

# Flow in small systems at the LHC

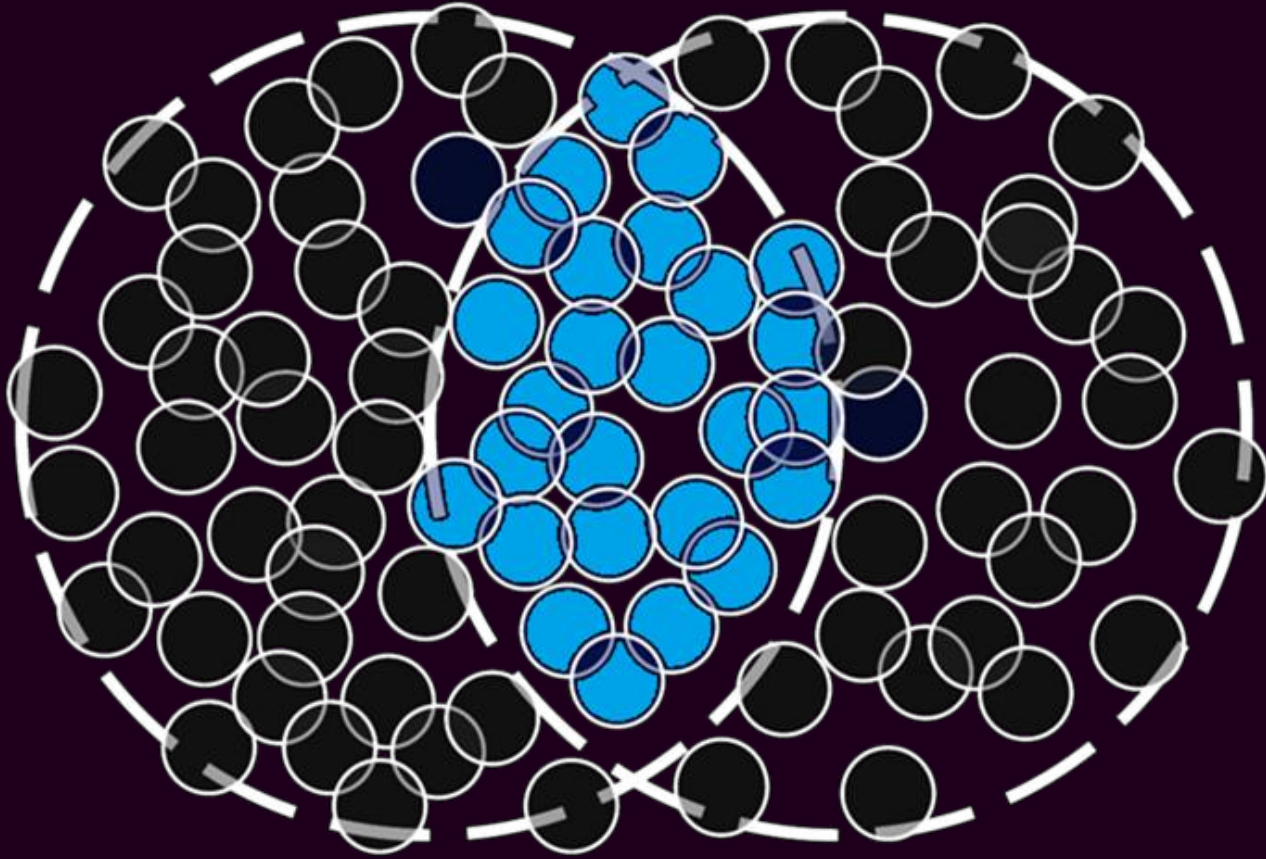
Blair Daniel Seidlitz

Columbia University



RHIC/AGS Users Meeting May. 21<sup>st</sup> , 2025

# Initial state of heavy ion collisions

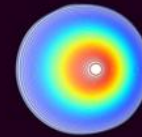


Parametrize spatial anisotropy

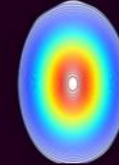
Geometric eccentricity

$$\varepsilon_n = \frac{\sqrt{\langle r_i^n \cos(n\phi_i) \rangle^2 + \langle r_i^n \sin(n\phi_i) \rangle^2}}{\langle r_i^n \rangle}$$

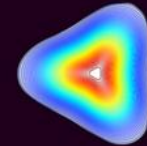
n=1



n=2



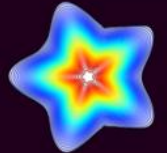
n=3

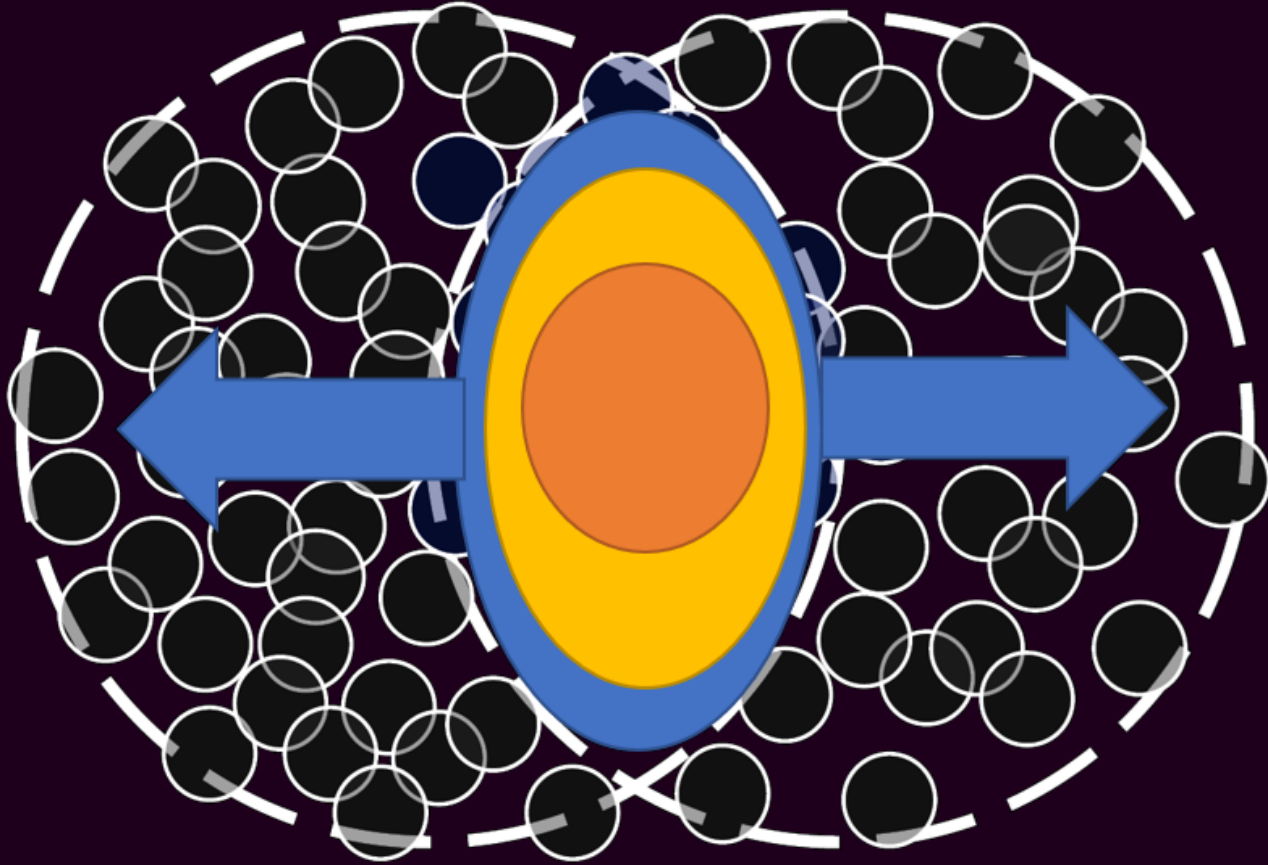


n=4



n=5





Parametrize spatial anisotropy

Geometric eccentricity

$$\varepsilon_n = \frac{\sqrt{\langle r_i^n \cos(n\phi_i) \rangle^2 + \langle r_i^n \sin(n\phi_i) \rangle^2}}{\langle r_i^n \rangle}$$

Viscous Hydrodynamics

$$T^{\mu\nu} = \epsilon u^\mu u^\nu + P[\epsilon] \Delta^{\mu\nu} - \eta[\epsilon] \sigma^{\mu\nu} - \zeta[\epsilon] \Delta^{\mu\nu} \nabla_\lambda^\perp u^\lambda$$

Equation of state  
transport coefficients  $P[\epsilon]$   $\eta[\epsilon]$   $\zeta[\epsilon]$ .

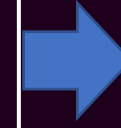
Initial state



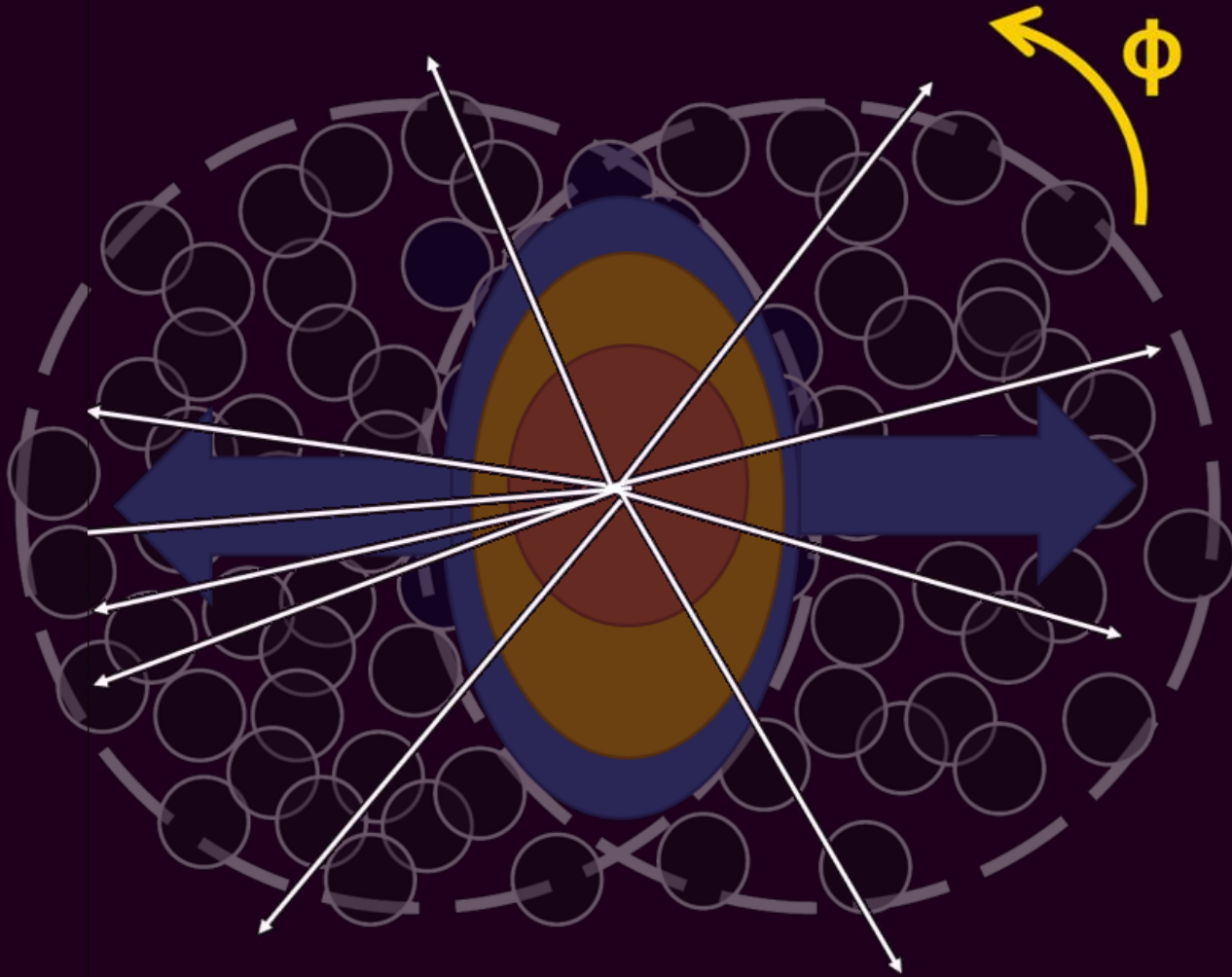
Pre-equilibrium  
dynamics



Hydrodynamic  
expansion



Momentum  
anisotropy



Parametrize spatial anisotropy

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Viscous Hydrodynamics

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Equation of state  
transport coefficients  $P[\epsilon]$   $\eta[\epsilon]$   $\zeta[\epsilon]$ .

Azimuthal Momentum anisotropy

$$\frac{dN}{d\phi} \propto 1 + 2v_1 \cos(\phi) + 2v_2 \cos(2\phi) + 2v_3 \cos(3\phi) + \dots$$

Initial state

$\varepsilon_n$

Pre-equilibrium  
dynamics

Hydrodynamic  
expansion

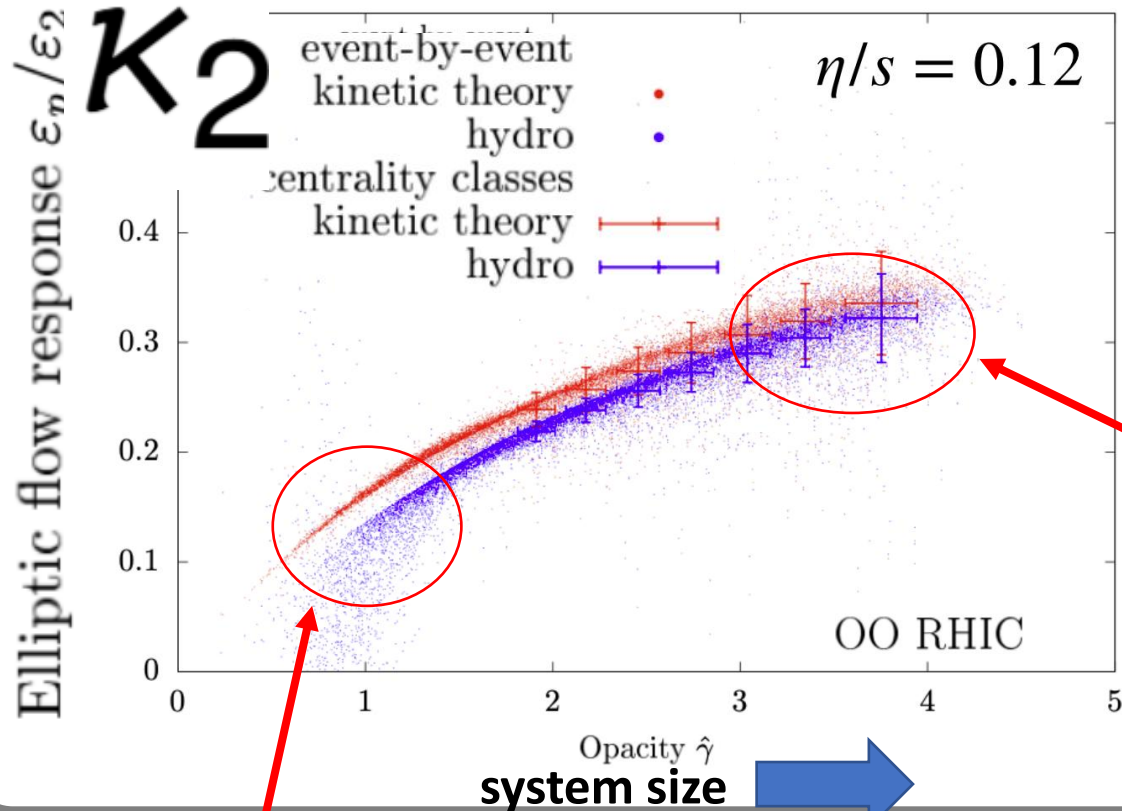
Momentum  
anisotropy

$v_n$

$$v_n \simeq \kappa_n \varepsilon_n$$

# Break down of hydrodynamics in small systems

*Schlichting, Werthmann*



Azimuthal anisotropy is a response to geometry

$$V_n \simeq K_n \varepsilon_n$$

Event by event calculation of kappa for **hydrodynamics** and **kinetic theory**

The two models converge at large system size

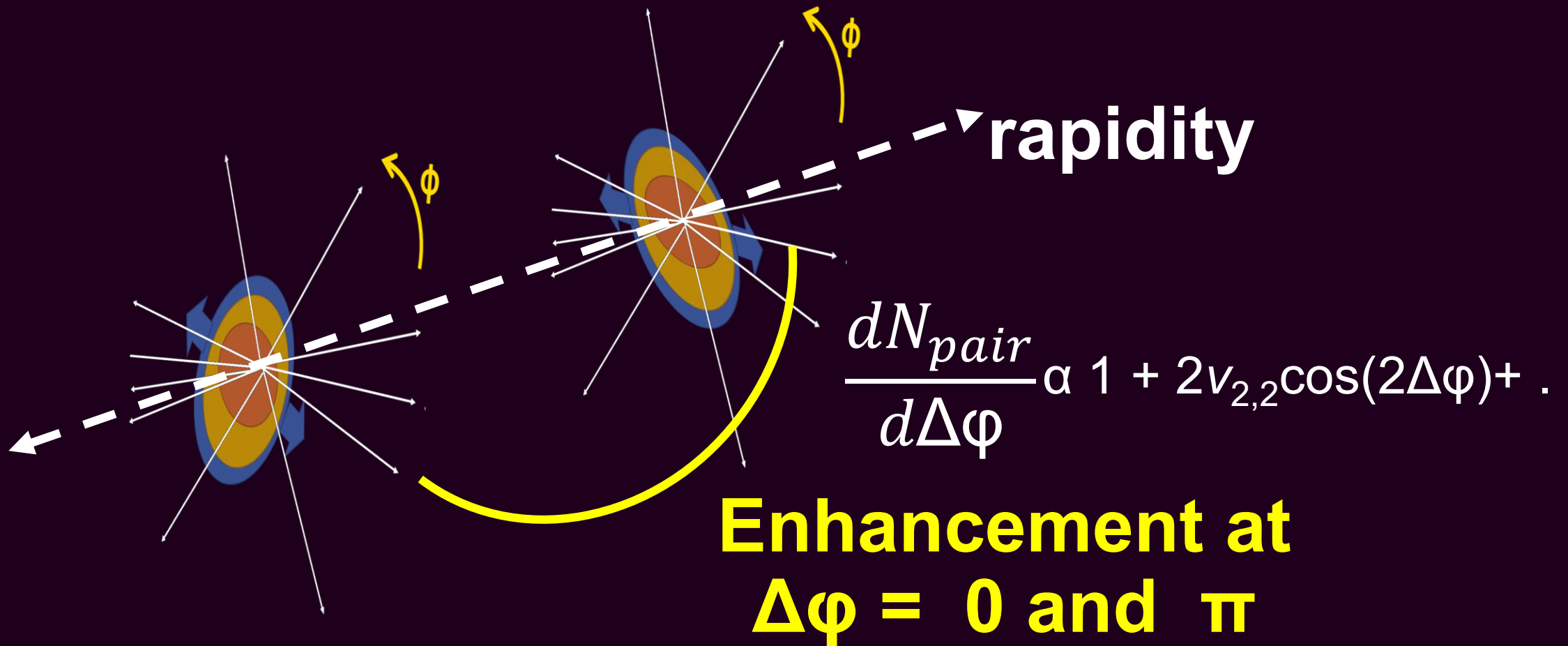
Break down of hydro at small system size

