

Loopfest 2026 – Brookhaven National Lab

Electroweak Corrections to $gg \rightarrow \gamma\gamma$

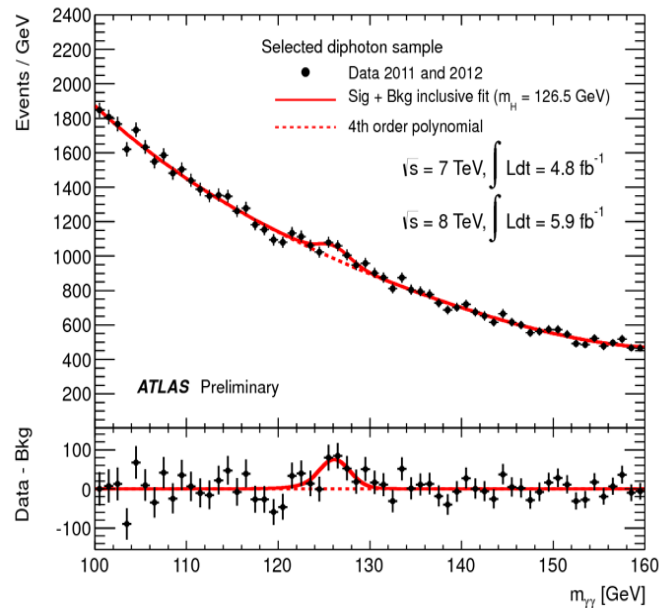
Gabriele Fiore

In collaboration with *Ciaran Williams*

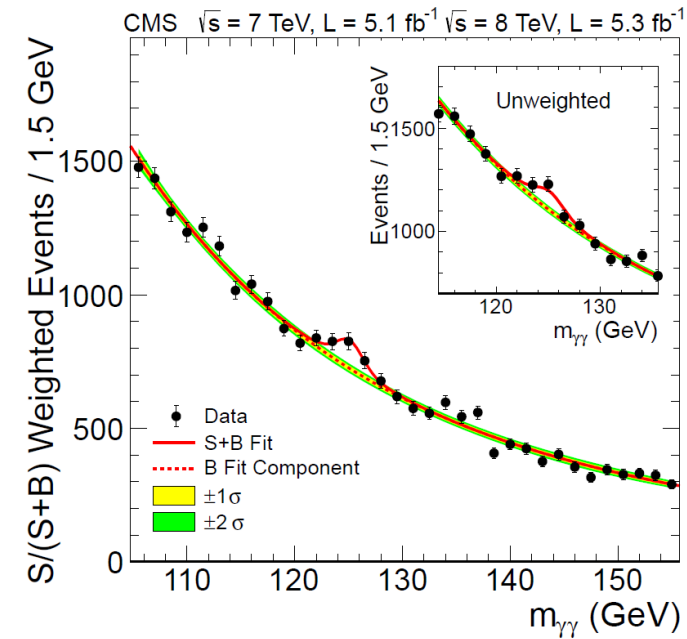
Based on: arXiv:2306.03956 and arXiv:2510.21643

Motivation: Diphoton channel

➤ Diphoton channel has a historical importance in the discovery of the Higgs boson.



ATLAS Collaboration, 2012



CMS Collaboration, 2012

➤ We distinguish between two different initial states: $q\bar{q} \rightarrow \gamma\gamma$ and $gg \rightarrow \gamma\gamma$.

Diphoton Predictions: $q\bar{q}$ Channel

$$q\bar{q} \rightarrow \gamma\gamma$$

Full NLO:

- *Binoth, Guillet, Pilon, Werlen, '99*

NNLO, QCD:

- *Catani, Cieri, de Florian, Ferrera, Grazzini, 2011*
- *Campbell, Ellis, Li, Williams, 2016*
- *Gehrmann, Glover, Huss, Whitehead, 2020*

NNLO, q_t -resummed:

- NNLL: *Cieri, Coradeschi, de Florian, 2015*
- N3LL: *Neumann, 2021*

NNLLO + NNLO_{QCD} + PS:

- *Alioli, Broggio, Gavardi, et. al 2021*
- *Gavardi, Oleari, Re, 2022*

NNLO, top-mass effects:

- *Becchetti, Bonciani, Cieri, Coro, Ripani, 2024*

N3LO amplitude:

- *Caola, Manteuffel, Tancredi, 2020*

NLO EW corrections:

- *Bierweiler, Kasprzik, Kühn, 2013*
- +2j: *Chiesa, Greiner, Schonherr, Tramontano, 2017*

NLO QCD + jets:

- *Del Duca, Maltoni, Nagy, Trocsanyi, 2003*
- *Gehrmann, Greiner, Heinrich, 2013*
- *Badger, Guffanti, Yundin, 2013*
- *Bern, Dixon, Cordero, Hoeche, Ita, Kosower, Lo Presti, Maitre, 2014*

Diphoton Predictions: gg Channel

$$gg \rightarrow \gamma\gamma$$

NNLO:

➤ *Dicus, Willenbrock, '88*

N3LO, QCD:

➤ Light quarks: *Bern, Dixon, Schmidt 2002*

➤ Heavy quarks: *Maltoni, Mandal, Zhao, 2018*

Chen, Heinrich, et al., 2020

Ahmed, Chakraborty, Chaubey, Kaur, 2025

Davies, Grau, Schönwald, Steinhauser, Stremmer, Vitti, 2025

N3LO + jet, QCD corrections:

➤ *Badger, Gehrmann, Marcoli, Moodie, 2021*

N4LO, QCD amplitude:

➤ *Bargiela, Caola, Manteuffel, Tancredi, 2022*

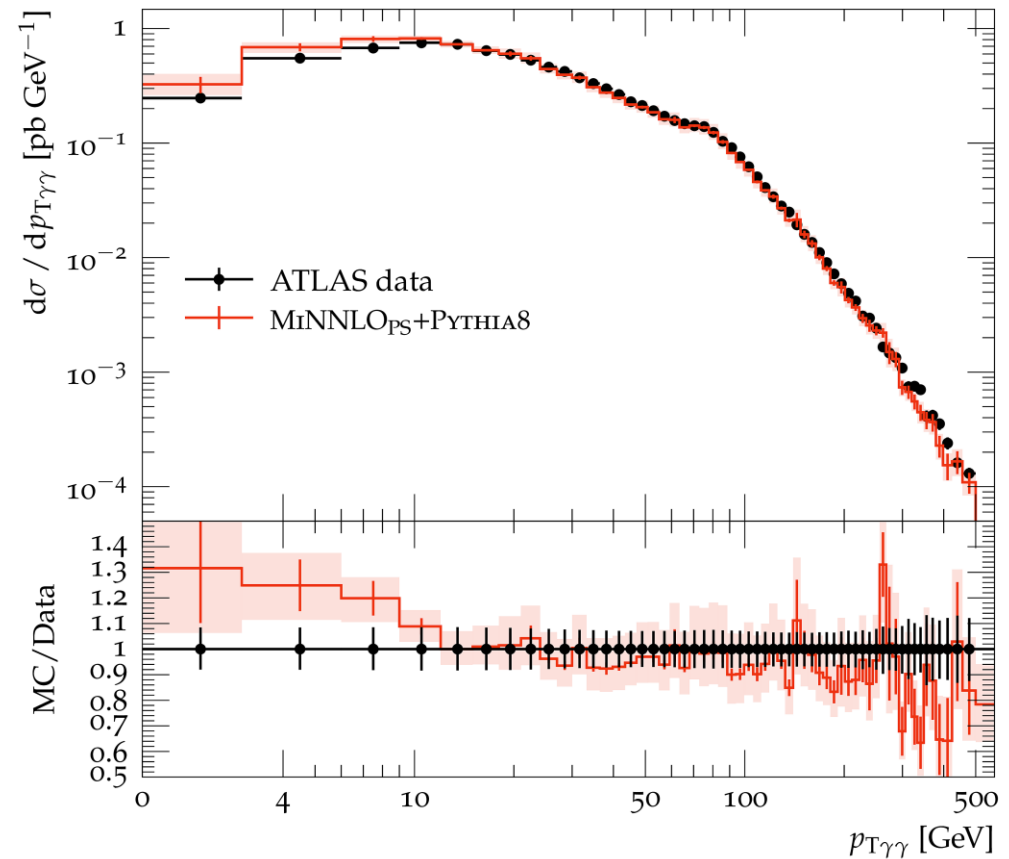
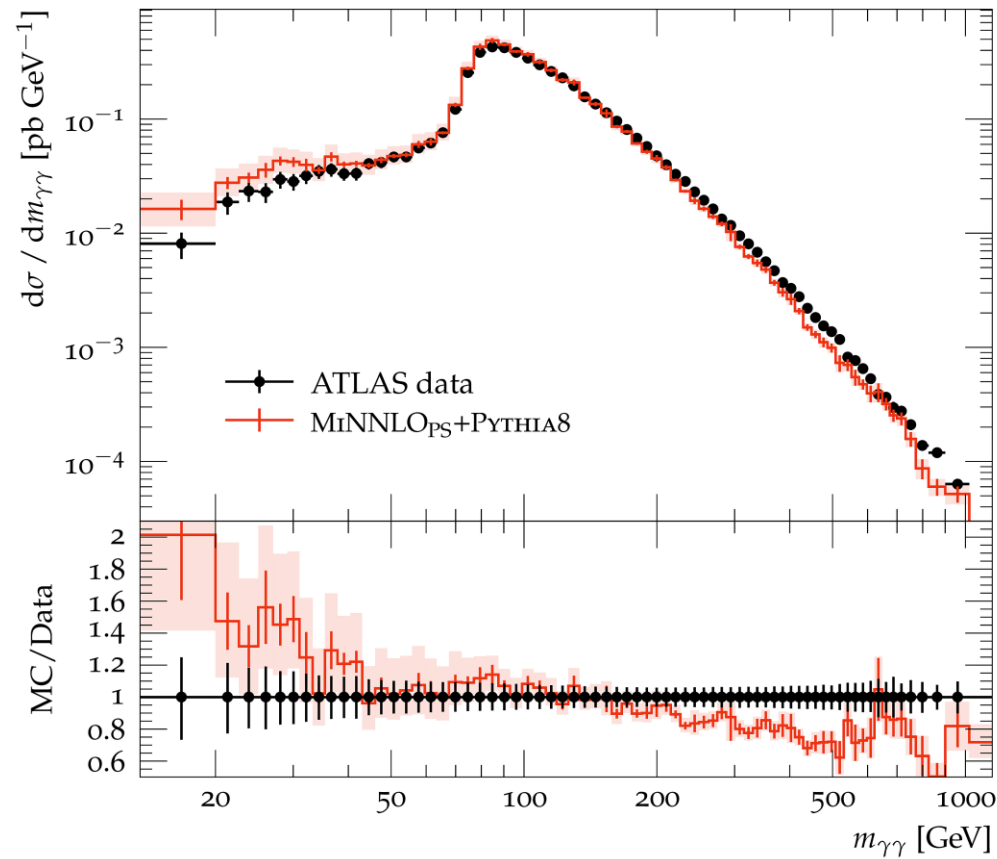
EW corrections missing



Objective of today's talk!

Motivation: Big Picture

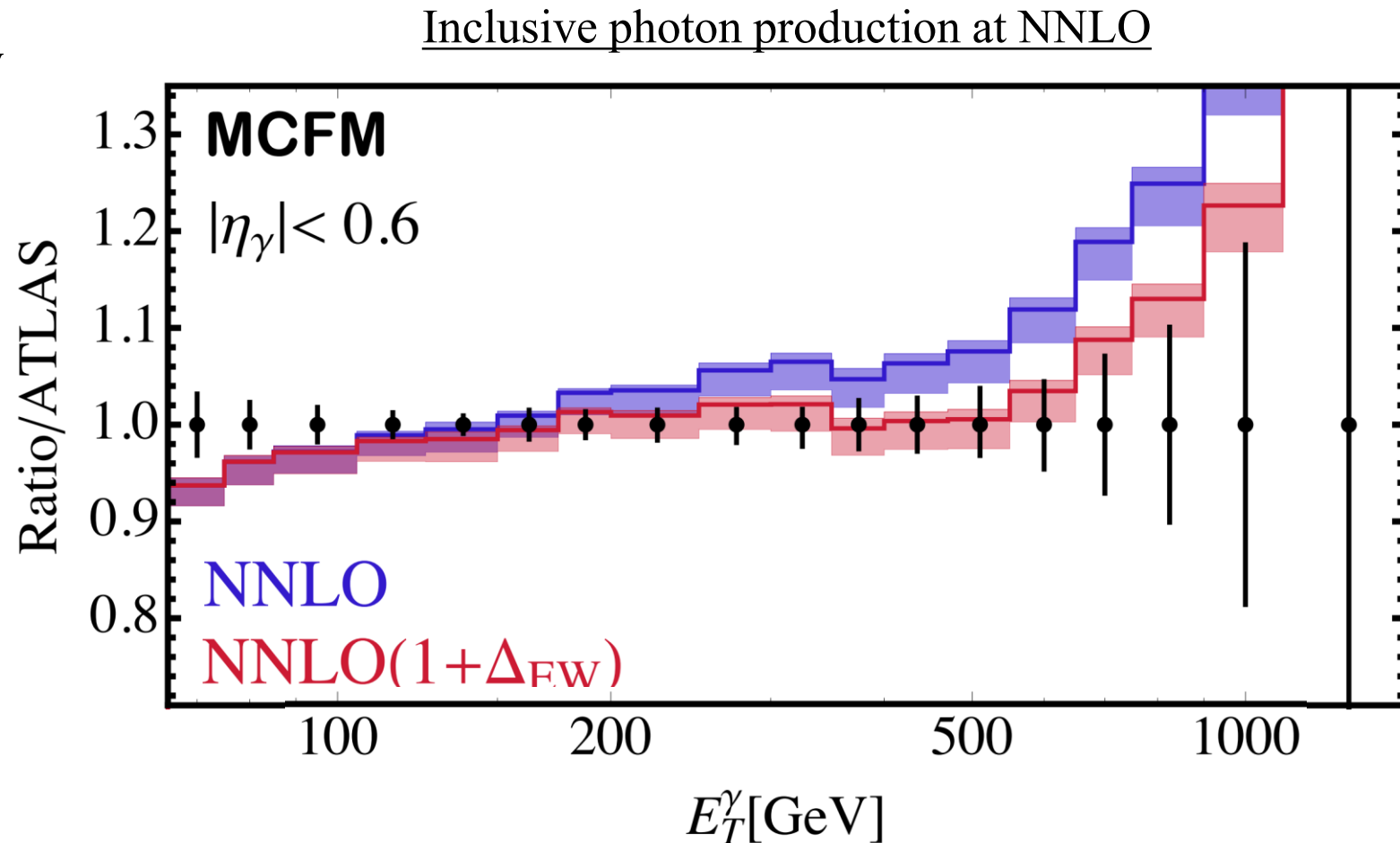
➤ Modern analysis, NNLO + PS matching in QCD compared with Atlas data.



