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# **Large amplitude characterization of CBETA arc cell**

## **using 3-D OPERA field maps**

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# 1 Introduction

◇ **Two things are addressed in these slides :**

**- the paraxial amplitude properties of CBETA arc cell.**

**This is in order essentially to make sure with tackle the next point with the correct hypotheses, namely**

**- large amplitude properties of CBETA arc cell.**

◇ **The investigation aims**

**- to validate, or to feedback on, the cell parameters and/or OPERA outcomes,**

**- with the criterion that the expectations are met (tune footprint, dynamical admittance, etc.)**

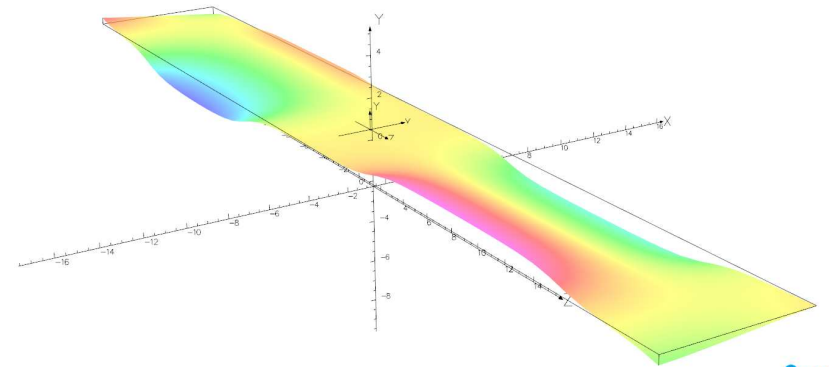
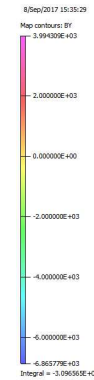
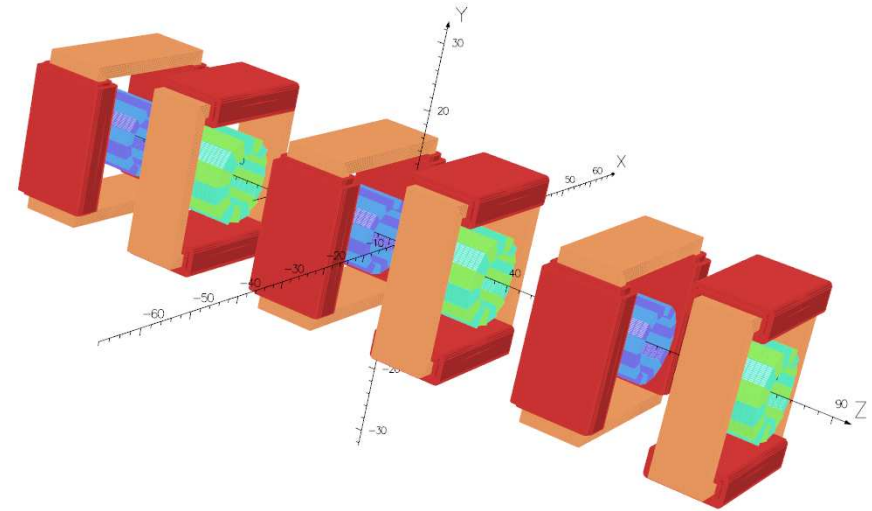
◇ **The exercise, back and forth between OPERA computation and ray-tracing, including the numerous changes in the lattice parameters, has kept us busy the past 2~3 years.**

## 2 Field maps

◇ The 3D field maps of concern have been produced by Nick Tsoupas

◇ As field maps are used, the study is *fully based on stepwise ray-tracing techniques* : particles are pushed across the field map step-by-step, typically millimeter steps.

- ◇ Two models of the cell can be used for ray-tracing :
  - either a single field map comprising QF and BD
  - or two separate field maps.



**UNITS**

Length	cm
Magn Flux Density	gauss
Magnetic Field	weber
Magn Scalar Pot	weber
Current Density	amp/cm
Power	W
Force	N

**MODEL DATA**

Scale: v6+H8sch\_1/2016-10-16  
 Magnetization: (T) (Gauss)  
 Nonlinear materials:  
 Simulation No: 1 of 1  
 5053904 elements  
 702024 nodes  
 24 conductors  
 NodeId interpolated fields  
 Activated in global coord

