

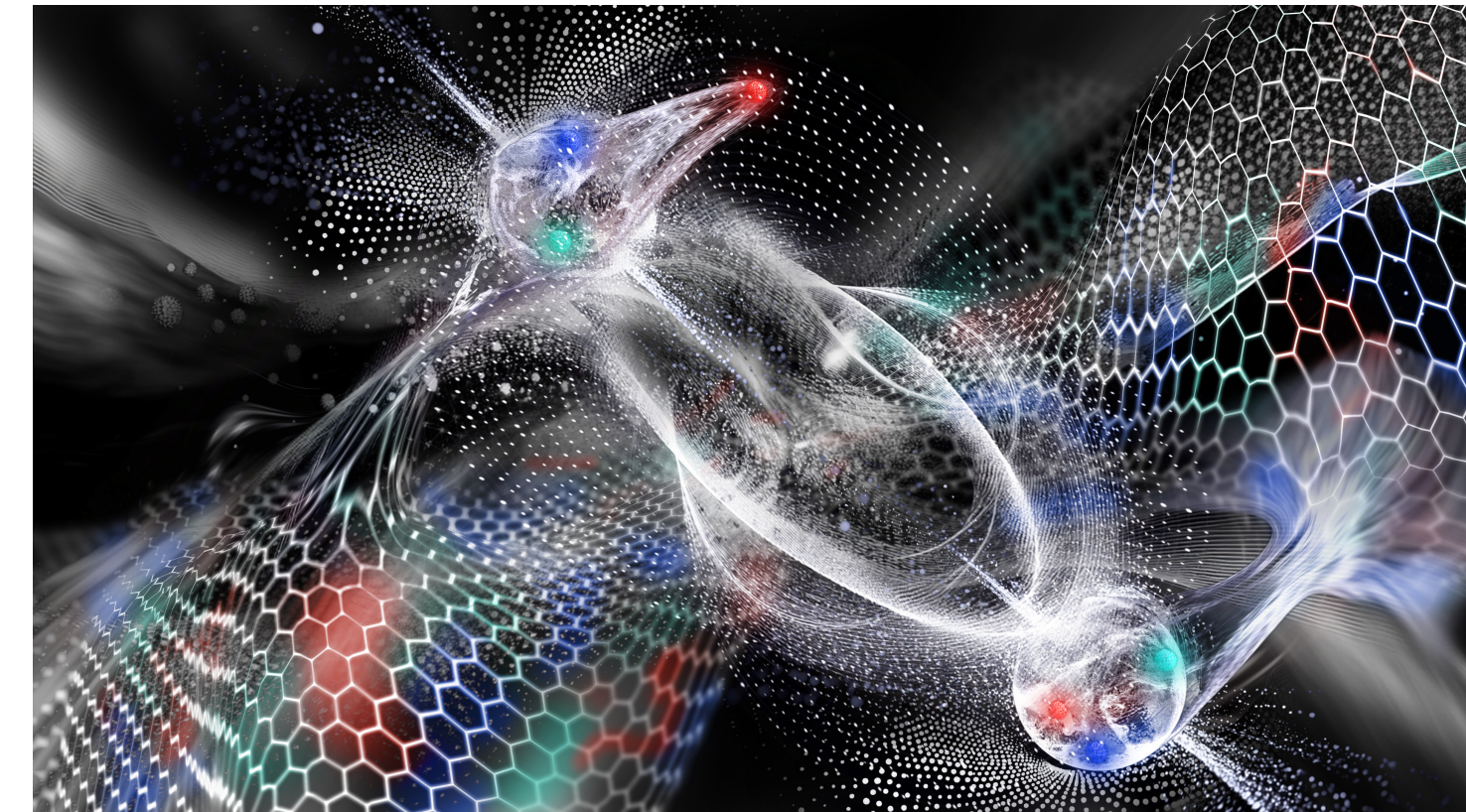
Theory of Quarkonia

Xiaojun Yao

University of Washington



InQubator for Quantum Simulation

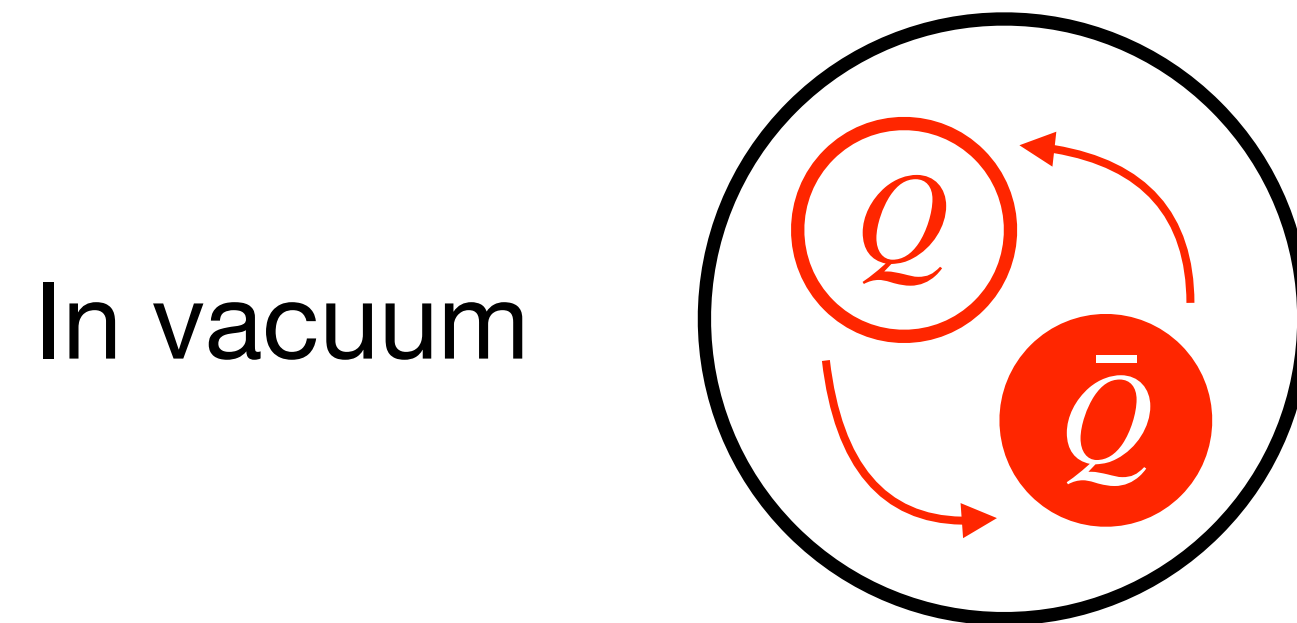


2025 RHIC/AGS Annual Users' Meeting
Brookhaven National Laboratory
May 20, 2025

Quarkonium Probes Deconfinement: Screening

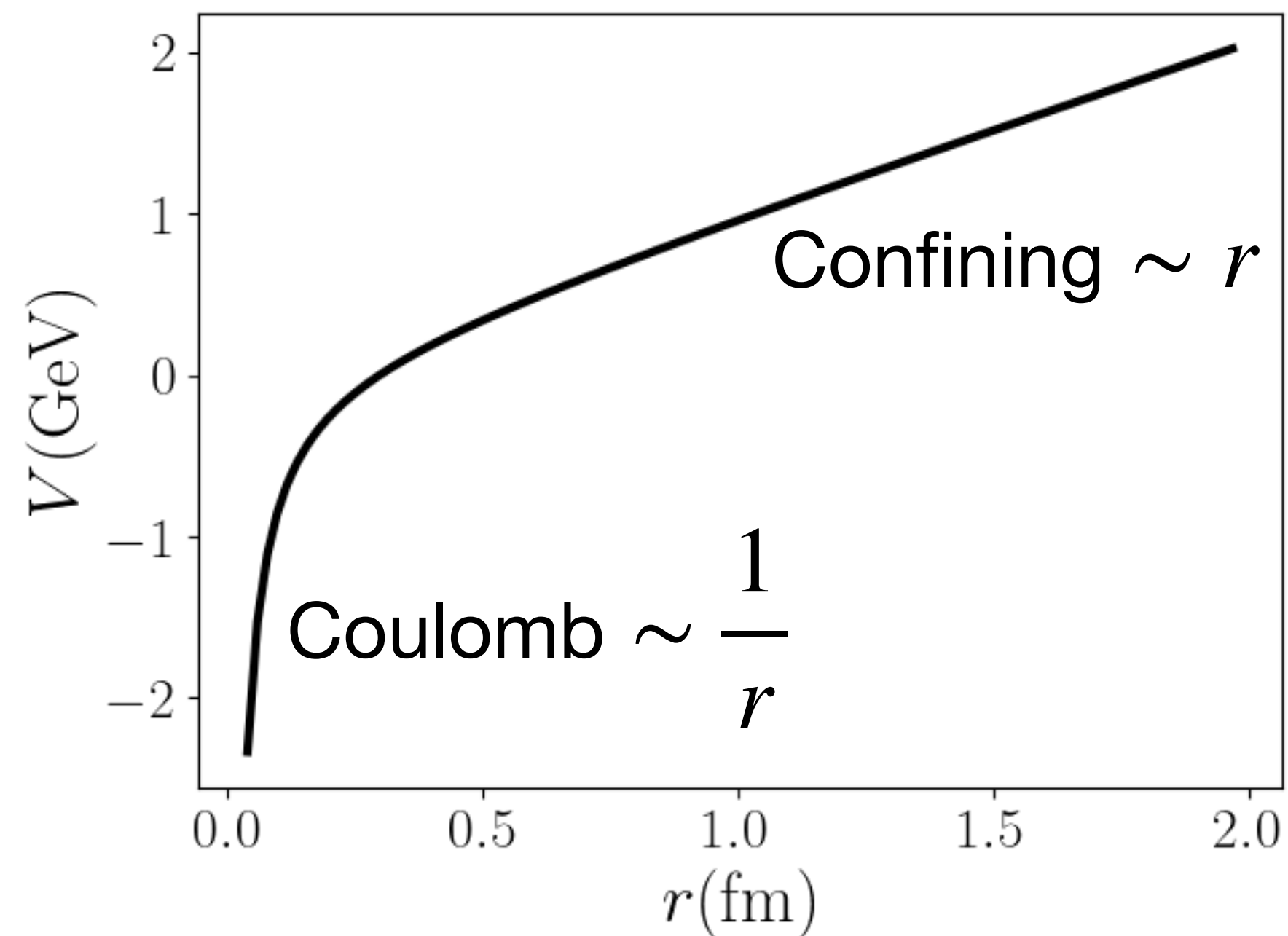
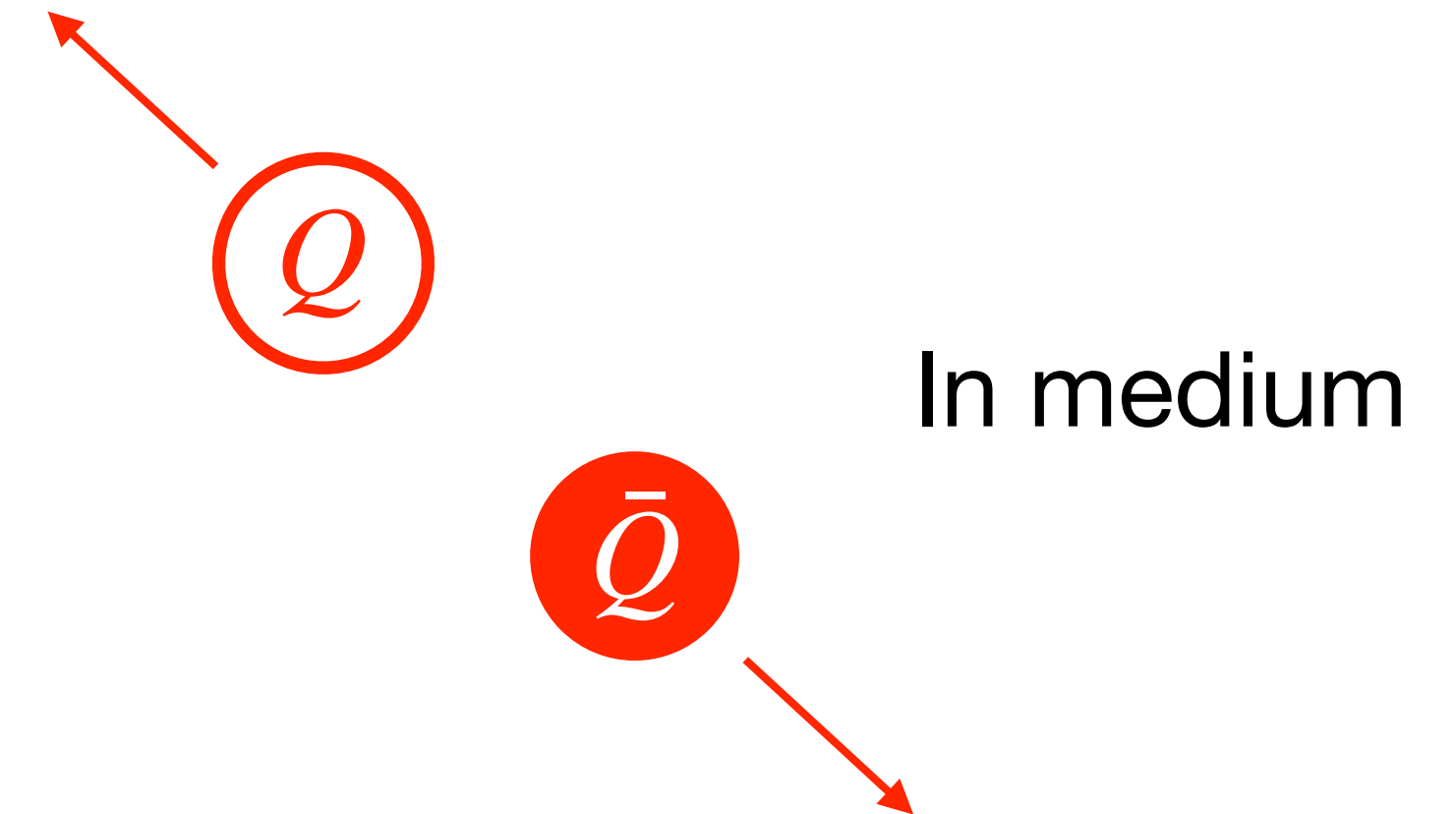
- Quarkonium: bound state of heavy $Q\bar{Q}$ pair

- Screening of potential at finite T

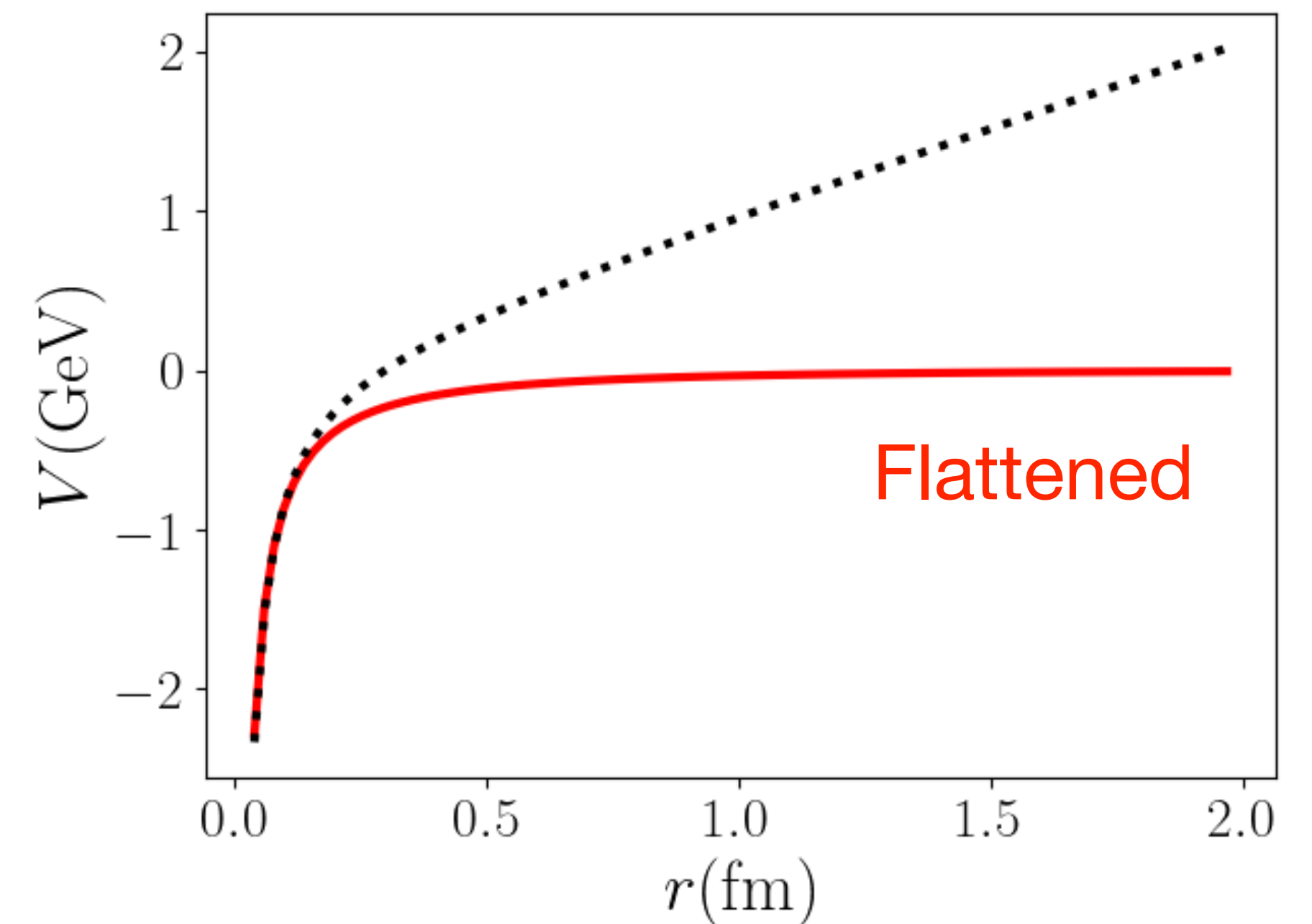


“Melting”

