

UCLA 2019 Santa Fe Jets and Heavy Flavor Workshop, Jan. 30, 2019, UCLA

Gluon TMDs from Quarkonium Production in proton collisions

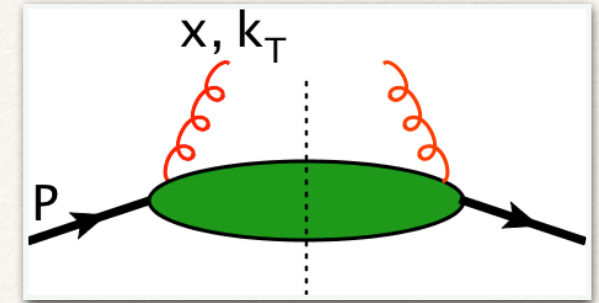
Marc Schlegel
Department of Physics
New Mexico State University

in collaboration with J.-P. Lansberg, C. Pisano, F. Scarpa
based on **PLB 784, 217 (2018)** ; **NPB 920, 192 (2017)**

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of gluons in nucleon**

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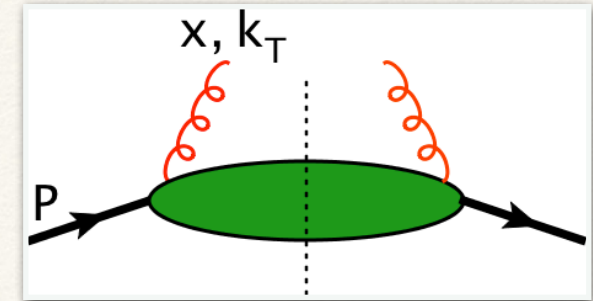
TMD gluonic matrix element



$$\Gamma^{\alpha\beta; [\mathcal{W}, \mathcal{W}']} (x, \mathbf{k}_T) = \int \frac{d\lambda d^2 z_T}{(2\pi)^3} e^{ix\lambda + i\mathbf{k}_T \cdot \mathbf{z}_T} \langle P, S | F^{n\alpha}(0) \mathcal{W} F^{n\beta}(\lambda n + z_T) \mathcal{W}' | P, S \rangle$$

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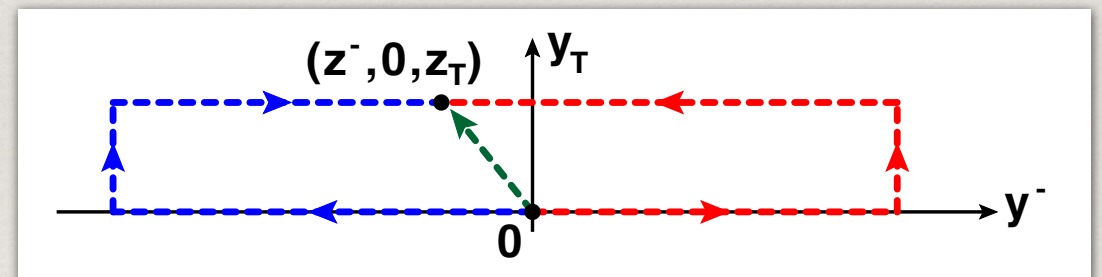
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Wilson line in fundamental representation

Color Gauge invariant definition of TMDs $\rightarrow$ Wilson line

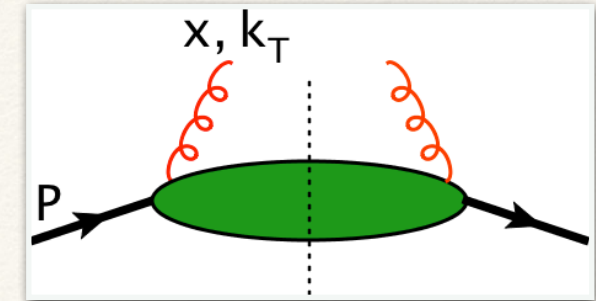
$$\mathcal{W}[a, b] = \mathcal{P} e^{-ig \int_a^b ds \cdot A^a(s) t^a}$$

$\rightarrow$ Wilson line for TMD:
nontrivial, process dependent:



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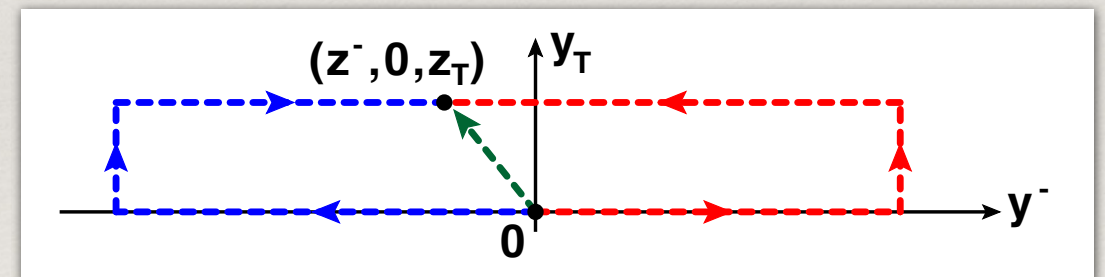
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→ Wilson line for TMD:
nontrivial, process dependent:



past-pointing WL →
Initial State Interactions:
pp → color singlet + X (DY-type)

$$\Gamma^{\alpha\beta}; [-, -]$$

future-pointing WLs →
Final State Interactions:
ep → jets + X (SIDIS-type)

$$\Gamma^{\alpha\beta}; [+, +]$$

Proper Definition & Soft function

[Collins; Sun, Xiao, Yuan; Aybat, Rogers; see also: Echevarria, Kasemats, Mulders, Pisano, JHEP 07, 158 (2015)]

Inclusion of Soft Function $\implies$

$$S(z_T) = \frac{\text{Tr}_c}{N_c} \langle 0 | \mathcal{W}_n^\dagger(-z_T/2) \mathcal{W}_{\bar{n}}(-z_T/2) \mathcal{W}_{\bar{n}}^\dagger(z_T/2) \mathcal{W}_n(z_T/2) | 0 \rangle$$

$$\tilde{\Gamma}^{\alpha\beta}(x, \mathbf{b}_T) \rightarrow \frac{\tilde{\Gamma}^{\alpha\beta}(x, \mathbf{b}_T)}{\sqrt{S(\mathbf{b}_T)}}$$

Modification needed in order to have:

- 1) Renormalizable Matrix Element
- 2) Finite matching coefficients at small b_T

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Parameterization

unpolarized & linearly polarized gluons :
helicity flip TMDs $\rightarrow$ azimuthal modulations

$$\Gamma^{\alpha\beta}(x, k_T) = \frac{1}{2x} \left[-g_T^{\alpha\beta} f_1^g(x, k_T^2) + \frac{k_T^\alpha k_T^\beta - \frac{1}{2} k_T^2 g_T^{\alpha\beta}}{M^2} h_1^{\perp g}(x, k_T^2) \right]$$

