

Beam-induced target depolarization at RHIC and EIC

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July 3, 2024

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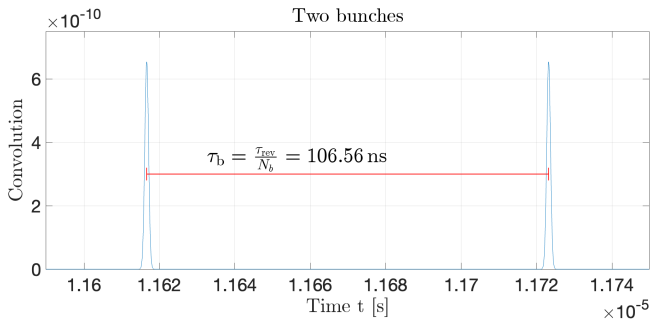
Introduction

- Development of polarized hydrogen jet (HJET) for RHIC finished 20 yrs ago.
- Many details on the technical structure and operation cannot be found in the literature, and there is also no comprehensive publication.
- Hardly anyone around to ask:
 - Tom Wise send me some unfinished paper drafts and other material.
 - I am in touch with Alexander Nass about BRP operation and most recent BRP measurement (from 2004).
- **Today's item:**
 1. **Beam-induced depolarization of target atoms**
 - Bunch repetition frequency will be much larger at EIC than at RHIC
 - Goal: Understand corresponding situation at EIC

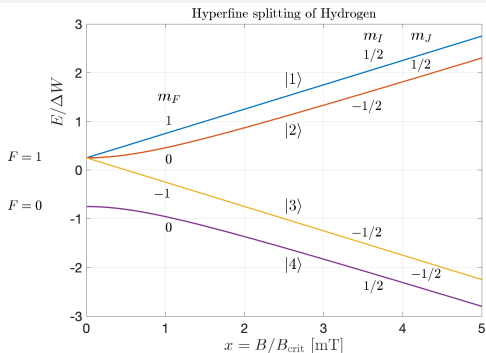
Bunch structure

RHIC situation:

- Time period between two adjacent bunches: $\tau_b = \frac{\tau_{\text{rev}}}{N_b} = 106.57 \text{ ns}$
- Number of stored bunches $N_b = 120$
- Bunch frequency $f_b = \frac{1}{\tau_b} = 9.3831 \text{ MHz}$
- Large number of harmonics contribute to induced magnetic high-frequency field close to RHIC beam, as bunches are short ($\sigma_t = 0.5 \text{ ns}$)



Hyperfine states of hydrogen



Critical field B_c (slide 23)

- Zeeman energy $g_J \mu_B B$ comparable to E_{hfs}
- $E_{\text{hfs}} \approx 5.874 \times 10^{-6} \text{ eV}$ ($\approx 1420 \text{ MHz}$ [1]):
- $B_c = 50.7 \text{ mT}$

Transition frequencies

- Transition frequency between two hyperfine states $|i\rangle$ and $|j\rangle$ given by:

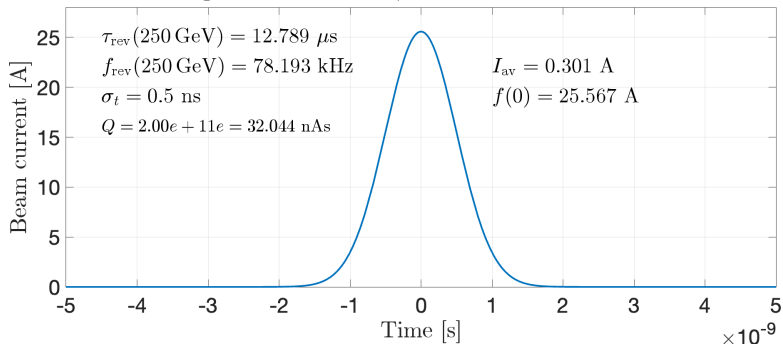
$$f_{ij} = \frac{E_{|i\rangle}(B) - E_{|j\rangle}(B)}{h} \quad (1)$$

- When f_{ij} matches one of the beam harmonics at a certain holding field B , resonant depolarization occurs [2].