Gavardi, Oleari, Re, 2022

Diphoton Predictions: Impact of EW corrections

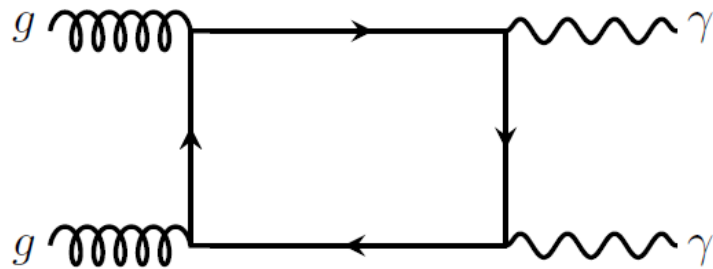
- With the QCD corrections reaching very high precision, the EW sector becomes more relevant.
- At high energy the Sudakov effects improve the comparison with the data.
- Similar trend is observed on different processes as well. e.g. DY with mixed QCD-EW (*Balossini, Calame, Montagna, Moretti, Nicrosini, Piccinini, Treccani, Vicini, 2009*).



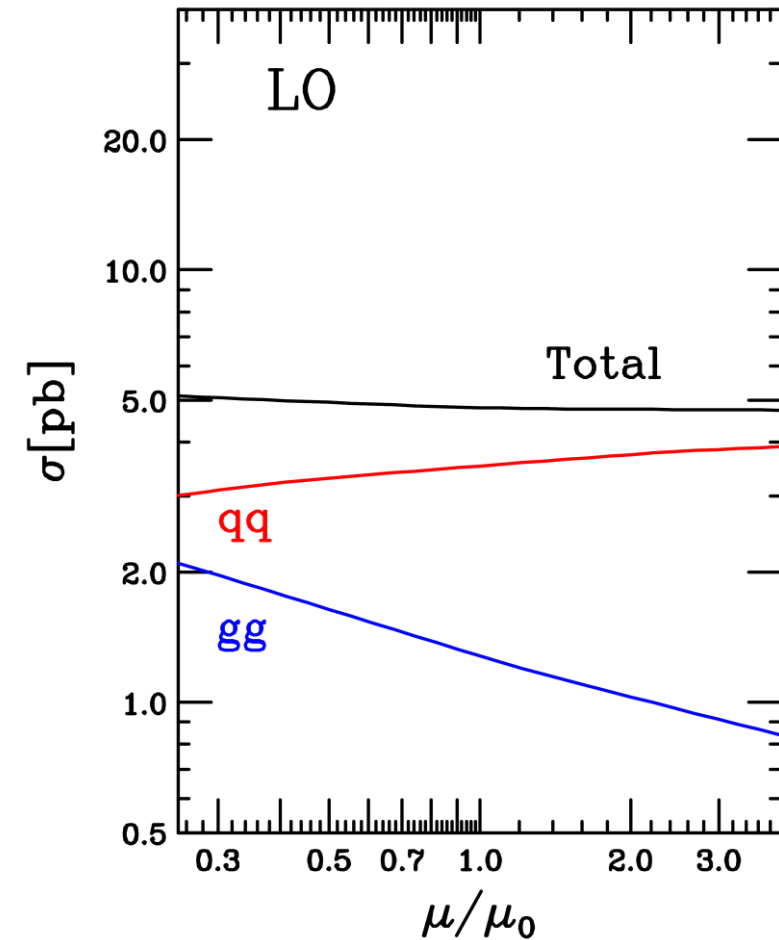
Campbell, Ellis, Williams, 2016

Motivation: gg Channel

- As part of the partonic process $pp \rightarrow \gamma\gamma$, the corrections enter at NNLO in QCD.



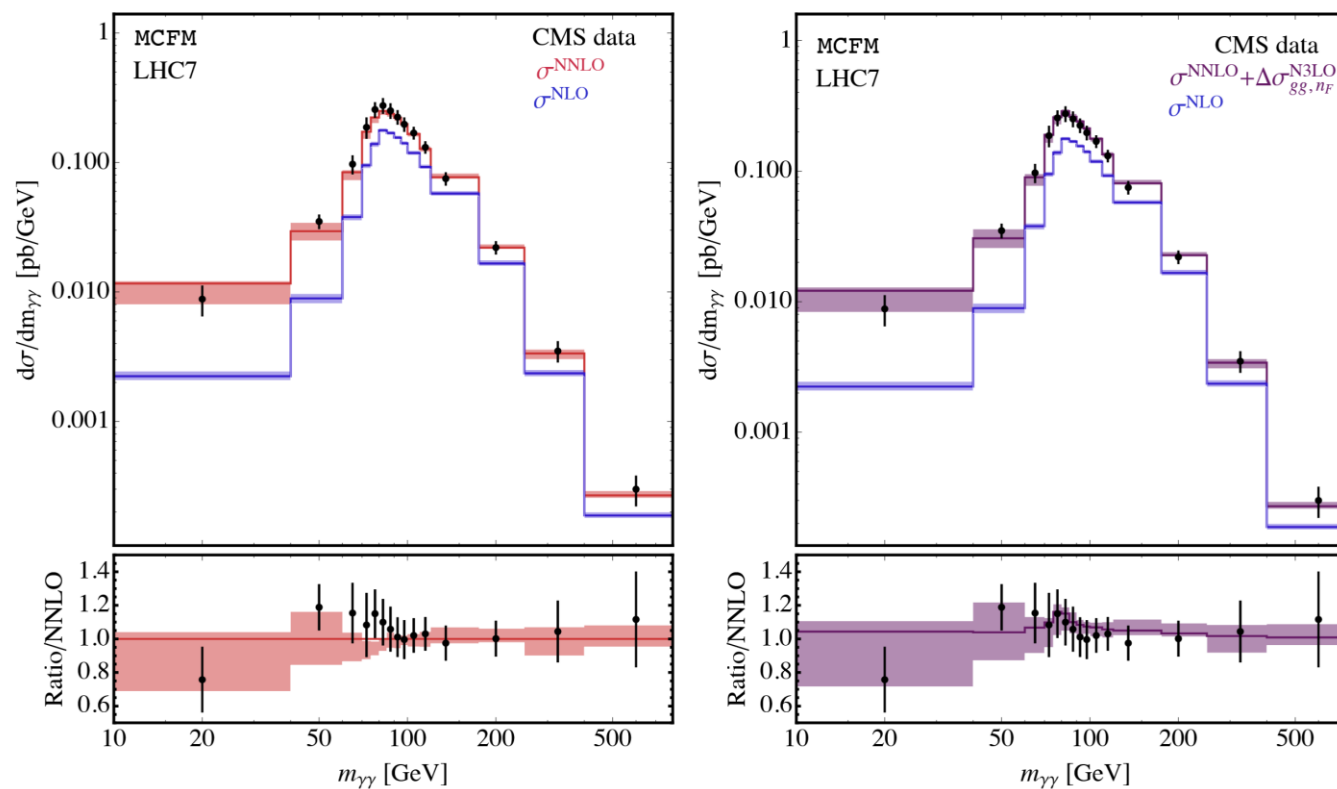
- Sizable contributions to the cross-section due to large initial-state gluon flux.
- Given the quality of recent LHC data, $\mathcal{O}(\alpha)$ corrections for the $gg \rightarrow \gamma\gamma$ channel are required to improve the precision our theoretical predictions.



Campbell, Ellis, Williams, 2011

Motivation: gg Channel

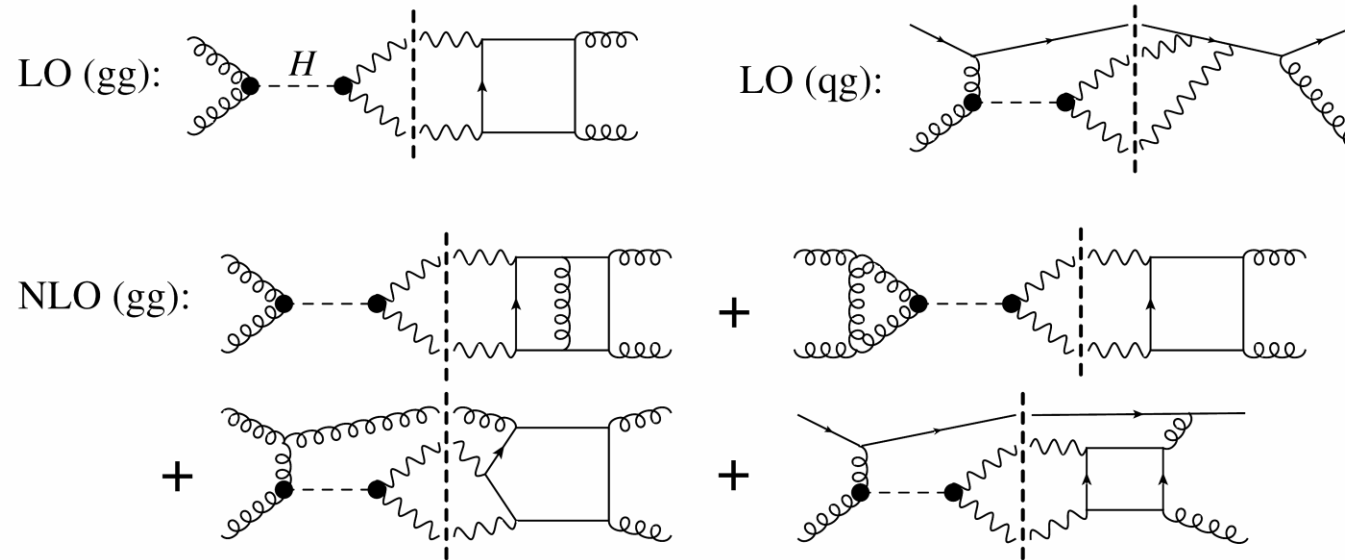
- Below, the impact of the gg channel in the NNLO implementation of $pp \rightarrow \gamma\gamma$ in MCFM. Data points taken from CMS(1405.7225).



Campbell, Ellis, Li, Williams, 2016

Motivation: Interference

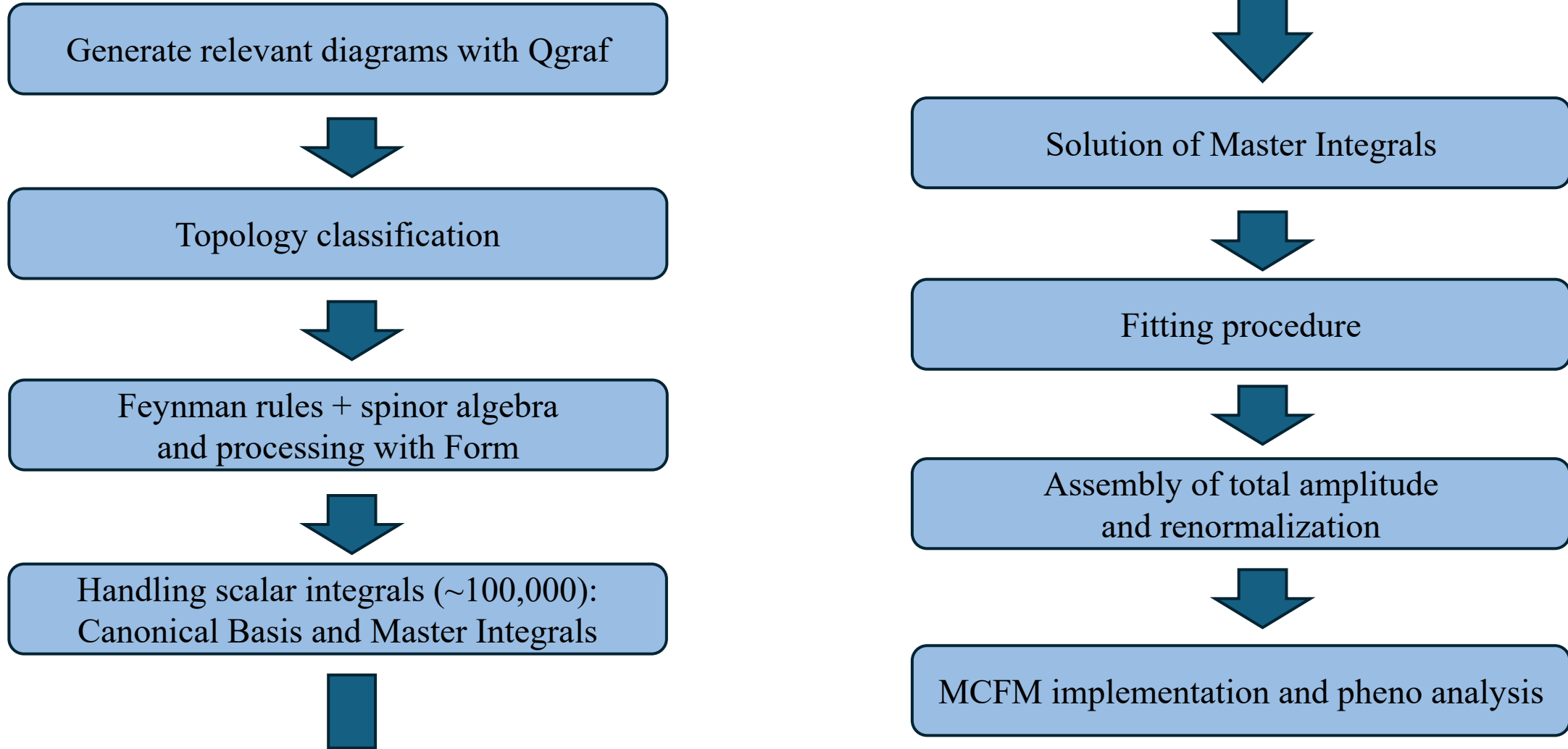
- $gg \rightarrow \gamma\gamma$ channel is crucial for studying the Higgs properties. A crucial example is the width, too small to be directly measured at LHC.
- The interference between the diphoton channel $gg \rightarrow H \rightarrow \gamma\gamma$ and the QCD background $gg \rightarrow \gamma\gamma$ can be used to extract better bounds on Higgs width (*Dixon, Siu, 2006; Martin, 2012; Dixon, Li, 2013*).



