

Brookhaven forum 2025: Launching the second Century of Quantum Physics

The strong force could have heated the Universe's first moment

Kim V. Berghaus

California Institute of Technology



Brookhaven
National Laboratory

A conceptual illustration of the universe's evolution. It features a central bright yellow-white point from which numerous thin white lines radiate outwards. To the left, a series of green, glowing, nebula-like structures are arranged in a sequence, suggesting the early stages of cosmic expansion. To the right, a sequence of blue and purple spheres, resembling galaxies or planets, are shown in a line, representing the formation and growth of celestial bodies. The background is a deep blue with various white and yellow circular and spherical shapes of different sizes, some with internal patterns, creating a sense of depth and complexity.

What are the laws of nature

that govern the evolution of our Universe

A conceptual diagram illustrating the evolution of the universe. It features a central bright yellow-white point from which numerous lines radiate outwards. To the left, these lines lead to a series of glowing green and blue particles, representing the early universe. To the right, the lines lead to a sequence of celestial bodies: a small blue sphere, a larger blue sphere with a grid, and a large, detailed blue sphere containing a complex internal structure, representing the formation of galaxies and the current universe. The background is a dark blue gradient with faint, glowing lines and particles.

What are the laws of nature
Standard Model of Particle Physics

Λ CDM Model of Cosmology
that govern the evolution of our Universe

What are the laws of nature
Standard Model of Particle Physics



Λ CDM Model of Cosmology

that govern the evolution of our Universe

Fundamental Physics in the Era of Precision Cosmology

Evolving dark energy



A Tantalizing 'Hint' That Astronomers Got Dark Energy All Wrong

Scientists may have discovered a major flaw in their understanding of that mysterious cosmic force. That could be good news for the fate of the universe.

New York Times,
April 4th , 2024

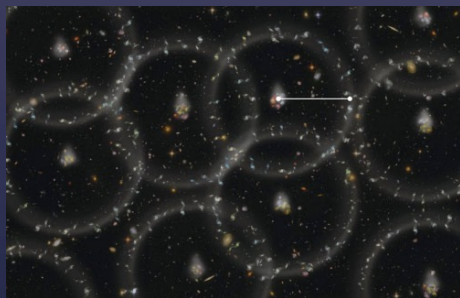
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DESI measures baryon acoustic oscillations and combined with other datasets sees hints of evolving dark energy

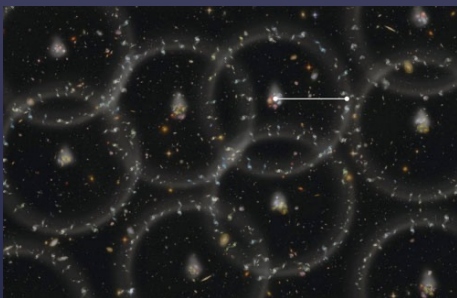
Evolving dark energy



A Tantalizing 'Hint' That Astronomers Got Dark Energy All Wrong

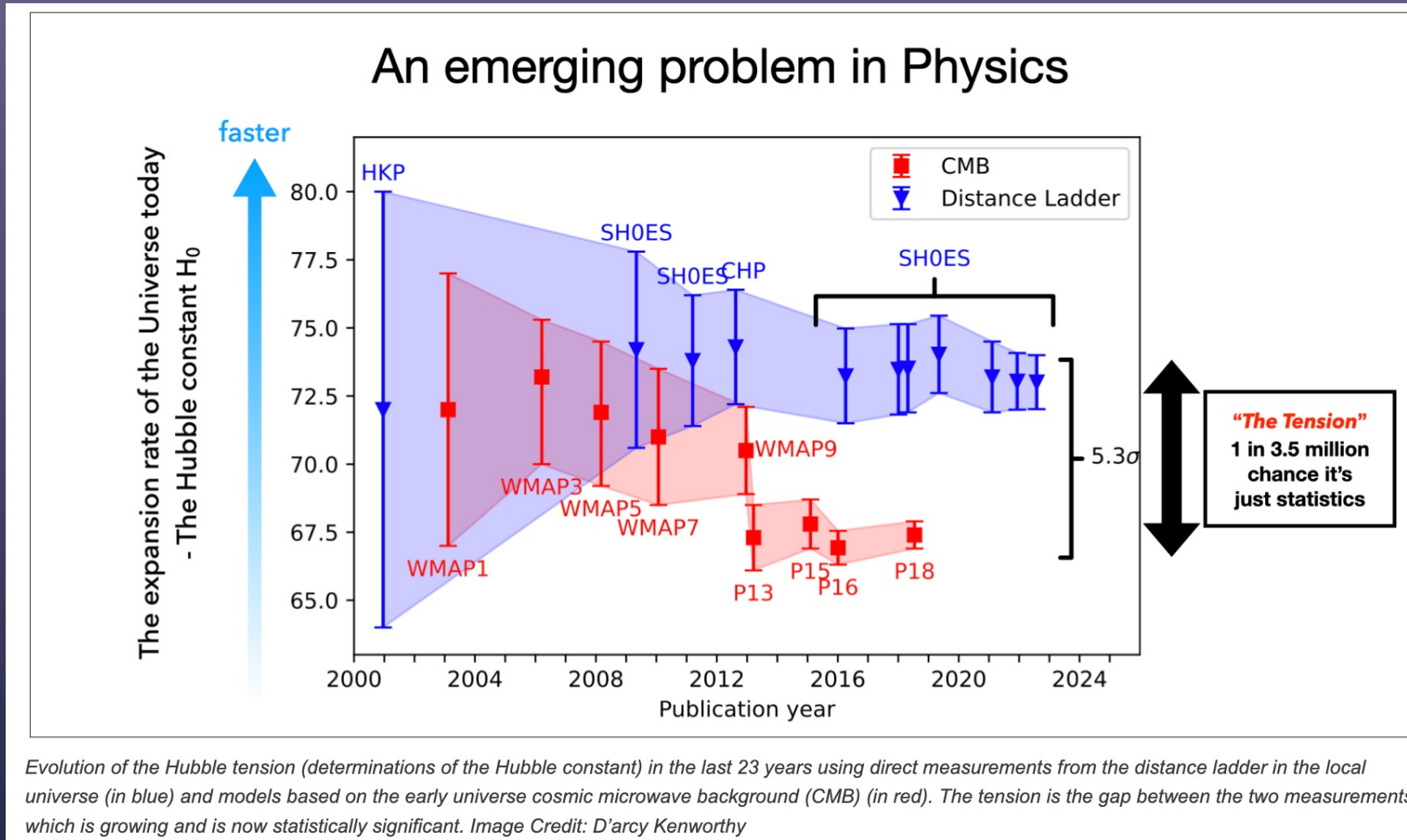
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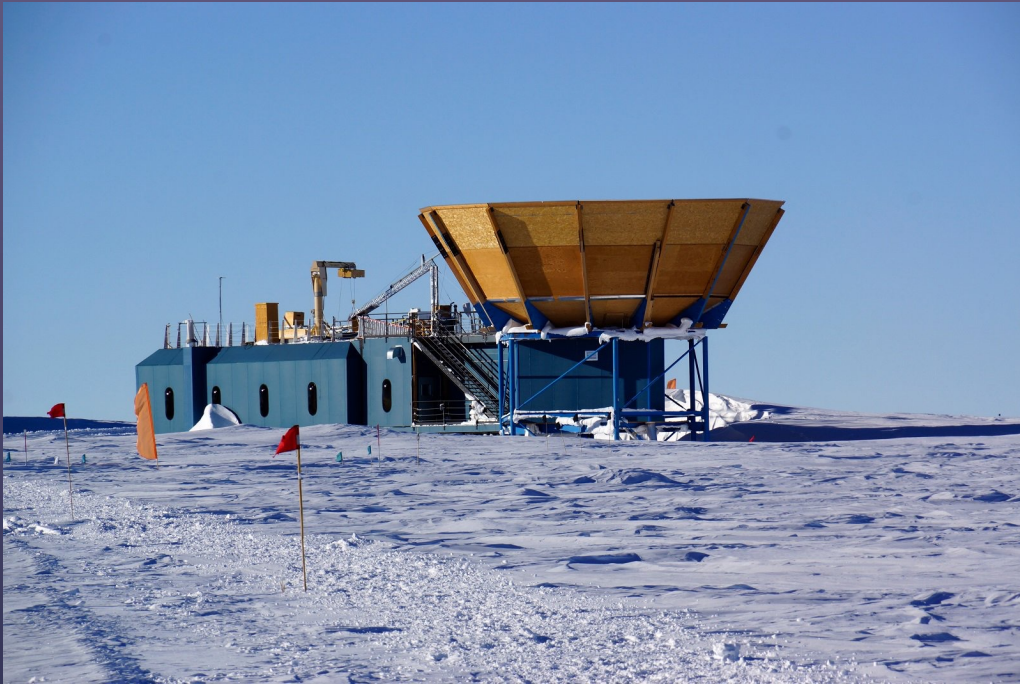


DESI measures baryon acoustic oscillations and combined with other datasets sees hints of evolving dark energy

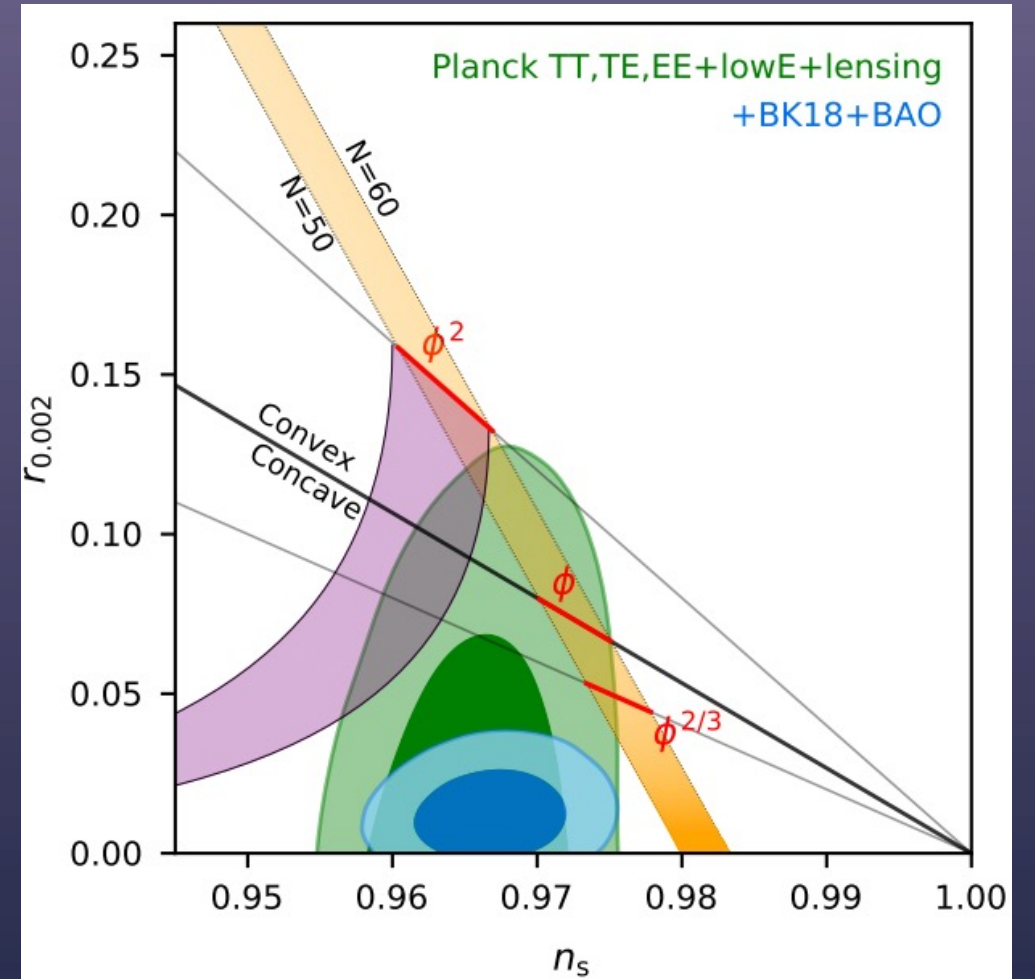
The Hubble Tension



(Lack of) primordial gravitational waves

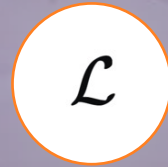


BICEP Array



What are the laws of nature

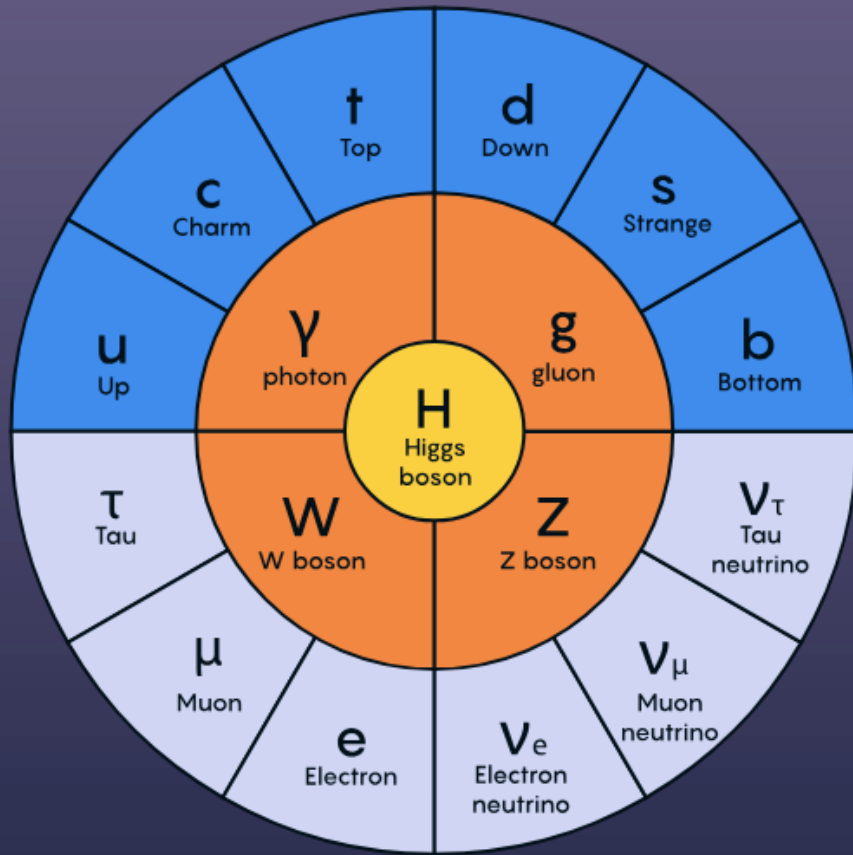
Standard Model of Particle Physics



Λ CDM Model of Cosmology

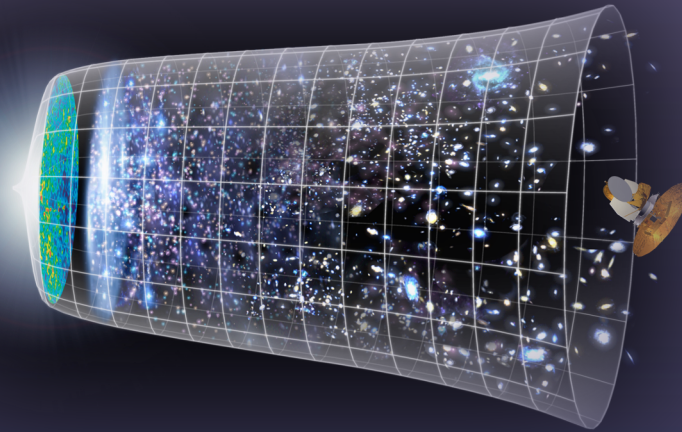
that govern the evolution of our Universe

The Standard Model



Λ CDM Model of Cosmology

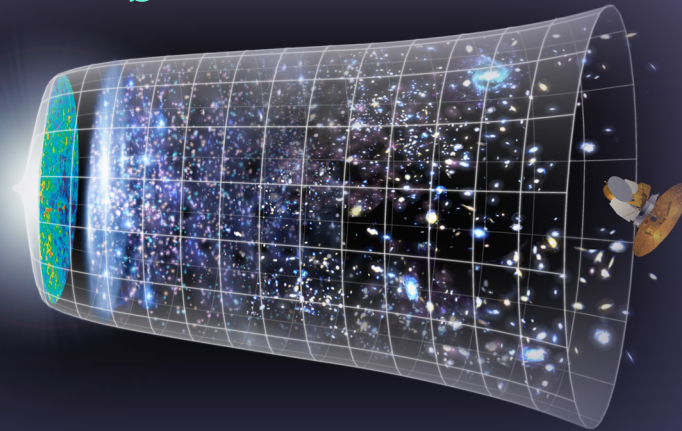
six parameters



Λ CDM Model of Cosmology

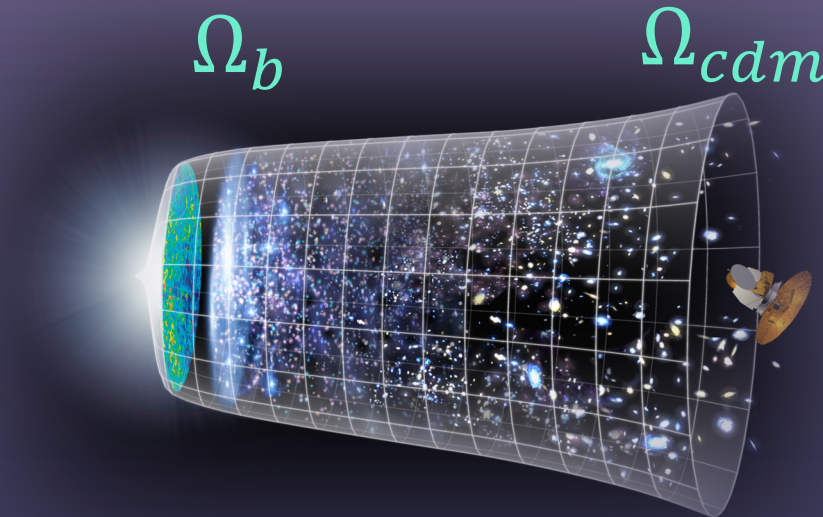
six parameters

Ω_b



Λ CDM Model of Cosmology

six parameters



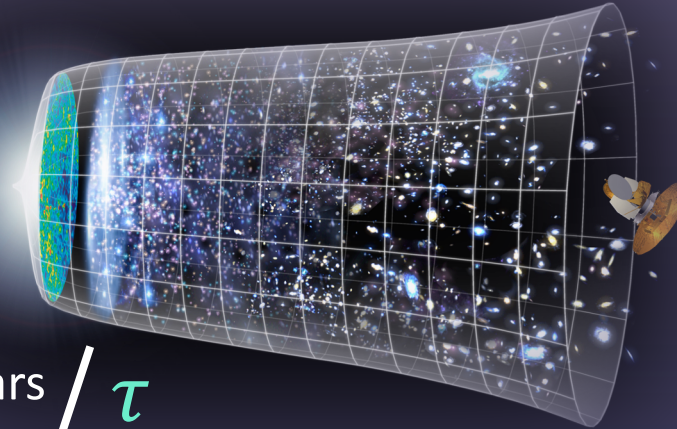
Λ CDM Model of Cosmology

six parameters

Ω_b

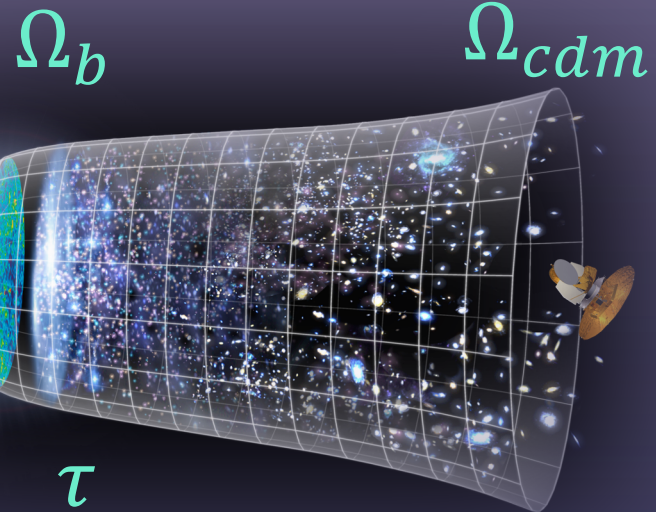
Ω_{cdm}

1 billion years / τ
star formation



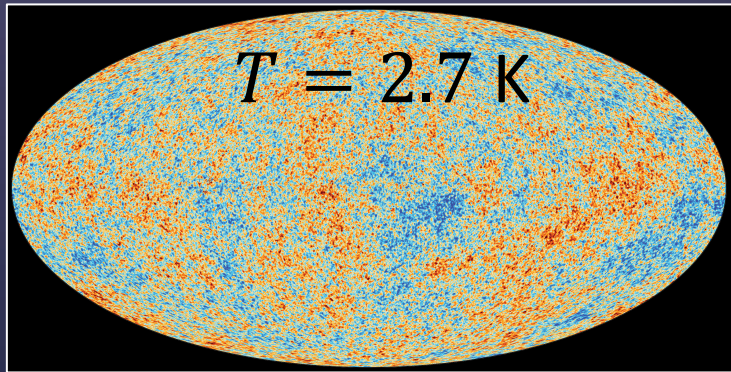
Λ CDM Model of Cosmology

six parameters



400.000 years

cosmic microwave background (CMB)