Matsui, Satz 1986



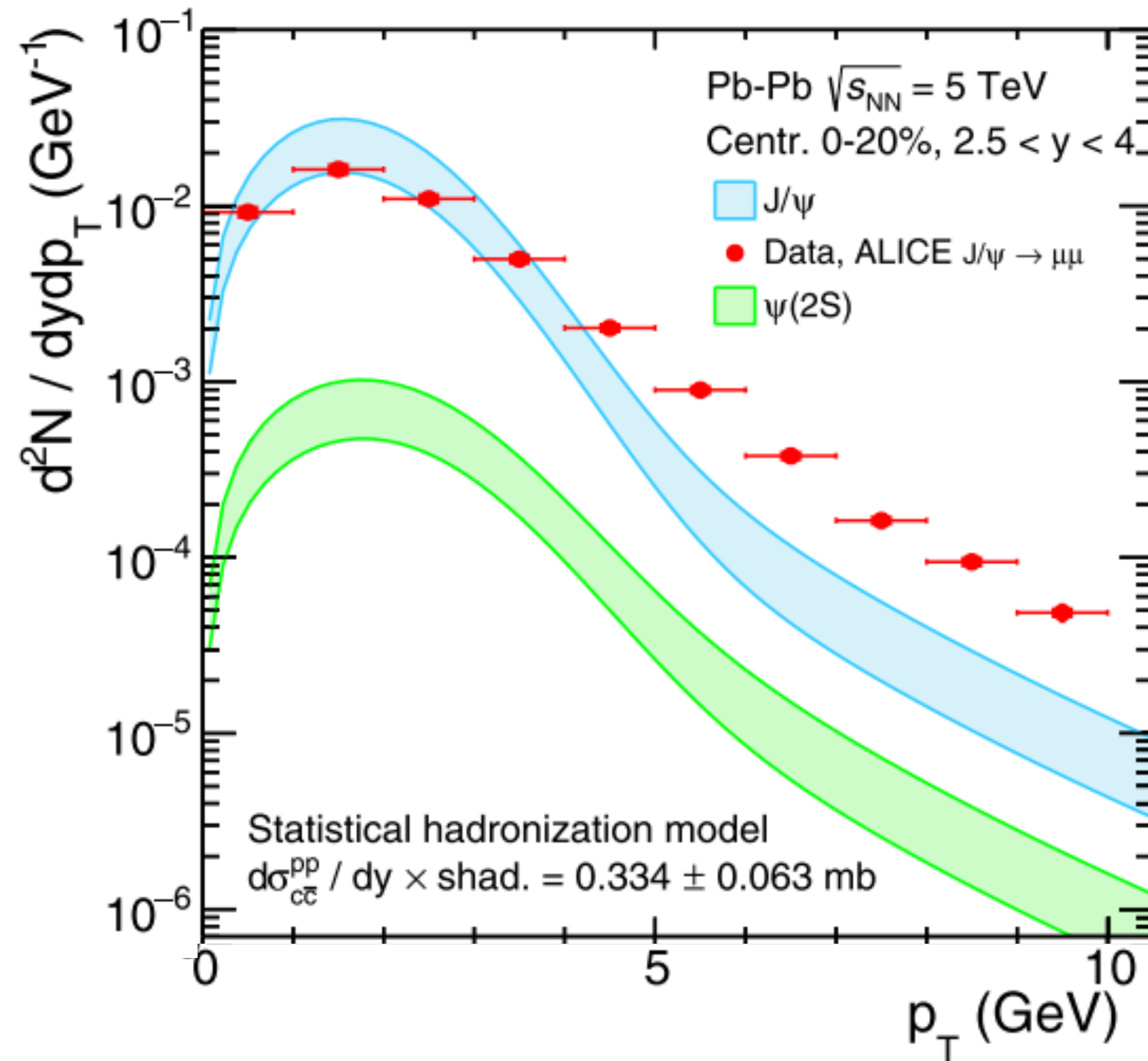
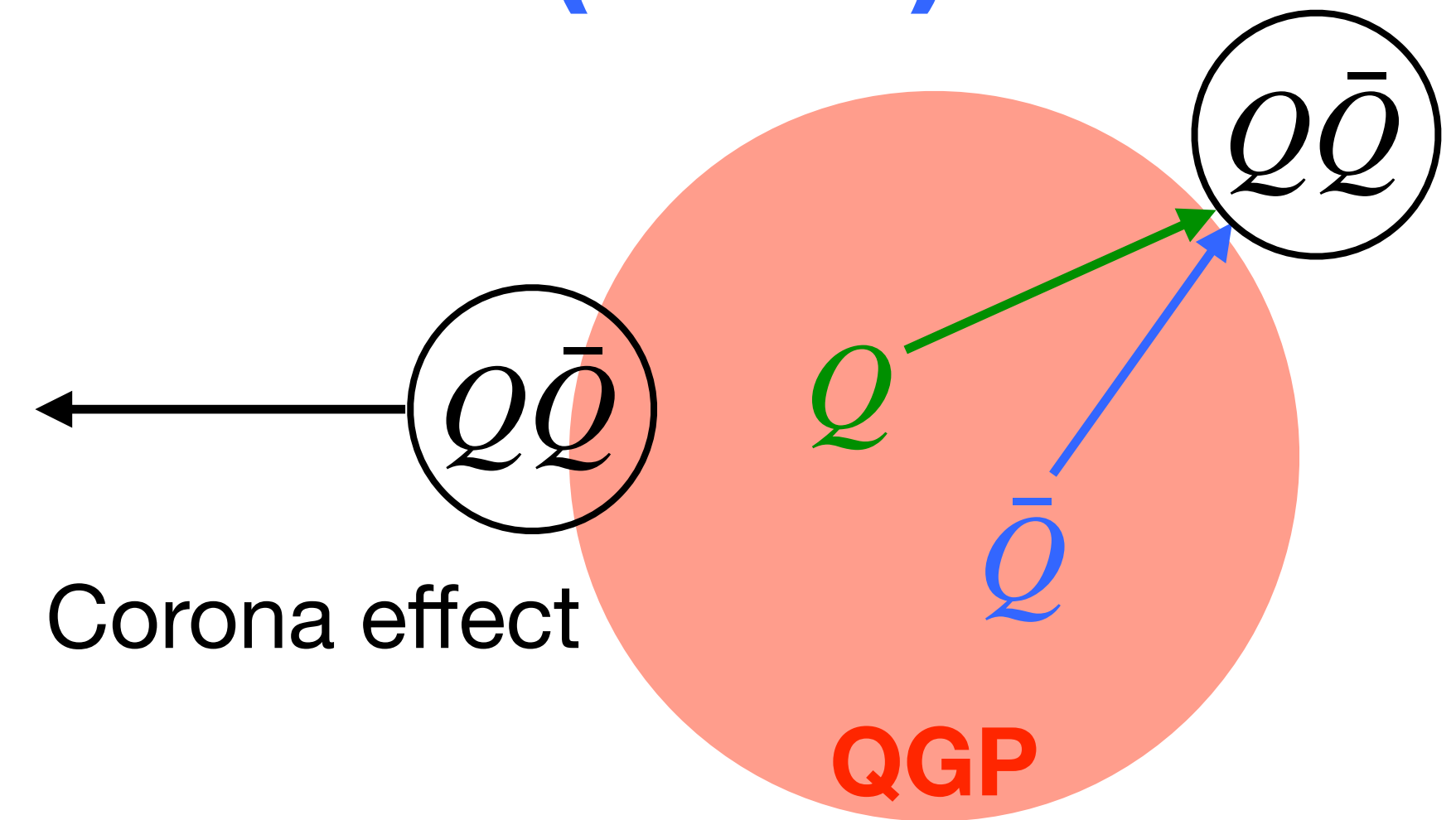
Plasma screening



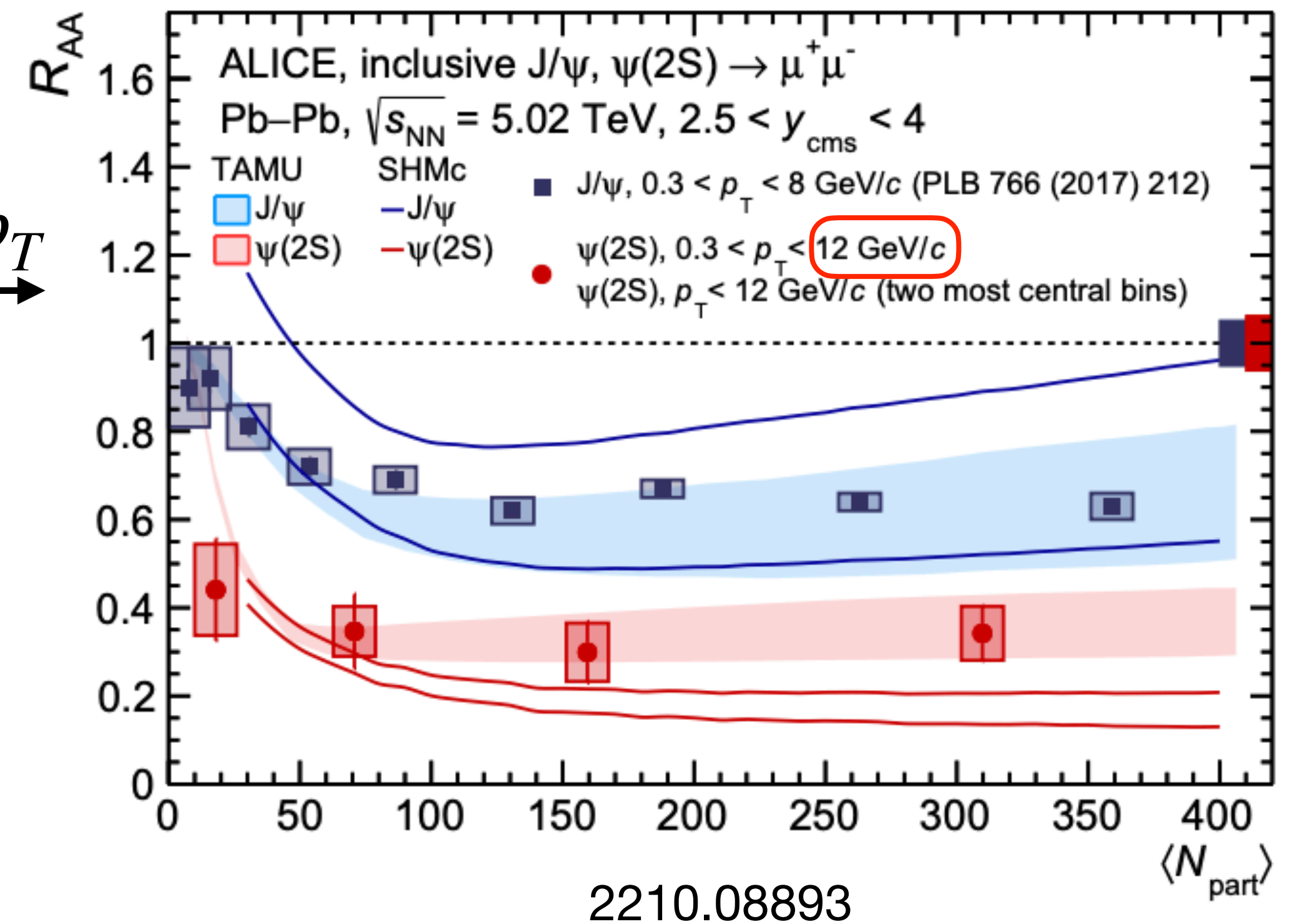
**Assume quarkonium not exist inside QGP,
only (re)generated at phase transition**

Statistical Hadronization Model (SHM)

- In QGP: charms unbound, reach **kinetic equilibrium** (large mass)
- Instantaneous hadronization at freezeout

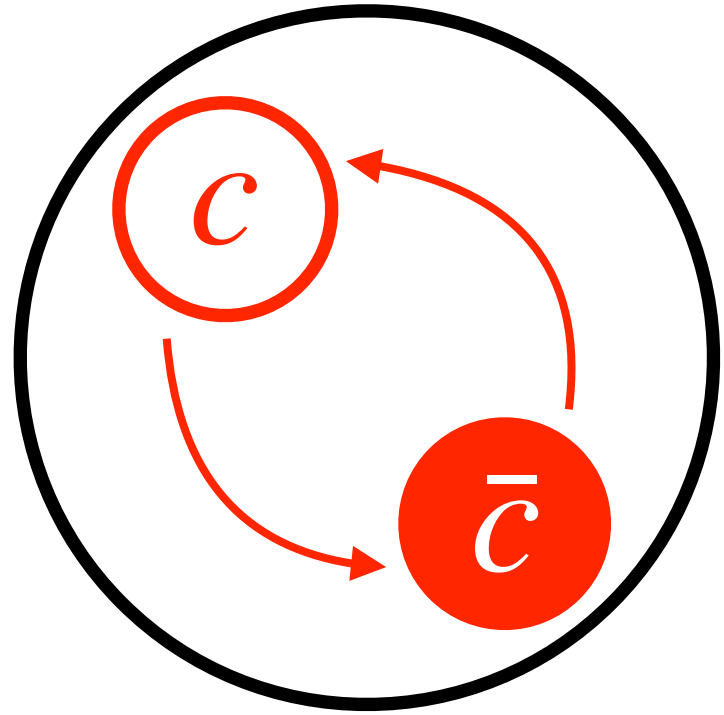


Valid at low p_T

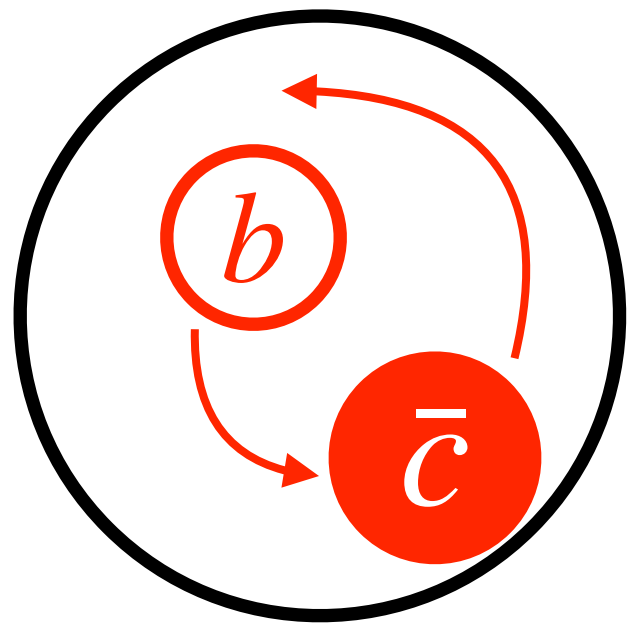


B_c Meson Production

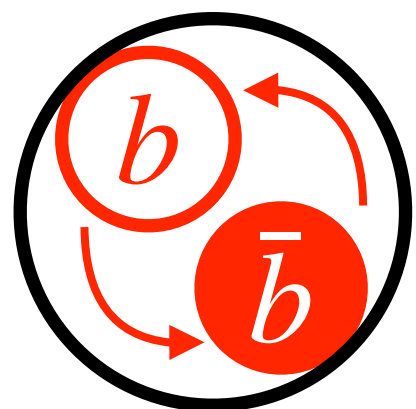
- B_c more “similar” to J/ψ than $\Upsilon(1S)$



