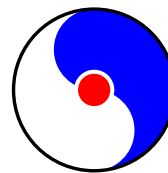
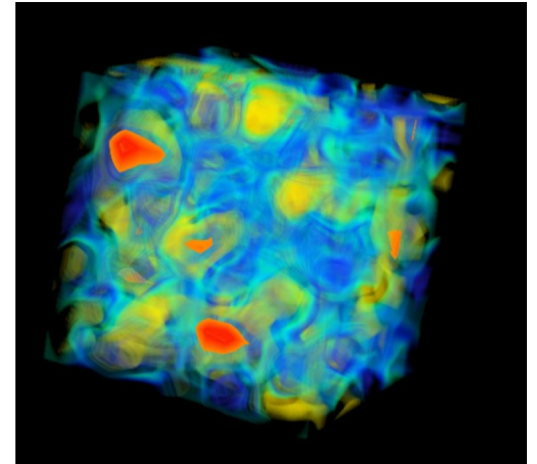


Introduction to Lattice QCD

Taku Izubuchi



RIKEN BNL
Research Center

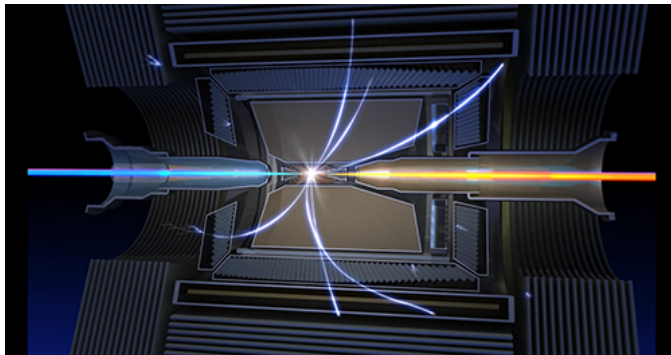
August 1st, 2019 Belle-II school, BNL

Text Book

- Gattringer, Lang, “Quantum Chromodynamics on the Lattice”
- DeGrand, DeTar, “Lattice Methods for Quantum Chromodynamics”
- Montvay, Munster, “Quantum Fields on Lattice”
- Creutz, “Quarks, Gluons and Lattices”

Particle & Nuclear Physics

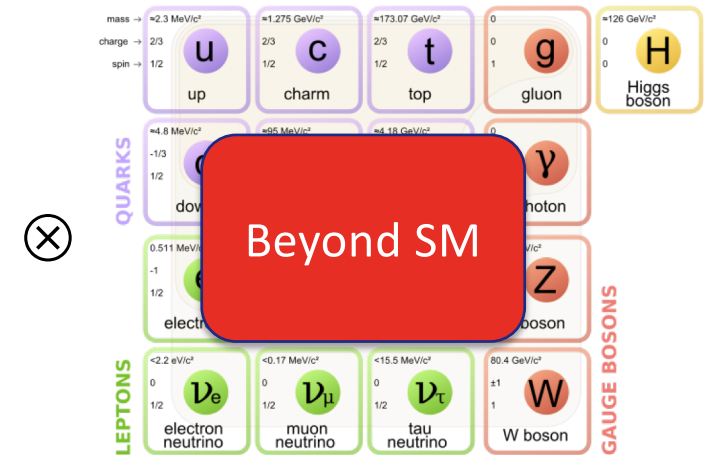
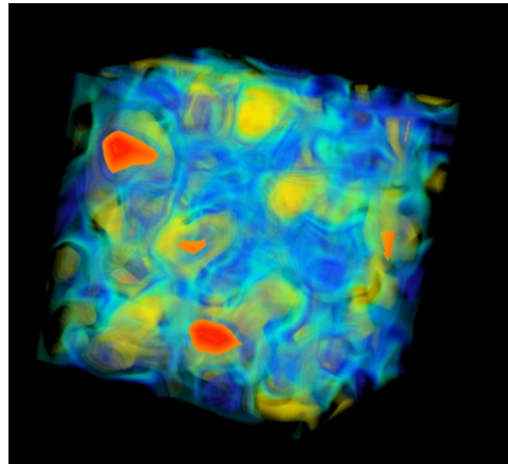
- [experiment]
theory]



[KEK.JP]

[QCD, Strong Interaction]

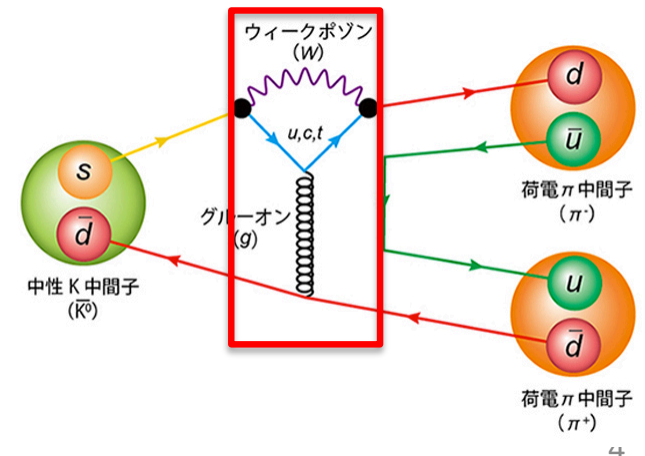
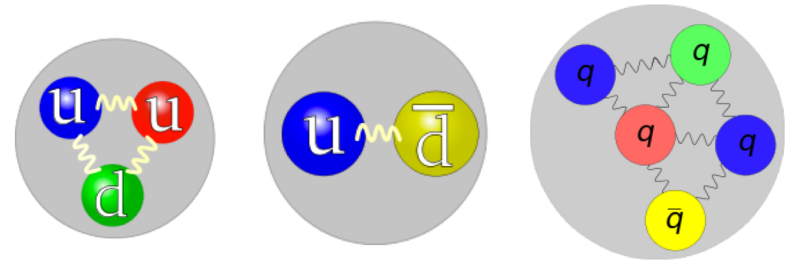
[Standard Model



- Goal : understand and check fundamental law of sub-atomic physics, Standard Model, through many interplay between experiments and theory
- To interpret experimental results, one often needs to solve how sub-atomic particles, quarks and gluons interacts each other.

Why study QCD ?

- Strong dynamics
- Make composite "stuff" :
Baryons, mesons, exotics
[spectroscopy , energy]
- To interpret experimental results
[hadronic matrix elements]
- Precision with reliable error



Quantum Chromo Dynamics (QCD)

- **Gluon** : 100 times stronger coupling than electromagnetism, strong coupling constant α_s
- **Quark** : quark mass m_f
- **confinement** → only color neutral, White, can exist at low energy and low temperature
- Make composite objects and phenomena, large hierarchy of scale
- Inter nuclear force is also from QCD → how nuclei exist

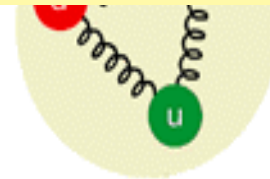
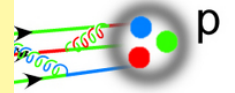
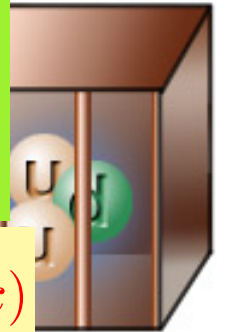
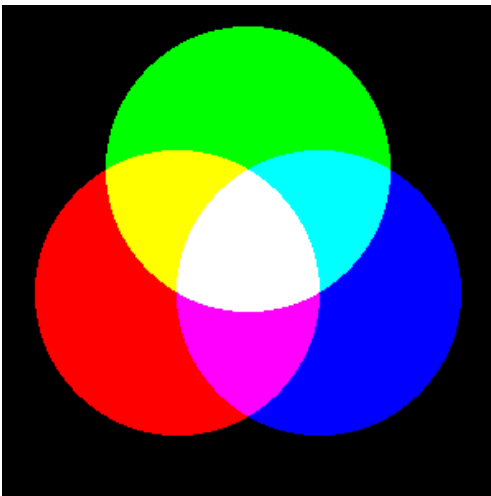
(so far) there is no further sub-structure nor more fundamental law known

If we solve QCD (in principle) we know (almost) **everything very precisely**

$$\mathcal{L}_{\text{QCD}} = -\frac{1}{4}F_{\mu\nu}^2(x) + \sum_f \bar{\psi}_f(i\gamma^\mu D_\mu - m_f)\psi_f(x)$$

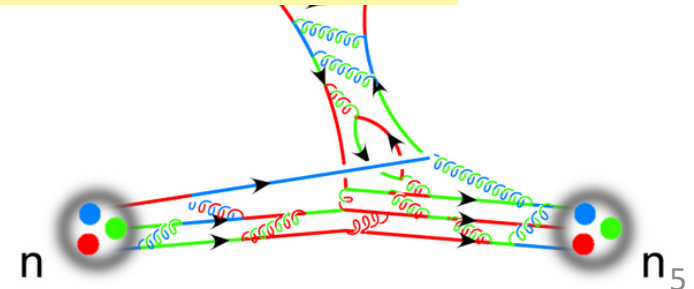
$$F_{\mu\nu}^a = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + gf^{abc}A_\mu^b A_\nu^c$$

$$D_\mu = \partial_\mu - igT^a A_\mu^a(x)$$



u-quark

Proton

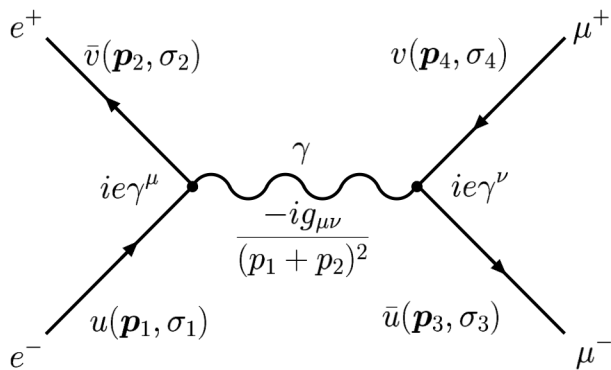


mass →	$\approx 2.3 \text{ MeV}/c^2$	$\approx 1.275 \text{ GeV}/c^2$	$\approx 173.07 \text{ GeV}/c^2$	0	$\approx 126 \text{ GeV}/c^2$
charge →	$2/3$	$2/3$	$2/3$	0	0
spin →	$1/2$	$1/2$	$1/2$	1	0
	u up	c charm	t top	g gluon	H Higgs boson
QUARKS	$\approx 4.8 \text{ MeV}/c^2$	$\approx 95 \text{ MeV}/c^2$	$\approx 4.18 \text{ GeV}/c^2$	0	
	$-1/3$	$-1/3$	$-1/3$	0	
	$1/2$	$1/2$	$1/2$	1	
	d down	s strange	b bottom	γ photon	
LEPTONS	$0.511 \text{ MeV}/c^2$	$105.7 \text{ MeV}/c^2$	$1.777 \text{ GeV}/c^2$	$91.2 \text{ GeV}/c^2$	
	-1	-1	-1	0	
	$1/2$	$1/2$	$1/2$	1	
	e electron	μ muon	τ tau	Z Z boson	
	$< 2.2 \text{ eV}/c^2$	$< 0.17 \text{ MeV}/c^2$	$< 15.5 \text{ MeV}/c^2$	$80.4 \text{ GeV}/c^2$	
	0	0	0	± 1	
	$1/2$	$1/2$	$1/2$	1	
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W W boson	
				GAUGE BOSONS	

GLUON

Non perturbative

- Perturbation theory is a powerful tool, it's origin of intuitive understandings.
- Feynman diagram is a perturbative expansion of amplitude



$$|\overline{\mathcal{M}}|^2 = \frac{e^4}{4(p_1 + p_2)^4} \text{Tr}[(\not{p}_1 + m_1)\gamma^\mu(\not{p}_2 - m_2)\gamma^\nu] \text{Tr}[(\not{p}_4 - m_4)\gamma_\mu(\not{p}_3 + m_3)\gamma_\nu]$$

- Convergence of perturbative series ? Precision of calculation ?

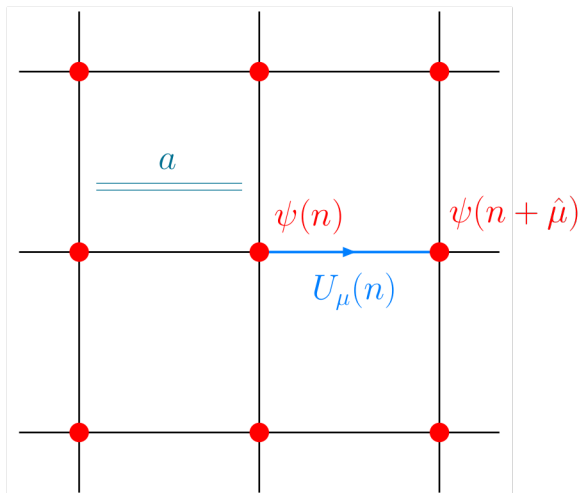
$$\mathcal{M} \propto \exp\left(-\frac{8\pi^2}{g^2}\right) = 0 + 0g^2 + 0g^4 \dots$$

- Non-perturbative effects are responsible for many important physics

Gluon on Lattice

- Lattice discretization for a non-perturbative analysis of strong dynamics, (K.G. Wilson 1974)

$\Psi(x), A_\mu(x), x \in \mathcal{R}^4$: *continuous infinity*
quantum divergences: needs regularization and renormalization



- Discretized Euclidean space-time
- lattice spacing $a \sim 0.1 \text{ fm}$
 (UV cut-off $|p| \leq \pi/a$)
- $\psi(n)$: Fermion field (Grassmann number)
- $U_\mu(n)$: Gauge field

gauge invariance $g(x) : \text{SU}(3)$

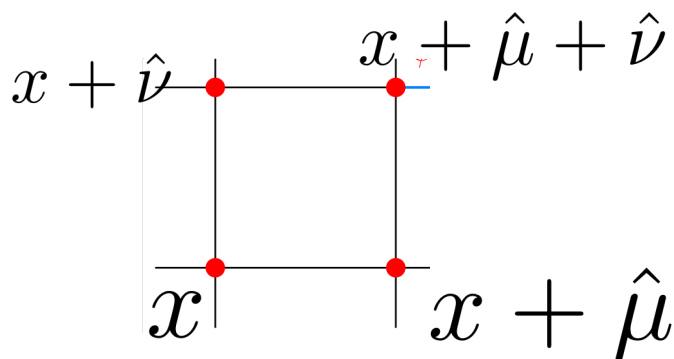
$$\psi(x) \rightarrow g(x)\psi(x), \bar{\psi}(x) \rightarrow \bar{\psi}(x)g^\dagger(x),$$

$$U(x, \mu) \rightarrow g(x)U(x, \mu)g^\dagger(x + \mu)$$

- SU(3)-valued gauge field $U(x, \mu)$ “link field”

$$U(x, \mu) = \exp(iaA_\mu(x)) = 1 + iaA_\mu(x) + \dots$$

- Field strength



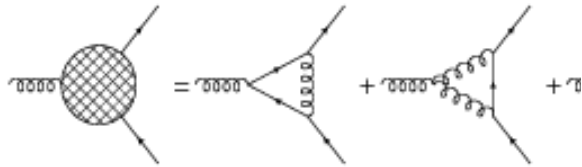
$$\begin{aligned} & \text{Tr} [U(x, \mu)U(x + \hat{\mu}, \nu)U^\dagger(x + \hat{\nu}, \mu)U^\dagger(x, \nu)] \\ &= \text{Tr} \left[e^{igaA_\mu(x)} e^{iga(A_\nu + a\partial_\mu A_\nu)} e^{-iga(A_\mu(x + \partial_\nu A_\mu))} e^{-igaA_\nu} \right] \\ &= \text{Tr} \exp \left[ig a^2 (\partial_\mu A_\nu - \partial_\nu A_\mu) - g^2 a^2 [A_\mu, A_\nu] \right] \\ &= 3 - \frac{1}{2} g^2 a^4 \text{Tr} [F_{\mu\nu}^2] \end{aligned}$$

- Wilson’s plaquette action, $\beta = 6 / g^2$

$$S_g = \sum_{\square} \beta \left\{ 1 - \frac{1}{3} \text{Re tr}(U_{\square}) \right\} \xrightarrow{a \rightarrow 0} \frac{1}{2g^2} a^4 \sum_x \text{tr}[F_{\mu\nu}(x)^2] + \mathcal{O}(a^6)$$

Strong coupling constant

- Strong coupling constant $\alpha_s = g^2 / 4 \pi$ runs due to the gluon's self coupling

