

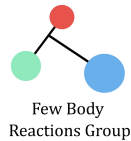
BAND Approach to Optical Potentials

Kyle Beyer

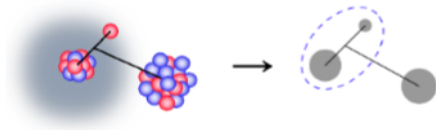
CSWEG 2025



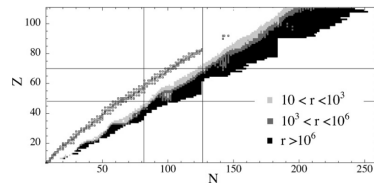
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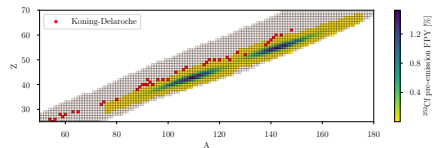
Nuclear reactions away from stability involves extrapolating the optical potential



- optical potential reduce many-body DOF in reactions to few
- phenomenological optical potentials are extrapolated away from β -stability
- along with NLDs, γ SF, etc., they are necessary for prediction and evaluation in data-sparse regime



(n, γ) rates for r-process ¹

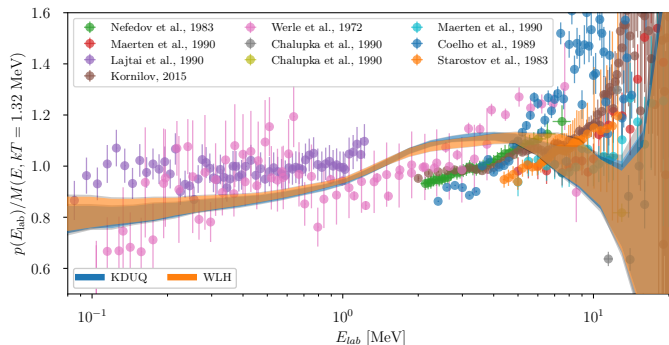


fission fragments ²

¹Goriely and Delaroche 2007, *Physics Letters B* 653(2-4), pp. 178–183

²Beyer et al. 2025, *Phys. Rev. C* 112, p. 024604

Past work: optical model uncertainty in CGMF fission observables



Optical model uncertainty in
CGMF Monte Carlo
Hauser-Feshbach calculations of
 ^{252}Cf PFNS³

³Beyer et al. 2025, *Phys. Rev. C* 112, p. 024604

Calibration of an uncertainty-quantified global optical potential

- Canonical KOH statistical approach to model calibration ⁴

$$\underbrace{y_i + \epsilon_i}_{\text{observation with uncertainty}} = \underbrace{\zeta(x_i)}_{\text{latent truth}} = \underbrace{\eta(x_i; \alpha) + \delta(x_i)}_{\text{parametric model with discrepancy}} + \epsilon_i \quad (1)$$

- given priors

$$\begin{aligned}\epsilon_i &\sim \mathcal{N}(0, \Sigma^{\text{exp}}) \\ \delta(x_i) &\sim \mathcal{N}(0, \Sigma^{\text{md}})\end{aligned}$$

- and data set(s) $\mathcal{D} = \{x_i, y_i\}_{i=0}^N$
- the likelihood is

$$p(\alpha, \beta | \mathcal{D}) = \left((2\pi)^N |\mathbf{V}(\beta)| \right)^{-1/2} \exp \left\{ -\frac{1}{2} \mathbf{\Delta}^\top \mathbf{V}^{-1}(\beta) \mathbf{\Delta} \right\},$$

where $\mathbf{\Delta} \equiv \mathbf{y} - \boldsymbol{\eta}(\mathbf{x}, \alpha)$ and $\mathbf{V}(\beta) \equiv \Sigma^{\text{exp}} + \Sigma^{\text{md}}(\beta)$

- Baye's rule provides posterior given priors $p(\alpha)$, $p(\beta)$
- for global nucleon-nucleus optical potentials, $x = \{A, Z, E, \theta, \dots\}$ and $y = d\sigma/d\Omega, A_y, \dots$

⁴Kennedy and O'Hagan 2001, *Journal of the Royal Statistical Society: Series B (Statistical Methodology)* 63(3), pp. 425–464

Statistical approach to optical potentials

- **1950s on:** Frequentist, least-squares fitting, a few parameters at a time. Parameter ambiguities observed. ^a
- **2003:** Koning-Delaroche ^b is state-of-the-art: 50 parameter fit with computational steering, local, global hierarchy
- **2017:** Bayesian calibration and UQ with local potentials (Lovell et al. ^c)
- **2023:** First global Bayesian calibration: CHUQ and KDUQ (Pruitt et al. ^d)

^aSatchler et al. 1964, *Physical Review* 136(3B), B637

^bKoning and Delaroche 2003, *Nuclear Physics A* 713(3-4), pp. 231–310

^cLovell et al. 2017, *Physical Review C* 95(2), p. 024611

^dPruitt et al. 2023, *Physical Review C* 107(1), p. 014602

Progress on uncertainty quantification and extrapolation of the optical potential

Statistical approach to optical potentials

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- systematic errors ignored!
- model discrepancy ignored (or treated with diagonal covariance)!
- non-identifiability not treated rigorously!

Can we trust the extrapolation into data-sparse regime?

Extrapolation off the line of stability constrained by charge-symmetry of the nuclear force

- Reasonable ansatz for isovector dependence of optical potential ⁵:

$$\begin{aligned} U(r) &= U_0(r) + 4 \frac{\boldsymbol{\tau} \cdot \mathbf{T}}{A} U_1(r) \\ &= U_0(r) + \underbrace{4 \frac{\tau_3 T_3}{A} U_1(r)}_{(N-Z)/A \text{ dependence in elastic scattering}} + \underbrace{4 \frac{\tau_+ T_- + \tau_- T_+}{2A} U_1(r)}_{(p, n) \text{ to isobaric analog}} \end{aligned}$$

Goal: uncertainty-quantified extrapolation in $(N-Z)/A$ using $(p, n)_{\text{IAS}}$ as a constraint and rigorous Bayesian approach

⁵Lane 1962, *Nuclear Physics* 35, pp. 676–685

(p, n) to the isobaric analog state

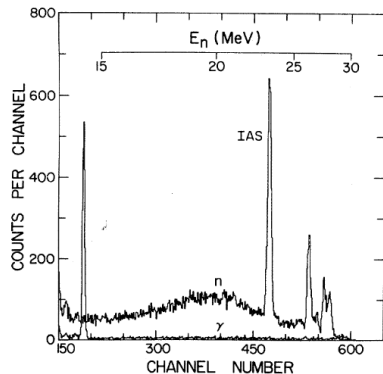
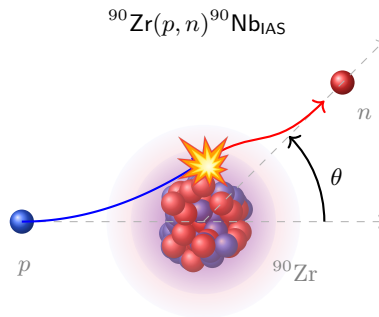


FIG. 1. Neutron (n) and γ -ray (γ) time-of-flight spectra at 0° for 35-MeV protons on ^{90}Zr . Each channel represents 0.1 nsec.