Motivation: Interference

- NLO analysis shows mass shift effect of $\sim 70\text{MeV}$ and destructive interference to the total cross-section of about 2% at NLO in QCD (Campbell, Carena, Harnik, Liu, 2017).
- Higgs width bounded to $20\text{-}40 \Gamma_{\text{H}}^{\text{SM}}$.
- These NLO QCD corrections are responsible for 40% effect on the mass shift, highlighting the relevance of higher order corrections.
- Beyond NLO (soft virtual approximation): reduction of the cross-section of about 1.7%, and an additional 30% for the mass shift on top of the NLO prediction (Bargiela, Buccioni, Caola, Devoto, Manteuffel, Tancredi, 2022).
- The EW corrections are still a missing piece in this picture!

Calculation Workflow



Calculation Workflow

Generate relevant diagrams with Qgraf

Topology classification

Feynman rules + spinor algebra
and processing with Form

Handling scalar integrals ($\sim 100,000$):
Canonical Basis and Master Integrals

Solution of Master Integrals

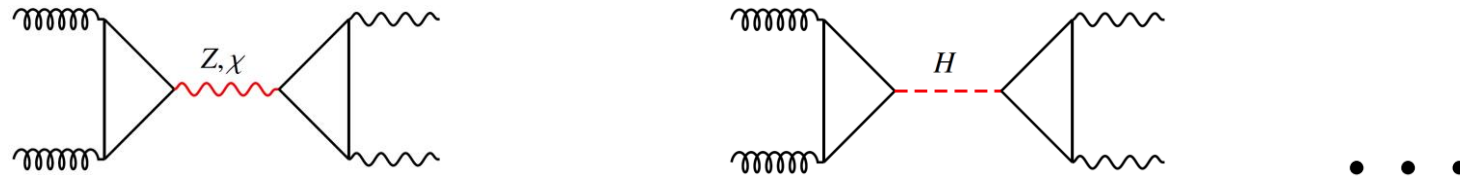
Fitting procedure

Assembly of total amplitude
and renormalization

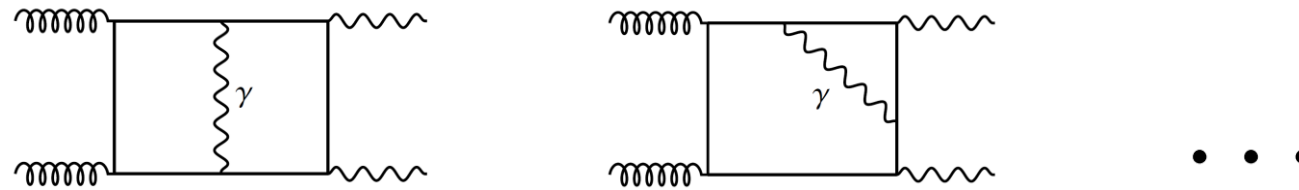
MCFM implementation and pheno analysis

Relevant Diagrams

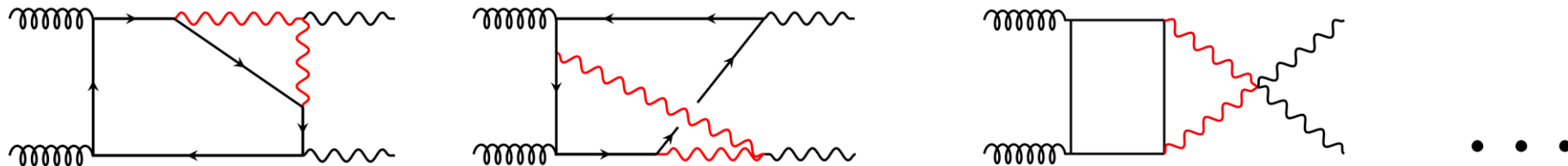
- Many of those are zero for the massless case (Furry theorem, charge conservation, massless limit,...).



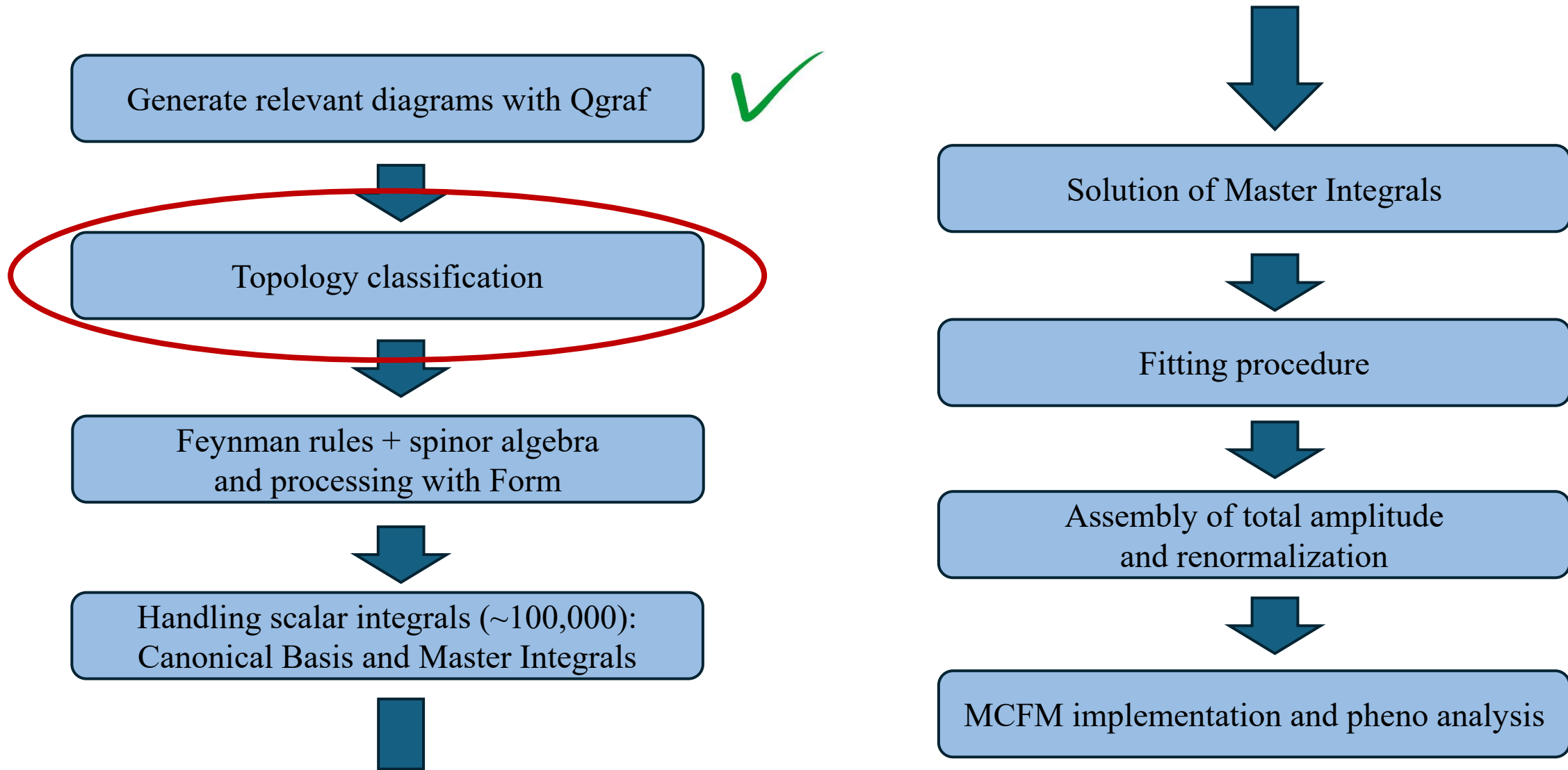
- QED diagrams already known. Use them as benchmark! (Bern, De Freitas, Dixon, 2001)



- New: pure non-abelian contributions with W -type insertion. Different behavior and charge structure.



Calculation Workflow



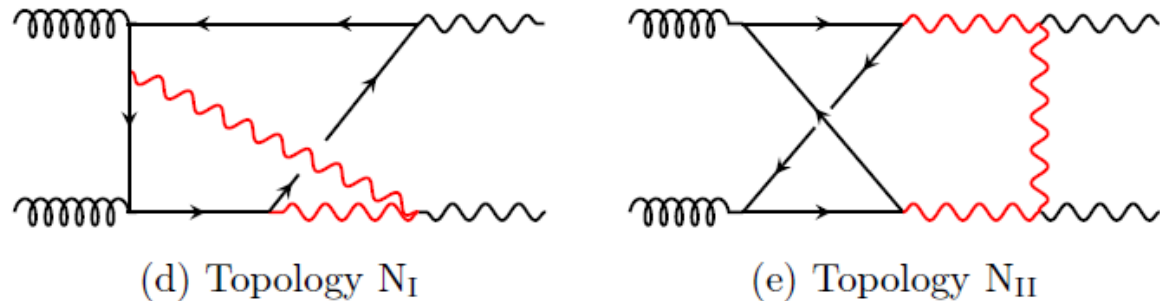
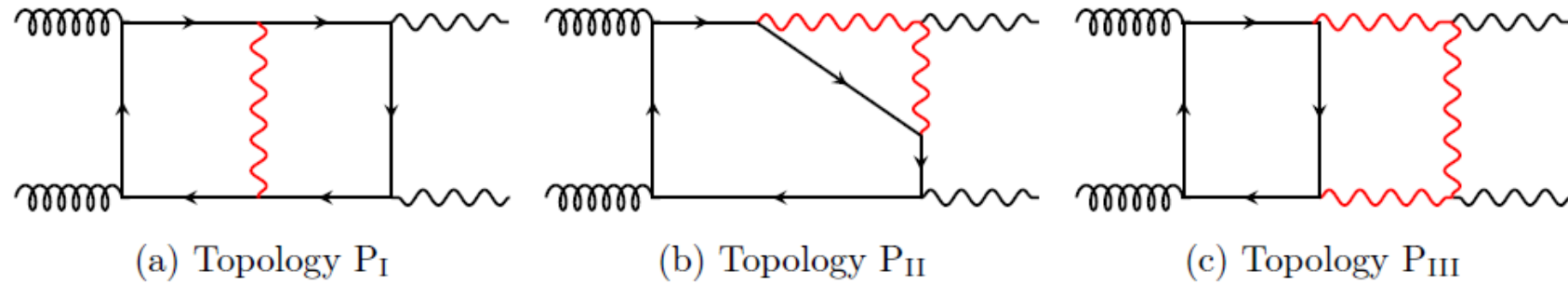
Topology Classification

- We identify 5 topologies:
 - 3 planar (P_I , P_{II} , P_{III}).
 - 2 non-planar (N_I , N_{II}).

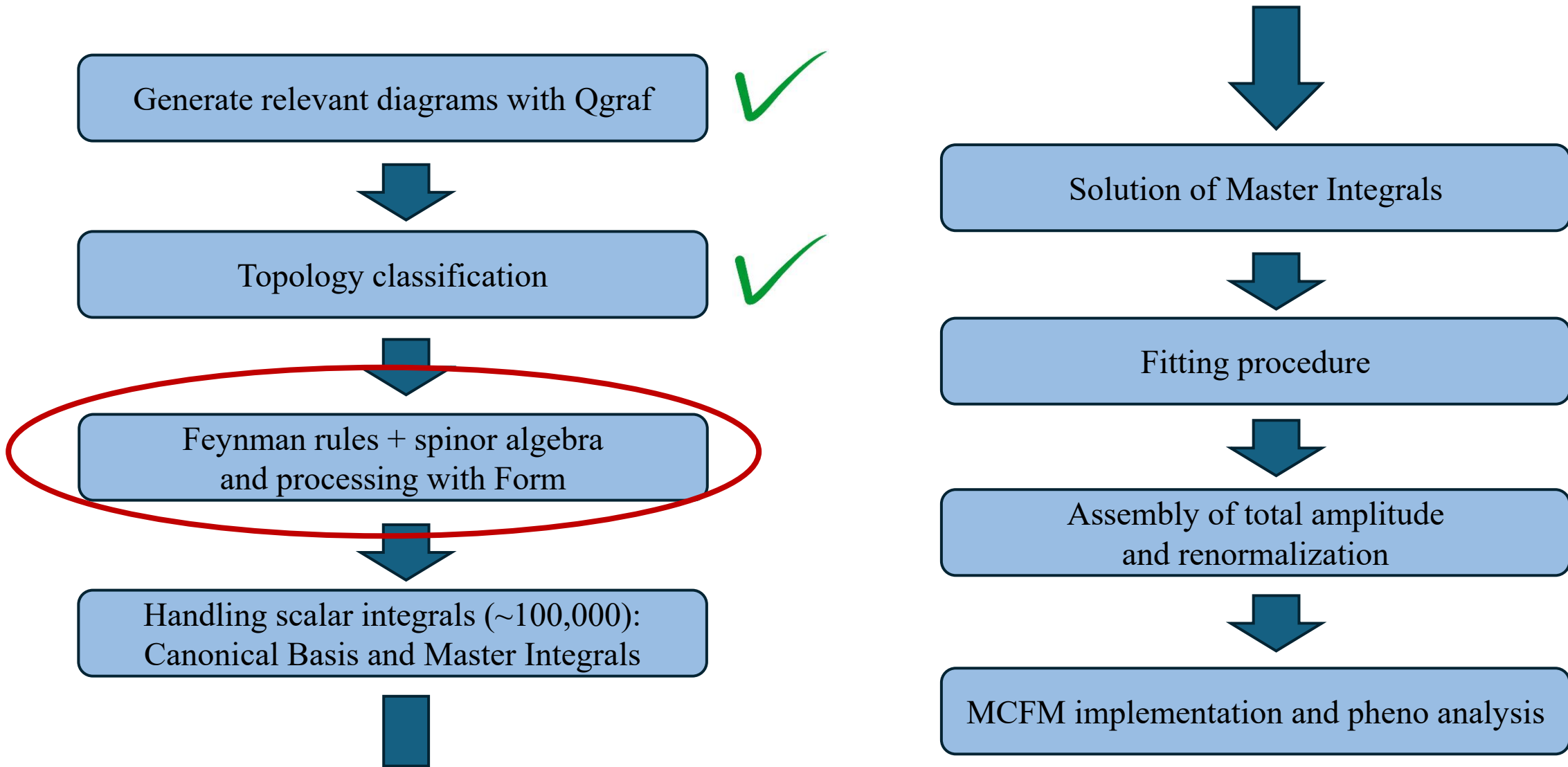
- Topology P_{II} is technically a subtopology of P_{III} but it's convenient to separate them.

