

Chiral magnetohydrodynamics & machine-learning CME

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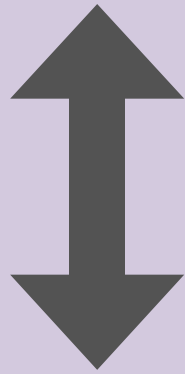


Chiral magnetic effect

$$\mathbf{j}_{\text{anom}} = \kappa_B \mathbf{B} + \kappa_\omega \boldsymbol{\omega}$$

$$\mathbf{j}_{5,\text{anom}} = \xi_B \mathbf{B} + \xi_\omega \boldsymbol{\omega}$$

Chiral Fluid



EM fields

Chiral MHD

MHD = Magnetohydrodynamics

Systems described by chiral MHD

- Heavy-ion collisions
 - For the CME search, reliable estimate of the lifetime of **B** is important
- Early Universe
- Weyl/Dirac semimetals

CME currents from magnetic reconnections

[Hirono-Kharzeev-Yin, PRL'16]

Magnetic & fermionic helicities

$$\partial_\mu \dot{j}_A^\mu = C_A \mathbf{E} \cdot \mathbf{B}$$



$$\frac{d}{dt} [\mathcal{H} + \mathcal{H}_F] = 0$$

$$\mathcal{H} = \int d^3x \mathbf{A} \cdot \mathbf{B}$$

Magnetic helicity

$$\mathcal{H}_F = \frac{2}{C_A} \int d^3x n_A$$

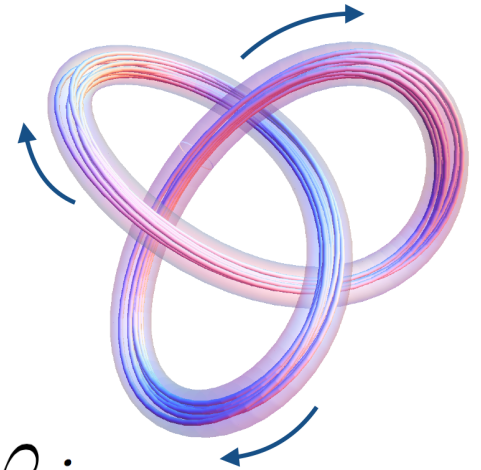
Fermionic helicity

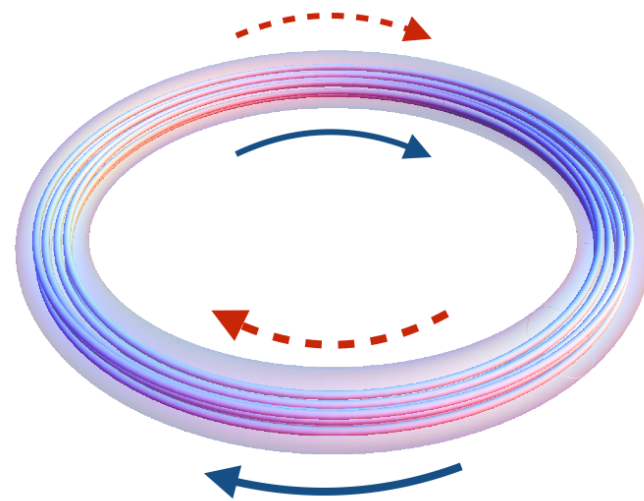
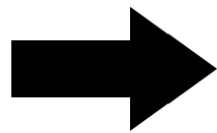
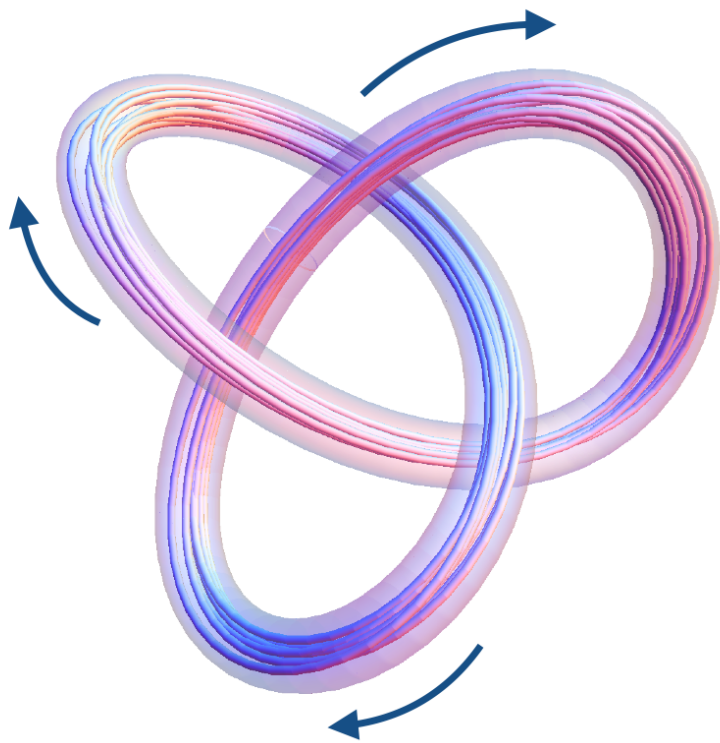
Magnetic helicity knows topology

$$\begin{aligned}\mathcal{H} &= \int d^3x \mathbf{A} \cdot \mathbf{B} \\ &= \sum_i \mathcal{S}_i \varphi_i^2 + 2 \sum_{i,j} \mathcal{L}_{ij} \varphi_i \varphi_j\end{aligned}$$

