### Status of Muon g-2

Luchang Jin
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Sept 26, 2019

Brookhaven Forum 2019: Particle Physics and Cosmology in the 2020's Brookhaven National Lab

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     Taku Izubuchi (BNL/RBRC)
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        Chulwoo Jung (BNL)
   Andreas Juttner (Southampton)
       Christoph Lehner (BNL)
     Antonin Portelli (Edinburgh)
         Cheng Tu (UConn)
and the RBC/UKQCD collaborations
```

Thomas Blum (UConn/RBRC)
Peter Boyle (Edinburgh)

Outline 3 / 39

- Introduction
- Phenomenology determination
- Lattice calculation
- Perspect of near future lattice results

$$\begin{cases} q = p' - p, \nu \\ p \end{cases}$$

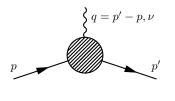
$$\vec{\mu} = g \frac{e}{2m} \vec{L}$$

Dirac equation implies:

$$\bar{u}(p')\gamma_{\nu}u(p)$$

$$g=2$$





$$\bar{u}(p')\left(F_1(q^2)\gamma_{\nu} + i\frac{F_2(q^2)[\gamma_{\nu},\gamma_{\rho}]q_{\rho}}{4m}\right)u(p)$$
(Euclidean space time)

q-2

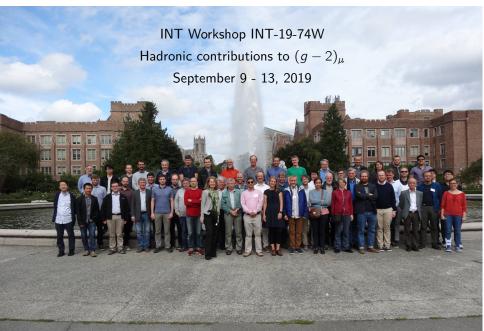
$$a = F_2(q^2 = 0) = \frac{g - 2}{2}$$

- The quantity a is called the anomalous magnetic moments.
- Its value comes from quantum correction.

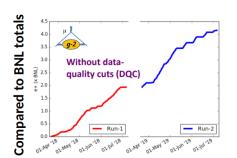
# Muon g-2: Fermilab E989, J-PARC E34 5/39

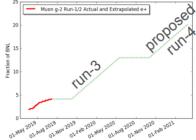
 $3.7\sigma$  deviations New Physics?





D. Hertzog's talk at Seattle INT workshop.





- Next Collaboration meeting is mid November
- MEANWHILE: RUN-3 Starts in October !!!



• J-PARC muon g-2 experiments in preparation.

	$a_{\mu}  imes 10^{10}$		
QED 5-loops	$11658471.8853 \pm 0.0036$	Aoyama, et al, 2012	
Weak 2-loops	$15.36 \pm 0.10$	Gnendiger et al, 2013	
HVP (LO)	$692.5 \pm 2.7$	RBC-UKQCD and FJ17 combined	
	$693.26 \pm 2.46$	KNT18	
	$693.9 \pm 4.0$	DHMZ19	
HVP (NLO)	$-9.93 \pm 0.07$	Fred Jegerlehner, 2017	
HVP (NNLO)	$1.22 \pm 0.01$	Fred Jegerlehner, 2017	
HLbL	$10.3 \pm 2.9$	Fred Jegerlehner, 2017	
	$10.5\pm2.6$	Glasgow Consensus, 2007	
SM Theory	$11659181.3 \pm 4.0$		
BNL E821 Exp	$11659208.9 \pm 6.3$		
Exp — SM	$27.6 \pm 7.5$		

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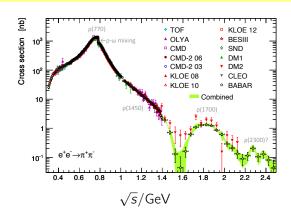
- Introduction
- Phenomenology determination
- Lattice calculation
- Perspect of near future lattice results

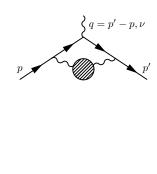
# HVP: dispersive approach

- C. Bouchiat and L. Michel, J. Phys. Radium22(1961) 121
- Z. Zhang's talk at Seattle INT workshop.