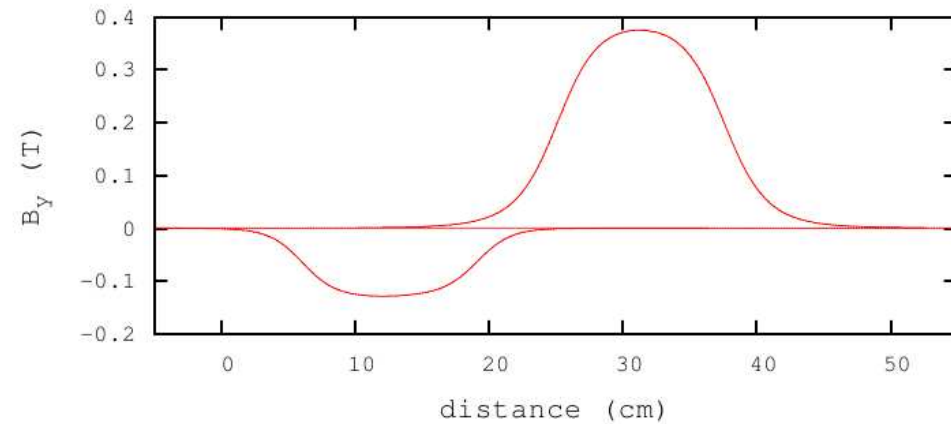
**Field Point Local Coord**  
 Local = Global

**FIELD EVALUATIONS (p)**  
 Cartesian: CARTESIAN (p)  
 x = -3.5 to 3.5

**Opera**  
COMSOL

The latter is the case in these slides, two OPERA file that we called  
 “QF\_v6 x=+-4p1y=+-1p4z=+-25 stp=1mm integral N35EH.table”  
 “BD\_v6 x=+-4p1y=+-1p4z=+-25 stp=1mm integral N35EH.table”

◇ The cell magnets being short, the field along the orbits does not feature any plateau.

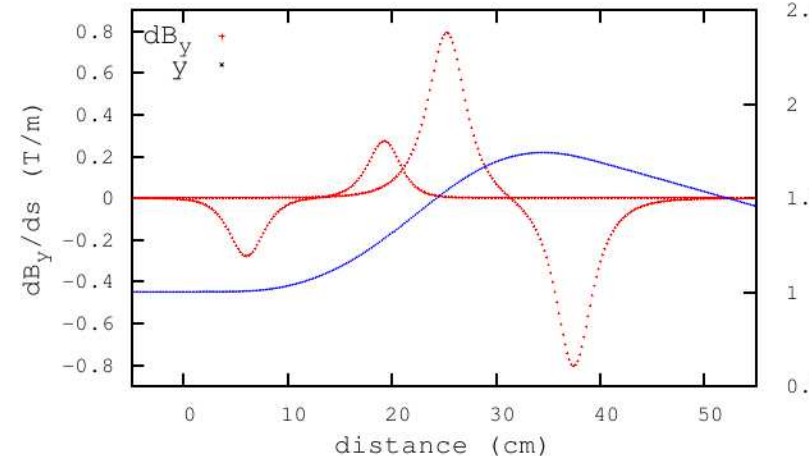


◇ As a consequence it has a rich content in derivatives, as results from

$$V_1(s, x, y) = G_1(s)y - \frac{G_1^{(2)}(s)}{8}(x^2 + y^2)y + \frac{G_1^{(4)}(s)}{192}(x^2 + y^2)^2y + \dots \quad (\text{dipole})$$

$$V_2(s, x, y) = G_2(s)xy - \frac{G_2^{(2)}(s)}{12}(x^2 + y^2)xy + \frac{G_2^{(4)}(s)}{384}(x^2 + y^2)^2xy + \dots \quad (\text{quadrupole})$$

◇ This is the origin of the huge reduction of the dynamical admittance, compared to that of linear magnets (slide #10) (the latter is limited as well though, a sufficient cause for that are kinematic nonlinearities)



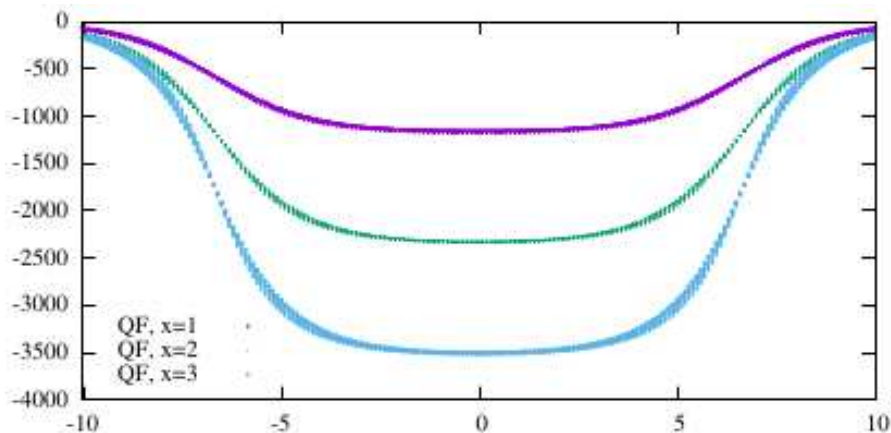
First order field derivative  $dB_y/ds$  (red) along a trajectory off-mid-plane (blue).

