

Groomed and ungroomed jet substructure from SCET

Kyle Lee

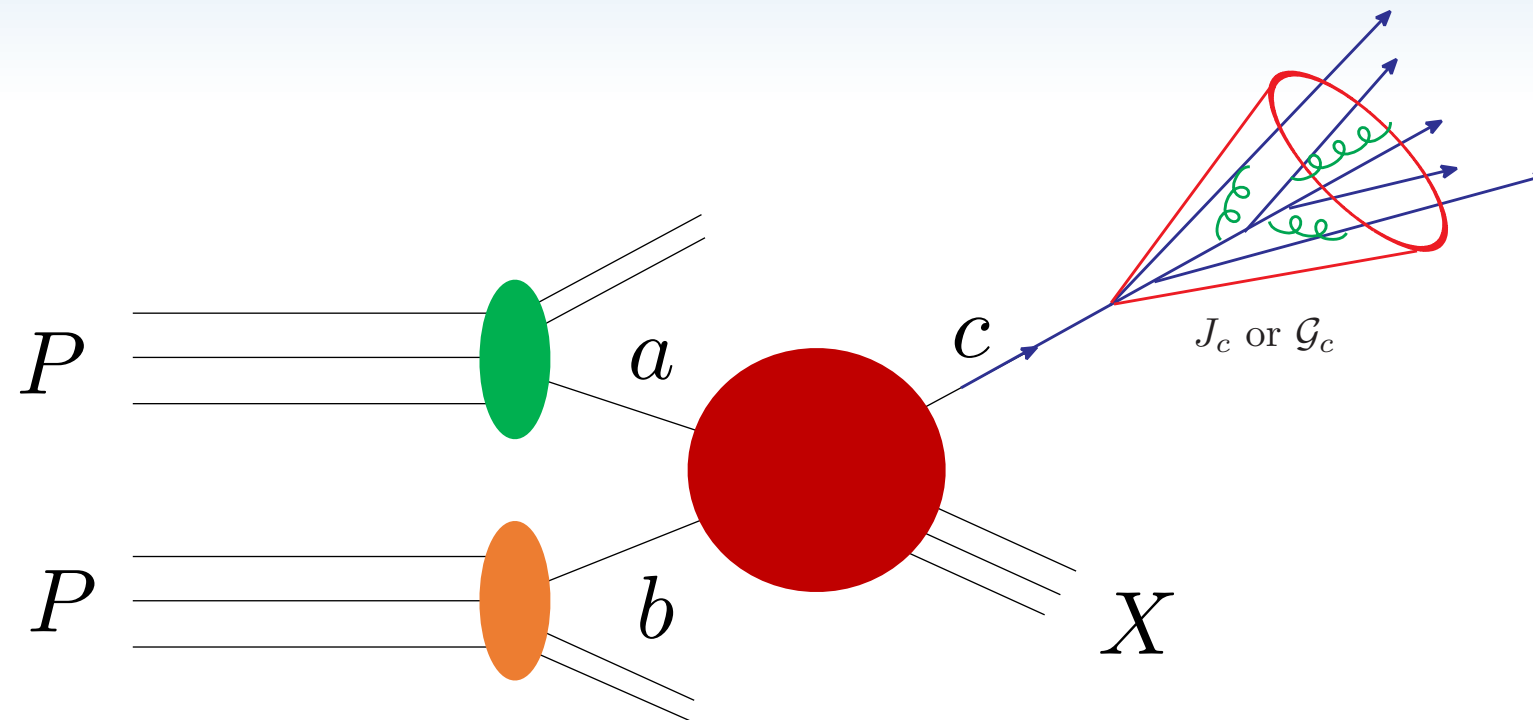
Stony Brook University

Works in collaboration with Zhong-Bo Kang, Xiaohui Liu, and Felix Ringer

Santa Fe Jets and Heavy Flavor Workshop
01/28/19 - 01/30/19



Processes of Interest



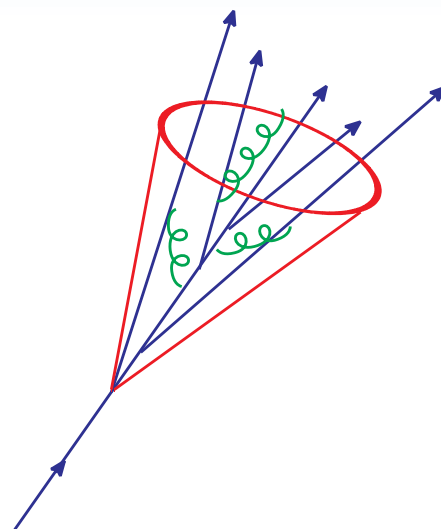
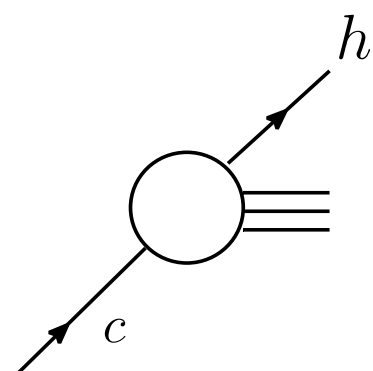
- ▶ We want to study semi-inclusive jet production
 $p + p \rightarrow \text{Jet}(\text{(with/without) substructure}) + X$
- ▶ More statistics. No veto on additional jets.

Plans of this talk

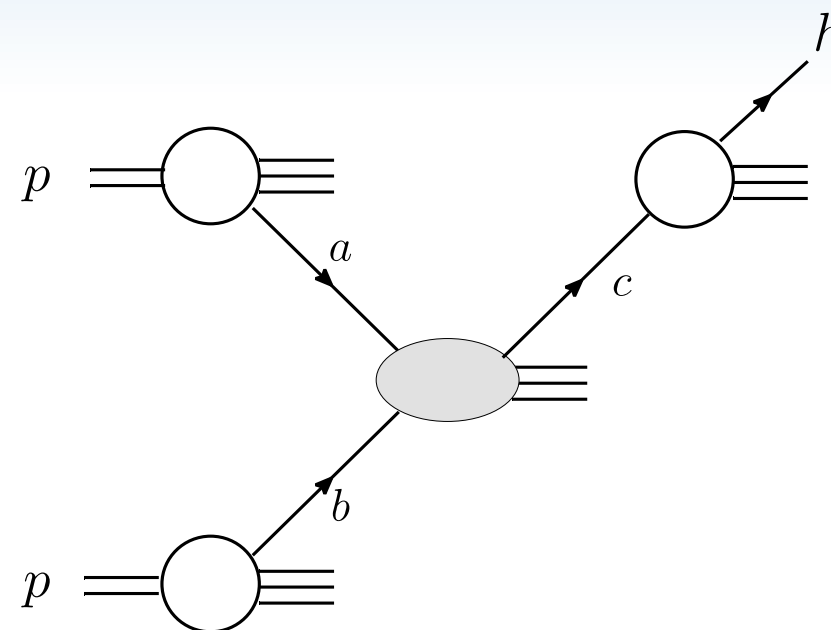
- Inclusive jet production at the LHC
- Ungroomed jet angularity measurements at the LHC
- Role of non-perturbative effects
- Groomed observables
- Conclusions

Factorization of Inclusive Jet Production

$$R \ll 1$$



from



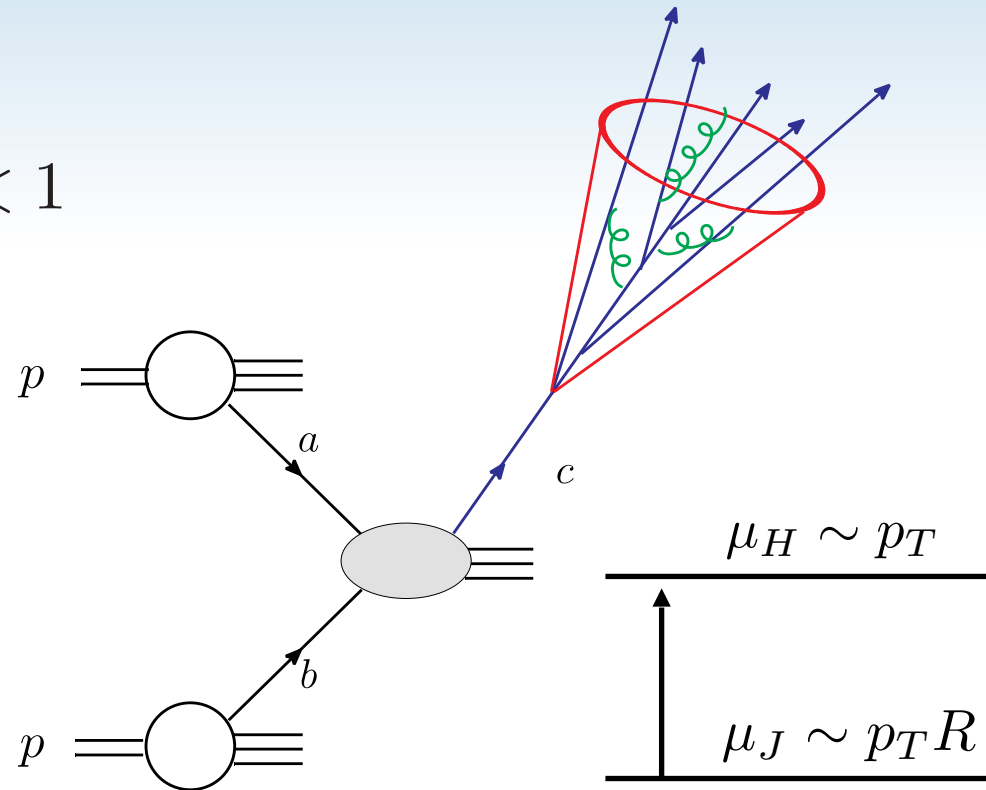
$$D_c^h \rightarrow J_c$$

$$pp \rightarrow hX$$

- Simple replacement of the fragmentation function by “semi-inclusive jet function”.

Comparison with the inclusive hadron production case

$$R \ll 1$$



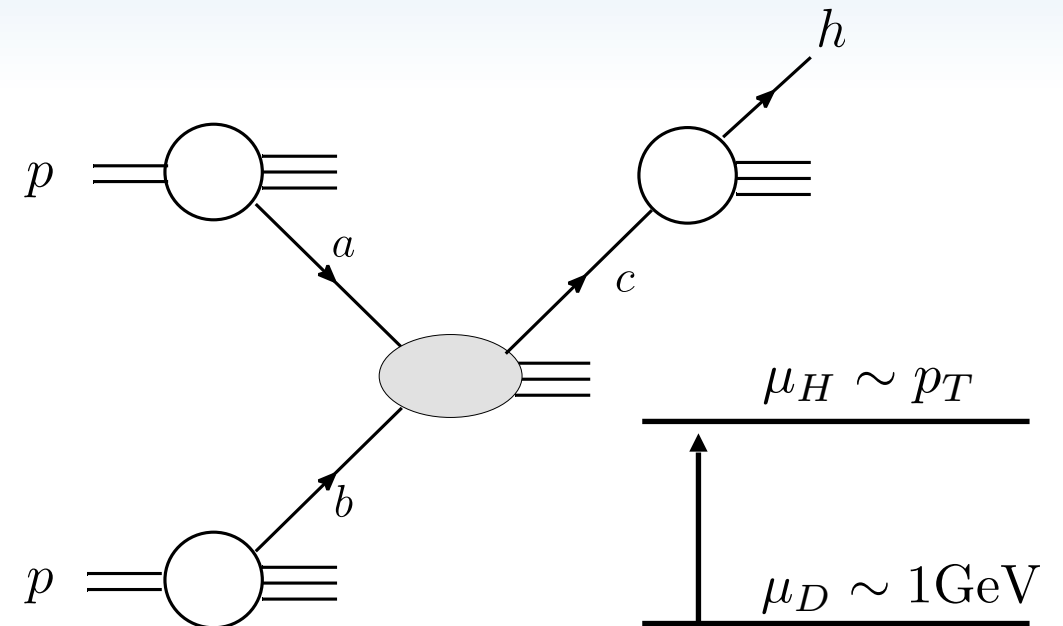
Factorization

Inclusive Jet

$$\frac{d\sigma^{pp \rightarrow \text{jet} X}}{dp_T d\eta} = \sum_{a,b,c} f_a \otimes f_b \otimes H_{ab}^c \otimes J_c + \mathcal{O}(R^2)$$

Hadron

$$\frac{d\sigma^{pp \rightarrow h X}}{dp_T d\eta} = \sum_{a,b,c} f_a \otimes f_b \otimes H_{ab}^c \otimes D_c^h$$

