

# P<sup>3</sup> : Precision in Polarized Pdfs

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Uncovering  
New Laws of  
Nature at the EIC

Brookhaven National Laboratory, Upton, NY USA  
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# Summary

- Brief intro to polarized pdfs/DIS
- Why Global and why NNLO
- BDSSV24 NNLO global fit      **Borsa, deF, Sassot, Stratmann, Vogelsang (2024)**
- QED and photon pdfs              **deF., Palma, Volonnino (2024)**
- Conclusions

# Brief intro polarized DIS

**cross section**

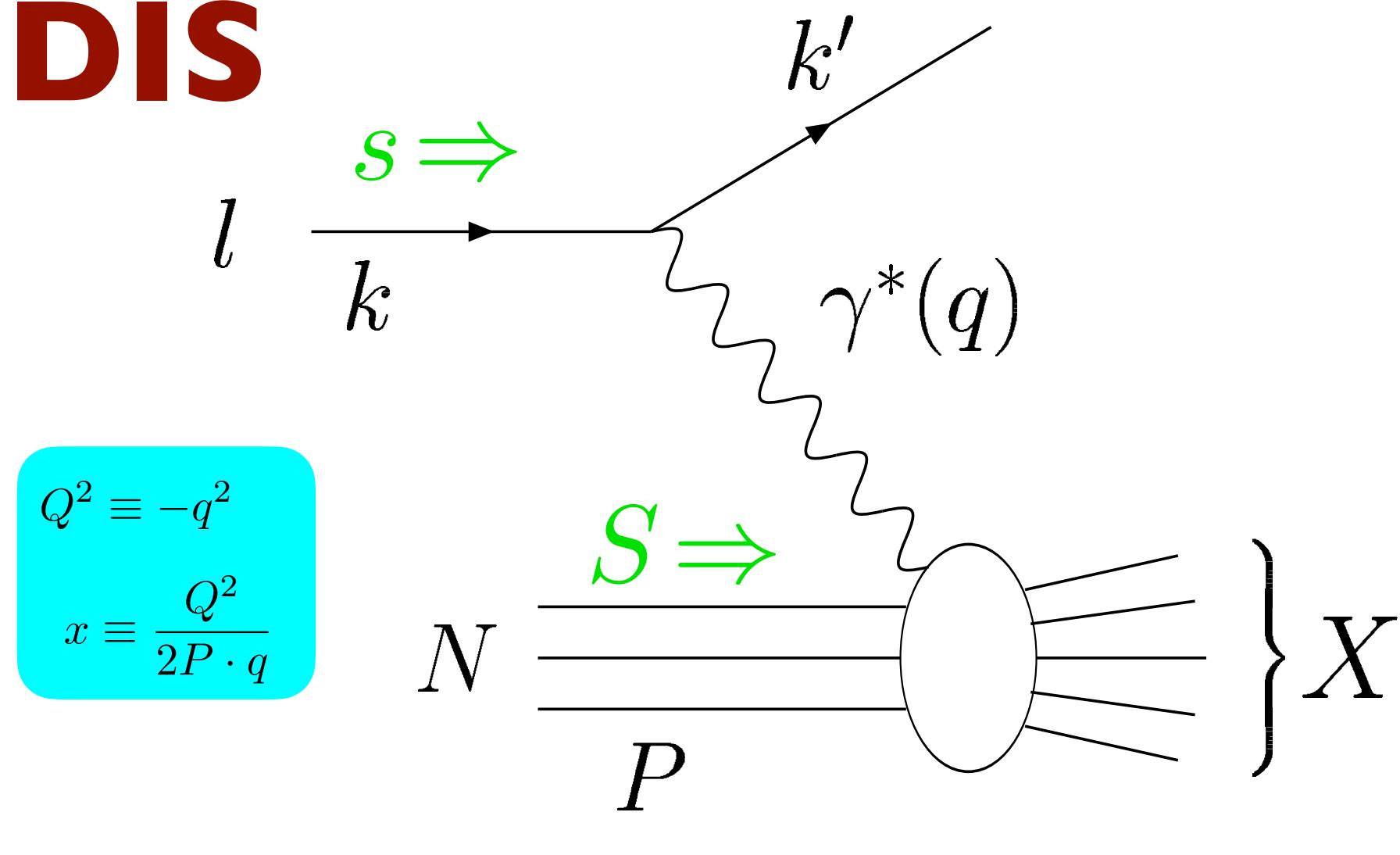
$$\frac{d\sigma}{dE' d\Omega} = \frac{\alpha^2}{Q^4} \frac{E'}{E} \underbrace{\mathcal{L}_{\mu\nu}(k, q, \textcolor{green}{S})}_{\text{leptonic}} \cdot \underbrace{\mathcal{W}^{\mu\nu}(P, q, \textcolor{green}{S})}_{\text{hadronic}}$$

- Hadronic tensor in terms of Structure Functions (non-perturbative)

$$\begin{aligned} \mathcal{W}^{\mu\nu}(P, q, S) &= \left( -g^{\mu\nu} + \frac{q^\mu q^\nu}{q^2} \right) F_1(x, Q^2) + \left( P^\mu - \frac{P \cdot q}{q^2} q^\mu \right) \left( P^\nu - \frac{P \cdot q}{q^2} q^\nu \right) F_2(x, Q^2) \\ &\quad + i M \varepsilon^{\mu\nu\rho\sigma} q_\rho \left[ \frac{S_\sigma}{P \cdot q} g_1(x, Q^2) + \frac{S_\sigma(P \cdot q) - P_\sigma(\textcolor{green}{S} \cdot q)}{(P \cdot q)^2} g_2(x, Q^2) \right] \end{aligned}$$

Asymmetry  $A_1(x, Q^2)$

$$A_{||} = \frac{d\sigma^{(\rightarrow\leftarrow)} - d\sigma^{(\rightarrow\rightarrow)}}{d\sigma^{(\rightarrow\leftarrow)} + d\sigma^{(\rightarrow\rightarrow)}} = D(y) \frac{g_1(x, Q^2)}{F_1(x, Q^2)}$$



$$\begin{aligned} Q^2 &\equiv -q^2 \\ x &\equiv \frac{Q^2}{2P \cdot q} \end{aligned}$$

$F_i$ : unpolarized structure functions

$g_i$ : polarized structure functions

# Brief intro polarized DIS

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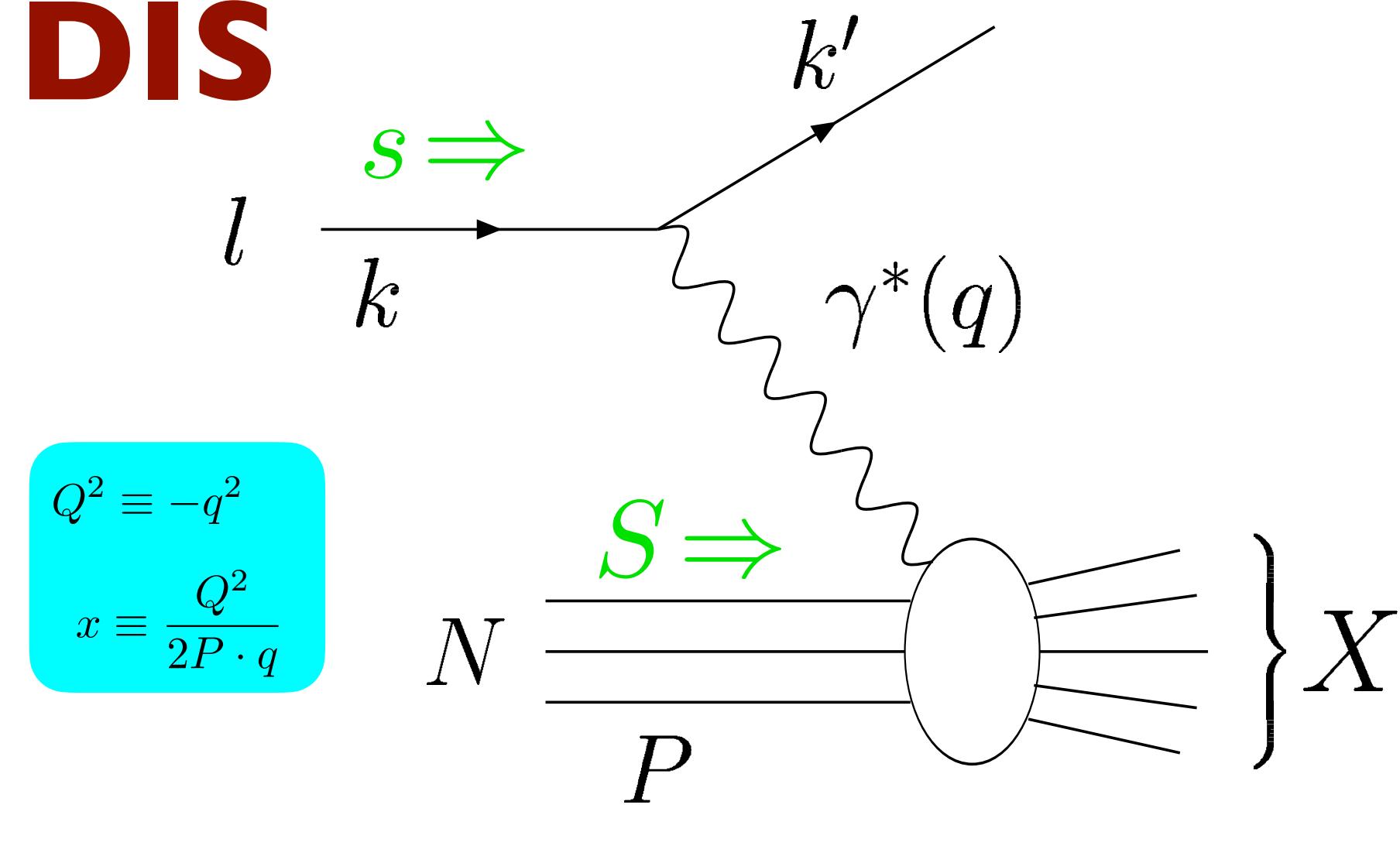
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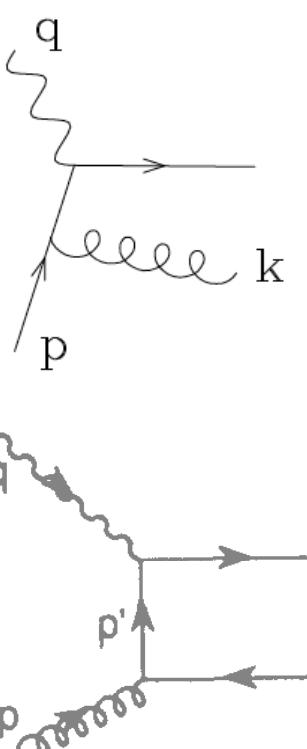
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$$g_1(x, Q^2) = \sum_i \frac{e_i^2}{2} \Delta q_i(x, Q^2) \quad \text{LO}$$

$$\begin{aligned} &+ \frac{\alpha_s}{2\pi} \sum_i \frac{e_i^2}{2} \int_x^1 \frac{dz}{z} \Delta C_i^q(z) \Delta q_i(x/z, Q^2) \\ &+ \frac{\alpha_s}{2\pi} \sum_i \frac{e_i^2}{2} \int_x^1 \frac{dz}{z} \Delta C^g(z) \Delta g(x/z, Q^2) \end{aligned}$$

**NLO**



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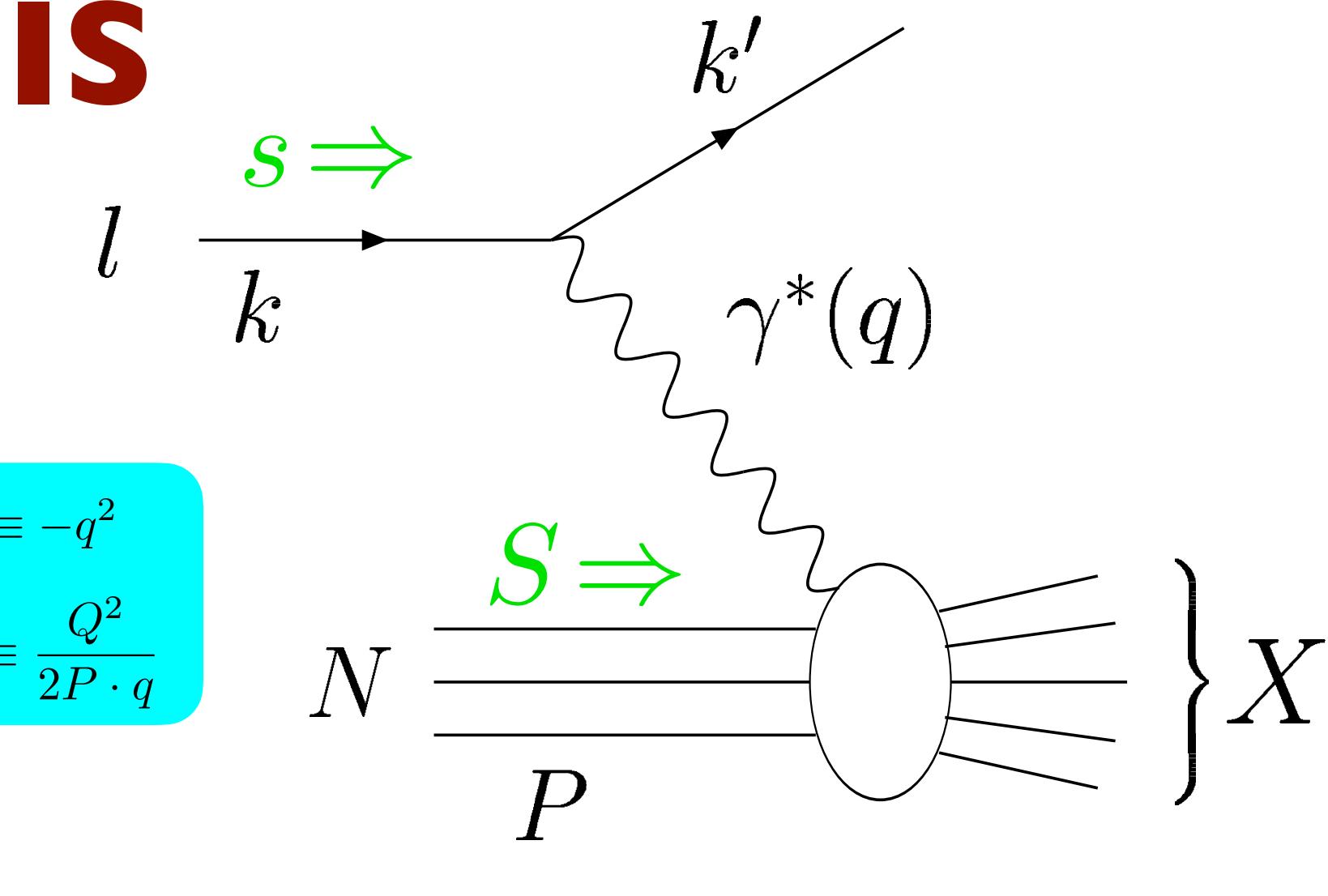
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- quark and gluon contribution to proton spin

$$\Delta\Sigma = \sum_i \int_0^1 \Delta q_i(x, Q^2) dx \quad \Delta G = \int_0^1 \Delta g(x, Q^2) dx$$

needs extrapolation to 0



$$\begin{aligned} Q^2 &\equiv -q^2 \\ x &\equiv \frac{Q^2}{2P \cdot q} \end{aligned}$$

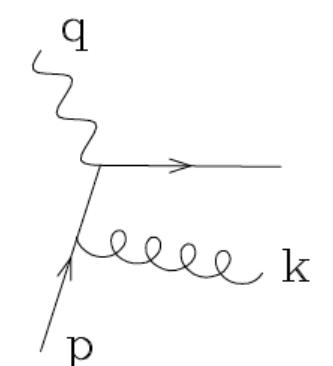
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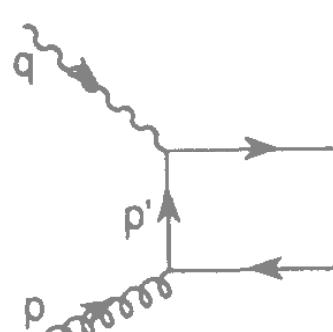
$$g_1(x, Q^2) = \sum_i \frac{e_i^2}{2} \Delta q_i(x, Q^2) \text{ LO}$$

$$+ \frac{\alpha_s}{2\pi} \sum_i \frac{e_i^2}{2} \int_x^1 \frac{dz}{z} \Delta C_i^q(z) \Delta q_i(x/z, Q^2)$$

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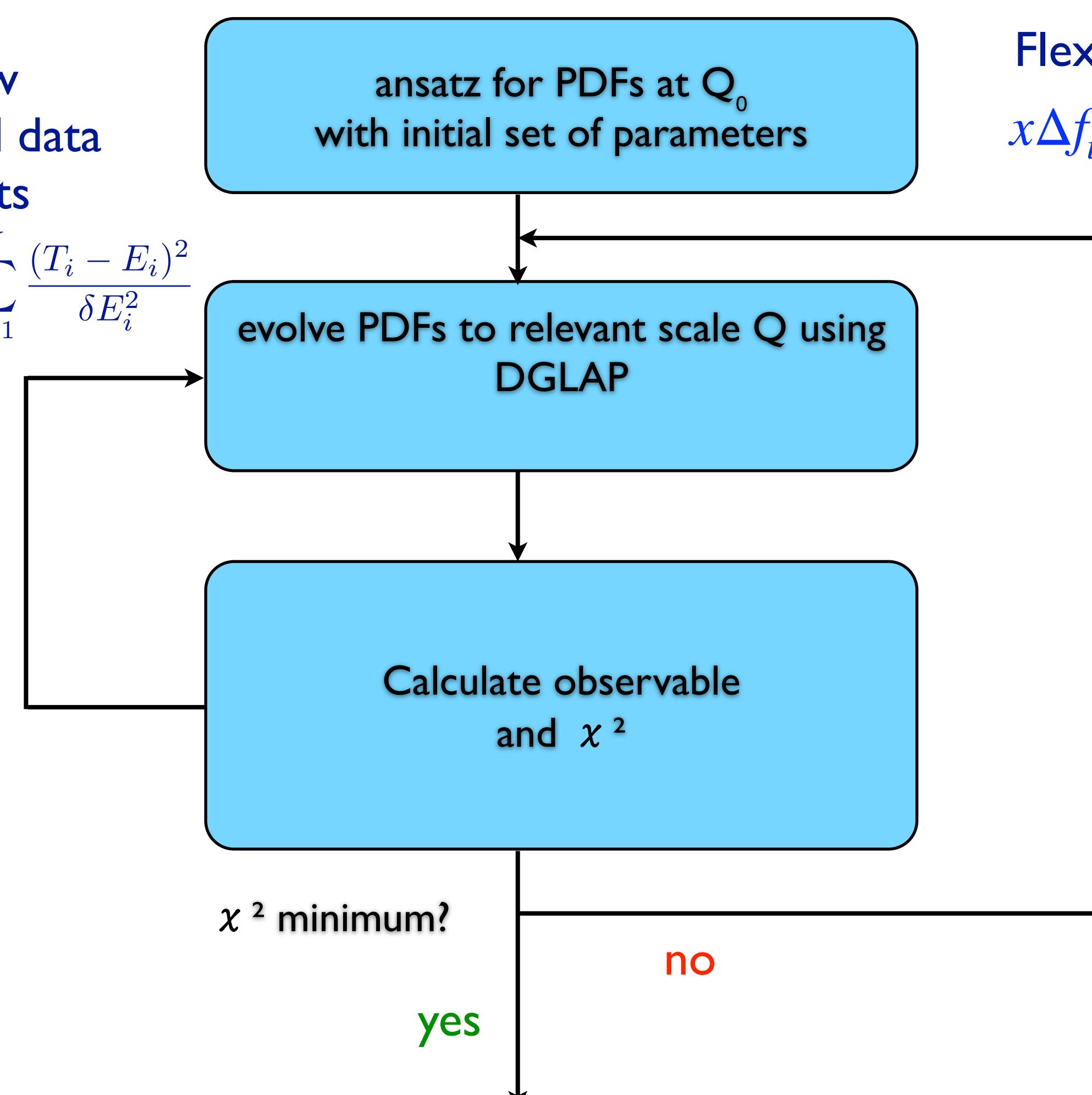
NLO



# PDFs obtained by **global fit** : $\chi^2$ minimization

~few  
hundred data  
points

$$\chi^2 = \sum_{i=1}^N \frac{(T_i - E_i)^2}{\delta E_i^2}$$



result : best fit + uncertainties  
(many fits)

Flexible parametrization at  $Q_0$

$$x \Delta f_i(x, Q_0^2) = N_i x^{\alpha_i} (1-x)^{\beta_i} (1 + \gamma_i x + \eta_i x^{\kappa_i})$$

change  
parameters  
~ few  
thousand  
times

several NLO analyses

DSSV  
NNPDF  
JAM

“partial” NNLO

MAP Bertone, Chiefa, Nocera (2024)

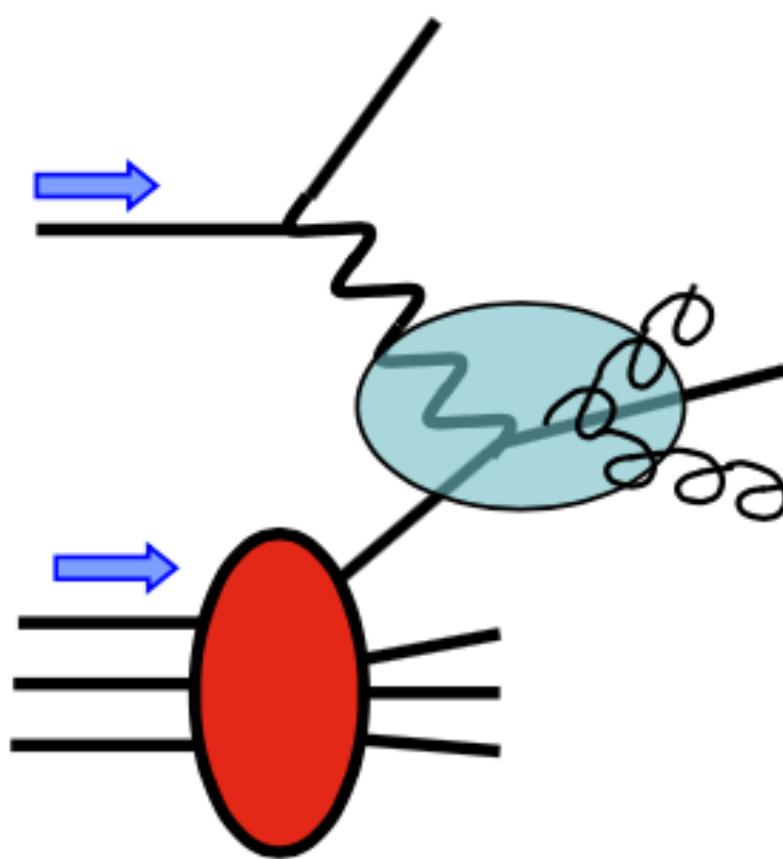
NNPDF coming: talk by Juan  
first global NNLO

**BDSSV24**

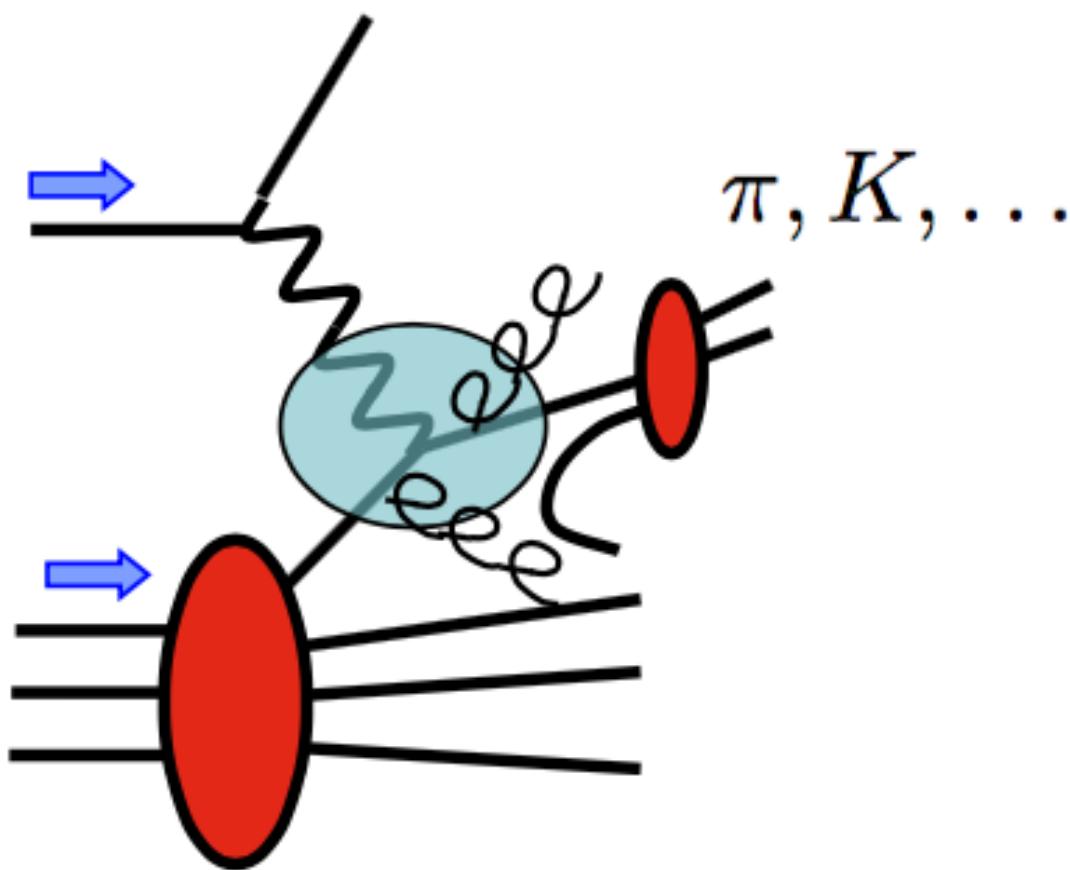
Borsa, deF, Sassot, Stratmann, Vogelsang (2024)

