

# Flow angle and flow magnitude factorization breaking

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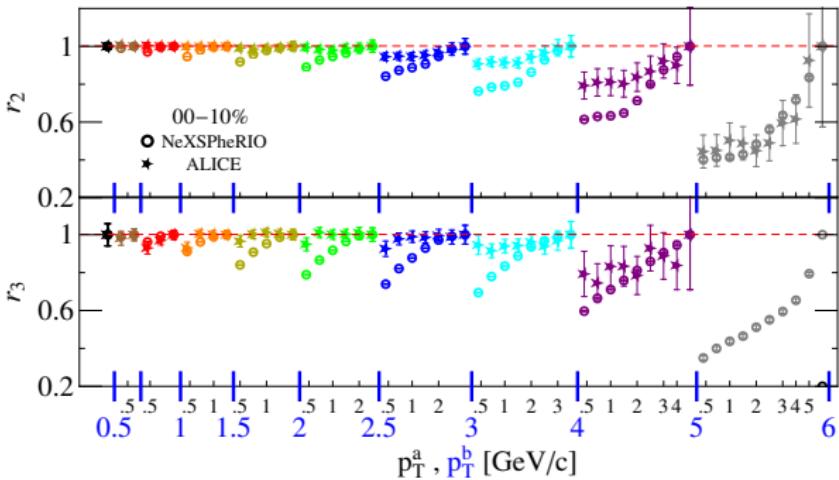
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with: W. Broniowski arXiv: 1711.03325  
and H. Mehrabpour *in preparation*



# Flow factorization breaking

$$r_n(p_a, p_b) = \frac{< v_n(p_a) v_n^*(p_b) >}{\sqrt{< v_n^2(p_a) >} \sqrt{< v_n^2(p_b) >}}$$



Gardim, Grassi, Luzum, Ollitrault: 1211.0989

small decorrelation of flow vectors due to fluctuations

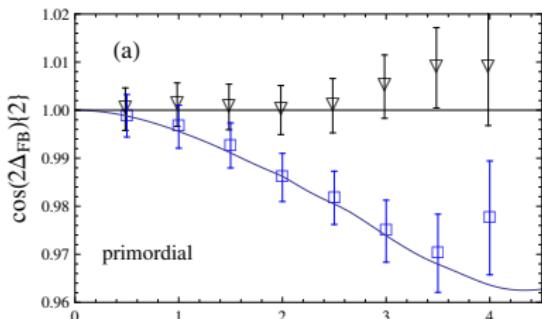
# Flow factorization breaking in rapidity

## forward-backward decorrelation

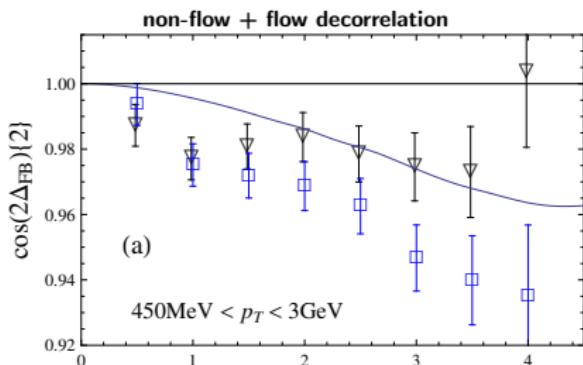
$$\frac{\langle v_2(F)v_2^*(B) \rangle}{\sqrt{\langle v_2^2(F) \rangle} \sqrt{\langle v_2^2(B) \rangle}}$$

PB, Broniowski, Moreira 1011.3354

events with decorrelation versus events without decorrelation



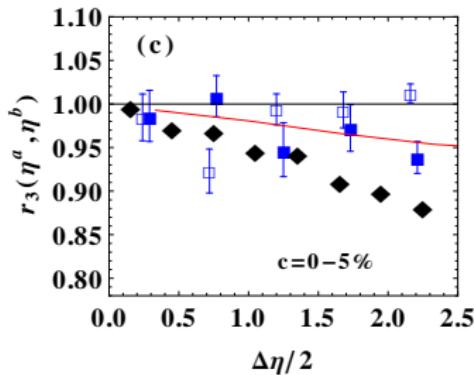
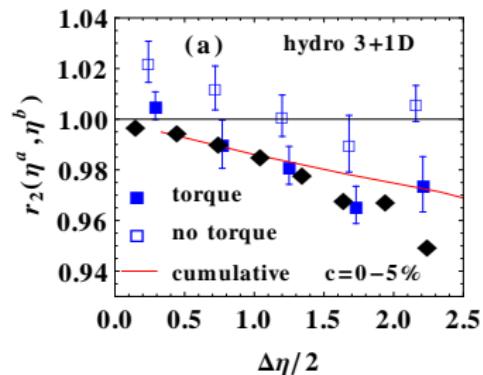
strong non-flow effects



