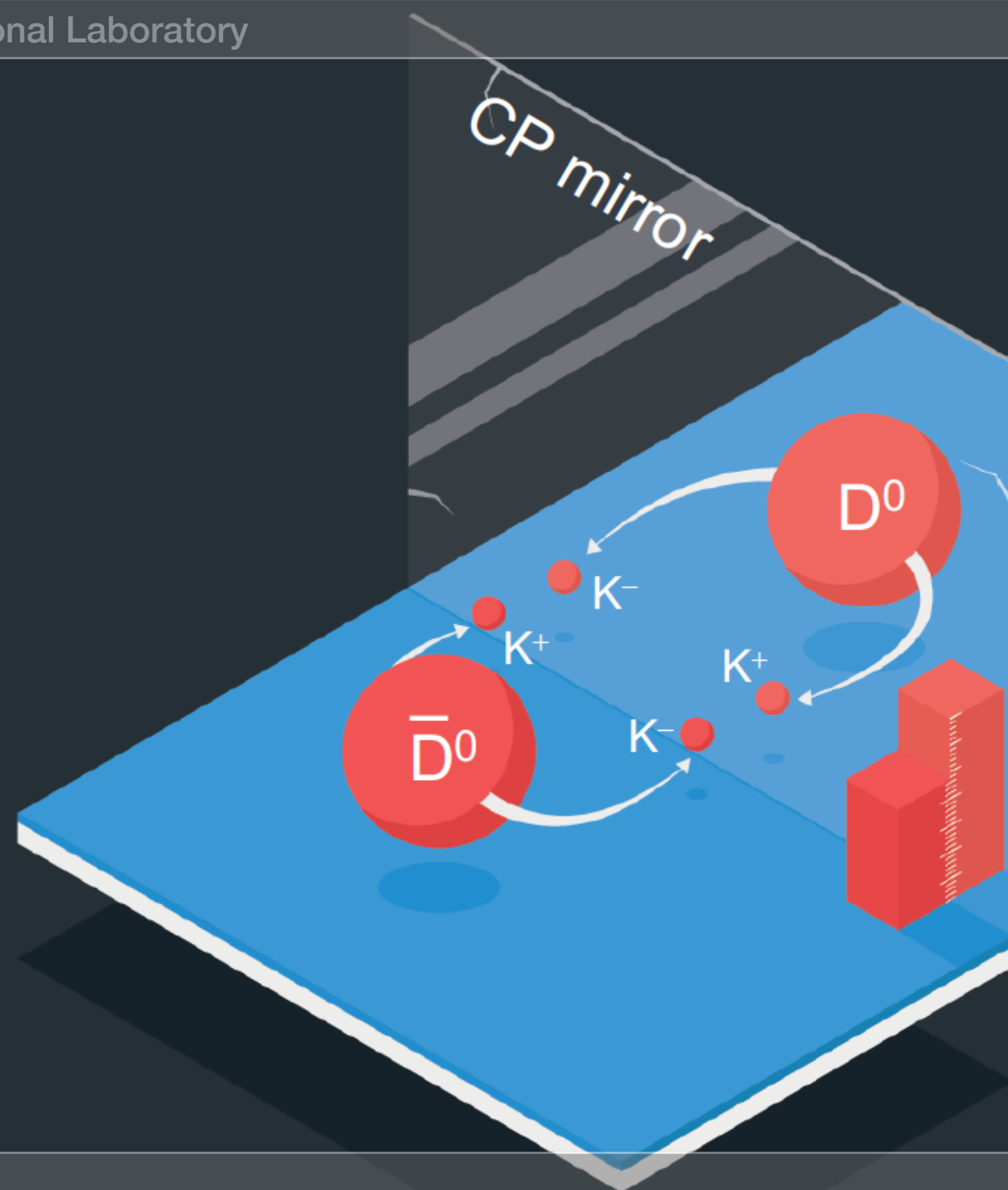




May 3rd, 2019

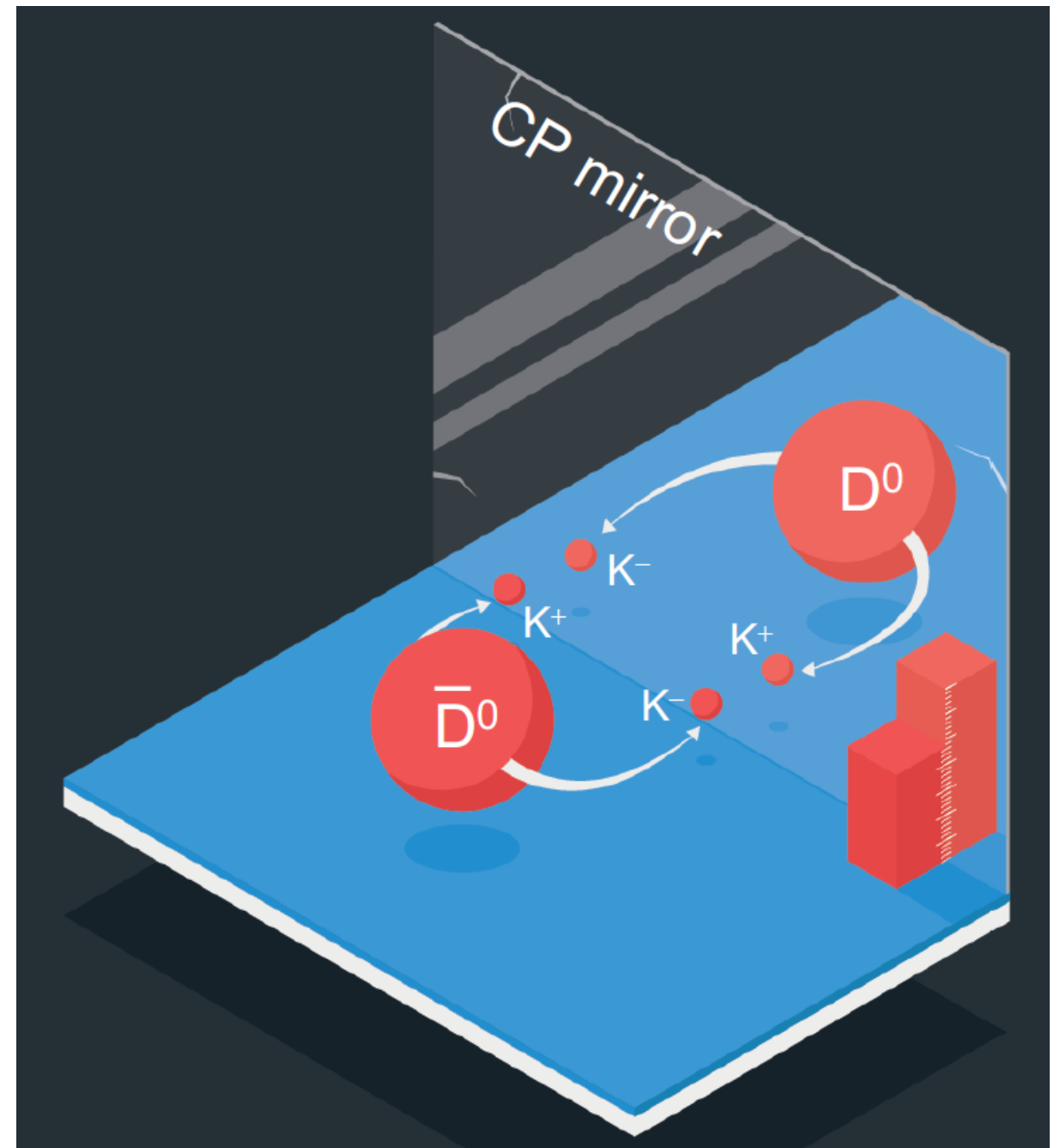
Observation of CP violation in charm decays

