

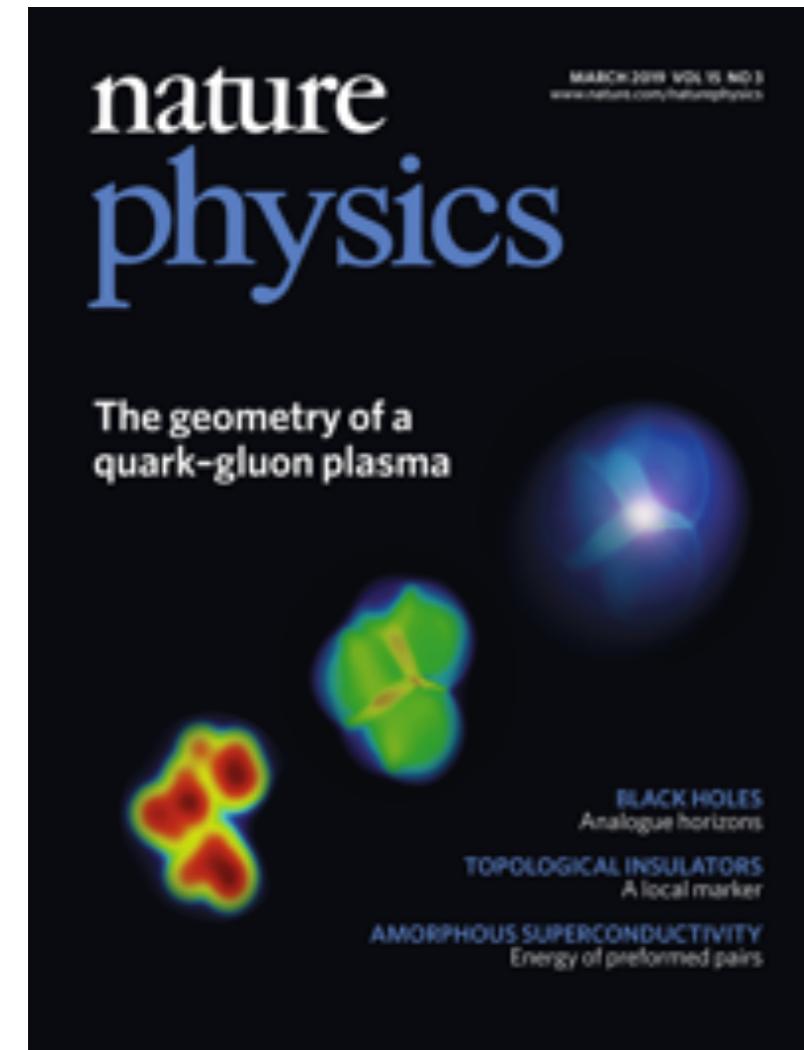
# Observation of collectivity in p+Au, d+Au and ${}^3\text{He}+\text{Au}$ collisions with PHENIX

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# Creation of quark-gluon plasma droplets with three distinct geometries

PHENIX Collaboration\*

Experimental studies of the collisions of heavy nuclei at relativistic energies have established the properties of the quark-gluon plasma (QGP), a state of hot, dense nuclear matter in which quarks and gluons are not bound into hadrons<sup>1–4</sup>. In this state, matter behaves as a nearly inviscid fluid<sup>5</sup> that efficiently translates initial spatial anisotropies into correlated momentum anisotropies among the particles produced, creating a common velocity field pattern known as collective flow. In recent years, comparable momentum anisotropies have been measured in small-system proton-proton (p+p) and proton-nucleus (p+A) collisions, despite expectations that the volume and lifetime of the medium produced would be too small to form a QGP. Here we report on the observation of elliptic and triangular flow patterns of charged particles produced in proton-gold (p+Au), deuteron-gold (d+Au) and helium-gold (<sup>3</sup>He+Au) collisions at a nucleon-nucleon centre-of-mass energy  $\sqrt{s_{NN}} = 200$  GeV. The unique combination of three distinct initial geometries and two flow patterns provides unprecedented model discrimination. Hydrodynamical models, which include the formation of a short-lived QGP droplet, provide the best simultaneous description of these measurements.

anisotropy to momentum correlations generated at the earliest stages of the collision, hence referred to as initial-state momentum correlation models (see refs <sup>12,13</sup> for recent reviews).

A projectile geometry scan utilizing the unique capabilities of the RHIC was proposed in ref. <sup>14</sup> to discriminate between hydrodynamical models that couple to the initial geometry and initial-state momentum correlation models that do not. Varying the collision system from p+Au, to d+Au, to <sup>3</sup>He+Au changes the initial geometry from dominantly circular, to elliptical, to triangular configurations, respectively, as characterized by the second- and third-order spatial eccentricities, which correspond to ellipticity and triangularity, respectively. The  $n$ th order spatial eccentricity of the system,  $\varepsilon_n$ , typically determined from a Monte Carlo (MC) Glauber model of nucleon-nucleon interactions (see for example ref. <sup>15</sup>), can be defined as

$$\varepsilon_n = \frac{\sqrt{\langle r^n \cos(n\phi) \rangle^2 + \langle r^n \sin(n\phi) \rangle^2}}{\langle r^n \rangle} \quad (2)$$

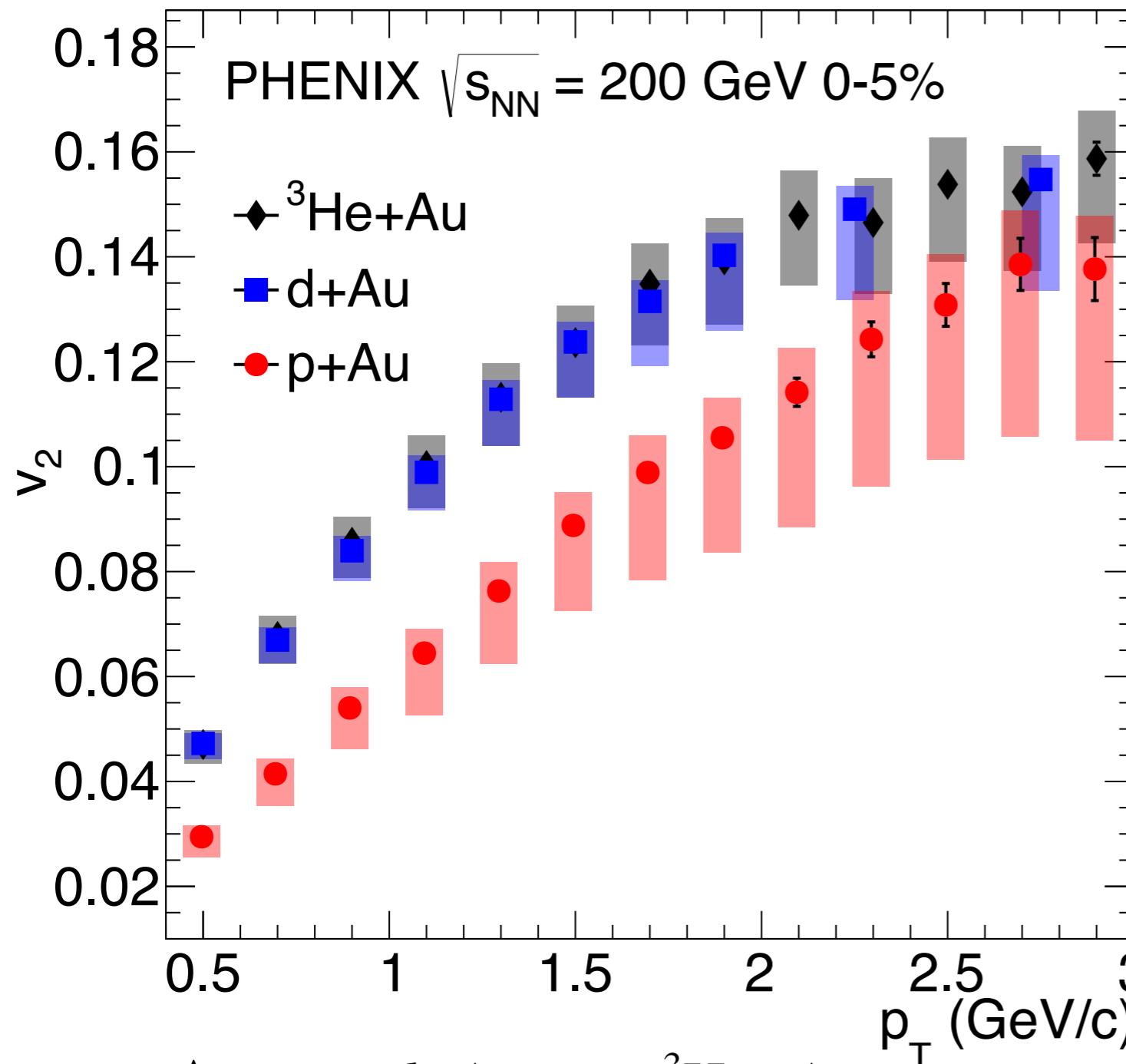
where  $r$  and  $\phi$  are the polar coordinates of participating nucleons<sup>16</sup>. The eccentricity fluctuates event-by-event and is generally dependent on the impact parameter of the collision and the number of

# STRONG EVIDENCE FOR QGP IN SMALL SYSTEMS

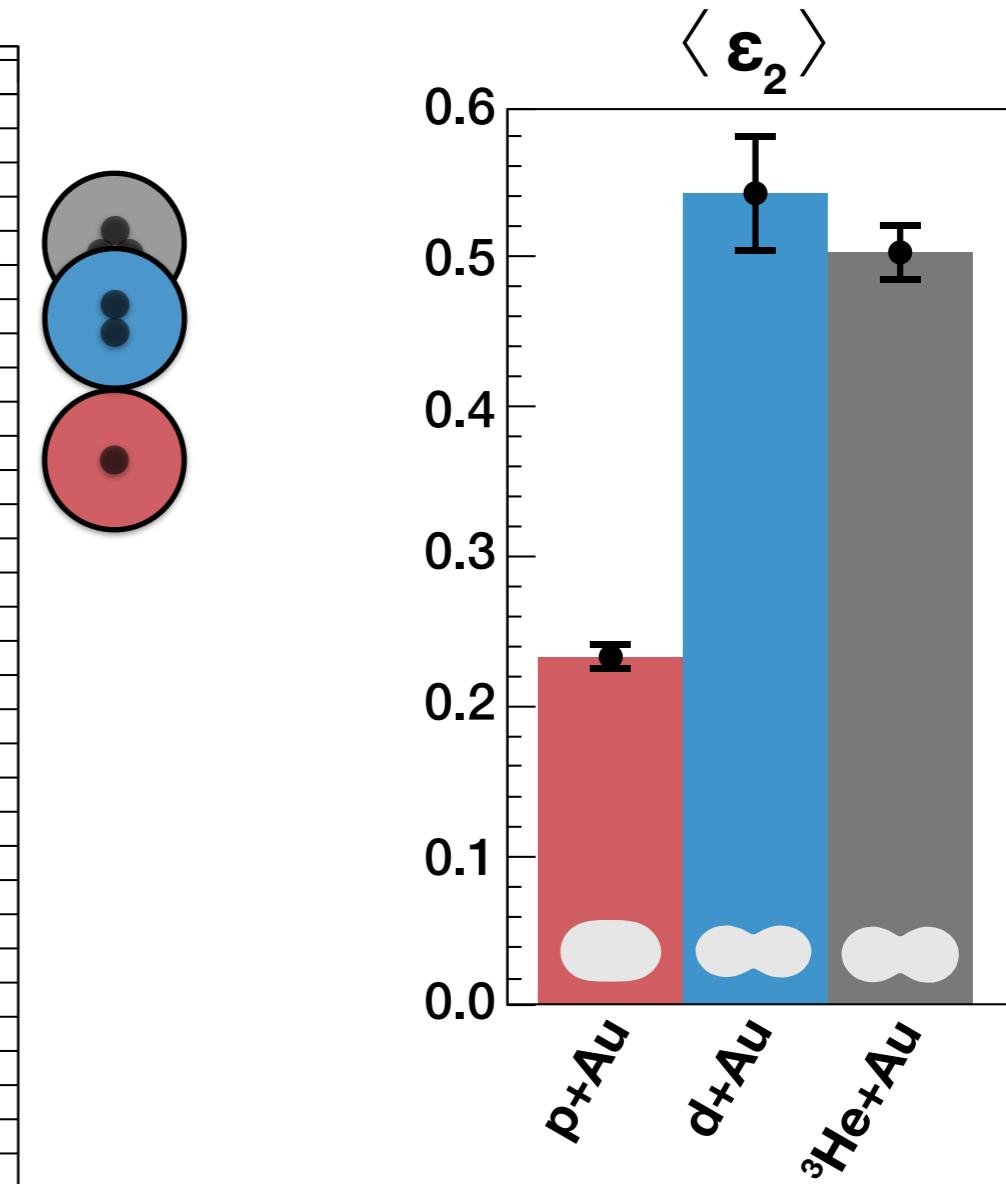


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# $v_2(p_T)$ measurement in small systems



$$v_2^{p+\text{Au}} < v_2^{d+\text{Au}} \approx v_2^{{}^3\text{He}+\text{Au}}$$



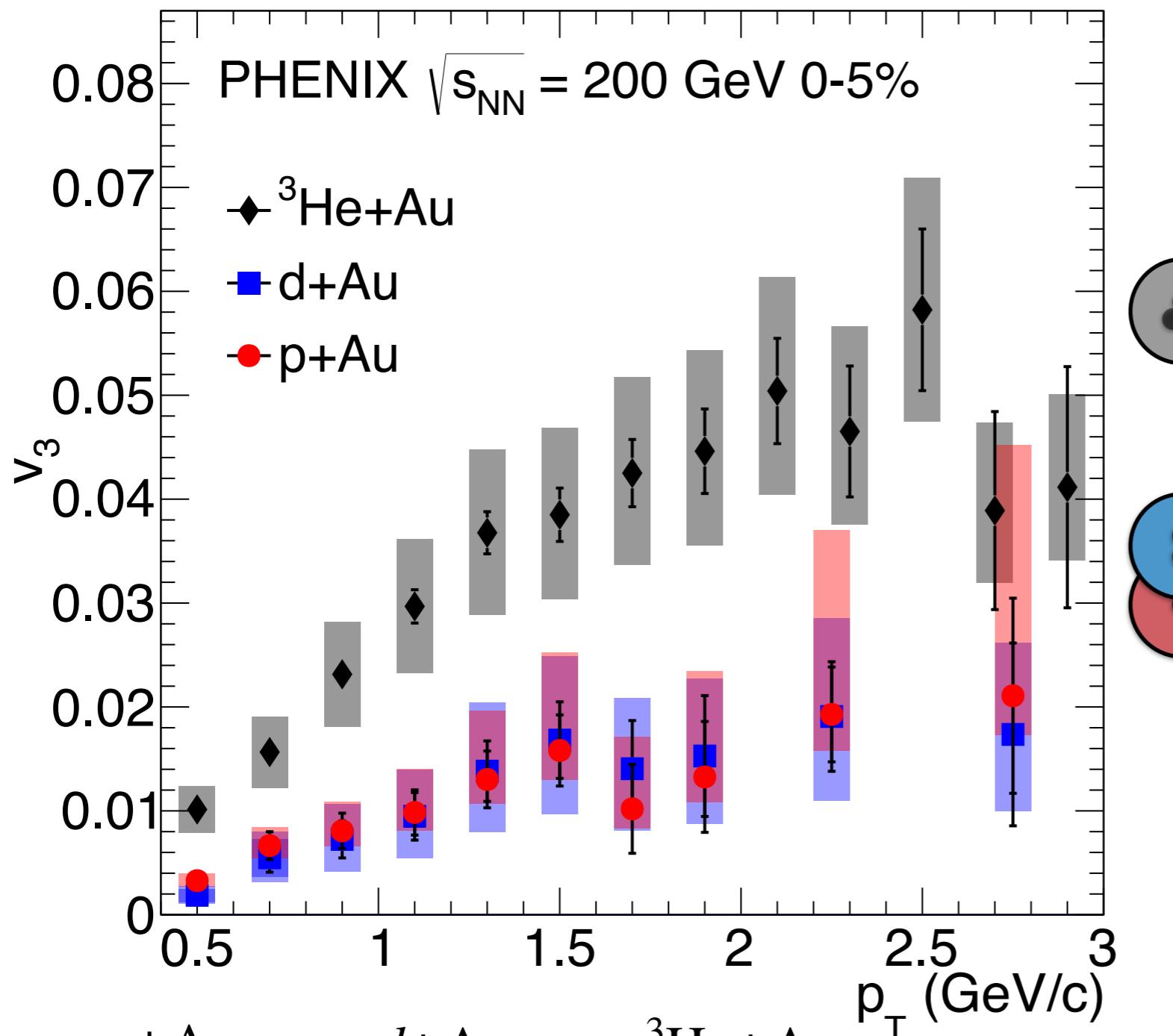
Glauber Model calculations

$$\epsilon_2^{p+\text{Au}} < \epsilon_2^{d+\text{Au}} \approx \epsilon_2^{{}^3\text{He}+\text{Au}}$$



