Initial state physics: An introduction to the Color Glass Condensate

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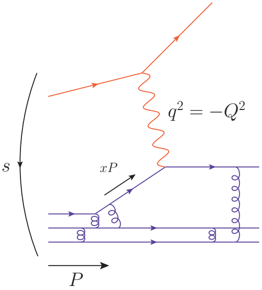
Hard Probe 2020

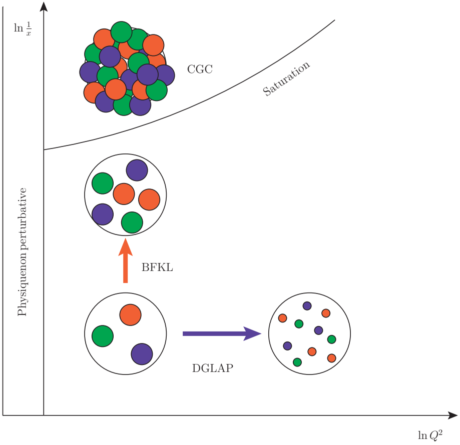
Student lectures

Overview

Accessing the partonic content of hadrons with an electromagnetic probe

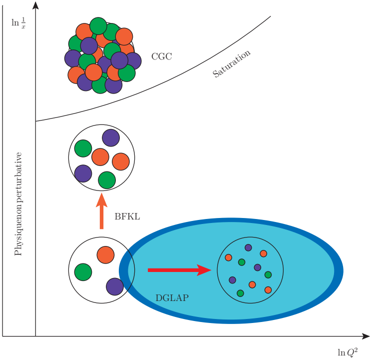


Electron-proton collision (parton model)



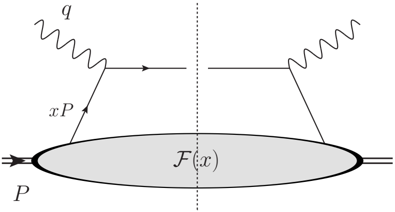
QCD at moderate $x_B = Q^2/s$

$$Q^2 \sim s$$



QCD factorization

processes with a hard scale $Q \gg \Lambda_{QCD}$



$$\sigma = \mathcal{F}(x,\mu) \otimes \mathcal{H}(x,\mu)$$

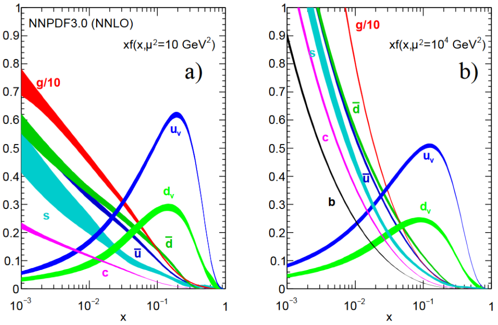
At a scale μ, the process is factorized into:

- A hard scattering subamplitude $\mathcal{H}(x,\mu)$
- A Parton Distribution Function (PDF) $\mathcal{F}(x,\mu)$

μ independence: DGLAP renormalization equation for $\mathcal{F}$

Parton Distribution Functions

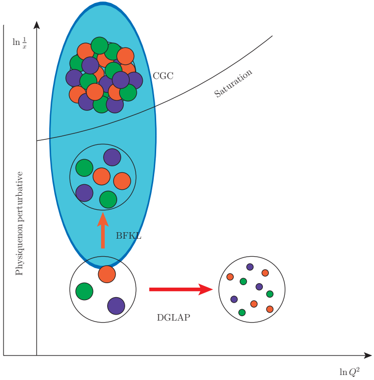
Gluon exchanges dominate at small x



[NNLO NNPDF3.0 global analysis, taken from PDG2018]

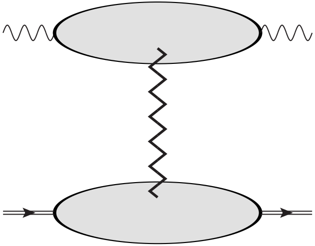
QCD at small $x_B = Q^2/s$

$$Q^2 \ll s$$



The Pomeron

Regge theory: for asymptotic values of s, an effective particle with the quantum numbers of the vacuum is exchanged

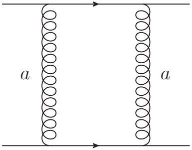


Positive C-parity: Pomeron exchange, negative C-parity: Odderon exchange

- How can we understand the Pomeron and the Odderon in perturbative QCD?
- How does it couple to hadrons?

