

Lightish but clumpy: scalar dark matter from inflationary fluctuations

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Based on: Gonzalo Alonso-Álvarez, J. Jaeckel, **arXiv:1807.09785**
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Goals & Motivation

Goal: Link high energy phenomena with low energy observables.

High Energy: Inflation \longleftrightarrow Low Energy: Dark matter

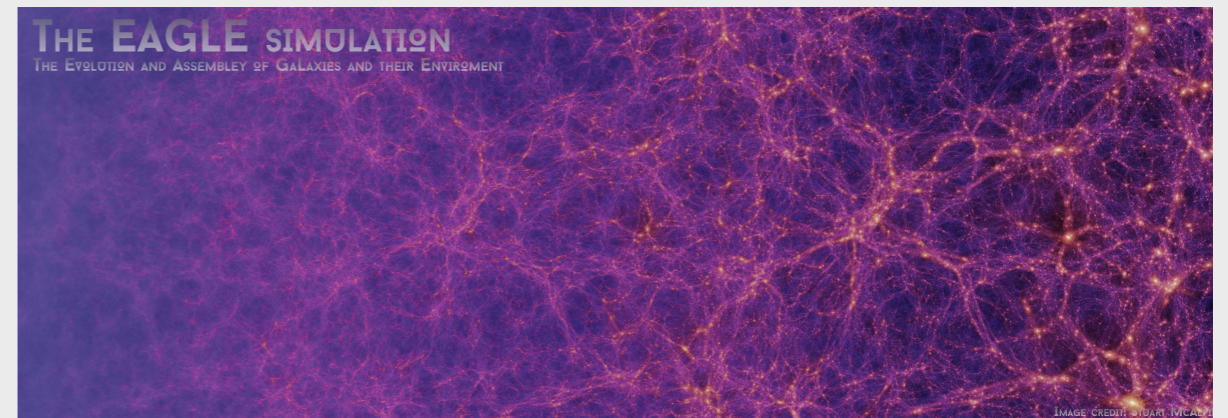
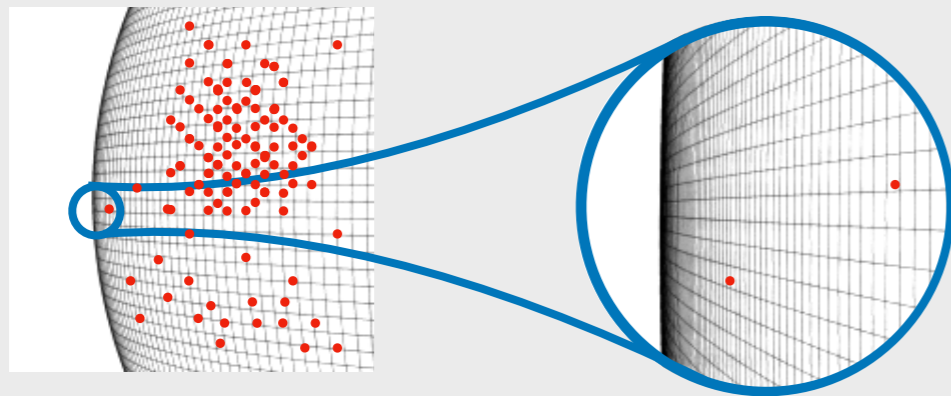


Image: <http://icc.dur.ac.uk/Eagle>

How: New mechanisms of non-thermal production of dark matter.

Main idea

Initial premise:

Dark matter can be generated **gravitationally** from quantum fluctuations of a scalar field during inflation.

Caveat:

Isocurvature perturbations are not observed.

Key observation:

A **non-minimal coupling to gravity** suppresses isocurvature and makes the scenario phenomenologically viable.

Different ideas in this direction

- Misalignment mechanism for scalars and vectors
Arias et al [1201.5902], *Nelson & Scholtz* [1105.2812], ...
- Resonant production at (p)reheating
 - WIMPZILLAs: *Chung et al* [hep-ph/9802238], ...
 - Non-minimally coupled scalars: *Fairbairn et al* [1808.08236], ...
- Vector dark matter from inflationary fluctuations
Graham et al [1504.02102]
- Thermal production via graviton exchange
Garny et al [1511.03278], *Ema et al* [1804.07471], ...

Relevance & applications

Most immediate and minimal scenario:

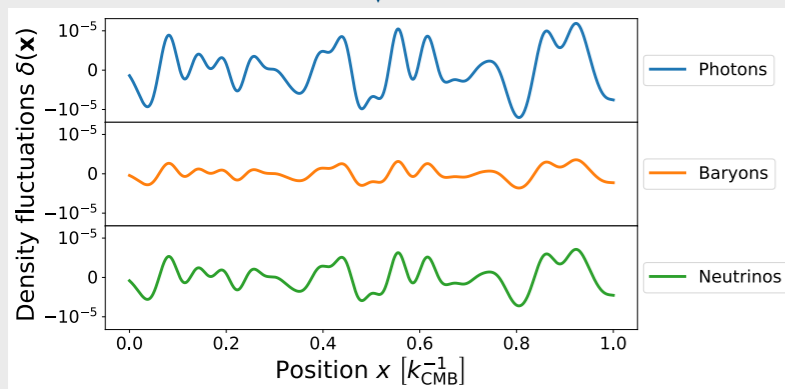
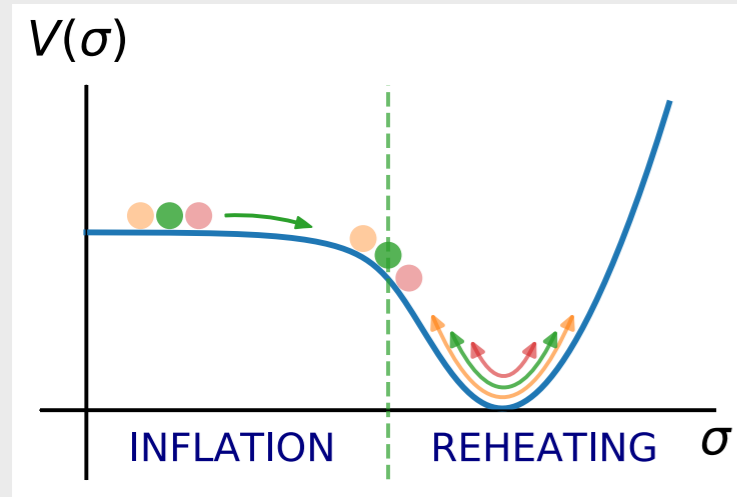
Gravitational generation of all the dark matter that we observe.