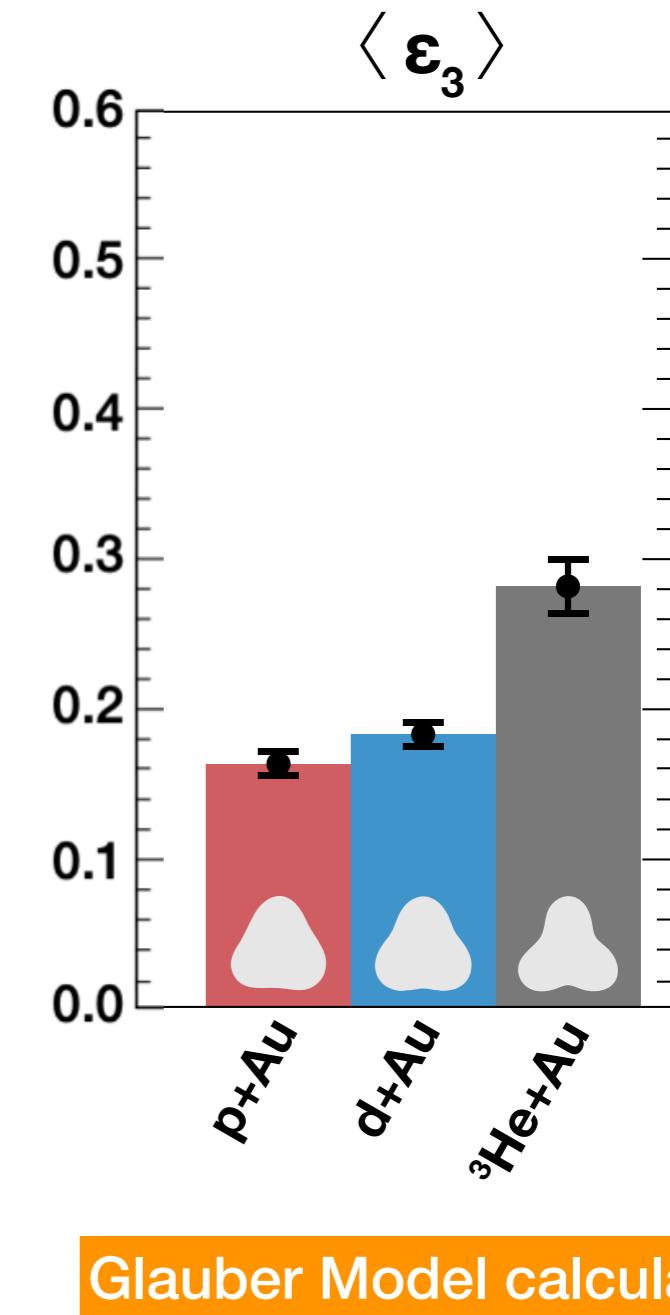
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# $v_3(p_T)$ measurement in small systems



$$v_3^{p+\text{Au}} \approx v_3^{d+\text{Au}} < v_3^{{}^3\text{He}+\text{Au}}$$

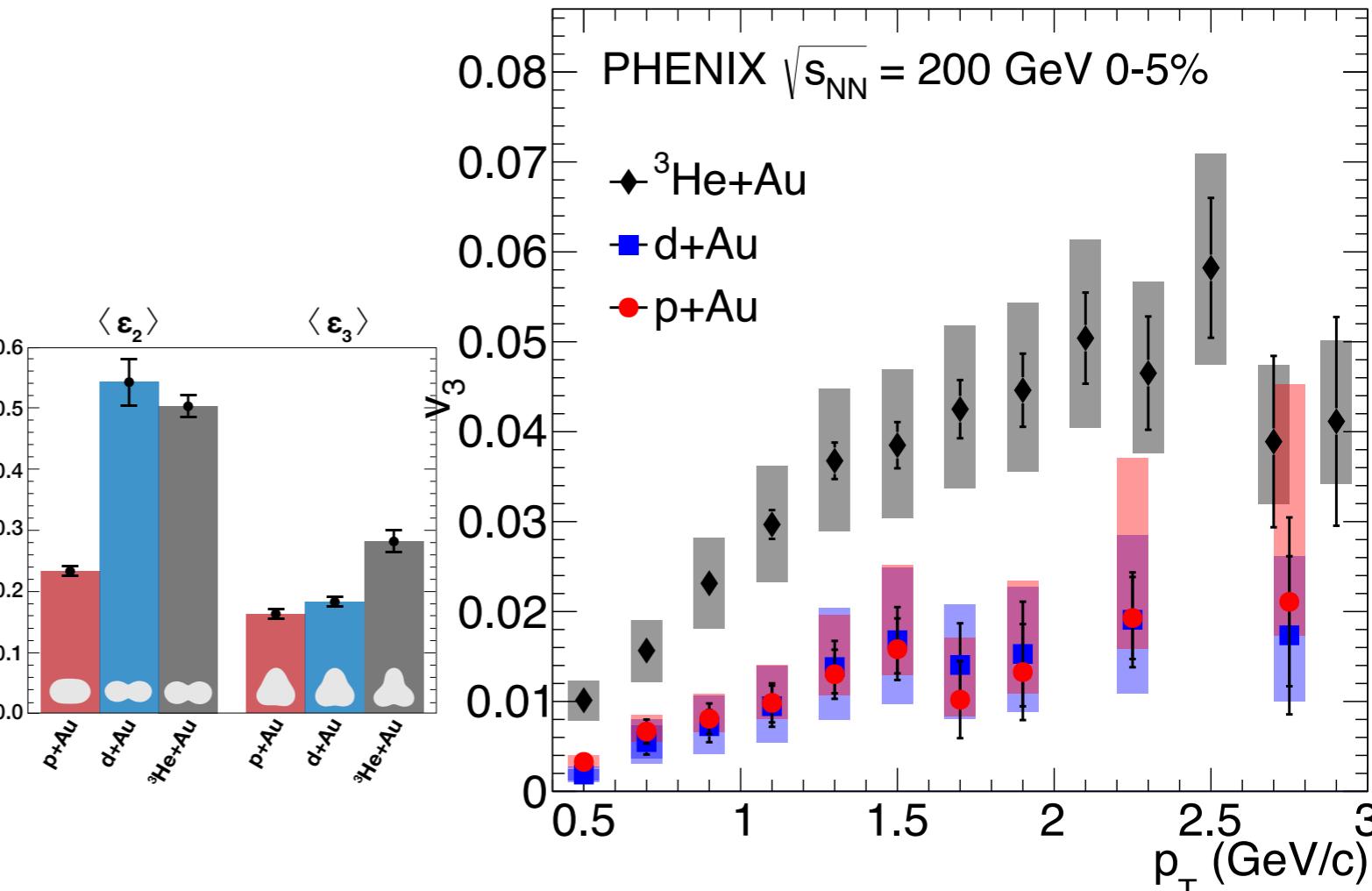
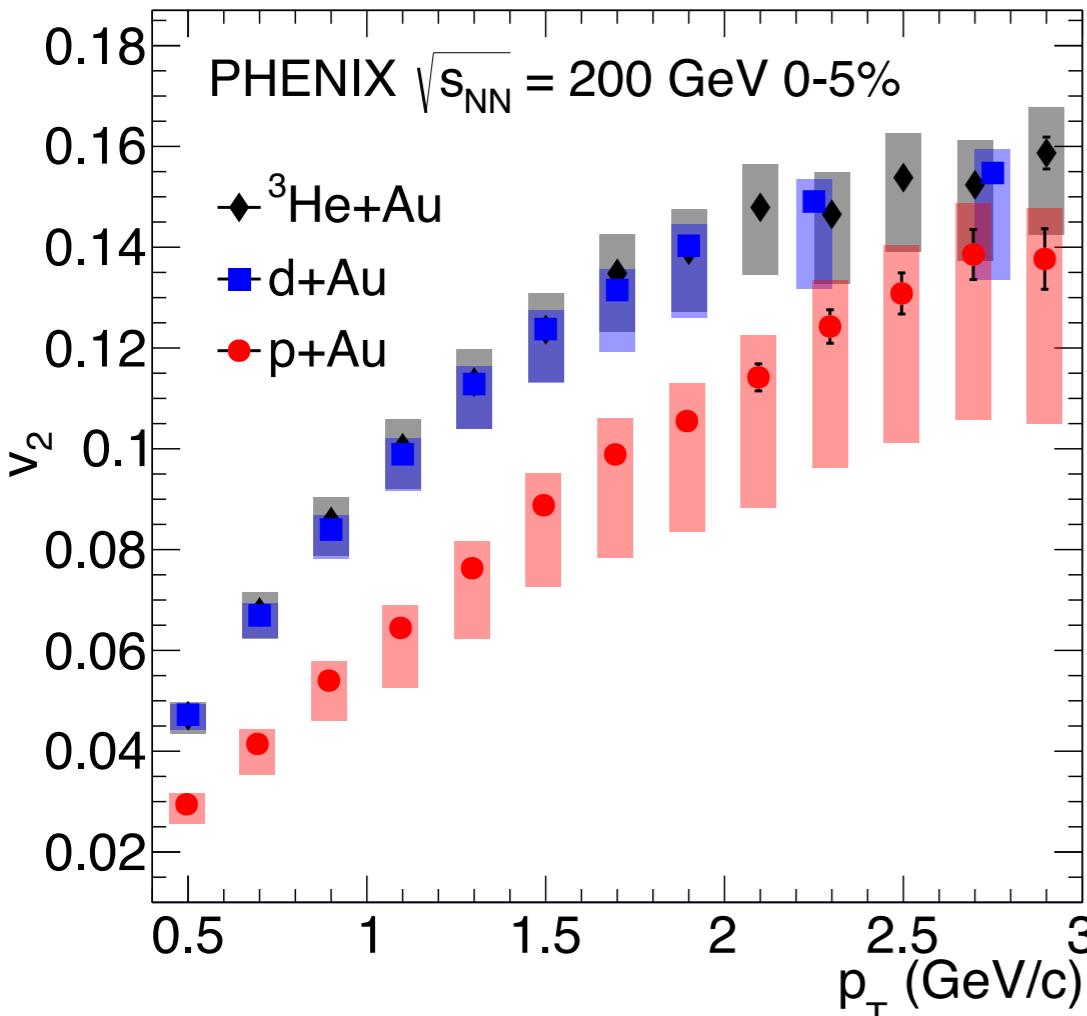
$$\varepsilon_3^{p+\text{Au}} \approx \varepsilon_3^{d+\text{Au}} < \varepsilon_3^{{}^3\text{He}+\text{Au}}$$





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# $v_n(p_T)$ measurement in small systems



$$v_2^{p+\text{Au}} < v_2^{d+\text{Au}} \approx v_2^{{}^3\text{He}+\text{Au}}$$

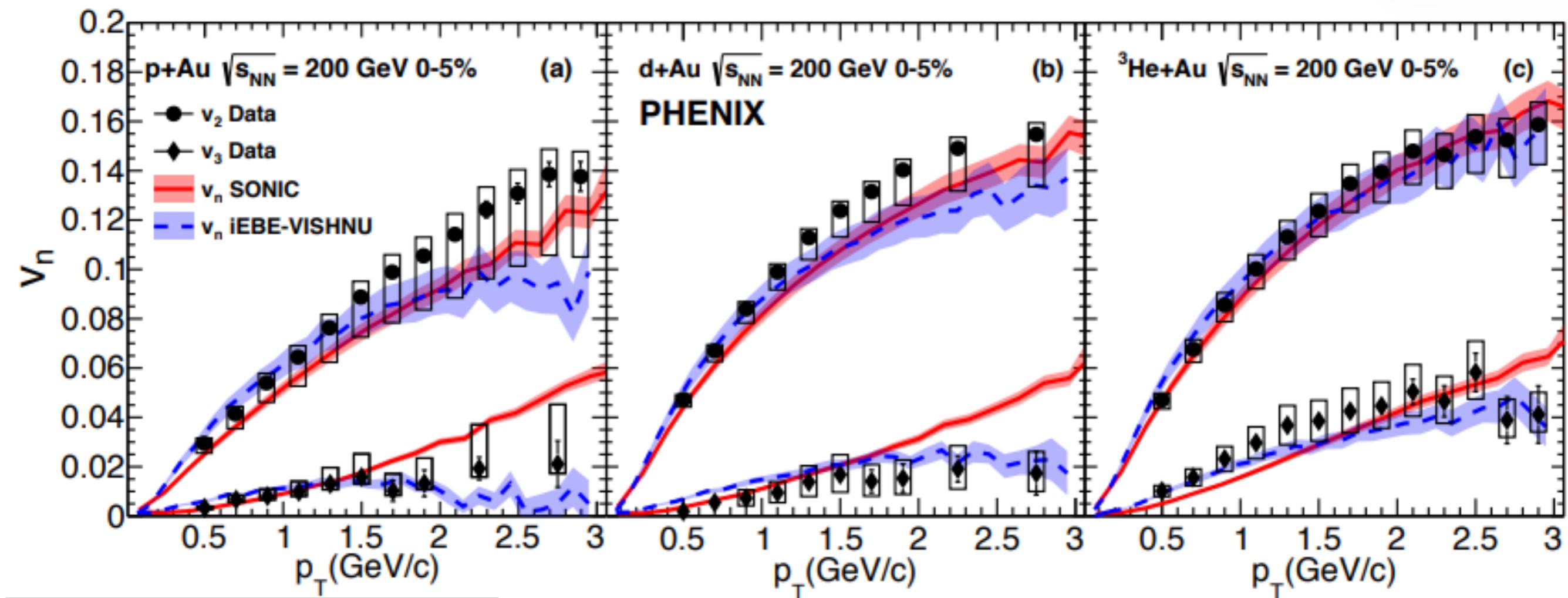
$$v_3^{p+\text{Au}} \approx v_3^{d+\text{Au}} < v_3^{{}^3\text{He}+\text{Au}}$$