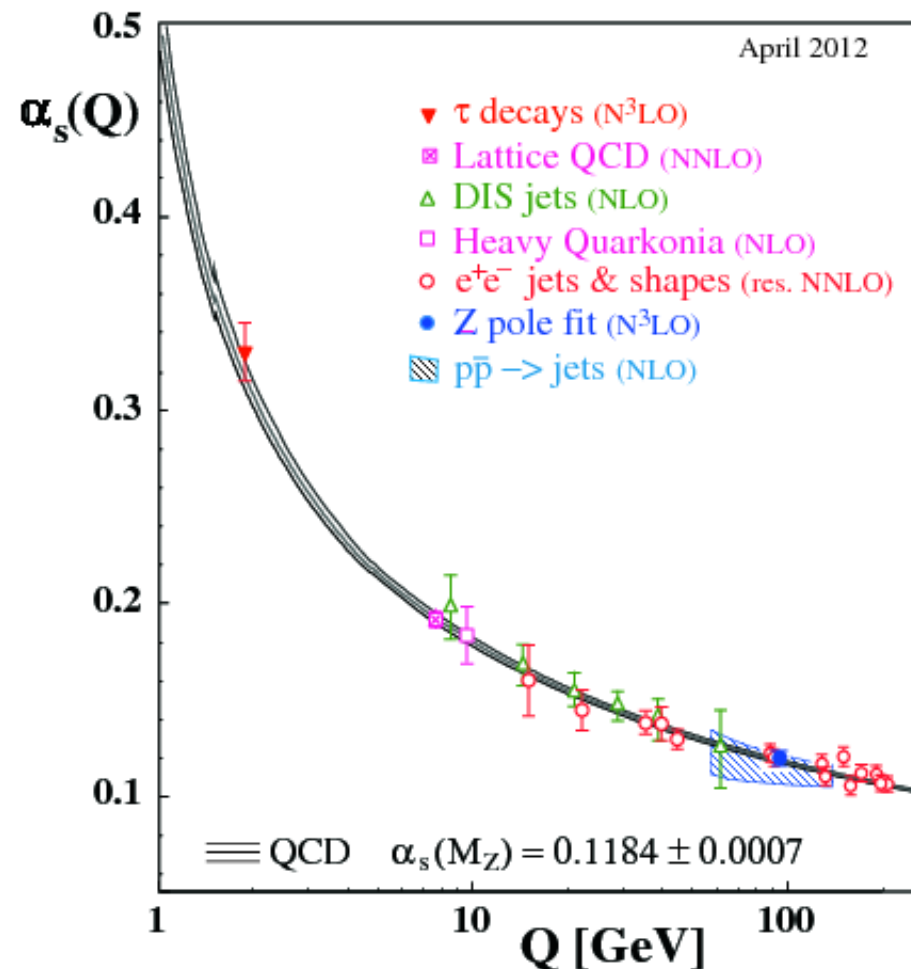


- Asymptotic freedom

- Lattice coupling

$$\beta = 6 / g^2$$

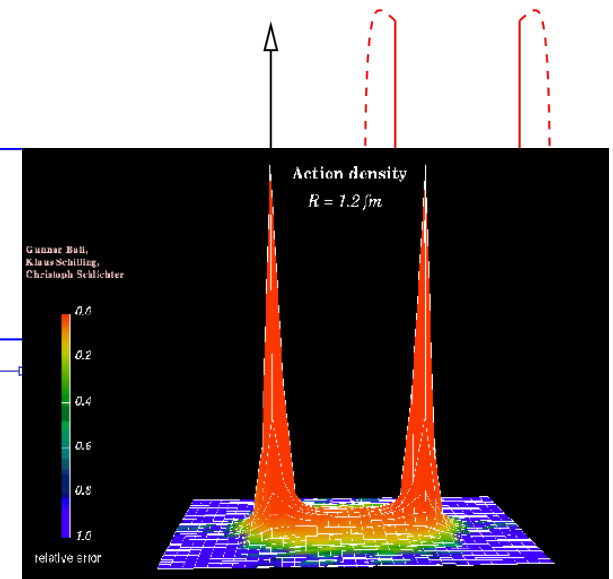
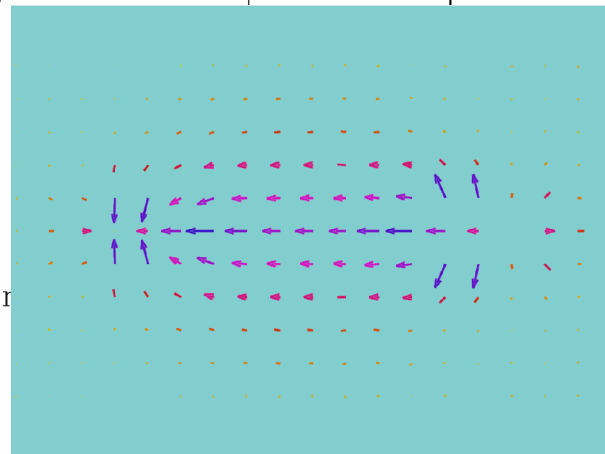
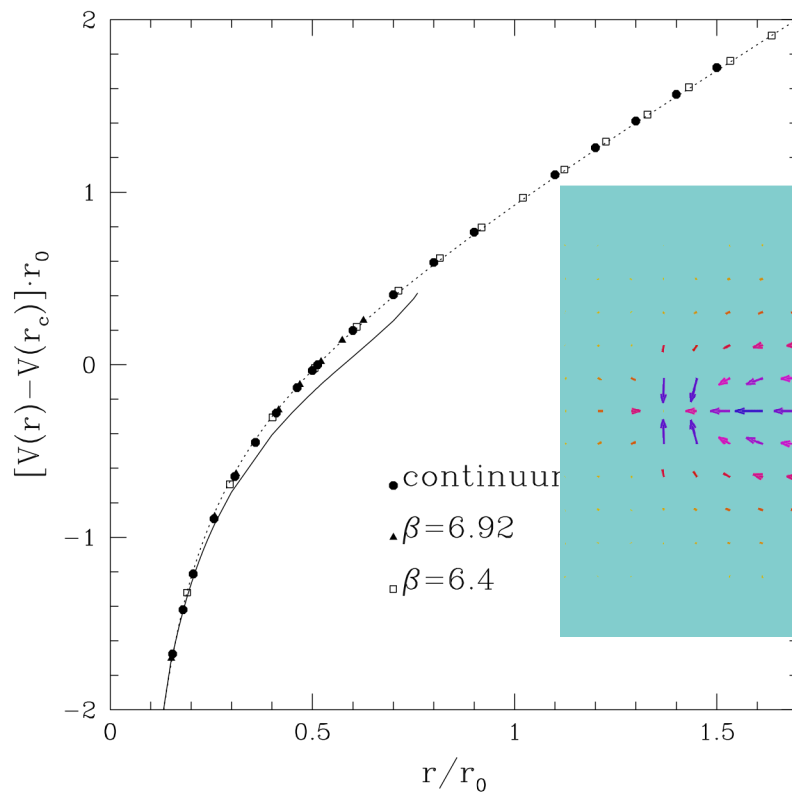
β large \leftrightarrow small a



Static quark potential

- Infinity heavy quark
- Cornell parametrization

$$V(r) = \sigma r + \frac{\gamma}{r} + \dots$$



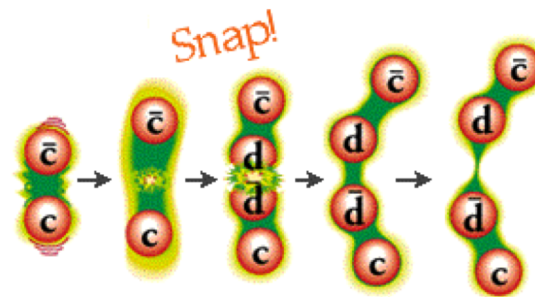
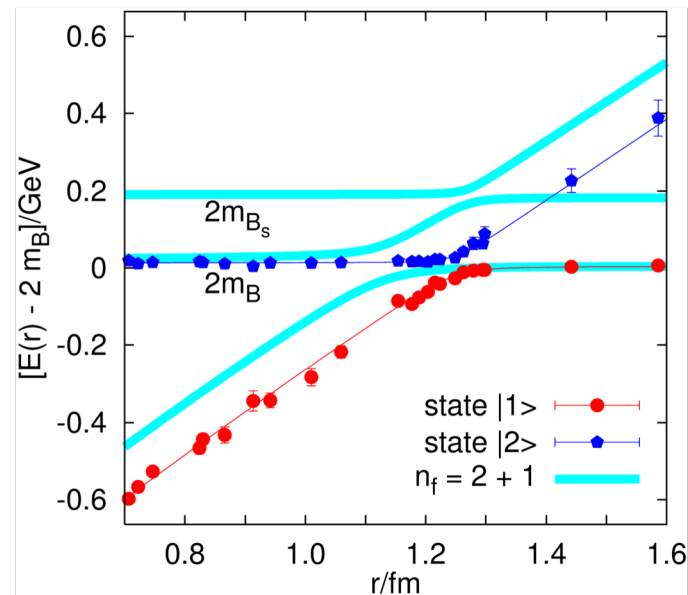
[G. Bali. Home page]

[compilation F. Knechtli, arxiv:1706.00282]

String breaking

- [G. Bali PRD71 (2005) 114513]

$$\left(\begin{array}{c} \boxed{} \\ \sqrt{N_f} \boxed{} \end{array} \right) \begin{array}{c} \sqrt{N_f} \boxed{} \\ -N_f \boxed{} + \text{wavy line} + \text{wavy line} \end{array}$$



mass →	$\approx 2.3 \text{ MeV}/c^2$	$\approx 1.275 \text{ GeV}/c^2$	$\approx 173.07 \text{ GeV}/c^2$	0	$\approx 126 \text{ GeV}/c^2$
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	u up	c charm	t top	g gluon	H Higgs boson
	d down	s strange	b bottom	γ photon	
	e electron	μ muon	τ tau	Z Z boson	
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W W boson	

QUARKS

LEPTONS

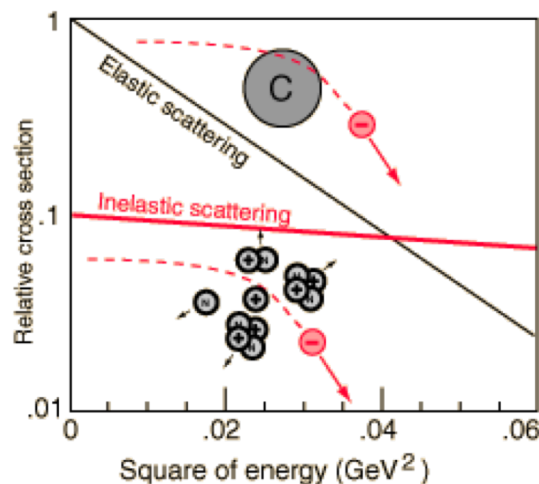
GAUGE BOSONS

QUARK

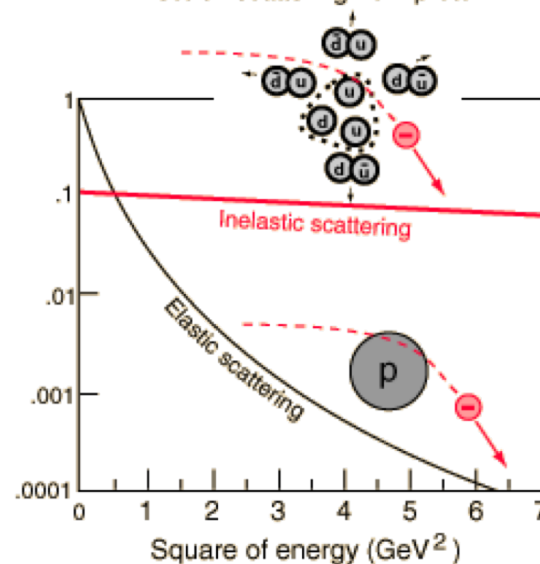
Proton/neutron are made of **Quarks**

- Electron - nucleon scattering (1968 SLAC)
- Proton has a substructure seen by energetic electrons
- “three quarks for Muster Mark”
James Joyce *Finnegan's Wake*

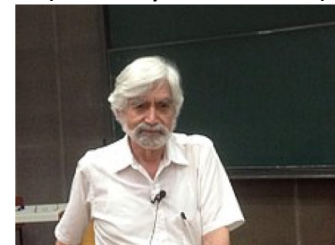
Electron scattering from carbon atom



Electron scattering from proton



(Murray Gell-Mann)



(George Zweig)

Proton
Charge:
+ 1



Neutron
Charge:
0



U mass 2.3(7)(5) MeV

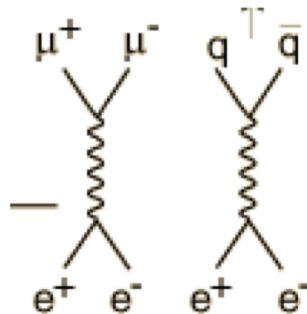
D mass 4.8(5)(3) MeV

U up quark, charge : $\frac{2}{3}$

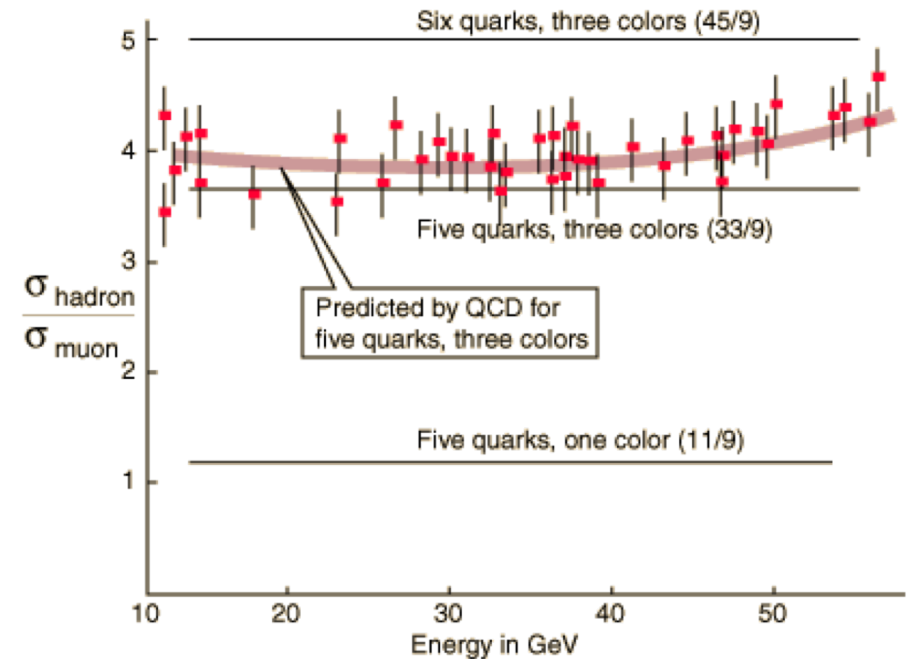
D down quark, charge: $-\frac{1}{3}$

R-ratio

$$\frac{\sigma(e^+e^- \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)}$$



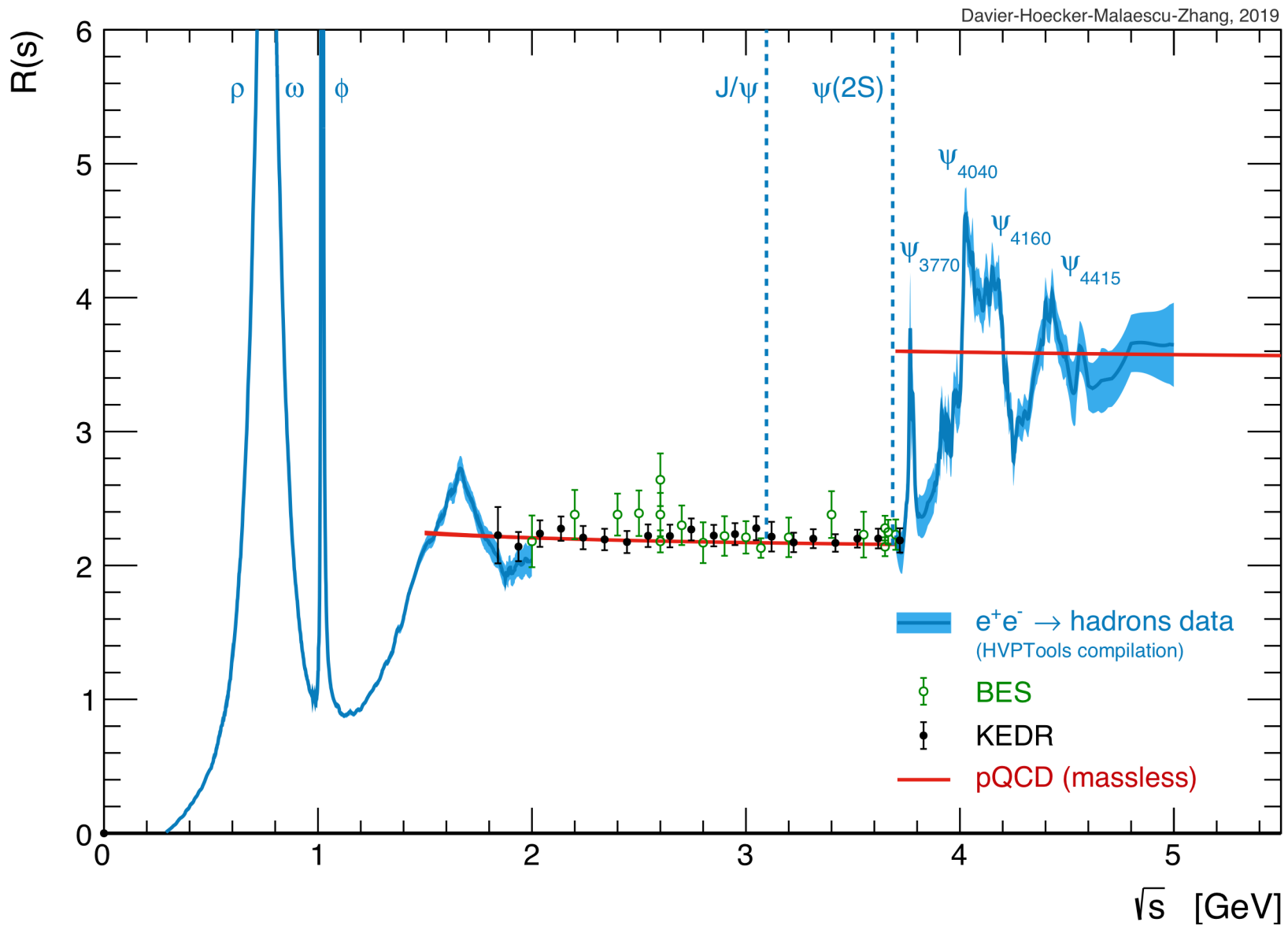
$$\left[\left(\frac{2}{3} \right)^2 + \left(-\frac{1}{3} \right)^2 + \dots \right] \times 3 = \frac{11}{3}$$



[Hyper Physics]

3 internal d.o.f : color SU(3)