Other potentially interesting applications:

- Freeze-in dark matter: [Hall et al \[0911.1120\]](#) ([Bernal et al \[1706.07442\]](#))
- Affleck-Dine baryogenesis: [Bettoni & Rubio \[1805.02669\]](#)
- Asymmetric dark matter: [Kaplan et al \[0901.4117\]](#) ([Zurek \[1308.0338\]](#))
- Higgs portal DM: [Cosme et al \[1802.09434\]](#)

(Dark) matter & inflation

Ordinary matter



Adiabatic perturbations

seen in the CMB

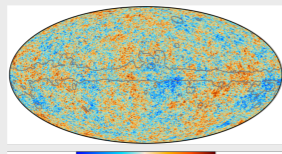
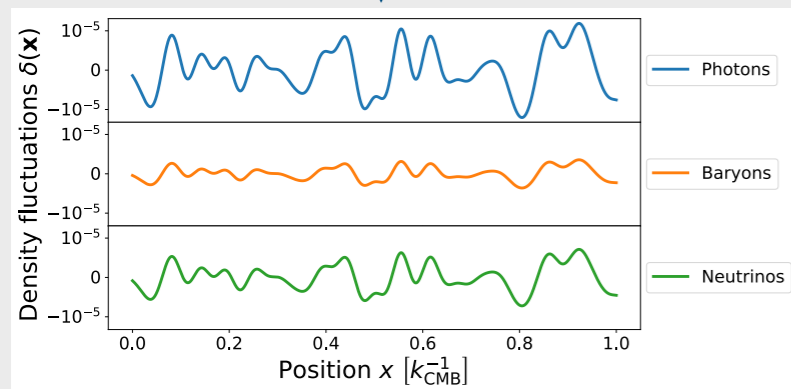
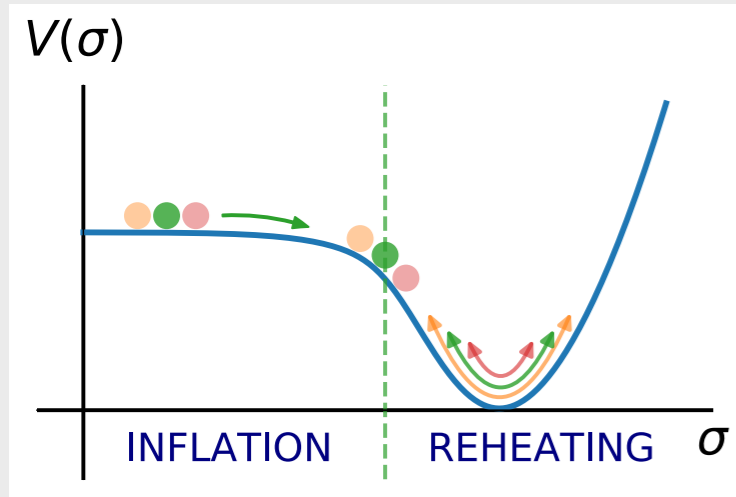


Image: <http://sci.esa.int/planck>

(Dark) matter & inflation

Dark matter?



Adiabatic perturbations
seen in the CMB

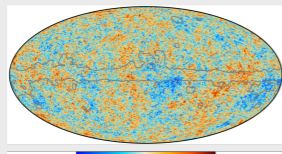
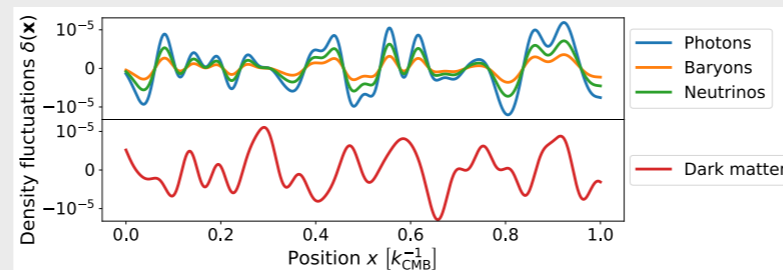
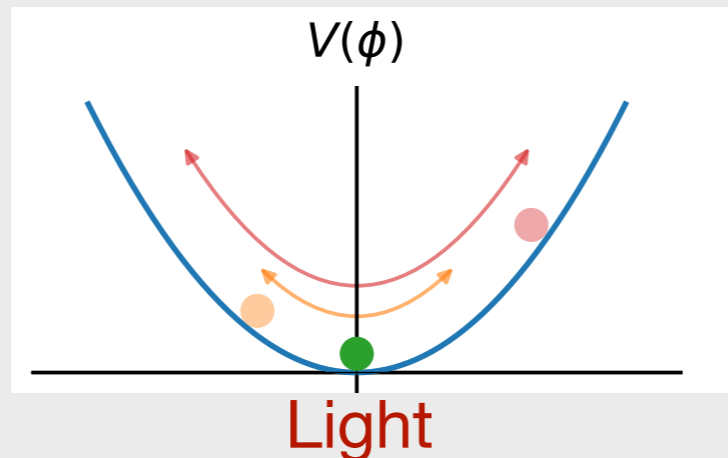