$$r[J/\psi] \sim 0.4 \text{ fm}$$

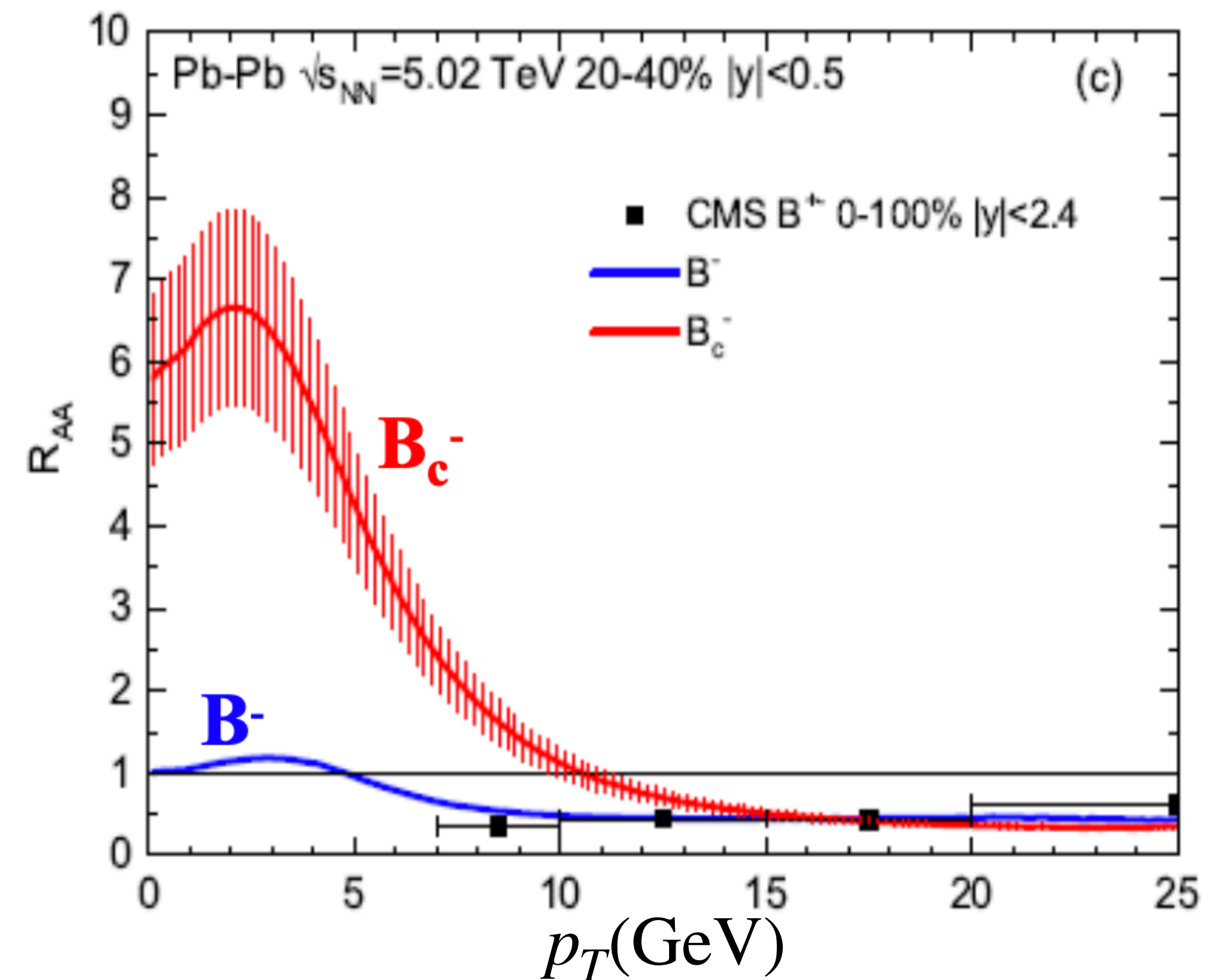


$$r[B_c] \sim 0.34 \text{ fm}$$



$$r[\Upsilon(1S)] \sim 0.2 \text{ fm}$$

- Langevin for unbound heavy quarks
- Instantaneous coalescence at freezeout: one bottom surrounded by multiple charms

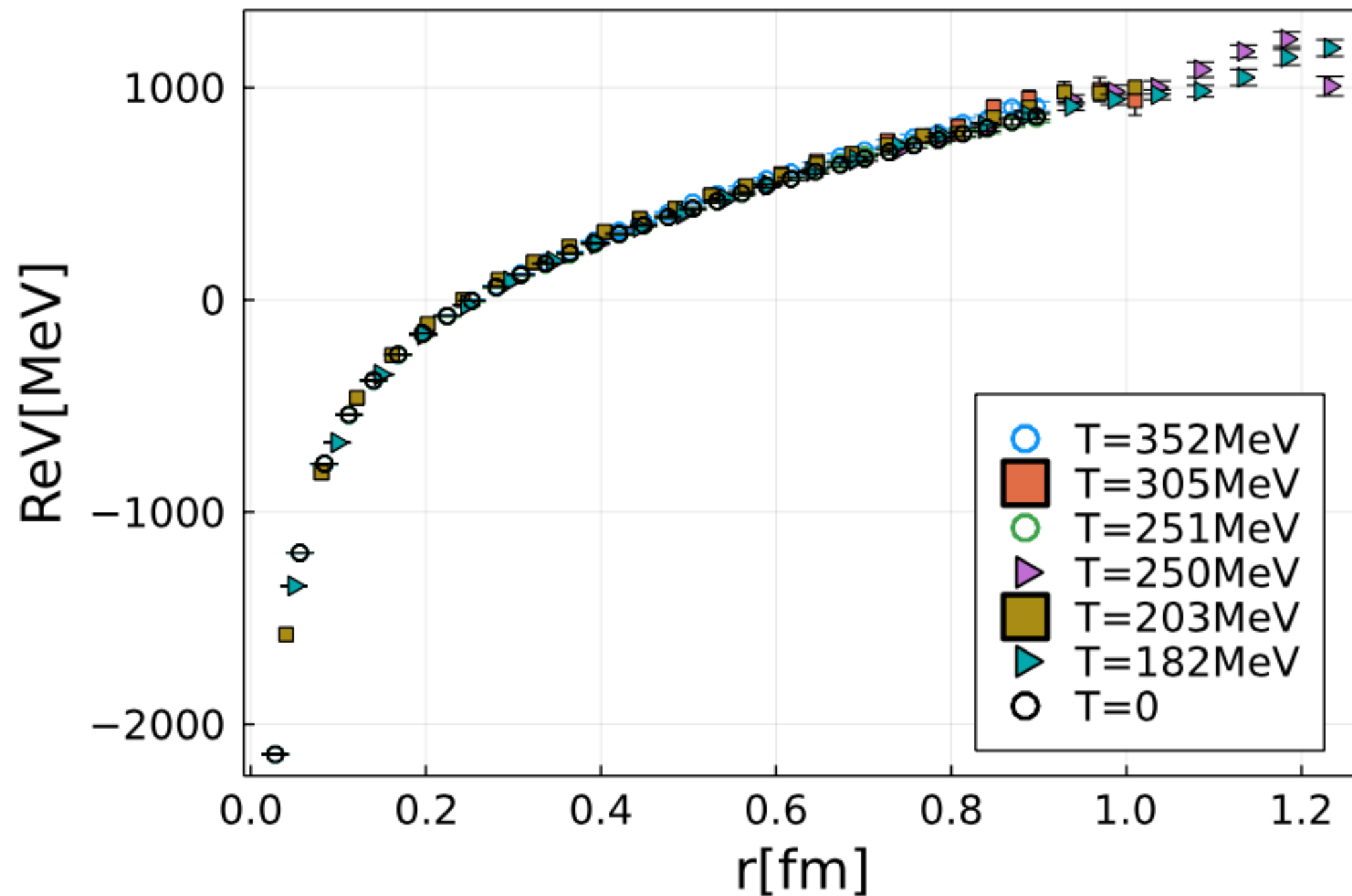


Zhao, He, 2407.05234

Can Quarkonium Exist inside QGP?

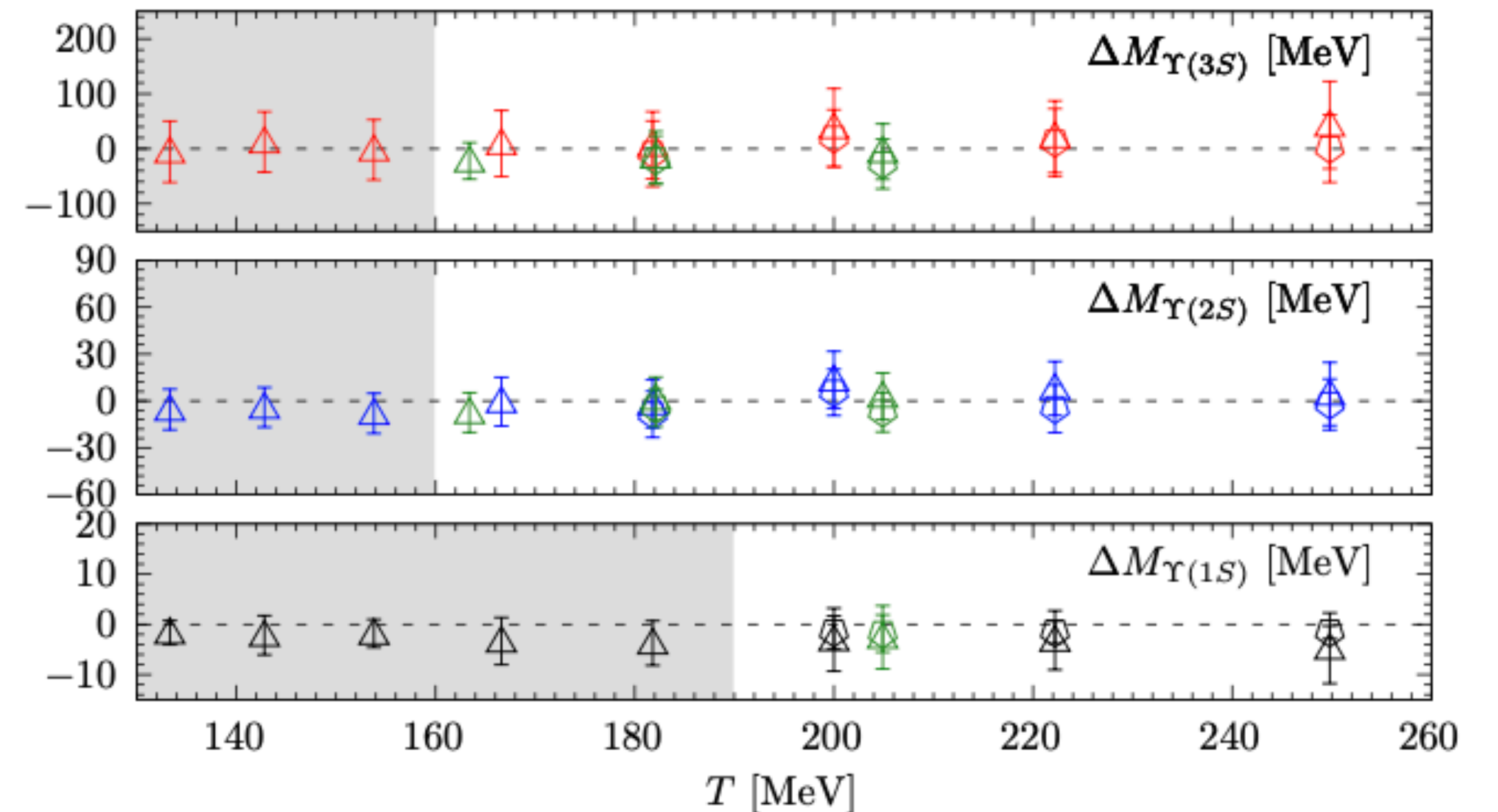
- Deeply bound state at low $T \approx$ small color dipole, e.g. $rT \sim 0.2$ for $\Upsilon(1S)$ at $T = 200$ MeV
- Lattice hint at static limit $M_Q \rightarrow \infty$:
no screening of potential

$$\langle \text{Tr}[W(\tau, r)W^\dagger(\tau, 0)] \rangle_T = \int_\omega \rho(\omega, r) e^{-\beta\omega}$$



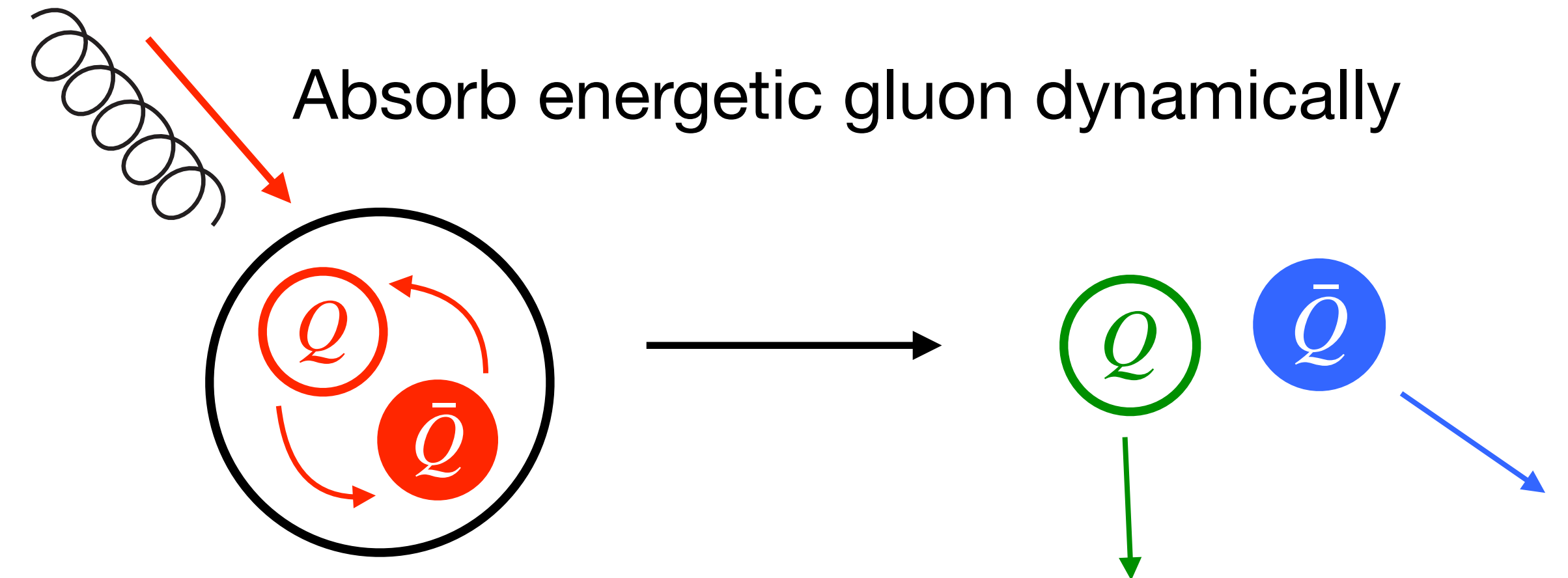
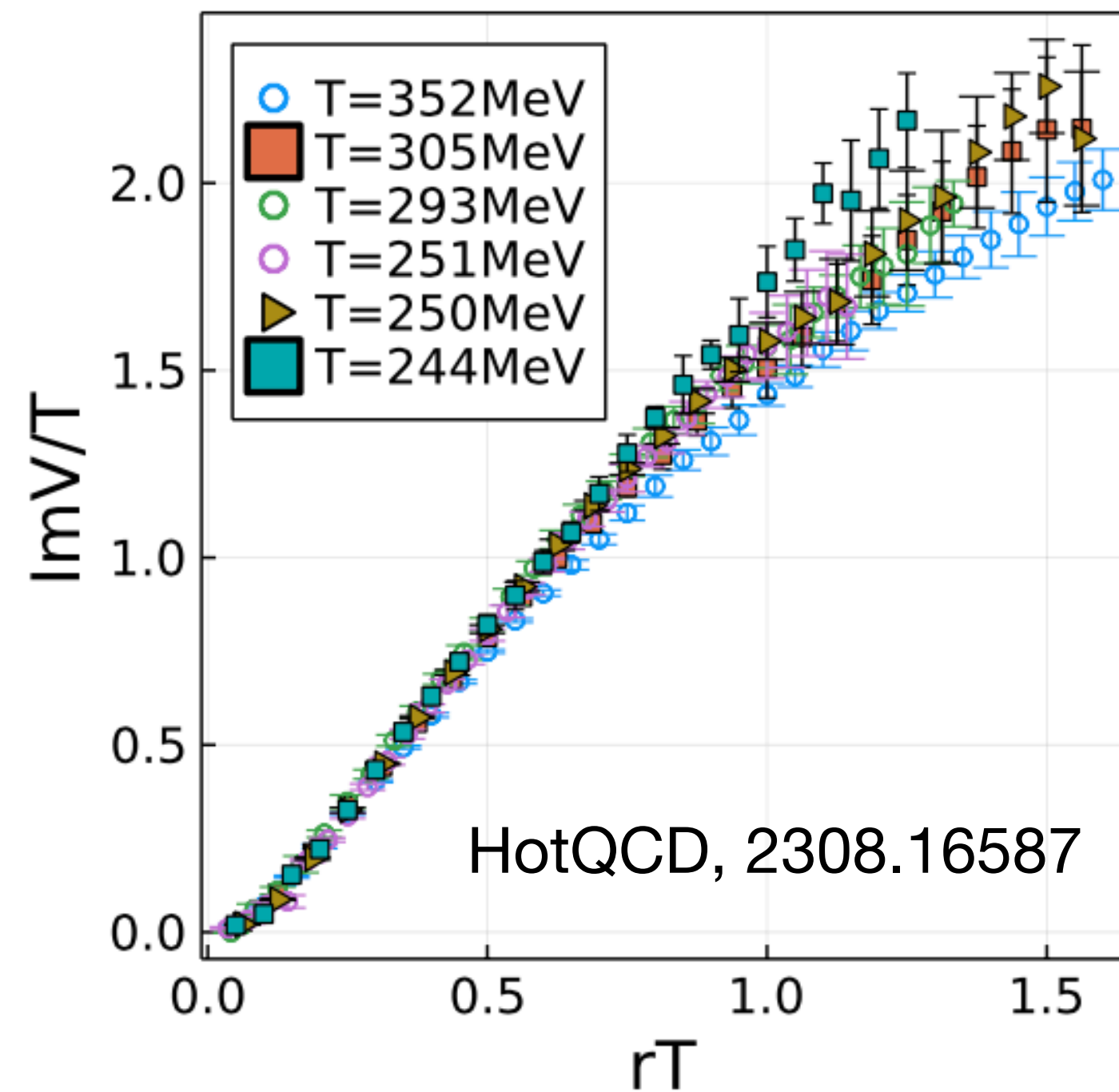
HotQCD, 2308.16587

