

Electroweak Boxes & Dispersion Relations

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My pronouns: he/him/his

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BNL, March 2019

Goals For This Talk

- *Provide a brief BSM context & basic formalism*
- *Introduce the dispersion relation framework*
- *Apply the DR framework to nucleon & nuclear contributions to the $W\gamma$ box correction*
- *Discuss possible tests*

Outline

- I. Context*
- II. W_γ Box: Dispersion Relations*
- III. W_γ Box: Free Nucleon*
- IV. W_γ Box: Nuclei*
- V. Electroweak Boxes More Generally*
- VI. Outlook*

I. Context

Weak Decays: CKM Unitarity

$$d \rightarrow u e^- \bar{\nu}_e$$

$$s \rightarrow u e^- \bar{\nu}_e$$

$$b \rightarrow u e^- \bar{\nu}_e$$

$$(u \quad c \quad t) \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

$$\Delta_{\text{CKM}} = (|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2)_{\text{exp}} - 1$$

$$0.94906 \pm 0.00041$$

$$0.05031 \pm 0.00022$$

$$0.00002$$

$$\Delta_{\text{CKM}} = -0.0006 \pm 0.0005$$

Precision ~ BSM Mass Scale

Precision Goal:

$$\delta \Delta_{CKM} \sim O(10^{-4})$$

Heavy BSM Physics:

$$\Delta_{CKM} \sim C \left(v/\Lambda \right)^2$$

$$\Lambda \sim 10 \text{ TeV (tree)}$$

$$\Lambda < 1 \text{ TeV (loop)}$$

Ultralight BSM Physics:


$$\Delta_{CKM} \sim \varepsilon^2 \left(\alpha/4\pi \right)$$

$$\varepsilon < 1 \text{ (loop)}$$


Error Budget

$$\Delta_{\text{CKM}} = (|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2)_{\text{exp}} - 1$$

0.94906 ± 0.00041



0.05031 ± 0.00022



0.00002



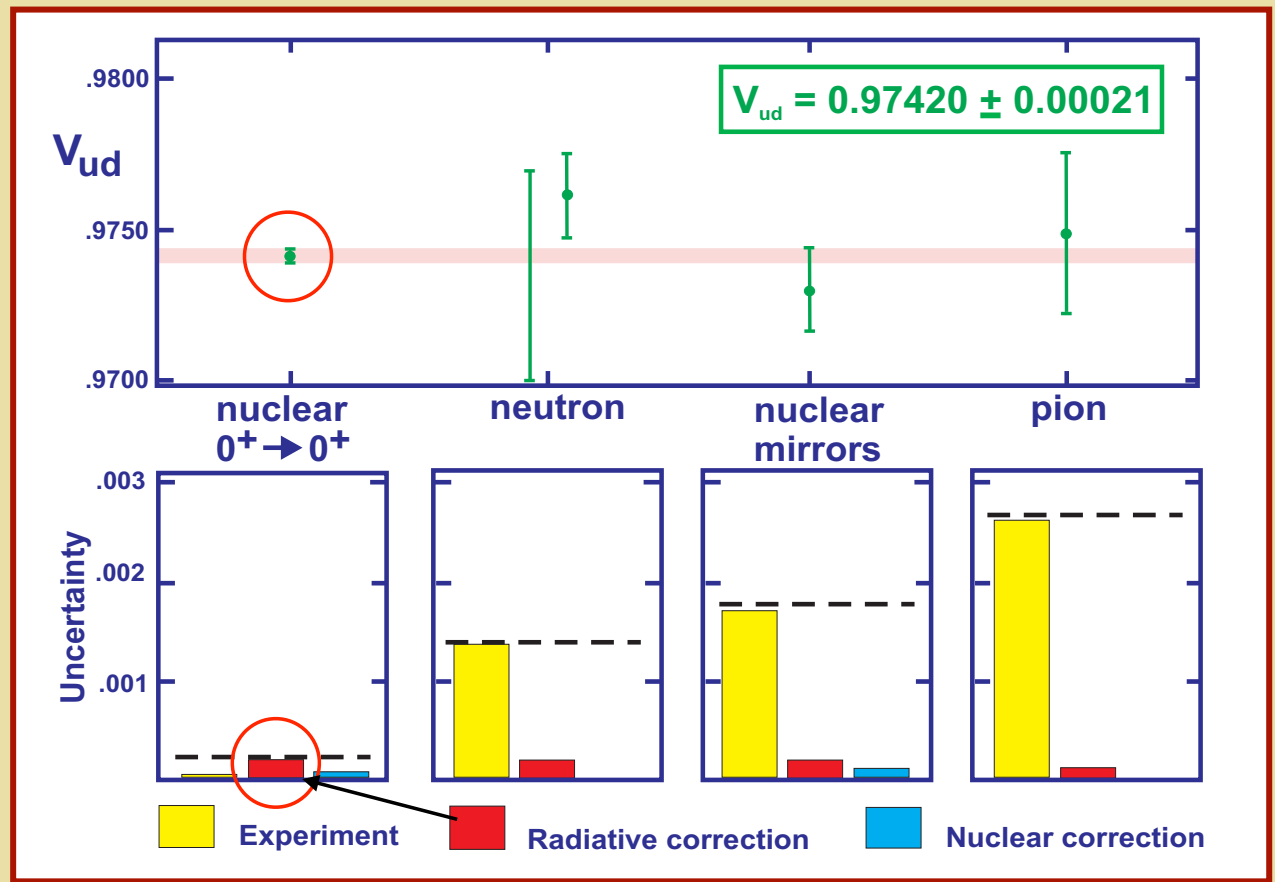
Error Budget

$$\Delta_{\text{CKM}} = (|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2)_{\text{exp}} - 1$$

0.94906 ± 0.00041

Radiative Correction

*Factor of 2 reduction
using disp relations*



Thanks: J. Hardy

Error Budget

$$\Delta_{\text{CKM}} = (|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2)_{\text{exp}} - 1$$

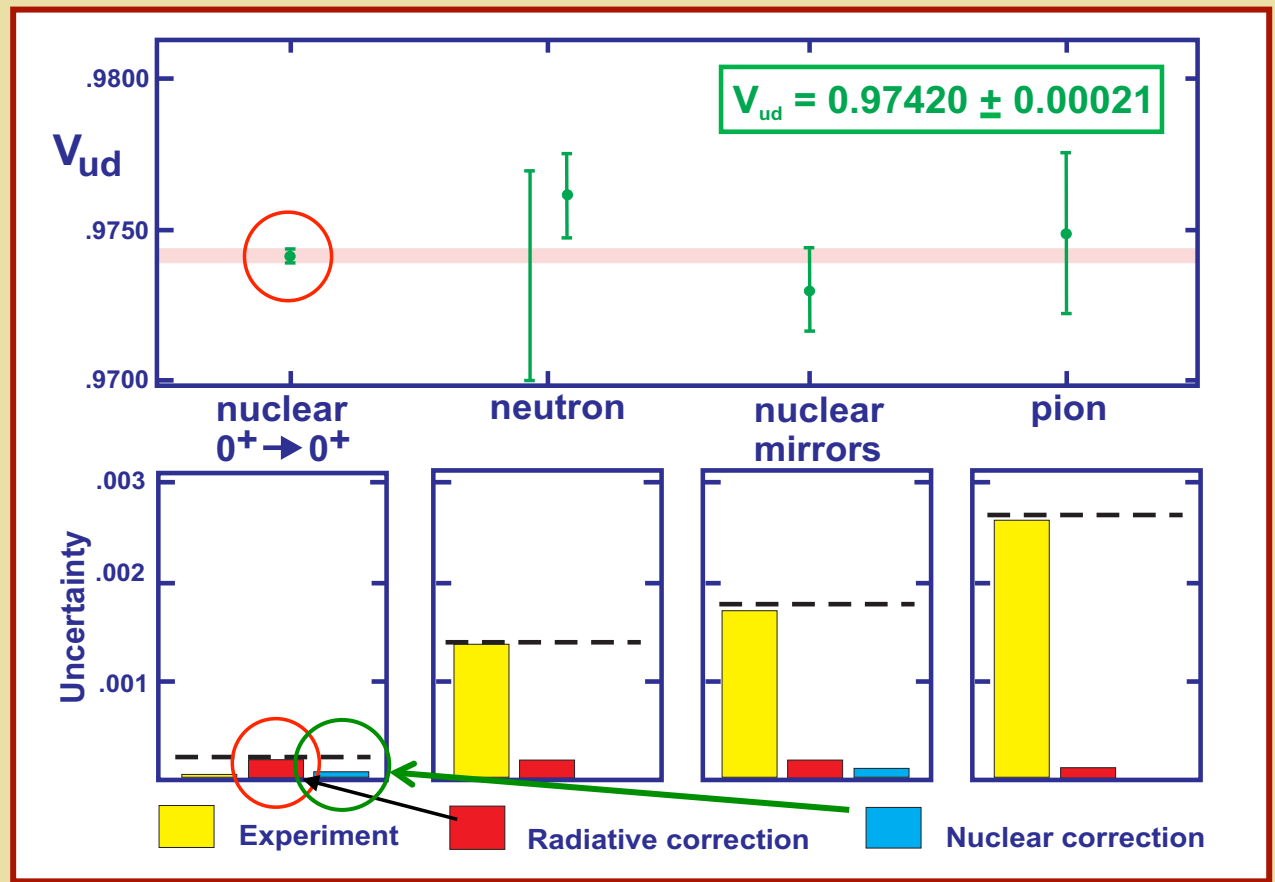
$$0.94906 \pm 0.00041$$

Radiative Correction

*Factor of 2 reduction
using disp relations*

Nuclear Correction

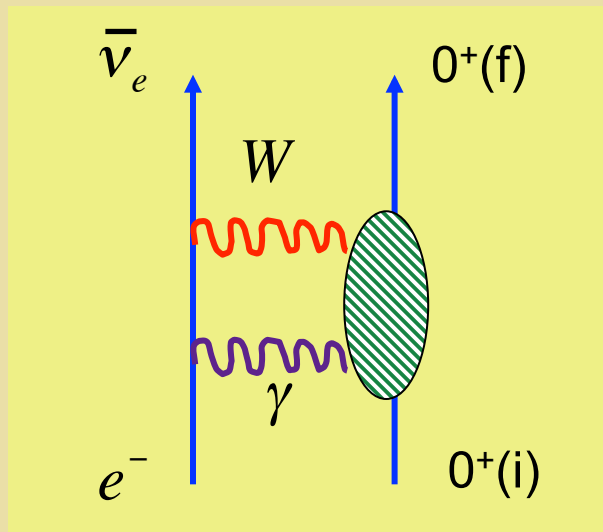
*Increase due to
previously omitted
contributions*



Thanks: J. Hardy

Radiative Corrections

Dominant source of uncertainty:



$$M_{\gamma W} = \frac{G_F V_{ud}}{\sqrt{2}} \frac{\alpha}{8\pi} \left[\ln \left(\frac{M_Z^2}{\Lambda^2} \right) + C_{\gamma W}(\Lambda) \right]$$

Short distance

Long distance

Long distance

Sensitive to hadronic & nuclear dynamics

Radiative Corrections & Ft Values

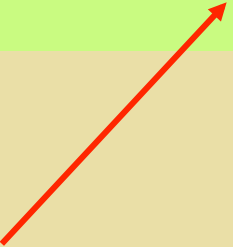
Corrected ft values:

$$\mathcal{F}t = ft (1 + \delta'_R) (1 + \delta_{NS} - \delta_C) (1 + \Delta_E^{NS})$$