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Isocurvature perturbations

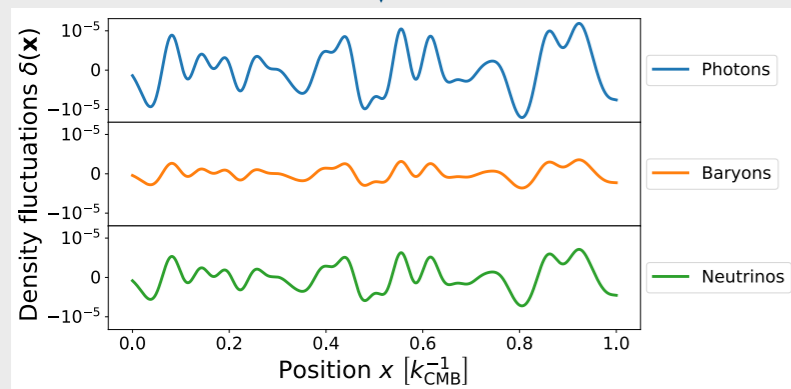
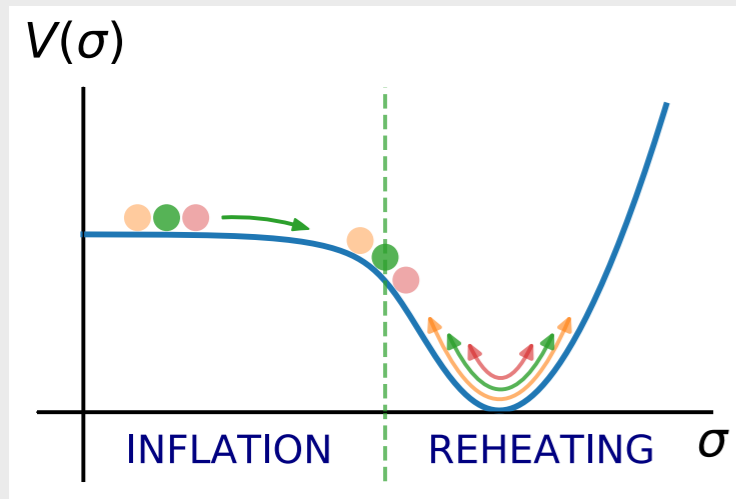
Planck 2018 [1807.06211]



$$\Rightarrow \frac{\delta\rho_{\text{iso}}(k_*)}{\delta\rho_{\text{ad}}(k_*)} \lesssim 0.03$$

at scales $k_{\text{CMB}}^{-1} \simeq 20 \text{ Mpc}$

(Dark) matter & inflation



Adiabatic perturbations
seen in the CMB

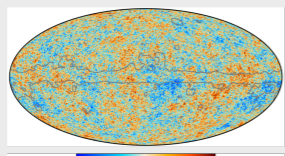
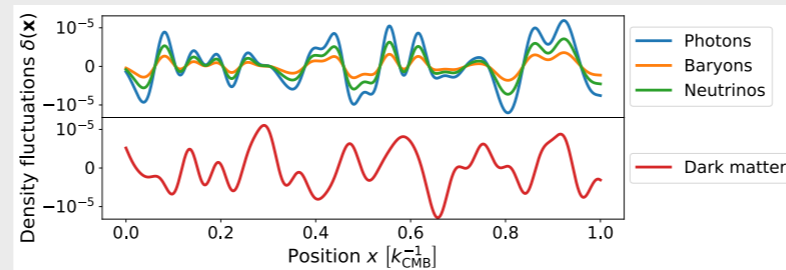
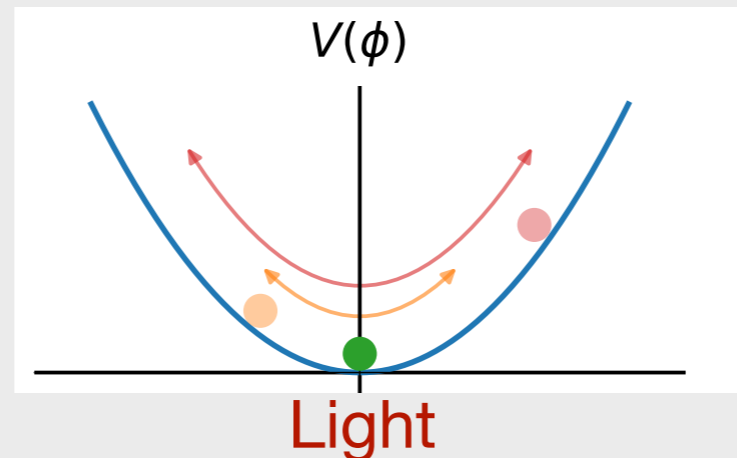


