

# Valence Pion PDF

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## Definitions

quark PDF

$$f_{\psi/h}(x, \mu^2) = \int \frac{d\xi^-}{4\pi} e^{-ixP^+z^-} \langle h(P) | [\bar{\psi}(\xi^-) \gamma^+ W_L(\xi^-, 0) \psi(0)]_{\mu} | h(P) \rangle$$

quark qPDF [arXiv:1305.1539]

$$q_{\psi/h}(x, P_z, \mu_R^2) = \int \frac{dz}{4\pi} e^{-ixP_z z} \langle h(P) | [\bar{\psi}(z) \Gamma W_L(z, 0) \psi(0)]_{\mu_R} | h(P) \rangle$$

$$\Gamma \in \{\gamma_t, \gamma_z\}$$

qPDF  $\rightarrow$  PDF [arXiv:1709.04933, arXiv:1706.01295]

$$q(x, P^z, \mu_R^2) = \int \frac{dy}{|y|} C\left(\frac{x}{y}, \frac{\mu_R}{P^z}, \frac{\mu}{|y|P^z}\right) f(y, \mu) + \mathcal{O}\left(\frac{m_h^2}{P_z^2}, \frac{\Lambda_{QCD}^2}{P_z^2}\right)$$

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$$\mathcal{O}_{\Gamma}^{\psi}(z, 0) = \psi(z) \Gamma W_L(z, 0) \psi(0)$$

# What it is we compute

## Isovector quark PDF

$$f_{\pi^+}^{u-d}(x, \mu) = f_{\pi^+}^u(x, \mu) - f_{\pi^+}^d(x, \mu)$$

Valence quark PDF  $x \in [0, 1]$

Matching quasi-quark PDF to PDF [arXiv:1709.04933, arXiv:1706.01295]

$$h_{u-d/\pi^+}(x, P^z, \mu_R) = f_{u-d/\pi^+}(x) + \frac{\alpha_s C_F}{2\pi} \int_{-1}^1 dy \left[ \frac{1}{|y|} \mathcal{F}_1\left(\frac{x}{y}\right)_+ + |\eta'| \mathcal{F}_2(1 + \eta'(x - y))_+ \right] f_{u-d/\pi^+}(y)$$

# Work Flow

## (1) Remove Excited States

$$\langle 0 | \pi_{P^z; \Delta t}^+ \mathcal{O}_\Gamma^{u-d}(z, 0; \tau) \bar{\pi}_{P^z; 0}^+ | 0 \rangle_a \rightarrow \langle \pi^+(P^z) | \mathcal{O}_\Gamma^{u-d}(z, 0) | \pi^+(P^z) \rangle_a$$

## (2) Renormalize qPDF

$$\langle \pi^+(P^z) | \mathcal{O}_\Gamma^{u-d}(z, 0) | \pi^+(P^z) \rangle_a \rightarrow \langle \pi^+(P^z) | \mathcal{O}_\Gamma^{u-d}(z, 0) | \pi^+(P^z) \rangle_{\mu_R}$$

## (3) qPDF $\rightarrow$ PDF

$$\langle \pi^+(P^z) | \mathcal{O}_\Gamma^{u-d}(z, 0) | \pi^+(P^z) \rangle_{\mu_R} \rightarrow f(y, \mu) \text{ via matching formula}$$

# Lattice Parameters & Computation Details

- Mixed Action
  - ▶ Sea: HotQCD HISQ gauge ensemble [arXiv:1407.6387]
  - ▶ Valence: Wilson Clover with 1HYP
- Lattice Spacing: 0.06 fm
- Lattice Size:  $48^3 \times 64$
- Pion Mass: 300 MeV
- Configuration Used:
  - ▶  $N_{CFG}$ : 100  $|z| > 16$
  - ▶  $N_{CFG}$ : 162  $16 \geq |z| > 8$
  - ▶  $N_{CFG}$ : 214  $8 \geq |z| > 0$
- Correlators per config computed using All-Mode Averaging (AMA) [arXiv:1402.0244]
  - ▶ 1 exact sample
  - ▶ 32 sloppy samples

