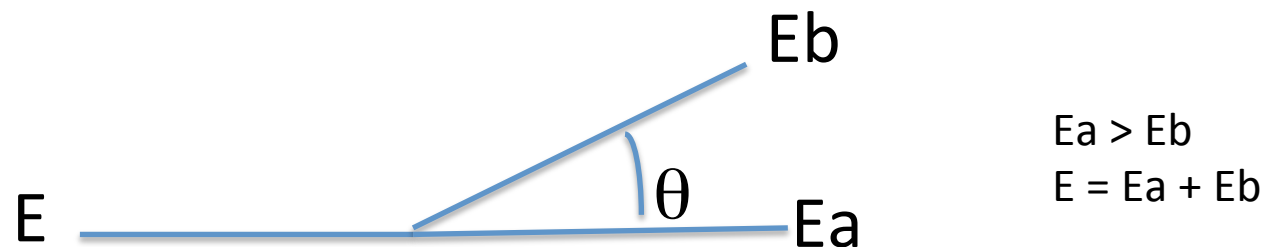


Searching for the dead cone effects with iterative declustering of heavy-flavor jets

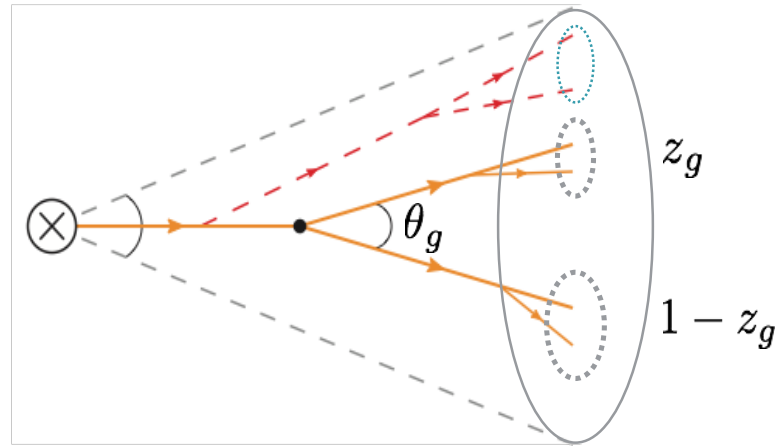
Leticia Cunqueiro (ORNL/UTK), MP

<https://arxiv.org/abs/1812.00102>

$$\Theta_{\text{d.c.}} < m / E$$



Y.Dokshitzer, V.A.Khoze and S.I.Troyan "On specific QCD properties of heavy quark fragmentation"



JET STRUCTURE

Jet Lund diagram

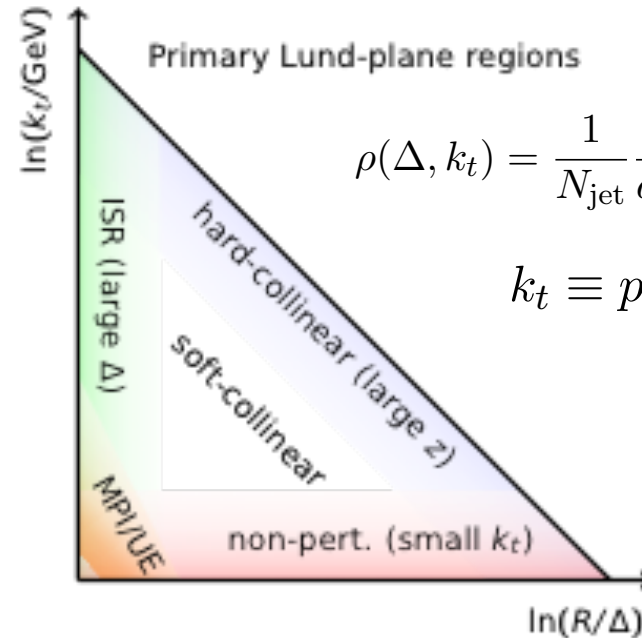
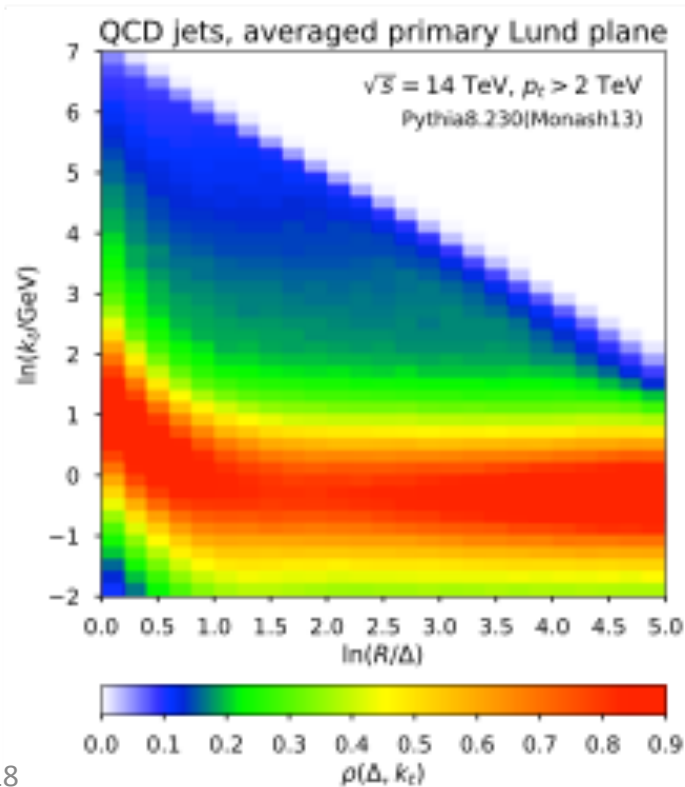


$$p_{T,a} > p_{T,b}, \quad \kappa = \frac{p_{T,b}}{p_{T,a} + p_{T,b}} \Delta_{ab}^3$$

$$\bar{\rho}(\Delta, \kappa) = \frac{1}{N_{\text{jet}}} \frac{dn_{\text{emission}}}{d \ln \kappa \, d \ln 1/\Delta}$$

Lund diagrams, a theoretical representation of the phase space within jets, have long been used in discussing parton showers and resummations. We point out that they can be created for individual jets through repeated Cambridge/Aachen declustering, providing a powerful visual representation of the radiation within any given jet.

arXiv:1807.04758



$$\rho(\Delta, k_t) = \frac{1}{N_{\text{jet}}} \frac{dn_{\text{emission}}}{d \ln k_t \, d \ln 1/\Delta}$$

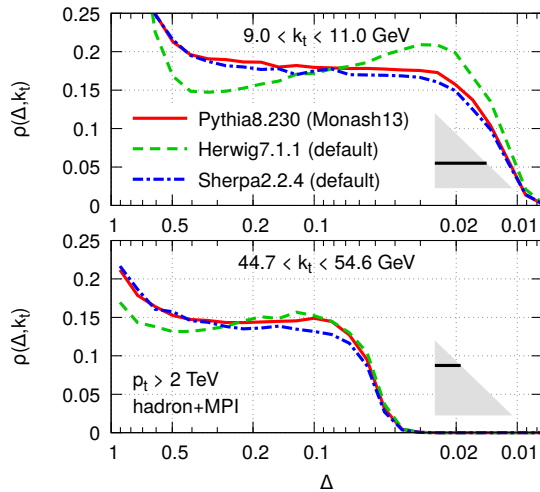
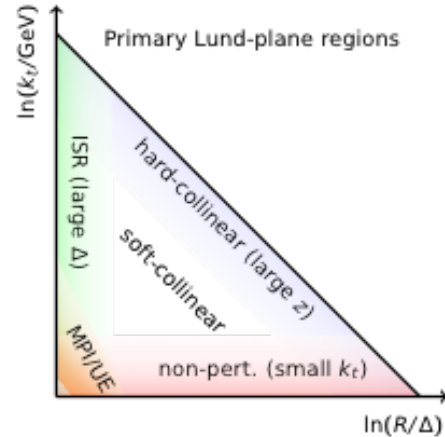
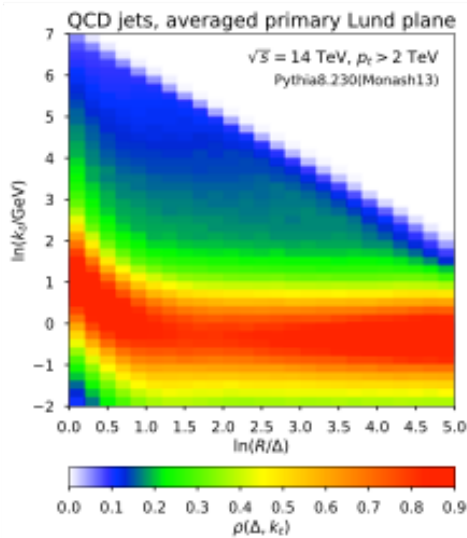
$$k_t \equiv p_{tb} \Delta_{ab},$$

Jet Lund diagram

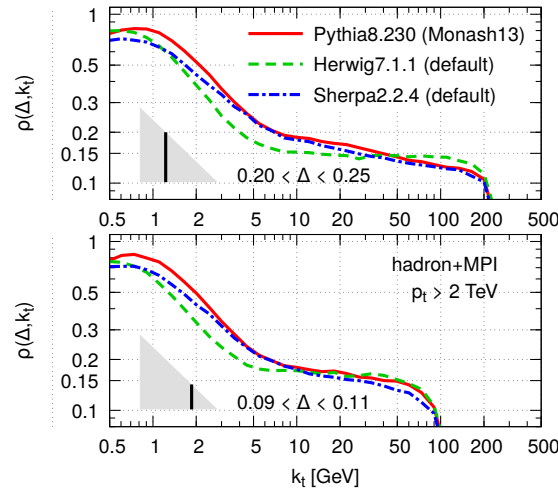
$$p_{T,a} > p_{T,b}, \kappa = \frac{p_{T,b}}{p_{T,a} + p_{T,b}} \Delta_{ab}^4$$

$$\bar{\rho}(\Delta, \kappa) = \frac{1}{N_{\text{jet}}} \frac{dn_{\text{emission}}}{d \ln \kappa d \ln 1/\Delta}$$

- Comparison of event generators
- Use for ML – jet ID (RHS below: boosted electroweak boson tagging at high momenta)



(a)



(b)

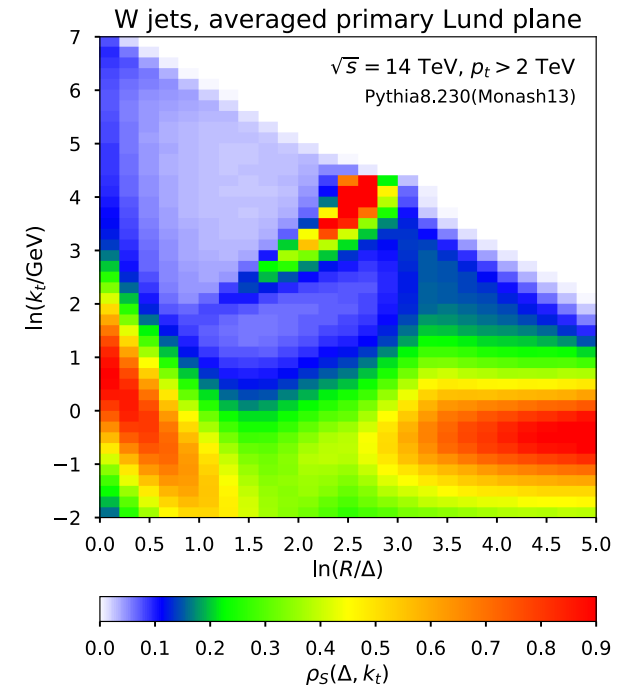
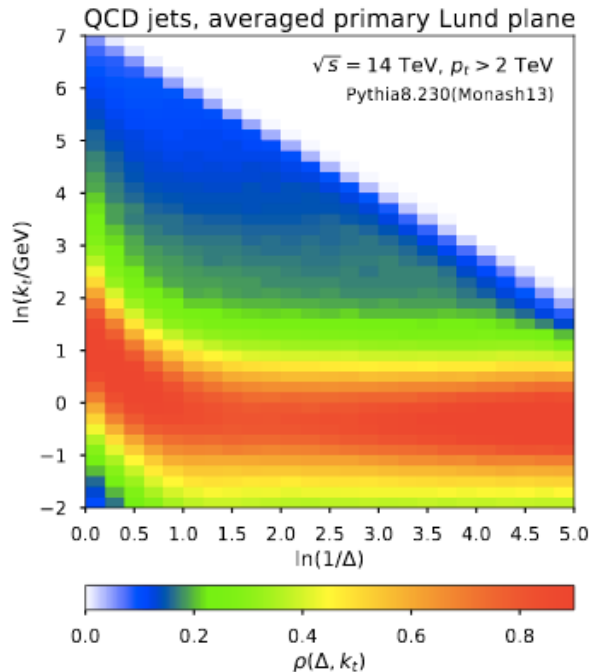


Figure 3: Emission density along slices of the Lund plane, at fixed k_t (top) and Δ (bottom), comparing three event generators.

The recipe... (1/2)



2D map of the jet tree filled via iterative declustering:

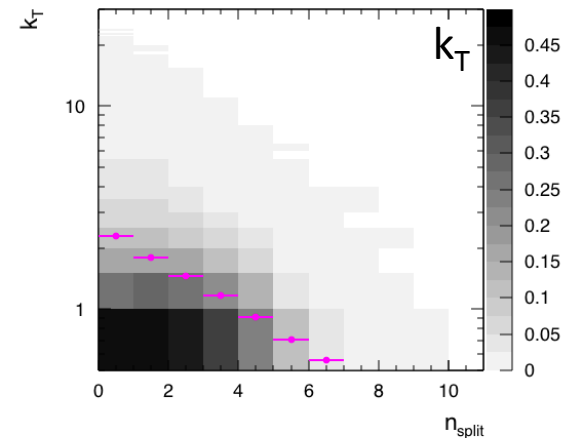
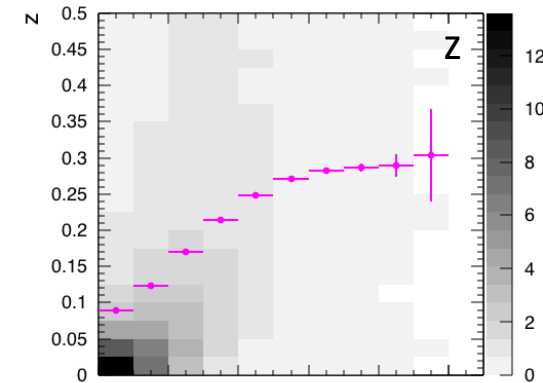
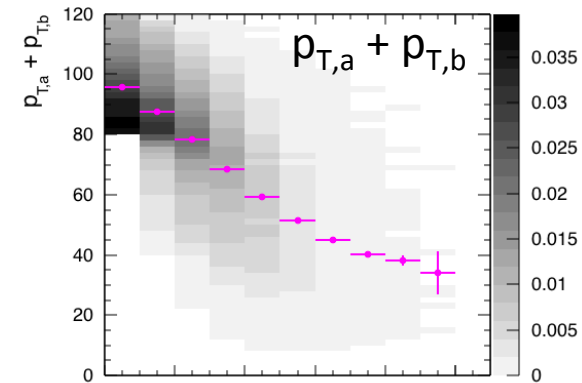
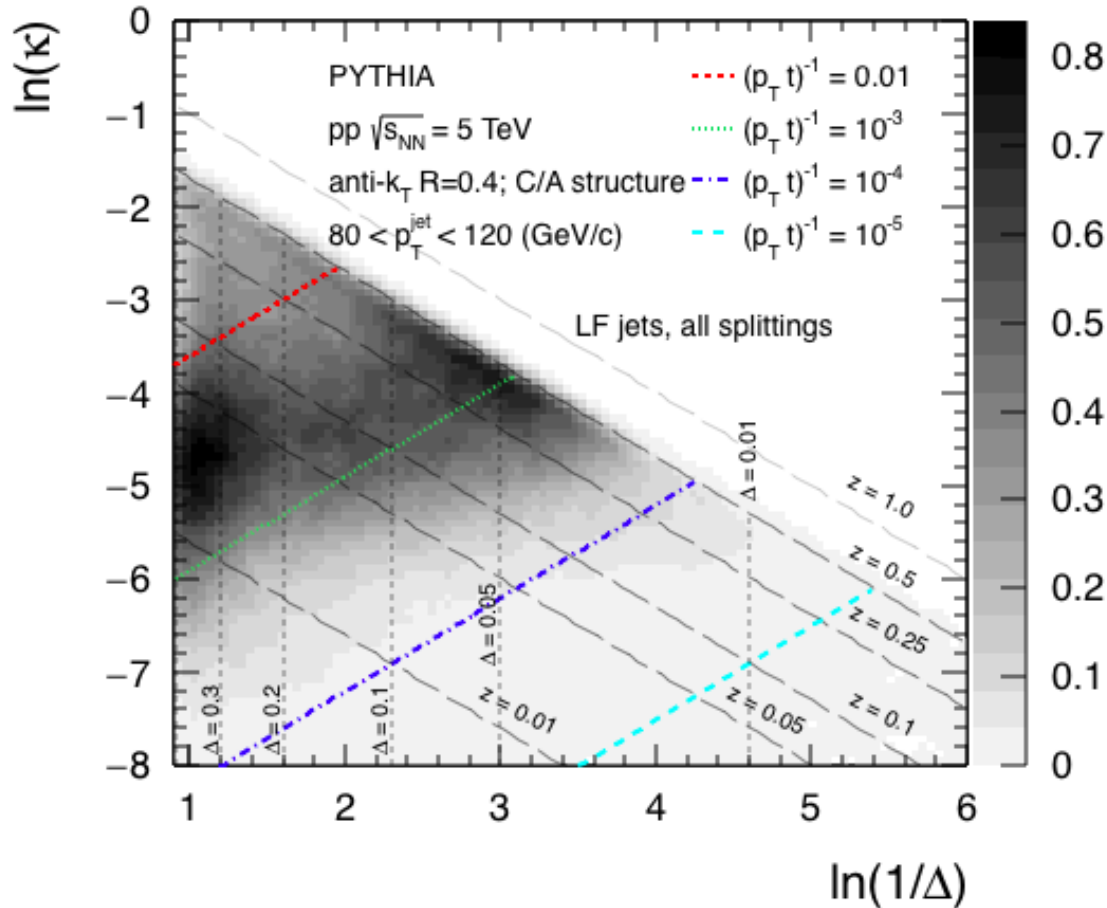
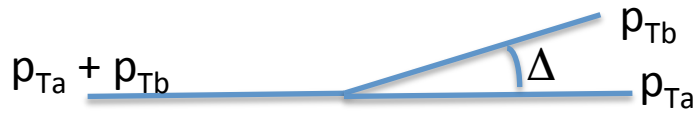
- Recluster the jet constituents with CA (a convenient metric for vacuum angular ordered showers)
- Uncluster it step by step, from the prongs that were merged last to the prongs that were merged first.
- Each time, register the scale k_T and the opening angle θ between the subleading and leading prong
- Always follow the hardest prong

The result is a flat 2D density map (except for the running of the coupling)

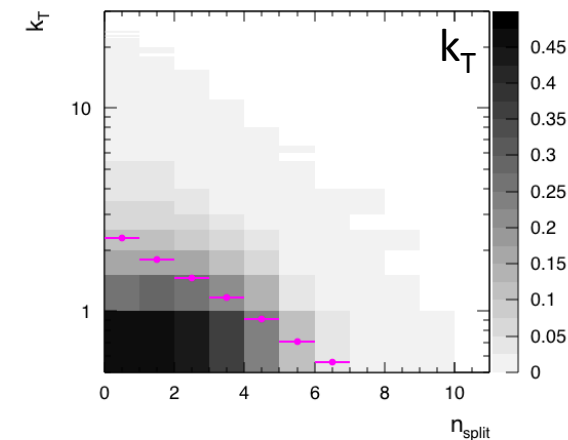
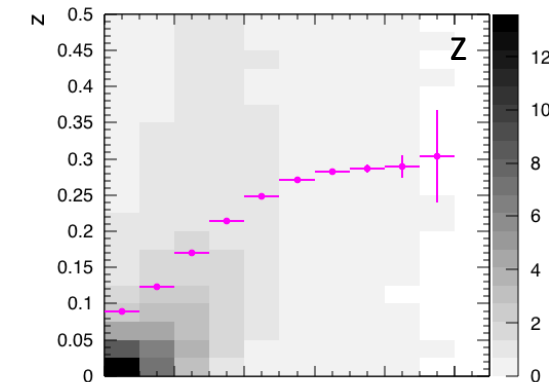
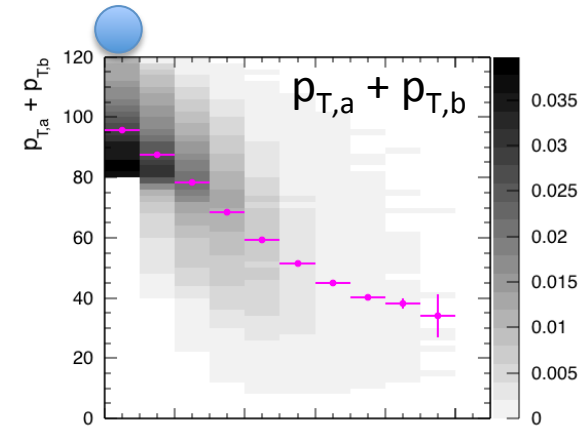
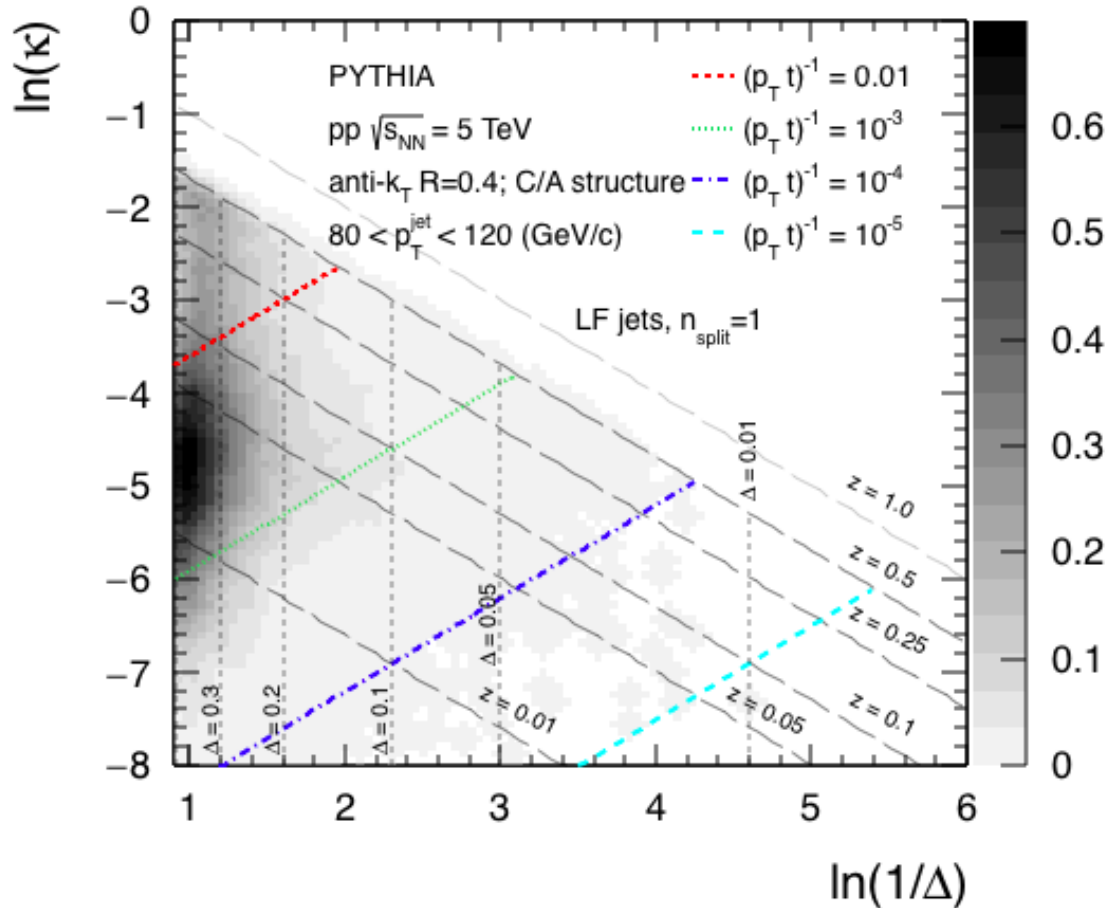
$$d\mathcal{P}_{\text{vac}} = 2 \frac{\alpha_s C_i}{\pi} d \log z \theta d \log \frac{1}{\theta}$$