# 3-bin measure of event-plane decorrelation (CMS)

$$r_n(\eta) = \frac{< q_n(-\eta) q_n^*(\eta_{ref}) >}{< q_n(\eta) q_n^*(\eta_{ref}) >}$$

only pairs with large rapidity gap  $\Delta\eta = \pm\eta - \eta_{ref}$

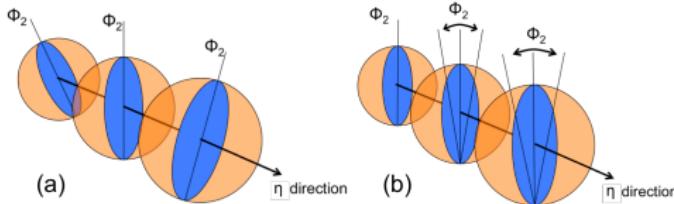


- nonflow under control
- torque effect seen in the CMS data
- semiquantitative agreement
- small factorization breaking observed

$$\text{F-slope} \quad r_n(\eta) \simeq 1 - 2F_n^\eta \eta$$

# Flow magnitude decorrelation + Twist angle

$$r_2(\eta) = 1 - 2F_n\eta = 1 - 2F_n^{\text{asy}}\eta - 2F_n^{\text{twi}}\eta$$



Jia, Huo 1402.6680, 1403.6077

the two can be separated using

3-bin and 4-bin correlators ATLAS 1709.02301

$$R_n(\eta) = \frac{\langle q_n(-\eta_{\text{ref}})q_n^*(\eta)q_n(-\eta)q_n^*(\eta_{\text{ref}}) \rangle}{\langle q_n(-\eta_{\text{ref}})q_n^*(-\eta)q_n(\eta)q_n^*(\eta_{\text{ref}}) \rangle} \simeq 1 - 2F_{n,2}^{\text{twi}}\eta$$

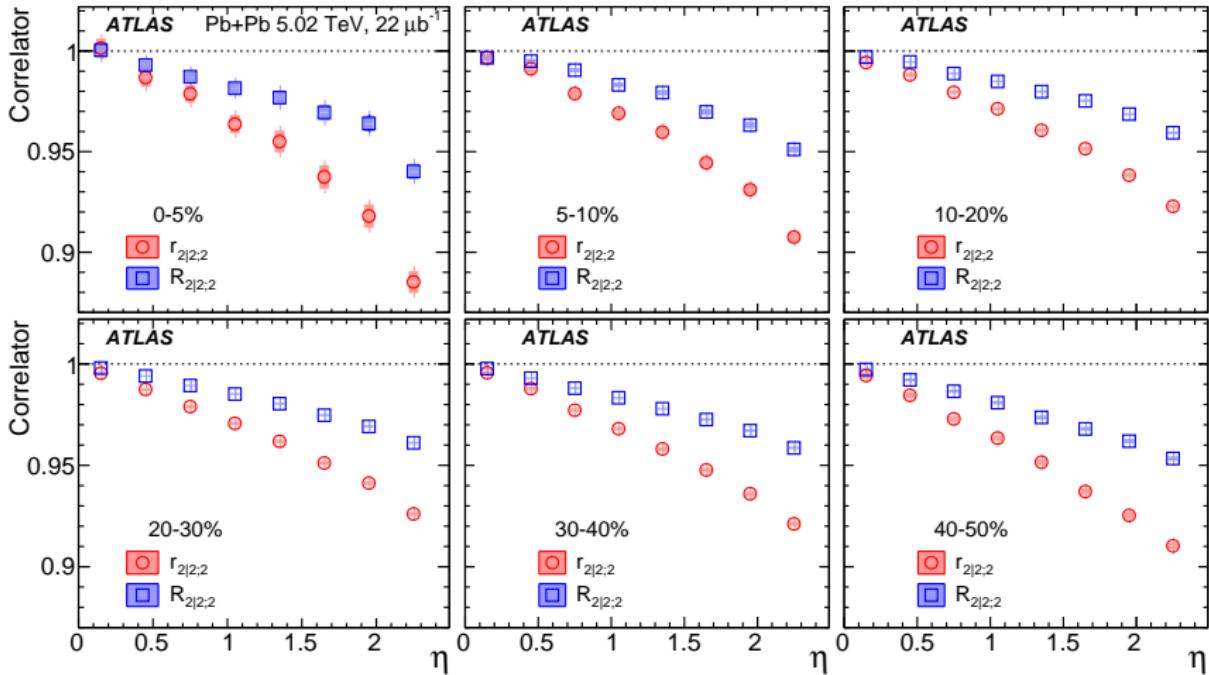
flow angle decorrelation

$$r_{n,2}(\eta) = \frac{\langle q(-\eta)^2 q^*(\eta_{\text{ref}})^2 \rangle}{\langle q(\eta)^2 q^*(\eta_{\text{ref}})^2 \rangle} \simeq 1 - 2F_{n,2}^{\text{asy}}\eta - 2F_{n,2}^{\text{twi}}\eta$$