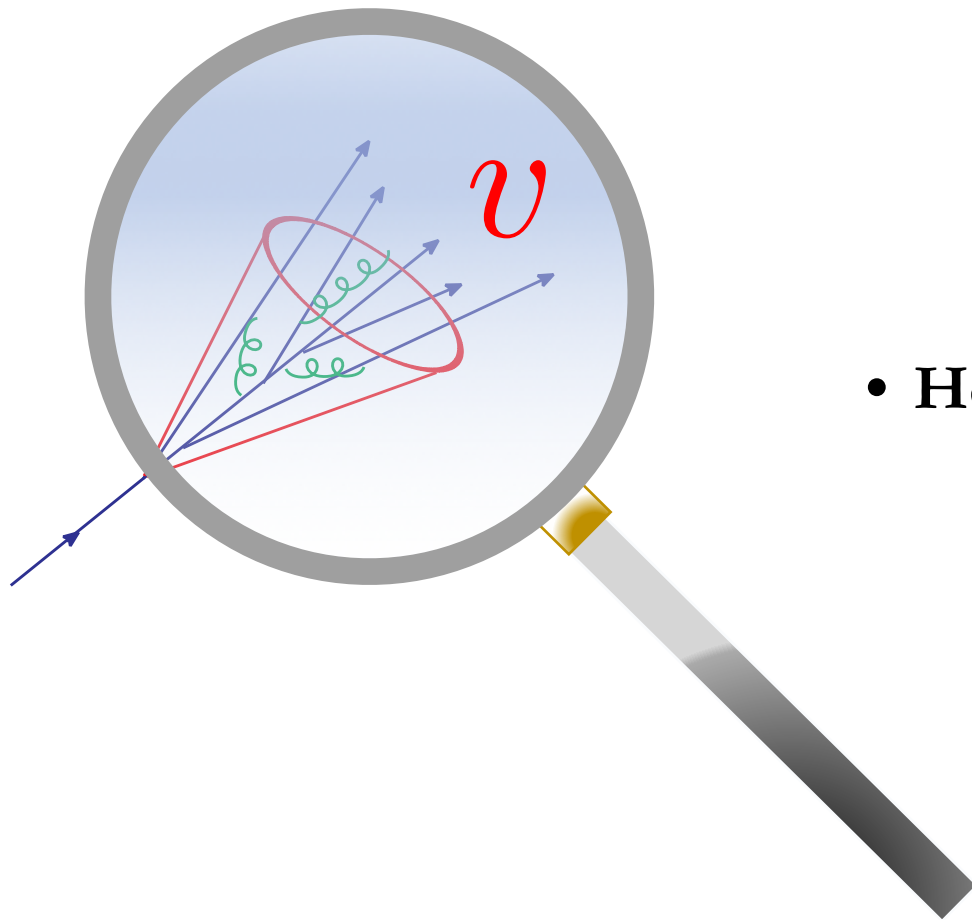


Evolution

$$\mu \frac{d}{d\mu} J_i = \sum_j P_{ji} \otimes J_j$$

$$\mu \frac{d}{d\mu} D_i^h = \sum_j P_{ji} \otimes D_j^h$$

Jet Substructure Measurements



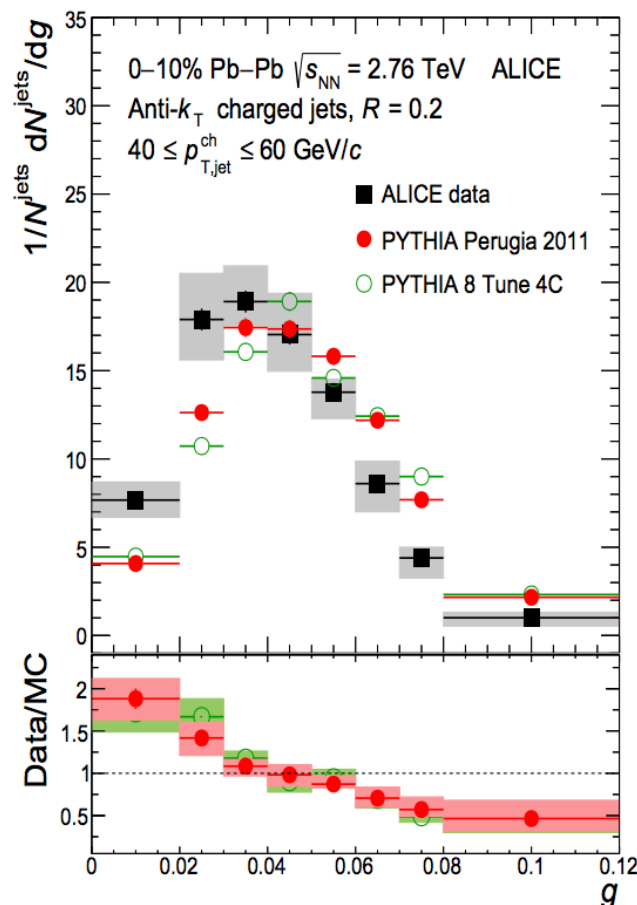
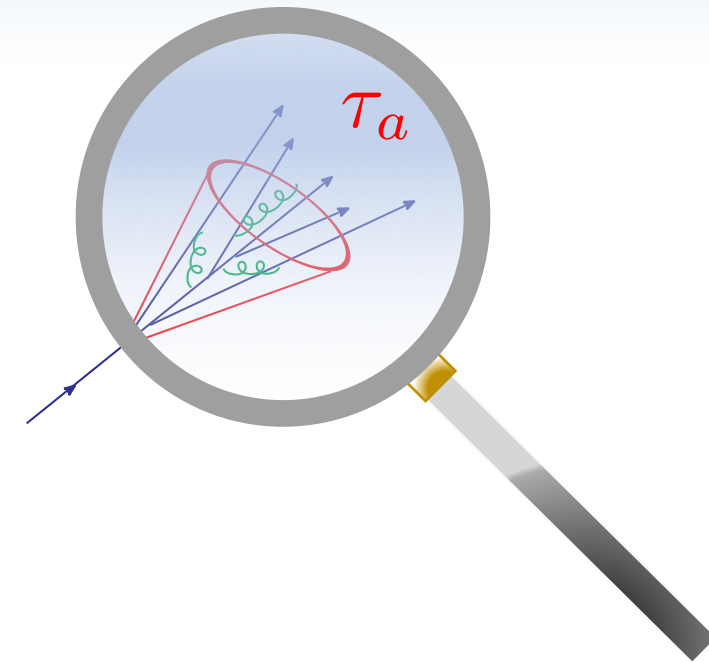
- How do we measure a substructure v inside the jet?

Jet angularity

- A generalized class of IR safe observables ($-\infty < a < 2$), angularity (applied to jet):

$$\tau_a^{pp} = \frac{1}{p_T} \sum_{i \in J} p_{T,i} (\Delta R_{iJ})^{2-a}$$

$$\tau_0^{pp} = \frac{m_J^2}{p_T^2} + \mathcal{O}((\tau_0^{pp})^2)$$

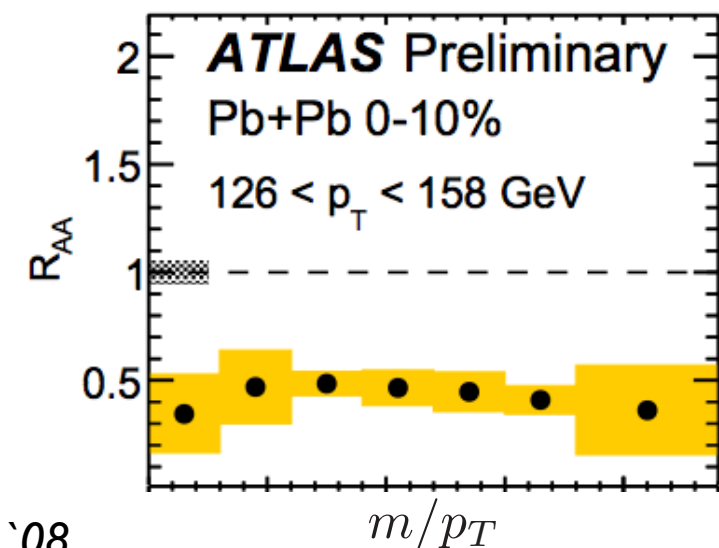


Medium modifications

$$a = 1$$

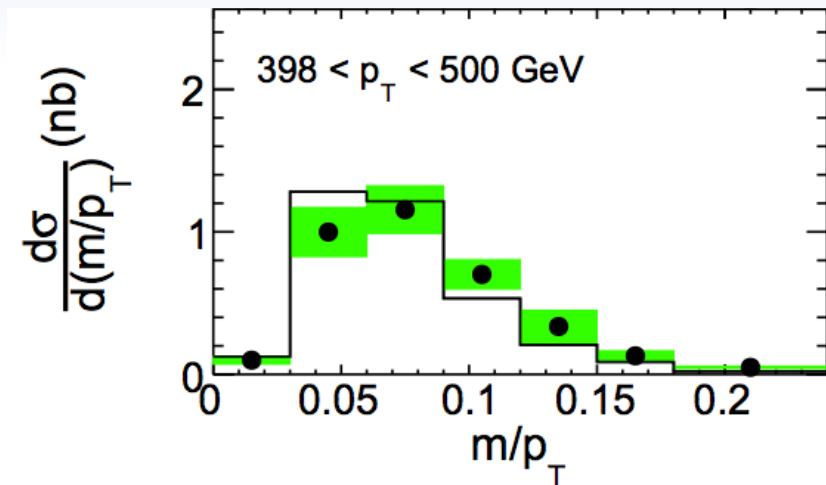
$$g(\text{girth}) = \frac{1}{p_T} \sum_{i \in J} p_{T,i} (\Delta R_{iJ})$$

$$a = 0$$

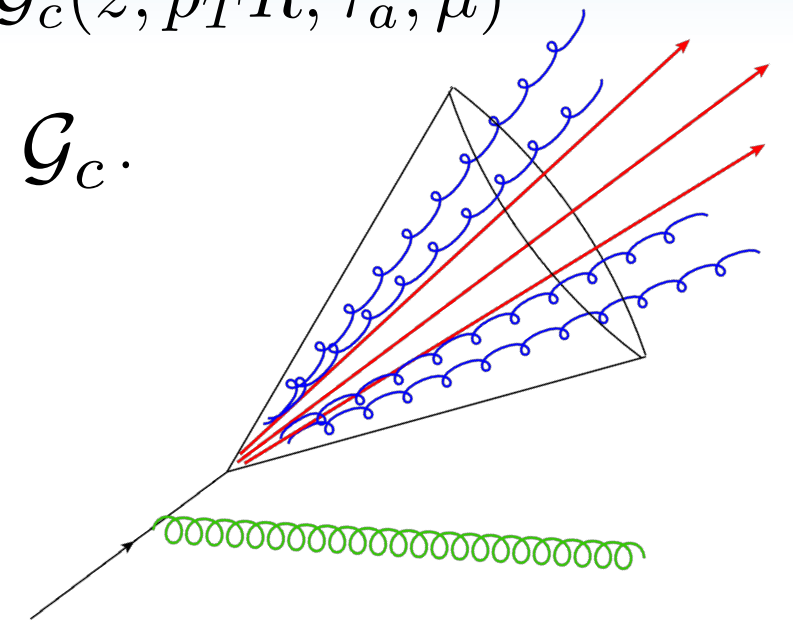


Sterman et al. '03, '08,
 Hornig, C. Lee, Ovanesyan '09, Ellis, Vermilion, Walsh, Hornig, C. Lee '10,
 Chien, Hornig, C. Lee '15, Hornig, Makris, Mehen '16, Kang, KL, Ringer '18

Factorization for jet angularity



- Replace $J_c(z, p_T R, \mu) \rightarrow \mathcal{G}_c(z, p_T R, \tau_a, \mu)$
- When $\tau_a \ll R^2$, refactorize $\mathcal{G}_c$.



Relevant modes for $\tau_a \ll R^2$

$$\tau_a \sim z \theta^{2-a}$$

Collinear

$$z_c \sim 1$$

$$\theta_c \sim \tau_a^{\frac{1}{2-a}}$$

Collinear-soft

$$\theta_s \sim R$$

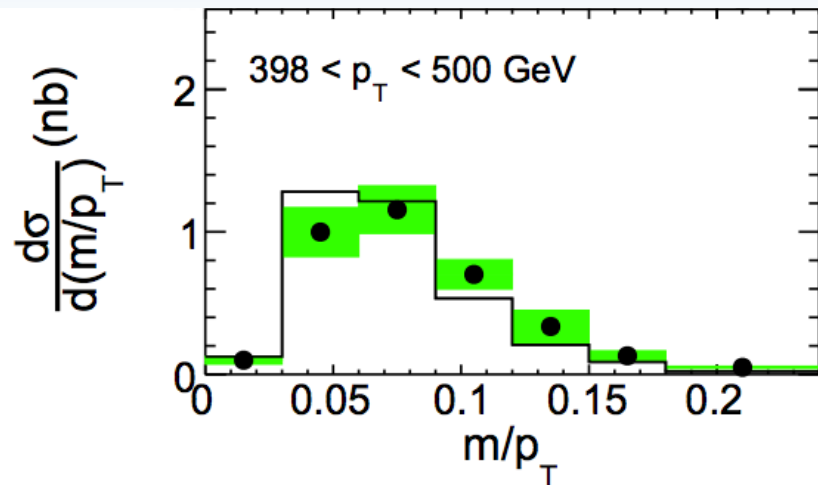
$$z_{cs} \sim \frac{\tau_a}{R^{2-a}}$$

Hard-collinear

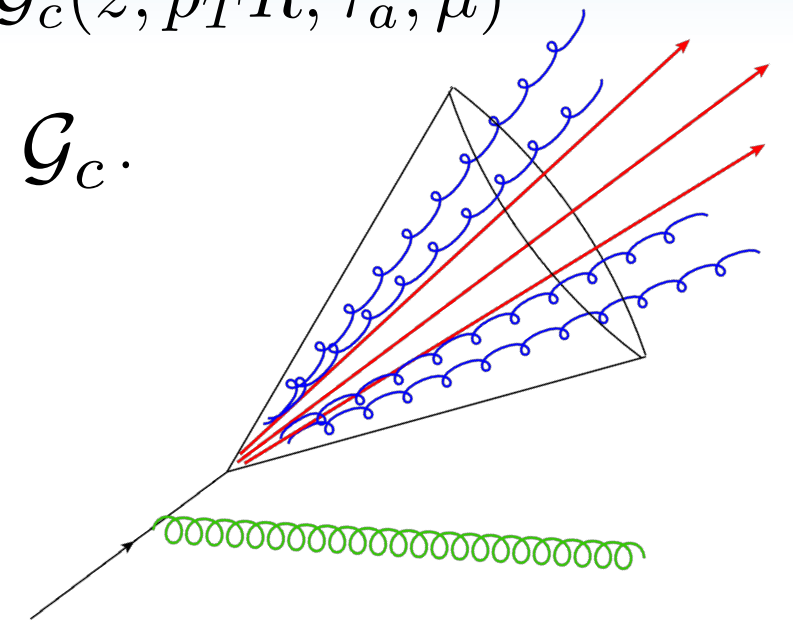
$$\theta_H \sim R$$

$$z_H \sim 1$$

Factorization for jet angularity



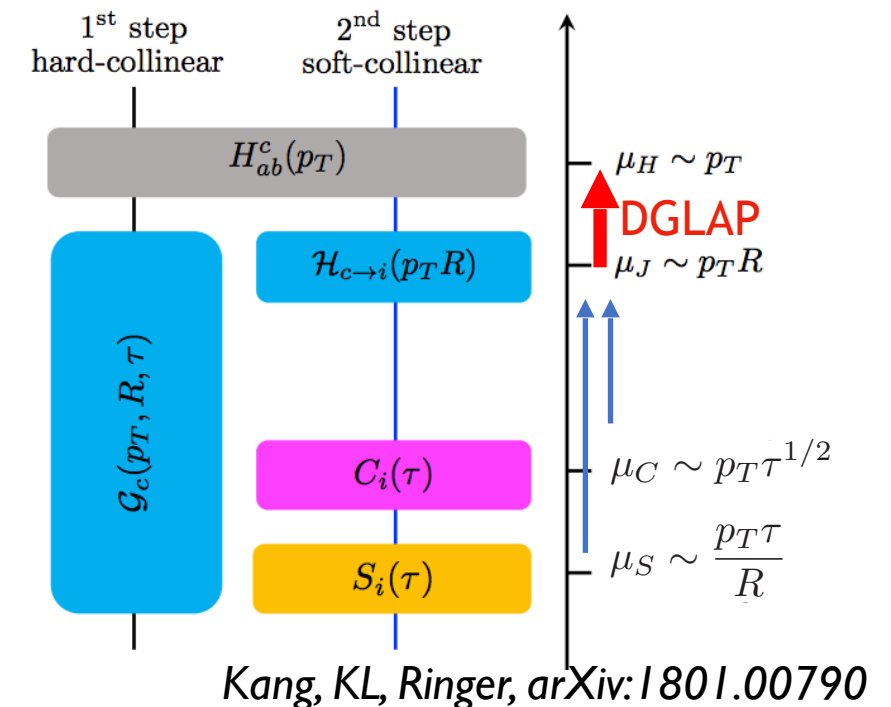
- Replace $J_c(z, p_T R, \mu) \rightarrow \mathcal{G}_c(z, p_T R, \tau_a, \mu)$
- When $\tau_a \ll R^2$, refactorize $\mathcal{G}_c$.