Dreyer, Soye, Salam JHEP 1812 (2018) 064

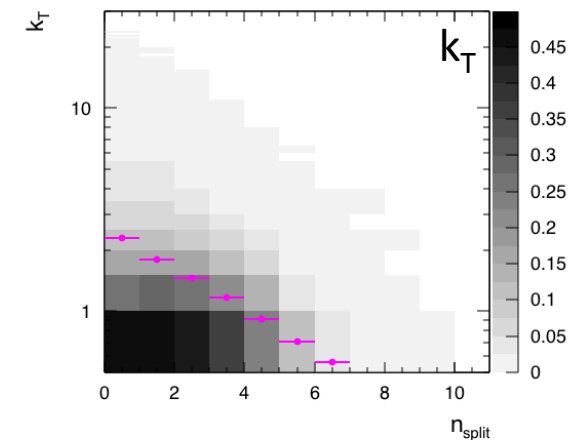
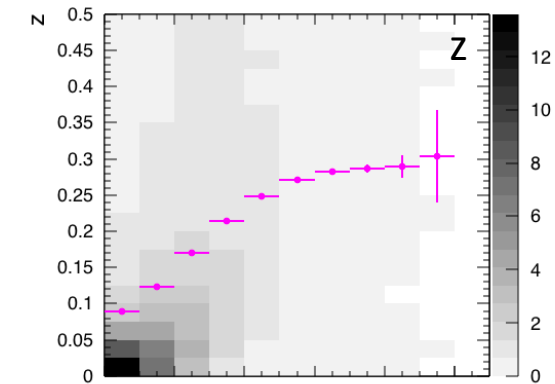
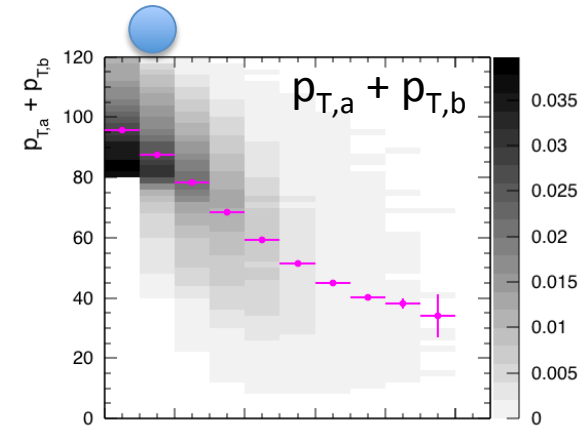
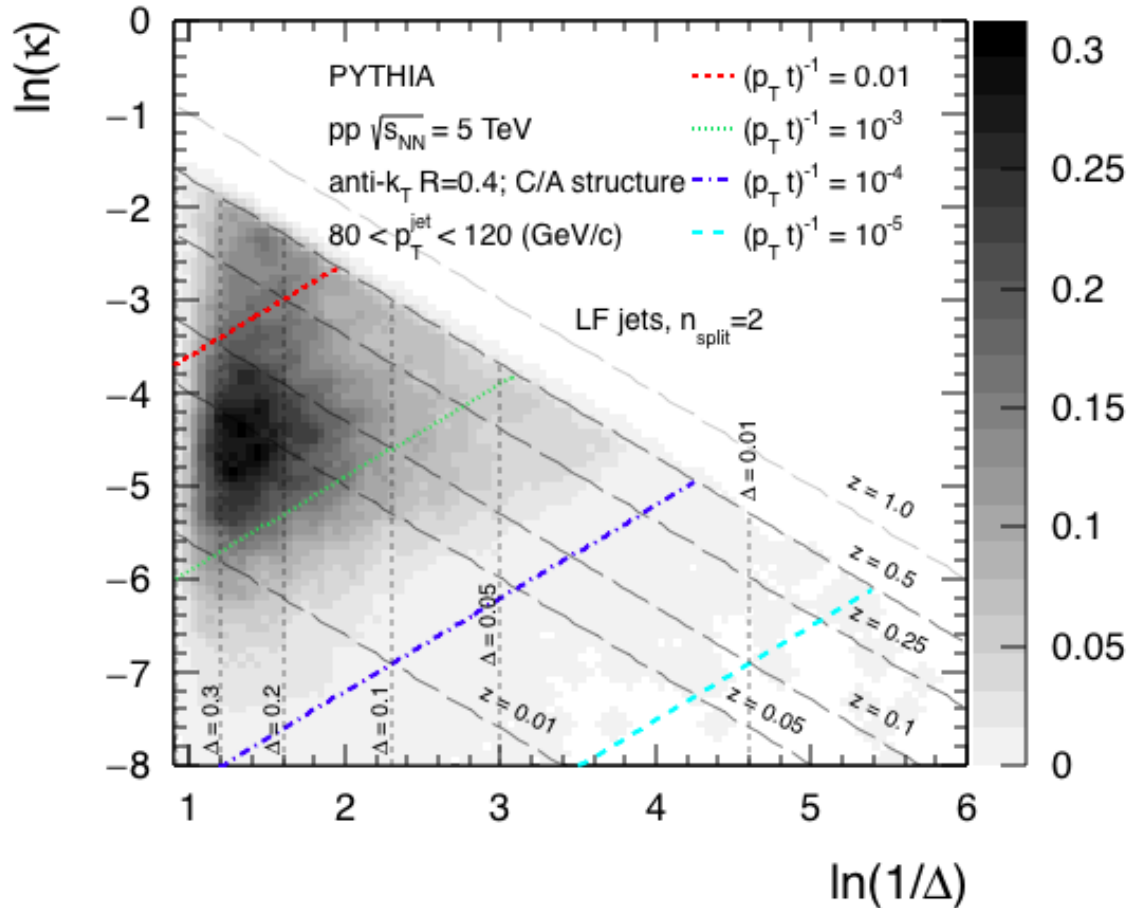
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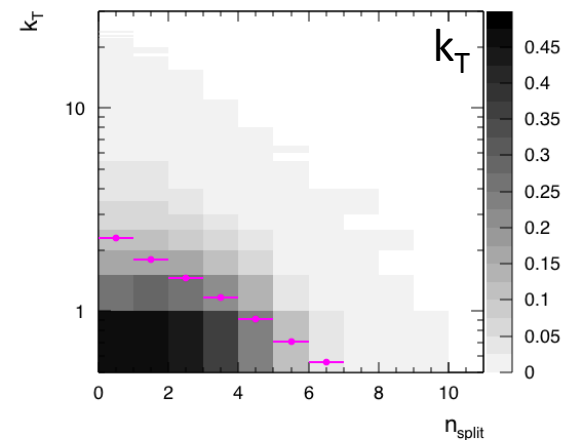
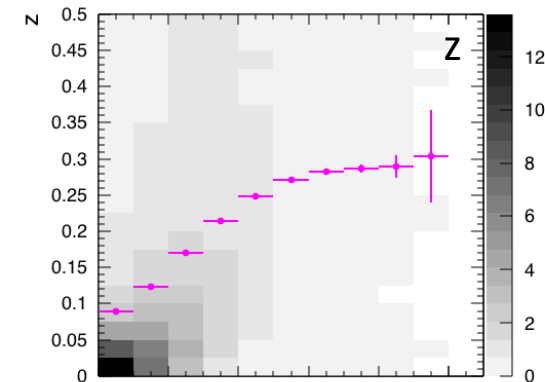
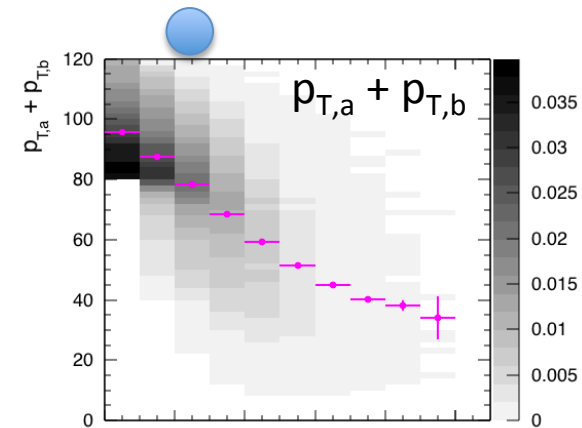
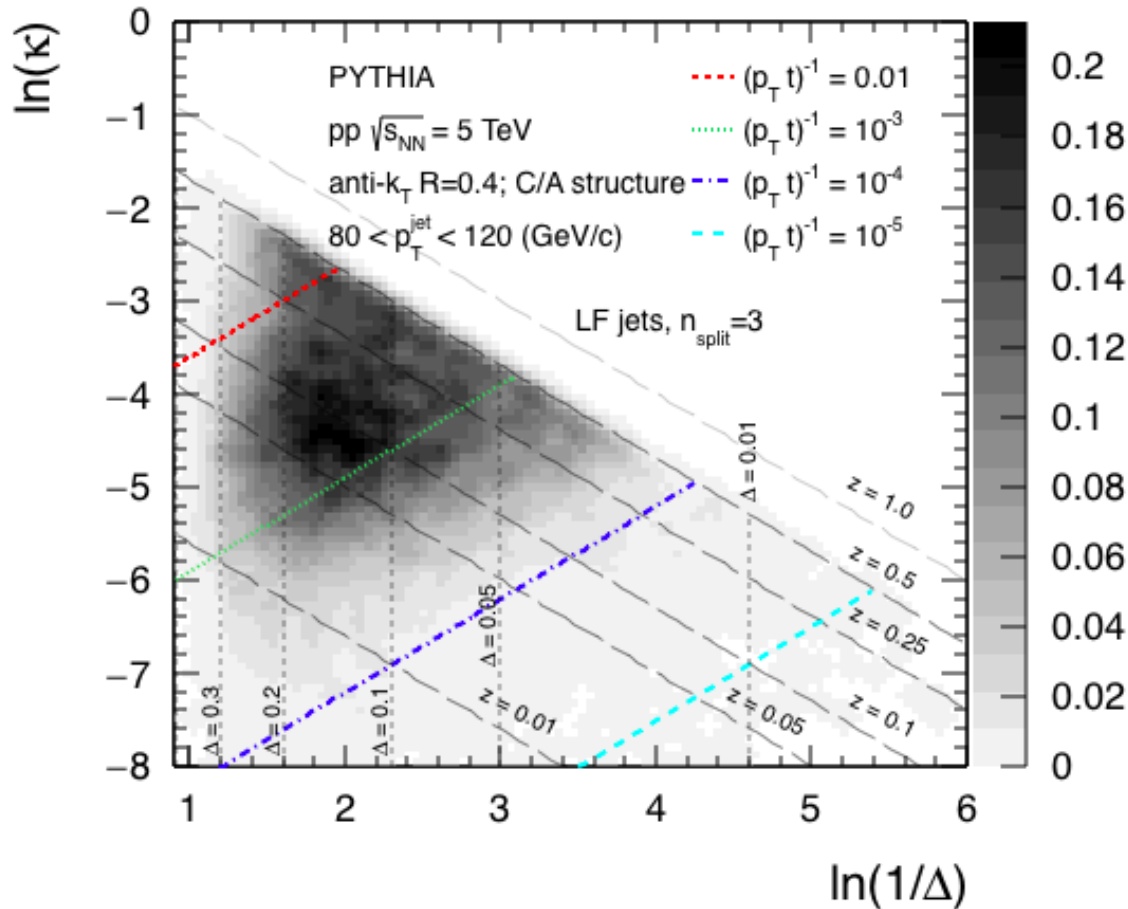
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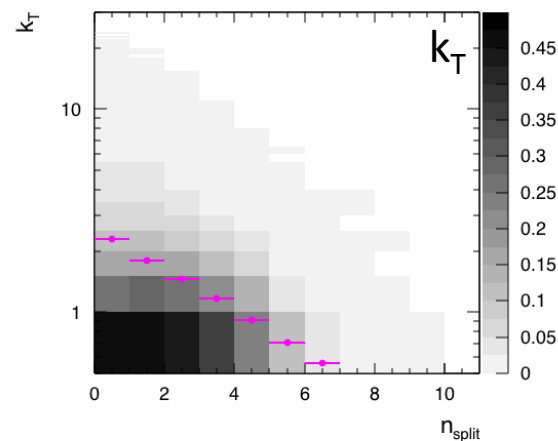
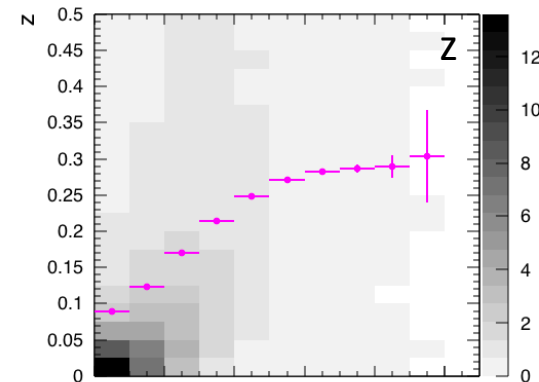
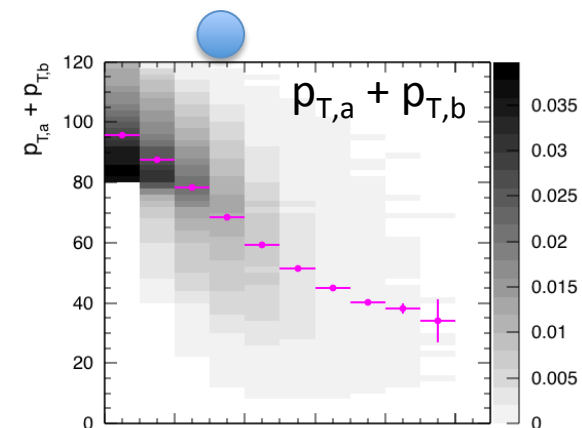
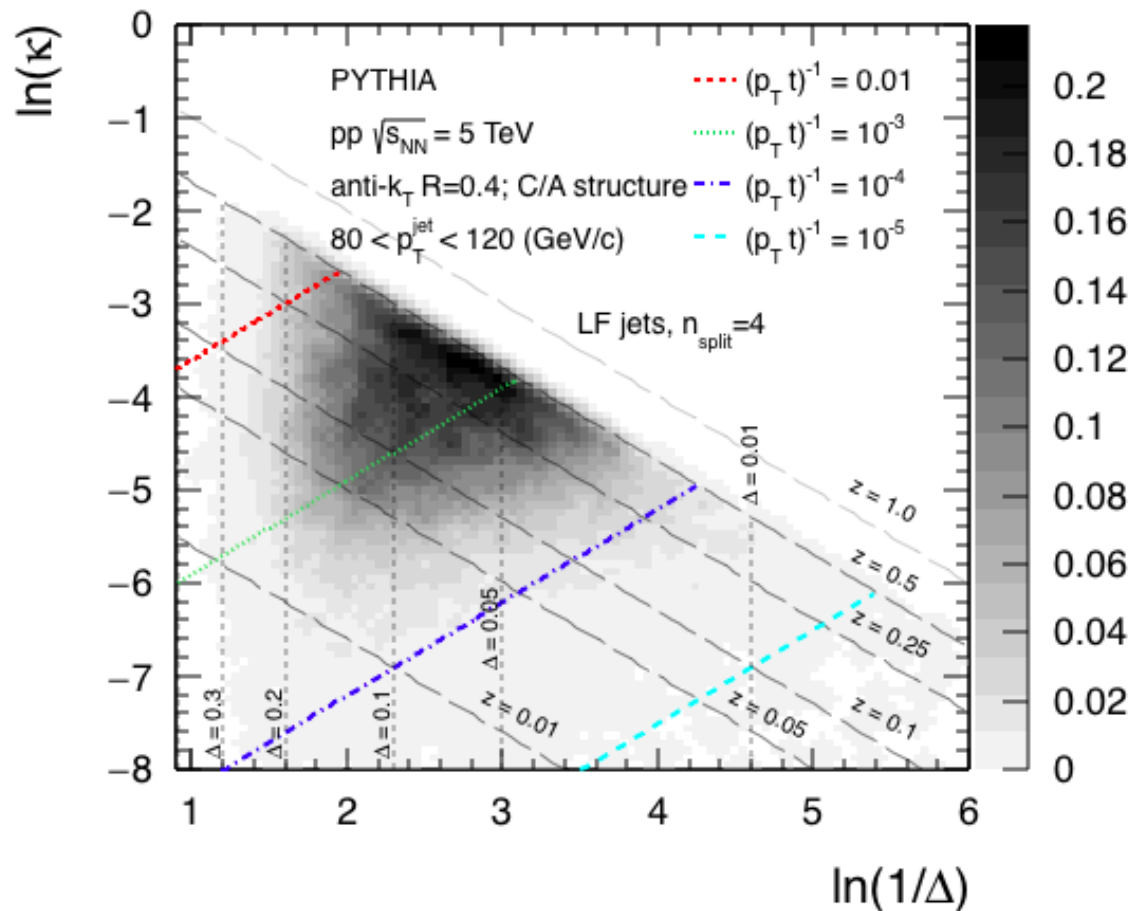
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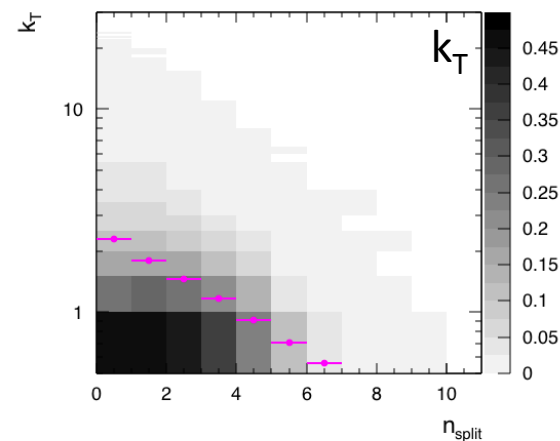
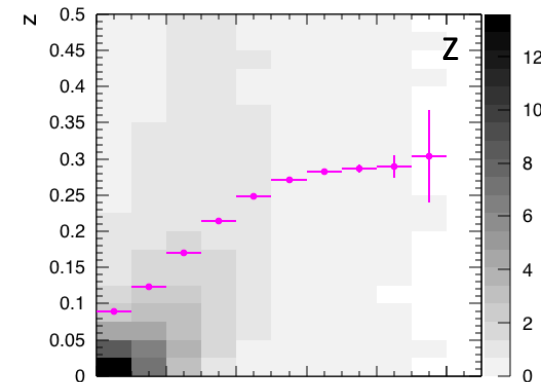
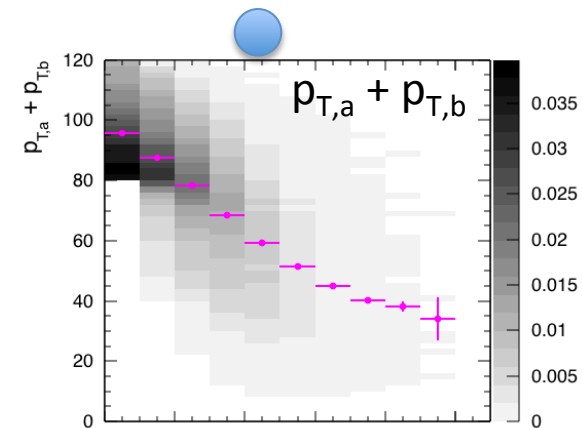
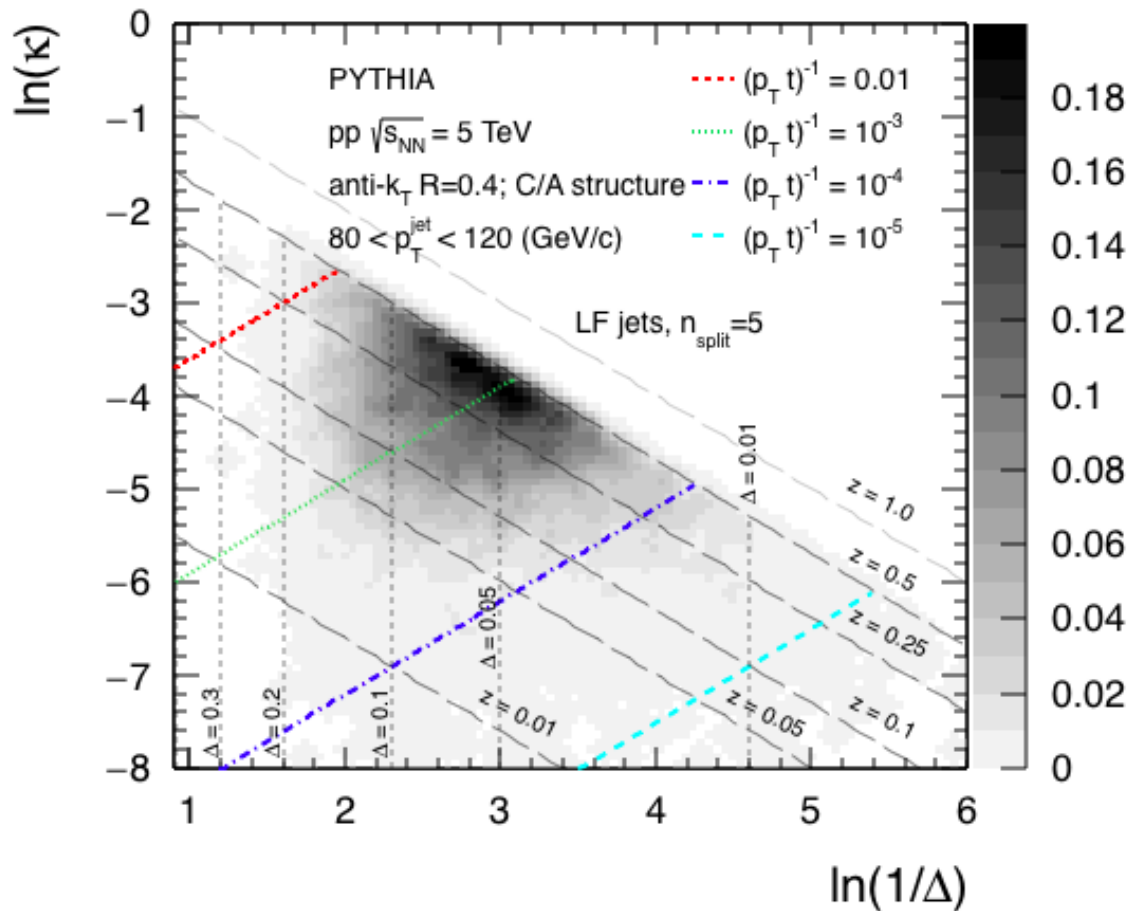
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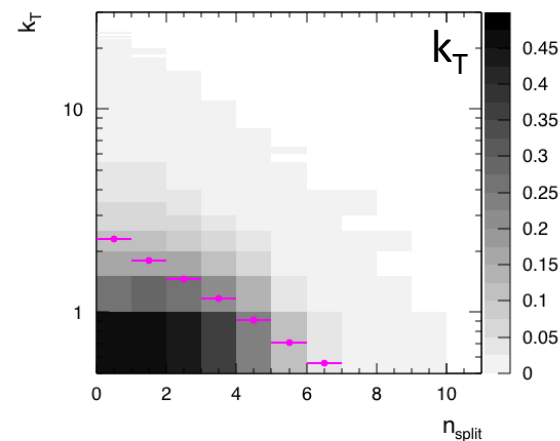
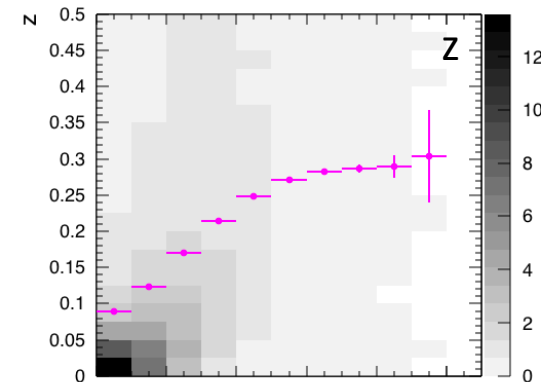
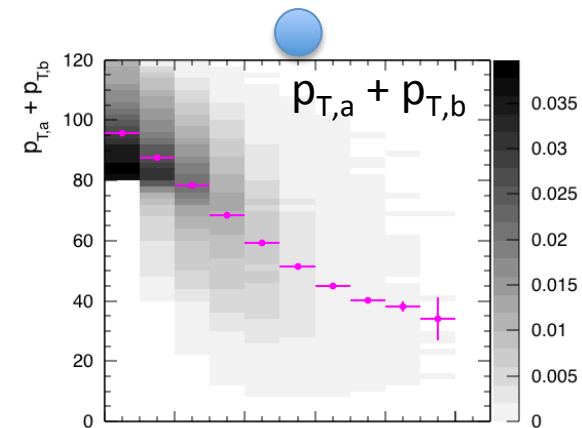
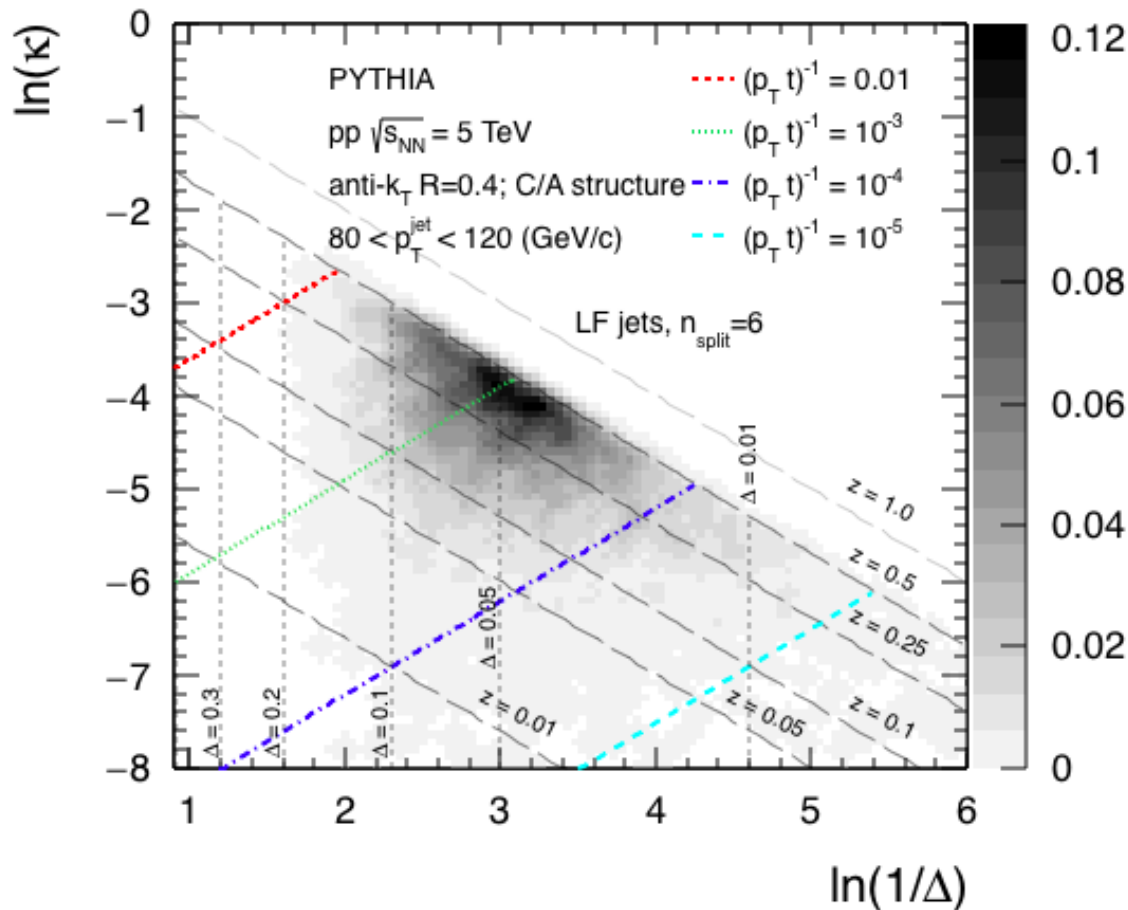
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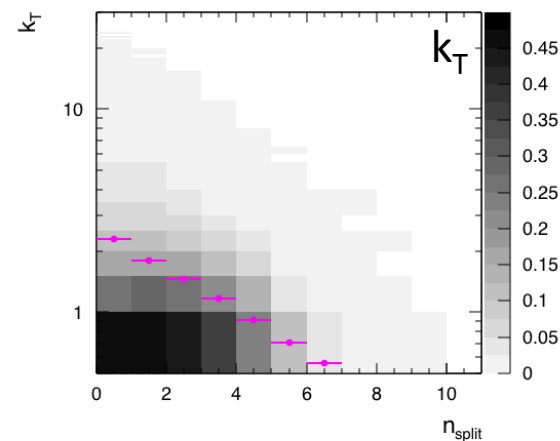
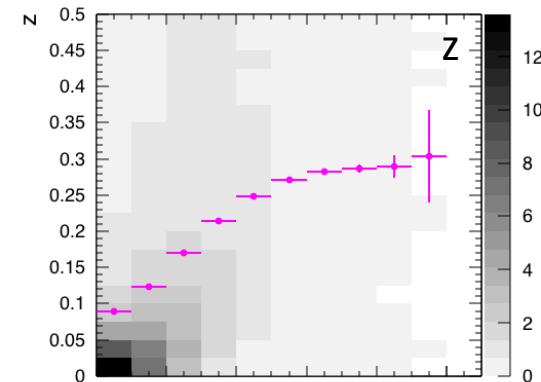
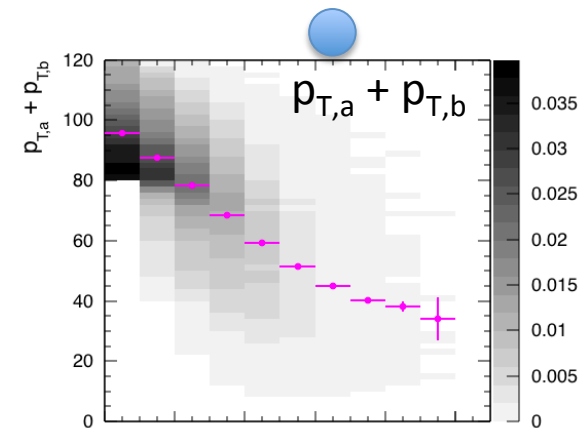
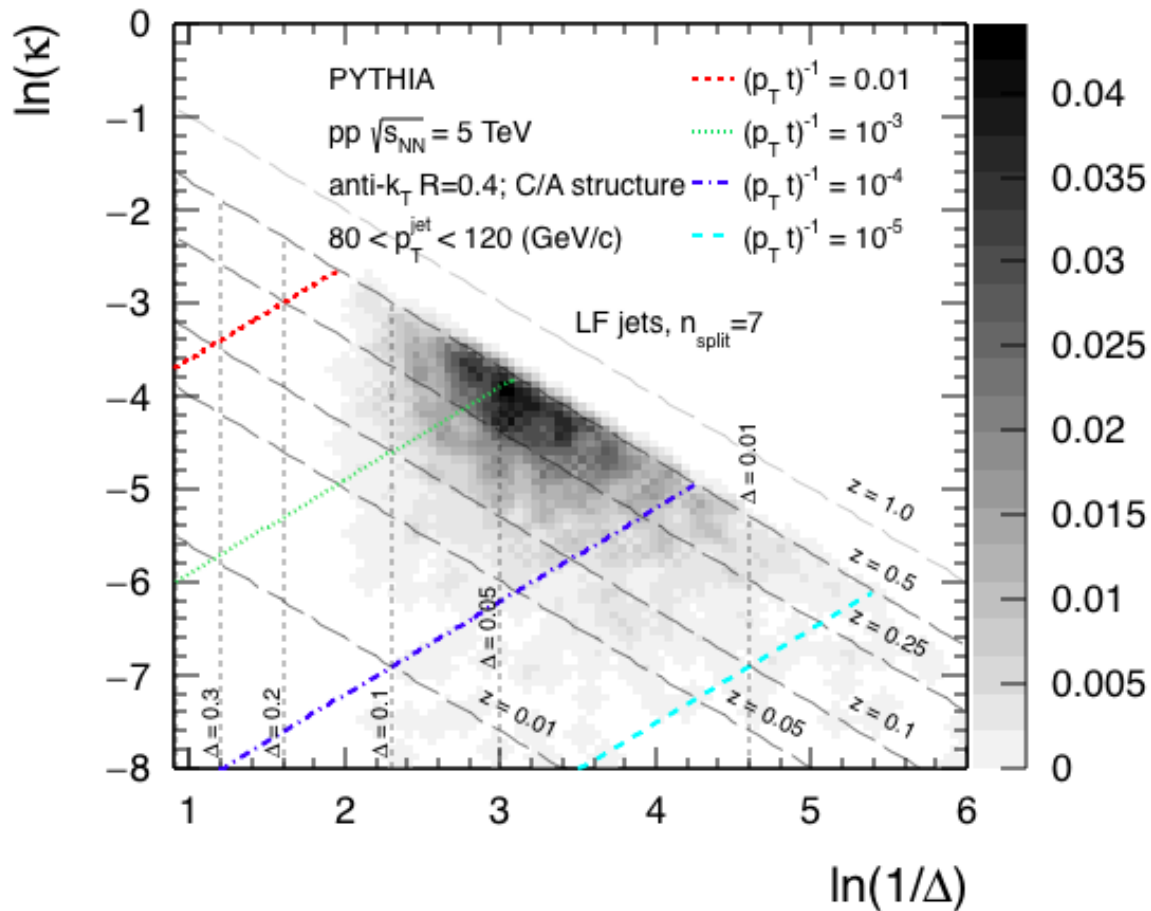
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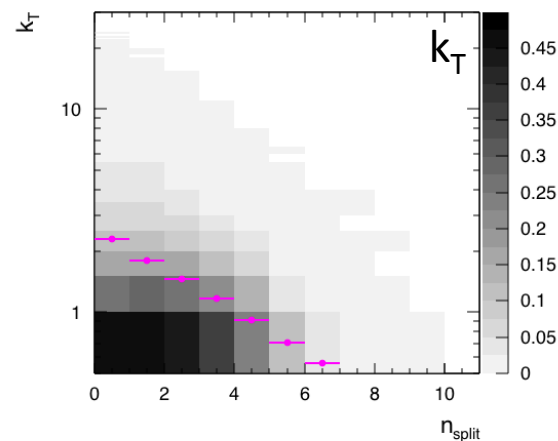
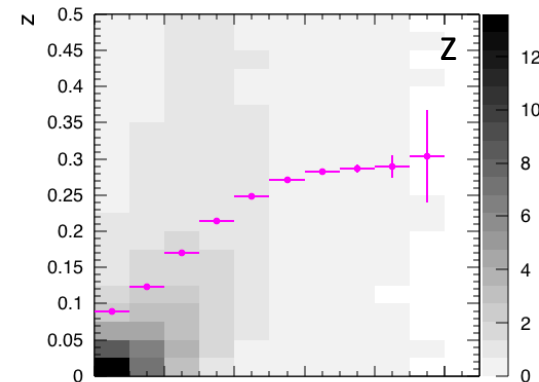
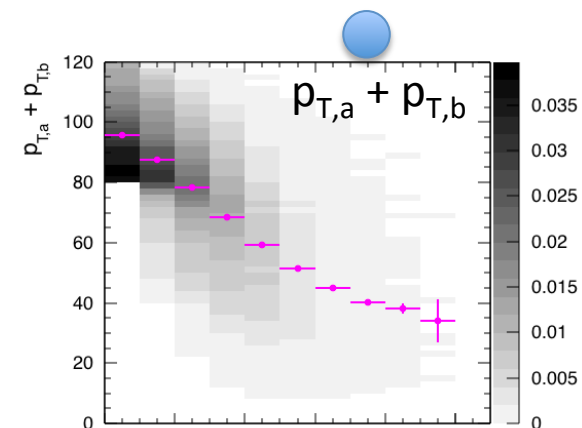
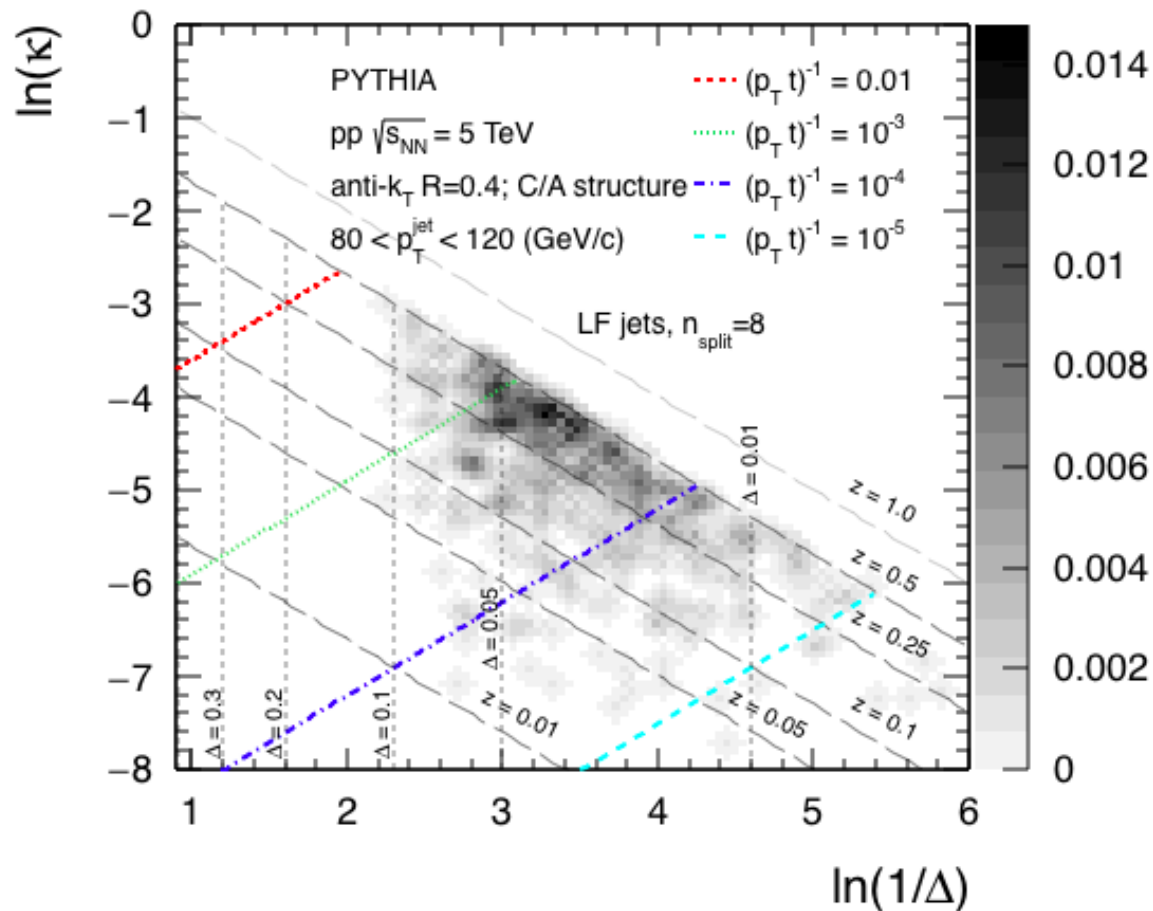
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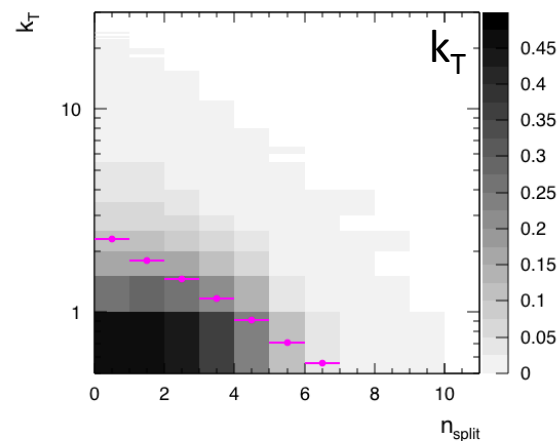
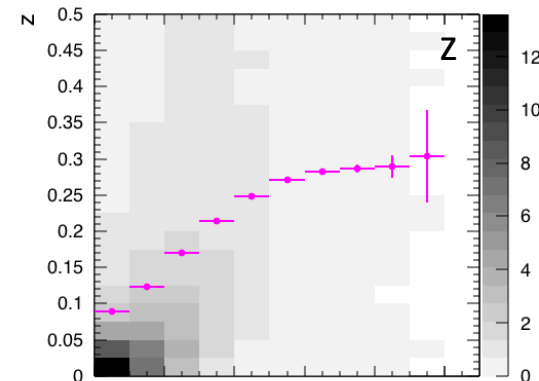
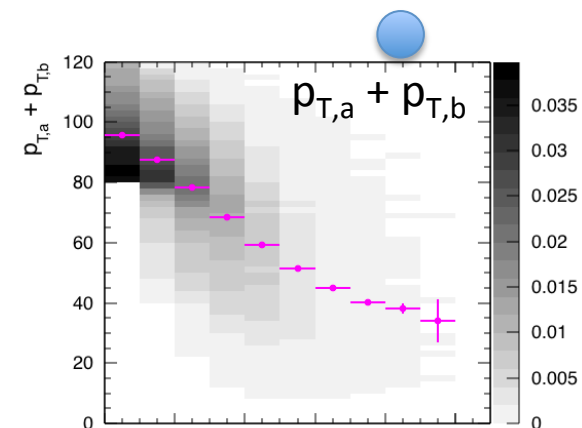
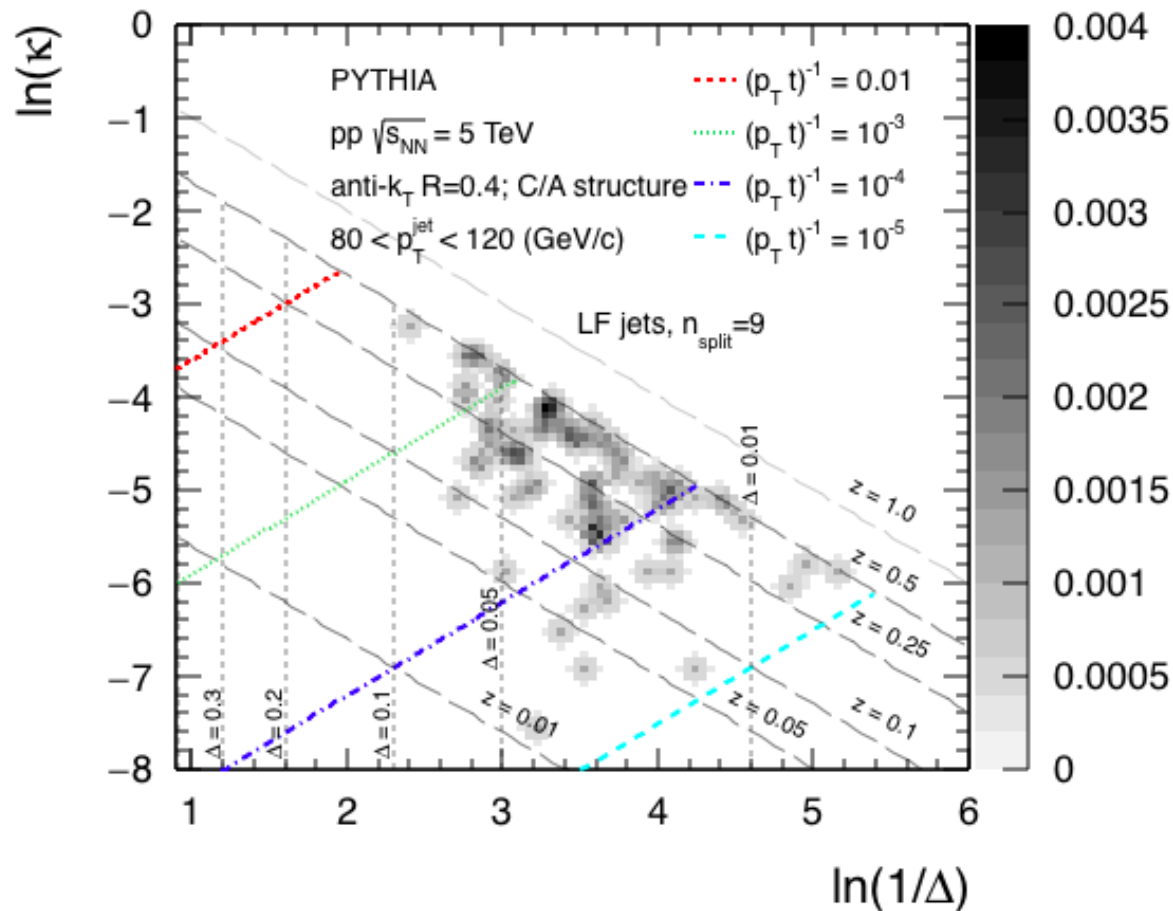
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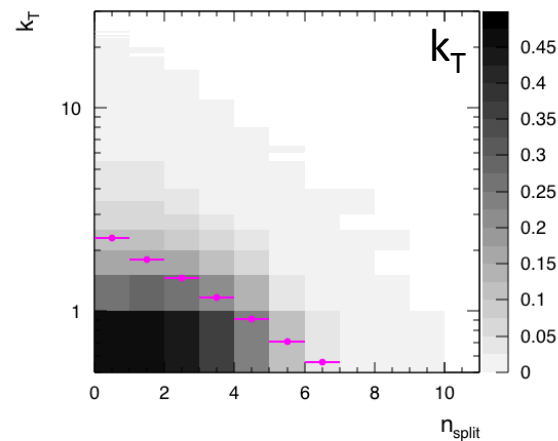
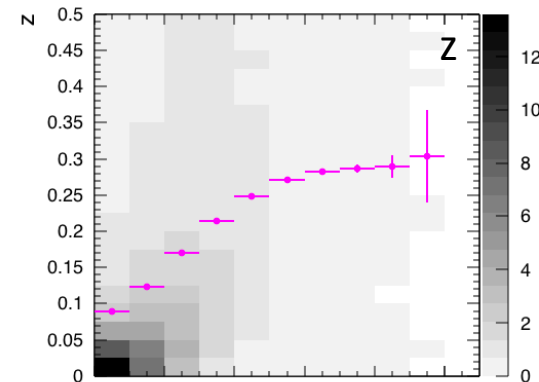
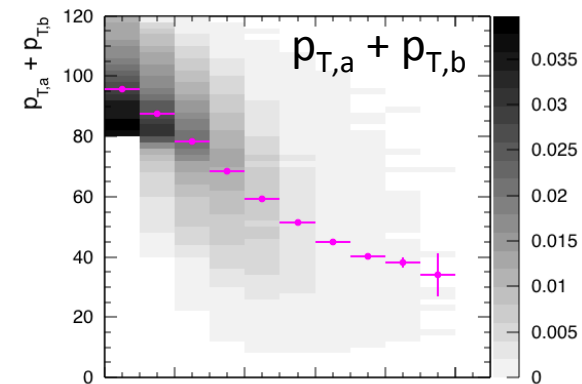
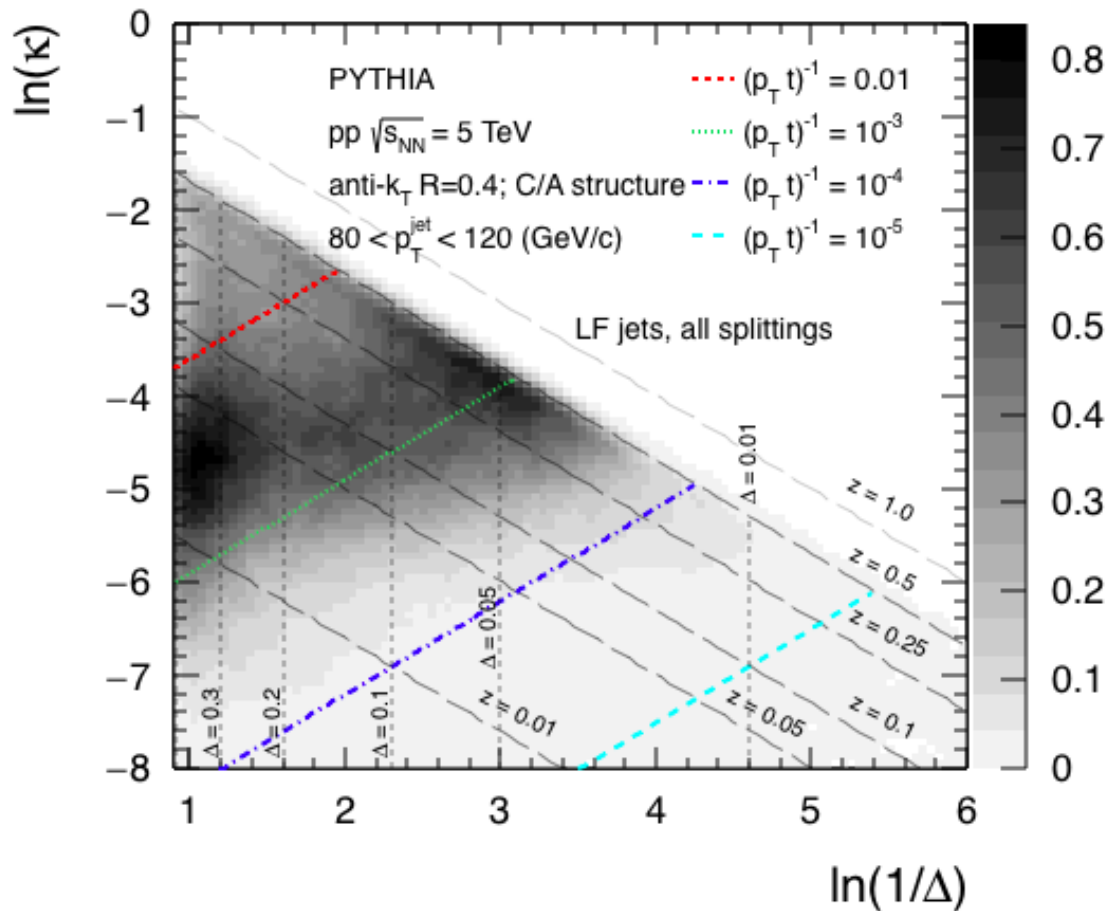
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Declusterization