Explore HOT QCD

Can we gain experimental sensitivity to far-from-equilibrium dynamics

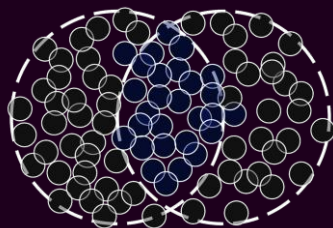


# Two-particle correlations

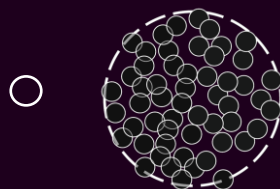


# System size

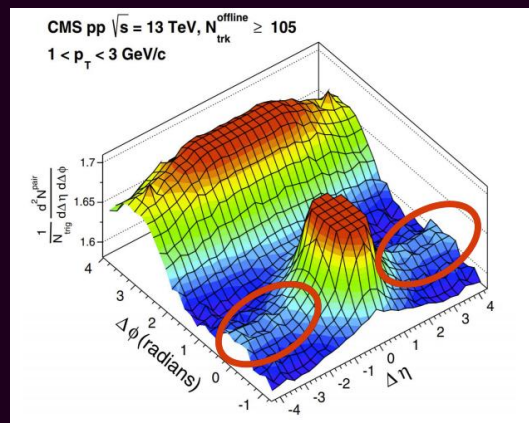
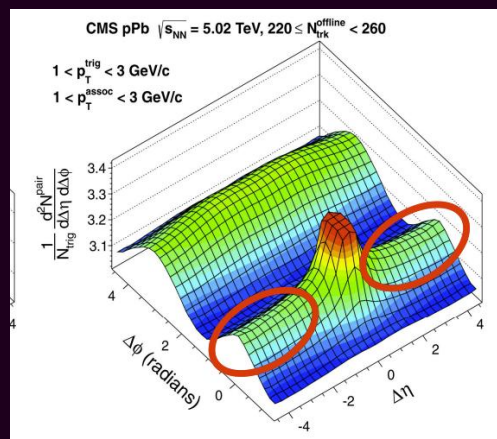
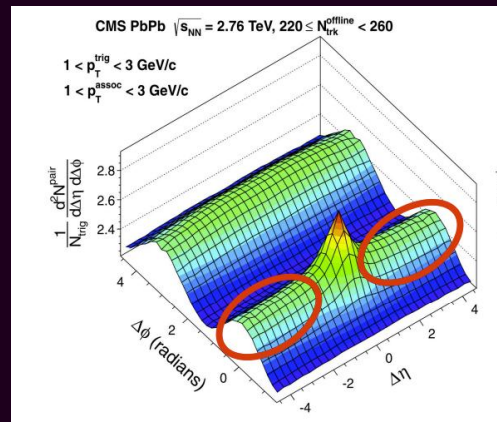
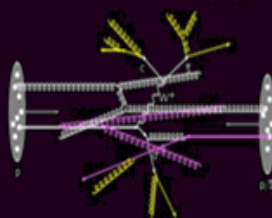
Pb+Pb



p+Pb



pp

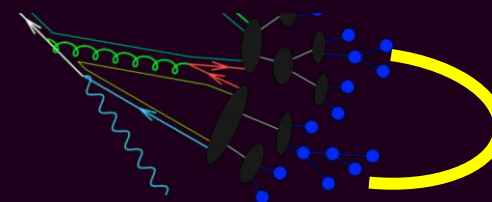
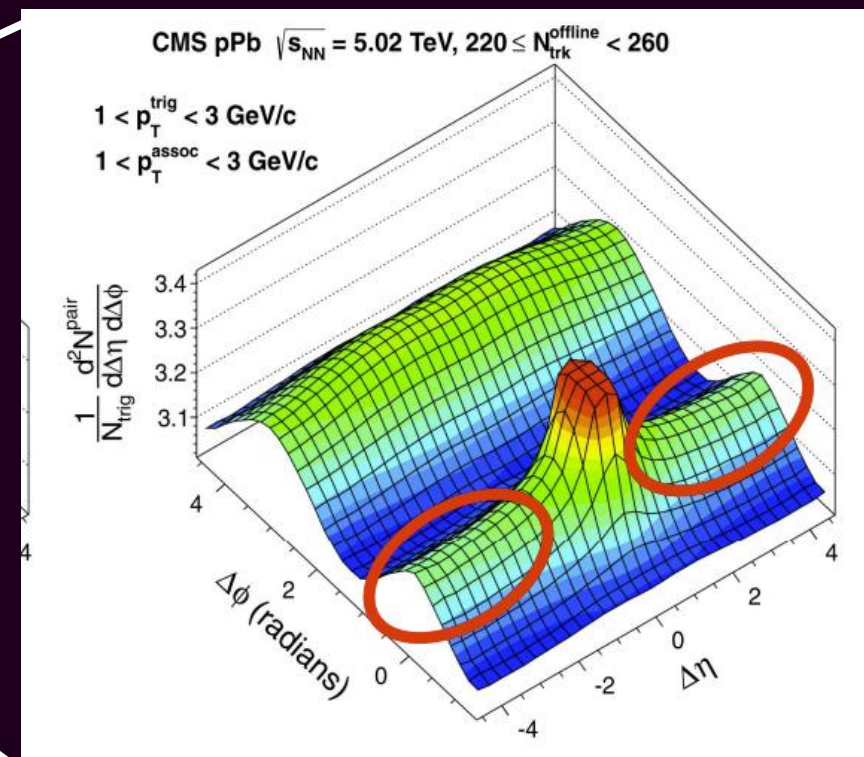
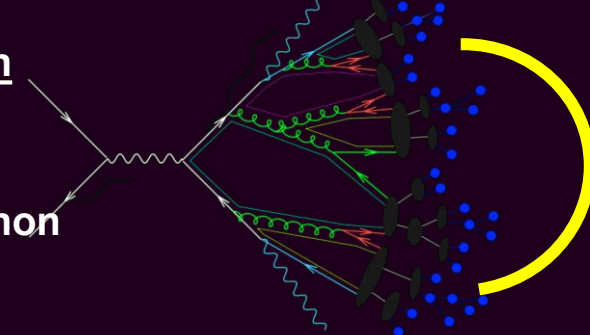


## Momentum conservation

Jets & particle decays

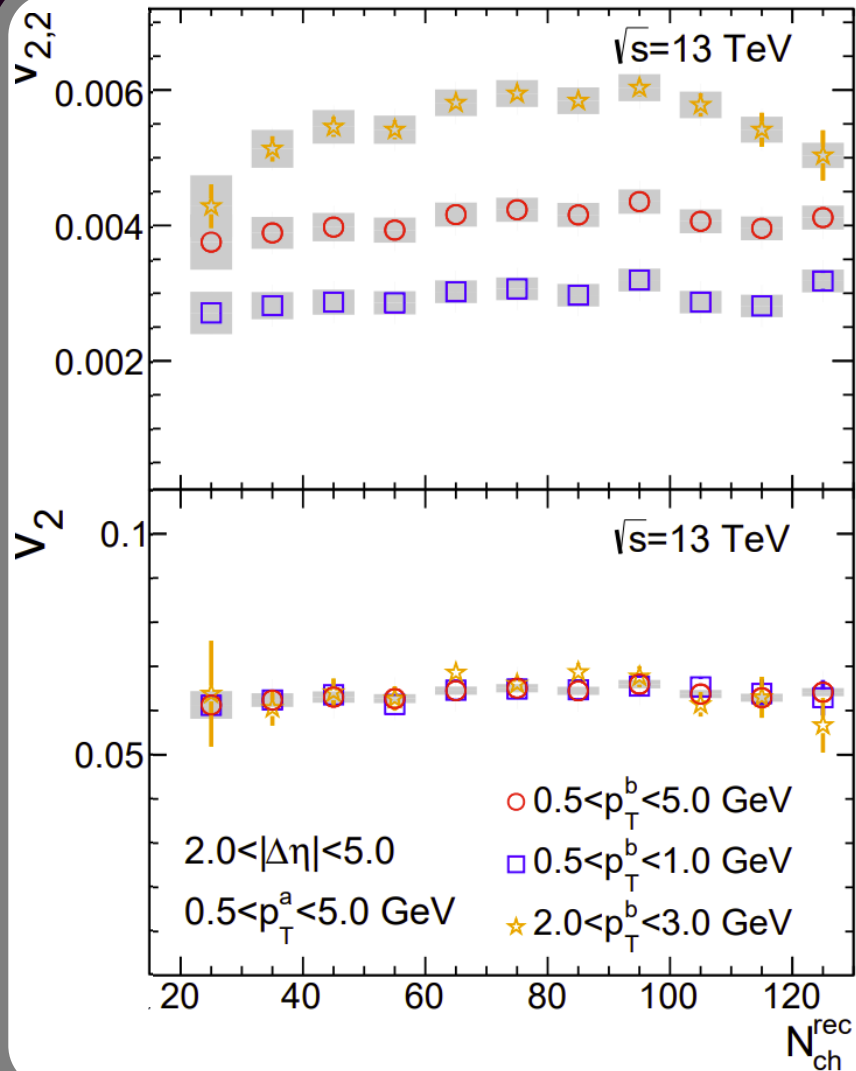
Termed “nonflow”

Not collective phenomenon



$\Delta\phi < 2$   
 $\Delta\eta < 2$

# First signs of collectivity in small systems



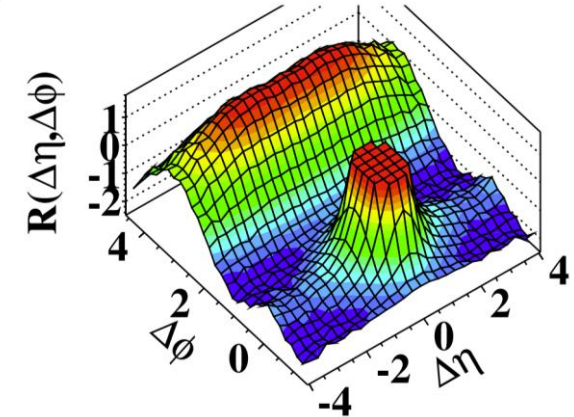
[arXiv:1509.04776](#)

## *CMS* ridge observation

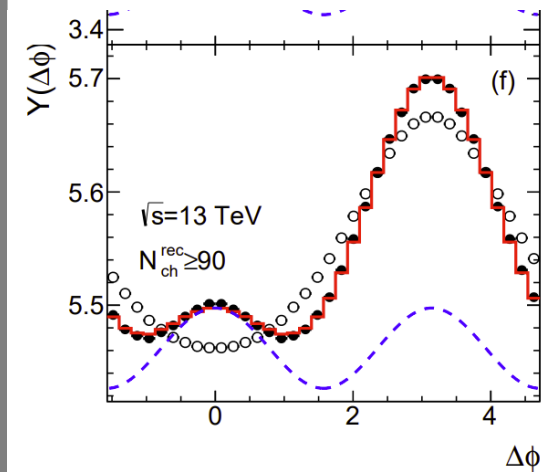
- First small systems ridge by
- Leveraged LHC's experiments,
  - high performance software trigger
  - larger  $N_{ch}$  reach

## *ATLAS* measurement

- Non-flow subtraction – first  $v_n$
- Demonstrated factorization
  - Consistent with anisotropy generated from universal (in  $p_T$ ) hydro flow vector
  - Inconsistent with CGC which has large decorrelation in  $p_T$



[arXiv:1009.4122](#)



[arXiv:1509.04776](#)



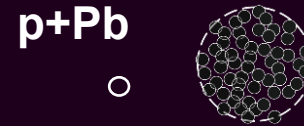
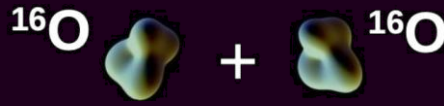
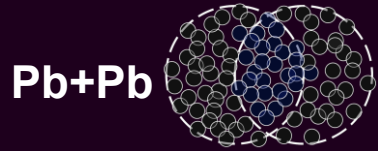
# Essential questions

- Can we break hydro by pushing into a regime governed, by what, far-out-of-equilibrium hydro, kinetic theory, parton escape, etc?
- How well do we know the initial-state geometry?
- Can we remove nonflow effects in a systematic way?
- What is the role of fluctuations?
- Are there other precision tools?

# What is the approach

An earnest precision-oriented march to lower and lower multiplicities and system size.

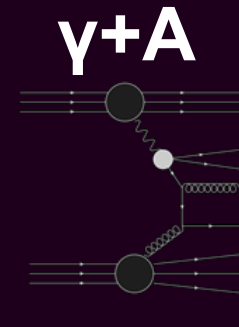
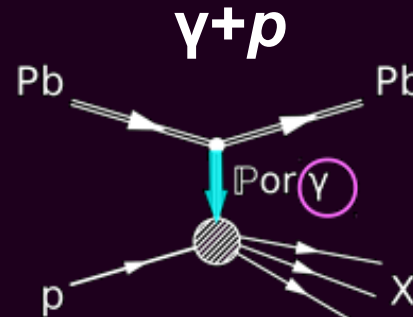
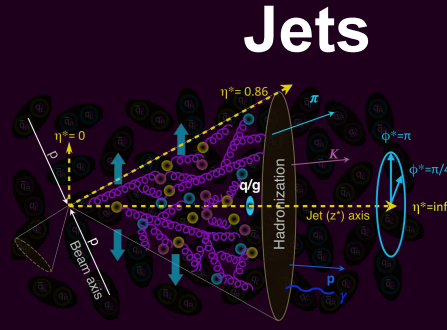
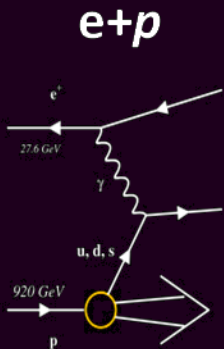
Too big?



Too small?

OR

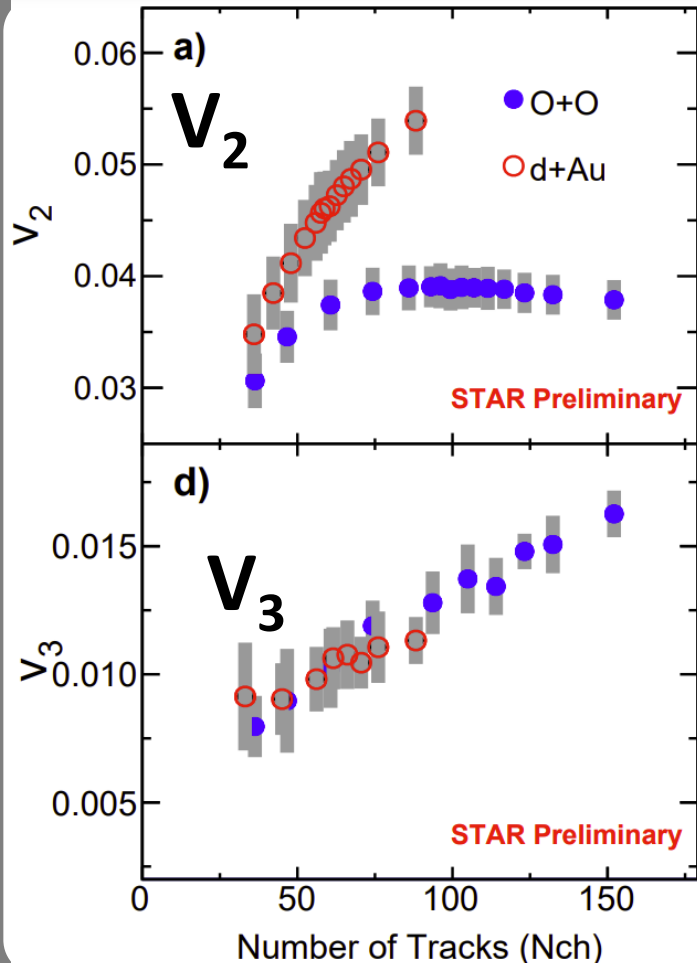
A creative explosion which interrogates the assumed building blocks of collectivity by completely removing them.



