

The Apex of RHIC Physics

Resolving the Strong Force

May 11–15, 2026

Workshop Report: Baryon Junction & Stopping

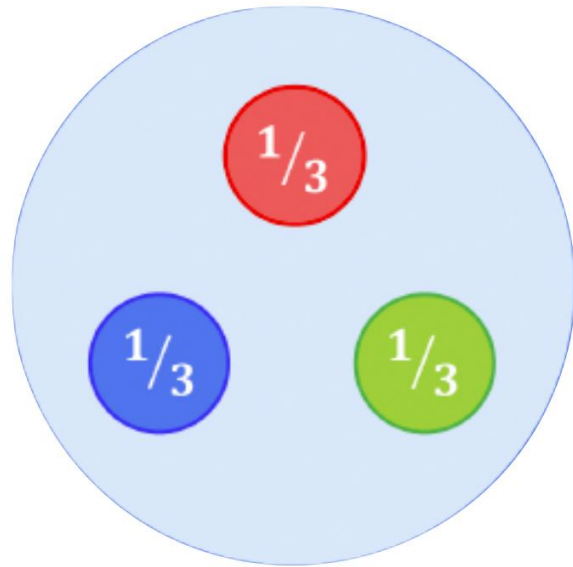
2026 RHIC/AGS Users' meeting

Chun Yuen Tsang (ANL) and Prithwish Tribedy

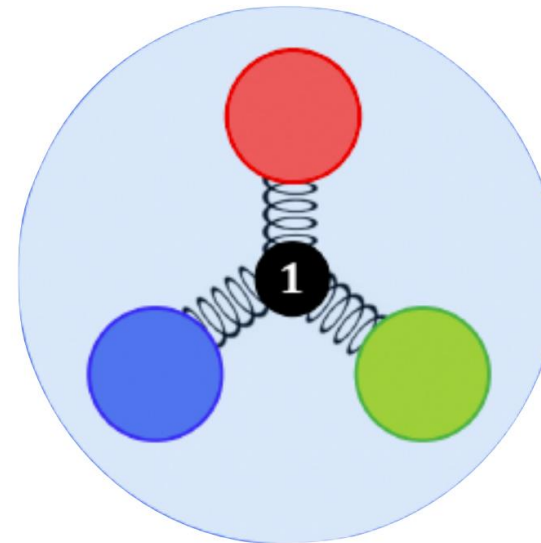
May 13, 2026

What carries baryon number, valence quark vs. baryon junction?

Naïve quark picture



Baryon junction picture [1, 2]



- [1]: Nucl. Phys. B 85, 442–460 (1975).
[2]: Nucl. Phys. B 123, 507–545 (1977)

7 talks in the Baryon Junction Section

Mon 11/05

09:00	Current status of the baryon junction <i>Dmitri Kharzeev et al.</i>	09:00 - 10:00
10:00	Generalized parton distributions and their implications for observables relevant to baryon junction dynamics <i>Simonetta Liuti</i>	10:00 - 10:30

11:00	Probing gluon density via deep inelastic scattering, and its implications on baryon junctions <i>Prithwish Tribedy</i>	
	Flavor-independent rapidity slope in baryon transport <i>Zhangbu Xu et al.</i>	11:30 - 12:00

14:00	Measuring baryon production at HERA, and its potential on baryon junctions <i>Gage Tustin</i>	
	Probing the Baryon Junction at the CLAS12 Experiment <i>Wenliang Li</i>	14:30 - 15:00
15:00	Finding the baryon number within the proton <i>Spencer Klein</i>	15:00 - 15:30

History of baryon junction: Gauge invariant baryon operator

How to construct a gauge-invariant wave function of a baryon?



G. Rossi, G. Veneziano 1977



Table IIa

Simplest mesons and baryons : colour structure and string picture

HADRON	GAUGE INVARIANT OPERATOR	STRING PICTURE
$M_2 = q\bar{q}$ meson	$\bar{q}^{j_2}(x_2) \left[P \exp \left(i g \int_{x_1}^{x_2} A_\mu dx^\mu \right) \right]_{j_2}^{j_1} q_{j_1}(x_1)$	
$M_0 =$ quarkless meson	$\text{Tr} \left[P \exp \left(i g \oint A_\mu dx^\mu \right) \right]$	
$B_3 = qqq$ baryon	$\epsilon^{j_1 j_2 j_3} \left[P \exp \left(i g \int_{x_1}^x A_\mu dx^\mu \right) q(x_1) \right]_{j_1} \left[P \exp \left(i g \int_{x_2}^x A_\mu dx^\mu \right) q(x_2) \right]_{j_2} \left[P \exp \left(i g \int_{x_3}^x A_\mu dx^\mu \right) q(x_3) \right]_{j_3}$	

Dmitri Kharzeev

Baryon-number – flavor separation in high energy collisions

Physics Letters B 378 (1996) 238–246

Can gluons trace baryon number?

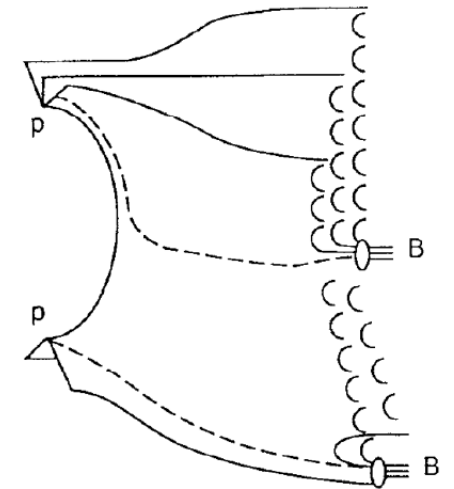
D. Kharzeev

Theory Division, CERN, CH-1211 Geneva, Switzerland
and Fakultät für Physik, Universität Bielefeld, D-33501 Bielefeld, Germany

Received 15 March 1996

Editor: R. Gatto

$$E_B \frac{d^3 \sigma^{(1)}}{d^3 p_B} = 8\pi G_p^M(0) G_p^P(0) f_B^{MP}(m_t^2) \left(\frac{\sqrt{s} m_t}{s_0} \right)^{\alpha_0^J(0) + \alpha_P(0) - 2} \times \left(\exp[y^*(\alpha_P(0) - \alpha_0^J(0))] + \exp[-y^*(\alpha_P(0) - \alpha_0^J(0))] \right).$$



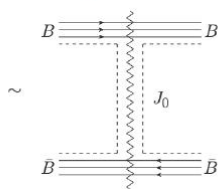
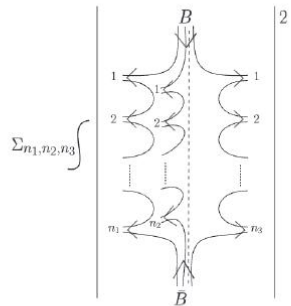
- The only way to construct a gauge-invariant baryon operator involves products of exponentials of path integrals of the gluon potential, all connecting to a common central point.
- D. Kharzeev then showed that such a topological configuration carries baryon number.