**Suggests flow is geometric in origin**



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# $v_n(p_T)$ measurement comparing to hydro models



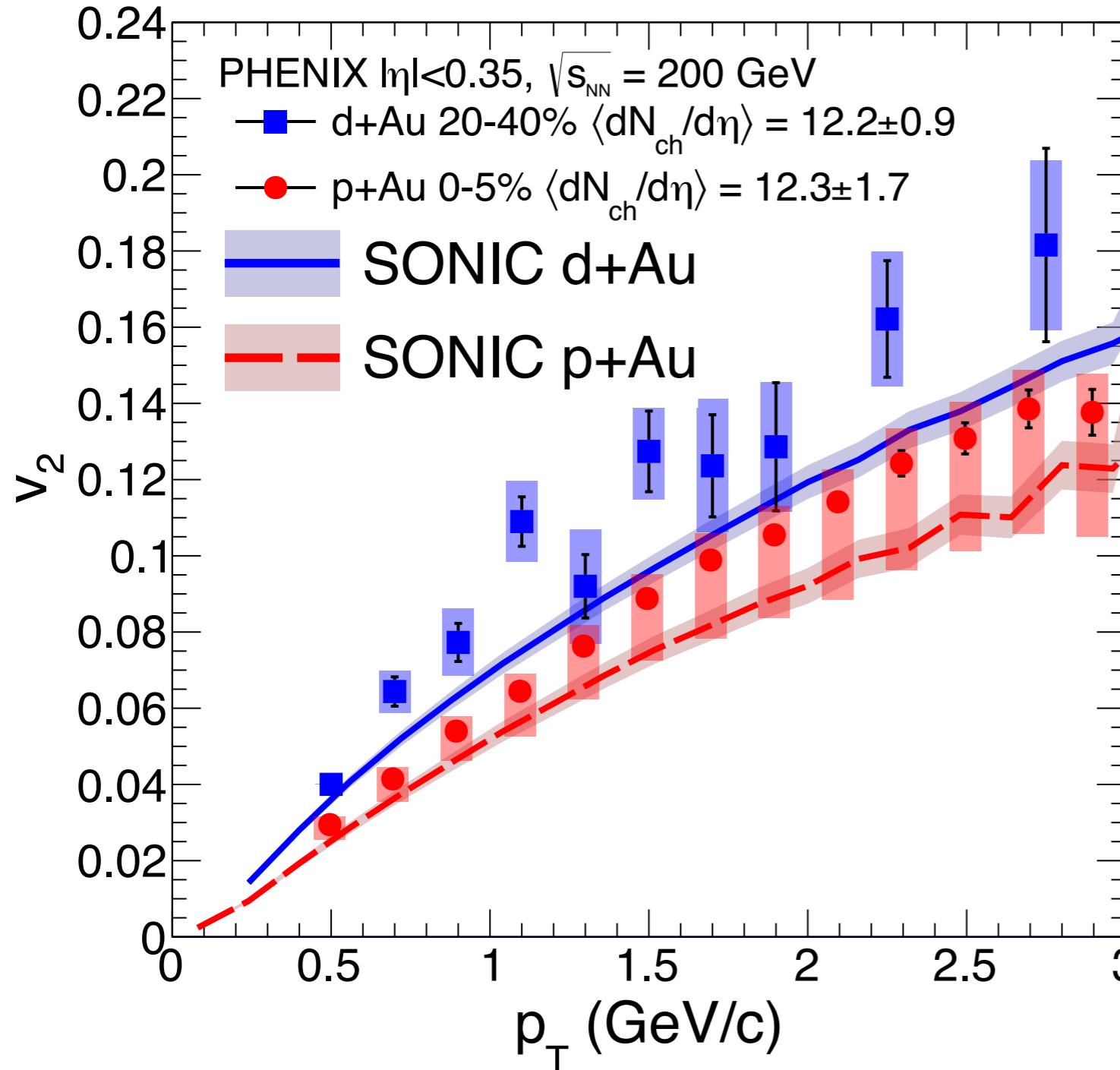
*iEBE-VISHNU*: C. Shen et al., Phys. Rev. C 95, 014906 (2017).  
*SONIC*: M. Habich et al., Eur. Phys. J. C 75, 15 (2015).

hydrodynamical models provide a simultaneous  
and quantitative description of the data in all three



# p/d+Au at same $\langle dN_{ch}/d\eta \rangle$

- d+Au 20-40%  $\langle dN_{ch}/d\eta \rangle = 12.2 \pm 0.9$  [PRC 96, 064905 \(2017\)](#)
- p+Au 0-5%  $\langle dN_{ch}/d\eta \rangle = 12.3 \pm 1.7$  [PRC 95, 034910 \(2017\)](#)



Good qualitative  
agreement of  
systematic  
difference between  
p+Au and d+Au

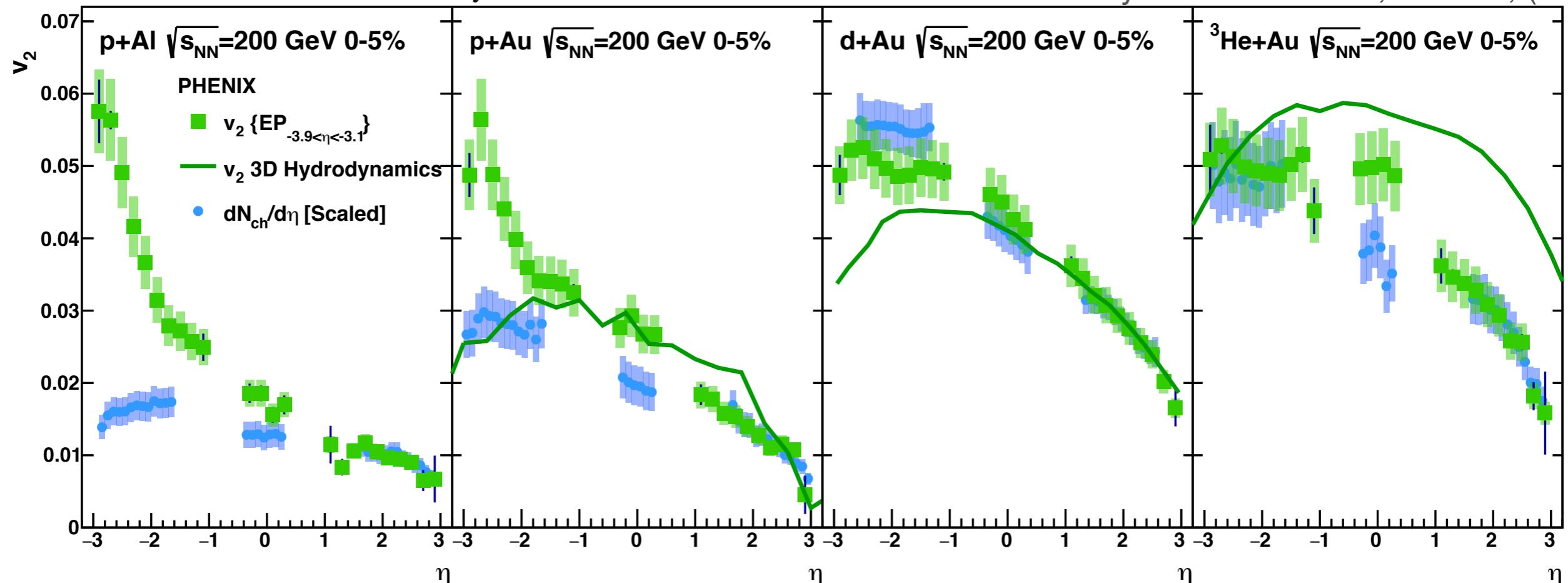


# $v_2(\eta)$ measurement comparing to hydro models

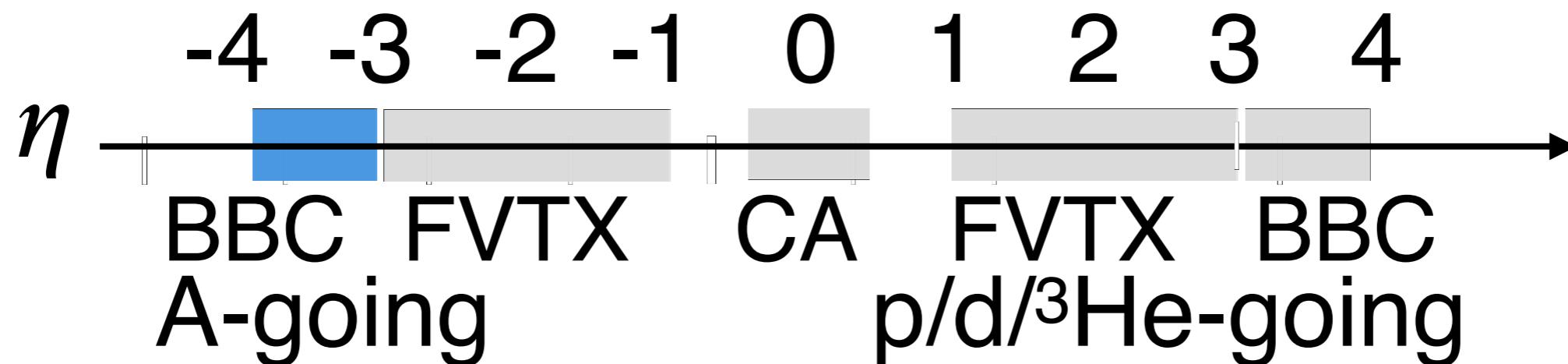
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P. Bozek etc.'s 3D hydro model

Phys. Rev. Lett. 121, 222301, (2018)



$v_2(\eta)$  is correlated with  $dN_{ch}/d\eta$  but not exact scaling  
3D Hydrodynamics predicts the trend

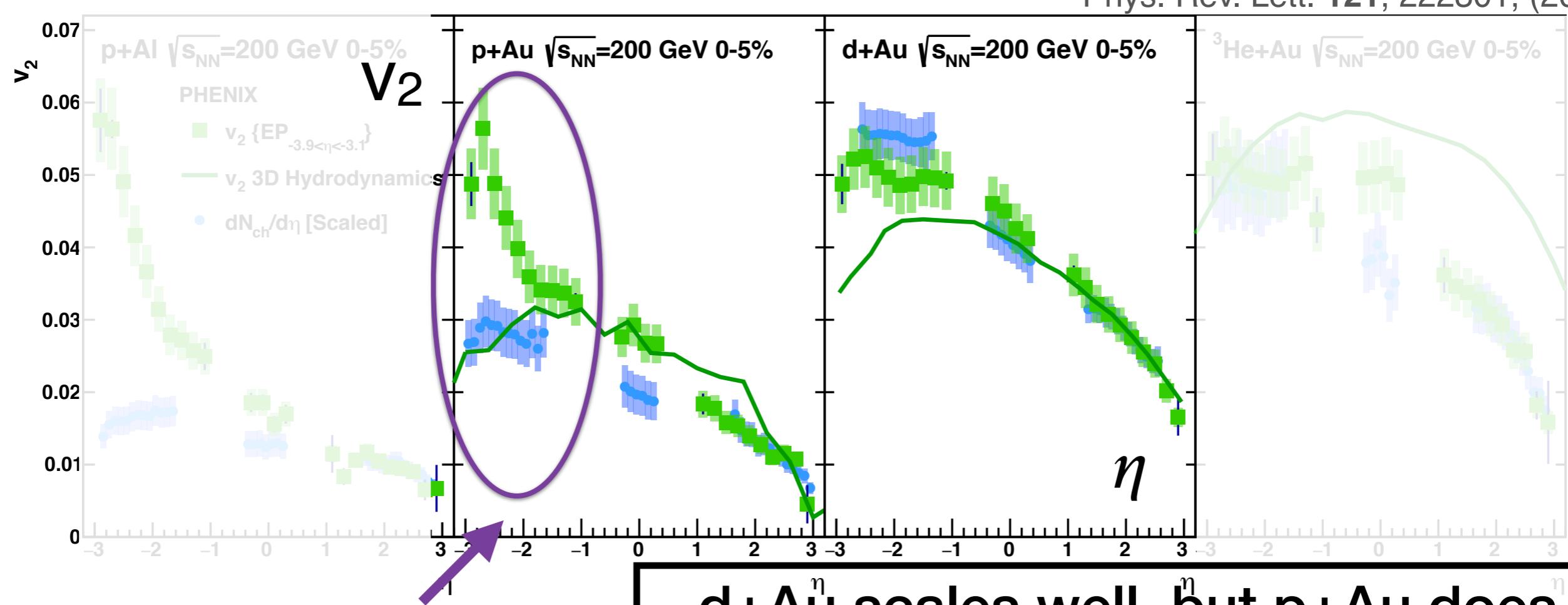




# $v_2(\eta)$ measurement comparing to hydro models

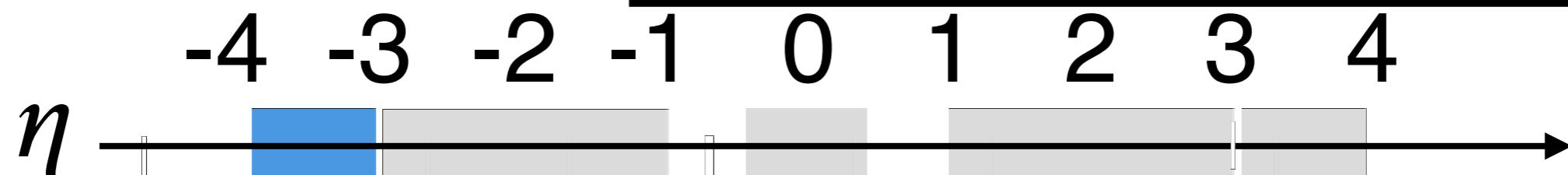
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Phys. Rev. Lett. 121, 222301, (2018)



Non-flow dominates near the EP detector

$d+Au$  scales well, but  $p+Au$  does not at backward rapidity, non-flow becomes more significant