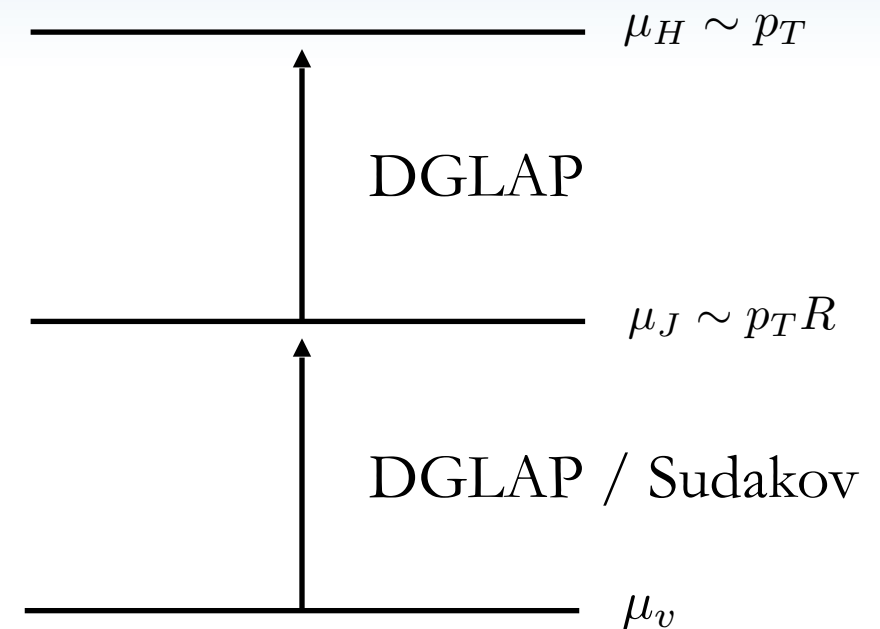
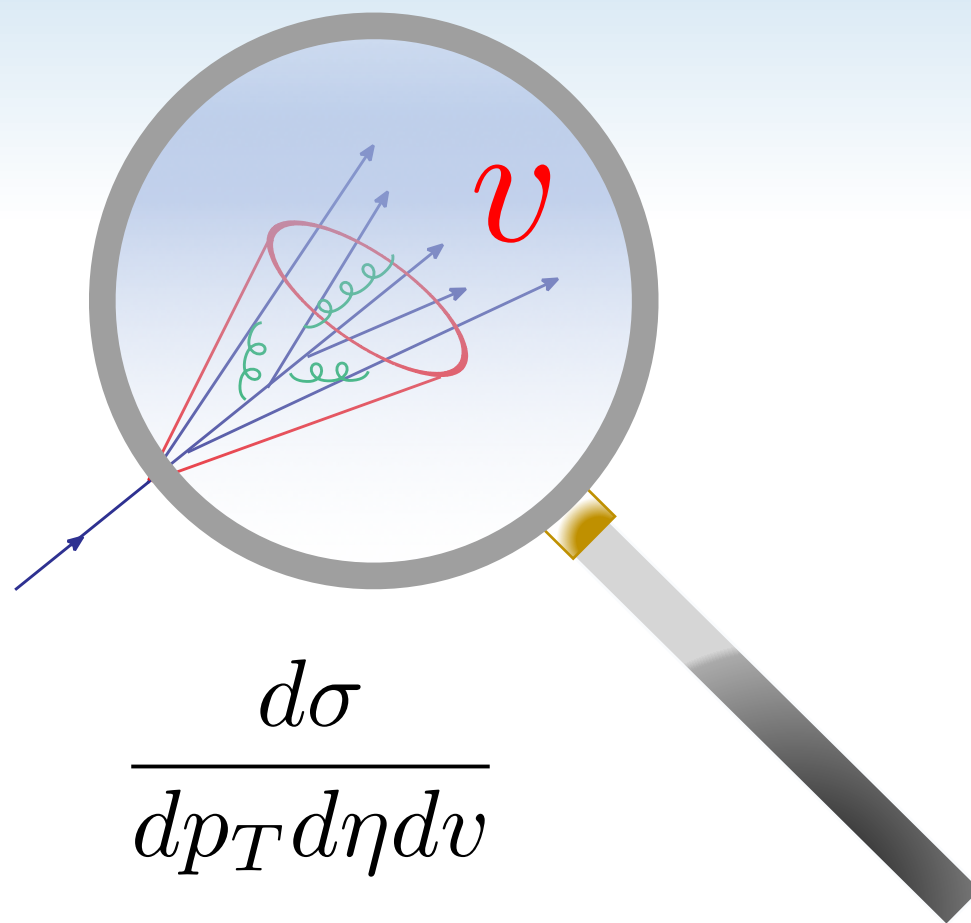
$$\mathcal{G}_c(z, p_T R, \tau_a, \mu) = \sum_i \mathcal{H}_{c \rightarrow i}(z, p_T R, \mu)$$

$$\times \int d\tau_a^{C_i} d\tau_a^{S_i} \delta(\tau_a - \tau_a^{C_i} - \tau_a^{S_i}) \mathcal{C}_i(\tau_a^{C_i}, p_T \tau_a^{\frac{1}{2-a}}, \mu) \mathcal{S}_i(\tau_a^{S_i}, \frac{p_T \tau_a}{R^{1-a}}, \mu) + \mathcal{O}(\tau_a/R^2)$$

- Each piece describes physics at different scales.
- $\mu_J \rightarrow \mu_H$ evolution follows DGLAP evolution equation again
- Resums $(\alpha_s \ln R)^n$ and $\left(\alpha_s \ln^2 \frac{R}{\tau^{1/2}}\right)^n$



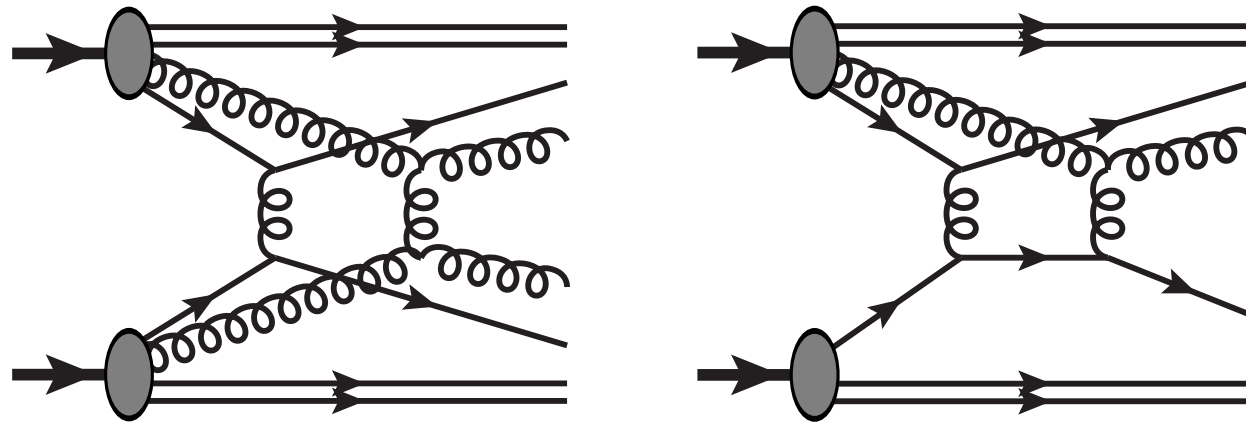
Patterns emerging



- When we measure a substructure v from the jet, once we evolve to μ_J the remaining evolution to μ_H is given by DGLAP evolution!
- Two step factorization:
 - a) production of a jet
 - b) probing the internal structure of the jet produced.

Non-perturbative Effects

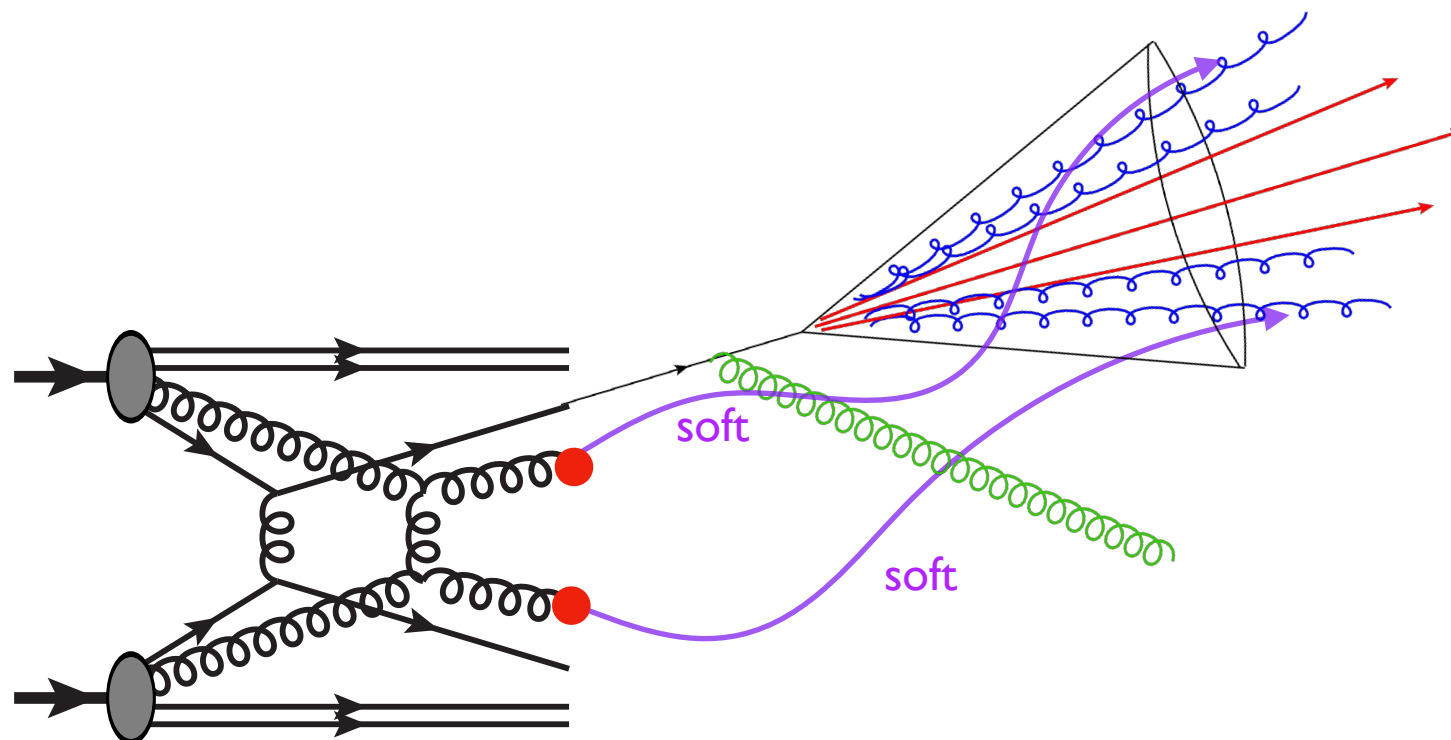
- Non-perturbative effects:



Figs from P. Bartalini et al. '11

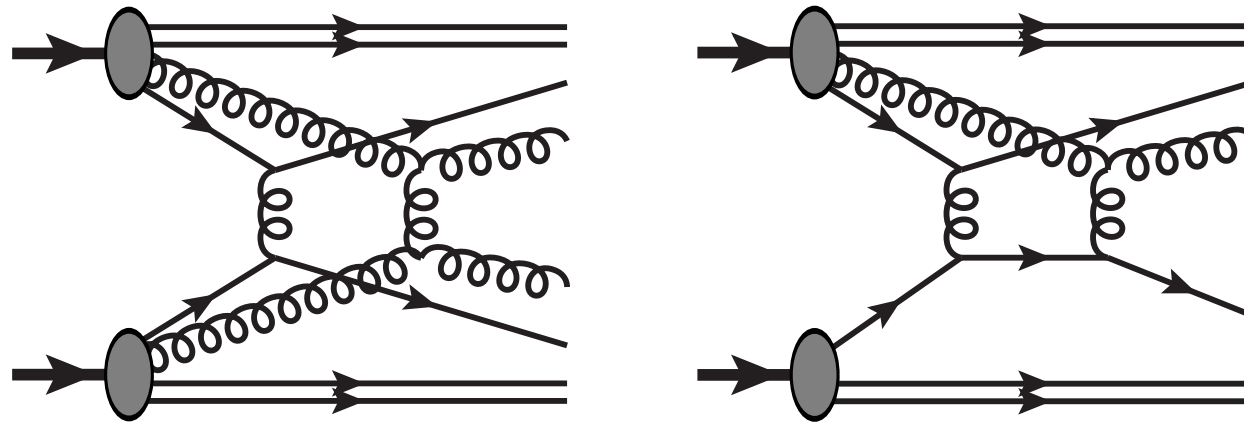
- **Multi-Parton Interactions (MPI) (Underlying Events (UE))**

Multiple secondary scatterings of partons within the protons may enter and contaminate jet.

