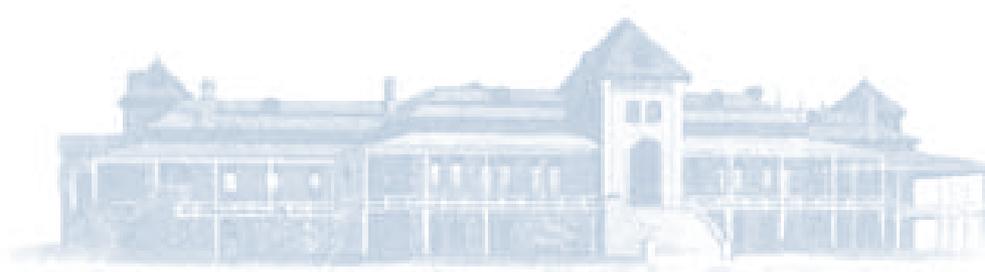


Jet Substructure in ATLAS at the LHC - a Tool for Discoveries and Measurements

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Seminar
**Brookhaven National
Laboratory**
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Roadmap for this Talk

Jet substructure

Introduction to jet substructure phenomenology
Selective review of most popular techniques

Jet reconstruction inputs

The ATLAS detector
Jet reconstruction inputs

Substructure reconstruction performance

Reconstructing observables in the presence of pile-up

Measurements & Searches

QCD jet structure
Application of substructure techniques in final states with top quarks

Conclusions and outlook

Jet Substructure



Kinematic reach at LHC allows production of (highly) boosted particles

All decay products of hadronically decaying Standard Model particles can be collected into single jet

$W \rightarrow qq$, Higgs bosons $H \rightarrow bb$, top quarks $t \rightarrow Wb \rightarrow qqb$

Searches for new heavy particles with boosted (SM) decay products

Single jet mass indicative observable for new particle production

Experimental challenges – detector granularity & pile-up

Limited granularity introduces limits in resolving sub-structure in a jet

E.g., signals from near-by particles can easily overlap in calorimeters

Presence of pile-up can disturb single jet mass reconstruction from internal flow measurements

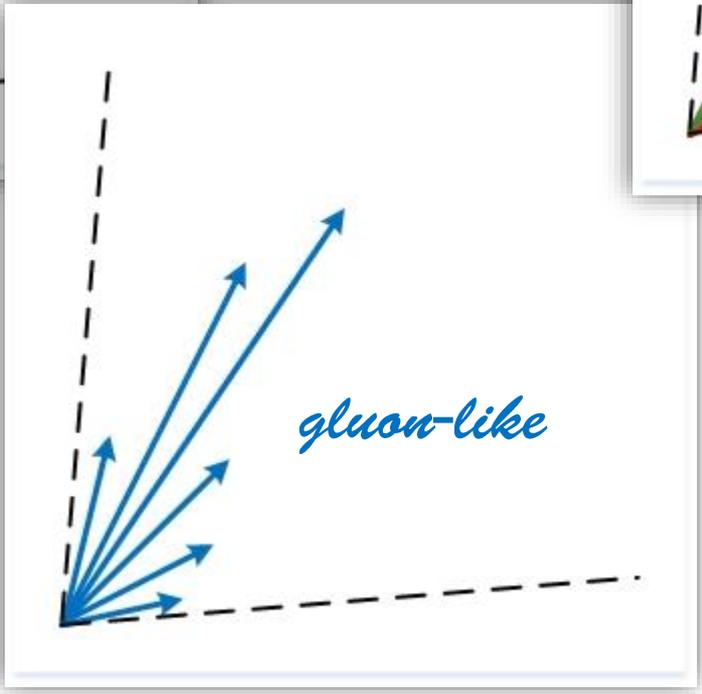
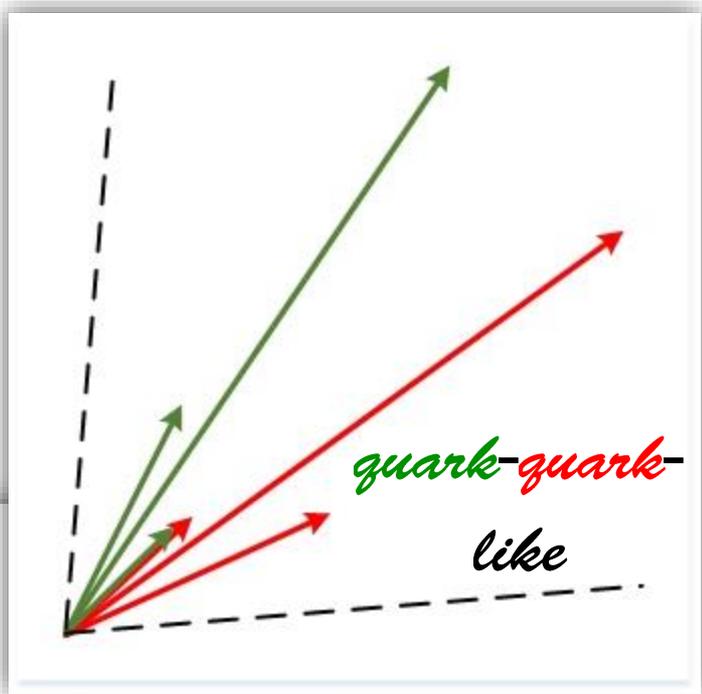
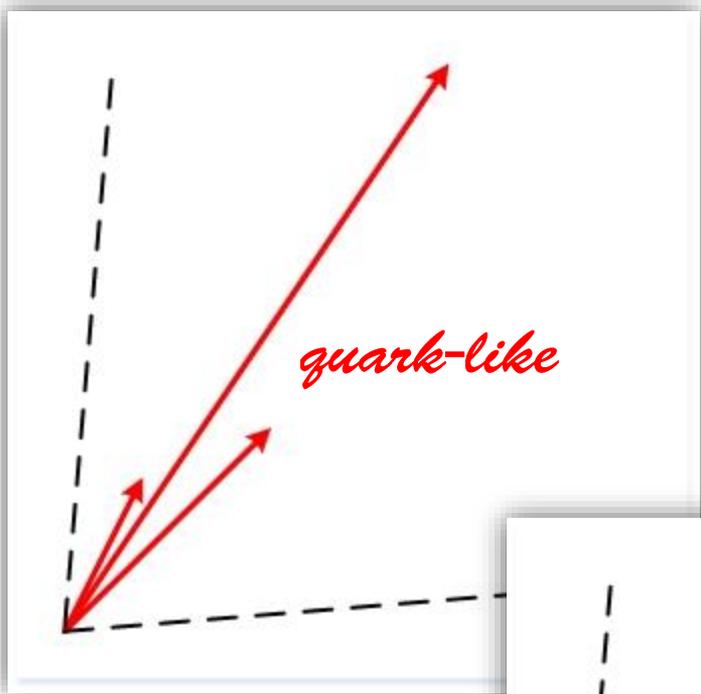
Needs to extract relevant internal jet energy flow structures for mass reconstruction etc. from diffuse pile-up contributions

Jet substructure analysis

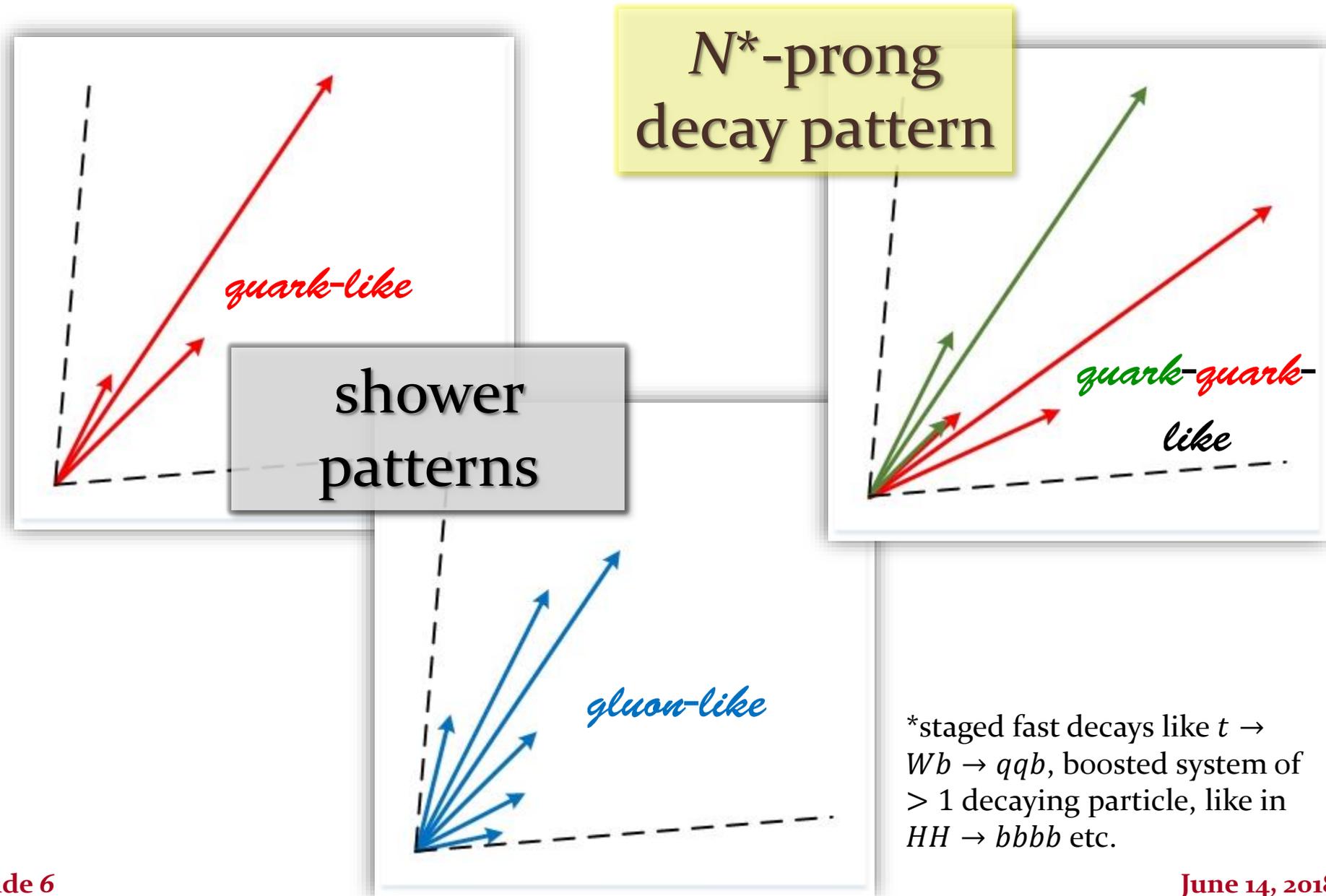
Collection of techniques aiming at enhancing two- or three-particle (“prongs”) decay patterns in single jets

Typically leads to suppression of QCD-like backgrounds from quark- and gluon jets with their typical parton shower and fragmentation driven internal flow structure

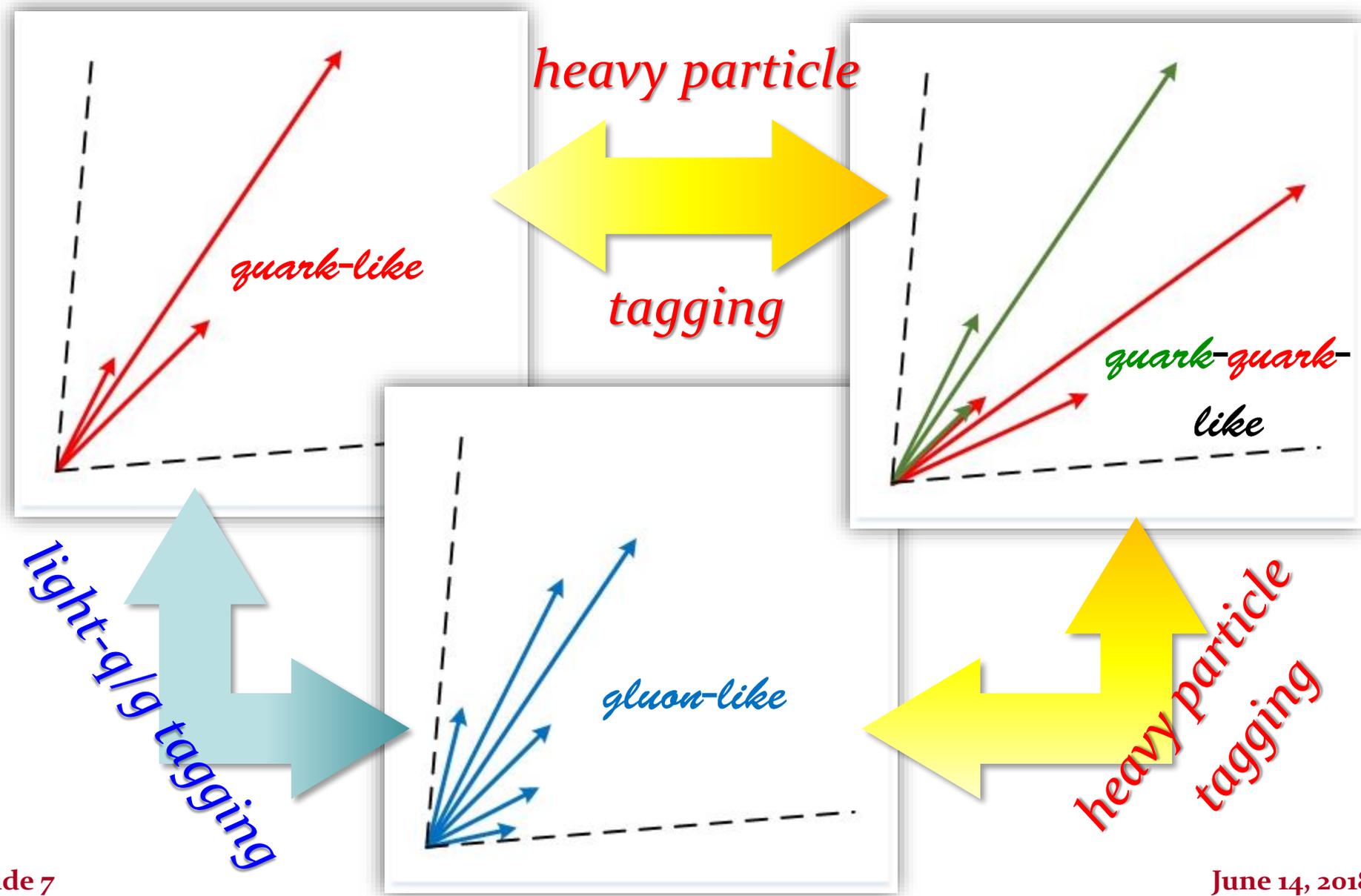
Separation of Flow Patterns



Separation of Flow Patterns



Separation of Flow Patterns



Particle flow inside a jet hints to source

Boosted (heavy) particle decay emits particles in typical jet cones

Standard Model as well as new particles

Usefulness depends on the ability to resolve decay structure

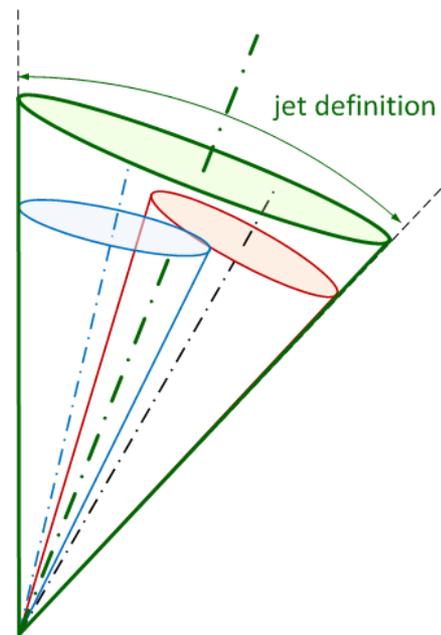
E.g., 2-prong (like W) or 3-prong (top) decays

Resolution scale given by (hypothetic) mass of particle – limited by detector capabilities

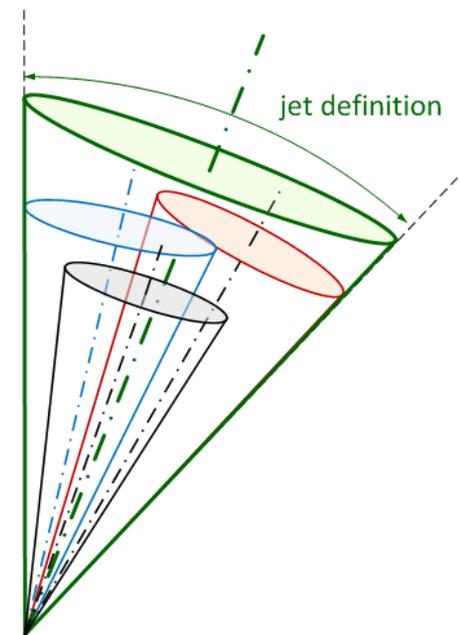
Jet size

“Rule of thumb” $R = 2m/p_T$

Cone of size R contains decay products of particle with mass m and transverse momentum p_T



2 – prong decay inside reconstructed jet, e.g. from $W \rightarrow q\bar{q}$ (SM) or heavy new object like $\phi \rightarrow gg$ or $Z' \rightarrow q\bar{q}$ (BSM)



3 – prong decay inside reconstructed jet, e.g. from $t \rightarrow q\bar{q}b$ (SM) or heavy new object like $\phi_{KK} \rightarrow Q\bar{Q}b + X$ or $t' \rightarrow q\bar{q}b$ (BSM)

Example for Boosted Particle Decays

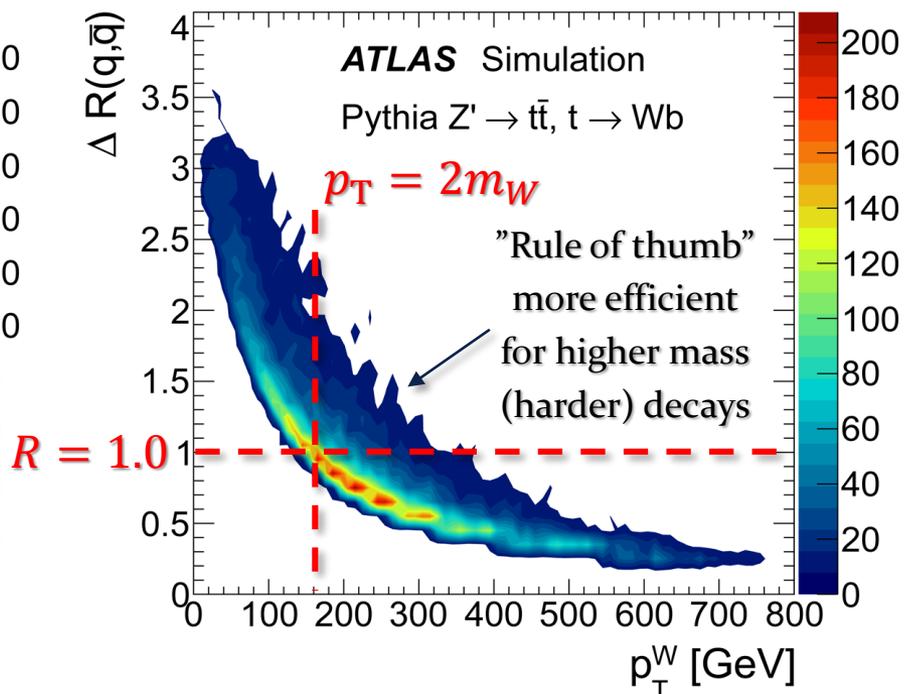
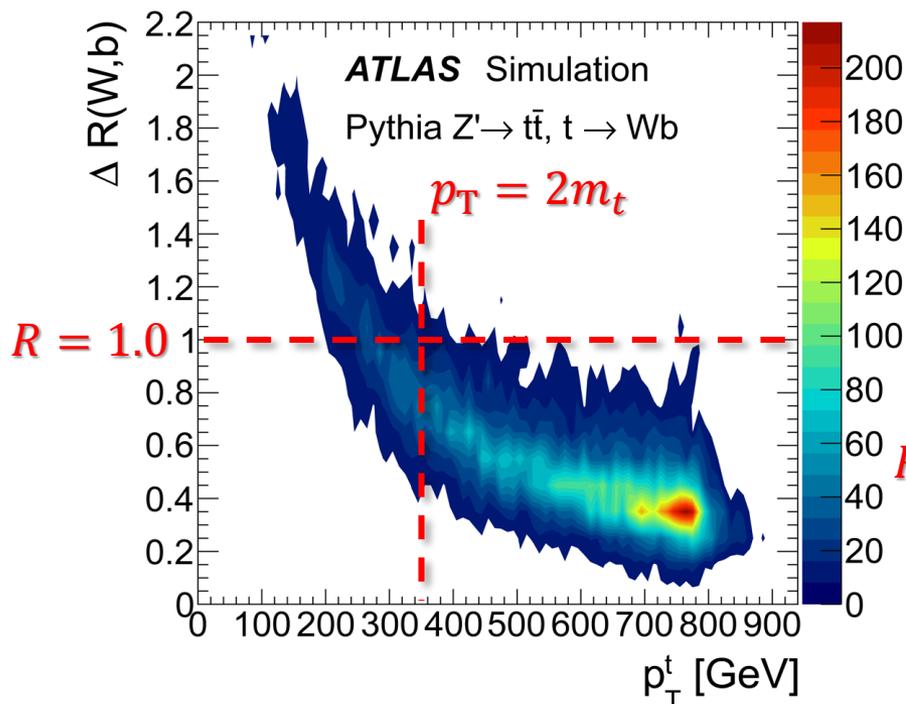
High p_T top quark hadronic decay

