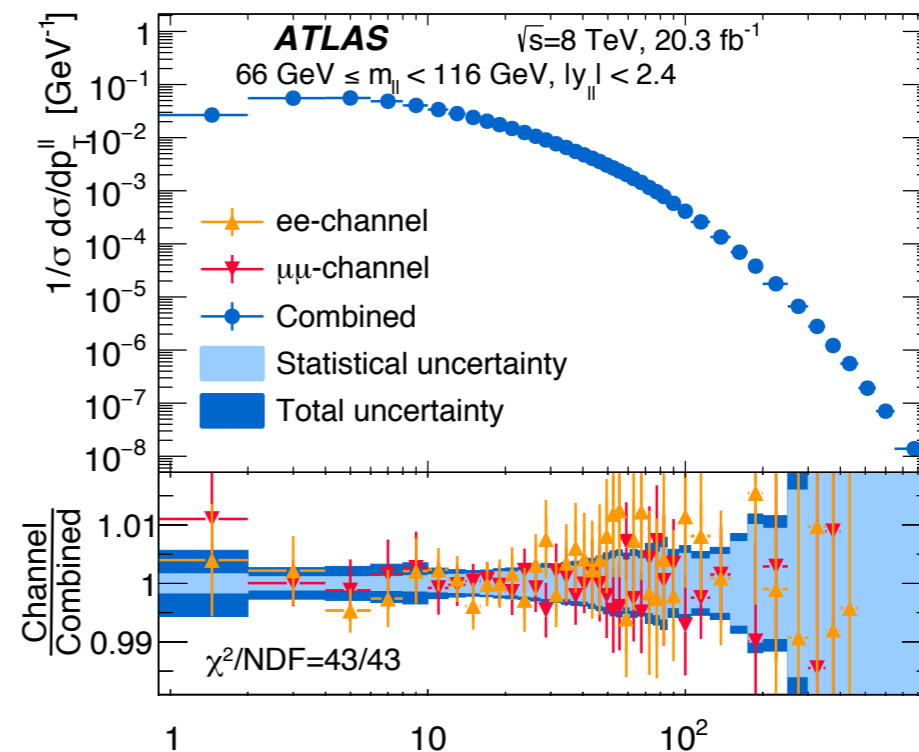




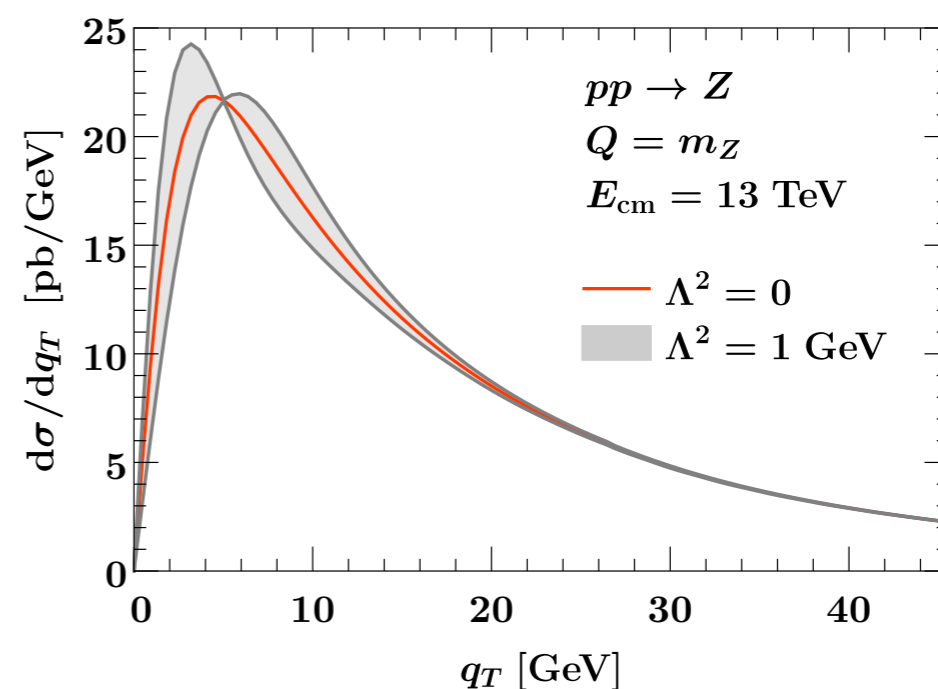
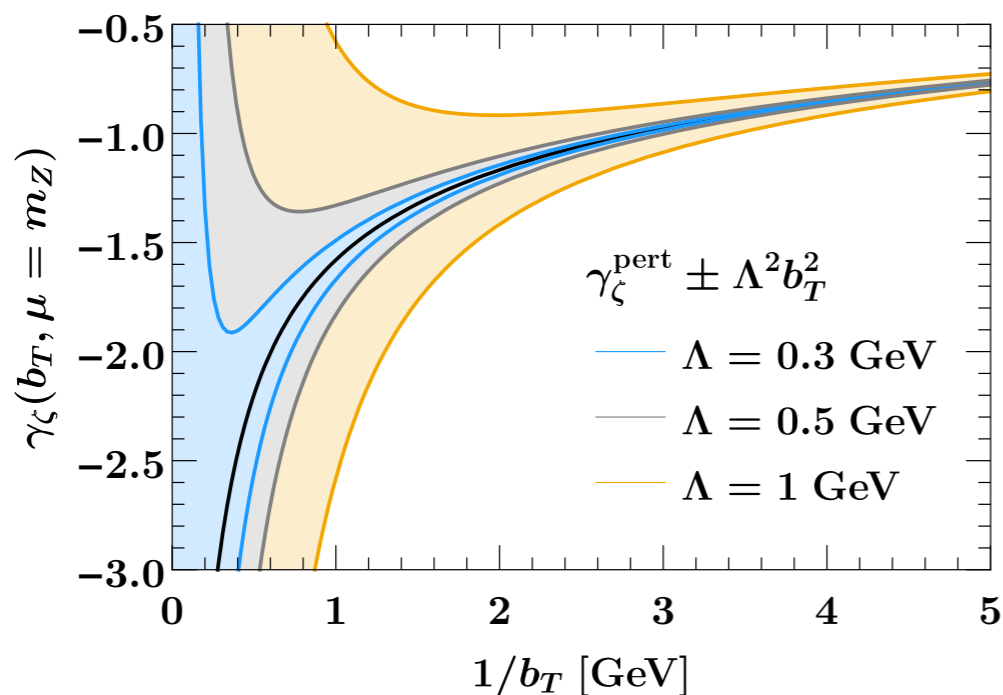
# TMD evolution

TMDPDFs encode 3D structure of hadrons,  $q_T$  dependence of hadronic cross-sections

Drell-Yan  $q_T$  dependence measured at LHC,  $<1\%$  precision



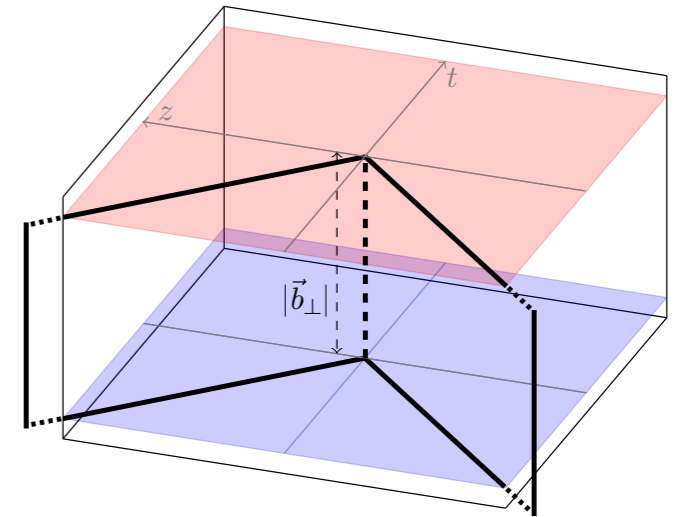
Nonperturbative contributions to rapidity anomalous dimension limit QCD theory predictions to much lower precision



# TMD evolution from LQCD

TMDPDFs inaccessible to LQCD with LaMET  
 — soft factor includes two light-light directions

Ebert, Stewart, Zhao, arXiv:1901.03685



Ratios of TMDPDFs free from soft factors, can be calculated with LQCD

Musch et al, PRD 85 (2012)

Engelhardt et al, PRD 93 (2016)

Yoon et al, PRD 96 (2017)

TMDPDF rapidity anomalous dimensions (Collins-Soper kernel)  
 calculable from ratios of quasi-TMDPDFs

Ebert, Stewart, Zhao, PRD 99 (2019)

$$\gamma_{\zeta}^{q, \overline{\text{MS}}}(b_T, \mu) = \zeta \frac{d}{d\zeta} f_q^{\overline{\text{MS}}}(x, b_T, \mu, \zeta)$$

$$= \frac{1}{\ln(p_1^z/p_2^z)} \ln \frac{C_{\text{TMD}}^{\overline{\text{MS}}}(\mu, xP_2^z) \int db^z e^{ib^z xp_1^z} \tilde{B}_q^{\overline{\text{MS}}}(b^z, b_T, \eta, \mu, p_1^z)}{C_{\text{TMD}}^{\overline{\text{MS}}}(\mu, xp_1^z) \int db^z e^{ib^z xp_2^z} \tilde{B}_q^{\overline{\text{MS}}}(b^z, b_T, \eta, \mu, p_2^z)}$$

LQCD-friendly quasi-beam function

# LQCD Setup

$$\gamma_{\zeta}^{q, \overline{\text{MS}}}(b_T, \mu) = \frac{1}{\ln(p_1^z/p_2^z)} \ln \frac{C_{\text{TMD}}^{\overline{\text{MS}}}(\mu, xP_2^z) \int db^z e^{ib^z xp_1^z} \tilde{B}_q^{\overline{\text{MS}}}(b^z, b_T, \eta, \mu, p_1^z)}{C_{\text{TMD}}^{\overline{\text{MS}}}(\mu, xp_1^z) \int db^z e^{ib^z xp_2^z} \tilde{B}_q^{\overline{\text{MS}}}(b^z, b_T, \eta, \mu, p_2^z)}$$

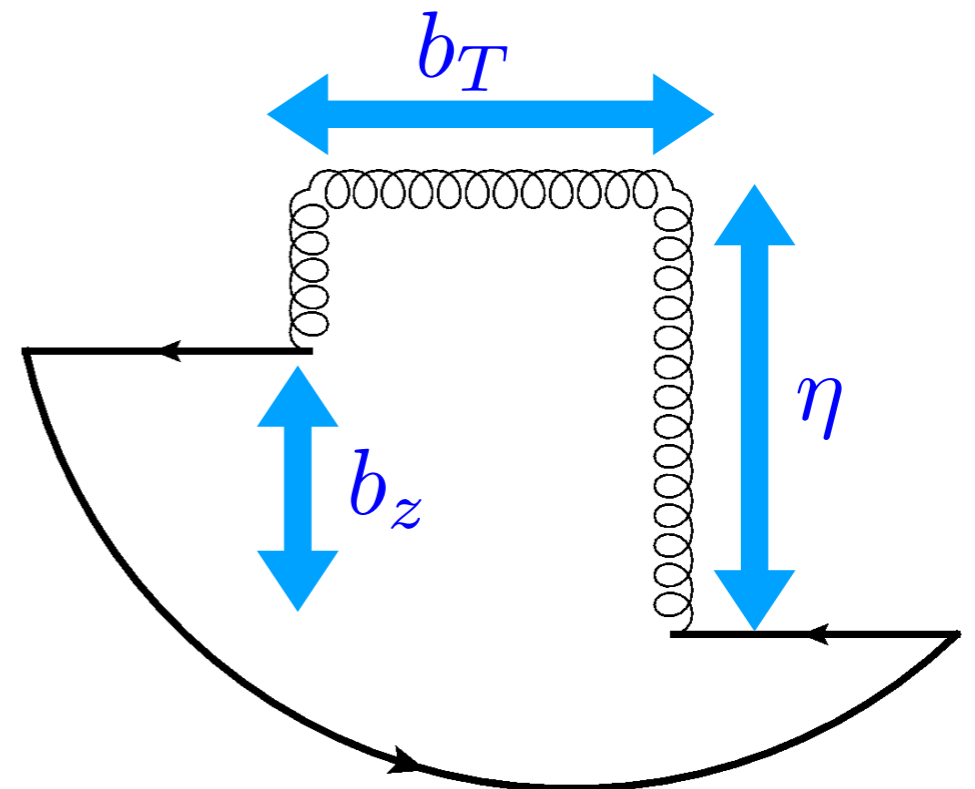
Independent of hadron state, choice of momenta, choice of  $x$

...up to power corrections:  $b_T/\eta$ ,  $1/(p^z b_T)$ ,  $M/p^z$

Exploit independence,  
calculate for valence pion  
with  $m_{\pi} \sim 1.2$  GeV