**Geometry**  
Fluctuations  
Transport  
Novel small systems  
Future

# $^{16}\text{O}+^{16}\text{O}$ with STAR

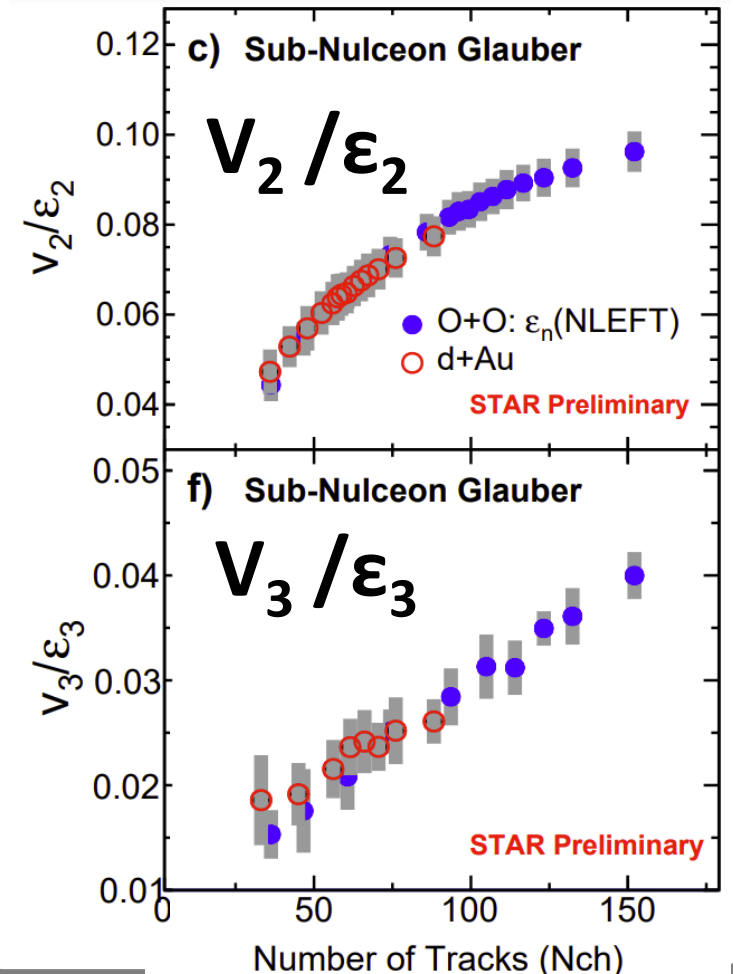
Zhengxi Yan



- $v_2$  and  $v_3$  in O+O and d+Au with nonflow subtraction
- 20-30% difference in  $v_2$  between O+O and d+Au

Divide by geometry

- Agreement of  $v_n/\epsilon_n$
- O+O  $\epsilon_2$  is less than a few percent sensitive to subnucleonic effects
- Larger systems transfer initial-state geometry more effectively



$v_n$  of equally sized systems are driven by geometry



Geometry  
**Fluctuations**  
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# Longitudinal decorrelation in small systems

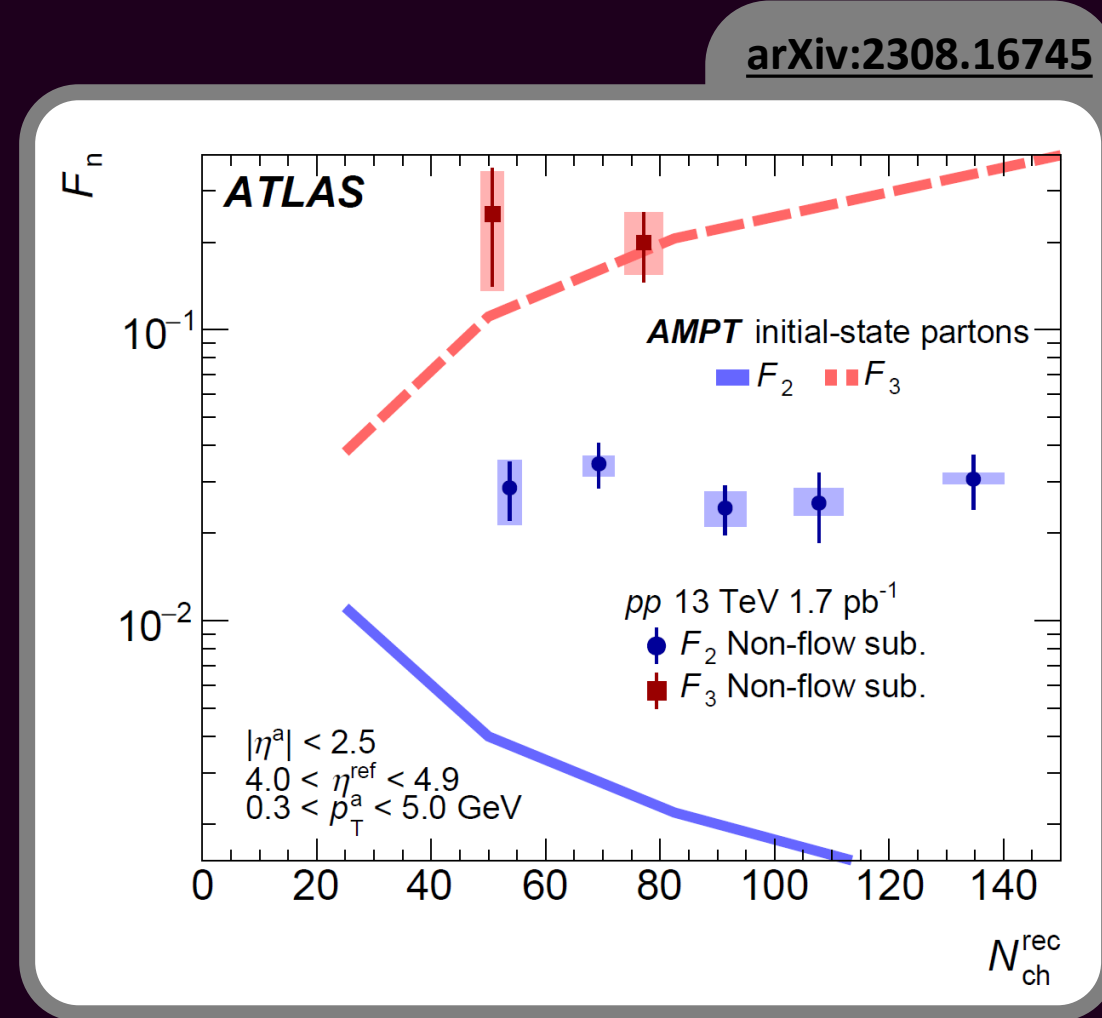
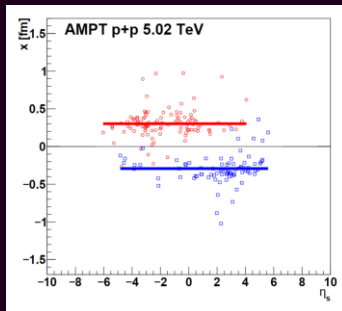
$F_n$  is the fractional change in  $v_{n,n}$  per a unit additional rapidity separation

$$v_{n,n} \propto 1 + F_n \eta_{separation}$$

Calculated by comparing  $v_{2,2}$  with  $|\Delta\eta|$  as low as  $|\Delta\eta| = 1.5$  to as large as  $|\Delta\eta| = 7.4$

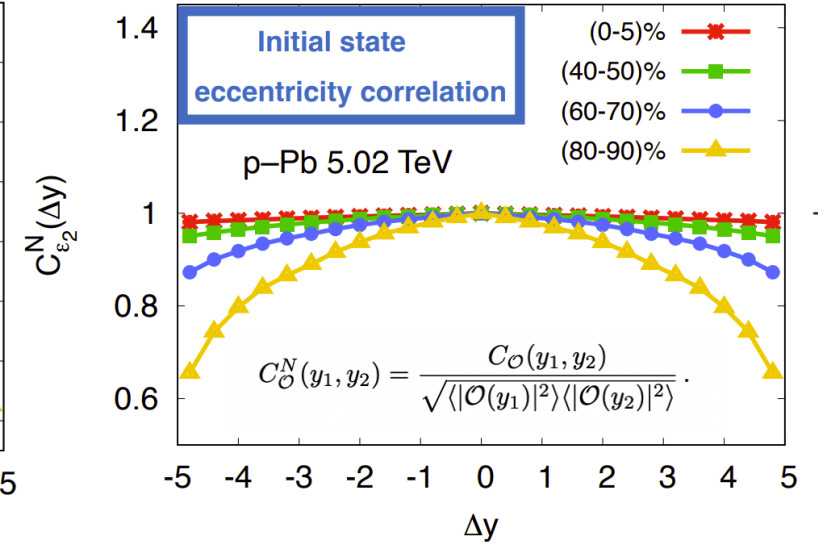
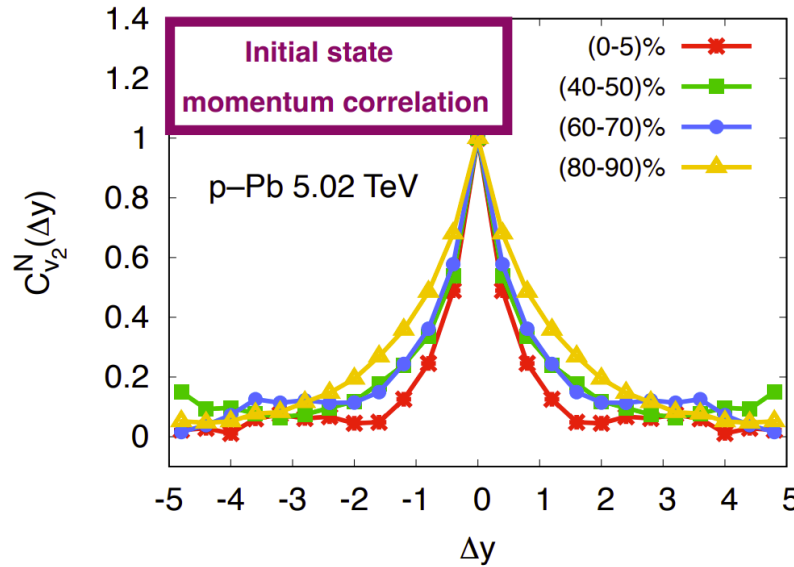
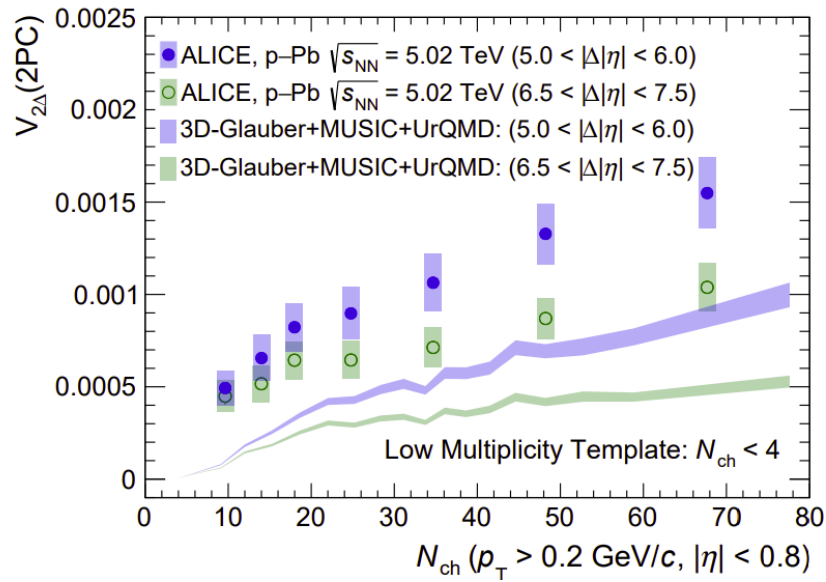
probes the shape of the initial state

- AMPT's simple 2-string model has insufficient longitudinal fluctuations



Nucleon-level geometry is inadequate to describe longitudinal dynamics

# How long is the long-range ridge?



arXiv:2504.02359

Phys. Rev. D 105, 094023

- $\sim 3\%$   $v_2$  for very large rapidity gap  $|\Delta\eta| > 7.5$
- Initial-state momentum anisotropy correlations die out with  $|\Delta\eta| > 2$
- Final-state effects (hydro+geometry) survive large separations in  $\eta$

Long range correlations are driven by final state effects

# Multi-particle cumulant high- $p_T$ $v_2$ in $pPb$

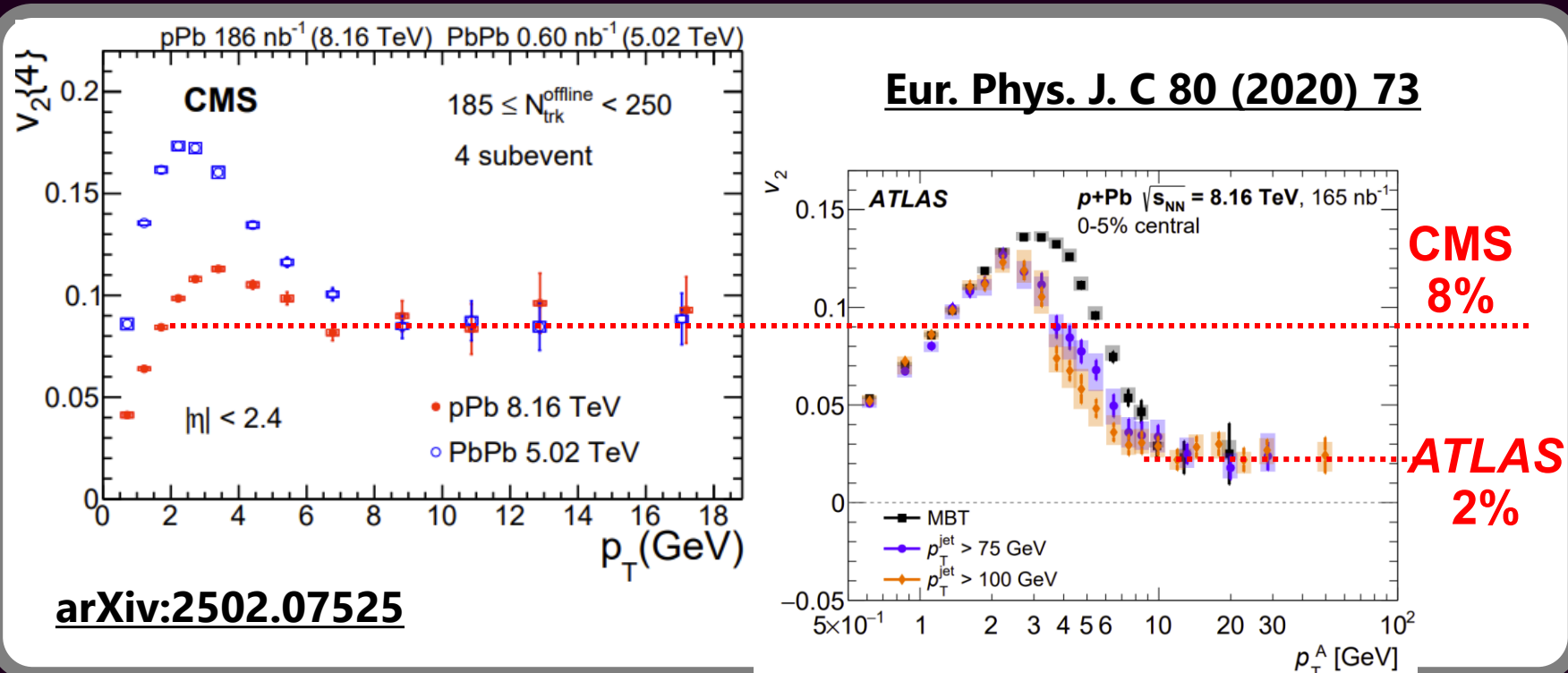
No Jet quenching in  $pPb$   
 $R_{AA} -1 < +/- 10\%$

High- $p_T$   $v_2$  observed by  
 ATLAS

Multiparticle cumulants  
 are less sensitive to  
 nonflow

Cumulants give access to  
 fluctuations

How  $v_2\{4\}=8\%$  and  
 $v_2\{2\}=2\%$ ?



$$v_n\{4\} = - \frac{d_n\{4\}}{(-c_n\{4\})^{3/4}}$$

$$d_n\{4\} = \langle\langle 4' \rangle\rangle - 2 \langle\langle 2' \rangle\rangle \langle\langle 2 \rangle\rangle$$

1 trigger particle      3 associated particles

↑ Fluctuations?

$$c_n\{4\} = \langle\langle 4 \rangle\rangle - 2 \langle\langle 2 \rangle\rangle^2$$

4 associated particles

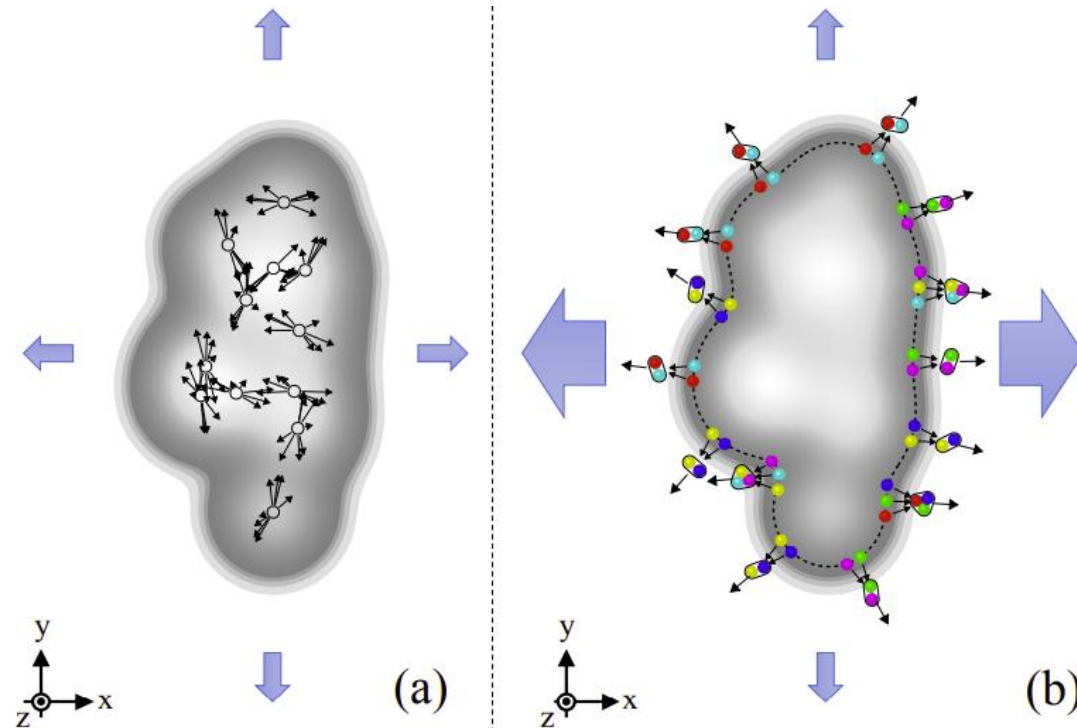
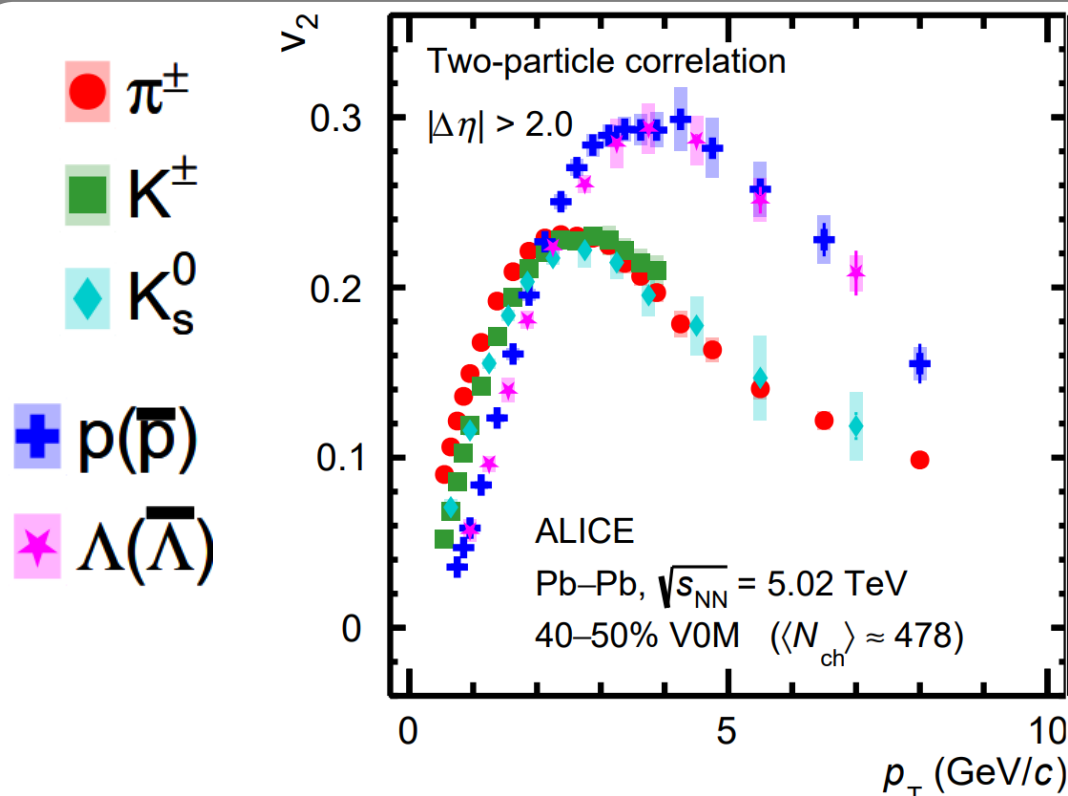
↓ Fluctuations?