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Isocurvature perturbations

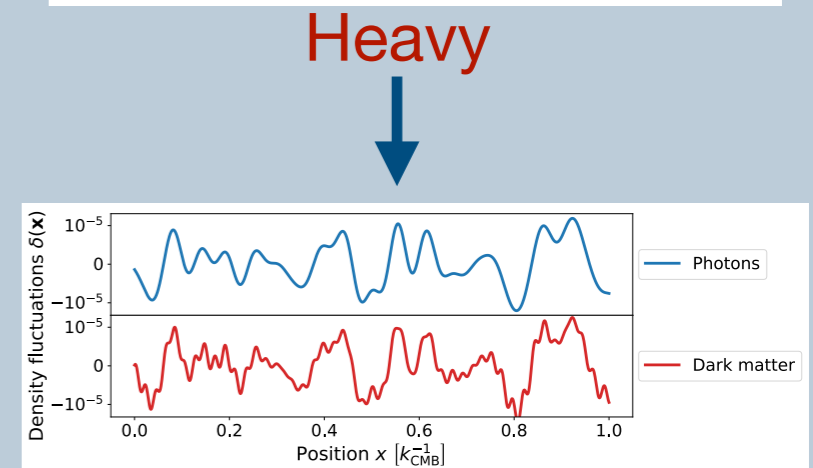
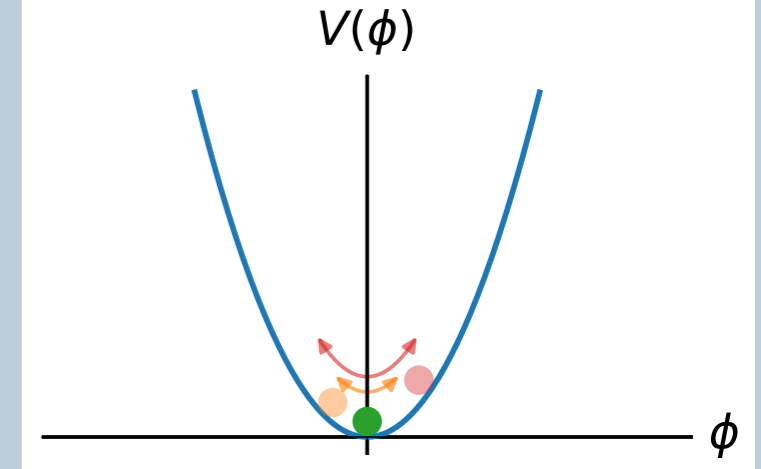
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$$\Rightarrow \frac{\delta\rho_{\text{iso}}(k_*)}{\delta\rho_{\text{ad}}(k_*)} \lesssim 0.03$$

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Dark matter!



Isocurvature perturbations
suppressed at large scales.

But only if $m \sim \mathcal{O}(H_I)$!

How may this happen?

Non-minimal coupling to gravity

Action in the Jordan frame:

$$S = \int d^4x \sqrt{-g} \left(\left(\tilde{M}_p^2 - \xi \phi^2 \right) R - \frac{1}{2} g^{\mu\nu} \nabla_\mu \phi \nabla_\nu \phi - m^2 \phi^2 \right)$$

Effective mass for ϕ :

$$m_{\text{eff}}^2 = m^2 + \xi R$$

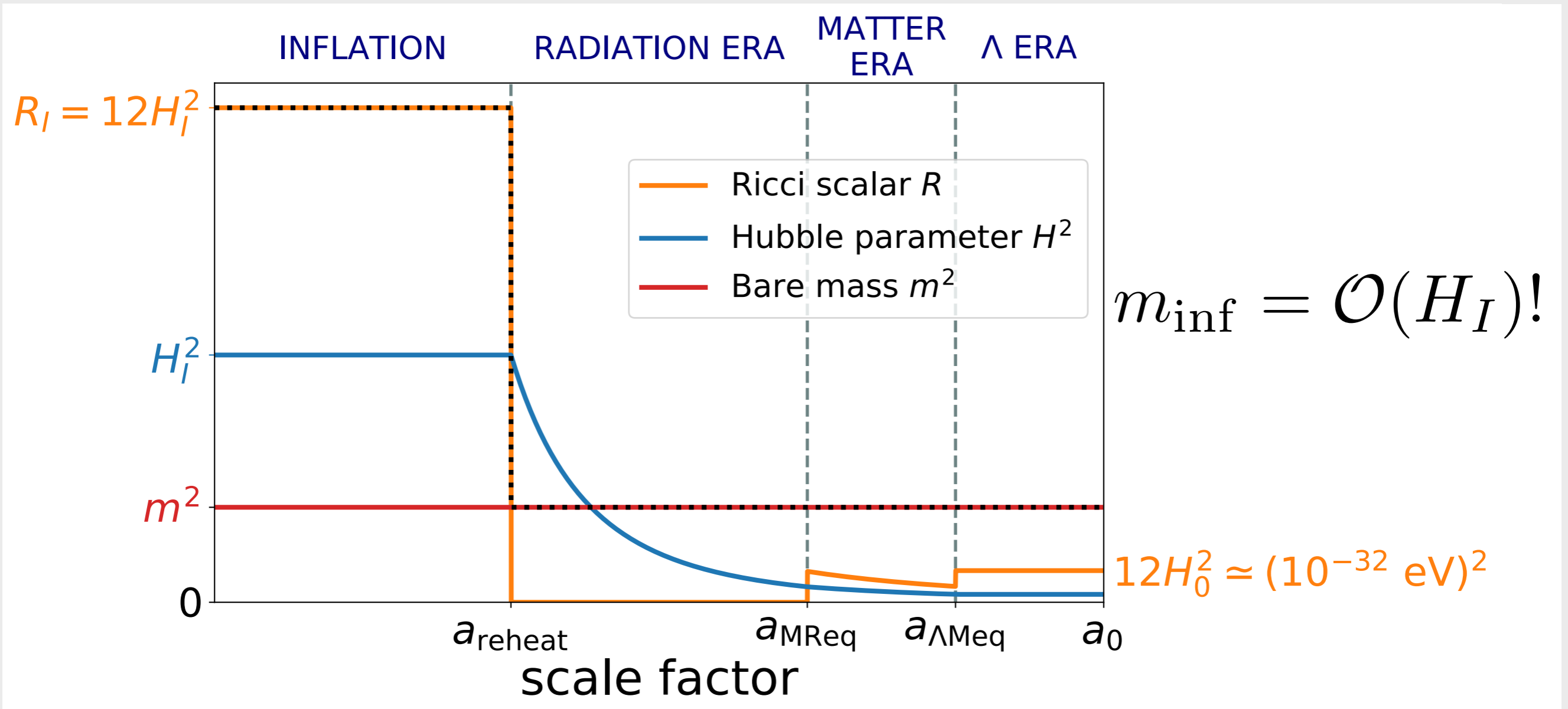
Bare mass
(constant)

“Gravitational” mass
(space-time dependent)

Negligible backreaction on the geometry.