Self-linking number

Linking number





CME currents from reconnections of \mathbf{B}

[Hirono-Kharzeev-Yin PRL'16]

$$\sum_i \oint_{C_i} \Delta \mathbf{J} \cdot d\mathbf{x} = -\frac{e^3}{2\pi^2} \Delta \mathcal{H}$$

Change of topology induces CME currents!

Formulation & waves of chiral MHD

[Hattori-Hirono-Yee-Yin, in preparation]

Chiral MHD

- MHD & chiral MHD can be understood as a low-energy effective theory basing on derivative expansion
- A new anomaly-induced instability

EOM of MHD

$$\partial_{\mu} T_{\text{fluid}}^{\mu\nu} = F^{\nu\rho} j_{\rho}$$

$$\partial_{\mu} F^{\mu\nu} = j^{\nu}$$

$$\partial_{\mu} \tilde{F}^{\mu\nu} = 0$$

$$\tilde{F}^{\mu\nu} \equiv \epsilon^{\mu\nu\alpha\beta} F_{\alpha\beta}$$

EOM of MHD

$$\partial_{\mu} T_{\text{tot}}^{\mu\nu} = 0 \quad \partial_{\mu} \tilde{F}^{\mu\nu} = 0$$

$$T_{\text{tot}}^{\mu\nu} = T_{\text{fluid}}^{\mu\nu} + T_{\text{EM}}^{\mu\nu}$$

$$T_{\text{EM}}^{\mu\nu} = -F^{\mu}_{\alpha} F^{\nu\alpha} + \frac{1}{4} \eta^{\mu\nu} F_{\alpha\beta} F^{\alpha\beta}$$

EOM of MHD

$$E^\mu \equiv F^{\mu\nu} u_\nu, \quad B^\mu \equiv \tilde{F}^{\mu\nu} u_\nu$$

$$F^{\mu\nu} = E^\mu u^\nu - E^\nu u^\mu - \epsilon^{\mu\nu\rho} B_\rho$$

$$\tilde{F}^{\mu\nu} = B^\mu u^\nu - B^\nu u^\mu + \epsilon^{\mu\nu\rho} E_\rho$$

$$\epsilon^{\mu\nu\alpha} \equiv \epsilon^{\mu\nu\alpha\beta} u_\beta$$

Hydrodynamic variables

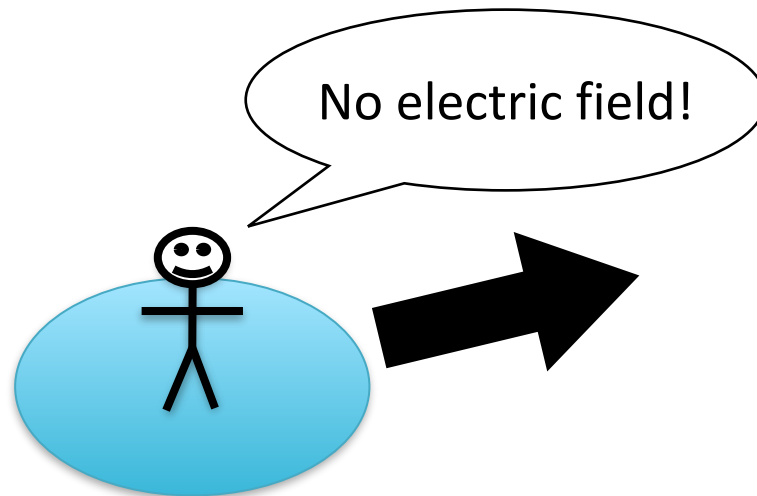
- Parameters characterizing local thermal equilibrium
- Neutral fluid with a conserved charge $\{T(\mathbf{x}), u^\mu(\mathbf{x}), \mu(\mathbf{x})\}$
- MHD $\{T(\mathbf{x}), u^\mu(\mathbf{x}), B^\mu(\mathbf{x})\}$

of hydro variables = # of equations = 7

No electric field in the fluid frame in ideal MHD

$$E_{(0)}^{\mu} = 0$$

Correspond to large conductivity limit



Constitutive relation for ideal MHD

$$T_{\text{tot}(0)}^{\mu\nu} = (e + p)u^\mu u^\nu - p\eta^{\mu\nu} + \mathbf{B}^2 \left[u^\mu u^\nu - b^\mu b^\nu - \frac{1}{2}\eta^{\mu\nu} \right]$$

$$B^\mu = |\mathbf{B}|b^\mu \quad b_\mu b^\mu = -1$$

$$F_{(0)}^{\mu\nu} = \epsilon^{\mu\nu\rho\sigma} u_\rho B_\sigma$$

CME doesn't play any role in ideal MHD!

Why?