Outer
correction




Nuclear struct
part of $M_{\gamma W}$
E-independent



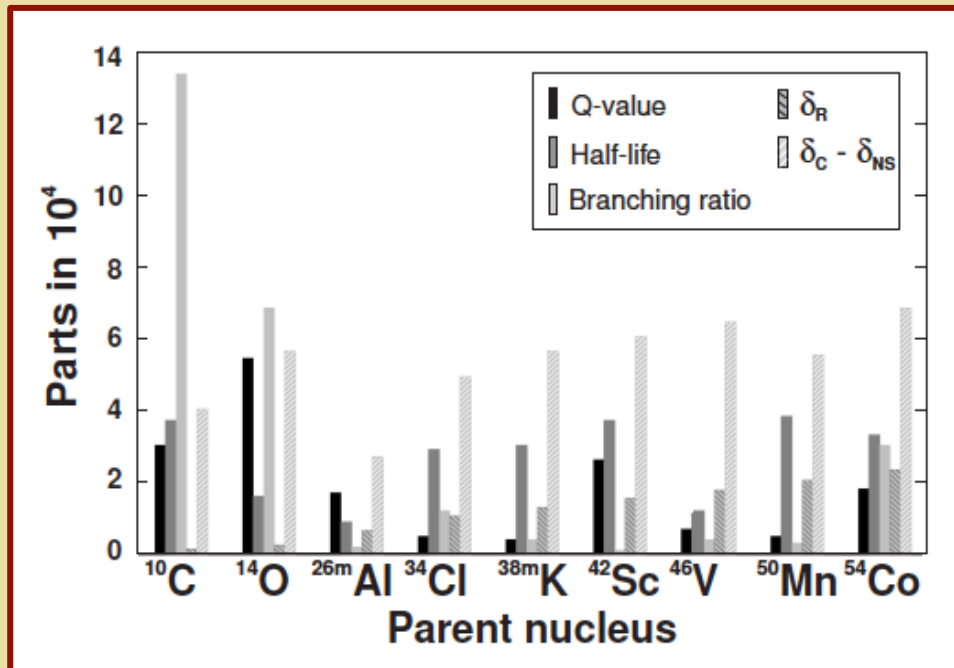
Nucl wavef'n
Not a RC



Nuclear struct
part of $M_{\gamma W}$
E-dependent



$0^+ \rightarrow 0^+$ Dispersion Corrections: δ_{NS}



b_F : scalar currents

Input for V_{ud} & CKM unitarity test

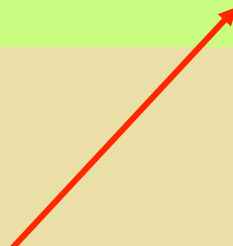
Towner & Hardy, PRC 91 (2015) 2, 025501

Radiative Corrections & V_{ud}

Superallowed

$$|V_{ud}|^2 = \frac{2984.43s}{\mathcal{F}t(1 + \Delta_R^V)}$$

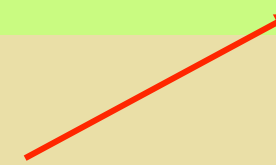
*Hadronic & short
distance part of $M_{\gamma W}$*



Neutron

$$|V_{ud}|^2 = \frac{5099.34s}{\tau_n(1 + 3\lambda^2)(1 + \Delta_R)}$$

$(g_A/g_V)^2$

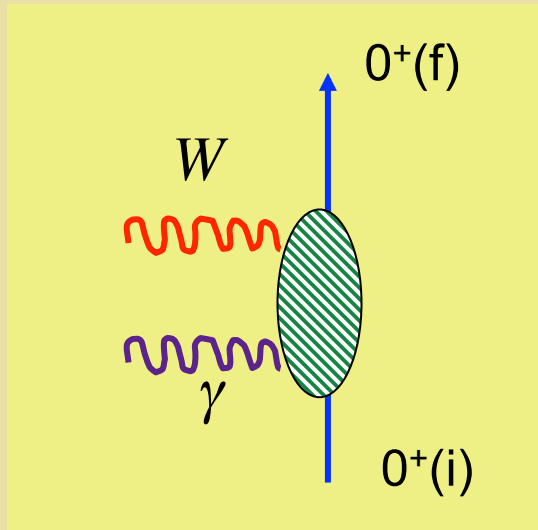


Contains Δ_R^V



II. W_γ Box: Dispersion Relations

Dispersion Relations



Electroweak virtual Compton amplitude:

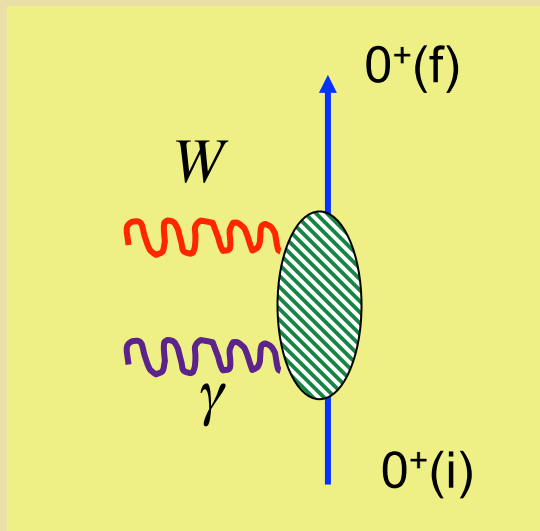
$$T_{\gamma W}^{\mu\nu} = \left[-g^{\mu\nu} + \frac{q^\mu q^\nu}{q^2} \right] T_1 + \frac{\hat{p}^\mu \hat{p}^\nu}{(p \cdot q)} T_2 + \frac{i\epsilon^{\mu\nu\alpha\beta} p_\alpha q_\beta}{2(p \cdot q)} T_3$$

Radiative correction:

$$\square_{\gamma W}^{VA} = 4\pi\alpha \text{Re} \int \frac{d^4 q}{(2\pi)^4} \frac{M_W^2}{M_W^2 + Q^2} \frac{Q^2 + \nu^2}{Q^4} \frac{T_3(\nu, Q^2)}{M\nu}$$

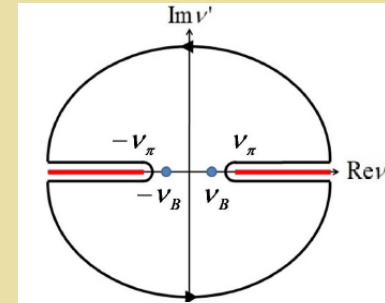
$$\square_{\gamma W}^{VA} = \frac{1}{2} (\Delta_R^V)_{\gamma W}^{VA}$$

Dispersion Relations



Dispersion relation:

Write T_3 as integral over discontinuity along cut



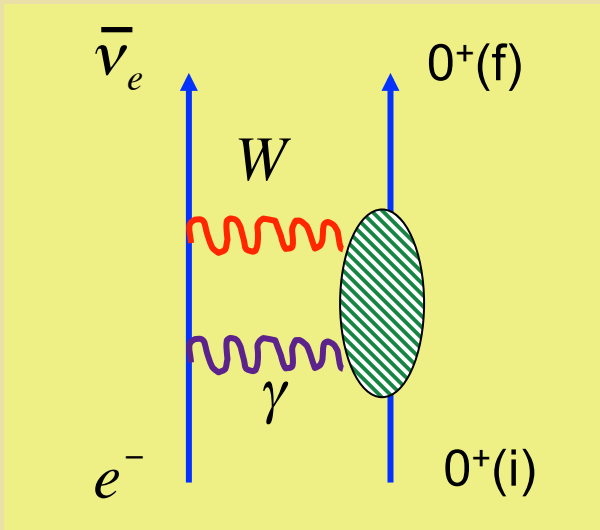
$$T_3^{(I)}(\nu, Q^2) = \frac{2}{i} \int_0^\infty d\nu' \left[\frac{1}{\nu' - \nu} + \frac{\xi^I}{\nu' + \nu} \right] F_3^{(I)}(\nu', Q^2)$$

Electroproduction structure functions:

$$\begin{aligned} W_{\gamma W}^{(I)\mu\nu} &= \frac{1}{8\pi} \sum_X (2\pi)^4 \delta^4(p + q - p_X) \langle p | J_{em}^{(I)\mu} | X \rangle \langle X | J_W^\nu | n \rangle \\ &= \left[-g^{\mu\nu} + \frac{q^\mu q^\nu}{q^2} \right] F_1^{(I)} + \frac{\hat{p}^\mu \hat{p}^\nu}{(p \cdot q)} F_2^{(I)} + \frac{i\epsilon^{\mu\nu\alpha\beta} p_\alpha q_\beta}{2(p \cdot q)} F_3^{(I)} \end{aligned}$$

$$d\sigma \propto L_{\mu\nu} W^{\mu\nu}$$

Dispersion Relations



Radiative Correction:

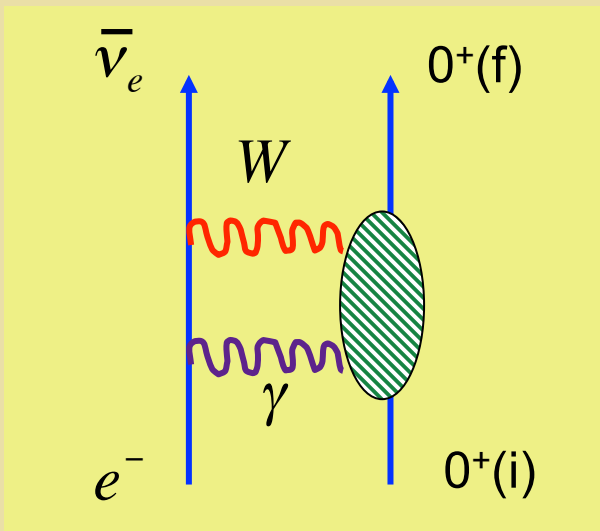
$$\begin{aligned}\square_{\gamma W}^{VA(0)} &= \frac{\alpha}{\pi M} \int_0^\infty \frac{dQ^2 M_W^2}{M_W^2 + Q^2} \int_0^\infty d\nu \frac{(\nu + 2q)}{\nu(\nu + q)^2} F_3^{(0)}(\nu, Q^2) \\ &= \frac{3\alpha}{2\pi} \int_0^\infty \frac{dQ^2 M_W^2}{Q^2 [M_W^2 + Q^2]} M_3^{(0)}(1, Q^2)\end{aligned}$$

Nachtmann Moments:

$$M_3^{(0)}(N, Q^2) = \frac{N+1}{N+2} \int_0^1 \frac{dx \xi^N}{x^2} \left[2x - \frac{N\xi}{N+1} \right] F_3^{(0)}$$

$$\xi = 2x \left(1 + \frac{4M^2 x^2}{Q^2} \right)^{-1}$$

Dispersion Relations



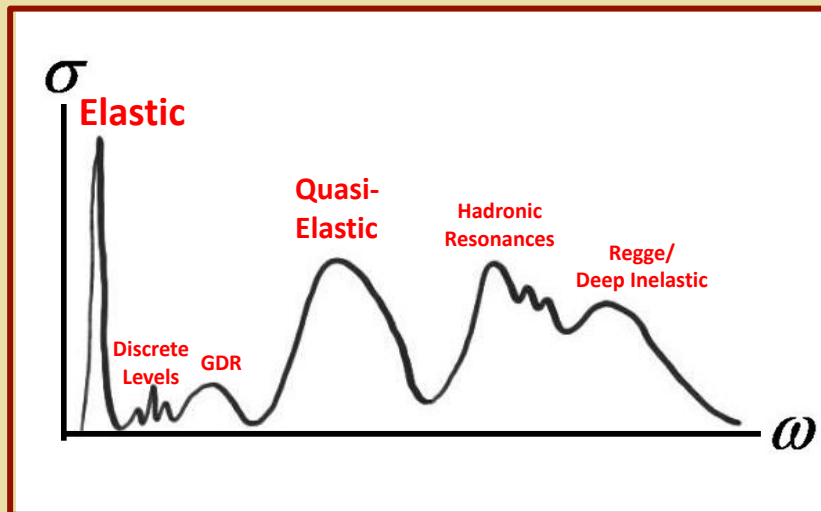
Radiative Correction:

$$\begin{aligned} \square_{\gamma W}^{VA(0)} &= \frac{\alpha}{\pi M} \int_0^\infty \frac{dQ^2 M_W^2}{M_W^2 + Q^2} \int_0^\infty d\nu \frac{(\nu + 2q)}{\nu(\nu + q)^2} F_3^{(0)}(\nu, Q^2) \\ &= \frac{3\alpha}{2\pi} \int_0^\infty \frac{dQ^2 M_W^2}{Q^2 [M_W^2 + Q^2]} M_3^{(0)}(1, Q^2) \end{aligned}$$