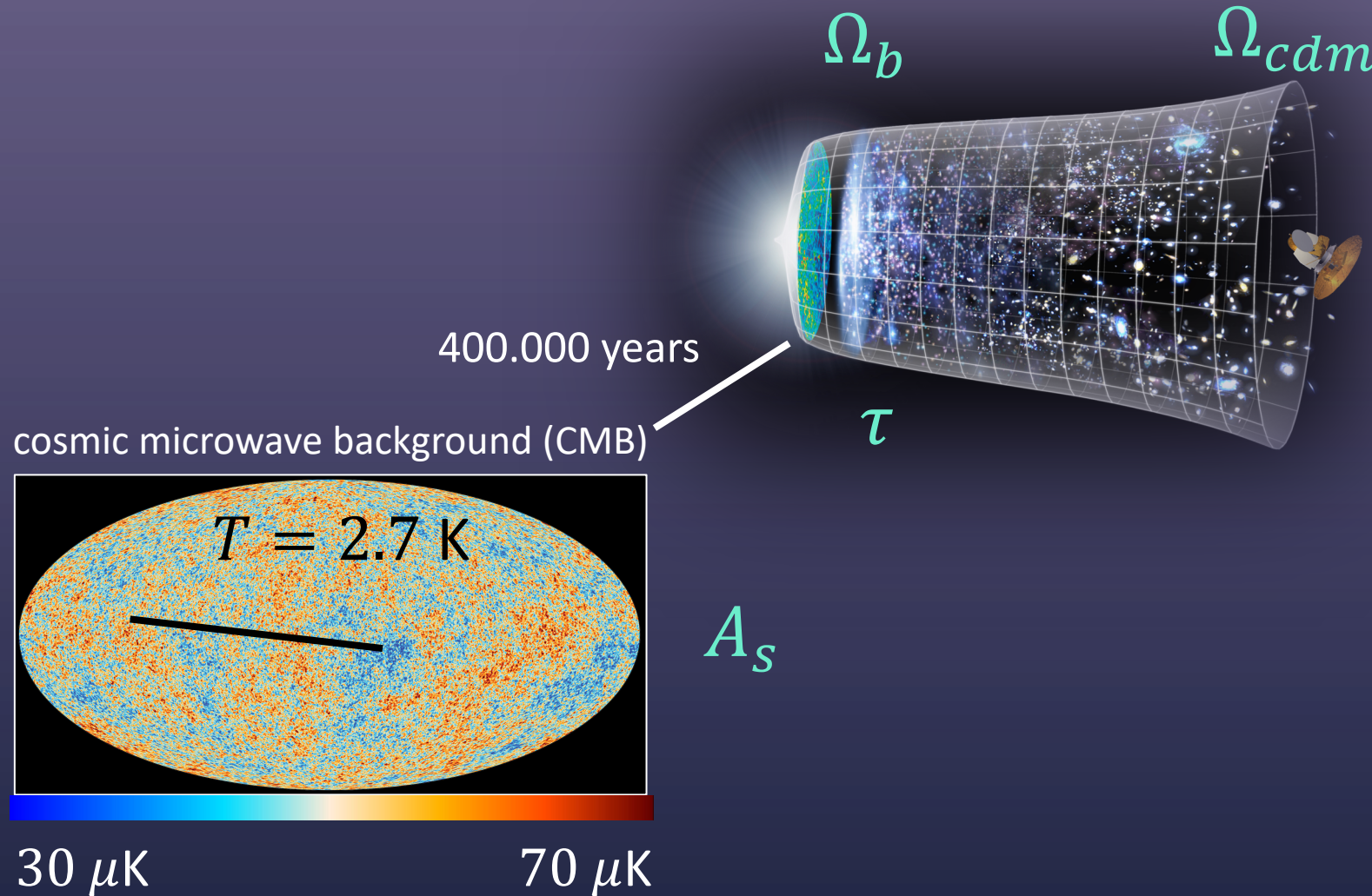


$30 \mu\text{K}$

$70 \mu\text{K}$

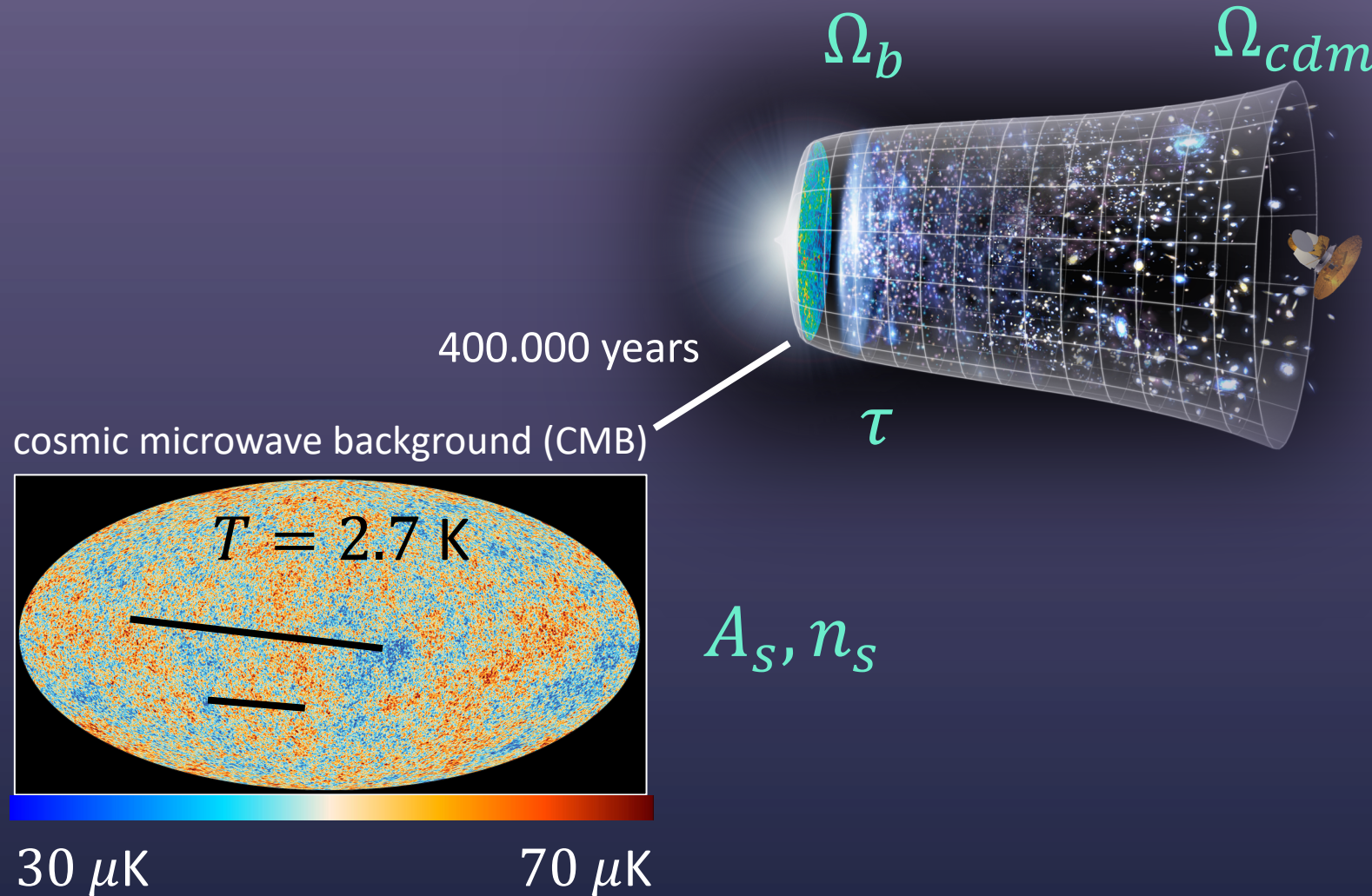
Λ CDM Model of Cosmology

six parameters



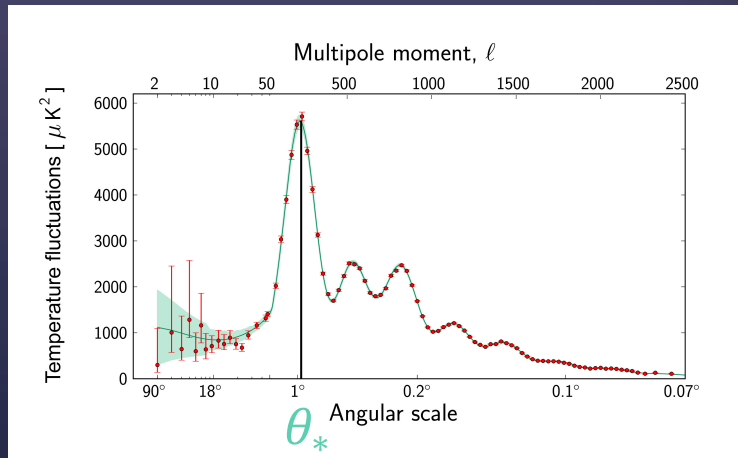
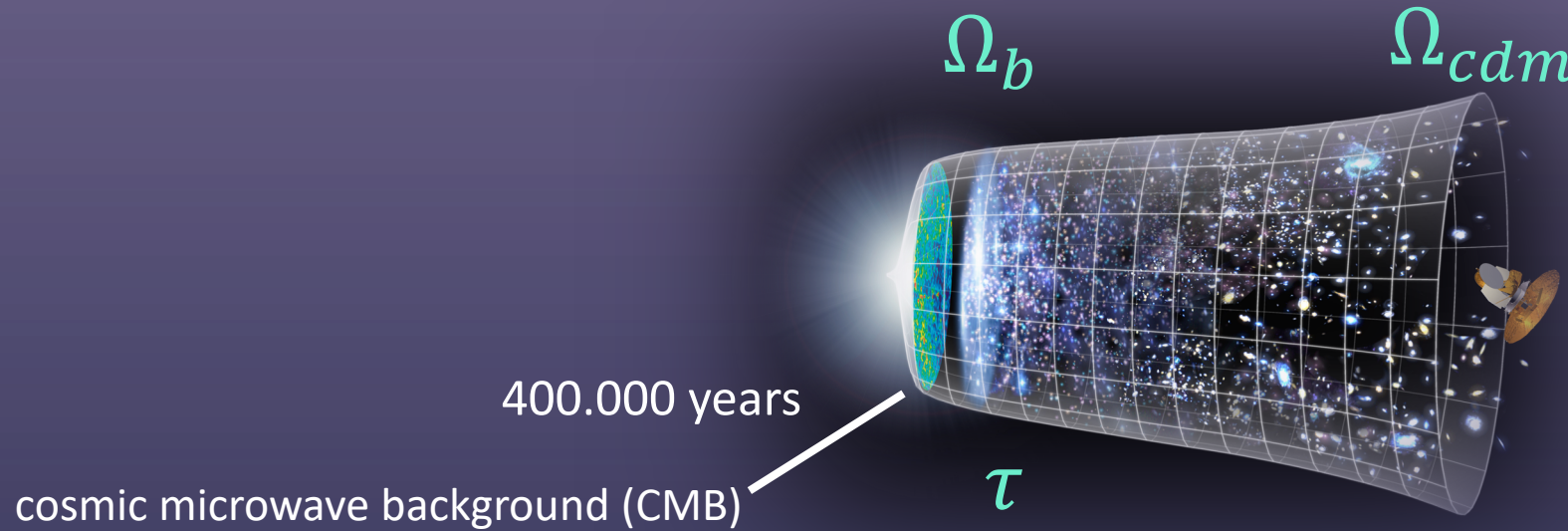
Λ CDM Model of Cosmology

six parameters



Λ CDM Model of Cosmology

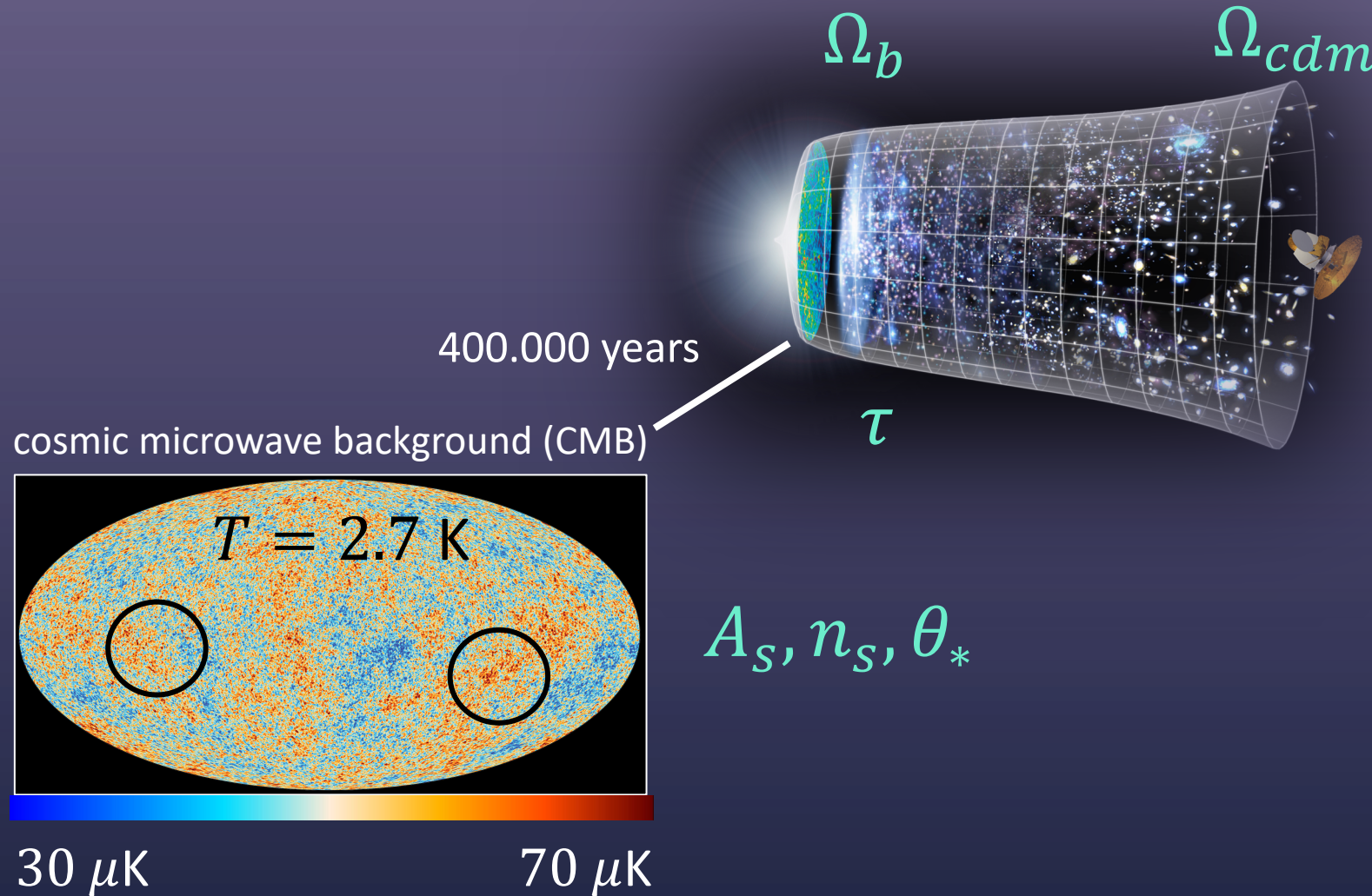
six parameters



A_s, n_s, θ_*

Λ CDM Model of Cosmology

six parameters



Λ CDM Model of Cosmology

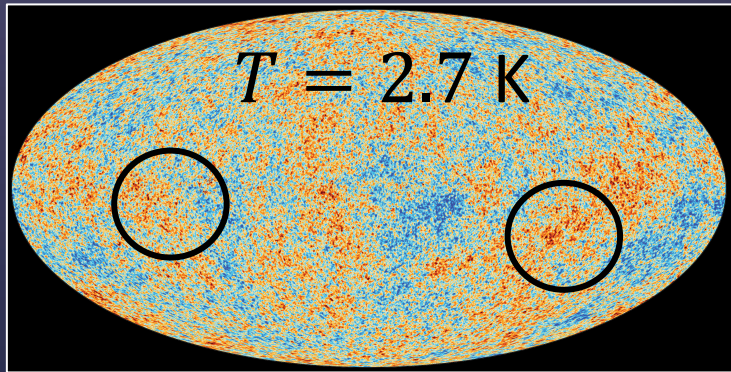
cosmic inflation

six parameters

$\ll s$

400.000 years

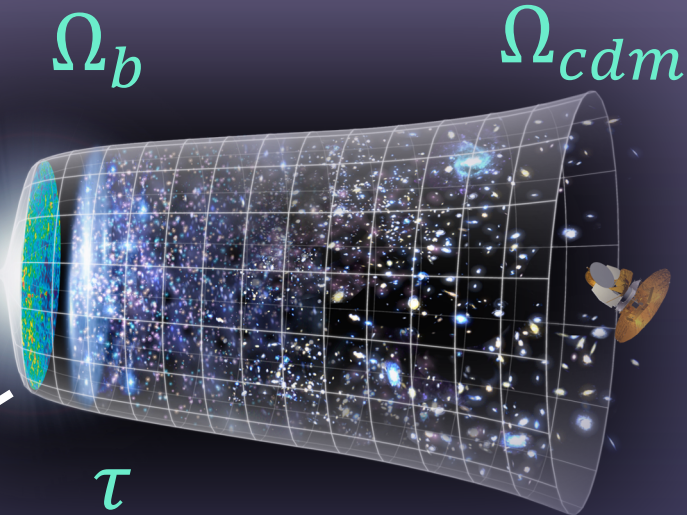
cosmic microwave background (CMB)