Naive perturbative Pomeron and Odderon

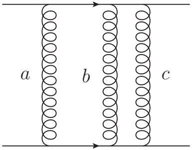
Naive perturbative description of the target hadron



Two gluons on a color singlet state

$$\mathrm{tr}(t^a t^a)$$

Leading Pomeron



Three gluons on a color singlet state

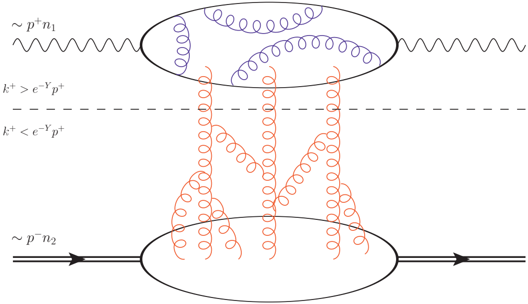
$$\mathrm{tr}(t^a t^b t^c) = \frac{1}{4}(d^{abc} + i f^{abc})$$

f^{abc}: subleading Pomeron
d^{abc}: leading Odderon

More involved but still for perturbative targets: BFKL, BKP, BLV...

Most general framework: small-x semiclassical effective theory

Effective semiclassical description of small x QCD

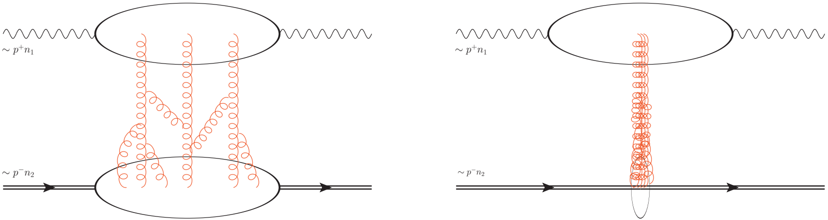


Let us split the gluonic field between "fast" and "slow" gluons

$$\mathcal{A}^{\mu a}(k^+, k^-, \vec{k}) = A^{\mu a}_{Y_c}(|k^+| > e^{-Y_c} p^+, k^-, \vec{k}) + b^{\mu a}_{Y_c}(|k^+| < e^{-Y_c} p^+, k^-, \vec{k})$$

$e^{-Y_c} \ll 1$

Large longitudinal boost to the projectile frame



$$b^+(x^+, x^-, \vec{x}) \qquad\qquad \frac{1}{\Lambda} b^+(\Lambda x^+, \frac{x^-}{\Lambda}, \vec{x})$$

$$b^-(x^+, x^-, \vec{x}) \quad \longrightarrow \quad \Lambda b^-(\Lambda x^+, \frac{x^-}{\Lambda}, \vec{x})$$

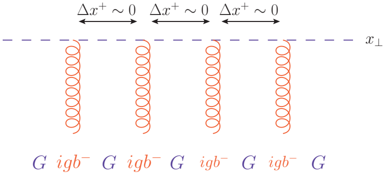
$$b^k(x^+, x^-, \vec{x}) \qquad \Lambda \sim \sqrt{\frac{s}{m_t^2}} \qquad b^k(\Lambda x^+, \frac{x^-}{\Lambda}, \vec{x})$$

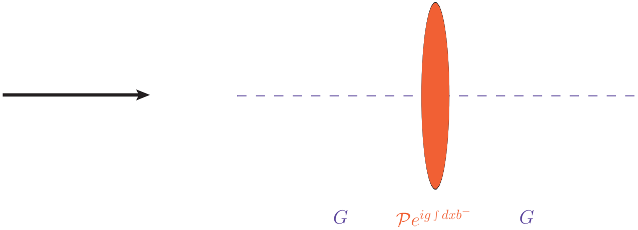
$$b^\mu(x) \to b^-(x)\, n_2^\mu = \delta(x^+)\, \mathbf{B}(\vec{x})\, n_2^\mu + O(\sqrt{\frac{m_t^2}{s}})$$