Constitutive relation doesn't determine the current

$$j^\mu = \underbrace{\sigma}_{\infty} E^\mu + \sigma_B B^\mu$$

$\swarrow \quad \searrow$
 $\infty \quad 0$

Current is determined by the Maxwell equation

$$j^\mu = \partial_\nu F^{\mu\nu} \quad F^{\mu\nu} = \epsilon^{\mu\nu\rho\sigma} u_\rho B_\sigma$$

Chiral magnetic conductivity never appears in EOM

$$\partial_\mu T_{\text{tot}}^{\mu\nu} = 0 \quad \partial_\mu \tilde{F}^{\mu\nu} = 0$$

Conservation of topology of B

- Flux is “frozen in” to the fluid
- Magnetic helicity is conserved in ideal MHD
 - **No reconnection**

$$h_{\mathbf{B}}^{\mu} = \tilde{F}^{\mu\nu} A_{\nu} \quad : \text{ helicity current}$$

$$\partial_{\mu} h_{\mathbf{B}}^{\mu} = 2\tilde{F}^{\mu\nu} F_{\mu\nu} = 8E^{\mu} B_{\mu} = 0$$

$$\mathcal{H} = \int d^3x h_{\mathbf{B}}^0 \quad \text{is conserved}$$

First order in derivative expansion

Using the second law,

$$T_{\text{tot}(1)}^{\mu\nu} = \zeta \Delta^{\mu\nu} \partial \cdot u + 2\eta \nabla \langle{}^\mu u^\nu\rangle$$

E^μ : C-odd, P-odd

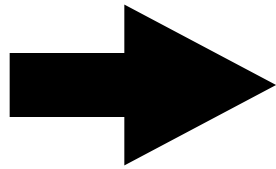
$$E_{(1)}^\mu = \frac{1}{\sigma\beta} \epsilon^{\mu\nu\alpha\beta} u_\nu \partial_\alpha (\beta B_\beta) - \underbrace{\epsilon_B B^\mu}_{\text{CME}}$$

σ : electric conductivity

$$\epsilon_B = \frac{\sigma_B}{\sigma} \quad \sigma_B : \text{chiral magnetic conductivity}$$

First order in derivative expansion

$$\tilde{F}_{(1)}^{\mu\nu} = \epsilon^{\mu\nu\rho\sigma} E_{(1)\rho} u_\sigma$$



$$\partial_\mu \left[T_{\text{tot}(0)}^{\mu\nu} + T_{\text{tot}(1)}^{\mu\nu} \right] = 0$$

$$\partial_\mu \left[\tilde{F}_{(0)}^{\mu\nu} + \tilde{F}_{(1)}^{\mu\nu} \right] = 0$$

Waves in chiral MHD

- Linear fluctuations - 6 modes in total

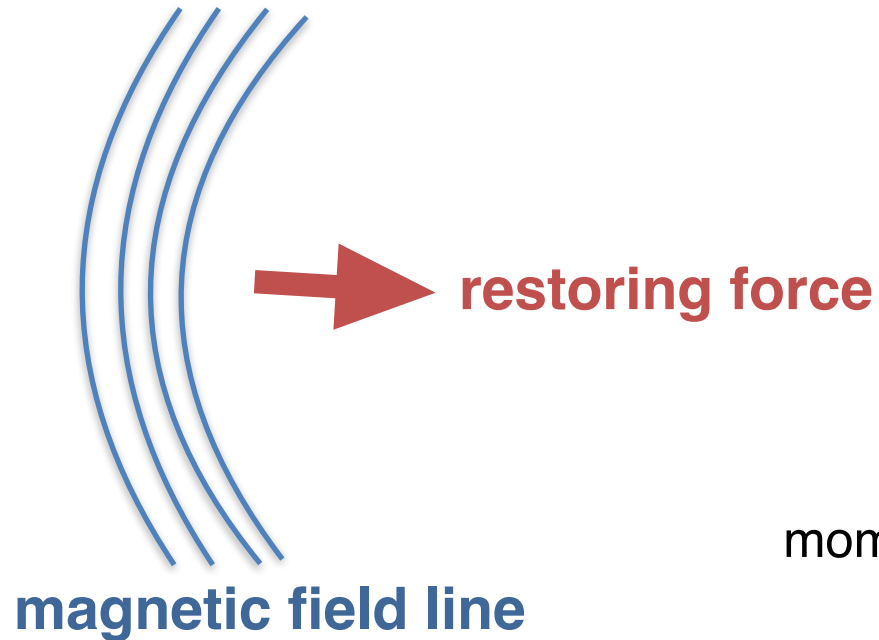
$$e \rightarrow e + \delta e,$$

$$B \rightarrow B + \delta B,$$

$$u^\mu \rightarrow u^\mu + \delta u^\mu,$$

$$b^\mu \rightarrow b^\mu + \delta b^\mu.$$

Alfven wave



momentum along the background \mathbf{B}
↓

Dispersion relation $\omega = \pm v_A k_{||}$

$$v_A^2 = \frac{B^2}{e + p + B^2} \quad \text{Alfven velocity}$$

Alfven wave in dissipative MHD

Dispersion relation

$$\omega = \pm v_A k_{||} - \frac{i}{2} [\bar{\eta} + \lambda] k^2$$

.....
damping

$$\lambda = \frac{1}{\sigma}$$

σ : electric conductivity

$$\bar{\eta} \equiv \frac{\eta}{e + p + \mathbf{B}^2}$$

η : shear viscosity

Alfven wave in chiral MHD

Including CME (when $\mathbf{k} \propto \mathbf{B}$)

$$\omega = \pm v_A k_{||} - \frac{i}{2} \left[(\bar{\eta} + \lambda) k_{||}^2 - \underbrace{s \epsilon_B k_{||}}_{\text{CME}} \right]$$

