

# LP<sup>3</sup> qPDF Results



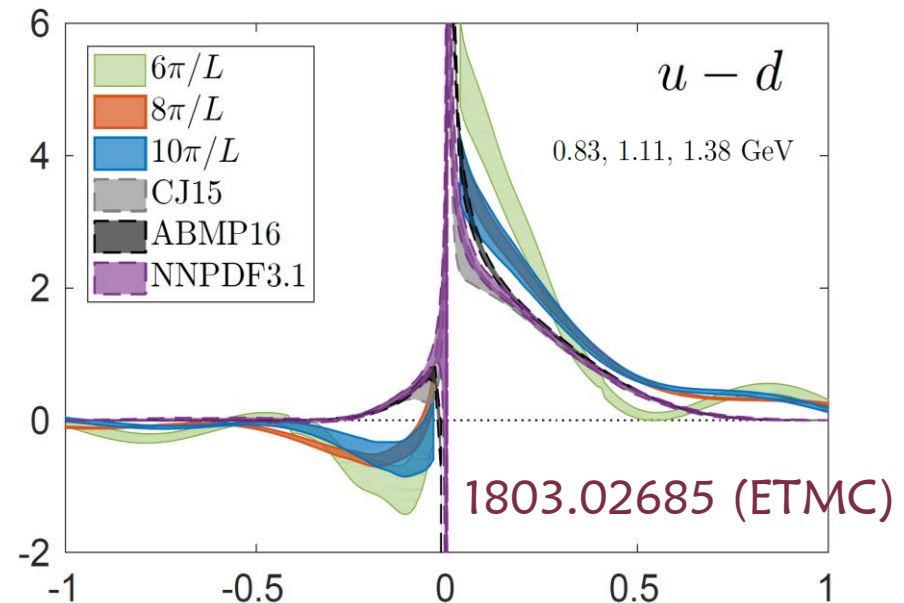
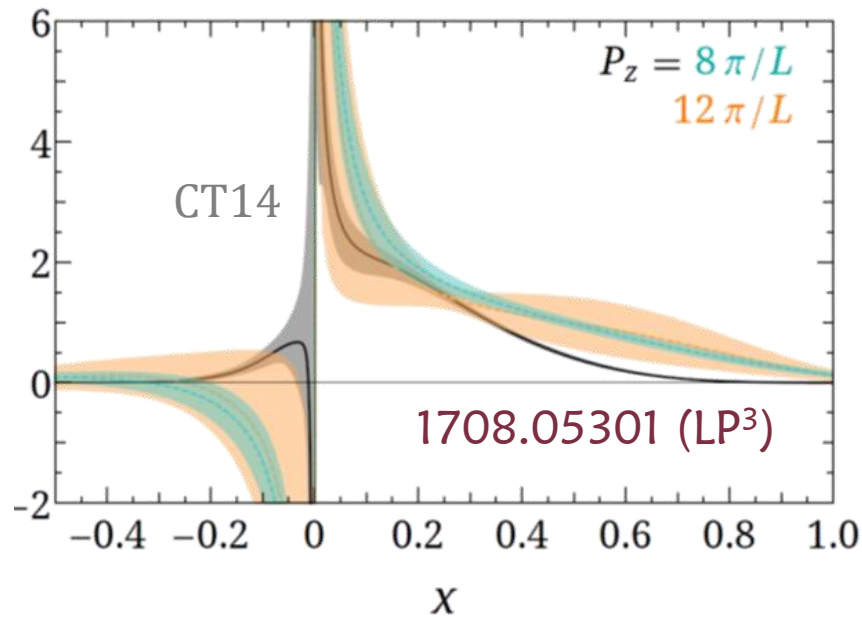
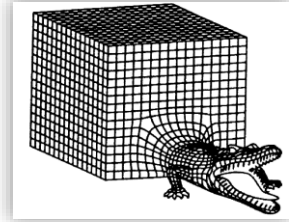
HUEY-WEN LIN

This work is supported by NSF under grant PHY 1653405  
 "CAREER: Constraining Parton Distribution Functions for New-Physics Searches"

# Physical Pion Mass Results

§ Exciting! Two collaborations' results at physical pion mass

- ∞ Boost momenta  $P_z \leq 1.4$  GeV
- ∞ Study of systematics still needed



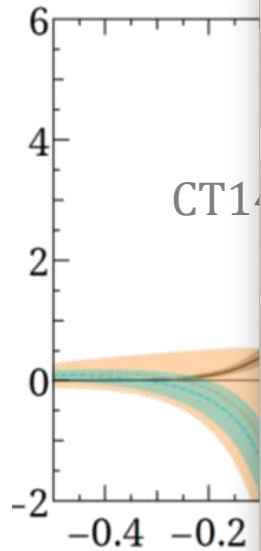
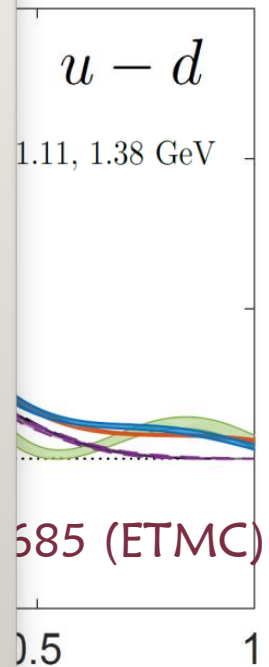
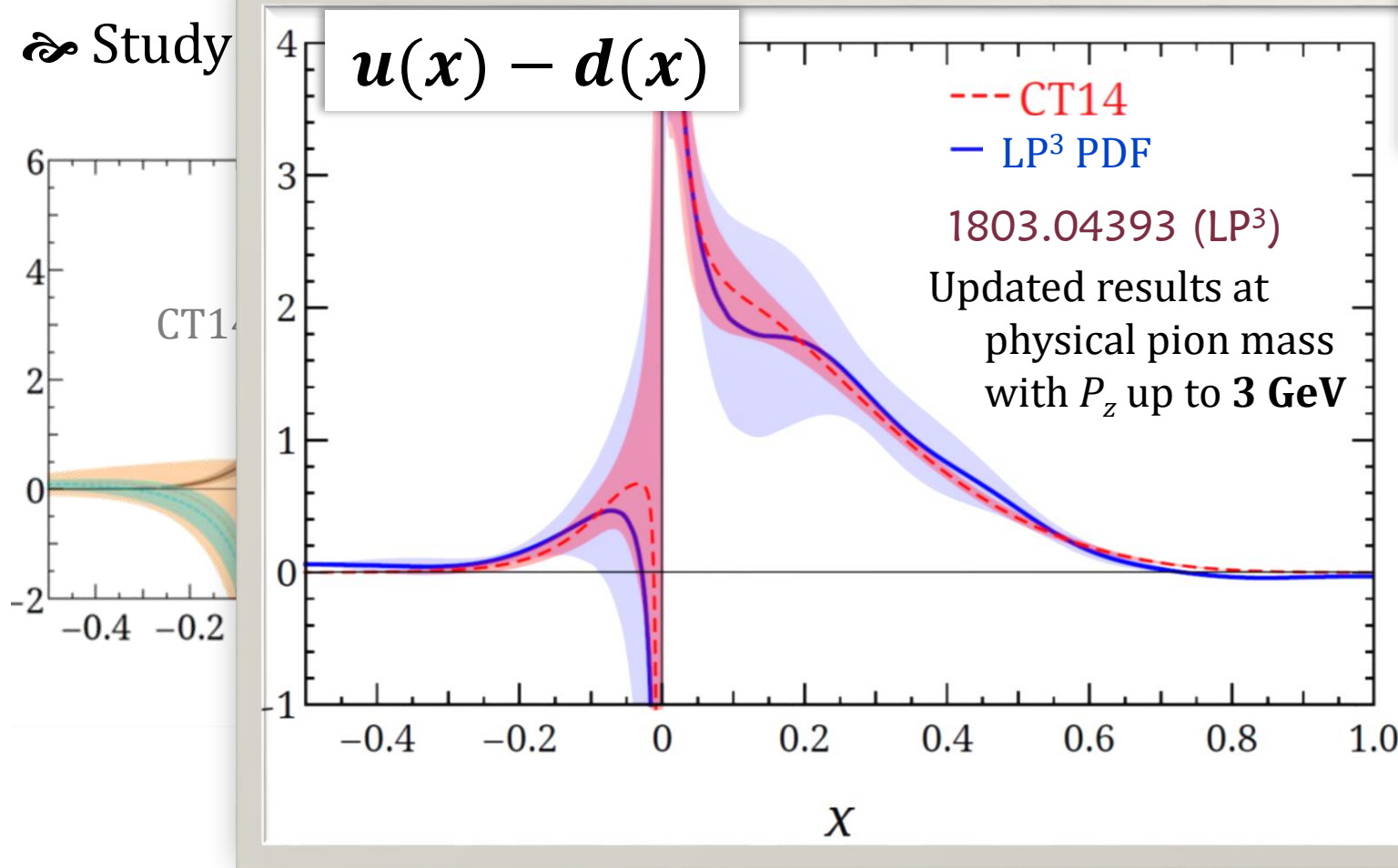
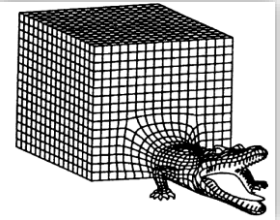
Not using parametrization (e.g.  $xf(x, \mu_0) = a_0 x^{a_1} (1-x)^{a_2} P(x)$ )  
Less pretty results; less likely to exactly coincide with global fits.

# Physical Pion Mass Results

§ Exciting! Two collaborations' results at physical pion mass

∞ Boost  $P_z \sim 1.4 \text{ GeV}$

∞ Study



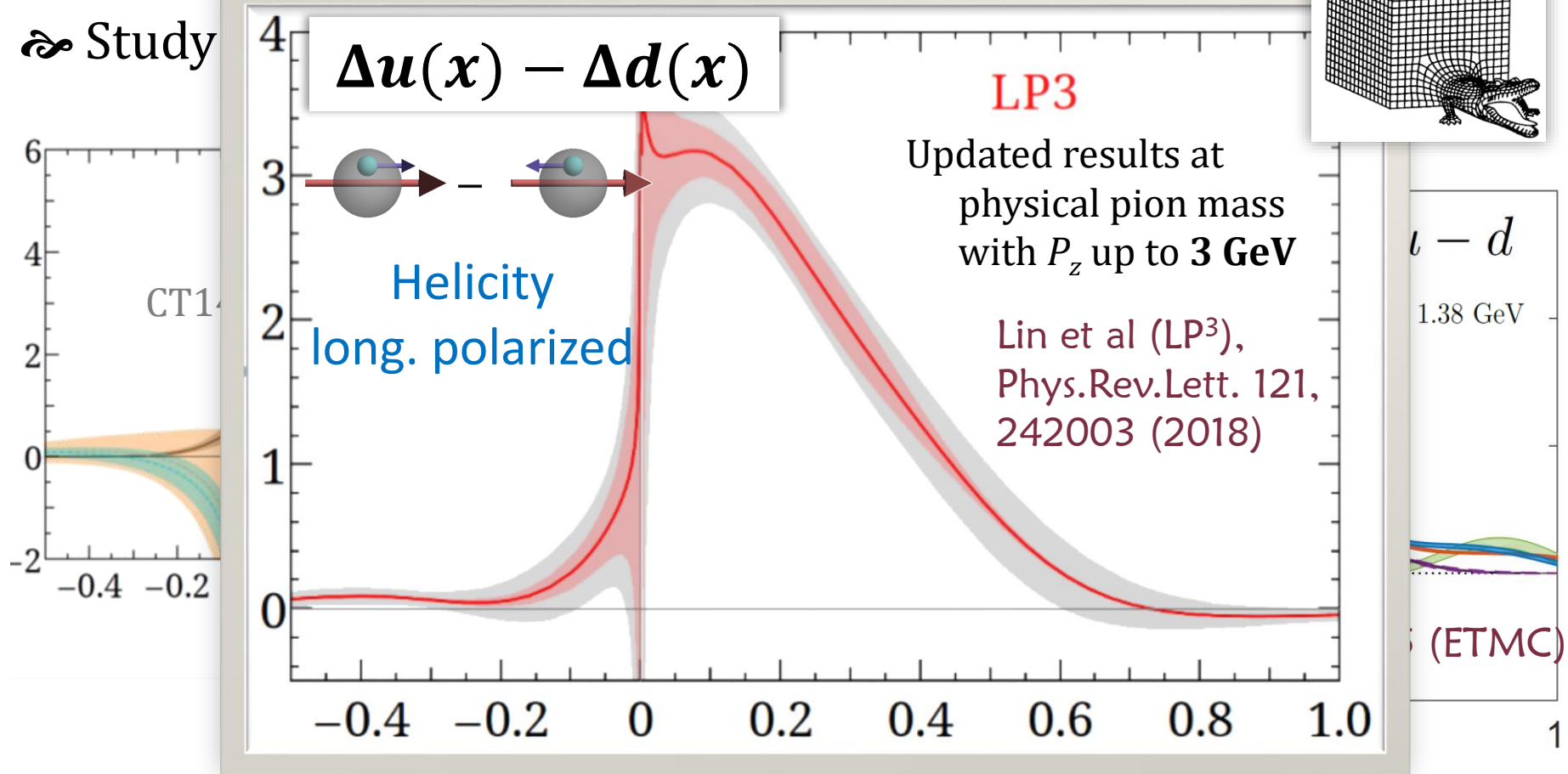
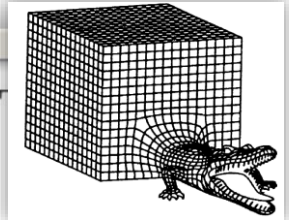


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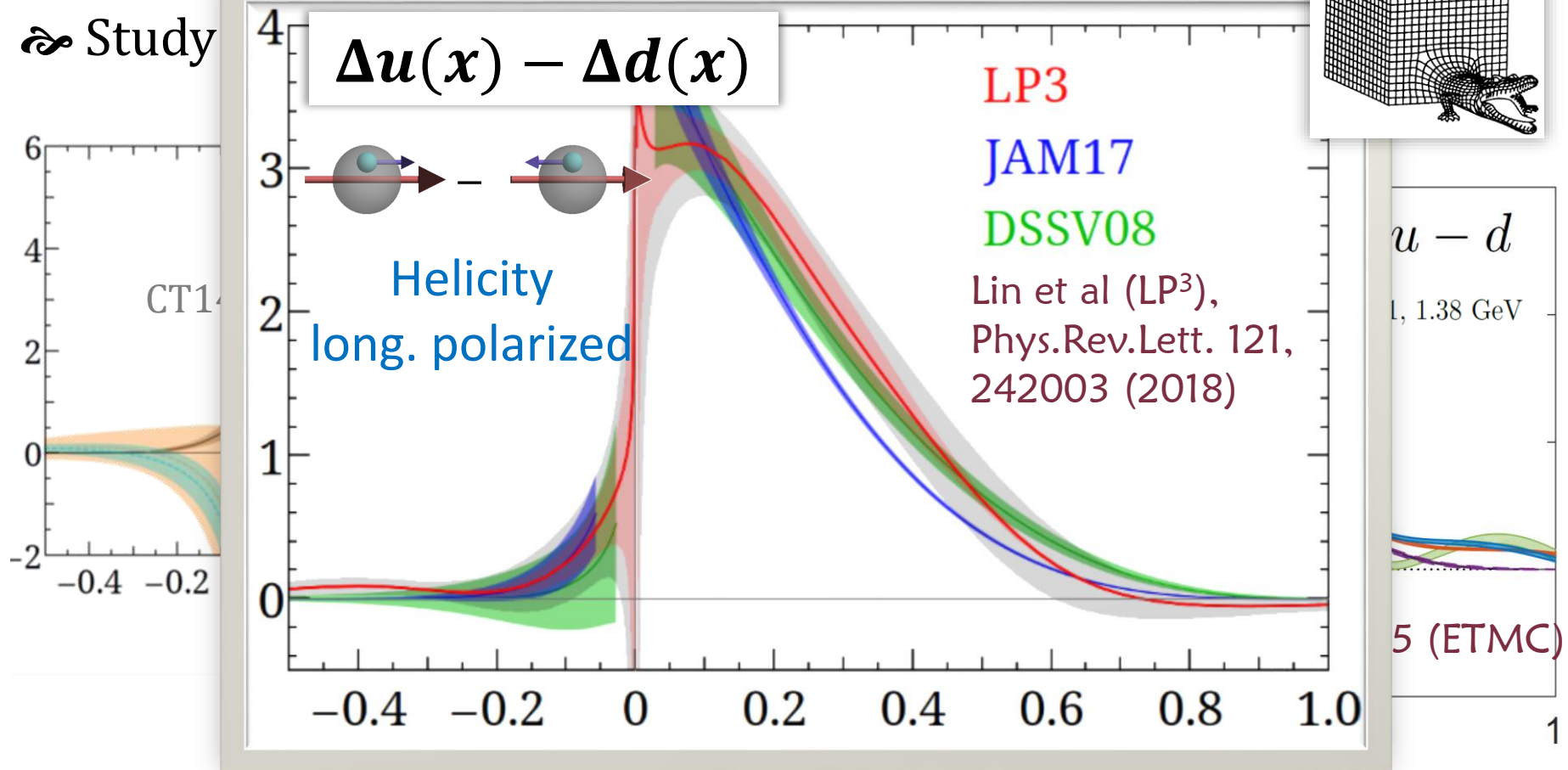
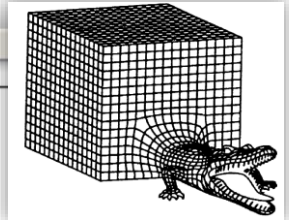