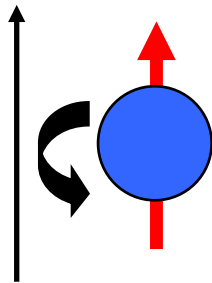
R-ratio



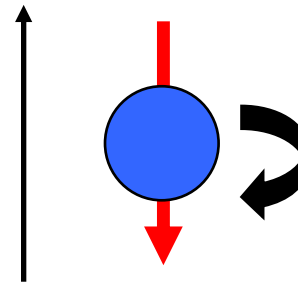
Quark spin & chiral symmetry

- Dirac equation: quark, electron have two degrees of freedom

SPIN



Right handed

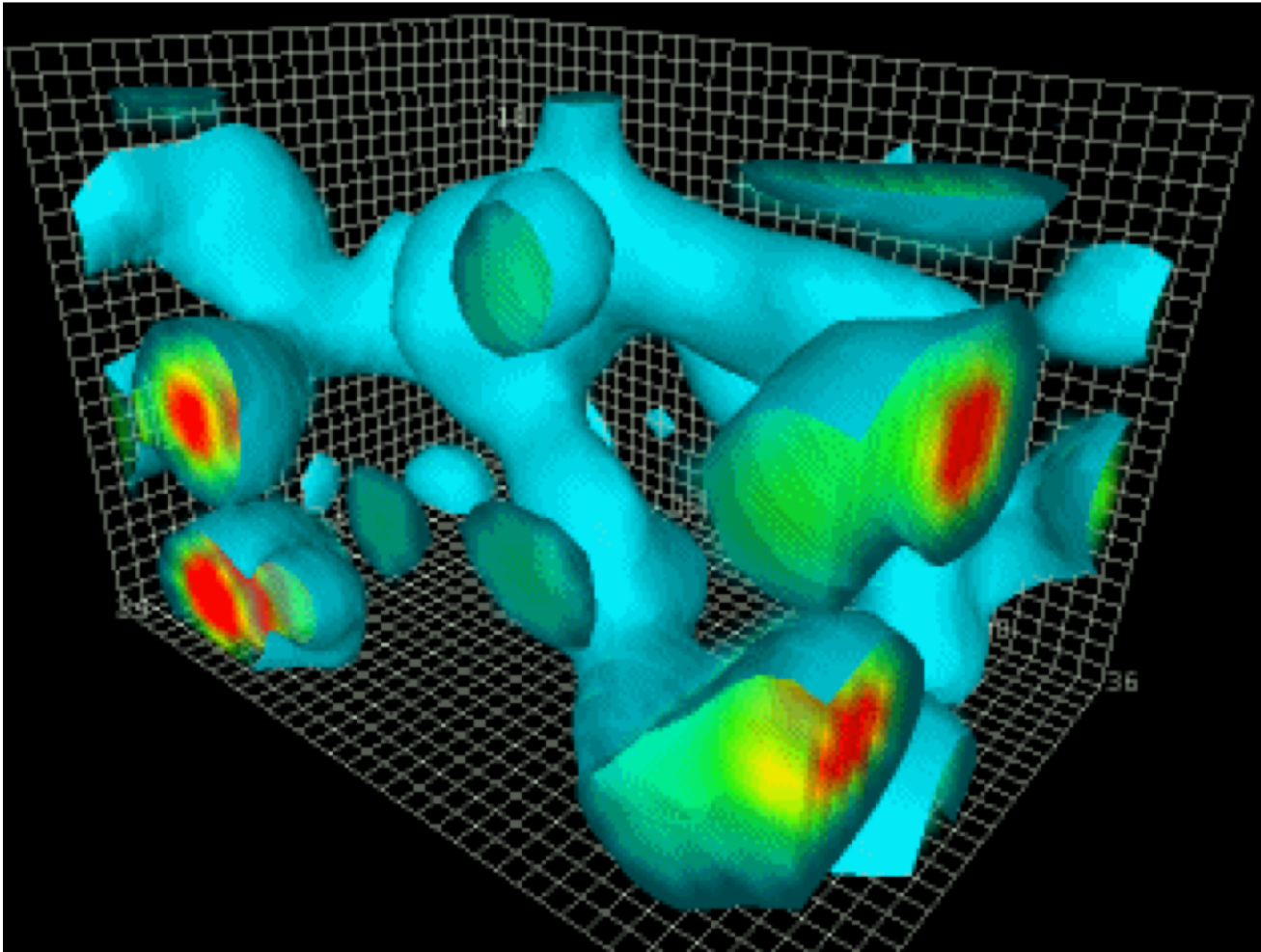


Left handed

- **Massive particles** (slower than speed of light)
⇔ mixing between left and right spins
- **massless particles** (speed of light)
⇔ NO mixing between left and right spins

Chiral symmetry

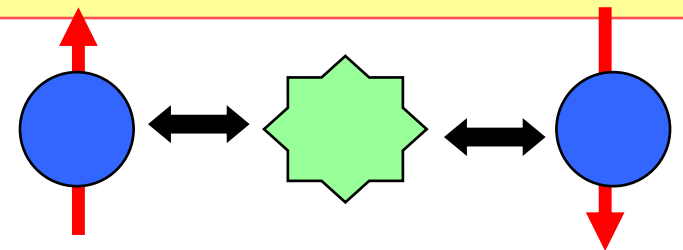
Fluctuation of QCD vacuum and spontaneous breaking of Chiral symmetry



QCD is creating / annihilating instantaneous “lump” of chiral charges

induce **effective mass** for quark

quark “picks” this chiral charge, which could be seen as **spontaneous breaking of chiral symmetry**



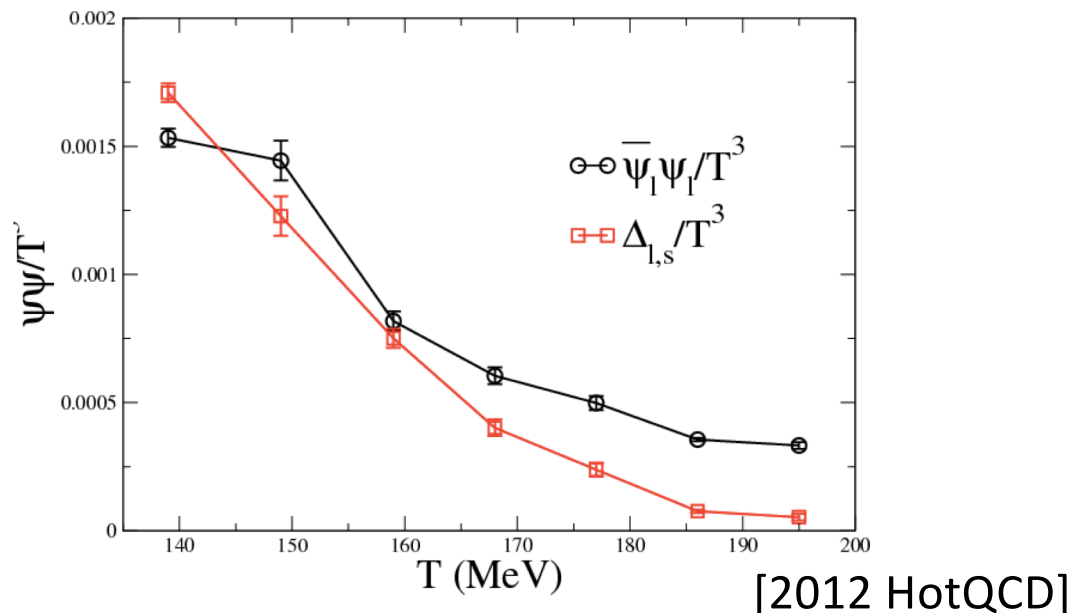
(Courtesy of Derek Leinweber, CSSM, University of Adelaide)

Origin of mass

- QCD (T=0): quark and anti-quark creates a pair (chiral condensate) , which is energetically stable than nothing !

$$\langle \bar{q}q \rangle = [242(04)(18) \text{ MeV}]^3 \quad (2010 \text{ JLQCD})$$

- Produced quark an effective mass roughly about a third of proton masses
 $3 \times 0.25 \text{ GeV} \sim 1 \text{ GeV}$
 c.f. up down quark mass ~ a few MeV
- At high temperature, chiral condensate decreases (Quark-Gluon Plasma)



Y. Nambu



J. Goldstone



[CERN] cf) Higgs mechanism, or superconductor

fermion action

Lattice fermion and **doublers**

$$S = \int d^4x \bar{\psi}(x) \not{D} \psi(x)$$

↓ on lattice

$$S = a^4 \sum_n \bar{\psi}_n \gamma_\mu \frac{1}{2a} (\psi_{n+\mu} - \psi_{n-\mu})$$

$$S_F = a^4 \sum_x \bar{\psi}(x) (\not{D} + m_0) \psi(x)$$

$$\not{D} = \frac{1}{2} \{ \gamma_\mu (\nabla_\mu^* + \nabla_\mu) \}$$

$$\nabla_\mu \psi(x) = \frac{1}{a} [U_\mu(x) \psi(x + a\hat{\mu}) - \psi(x)]$$

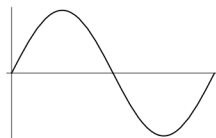
Wilson fermion (breaks chiral sym)

$$S = \sum_n \bar{\psi}_n \left[\gamma_\mu D_\mu - \frac{a}{2} D^2 \right] \psi_n,$$

$$S = \int d^4p \bar{\psi} \left[i\gamma_\mu \sin p_\mu + \frac{1}{a} \sum_\mu (1 - \cos p_\mu) \right] \psi.$$

- **Doubler problem**
(Nielsen-Ninomiya no-go theorem)

$$S = \int_{-\pi}^{\pi} d^4p \bar{\psi}(-p) i\gamma_\mu \sin p_\mu \psi(p)$$



$$\not{D}(p) = i\gamma_\mu \sin p_\mu \rightarrow i\gamma_\mu p_\mu$$

: Particle

$$\not{D}(\pi - p) = i\gamma_\mu \sin(\pi - p_\mu) \rightarrow -i\gamma_\mu p_\mu$$

: Doublers

↓

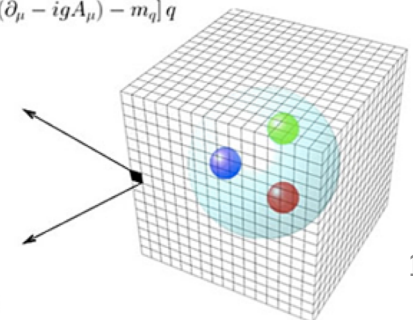
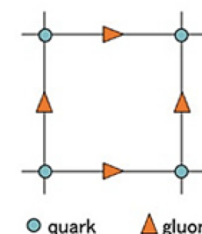
16 fermions. (4 dim.)

$$\not{D}(p) \rightarrow i\gamma_\mu p_\mu + \mathcal{O}(a),$$

$$\not{D}(\pi - p) \rightarrow -i\gamma_\mu p_\mu + \frac{2}{a} + \mathcal{O}(a).$$

QCD Lagrangian

$$\mathcal{L} = -\frac{1}{4} F^{\mu\nu} F_{\mu\nu} + \sum_{q=u,d,s,c,b,t} \bar{q} [i\gamma^\mu (\partial_\mu - igA_\mu) - m_q] q$$



Lattice quark types

- Staggered-types [MILC, HPQCD, BMWc,]

Fast, high statistics, Highly Improved Staggered action (HISQ)

4 “taste”, taking $[\det D]^{1/4}$, Non-local action, Unitarity ?

- 4D Wilson-types

[Alpha, PACS, BMWc, NPLQCD, PDNME,]

No taste problem, exact flavor symmetry, chiral symmetry broken, need to tune to remove $O(a)$ effects (clover term)

Twisted Wilson [ETMC]

Fast, automatic $O(a)$, breaks isospin symmetry

- Domain-Wall Fermion (DWF) / Overlap

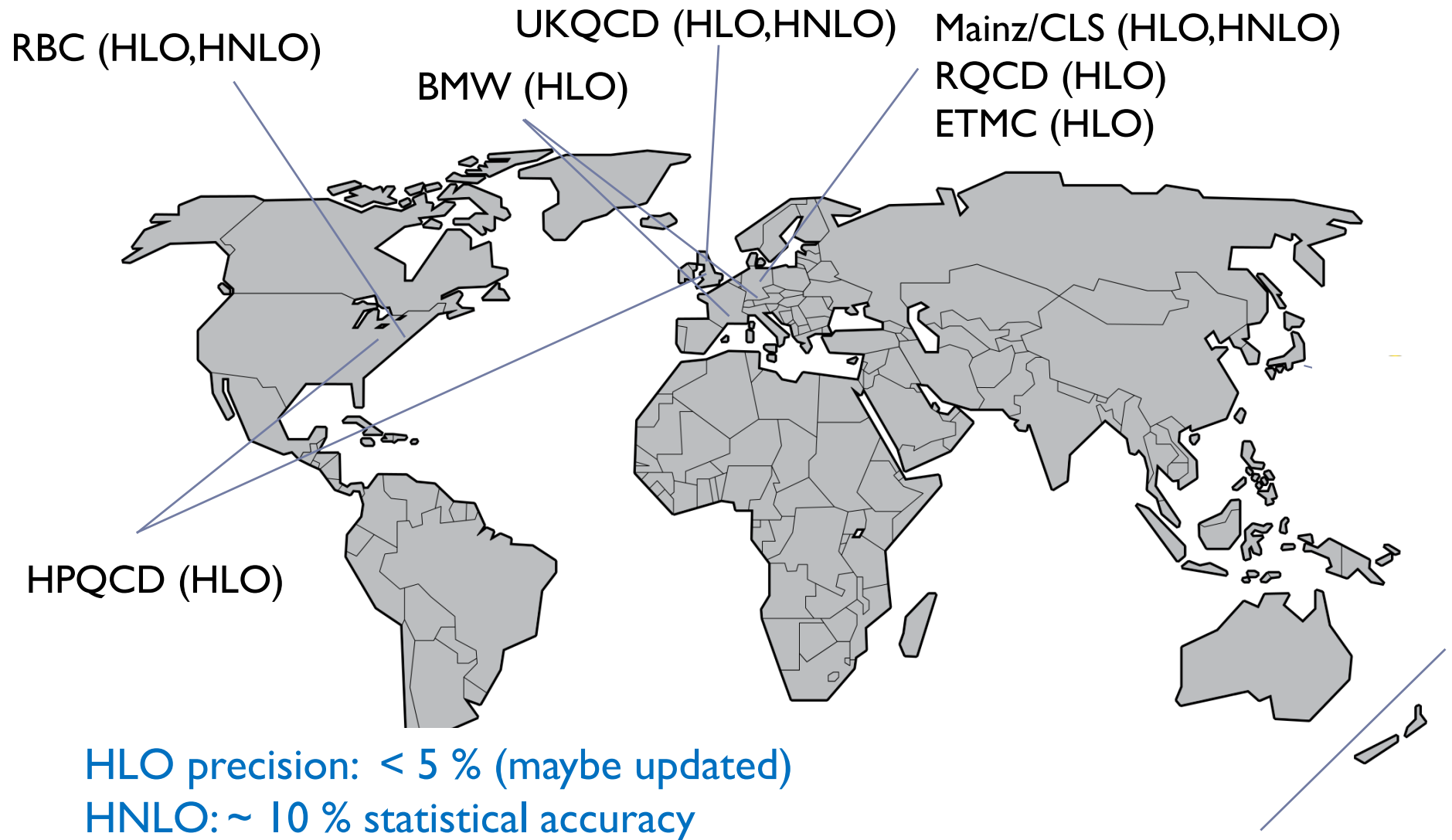
[RBC/UKQCD, JLQCD, xQCD....]

(almost) exact chiral symmetry, automatic off-shell $O(a)$ improved, expensive

1. Introduction

Lattice works

[Slide from Eigo Shintani]



Domain Wall Quarks (for up, down, and strange)

[Kaplan, Shamir, Blum & Soni]

- 4D lattice quark utilizing an "extra dimension", L_s . (expensive)

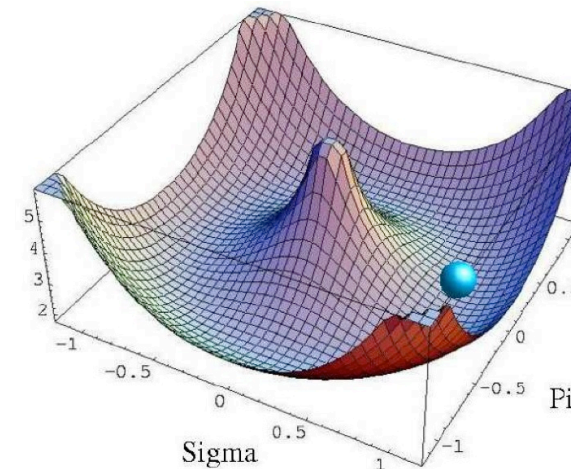
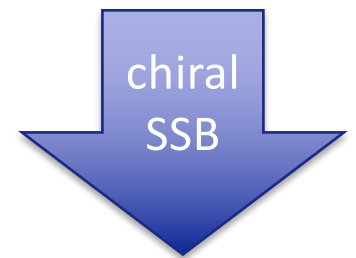
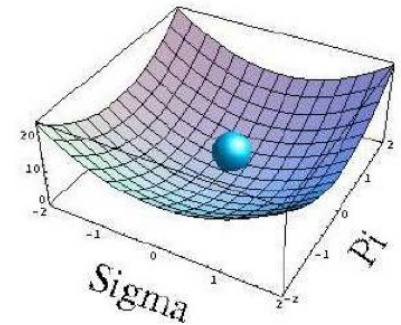
- Almost perfect chiral symmetry

Small unphysical mixing for the Weak Matrix Elements

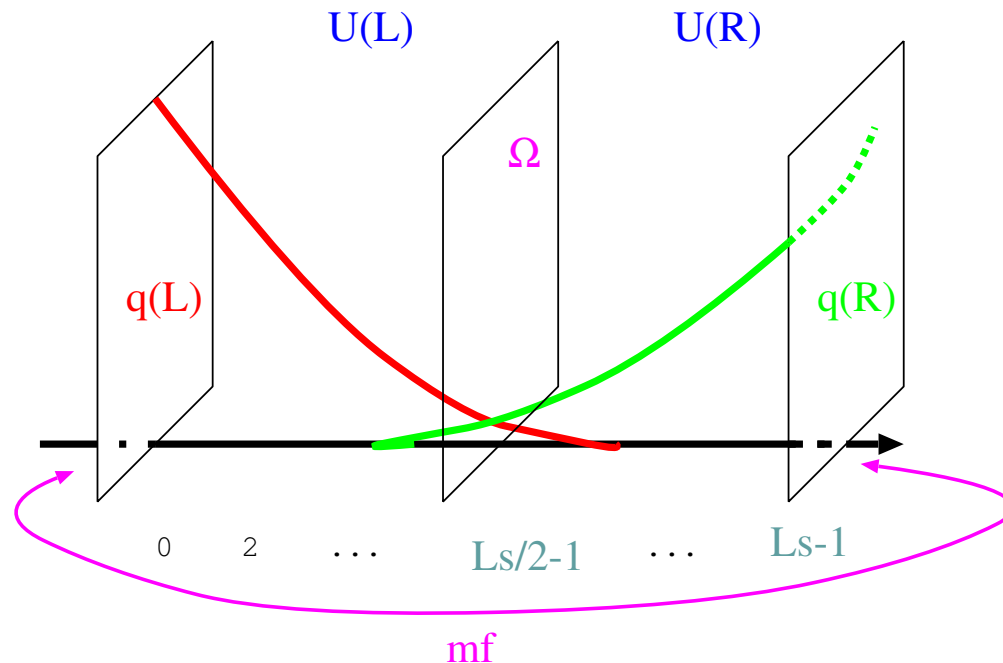
Error from discretization is small $\mathcal{O}(a^2 \Lambda_{QCD}^2) \sim \text{a few } \%$.