CME

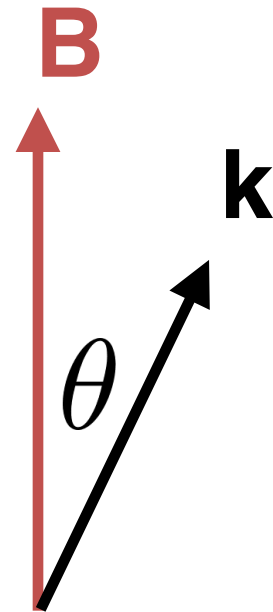
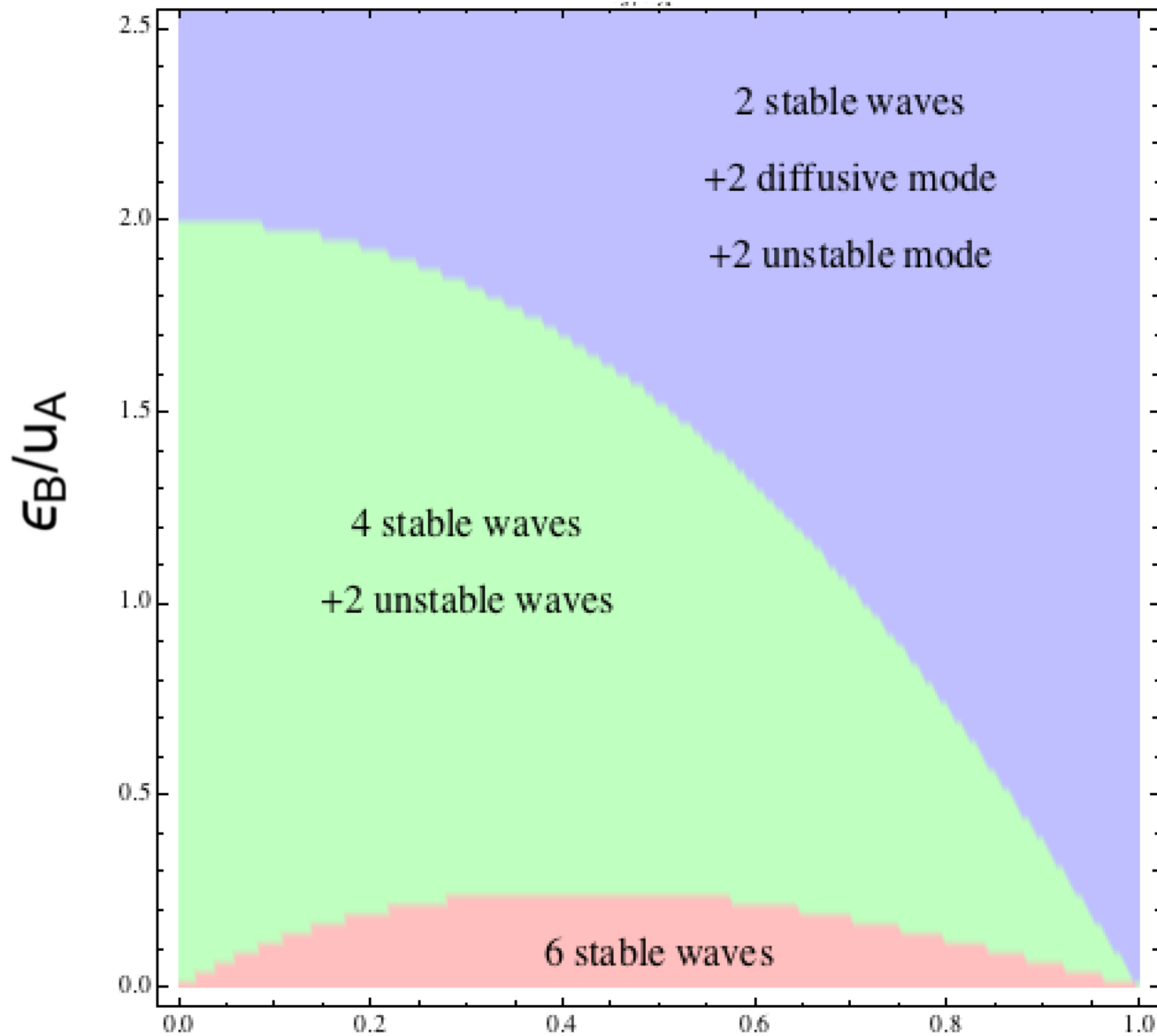
$s = \pm 1$ indicates the helicity of the mode