From measurements at MSU Cyclotron Lab ^a

^aDoering et al. 1975, *Physical Review C* 12(2), p. 378



- good single-particle description: replacement of valence neutron orbital with incident proton
- probe of N, Z symmetry energy and neutron skins

The East Lansing Model

$$U(r; E, A, Z) = U_0(r; E, A) + \left(\frac{N - Z}{N + Z} \right) U_1(r; E, A)$$

- New global, uncertainty quantified optical potential
- Lane consistent
- 24 parameters (compared to 48 for Koning-Delaroche)
- Coulomb correction to all orders
- open-source data curation (`exfor-tools`), solver (`jitR`), and statistical modeling software (`rxmc`) as part BAND framework⁶



visitlearn.msu.edu/resources/facility-for-rare-isotope-beams
photos.msu.edu

github.com/beykyle/elm

⁶Beyer et al. 2025, (Version 0.5.0)

jit $\mathcal{R}$ allows for global calibration without needing emulators

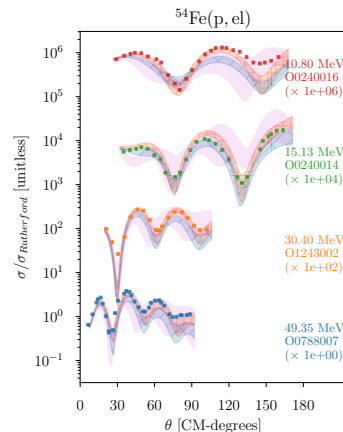
philosophy:

Modular library for building workspaces which take in interaction models and spit out reaction observables

- calculable $\mathcal{R}$ -matrix on Lagrange mesh ^a
- modular python library
- precompute everything you can
- numba just-in-time compilation, heavy numpy vectorization
- **2 orders of magnitude faster than comparable fortran codes**, just as fast as emulators ^b in UQ setting

^aBaye 2015, *Physics reports* 565, pp. 1–107

^bOdell et al. 2024, *Physical Review C* 109(4), p. 044612



Open source, part of BAND framework ^a

^aBeyer et al. 2025, (Version 0.5.0)

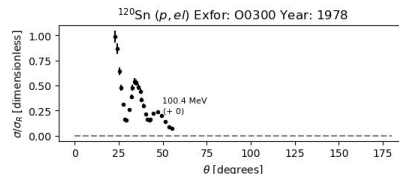
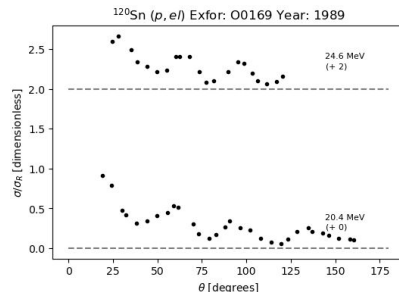


Curation of a body of experimental data

- ~ 100 EXFOR entries, $\sim 10K$ data points
- sparse coverage of input space
- $40 \leq A \leq 208$
- $10 \text{ MeV} < E_{\text{lab}} < 200 \text{ MeV}$
- $d\sigma/d\Omega$ and A_y from elastic (p, p) , (n, n)
- $d\sigma/d\Omega$ with clean $\Delta J^\pi = 0^+$ from $(p, n)_{\text{IAS}}$
- manual ID & cleaning of outliers
- types of error reported inconsistent

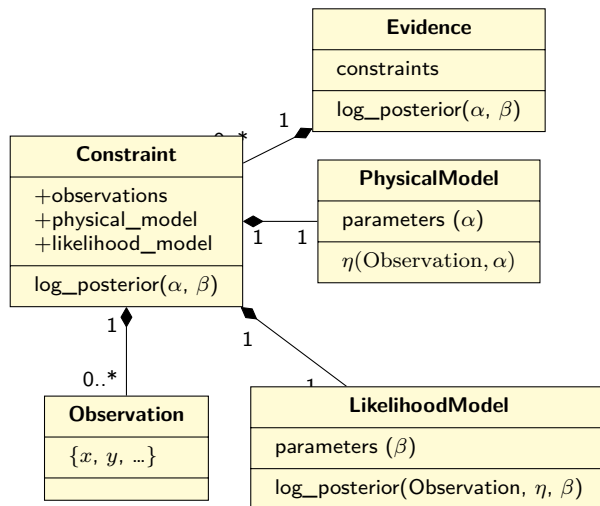
github.com/beykyle/exfor_tools, built on top of David Brown's x4i package.

```
all_entries_pp = exfor_tools.get_exfor_differential_data(  
    target=target,  
    projectile=proton,  
    quantity="dXS/dA",  
    product="EL",  
    energy_range=[10, 200], # MeV  
)
```



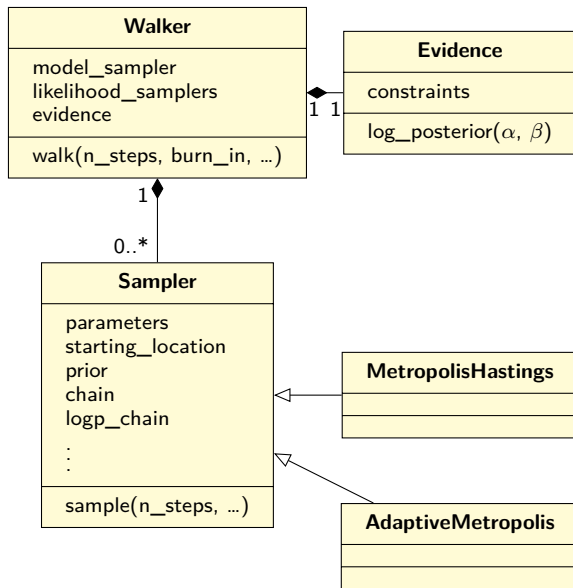
rxmc: flexible, hierarchical Bayesian calibration to heterogeneous data