# Why global?

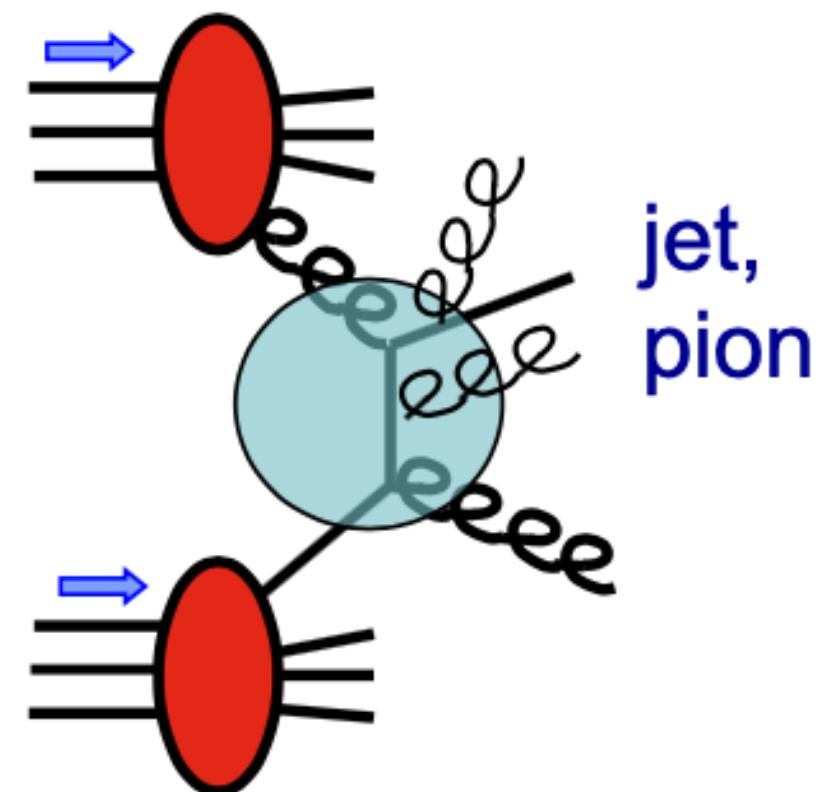
- Observables provide sensitivity on different pdfs



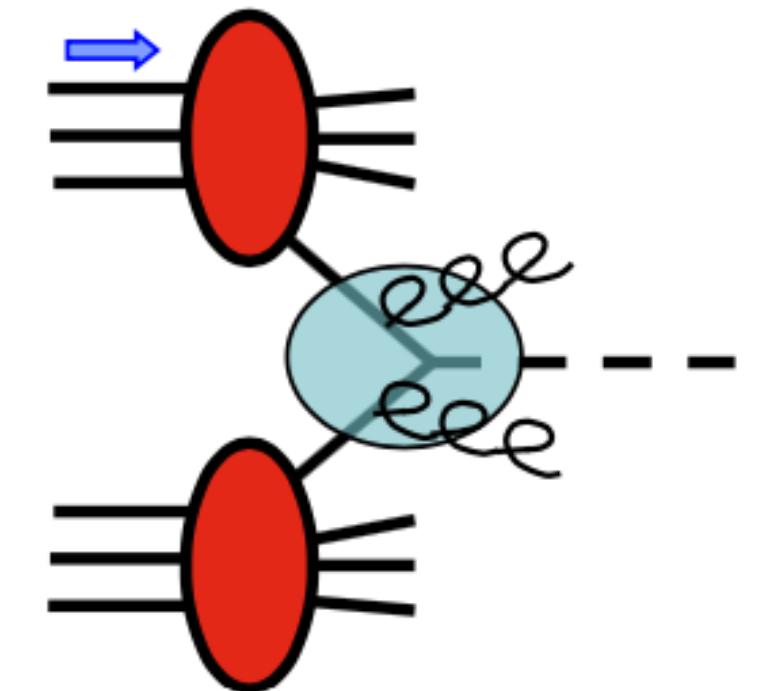
DIS



SIDIS



pp high- $p_T$



W bosons

$$\Delta\Sigma, (\Delta q + \Delta\bar{q}), \Delta g$$
$$q = u, d, s$$

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$$q = u, d, s$$

$$\Delta g, \Delta q, \Delta\bar{q}$$

$$\Delta q, \Delta\bar{q}$$

# Why NNLO?

TH needed to match experimental accuracy/understanding

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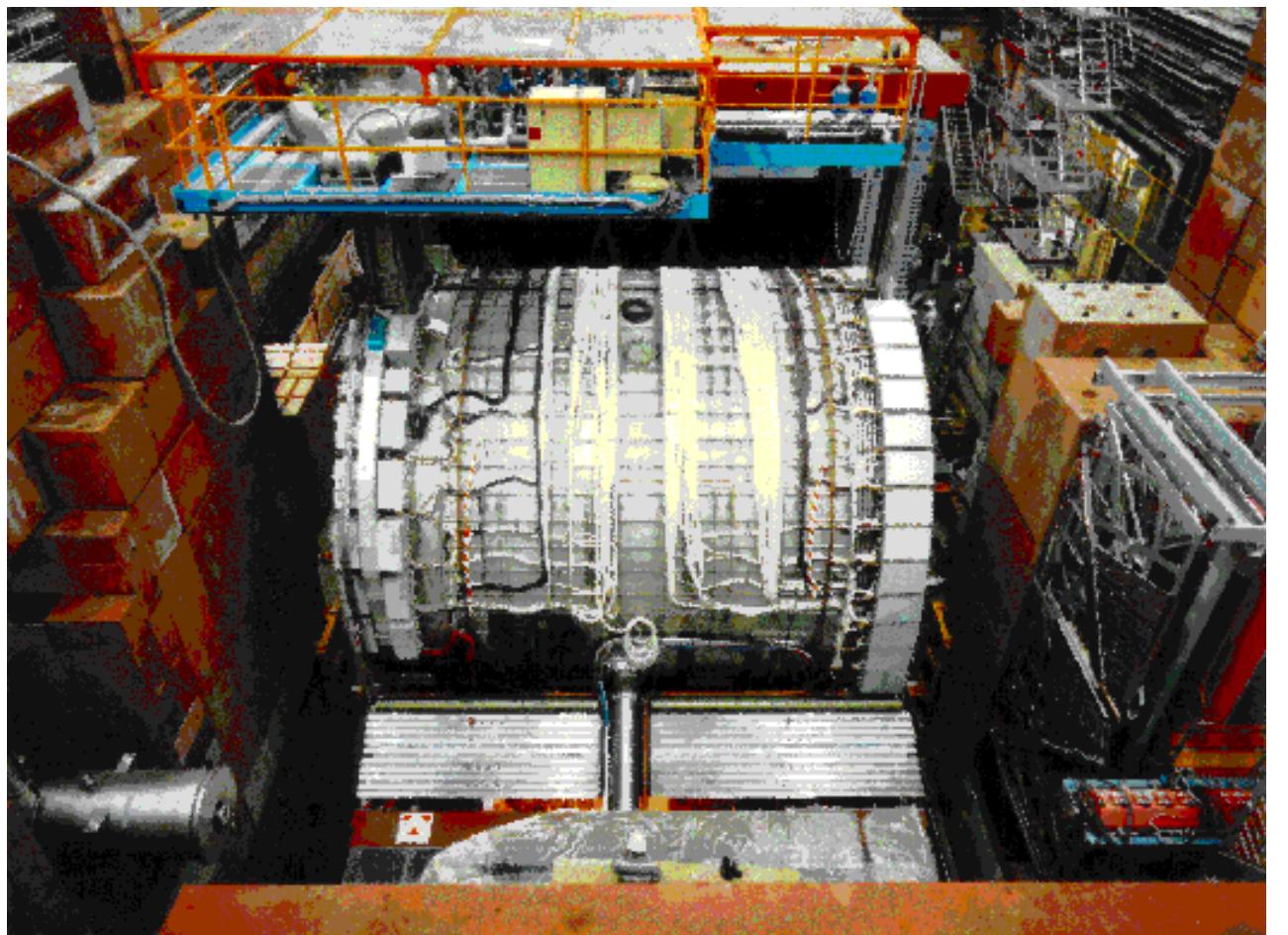
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70's, 80's LO  
(Born Level)



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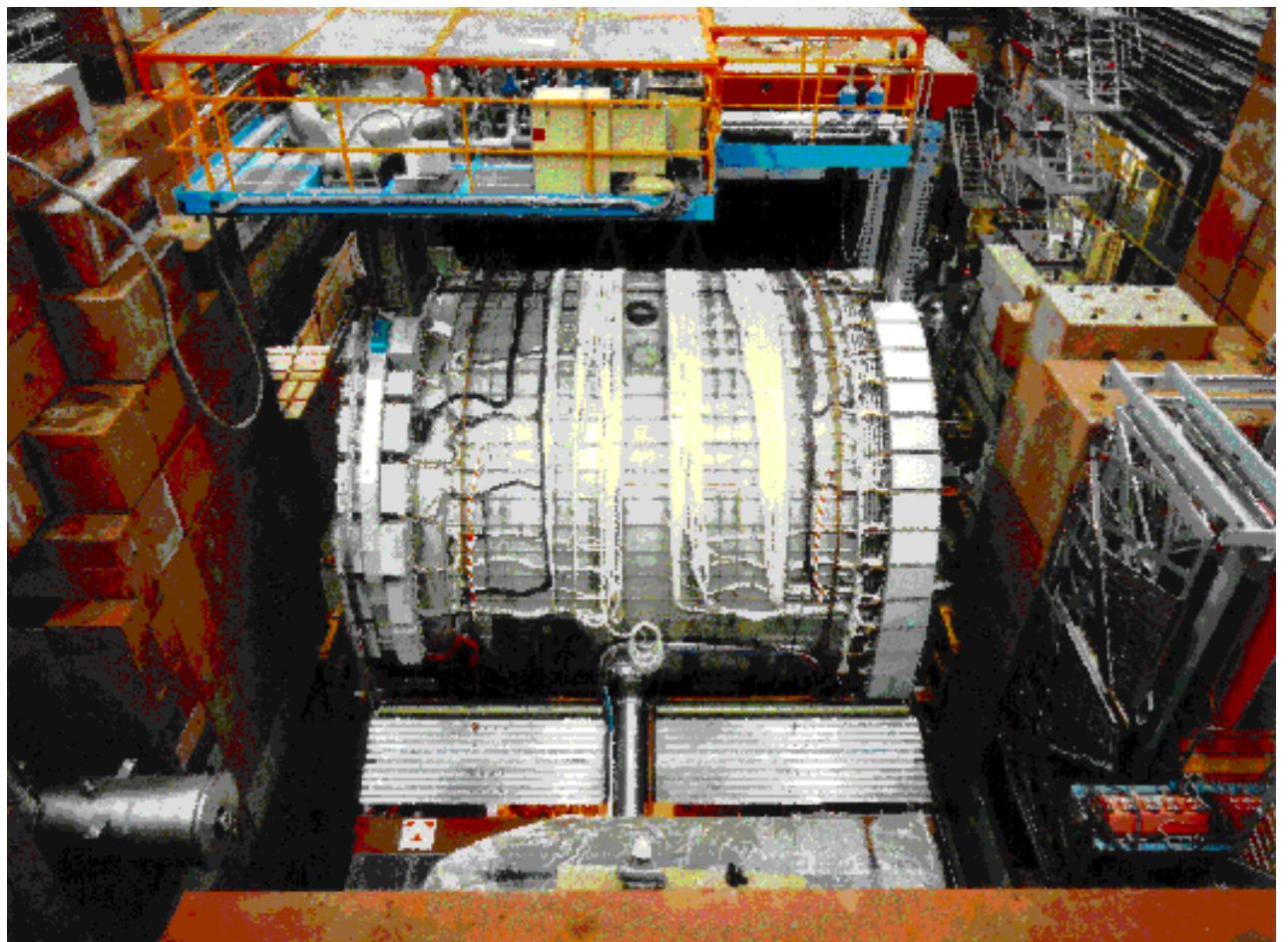
LEP- HERA 90's -00's  
NLO+

70's, 80's LO  
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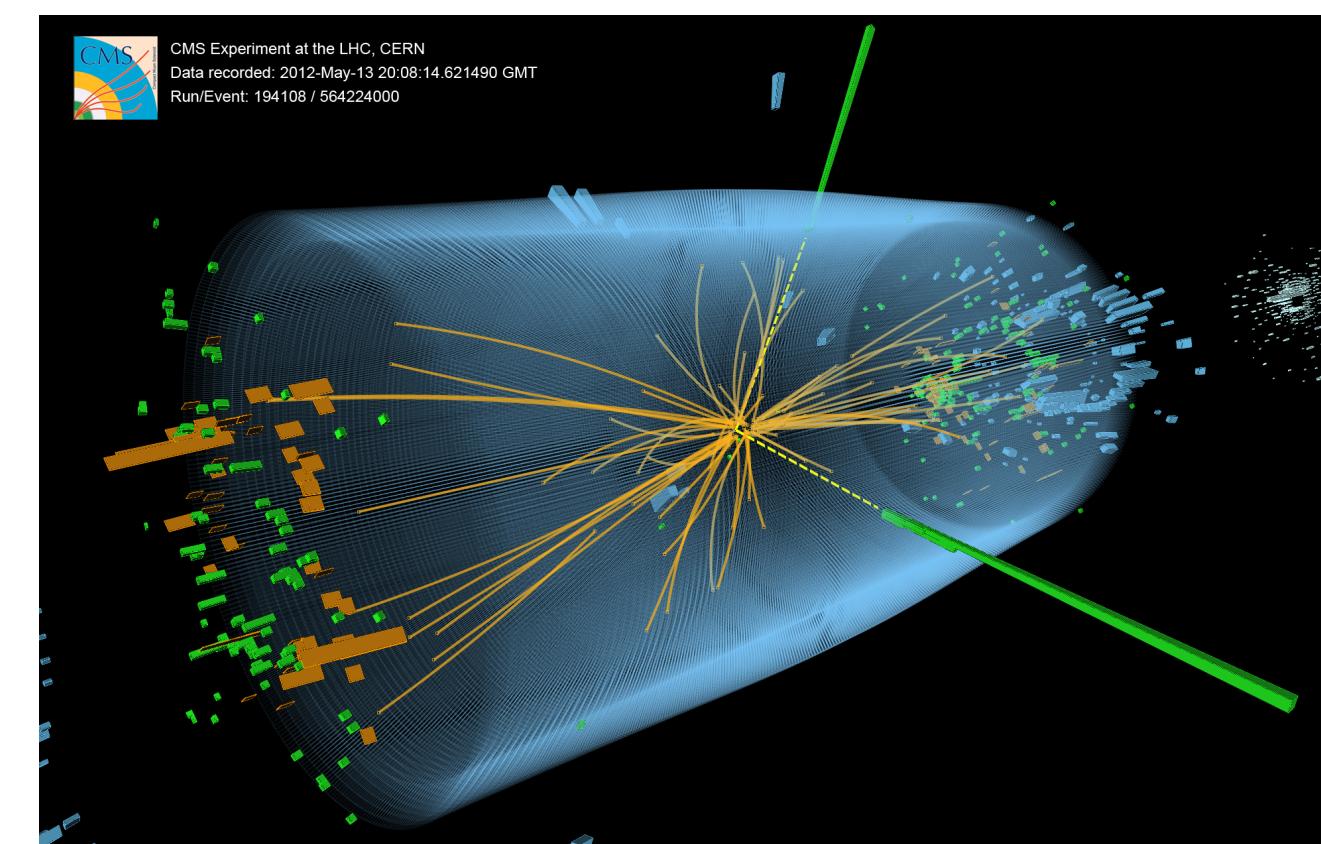


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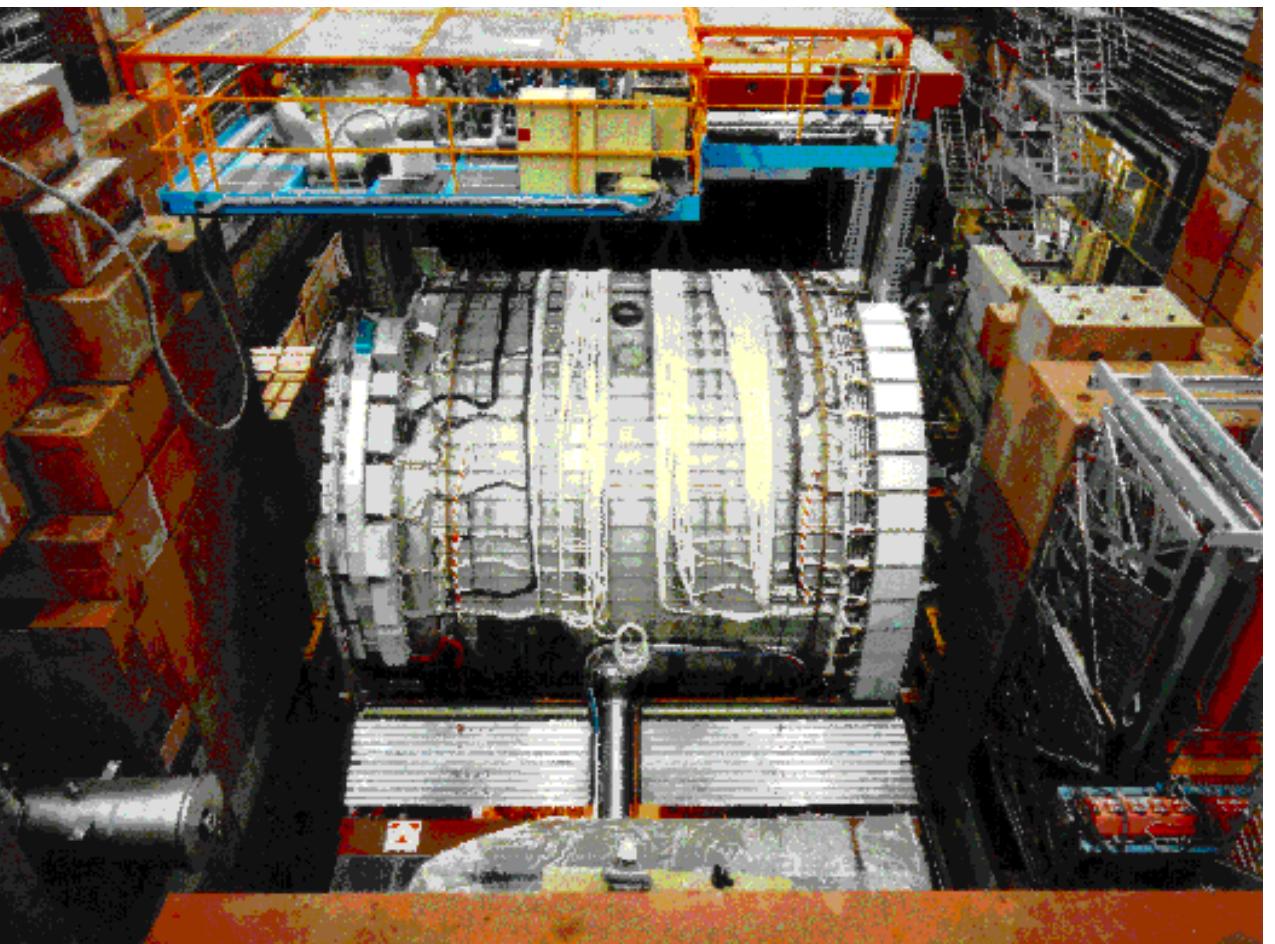


LHC 10's-20's-  
NLO+PS  
NNLO and beyond



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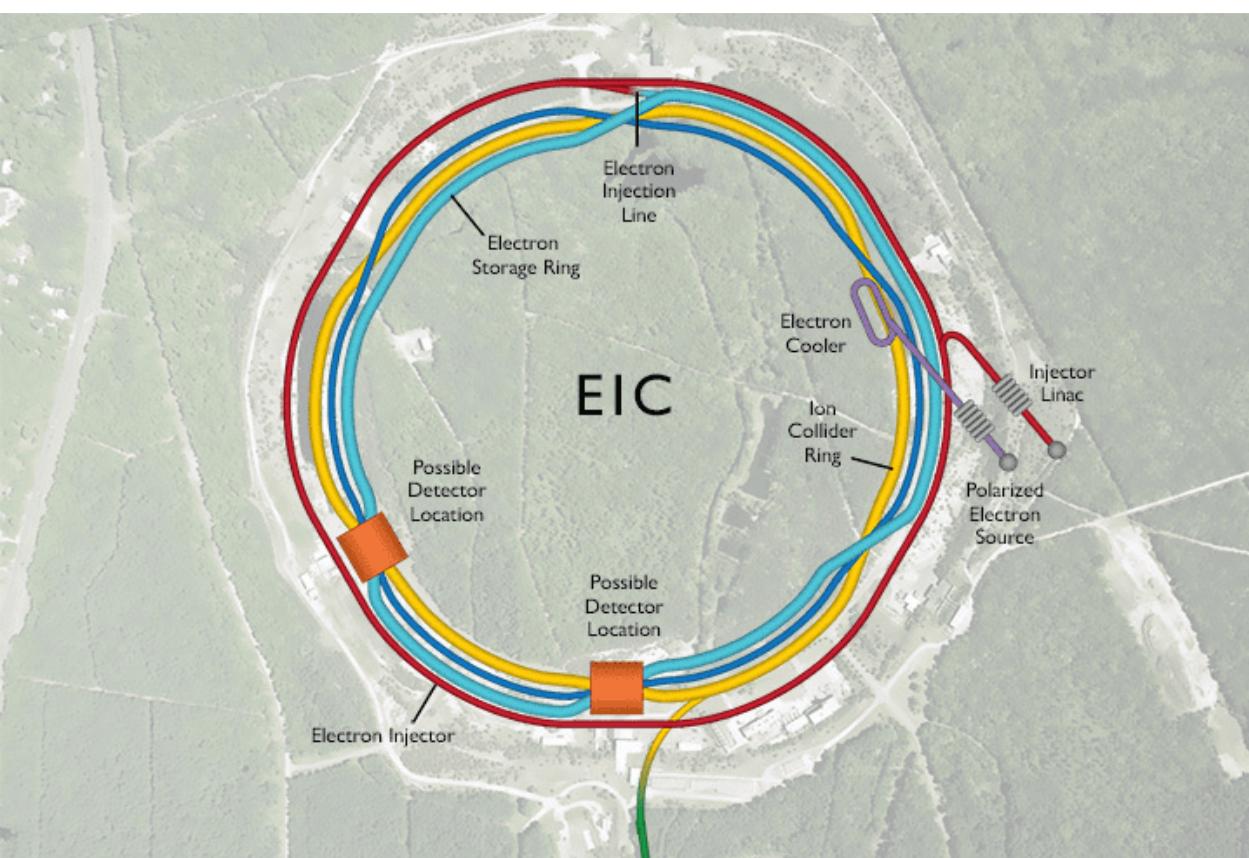