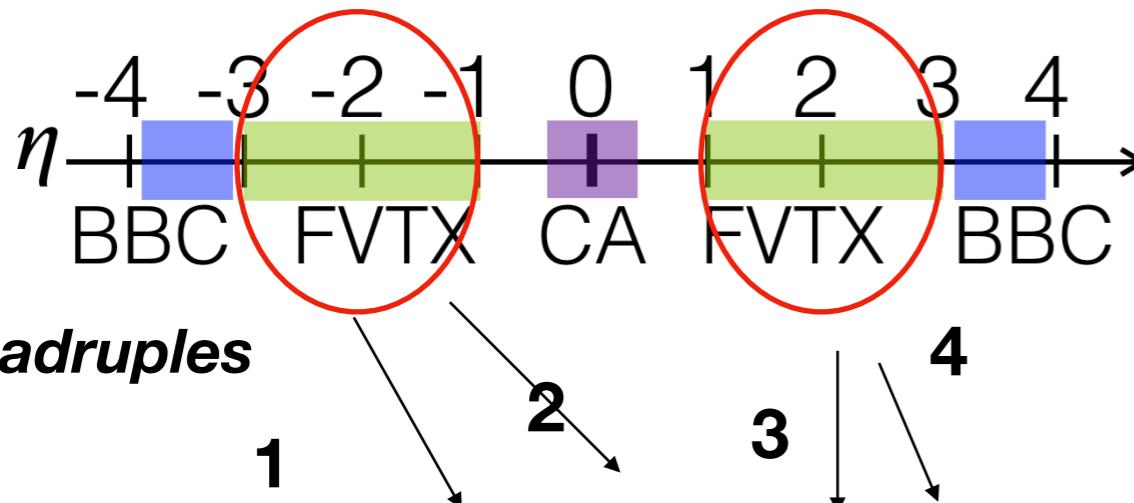
A-going

p/d/ ${}^3\text{He}$ -going



# Multi-particle correlation method

4-p correlation as an example:



$$\langle\langle 4 \rangle\rangle = \langle\langle e^{in(\phi_1 + \phi_2 - \phi_3 - \phi_4)} \rangle\rangle$$

$$\langle\langle 2 \rangle\rangle = \langle\langle e^{in(\phi_1 - \phi_2)} \rangle\rangle$$

$$c_2\{4\} = \langle\langle 4 \rangle\rangle - 2 * \langle\langle 2 \rangle\rangle^2$$

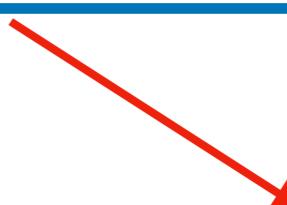
$$v_2\{4\} = (-c_2\{4\})^{-1/4}$$

$$v_2\{4\}^2 \approx v_2\{6\}^2 \approx \langle v_2 \rangle^2 - \sigma_{v2}^2$$

$$v_2\{2\}^2 \approx \langle v_2 \rangle^2 + \sigma_{v2}^2 + \delta^2$$

Multi-particle correlations suppress non-flow effectively

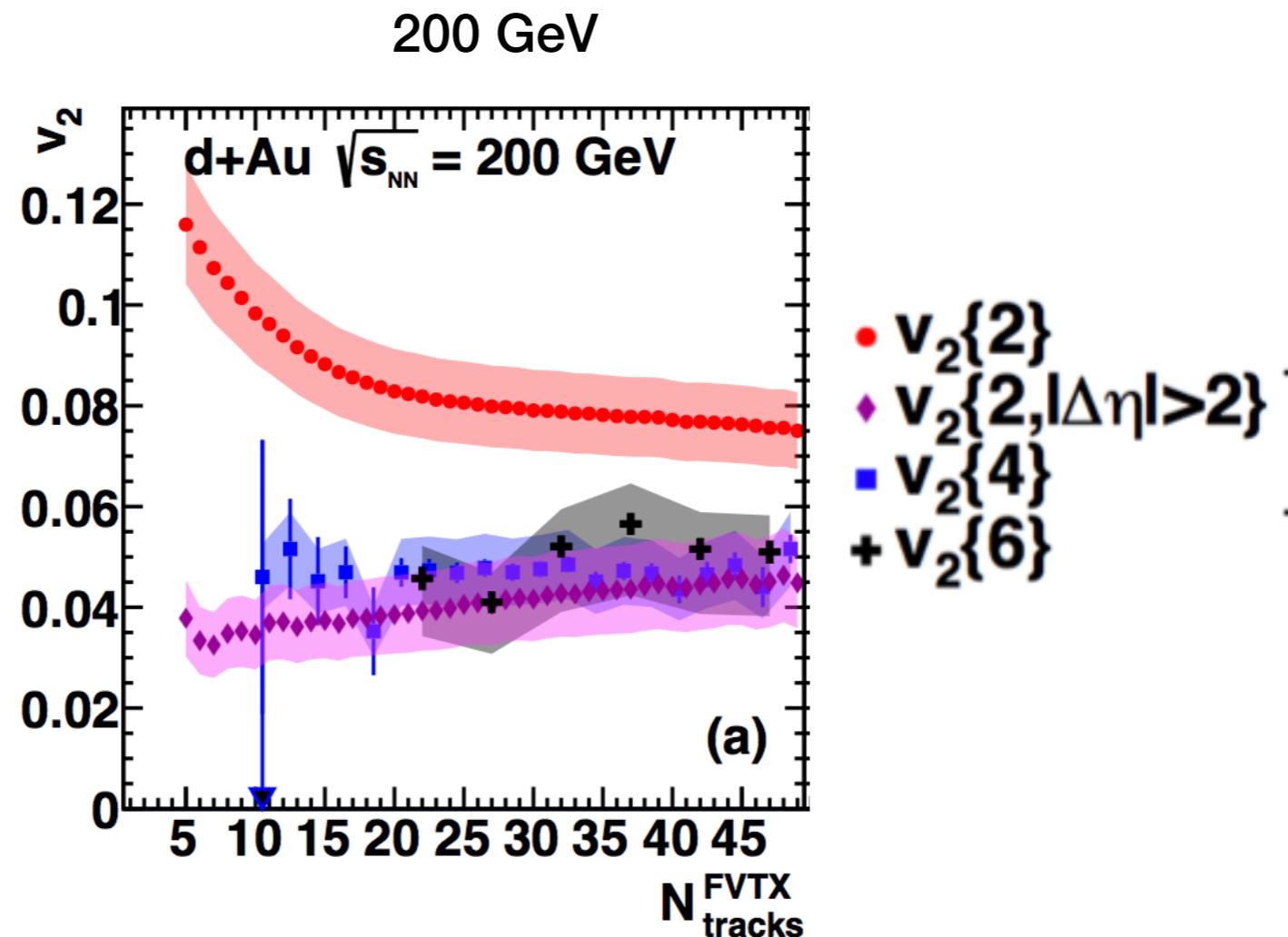
Multi-particle correlations also provide information about how flow fluctuate event-by-event





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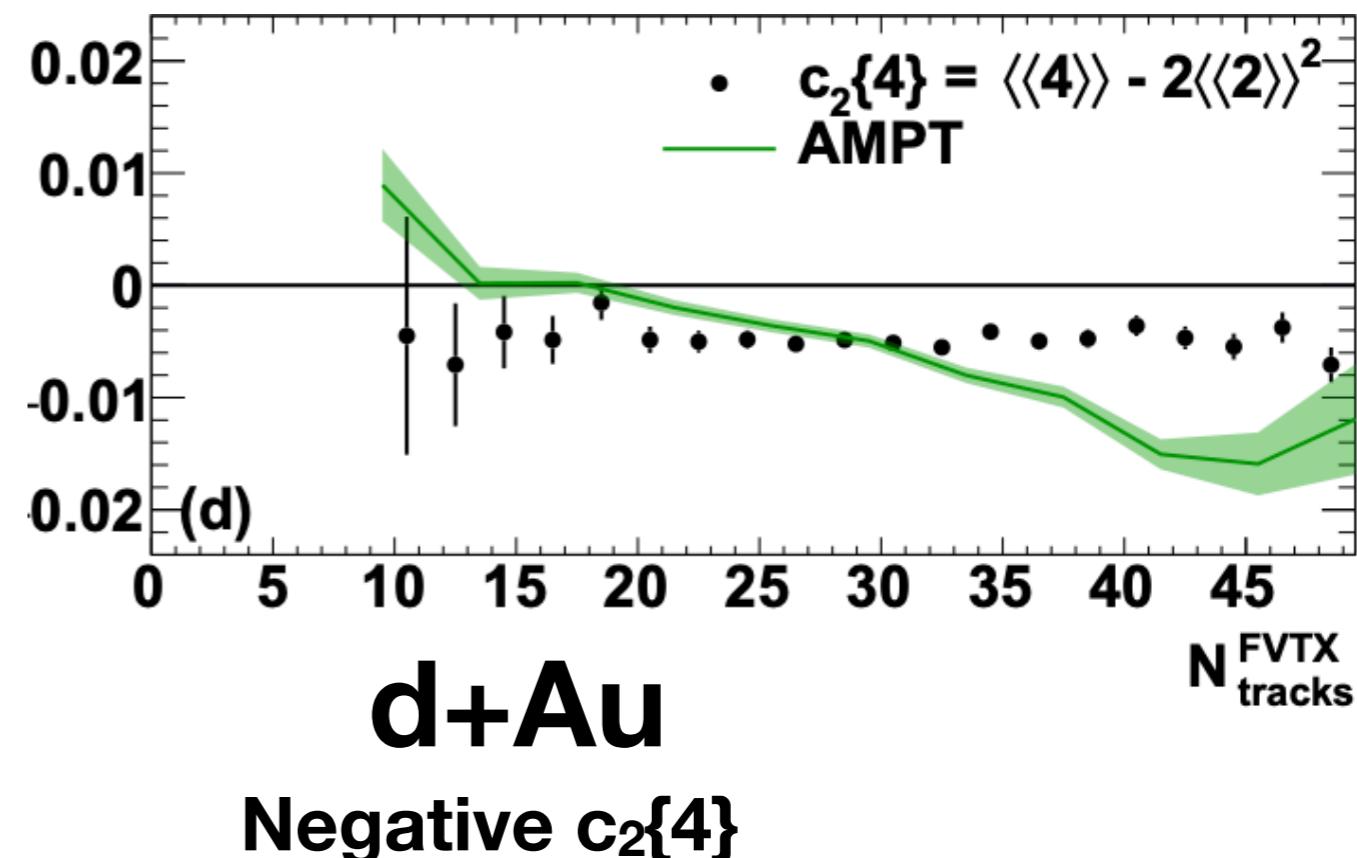
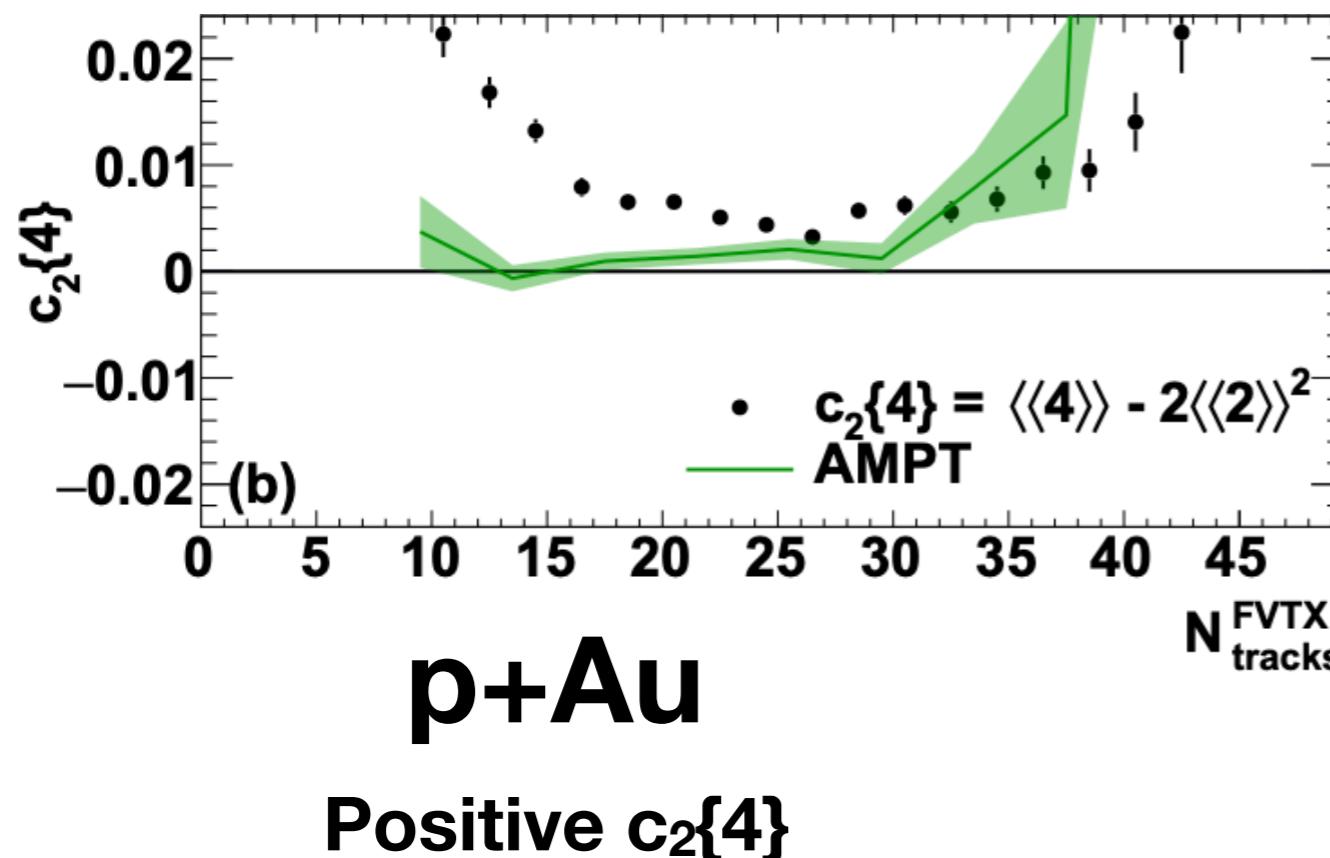
# Multi-particle correlations in d+Au@200GeV



- 200 GeV  $v_2\{2\}[\eta \text{ gap}] \approx v_2\{4\} \approx v_2\{6\}$ , indication of strong collectivity
- 200 GeV  $v_2\{4\}$  dominated by flow, not non-flow