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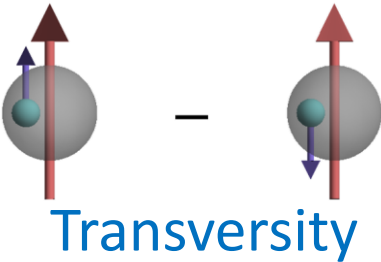


# Physical Pion Mass Results

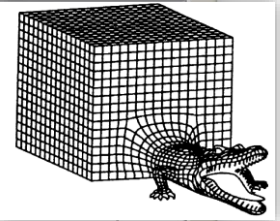
§ Exciting! Two collaborations' results at physical pion mass

∞ Bo  
∞ Stu

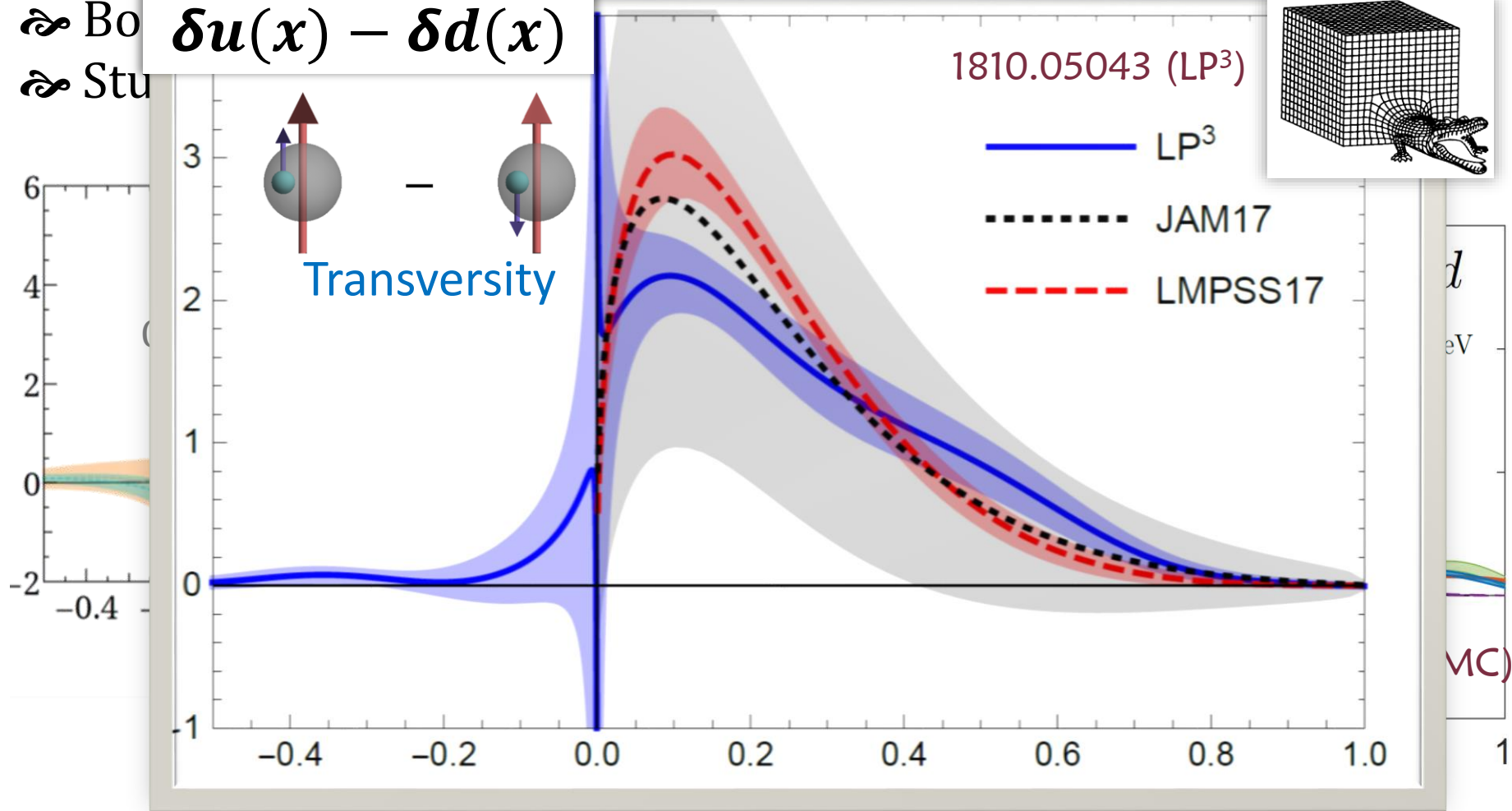
$$\delta u(x) - \delta d(x)$$



1810.05043 (LP<sup>3</sup>)



— LP<sup>3</sup>  
- - - JAM17  
- - - LMPSS17



$d$   
eV  
MC)

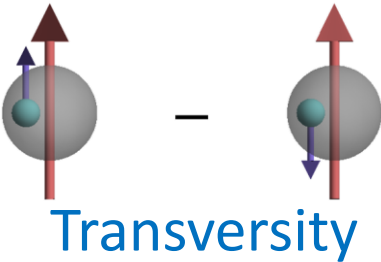
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# Physical Pion Mass Results

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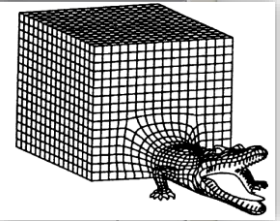
∞ Bo  
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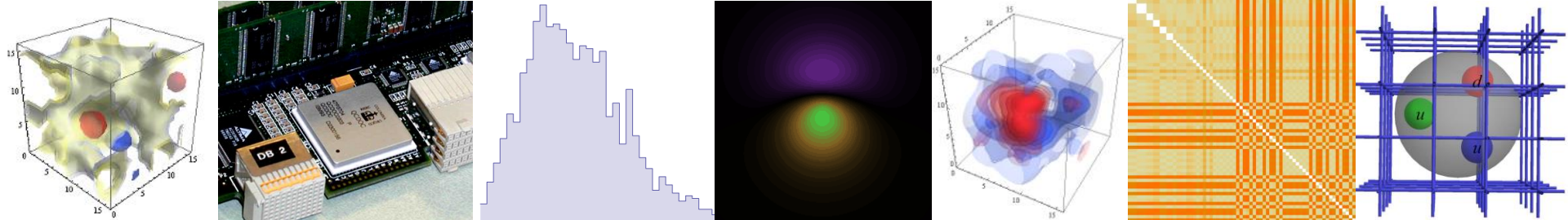


Very encouraging results  
More systematics studies are underway:  
multiple lattice spacings, volumes, etc.

$d$   
eV  
MC)

1





# MythBusting Criticism of $LP^3$ qPDF Calculations



HUEY-WEN LIN

This work is supported by NSF under grant PHY 1653405  
“CAREER: Constraining Parton Distribution Functions for New-Physics Searches”



# Outline

Apologize in advance for a rather technical talk

Need to clear up many false rumors

§ The need for large boost momentum

⇒ Nothing to do with lattice

§ Please stop using plateau fit

⇒ What's the point of getting many wrong answers?

§ I will take more criticism

⇒ As time allows



# *The Necessity of Large Boosted Momentum without modeling*



# Myth

One can recover the full PDF using  
boost momentum below 2 GeV.





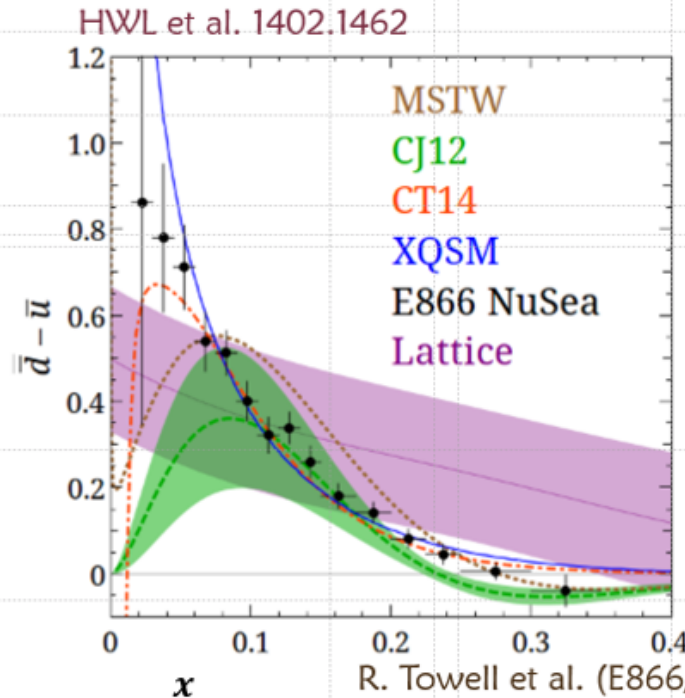
# Backstory

§ Many of you are old enough to remember this:

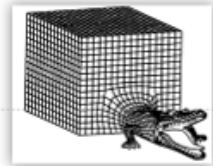
## Sea Flavor Asymmetry

§ First time in LQCD history to study antiquark distribution!

$\propto M_\pi \approx 310 \text{ MeV}, a \approx 0.12 \text{ fm}$



$$\bar{q}(x) = -q(-x)$$



Lost resolution in  
small- $x$  region

Future improvement:  
larger lattice volume

$$\int dx (\bar{u}(x) - \bar{d}(x)) \approx -0.16(7)$$

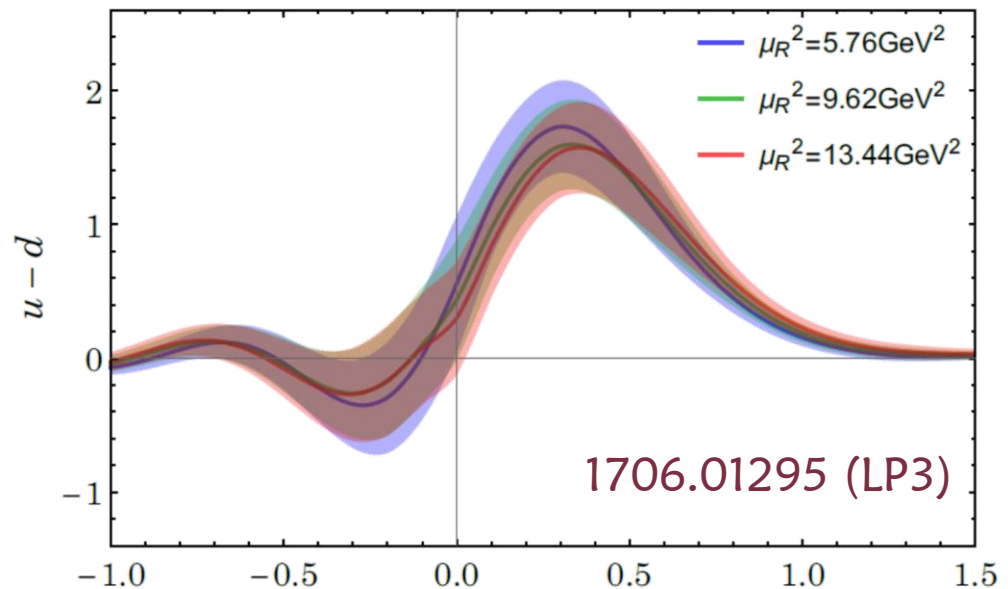
Experiment	$x$ range	$\int_0^1 [\bar{d}(x) - \bar{u}(x)] dx$
E866	$0.015 < x < 0.35$	$0.118 \pm 0.012$
NMC	$0.004 < x < 0.80$	$0.148 \pm 0.039$
HERMES	$0.020 < x < 0.30$	$0.16 \pm 0.03$

**Caveat:** These matrix elements are not properly renormalized

# Backstory

§ Efforts by multiple collaborations have been devoted into working on lattice renormalization

∞ We finally obtained the renormalized ME, and the renormalized PDF results puzzled us for months!

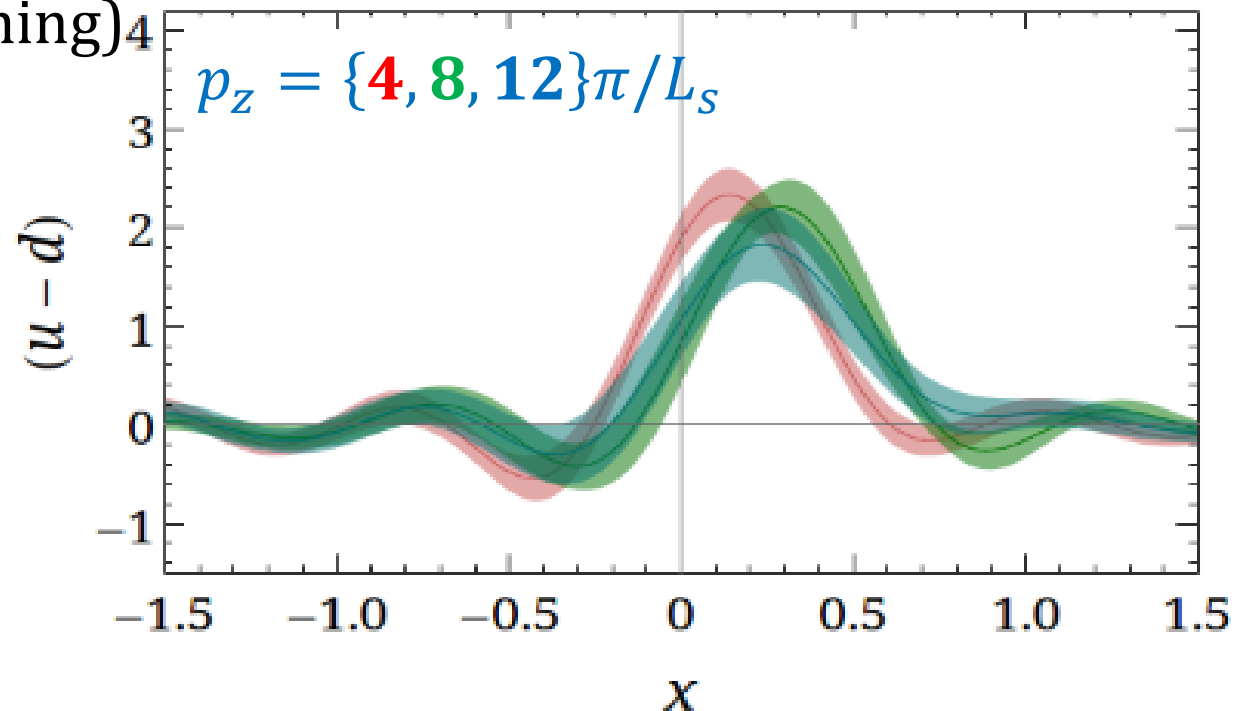


∞ We finally posted the results to arXiv, since<sup>x</sup> others had already posted their renormalized result

# Backstory

§ Immediately, we checked a different lattice and observed the same thing and worse, since we had more “z” data!