Angelo Di Canto

on behalf of the LHCb collaboration

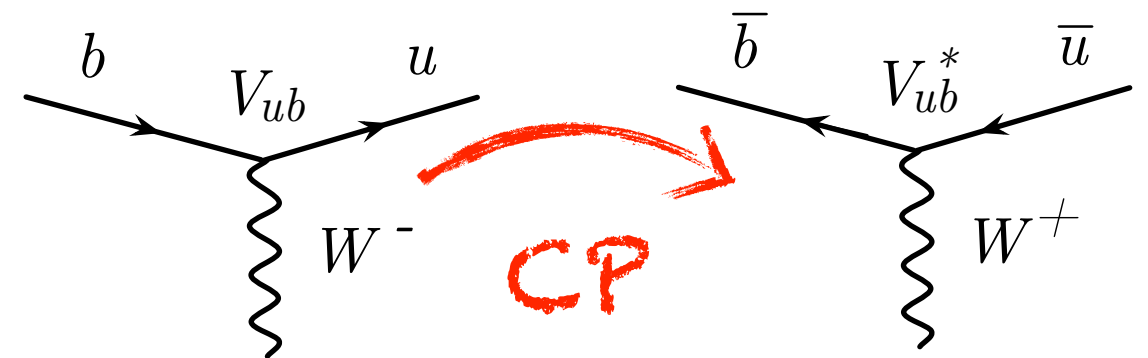
What is CP violation?

- Combination of mirror-reversal of the spatial arrangement (P) and replacement of particles with antiparticles (C)
- Dynamics of quarks (and maybe also leptons) is not invariant under CP transformation



From where does it come?

- Irreducible phase of the Cabibbo-Kobayashi-Maskawa matrix describing the charged currents of quarks
- Potentially large effects in transitions involving the third generation
- Highly suppressed for charm, which has small couplings to the third generation



$$V \approx \begin{pmatrix} \boxed{d} & \boxed{s} & \boxed{b} \\ \begin{pmatrix} 1 & \lambda & \lambda^3 e^{i\varphi} \\ -\lambda & 1 & \lambda^2 \\ -\lambda^3 e^{-i\varphi} & -\lambda^2 & 1 \end{pmatrix} & \begin{matrix} \boxed{u} \\ \boxed{c} \\ \boxed{t} \end{matrix} \end{pmatrix}$$

$$\lambda \approx 0.22$$

Brief history of CP violation

P violation in weak interactions postulated by Lee and Yang

CP violation in strange mesons by Cronin, Fitch et al

CP violation in beauty mesons by BaBar and Belle

1956 1957 ... 1963 1964 ... 1973 2001 ... 2019

P and C violation by Wu *et al* and Lederman *et al*

Cabibbo mixing angle

Kobayashi and Maskawa propose third generation to explain CP violation

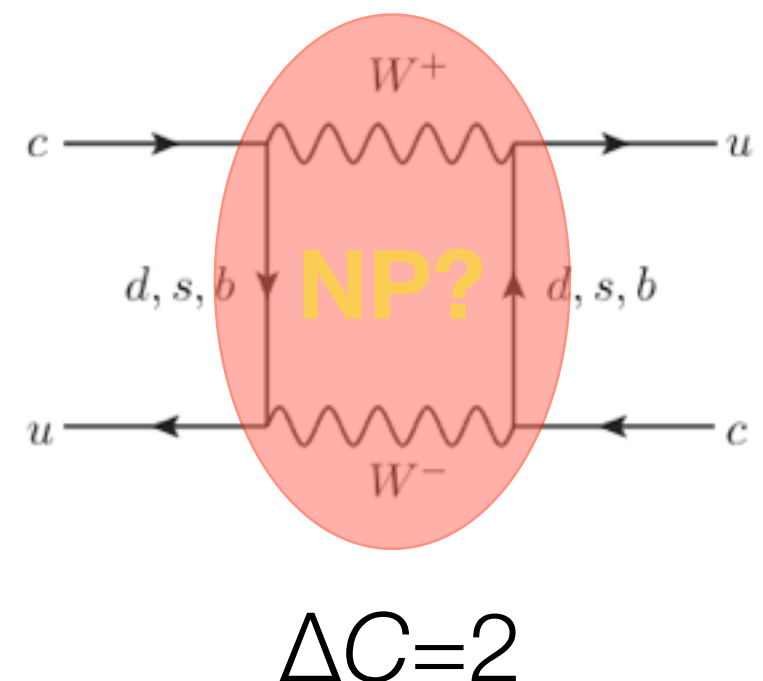
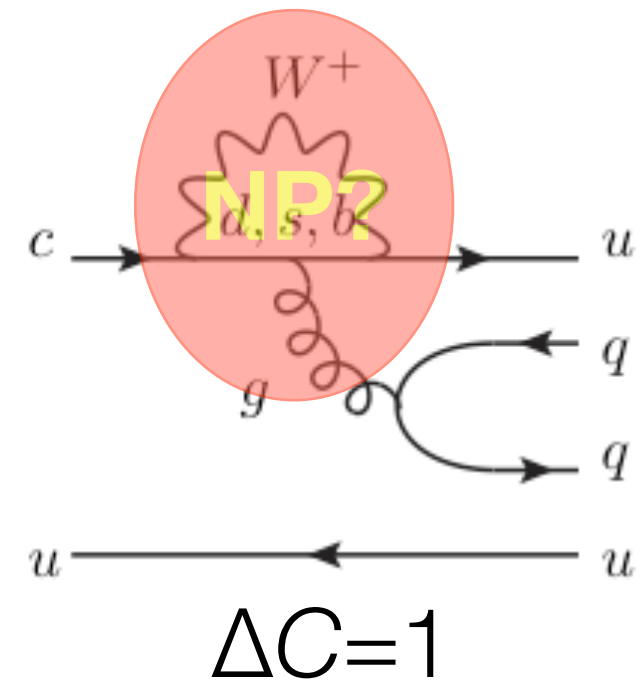
CP violation in charm mesons by LHCb

Why shall we bother about charm?