$$a_{\mu}^{\mathrm{HVP\ LO}} = rac{lpha^2}{3\pi^2} \int_{4m_{\pi}^2}^{\infty} ds rac{K(s)}{s} R(s), \qquad R(s) = rac{\sigma(e^+e^- o \mathrm{hadroncs})}{\sigma(e^+e^- o \mu^+\mu^-)}$$

#### The Dominant $\pi^+\pi^-$ Channel (1)



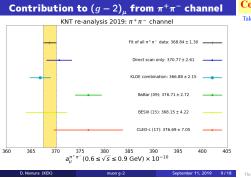


### HVP: KNT18 & DHMZ19

KNT18; D. Nomura's talk at Seattle INT workshop. DHMZ19; Z. Zhang's talk at Seattle INT workshop.

- KNT18:  $a_{\mu}^{\text{HVP LO}} = 693.26(2.46) \times 10^{-10}$ .
- DHMZ19:  $a_{\mu}^{\text{HVP LO}} = 693.9(4.0) \times 10^{-10}$ .

#### BaBar & KLOE tension



#### Combined Results Fit [<0.6 GeV] + Data [0.6-1.8 GeV]

Take into account the correlation of 62% (based on pseudo-data samples) of the two regions



⇒ The difference "All but BABAR" and "All but KLOE" = 5.6 to be compared with 1.9 uncertainty with "All data"

KLOE

- ➤ The local error inflation is not sufficient to amplify the uncertainty
- ➤ Global tension (normalisation/shape) not previously accounted for
- ➤ Potential underestimated uncertainty in at least one of the measurements?
- ➤ Other measurements not precise enough and are in agreement with BABAR or
- ⇒ Given the fact we do not know which dataset is problematic, we decide to
  - ➤ Add half of the discrepancy (2.8) as an additional uncertainty (correcting the local PDG inflation to avoid double counting)
- ➤ Take the mean value "All but BABAR" and "All but KLOE" as our central value

The muon o.2 theory initiatives. Sentile. Sent 9.13. 2019. Zhiqing Zhang (LAL, Orsay)

Future experiments at Belle II, CMD-3, SND can provide more insight.

### **HVP**: Seattle consensus

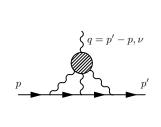
Bogdan Malaescu's talk at Seattle INT workshop.

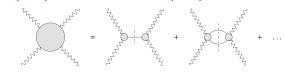
### Proposal: conservative merging of model-independent HVP combination results

- Basic requirements for the *merging* procedure:
- *Conservative* (see tensions between experimental data and differences between combinations based on same datasets )
- Accounting for correlations between different channels (understood meaning of systematic uncertainties and identified 15 common ones, DHMZ since arXiv:1010.4180) Yields unavoidable increase of total uncertainty  $693.9 \pm 1.0 \pm 3.4 \pm \boxed{1.6} \pm 0.1_{\psi} \pm 0.7_{\rm QCD}$
- Proposed merging procedure:
- <u>Central value</u>: simple average of the <u>DHMZ</u> and <u>KNT</u> sums of channels (the <u>DHMZ</u> and <u>KNT</u> central values are, *by chance*, very similar)
- Experimental uncertainties: in each channel/mass range use max(DHMZ, KNT) and see by how much to increase the corresponding DHMZ uncertainty (sq. difference); enhance the DHMZ sum of channels (with correlations) by these amounts (sq. sum)
- Use  $|\underline{DHMZ}(\text{ch.})-\underline{KNT}(\text{ch.})|/2$  as extra systematic in each channel; independent between channels (sign of algebraic difference fluctuates for various channels) o)  $\pi\pi$  BABAR/KLOE systematic: max(DHMZ B./K. syst.,  $|\underline{DHMZ}(\pi\pi)-\underline{KNT}(\pi\pi)|/2$ ) (stay conservative, but avoid double-counting the effect of this B./K. tension)
- o)  $\pi^+\pi^-\pi^0$ : do not include this systematic (difference understood:  $1^{st}/2^{nd}$  order interp.)