- Z-box type diagrams are expressed in terms of topology P_I .

- QED type diagrams are analyzed separately.



Calculation Workflow



Helicity Amplitudes: Tensor Structures

- The idea is to facilitate the calculation by using the well-known spinor-helicity formalism.

$$\longrightarrow \mathcal{A}_{gg \rightarrow \gamma\gamma} = \varepsilon_\mu \varepsilon_\nu \varepsilon_\rho \varepsilon_\sigma \mathcal{A}_{gg \rightarrow \gamma\gamma}^{\mu\nu\rho\sigma}$$

- We write the amplitude in terms of a rank-4 Lorentz tensor, and we expand it into a tensor structure basis.

$$\longrightarrow \mathcal{A}_{gg \rightarrow \gamma\gamma}^{\mu\nu\rho\sigma} = \sum_{i=1}^{N_T} \mathcal{F}_i T_i^{\mu\nu\rho\sigma}$$

- We construct Projectors based on those tensor structures and we extract the relevant form factors.

$$\longrightarrow \mathcal{F}_{\lambda_i} = \sum_{pol} \mathcal{P}_{\lambda_i}^{\mu\nu\rho\sigma} (\varepsilon_{1,\mu})^* (\varepsilon_{2,\nu})^* (\varepsilon_{3,\rho})^* (\varepsilon_{4,\sigma})^* \mathcal{A}_{\lambda_i}$$

Physical Helicity Amplitudes

➤ From the Form Factors we can finally reconstruct the physical helicity amplitudes. We show one for brevity:

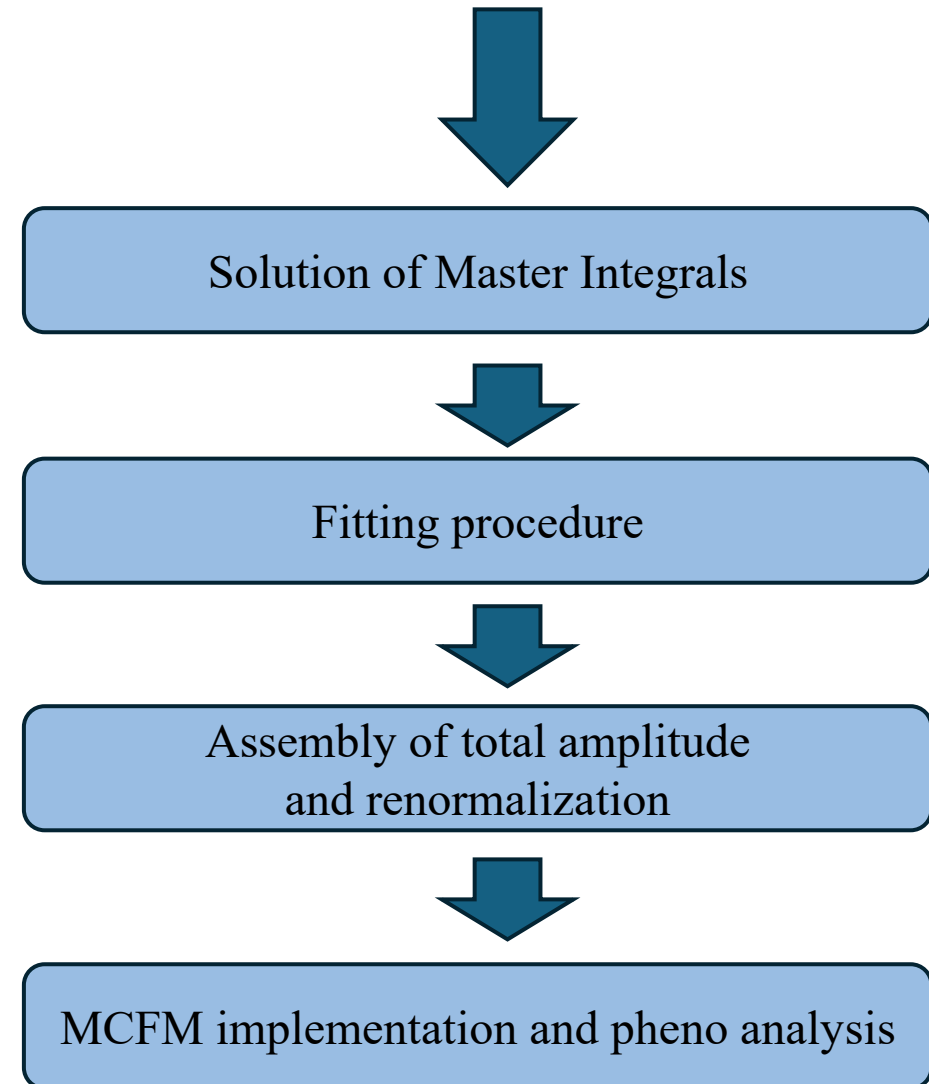
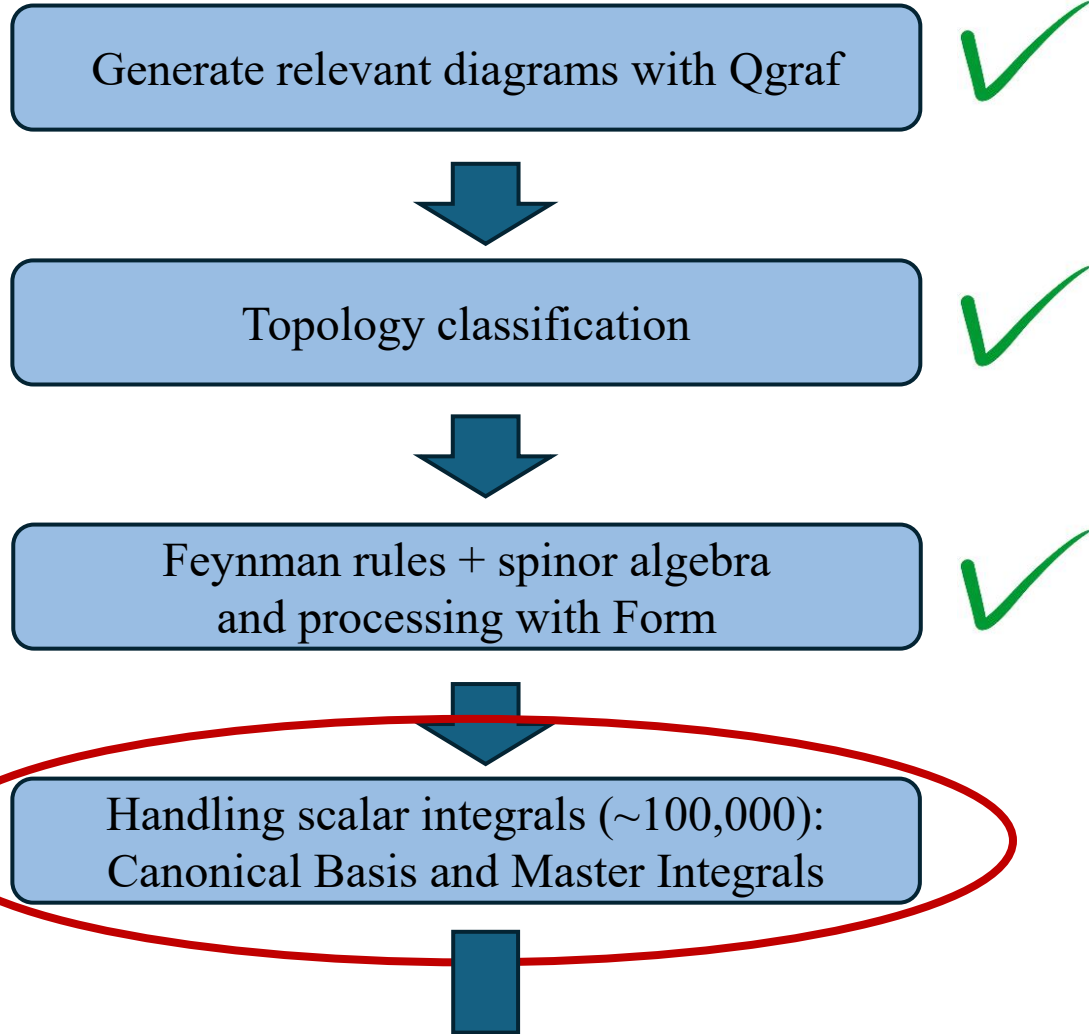
$$A_{\lambda_1} = \frac{t^2}{4} \left(\frac{2\mathcal{F}_6}{u} - \frac{2\mathcal{F}_3}{s} - \mathcal{F}_1 \right) + \mathcal{F}_8 \left(\frac{s}{u} + \frac{u}{s} + 4 \right) + \frac{t}{2} (\mathcal{F}_2 - \mathcal{F}_4 + \mathcal{F}_5 - \mathcal{F}_7)$$

➤ We have 8 independent helicity amplitudes. One for each of the following helicity configuration:

$$\begin{aligned} \lambda_1 &= \{+, +, +, +\}, & \lambda_2 &= \{-, +, +, +\} \\ \lambda_3 &= \{+, -, +, +\}, & \lambda_4 &= \{+, +, -, +\} \\ \lambda_5 &= \{+, +, +, -\}, & \lambda_6 &= \{-, -, +, +\} \\ \lambda_7 &= \{-, +, -, +\}, & \lambda_8 &= \{+, -, -, +\} \end{aligned}$$

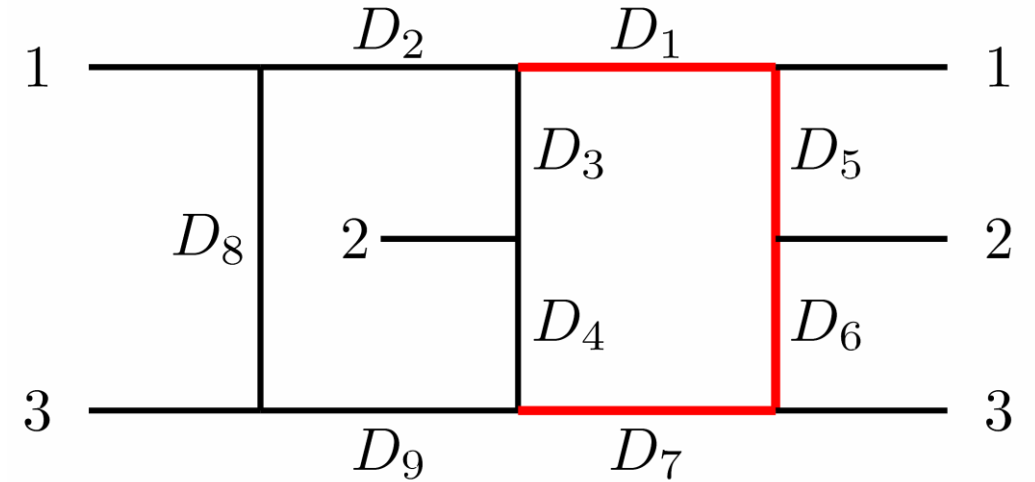
➤ The form factors contain thousands of loop integrals → MIs reduction!

Calculation Workflow

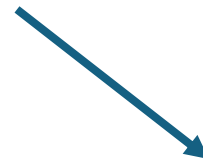


Auxiliary Topology

➤ We introduce an auxiliary topology to account for all the possible scalar products involving the loop momenta.



➤ All the scalar integrals can be written in terms of the auxiliary topology's propagators.



➤ This preliminary scalar integral reduction produces a multitude ($\sim 100,000$) of scalar integrals! ➔ Master Integral reduction

$$I_{a_1 \dots a_9}^{\text{NII}} = \mathcal{C}(\epsilon) \int \frac{d^d \ell_1}{(2\pi)^d} \frac{d^d \ell_2}{(2\pi)^d} \frac{1}{D_1^{a_1} \dots D_9^{a_9}}$$

DEQ: Overview

- General strategy to evaluate MIs: Differential Equation method (DEQ)

$$\text{Laporta Basis } \vec{\mathcal{I}}(\{x_i\}) \longrightarrow \frac{\partial}{\partial x_i} \vec{\mathcal{I}}(\{x_i\}) = \underbrace{\mathcal{A}(\{x_i\}, \epsilon)} \vec{\mathcal{I}}(\{x_i\})$$

- The complexity goes up pretty quickly for the Laporta Basis. Solution: Canonical Basis.

$$\text{Canonical Basis } \vec{\mathcal{G}}(\{x_i\}) \longrightarrow \frac{\partial}{\partial x_i} \vec{\mathcal{G}}(\{x_i\}) = \epsilon \underbrace{\mathcal{A}(\{x_i\})} \vec{\mathcal{G}}(\{x_i\})$$

- The factorization of the space-time parameter allows to write the solution in an iterative way:

$$\vec{\mathcal{G}}(\{x_i\}) = \left(\mathbb{1} + \epsilon \int_{\gamma} d\mathcal{A} + \epsilon^2 \int_{\gamma} d\mathcal{A} d\mathcal{A} + \dots \right) \vec{\mathcal{G}}(\{x_i^0\})$$

DEQ: Canonical Basis

- The kinematic information is encoded in the alphabet $\{\eta_k\}$ of the differential equation:

$$d\mathcal{A} = \sum_k \mathcal{M}_k d \log \eta_k$$

- The solution for rational letters is expressed in terms of Goncharov Polylogarithms (GPLs) (*Goncharov, 1998*).

$$G(a_n, \dots, a_1; x_0) = \int_0^{x_0} G(a_{n-1}, \dots, a_1; x_0) \frac{dt}{t - a_n},$$

- Efficient numerical evaluation of GPLs using GiNaC (*Vollinga, 2005*)/HandyG (*Naterop, Signer, Ulrich, 2016*).
- Non-rational letters may or may not be represented in terms of GPLs. We implement a generalized approach based on Chen's Iterated Integrals.