- curating a body of evidence from multiple constraints
- Bayesian calibration with flexible likelihood models
- designed with `jitR` and `exfor-tools` in mind
- built-in sampler, or interface with external MCMC package
- handles parametric likelihood models with batched Metropolis-in-Gibbs sampling
- in alpha stage:
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Likelihood model

$$\ln p(\mathcal{D}|\alpha, \mathcal{M}) = -\frac{1}{2} \left(\mathbf{\Delta}^\top \cdot \mathbf{V}^{-1} \cdot \mathbf{\Delta} + \ln \left((2\pi)^k |\mathbf{V}| \right) \right)$$

$$\Delta_i(\alpha) \equiv \eta(x_i, \alpha) - y_i$$

$$\mathbf{V}(\beta) = \mathbf{\Sigma}_{ij}^{\text{exp}} + \mathbf{\Sigma}_{ij}^{\text{MD}}(\beta)$$

- diagonal model discrepancy, learned from data (following ⁷).

$$\mathbf{\Sigma}_{ij}^{\text{exp}} \equiv \delta_{ij} \sigma_{\text{exp},i}^2$$

$$\mathbf{\Sigma}_{ij}^{\text{MD}}(\beta) \equiv \delta_{ij} (\beta \tilde{y}_i)^2, \quad \tilde{y}_i \equiv (\eta(x_i; \alpha) + y_{i,\text{exp}}) / 2$$

- β are free parameters (one per data sector)

Prior

- estimated from CHUQ

Sampling

- Adaptive Metropolis-in-Gibbs with rxmc
- Goodman-Weare from emcee⁸
- massively parallel

⁷Pruitt et al. 2023, *Physical Review C* 107(1), p. 014602

⁸Foreman-Mackey et al. 2013, *PASP* 125(925), p. 306; Goodman and Weare 2010, *Communications in applied mathematics and computational science* 5(1), pp. 65–80

Likelihood model

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- β are free parameters (one per data sector)

To re-scale or not to re-scale the log-likelihood?

- Pruitt et al. re-scaled the log likelihood by k/N (# of free parameters / # of experimental data points $\sim 20/10000$):

$$\ln p(\mathcal{D}|\alpha, \mathcal{M}) \rightarrow$$

$$\frac{k}{N} \ln p(\mathcal{D}|\alpha, \mathcal{M})$$

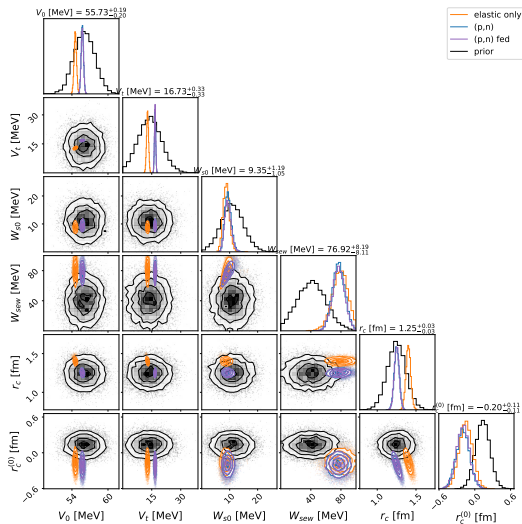
- re-scaling is equivalent to “cold” posterior with $\lambda = k/N < 1$ (post-Bayesian approach) ⁸

$$p_\lambda(\alpha|\mathcal{D}, \mathcal{M}) \propto p(\mathcal{D}|\alpha, \mathcal{M})^\lambda p(\alpha)$$

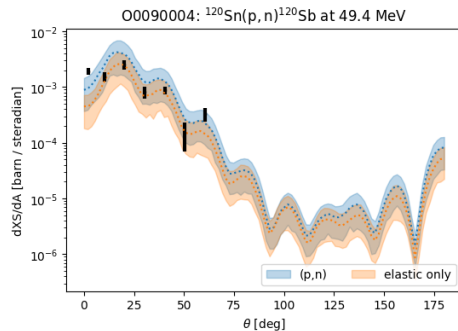
⁷Pruitt et al. 2023, *Physical Review C* 107(1), p. 014602

⁸Grünwald 2011, vol. 19, pp. 397–420

Effect of (p, n)

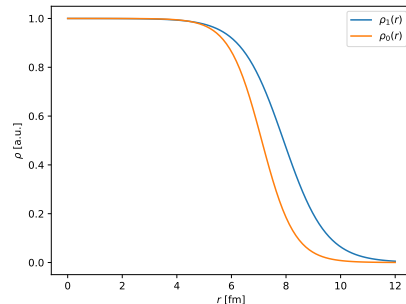
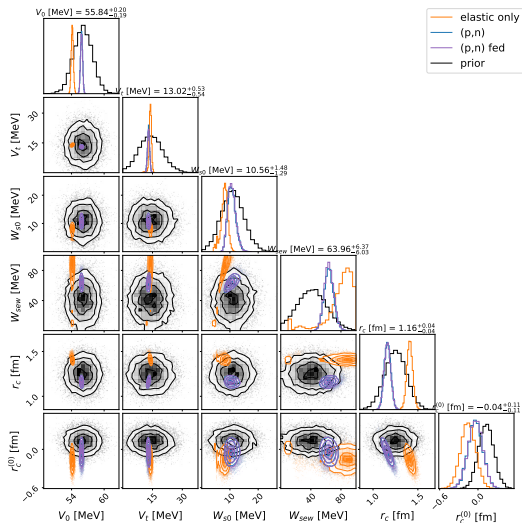


- Introduction of (p, n) constraint doesn't make a huge difference ...



Corner plot for subset of posterior with and without (p, n)

Effect of independent isovector geometry

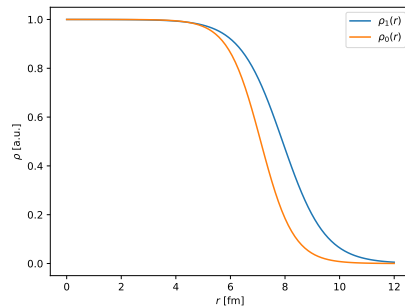
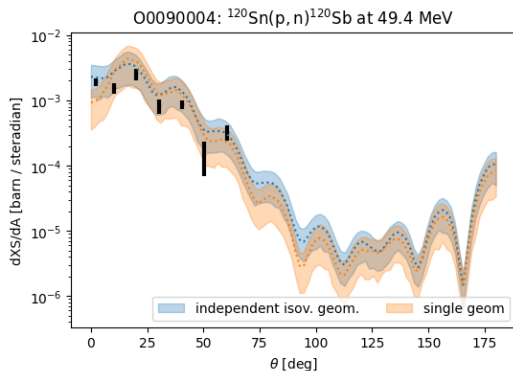


- Independent isoscalar and isovector geometries suggested by Danieliwick et al. ^a
- Convergence only reached when (p, n) constraint is included

Corner plot for subset of posterior with and without (p, n)

