

# Lecture 3: Factorization & Resummation in SCET

Tues. Aug 6, 2019 (pm)

Have to compute matching condition:

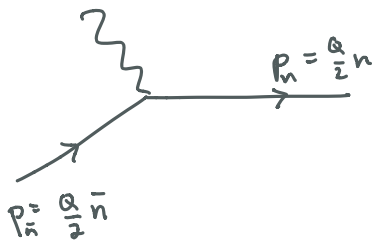
$$j_{\text{QCD}}^\mu = \bar{q} \gamma^\mu q \rightarrow j_{\text{SCET}}^\mu \sum_{\substack{n_1, n_2 \\ \tilde{p}_1, \tilde{p}_2}} C_2(p_1, p_2, \mu) \mathcal{O}_{n_1 n_2}^\mu(\tilde{p}_1, \tilde{p}_2)$$

$$\mathcal{O}_{2, n_1 n_2}^\mu(\tilde{p}_1, \tilde{p}_2) = \bar{\chi}_{n_1, p_1} \gamma_\perp^\mu \chi_{n_2, p_2} \xrightarrow{\text{red}} \bar{\chi}_{n_1, p_1} \gamma_{\perp 1}^\mu \gamma_{\perp 2}^\mu \chi_{n_2, p_2}$$

where  $\chi_{n, p} = [W_n^\dagger \xi_n]_p$

Match matrix elements using any external states w/ nonzero overlap w/ operator.

e.g. DIS



Tree level:

$$\langle q(p_n) | j^\mu | q(p_n) \rangle = \bar{u}_n \gamma_\perp^\mu u_n$$

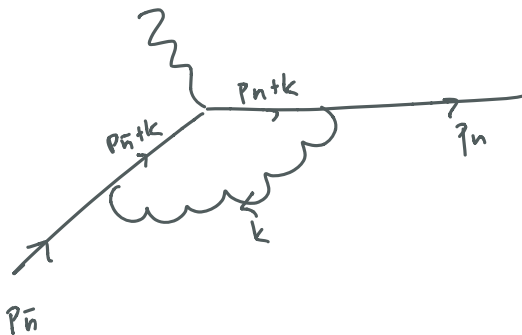
SCET:

$$\langle q_n | \bar{\chi}_n \gamma_\perp^\mu \chi_n | q_n \rangle = \text{same}$$

$$\Rightarrow C_2 = 1.$$

1-loop:

QCD:



$$= \bar{u}_n i g_s \mu^{2\epsilon} \gamma^\alpha t^a \int \frac{d^d k}{(2\pi)^D} \frac{i(p_n+k)}{(p_n+k)^2 + i\epsilon} \gamma^\mu \frac{i(p_n+k)}{(p_n+k)^2 + i\epsilon} i g_s \mu^{2\epsilon} \gamma_\alpha t^a u_n \frac{-i}{k^2 + i\epsilon}$$

$$= -i g_s^2 \mu^{2\epsilon} C_F \int \frac{d^d k}{(2\pi)^D} \frac{\bar{u}_n \gamma^\alpha (p_n+k) \gamma^\mu (\tilde{p}_n+k) \gamma_\alpha u_n}{(k^2 + i\epsilon) (k^2 + 2p_n \cdot k + i\epsilon) (k^2 + 2\tilde{p}_n \cdot k + i\epsilon)}$$