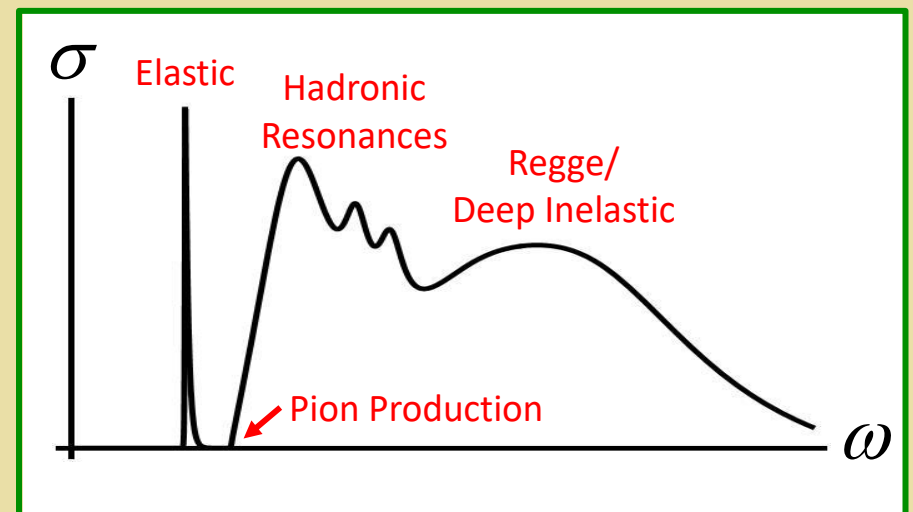
- **Relate $F_3^{(0)}$ and $M_3^{(0)}$ to data and/or**
- **Compute $F_3^{(0)}$ and $M_3^{(0)}$ using same methods used to describe semi-leptonic scattering processes with nucleon & nuclear targets**

Leptoproduction: Had & Nuc Response

Nuclei



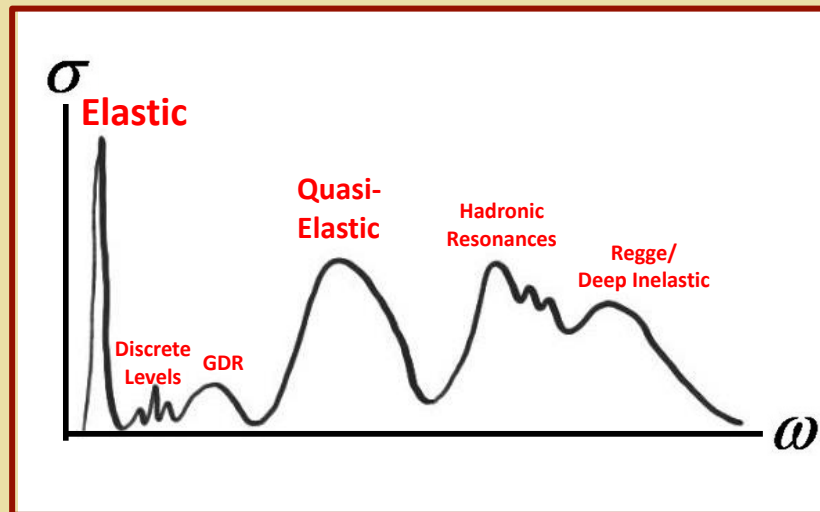
Free nucleons



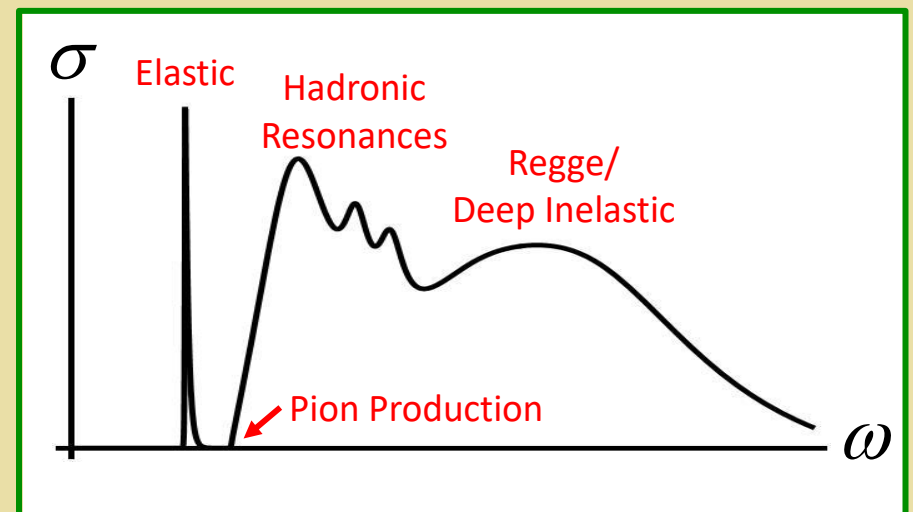
III. W_γ Box: Free Neutron

Leptoproduction: Had & Nuc Response

Nuclei



Free nucleons

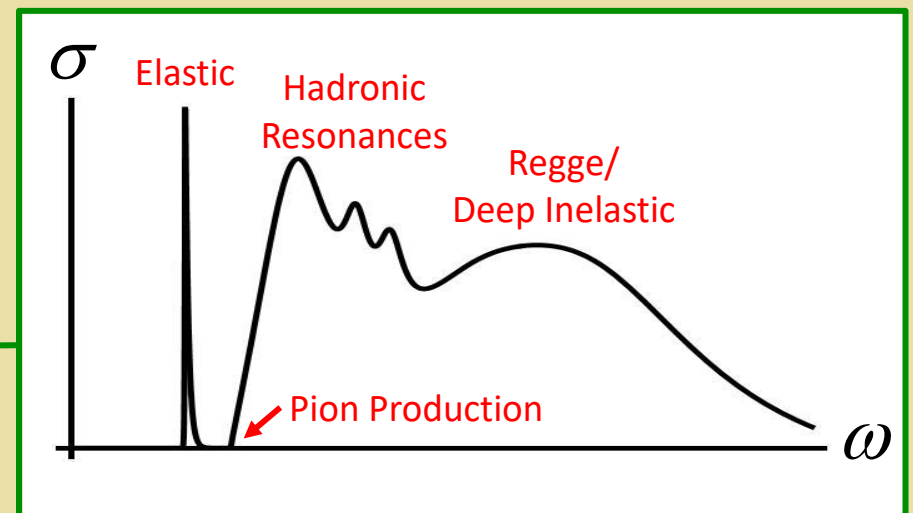
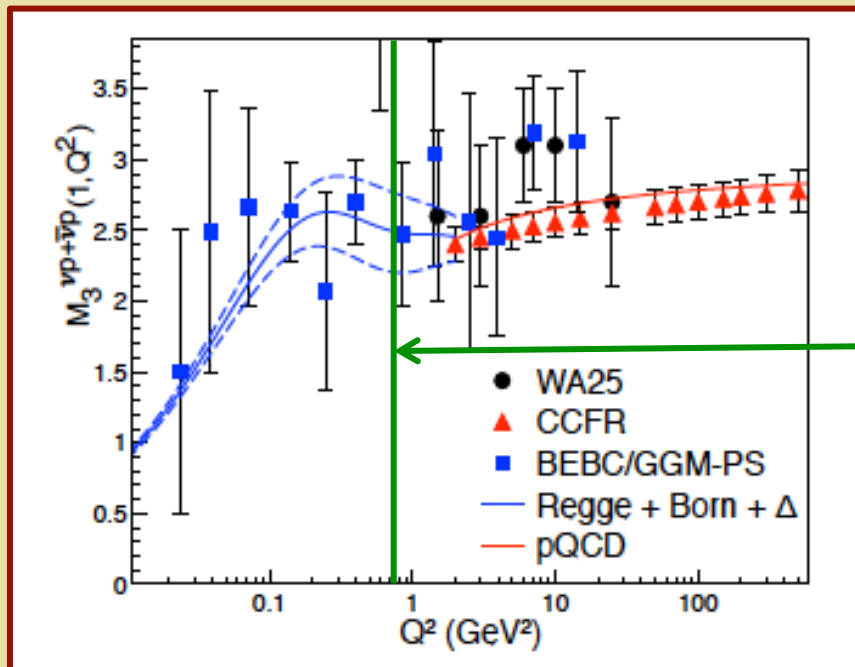


Single nucleon: PRL 121 (2008) 241804

$$\Delta_R^V = 0.02361(38) \rightarrow 0.04267(22)$$

Neutrino Scattering

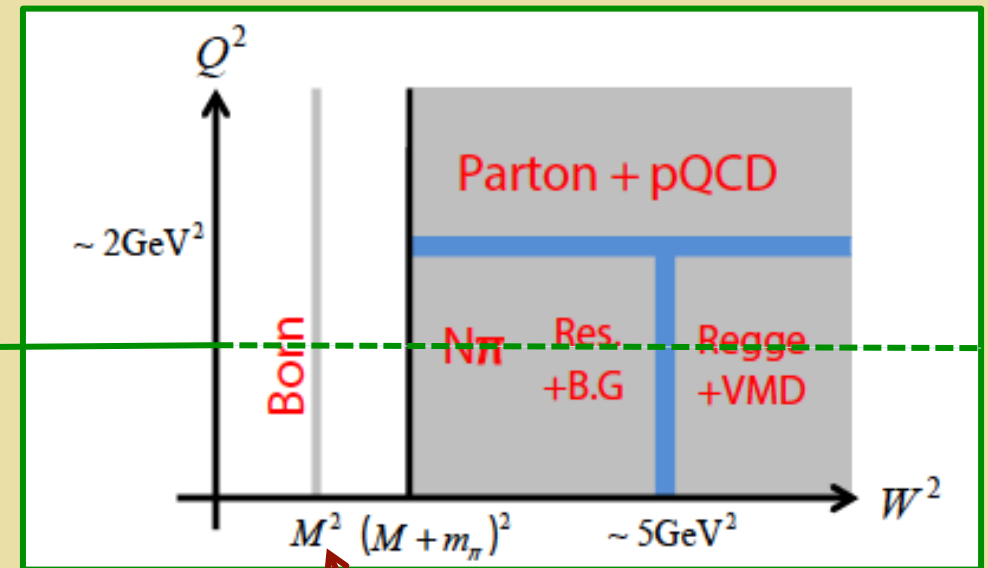
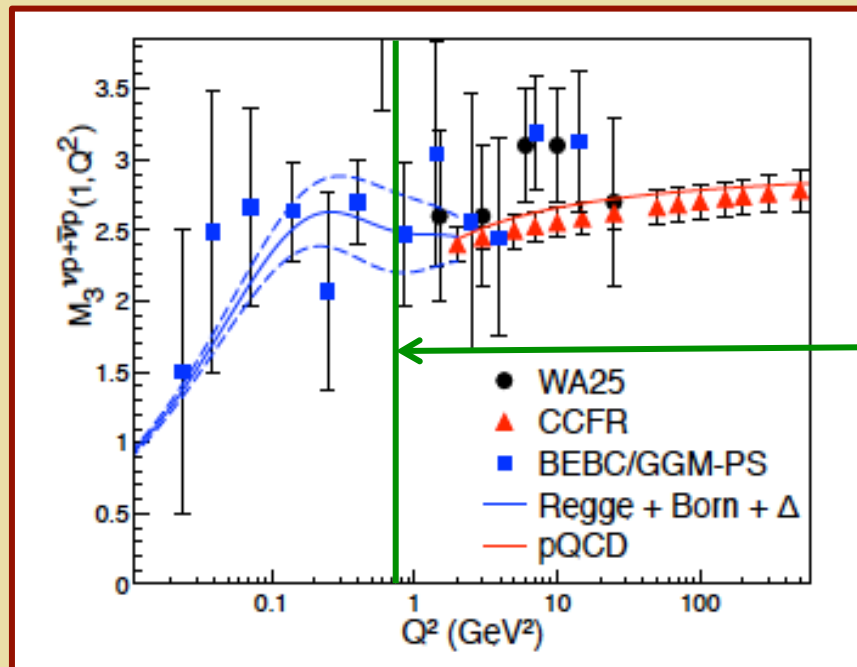
Free nucleons



- Compute contributions to $M_3^{vp+\bar{\nu}p}$ at each Q^2 from different ω regions
- Isospin rotate to $M_3^{(0)}$