Theoretical considerations that give rise to baryon junction

Dmitri Kharzeev

Baryon-number – flavor separation from the topological expansion of QCD

D. Frenklakh, DK, G. Rossi, G. Veneziano, arXiv:2405.04569



Three ingredients:
Duality
Topological expansion
Feynman-Wilson gas

Baryons junctions and duality

A POSSIBLE DESCRIPTION OF BARYON DYNAMICS IN DUAL AND GAUGE THEORIES

G.C. Rossi ^{*)}

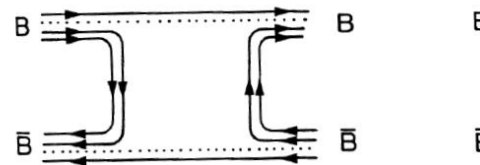
and

1977

G. Veneziano ^{**)}

CERN - Geneva

No annihilation:



(b)

The junction-anti-junction intercept

$$\alpha_{\mathbb{P}} = 1 + C_{RL}$$

$$\alpha_{\mathbb{J}_0} = (2\alpha_{\mathbb{B}} - 1) + 3(1 - \alpha_{\mathbb{P}}) + 3(1 - \alpha_{\mathbb{R}}) - C_3 \sim 0.5 - 3C_{RL} - C_3$$

$$\alpha_{\mathbb{J}_0} \simeq 0.26$$

beam rapidity dependence $e^{(\alpha_{\mathbb{J}_0} + \alpha_{\mathbb{P}} - 2)Y/2} = e^{-0.66 Y/2}$

D. Frenklakh, DK, G. Rossi, G. Veneziano, arXiv:2405.04569

- Difficulty in the duality of t- and u-channels for baryons implies that baryon number is carried by a separate entity.
- Derived cross-section dependence using an analogy between high-energy interactions and a classical gas.

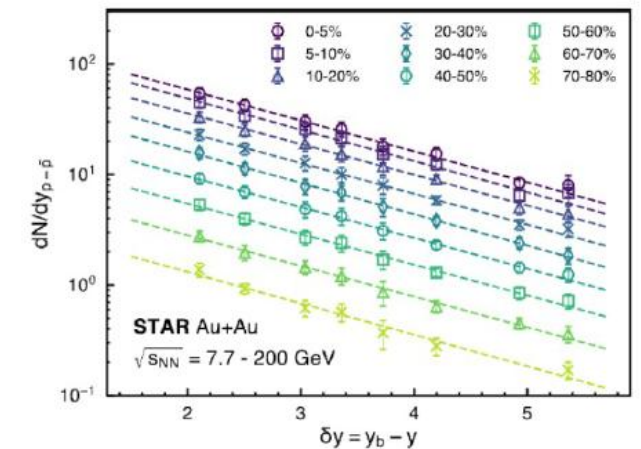
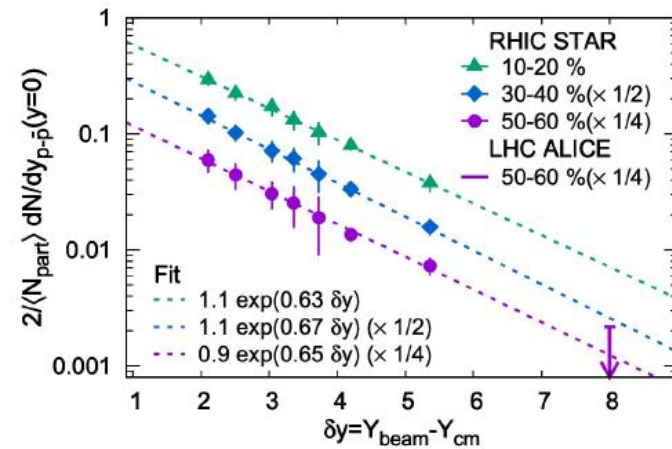
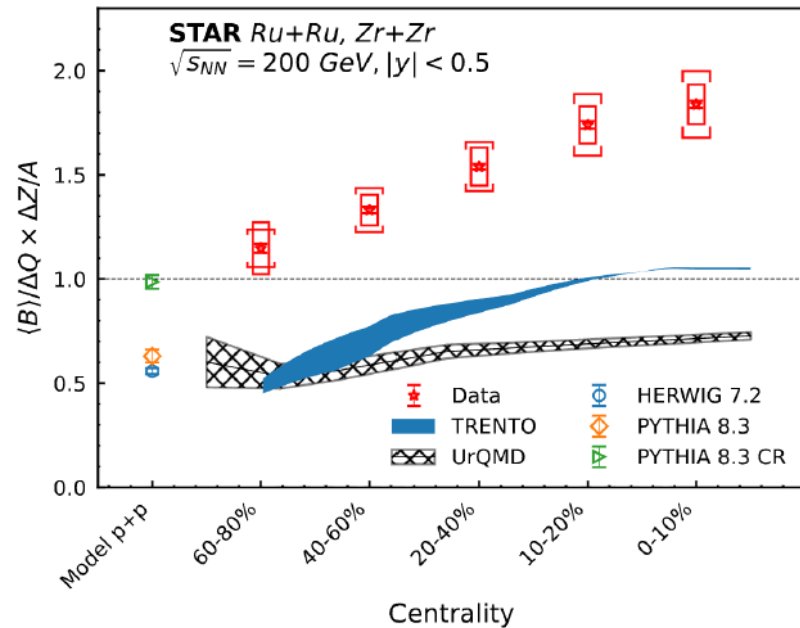
Current evidence for baryon junction transport

STAR Collaboration, arXiv:2408.15441;
Science, to appear

Dmitri Kharzeev

Search for baryon junctions in photonuclear processes and isobar collisions at RHIC

James Daniel Brandenburg,¹ Nicole Lewis,^{1,*} Prithwish Tribedy,¹ and Zhangbu Xu¹
¹Physics Department, Brookhaven National Laboratory, Upton, NY 11973, USA
(Dated: August 25, 2022)

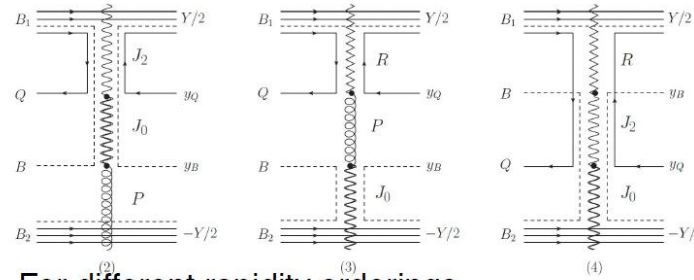


- Net-baryon/Net-charge difference between isobars at mid-rapidity: baryon and electric charge transport are decoupled.
- Net-proton vs rapidity difference between beam and fragment: Consistent with predictions of baryon junction stopping from Regge theory.