Discovery tool, standard model effects are $\sim 10^{-3}$ or smaller, potential room for new physics to show up

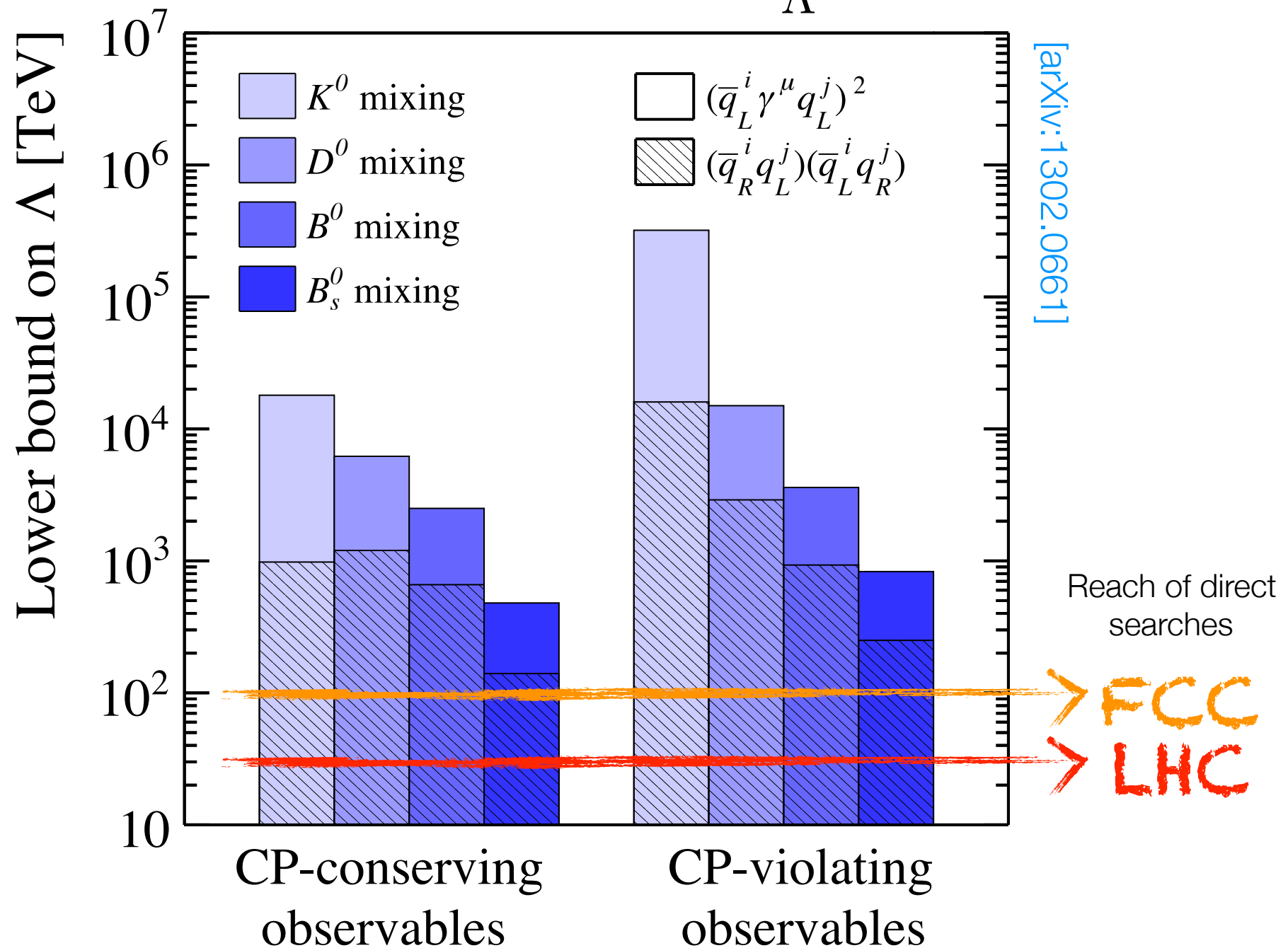
Unique, gives sensitivity to new physics coupling to up-type quarks (complementary to K and $B_{(s)}$ decays)

Challenging, need $O(1-100M)$ yields and control over systematic uncertainties. Predictions are difficult, not a precision probe

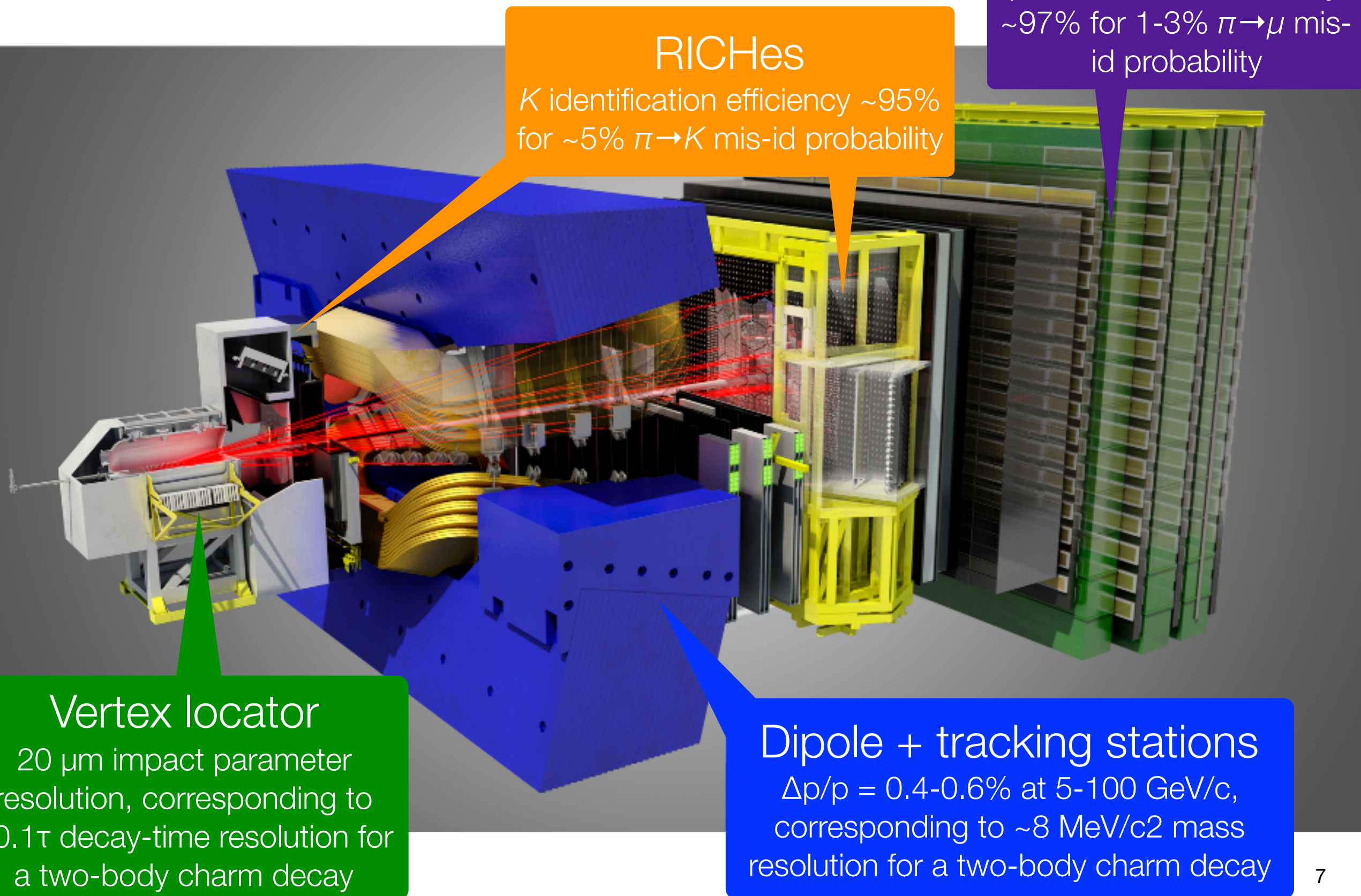


Reach/impact

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \frac{1}{\Lambda^2} \mathcal{O}_{\Delta F=2}$$