Lattice NRQCD at finite bottom quark mass



Dissociation and Regeneration

- Same lattice calculation shows thermal width $\longrightarrow$ Quarkonium can dissociate

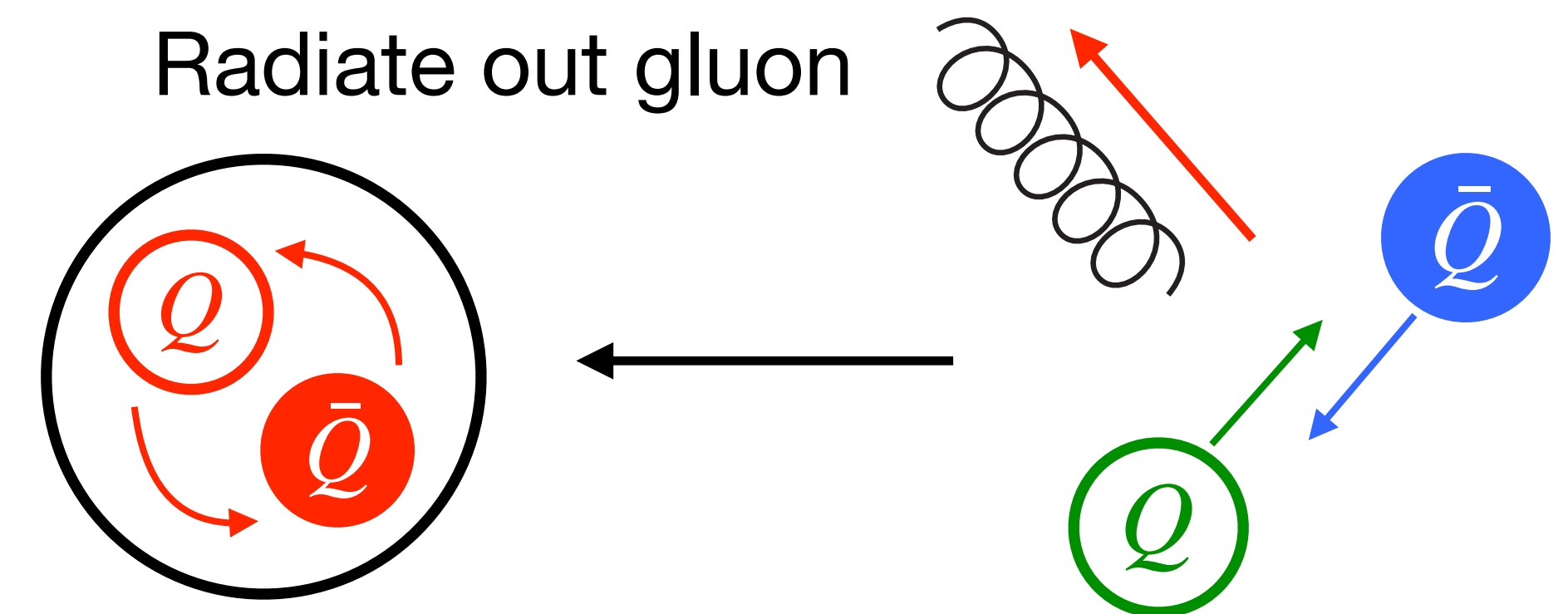


Laine, Philipsen, Romatschke, Tassler, hep-ph/0611300

Beraudo, Blaizot, Ratti, 0712.4394

Brambilla, Ghiglieri, Petreczky, Vairo, 0804.0993

- If quarkonium exists inside QGP, *recombination* can also happen, as inverse process of dissociation



Thews, Schroedter, Rafelski, hep-ph/0007323

Phenomenology with Dissociation and Regeneration

- **Texas A&M University (TAMU) model**

$$\frac{dN_{\mathcal{B}}(\tau)}{d\tau} = -\Gamma_D[T(\tau)] \left(N_{\mathcal{B}}(\tau) - N_{\mathcal{B}}^{\text{eq}}[T(\tau)] \right)$$

#(bound state)
change rate

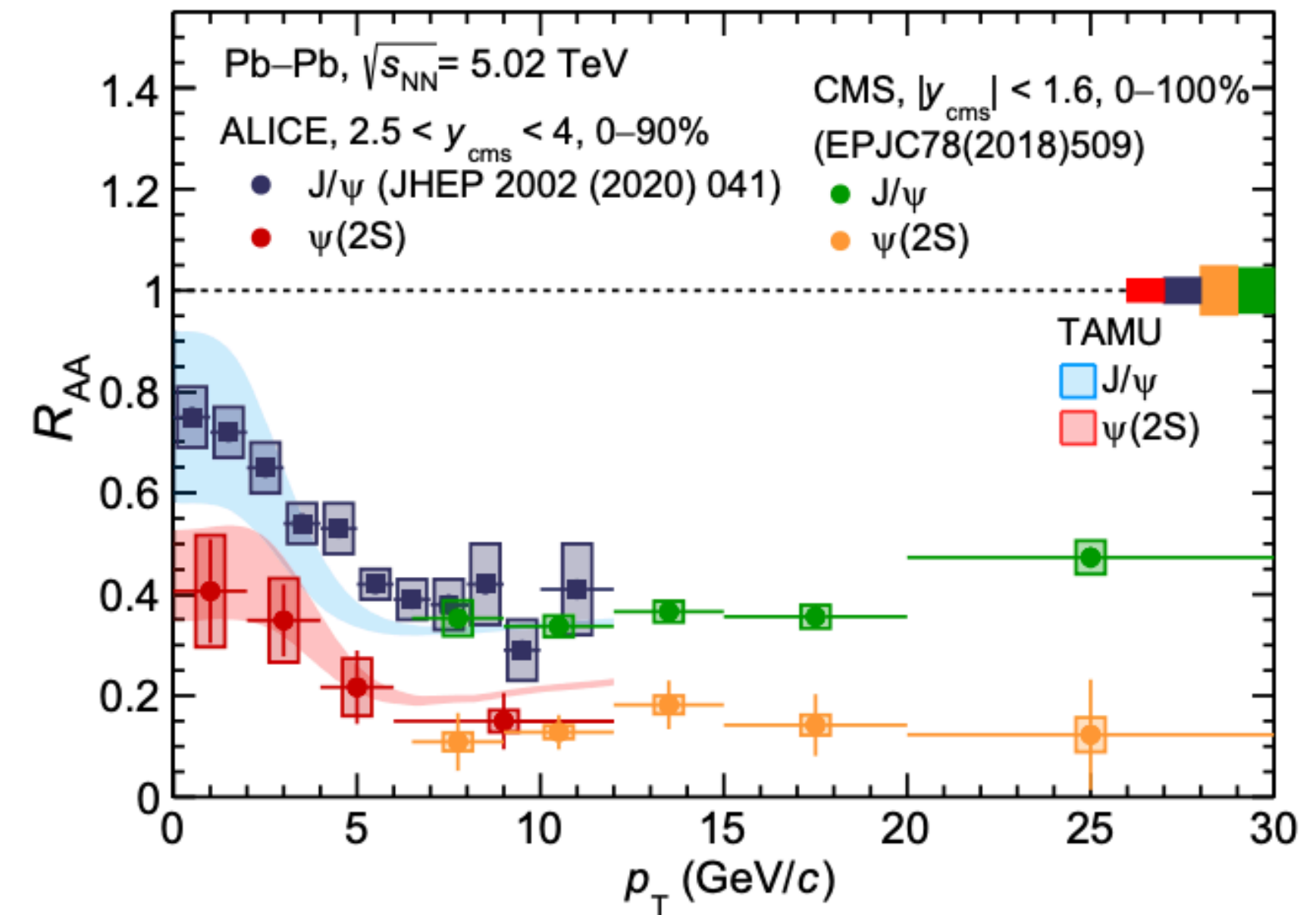
Equilibrium number,
time-dependent

- Dissociation rate from $g + \mathcal{B} \rightarrow Q + \bar{Q}$
 $2 \times (q/g + Q \rightarrow q/g + Q)$

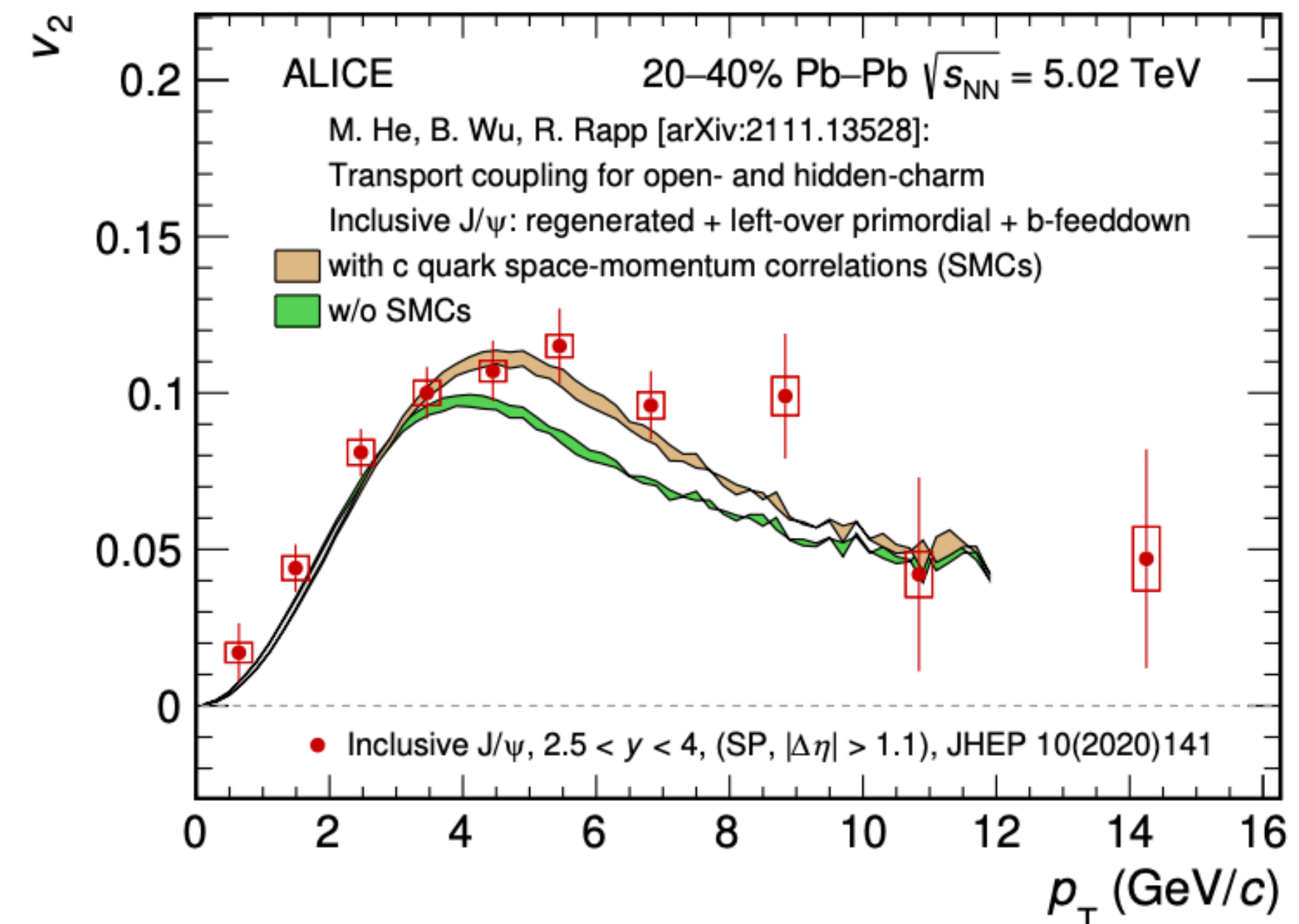
- Phenomenological factors, interference

- New development:
T-matrix summation + lattice input

Wu, Tang, Rapp, 2503.10089



2210.08893



2005.14518

Phenomenology with Dissociation and Regeneration

- Santiago comover interaction model**

$$\tau \frac{d\rho_{\mathcal{B}}(\tau)}{d\tau} = -\sigma_{\text{co}-\mathcal{B}} \left(\rho_{\text{co}} \rho_{\mathcal{B}} - \rho_Q \rho_{\bar{Q}} \right) \quad \left[\cosh(y - \eta) \frac{\partial}{\partial \tau} + \frac{\sinh(y - \eta)}{\tau} \frac{\partial}{\partial \eta} + \mathbf{v}_T \cdot \nabla_T \right] f_{\mathcal{B}} = -\alpha f_{\mathcal{B}} + \beta$$

