

# SiPM Noise Simulation

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# The Challenge

SiPMs are noisy

- Single-photon signals at  $O(\text{MHz})$
- Even worse once irradiated

Since the SciFi readout is longitudinally segmented, calibration of individual channels can be challenging

The MIP peak provides an ideal signal to use for a channel-by-channel calibration

Want to keep the signal amplitude of the MIP above the noise floor from the SiPMs

# The Mitigation Strategies

Reduce the SiPM noise per channel, or increase the MIP signal per channel

Noise Reduction Option	Pro	Con
S13 vs. S14	Easy integration	Less Noise = Worse PDE
Reduce Temperature	Strong noise reduction	Hard integration
Replace SiPMs	Fresh sensors	\$\$\$
Anneal SiPMs	Undo radiation damage	Hard integration(?)
Smaller readout channels	Better granularity	MIP shrinks too

MIP Improvement Option	Pro	Con
S13 vs. S14	Easy integration	More PDE = More Noise
Silicone Cookie	Easy integration	?
Better SciFis	Less attenuation, more light	\$\$\$\$\$\$\$
Larger readout channels	More light	Noise $\propto$ SiPM Area

For more detail on these, see [our session in the Jan. collaboration meeting](#)

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Let's investigate these options a bit more...

# Readout Channel Splitting

Investigate the effect of splitting our current  $1.2 \times 1.2 \text{ cm}^2$  channels into 2 or 4 DCR scales as SiPM area while MIP signal scales as path length

- Splitting in four reduces DCR by a factor of 4, signal by a factor of 2
- Splitting in two reduces DCR by a factor of 2, keeps MIP signal  $\sim$  the same

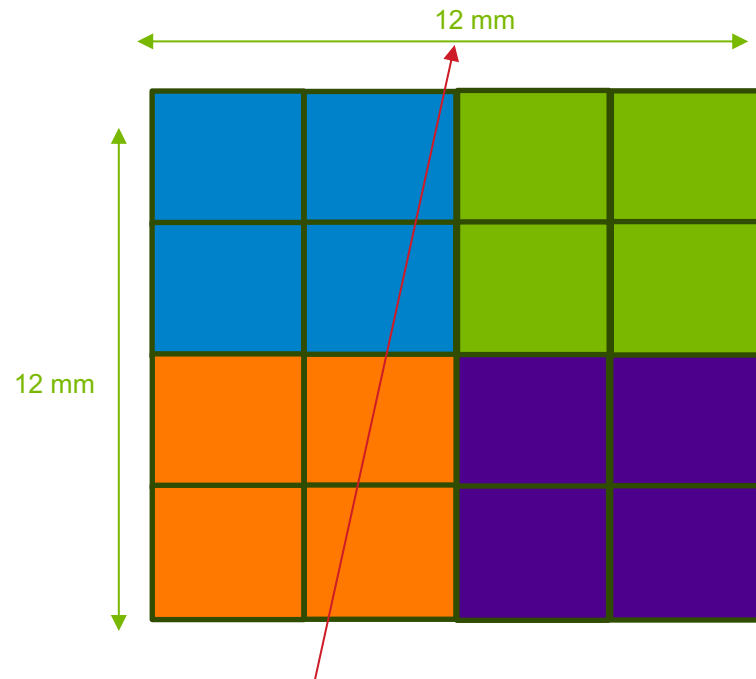
DCR is *not* linearly proportional to where can you set your threshold, needs to be simulated!

- DCR proportional to mean of pedestal, but we need to know the sigma!

Additional benefits of smaller channels:

- Reduces pulse size going into HGCROC
- Improves granularity for e/pi separation
- Reduces needed light guide length (?)

Splitting in two = sum orange + blue and green+purple



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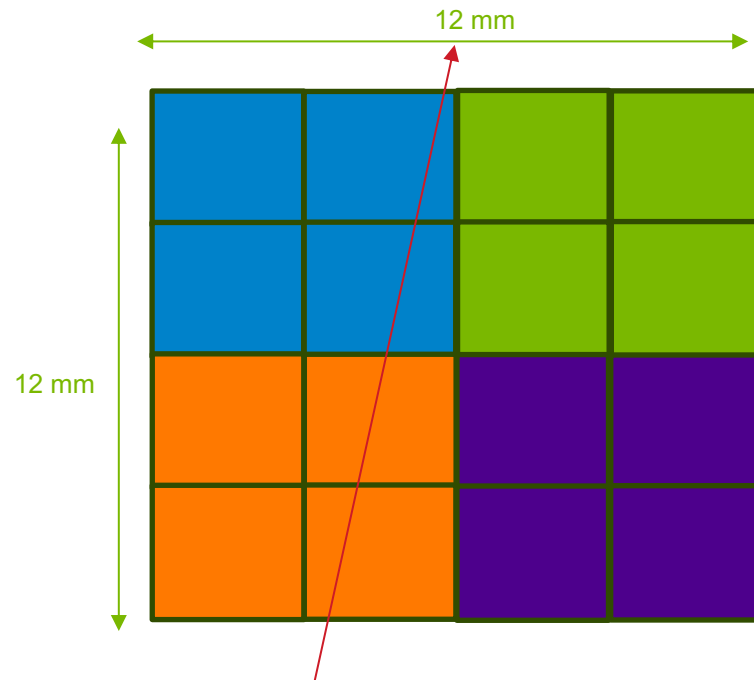
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This is probably worth a whole other talk...

# A word on annealing...

INFN Bologna successfully tested in-situ annealing via forward-bias of the SiPM

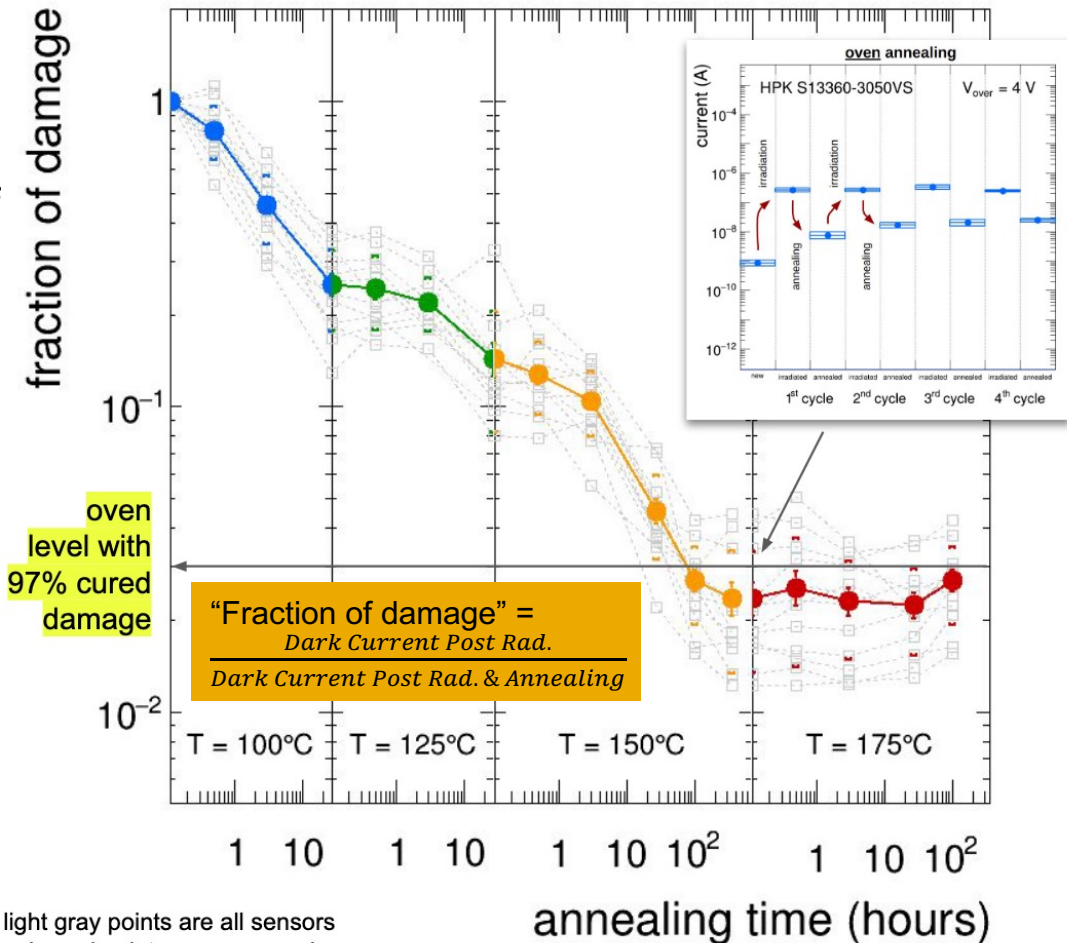
- Run a current through the diode until it heats to the desired temperature

This provides a convenient alternative to removing the SiPMs from the detector and putting them in an oven!

Achieved a factor ~50 (!!!) reduction in dark current after several days of heating at 150° C

- **Very** attractive option for us
- **IF** we can integrate it in the ESB

## online self-annealing with forward bias



[Click here for Roberto's talk](#)

# Simulation Technique

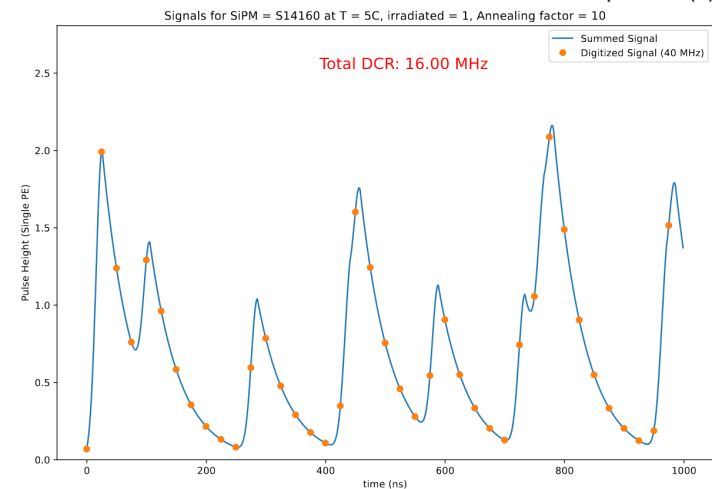
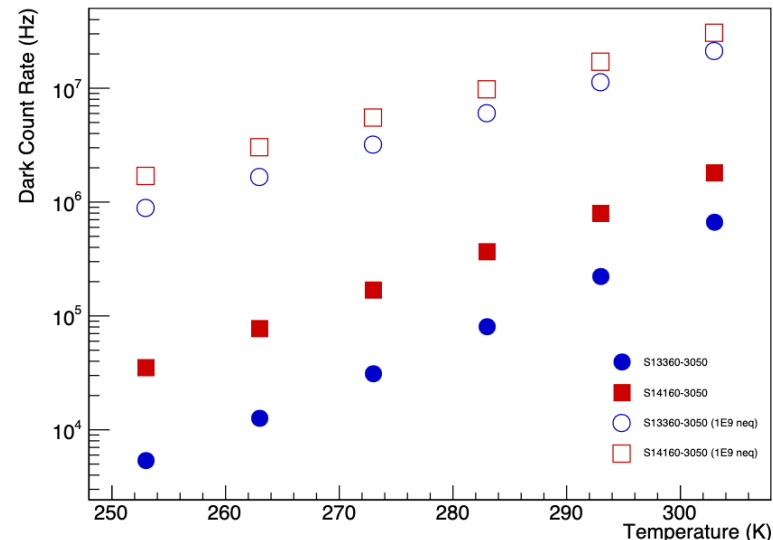
Throw waveforms at the dark count rates supplied to us by INFN Bologna for S14160-3050 SiPMs

- Various temperatures provided
- Both unirradiated and irradiated with  $1E9$  1 MeV neutron equivalent dose

Include the crosstalk numbers supplied by Hamamatsu

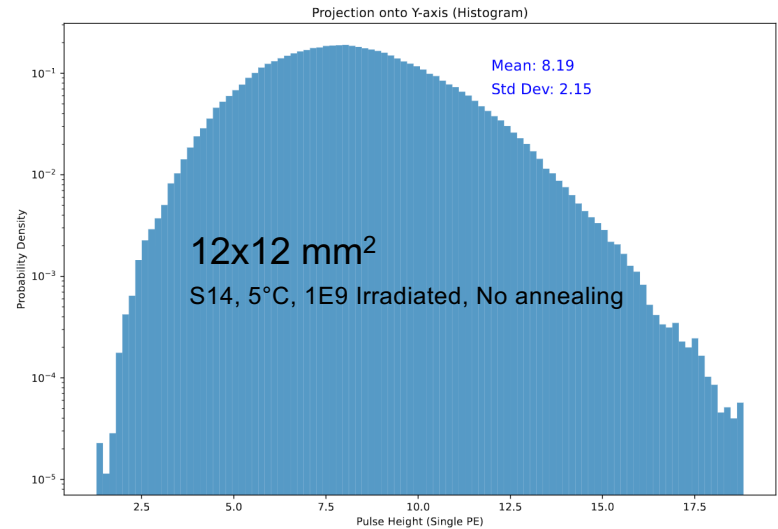
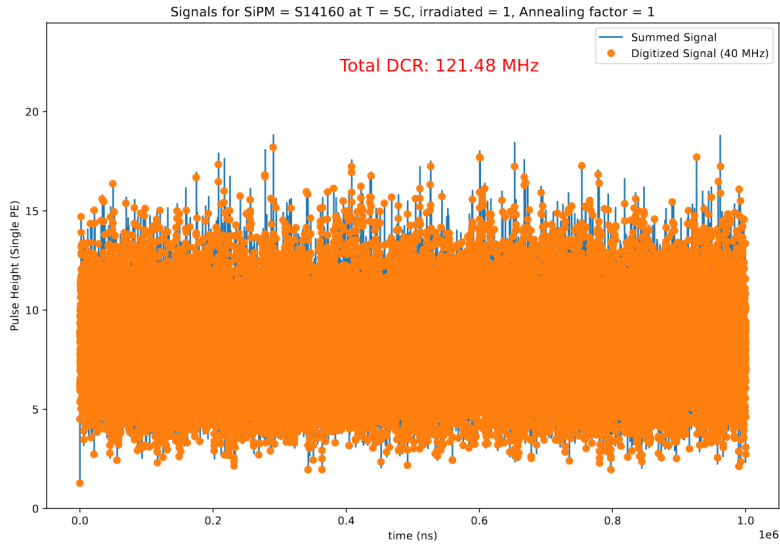
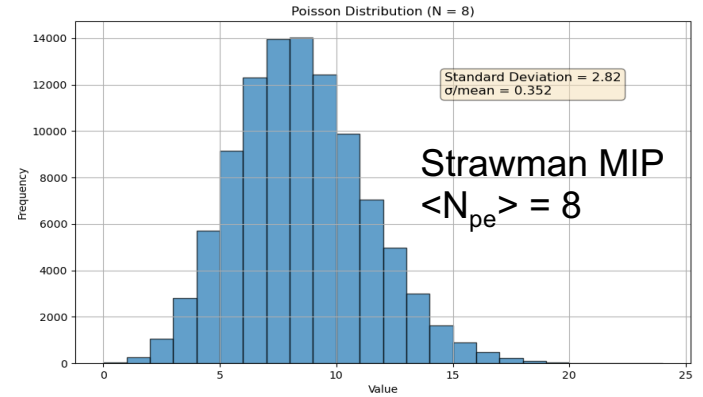
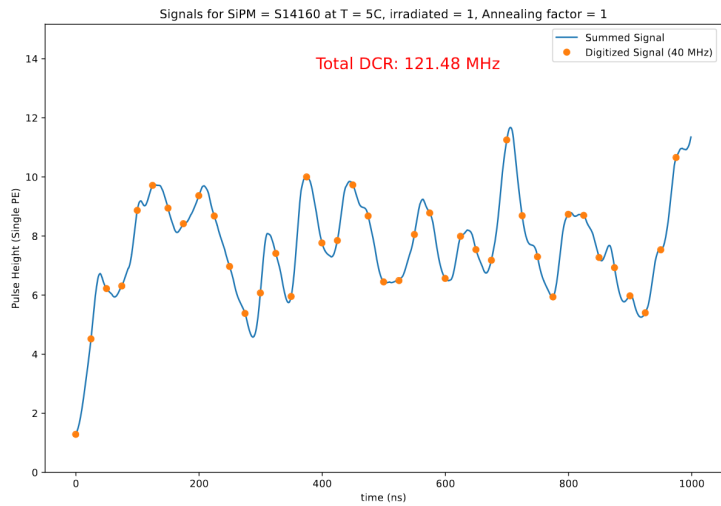
- This is important to determine the sigma of the noise!

Focus on irradiated sensors, since this is the concerning case for the MIP

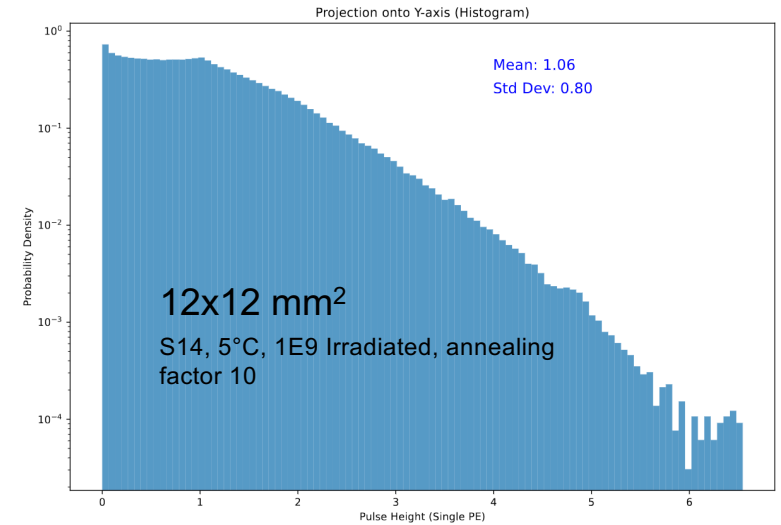
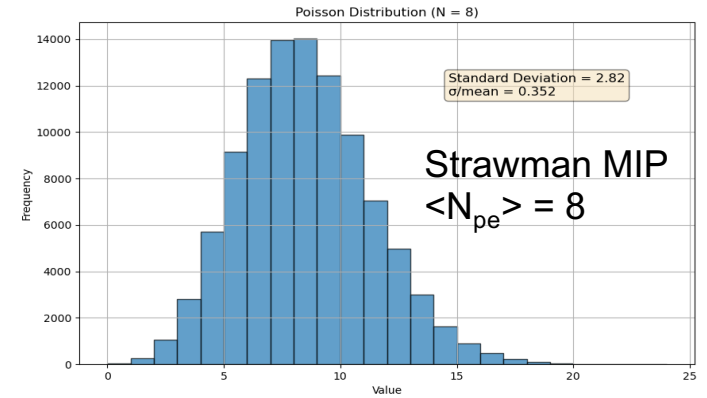
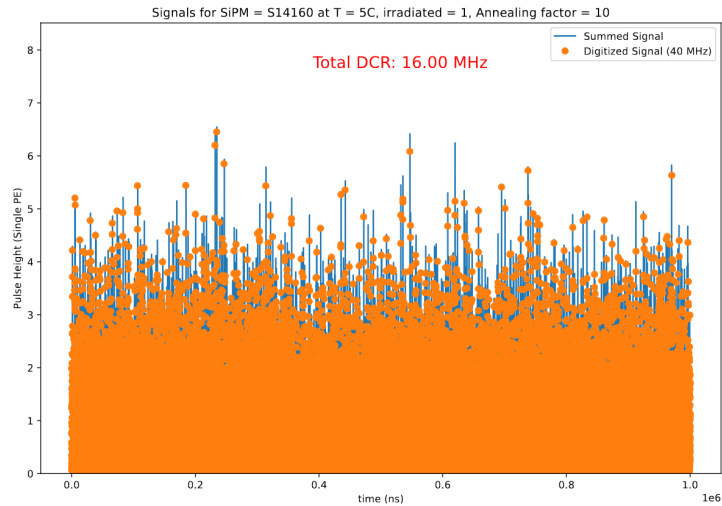
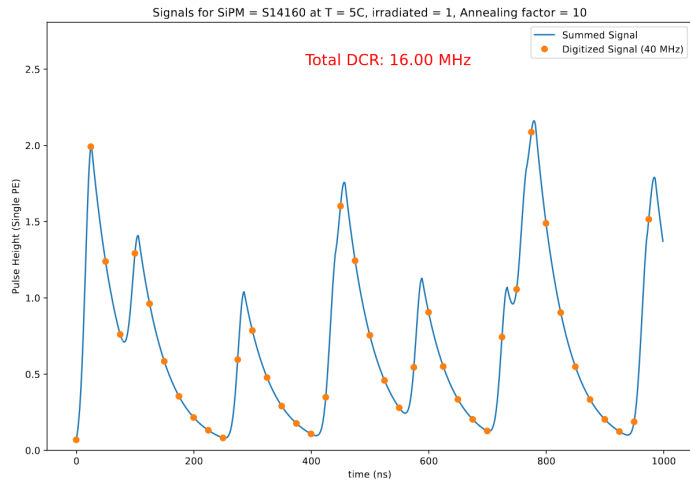