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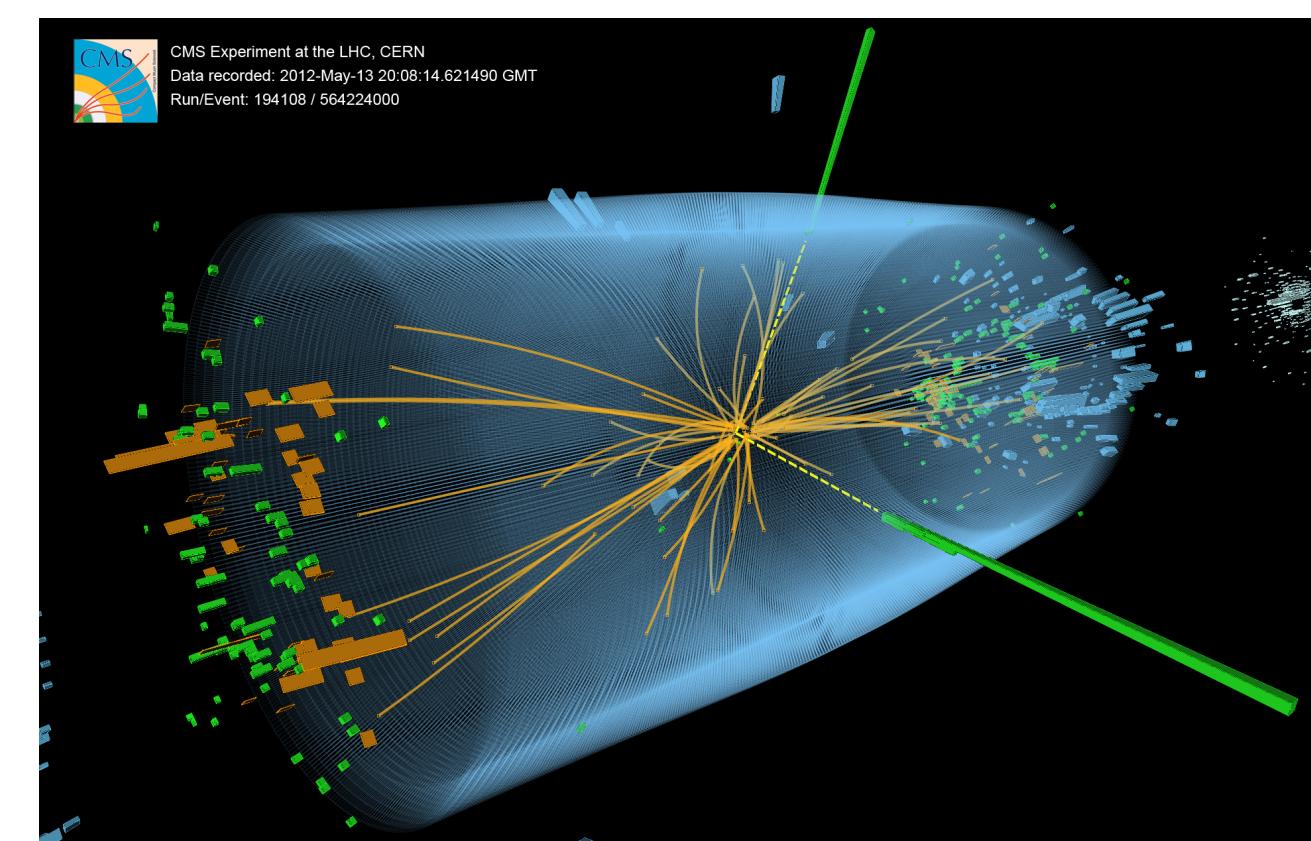
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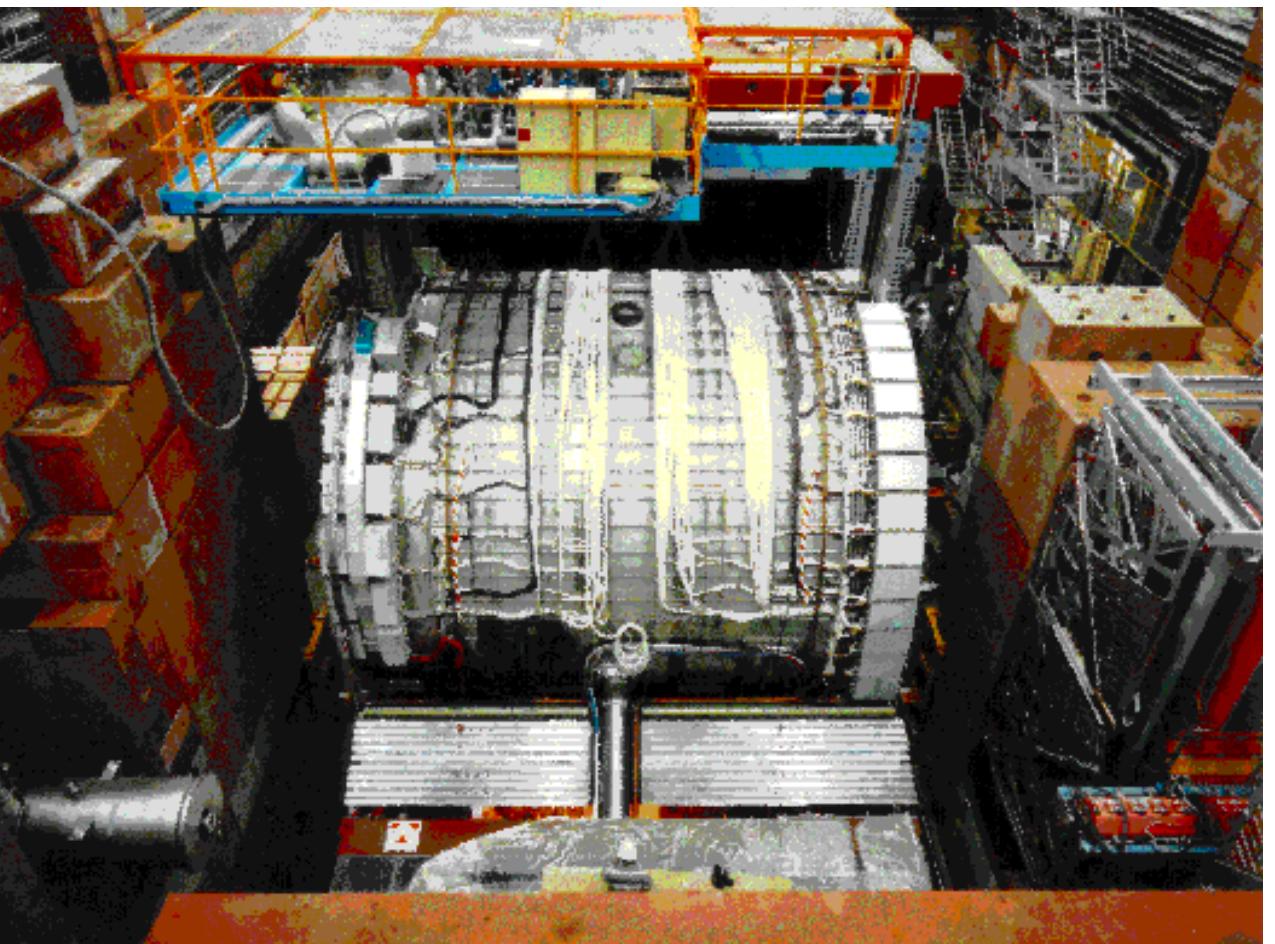


EIC 30's-  
 $\geq$ NNLO+ and beyond



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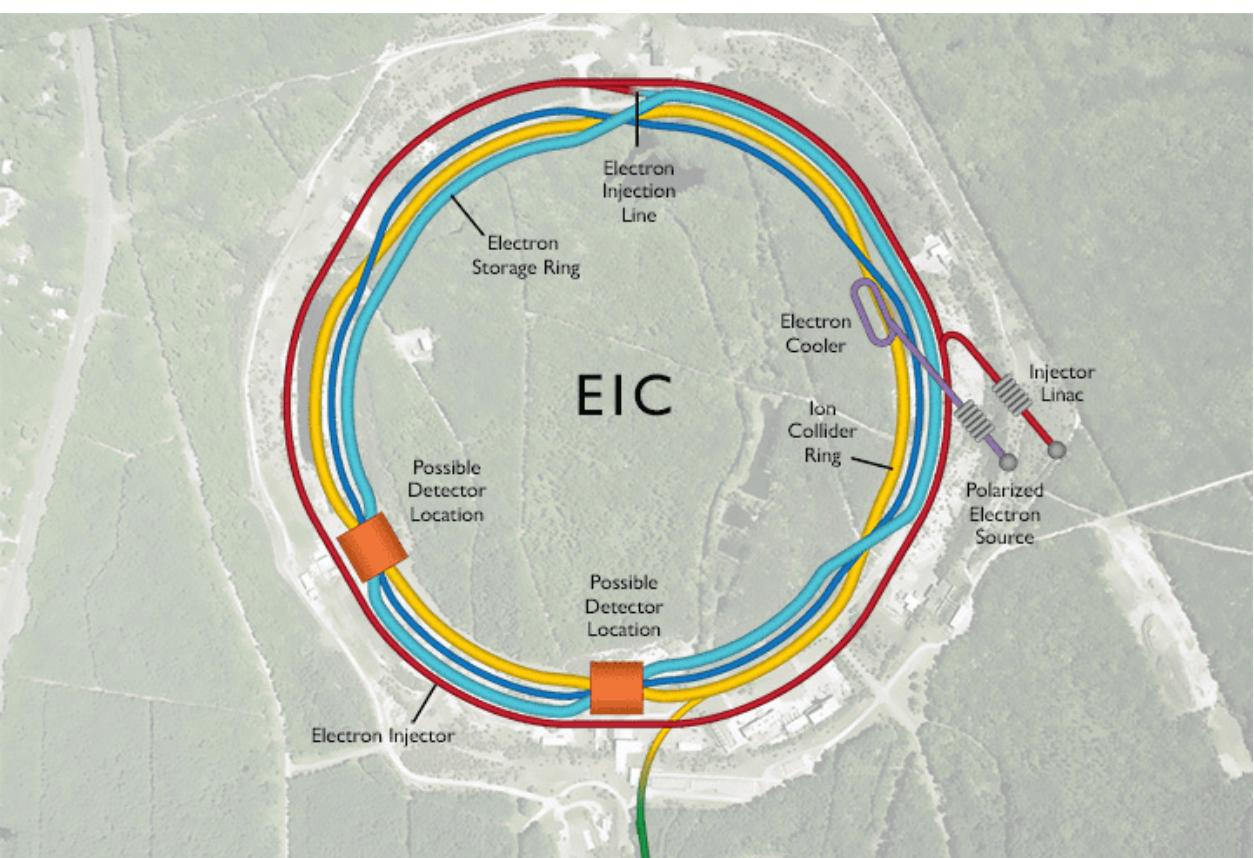


LEP- HERA 90's -00's  
NLO+

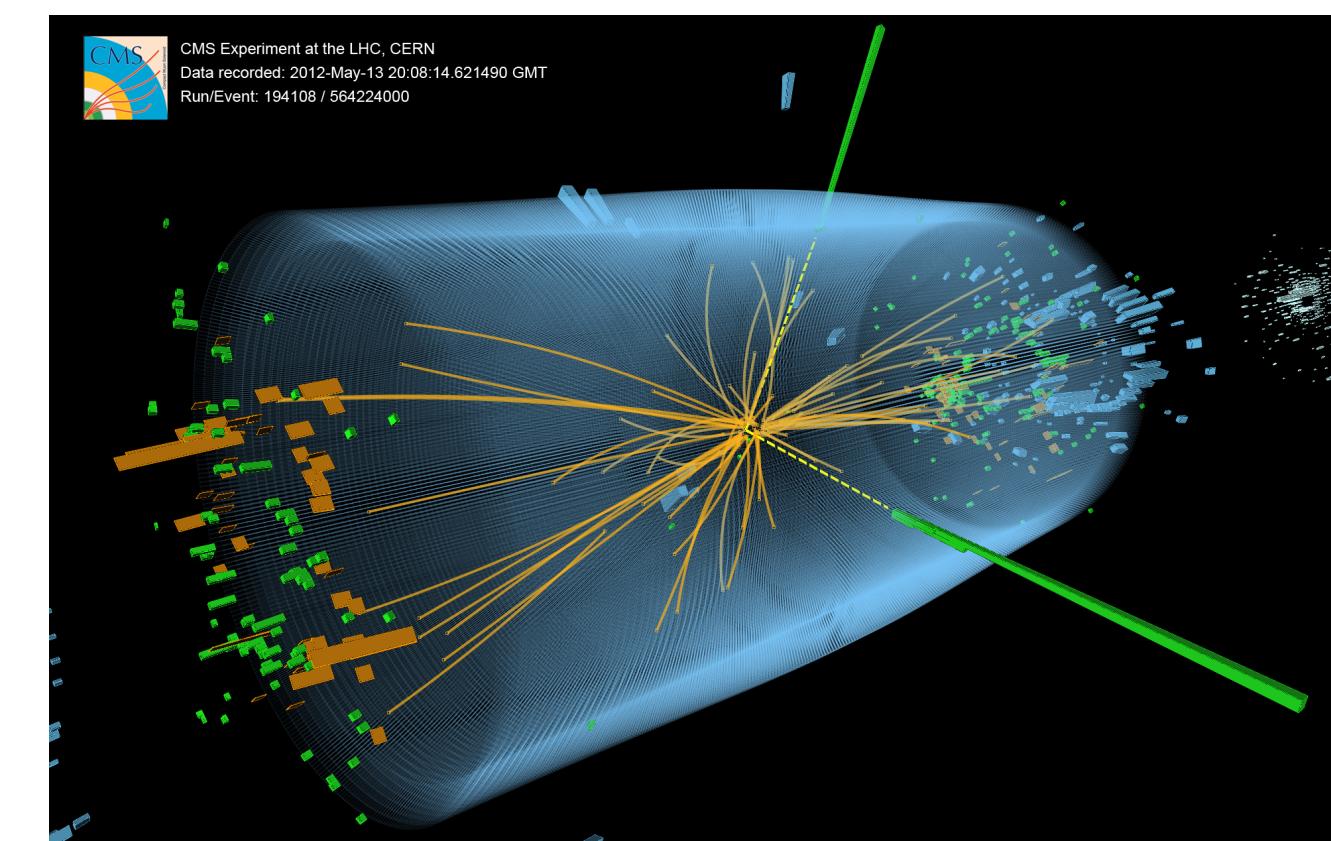
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LHC 10's-20's-  
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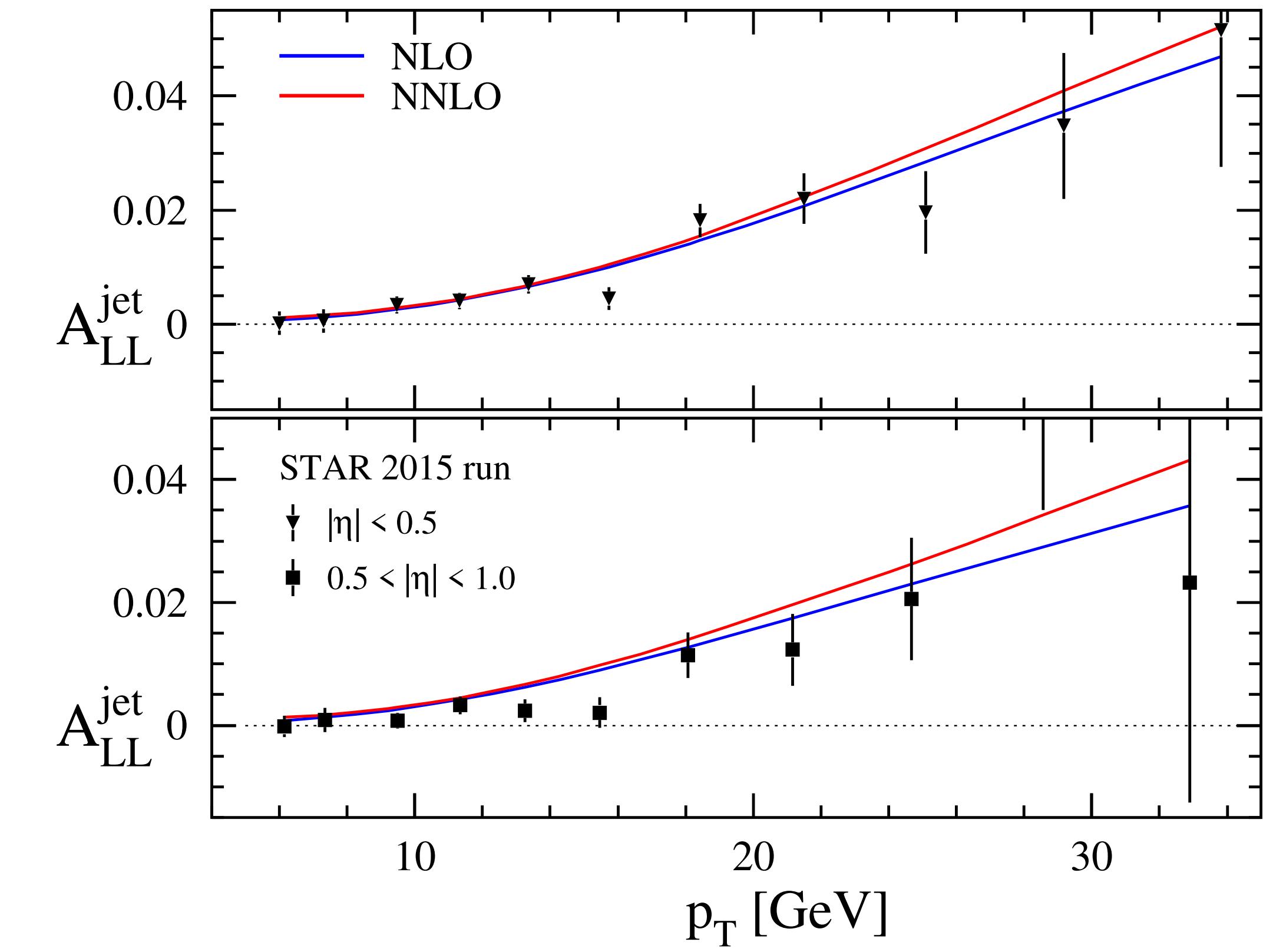
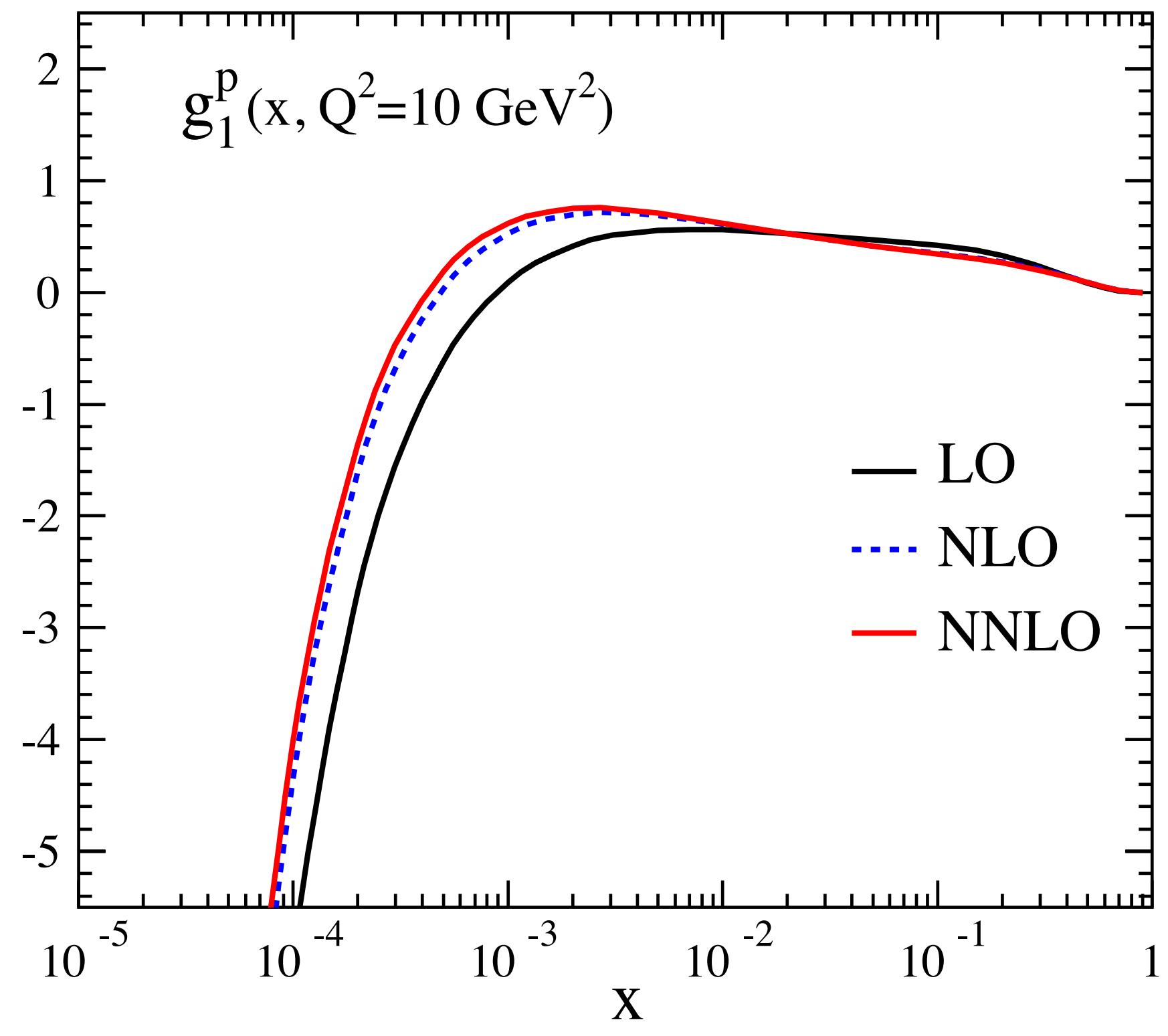
EIC 30's-  
 $\geq$ NNLO+ and beyond



TH effort driven by  
experimental possibilities!

# Why NNLO?

► Precision!



- Not all observables known at NNLO accuracy yet in polarized case



**DIS**

**NNLO  $g_1$  coefficients**

E. B. Zijlstra and W. L. van Neerven (1994)



**DY**

**NNLO  $p p \rightarrow W \rightarrow e^\pm \nu^{(-)}$**

R. Boughezal, Hai Tao Li, F. Petriello (2021)



**Evolution**

**NNLO evolution kernels**

S. Moch, J.A.M. Vermaseren, A. Vogt (2014, 2015)

A. Vogt, S. Moch, M. Rogal, J.A.M. Vermaseren (2008)



**HQ**

**Heavy quark matching coefficients**

Bierenbaum et al. (2022)



**SIDIS**

**NNLO SIDIS coefficients**

not ready for use in global fit

L.Bonino, T.Gehrmann, M.Löchner, K.Schonwald (2024)

hard to implement in Mellin space

(became available while performing fit)

S.Goyal, R.Lee, S.Moch, V.Pathak, N.Rana, V.Ravindran (2024)



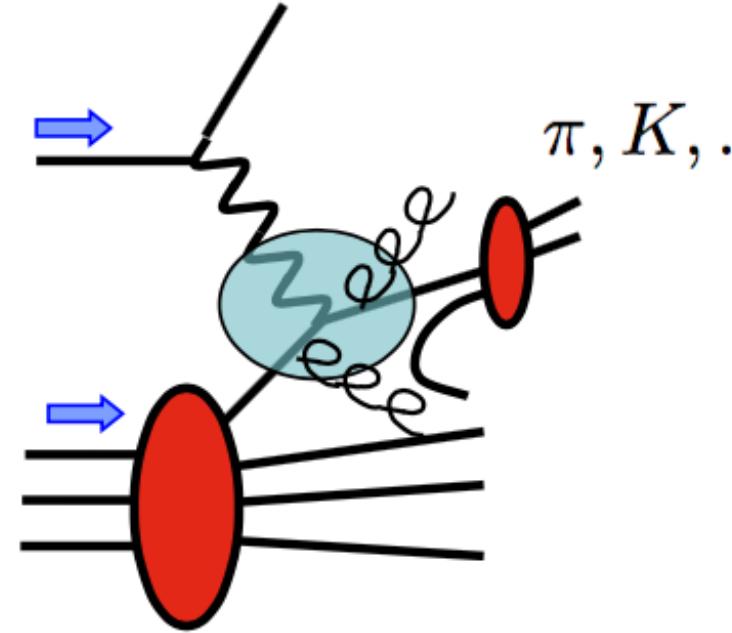
**RHIC jets, dijets, hadrons**

**$p p \rightarrow \text{jets}, \pi^{\pm,0}$**

only NLO known

► Rely on Soft approximation at NNLO : dominant terms in the threshold limit (on top of full NLO)

Use threshold resummation at NNLL to obtain the leading terms in the perturbative expansion



SIDIS

$$\alpha_s^k \delta(1 - \hat{x}) \left( \frac{\ln^m(1 - \hat{z})}{1 - \hat{z}} \right)_+, \quad \alpha_s^k \delta(1 - \hat{z}) \left( \frac{\ln^m(1 - \hat{x})}{1 - \hat{x}} \right)_+, \text{ with } m \leq 2k - 1.$$

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$$\hat{x} = \frac{Q^2}{2p_a \cdot q}$$

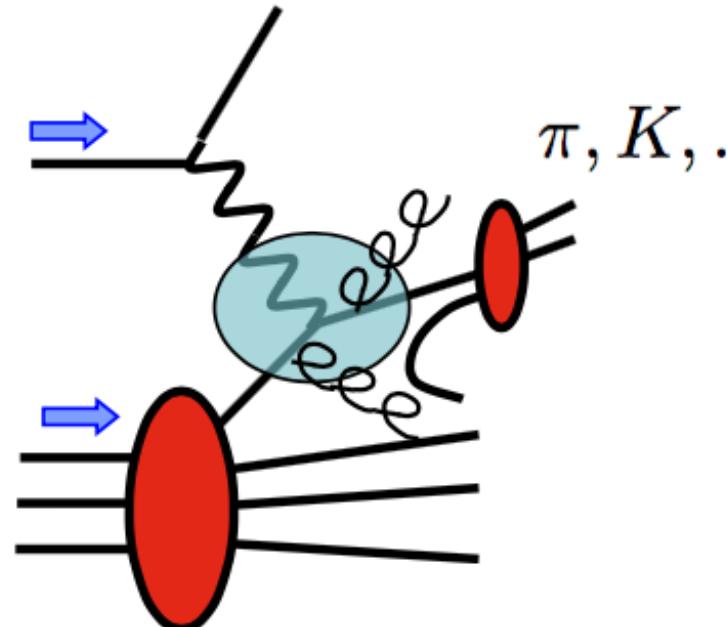
$$\hat{z} = \frac{p_a \cdot p_b}{p_a \cdot q}$$

All distributions known (up to N3LO), missing regular terms

M.Abele, D.de Florian,W.Vogelsang (2022,2023)

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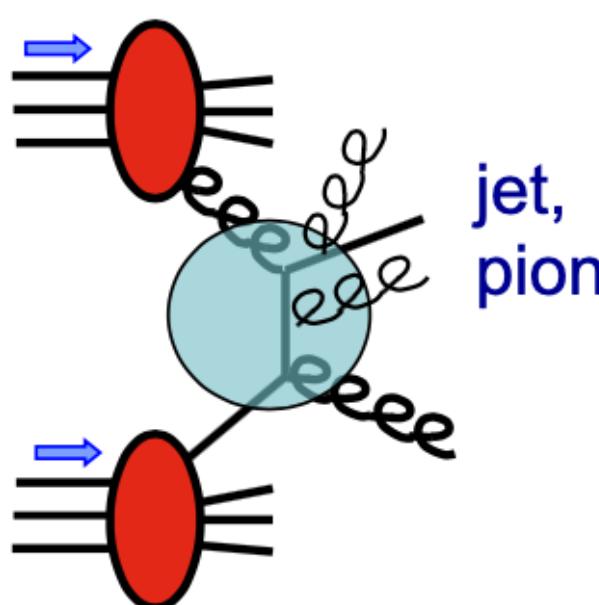
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pp high- $p_T$