$30 \mu\text{K}$

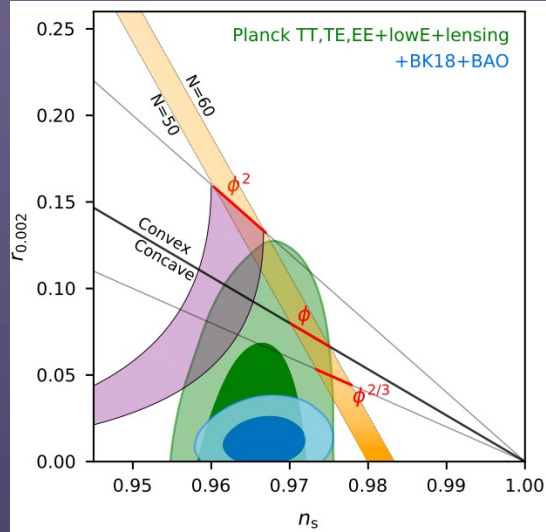
$70 \mu\text{K}$

A_s, n_s, θ_*



lack of inflationary
gravitational waves

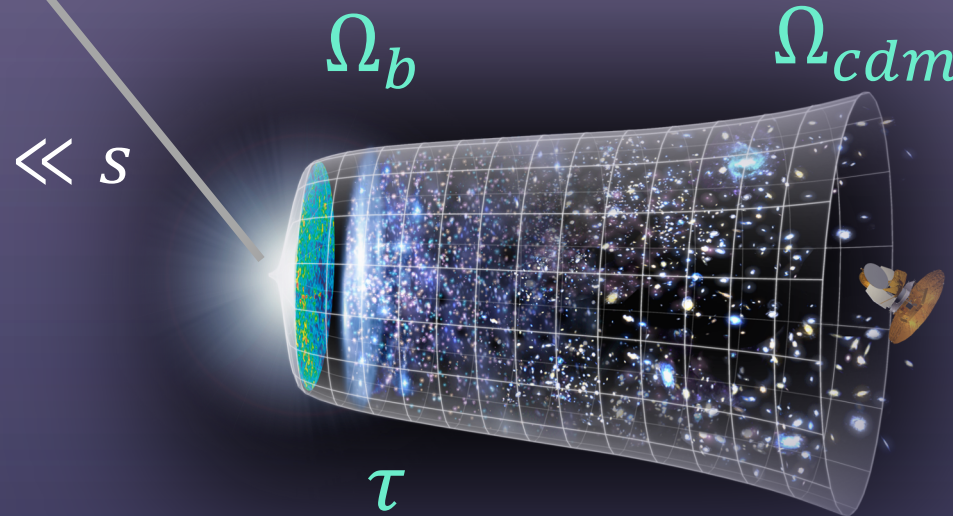
cosmic inflation



Bicep/Keck XIII, 2021

Λ CDM Model of Cosmology

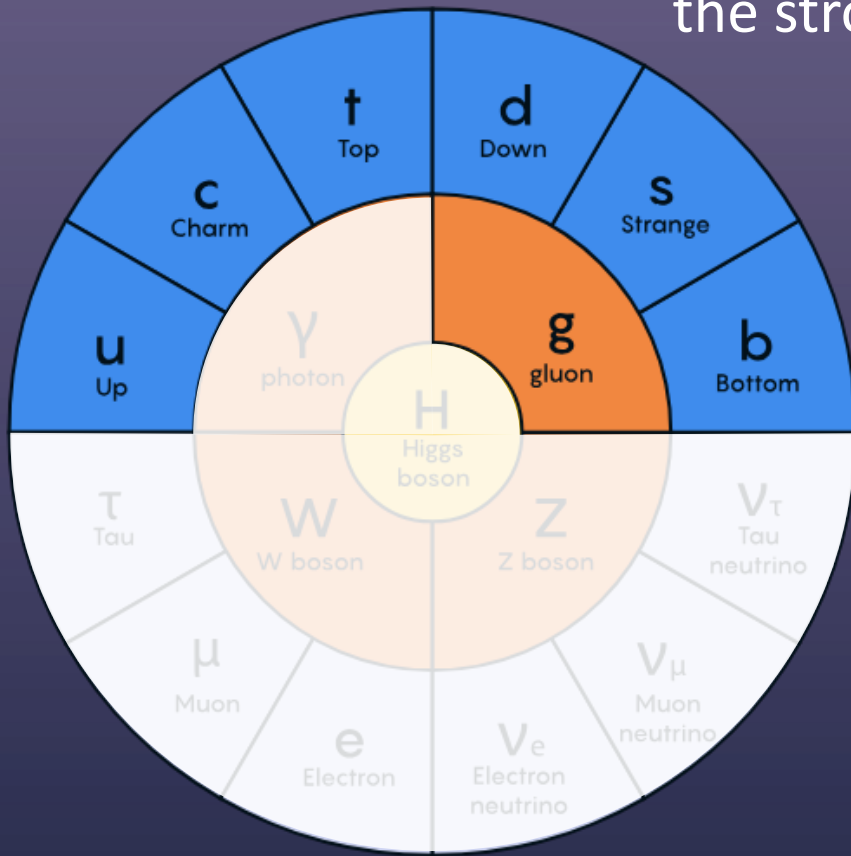
six parameters



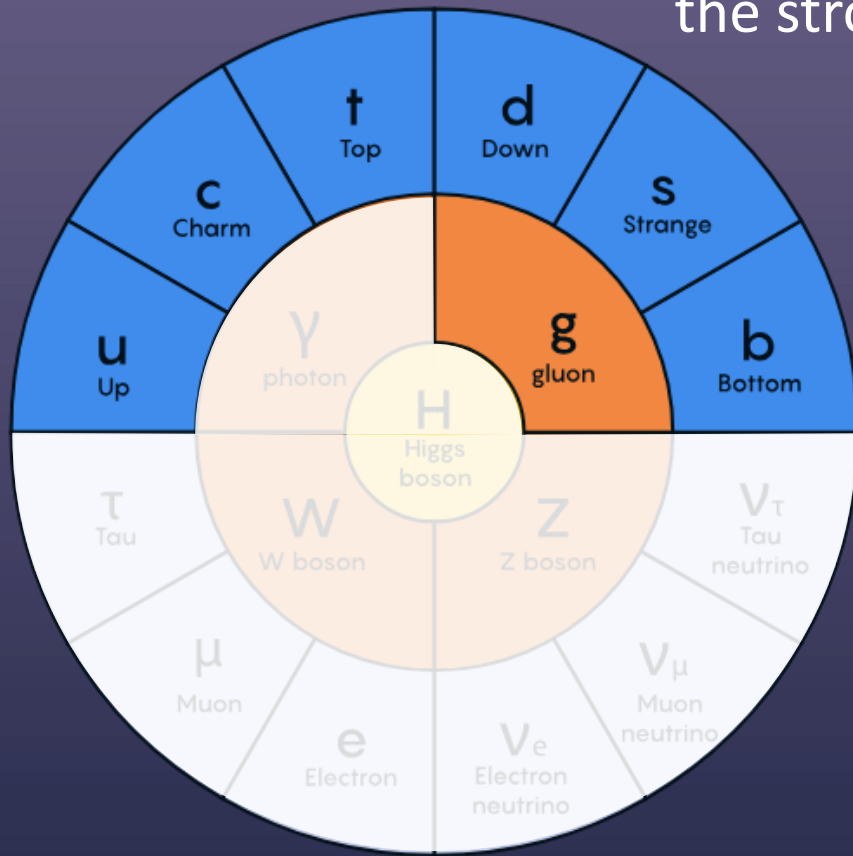
$$A_s, n_s, \theta_*$$

The Standard Model

the strong force

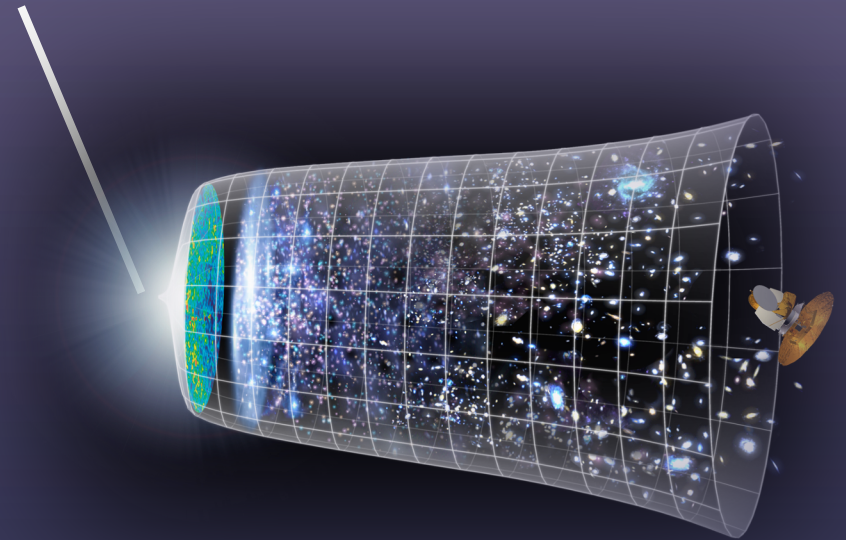


The Standard Model

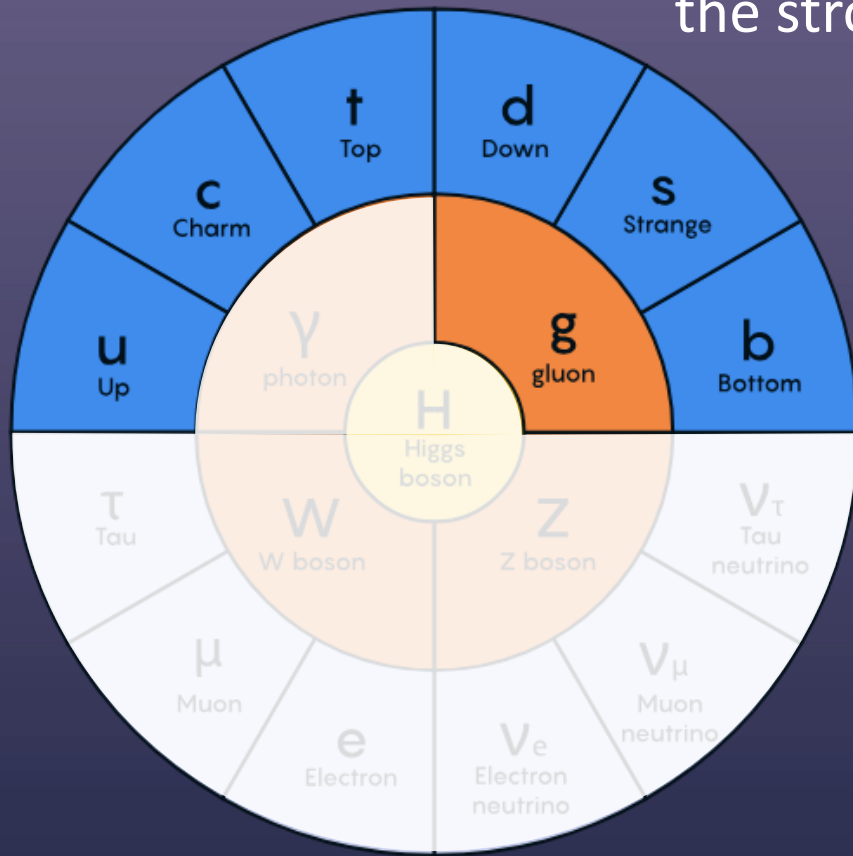


the strong force

cosmic inflation

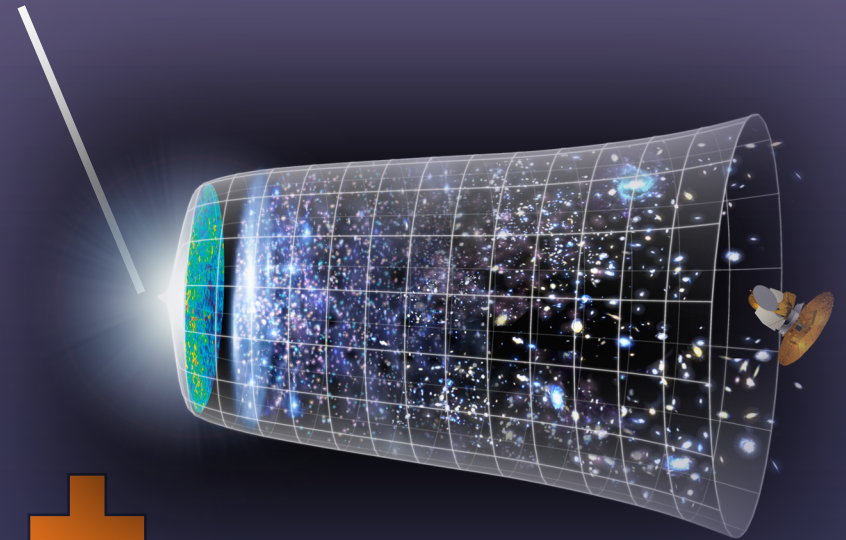
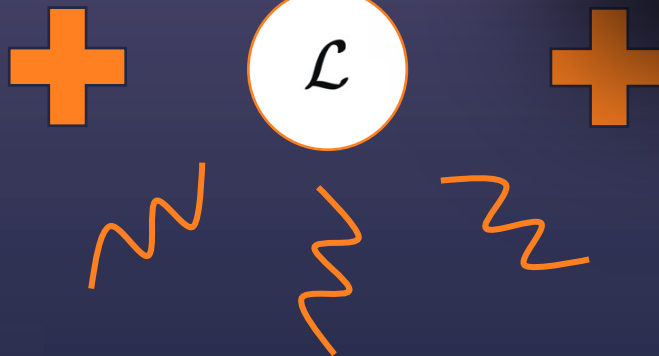


The Standard Model

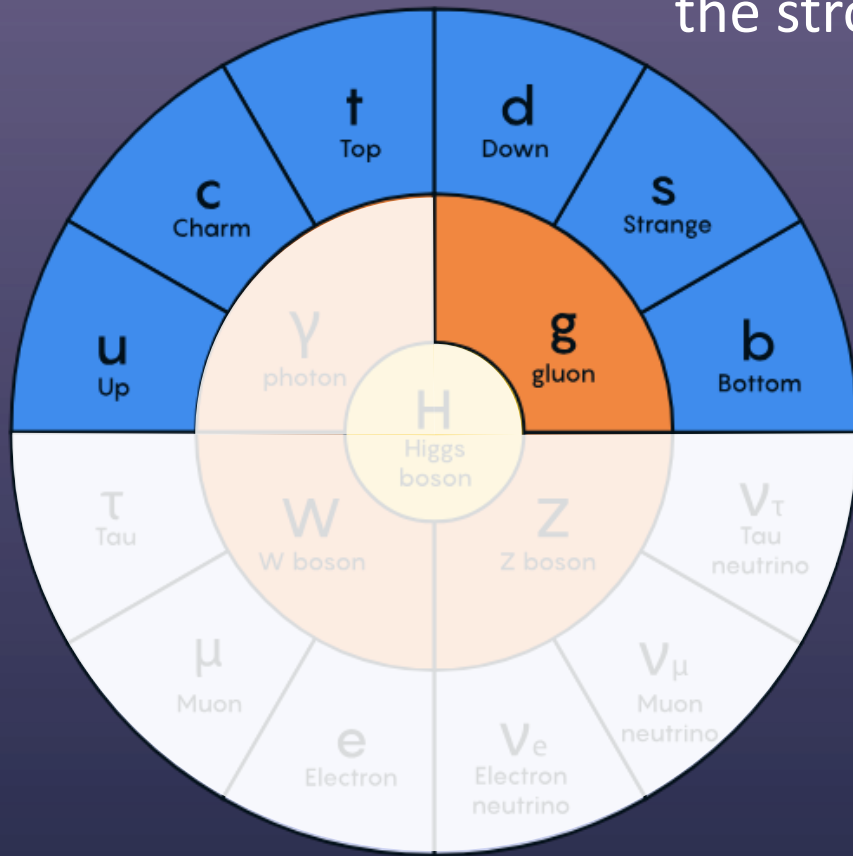


the strong force

cosmic inflation



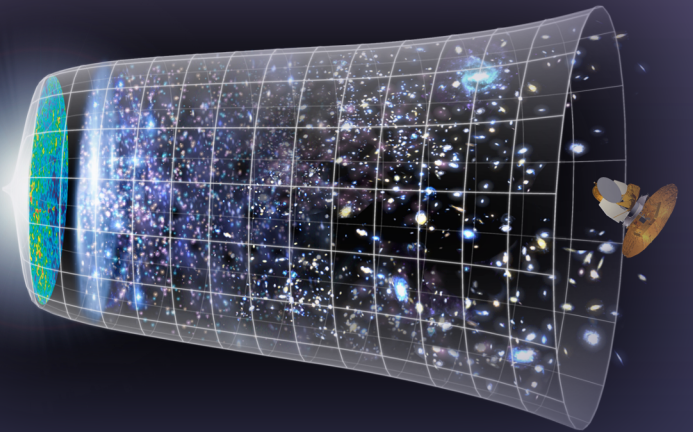
The Standard Model



the strong force

cosmic inflation

$$\mathcal{L}_{\text{int}} = -\phi \frac{\alpha_s}{16\pi f} \tilde{G} G$$



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Warm Inflation with the Standard Model

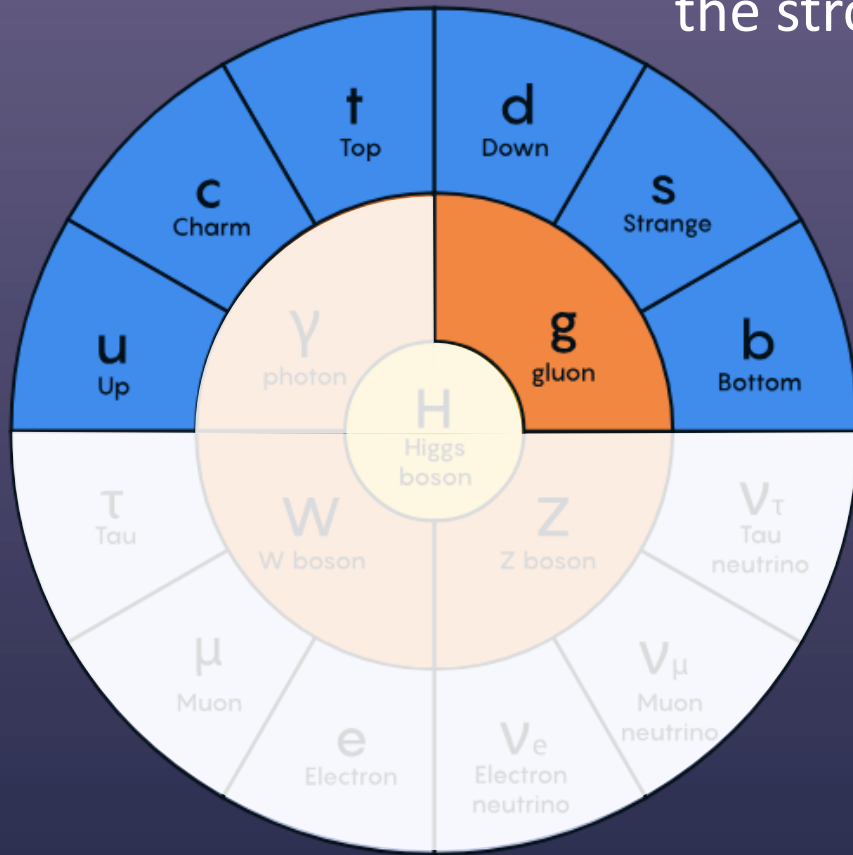
[Kim V. Berghaus](#) ¹, [Marco Drewes](#) ^{2,3}, and [Sebastian Zell](#) ^{4,5,2}