DEQ: Solution

In a more general scenario, the solution can be evaluated (up to boundary constants) at any fixed order:

$$\vec{\mathcal{G}}^{(1)}(\{x_i\}) = \int_{\gamma} d\mathcal{A} \vec{\mathcal{G}}^{(0)}(\{x_i^0\}) + \vec{\mathcal{G}}^{(1)}(\{x_i^0\})$$

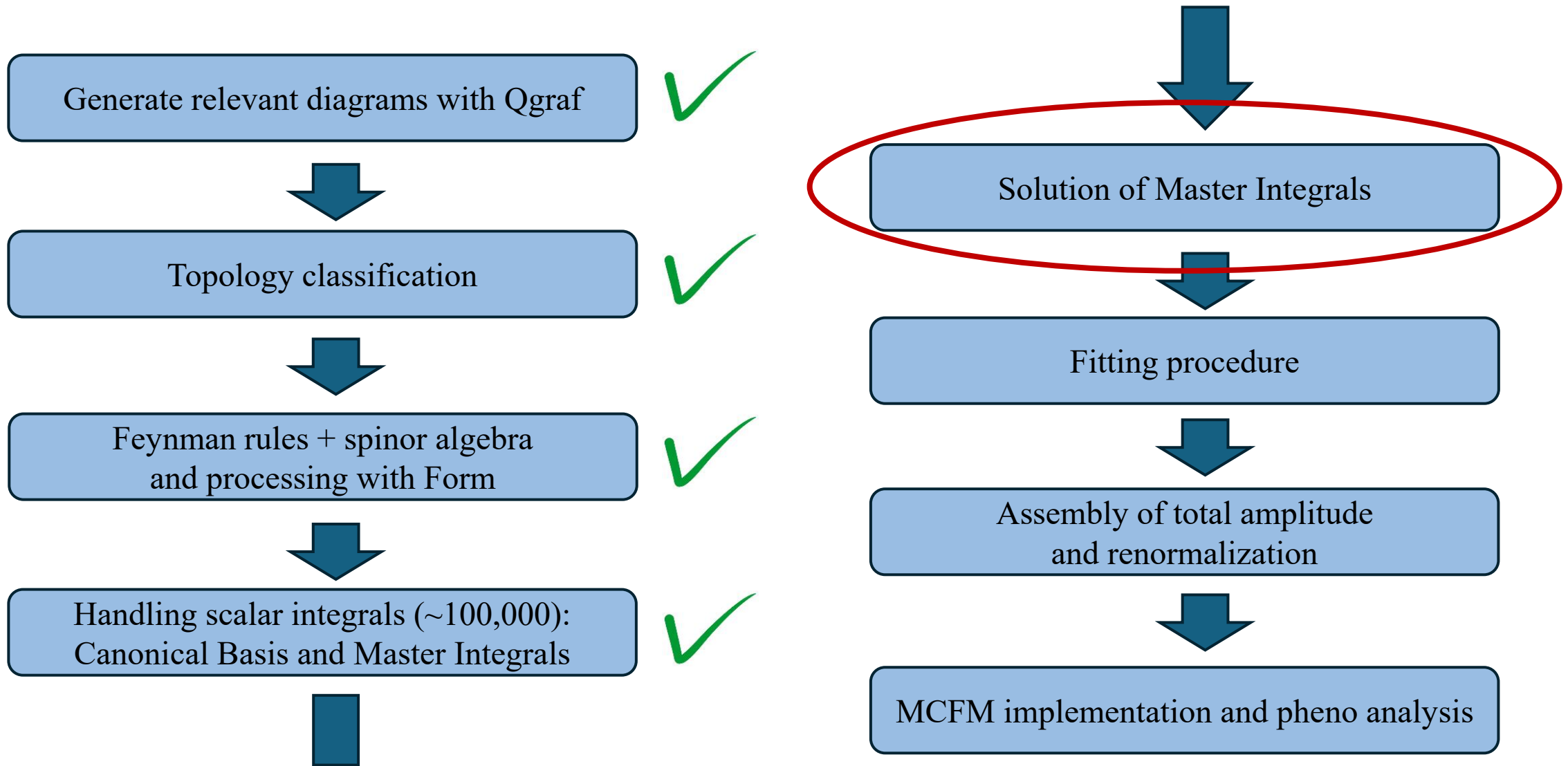
$$\vec{\mathcal{G}}^{(2)}(\{x_i\}) = \int_{\gamma} d\mathcal{A} d\mathcal{A} \vec{\mathcal{G}}^{(0)}(\{x_i^0\}) + \int_{\gamma} d\mathcal{A} \vec{\mathcal{G}}^{(1)}(\{x_i^0\}) + \vec{\mathcal{G}}^{(2)}(\{x_i^0\})$$

$$\vec{\mathcal{G}}^{(3)}(\{x_i\}) = \int_{\gamma} d\mathcal{A} d\mathcal{A} d\mathcal{A} \vec{\mathcal{G}}^{(0)}(\{x_i^0\}) + \int_{\gamma} d\mathcal{A} d\mathcal{A} \vec{\mathcal{G}}^{(1)}(\{x_i^0\}) + \int_{\gamma} d\mathcal{A} \vec{\mathcal{G}}^{(2)}(\{x_i^0\}) + \vec{\mathcal{G}}^{(3)}(\{x_i^0\})$$

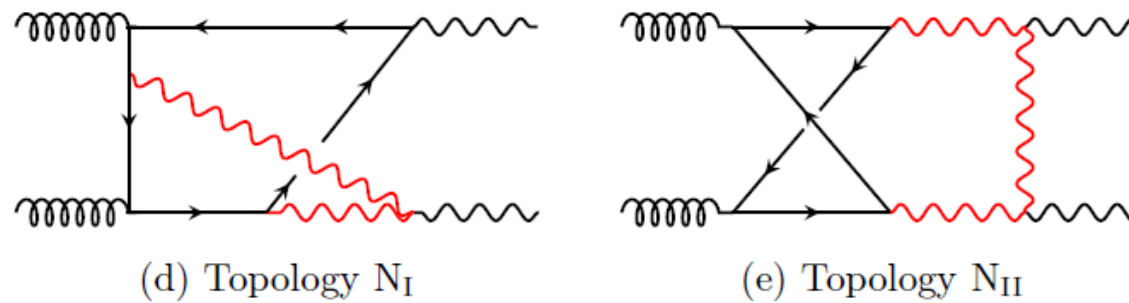
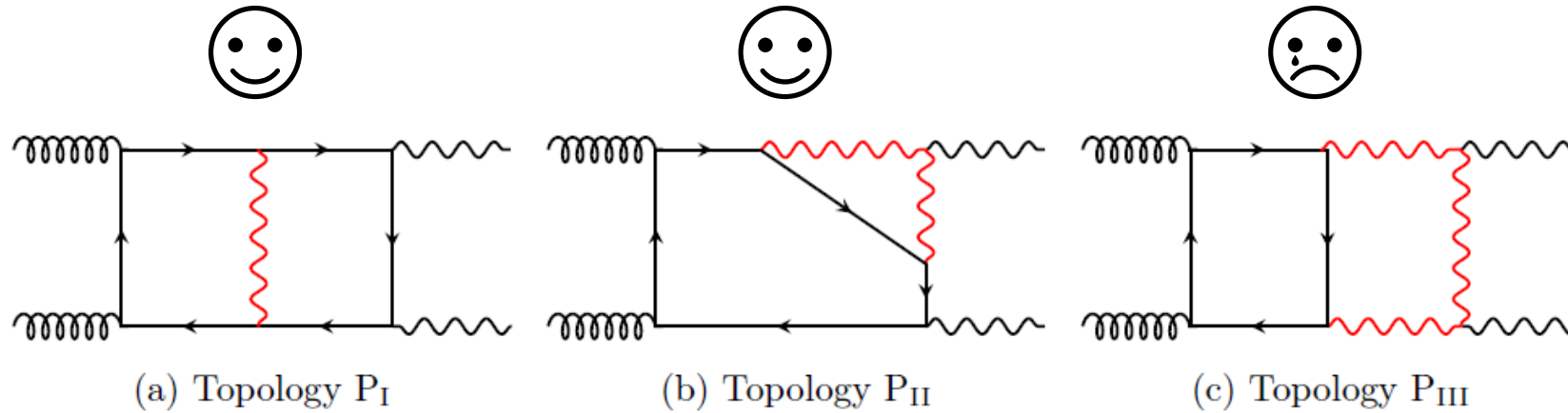
$$\begin{aligned} \vec{\mathcal{G}}^{(4)}(\{x_i\}) &= \int_{\gamma} d\mathcal{A} d\mathcal{A} d\mathcal{A} d\mathcal{A} \vec{\mathcal{G}}^{(0)}(\{x_i^0\}) + \int_{\gamma} d\mathcal{A} d\mathcal{A} d\mathcal{A} \vec{\mathcal{G}}^{(1)}(\{x_i^0\}) + \int_{\gamma} d\mathcal{A} d\mathcal{A} \vec{\mathcal{G}}^{(2)}(\{x_i^0\}) \\ &+ \int_{\gamma} d\mathcal{A} \vec{\mathcal{G}}^{(3)}(\{x_i^0\}) + \vec{\mathcal{G}}^{(4)}(\{x_i^0\}). \end{aligned}$$

Efficient semi-numerical implementation thanks to the recursive nature of the iterated integrals.

Calculation Workflow



Rational vs Non-Rational alphabet



Topology P_{III}: Canonical Basis

The Canonical basis is obtained again via Magnus Series Expansion. Here we show the 7-propagators family:

$$\mathcal{G}_{29}^{\text{P}_{\text{III}}} = \epsilon^2 y \sqrt{r_2} \mathcal{F}_{29},$$

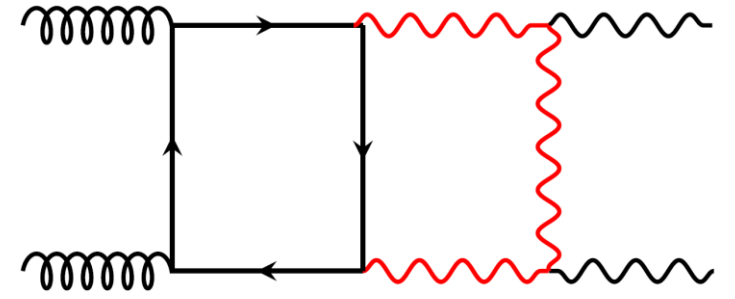
$$\mathcal{G}_{30}^{\text{P}_{\text{III}}} = \epsilon^2 \sqrt{r_1} (2x \mathcal{F}_{28} + xy \mathcal{F}_{29} + y \mathcal{F}_{30}),$$

$$\mathcal{G}_{31}^{\text{P}_{\text{III}}} = \epsilon^2 y^2 (\mathcal{F}_{29} + \mathcal{F}_{31}),$$

$$\mathcal{G}_{32}^{\text{P}_{\text{III}}} = -\frac{1}{2} \epsilon \left(2x(1 - 2\epsilon) \mathcal{F}_{21} + y^2 \epsilon (-\mathcal{F}_{27} + x \mathcal{F}_{29} + \mathcal{F}_{30}) \right. \\ \left. + 2y \epsilon (2\mathcal{F}_{17} + \mathcal{F}_{22} + x \mathcal{F}_{28} + x \mathcal{F}_{29} - \mathcal{F}_{32}) \right).$$

The differential equation is expressed in terms of the variables:

$$x = -\frac{2p_1 \cdot p_2}{M_W^2} \quad y = -\frac{2p_2 \cdot p_3}{M_W^2}$$



➤ Non-Rational letters:

$$r_1 = y(4 + y),$$

$$r_2 = y(y + 2xy + x^2(4 + y)),$$

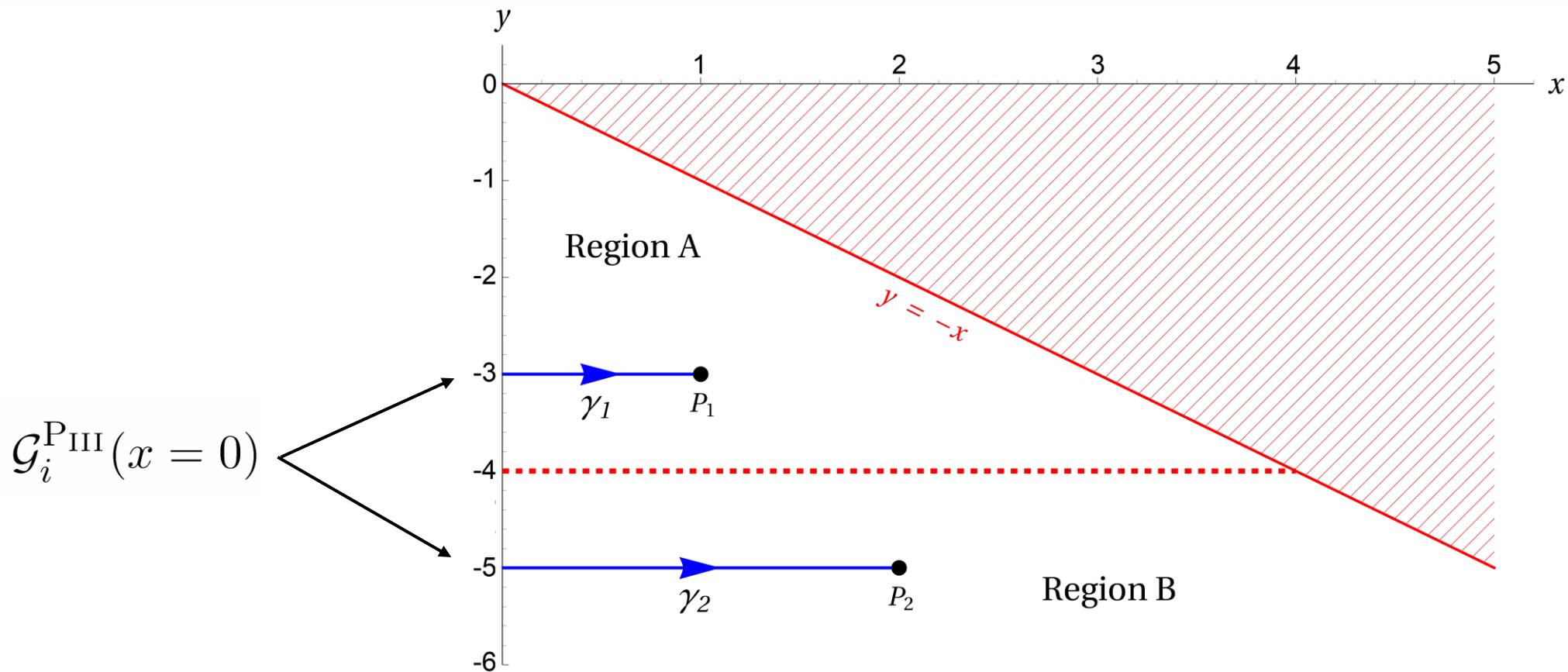
$$r_3 = xy(4y + x(4 + y)).$$

➤ For efficiency reasons, we only rationalize r_1 .