Simulated decay of new heavy particle $Z' \rightarrow t\bar{t}$

Topology is two sequential 2-prong decays with high p_T

Top quarks have high p_T – decay into $t \rightarrow Wb$ with high- p_T W boson

$$\Delta R(W, b) = \sqrt{\Delta y_{W \leftrightarrow b}^2 - \Delta \phi_{W \leftrightarrow b}^2} \quad \text{and} \quad \Delta R(q, \bar{q}) = \sqrt{\Delta y_{q \leftrightarrow \bar{q}}^2 - \Delta \phi_{q \leftrightarrow \bar{q}}^2}$$



Jet grooming – groomed jet

Means **any** substructure technique has been applied to a jet
Including combinations of several techniques

Fat jet

Reconstructed large- R jet

Often configured as a search tool

Popular are anti- k_t and C/A jets in filter-type applications

Typically $R \geq 1$

Sub-jet

Narrow jet inside a fat jet

Result of filter-type applications

Selected sub-jets are recombined to reconstruct observables of interest

Taggers

Characterization of a specific jet source

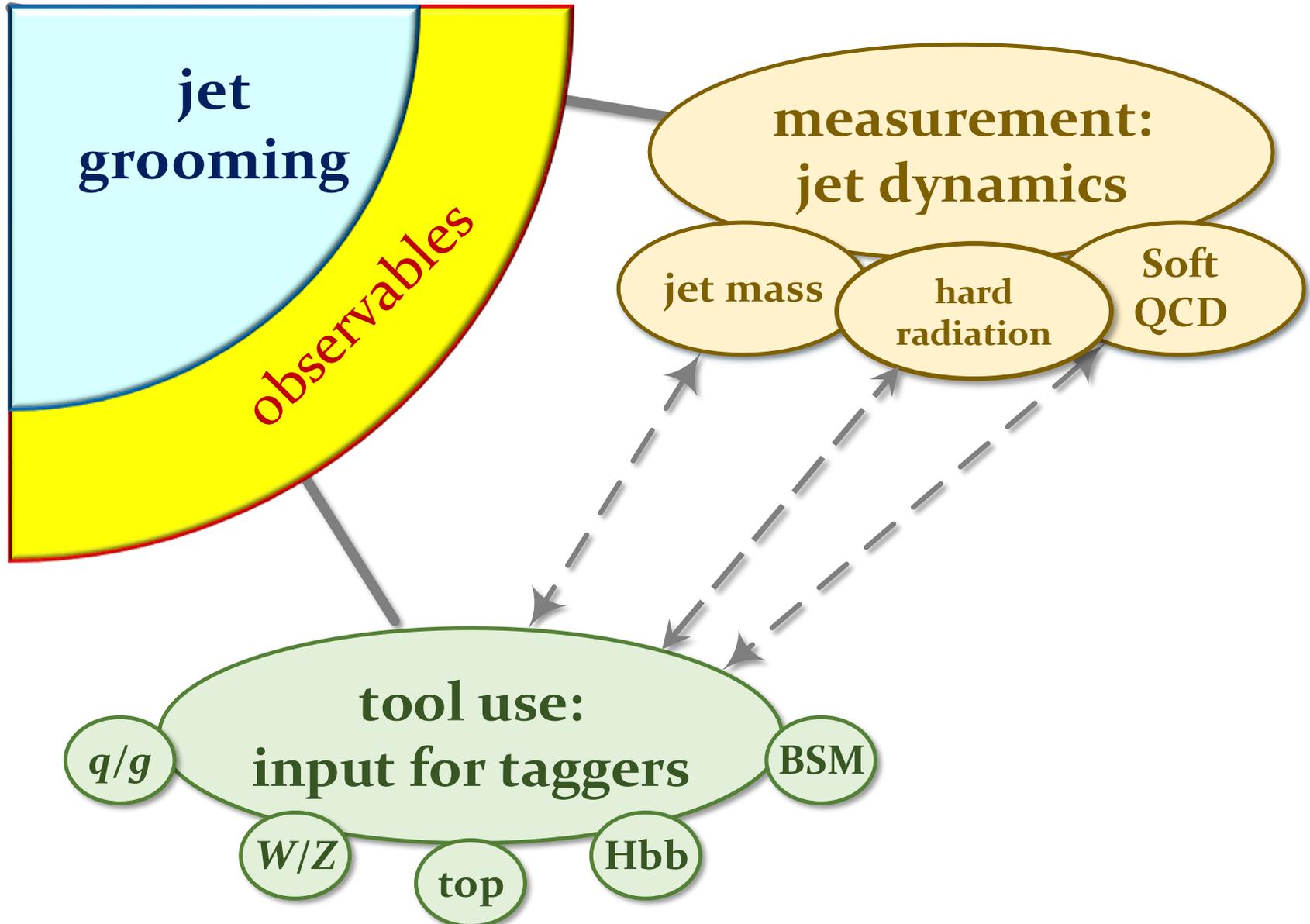
Implemented by (optimized) combinations of substructure techniques

E.g., HEPTOPTAGGER⁽¹⁾, W boson tagging⁽²⁾ etc.

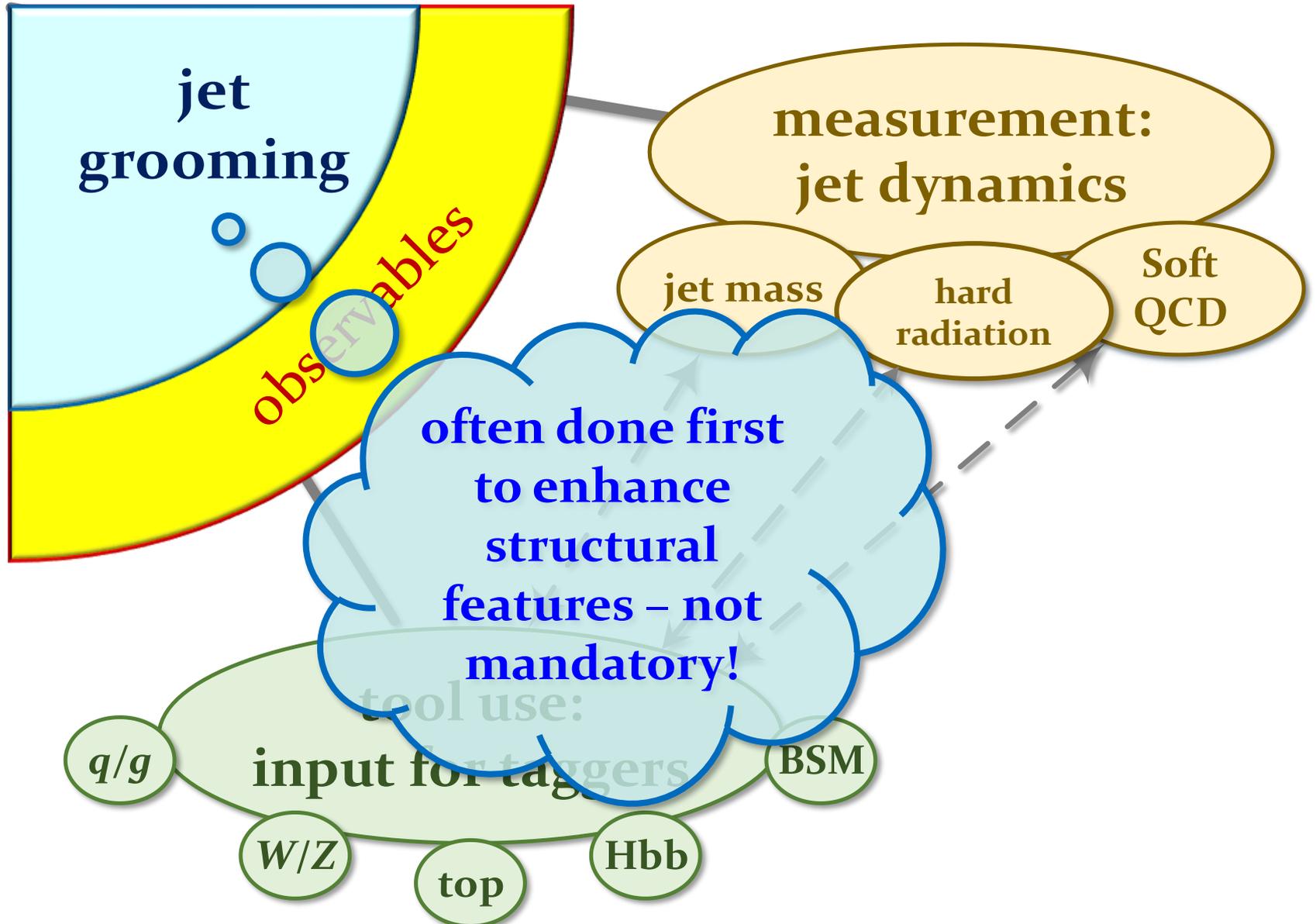
⁽¹⁾ T. Plehn, M.Spannowsky, M.Takeuchi, & D.Zerwas, JHEP 1010, 078 (2010)

⁽²⁾ Y.Cui, Z.Han, & M.D.Schwartz, Phys.Rev. D83 (2011) 074023

Jet Substructure in Physics Analysis



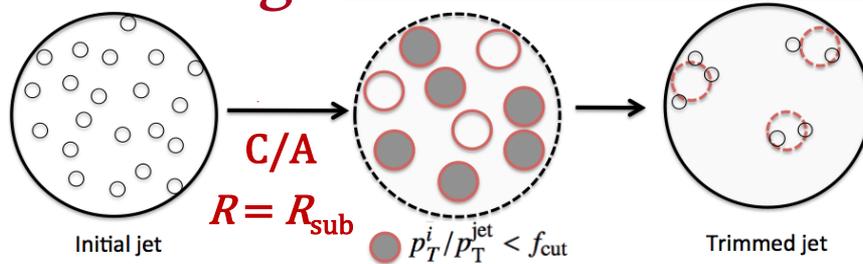
Jet Substructure in Physics Analysis



Jet Grooming Techniques (1)

Trimming

D.Krohn, J.Thaler, L.Wang, JHEP 02 (2010) 84

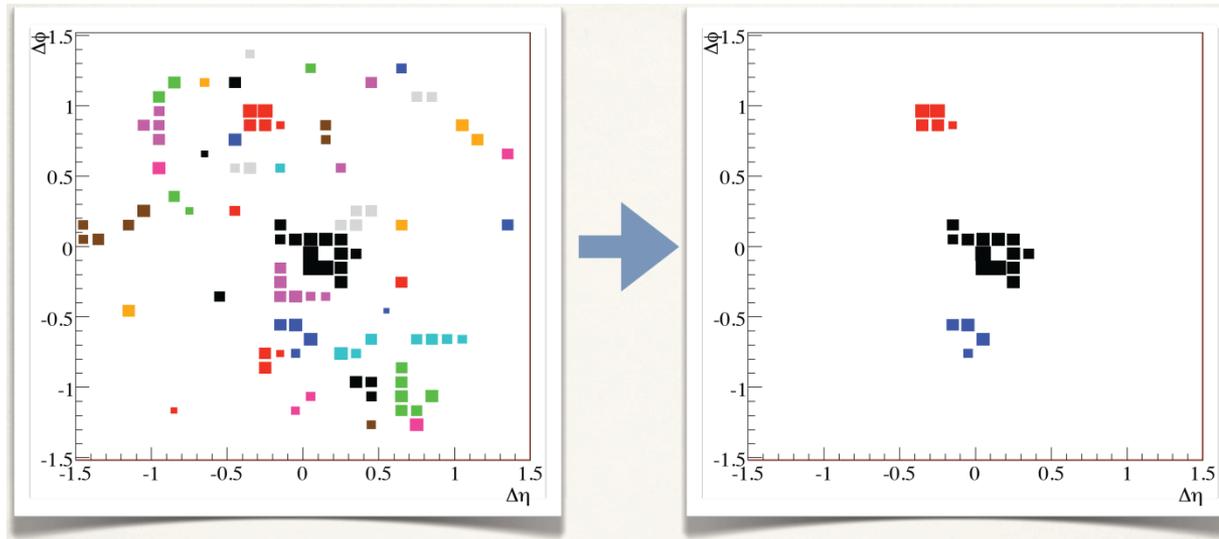


$$R_{\text{sub}} = \{0.2, 0.3\}$$

$$f_{\text{sub}} = \{0.01, 0.03, 0.05\}$$

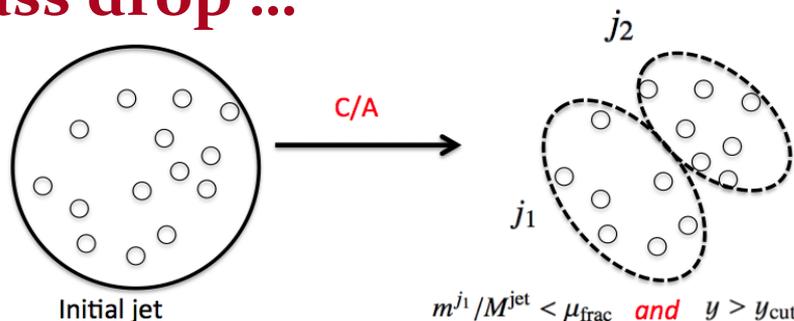
$$p_T^{\text{sub}} > f_{\text{sub}} \times p_T^{\text{jet}} \text{ (accepted sub-jets)}$$

Recombine surviving sub-jets to form groomed jet
 Works with any meaning full initial jet definition



Jet Grooming Techniques (2)

Mass drop ...



Butterworth, Davison, Rubin, Salam,
 Phys.Rev.Lett. 100 (2008) 242001

**REQUIRES MEANINGFUL
 RECURSIVE RECOMBINATION
 JET CLUSTER SEQUENCE**

Split jet into two sub-jets by undoing last recombination

Accept if masses of sub-jets considerably smaller than initial jet mass ...

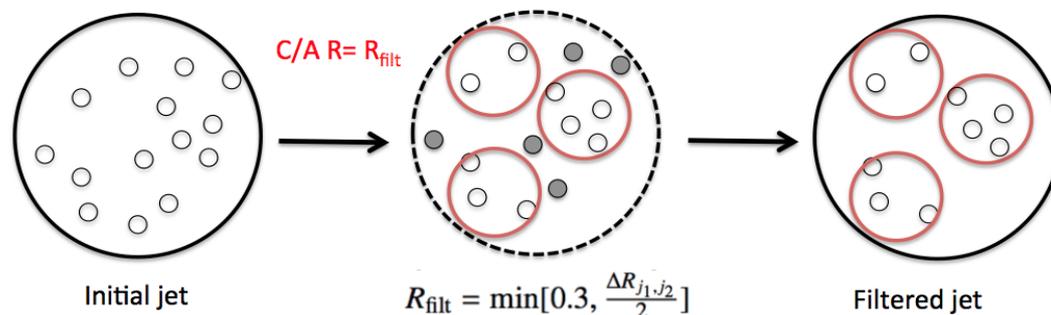
$$\min(m_1, m_2)/m < \mu_{\text{frac}} \text{ with } \mu_{\text{frac}} = \left\{ \frac{1}{5}, \frac{1}{3}, \frac{2}{3} \right\}$$

... and p_T sharing fairly symmetric

$$y = \min(p_{T,1}^2, p_{T,2}^2)/m^2 \times \Delta R_{12} < y_{\text{cut}} \text{ with } y_{\text{cut}} = 0.09$$

BDRS

... and filter



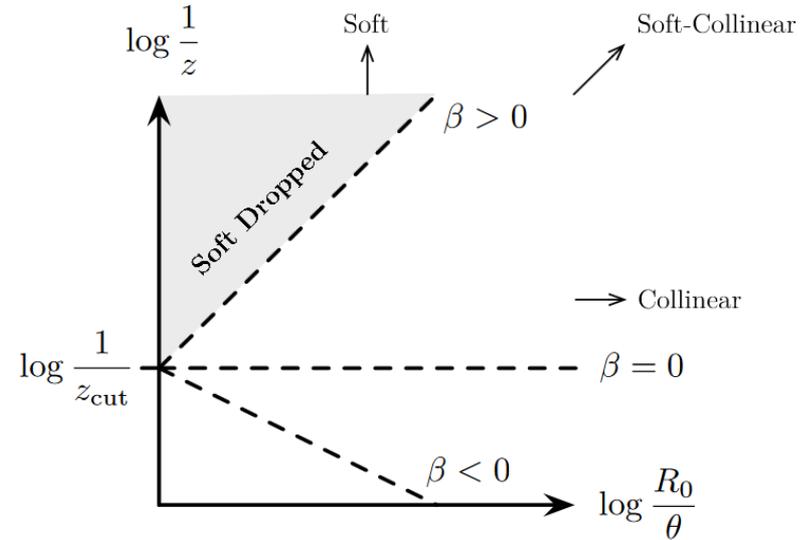
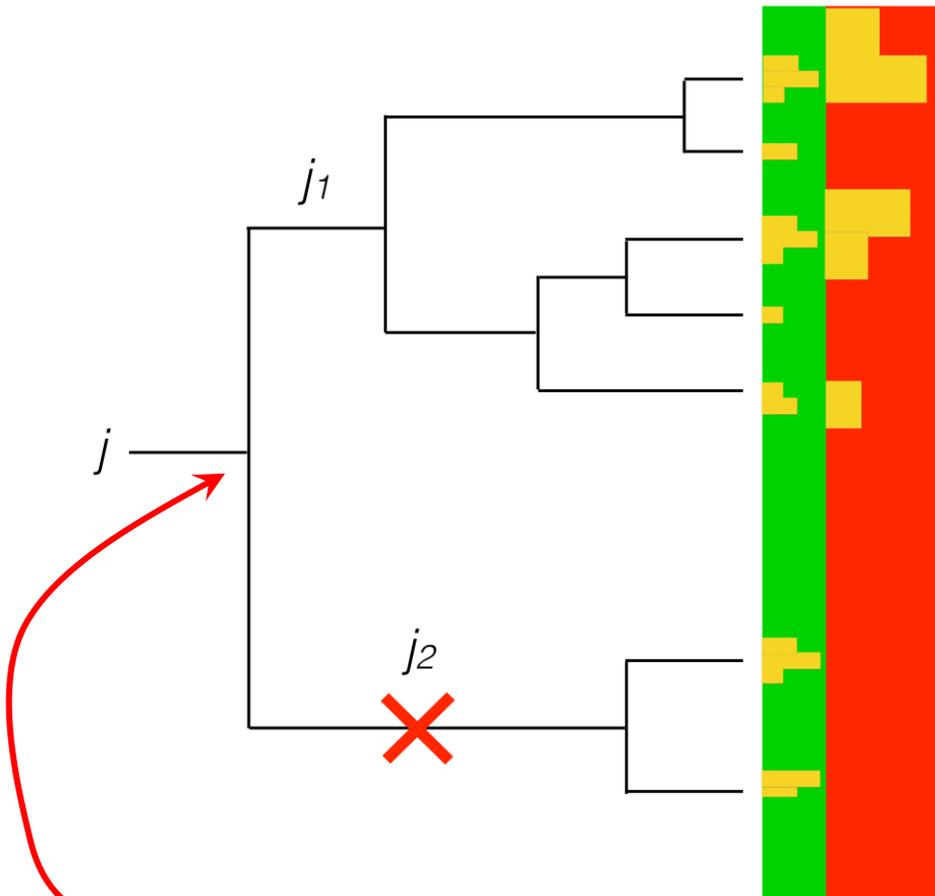
$$R_{\text{filt}} = \min(0.3, \Delta R_{12}/2)$$

Jet Grooming Techniques (3)

Soft Drop

Larkoski, Marzani, Soyez, Thaler, JHEP 1405 (2014) 146

De-clustering algorithm removes wide angle soft emissions inside any IRC safe clustered jet



Controlling soft drop phase space

- $\beta \rightarrow \infty$ yields ungroomed jet
- $\beta = 0$ corresponds to modified mass drop
- $\beta > 0$ removes soft radiation
- $\beta < 0$ removes soft-collinear radiation

fails soft drop condition $\min(p_{T,1}, p_{T,2}) / (p_{T,1} + p_{T,2}) > z_{\text{cut}} (\Delta R_{12} / R_0)^\beta$

Soft Drop

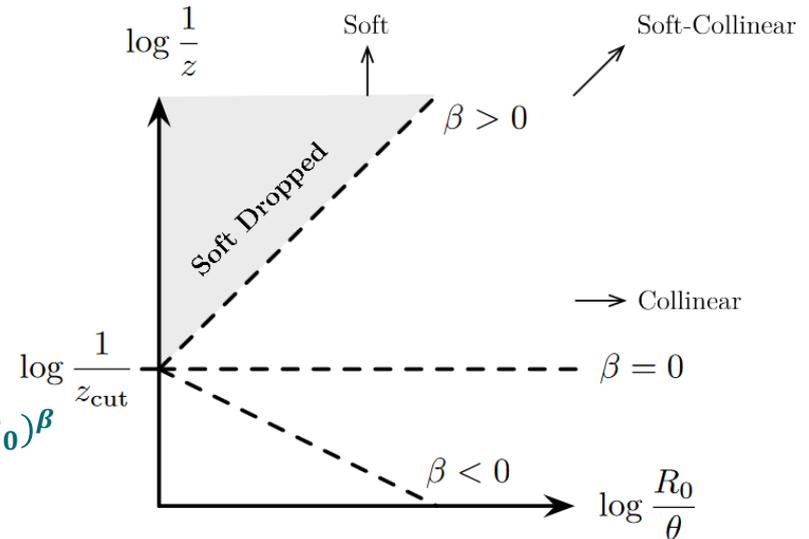
Larkoski, Marzani, Soyez, Thaler, JHEP 1405 (2014) 146

De-clustering algorithm removes wide angle soft emissions inside any IRC safe clustered jet

- Needs physically meaningful cluster sequence – recluster jet with C/A recursive recombination algorithm
- Undo the last stage of C/A clustering (largest angular distances!) → check if sub-jets j_1, j_2 pass the soft drop condition

$$\min(\mathbf{p}_{T,1}, \mathbf{p}_{T,2}) / (\mathbf{p}_{T,1} + \mathbf{p}_{T,2}) > z_{\text{cut}} (\Delta R_{12} / R_0)^\beta$$
 with parameters $z_{\text{cut}} \sim 0.1$, $\beta = R_0$ is nominal jet size

- If condition passed → original jet is soft dropped jet
- If softer sub-jet fails, harder sub-jet is redefined as the new jet → procedure is iterated until no more sub-jets are dropped
- If jet is a singleton and cannot be de-clustered anymore → **removed** from considerations (**tagging** mode) or considered as the **final soft-dropped jet** (**grooming** mode)



Controlling soft drop phase space

- $\beta \rightarrow \infty$ yields ungroomed jet
- $\beta = 0$ corresponds to modified mass drop
- $\beta > 0$ removes soft radiation
- $\beta < 0$ removes soft-collinear radiation

Relatively long list ...

