

Extracting the in-Medium Color Force from Heavy-Ion Collisions

Xiaojian Du, Shuai Liu, Ralf Rapp

Cyclotron Institute
Department of Physics and Astronomy
Texas A&M University

Santa Fe Jets and Heavy Flavor Workshop 2019

2019.01.29

University of California at Los Angeles, CA



Outline

- Introduction: Why Quarkonium
- Quarkonium Transport Approach
 - Rate Equation
- Advantage of Bottomonium
 - Comparison to Charmonium
- Statistical Extraction of the Heavy Quark Potential
 - Assembly Line
 - Ansatz for Potential
 - χ^2 Tests and Results
- Summary

Quarkonium in Heavy-Ion Collisions

Heavy Quarkonium, J/ψ , $\psi(2S)$, $\Upsilon(1S)$, $\Upsilon(2S)$,as a probe:

1. Survive in QGP ($E_{\text{BINDING}} > T_c$),
2. Large masses (potential picture works),
3. Various species (hierarchy), ...

Hierarchy:

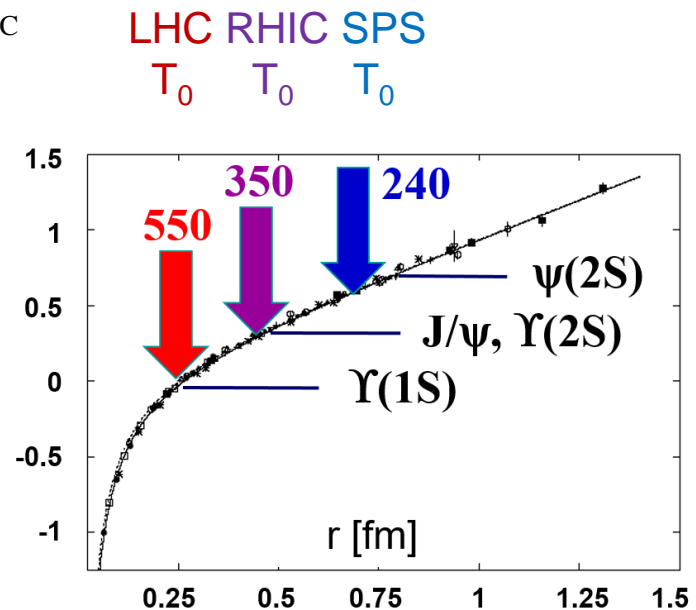
$$T_{\text{melt}}[\psi(2S)] < T_0^{\text{SPS}} < T_{\text{melt}}[J/\psi, \Upsilon(2S)] < T_0^{\text{RHIC}} < T_{\text{melt}}[\Upsilon(1S)] < T_0^{\text{LHC}}$$

For J/ψ , $\psi(2S)$, $\Upsilon(1S)$, $\Upsilon(2S)$:

- Bound at different parts of potential
- Different melting temperatures
- Systematic excitation functions: SPS, RHIC, LHC...

Ideal for probing:

strong force / $Q\bar{Q}$ potential in medium



Transport Approach

Kinetic Rate Equation

$$\frac{dN_Y(\tau)}{d\tau} = -\Gamma_Y(T(\tau)) [N_Y(\tau) - N_Y^{\text{eq}}(T(\tau))]$$

primordial

regeneration

Transport Coefficients

- Reaction Rates

$\Gamma_Y(T(\tau))$



NLO Quasi-free (dominates small binding)
LO Gluo-dissociation (dominates large binding)

- Equilibrium Limits

$N_Y^{\text{eq}}(T(\tau))$



From heavy quark
conservation

$$N_Y^{\text{eq}} = V_{FB} \gamma_b^2 d_Y \int \frac{d^3p}{(2\pi)^3} f_Y^{\text{eq}}(E_p; T)$$

Key Parameters

- Coupling α_s



Affects reaction rates

Fixed from previous calculations compared with data

- Thermal Relaxation Time
of Heavy Quarks



Modifies equilibrium limit

Extracted from heavy quark diffusion simulations

Key Inputs

- In-Medium $Q\bar{Q}$ Potential/Binding Energy → Reaction Rates

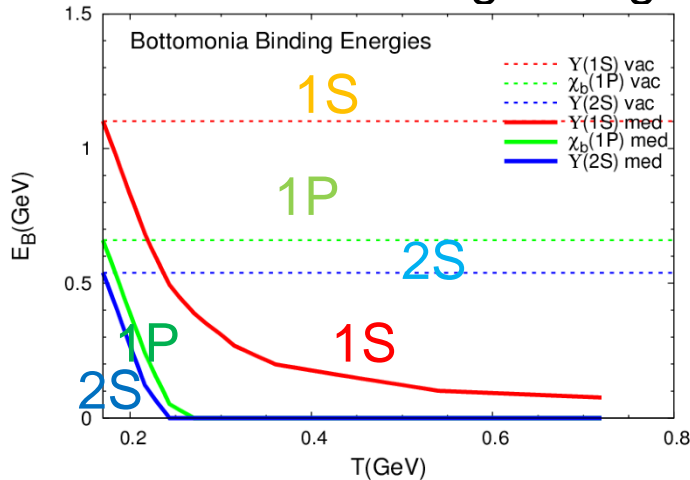
- Heavy Quark/-onium pp Cross Section → fugacity factor γ_b → Equilibrium Limit

- Initial State Effects (nPDF)