Declusterization



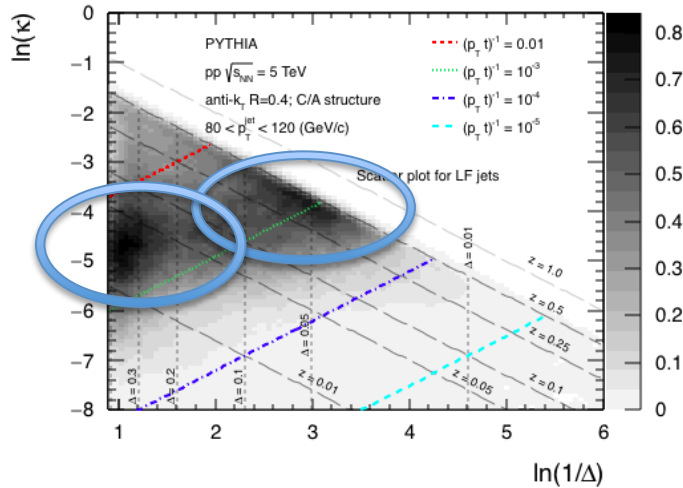
Jet Lund diagram

$$p_{T,a} > p_{T,b}, \kappa = \frac{p_{T,b}}{p_{T,a} + p_{T,b}} \Delta_{ab}^{17}$$

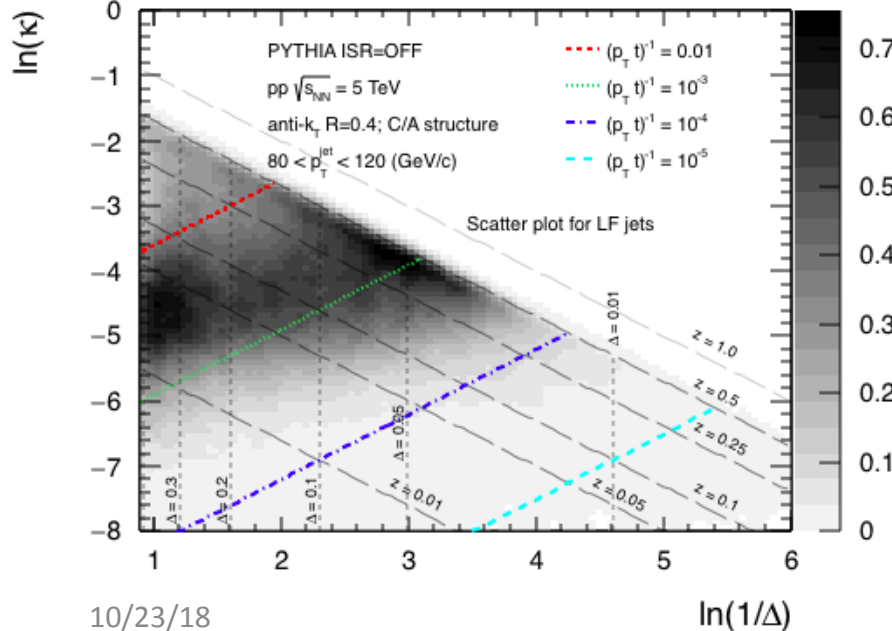
$$\bar{\rho}(\Delta, \kappa) = \frac{1}{N_{\text{jet}}} \frac{dn_{\text{emission}}}{d \ln \kappa d \ln 1/\Delta}$$

Expectation from theory: for every R and p_T , the phase space should be uniformly filled => the peak at the maximal angle is not expected.. The uniform distribution should be a feature of FSR.