Future tests for Baryon Junction

Dmitri Kharzeev

Combined Baryon-number-charge distribution:



For different rapidity orderings

$$F_{(B,Q)}^{(1)} \propto e^{(\alpha_P + \alpha_{J_2} - 2)\frac{Y}{2}} e^{(\alpha_R - \alpha_{J_2})y_B} e^{(\alpha_P - \alpha_R)y_Q},$$

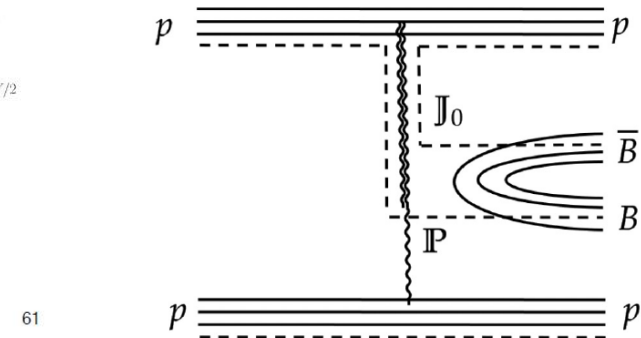
$$F_{(B,Q)}^{(2)} \propto e^{(\alpha_P + \alpha_{J_2} - 2)\frac{Y}{2}} e^{(\alpha_{J_0} - \alpha_{J_2})y_Q} e^{(\alpha_P - \alpha_{J_0})y_B},$$

$$F_{(B,Q)}^{(3)} \propto e^{(\alpha_R + \alpha_{J_0} - 2)\frac{Y}{2}} e^{(\alpha_P - \alpha_R)y_Q} e^{(\alpha_{J_0} - \alpha_P)y_B},$$

$$F_{(B,Q)}^{(4)} \propto e^{(\alpha_R + \alpha_{J_0} - 2)\frac{Y}{2}} e^{(\alpha_{J_0} - \alpha_{J_2})y_Q} e^{(\alpha_{J_2} - \alpha_R)y_B}.$$

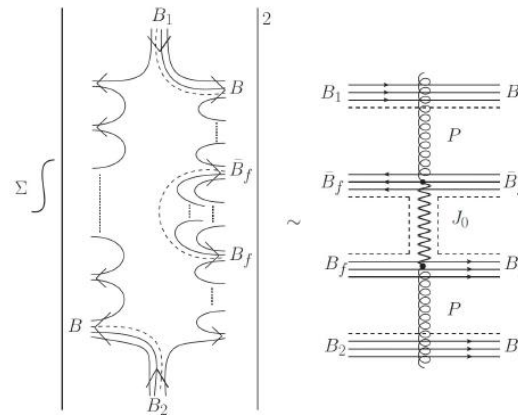
Future tests at RHIC and LHC

Extract the slope of J_0 trajectory in diffractive production of baryon-antibaryon pairs (or tetraquarks):



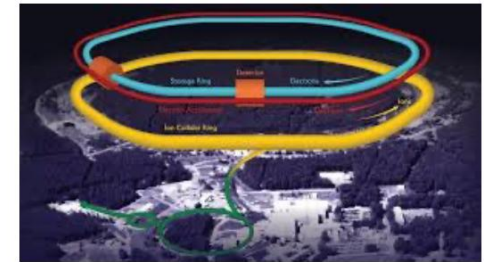
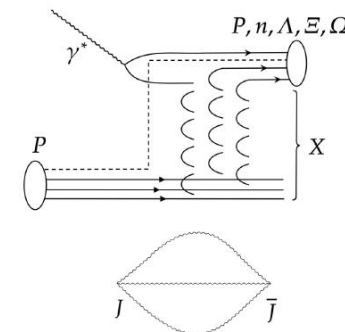
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Baryon-antibaryon correlations in rapidity:



$$\rho_{B,\bar{B}}(\Delta y) \sim e^{-|\Delta y|(\alpha_P - \alpha_{J_0})}$$

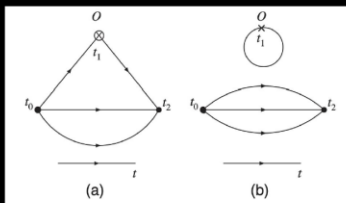
Baryon junctions at EIC



Generalized Parton Distributions, disconnected sea quarks, and baryon junction phenomenology

Simonetta Liuti

Including connected and disconnected contributions in GPD phenomenology



$$F_1^q(t) = A_{10}^q(t) = \int_{-1}^1 dx H_q(x, \xi, t),$$

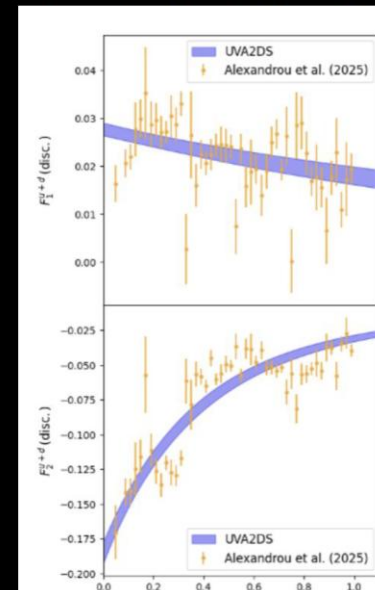
$$F_2^q(t) = B_{10}^q(t) = \int_{-1}^1 dx E_q(x, \xi, t)$$

$$F_1^q(t) = F_1^{q,C} + F_1^{q,D} = \int_{-1}^1 dx H_q^C + \int_{-1}^1 dx H_q^D$$

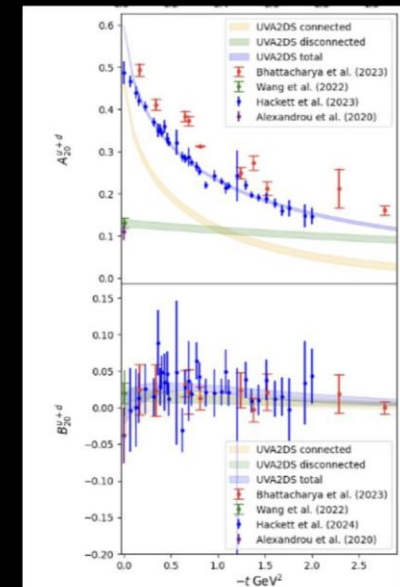
$$H_q^C = \frac{1}{2}(H_{q_v} + H_{q_{sea}}^C) = \frac{1}{2}(H_q^{-,C} + H_q^{+,C})$$

