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# Synergies between a U.S.-based Electron-Ion Collider and the European research in Particle Physics

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## Abstract

This document is submitted as input to the European Strategy for Particle Physics Update (ESPPU). In April 2024, the U.S.-based Electron-Ion Collider (EIC) project received approval for critical-decision 3A (CD-3A) allowing for Long-Lead Procurement, bringing its realization another step closer. since July 2022 the experimental ePIC Collaboration has been established around the realization of a general purpose detector. A large group of European scientists is already involved in EIC. Currently, about a quarter of the EIC User Group (consisting of over 1500 scientists) and of the ePIC Collaboration (consisting of 900 members) is based in Europe. This European involvement is not only an important driver of the EIC, but can also be beneficial for a number of related ongoing and planned particle physics experiments at CERN. In this document, the connections between the scientific questions addressed at CERN and at the EIC are outlined. The aim is to highlight how the synergies between the European Particle Physics research and the EIC project offer ample opportunities to foster progress at the forefront of collider physics.

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# 1 Introduction

23

The Electron-Ion Collider (EIC) is a major new research facility to advance the long term vision for Nuclear Physics to discover and understand the emergent phenomena of Quantum Chromo-Dynamics (QCD).

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The worldwide fast growing community of scientists interested in the EIC organized itself under the EIC Users Group (EICUG) [web site <http://www.eicug.org/>]. As of February 2018, the EICUG consists of 1549 scientists (including 377 theorists) from 303 institutions of 40 countries in all world regions, with a large European involvement consisting of 412 scientists (27%) from 86 institutions.

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In December 2019, following the extremely positive assessment by the U.S. National Academies of Sciences, Engineering, and Medicine (NAS) [1], the US Department of Energy (DoE) established EIC Critical Decision 0 (CD0), a “mission need” declaration, formally starting the EIC Project.

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The ePIC Collaboration has been established in July 2022, with the goal of realizing a general purpose detector designed to deliver the whole EIC core science program. As of February 2018, the ePIC Collaboration consists of 901 members, of whom 27% are based in Europe.

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In April 2024, the EIC Project received approval for critical-decision 3A (CD-3A) allowing the Project to start Long-Lead Procurement. This brings the realization of such a collider another step closer.

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Given the substantial European interest and involvement in the EIC and its physics program, it seems appropriate to outline to the Panel of the European Strategy for Particle Physics Update (ESPPU) the current European involvement in the EIC endeavor and the mutual benefits for particle physics experiments at CERN and EIC experiments.

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The general need for and uses of high-energy electron-proton and electron-ion collisions, often referred to as Deep Inelastic Scattering (DIS), are outlined in a separate document submitted to this Panel (is this still correct?). Also, detailed accelerator (R&D) aspects of the U.S.-based EIC are outlined in a dedicated document. Here, the focus will be on the EIC physics. The detector R&D program is outlined in a dedicated document, the synergies with present and planned CERN experiments, and outlining the European involvement in the EIC.

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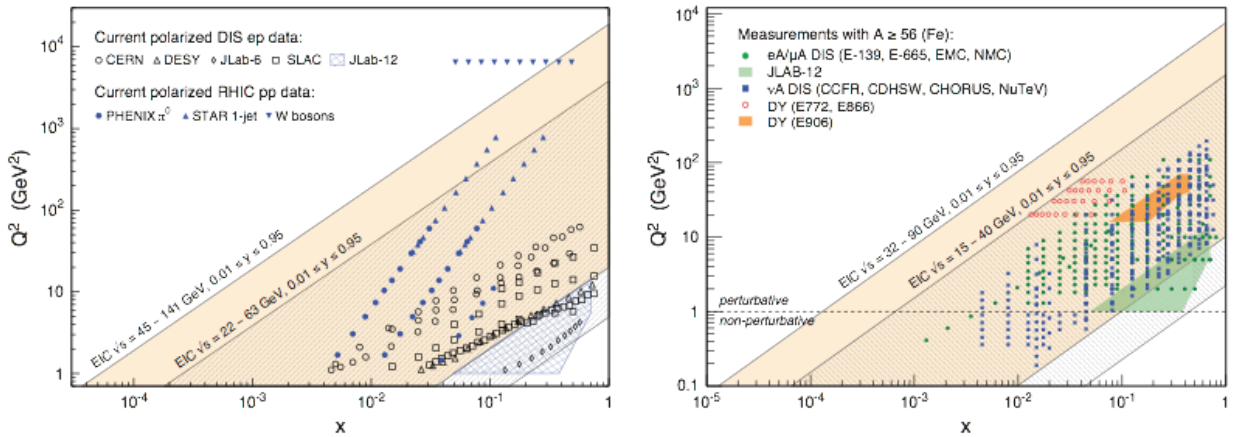


Figure 1: **Old plots! Do we need them?** Left: the phase space in  $(x, Q^2)$  covered in polarized electron-proton DIS by two different setups for the EIC, in comparison with past and current DIS machines and RHIC (updated version of Ref. [3]). A center-of-mass energy in the range of 20-100 GeV is foreseen for the EIC, with 45 GeV having maximum luminosity and 140 GeV being the maximal energy after a possible future upgrade. Right: EIC kinematical reach for nuclei, compared to earlier nuclear DIS experiments.

