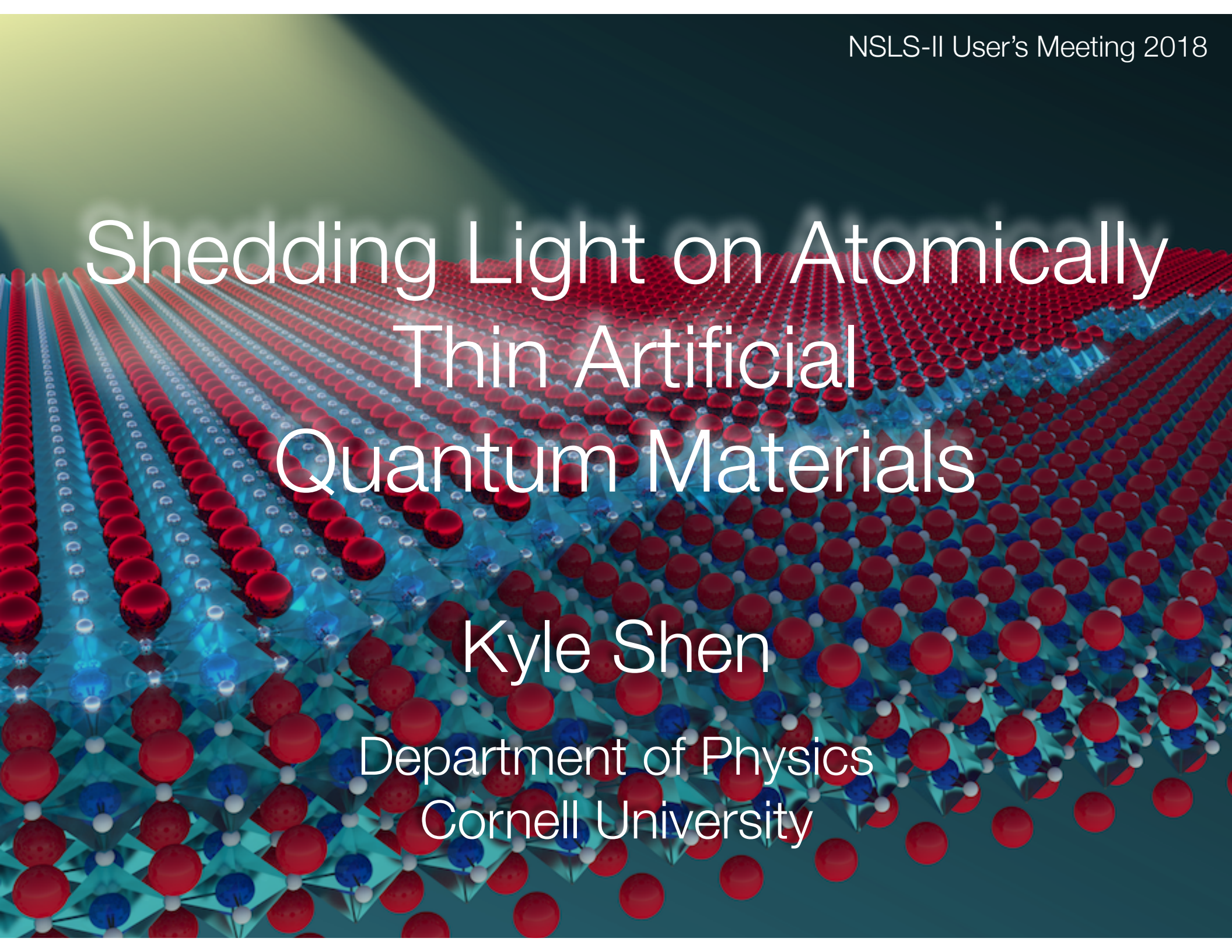


Shedding Light on Atomically Thin Artificial Quantum Materials

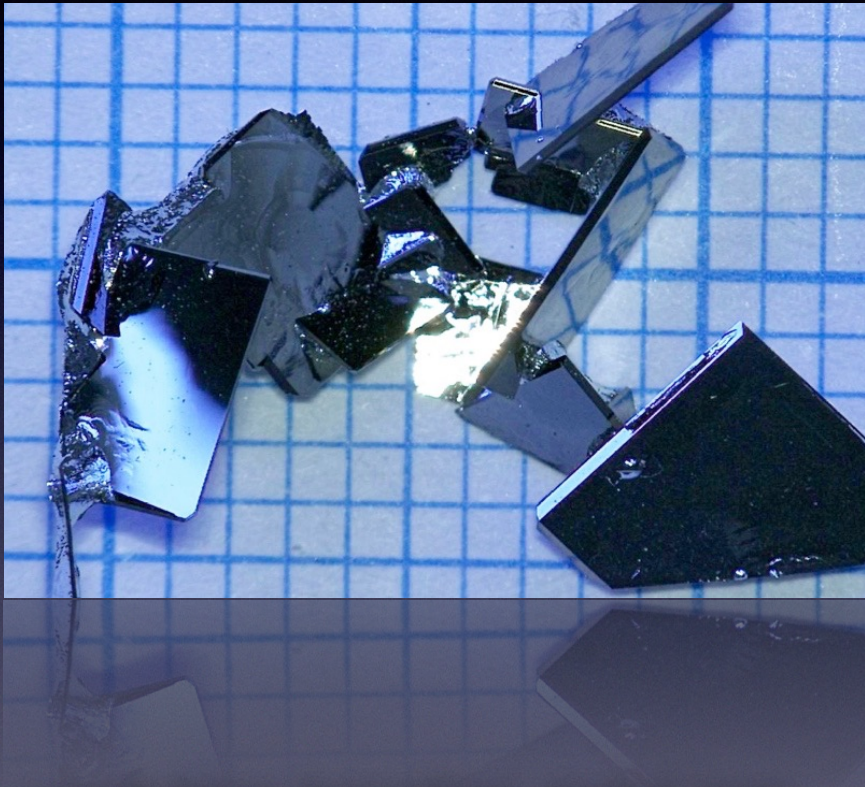
A 3D visualization of an atomically thin artificial quantum material. The structure consists of a lattice of red and blue spheres, with a central channel of blue spheres. The spheres are arranged in a regular, repeating pattern, and the overall structure is shown in a perspective view, receding into the distance. The background is a dark, gradient color, transitioning from light blue on the left to dark blue on the right.

Kyle Shen

Department of Physics
Cornell University

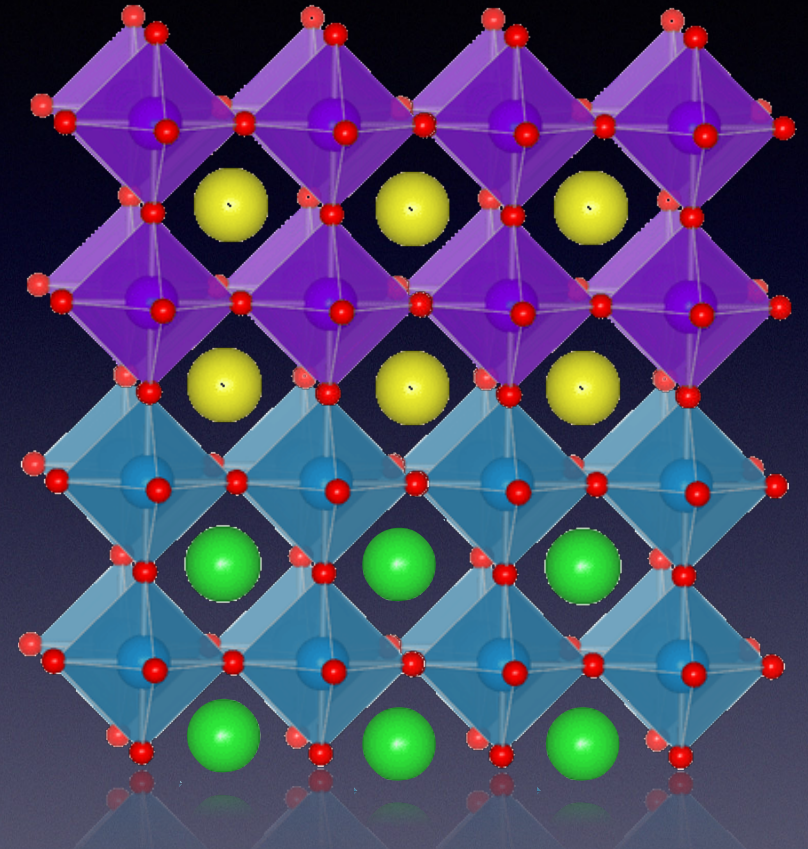
transition metal oxide interfaces : a new frontier

bulk materials



- switching transition metal ions
- chemical doping (cation substitution)
- lattice parameters (cation radius)

interfaces



- broken symmetry @ interface
- lowered dimensionality
- chemical potential offsets
- epitaxial strain (lattice constants)

interfaces made from complex quantum materials

- control of electron density
- dimensional confinement
- strength of interactions
- mass of the carriers
- magnetic interactions
- relativistic or classical

} conventional materials (Si, GaAs, Ge...)