Variation of  $m_{\pi}$  probes  
power corrections

**Not** independent of sea quark mass,  
quenched gauge fields used for  
exploratory calculation



# Challenges

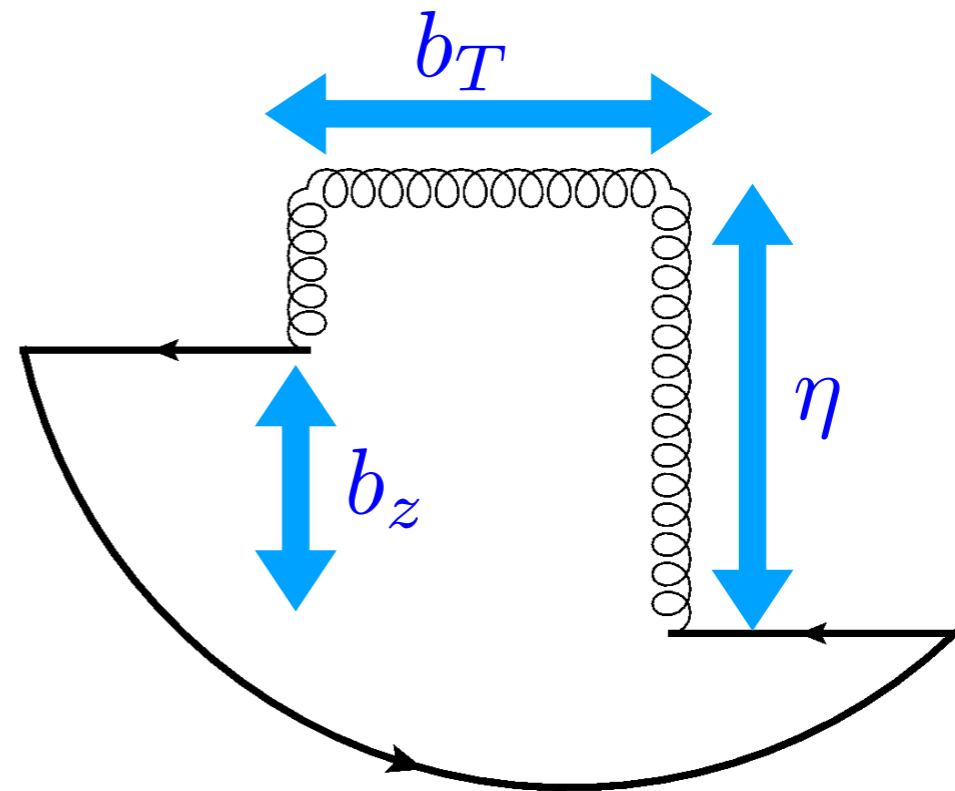
$$\tilde{B}_q^{\overline{\text{MS}}}(b^z, b_T, \eta, \mu, p^z) = \lim_{t, \tau, |t-\tau| \rightarrow \infty} 2E(p) \frac{G^{3pt}(b^z, b_T, \eta, \mu, p^z, t, \tau)}{G^{2pt}(p^z, t)}$$

- Large momentum. Noisy, need fine lattice spacing, ...

- Many bare matrix elements useful

$$B(b_z, b_T, \eta, p_z)$$

*Redundant  $\eta$  and  $p_z$  choices probe systematics*



- And many nonperturbative renormalization factors

$$Z(b_z, b_T, \eta, \mu)$$

$$0 \leq b_z, b_T \leq \eta \leq 10$$

$$p_z = 2, 3, 4$$

$$p_z^{max} \sim 2.6 \text{ GeV}$$

Analyze 2838  
3pt functions  
and NPR  
factors

- Fourier transform over  $b_z$

# Exploratory ensembles

Quenched Wilson gauge configurations — *generated by Michael Endres*

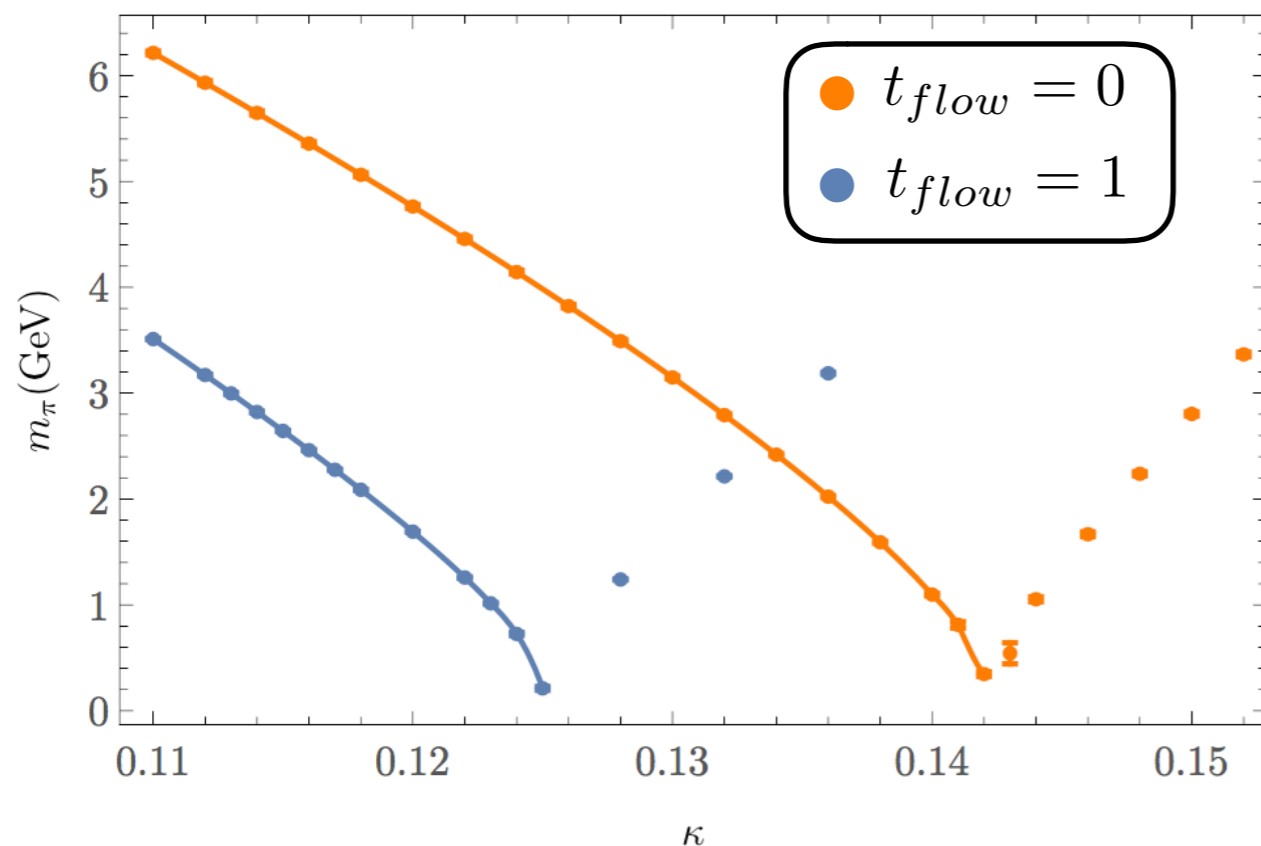
$$32^3 \times 64 \quad \beta = 6.30168$$

Lattice spacing determined from  
gradient flow scale-setting

$$a = 0.06(1) \text{ fm}$$

Lüscher, JHEP 1008 (2010)

Borsanyi et al, JHEP 1209 (2012)



Gradient flow also used as link  
smearing to reduce noise

— Bare quark mass tuned after  
smearing

Low statistics:

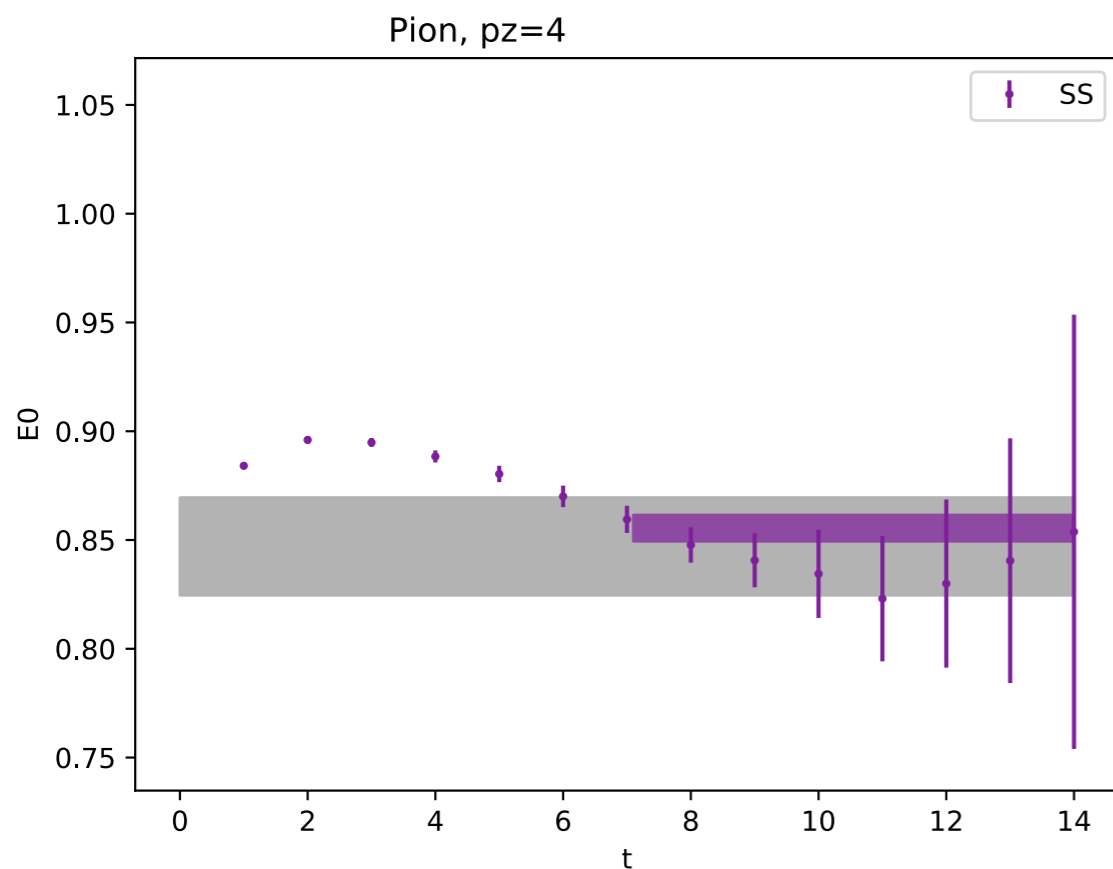
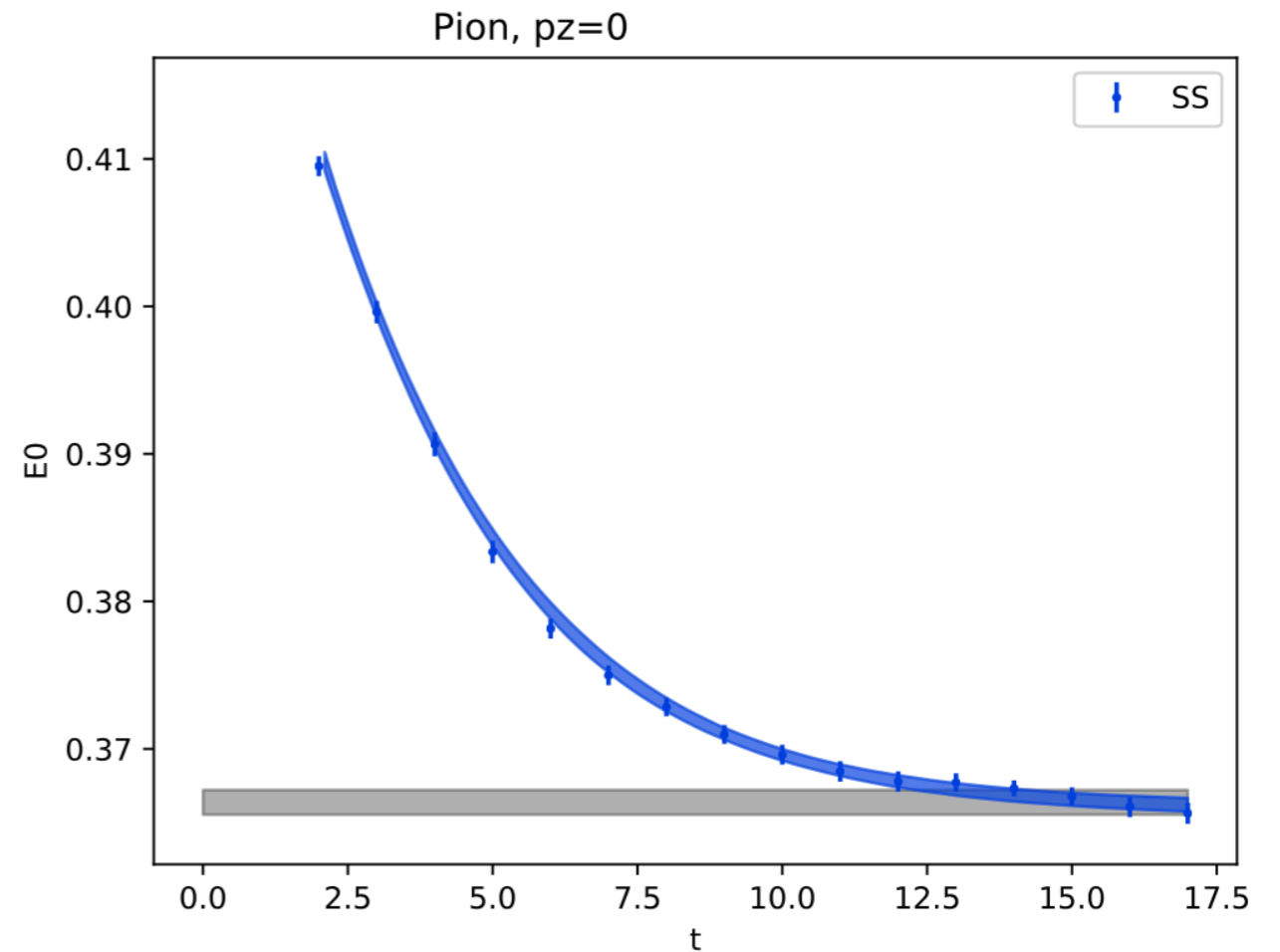
$N = 197$  cfigs x 2 sources/cfg