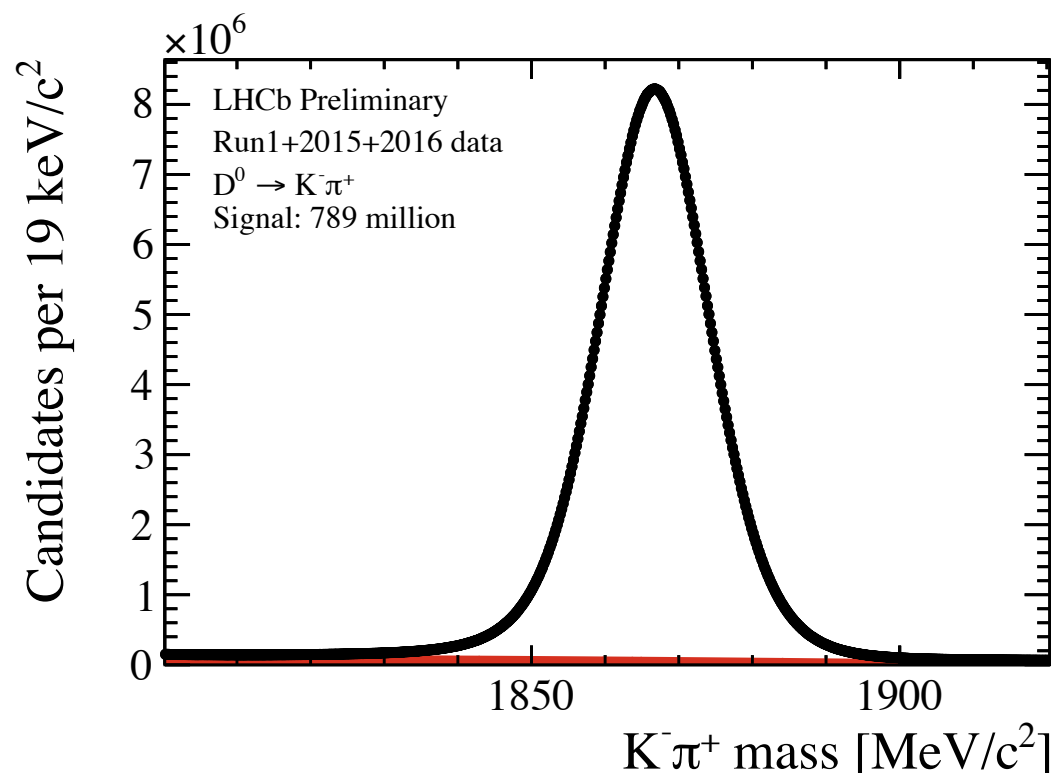


The beauty of the LHC



The trigger

- Real-time calibration of the detector allows offline-quality reconstruction
- Trigger candidates are used directly in the analysis



LHCb 2015 Trigger Diagram

40 MHz bunch crossing rate

L0 Hardware Trigger : 1 MHz readout, high E_T/P_T signatures

**450 kHz
 $h^\pm$**

**400 kHz
 $\mu/\mu\mu$**

**150 kHz
 e/γ**

Software High Level Trigger

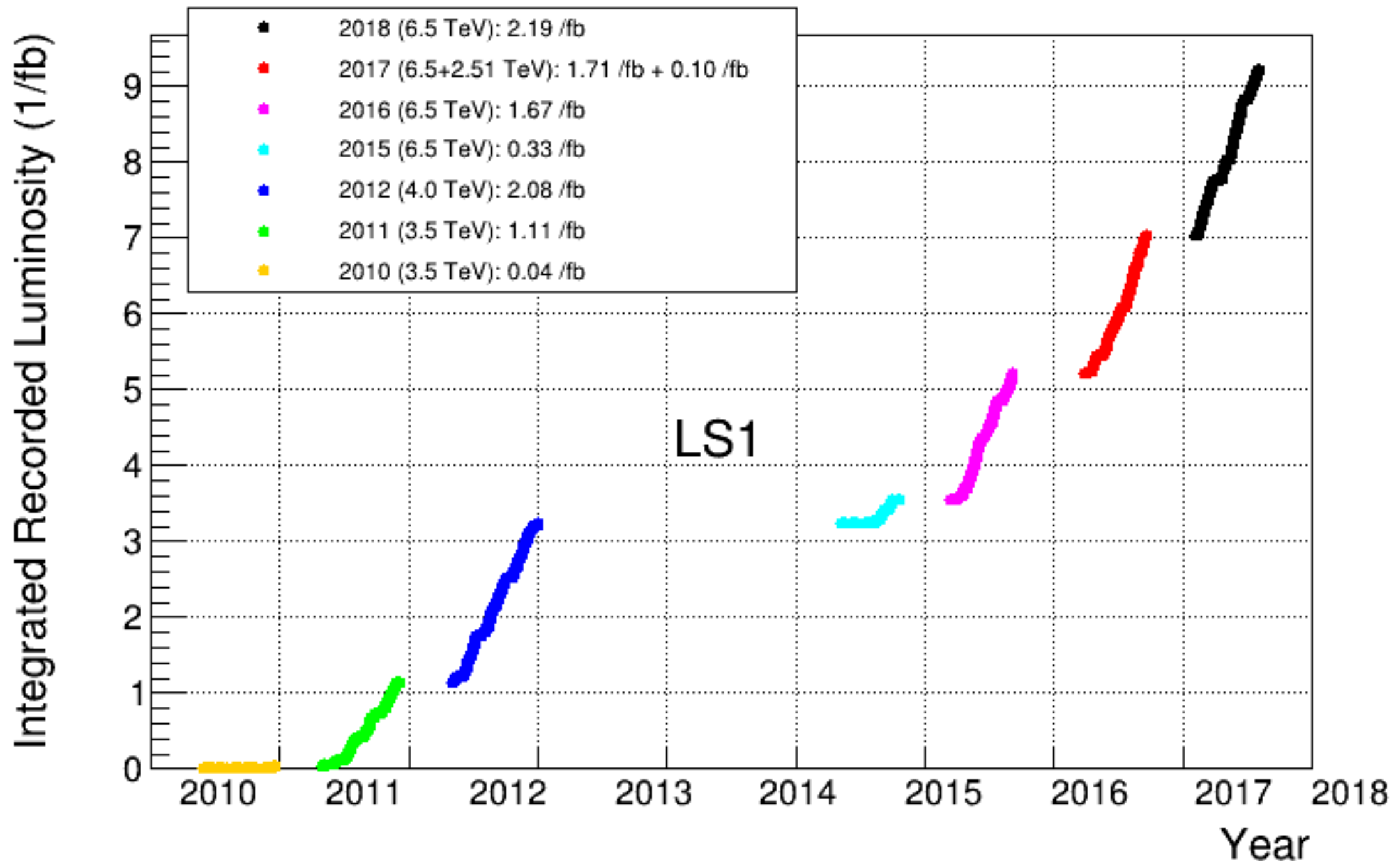
Partial event reconstruction, select displaced tracks/vertices and dimuons