^aDanielewicz et al. 2017, *Nuclear Physics A* 958, pp. 147–186

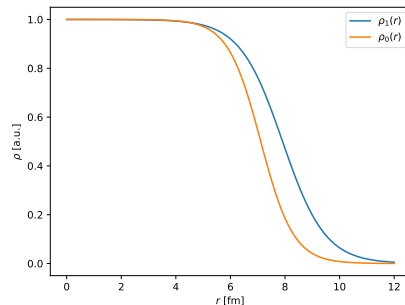
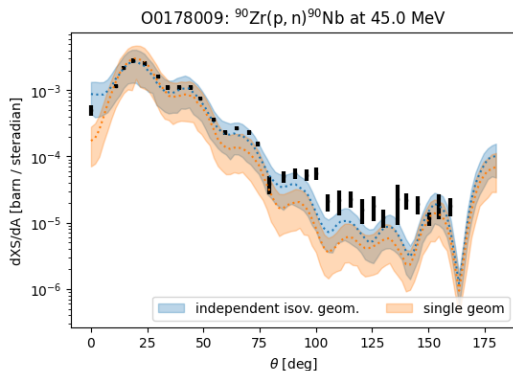
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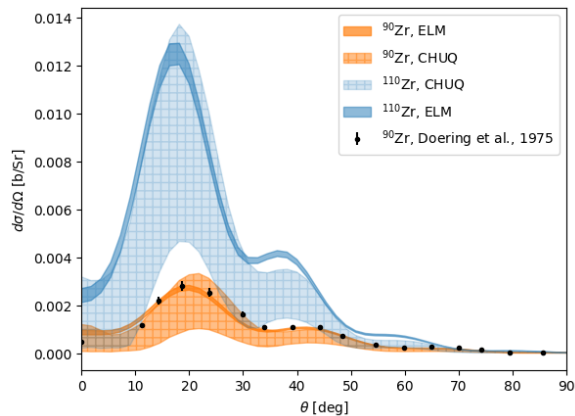
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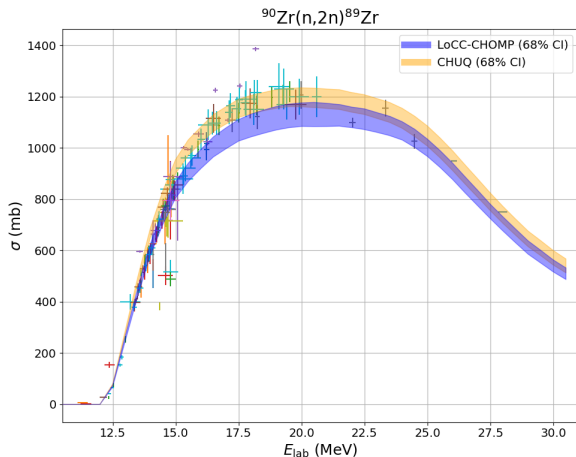
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Extrapolation in $(N - Z)/A$ with reduced uncertainties compared to CHUQ



$(p, n)_{IAS}$ for Zr isotopes

We've begun propagating uncertainties to Hauser-Feshbach calculations



BAND Fellow, Samuel Sullivan
(PhD student at University of Surrey)

- Lane-consistent potential calibrated to $d\sigma_{el}/d\Omega$ and A_y for (n,n) , (p,p) along isotopic chains from 10-50 MeV using `rxmc`
- propagation to compound channels using TALYS
- covariances can be obtained empirically from posterior predictive

ELM coming soon!

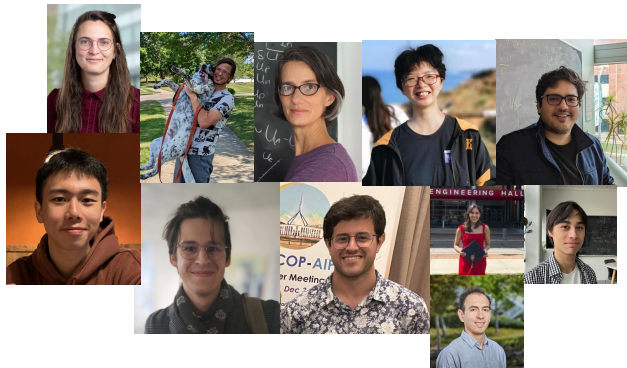
- need both (p, n) constraints and independent isoscalar and isovector geometries
- developed open-source software ecosystem for Bayesian inference in nuclear reactions

Future work:

- systematic experimental error, model discrepancy, errors in angle
- handle non-identifiability of spin-orbit potential, geometric parameters
- dispersion relation and structure constraints
- UQ of nuclear level density

Thank you for your time!

In collaboration with Filomena Nunes and the Few Body Reactions Group



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Few Body
Reactions Group

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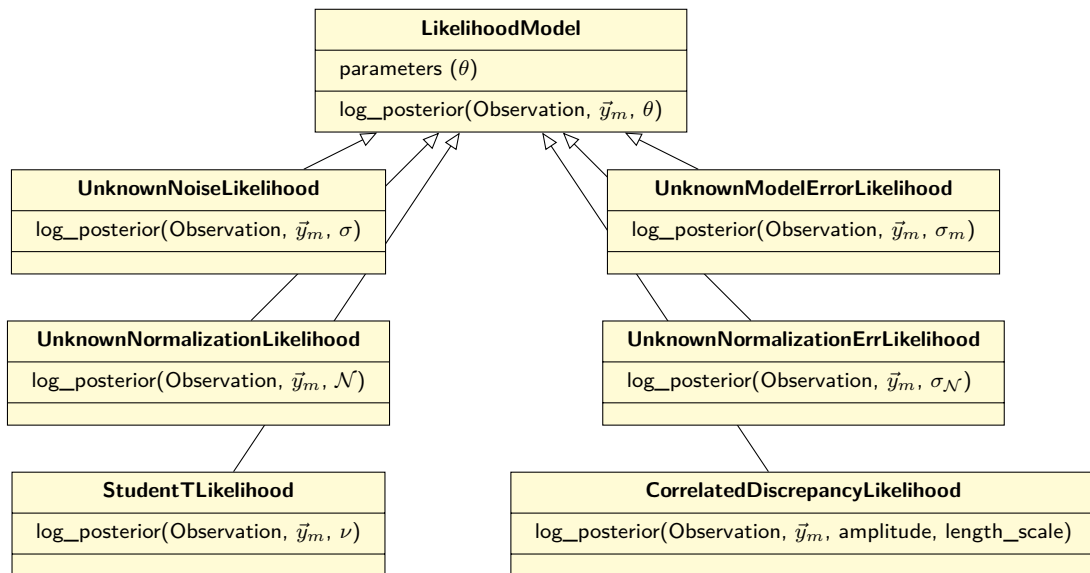
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-  Doering, RR et al. (1975). "Microscopic description of isobaric-analog-state transitions induced by 25-, 35-, and 45-MeV protons". In: *Physical Review C* 12.2, p. 378.
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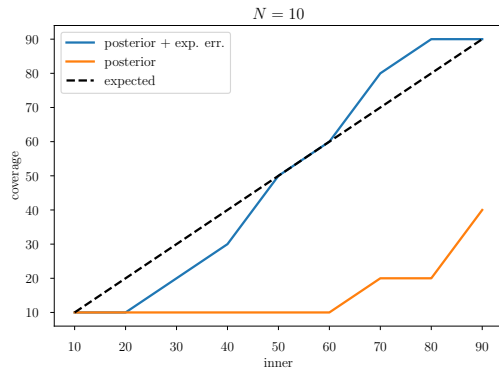
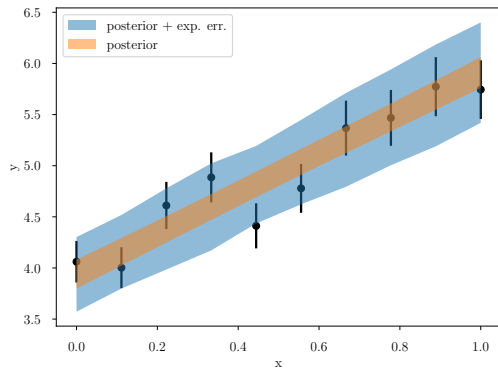
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-  Satchler, GR et al. (1964). "Optical-Model Analysis of" Quasielastic"(p, n) Reactions". In: *Physical Review* 136.3B, B637.
-  Schery, SD et al. (1974). "The (p, n) reaction to the isobaric analogue state of high-Z elements at 25.8 MeV". In: *Nuclear Physics A* 234.1, pp. 109–129.