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The EIC is a machine that is unique compared to any previous DIS experiment because of the combined availability of high energy, high luminosity, ion versatility, and polarization. It is the first ever machine with the capability to collide highly polarized electrons on polarized protons and light ions, as well as on unpolarized heavier ions up to uranium. The EIC has a large reach in  $x$  and  $Q^2$  (see Fig. 1). High energy scattering of polarized electrons and ions, including both longitudinally and transversely polarized light ions, is crucial to a full understanding of the quark-gluon structure and dynamics of baryons, mesons, and nuclei. Compared to the HERA collider at DESY, the EIC will have lower energy but a much higher luminosity **put lumi plot?**. It can conduct measurements that

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53 have never been feasible before. The NAS committee finds the scientific case for EIC compelling, unique, timely.  
54 According to the NAS report [1]: “*The science questions that an EIC will answer are central to completing an*  
55 *understanding of atoms as well as being integral to the agenda of nuclear physics today. In addition, the development*  
56 *of an EIC would advance accelerator science and technology in nuclear science; it would as well benefit other fields*  
57 *of accelerator-based science and society, from medicine through materials science to elementary particle physics.*”

58 Since 1999, there have been many dedicated scientific meetings to shape the physics case for the EIC, most  
59 notably the yearly POETIC and EICUG meetings. In 2010 there was a ten-week program at the Institute of  
60 Nuclear Theory (INT) in Seattle that resulted in a 547-page document [4] detailing the science case. This was  
61 updated and condensed in the EIC white paper [3].

62 The NAS report summarizes the physics objectives as follows: “*An EIC can uniquely address three profound*  
63 *questions about nucleons - neutrons and protons - and how they are assembled to form the nuclei of atoms:*

- 64 - *How does the mass of the nucleon arise?*
- 65 - *How does the spin of the nucleon arise?*
- 66 - *What are the emergent properties of dense systems of gluons?”*

67 To answer these fundamental questions one needs to probe with high precision and high energy the quark and  
68 gluon structure of nucleons and nuclei. It is of great advantage to do this with a simple and well-known probe,  
69 such as the electron or the photon. The quark and gluon structure of nucleons is expressed theoretically in terms of  
70 parton distributions of various levels of detail and sophistication. For inclusive DIS, where one ignores details of the  
71 final state, the scattering process is described in terms of collinear (i.e., transverse-momentum integrated) Parton  
72 Density Functions (PDFs). When more aspects of the final state are measured, one can become sensitive to the  
73 three-dimensional momentum distributions (Transverse Momentum Dependent PDFs or TMDs). Exclusive and  
74 diffractive processes allow one to probe also transverse spatial distributions, given in terms of Generalized Parton  
75 Densities (GPDs). The information contained in TMDs and GPDs is truly complementary since the position  
76 of partons in a transverse plane is not Fourier conjugate to their transverse momentum. Hence, the combined  
77 investigation of PDFs, TMDs, and GPDs, will allow one to arrive at a more complete picture of how nucleons are  
78 composed at the level of quarks and gluons. These studies go far beyond global (and scale-dependent) observations  
79 like “*50% of the proton’s momentum is carried by gluons*” or “*only about 30% of the proton’s spin is carried by*  
80 *quarks*” (a well-known conclusion from the EMC, SMC, and later experiments). More specifically, the three major  
81 themes of the EIC physics program are:

- 82 - the flavor and spin structure of the proton
- 83 - three-dimensional structure of nucleons and nuclei in momentum and configuration space
- 84 - QCD in nuclei

85 In addition, the high luminosity and cleaner environment (with respect to hadron colliders) will enable precision  
86 studies in electroweak physics and some specific searches of physics beyond the Standard Model.

87 The EIC is expected to start operating toward the mid of the 2030’s. It will likely run concurrently with LHC  
88 after its high-luminosity upgrade. Hence, it seems appropriate and timely to outline below how investigations at  
89 the EIC of each of the above topics could benefit CERN experiments and vice versa.

## 90 **2 The flavor and spin structure of the proton**

91 This topic centers around the accurate determination of collinear parton densities for both unpolarized and polarized  
92 protons (and neutrons and deuterons). The unpolarized collinear PDFs currently used for LHC studies are primarily  
93 determined by the HERA data and have an (NNLO) uncertainty of 2.4% for up quarks at  $x = 0.5$  and  $Q^2 = 100$   
94 GeV, but 12% for down quarks, 140% for strange quarks and 34% for gluons [5]. The uncertainty quickly rises for  
95 all flavors at  $x \gtrsim 0.6 - 0.7$  (see Ref. [6]).

96 It is important to constrain PDFs in the limit of large  $x$  over a wide range of  $Q^2$  because they influence the  
97 production rate of high transverse momentum  $W$  and  $Z$  bosons and jets, as well as the possible production of  
98 new heavier partners that are predicted in several BSM extensions (see the DIS document for more details). LHC  
99 data will help to decrease these uncertainties, particularly after the high-luminosity upgrade [7], but for the search

100 for BSM physics at the LHC it is essential that the employed PDFs are obtained from data that are insensitive  
 101 to that BSM physics. Proposed future experiments at CERN, such as a fixed-target experiment [8] or the LHeC  
 102 experiment (see the related document submitted to the ESPPU Panel), could provide such data with high precision.  
 103 An overall 1% uncertainty in the PDFs would be the desired goal at the LHC to confront with theory [5]. High  
 104 statistics data obtained with the EIC from neutral and charge currents in electroweak DIS would also help to reach  
 105 a better precision at large  $x$ , especially for the EIC configuration with maximum center-of-mass energy of 140 GeV.  
 106 For instance, a projection study shows that charged current DIS at the EIC would have very strong impact on the  
 107  $x\bar{d} + x\bar{s}$  combination of quark distributions [6].