Buffer events to disk, perform online detector calibration and alignment

Full offline-like event selection, mixture of inclusive and exclusive triggers

12.5 kHz (0.6 GB/s) to storage

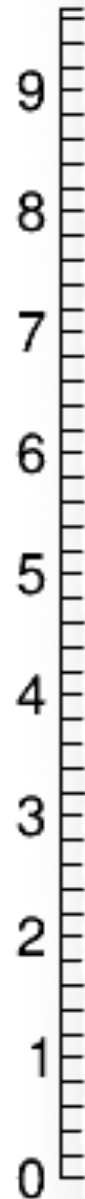
LHCb Cumulative Integrated Recorded Luminosity in pp, 2010-2018





CERN-EP-2019-042
LHCb-PAPER-2019-006
March 21, 2019

Integrated Recorded Luminosity (1/fb)



[arXiv:1903.08726]

Observation of CP violation in charm decays

LHCb collaboration[†]

Abstract

A search for charge-parity (CP) violation in $D^0 \rightarrow K^- K^+$ and $D^0 \rightarrow \pi^- \pi^+$ decays is reported, using pp collision data corresponding to an integrated luminosity of 6 fb^{-1} collected at a center-of-mass energy of 13 TeV with the LHCb detector. The flavor of the charm meson is inferred from the charge of the pion in $D^{*}(2010)^+ \rightarrow D^0 \pi^+$ decays or from the charge of the muon in $\bar{B} \rightarrow D^0 \mu^- \bar{\nu}_\mu X$ decays. The difference between the CP asymmetries in $D^0 \rightarrow K^- K^+$ and $D^0 \rightarrow \pi^- \pi^+$ decays is measured to be $\Delta A_{CP} = [-18.2 \pm 3.2 (\text{stat.}) \pm 0.9 (\text{syst.})] \times 10^{-4}$ for π -tagged



First results with full statistics!

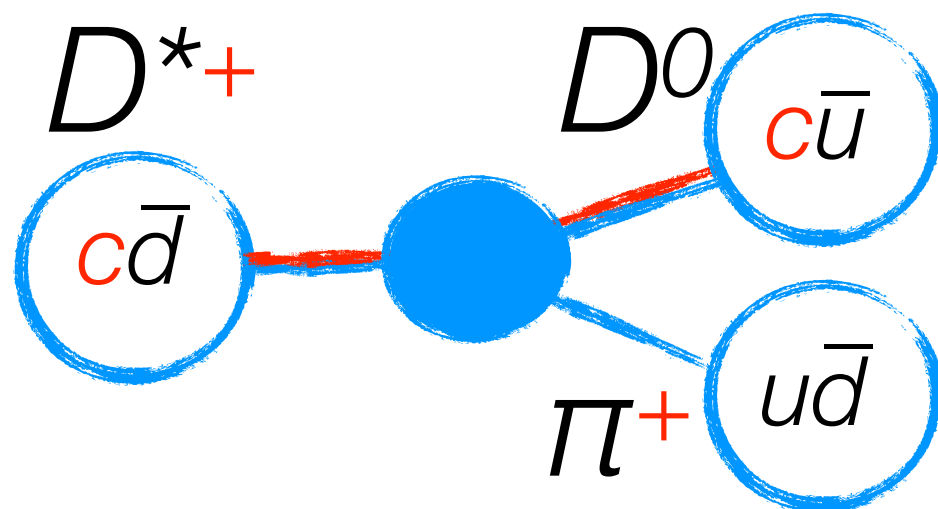
What do we measure?

$$A_{CP}(D^0 \rightarrow h^+ h^-) = \frac{\Gamma(D^0 \rightarrow h^+ h^-) - \Gamma(\bar{D}^0 \rightarrow h^+ h^-)}{\Gamma(D^0 \rightarrow h^+ h^-) + \Gamma(\bar{D}^0 \rightarrow h^+ h^-)}$$

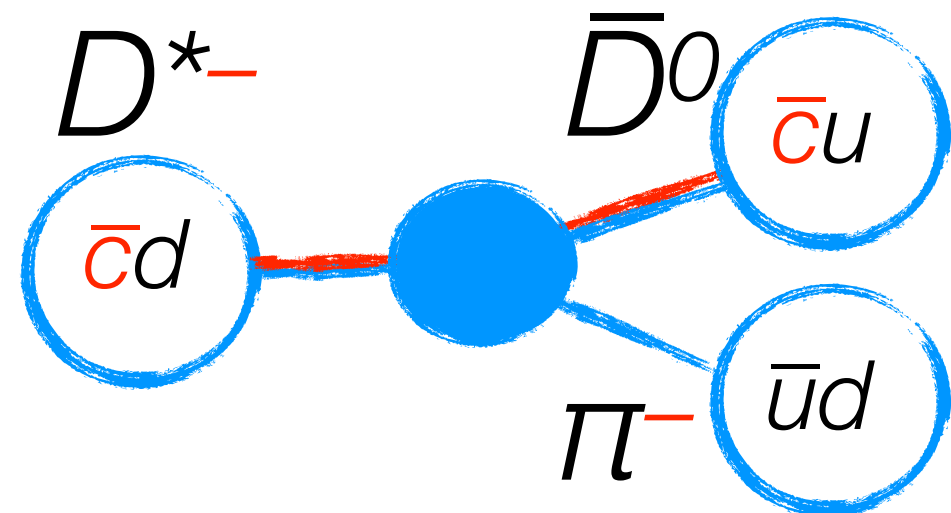
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$$A_{CP}(D^0 \rightarrow h^+ h^-) = \frac{\Gamma(D^0 \rightarrow h^+ h^-) - \Gamma(\bar{D}^0 \rightarrow h^+ h^-)}{\Gamma(D^0 \rightarrow h^+ h^-) + \Gamma(\bar{D}^0 \rightarrow h^+ h^-)}$$

How do we tell if it's a D^0 or a $\bar{D}^0$?