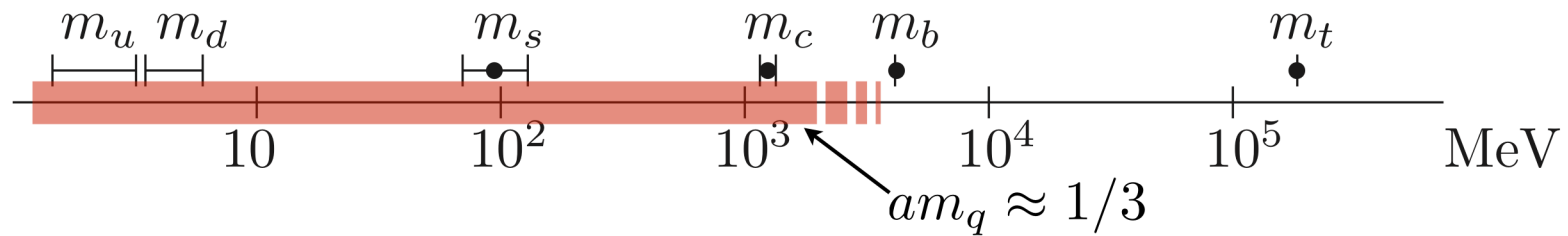
Chiral extrapolation is simpler, continuum like.

- Unitary theory (at long distance).

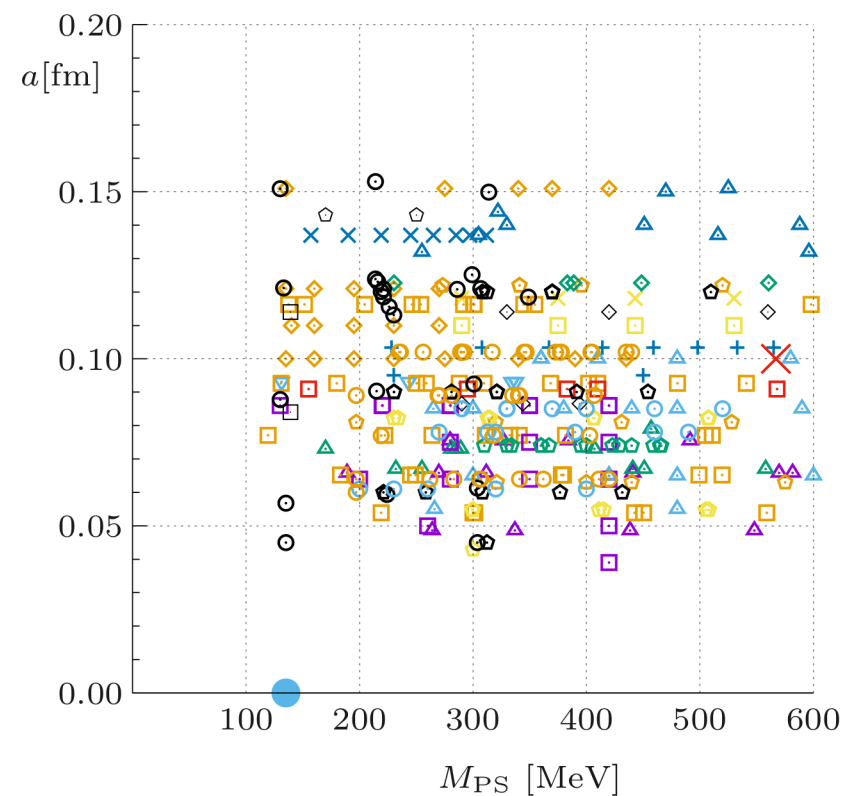


Origin of Proton mass
Emergence of NG boson



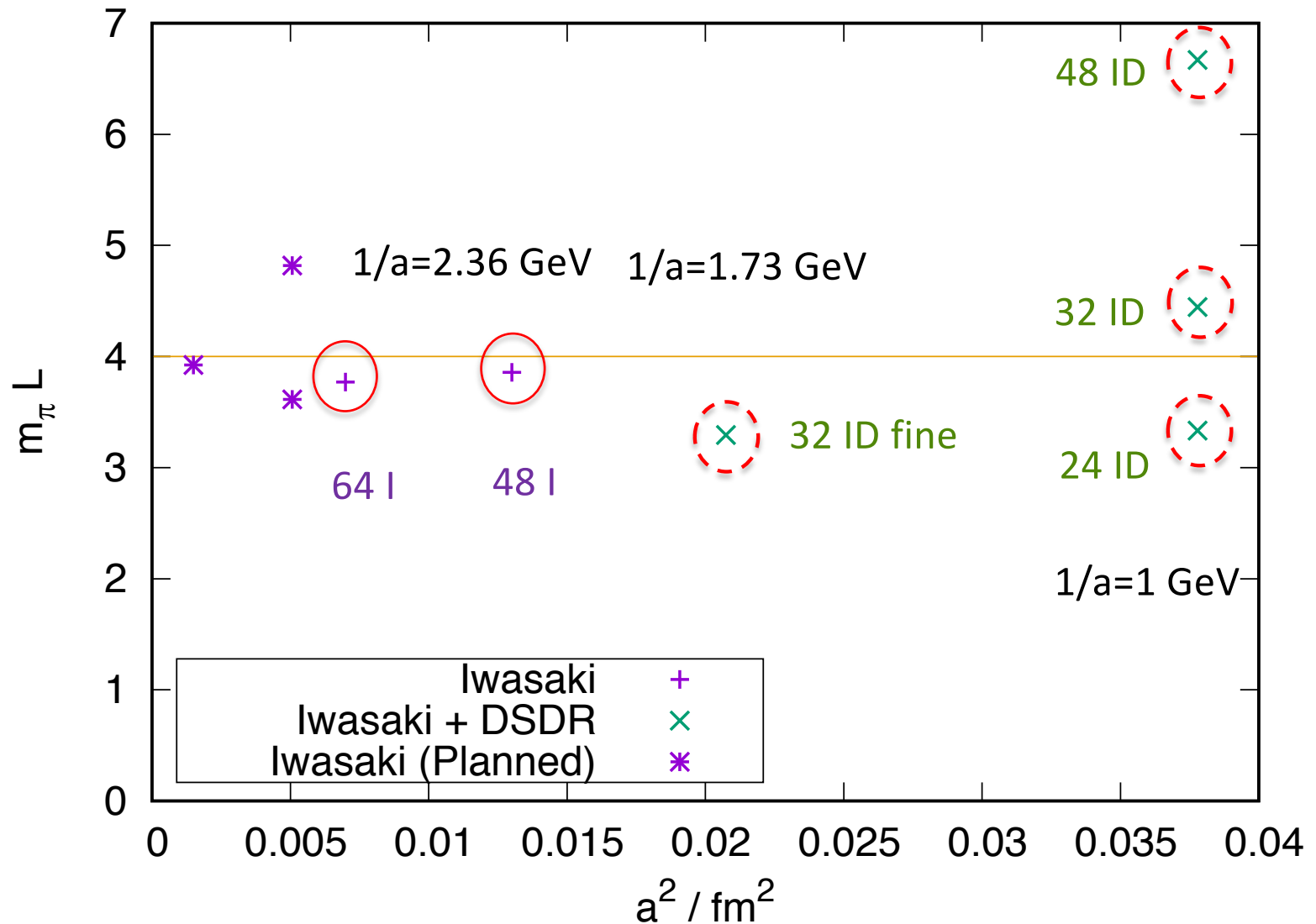


CLS	$N_f = 2$	\blacktriangle
ETMC	$N_f = 2$	\triangle
(clover) ETMC	$N_f = 2$	∇
QCDSF	$N_f = 2$	\blacktriangle
BGR	$N_f = 2$	\triangle
JLQCD	$N_f = 2$	\times
(plaq) TWQCD	$N_f = 2$	$+$
(Iwa) TWQCD	$N_f = 2$	\times
(HEX) BMW	$N_f = 2 + 1$	\square
(stout) BMW	$N_f = 2 + 1$	\diamond
(stout-stag) BMW	$N_f = 2 + 1$	\diamond
CLS	$N_f = 2 + 1$	\square
HSC	$N_f = 2 + 1$	\diamond
PACS-CS	$N_f = 2 + 1$	\square
QCDSF	$N_f = 2 + 1$	\diamond
JLQCD	$N_f = 2 + 1$	\square
(Möbius) JLQCD	$N_f = 2 + 1$	\diamond
RBC-UKQCD	$N_f = 2 + 1$	\diamond
(DSDR) RBC-UKQCD	$N_f = 2 + 1$	\diamond
(Möbius) RBC-UKQCD	$N_f = 2 + 1$	\square
MILC	$N_f = 2 + 1$	\diamond
MILC	$N_f = 2 + 1 + 1$	\circ
ETMC	$N_f = 2 + 1 + 1$	\circ
BMW	$N_f = 1 + 1 + 1 + 1$	\circ
LQCD/CP-PACS (2001)	$N_f = 2$	\times
M_π (experiment)		\bullet



[Herdoíza summer 2015+partial updates]

Nf=2+1 DWF QCD ensemble at physical quark mass



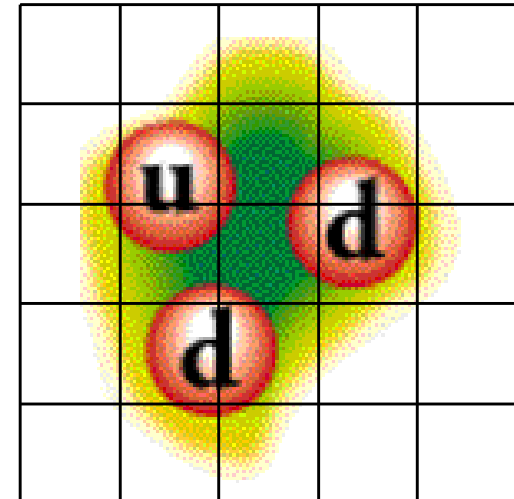
Lattice QCD calculation

- Solve QCD numerically, and enable direct comparison to experimental results
- Perform Feynman's Path Integral to compute quantum expectation value of an observable \mathcal{O}

$$\langle \mathcal{O} \rangle = \frac{1}{Z} \int \mathcal{D}U_\mu(x) \mathcal{D}\psi(x) \mathcal{D}\bar{\psi}(x) e^{-S_G + \bar{\psi}(\not{D} + m)\psi} \mathcal{O}[\psi, \bar{\psi}, U_\mu]$$

- To fit protons
space-time : $L \sim 5 \text{ fm}$,
- To avoid large systematic error
Lattice spacing : $a \sim 0.1 \text{ fm}$
→ Integration over $\sim (50)^4 \times 4 \times 8 \sim 100 \text{ Million variables}$
- To compute proton mass, typically needs, at least
 - $\sim 100\text{-}1,000 \text{ Exa floating point operations}$
 - $\sim 10^{20}\text{-}10^{21} \text{ floating operations}$
 - $\sim 100 \text{ days on one rack of Blue Gene / Q}$
 - $\sim 100\text{-}1000 \text{ years on one desktop (if fits)}$

[© AICS, Y. Kuramashi]

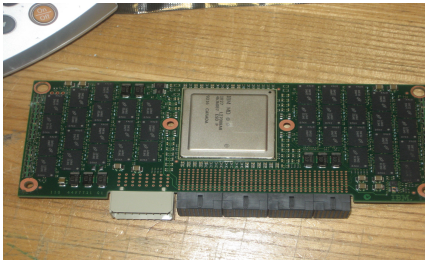


- QCD and SM is the first principle (known to date) : **stable problem**
- Need **precise numerical results** to confront big experimental results

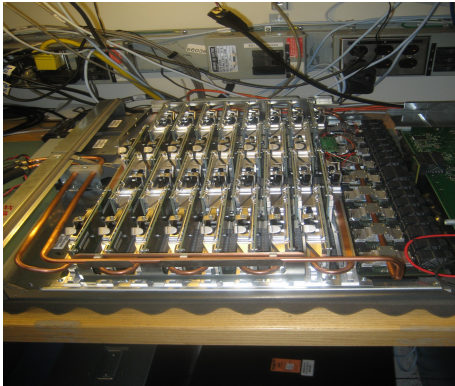
Lattice QCD super computers

- ~ massively parallel computer (1K-10K CPUs)
- ~1 - 10 PFLOPS (1-10 million times faster than PC)
- needs very fast interconnections

QCDCQ at BNL (pre-commercial Blue Gene / Q) 3.5 racks 0.7 Pflops total, 2012-

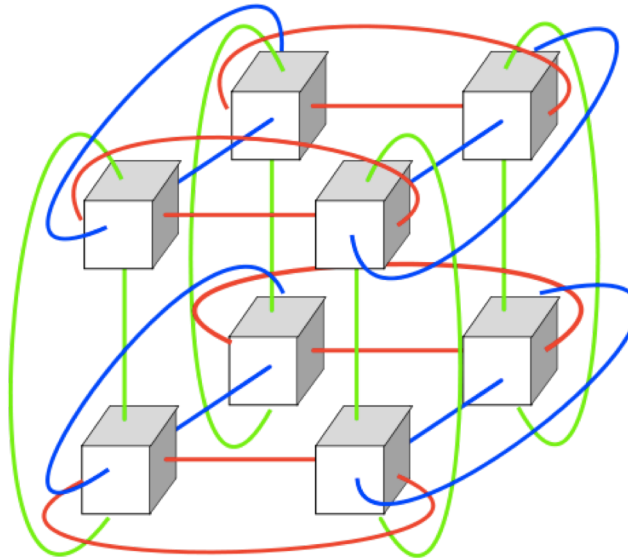


Compute card (16 cores + 16GBmemory)



Node Card (32 compute cards, 5D optical link)

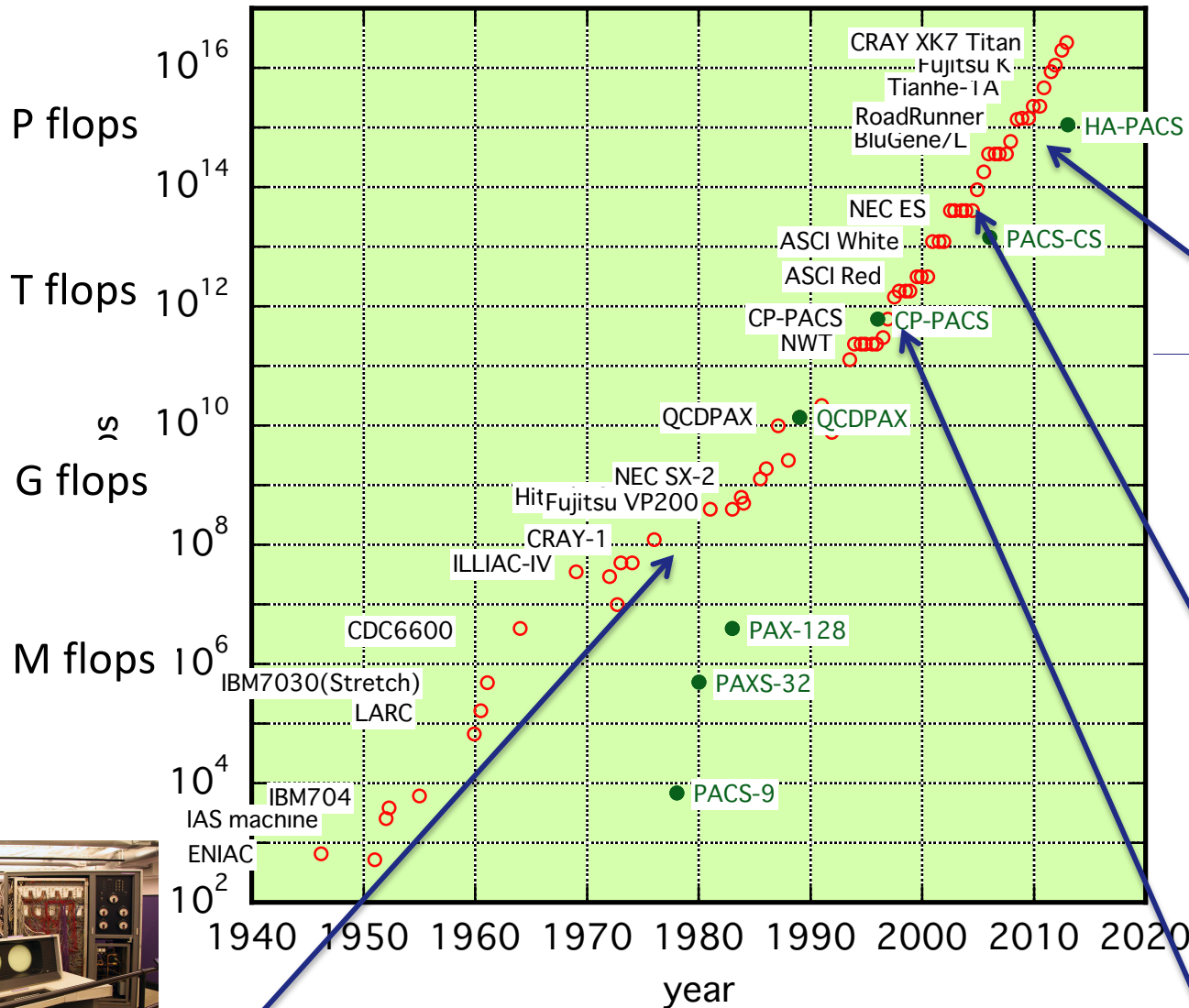
BNL, Columbia, RIKEN, U of Edinburgh
+ IBM Watson Lab.