# (1) Remove Excited States

## 2-Point Function

$$C_{2pt}(t) = \sum_n |A_n| e^{-E_n t}$$

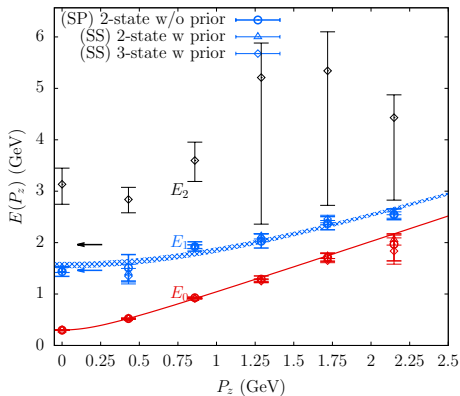
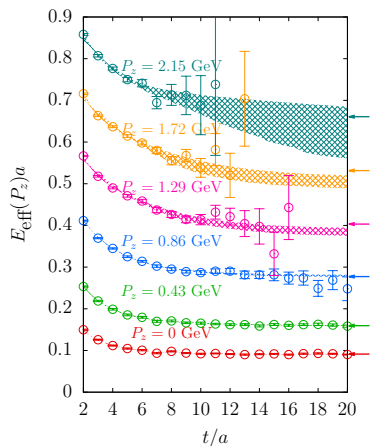
Fit two-point function to sum of exponentials to acquire parameters

## 3-Point Function

$$C_{3pt}^{\vec{P}, \Gamma}(\Delta t, \tau; z) = \sum_{n,m} e^{-E_n \Delta t} e^{E_n \tau} e^{-E_m \tau} \times A_n^* \langle n | \mathcal{O}_\Gamma^{u-d}(z, 0) | m \rangle A_m$$

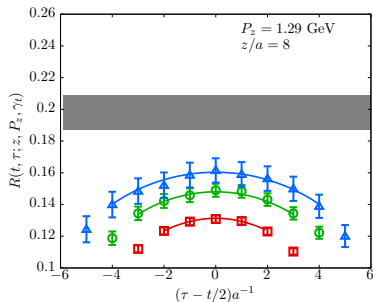
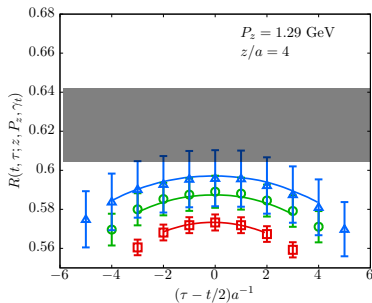
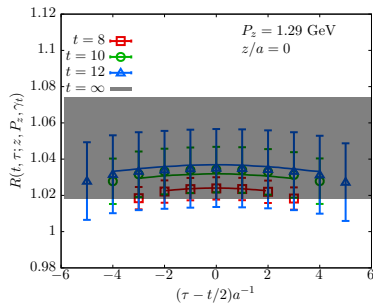
- Define  $R(\Delta t, \tau, z) = C_{3pt}(\Delta t, \tau, z) / C_{2pt}(\Delta t)$   
 $R_{sum}(\Delta t, z) = \sum_{\tau=t_0}^{\Delta t-t_0} R(\Delta t, \tau, z) = C + (\Delta t - t_0)\mathcal{M}$
- Fit  $C_{3pt}$  across all  $\Delta t$  and  $\tau$  to get  $\mathcal{M}$ , fixing  $A_n$  and  $E_n$  from 2-point fit