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Phys. Rev. Lett. **135**, 171002 – Published 23 October, 2025

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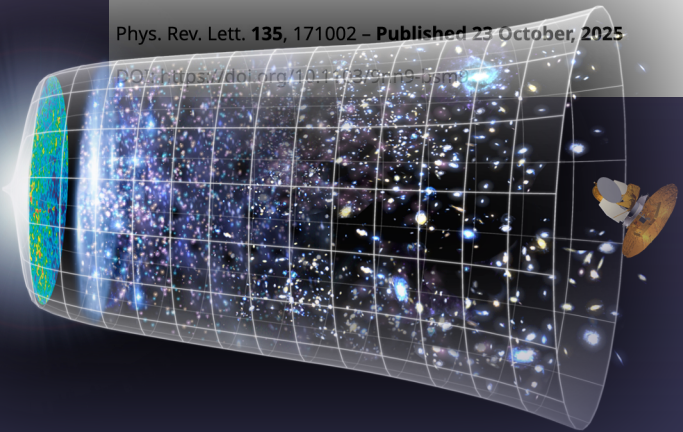
The Standard Model



the strong force

cosmic inflation

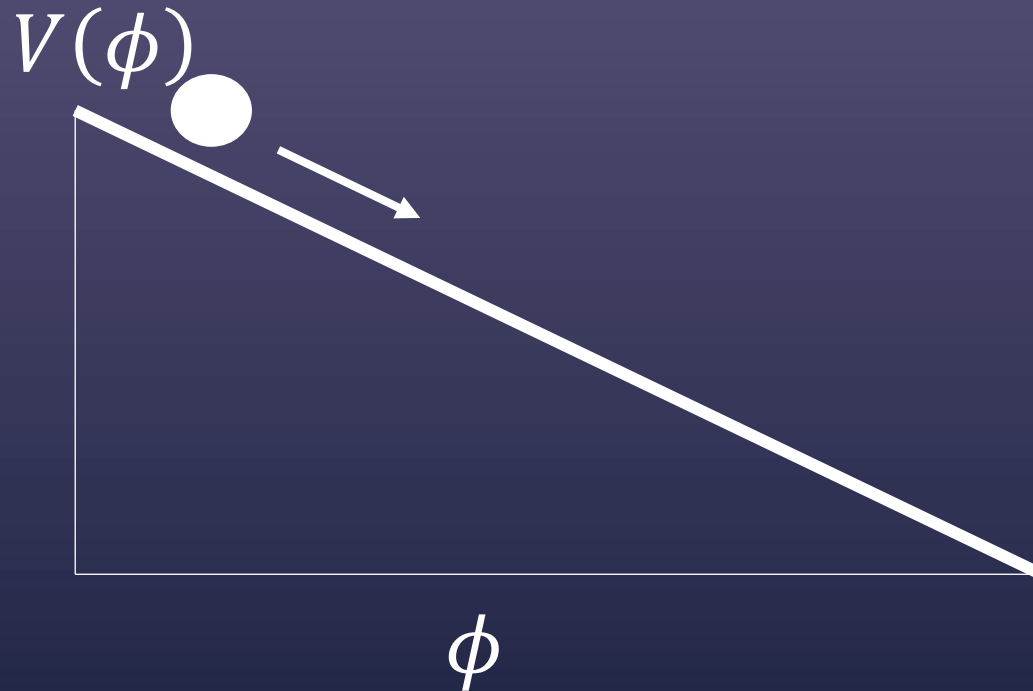
$$\mathcal{L}_{\text{int}} = -\phi \frac{\alpha_s}{16\pi f} \tilde{G} G$$



Inflation

General relativity: content of our Universe dictates its expansion

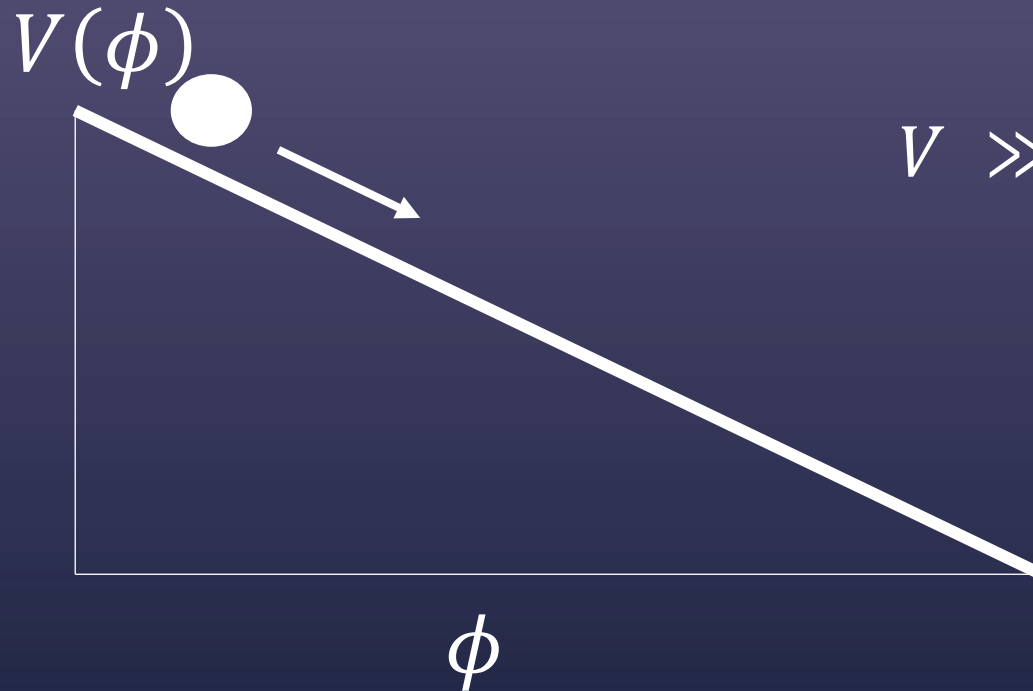
Condition for accelerated expansion: $\frac{P}{\rho} < -\frac{1}{3}$ (negative pressure, quite exotic)



Inflation

General relativity: content of our Universe dictates its expansion

Condition for accelerated expansion: $\frac{P}{\rho} < -\frac{1}{3}$ (negative pressure, quite exotic)



$$V \gg \frac{1}{2} \dot{\phi}^2$$

$$\ddot{\phi} + 3H\dot{\phi} + V' = 0$$



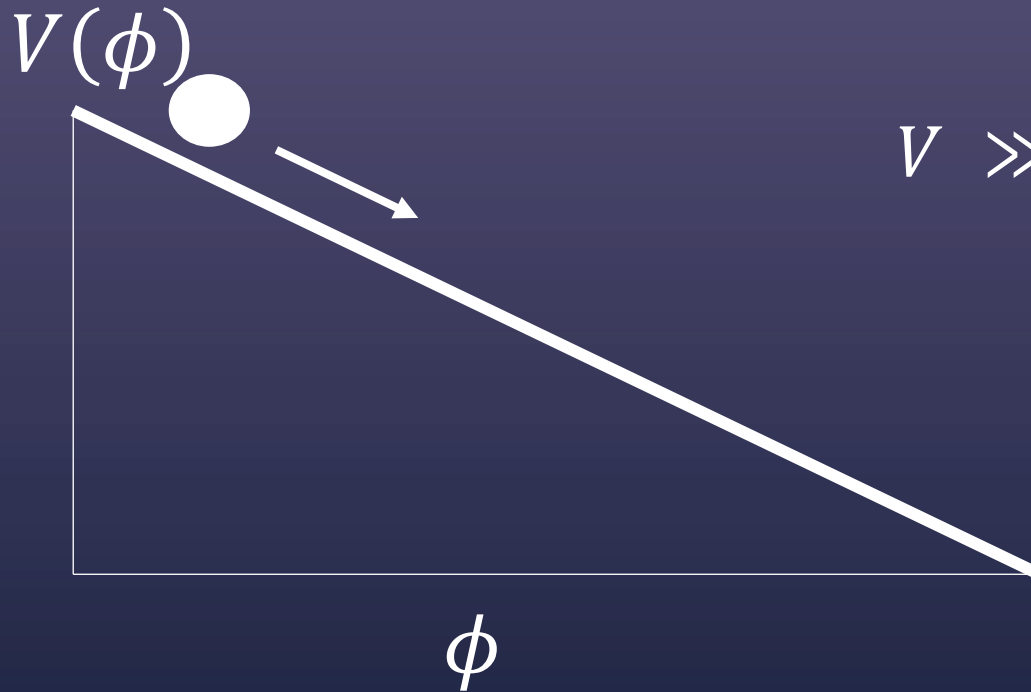
Hubble friction
from expansion

Inflation

$$\varepsilon_V = \frac{M_{Pl}^2}{2} \left(\frac{V'}{V} \right)^2 \ll 1$$

General relativity: content of our Universe dictates its expansion

Condition for accelerated expansion: $\frac{P}{\rho} < -\frac{1}{3}$ (negative pressure, quite exotic)



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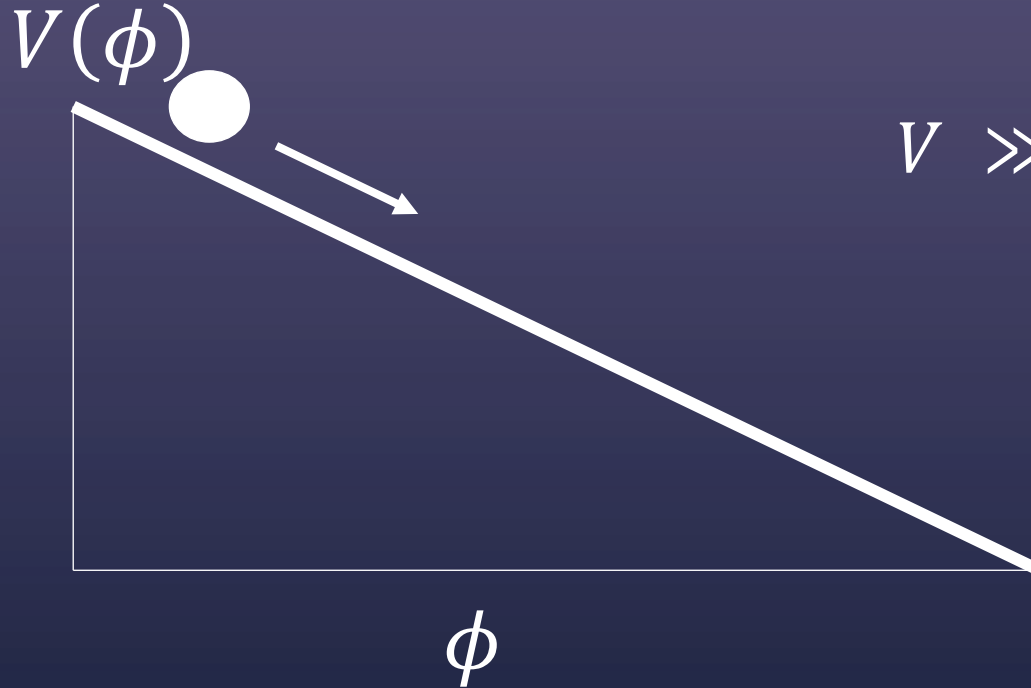
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$$V \gg \frac{1}{2} \dot{\phi}^2$$

$$\ddot{\phi} + 3H\dot{\phi} + V' = 0$$



Hubble friction
from expansion

$$\langle \delta\phi \delta\phi \rangle \propto A_s$$

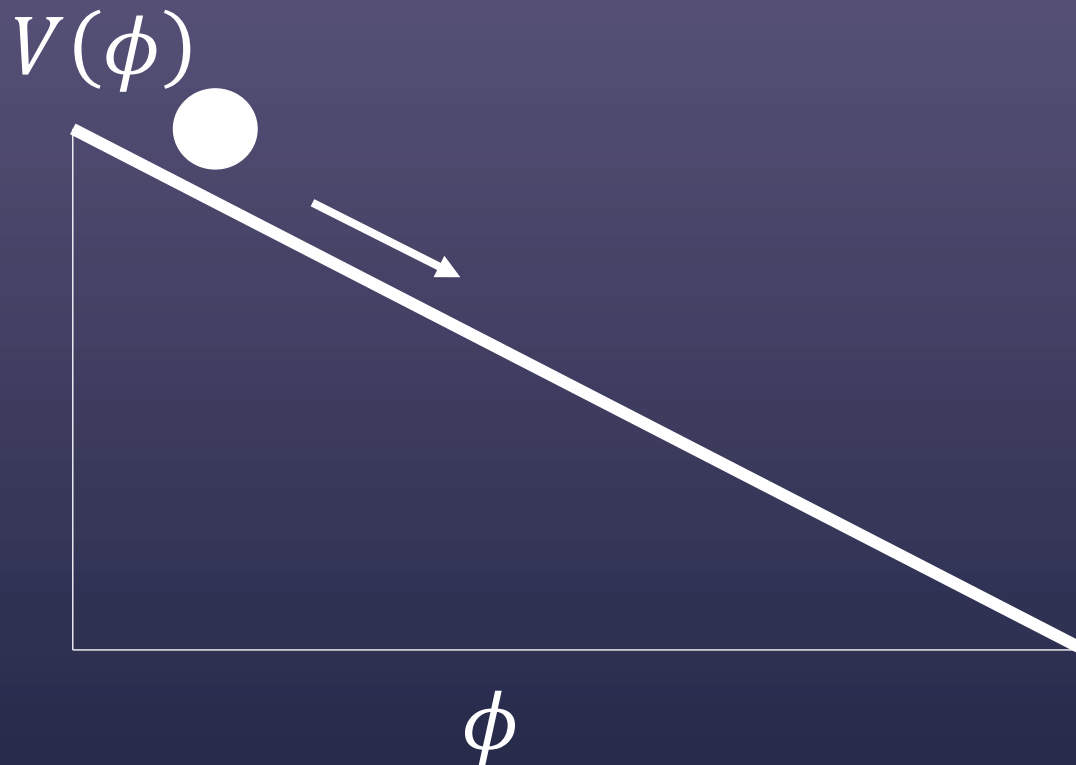
Inflation

$$\varepsilon_V = \frac{M_{Pl}^2}{2} \left(\frac{V'}{V} \right)^2 \ll 1$$

$$\ddot{\phi} + 3H\dot{\phi} + V' = 0$$

$$\delta\phi \sim H$$

quantum fluctuations
seed structure



Inflation

$$\varepsilon_V = \frac{M_{Pl}^2}{2} \left(\frac{V'}{V} \right)^2 \ll 1$$

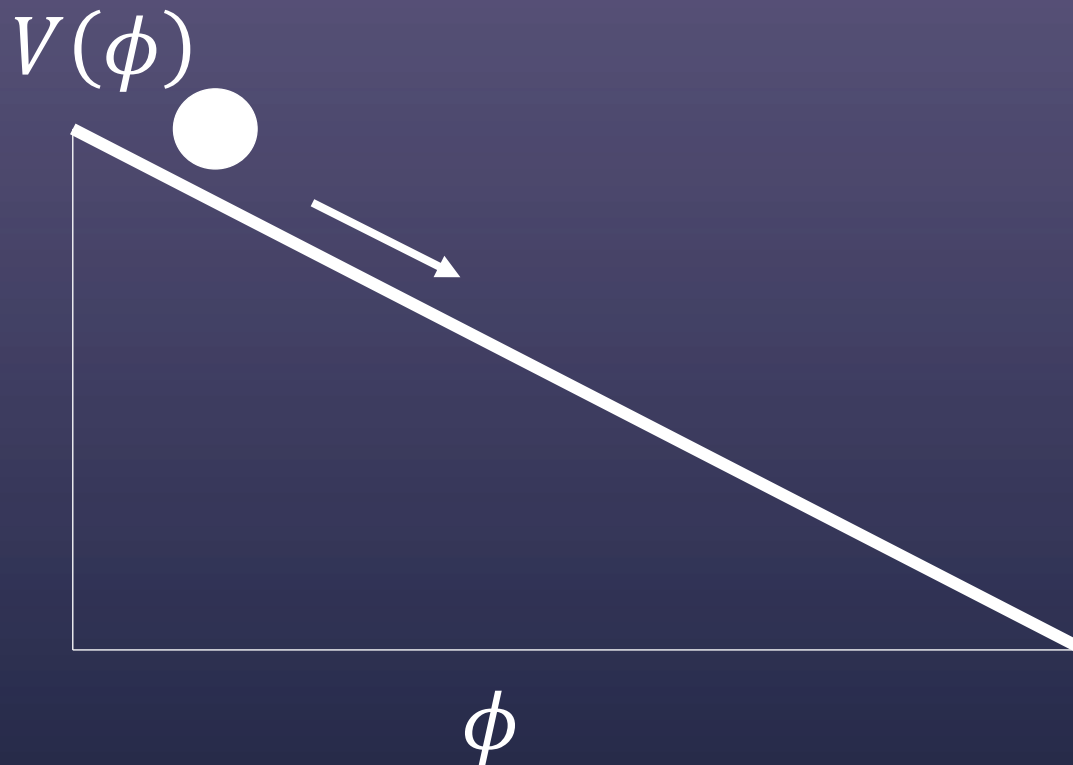
$$\ddot{\phi} + 3H\dot{\phi} + V' = 0$$

$$\delta\phi \sim H$$

quantum fluctuations
seed structure

$$h \sim H$$

quantum fluctuations
seed gravitational waves



Inflation

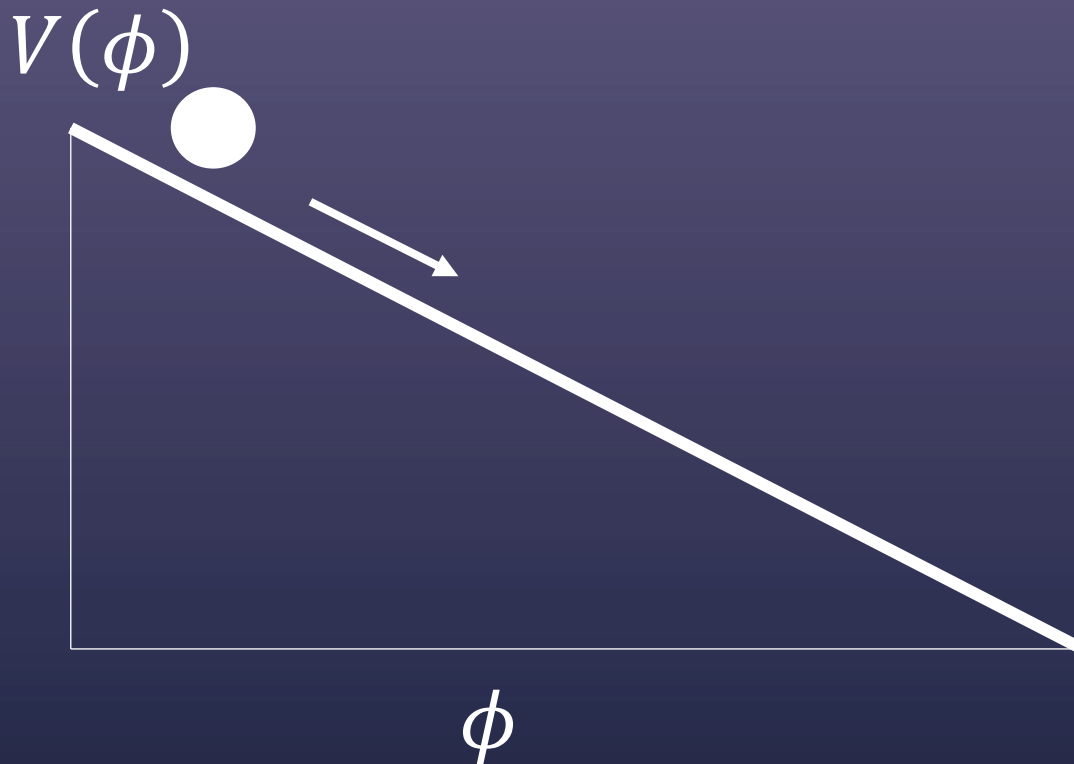
$$\varepsilon_V = \frac{M_{Pl}^2}{2} \left(\frac{V'}{V} \right)^2 \ll 1$$