} more exotic building blocks with tunable properties and interactions!

conventional
semiconductors



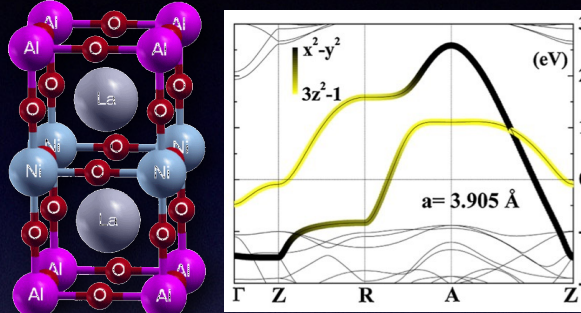
“correlated”
quantum materials



the grand challenge : controlling & optimizing quantum materials

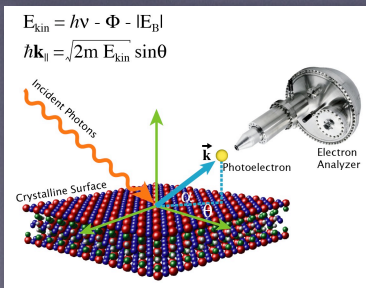
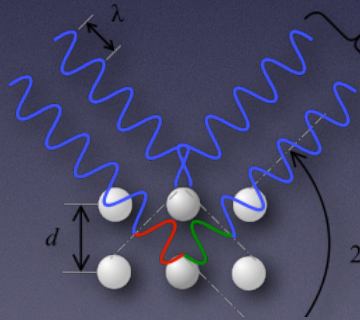
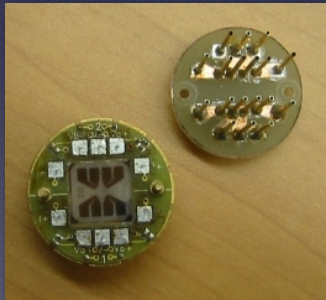
prediction

Superconductivity?

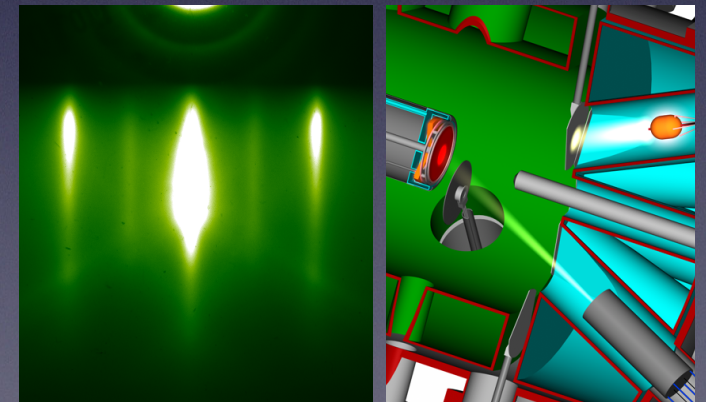


P. Hansmann *et al.*, *Phys. Rev. Lett* **103**, 016401 (2009)

characterization



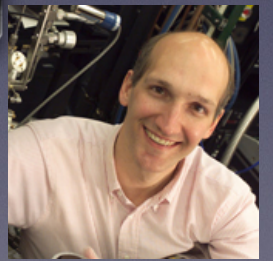
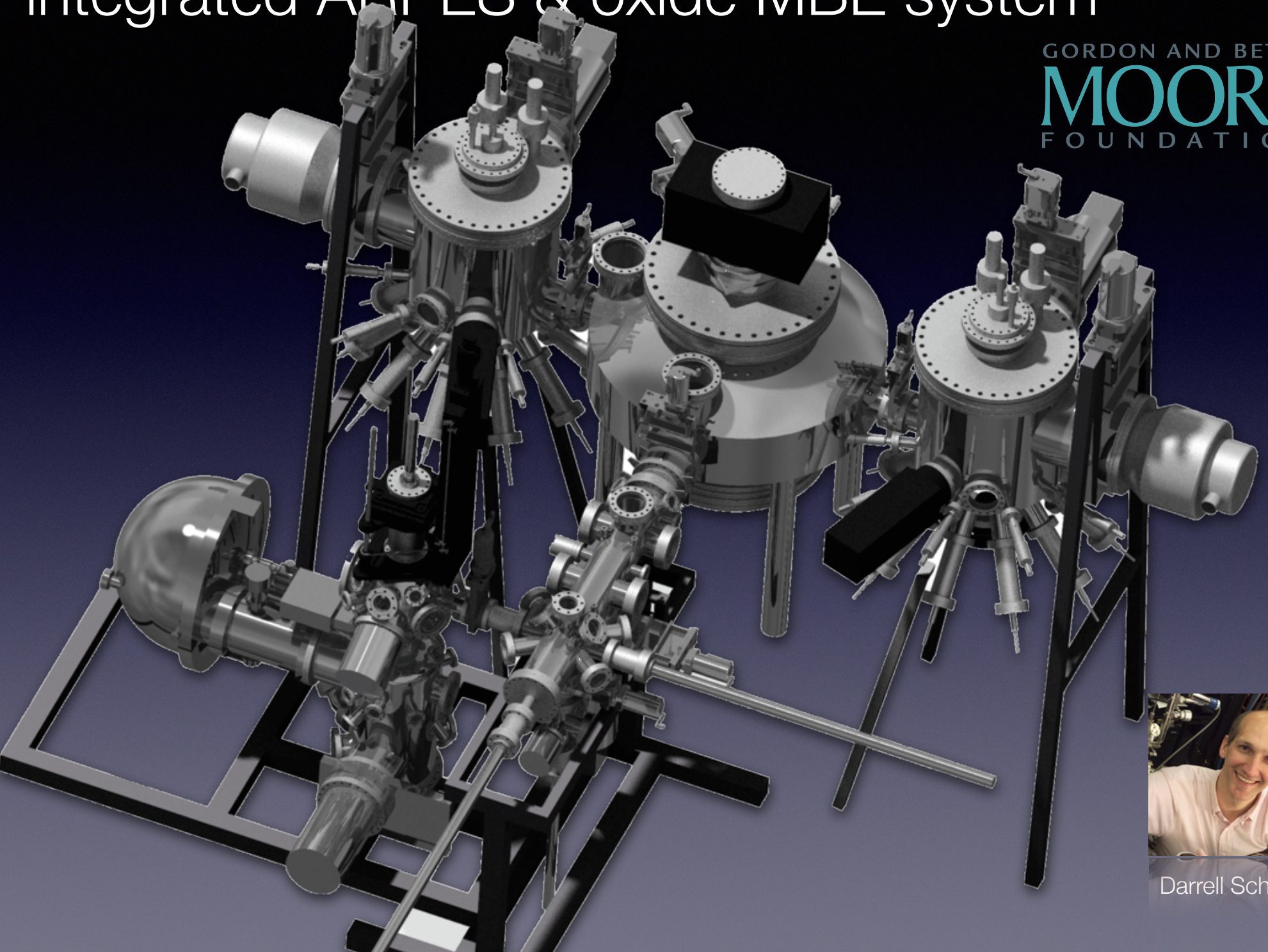
synthesis



Courtesy of Schlom Group

integrated ARPES & oxide MBE system

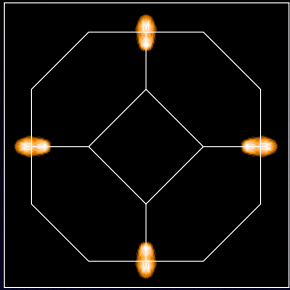
GORDON AND BETTY
MOORE
FOUNDATION



Darrell Schlom

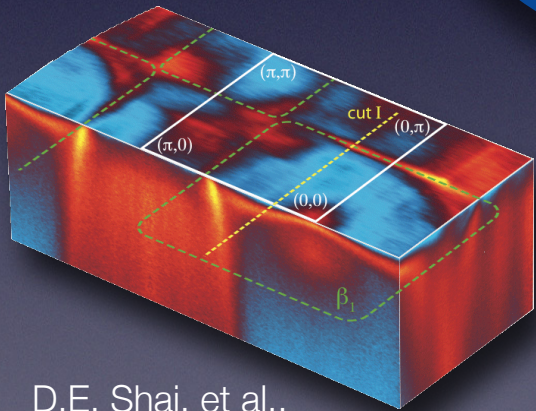
expanding the palette of available materials

EuO



D.E. Shai, et al.,
PRL 108, 267003
(2012)

SrRuO₃



D.E. Shai, et al.,
PRL 110, 087004
(2013)

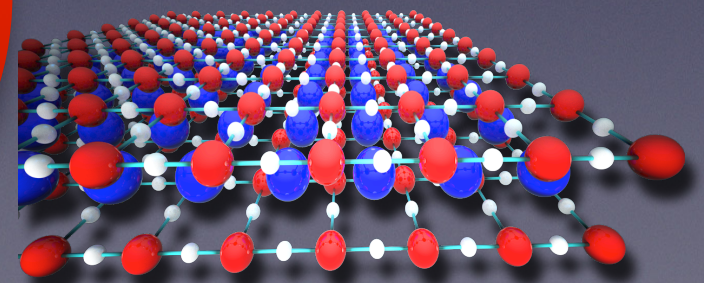
Bulk materials

Cleavable
single
crystals

Artificial
Heterostructures
&
Thin Films

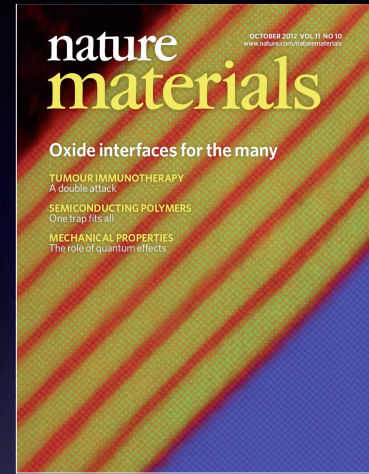
Metastable
Phases not
available
in bulk

SrCuO₂



J.W. Harter et al., *PRL* 109, 267001 (2012)

oxide
superlattices



E.J. Monkman et al.,
Nature Materials 11, 855
(2012)

recent examples in atomically thin interface materials

- quantum materials in the ultrathin limit : atomically thin LaNiO_3
- enhanced high- T_c superconductivity in monolayer FeSe grown on SrTiO_3