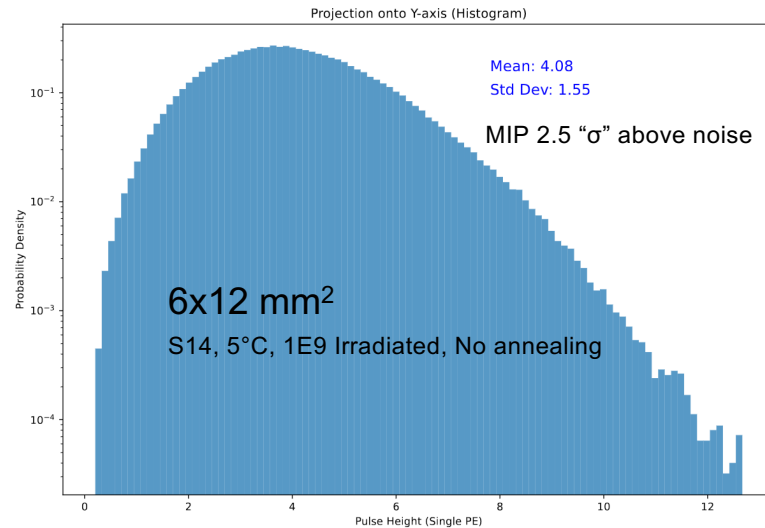
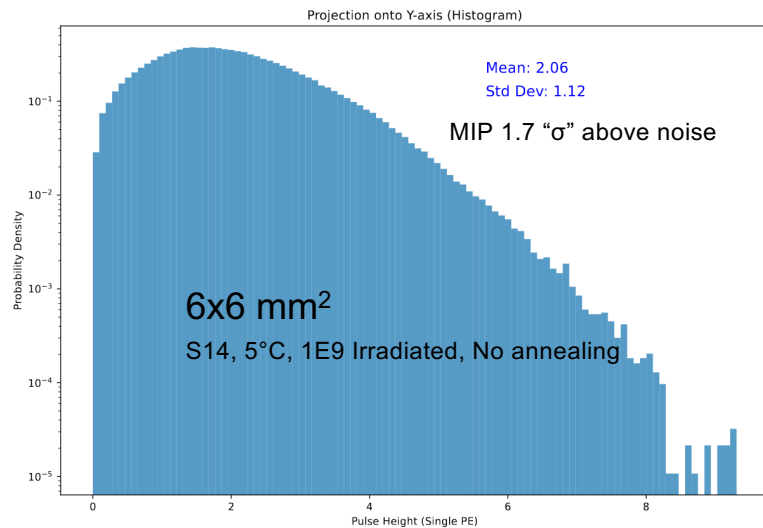
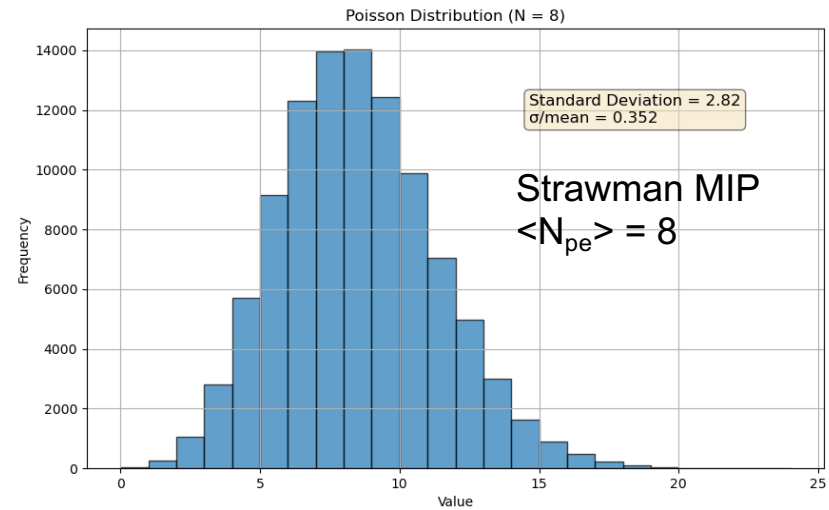
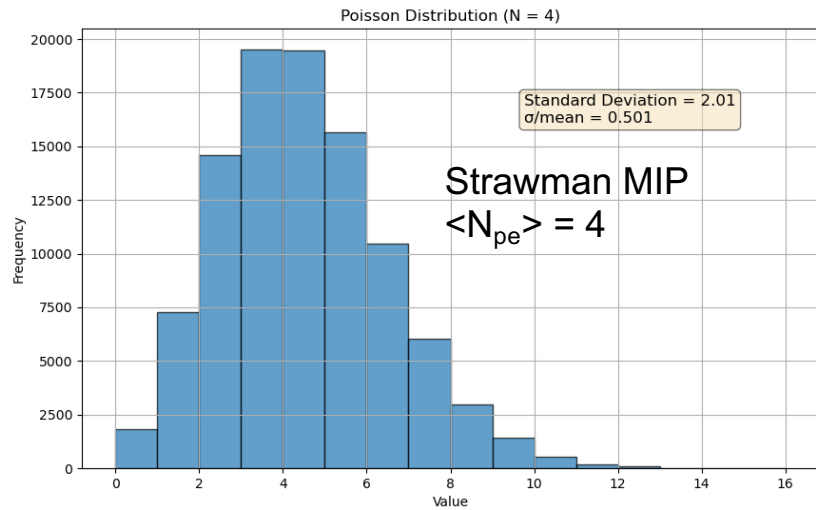


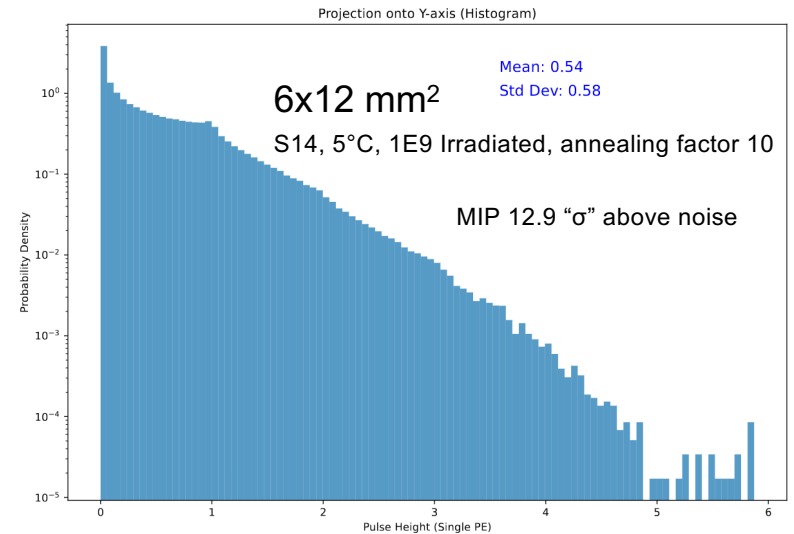
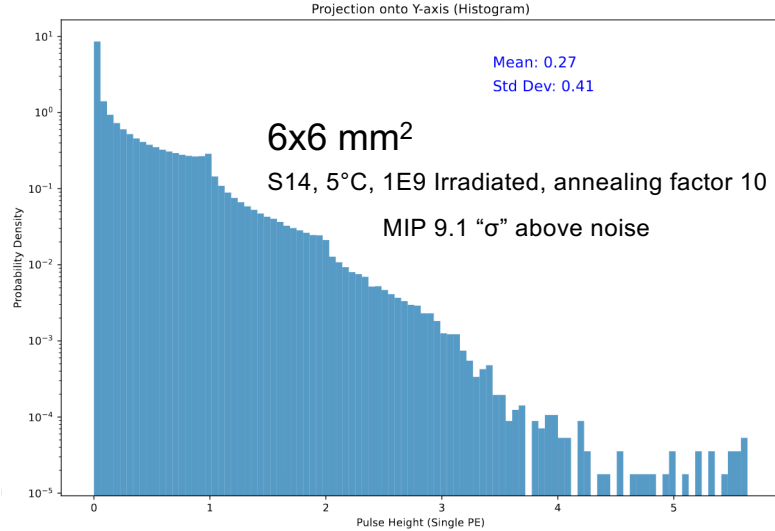
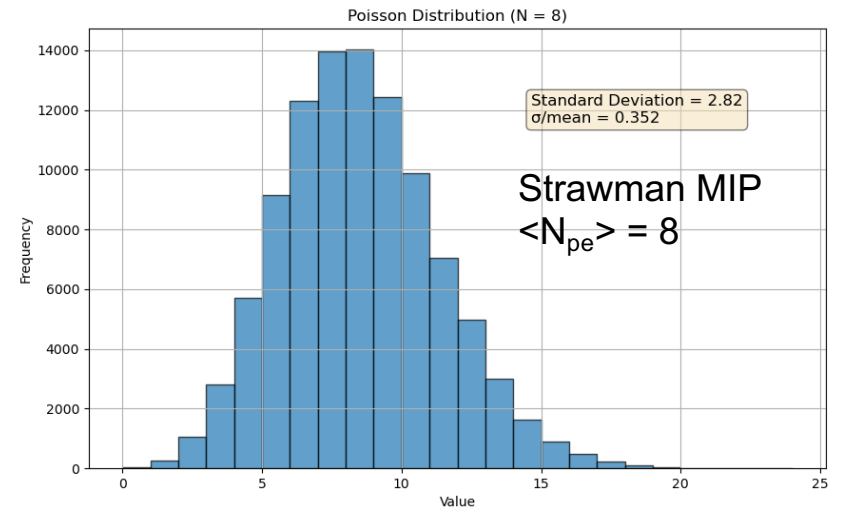
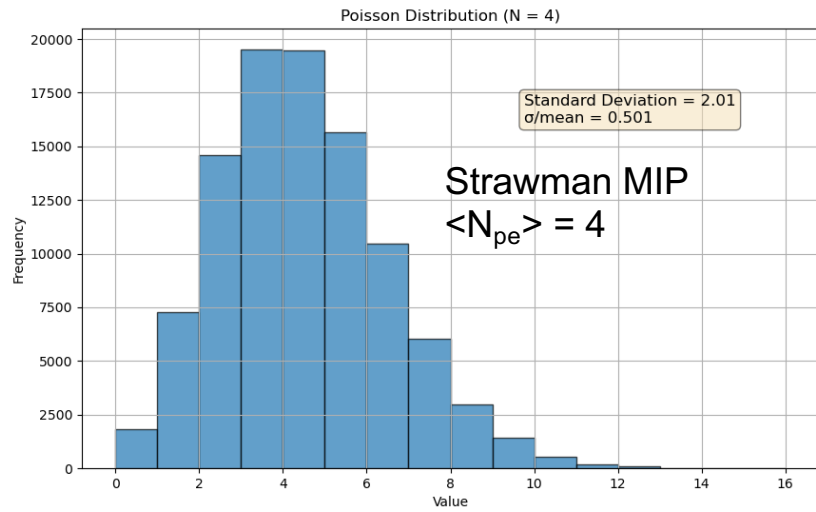
# Baseline – 12 x 12 mm<sup>2</sup>

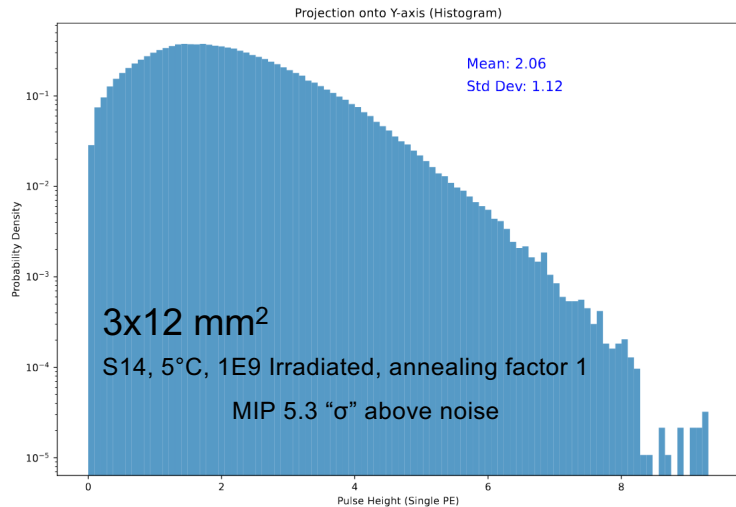
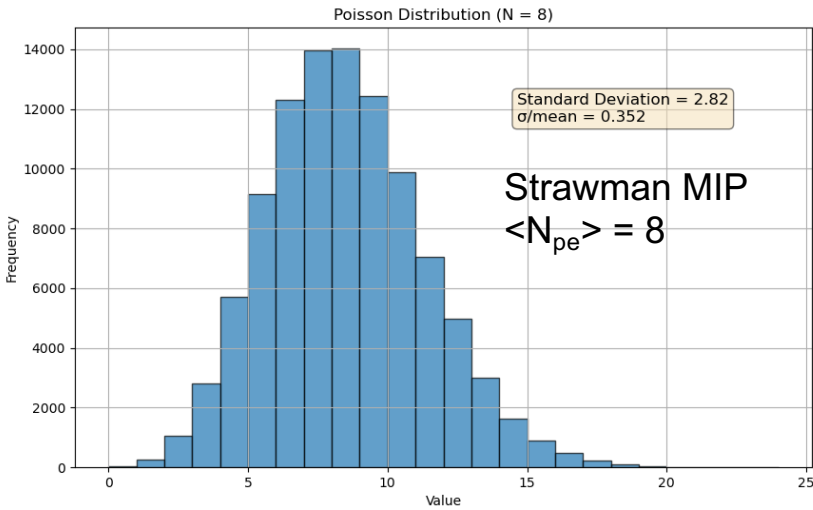


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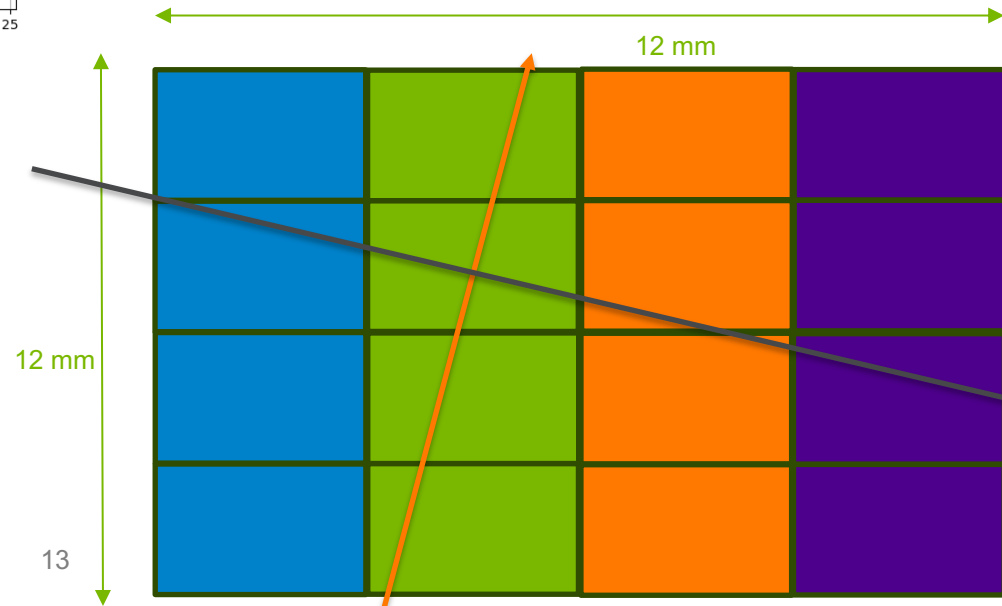


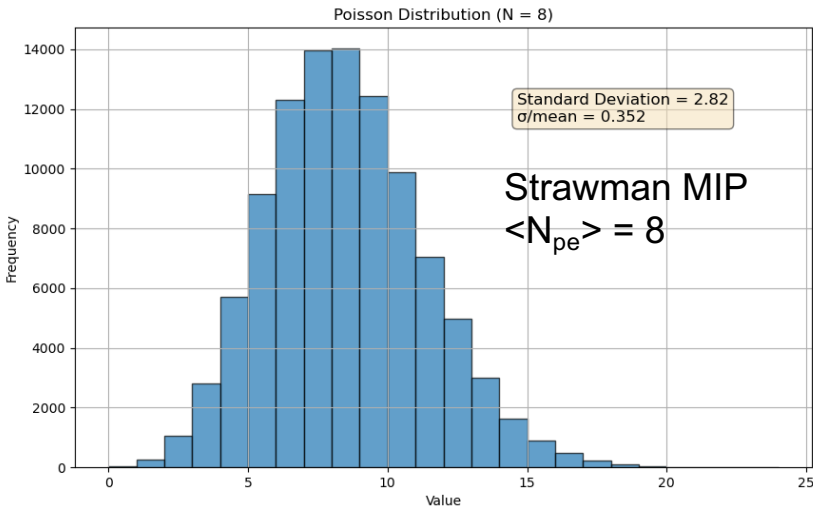


# To go even further...

Note that this case will only really be true for field-off runs of the collider

- Cosmics will rarely hit in the “aligned” orientation

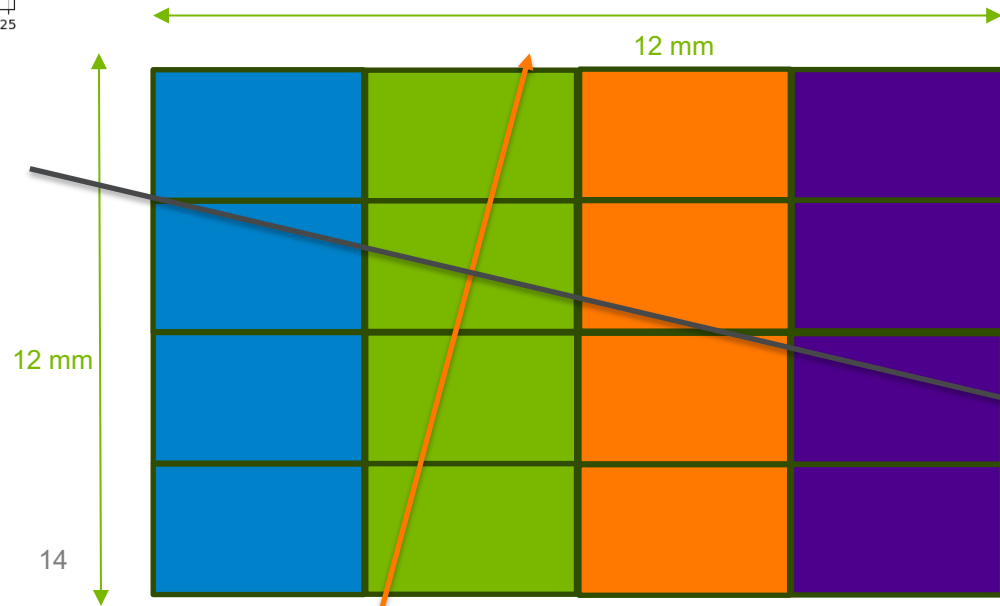
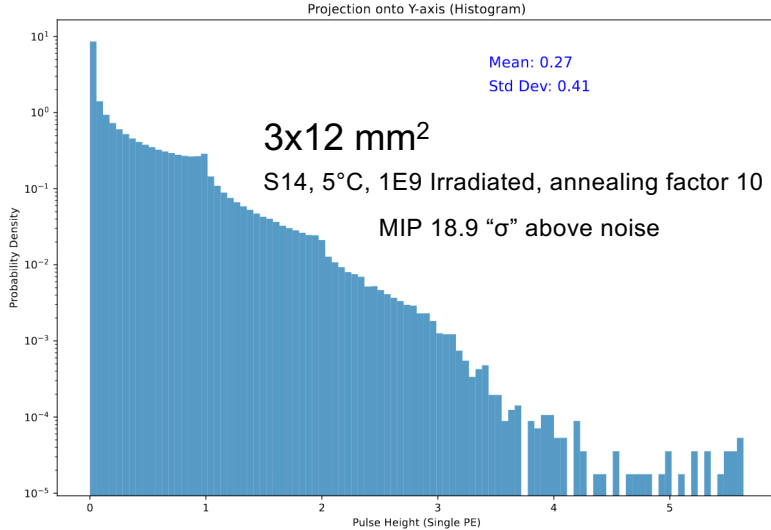




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# Conclusions

- If the annealing factor of 10 can be achieved, all schemes seem acceptable
- If not, the 12 mm x 12 mm and 6 mm x 6 mm are in trouble, 6 mm x 12 mm looks okay still
  - The 6x12 mm option provides decent headroom at little cost
    - Requires individual 6x6mm SiPMs and not arrays