Shockwave approximation

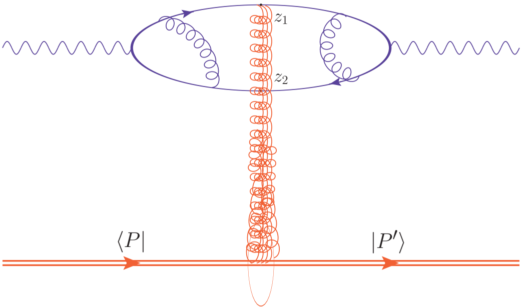
Effective Feynman rules in the slow background field

The interactions with the background field can be exponentiated





Factorized picture



Factorized amplitude

$$\mathcal{A}^{Y_c} = \int d^{D-2}\vec{z}_1 d^{D-2}\vec{z}_2 \, \Phi^{Y_c}(\vec{z}_1, \vec{z}_2) \, \langle P'|[\mathrm{Tr}(U^{Y_c}_{\vec{z}_1} U^{Y_c\dagger}_{\vec{z}_2}) - N_c]|P\rangle$$

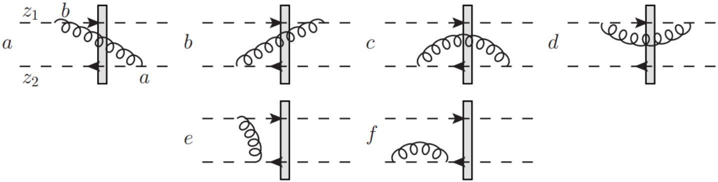
Dipole operator $\mathcal{U}^{Y_c}_{ij} = \frac{1}{N_c}\mathrm{Tr}(U^{Y_c}_{\vec{z}_i} U^{Y_c\dagger}_{\vec{z}_j}) - 1$

Written similarly for any number of Wilson lines in any color representation!

Y_c independence: B-JIMWLK hierarchy of equations
[Balitsky, Jalilian-Marian, Iancu, McLerran, Weigert, Leonidov, Kovner]

Evolution for the dipole operator

$$\mathcal{U}_{12}^{Y_c+\delta Y} - \mathcal{U}_{12}^{Y_c}$$



B-JIMWLK hierarchy of equations
[Balitsky, Jalilian-Marian, Iancu, McLerran, Weigert, Leonidov, Kovner]

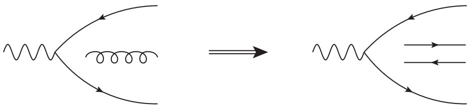
$$\frac{\partial \mathcal{U}_{12}^{Y_c}}{\partial Y_c} = \frac{\alpha_s N_c}{2\pi^2}\int d\vec{z}_3 \frac{\vec{z}_{12}^{\,2}}{\vec{z}_{13}^{\,2}\vec{z}_{23}^{\,2}}\left[\mathcal{U}_{13}^{Y_c}+\mathcal{U}_{32}^{Y_c}-\mathcal{U}_{12}^{Y_c}+\mathcal{U}_{13}^{Y_c}\,\mathcal{U}_{32}^{Y_c}\right]$$

$$\frac{\partial \mathcal{U}_{13}^{Y_c}\,\mathcal{U}_{32}^{Y_c}}{\partial Y_c} = \ldots$$

Evolves a dipole into a double dipole

The BK equation

Mean field approximation, or 't Hooft planar limit $N_c \to \infty$ in the dipole B-JIMWLK equation



[Mueller]