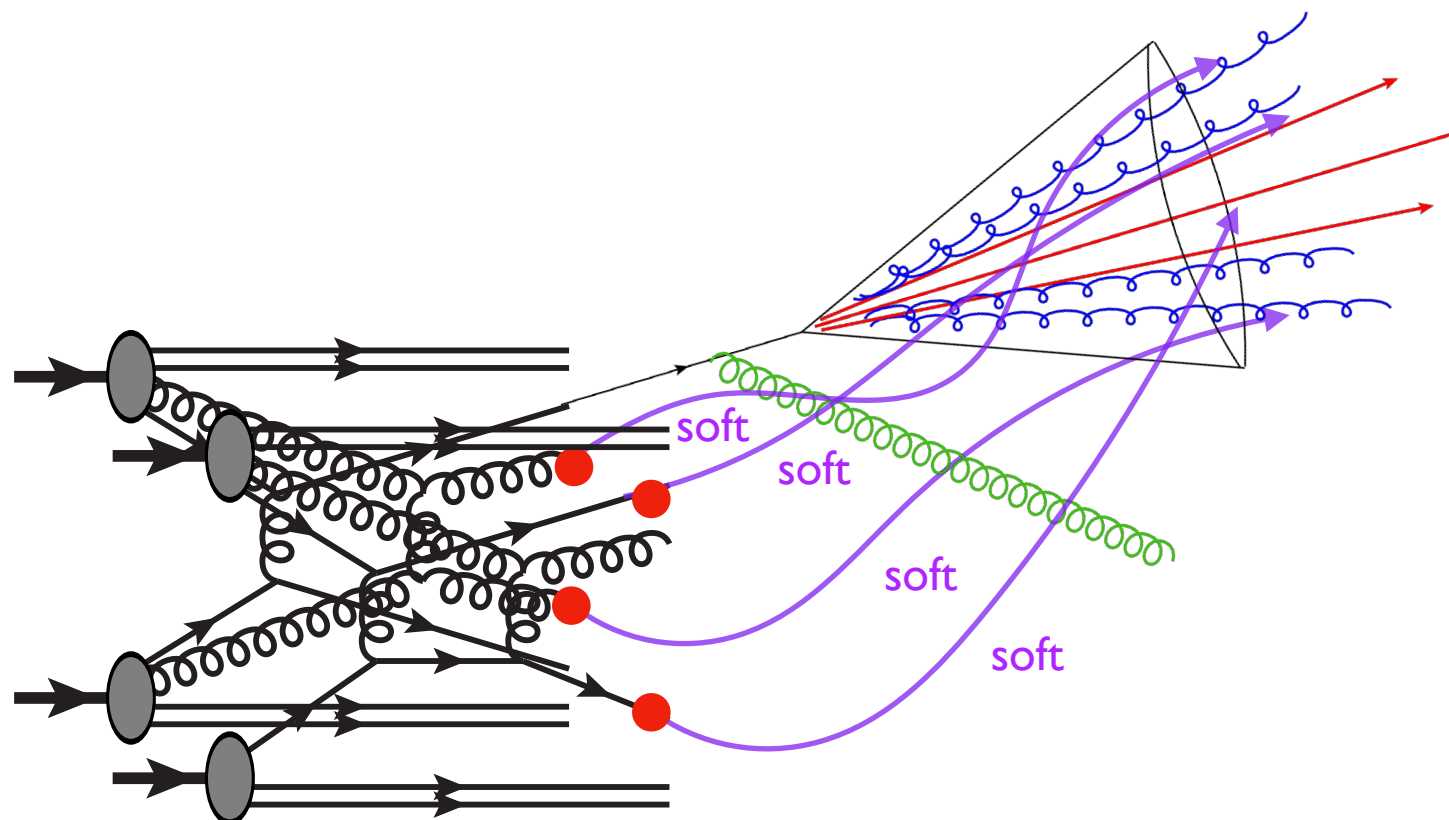


Non-perturbative Effects

- Non-perturbative effects:



Figs from P. Bartalini et al. '11



- **Multi-Parton Interactions (MPI) (Underlying Events (UE))**

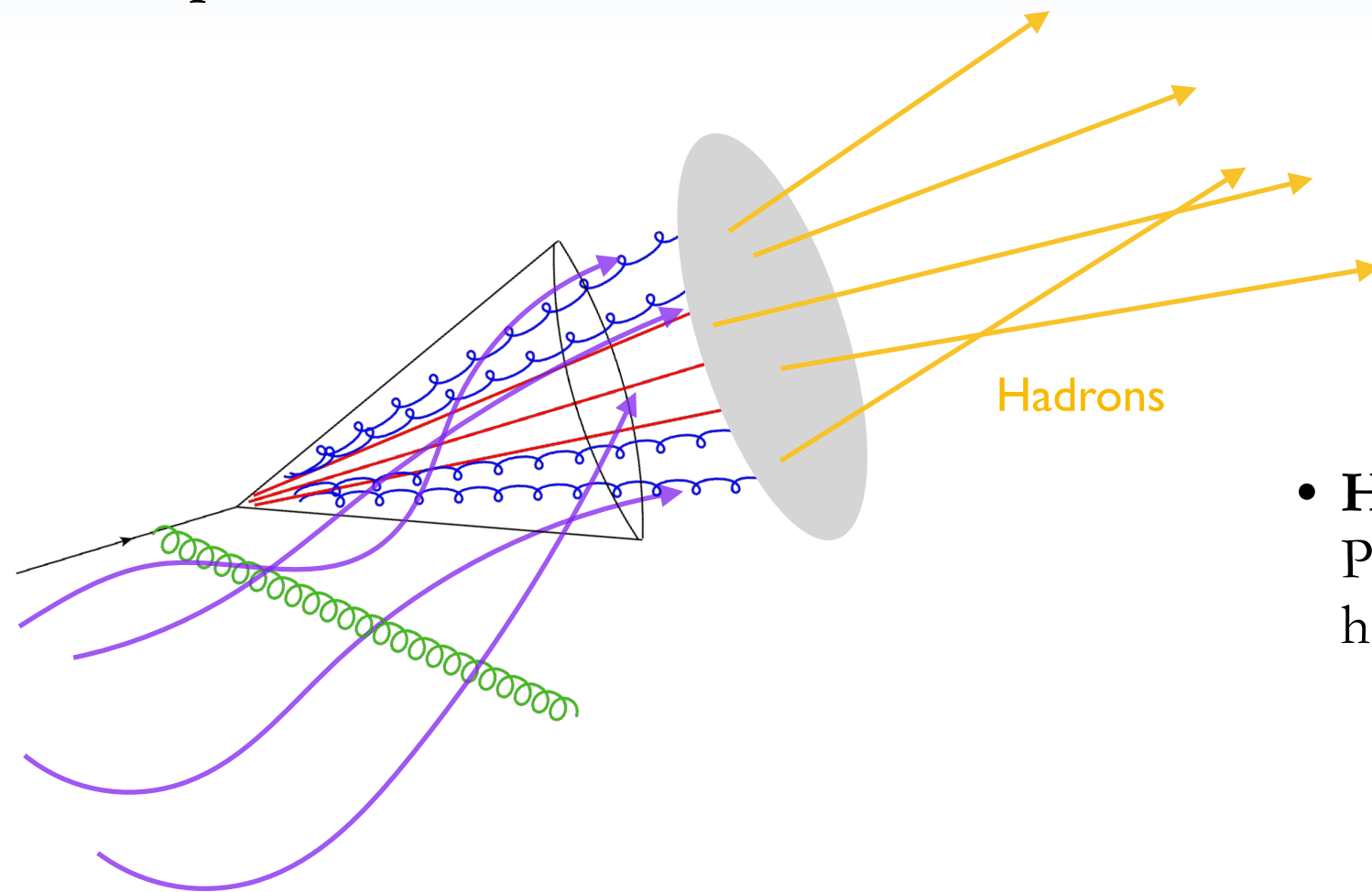
Multiple secondary scatterings of partons within the protons may enter and contaminate jet.

- **Pileups**

Secondary proton collisions in a bunch may enter and contaminate jet.

Non-perturbative Effects

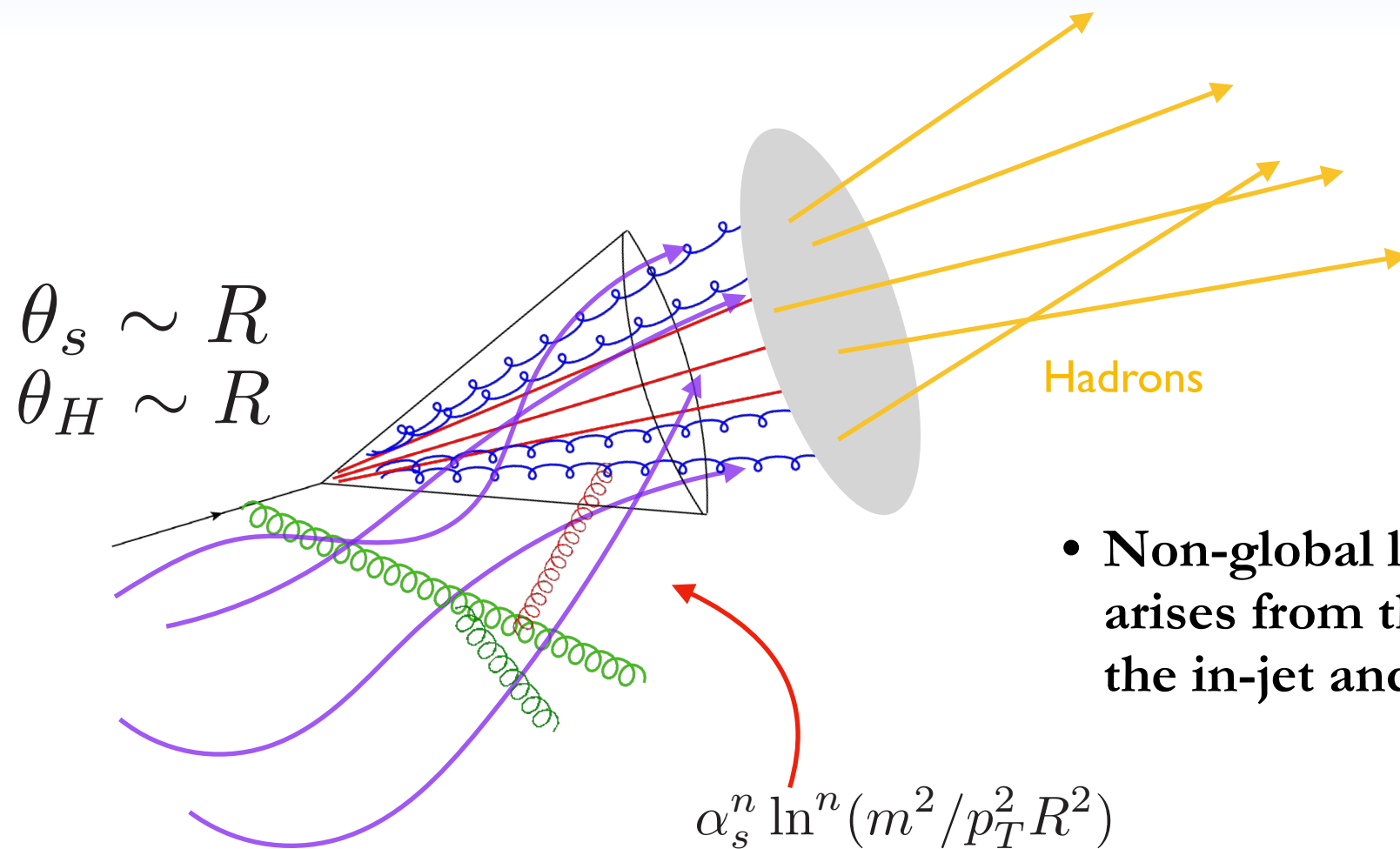
- Non-perturbative effects:



- **Hadronization**
Partons forming the jet eventually hadronizes.

Non-global logarithms

Dasgupta, Salam '01
 Banfi, Marchesini, Smye '02
 Larkoski, Moulton, Neill '15
 Becher, Neubert, Rothen, Shao '15, '16 ...



- **Non-global logarithms (NGLs):** arises from the correlation between the in-jet and out-of-jet radiation.

rather small effect for jet mass

Non-perturbative Model

- As τ gets smaller, $\mu_S \sim \frac{p_T \tau}{R}$ (smallest scale) can approach a non-perturbative scale.

We shift our perturbative results by convolving with non-perturbative shape function to smear

$$\frac{d\sigma}{d\eta dp_T d\tau} = \int dk F_\kappa(k) \frac{d\sigma^{\text{pert}}}{d\eta dp_T d\tau} \left(\tau - \frac{R}{p_T} k \right)$$

- Single parameter NP soft function :

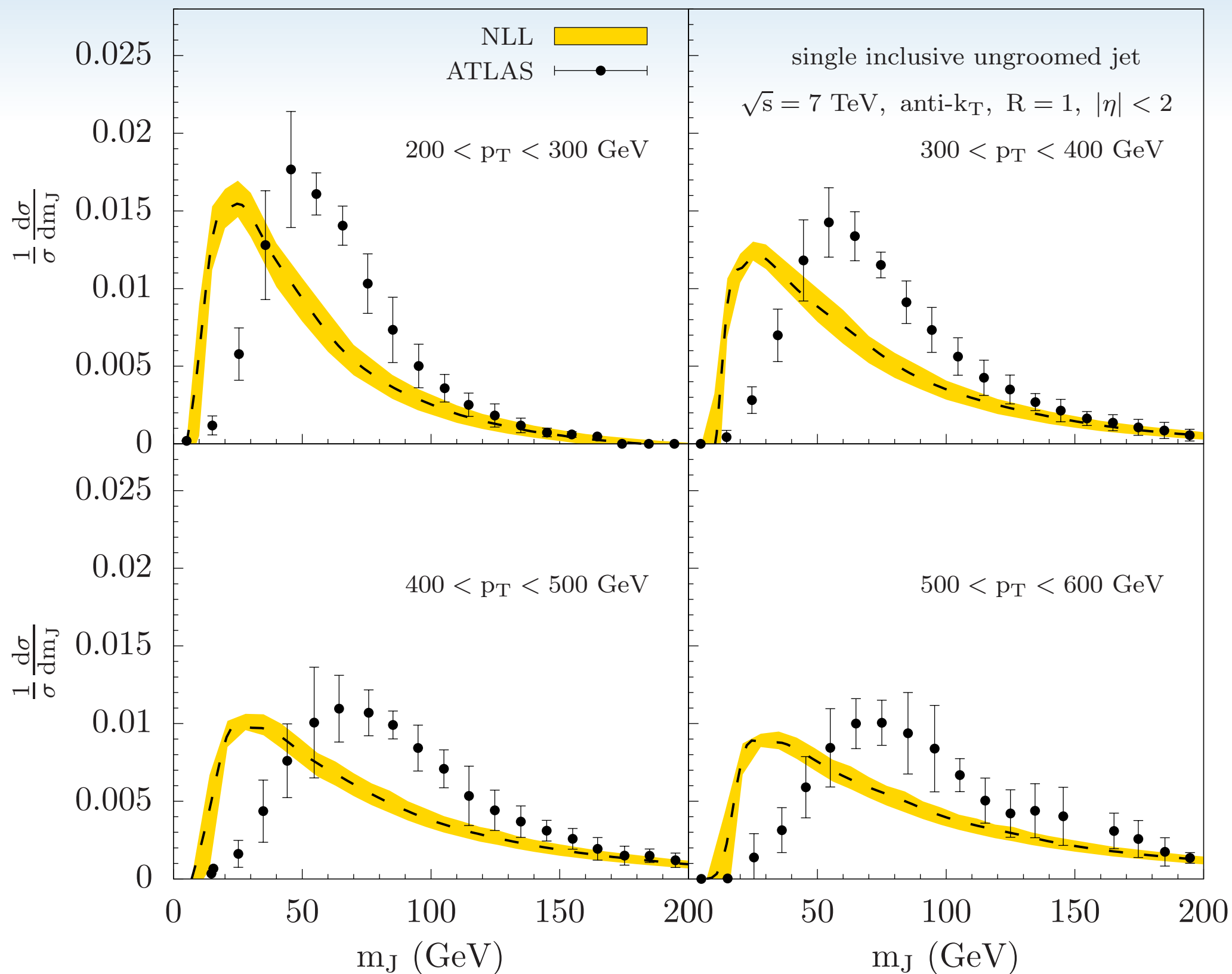
$$F_\kappa(k) = \left(\frac{4k}{\Omega_\kappa^2} \right) \exp \left(-\frac{2k}{\Omega_\kappa} \right) \quad \text{Stewart, Tackmann, Waalewijn '15}$$