# Pion correlation functions

Pion energy fit from 2pt function

$$G_\pi(p, t) = \sum_x e^{ip \cdot x} \langle \pi(x, t) \pi(0)^\dagger \rangle$$

$$= \sum_n Z_n(p) e^{-E_\pi(p)t}$$



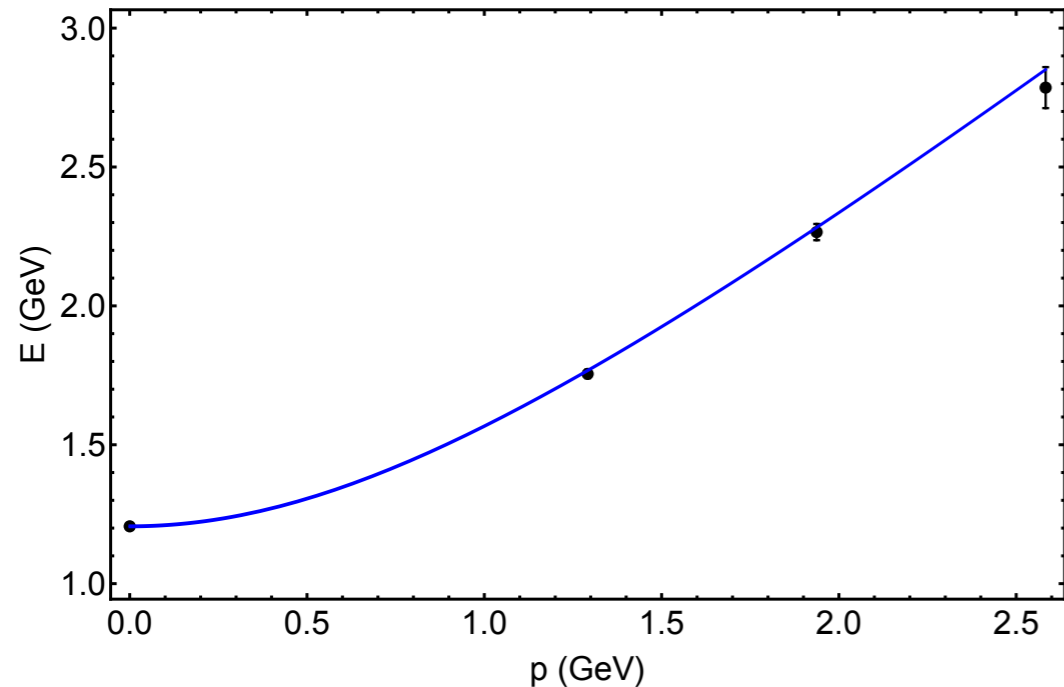
Boosted pion has exponential signal-to-noise problem

$$\text{StN}(t) \propto \sqrt{N} e^{-(E_\pi(p) - m_\pi)t}$$

Momentum smearing increases overlap and reduces noise

Bali, Lang, Musch, Schäfer, PRD 93 (2016)

# Boosted pion energies



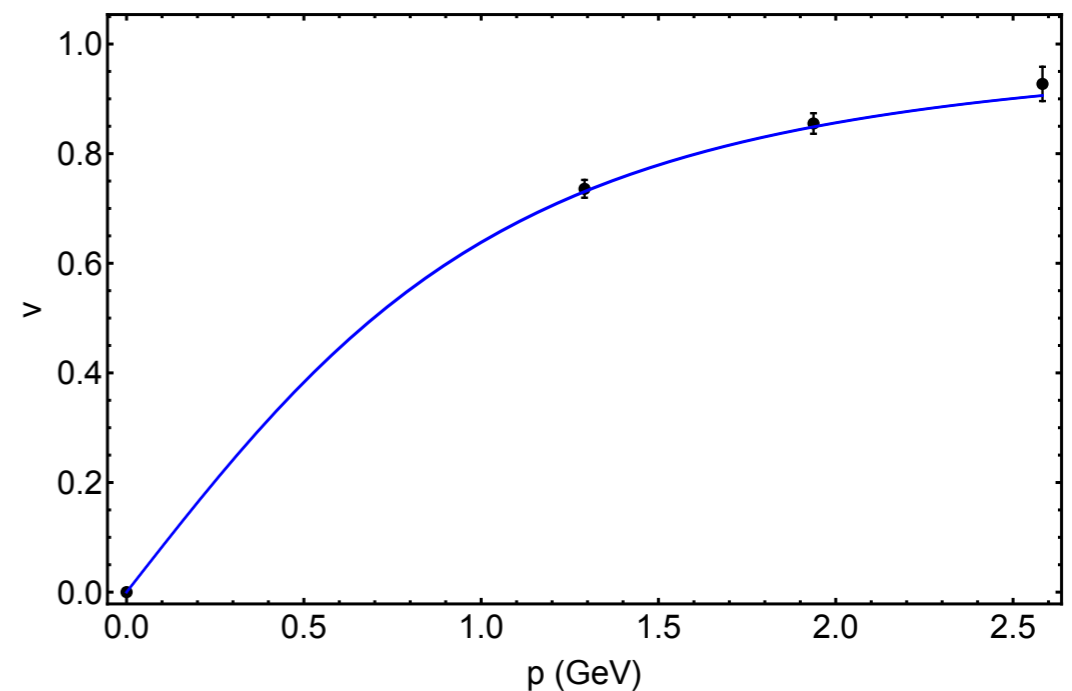
Results described well by  
continuum dispersion relation

$$p_z^{max} \sim 2.6 \text{ GeV}$$

$$a^{-1} \sim 3.3 \text{ GeV}$$

Boosted pions highly relativistic

$$v^{max} \sim 0.93$$





# Three-point functions

Ground-state matrix elements extracted from ratios of 3pt/2pt functions

$$\frac{G^{3pt}(b^z, b_T, \eta, \mu, p^z, t, \tau)}{G^{2pt}(p^z, t)} = \frac{\sum_{n,m} \mathcal{M}_{nm} \sqrt{Z_n Z_m} e^{-E_n(p^z)\tau} e^{-E_n(p^z)(t-\tau)}}{\sum_n Z_n e^{-E_n(p^z)\tau}}$$

Rearranging, matrix elements accessible from **linear** fit

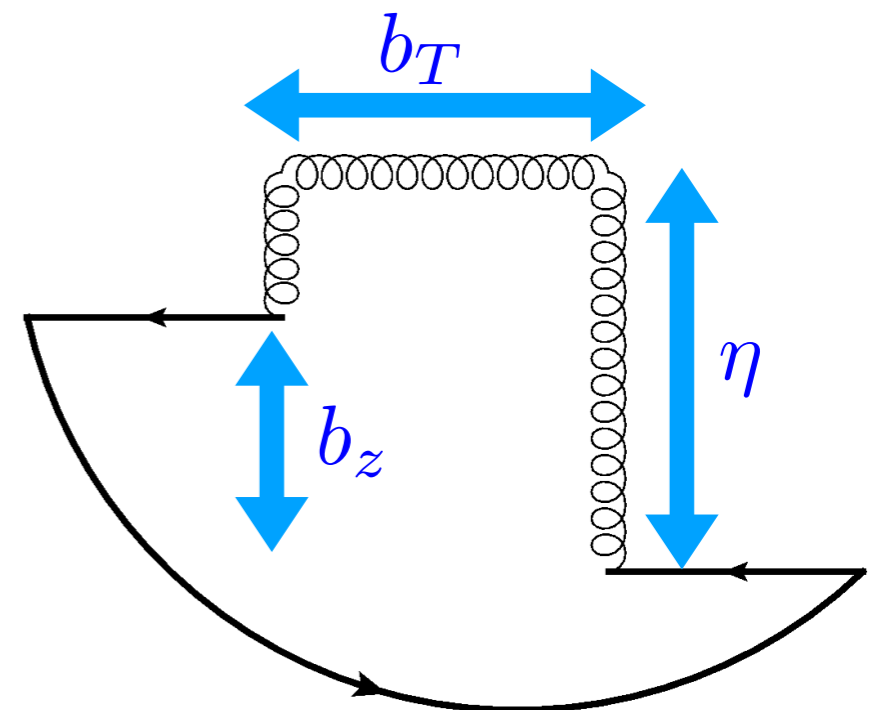
$$\left( 1 + \sum_{n>0} \frac{Z_n}{Z_0} e^{-[E_n(p^z) - E_0(p^z)]t} \right) \frac{G^{3pt}(b^z, b_T, \eta, \mu, p^z, t, \tau)}{G^{2pt}(p^z, t)} =$$

$$\mathcal{M}_{00} + \sum_{n,m>0} \tilde{\mathcal{M}}_{nm} e^{-[E_n(p^z) - E_0(p^z)]\tau} e^{-[E_n(p^z) - E_0(p^z)](t-\tau)}$$

Three source/sink separations:

$$t = 0.48, 0.6, 0.72 \text{ fm}$$

Simultaneous fit to all  $t, \tau$



# Fitting sums of exponentials

**What fit range(s)?**

tmax: signal-to-noise  $> 2$  (results mostly insensitive)

tmin: use all choices where operators are separated by width of transfer matrix

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Start with 1.

Try 2. Keep if preferred by information criterion, e.g. AIC

Try 3....

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Bootstrap

Optimal shrinkage (interpolates between correlated and uncorrelated fit)

Ledoit, Wolf, Journal of Multivariate Analysis 88 (2004)

Rinaldi, Syritsyn, MW et al, arXiv:1901.07519

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Weighted average, e.g. weight = p-value/variance

