







TRACES OF NON-EQUILIBRIUM DYNAMICS IN RELATIVISTIC HEAVY-ION COLLISIONS

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Description of relativistic heavy-ion collisions

Two types of model have been successful in describing relativistic heavyion collisions:

 Hybrid approaches: Hydro description of the QGP + microscopic transport for the hadronic sector



- Simplified dynamics, assumption of local equilibrium
- Direct access to QGP properties in equilibrium
- **Transport approaches:** full microscopic description of heavy-ion collisions
- No assumptions about local equilibrium
- Can we identify which model features reflect the actual physical nature of the QGP ?
- Access to the QGP degrees of freedom
- Is there any differences in the dynamical evolution of the system that can be attribtuted to non-equilibrium effects ?



2D+1 viscous hydrodynamics

Space-time evolution of the QGP via conservation equations:

$$\partial_{\mu}T^{\mu\nu} = 0 \qquad T^{\mu\nu} = e \, u^{\mu}u^{\nu} - \Delta^{\mu\nu}(P + \Pi) + \pi^{\mu\nu}$$

 u^{μ} : cell 4-velocityP: local isotropic pressureΠ: bulk viscous pressuree: local energy density $\pi^{\mu\nu}$: shear stress tensor $\Delta^{\mu\nu} = g^{\mu\nu} - u^{\mu}u^{\nu}$

- For this study we use VISH2+1: time evolution of the viscous corrections through the 2nd order Israel-Stewart equations
 - $\square \quad \eta: \text{ shear viscosity}$
 - $\Box \quad \zeta : \text{ bulk viscosity}$

 $\tau_{\Pi}\dot{\Pi} + \Pi = -\zeta\theta - \delta_{\Pi\Pi}\Pi\theta + \phi_{1}\Pi^{2} + \lambda_{\Pi\pi}\pi^{\mu\nu}\sigma_{\mu\nu} + \phi_{3}\pi^{\mu\nu}\pi_{\mu\nu}$

$$\tau_{\pi} \dot{\pi}^{\langle \mu\nu\rangle} + \pi^{\mu\nu} = 2\eta \sigma^{\mu\nu} + 2\pi_{\alpha}^{\langle \mu} w^{\nu\rangle\alpha} - \delta_{\pi\pi} \pi^{\mu\nu} \theta + \phi_7 \pi_{\alpha}^{\langle \mu} \pi^{\nu\rangle\alpha} - \tau_{\pi\pi} \pi_{\alpha}^{\langle \mu} \sigma^{\nu\rangle\alpha} + \lambda_{\pi\Pi} \Pi \sigma^{\mu\nu} + \phi_6 \Pi \pi^{\mu\nu}$$

C. Shen, Z. Qiu, H. Song, J. Bernhard, S. Bass and U. Heinz, Comput. Phys. Commun. 199 (2016) 61-85









Stages of a collision in VISHNU



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The PHSD transport approach



- **Goal:** Study the properties of strongly interacting matter under extreme conditions from a microscopic point of view
- **Realization:** dynamical many-body transport approach

The PHSD transport approach



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Parton-Hadron-String-Dynamics (PHSD)

- Explicit parton-parton interactions, explicit phase transiton from hadronic to partonic degrees of freedom
- Transport theory: off-shell transport equations in phase-space representation based on Kadanoff-Baym equations for the partonic and hadronic phase



W.Cassing, E.Bratkovskaya, PRC 78 (2008) 034919; NPA831 (2009) 215; W.Cassing, EPJ ST 168 (2009) 3

Traces of non-equilibrium dynamics in relativistic heavy-ion collisions

Dynamical Quasi-Particle Model (DQPM)

Medium evolution: Hydro vs PHSD

The QGP phase is described in terms of interacting quasiparticles: quarks and gluons with Lorentzian spectral functions:

$$\rho_i(\omega,T) = \frac{4\omega\Gamma_i(T)}{(\omega^2 - \mathbf{p}^2 - M_i^2(T))^2 + 4\omega^2\Gamma_i^2(T)} \qquad (i = q, \bar{q}, g)$$

 Properties of quasiparticles (large widths and masses) are fitted to the lattice QCD results