Single bunch distribution

- (Gaussian) bunch in RHIC

Single Gaussian bunch, RHIC nominal conditions



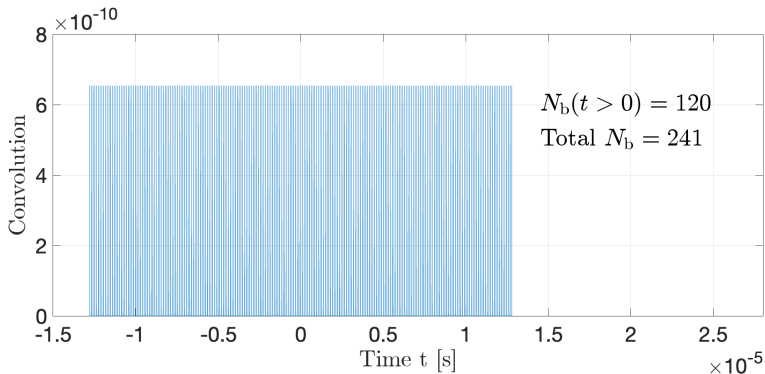
Pulse shape described by

$$f(t) = \frac{Q}{\sqrt{2\pi}\sigma_t} \exp\left(-\frac{t^2}{2\sigma_t^2}\right) \quad (2)$$

Gaussian convoluted with (finite) series of delta functions

Total beam current as function of time t given by

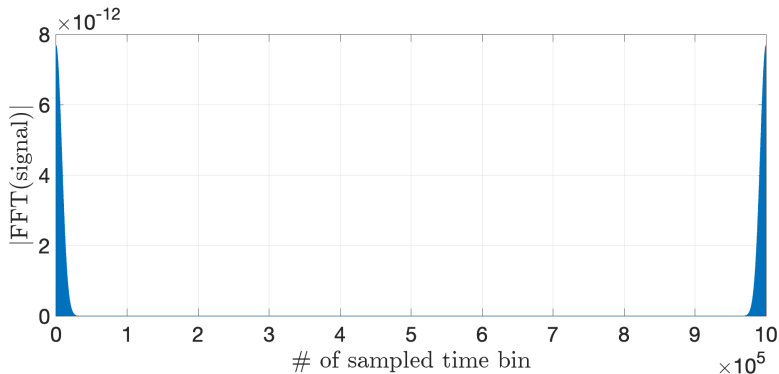
$$I(t) = \int_{-\infty}^{\infty} f(t - \xi) \sum_{k=-\infty}^{\infty} \delta\left(\xi - k \frac{\tau_{\text{rev}}}{N_b}\right) d\xi \quad (3)$$



Radiofrequency-fields

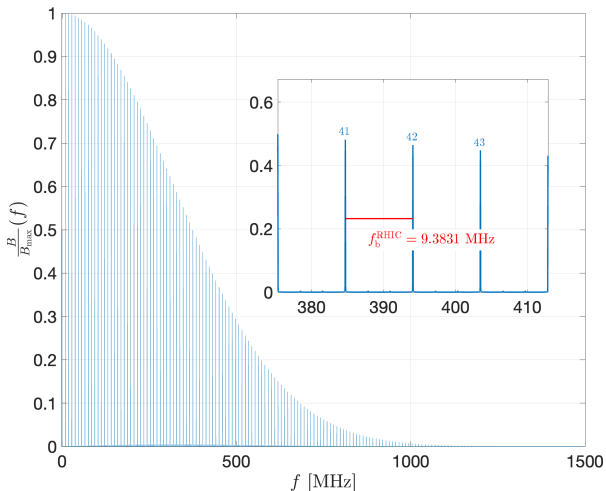
FFT of convolution

- Two-sided amplitude spectrum of FFT of the convolution



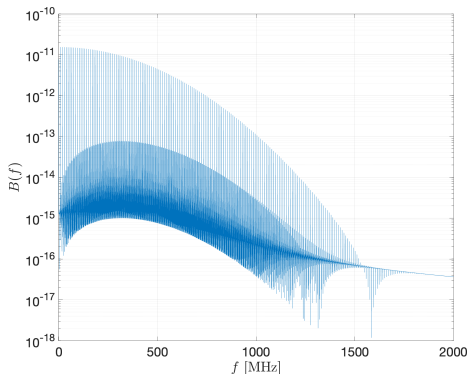
Produced radio-frequency fields

- Single-sided amplitude spectrum of FFT
- x-axis converted to frequency



Amplitudes of magnetic RF fields

- Same, but logy
- FFT background $\leq 1\%$
 - not at f_{rev}
 - not from finite set of δ fcts
 - \rightarrow probably numerical from FFT



Transition frequencies between hyperfine states of H

Based on Zeeman splitting, shown on slide 5, using Eq. (1)