- ◇ **An ideal cell without any random errors (neither field not alignment) is considered.**
  
- ◇ **It does include systematic non-linearities, namely,**
  - (i) the multipole content intrinsic to the magnet geometry in the OPERA simulation, and**
  
  - (ii) the non-linearity content resulting from the longitudinal form factor as addressed above. Note that the latter includes all orders in  $x, y$ .**
  
- ◇ **In addition, as stepwise ray-tracing is used the kinematic non-linearities are present in the motion, always.**

- Check the positioning of the magnet with respect to the map frame (which is also the mesh frame). They are centered on the center of the field map.

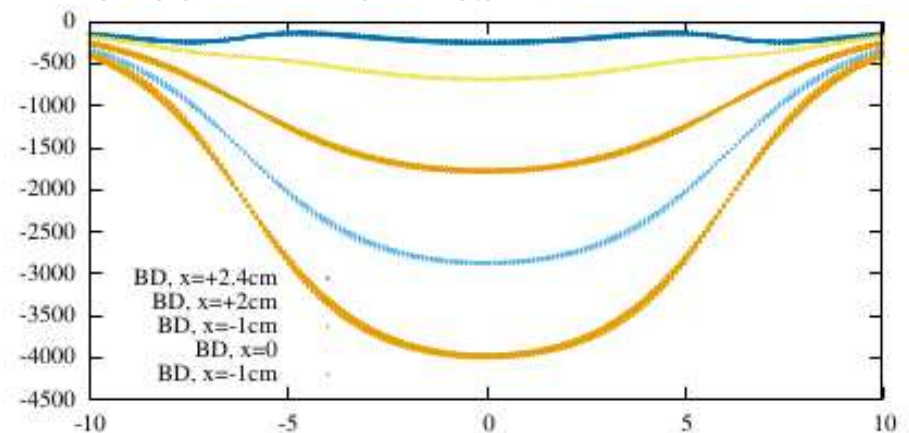
◇ Field from QF field map file, along  $x=1, 2$  and  $3$  cm.

The field is symmetric wrt.  $s=0$  plane  $\rightarrow$  magnet centered in the mesh.



◇ Field from BD field map file, along lines at various constant- $x$  values.

The field is symmetric wrt.  $s=0$  plane  $\rightarrow$  magnet centered in the mesh.



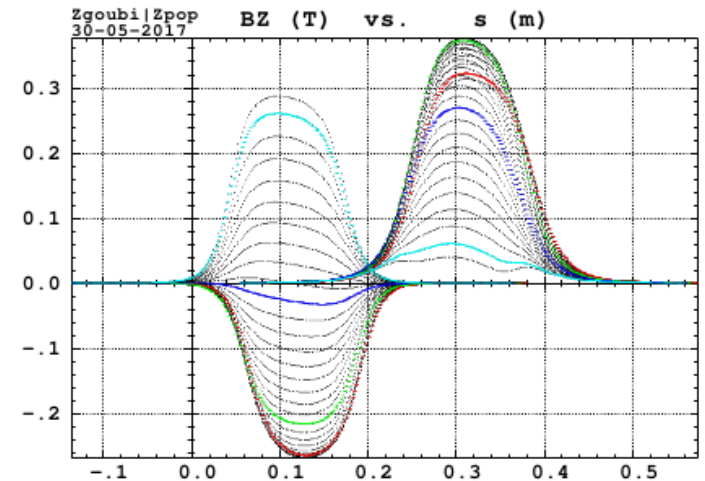
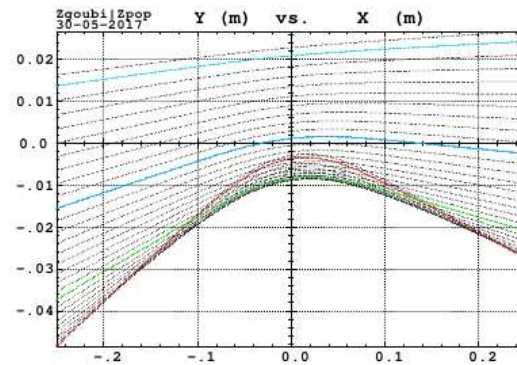
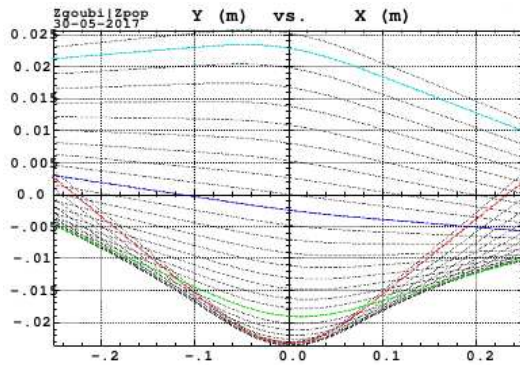
- Check the centering of extreme orbits (42 and 150 MeV) in QF and BD maps.