Volatility, jet width/angularities, jet planarity, ...

Application depends on signal extraction efficiency and background rejection for a given search

Some are very sensitive to e.g. pile-up – performance can be largely degraded in the experiment

Some are hard to control experimentally – systematic uncertainties can be large

... recently most popular with experiments:

Single jet mass

Generalized angularities

Splitting scale

N -subjettiness

Energy-energy correlations

Soft drop mass

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Single jet mass

CAN BE USED WITH ANY
MEANINGFUL JET ALGORITHM

$$m^{\text{jet}} = \sqrt{(E^{\text{jet}})^2 - |\vec{p}^{\text{jet}}|^2}$$

Powerful and indicative observable in (resonant) searches

Useful in tag direct or subsequent decays of (boosted) SM particles

Rule of thumb for light quark/gluon jets (from NLO)

$$\sqrt{\langle (m^{\text{jet}})^2 \rangle_{\text{NLO}}} \approx (0.1 - 0.2) \cdot p^{\text{jet}} R$$

Ellis, Huston, Loch, Tönnemann,
Prog.Part.Nucl.Phys. 60 (2008) 484-551

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MEANINGFUL JET ALGORITHM

$$\lambda_{\beta}^{\kappa} = \sum_{i \in \text{jet}} z_i^{\kappa} \times (R_i/R_0)^{\beta}, \text{ with } z_i^{\kappa} = \left(p_{T,i} / \sum_{k \in \text{jet}} p_{T,k} \right)^{\kappa}$$

E.g., **jet width** ($\kappa = 1, \beta = 1$) useful for quark-gluon tagging

Larkoski, Thaler, Waalewijn, JHEP 11 (2014) 129

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**REQUIRES MEANINGFUL
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$$d_{ij} = \min(p_{T,i}^2, p_{T,j}^2) \times \Delta R_{ij}^2 \Rightarrow y_{\text{scale}}^{i \rightarrow j} = \sqrt{d_{ij}} / \Delta R_{ij}$$

Has expectation value for pronged decays, e.g. $d_{12} \approx (m/2)^2$

d_{12} ($t \rightarrow Wb$) and d_{23} ($W \rightarrow qq'$) used in hadronic top decay tagging

Butterworth, Cox, Forshaw, Phys.Rev. D65(2002) 096014

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$$\tau_N = \sum_k p_{T,k} \times \min\{\Delta R_{1,k}, \Delta R_{2,k}, \dots, \Delta R_{N,k}\} / (R_0 \sum_k p_{T,k})$$

Likelihood measure for jet to have 1 ... N subjets

Ratios like $\tau_{21} = \tau_2/\tau_1$ and $\tau_{32} = \tau_3/\tau_2$ highly indicative of 2- or 3-prong decay structure – used in hadronic top decay tagging

Relatively long list ...

Volatility, jet width/angularities, jet planarity, ...

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Energy-energy correlations

Employs N -point energy correlation functions between jet constituents

$$\text{ECF}(N, \beta) = \sum_{i_1 < i_2 < \dots < i_N \in \text{Jet}} \left(\prod_{a=1}^N p_{Ti_a} \right) \left(\prod_{b=1}^{N-1} \prod_{c=b+1}^N R_{i_b i_c} \right)^\beta$$

CAN BE USED WITH ANY
MEANINGFUL JET ALGORITHM

Relatively long list ...

Volatility, jet width/angularities, jet planarity, ...

Application depends on signal extraction efficiency and background

$$0\text{-point corr. } ECF(0, \beta) = 1$$

$$1\text{-point corr. } ECF(1, \beta) = \sum_{i \in \text{jet}} p_{Ti}$$

$$2\text{-point corr. } ECF(2, \beta) = \sum_{i < j \in \text{jet}} p_{Ti} p_{Tj} (R_{ij})^\beta$$

$$3\text{-point corr. } ECF(3, \beta) = \sum_{i < j < k \in \text{jet}} p_{Ti} p_{Tj} p_{Tk} (R_{ij} R_{ik} R_{jk})^\beta$$

$$4\text{-point corr. } ECF(4, \beta) = \sum_{i < j < k < l \in \text{jet}} p_{Ti} p_{Tj} p_{Tk} p_{Tl} (R_{ij} R_{ik} R_{il} R_{jk} R_{jl} R_{kl})^\beta$$

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Application depends on signal extraction efficiency and background

Double-ratios indicate jet source

Generally $C_N^{(\beta)} = (ECF(N + 1, \beta)ECF(N - 1, \beta))/ECF(N, \beta)^2$

$C_1^{(\beta)}$ is useful for quark/gluon separation for small $\beta \simeq 0.2$ – exploits different soft radiation patterns (uses 2-point correlations)

$C_2^{(\beta)}$ helps with boosted $W/Z/H$ identification with $\beta \simeq 0.5$ better for high mass and $\beta \simeq 2$ better for lower mass resonances – at a fixed p_T (uses 3-point correlations)

$C_3^{(\beta)}$ distinguishes QCD jets from boosted top quarks best for $\beta \approx 1 - 2$
(uses 4-point correlations)

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Relatively long list ...

Volatility, jet width/angularities, jet planarity, ...

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Eur.Phys.J. C76 (2016), no.3, 154

Boosted boson tagging in ATLAS

$$C_2^{(\beta)} = (\text{ECF}(3, \beta) \text{ECF}(1, \beta)) / (\text{ECF}(3, \beta))^2$$

and optimized for 2-prong decays

$$D_2^{(\beta)} = (\text{ECF}(3, \beta) \text{ECF}(1, \beta)) / (\text{ECF}(3, \beta))^3$$

with $\beta = 1, 2$ studied

Larkoski, Mout, Neill,
JHEP 1409 (2014)

Energy-energy correlations

Employs N -point energy correlation functions between jet constituents

$$\text{ECF}(N, \beta) = \sum_{i_1 < i_2 < \dots < i_N \in \text{Jet}} (\prod_{a=1}^N p_{Ti_a}) (\prod_{b=1}^{N-1} \prod_{c=b+1}^N R_{i_b i_c})^\beta$$

Larkoski, Salam, Thaler, JHEP 1306 (2013) 108

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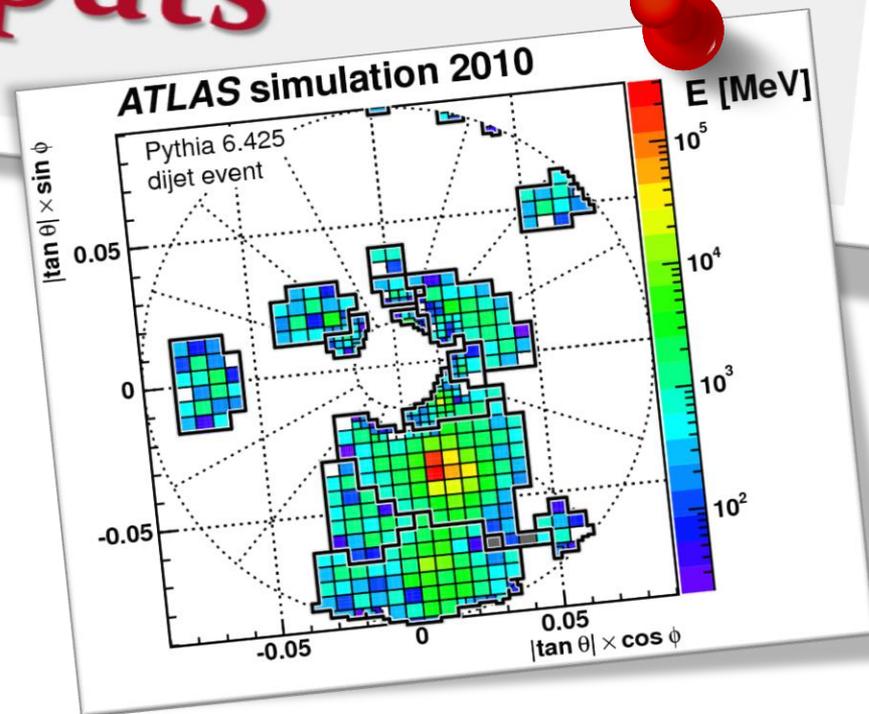
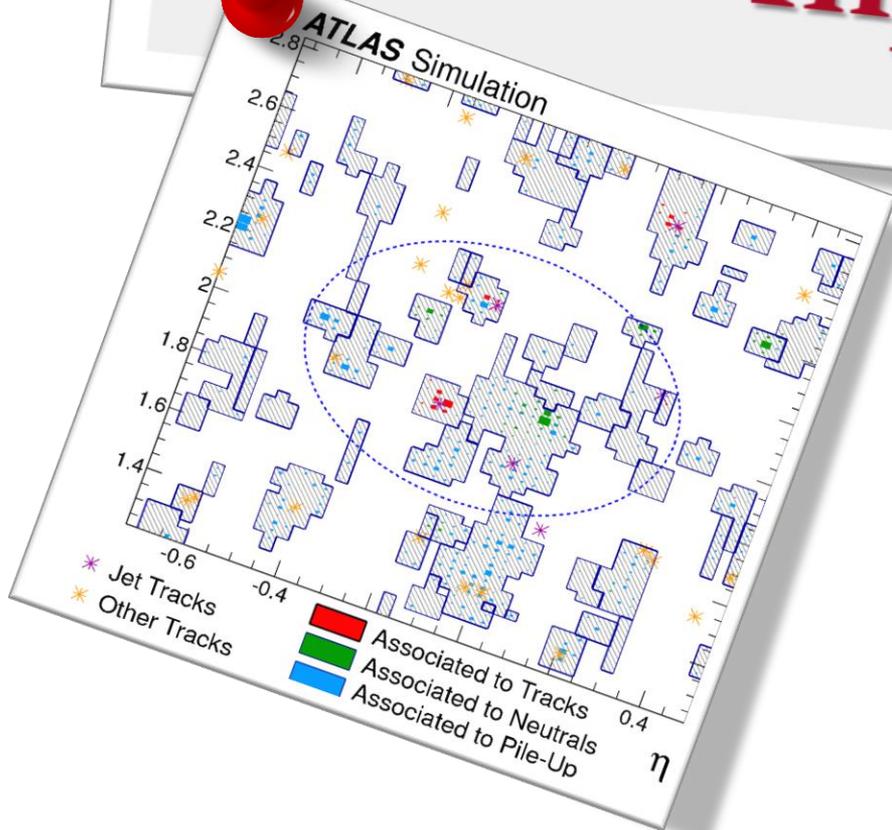
N -subjettiness

Energy-energy correlations

Soft drop mass

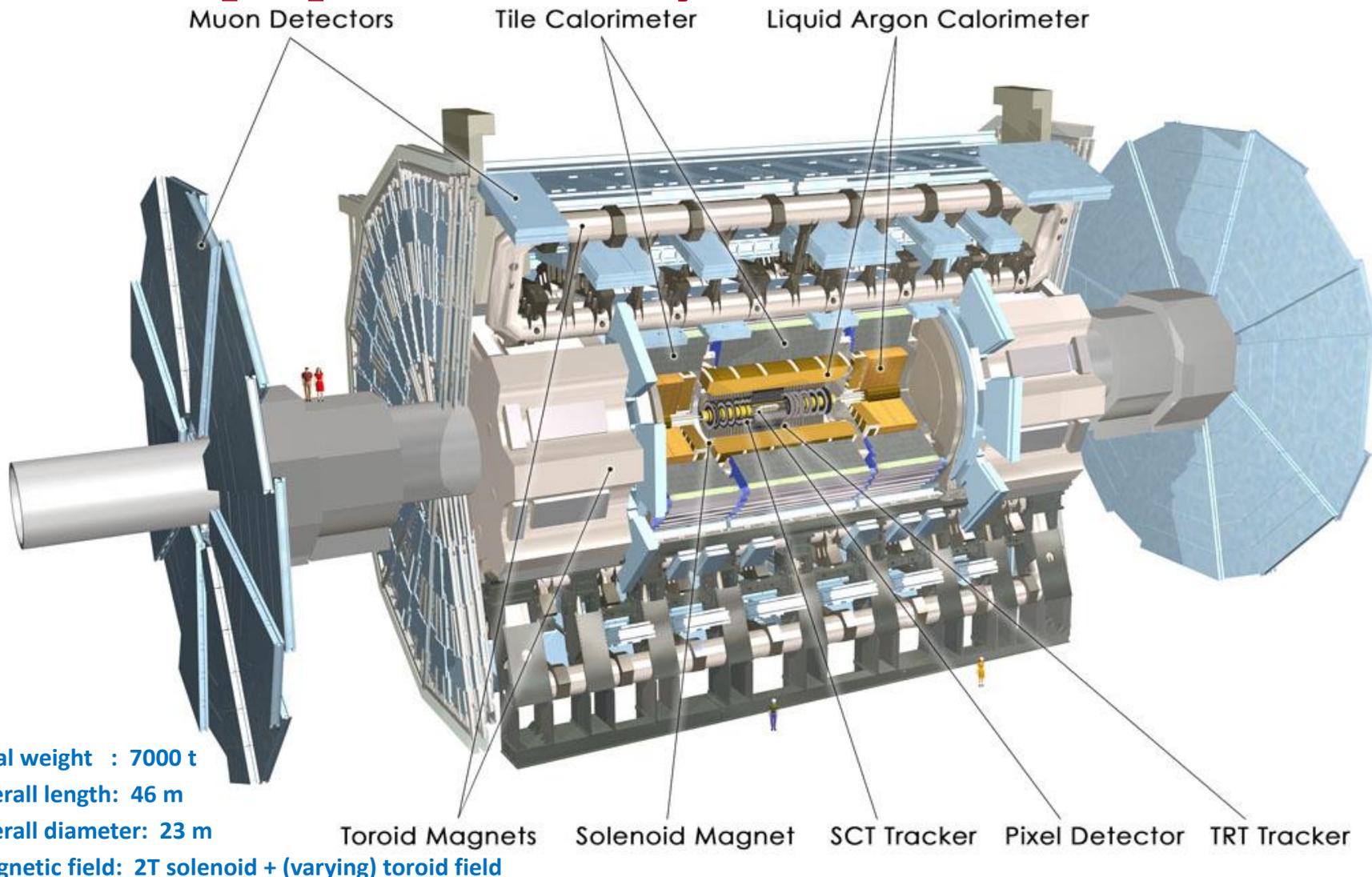
Recalculate mass of soft dropped jet – e.g. gives experimental access to evaluation of calculated internal dynamics of QCD jets

Jet Reconstruction Inputs



ATLAS at the LHC

A multi-purpose detector system



Calorimeters (Calo)

Provides principal signals for jet kinematics and substructure measurement

Full coverage within $|\eta| < 4.9$ with depth $\gtrsim 10 \lambda_{\text{int}}$

Highly segmented for energy flow measurements ($\sim 188,000$ channels)

Granularity in $\Delta\eta \times \Delta\varphi = 0.025 \times \pi/128$ (central EM, $|\eta| \lesssim 2.5$)

Up to seven depth layers (*samplings*)

Inner detector (ID)

Provides charged particle tracks and vertices

Coverage $|\eta| < 2.5, p_{\text{T}} > 500$ MeV

Pile-up mitigation and overall jet calibration refinement by associating vertex with jets and/or sub-jets

Track-assisted measurement of flow structures – improved angular component in substructure and mass reconstruction

Particle flow

Replace charged response in calorimeter with kinematics from well-measured tracks

Topological clustering

Collects signals from individual or close-by particles into 3-dimensional *energy blobs*

Connect cells by following spatial signal significance patterns with seed and growth control

Massless pseudo-particle representation

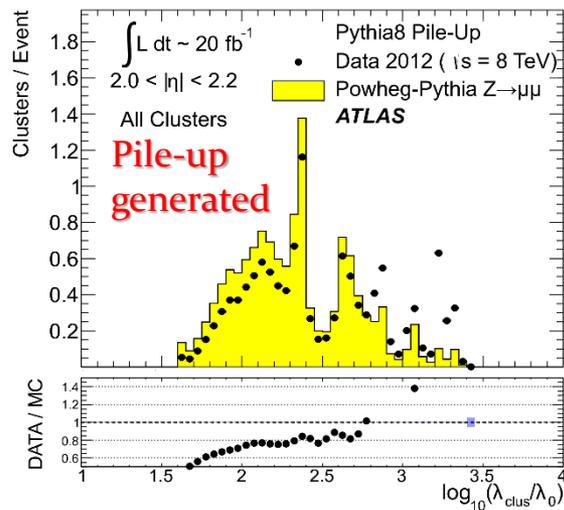
Recombines (weighted) cell energies to four-momentum $P_{clus} = (E_{clus}, \eta_{clus}, \varphi_{clus}, m_{clus} = 0)$

Cluster direction with respect to nominal collision vertex

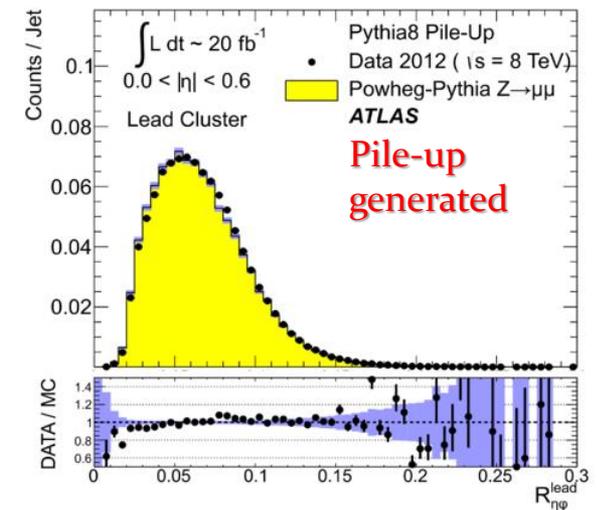
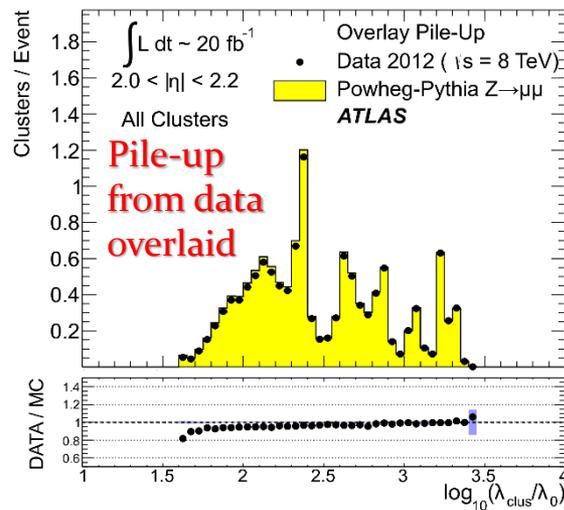
Cluster signal on basic (EM) or locally calibrated (LCW) scale

$$P_{clus} = P_{clus}^{EM} = (E_{clus}^{EM}, \eta_{clus}^{EM}, \varphi_{clus}^{EM}, 0) \text{ or } P_{clus} = P_{clus}^{LCW} = (E_{clus}^{LCW}, \eta_{clus}^{LCW}, \varphi_{clus}^{LCW}, 0)$$

LCW uses cluster shapes – preferred for substructure but some sensitivity to pile-up