Rare-earth nickelates $RENiO_3$

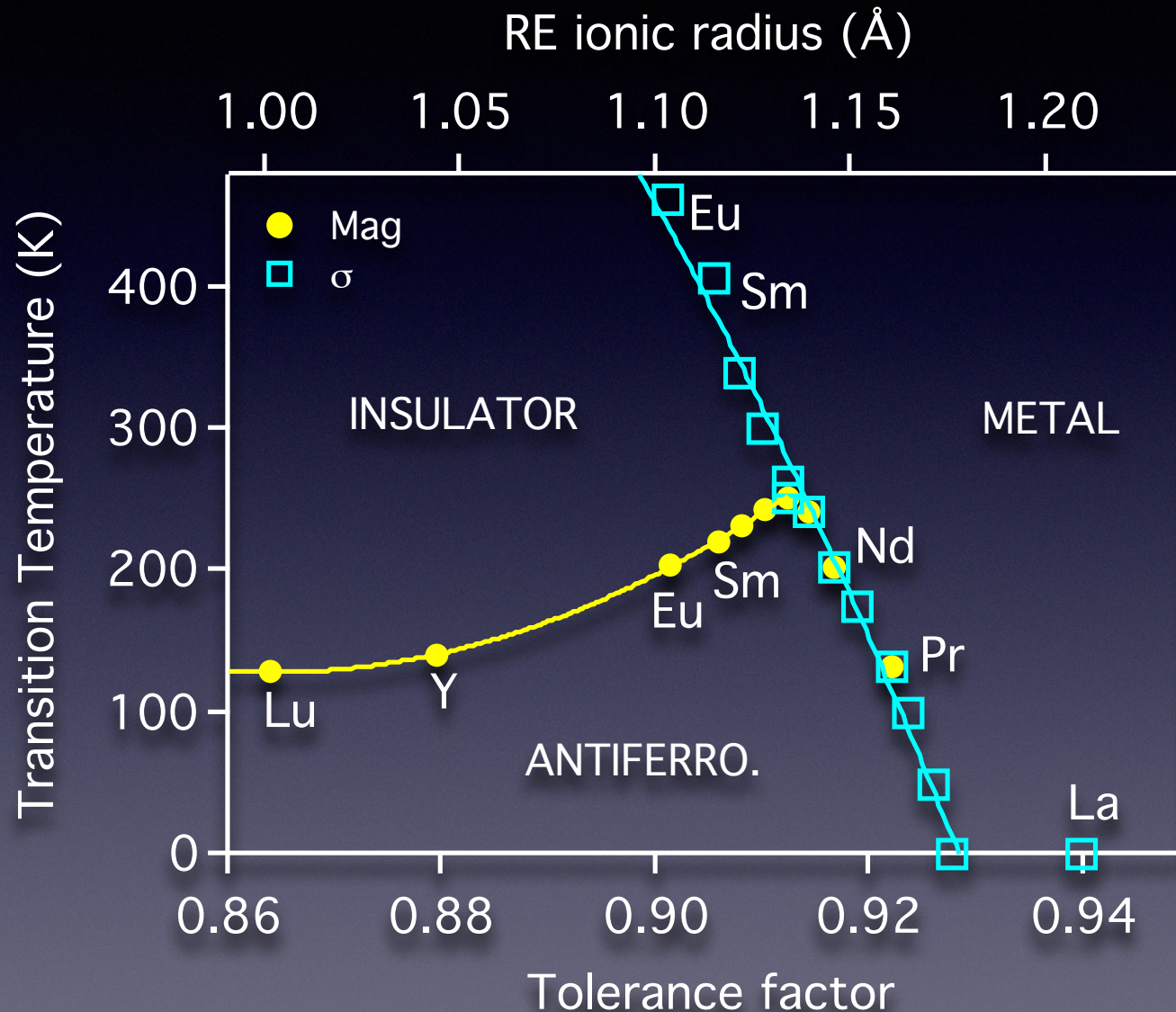
Phil King
(Kavli Fellow;
now faculty @ St.
Andrews Univ)



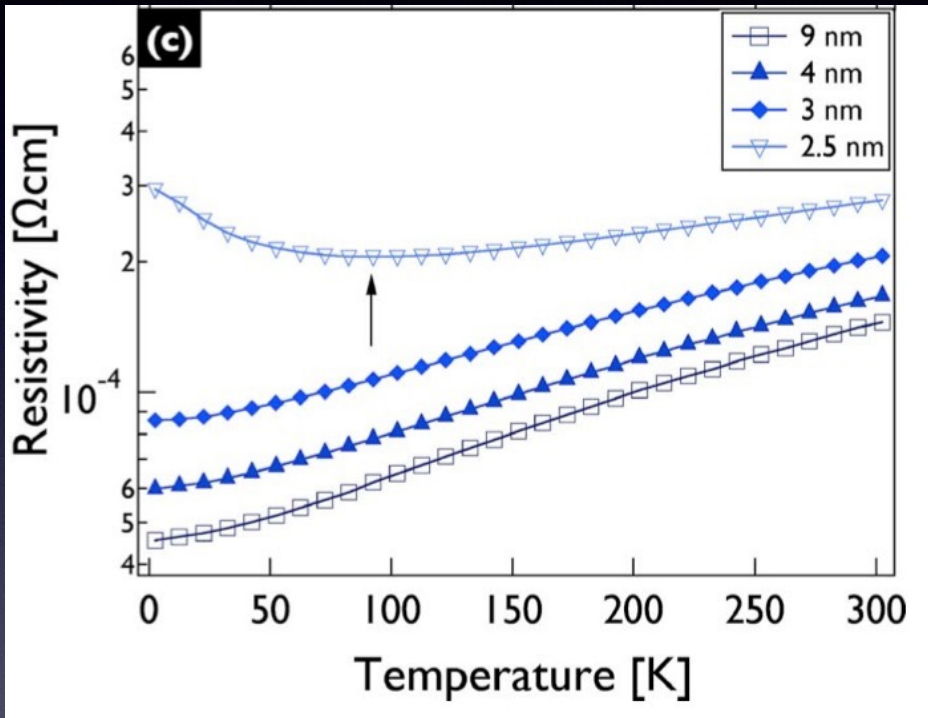
Masaki Uchida
(JSPS Fellow;
now Asst. Prof.
U. Tokyo)