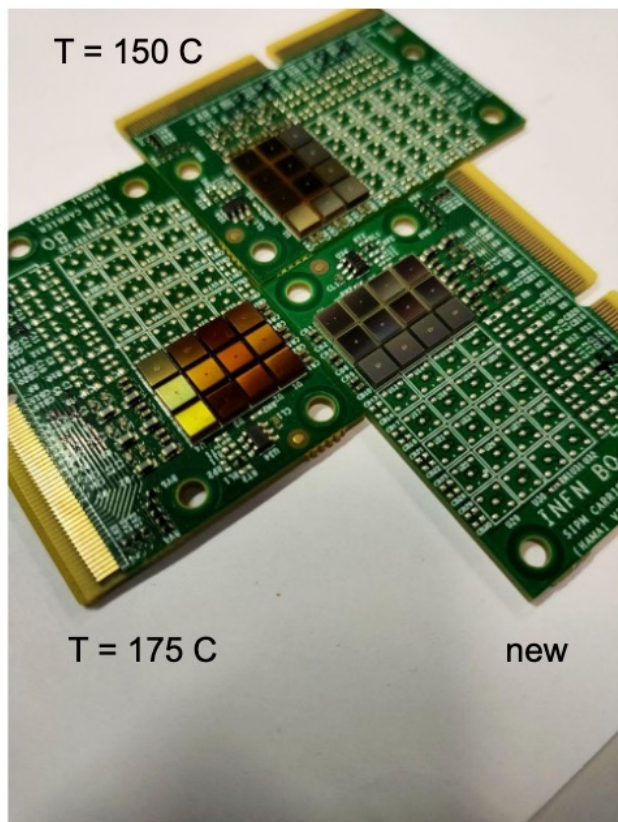
## Caveats

- We assumed as our strawman an average of 8 p.e. for the MIP
  - Single-clad Kuraray was something like 6.4 p.e. on average in Maria's previous sims, Luxium was 5.1 p.e.
  - Tegan's light guide studies suggest we gain more from the silicone cookie than we assumed at that time
  - 6.4 would be okay for the 2x1 split
- No investigation of splitting effect on signal shape!
  - For fixed external circuitry, the smaller arrays will have faster signals (SiPM capacitance)
  - This will reduce their sigma somewhat, but shorter signals will have fewer useful ADC samples in HGCROC

# Backup

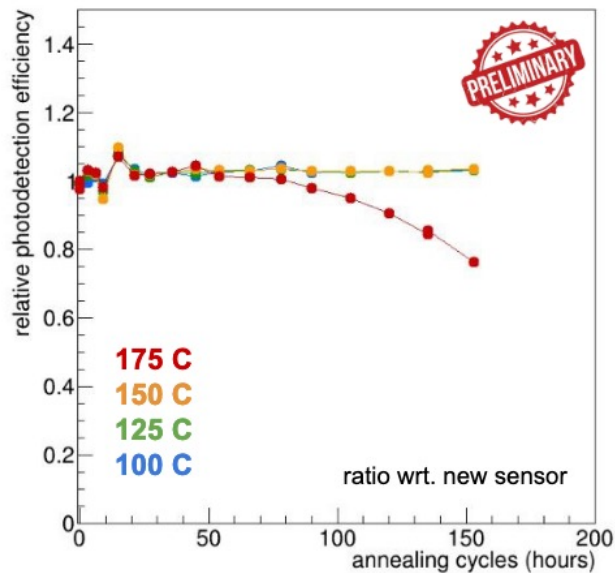


# Detailed studies of SiPM online self-annealing



## after many hours of online annealing

we noticed alterations on the SiPM windows  
in particular in one board that underwent  
500 hours of online annealing at  $T = 175\text{ C}$   
the sensors appear "yellowish" when compared to new



detailed studies are ongoing, preliminary results indicate efficiency loss after 100 hours of annealing at  $T = 175\text{ C}$ . **lower temperatures unaffected up to 150 hours**

# Backup Backup (January)

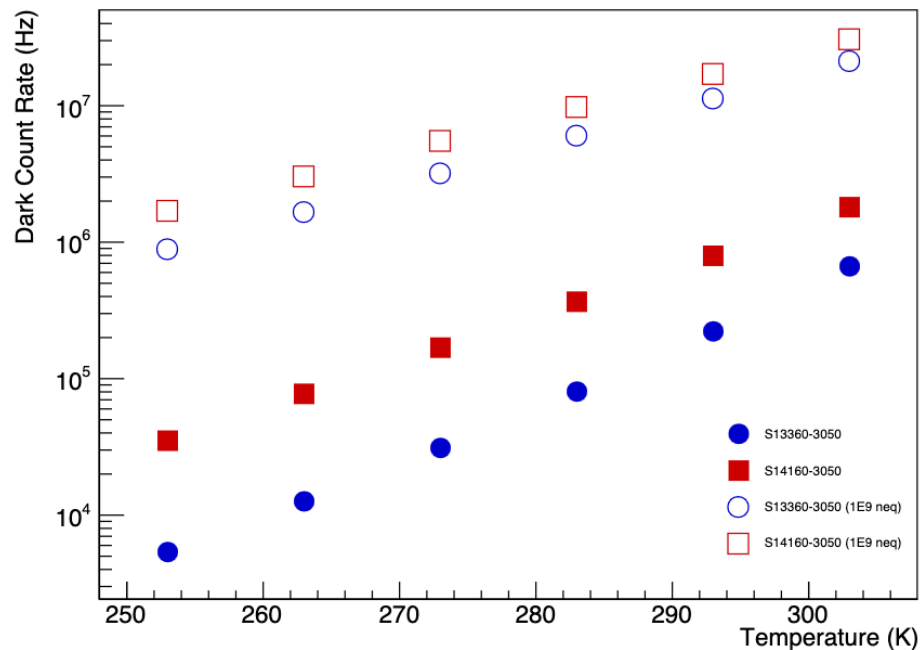
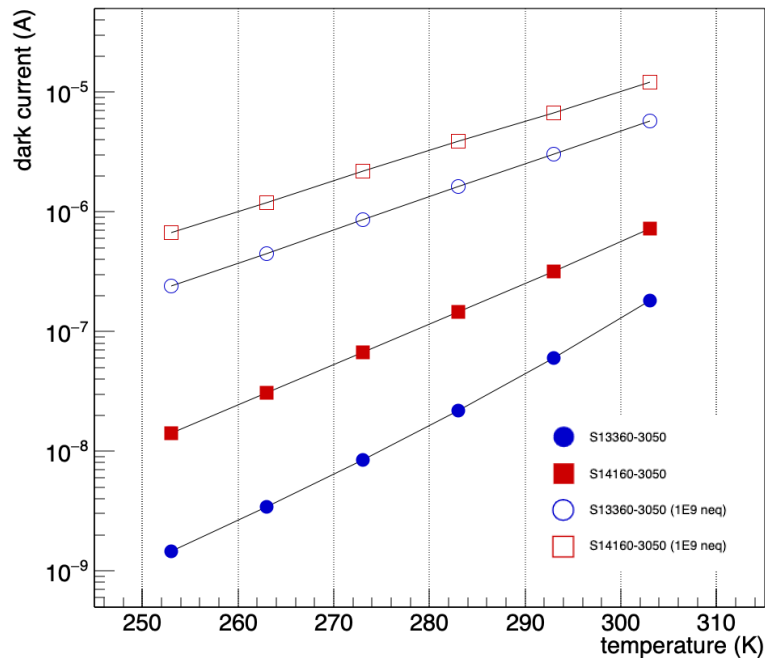
# Dark Currents from INFN Measurements

Scale currents by gain &  $q_e$  to get a dark count rate

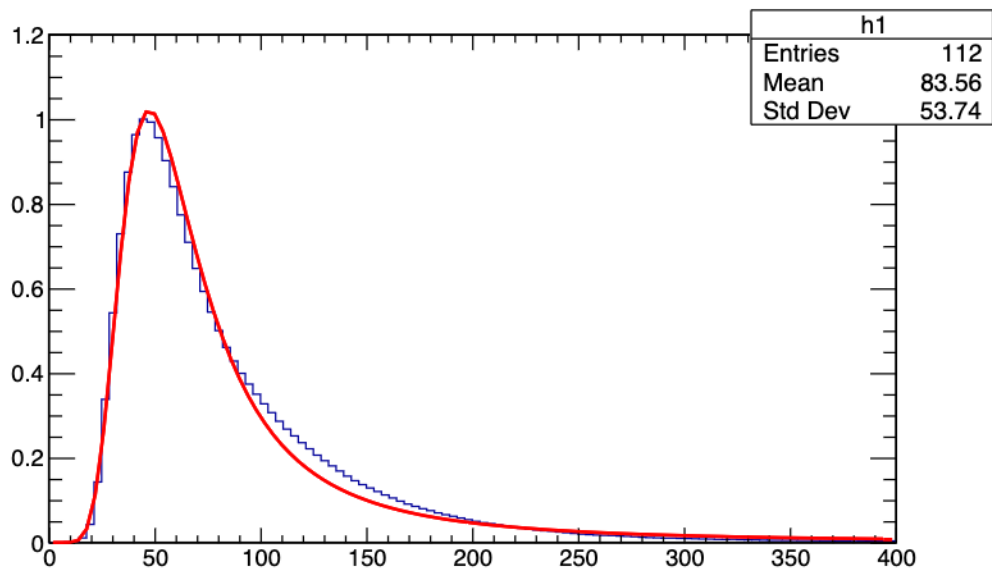
- This is a bit naïve, but Roberto indicated it should be alright

Order of magnitude increase after  $1E9$  1 MeV neq dose

- Our conservative estimate of dose for lifetime of experiment was  $3E10$



# HGCROC Signals



Pulse from charge injection tests on HGCROC

- Kindly supplied by Norbert
- Digitized much faster than HGCROC actually digitizes
- Landau fit
- Peak set to 1 photoelectron

For now, I just do the stupid thing and throw these distributions with the frequency of the DCR

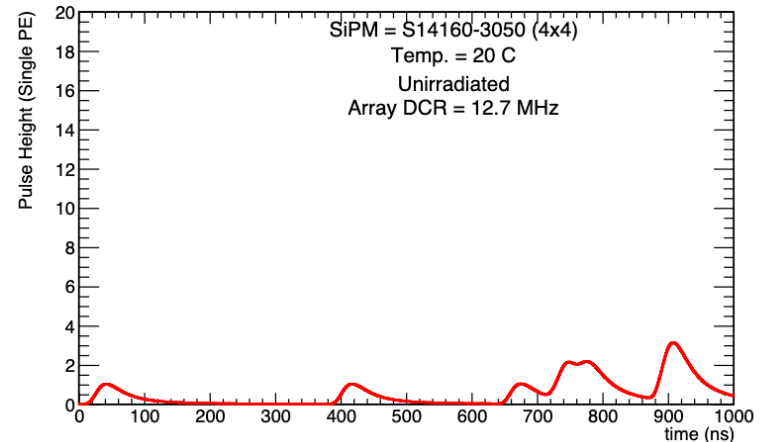
The reality of how often we will see hits from noise will depend on what ADC value we call a “hit”

# Waveforms

Created by Monte Carlo which throws Landau distributions with parameters taken from the HGCROC charge injection signal at different time slices with probability defined by the DCR

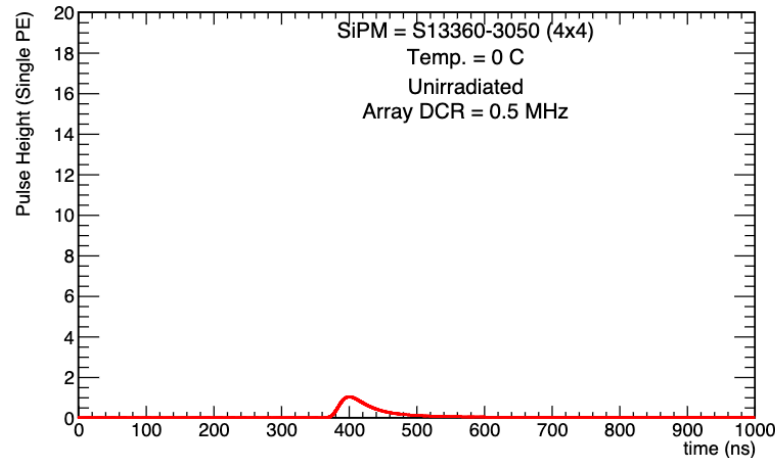
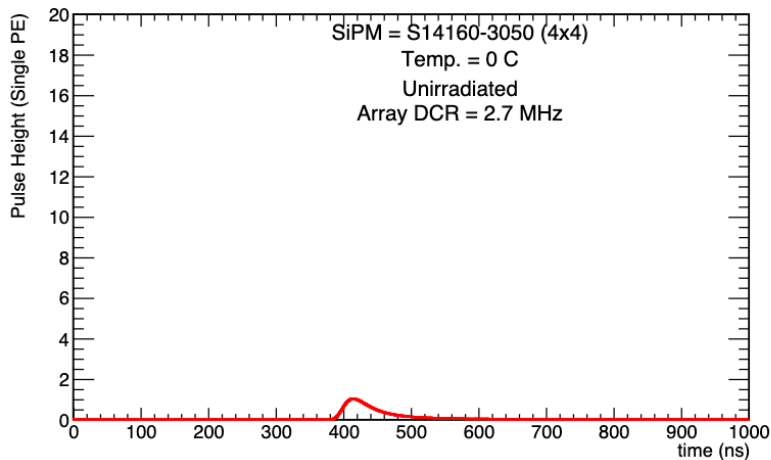
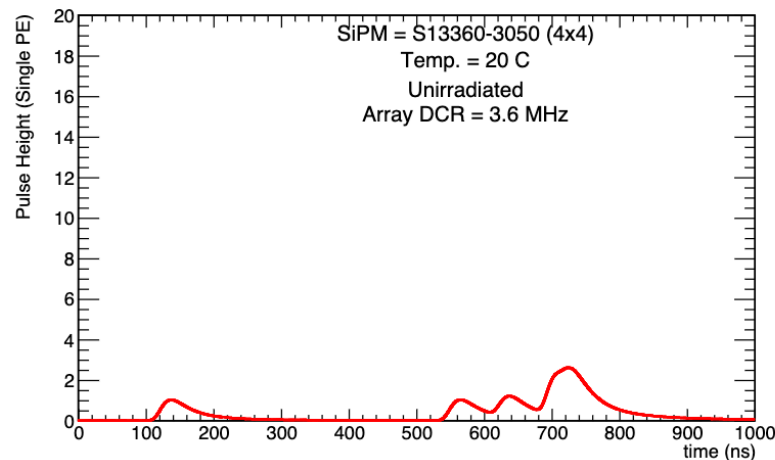
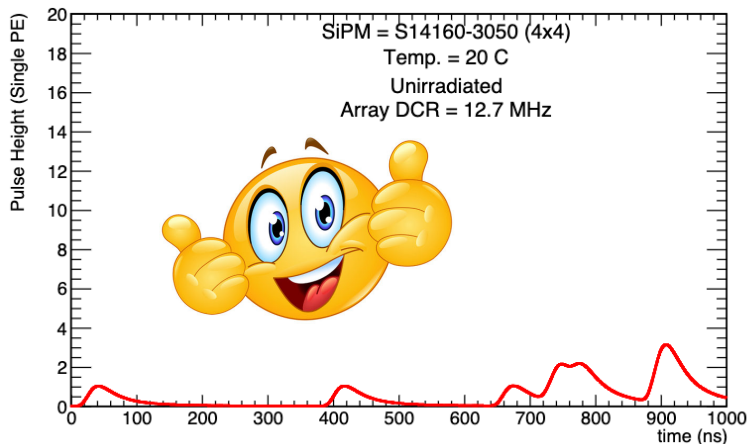
Crosstalk produces pulses 2x or 3x as high as the standard single-photon pulse

- Here a 7% crosstalk probability is assumed for both S14160 and S13360
- Final crosstalk numbers will depend on the optical coupling to the lightguide



# Temperature Dependence

Both sensors look reasonably good before irradiation, could easily set a threshold at 4 or 5 p.e.

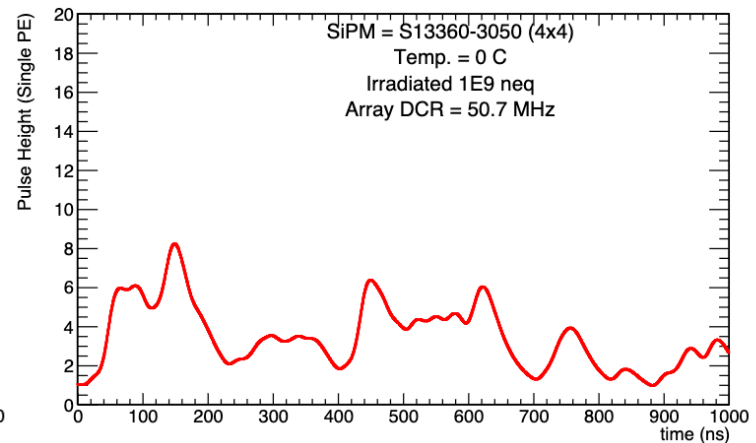
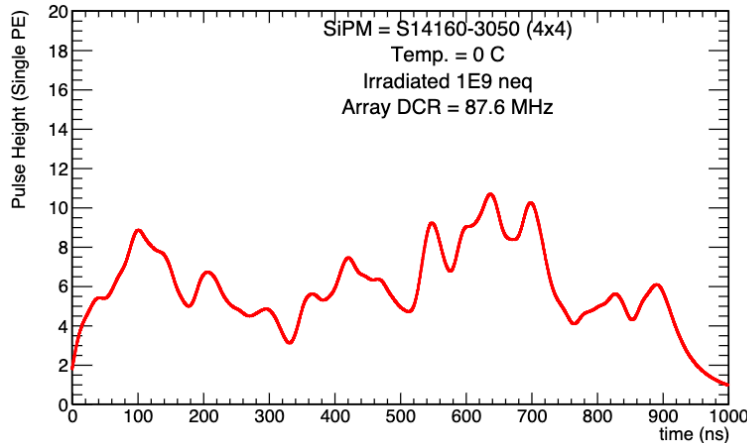
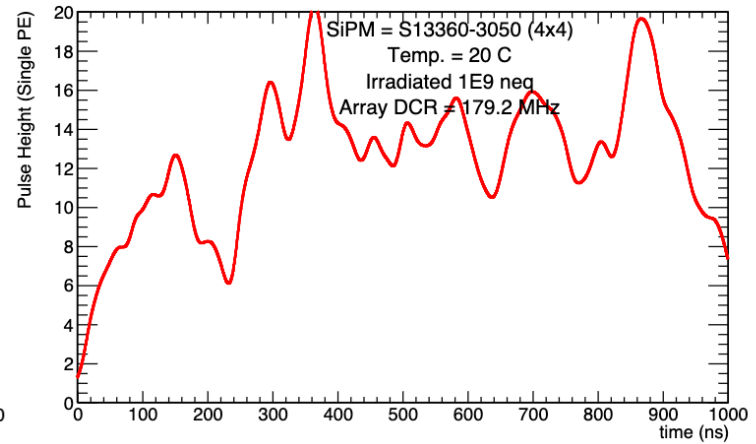
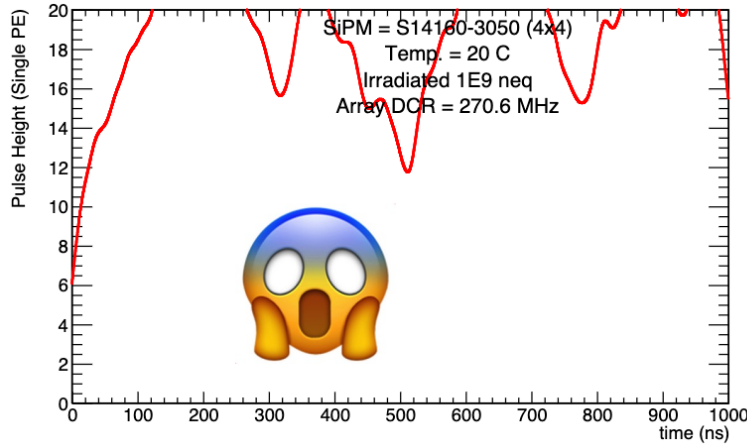


# Radiation Damage

0 °C not cool enough to save the MIP for either SiPM after 1E9 1 MeV neq dose

This assumes that the dark *current* produced by irradiation translates directly into dark *counts*, which is the **worst case** scenario

Keep in mind that this represents only ~1/30<sup>th</sup> of the 10-year expected dose