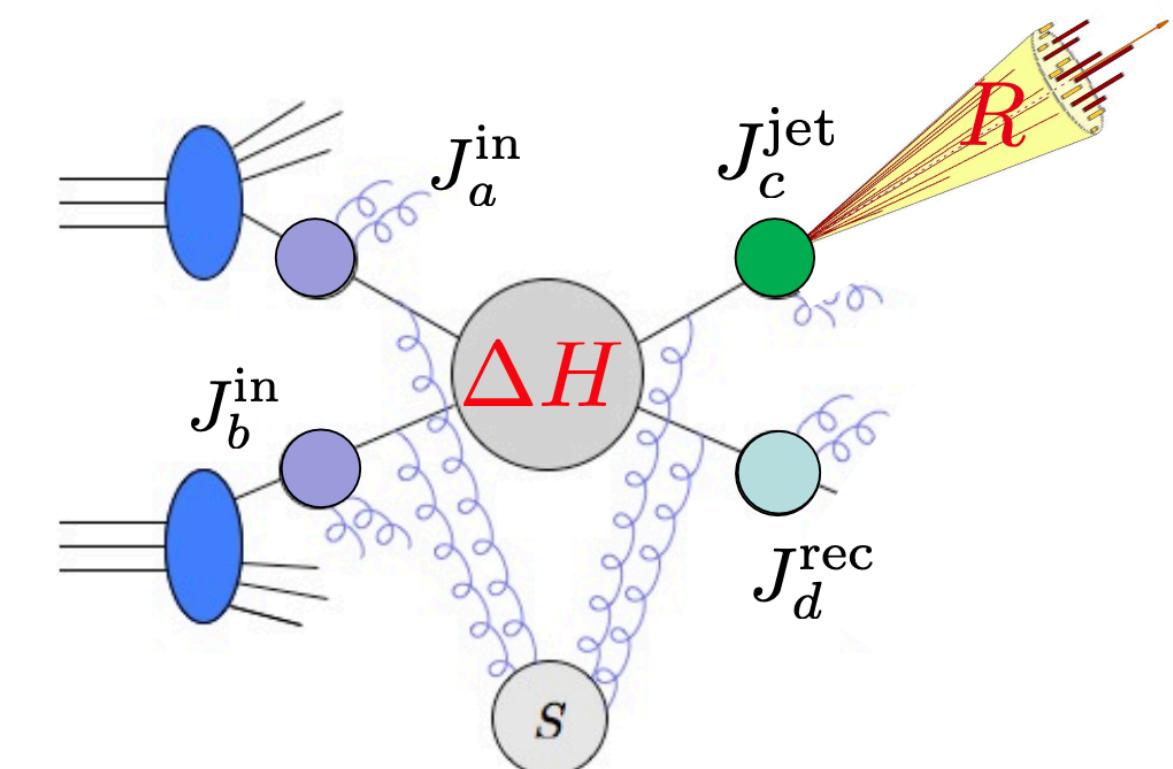
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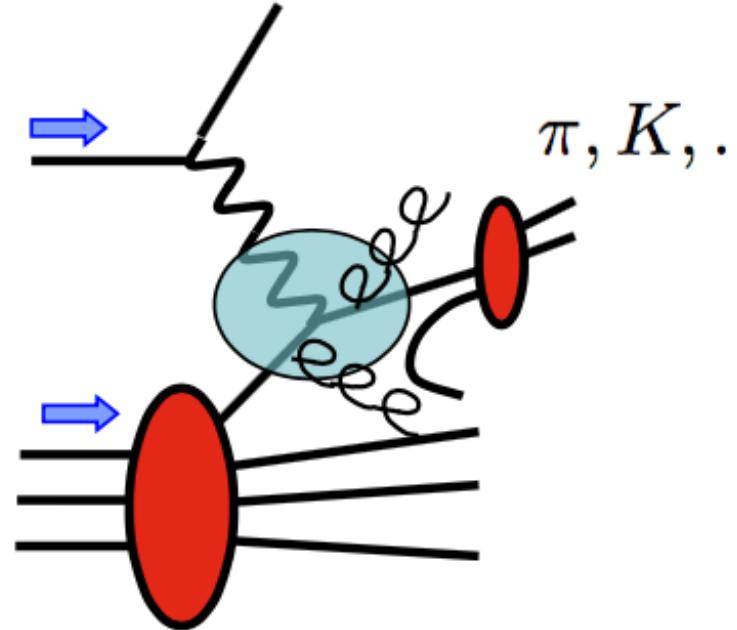
$$\Delta\hat{\sigma}^{ab \rightarrow cd} \sim J_a^{\text{in}} \times J_b^{\text{in}} \times J_c^{\text{jet}} \times J_d^{\text{rec}} \times \text{Tr}[\Delta H \mathcal{S}^\dagger S \mathcal{S}]_{ab \rightarrow cd}$$

D.deF,W.Vogelsang (in preparation)



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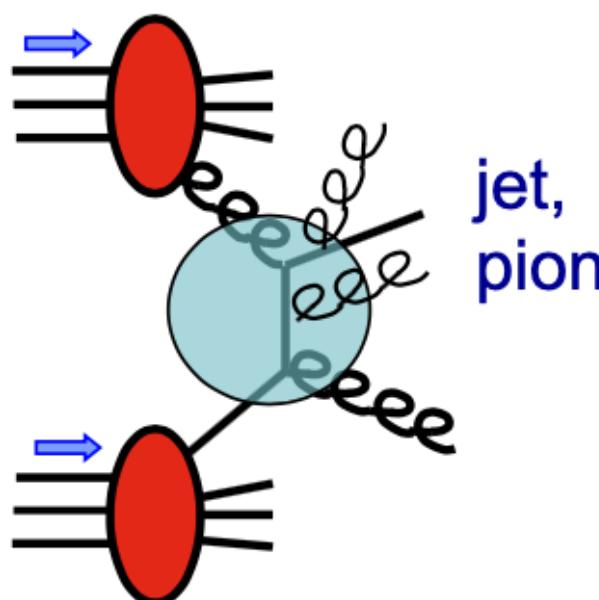
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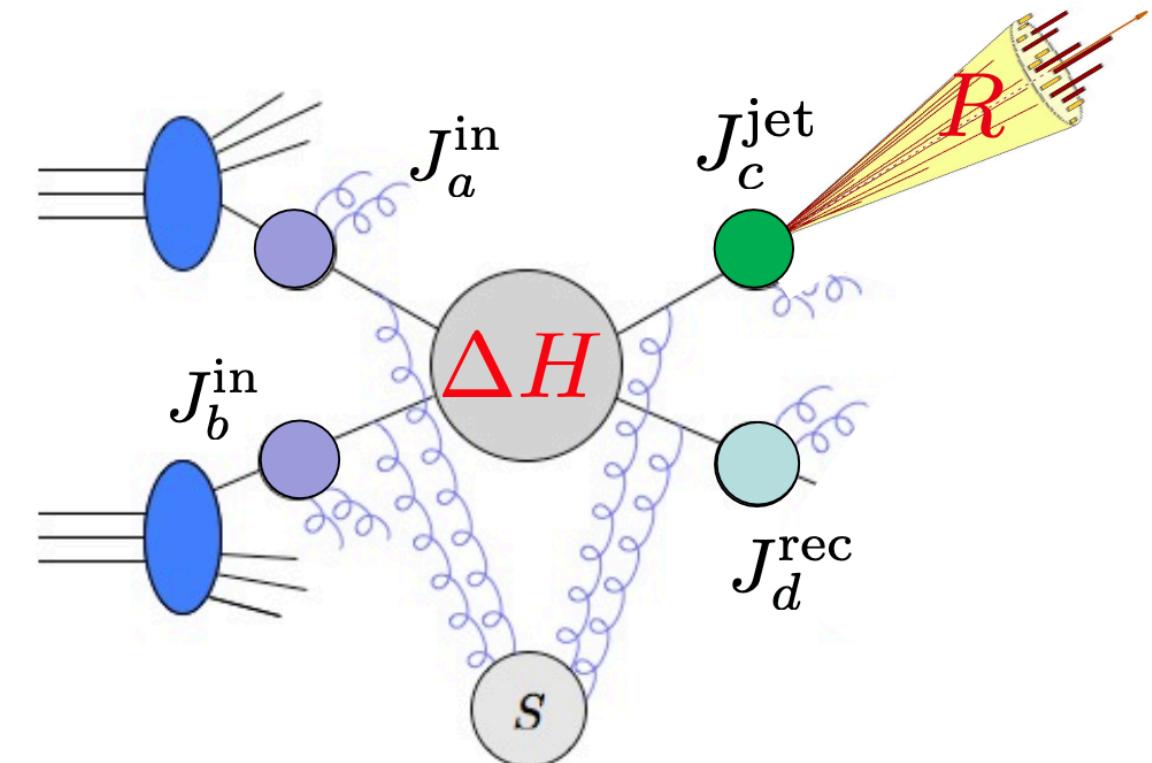


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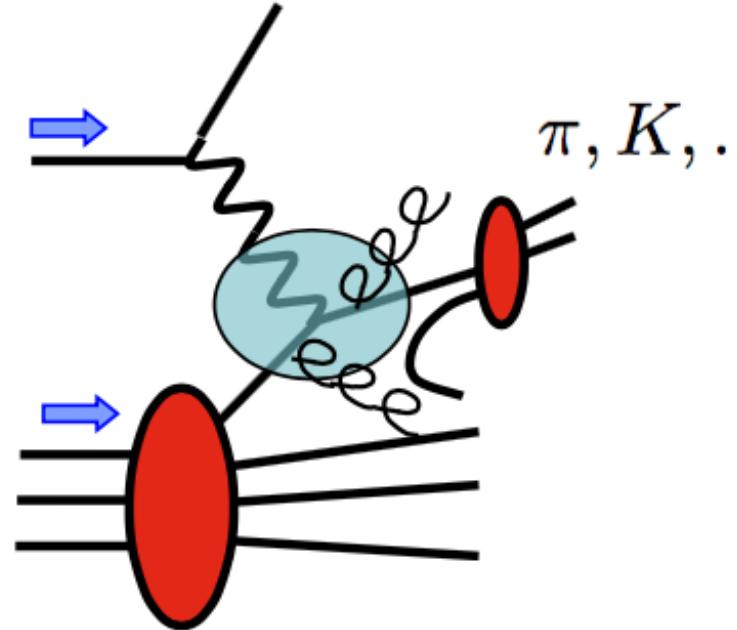
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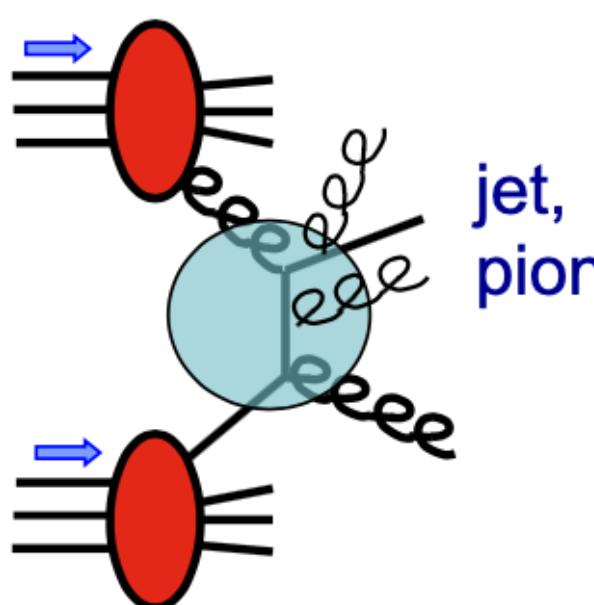
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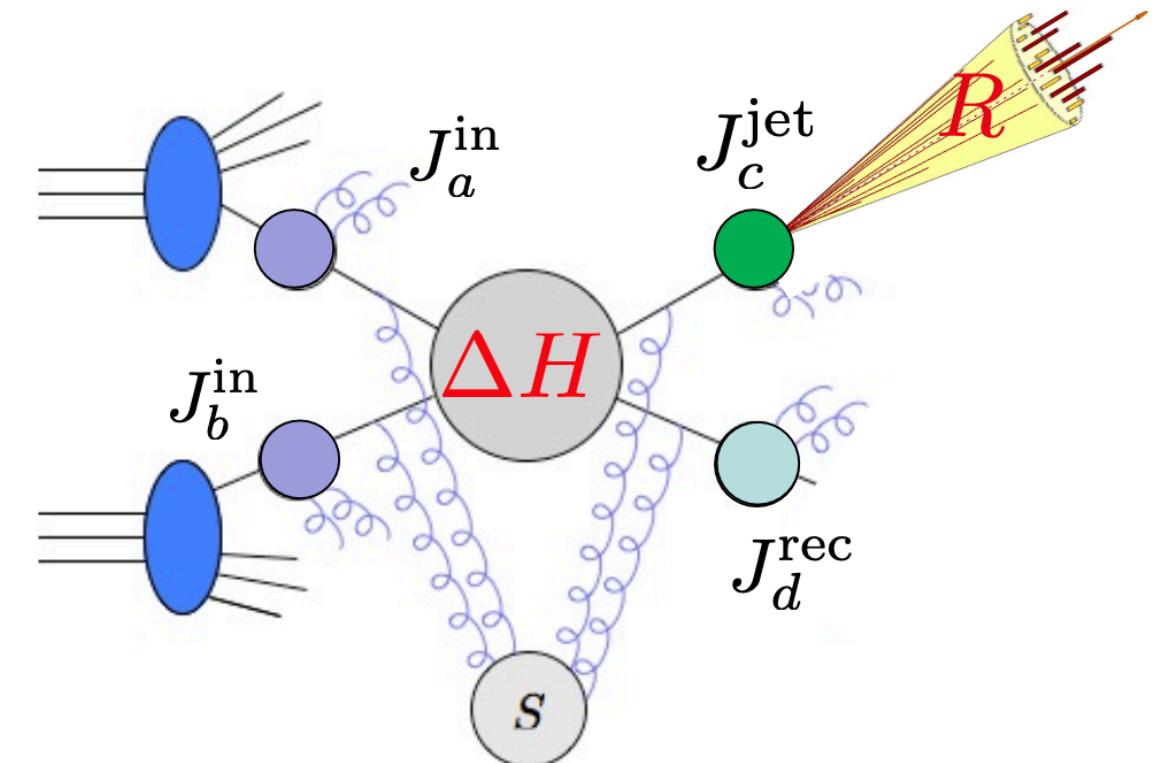
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D.deF,W.Vogelsang (in preparation)



- Another issue: need fragmentation functions, not yet fully global set available at NNLO (use NLO set)
- To be safe: select data in a restricted region of phase space (cuts in  $x_{\text{SIDIS}} > 0.12$  and  $p_T > 1.5 \text{ GeV}$ )

# Ingredients

## ► Parameterizations:

$$(\Delta q + \Delta \bar{q})(x, Q_0^2) = N_q x^{\alpha_q} (1 - x)^{\beta_q} \left( 1 + \gamma_q x^{\delta_q} + \eta_q x \right) \quad (u, d) \quad Q_0^2 = 1 \text{ GeV}^2$$

$$\Delta \bar{q}(x, Q_0^2) = N_{\bar{q}} x^{\alpha_{\bar{q}}} (1 - x)^{\beta_{\bar{q}}} \left( 1 + \gamma_{\bar{q}} x^{\delta_{\bar{q}}} \right) \quad (u, d, s)$$

$$\Delta g(x, Q_0^2) = N_g x^{\alpha_g} (1 - x)^{\beta_g} \left( 1 + \gamma_g x^{\delta_g} \right)$$

(32 parameters)

## ► Evolution: ZMVFNS HQ matching coefficients

QCD-PEGASUS framework  
A. Vogt (2004)

checked EKO and APFEL

## ► Assumptions: no $SU(2)/SU(3)$ constraints

positivity relative to MSHT20

S. Bailey et al.

## ► Fragmentation: BDSS20 BDSS24 NLO FFs jets/ $\pi^0/\pi^\pm$ (cuts in $x$ and $p_T$ ) I. Borsa et al. (2021,2023)

## ► Errors: Montecarlo Error Sampling Systematic Implementation

D. deF et al. (2019)

data replicas → 600 PDFs replicas

→ error propagation  
→ reweighting

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## ► Parameterizations:

$$(\Delta q + \Delta \bar{q})(x, Q_0^2) = N_q x^{\alpha_q} (1 - x)^{\beta_q} \left( 1 + \gamma_q x^{\delta_q} + \eta_q x \right) \quad (u, d) \quad Q_0^2 = 1 \text{ GeV}^2$$

$$\Delta \bar{q}(x, Q_0^2) = N_{\bar{q}} x^{\alpha_{\bar{q}}} (1 - x)^{\beta_{\bar{q}}} \left( 1 + \gamma_{\bar{q}} x^{\delta_{\bar{q}}} \right) \quad (u, d, s)$$

$$\Delta g(x, Q_0^2) = N_g x^{\alpha_g} (1 - x)^{\beta_g} \left( 1 + \gamma_g x^{\delta_g} \right)$$

(32 parameters)

## ► Evolution: ZMVFNS HQ matching coefficients

QCD-PEGASUS framework  
A. Vogt (2004)

checked EKO and APFEL

## ► Assumptions: no $SU(2)/SU(3)$ constraints

positivity relative to MSHT20

S. Bailey et al.

## ► Fragmentation: BDSS20 BDSS24 NLO FFs jets/ $\pi^0/\pi^\pm$ (cuts in $x$ and $p_T$ ) I. Borsa et al. (2021,2023)

## ► Errors: Montecarlo Error Sampling Systematic Implementation

data replicas  $\rightarrow$  600 PDFs replicas

D. deF et al. (2019)

→ error propagation  
→ reweighting



# Ingredients and Results

## Data Selection:

**DIS:** EMC, SMC, E142, E143, E154, E155,  
HERMES, COMPASS, HALL-A, CLAS  
(p, n, d, He)

**SIDIS:** SMC, HERMES, COMPASS  
 $(p, d) \rightarrow (\pi^\pm, K^\pm, h^\pm)$

**PP-JETS:** STAR run 5, 6, 9, 12, 13, 15  
( $\sqrt{s} = 200, 510 \text{ GeV}$ )

**PP- $\pi^0/\pi^\pm$ :** PHENIX, STAR

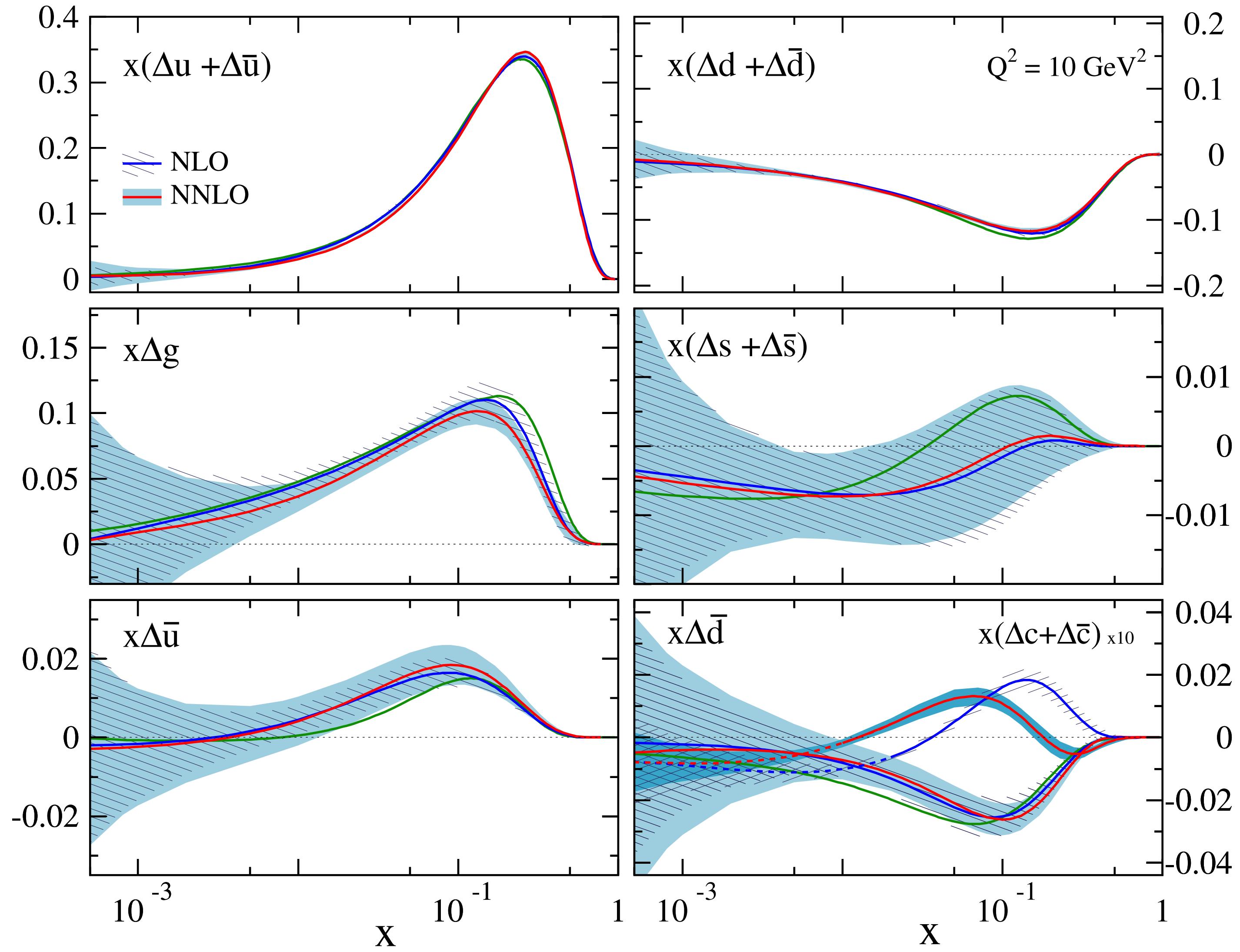
**PP  $W^\pm$ :** PHENIX, STAR

**Total:**

#data	NLO	NNLO
368	302.7	294.3
114	127.6	122.9
91	111.1	104.7
78	63.5	66.0
22	22.3	20.3
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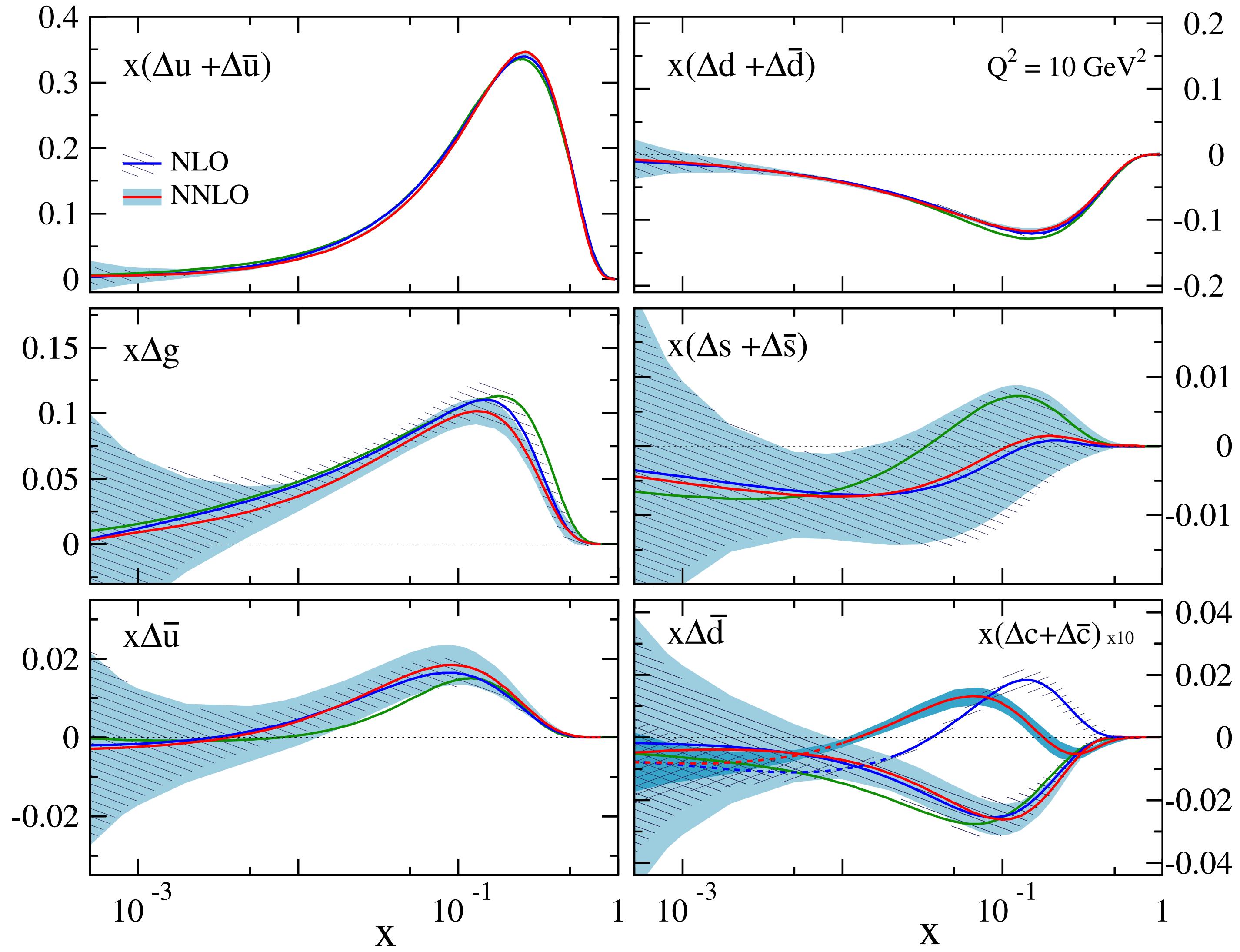
$$x_{SIDIS} > 0.12 \quad p_T > 1.5 \text{ GeV}$$

# Distributions



BDSSV22: NLO with dijets and no cuts on sidis  
( $\Delta u + \Delta \bar{u}$ ) and ( $\Delta d + \Delta \bar{d}$ ) well constrained  
no significant NNLO/NLO differences