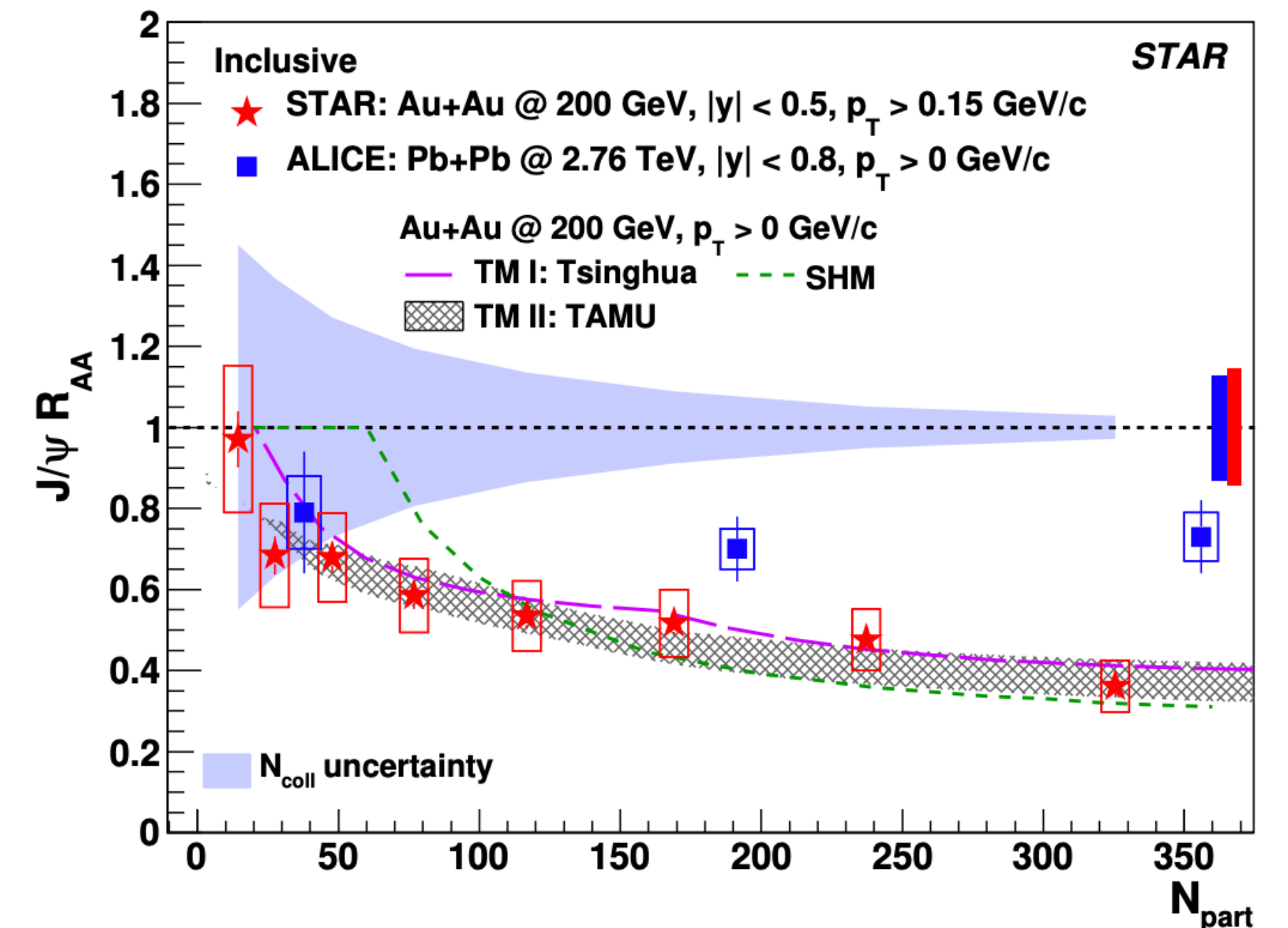
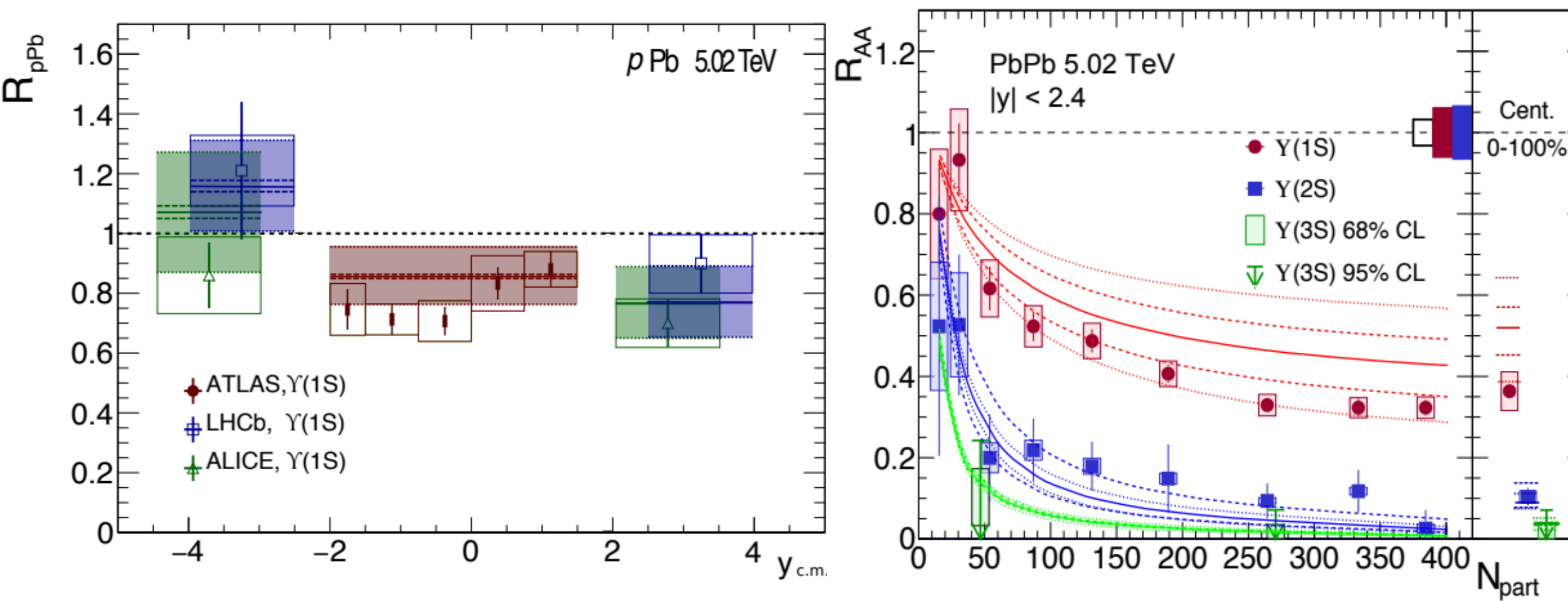
Input: comover-quarkonium cross section

No regeneration for bottomonium

- Tsinghua model: relativistic Boltzmann**

Only leading-order processes included

$$g + \mathcal{B} \leftrightarrow Q + \bar{Q}$$



What Properties of QGP are Probed?

- **Temperature?** $\longrightarrow$ Input from well-studied hydro
- **Dissociation/recombination rate?** $\longrightarrow$ Property of quarkonium

**Separate intrinsic properties of quarkonium
from intrinsic properties of QGP**

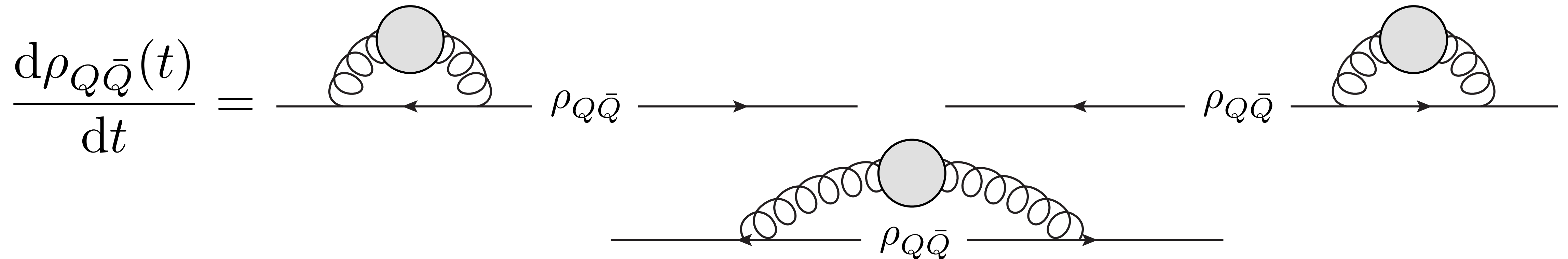


Open quantum system (OQS) + effective field theory (EFT)

Open Quantum System for Quarkonium

- **QGP + $Q\bar{Q}$ = closed quantum system; $Q\bar{Q}$ = open quantum system**

$$\rho_{Q\bar{Q}}(t) = \text{Tr}_{\text{QGP}}[U(t)\rho_{\text{tot}}(0)U^\dagger(t)]$$



Akamatsu, Rothkopf, 1110.1203; Akamatsu, 1403.5783

Brambilla, Escobedo, Soto, Vairo, 1612.07248, 1711.04515; Blaizot, Escobedo, 1711.10812;
Miura, Akamatsu, Asakawa, Rothkopf, 1908.06293, Yao, Mehen, 1811.07027, 2009.02408

Reviews of OQS for quarkonia: Rothkopf, 1912.02253; Akamatsu, 2009.10559, Sharma, 2101.04268; Yao, 2102.01736

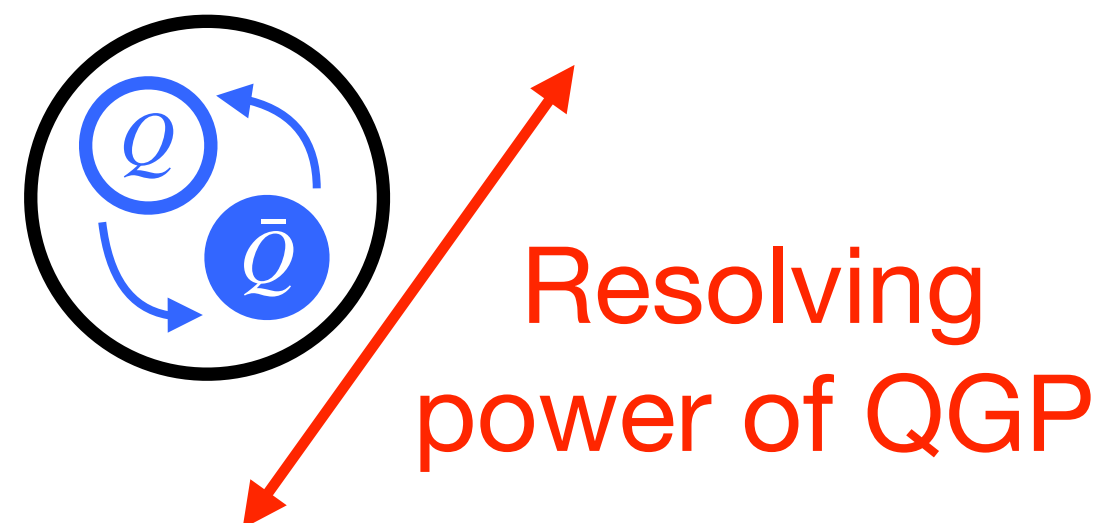
Separation of Energy Scales and Nonrelativistic EFT

In vacuum $M \gg M_v \gg M_v^2$

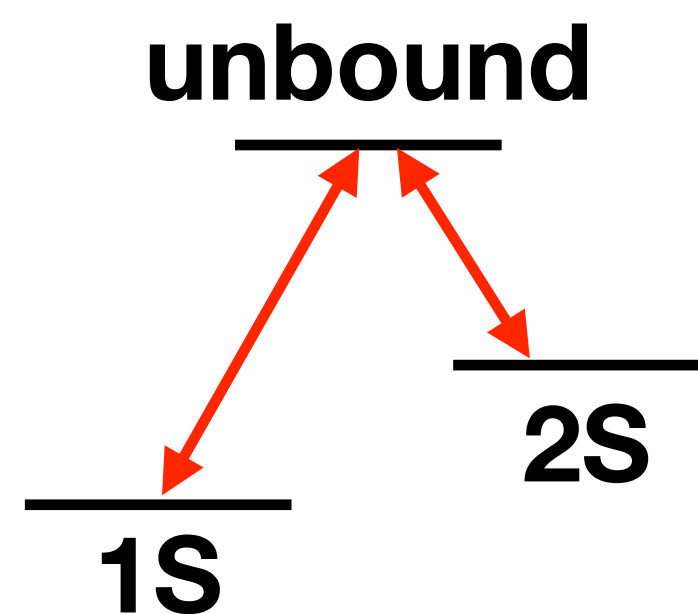
QCD	M	heavy quark mass	$c\bar{c}$	$b\bar{b}$
			1.3	4.5 GeV
NRQCD	M_v	inverse of quarkonium size r^{-1}	0.7	1.1 GeV
pNRQCD	M_v^2	quarkonium binding energy	0.5	0.5 GeV

Medium case depends on where T fits

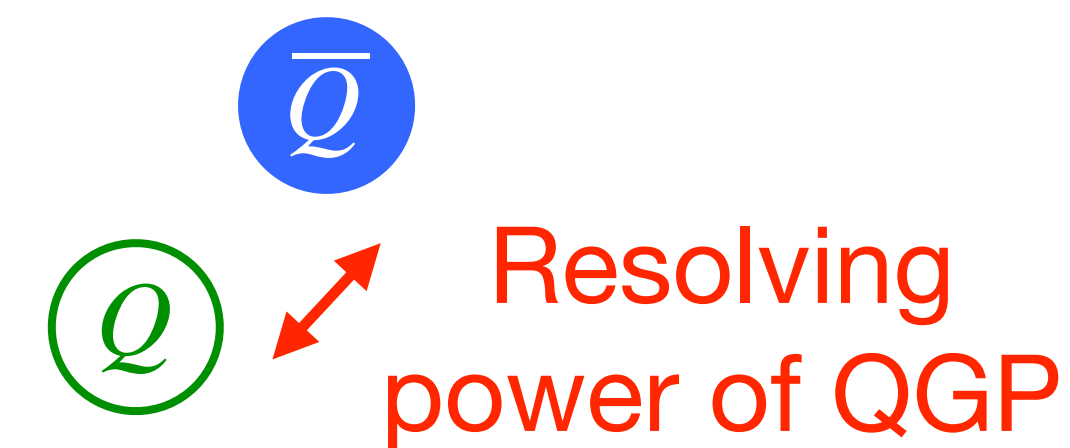
- Quantum optical limit: low T



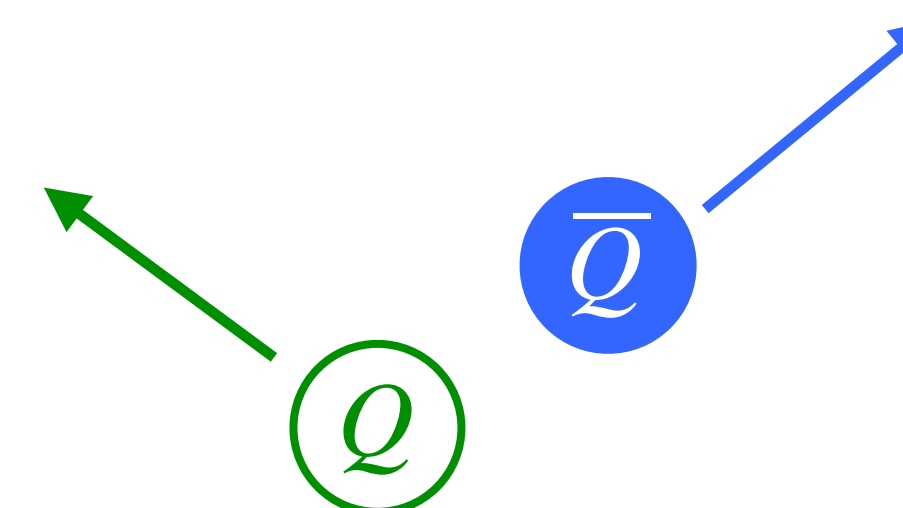
Transitions between levels



- Quantum Brownian motion: high T



Decoherence of $Q\bar{Q}$ pair, “diffusion”



Separation of Energy Scales and Physical Scenarios

Hierarchy of energy scales	EFT	Quantum Description	Classical Limit
$M \gg T \gg Mv^2, \Lambda_{\text{QCD}}$	NRQCD $\alpha_s(T)$ small	Lindblad (quantum Brownian motion)	Diffusion equation (semiclassical limit)
$M \gg Mv \gg T, \Lambda_{\text{QCD}}$ $T \gg Mv^2$, expand Mv^2/T	pNRQCD rT small	Lindblad (quantum Brownian motion)	
$M \gg Mv \gg T, \Lambda_{\text{QCD}}$ No expansion of Mv^2/T	pNRQCD rT small		Boltzmann equation (quantum optical and semiclassical limits)

Case 1: Lindblad Equation in Quantum Brownian Motion

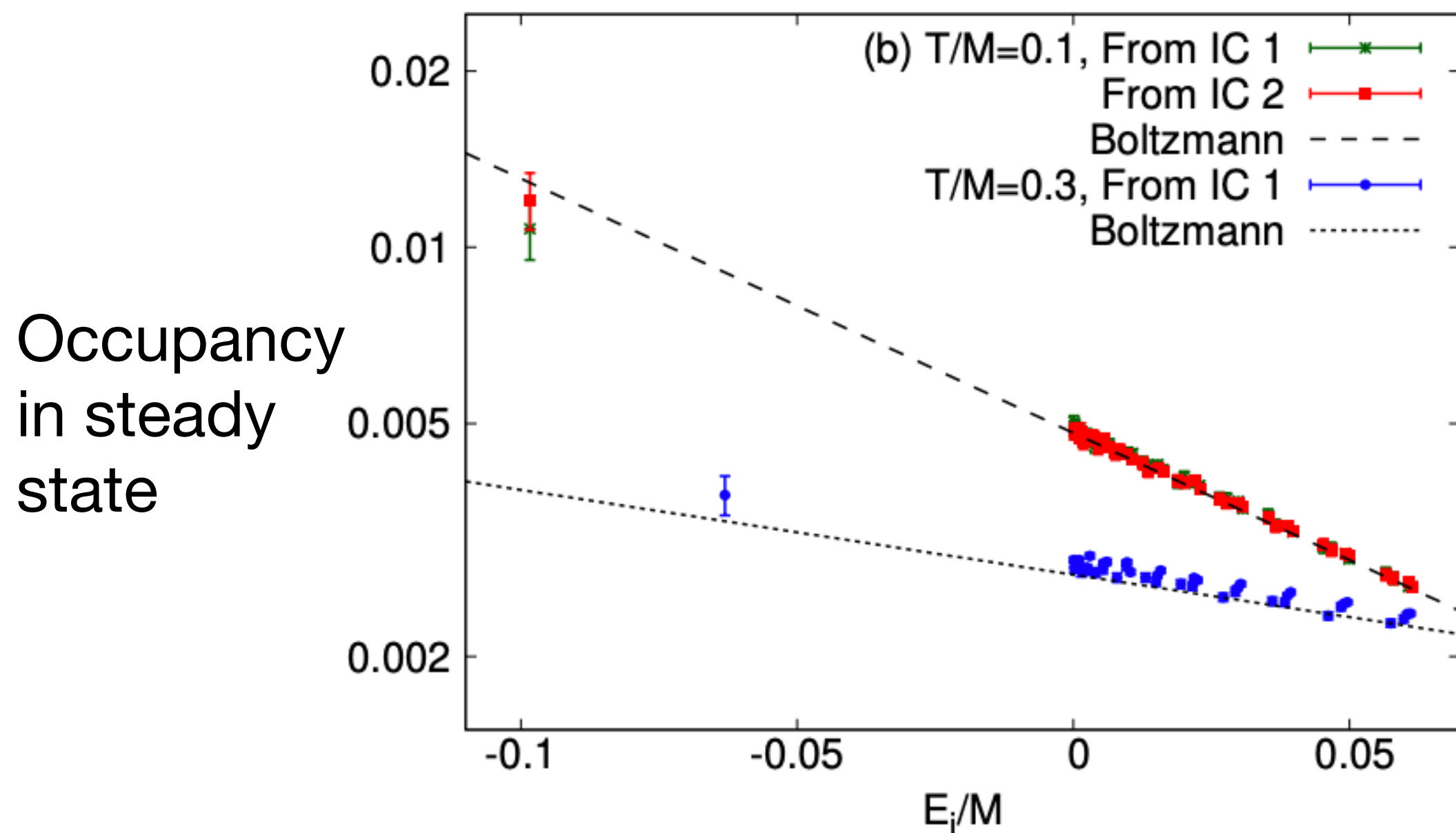
- $M \gg T \gg Mv^2, \Lambda_{\text{QCD}}, 1\text{D numerics}$

$$\frac{d\rho_{Q\bar{Q}}(t)}{dt} = -i[H, \rho_{Q\bar{Q}}(t)] + \sum_{ij} D_{ij} \left(L_i \rho_{Q\bar{Q}}(t) L_j^\dagger - \frac{1}{2} \{ L_j^\dagger L_i, \rho_{Q\bar{Q}}(t) \} \right)$$