Neutrino Scattering

Free nucleons



Illustrate w/ Born

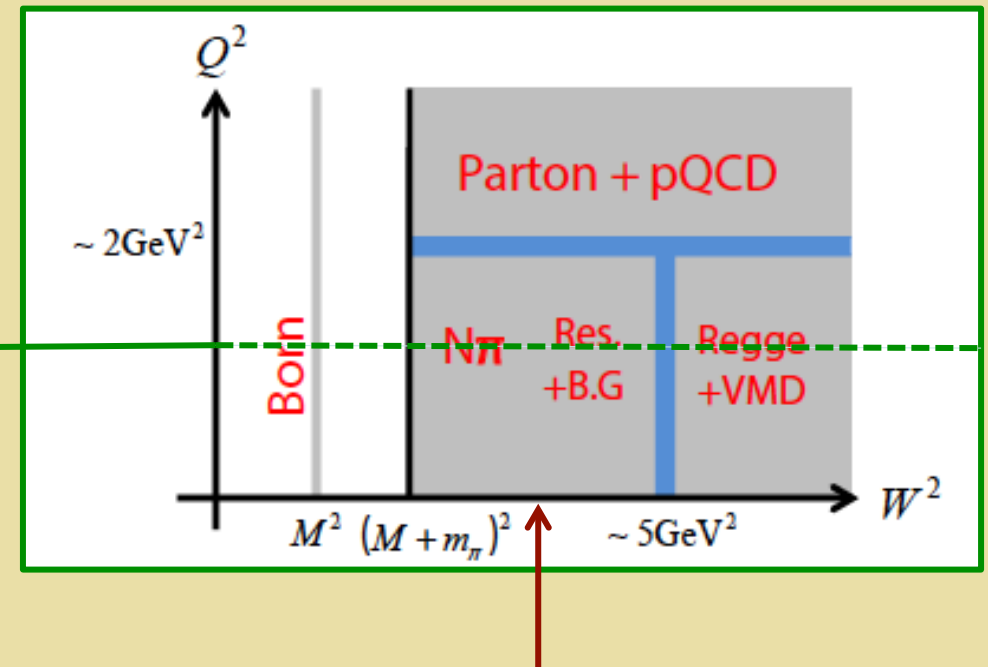
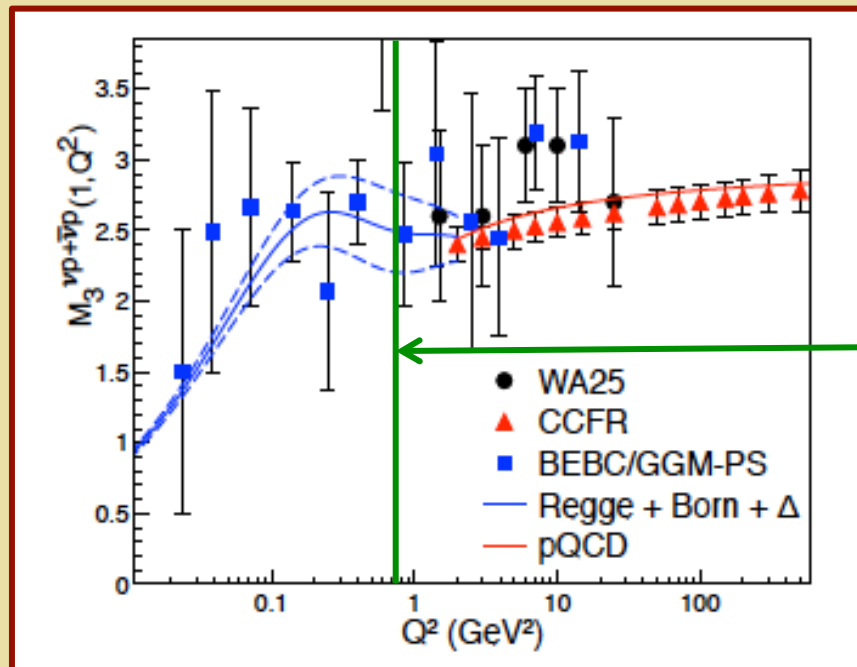
- Compute contributions to $M_3^{vp+\bar{\nu}p}$ at each Q^2 from different ω regions
- Isospin rotate to $M_3^{(0)}$

Born Contribution

$$\begin{aligned} F_{3, \text{Born}}^{\nu p + \bar{\nu} p} &= -G_A(Q^2) G_M^V(Q^2) \delta(1 - x), \\ F_{3, \text{Born}}^{(0)} &= -\frac{1}{4} G_A(Q^2) G_M^S(Q^2) \delta(1 - x), \end{aligned}$$

Neutrino Scattering

Free nucleons

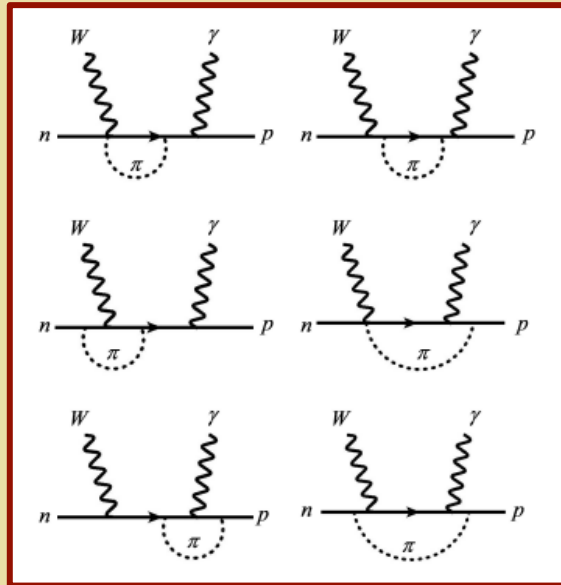


πN + Resonance

- Compute contributions to $M_3^{vp+\bar{\nu}p}$ at each Q^2 from different ω regions
- Isospin rotate to $M_3^{(0)}$

$\pi N + \text{Resonance}$

Non-resonant



$$F_{3,\pi N}^{(0)} = \sum_{i=1}^6 F_{3,i}^{(0)\pi N}$$

Resonant

$\Delta(1232): \nu p \text{ only}$

Form factors

$$F_{3,\Delta}^{\nu p + \bar{\nu} p} = -\frac{2\nu}{M} \frac{m_{\Delta} \Gamma_{\Delta}}{\pi} \frac{1}{(W^2 - m_{\Delta}^2)^2 + m_{\Delta}^2 \Gamma_{\Delta}^2} \frac{V_3}{3}$$

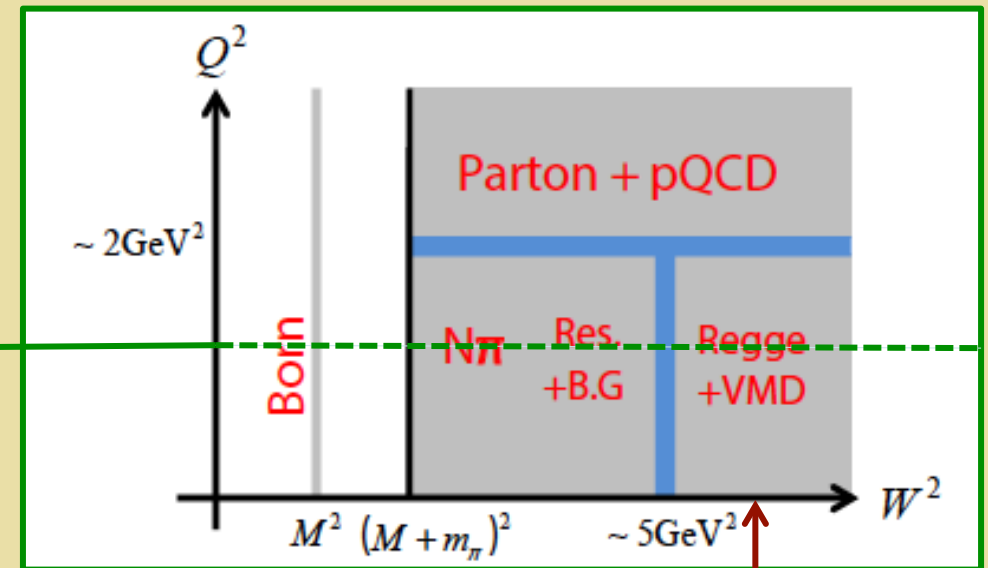
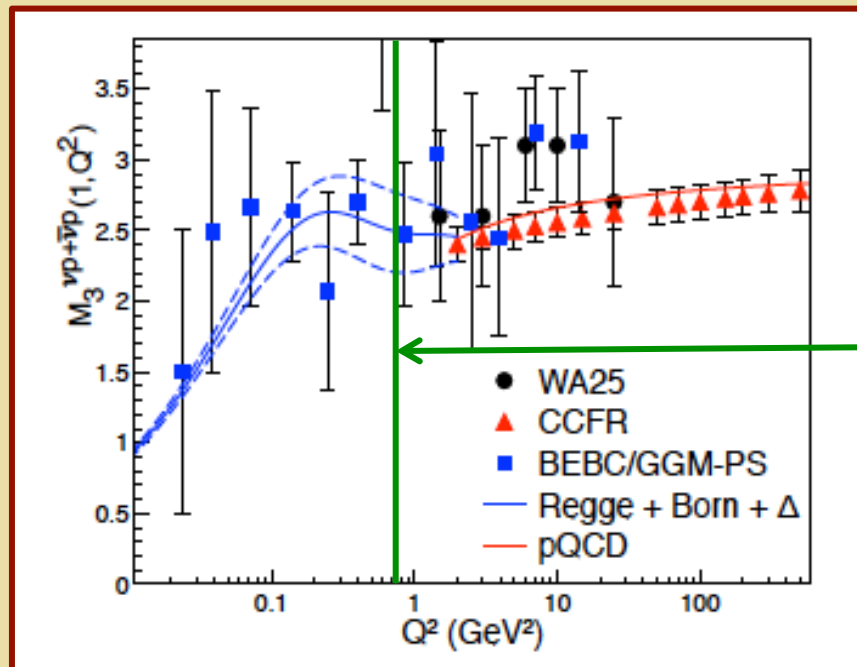
$$F_{3,\text{res}} = \sum_R \frac{\frac{\nu}{q} \Gamma_R m_R}{(W^2 - m_R^2)^2 + \Gamma_R^2 m_R^2} \sqrt{\frac{M(m_R^2 - M^2)}{16\pi^3 \alpha}} \times \sum_{J_z=1/2, 3/2} (A_{em,J_z}^{R,p} + A_{em,J_z}^{R,n})^* A_{w,J_z}^R$$

$F_3^{\nu p} : \text{both } 1/2, 3/2$

$F_3^{(0)} : \text{only } 1/2$

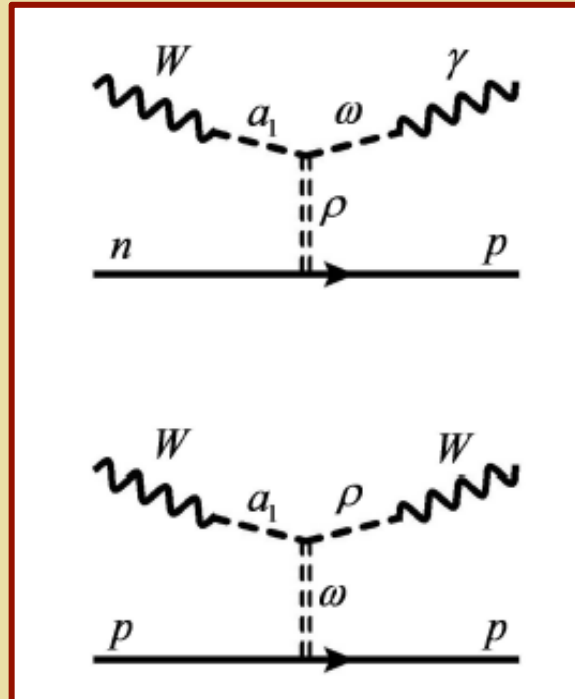
Neutrino Scattering

Free nucleons



- Compute contributions to $M_3^{vp+\bar{\nu}p}$ at each Q^2 from different ω regions
- Isospin rotate to $M_3^{(0)}$