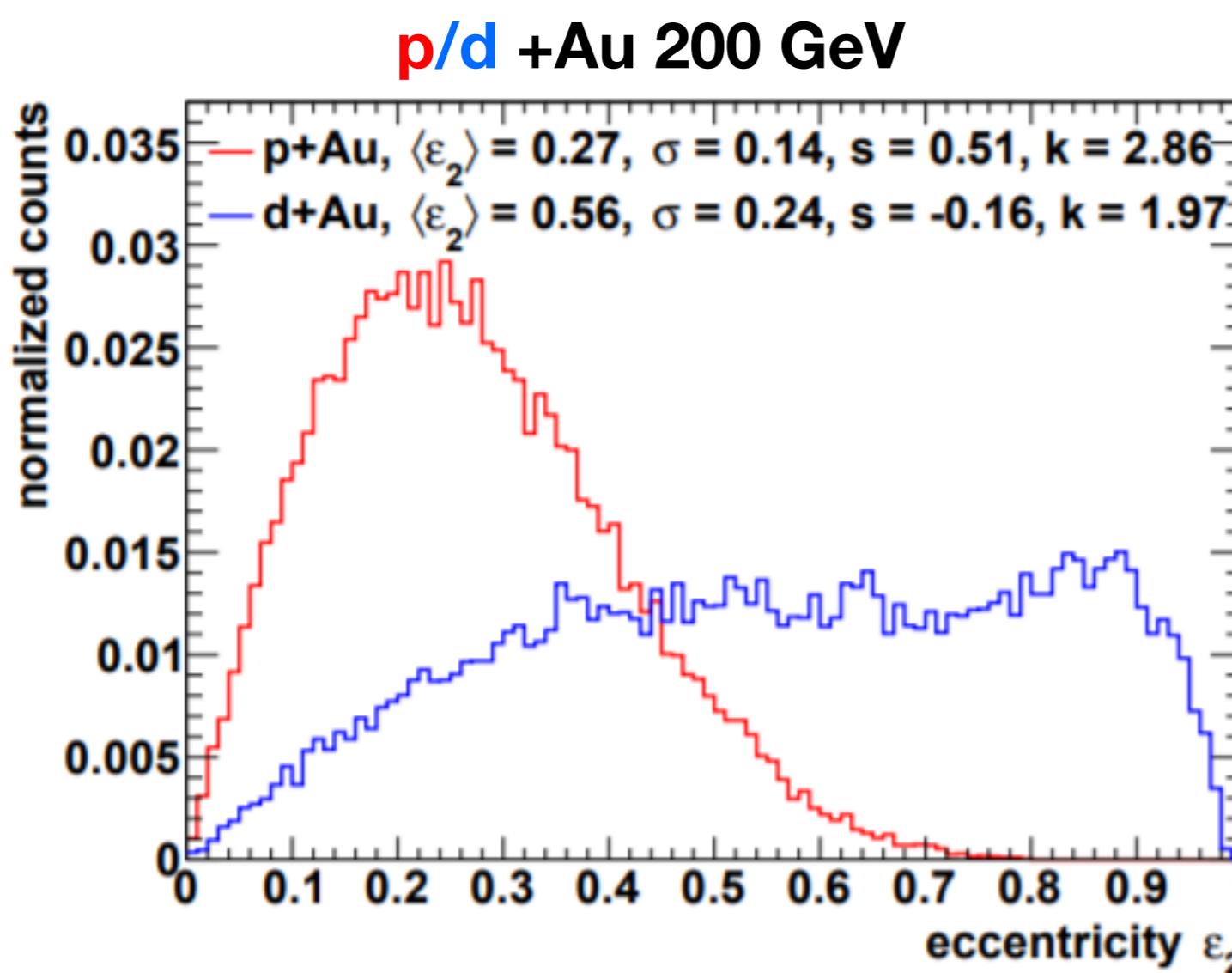
# Flow fluctuations in p/d+Au@200GeV

*Phys. Rev. Lett. 120, 062302 (2018)*



- $v_2\{4\}^2 \approx \langle v_2 \rangle^2 - \sigma_{v2}^2$ , If fluctuation  $\sigma_{v2} >$  mean  $v_2$ ,  $c_2\{4\}$  is positive
- Implies  $c_2\{4\}$  in  $p+\text{Au}$  is dominated by fluctuations
- AMPT (A Multi-phase transport model) describes the sign

# Initial eccentricity distribution



*Monte Carlo Glauber*

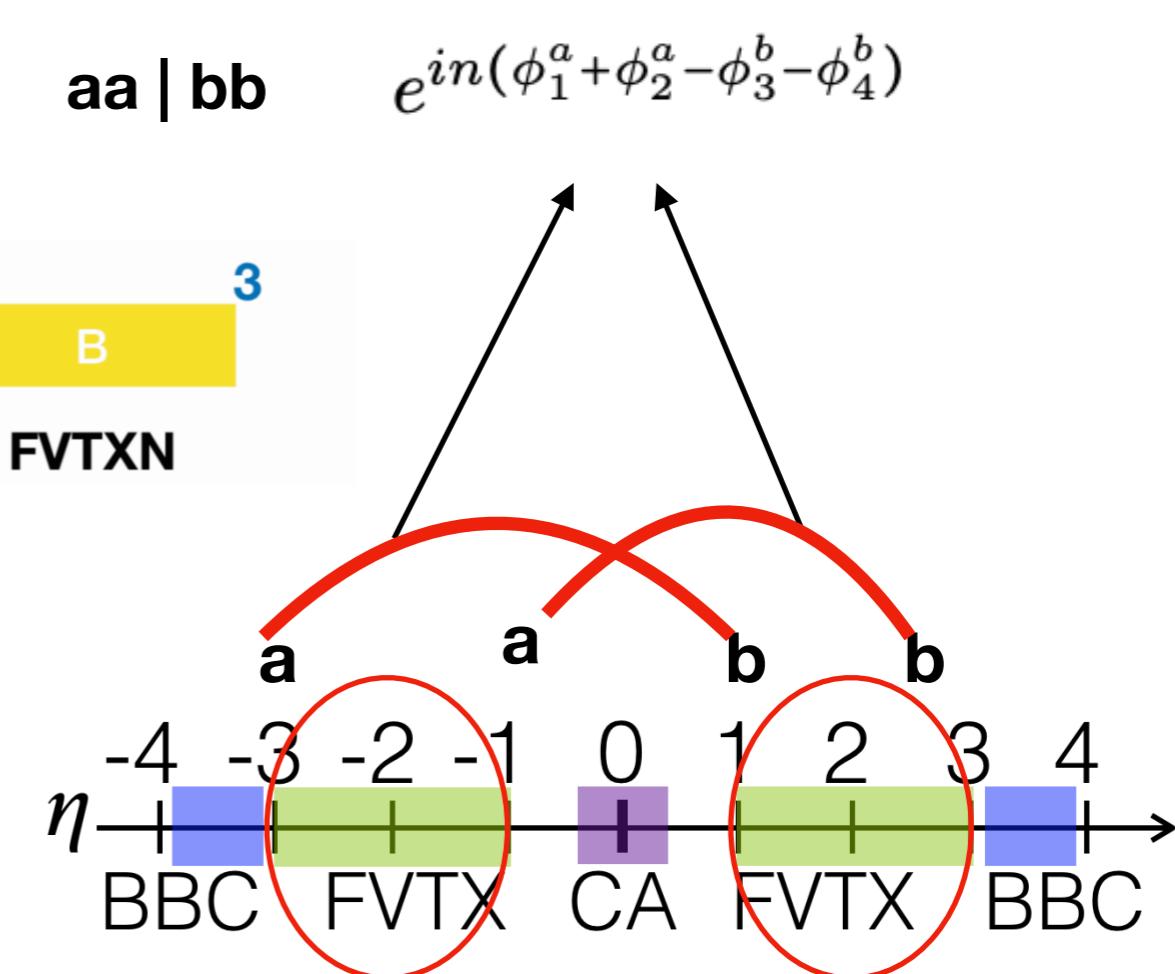
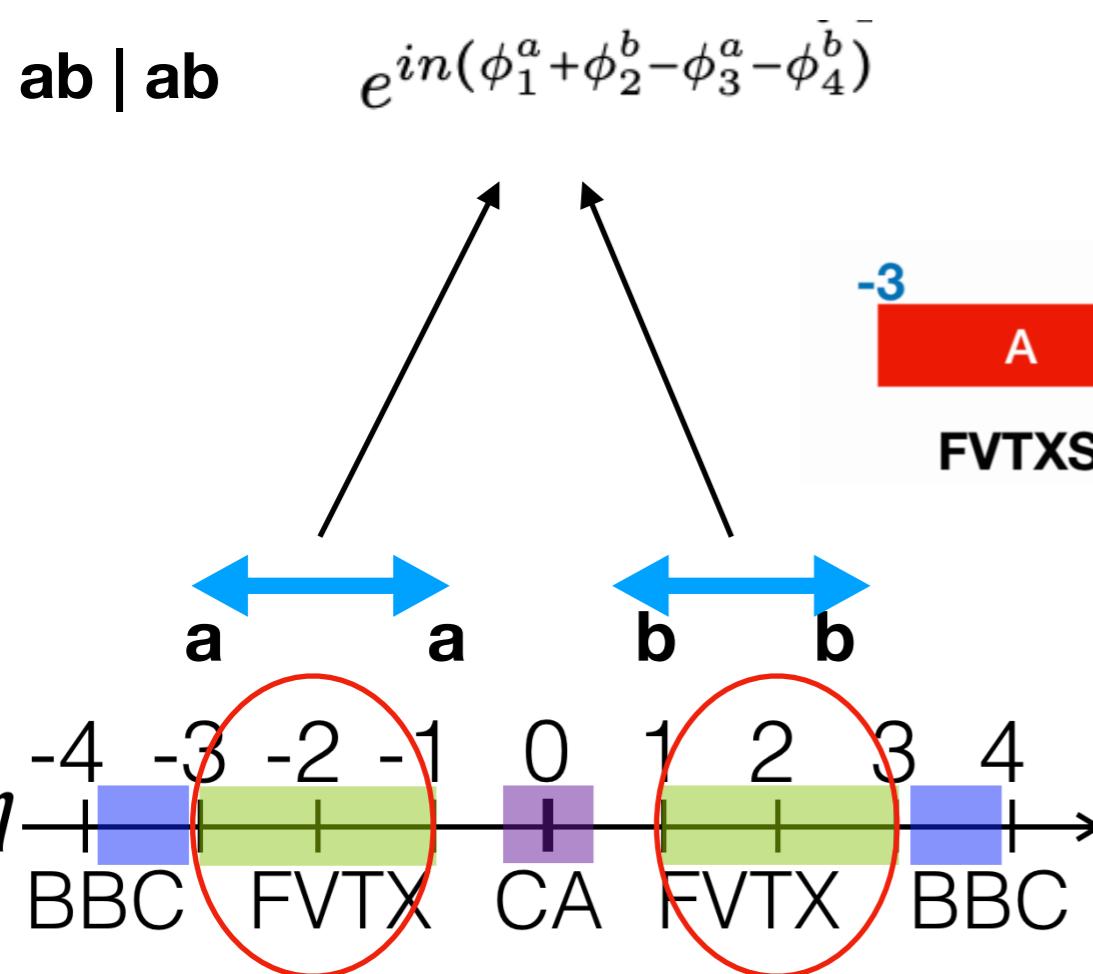
High positive skewness in p+Au  
 Significantly platykurtic in d+Au

- Initial eccentricity distribution is highly non-Gaussian.
- Fluctuations are highly non-trivial in small systems

*Phys. Rev. Lett. 120, 062302 (2018)*

# Subevent cumulant method

- To further suppress non-flow and investigate the role of fluctuations, 2-subevent cumulant method is used
- These 2 subevents already has an  $\eta$  gap  $> 2$
- Two types depending on how we correlate particles are studied
- Typically non-flow contained: **aa|bb < ab|ab < standard method**

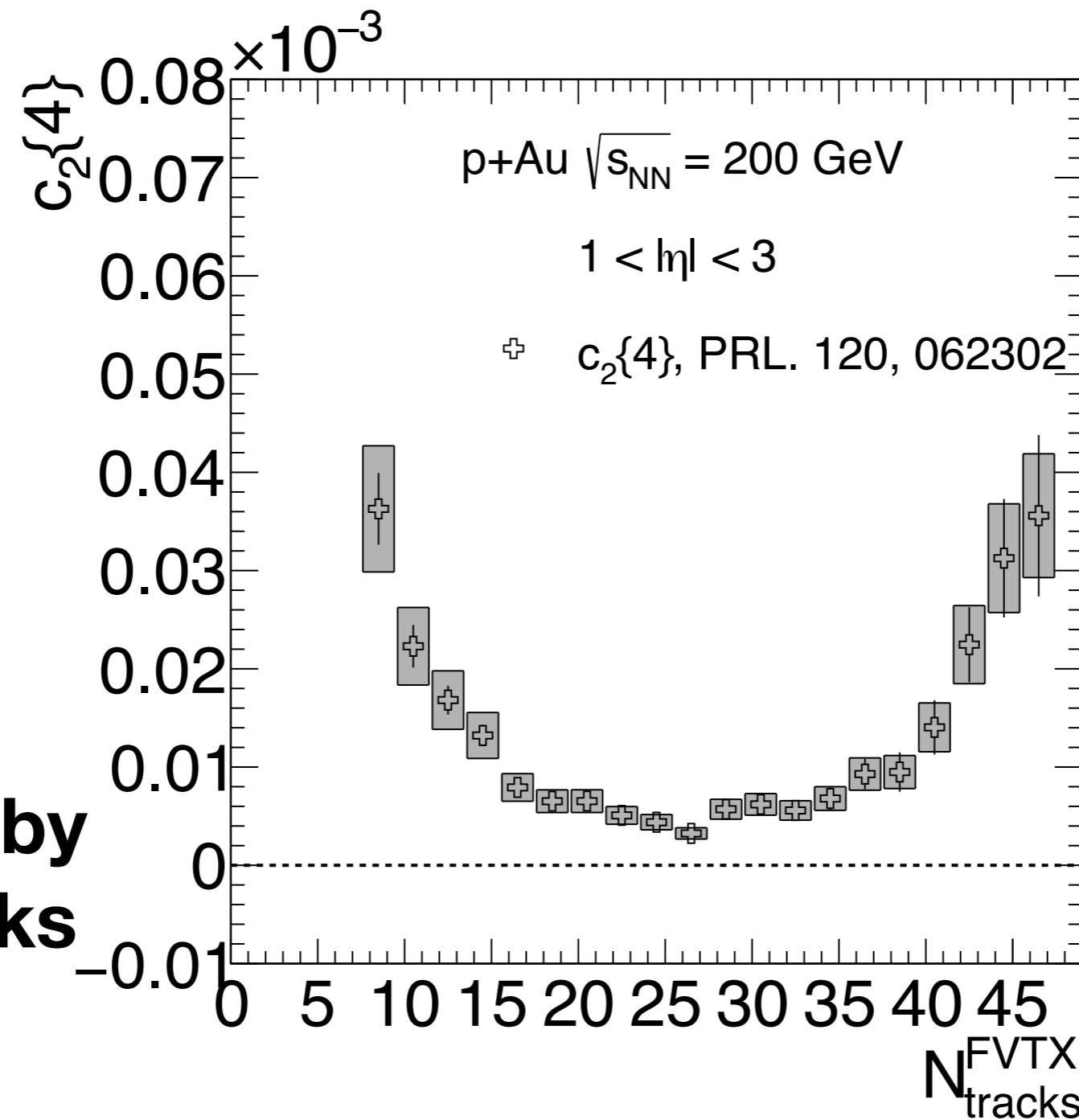




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# Subevent cumulants in p+Au@200GeV