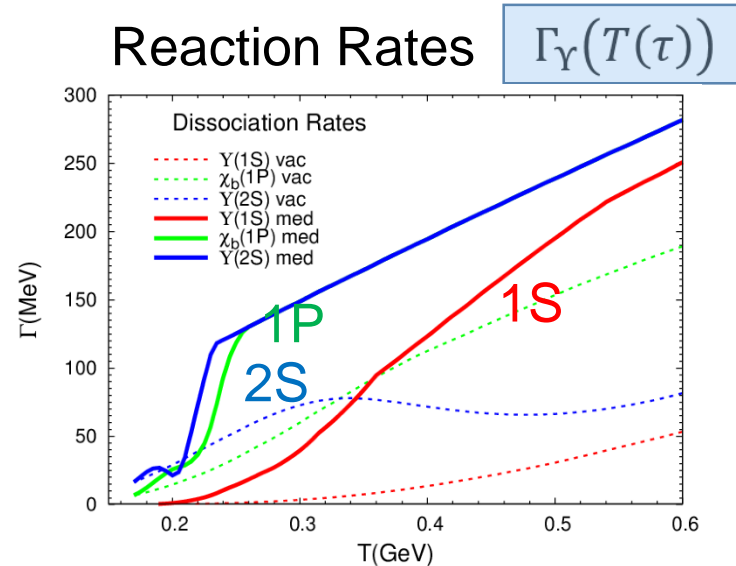
- Fireball Evolution

Binding Energy and Reaction Rates

Bottomonium Binding Energies

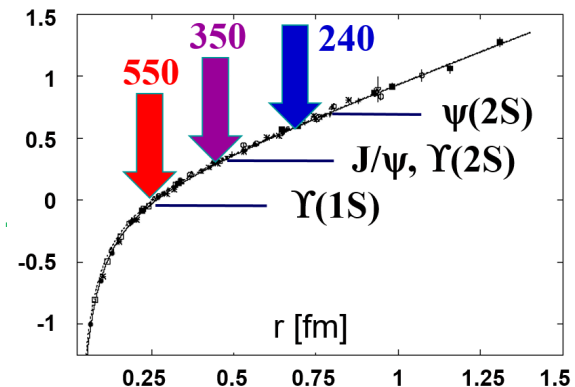


Reaction Rates



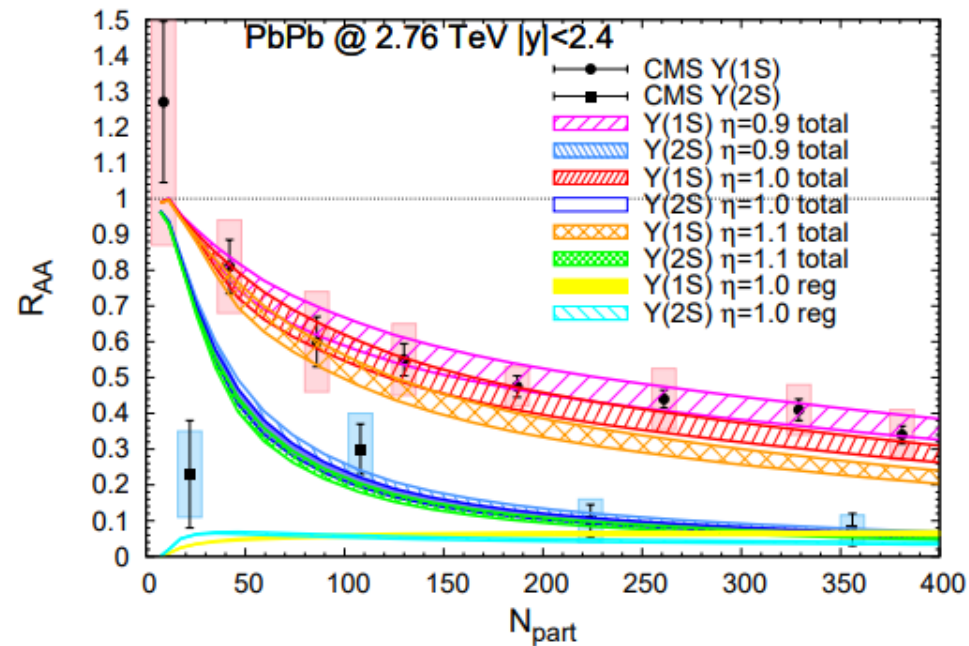
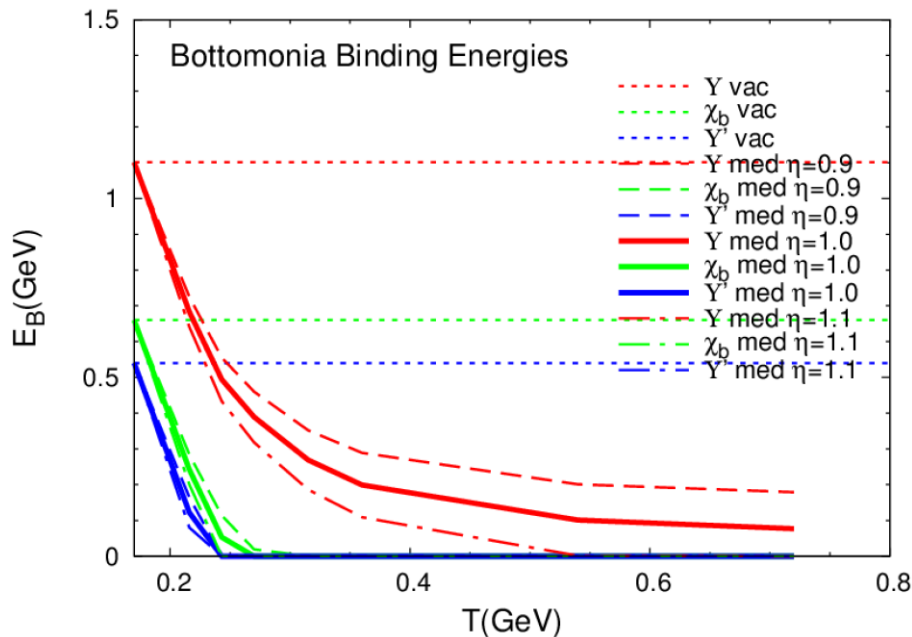
X. Du, M. He, R. Rapp, PRC96 (2017) 054901