$$i\mathbf{k} \times \underbrace{\mathbf{e}^{(s)}}_{\text{helicity eigenstate}} = s k \mathbf{e}^{(s)}$$

helicity eigenstate

Instability in one of the helicity modes

Stability of the waves

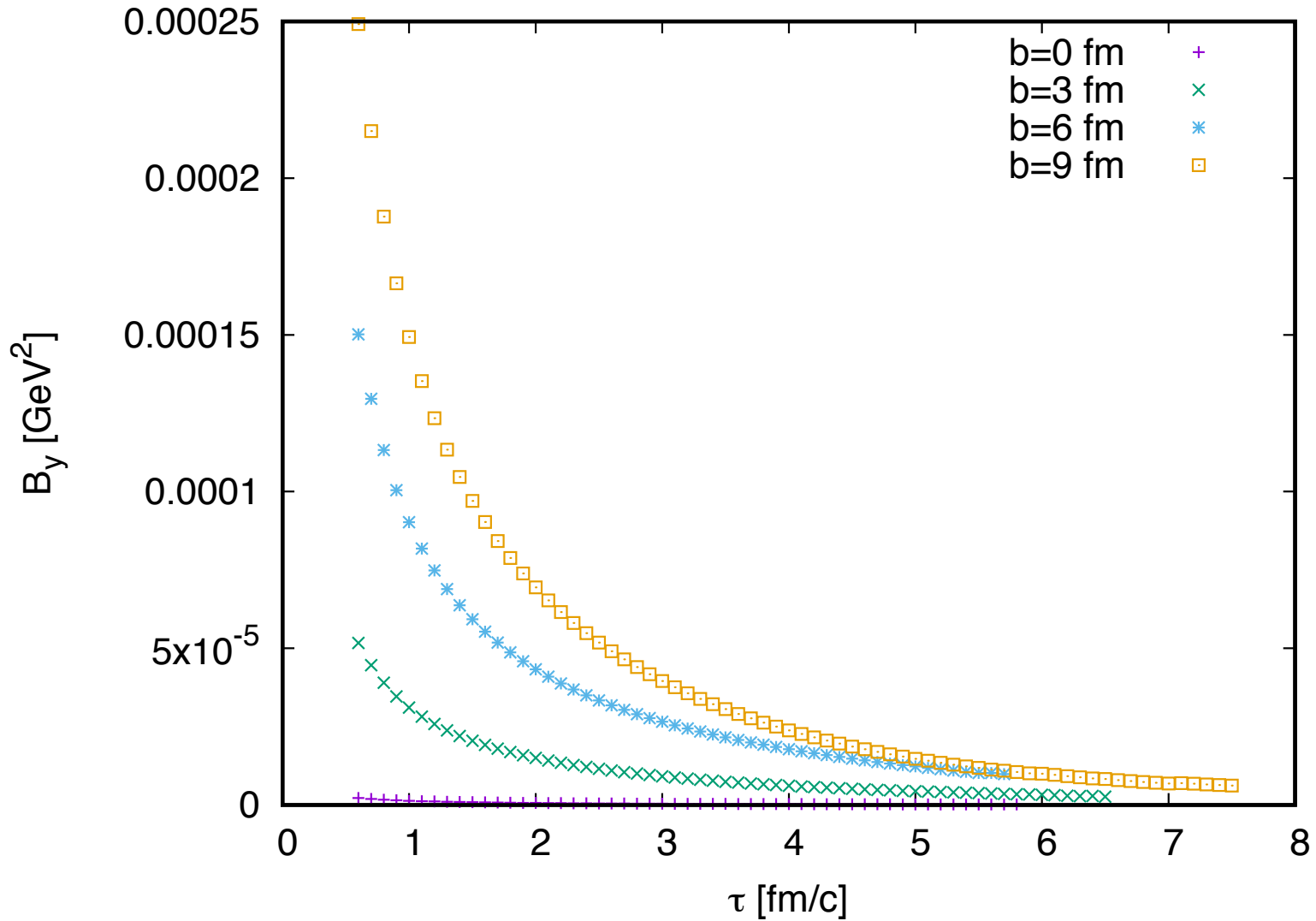


Toward numerical impl. of chiral MHD

with Mace, Kharzeev, Inghirami et.al.