∞ Results from 2017 Summer at physical pion mass  
(before matching)



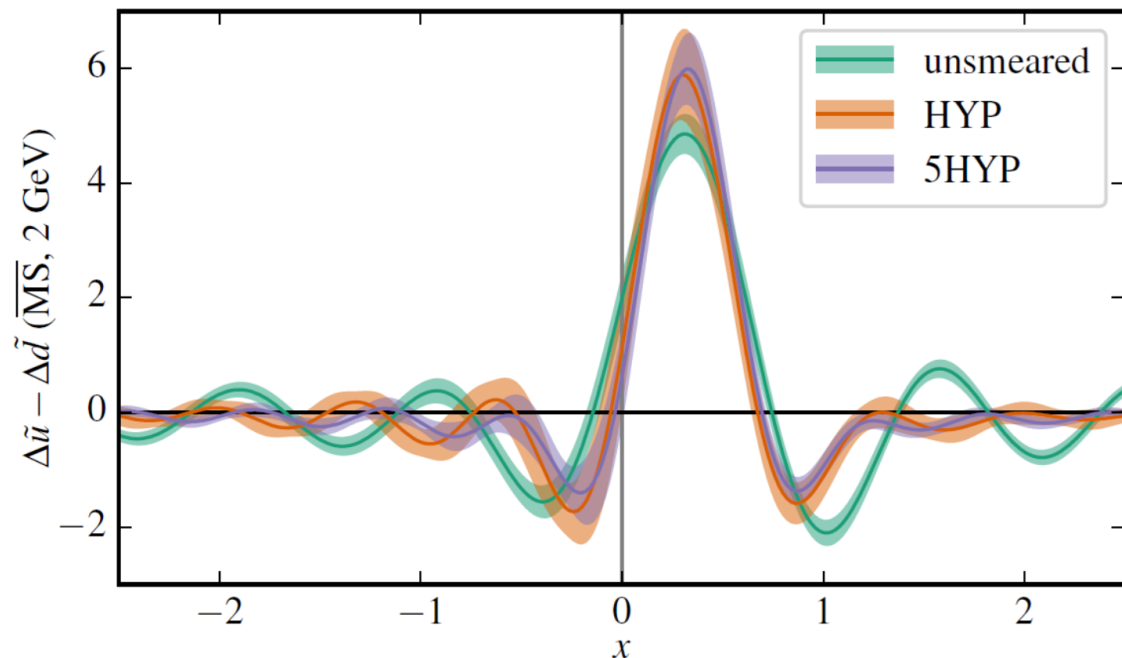


# Backstory

## § Something is obviously wrong!

- ⌘ This effect was missed by earlier ETMC work due to the small  $z$  and momentum combination

J. Green et al 1707.07152



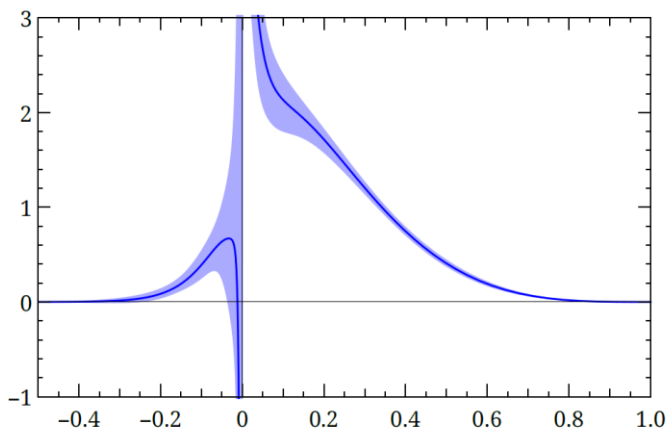
- ⌘ We called it “oscillation” and F.T. truncation artifacts
- ⌘ For a while, people had no idea what we are talking about

# Focus on Continuum

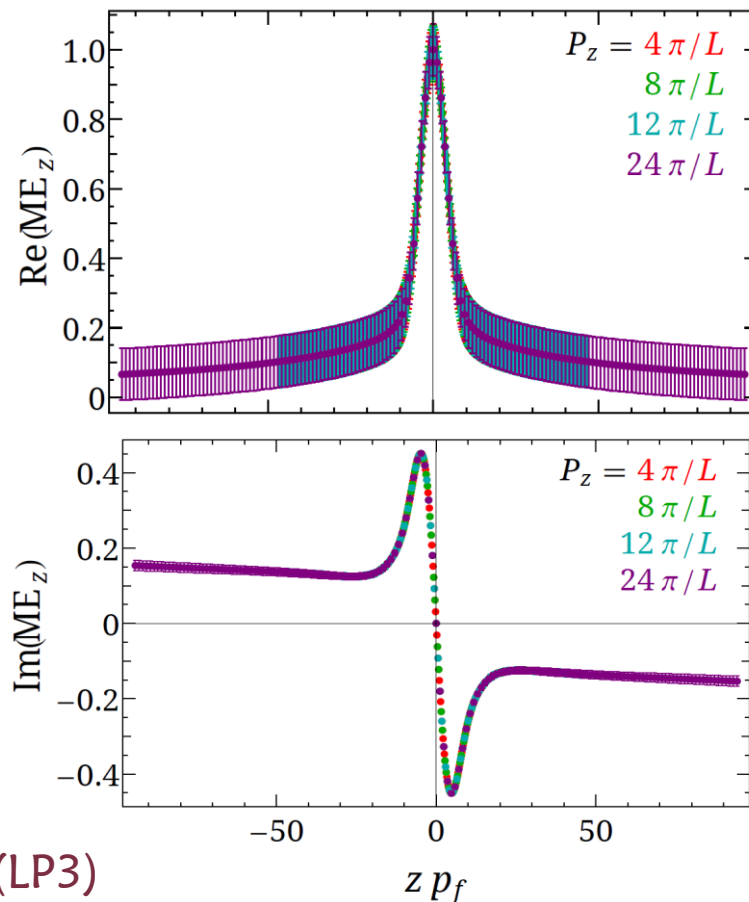
§ Not a lattice problem but a Fourier-transform issue

§ Simple exercise with CT14 PDF 1506.07443

Pick your favorite PDF  
(CT14 here)



Fourier to  
coordinate space  
at some momentum



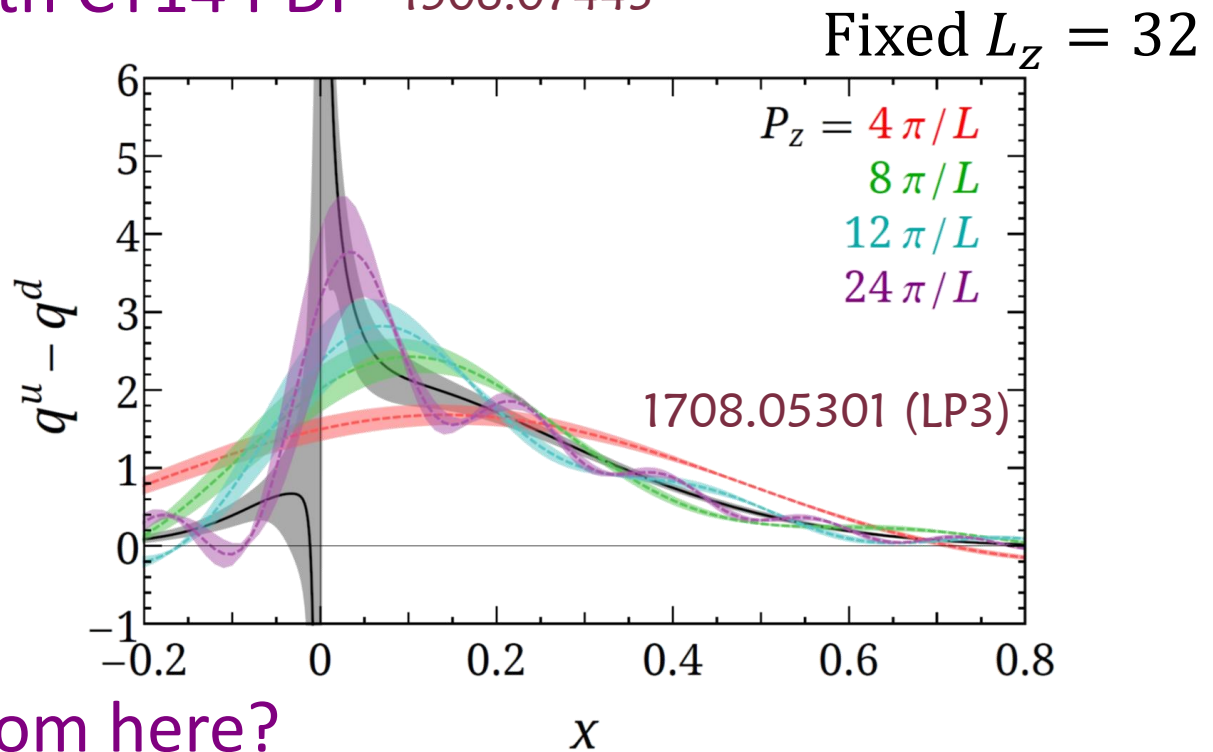
1708.05301 (LP3)

# Focus on Continuum

§ Not a lattice problem but a Fourier-transform issue

§ Simple exercise with CT14 PDF 1506.07443

Inverse Fourier  
transform back to  
momentum space



§ What do you do from here?

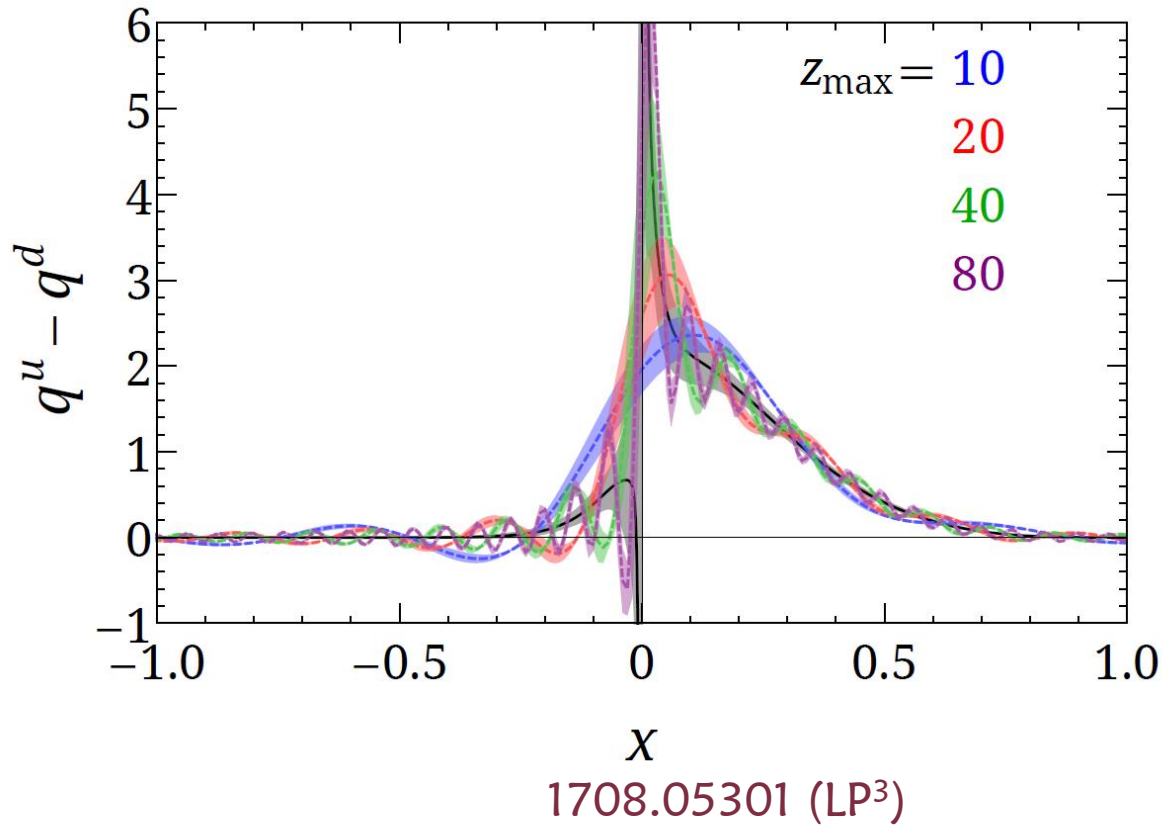
- ⌘ Throw up your hands and add 100% errorbars across all  $x$
- ⌘ Find some way to salvage as much information as possible



# Continuum

§ Fix  $P_z$  and use large  $z_{\max}$ ?

Fixed  $P_z = 24\pi/L$



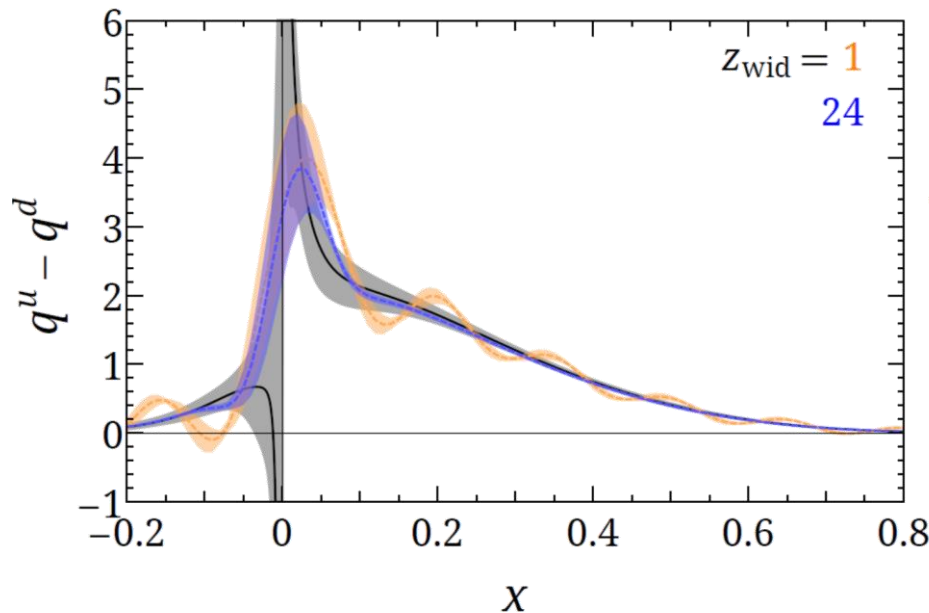
# Continuum

§ Luckily, we know the answer

§ Two possible solutions proposed (likely more) <sup>1708.05301 (LP<sup>3</sup>)</sup>