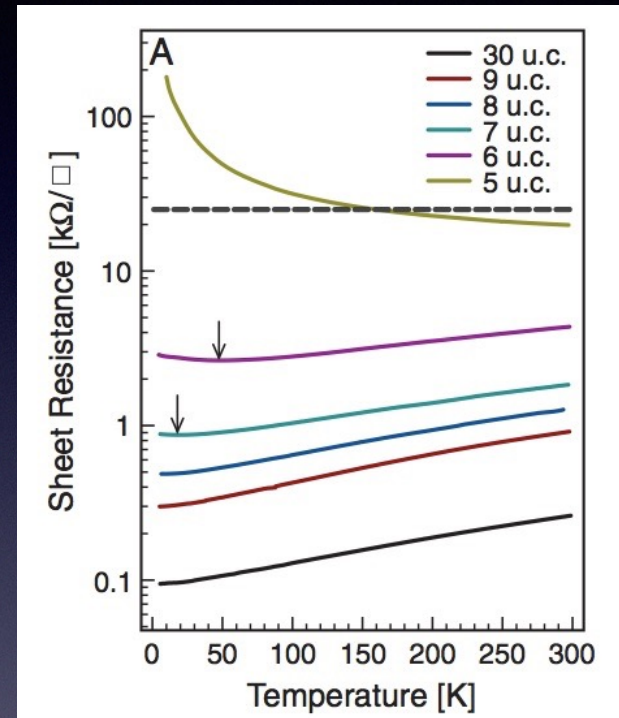
Haofei Wei



metal-insulator transitions in ultrathin LaNiO_3 films



J. Son *et al*, Appl. Phys. Lett. 96, 062114 (2010)

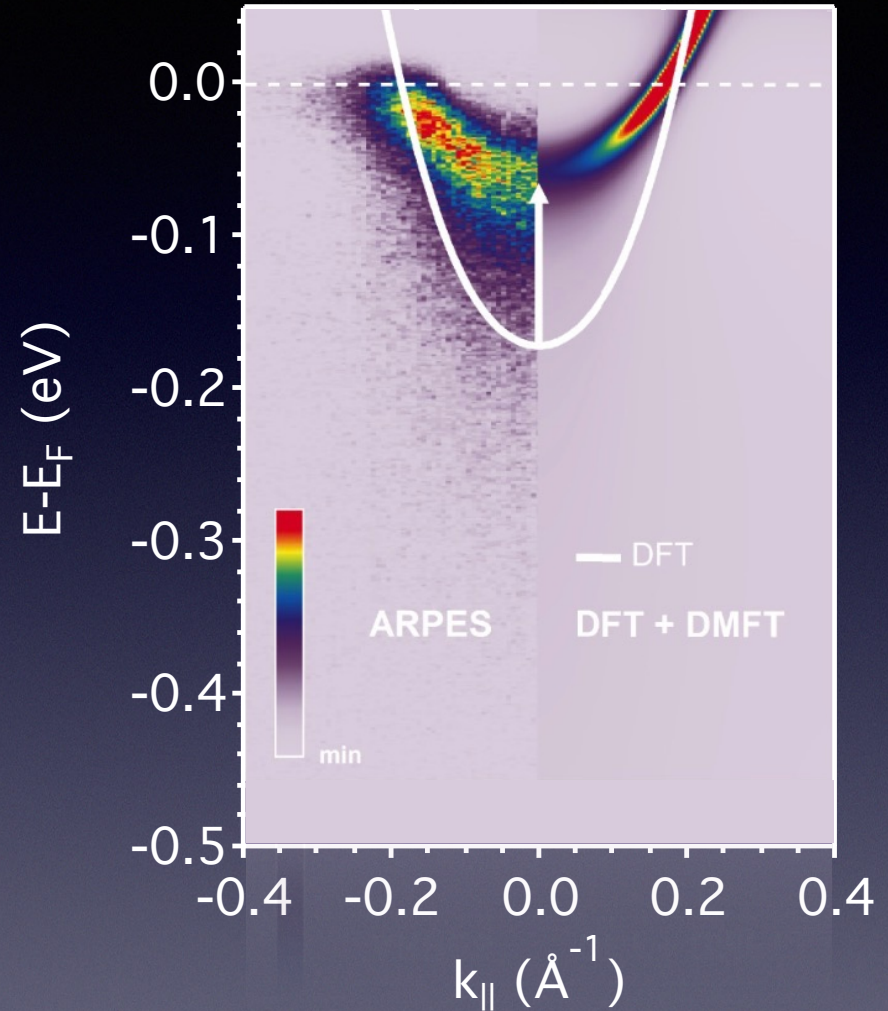
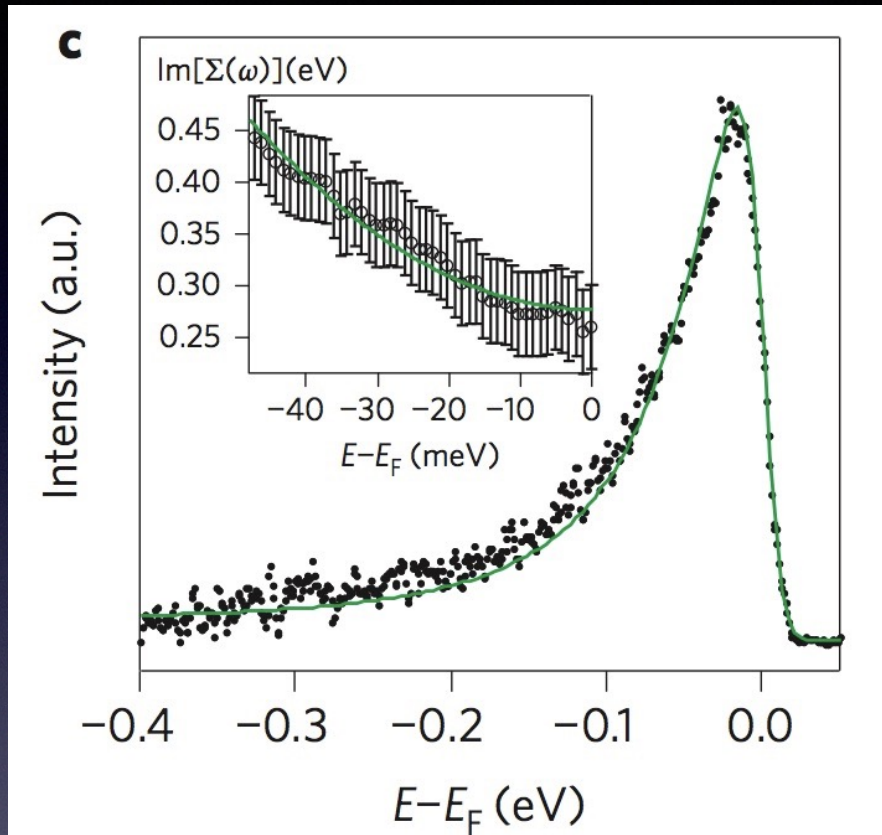


J. Scherwitzl *et al*, Phys. Rev. Lett. 106, 246403 (2011)

LaNiO_3 is a correlated metal in bulk, but reducing its thickness below a few unit cells turns the system progressively more insulating

bulk LaNiO_3 exhibits a large mass renormalization

LaNiO_3

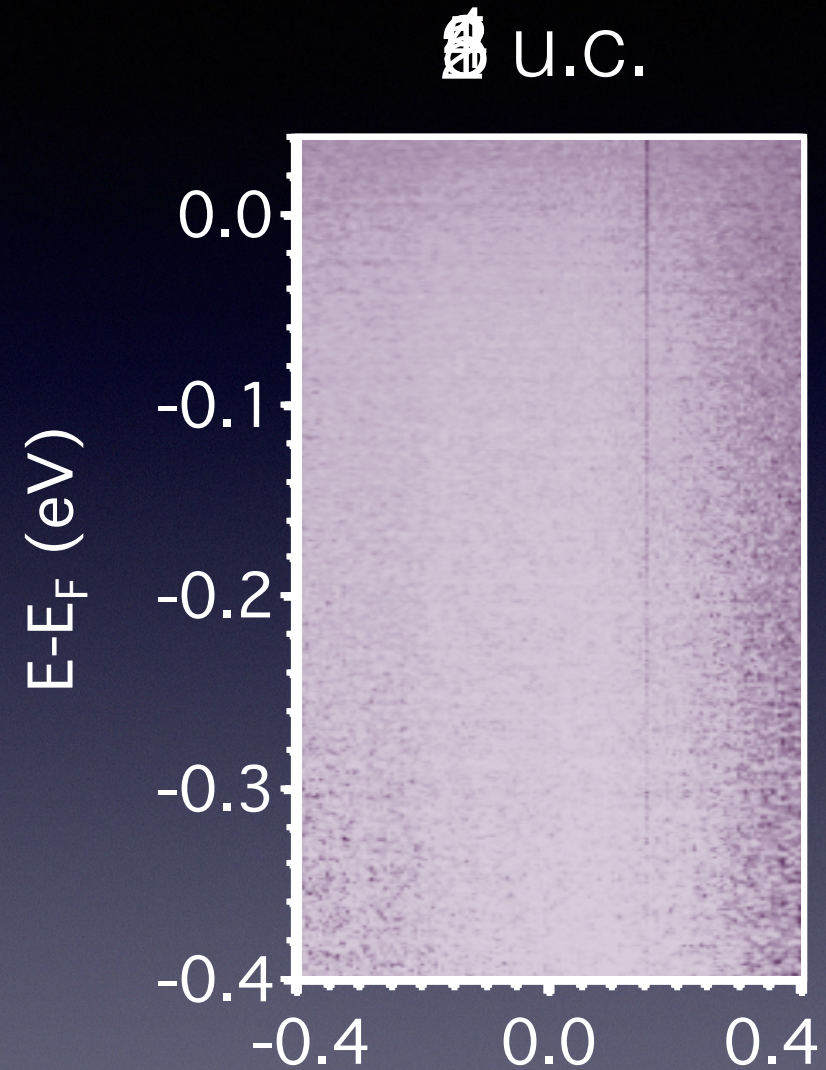
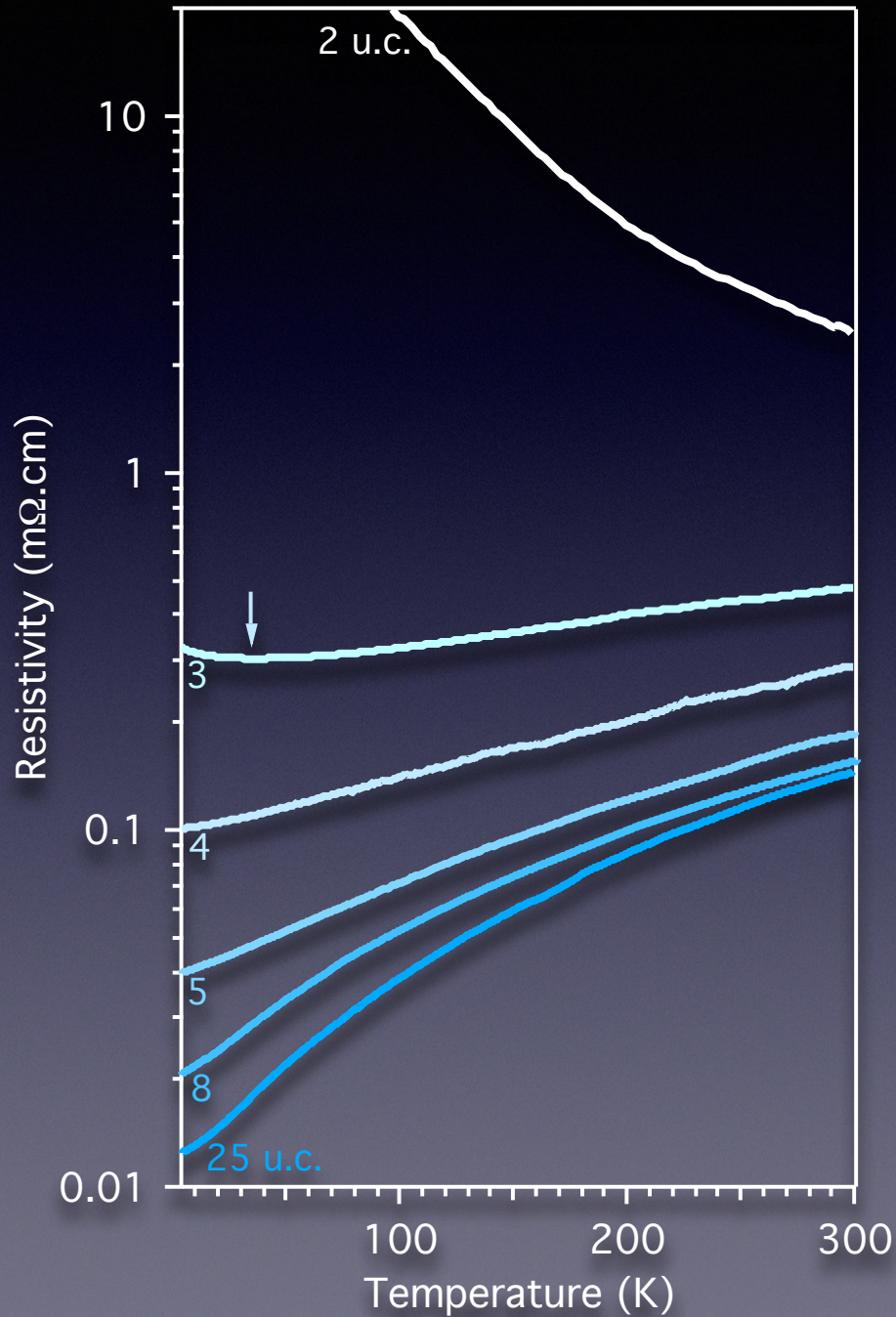