$H = \text{kinetic} + \text{potential}$

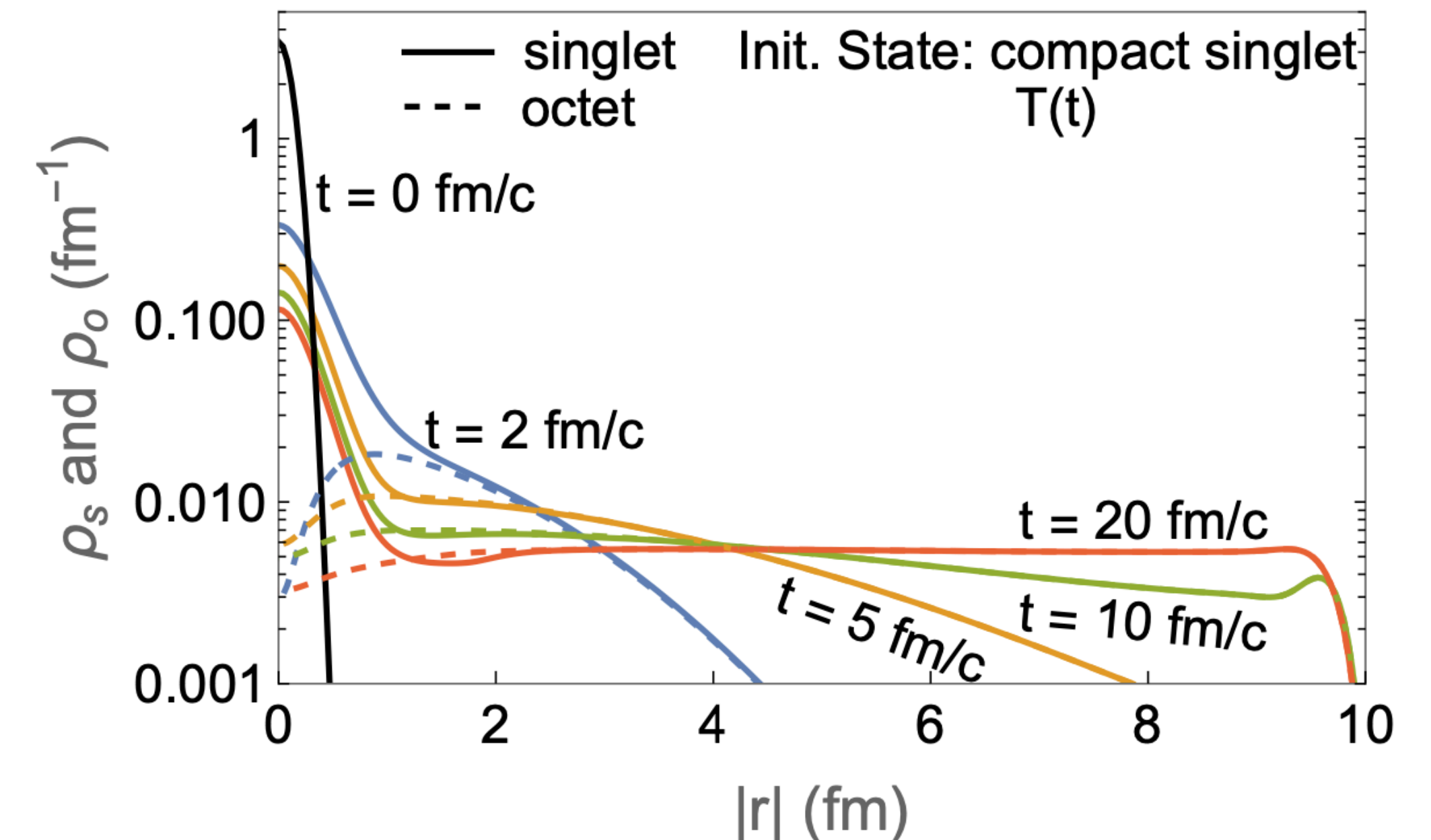
$L_i : \text{singlet-octet, octet-octet transitions}$

- **Osaka group: momentum basis $D_{ij} \propto \delta_{ij}$**



Miura, Akamatsu, Asakawa, Kaida, 2205.15551

- **Nantes: position basis D_{ij} has off-diagonal**



Delorme, Katz, Gousset, Gossiaux, Blaizot, 2402.04488

Case 2: Lindblad Equation in Quantum Brownian Motion

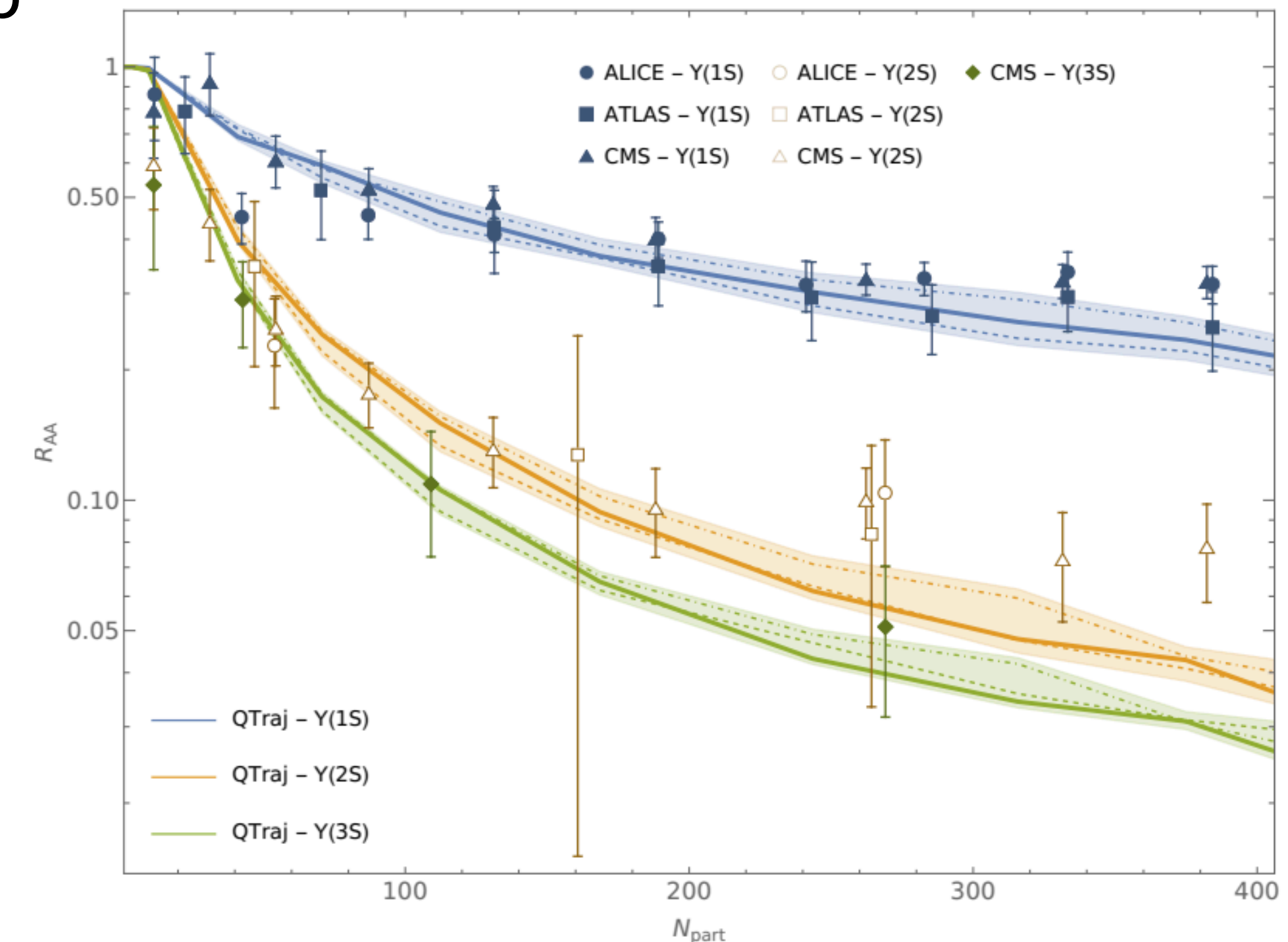
- **Munich-Kent approach:** $M \gg Mv \gg T, \Lambda_{\text{QCD}}$ and $T \gg Mv^2$

$$\frac{d\rho_{Q\bar{Q}}(t)}{dt} = -i[H + \gamma_{\text{adj}}\Delta h, \rho_{Q\bar{Q}}(t)] + \kappa_{\text{adj}}(L_{\alpha i}\rho_{Q\bar{Q}}(t)L_{\alpha i}^\dagger - \frac{1}{2}\{L_{\alpha i}^\dagger L_{\alpha i}, \rho_{Q\bar{Q}}(t)\})$$

- Perturbative at Mv : **3-loop QCD potential**, beyond Coulomb
Nonperturbative transport parameters: κ_{adj} and γ_{adj} from
chromoelectric correlator

$$\kappa_{\text{adj}} + i\gamma_{\text{adj}} \equiv \frac{g^2 T_F}{3N_c} \int dt \langle \mathcal{T} E_i^a(t) \mathcal{W}^{ab}(t, 0) E_i^b(0) \rangle_T$$

- New development: perturbative beyond dipole limit



Brambilla, Magorsch, Strickland,
Vairo, Griend, 2403.15545

Case 3: Boltzmann Equation in Quantum Optical & Classical

- Rigorous derivation** $M \gg Mv \gg T, \Lambda_{\text{QCD}}$
- Duke-MIT: coupled Boltzmann equations for open heavy quarks and quarkonia**

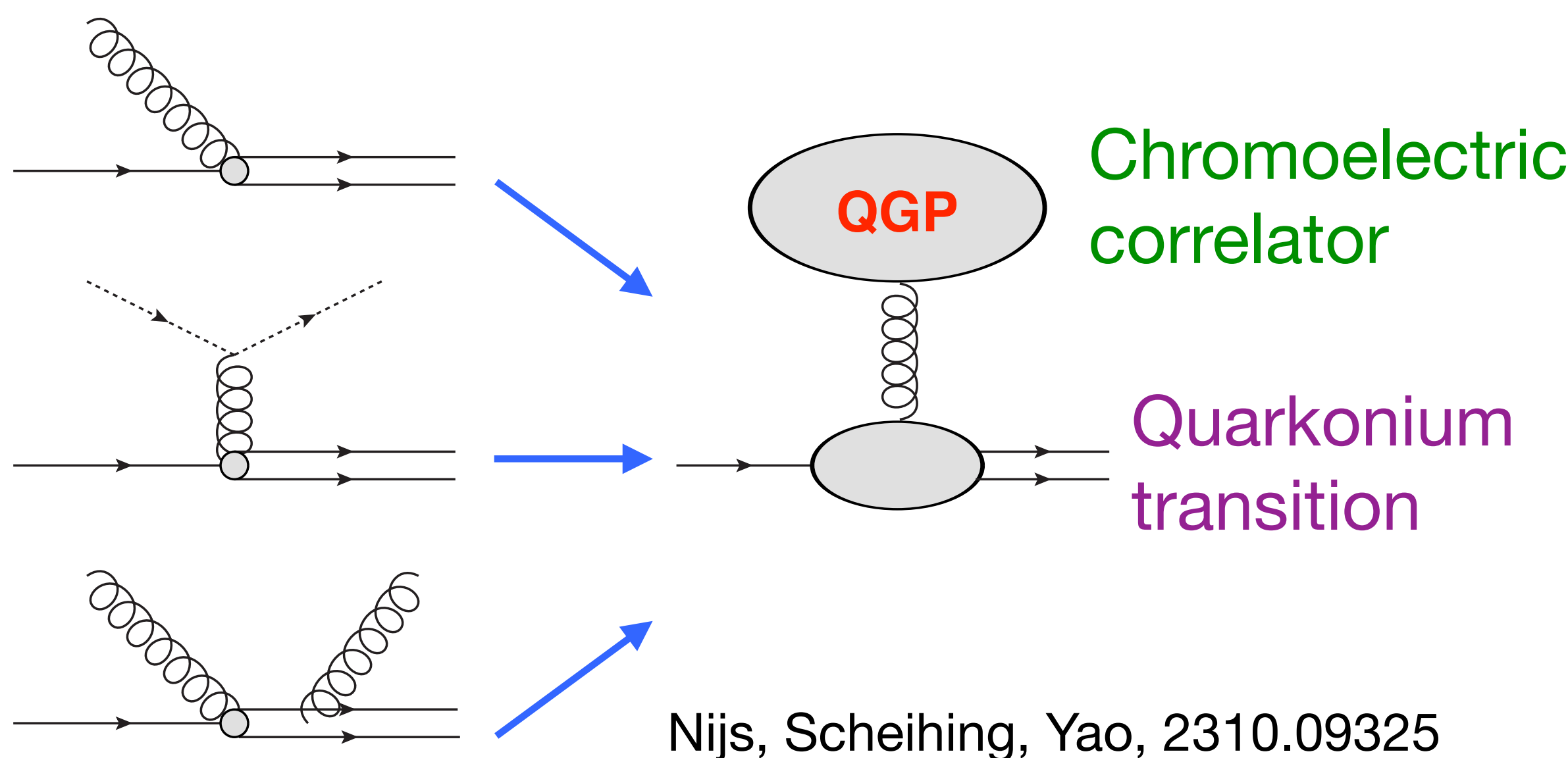
$$\frac{dn_{\mathcal{B}}(t, \mathbf{x})}{dt} = -\Gamma n_{\mathcal{B}}(t, \mathbf{x}) + F(t, \mathbf{x})$$