$$H_q^D = \frac{1}{2}(H_q^{-,D} + H_q^{+,D}).$$

"-" disconnected n=1 Mellin moment



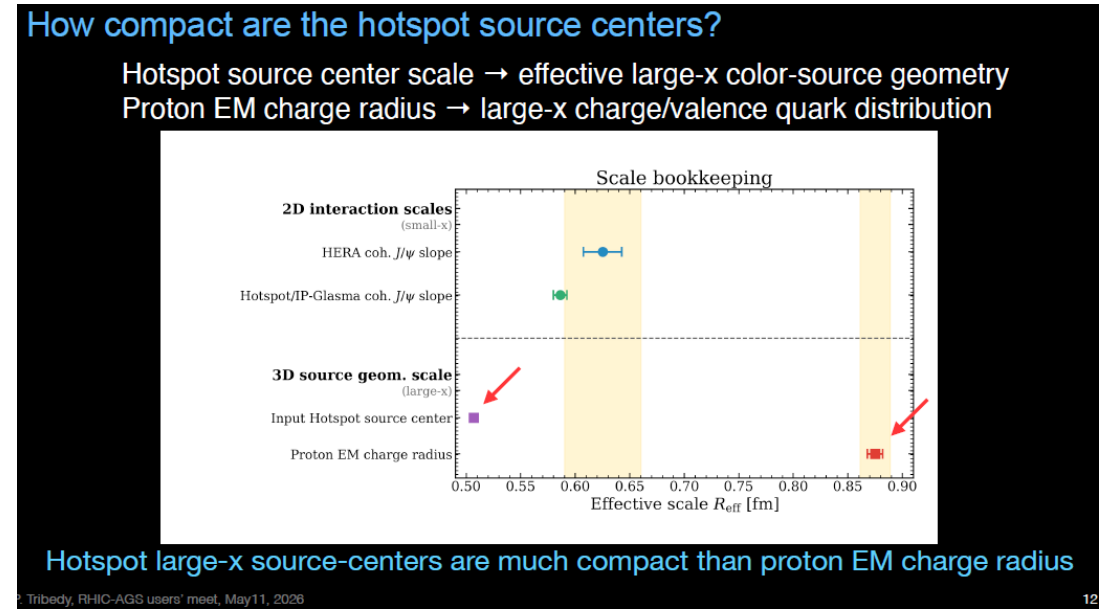
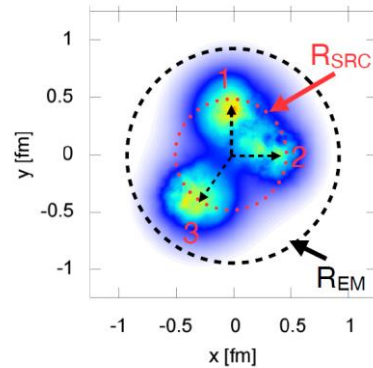
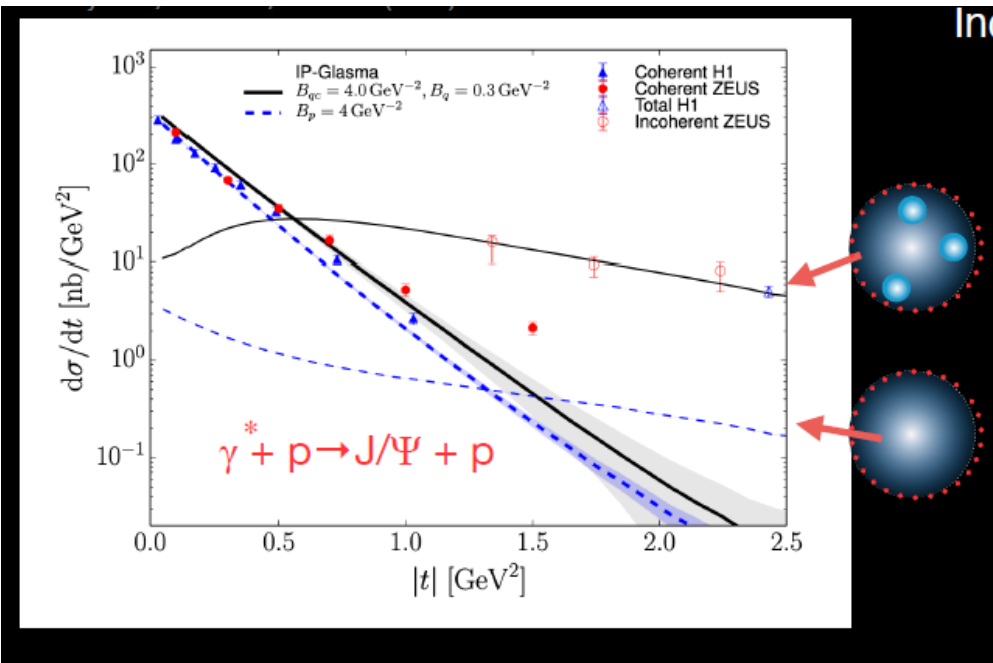
"+" connected and disconnected n=2 Mellin moment



- First quantitative study of the impact of disconnected contributions to the parton distribution functions.
- Junction carries baryon number. Provide sea quarks input for recombination into mid-rapidity baryons.

Phenomenological model inspired by baryon junction

Prithwish Tribedy



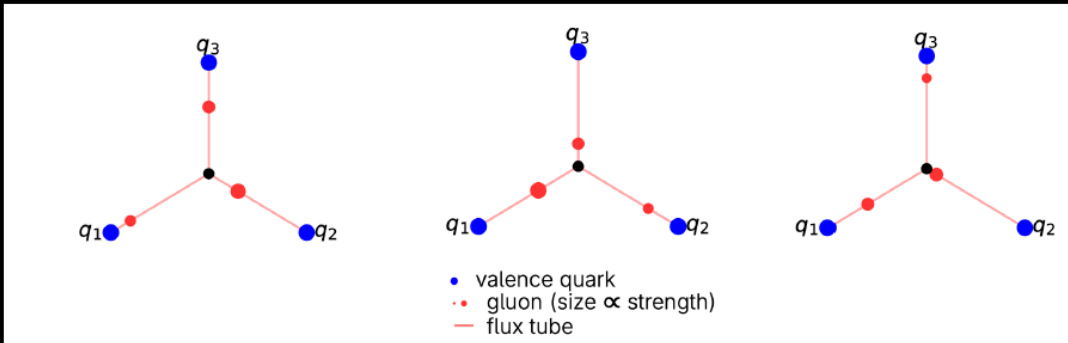
- Cross-section data of coherent small x $\gamma+p$ interactions agree with the hotspot model: color source density peaks at three points.
- Problem: Effective size of the same model at large-x \ll proton radius?

Phenomenological model inspired by baryon junction (cont.)

Prithwish Tribedy

How the stringy-proton geometry is built

Work in progress: Roch, Mäntysaari, Tribedy, Schenke, Shen & Zhao



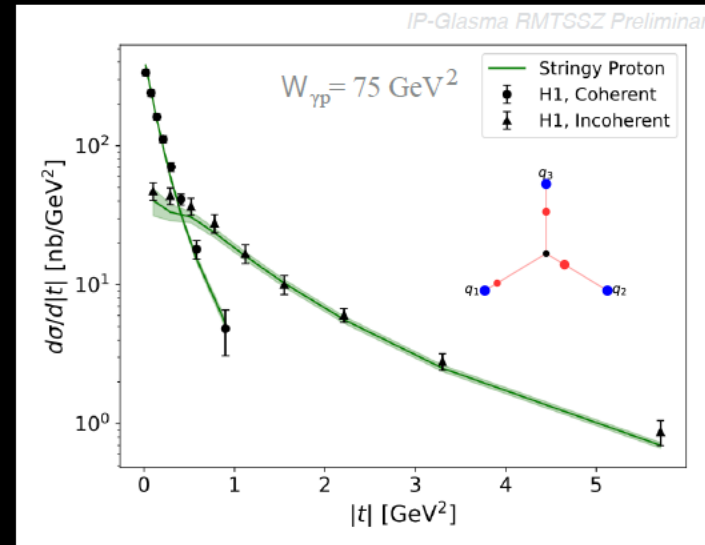
- 1) Sample three source/charge centers, 2) Connect them through a Y-like geometry
- 3) Place gluonic support along the connections, 4) Add strength fluctuations