Rack
(32 Node cards + I/O drawers)

pictures : BNL machines

computer speed



QCDCQ 0.7 P flops 2012



QCDOC 20 T flops 2005



12 racks (1024 DB per rack) ~ 600GFlops @ BNL

QCDSP 0.6T flops 1998



CDC 6600/7600 40 M flops 1970s

[A. Ukawa]

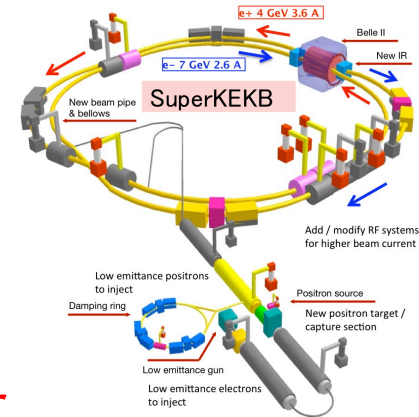
Lattice Gauge Theory Receipt

- QCD vacuum ensemble generation \sim Accelerator

- choice of gauge / sea quark actions

- Algorithms / Machines capability machines

- for each parameters $(a^{-1}, V, m^{\text{sea}})$



- Physical observable measurements \sim Detector

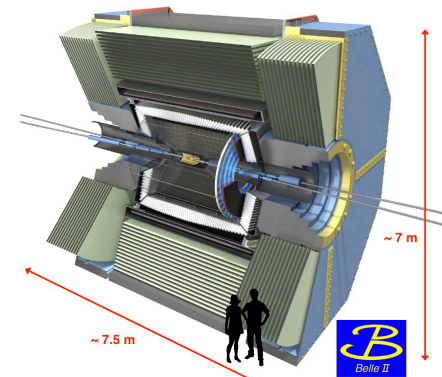
- valence quark propagators (low eigenvectors), m_f

- Hadron n -point green's functions, matrix elements

$$\langle \mathcal{O} \rangle = \int \mathcal{D}U_\mu \mathcal{D}\bar{q}_i \mathcal{D}q_i \mathcal{O} e^{-S_{\text{LGT}}} / \langle 1 \rangle$$

- Renormalize and Chiral/Volume/continuum extrapolations

- Algorithms / Machines capacity machines



\Rightarrow The final answers

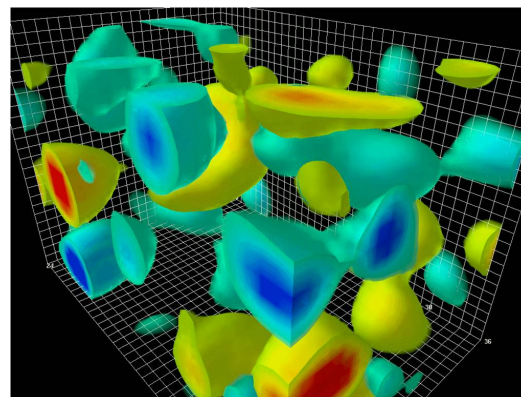
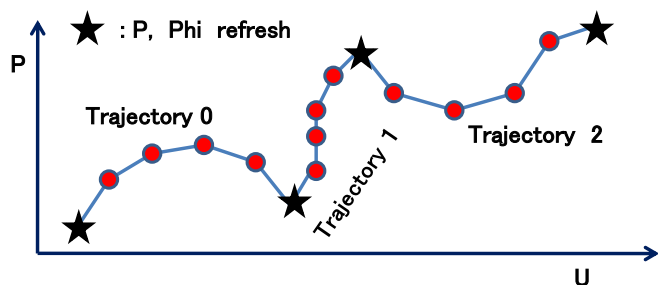
LGT ensemble generation

(\leftrightarrow Ising : spin configuration generation)

- Monte Carlo to Sample Important configurations of QCD action $e^{-S_{\text{QCD}}}$
- Accumulate samples of **QCD vacuum**, typically $\mathcal{O}(100) \sim \mathcal{O}(1,000)$ files of gauge configuration $U_\mu(n)$ on disk (1 \sim 10 GB/conf).
- By solving a classical QCD, with an occasional stochastic "hit": **exactly** $\propto e^{-S_{\text{QCD}}}$
- Must generate sequentially $\{U_\mu^{(0)} \rightarrow U_\mu^{(1)} \rightarrow \dots\}$, which **needs capable machines**.

$$\langle \mathcal{O} \rangle = \frac{1}{Z} \int \mathcal{D}U_\mu(x) \prod_{f=1}^{N_f} \det(D + m_f) e^{-S_G} \mathcal{O}$$

$$\text{Prob}(U_\mu) \propto \det D_{u,d,s}[U] e^{-S_g}$$



[D. Leinweber]

- RHMC for odd flavor [Clark Kennedy]
- Solve **short** (**long**) modes **more** (**less**) frequently [Hasenbush's trick]

LGT observable measurement

(<->detector)

- Measurements **physical observables** on the vacuum ensemble.

$$\langle \mathcal{O} \rangle = \int \frac{\mathcal{D}U_\mu \text{ Prob}[U_\mu]}{\mathcal{Z}} \times \mathcal{O}[U_\mu]$$

- Could do Analysis on many configurations **independently** (trivial parallel jobs) → could also use PC Clusters
- We made hadron **operator** (EW operators) from quark, and let the quark propagates on each of the generated QCD configuration (by solving the Dirac Eq)
- Obtain **hadron mass** or **QCD matrix elements** of operators

Euclidean 2pt Green function

O: hadron operator of target quantum #

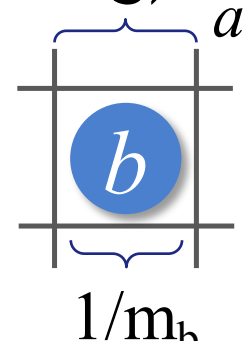
$$C_{XY}(t; \vec{p}) = \sum_{\vec{x}} \langle O_X(t, \vec{x}) O_Y^\dagger(0) \rangle e^{-i\vec{p} \cdot \vec{x}}.$$

$$C_{XY}(t) \xrightarrow{t \rightarrow \infty} \frac{1}{2E_0(\vec{p})} \langle 0 | O_X(0) | H_0(\vec{p}) \rangle \langle H_0(\vec{p}) | O_Y^\dagger(0) | 0 \rangle e^{-E_0(\vec{p})t}$$

$$C_{XY}(t; \vec{p}) = \sum_{i=0}^{\infty} \frac{1}{2E_i(\vec{p})} \langle 0 | O_X(0) | H_i(\vec{p}) \rangle \langle H_i(\vec{p}) | O_Y^\dagger(0) | 0 \rangle e^{-E_i(\vec{p})t}$$

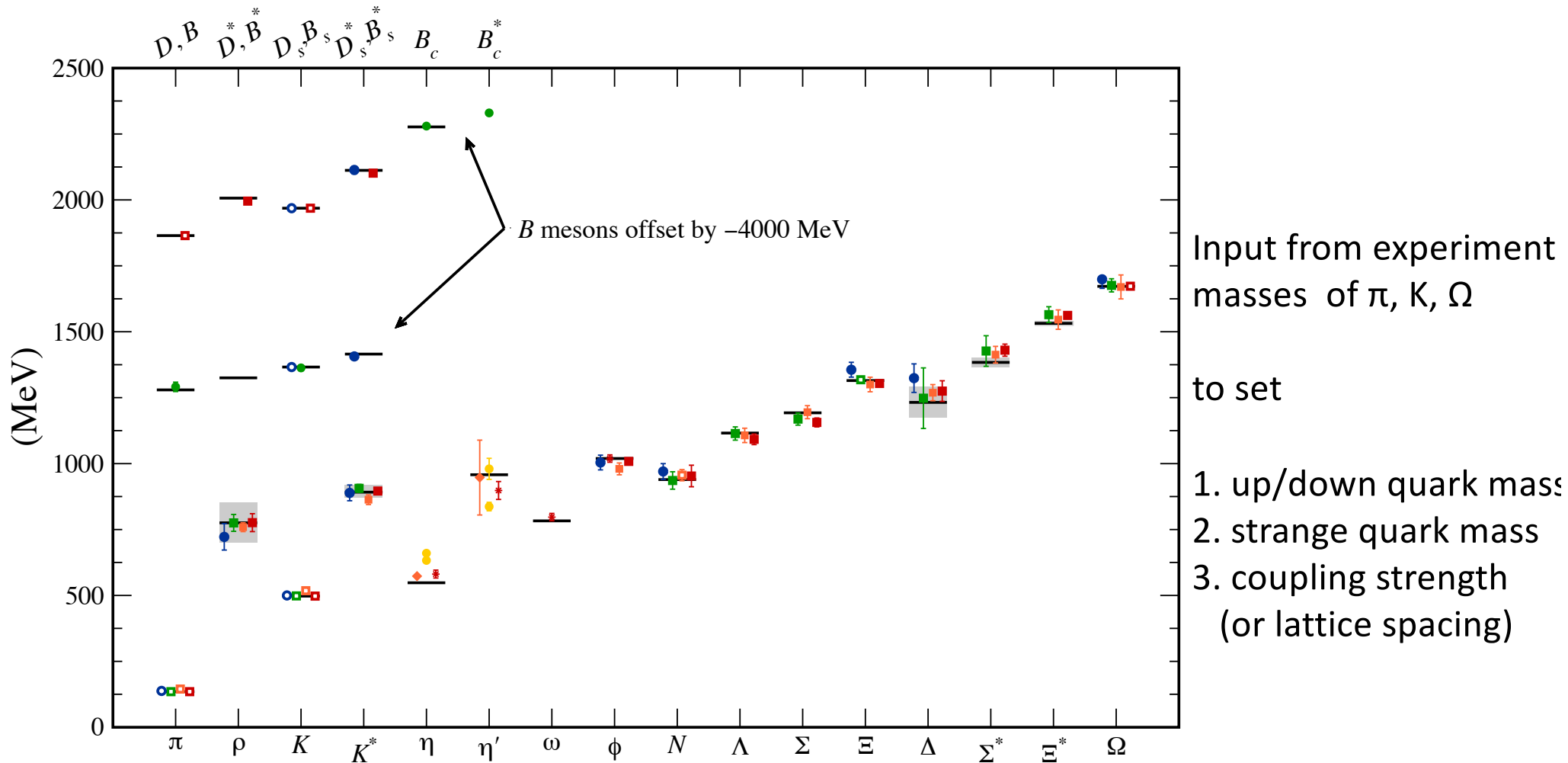
Input parameters

- Same as the continuum theory (coupling, quark mass)
- coupling (lattice spacing a): for given β , measure one hadron mass (e.g. Omega baryon) ($a m_\Omega$) = $a \times 1.672.45 \text{ GeV}$
- quark mass m_f : measure hadron masses (e.g. charged/neutral π and K)
 $a M_\pi(a m_u, a m_d) = a \times 0.139 6 \text{ GeV}$
- Bottom (& charm) quark needs special action for a large $a m_b \rightarrow$ effective quark actions
 HQET, NRQCD, Fermilab action (needs matching)



APPLICATION – SPECTRUM -

Lattice calculation vs experiments : hadron masses

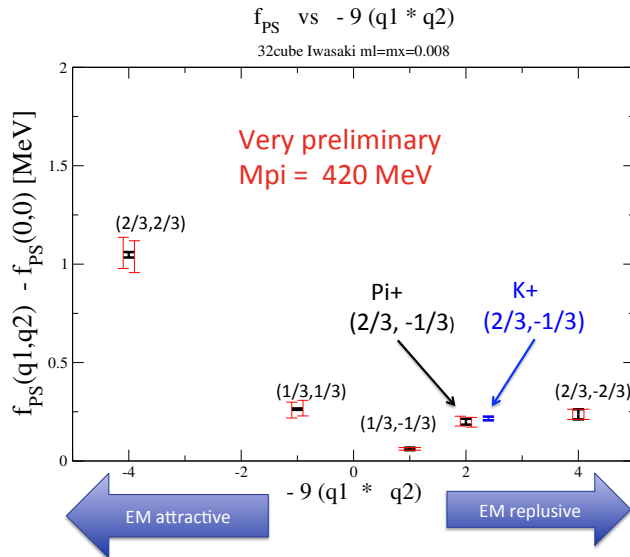


$\pi \dots \Omega$: BMW, MILC, PACS-CS, QCDSF; η - η' : RBC, UKQCD, Hadron Spectrum (ω);
 D, B : Fermilab, HPQCD, Mohler-Woloshyn

QCD+QED simulation

[T. Blum et al.]

EM effects on PS decay



Statistically well resolved by +e/-e averaging.

c.f. [Bijnens Danielsson 2006]

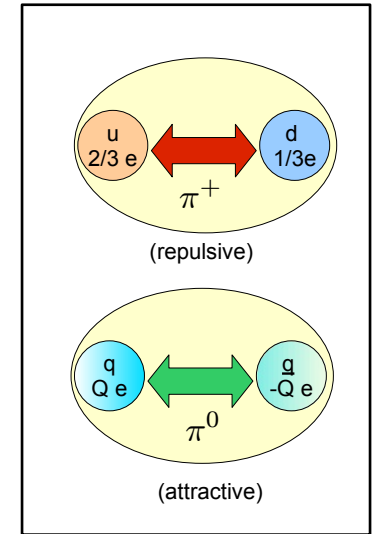
$$f_{\pi^+, NLO}/F_0 = 0.0039$$

$$f_{K^+, NLO}/F_0 = 0.0056$$

EM turned on, but $m_u = m_d$

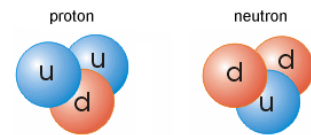
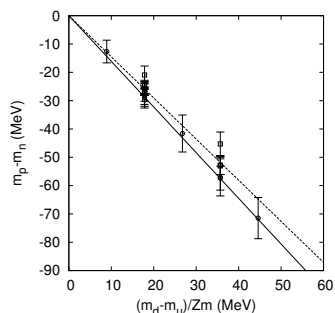
Iwasaki-DWF Nf=2+1,

$(2.7 \text{ fm})^3$, $a^{-1} \sim 2.3 \text{ GeV}$



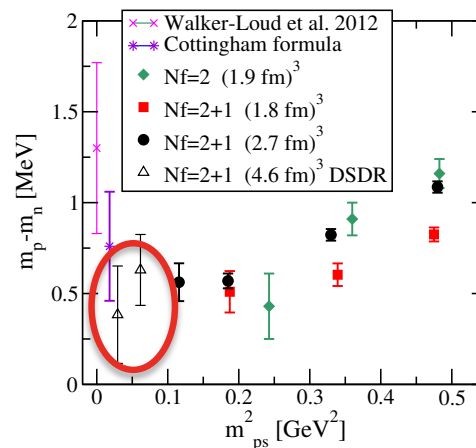
Proton / Neutron mass difference

$(m_u - m_d)$ effect



DSDR DWF Nf=2+1
 $(4.6 \text{ fm})^3$,
 $a^{-1} \sim 1.4 \text{ GeV}$

EM effect

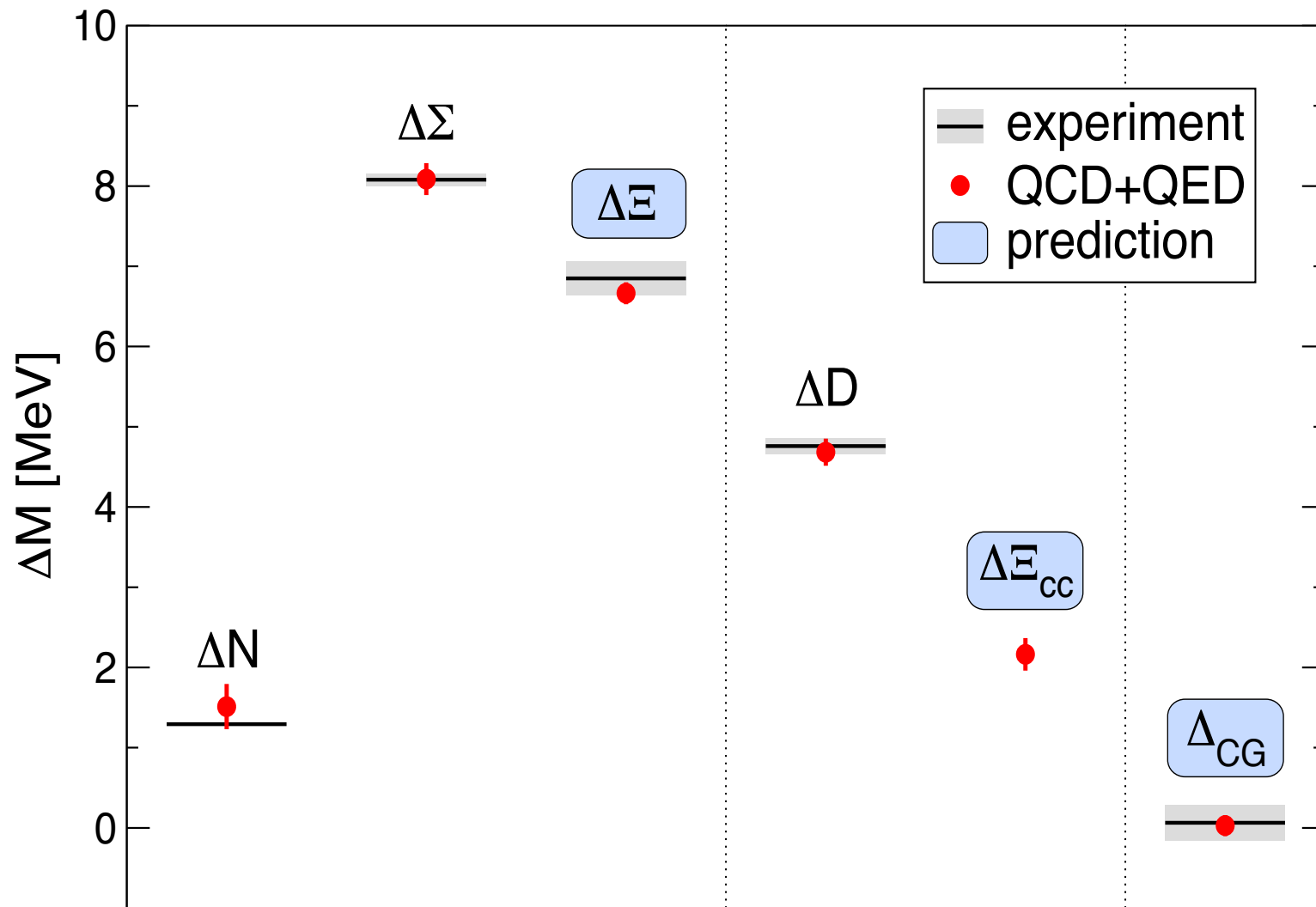


	$m_u - m_d$	EM
NPLQCD	2.26(72)	
BLUM	2.51(71)	0.54(24)
RM123	2.80(70)	
QCDSF-UKQCD	3.13(77)	