$\Rightarrow$ BK equation [Balitsky, 1995] [Kovchegov, 1999]

$$\frac{\partial \left\langle \mathcal{U}_{12}^{Y_c} \right\rangle}{\partial Y_c} = \frac{\alpha_s N_c}{2\pi^2} \int d\vec{z}_3 \, \frac{\vec{z}_{12}^{\,2}}{\vec{z}_{13}^{\,2}\, \vec{z}_{23}^{\,2}} \left[\left\langle \mathcal{U}_{13}^{Y_c} \right\rangle + \left\langle \mathcal{U}_{32}^{Y_c} \right\rangle - \left\langle \mathcal{U}_{12}^{Y_c} \right\rangle + \left\langle \mathcal{U}_{13}^{Y_c} \right\rangle \left\langle \mathcal{U}_{32}^{Y_c} \right\rangle \right]$$

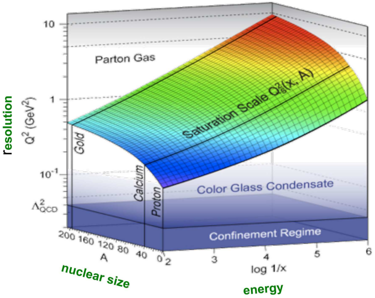
BFKL/BKP part — Triple pomeron vertex

Non-linear term : one type of saturation

Non-perturbative elements are compatible with CGC-type models

Saturation scale: a quick estimate

The saturation scale Q_s



Gluons per unit area

$$\rho \propto \frac{xG_A(x,Q^2)}{\pi R_A^2}$$

Recombination cross section

$$\sigma_{gg\to g} \propto \frac{\alpha_s}{Q^2}$$

Saturation starts when $\rho\sigma \simeq 1$, which means Q_s^2 solves

$$Q_s^2 \propto \alpha_s \frac{xG_A(x, Q_s^2)}{\pi R_A^2}.$$

$$Q_s^2 \propto A^{1/3} x^{-0.3}$$

Loop corrections: probing a single hadron

One-loop corrections with saturation effects: state of the art

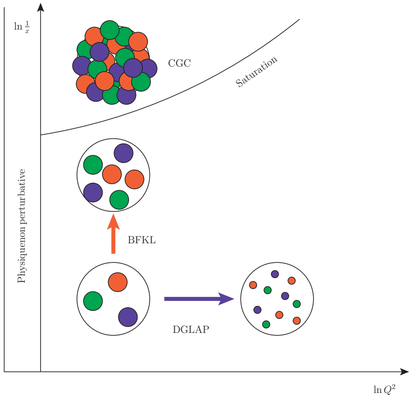
Evolution

- Dipole evolution [Balitsky, Chirilli]
- 3-point operator evolution [Balitsky, Gerasimov, Grabovsky]
- 4-point operator evolution [Grabovsky]
- Full JIMWLK Hamiltonian [Kovner, Lublinsky, Mulian, 2014]

Observables

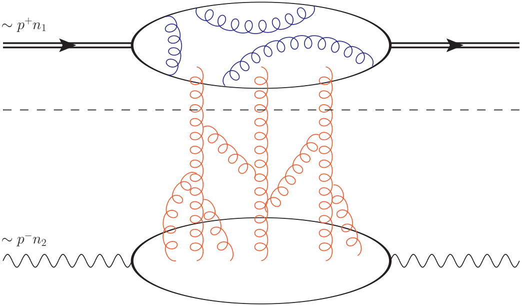
- Fully inclusive Deep Inelastic Scattering [Balitsky, Chirilli], [Beuf], [Hänninen, Lappi, Paatelainen]
- (Semi-inclusive) Photon+dijet in for *ep* and *eA* [Roy, Venugopalan]
- Exclusive dijet in *ep*, *eA*, γp or γA [RB, Grabovsky, Szymanowski, Wallon]
- Exclusive light vector meson in *ep* and *eA* [RB, Ivanov, Grabovsky, Szymanowski, Wallon]