Geometry  
Fluctuations  
**Transport**  
Novel small systems  
Future

# What is flowing?

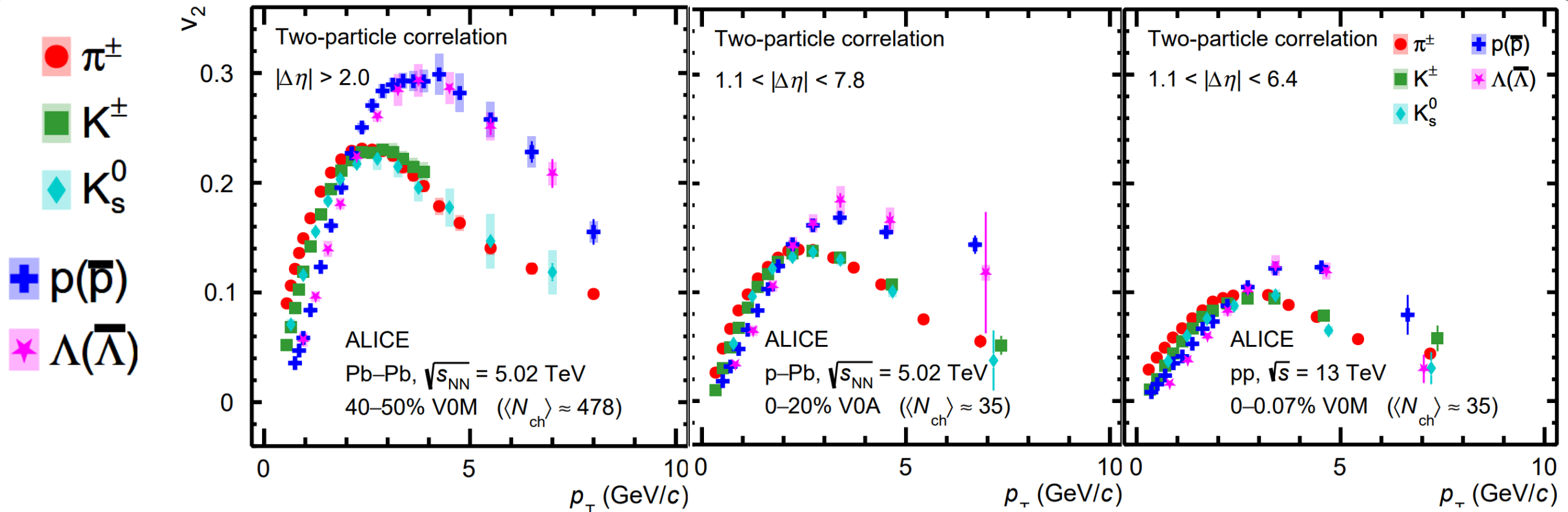
arXiv:2411.09323



- $p_T < 3$  GeV: Mass ordering arises from the interplay between radial expansion and anisotropic flow of the velocity field which pushes heavier particles to higher  $p_T$
- $3 < p_T < 5$  GeV: Baryon and meson grouping arises from parton coalescence
- Non-flow subtraction with the template fit

# Partonic flow in $pp$

arXiv:2411.09323

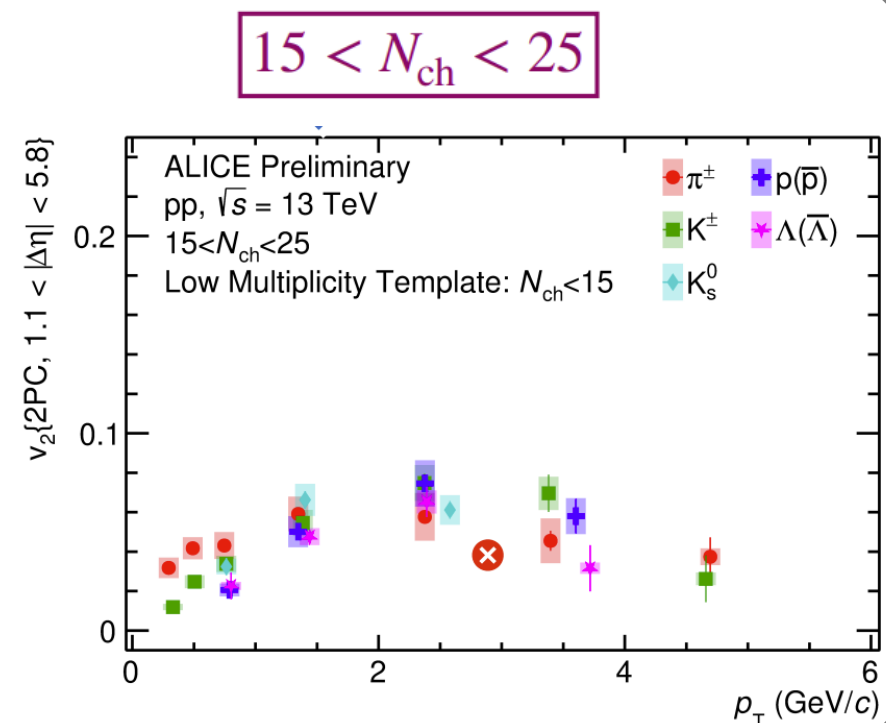
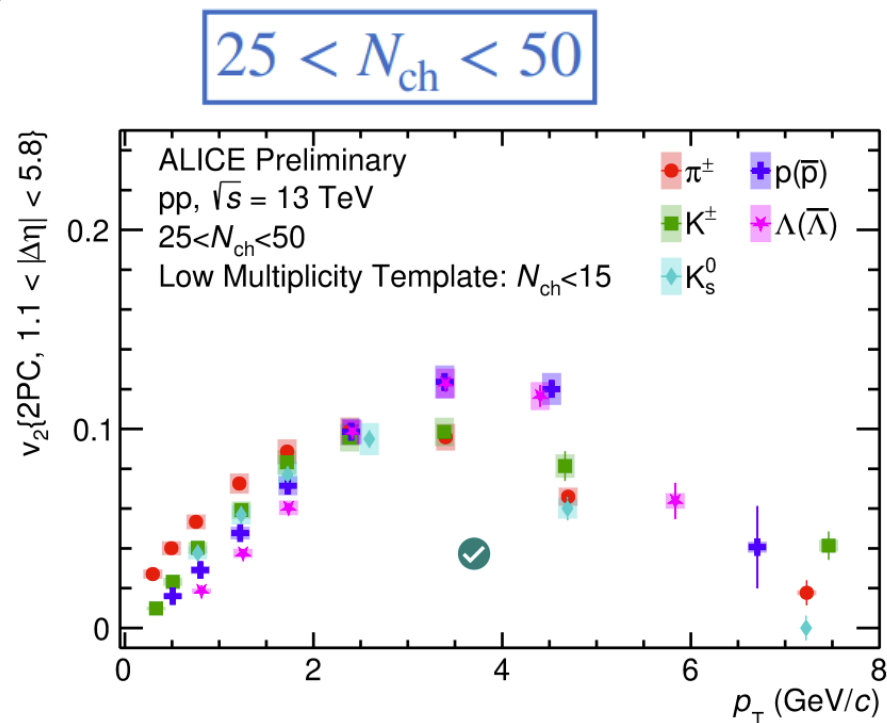


- $p_T < 3$  GeV: Mass ordering arises from the interplay between radial expansion and anisotropic flow of the velocity field which pushes heavier particles to higher  $p_T$
- $3 < p_T < 5$  GeV: Baryon and meson grouping arises from parton coalescence
- Non-flow subtraction with the template fit

Strong evidence of a deconfined flowing parton medium in  $pp$  and  $pPb$  19

# The breakdown of partonic flow?

- Breakdown of baryon-meson grouping in low multiplicity  $pp$
- Mass ordering persists
- Nonflow subtract issues are largest at high  $p_T$  and low multiplicity
- Are probes like these **less sensitive to uncertainties in initial state geometry?**



small collision systems presented either a single baryon or meson  $v_2$ , limiting the opportunity to explore potential groupings among baryons and mesons [45, 46]. Other similar measurements have not shown a clear grouping and splitting of baryon and meson  $v_2$  at intermediate  $p_T$  in small systems, differing from similar measurements in heavy-ion collisions [24]. This difference may arise from difficulties in accounting for non-flow effects in small systems. The results presented in Fig. 2 show, for the first time,

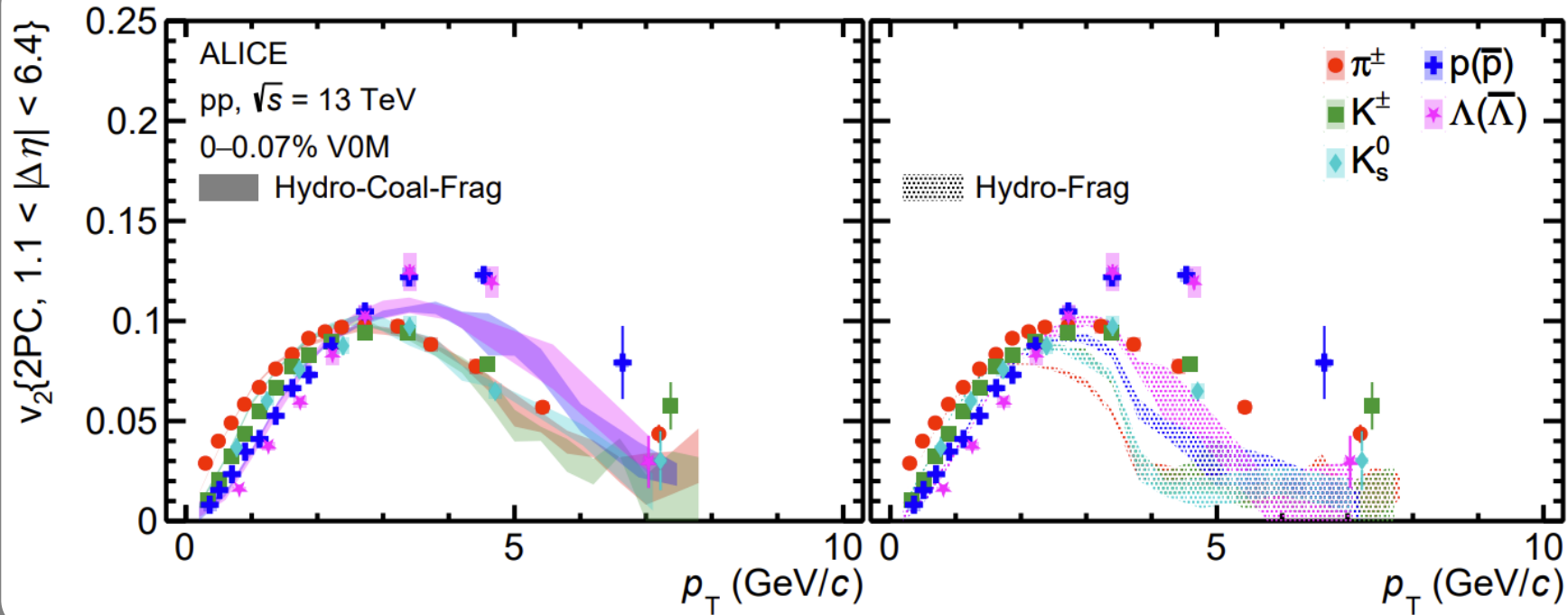
**Nonflow subtraction is imperative for precision measurements**

**Are there tools like this to probe the breakdown of hydro?**



# Partonic flow in $pp$

arXiv:2411.09323

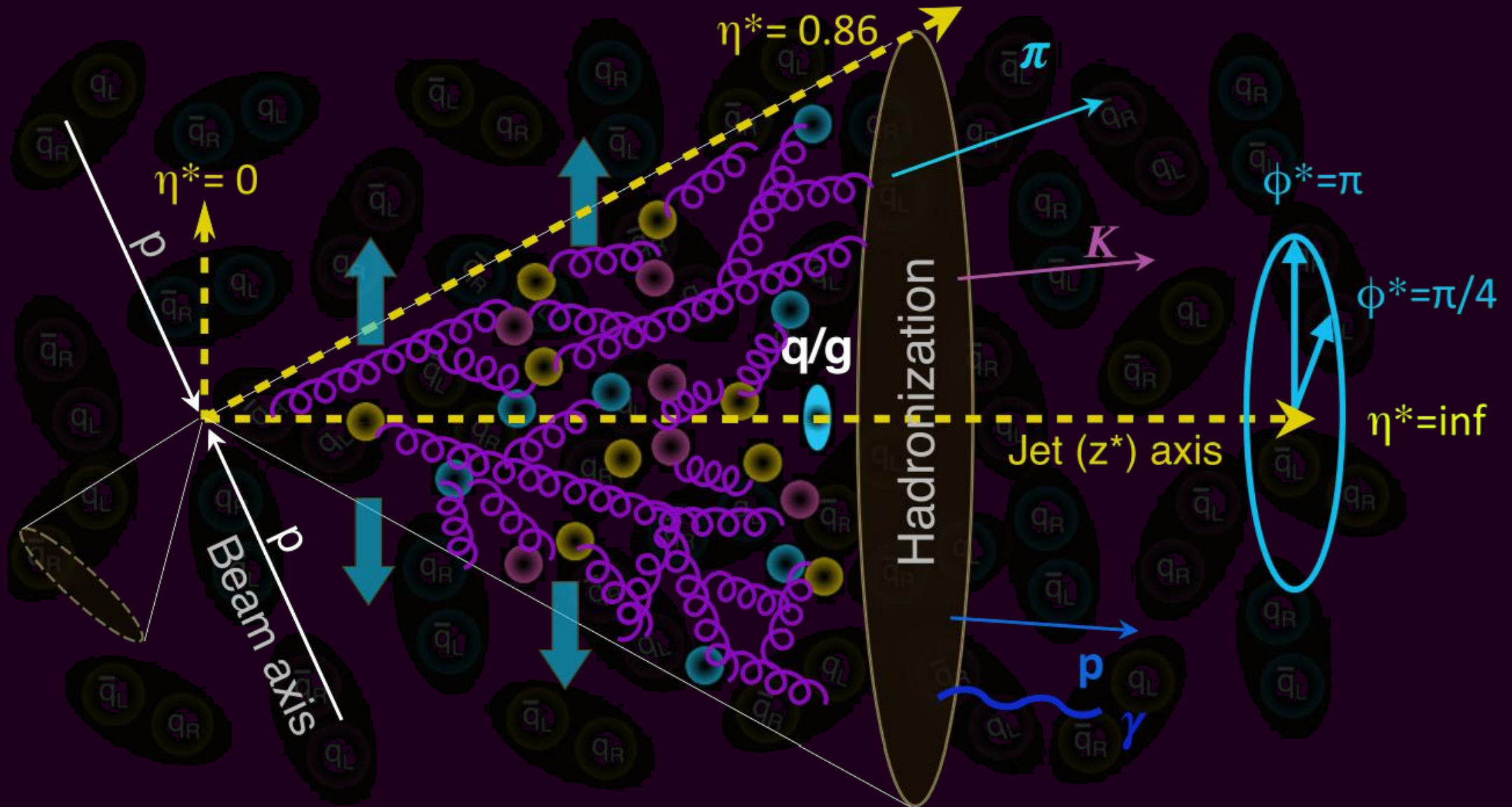


- Coalescence between medium quarks and fragmentation quarks is necessary for adequately describing  $pp$  results

Coalescence is necessary but not sufficient to reproduce baryon-meson grouping

Geometry  
Fluctuations  
Transport  
**Novel small systems**  
Future

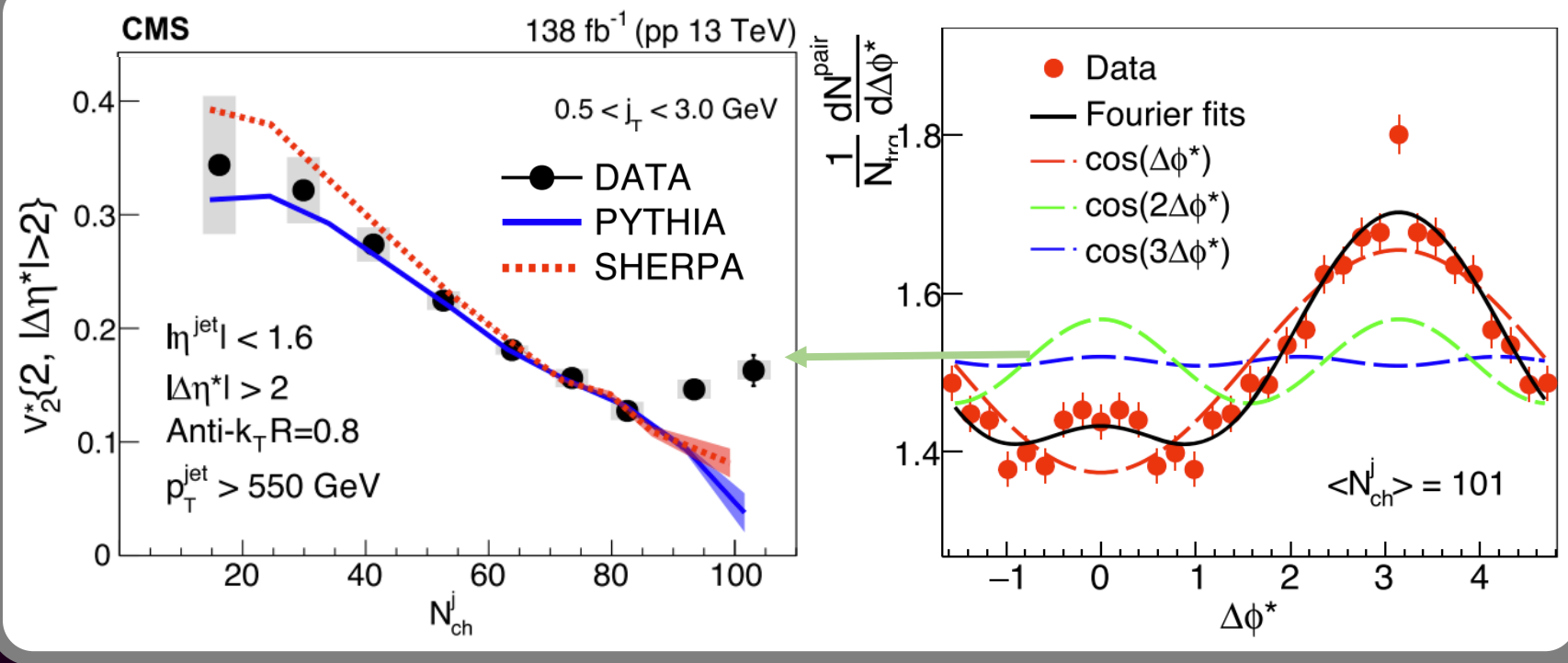
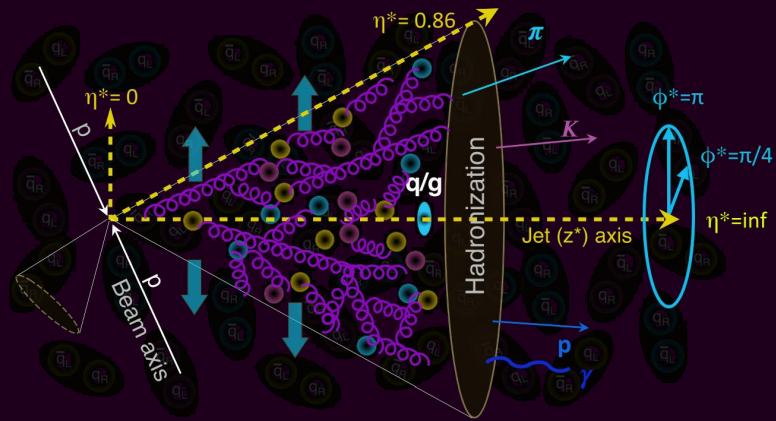
# High multiplicity jet system