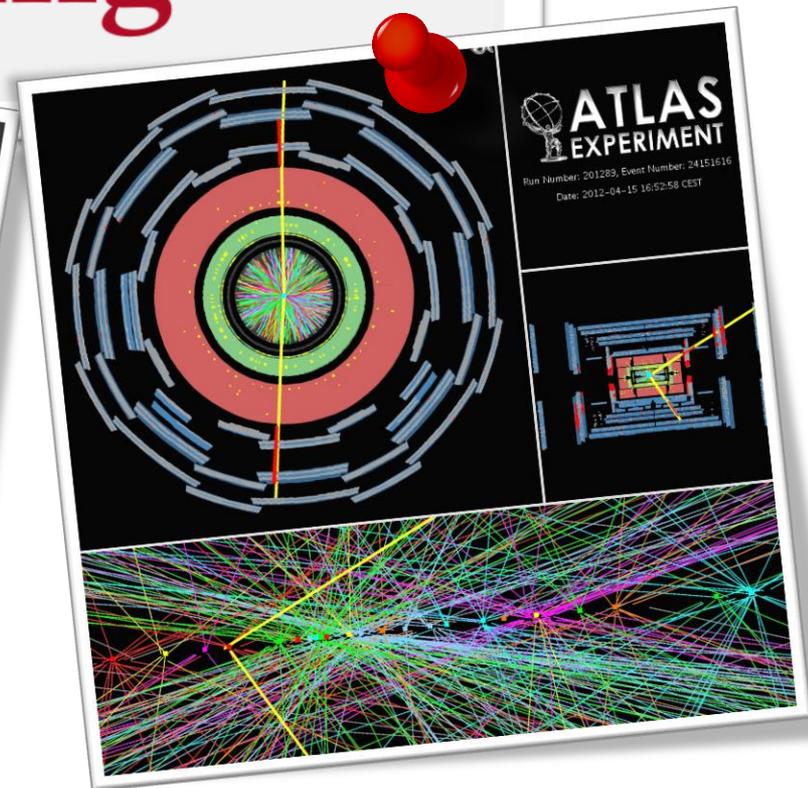
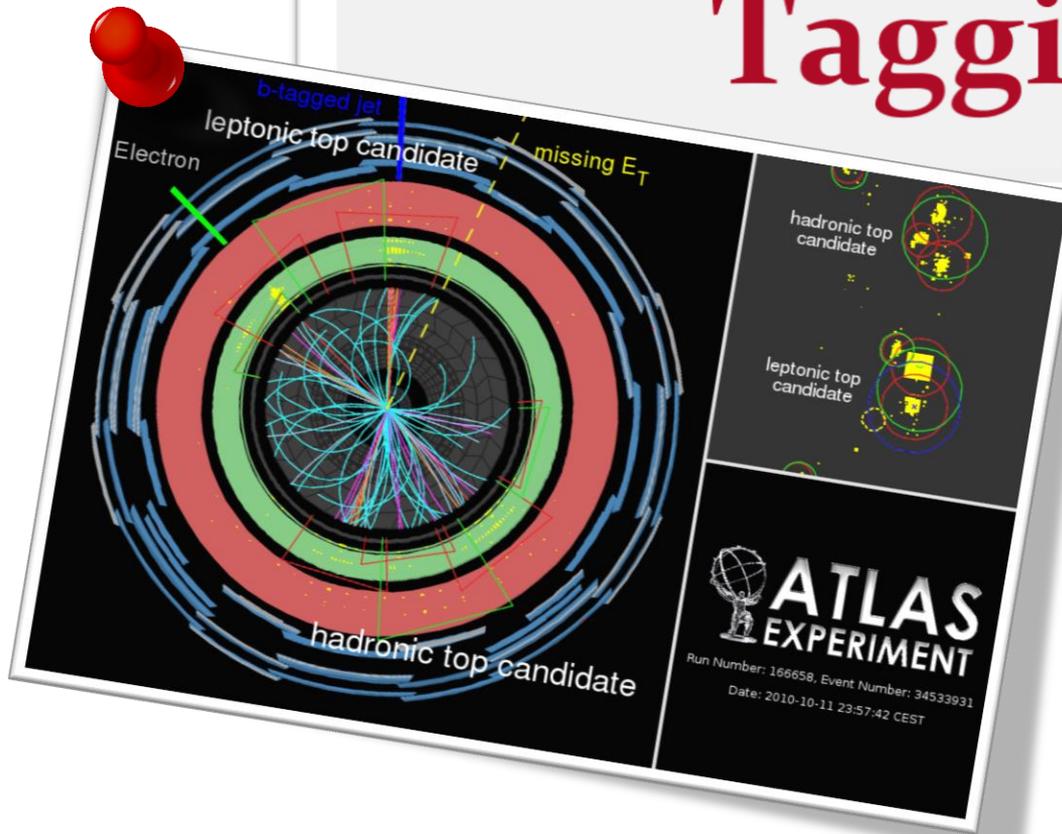


Topo-cluster depth location measure $\log_{10}(\lambda_{clus}/1 \text{ mm})$



Leading topo-cluster size R_{η}^{lead}

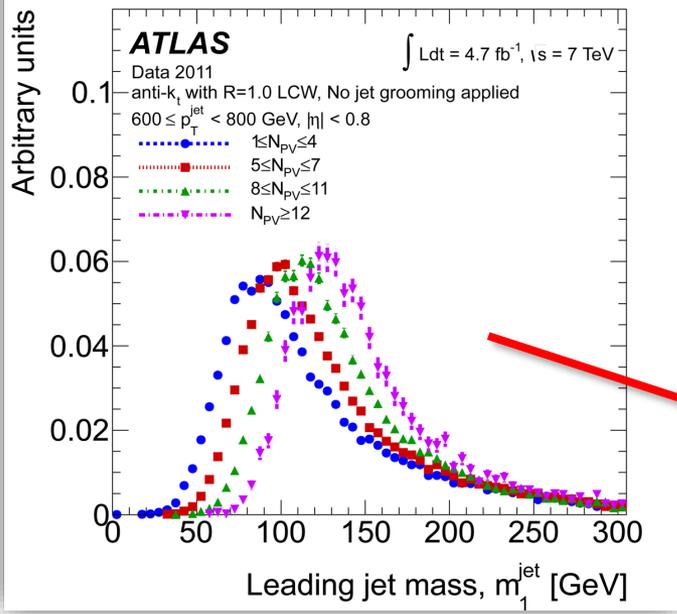
Substructure Performance & Tagging



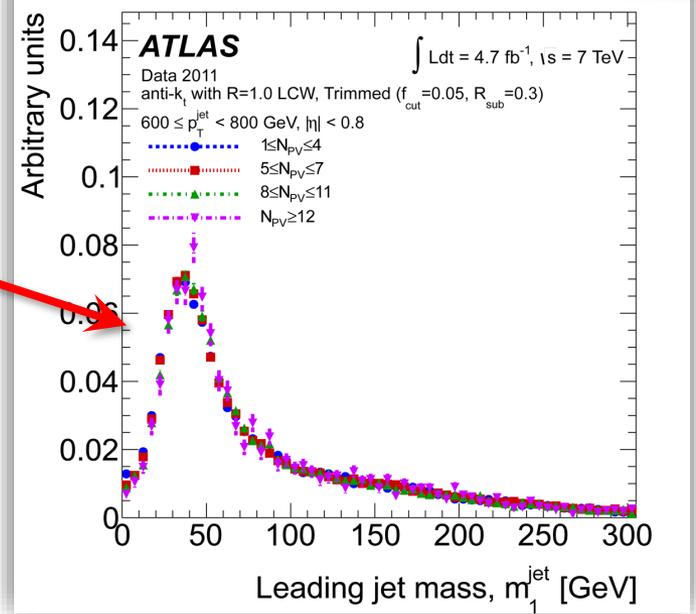
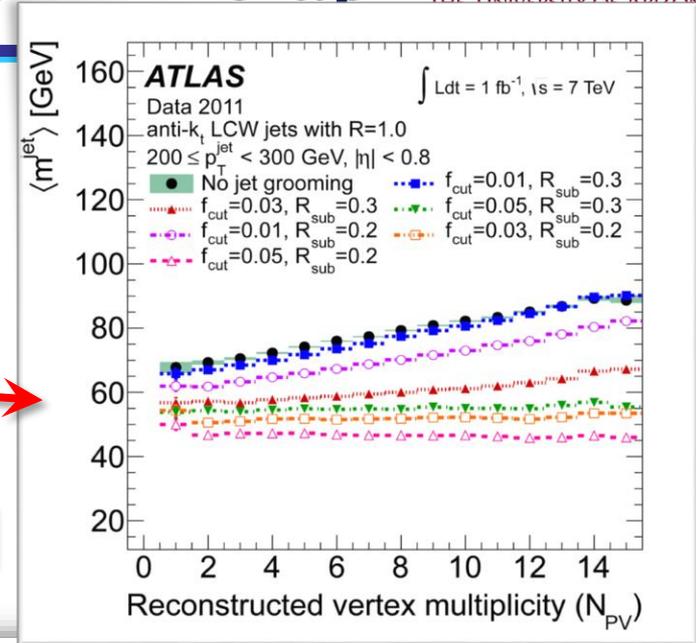
Single Jet Mass Measurement in Pile-up

Apply grooming techniques to mitigate pile-up effects

Pile-up affects average reconstructed mass for jets



JHEP 1309 (2013) 76



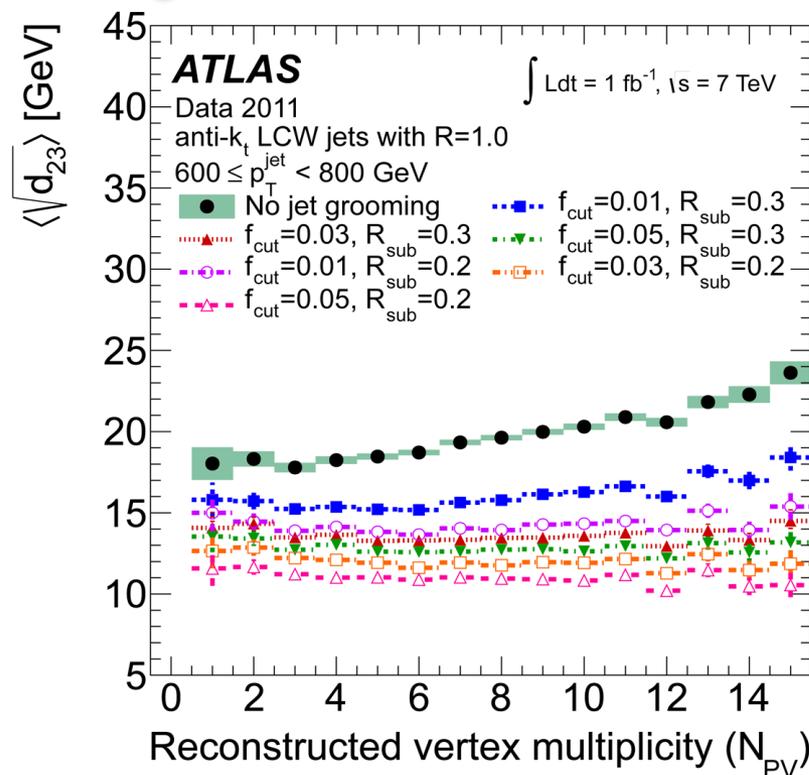
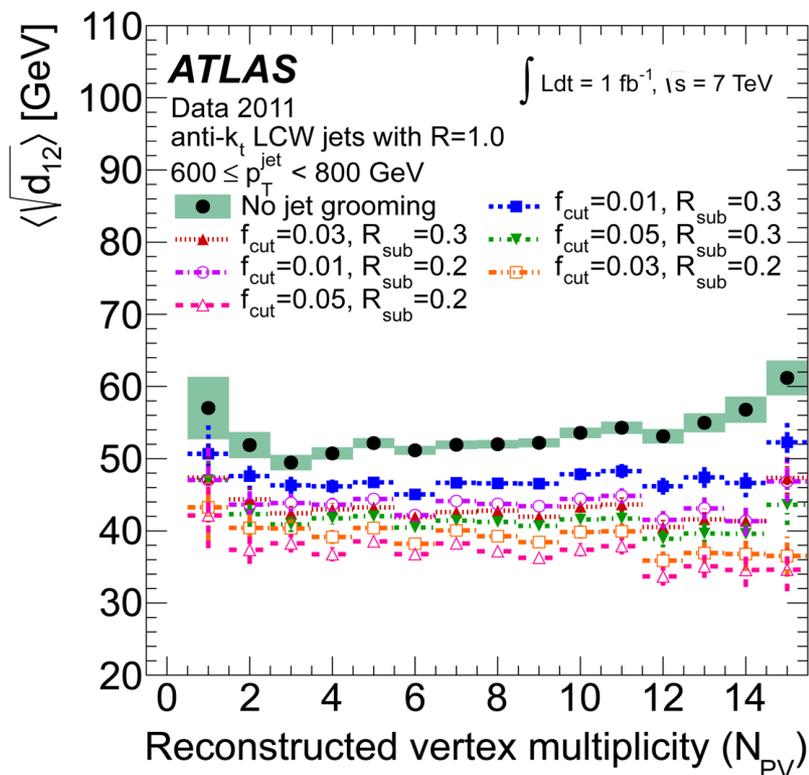
Anti- k_t , $R = 1.0$
 $600 \leq p_T^{\text{jet}} < 800 \text{ GeV}$

Substructure in Pile-up (1)

Pile-up mitigation by trimming

Splitting scales

Anti- k_t , $R = 1.0$
 $600 \leq p_T^{\text{jet}} < 800 \text{ GeV}$



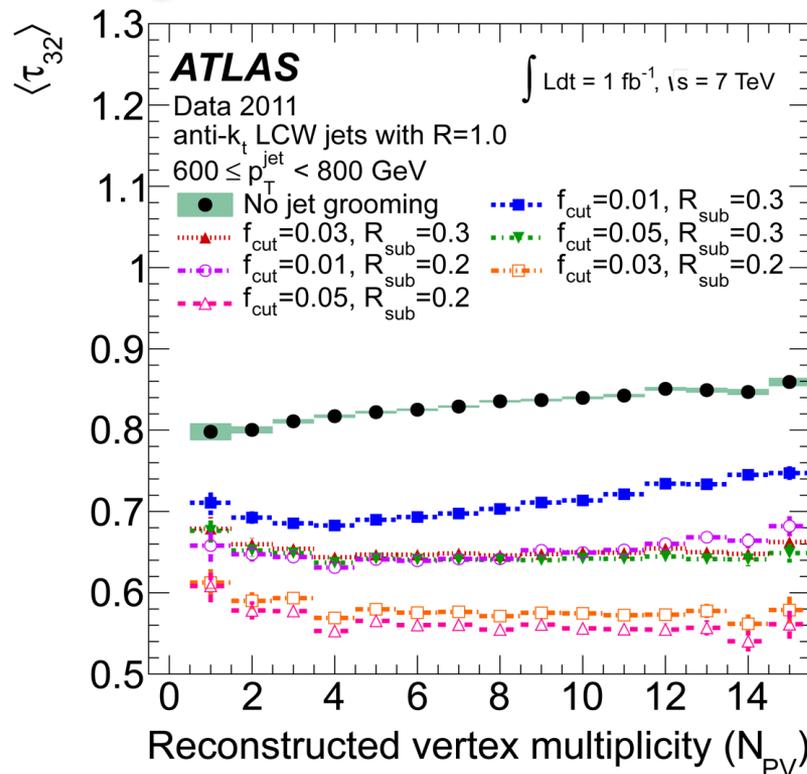
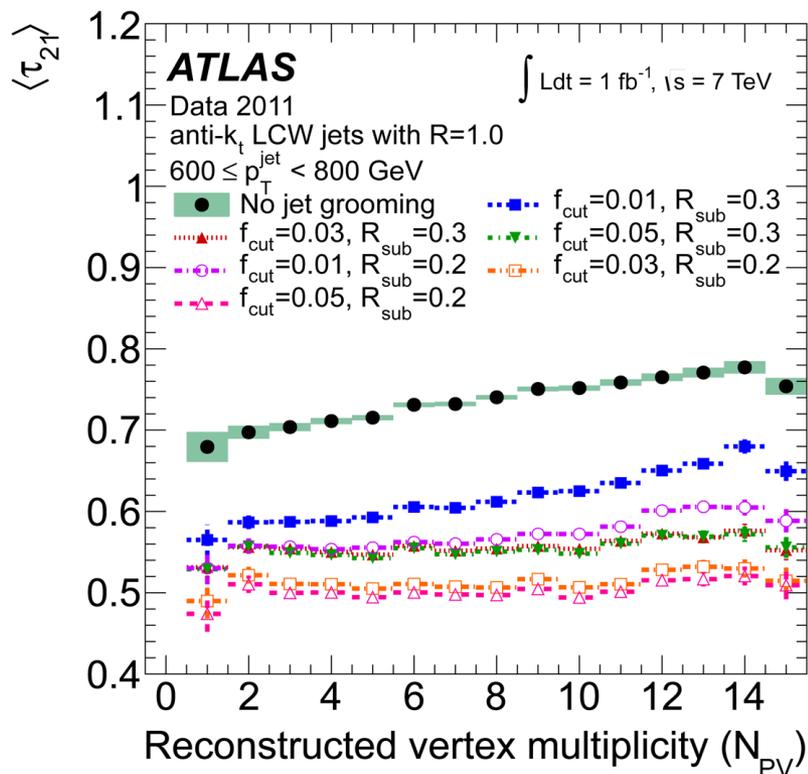
JHEP 1309 (2013) 76

Substructure in Pile-up (2)

Pile-up mitigation by trimming

N -subjettiness

Anti- k_t , $R = 1.0$
 $600 \leq p_T^{\text{jet}} < 800 \text{ GeV}$



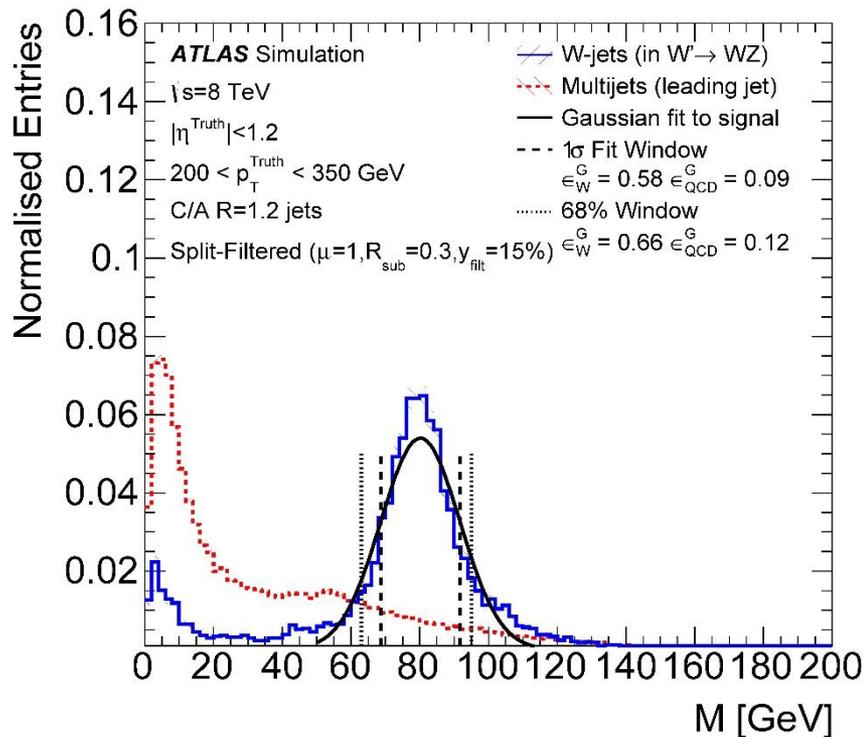
JHEP 1309 (2013) 76

Grooming techniques/configurations need optimization

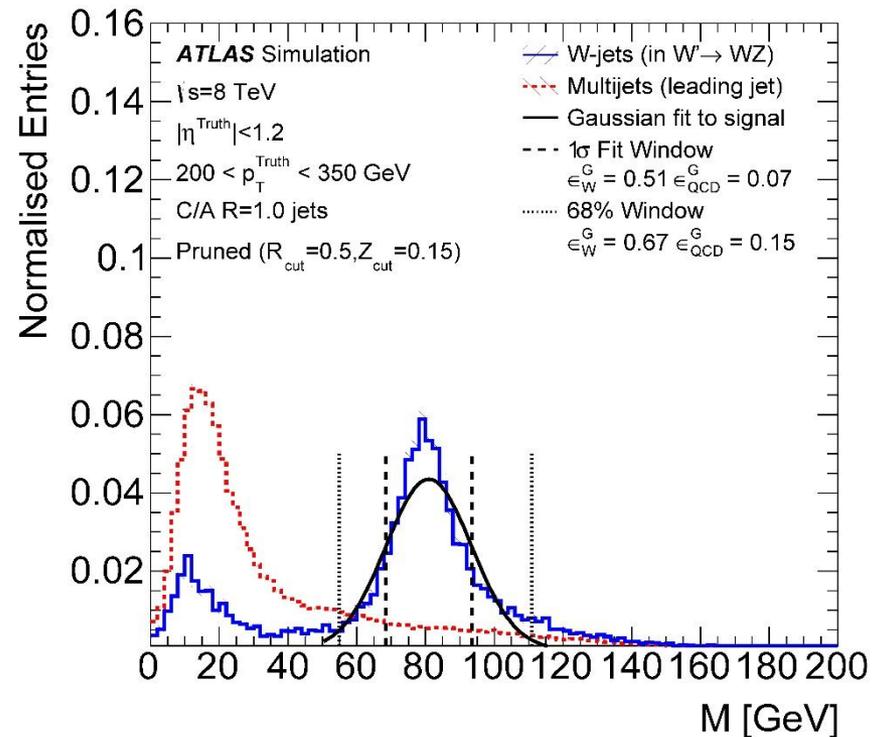
Adapt to search goal – signal efficiency and background rejection

Example: grooming in simulated $W' \rightarrow WZ$ with $W \rightarrow qq$

EPJC 76 (2016) No 3. 154



split-filtered C/A R = 1.2 jets



pruned C/A R = 1.0 jets

SM $t\bar{t}$ final state with one full hadronic top decay

Efficient event selection from semi-leptonic decay

Lepton and missing transverse momentum triggers

Offline event selection based on leptons, missing transverse momentum and number of (small radius) jets (e.g., $N_{\text{jet}} \geq 4$)

Three-prong decay $t \rightarrow Wb \rightarrow qq'b$

Measure performance of substructure observables in top decay dynamics – useful for searches for new particles decaying into top quarks

Enhanced sample for two-prong decay $W \rightarrow qq'$

Most efficient source for hadronic W decays

Allows evaluation of W or top tagging configurations

Momentum and topology selection of large radius jet ($R_{\text{jet}} = 1$)

E.g., $p_{\text{T}}^{\text{jet}} > 200$ GeV, large angular distance $\Delta R(W, b) > R_{\text{jet}}$, two-prong decay indication (τ_{21} small, $\sqrt{d_{12}} \approx m_W, \dots$) \rightarrow contains (only) hadronic W decay

E.g., $p_{\text{T}}^{\text{jet}} > 350$ GeV, small angular distance $\Delta R(W, b) < R_{\text{jet}}$, three-prong decay indication ($\sqrt{d_{12}} \approx m_{\text{top}}/2, \sqrt{d_{23}} \approx m_W/2, \dots$) \rightarrow contains hadronic top decay

W can be reconstructed in large radius jets in enhanced top sample!