- Collaborating with ECHO-QGP
- Implement:
 - Finite resistivity + anomalous effects
 - Pre-equilibrium CME currents

B-field evolution



Machine-learning CME

[Hirono-Kharzeev-Mace, in progress]

Observables for CME

- gamma correlation $\gamma_{\alpha\beta} = \left\langle \cos(\phi_1^\alpha + \phi_2^\beta - 2\Psi_{\text{RP}}) \right\rangle$
- Can pick up two-particle correlations unrelated to charge separation (“non-flow”)

$$f_{\alpha\beta}(\mathbf{p}_1, \mathbf{p}_2) = f_\alpha(\mathbf{p}_1) f_\beta(\mathbf{p}_2) + \underbrace{f_{\alpha\beta}^c(\mathbf{p}_1, \mathbf{p}_2)}_{\sim \frac{1}{N}}$$

Backgrounds

- Flowing resonances [Voloshin PRC'04]
[F. Wang PRC'10]
- Local charge conservation + flow
[Schlichting-Pratt, PRC'11]
[Pratt-Schlichting-Gavin PRC'11]
- Transverse momentum conservation
[Bzdak-Koch-Liao PRC'10, PRC'11]
[Pratt-Schlichting-Gavin PRC'11]

Those all contribute as $\sim \frac{v_2}{N}$

Machine-learning CME

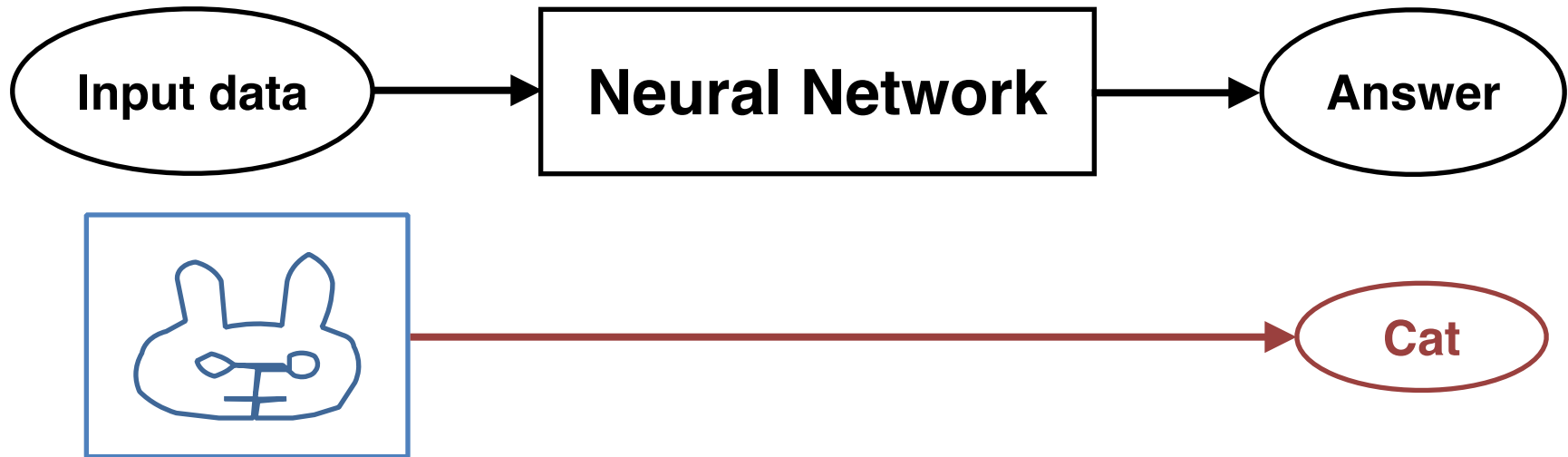
[Hirono-Kharzeev-Mace, in progress]

- How to separate of signal & background is a big problem
- Gammas - include contributions from background effects
- There can be better observables
 - Less sensitive to background effects

Machine-learning CME

[Hirono-Kharzeev-Mace, in progress]

- Use **neural network** to learn the feature of CME
 - Training (parameter tuning)