positive $\pi \Rightarrow D^0$

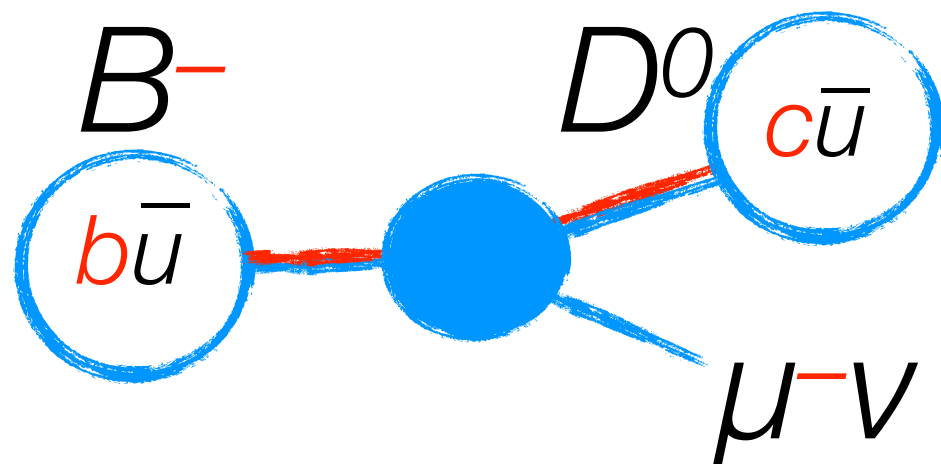


negative $\pi \Rightarrow \bar{D}^0$

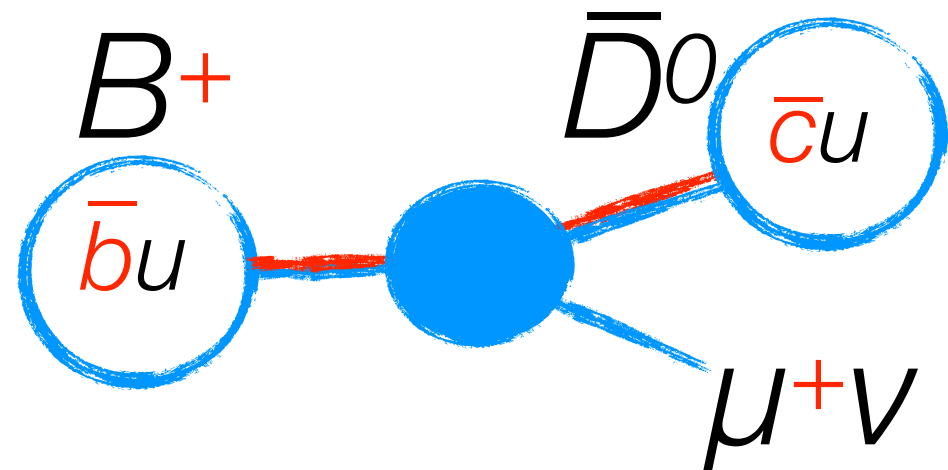
What do we measure?

$$A_{CP}(D^0 \rightarrow h^+ h^-) = \frac{\Gamma(D^0 \rightarrow h^+ h^-) - \Gamma(\bar{D}^0 \rightarrow h^+ h^-)}{\Gamma(D^0 \rightarrow h^+ h^-) + \Gamma(\bar{D}^0 \rightarrow h^+ h^-)}$$

How do we tell if it's a D^0 or a $\bar{D}^0$?



negative $\mu \Rightarrow D^0$

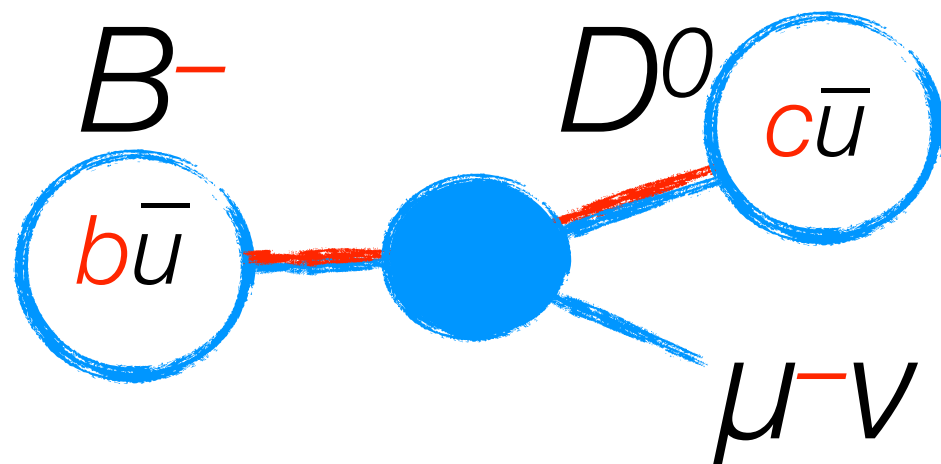


positive $\mu \Rightarrow \bar{D}^0$

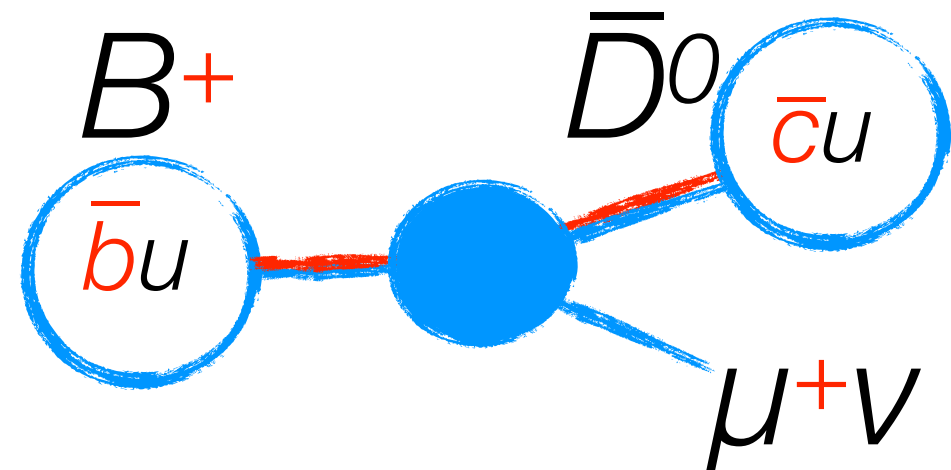
What do we measure?

$$A(D^0 \rightarrow h^+ h^-) = \frac{N(D^0 \rightarrow h^+ h^-) - N(\bar{D}^0 \rightarrow h^+ h^-)}{N(D^0 \rightarrow h^+ h^-) + N(\bar{D}^0 \rightarrow h^+ h^-)}$$

How do we tell if it's a D^0 or a $\bar{D}^0$?