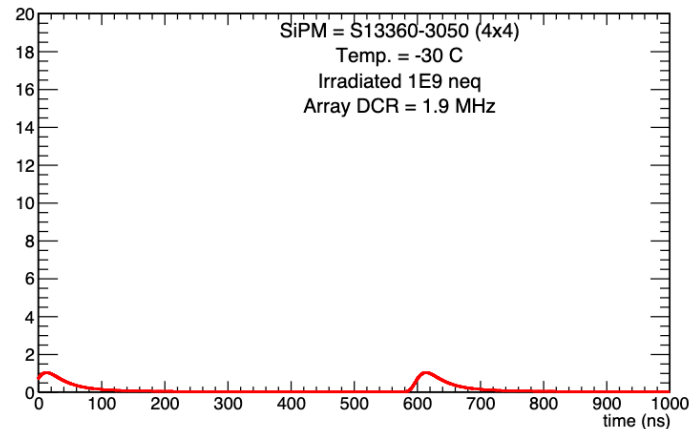
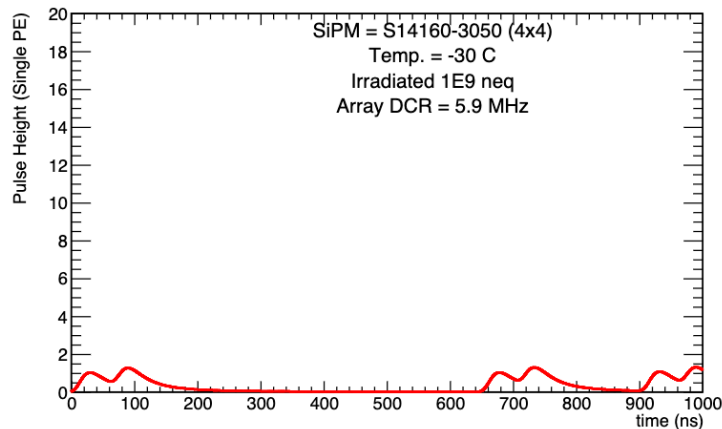
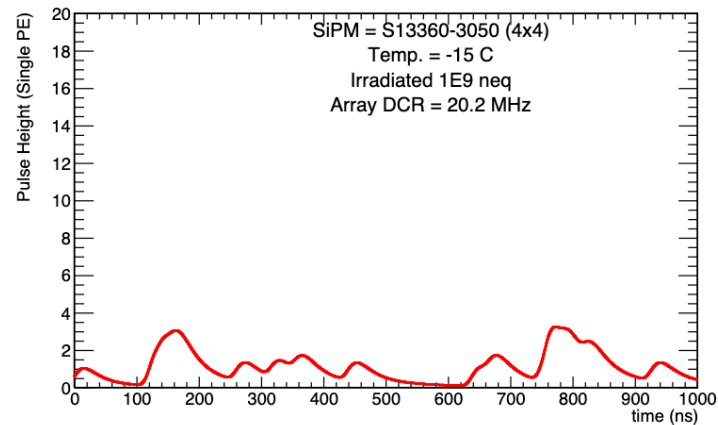
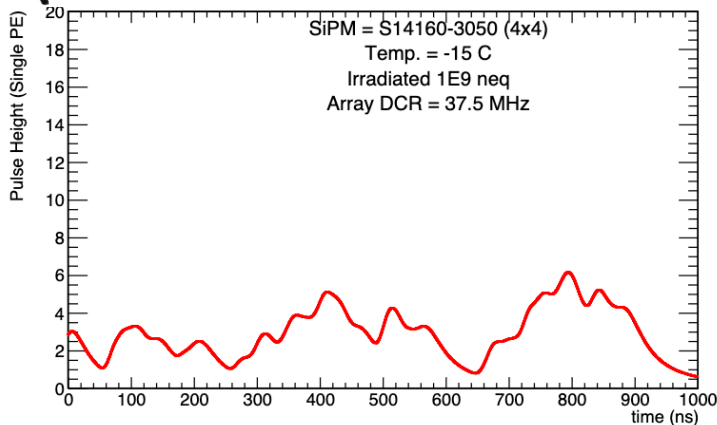
# Sub-Zero



Looks acceptable, as expected since -30C is the nominal temperature for the single-photon counting in the dRICH

DCR is an exponential function of temperature

Integration of sub-zero cooling would likely be very challenging, but in principle possible since the dRICH is already doing it



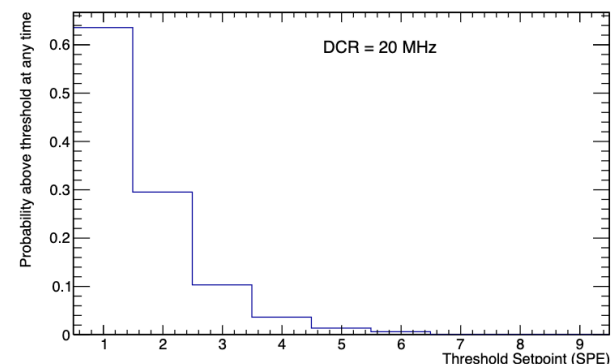
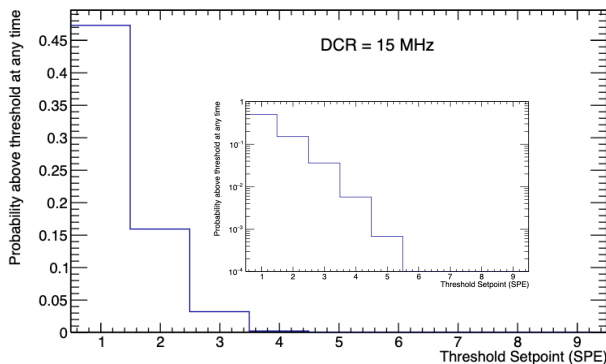
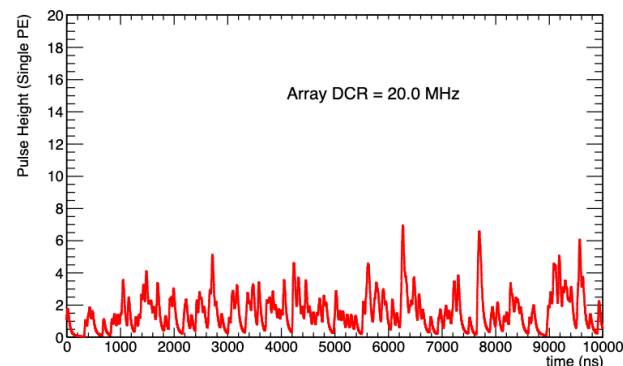
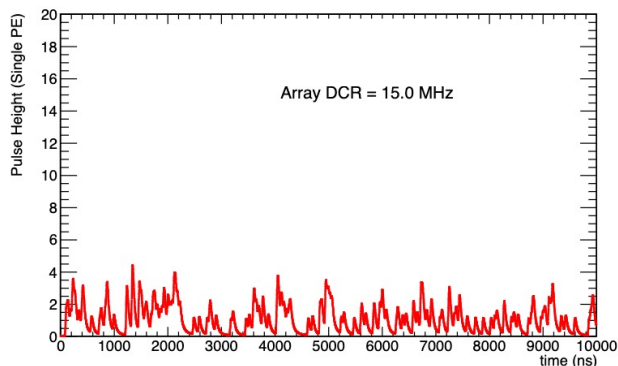


# Threshold

HGCROC will report information whenever the analog input is above a threshold

Given HGCROC pulse shapes, seems that something like ~15 MHz could be a reasonable goal for maximum DCR

- In 100  $\mu$ s, expect analog signal above a 5  $N_{pe}$  threshold ~0.05% of the time



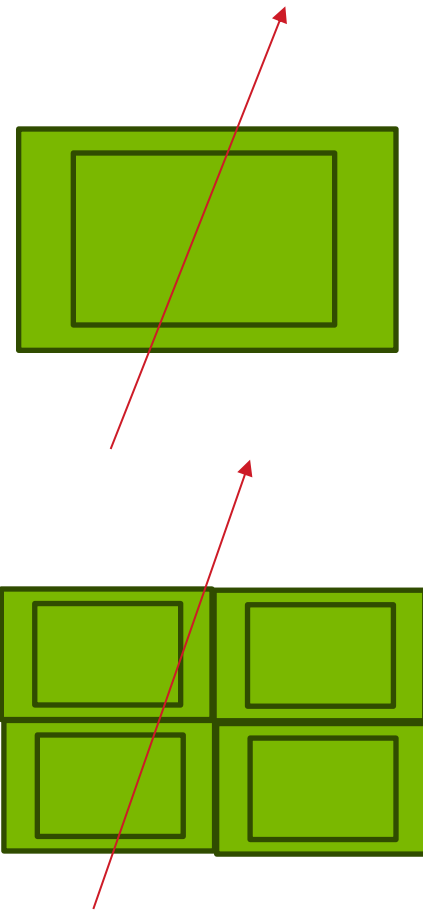
# Summing Channels

How can we reach 15 MHz?

- What matters for keeping the MIP is that there are less than 15 MHz of noise in the area that the MIP photons are hitting

Splitting the 4x4 array into four optically isolated 2x2 arrays could ~half the noise on the MIP

- Set thresholds at  $\sim \frac{1}{4}$  the value of the 4x4 array
- Four light guides covering smaller areas
- Unlike a shower, the MIP energy deposit in the SciFi can be approximated as a line
  - Will typically send all its photons to  $\sim$ two of the four SiPMs!
- Half the surface area of active SiPM means half the noise



# Discussion

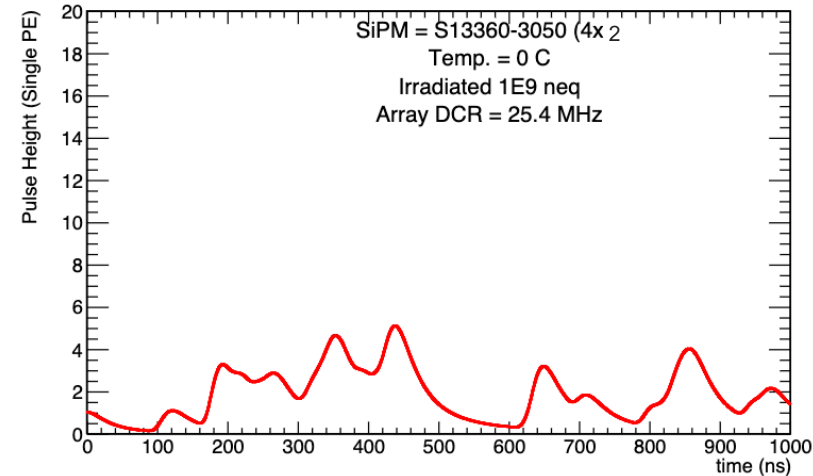
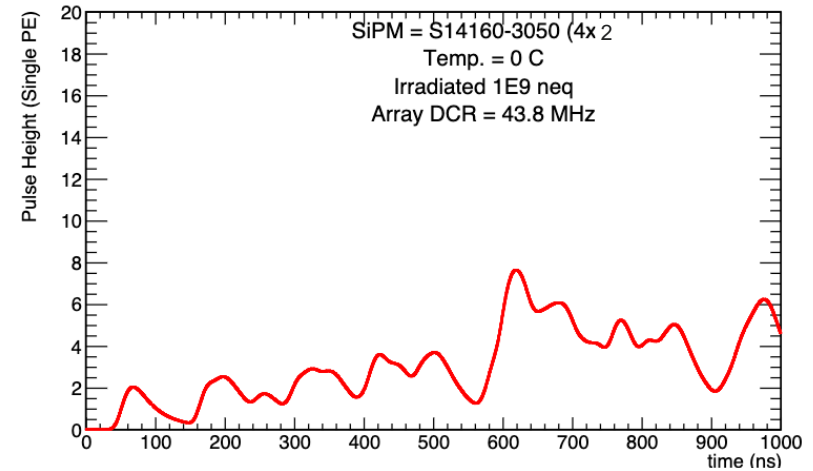
Even with splitting readout channels from 4x4 summed to 2x2 summed, can't reach 15 MHz for irradiated SiPMs even at 0 °C

- 1x1 channels?  $5760 \times 16 = 92\text{k}$  readout channels

Effects of annealing between runs not considered

- Time heals all wounds
- Some studies show 1 month of sitting at room temp. after proton irradiation can decrease dark current by a factor of 2

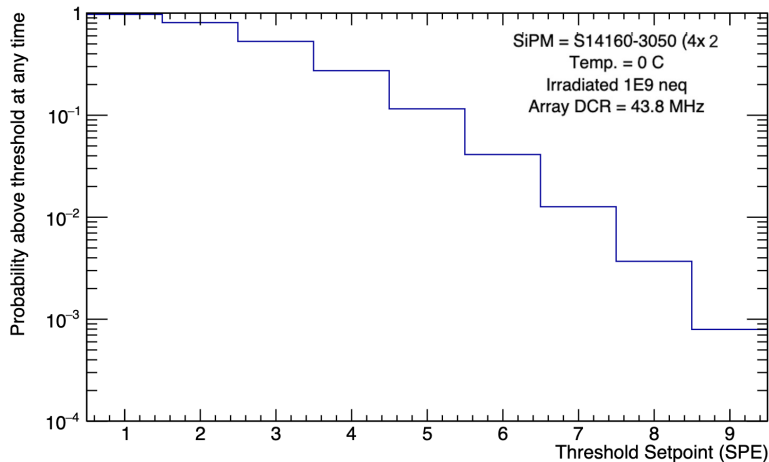
Or do as Craig Woody said and use the MIP to cross-calibrate in the early stages of the experiment, then trust that calibration once the MIP is below threshold



# Fraction of time spent above a threshold of $N_{\text{SPE}}$ (1 ms simulated)

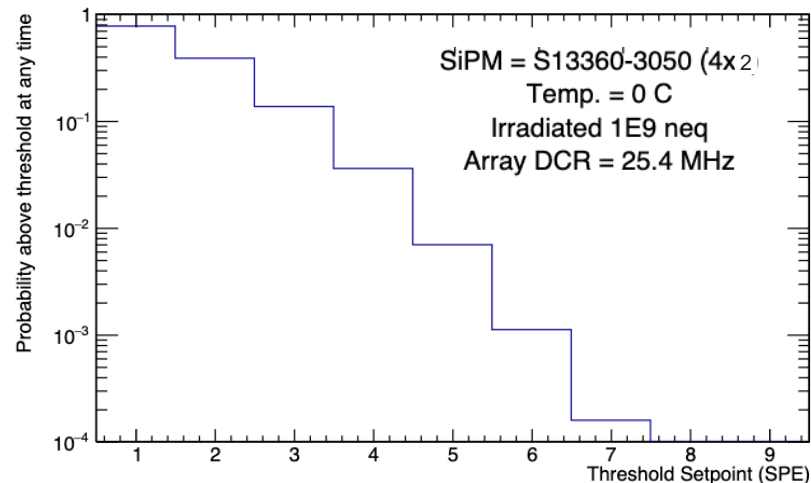
```
percentage above 2 SPE = 0.805159
percentage above 3 SPE = 0.526815
percentage above 4 SPE = 0.272328
percentage above 5 SPE = 0.115044
percentage above 6 SPE = 0.040986
percentage above 7 SPE = 0.012618
percentage above 8 SPE = 0.003681
percentage above 9 SPE = 0.000791
```

```
percentage above 2 SPE = 0.38593
percentage above 3 SPE = 0.136293
percentage above 4 SPE = 0.035822
percentage above 5 SPE = 0.006938
percentage above 6 SPE = 0.001113
percentage above 7 SPE = 0.000158
percentage above 8 SPE = 0
percentage above 9 SPE = 0
```



Per channel, expect 11% of samples to be above a 5 p.e. threshold (4.4 MHz)

Ignoring time binning effects



Per channel, expect 0.6% of samples to be above a 5 p.e. threshold (0.24 MHz)

# HGCROC Bandwidth

2x2 array in 100  $\mu$ s

HGCROC sends 32 bits/hit

Currently each HGCROC has 2 x 1.28 Gbps links

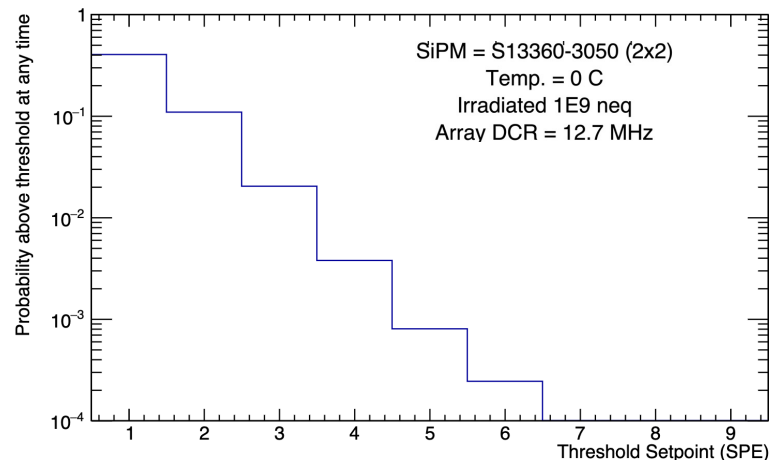
Max hit rate above threshold: 80 MHz for all channels on that HGCROC

Anticipate 60 Channels/HGCROC

Max hit rate of  $\sim 1.3$  MHz/channel can be passed on by HGCROC

- For irradiated S13360 with readout split into 2x2, expect total DCR/channel of 12.7 MHz, want a threshold at  $\sim 2$ -3 SPE
  - Above 2 PE 10% of the time  $\sim 4$  MHz
  - above 3 PE 2% of the time  $\sim 0.8$  MHz
  - Cross-talk has much larger effect on operation at lower thresholds
- Rate/channel situation is slightly worse for split readout due to cross-talk

percentage above 1	SPE = 0.403242
percentage above 2	SPE = 0.109232
percentage above 3	SPE = 0.020353
percentage above 4	SPE = 0.00378
percentage above 5	SPE = 0.000802
percentage above 6	SPE = 0.000244
percentage above 7	SPE = 2.4e-05
percentage above 8	SPE = 0
percentage above 9	SPE = 0



# Conclusion

With measurements from INFN and HGCROC signals, made some first estimates for what waveforms will look like

Irradiated sensors very likely to lose the MIP without intervention

- Additional cooling
- Annealing
- Further splitting of channels

Caveats:

- Dark Current to DCR conversion
- Pulse shapes will depend on how many SiPMs are ganged together
- No room-temperature annealing effect

# BACKUP

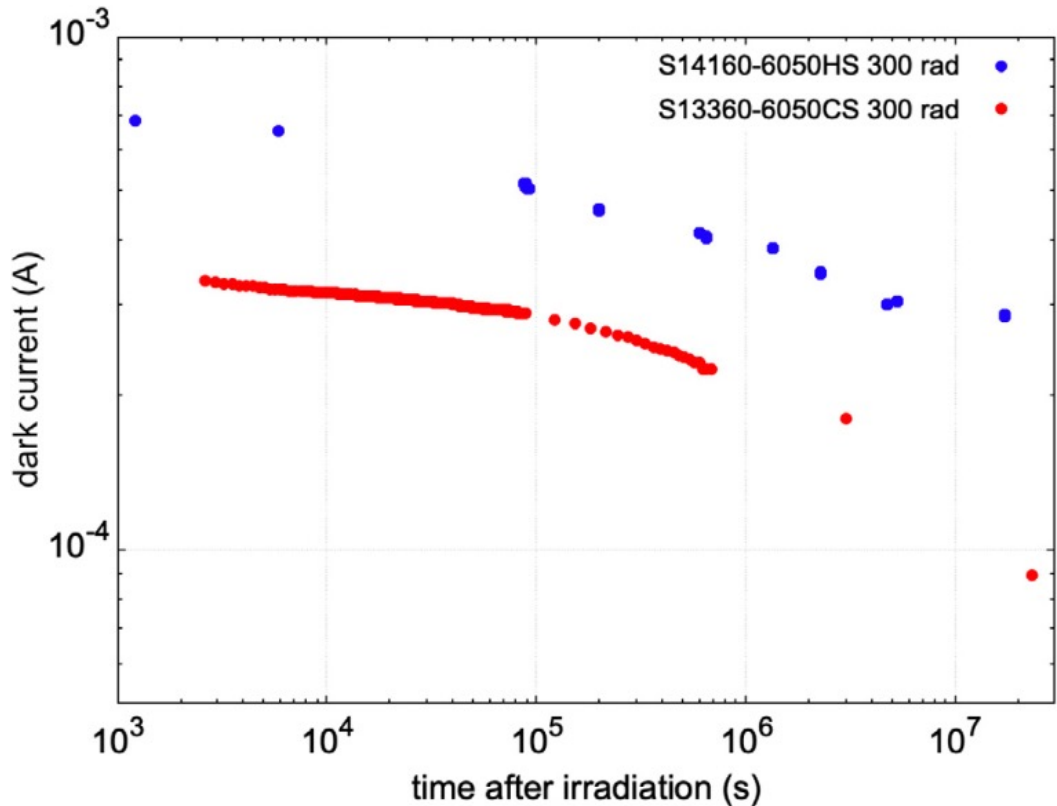
# HGCROC ADC Resolution

Resolution of 0.4 fC ~ 2500 electrons

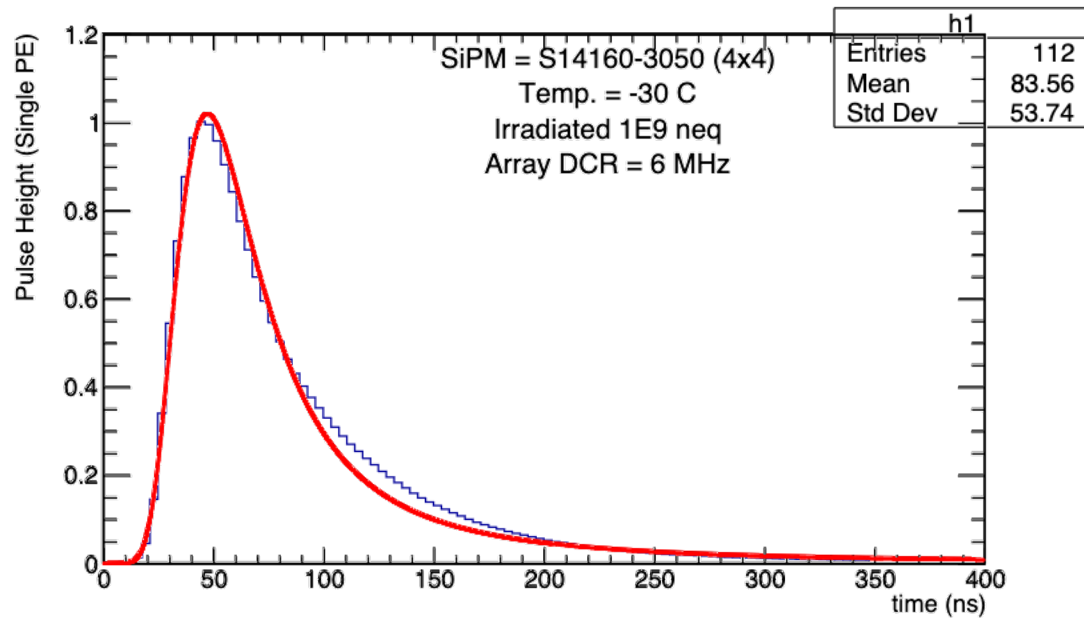
Much smaller than the signal from a single SPAD firing in the SiPM