Regge Contribution



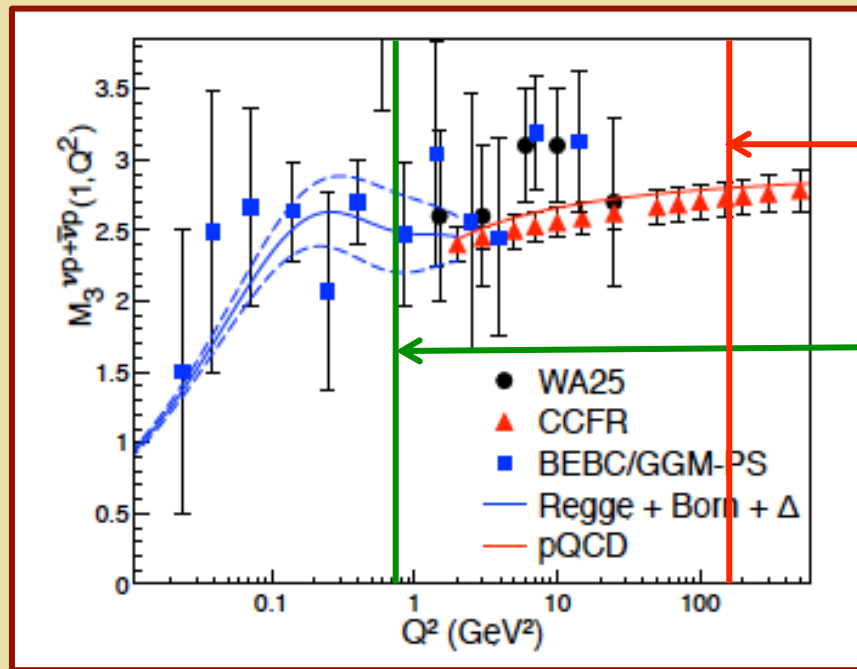
$$F_{3,\mathbb{R}}^{\nu p + \bar{\nu} p} = \frac{C(Q^2) f_{th}(W)}{[1 + Q^2/m_\rho^2] [1 + Q^2/m_{a_1}^2]} \left(\frac{\nu}{\nu_0} \right)^{\alpha_0}$$

$$\downarrow$$

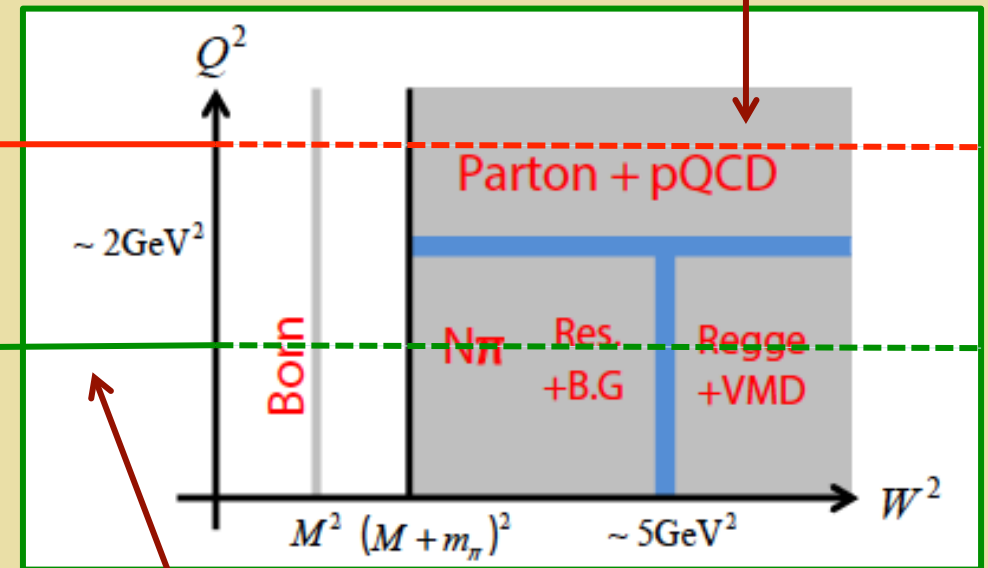
$$F_{3,\mathbb{R}}^{(0)} = \frac{1}{36} \frac{C_{\gamma W}(Q^2) f_{th}(W)}{[1 + Q^2/m_\rho^2] [1 + Q^2/m_{a_1}^2]} \left(\frac{\nu}{\nu_0} \right)^{\alpha_0}$$

- **Matching at $Q^2 = 0$ and $Q^2 = 2 \text{ (GeV)}^2$ [pQCD regime] $\rightarrow C_{\gamma W}(Q^2) = C(Q^2)$**
- **Factor of 1/36: matching at pQCD scale**

Neutrino Scattering



Free nucleons



- Compute contributions to $M_3^{vp+\bar{\nu}p}$ at each Q^2 from different ω regions
- Isospin rotate to $M_3^{(0)}$

Born + (πN + Res) + Regge

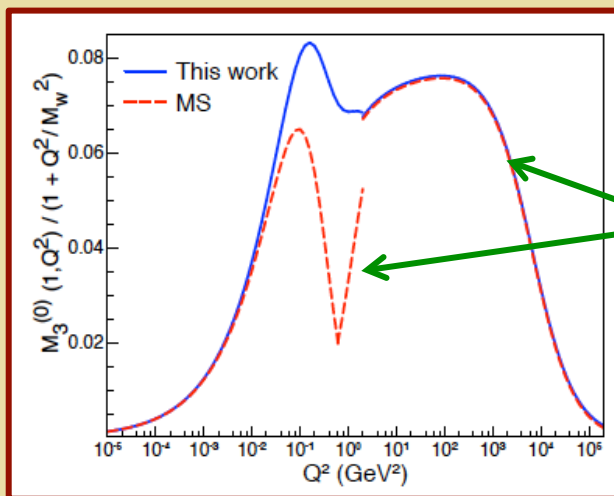
W_γ Box: Update from 2006

$$\square_{\gamma W}^{VA} = \frac{\alpha}{2\pi} [C_{DIS} + C_B + C^{Regge} + C^{\pi N} + C^{Res}]$$

$$C_{DIS}^{MS} = 1.84 \rightarrow C_{DIS}^{new} = 1.87$$

$$C_B^{MS} = 0.829(83) \rightarrow C_B^{new} = 0.91(5)$$

$$C_{INT}^{MS} = 0.14(14) \rightarrow C^{Regge} + C^{\pi N} + C^{Res} = 0.48(7)$$



$$F_{MS}(Q^2) = \frac{12}{Q^2} M_3^{(0)}(1, Q^2)$$

***See W. Marciano
talk today***

Future Tests

- *Lattice computation of $M_3^{(0)}(Q^2)$*
- *PV electron scattering*

Isospin relation

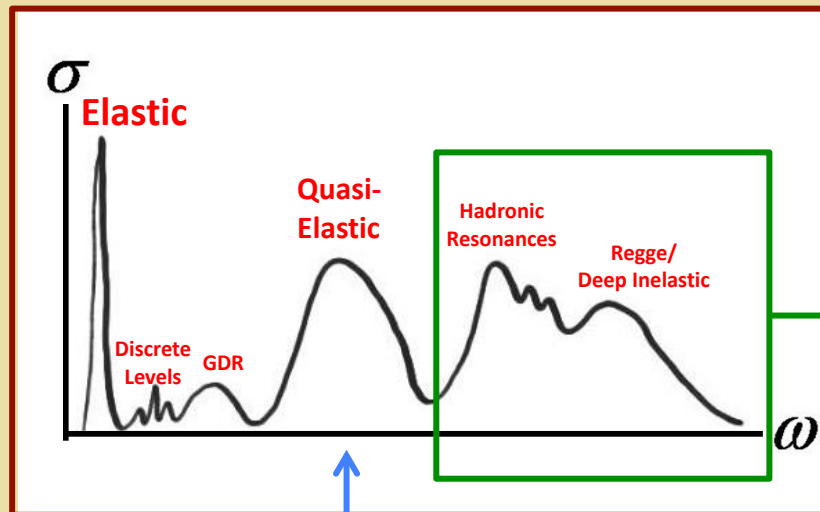
$$4F_3^{(0)} = F_{3,\gamma Z}^p - F_{3,\gamma Z}^n$$

- ***SoLID ?***
- ***EIC ?***

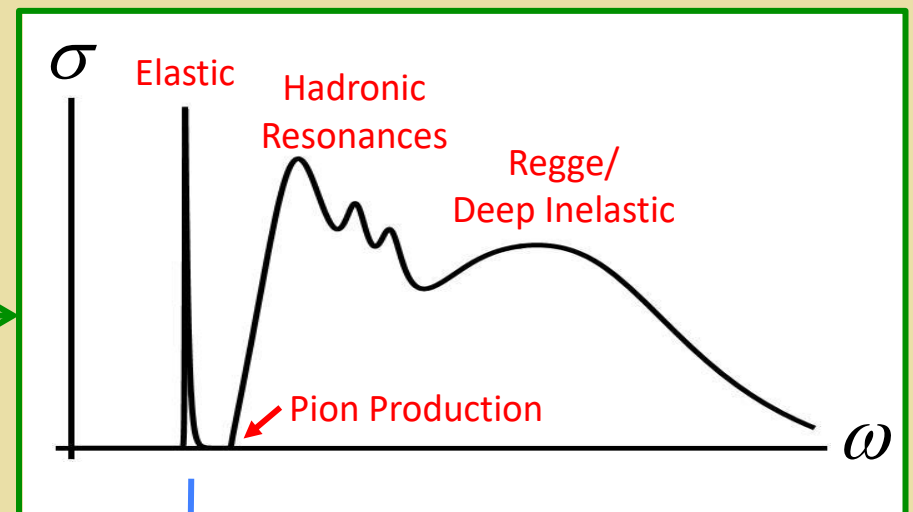
IV. W_γ Box: Nuclei

Leptonproduction: Had & Nuc Response

Nuclei



Free nucleons

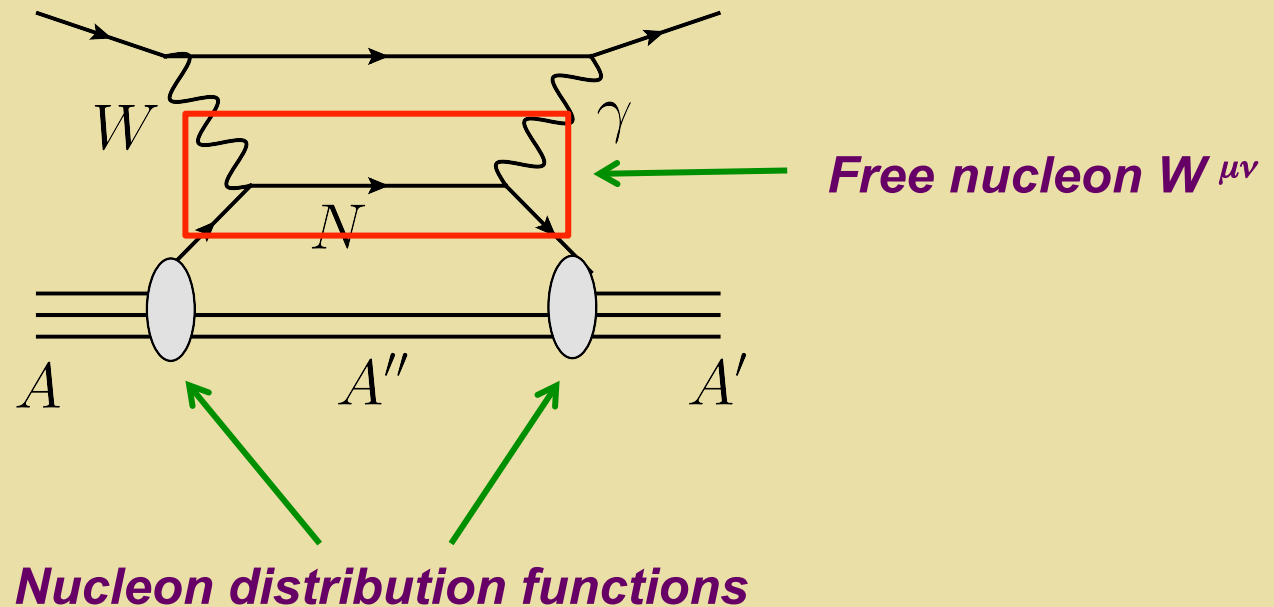


Quasielastic response

Part of δ_{NS} : “ C_B^{Nucl} ”