# Enhanced $v_2$ within high-multiplicity jets

Use jet axis as beam axis

$$\phi, \eta, p_T \rightarrow \phi^*, \eta^*, j_T$$



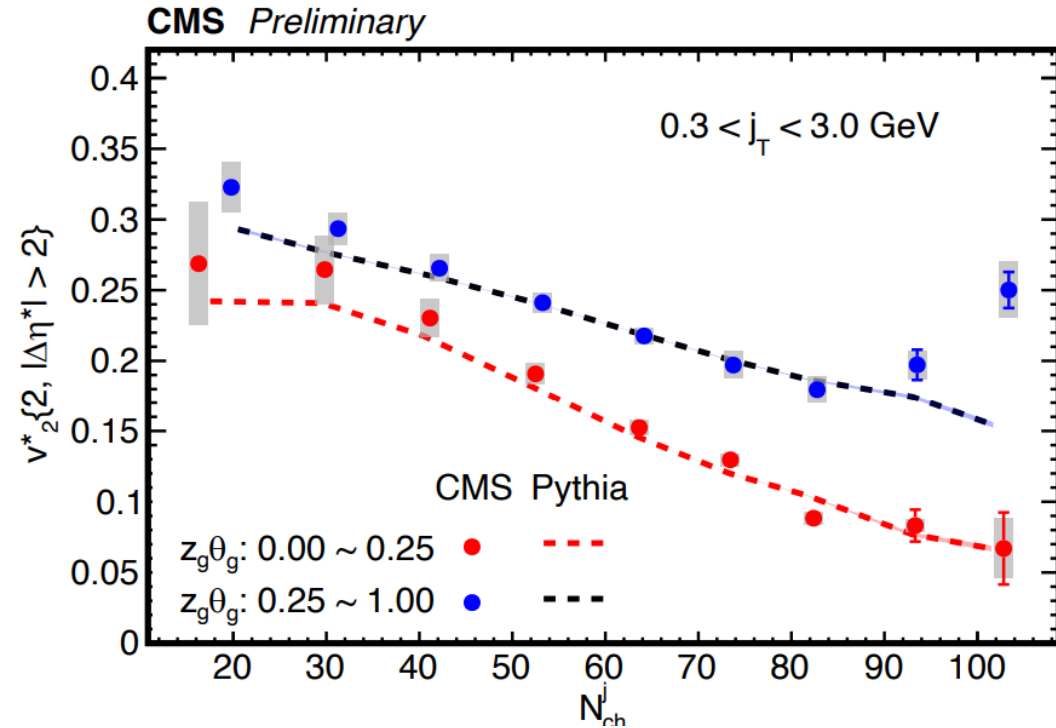
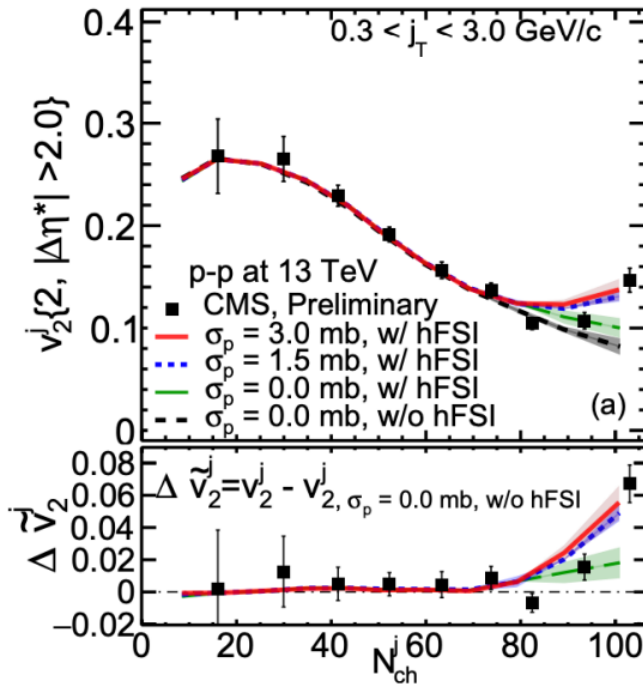
***PRL 133 (2024) 142301***

- 138 fb<sup>-1</sup>: Full Run 2 statistics used
- In the highest multiplicity jets a flow-like elliptic enhancement is observed
- Measured a very large  $v_2 = 16.29\%$
- **What's next:**  $j_T$  dependence, flow factorization, ...?

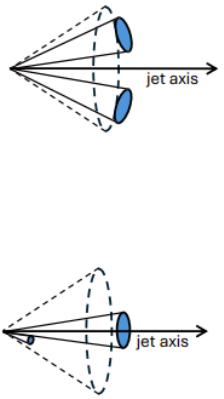
# Enhanced $v_2$ within high-multiplicity jets

arXiv:2401.13137

Theory comparison using PYTHIA jet shower model combine with parton re-scattering ZPC



Xiaoyu Liu



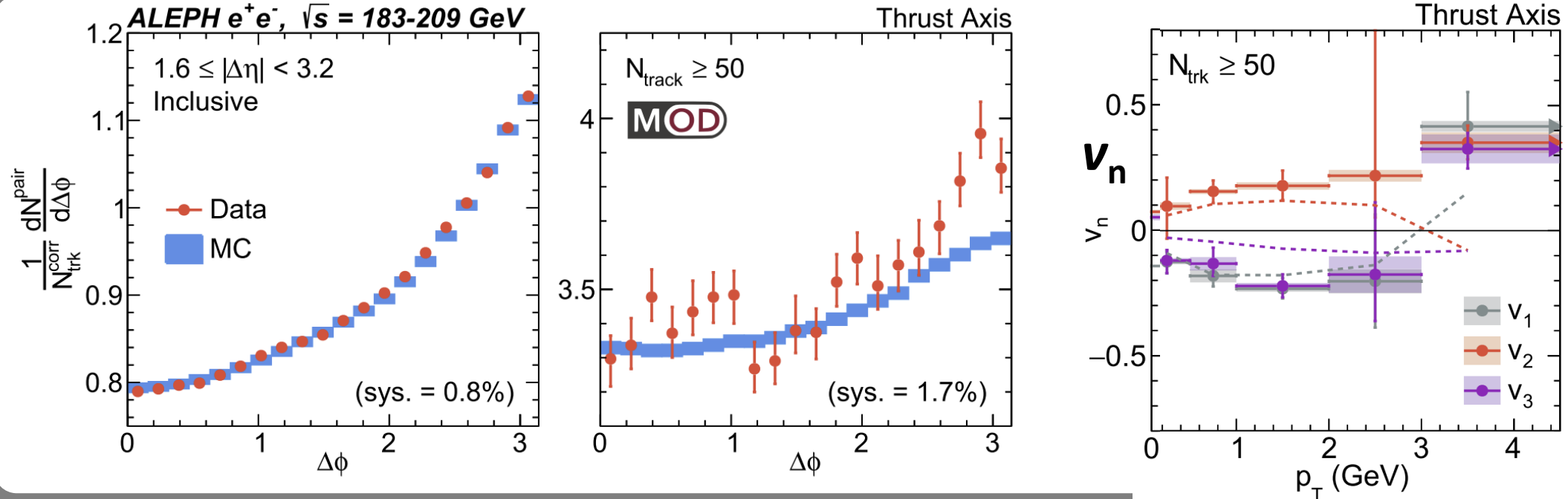
PYTHIA AMPT-like

Does pQCD provide us a well-constrained geometry?

- New results: grouping jets in terms of substructure
- 2-prong jets show much larger collective effects



# Azimuthal anisotropy in $e^+e^-$



PLB 856, 138957 (2024)

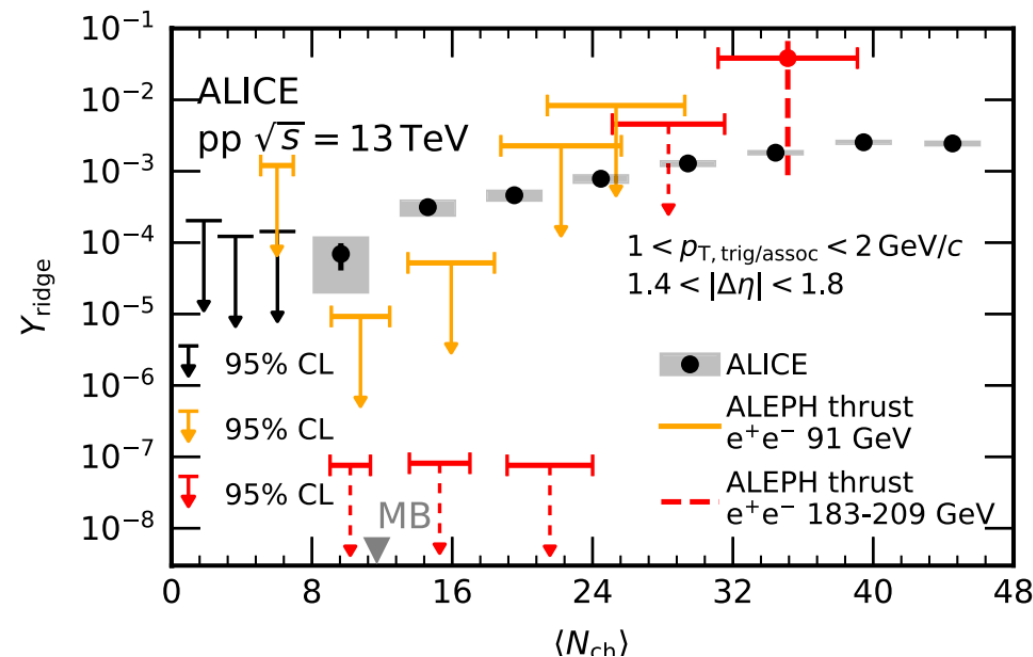
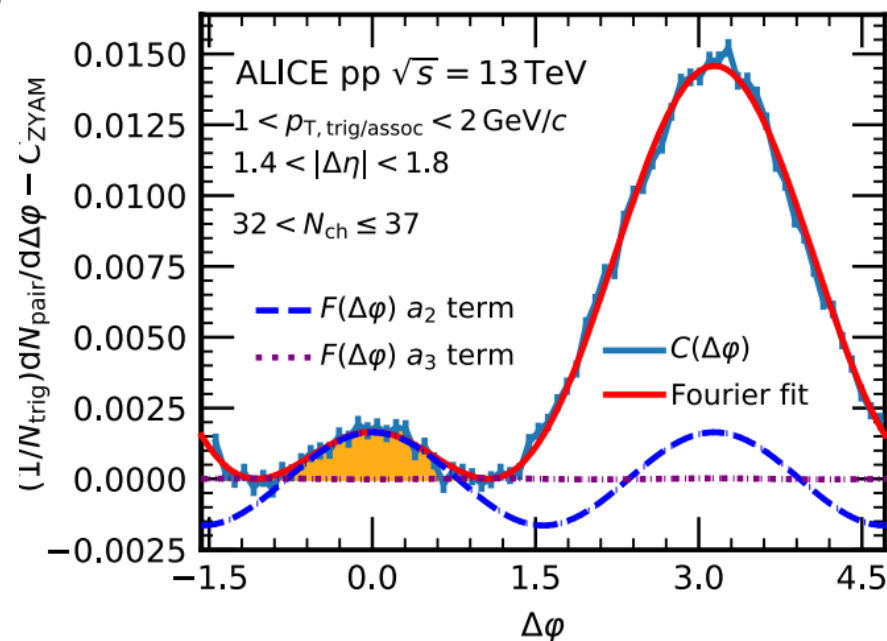
- 2-particle correlation performed in the thrust axis (much like the high multiplicity jets)
- The  $V_n$  coefficients have poor data-MC agreement when pushed to the highest multiplicities
- One can interpret this as signs of other effects in data not simulated and perform nonflow subtraction
- $v_n^{\text{flow}} = v_n^{\text{data}} - v_n^{\text{MC}}$ : this method yields a positive  $v_2$  but a negative  $v_3$  which is not interpretable in the geometry-driven interpretation

Issues with nonflow sub. make qualitative conclusions of collectivity hard



# Ridge yields in the smallest systems: $pp$ & $e^+e^-$

Extraction of nearside per-a-trigger yields: integral of **Orange** region



Phys. Rev. Lett. 132

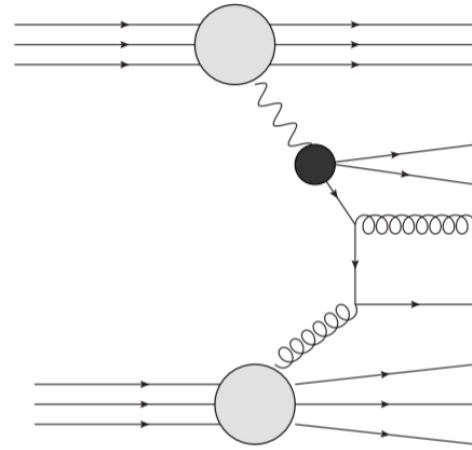
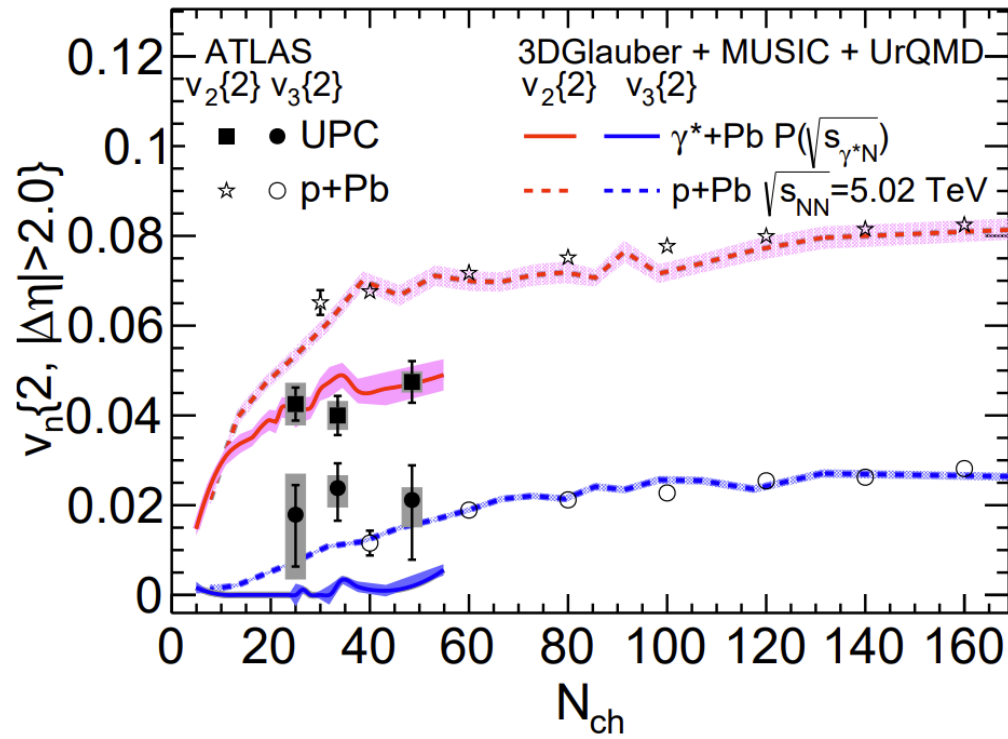
Nearside per-a-trigger yields are challenging to interpret

- yield depends on shape and scale of nonflow
- Per-trigger yield grows with growing  $N_{ch}$  – results don't necessarily demonstrate a increasing  $v_2$
- have to be compared to theory with flow and nonflow

Can make direct comparisons to  $e^+e^-$

**Only quantitative statements can be made about nonflow**

# $\gamma$ +Pb azimuthal anisotropy at the LHC



## $\gamma$ +Pb collisions

$\langle W_{\gamma N} \rangle$  range 0.5 – 1.0 TeV  
 Single-sided nuclear breakup  
 Mostly hadronic-like collisions  
 Selected with rapidity gap