(Orbits over a 40-153 MeV range are plotted here, the design energies are colored.)

◇ QF : The 42 and 150 MeV orbit excursions reach  $\pm 23.33$  mm : the beam is correctly centered.

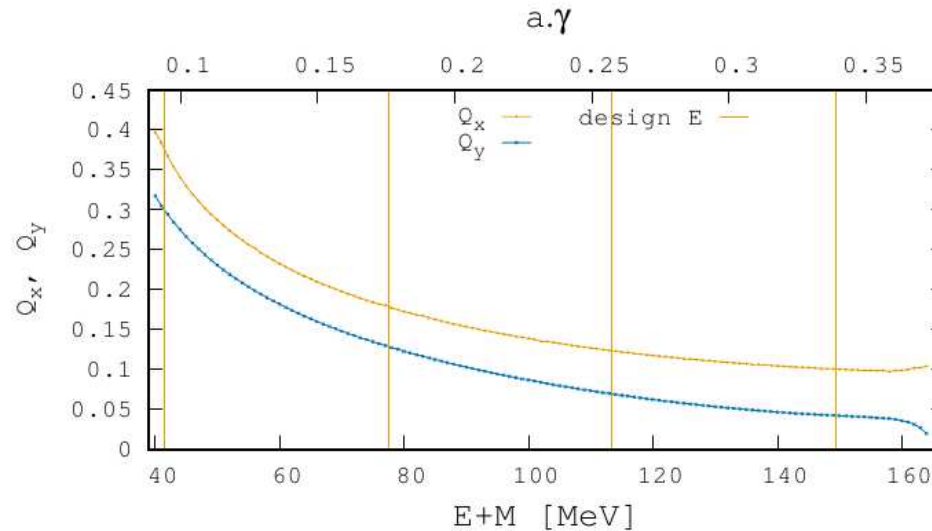
◇ Orbit excursion in BD magnet (centered on 0, length 12 cm).

◇ Field along 30 different orbits in QF and BD (40 to 153 MeV).

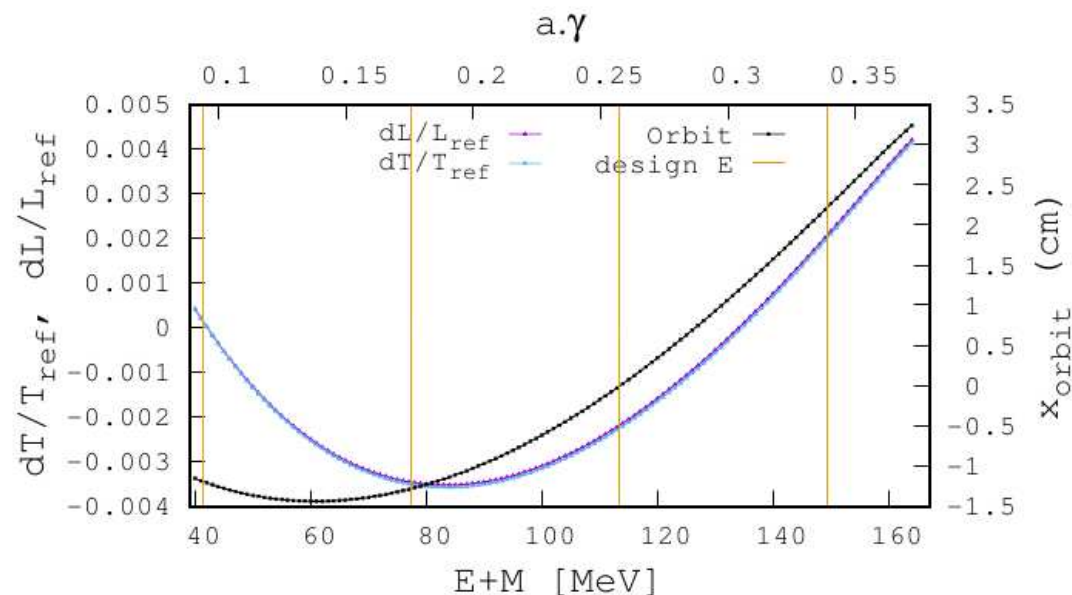




- Energy dependence of cell tunes (from first order mapping), from 40 MeV up to maximum stable energy ( $\approx 166$  MeV).