- Hierarchy:
→ Different Melting temperatures for $Y(1S)$, $Y(2S)$..
- Probing In-Medium Potential



Bottomonium: Sensitivity to Binding Energy

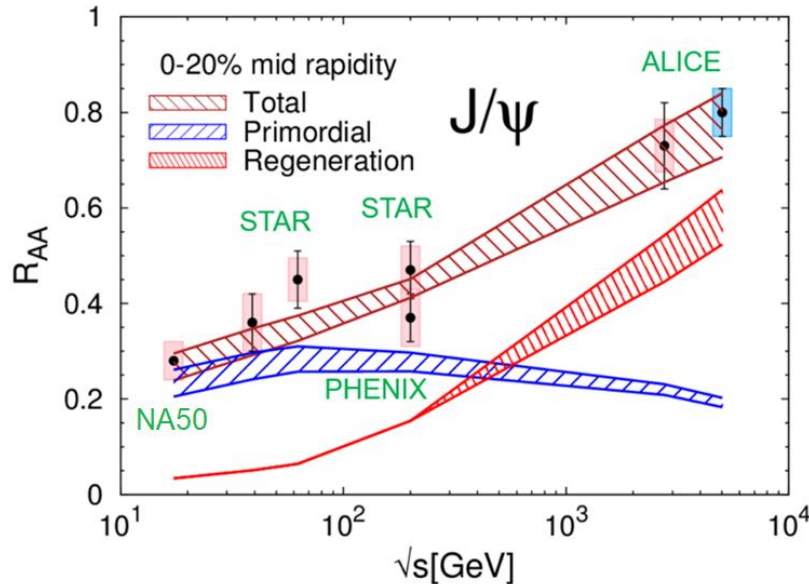
Binding Energies \rightarrow Reaction Rates $\xrightarrow{\text{Transport}}$ R_{AA} Observables



- Significant binding energy dependence of the $Y(1S)$ R_{AA}

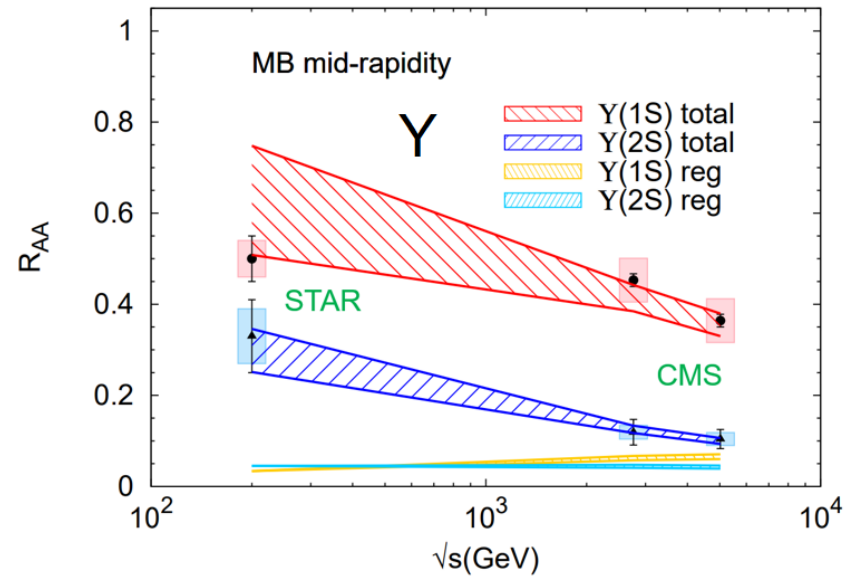
X. Du, M. He, R. Rapp, PRC96 (2017) 054901