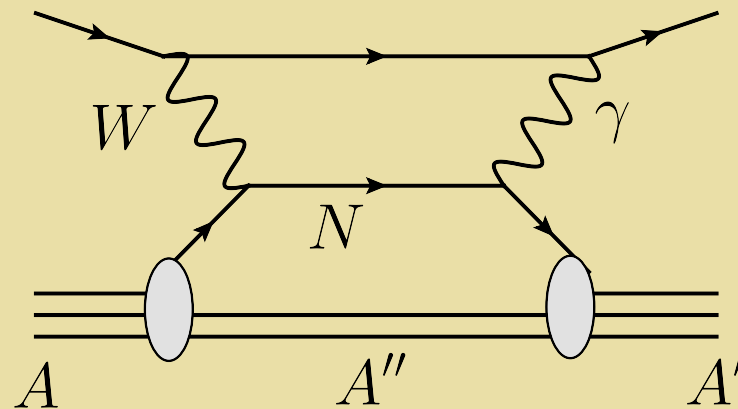
New work

Quasielastic Contribution to δ_{NS}



$$W^{\mu\nu} = \frac{3\mathcal{N}m_N^2}{4\pi p_F^3} \int \frac{d^3p}{E(p)E(p+q)} \delta\{\omega - [E(p+q) - E(p)]\} \\ \times \theta(p_F - |p|) \theta(|p+q| - p_F) f^{\mu\nu}(P+Q, P),$$

Quasielastic Contribution to δ_{NS}



$$\delta_{NS}^{QE} = \frac{\alpha}{\pi} C_{QE}$$

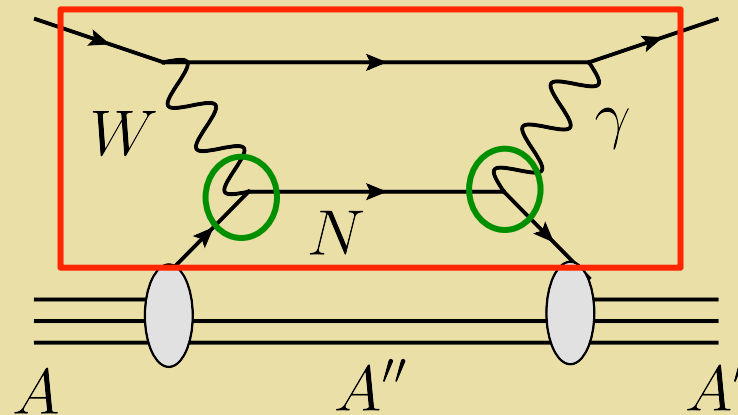
$$C_{QE} = 2 \int_0^\infty dQ^2 \int_{\nu_{min}}^{\nu_\pi} \frac{d\nu(\nu + 2q)}{M\nu(\nu + q)^2} F_{3,\gamma W}^{(0),QE}(\nu, Q^2)$$

$$\frac{1}{Z} F_{3, \gamma W}^{(0), QE}(\nu, Q^2) = -G_A G_M^S \frac{3Q^2}{32q} F_P \frac{\left((\tilde{k}_+)^2 - (\tilde{k}_-)^2 \right)}{k_F^3} \quad \leftarrow \text{Functions of B.E., } M_A \dots$$

Nucleon form factors

Pauli blocking

TH: Nucleon Born Contribution to δ_{NS}

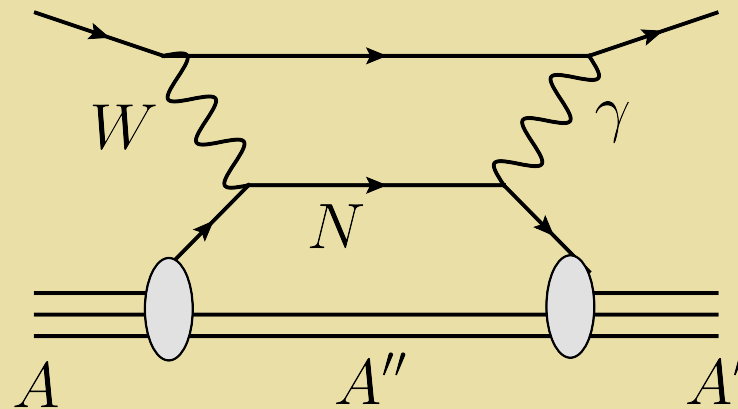


TH** approach: $N \times$ free nucleon loop computation but with “quenched” nucleon ff’s

$$\begin{aligned} G_M^S &\longrightarrow q_S^{(0)} G_M^S \\ G_A &\longrightarrow q_A G_A \end{aligned}$$

***** Towner & Hardy***

Impact on δ_{NS}



$$\Delta \delta_{NS} = \frac{\alpha}{\pi} \left(C_{QE} - q_S^{(0)} q_A C_B \right) = -(4.6 \pm 0.9) \times 10^{-4}$$

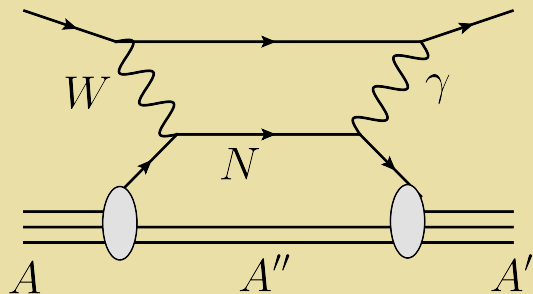
**Neglects A-
dep variations**

**Ave over 20
transitions**

Optimistic:

- Correlations
- 2-body currents
- Rel corrections

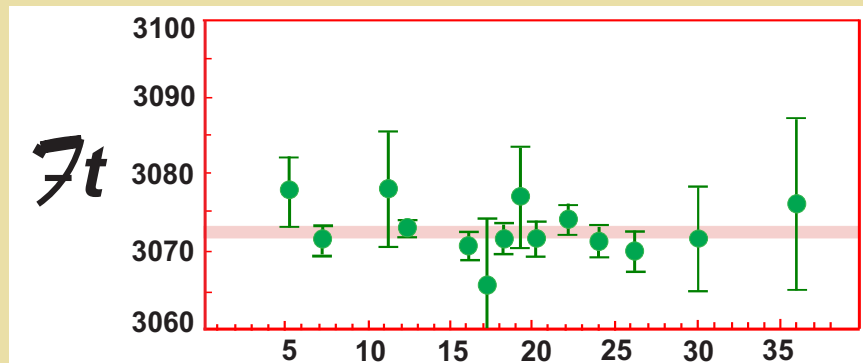
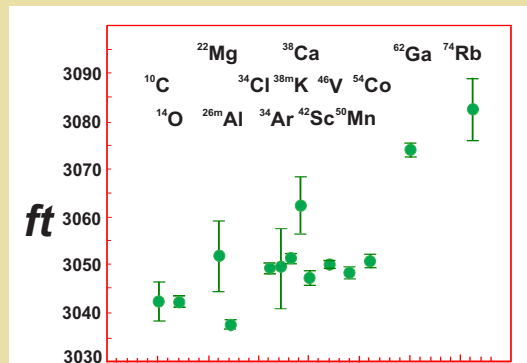
Refinements



Apply state-of-art methods

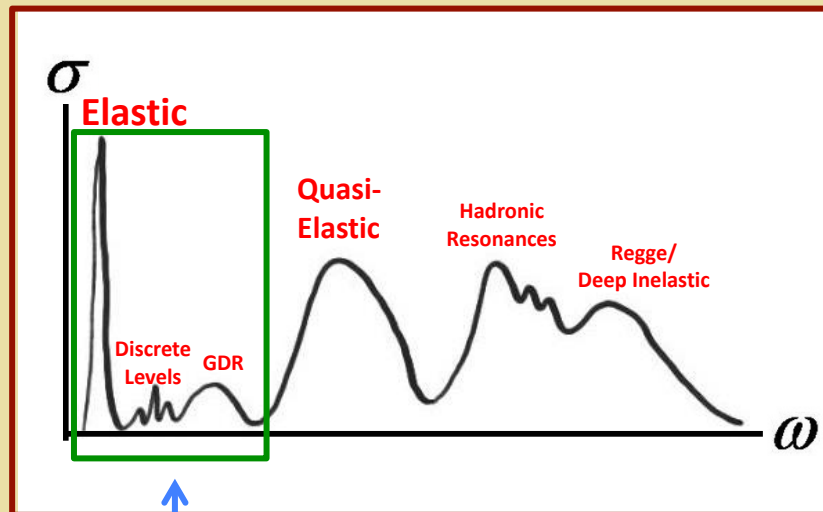
- Correlations
- 2-body currents
- Rel corrections

Consistency w/ CVC ?