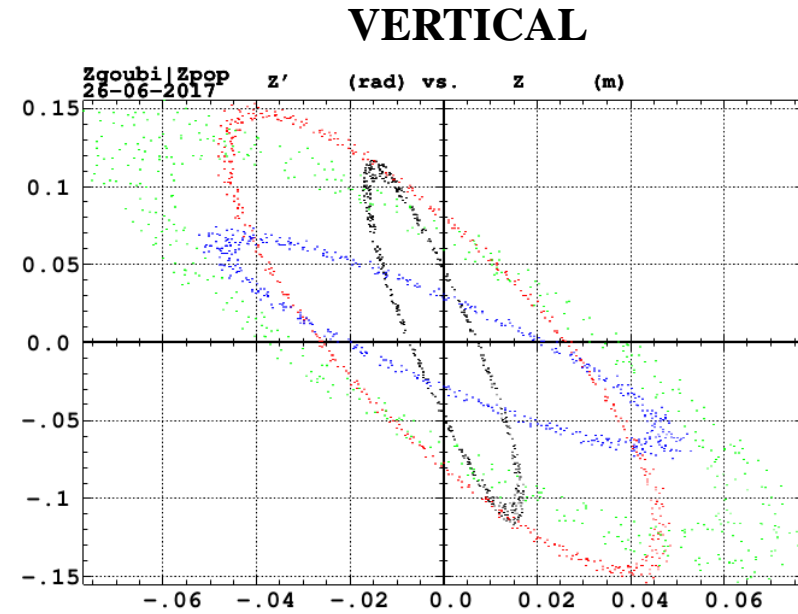
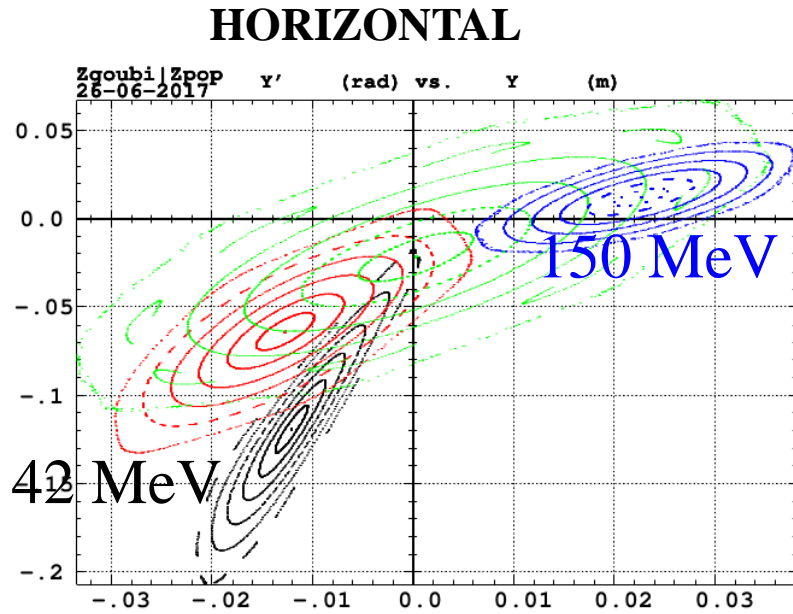


- Energy dependence of orbit position (at center of long drift), path length and time of flight.



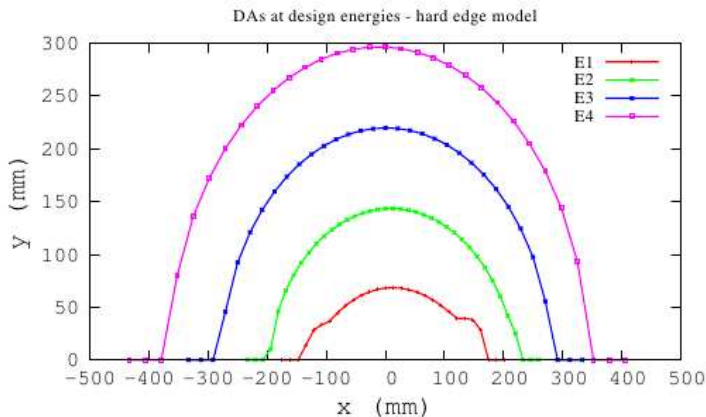
### 3 Large amplitude

- Maximum stable invariants, for each one of the 4 design energies.



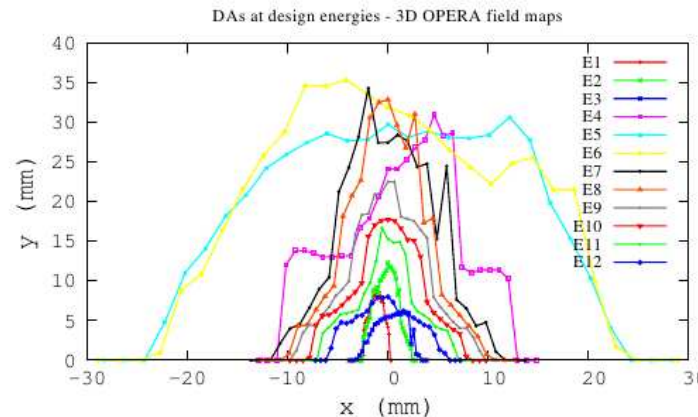
- A parenthesis concerning the “dynamical admittance”.
- Typical behavior of hard-edge model versus realistic fields :