negative $\mu \Rightarrow D^0$

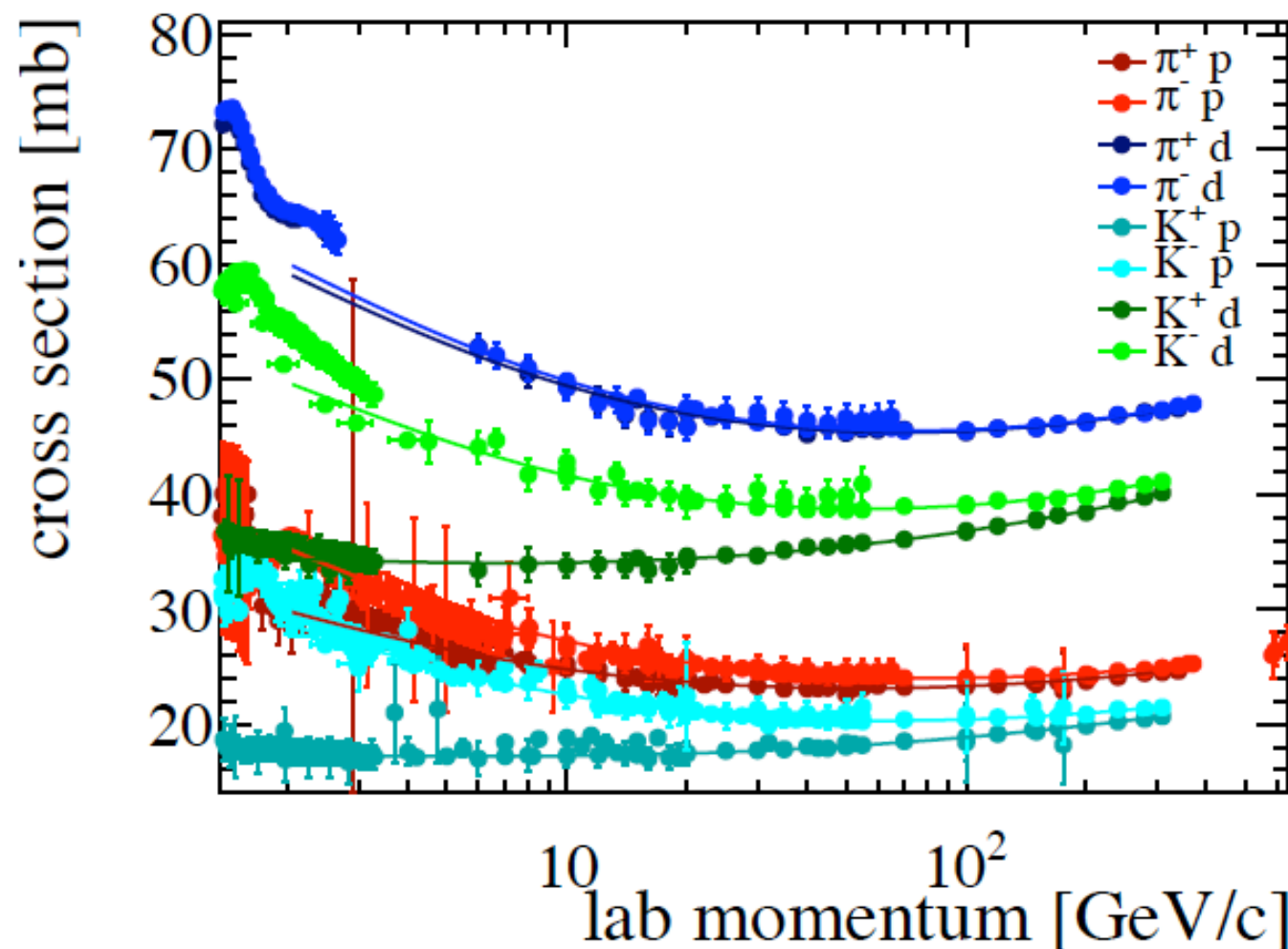


positive $\mu \Rightarrow \bar{D}^0$

Detection asymmetries

- Detector layout and/or differences in interaction cross section result in different reconstruction efficiencies for positively and negatively charged particles

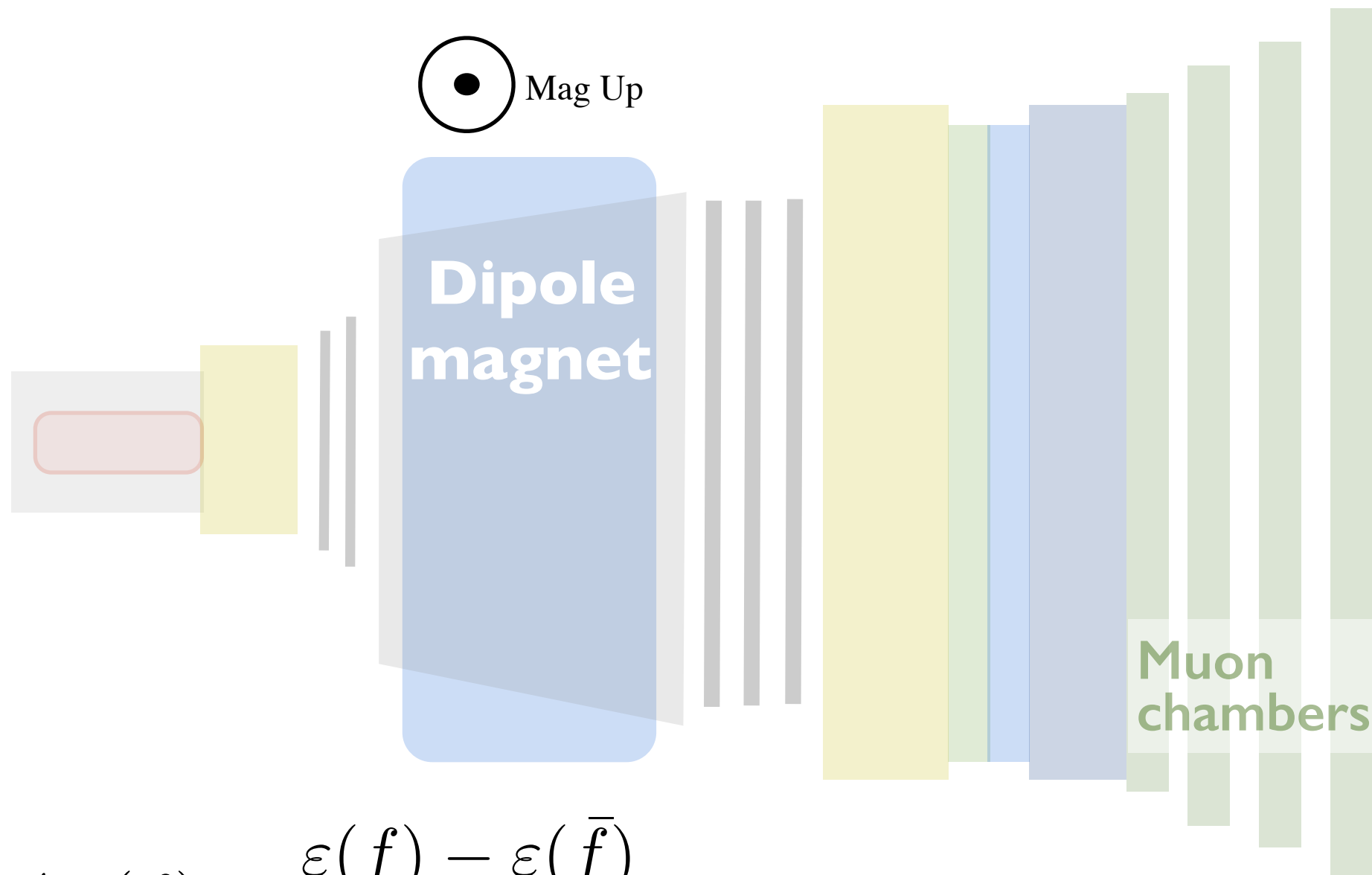
Chin. Phys. C 38 (2014) 090001



$$A_D(f) = \frac{\varepsilon(f) - \varepsilon(\bar{f})}{\varepsilon(f) + \varepsilon(\bar{f})}$$