Same large source scale, different place for the glue

P. Tribedy, RHIC-AGS users' meet, May 11, 2026

Stringy proton describes HERA J/ψ coherent/incoherent data

Work in progress: Roch, Mäntysaari, Tribedy, Schenke, Shen & Zhao



Coherent: average geometry described

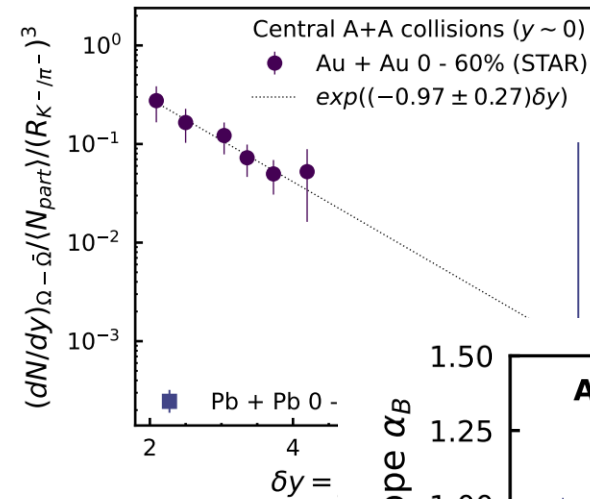
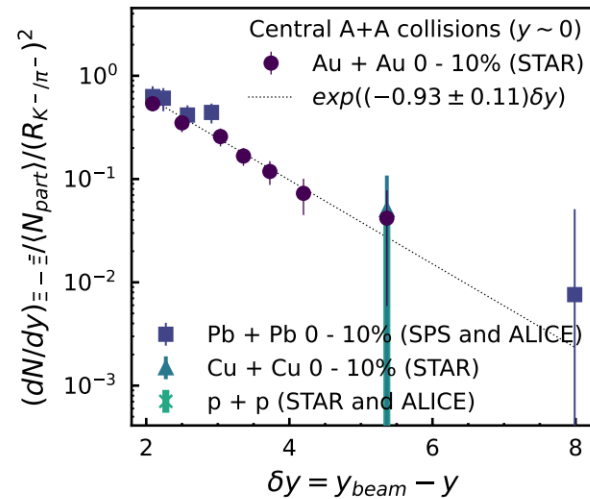
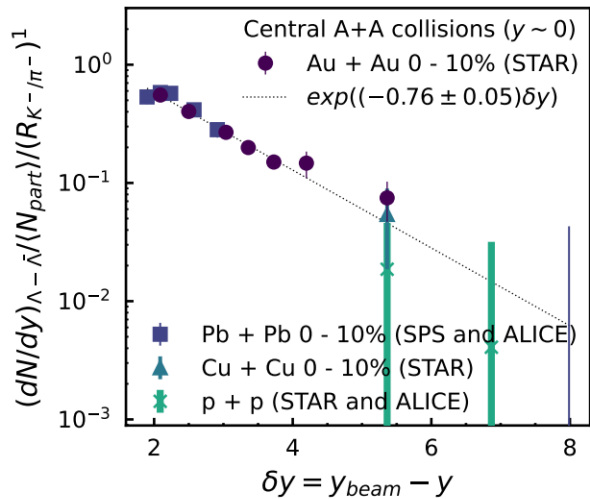
Incoherent: fluctuations also described

Changing where the glue lives does not destroy the HERA description

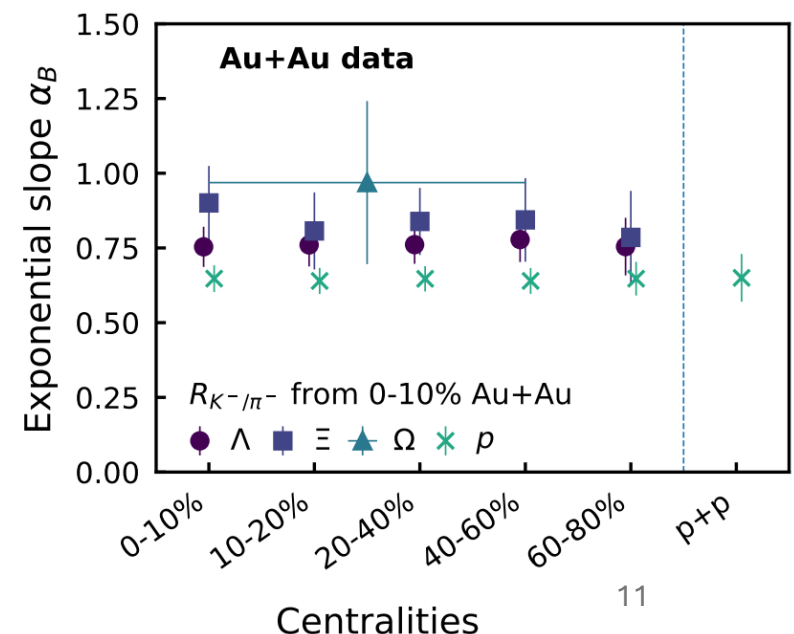
- Stringy proton model: Color sources are distributed along the three lines of the Y-shape.
- This model describes both large and small x data.

Another test of junction picture: Flavor independent

Zhangbu Xu

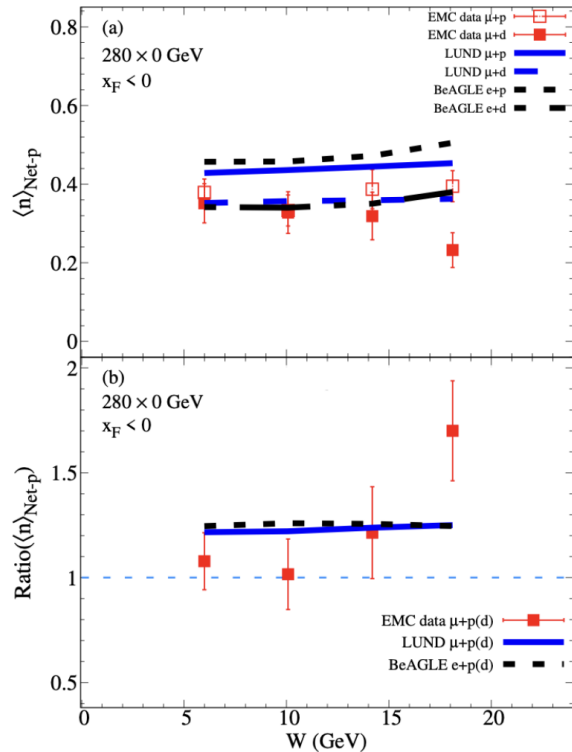


- Net- Λ , net- Ξ and net- Ω yields (after correcting for s-quark production rate) follows an exponential dependence on rapidity.
- Exponential slopes are consistent with each other and net-proton.
- Baryon stopping is independent of quark flavor: Consistent with the baryon junction picture.