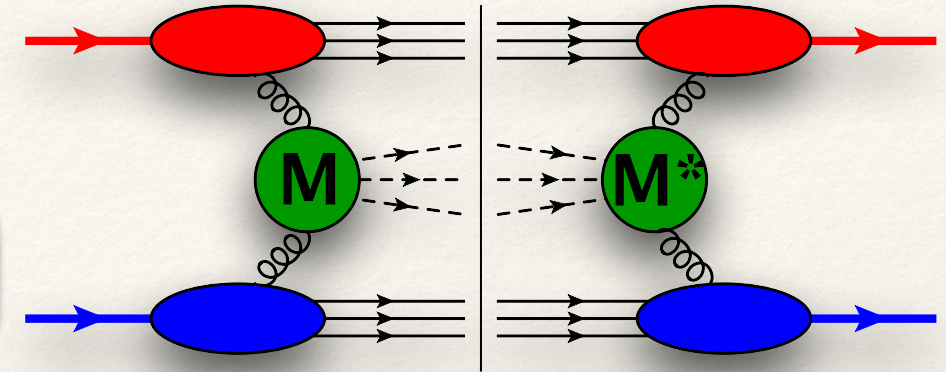
Gluon TMDs from pp - collisions

[J.-P. Lansberg, C. Pisano, M.S., NPB 920, 192 (2017)]

General TMD expression for gluon fusion:

$$\frac{d\sigma}{d^4q...}(q_T \ll Q) \propto \mathcal{C}[\Gamma^{\alpha\alpha'}(x_a, k_{aT}) \Gamma^{\beta\beta'}(x_b, k_{bT})] (\mathcal{M}_{\alpha\beta} (\mathcal{M}_{\alpha'\beta'})^*)$$

$$\mathcal{C}[w \textcolor{blue}{f} \textcolor{red}{g}] = \int d^2k_{aT} \int d^2k_{bT} \delta^{(2)}(k_{aT} + k_{bT} - q_T) w(k_{aT}, k_{bT}) \textcolor{blue}{f}(x_a, k_{aT}) \textcolor{red}{g}(x_b, k_{bT})$$

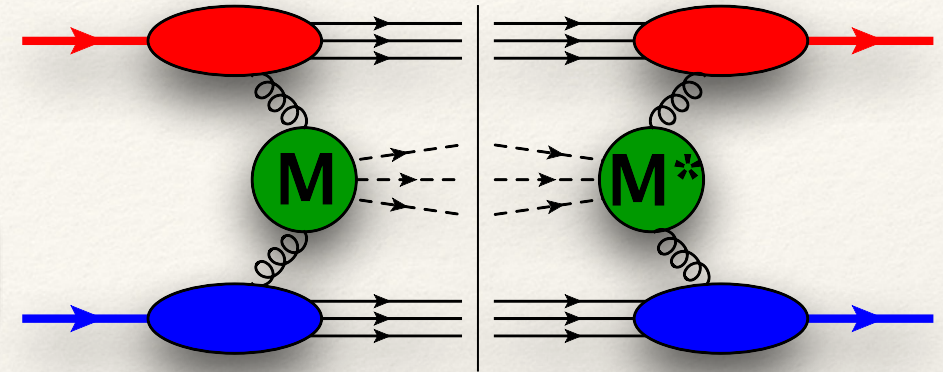


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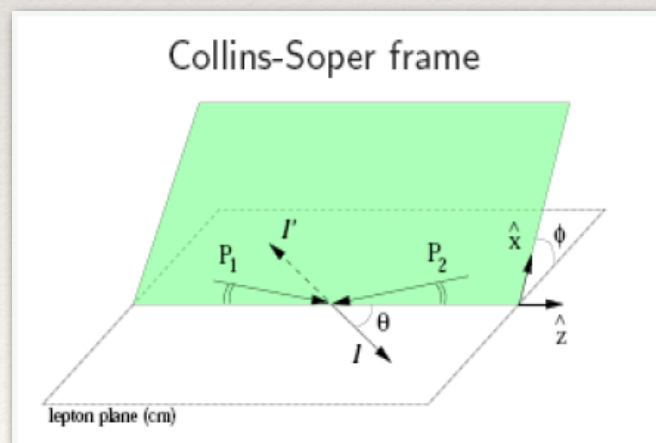
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2-particle (pair) final states in pp - collisions

$$\left. \frac{d\sigma^{gg}}{d^4q d\Omega} \right|_{q_T \ll Q} = \hat{F}_1 [f_1^g \otimes f_1^g] + \hat{F}_2 [h_1^{\perp g} \otimes h_1^{\perp g}] + \cos(2\phi) \hat{F}_3 [h_1^{\perp g} \otimes f_1^g + f_1^g \otimes h_1^{\perp g}] + \cos(4\phi) \hat{F}_4 [h_1^{\perp g} \otimes h_1^{\perp g}]$$



Evaluate cross section in
c.m. frame of the produced pair
Collins - Soper angles θ , ϕ

Gluon TMDs do not appear in Drell-Yan

Double $J/\psi(\Upsilon)$ - production

[Lansberg, Pisano, Scarpa, PLB 784, 217 (2018)]

TMD - formalism: J/ψ pair back-to-back; Color singlet

Here: Sample diagram (of about $O(40)$)

