Fluctuations and anisotropy in heavy-ion collisions: A new paradigm

by

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Based on:

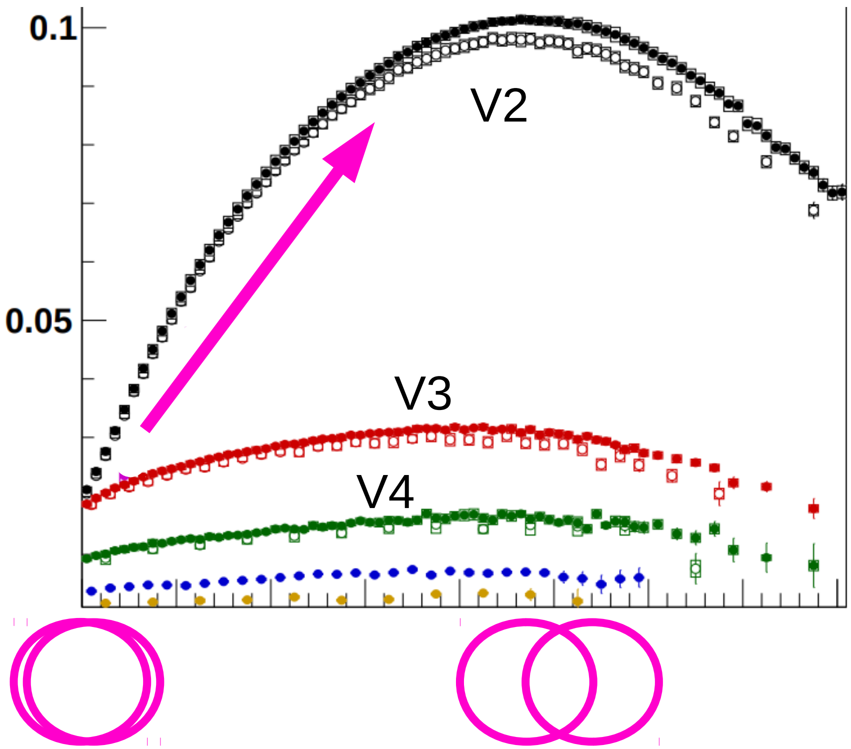
- Giacalone, Guerrero-Rodriguez, Luzum, Marquet, Ollitrault 1902.07168

Message from data: Anisotropy from anisotropy.

Final-state harmonic ← $$v_n = \kappa_n \varepsilon_n$$ → Initial-state harmonic

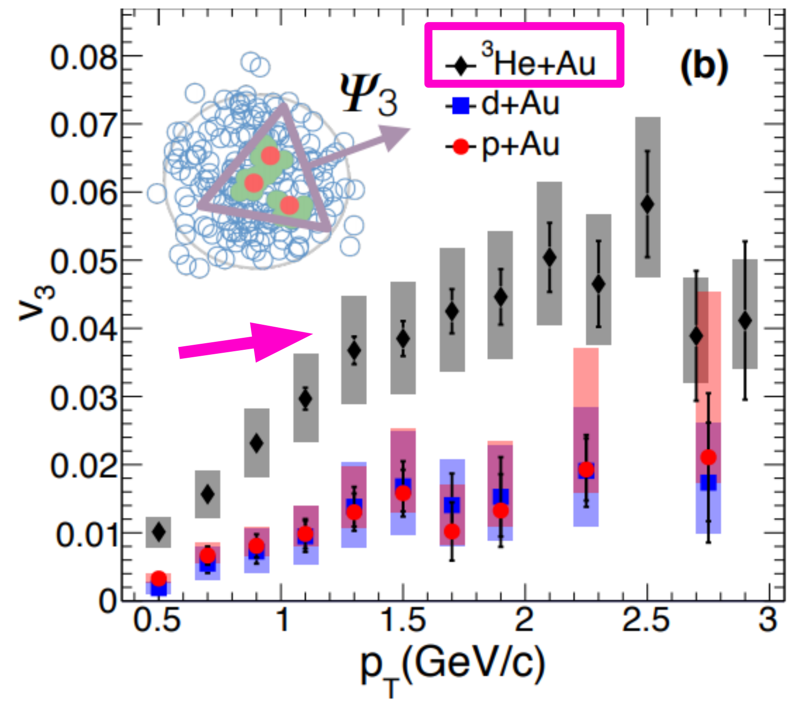
[ALICE collaboration, 1804:02944]

n=2



[PHENIX collaboration, 1805:02973]

n=3



This talk: AA + ‘fluid paradigm’

κ_n → Just a number. Rather independent of centrality up to ~20-30%.
[Noronha-Hostler, Yan, Gardim, Ollitrault, arXiv 1511:03869]

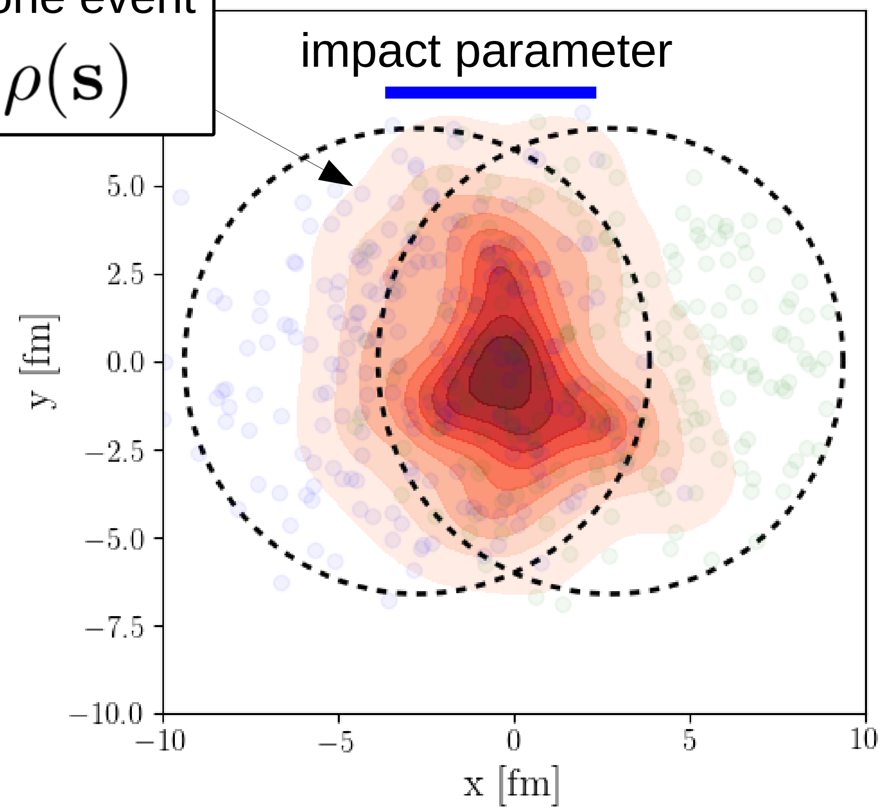
How do we calculate the initial anisotropy?

[Teaney, Yan 1010.1876]

$$\varepsilon_n \equiv \frac{\int_{\mathbf{s}} \mathbf{s}^n \rho(\mathbf{s})}{\int_{\mathbf{s}} |\mathbf{s}|^n \rho(\mathbf{s})}$$

$$\mathbf{s} = x + iy$$

density field in one event $\rho(\mathbf{s})$



Origin of anisotropy:

n=2
elliptic flow ⟶ **geometry** + **fluctuations**
[PHOBOS Collaboration **nucl-ex/0610037**]