flow angle+flow magnitude decorrelation

$(r_{n,1}(\eta)$  first measured by CMS 1503.01692)

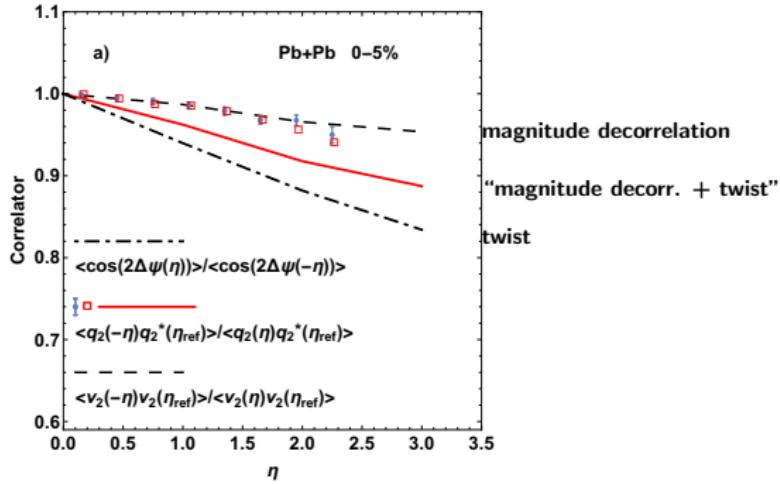
# Observation of separate flow magnitude decorrelation and twist angle



$$r_2 \simeq 1 - 2F_2^{\text{asy}}\eta - 2F_2^{\text{twi}}\eta$$

$$R_2 \simeq 1 - 2F_2^{\text{twi}}\eta$$

# twist angle and flow magnitude decorrelation 3+1D hydro model

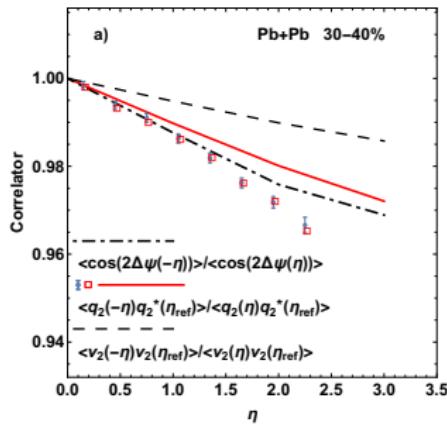


**surprising result:** “inverted hierarchy”

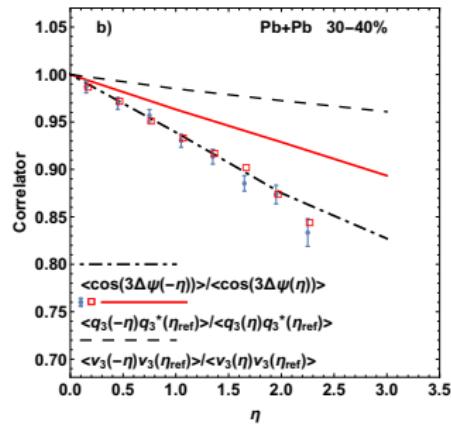
**magnitude decorr. < “magnitude decorrelation + twist” < twist**