- S13360 Gain 1.7E6
- S14160 Gain 2.5E6





**Fig. 2.** Dark current in S13 (red) and S14 (blue) as a function of time after 300 rad irradiation at room temperature. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)



# CONSIDERATIONS

Photon detection efficiency

Noise

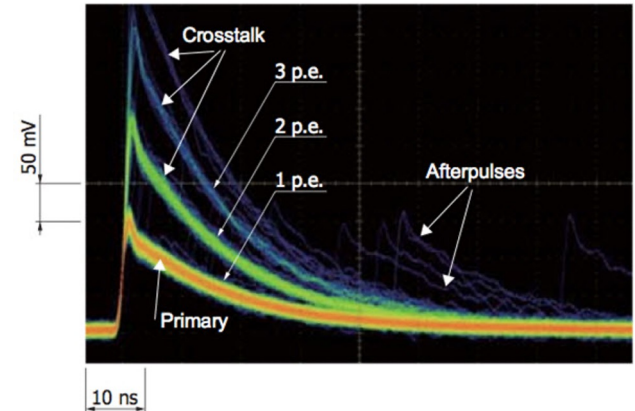
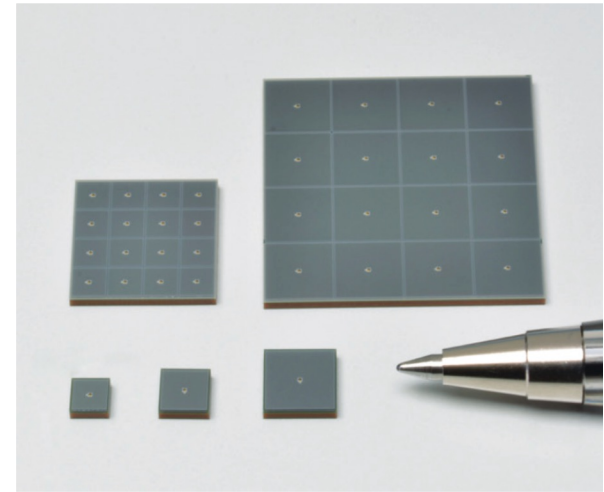
- Low dark count rate necessary to see small signals from e.g. MIPs
- Low crosstalk & afterpulsing preferable

Pulse shape

- Fast rise time for z-position resolution
  - Time-projection Calorimeter (TPC)
- Short tail to reduce signal pileup
- Consistent over time
- Proportional to  $N_{pe}$

Consider the performance in the BIC of Hamamatsu S13660 and S14160 Series

- Biggest challenge is seeing the MIP at midrapidity, use this as a benchmark



# PDE

PDE important to minimize statistical error on energy measurement & maximize efficiency for small signals

50 micron pixel size chosen

- Trade off between geometric efficiency & saturation point

Relevant wavelengths determined by emission spectrum of scintillating fibers

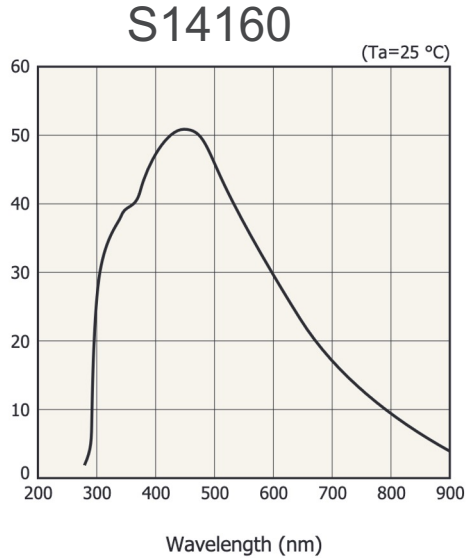
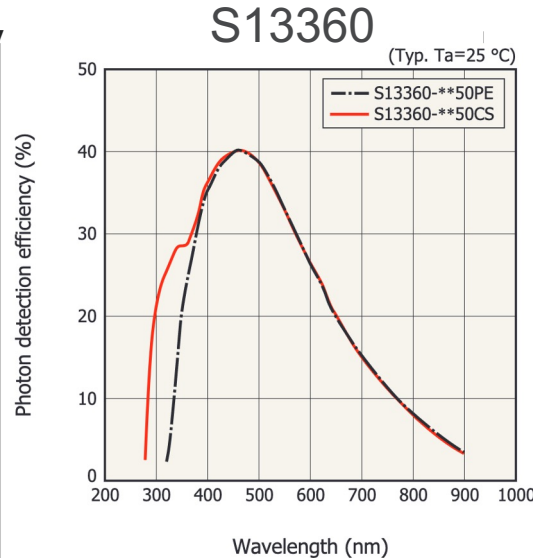
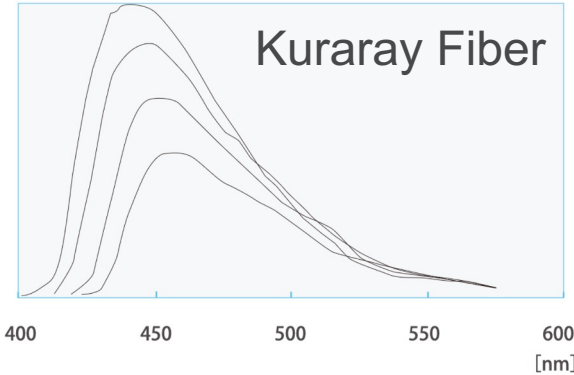
- Both SiPMs peak in PDE near the emission peak around 450 nm

**S14160 peaks at ~50% PDE**

**S13360 peaks at ~40% PDE**

In PDE, S14160 series wins

SCSF-78



# NOISE

Dark count rate (DCR) determines threshold

- MIPs at midrapidity will generate 3-6  $N_{pe}$  on average
  - Would be good to have threshold slightly below MIP

DCR above a few 10s of MHz will endanger the MIP

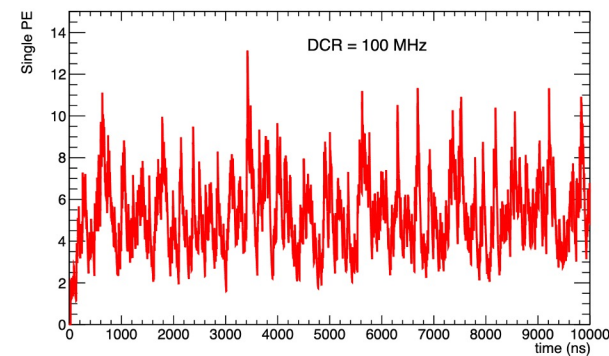
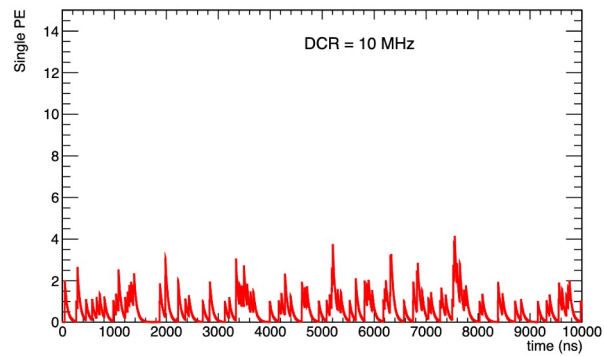
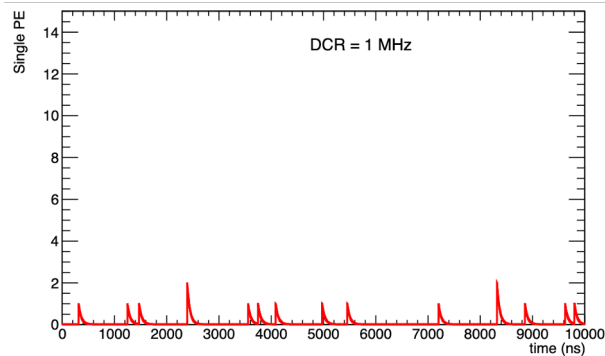
Specification	S13360-3050 (3x3 mm)	S14160-3050 (3x3 mm)
DCR (Typ.)	500 kHz	1 MHz***
Crosstalk (%)	3	7

\*\*\* Estimated, differing values in literature

Signal will gang 1.2 cm x 1.2 cm area (16 3x3 mm or 4 6x6 mm)

- DCR for one BIC channel will be ~16x value in table

Plan to test S14160 SiPMs at ANL & Regina



# PULSE SHAPE

Fast rise time improves position resolution in z-direction

- If not limited by other factors like readout or scintillation decay

Fast fall time reduces pileup of signals (and dark counts)

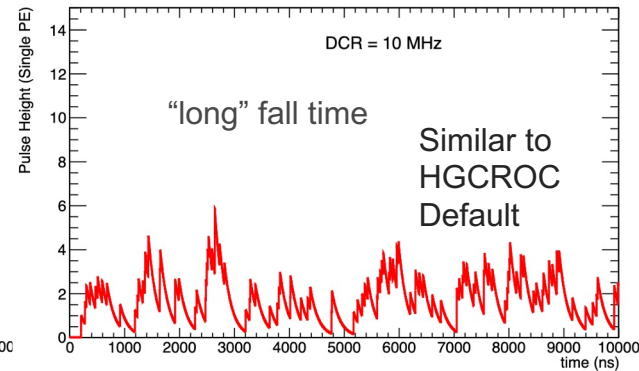
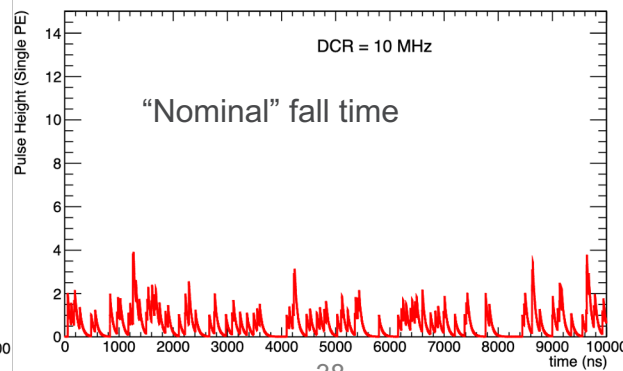
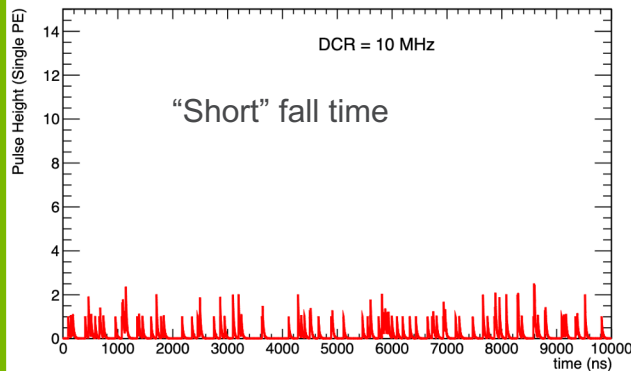
- Shape will depend heavily on the readout circuit

*Appears that S14160 has fast rise time, slower fall time than S13360*

- Challenging to compare between papers due to differences in readout
- Will soon compare the two in an identical setup at ANL

Should converge on a reasonable target for these parameters based on physics

- Keeping the MIP, low energy  $\gamma$



**TABLE 1: Barrel Imaging Calorimeter SiPM Specs**

Parameter	Specification	Notes
Active Area	3 mm x 3 mm (4 x 4 array)	Preassembled array covering 1.2cm x 1.2cm
Pixel Size	50 μm	
Package Type	Surface Mount	
Peak Sensitivity	450 nm	
PDE	~ 50%	
Gain	> ~ 2 x 10 <sup>6</sup>	
DCR	Typ.: ~ 500kHz / SiPM Max: < 1.5 MHz / SiPM	DCR applies to each SiPM in the 4 x 4 array
Temperature coefficient of Vop	< 40mV/C	
Direct crosstalk probability	< ~ 7%	
Terminal capacity	~ 500pF / SiPM	Applies to each SiPM in the 4 x 4 array
Packing granularity		
Vop variation within a tray	< 200 mV	
Recharge Time	< 100 ns	
Fill Factor	> 70%	
Protective Layer	Silicone (n ~ 1.5-1.6)	

S13360



~



S14160



# OPEN/DISCUSSION QUESTIONS

What DCR can we really tolerate?

Depends on signal shape, shorter fall time better

With HGCROC-length signals, 10 MHz too large if threshold is  $3 N_{pe}$

Do we want to actively temperature control the SiPMs?

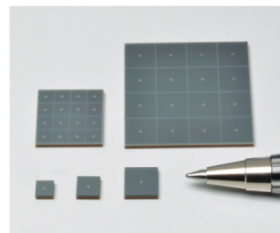
- Maintain a constant DCR by decreasing temp as rad damage accumulates
- More effective with S13360 series than S14160

If we go with the S13360 series, how can we compensate the loss in PDE?

Does the lower operating voltage of the S14160 series benefit us?



## MPPC® (Multi-Pixel Photon Counter)



S14160/S14161 series

### Low breakdown voltage type MPPC for scintillation detector

The S14160/S14161 series achieve higher PDE (photon detection efficiency) and lower operation voltage than other MPPC to adapt for PET and radiation monitor application. They achieve small dead space in a photosensitive area with HWB (hole wire bonding) technology (Patent pending). And the gap from the photosensitive area edge to the package edge is only 0.2 mm. This package realizes the four-side buttable arrangement.

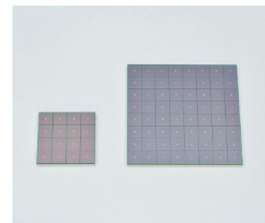
#### Features

- Higher PDE (50% at  $\lambda_p$ ,  $V_{op}=V_{BR} + 2.7 V$ )
- Lower voltage ( $V_{BR}=38 V$  typ.) operation
- Small dead space in photosensitive area
- Low afterpulses and crosstalk
- High gain:  $10^6$  order
- Excellent time resolution
- Immune to effects of magnetic fields

#### Applications

- PET (positron emission tomography)
- Radiation monitor

## MPPC® (Multi-Pixel Photon Counter) arrays



S13361-3050 series

### MPPC arrays in a chip size package miniaturized through the adoption of TSV structure

The S13361-3050 series is a MPPC array for precision measurement miniaturized by the use of TSV (through-silicon via) and CSP (chip size package) technologies. The adoption of a TSV structure made it possible to eliminate wiring on the photosensitive area side, resulting in a compact structure with little dead space compared with previous products. The four-side buttable structure allows multiple devices to be arranged side by side to fabricate large-area devices.

They are suitable for applications, such as medical, non-destructive inspection, environmental analysis, and high energy physics experiment, that require photon counting measurement.

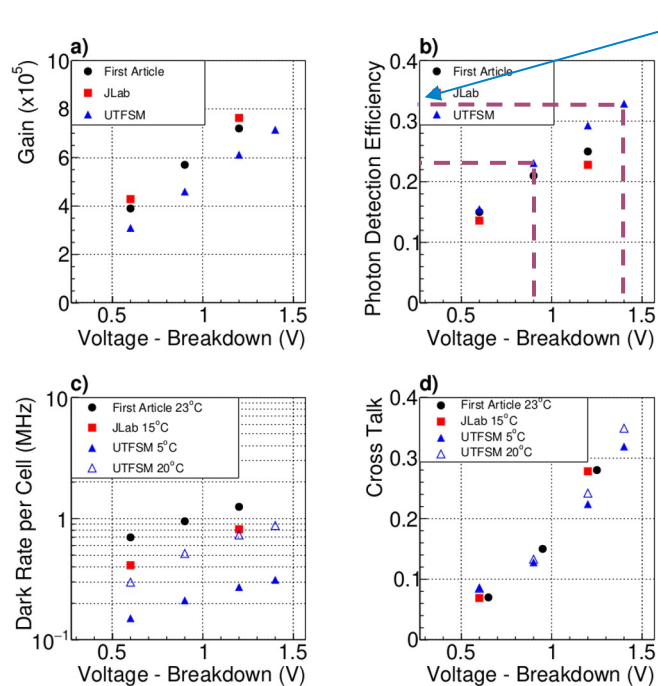
#### Features

- Low crosstalk
- Low afterpulses
- Outstanding photon counting capability (outstanding photon detection efficiency versus numbers of incident photons)
- Compact chip size package with little dead space
- Low voltage ( $V_{BR}=53 V$  typ.) operation
- High gain:  $10^5$  to  $10^6$

#### Applications

- Astro physical application
- High energy physics experiment
- Nuclear medicine
- PET
- Environmental analysis

# PDE: GlueX SiPM Parameters



PDE ~33%

The Hamamatsu specification sheets provide the recommended operating voltage for a nominal gain of  $7.5 \times 10^5$ , although our measurements indicate lower gains (Fig. 4a). We determined that this operational voltage on average corresponds to 0.9 V above breakdown; to obtain our setting at an overvoltage of 1.4 V, we added 0.5 V and then adjusted for temperature.