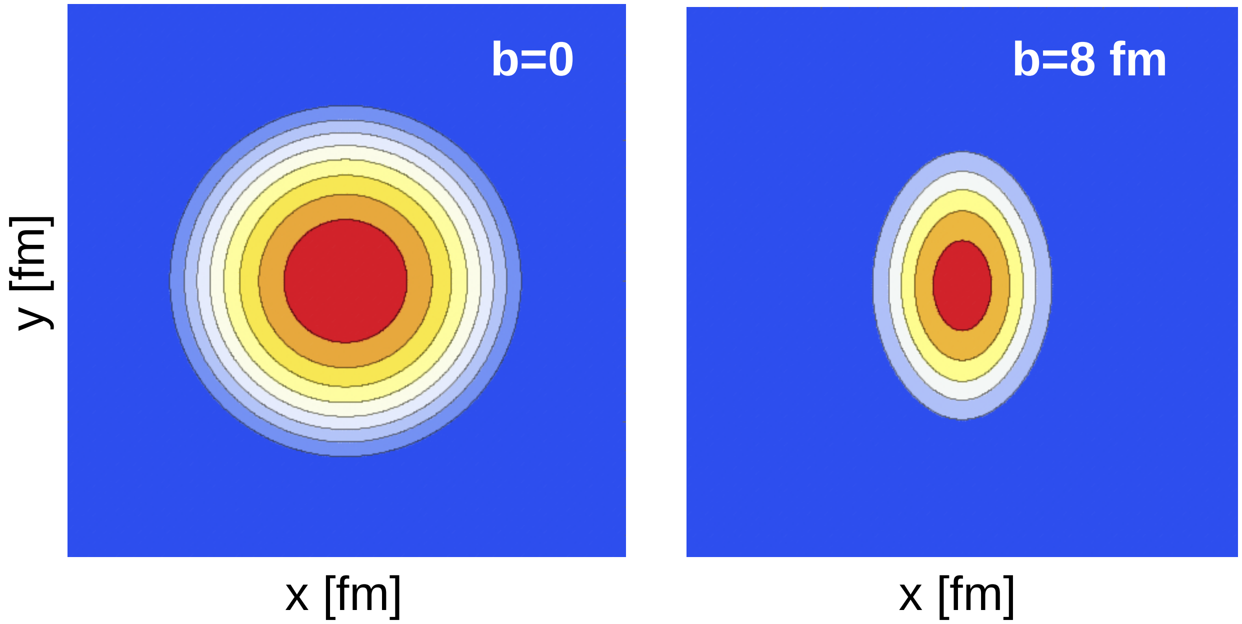
n=3
triangular flow ⟶ **fluctuations only**
[Alver, Roland **1003.0194**]

The theoretical input is a model for $\rho(\mathbf{s})$ and its fluctuations.

What do we need?

- $\langle \rho(\mathbf{s}) \rangle$

The average density.

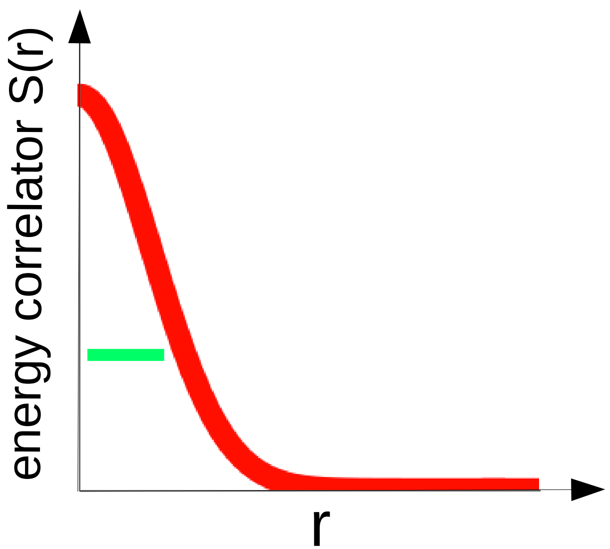


- **The fluctuations** around the average.

$$S(\mathbf{s}_1, \mathbf{s}_2) = \langle \rho(\mathbf{s}_1) \rho(\mathbf{s}_2) \rangle - \langle \rho(\mathbf{s}_1) \rangle \langle \rho(\mathbf{s}_2) \rangle$$

Density of variance.
Quantifies the correlation of fluctuations.

$$\mathbf{r} = \mathbf{s}_1 - \mathbf{s}_2$$



Input from high-energy QCD (or CGC). Nuclei characterized by a scale:

$$Q_s^2(\mathbf{s}) \propto T(\mathbf{s})$$

← Nuclear thickness

Proportional to the density of ‘color charges’, that source the energy density.
Average energy density after the collision known for a long time:

$$\langle \rho(\mathbf{s}) \rangle \propto Q_A^2(\mathbf{s}) Q_B^2(\mathbf{s})$$

[Lappi **hep-ph/0606207**]
[Lappi, Venugopalan **nucl-th/0609021**]

Fluctuation calculated recently:

[Albacete, Guerrero-Rodriguez, Marquet **1808.00795**]

The leading order of the expansion, of order N_c^0, reads:

$$\begin{aligned}
&[\mathrm{Cov}[\epsilon_{\rm MV}](0^+;x_\perp,y_\perp)]_{N_c^0} \\
&=\Bigg[\frac{1}{g^4 r^8} e^{-\frac{r^2}{2}(Q_{s1}^2+Q_{s2}^2)}\Bigg(16+32e^{\frac{Q_{s1}^2 r^2}{2}} \\
&-64e^{\frac{Q_{s1}^2 r^2}{4}}-4e^{\frac{r^2}{4}(2Q_{s1}^2+Q_{s2}^2)}\left(Q_{s2}^4 r^4-2\left(4\pi\partial^2 L(0_\perp)\right)^2\bar{Q}_{s1}^4 r^4+8Q_{s2}^2 r^2+48\right) \\
&+\frac{1}{8}e^{\frac{r^2}{4}(Q_{s1}^2+Q_{s2}^2)}\left(Q_{s1}^4 Q_{s2}^4 r^8+(4Q_{s1}^2 Q_{s2}^2 r^6+128r^2)(Q_{s1}^2+Q_{s2}^2)+16r^4\left(Q_{s1}^2+Q_{s2}^2\right)^2+1024\right) \\
&+2e^{\frac{r^2}{2}(Q_{s1}^2+Q_{s2}^2)}\left(\bar{Q}_{s1}^4 r^4(Q_{s2}^2 r^2-4)\left(4\pi\partial^2 L(0_\perp)\right)^2+40\right)\Bigg)\Bigg]+[1\leftrightarrow 2]. \qquad (4.48)
\end{aligned}$$

The next term, of order N_c^{-2}, reads:

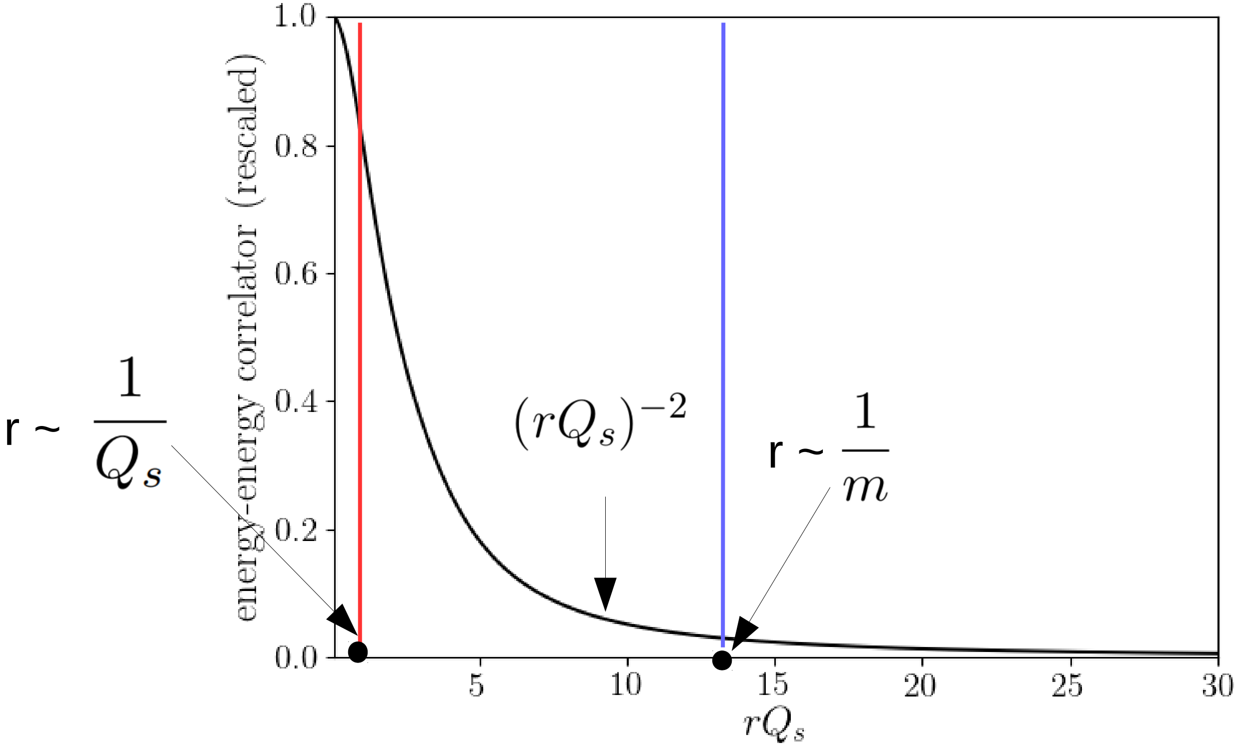
$$\begin{aligned}
&[\mathrm{Cov}[\epsilon_{\rm MV}](0^+;x_\perp,y_\perp)]_{N_c^{-2}} \\
&=\Bigg[\frac{1}{N_c^2 g^4 r^8} e^{-\frac{r^2}{2}(Q_{s1}^2+Q_{s2}^2)}\Bigg(2\left(Q_{s1}^2 r^2+Q_{s2}^2 r^2+8\right)^2 \\
&+4Q_{s1}^2 r^2(8+Q_{s1}^2 r^2)e^{\frac{Q_{s2}^2 r^2}{2}}-8(8+Q_{s1}^2 r^2)(4+Q_{s1}^2 r^2)e^{\frac{Q_{s2}^2 r^2}{4}} \\
&+4e^{\frac{r^2}{4}(2Q_{s1}^2+Q_{s2}^2)}\left(Q_{s2}^4 r^4-2(4\pi\partial^2 L(0_\perp))^2\bar{Q}_{s1}^4 r^4+8Q_{s2}^2 r^2+16Q_{s1}^2 r^2\right) \\
&-\frac{1}{8}e^{\frac{r^2}{4}(Q_{s1}^2+Q_{s2}^2)}\left(Q_{s1}^4 Q_{s2}^4 r^8+(4Q_{s1}^2 Q_{s2}^2 r^6+128r^2)(Q_{s1}^2+Q_{s2}^2)+16r^4\left(Q_{s1}^2+Q_{s2}^2\right)^2-1024\right) \\
&-2e^{\frac{r^2}{2}(Q_{s1}^2+Q_{s2}^2)}\left(\bar{Q}_{s1}^4 r^4(Q_{s2}^2 r^2-4)(4\pi\partial^2 L(0_\perp))^2+32Q_{s1}^2 r^2-4Q_{s1}^2 Q_{s2}^2 r^4\right)\Bigg)\Bigg]+[1\leftrightarrow 2]. \qquad (4.49)
\end{aligned}$$