Observe Nonzero  $\gamma$ Pb  $v_2$

comparison to 3D Glauber + MUSIC + UrQMD

Why is  $v_2(\gamma^* + \text{Pb}) < v_2(\text{p} + \text{Pb})$

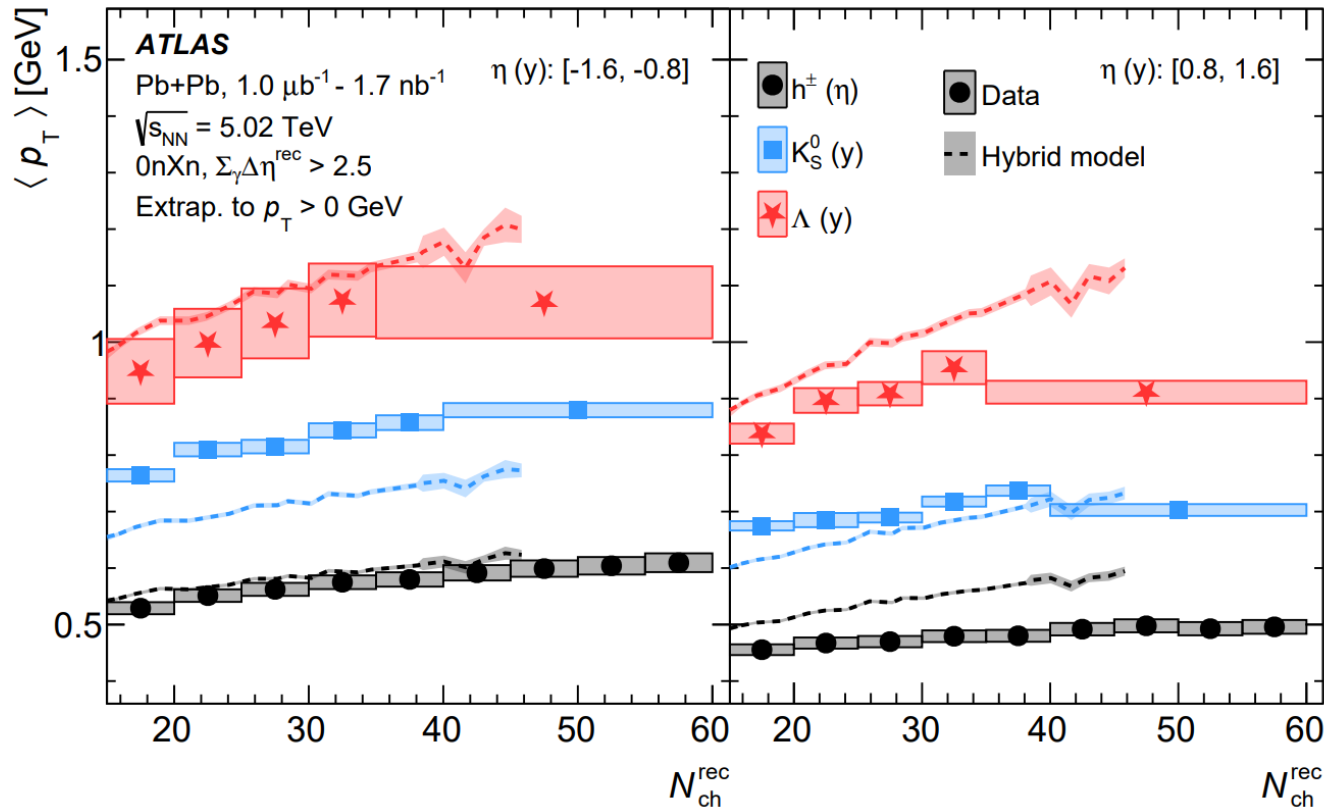
$\gamma$ +Pb correlations performed at forward rapidities  
 which suppresses observed  $v_2$  due to enhanced longitudinal decorrelation

Phys. Rev. Lett. 129, 252302  
 Phys. Rev. C 104, 014903



# Identified particle $\langle p_T \rangle$ in $\gamma$ +Pb

arXiv:2503.08181



➤ Why  $\langle p_T \rangle$  with  $N_{ch}$

- Higher energy density achieved in higher multiplicity collisions leads to stronger radial expansion.
- Signature of collectivity

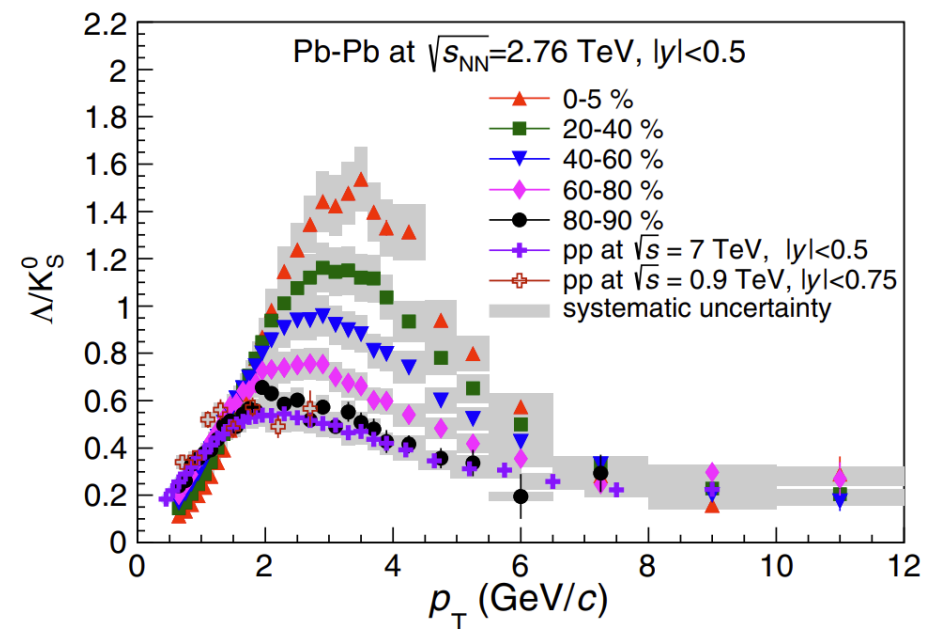
➤ Larger  $\langle p_T \rangle$  in the Pb-going direction

➤ Qualitative agreement with the Hydro model

Behavior in data is consistent with qualitative picture of radial flow

# Baryon anomaly in $\gamma A$

[Phys. Rev. Lett. 111, 222301](#)



- Evidence of hadronization in a dense environment
- Observe large baryon enhancement at mid- $p_T$  in  $\gamma A$ , similar to pPb
- Possibly see larger baryon enhancement in the Pb going direction

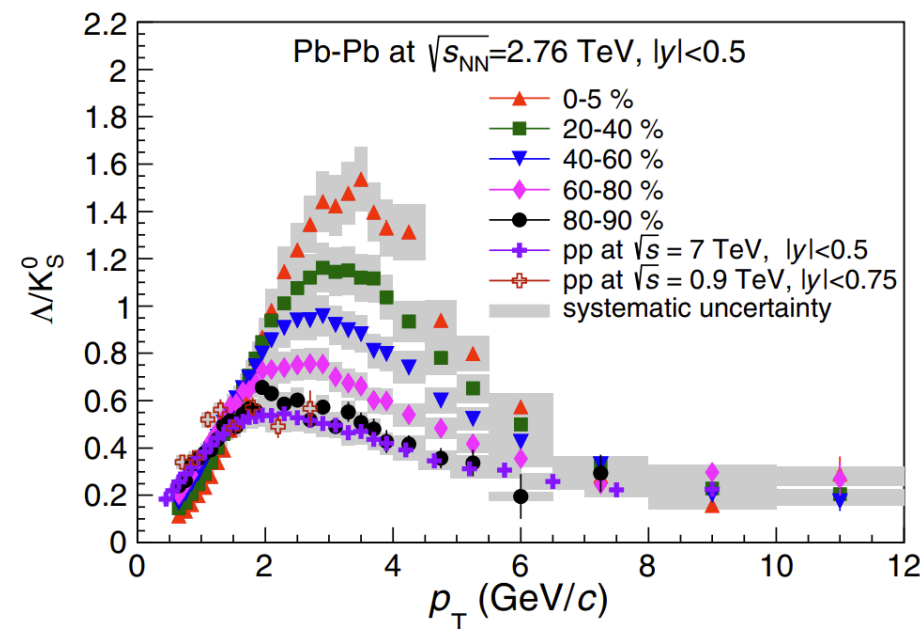
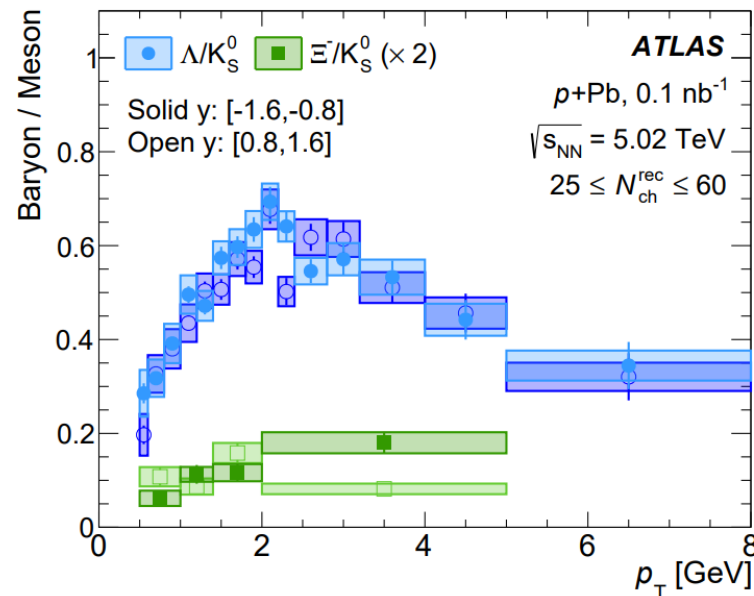
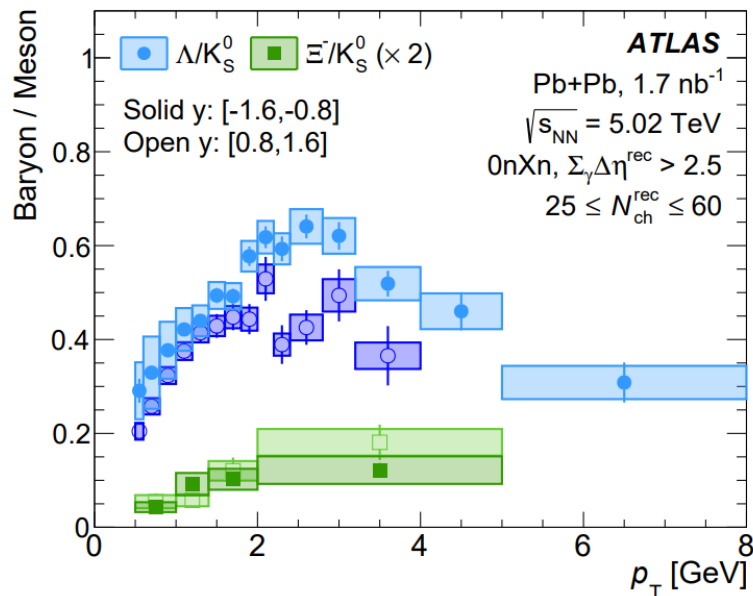
# Baryon anomaly in $\gamma$ +Pb

$\gamma$ +Pb

arXiv:2503.08181

$p$ +Pb

Phys. Rev. Lett. 111, 222301



- Evidence of hadronization in a dense environment
- Observe large baryon enhancement at mid- $p_T$  in  $\gamma$ Pb, similar to pPb
- Possibly see larger baryon enhancement in the Pb-going direction

Observe large baryon enhancement at mid- $p_T$  in  $\gamma A$ , similar to pPb

# The smallest systems

Observing the minimal criteria for QGP formation?

Uncertainty in  
nonflow subtraction  
→ no firm  
conclusion

Maybe?

$N_{ch}$  threshold observed

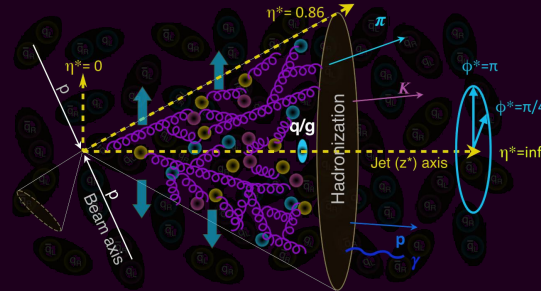
So much like previously explored  
systems, how is  $\gamma$ Pb unique?

1. Boast / decorrelations
2. Unique quantum numbers

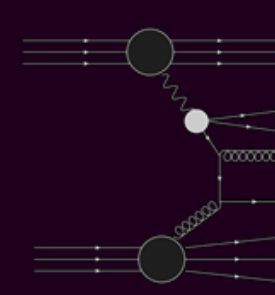
$\gamma+p$



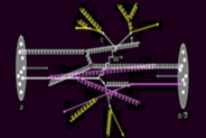
Jets



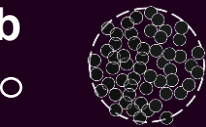
$\gamma+A$



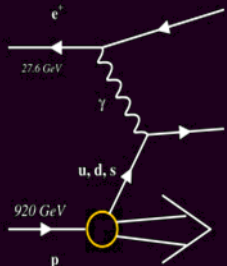
$pp$



$p+Pb$



$e+p$



$e^+e^-$





Geometry  
Fluctuations  
Transport  
Novel small systems  
Future

# Future: LHC $^{16}\text{O}+^{16}\text{O}$ and $^{20}\text{Ne}+^{20}\text{Ne}$ runs

## The plan in July

- $\sim 2 \text{ nb}^{-1}$  of 5.36 TeV  $^{16}\text{O}+^{16}\text{O}$
- $< 1$  day of 5.36 TeV  $^{20}\text{Ne}+^{20}\text{Ne}$

## Advantages of O+O and over pPb

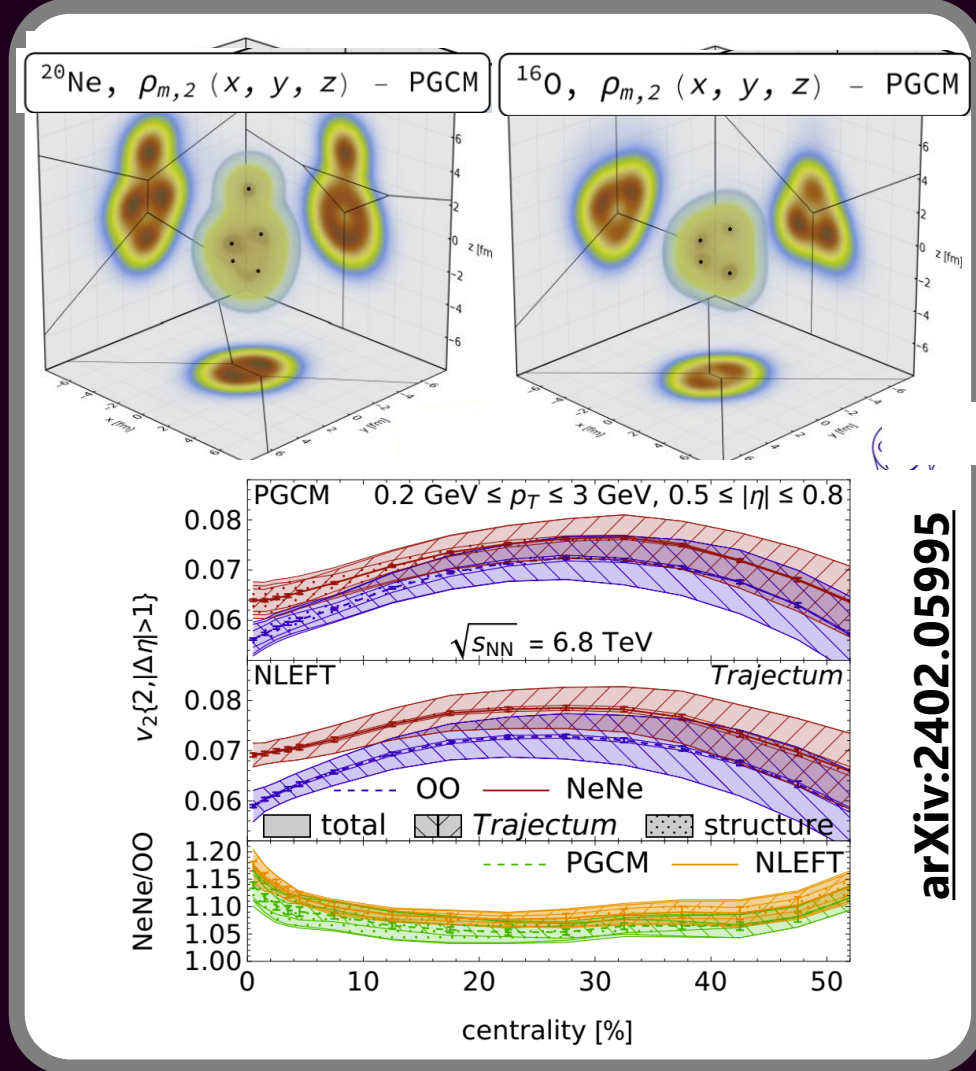
- Initial-state geometry better under control
- less sensitivity to proton substructure
- Symmetric collision system
- First identical system at RHIC and LHC

## Comparisons to Ne+Ne

- We can compare 2 small systems of similar size but different initial-state geometries
- Ratios of O+O to Ne+Ne can test linear response to geometry

Density  
map of  $^{16}\text{O}$   
and  $^{20}\text{Ne}$

$v_2$  of OO and  
NeNe for  
geometry  
models PGCM  
(top) and  
NLEFT (middle)

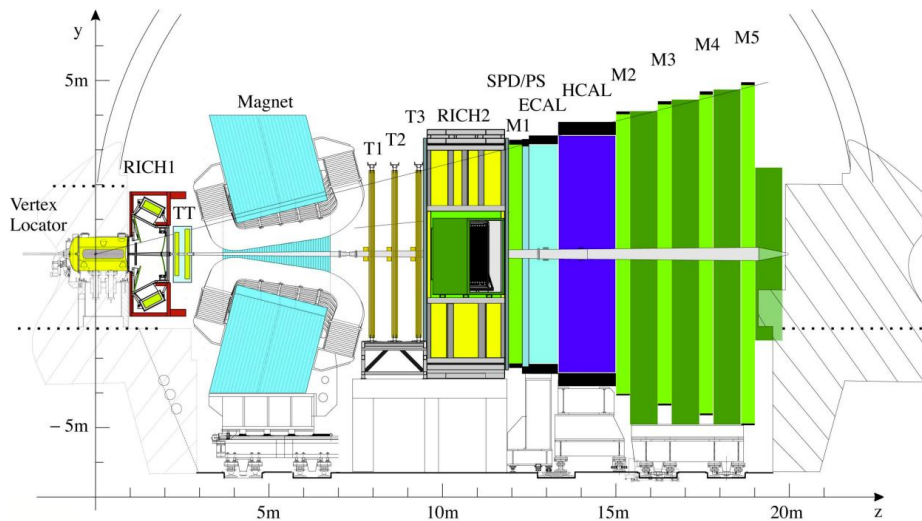
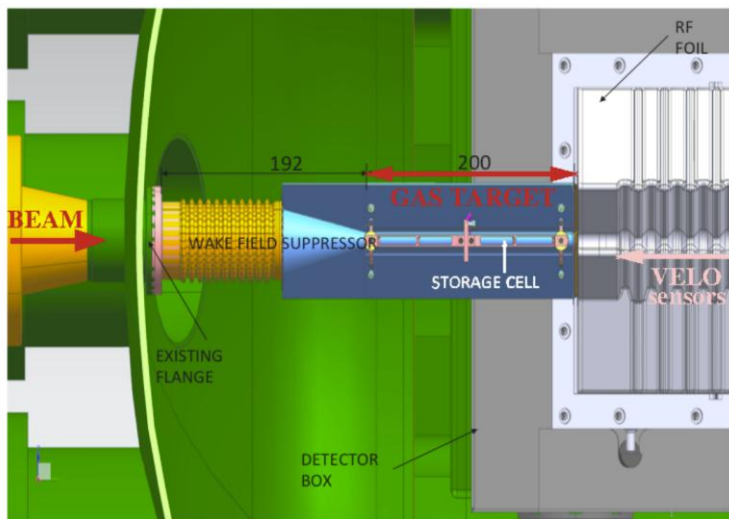


arXiv:2402.05995

**Geometry better constrained → sharper test of hydrodynamics**

# LHCb + SMOG2

[More details here](#)