108 The EIC can measure various processes from which the contribution from different quark and antiquark flavors  
 109 can be separately extracted over a very broad range in  $x$ . These processes include not only neutral and charged  
 110 current electroweak DIS but also semi-inclusive DIS (SIDIS), where a hadron in the final state is identified. SIDIS  
 111 at the EIC is particularly useful for a precise determination of the strange quark distribution [6]. Such analyses  
 112 go hand in hand with the extraction of collinear fragmentation functions (FFs), where light flavor contributions  
 113 can be individually tagged. Therefore, our knowledge of PDFs is influenced also by the accuracy at which FFs are  
 114 known. The FFs are usually determined in  $B$ -factory experiments by measuring electron-positron collisions, but  
 115 this yields only a limited knowledge of each individual flavor contribution, and the gluon channel is reachable only  
 116 at subleading order. Ideally, one should extract FFs by performing a global fit to data of all available reactions,  
 117 including hadron collider data. At present, this was done only by one group [9, 10]. More recently, the NNPDF  
 118 collaboration found that hadron colliders data (CDF, CMS, ALICE) can significantly constrain the gluon FFs [11].  
 119 On the other hand, the (anti)strange FFs still suffer from large uncertainties, and this reflects in a large uncertainty  
 120 on the (anti)strange PDFs as well. More generally, the limitations of the current SIDIS data hinder a complete  
 121 flavor separation of unfavored channels. At the EIC, the high luminosity, combined with the large lever arm in  
 122 the hard scale  $Q$  and the purposefully planned detector capabilities, will allow for very precise studies of the flavor  
 123 dependence of FFs over a large phase space. Current studies on the projected relative error indicate that significant  
 124 improvements can be achieved also for PDFs of light flavors over a wide range of low to medium  $x$ , particularly  
 125 for the strange component [6]. Hence, combining inputs from the EIC, hadron colliders and  $B$ -factories, will allow  
 126 to drastically reduce the uncertainties on FFs, and could make it possible also to reach the ultimate goal of a  
 127 simultaneous extraction from data of both PDFs and FFs. Additionally, it could help in clarifying if the intrinsic  
 128 flavor content of the proton (i.e., Fock components in its wave function) receives contributions also from charm [12].

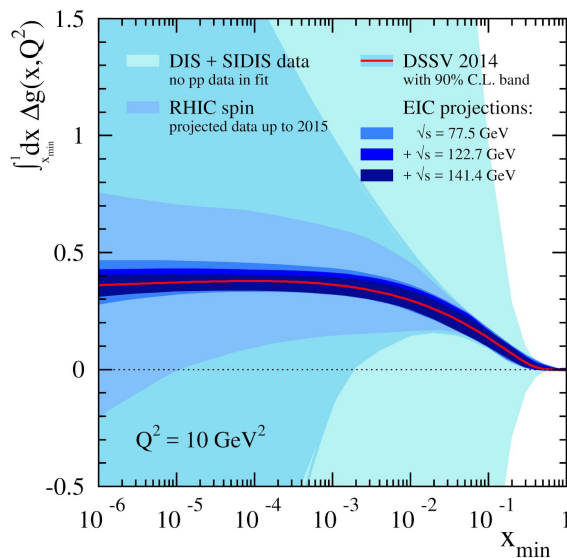


Figure 2: Integrated gluon helicity as a function of the attainable  $x_{\min}$  for various EIC configurations [13].

129 The availability of polarized proton beams at the EIC allows to do a similar analysis for the polarized quark and  
 130 gluon PDFs and thereby to shed further light on their contribution to the proton spin. In particular, the strange

131 quark and gluon PDFs still have large uncertainties. Recent results obtained at RHIC give evidence that the gluon  
 132 contribution is nonvanishing and positive, although the uncertainty is large because the result is very sensitive to  
 133 the minimum attainable  $x$  ( $x_{\min}$ ). With its unique capability of colliding longitudinally polarized electrons and  
 134 protons, while spanning small  $x$  even below  $10^{-4}$ , the EIC will drastically reduce this uncertainty (see Fig. 2). Both  
 135 DIS and SIDIS data will be important here. The development of reliable and versatile Monte Carlo generators to  
 136 analyze them is also of vital importance, as are the collaborations based on the European leadership in this field.