$$\ddot{\phi} + 3H\dot{\phi} + V' = 0$$

$\delta\phi \sim H$ quantum fluctuations
seed structure

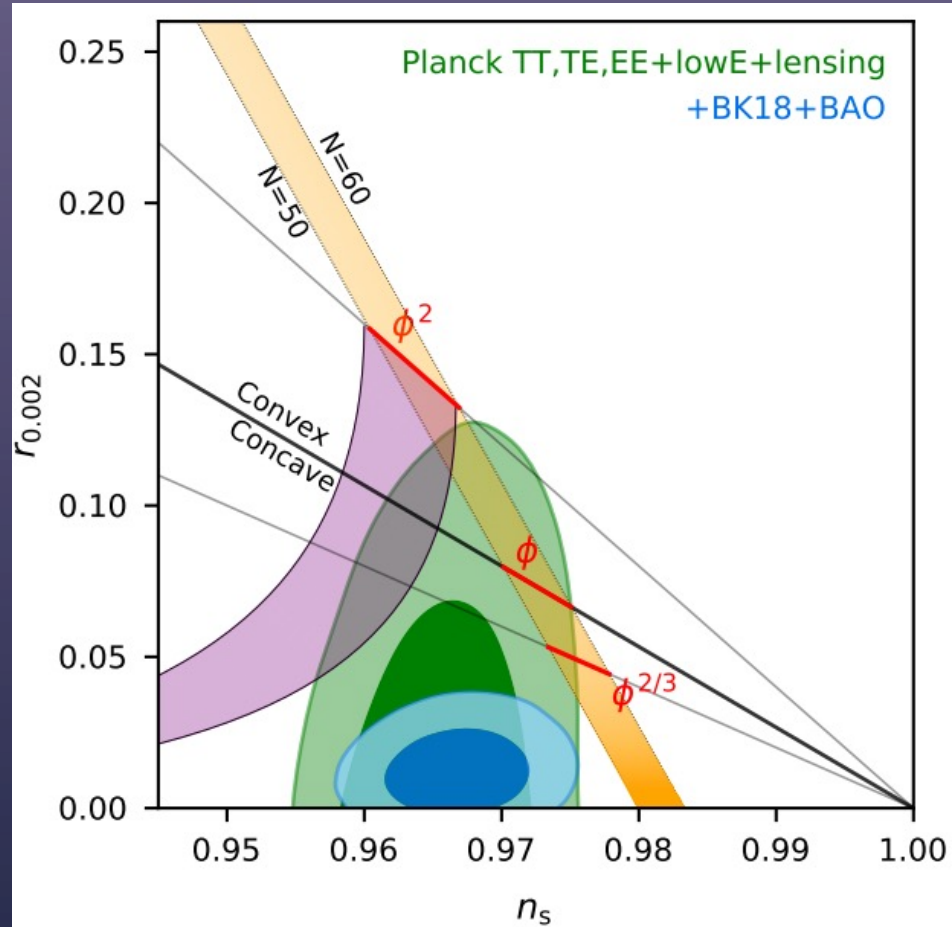
$h \sim H$ quantum fluctuations
seed gravitational waves

$$r \propto \frac{\langle hh \rangle}{\langle \delta\phi \delta\phi \rangle} \approx 16\varepsilon_V$$



Inflation

$$\varepsilon_V = \frac{M_{Pl}^2}{2} \left(\frac{V'}{V} \right)^2 \ll 1$$



Bicep/Keck XIII, 2021

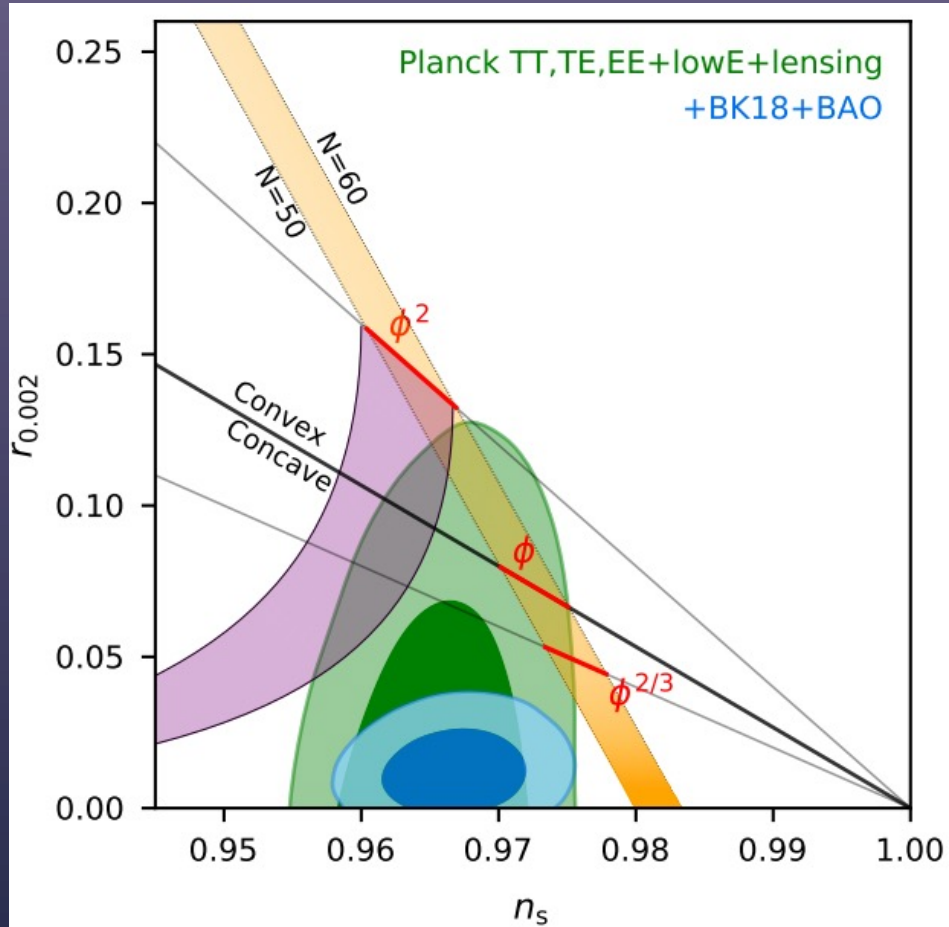
$$\ddot{\phi} + 3H\dot{\phi} + V' = 0$$

$\delta\phi \sim H$ quantum fluctuations
seed structure

$h \sim H$ quantum fluctuations
seed gravitational waves

$$r \propto \frac{\langle hh \rangle}{\langle \delta\phi \delta\phi \rangle} \approx 16\varepsilon_V$$

Inflation



Bicep/Keck XIII, 2021

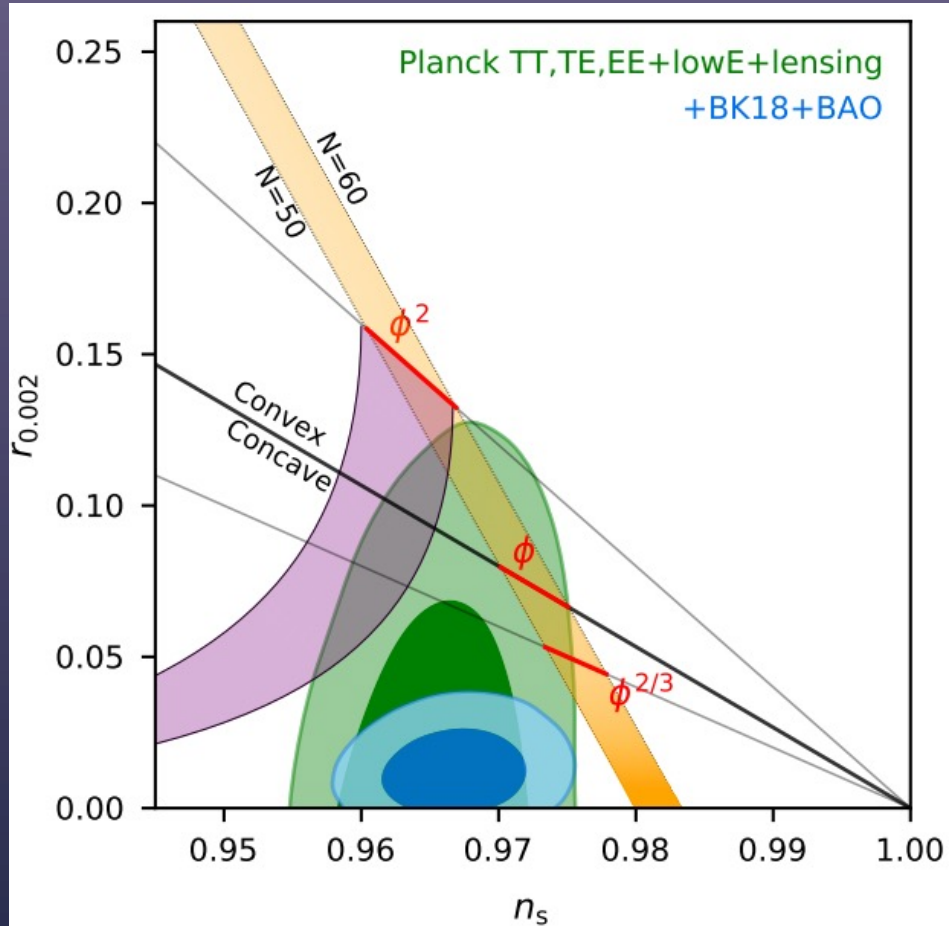


modified gravity

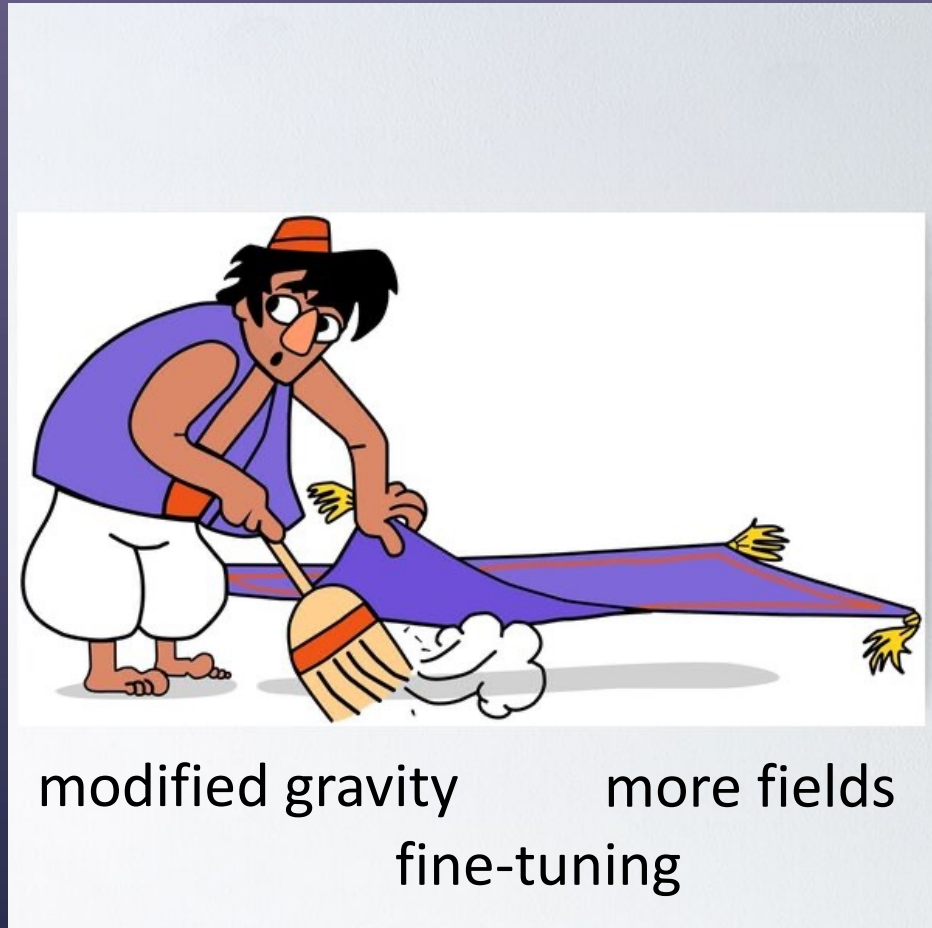
more fields

fine-tuning

Inflation



Bicep/Keck XIII, 2021



Let's include interactions

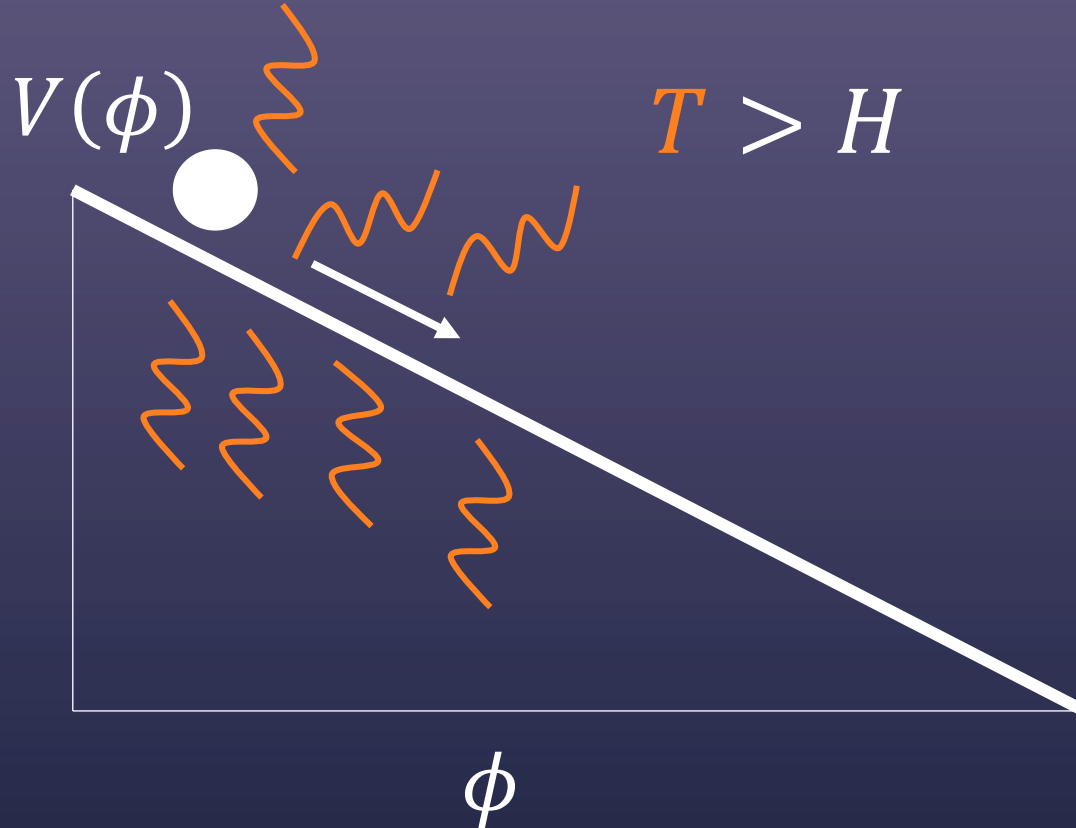
Warm Inflation

Warm Inflation

$$\varepsilon_V = \frac{M_{Pl}^2}{2(1 + \gamma/3H)} \left(\frac{V'}{V} \right)^2 \ll 1$$

consider particle interactions

$$\ddot{\phi} + 3(H + \gamma)\dot{\phi} + V' = 0$$



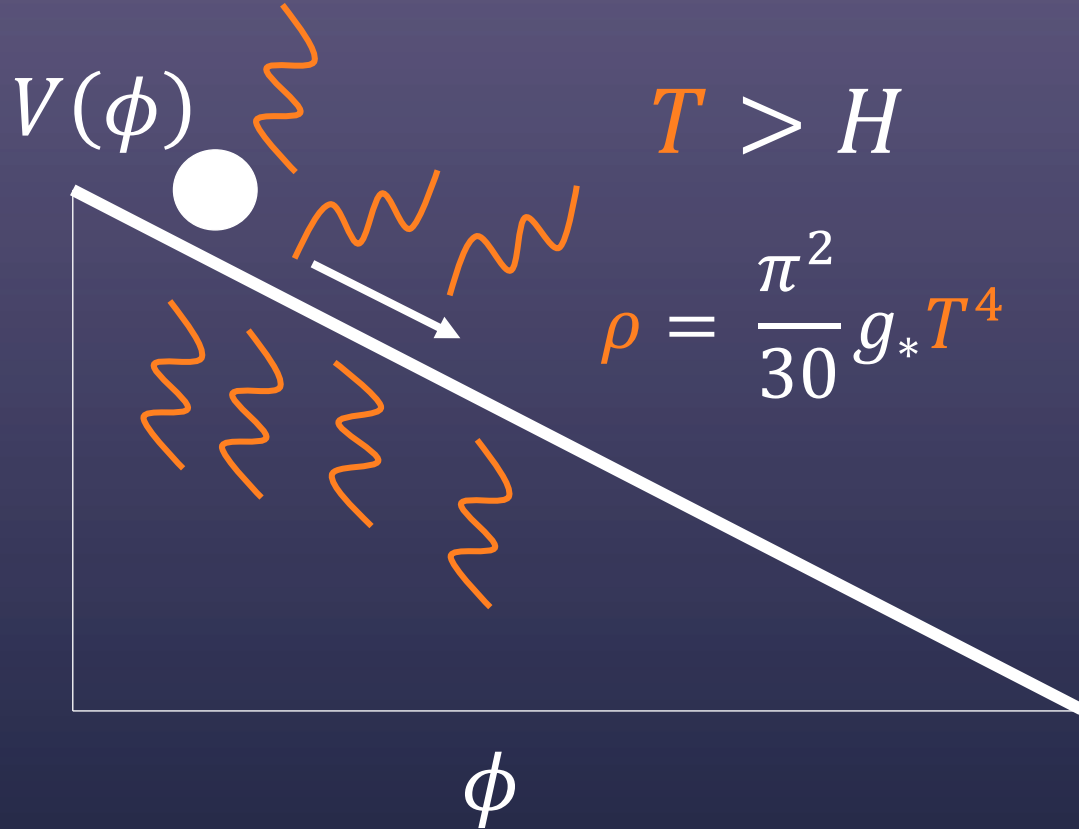
Warm Inflation

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consider particle interactions