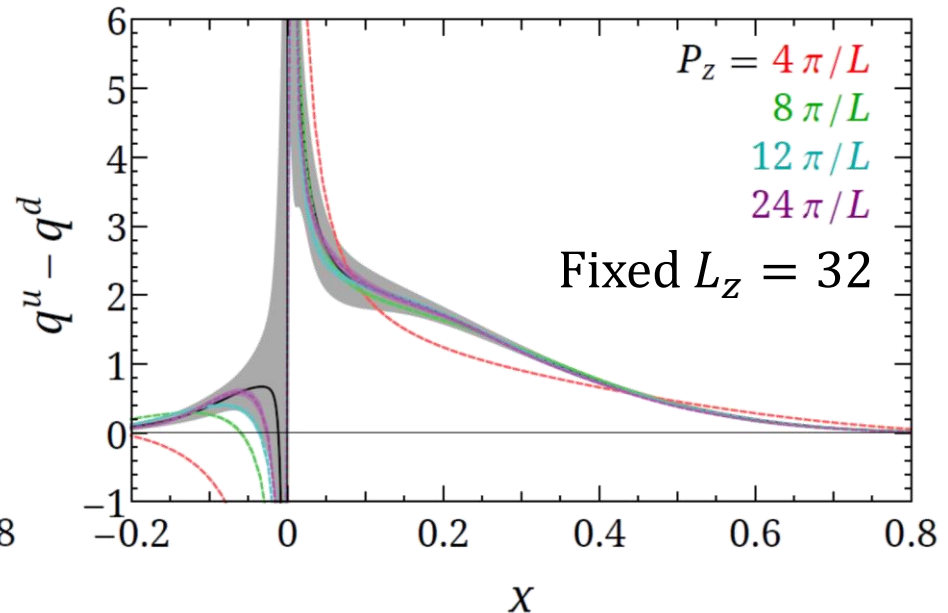
Filter approach

$$F(z_{\text{lim}}, z_{\text{wid}}) = \frac{1 + \text{erf}\left(\frac{z + z_{\text{lim}}}{z_{\text{wid}}}\right)}{2} \frac{1 - \text{erf}\left(\frac{z - z_{\text{lim}}}{z_{\text{wid}}}\right)}{2}$$

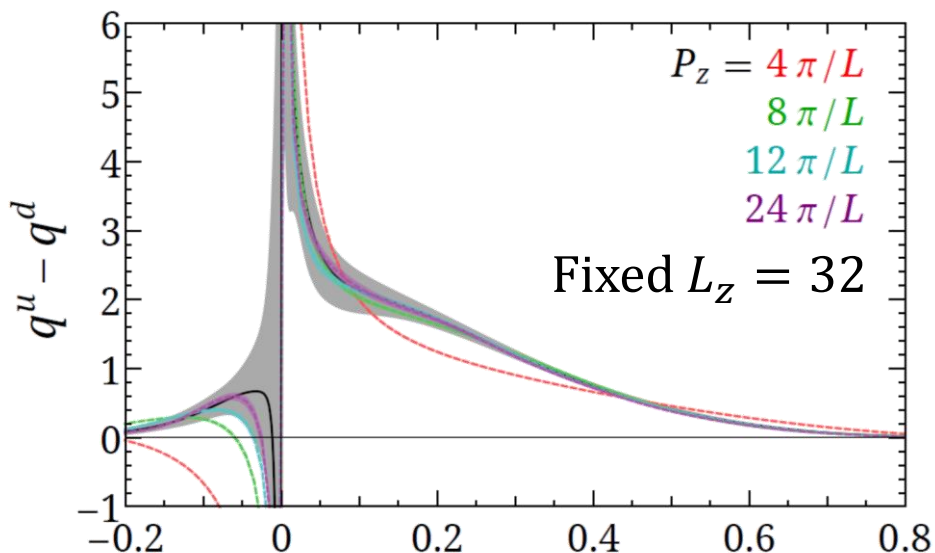


Derivative approach

$$q(x) = \int_{-z_{\text{max}}}^{+z_{\text{max}}} dz \frac{-P_z}{2\pi} \frac{e^{ixP_z z}}{iP_z x} h'(z)$$

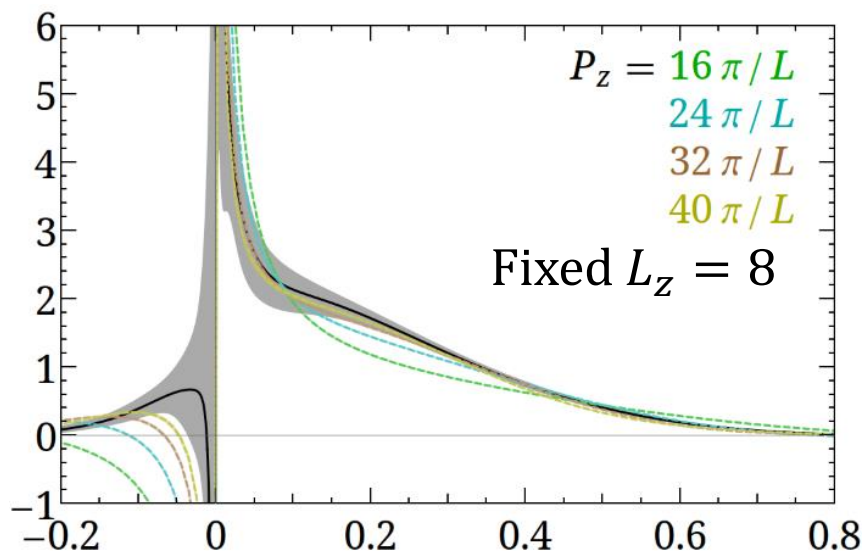


# Continuum



## § What lessons learned here?

- Given  $L_z \approx 15$ , you need large momentum to just get the sign of the antiquark correct!
- With small  $zP_z$ , you will miss over the majority of  $x$
- Not just a quasi-PDF problem
- Going for large  $P_z$  is an unavoidable direction for any method that requires steps similar to Fourier transformation



# Continuum

## § What lessons learned here?

➤ Given  $L_z \approx 15$ , you need large momentum to just get

mark

will miss

problem

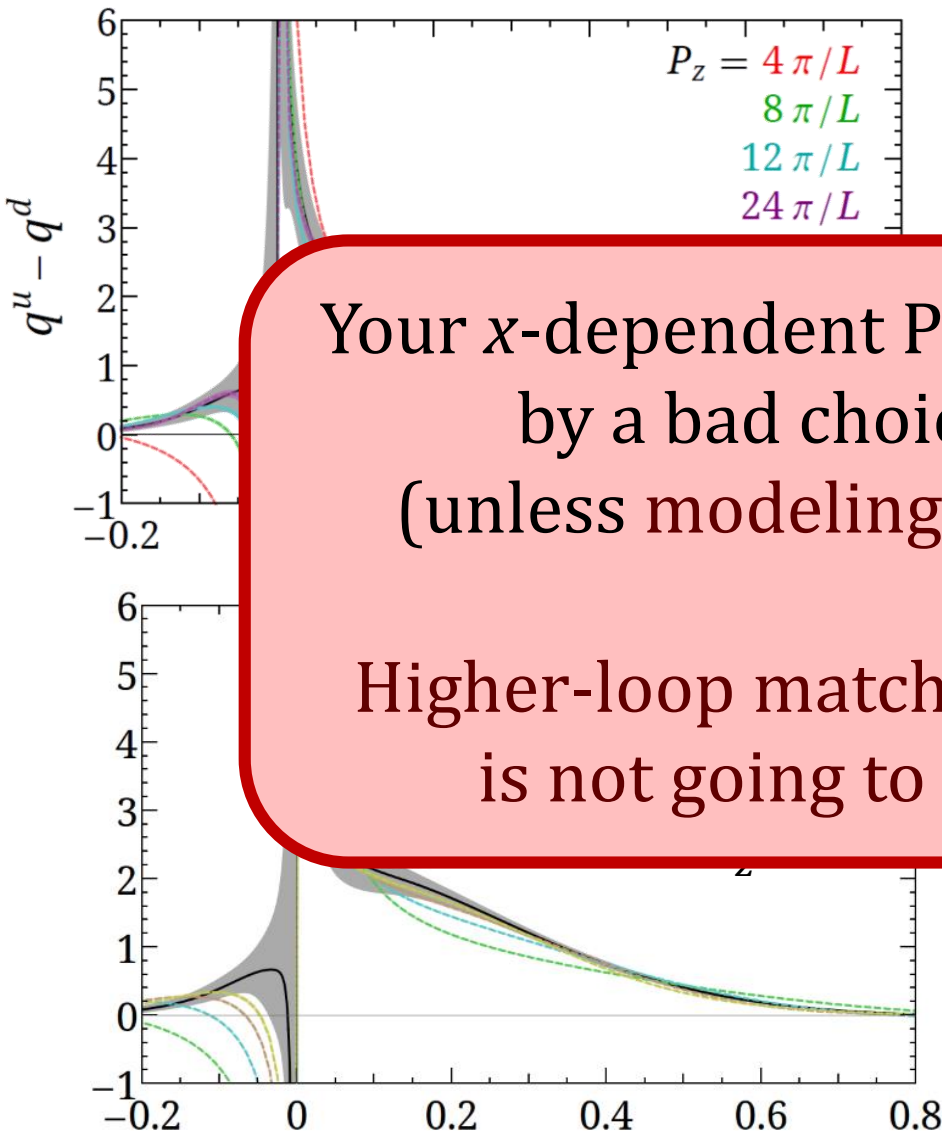
in

and results are not the same for any

method that requires steps similar to Fourier transformation

Your  $x$ -dependent PDFs will be doomed by a bad choice of  $\max z P_z$ ! (unless modeling  $z P_z$  dependence)

Higher-loop matching in LaMET later is not going to do much for it!



# Alternative Solution

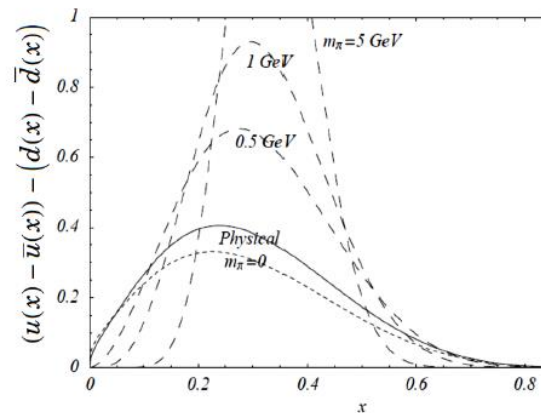
## § One can model the lattice data?

- ⌘ Replaces cutoff systematic by model dependence
- ⌘ No obvious advantage for direct- $x$  approach from moments  
(*already did this in 2001!*)
- ⌘ Nor guide the global PDF with correct  $x$ -forms or improve that fit-form dependence

## § What can we learn about the $x$ -distribution?

- ⌘ Make an ansatz of some smooth form for the distribution and fix the parameters by matching to the lattice moments

$$xq(x) = \alpha x^b (1-x)^c (1 + \epsilon \sqrt{x} + \gamma x)$$



Cannot separate valence-quark contribution from sea

New idea needed to access the sea!

W. Detmold et al, Eur.Phys.J.direct C3 (2001) 1-15



Huey-Wen Lin — Lattice 2013 @ Mainz

3

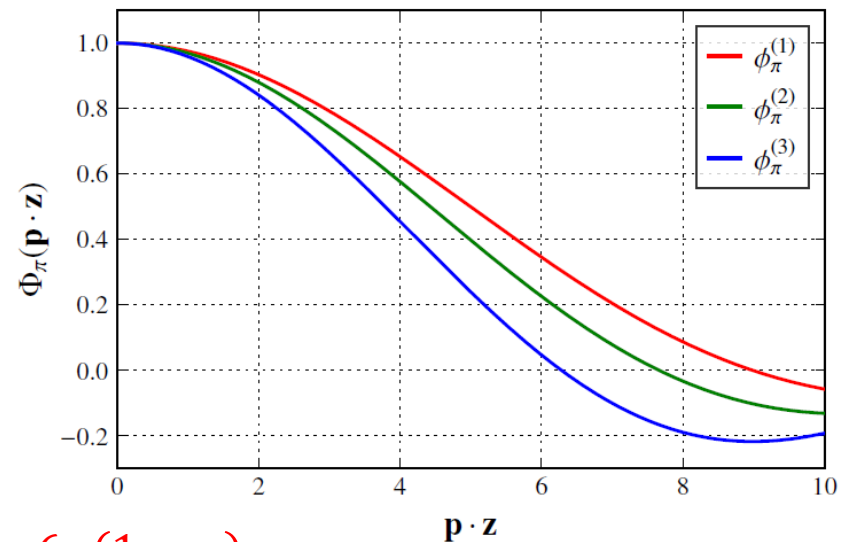
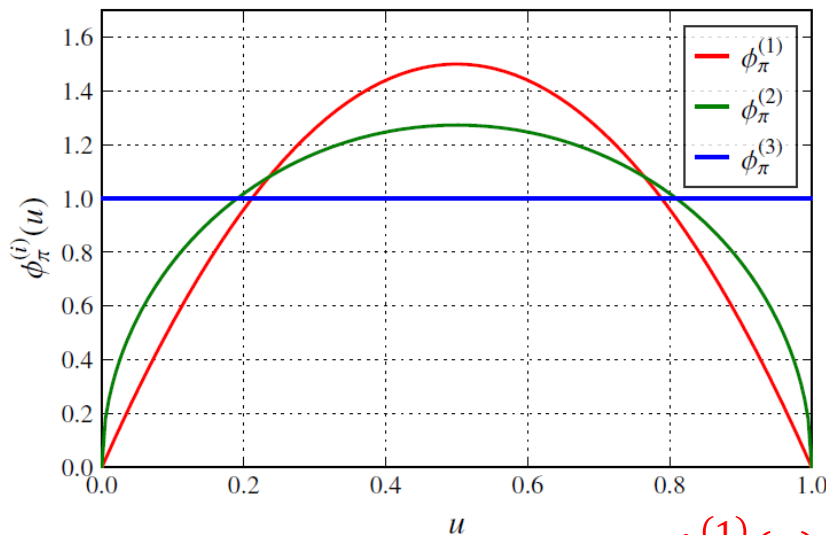


# Alternative Solution

## § Coordinate-space PDFs?

∞ Lose sensitivity in constraining PDF forms

1709.04325(RQCD)



$$\begin{aligned}\phi_\pi^{(1)}(u) &= 6u(1-u) \\ \phi_\pi^{(2)}(u) &= \frac{8}{\pi}\sqrt{u(1-u)} \\ \phi_\pi^{(3)} &= 1\end{aligned}$$

# *Please Stop Using Plateau Fit*



# Myth

We should all use the plateau fit!  
We can control excited states using large  
source-sink separation.



# What is Plateau Fit?

## § A “ruler” fit from a ground-state dominated world

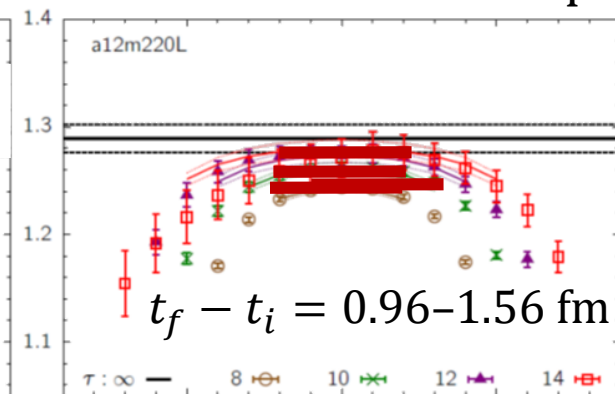
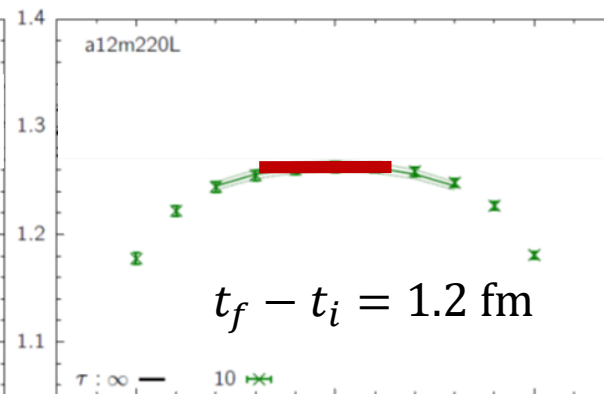
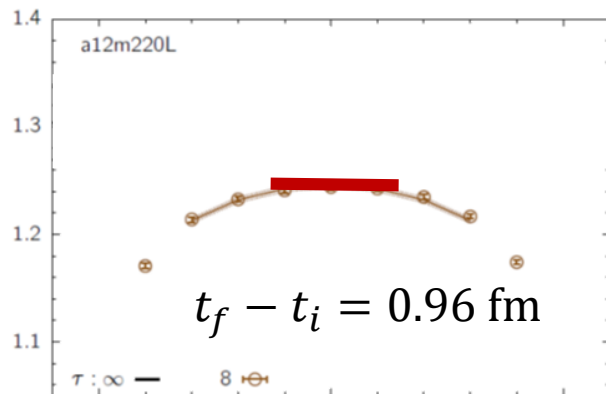
⌘ Advantage: simple to eyeball a straight line

$$C^{3\text{pt}}(t_f, t, t_i) = |\mathcal{A}_0|^2 \langle 0 | \mathcal{O}_T | 0 \rangle e^{-M_0(t_f - t_i)}$$

$$C^{2\text{pt}}(t_f, t_i) = |\mathcal{A}_0|^2 e^{-M_0(t_f - t_i)}$$



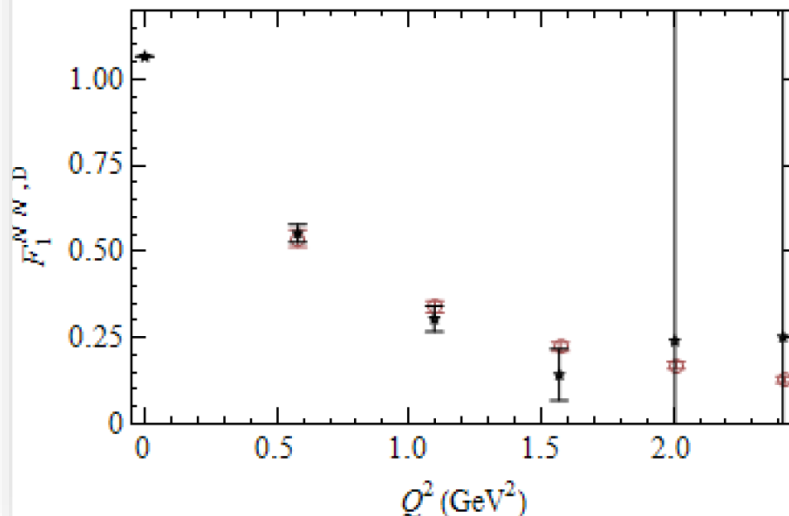
⌘ You get **wrong** answers many times until you get to large  $t_{\text{sep}}$



# No More Plateau Fit!

- § Since 2008, I have advocated abandoning plateau fits at large source-sink separation
- § Why waste precious computing resources in the regions where noise dominates?