Non-minimal coupling in cosmology

Effective mass for ϕ : $m_{\text{eff}}^2 = m^2 + \xi R$, $R \propto H^2$



The field is heavy during inflation, but it can be light from then on.

Cosmological evolution

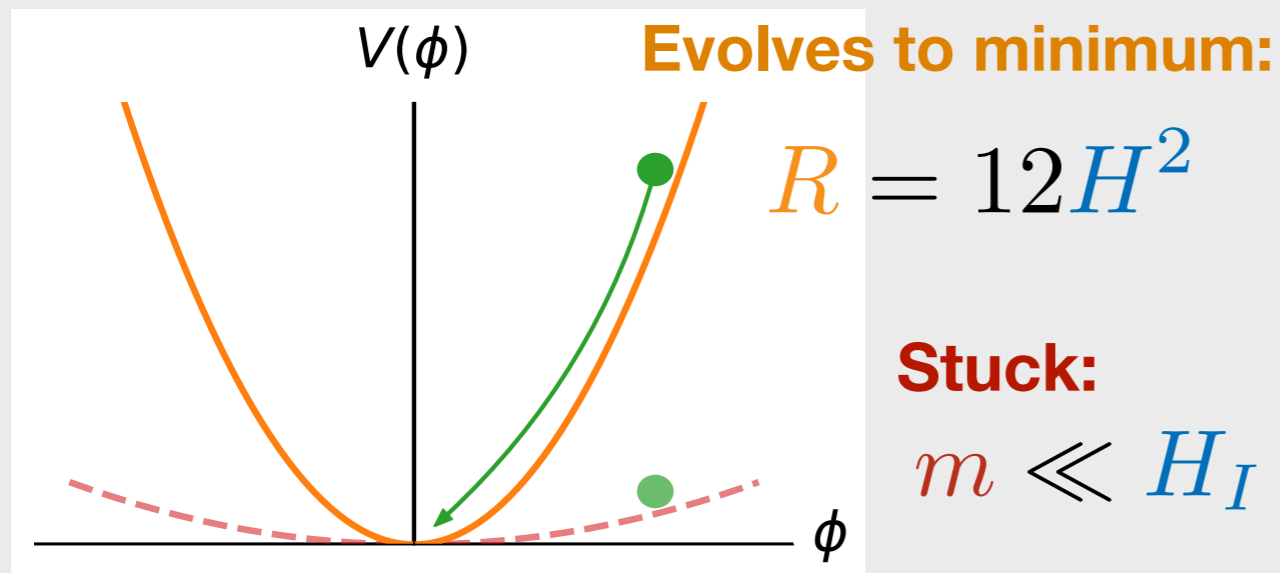
- I. Quantum evolution during inflation
 - (i) Homogeneous field
 - (ii) Non-zero momentum modes (fluctuations)
- II. Classical evolution after inflation
- III. Late times power spectrum and relic abundance

I. Evolution during inflation

Classically

The homogeneous field is damped away during inflation.

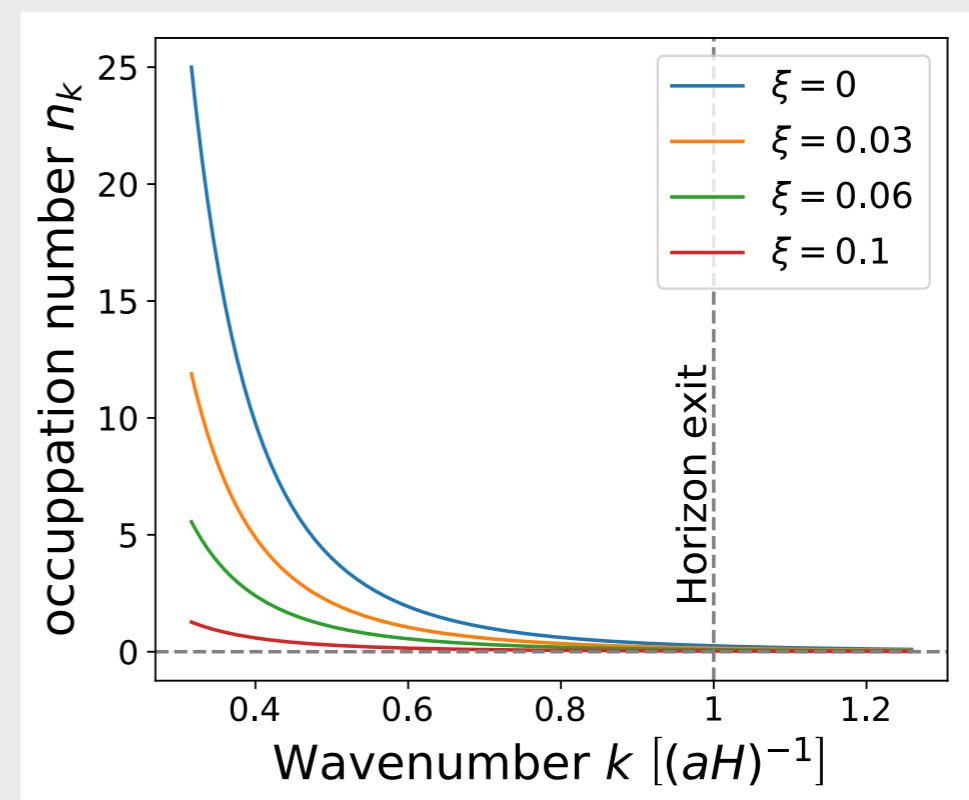
$$\ddot{\phi} + 3H\dot{\phi} + (m^2 + \xi R)\phi = 0.$$