- Determine transition frequencies f_{ij} between hyperfine states $|i\rangle$ and $|j\rangle$.
- Classification refers to change of quantum numbers (see Ramsey [3]):
 - B_0 is static field, B_1 is RF field that exerts torque on magnetic moment μ :
 - π ($B_1 \perp B_0$) transitions *within one F* multiplet:

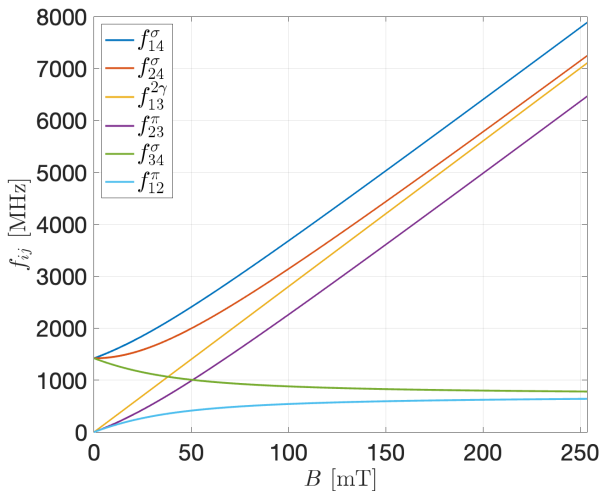
$$\Delta F = 0, \quad \Delta m_F = \pm 1. \quad (4)$$

- σ ($B_1 \parallel B_0$) transitions *between different F* multiplets:

$$\Delta F = \pm 1, \quad \Delta m_F = 0, \pm 1. \quad (5)$$

- Single photon transitions in H: f_{12}^π , f_{23}^π , f_{14}^σ , f_{24}^σ , and f_{34}^σ .
- Transition $f_{13}^{2\gamma}$ with $\Delta m_F = 2$ requires two photons.

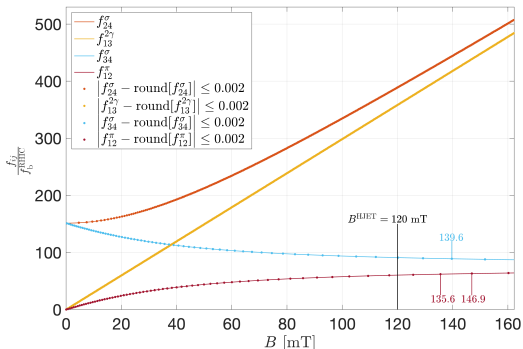
Transition frequencies between hyperfine states of H



Hydrogen hyperfine transitions from bunch fields

Depolarization when f_{ij} multiple of bunch frequency f_b^{RHIC}

- HJET injects states $|1\rangle + |4\rangle$ (p^\uparrow) and $|2\rangle + |3\rangle$ (p^\downarrow).
 - What is actual orientation of B^{HJET} ?

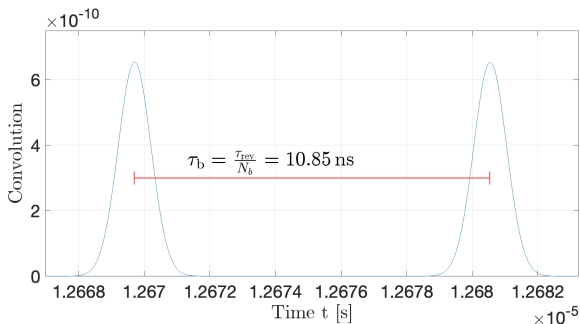


- $\left| \frac{f_{ij}}{f_b^{\text{RHIC}}} - n \right| \leq 0.002, n \in \mathbb{N}$
- Same $m_l \Rightarrow f_{14}^\sigma, f_{23}^\pi$ omitted

Bunch structure

EIC situation:

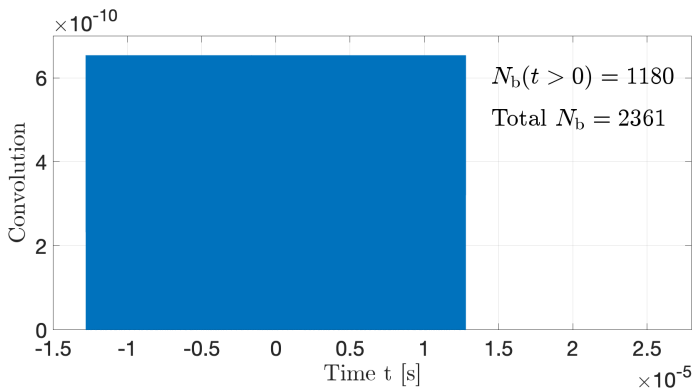
- Time period between two adjacent bunches: $\tau_b = \frac{\tau_{\text{rev}}}{N_b} = 10.85 \text{ ns}$
- Number of stored bunches $N_b = 1180$
- Bunch frequency $f_b = \frac{1}{\tau_b} = 92.2081 \text{ MHz}$



Gaussian convoluted with (finite) series of delta functions

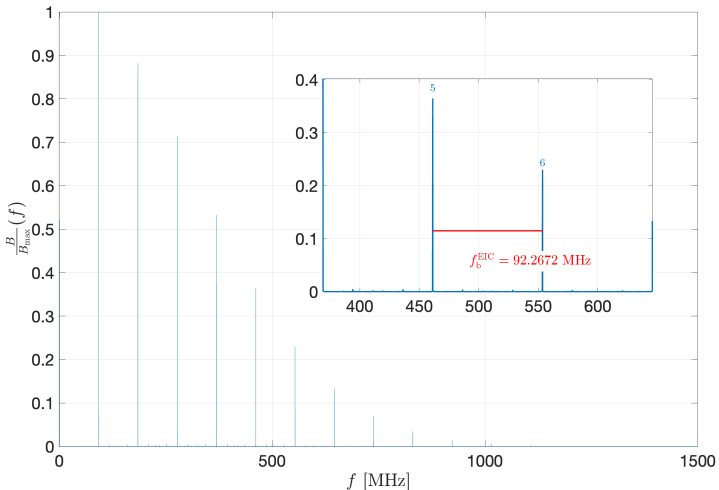
Total beam current as function of time t given by

$$I(t) = \int_{-\infty}^{\infty} f(t - \xi) \sum_{k=-\infty}^{\infty} \delta\left(\xi - k \frac{\tau_{\text{rev}}}{N_b}\right) d\xi \quad (6)$$



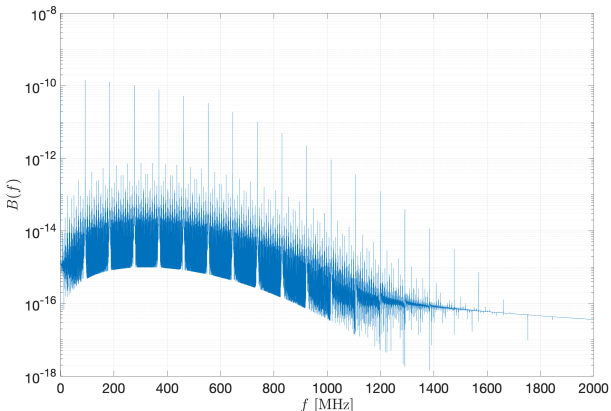
Produced radio-frequency fields

- Single-sided amplitude spectrum of FFT



Amplitudes of radio-frequency fields

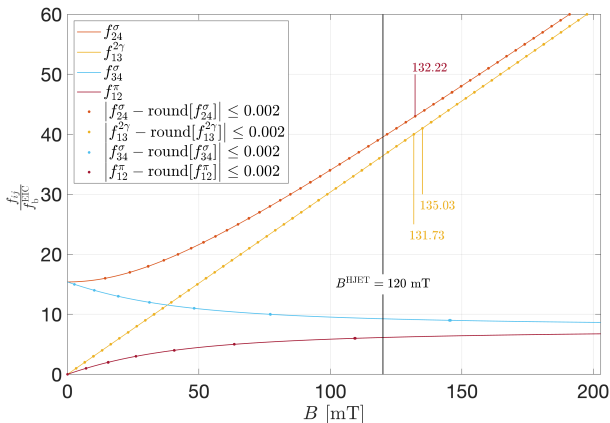
- RF field amplitudes at EIC $\approx 10\times$ larger compared to RHIC
 \Rightarrow increased transition probability due more photons ($n_\gamma \propto B^2$).
- Frequency spacing in spectrum will become much larger
 \Rightarrow fewer contributing resonances.