## ◆ Results



- ◆ Consistent with conventional ratio approach
- ◆ Smaller errorbar at larger  $Q^2$
- ◆ Normally smearings are tuned to eliminate higher-energy contribution

HL@Lattice2008

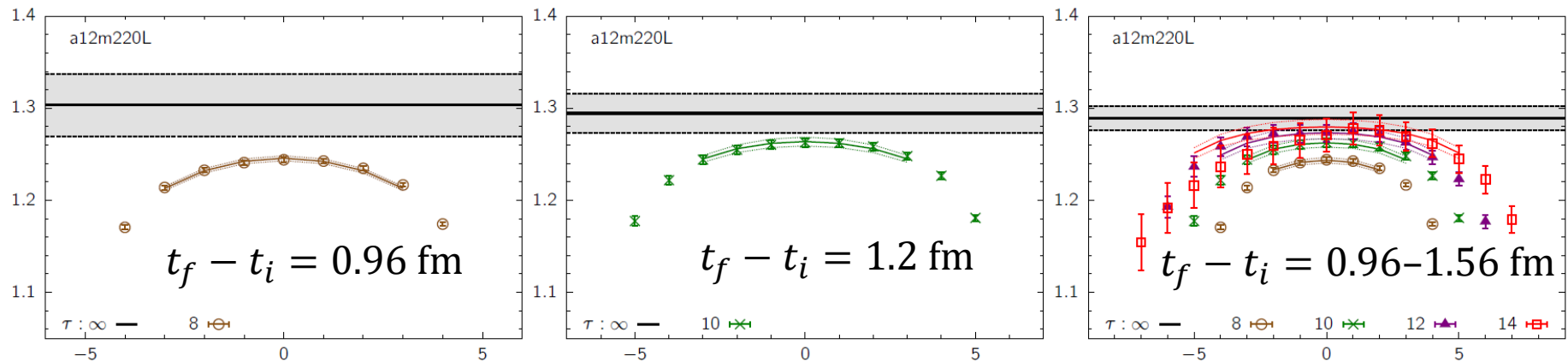
∞ The only people who listens mainly to claim credits:  
We have been doing this for meson cases for ages! This is nothing new!



# What is Plateau Fit?

## § What isn't plateau fit?

- ↻ Multistate fit to multiple separations
- ↻ Converts excited-state systematic into statistical uncertainty

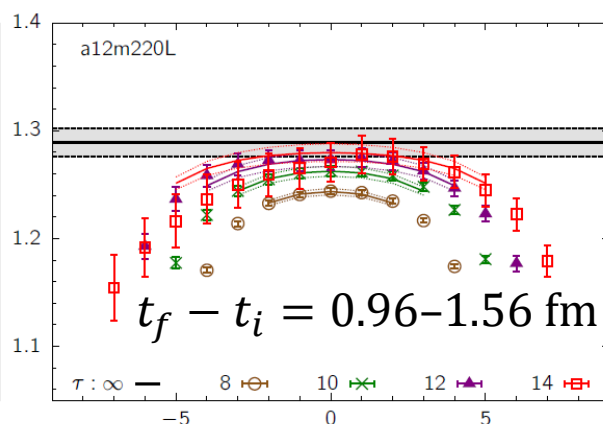
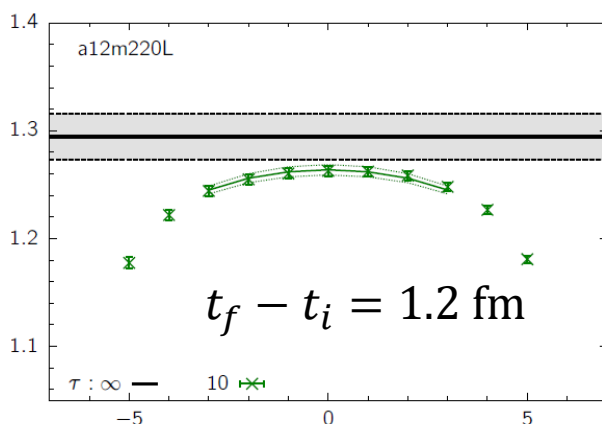
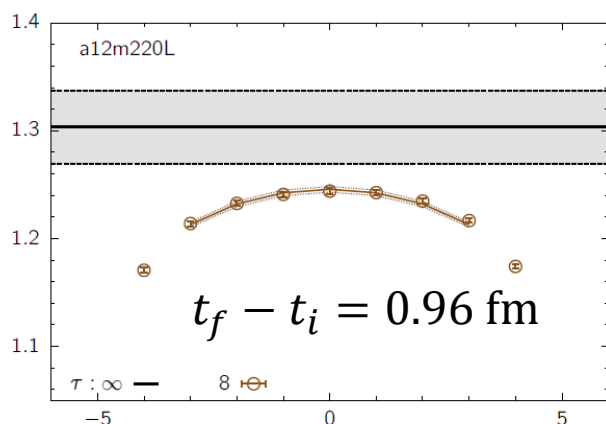


$$\begin{aligned}
 C^{3\text{pt}}(t_f, t, t_i) = & |\mathcal{A}_0|^2 \langle 0 | \mathcal{O}_\Gamma | 0 \rangle e^{-M_0(t_f - t_i)} \\
 & + \mathcal{A}_0 \mathcal{A}_1^* \langle 0 | \mathcal{O}_\Gamma | 1 \rangle e^{-M_0(t - t_i)} e^{-M_1(t_f - t)} + \mathcal{A}_0^* \mathcal{A}_1 \langle 1 | \mathcal{O}_\Gamma | 0 \rangle e^{-M_1(t - t_i)} e^{-M_0(t_f - t)} \\
 & + |\mathcal{A}_1|^2 \langle 1 | \mathcal{O}_\Gamma | 1 \rangle e^{-M_1(t_f - t_i)} \\
 C^{2\text{pt}}(t_f, t_i) = & |\mathcal{A}_0|^2 e^{-M_0(t_f - t_i)} + |\mathcal{A}_1|^2 e^{-M_1(t_f - t_i)} + \dots
 \end{aligned}$$

# No More Plateau Fit!

§ If one can get to the right answer sooner, why not?

- ∞ Errors are large at small  $t_{\text{sep}}$ ; errors are bigger than plateau fit
- ∞ They are consistent
- ∞ More  $t_{\text{sep}}$  data helps to further reduce the errors



$$\begin{aligned}
 C^{3\text{pt}}(t_f, t, t_i) &= |\mathcal{A}_0|^2 \langle 0 | \mathcal{O}_\Gamma | 0 \rangle e^{-M_0(t_f - t_i)} \\
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 \end{aligned}$$

# *Myth*

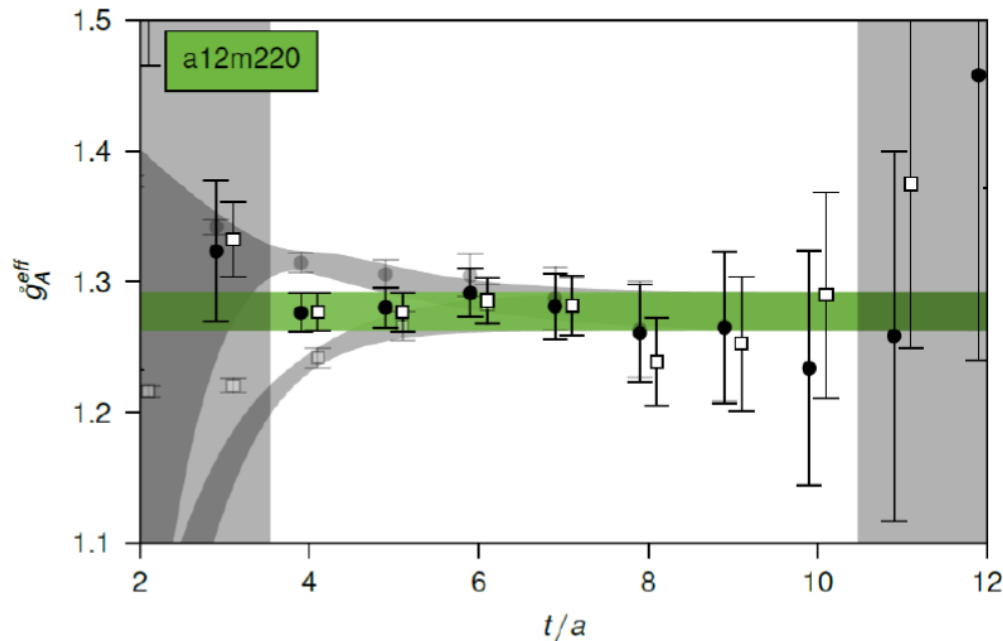
The two-state ground-state extraction  
will be dominated by the smallest  
source-sink separation



# No More Plateau Fit!

§ When the smallest source-sink separation dominates the ground-state errors

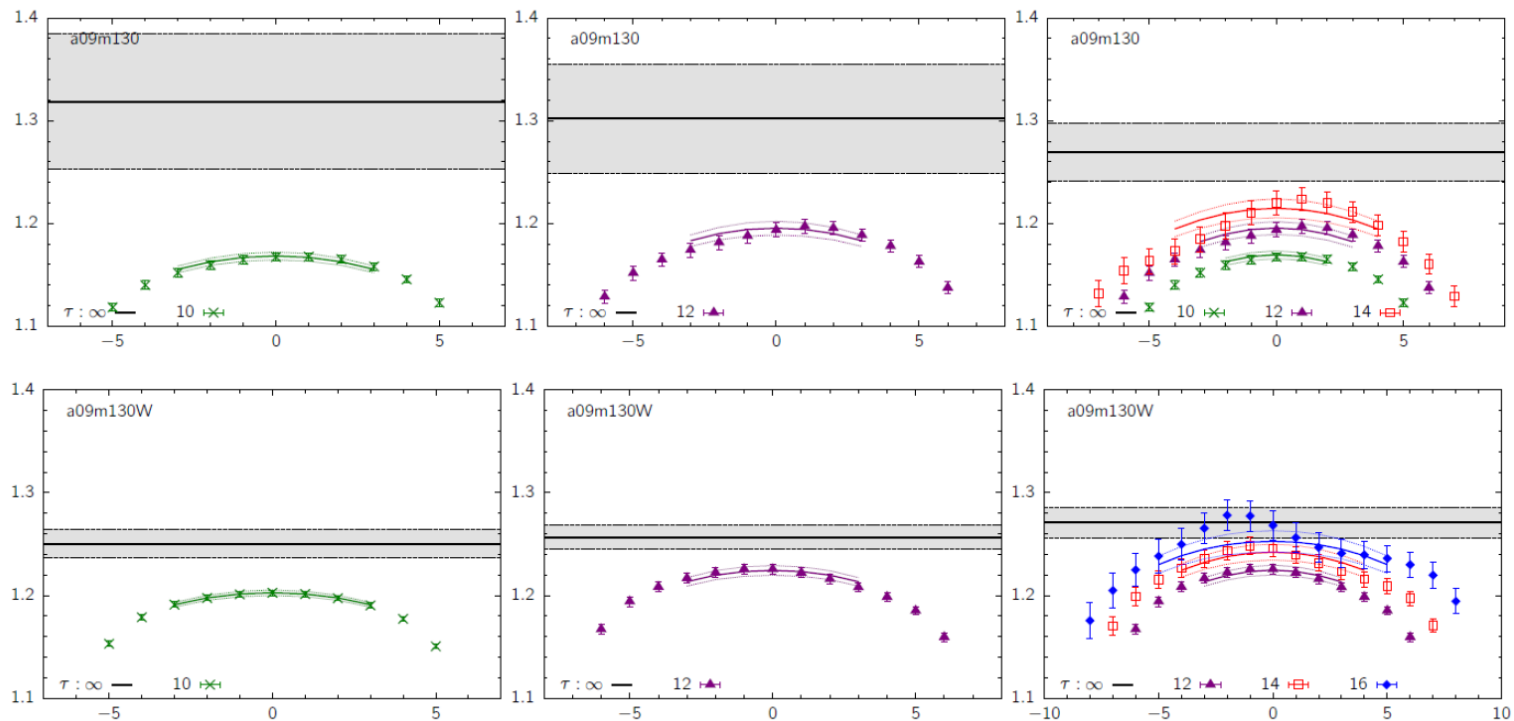
CalLat, *Nature* 558 (2018) no.7708, 91-94



# No More Plateau Fit!

§ When the smallest source-sink separation DO NOT dominate the errors

PNDME, Phys. Rev. D98 (2018) 034



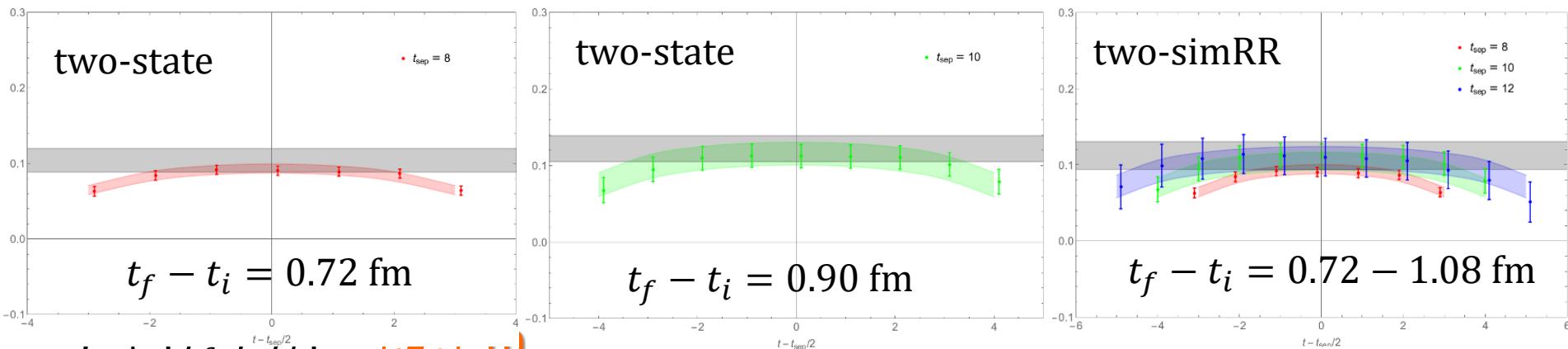


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$P_z = 2.6 \text{ GeV}$ ,  $z = 4$ , real (plot by Zhouyou Fan)



$$C^2\text{pt}(t_f, t_i) = |\mathcal{A}_0|^2 e^{-M_0(t_f - t_i)} + |\mathcal{A}_1|^2 e^{-M_1(t_f - t_i)} + \dots$$

