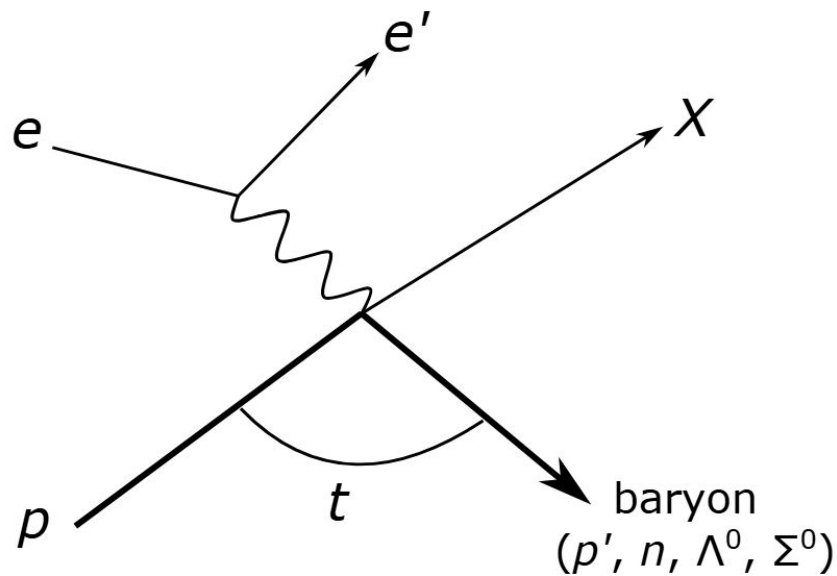


The t RECO Convention



- p_{BA}^μ : outgoing baryon 4-momentum.
- p_{BE}^μ : incoming proton beam momentum.
- $p_{\gamma^*}^\mu$: difference between incoming and scattered electron 4-momentum.
- p_X^μ : 4-momentum of the rest of final state.

Why is t reconstruction not trivial?

Overconstrained system, with momentum conservation, t can be measured in many different ways. Ideal way to compute t might depend on particular channel and particular kinematics

Important factors to consider:

- Beam divergence of either beam
- QED radiation. Initial and Final state radiation
- Detector resolution (in both transverse momentum, and longitudinal momentum)
- Subtraction of large numbers, etc.

Quite a bit of
proliferation of
methods in
inclusive DIS in
HERA days...

No naming
convention!!!

Method name	Observables	y	Q^2	$x \cdot E_p$
Electron (e)	$[E_0, E, \theta]$	$1 - \frac{\Sigma_e}{2E_0}$	$\frac{E^2 \sin^2 \theta}{1-y}$	$\frac{E(1+\cos \theta)}{2y}$
Double angle (DA) [6, 7]	$[E_0, \theta, \gamma]$	$\frac{\tan \frac{\gamma}{2}}{\tan \frac{\gamma}{2} + \tan \frac{\theta}{2}}$	$4E_0^2 \cot^2 \frac{\theta}{2} (1-y)$	$\frac{Q^2}{4E_0 y}$
Hadron (h , JB) [4]	$[E_0, \Sigma, \gamma]$	$\frac{\Sigma}{2E_0}$	$\frac{T^2}{1-y}$	$\frac{Q^2}{2\Sigma}$
ISigma ($\text{I}\Sigma$) [9]	$[E, \theta, \Sigma]$	$\frac{\Sigma}{\Sigma + \Sigma_e}$	$\frac{E^2 \sin^2 \theta}{1-y}$	$\frac{E(1+\cos \theta)}{2y}$
IDA [7]	$[E, \theta, \gamma]$	y_{DA}	$\frac{E^2 \sin^2 \theta}{1-y}$	$\frac{E(1+\cos \theta)}{2y}$
$E_0 E \Sigma$	$[E_0, E, \Sigma]$	y_h	$4E_0 E - 4E_0^2 (1-y)$	$\frac{Q^2}{2\Sigma}$
$E_0 \theta \Sigma$	$[E_0, \theta, \Sigma]$	y_h	$4E_0^2 \cot^2 \frac{\theta}{2} (1-y)$	$\frac{Q^2}{2\Sigma}$
$\theta \Sigma \gamma$ [8]	$[\theta, \Sigma, \gamma]$	y_{DA}	$\frac{T^2}{1-y}$	$\frac{Q^2}{2\Sigma}$
Double energy (A4) [7]	$[E_0, E, E_h]$	$\frac{E-E_0}{(xE_p)-E_0}$	$4E_0 y (xE_p)$	$E + E_h - E_0$
$E \Sigma T$	$[E, \Sigma, T]$	$\frac{\Sigma}{\Sigma + E \pm \sqrt{E^2 + T^2}}$	$\frac{T^2}{1-y}$	$\frac{Q^2}{2\Sigma}$
$E_0 E T$	$[E_0, E, T]$	$\frac{2E_0 - E \mp \sqrt{E^2 - T^2}}{2E_0}$	$\frac{T^2}{1-y}$	$\frac{Q^2}{4E_0 y}$
Sigma (Σ) [9]	$[E_0, E, \Sigma, \theta]$	$y_{\text{I}\Sigma}$	$Q_{\text{I}\Sigma}^2$	$\frac{Q^2}{4E_0 y}$
$e\Sigma$ ($e\Sigma$) [9]	$[E_0, E, \Sigma, \theta]$	$\frac{2E_0 \Sigma}{(\Sigma + \Sigma_e)^2}$	$2E_0 E (1 + \cos \theta)$	$\frac{E(1+\cos \theta)(\Sigma + \Sigma_e)}{2\Sigma}$

Table 1. Summary of basic reconstruction methods that employ only three out of five quantities: E_0 (electron-beam energy), E and θ (scattered electron energy and polar angle), Σ and γ (longitudinal energy-momentum balance, $\Sigma = \sum_{\text{HFS}} (E_i - p_{z,i})$, and the inclusive angle of the HFS).