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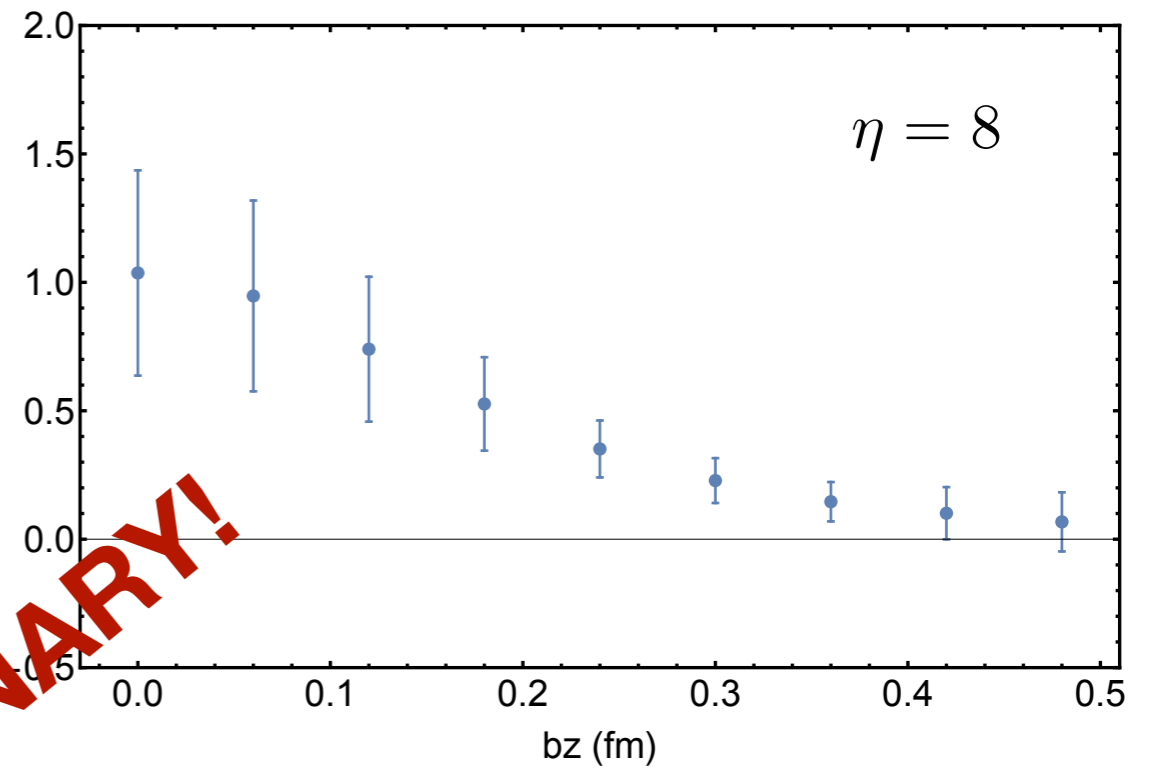
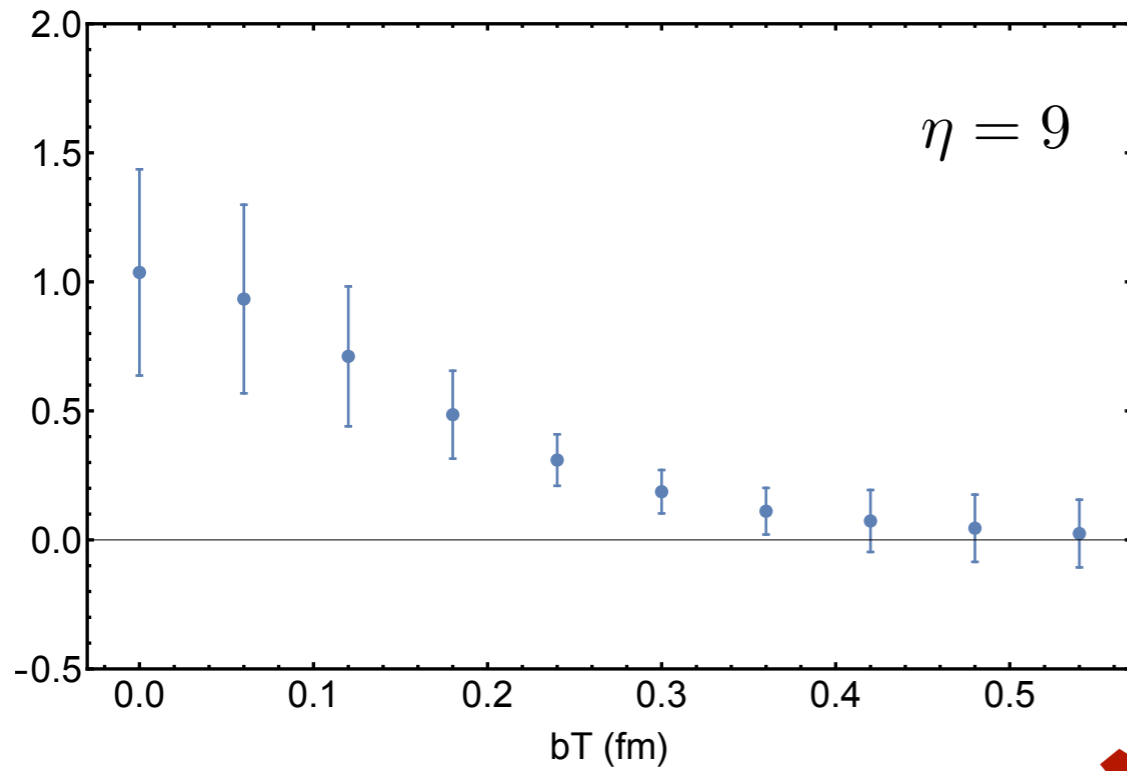
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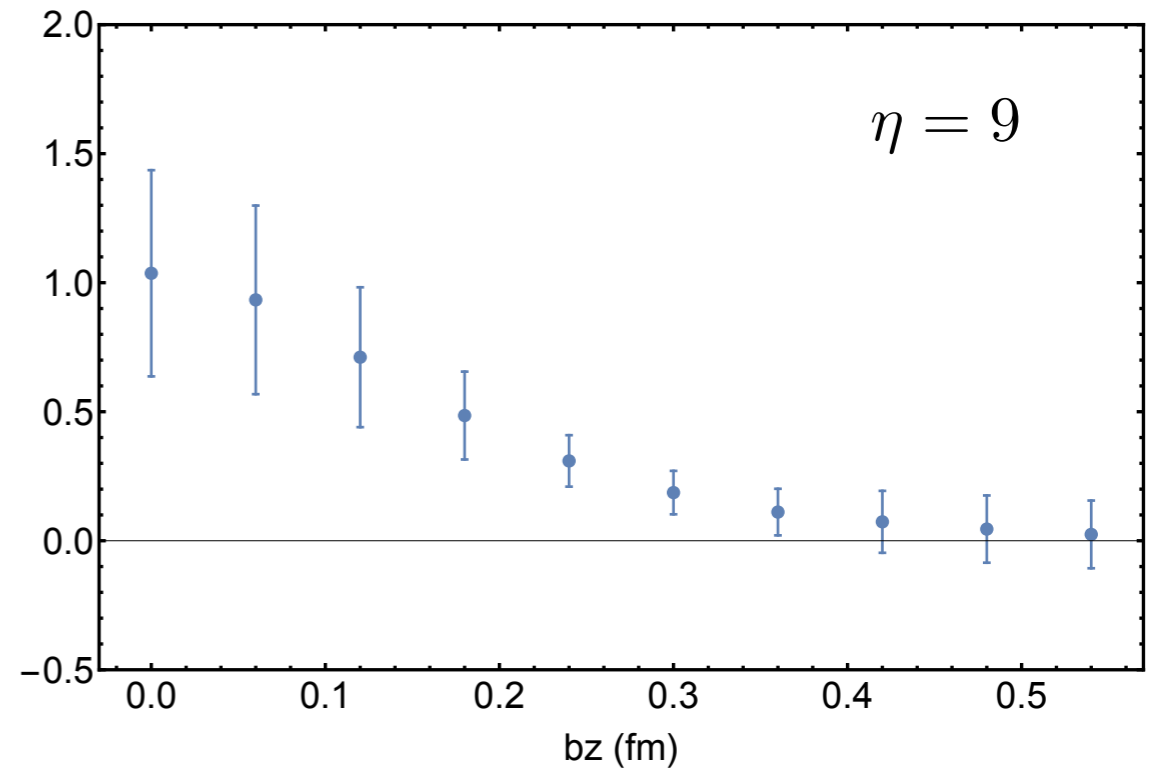
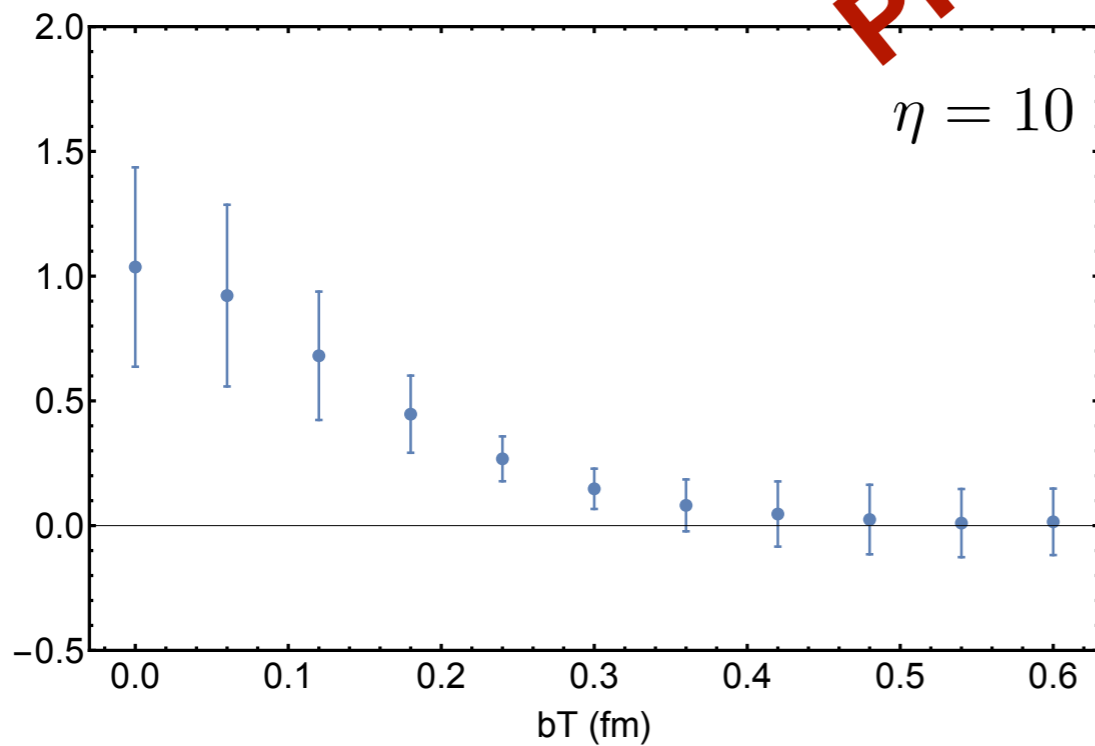


We're not using  
plateau fits!

# Bare beam functions



**PRELIMINARY!**



# Quark field renormalization

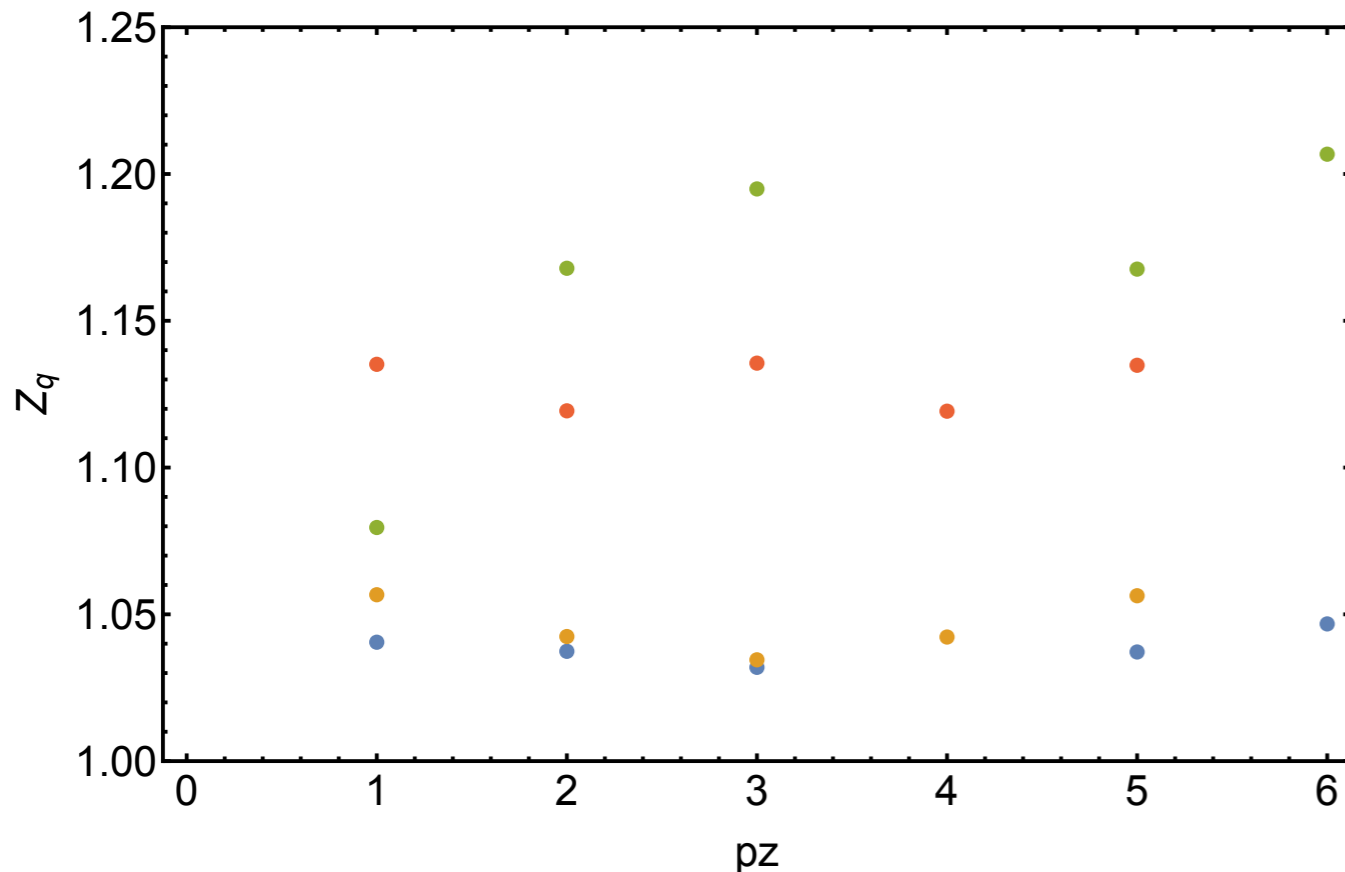
(Landau) gauge-fixed quark propagators accessible to LQCD and perturbation theory

$$S_{\alpha\beta}(p) = \sum_{x,y} e^{ip \cdot (x-y)} \langle q_\alpha(x) \bar{q}_\beta(y) \rangle$$