- Potential feature of UE (MPIs) ? – TBD but not a show stopper at this point – treat MC as real data

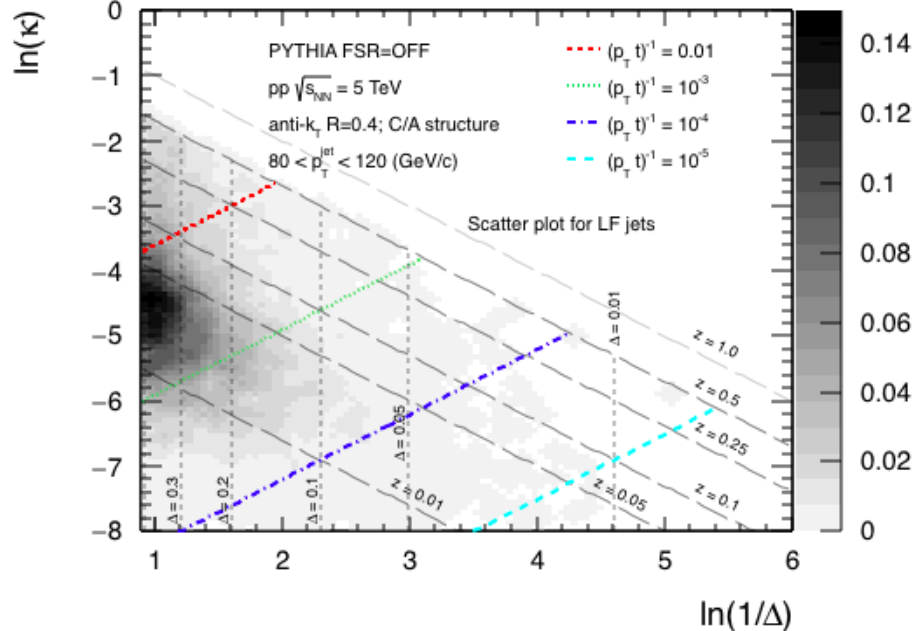


Initial State Radiation OFF



10/23/18

Final State Radiation OFF

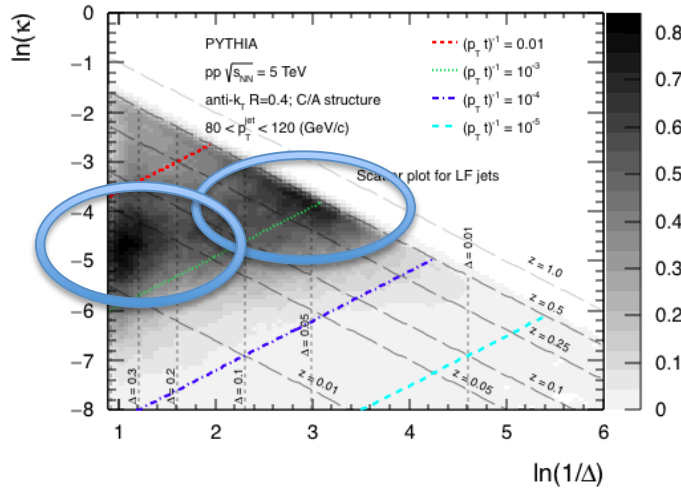


Tue 23/10/2018 09:32:18 CEST

Jet Lund diagram

$$p_{T,a} > p_{T,b}, \kappa = \frac{p_{T,b}}{p_{T,a} + p_{T,b}} \Delta_{ab}^{18}$$

$$\bar{\rho}(\Delta, \kappa) = \frac{1}{N_{\text{jet}}} \frac{dn_{\text{emission}}}{d \ln \kappa d \ln 1/\Delta}$$



Structures vary with the algorithm used

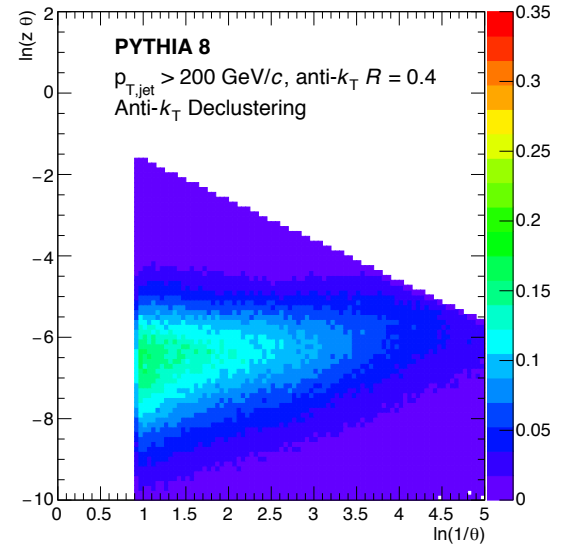
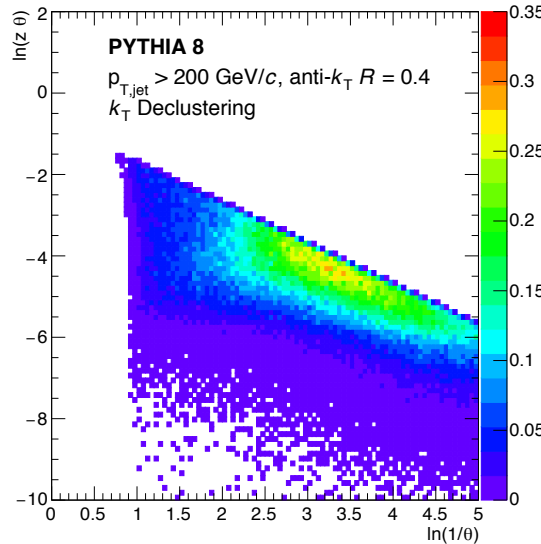
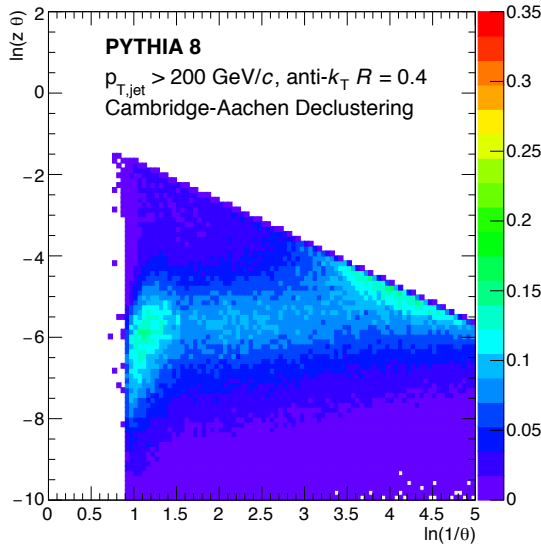
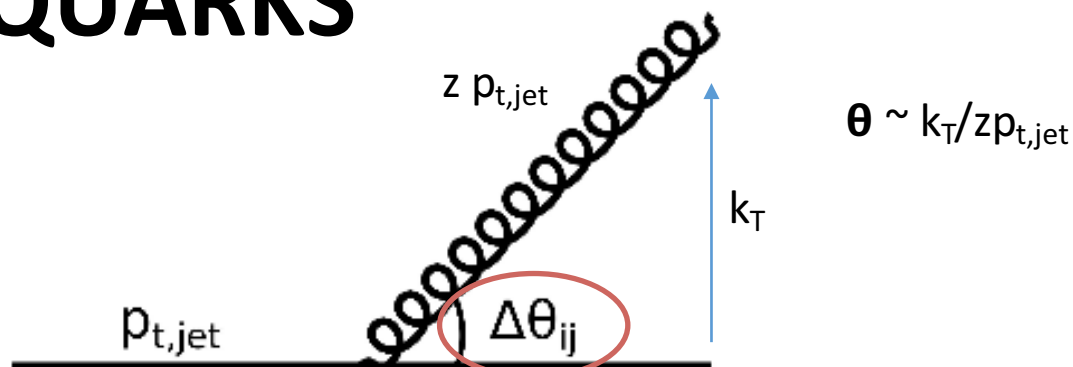


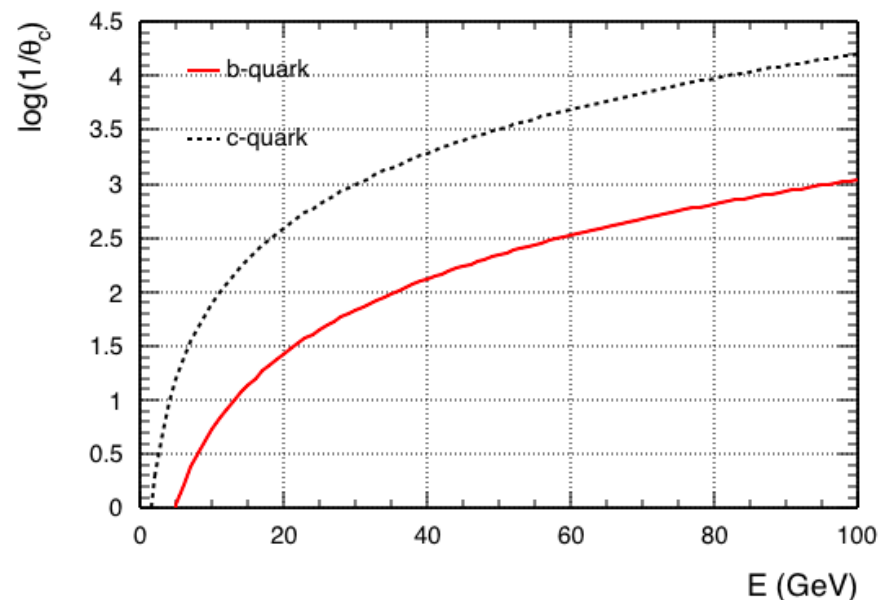
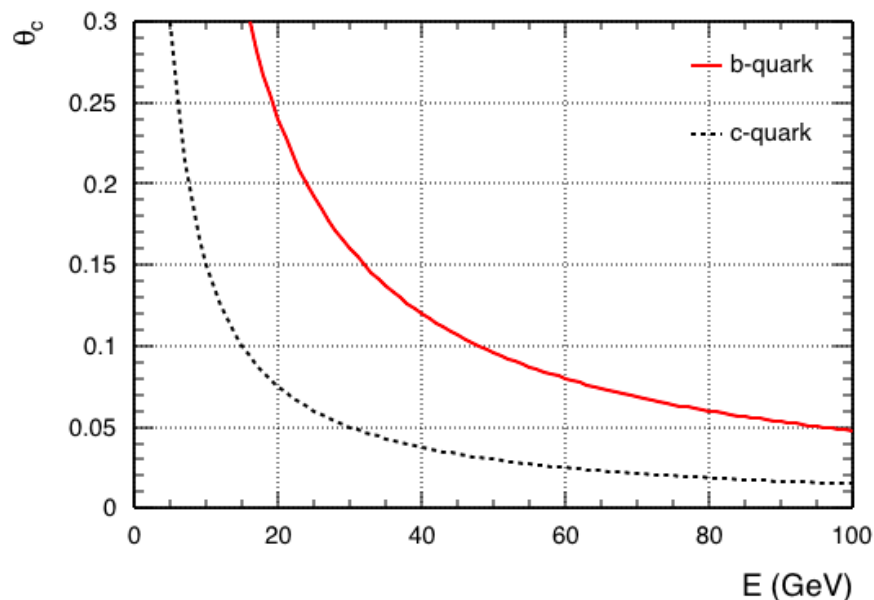
Figure 3: Lund diagrams reconstructed from a sample anti- k_T $R = 0.4$ jets generated by PYTHIA8. Three reclustering strategies were considered: C/A (left), k_T (middle), and anti- k_T (right).

ONTO HEAVY QUARKS

Hard gluons with restricted to radiate outside the dead cone (small k_T)



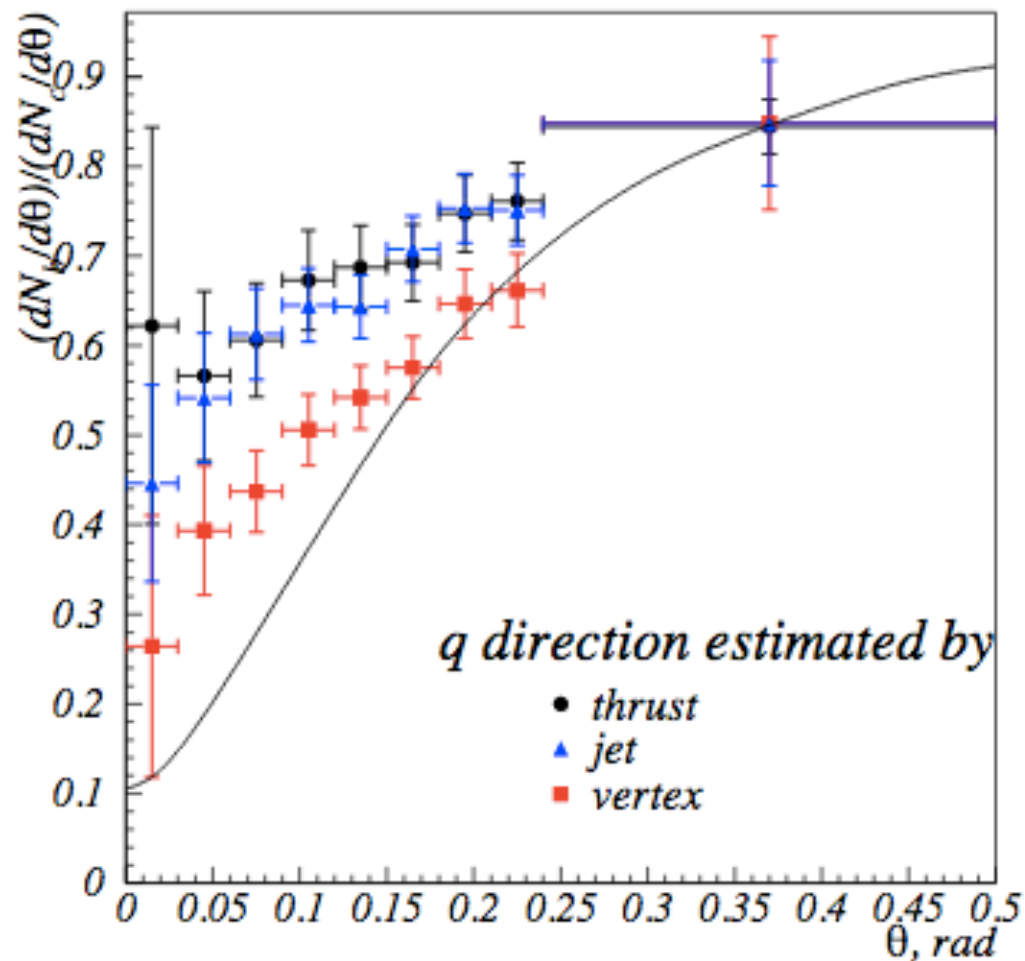
Dead cone for $\theta < m / E$



Dead cone measured?

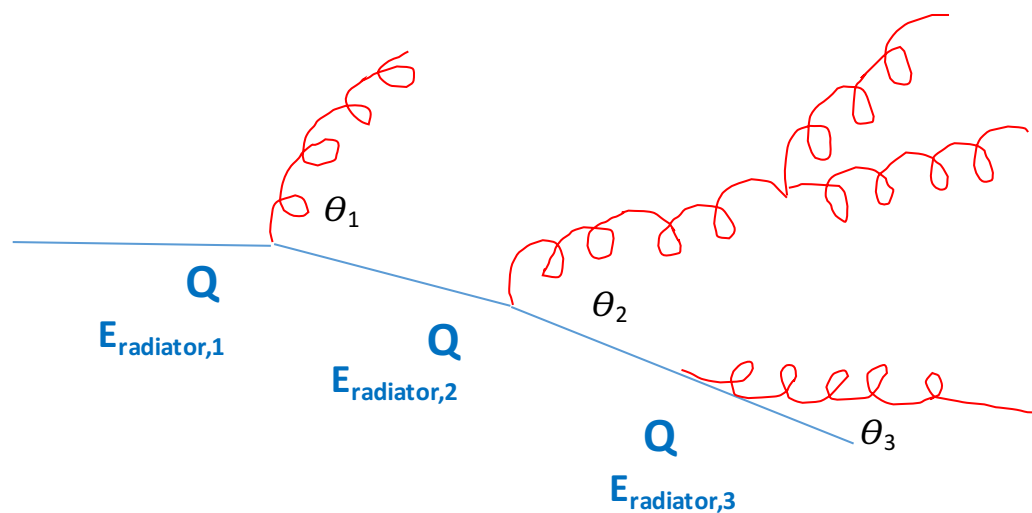
Battaglia, Orava, Sami "A study of depletion of fragmentation particles at small angles in b-jets with the DELPHI detector at LEP"

The first and only attempt of direct dead cone measurement



Recipe... 2/2

“tag & follow” heavy-flavors (hardest branch)...



We do the standard declustering process

We follow the prong containing the heavy flavor at each step

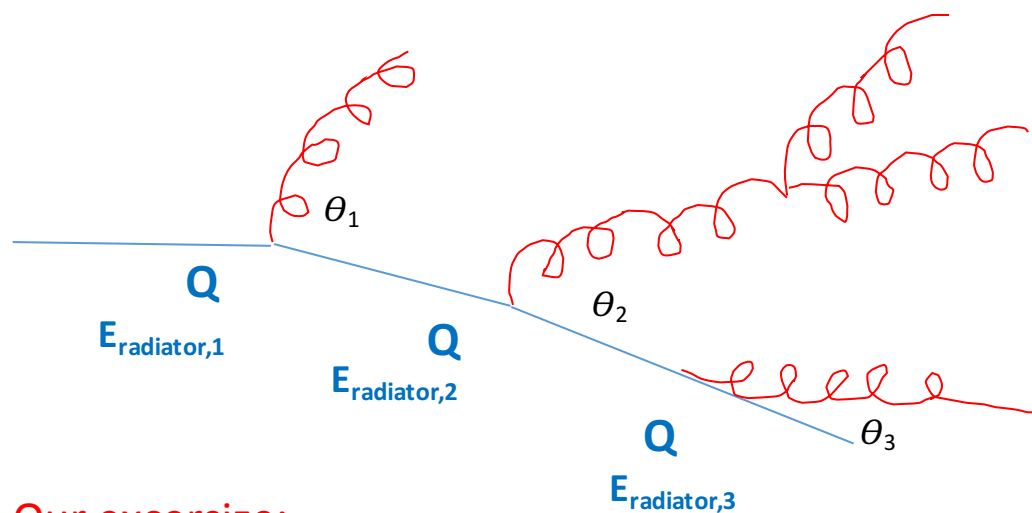
Negligible number of cases where the prong containing the heavy flavor is not the hardest, no ambiguity when comparing to light quarks

At each step of the jet tree we can access the angle of the splitting and the energy of the radiating parent subjet, E_{radiator} . For $\theta < m_Q/E_{\text{radiator}}$ we expect to see a suppression of splittings.

For low E_{radiator} , which means penetrating the jet to deep levels, the phase space for the dead cone observation grows.

Recipe... 2/2

“tag & follow” heavy-flavors (hardest branch)...



We do the standard declustering process

We follow the prong containing the heavy flavor at each step

Negligible number of cases where the prong containing the heavy flavor is not the hardest, no ambiguity when comparing to light quarks

Our exercise:

- Pythia8 Tune 4C
- pp collisions at $\sqrt{s}=13$ TeV
- Select $c\bar{c}$, $b\bar{b}$ production channels
- At parton level the leading prong at each declustering step is the one containing the heavy parton
- At hadron level, we inhibit the decay of the B and D mesons and the leading prong at each declustering step the one containing the heavy meson.
- Antik_T jets $R=0.4$, CA reclustering algo

LF vs HQ – parton level – no UE, no ISR

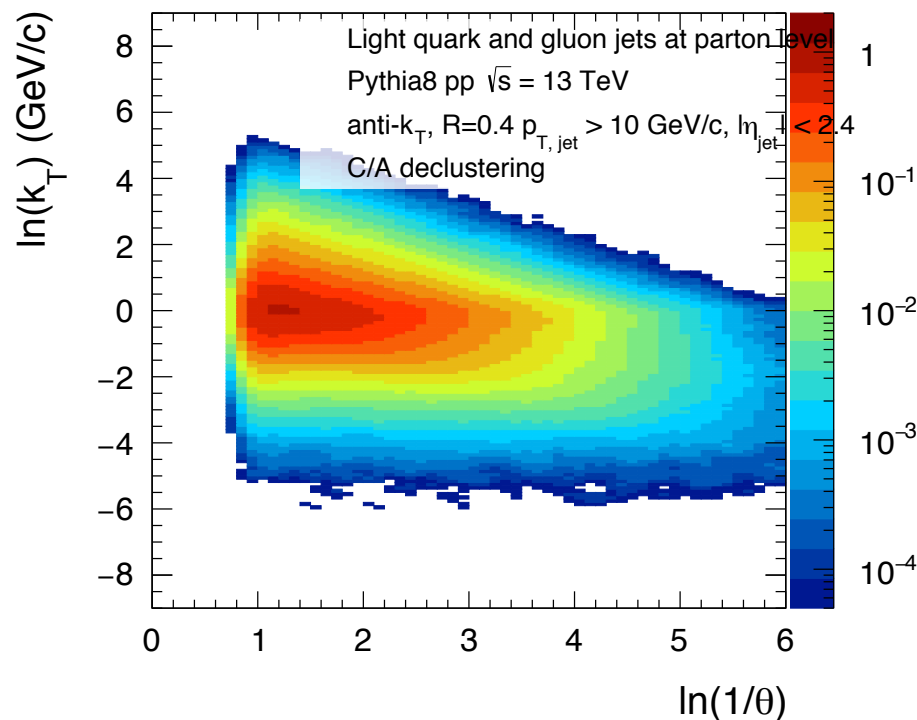
All jets (light quark and gluon)

$$\log(kT) = -2 \Leftrightarrow kT = 0.135 \text{ GeV}$$

$$\log(kT) = 0 \Leftrightarrow kT = 1 \text{ GeV}$$

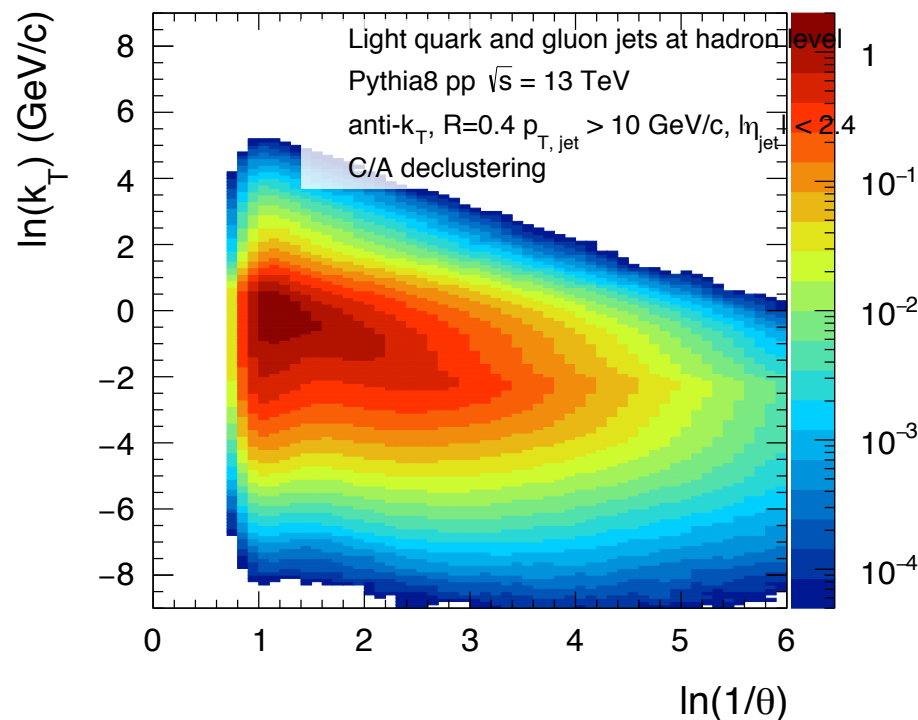
Parton level

Lund diagram $\rho(k_T, \theta) = 1/N_{\text{jet}} d^2N / d \ln(k_T) / d \ln(1/\theta) (\text{GeV/c})^{-1}$



Hadron level

Lund diagram $\rho(k_T, \theta) = 1/N_{\text{jet}} d^2N / d \ln(k_T) / d \ln(1/\theta) (\text{GeV/c})^{-1}$



Tue 29/01/2019 11:48:36 PST

Tue 29/01/2019 11:48:36 PST

Hadron level – ALL processes

LF vs HQ – parton level – no UE, no ISR

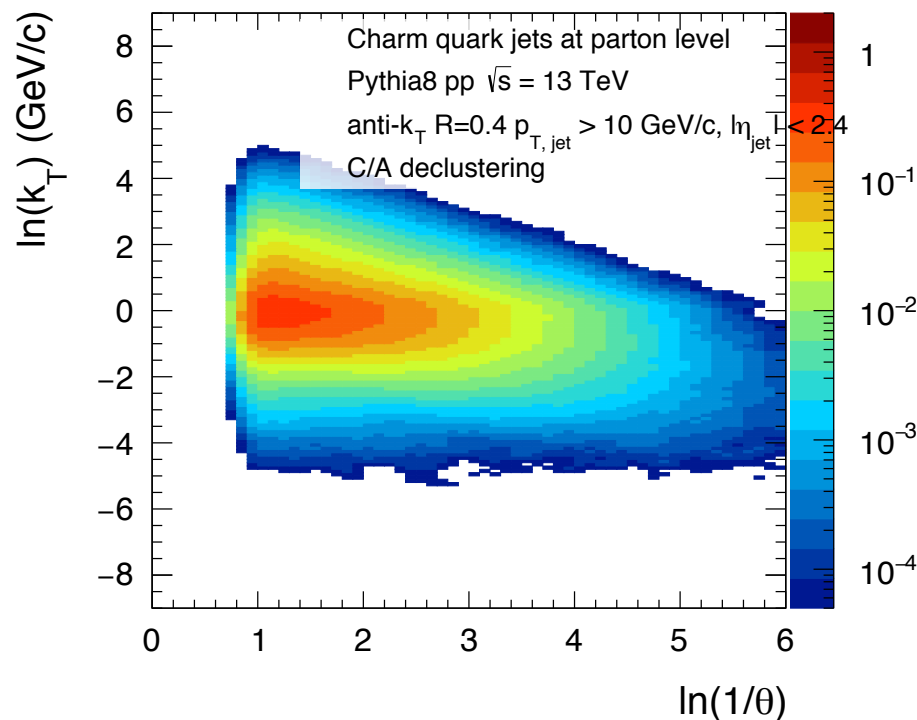
Charm

$$\log(kT) = -2 \Leftrightarrow kT = 0.135 \text{ GeV}$$

$$\log(kT) = 0 \Leftrightarrow kT = 1 \text{ GeV}$$

Parton level

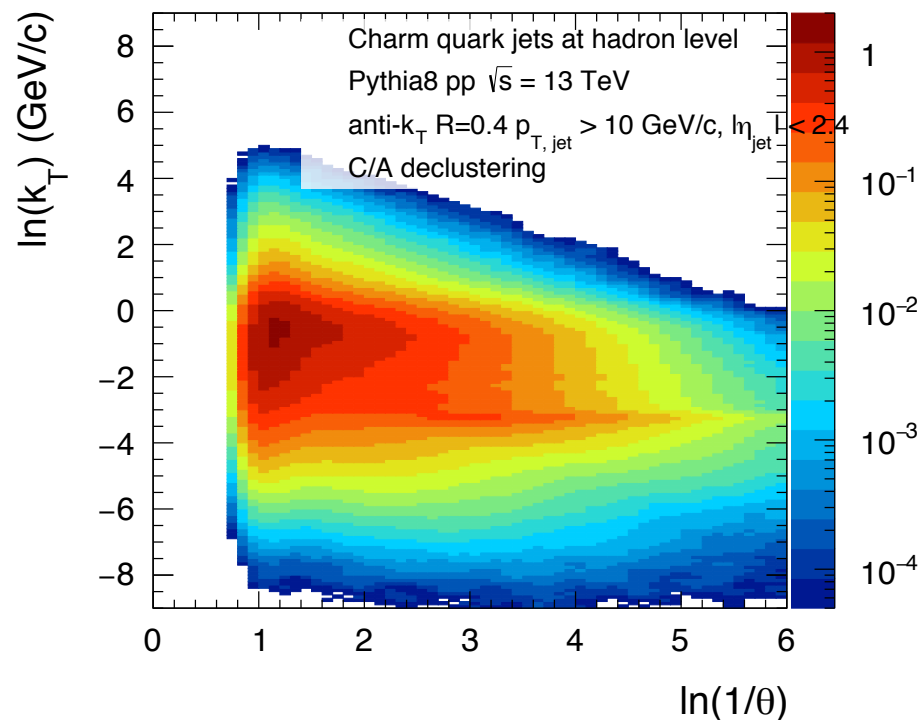
Lund diagram $\rho(k_T, \theta) = 1/N_{\text{jet}} d^2N / d \ln(k_T) / d \ln(1/\theta) (\text{GeV}/c)^{-1}$