Topology P_{III} : Kinematic Sub-Regions

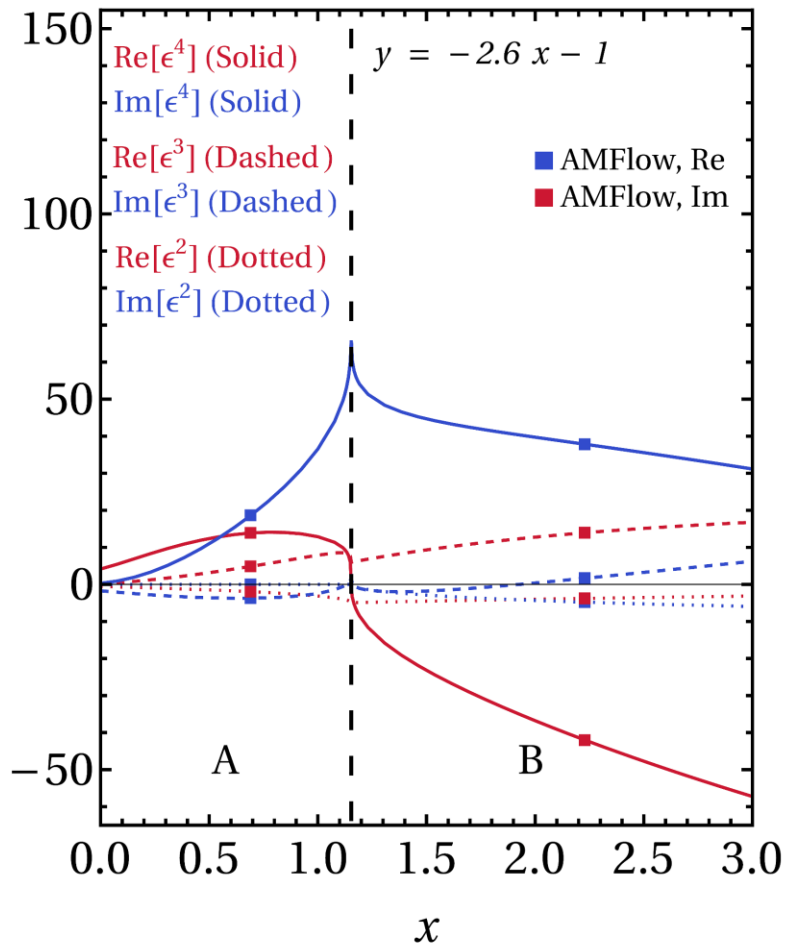
Region A : $[0 \geq y \geq -4, x \leq -y]$, $y = -\frac{4z_A^2}{1+z_A^2}$

Region B : $[y \leq -4, x \leq -y]$, $y = -\frac{4}{1-z_B^2}$

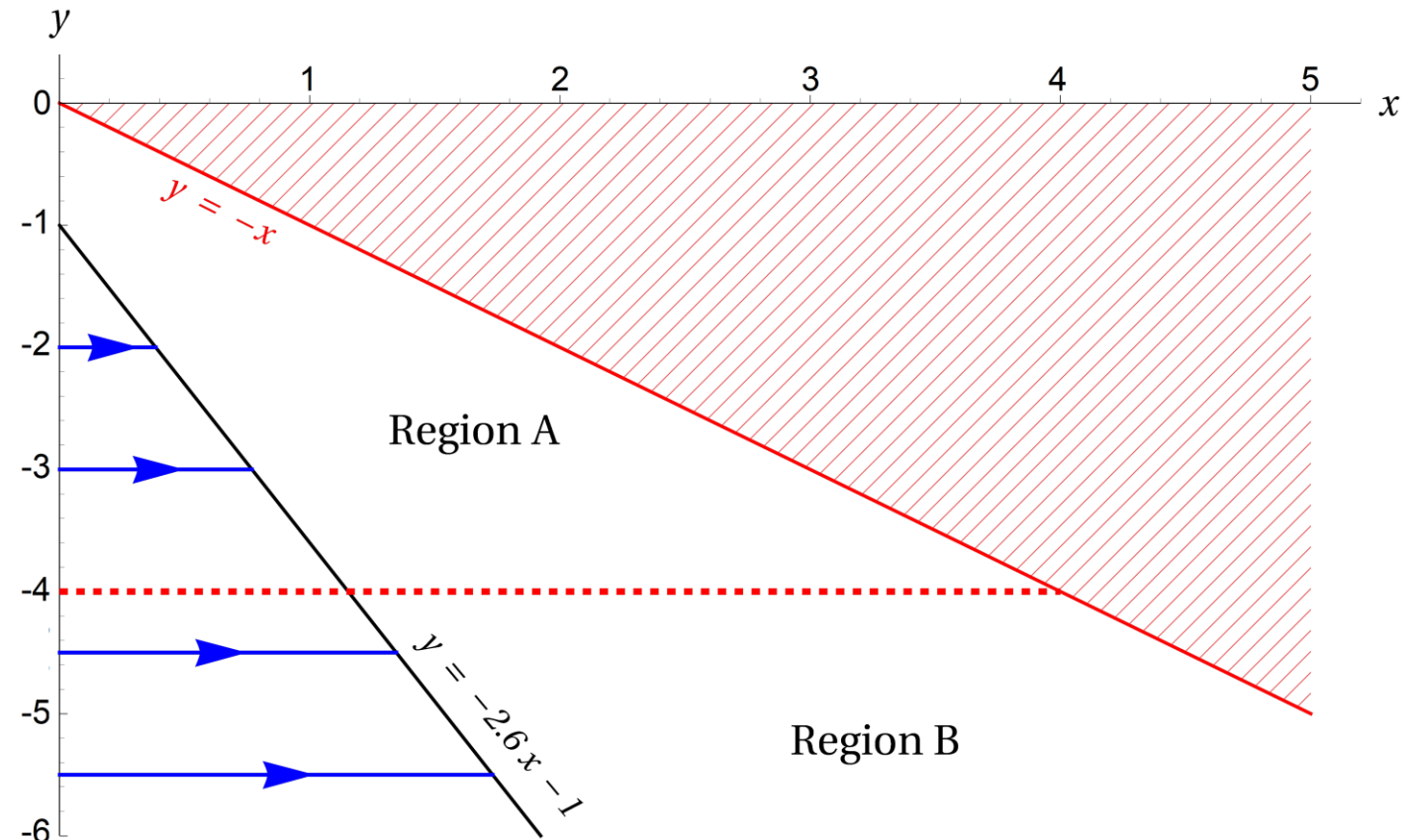


Topology \mathcal{P}_{III} : Numerical Evaluation and AMFlow Check

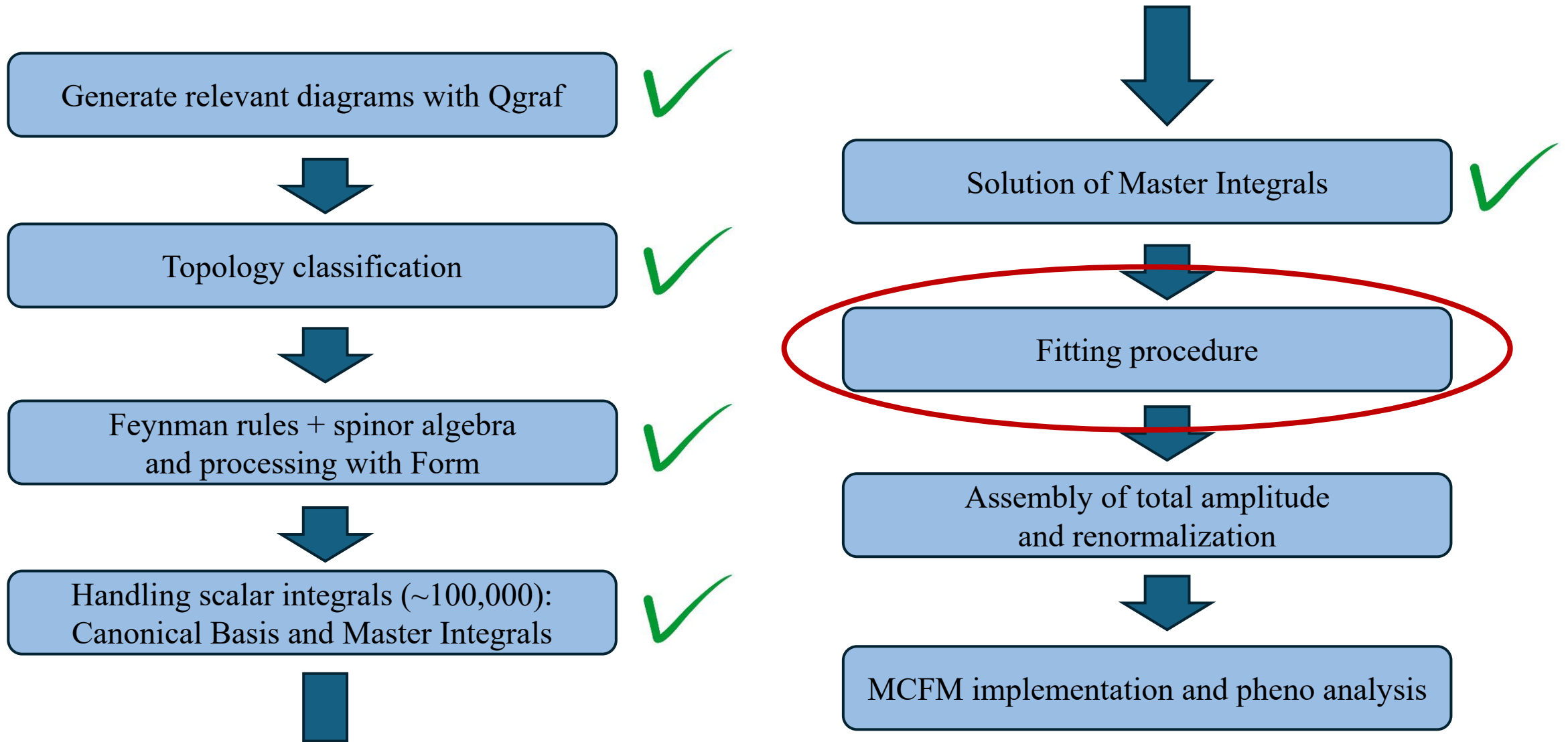
$$\mathcal{G}_{31}^{\mathcal{P}_{\text{III}}}$$



Integration Path:



Calculation Workflow



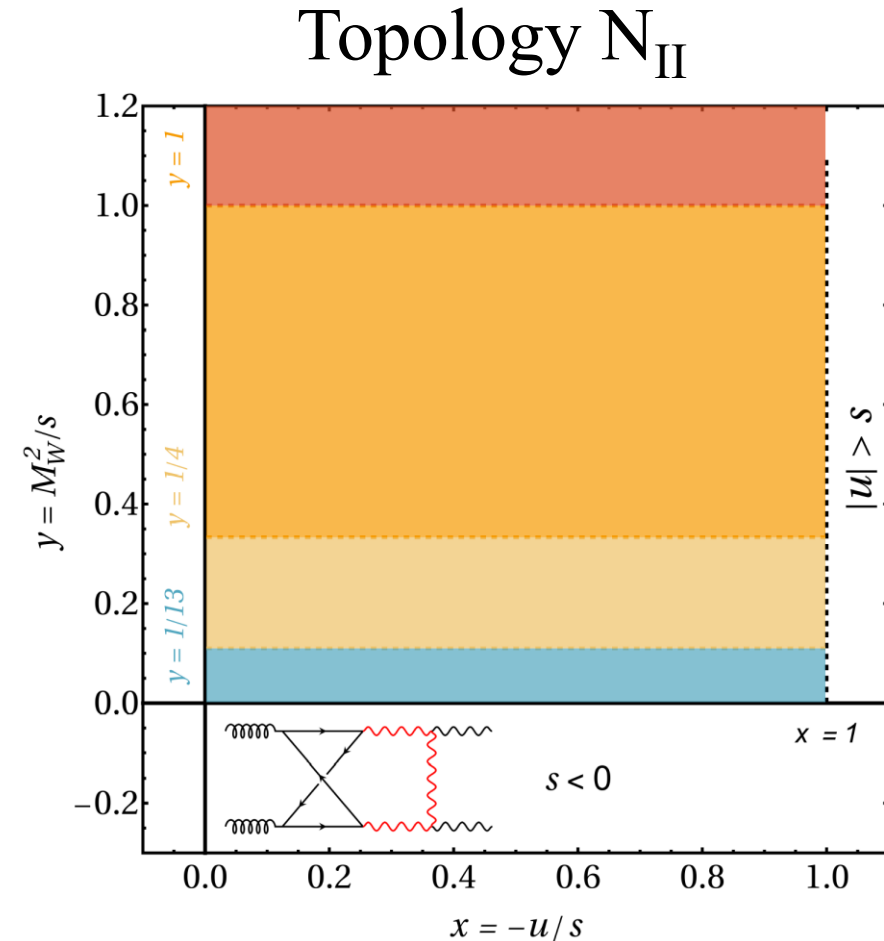
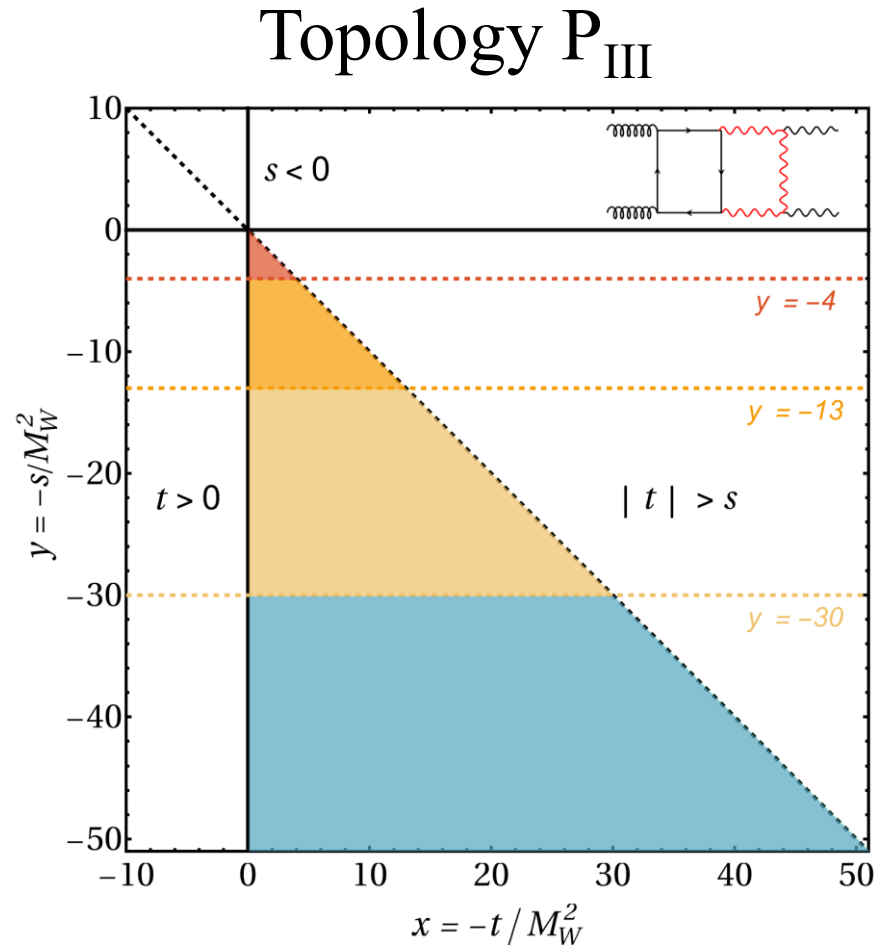
Fitting Procedure: Methodology

- Once we have the MIs we can evaluate them. GPL vs Chen.
- Overall goal: MCFM implementation. It is Crucial to find an efficient evaluation over the phase space.
- Idea: generate a grid of values in Mathematica and try a fitting strategy. Very efficient and reliable if successful, but control over the errors introduced by the approximation is mandatory.
- Exploit the underlying logarithmic structure of Chen Integrals. Fit real and Imaginary part separately:

$$\mathcal{K}_n^{(w)}(x, y) = \frac{\sum_{i+j=6} a_{i,j} x^i y^j}{\sum_{i+j=6} b_{i,j} x^i y^j} + i \frac{\sum_{i+j=6} c_{i,j} x^i y^j}{\sum_{i+j=6} d_{i,j} x^i y^j}$$

Fitting Procedure: Regions

- To prevent numerical instabilities given by real and pseudo thresholds, it's more efficient to break down the fitting by regions.

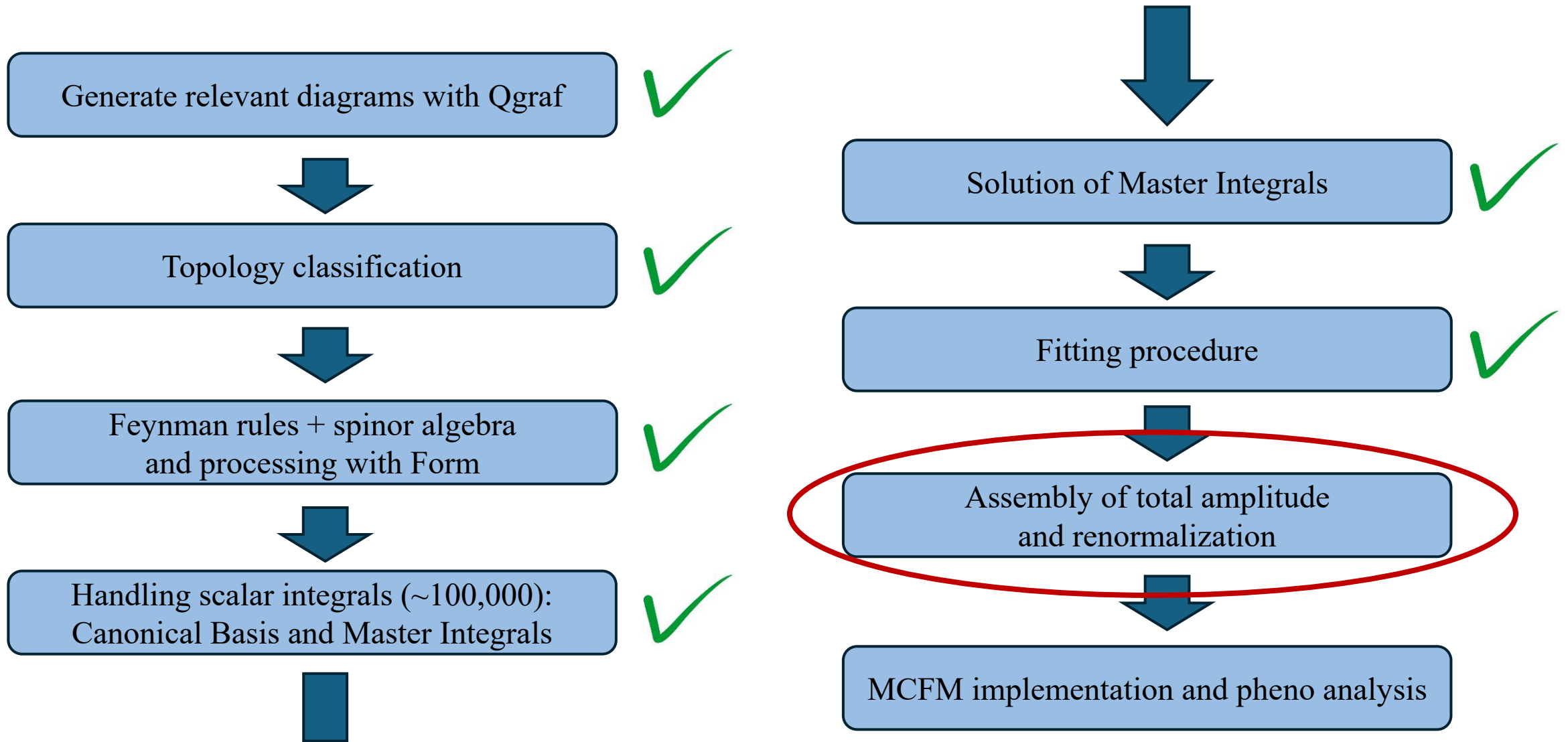


Fitting Procedure: Validation

- Preliminary analysis shows very good accuracy across the phase-space. Do the pseudo-thresholds cause any real issue? We find an estimate on the cross-section for the fully integrated effects!
- Idea: use NI as a testing ground. Most complex GPL topology $\sim 20,000$ GPL!
- We compare analytic $\sigma_{\text{NI}}^{(\mathcal{G})}$ vs fitting $\sigma_{\text{NI}}^{(\mathcal{K})}$ and extract a comparison from there. Expect similar behaviour from PIII and NII.

Ratio $\sigma_{\text{NI}}^{(\mathcal{K})} / \sigma_{\text{NI}}^{(\mathcal{G})}$	Uncertainty
1.00001	± 0.007

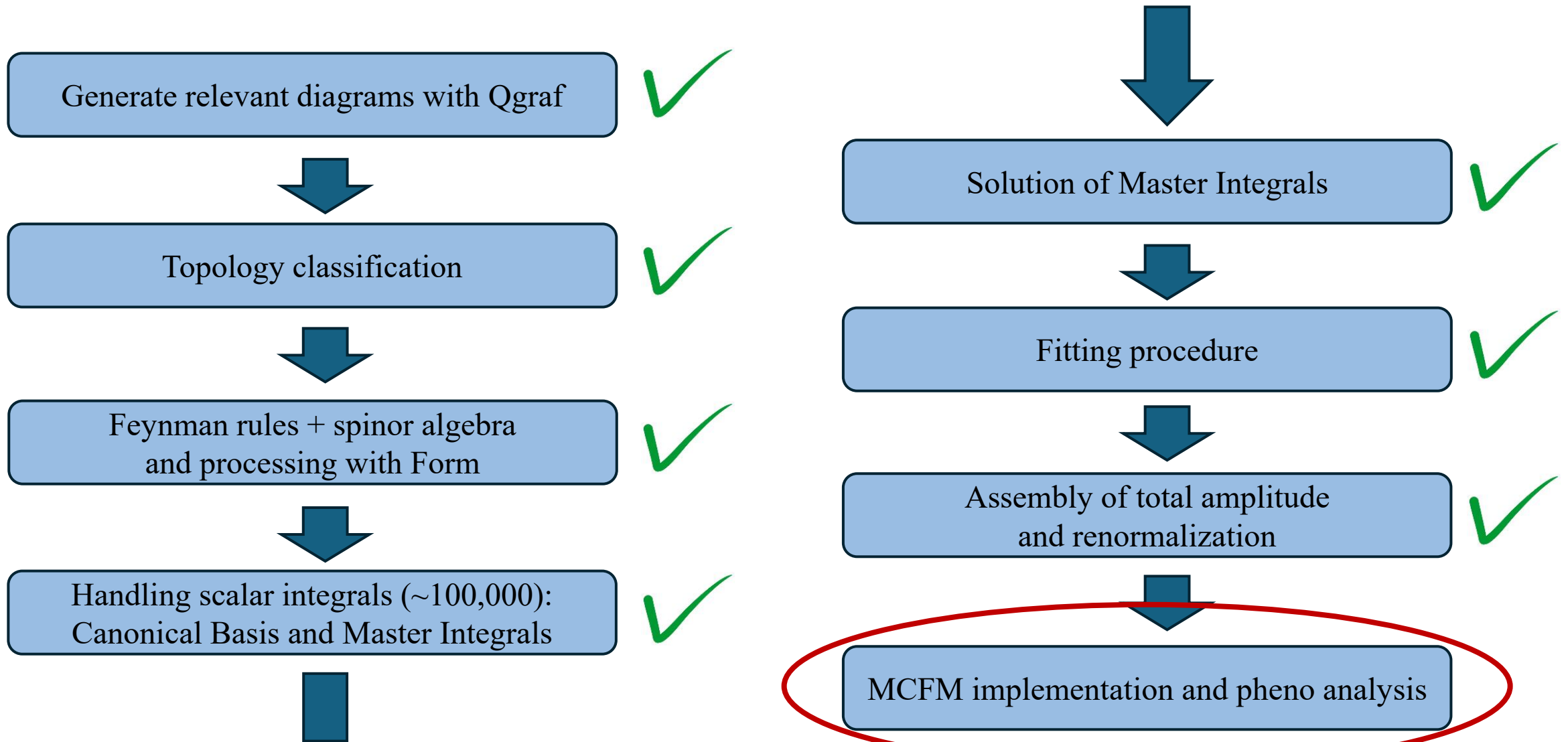
Calculation Workflow



Various validations for our results:

- Extensive numerical check for MIs with AMFlow.
- IR and UV analytic cancellation.
- Comparison for LO and NLO QED corrections with the literature.
- Scheme independence for γ_5 .
- Invariance under $u \leftrightarrow t$ and symmetry of rapidity distribution.

Calculation Workflow



MCFM Setup

➤ We present the results of our implementation for:

- 1) Total cross-section,
- 2) Invariant mass $m_{\gamma\gamma}$ distribution,
- 3) Transverse momentum p_T distribution,
- 4) Rapidity of the system $Y_{\gamma\gamma}$ distribution.