# Distributions



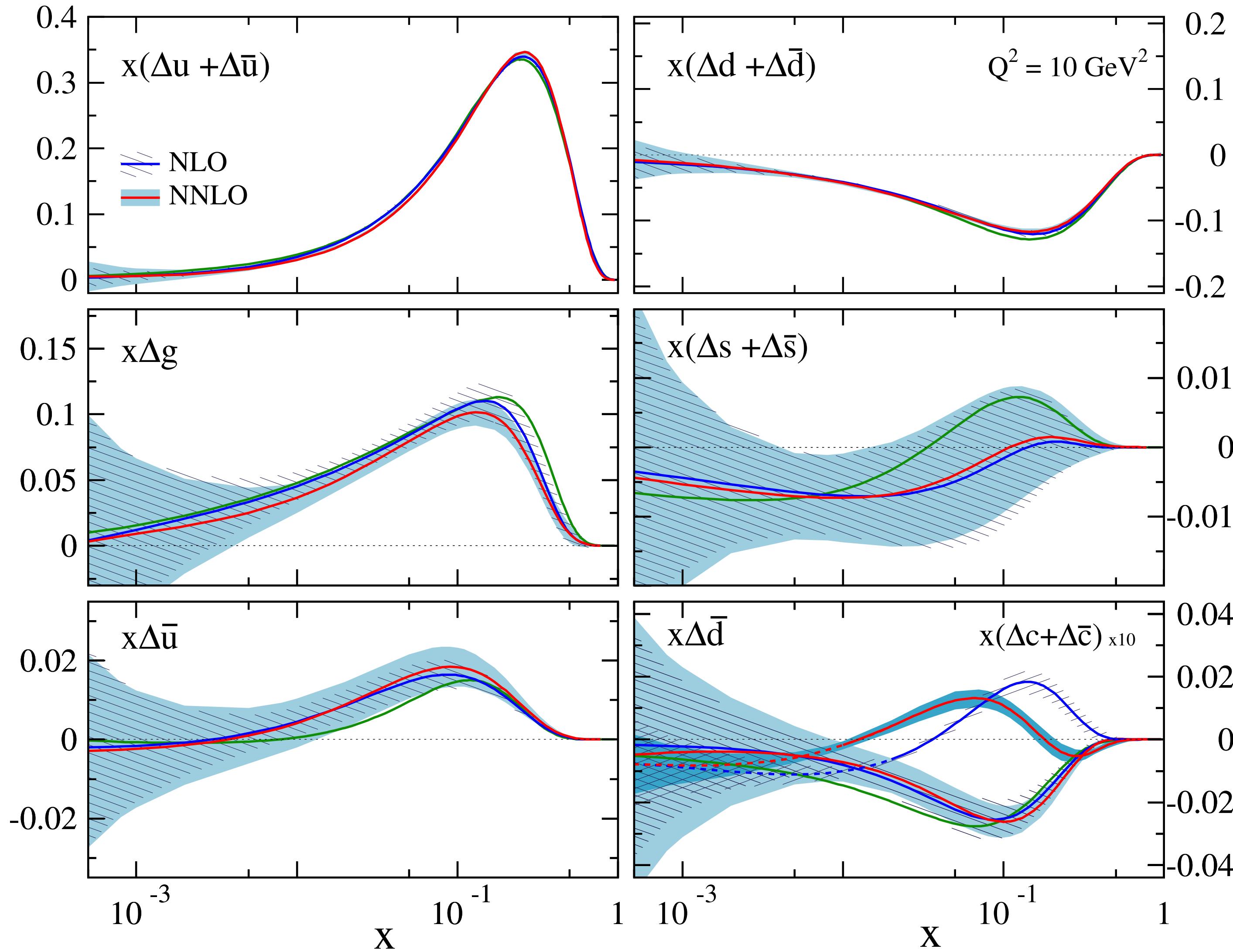
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NNLO/NLO differences within uncertainties

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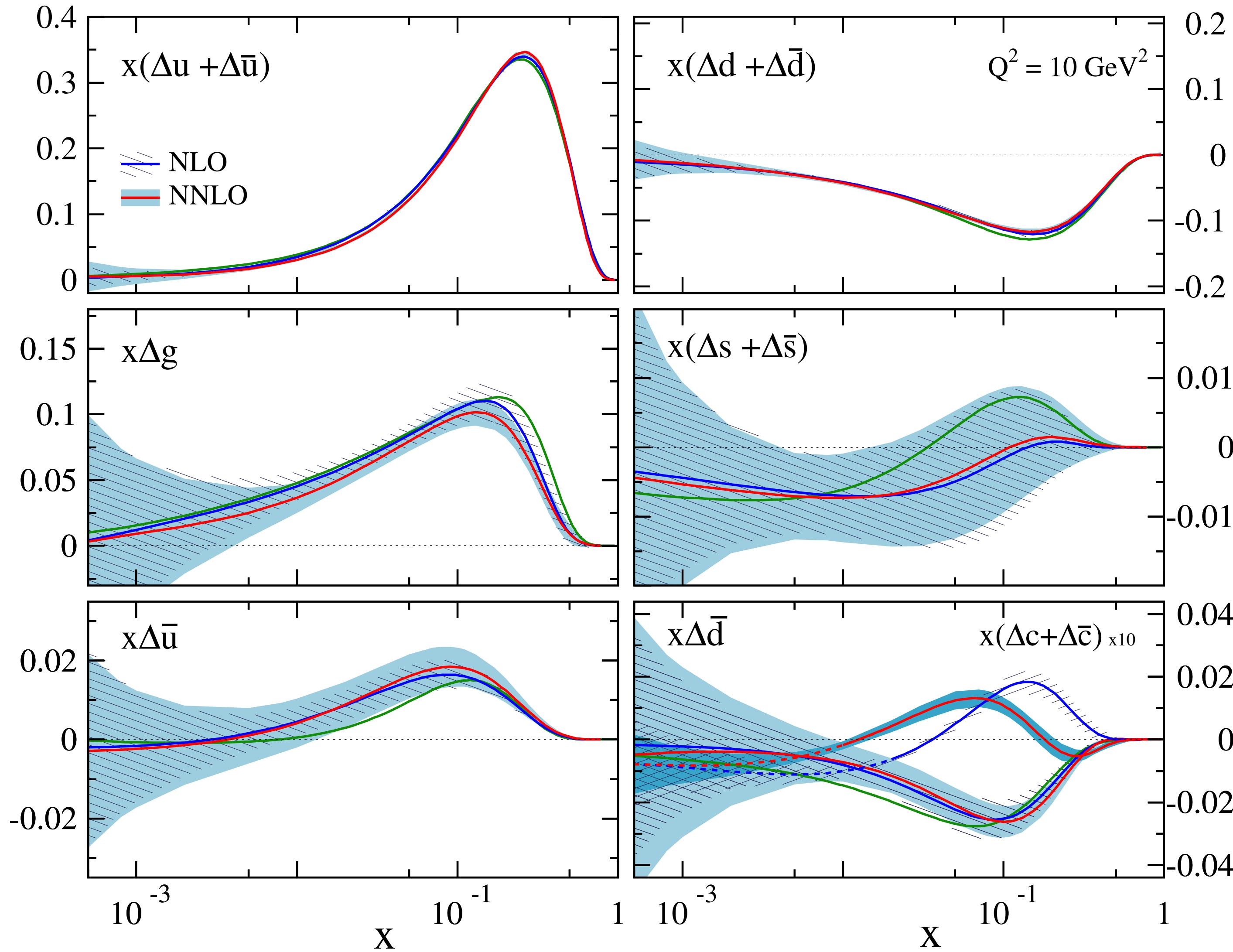
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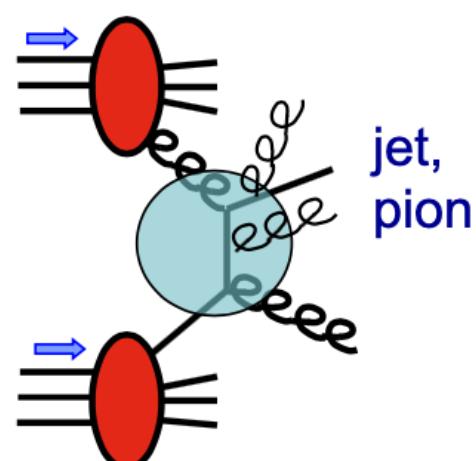
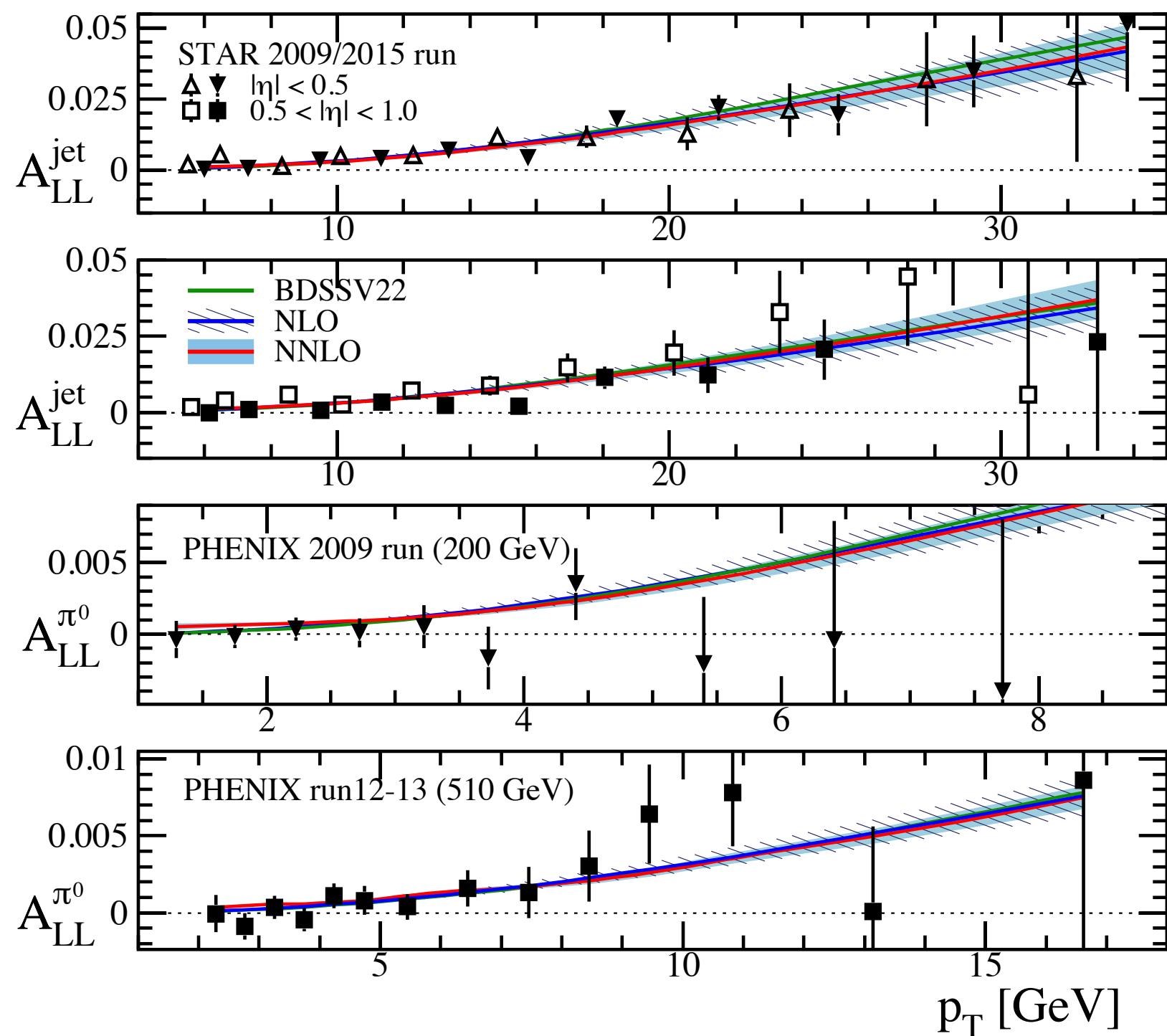
$(\Delta s + \Delta \bar{s})$  consistent with zero  
suffers the cut and lack of  $F, D$  constraints

$\Delta \bar{u}$  and  $\Delta \bar{d}$  opposite signs  
constrained by  $W$

$(\Delta c + \Delta \bar{c})$  small!  
inherits the gluon uncertainties

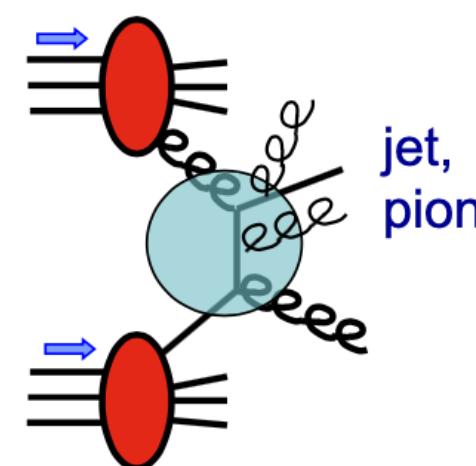
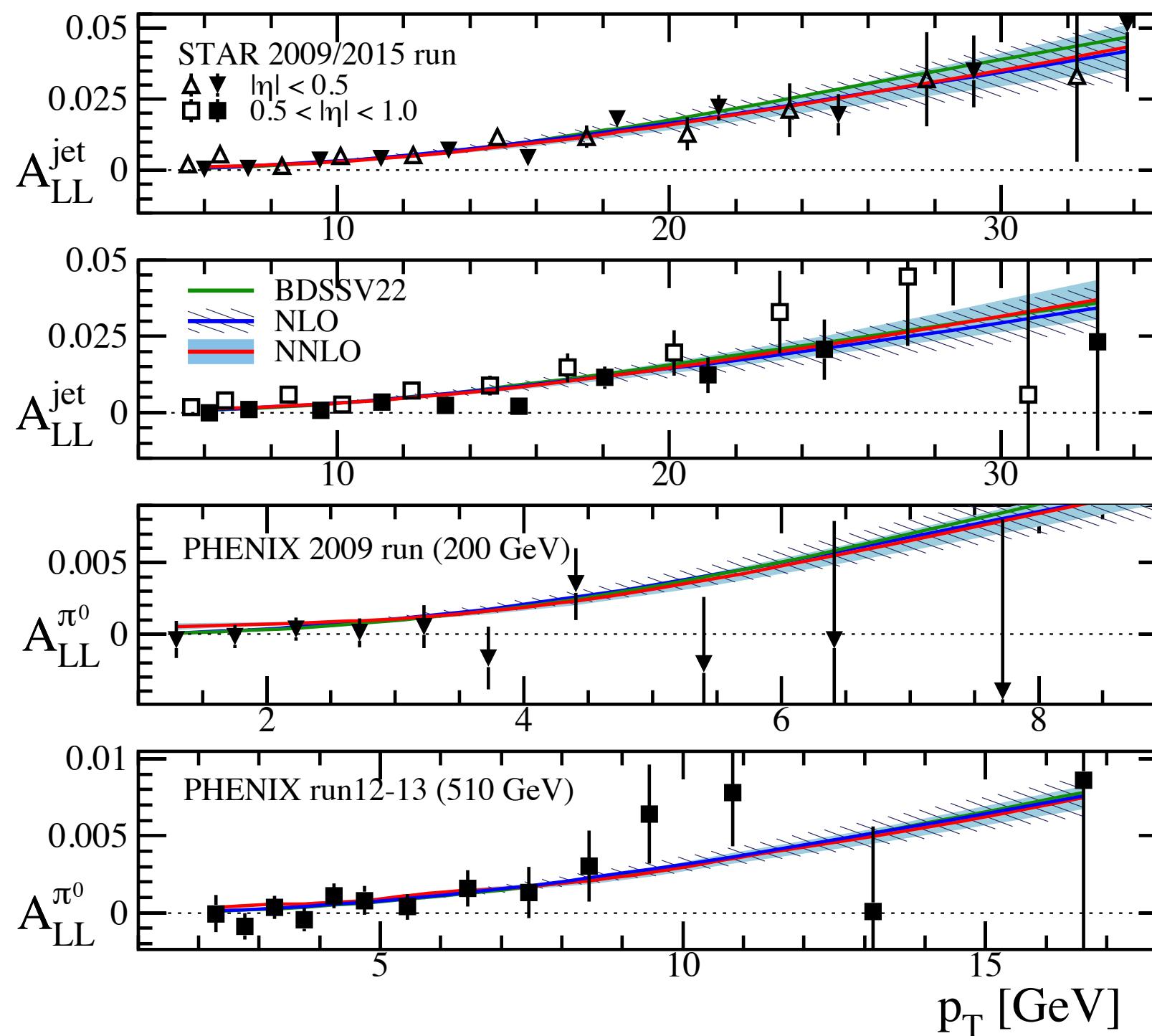
strongly dependent on perturbative order

# Selected Data Sets

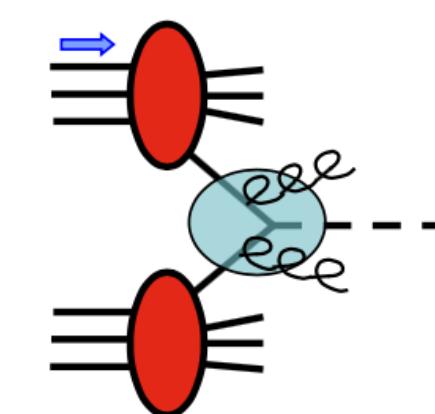
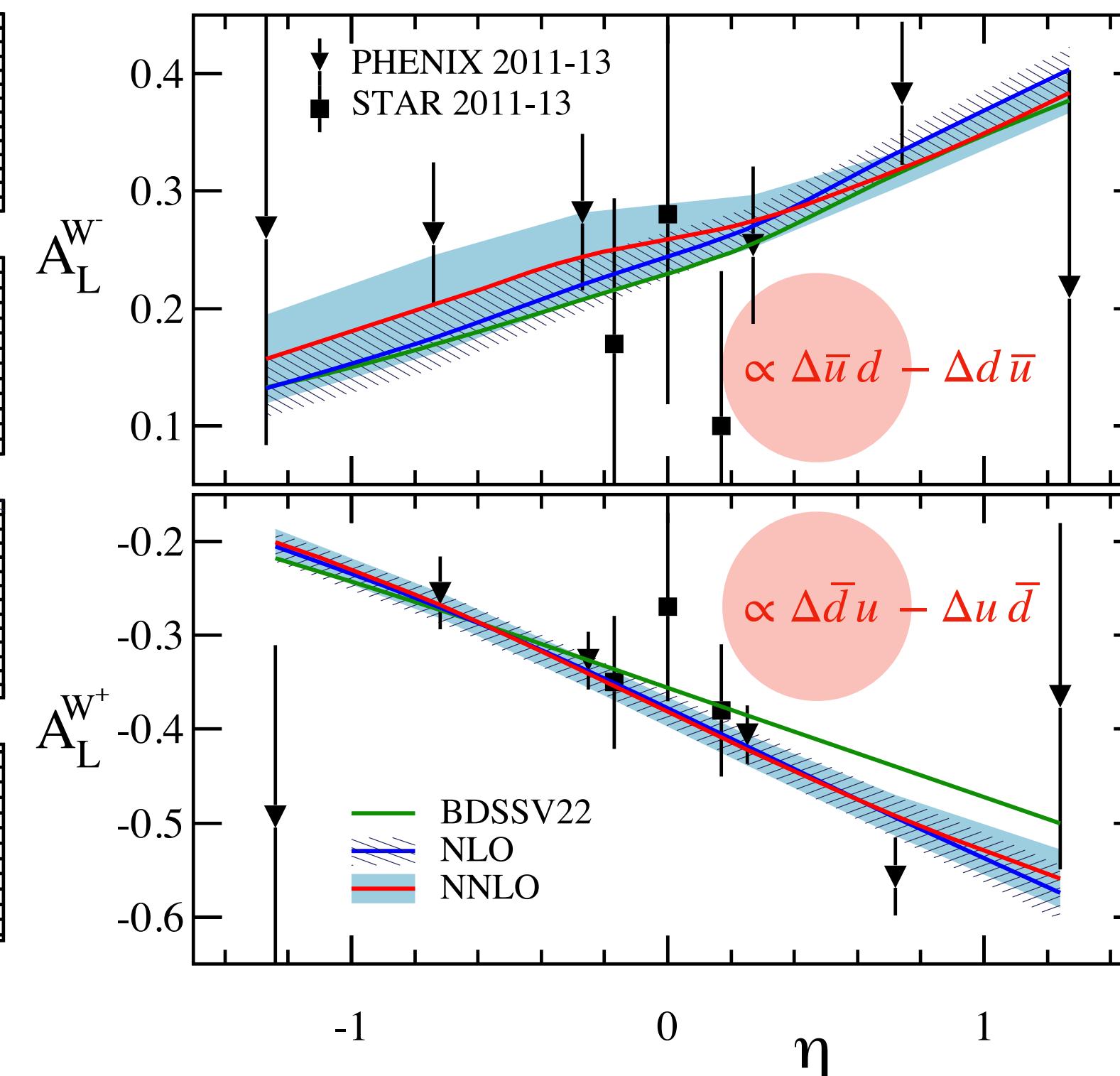


pp high- $p_T$

# Selected Data Sets

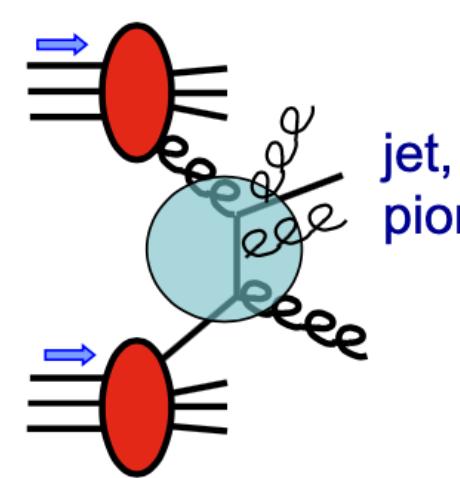
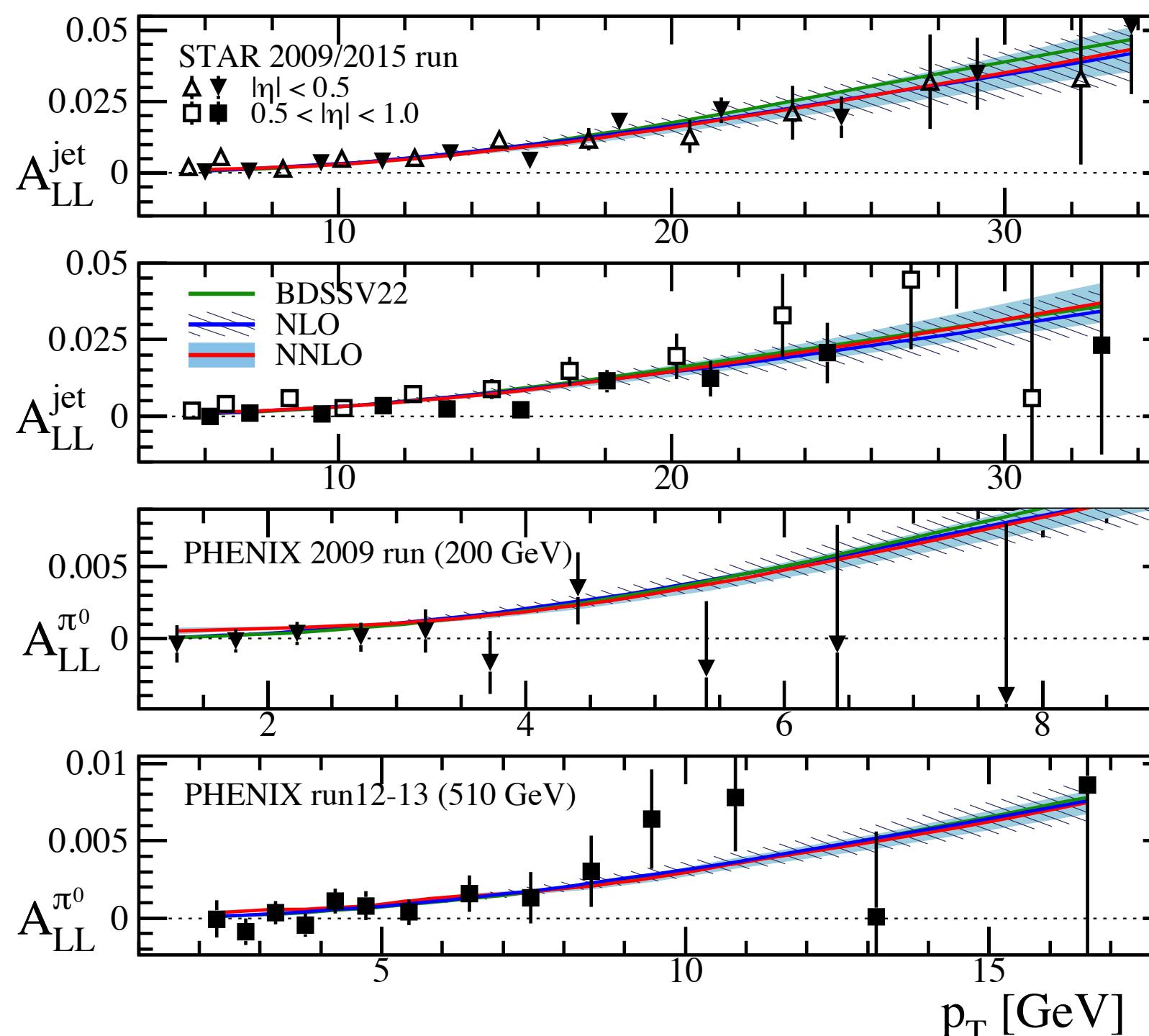