Hamamatsu Multi-Pixel  
Photon Counter (MPPC)  
S12045(X):  
16 x 3600 pixels (50  $\mu\text{m}$ )

Construction and Performance  
of the Barrel Electromagnetic Calorimeter  
for the GlueX Experiment

T.D. Beattie<sup>a</sup>, A.M. Foda<sup>a</sup>, C.L. Henshel<sup>a</sup>, S. Katsaganis<sup>a</sup>, S.T. Krueger<sup>a</sup>,  
G.J. Luke<sup>a</sup>, Z. Papandreou<sup>a\*</sup>, E.L. Plummer<sup>a</sup>, J.A. Serrano<sup>a</sup>,  
A.M. Sorenson<sup>a</sup>, F. Barboza<sup>b</sup>, E. Chudobov<sup>c</sup>, M.M. Dillou<sup>d</sup>, D. Lawrence<sup>b</sup>,  
Y. Qiang<sup>e</sup>, N. Sandwell<sup>b</sup>, E.S. Sattler<sup>b</sup>, C. Stauden<sup>b</sup>, J.R. Stevens<sup>b</sup>,  
S. Taylor<sup>a</sup>, T. Whited<sup>b</sup>, B. Zihlmann<sup>b</sup>, W. Levine<sup>a</sup>, W. McKinley<sup>a</sup>,  
C.A. Meyer<sup>a</sup>, M.J. Stubb<sup>a</sup>, E.G. Anagnostis<sup>a</sup>, C. Kourkomezis<sup>a</sup>,  
G. Vasilakidis<sup>a</sup>, G. Voglmeier<sup>a</sup>, W.K. Brooker<sup>a</sup>, H. Hakobyan<sup>a</sup>, S. Kolshrub<sup>a</sup>,  
R. Rojas<sup>a</sup>, C. Romero<sup>a</sup>, O. Soto<sup>a</sup>, A. Toró<sup>a</sup>, I. Vega<sup>a</sup>, M.R. Shepherd<sup>d</sup>

<sup>a</sup>Department of Physics, University of Regina, Regina, Saskatchewan, Canada S4S 0A2  
<sup>b</sup>Jefferson Laboratory, Newport News, Virginia 23060, USA  
<sup>c</sup>Cornell University, Ithaca, Pennsylvania 14853, USA  
<sup>d</sup>National and Kapodistrian University of Athens, 15701 Athens, Greece  
<sup>e</sup>University of Tennessee, Knoxville, Tennessee 37996, USA  
<sup>f</sup>Indiana University, Bloomington, Indiana 47405, USA

## Abstract

The barrel calorimeter is part of the new spectrometer installed in Hall D at Jefferson Lab for the GlueX experiment. The calorimeter was installed in 2013, commissioned in 2014 and has been operating routinely since early 2015. The detector configuration, associated Monte Carlo simulations, calibration and operational performance are described herein. The calorimeter records the time and energy deposited by charged and neutral particles created by a multi-GeV photon beam. It is constructed as a lead and scintillating-fiber calorimeter and read out with 3840 megaelectron-silicon photomultiplier arrays. Particles impinge on the detector over a wide range of angles, from normal incidence at 90 degrees down to 11.5 degrees, which defines a geometry that is fairly unique among calorimeters. The response of the calorimeter has been measured during a running experiment and performs as expected for electromagnetic showers below 2.5 GeV. We characterize the performance of the BCal using the energy resolution integrated over typical angular distributions for  $e^+e^-$  and  $\eta$  production of  $\sigma_E/E=2\%/\sqrt{E(\text{GeV})}$ , 3.6% and a timing resolution of  $\sigma_t=10$  ps at 1 GeV.

arXiv:1801.03088v2 [physics.ins-det] 20 Apr 2018

Figure 4: Measurements of the first-article samples (black circles) [20, 23], production samples at JLab (red squares) and production samples at UTFSM (triangles) [21, 22] of four basic SiPM parameters as a function of the voltage over breakdown. a) gain, b) photon detection efficiency, c) dark rate per tile (the dark rate for the array is 16 times higher) and d) cross talk determined from deviations of the single-pixel distributions from a pure Poisson function. As long as the voltage over breakdown is kept constant, the dark rate is the only parameter that has a significant temperature dependence. The nominal operating voltage for the GlueX experiment is 1.4 V above breakdown. (Color online)

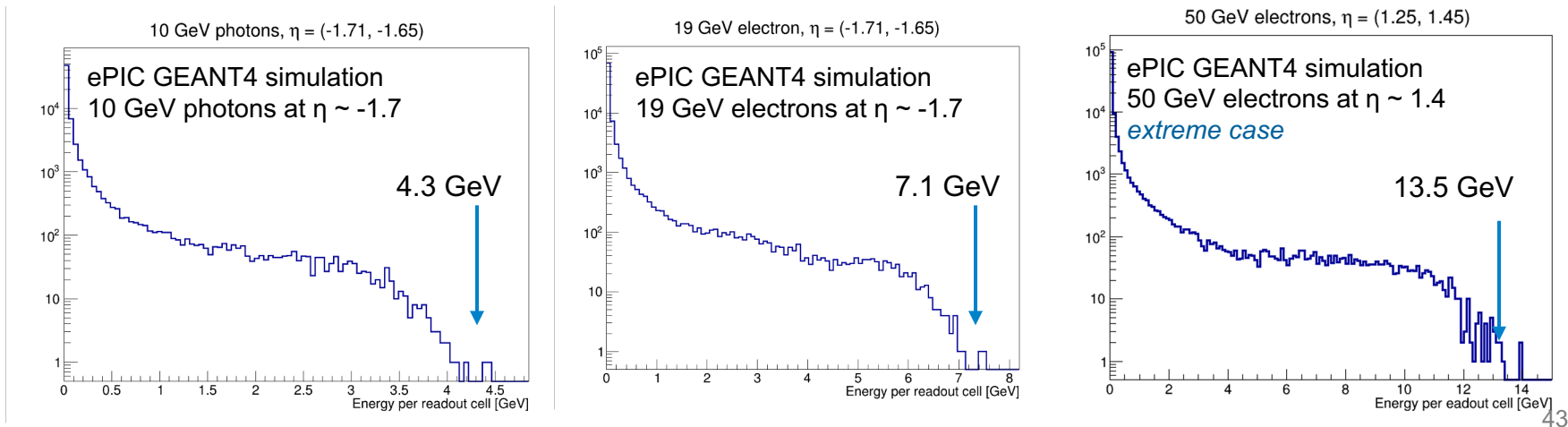
# Pixel Size and Number of Pixels

Defined by **photoelectron statistics** and **energy range** to be measured

## Energy measurement ranges in BECal:

- Shall provide photon measurements up to 10 GeV (F-DET-ECAL-BAR.2) and down to 100 MeV (F-DET-ECAL.9)
- Shall provide electron ID up to 50 GeV and down to 1 GeV and below (F-DET-ECAL-BAR.1)
  - Electron energy measurement needed for  $e/\pi$  separation only (straightforward at high energies)
- Reasonable performance for MIPs needed for calibration and for muon ID

Largest energy deposit occurs for particles at large  $\eta$  (steep angle) where the pathlength in a cell is maximal and the attenuation is minimal.



# Photoelectron statistics

From our 2023 Hall D tests using GlueX SiPMs and double-clad Kuraray fibers: **1000 phe/GeV** per side for showers at the center of the Baby BCAL prototype

- Corrected for attenuation: **1077 phe/GeV\*** per side

We can scale these results for the **ePIC Barrel ECal\***:

- x 1.5 factor improvement in **SiPM photon detection efficiency**
- x 1.16 factor to account for **better optical coupling**
- x 0.69 reduction accounting for **single-clad Kuraray fibers**

This gives **~ 1239 phe/GeV** per side (fully corrected for attenuation)

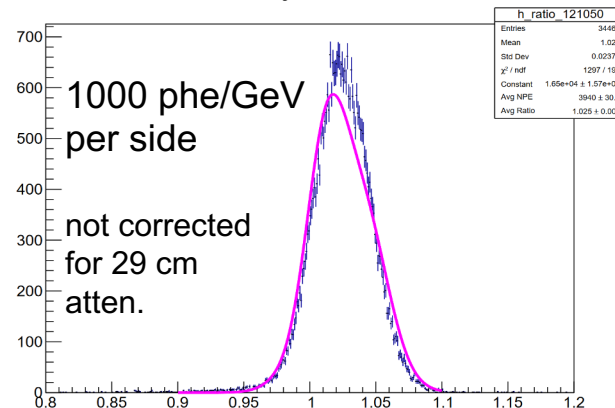
- **10 GeV  $\gamma$  at  $\eta \sim -1.7$ : 5560 phe  $\rightarrow$  **9.8 % max SiPM occupancy****
- **19 GeV  $e^-$  at  $\eta \sim -1.7$ : 9181 phe  $\rightarrow$  **16.1 % max SiPM occupancy****
- **50 GeV  $e^-$  at  $\eta \sim 1.4$  (most extreme case): 17456 phe  $\rightarrow$  **30.1% max SiPM occupancy****

Well below the region where large nonlinearities in the SiPM response are expected in almost all cases.

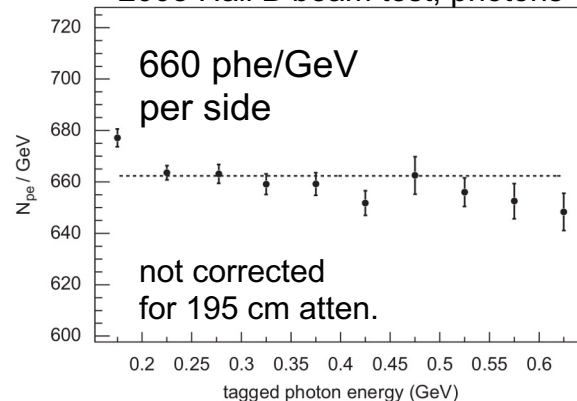
Small non-linear effects possible for some ultra-high energy electrons, which is acceptable ( $e-\pi$  separation straightforward).

\* See backup slide for the attenuation length measurement and extraction of those factors

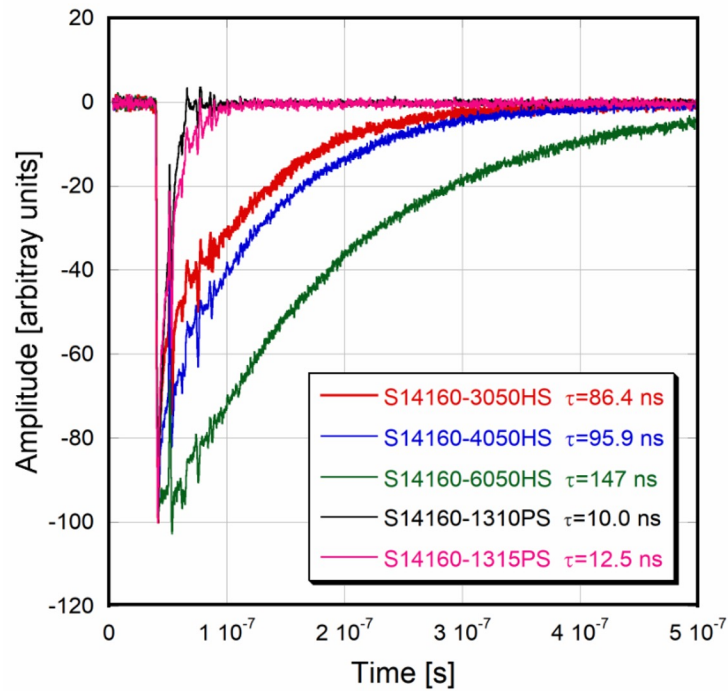
2023 Hall D, Baby BCal, 3.9 GeV  $e^+$



2008 Hall B beam test, photons



**Fig. 16.** The number of photoelectrons per GeV per end of the BCAL module is shown as a function of energy. A one parameter fit is plotted (dashed line). For more details see the text.



# NOISE

Dark count rate (DCR) determines threshold

- MIPs at midrapidity will generate 3-6  $N_{pe}$  on average
  - Would be good to have threshold slightly below MIP

DCR above a few 10s of MHz will endanger the MIP

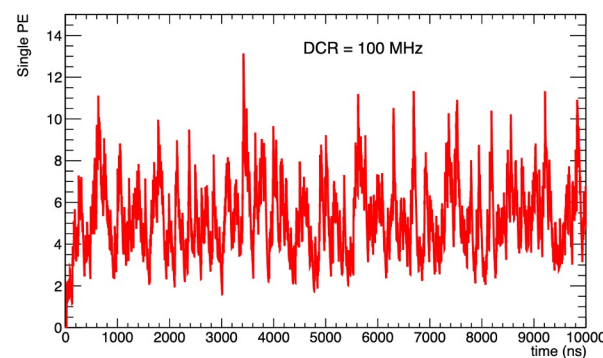
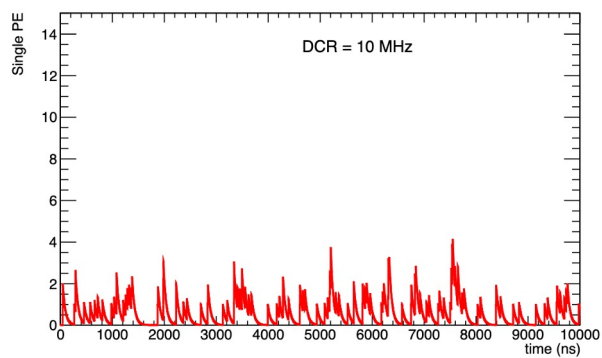
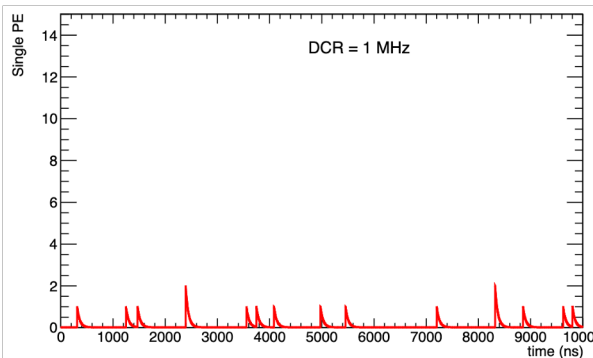
Specification	S13360-3050 (3x3 mm)	S14160-3050 (3x3 mm)
DCR (Typ.)	500 kHz	1 MHz <sup>***</sup>
Crosstalk (%)	3	7

\*\*\* Differing values in literature

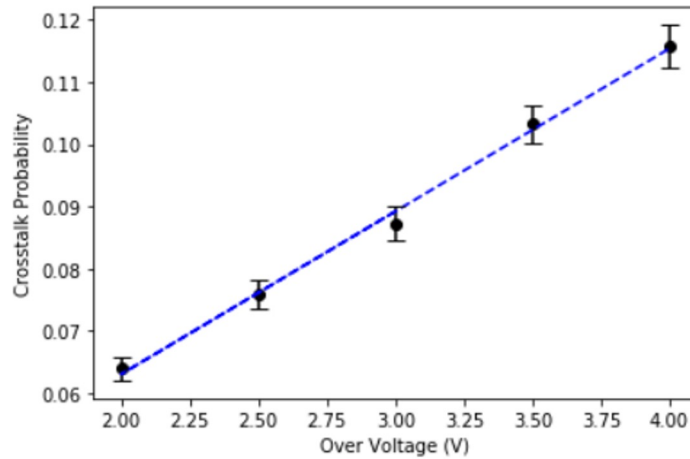
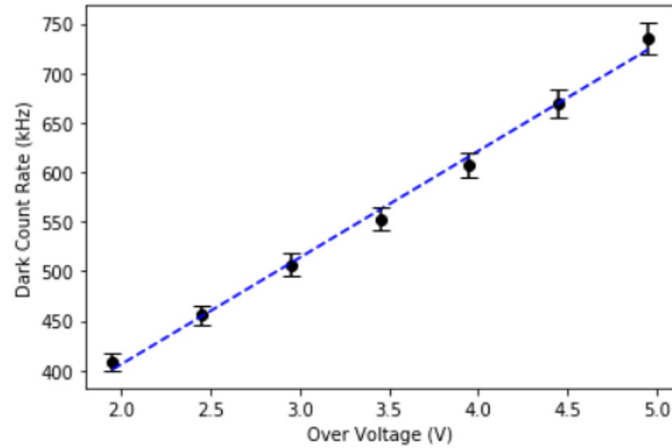
Signal will gang 1.2 cm x 1.2 cm area (16 3x3 mm or 4 6x6 mm)

- DCR for one BIC channel will be ~16x value in table

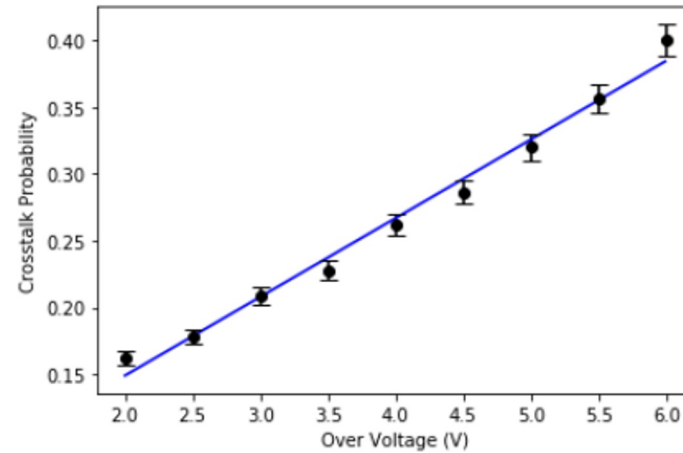
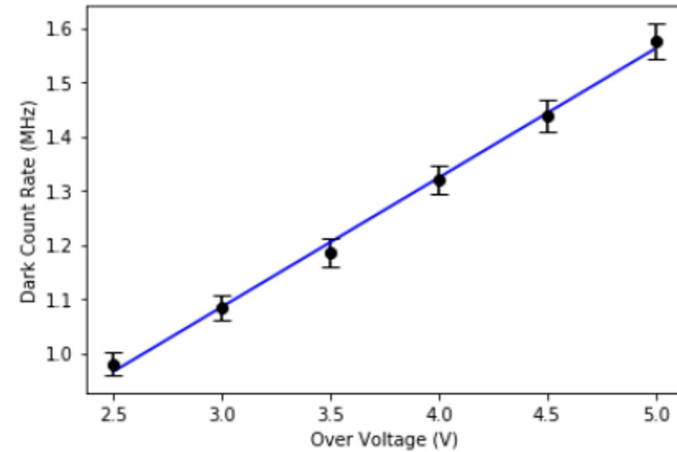
Plan to test S14160 SiPMs at ANL & Regina



## S13360-3050 (3x3 mm)



## S14160-3050 (3x3 mm)



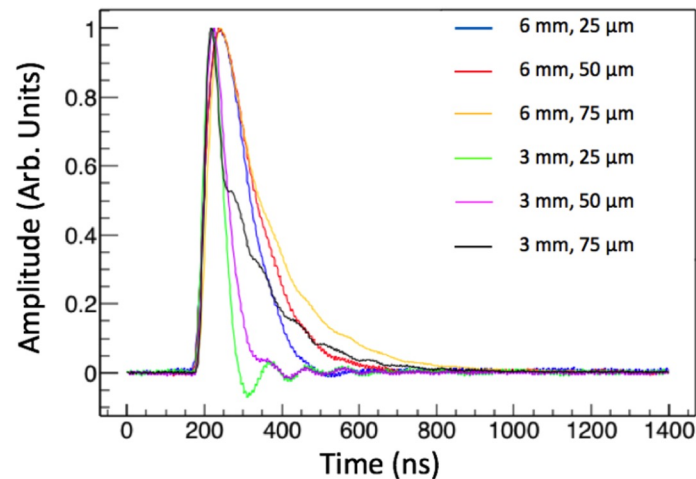
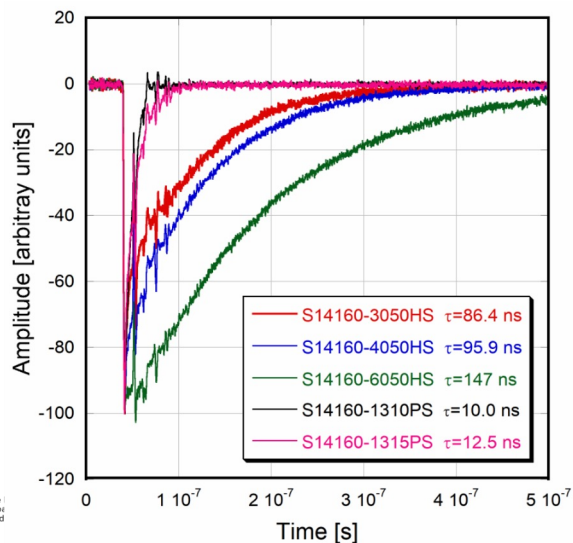
# PULSE SHAPE

Pulse shape strongly defined by how signals are handled

S14160 has faster rise time, slower fall time

Larger SiPMs (6x6 mm) have ~2x longer fall times due to capacitance

- Can we mitigate this in our ganging scheme?