## elliptic versus triangular

$V_2$

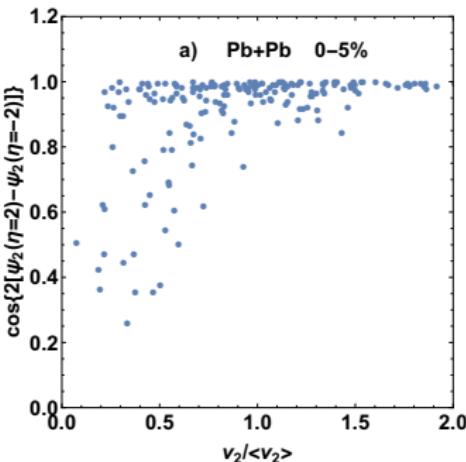


$V_3$



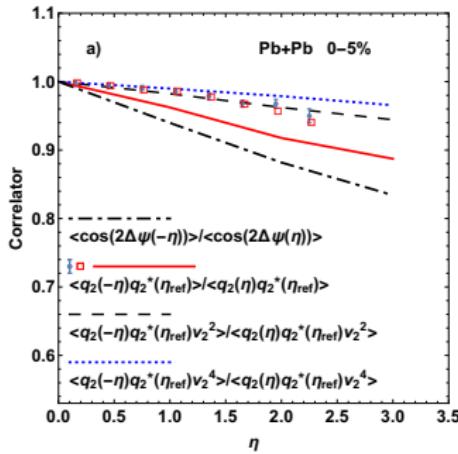
the “inverted hierarchy” effect is stronger for  $v_3$   
triangular flow - fluctuation dominated

## Correlation between flow magnitude and twist angle



- ▶ **strong** correlation between flow magnitude and twist angle
- ▶ events with large flow have smaller twist angle
- ▶ twist angle measure  $\langle \cos(\Delta\Psi_2) \rangle \propto (v_2)^0$   
“magnitude decorr.+twist”  $\langle q_2(\eta)q_2^*(\eta_{ref}) \rangle \propto (v_2)^2$
- ▶ different weighting by ( $v_2$ ) powers explains “inverted hierarchy”

## Correlators weighted by powers of $v_n$

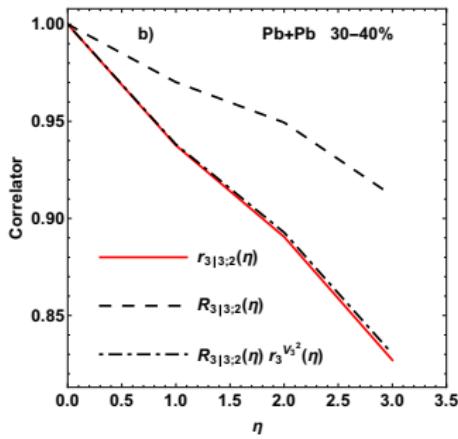


- hierarchy of correlators consistent with expectations

$$\frac{\langle q_n(-\eta) q_n^*(\eta_{ref}) \rangle}{\langle q_n(\eta) q_n^*(\eta_{ref}) \rangle} < \frac{\langle q_n(-\eta) q_n^*(\eta_{ref}) v_n^2 \rangle}{\langle q_n(\eta) q_n^*(\eta_{ref}) v_n^2 \rangle} < \frac{\langle q_n(-\eta) q_n^*(\eta_{ref}) v_n^4 \rangle}{\langle q_n(\eta) q_n^*(\eta_{ref}) v_n^4 \rangle}$$

- the correlation between flow magnitude and twist can be measured experimentally

# Factorization of factorization breaking



- ▶ flow magnitude factorization breaking (square)

$$r_n^2(\eta) = \frac{\langle v_n^2(-\eta) v_n^2(\eta_{ref}) \rangle}{\langle v_n^2(\eta) v_n^2(\eta_{ref}) \rangle}$$

- ▶ twist angle factorization breaking (weighted with  $v_n^4$ )

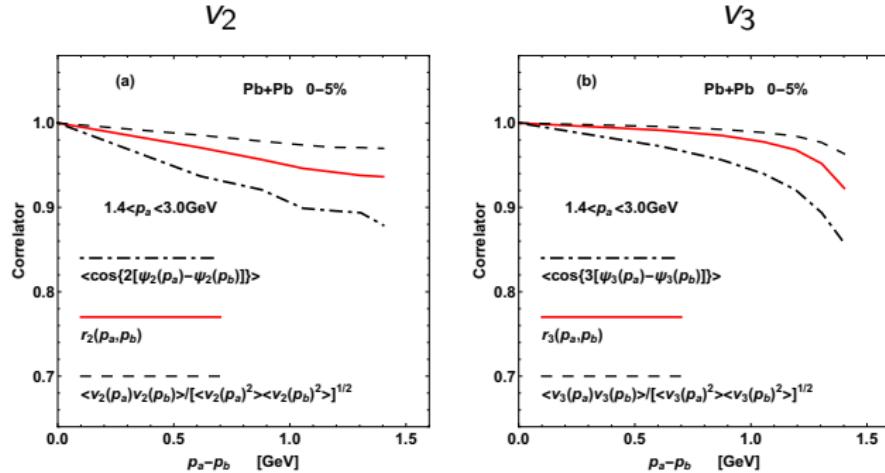
$$R_{n|n;2}(\eta) = \frac{\langle q_n(-\eta_{ref}) q_n^*(\eta) q_n(-\eta) q_n^*(\eta_{ref}) \rangle}{\langle q_n(-\eta_{ref}) q_n^*(-\eta) q_n(\eta) q_n^*(\eta_{ref}) \rangle}$$