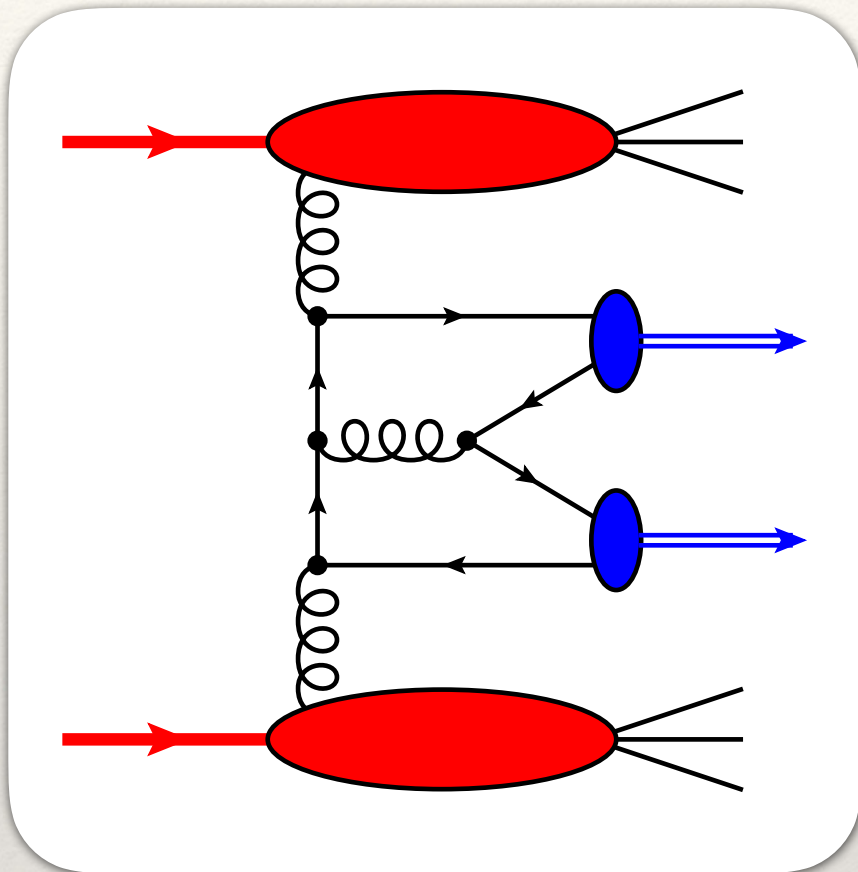
Full analytical LO amplitude

$$g + g \rightarrow J/\psi + J/\psi$$

(J/ψ in color singlet configuration)

Presented in Appendix of
[Qiao, Sun, Sun; 0903.0954] (Kudos!)

Helicity formalism: contract with polarization vectors
→ perturbative factors F_1, F_2, F_3, F_4



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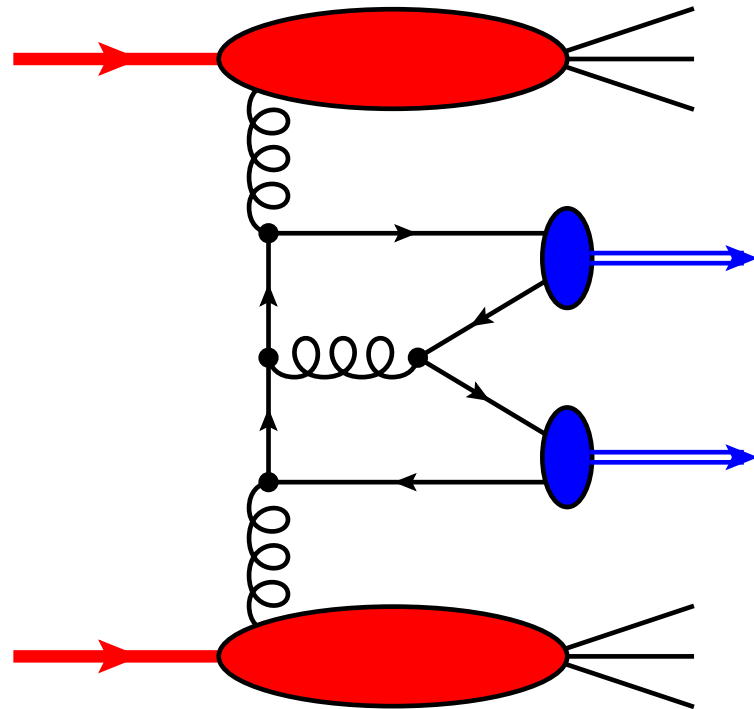
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$$F_i(Q, \cos^2 \theta) \propto \frac{\sum_{k=0}^{N_i} c_k(\alpha) (\cos^2 \theta)^k}{(1 - (1 - \alpha^2) \cos^2 \theta)^4}$$

Degree of polynomial $N_i = 6, 4, 5, 6$

Q : invariant mass of Quarkonium pair

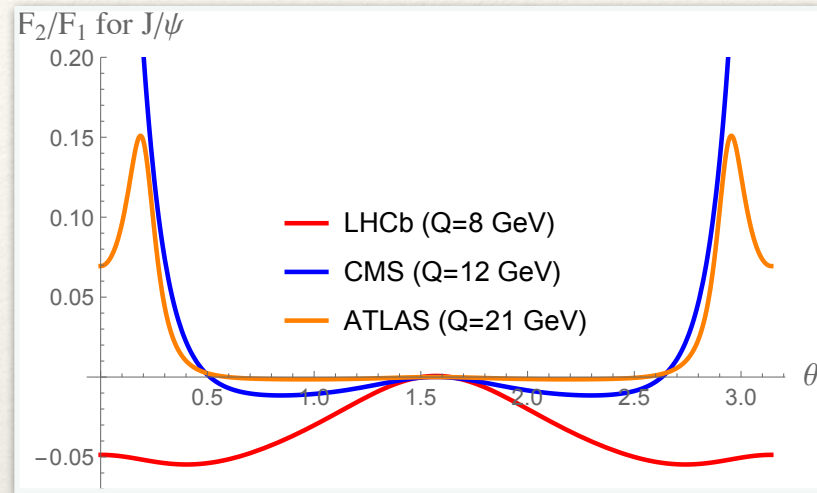
$$\alpha = \frac{2M_{J/\psi}}{Q}$$

$\alpha = 1$ threshold limit

$\alpha = 0$ high-energy limit

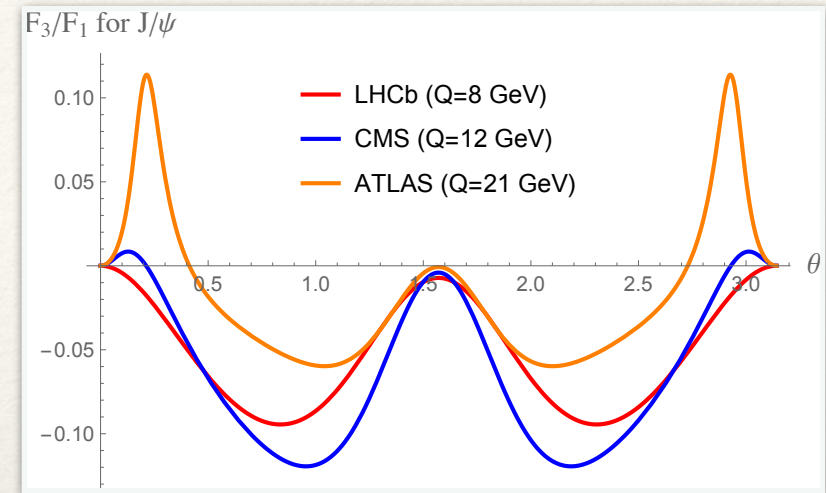
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Ratios:



q_T - spectrum:

Modification from lin. pol. glue mostly negligible

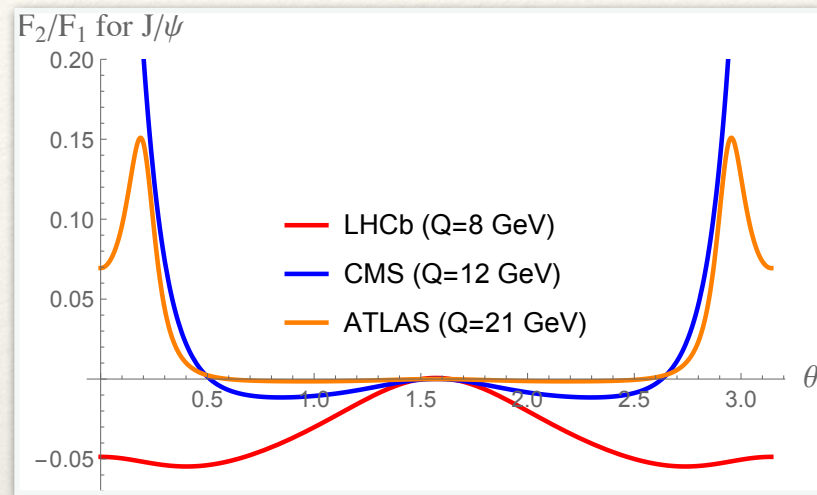


$\cos(2\phi)$:

about -10% at $\theta = \pi/4, 3\pi/4$

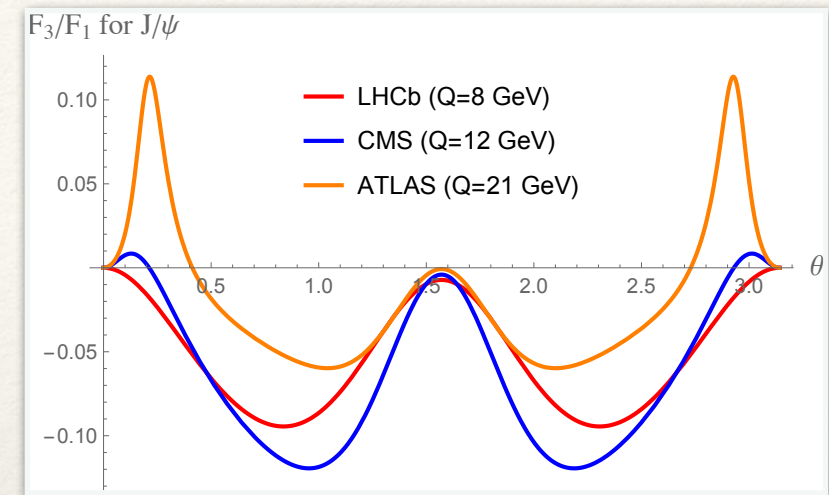
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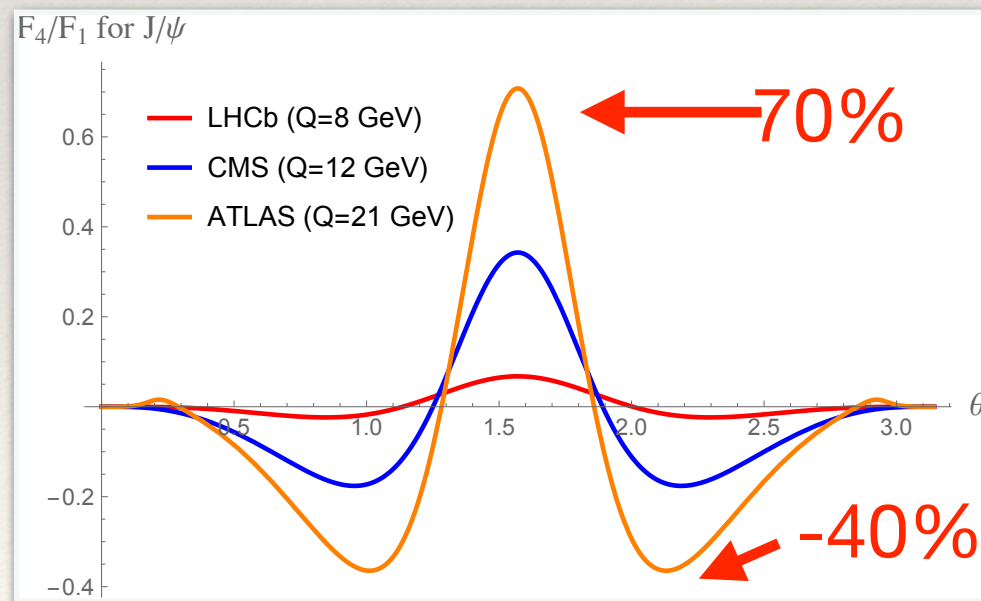
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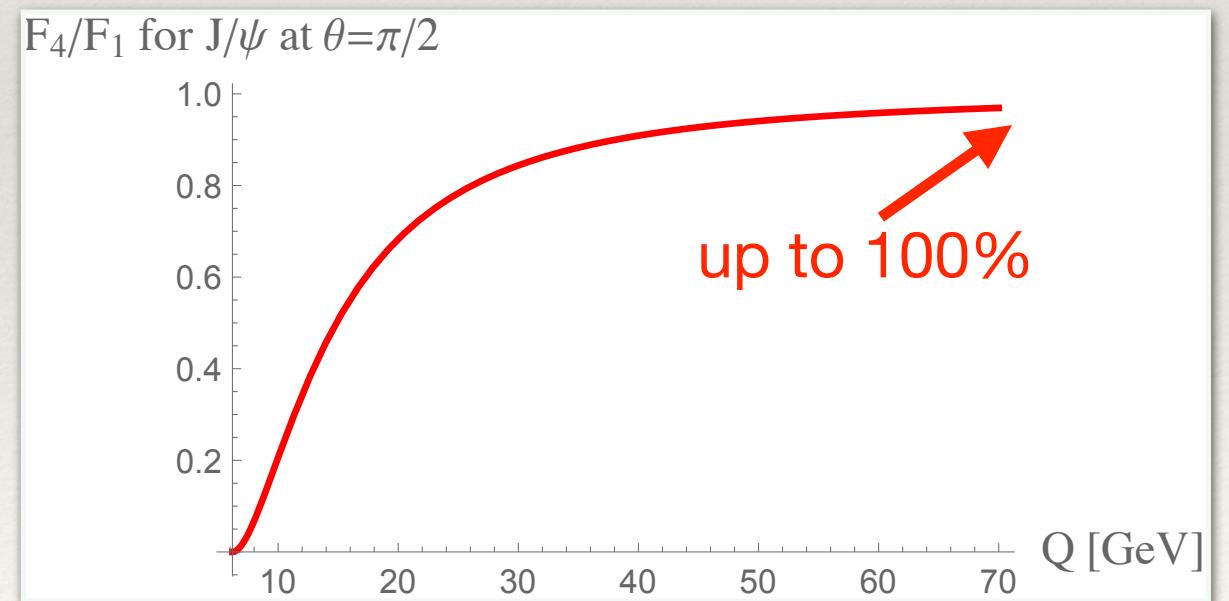
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$\cos(4\phi)$: Large at $\theta = \pi/2$!