- Applied for detection of first-order PT

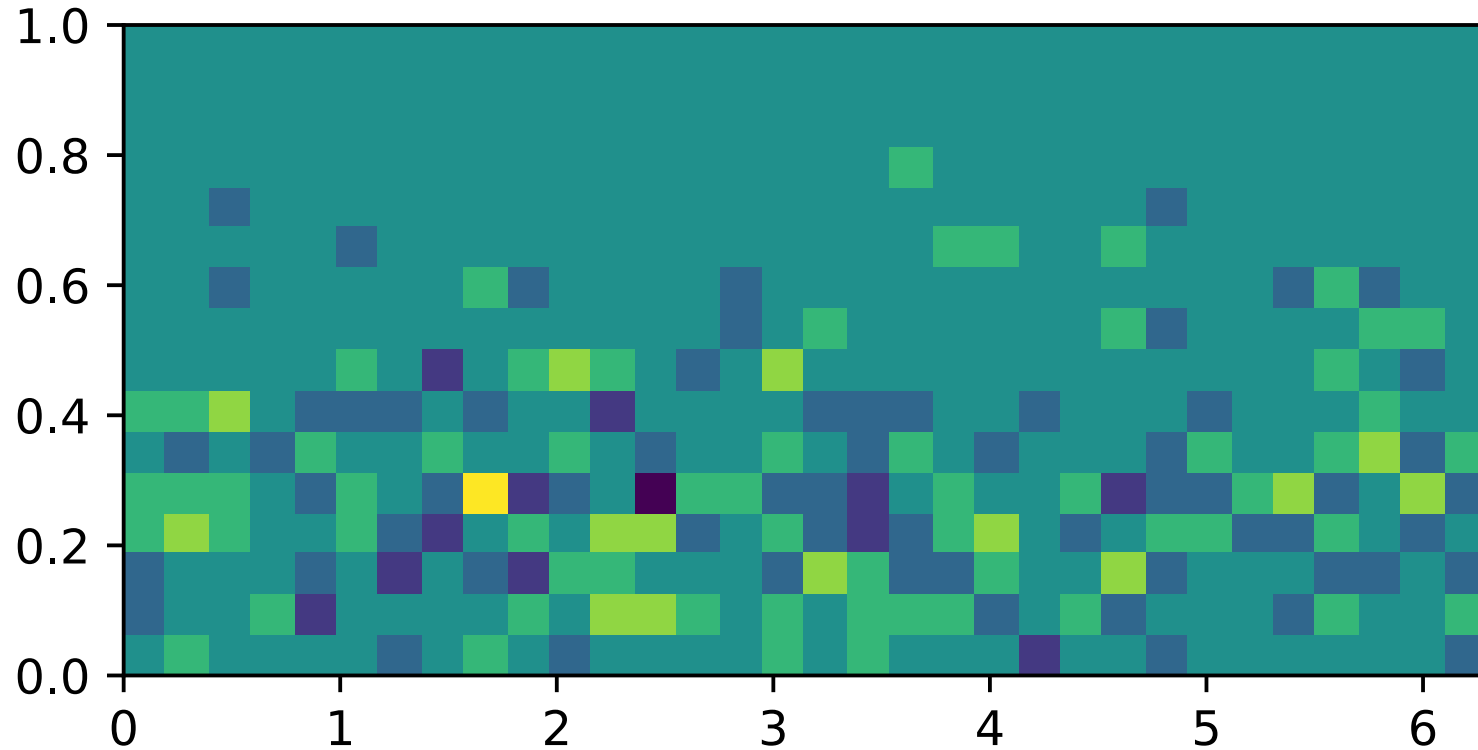
[Pang-Zhou-Su-Petersen-Stöcker-Wang 1612.04262]

Data preparation & learning

- Supervised learning
- Raw event data (momenta of each particle)
 - Fireball
 - With charge separation (about 1% in a1)
 - Without charge separation
 - Produce particle via Cooper-Frye prescription
 - Perform resonance decay
- Charge-dependent distributions
 - One-particle distributions
 - Multi-particle distributions
- Use Keras + Tensorflow

$$N^+(p_T, \phi) - N^-(p_T, \phi)$$

p_T [GeV]



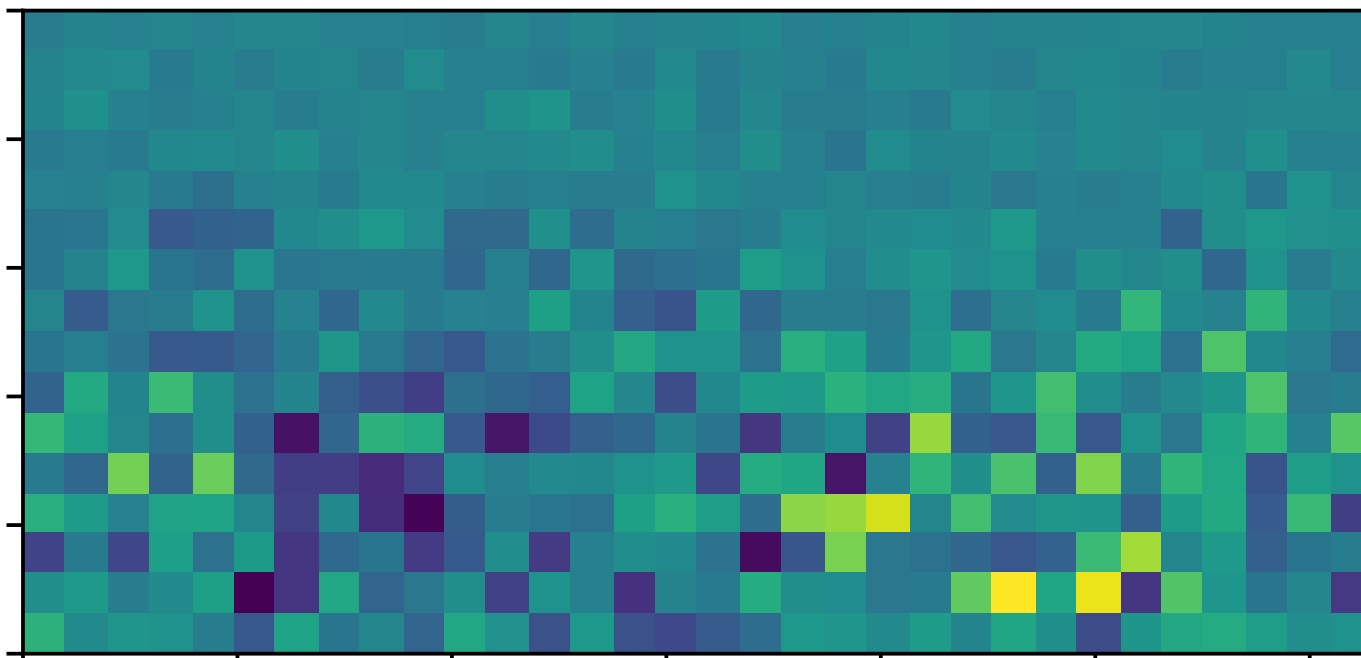
- One event

$\phi - \psi_{RP}$

$$N^+(p_T, \phi) - N^-(p_T, \phi)$$

p_T [GeV]