Tue 29/01/2019 11:48:36 PST

Hadron level

Lund diagram $\rho(k_T, \theta) = 1/N_{\text{jet}} d^2N / d \ln(k_T) / d \ln(1/\theta) (\text{GeV}/c)^{-1}$



Tue 29/01/2019 11:48:37 PST

Hadron level – ALL processes

LF vs HQ – parton level – no UE, no ISR

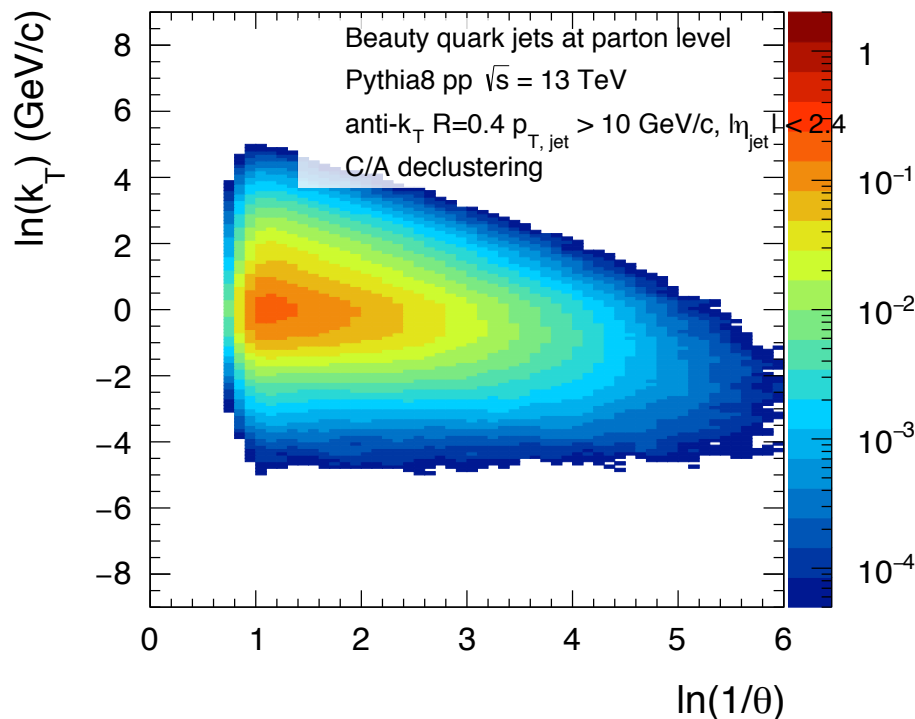
Beauty

$$\log(kT) = -2 \Leftrightarrow kT = 0.135 \text{ GeV}$$

$$\log(kT) = 0 \Leftrightarrow kT = 1 \text{ GeV}$$

Parton level

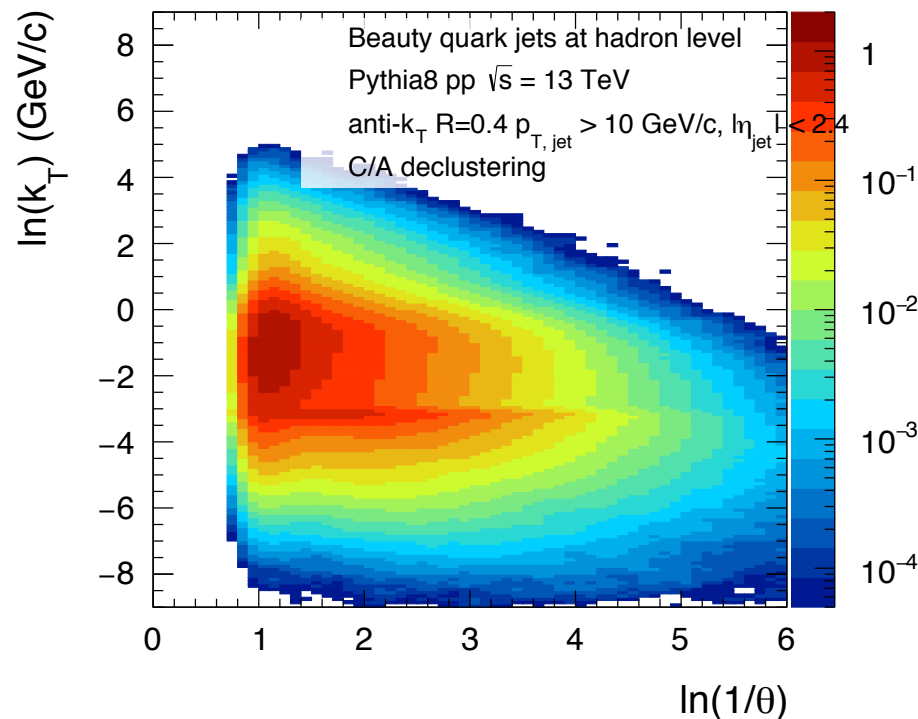
Lund diagram $\rho(k_T, \theta) = 1/N_{\text{jet}} d^2N / d \ln(k_T) / d \ln(1/\theta) (\text{GeV}/c)^{-1}$



Tue 29/01/2019 11:48:36 PST

Hadron level

Lund diagram $\rho(k_T, \theta) = 1/N_{\text{jet}} d^2N / d \ln(k_T) / d \ln(1/\theta) (\text{GeV}/c)^{-1}$



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Hadron level – ALL processes

LF vs HQ – parton level – no UE, no ISR

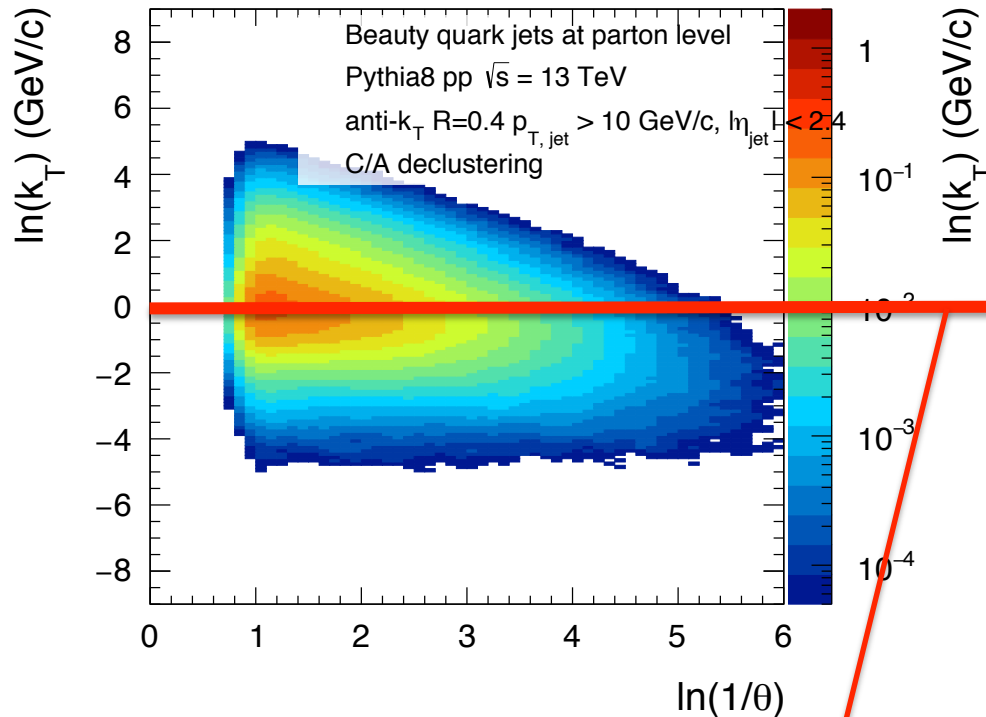
Beauty

$$\log(kT) = -2 \Leftrightarrow kT = 0.135 \text{ GeV}$$

$$\log(kT) = 0 \Leftrightarrow kT = 1 \text{ GeV}$$

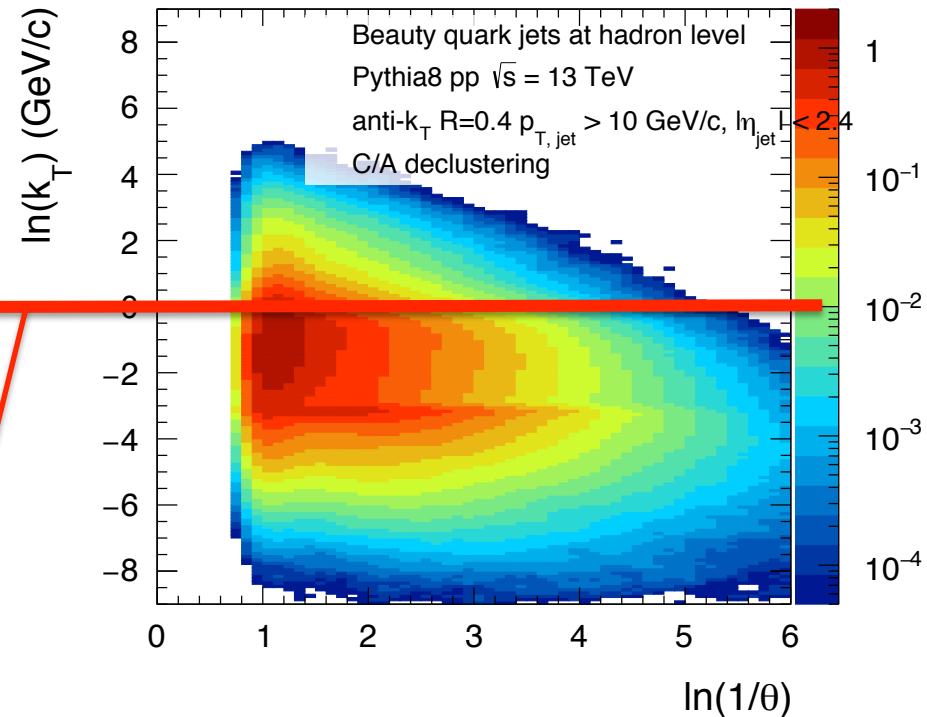
Parton level

Lund diagram $\rho(k_T, \theta) = 1/N_{\text{jet}} d^2N / d \ln(k_T) / d \ln(1/\theta) (\text{GeV}/c)^{-1}$



Hadron level

Lund diagram $\rho(k_T, \theta) = 1/N_{\text{jet}} d^2N / d \ln(k_T) / d \ln(1/\theta) (\text{GeV}/c)^{-1}$



Hadron level – ALL processes

Non-perturbative effects can be removed/isolated by cutting the region $\ln(k_T) < 0$

Dead cone on parton level

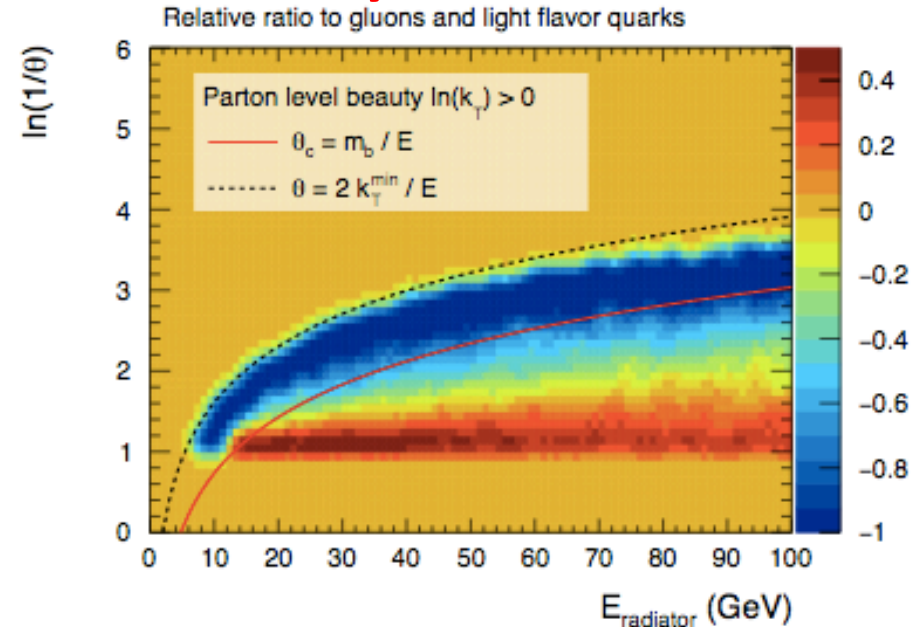
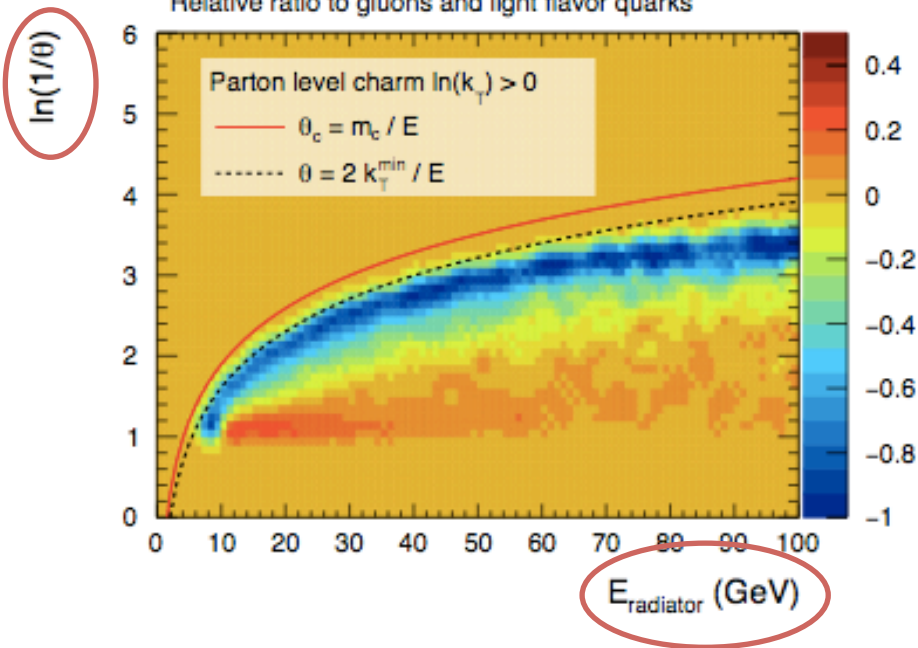
Parton level – no UE, no ISR

Angle vs. Energy of the radiator

Cut on $\log(k_T) > 0$
 $E_{\text{radiator}} > 1/(z\theta)$

charm

beauty



$$Q = \frac{P^Q(\log(1/\theta), E_{\text{radiator}}) - P^{\text{inc}}(\log(1/\theta), E_{\text{radiator}})}{P^{\text{inc}}(\log(1/\theta), E_{\text{radiator}})}$$

The cut $\log(k_T) > k_T^{\min}$ translates into $E_{\text{radiator}} > k_T^{\min} / z\theta$. The black dashed line corresponds to the kinematic limit of $z=0.5$, above which there are no more entries in the inclusive reference

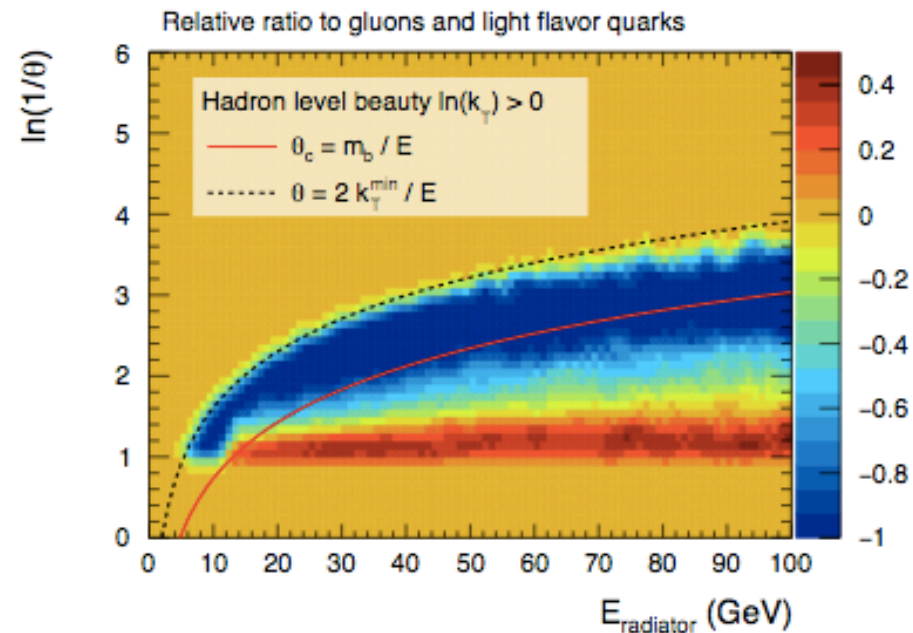
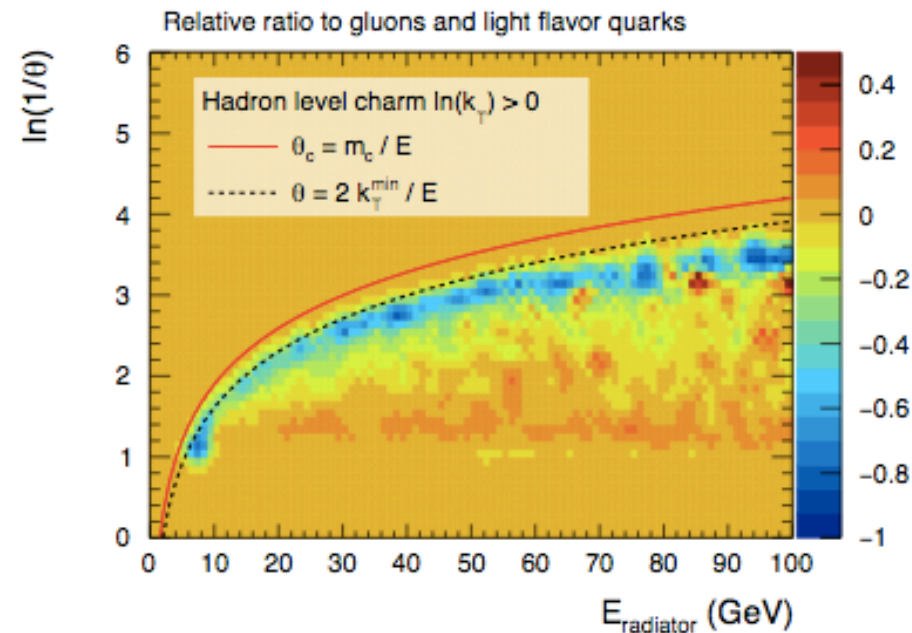
We can identify regions of phase space where $P^Q = -1$, meaning no radiation off the heavy quark radiator. **The suppression of large angles couples to the suppression of large z**

Dead cone at hadron level

Hadron level – ALL processes

charm

beauty



$$Q = \frac{P^Q(\log(1/\theta), E_{\text{radiator}}) - P^{\text{inc}}(\log(1/\theta), E_{\text{radiator}})}{P^{\text{inc}}(\log(1/\theta), E_{\text{radiator}})}$$

At hadron level the effects are smeared but not washed out

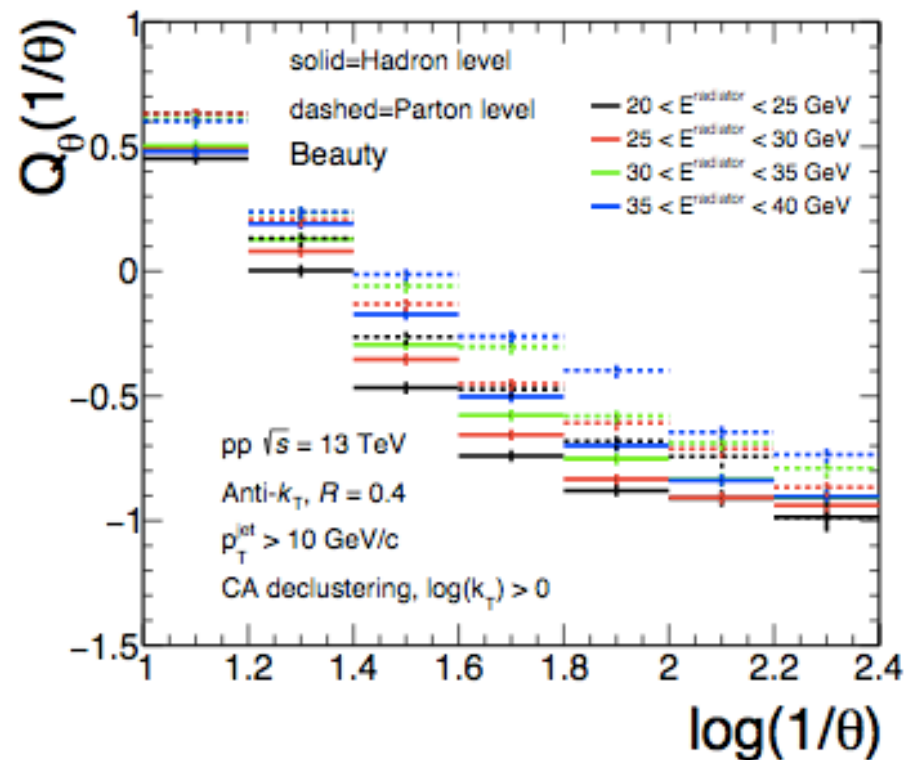
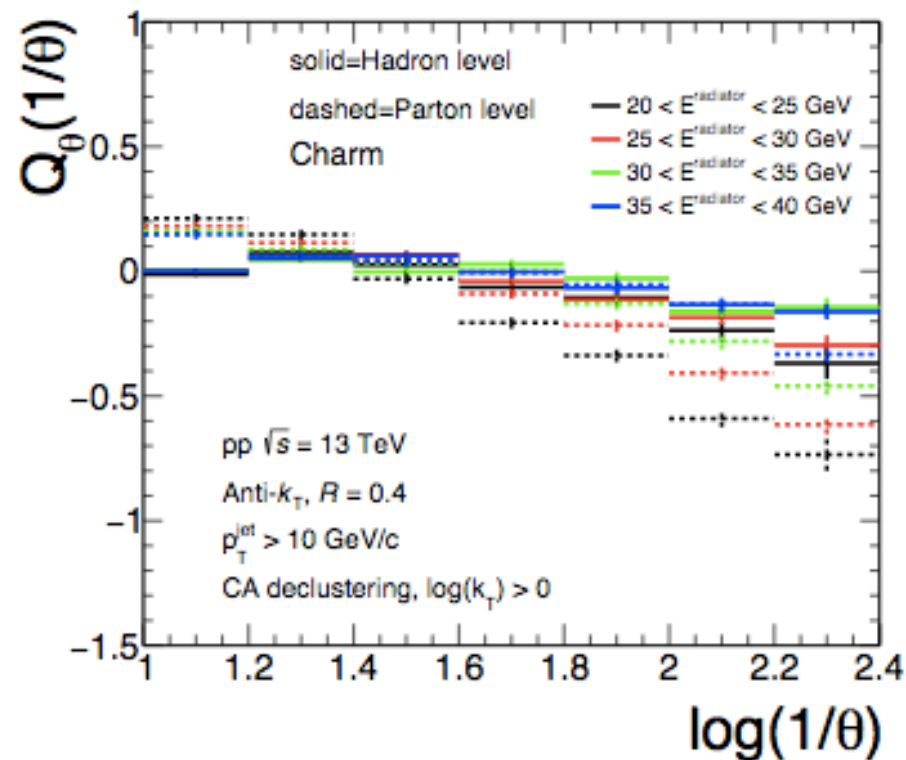
As expected: the higher E_{radiator} is, the dead cone effects appear at smaller angles.

For D jets, the effects appear at measurable angles of ~ 0.1 rad for radiator energies of 10-30 GeV

For B jets, one can go higher in radiator energy and still have effects at angles of the order of 0.1 rad

The observable – projections for E

$$Q_\theta = \frac{P^Q(1/\theta) - P^{inc}(1/\theta)}{P^{inc}(1/\theta)}, E_{\text{radiator}} \in (E_{\text{min}}, E_{\text{max}})$$



For $E_{\text{radiator}} = 20$ GeV,

D-jets are suppressed by 30% relative to inclusive at $\theta \sim 0.1$ rad

B-jets are suppressed by nearly by 100% relative to inclusive at $\theta \sim 0.1$ rad

SOME NOTES...

Heavy-flavor and the dead cone...