### 137 **3 Three-dimensional structure of nucleons and nuclei in momentum** 138 **and configuration space**

139 One of the main topics to be studied with the EIC is that of transverse momenta and positions of quarks and gluons  
 140 inside hadrons, as quantified by TMDs and GPDs, respectively. The ultimate aim of such studies is to gain much  
 141 deeper insight into the dynamics of quarks and gluons in hadrons than what can be inferred from PDFs alone.  
 142 TMDs and GPDs have been central to the investigations conducted earlier by the HERMES experiment at DESY and  
 143 currently by the COMPASS experiment at CERN and at Jefferson Lab with both polarized and unpolarized targets.  
 144 Quark TMDs are typically studied using SIDIS and in vector boson production in proton-proton collisions. Quark  
 145 GPDs are typically studied through the Deeply Virtual Compton Scattering (DVCS) and Deeply Virtual Meson  
 146 Production (DVMP) processes. Polarized light ions at the EIC (both longitudinal and transverse) are crucial for  
 147 fully resolving the diverse spin-spatial and spin-momentum correlations of hadronic structure.

#### 148 **TMD factorization and evolution**

149 Many critically important aspects of TMDs can be studied experimentally with the EIC. Apart from measuring  
 150 the various TMDs for polarized and unpolarized quarks, gluons and/or hadrons, the TMD formalism itself needs  
 151 thorough investigations.

152 First of all, TMD factorization, which allows for the theoretical description of particle spectra at small or  
 153 intermediate transverse momentum, needs to be tested. A main objective is to demonstrate to what extent the  
 154 TMDs are universal (like the collinear PDFs and FFs). Although the TMD for unpolarized quarks and hadrons  
 155 is expected to be the same in SIDIS and Drell-Yan (DY), for some polarized cases there is a calculable process  
 156 dependence. In 2015 and 2018 the COMPASS experiment acquired DY data with the objective of testing this (more  
 157 specifically, determining the predicted overall sign difference between the SIDIS and DY measurements of the Sivers  
 158 TMD effect) and the results so far are consistent with the expectations from the TMD formalism, although the  
 159 uncertainties are still large. The same applies to the asymmetry measurements at RHIC in  $W$  production [14].  
 160 Furthermore, for gluon TMDs even the unpolarized case is expected to be non-universal. Hence, it is important  
 161 to compare observables at the LHC, e.g., for quarkonium production [15], to related observables at the EIC. The  
 162 specific predictions from the TMD formalism can be safely tested in a high  $Q^2$  regime ( $Q^2 \gg M_p^2$ ), in order to  
 163 avoid complications from power suppressed quark-gluon correlation effects. Secondly, TMD evolution, i.e., the  $Q^2$   
 164 dependence of the predictions, needs to be tested, as was done extensively for the scaling violations in DIS with  
 165 HERA, showing that Altarelli-Parisi equations (DGLAP evolution) for collinear PDFs work over many orders of  
 166 magnitude, but at the same time hinting at possible deviations at small  $x$  and moderate  $Q^2$  [16]. Investigation of  
 167 the latter regime is another objective of the EIC (covered in Sec. 4).

#### 168 **Transition between TMD and collinear frameworks**

169 Besides testing its predictions, the limitations of the TMD formalism also need to be clarified. For this purpose  
 170 proton-proton collisions, even unpolarized ones, are particularly useful. Currently, it is expected on theoretical  
 171 grounds that azimuthal correlations in dijet production in  $p-p$  collisions do not factorize [17]. The size of the  
 172 factorization breaking effects is entirely unknown and can only be assessed from a comparison of high energy  $p-p$   
 173 and  $e-p$  collisions. The data from HERA on dijet production in  $e-p$  collisions is too limited to do this already  
 174 now. Recent studies of various dijet measurements at the EIC indicate that they are feasible [18, 19]. The same  
 175 applies to  $D$ -meson pair production, which is another process of interest for TMD factorization breaking tests.

176 The TMD formalism applies to the region of low transverse momenta (much smaller than the large scale in the  
 177 process). Therefore, sufficient momentum resolution is required in the full  $p_T$  range in order to test the limits of  
 178 applicability and the transition to the collinear formalism that applies at large  $p_T$ . In many of the observables of  
 179 interest, such as azimuthal correlations in the Drell-Yan process (studied by, e.g., CMS [20] and ATLAS [21]), the

180 momentum resolution is currently too limited, but that will improve over time. The same applies to di-photon and  
181 quarkonium (pair) production. The transition region of intermediate  $p_T$  has recently attracted much attention from  
182 the theoretical point of view [22, 23]. Because of the additional degrees of freedom (observables differential in more  
183 variables) it is harder to reach the same level of precision as for observables calculable in collinear factorization.  
184 Nevertheless, at present NNLL expressions are available for a number of processes and more data, from LHC and  
185 other experiments, are needed to confront with theoretical predictions.

## 186 GPDs

187 GPDs enter collinear factorization expressions and are expected to be universal. While the associated theory has  
188 reached a high level of sophistication, progress on relevant measurements has been comparatively slow because  
189 the experiments require high statistics and excellent detector coverage. One is dealing with exclusive processes in  
190 this case. HERA, COMPASS and Jefferson Lab experiments have already obtained experimental information on some  
191 of the GPDs, but this is only the beginning when it comes to extractions of the GPDs with precision and in a  
192 sufficiently large kinematic domain. The capabilities of the EIC will open the way to a thorough exploration of GPD  
193 properties. In particular, the large  $Q^2$  coverage will enable tests of GPD evolution and studies of power-suppressed  
194 higher-order correlations.