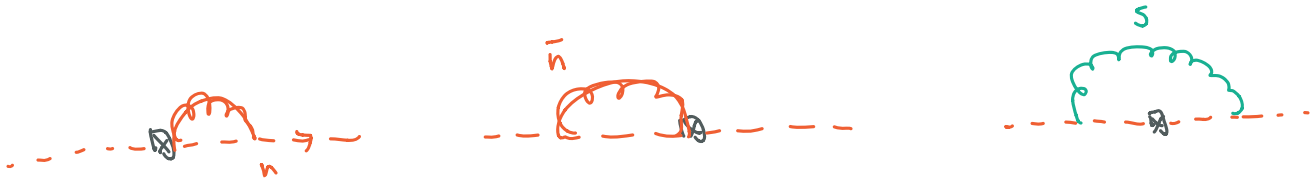
... [exercise; remember  $(\frac{1}{\sqrt{2L}})^2$    $Z_2 = 1 + \frac{d_{\text{SCET}}}{4\pi} (\frac{1}{\epsilon_{UV}} - \frac{1}{\epsilon_{IR}})$ ]

$$\overline{\text{MS}} = \frac{\alpha_s C_F}{4\pi} \left[ -\frac{2}{\epsilon_{IR}^2} - \frac{2 \ln \frac{\mu^2}{Q^2} + 3}{\epsilon_{IR}} - \ln^2 \frac{\mu^2}{Q^2} - 3 \ln \frac{\mu^2}{Q^2} - 8 + \frac{\pi^2}{6} \right]$$

D=4-2ε

no  $\frac{1}{\epsilon_{UV}}$ :  $j^\mu$  not renormalized due to current conservation

SCE1: 3 diagrams to compute:



(w/ renorm ( $Z_2$  same as in QCD))

In pure DR all dim scales  $\Rightarrow 0$

Technically

$$A_n + A_{\bar{n}} + A_5 \propto \left( \frac{2}{\epsilon_{UV}^2} - \frac{2}{\epsilon_{IR}^2} + (3 - \ln \frac{\mu^2}{Q^2}) \left( \frac{1}{\epsilon_{UV}} - \frac{1}{\epsilon_{IR}} \right) \right)$$

$\Rightarrow$  QCD - SCE1  $\Rightarrow$

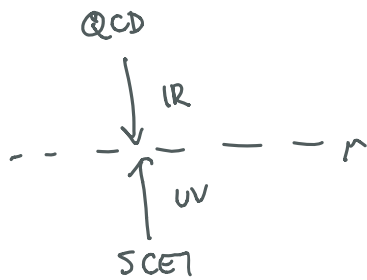
$$Z_2 = 1 + \frac{\alpha_S C_F}{4\pi} \left( -\ln^2 \frac{\mu^2}{Q^2} - 3 \ln \frac{\mu^2}{Q^2} - 8 + \frac{\pi^2}{6} \right)$$

( $Q^2 \rightarrow -Q^2$  etc or pp)

$$Z_2 = 1 + \frac{\alpha_S C_F}{4\pi} \left( -\frac{2}{\epsilon_{UV}^2} - \frac{3}{\epsilon_{UV}} - \frac{2}{\epsilon_{UV}} \ln \frac{\mu^2}{Q^2} \right)$$

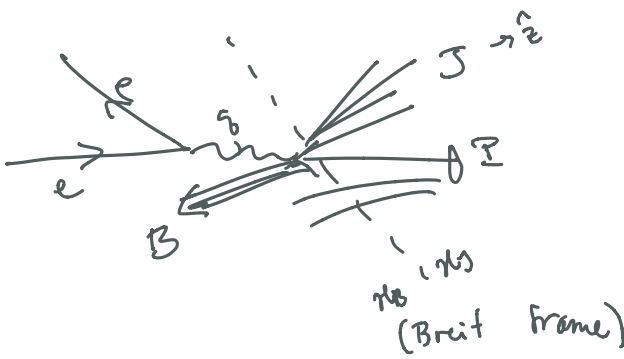
$\hookrightarrow$  IR div in QCD turns into UV div in SCE1

$\Rightarrow$  RGE to resum IR logs



Exercise: put in explicit IR regulator and repeat,  
e.g. off shellness  $p_n^2, p_{\bar{n}}^2 \neq 0$ .