➤ It's phenomenologically interesting to distinguish between two different invariant mass regions:

“Higgs” window

$$M_Z < m_{\gamma\gamma} < 2M_W$$

“High-energy” window

$$2M_W < m_{\gamma\gamma} < 500\text{GeV}$$

Results: Cross-section

Correction of $\sim 2.43\%$ in the Higgs window

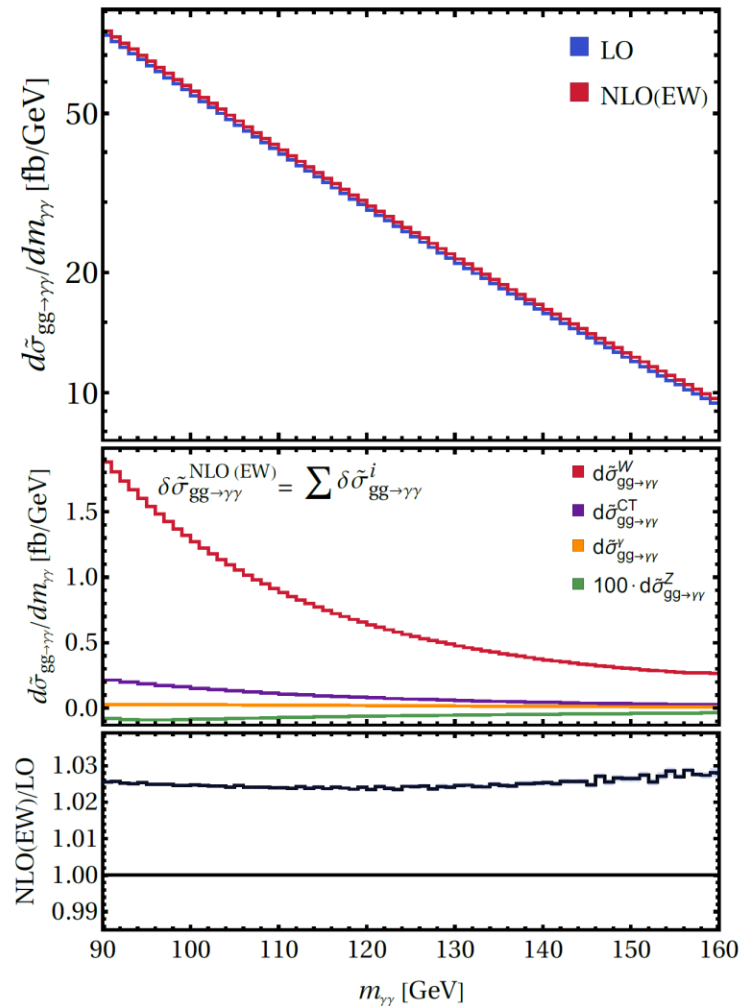
Order	Cross Section [pb]	Uncertainty [pb]
$\tilde{\sigma}_{gg \rightarrow \gamma\gamma}^{(\text{LO})}$	2.176	± 0.002
$\tilde{\sigma}_{gg \rightarrow \gamma\gamma}^{(\text{NLO-EW})}$	2.229	± 0.003

Correction of $\sim 1.62\%$ in the high energy region

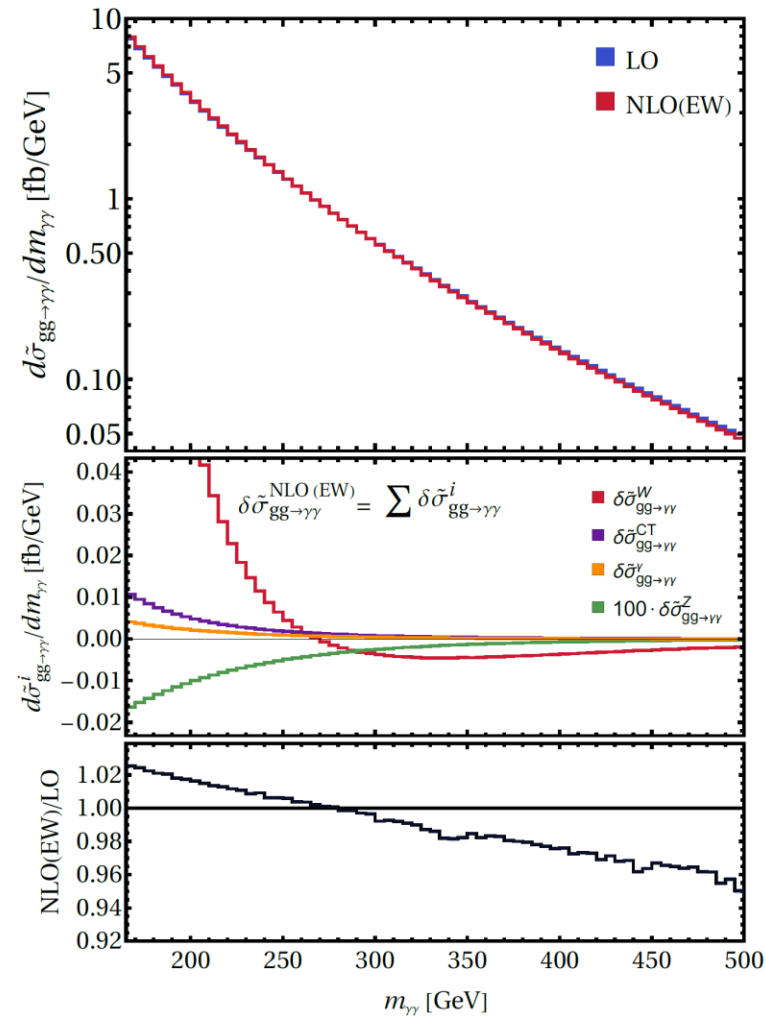
Order	Cross Section [pb]	Uncertainty [pb]
$\tilde{\sigma}_{gg \rightarrow \gamma\gamma}^{(\text{LO})}$	0.39119	± 0.00001
$\tilde{\sigma}_{gg \rightarrow \gamma\gamma}^{(\text{NLO-EW})}$	0.39751	± 0.00006

Results: Invariant Mass Distributions

“Higgs” window

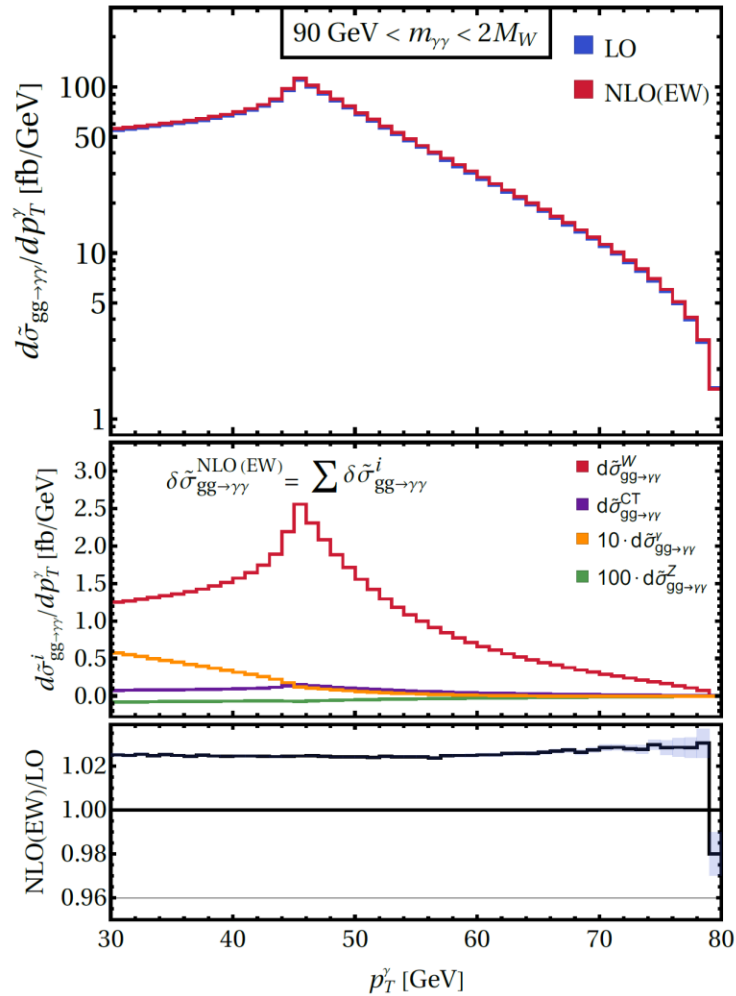


“High-energy” window

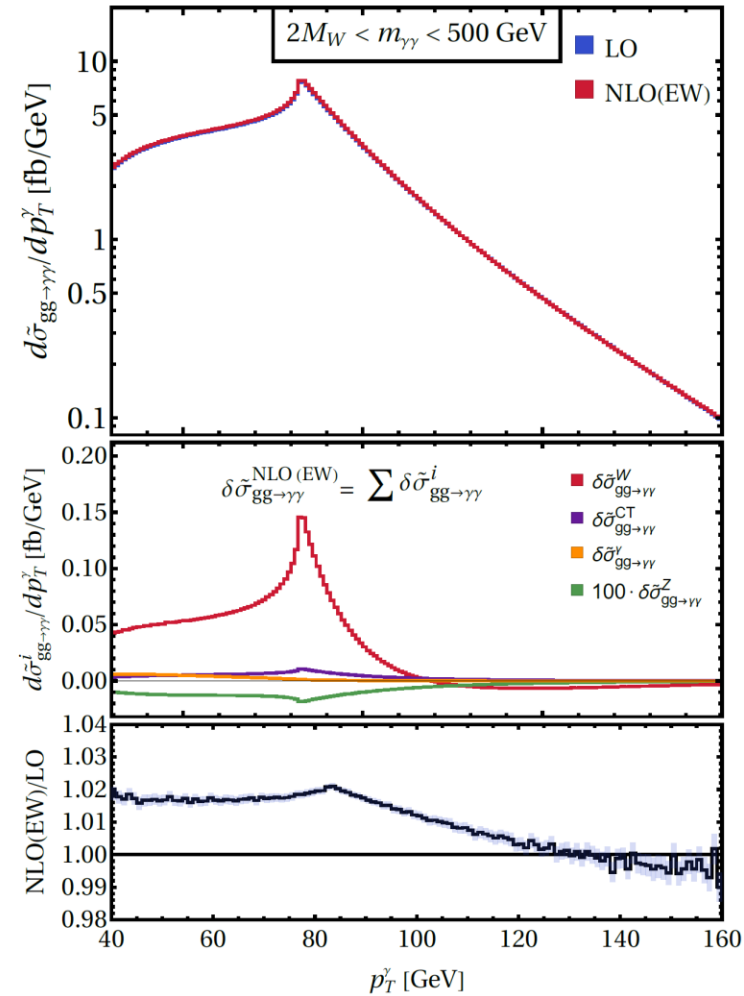


Results: Transverse Momentum Distributions

“Higgs” window

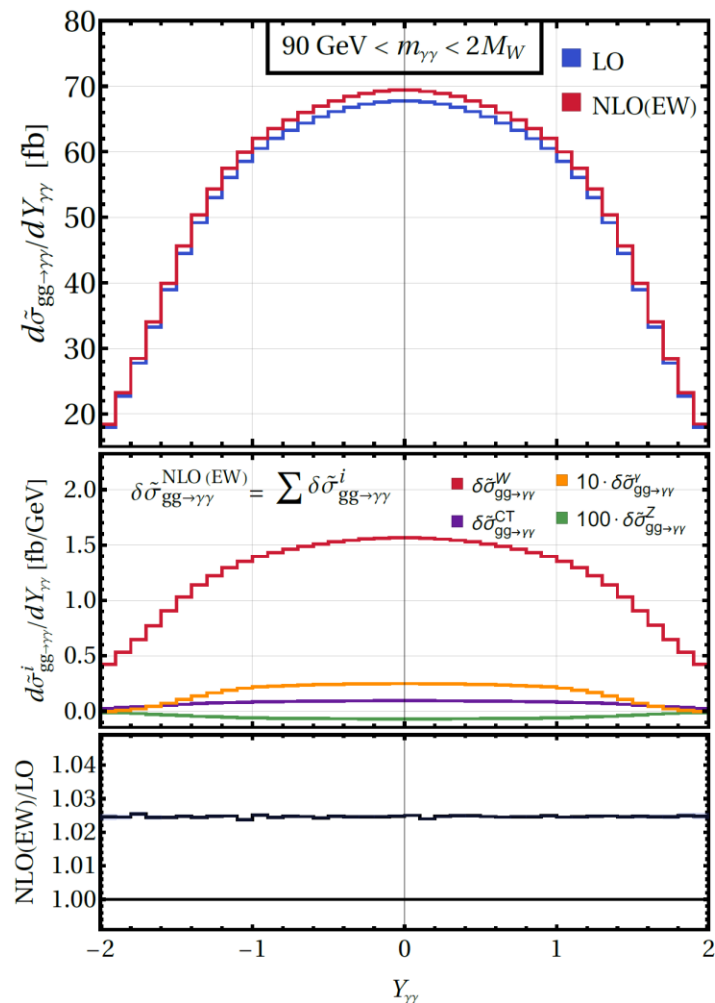


“High-energy” window

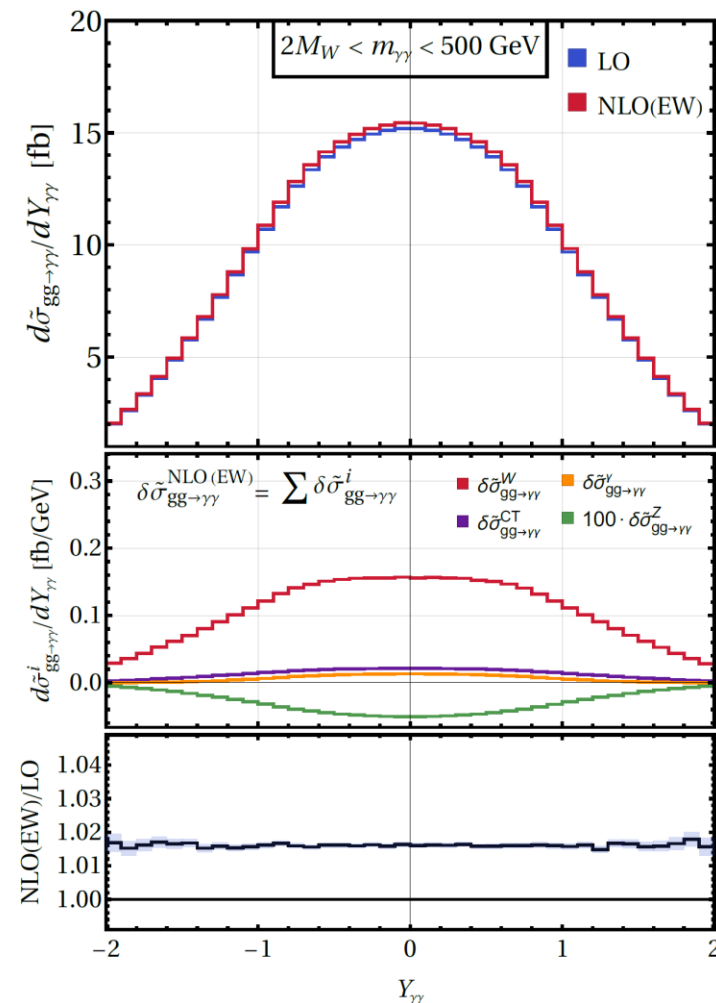


Results: Rapidity Distributions


“Higgs” window



“High-energy” window



Summary and Future Plans

- Evaluation for the first time of EW correction to $gg \rightarrow \gamma\gamma$ for light quarks. We observed impact of 2.43% in the Higgs central region, and a more modest impact of 1.62% in the high energy region.
- Differential distributions shows corrections at the percent level, with more structure happening in the high energy region. Symmetry of the rapidity distribution continues a non-trivial check for our results.
- We expect the high energy region to be sensitive to the top-quark correction  Work in progress!
- With the inclusion of the top-quark corrections, we can calculate the EW correction to the Higgs interference and ultimately extrapolate the corresponding corrections to the Higgs Width

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Thank you for listening!