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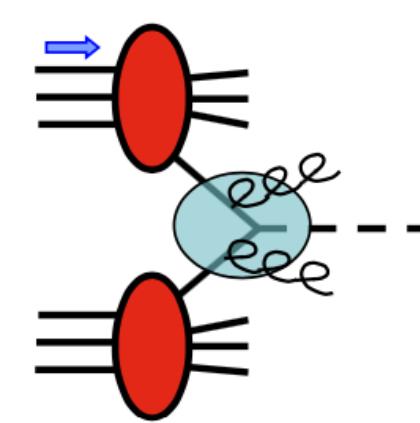
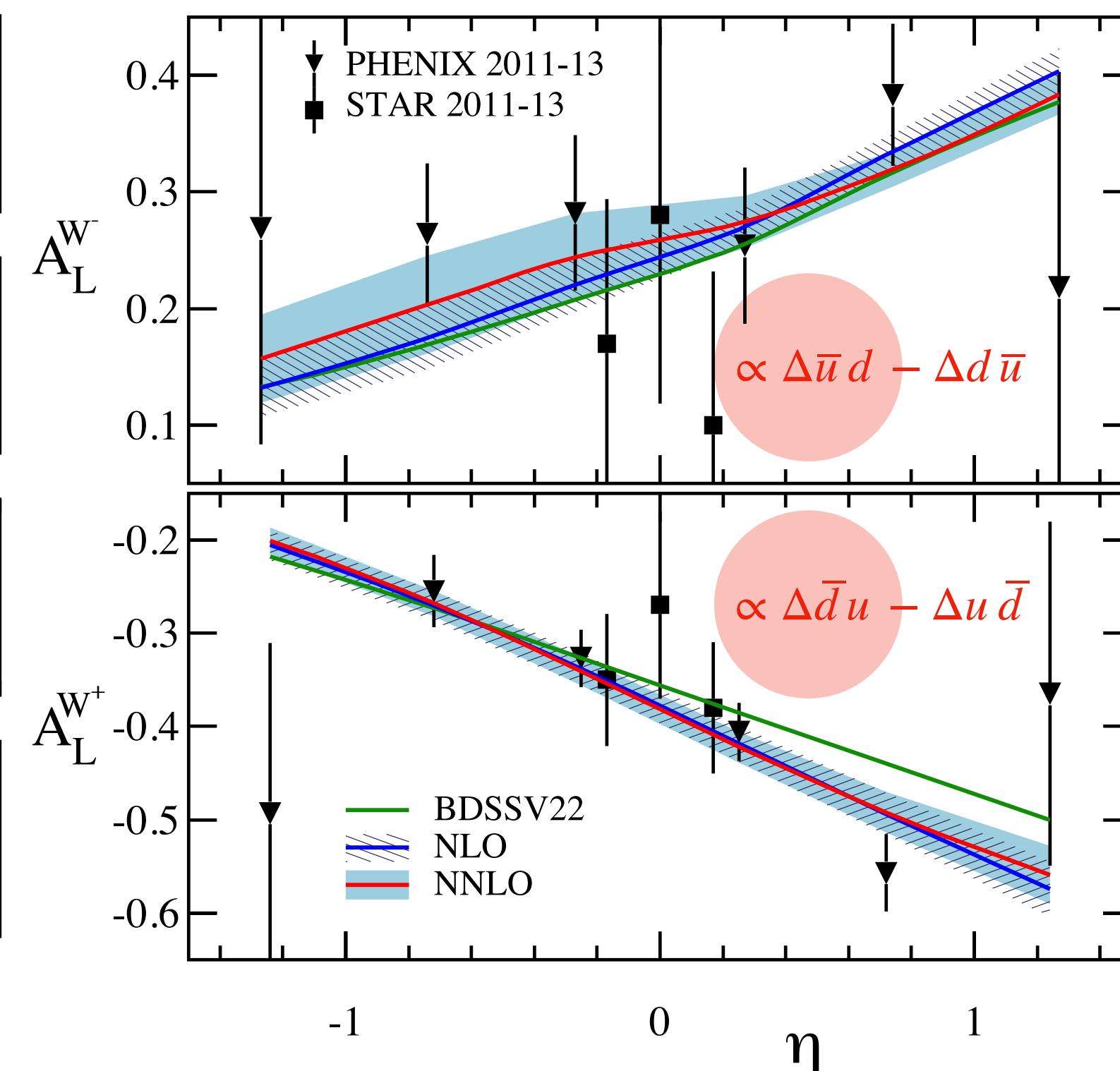


$W$  bosons

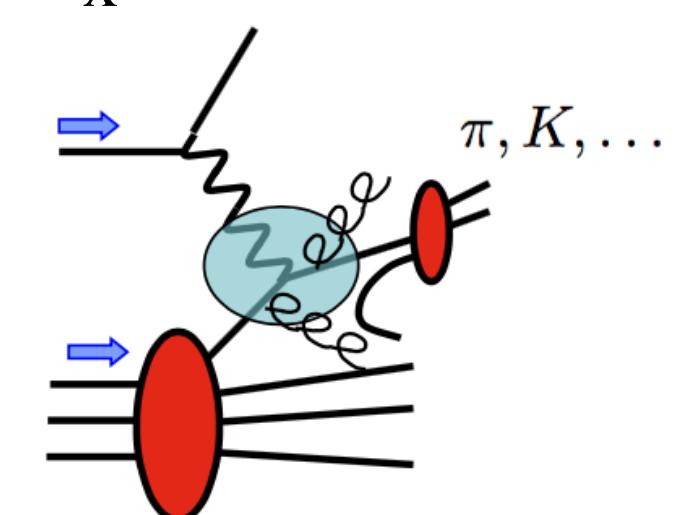
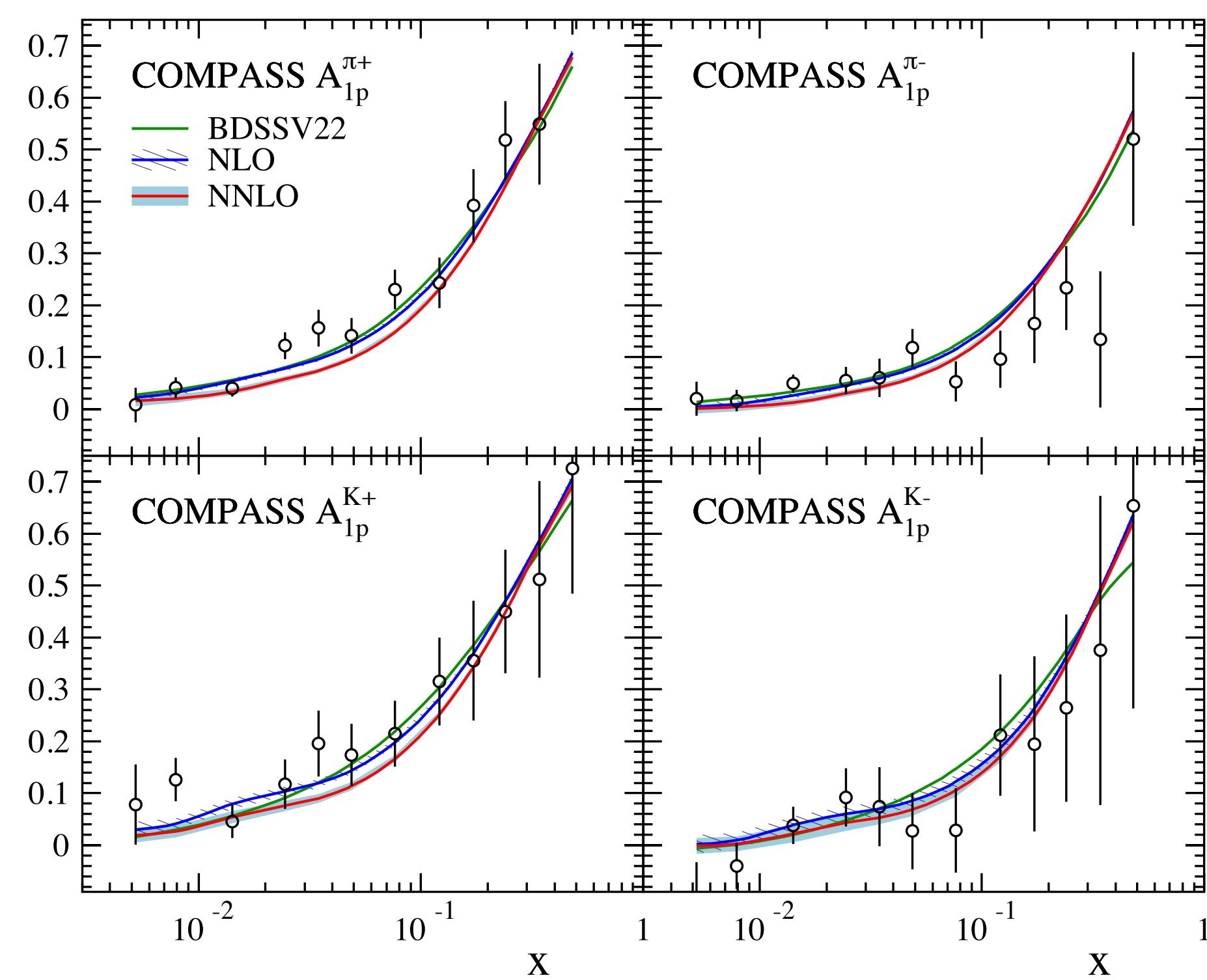
# Selected Data Sets



$\text{pp}$  high- $p_T$

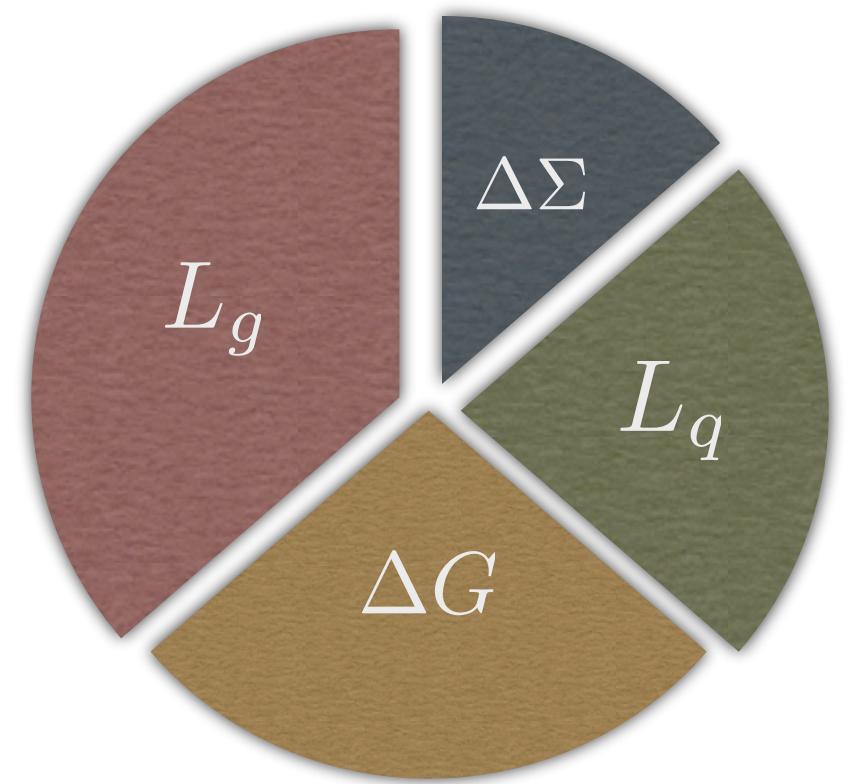


$W$  bosons



SIDIS

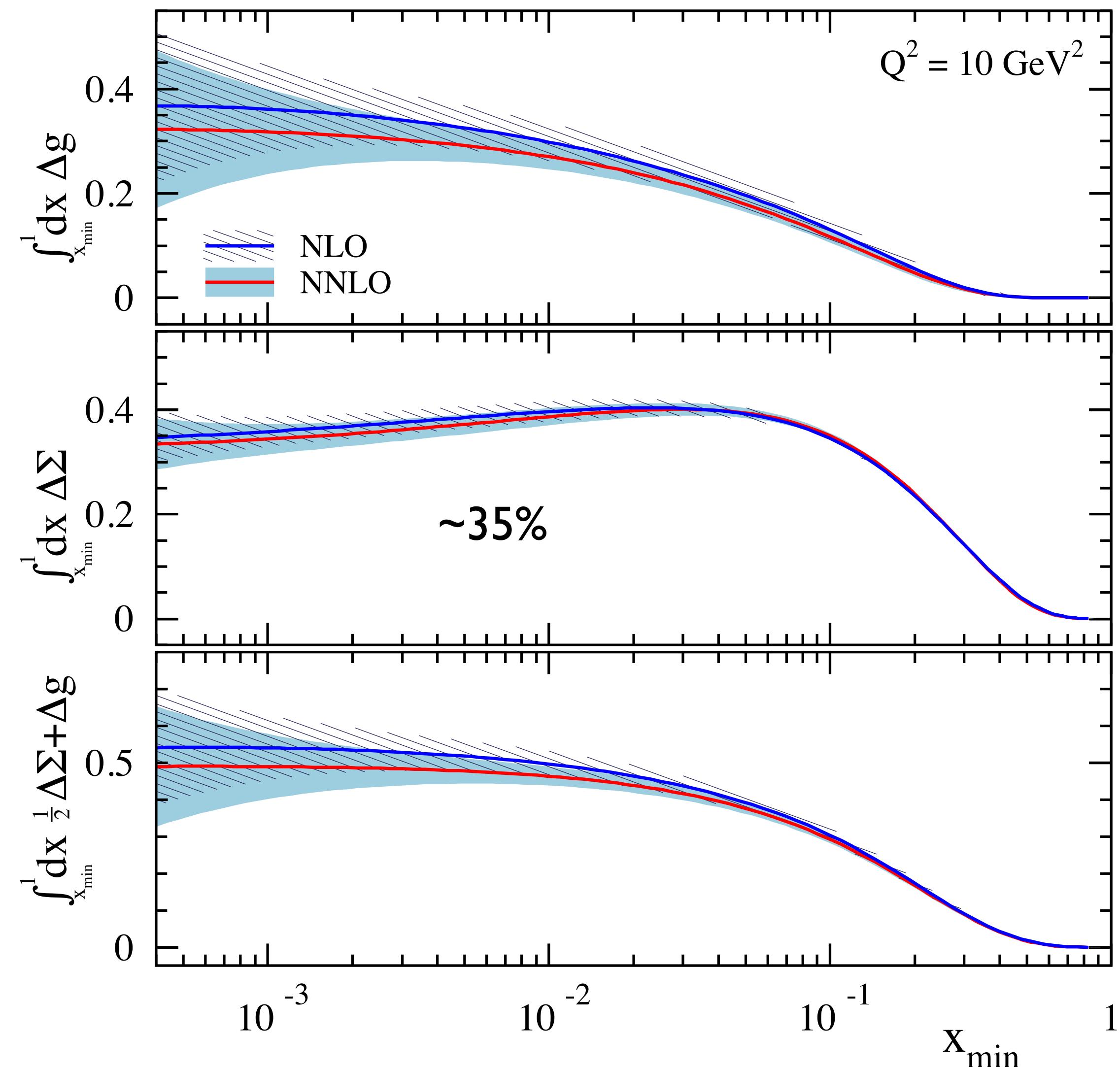
# Spin Sum Rule



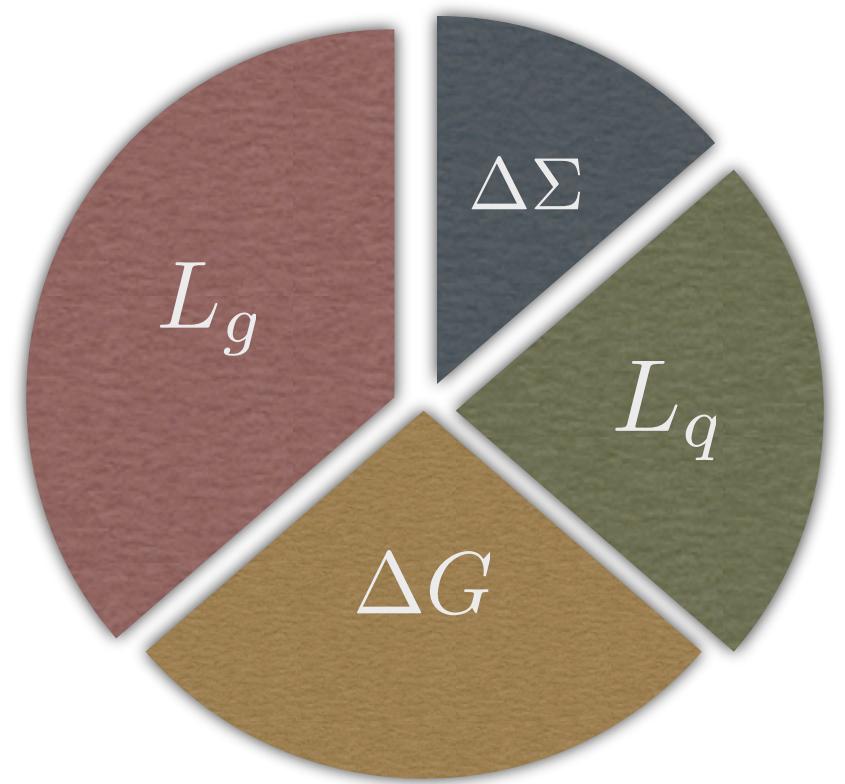
$$\frac{1}{2}\Delta\Sigma + \Delta G + L_q + L_g = \frac{1}{2}$$

Extrapolation at small  $x$  very uncertain...

but might saturate the spin sum rule just from spin of quarks and gluons (no orbital angular momentum)



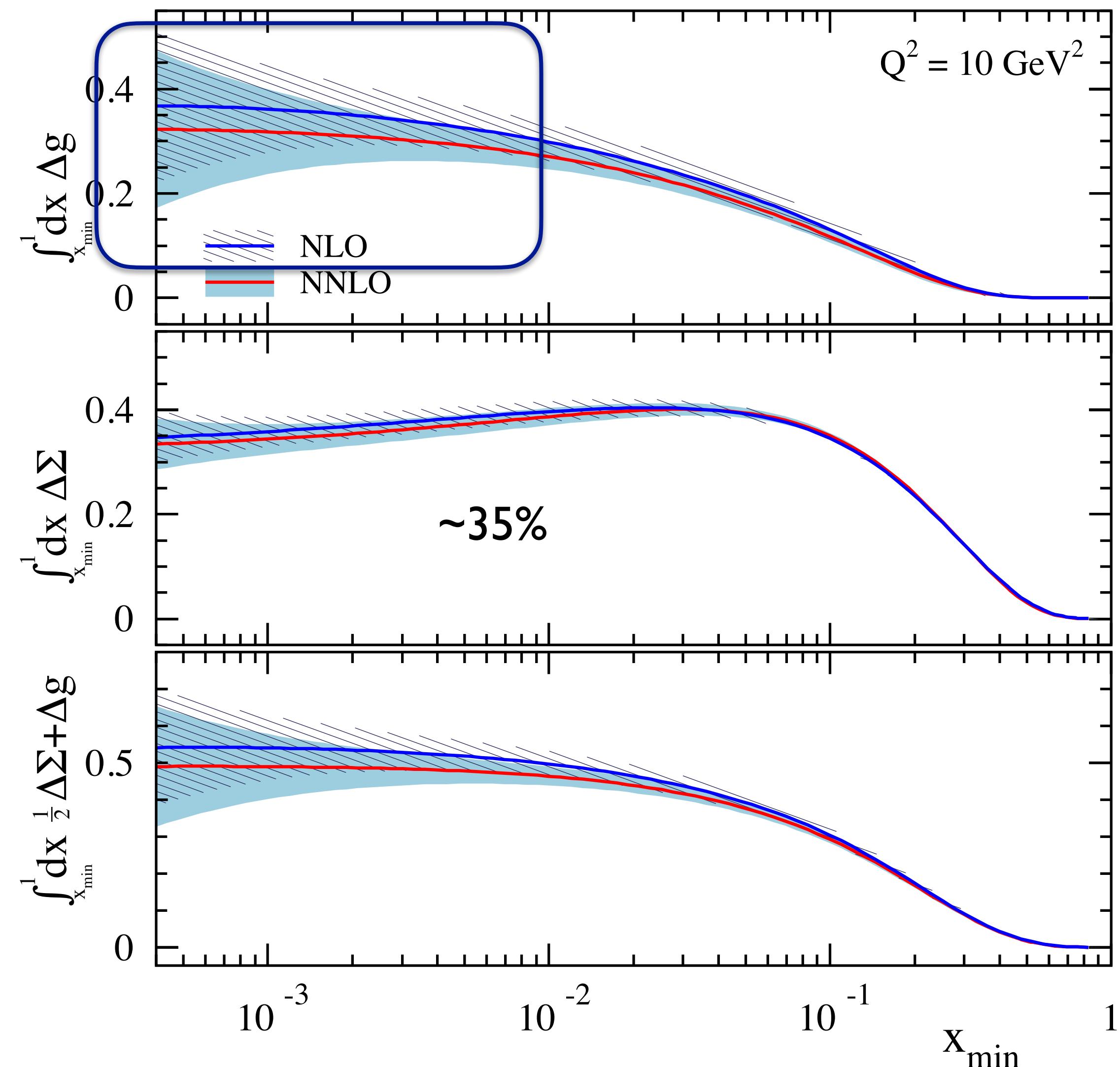
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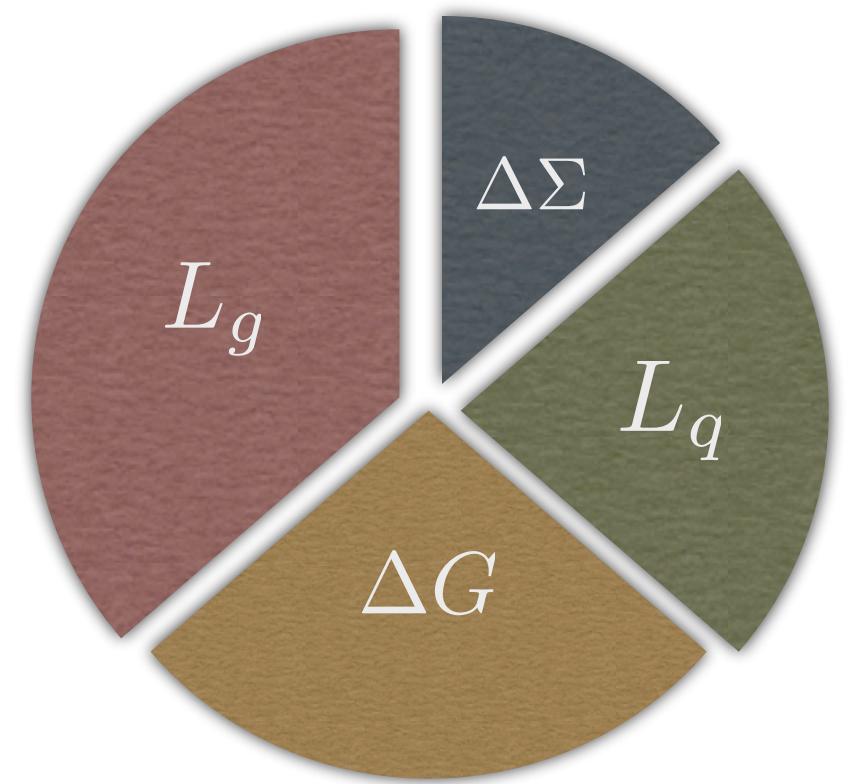
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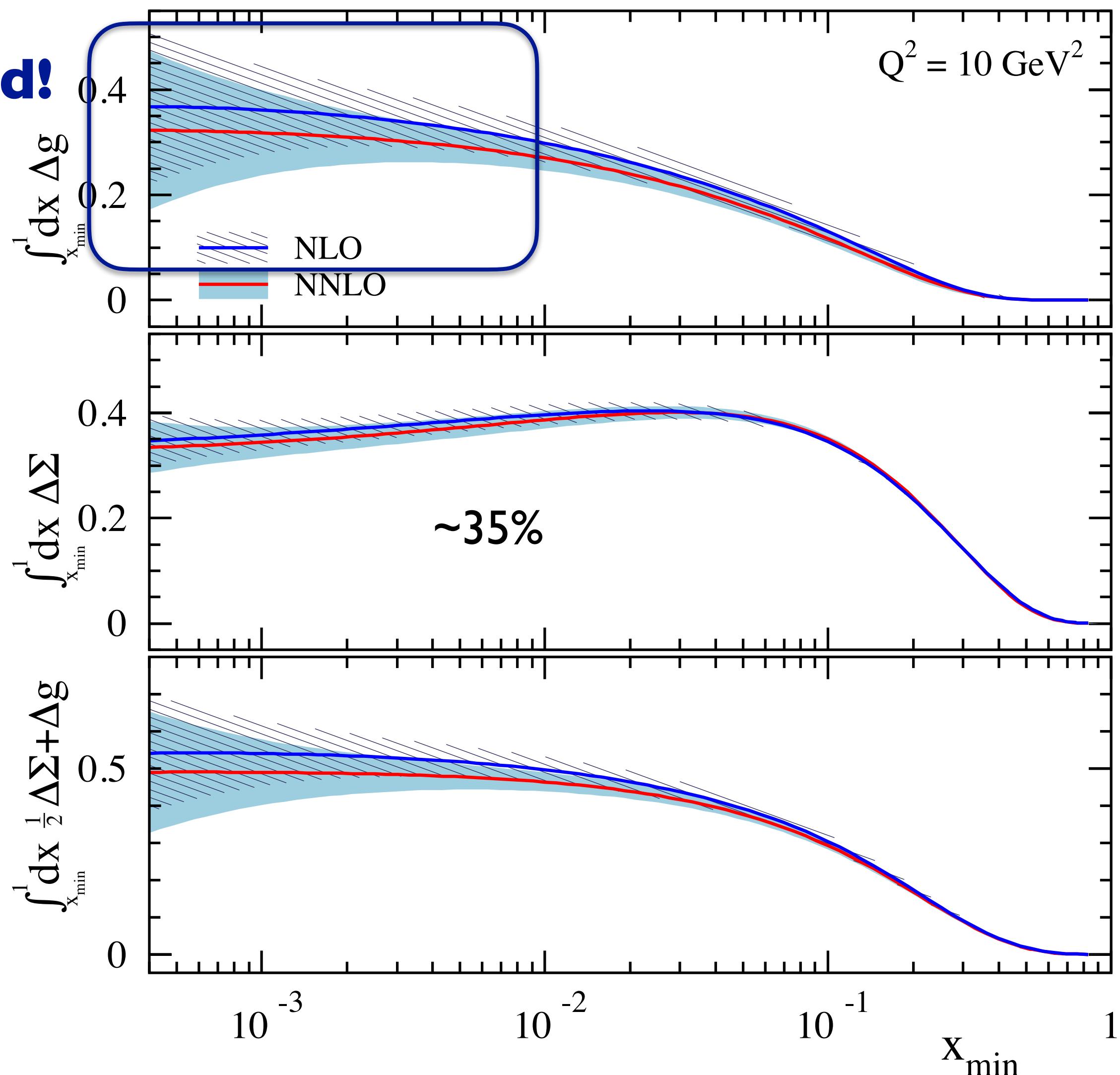