These expressions give:

$$S(\mathbf{s}_1,\mathbf{s}_2) = \langle \rho(\mathbf{s}_1)\rho(\mathbf{s}_2)\rangle - \langle \rho(\mathbf{s}_1)\rangle\langle \rho(\mathbf{s}_2)\rangle$$

in collisions of large nuclei.

what are the relevant features?



$$\frac{1}{Q_s} \ll \frac{1}{m}$$

- It is very sharp compared to the system size.

  Short-range correlations: $S(\mathbf{s}_1, \mathbf{s}_2) \approx \xi(\mathbf{s})\delta(\mathbf{r})$, $\mathbf{s} = \frac{\mathbf{s}_1 + \mathbf{s}_2}{2}$

- Its integral is divergent and dominated by r^-2 tail.

An infrared cutoff naturally emerges, around the size of the nucleon.

Scales are separated: we can isolate the leading contribution.

$$\xi(\mathbf{s}) \equiv \int_{\mathbf{r}} S\left(\mathbf{s} + \frac{\mathbf{r}}{2}, \mathbf{s} - \frac{\mathbf{r}}{2}\right) \propto Q_A^2(\mathbf{s}) Q_B^2(\mathbf{s}) \left[Q_A^2(\mathbf{s}) \ln\left(1 + \frac{Q_B^2(\mathbf{s})}{m^2}\right) + Q_B^2(\mathbf{s}) \ln\left(1 + \frac{Q_A^2(\mathbf{s})}{m^2}\right)\right]$$

How do we use it in practice ? We follow [Blaizot, Broniowski, Ollitrault, 1405.3572].

$$\rho(\mathbf{s}) = \langle\rho(\mathbf{s})\rangle + \delta\rho(\mathbf{s}), \quad \langle\rho(\mathbf{s})\rangle \gg \delta\rho(\mathbf{s})$$

(on long wavelengths)

Perturbative expansion of the anisotropy: $\varepsilon_n \equiv \dfrac{\int_{\mathbf{s}} \mathbf{s}^n \rho(\mathbf{s})}{\int_{\mathbf{s}} |\mathbf{s}|^n \rho(\mathbf{s})}$

To first nontrivial order:

$$\langle\varepsilon_2\varepsilon_2^*\rangle = \varepsilon_2\{2\}^2 = \sigma^2 + \bar{\varepsilon}_2^2$$

and

$$\varepsilon_3\{2\}^2 = \frac{\int_{\mathbf{s}} |\mathbf{s}|^6\, \xi(\mathbf{s})}{\left(\int_{\mathbf{s}} |\mathbf{s}|^3 \langle\rho(\mathbf{s})\rangle\right)^2}$$

$$\sigma^2 = \frac{\int_{\mathbf{s}} |\mathbf{s}|^4\, \xi(\mathbf{s})}{\left(\int_{\mathbf{s}} |\mathbf{s}|^2 \langle\rho(\mathbf{s})\rangle\right)^2}, \quad \bar{\varepsilon}_2 = \frac{\int_{\mathbf{s}} \mathbf{s}^2 \langle\rho(\mathbf{s})\rangle}{\int_{\mathbf{s}} |\mathbf{s}|^2 \langle\rho(\mathbf{s})\rangle}$$

From the CGC we have:

$$\langle\rho(\mathbf{s})\rangle = \frac{4}{3g^2} Q_A^2(\mathbf{s}) Q_B^2(\mathbf{s})$$

Prefactors from:
[Albacete, Guerrero-Rodriguez, Marquet 1808.00795]

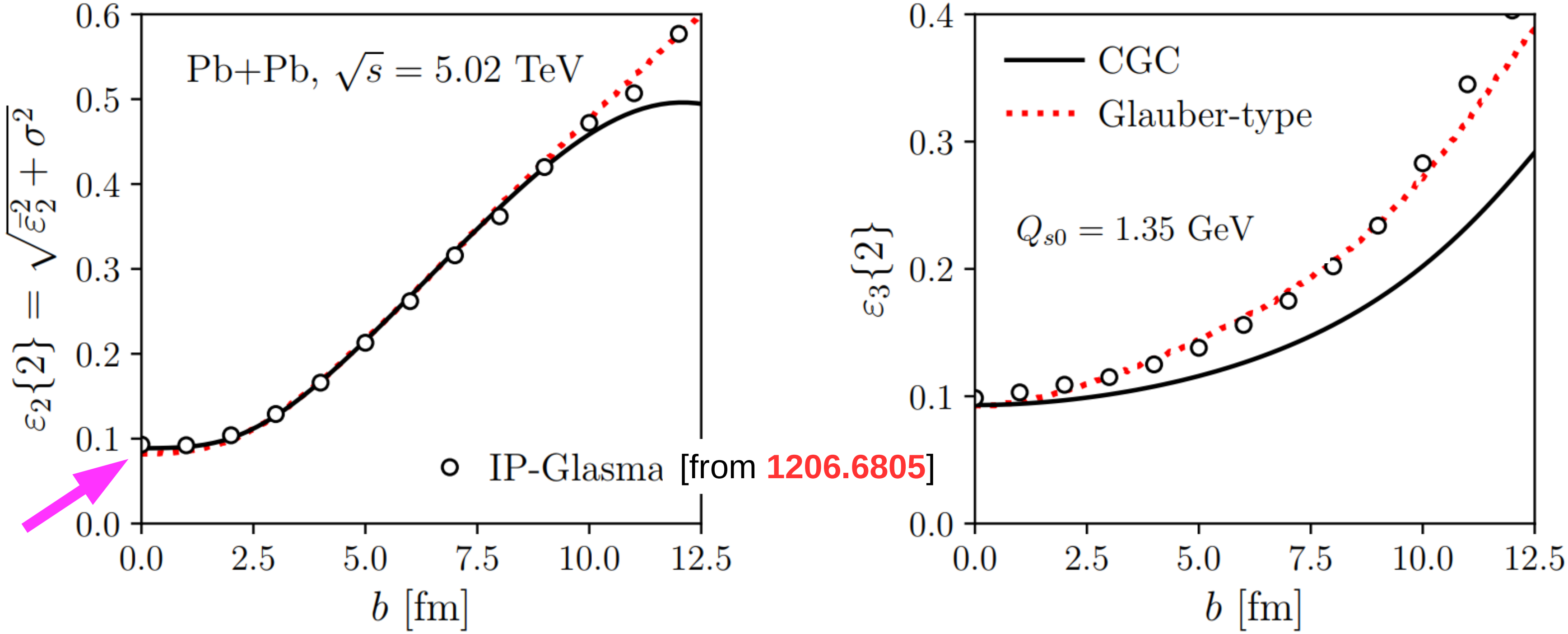
$$\xi(\mathbf{s}) = \frac{16\pi}{9g^4} Q_A^2(\mathbf{s}) Q_B^2(\mathbf{s}) \left(Q_A^2(\mathbf{s}) \ln\left(1 + \frac{Q_B^2(\mathbf{s})}{m^2}\right) + Q_B^2(\mathbf{s}) \ln\left(1 + \frac{Q_A^2(\mathbf{s})}{m^2}\right)\right)$$