$$F_4(\theta = \frac{\pi}{2}, Q \gg 2M_{J/\psi}) \rightarrow F_1$$



unique to this process!

q_T - spectrum: Data on J/ψ - pairs available from LHC

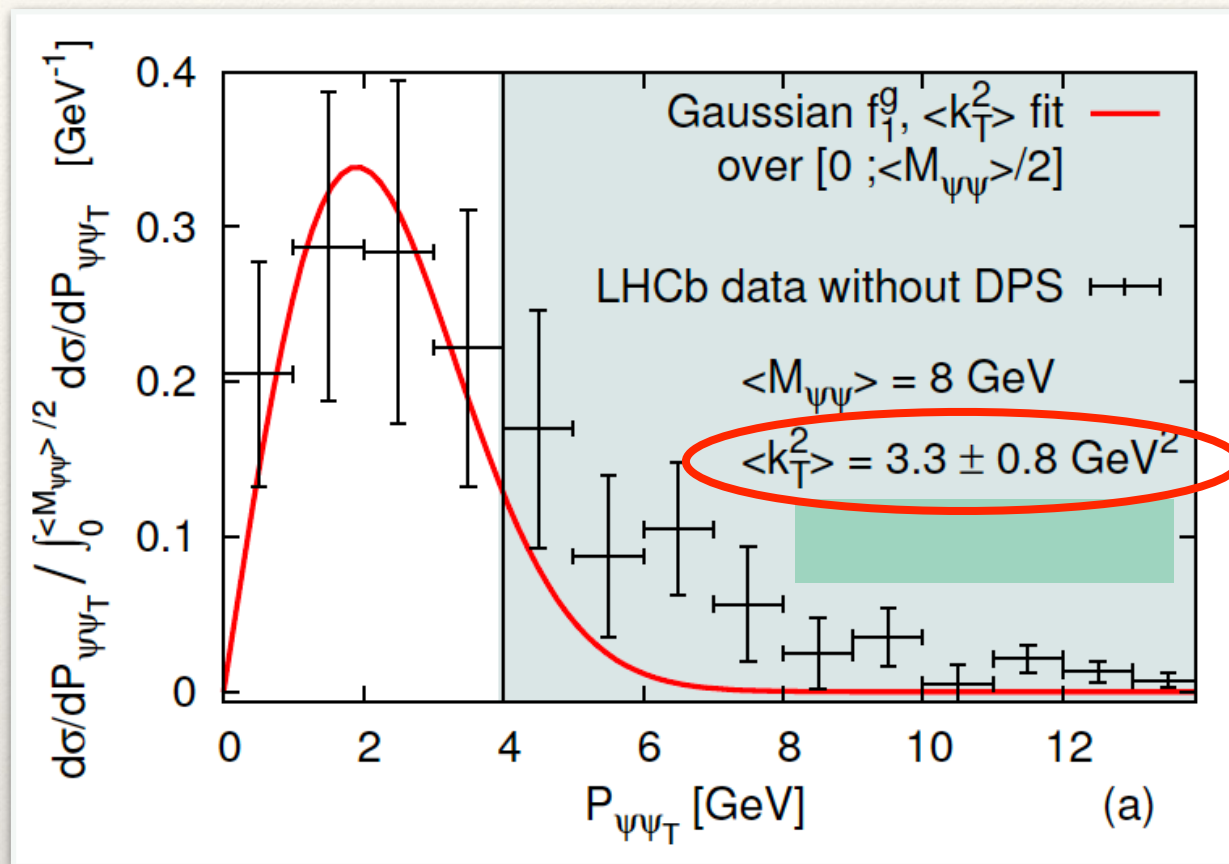
Measurements from CMS and ATLAS at 7,8 TeV, but not suitable for TMD exploration

LHCb data (2017) at 13 TeV slightly above threshold, $Q = 8$ GeV

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Fit observable

$$\frac{\left[\frac{d\sigma}{dq_T} \right]_{\text{bins}}}{\int_0^{Q/2} dq_T \left[\frac{d\sigma}{dq_T} \right]_{\text{bins}}} \propto \frac{\mathcal{C}[f_1^g f_1^g]}{\int_0^{Q/2} dq_T \mathcal{C}[f_1^g f_1^g]}$$

First glimpse on unpol. gluon TMD:

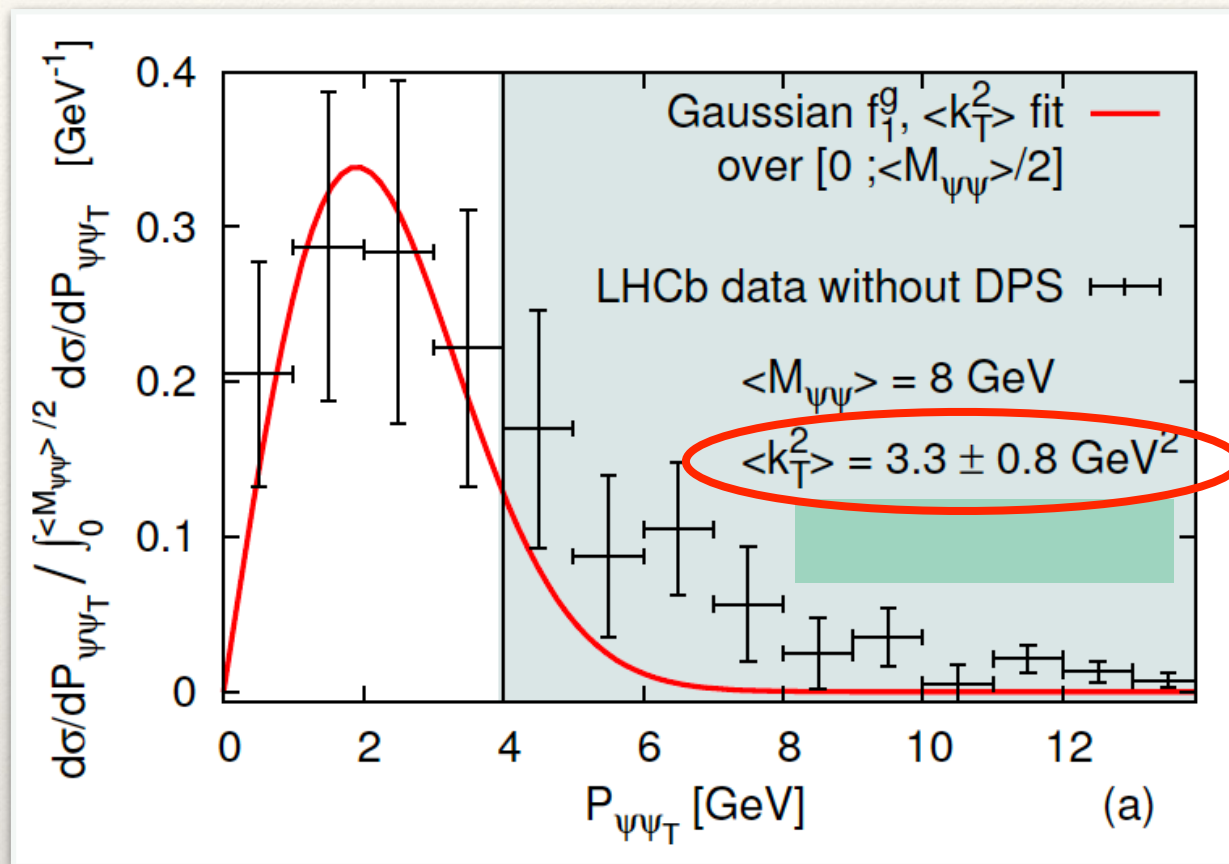
No fits for f_1 so far, Gaussian ansatz

$$f_1^g(x, k_T^2; Q = 8 \text{ GeV}) = \frac{g(x)}{\pi \langle k_T^2 \rangle} \exp(-k_T^2 / \langle k_T^2 \rangle)$$

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- One-parameter fit for $\langle k_T^2 \rangle$, $\chi^2 = 1.08$: *effective*, but not *intrinsic* width
- TMD evolution not included, needed if more data sets at different Q are available
- LHCb data corrected for double-parton scattering
- Expectation: Color Singlet Mode dominant

[Lansberg, Shao, PLB 2015; Ko, Yu, Lee, JHEP 2011; Li, Xu, Liu, Zhang, JHEP 1023]

Azimuthal modulation: no data available → predictions

Suggested observables: weighted cross section ratios

$$\frac{\int d\phi \cos(2\phi) \left[\frac{d\sigma d\phi}{dq_T} \right]_{\text{bins}}}{\int_0^{Q/2} dq_T \left[\frac{d\sigma}{dq_T} \right]_{\text{bins}}} \propto \frac{\mathcal{C}[w_3 h_1^{\perp g} f_1^g] + \mathcal{C}[w_3 f_1^g h_1^{\perp g}]}{\int_0^{Q/2} dq_T \mathcal{C}[f_1^g f_1^g]}$$

$$\frac{\int d\phi \cos(4\phi) \left[\frac{d\sigma d\phi}{dq_T} \right]_{\text{bins}}}{\int_0^{Q/2} dq_T \left[\frac{d\sigma}{dq_T} \right]_{\text{bins}}} \propto \frac{\mathcal{C}[w_4 h_1^{\perp g} h_1^{\perp g}]}{\int_0^{Q/2} dq_T \mathcal{C}[f_1^g f_1^g]}$$

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Input for linearly pol. gluon TMD: models

Model 1

$$h_1^{\perp g}(x, k_T^2) = \frac{M^2}{\langle k_T^2 \rangle} \frac{g(x)}{\pi \langle k_T^2 \rangle} \exp \left(1 - \frac{3 k_T^2}{2 \langle k_T^2 \rangle} \right)$$