### HLbL: dispersive approach

G. Colangelo et al., Phys.Rev.Lett. 118 (2017) no.23, 232001M. Hoferichter et al., JHEP 1810 (2018) 141; Mainz PRD 100 (2019) no.3, 034520





 $\label{eq:Figure 2.} \textbf{ Intermediate states in the direct channel: pion pole and two-pion cut.}$ 

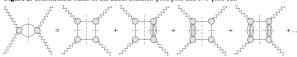


Figure 3. Two-pion contributions to HLbL. Further crossed diagrams are not shown explicitly.

- Pion pole + pion box + rescattering + · · ·
   Estimate/bound short distance contribution.
- "Not in a form as compact as for HVP".
   Need many inputs from experiments for lattice.

## HLbL: model & dispersive

### G. Colangelo's talk at Seattle INT workshop.

Contribution	PdRV(09)	N/JN(09)	J(17)	White Paper
$\pi^0, \eta, \eta'$ -poles	$114\pm13$	$99\pm16$	$95.45 \pm 12.40$	$93.8 \pm 4.0$
$\pi$ , K-loop/box	$-19\pm19$	$-19\pm13$	$-20\pm5$	$-16.4 \pm 0.2$
S-wave $\pi\pi$	_	_	_	$-8 \pm 1$
scalars	$-7\pm7$	$-7\pm2$	$-5.98 \pm 1.20$	$\left.\right\}$ - 2 ± 3
tensors	_	_	$1.1 \pm 0.1$	] -2 ± 3
axials	$15\pm10$	$22\pm 5$	$7.55 \pm 2.71$	8 ± 8
q-loops / SD	2.3	$21\pm3$	$22.3 \pm 5.0$	$10\pm10$
total	$105\pm26$	$116\pm39$	$100.4 \pm 28.2$	85 ± <i>XX</i>

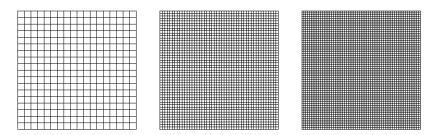
- In unit of 10<sup>-11</sup> for this slides.
- PdRV = Prades, de Rafael, Vainshtein ("Glasgow consensus").
- N = Nyffeler, J = Jegerlehner.
- Uncertainties added in quadrature: XX=12.
- Uncertainties added linearly: XX=21.

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The QCD partition function in Euclidean space time:

$$Z = \int [\mathcal{D}U_{\mu}]e^{-S_{\mathcal{O}}[U]}\det(D[m_l, U])^2 \det(D[m_s, U])$$
(1)



(Left)  $19 \times 19$  Go board (Middle)  $48 \times 48$  (Right)  $64 \times 64$ 

The configuration is stored in position space. The reason is that the action is local in position space. Working in position makes the calculation simpler.

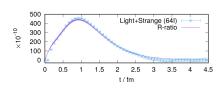
This is in contrast to analytical perturbative calculation, where interaction only happens occasionally. So it is advantagous to work in momentum space, where the propagator can be diagonalized.

# HVP: lattice approach

T. Blum 2003; D. Bernecker, H. Meyer 2011.

$$C(t) = \frac{1}{3} \sum_{\vec{x}} \sum_{j=0,1,2} \langle J_j(\vec{x}, t) J_j(0) \rangle$$

$$a_{\mu}^{\mathsf{HVP LO}} = \sum_{t=0}^{+\infty} w(t) C(t)$$



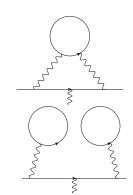
 In Euclidean space-time, C(t) decreases exponentially as t increases.

For 
$$t\sim 1$$
 fm,  $C(t)\sim e^{-m_{
ho}t}$ .

For 
$$t \to \infty$$
,  $C(t) \sim e^{-2m_{\pi}t}$ .

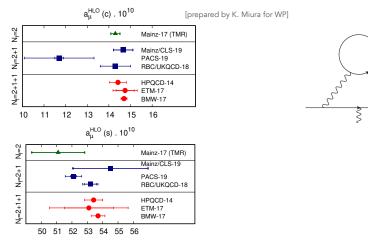
Lattice statistical error:  $\delta C(t) \sim e^{-m_\pi t}$ .

• For  $t\lesssim 1$  fm,  $w(t)\sim t^4$ . For  $t\to\infty$   $(m_\mu t\gg 1)$ ,  $w(t)\sim t^2$ .