195 GPDs encode information about the spatial distribution of partons inside a hadron, correlated with their  
196 distribution in longitudinal momentum. The spatial distribution is obtained in a rather direct way by Fourier  
197 transforming the transverse momentum spectrum of the scattered hadron in suitable exclusive reactions like DVCS  
198 or DVMP. In an indirect manner, this distribution also influences the dynamics of  $p-p$  as well as  $p-A$  and  $A-A$   
199 collisions, namely in the context of multiparton interactions (MPI). In such interactions, several partons in the  
200 colliding hadrons take place in independent hard scatters, and the relative transverse distance between the partons  
201 is of crucial importance for this mechanism. Information on the spatial distribution of single partons from GPDs  
202 provides a quantitative baseline expectation, on top of which one can then attempt to assess correlation effects  
203 between different partons. In this sense, MPI and underlying events description will considerably benefit from a  
204 fully developed GPD picture.

205 Containing simultaneous information about longitudinal momentum and transverse position, GPDs give access  
206 to the orbital angular momentum carried by quarks and gluons. The most prominent (but not the only) expression  
207 of this is Ji's sum rule [24], which relates integrals over GPDs with the total (spin and orbital) angular momentum  
208 carried by a parton species in the proton. The detailed measurements that will become possible at the EIC will  
209 bring us much closer to a quantitative understanding of the role of orbital angular momentum and hence to the  
210 question of how the nucleon spin arises from its constituents at parton level. Moreover, GPDs will enable detailed  
211 studies of the spatial distribution (tomography) of several interesting observables like charge, pressure, energy and  
212 number densities.

213 TMDs and GPDs have been studied extensively. Various computer codes and tools, to a large extent developed  
214 by European groups, are available, such as for example TMDlib and TMDplotter (<https://tmdlib.hepforge.org/doxy/html/index.html/>) and for GPD model studies the package PARTONS (<https://arxiv.org/abs/1512.06174>). There are also proposals for future TMD and GPD studies at CERN. The AMBER experiment at the M2  
217 beam line of the CERN SPS started operation in 2023 [25], planning Drell-Yan measurements with pion and kaon  
218 beams. Another initiative that is being put forward is that of a fixed-target program which would allow scattering  
219 of an LHC proton or lead beam on a polarized fixed-target. The physics case for such an experiment has been  
220 developed over recent years under the name "AFTER@LHC" [8] and a particular proposal for such a fixed-target  
221 experiment at LHCb is called "LHC-spin". There would of course be ample synergies between such experiments and  
222 the EIC as well.

## 223 4 QCD in nuclei

224 The EIC could be the first ever collider experiment for deeply inelastic scattering with nuclei, opening up a large  
225 new window of high atomic number  $A$  at small  $x$  and large  $Q^2$  that has never before been experimentally accessed  
226 (see Fig. 1). Electron-ion collisions allow to study the internal structure of heavy ions in terms of elementary  
227 partonic constituents, quarks and gluons. In addition to a new understanding of QCD in large nuclei being a  
228 fundamental topic in itself, it is also complementary to the program of heavy-ion collisions at the LHC. Probing

229 the partonic structure of the colliding nuclei by DIS experiments is important for understanding the production  
230 of both the matter that then becomes a quark-gluon plasma, and the jets and other hard probes that are used to  
231 explore its properties.

### 232 **Nuclear PDFs**

233 At the LHC kinematics, one is mostly sensitive to the region of relatively small  $x$ , where the nuclear PDFs are  
234 very poorly known since they are hardly constrained by any currently existing data. Studies of modifications to  
235 the structure of jets as they pass through deconfined QCD matter are becoming very sophisticated and will form  
236 an increasingly important part of the nuclear collision program at the LHC when it moves to higher luminosities.  
237 These studies will require a reduction of the large uncertainty of nuclear PDFs at small  $x$ , and the EIC will have  
238 a significant impact here [26]. In particular, one should note that accurate measurements of charm structure  
239 functions at the EIC would have an even larger effect than total cross section measurements, also improving the  
240 determination of the nuclear gluon PDF at large  $x$  (aiding the further study of the EMC effect and of (anti-)  
241 shadowing). Semi-inclusive DIS at large  $x$  in nuclei also provides a laboratory to study the energy loss of high-  
242 energy partons passing through ordinary, confined nuclear matter. Such measurements provide the comparison  
243 that is necessary to calibrate and check theoretical approaches to jet energy loss. Last but not least, the exclusive  
244 and inclusive measurements will give access to the GPDs and TMDs of nuclei, respectively.

### 245 **Saturation phenomena**

246 At high collision energies, or equivalently small  $x$ , the phase space available for emitting soft gluons is very  
247 large. Since every emitted gluon is itself a source of further radiation, a fast growing cascade is created that  
248 leads to a strong growth of the gluon distribution. Eventually unitarity is expected to be preserved due to the  
249 nonlinear interactions, i.e., gluon mergings, that start to play an important role, leading to the phenomenon of  
250 gluon saturation. Gluon saturation then controls the physics in DIS collisions in the small  $x$  limit at moderate  $Q^2$ .  
251 The gluons in the same regime are also responsible for the production of deconfined QCD matter – the quark-gluon  
252 plasma – in heavy ion collisions, studied at the LHC and in particular by the ALICE collaboration.