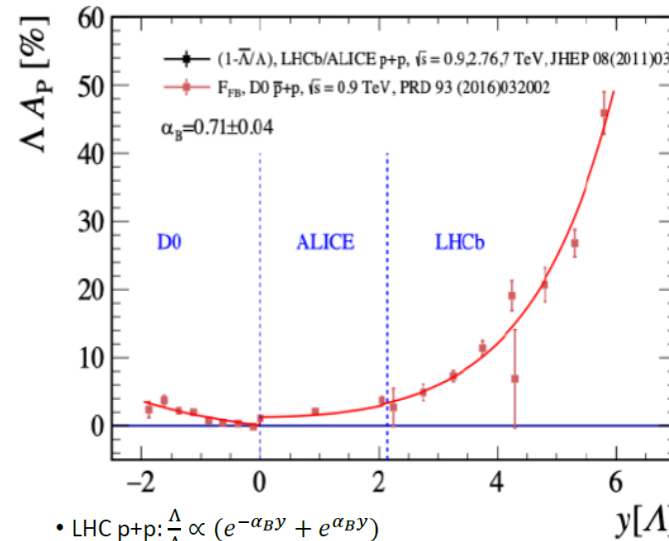


Another test of junction picture: Flavor independent (Cont.)

Zhangbu Xu



Combined ALICE/LHCb/D0 $\bar{\Lambda}/\Lambda$



LHCb Λ_b

Anantha Nair (KSU)
Existing high-statistic data at 13TeV

- LHC p+p: $\frac{\Lambda}{\bar{\Lambda}} \propto (e^{-\alpha_B y} + e^{\alpha_B y})$
- Tevatron $\bar{p}+p$: $\frac{\bar{\Lambda}}{\Lambda} \propto (e^{-\alpha_B y} - e^{\alpha_B y})$

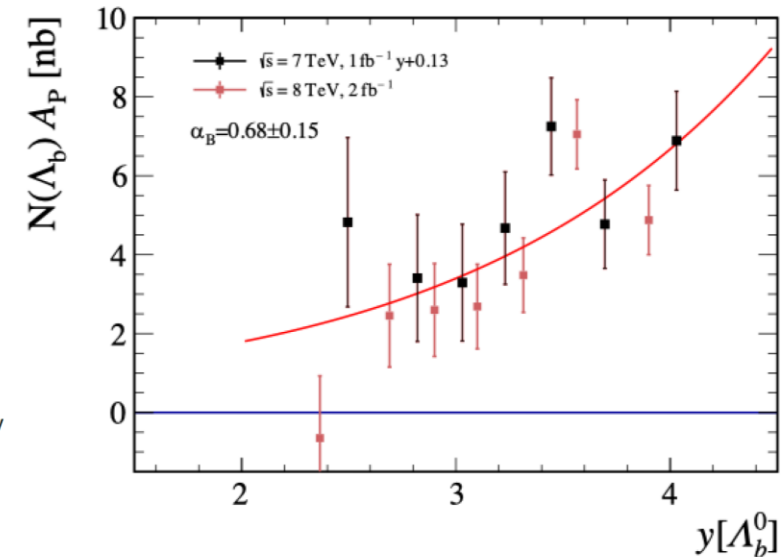


FIG. 7. The W dependence of the net-proton for $\mu+p$ and $\mu+d$ at 280-GeV muon on fixed targets are shown in panel (a). The ratios between $\mu+p$ and $\mu+d$ are presented in panel (b). Data and LUND model calculations are extracted from Ref. [15].

- Measured net-proton yield of e+p and e+deuteron are consistent at $x_F < 0$: Contradicting the Lund model!
- Exponential slope of net- Λ and Λ_b from various experiments agree with that from STAR.
- Conclusion: Baryon stopping is independent of quark flavor.

Another source of data: HERA

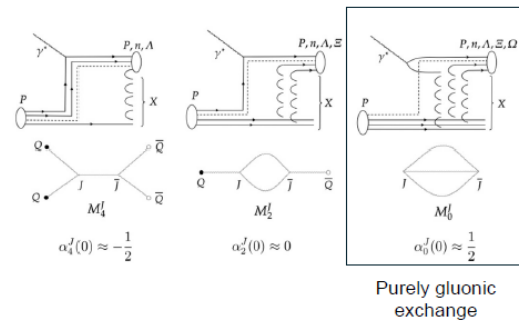
Baryon Junction Motivation

- David Frenklakh and Dmitri Kharzeev gave us a quick calculation based on their gluon junction framework in this publication and estimate a baryon asymmetry in DIS of 0.03 ± 0.02 within the accessible kinematic range at HERA

Signatures of baryon junctions in semi-inclusive deep inelastic scattering

David Frenklakh ^{a,b}, Dmitri E. Kharzeev ^{a,b,c}, Wenliang Li ^b

<https://doi.org/10.1016/j.physletb.2024.138680>



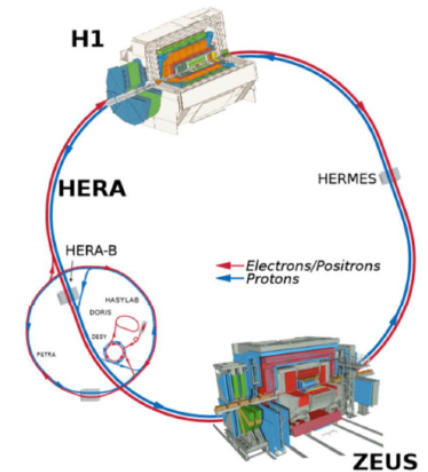
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- Based on a $3 \rightarrow 3$ Forward scattering in Regge Theory framework
- Regge intercept for purely gluonic junction exchange cannot be calculated from first principles, but they estimated it from the Au-Au RHIC Beam Energy Scan data, hence the large uncertainty
- Not a precise estimate but gives us a ballpark expectation

HERA

- The only high energy ep collider accelerator ever operated
- Ran from 1992-2007 at DESY in Hamburg, Germany with nominal electron/positron energy of 27.6 GeV and proton energy of 920 GeV ($\sqrt{s} = 319$ GeV)
- Served four experiments
 - Hermes: Fixed target with electron beam
 - HERA-B: Fixed target with proton beam
 - ZEUS: General purpose hermetic detector
 - H1: General purpose hermetic detector
- Delivered $\sim 1 \text{ fb}^{-1}$ of total integrated luminosity to H1/ZEUS, perfect starting point to get a head start on what baryon junction signatures in EIC may look like

Gage Tustin



Hadron-Electron Ring Accelerator (HERA) complex, figure from [1]