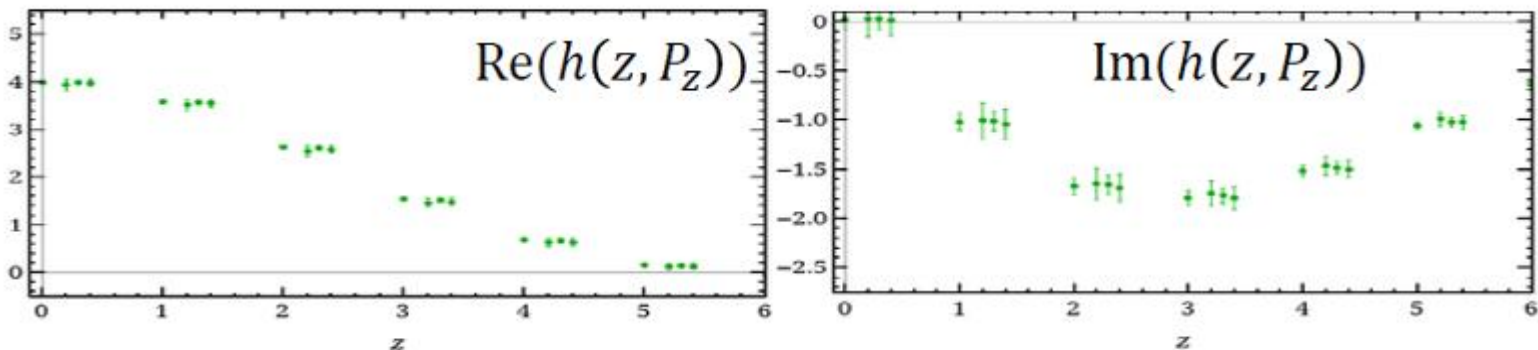
$$+ \mathcal{A}_0 \mathcal{A}_1^* \langle 0 | \mathcal{O}_\Gamma | 1 \rangle e^{-M_0(t - t_i)} e^{-M_1(t_f - t)} + \mathcal{A}_0^* \mathcal{A}_1 \langle 1 | \mathcal{O}_\Gamma | 0 \rangle e^{-M_1(t - t_i)} e^{-M_0(t_f - t)} \\ + |\mathcal{A}_1|^2 \langle 1 | \mathcal{O}_\Gamma | 1 \rangle e^{-M_1(t_f - t_i)}$$

$$C^2\text{pt}(t_f, t_i) = |\mathcal{A}_0|^2 e^{-M_0(t_f - t_i)} + |\mathcal{A}_1|^2 e^{-M_1(t_f - t_i)} + \dots$$

# We Do Not Use Plateau Fit

1) Calculate nucleon matrix elements on the lattice

$$h(z, P_z) = \left\langle P \left| \bar{\psi}(z) \gamma_z \exp\left(-i g \int_0^z dz' A_z(z')\right) \psi(0) \right| P \right\rangle$$



$$P_z = 2.6 \text{ GeV}$$

$$M_\pi \approx 135 \text{ MeV}, a \approx 0.09 \text{ fm}$$

LP<sup>3</sup> 1804.01483



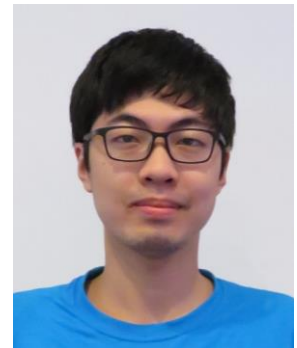
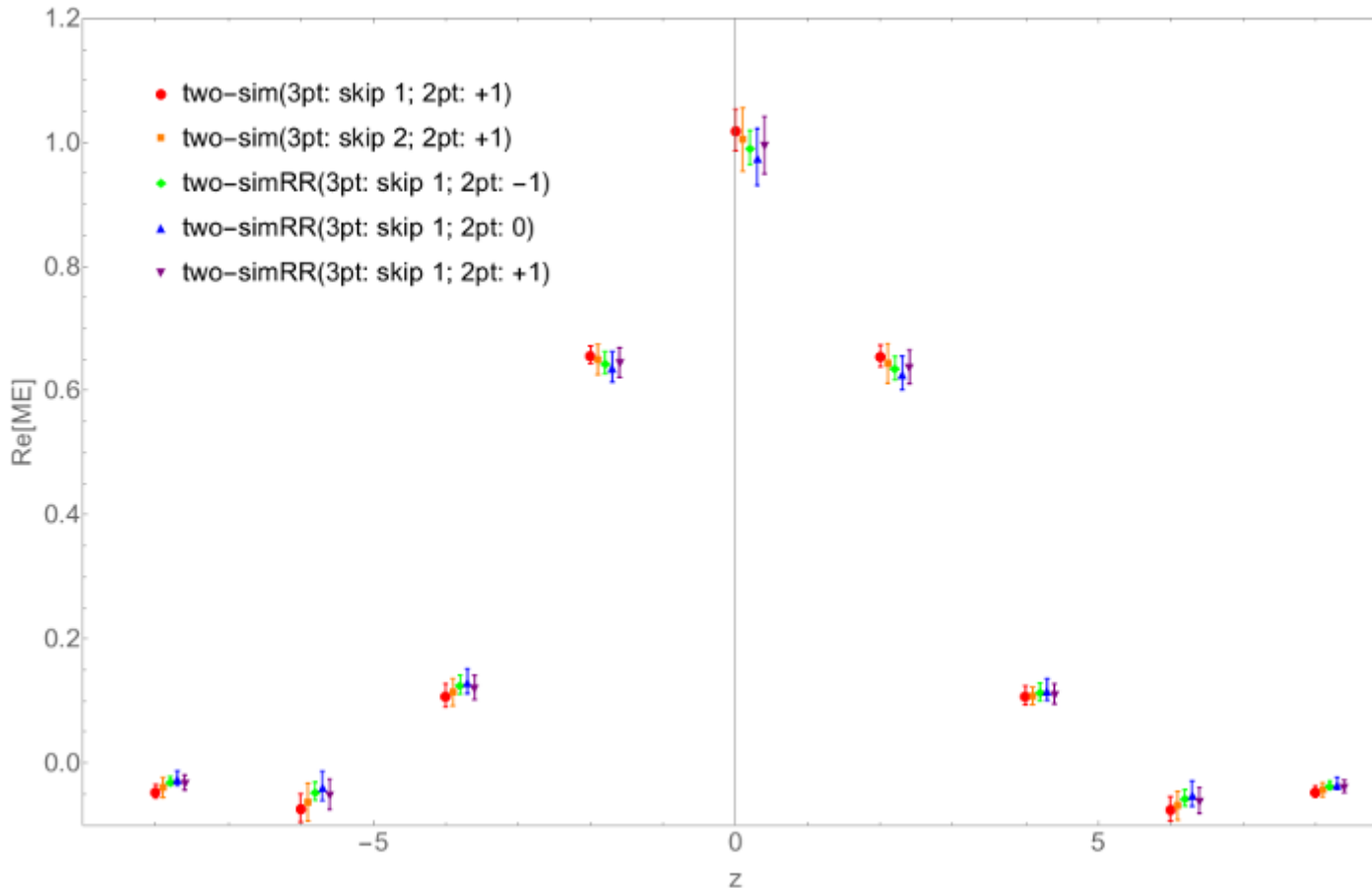
Ruizi Li

Blinded 3-state fits  
produced consistent  
results

HL@INT2018, DNP2018,...

# *We Do Not Use Plateau Fit*

$P_z = 2.6$  GeV, transversity (plot by Zhouyou Fan)



Zhouyou Fan

# No More Plateau Fit!

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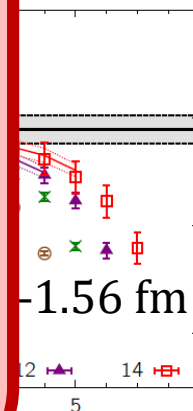
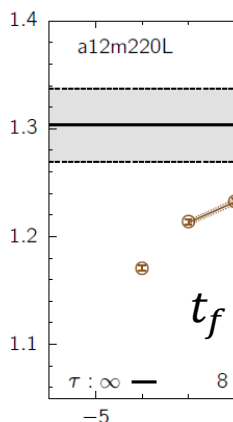
∞ More

Don't like multiple-state fits?

Use another method:

summation, GEV, or invent your own to  
take account into excited states

Anything but plateau fit!



$$C^{3\text{pt}}(t_f, t, t_i) = |\mathcal{A}_0|^2 \langle 0 | \mathcal{O}_\Gamma | 0 \rangle e^{-M_0(t-t_i)} e^{-M_0(t_f-t)} + \mathcal{A}_0 \mathcal{A}_1^* \langle 0 | \mathcal{O}_\Gamma | 1 \rangle e^{-M_0(t-t_i)} e^{-M_1(t_f-t)} + \mathcal{A}_0^* \mathcal{A}_1 \langle 1 | \mathcal{O}_\Gamma | 0 \rangle e^{-M_1(t-t_i)} e^{-M_0(t_f-t)} + |\mathcal{A}_1|^2 \langle 1 | \mathcal{O}_\Gamma | 1 \rangle e^{-M_1(t_f-t_i)}$$

$$C^{2\text{pt}}(t_f, t_i) = |\mathcal{A}_0|^2 e^{-M_0(t_f-t_i)} + |\mathcal{A}_1|^2 e^{-M_1(t_f-t_i)} + \dots$$

# *No More Plateau Fit!*

§ Avoid getting signal from noisy areas!

↻ Many people are doing this!

↻ FIND EVAN'S SLIDE AND SOME NPLQCD



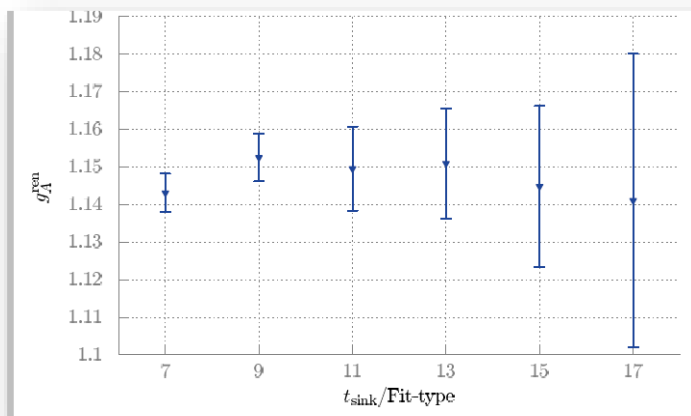
# Myth

One needs 1 fm source-sink separation  
to control the excited state  
for nucleon matrix elements



# Smearing Dependent

## § Even with the plateau method



### ① Plateau Method: single-state

⇐ RQCD (2014):

[G.Bali et al. (RQCD), 2014]

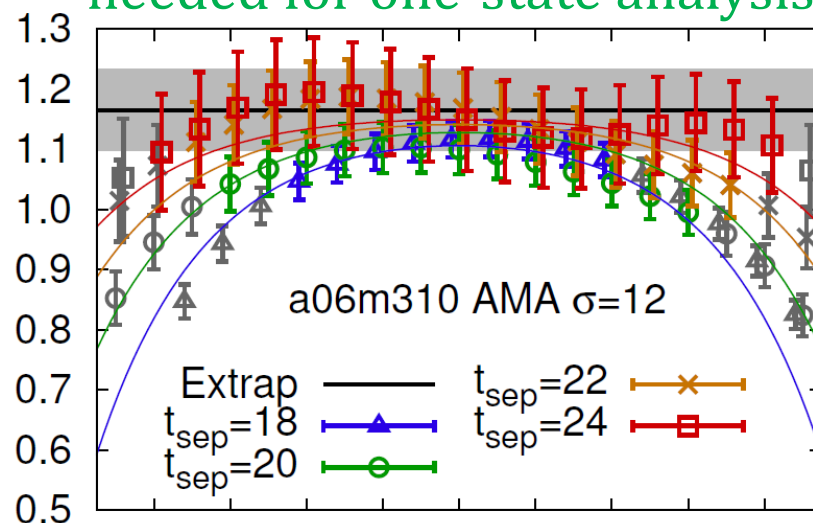
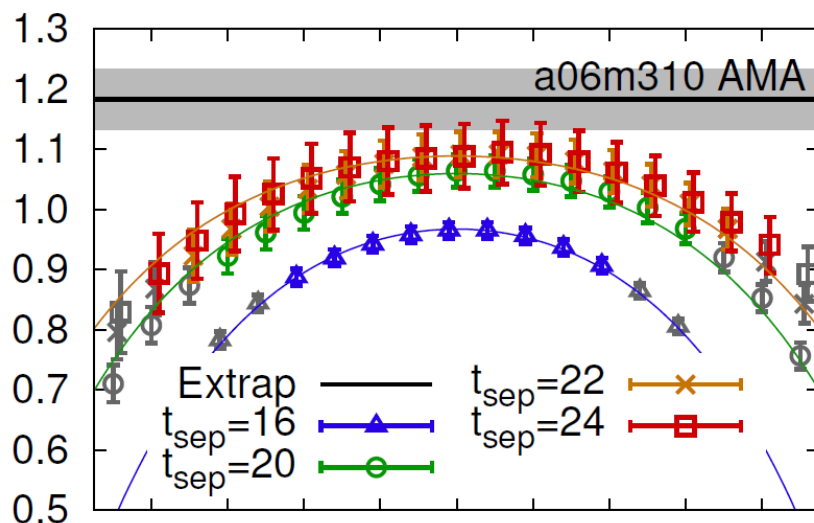
- ▶  $m_\pi = 285 \text{ MeV}$
- ▶  $g_A$  not sensitive on  $T_{\text{sink}}$ : 0.49-1.19 fm

M. Constantinou @ Lattice 2014

# Smearing Dependent

§ An example study on excited-state from PNDME work (42K)