253 Since a high atomic number  $A$  also increases the gluon density at a given  $Q^2$ , gluon saturation is accessible in  
254 DIS experiments at lower energies with nuclei than with protons. The EIC will provide a versatile experimental  
255 program to investigate in detail this new regime of QCD. In order to fully understand this regime, it is important  
256 to simultaneously measure inclusive and semi-inclusive cross sections, inclusive diffraction (diffractive structure  
257 functions) and exclusive reactions, such as vector meson production and diffractive dijets. The EIC is being  
258 designed to be a facility that can perform this broad set of measurements that are necessary for a full picture of  
259 the gluonic structure of nuclei.

260 The small- $x$  physics at the EIC is complemented by several measurements that are performed in hadron colliders  
261 (RHIC and LHC) to probe aspects of the same physics. Inclusive particle production and two-particle correlations  
262 in  $p-p$  and  $p-A$  collisions, especially at forward rapidities, directly probe the small- $x$  gluons in nuclei. At CERN,  
263 these forward measurements are performed at LHCb and, especially with future instrumentation upgrades, at CMS  
264 and ATLAS. Furthermore, as mentioned, the small- $x$  and high- $A$  program at the EIC is very closely connected to  
265 studies of deconfined QCD matter in heavy-ion collisions. The correct interpretation of several aspects of heavy-ion  
266 collisions, such as the ridge, requires data from high-energy electron-ion experiments.

### 267 **Initial conditions for Quark-Gluon Plasma studies**

268 One of the most challenging problems in heavy-ion physics is to understand how the gluons and quarks from the  
269 colliding nuclei form a thermalized plasma. Data from  $A-A$  collisions seems to be well described by models  
270 assuming a very quick formation of an equilibrated medium. It is, however, very difficult to get direct experimental  
271 access to the earliest stages of a heavy-ion collision, and the theoretical understanding of the thermalization process  
272 is still quite incomplete. Here, the picture of the small- $x$  degrees of freedom in the nucleus obtained from electron-  
273 ion collisions is crucial, as it provides the starting point – the initial condition – from which one evolves towards a  
274 deconfined matter.

275 In recent studies at the LHC, it has become clear that effects usually attributed to collective behavior in  $A-A$   
276 collisions, such as elliptic flow and the ridge, are also visible in  $p-A$  and  $p-p$  collisions. There is currently an  
277 intense debate in the field concerning the correct interpretation of these results. They have been explained either in  
278 terms of multiparticle correlations already present in the colliding protons or nuclei, or alternatively ones generated  
279 by collective interactions if deconfined QCD matter is present. For a resolution of the puzzle posed by these results,

280 a baseline measurement of such initial state correlations in a more tractable collision system is essential. The ideal  
281 experiments for this are provided by  $e - p$  and  $e - A$  collisions.

## 282 UPCs

283 Besides being a source for many small- $x$  gluons, ultra-relativistic heavy ions also form a source of strong electro-  
284 magnetic fields, which can be probed in Ultra-Peripheral Collisions (UPCs) at high energies. This effectively leads  
285 to high-energy photon-ion scattering, where partons at small  $x$  are probed both in nuclei (in  $A - A$  collisions) and  
286 protons (in  $p - A$  collisions). Since in these collisions the photon is always quasi-real, one does not have the same  
287 ability to vary  $Q^2$  as in DIS experiments, but on the other hand the higher collision energy gives access to smaller  
288 values of  $x$  than available at the EIC. To date, exclusive vector meson production in  $\gamma - A$  collisions measured  
289 by ALICE [27] have been used to probe nuclear gluons at small  $x$ . The possibility to separately perform coher-  
290 ent (nucleus stays intact) and incoherent (nucleus breaks up into smaller color neutral fragments) measurements  
291 gives an additional handle on probing the nuclear geometry and its fluctuations, which are important features for  
292 understanding the initial state of  $A - A$  collisions. Exclusive  $J/\psi$  and  $\Upsilon$  productions in UPCs have been studied  
293 at ALICE, LHCb and CMS. The results (e.g., see Refs. [28, 29, 30]) are consistent with HERA measurements of the  
294 same process at lower energies, and can be used to constrain the  $x$  dependence in different theory calculations.

## 295 5 Electroweak physics and the search for physics beyond the Standard 296 Model

297 Although not the main objective of the EIC, certain EW precision measurements, despite  $Q$  being well below  
298  $M_W$ , can be done at the EIC thanks to the high luminosity. For example, precision determinations of the weak  
299 mixing angle  $\sin^2 \theta_W$  can be performed in collisions of polarized electrons and unpolarized deuterons (see Fig. 3).  
300 This measurement as a function of energy is of interest because deviations from the SM running would signal  
301 BSM physics. These modifications can not only affect the high energy part, but also the low energy one, e.g.,  
302 through effects induced by “dark”  $Z$  bosons [31]. The EIC data is projected to be in the intermediate  $Q$  range,  
303 which is beyond the reach of fixed-target experiments. Apart from the weak mixing angle, the EIC can put  
304 further constraints on electron  $\rightarrow \tau$  flavor violating transitions, improving on an earlier limit set by HERA and  
305 complementing the limits from  $\tau$  decays at  $B$ -factories (see Ref. [4] for further discussion of these and related  
306 topics).