Saturation scale proportional to the integrated nuclear density:

$$Q_s^2(\mathbf{s}) = Q_{s0}^2 T(\mathbf{s})/T(\mathbf{0})$$

RESULTS

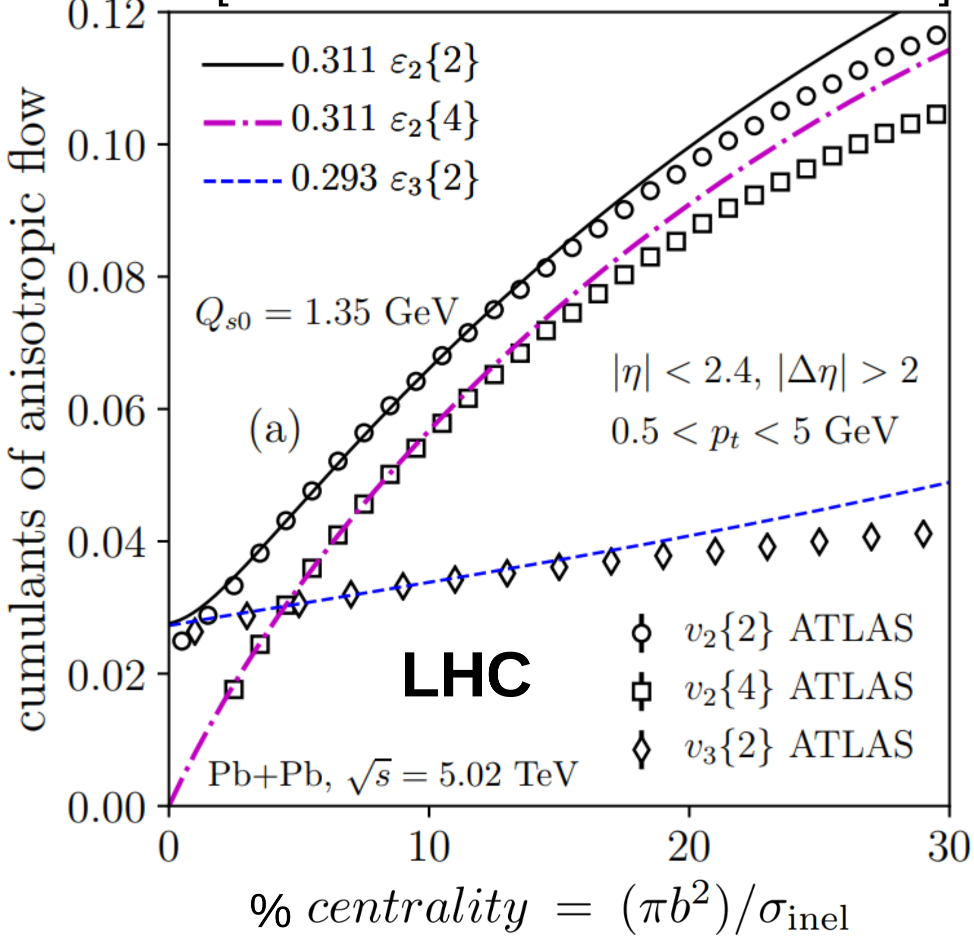
Giacalone, Guerrero-Rodriguez, Luzum, Marquet, Ollitrault
[1902:07168]



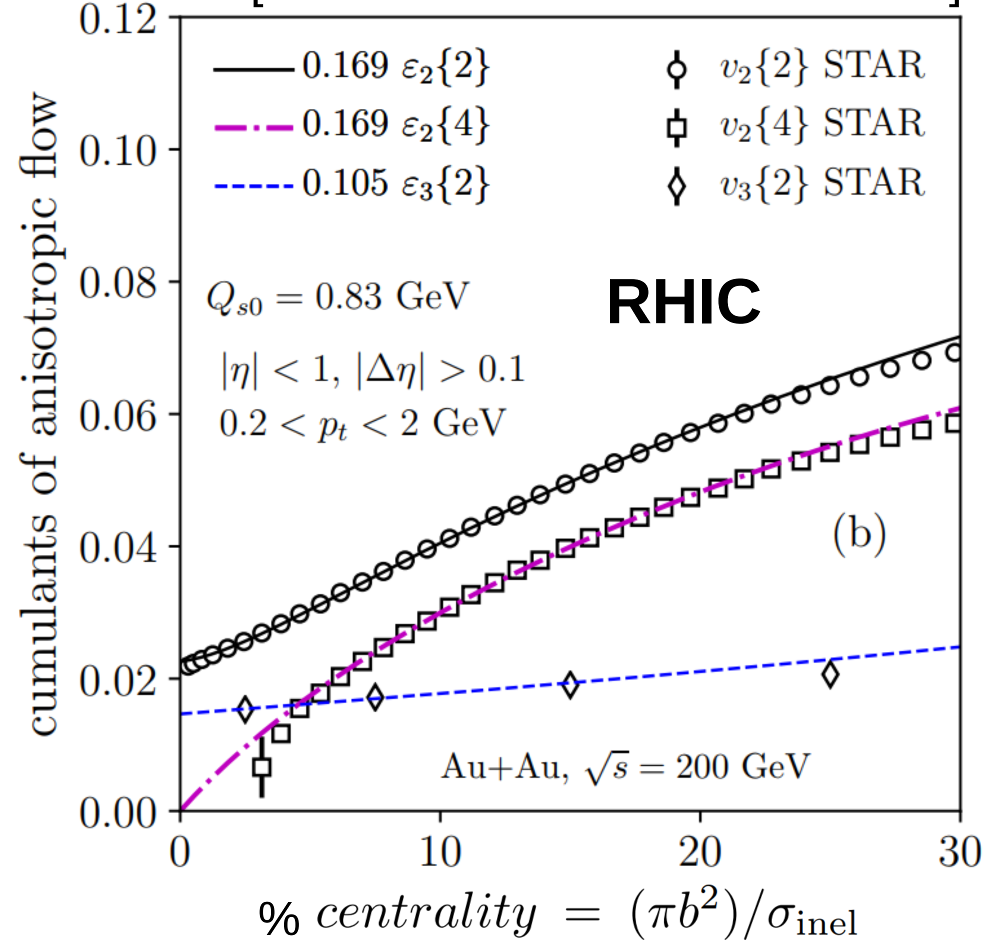
We reproduce the Glauber results. Not a lucky coincidence!

$$\varepsilon_2\{2\}^2 \approx \varepsilon_3\{2\}^2 \approx \frac{\log(Q^2/m^2)}{R^2Q^2} \longrightarrow \varepsilon_2\{2\} \approx 0.1 \text{ for } \begin{array}{l} m = 0.1 \text{ GeV} \\ Q = 1 \text{ GeV} \\ R = 5 \text{ fm} \end{array}$$

[ATLAS Collaboration 1904.04808]



[STAR Collaboration 1505.07812]



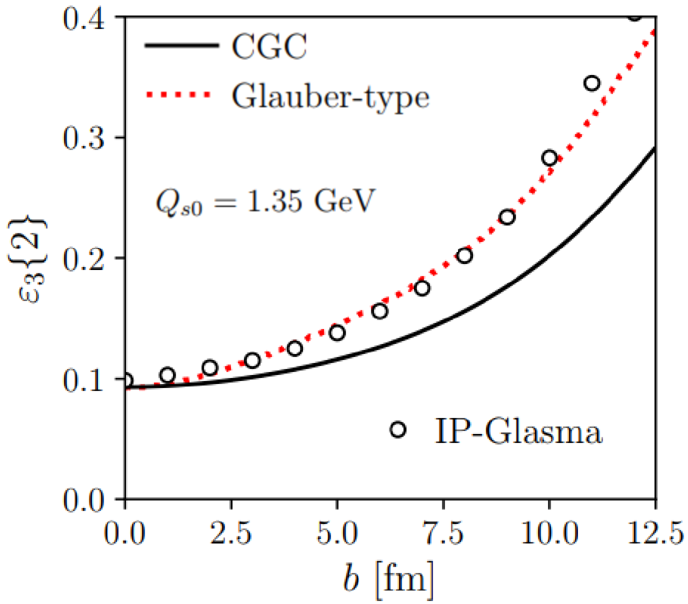
- We use $v_2\{4\} = \kappa_2 \varepsilon_2\{4\} \approx \kappa_2 \bar{\varepsilon}_2$ [Voloshin, Poskanzer, Tang, Wang, 0708.0800]
  $v_2\{4\}$ fixes the response coefficient k_2.