Bottomonium: Small Regeneration



R. Rapp, X. Du, NPA967 (2017) 216

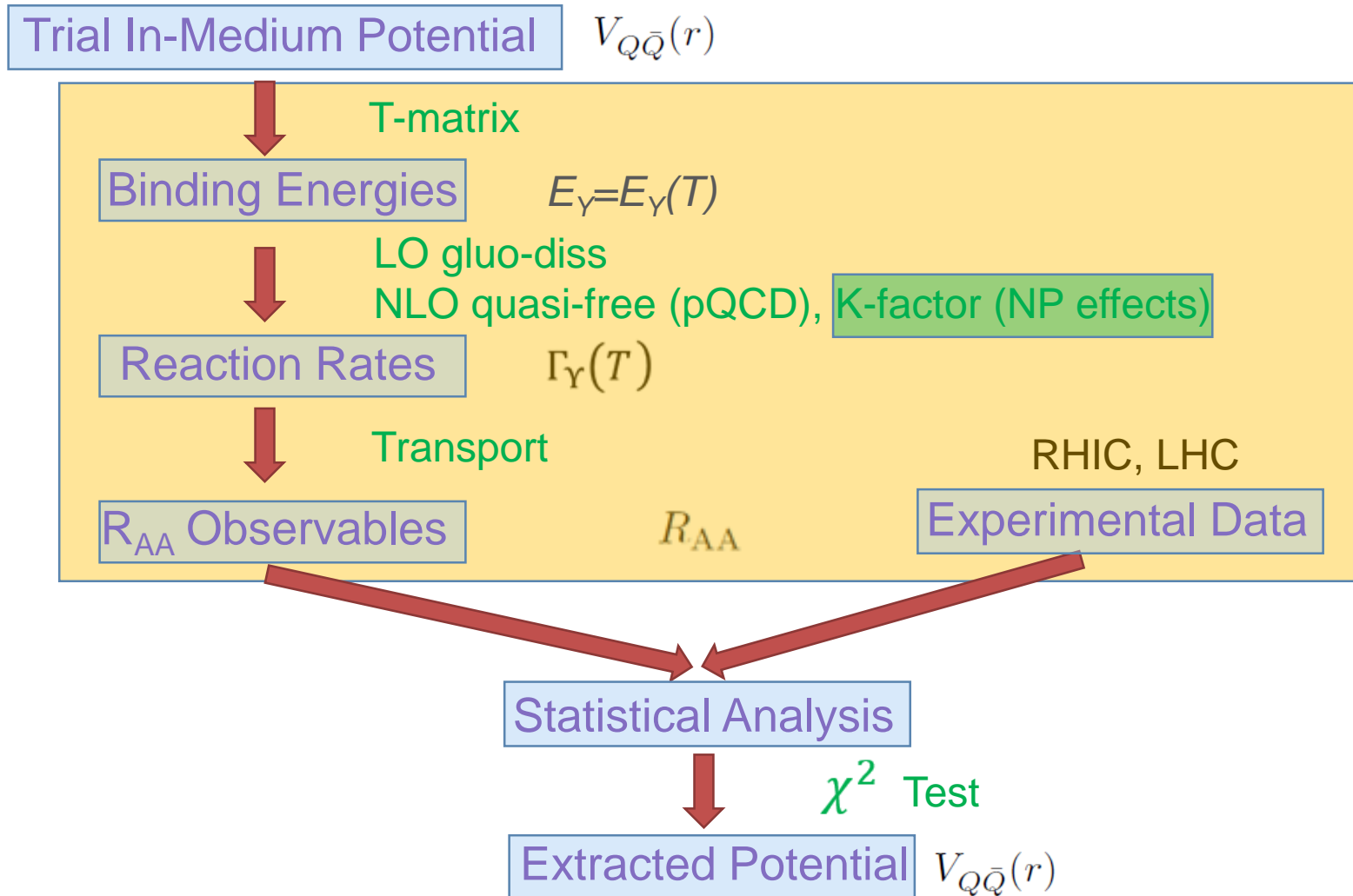
- J/ψ Excitation Function:
→ Demonstration of regeneration
- Bottomonia have species



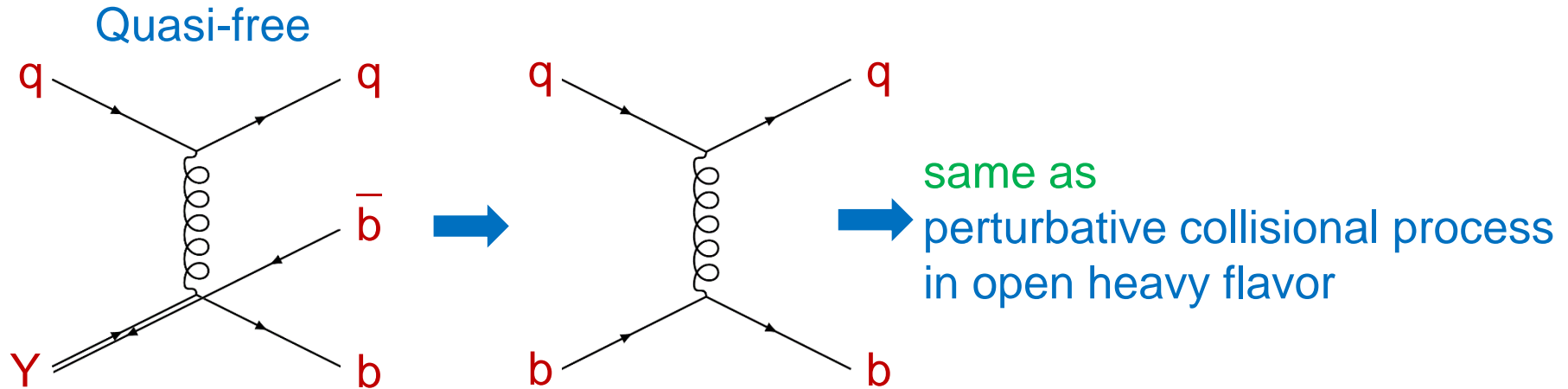
- J/ψ and $Y(2S)$:
 $E_{\text{bind}}[J/\psi] \approx E_{\text{bind}}[Y(2S)]$
but
Different excitation functions
→ Due to larger regeneration for J/ψ

- Bottomonium:
→ Clean and suitable for probing in-medium potential !