- ▶ angle+magnitude factorization breaking

$$r_{n|n;2}(\eta) = \frac{\langle q_n(-\eta)^2 q_n^*(\eta_{ref})^2 \rangle}{\langle q_n(\eta)^2 q_n^*(\eta_{ref})^2 \rangle}$$

(angle+magnitude f. b.)  $\simeq$  (twist angle f. b.)(flow magnitude f. b.)  
measured by ATLAS

# flow magnitude and flow angle decorrelation $p_T$ factorization breaking

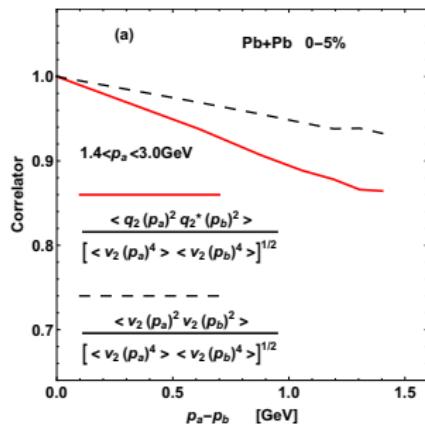


PB : 1808.09840

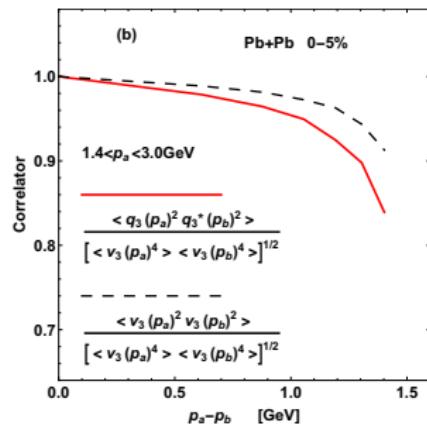
Similar effect for  $p_T$  and rapidity factorization breaking

# Measuring flow magnitude and flow angle decorrelation $p_T$ factorization breaking

$V_2$



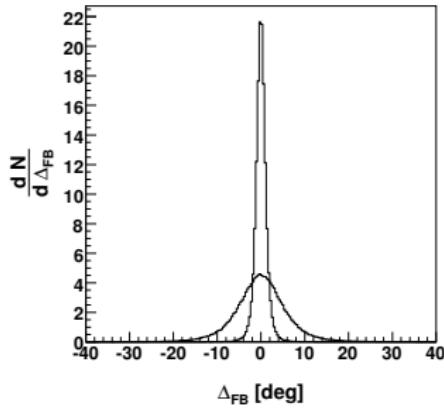
$V_3$



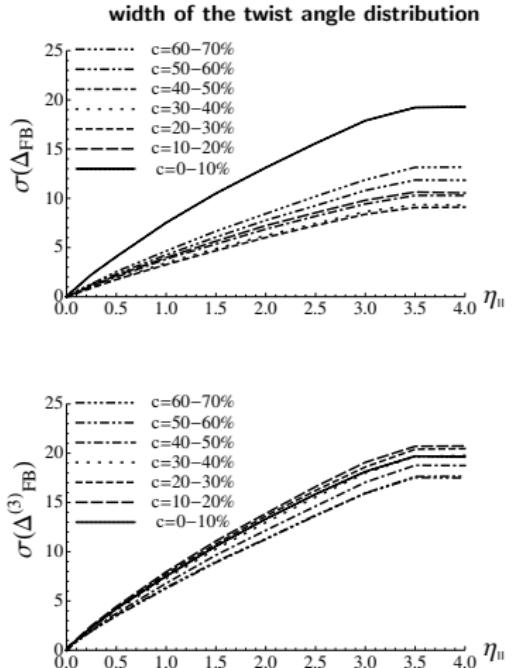
Flow magnitude and flow angle decorrelation in  $p_T$   
can be separately measured in experiment