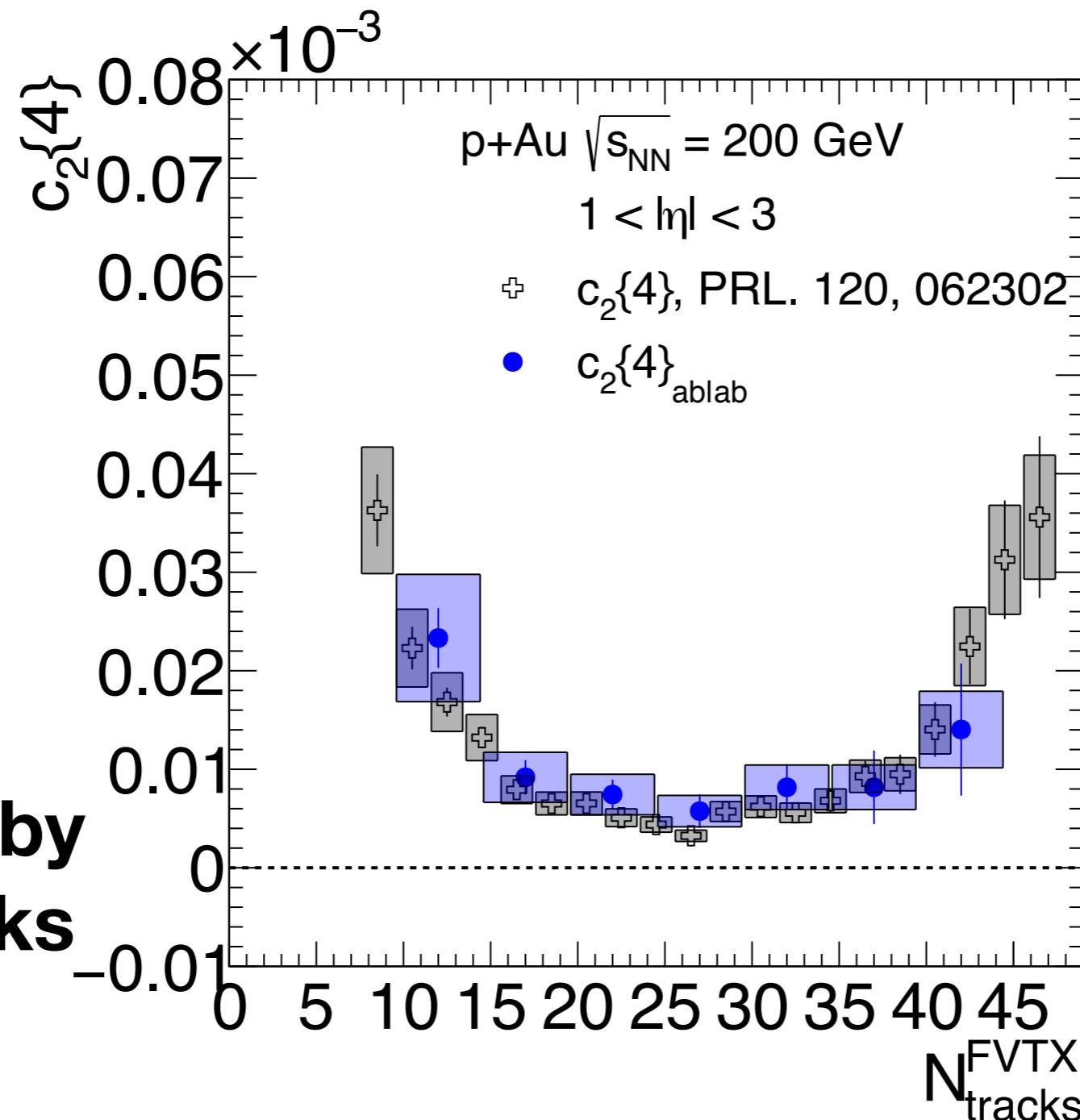
**dominated by  
low  $p_T$  tracks**



- Standard  $c_2\{4\}$  in p+Au is positive

# Subevent cumulants in p+Au@200GeV

**dominated by  
low  $p_T$  tracks**



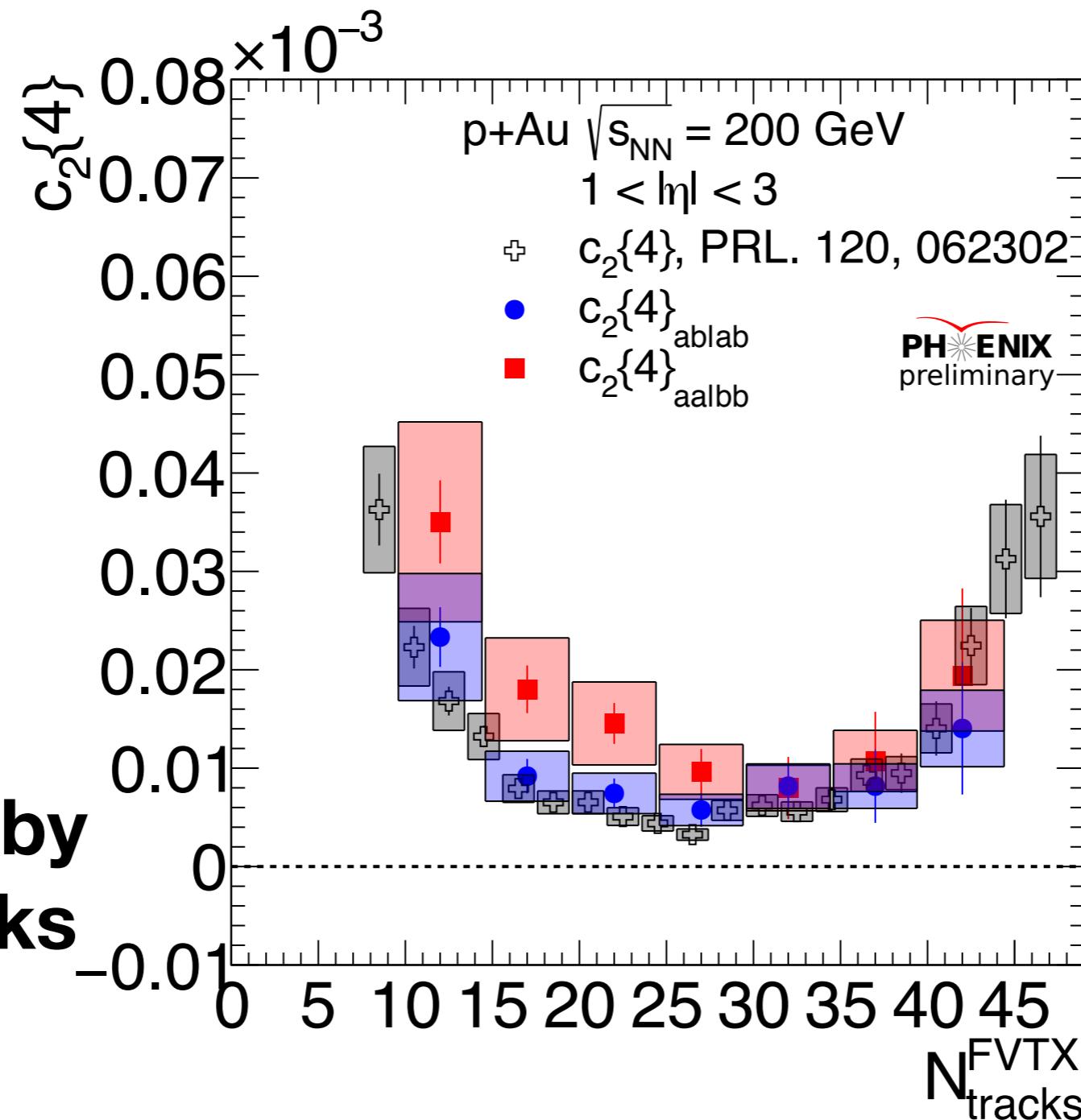
- Two sub-event  $c_2\{4\}$  in p+Au are still positive
- Consistent with standard  $c_2\{4\}$  without subevent



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# Subevent cumulants in p+Au@200GeV

dominated by  
low  $p_T$  tracks



Our PRL results  
are very robust!

- Two sub-event  $c_2\{4\}$  in p+Au are still positive
- Further confirm that flow fluctuations are important in p+Au

# Conclusions

- **Initial Geometry** is the driving force of the final correlations
- **Small variance limit** breaks in p+Au and in d+Au
- **Flow fluctuations are significant in  $c_2\{4\}$**  in p+Au collisions confirmed by **sub-event cumulant analysis**



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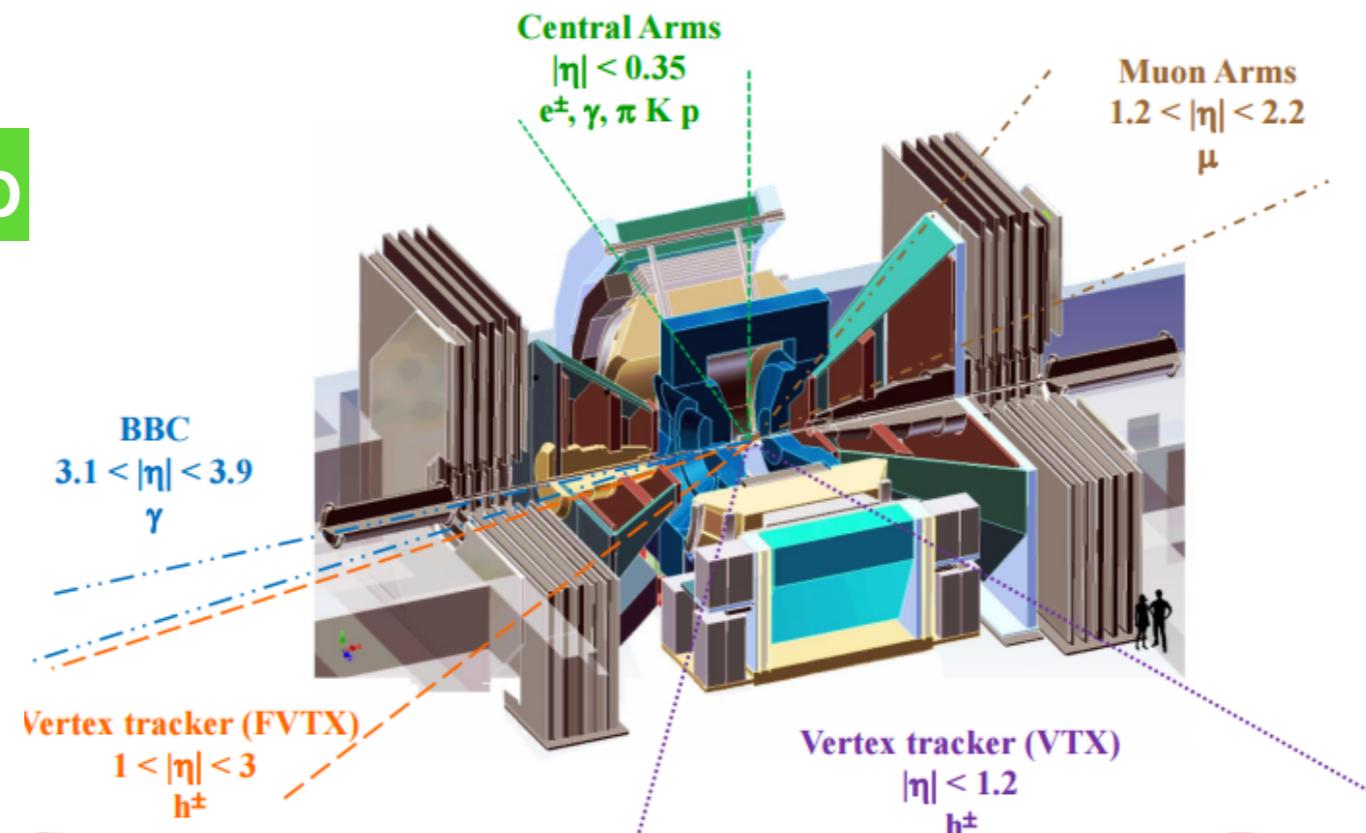
# Back Up



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# Small system measurements in PHENIX

## Detector Setup



- Midrapidity: DC, PC, TOF  $\rightarrow$  tracking and PID
- Forward: BBC, FVTX  $\rightarrow$  triggering, event selection, correlations with midrapidity particles, event plane determination
- Muon arm: FVTX, Muon Tracking, Muon ID  $\rightarrow$  heavy flavor tracking and identification

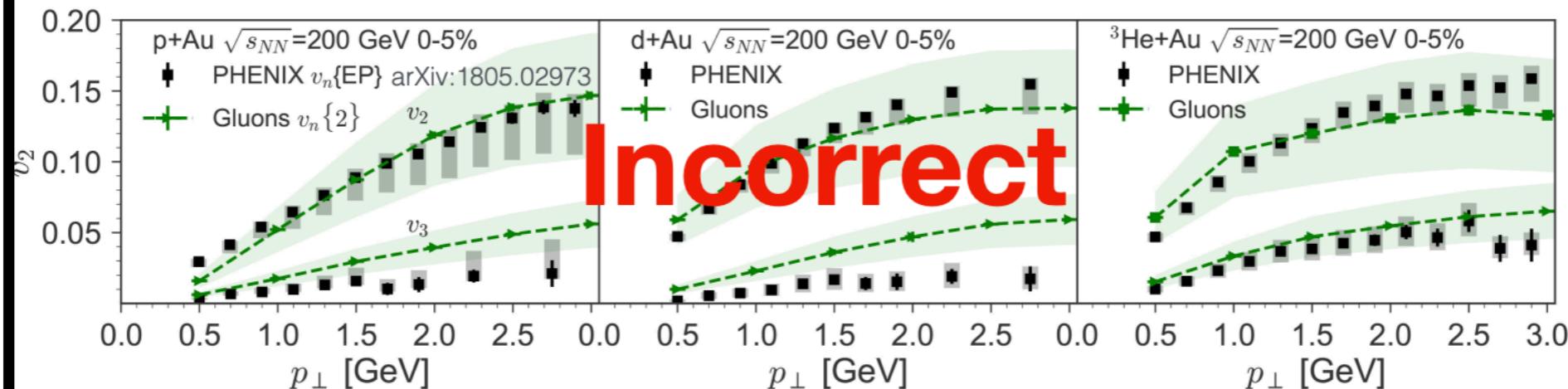


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# From Latest development for CGC models

from M.Mace at INT - May 2, 2019 Origins of Correlations in High Energy Collisions