307 The EIC can contribute to reducing the uncertainties in the determination of the  $W$  mass from its  $p_T$  spectrum.  
308 Currently, the best determination of the  $W$  mass is deduced from a global electroweak fit [33] ( $\delta M_W = 8$  MeV), but  
309 precise extractions of this SM parameter have been obtained also by fitting the transverse mass and momentum  
310 distributions of the  $W$  decay products in proton-proton collisions at ATLAS [34] and in proton-antiproton collisions  
311 at D0 and CDF [35] ( $\delta M_W \approx 20$  MeV). By introducing a flavor dependence in the intrinsic transverse momenta of  
312 the partons entering the collision, it was recently found that this nonperturbative effect (already observed in the  
313 HERMES data at low energy) has a statistically significant impact on the extracted values of the  $W^\pm$  masses [36].  
314 The details of this flavor dependence are best studied in SIDIS by looking at unintegrated multiplicities. In this  
315 way the measurements at the EIC will indirectly affect the accuracy level in the determination of the  $W$  mass,  
316 consequently influencing also searches of new BSM physics.

317 The Standard Model does not contain fundamental tensor interactions. However, loop effects and new interac-  
318 tions at the TeV scale can effectively generate such interactions that can be indirectly constrained either by data  
319 at the LHC or in hadronic decays at low energy (see, e.g., Ref. [37] and references therein). In the latter case,  
320 helicity-flip parameters in the neutron  $\beta$ -decay cross section can be combined with information on the isovector  
321 component of the nucleon tensor charge to put constraints on novel tensor interactions at the TeV scale. Infor-  
322 mation on the nucleon tensor charge can come from very precise lattice computations, in particular regarding  
323 its isovector component. But it can also be obtained from SIDIS (HERMES, COMPASS) or proton-proton collisions  
324 (RHIC) with transversely polarized hadrons by extracting the chiral-odd transversity distribution, whose first  
325 Mellin moment is the tensor charge. The EIC can significantly contribute in reducing the current uncertainty on  
326 the phenomenological extraction of the tensor charge by widely enlarging the covered  $x$  and  $Q^2$  range.



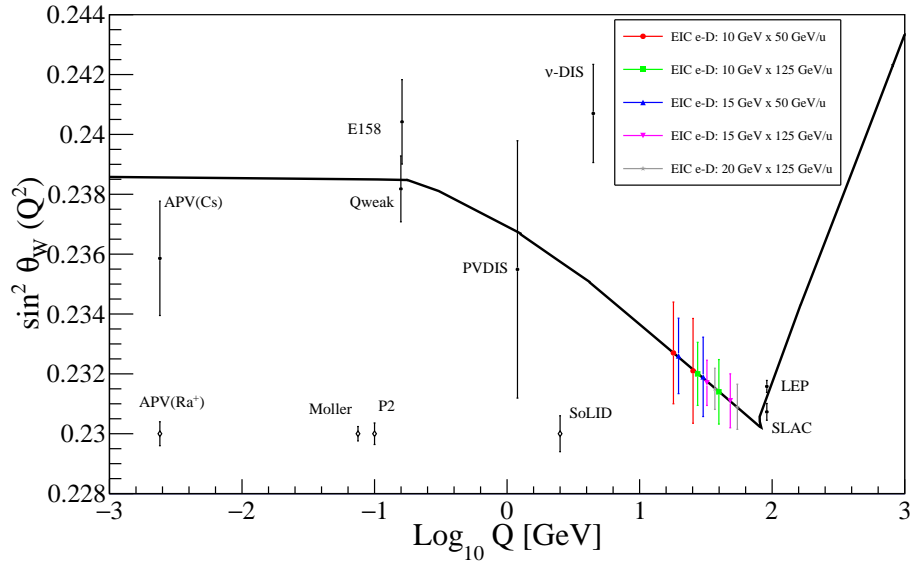


Figure 3: The dependence of  $\sin^2 \theta_W$  on  $Q$  calculated in the Standard Model (solid line) compared with existing data and projected fixed-target and EIC measurements. For the EIC projections 200 days of electron-deuteron running is assumed (with luminosity of  $267 \text{ fb}^{-1}$  per nucleon and electrons 80% polarized). QED and QCD radiative corrections and effects of detector smearing are taken into account [32].

## 6 Conclusions

In this document for the ESPPU, we outline the European involvement in the U.S.-based Electron-Ion Collider (EIC) project, which recently received strong endorsement by the U.S. National Academies of Sciences, Engineering, and Medicine. More than a quarter of the EIC User Group consists of European scientists, indicating a large European interest in the EIC, in both its science case and its detector and accelerator development. Furthermore, this document reviews the large variety of mutual benefits for CERN and EIC experiments, ranging from the connections between the scientific questions addressed to the shared interest regarding detector R&D. We conclude that strengthening the ties between the particle physics community in Europe and the EIC project, would be very beneficial to all parties involved and could foster important progress in research at the forefront of collider physics.