- High-energy HF – little dead-cone effect because of m/E small – radiative quark (a la LF) e-loss dominant
- I. Vitev et al. study HF-jets with z_g (standard grooming tech.) => no significant differences as compared to LF jets (high-momentum 140-160 GeV) for HF tagged jet $Q \rightarrow Qg$
- QQbar splits in parton shower:
 - I. Vitev – little / no in medium modification
 - Novel techniques in pp on Disentangling Heavy Flavor at Colliders [arXiv:1702.02947] – potentially interesting for AA

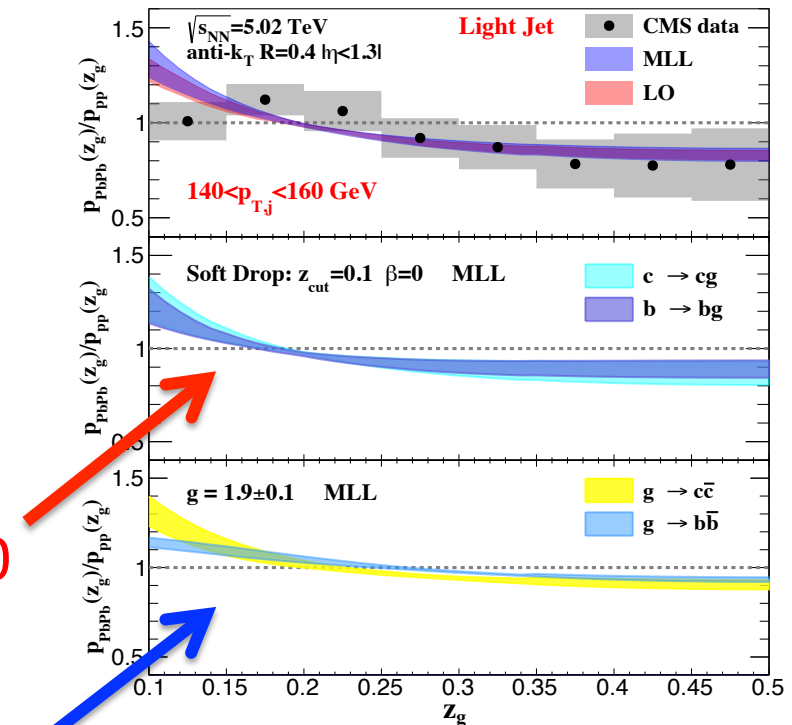
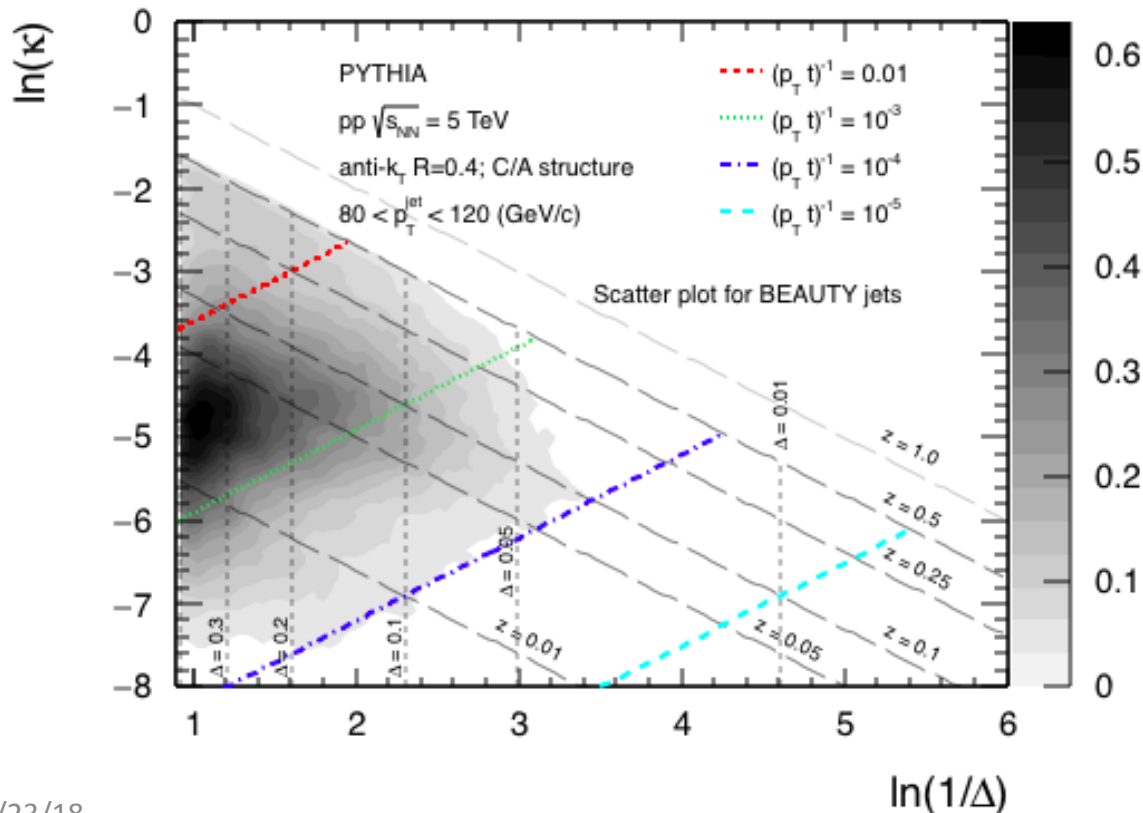


Figure 2. The modification of the jet splitting functions in 0-10% central Pb+Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV for the p_T bin $140 < p_{T,j} < 160$ GeV. The upper panels compare the LO and MLL predictions to CMS light jet substructure measurements [12]. The middle and lower panels present the MLL modifications for heavy flavor tagged jet - the $Q \rightarrow Qg$ and $\rightarrow Q\bar{Q}$, respectively.

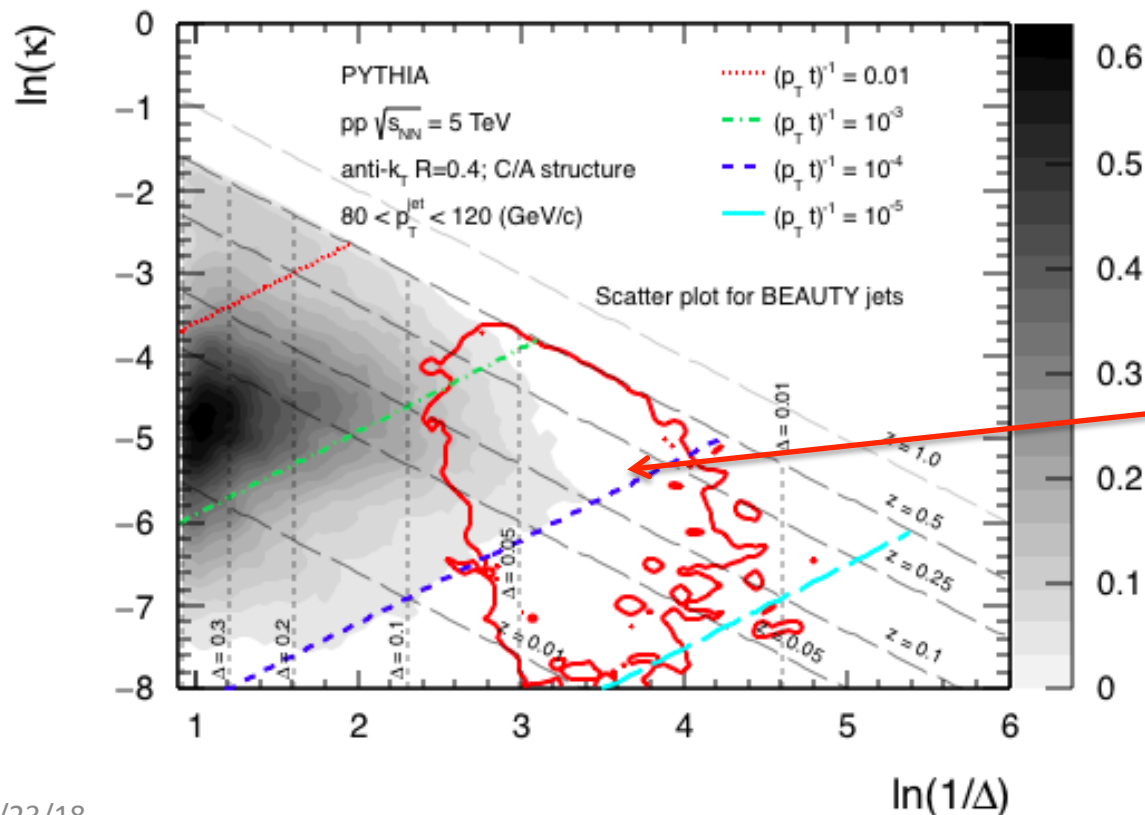
Heavy-flavor and the dead cone...

- Can we take a look with Lund diagram?
 - Use leading (high- p_T) HF-hadron (lepton) for the tag & follow declusterization
- Why z_g (used so far) not good?



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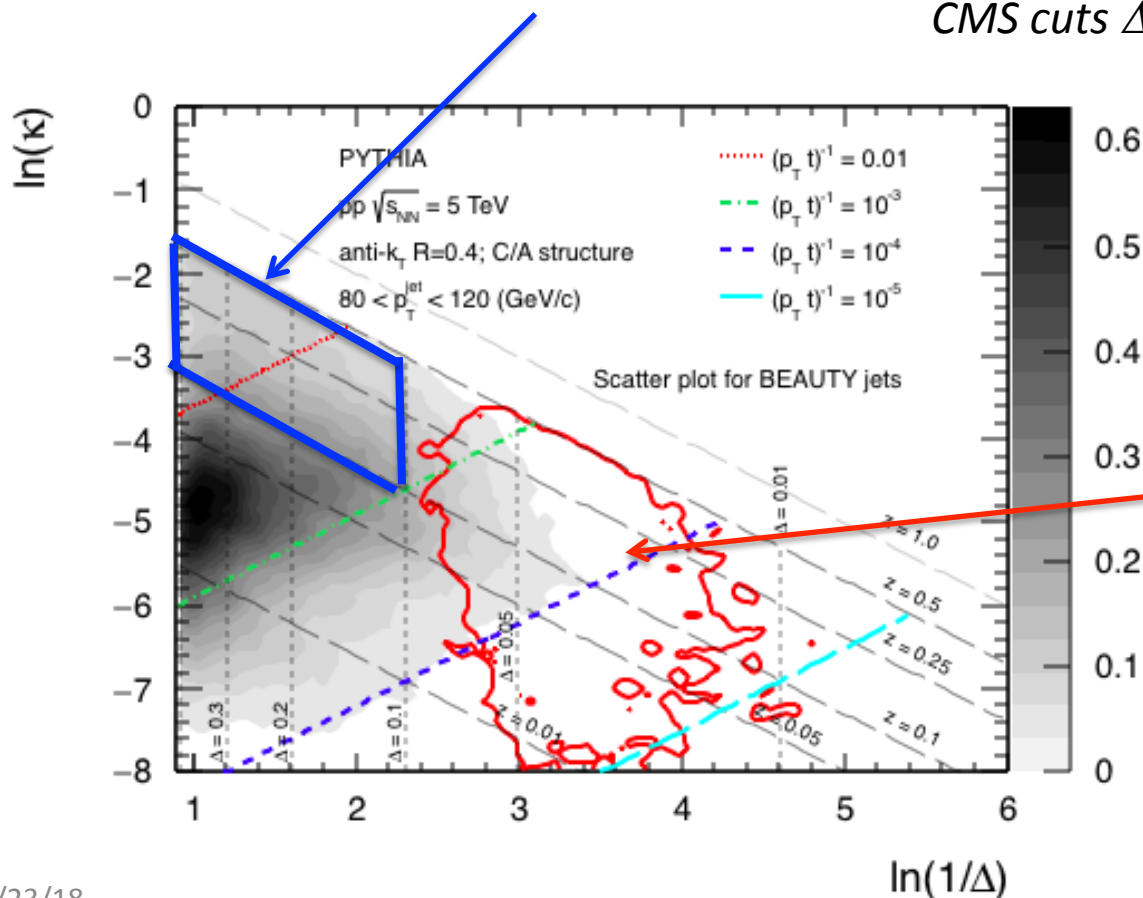
Dead cone

- An approximation
 - $\Delta < m/E$
- Entries that dead cone ought to suppress

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CMS cuts $\Delta < 0.1$ – ALICE does not



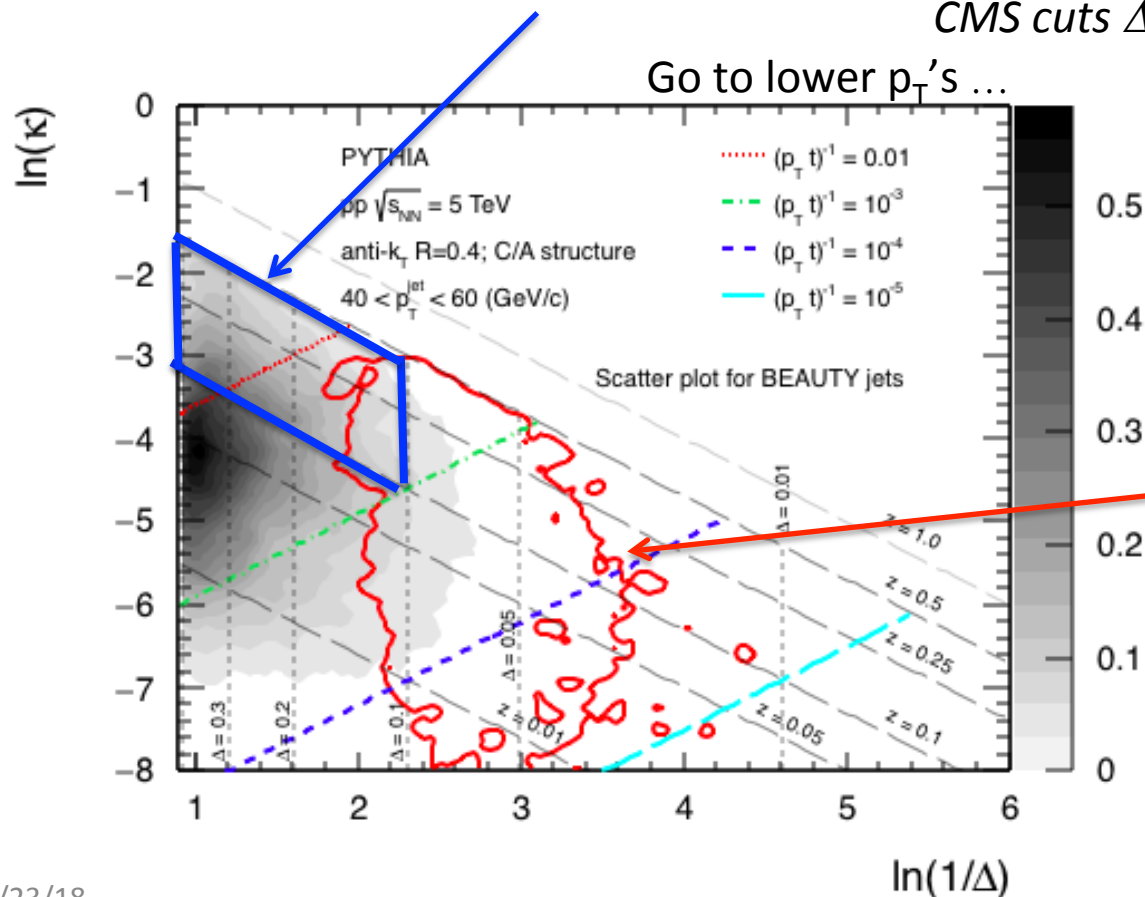
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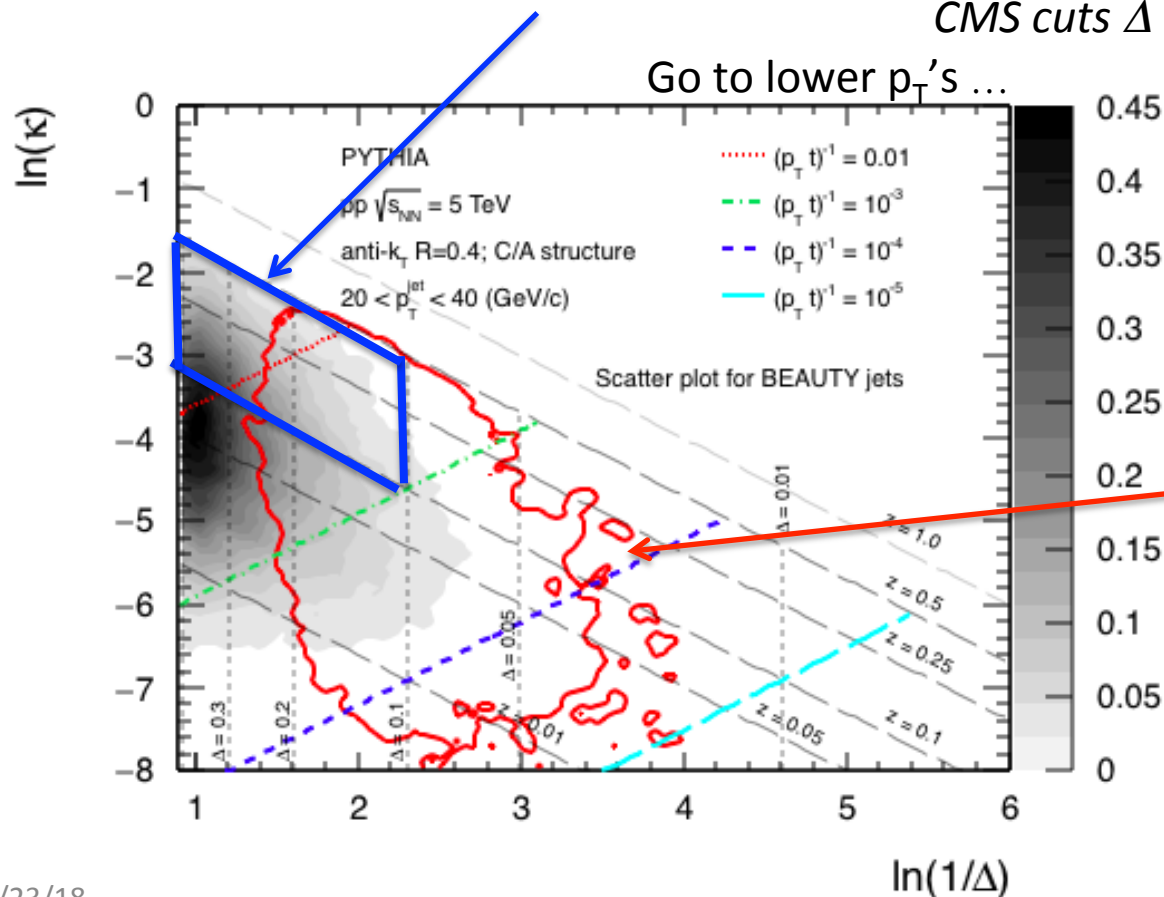
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Dead cone

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Lower pT

High- p_T

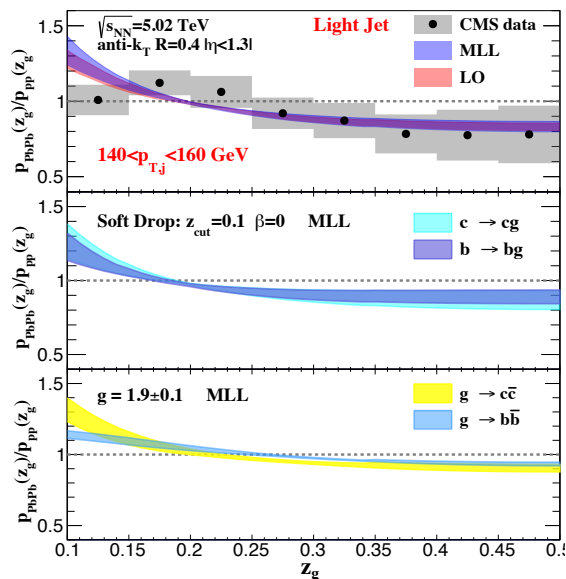
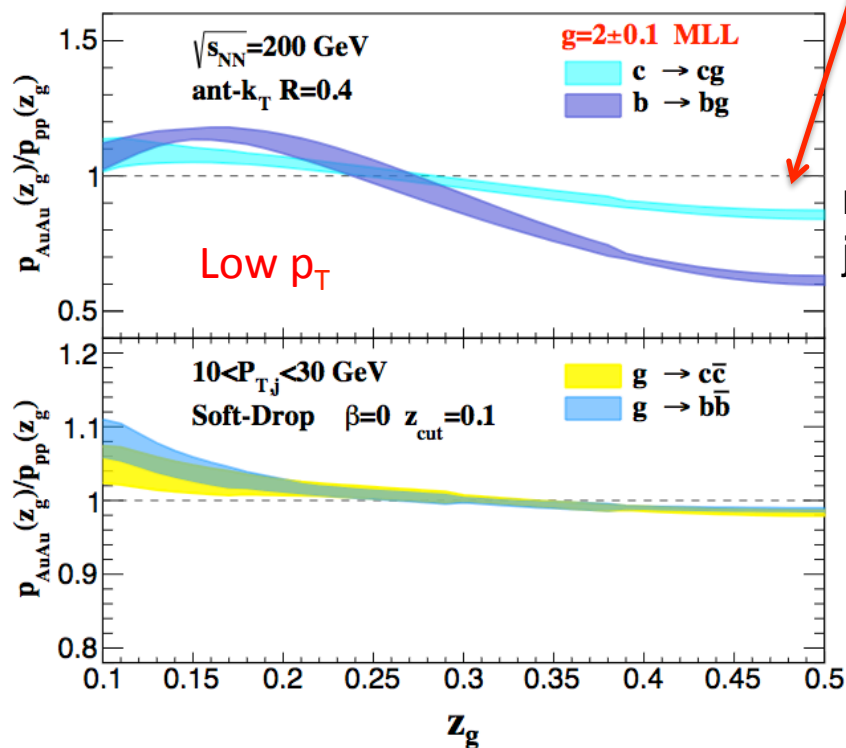


Figure 2. The modification of the splitting functions in 0-10% central Pb-collisions at $\sqrt{s_{NN}} = 5.02$ TeV for the p_T $140 < p_{T,j} < 160$ GeV. The upper panel compares the LO and MLL predictions to the light jet substructure measurements [12]. The middle and lower panels present the MLL modifications for heavy flavor tagged jet - $Q \rightarrow Qg$ and $\rightarrow Q\bar{Q}$, respectively.

Modification of fragmentation functions for gluon and quark
Kang et al 2014



Hai Tao Li, yesterday...
There it is!
Normalization?(!)

modifications of the jet splitting functions

HTL, Vitev 2018

https://indico.bnl.gov/event/5039/contributions/26256/attachments/21619/29615/b_jet_haitao.pdf

kT cut vs. Soft Drop ?

Soft Drop Condition:

$$\frac{\min(p_{t1}, p_{t2})}{p_{t1} + p_{t2}} > z_{\text{cut}} \left(\frac{\Delta R_{12}}{R_0} \right)^\beta$$

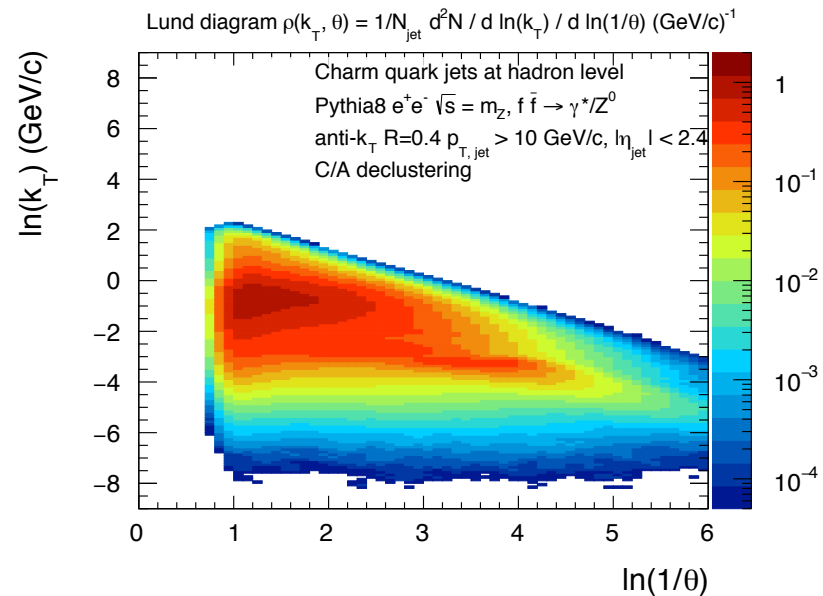
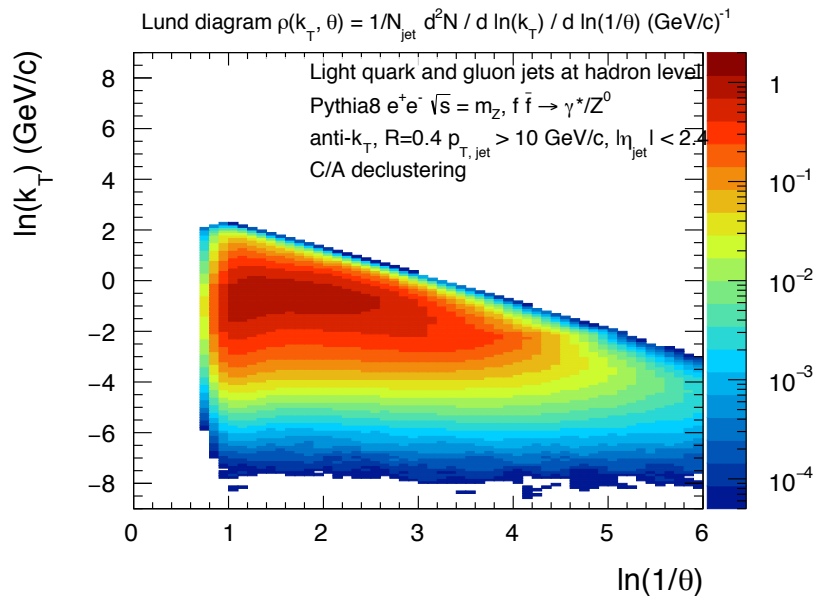
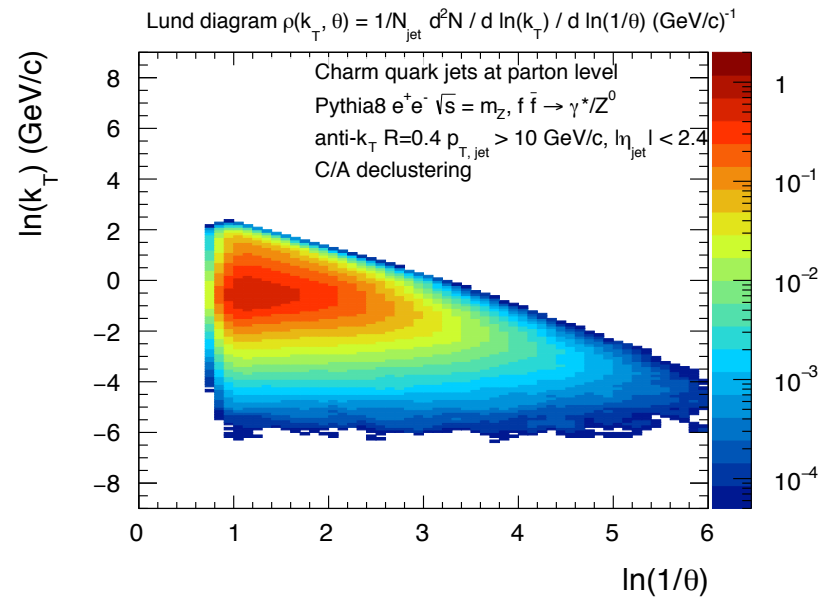
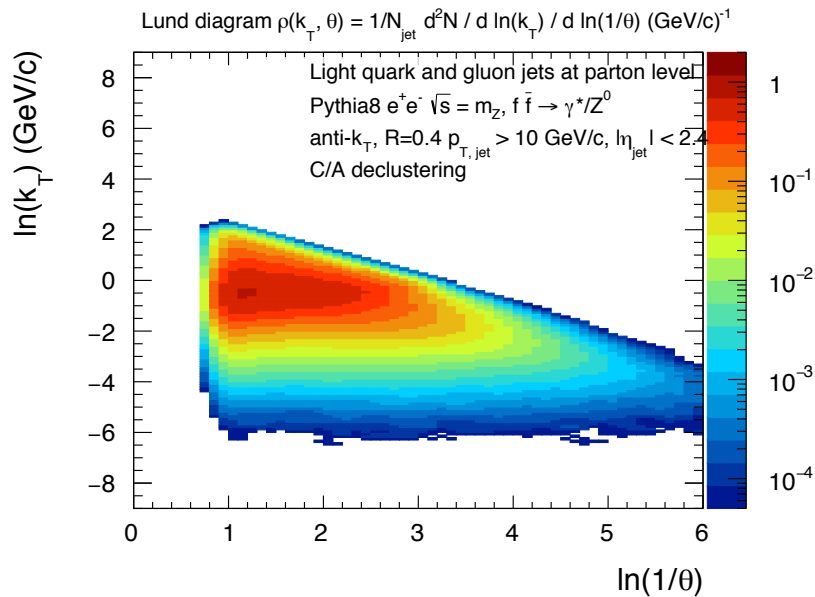
$$z = \frac{p_{t2}}{p_{t1} + p_{t2}}$$