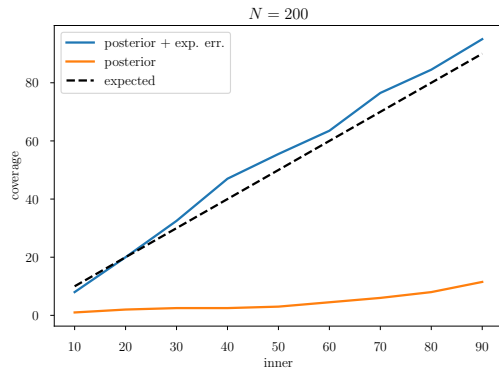
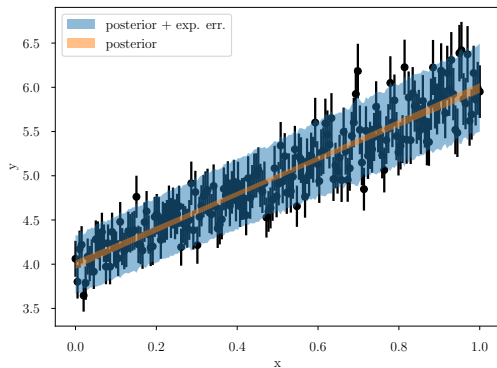
rxmc: building and composing likelihood models



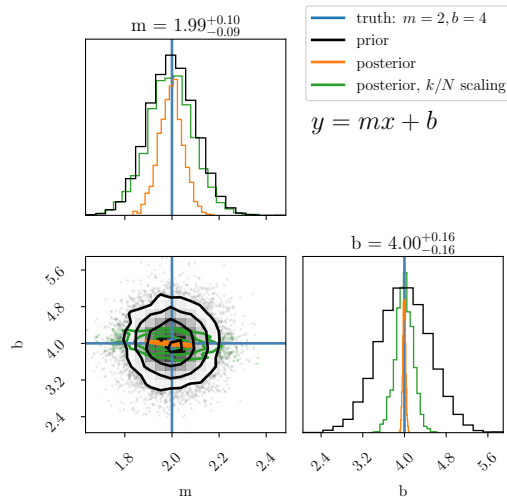
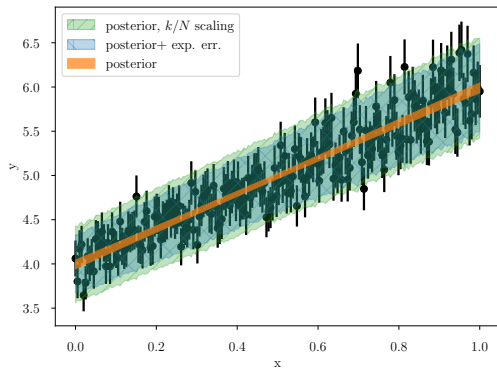
Empirical coverage



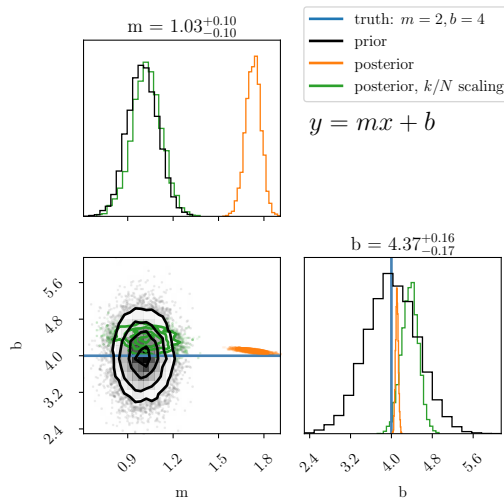
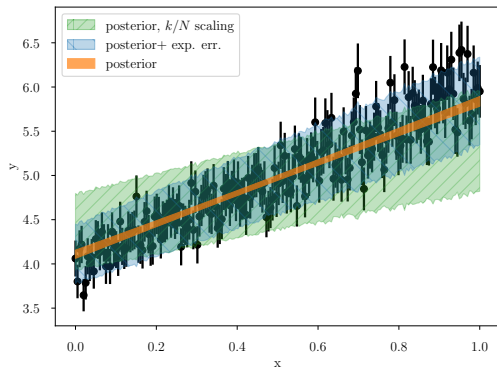
Empirical coverage



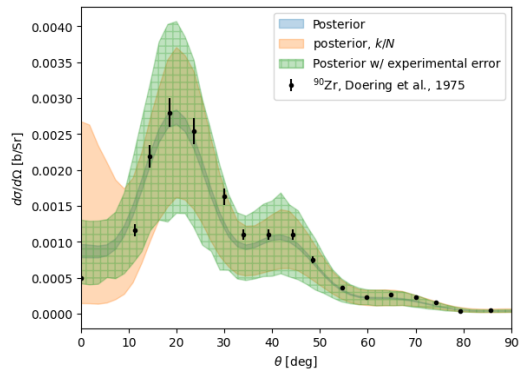
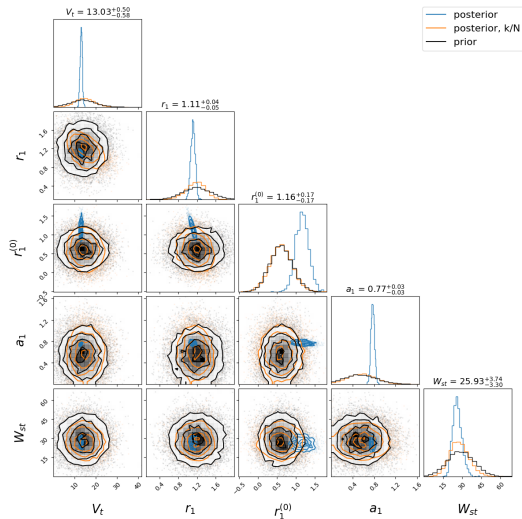
“Cold” posterior