- Both hadronization and MPI effects in jet mass is well-represented by just shifting first-moments.
- The parameter Ω_κ is related to shift in the distribution:

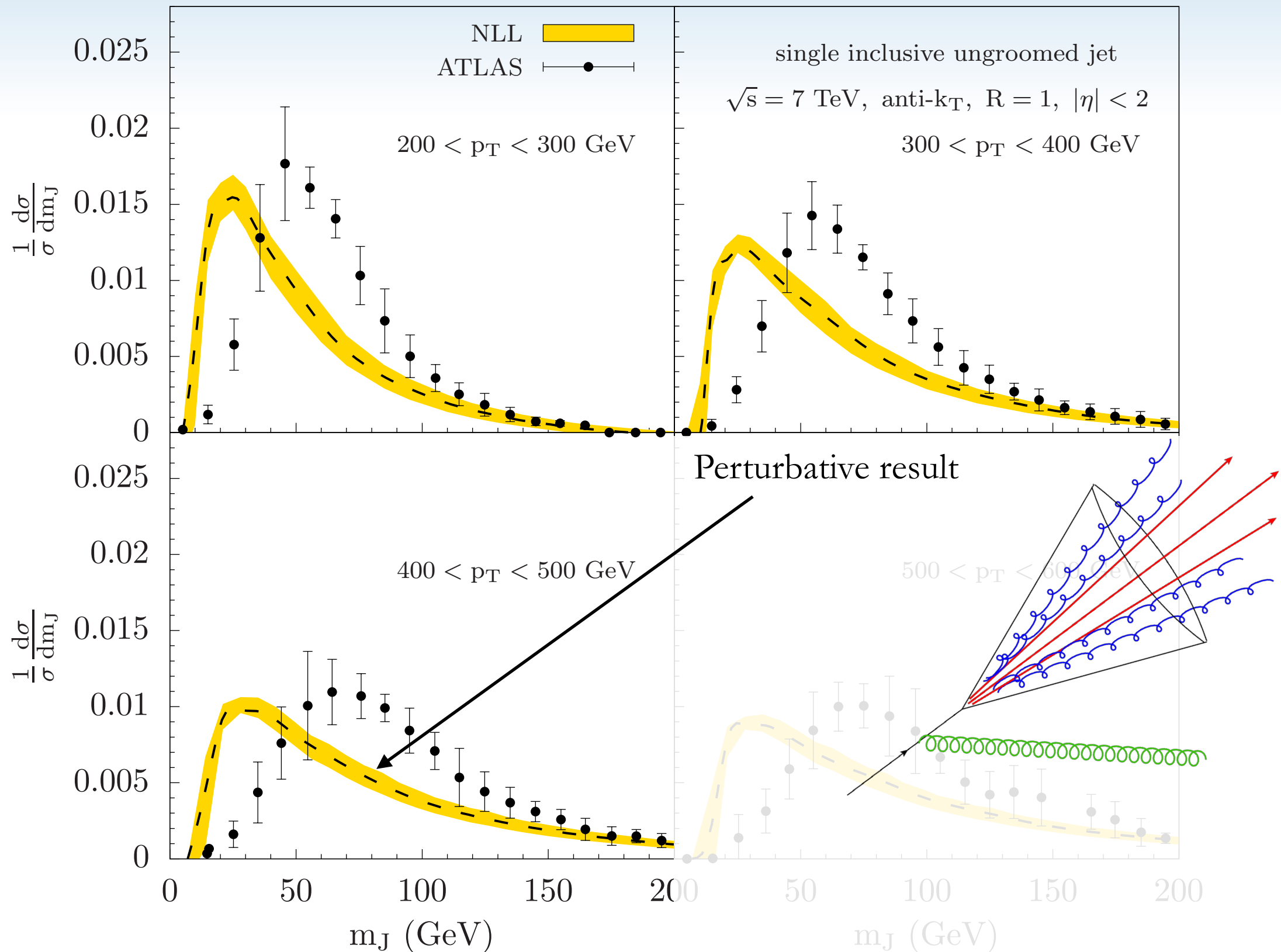
$$\tau = \tau_{\text{pert}} + \tau_{\text{NP}} = \tau_{\text{pert}} + \frac{R\Omega_\kappa}{p_T} = \tau_{\text{pert}} + \frac{R(\Lambda_{\text{hadro.}} + \Lambda_{\text{MPI}})}{p_T}$$

$\Omega_\kappa \sim \Lambda_{\text{had}} \sim 1 \text{ GeV}$ corresponds to non-perturbative effects coming primarily from the hadronization alone.

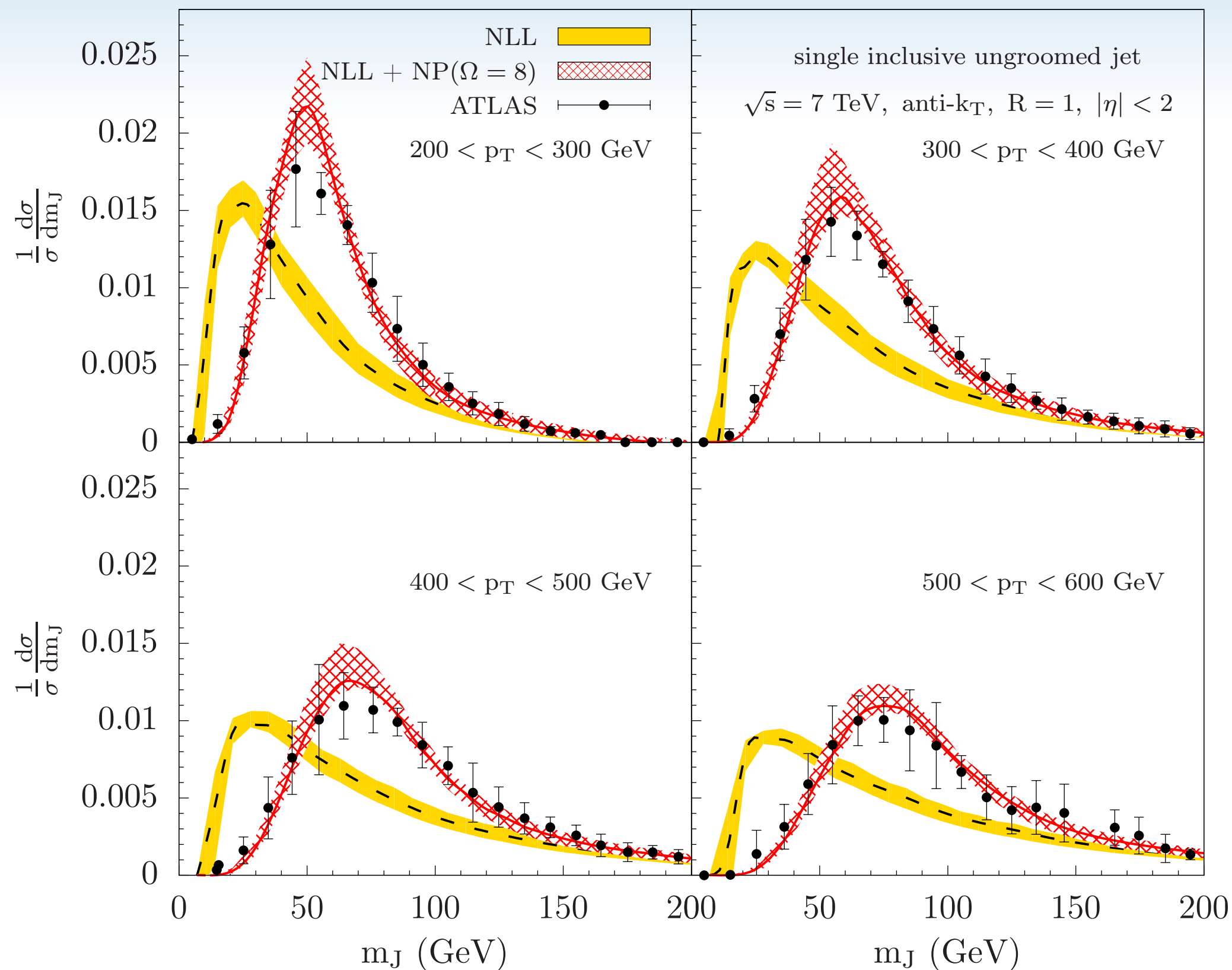
Phenomenology



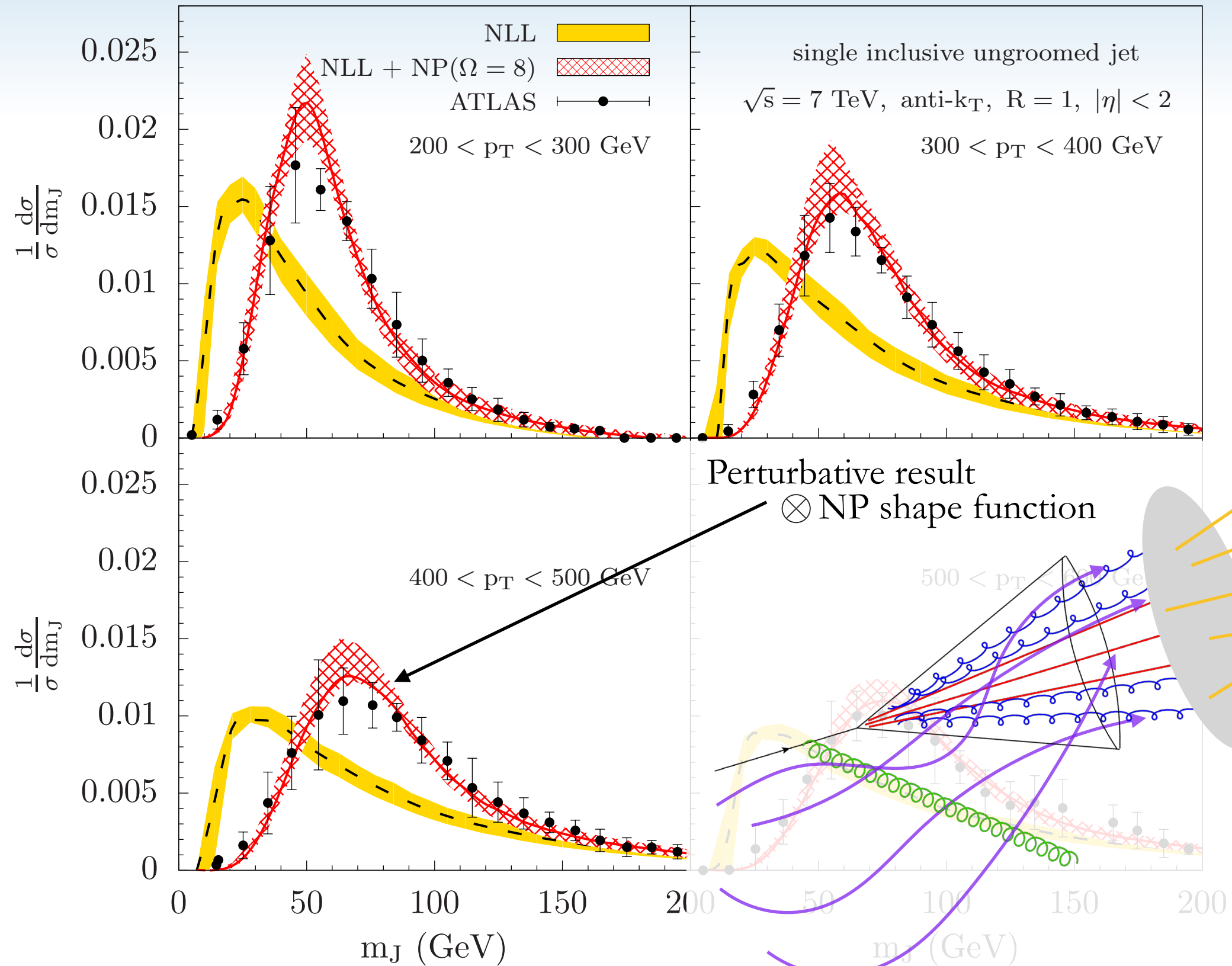
Phenomenology



Phenomenology



Phenomenology



Getting a better hold of MPI

- Underlying Events (UE) are difficult to understand.

How do we get a better hold of these contaminations in the jet?

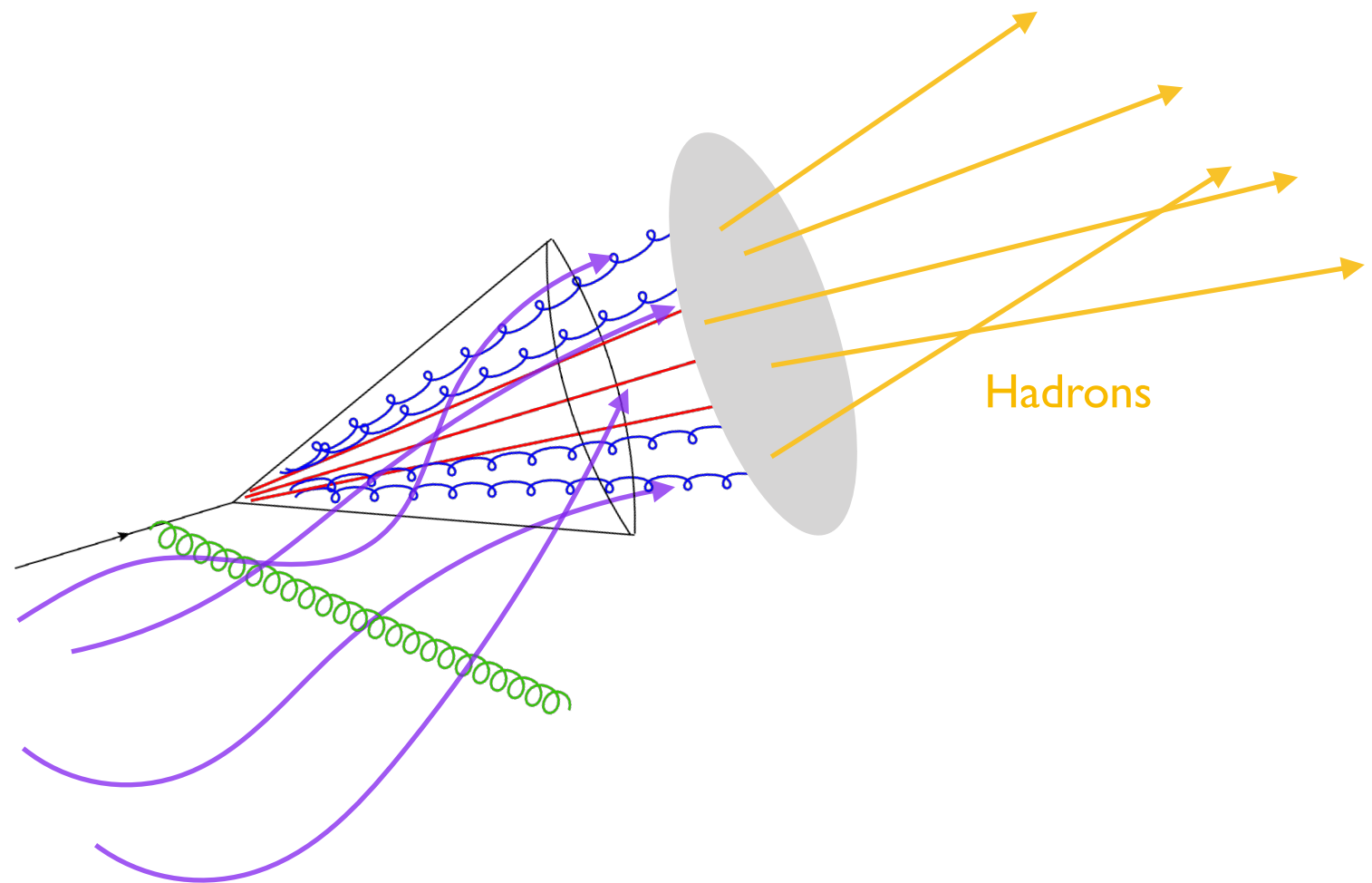
- Hint : contamination generally from soft radiations.
 1. Subtracted jet mass moments (See Yiannis's talk)
Chien, Kang, KL, Makris, arXiv:1812.06977
 2. Grooming

Soft Drop Grooming

- Underlying Events (UE) are difficult to understand.