$$\frac{m^*}{m} = 3.3 \pm 0.5$$

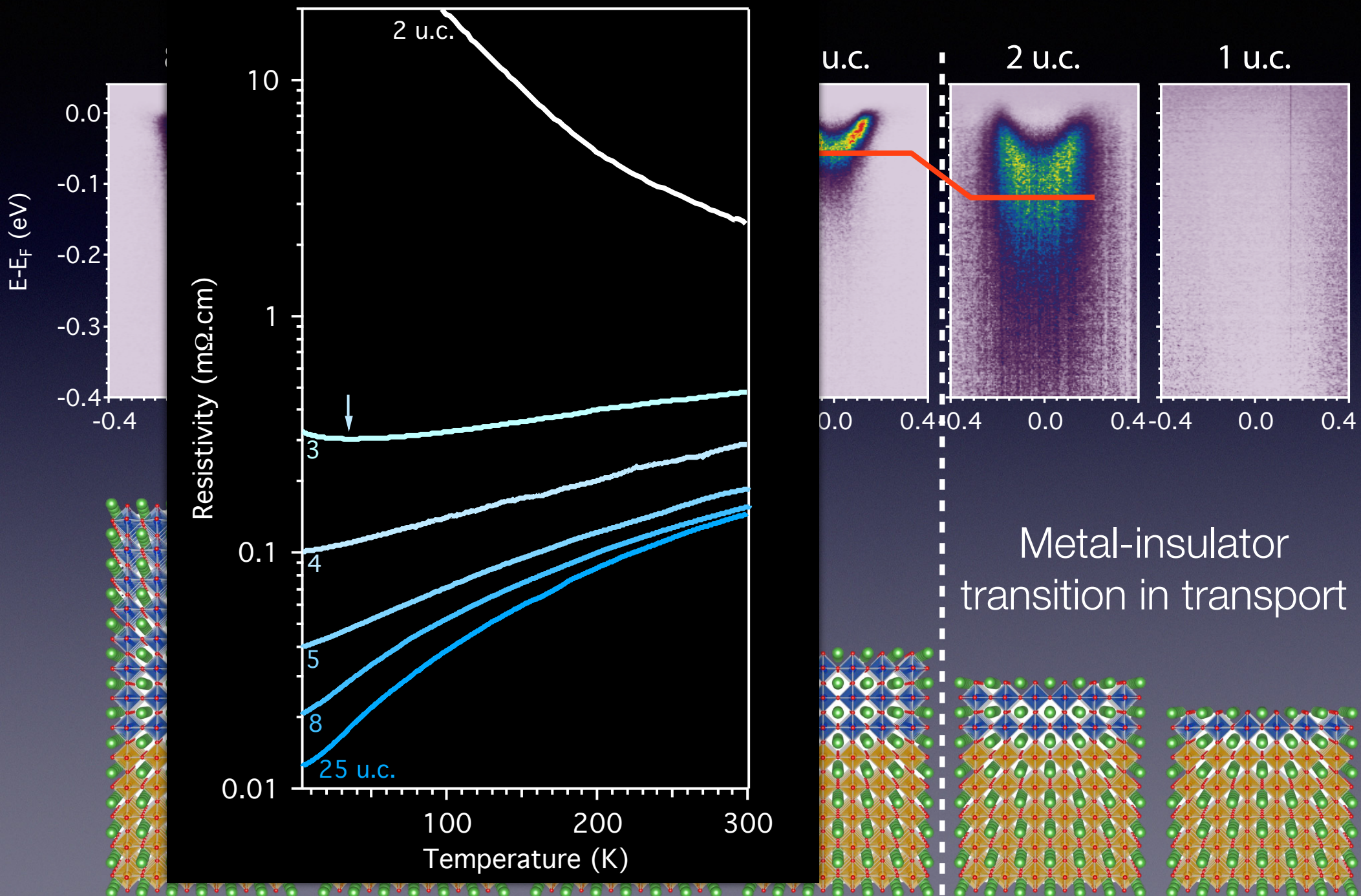
in agreement with optical & thermodynamic measurements

Metal-insulator transition at 2 unit cells

Decreasing film thickness

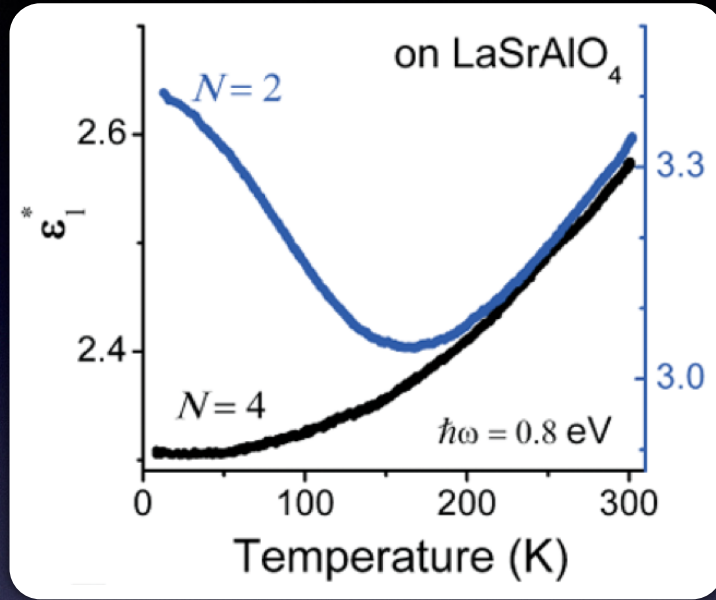


Electronic structure evolution with LaNiO_3 thickness



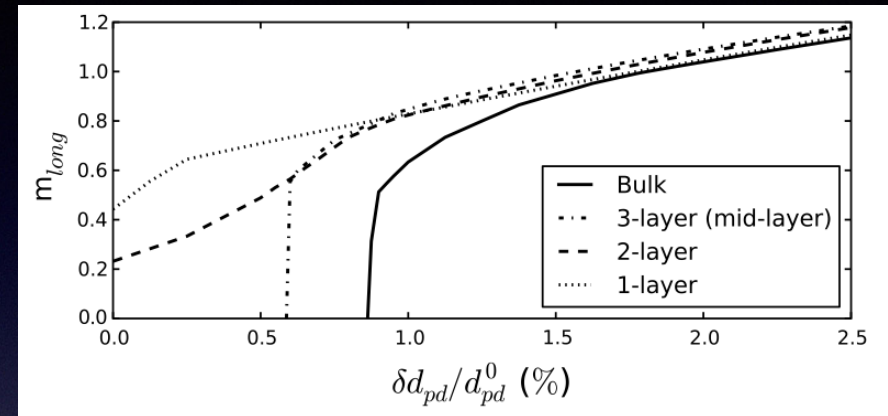
ordering tendencies in low-dimensional nickelates