Hydrogen hyperfine transitions from bunch fields

Depolarization occurs when $\left| \frac{f_{ij}}{f_b^{\text{EIC}}} - n \right| \leq 0.002$, $n \in \mathbb{N}$

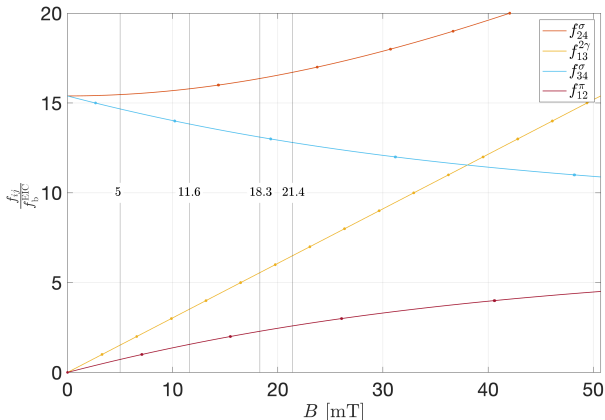
- For $B < 200$ mT, all transitions below harmonic number 50 will contribute!



Hydrogen hyperfine transitions from bunch fields

How about the small B region?

- At RHIC, this region was inaccessible, as spacing of $f_{13}^{2\gamma} \approx 0.3$ mT.
- At EIC, at ≈ 5 mT, spacing is of $f_{13}^{2\gamma} \approx 3.3$ mT.



Conclusion and outlook

Conclusion

- Depolarization:
 - RHIC: harmonic numbers > 100 were ignored.
 - EIC: All depolarization transitions of H appear at harmonic numbers < 50 .
 - EIC: If we stay in 120 mT region, field accuracy $\lesssim 1$ mT required.
 - Contrary to situation at RHIC, smaller B may be an option for EIC.
- Extend study of beam-induced depolarization to D and ^3He atoms.

Outlook: Other things that need to be looked into

1. How does beam-induced \vec{B} field affect target polarization? \rightarrow in progress
2. Zero-crossings along the vertical axis of the source.
 - Revisit magnetic field calculations of HJET holding field \rightarrow in progress.
3. Polarization measurements with all transition units in ABS and BRP.
4. Tracking of atoms in HJET & sextupole magnet systems
 - Recuperated tracking code used for HJET design from Michelle/Paolo \rightarrow to make it work is time consuming

References I

- [1] M. Diermaier et al., *Nature Commun.* **8**, 5749 (2017).
- [2] A. Airapetian et al., *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment* **540**, 68 (2005).
- [3] N. Ramsey, *Molecular Beams*, Oxford University Press, 1956.

Spare slides

Critical field for hydrogen hyperfine splitting I

Zeeman region:

- magnetic flux density at which energy separation between different hyperfine levels becomes comparable to Zeeman splitting.
- referred to as *critical magnetic field* or *Breit-Rabi field* B_c
- Breit-Rabi formula (energy levels of hydrogen atom in external magnetic field:

$$E_{F,m_F} = -\frac{E_{\text{hfs}}}{2(2I+1)} + g_J \mu_B m_J B \pm \frac{E_{\text{hfs}}}{2} \sqrt{1 + \frac{2m_F x}{F} + x^2}, \text{ where} \quad (7)$$

- E_{hfs} is hyperfine splitting energy
- I is nuclear spin (for H, $I = \frac{1}{2}$)
- g_J is Landé g-factor
- μ_B is Bohr magneton
- m_J is magnetic quantum number
- m_F is total angular momentum quantum number
- $x = \frac{g_J \mu_B B}{E_{\text{hfs}}}$
- $F = I + J$ is total angular momentum (for H, $J = \frac{1}{2}$)

Critical field for hydrogen hyperfine splitting II

For H:

- hyperfine splitting energy E_{hfs} (1420 MHz):

$$E_{\text{hfs}} \approx 5.874 \times 10^{-6} \text{ eV} \quad (8)$$

- Critical field B_c is when Zeeman energy $g_J \mu_B B$ is comparable to E_{hfs} . With $g_J \mu_B B_c \approx E_{\text{hfs}}$, we get:

$$B_c \approx \frac{E_{\text{hfs}}}{g_J \mu_B} \quad (9)$$

- For H, $g_J \approx 2$ (approximately for electron), and $\mu_B \approx 5.788 \times 10^{-5} \text{ eV/T}$. Thus,

$$B_c \approx \frac{5.874 \times 10^{-6} \text{ eV}}{2 \times 5.788 \times 10^{-5} \text{ eV/T}} \approx 50.7 \text{ mT} \quad (10)$$