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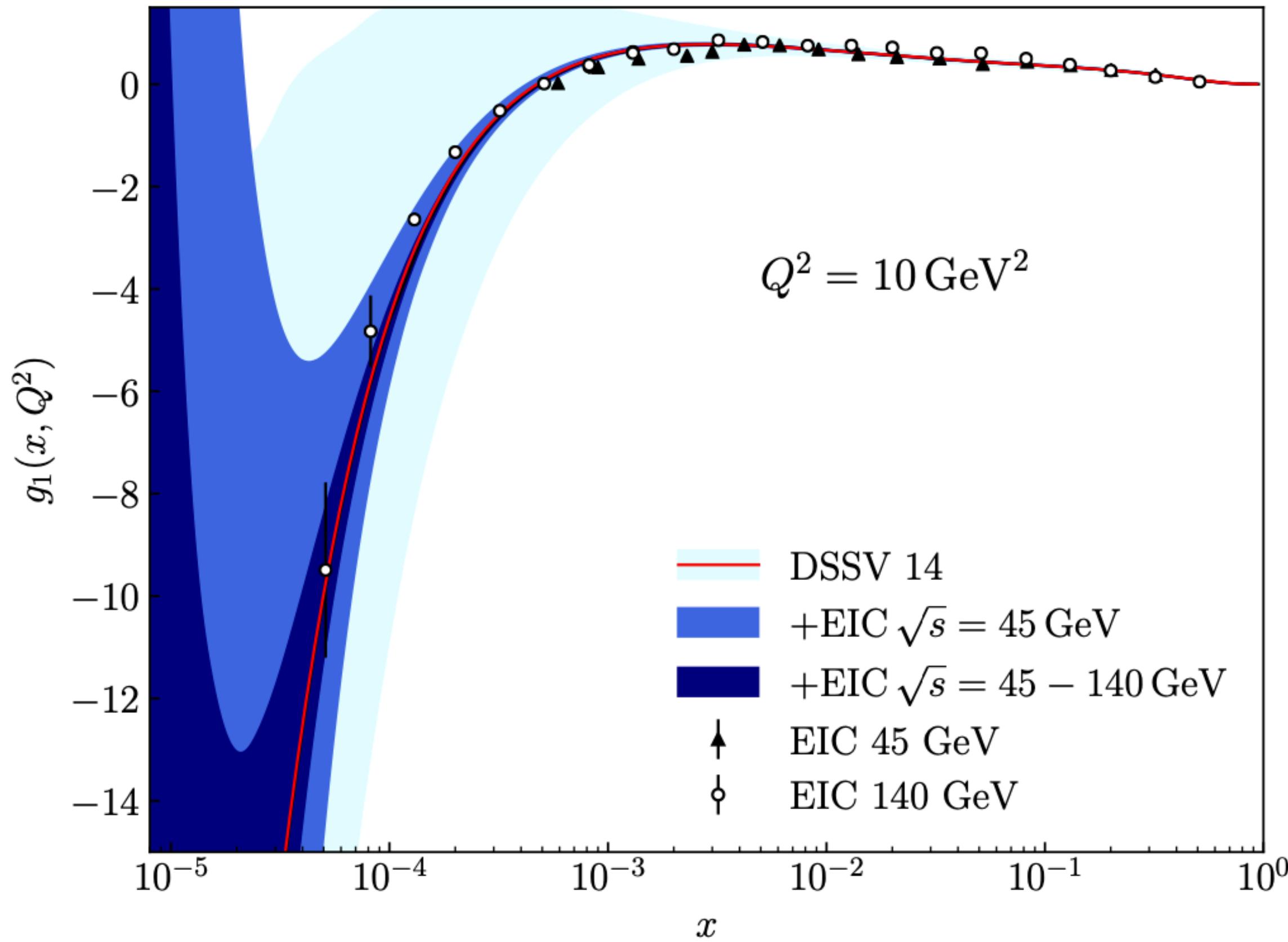
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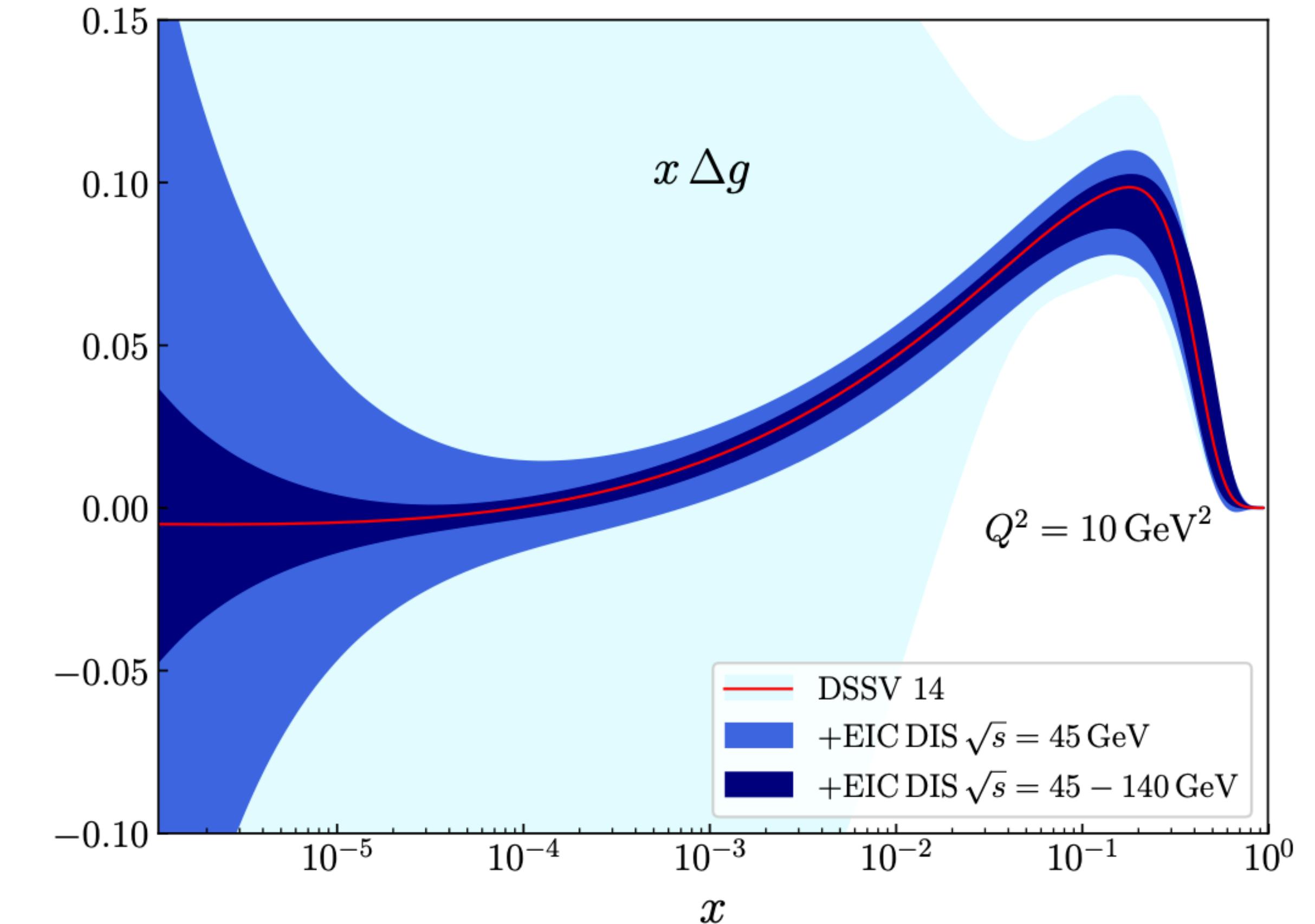
**EIC  
needed!**



# Key measurement for polarized gluon distribution : $g_1(x, Q^2)$ at small $x$ at EIC (coefficient + evolution)



at small  $x$   $\frac{\partial g_1(x, Q^2)}{\partial \ln Q^2} \approx -\Delta g(x, Q^2)$



Aschenauer, Borsa, Lucero, Nunes, Sassot (2020)

# QCD+QED/EW effects

Notice that  $\alpha_s^2 \sim \alpha \sim 0.01$  and EW effects can be enhanced..

$$\sigma(\alpha_s, \alpha) = \sigma^{(0,0)} + \alpha_s \sigma^{(1,0)} + \alpha_s^2 \sigma^{(2,0)} + \dots + \alpha \sigma^{(0,1)} + \dots + \alpha_s \alpha \sigma^{(1,1)} + \dots$$

LO	NLO QCD	NNLO QCD	NLO EW	NNLO mixed QCD-EW
----	---------	----------	--------	-------------------

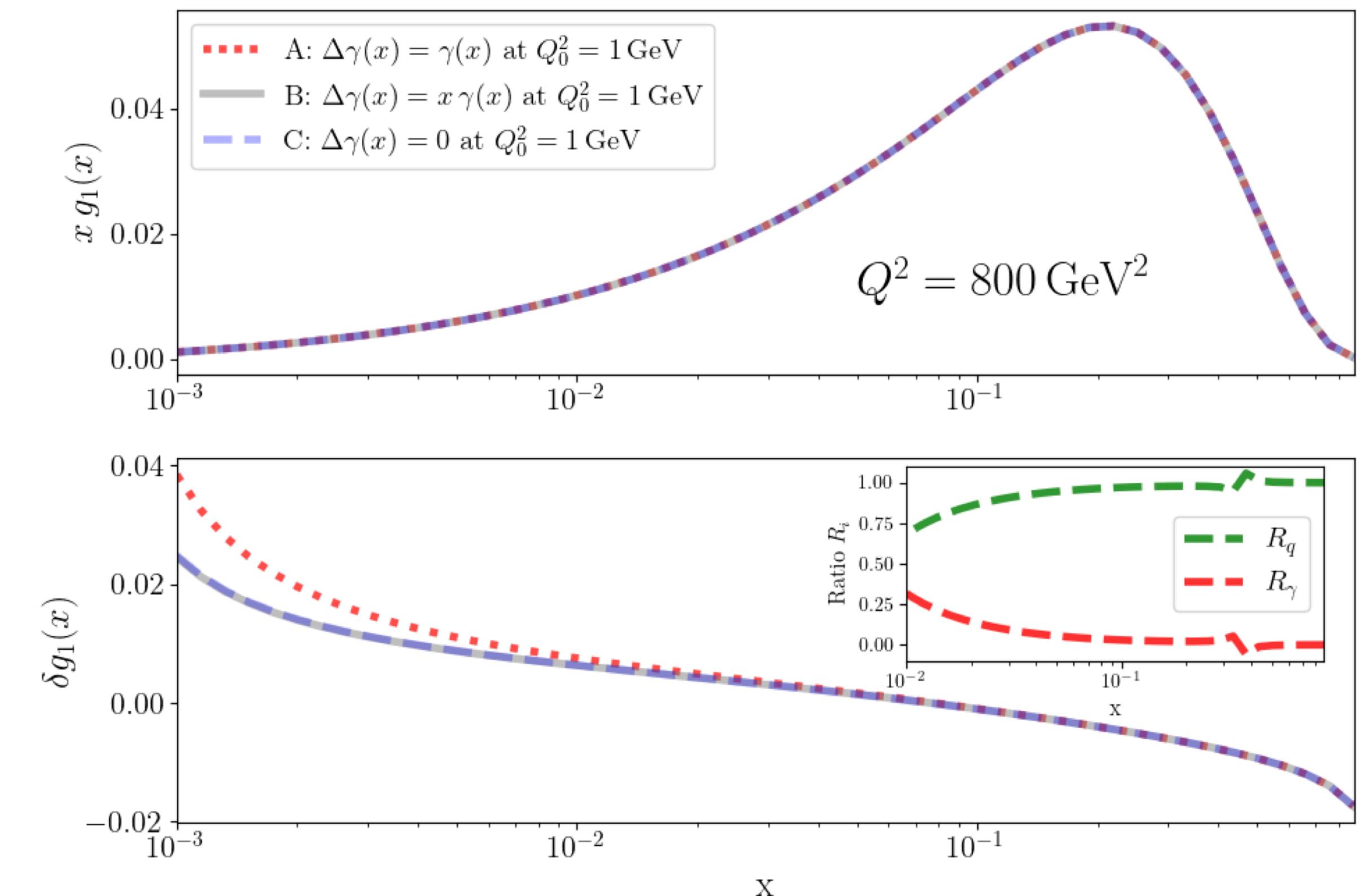
- mixed QCD-QED splitting functions known

unpolarized	<a href="#">deF. Rodrigo, Sborlini (2016)</a>
polarized	<a href="#">deF., Palma (2023)</a>

- QED corrections to  $g_1(x, Q^2)$  [deF., Palma \(2023\)](#)

- Corrections typically  $\sim 1\%$  but can be enhanced in certain kinematical regions

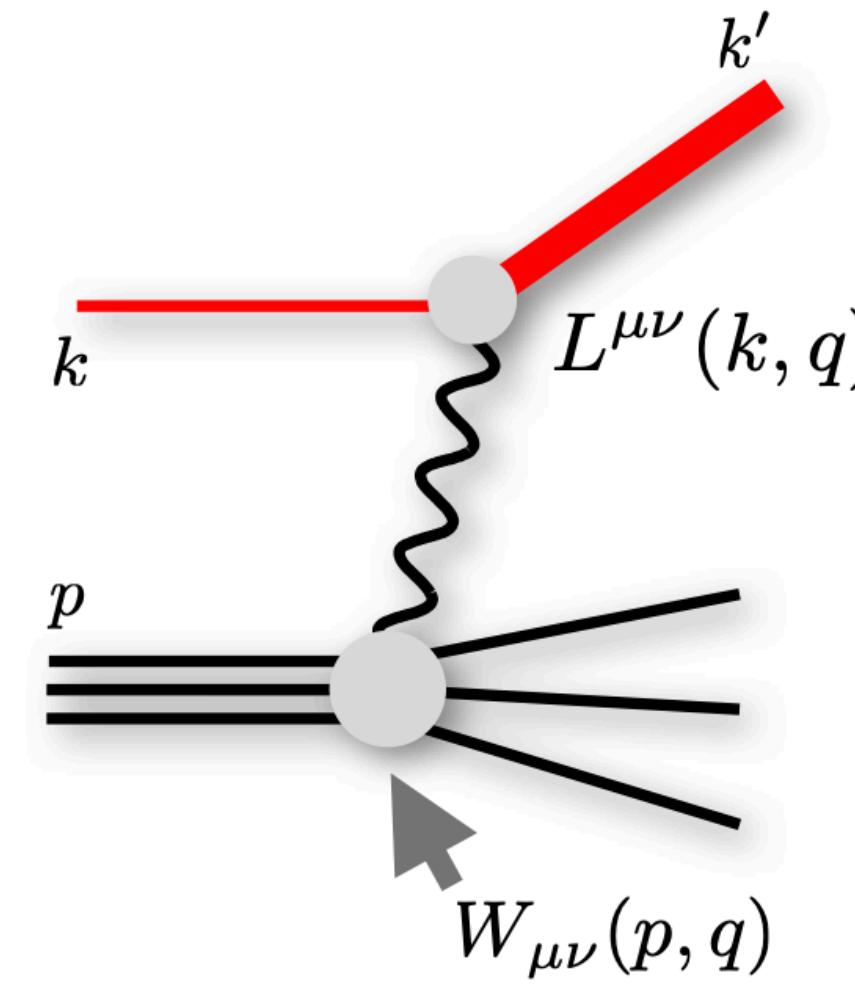
new distributions: photon and lepton in the proton



# photon distribution :LUXQED

“Hypothetical DIS process”  $l + p \rightarrow L + X$  calculated in two ways Manohar, Nason, Salam, Zanderighi (2016)

I. total cross section in terms of structure functions

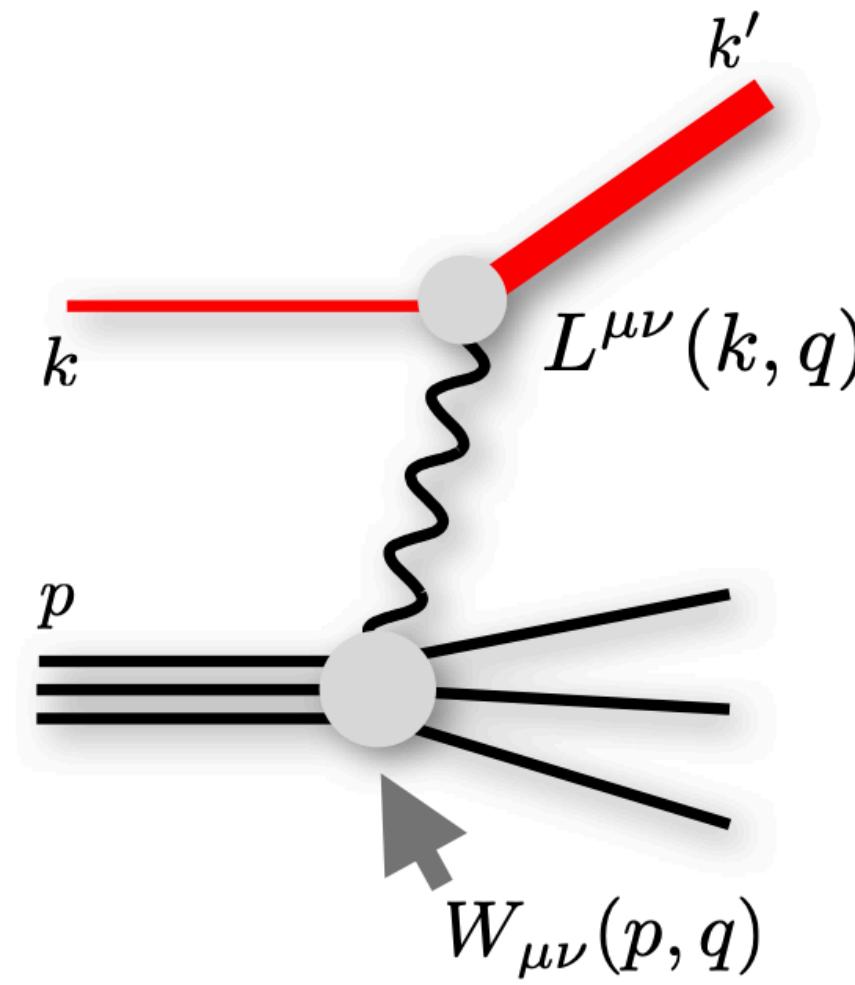


$$\Delta\sigma_{lp}(p) = \frac{1}{2\pi\alpha(\mu^2)}\sigma_0 \int_x^1 \frac{dz}{z} \int_{Q_{\min}^2}^{Q_{\max}^2} \frac{dQ^2}{Q^2} \alpha_{\text{ph}}^2(-Q^2)$$
$$\left\{ H \left( 4 - 2z - \frac{4m_p^2 x^2}{Q^2} - \frac{4m_p^2 x^2 Q^2}{M^4} - \frac{8m_p^2 x^2}{M^2} - \frac{2zQ^2}{M^2} \right) x g_1(x/z, Q^2) \right.$$
$$\left. - H \left( \frac{8m_p^2 x^2}{zM^2} + \frac{8m_p^2 x^2}{zQ^2} \right) x g_2(x/z, Q^2) \right\}$$

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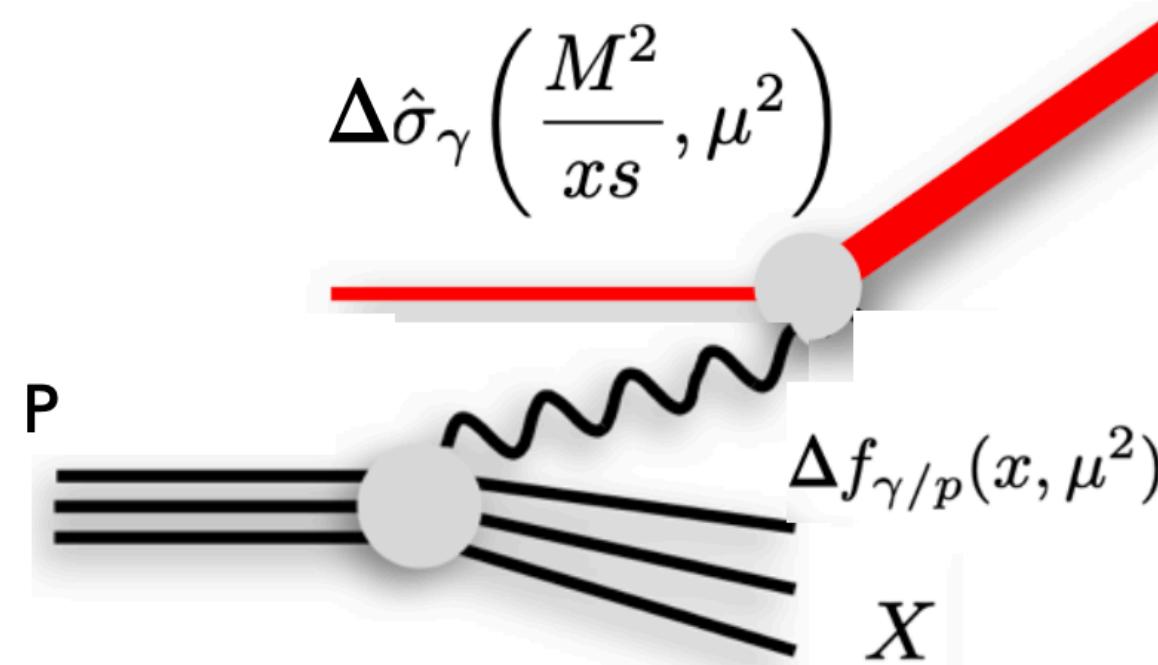


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2. from collinear factorization in terms of photon distribution



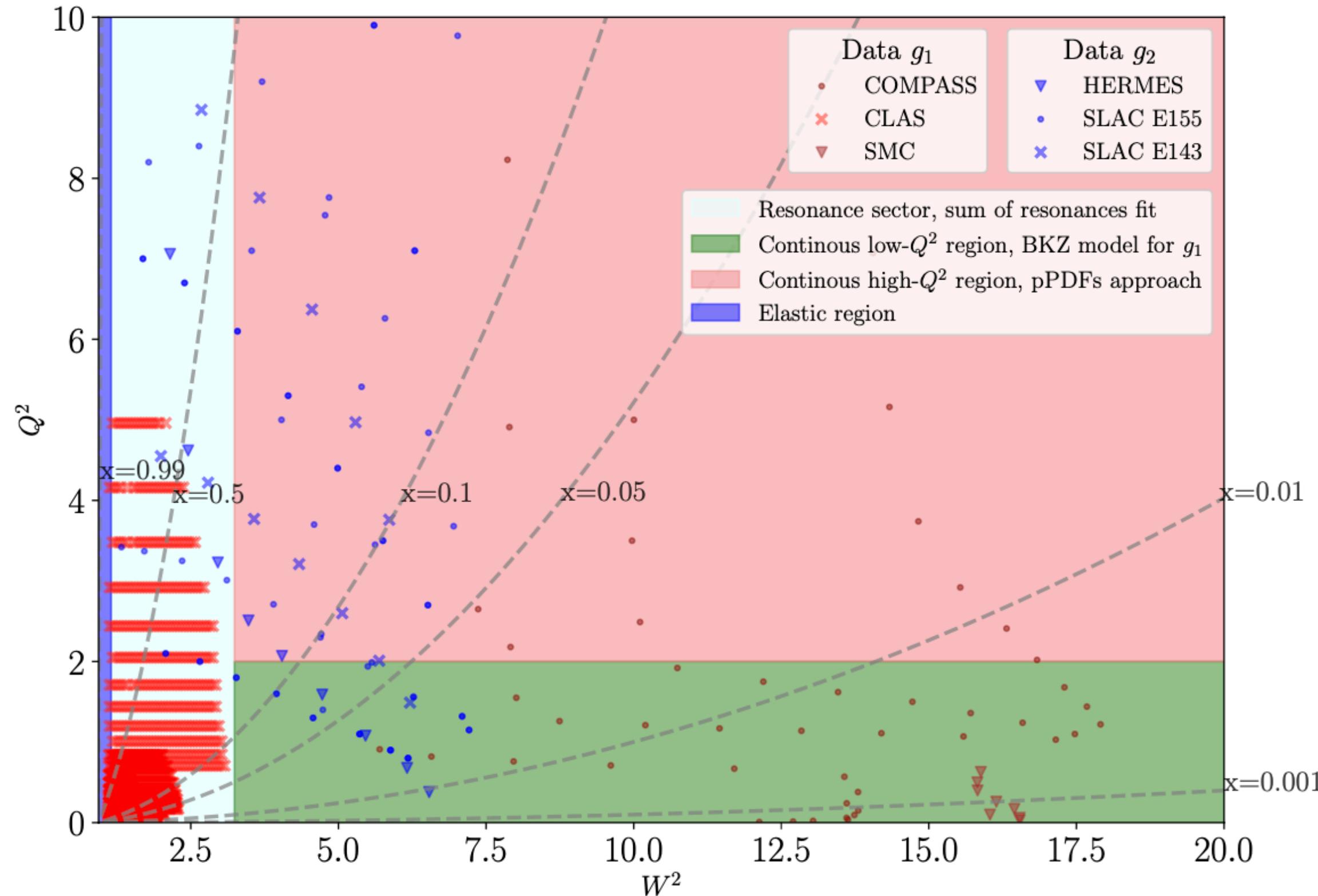
$$\Delta\sigma_{lp}(p) = \int dy \Delta\hat{\sigma}_{l\gamma}^{(0,0)}(yp) \Delta f_\gamma(x, \mu^2) + \cdot$$

equate and deduce photon distribution

# Polarized photon pdf

Allows to compute the polarized photon distribution

$$\Delta f_\gamma(x, \mu^2) = \frac{1}{2\pi\alpha(\mu^2)} \int_x^1 \frac{dz}{z} \left\{ \int_{Q_{\min}^2}^{\frac{\mu^2}{1-z}} \frac{dQ^2}{Q^2} \alpha_{\text{ph}}^2(Q^2) \right. \\ \left[ \left( 4 - 2z - \frac{4m_p^2 x^2}{Q^2} \right) g_1(x/z, Q^2) - \left( \frac{8m_p^2 x^2}{zQ^2} \right) g_2(x/z, Q^2) \right] \\ \left. + \alpha^2(\mu^2) 4(1-z) g_1(x/z, \mu^2) \right\}$$



Requires knowledge of structure functions over full kinematical regime : Elastic, Resonance, Continuum

Elastic: E/M Sach Form factors

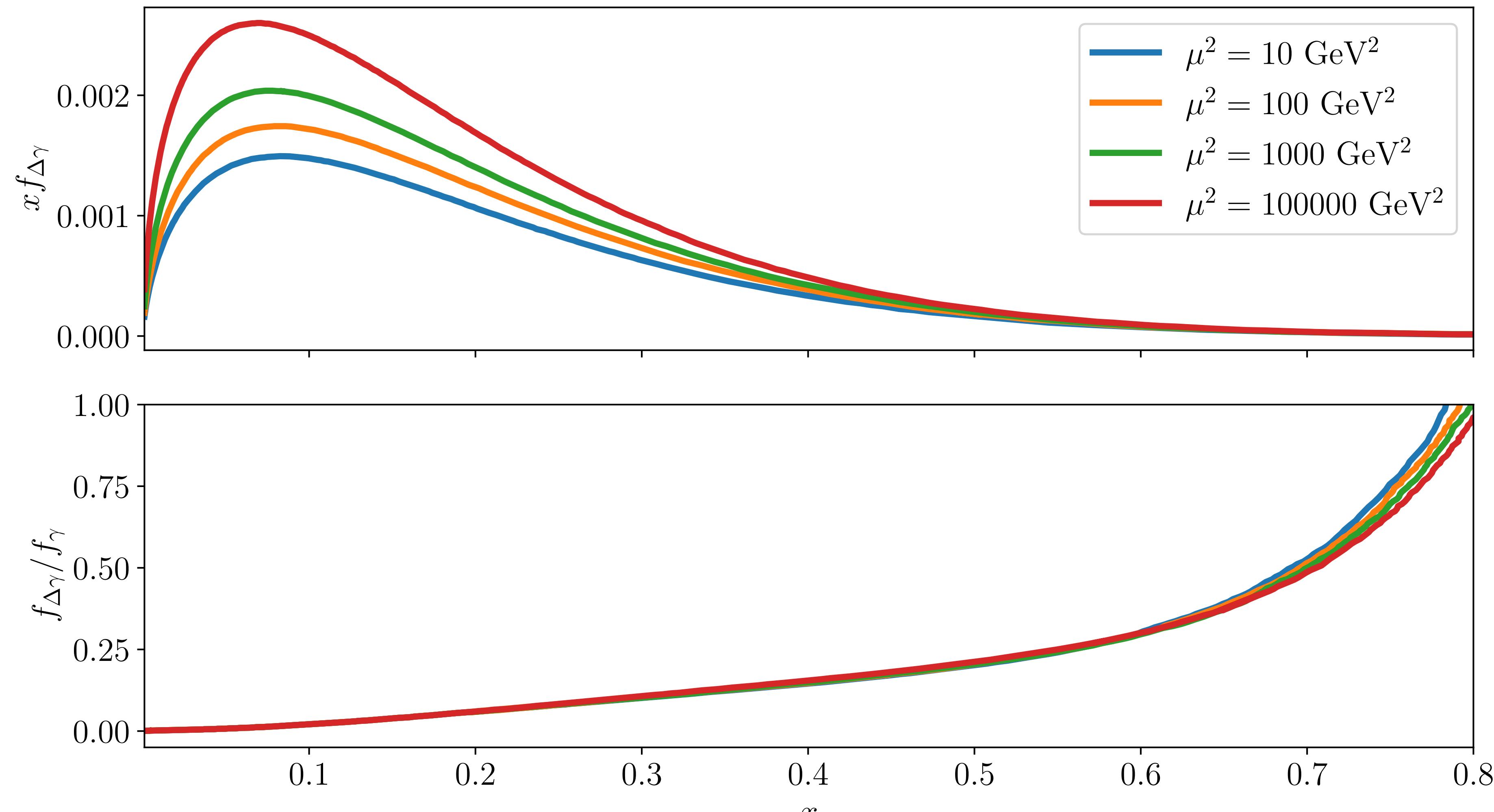
Resonance: Fits CLAS data

Low Q continuum: VMD model/Fits

High Q continuum: Perturbative QCD

deF., Palma, Volonnino (2024)

several sources of uncertainties...