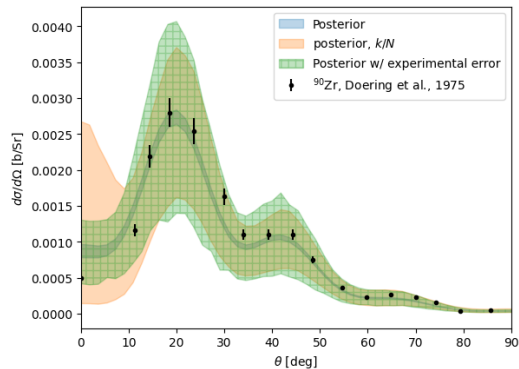
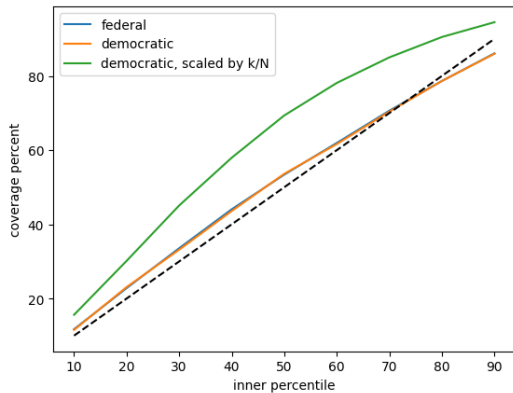
“Cold” posterior: over-emphasis on the prior



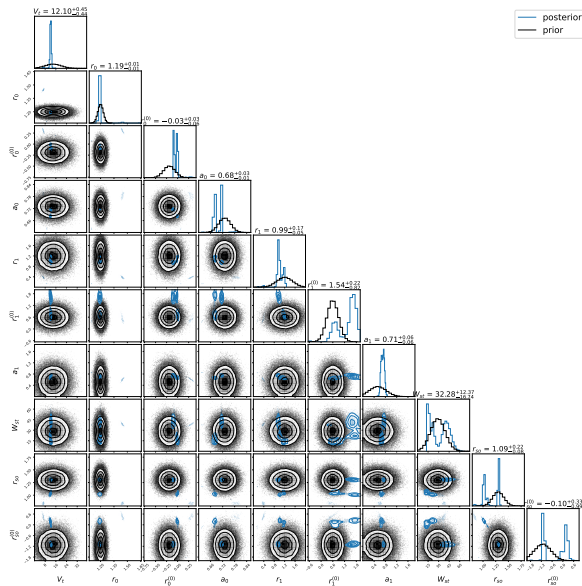
“Cold” posterior: over-emphasis on the prior



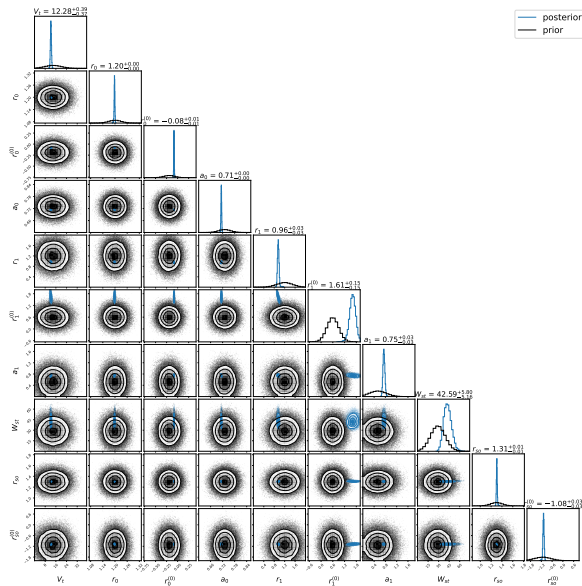
Predictive posteriors and empirical coverage



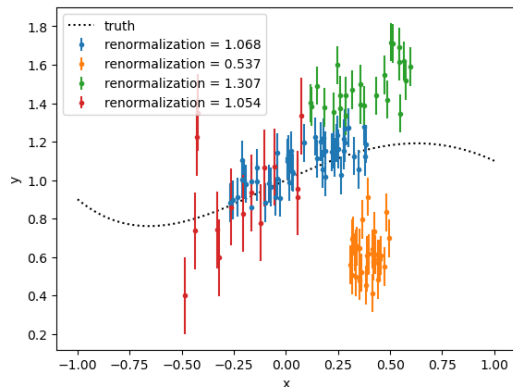
Multimodality, parameter ambiguity in radii



Multimodality, parameter ambiguity in radii



Handling multiple constraints with systematic errors



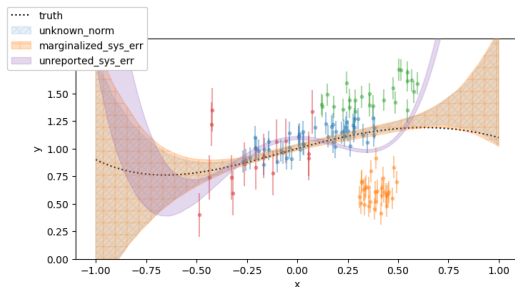
Each synthetic data set j has it's own systematic error σ_j

- treat data normalization as free parameter:
 $y_{i \in \mathcal{D}_j} = \zeta(x_i) + \epsilon_i = \rho_j \eta(x_i; \alpha) + \delta(x_i)$
- explicitly marginalize over $\rho_j \sim \mathcal{N}(1, \sigma_j)$,
leads to *model-dependent* covariance matrix:
 $\Sigma_{lm} = \delta_{lm} \sigma_{stat,l}^2 + \sigma_j^2 \eta(x_l; \alpha) \eta(x_m; \alpha)$.
Note: doing this wrong leads to Peelle's Pertinent Puzzle ⁹!