Consider measurement of DIS "thrust":



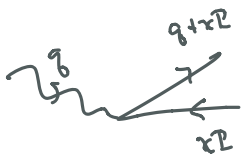
"y-tilde thrust":

$$\tau_1 = \frac{2}{Q^2} \sum_{i \in X} \min \{ q_{B \cdot p_i}, q_{J \cdot p_i} \}$$

here  $q_B = \kappa P$  (other choices possible) [1303.4952]

$$q_J = q + \kappa P$$

$$\tau_1 = \frac{1}{Q} \sum_{i \in X} \max \{ \tilde{n}_z \cdot p_i, n_z \cdot p_i \} = \text{DIS thrust}$$



exact partonic limit  $\tau_1 \rightarrow 0$ .

$$= 1 - \frac{2}{Q} \sum_{i \in X} p_z^i$$

[exercise]

(analogous to back-to-back jets in  $e^+e^-$  )

full QCD:

$$\frac{d\sigma}{dx dQ^2 dt} = L_{\mu\nu}(x, Q^2) W^{\mu\nu}(x, Q^2, \tau)$$

$$W^{\mu\nu} = \sum_X \langle P | J^{\mu\dagger} | X \rangle \langle X | J^\nu | P \rangle (2\pi)^4 (\mathbb{1} + \not{q} - \not{p}_X) \delta(\tau - z(x))$$

$$J^\mu(x) = \bar{q} \gamma^\mu q$$

$$\sum_{n_1, n_2} \int d^3 \tilde{p}_1 d^3 \tilde{p}_2 e^{i(\tilde{q}_1 - \tilde{p}_2) \cdot x} C_{\alpha\beta}^\mu(\tilde{p}_1, \tilde{p}_2) \Theta_{\tilde{q}\tilde{p}}^{\alpha\beta}(\tilde{p}_1, \tilde{p}_2; x)$$

where  $\Theta_{\tilde{q}\tilde{p}}^{\alpha\beta} = \bar{\chi}_{n_1, \tilde{p}_1}^\alpha(x) \chi_{n_2, \tilde{p}_2}^\beta(x)$

↓ redef

$$\bar{\chi}_{n_1, \tilde{p}_1}^{\alpha'} \chi_{n_1}^\dagger \chi_{n_2} \chi_{n_2, \tilde{p}_2}^\beta$$

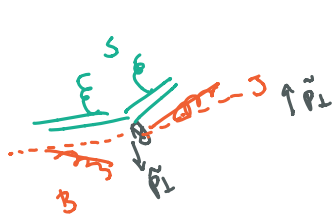
$$W^{\mu\nu} = \int d^4 x e^{i\tilde{q} \cdot x} \sum_{\substack{n_1, n_2 \\ n_1' n_2'}} \int d^3 \tilde{p}_1 d^3 \tilde{p}_2 d^3 \tilde{p}_1' d^3 \tilde{p}_2' e^{i(\tilde{q}_1 - \tilde{p}_1') \cdot x} \int dt_3 dt_4 dt_5 \delta(\tau - \tau_3 - \tau_4 - \tau_5) \\ \times \langle P_{n_B} | \bar{C}_{\nu}^{\alpha\beta}(\tilde{p}_1, \tilde{p}_2) \bar{\chi}_{n_2, \tilde{p}_2}^\beta \chi_{n_1}^\dagger \chi_{n_1}^\alpha \chi_{n_1, \tilde{p}_1}^\alpha(x) \delta(\tau_3 - \frac{n_3 \cdot \hat{p}_3}{Q}) \delta(\tau_4 - \frac{n_4 \cdot \hat{p}_4}{Q}) \\ \times \delta(\tau_5 - \frac{n_5 \cdot \hat{p}_5}{Q} - \frac{n_5 \cdot \hat{p}_5'}{Q}) C_{\nu}^{\alpha'\beta'}(\tilde{p}_1', \tilde{p}_2') \bar{\chi}_{n_1, \tilde{p}_1'}^{\alpha'} \chi_{n_1'}^\dagger \chi_{n_2} \chi_{n_2, \tilde{p}_2'}^{\beta'}(0) | P_{n_A} \rangle$$