Assembly Line



K-factor



K-factor:

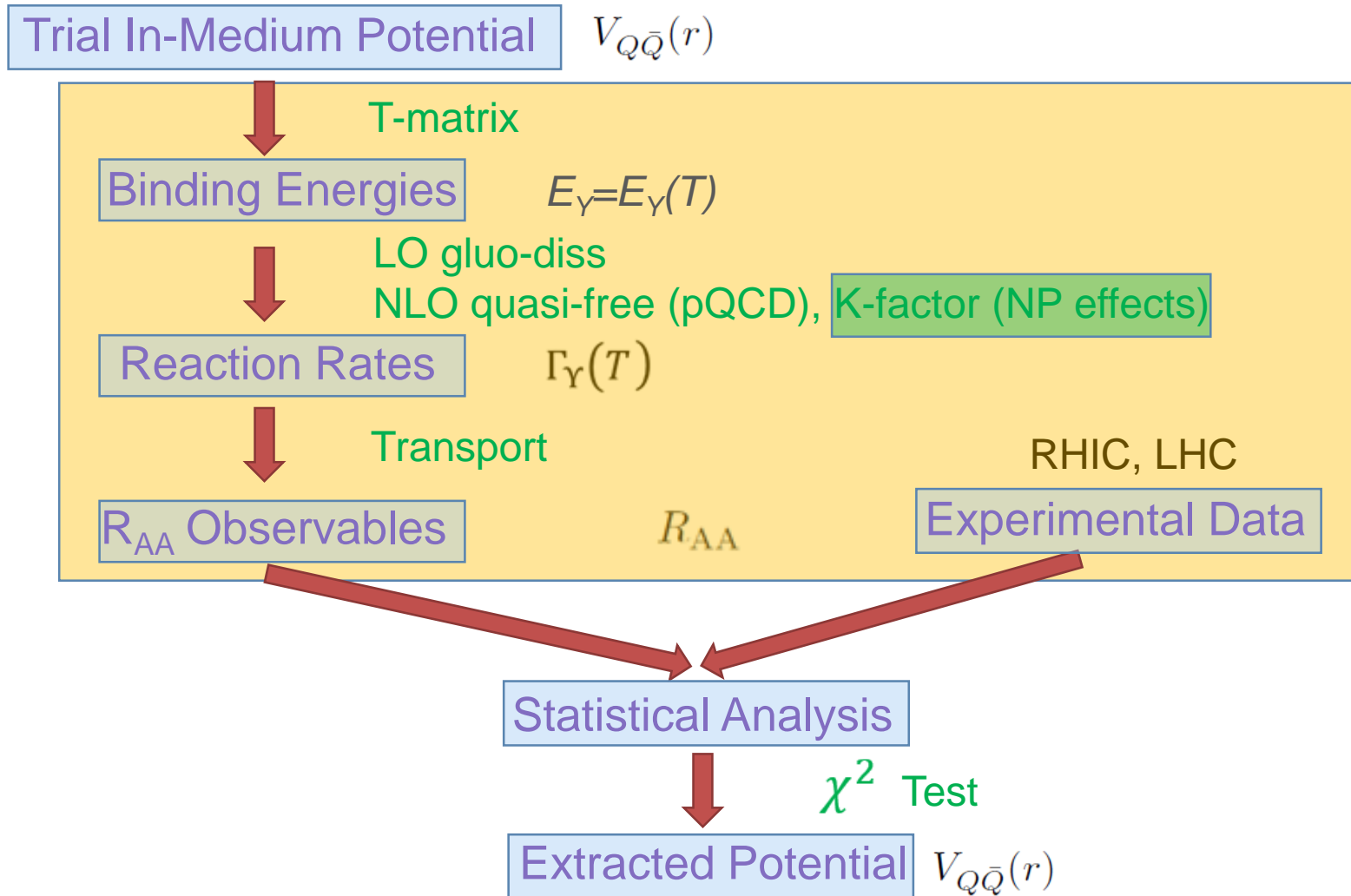
- Multiplication to pQCD calculated Reaction Rates
- Accounts for nonperturbative effects
- Open heavy-flavor suggests **K-factor** > 5

R. Rapp, et al, NPA979 (2018) 21-86
S. Cao, et al, arXiv: 1809.07894
R. Rapp, arXiv: 1901.06440

$D_s(2\pi T)$:

- pQCD: ~ 30
- Required by phenomenology: $2 \sim 5$

Assembly Line



Trial Potential

$$V_{Q\bar{Q}}(r) = -\frac{4}{3}\alpha_s \left(\frac{e^{-\boxed{m_D} r}}{r} + m_D \right) + \min \left(\sigma r, \frac{\sigma}{\boxed{m_S}} \right)$$

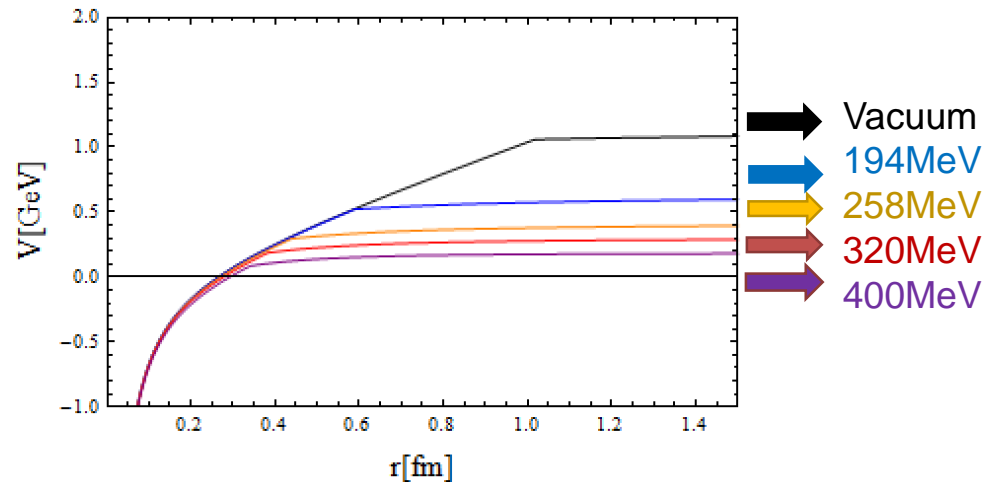
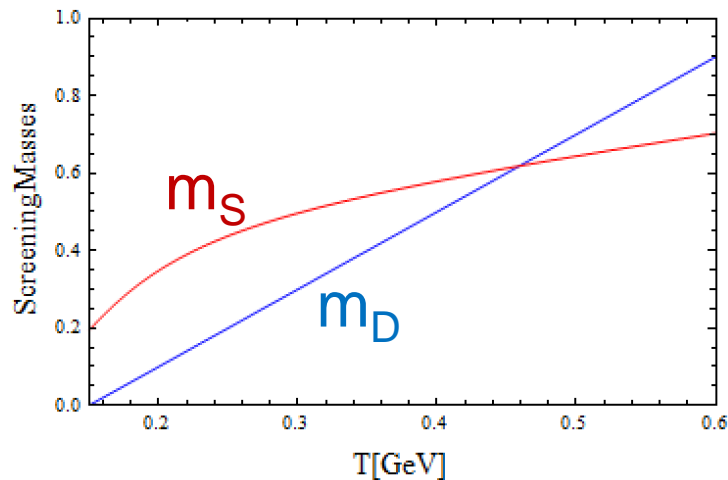
Coulomb