# Twist angle distribution - Glauber model

$$\Psi_2(\eta) - \Psi_2(-\eta), \quad \Delta\eta = 1, 5$$

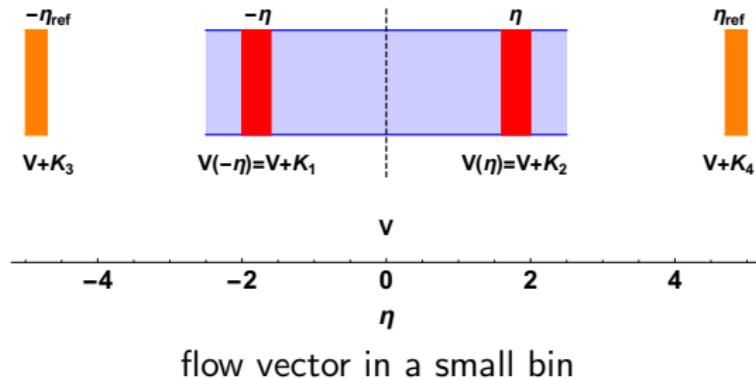


- very forward (backward), maximal decorrelation
- in between, intermediate
- linear around  $\eta = 0$



- $n = 2$ , largest decorrelation for central collisions
- $n = 3$ , similar decorrelation for all centralities

## Simple model of flow factorization breaking



$$Q_i = V_i + K_i$$

has a random component + global flow

random components correlated (depending on bin separation)

random component independent from global flow constraints on global flow, flow moments in small bins

PB, H. Mehrabpour

# Predicts flow angle, flow magnitude, and flow vector decorrelation

decorrelation of flow vectors

$$\frac{\langle q_n^m(\eta)q_n^m * (-\eta) \rangle}{\sqrt{\langle q_n^m(\eta)q_n^m * (\eta) \rangle \langle q_n^m(-\eta)q_n^m * (-\eta) \rangle}} \simeq 1 - 2m \frac{\langle \delta^2 \rangle}{\langle v_n^2 \rangle}$$

decorrelation of flow angles

$$\frac{\langle v_n^{2m} \cos(mn(\Psi_n(\eta) - \Psi_n(-\eta))) \rangle}{\langle v_n^{2m} \rangle} \simeq 1 - m \frac{\langle \delta^2 \rangle}{\langle v_n^2 \rangle}$$

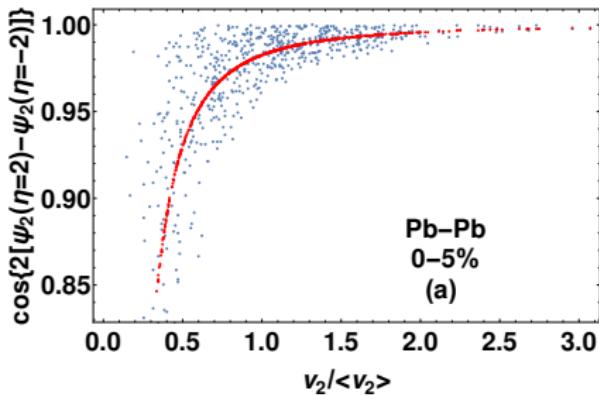
decorrelation of flow magnitudes

$$\frac{\langle v_n^m(\eta)v_n^m(-\eta) \rangle}{\sqrt{\langle v_n^{2m}(\eta) \rangle \langle v_n^{2m}(-\eta) \rangle}} \simeq 1 - m \frac{\langle \delta^2 \rangle}{\langle v_n^2 \rangle}$$

$$\delta = K(\eta) - K(-\eta)$$

increases with bin separation

## global flow - twist angle anticorrelation

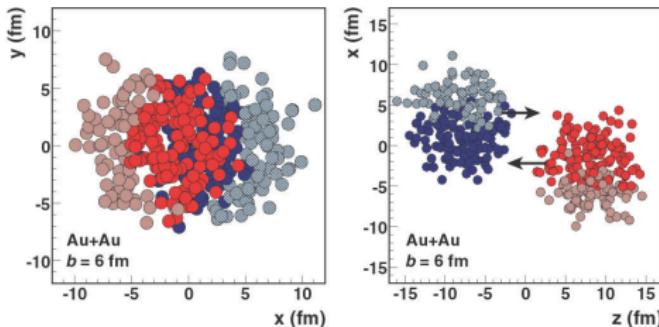


$$\cos(n(\Psi_n(\eta) - \Psi_n(-\eta))) \simeq 1 - \frac{\delta_n^2}{V_n^2}$$

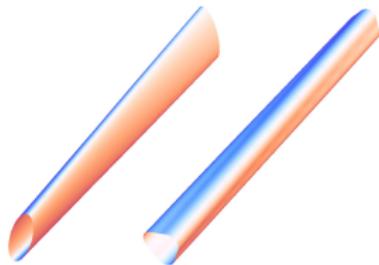
explains the relation between twist angle and global flow magnitude  
explains differences between  $v_2$  and  $v_3$  (smaller global flow)  
**also : explains qualitatively 3- and 4-bin correlators**  
**explains relation between correlators of different powers of flow**