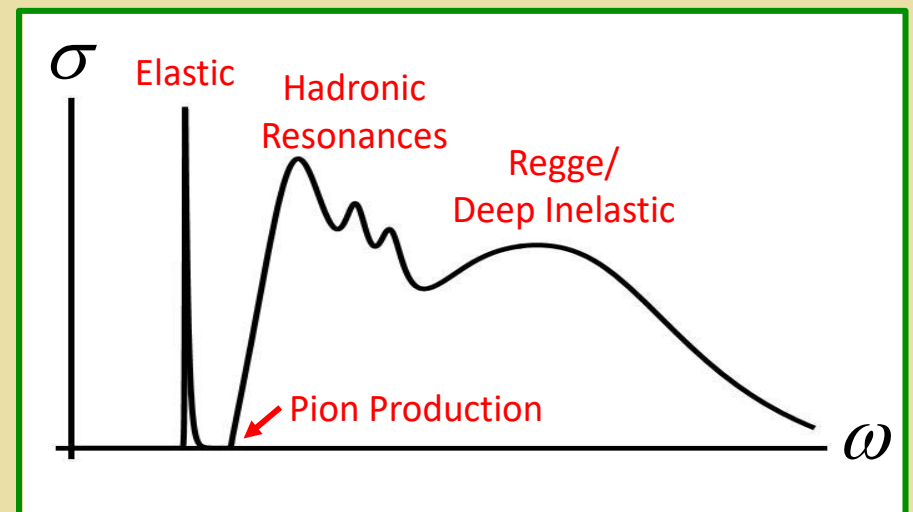


Other Nuclear Corrections

Nuclei



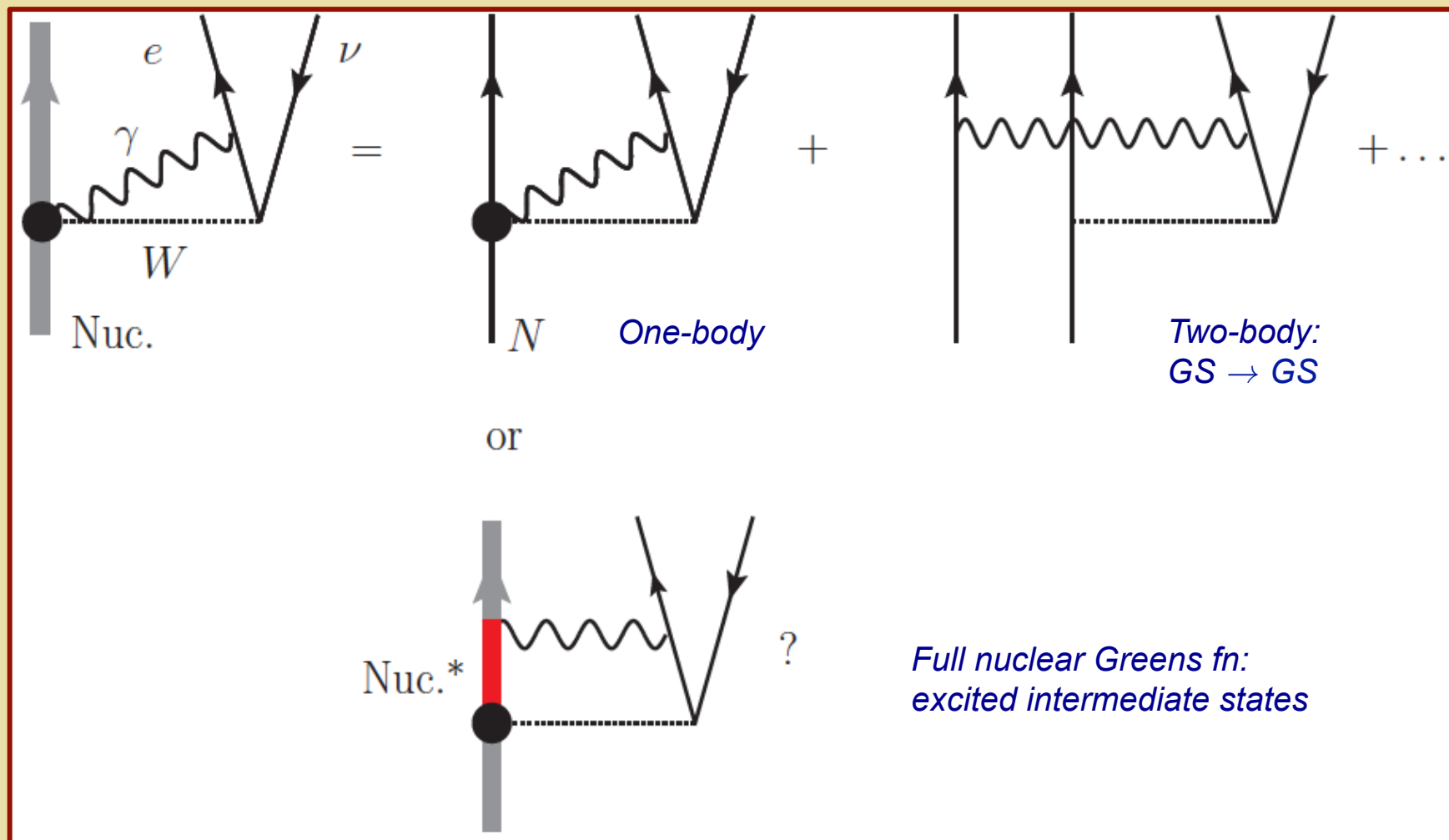
Free nucleons



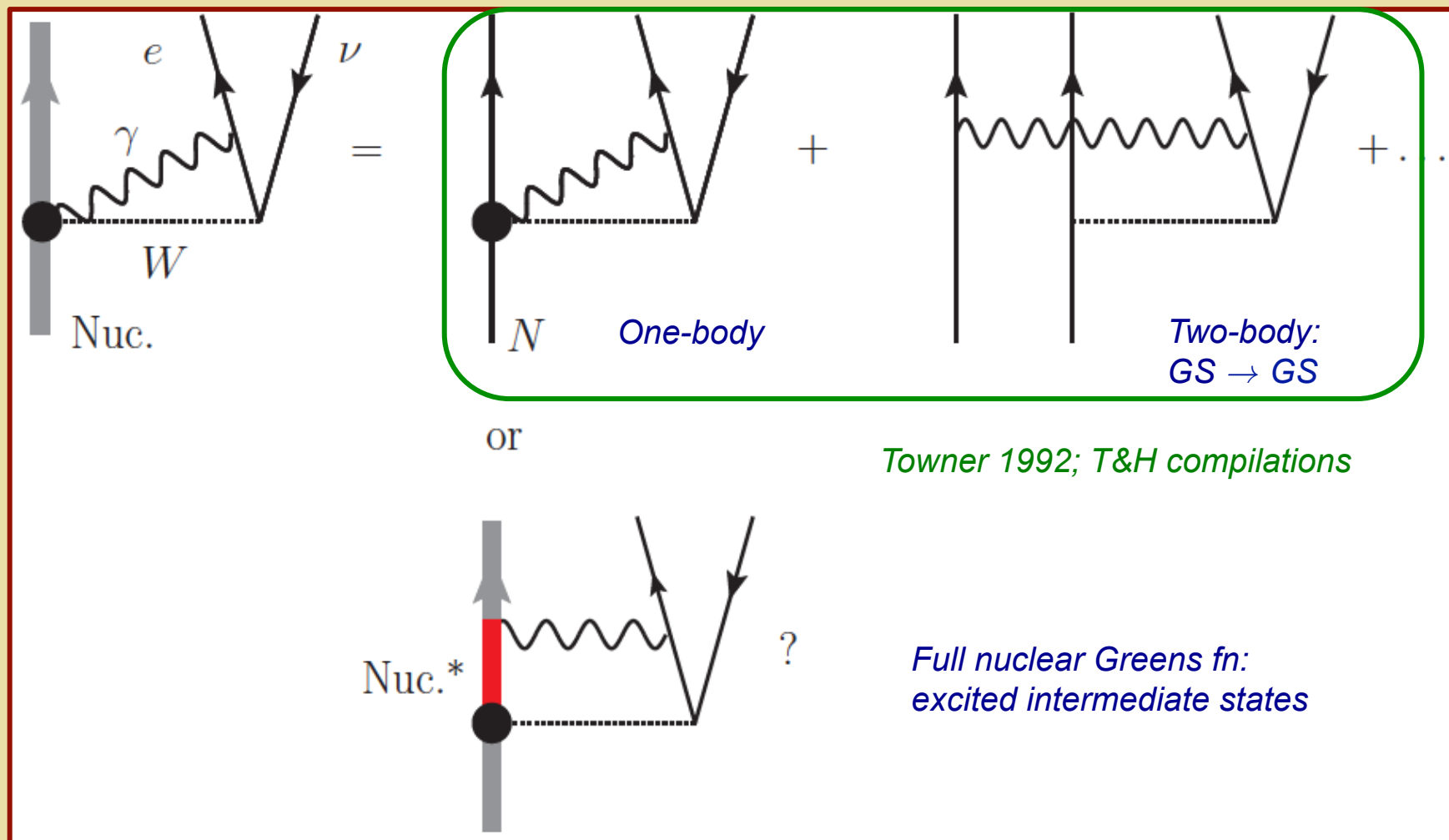
Low-lying transitions

Part of δ_{NS}

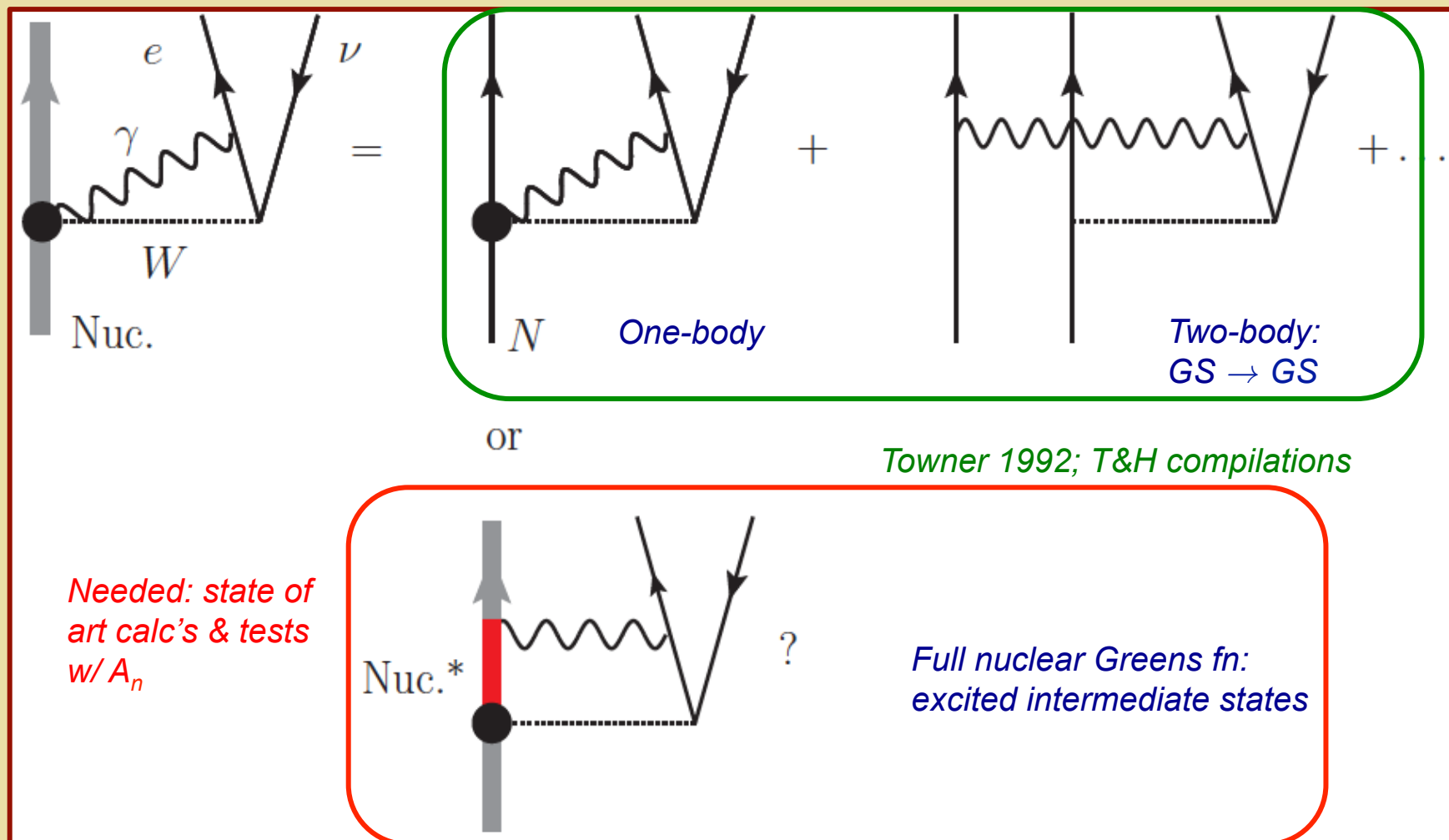
$0^+ \rightarrow 0^+$ Decay: δ_{NS}



$0^+ \rightarrow 0^+$ Decay: δ_{Ns}



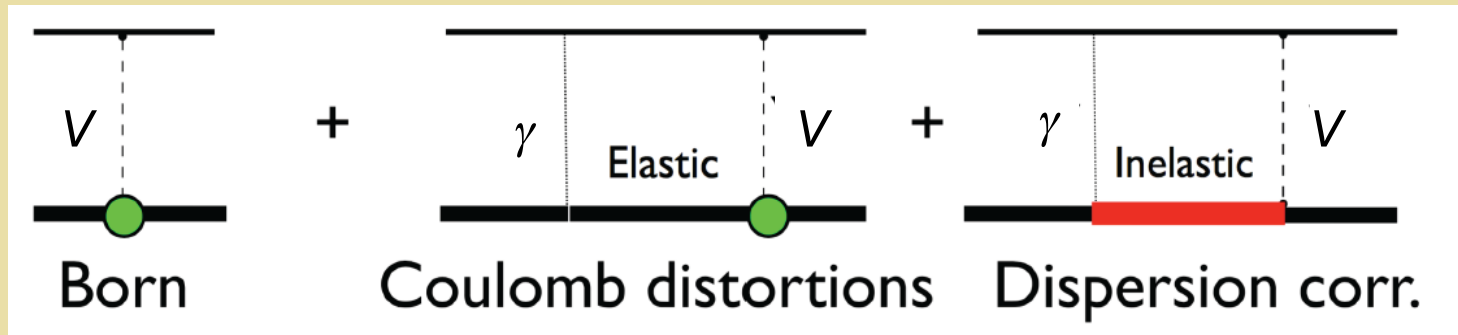
$0^+ \rightarrow 0^+$ Decay: δ_{Ns}



V. EW Boxes More Generally

Dispersion Corrections

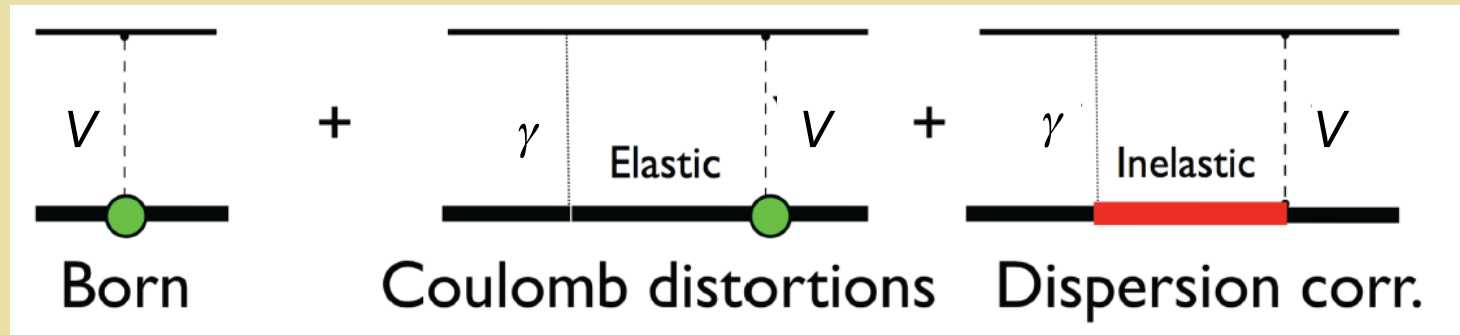
Two-boson exchange in semileptonic processes: important for elastic PV eN & eA scattering (^{12}C) & nuclear β -decay; beam normal asymmetry, Olympus... provide tests