Model 2: Saturation of positivity bound

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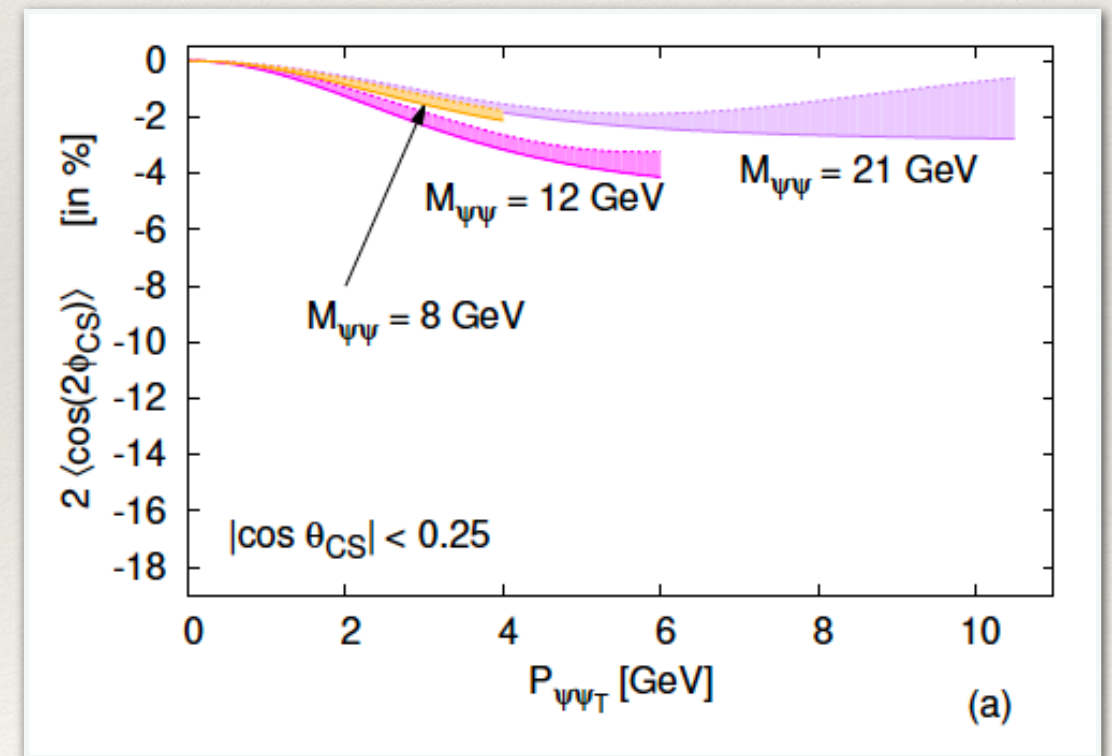
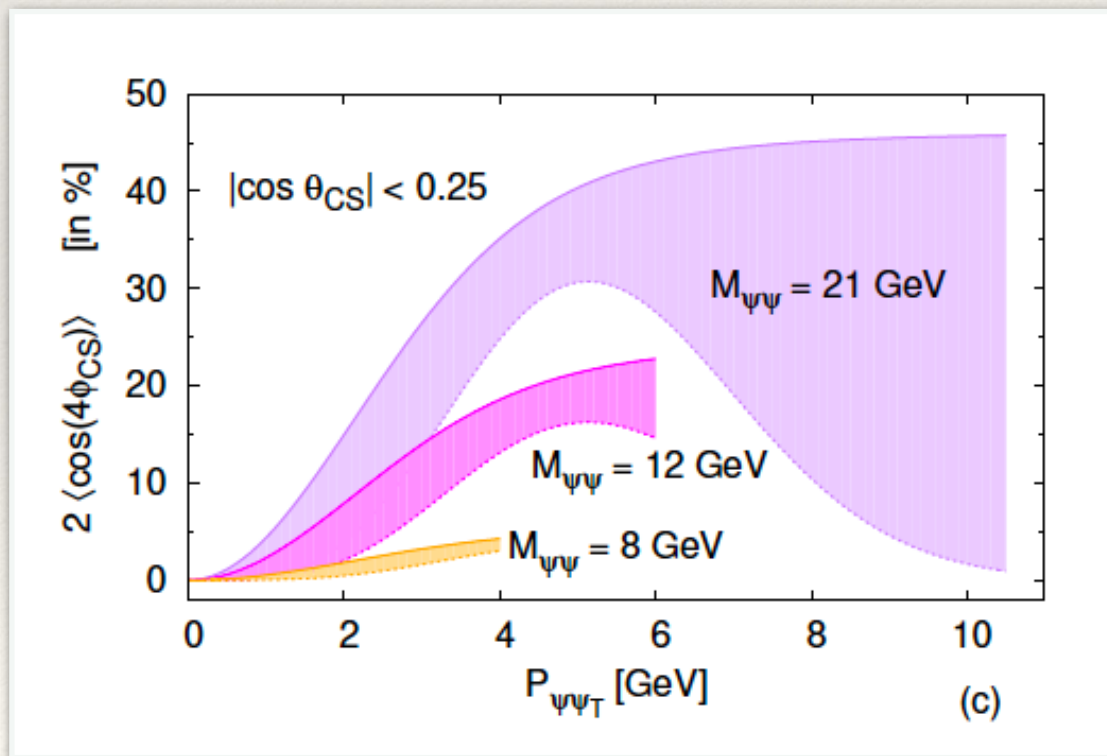
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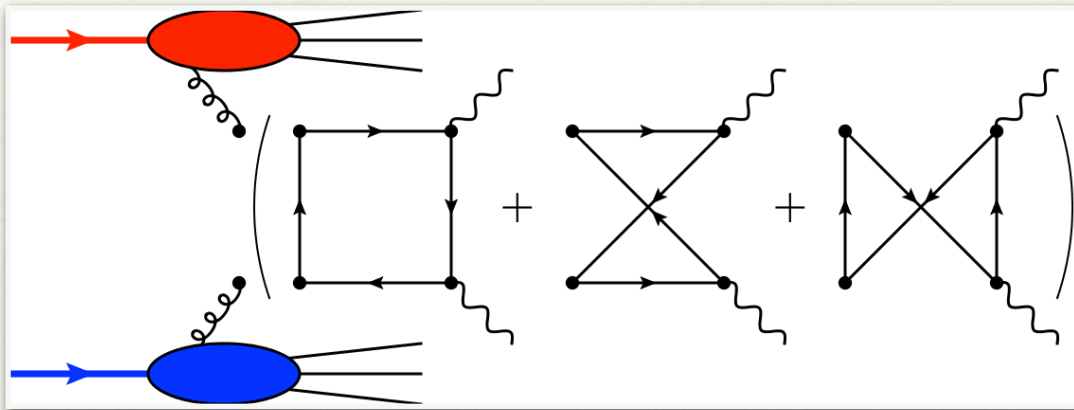
- Large effects for $\cos(4\phi)$ for Q larger than threshold: 10% - 20%, $\cos(2\phi)$ few percent

Summary & Outlook

- ❖ Gluon TMDs → new aspects on the 3D gluonic structure of the nucleon → linear gluon polarization
- ❖ Several final states in pp - collisions: Leptons, Photons, Heavy Quarkonium states
- ❖ Promising: J/ψ pairs at the LHC, particularly large: $\cos(4\phi)$ azimuthal modulation
- ❖ Data is coming: first extraction of unpolarized gluon TMD
- ❖ Important for EIC: An idea what to expect for gluon TMDs
- ❖ Outlook:
 - Improve fit for unpolarized gluon TMD including ATLAS, CMS data and TMD evolution
 - Extend helicity formalism to spin-dependent observables and gluon TMDs (sPHENIX)

Diphoton production

[Qiu, M.S., Vogelsang, PRL 2011]



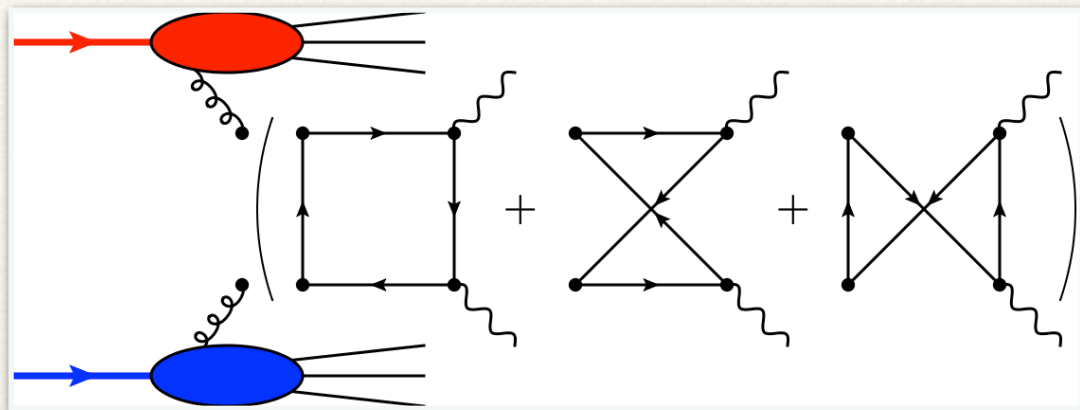
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only $\cos(4\phi)$ -modulation purely gluonic
- ❖ Photons are difficult to measure:
 $\gamma\gamma$ - production: background from π^0 - decays

Linearly polarized gluons potentially useful for Higgs - production

[Boer, den Dunnen, Pisano, MS, Vogelsang, PRL 2012, 2013]