Experimental

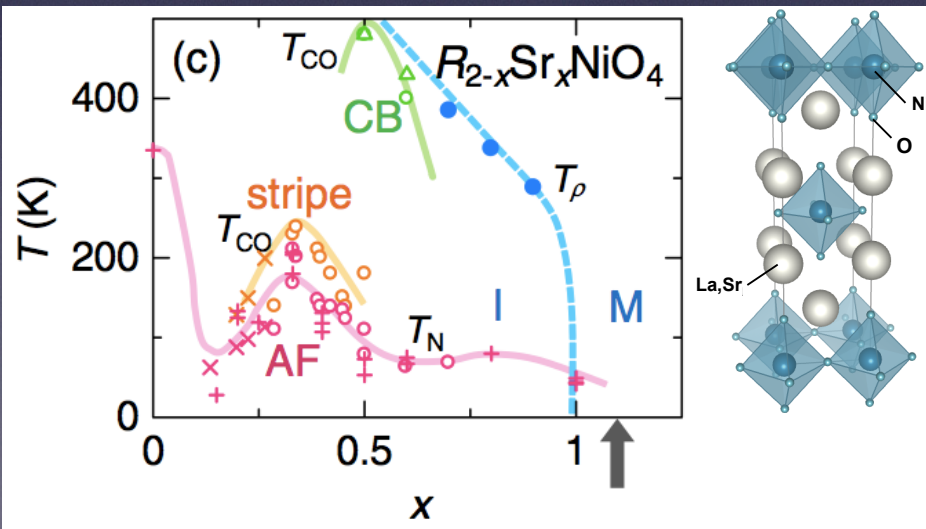


Boris et al., *Science* 332 (2011) 937

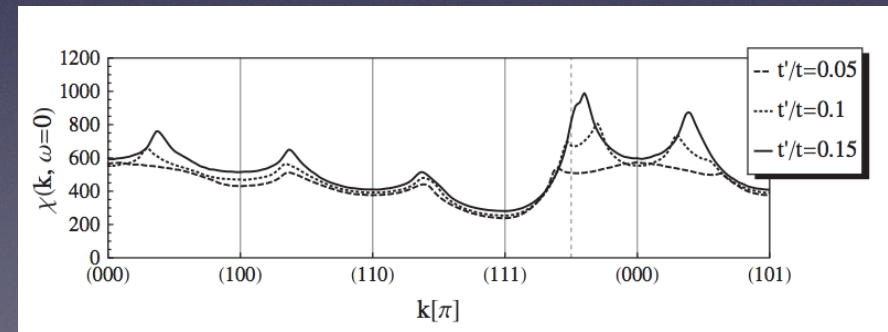
Theoretical



B. Lau & A.J. Millis, *Phys. Rev. Lett.* **110**, 126404 (2013)



M. Uchida et al., *Phys. Rev. Lett.* **106**, 027001 (2011)

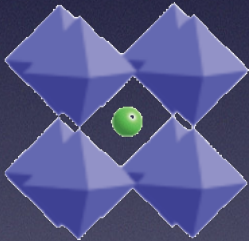


S. Lee, R. Chen & L. Balents, *Phys. Rev. Lett.* **106**, 016405 (2011)

dimensionality plays an important role in layered nickelates

Ruddlesden-Popper Series : $\text{La}_{n+1}\text{Ni}_n\text{O}_{3n+1}$

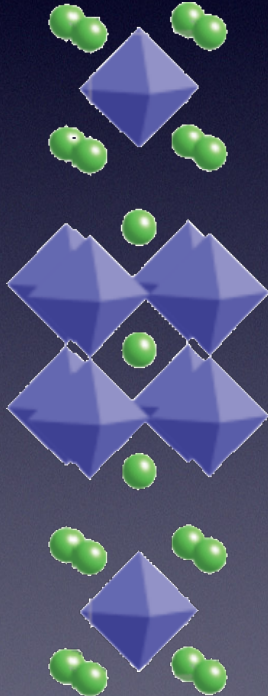
$n = \infty$



paramagnetic,
correlated metal

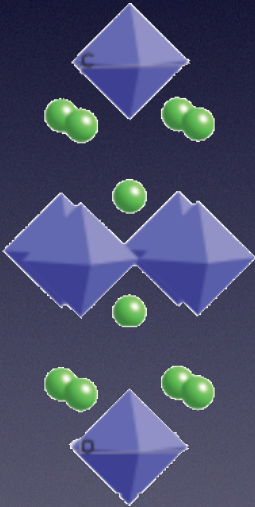


$n = 2$



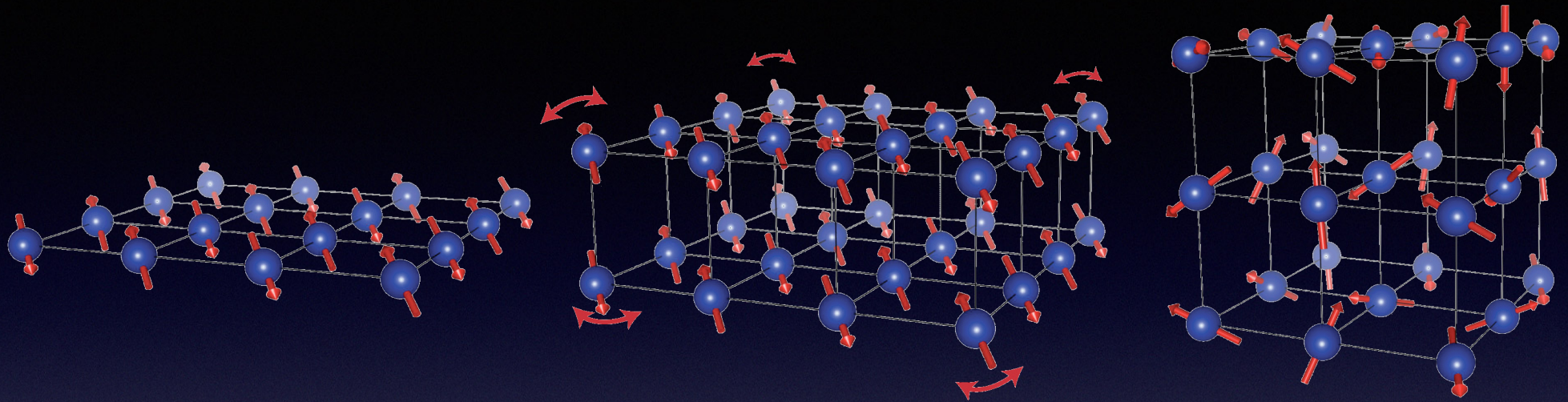
paramagnetic
“bad metal”

$n = 1$

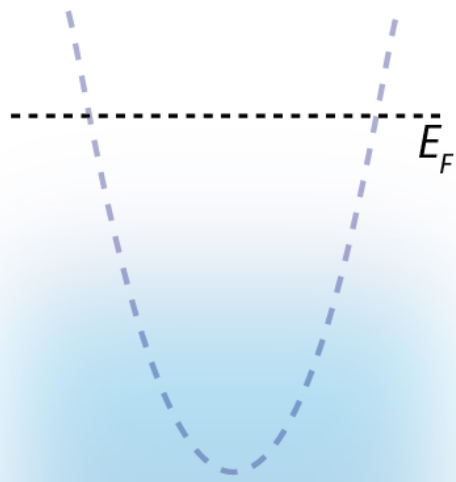


correlated
antiferromagnetic
insulator, insulating
charge / spin stripes

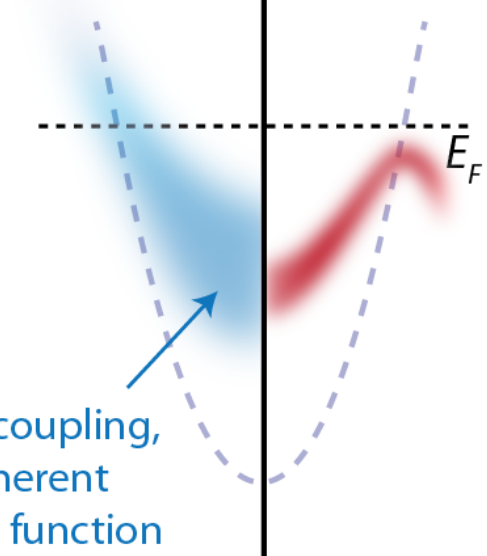
Ordering instability likely onsets below 3 unit cells



Static charge/spin order
Fully-gapped insulator

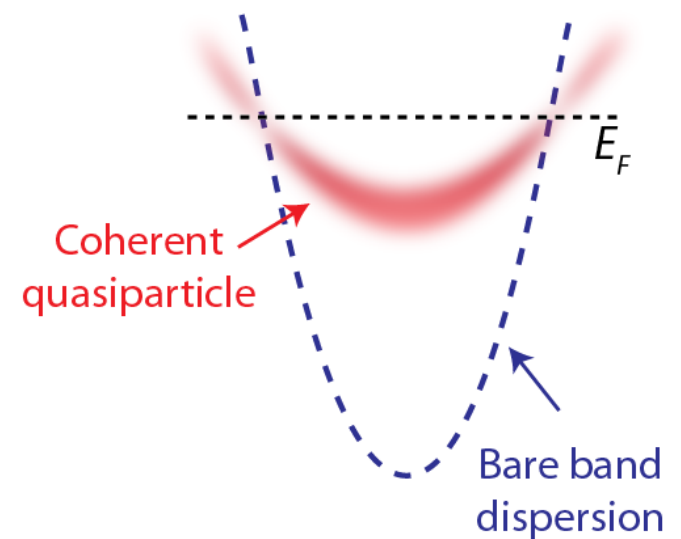


Incipient order
Fluctuating | Weak coupling



Strong coupling,
incoherent
spectral function

Strongly-interacting
Fermi liquid metal



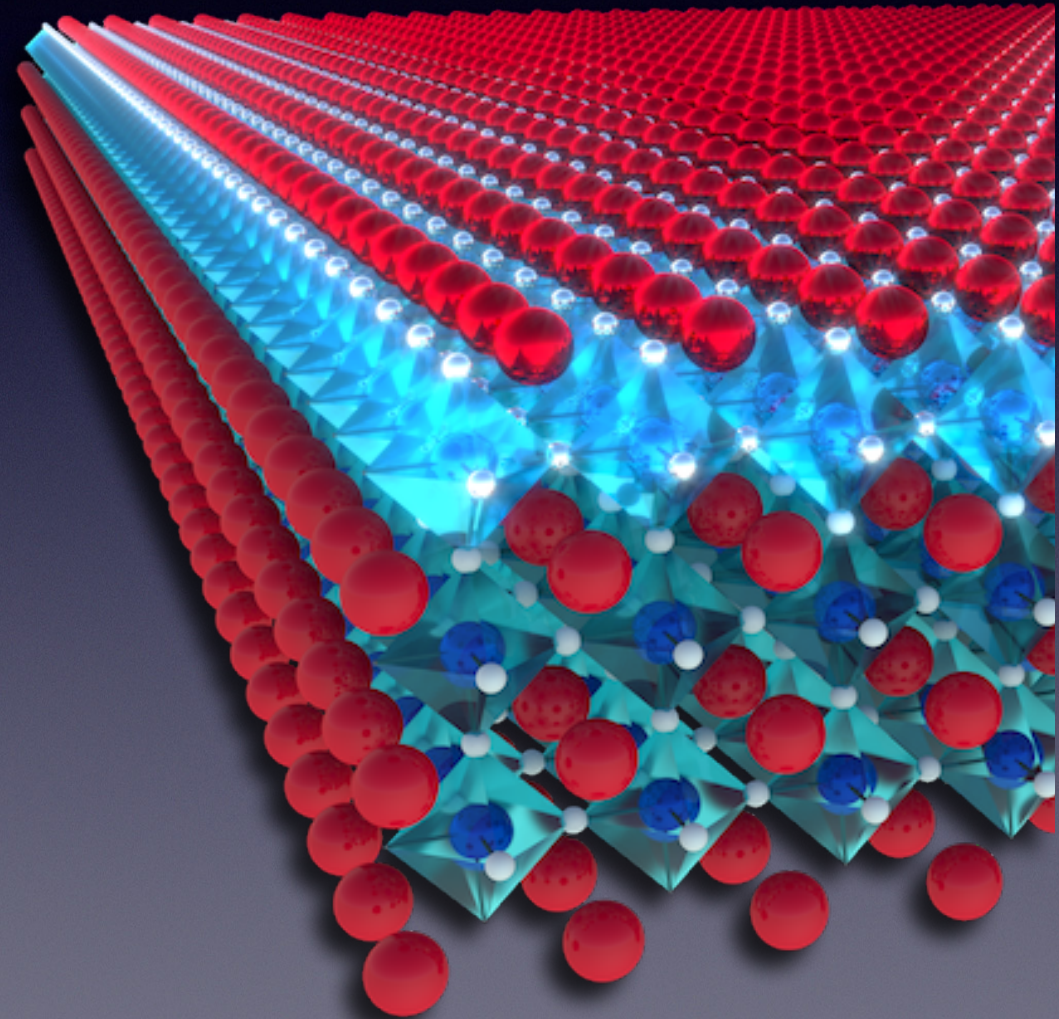
Coherent
quasiparticle

Bare band
dispersion

Scientific Needs from NSLS-II

- can we directly detect static spin/charge ordering stabilized in the ultrathin limit?
- can we observe short-ranged spin fluctuations before static long-ranged order onsets?
- what is the effect of octahedral rotations and/or polar distortions near the film-substrate interface?