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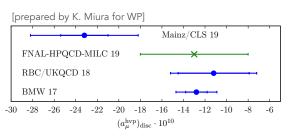
### A. El-Khadra's talk at Seattle INT workshop.



Charm and Strange quark connected is very accurate.

# HVP: disconnected diagram

#### A. El-Khadra's talk at Seattle INT workshop.



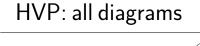


- Include *u*, *d*, *s* quark contribution.
- Strong suppression from charge factor.

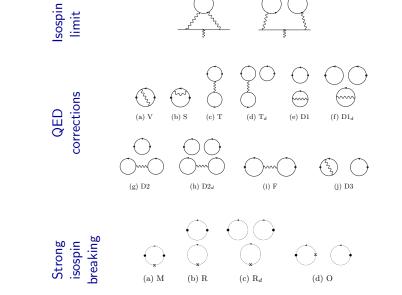
- con: 
$$e_u^2 + e_d^2 = 5/9$$

- discon: 
$$(e_u + e_d)^2 = 1/9$$

• Vanishes in the flavor SU(3) limit  $=> \sim (m_{u,d}-m_s)^2$ .

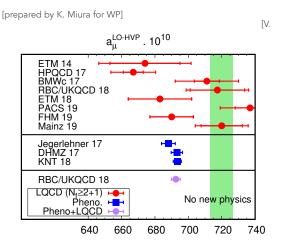


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### HVP: pure lattice

### A. El-Khadra's talk at Seattle INT workshop.



# HVP: hybrid (lattice & dispersive)

PHYSICAL REVIEW LETTERS 121, 022003 (2018)

Editors' Suggestion

#### Calculation of the Hadronic Vacuum Polarization Contribution to the Muon Anomalous Magnetic Moment

T. Blum, P. A. Boyle, V. Gülpers, T. Izubuchi, J. L. Jin, J. C. Jung, A. Jüttner, C. Lehner, A. Portelli, and J. T. Tsang

(RBC and UKQCD Collaborations)

<sup>1</sup>Physics Department, University of Connecticut, Storrs, Connecticut 06269-3046, USA <sup>2</sup>School of Physics and Astronomy, The University of Edinburgh, Edinburgh EH9 3FD, United Kingdom <sup>3</sup>School of Physics and Astronomy, University of Southampton, Southampton SO17 1BJ, United Kingdom <sup>4</sup>Physics Department, Brookhaven National Laboratory, Upton, New York 11973, USA <sup>5</sup>RIKEN-BNL Research Center, Brookhaven National Laboratory, Upton, New York 11973, USA



(Received 25 January 2018; published 12 July 2018)

We present a first-principles lattice QCD + QED calculation at physical pion mass of the leading-order hadronic vacuum polarization contribution to the muon anomalous magnetic moment. The total contribution of up, down, strange, and charm quarks including QED and strong isospin breaking effects is  $a_{\mu}^{\text{HVP LO}} = 715.4(18.7) \times 10^{-10}$ . By supplementing lattice data for very short and long distances with R-ratio data, we significantly improve the precision to  $a_{\mu}^{\text{HVP LO}} = 692.5(2.7) \times 10^{-10}$ . This is the currently most precise determination of a HVP LO.

Pure lattice result and dispersive result with reduced  $\pi\pi$  dependence (window method)

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RBC-UKQCD PRL 121, 022003 (2018)

$$a_{\mu}^{\text{HVP LO}} = \sum_{t=0}^{+\infty} w(t)C(t) \qquad w(t) = w^{\text{SD}}(t) + w^{\text{W}}(t) + w^{\text{LD}}(t)$$

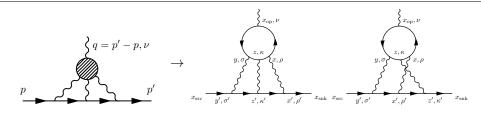
Most of  $\pi\pi$  peak is captured by window from  $t_0 = 0.4$  fm to  $t_1 = 1.5$  fm, so replacing this region with lattice data reduces the dependence on BaBar versus KLOE data sets.

Babar & KLOE tension => syst error  $5.6/2 = 2.8 \times 10^{-10}$ .