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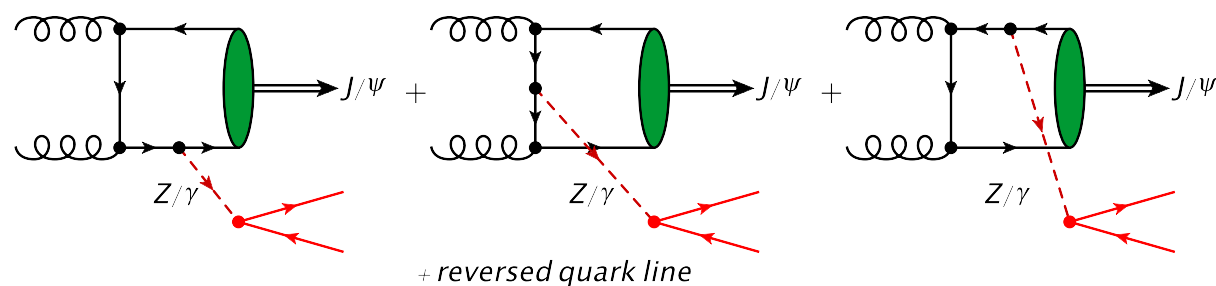
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Associated $\Upsilon(J/\psi) + \gamma/Z$ - production at the LHC

[den Dunnen, Lansberg, Pisano, M.S., PRL 2014 / Lansberg, Pisano, M.S., NPB 2017]

Quarkonium in color singlet configuration compatible with TMD factorization

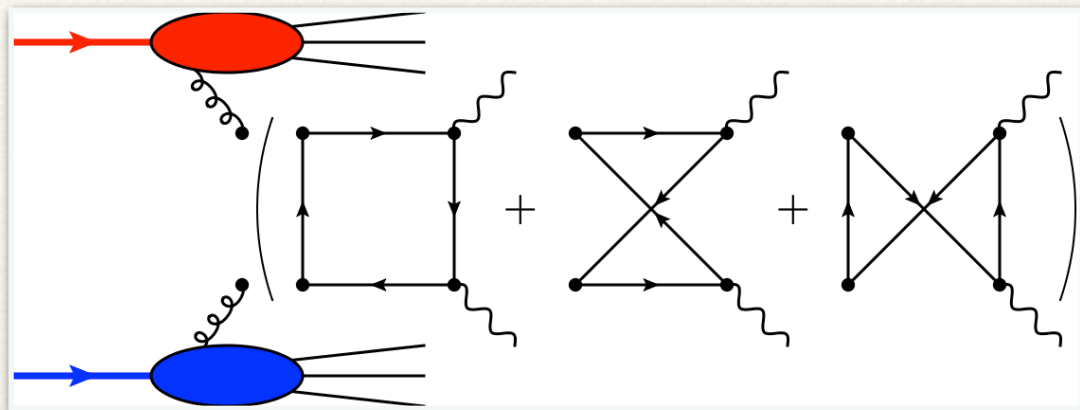


- ❖ main contribution from gluon fusion
- ❖ Υ - production:
Color Octet (Fragm.) $\ll$ Color Singlet ✓
- ❖ Real γ : $F_2 = 0$, $\cos(2\phi) \sim 3\%$, $\cos(4\phi) \sim 1\%$!
- ❖ Real Z: $F_1 \gg F_2, F_3, F_4$

first glimpse of linearly polarized gluons may be possible at the LHC

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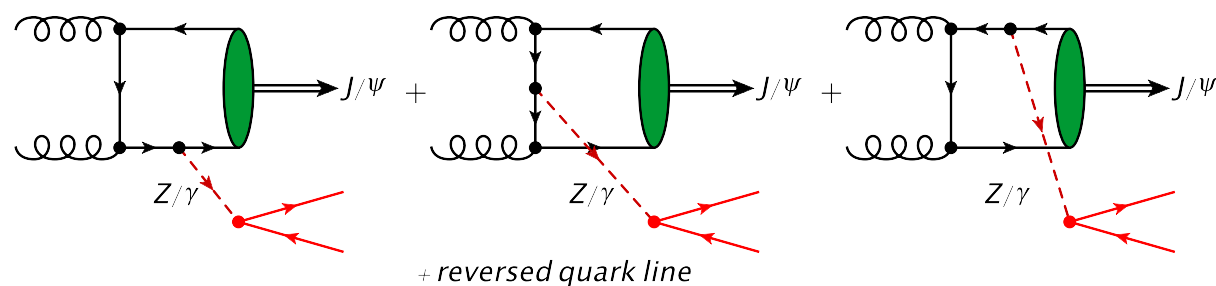
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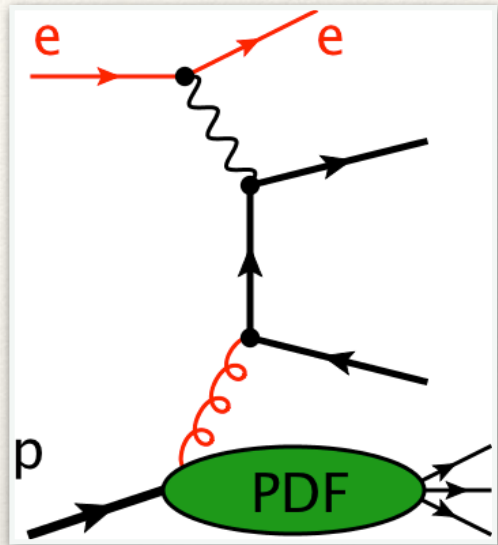


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Processes sensitive to gluon TMDs

Processes sensitive to gluon TMDs



Heavy Quark production in e + p - collisions

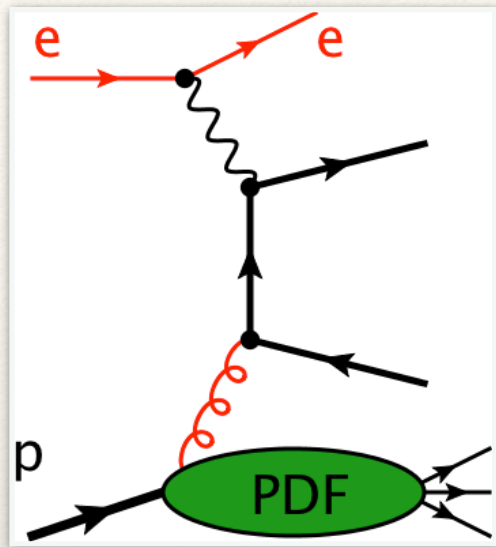
[Boer, Brodsky, Mulders, Pisano, PRL 2011]

$$e + p^{(\uparrow)} \rightarrow e' + \text{jet}(c, b) + \text{jet}(\bar{c}, \bar{b}) + X$$

TMD factorization ok!

Spin dependent (+independent...) gluon TMDs: EIC would be ideal!

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[Boer, Brodsky, Mulders, Pisano, PRL 2011]

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unpolarized cross section:

$$\frac{d\sigma_{UU}}{dq_T} \propto (F_1 + F_2 \cos(2\phi))$$

azimuthally independent term:

$$F_1 \propto f_1^g(x, q_T)$$

→ unpol. gluon distribution

azimuthally dependent term:

$$F_2 \propto h_1^{\perp g}(x, q_T)$$

→ linearly pol. gluons