## A DARK FORCE FOR BARYONS

Ian Shoemaker<br>BF2OII<br>October 19th, 201 |

with Michael Graesser and Luca Vecchi

## OR:

# WHY YOU SHOULD LEARN TO STOP WORRYING AND GAUGE BARYON NUMBER 

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Dark matter à la Occam

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THE STANDARD MODEL

|  |  | ermions |  | Bosons |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\begin{aligned} & U \\ & \text { up } \end{aligned}$ |  <br> charm | $\underset{\text { top }}{t}$ | $\mathcal{P}$ |  |
|  | down | strange |  |  |  |
|  |  |  |  |  | \% |
|  |  |  | $7$ <br> tau | $\underset{\text { gluon }}{9}$ |  |
|  |  |  |  |  |  |
| *Yet to be confirmed |  |  |  |  |  |

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lex parsimoniae

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THE STANDARD MODEL


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## What do we really know about DM?

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2. It's stable (or at least very long-lived).

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3. An indication of an underlying origin.

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The only problem is...

## Super Kamiokande says:



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## The proton is stable.



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- Baryogenesis requires a DM asymmetry.
- Shared gauge interactions with baryons predict novel signatures: monojets and low mass DD.


## Gauging baryon number

- Older examples:
- Carone and Murayama 1994; Bailey and Davidson 1995; Aranda and Carone 1998.
- More recently:
- Dulaney, Fileviez-Perez and Wise (2010); Buckley, Fileviez-Perez, Hooper, and Neil (201I).


## An anomaly-free example



- New chiral states

|  | $S U(3)_{C}$ | $S U(2)_{W}$ | $U(1)_{Y}$ | $U(1)_{B}$ |
| :---: | :---: | :---: | :---: | :---: |
| $Q_{i}^{\prime}$ | 3 | 2 | $+\frac{1}{6}$ | $-\frac{1}{N}$ |
| $u_{c i}^{\prime}$ | $\overline{3}$ | 1 | $-\frac{2}{3}$ | $+\frac{1}{N}$ |
| $d_{c i}^{\prime}$ | $\overline{3}$ | 1 | $+\frac{1}{3}$ | $+\frac{1}{N}$ |
| $L_{i}^{\prime}$ | 1 | 2 | $-\frac{1}{2}$ | 0 |
| $\nu_{c i}^{\prime}$ | 1 | 1 | 0 | 0 |
| $e_{c i}^{\prime}$ | 1 | 1 | +1 | 0 |

$N$ dark generations

- Spontaneously break $U(1)_{B}$

| $S^{+}$ | 1 | 1 | 0 | $+B(S)$ |
| :--- | :--- | :--- | :--- | :--- |
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- Unlike conventional ADM, the asymmetries are generated simultaneously.
- Recent work by: Bell, Petraki, IMS, Volkas [I I 05.3730].


## DIRECT DETECTION BOUNDS



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annihilation physics

$$
\begin{gathered}
\mathfrak{\imath} \\
\text { DM-quark } \\
\text { scattering }
\end{gathered}
$$

## RECOIL SPECTRUM

$$
\frac{d R}{d E_{R}}=\frac{N_{T} \rho_{\odot}}{m_{X}} \int_{|\vec{v}|>v_{m i n}} d^{3} v v f\left(\vec{v}, \vec{v}_{\oplus}\right) \frac{d \sigma}{d E_{R}}
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particle physics

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f(v) \propto\left[\exp \left(\frac{v_{e s c}^{2}-v^{2}}{k v_{0}^{2}}\right)-1\right]^{k}
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[Lisanti, Strigari, Wacker, Wechsler (20 10 )]

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High-velocity tail is important for light DM.

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VECTOR CASE:

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\frac{d \sigma}{d E_{R}}=\frac{m_{N} A^{2}}{2 \pi v^{2}}\left(\frac{q_{V} g_{B}^{2}}{m_{B}^{2}}\right)^{2} F^{2}\left(E_{R}\right)
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\text { DD imposes: } \\
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## Monojets at the Tevatron

- For light DM, the Tevatron and the LHC are the world's best DD experiments [Goodman, et al. (20 I 0); Bai, Fox, Harnik (20 I 0)].


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$$
p \bar{p} \rightarrow \not \oiint_{T}+j
$$

See Luca's talk.

## Combined constraints:

## axial case

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D^{\mu} X=\partial^{\mu} X+i g_{B}\left(q_{V}^{0}+q_{A} \gamma^{5}\right) Z_{B}^{\mu} X
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## cONCLUSIONS

- Gauging baryon number saves the proton + automatic DM candidate charged under baryonic force.
- Simultaneous generation of dark and visible asymmetries.
- Consistent with bounds from B-factories, LEP, mono-jet Tevatron searches, and direct detection for:
- GeV-scale DM with a GeV-scale mediator.
- LHC and direct detection will probe much of the remaining parameter space.


## EXTRAS

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Decay operator $\leftrightarrow$ asymmetry transfer operator