## References

- [1] NAS, *An Assessment of U.S.-based Electron-Ion Collider Science* (2018) <https://www.nap.edu/catalog/25171/an-assessment-of-us-based-electron-ion-collider-science>
- [2] NSAC, *Reaching for the Horizon: the 2015 US Long Range Plan for Nuclear Science* (2015) <https://science.energy.gov/np/nsac>
- [3] A. Accardi *et al.*, *Electron Ion Collider: The Next QCD Frontier*, Eur. Phys. J. **A52** (2016) 268 [arXiv:1212.1701]
- [4] D. Boer *et al.*, *Gluons and the quark sea at high energies: Distributions, polarization, tomography*, arXiv:1108.1713
- [5] G. Salam, talk given at POETIC 7 (2016) <https://indico.bnl.gov/event/2095/contributions/4437/attachments/4027/4779/Poetic-LHC-Theory-precision.pdf>
- [6] E.C. Aschenauer, talk given at INT program INT-1803 *Probing Nucleons and Nuclei in High Energy Collisions* (2018) [http://www.int.washington.edu/talks/WorkShops/int\\_18\\_3/People/Aschenauer\\_EC/Aschenauer2.pdf](http://www.int.washington.edu/talks/WorkShops/int_18_3/People/Aschenauer_EC/Aschenauer2.pdf)

- 350 [7] R.A. Khalek *et al.*, (2018) arXiv:1810.03639
- 351 [8] C. Hadjidakis *et al.*, (2018) arXiv:1807.00603
- 352 [9] D. de Florian *et al.*, Phys. Rev. **D91** (2015) 014035 [arXiv:1410.6027]
- 353 [10] D. de Florian *et al.*, Phys. Rev. **D95** (2017) 094019 [arXiv:1702.06353]
- 354 [11] V. Bertone *et al.* (NNPDF), Eur. Phys. J. **C78** (2018) 651 [arXiv:1807.03310]
- 355 [12] R.D. Ball *et al.* (NNPDF), Eur. Phys. J. **C77** (2017) 663 [arXiv:1706.00428]
- 356 [13] E.C. Aschenauer, R. Sassot, and M. Stratmann, Phys. Rev. **D92** (2015) 094030 [arXiv:1509.06489]
- 357 [14] L. Adamczyk *et al.* (STAR), Phys. Rev. Lett. **116** (2016) 132301 [arXiv:1511.06003]
- 358 [15] D. Boer and C. Pisano, Phys. Rev. **D86** (2012) 094007 [arXiv:1208.3642]
- 359 [16] R.D. Ball *et al.*, Eur. Phys. J. **C78** (2018) 321 [arXiv:1710.05935]
- 360 [17] T.C. Rogers and P.J. Mulders, Phys. Rev. **D81** (2010) 094006 [arXiv:1001.2977]
- 361 [18] A. Dumitru, V. Skokov, and T. Ullrich, (2018) arXiv:1809.02615
- 362 [19] L. Zheng *et al.*, Phys. Rev. **D98** (2018) 034011 [arXiv:1805.05290]
- 363 [20] V. Khachatryan *et al.* (CMS), Phys. Lett. B **750** (2015) 154 [arXiv:1504.03512]
- 364 [21] G. Aad *et al.* (ATLAS), JHEP **1608** (2016) 159 [arXiv:1606.00689]
- 365 [22] J. Collins *et al.*, Phys. Rev. **D94** (2016) 034014 [arXiv:1605.00671]
- 366 [23] M.G. Echevarria *et al.*, Phys. Lett. **B781** (2018) 161 [arXiv:1801.01480]
- 367 [24] X.-D. Ji, Phys. Rev. Lett. **78** (1997) 610 [hep-ph/9603249]
- 368 [25] B. Adams, C. A. Aidala, G. D. Alexeev *et al.*, COMPASS++/AMBER: Proposal for Measurements at the  
369 M2 beam line of the CERN SPS Phase-1: 2022-2024 CERN-SPSC-2019-022 ; SPSC-P-360
- 370 [26] E.C. Aschenauer *et al.*, Phys. Rev. **D96** (2017) 114005 [arXiv:1708.05654]
- 371 [27] E. Abbas *et al.* (ALICE), Eur. Phys. J. **C73** (2013) 2617 [arXiv:1305.1467]
- 372 [28] B.B. Abelev *et al.* (ALICE), Phys. Rev. Lett. **113** (2014) 232504 [arXiv:1406.7819]
- 373 [29] R. Aaij *et al.* (LHCb), JHEP **09** (2015) 084 [arXiv:1505.08139]
- 374 [30] A.M. Sirunyan *et al.* (CMS), JHEP **1806** (2018) 128 [arXiv:1803.11133]
- 375 [31] H. Davoudiasl, H.-S. Lee, and W.J. Marciano, Phys. Rev. **D92** (2015) 055005 [arXiv:1507.00352]
- 376 [32] Y.X. Zhao *et al.*, Eur. Phys. J. **A53** (2017) 55 [arXiv:1612.06927]
- 377 [33] M. Baak *et al.* (Gfitter Group), Eur. Phys. J. **C74** (2014) 3046 [arXiv:1407.3792]
- 378 [34] M. Aaboud *et al.* (ATLAS), Eur. Phys. J. **C78** (2018) 110 [arXiv:1701.07240]
- 379 [35] T.A. Aaltonen *et al.* (CDF), Phys. Rev. **D89** (2014) 072003 [arXiv:1311.0894]
- 380 [36] A. Bacchetta *et al.*, Phys. Lett. **B** (2018) in press [arXiv:1807.02101]
- 381 [37] R. Gupta *et al.*, Phys. Rev. **D98** (2018) 034503 [arXiv:1806.09006]