Screening
Masses

$$m_D = a(T - T_0)$$

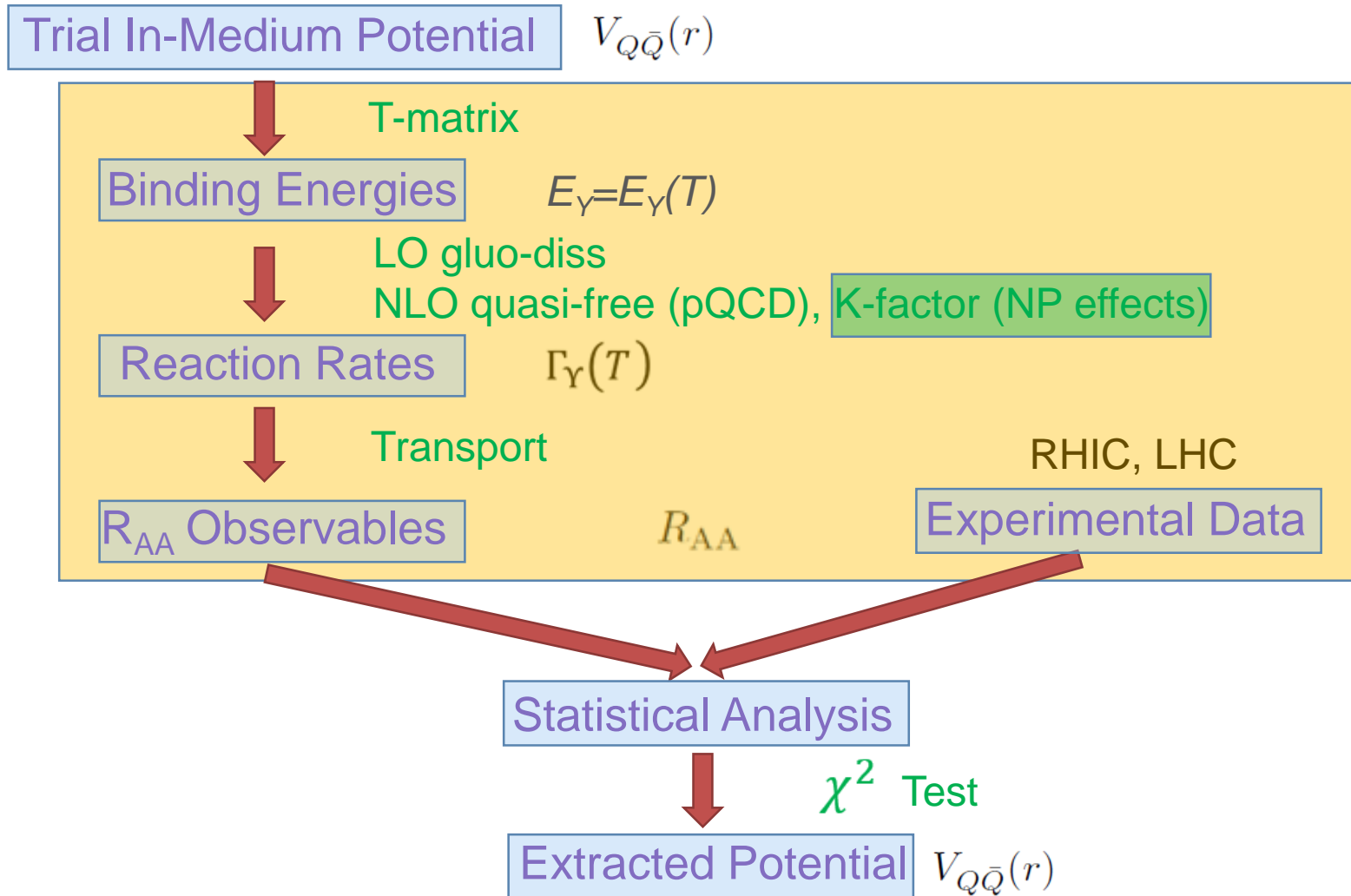
String

$$m_S = m_v + b \left(\sqrt{(T - T_0)^2 + (dT_0)^2} - dT_0 \right) - c \left(\sqrt{(T - T_0)^2 + (dT_0)^2} - (T - T_0) - dT_0 \right)$$



In relevant parameter space, sensitive parameters: a ($\sim m_D/T$) & b ($\sim m_S/T$ at high- T)

Assembly Line



Statistical Analysis: χ^2 Test

Model Calculations
with trial potential

$$\chi^2 = \sum_{i=1}^N \left(\frac{R_{AA}^{\text{mod}}(a, b, c, d) - R_{AA}^{\text{exp}}}{\sigma_{\text{exp}}} \right)^2$$

Data selection (Bottomonium data):
53 R_{AA} data points from RHIC to LHC
energies

Multiple Collision Systems:
193GeV U-U, 200GeV Au-Au,
2.76TeV Pb-Pb, 5.02TeV Pb-Pb.