Summary: probing a single hadron



Overview

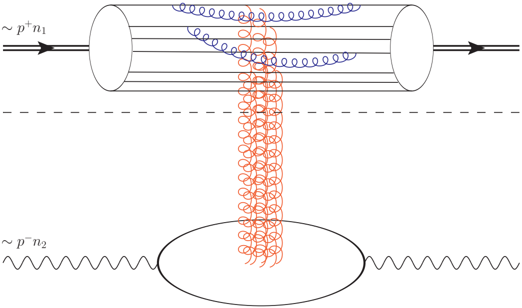
The Color Glass Condensate



fast partons $\leftrightarrow$ valence partons

slow gluons $\leftrightarrow$ wee gluons

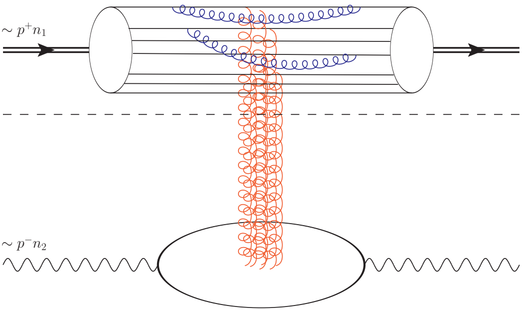
The Color Glass Condensate



Hadron wave function = collection of static color sources

Color sources ρ are classical random variables, treated with a weight function $W_Y[\rho]$

The Color Glass Condensate



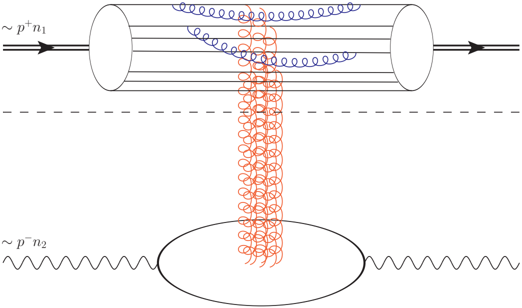
Static source = static current of color charge

$$J_a^\mu = \delta^{\mu +} \rho_a(x)$$

Wee gluons: solutions to the classical Yang-Mills equation with the source

$$[D_\nu, F^{\mu\nu}] = \delta^{\mu +} \rho_a(x) T^a$$

The Color Glass Condensate



Target matrix elements → averages over configurations of sources and dynamical fields A^μ

$$\frac{\langle P|\mathcal{O}|P\rangle}{\langle P|P\rangle} \to \langle \mathcal{O} \rangle = \int \mathcal{D}\rho\, \mathcal{D}A^\mu\, W[\rho]\, e^{iS[\rho,A]}\, \mathcal{O}[\rho, A]$$

The MV model

McLerran-Venugopalan model

- Sources $\simeq$ valence quarks $\Rightarrow$ number of sources $\sim N_c A$
- Transverse radius $R_A \sim A^{1/3}\Lambda_{\rm QCD}^{-1}$
- Transverse resolution of the probe $1/Q^2$
- Number of sources seens by the probe $\Delta N = \frac{\Lambda_{\rm QCD}^2}{Q^2}\frac{N_c A^{1/3}}{\pi}$

If $Q^2 \ll \Lambda_{\rm QCD}^2 A^{1/3}$, a large number of sources is probed

Random distribution of sources, total color charge probed is 0:

$$\langle \mathcal{Q} \rangle = \int_{1/Q^2} d^2\vec{x} \int dx^- \rho(x^-, \vec{x}) = 0$$

The MV model

McLerran-Venugopalan model

- Assume that $\langle \rho_a(x^-, \vec{x}) \rangle = 0$
- Write that
  $\langle \rho_a(x^-, \vec{x}) \rho_b(y^-, \vec{y}) \rangle = g_s^2 \delta_{ab} \delta(x^- - y^-) \delta(\vec{x} - \vec{y}) \lambda(x^-)$
- Assume that higher-point functions vanish

Correlators are generated from a Gaussian weight function

$$\Phi[\rho] \propto \exp\left(-\frac{1}{2} \int d^2\vec{x} \, \frac{\rho_a \, \rho_a}{\mu^2} \right), \quad \mu \propto \int dx^- \lambda(x^-)$$

Target matrix elements:

$$\frac{\langle P | \mathcal{O} | P \rangle}{\langle P | P \rangle} \to \frac{\int \mathcal{D}\rho \, \Phi[\rho] \, \mathcal{O}}{\int \mathcal{D}\rho \, \Phi[\rho]}$$