Adapted from open source demos in <https://github.com/beykyle/rxmc>

⁹Barlow 2021, *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment* 987, p. 164864; D'Agostini 1994, *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment* 346(1-2), pp. 306–311; Neudecker et al. 2012, *Nuclear science and engineering* 170(1), pp. 54–60

Handling multiple constraints with systematic errors



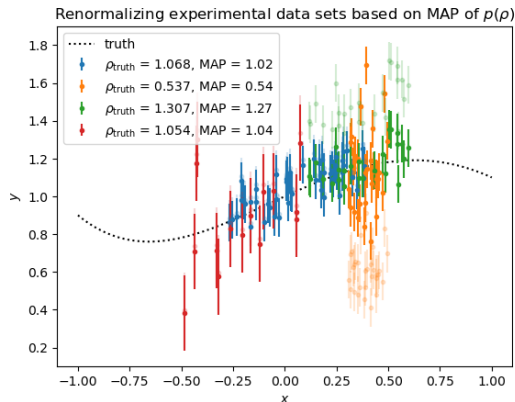
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- compare to ignoring systemtic error

Adapted from open source demos in <https://github.com/beykyle/rxmc>

⁹Barlow 2021, *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment* 987, p. 164864; D'Agostini 1994, *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment* 346(1-2), pp. 306–311; Neudecker et al. 2012, *Nuclear science and engineering* 170(1), pp. 54–60

Handling multiple constraints with systematic errors



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 $\Sigma_{lm} = \delta_{lm} \sigma_{\text{stat}, l}^2 + \sigma_j^2 \eta(x_l; \alpha) \eta(x_m; \alpha)$.
Note: doing this wrong leads to Peelle's Pertinent Puzzle ⁹!
- advantage of keeping ρ_j free is the ability to renormalize experimental data to be consistent with each other

Adapted from open source demos in <https://github.com/beykyle/rxmc>

⁹Barlow 2021, *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment* 987, p. 164864; D'Agostini 1994, *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment* 346(1-2), pp. 306–311; Neudecker et al. 2012, *Nuclear science and engineering* 170(1), pp. 54–60

- statistical only:

$$\Sigma_{ij} = \delta_{ij} \sigma_i^2$$

- fit (fractional) statistical error:

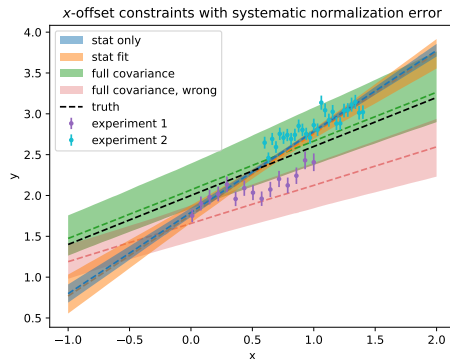
$$\Sigma_{ij} = \delta_{ij} (\eta y_{\mathcal{M}}(x_i; \alpha))^2$$

- full covariance:

$$\Sigma_{ij} = \delta_{ij} \sigma_i^2 + \sigma_N^2 y_{\mathcal{M}}(x_i; \alpha) y_{\mathcal{M}}(x_j; \alpha)$$

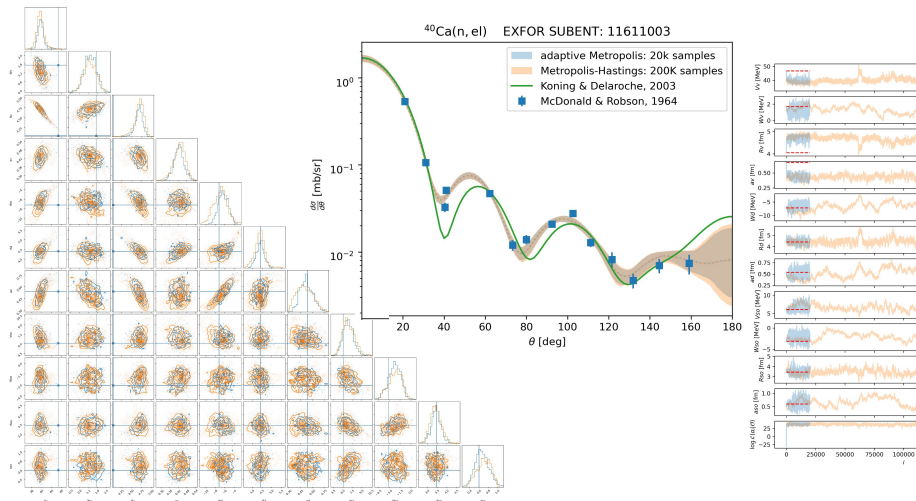
- full covariance, wrong:

$$\Sigma_{ij} = \delta_{ij} \sigma_i^2 + \sigma_N^2 y_i y_j$$



https://github.com/beykyle/rxmc/blob/main/examples/systematic_err_demo.ipynb

rxmc demos and tutorials: comparison of sampling algorithms for fitting an optical potential



https://github.com/beykyle/rxmc/blob/main/examples/fitting_an_optical_potential.ipynb

$$U(r, E, A, Z) = U_0(r, E, A) \pm \frac{N - Z}{N + Z} U_1(r, E, A) + V_C$$

$$U_0(r, E, A) = V_0(E) f((r - R_0(A))/a_0) + iW_0(E) f((r - R_w(A))/a_w) \\ - 4ia_w W_{s,0}(E) \frac{d}{dr} f((r - R_w(A))/a_w) + 2(\boldsymbol{\sigma} \cdot \boldsymbol{\ell}) V_{so} \frac{1}{r} \frac{d}{dr} f((r - R_{so}(A))/a_{so})$$

$$U_1(r, E, A) = V_1(E) f((r - R_1(A))/a_1) - 4ia_w W_{s,1}(E) \frac{d}{dr} f((r - R_w(A))/a_w)$$

$$f(x) = \frac{1}{1 + \exp(x)}$$

$$R_i(A) = r_{i,0} + r_{i,A} A^{1/3}$$

$$E_C = 6Z\alpha\hbar c/5R_C$$

$$V_{0(1)}(E) = V_{0(1)} - \alpha(E - E_C)$$

$$W(E) = W \left/ \left(1 + \exp\left(\frac{\eta_W - (E - E_C)}{\lambda_W}\right) \right) \right.$$

$$W_{s,0(1)}(E) = W_{s,0(1)} \left/ \left(1 + \exp\left(\frac{(E - E_C) - \eta_s}{\lambda_s}\right) \right) \right.$$

- Goodman-Weare¹⁰ affine-invariant stretch move sampler from emcee¹¹, using default step-size ($a = 2$)
- 256 chains, 200,000 samples each
- massively parallel calibration requires ~ 3 days to run
- priors taken from CHUQ posterior
- priors for r_1^0 , r_1^A and a_1 taken from Danielewicz¹²
- each data sector re-scaled to be equally important
- fractional uncertainties reasonably consistent with CHUQ:
 $\frac{d\sigma^{(n,n)}}{d\Omega}$: 35%, $\frac{d\sigma^{(p,p)}}{d\Omega}$: 44%, $A_y^{(n,n)}$: 140%, $A_y^{(p,p)}$: 155%, $\frac{d\sigma^{(p,n)}}{d\Omega}$: 55%