- SMOG2 provides fixed target capabilities which allows for many targets
- Could repeat the RHIC small system energy scan but at similar energies as RHIC
  - Pb+p, Pb+d, Pb+He
- Ratios of similar collisions systems at different energies can cancel initial state geometry effects and yield precision measurements of transport
- Important collective studies with the closing of RHIC

# Summary

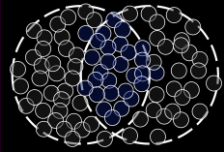
- We have learned a lot:
  - hydro qualitatively works
  - fluctuations are important
  - sub-nucleonic structure is necessary
  - Nonflow has to be systematically considered
- The findings of the field are largely qualitative
- In many small systems, the uncertainties in and the general complexity of the initial state prevents quantitative statements on the nature of collectivity in small systems and how for example it differs from large systems
- Attempts to find the smallest systems or systems with very simple initial state geometry (e.g.  $e^+e^-$ ) do not produce enough particles to suppress uncertainties of nonflow removal
- Our tool kit is growing rapidly
- Light ion collisions maybe be a sweet spot, for our current level of understanding, because of its sensitivity to out-of-equilibrium effects, relatively small amount of nonflow, and its better constrained geometry

# Backup

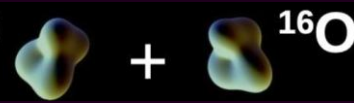
# A push towards precision

The near future lies in a regime constrained theoretical uncertainties and well understood removal of non-flow

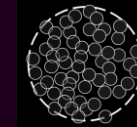
Pb+Pb



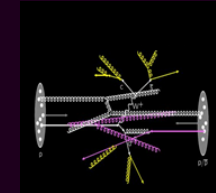
$^{16}\text{O}$



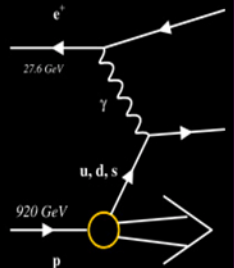
p+Pb



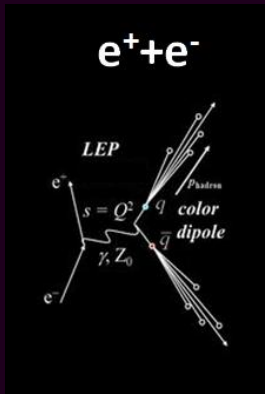
pp



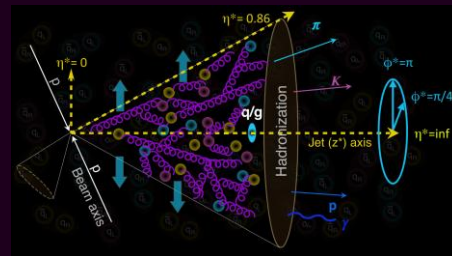
e+p



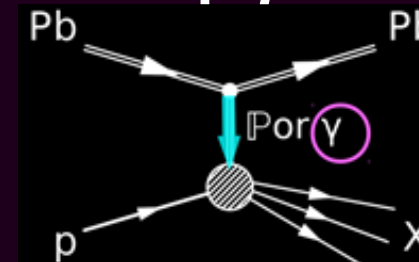
e<sup>+</sup>e<sup>-</sup>



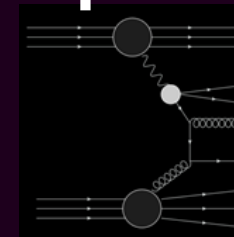
Jets



$\gamma$ +p

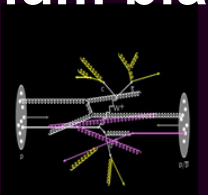


$\gamma$ +A



Minimum bias

pp

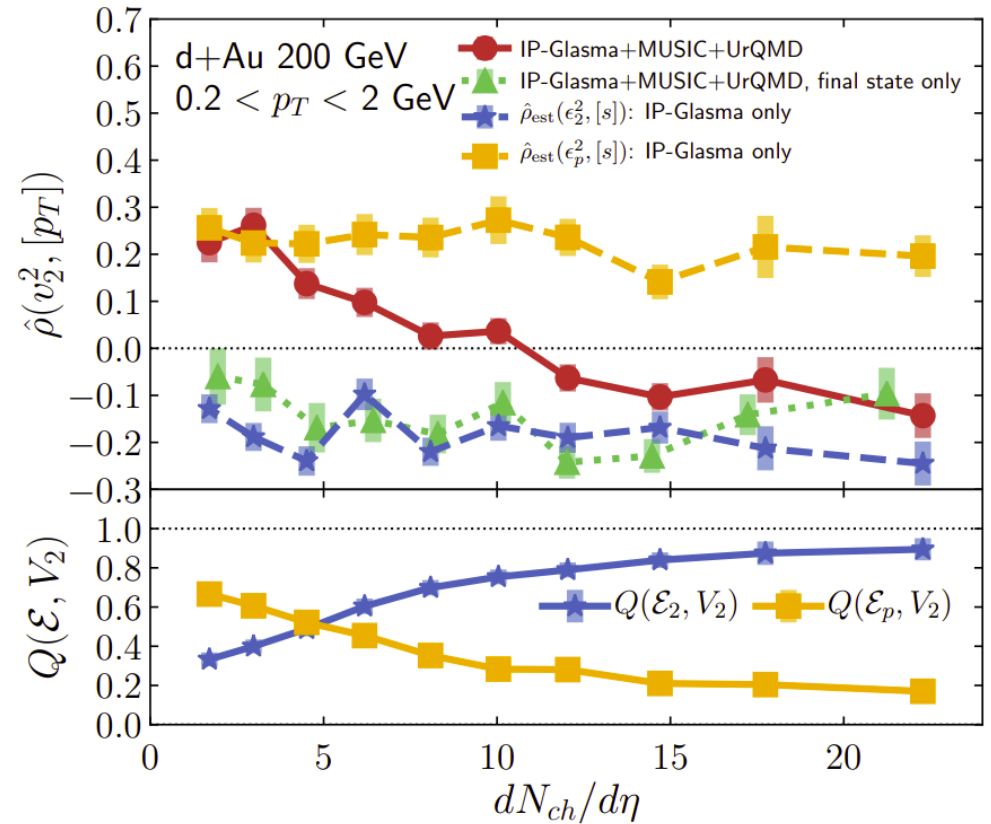
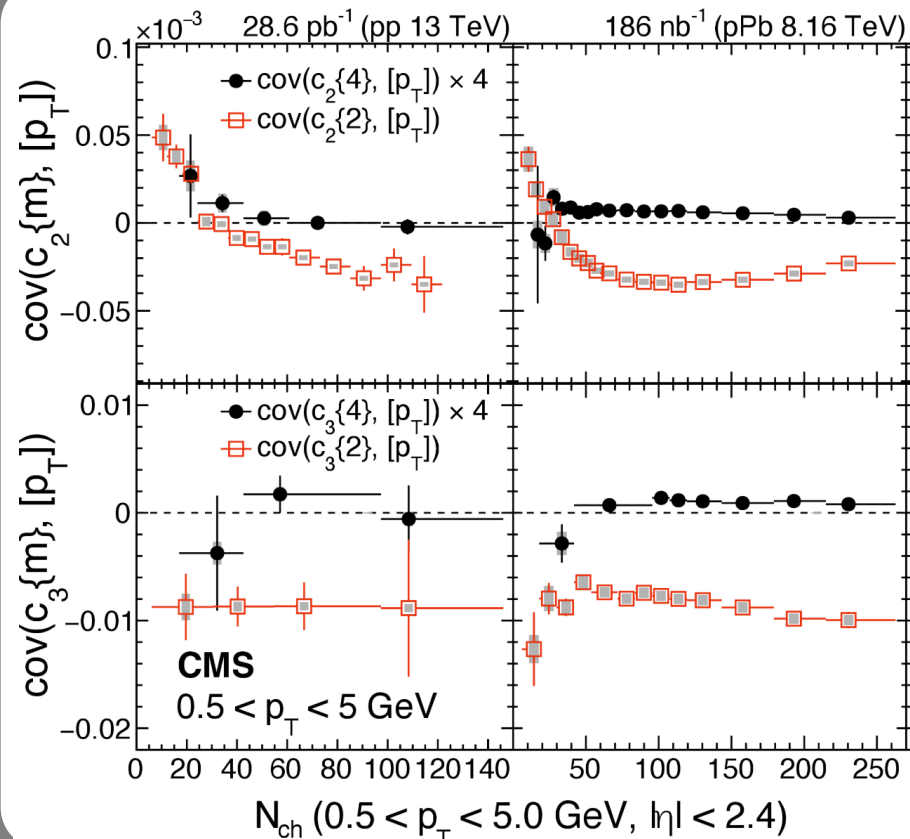


Without this these will remain a disconnected set of facts



# Vn-pT with CMS in small systems

$$\text{cov}(c_n\{2\}, [p_T]) = \Re \left\langle \sum_{a,b} \exp^{in(\phi_a - \phi_b)} ([p_T] - \langle [p_T] \rangle) \right\rangle$$



Phys. Rev. C 110

- contributions from nonflow can create sign change at low  $N_{ch}$
- Without subevents, it's hard to interpret the effects from nonflow even on a qualitative level
- Vn pt [arXiv:2410.04578](https://arxiv.org/abs/2410.04578)

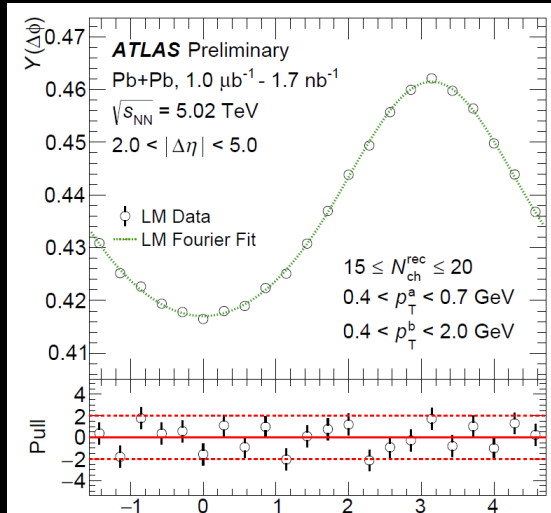
# Template fit in photonuclear



High-multiplicity (HM) correlation data



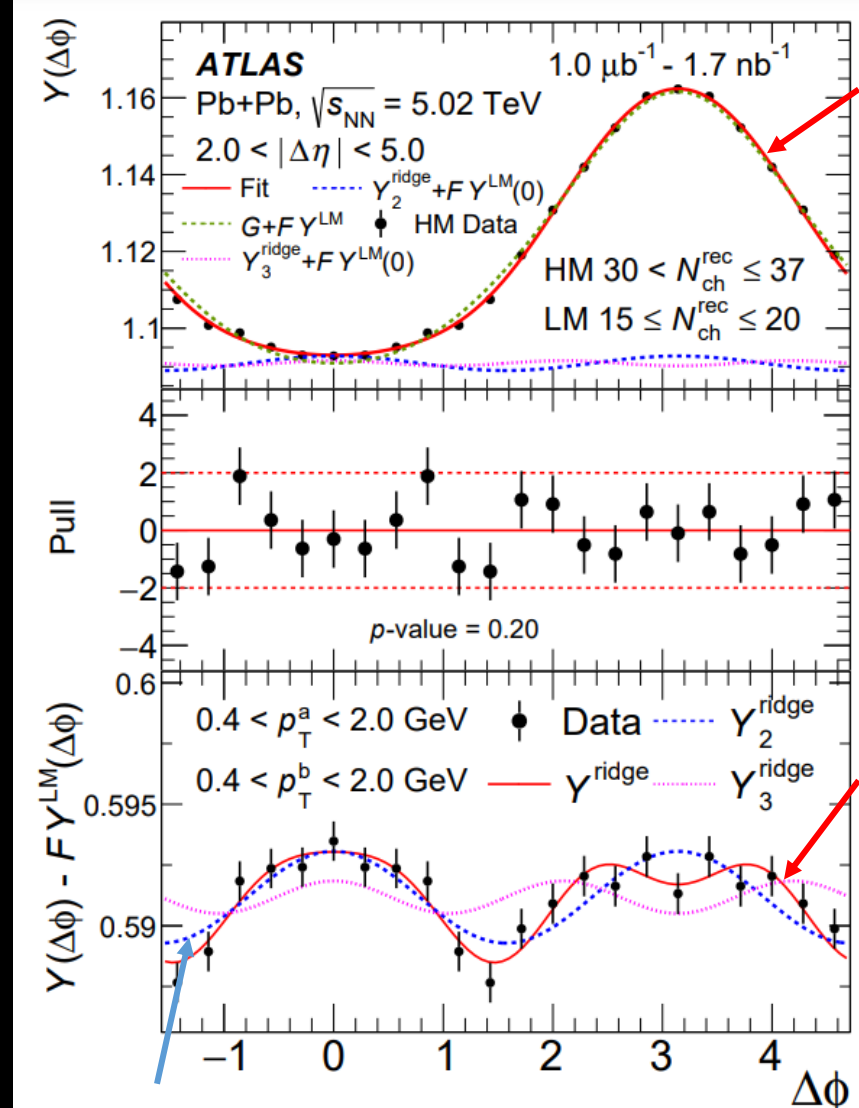
Low multiplicity (LM)  
template for jet/non-flow correlation



## Nonflow subtraction

- HM fit with LM data and flow coef.
- HM and LM assumed to have same flow shape
- Different LM selection leads to similar results

$$Y^{\text{HM}}(\Delta\phi) = FY^{\text{LM}}(\Delta\phi) + G \left\{ 1 + 2 \sum_{n=2}^3 v_{n,n} \cos(n\Delta\phi) \right\}$$



fit

fit

arXiv:2101.1077

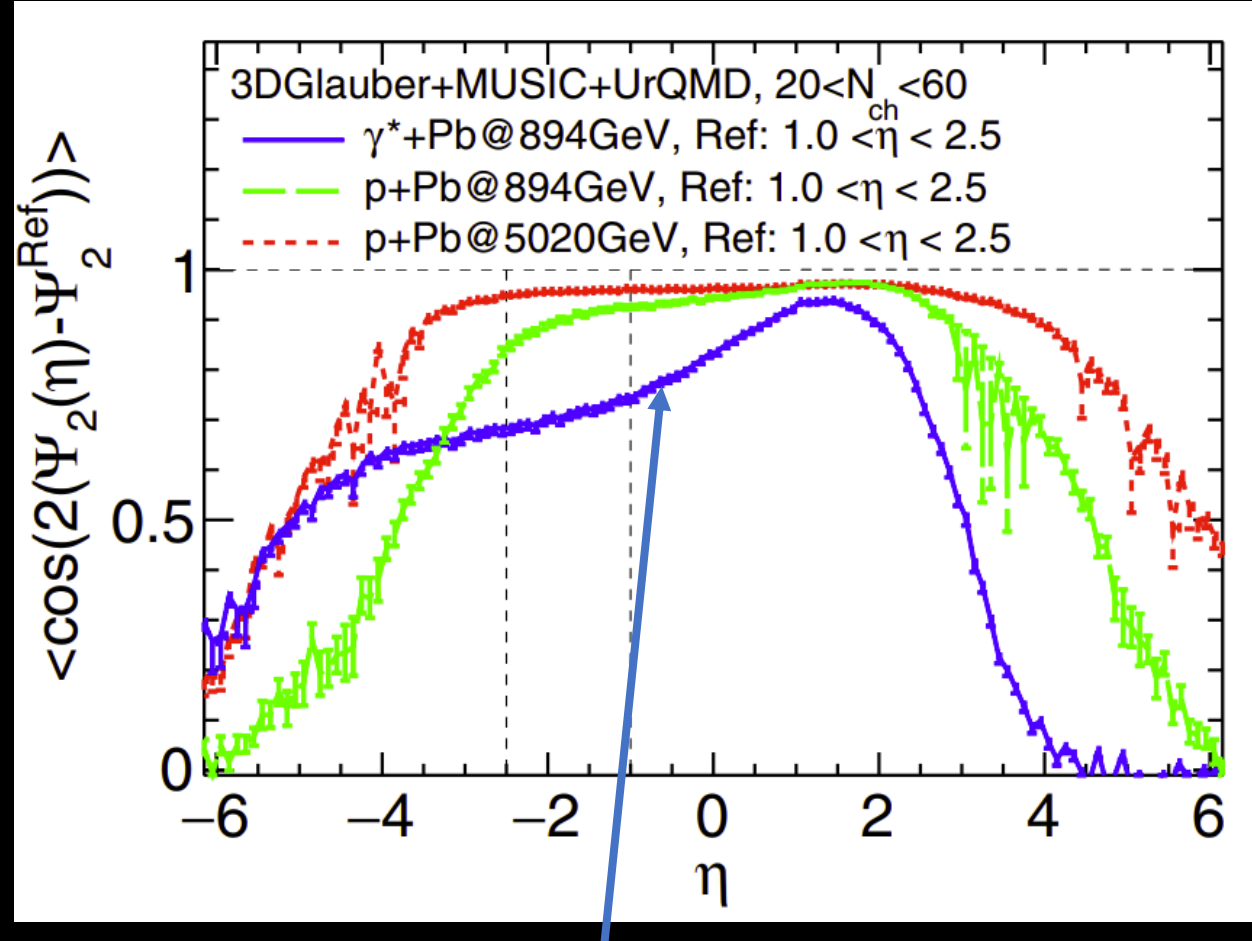
1

After nonflow subtraction clear  $\cos(2\Delta\phi)$  modulation

# Why is $\gamma$ Pb $v_2$ smaller

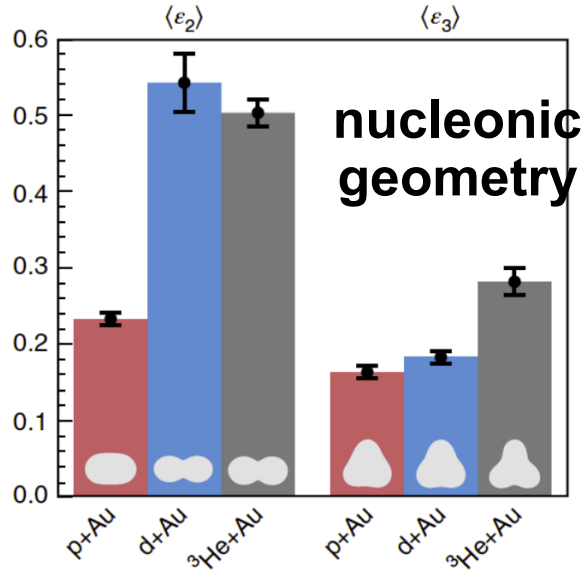
EPJ Web of Conferences 276, 01002 (2023) SQM22

- Correlations in small systems are performed with a rapidity gap between the particles
- The event plan can fluctuate between these rapidities and decreasing the observed  $v_2$
- This effect is larger at forward rapidities.
- Because  $\gamma$ Pb is so boosted the "forward rapidities" are probes relative to other systems with the ATLAS detector.