Ultrathin
 LaNiO_3



recent examples in atomically thin interface materials

- quantum materials in the ultrathin limit : atomically thin LaNiO_3
- enhanced high- T_c superconductivity in monolayer FeSe grown on SrTiO_3

characteristics of superconductors

nature of the superconducting wavefunction

$$\Psi(r_1, s_1; r_2, s_2) = \phi(r_1, r_2) \chi(s_1, s_2)$$

pair wavefunction must be antisymmetric

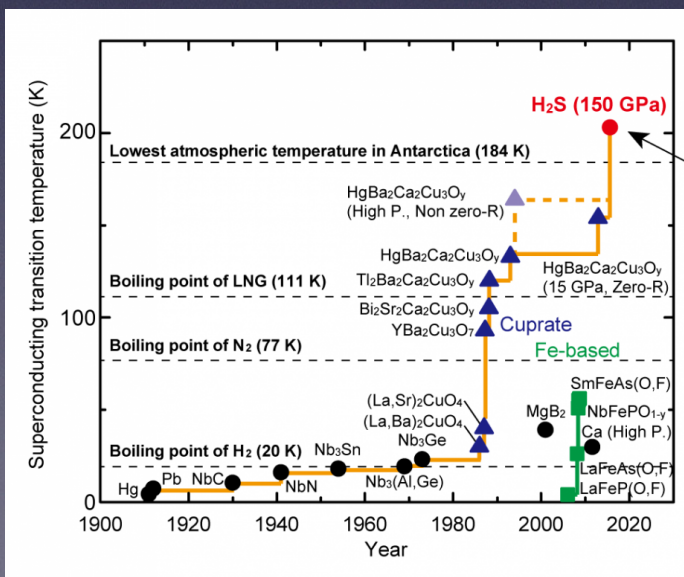
$$\left\{ \begin{array}{l} \phi(r_1, r_2) \text{ spatial part} \\ \chi(s_1, s_2) \text{ spin part} \end{array} \right.$$

~ 99% of SC : $\chi(s_1, s_2) \propto (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$
antisymmetric : spin singlet

conventional BCS (Pb, Al, Nb), high Tc cuprates, Fe-SC

very rare : $\chi(s_1, s_2) \propto |\uparrow\uparrow\rangle$
symmetric : spin triplet

superconducting transition temperature (T_c)



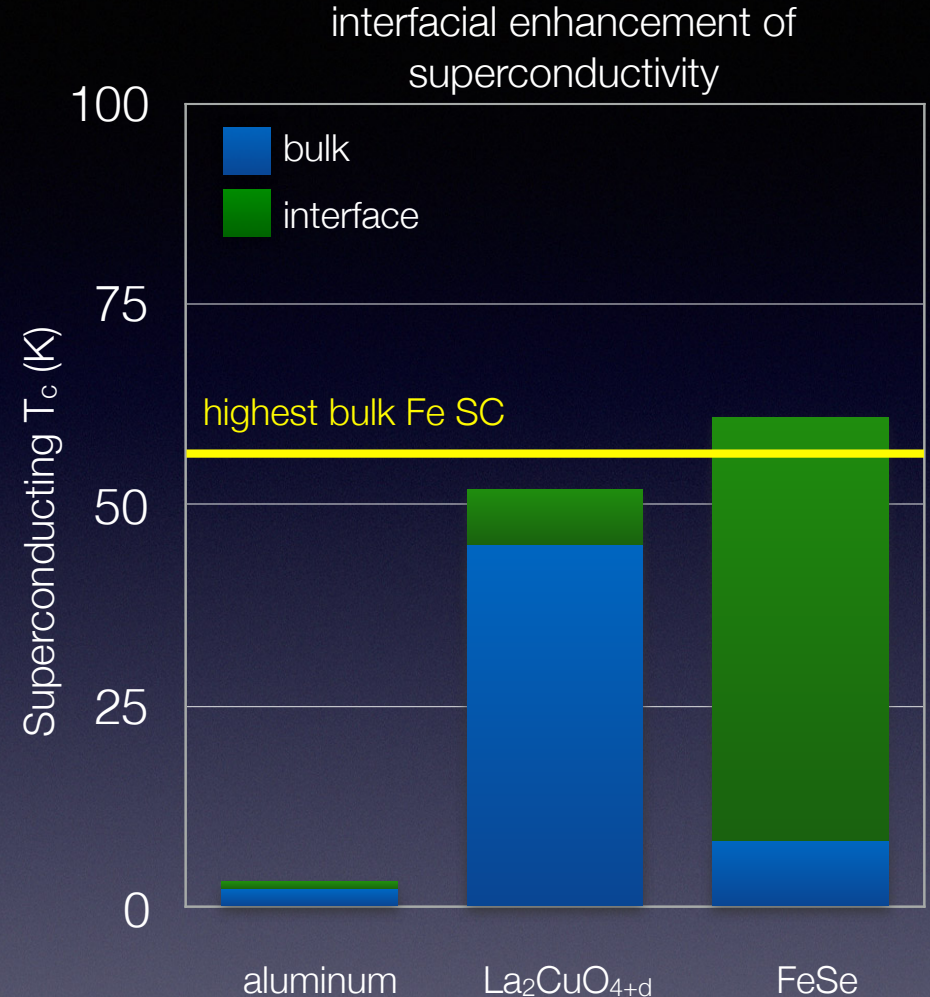
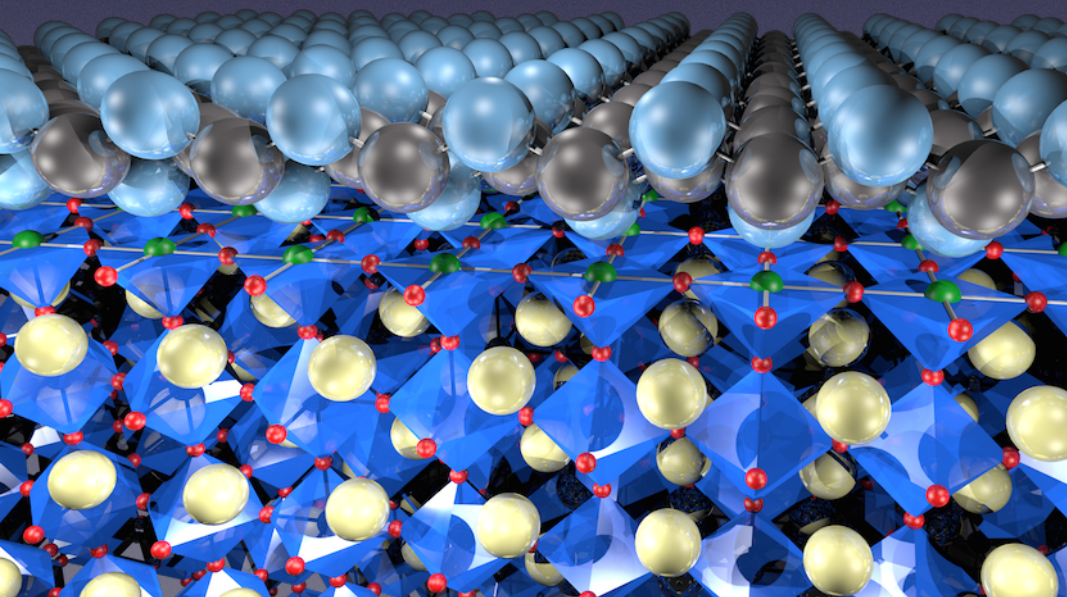
highest T_c known superconductor : H₂S @ 203 K

highest T_c ambient SC : HgBa₂Ca₂Cu₃O₈ @ 134 K

highest T_c Fe-based SC : SmFeAsO ~ 56 K

interfacial enhancement of superconductivity

- idea of interfacial enhancement of SC dates back to Ginzburg (1964)
- small (10%) effects observed in metal films and cuprate thin films
- in bulk, FeSe has a T_c of 8 K, and is structurally the simplest of the Fe-based SCs
- initial measurements of $\sim 60\text{-}70$ K T_c 's of FeSe / SrTiO₃ performed by STM and ARPES