How do we get a better hold of these contaminations in the jet?

- Hint : contamination generally from soft radiations.



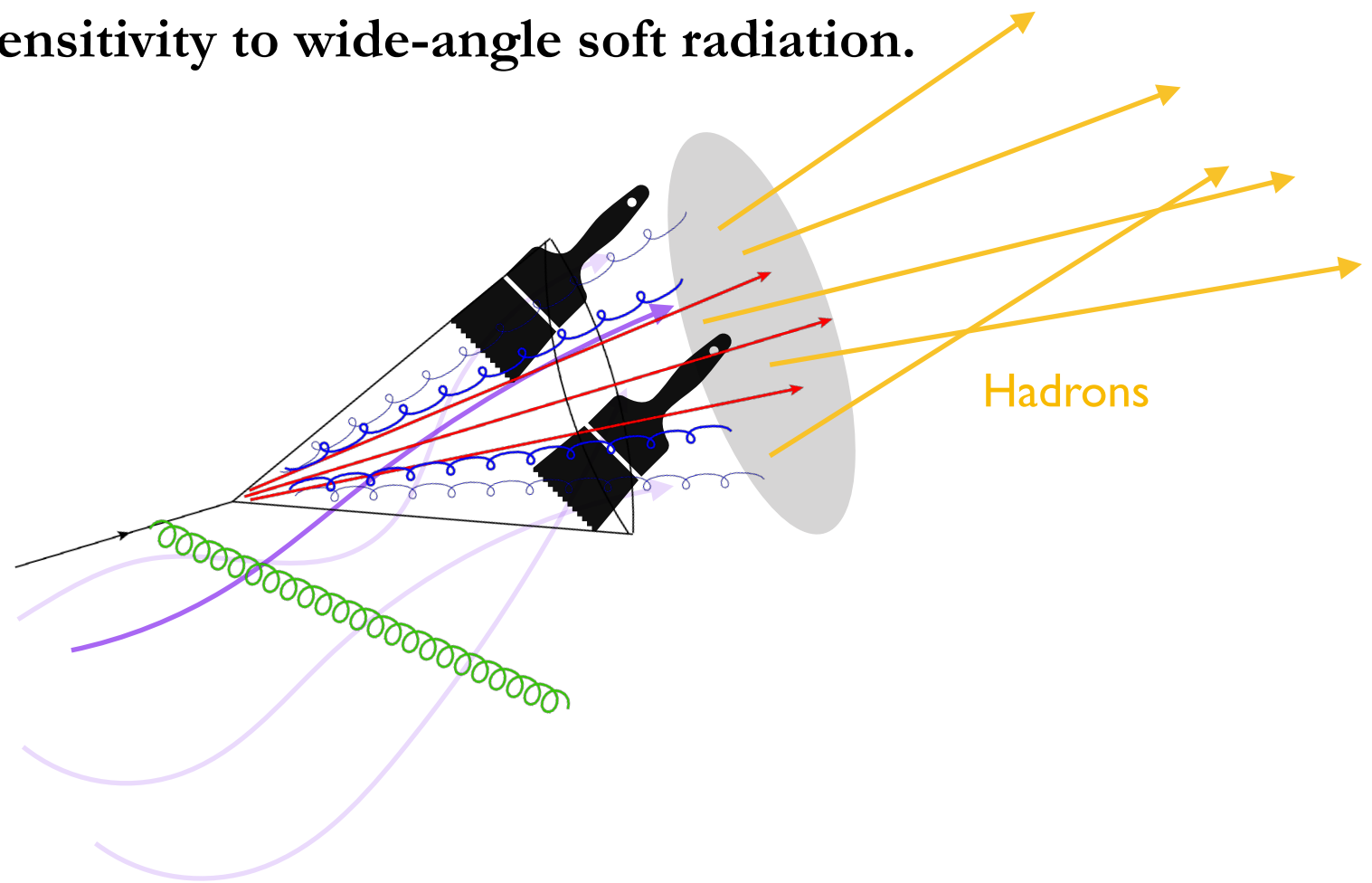
Soft Drop Grooming

- Underlying Events (UE) are difficult to understand.

How do we get a better hold of these contaminations in the jet?

- Hint : contamination generally from soft radiations.

Groom jets to reduce sensitivity to wide-angle soft radiation.



Soft Drop Grooming

- Underlying Events (UE) are difficult to understand.

How do we get a better hold of these contaminations in the jet?

- Hint : contamination generally from soft radiations.

Groom jets to reduce sensitivity to wide-angle soft radiation.

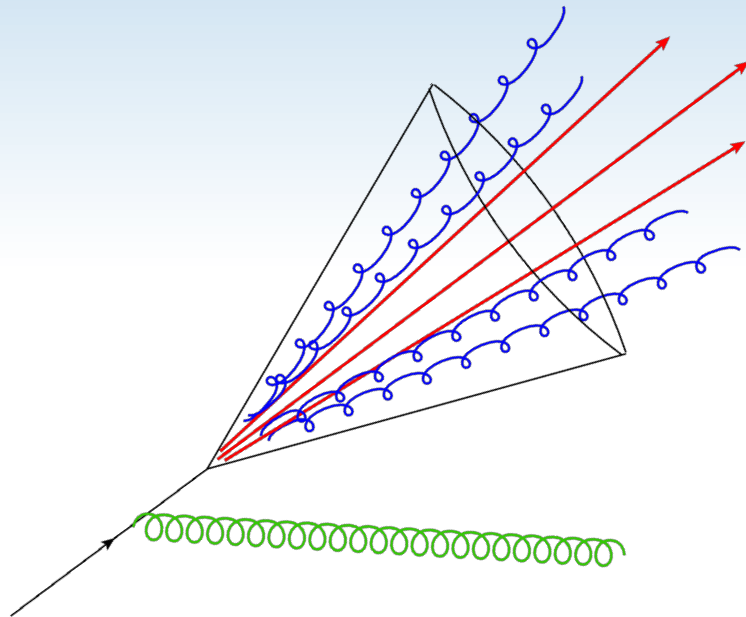


- Soft drop grooming algorithms:

- Reorder emissions in the identified jet according to their relative angle using C/A jet algorithm.
- Recursively remove soft branches until soft drop condition is met:

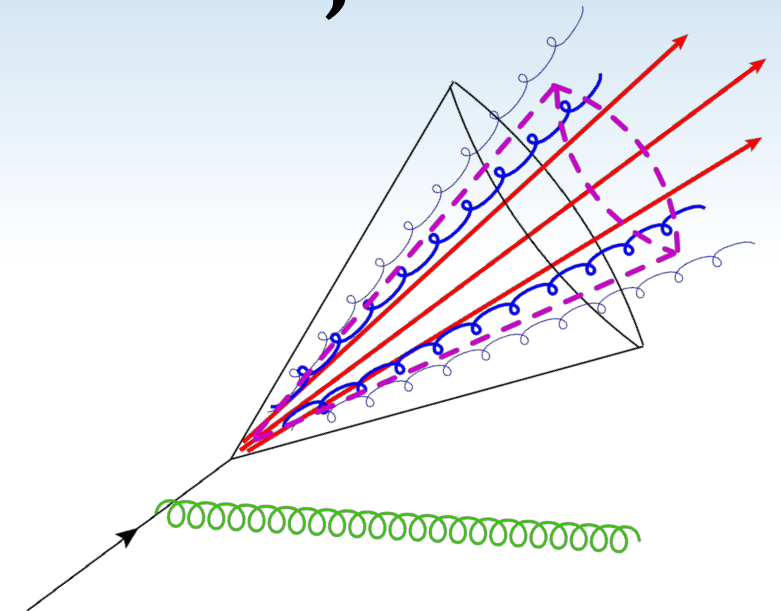
$$\frac{\min[p_{T,i}, p_{T,j}]}{p_{T,i} + p_{T,j}} > z_{\text{cut}} \left(\frac{R_{ij}}{R} \right)^\beta$$

Relevant modes in the groomed jet



$$\tau_a \sim z \theta^{2-a}$$

$$z > z_{\text{cut}} (\theta/R)^\beta$$



- The ungroomed case ($\tau \ll R^2$)

Hard-collinear

$$\theta_H \sim R \quad z_H \sim 1$$

Collinear

$$z_c \sim 1 \quad \theta_c \sim \tau_a^{\frac{1}{2-a}}$$

Collinear-soft

$$\theta_s \sim R \quad z_{cs} \sim \frac{\tau_a}{R^{2-a}}$$

- The groomed case ($\tau_{\text{gr}}/R^2 \ll z_{\text{cut}} \ll 1$)

Hard-collinear

$$\theta_H \sim R \quad z_H \sim 1$$

Collinear

$$z_c \sim 1 \quad \theta_c \sim \tau_a^{\frac{1}{2-a}}$$

$\notin gr \text{ soft}$

$$\theta_{\notin gr} \sim R \quad z_{\notin gr} \sim z_{\text{cut}} \left(\frac{\theta}{R} \right)^\beta = z_{\text{cut}}$$

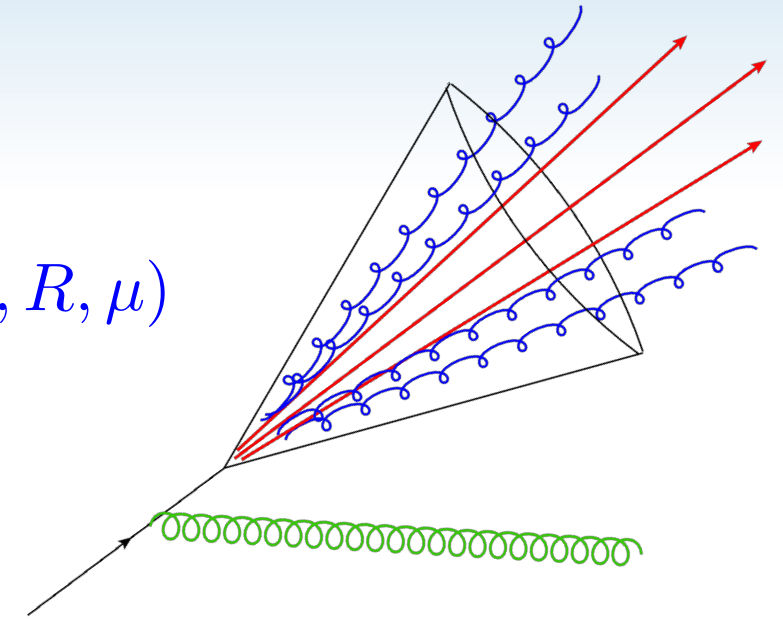
$\in gr \text{ soft}$

$$z_{\in gr} \sim z_{\text{cut}} \left(\frac{\theta}{R} \right)^\beta = z_{\text{cut}}^{\frac{2-a}{2-a+\beta}} \left(\frac{\tau_a}{R^{2-a}} \right)^{\frac{\beta}{2-a+\beta}} \quad \theta_{\in gr} \sim \left(\frac{\tau_a R^\beta}{z_{\text{cut}}} \right)^{\frac{1}{2-a+\beta}}$$

Groomed jet mass factorization

- The ungroomed case ($\tau \ll R^2$)

$$\mathcal{G}_i(z, p_T R, \tau, \mu) = \sum_j \mathcal{H}_{i \rightarrow j}(z, p_T R, \mu) C_j(\tau, p_T, \mu) \otimes S_j(\tau, p_T, R, \mu)$$

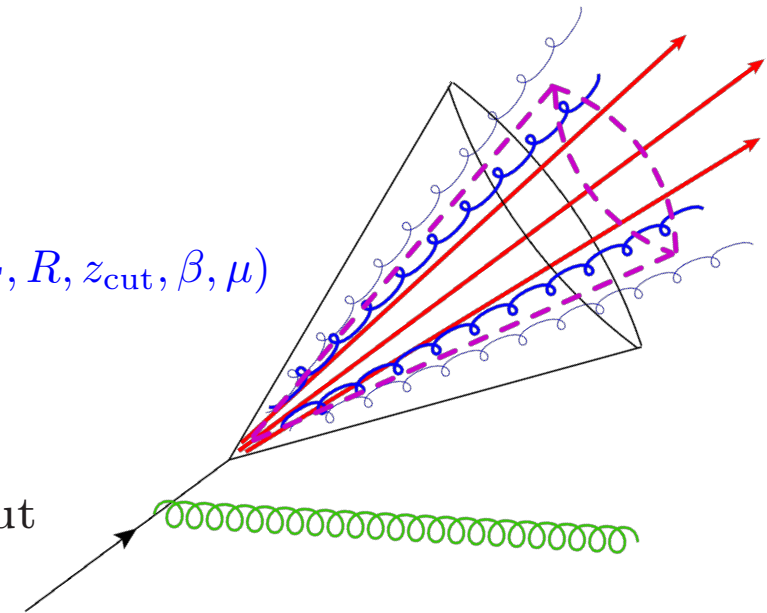


- Resums global logs $\alpha_s^n \ln^n R$ and $\alpha_s^n \ln^{2n} \tau / R^2$

- The groomed case ($\tau_{\text{gr}} / R^2 \ll z_{\text{cut}} \ll 1$)

$$\mathcal{G}_i(z, p_T R, \tau_{\text{gr}}, z_{\text{cut}}, \beta, \mu) = \sum_j \mathcal{H}_{i \rightarrow j}(z, p_T R, \mu) S_j^{\not\in \text{gr}}(p_T, R, z_{\text{cut}}, \beta, \mu) C_j(\tau, p_T, \mu) \otimes S_j^{\in \text{gr}}(\tau, p_T, R, z_{\text{cut}}, \beta, \mu)$$