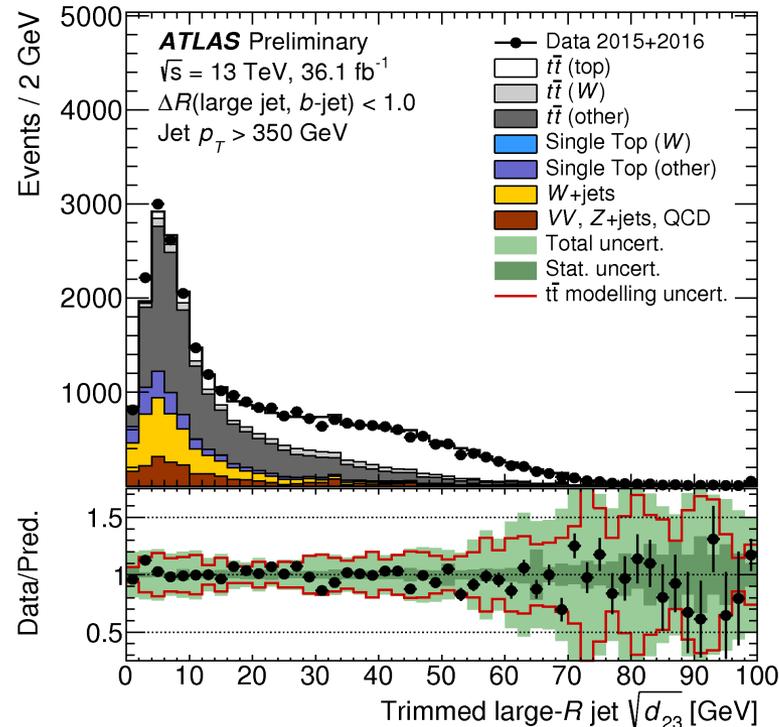
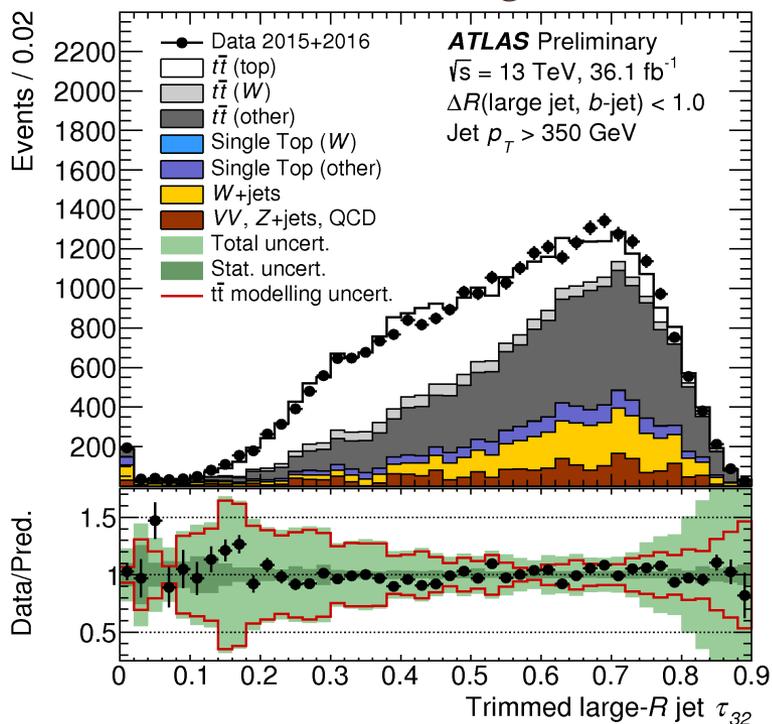
N-subjettiness & Splitting Scale

Modeling in $t\bar{t}$ decays

Complex measures in a physics analysis application

Understanding non-trivial distributions in data \rightarrow input to BDTs, (deep) neural networks (DNN), machine learning (ML), ...

$\text{Anti-}k_t, R = 1.0$
 $p_T^{\text{jet}} > 350 \text{ GeV}$
(top enhanced)

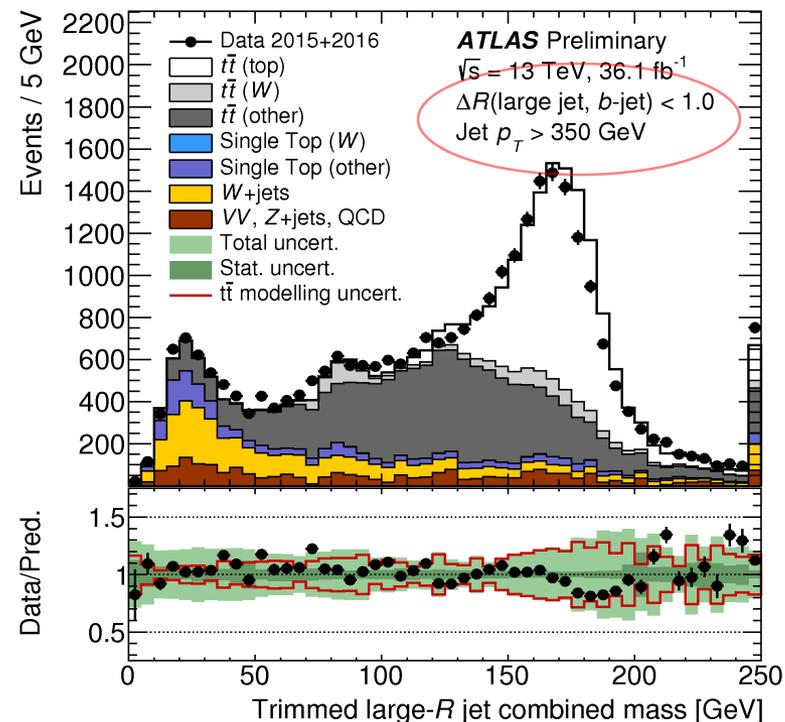
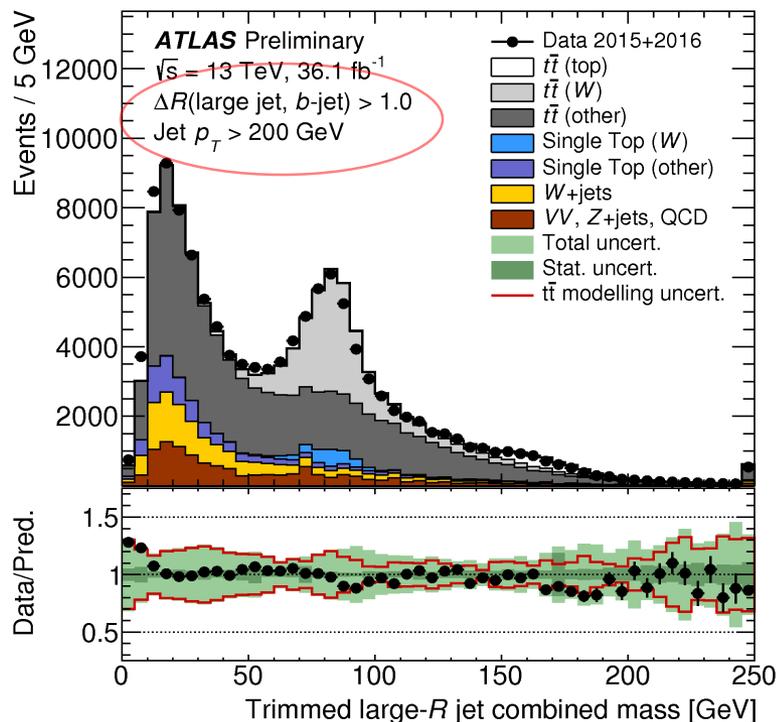


ATLAS-CONF-2017-064 (2018)

Decay Tagging in $t\bar{t}$ Final State

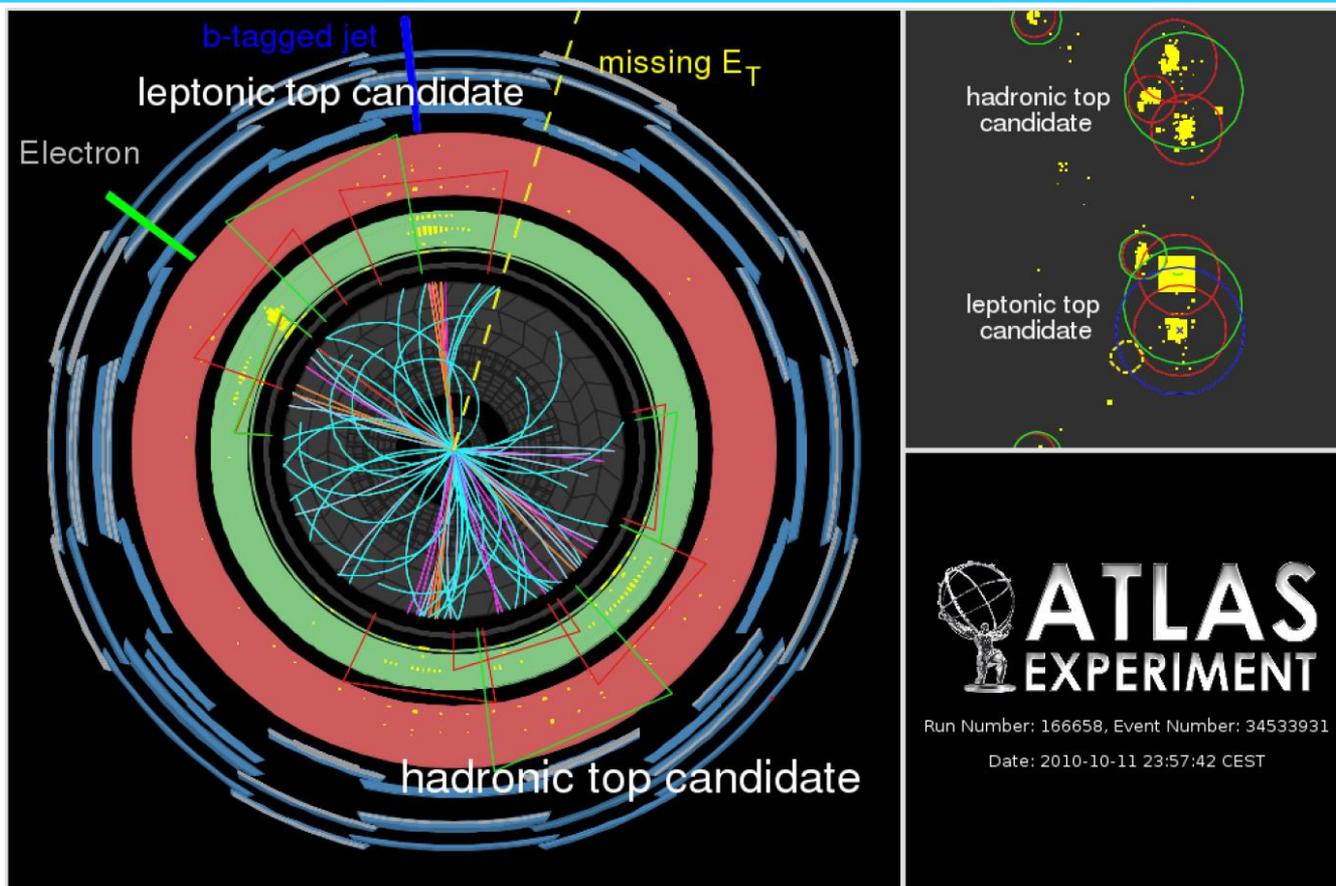
Anti- k_t , $R = 1.0$
 $p_T^{\text{jet}} > 200 \text{ GeV}$
 (W enhanced)

Anti- k_t , $R = 1.0$
 $p_T^{\text{jet}} > 350 \text{ GeV}$
 (**top** enhanced)



ATLAS-CONF-2017-064 (2018)

Boosted Top Decay in ATLAS



ATLAS
 EXPERIMENT

Run Number: 166658, Event Number: 34533931

Date: 2010-10-11 23:57:42 CEST

Leptonic top	E_T^{miss} : $E_T = 36$ GeV, $\phi = -1.5$ electron: $p_T = 145$ GeV, $\eta = 1.1$, $\phi = 2.5$ jet: index = 1, $E_T = 194$ GeV, $\eta = 1.2$, $\phi = 1.7$, $m_j = 17$ GeV
Hadronic top ($R = 0.4$ clustering)	jet 2, $E_T = 155$ GeV, $\eta = 1.1$, $\phi = -0.7$ rad, $m_j = 22.7$ GeV + jet 3, $E_T = 113$ GeV, $\eta = 1.3$, $\phi = -1.7$ rad, $m_j = 14$ GeV + jet 4, $E_T = 54$ GeV, $\eta = 0.6$, $\phi = -1.7$ rad, $m_j = 8$ GeV
Hadronic top ($R = 1.0$ clustering)	jet 1, $E_T = 356$ GeV, $\eta = 1.3$, $\phi = -1.1$ rad, $m_j = 197$ GeV $\sqrt{d_{12}} = 110$, $\sqrt{d_{23}} = 40$

Improving jet substructure

Mass measurement improvement

$$\text{Track-assisted mass } m_{\text{TA}} = p_{\text{T}}^{\text{calo}} / p_{\text{T}}^{\text{track}} \times m_{\text{track}}$$

Combining m_{calo} and m_{TA}

Track-Calo-Cluster (TCC)

Using calorimeter response and track direction – efficient at high jet p_{T}

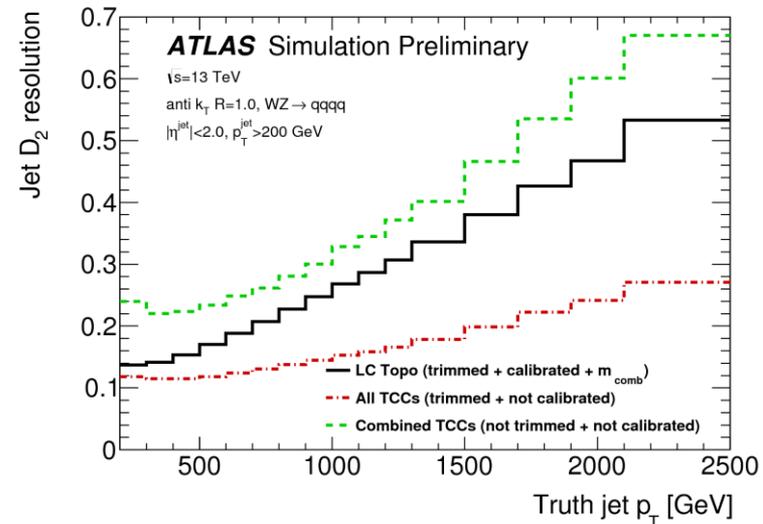
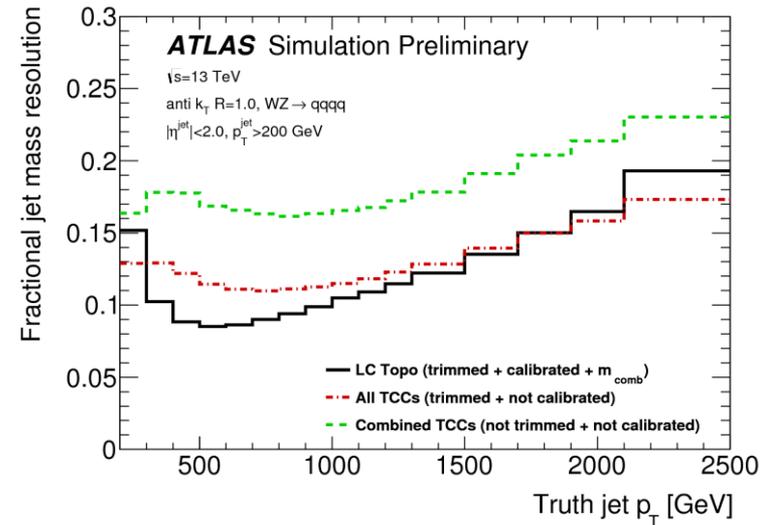
Track-assisted & TCC substructure

E.g., energy-energy correlation D_2

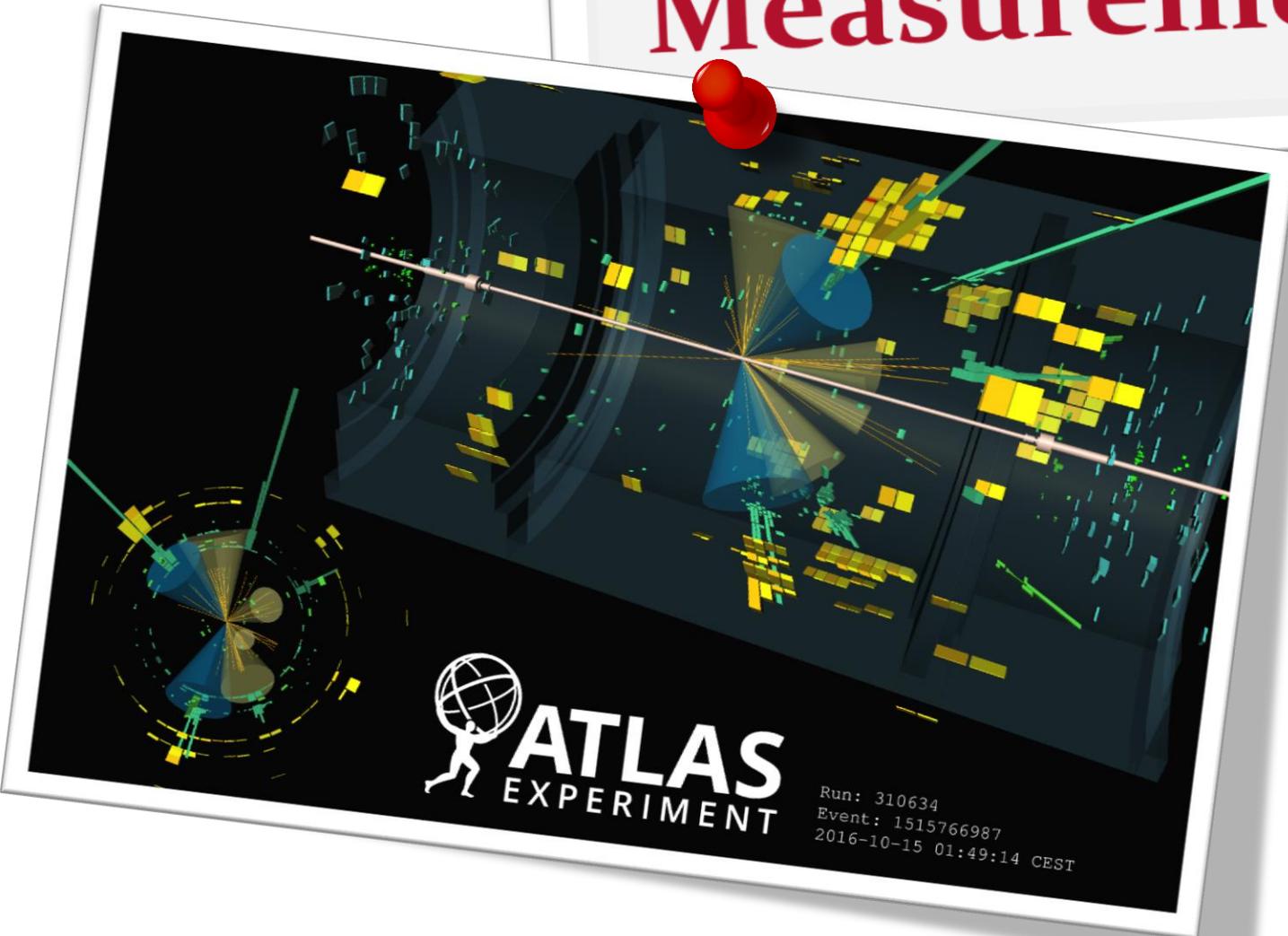
Reclustering

Cluster fat jets from narrow jets

Fully calibrated and pile-up corrected



Measurements



Predictions for $\rho = m^{\text{soft-drop}} / p_T^{\text{ungroomed}}$

Available (recently) NLO (order α_s) + NLL (resummation), LO + NNLL

Measurements

Evaluate calculations of ρ using data

$\log_{10}(\rho^2)$ distributions reflect internal jet dynamics

Algorithm and selections

Anti- k_t , $R = 0.8$ jets

Di-jet topology with $p_{T,1} > 600$ GeV and $p_{T,1}/p_{T,2} < 1.5$

Soft drop configuration $z_{\text{cut}} = 0.1$, $\beta = \{0,1,2\}$

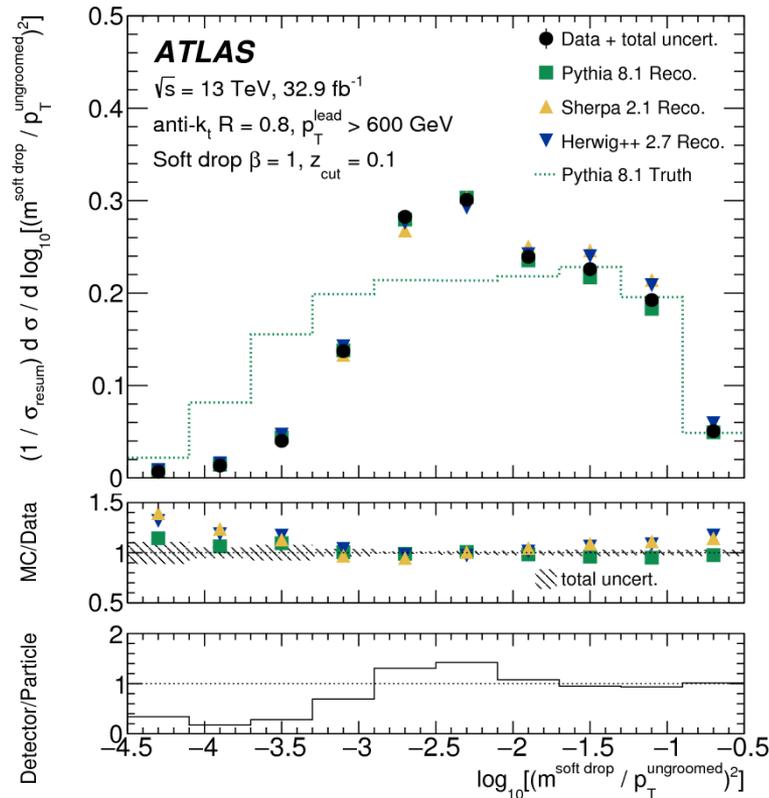
Dominating systematic uncertainties

Calorimeter (topo-cluster) energy scale (*bottom-up* approach)

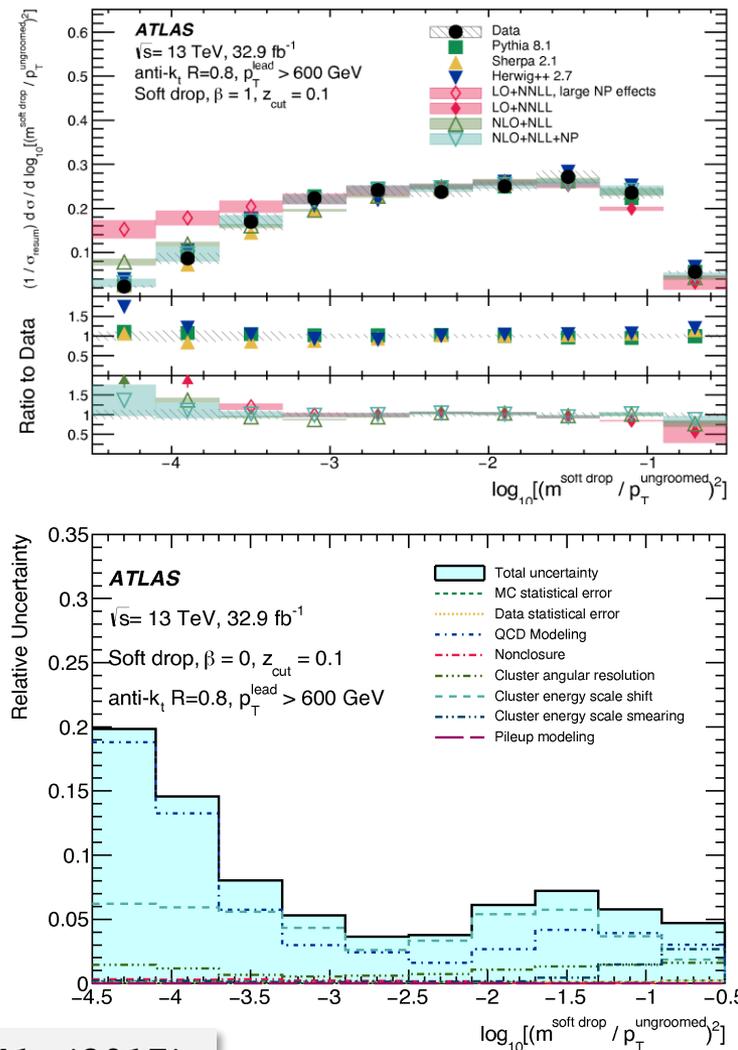
Modeling – limitations in fixed-order (particle) generators (PYTHIA and SHERPA)

Soft Drop Mass Measurement - Results

Detector-level ...