( $\bar{c} = \gamma^0 \gamma^c$ )

(momentum conservation,  $n_{J,B}$  coll & soft decoupling)

$$\begin{aligned}
 &= Q^2 \int dt_J dt_B dt_S \delta(\tau - t_J - t_B - t_S) \bar{C}_{\mu}^{\beta\alpha} C_{\nu}^{\alpha'\rho'} \int d^2 \vec{p}_B \\
 &\quad \times \langle 0 | [\chi_{n_B}^{\dagger} \chi_{n_J}] (0) \delta(Q t_S - n_J \cdot \hat{p}_J^S - n_B \cdot \hat{p}_B^S) (\chi_{n_J}^{\dagger} \chi_{n_B}) (0) | 0 \rangle \\
 &\quad \times \left\{ \langle \mathbb{P}_{n_B} | \bar{\chi}_{n_B}^{\alpha} (0) \delta(Q t_B - n_B \cdot \hat{p}_B^{\alpha}) \delta(\bar{n}_B \cdot \hat{q} + \bar{n}_B \cdot \mathcal{P}) \delta^2(\vec{p}_B - \vec{p}_B) \chi_{n_B}^{\alpha'} (0) | \mathbb{P}_{n_B} \rangle \right. \\
 &\quad \times \langle 0 | \chi_{n_J}^{\alpha} (0) \delta(Q t_J - n_J \cdot \hat{p}_J^{\alpha}) \delta(\bar{n}_J \cdot \hat{q} + \bar{n}_J \cdot \mathcal{P}) \delta^2(\vec{p}_J + \vec{p}_J) \chi_{n_J}^{\alpha'} (0) | 0 \rangle \\
 &\quad \left. + \chi \leftrightarrow \bar{\chi} \right\}
 \end{aligned}$$

$$= \frac{d\sigma_0}{d\tau d\alpha^2} \int dt_J dt_B dt_S \delta\left(\tau - \frac{t_J}{Q^2} - \frac{t_B}{Q^2} - \frac{t_S}{Q^2}\right) S(k_S, \mu)$$



$$\times \int d^2 p_{\perp} J_g(t_J - \vec{p}_{\perp}^2, \mu) \left[ H_g(\alpha^2, \mu) B_g(t_B, \tau, \vec{p}_{\perp}^2, \mu) + H_{\bar{g}}(\alpha^2, \mu) B_{\bar{g}}(t_B, \tau, \vec{p}_{\perp}^2, \mu) \right]$$

where  $\frac{d\sigma_0}{d\tau d\alpha^2} = \frac{2\pi \alpha_{em}^2}{Q^4} [(1-y)^2 + 1]$  where  $y = \frac{\alpha^2}{xS}$

Hard function:  $H_{g, \bar{g}}(\alpha^2, \mu) = |C(\alpha^2, \mu)|^2 \cdot Q_f^2$  ( $\gamma^*$  channel only)

Soft function:  $S(k_S, \mu) = \frac{1}{N_c} \text{Tr} \sum_{X_S} |\langle X_S | \chi_{n_J}^{\dagger} \chi_{n_B} | 0 \rangle|^2 \delta(k_S - \sum_{i \in X_S} [\theta(n_{J_i} \cdot k - n_{J_i} \cdot k) n_{J_i} \cdot k + \theta(n_{B_i} \cdot k - n_{B_i} \cdot k) n_{B_i} \cdot k])$

Jet function:  $J_g(\vec{p}^- \cdot k^{\dagger} + \vec{p}_{\perp}^2, \mu) = \frac{(2\pi)^2}{N_c} \int \frac{dy^-}{2\vec{p}^-} e^{ik^{\dagger} y^- / 2} \text{Tr} \langle 0 | \frac{\bar{X}}{2} \chi_n(y^- \frac{n}{2}) \delta(\vec{p}^- + \bar{n} \cdot \mathcal{P}) \delta^2(\vec{p}_{\perp}^- + \vec{p}_{\perp}) \bar{\chi}_n(0) | 0 \rangle$