$$\phi_E \simeq \phi_0 e^{-\frac{4}{3}\xi N}$$

Quantum mechanically

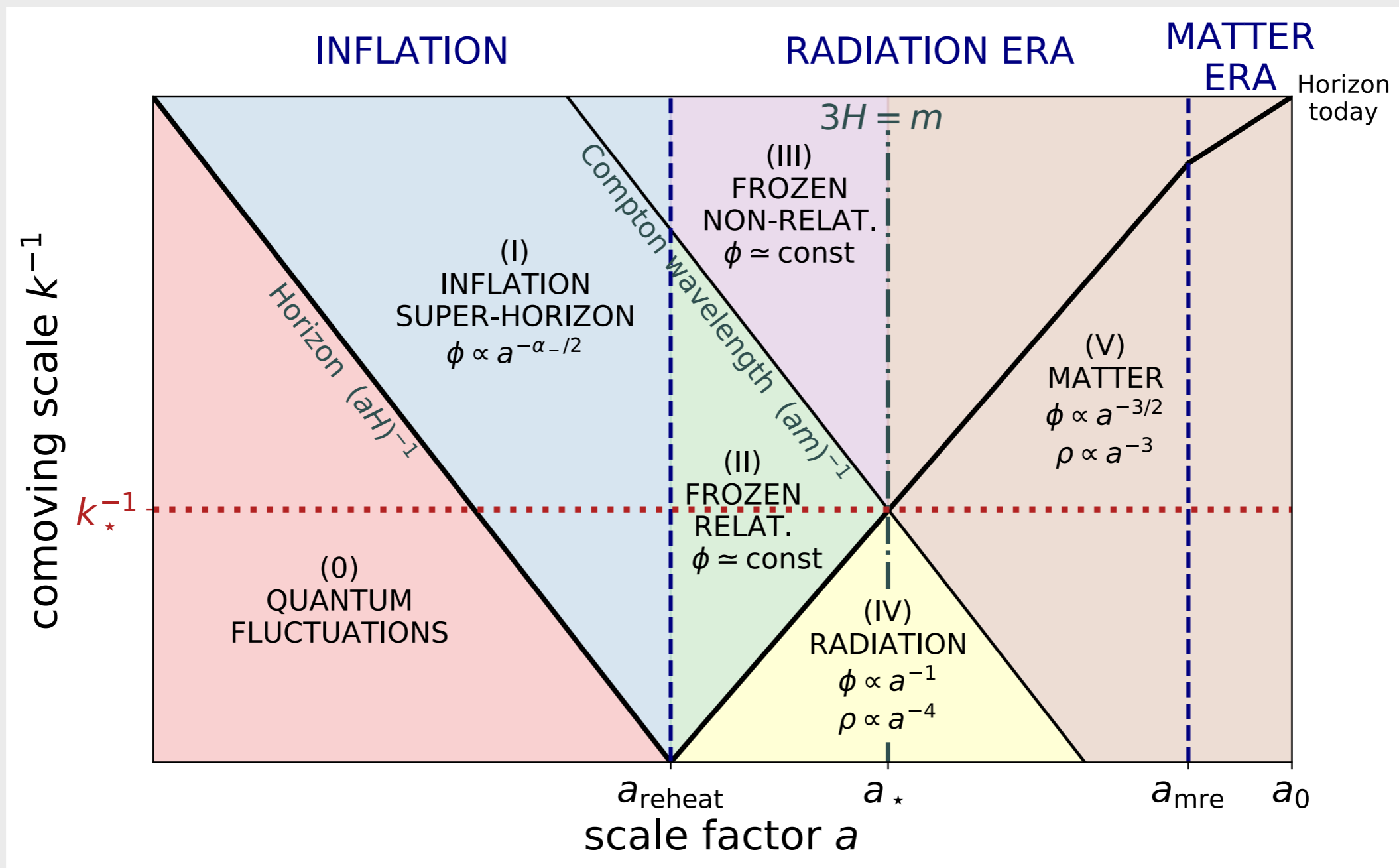
Higher momentum modes are excited due to the time-dependent gravitational background.



A larger ξ suppresses the occupation number.