2.68(35) 0.54(24)

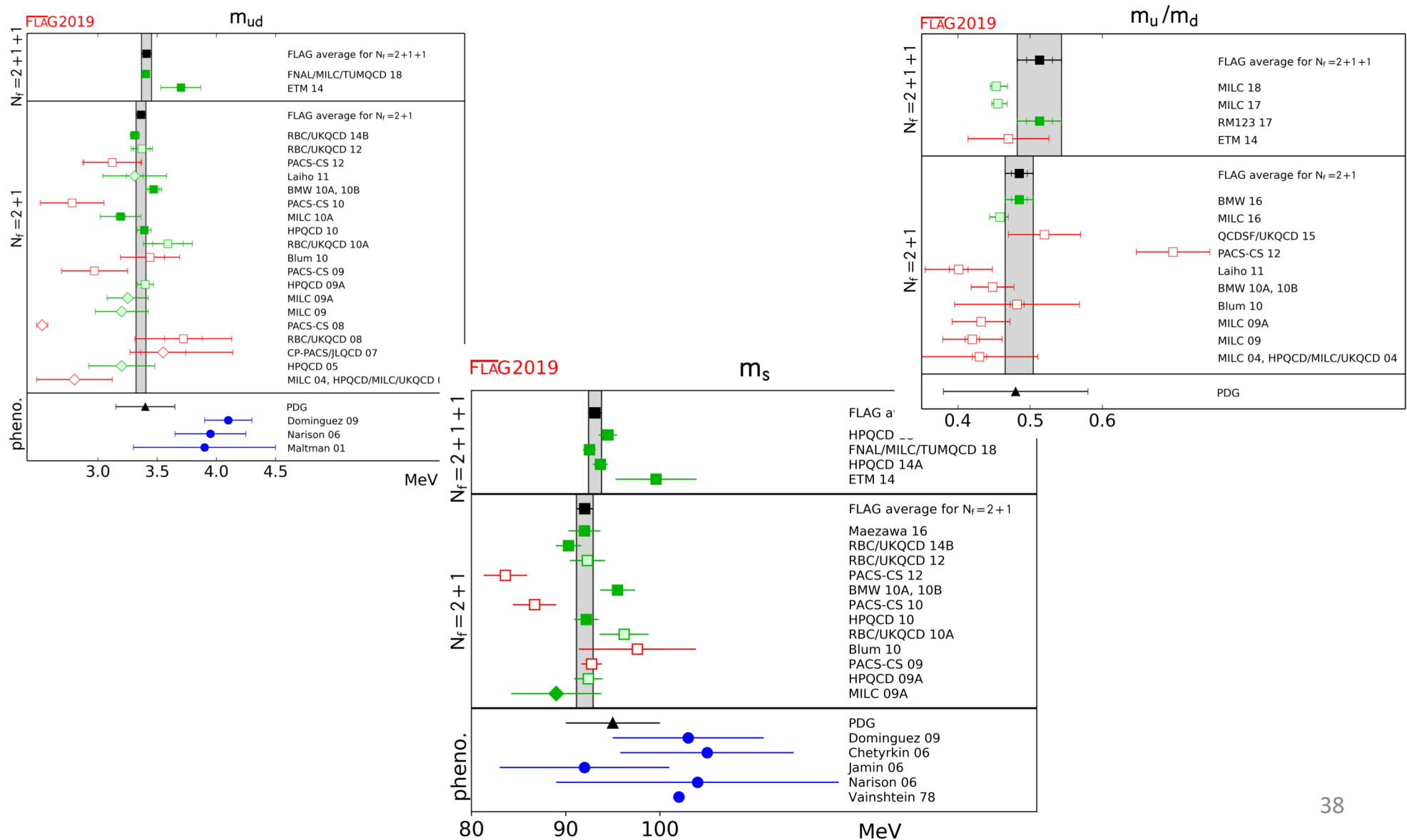
$\Rightarrow |M_N - M_p| = 2.14(42) \text{ MeV}$
(experiment: 1.2933321(4) MeV)

Isospin-breaking spectrum



quark mass results

■ FLAG (Flavor Lattice Averaging Group)



rating for systematic errors

S. Aoki et al.

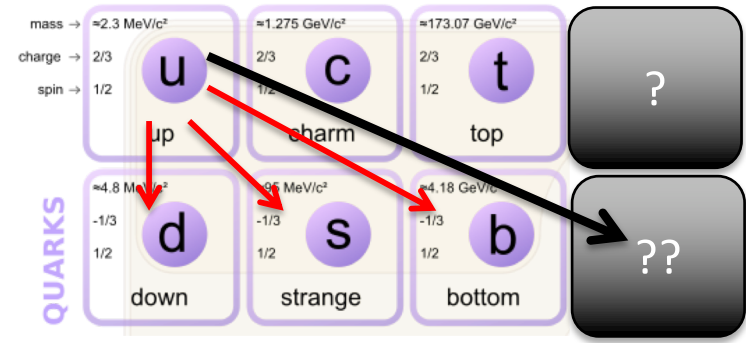
1902.08191

Collaboration	Ref.	publication status	chiral extrapolation	continuum extrapolation	finite volume	renormalization	running	m_{ud}	m_s
Maezawa 16	[33]	A	★	★	★	★	d	–	92.0(1.7)
RBC/UKQCD 14B [⊖]	[34]	A	★	★	★	★	d	3.31(4)(4)	90.3(0.9)(1.0)
RBC/UKQCD 12 [⊖]	[32]	A	★	○	★	★	d	3.37(9)(7)(1)(2)	92.3(1.9)(0.9)(0.4)(0.8)
PACS-CS 12*	[35]	A	★	■	■	★	b	3.12(24)(8)	83.60(0.58)(2.23)
Laiho 11	[36]	C	○	★	★	○	–	3.31(7)(20)(17)	94.2(1.4)(3.2)(4.7)
BMW 10A, 10B ⁺	[37, 38]	A	★	★	★	★	c	3.469(47)(48)	95.5(1.1)(1.5)
PACS-CS 10	[39]	A	★	■	■	★	b	2.78(27)	86.7(2.3)
MILC 10A	[40]	C	○	★	★	○	–	3.19(4)(5)(16)	–
HPQCD 10**	[41]	A	○	★	★	–	–	3.39(6)	92.2(1.3)
RBC/UKQCD 10A	[42]	A	○	○	★	★	a	3.59(13)(14)(8)	96.2(1.6)(0.2)(2.1)
Blum 10 [†]	[9]	A	○	■	○	★	–	3.44(12)(22)	97.6(2.9)(5.5)
PACS-CS 09	[43]	A	★	■	■	★	b	2.97(28)(3)	92.75(58)(95)
HPQCD 09A [⊕]	[44]	A	○	★	★	–	–	3.40(7)	92.4(1.5)
MILC 09A	[45]	C	○	★	★	○	–	3.25 (1)(7)(16)(0)	89.0(0.2)(1.6)(4.5)(0.1)
MILC 09	[46]	A	○	★	★	○	–	3.2(0)(1)(2)(0)	88(0)(3)(4)(0)
PACS-CS 08	[47]	A	★	■	■	■	–	2.527(47)	72.72(78)
RBC/UKQCD 08	[48]	A	○	■	★	★	–	3.72(16)(33)(18)	107.3(4.4)(9.7)(4.9)
CP-PACS/ JLQCD 07	[49]	A	■	★	★	■	–	3.55(19)(⁺⁵⁶ ₋₂₀)	90.1(4.3)(^{+16.7} _{-4.3})
HPQCD 05	[50]	A	○	○	○	○	–	3.2(0)(2)(2)(0) [‡]	87(0)(4)(4)(0) [‡]
MILC 04, HPQCD/ MILC/UKQCD 04	[51, 52]	A	○	○	○	■	–	2.8(0)(1)(3)(0)	76(0)(3)(7)(0)

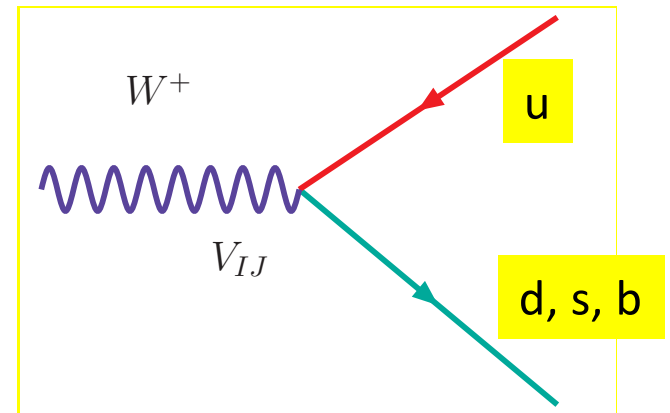
– For QCD:

- ★ $[M_{\pi,\min}/M_{\pi,\text{fid}}]^2 \exp\{4 - M_{\pi,\min}[L(M_{\pi,\min})]_{\max}\} < 1$, or at least three volumes
- $[M_{\pi,\min}/M_{\pi,\text{fid}}]^2 \exp\{3 - M_{\pi,\min}[L(M_{\pi,\min})]_{\max}\} < 1$, or at least two volumes
- otherwise

Quark flavor physics



- Cabibbo, Kobayashi and Maskawa (CKM) advocated quark “flavor” could change by emitting W boson (flavor mixing).
- up quark could change into the three down type quarks with mixing amplitude, called CKM matrix element V_{ij}
 down : V_{ud} ,
 strange : V_{us} ,
 bottom : V_{ub}
- Due to preservation of total probability,
 $|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1$
 (CKM unitarity condition)
- If this equation is not fulfilled, something **new and unknown**, such as 7th quark (or 4th generation) may be responsible

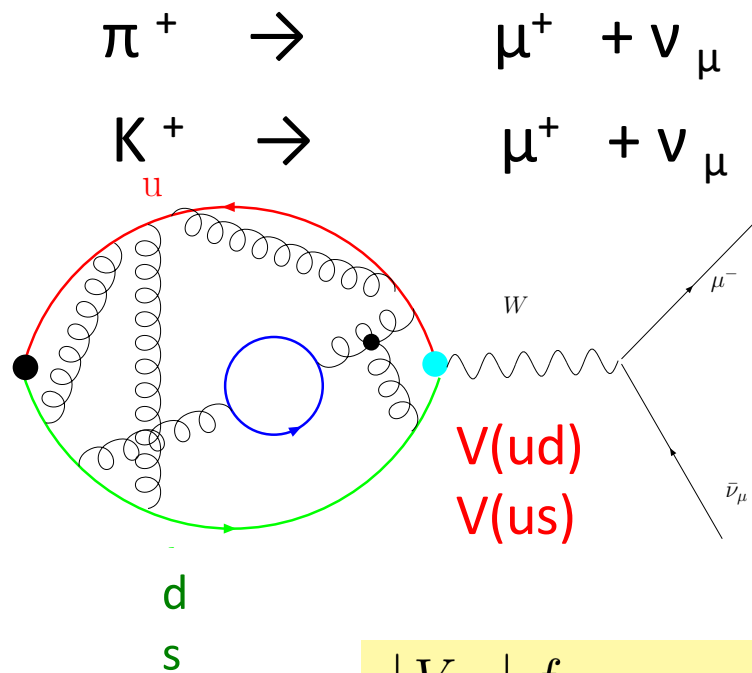


Sub-percent accuracy on Physical point

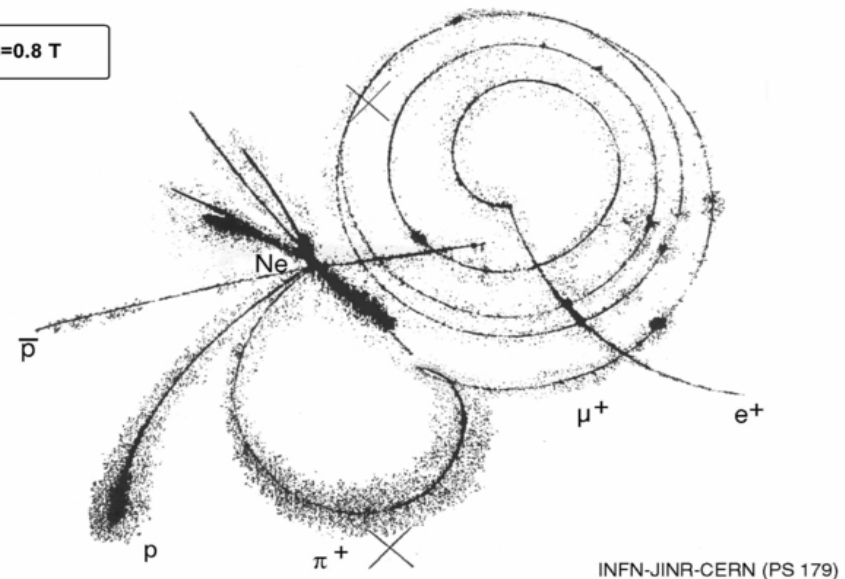
- charged π and K meson decay into muon and neutrino

$$\Gamma(P \rightarrow \ell \nu) = \frac{G_F^2}{8\pi} f_P^2 m_\ell^2 M_P \left(1 - \frac{m_\ell^2}{M_P^2}\right)^2 |V_{q_1 q_2}|^2 \quad (\text{experiment})$$

$$\langle 0 | \bar{q}_1 \gamma_\mu \gamma_5 q_2 | P(p) \rangle = i p_\mu f_P \quad (\text{lattice})$$