# NOISE

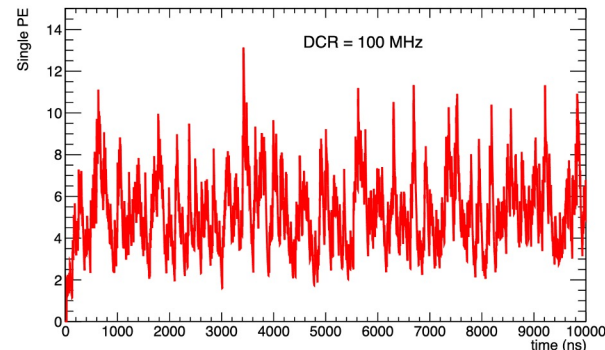
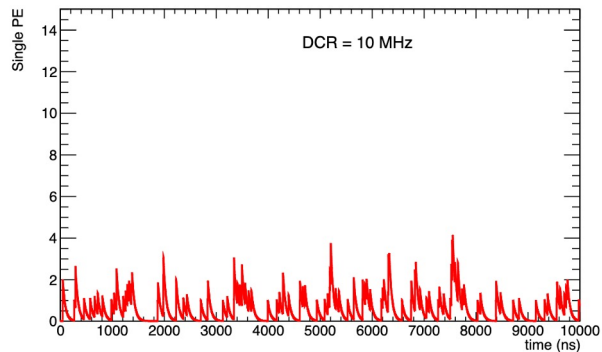
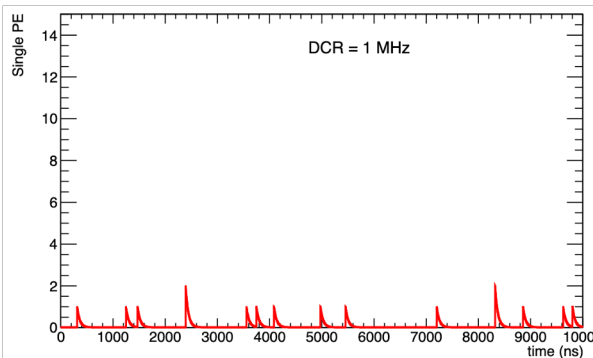
Dark count rate (DCR) determines threshold

- MIPs at midrapidity will generate 3-6  $N_{pe}$  on average
  - Would be good to have threshold slightly below MIP

DCR above a few 10s of MHz at the readout will swamp the MIP

Literature seems divided on noise characteristics of S14160 series

- Plan to test S14160 SiPMs at ANL & Regina



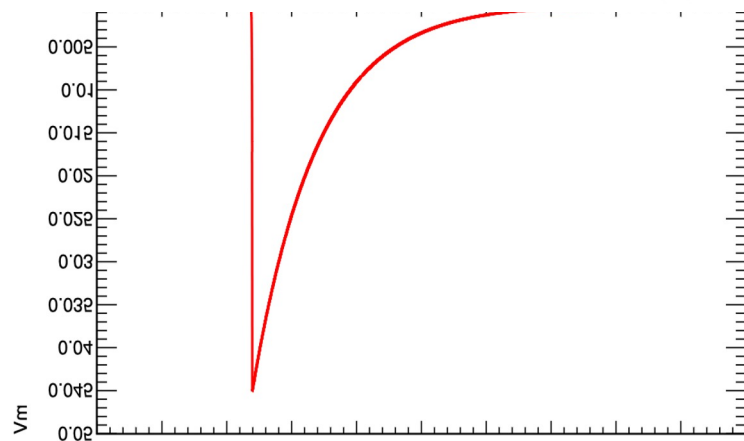
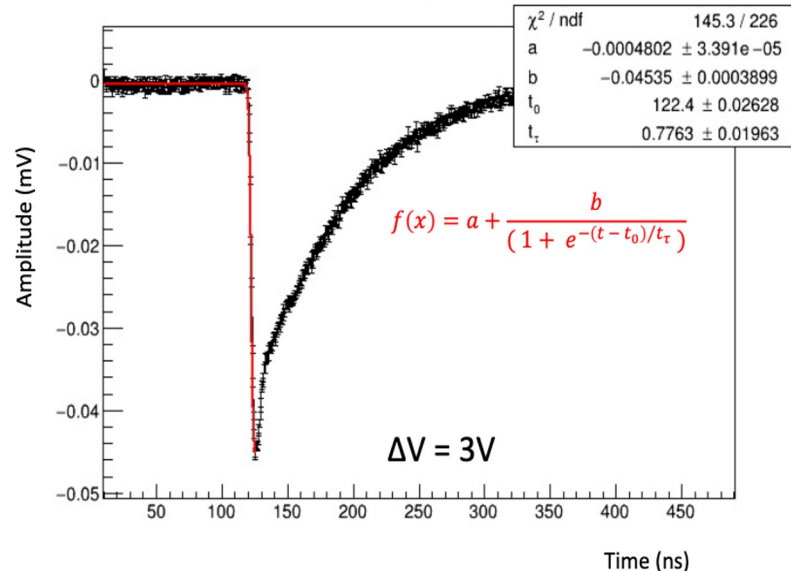
# WAVEFORMS

Parameterize S14161 waveform based on presentation from AMS-100 ([here](#))

Exponential rise and exponential fall  
– Different time constants

Pulse height around 0.045 mV  
– Take this as 1 Npe

SiPM1 ( $V_{bd} = 156.9V$ )

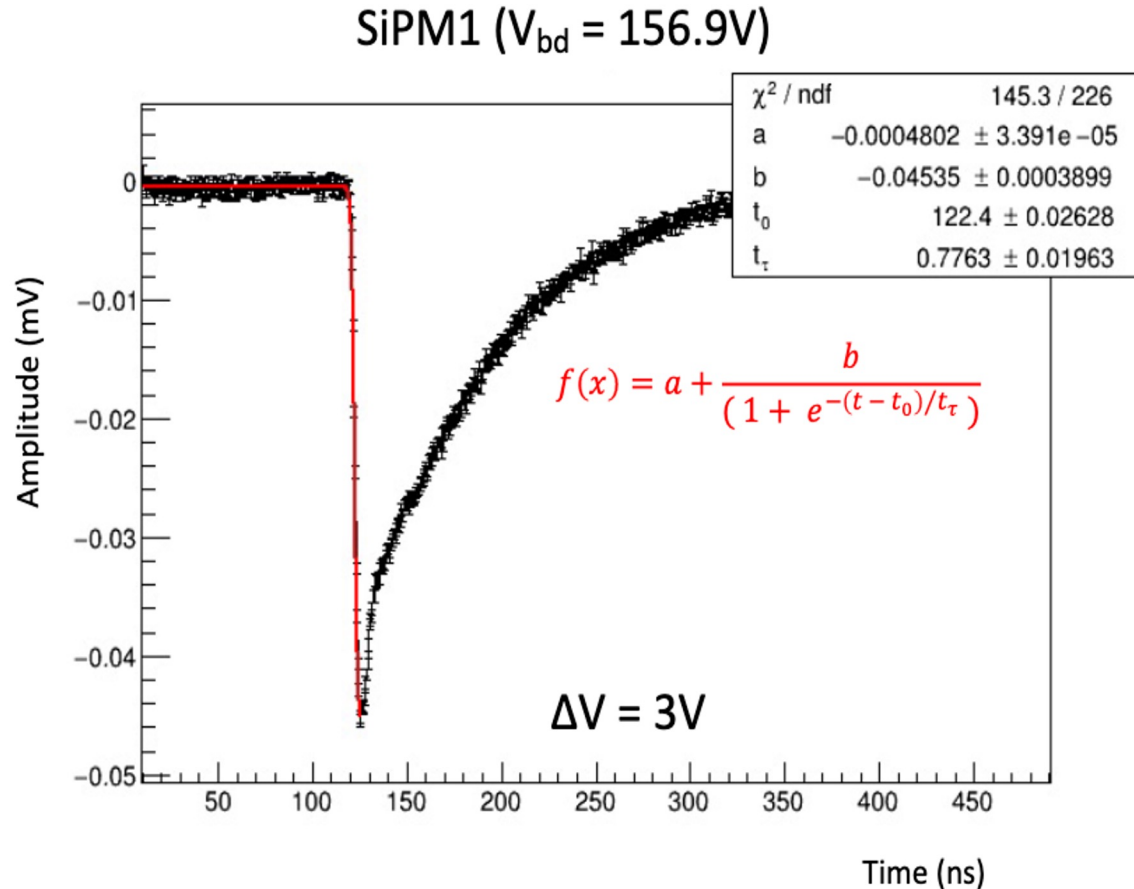


# WAVEFORMS

Agreement not good but also not so terrible

Tail a bit wider in the data

Good enough for now

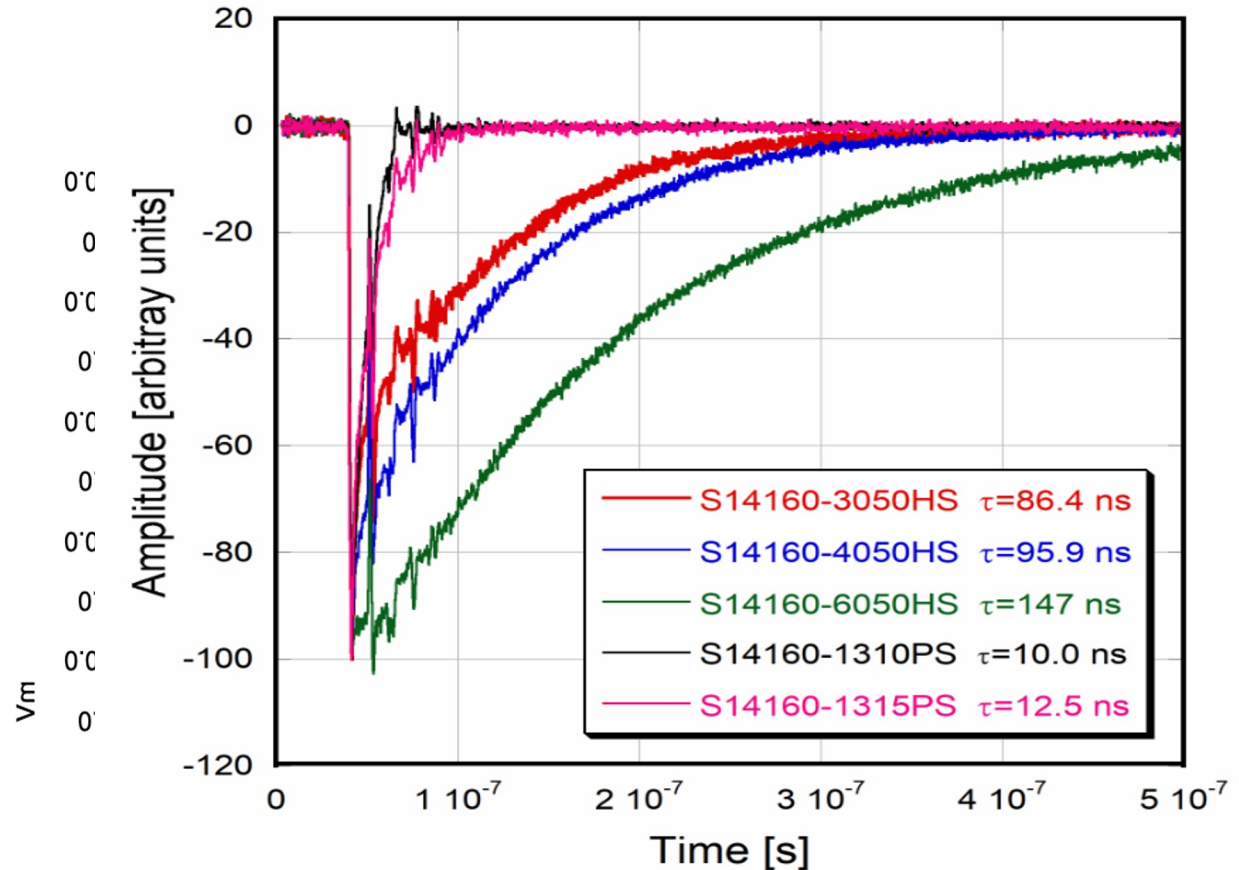


# WAVEFORMS

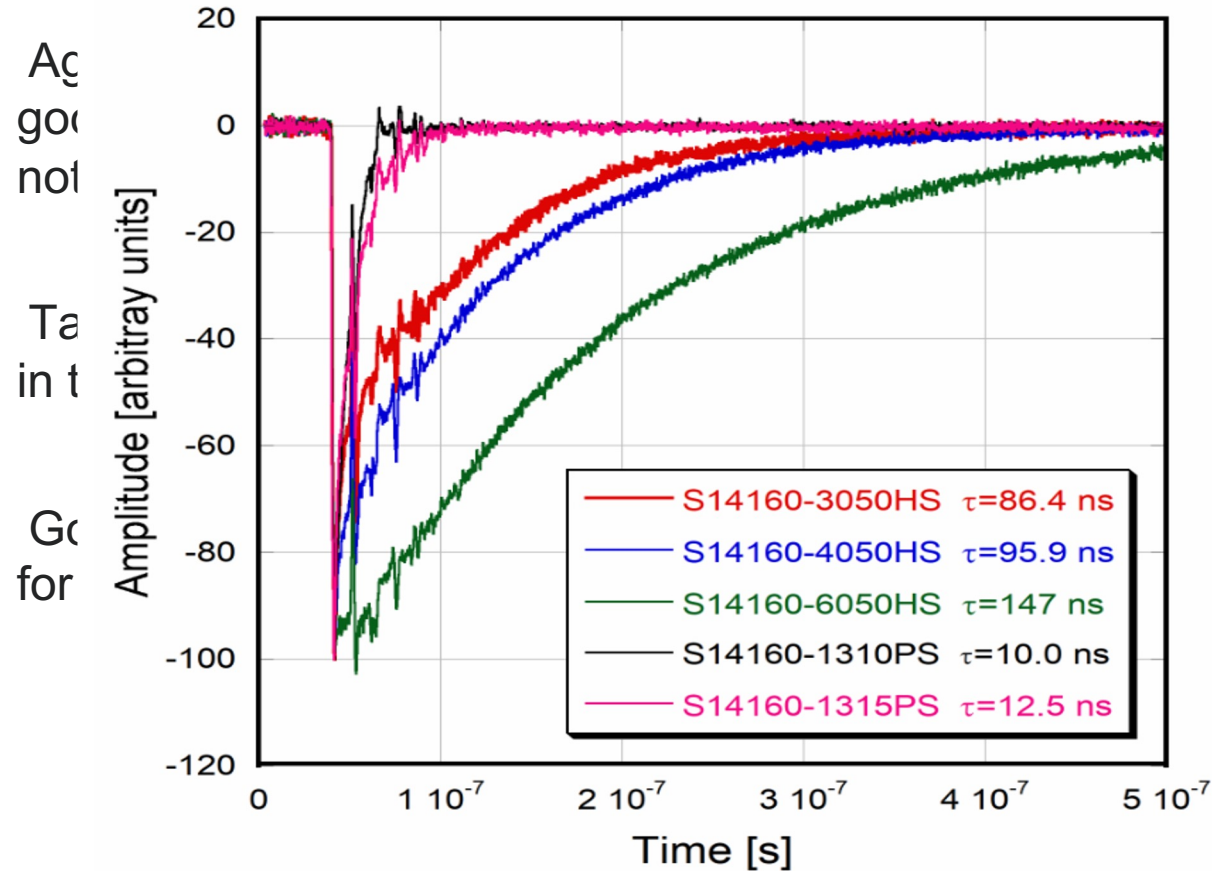
Agreement not good but also not so terrible

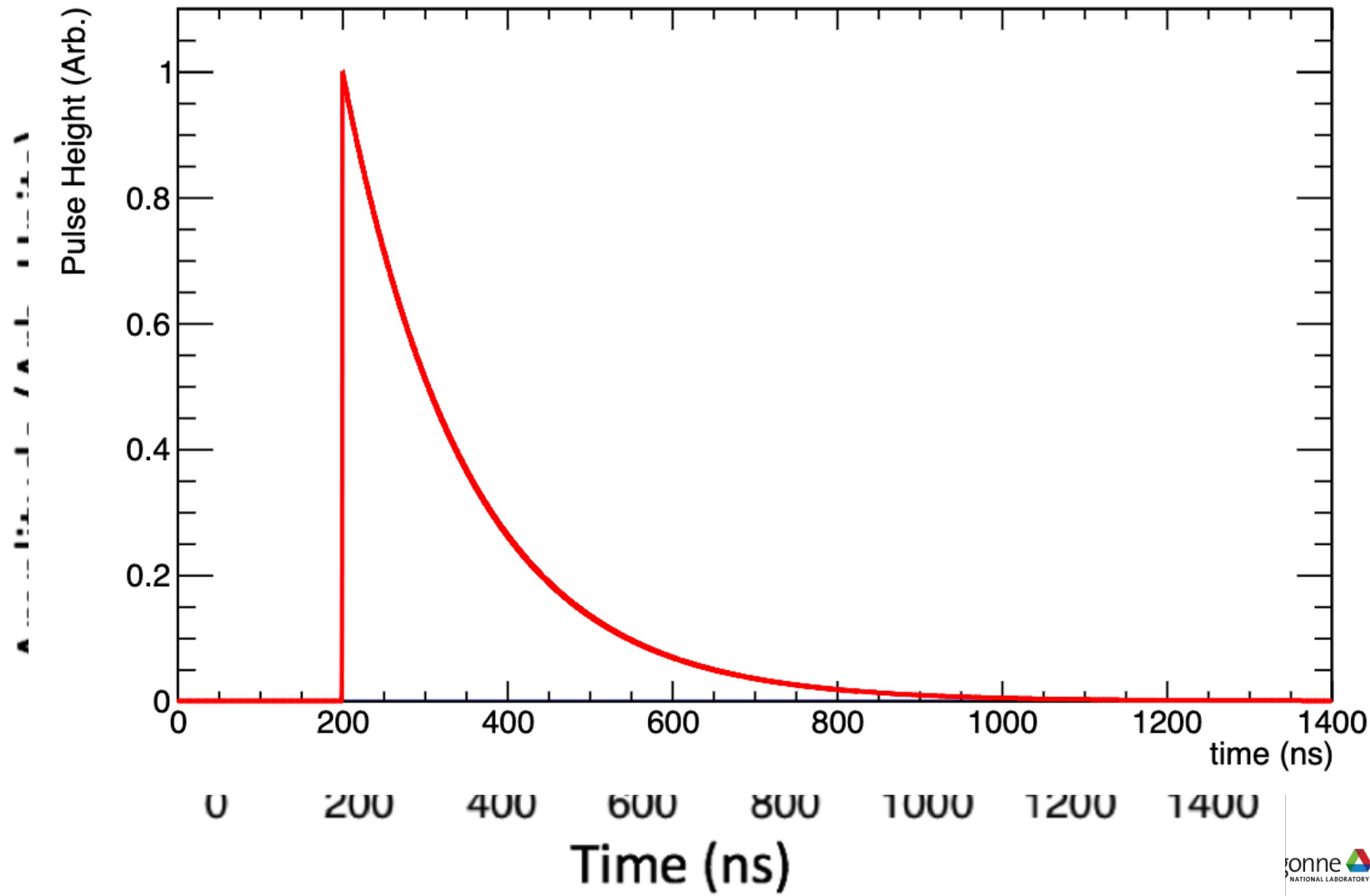
Tail a bit wider in the data

Good enough for now



# WAVEFORMS





# WAVEFORMS

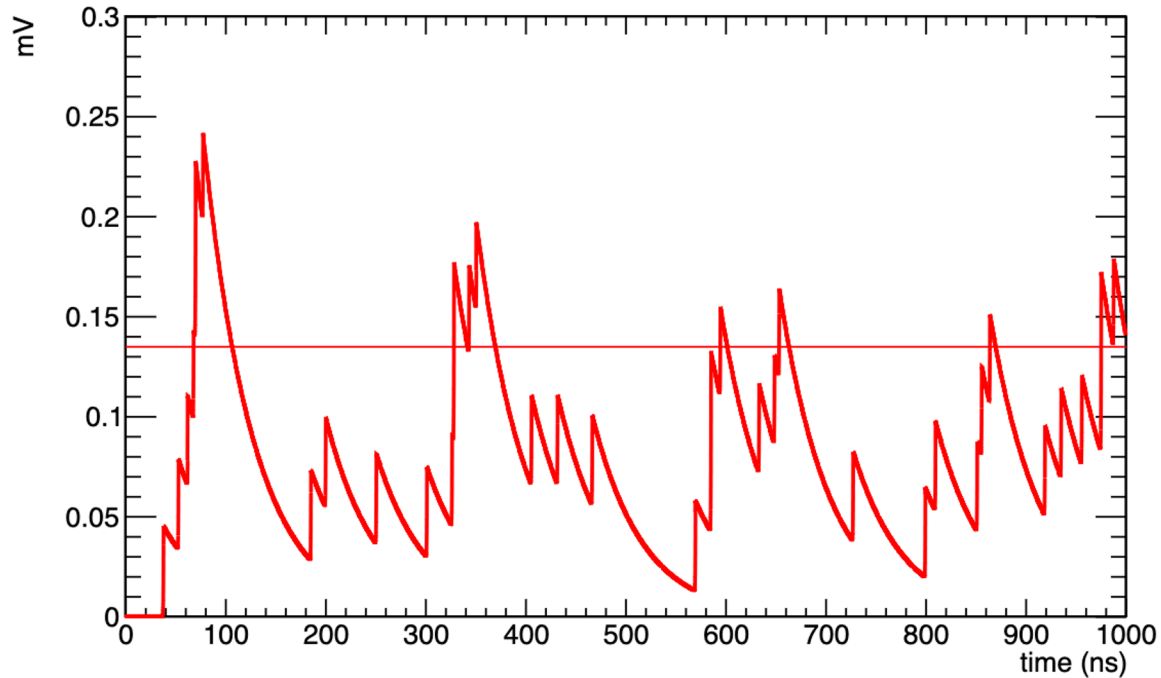
Monte Carlo throwing signals  
with expected rate

- 32 MHz
- 1 microsecond

Crosstalk probability of 7%  
included (should it be, or is it  
included in the number from  
Hamamatsu?)