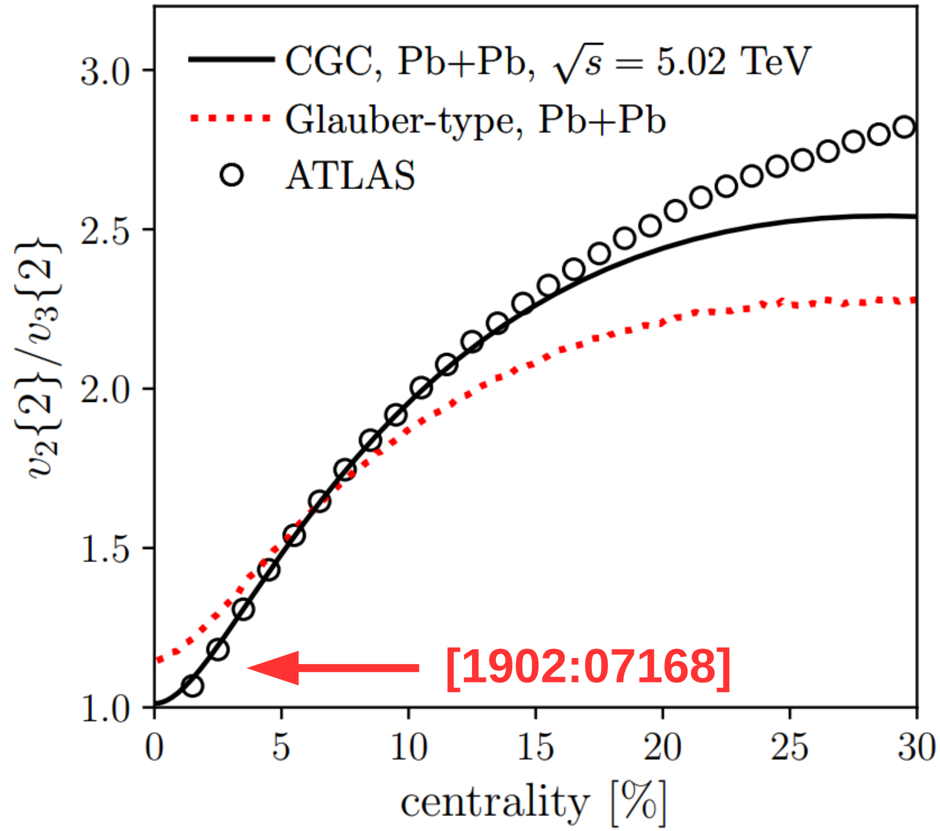
- The splitting between v2{2} and v2{4} is due to fluctuations:

$$Q_s[\text{LHC}] > Q_s[\text{RHIC}] \implies$$ smaller splitting at LHC.

TRIANGULAR FLOW

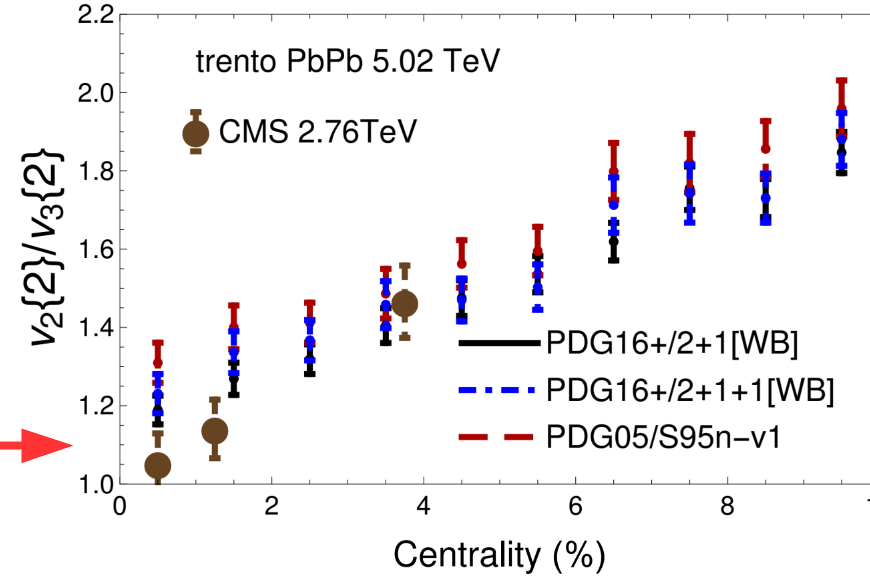
In the CGC, triangular flow grows more mildly than in a Glauber calculation.





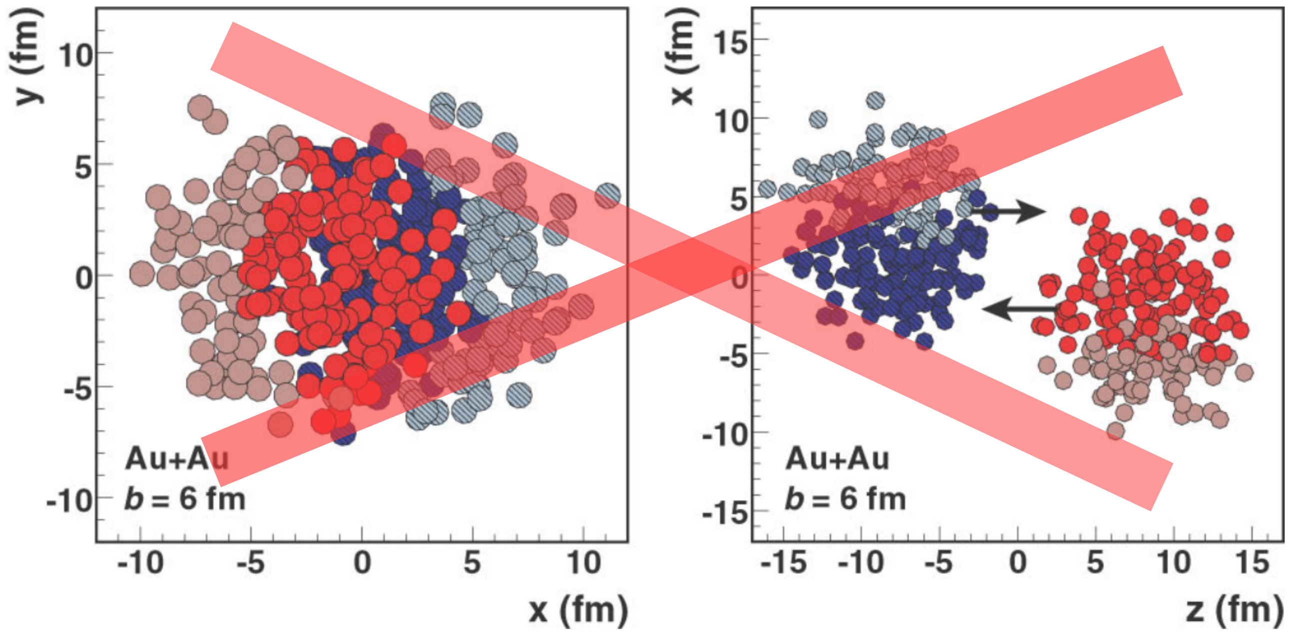
This fixes a longstanding problem of hydro-to-data comparisons:
The ratio $v_2\{2\}/v_3\{2\}$ grows quickly with centrality.

e.g. [Shen, Qiu, Heinz, 1502.04636]



[Alba, Mantovani Sarti, Noronha, Noronha-Hostler, Parotto, Portillo Vasquez, Ratti 1711.05207]

CONCLUSION: A NEW PARADIGM FOR FLUCTUATIONS.



[Miller, Reygers, Sanders, Steinberg **nucl-ex/0701025**]

We free the description from the Glauber Monte Carlo Ansatz:

- No random sampling of nucleons.