1.0
0.8
0.6
0.4
0.2
0.0



0.06

0.04

0.02

0.00

-0.02

-0.04

- Averaged over 2000 events
with fixed direction of charge separation

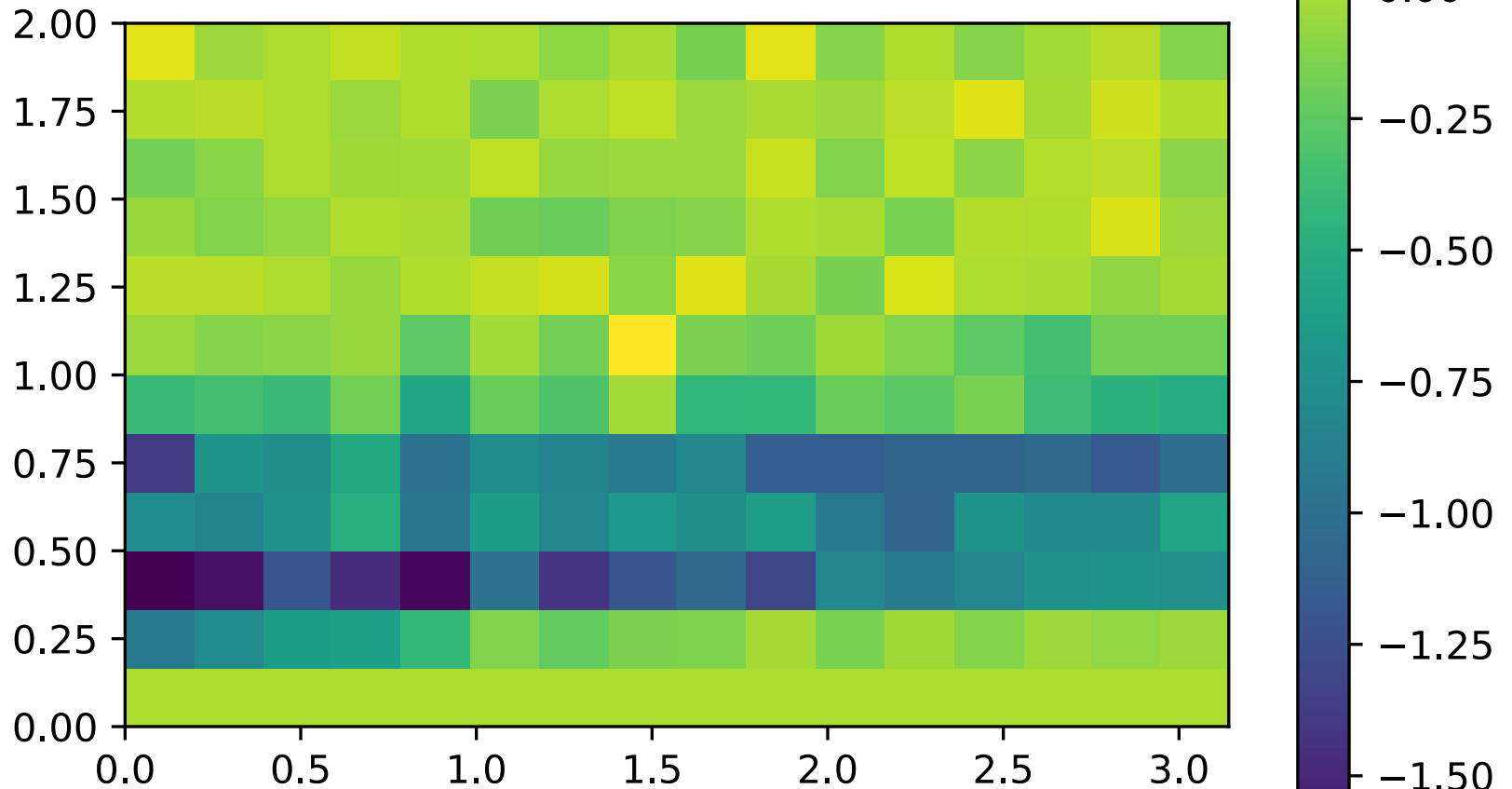
$\phi - \psi_{RP}$

Machine-learning CME

- Machines can detect CME well for the single-particle net-charge distribution (averaged over a certain number of events), **when the direction of charge separation is fixed**
- We are testing this method for two-particle charge-dependent corr.

$$N^{++} + N^{--} - N^{+-} - N^{-+}$$

m_{inv} [GeV]



$\Delta\phi$

Summary

1. B-field topology & CME currents

- Reconnections generate CME currents

2. Chiral MHD

- CME appears in the first order P-odd correction
- Helicity dependent instability
- Numerical implementation is under way

3. Machine learning CME

- Use of neural network for a better observable