Statistics of energy-averaged cross sections: "<u>unknown unknowns</u>" or neglected "<u>unknown knowns</u>"?

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Theoretical estimate of cross section fluctuations

• Measure of importance of fluctuations:

$$R_{ab} = \frac{\sqrt{\text{cov}(\sigma_{ab}, \sigma_{ab})}}{\langle \sigma_{ab} \rangle} = \frac{\Delta(\sigma_{ab})}{\langle \sigma_{ab} \rangle}$$

where

$$\mathbf{cov}(\sigma_{ab}, \sigma_{cd}) = \langle \sigma_{ab} \sigma_{cd} \rangle - \langle \sigma_{ab} \rangle \langle \sigma_{cd} \rangle$$

Brown-Kawano result:

Average resonance separation

$$\frac{R_{ab}^2}{\Gamma_W} = \frac{\pi a}{\Gamma_W} \frac{W_{abab}}{W_{ab}^2} - 1$$

- \mathcal{W}_{ab} : standard WFC
- \mathcal{W}_{abab} : quartic generalization of WFC
- Γ_W : Weisskopf correlation width

$$\Gamma_W = \frac{d}{2\pi} \sum_c T_c$$
Transmission
coefficients



Fluctuations substantial!

11/8/17

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Variance of energy-averaged cross section: an identity

• Model energy-averaged cross section as

$$\langle \sigma_{ab}(E) \rangle_{\Delta} = \frac{1}{\Delta} \int_{E-\frac{\Delta}{2}}^{E+\frac{\Delta}{2}} \sigma_{ab}(E')dE'$$

• Variance about theoretical (GOE) average $\bar{\sigma}_{ab}$:

Gaussian orthogonal ensemble of random Hamiltonian matrices

$$v_{ab}(\Delta) = \overline{(\langle \sigma_{ab}(E) \rangle_{\Delta} - \overline{\sigma}_{ab})^2} = \frac{2}{\Delta} \int_{0}^{\Delta} (\Delta - \varepsilon) C_{ab}(\varepsilon) d\varepsilon$$

where the auto-correlation function

$$C_{ab}(2\varepsilon) = \overline{\sigma_{ab}(E-\varepsilon)\sigma_{ab}(E+\varepsilon)} - (\overline{\sigma}_{ab})^2$$

What's known about autocorrelation functions?

• Numerical studies able to relate $C_{ab}(2\varepsilon)$ to $F_{ab}(2\varepsilon)$

$$=\overline{(S_{ab}(E-\varepsilon-\delta_{ab})^2(S_{ab}^*(E+\varepsilon)-\delta_{ab})^2}$$

Exact reduction possible of GOE average in F_{ab} to
 3-dimensional integral (Phys Lett B <u>211</u>, 379)

• Simple relations inferred – e.g., for $a \neq b$, $C_{ab}(\varepsilon) \approx \frac{R_{ab}^2}{R_{ab}^2 + 1} |F_{ab}(\varepsilon)|$ (*)

•
$$R_{ab}$$
 calculated with $F_{ab}(0)$ and $\overline{\sigma}_{ab}$





Solid line: $C_{ab}(\varepsilon)$ [GOE-MC calc's] Dashed line: rhs of (\bigstar) [analytical] (32 equal T_c 's) 5

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 $F_{ab}(\varepsilon) \quad (a \neq b)$

Similar to result for $\bar{\sigma}_{ab}$ in Phys Rep <u>129</u>, 367

$$=\frac{1}{8}\int_{0}^{\infty}d\lambda_{1}\int_{0}^{\infty}d\lambda_{2}\int_{0}^{1}d\lambda e^{-i\varphi(\lambda_{1},\lambda_{2},\lambda)}\mu(\lambda_{1},\lambda_{2},\lambda)\Pi(\lambda_{1},\lambda_{2},\lambda)f_{a}(\lambda_{1},\lambda_{2},\lambda)f_{b}(\lambda_{1},\lambda_{2},\lambda)$$

•
$$\mu(\lambda_1, \lambda_2, \lambda) = \frac{\lambda(1-\lambda)|\lambda_1-\lambda_2|}{(\lambda+\lambda_1)^2(\lambda+\lambda_2)^2\sqrt{\lambda_1(1+\lambda_1)\lambda_2(1+\lambda_2)}}$$

Not the most physically transparent result

(Useful comments on numerical evaluation in Phys Lett B <u>685</u>, 263)

$$= f_c(\lambda_1, \lambda_2, \lambda) = \frac{\lambda_1(1+\lambda_1)}{(1+T_c\lambda_1)^2} + \frac{\lambda_2(1+\lambda_2)}{(1+T_c\lambda_2)^2} + \frac{2\lambda(1-\lambda)}{(1-T_c\lambda)^2} + \frac{1}{2}(1-T_c)\left(\frac{\lambda_1}{(1+T_c\lambda_1)} + \frac{\lambda_2}{(1+T_c\lambda_2)} + \frac{2\lambda}{(1-T_c\lambda_2)}\right)^2$$

•
$$\boldsymbol{\varphi}(\lambda_1, \lambda_2, \lambda) = \pi \frac{\varepsilon}{d} (\lambda_1 + \lambda_2 + 2\lambda)$$

• $\Pi(\lambda_1, \lambda_2, \lambda) = \prod_c \frac{1 - T_c \lambda}{\sqrt{(1 + T_c \lambda_1)(1 + T_c \lambda_2)}}$

Alternative calculation of autocorrelation functions

Adopt statistical Breit-Wigner (SBW) model for *S*-matrix

S-matrix = sum of Breit-Wigners with random parameters (distributed as in GOE)



A simple result for autocorrelation functions (lumping approximation) Nucl Sci Eng <u>62</u>, 756



<u>Implication</u>: ergodicity of σ_{ab} [: $C_{ab}(\varepsilon) \rightarrow 0 \text{ as } \varepsilon \rightarrow \infty$]

For large enough Δ ,

$$\langle \sigma_{ab}(E) \rangle_{\Delta} = \bar{\sigma}_{ab}$$



Closing thoughts on "unknown knowns"

- In unresolved resonance regime, unless averaging interval $\Delta > 10^4 \Gamma_W$, there is a sizeable probability (~30%) that *energy-averaged* cross section inferred from experiment and *theoretical ensemble* average should differ by > 1%
- Distinction between ensemble and running averages impacts on uncertainties in the self-shielding correction factors used in the extraction of cross sections
 - Concern about MC estimates of self-shielding corrections is not new: in "[c]omparing calculation of measurement, we must keep in mind that the Monte Carlo distribution is an expectation value, i.e., an average over many sampled ladders, whereas the observed distribution comes from the one ladder realized in nature". (Fröhner & Larson, 1995)
 - Current view (EJP A <u>52</u>, 170): "the uncertainty on [the self-shielding and multiple interaction] correction factor is ≤ 0.5%, as demonstrated in ref. [EJP A <u>50</u>, 124]". This error estimate is based on the agreement of ratios of expectation values calculated with SESH and MCNP.

Thank you for your attention!