➤ Robustness of the 2-state fit  $a = 0.06 \text{ fm}$ , 220-MeV pion  $1606.07049$   
Small smearing param.  $\sigma_G=6.5$   $g_A^{\text{bare}}$  Bigger smearing param.  $\sigma_G=10$   
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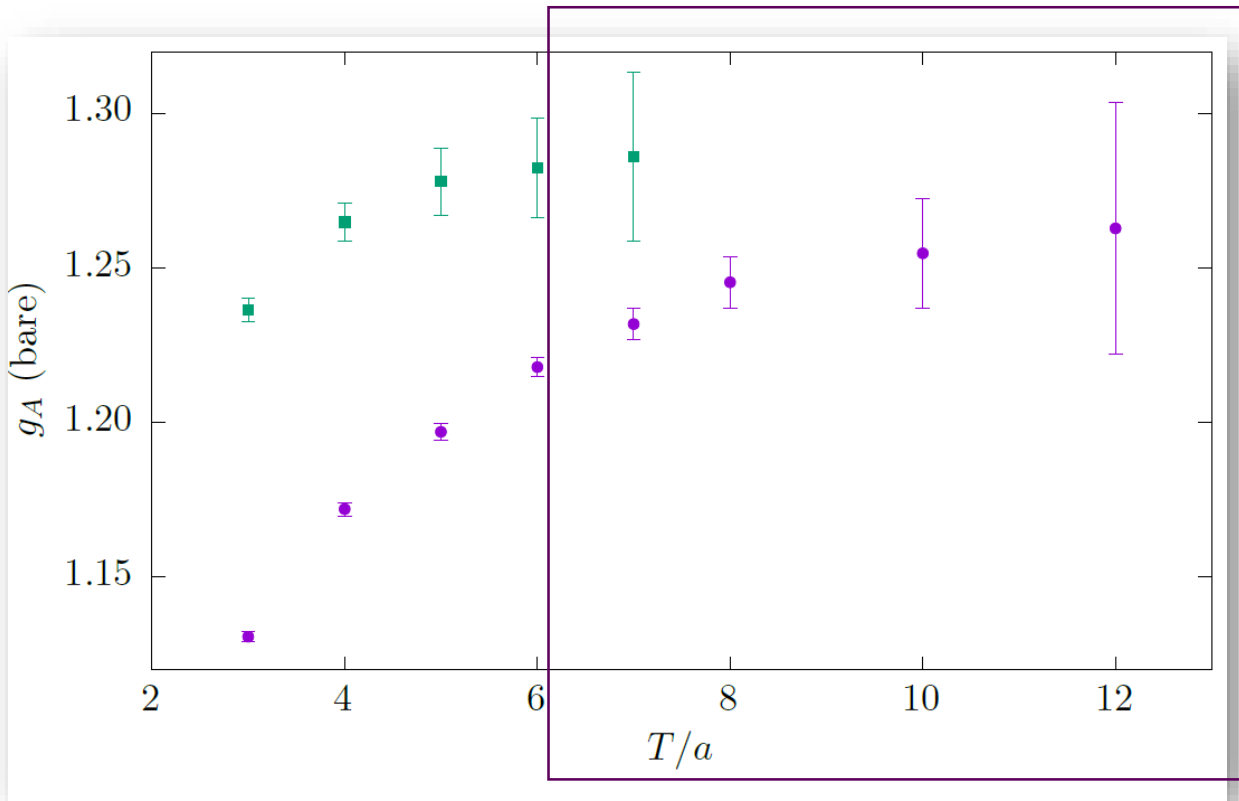
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# Analysis Method Dependent

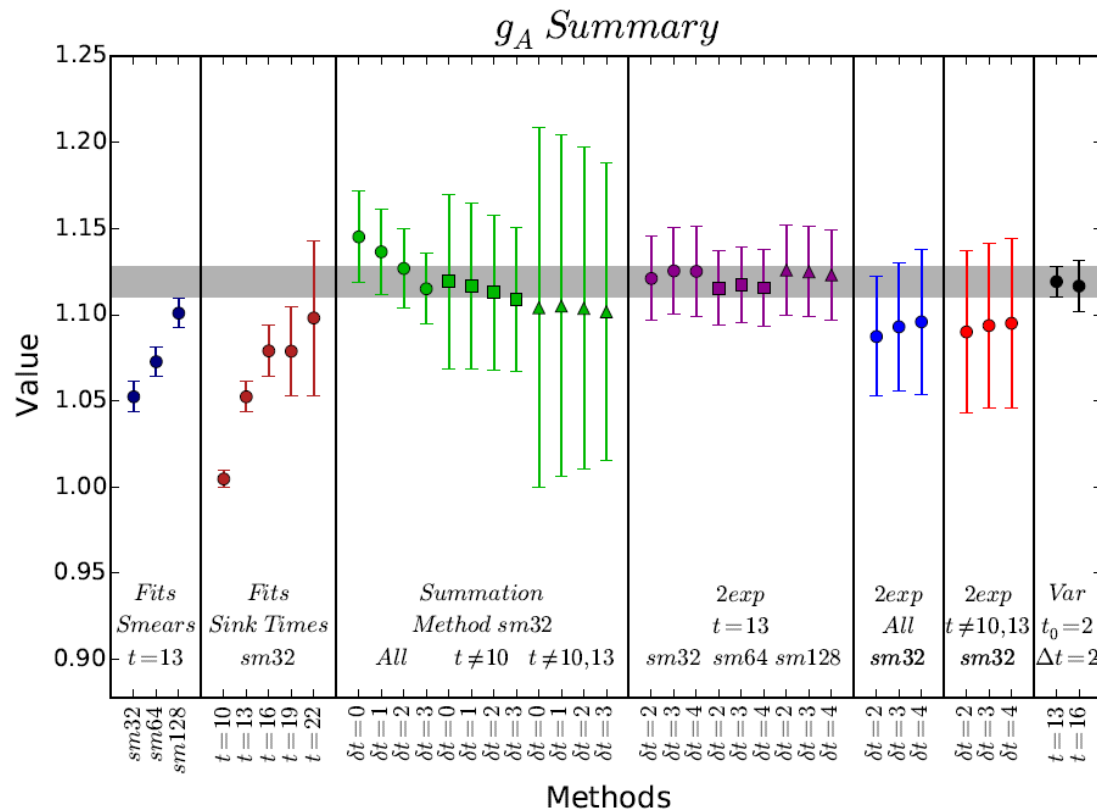
## § Summation method

≈ 0.116 fm at physical pion mass

1812.10574 J. Green



# Analysis Method Dependent



J. Green@Lattice 2018

$$m_\pi = 460 \text{ MeV}, a = 0.074 \text{ fm}$$

J. Dragos *et al.*, PRD **94**, 074505 (2016) [1606.03195]



# *Myth*

You need way too much computer to reach high boost momentum



# The Case of 3-GeV $P_z$

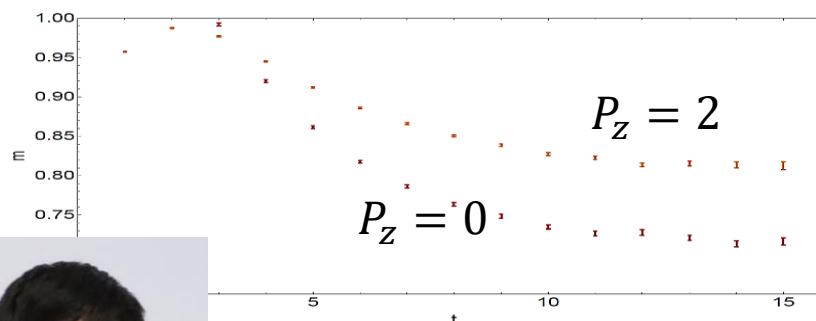
§ The signal-to-noise ratio depends on many parameters

§ Be careful what one extrapolates

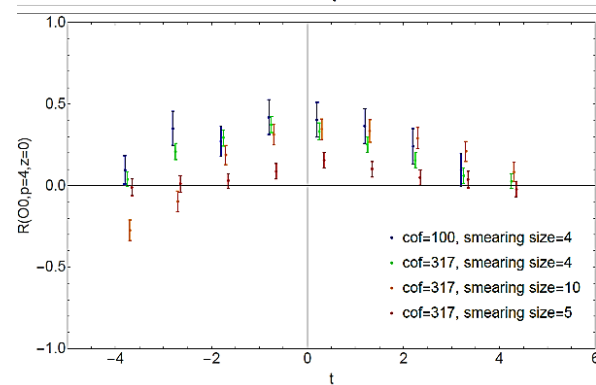
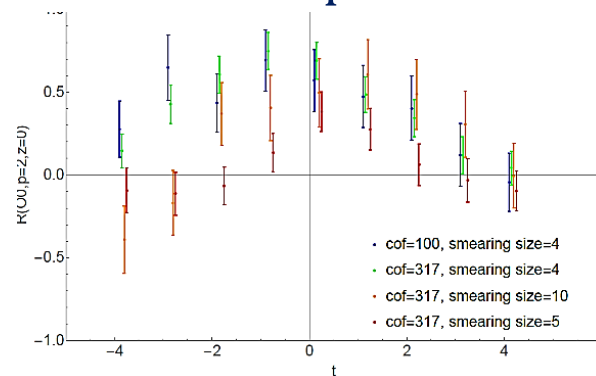
∞ Here is an example from my student

∞ If one simply extrapolates from this data, the conclusion:

∞ One needs more statistics to get zero-momentum proton mass and moments; obviously NOT true



Zhouyou Fan

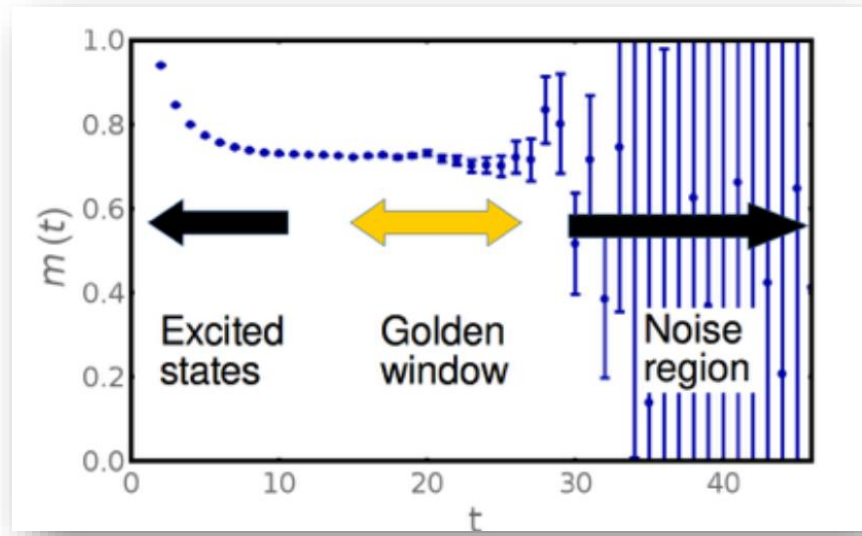


# The Case of 3-GeV $P_z$

## § Our original target was not 3 GeV

- ⌘ We were hoping to get *some* signal at 2.6 GeV with momentum smearing fine-tuned to this point
- ⌘ We thought that we would have 2.2 GeV to fall back, since it's not far from 2.6 GeV
- ⌘ 3.0 GeV was a surprise to me too

No one thought NPLQCD could calculate  $NN$  interactions when they first started, but now they are the main force in lattice nuclear physics

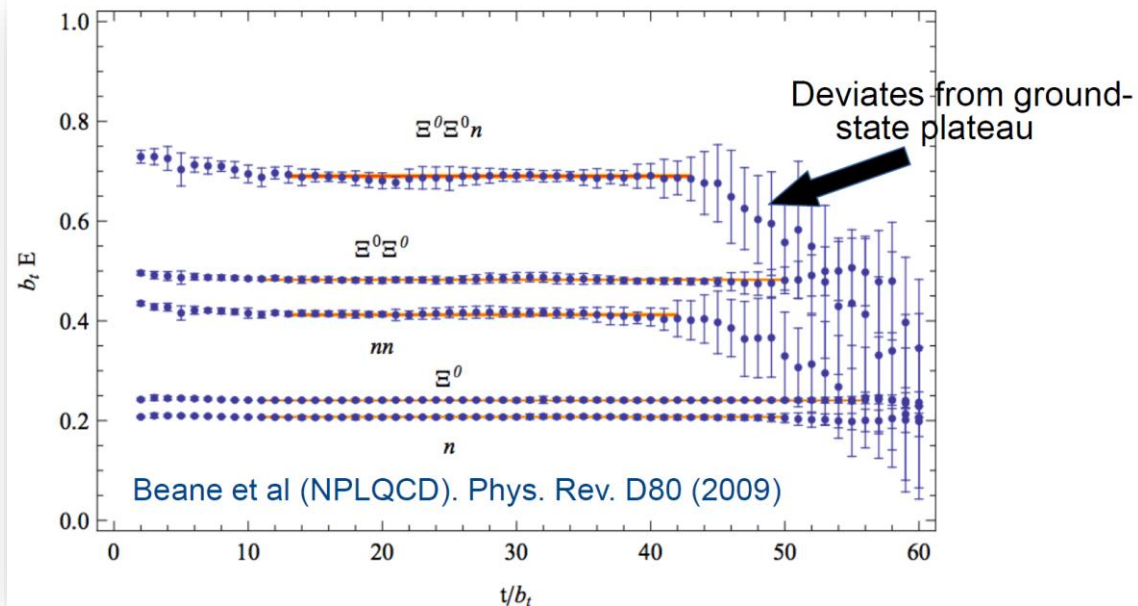


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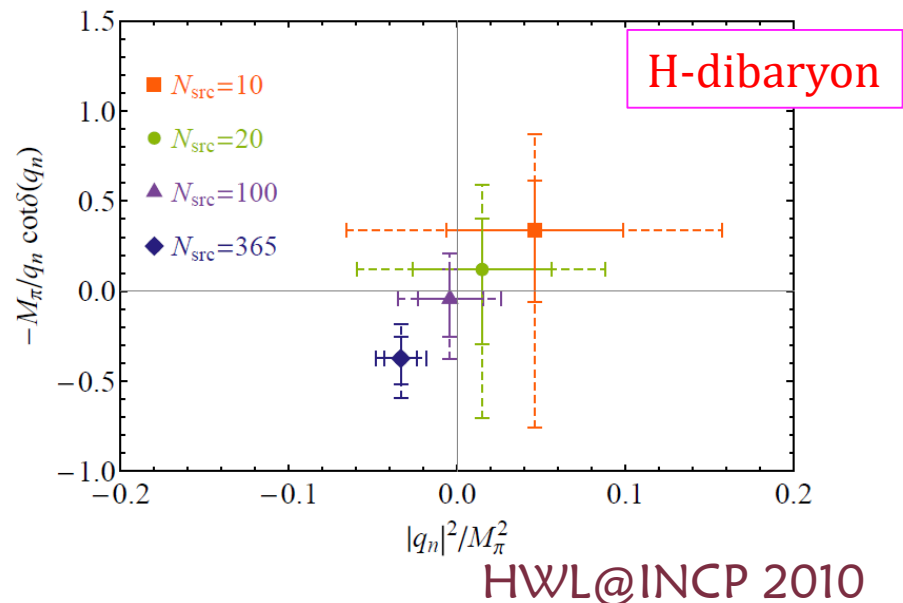


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Words of advice:  
Similar to hadron-  
interaction calculations,  
one can easily reach  
wrong conclusions with  
insufficient statistics





# *Excited States at 3-GeV $P_z$*

§ How about excited states?