 DQPM provides mean-fields (1P1) for quarks and gluons as well as effective 2-body interactions (2P1)

Traces of non-equilibrium dynamics in relativistic heavy-ion collisions

Peshier, Cassing, PRL 94 (2005) 172301; Cassing, NPA 791 (2007) 365: NPA 793 (2007)





4 ³ k [GeV]

5 5

 ρ [GeV²]

0.1



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Introduction Initial conditions from PHSD Medium evolution

Medium evolution: Hydro vs PHSD

Summary

Stages of a collision in PHSD



- String formation in primary NN collisions
- → decays to pre-hadrons (baryons and mesons)



Stages of a collision in PHSD



Stages of a collision in PHSD



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Stages of a collision in PHSD



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Stages of a collision in PHSD



Stages of a collision in PHSD

t = 7.31921 fm/c

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Summary

Stages of a collision in PHSD

Stages of a collision in PHSD

Coarse graining of PHSD

- Goal: Initialize the hydro with a non-equilibrium profile from PHSD to compare the both evolutions
- Energy-momentum tensor should have the form:

$$T^{\mu\nu} = e \, u^{\mu} u^{\nu} - \Delta^{\mu\nu} (P + \Pi) + \pi^{\mu\nu}$$

- The viscous corrections should be small
- Spatial gradients should not be too large
- Coarse graining of PHSD medium in the transverse plane at the midrapidity region ($z \approx 0$. fm):

$$T^{\mu\nu}(x) = \sum_{i} \int_{0}^{\infty} \frac{d^{3}p_{i}}{(2\pi)^{3}} f_{i}(E_{i}) \frac{p_{i}^{\mu}p_{i}^{\nu}}{E_{i}} = \frac{1}{V} \sum_{i} \frac{p_{i}^{\mu}p_{i}^{\nu}}{E_{i}}$$
$$\Delta x = \Delta y = 1 \text{ fm}$$

Evaluation of the energy momentum tensor

Diagonalization of the energy-momentum tensor

$$T^{\mu\nu} (x_{\nu})_i = \lambda_i (x^{\mu})_i = \lambda_i g^{\mu\nu} (x_{\nu})_i$$

Landau-matching condition:

$$T^{\mu\nu}u_{\nu} = eu^{\mu} = (eg^{\mu\nu})u_{\nu}$$

Evaluation of the characteristic polynomial:

$$P(\lambda) = \begin{vmatrix} T^{00} - \lambda & T^{01} & T^{02} & T^{03} \\ T^{10} & T^{11} + \lambda & T^{12} & T^{13} \\ T^{20} & T^{21} & T^{22} + \lambda & T^{23} \\ T^{30} & T^{31} & T^{32} & T^{33} + \lambda \end{vmatrix}$$

The four solutions λ_i are identified to $(e, -P_1, -P_2, -P_3)$

The pressure components P_i do not necessarily correspond to (P_x, P_y, P_z)

Evaluation of the energy momentum tensor

Using the Landau-matching condition we have :

$$\begin{array}{l} (T^{00}-e)+T^{01}X+T^{02}Y+T^{03}Z=0\\ T^{10}+(T^{11}+e)X+T^{12}Y+T^{13}Z=0\\ T^{20}+T^{21}X+(T^{22}+e)Y+T^{23}Z=0\\ T^{30}+T^{31}X+T^{32}Y+(T^{33}+e)Z=0 \end{array}$$

• With the 4-velocity $u_{\nu} = \gamma (1, X, Y, Z) = \gamma (1, -\beta_x, -\beta_y, -\beta_z)$

Evaluation of viscous corrections:

$$\Pi = -\frac{1}{3} \,\Delta_{\mu\nu} T^{\mu\nu} - \mathcal{P}_{\rm EoS}$$

$$\pi^{\mu\nu} = T^{\mu\nu} - eu^{\mu}u^{\nu} - \Delta^{\mu\nu}(\mathcal{P}_{\rm EoS} + \Pi)$$

Liu et al., Phys. Rev. C 91, 064906 (2015)

Dependance on the parallel ensembles

- In order to have a smooth mean-field potential to propagate the particles and to initialize the hydro evolution, we need to average PHSD events over N parallel ensembles.
- The more parallel ensembles (NUM) are considered, the smoother the obtained profile is:

• For this study we take NUM = 30 parralel events

Pressure isotropization

- We look at the pressure components P_i of several cells along the x-axis
- The longitudinal pressure dominates when the collisions happens and then decrease rapidly towards 0
- The transverse pressure increases with time and reach the equilibrated pressure at 0.5 – 1. fm/c

PHSD profile averaged over 100*NUM events

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- We look at the pressure components P_i of several cells along the x-axis
- The longitudinal pressure dominates when the collisions happens and then decrease rapidly towards 0
- The transverse pressure increases with time and reach the equilibrated pressure at 0.5 – 1. fm/c
- We choose t₀ = 0.6 fm/c to initialize the hydro evolution

PHSD profile averaged over 100*NUM events

Pressure isotropization

Evaluation of the relative value between the transverse pressure and the one given by the QCD EoS

PHSD profile averaged over 100*NUM events

In averaged, we can see that the PHSD medium reach the EoS pressure in the transverse direction after a time of 0.5 – 1. fm/c

Adjustment of hydro parameters

In order to compare the two models, the temperature-dependant shear viscosity from PHSD is used in the hydro code

PHSD results from: Ozvenchuk et al., **Phys. Rev. C 89, 064903 (2013)** Duke Bayesian analysis for the bulk viscosity: Bernhard et al, **arXiv:1704.04462**

PHSD

Summary

Space-time evolution (cell velocity β)

Space-time evolution (cell velocity β)

- The PHSD evolution remains chaotic for all times
- The hydro code seems to smooth the initial PHSD profile during the evolution in time

Fourier images of energy density

- How to quantify this change in the hydro evolution ?
- Fourier transformation of the energy density profile:

$$\tilde{e}(k_x, k_y) = \frac{1}{mn} \frac{1}{n} \sum_{x=0}^{m-1} \sum_{y=0}^{n-1} e(x, y) e^{2\pi i \left(\frac{xk_x}{m} + \frac{yk_y}{n}\right)}$$

Radial distribution of the Fourier modes of the energy density

Shorter wavelength modes survive only in PHSD which indicates the constant inhomogeneity of the QGP medium

Spatial eccentricity

- □ We compare 100 events from PHSD and the event-by-event hydro code
- The spatial eccentricity is calculated as a function of time by the formula:

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$$\epsilon_2 = \frac{\sqrt{\{r^2 \cos(2\phi)\}^2 + \{r^2 \sin(n\phi)\}^2}}{\{r^2\}}$$

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- The green dots represent each
 PHSD event: large fluctuations
- Each hydro event has a very smooth evolution contrary to PHSD
- Both PHSD and hydro agree with each other when the full initial state information is taken into account (initial velocity β_i + viscous corrections $\pi_i^{\mu\nu}$)

PHSD profile averaged over 100*NUM events

Momentum eccentricity

- Momentum anisotropies reflect the medium's response to initial spatial anisotropies:
- □ Including the initial flow velocity β_i in the hydro has a huge effect in ϵ_p
- The initial viscous corrections $\pi_i^{\mu\nu}$ play only a role at small times

$$\epsilon_p = \frac{\int dx dy \ (T^{xx} - T^{yy})}{\int dx dy \ (T^{xx} + T^{yy})}$$

 The increase of the bulk viscosity produces a bump during the hadronization process

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PHSD profile averaged over 100*NUM events

- We have compared two descriptions of the QGP medium evolution in heavy-ion collisions: PHSD and 2+1D hydro
- We matched the hydrodynamical evolution as closely as possible with the PHSD medium
- Similar QCD EoS Same shear viscosity η/s Flexible bulk viscosity ζ/s

In average, both QGP mediums evolve in a similar way, but:

- Very strong fluctuations are observed in the PHSD medium at any time of the evolution while the hydrodynamical medium evolves smoothly
- Strong response of the hydrodynamical medium to transport coefficients
- In the future, non-equilibrium traces on other observables (hard probes...)

Medium evolution: Hydro vs PHSD Summary

Thank you for your attention!

FIAS Frankfurt Institute for Advanced Studies

HGS-HIRe for FAIR Helmholtz Graduate School for Hadron and Ion Research

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DAAD

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