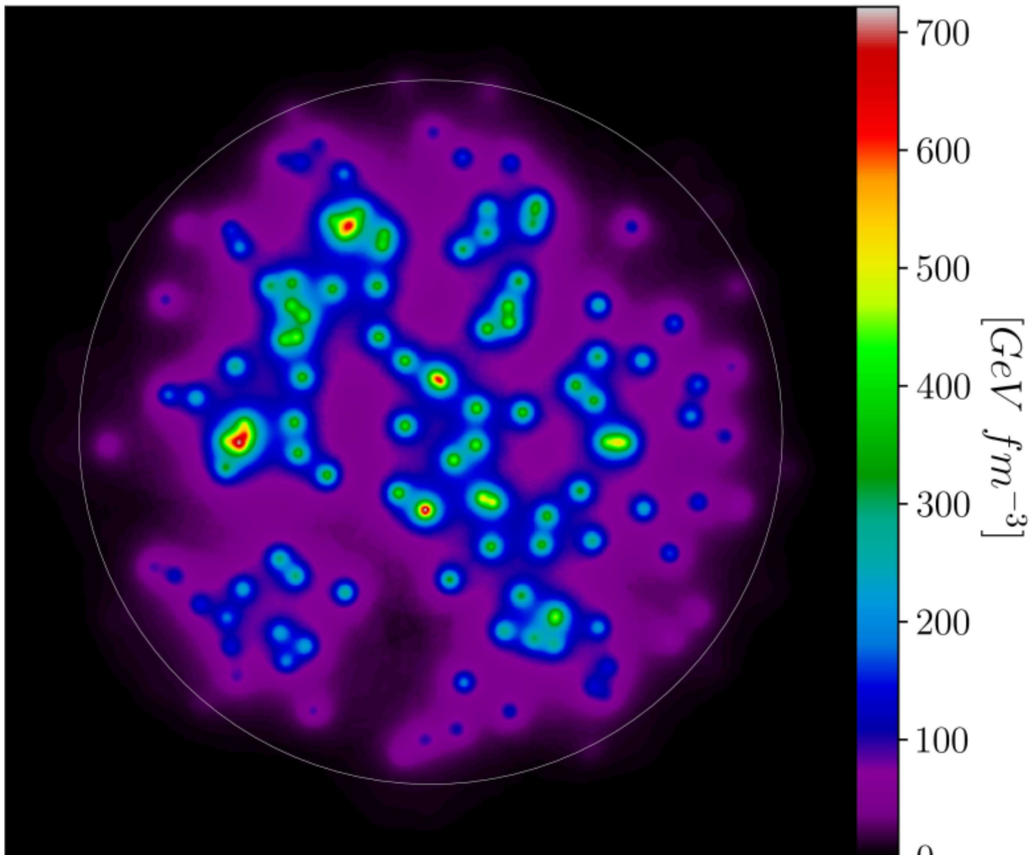
- No ad hoc prescriptions about the deposition of energy.

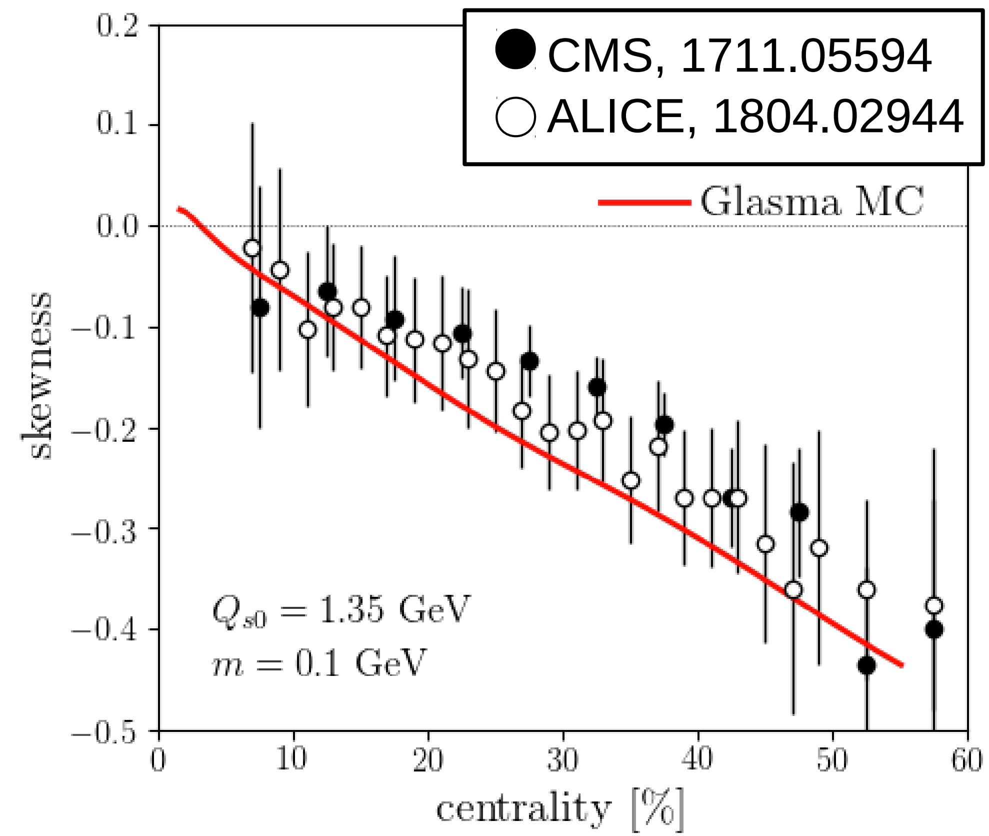
- Nonperturbative physics only through the mass parameter.

BACKUP

Beyond small-and-local fluctuation approximation?

Event-by-event energy density profiles.
Non-Gaussianity generated by positive sign
of energy at a transverse point.

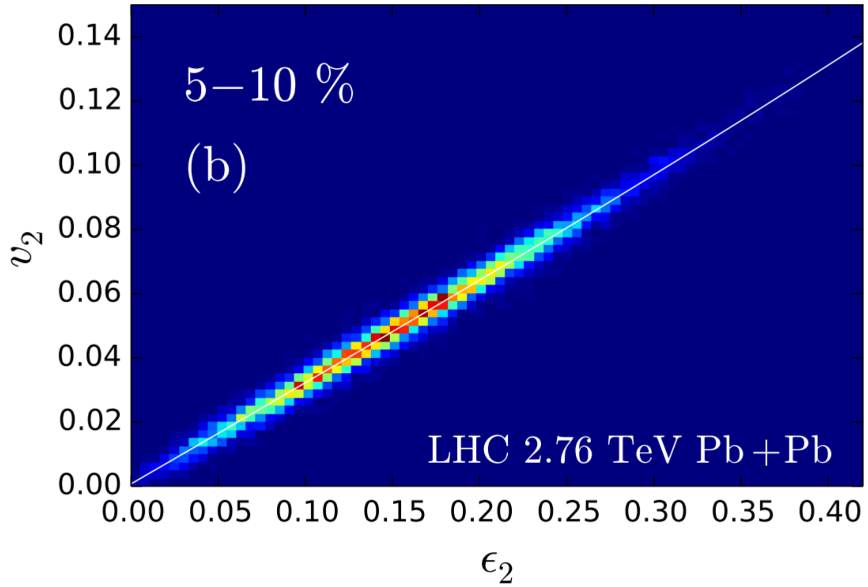




[Gelis, Giacalone, Guerrero-Rodriguez, Marquet, Ollitrault, **to appear**]

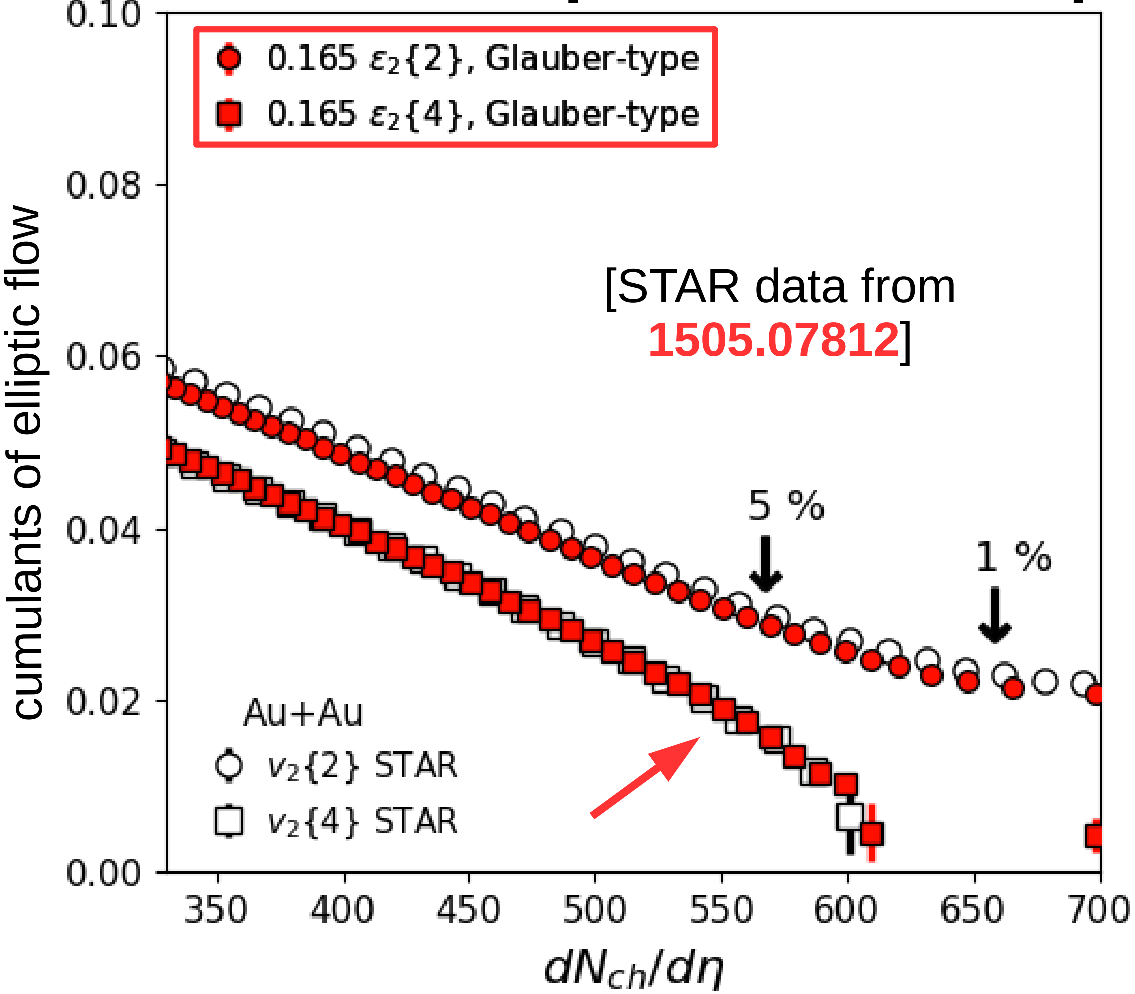
Stay tuned (or ask for details).