II. Classical evolution after inflation

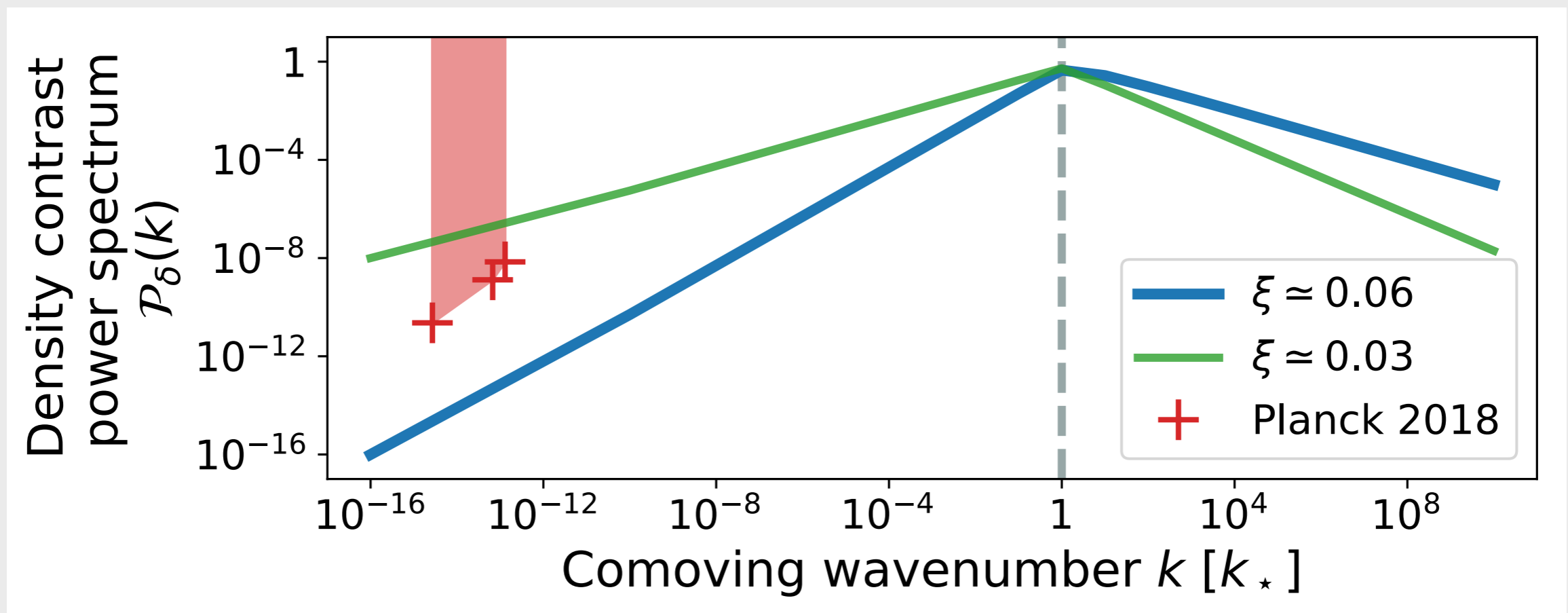
$$\ddot{\phi}_k + 3H\dot{\phi}_k + \left(\frac{k^2}{a^2(t)} + m^2 + \xi R \right) \phi_k = 0.$$



III. Isocurvature power spectrum

The density power spectrum is peaked at the comoving scale k_{\star}^{-1} .

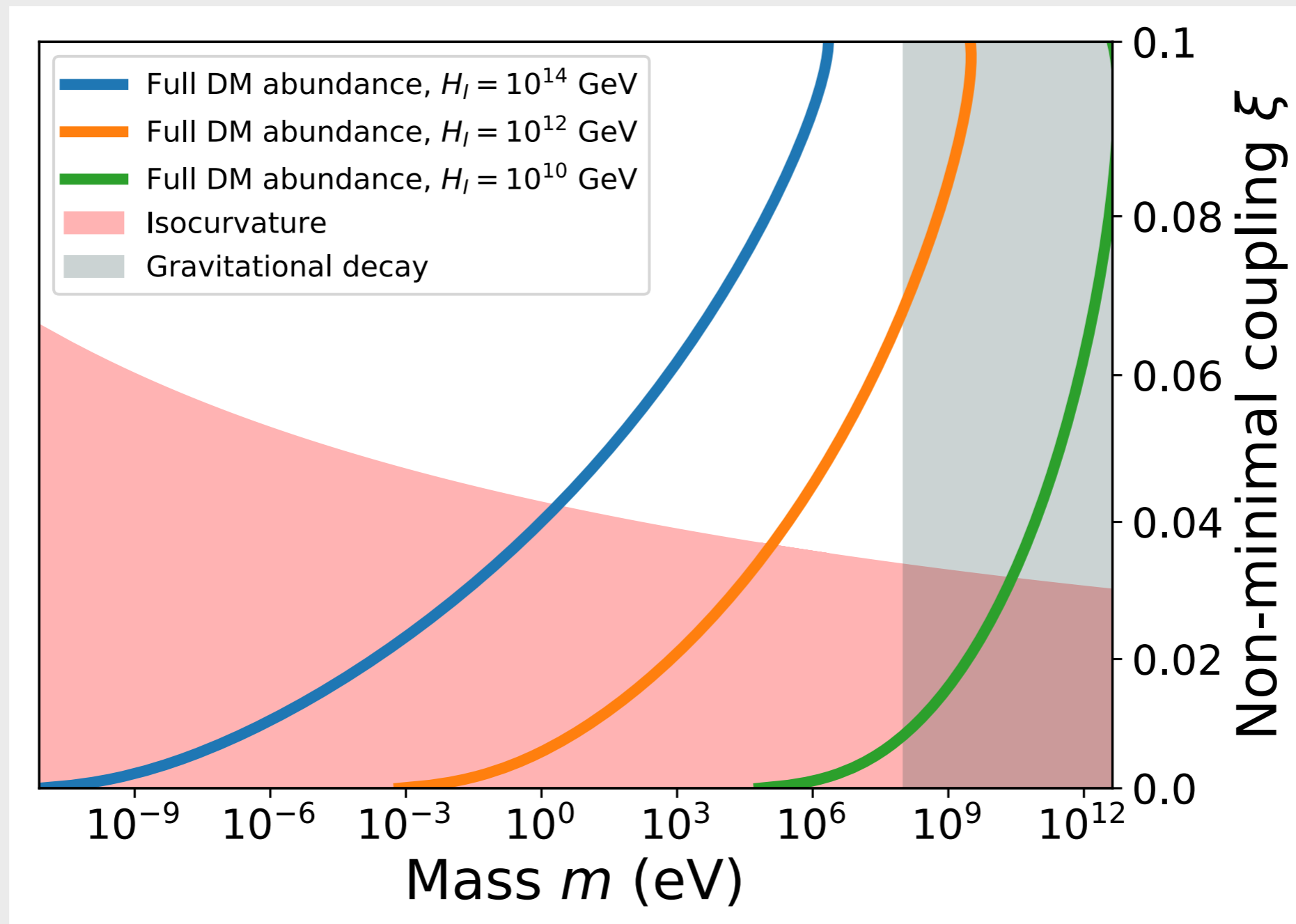
$$k_{\star}^{-1} \simeq 4 \cdot 10^7 \text{ km} \sqrt{\frac{\text{eV}}{m}} \\ (1 \mu\text{pc})$$



Isocurvature fluctuations are small enough at the CMB scales.