## An error



Digging deeper in the code in answer questions raised in small systems workshop in Houston and preparing to make code public, a bug was found in the units

$$\frac{dN}{d^2k} [\text{GeV}^{-2}] \stackrel{!}{=} f(k') [\text{fm}^2]$$

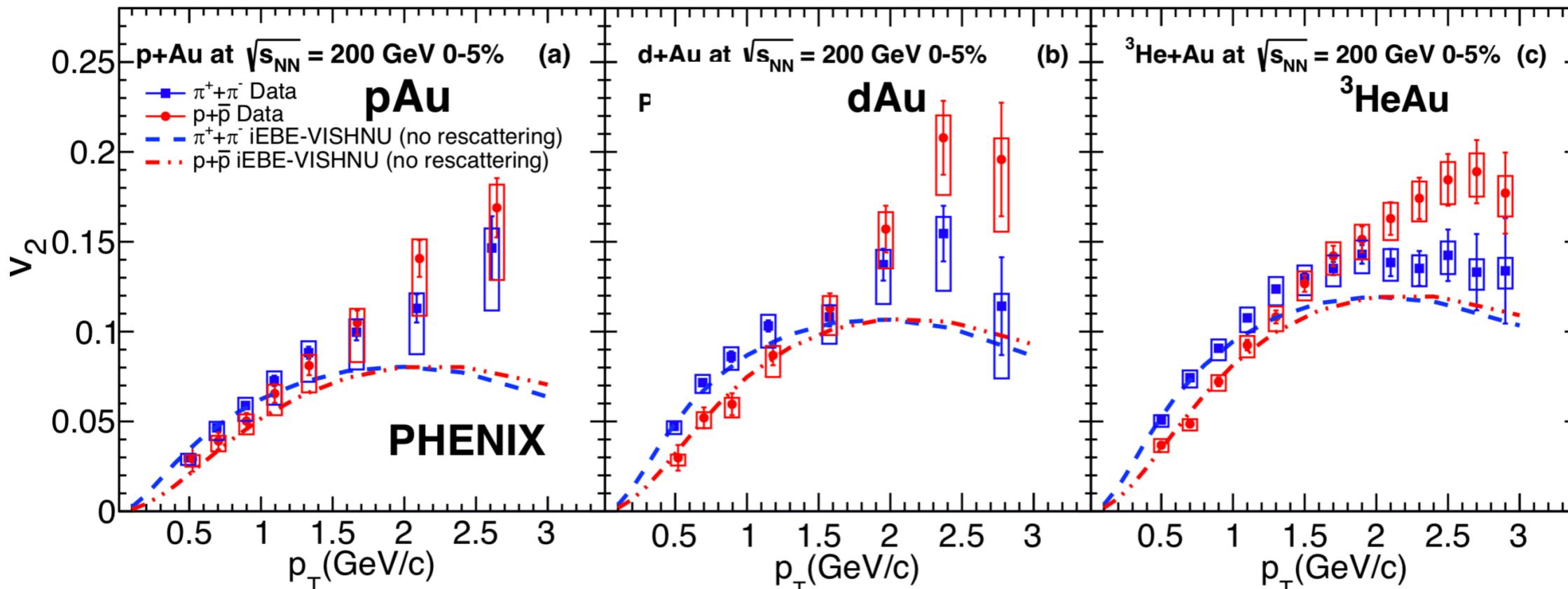
Effectively, all results were off by in momenta factor of  $\hbar c$



# $v_2(p_T)$ for identified $\pi^\pm$ and protons

System Size Increases

Phys. Rev. C 97, 064904 (2018)



- $v_2(\pi^\pm) > v_2(\text{proton})$  at  $p_T < 1.5$  GeV/c, reversed at higher  $p_T$
- The hydro model describes the low- $p_T$  mass-ordering in  $v_2(p_T)$  well
- hydro doesn't get the intermediate pT splitting so well because of the hadronization

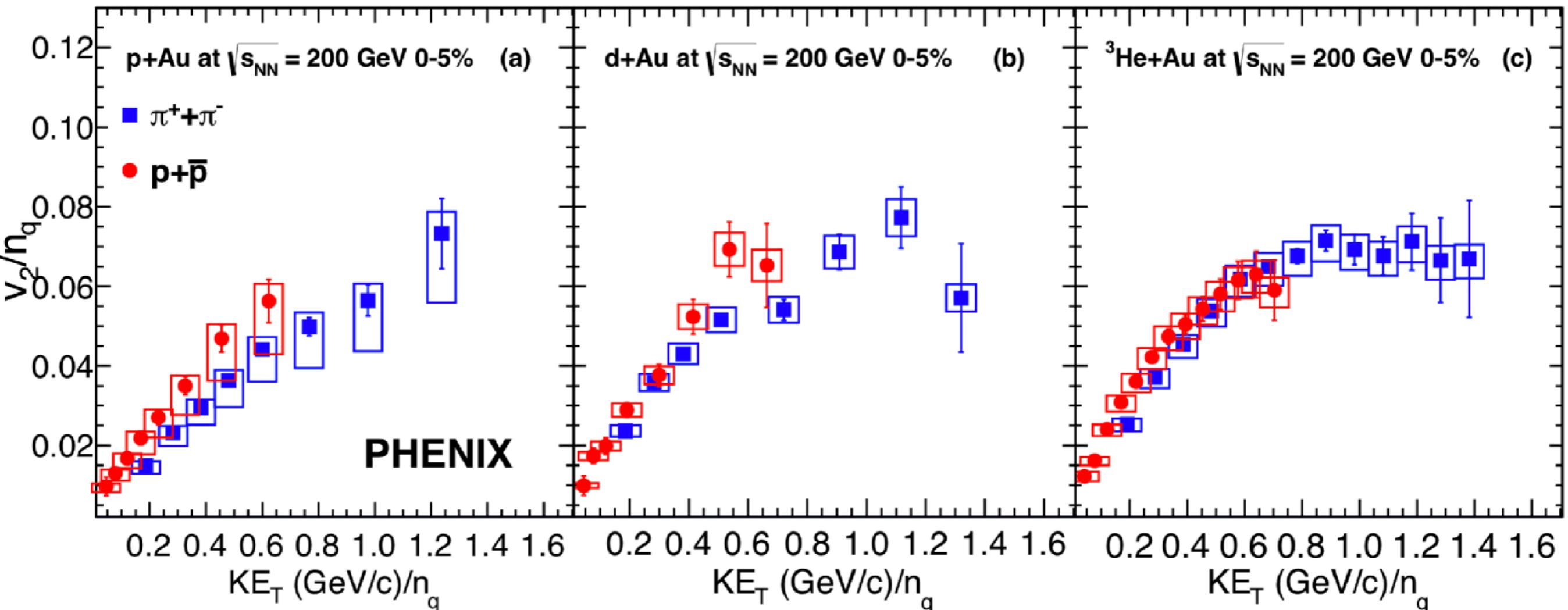


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# Quark number scaling of $v_2$ in p/d/ $^3\text{He}$ +Au

## Is there a common velocity field?

PRC 97, 064904 (2018)



Quark number scaling observed similar to AA  
holds better as the system size increases

# Analysis methods for Flow

## Two – particle correlation method

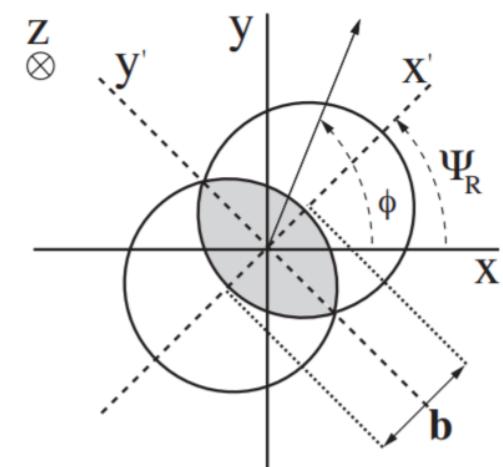
Pairs:  $\frac{dN}{d\Delta\phi} \propto 1 + \sum_n 2v_n^a v_n^b \cos(n\Delta\phi)$  **2PC method**

## Event plane method:

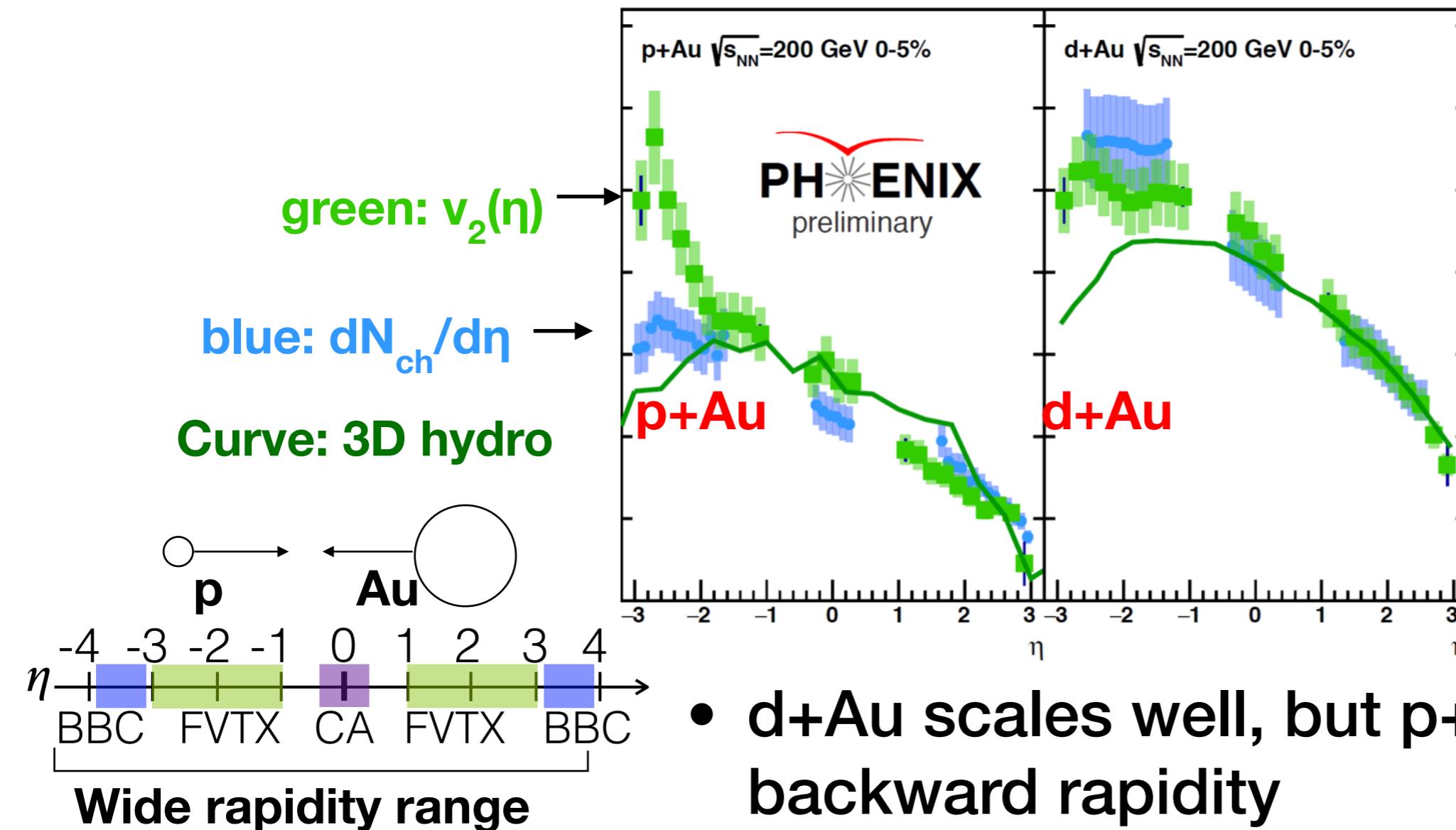
$$dN / d\phi = 1 + \sum_n 2v_n \cos(n(\phi - \Psi_n))$$

## Multi-particle correlation method:

$$\begin{aligned} \langle 2 \rangle &\equiv \left\langle e^{in(\phi_1 - \phi_2)} \right\rangle \equiv \frac{1}{P_{M,2}} \sum'_{i,j} e^{in(\phi_i - \phi_j)}, \\ \langle 4 \rangle &\equiv \left\langle e^{in(\phi_1 + \phi_2 - \phi_3 - \phi_4)} \right\rangle \\ &\equiv \frac{1}{P_{M,4}} \sum'_{i,j,k,l} e^{in(\phi_i + \phi_j - \phi_k - \phi_l)}, \end{aligned}$$



# $v_2(\eta)$ vs $dN_{ch}/d\eta$ in Geometry Control Scan



- d+Au scales well, but p+Au does not at backward rapidity
- 3D hydrodynamics quantitatively describes the data in p+Au and d+Au

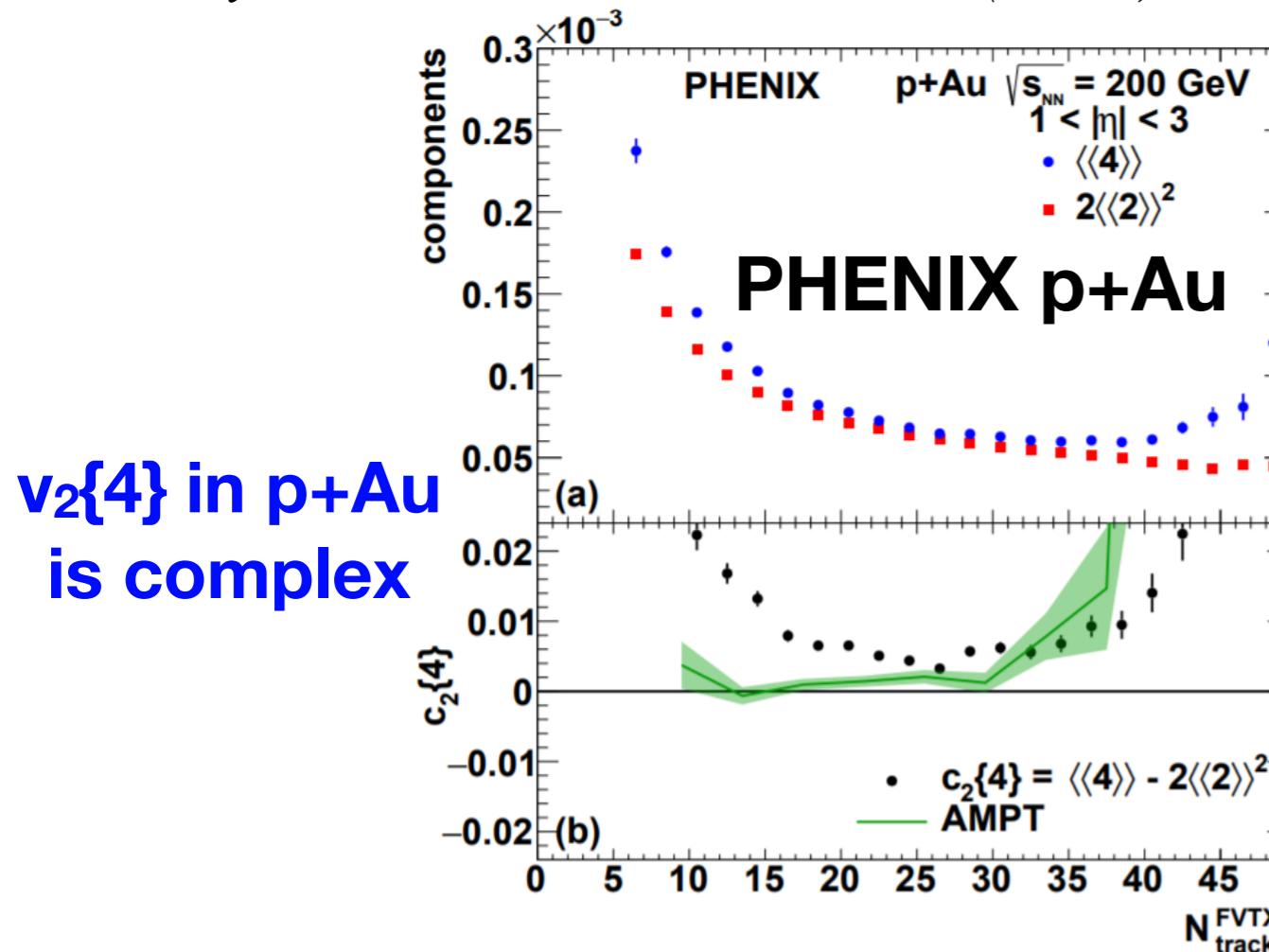
The event plane is measured in  $-3.9 < \eta < -3.1$