$$k_T = p_{t2} \cdot \Delta R_{12}$$

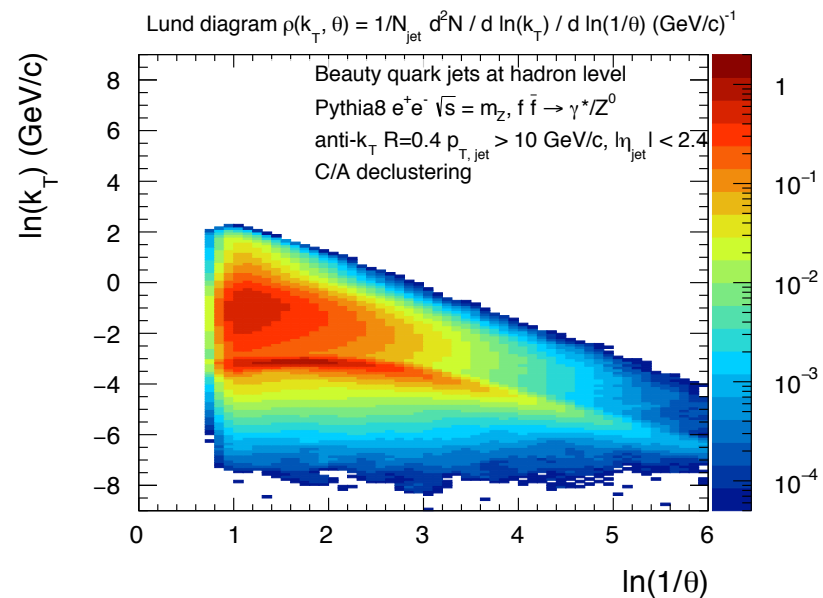
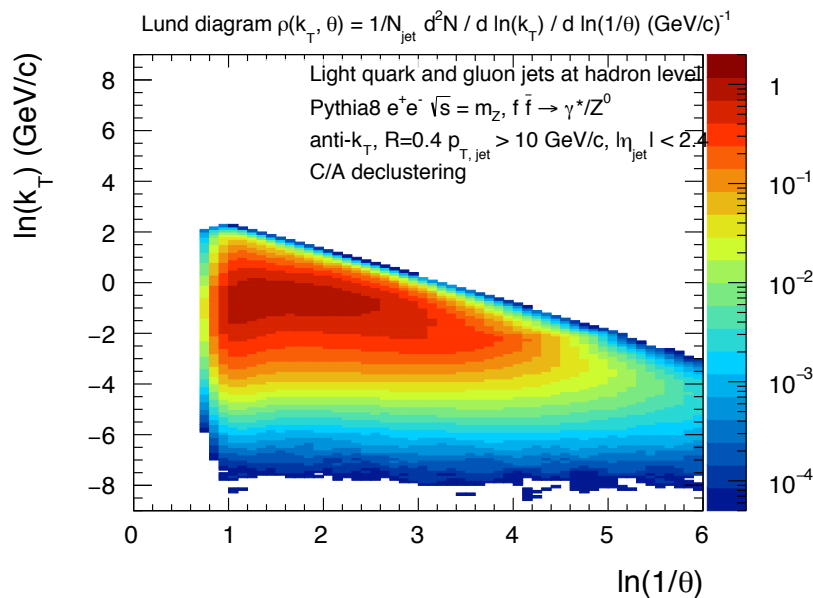
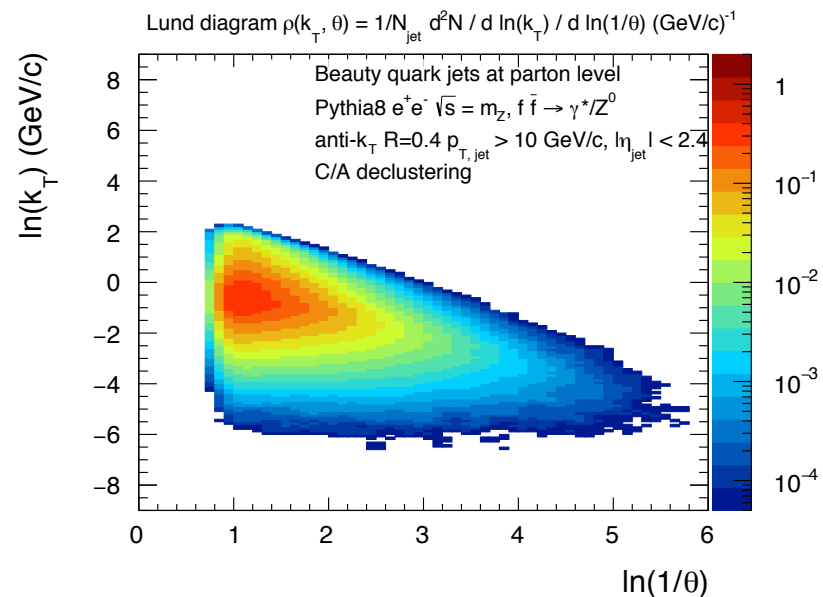
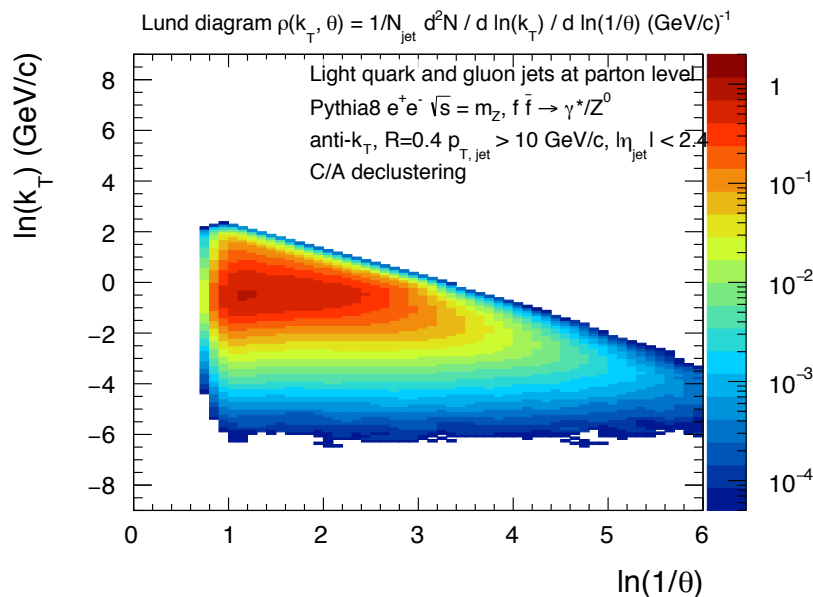
$$p_{t2} > k_{T\text{cut}} \quad \text{for } R_0 = 1, \beta = 1$$

**OUTLOOK: REVISIT LEP, MEASURE AT
LHC, ... EIC? DETAILED HADRONIZATION
STUDIES?**

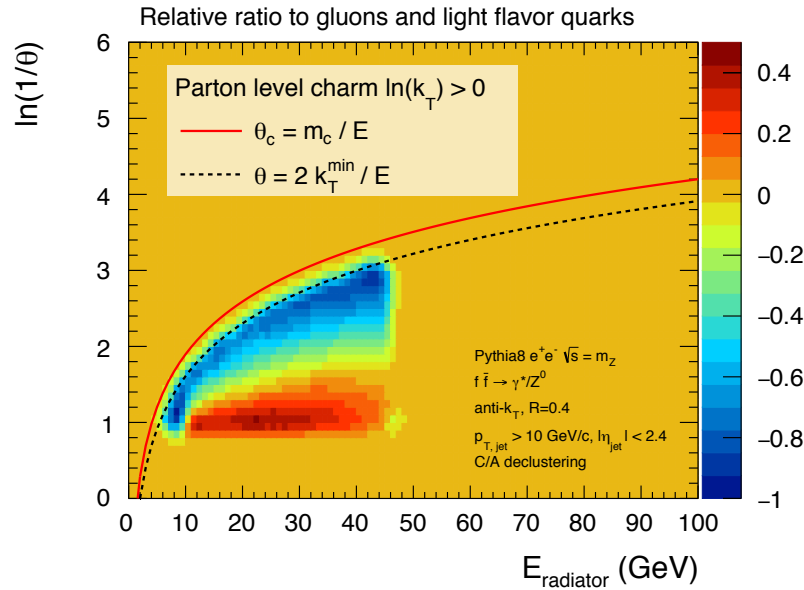
Outlook : LEP, (eIC...) - charm



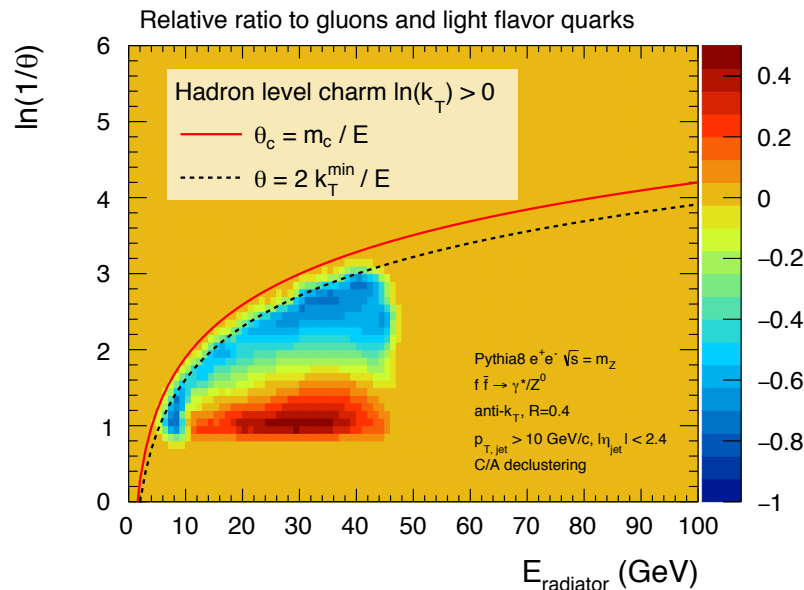
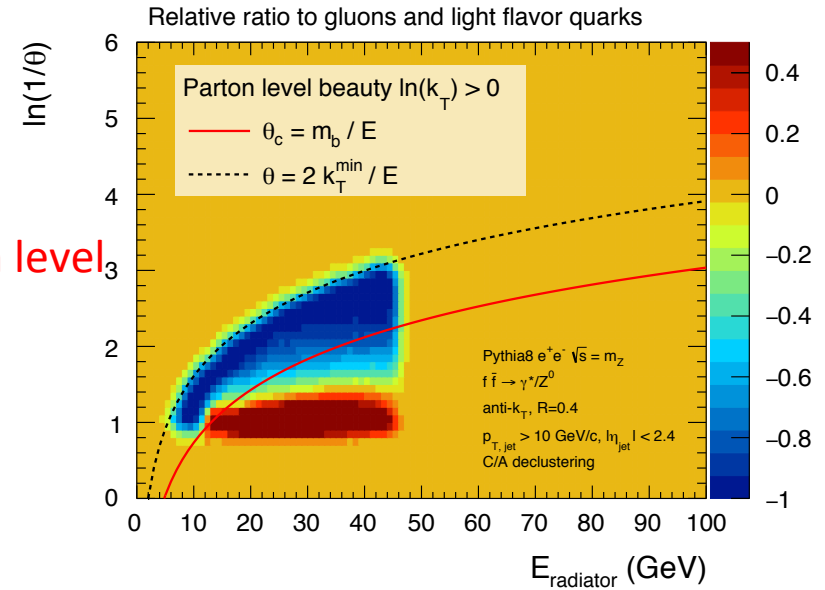
Outlook : LEP, (eIC...) - beauty



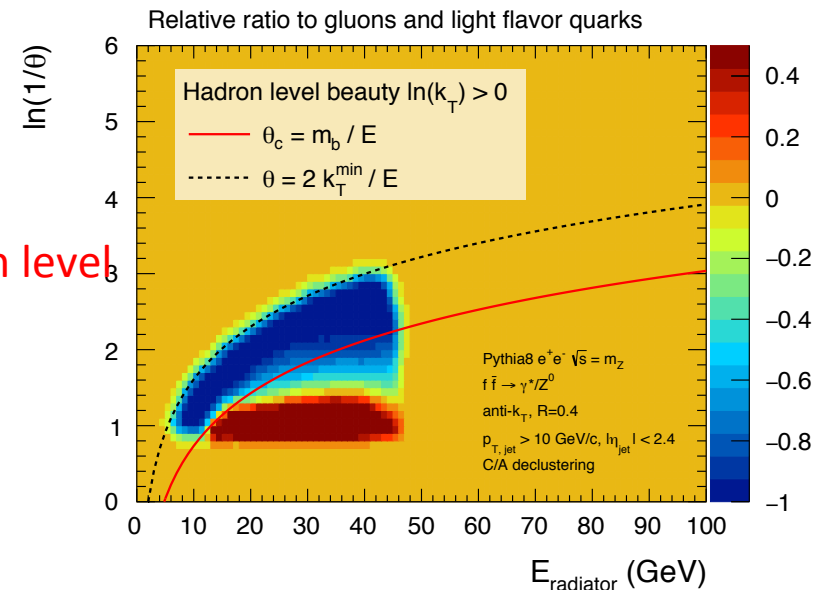
Outlook : LEP, (eIC...)



Parton level



Hadron level



Summary

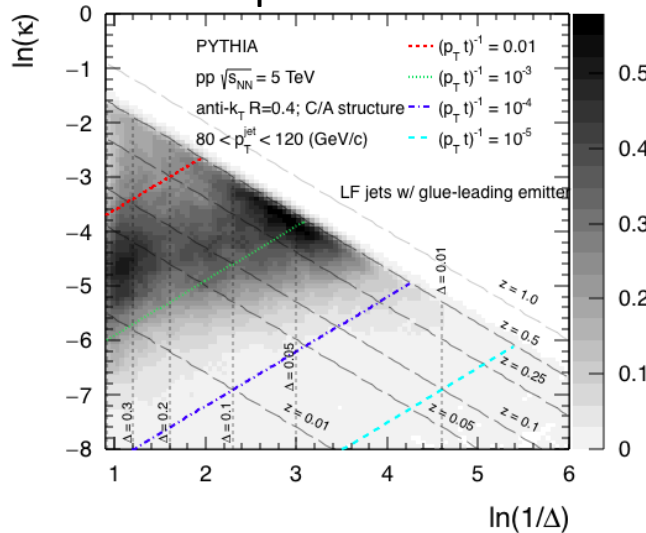
- Dead cone at the subjet level:
 - rather “simple” conceptually
 - Cut on k_T : opportunity to study the effect systematically => constraints on hadronization (LEP, eIC; hadronic collisions more difficult ISR, MPI,...)
- More applications for declustering in considerations (jet classification based on Lund plane - k_T)
- Last but not least: no grooming, no soft drop, no ΔR cuts... => opportunity for heavy-ion collisions => jet modifications, medium response(?), ...

Notes on heavy-quarks

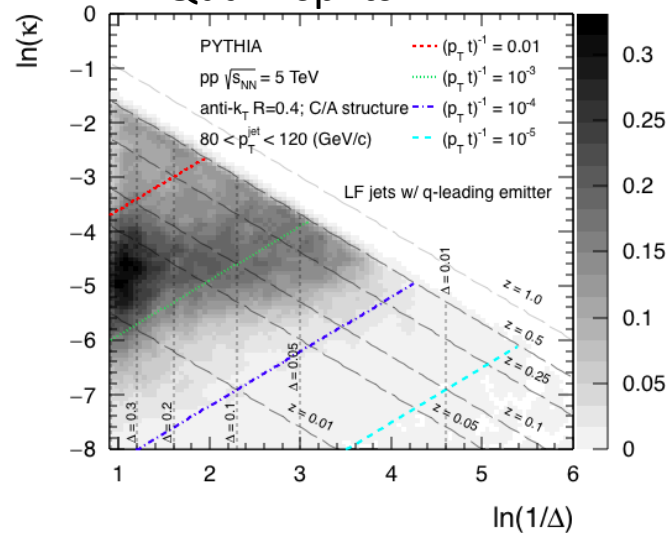
$$\bar{\rho}(\Delta, \kappa) = \frac{1}{N_{\text{jet}}} \frac{dn_{\text{emission}}}{d \ln \kappa d \ln 1/\Delta}^{45}$$

$$p_{T,a} > p_{T,b}, \kappa = \frac{p_{T,b}}{p_{T,a} + p_{T,b}} \Delta_{ab}$$

Gluon splits



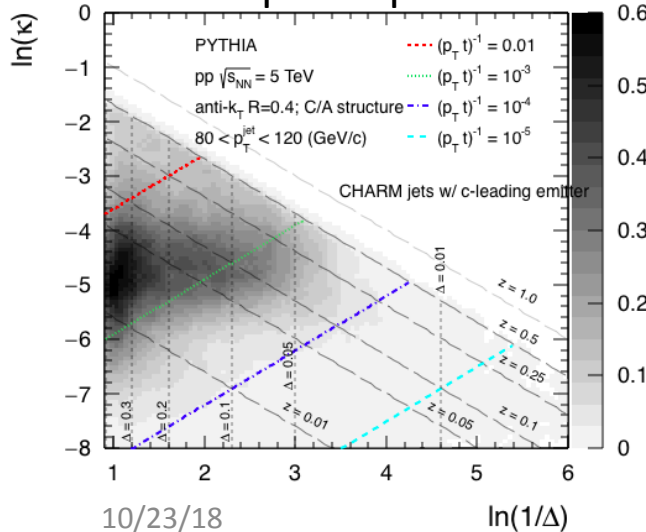
LF Quark splits



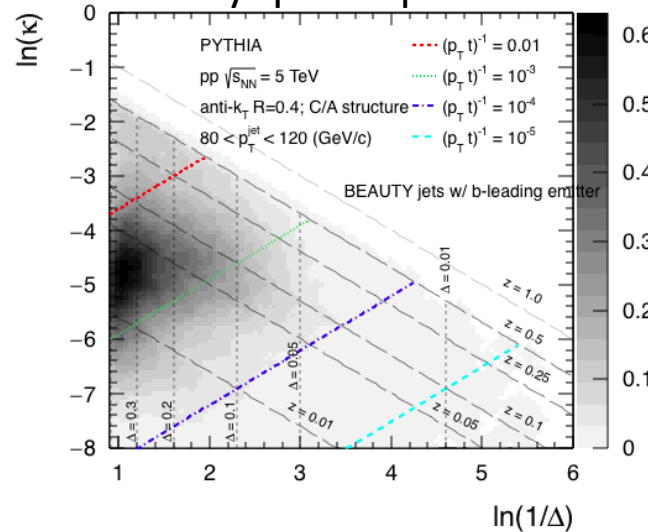
Purely MC exercise – compare LF to HF jets

An approximation:
- use the leading parton within the emitter (follow Q)

Charm quark splits



Beauty quark splits



Opportunity? – explore with Machine Learning

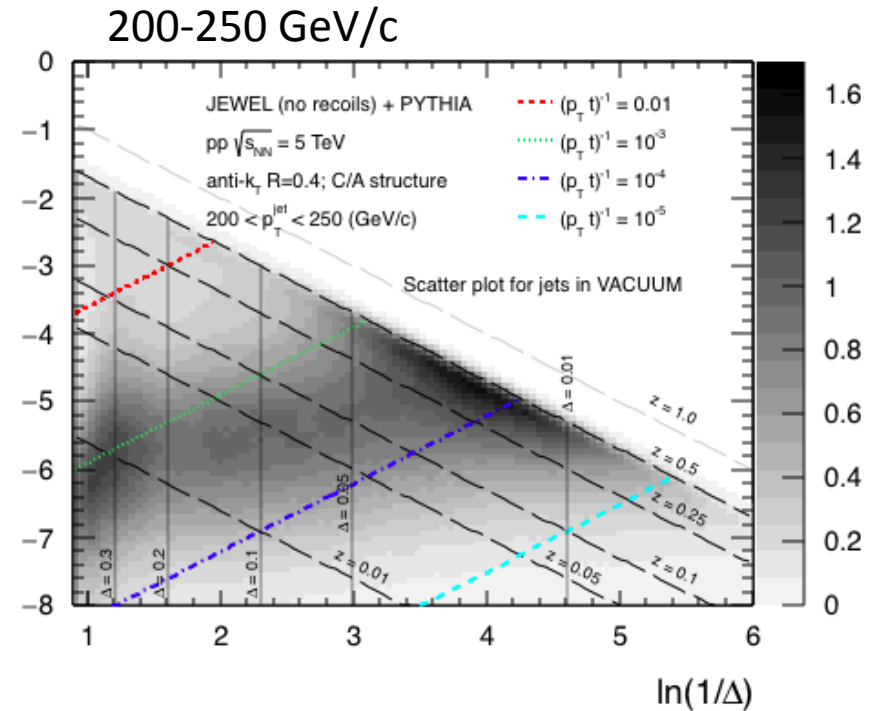
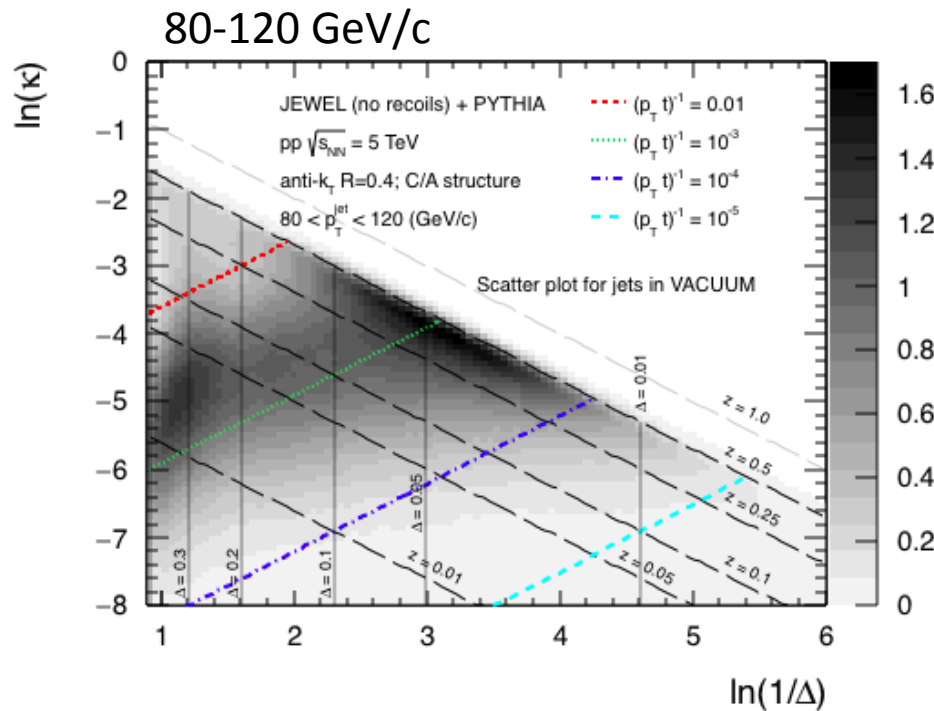
JET QUENCHING – MC STUDY

Jet Lund diagram

JEWEL + PYTHIA (RECOIL=OFF)
10% most central PbPb at 5 TeV
VACUUM

$$p_{T,a} > p_{T,b}, \kappa = \frac{p_{T,b}}{p_{T,a} + p_{T,b}} \Delta_{ab}^{47}$$

$$\bar{\rho}(\Delta, \kappa) = \frac{1}{N_{\text{jet}}} \frac{dn_{\text{emission}}}{d \ln \kappa d \ln 1/\Delta}$$