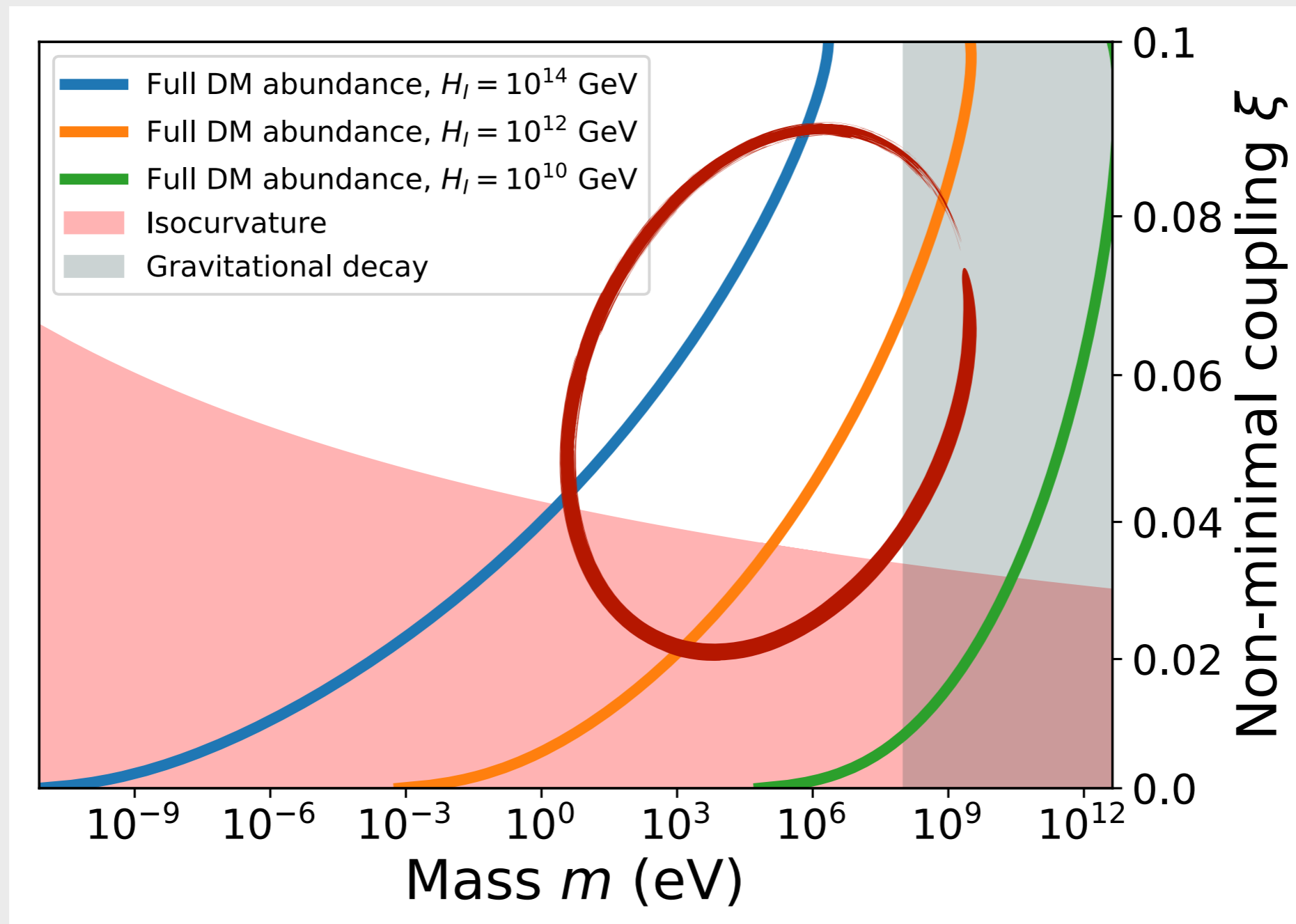
Relic density

For a given m and ξ , fixing the abundance to the dark matter one selects a scale of inflation.



Relic density

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Observables: cosmology

Two main particular features of the mechanism:

1. Presence of a small component of isocurvature perturbations.
 - CDI “axion II” scenario in Planck 2018 analysis.
2. Very strong clumping at small scales.
 - Fluctuations at the peak are large and they collapse around matter-radiation equality.

$$\ell_{\text{today}} \simeq \frac{1}{z_{\text{eq}}} k_{\star}^{-1} \simeq 10^4 \text{ km} \sqrt{\frac{\text{eV}}{m}}$$

- Potential signals in gravitational lensing, astrophysical processes and direct & indirect detection.

Observables: particle physics

So far, we haven't considered any non-gravitational interactions with other fields.

The field has a \mathbb{Z}_2 symmetry that impedes the decay.

$$\phi \longleftrightarrow -\phi$$

We can assume that at least gravity breaks this global symmetry.

\Rightarrow Gravitational-mediated decay, eg $\frac{\phi}{M_p} \mathcal{F} \mathcal{F}$

$$\text{Stability: } \tau = \Gamma^{-1} \sim \frac{64\pi M_p^2}{m^3} \lesssim \tau_{\text{univ}} \sim 10^{-33} \text{ eV}$$

Other couplings are compatible with the production mechanism.

Conclusions

1. Dark matter can be generated from quantum fluctuations of a light scalar field during inflation.
 2. A small non-minimal coupling to gravity suppresses isocurvature perturbations at CMB scales.
- In the future: Study the phenomenology of the scenario
 - Cosmology: clumping at small scales.
 - Particle physics: coupling to other fields? Detection & stability.

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Thanks!



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