$(J_{\bar{g}} = J_g \quad \chi \leftrightarrow \bar{\chi})$

Beam function:

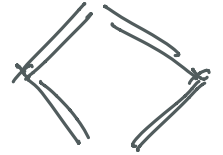
$$B_g(\vec{p}^- \cdot k^{\dagger}, \frac{\vec{p}_{\perp}^-}{P^-}, \vec{p}_{\perp}^2, \mu) = \frac{\theta(\vec{p}^-)}{\vec{p}^-} \int \frac{dy^-}{4\pi} e^{ik^{\dagger} y^- / 2} \times \langle \mathbb{P}_n(P^-) | \bar{\chi}_n(y^- \frac{n}{2}) \frac{\bar{X}}{2} [\delta(\vec{p}^- - \bar{n} \cdot \mathcal{P}) \frac{1}{\pi} \delta(\vec{p}_{\perp}^2 - \mathcal{P}_{\perp}^2) \chi_n(0)] | \mathbb{P}_n(P^-) \rangle$$

$(B_{\bar{g}} \quad \chi \leftrightarrow \bar{\chi})$

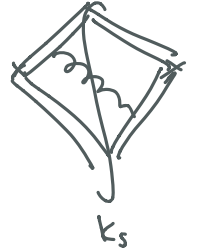
1-loop results:

$$H(Q^2, \mu) = 1 + \frac{\alpha_s(\mu) C_F}{2\pi} \left( -\ln^2 \frac{\mu^2}{Q^2} - 3 \ln \frac{\mu^2}{Q^2} - 8 + \frac{\pi^2}{6} \right)$$

$$S(k_s, \mu) = \delta(k_s)$$

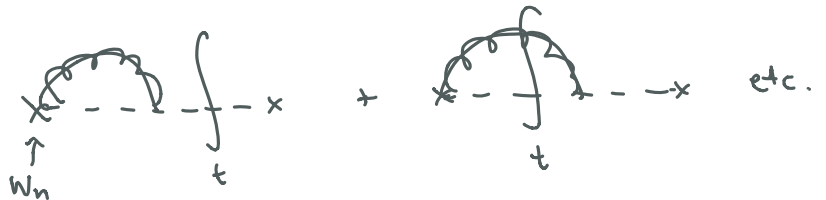


$$+ \frac{\alpha_s(\mu) C_F}{4\pi} \left\{ \frac{\pi^2}{3} \delta(k) - \frac{16}{\mu} \left[ \frac{\theta(k) g_n(k/\mu)}{k/\mu} \right]_+ \right\}$$



$$J(t, \mu) = \delta(t) \quad x \text{ --- } \int \text{---} x$$

$$+ \frac{\alpha_s(\mu) C_F}{4\pi} \left\{ (7 - \pi^2) \delta(t) - \frac{3}{\mu^2} \left[ \frac{\theta(t)}{t/\mu^2} \right]_+ + \frac{4}{\mu^2} \left[ \frac{\theta(t) \ln(t/\mu^2)}{t/\mu^2} \right]_+ \right\}$$



Beam functions  $\mathcal{B}_q(t, x, \vec{p}_\perp^2, \mu)$  match onto ordinary PDFs:

$$f_q \left( \frac{\vec{p}_\perp^2}{z}, \mu \right) = \theta(\vec{p}_\perp^2) \langle \mathcal{I}_n(\mathbb{R}^-) | \bar{\chi}_n(0) \frac{\vec{p}_\perp^2}{2} [\mathcal{S}(\vec{p}_\perp^2 - \vec{n} \cdot \mathbb{P}) \chi_n(0)] | \mathcal{I}_n(\mathbb{R}^-) \rangle$$