**shape consistent with  $x^* f$**

► Uncertainties on distribution (10-20% level)

$$\int_{0.001}^1 \Delta\gamma dx \simeq 0.0049 \pm 0.0008$$

proton spin carried  
by photons ( $\sim 1\%$ )

# Conclusions

- ▶ First NNLO global analysis of polarized PDFs:
  - Good perturbative stability going from NLO to NNLO
  - Slight improvement in the description of data (after imposing cut on  $x_{\text{SIDIS}}$ )
- ▶ Outlook:
  - Include full NNLO SIDIS results → First stage: new NNLO analysis of FFs.
- ▶ Percent level accuracy required both for EXP and TH (even QED can contribute)
- ▶ Still rather incomplete picture of the proton's spin in terms of the contribution from quarks, anti-quarks and gluons.
- ▶ The spin program at the future EIC expected to give unique access to the proton's spin structure.





Who carries  
the spin of the  
proton?



I don't know, ask  
the americans, they  
are building a new  
collider

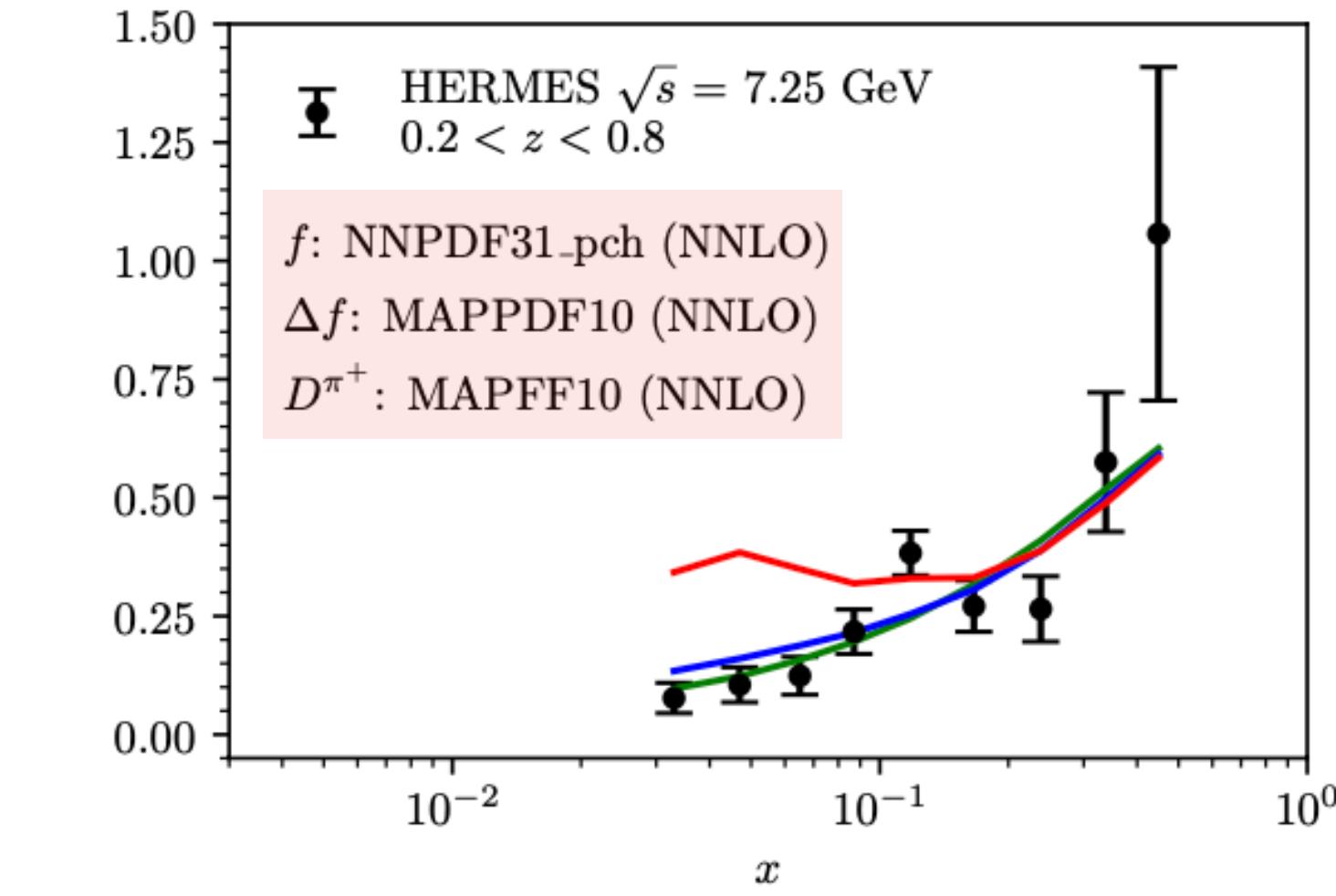
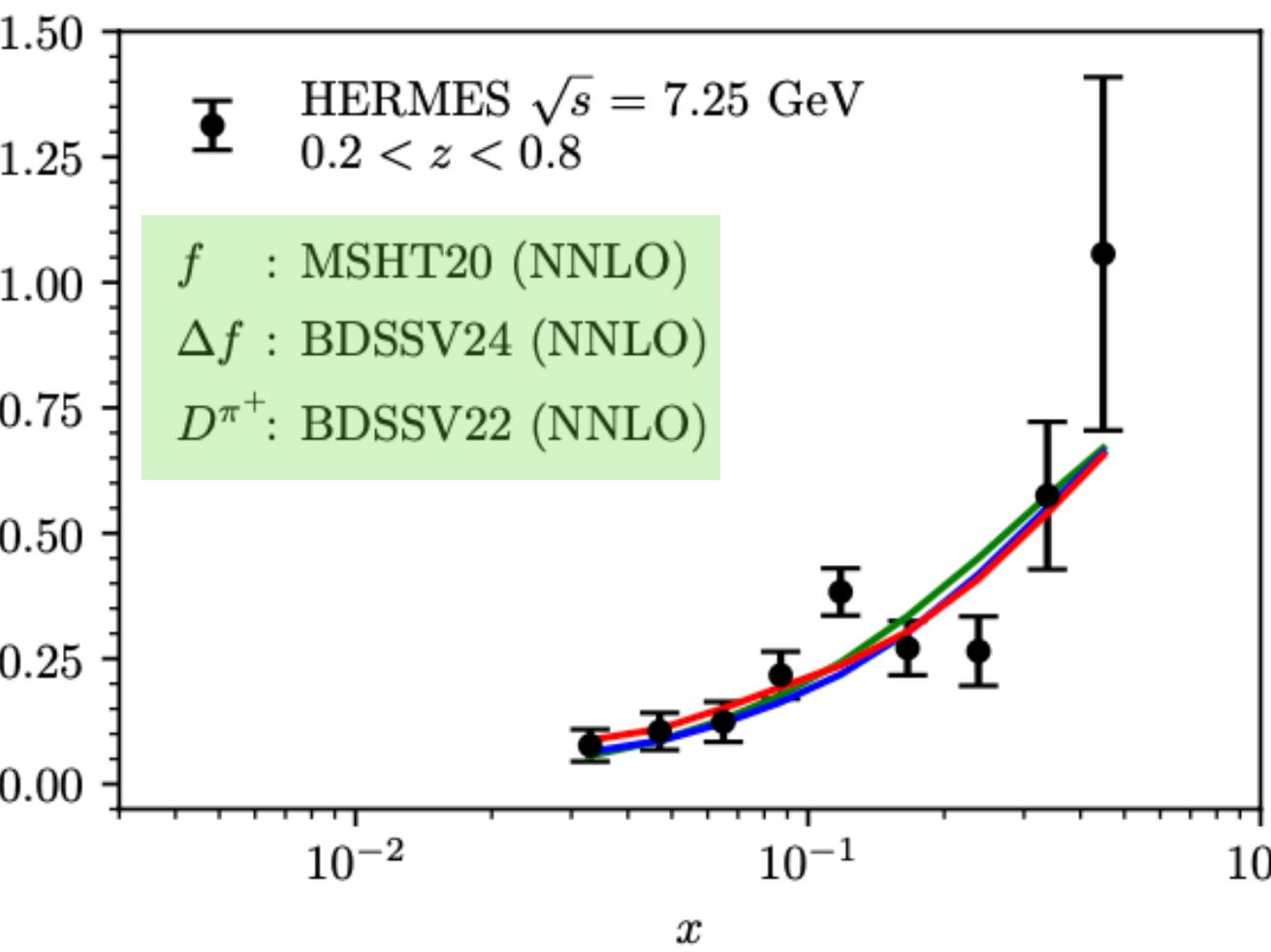
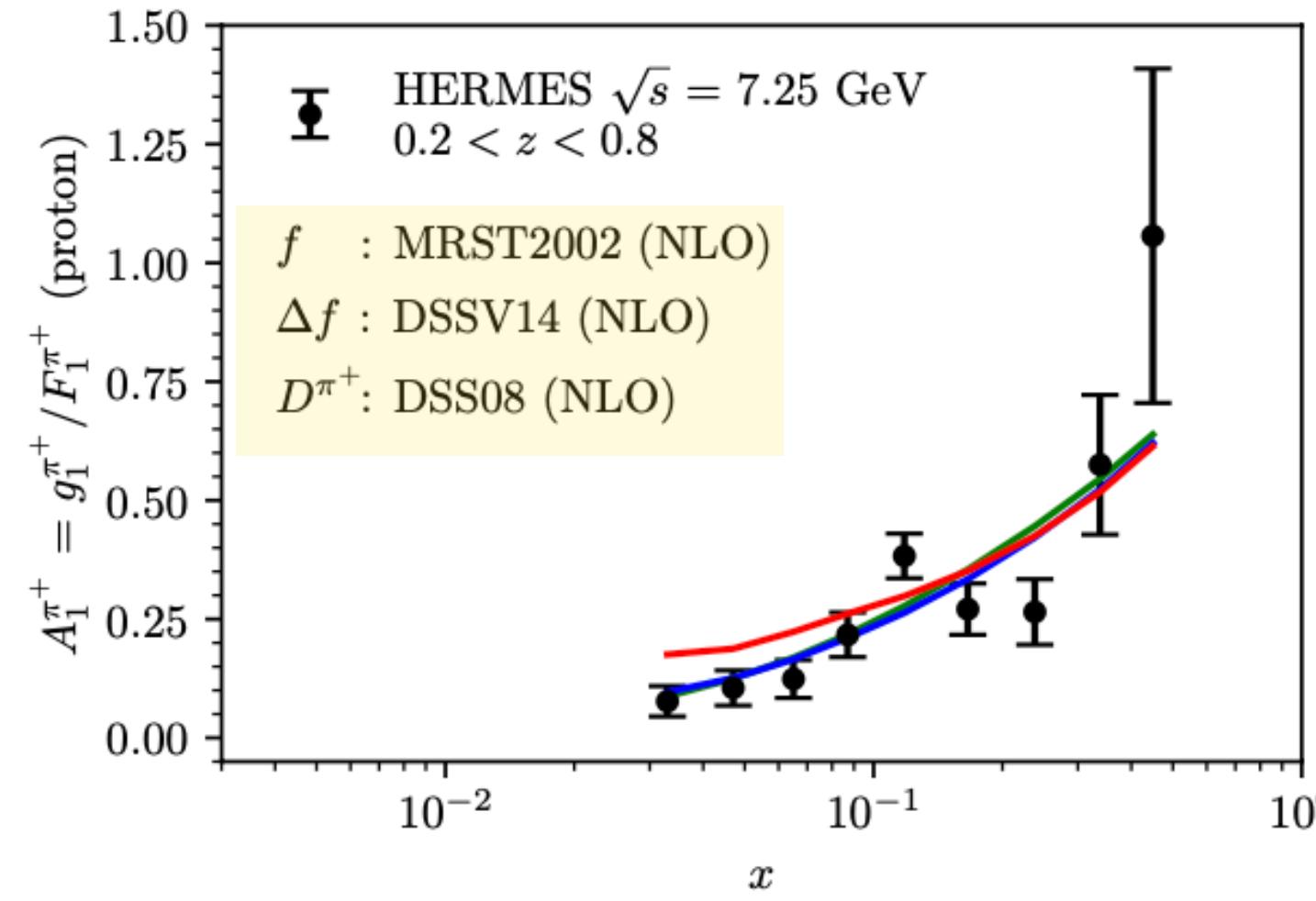
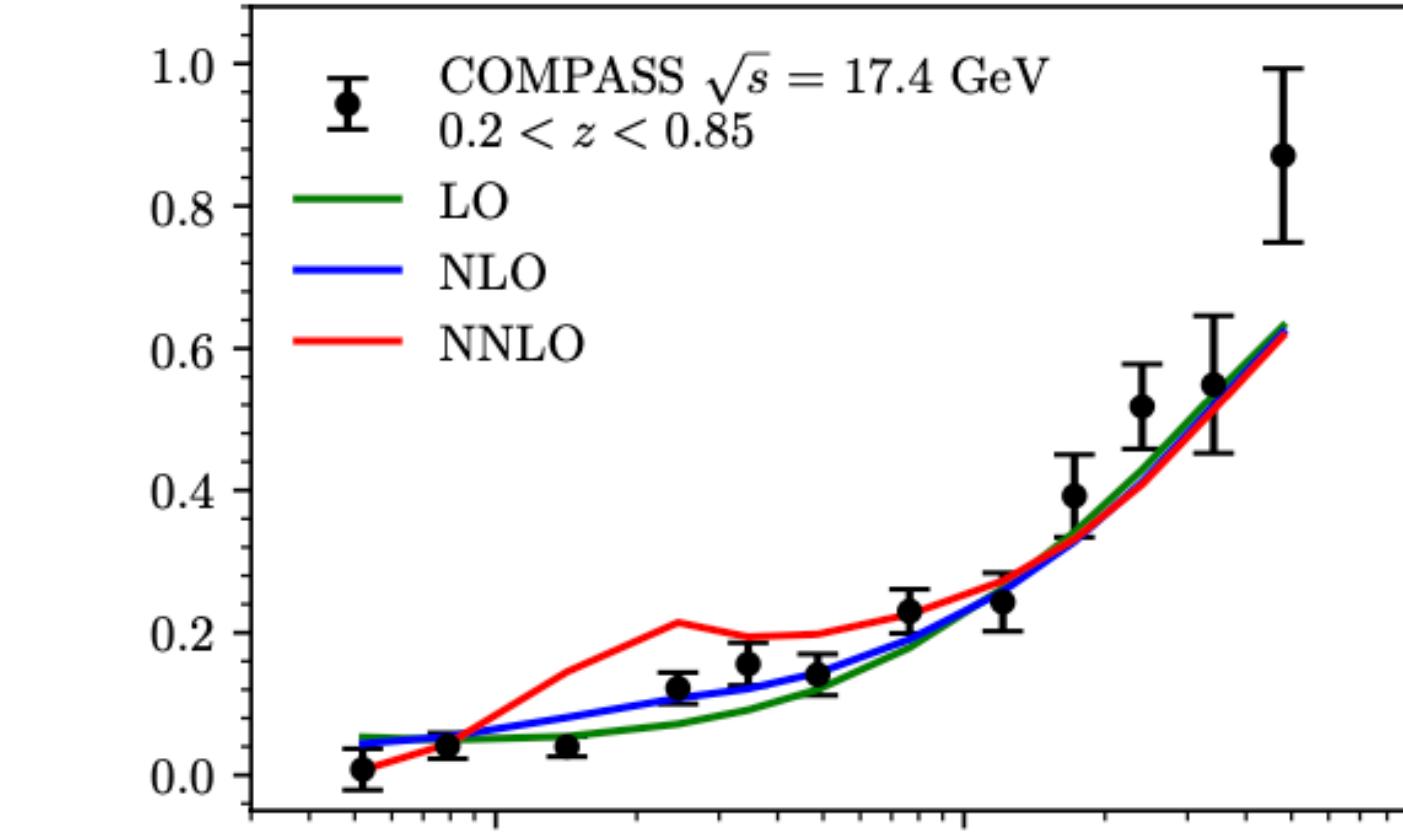
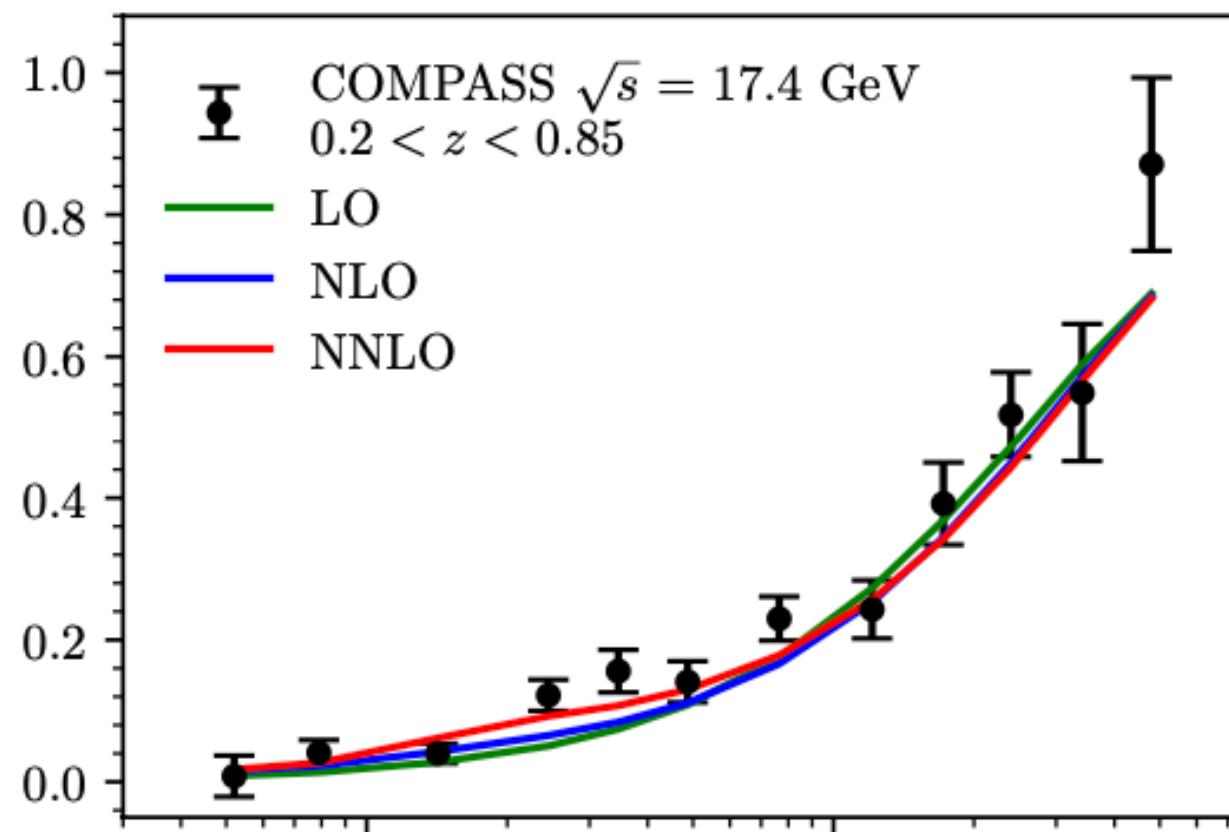
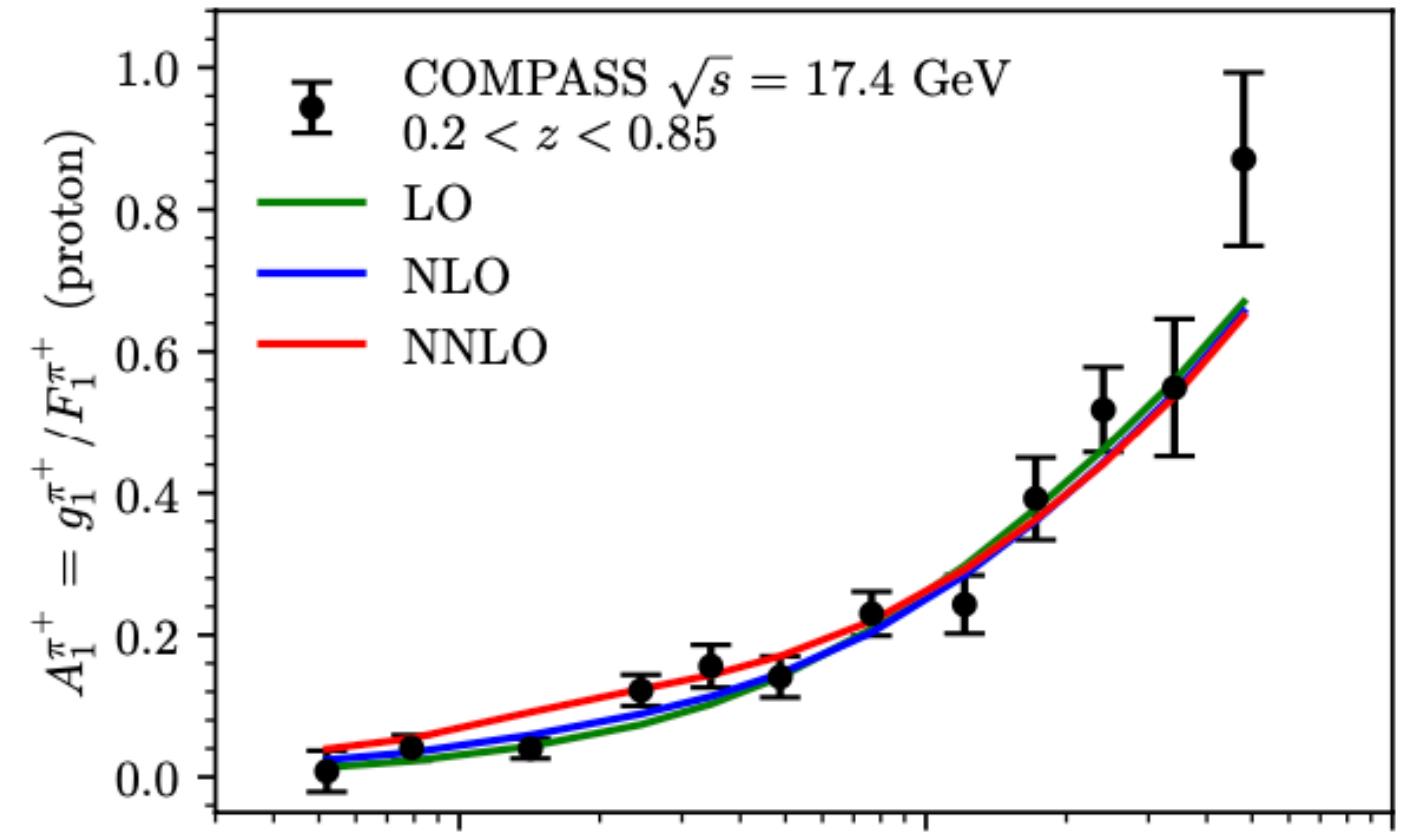
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the spin of the  
proton?

# THANKS

# Backup

# Post validation: Exact NNLO SIDIS coefficients

L.Bonino, T.Gehrmann, M.Löchner, K.Schonwald (2024)



NNLO matter for SIDIS  
 NNLO matter for EIC

FFs matter for SIDIS  
 No NNLO FFS for RHIC yet

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