... unfolded measurements



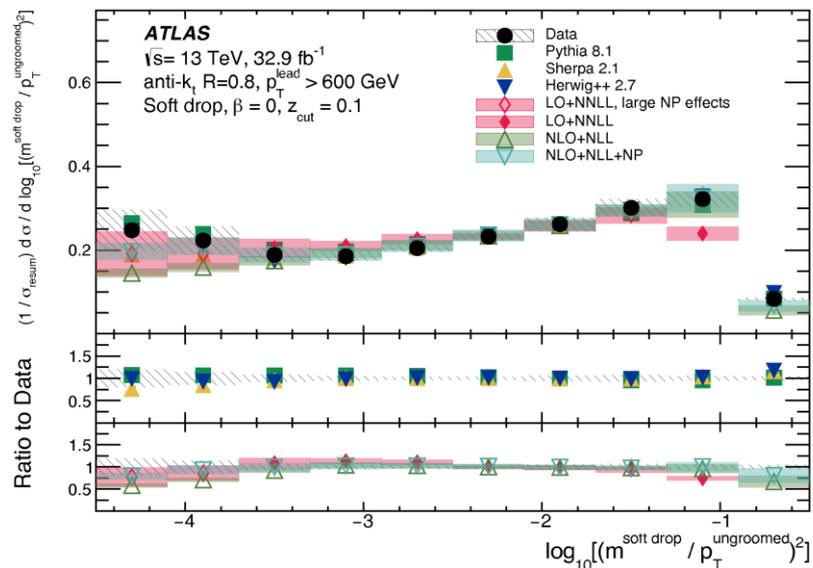
Systematic uncertainties →

arXiv:1711.08341 (2017)

Soft Drop Mass – Interpretation

$\beta = 0$

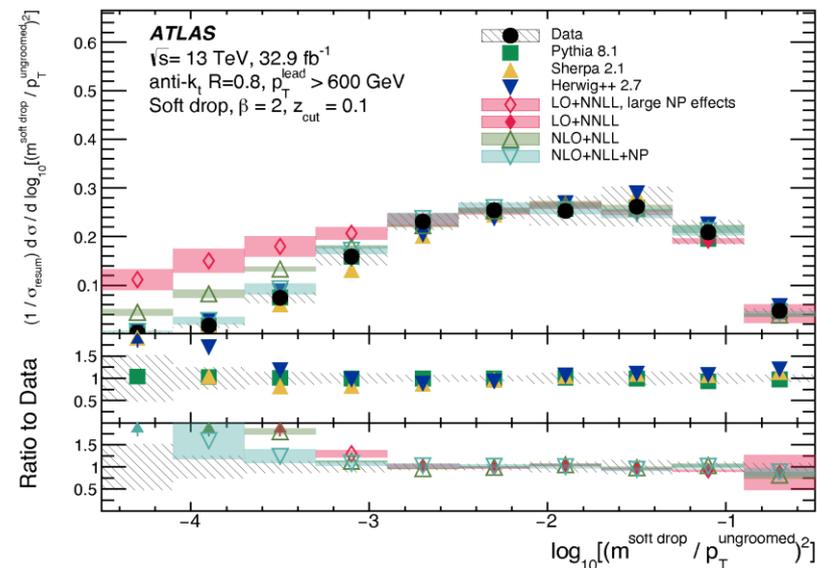
Removes soft \rightarrow collinear radiation



arXiv:1711.08341 (2017)

$\beta = 2$

Removes soft radiation

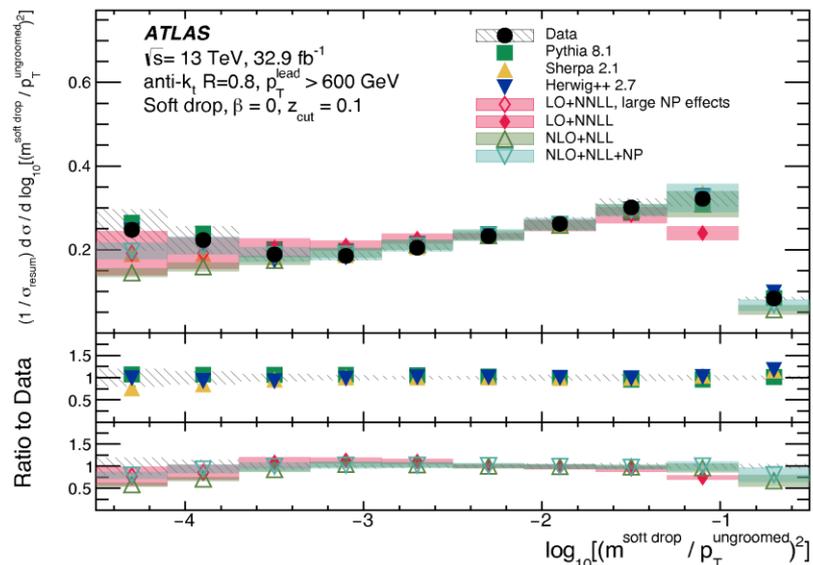


arXiv:1711.08341 (2017)

Soft Drop Mass – Interpretation

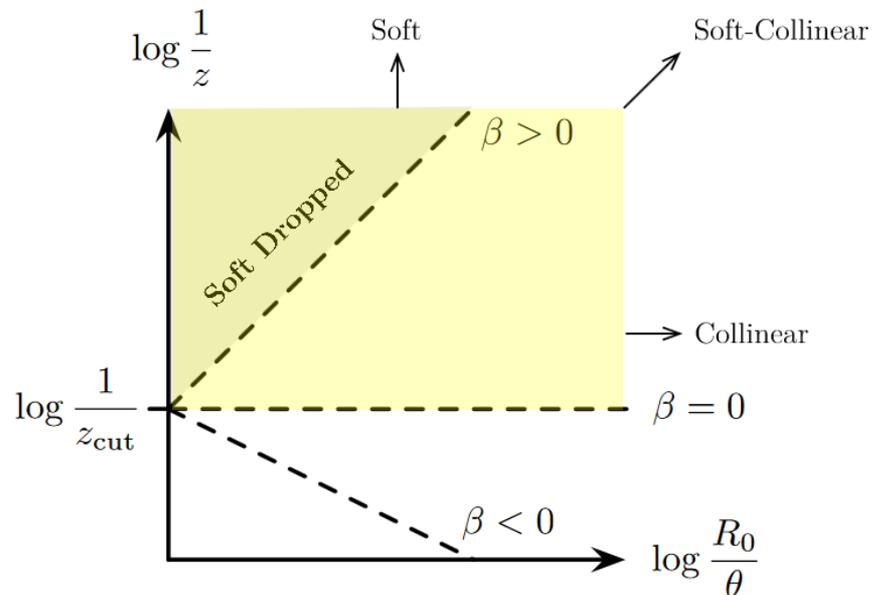
$\beta = 0$

Removes soft \rightarrow collinear radiation



$\beta = 2$

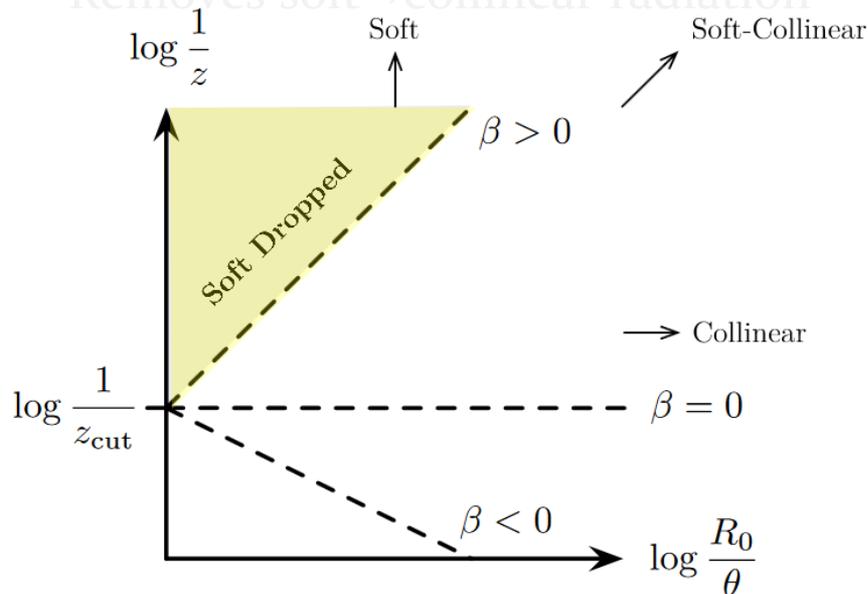
Removes soft radiation



Soft Drop Mass – Interpretation

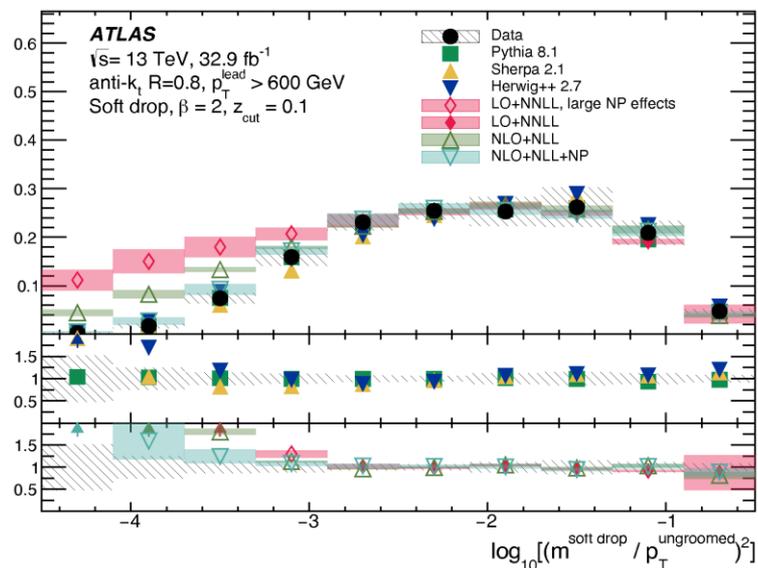
$\beta = 0$

Removes soft \rightarrow collinear radiation



$\beta = 2$

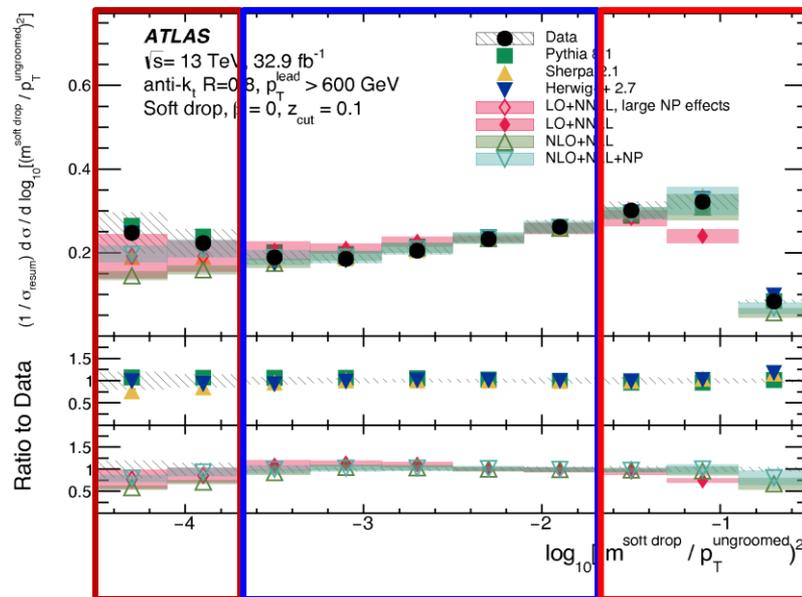
Removes soft radiation



Soft Drop Mass – Interpretation

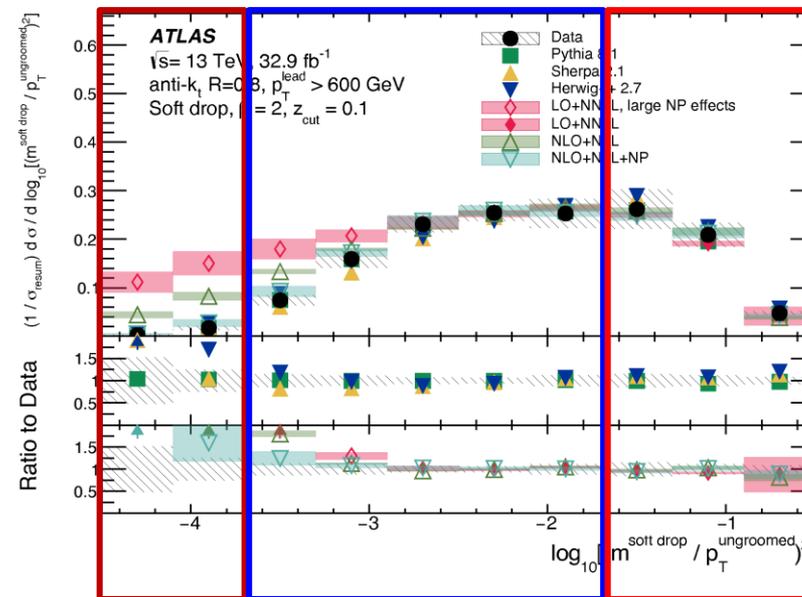
$\beta = 0$

Removes soft \rightarrow collinear radiation



$\beta = 2$

Removes soft radiation



< -3.7 non-perturbative regime

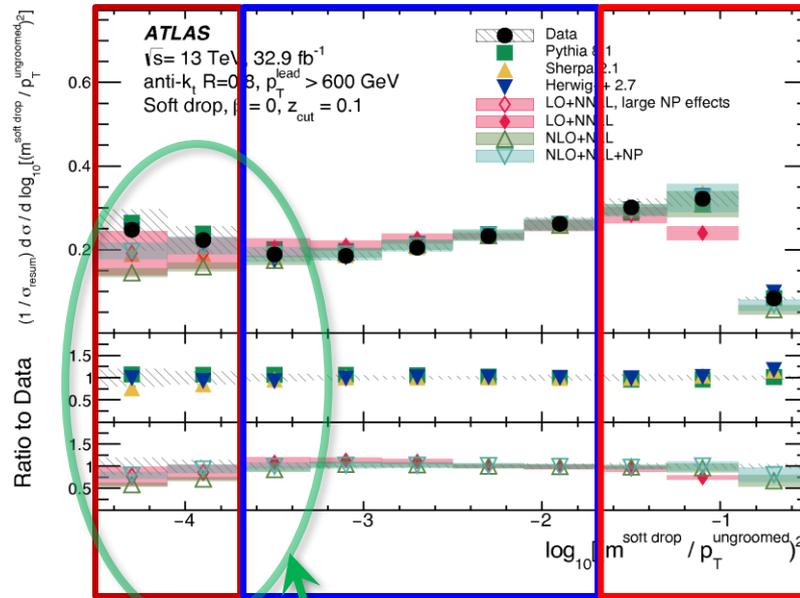
$-3.7 \rightarrow -1.7$ resummation region

> -1.7 fixed order regime
 (large angle gluon emission)

Soft Drop Mass – Interpretation

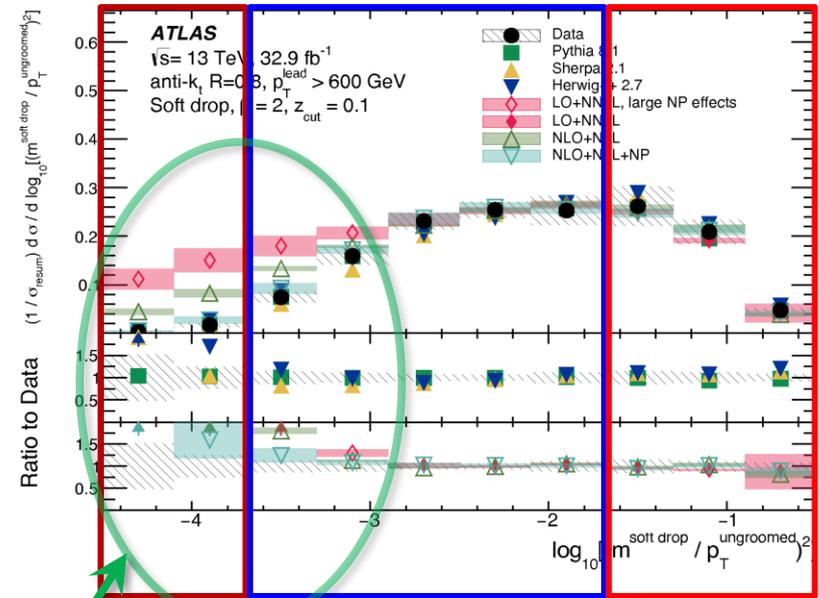
$\beta = 0$

Removes soft \rightarrow collinear radiation



$\beta = 2$

Removes soft radiation



Larger disagreements for less inclusive
 $(\beta = 2)$ soft contribution removal

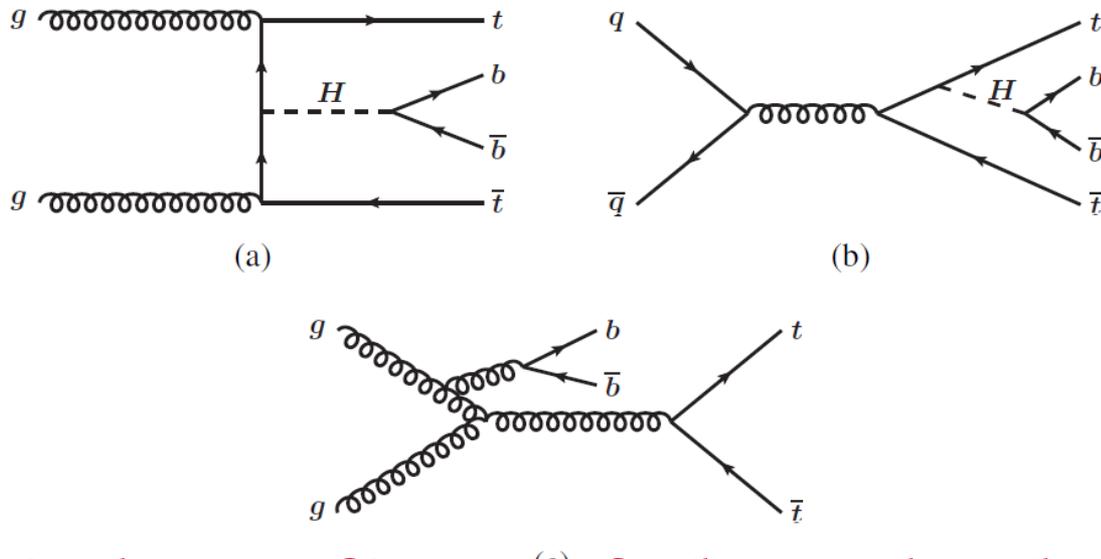
(large angle gluon emission)