- ▶ Longitudinal flow decorrelation - **angle+magnitude has been measured by ATLAS**
- ▶ Additional event-by-event correlations  
Large flow magnitude  $\leftrightarrow$  Small twist angle  
Small flow magnitude  $\leftrightarrow$  Large twist angle  
(cannot be easily separated)
- ▶ Factorization breaking in  $p_T$  - **flow+magnitude can be measured**
- ▶ simple random model of flow decorrelation  
explains qualitatively separate flow magnitude and flow angle decorrelation

## Twisted event-plane angles - torque effect



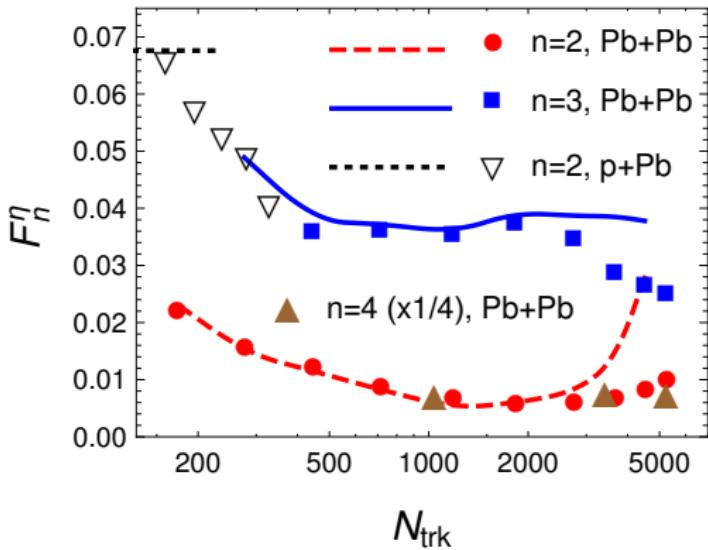
Ann.Rev.Nucl.Part.Sci. 57 (2007) 205



- due to fluctuations of F-B participants
- left-right orientation and angle magnitude are random
- only “smooth” long range twist
- random decorrelations on small scale, difficult to observe

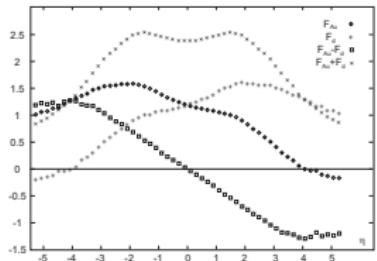
also other models of initial state: fluctuating strings, hybrid models ... → additional fluctuations

$$\text{F slope} \quad (r_n(\eta) \simeq 1 - 2F_n^\eta \eta)$$



- fair description of mid-central collisions
- overestimates decorrelation in central collisions
- $F_4 \simeq 4F_2$

## Matter distribution in space-time rapidity



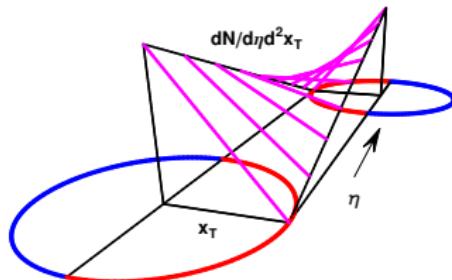
Asymmetric emission functions

(Białas, Czyż, 2005)

signatures of tilt:

charged particles  $v_1$ , tilted source HBT, D meson directed flow

$$\rho(\eta, x, y) \propto f_+(\eta)N_+(x, y) + f_-(\eta)N_-(x, y)$$



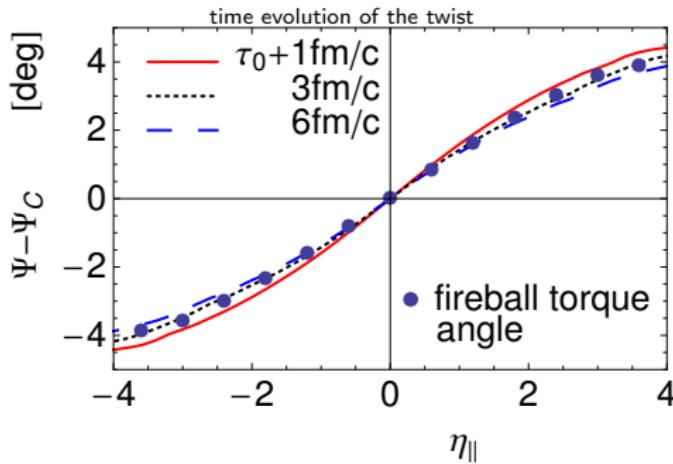
bremsstrahlung Adil Gyulassy, 2005

# One-shot 3+1D hydro evolution (2010)

initial density with a twist

$$s(x, y, \eta) \propto \rho_+(Rx, Ry)f_+(\eta) + \rho_-(R^T x, R^T y)f_-(\eta)$$