Multiple Rapidities:
mid-rapidity, forward rapidity.

Multiple Species:
1S, 2S, 3S, 1S+2S+3S, 2S+3S, ...

Chi-square distribution degree of freedom:
53 data – 4 parameters = 49

p-value

Assuming the model is correct, there is p-value chance
that the data is at least this discrepancy from the model

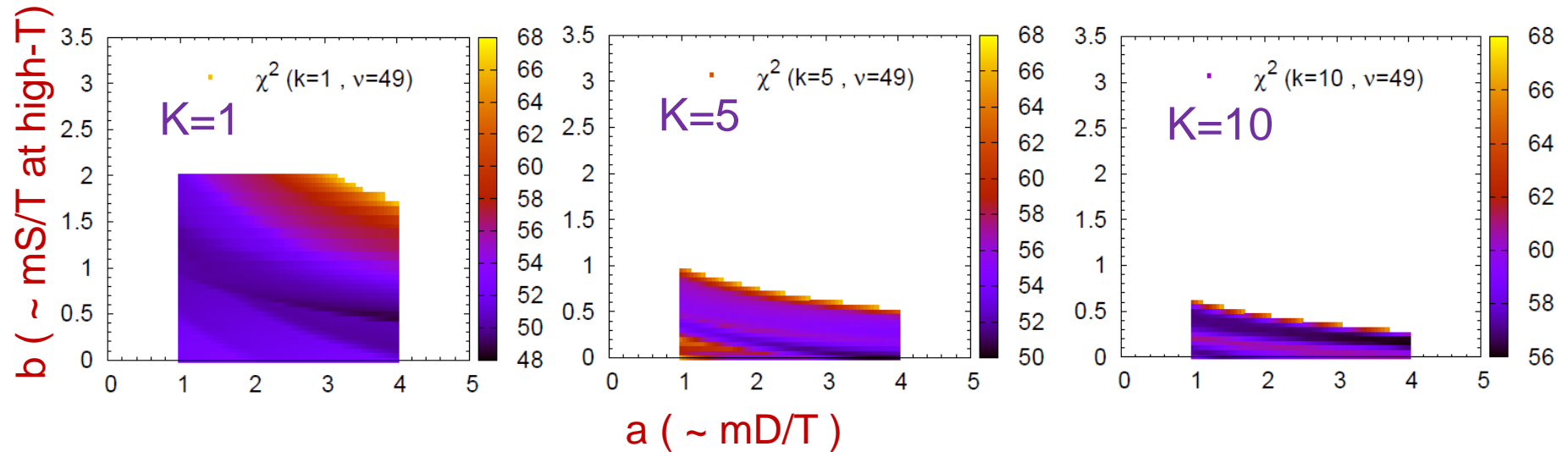
- Larger χ^2 smaller p-value
- Smaller χ^2 larger p-value