Boosted $t\bar{t}H$ Production

Phys. Rev. D97 (2018) 072016

Motivation

Direct experimental access to top-quark Yukawa coupling

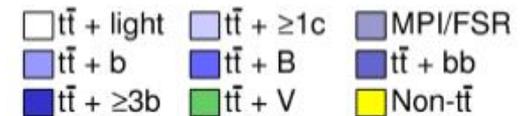
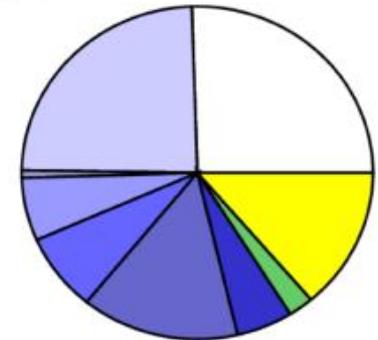


Final state of interest^(c) for boosted analysis

$H \rightarrow b\bar{b}$ reconstructed as anti- k_t jet ($R = 1.0$)
 with two b -tagged sub-jets

$t \rightarrow Wb$ reconstructed as anti- k_t jet ($R = 1.0$)
 with one b -tagged sub-jets and ≥ 1 sub-jet not b -tagged

SR^{boosted}



Boosted $t\bar{t}H$ Production

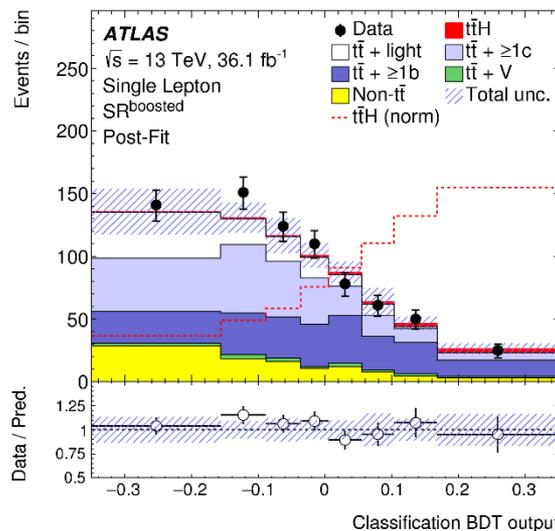
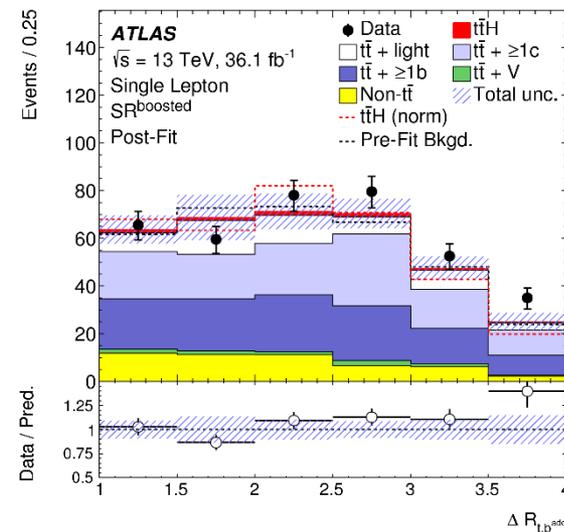
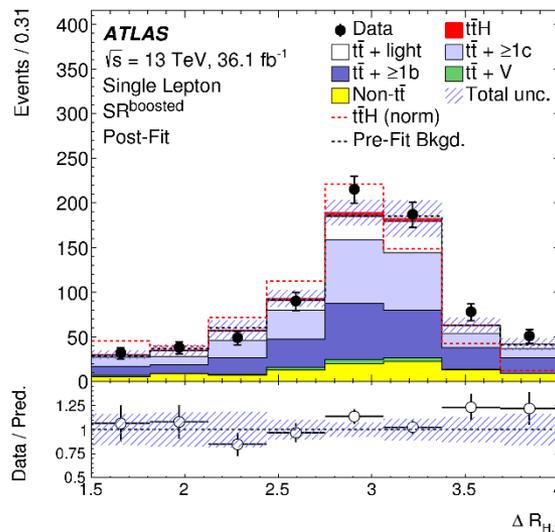
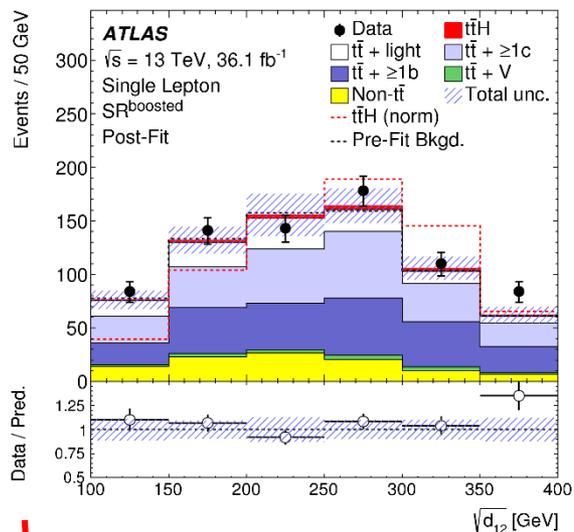
Highest discrimination in BDT

Phys. Rev. D97 (2018) 072016

Splitting scale $\sqrt{d_{12}}$

angular distance $H \leftrightarrow t$

angular distance $t \leftrightarrow b$



Jet substructure is mainstream analysis tool

Many relevant observables well understood

Powerful observables to extract hard internal jet structure often little dependent on pile-up

Residual or fundamental dependences on pile-up can be mitigated by appropriate grooming techniques

Well applicable both in searches and measurements

Measurement of mass of light quark/gluon jets provides guidance for theory for calculations in soft and soft-collinear regimes

Contribution to important measurements of Higgs production mechanisms

Direct searches for new heavy particles decaying into gluons or light quarks

Indirect searches for new heavy particles decaying into heavy quarks and/or gauge bosons

Exiting field poses significant experimental challenges

Limited detector capabilities

Resolution of energy flow in highly boosted regimes in calorimeters

Effects of pile-up at high-luminosity LHC operations – calorimeters and tracking (e.g. vertex resolution)

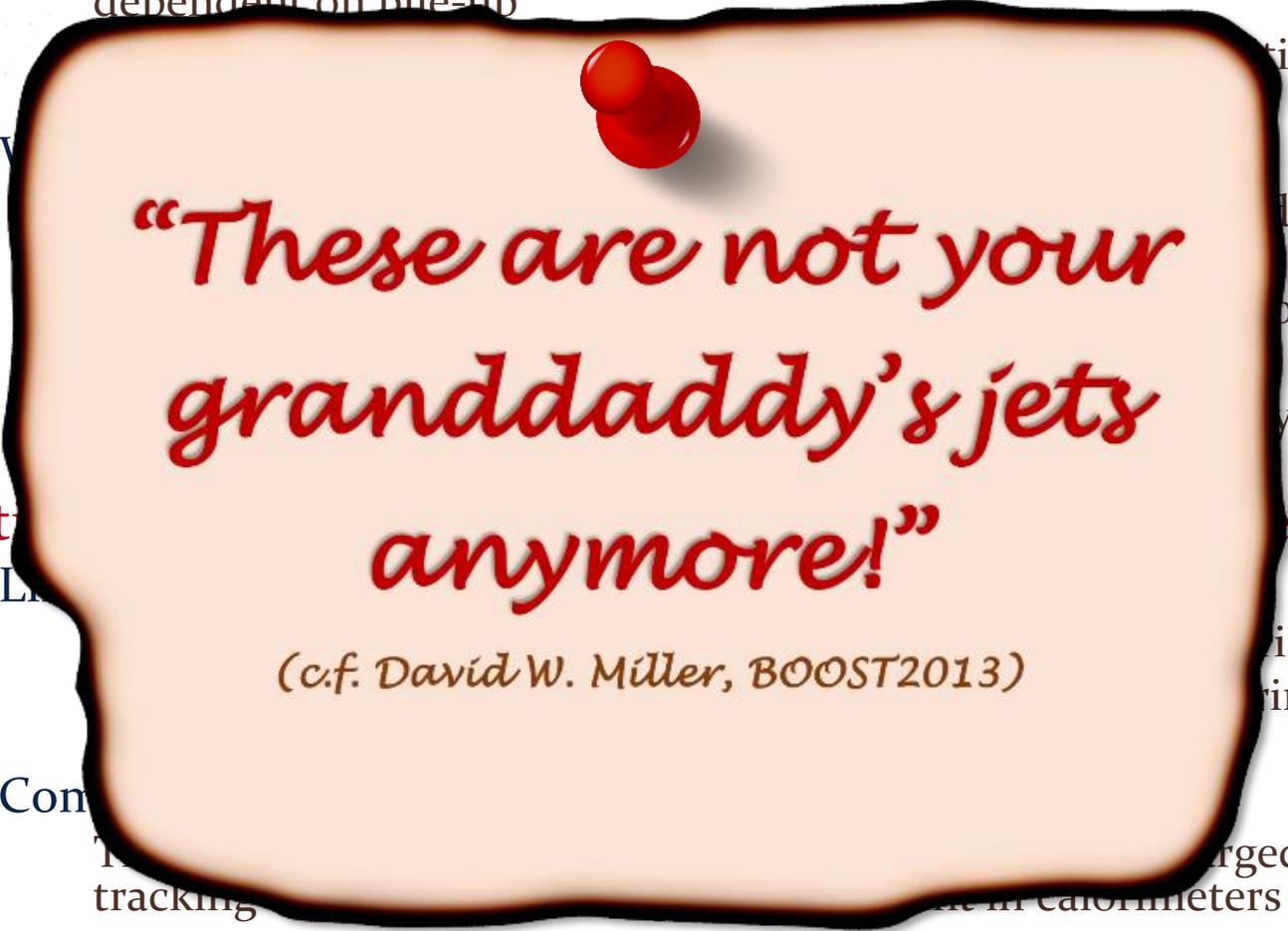
Combining reconstructed objects

Track-assisted reconstruction – combining direction of charged flow from tracking with complete energy measurement in calorimeters

Jet substructure is mainstream analysis tool

Many relevant observables well understood

Powerful observables to extract hard internal jet structure often little dependent on pile-up



*“These are not your
granddaddy’s jets
anymore!”*

(c.f. David W. Miller, BOOST2013)

Exit

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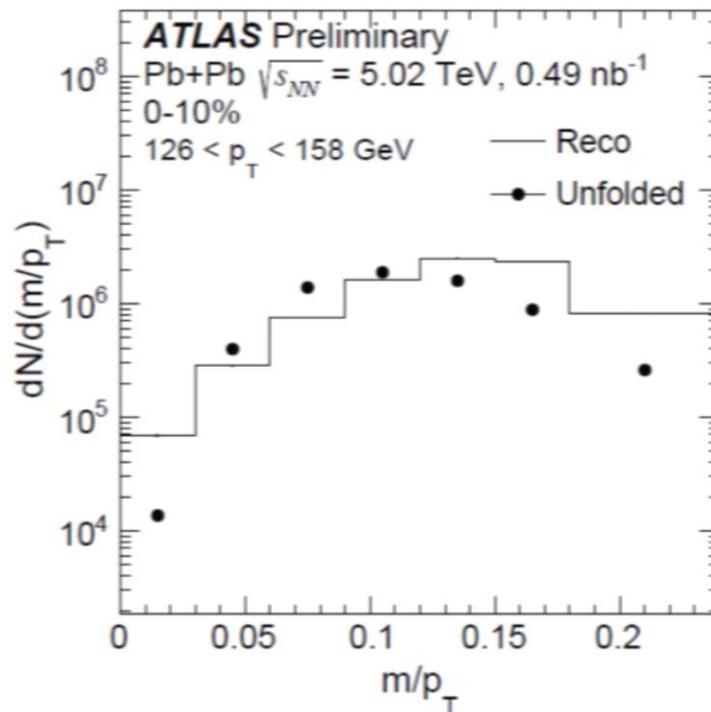
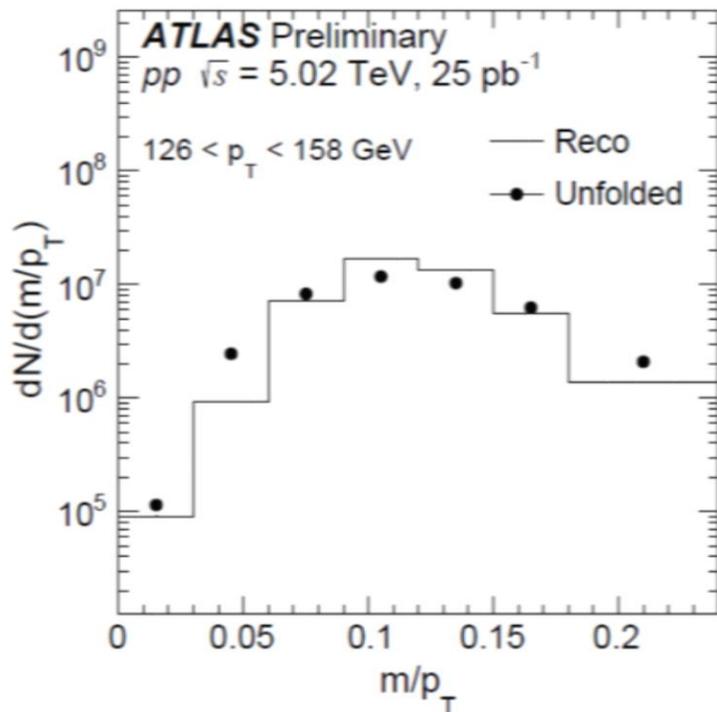
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arged flow from
e in calorimeters

Additional Material

Jet Substructure In Heavy Ion Collisions

- Measurement of m/p_T in pp and PbPb collisions at $\sqrt{s} = 5.02$ TeV
ATLAS-CONF-2018-014
- Jet mass measurement in PbPb gives insights in the modification of the jet while passing through quark-gluon plasma
- Jets are reconstructed with anti- k_t $R = 0.4$ algorithm from 0.1×0.1 towers
 - Underlying event is subtracted from towers before calibrations are applied

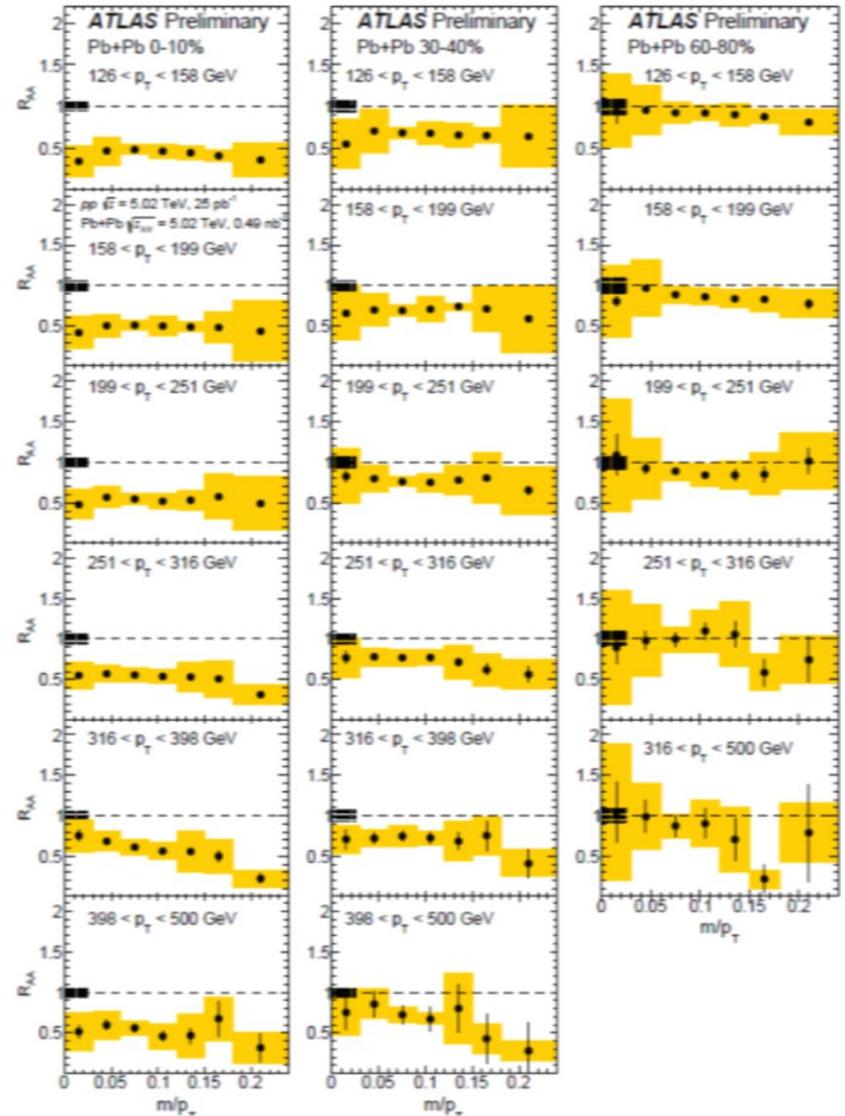


Jet Substructure In Heavy Ion Collisions

- Nuclear modification factor R_{AA} :
 - Quantifies modification of jet yields PbPb collisions with respect to pp

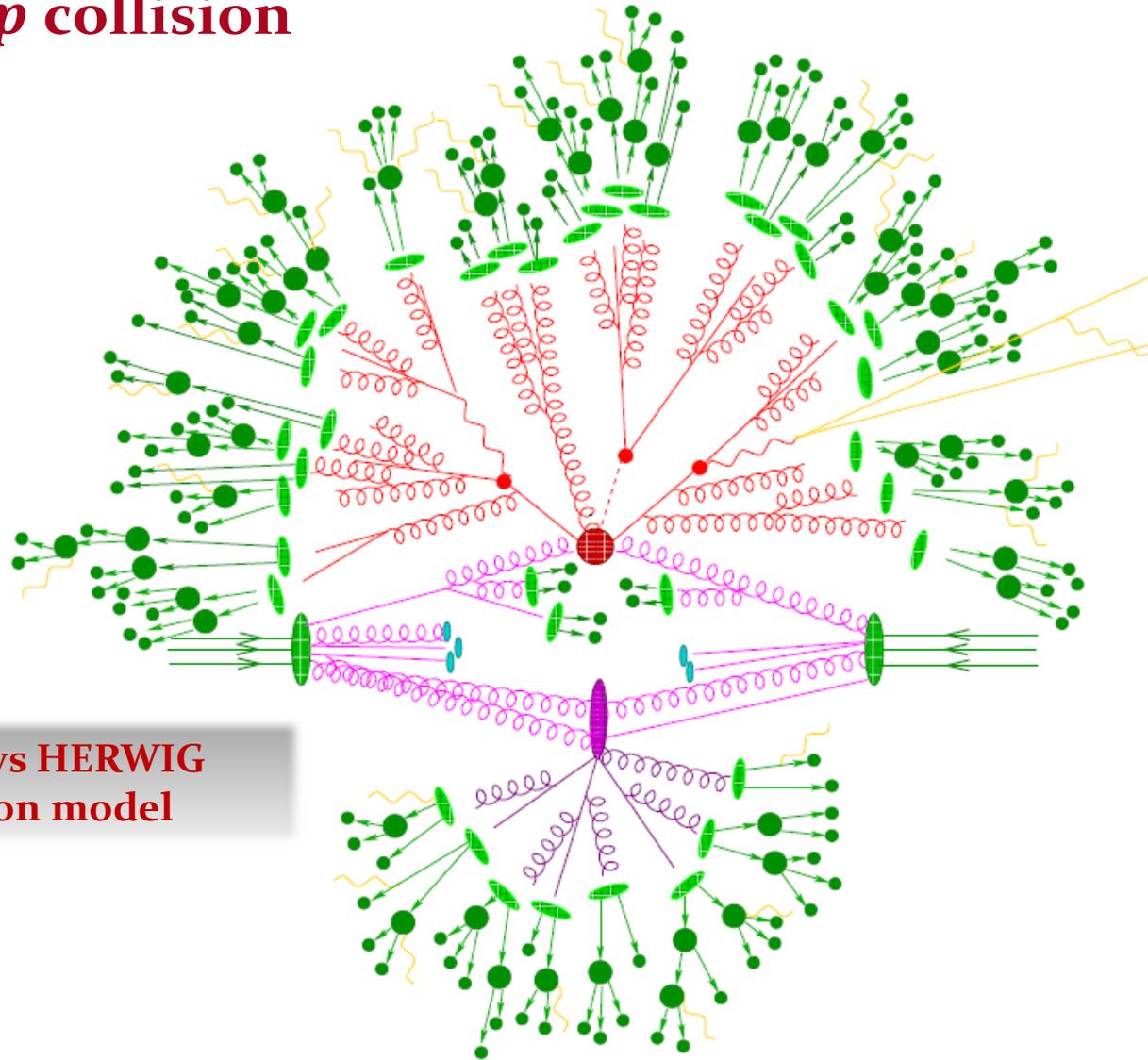
$$R_{AA}(m/p_T, p_T) = \frac{1}{N_{evt}} \frac{dN_{jet}^{Pb+Pb}}{d(m/p_T)}(p_T) \Big|_{cent}}{\langle T_{AA} \rangle \frac{d\sigma_{jet}^{PP}}{d(m/p_T)}(p_T)}$$

- $\langle T_{AA} \rangle$: nuclear thickness function
- R_{AA} measured for various p_T and centralities (based on E_T in FCal)
- Modification factor is relatively flat in m/p_T but shows suppression of jets in PbPb collisions