[Niemi, Eskola, Patelaainen, arXiv 1505:02677]

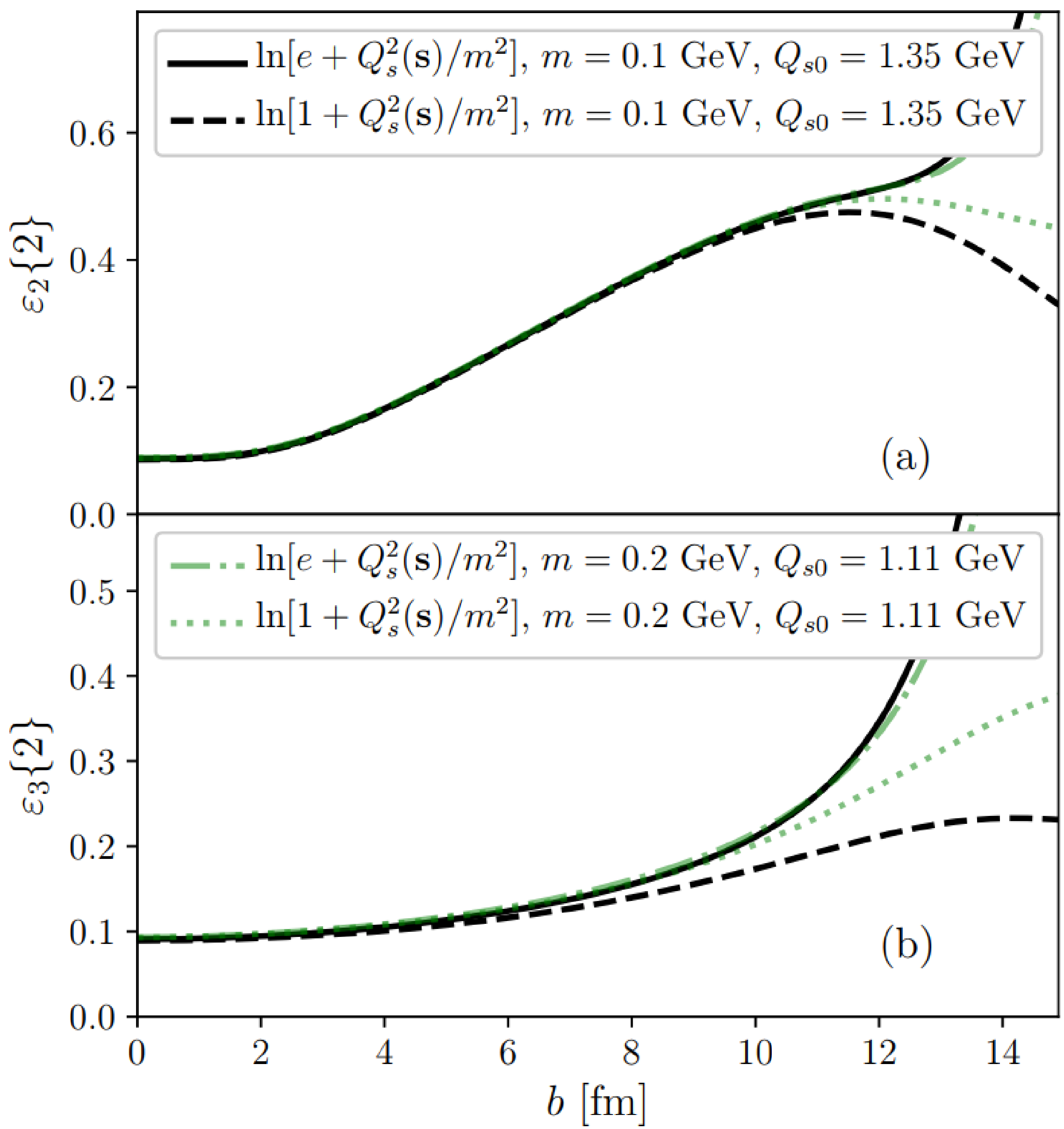


Calculations suggest a linear relation.

[Giacalone, 1811.03959]

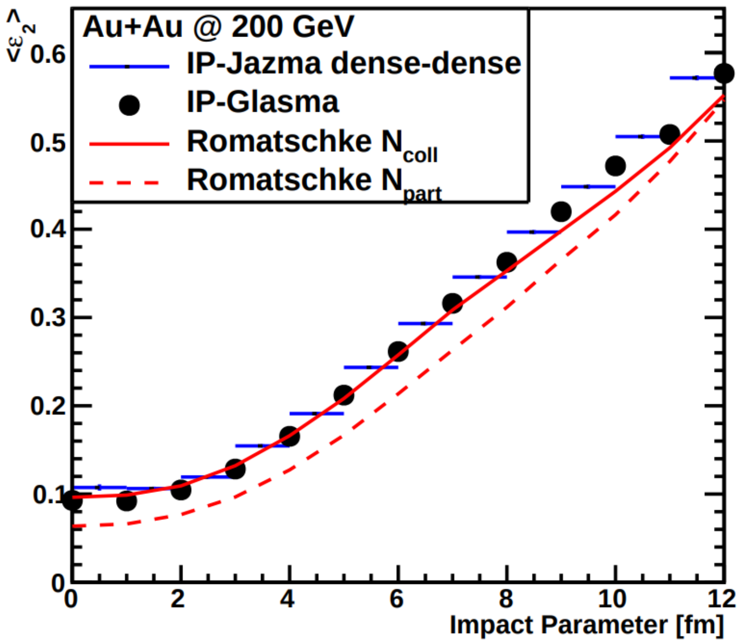


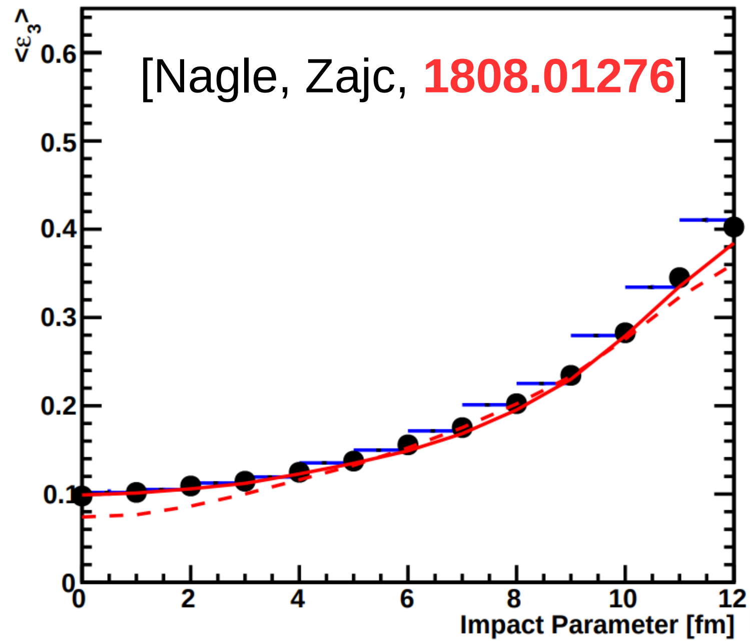
How robust is the formalism ?