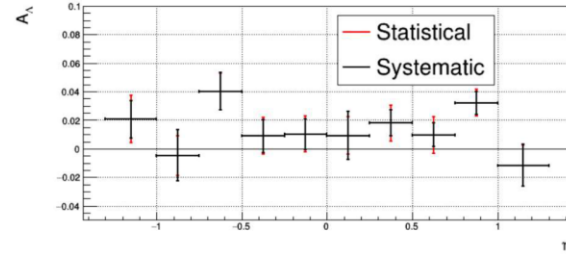
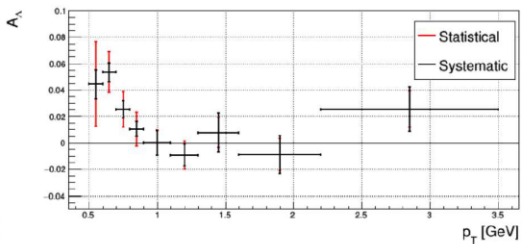
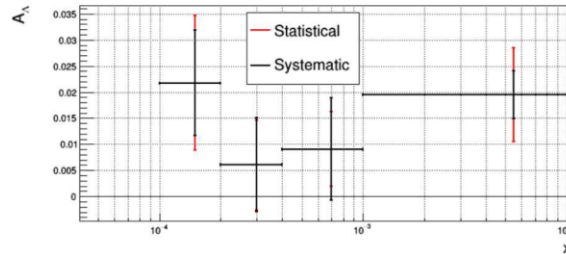
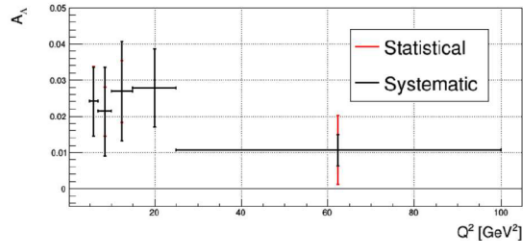
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- HERA: $e+p$ data at high beam energy.
- Net- Λ can be reconstructed: complement aforementioned flavor independent test on baryon stopping.

Another source of data: HERA

Gage Tustin

Asymmetry Distributions



- $\Lambda - \bar{\Lambda}$ Asymmetry defined as this ratio of cross sections, many factors cancel out

$$A_{\Lambda} = \frac{\sigma_{vis}(ep \rightarrow e'\Lambda X) - \sigma_{vis}(ep \rightarrow e'\bar{\Lambda} X)}{\sigma_{vis}(ep \rightarrow e'\Lambda X) + \sigma_{vis}(ep \rightarrow e'\bar{\Lambda} X)} = \frac{\frac{N_{\Lambda}}{\epsilon_{rec,\Lambda}} - \frac{N_{\bar{\Lambda}}}{\epsilon_{rec,\bar{\Lambda}}}}{\frac{N_{\Lambda}}{\epsilon_{rec,\Lambda}} + \frac{N_{\bar{\Lambda}}}{\epsilon_{rec,\bar{\Lambda}}}}$$

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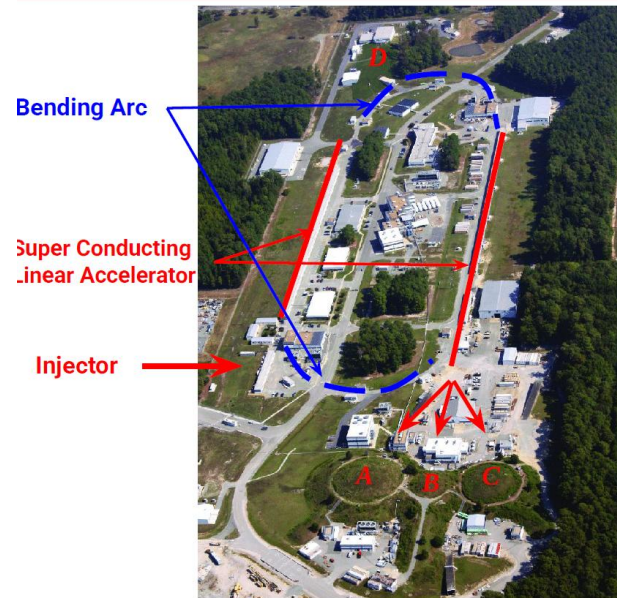
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- Data analysis in progress. Preliminary results show non-zero net- Λ .
- Efficiency correction TBD.

Proposal to use CLAS12 experiment for baryon junction test.

Wenliang Li

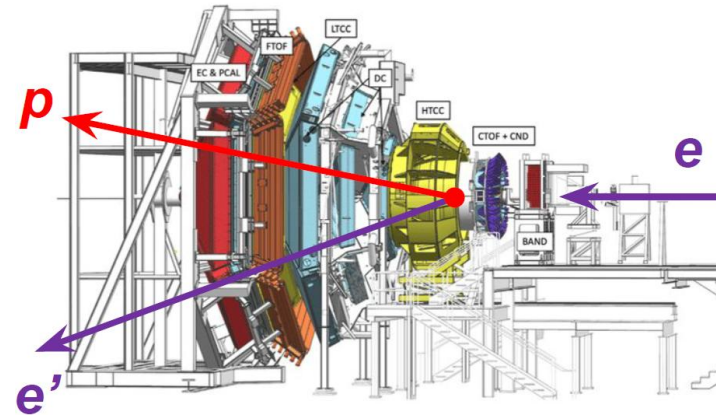
Jefferson Lab and Four Experimental Halls



- **Facility cost:**
 - Two Superconducting LINAC
 - Four experimental Halls
 - Electron beam energy up to 12 GeV
 - Duty cycle~99%
- **Peak Power Consumption: 1.2 M Watts**
- **Standard 6 month/year operation**
- **Four halls with different objectives:**
 - Hall A: preparing Moller experiment
 - Hall B: large acceptance spectrometer
 - Hall C: Narrow acceptance high precision spectrometers
 - Hall D: real photon beam,

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Experimental configuration



Beam and target configuration:

- Upto 12 GeV e beam
- unpolarized LH2 or LD target

Detector stack

- Time-of-flight: CTOF, FTOF
- PID: HTCC, LTCC
- Calorimeter: EC, PCAL
- Tracker: DC
- Neutron detection: CND

Coverage:

- **Charged hadrons:** 5-125° coverage
- **Neutrons:** 40-120°, with momentum 0.2-1 GeV/c

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- JLab provides another source of e+p and e+d collision data.
- Test the consistency of net-proton yields between e+p and e+d collisions with improved statistical significance.

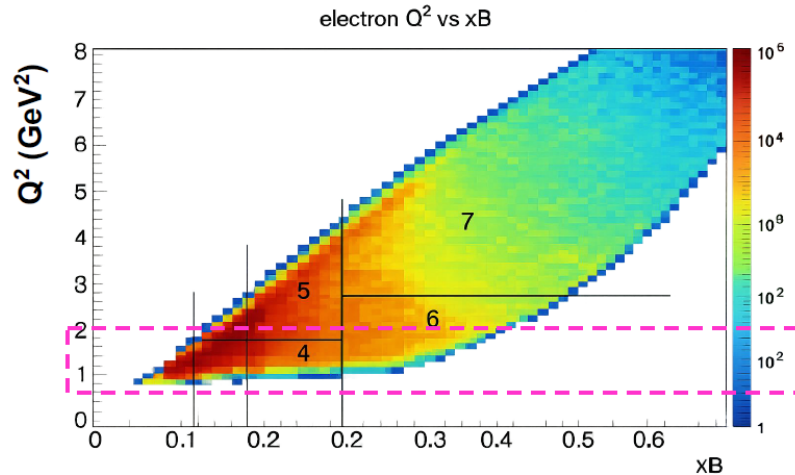
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Details of CLAS12 experiment