Yao, Mehen,
 1811.07027,
 2009.02408

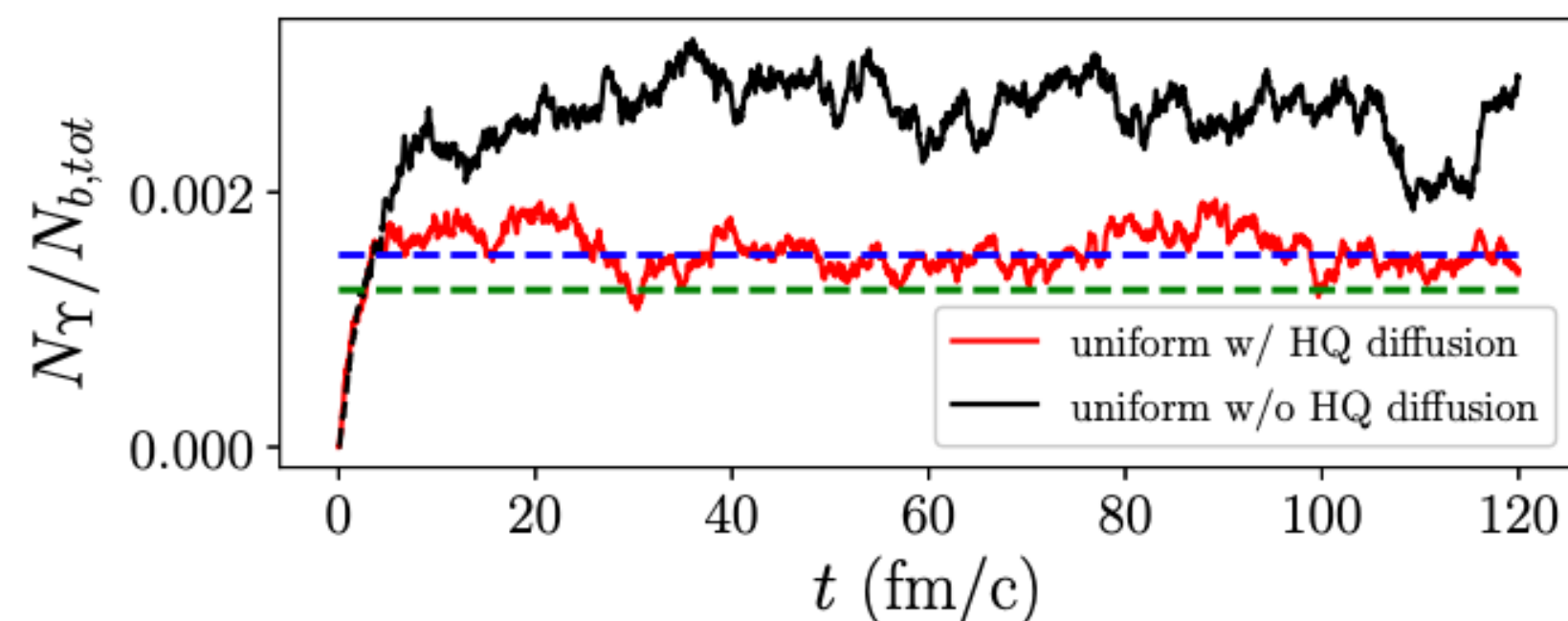
$$\Gamma = \int \frac{d^3p_{\text{rel}}}{(2\pi)^3} |\langle \psi_b | \mathbf{r} | \Psi_{\mathbf{p}_{\text{rel}}} \rangle|^2 [g_{\text{adj}}^{++}]^> (-\Delta E)$$

$$F = \int \frac{d^3p_{\text{cm}}}{(2\pi)^3} \frac{d^3p_{\text{rel}}}{(2\pi)^3} f_{Q\bar{Q}} |\langle \psi_b | \mathbf{r} | \Psi_{\mathbf{p}_{\text{rel}}} \rangle|^2 [g_{\text{adj}}^{--}]^> (\Delta E)$$

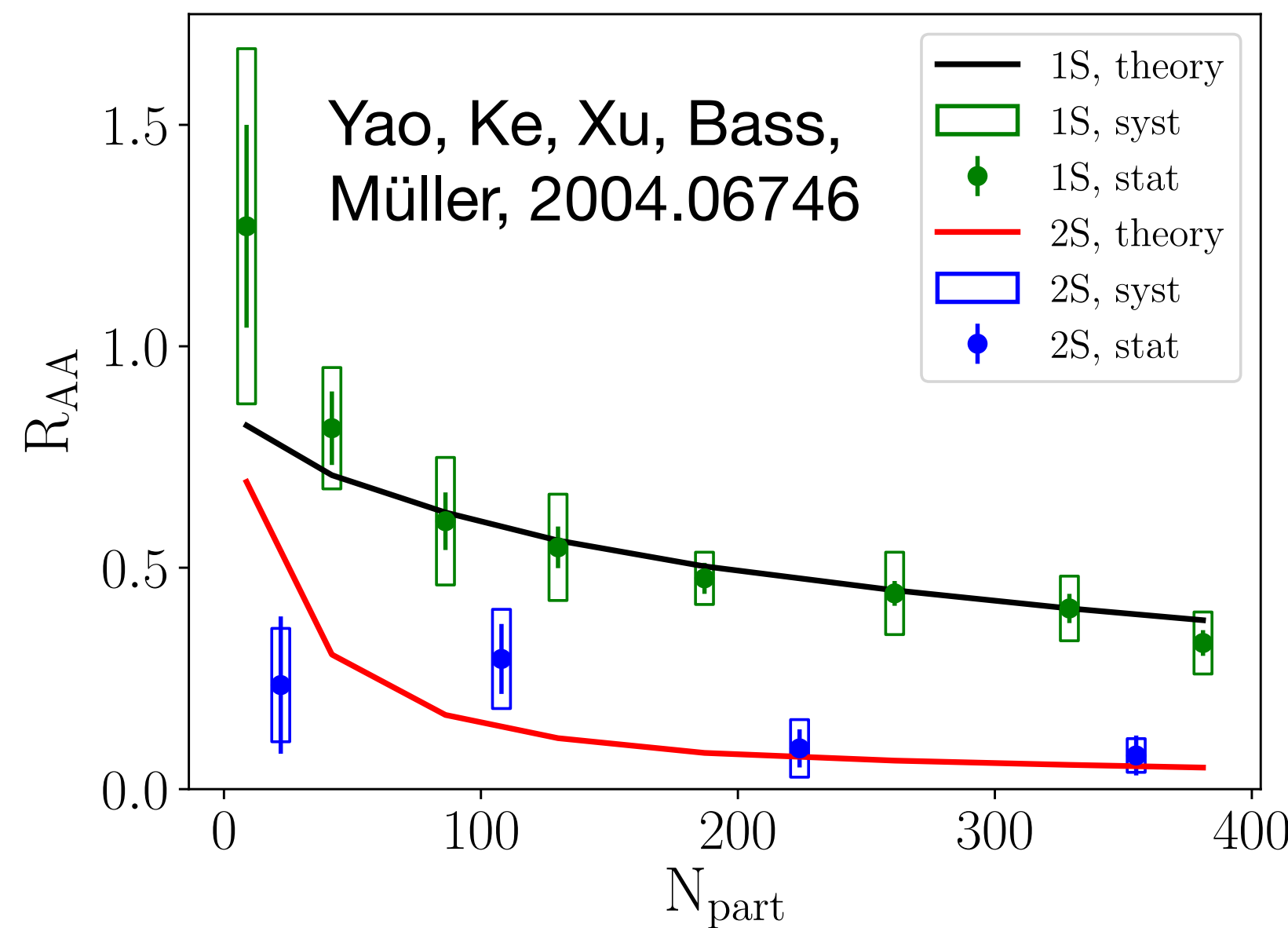
$$f_{Q\bar{Q}} \neq f_Q f_{\bar{Q}}$$



Open heavy quark dynamics
crucial for quarkonium equilibrium



Yao, Müller,
1709.03529

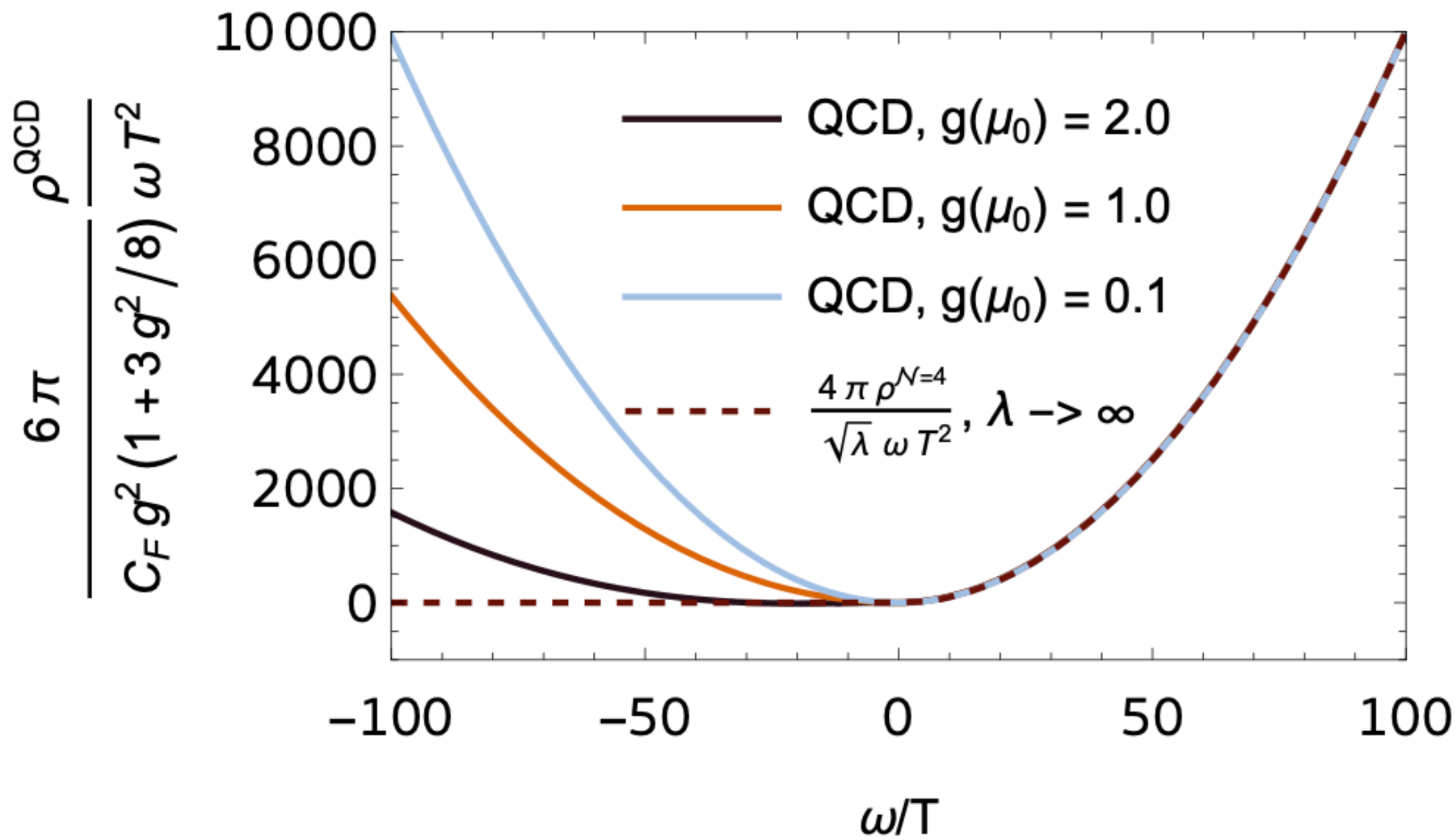
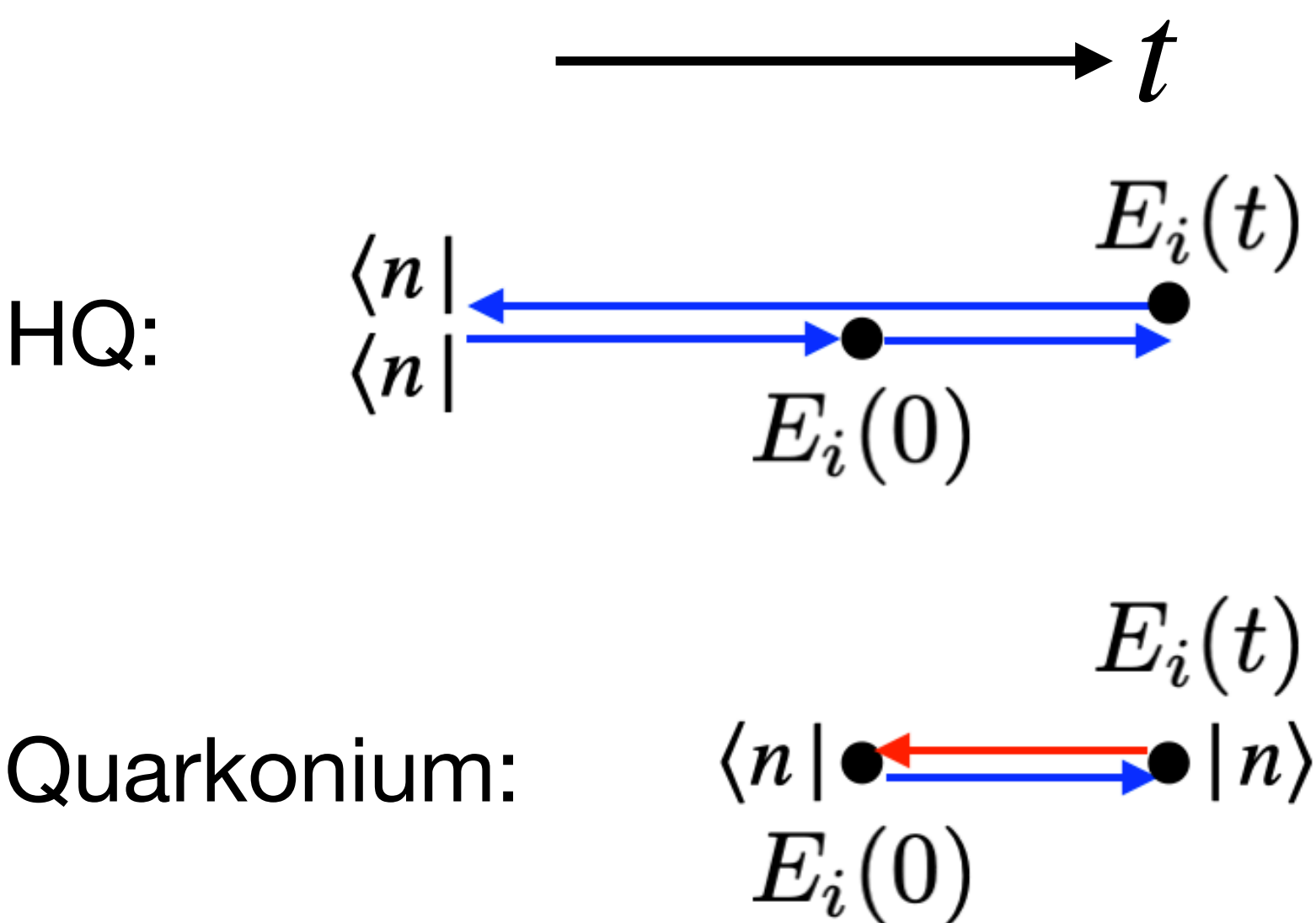


Chromoelectric Correlators for Heavy Quark and Quarkonium

- HQ diffusion** $\kappa_{\text{fund}} = \frac{g^2}{3N_c} \text{Re} \int dt \langle \text{Tr}_c [U(-\infty, t) E_i(t) U(t, 0) E_i(0) U(0, -\infty)] \rangle_{T, Q}$

- Quarkonium** $\kappa_{\text{adj}} = [g_{\text{adj}}^{++}]^>(\omega = 0) = \frac{g^2 T_F}{3N_c} \int dt \langle E_i^a(t) \mathcal{W}^{ab}(t, 0) E_i^b(0) \rangle_T$

- Operator orderings are different; evident difference seen in lattice calculation** Mayer-Steudte, PoS LATTICE2024 (2025) 205



Nijs, Scheihing, Yao, 2310.09325

Eller, Ghiglieri, Moore, 1903.08064; Binder, Mukaida, Scheihing, Yao, 2107.03945; Scheihing, Yao, 2205.04477, 2306.13127

Need new dynamical description to access strong coupling effect

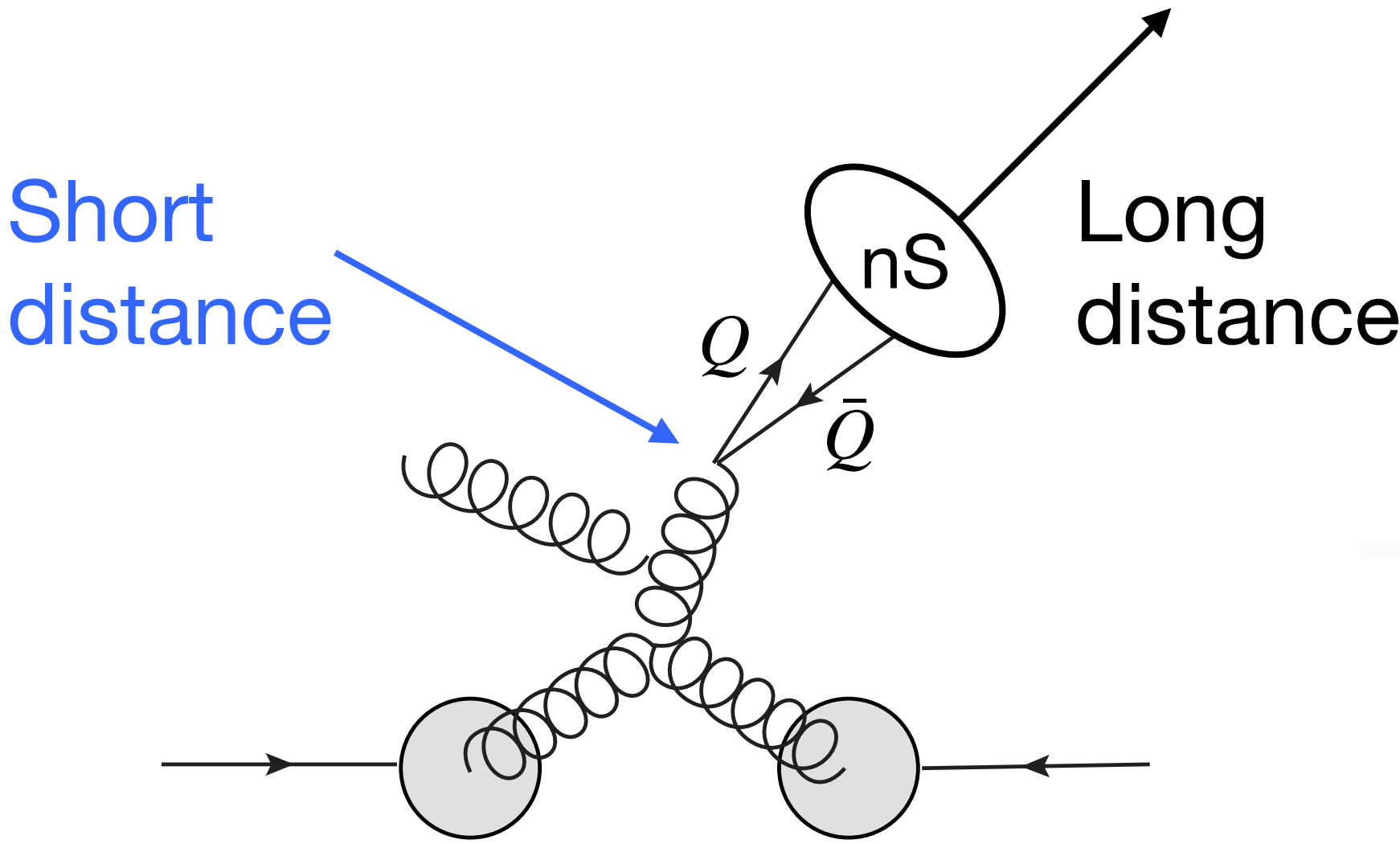
What Properties of QGP are Probed?