- Resums global logs $\alpha_s^n \ln^n R$, $\alpha_s^n \ln^{2n} \tau / R^2$, and $\alpha_s^n \ln^{2n} z_{\text{cut}}$



Non-global Logarithms

Dasgupta, Salam '01 and many more

- The ungroomed case ($\tau \ll R^2$)

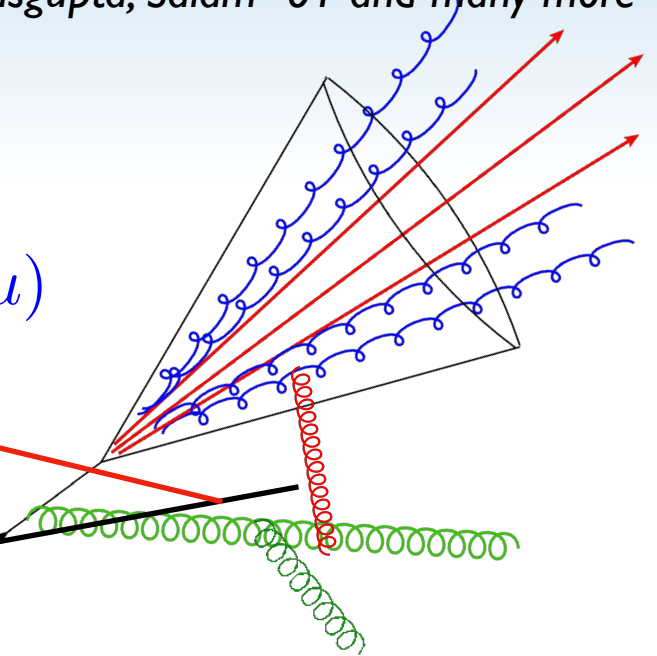
$$\mathcal{G}_i(z, p_T R, \tau, \mu) = \sum_j \mathcal{H}_{i \rightarrow j}(z, p_T R, \mu) C_j(\tau, p_T, \mu) \otimes S_j(\tau, p_T, R, \mu)$$

$$\theta_H \sim R$$

$$\theta_s \sim R$$

- Non-global logs directly affect the jet mass spectrum.

$$\alpha_s^n \ln^n(\tau/R^2) \quad n \geq 2$$



Non-global Logarithms

Dasgupta, Salam '01 and many more

- The ungroomed case ($\tau \ll R^2$)

$$\mathcal{G}_i(z, p_T R, \tau, \mu) = \sum_j \mathcal{H}_{i \rightarrow j}(z, p_T R, \mu) C_j(\tau, p_T, \mu) \otimes S_j(\tau, p_T, R, \mu)$$

$$\theta_H \sim R$$

$$\theta_s \sim R$$

- Non-global logs directly affect the jet mass spectrum.

$$\alpha_s^n \ln^n(\tau/R^2) \quad n \geq 2$$

- The groomed case ($\tau_{\text{gr}}/R^2 \ll z_{\text{cut}} \ll 1$)

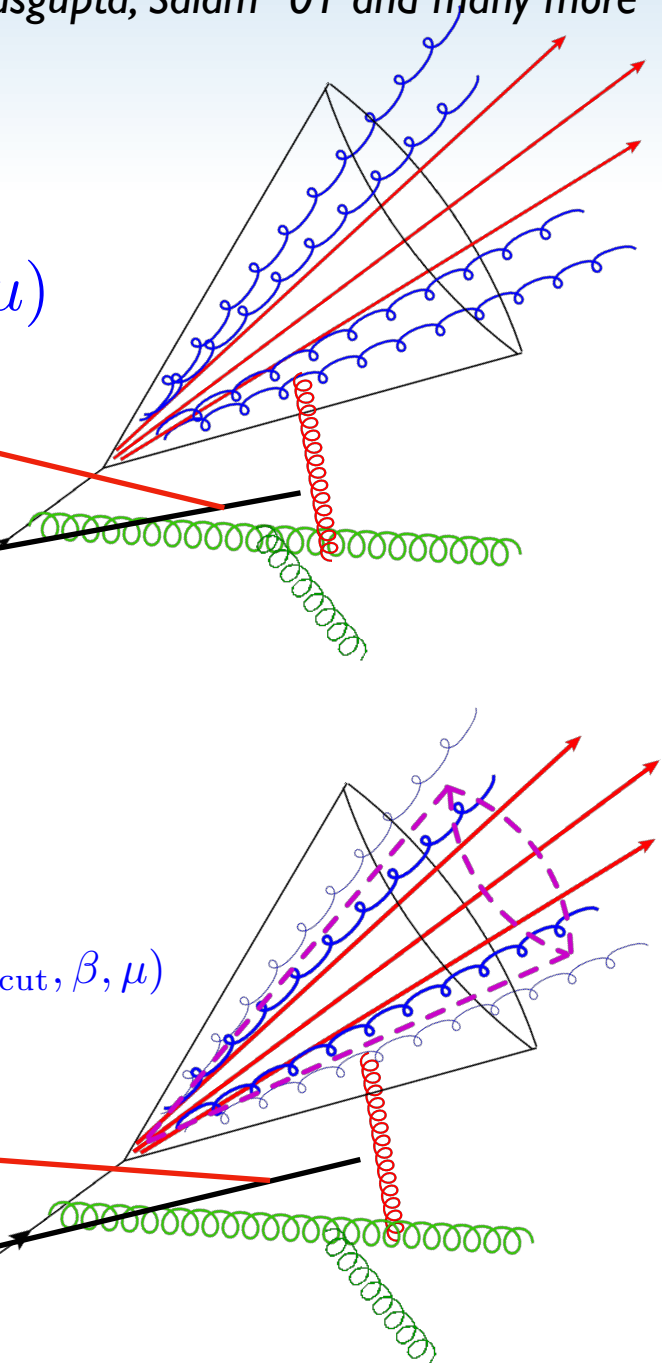
$$\mathcal{G}_i(z, p_T R, \tau_{\text{gr}}, z_{\text{cut}}, \beta, \mu) = \sum_j \mathcal{H}_{i \rightarrow j}(z, p_T R, \mu) S_j^{\not\in \text{gr}}(p_T, R, z_{\text{cut}}, \beta, \mu) C_j(\tau, p_T, \mu) \otimes S_j^{\in \text{gr}}(\tau, p_T, R, z_{\text{cut}}, \beta, \mu)$$

$$\theta_H \sim R$$

$$\theta_{\not\in \text{gr}} \sim R$$

- Non-global logs only indirectly affects the jet mass spectrum through normalization.

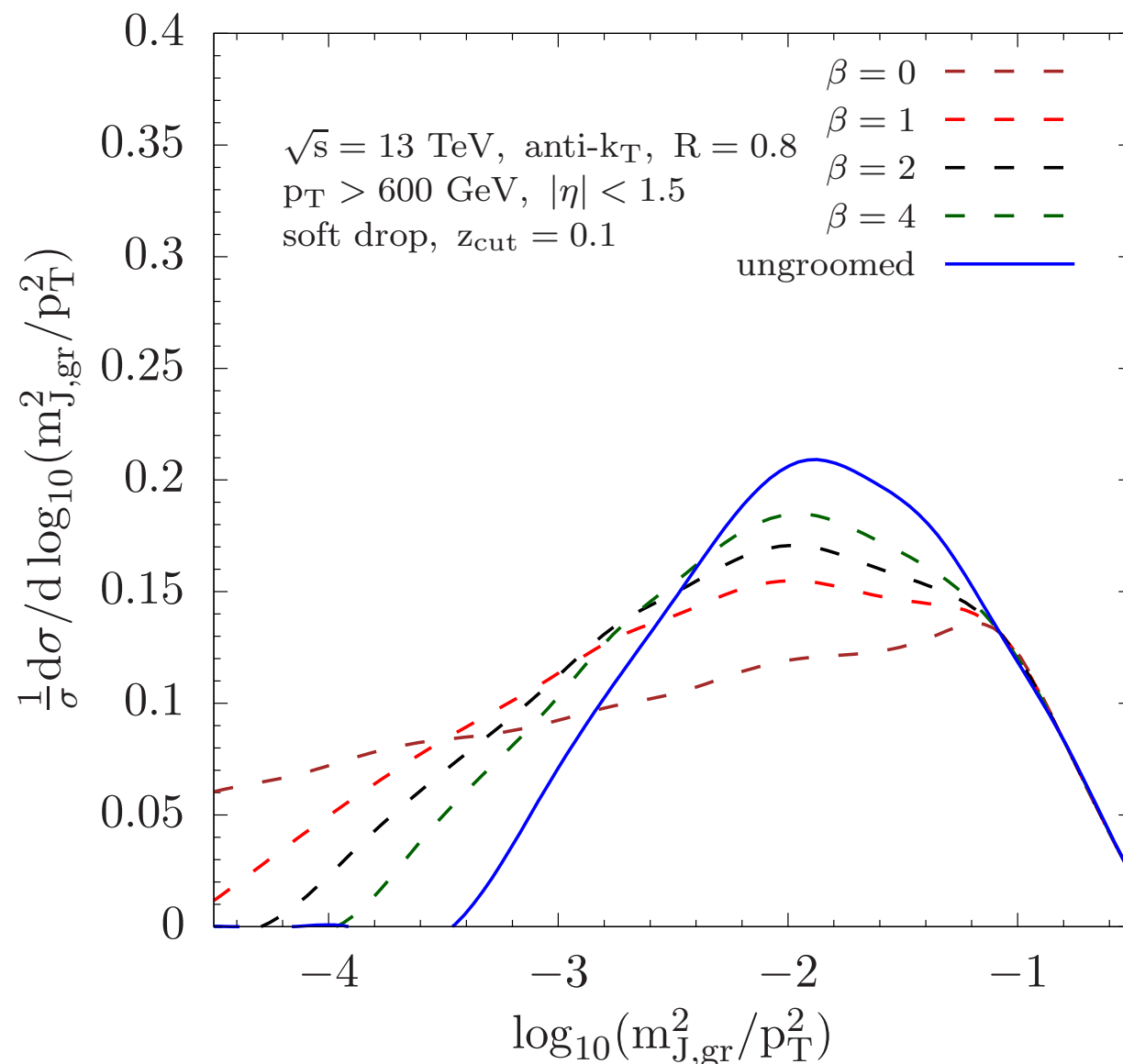
$$\alpha_s^n \ln^n(z_{\text{cut}}) \quad n \geq 2$$



Limit to the ungroomed case

- Soft drop condition is passed trivially when $\beta \rightarrow \infty \Leftrightarrow$ returns ungroomed case.

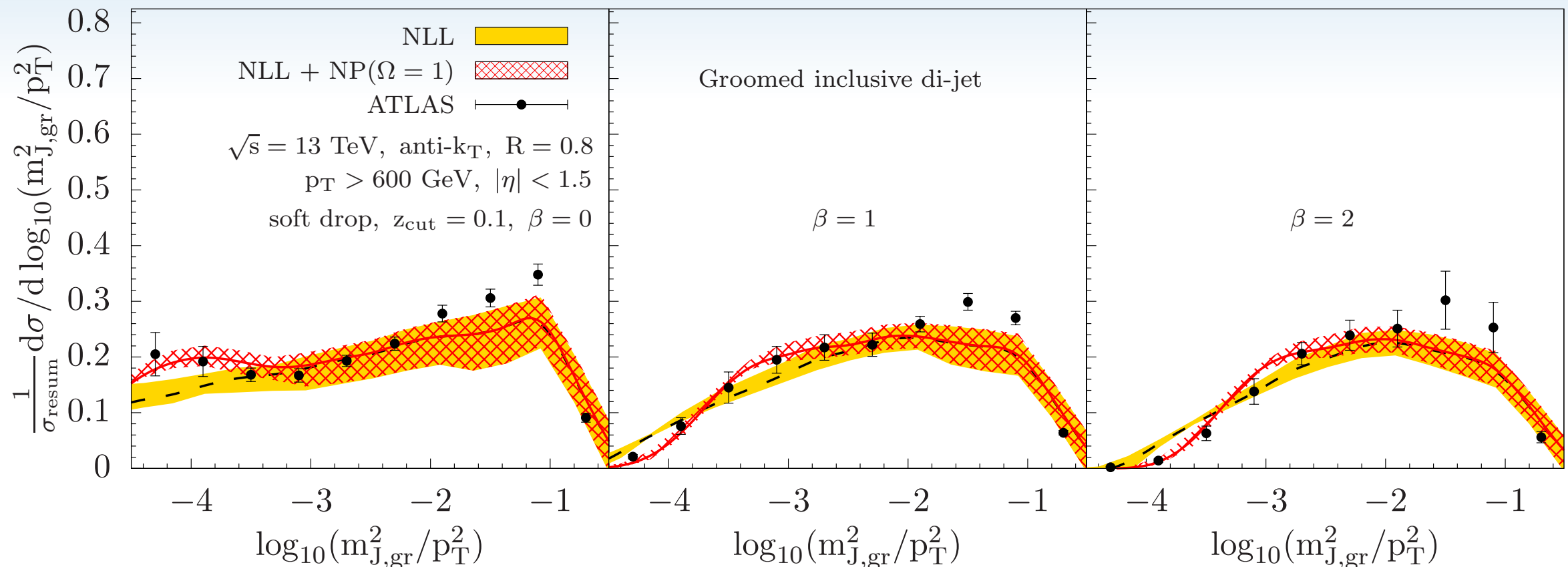
$$\frac{\min[p_{T,i}, p_{T,j}]}{p_{T,i} + p_{T,j}} > z_{\text{cut}} \left(\frac{R_{ij}}{R} \right)^\beta \rightarrow 0 \quad \text{when} \quad \beta \rightarrow \infty$$



Checked both numerically and analytically.

- At $\tau_{\text{gr}} = z_{\text{cut}} R^2$, the groomed result transitions to the ungroomed case.

Phenomenology (groomed jet mass)



- Developed the formalism for single inclusive groomed jet mass cross-section.
- Shows very good agreement with the data.
- $\Omega_k = 1 \text{ GeV} \implies$ Reduced contamination as expected.
NP effects mostly from hadronization.