Wenliang Li

Data Set Information

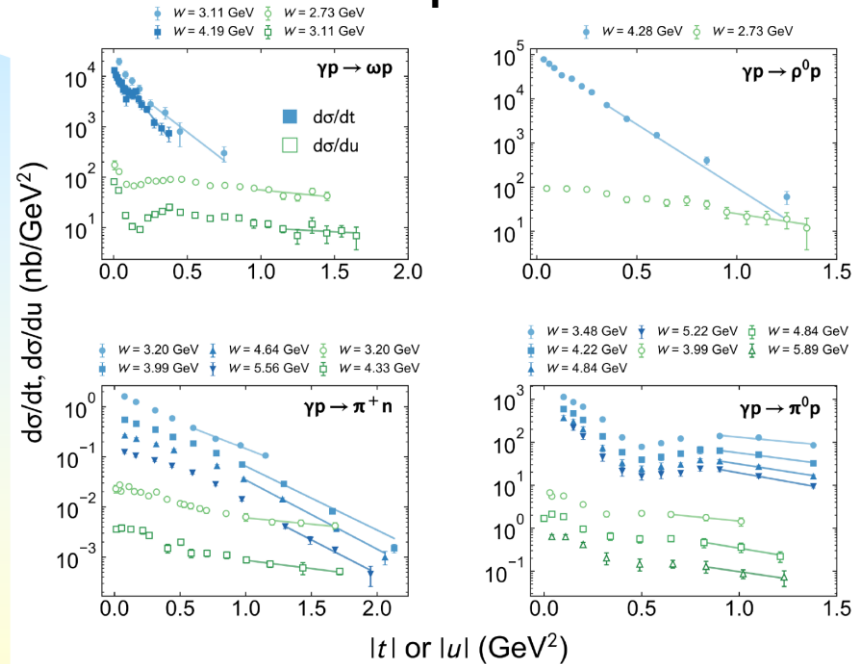


Q^2 vs x_B Phasespace coverage for RGA
(Image credit to the CLAS12 collaboration)

- **CLAS 12 Run group A (RGA)**
 - ~10.6 GeV Electron on unpolarized LH2
 - Completed data taking 2018–2019
 - 300 mC of integrated beam charge were collected, 500T of data size (50% completion)
- **CLAS 12 Run group B (RGB)**
 - ~ 10.6 GeV Electron on unpolarized LD2
 - Completed data taking 2019-2020
 - 136 mC of integrated beam charge were collected (50% completion)
- **More data taking is possible.**
- **Primary goal is to focus on $1 < Q^2 < 2 \text{ GeV}^2$**
 - Better comparison with the prediction

Imaging baryon number within protons

Input data

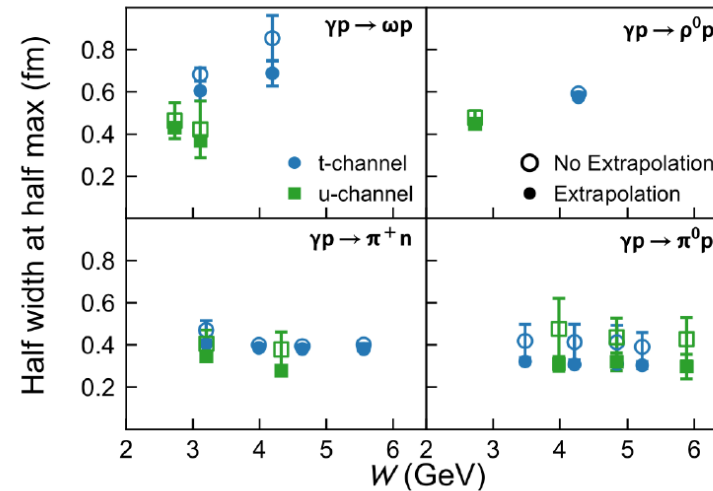


Some data exhibit diffractive minima, some do not
Diffractive minima are associated with sharp edges

Summarizing the data - FWHMs

Spencer Klein

On average, the u-channel radii are a bit smaller than the t-channel equivalents. For the ω , the differences are significant.



- Fourier transform of $f(p_T)$ gives impact parameter distribution.
- Perform transform on $\gamma p \rightarrow \omega p, \gamma p \rightarrow \rho p, \gamma p \rightarrow \pi^0 p, \gamma p \rightarrow \pi^+ n$
- Why are the spatial sizes of protons from the four channels inconsistent?

Hypothesis: Different channels probe different entities

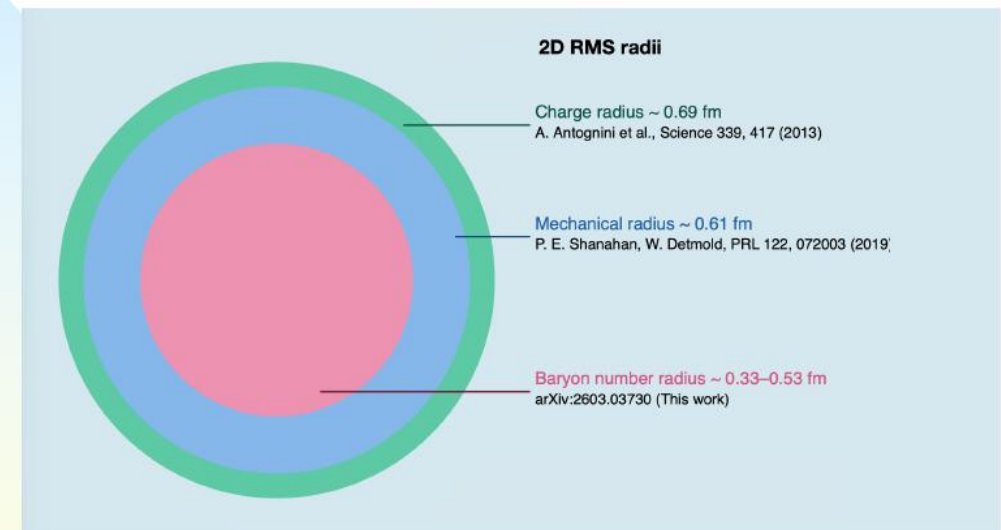
Spencer Klein

How do the sizes compare?

Reaction	Sensitive to	RMS Radius (fm)
Low-energy forward VM photoproduction	High- x quarks and gluons	0.67 – 0.77
Forward π^0 and π^+ photoproduction	High- x quarks	0.33 – 0.45
Backward meson production	Baryon number	0.33 – 0.53
High-energy forward VM photoproduction	Gluons	0.86 – 0.99
Electron scattering	Net charge	0.69 – 0.72
Hydrogen spectroscopy	Charge	0.69

- The baryon number size is in the range 0.33-0.53 fm, depending on reaction.
- This is significantly smaller than the proton size measured via reactions sensitive to gluons, net charge or charge.
- We conclude that the baryon number is somewhat concentrated in the center of the proton

The proton through different microscopes



- Baryon number is concentrated near the center of the proton.
- Consistent with it residing in either high- x quarks or baryon junctions.

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Graphic by Mathias Labonte

Conclusion & Outlook

- Multiple independent lines of evidence now support the baryon junction picture.
- Active experimental programs across RHIC, HERA, and CLAS12 are providing complementary tests.
- New probes of baryon number (e.g. spatial imaging) and theories of transported baryon creation (e.g. GPDs) are being developed.
- The EIC will be indispensable: enable accurate measurement of baryon transport in $e+A$, decisive tests on baryon junction picture.
- **A long-term program with a bright future ahead.**