The MV model

Beyond the McLerran-Venugopalan model

Possible extensions

- Add a transverse dependence
  $\langle \rho_a(x^-,\vec{x})\rho_b(y^-,\vec{y})\rangle = g_s^2 \delta_{ab}\delta(x^- - y^-)\delta(\vec{x}-\vec{y})\lambda(x^-,\vec{x})$
- Include higher-point functions, or use non-Gaussian weight functions

$$\Phi[\rho] \propto \exp\left(-\frac{1}{2}\int d^2\vec{x}\left[\frac{\rho_a\,\rho_a}{\mu^2} - \frac{d^{abc}\rho_a\rho_b\rho_c}{\kappa}\right]\right)$$

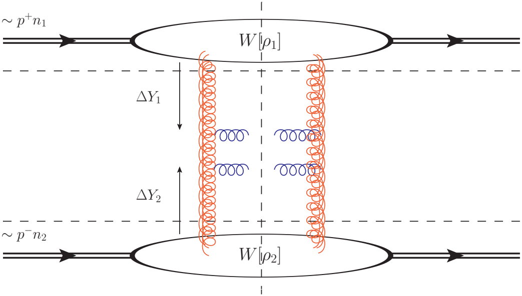
[Jeon, Venugopalan]

$\rho_a\rho_a$: Pomeron term, $d^{abc}\rho_a\rho_b\rho_c$: Odderon term

Overview

Hadron-hadron collisions in the CGC

Collisions of two distributions of color sources



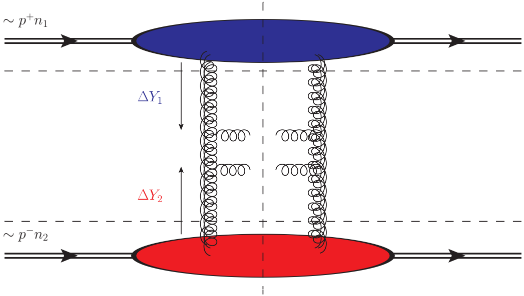
Expectation value of an operator

$$\langle \mathcal{O} \rangle = \int \mathcal{D}\rho_1 \, \mathcal{D}\rho_2 \, W_{Y_1}[\rho_1] \, W_{Y_2}[\rho_2] \, \mathcal{O}[\rho_1, \rho_2]$$

Source terms in both light cone directions $J_1^\mu = \delta^{\mu +} \rho_1$ and $J_2^\nu = \delta^{\nu -} \rho_2$...

Hadron-hadron collisions in the CGC

Two different saturation scales

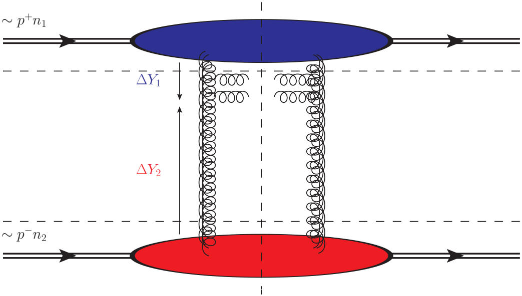


Projectile saturation scale $Q_{s1}^2 \propto (A_1/x_1)^{1/3}$

Target saturation scale $Q_{s2}^2 \propto (A_2/x_2)^{-1/3}$

Hadron-hadron collisions in the CGC

Hybrid factorization ansatz [Dumitru, Hayashigaki, Jalilian-Marian]



At forward rapidities, we can use the CGC to describe the target, while using colinear factorization to describe the projectile.

Allows to study the target with well-understood descriptions of the projectile.