$$\ddot{\phi} + 3(H + \gamma)\dot{\phi} + V' = 0$$

$$\dot{\rho} + 4H\rho = \gamma\dot{\phi}^2$$

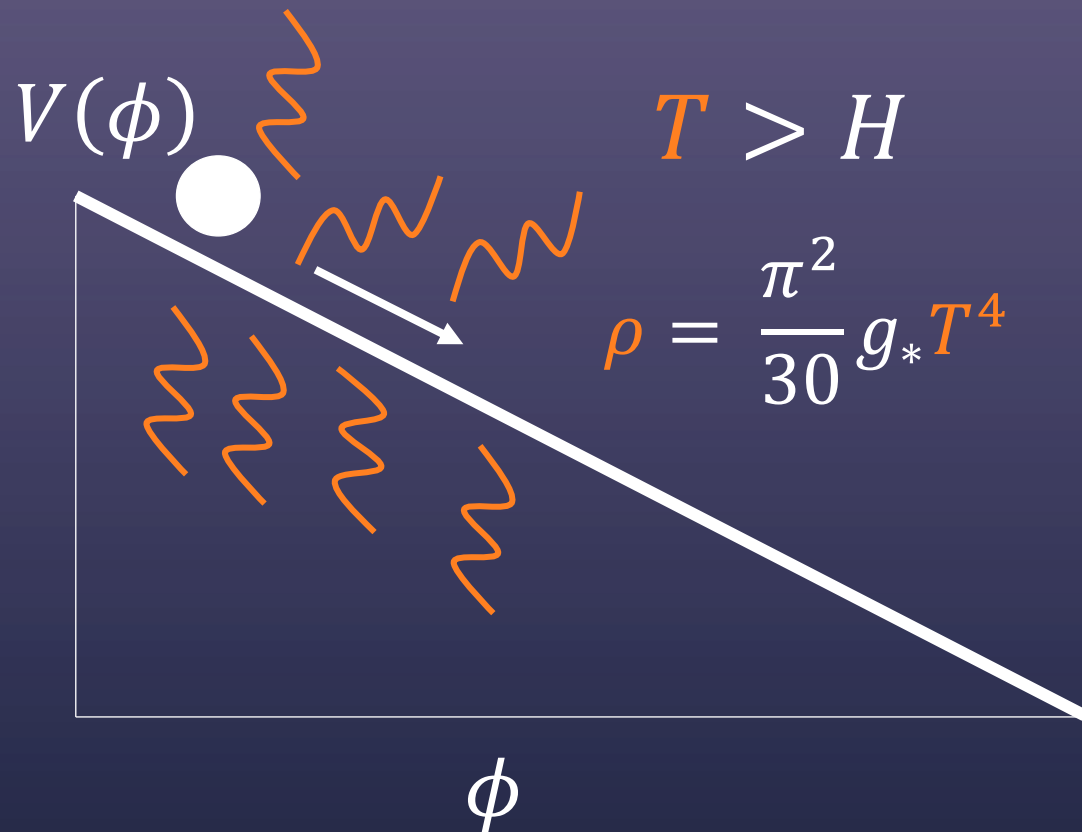


Warm Inflation

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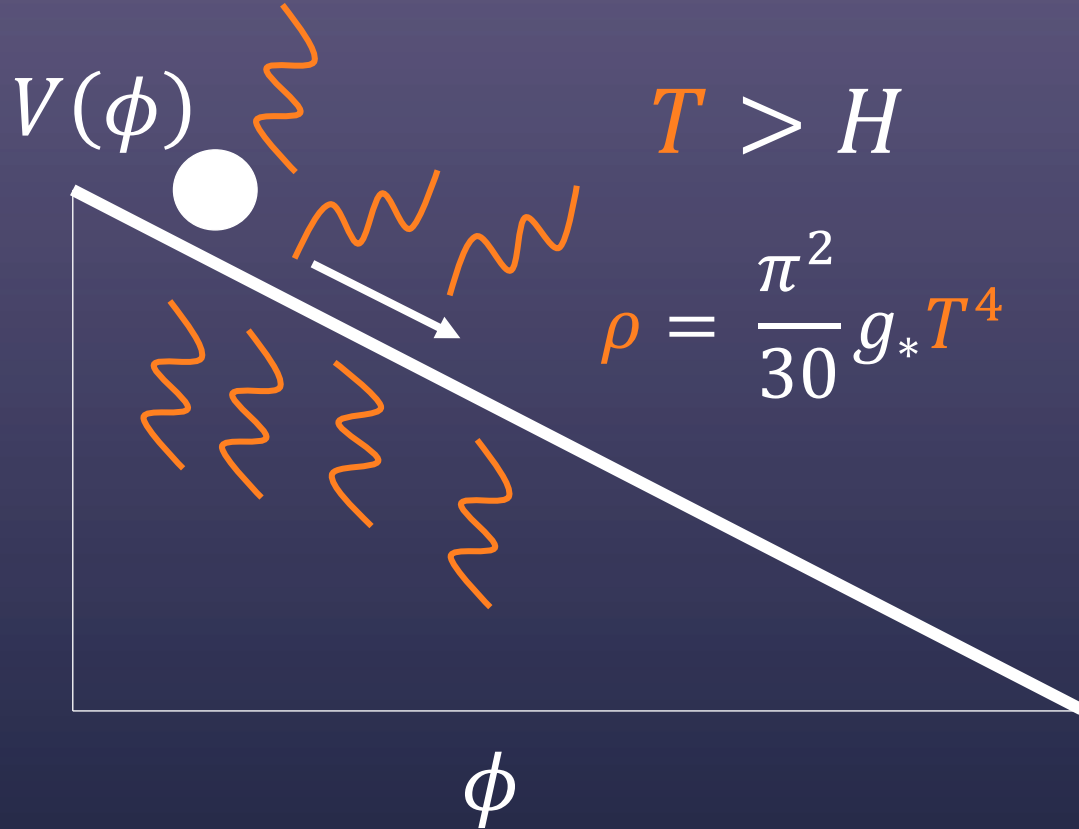
$$\ddot{\phi} + 3(H + \gamma)\dot{\phi} + V' = 0$$



$$\underbrace{\dot{\rho} + 4H\rho}_{\text{sink term}} = \underbrace{\gamma\dot{\phi}^2}_{\text{source term}}$$

Warm Inflation

consider particle interactions



$$\varepsilon_V = \frac{M_{Pl}^2}{2(1 + \Upsilon/3H)} \left(\frac{V'}{V} \right)^2 \ll 1$$

$$\ddot{\phi} + 3(H + \Upsilon)\dot{\phi} + V' = 0$$

$$\dot{\rho} + 4H\rho = \Upsilon\dot{\phi}^2$$

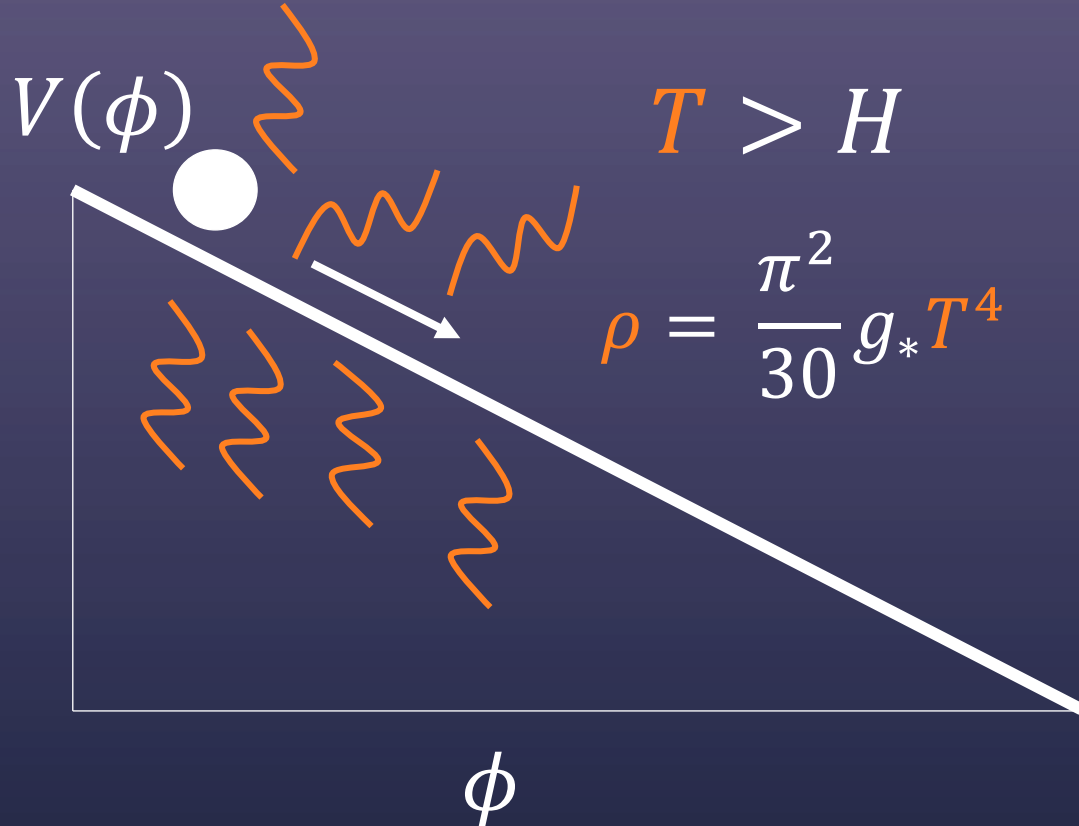
sink term \approx source term

$$\delta\phi \sim \sqrt{HT}$$

classical fluctuations
seed structure

Warm Inflation

consider particle interactions



$$\varepsilon_V = \frac{M_{Pl}^2}{2(1 + \gamma/3H)} \left(\frac{V'}{V} \right)^2 \ll 1$$

$$\ddot{\phi} + 3(H + \gamma)\dot{\phi} + V' = 0$$

$$\underbrace{\dot{\rho} + 4H\rho}_{\text{sink term}} = \underbrace{\gamma\dot{\phi}^2}_{\text{source term}}$$

sink term \approx source term

$$\delta\phi \sim \sqrt{HT}$$

classical fluctuations
seed structure

$$r \propto \frac{\langle hh \rangle}{\langle \delta\phi \delta\phi \rangle} \ll 16\varepsilon_V$$

Warm Inflation from first principles

- couple inflaton to light degrees of freedom $\mathcal{L}_{\text{int}} = -\phi J_{\text{int}}$

$$\ddot{\phi} + 3H\dot{\phi} + V' = -\langle J_{\text{int}} \rangle$$

Warm Inflation from first principles

- couple inflaton to light degrees of freedom $\mathcal{L}_{\text{int}} = -\phi J_{\text{int}}$

$$\ddot{\phi} + 3H\dot{\phi} + V' = -\langle J_{\text{int}} \rangle$$

back-reacts

$$\langle J_{\text{int}} \rangle \approx m_{th}^2 \phi + \Upsilon \dot{\phi} + O(\ddot{\phi})$$

Warm Inflation from first principles

- couple inflaton to light degrees of freedom $\mathcal{L}_{\text{int}} = -\phi J_{\text{int}}$

$$\ddot{\phi} + 3H\dot{\phi} + V' = -\langle J_{\text{int}} \rangle$$

back-reacts

$$\langle J_{\text{int}} \rangle \approx m_{th}^2 \phi + \Upsilon \dot{\phi} + O(\ddot{\phi})$$

Usually $m_{th}^2 \phi \gg \Upsilon \dot{\phi}$

Is warm inflation possible?



Minimal Warm Inflation

- couple inflaton to light degrees of freedom $\mathcal{L}_{\text{int}} = -\phi \frac{\alpha}{16\pi f} \tilde{G} G$



$$\ddot{\phi} + 3H\dot{\phi} + V' = - \left\langle \frac{\alpha}{16\pi f} \tilde{G} G \right\rangle$$

Minimal Warm Inflation

- couple inflaton to light degrees of freedom $\mathcal{L}_{\text{int}} = -\phi \frac{\alpha}{16\pi f} \tilde{G} G$

 g

$$\ddot{\phi} + 3H\dot{\phi} + V' = - \left\langle \frac{\alpha}{16\pi f} \tilde{G} G \right\rangle$$

$$\left\langle \frac{\alpha}{16\pi f} \tilde{G} G \right\rangle \approx \cancel{m_{th}^2} \phi + \Upsilon \dot{\phi} + O(\ddot{\phi})$$

not allowed
by symmetry

Minimal Warm Inflation

- couple inflaton to light degrees of freedom $\mathcal{L}_{\text{int}} = -\phi \frac{\alpha}{16\pi f} \tilde{G} G$

$\gtrsim g$

$$\ddot{\phi} + 3H\dot{\phi} + V' = - \left\langle \frac{\alpha}{16\pi f} \tilde{G} G \right\rangle$$

enhanced by sphalerons

$$\left\langle \frac{\alpha}{16\pi f} \tilde{G} G \right\rangle \approx \cancel{m_{th}^2 \phi} + \Upsilon \dot{\phi} + O(\ddot{\phi})$$

not allowed
by symmetry

Sphalerons



The Editor
Physical Review D
1 Research Road
Box 1000
Ridge, NY 11961

July 17, 1984

Dear Sir:

Please find enclosed our manuscript "A Saddle-Point Solution in the Weinberg-Salam Theory", which we have revised in response to some of the referee's criticisms. In particular, we have substantially rewritten the Introduction, and added few sentences at the end of Sec. III, in order to clarify that our solution is a saddle-point.

We have continued to use the word "sphaleron" in our paper, for which we think there is a real need. It describes a static unstable solution, localized in space, in any field theory, and not just the solution in the Weinberg-Salam theory which we discuss in this paper. Such solutions are very common in a variety of field theories, and they usually reflect the topological complexity of the theory. We think "sphaleron" is quite a nice word, and we spent some time seriously looking for alternatives, but were less happy with those. Sphaleron has the same root (sphaleros) as sphalerite, which is the name of a fairly common mineral. We noted that the referee had to refer to "this thing" when he didn't use the word sphaleron, which indicates that there is no obvious alternative word.

We would have liked to be more specific about the physical significance of our solution, as the referee requested, but this is a subject of ongoing research by us, and we have nothing more to say in the present paper.

We hope that the paper can be published in its present form.

Yours sincerely,



Frans Klinkhamer

Nicholas Manton

Sphalerons

We have continued to use the word "sphaleron" in our paper, for which we think there is a real need. It describes a static unstable solution, localized in space, in any field theory, and not just the solution in the Weinberg-Salam theory which we discuss in this paper. Such solutions are very common in a variety of field theories, and they usually reflect the topological complexity of the theory. We think "sphaleron" is quite a nice word, and we spent some time seriously looking for alternatives, but were less happy with those. Sphaleron has the same root (sphaleros) as sphalerite which is the name of a fairly common mineral. We noted that the referee had to refer to "this thing" when he didn't use the word sphaleron, which indicates that there is no obvious alternative word.



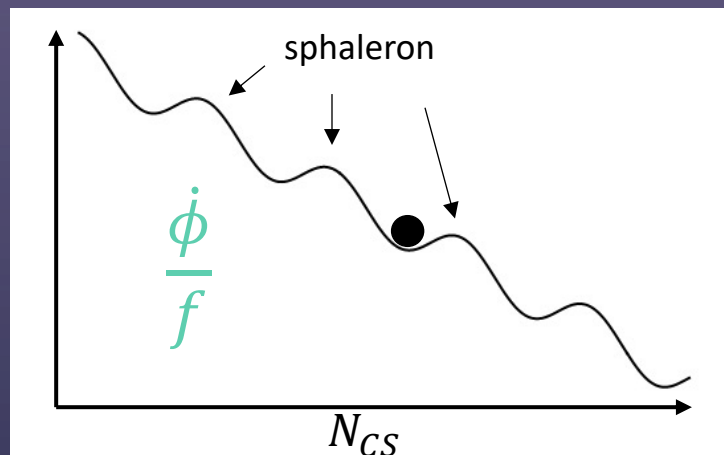
Frans Klinkhamer

Nicholas Manton

Sphaleron friction

$$\mathcal{L}_{\text{int}} = -\phi \frac{\alpha}{8\pi f} \tilde{G} G \approx \dot{\phi}/f N_{CS}$$

Inflaton acts as “chemical potential” for N_{CS}

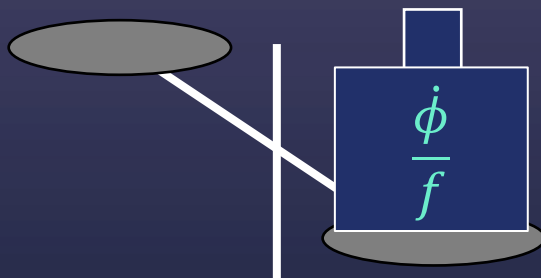
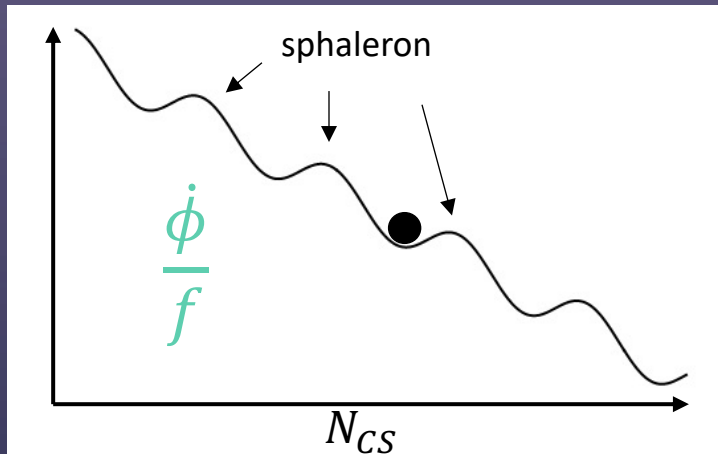


Not drawn to scale

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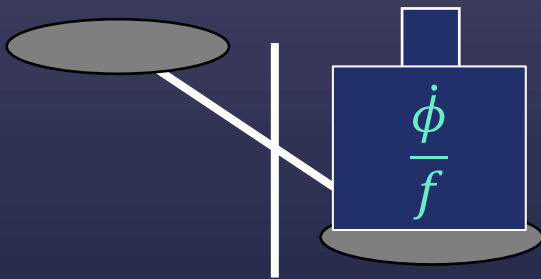
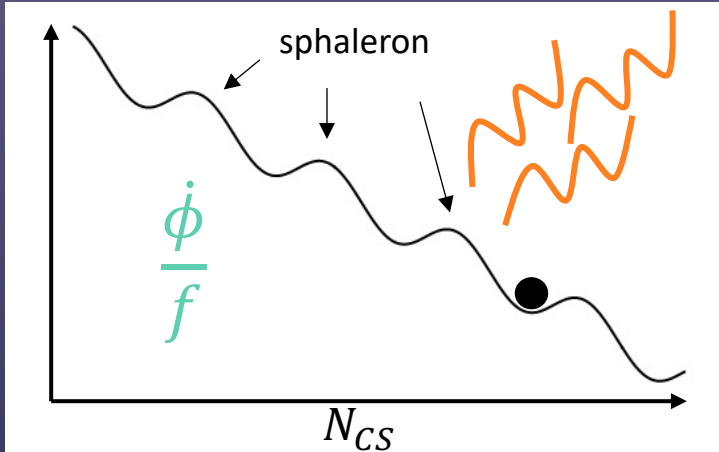
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induces CS-friction

$$\Upsilon \approx N_c^5 \alpha^5 \frac{T^3}{f^2}$$

$$\ddot{\phi} + (3H + \Upsilon)\dot{\phi} + V' = 0$$

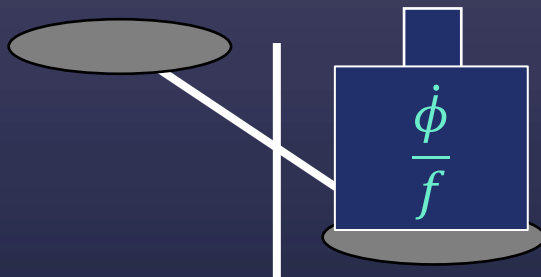
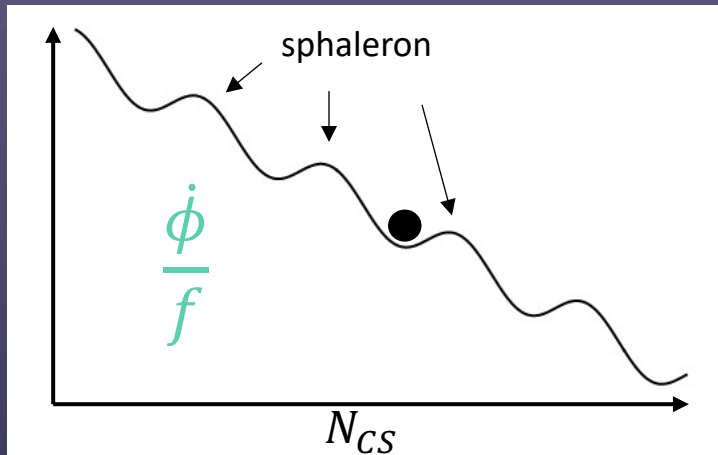
$$\rho + 4H\rho = \Upsilon \dot{\phi}^2$$