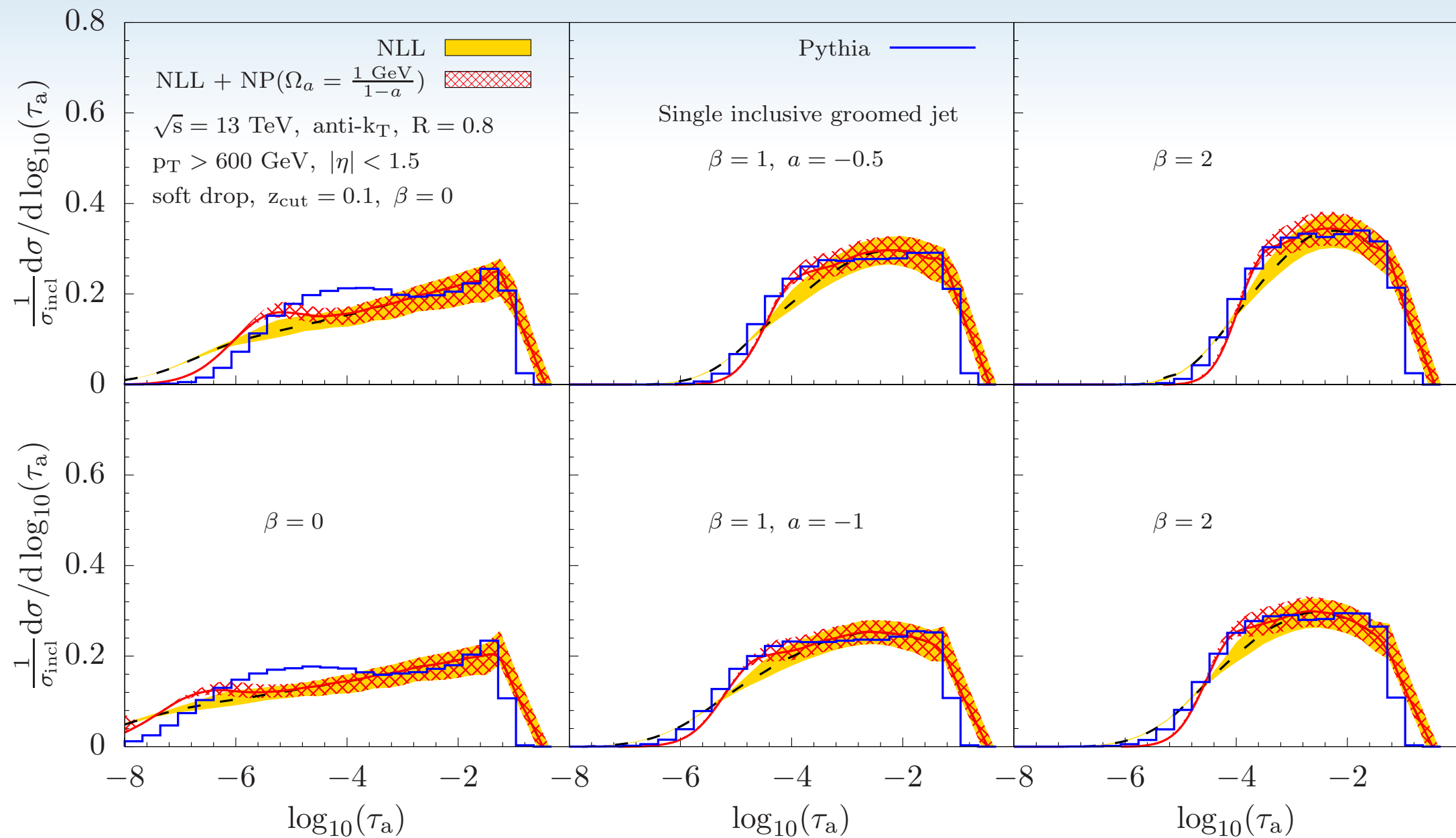
See also

ATLAS, *arXiv:1711.08341*

Larkoski, Marzani, Soyez, Thaler '14

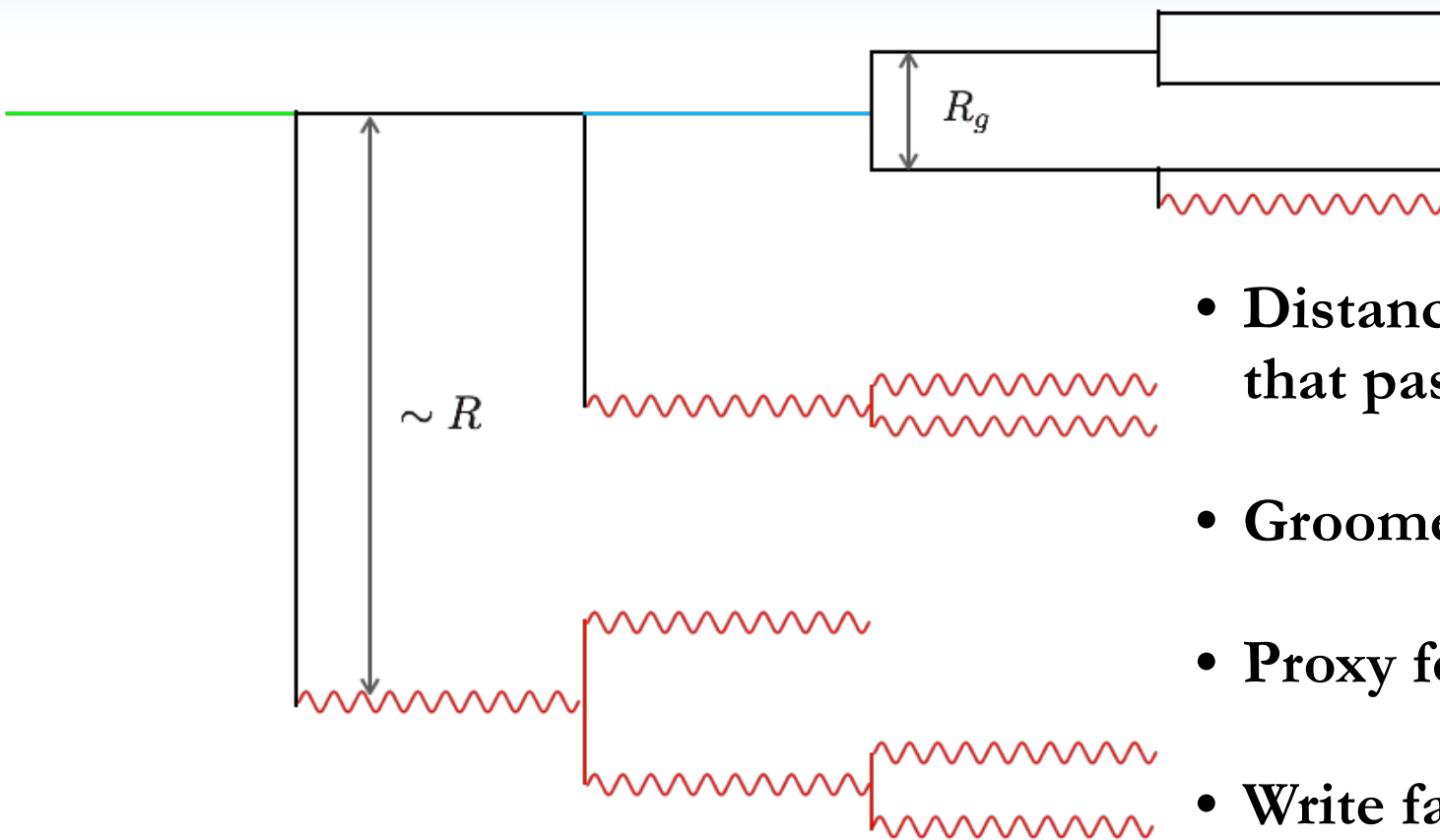
Frye, Larkoski, Schwartz, Yan '16

Phenomenology



- General angularities show decent agreement with Pythia with reduced contamination from UE/PU.

Groomed Jet Radius, R_g



Larkoski, Marzani, Soyez, Thaler '14
 Tripathy, Xue, Larkoski, Marzani, Thaler '17

- Distance between the two branches that passes the soft drop condition.
- Groomed jet area is approximately $\approx \pi R_g^2$
- Proxy for the sensitivity to contamination from pileup.
- Write factorization for cumulative distribution.

$$\frac{d\Sigma(\theta_g)}{d\eta dp_T} = \sum_{abc} f_a(x_a, \mu) \otimes f_b(x_b, \mu) \otimes H_{ab}^c(x_a, x_b, \eta, p_T/z, \mu) \otimes \mathcal{G}_c(z, p_T, \theta_g, \mu; z_{\text{cut}}, \beta)$$

Factorization for R_g

- $z_{\text{cut}} \ll 1$ and $R_g \ll R$
- For $z_{\text{cut}} \ll 1$, groomed jet axis and ungroomed jet axis

Hard-collinear

$$z_H \sim 1 \quad \theta_H \sim R \quad \text{correlation}$$

$\notin gr \text{ soft}$

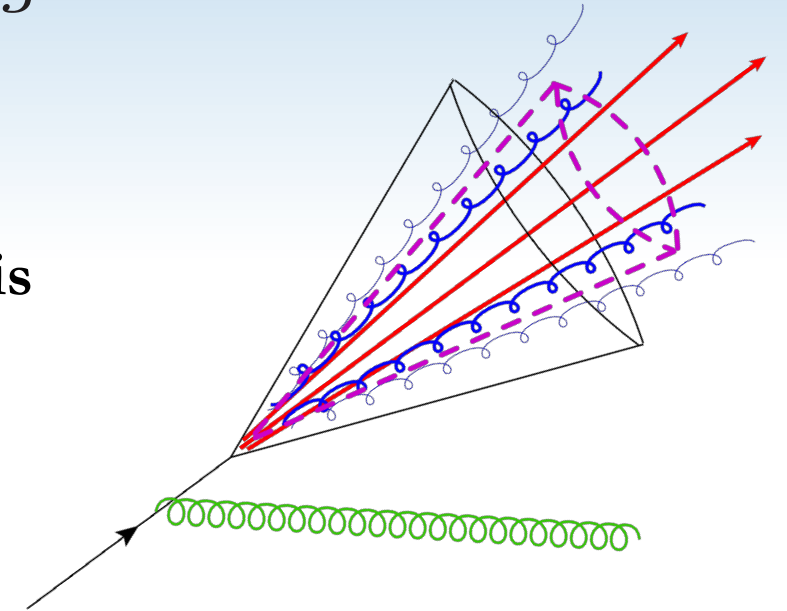
$$z_{\notin gr} \sim z_{\text{cut}} \left(\frac{\theta}{R} \right)^\beta = z_{\text{cut}} \quad \theta_{\notin gr} \sim R$$

Collinear

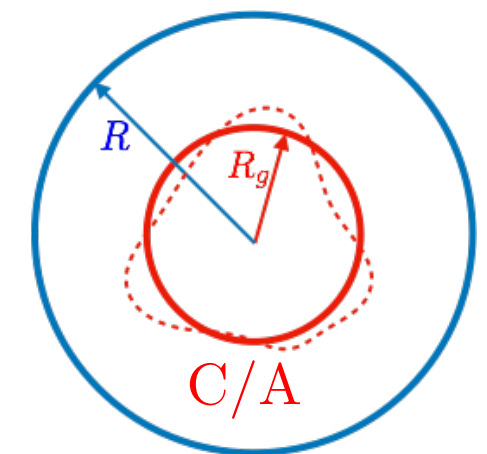
$$z_c \sim 1 \quad \theta_c \sim R_g \quad \text{correlation}$$

$\in gr \text{ soft}$

$$z_{\in gr} \sim z_{\text{cut}} \left(\frac{\theta}{R} \right)^\beta = z_{\text{cut}} \left(\frac{R_g}{R} \right)^\beta \quad \theta_{\in gr} \sim R_g$$



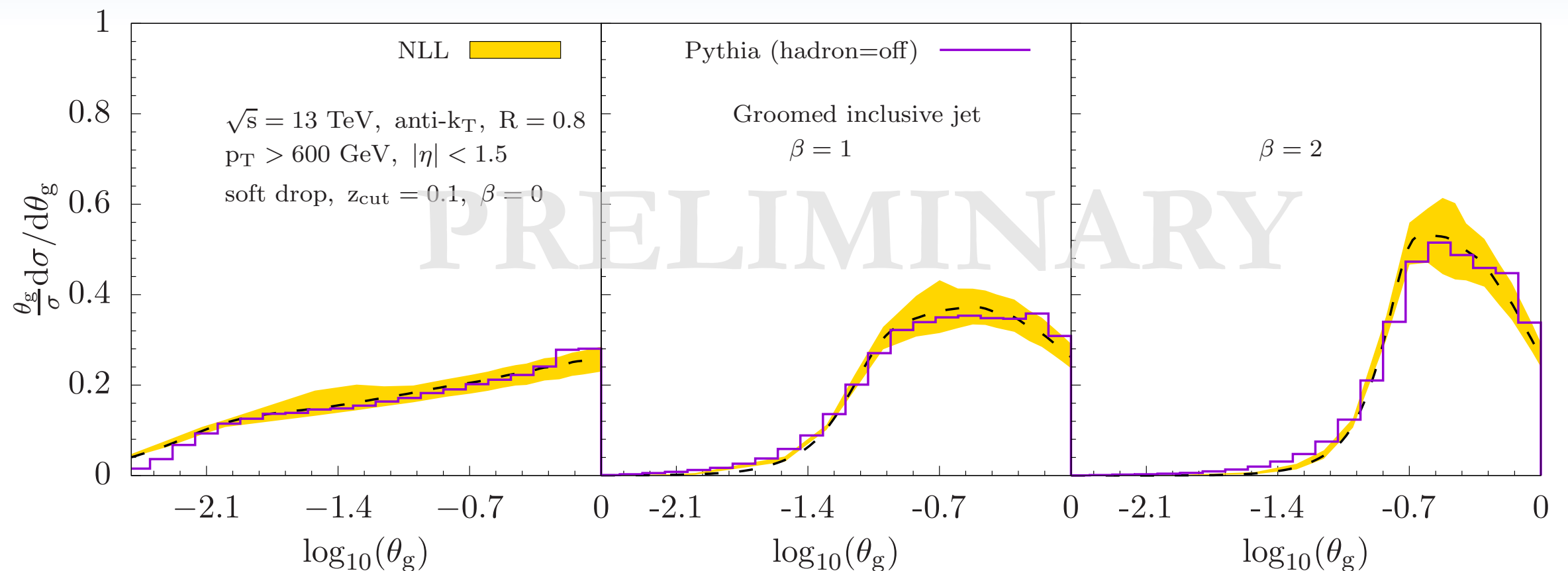
anti- k_T



Kelley, Walsh, Zuberi '12
Becher, Neubert, Rothen, Shao '15, '16

$$\mathcal{G}_c(z, p_T, R, R_g, \mu; z_{\text{cut}}, \beta) = \sum_i \sum_n \mathcal{H}_{c \rightarrow i}^n(z, p_T R, \mu) \otimes_\Omega S_{i,n}^{\notin gr}(z_{\text{cut}} p_T R, \mu) \sum_m C_i^m(p_T R_g, \mu) \otimes_\Omega S_{i,m}^{\in gr}(\theta_g z_{\text{cut}} p_T R, \mu; \beta)$$

Phenomenology (groomed jet radius)



$$\theta_g = \frac{R_g}{R}$$

- Shows very good agreement without having to account for NGLs and clustering logs.

Conclusions

- Formalisms for studying semi-inclusive jet production with and without a substructure measurement were introduced.
- Discussed phenomenology of ungroomed jet mass in the LHC.
- Discussed various non-perturbative effects.
- Discussed soft drop grooming which reduce contaminations from the uncorrelated radiations.
- Discussed factorization of groomed angularities and groomed jet radius, and demonstrated reduced contamination and good agreements with Pythia.