**Hard edge model**



E1 to E4 :  
42, 78,  
114 and  
150 MeV

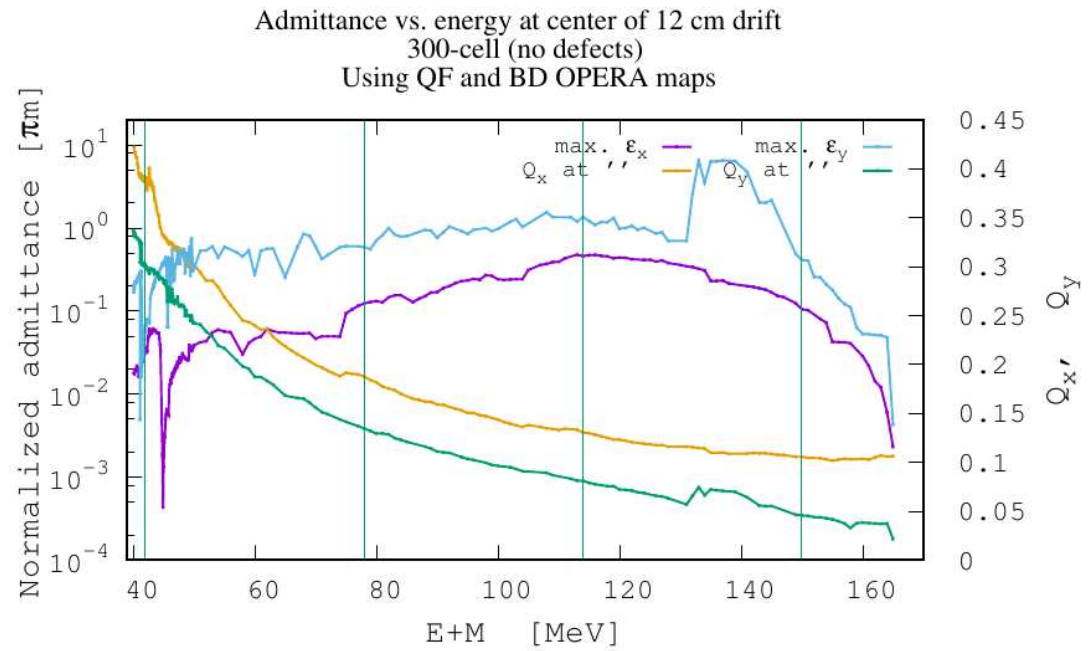
**3D OPERA field maps**



E1 to  
E12 : 40,  
41, 42,  
78 114,  
119, 150,  
152, 154,  
156, 158,  
160 MeV.

- A scan of the H and V dynamical admittances of a **300-cell beam line**, over **40-166 MeV** (this is the surface of the phase space invariant at maximum stable amplitude, respectively H and V, previous slide).

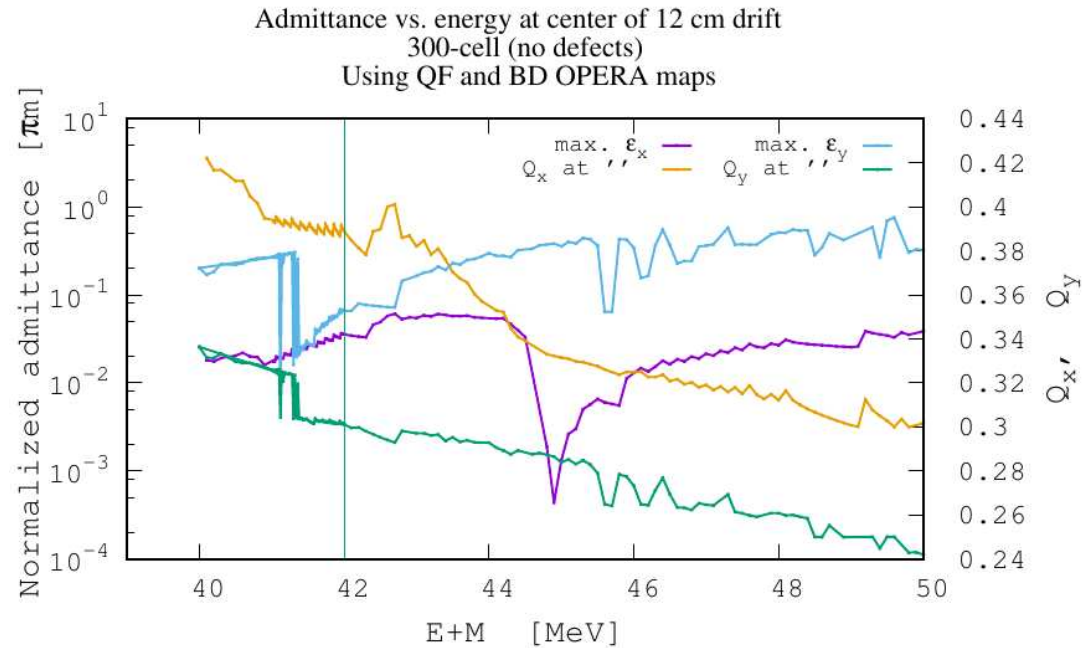
- ◊ The vertical bar is at **42 MeV**.



