#### Anson Hook UMD

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# Strong CP Problem

Axions are the simplest and most minimal solution to the Strong CP problem

# Classical Strong CP problem

# Neutron contains an up quark and two down quarks



# Classical Strong CP problem

Electric Dipole moment

 $d_n = qx$  $\frac{1}{3}$ 2 U D 3

# Expected Dipole moment

$$|d_n| \approx ex\sqrt{1 - \cos\theta}$$
  
 $\approx 10^{-14} e\sqrt{1 - \cos\theta} \text{ cm}$ 



## Measured EDM

$$|d_n| \approx ex\sqrt{1 - \cos\theta}$$
  
 $\approx 10^{-14} e\sqrt{1 - \cos\theta} \text{ cm}$ 



 $|d_n| < 2.9 \times 10^{-26} \, e \, \mathrm{cm}$ 

Baker et. al. hep-ex/0602020 : Institut Laue-Langevin, Grenoble

# Classical Strong CP problem

Measurement indicates a small theta



 $\theta < 10^{-12}$ 

Must be a reason!

# Strong CP Problem

Axion solution is the same solution as why carbon dioxide lives on a line

Angle is dynamical and relaxes to minimum

# Strong CP Problem

- Axions are the simplest and most minimal solution to the Strong CP problem
- Solves a problem and can be dark matter
- Axion dark matter obtains its number abundance through the misalignment mechanism
  - Produces cold dark matter regardless how light the axion is

- If it is dark matter, how can we look for it?
- The axion is a classical field due to large number abundance
  - If mass is less than eV, then many particles per Compton wavelength

$$\phi(x,t) = \sum \frac{1}{l^{3/2}} \frac{1}{\sqrt{2\omega_n}} (a_n e^{ip \cdot x} + a_n^{\dagger} e^{-ip \cdot x})$$

- First treat dark matter as a particle in a box
- Take velocity profile of dark matter from simulation and find the quantum state that reproduces it
  - Isothermal Profile

$$\rho = \int d^3 v f(v) \omega_v \overline{n}$$

Just like how a laser has Poisson statistics for number of photons, Axion state also has Poisson statistics for number of axions

$$N_n = \left(\frac{2\pi}{m}\right)^3 f(v_n)\overline{n}$$

$$\mid N \rangle = \alpha e^{\sqrt{N_n} a_n^{\dagger}} \mid 0 \rangle$$

$$\phi(t) = \frac{\sqrt{\rho}}{m} \sum_{i_r} \alpha_r \sqrt{f_r 4\pi v^2 \Delta v} \cos\left(\omega_r t + \phi_r\right)$$

# Crank through and calculate a(t) in this background

Sum of many sines with random phases

Basically the square root of an integral



# Qualitative Features

 Axion acts like a cosine function over distance scales

 $L \sim \frac{1}{m_a v}$ 

Afterwards amplitude and phase randomly scrambled

• Coherence length of order de Broglie wavelength

# Qualitative Features

 Axion acts like a cosine function over time scales



Afterwards amplitude and phase randomly scrambled

Coherence time is large

### Qualitative Features

- Frequency of the axion sine wave  $\omega \sim m_a \pm 10^{-6} m_a$
- Thus the quality factor of axion dark matter is very large

$$Q = \frac{1}{\omega\tau} \sim 10^6$$

### Looking for the axion

 $\mathcal{L} \supset \frac{a}{4f} F\tilde{F}$ 

- Looking for the axion through the coupling to gluons is HARD
  - Very few experiments can reach the QCD axion line
- Instead look for the axion through its coupling with the photon

# DISCLAIMER

- QCD Axion
  - Solves the Strong CP problem
  - Couples to photons and gluons and fermion spin
- ALP (Axion like particles)
  - Does NOT solve the Strong CP problem
  - Couples to photons and/or fermion spin
- Axions
  - Can be either
  - Figure it out from context

#### Effect of photon coupling

$$\mathcal{L} \supset \frac{a}{4f} F\tilde{F}$$

 $v_{\rm phase} \approx 1 \pm \frac{\dot{a}}{2kf}$ 

For circularly polarized light

$$-\omega^2 + k^2 \mp \frac{da}{dt}\frac{k}{f} = 0$$

### Effect of photon coupling

$$v_{\rm phase} \approx 1 \pm \frac{a}{2kf}$$

- Phase velocity of circularly polarized light is different depending on which polarization it is
- Device most sensitive to differences in phase velocities is an interferometer

# Axion interferometry

- One-to-one mapping between axion interferometry and gravity wave interferometry
- An axion interferometer can double as a gravity wave detector
- Axion dark matter appears in the same manner as a continuous gravity wave signal with a quality factor of 10<sup>6</sup>

## Gravity wave interferometry

#### Mirror

Consider a plus polarized gravity wave incident perpendicular to the interferometer

Mirror



# Axion interferometry







- Only difference is the presence of wave plates
- Needed to maintain polarization

# Axion wave



1011111



Mirror

#### No Axion DM

# Exactly the same as a gravity wave interferometer

# Experiment doubles as a gravity wave detector

No need to send the legs in different directions otherwise



# Resonant interferometry

#### **Resonant Detector instead!**



# Resonant interferometry

1. Optimal Length is as expected  $L = \lambda_g/2$ 

2. Resonant detector







Mirror

**Detector** 

Same Mapping as before Otherwise identical to Gravity wave detector

#### Parameters

- What are reasonable parameters?
  - Similar to gravity wave interferometers : Maybe do as part of setting up and testing a gravity wave interferometer
  - Cost is all in man power
- Assumption : Shot noise and radiation noise limit until 10 Hz where seismic noise becomes an issue
  - 40 m arm length
  - 10 kg mirror
  - 30 days run time : factor of 3 worse limits if you run for 6 hours



Seismic Noise becomes an issue

 Large Finesse/power not needed to probe new regions of parameter space!

- If detector is dedicated to an axion search and not gravity wave search, can do better!
- Radiation pressure can be mitigated if same mirror is used for both arms!

#### Radiation Pressure replaced by Radiation Torque





10 kg mirror 10 cm diameter 1 cm between beams Red: 1 MW power Black: 1 kW power Dotted :  $F = 10^6$ Solid :  $F = 10^{2}$ 

# Conclusion

- Axion dark matter changes the phase velocity of circularly polarized light
- Can look for this effect in an interferometer
- Can extend bounds by up to 2-3 orders of magnitude over some range of parameters
- Do not need the newest fanciest technology
  - Need to make sure that birefringent backgrounds are under control!