RI/MOM condition: full propagator = tree-level propagator (at one scale)

$$Z_q S(p) \Big|_{p^2=\mu_R^2} = S^{\text{tree}}(p)$$

$$\implies Z_q = \frac{1}{12} \text{Tr} [S^{-1}(p) S^{\text{tree}}(p)] \Big|_{p^2=\mu_R^2}$$



Momentum dependence  
O(10%) 1-loop effect

$Z_q^{\text{prop}}(p_1)$

$Z_q^{\text{prop}}(p_2)$

$Z_q^{\text{VC}}(p_1)$

$Z_q^{\text{VC}}(p_2)$

Alternative definition from conserved current

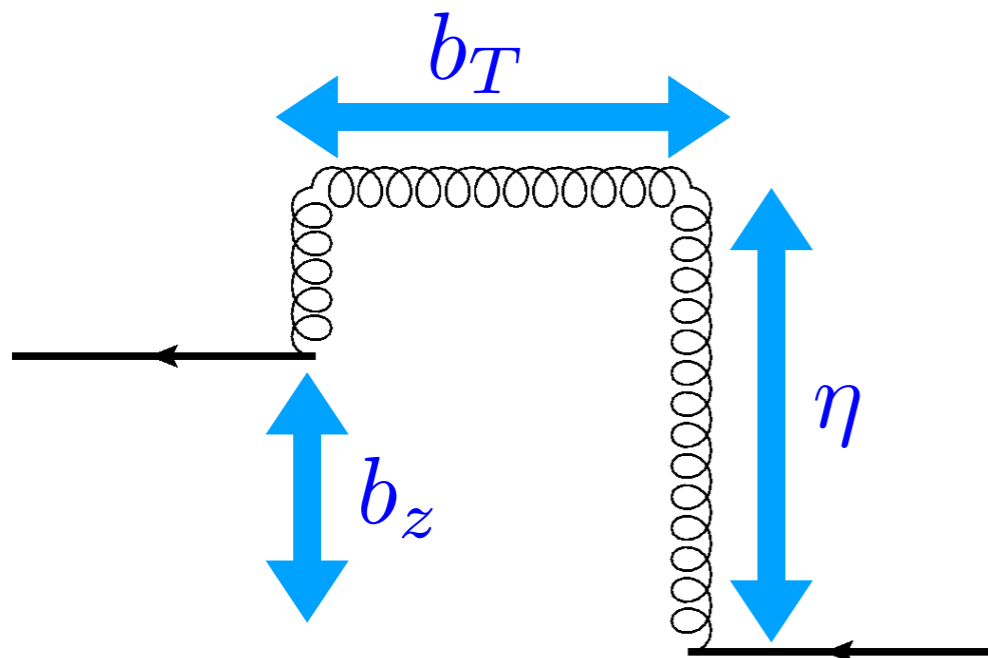


# Staple renormalization

Correlation function for nonlocal operator in gauge-fixed quark state

$$G_{\alpha\beta}(p) = \sum_{x,y,z} e^{ip \cdot (x-y)} \langle q_\alpha(x) \mathcal{O}(z+b, z) \bar{q}_\beta(y) \rangle$$

$$= \sum_z \langle \gamma_5 S^\dagger(p, b+z) \gamma_5 \tilde{W}(\eta; b+z, z) \frac{\Gamma}{2} S(p, z) \rangle_{\alpha\beta}$$



Vertex function accessible to LQCD and perturbation theory

$$\Lambda(p) = \left( \gamma_5 [S^{-1}(p)]^\dagger \gamma_5 \right) G(p) S^{-1}(p)$$

## RI/MOM condition

$$Z_q^{-1} Z_{\mathcal{O}} \text{Tr} [P_\Gamma \Lambda(p)] = \text{Tr} [P_\Gamma \Lambda^{\text{tree}}(p)] = 6e^{ip \cdot b}$$

# One-loop matching

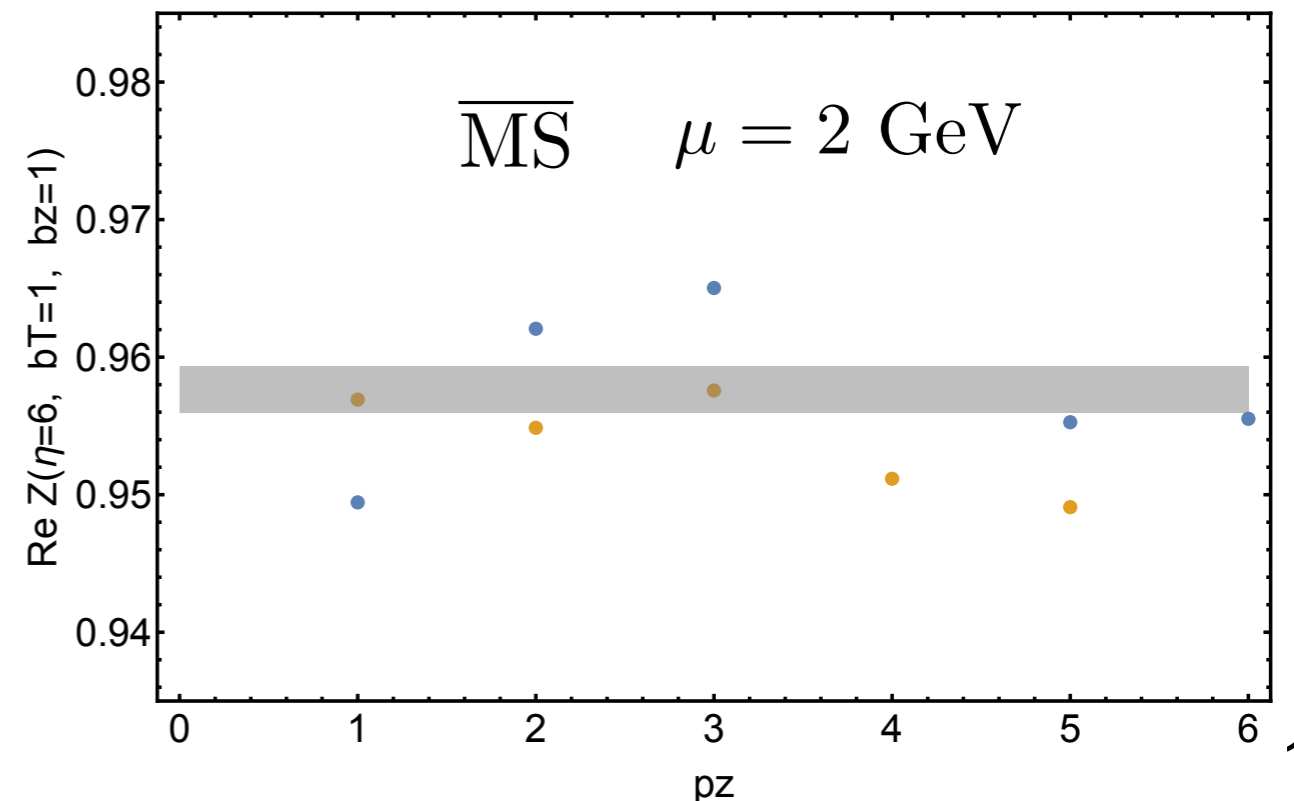
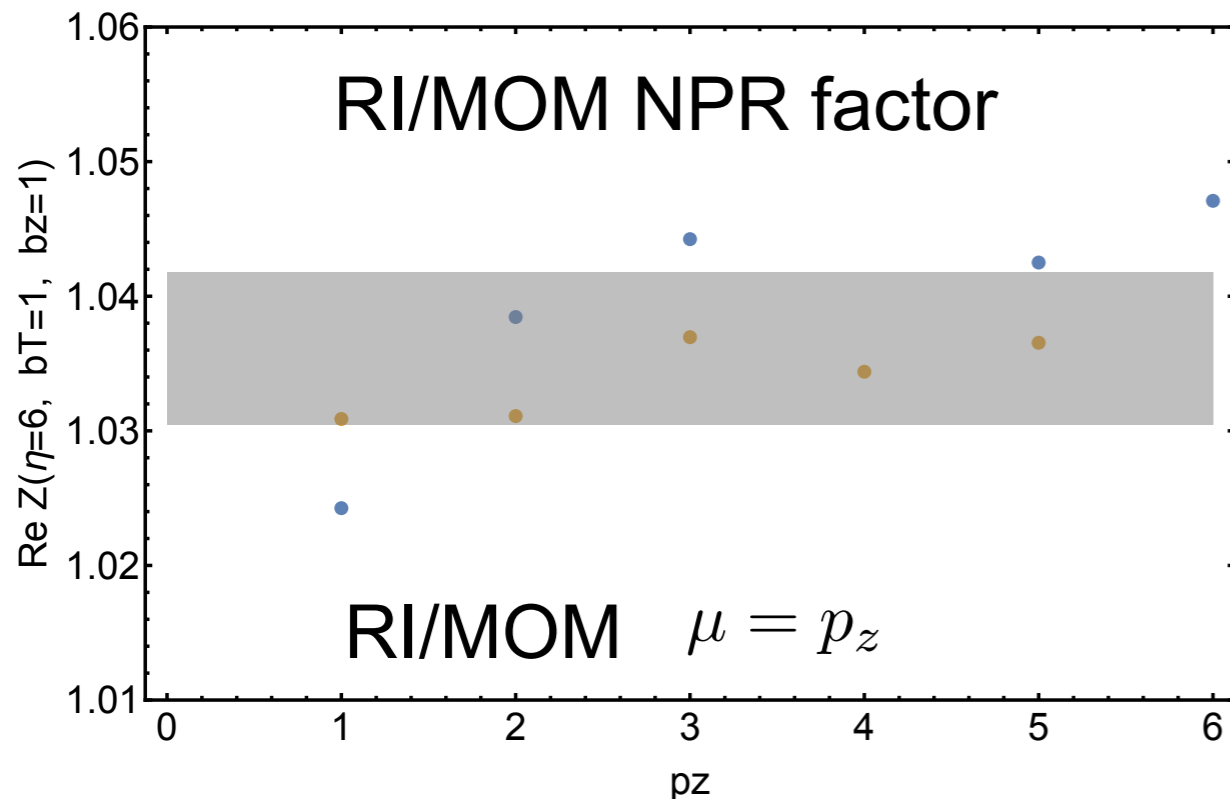
Dependence on  $p_z$  arises from perturbative RG and lattice artifacts

One-loop matching from RI/MOM  $\mu = p_z$  to  $\overline{\text{MS}}$   $\mu = 2 \text{ GeV}$  reduces  $p_z$  dependence

1-loop matching for asymmetric staple: Ebert, Stewart, Zhao *in preparation*

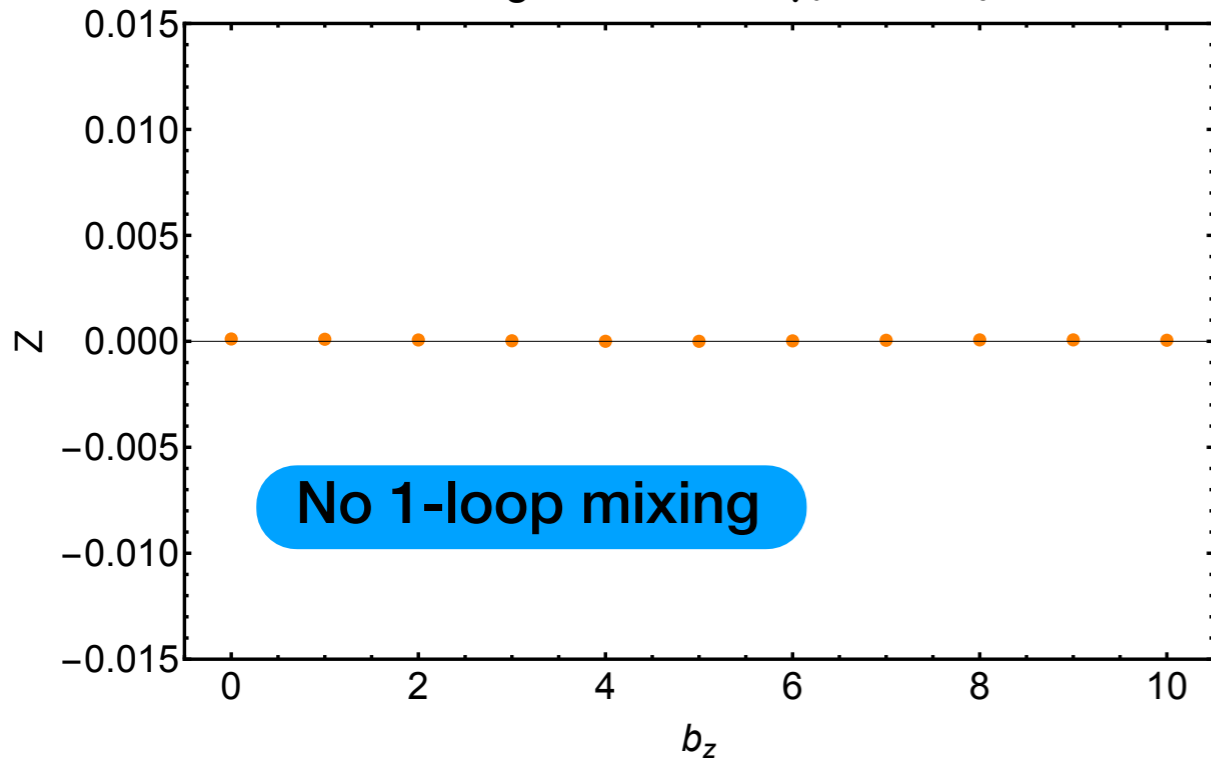
Residual  $p_z$  dependence from 2-loop RG and lattice artifacts