Sphaleron friction with the SM

$$\mathcal{L}_{\text{int}} = -\phi \frac{\alpha}{8\pi f} \tilde{G} G \approx \frac{\dot{\phi}}{f} N_{CS}$$

QCD $\gtrsim g, q$



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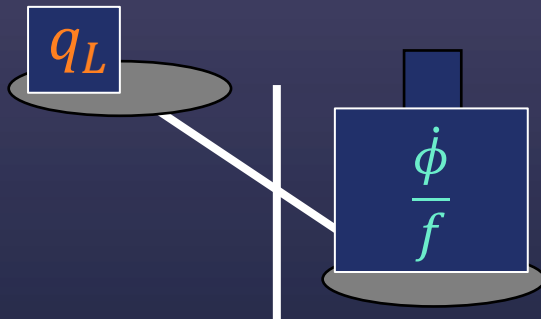
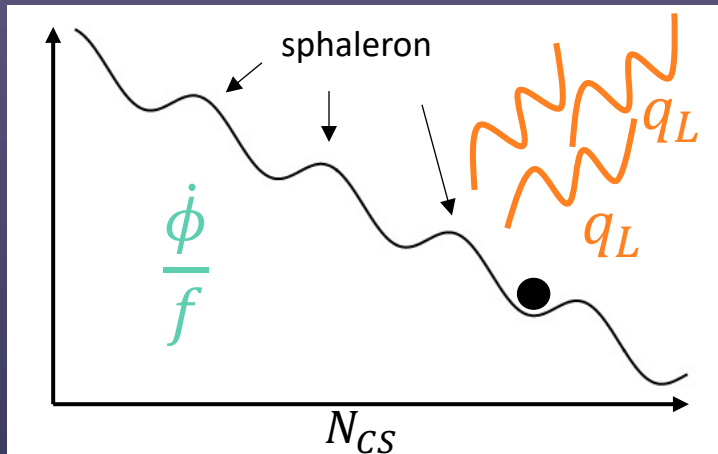
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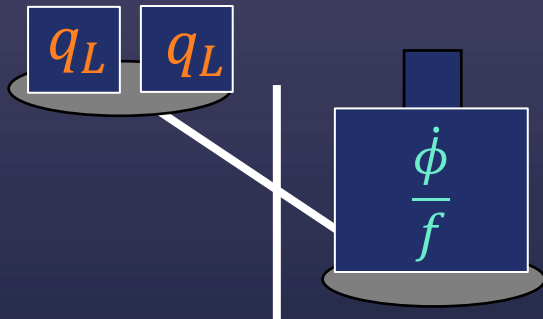
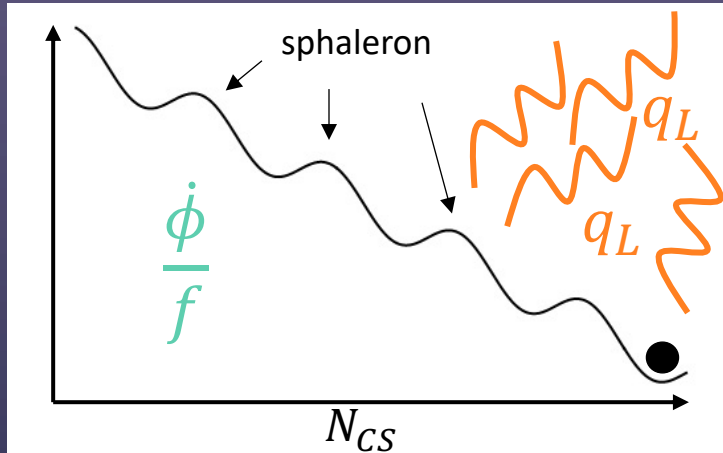
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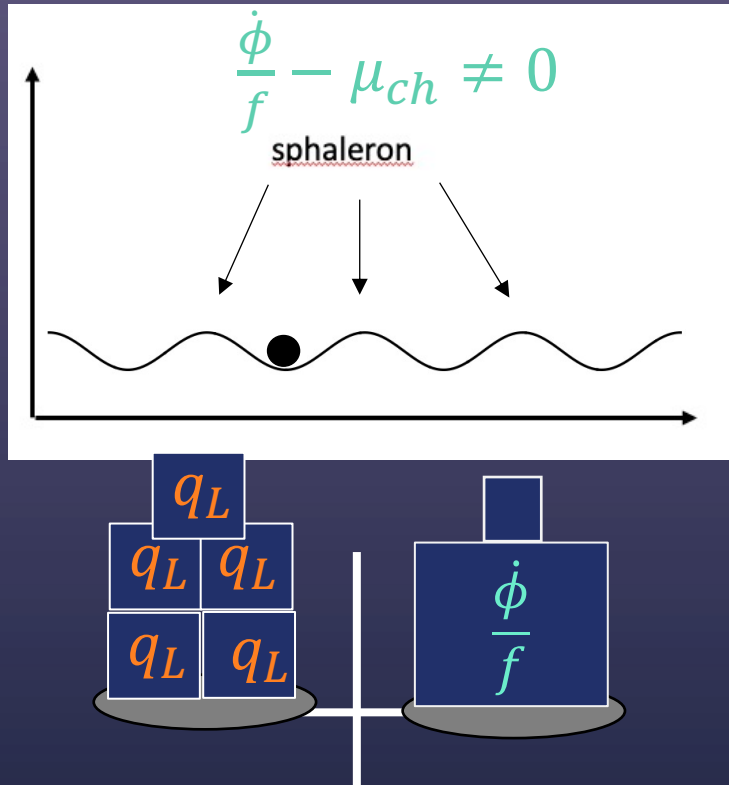
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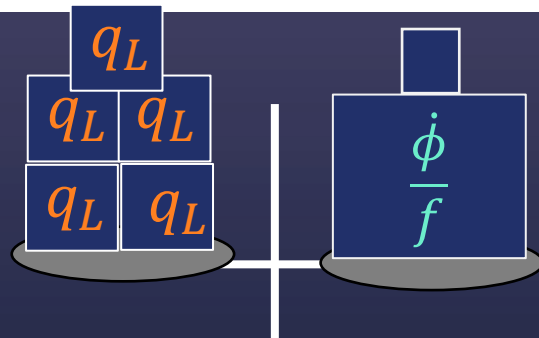
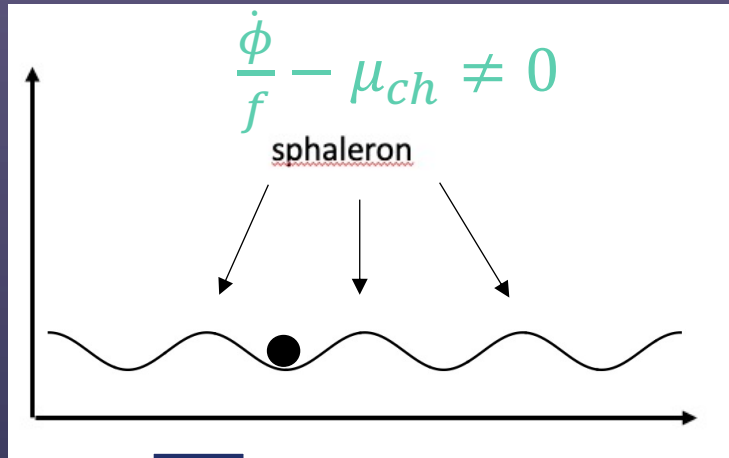
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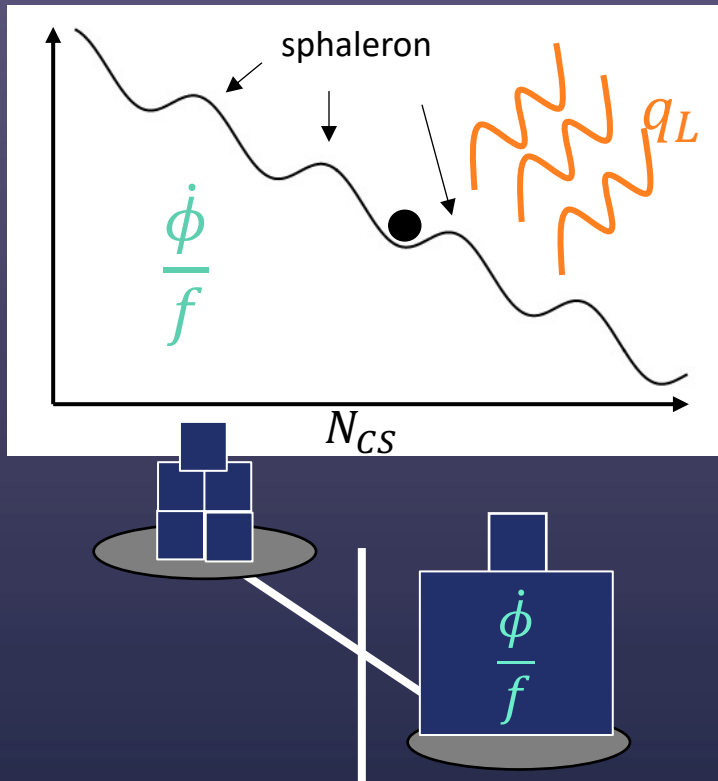
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Berghaus, Drewes, Zell, 2025

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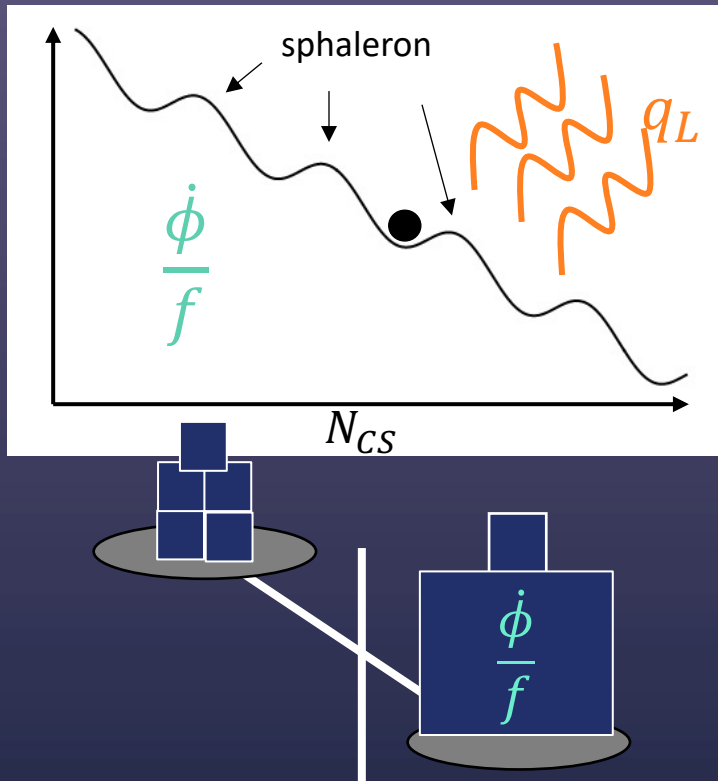
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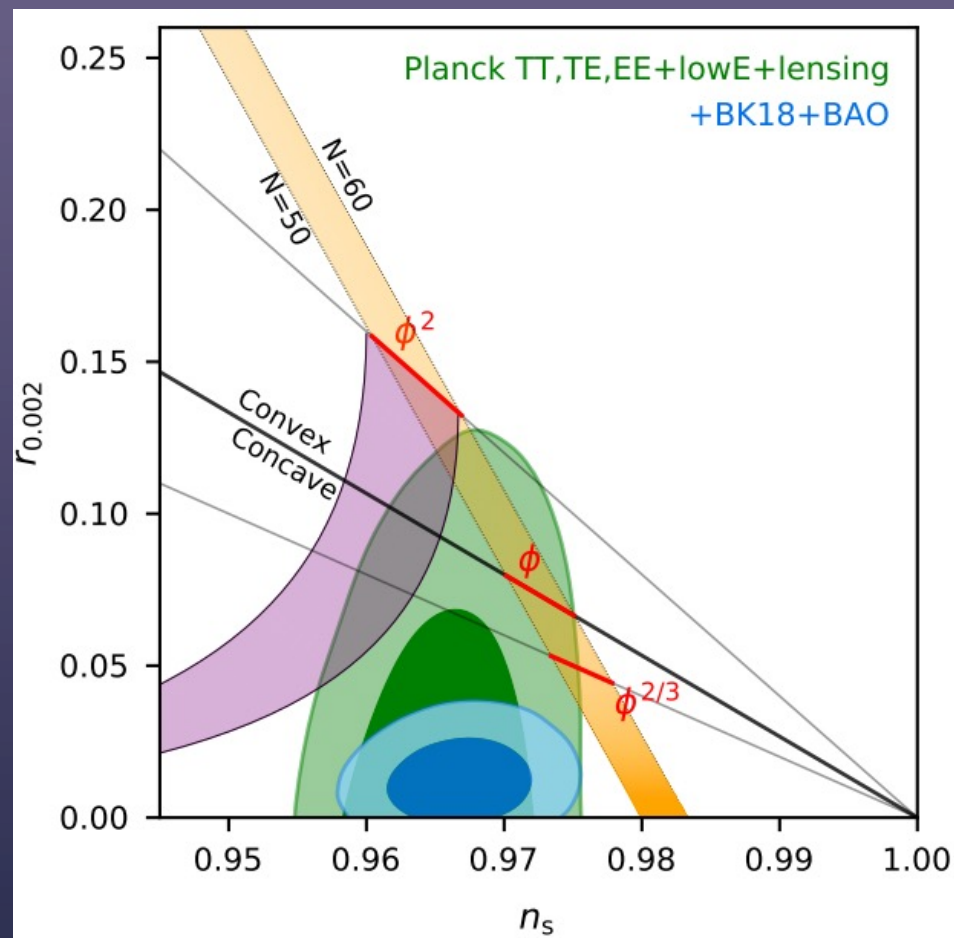
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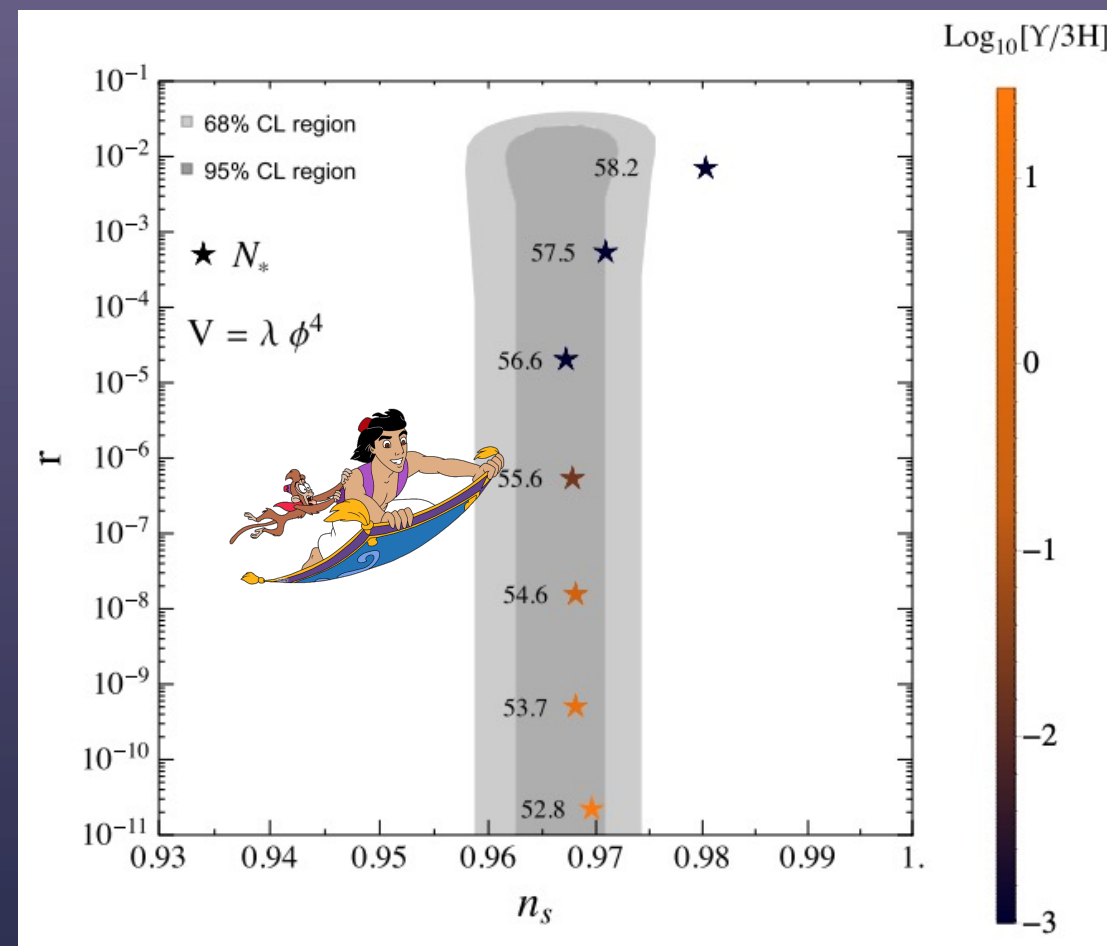
$$\Upsilon_{eff} = \Upsilon / \left(1 + \frac{N_f}{N_c} \frac{12f^2}{T^2} \frac{\Upsilon}{3H} \right)$$