The t -Reconstruction (t RECO) Convention establishes a standardized nomenclature for methods used to reconstruct the Mandelstam variable t in electron-proton (or ion) scattering. This convention aims to provide a unified framework for describing a broad class of experimental techniques to be used for t reconstruction at the Electron-Ion Collider.

TABLE I. The t RECO convention: classes of methods used to reconstruct the Mandelstam variable t are defined by the observables that are at least partially utilized. Here, p_{BA}^μ represents the outgoing baryon 4-momentum; p_{BE}^μ denotes the incoming proton beam momentum; $p_{\gamma^*}^\mu$ refers to the difference between the incoming and scattered electron 4-momenta; and p_{X}^μ represents the 4-momentum of the rest of the final

Class name	Observables used	Example(s)
BABE	$p_{\text{BA}}^\mu, p_{\text{BE}}^\mu$	$-t = (p_{\text{BA}}^\mu - p_{\text{BE}}^\mu)^2$ or $t = \vec{p}_{\text{BA}}^{\text{T}} ^2$
eX	$p_{\gamma^*}^\mu, p_{\text{X}}^\mu$	$-t = (p_{\gamma^*}^\mu - p_{\text{X}}^\mu)^2$ or $t = \vec{p}_{\text{X}}^{\text{T}} + \vec{p}_{e'}^{\text{T}} ^2$
eXBA	$p_{\gamma^*}^\mu, p_{\text{X}}^\mu, p_{\text{BA}}^\mu$	$-t = (p_{\text{corr}}^\mu - p_{\text{X}}^\mu)^2$ $p_{\text{corr}}^\mu = (\frac{\Sigma_{\text{XBA}}}{2} - \frac{E_{e'}}{2}(1 + \cos \theta_{e'}), p_{\gamma^*}^\mu[1], p_{\gamma^*}^\mu[2], -\frac{\Sigma_{\text{XBA}}}{2} - \frac{E_{e'}}{2}(1 + \cos \theta_{e'}))$ $\Sigma_{\text{XBA}} := (E_{\text{X}} - p_{\text{X}}^z) + (E_{\text{BA}} - p_{\text{BA}}^z)$
eXBE	$p_{\gamma^*}^\mu, p_{\text{X}}^\mu, p_{\text{BE}}^\mu$	$-t = (p_{\text{corr}}^\mu - p_{\text{BE}}^\mu)^2$ or $t = \vec{p}_{\text{miss}}^{\text{T}} ^2$ $p_{\text{corr}}^\mu = \left[\sqrt{ \vec{p}_{\text{miss}} ^2 + m_{\text{BA}}^2}, \vec{p}_{\text{miss}} \hat{n}(\theta_{\text{miss}}, \phi_{\text{miss}}) \right]$ $p_{\text{miss}}^\mu = p_{\gamma^*}^\mu + p_{\text{BE}}^\mu - p_{\text{X}}^\mu = [E_{\text{miss}}, \vec{p}_{\text{miss}} \hat{n}(\theta_{\text{miss}}, \phi_{\text{miss}})]$
eBABE	$p_{\gamma^*}^\mu, p_{\text{BA}}^\mu, p_{\text{BE}}^\mu$	$-t = (p_{\text{corr}}^\mu - p_{\text{BE}}^\mu)^2$ $p_{\text{corr}}^\mu = (p_{\text{BA}}^\mu[0], -p_{\gamma^*}^\mu[1], -p_{\gamma^*}^\mu[2], p_{\text{BA}}^\mu[3])$
XBABE	$p_{\text{X}}^\mu, p_{\text{BA}}^\mu, p_{\text{BE}}^\mu$	$-t = (p_{\text{corr}}^\mu - p_{\text{BE}}^\mu)^2$ $p_{\text{corr}}^\mu = (p_{\text{BA}}^\mu[0], -p_{\text{X}}^\mu[1], -p_{\text{X}}^\mu[2], p_{\text{BA}}^\mu[3])$
eXBABE	$p_{\gamma^*}^\mu, p_{\text{X}}^\mu, p_{\text{BA}}^\mu, p_{\text{BE}}^\mu$	$-t = (p_{\text{corr}}^\mu - p_{\text{BE}}^\mu)^2$ $p_{\text{corr}}^\mu = \left[\sqrt{ \vec{p}_{\text{miss}} ^2 + m_{\text{BA}}^2}, \vec{p}_{\text{miss}} \hat{n}(\theta_{\text{BA}}, \phi_{\text{BA}}) \right]$ $p_{\text{miss}}^\mu = p_{\gamma^*}^\mu + p_{\text{BE}}^\mu - p_{\text{X}}^\mu = [E_{\text{miss}}, \vec{p}_{\text{miss}} \hat{n}(\theta_{\text{miss}}, \phi_{\text{miss}})]$

The *t*RECO Convention

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You are welcome to comment, edit, and sign the convention