See e.g. Blossier et al (ETM), PRD 91 (2015)

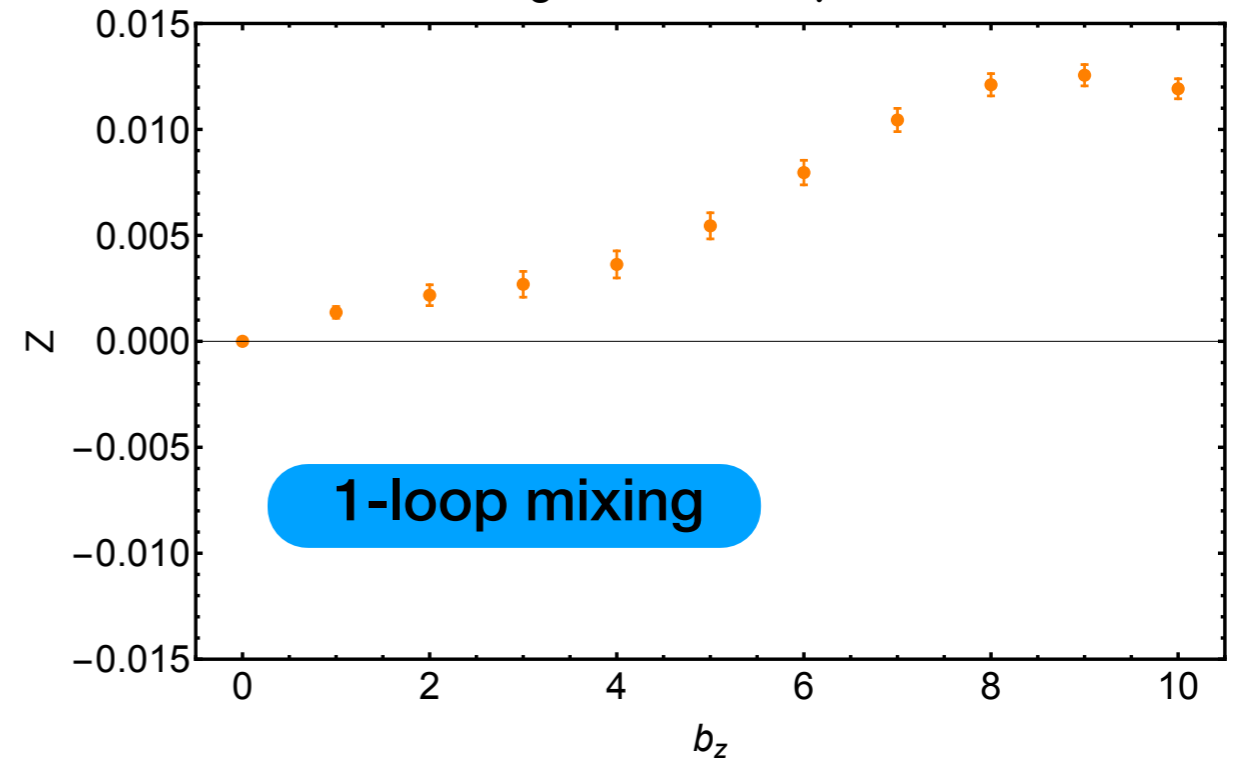


# Operator mixing results

Straight line:  $\Lambda = \gamma_t, P = \sigma_{zt}$



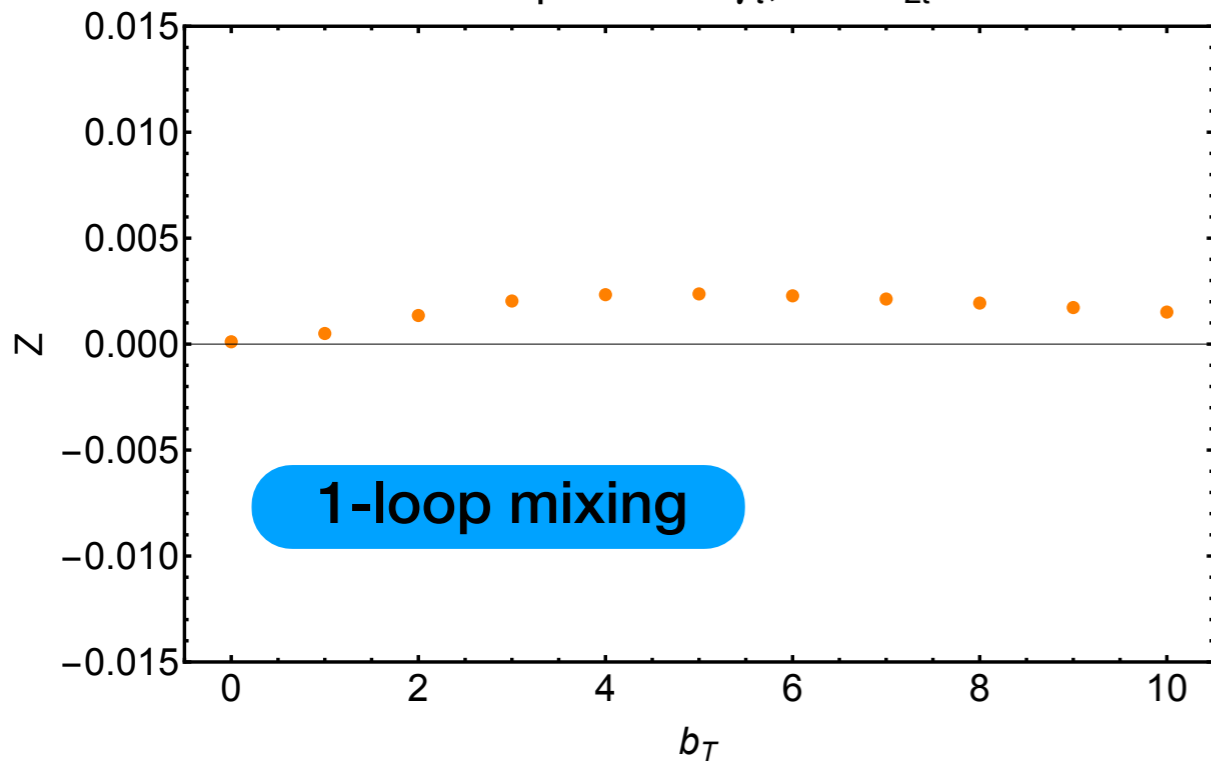
Straight line:  $\Lambda = \gamma_z, P = 1$



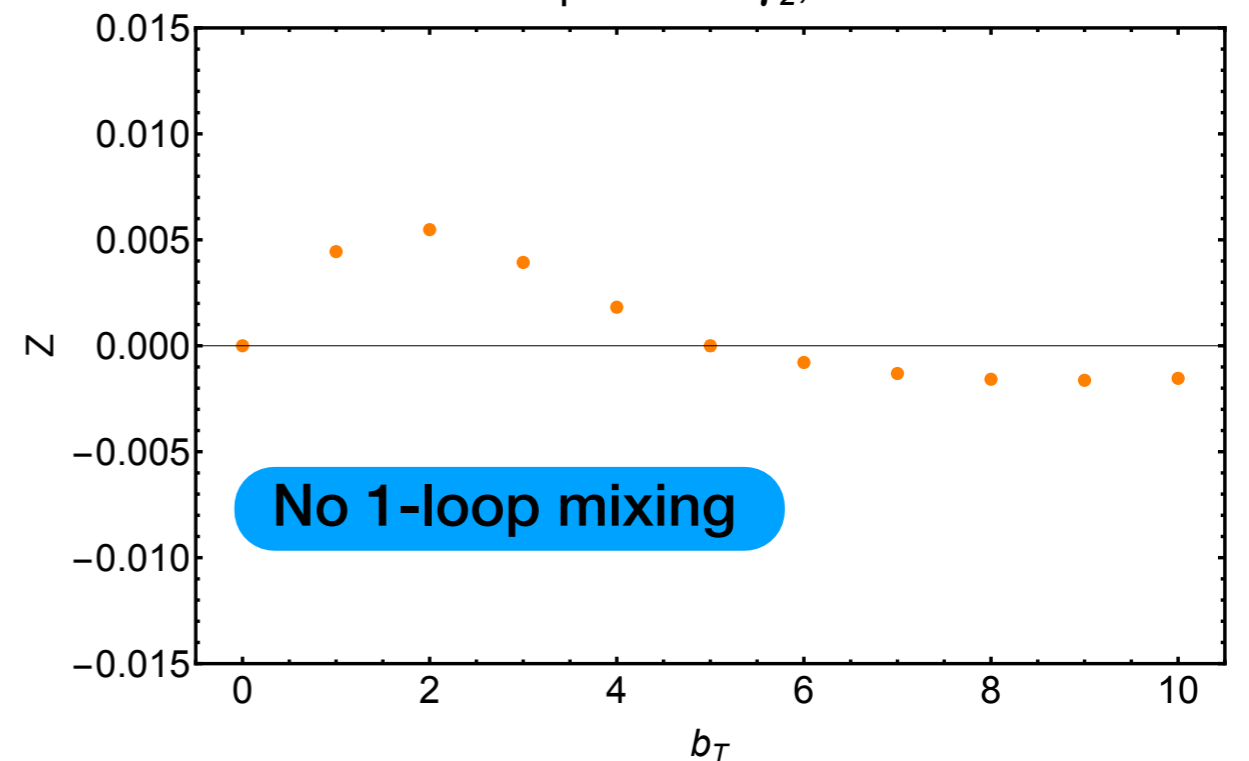
Lattice perturbation theory: Xiong, Ji, Zhang, Zhao, PRD 90 (2014)  
Constantinou, Panagopoulos, PRD 96 (2017)

Chen et al, arXiv:1710.01089  
Green, Jansen, Steffense, PRL 121 (2018)

Staple:  $\Lambda = \gamma_t, P = \sigma_{zt}$



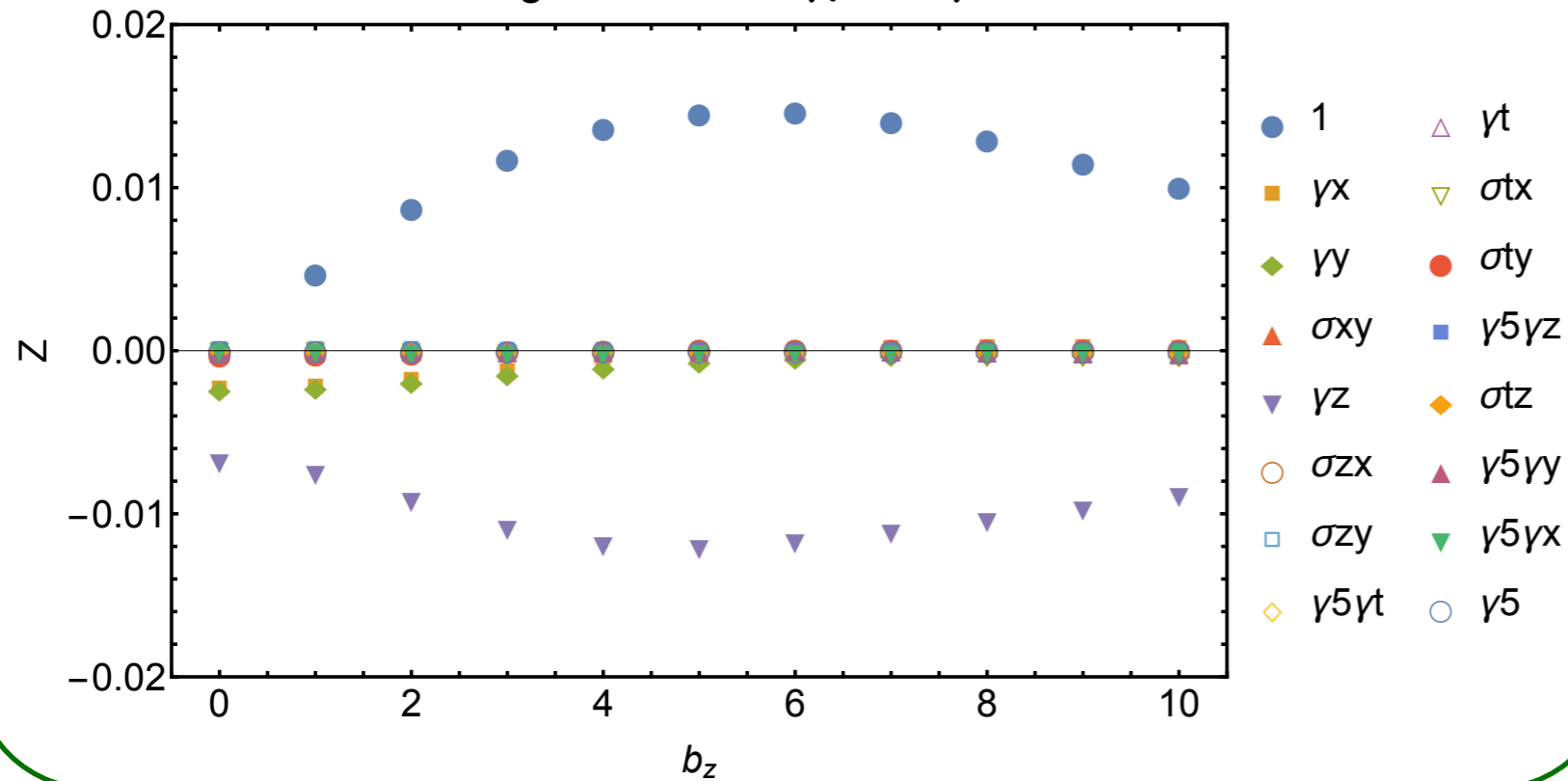
Staple:  $\Lambda = \gamma_z, P = 1$



Lattice perturbation theory: Constantinou, Panagopoulos and Spanoudes, arXiv:1901.03862