# Two-Point Fit Results

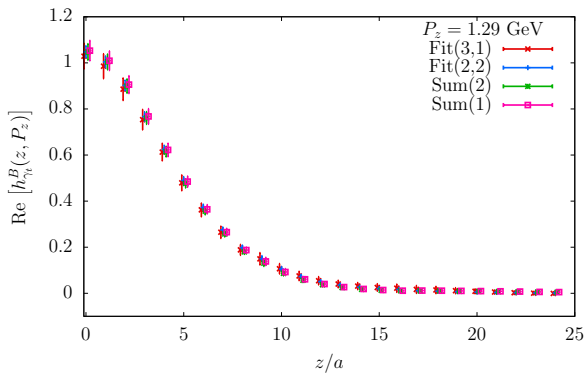




# Three-Point Fit Results



## Pion qPDF matrix element after extrapolation



## (2) qPDF Renormalization

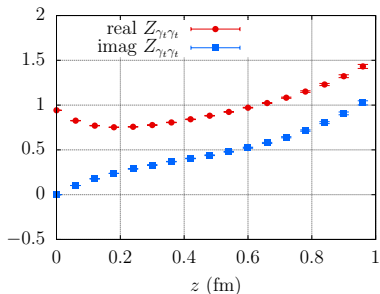
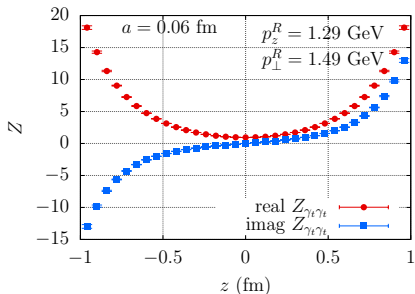
qPDF matrix element can be multiplicatively renormalized

$$h_{\gamma_t}^R = Z_{\gamma_t\gamma_t}(z; p^R) h_{\gamma_t}^o(z; P^z, a)$$

Renormalization coefficient renormalized under RI-MOM scheme

$$Z_{\gamma_t\gamma_t}(z, p^R) q_{\gamma_t}(z, p)|_{p=p^R} = e^{ip_z^R z}$$

where  $q$  is the quark-in-quark PDF



LHS: qPDF RI-MOM Z-factor

Divergent contribution  $\sim e^{c|z|}$

RHS: qPDF Z-factor after factoring out Wilson Line UV-divergent term

$c =$  number evaluated from heavy-quark effective potential [arXiv:1804.10600]

# Comparison to 1-Loop Perturbation Theory

How well does  $q(z, P^z, \mu_R^2)_{\text{lat}}$  match with  $q^{1\text{-Loop}}(z, P^z, \mu_R^2)$ ?

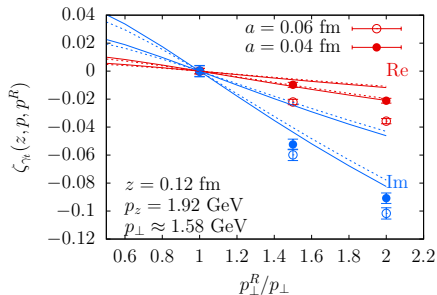
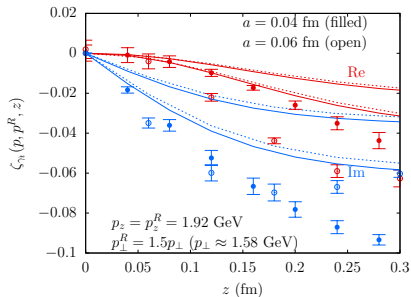
Study the behavior of

$$\zeta_{\Gamma}(z, p, p^R) = \frac{q^R(z, p, p^R)}{q^R(z, p, p)} - 1$$

In both 1-Loop Perturbation Theory and on the the Lattice.

No running of  $\alpha_s$  for  $\zeta$  in 1-Loop Perurbation theory.  
Compute  $\zeta^{1\text{-Loop}}$  using  $\alpha_s$  from scales  $\mu = \frac{1}{2}p^R$  to  $2p^R$

# $\zeta$ $z$ and $p^R$ dependence



# Identifying Perturbative/Nonperturbative Regime in Matrix Element

For what quark-antiquark separation do we expect to see nonperturbative physics to be prominent in real-space qPDF?

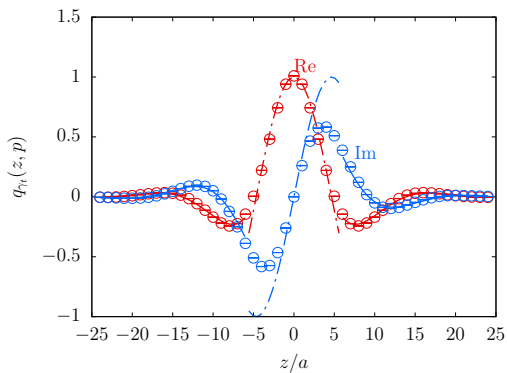
Model bare quark qPDF using the following ansatz:

$$q(z, P) = A e^{i\omega z} e^{-m_{\text{scr}}|z|} e^{-c|z|}$$

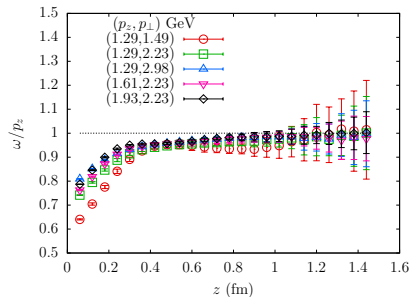
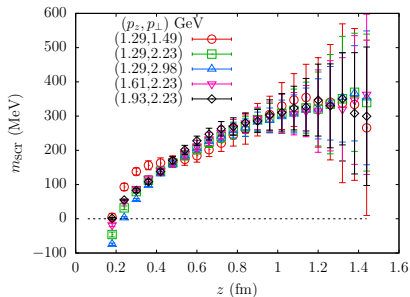
- $e^{i\omega z}$  represents tree-level bare quark-qPDF
- $e^{-c|z|}$  removes UV divergent piece of qPDF;  $c$  determined from heavy-quark potential calculation [arXiv:1804.10600]
- $e^{-m_{\text{scr}}|z|}$  models the damping of plane-wave quark via inverse screening length  $m_{\text{scr}} \sim \mathcal{O}(\Lambda_{\text{QCD}})$



How well does it fit the data?



# $m_{\text{scr}}$ and $\omega$

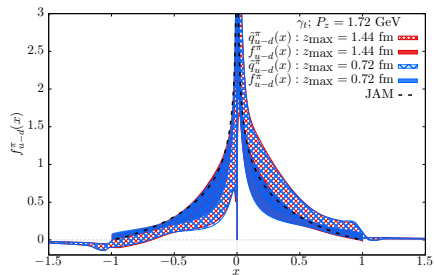
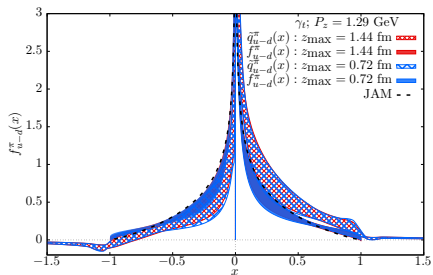


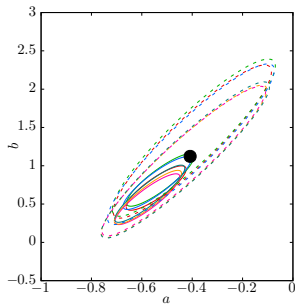
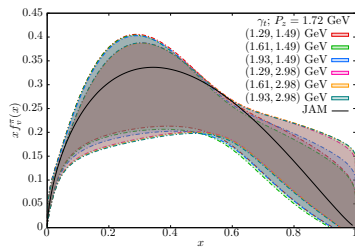
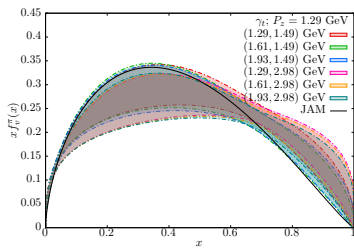
# (3) qPDF $\rightarrow$ PDF Matching

# Strategy

- A possible strategy. One could parameterize the real-space data  $h^R(z, P^z, p^R)$  over available  $z$  then model the behavior for larger  $z$ . Then perform Fourier Transform
  - ▶ Advantages: No model dependence when dealing with available data
  - ▶ Disadvantages: Must model real-space qPDF beyond available  $z$ -data. Performing Fourier Transform involves performing an integral in non-perturbative region of qPDF
- What we do. Use PDF ansatz  $f(x; a, b) = Ax^a(1-x)^b$ , transform the data into real-space quasiPDF, then fit that to data, minimizing.
  - ▶ Disadvantages: Introduce model dependence
  - ▶ Advantages: Model already used by phenomenologists fitting PDF's from dataCan control which  $z$ -data we can fit to

# Pion qPDF and PDF





# Conclusions

- We presented a calculation of the Pion valence-quark PDF using the quasi-PDF approach

$$m_\pi = 300 \text{ MeV}, a = 0.06 \text{ fm}, P^z = 1.29 \text{ GeV}, 1.72 \text{ GeV}$$

- We studied the running of the renormalization coefficient
  - ▶ Qualitative agreement between lattice renormalization and 1-loop
  - ▶ Smaller Lattice Spacing improves quantitative agreement
  - ▶ 2-Loops may be necessary
- More statistics are necessary to overcome noise in large-momentum data