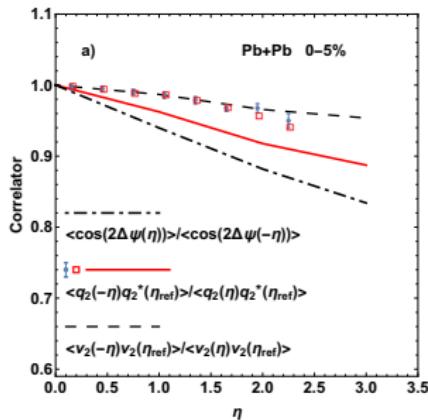
forward (backward) participants rotated in the transverse plane



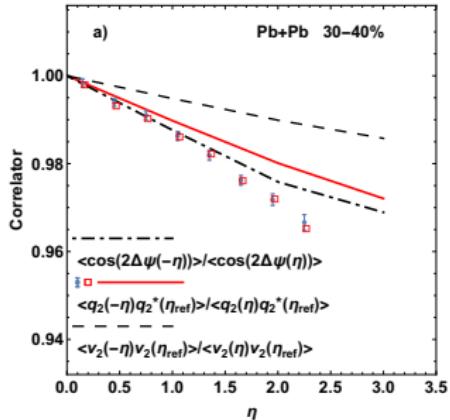
- the twist survives the hydrodynamic evolution

## central versus peripheral

0 – 5%



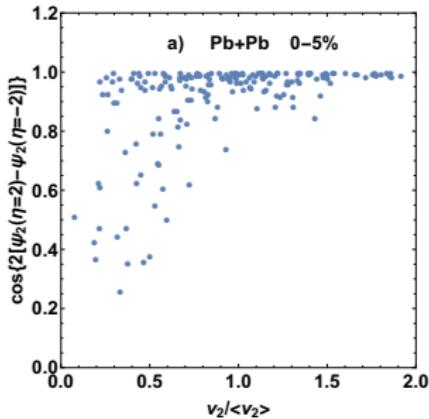
30 – 40%



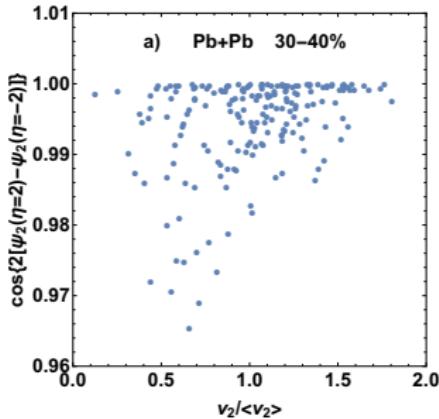
the “inverted hierarchy” effect is stronger in central collisions  
large elliptic flow in semi-central collisions → less fluctuations

## central versus peripheral

0 – 5%



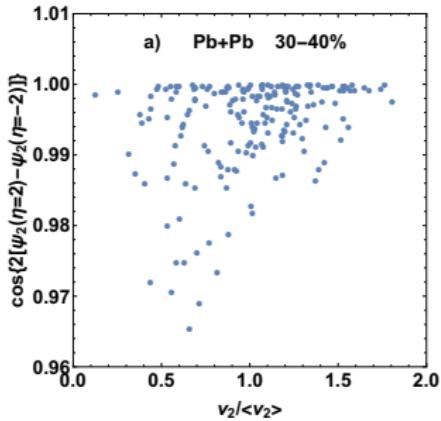
30 – 40%



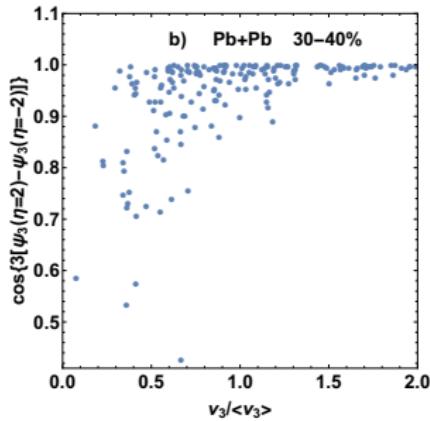
stronger correlation in central collisions

## elliptic versus triangular

$v_2$



$v_3$



stronger correlation for  $v_3$

## similar for triangular flow