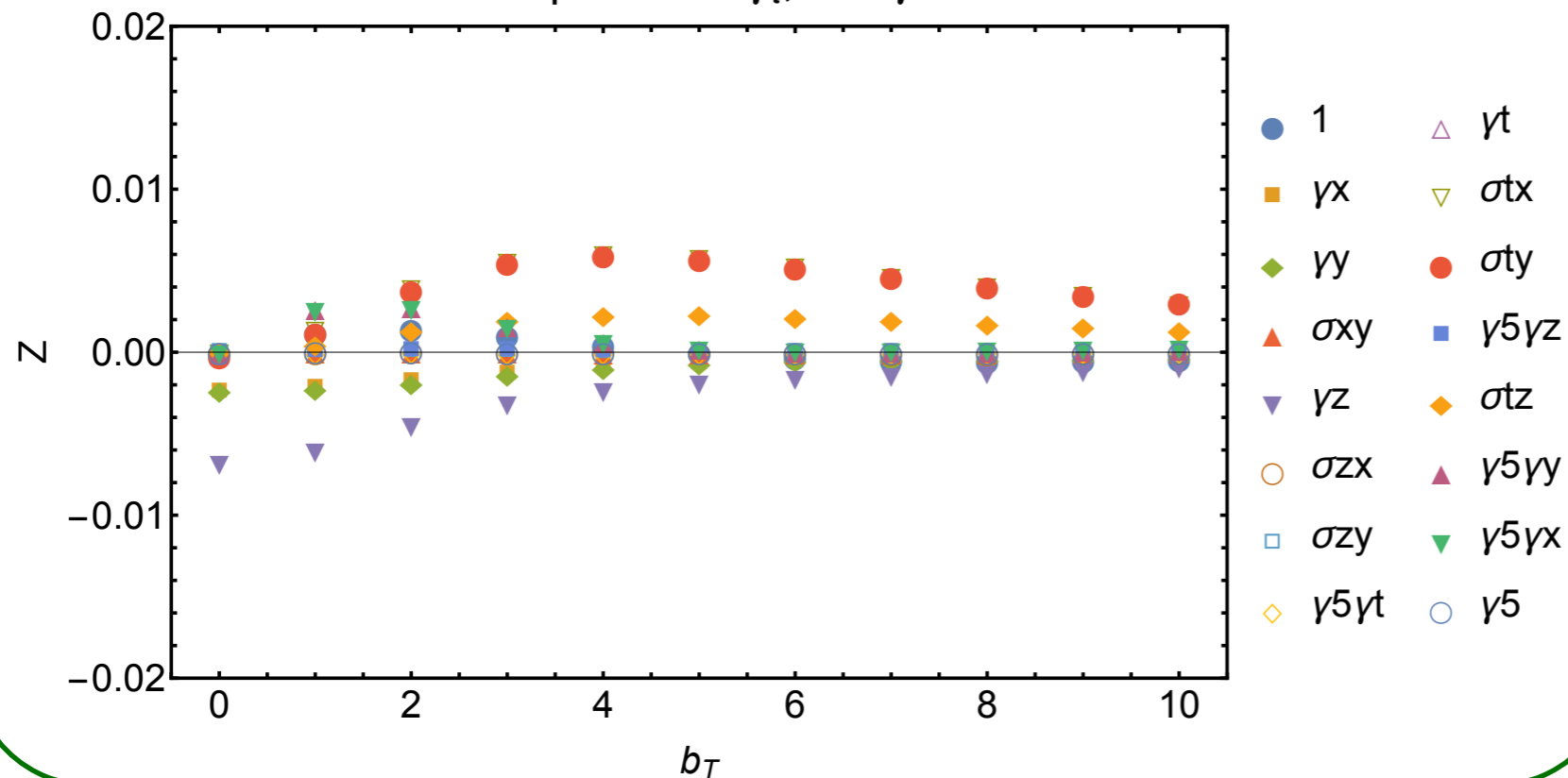
# Operator mixing is generic

Straight line:  $\Lambda = \gamma_t, P = \gamma$



Nonperturbative mixing between different currents generic for nonlocal operators

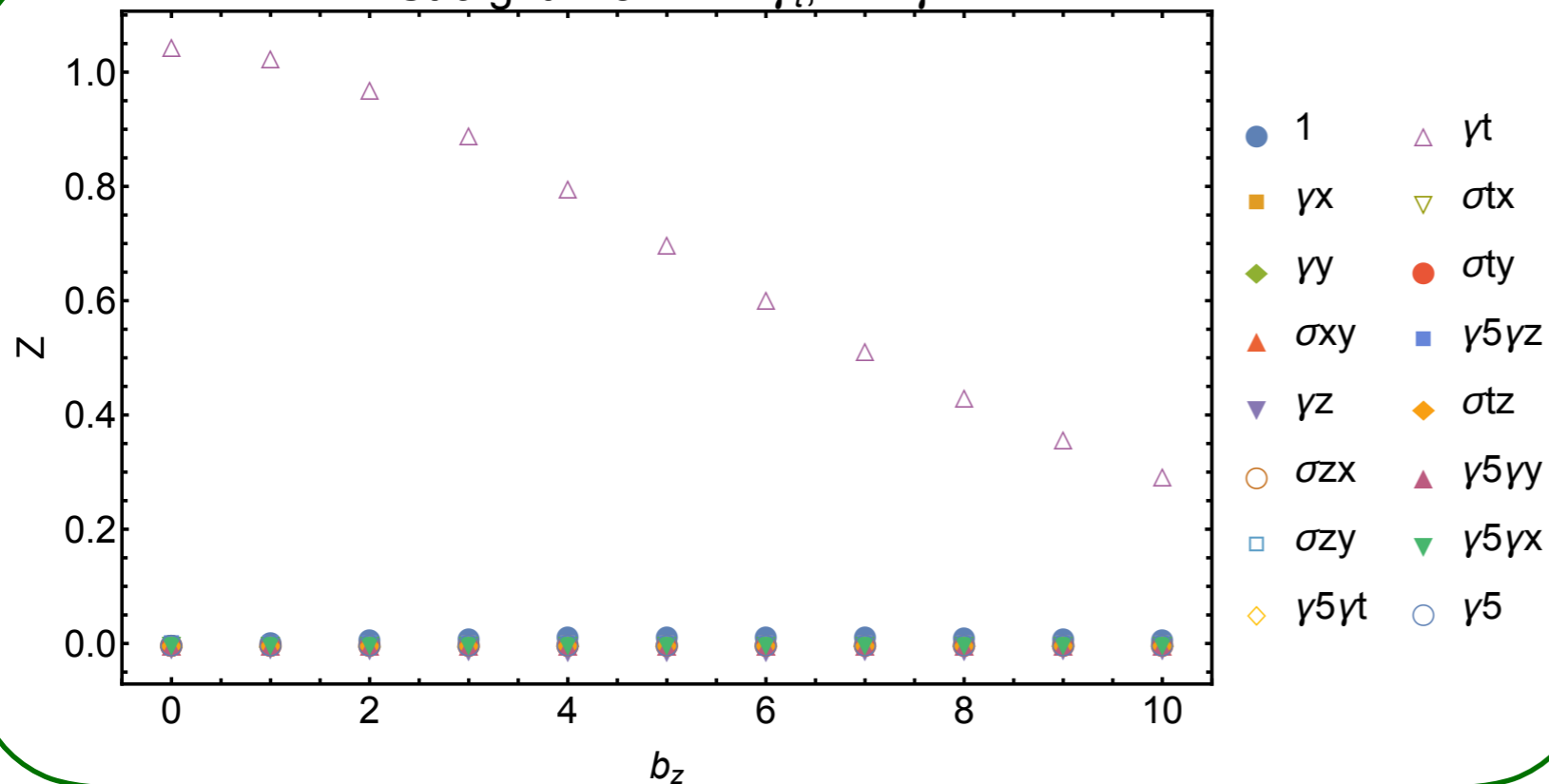
Staple:  $\Lambda = \gamma_t, P = \gamma$