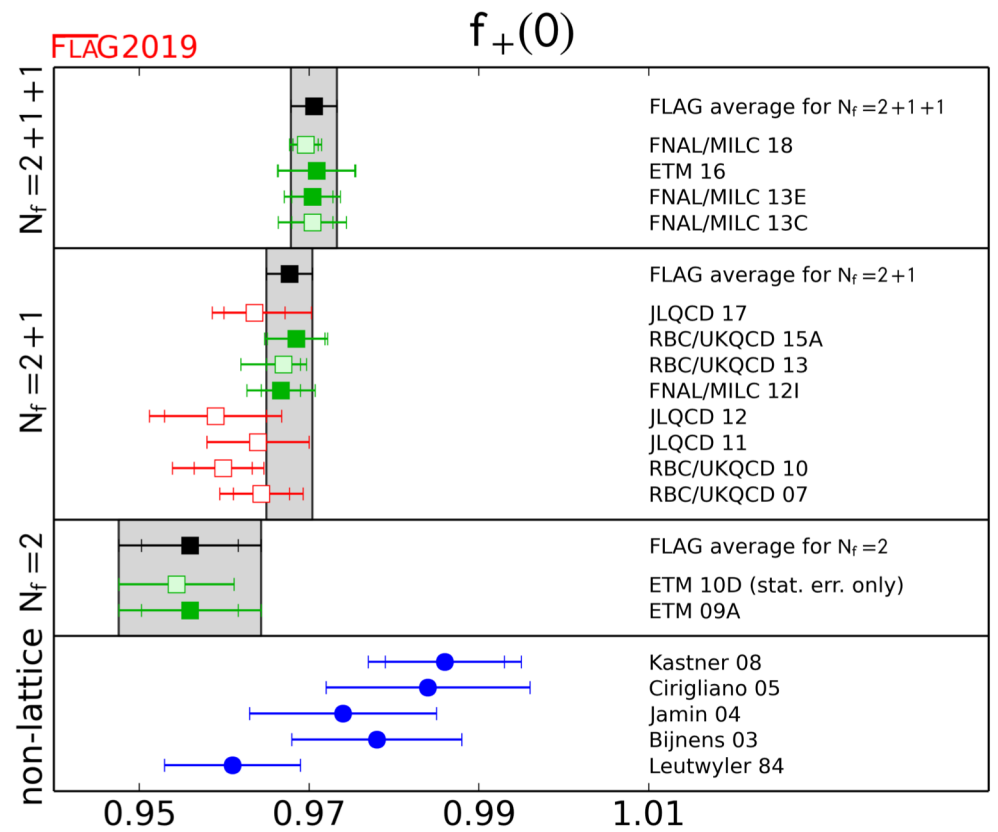
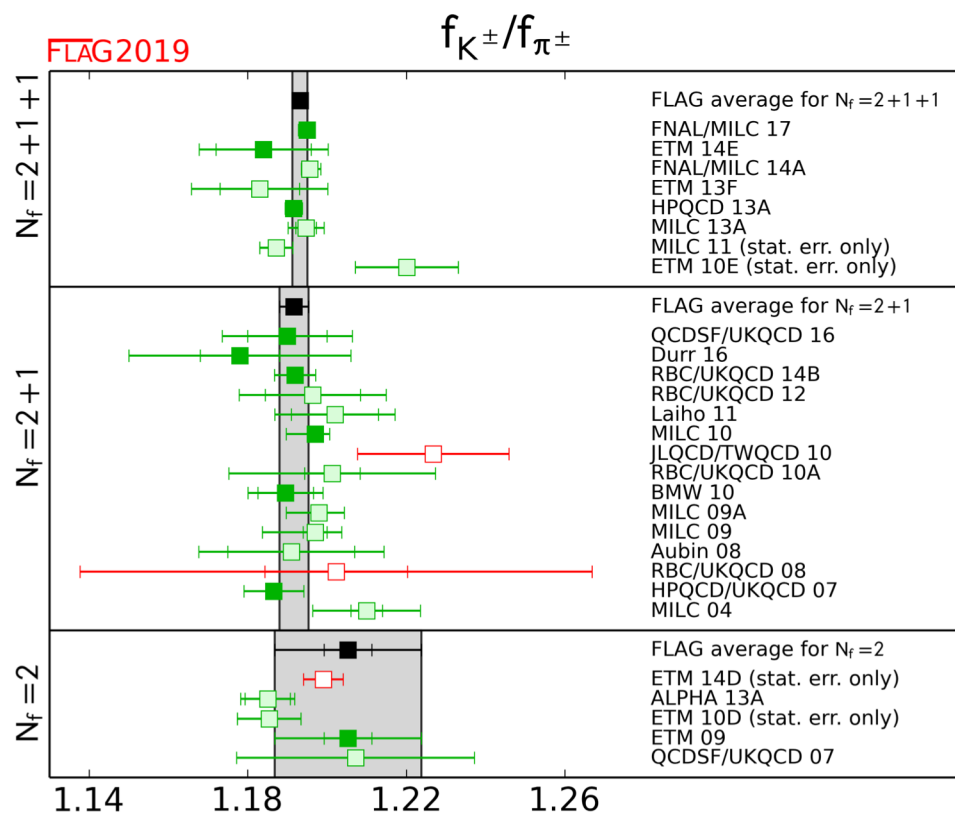


B=0.8 T



$$\left| \frac{V_{us}}{V_{ud}} \right| \frac{f_K}{f_\pi} = 0.27673 (29)_{\text{exp}} (23)_{\text{th}}$$

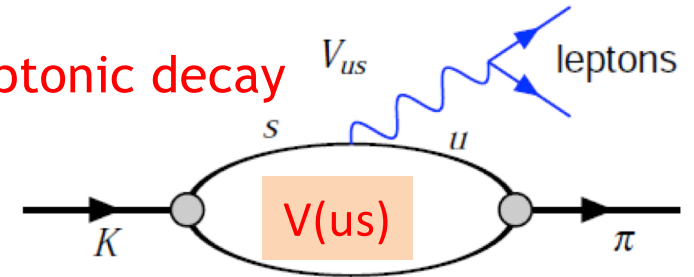
Experiment(FLAG2019)



A stringent check of Standard Model

- Together with similar analysis on K meson **semileptonic decay**

$$K^+ \rightarrow \pi^0 + l + \nu_l$$



- Super allowed Beta decay for [Hardy&Towner]

$$|V_{ud}| = 0.97420(21)$$

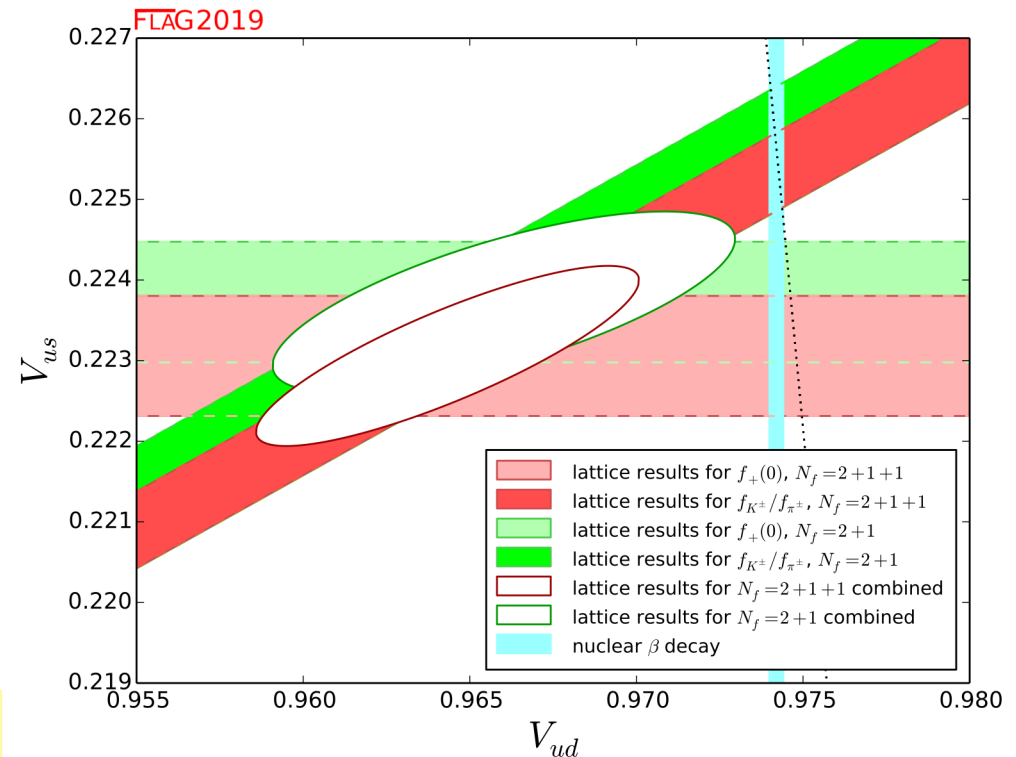
- V(ub) is negligible

- Very precise check for Standard Model, CKM unitarity [FLAG 2019]

$$\begin{aligned} & |V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 \\ &= 0.9797(74) \quad [2.7 \sigma] \quad (\text{Kl2+Kl3}) \\ &= 0.99884(53) \quad [2.2 \sigma] \quad (\text{Kl3+ B}) \\ &= 0.99986(46) \quad [0.3 \sigma] \quad (\text{Kl2+ B}) \end{aligned}$$

$$C_{KV_\mu D}(t_x, t_y; \vec{p}) = \sum_{i,j} \frac{1}{2m_{D_i} 2E_{K_j}(\vec{p})} e^{-m_{D_i} t_x - E_{K_j}(\vec{p}) |t_y|} \times$$

$$\times \langle 0 | O_K(t_x, \vec{x}) | K_i(\vec{p}) \rangle \langle K_i(\vec{p}) | V_\mu(0) | D_j(\vec{0}) \rangle \langle D_j(\vec{0}) | O_D^\dagger(0) | 0 \rangle.$$



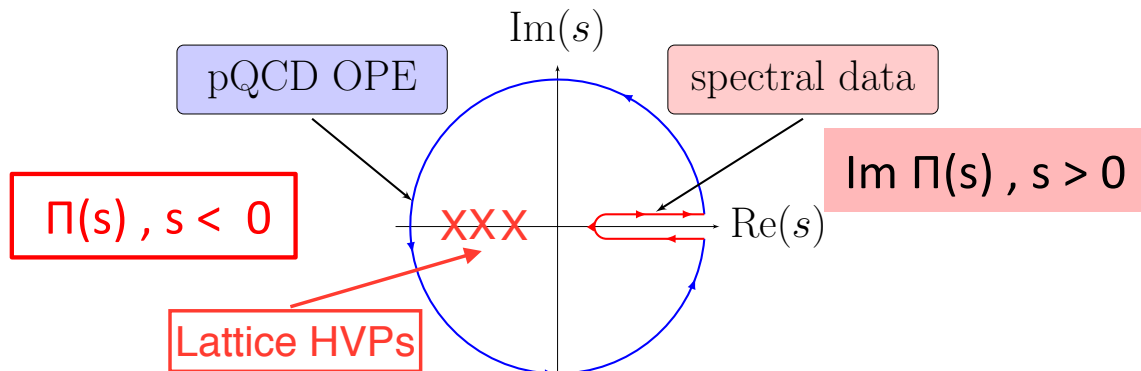
Inclusive tau decay [PRL 18]

[Hiroshi Ohki]

■ A new way to combine

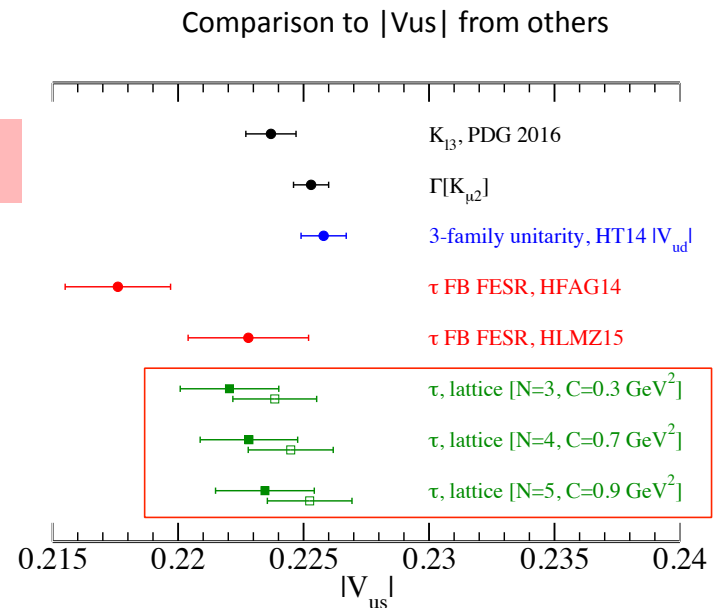
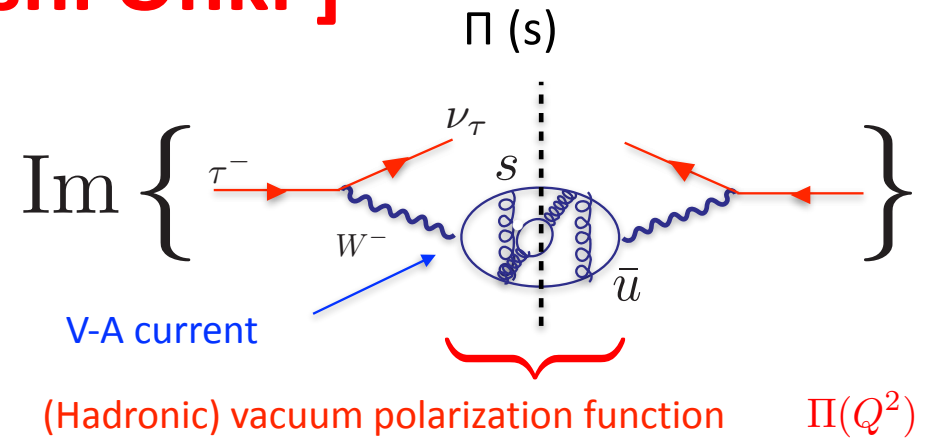
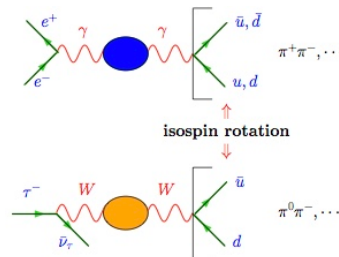
- $\tau \rightarrow \text{had} + \nu$ experimental spectrum
- Lattice vector/axial vacuum polarization **suppressing** quark-hadron duality violation

Unitarity and analytic continuation at work !



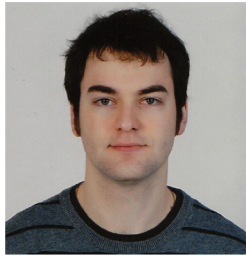
- Traditional analysis currently gives - 3.4 σ tension in V_{us}
- All our new results are consistent with CKM unitarity and has smaller errors

- Also τ input for muon g-2



Our result

All our results are consistent with each other within 1 σ error, as well as to CKM unitarity.

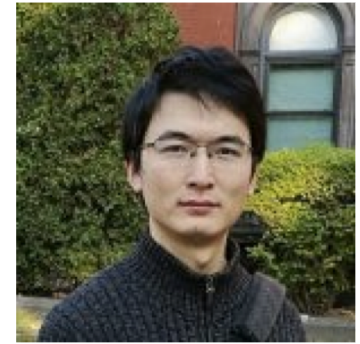


[Lehner, Blum, et al.]

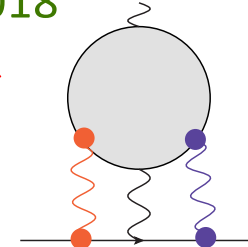
Muon g-2

3 PRLs including an Editors' suggestion
BNL Press Release on 6/13/2018

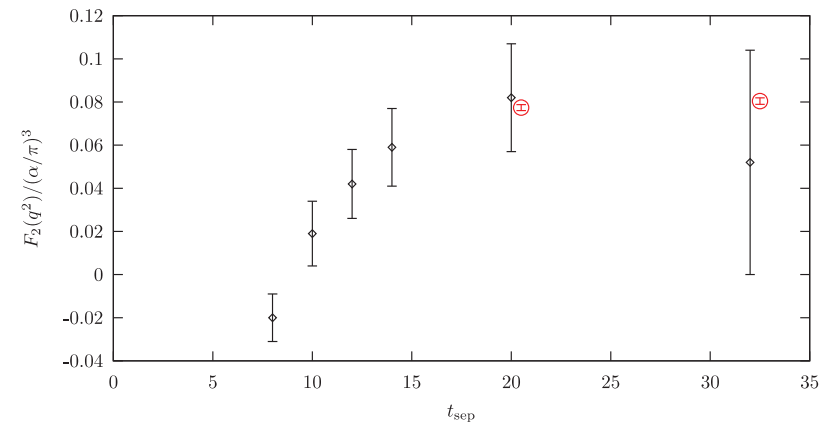
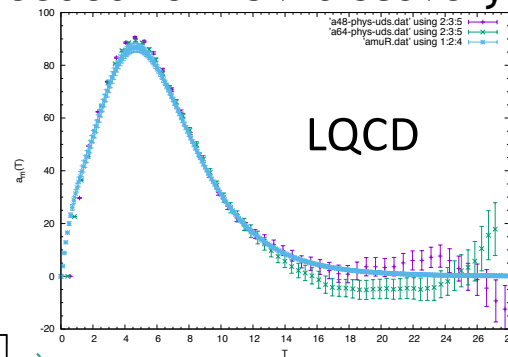
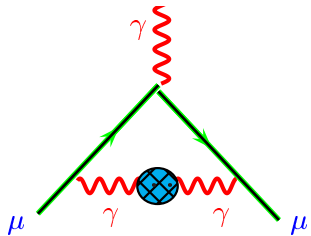
New Muon g-2 @ FNAL, J-PARC
to confirm $\sim 4 \sigma$ difference
from SM prediction