Accepted band
95% confidence: p-value > 0.05



Best fit
Smallest χ^2
Largest p-value

Resultant Parameter Space

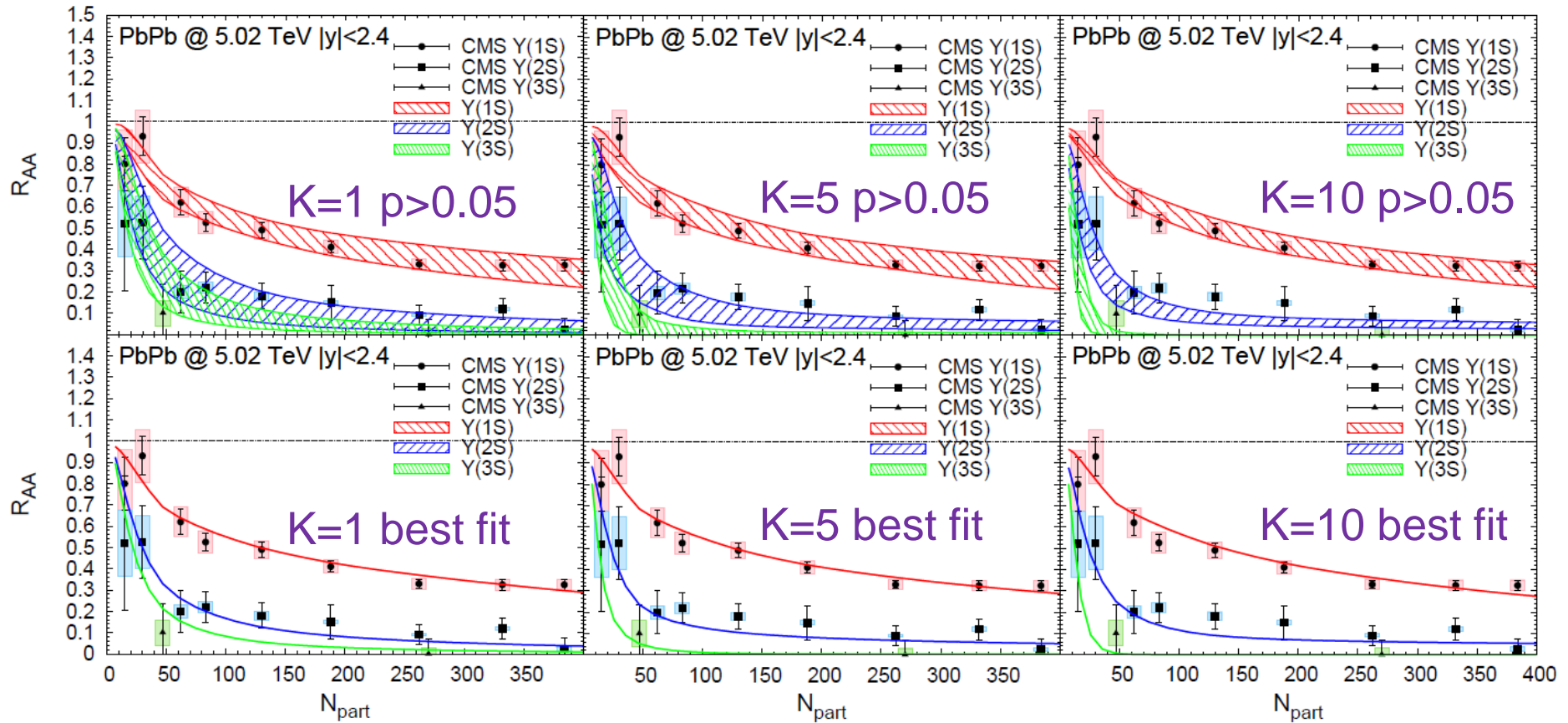


Smaller $b \rightarrow$ smaller $m_s \rightarrow$ stronger binding
Larger $b \rightarrow$ larger $m_s \rightarrow$ weaker binding

In order to describe data, rates are approximately constant at relevant temperature ($<400\text{MeV}$)

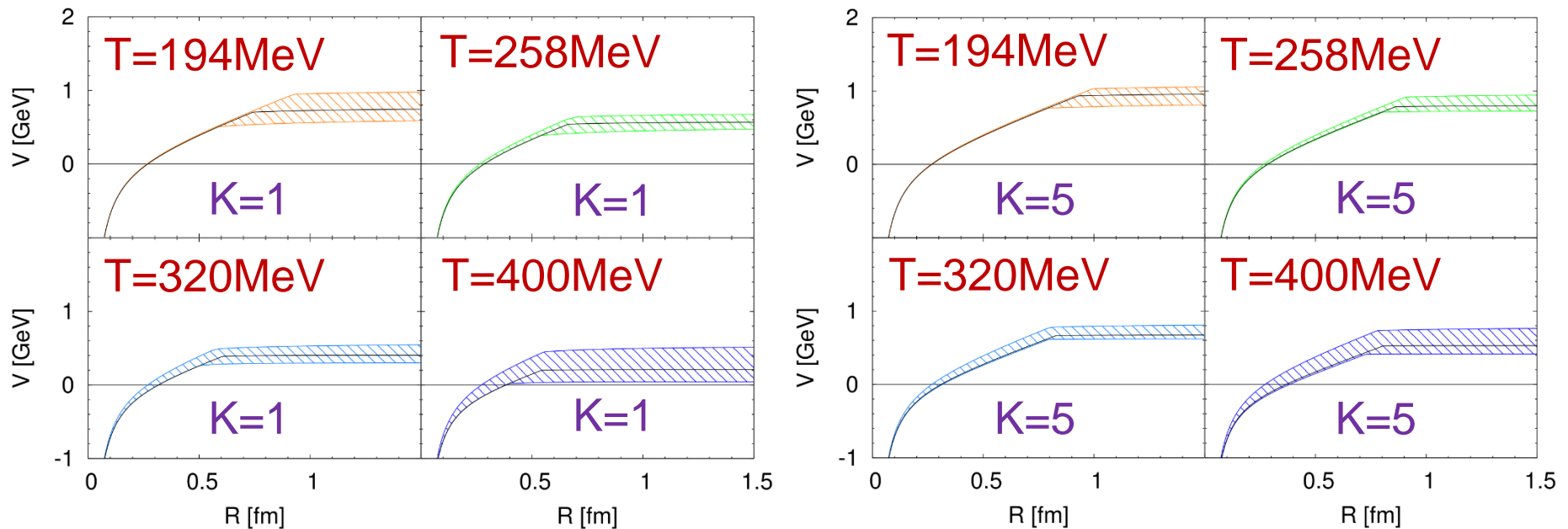
Thus Increasing K-factor requires stronger binding

Resultant Bands of R_{AA} s



Good description of data with parameters selected by statistics

Resultant Bands of In-Medium Potential



$V^{\text{best-fit}}(r=1\text{fm})$	$K=1$	$K=5$
$T=194\text{MeV}$	730MeV	940MeV
$T=320\text{MeV}$	400MeV	670MeV

$K=1$: allows for both weak/strong binding

$K=5$: suggests strong binding

$K=10$: even stronger..

Summary

- Bottomonium is ideal for probing the in-medium potential
 - R_{AA} observable is sensitive to in-medium potential
 - Has small regeneration contribution
 - Allows for more species (potentially more data constraints)
- Integration of open/hidden heavy flavor is important
 - A large K-factor is needed accounting for nonperturbative effects
- In-medium heavy-quark potential is strong
 - Statistics tell the story

Thank you!