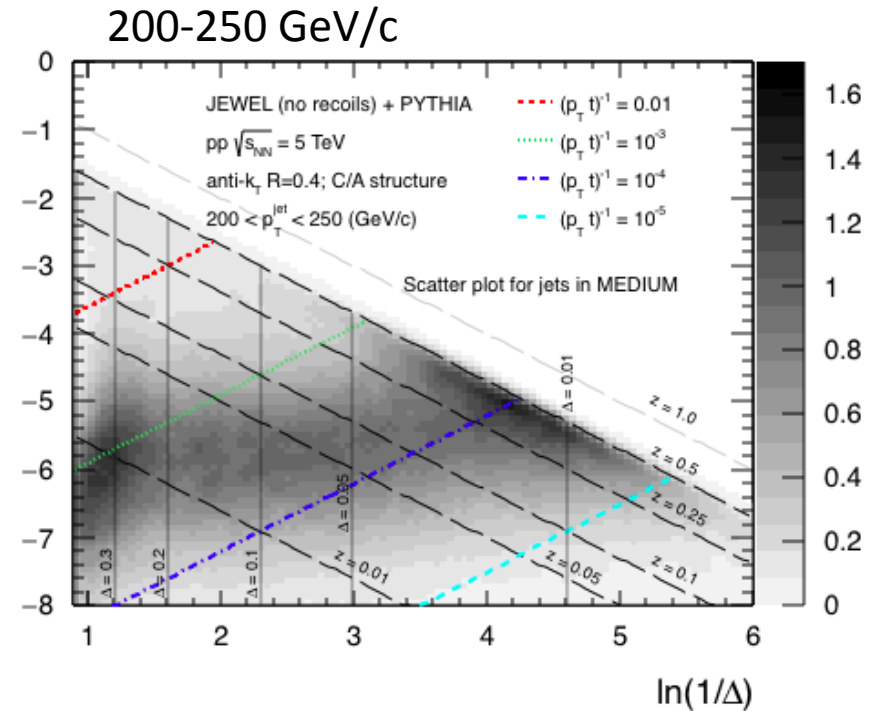
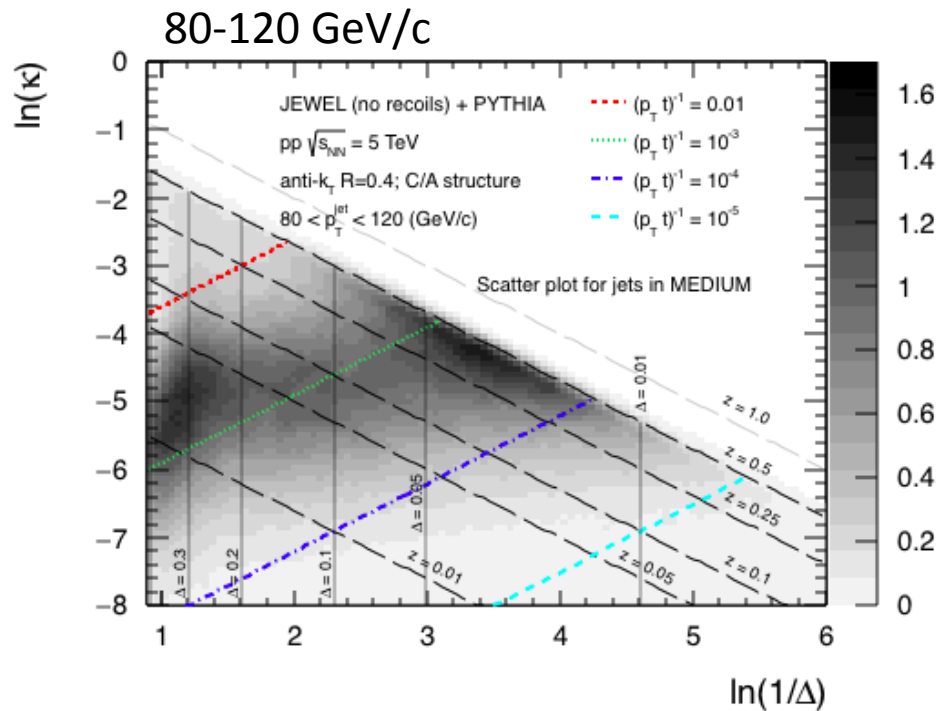


Jet Lund diagram

JEWEL + PYTHIA (RECOIL=OFF)
10% most central PbPb at 5 TeV
MEDIUM

$$p_{T,a} > p_{T,b}, \kappa = \frac{p_{T,b}}{p_{T,a} + p_{T,b}} \Delta_{ab}^{48}$$

$$\bar{\rho}(\Delta, \kappa) = \frac{1}{N_{\text{jet}}} \frac{dn_{\text{emission}}}{d \ln \kappa d \ln 1/\Delta}$$

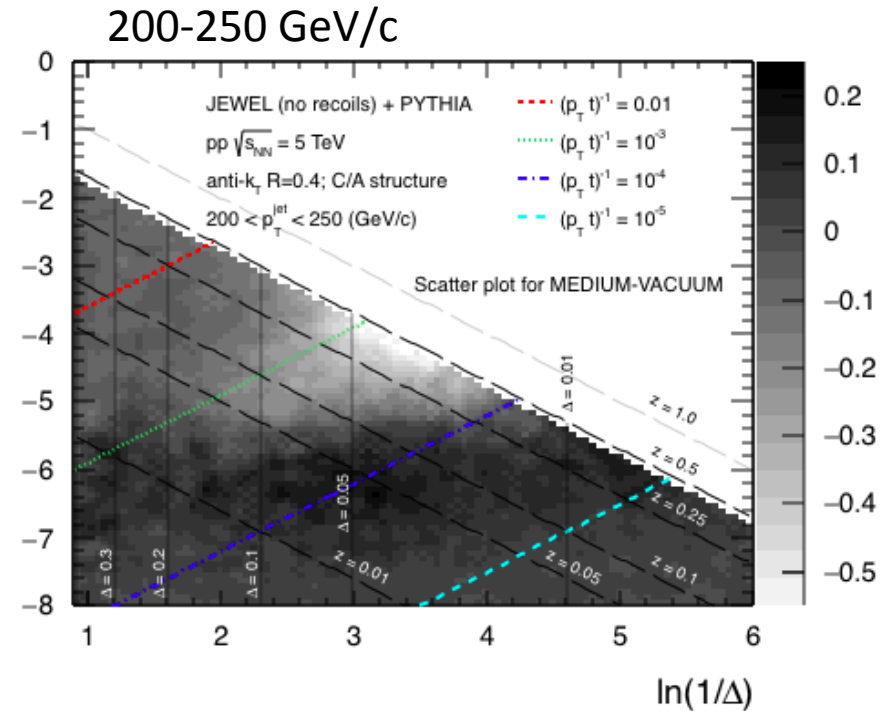
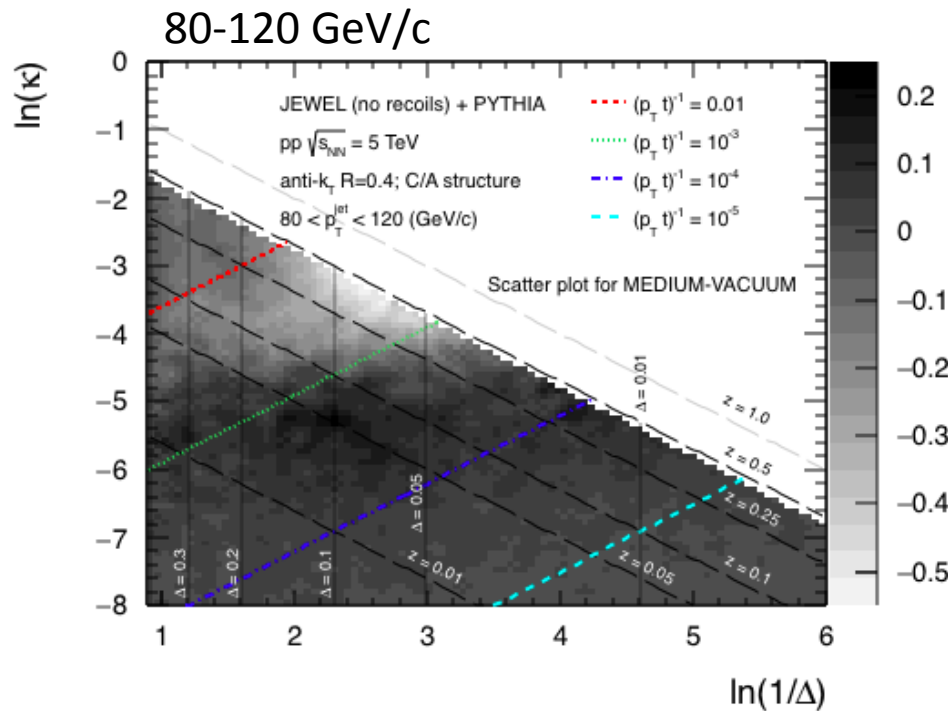


Jet Lund diagram

JEWEL + PYTHIA (RECOIL=OFF)
10% most central PbPb at 5 TeV
MEDIUM - VACUUM

$$p_{T,a} > p_{T,b}, \quad \kappa = \frac{p_{T,b}}{p_{T,a} + p_{T,b}} \Delta_{ab}^{49}$$

$$\bar{\rho}(\Delta, \kappa) = \frac{1}{N_{\text{jet}}} \frac{dn_{\text{emission}}}{d \ln \kappa d \ln 1/\Delta}$$

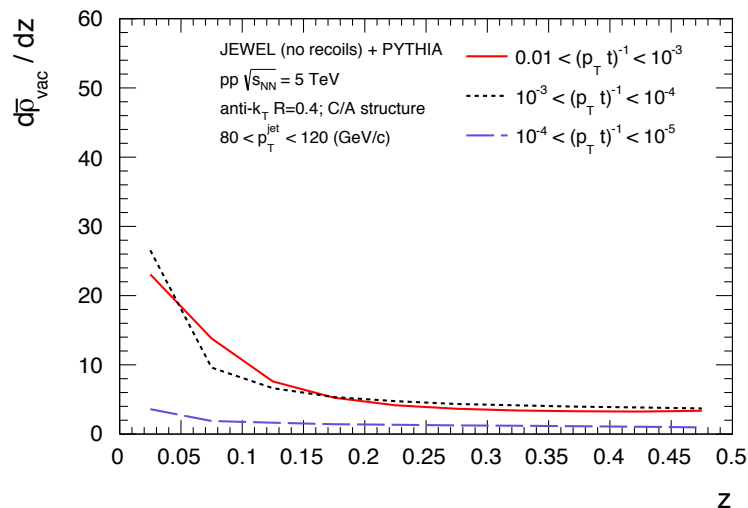


Jet Lund diagram

$$p_{T,a} > p_{T,b}, \kappa = \frac{p_{T,b}}{p_{T,a} + p_{T,b}} \Delta_{ab}^{50}$$

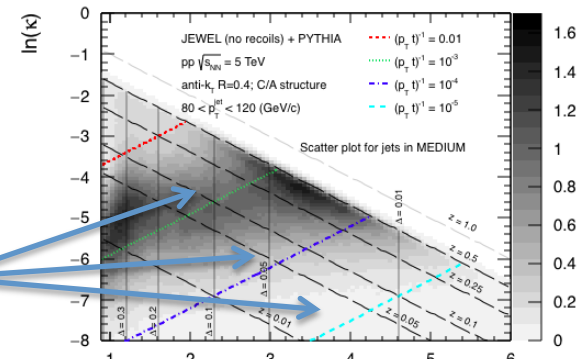
slicing through time

$$\bar{\rho}(\Delta, \kappa) = \frac{1}{N_{\text{jet}}} \frac{dn_{\text{emission}}}{d \ln \kappa d \ln 1/\Delta}$$

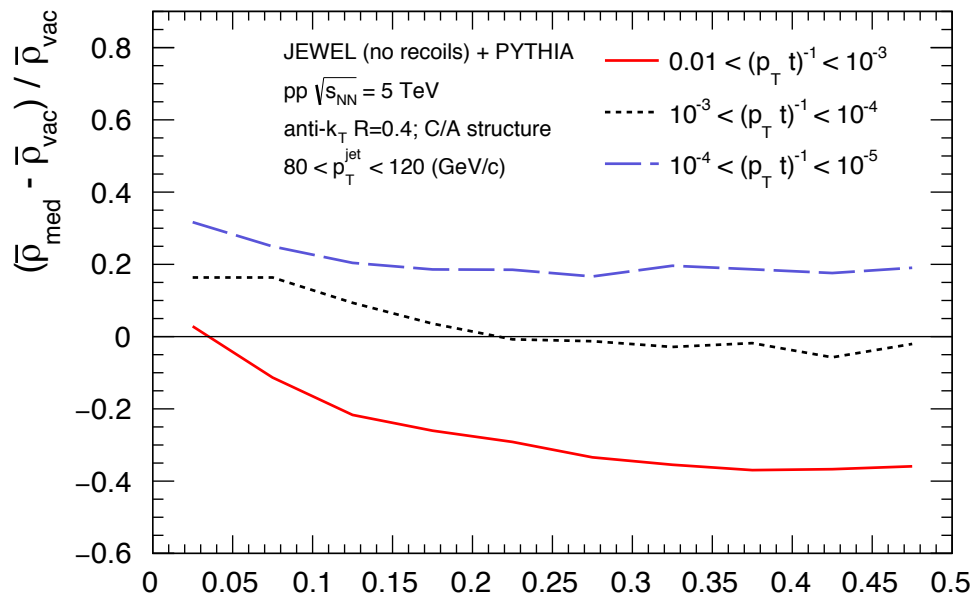
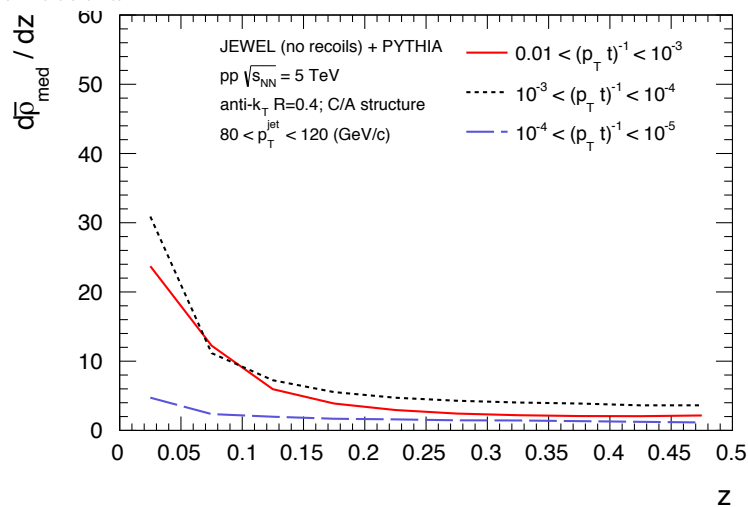


80 – 120 GeV/c

Slicing along $1/(p_T \times t)$
and projecting along z



Mon 22/10/2018 20:24:09 PDT



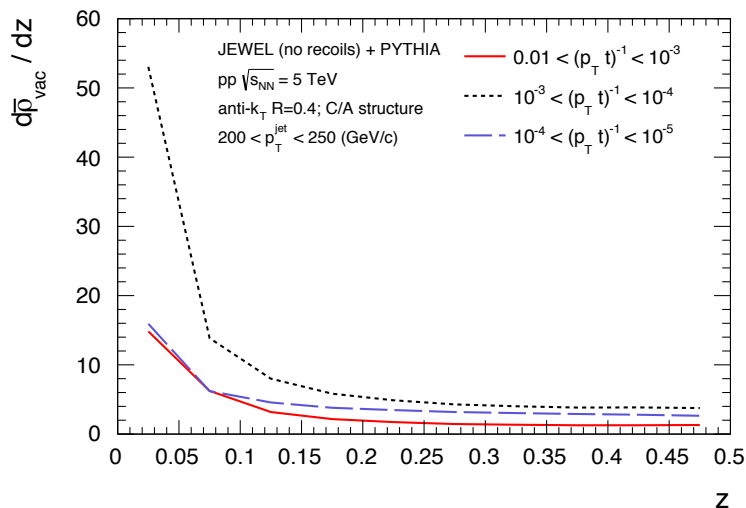
Mon 22/10/2018 20:24:09 PDT

Jet Lund diagram

$$p_{T,a} > p_{T,b}, \kappa = \frac{p_{T,b}}{p_{T,a} + p_{T,b}} \Delta_{ab}^{51}$$

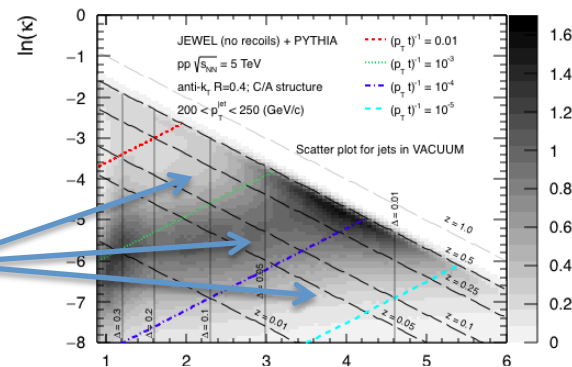
slicing through time

$$\bar{\rho}(\Delta, \kappa) = \frac{1}{N_{\text{jet}}} \frac{dn_{\text{emission}}}{d \ln \kappa d \ln 1/\Delta}$$

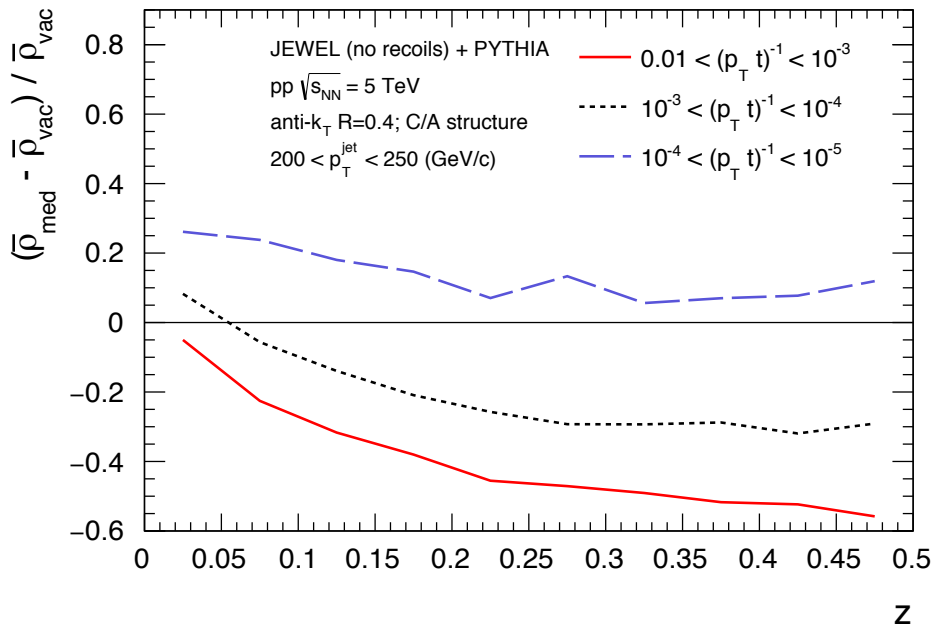
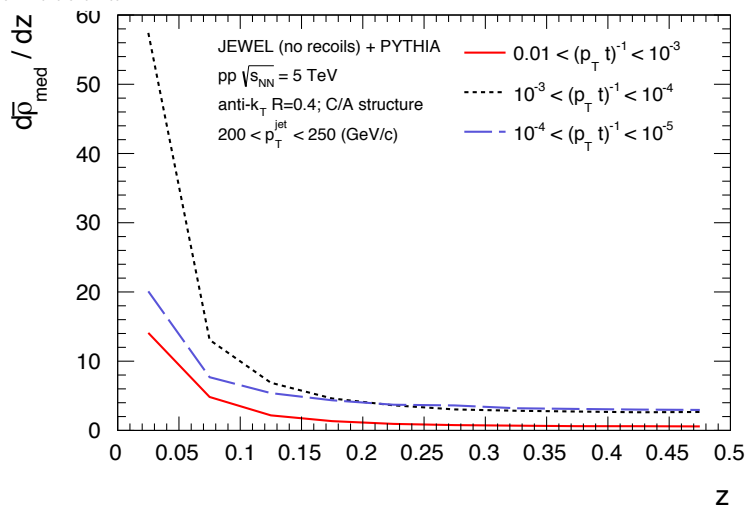


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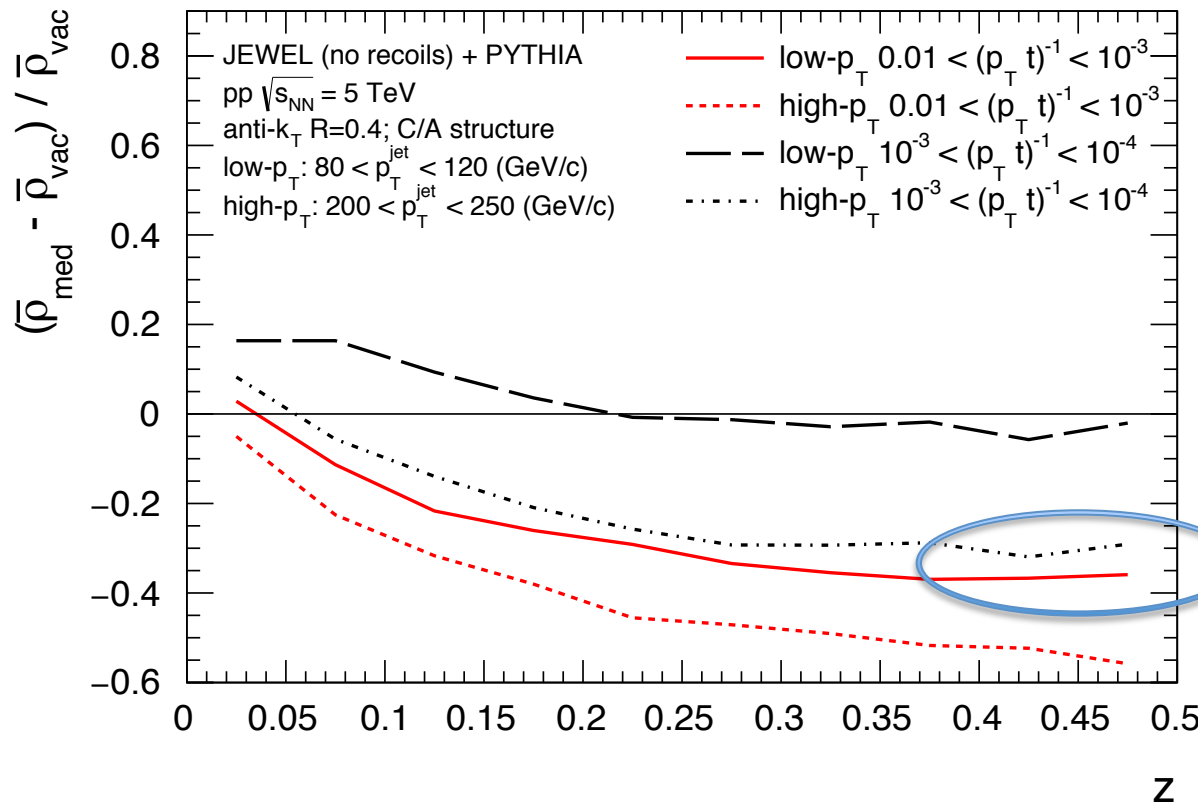
Z

Jet Lund diagram

$$p_{T,a} > p_{T,b}, \kappa = \frac{p_{T,b}}{p_{T,a} + p_{T,b}} \Delta_{ab}^{52}$$

slicing through time

$$\bar{\rho}(\Delta, \kappa) = \frac{1}{N_{\text{jet}}} \frac{dn_{\text{emission}}}{d \ln \kappa d \ln 1/\Delta}$$



Low- with high- p_T for
the similar t

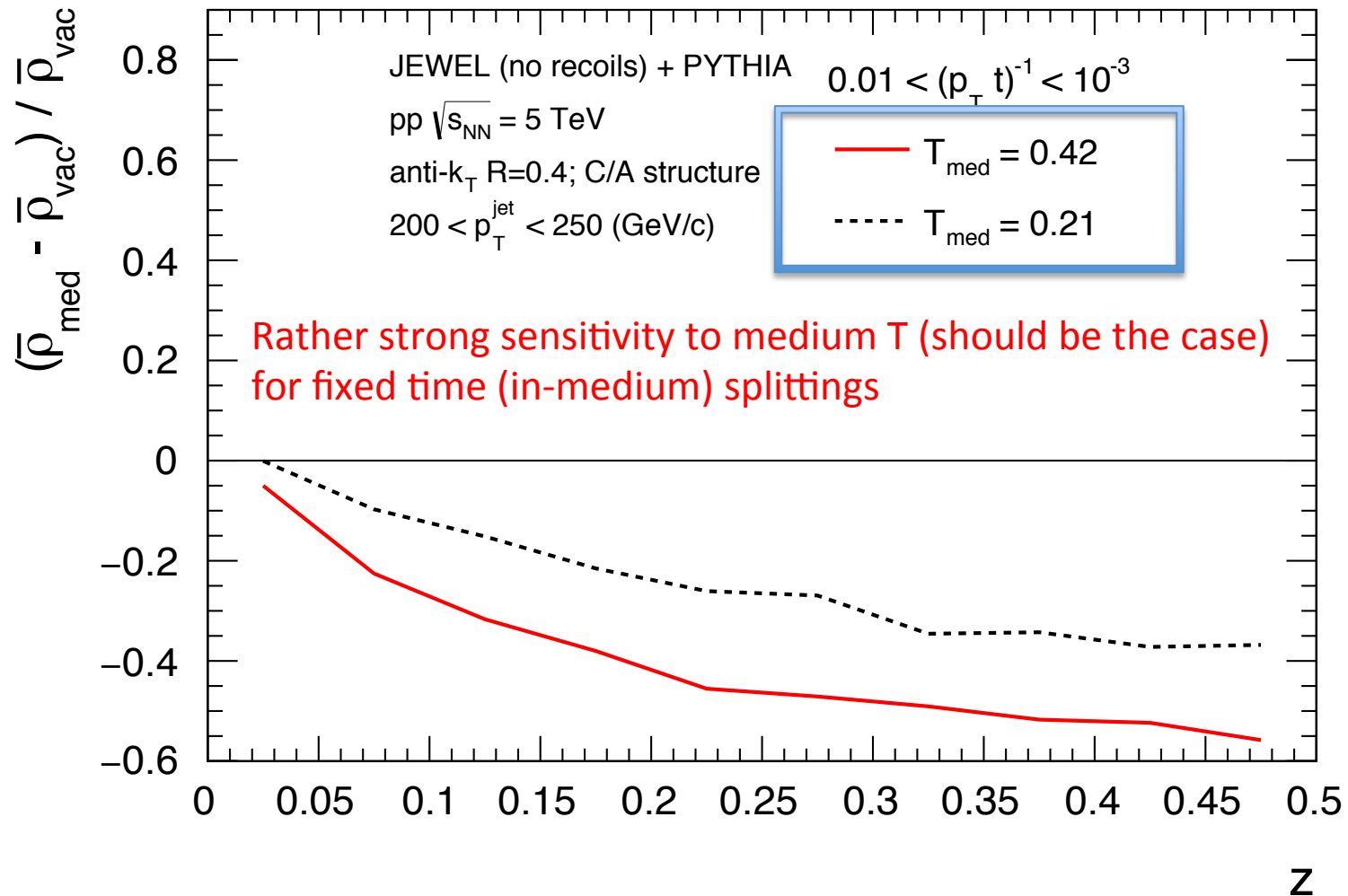


similar suppression
(scaling factor for $p_T \times t$
roughly 80/200)

Work on slicing on
formation time directly
ongoing

Sensitivity to medium's T

An example...



Sensitivity to medium's L (centrality)

An example...