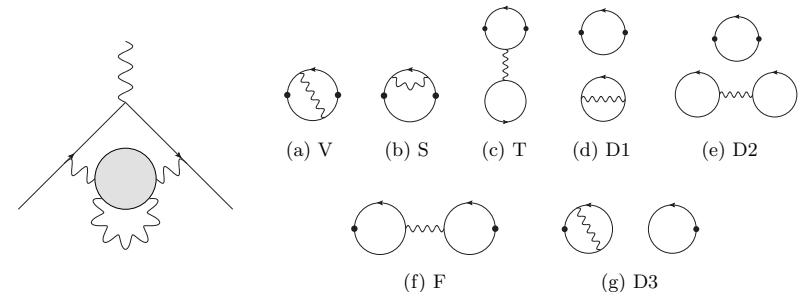
[Luchang's talk]



hadronic contributions needed for BSM discovery

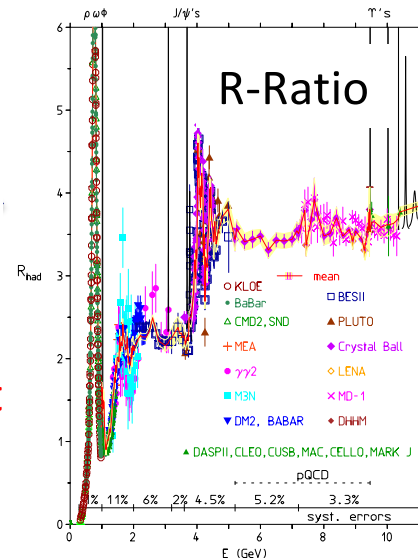


HLbL: x 1000 speed up
at physical quark mass



Most precise published result

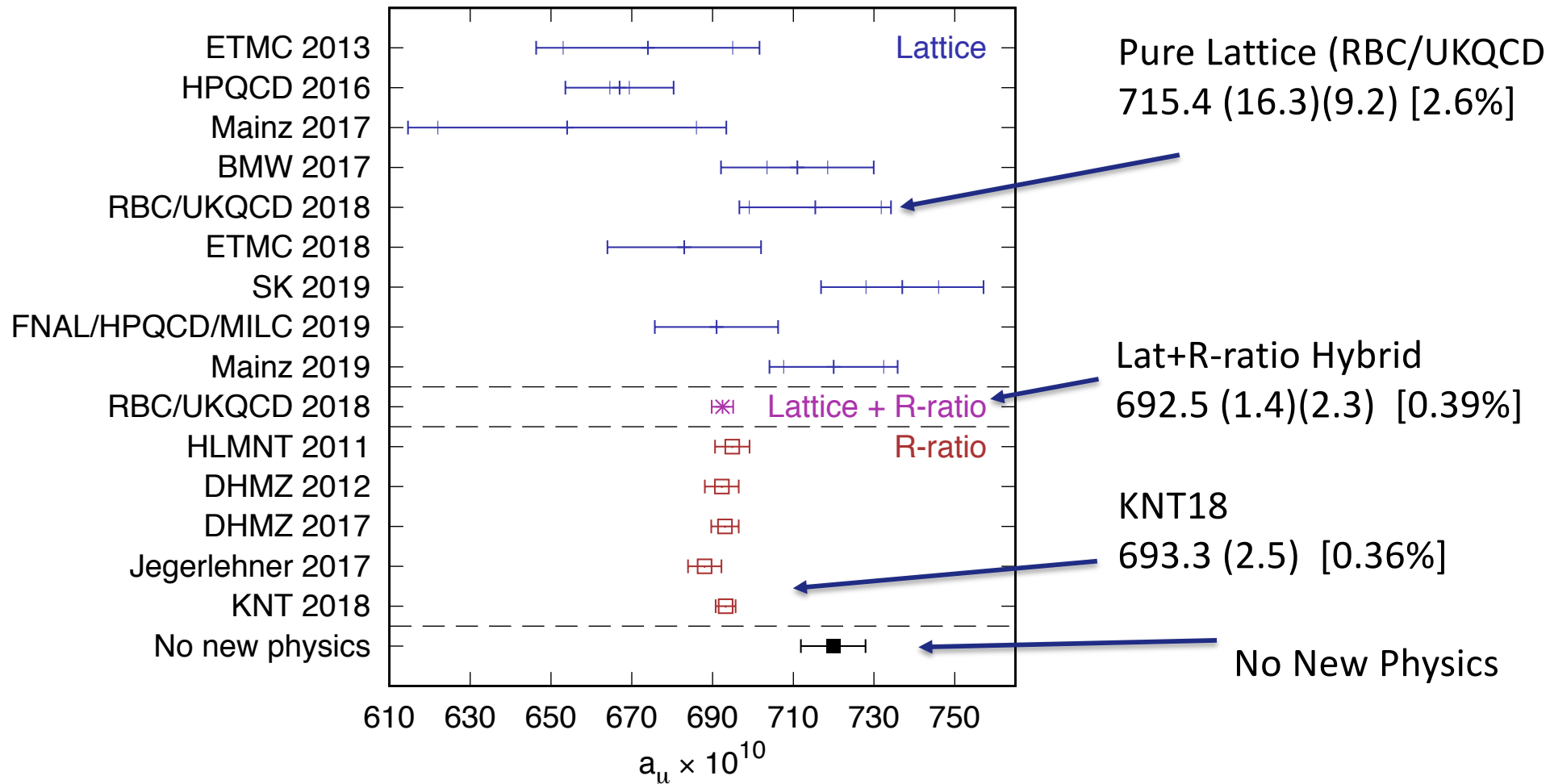
Cross-check each other
and get the precise and reliable
theoretical guidance to g-2 experiments



HVP: disconnected, QED & strong isospin
breaking in the first principle calculation

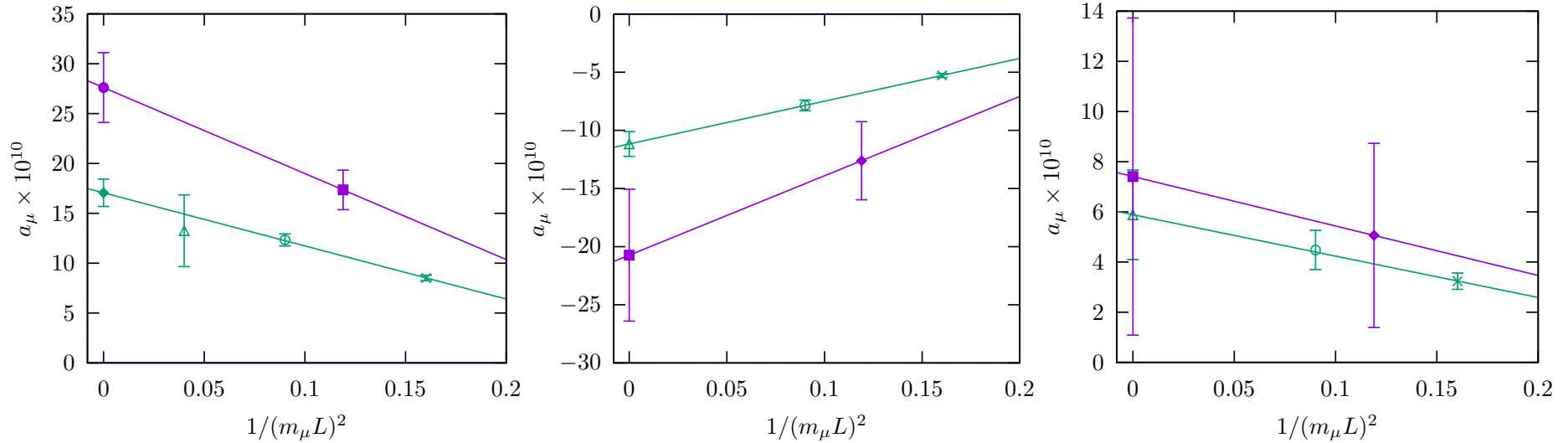
HVP results

[Christoph Lehner Lat19]

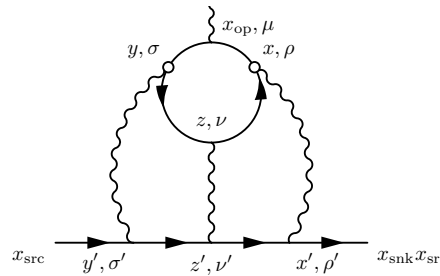


- Significant improvements is in progress for statistical error using 2π and 4π (!) states in addition to EM current (GEVP, GS-parametrization)
- Pure Lattice HVP 5×10^{-10} this year, 1×10^{-10} for long term
- Check BABAR-KLOE tension by window method, consolidate error at 3×10^{-10}

QED_L continuum and infinite volume extrapolation [Blum et al., 2019] (preliminary)



- Iwasaki ensembles: $a \rightarrow 0$ ($c_2 = 0$, conn. extrap.: up to 1 fm, 48^3 for $r > 1$ fm)
- I-DSDR ensembles: $L \rightarrow \infty$ ($b_2 = 0$)

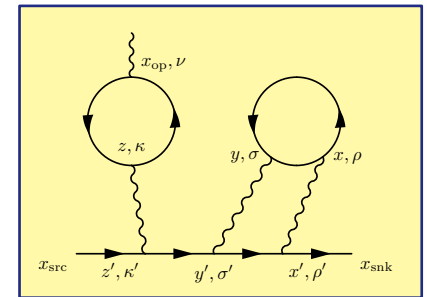


$$a_\mu^{cHLbL} = (27.61 \pm 3.51_{\text{stat}} \pm 0.32_{\text{sys}, a^2}) \times 10^{-10}$$

$$a_\mu^{dHLbL} = -20.20 \pm 5.65_{\text{stat}} \times 10^{-10}$$

$$a_\mu^{HLbL} = 7.41 \pm 6.32_{\text{stat}} \pm 0.32_{\text{sys}, a^2} \times 10^{-10}$$

$$F_2(a, L) = F_2 \left(1 - \frac{c_1}{(m_\mu L)^2} \right) (1 - c_2 a^2)$$



$K \rightarrow \pi \pi$ Decays

- quark flavor mixing (strange quark \rightarrow down quark) in **Cabibbo Kobayashi Maskawa Theory** (2008 Nobel prize)

- First direct methods (Lellouch Luscher formula)  **Nobelprize**
PRL 108, 141601 (2012)

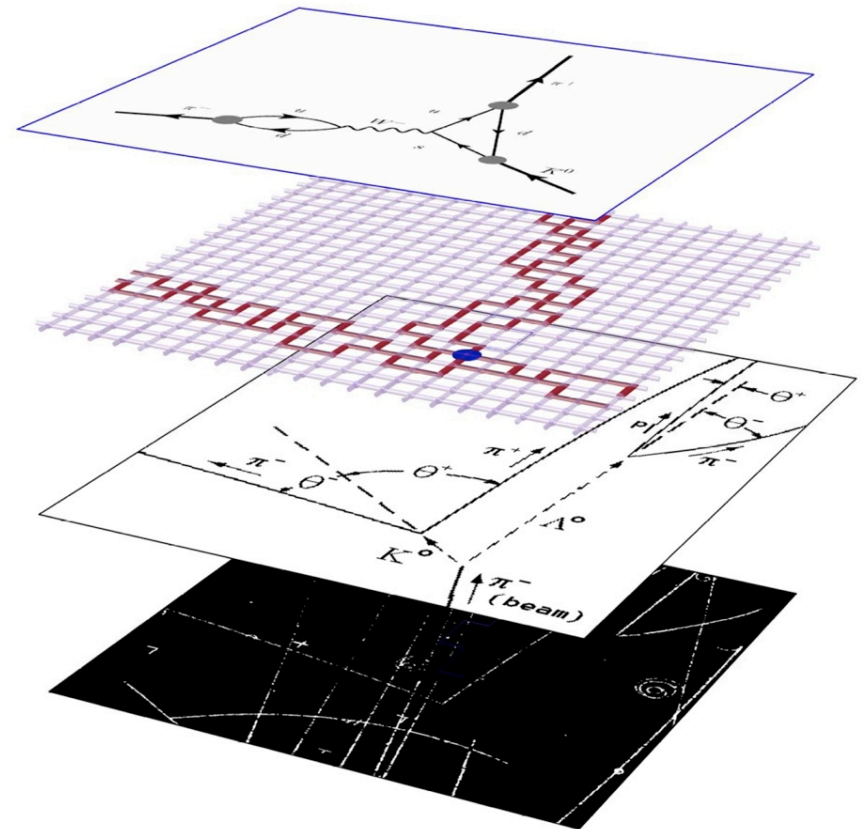
- Finer lattice spacing
~15% error \rightarrow ~5% error

- Realistic kinematics for $I=0$ by G-parity boundary condition

40 years old home work : ϵ'/ϵ

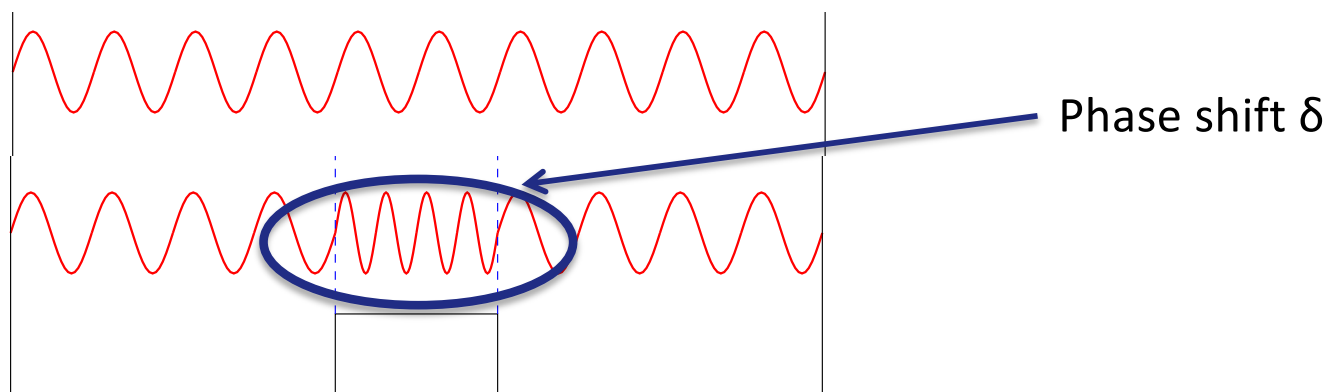
- K_L K_S mass difference rare Kaon decay ?

2012 Ken Wilson Lattice Award



K \rightarrow $\pi\pi$ decay on lattice

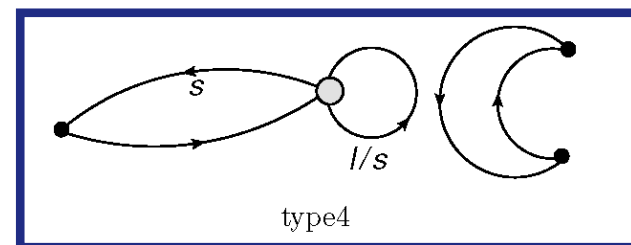
- Relates **energy on finite volume** $E_{\pi\pi}$ (V) to **phase shift** δ to obtain **complex** $\text{Amp}(K \rightarrow \pi\pi) = |A_l| e^{i\delta}$ (Lüscher, Lellouch-Lüscher)



- Momentum of pions are controlled by boundary condition (anti-periodic or G-parity b.c.)
- Mixing and Renormalization of operator is done using non-perturbative renormalization (NPR)
- Chiral Symmetry is curtail
- $l=2$ channel is under control, $l=0$ is still a challenge due to disconnected diagrams. [PRL (2015). update this year]

$$\text{Re} \left(\frac{\epsilon'}{\epsilon} \right)_{\text{SM}} = 1.38(5.15)(4.43) \times 10^{-4}$$

cf. $\text{Re} \left(\frac{\epsilon'}{\epsilon} \right)_{\text{exp}} = 16.6(2.3) \times 10^{-4}$



B physics

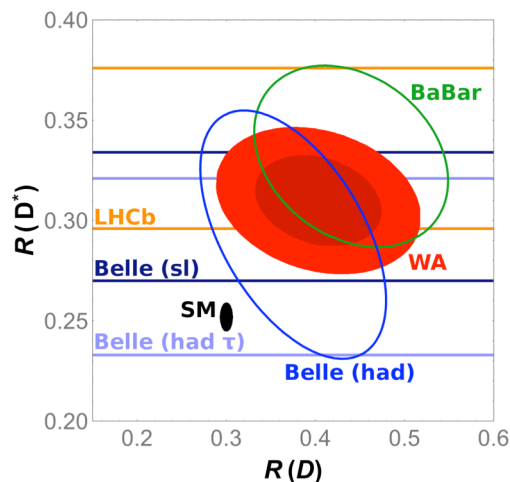
- Provide crucial theoretical guidance for B factories (LHCb and Belle-II), CKM physics

- Semileptonic form factor $B \rightarrow \pi l \nu$
- $1 \rightarrow 2$ transition (FCNC)

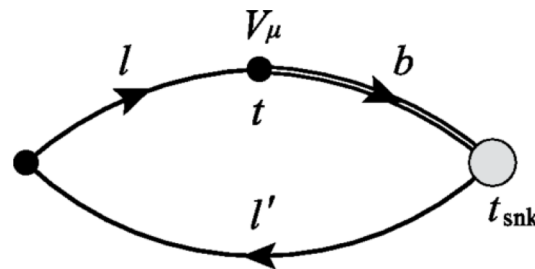
$$B \rightarrow K\pi l^+ l^-, \pi\pi l^+ \nu$$

- $|V_{cs}|$ from $\Lambda_c \rightarrow \Lambda$ [Meinel, PRL (2017)]
- $|V_{ub}|$ from $\Lambda_b \rightarrow p$ decay

V_{ud} $\pi \rightarrow l\nu$	V_{us} $K \rightarrow l\nu$ $K \rightarrow \pi l\nu$	V_{ub} $B \rightarrow \pi l\nu$
V_{cd} $D \rightarrow l\nu$ $D \rightarrow \pi l\nu$	V_{cs} $D_s \rightarrow l\nu$ $D \rightarrow K l\nu$	V_{cb} $B \rightarrow D l\nu$ $B \rightarrow D^* l\nu$
V_{td} $B_d \leftrightarrow \bar{B}_d$	V_{ts} $B_s \leftrightarrow \bar{B}_s$	V_{tb}

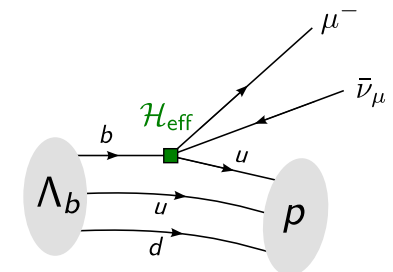


$B \rightarrow D\tau\nu$:



$$\mathcal{R}_D^{SM} = \frac{\mathcal{B}(\bar{B} \rightarrow D\tau^- \bar{\nu}_\tau)}{\mathcal{B}(\bar{B} \rightarrow De^- \bar{\nu}_e)} = 0.300 \pm 0.008$$

$$\mathcal{R}_{D^*}^{SM} = \frac{\mathcal{B}(\bar{B} \rightarrow D^*\tau^- \bar{\nu}_\tau)}{\mathcal{B}(\bar{B} \rightarrow D^*e^- \bar{\nu}_e)} = 0.252 \pm 0.003$$



Summary & Conclusions

- LQCD provides non-perturbative ab-initio analysis of QFT. Systematically improvable.
- LQCD now delivers a lot of promised results essential to make experiments meaningful, enables a lot of important physics analysis which has not been possible before.
- Wide spectrum of important studies in Particle and Nuclear Physics. Many basic quantities are determined $< 1\%$! Precision with reliable error/
- Hot topics : B physics, muon g-2, multi-hadron including $K \rightarrow \pi\pi$, EW box, CP violation in Nucleon (EDM), ...

Stay tuned ! Thank you !