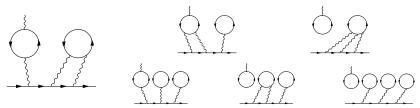
The window method is much less affected by this tension.

t / fm

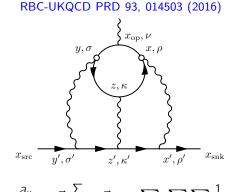
# HLbL: diagrams



- Gluons and sea quark loops (not directly connected to photons) are included automatically to all orders!
- There are additional four different permutations of photons not shown.



# HLbL: Exact photon and moment



 $ec{\mu} = \sum rac{1}{2} (ec{x}_{\mathsf{op}} - ec{x}_{\mathsf{ref}}) imes ec{J}(ec{x}_{\mathsf{op}})$ 

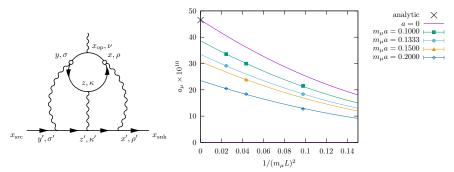
- Two point sources at x, y: randomly sample x and y.
- Importance sampling:
   focus on small |x y|.
- Complete sampling for  $|x y| \le 5a$  upto discrete symmetry.

$$\frac{a_{\mu}}{m_{\mu}}\bar{u}_{s'}(\vec{0})\frac{\sum}{2}u_{s}(\vec{0}) = \sum_{r=x-y}\sum_{z}\sum_{x_{op}}\frac{1}{2}(\vec{x}_{op} - \vec{x}_{ref}) \times \bar{u}_{s'}(\vec{0})i\vec{\mathcal{F}}^{C}(\vec{0}; x, y, z, x_{op})u_{s}(\vec{0})$$

- ,
- Only sum over r = x y.

### HLbL: muon leptonic LbL test

### RBC-UKQCD PRD 93, 014503 (2016)

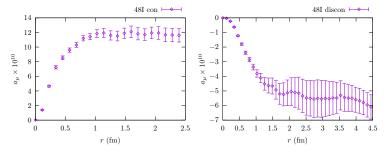


$$F_2(a,L) = F_2 \left( 1 - \frac{c_1}{(m_\mu L)^2} + \frac{c_1'}{(m_\mu L)^4} \right) (1 - c_2 \, a^2 + c_2' \, a^4) \ \rightarrow \ F_2 = 46.6(2) \times 10^{-10} \ (19)$$

- Pure QED computation. Muon leptonic light by light contribution to muon g-2. Phys.Rev. D93 (2016) 1, 014503. arXiv:1510.07100.
- Analytic results:  $0.371 \times (\alpha/\pi)^3 = 46.5 \times 10^{-10}$ .
- ullet  $\mathcal{O}(1/L^2)$  finite volume effect, because the photons are emitted from a conserved loop.

## HLbL: 48I $m_{\pi}=139$ MeV

RBC-UKQCD PRL 118, 022005 (2017)



- Left: connected diagrams. Right: leading disconnected diagrams.
- $48^3 \times 96$  lattice, with  $a^{-1} = 1.73$  GeV,  $m_{\pi} = 139$  MeV,  $m_{\mu} = 106$  MeV.
- We use Lanczos, AMA, and zMobius techniques to speed up the computations.
- 65 configurations are used. They each are separated by 20 MD time units.

$$a_{\mu}^{\text{cHLbL}} = (11.60 \pm 0.96) \times 10^{-10}$$
 (30)  
 $a_{\mu}^{\text{dHLbL}} = (-6.25 \pm 0.80) \times 10^{-10}$  (31)

$$a_{\mu}^{\text{CHLbL}} = (5.35 \pm 1.35) \times 10^{-10}$$
 (3)

(32)

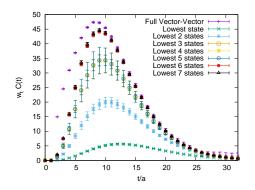
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C. Lehner's talk at Seattle INT workshop. RBC-UKQCD prelim.