Berghaus, Drewes, Zell, 2025

Warm Inflation with the Standard Model



Bicep/Keck XIII, 2021

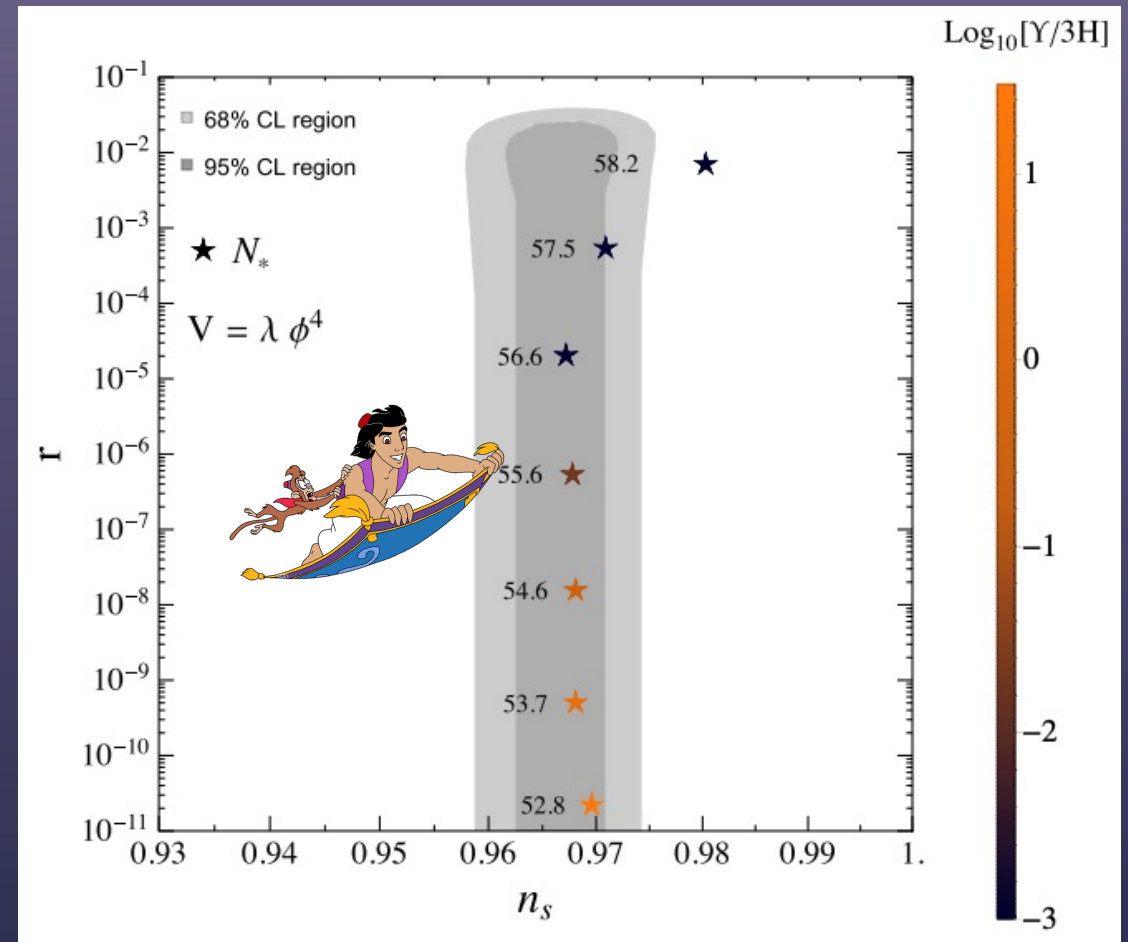


Berghaus, Drewes, Zell, 2025

Warm Inflation with the Standard Model

$$V = \lambda \phi^4$$

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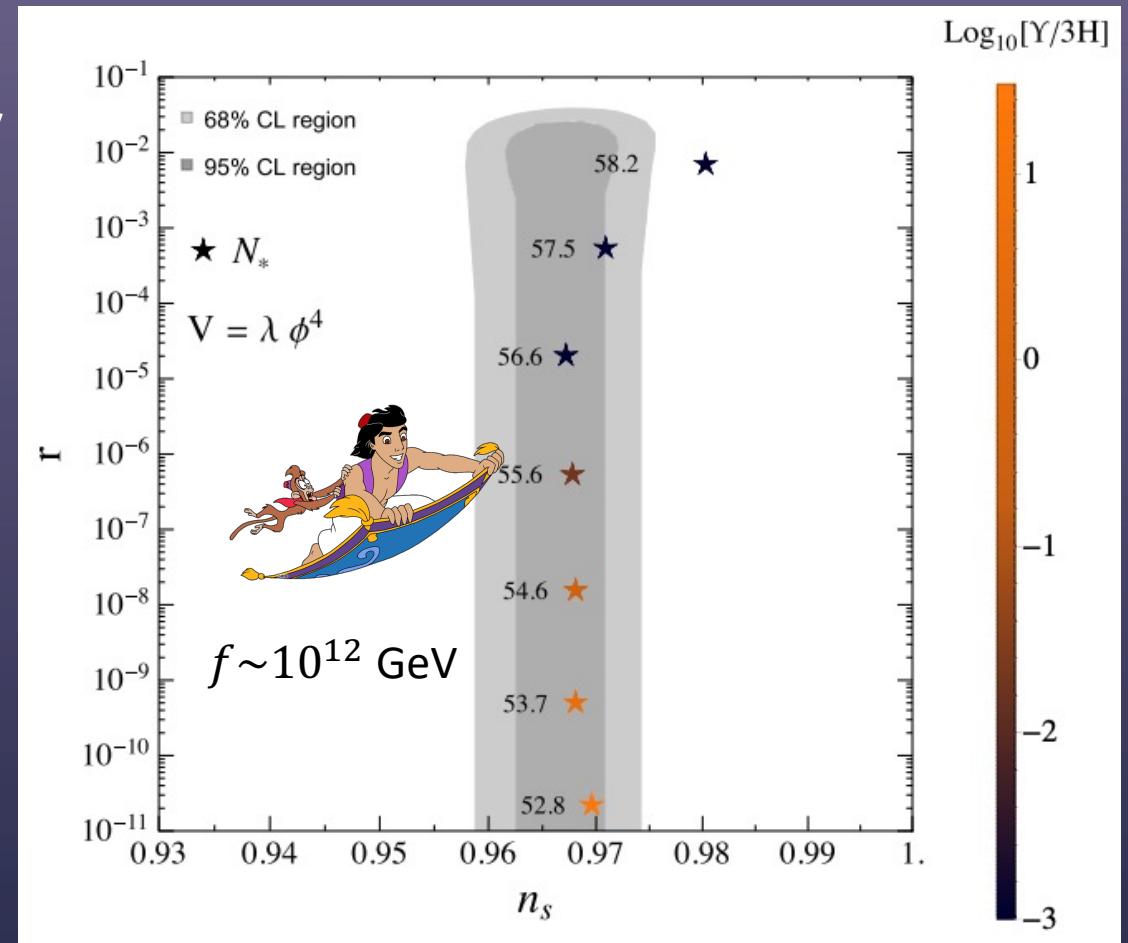
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accessible by
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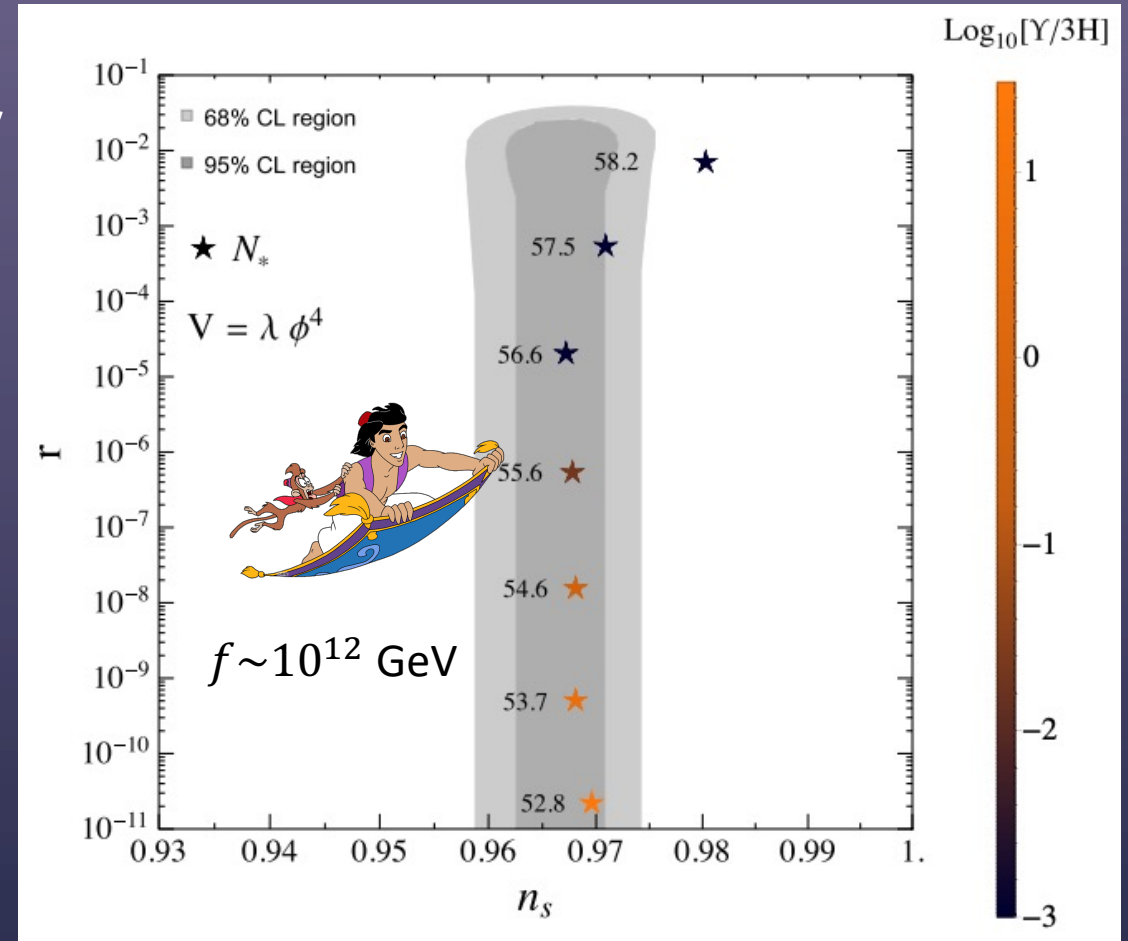
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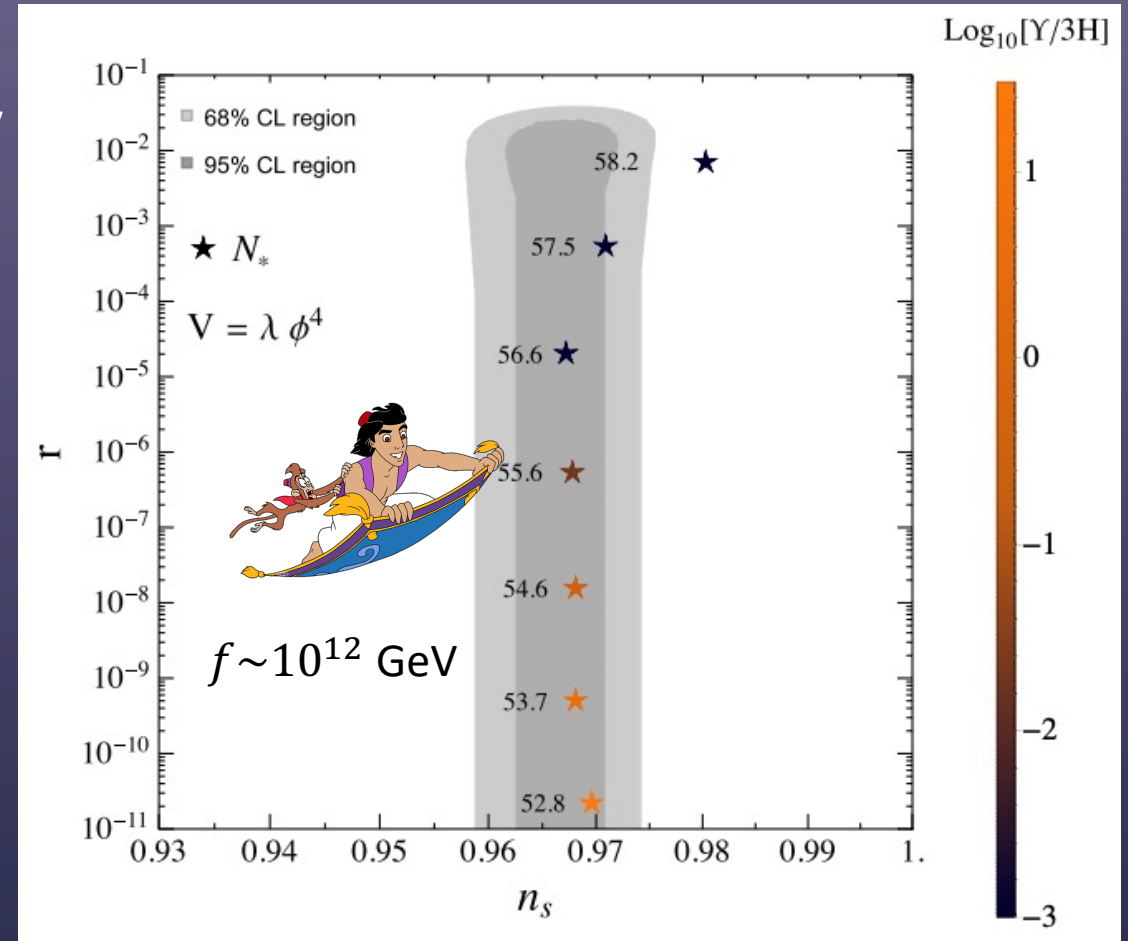
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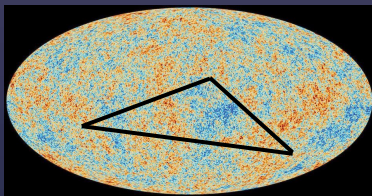
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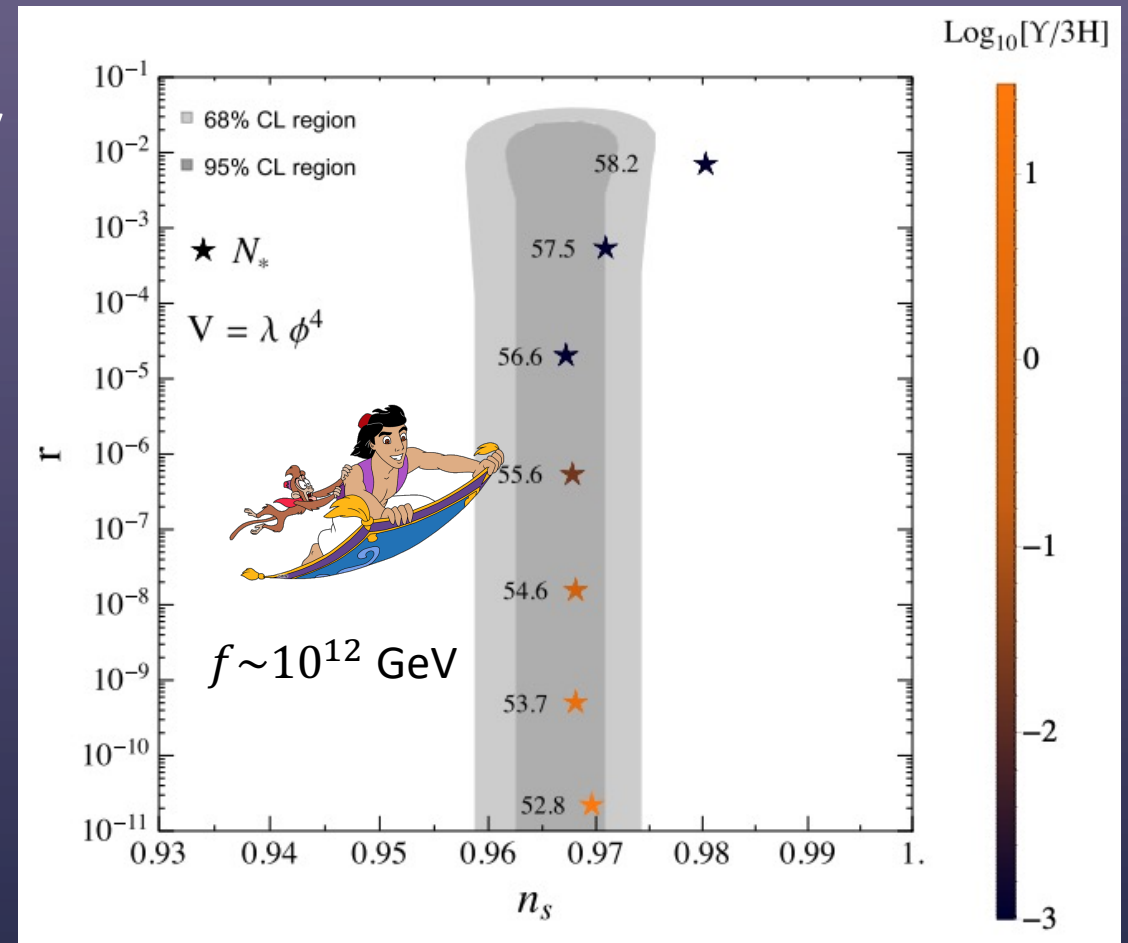
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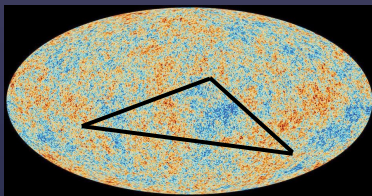
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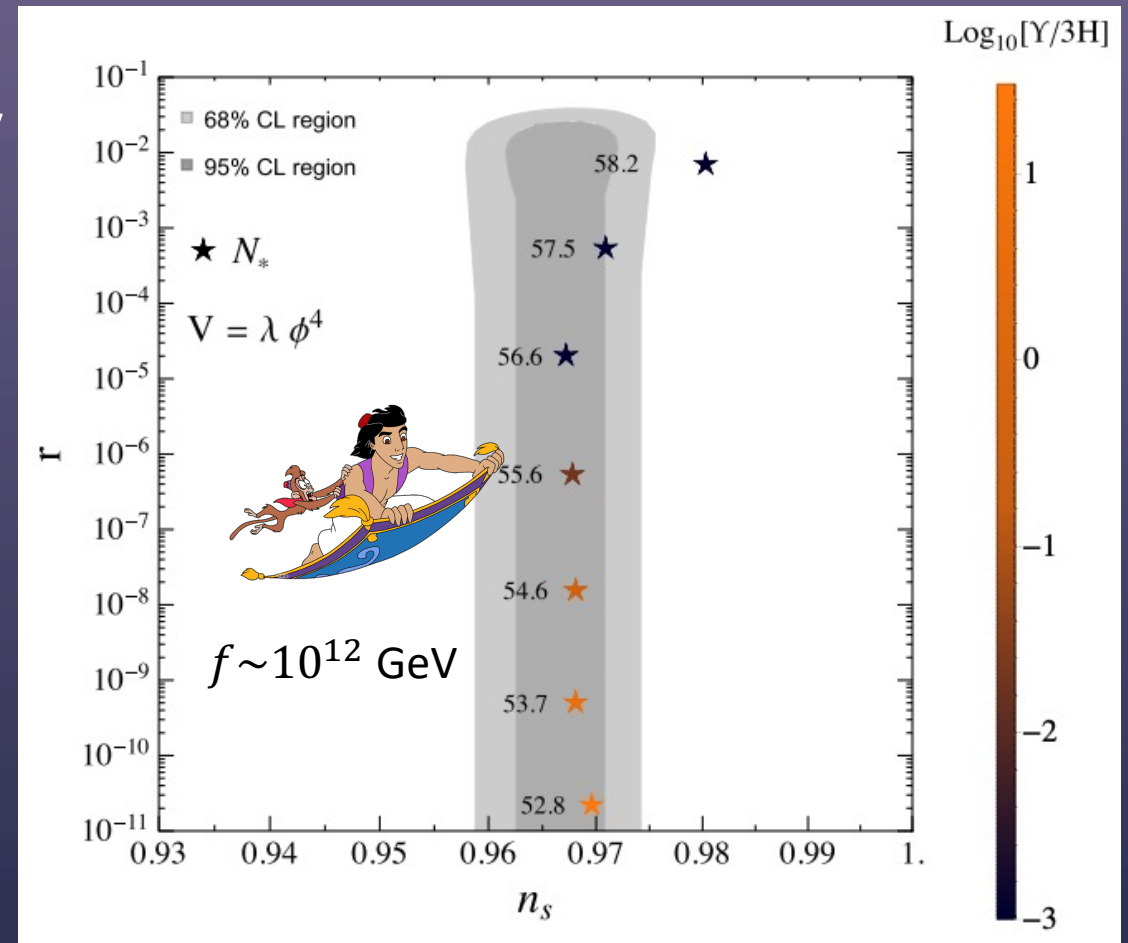
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What are the laws of nature

heating the early Universe

the strong CP problem

the nature of the inflaton

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Thank you!