- Up to two crosstalk hits
  - 3 and greater is a less  
than 1% effect

Line drawn at 3 \* single  
photoelectron peak



# HITS IN HGCROC WINDOW

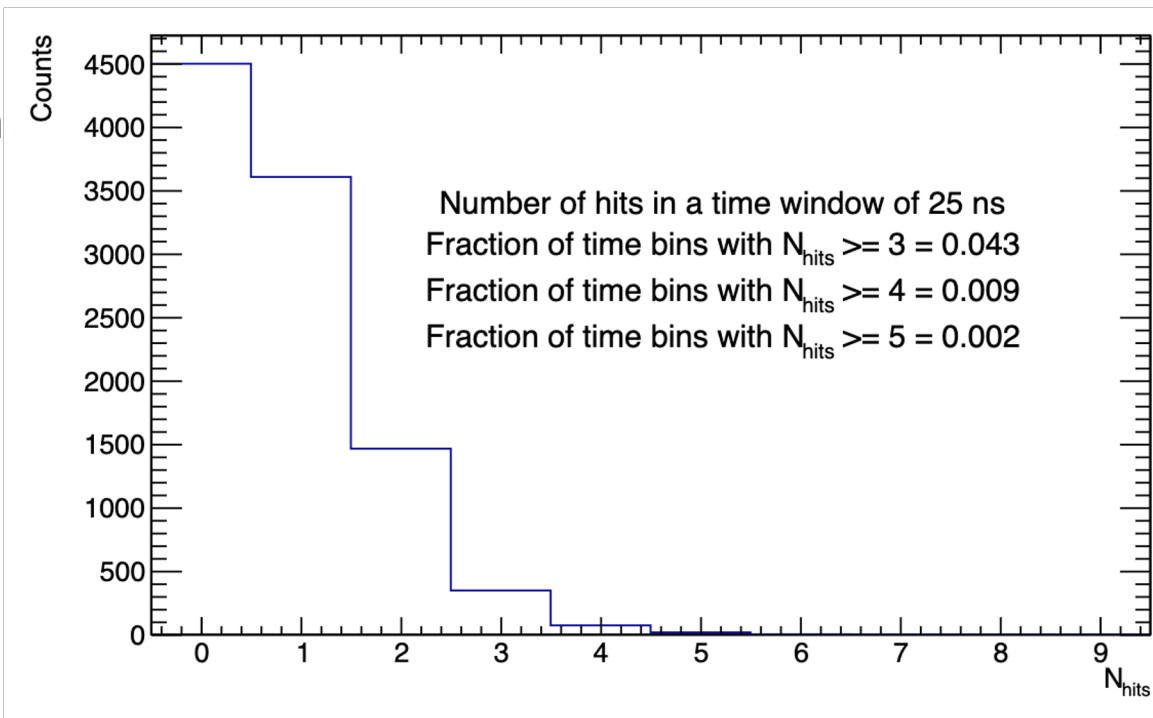
Take 25 ns window of HGCROC

Poissonian distribution with a mean of 0.8

- 25 ns \* 32 MHz

If threshold is set to 3 \* SPE pulse height, 4% of time bins will be triggered

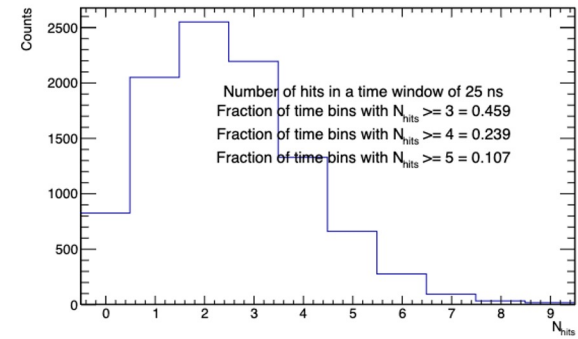
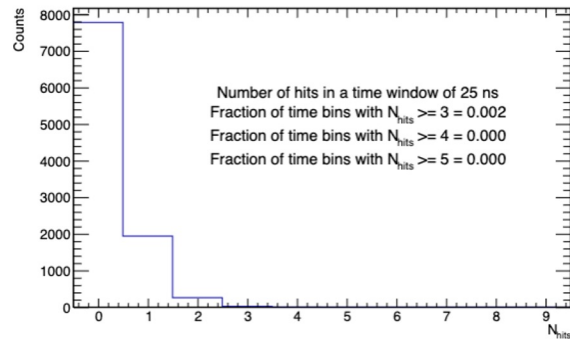
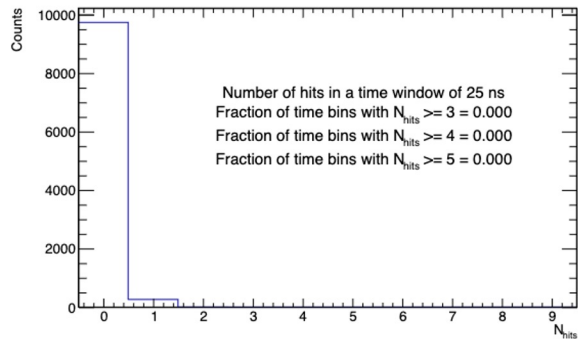
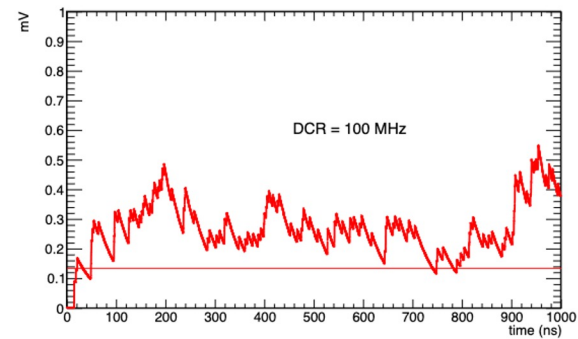
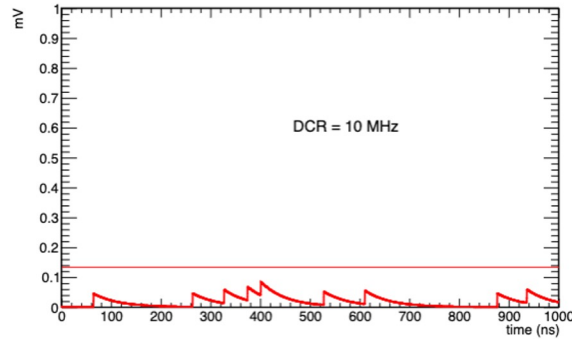
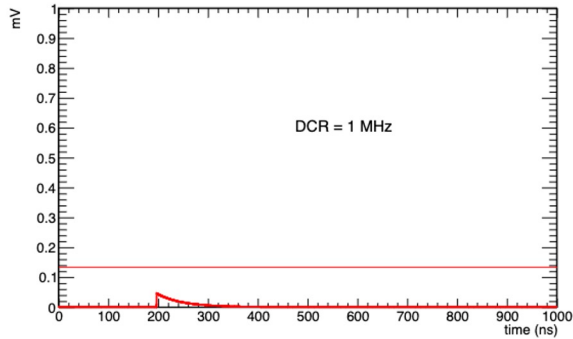
- 4% of the channels of the detector will be active in ToT mode at any given time





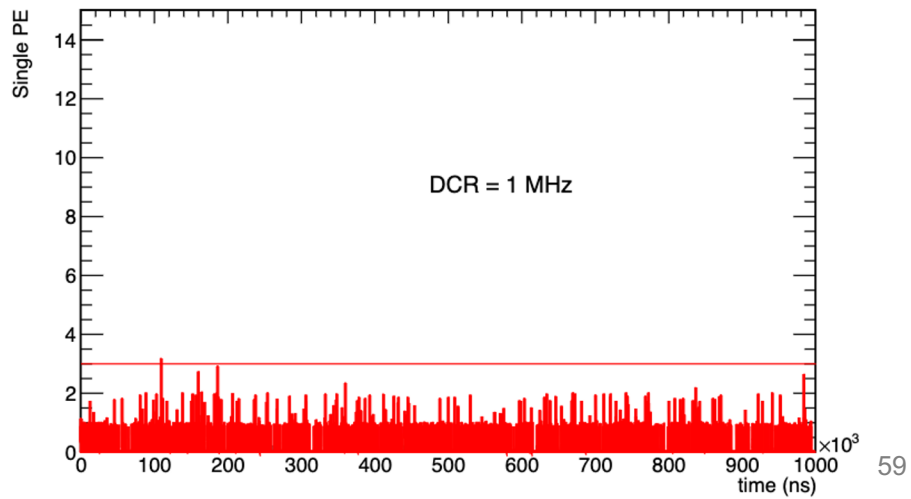
# Threshold of 3 p.e. likely excluded if DCR reaches 100 MHz

- This also poses an issue because we can't get around it with timing cuts in the same way as the dRICH, the detector could have a signal at any time
  - Bunch crossings every 10 ns, shorter than light propagation time



Threshold (p.e.)	Prob. Above threshold @ 1 MHz	Prob. Above threshold @ 10 MHz	Prob. Above threshold @ 30 MHz	Prob. Above threshold @ 50 MHz	Prob. Above threshold @ 100 MHz
2	0.01%	2%	29%	69%	99%
3	0.0005%	0.3%	8%	36%	94%
4	0%	0.04%	1.7%	14%	79%
5	0%	0.005%	0.2%	4%	57%
6	0%	0%*	0.03%	1.1%	35%
7	0%	0%*	0.005%	0.4%	17%

Numerical uncertainty of 0.03%



# TEMPERATURE DEPENDENCE

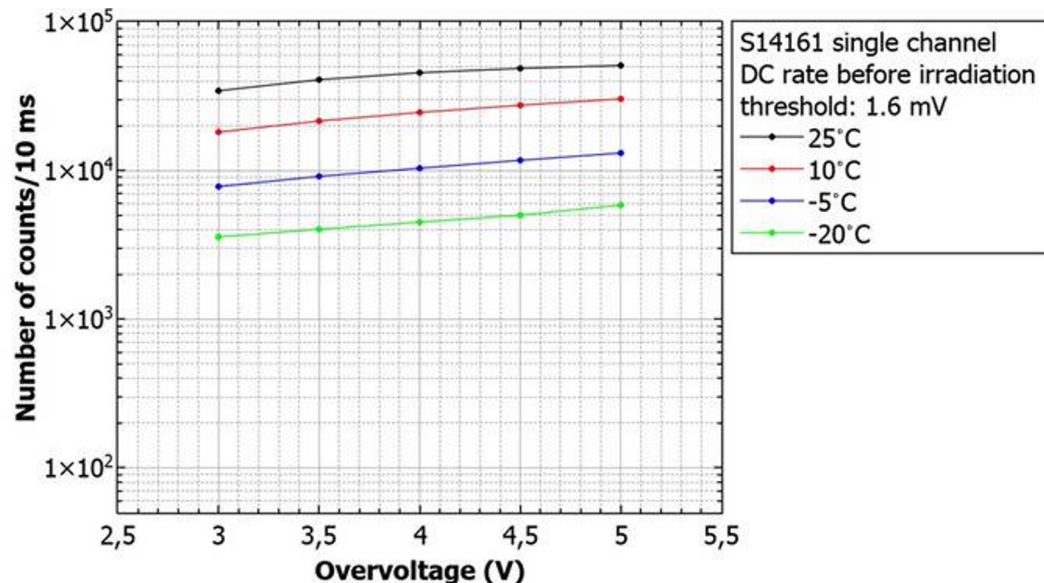
Conventional wisdom is that DCR is halved for every decrease of  $10^{\circ}\text{C}$

“Single-channel” here refers to 1/16 of a 4x4 array (S14161-6050HS-04)

- DCR numbers for ganged array should be 16x higher

To reach the  $\sim 4\text{ MHz}$  of GlueX with S14161-6050HS-04, need to go to  $-20^{\circ}\text{C}$  or lower

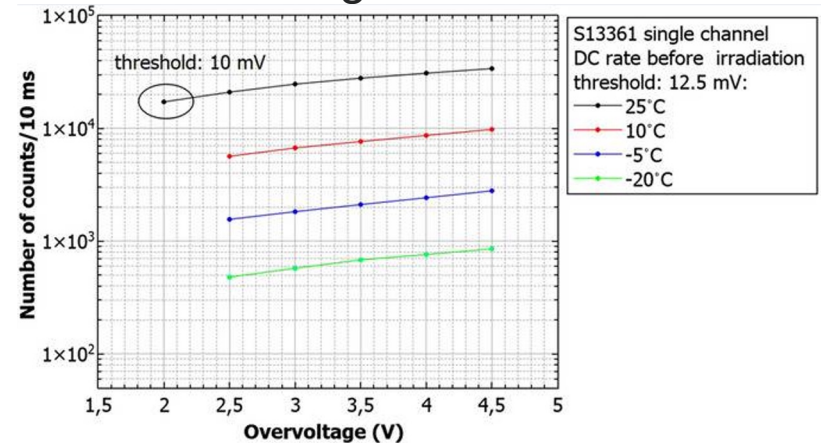
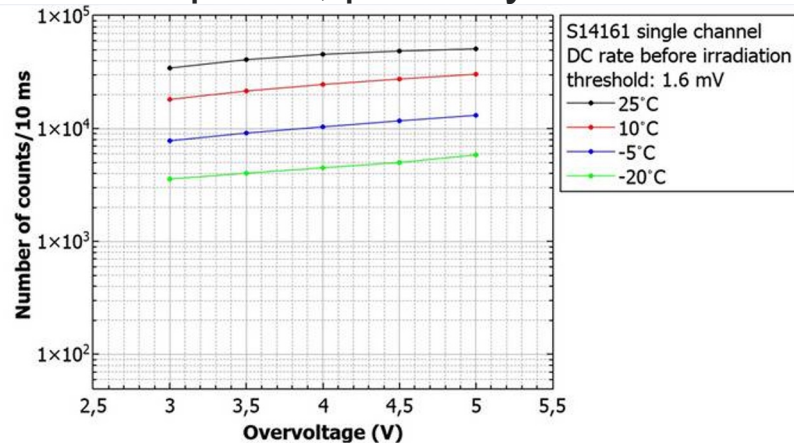
## Proton irradiation of SiPM arrays for POLAR-2



# TEMPERATURE DEPENDENCE

The authors of this paper report that DCR of the S14161 is 60% higher than S13361 at 25° C, and a factor of 5 higher at -20° C

- The DCR of the S14161 is apparently a much slower function of temperature
- This is bad, because it renders less effective the only handle we have over the DCR
- On the flip side, probably the DCR increases less if we go above 25° C...



On such a critical point, should consult an expert (Hamamatsu directly?) to see if this behavior is expected or not

# RADIATION DAMAGE

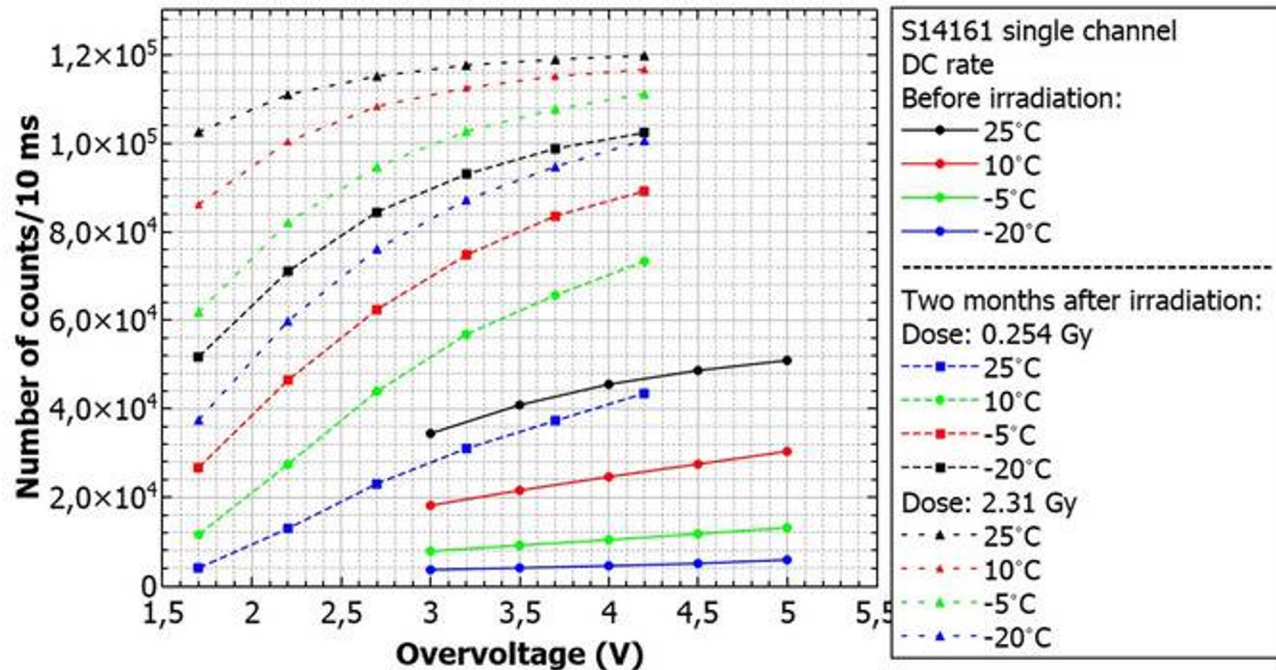
## Proton irradiation of SiPM arrays for POLAR-2

Pre-radiation DCR  
around 3 MHz (single  
channel)

- At 3V overvoltage

After ~200 Rad of  
proton radiation (and  
two months of  
waiting), DCR larger  
by factor of 4

- Half our expected  
dose

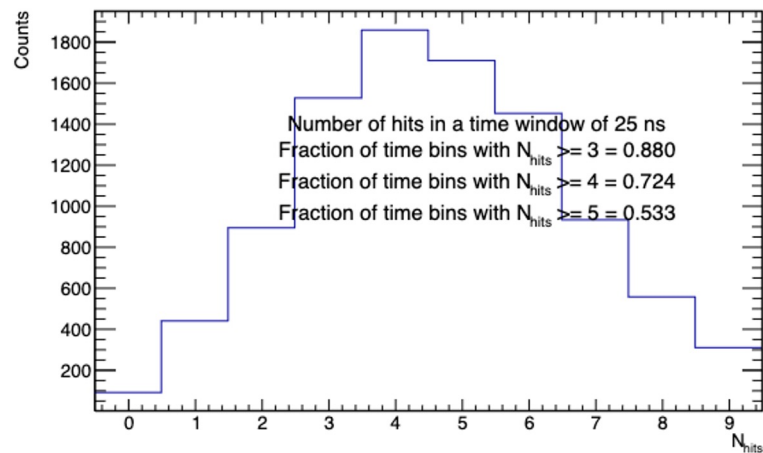
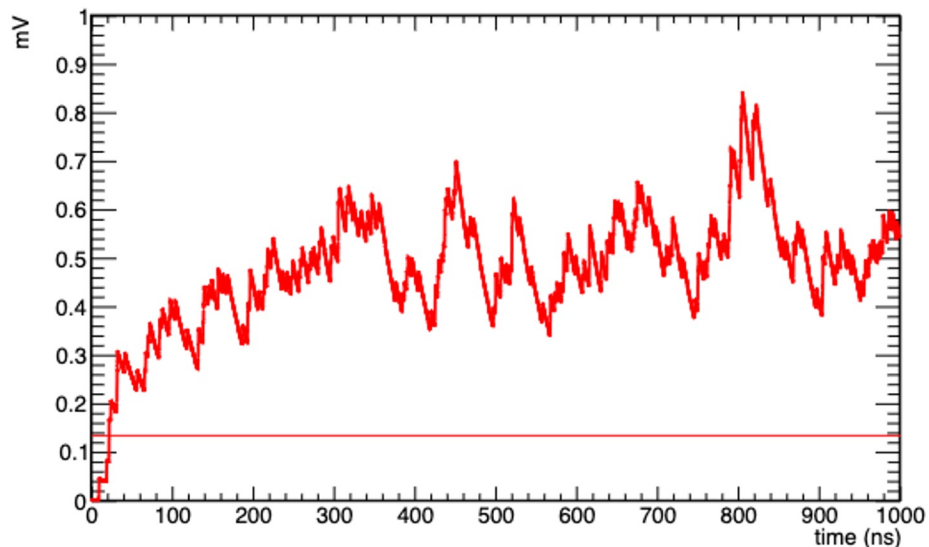


# RADIATION DAMAGE

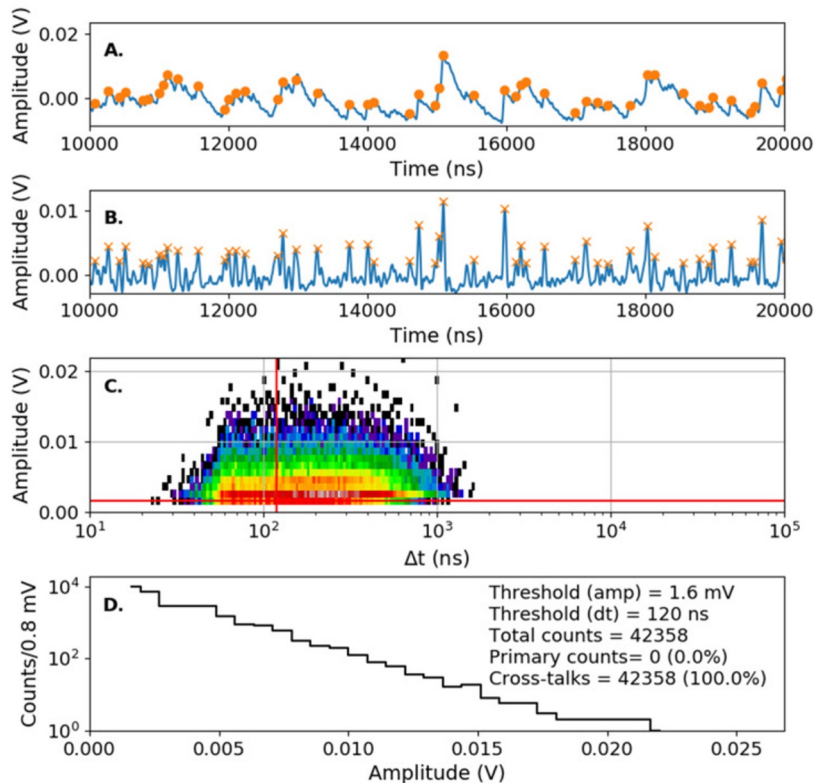
If we take the numbers provided in this paper seriously, expect 192 MHz of DCR after 200 Rad of radiation damage at room temperature

This is clearly too large, likely would swamp the MIP

- Threshold would need to be set at something like 9 Npe or higher



# S14161



# S13361