- | | | |
|---|---|---|
| • Light hadrons v_n and $\langle p_T \rangle$ | → | • Shear and bulk viscosities η, ζ |
| • High- p_T particles suppression | → | • Jet quenching parameter $\hat{q}$ |
| • Open heavy flavors R_{AA} and v_n | → | • Heavy quark drag and diffusion coefficients $\eta_D, \kappa_{\text{fund}}$ |
| • Low-p_T quarkonium R_{AA} and v_n | → | • Quarkonium transport parameters κ_{adj} and γ_{adj}, and more generally, chromoelectric correlator |

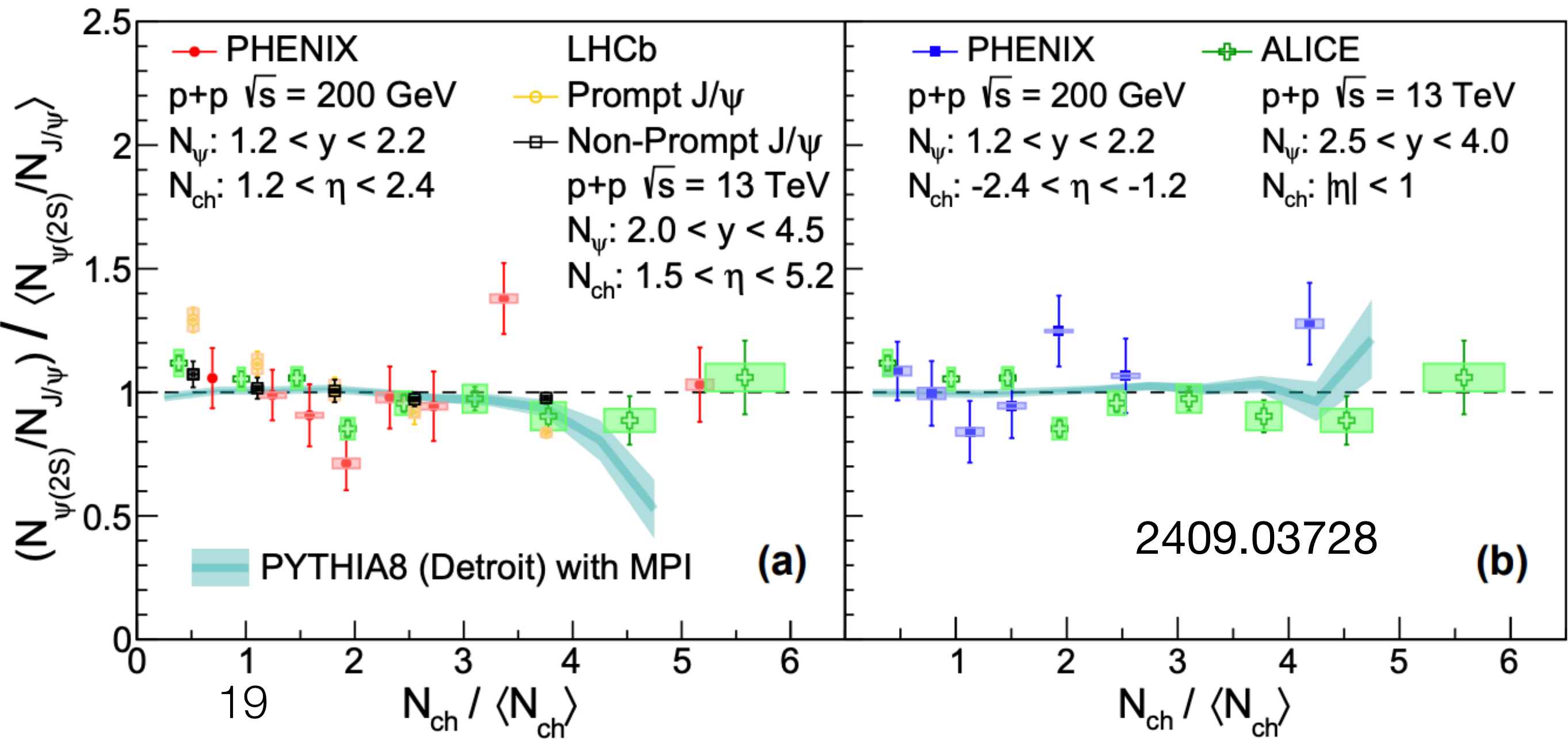
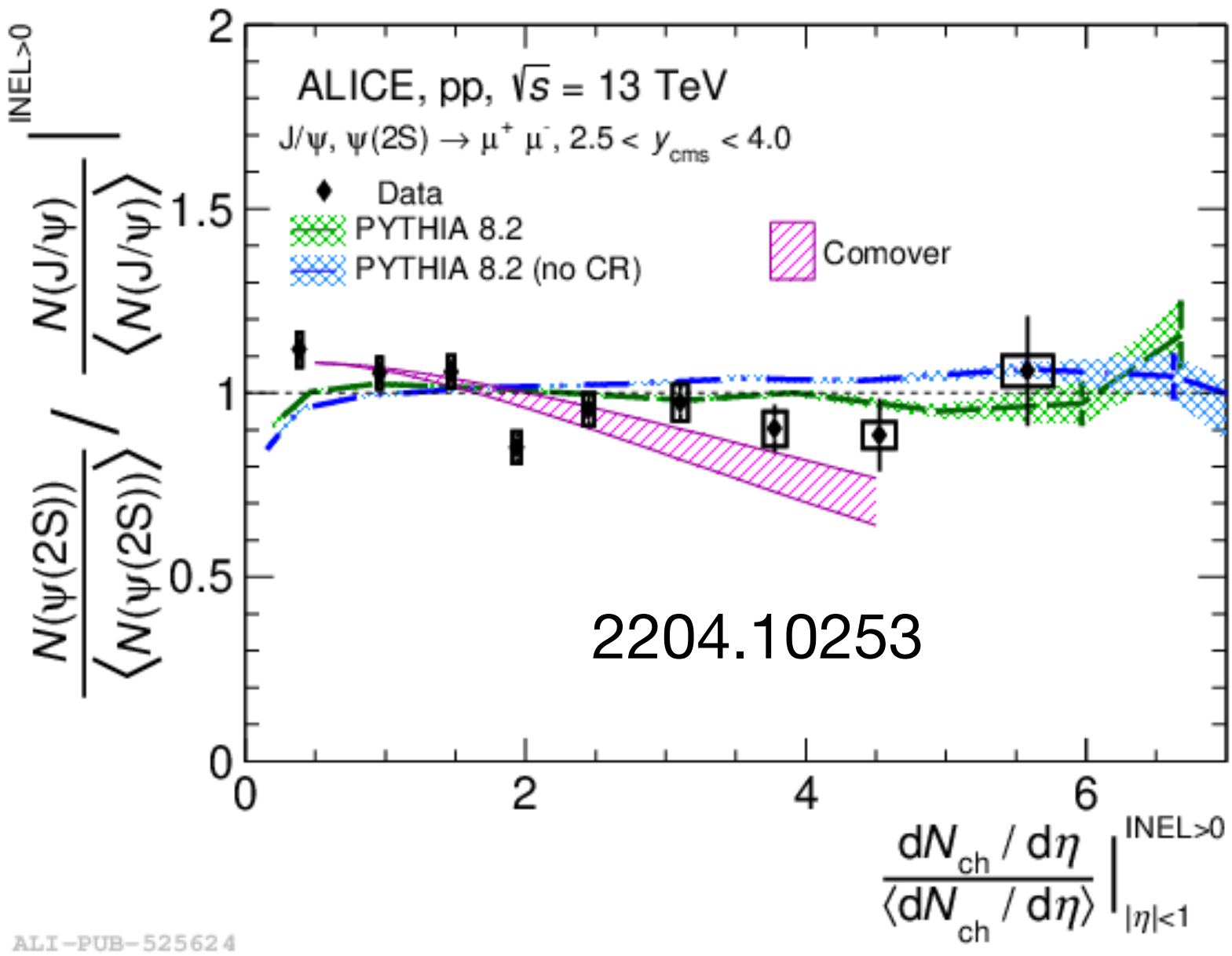
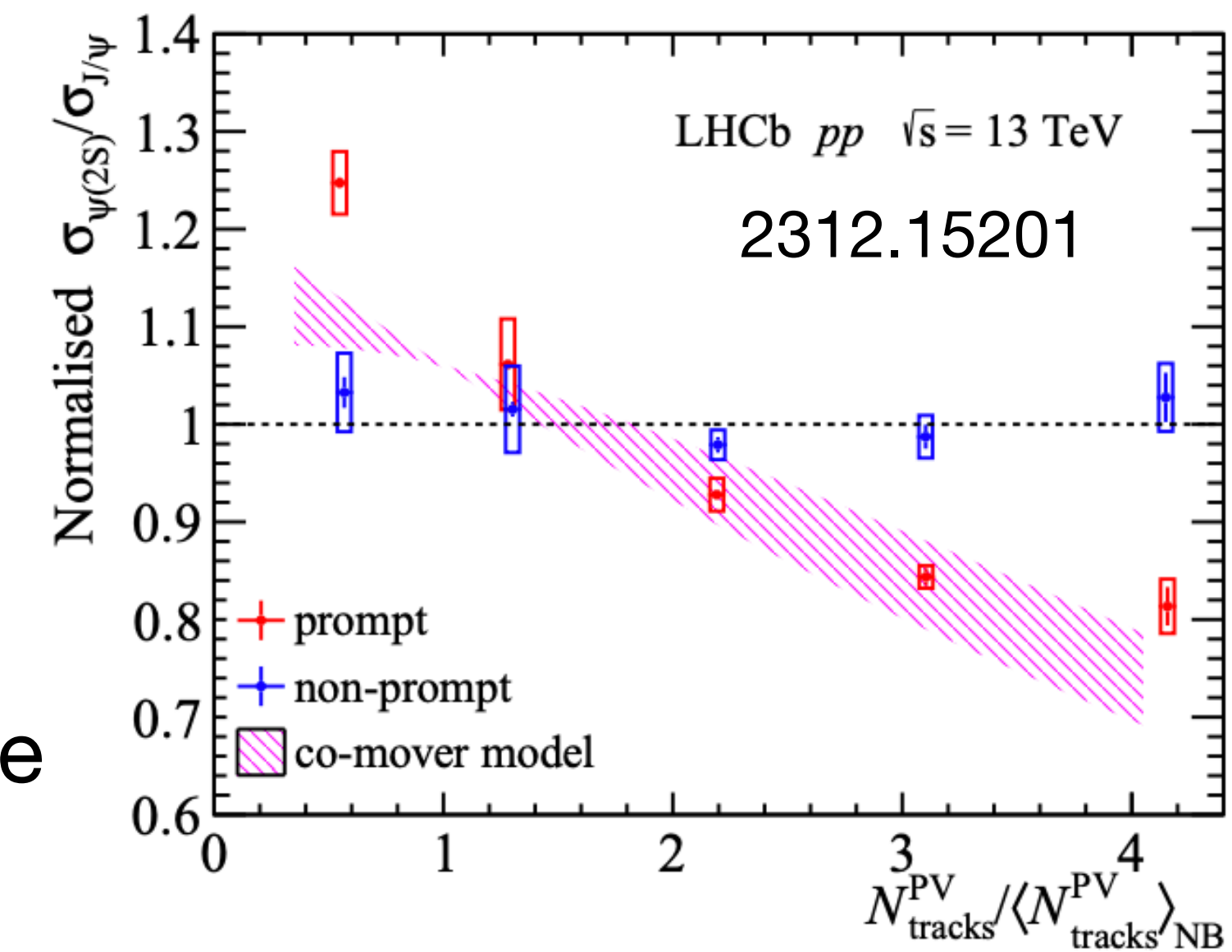
$$\langle E_i^a \mathcal{W}^{ab} E_i^b \rangle_T(\omega)$$

Quarkonium Production from AA to pp, pA

- $\frac{\psi(2S)}{J/\psi}$ v.s. event activity in pp

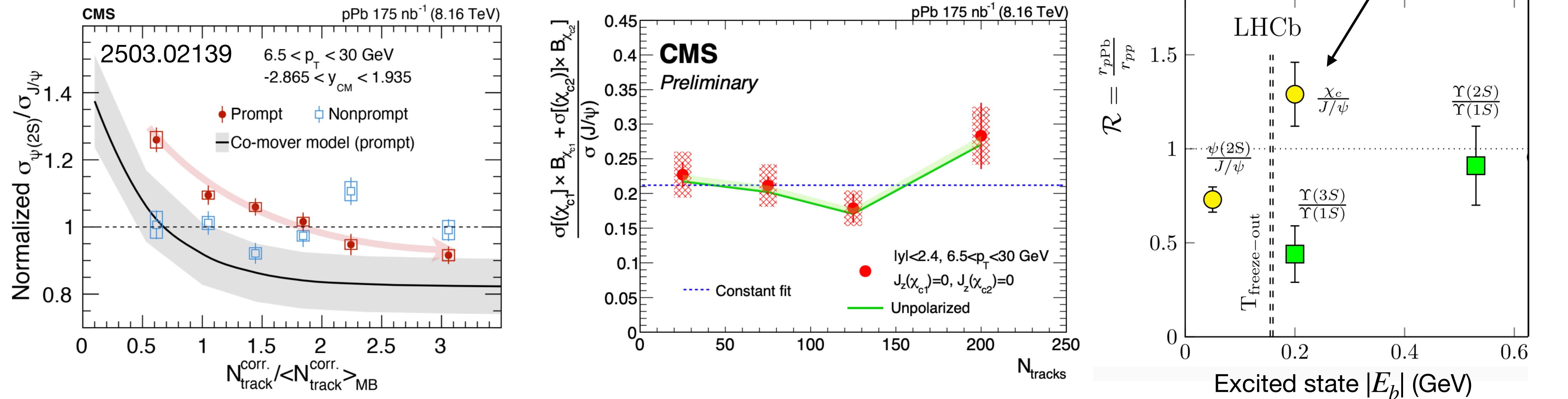


Universality of long-distance may break down in pp?



Quarkonium Production from AA to pp, pA

- Final-state interaction effect in quarkonium production in pA collisions



- Consistent theoretical description needed

Nature of “environment”, how its property depends on multiplicity

How formation of quarkonium is affected by “environment”, dependence on size & binding energy

Interaction time with “environment”, p_T/m dependence

Summary and Future Prospects

- Quarkonium probes deconfinement in AA
- **Uniquely** probes **chromoelectric correlator** of QGP in the dipole limit
- From AA to pp, pA
- Details of model descriptions: **Andronic, et al, 2402.04366**, summarizing efforts of EMMI Rapid Reaction Task Force started at GSI 2019
- Other interesting directions: (1) quarkonium spin alignment/polarization
 - (2) novel techniques to solve Lindblad equations
 - Neutral network: Lin, Luo, Yao, Shanahan, 2402.06607
 - Hydro formulation: Rais, van Hees, Greiner, 2502.21047
 - (3) quarkonium production at intermediate/high p_T , within jet