The Euclidean space-time correlator in finite volume

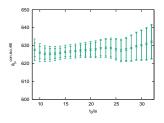
$$C(t) = \sum_{n} |\langle 0|V|n\rangle|^2 e^{-E_n t}$$



48I, ud,conn,iso:  $626.5(2.6)_{\text{stat}} \times 10^{-10}$ 

Previous stat error:  $14.2 \times 10^{-10}$ 

- At large t large statistical error dominated by a few states.
- SD: Full Vector-VectorLD: Sum of states
- Systematic can be strictly bounded.



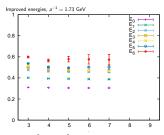
## HVP: finding the states

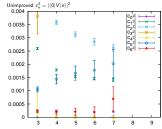
C. Lehner's talk at Seattle INT workshop. RBC-UKQCD prelim.

The Euclidean space-time correlator in finite volume

$$C(t) = \sum_{n} |\langle 0|V|n\rangle|^2 e^{-E_n t}$$

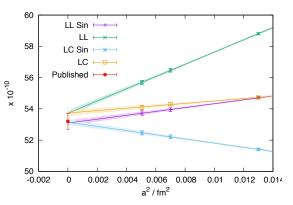
- Study the N × N correction matrix formed by N operators with the same quantum number:
  - local and smeared vector currents
  - 2 pion operator with 1,2,3,4 units of momentum
  - 4 pion operator
- Both  $E_n$  and  $\langle 0|V|n\rangle$  can be extracted.





C. Lehner's talk at Seattle INT workshop. RBC-UKQCD prelim.

Third lattice spacing for strange data ( $a^{-1} = 2.77$  GeV with  $m_{\pi} = 234$  MeV with sea light-quark mass corrected from global fit):



For light quark need new ensemble at physical pion mass. Started run on Summit Machine at Oak Ridge this year ( $a^{-1} = 2.77$  GeV with  $m_{\pi} = 139$  MeV).

#### RBC-UKQCD prelim

- MDWF+Iwasaki: continuum limit (5.4fm)
- MDWF+DSDR:  $a^{-1} = 1.015 \,\text{GeV}$ :  $24^3 \times 64$  (4.8 fm),  $32^3 \times 64$  (6.4 fm),  $48^3 \times 64$  (9.6 fm).
- MDWF+DSDR:  $a^{-1} = 1.371 \,\text{GeV}$ :  $32^3 \times 64 \,\text{(4.6 fm)}$ .

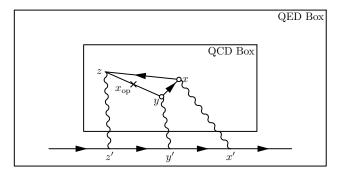
$$F_{2}(a,L) = F_{2}\left(1 - \frac{c_{1}}{(m_{\mu}L)^{2}}\right)(1 - c_{2}a^{2}) \tag{43}$$

### Disconnected diagrams

$$a_{\mu}^{\text{cHLbL}} = (27.61 \pm 3.51_{\text{stat}} \pm 0.32_{\text{sys},a^2}) \times 10^{-10}$$
 (44)  
 $a_{\mu}^{\text{dHLbL}} = (-20.20 \pm 5.65_{\text{stat}}) \times 10^{-10}$  (45)  
 $a_{\mu}^{\text{HLbL}} = (7.41 \pm 6.32_{\text{stat}} \pm 0.32_{\text{sys},a^2}) \times 10^{-10}$  (46)

# HLbL: $QED_{\infty}$ kernel & subtraction

Mainz PoS LATTICE2016 164, RBC-UKQCD PRD 96, 034515 (2017)



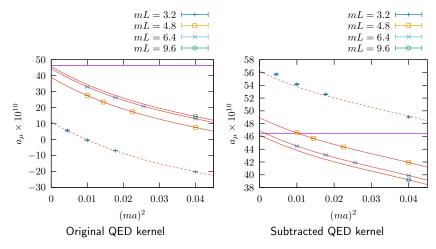
Subtracted QED kernel (use the current conservation condition) reduce discretization & finite volume & statistical error

$$\mathfrak{G}_{\rho,\sigma,\kappa}^{(2)}(x,y,z) = \mathfrak{G}_{\rho,\sigma,\kappa}^{(1)}(x,y,z) - \mathfrak{G}_{\rho,\sigma,\kappa}^{(1)}(y,y,z) - \mathfrak{G}_{\rho,\sigma,\kappa}^{(1)}(x,y,y) + \mathfrak{G}_{\rho,\sigma,\kappa}^{(1)}(y,y,y)$$

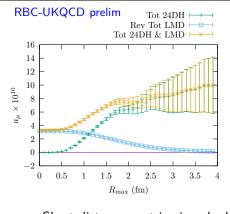
# HLbL: $QED_{\infty}$ muon leptonic LbL test

#### RBC-UKQCD PRD 96, 034515 (2017)

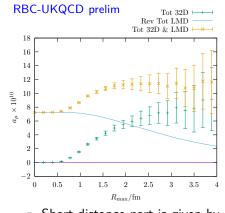
• Compare the two  $\mathfrak{G}_{\rho,\sigma,\kappa}(x,y,z)$  in pure QED computation.