§ We do what we can on our ensembles

↻ If only we were not bounded by finite computer resources

§ Future direction in progress

↻ Smaller lattice spacing

↻ Variational method to catch more excited states

↻ Happy to hear more suggestions

↻ (...unless they involve plateau fits)

# Conclusions

Apologize in advance for a rather technical talk

Need to clear up many false rumors

§ We still need large boost momentum

∞ The most non-model dependent way to access  
small-x physics for EIC

∞ Let's work together to  
figure out the best way to do it

§ Please stop using plateau fit

∞ Just stop



# *Further Criticism?*



# *Backup Slides*



# Control Excited States

§ Same conclusion from another study on excited-state from NME (Jlab,PNDME,LHPC) work  $a = 0.08 \text{ fm}$ , 312-MeV pion

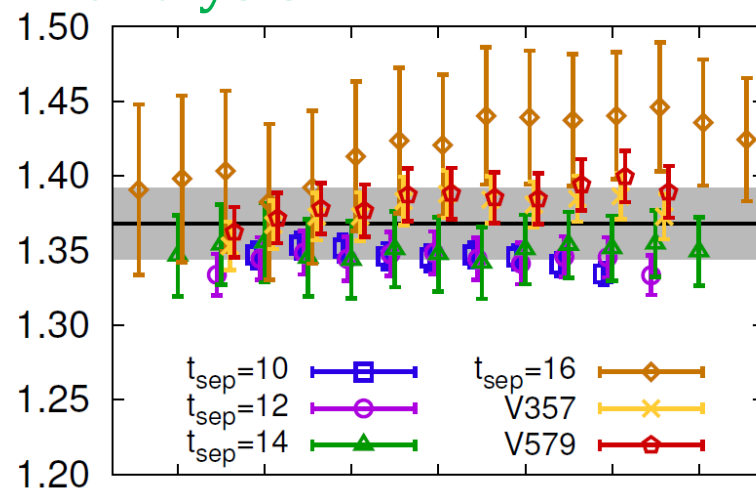
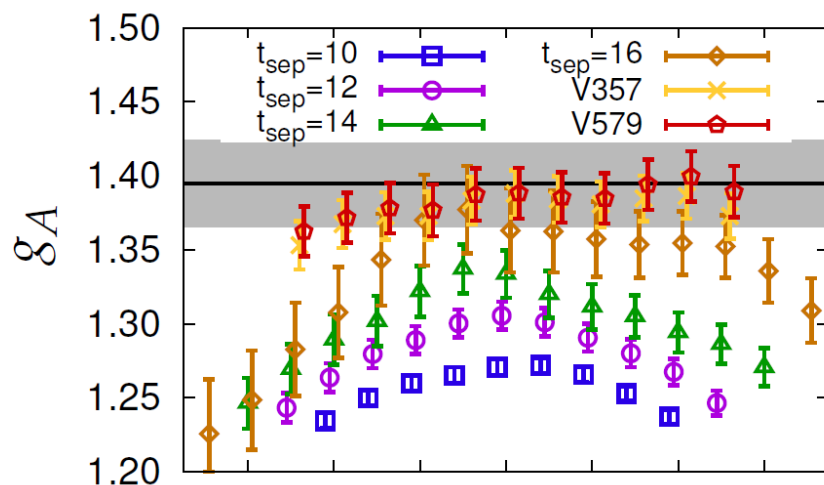
1602.07737

∞ Robustness of the 2-state fit; consistent with vibrational method

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§ V357 and V579 are results from different 3x3 correlator matrix analysis

§ Ground-state stat. error is NOT dominated by smallest tsep

# Many Such Examples

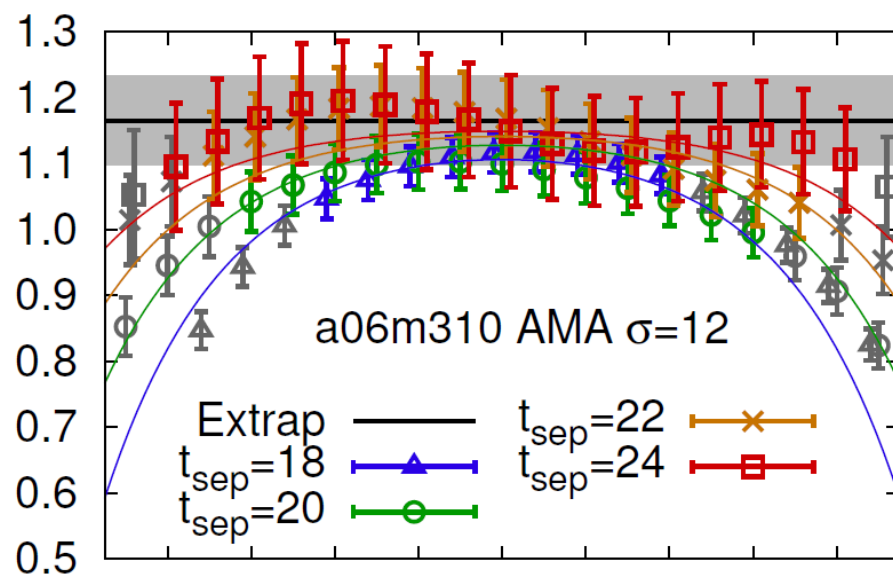
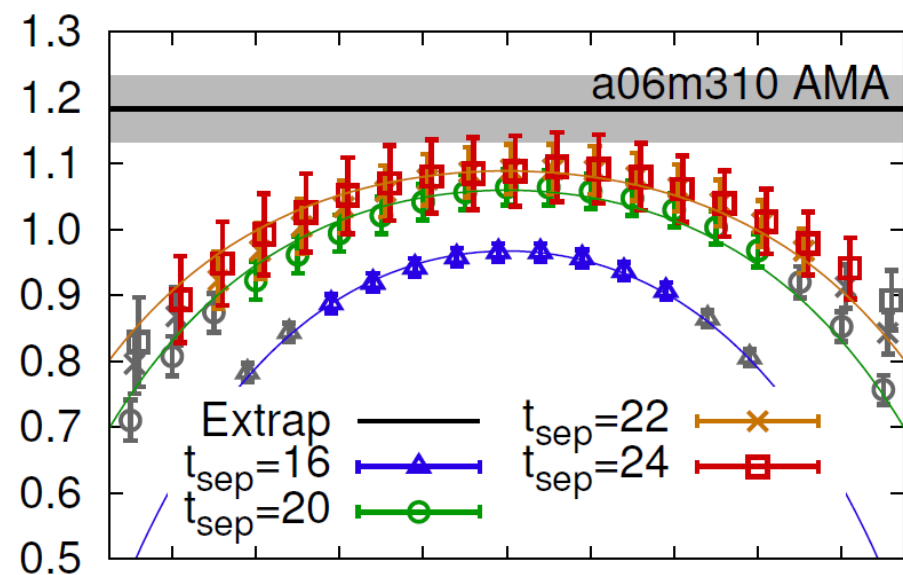
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1606.07049, Plots by Boram Yoon

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§ Very small gain by focusing on analysis with only ground state



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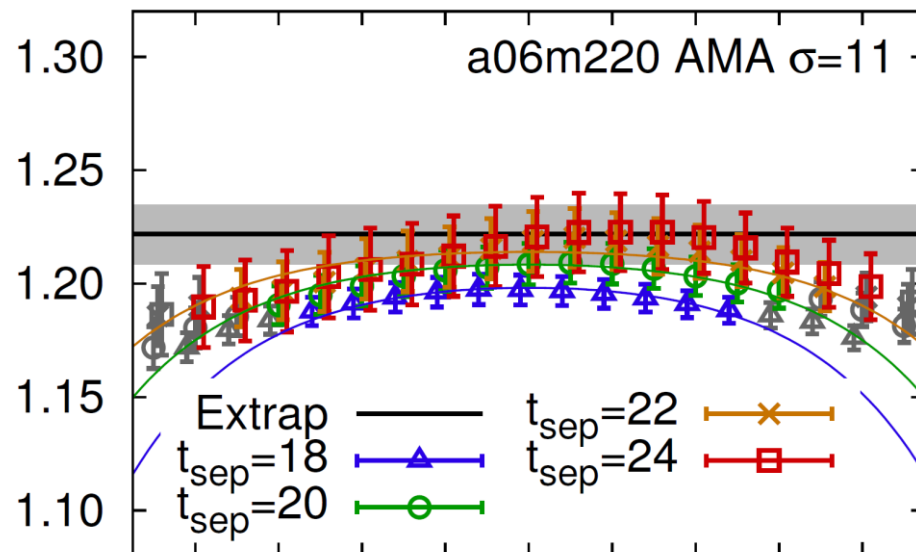
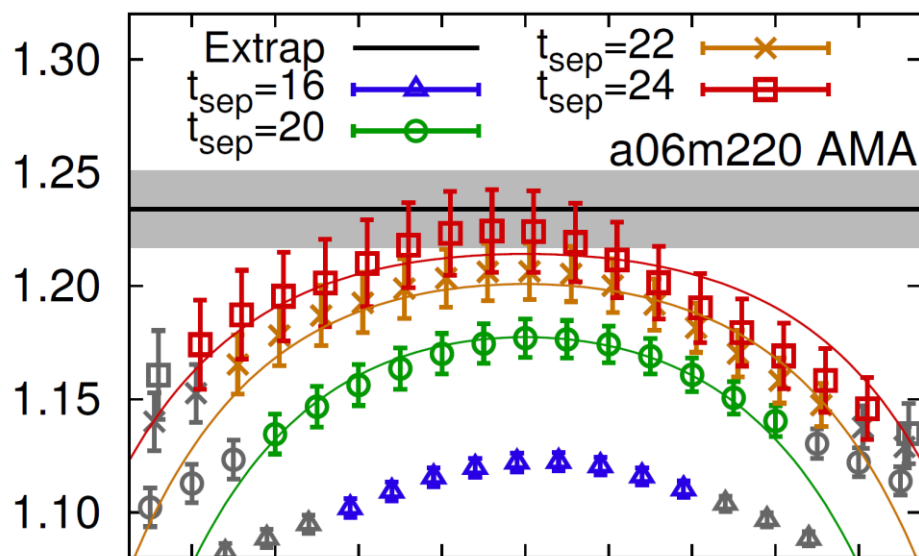
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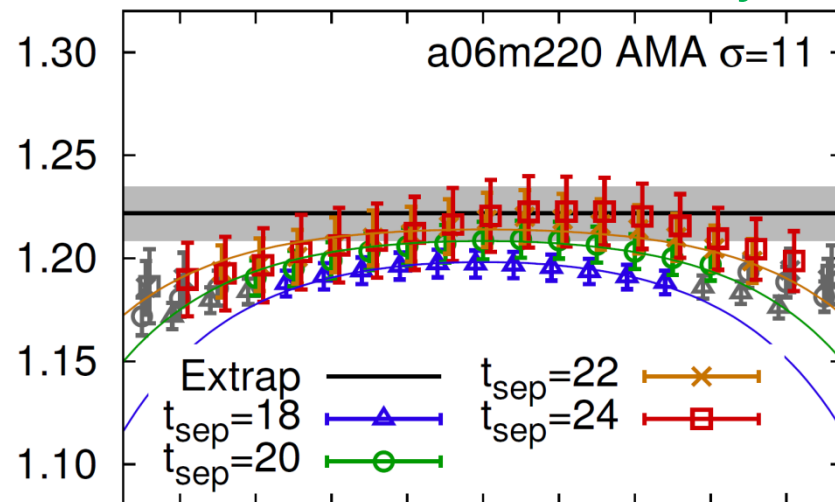
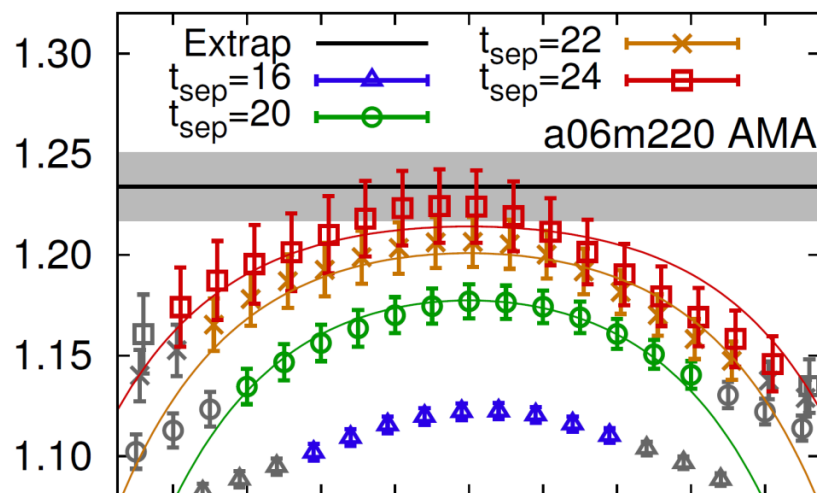
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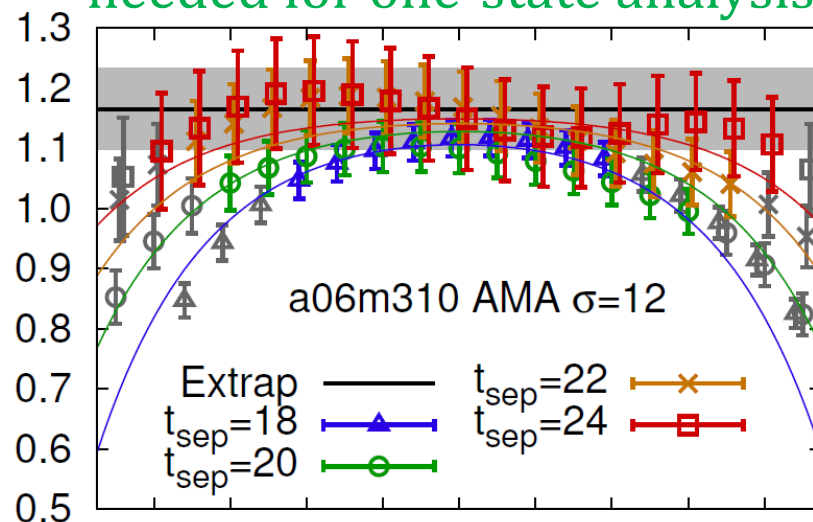
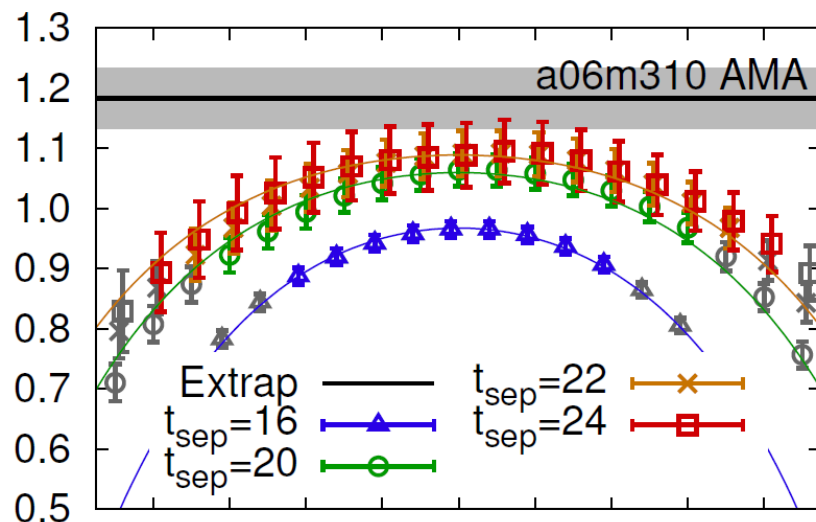
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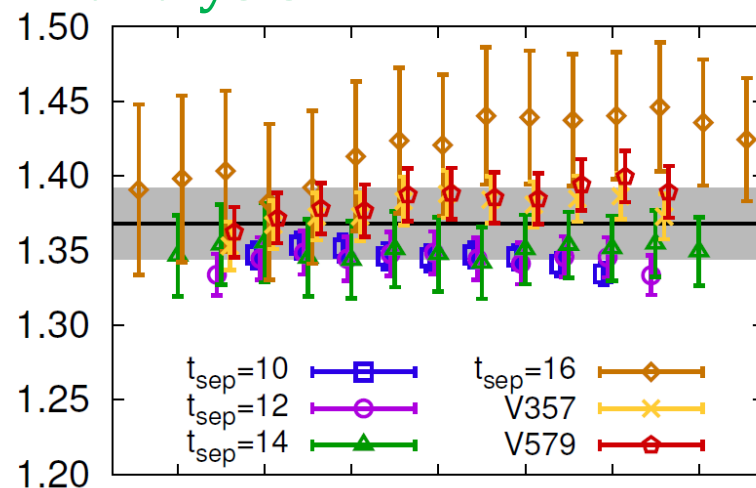
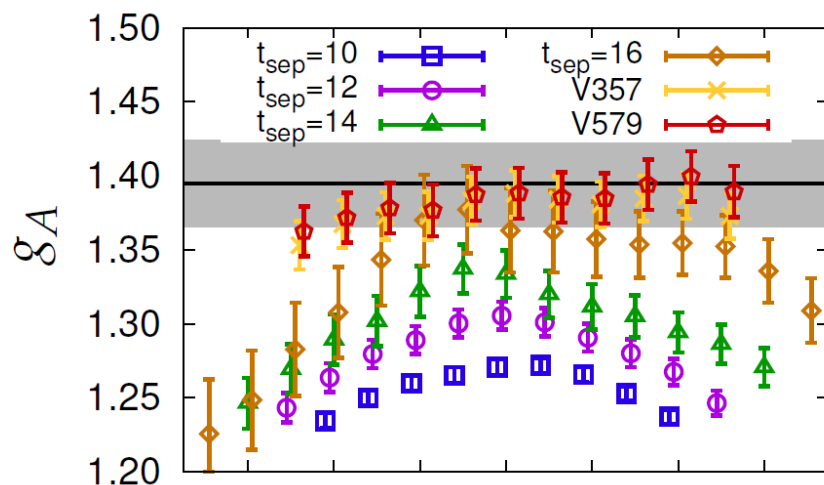
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