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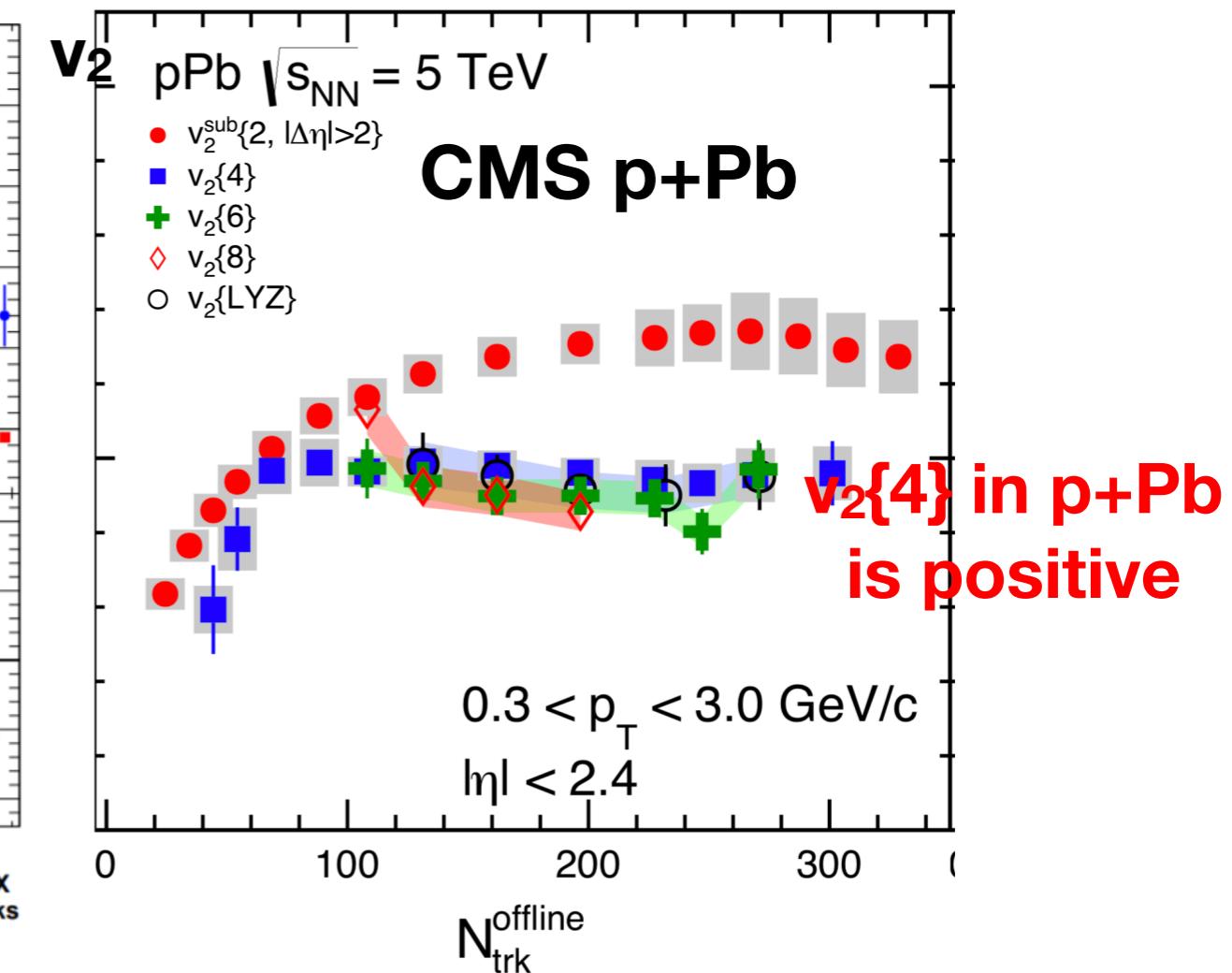
# Multi-particle correlations in p+A comparing RHIC and LHC results

*Phys. Rev. Lett. 120, 062302 (2018)*



**$v_2\{4\}$  in p+Au is complex**

*Phys. Lett. B 765 (2017) 193-220*

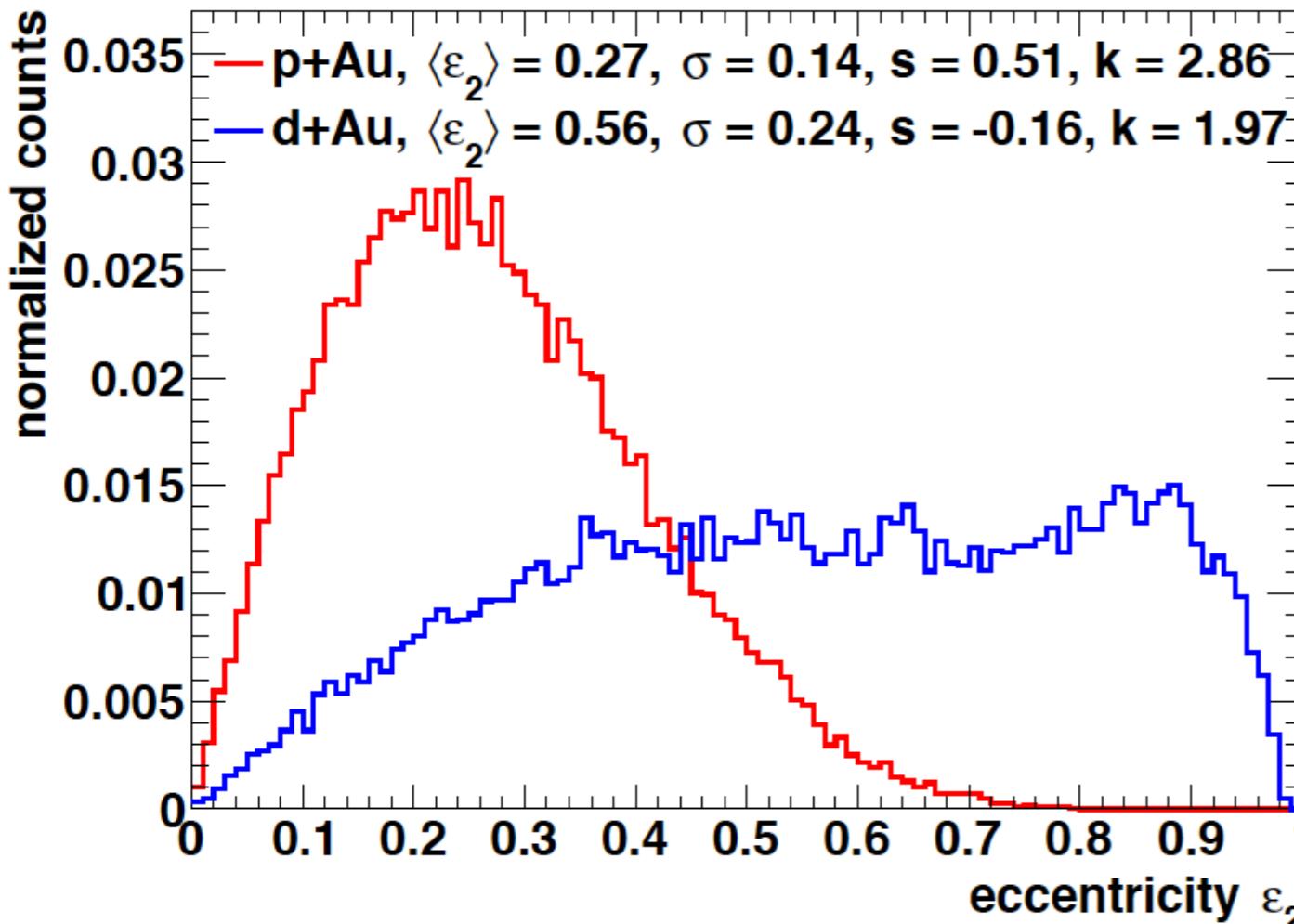


**$v_2\{4\}$  in p+Pb is positive**

- p+A at RHIC energy has more contributions from fluctuations than LHC energy



## Figure from Ronald Belmont



	p+Au	d+Au
$\varepsilon_2^4$	0.00531	0.0983
$2\varepsilon_2^2\sigma^2$	0.00277	0.0370
$4\varepsilon_2 s \sigma^3$	0.00147	-0.0053
$(k - 2)\sigma^4$	0.00031	-0.0001

$$\varepsilon_2\{4\} = (\varepsilon_2^4 - 2\varepsilon_2^2\sigma^2 - 4\varepsilon_2 s \sigma^3 - (k - 2)\sigma^4)^{1/4}$$

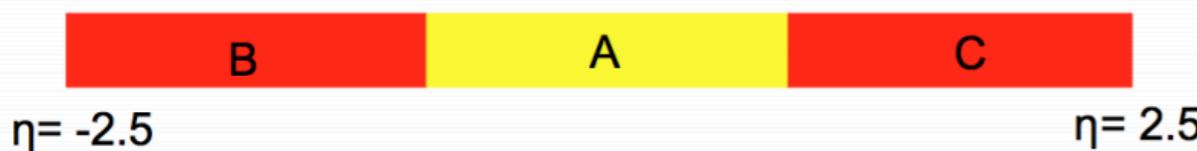
- the variance brings  $\varepsilon_2\{4\}$  down (this term gives the usual  $\sqrt{\nu_2^2 - \sigma^2}$ )
- positive skew brings  $\varepsilon_2\{4\}$  further down, negative skew brings it back up
- kurtosis  $> 2$  brings  $\varepsilon_2\{4\}$  further down, kurtosis  $< 2$  brings it back up
  - recall Gaussian has kurtosis = 3



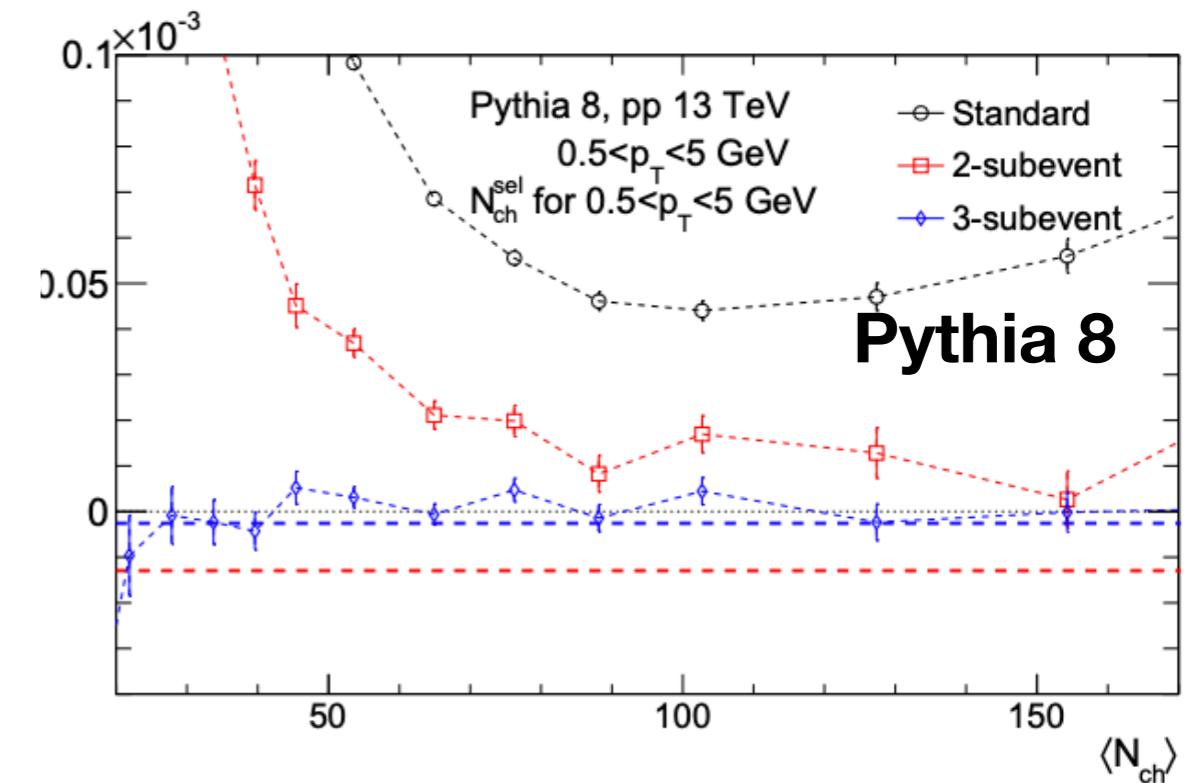
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# Sub-event cumulant method

(a) Three subevents



(b) Two subevents



The second type of four-particle cumulant has a larger contribution from non-flow

**Introduced to further reduce non-flow correlations in multi-particle cumulant in small systems**

**Pythia does not have flow correlations  
—> expectation is  $c_2\{4\} \sim 0$**

**Standard  $c_2\{4\} > 0 \rightarrow$  non-flow remains**

**2-subevent closer to 0 , but still positive**

**3-subevent - close to 0**

Reference: Revealing long-range multi-particle collectivity in small collision systems via subevent cumulants. Phys. Rev. C 96, 034906 (2017)

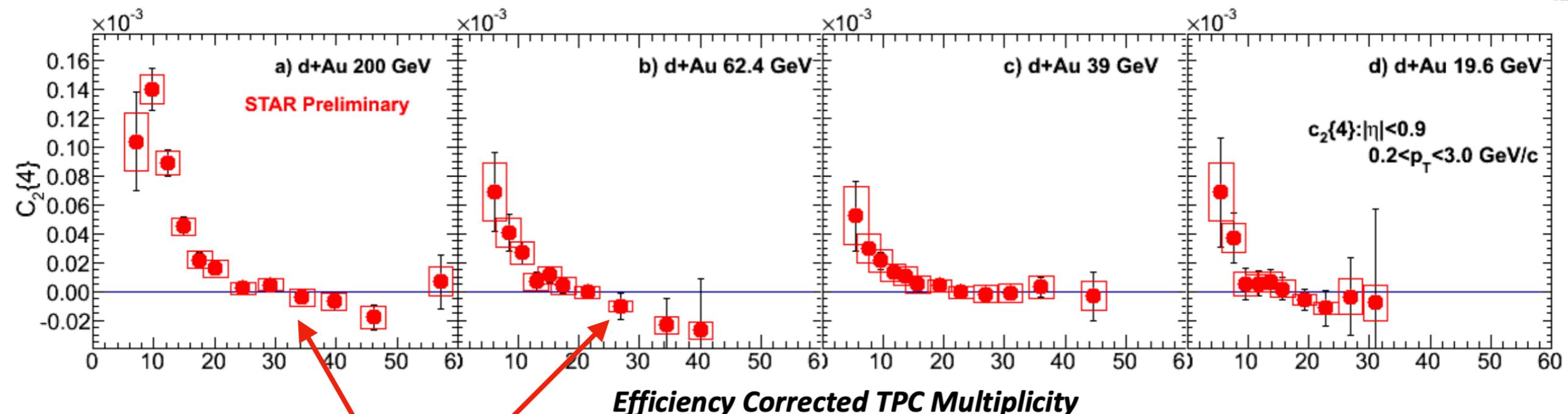


# Small system multi-particle correlations in STAR



$c_2\{4\}$  vs.  $\langle dN/d\eta \rangle$

Standard cumulant



indication of negative  $c_2\{4\}$  in d+Au

200 GeV d+Au 3sub-events