- A zoom in the **40-50 MeV region**.

- ◊ The  $A_y$  dip at **41+ MeV** is attributed to  $Q_x + 2Q_y = \text{int. sum resonance line}$ , TBC ! source : skew sextupole excitation (see middle and bot. Q-diagrams in slide #4).

- ◊ The  $A_x$  dip at **45- MeV** is due to  $3Q_x = \text{integer}$  (upright sextupole).

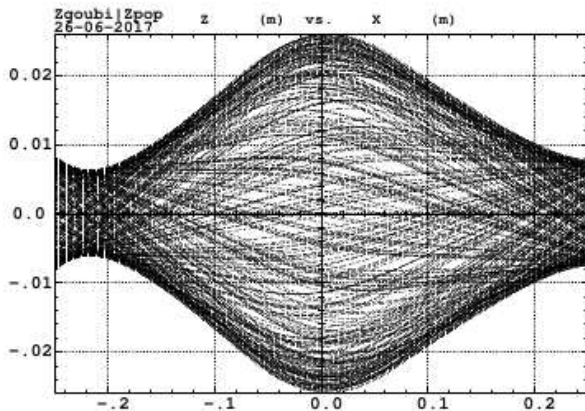


- $Q_x$  and  $Q_y$  tunes in these plots are for the motion at maximum stable amplitude (*i.e.*, the tunes for the maximum invariants, previous slide).

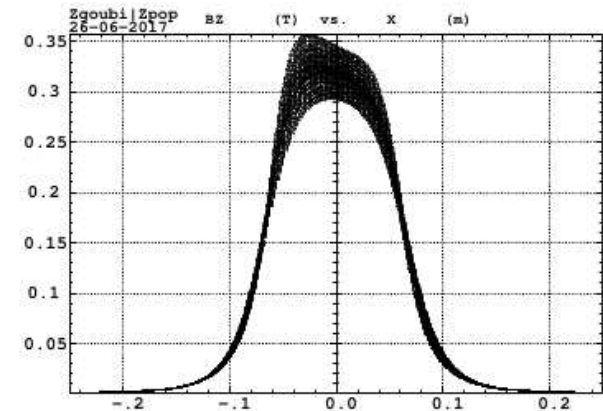
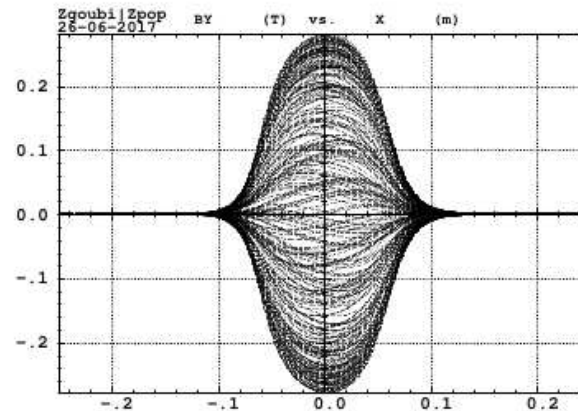
- Because the maximum stable vertical invariant *really is* at large amplitude ( $\pm 2.6$  cm, see left plot below), need to check that the field at such large amplitude is still correctly computed (at least, “looks realistic”), otherwise DA computations are questionable...

A 42 MeV electron is launched on maximum vertical stable invariant, and observed across BD over a few 10s of passes (that’s sufficient for the electron to reach large amplitudes/fields) :

◇ y-projection of the vertical motion.



◇ Radial (left) and vertical (right) components of the field experienced. Does not look to badly computed...