"   
 x

matching:  $\mathcal{B}_q(t, x, \vec{p}_\perp^2, \mu) = \sum_{j=q, g} \int_x^1 \frac{dz}{z} \mathcal{C}_{qj}(t, z, \vec{p}_\perp^2, \mu) f_j \left( \frac{x}{z}, \mu \right)$



$$\mathcal{C}_{qq}(t, z, \vec{p}_\perp^2, \mu) = \frac{1}{\pi} \delta(t) \delta(1-z) \delta(\vec{p}_\perp^2)$$

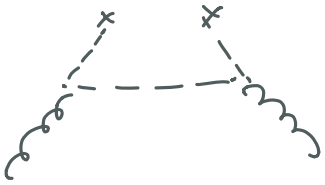
$$+ \frac{\alpha_s C_F}{2\pi^2} \theta(z) \left\{ \frac{2}{\mu^2} \left[ \frac{\theta(t) \ln(t/\mu^2)}{t/\mu^2} \right]_+ \delta(1-z) \delta(\vec{p}_\perp^2) \right.$$

$$\left. + \frac{1}{\mu^2} \left[ \frac{\theta(t)}{t/\mu^2} \right]_+ \left[ \mathcal{P}_{qq}(z) - \frac{3}{2} \delta(1-z) \right] \delta \left( \vec{p}_\perp^2 - \frac{(1-z)t}{z} \right) \right\}$$

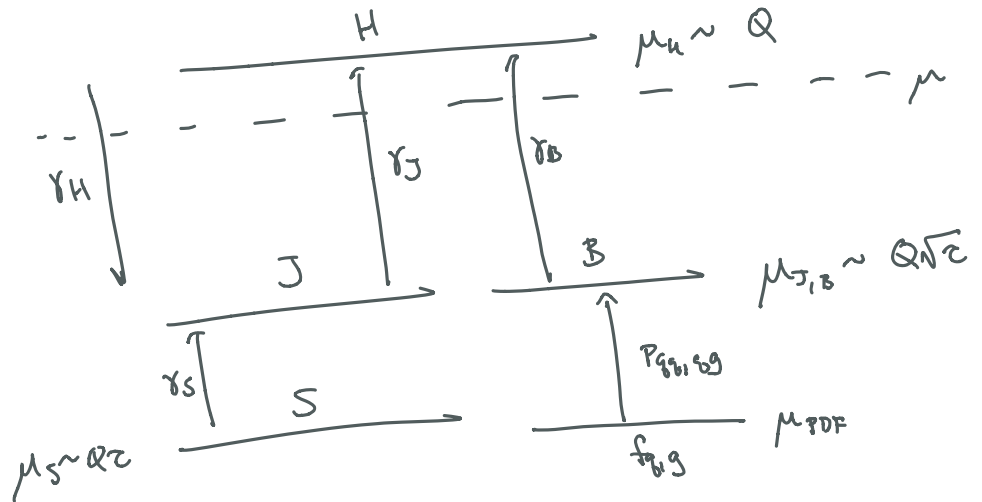


$$+ \delta(t) \delta(p_\perp^2) \left( \left[ \frac{\theta(1-z) \ln(1-z)}{1-z} \right]_* (1+z^2) - \frac{\pi^2}{6} \delta(1-z) + \theta(1-z) \left( 1-z - \frac{1+z^2}{1-z} \ln z \right) \right)$$

$$d\mathcal{G}_g(t, z, \vec{p}_\perp, r) = \frac{\alpha_s(\mu) T_F}{2\pi^2} \theta(z) \left\{ \frac{1}{\mu^2} \left[ \frac{\theta(t)}{t/r^2} \right]_* P_{gg}(z) \delta(p_\perp^2 - \frac{(1-z)}{z} t) + \delta(t) \delta(p_\perp^2) \left[ P_{gg}(z) \ln \frac{1-z}{z} + 2\theta(1-z) z(1-z) \right] \right\}$$