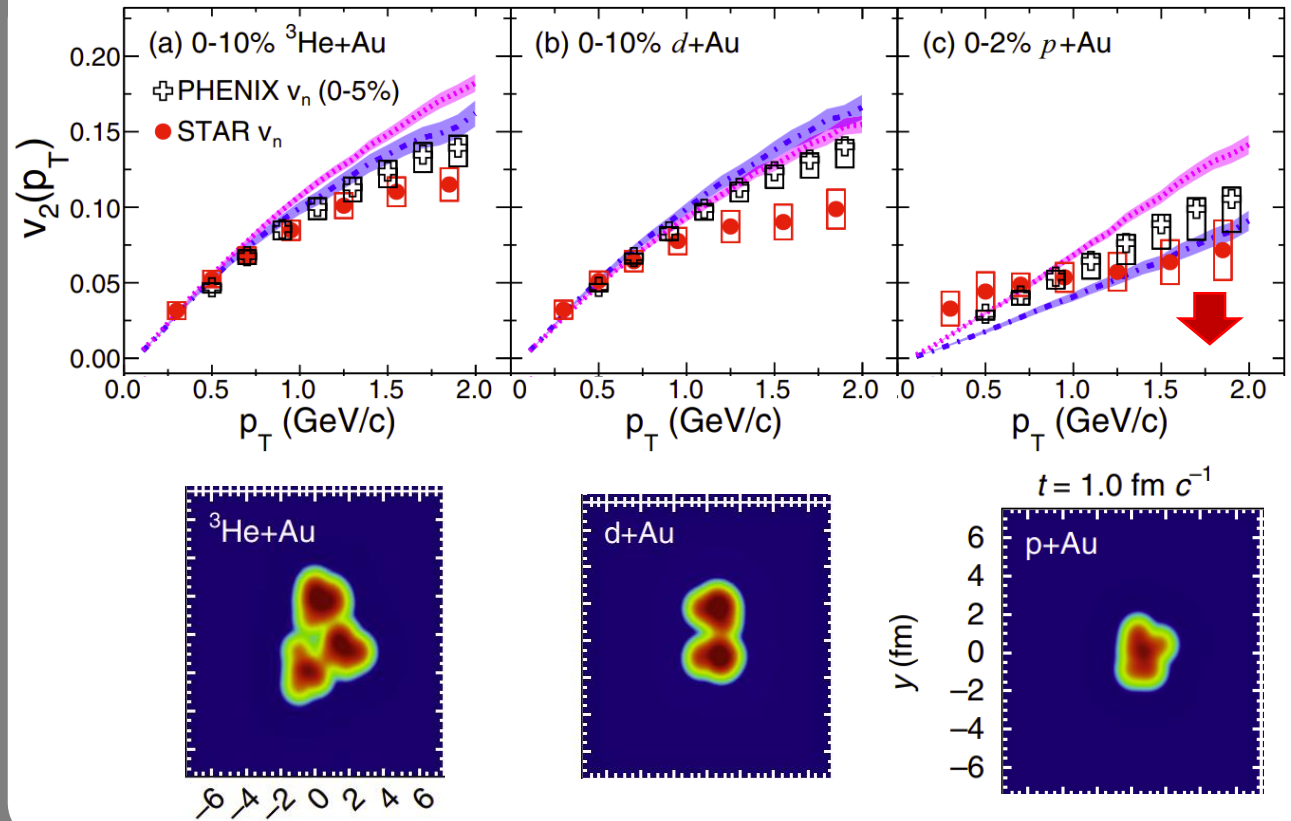


**I will show measurements that reflects the slope of these lines next!**

# Geometry: flow in p+Au d+Au He+Au



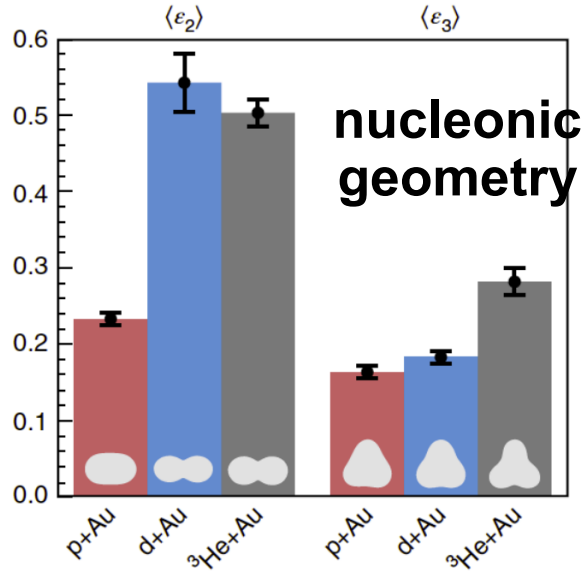
Nature Physics 15, (2019)



Phys. Rev. C 110

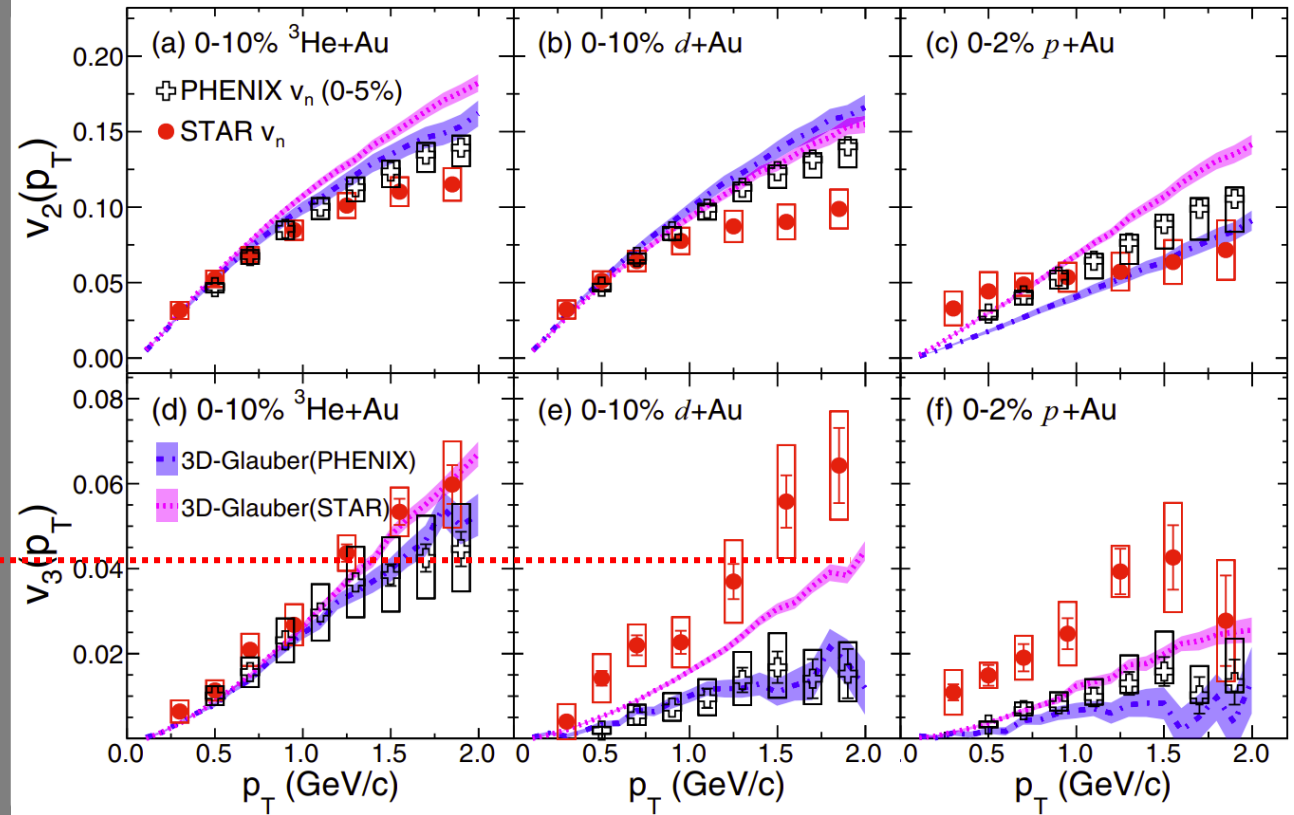
Response to geometry in small systems!

# Geometry: flow in p+Au d+Au He+Au



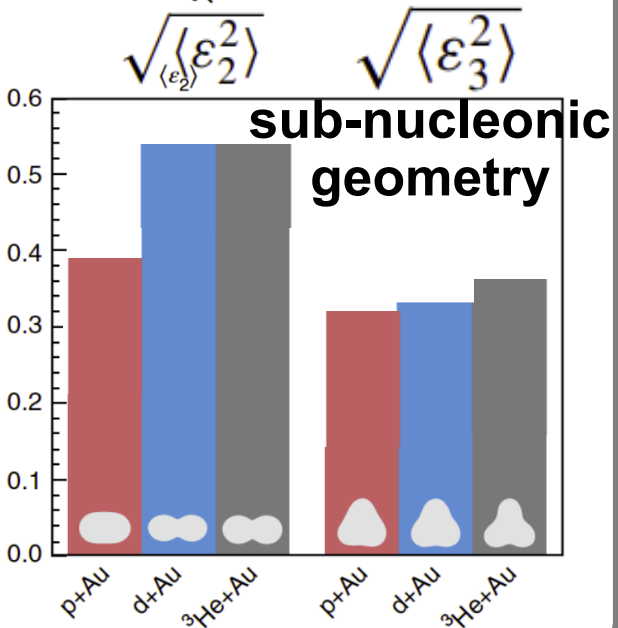
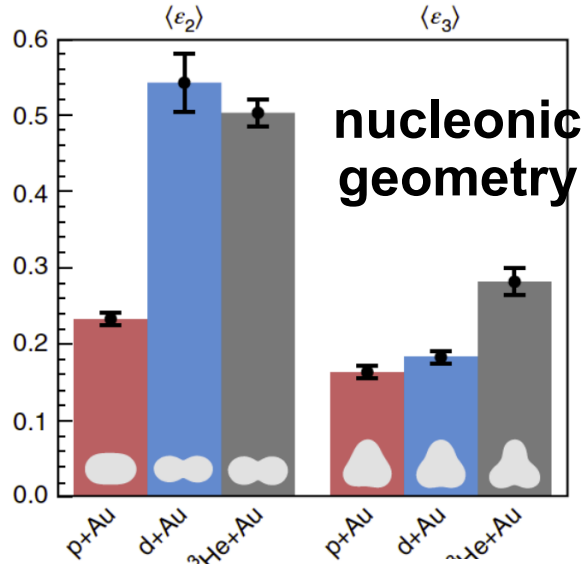
Nature Physics 15, (2019)

Similar  $v_3$  in STAR  
d+Au and  $^3\text{He}+\text{Au}$



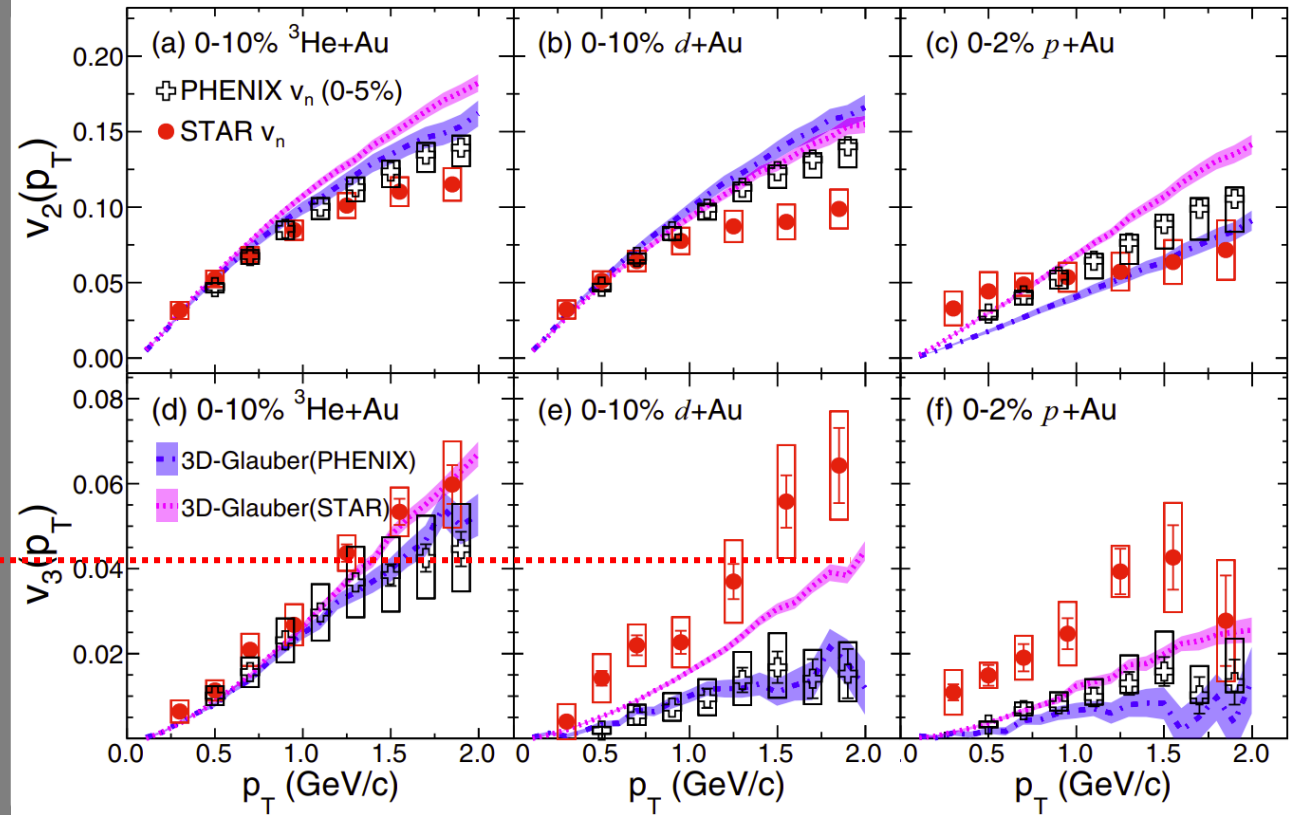
Phys. Rev. C 110

# Geometry: flow in p+Au d+Au He+Au



Nature Physics 15, (2019)

Similar  $v_3$  in STAR  
d+Au and  $^3\text{He}+\text{Au}$



Sub-nucleonic geometry + fluctuations predicts smaller  
difference in geometry between d and  $^3\text{He}$

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Precise sub-nucleonic geometry needed to describe  
small systems especially p+Au



# PHENIX vs. STAR: $p$ +Au $d$ +Au $^3\text{He}$ +Au

STAR

TPC

-1  $|\Delta\eta| > 1$  +1

PHENIX

-3.9 -3.1 -3.0 -1 -0.35 +0.35

Larger  $\eta$  separation  $\rightarrow$  more event plane fluctuation  
PHENIX has a larger  $\Delta\eta$  gap  $\rightarrow$  more fluctuations

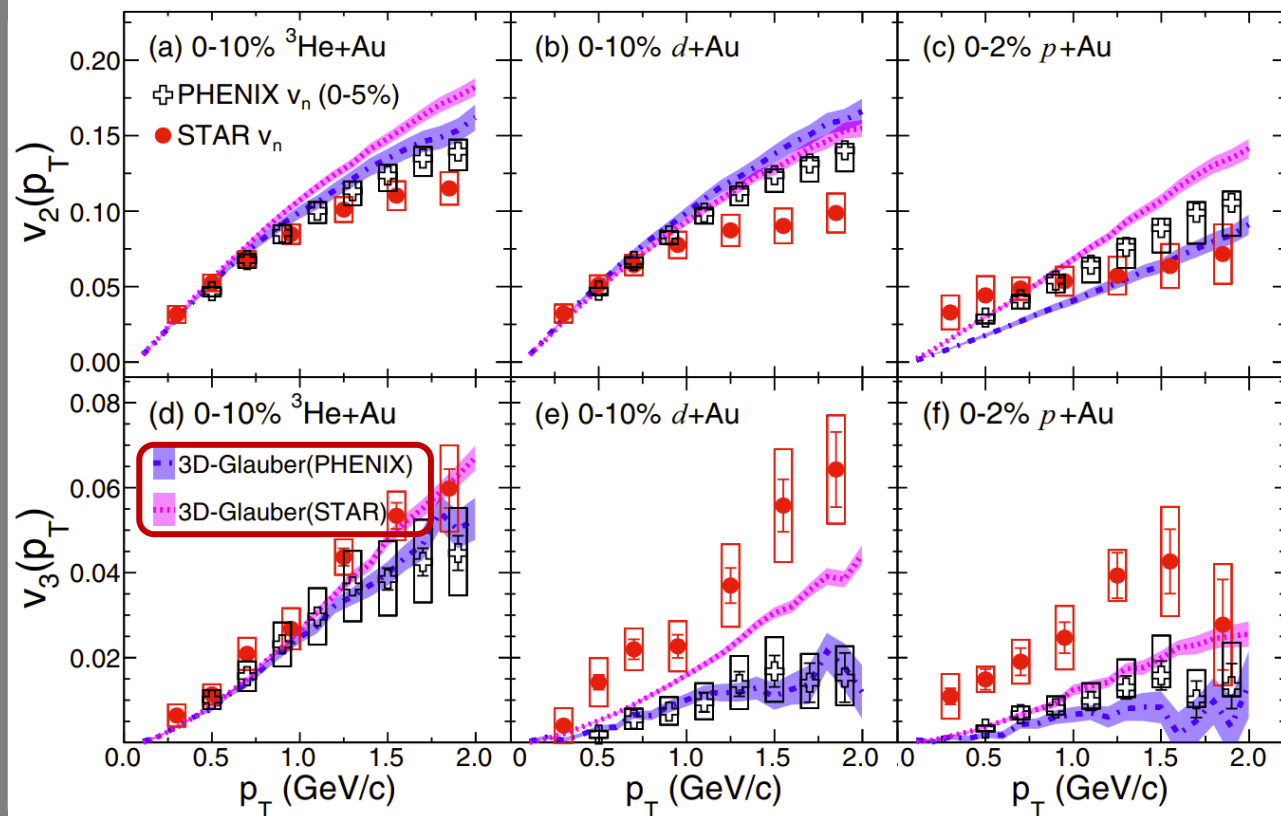
Model: 3D sub-nucleon Glauber+MUSIC+UrQMD

Still some 20-30% discrepancies between data-model

subnucleon fluctuations can also reproduce the measured  $v_3$  values.

In summary, our results highlight the importance of considering subnucleon fluctuations and longitudinal decorrelations in interpreting flow measurements in small collision systems, and they underscore the need for continued refinement of both additional measurements and theoretical models.

ACKNOWLEDGMENTS



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Longitudinal decorrelations is  
larger in smaller systems



Longitudinal decorrelations are important in the push to precision