Hierarchies between mixings visible but differ from 1-loop mixing pattern

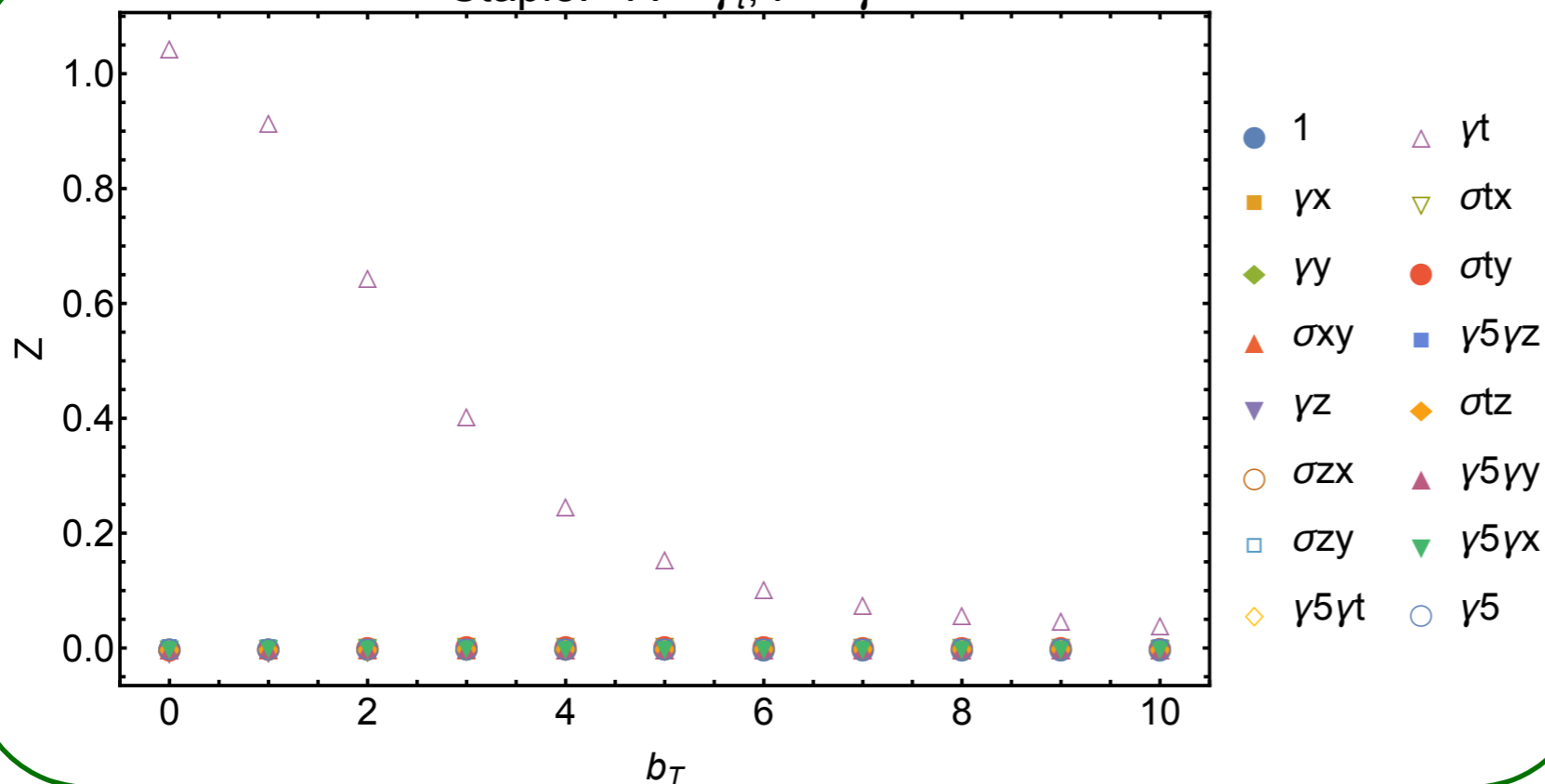
# Operator mixing is small

Straight line:  $\Lambda = \gamma_t, P = \gamma$



**Small** nonperturbative mixing between different currents generic for nonlocal operators

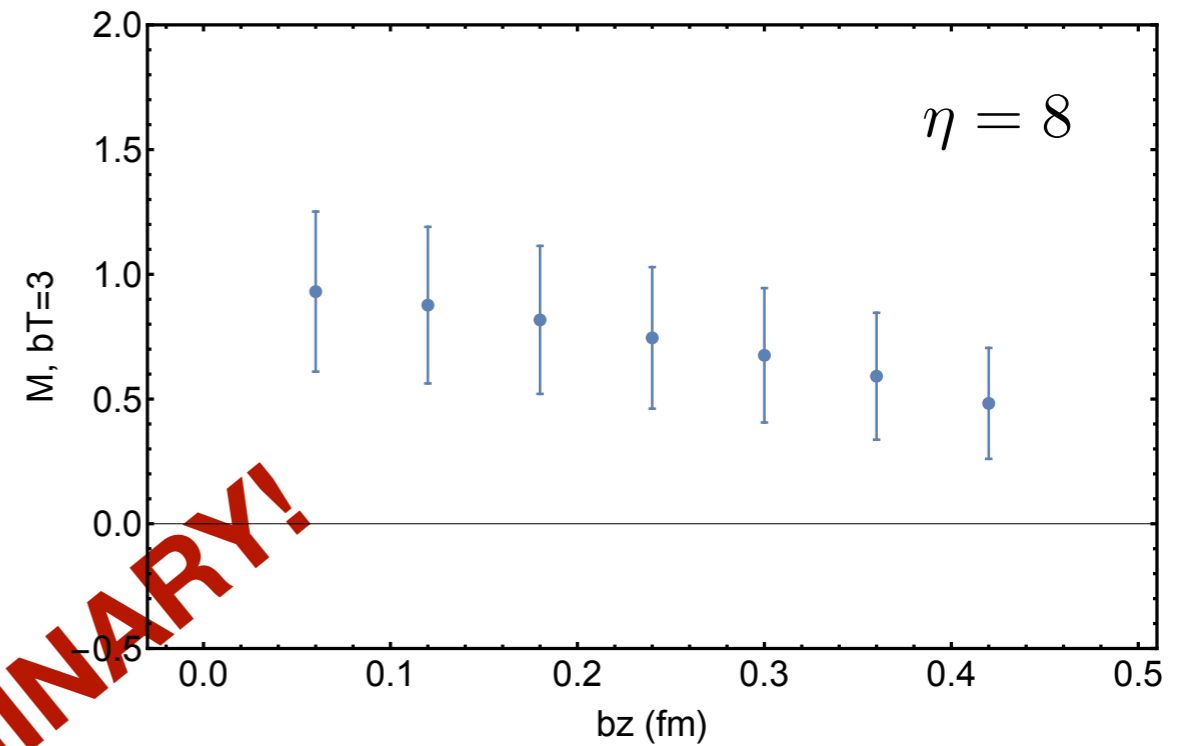
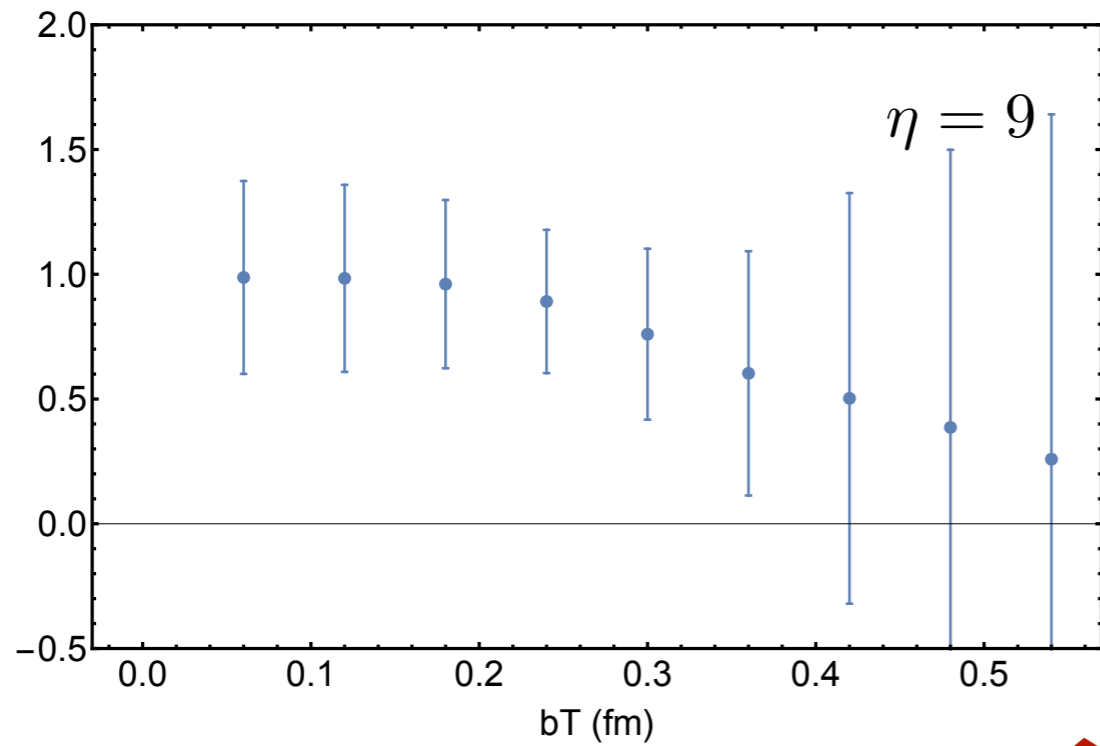
Staple:  $\Lambda = \gamma_t, P = \gamma$



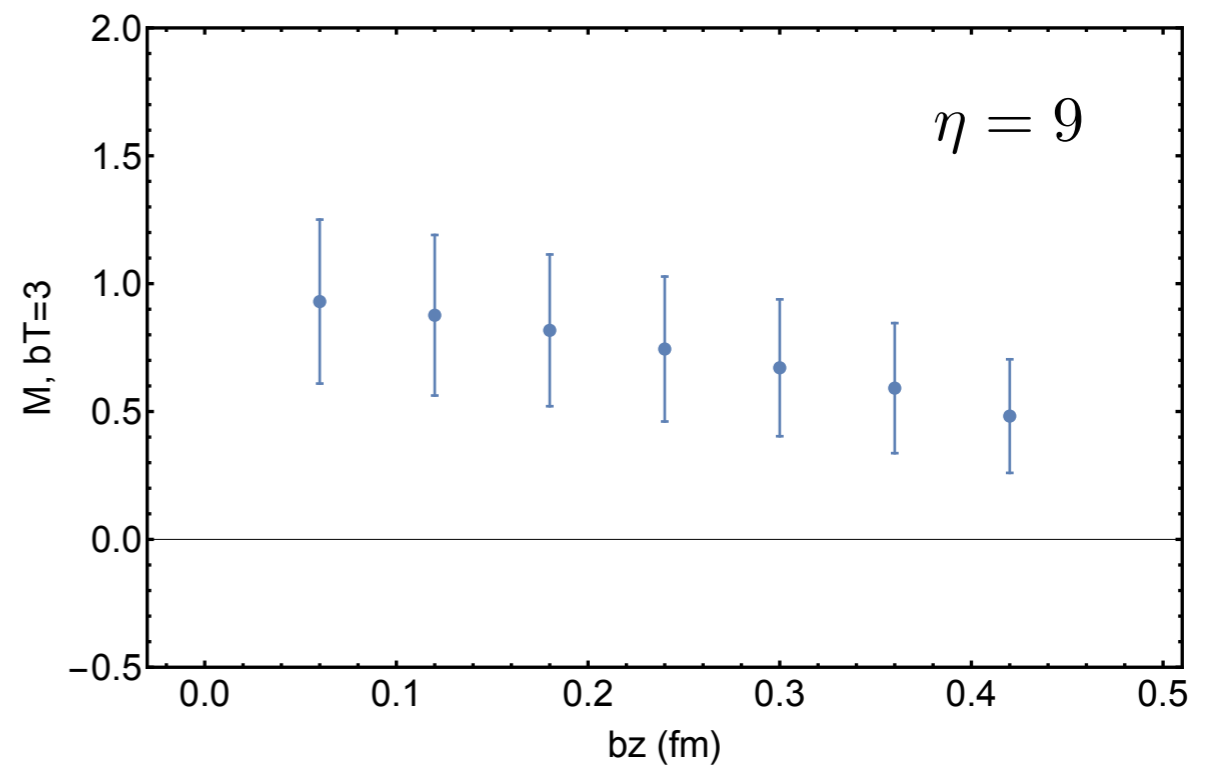
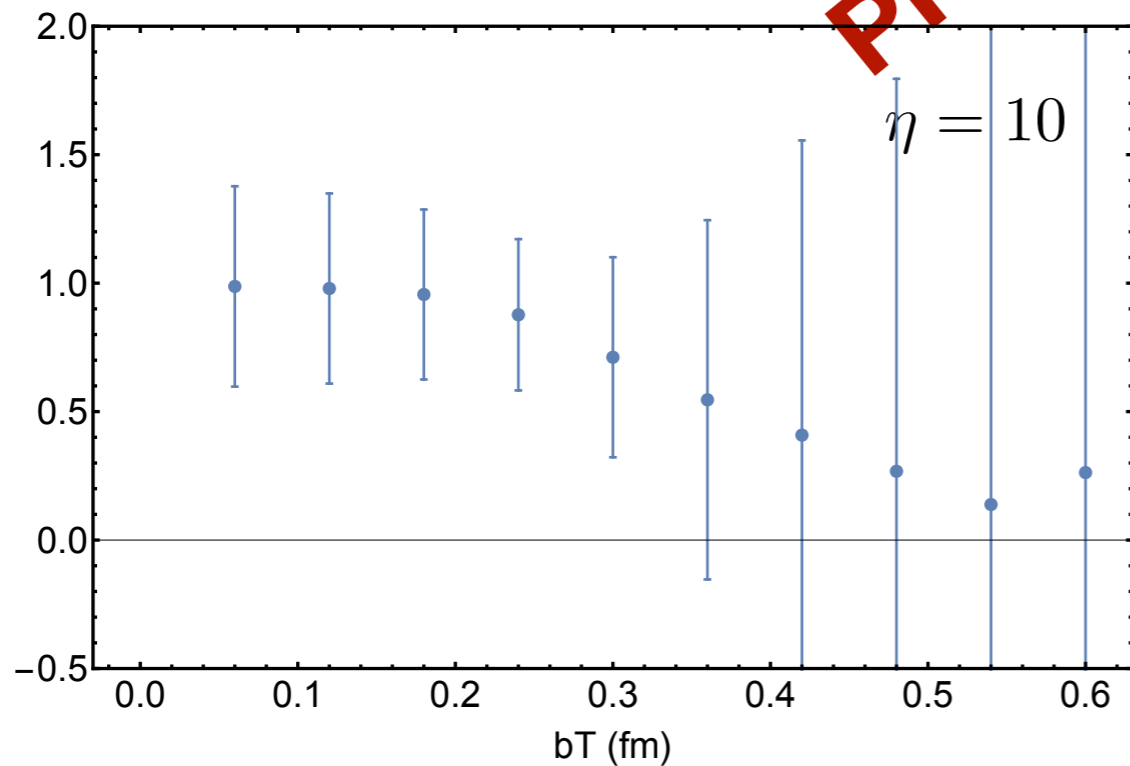
O(1%) effects negligible for exploratory calculations

Present in precision PDF calculations

# Beam functions



**PRELIMINARY!**



# Towards TMD evolution from LQCD

Next: Fourier transform beam function and form ratios

$$\gamma_{\zeta}^{q, \overline{\text{MS}}}(b_T, \mu) = \frac{1}{\ln(p_1^z/p_2^z)} \ln \frac{C_{\text{TMD}}^{\overline{\text{MS}}}(\mu, xP_2^z) \int db^z e^{ib^z xp_1^z} \tilde{B}_q^{\overline{\text{MS}}}(b^z, b_T, \eta, \mu, p_1^z)}{C_{\text{TMD}}^{\overline{\text{MS}}}(\mu, xp_1^z) \int db^z e^{ib^z xp_2^z} \tilde{B}_q^{\overline{\text{MS}}}(b^z, b_T, \eta, \mu, p_2^z)}$$

Detailed study of Fourier transform systematics important

Next-to-next: dynamical sea quarks, continuum extrapolation, finite volume, ...

**Stay tuned**