RG Evolution



Each function obeys an RGE eq:

$$\mu \frac{d}{d\mu} H(Q^2, r) = \gamma_H(r) H(Q^2, r)$$

$$\gamma_H(r) = 2\Gamma_{\text{cusp}}[\alpha_s(r)] \ln \frac{Q^2}{r^2} + \gamma_H[\alpha_s(r)]$$

[Laplace space  $\tilde{f}(v) = \int dt e^{vt} f(t)$ ]:

$$\mu \frac{d}{d\mu} \tilde{B}(v, r) = \gamma_B^{\sim}(v, r) \tilde{B}(v, r)$$

$$\gamma_B^{\sim}(r) = 2\Gamma_{\text{cusp}}[\alpha_s(r)] \ln(\mu^2 v e^{\epsilon}) + \gamma_B^{\sim}[\alpha_s(r)] = \gamma_B^{\sim}(r)$$

$$\mu \frac{d}{d\mu} \tilde{J}(v, r) = \gamma_J^{\sim}(v, r) \tilde{J}(v, r)$$

$$\gamma_J^{\sim}(r) = -4\Gamma_{\text{cusp}}[\alpha_s(r)] \ln(\mu v e^{\epsilon}) + \gamma_J^{\sim}[\alpha_s(r)]$$

$$\mu \frac{d}{d\mu} \tilde{S}(v, r) = \gamma_S^{\sim}(v, r) \tilde{S}(v, r)$$

$$\text{RG inv: } \gamma_H + \gamma_B + \gamma_J + \gamma_S = 0.$$

Solutions:

$$H(Q^2, \mu) = H(Q^2, \mu_H) \exp\left(\int_{\mu_H}^{\mu} \frac{d\mu'}{\mu'} \gamma_H(\mu')\right)$$

$$\tilde{F}(v, \mu) = \tilde{F}(v, \mu_F) \exp\left(\int_{\mu_F}^{\mu} \frac{d\mu'}{\mu'} \gamma_F(\mu')\right) \quad F = J, B, S$$

Achieve:

$$\tilde{\sigma}(v) \sim \left| 1 + \frac{d_s}{4\pi} (C_{12} \ln^2 v + C_{11} \ln v + C_{10}) + \left(\frac{d_s}{4\pi}\right)^2 (C_{24} \ln^4 v + C_{23} \ln^3 v + \dots + C_{20}) \right.$$

$$\sim \exp \left\{ \frac{d_s}{4\pi} (C_{12} \ln^2 v + C_{11} \ln v + C_{10}) + \left(\frac{d_s}{4\pi}\right)^2 (C_{23} \ln^3 v + C_{22} \ln^2 v + C_{21} \ln v + C_{20}) + \dots \right\}$$

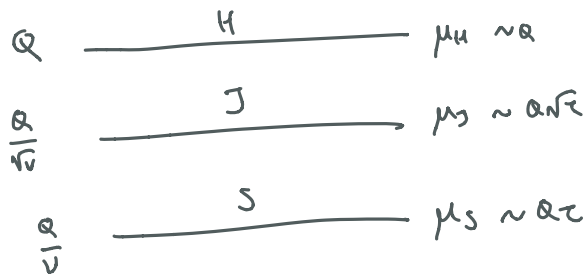
$$\ln v \sim \frac{1}{d_s}$$

$$\sim \frac{1}{d_s} \quad 1 \quad d_s \quad d_s^2 \dots$$

Compare most "traditional" resummations vs. SCET:

equiv of  $\mu_{H,B,S}$  usually pre-fixed

→ remain variable  
 ⇒ improved results in mom space & uncertainty estimates



hard to transfer  $d_s\left(\frac{Q}{v}\right)$

similar issue in  $q \leftrightarrow b$  space transfers for TMD cross sections