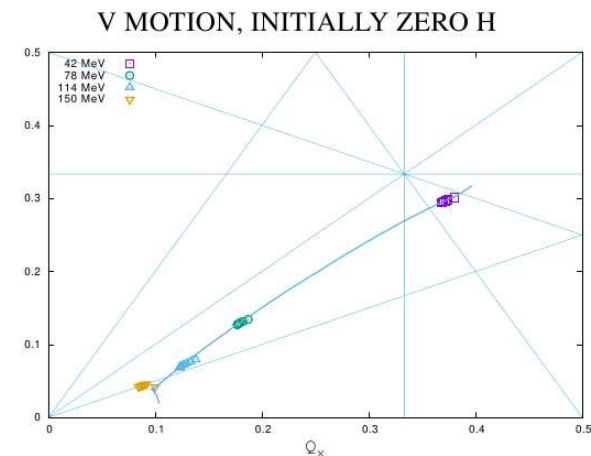
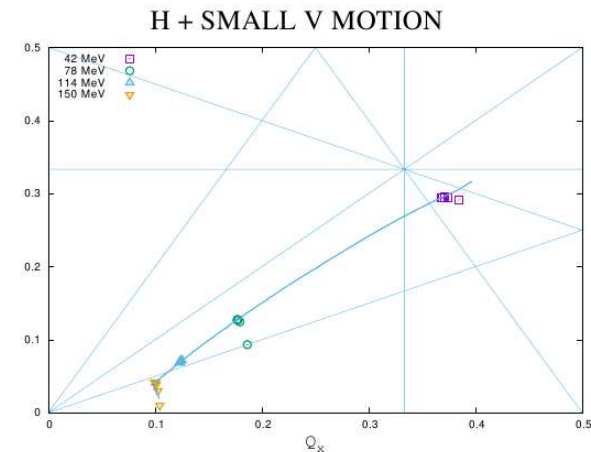
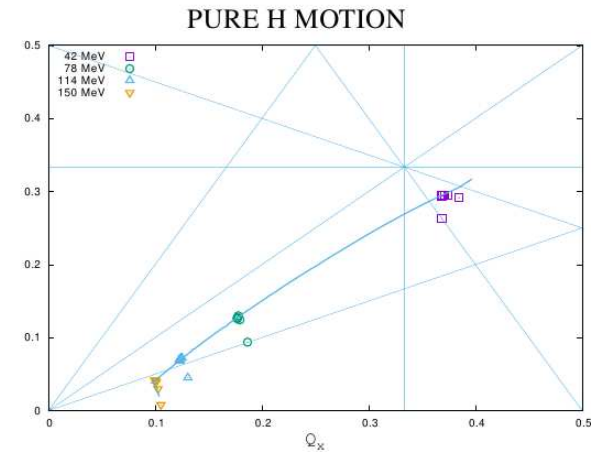
## 4 Tune footprint

- Spreading of  $Q_x$ ,  $Q_y$  under the effect of horizontal ( $dQ_x/dax$ ,  $dQ_x/day$ ) and vertical ( $dQ_y/dax$ ,  $dQ_y/day$ ) anharmonicities.

◇ For each of the 4 design energies, 7 invariants from paraxial to maximum stable are tracked and the tunes are computed.

- Top : case of pure H motion
- Middle : initial H + subliminal V
- Bottom : initial V + subliminal H

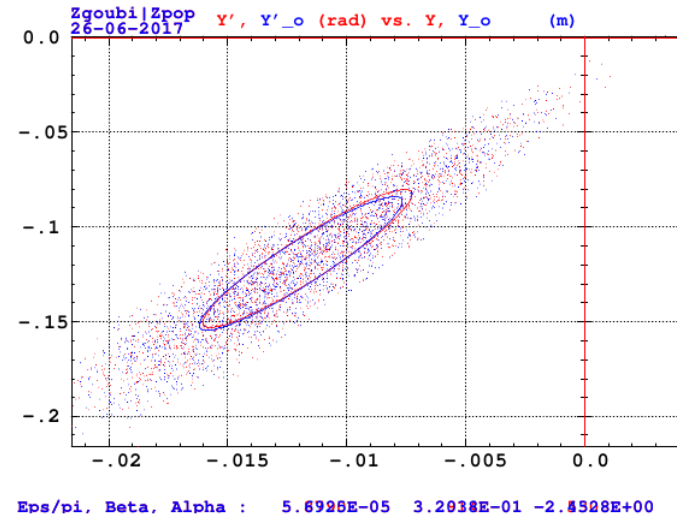
The solid line spans the 42-150 MeV range, case of paraxial tunes.



## 5 Transmission

- The focus here is on 42 MeV, the region of smaller admittance (previous slides).
- 10000 particles are launched in a wide  $\epsilon_x/\pi = 10^{-4}$  m,  $\epsilon_y/\pi = 10^{-3}$  m, geometric.
- Of these, 21% make it to the end of the 300-cell channel. This determines the H, V admittance of the 42 MeV channel, at its entrance.

- $x, x'$  coordinates (initial (blue) and final (red)) of the particles that make it thru, and matching ellipses. The geometrical admittance is  $\epsilon_x/\pi = 5.7 \cdot 10^{-5}$  m, i.e.,  $\epsilon_x/\pi = 4.8$  mm, normalized.



- $y, y'$  coordinates (initial (blue) and final (red)) of the particles that make it thru, and matching ellipses. The geometrical admittance is  $\epsilon_y/\pi = 1.7 \cdot 10^{-4}$  m, i.e.,  $\epsilon_y/\pi = 14.3$  mm normalized.