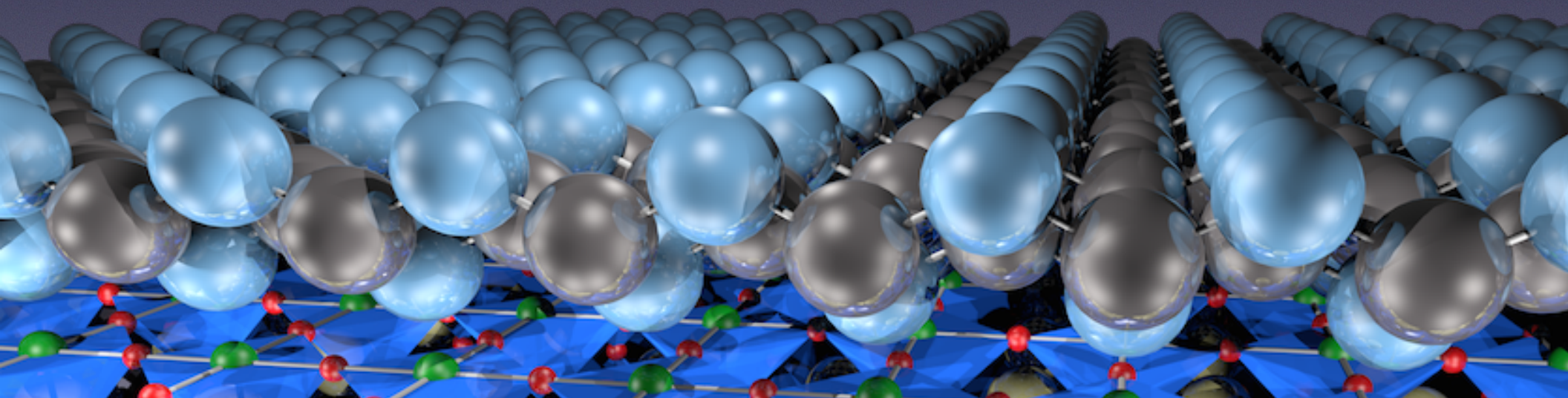
- V. Ginzburg. *Phys. Lett.* **13**, 101 (1964)
M. Strongin *et al.*, *Phys. Rev. Lett.* **21**, 1320 (1968)
I. Bozovic *et al.*, *Phys. Rev. Lett.* **89**, 107001 (2002)
Q.Y. Wang *et al.*, *Chin. Phys. Lett.* **29**, 017401 (2012)

central questions about monolayer FeSe / SrTiO₃

- what is the mechanism of the interfacial enhancement of T_c ?
- how does the T_c measured spectroscopically compare to other probes?
- what happens to the superconductivity as the FeSe film gets thicker?

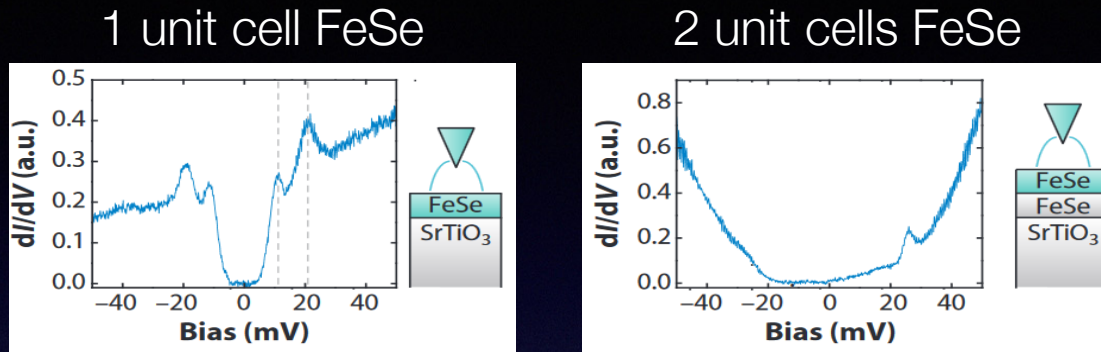
challenges

- superconductivity of FeSe / SrTiO₃ monolayer does not survive removal into air (unless samples are capped)
- capped samples appear to exhibit lower T_c 's when measured *ex situ* than reported by *in situ* spectroscopic probes
- a given sample is not measured by different techniques; T_c 's cannot be compared



what happens as the FeSe film gets thicker?

in situ probes (ARPES, STM) see no gap for films thicker than 1 UC



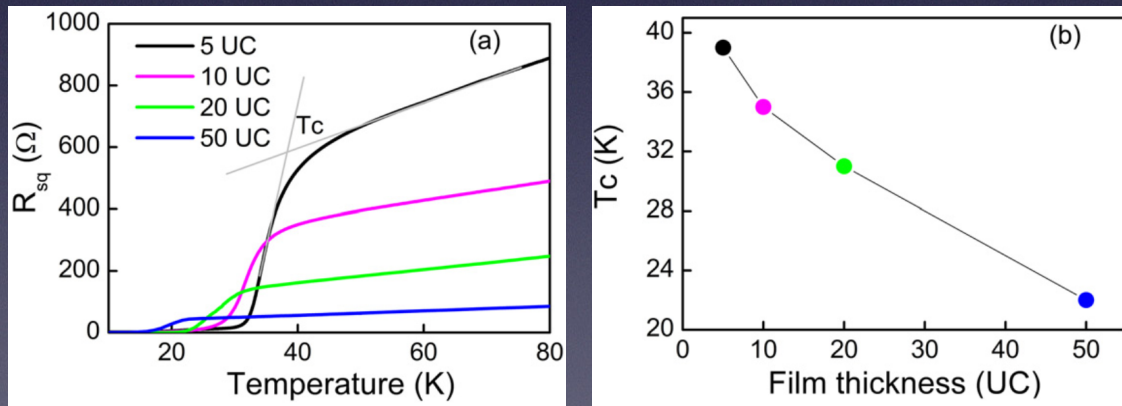
Wang et al., *Chin. Phys. Lett.* **29**, 037402 (2012)

- superconducting gap disappears for all films thicker > 1 uc

surface-sensitivity cannot see buried interface?

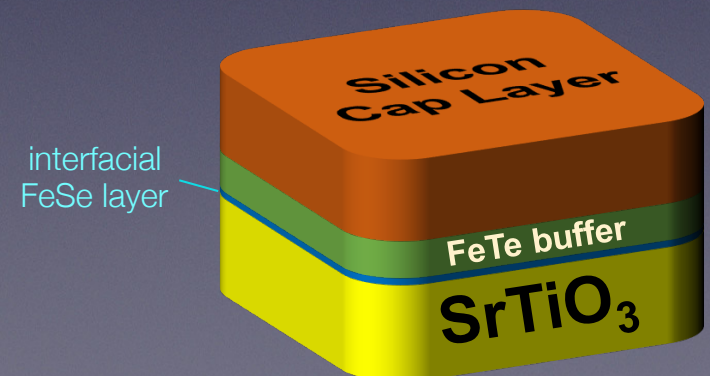


ex situ measurements on capped films



Wang et al., *2D Materials* **2**, 044012 (2015)

- SC observed for thicker capped films, but T_c appears to decrease with film thickness
- capping layers believed to reduce T_c

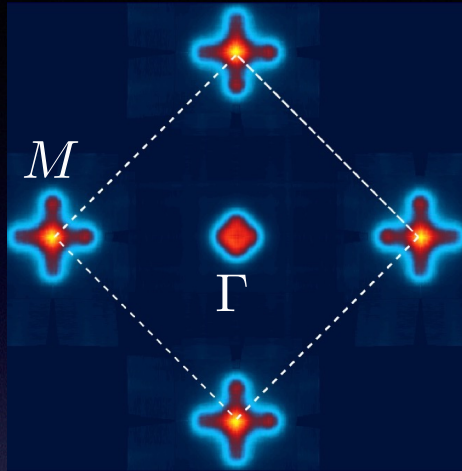


ARPES measurements as a function of film thickness

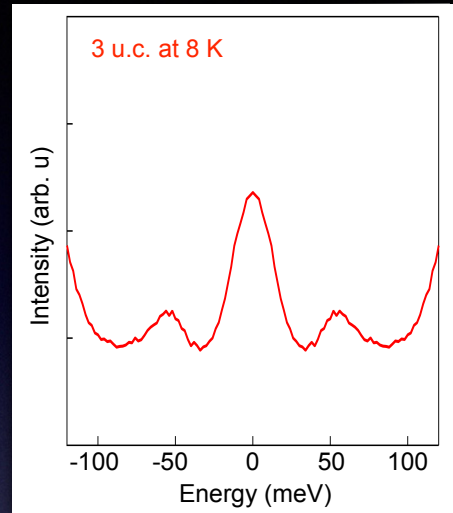
3 unit cell
FeSe



bulk-like Fermi surface



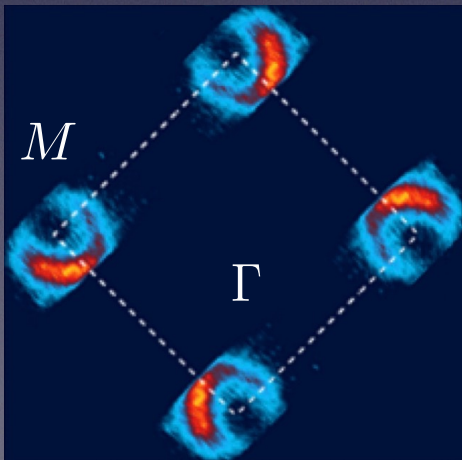
no detectable SC gap



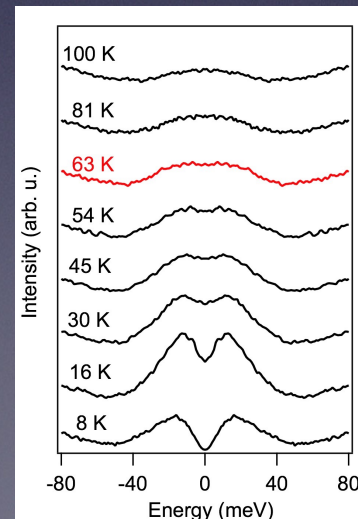
- ARPES data for 3 unit cell sample nearly identical to bulk FeSe
- No observable SC gap measured down to base temperatures

missing hole pocket at Γ

1 unit cell
FeSe



15 meV gap



gap closes ~ 60 K