C. Delitzsch, LHCP 2018

Schematic *pp* collision

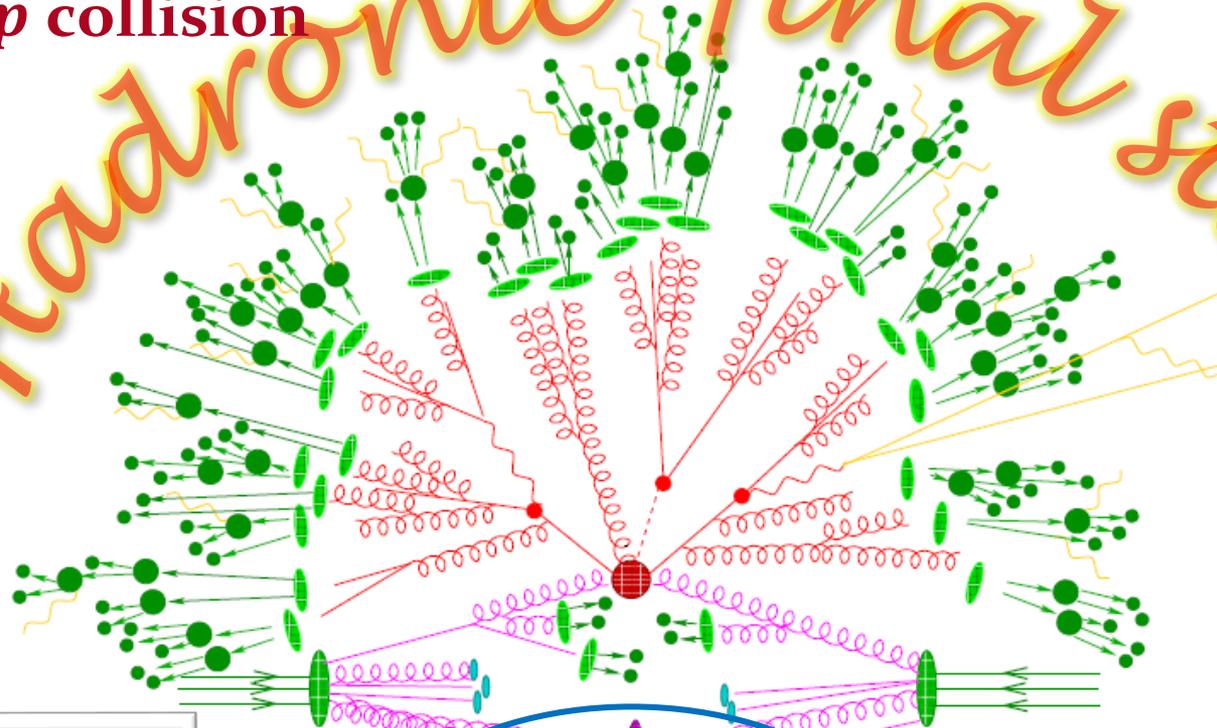


Graphics shows HERWIG
fragmentation model

Proton-Proton Collisions @ LHC

Schematic pp collision

Hadronic final state



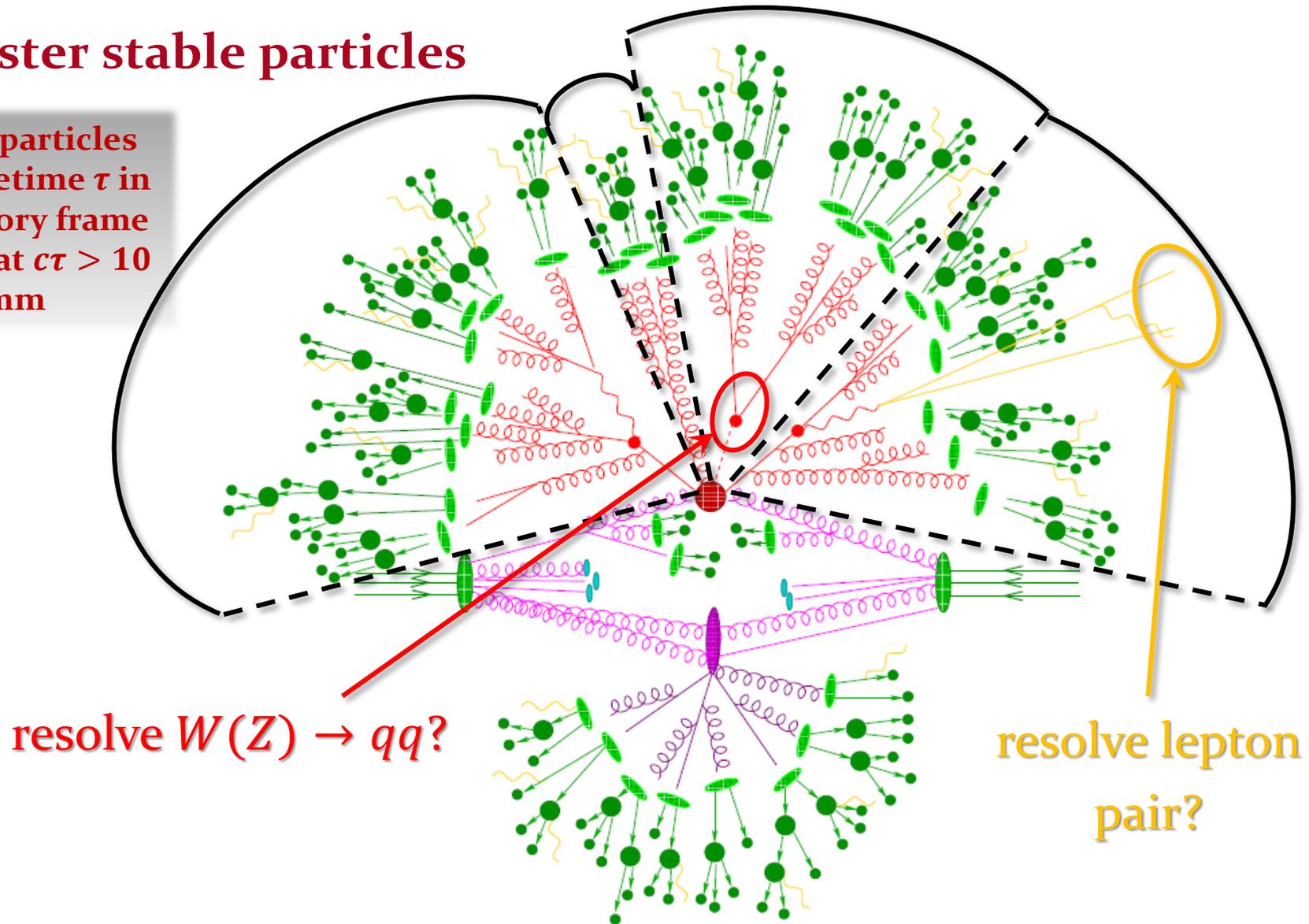
Interacting partons
Hard interaction
Parton Shower
Fragmentation
Hadronization

*Multiple
parton
interactions
(MPI)*

Particle Jet Reconstruction (MC only)

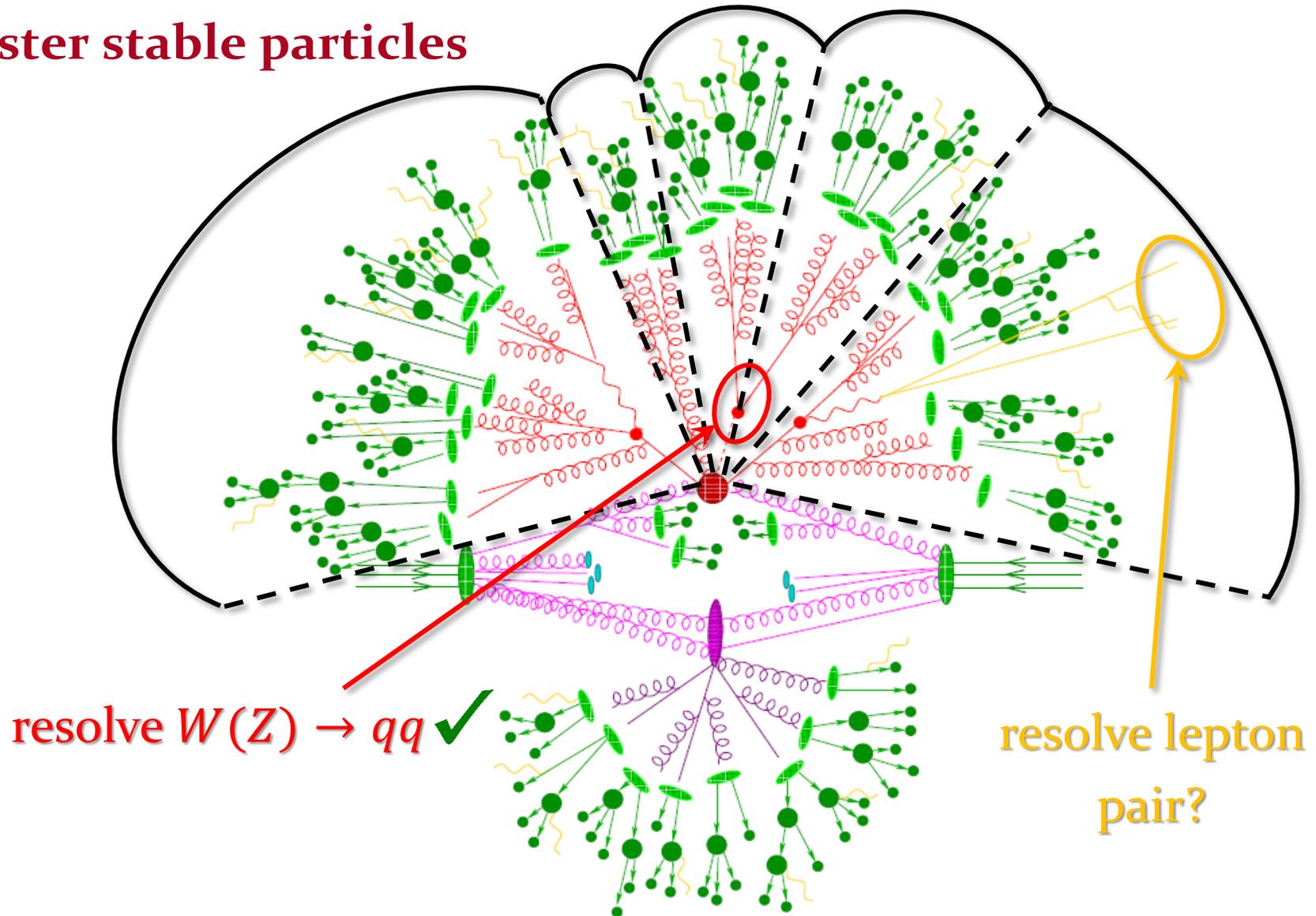
Cluster stable particles

Stable particles
have lifetime τ in
laboratory frame
such that $c\tau > 10$
mm



... or is this better?

Cluster stable particles



Final state component

Proxy for parton-level direction and energy

Partonic energy/momentum flow → indicative of process

Final states needs to be calculable in perturbation theory (pQCD)

Safe against soft emissions (**infrared safety**)

Stable against collinear splitting (**collinear safety**) – total energy carried by one or more particles in the same direction should have no effect

Algorithm choices

LHC uses recursive recombination algorithms (nearly) exclusively

Infrared safe – soft emissions may only affect jet shapes but not global kinematics

Collinear safe – no seeds used

Improvement from Tevatron

Computationally affordable implementations available – avoid seeded cone algorithms and workarounds to fix infrared and collinear issues

Allow comparisons experiment ↔ models (calculations)

Jet inputs

LHC – experiments correct jets to **particle level**

Jet algorithm type

Recursive recombination, seedless cones, ...

Recombination scheme

Two *proto-jets* recombined exclusively by **energy scheme** (adding four-momenta) – massive proto-jets after first recombination

Distance parameter (“radius”)

Defines maximum angular distance $\sqrt{\Delta y^2 + \Delta \phi^2}$ considered in clustering
Controls catchment area and jet shape – **spatial resolution** in energy flow

Angular scale for jets from gluon radiation – collinear limit in calculations

Minimum transverse momentum requirement

Sets acceptance/sensitivity in energy flow reconstruction

Agreed-upon concept at LHC ...

... since it's early days!

Favorite Jet Algorithms

k_T algorithm

Use intrinsic transverse momentum (k_T) between particles to decide on combination

Originally motivated by inversion of gluon splitting kernel $g \rightarrow gg$ Implements physically meaningful clustering sequence (“undoing” fragmentation)

Combination starts with the smallest available four-momentum

Large and irregularly shaped (y, φ) catchment area – jet four-momentum “walks” around a lot before final jet is formed

Cambridge/Aachen (C/A) flavor of k_T

Uses angular distance scale

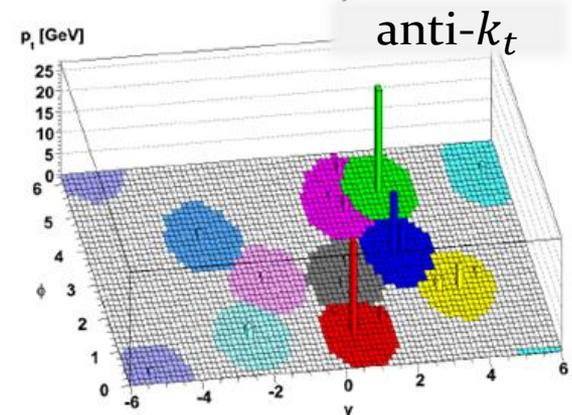
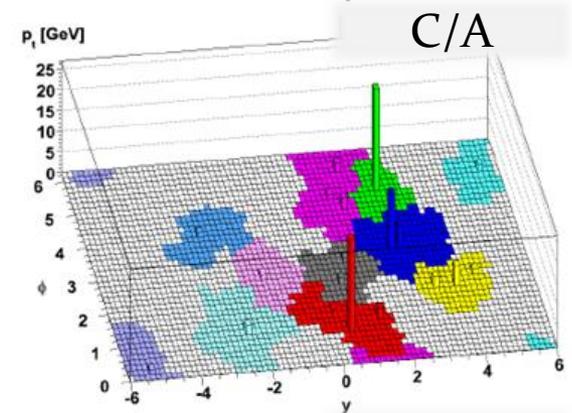
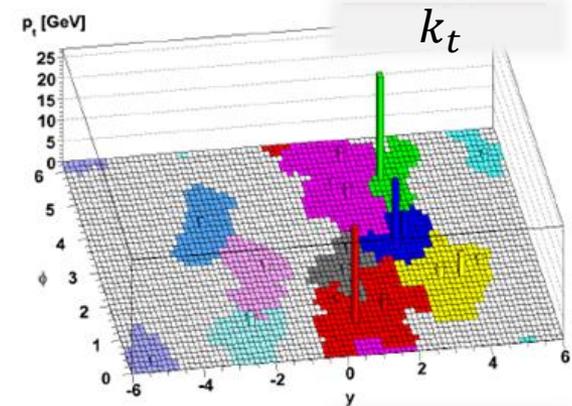
Similar jets than k_T

Anti- k_T algorithm

Inversion of k_T algorithm – start with largest four-momentum

Smaller and regular (circular) (y, φ) catchment area

Final jet physically meaningful but clustering sequence not physically meaningful (!)



Underlying event (experimental view)

Stochastic component

MPI generated independently from hard scatter interaction – but maybe reconnected by color flow

Total phase space limited by hard scatter

Correlated component

Initial (ISR) and final (FSR) state radiation off quarks and gluons involved in the hard scatter interaction

But we cannot disentangle them in the experiment

Suggestions for FSR jet tagging have been tested – low efficiency...

Pile-up

Diffuse emission of low-energy particles into a jet

Purely stochastic – no correlation to hard scatter interaction

Generated by additional (soft) proton-proton interactions occurring in the same bunch-crossing as the hard scatter

Scales with instantaneous luminosity – average number of additional interactions $\langle \mu \rangle$

Pile-up

Affects jet kinematics and shapes

Introduces additional energy/transverse momentum in jets

Deteriorates reconstruction of jet mass & internal particle flow features

Number of pile-up collisions μ

$$\mu = (\mathcal{L}_{\text{inst}} \times \sigma_{\text{inel}}^{pp}) / (N_b \times f_{\text{LHC}})$$

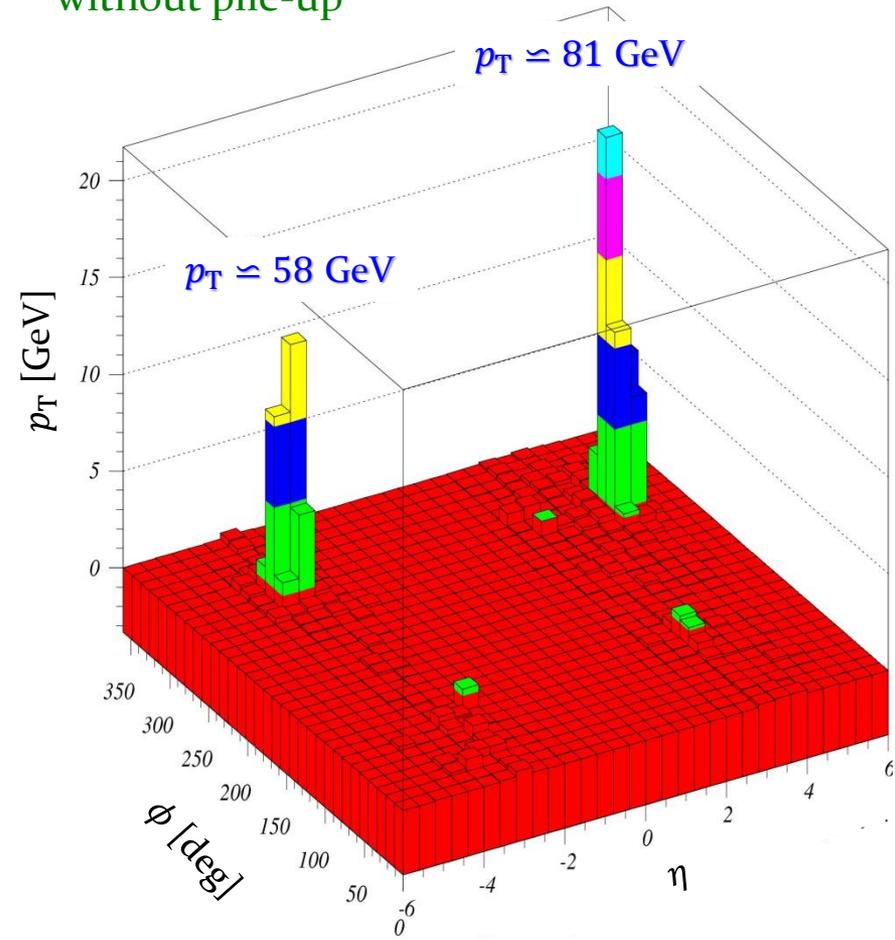
$\mathcal{L}_{\text{inst}}$ – instantaneous luminosity

$\sigma_{\text{inel}}^{pp}$ – inelastic pp cross section

N_b – number of colliding bunches in LHC

f_{LHC} – LHC revolution frequency

without pile-up



$qq \rightarrow qqWW \rightarrow qqH \rightarrow qqvvvv$

Pile-up

Affects jet kinematics and shapes

Introduces additional energy/transverse momentum in jets

Deteriorates reconstruction of jet mass & internal particle flow features

Number of pile-up collisions μ

$$\mu = (\mathcal{L}_{\text{inst}} \times \sigma_{\text{inel}}^{pp}) / (N_b \times f_{\text{LHC}})$$

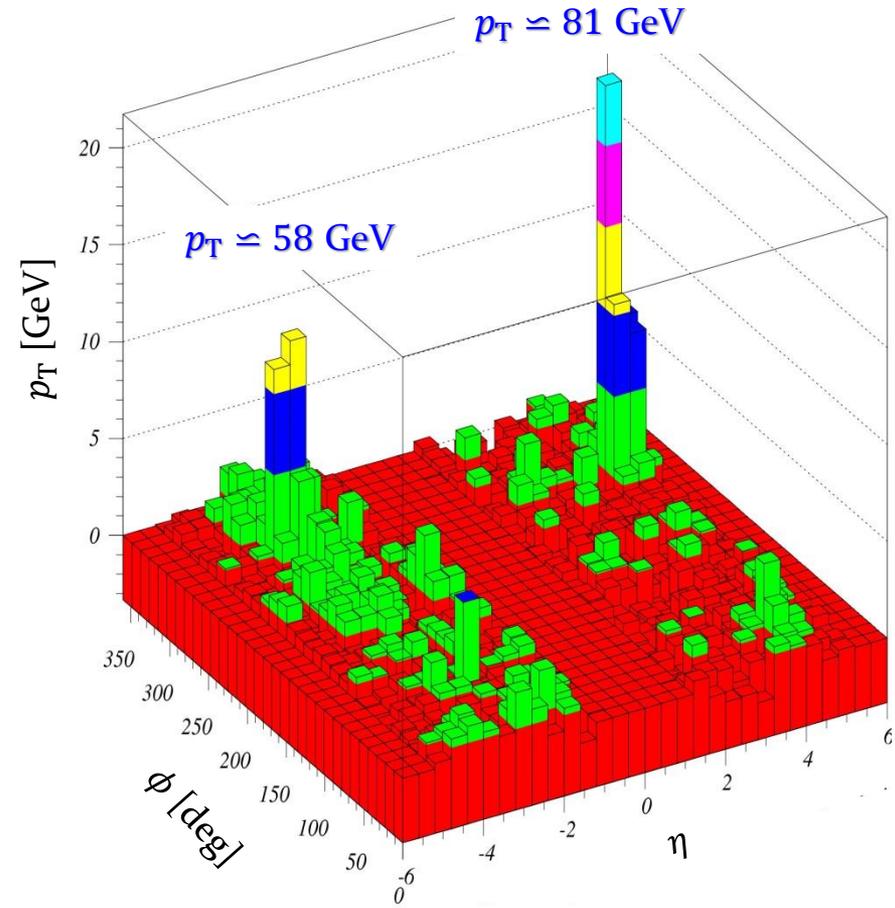
$\mathcal{L}_{\text{inst}}$ – instantaneous luminosity

$\sigma_{\text{inel}}^{pp}$ – inelastic pp cross section

N_b – number of colliding bunches in LHC

f_{LHC} – LHC revolution frequency

with pile-up @ $\mathcal{L} = 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$



$qq \rightarrow qqWW \rightarrow qqH \rightarrow qqvvvv$