Notice the vertical scales in the two plots are different.



- $R_{\text{max}} = \max(|x-y|, |x-z|, |y-z|).$
- $a^{-1} = 1.015$  GeV.
- 24DH: partial sum upto  $R_{\text{max}}$ .
- Rev LMD: reverse partial sum down to R<sub>max</sub>.
- 24DH & LMD: the sum of the above two curves.
- Short distance part is given by lattice data.
- Long distance part is given by LMD model of  $\pi^0$  pole.
- At 2.0 fm, the combination gives:  $a_{\mu}^{\text{tot}} = 7.46(62)_{\text{stat}} \times 10^{-10}$ .  $a_{\mu}^{\text{con}} = 13.44(10)_{\text{stat}} \times 10^{-10}$   $a_{\mu}^{\text{discon}} = -5.70(58)_{\text{stat}} \times 10^{-10}$



- $R_{\text{max}} = \max(|x-y|, |x-z|, |y-z|).$
- $a^{-1} = 1.015$  GeV.
- 32D: partial sum upto  $R_{\text{max}}$ .
- Rev LMD: reverse partial sum down to R<sub>max</sub>.
- 32D& LMD: the sum of the above two curves.
- Short distance part is given by lattice data.
- Long distance part is given by LMD model of  $\pi^0$  pole.
- At 2.5 fm, the combination gives:  $a_{\mu}^{\rm tot} = 11.40(1.27)_{\rm stat} \times 10^{-10}$ .  $a_{\mu}^{\rm con} = 29.19(0.73)_{\rm stat} \times 10^{-10}$   $a_{\mu}^{\rm discon} = -17.79(1.13)_{\rm stat} \times 10^{-10}$

RBC-UKQCD prelim (QED $_{\infty}$  & LMD)

Accumulate stat & Estimate syst

HVP  $692.5 \pm 2.7$ RBC-UKQCD and FJ17 combined (LO)  $693.26 \pm 2.46$ KNT18  $693.9 \pm 4.0$ DHM719  $??? \pm 5.0$ Pure lattice results  $??? \pm 2.0$ Hybrid approach (KNT & RBC-UKQCD) Weigh in on BaBar/KLOE HI bl  $10.3 \pm 2.9$ Fred Jegerlehner, 2017  $10.5 \pm 2.6$ Glasgow Consensus, 2007  $8.5 \pm 2.1$ Dispersive Seattle INT 2019

 $\pm 1.2$  if add error in quadrature. RBC-UKQCD prelim (QED, )  $7.41 \pm 6.32_{\rm stat} \pm 0.32_{\rm sys.a^2}$ 

 $11.40 \pm 1.27_{\rm stat} \pm ???_{\rm sys}$ 

# Timeline for the White Paper

Earliest possible release date for Fermilab g-2 measurement:

#### 15-20 December 2019

Post the WP on arXiv by:

1 Dec. 2019

Deadline for finalizing individual WP chapters:

1 Nov 2019

At this date the Overleaf chapters will be frozen.

• Editorial board will release complete WP to authors for feedback on:

15 Nov. 2019

will need to receive feedback from authors within a week

Experimental and theoretical inputs used in WP must be published by:

15 Oct 2019

To make sure to be included in WP discussion, a paper to be posted in arXiv by same date.

Note: The WP will be posted on arXiv in December, even if the Fermilab experiment's release date is delayed.

Seattle INT workshop summary

# Thank You!