¹⁰Goodman and Weare 2010, *Communications in applied mathematics and computational science* 5(1), pp. 65–80

¹¹Foreman-Mackey et al. 2013, *PASP* 125(925), p. 306

¹²Danielewicz et al. 2017, *Nuclear Physics A* 958, pp. 147–186

Charge exchange to isobaric analog states

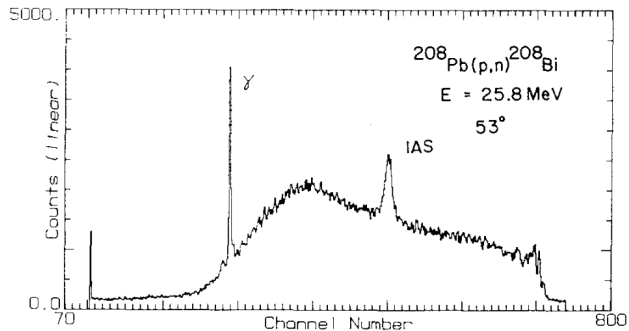
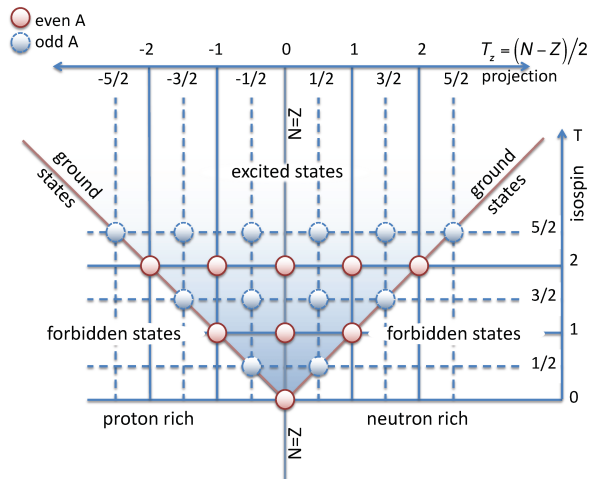


Fig. 1. Time-of-flight neutron spectrum for the reaction $^{208}\text{Pb}(p,n)^{208}\text{Bi}$ at 25.8 MeV showing the broad peak of the isobaric analogue state (IAS) and a narrow peak caused by incomplete rejection of γ -rays. The abscissa gives the time with units of 0.50 nsec per channel; higher energy neutrons are to the right in the spectrum. The flight path was approximately 9 m.

Reprinted from¹³

¹³SD Schery et al. (1974). "The (p, n) reaction to the isobaric analogue state of high-Z elements at 25.8 MeV". In: *Nuclear Physics A* 234.1, pp. 109–129.

Charge exchange to isobaric analog states



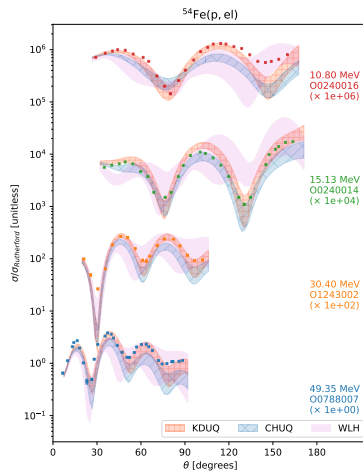
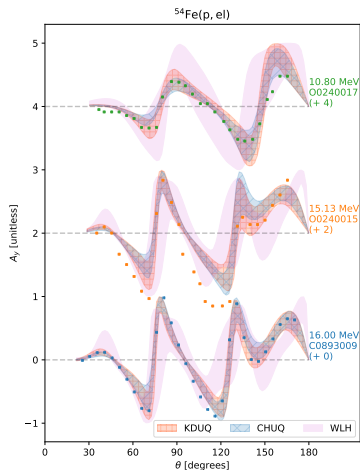
Reprinted from¹⁴

¹⁴Michael A Bentley (2022). "Excited States in Isobaric Multiplets—Experimental Advances and the Shell-Model Approach". In: *Physics* 4.3, pp. 995–1011. ISSN: 2624-8174. DOI: [10.3390/physics4030066](https://doi.org/10.3390/physics4030066). URL: <https://www.mdpi.com/2624-8174/4/3/66>

jit $\mathcal{R}$ demos and tutorials: built-in optical potentials

```
[2]: from jitr.optical_potentials import chuq, kduq, wlh
```

https://github.com/beykyle/jitr/blob/main/examples/notebooks/builtin_omps_uq.ipynb



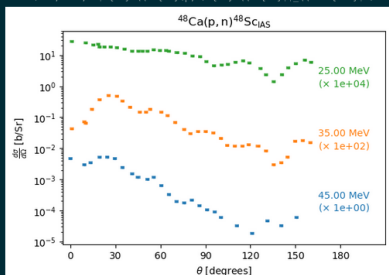
Curating a body of evidence

```
entry_pn = ExforEntry(
    entry="00178",
    reaction=ca48_pn,
    quantity="dXS/dA",
    vocal=True,
    parsing_kwargs={
        "statistical_err_labels": ["DATA-ERR"],
        "systematic_err_labels": ["ERR-1"],
    },
)

Found subentry 00178006 with the following columns:
['ERR-1', 'ERR-2', 'ERR-3', 'ERR-4', 'EN', 'Q-VAL', 'ANG-ERR', 'DATA-ERR', 'ANG-CM', 'DATA', 'ERR-T']
Found subentry 00178007 with the following columns:
['ERR-1', 'ERR-2', 'ERR-3', 'ERR-4', 'EN', 'Q-VAL', 'ANG-ERR', 'DATA-ERR', 'ANG-CM', 'DATA', 'ERR-T']
Found subentry 00178008 with the following columns:
['ERR-1', 'ERR-2', 'ERR-3', 'ERR-4', 'EN', 'Q-VAL', 'ANG-ERR', 'DATA-ERR', 'ANG-CM', 'DATA']

fig, ax = plt.subplots(1, 1, figsize=(6, 4))
entry_pn.plot(
    ax, offsets=100, label_kwargs={"label_xloc_deg": 165, "label_offset_factor": 0.0001}
)
ax.set_xlim([-5, 210])
ax.set_title(ax.title.get_text() + r"${\text{IAS}}$")

Text(0.5, 1.0, "${\rm{Ca}}\{{p,n}\}^{\rm{Sc}}_{\rm{IAS}}$")
```



https://github.com/beykyle/exfor_tools

- reproducible curation
- systematically improvable with new evaluations