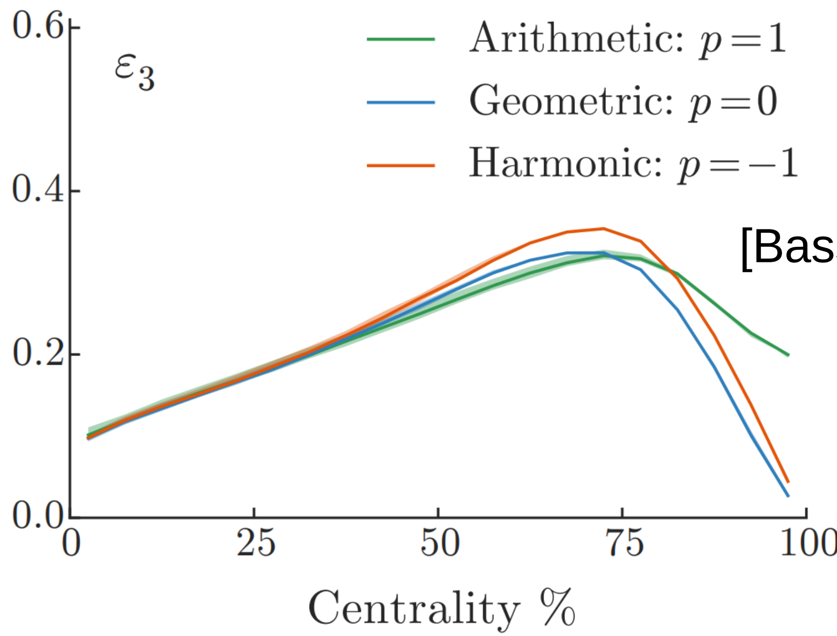


Breaks down at about b=12 fm.

All MC-Glauber-based models have the same eccentricities.







[Bass, Bernhard, Moreland 1412.4708]

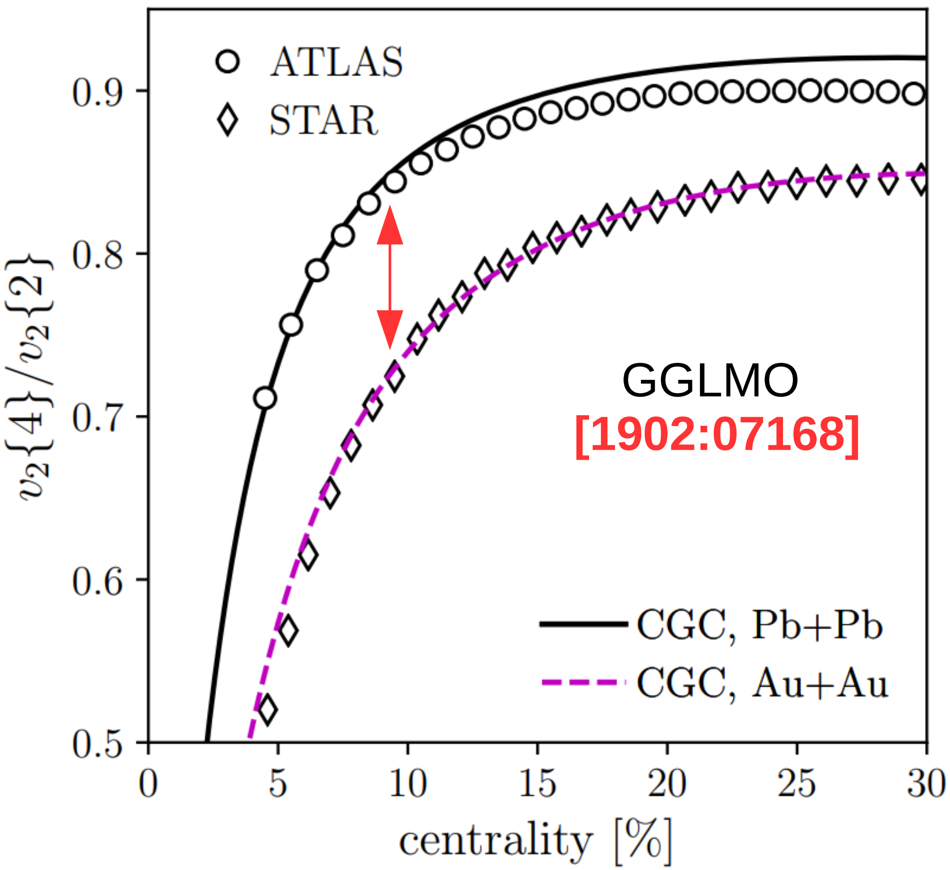
ELLIPTIC FLOW FLUCTUATIONS

Fluctuations of elliptic flow produce the **splitting between $v_2\{2\}$ and $v_2\{4\}$.**
Experimental data indicate that fluctuations are larger at RHIC energy.

Energy dependence of the saturation scale from fits of DIS data:

$$\frac{Q_s^2(x_1)}{Q_s^2(x_2)} = \left(\frac{\sqrt{s_1}}{\sqrt{s_2}}\right)^{0.28}$$

See e.g. [Albacete, Marquet, **1401.4866**]



Increase of **~1.6** from RHIC to LHC energy.

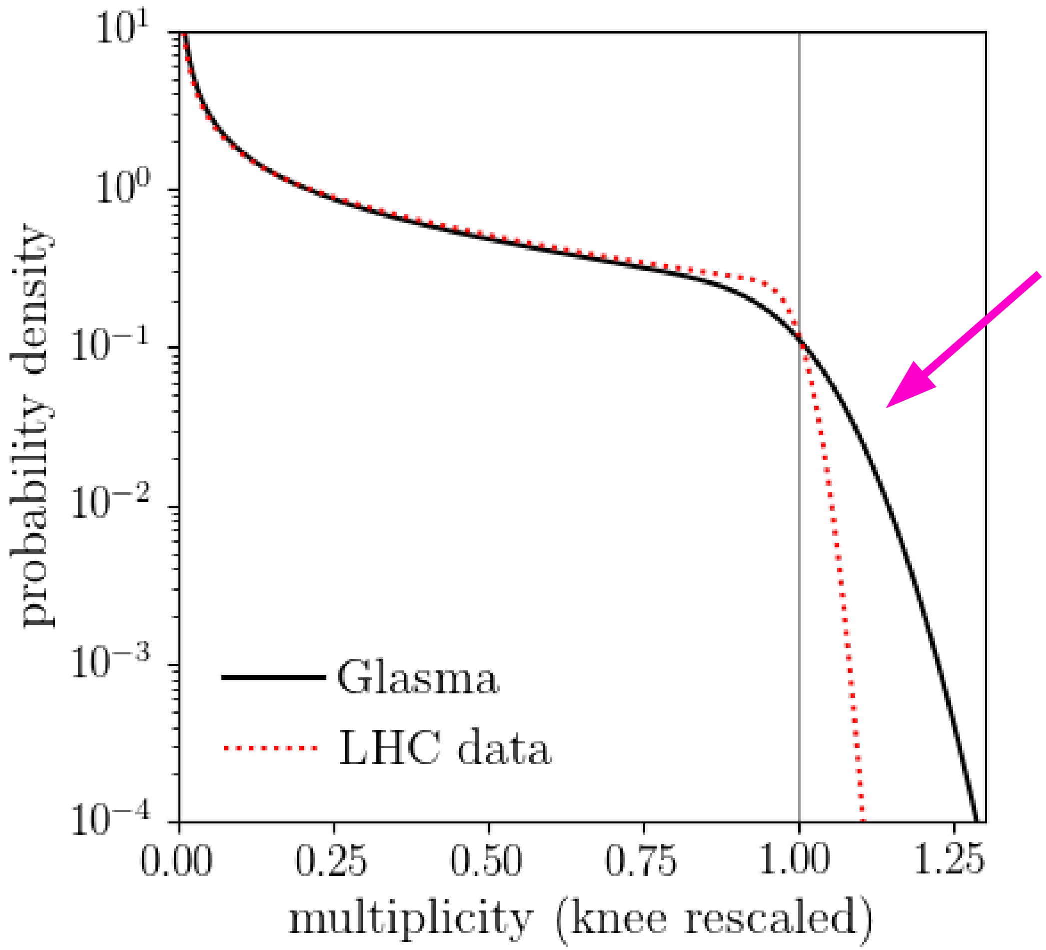
Compatible with the evolution of Q_{s0} found in our fit of anisotropic flow data:

Q_{s0} (LHC) ~ 1.3 GeV
Q_{s0} (RHIC) ~ 0.8 GeV

Very transparent physical explanation!

NB: the Glauber-type calculation does not make any specific predictions for this ratio.

The fluctuations of the primordial energy density are too large compared to the fluctuations of the final-state multiplicity observed at LHC.



Need full pre-equilibrium dynamics. Nontrivial task.