Loop corrections with the hybrid factorization ansatz

One-loop corrections with saturation effects: state of the art

Evolution

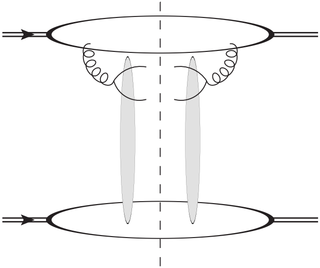
- Dipole evolution [Balitsky, Chirilli]
- 3-point operator evolution [Balitsky, Gerasimov, Grabovsky]
- 4-point operator evolution [Grabovsky]
- Full JIMWLK Hamiltonian [Kovner, Lublinsky, Mulian, 2014]

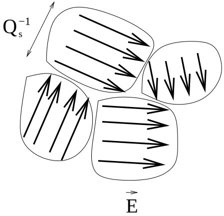
Observable

- Semi-inclusive hadron production in hybrid factorization [Chirilli, Xiao, Yuan], [Altinoluk, Armesto, Beuf, Kovner, Lublinsky]

Hybrid processes

Example: domain structure and the correlation limit





Picture from [Kovner, Lublisnky]

A pair of partons from the splitting of a colinear gluon from the projectile probes the target as a dipole of size $r_\perp$.

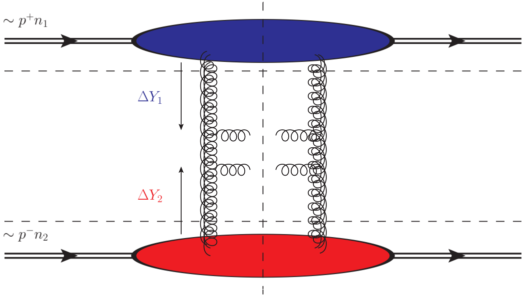
Domain structure: the target contains domains of oriented chromo-electric fields of size $1/Q_s$.

Small dipoles $|r_\perp| \ll 1/Q_s$ will probe a single domain. In momentum space, small dipole = back-to-back dijet.

Thus local correlations in the target lead to momentum correlations in the outgoing state

Hadron-hadron collisions in the CGC

Two different saturation scales

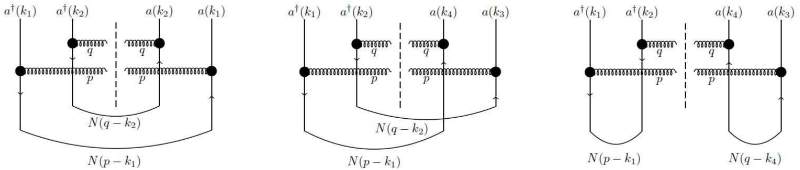


Projectile saturation scale $Q_{s1}^2 \propto (A_1/x_1)^{1/3}$

Target saturation scale $Q_{s2}^2 \propto (A_2/x_2)^{-1/3}$

Glasma graphs

An example of glasma graphs: double inclusive gluon production



Picture from [Altinoluk, Armesto]

- 2 gluons from the projectile: need to compute $\langle AAAA \rangle_{A_1}$. Assumption: each gluon comes from a different color charge density
- Scattering with the dense target via Wilson line operators: adjoint dipoles $N(p-k_1)$ and $N(q-k_2)$.
- Eikonal coupling between the t-channel gluons and the measured gluons via Lipatov vertices

Glasma graphs

Beyond glasma graphs

- Assumption: each gluon comes from a different color charge density:
  - A single charge could emit a gluon which splits into a gluon pairs [Kovner, Lublinsky, Skokov], [Kovchegov, Skokov]
  - Other assumption to relax: mean field $\langle AAAA\rangle_{A_1} \to \langle AA\rangle_{A_1}\langle AA\rangle_{A_1}$
- adjoint dipoles $N(p-k_1)$ and $N(q-k_2)$.
  - Also a possibility to relax the mean field approximation: $\langle N(p-k_1)N(q-k_2)\rangle_{A_2} \to \langle N(p-k_1)\rangle_{A_2}\langle N(q-k_2)\rangle_{A_2}$
- Eikonal coupling between the t-channel gluons and the measured gluons via Lipatov vertices
  - Posibility to include sub-eikonal corrections [Agostini, Altinoluk, Armesto]

Overview

Summary

A few typical CGC topics:

- One-loop corrections and precision phenomenology
- Target models beyond MV
- Correlations from the domain structure, from glasma graphs and beyond
- Spin effects in the CGC?
- Odd harmonics in the CGC?