$$V = Z^0, W, \gamma$$

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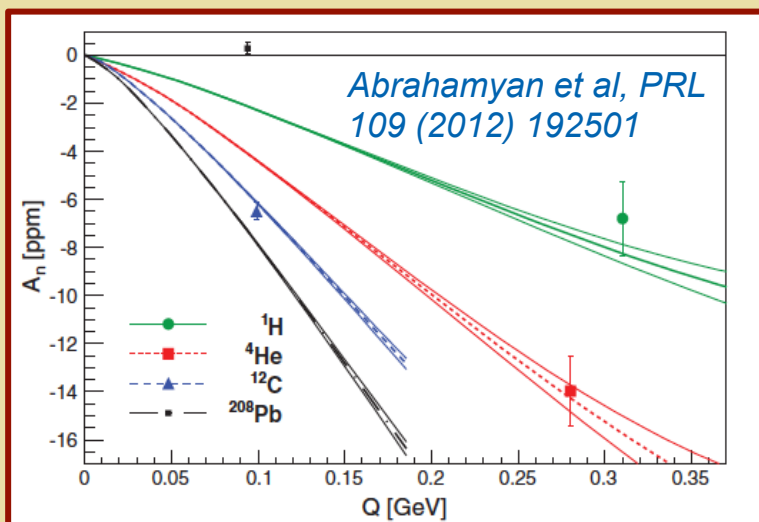
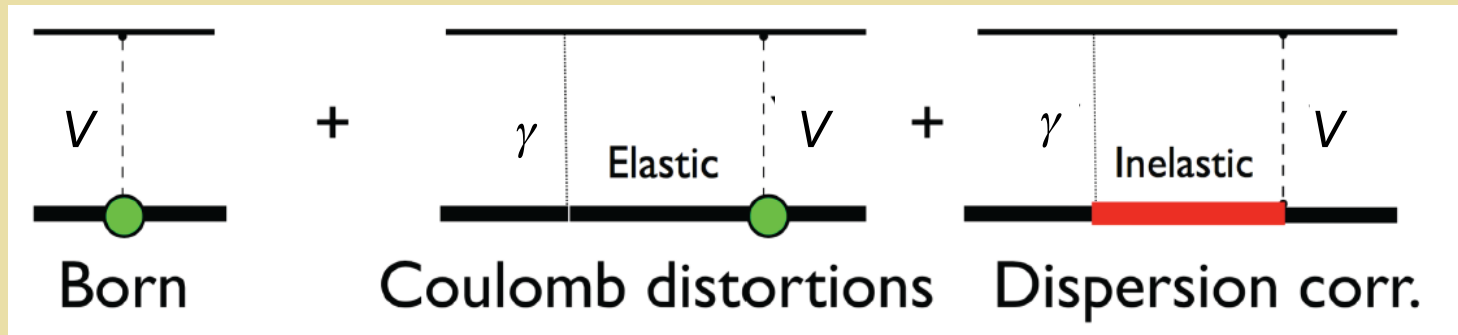
$$V = Z^0, W, \gamma$$



$V = \gamma$ *Beam normal asymmetry*

Dispersion Corrections

Two-boson exchange in semileptonic processes: important for elastic PV eN & eA scattering (^{12}C) & nuclear β -decay; beam normal asymmetry, Olympus... provide tests

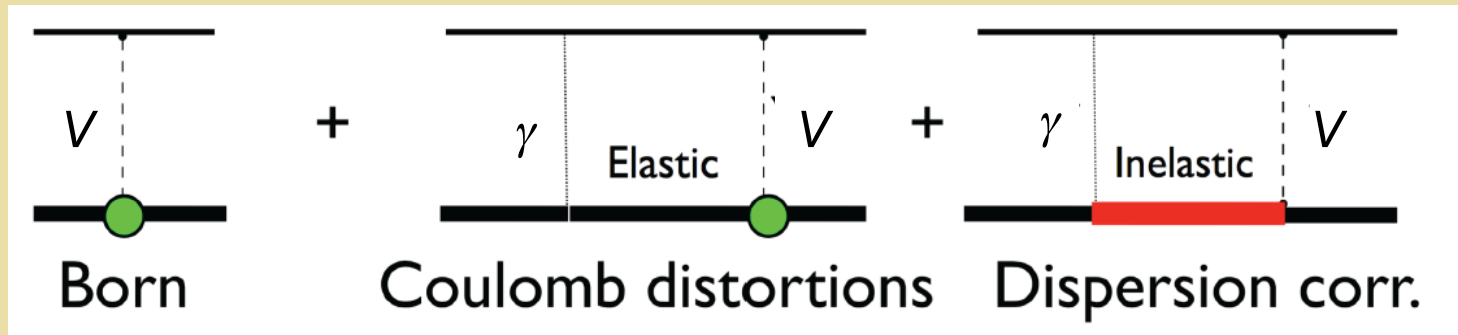


$V = \gamma$ *Beam normal asymmetry*

- J Lab Hall A
- Future: Mainz, J Lab

Dispersion Corrections

Two-boson exchange in semileptonic processes: important for elastic PV eN & eA scattering (^{12}C) & nuclear β -decay; beam normal asymmetry provides, Olympus... provide tests



$$V = Z^0, W, \gamma$$

Important for $O(0.1\%)$ probes of PV $^{12}\text{C}(e,e')$ & superallowed β -decay

$$V = \gamma$$

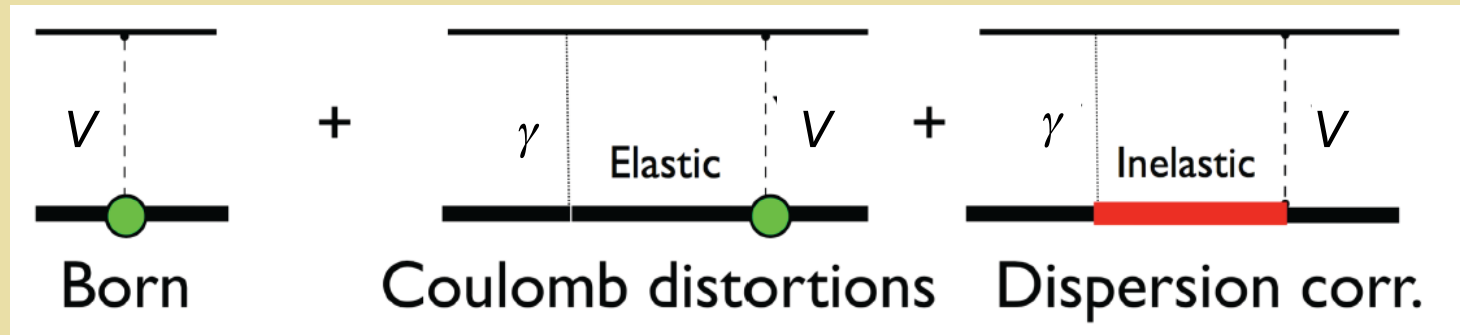
Beam normal asymmetry

$$V = Z^0, W$$

Nucleus-dependent QED & EW corrections

Dispersion Corrections

Proposal: (1) carry out a consistent set of computations for A_n , PV asymmetry, & δ_{NS} using different methods (2) develop a program of A_n measurements to test computations



$$V = Z^0, W, \gamma$$

$$V = \gamma$$

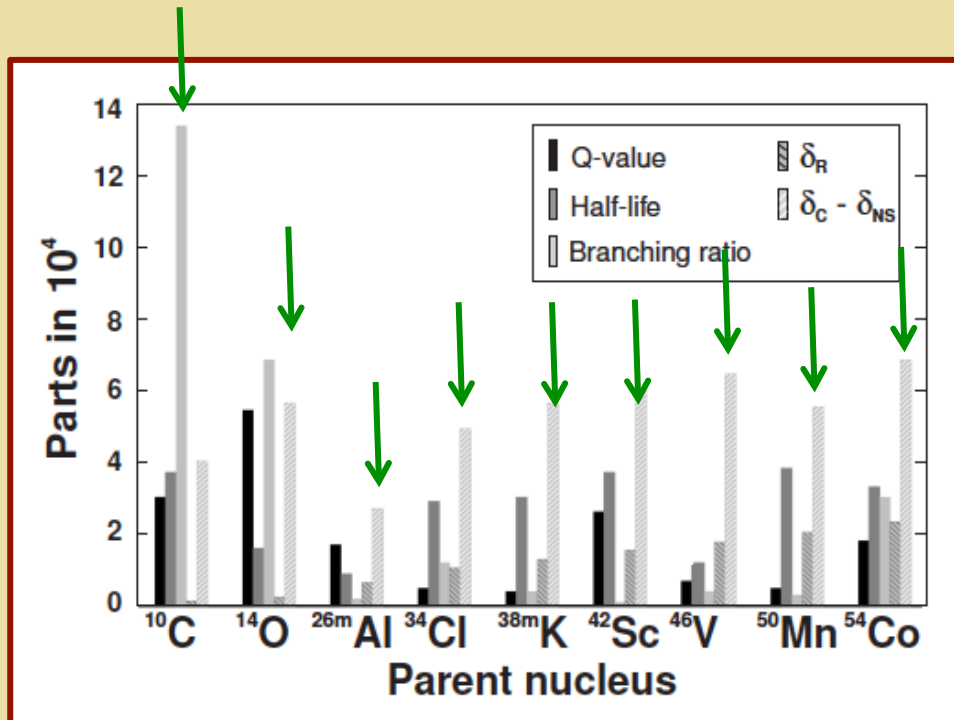
Beam normal asymmetry

Important for $O(0.1\%)$ probes of PV $^{12}\text{C}(e,e')$ & superallowed β -decay

$$V = Z^0, W$$

Nucleus-dependent QED & EW corrections

$0^+ \rightarrow 0^+$ Dispersion Corrections: δ_{NS}



b_F : scalar currents

Input for V_{ud} & CKM unitarity test

Towner & Hardy, PRC 91 (2015) 2, 025501

- Re-compute with state-of-the-art many-body methods
- Test w/ A_n predictions & expt for ^{10}B , ^{14}N , ^{26}Mg , ^{34}S , ^{38}Ar , ^{42}Ca , ^{46}Ti , ^{50}Cr , ^{54}Fe
- Investigate strategy for obtaining reduced error bars

IV. Outlook

- *Studies of neutron and nuclear β -decay are heading to a new era of precision, with a goal $\delta\Delta_{CKM} \sim O(10^{-4})$*
- *Hadronic and nuclear uncertainties in computing the $W\gamma$ box radiative correction remain one of the key challenges to reaching this goal*
- *Recent developments using dispersion relations open a new path toward reducing this uncertainty with an opportunity for new experimental tests using leptonproduction & theoretical tests with lattice QCD*
- *There exists an exciting opportunity to implement a unified, comprehensive program EW box computations (β -decay, PV electron scattering) and experimental tests with polarized electron-nucleus scattering (A_n).*