Reversing the magnet polarity

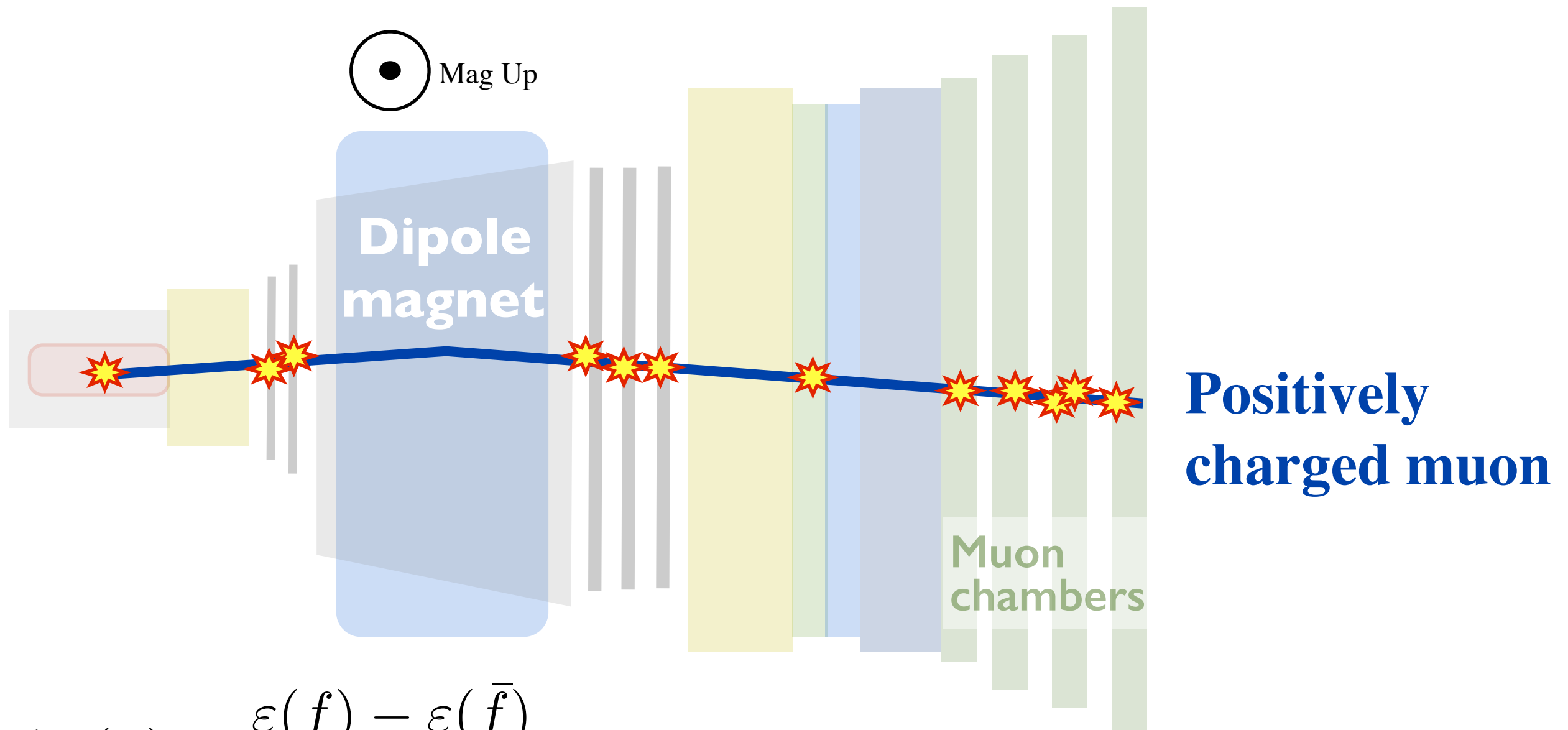
[LHCb-PUB-2014-006]



$$A_D(f) = \frac{\varepsilon(f) - \varepsilon(\bar{f})}{\varepsilon(f) + \varepsilon(\bar{f})}$$

Reversing the magnet polarity

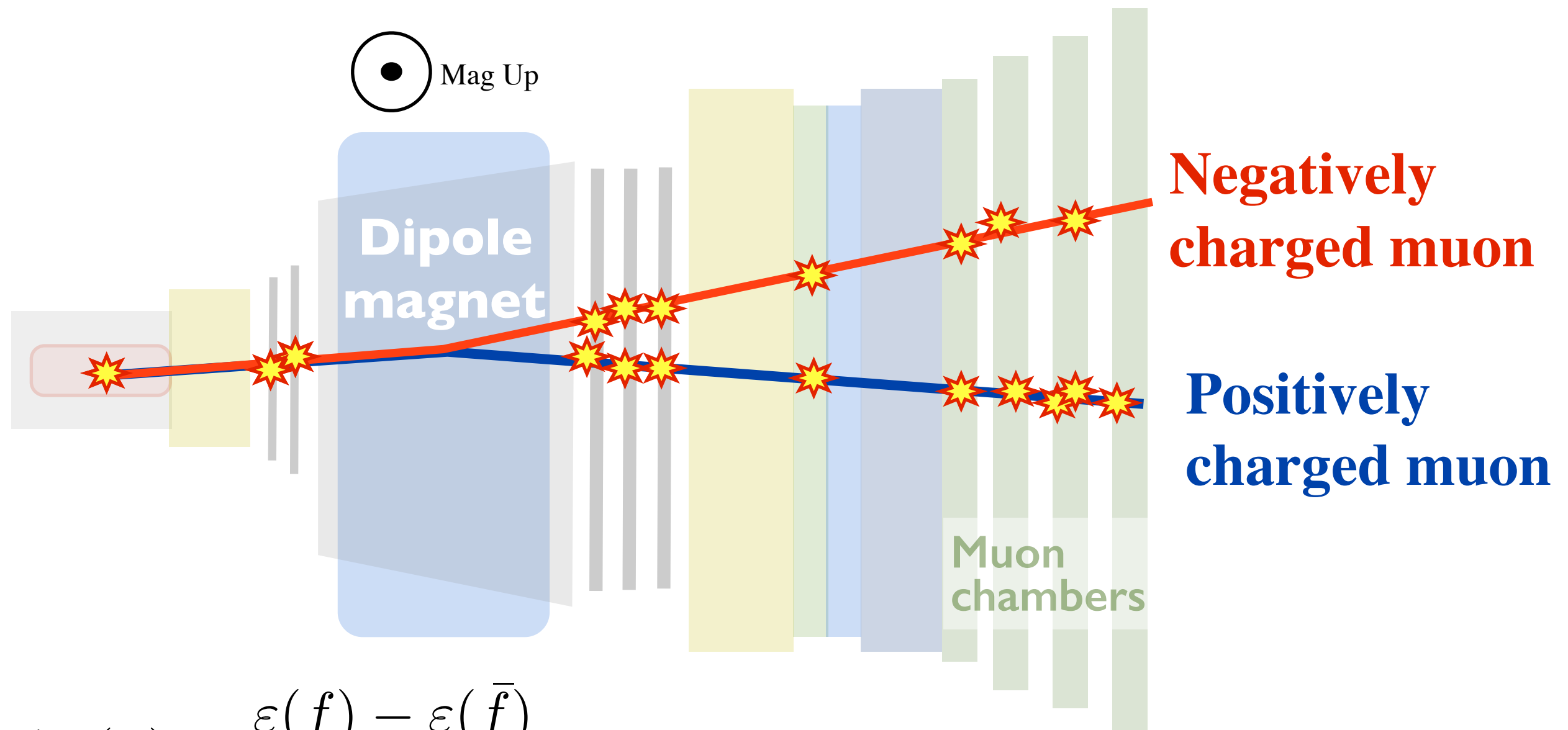
[LHCb-PUB-2014-006]



$$A_D(f) = \frac{\varepsilon(f) - \varepsilon(\bar{f})}{\varepsilon(f) + \varepsilon(\bar{f})}$$

Reversing the magnet polarity

[LHCb-PUB-2014-006]

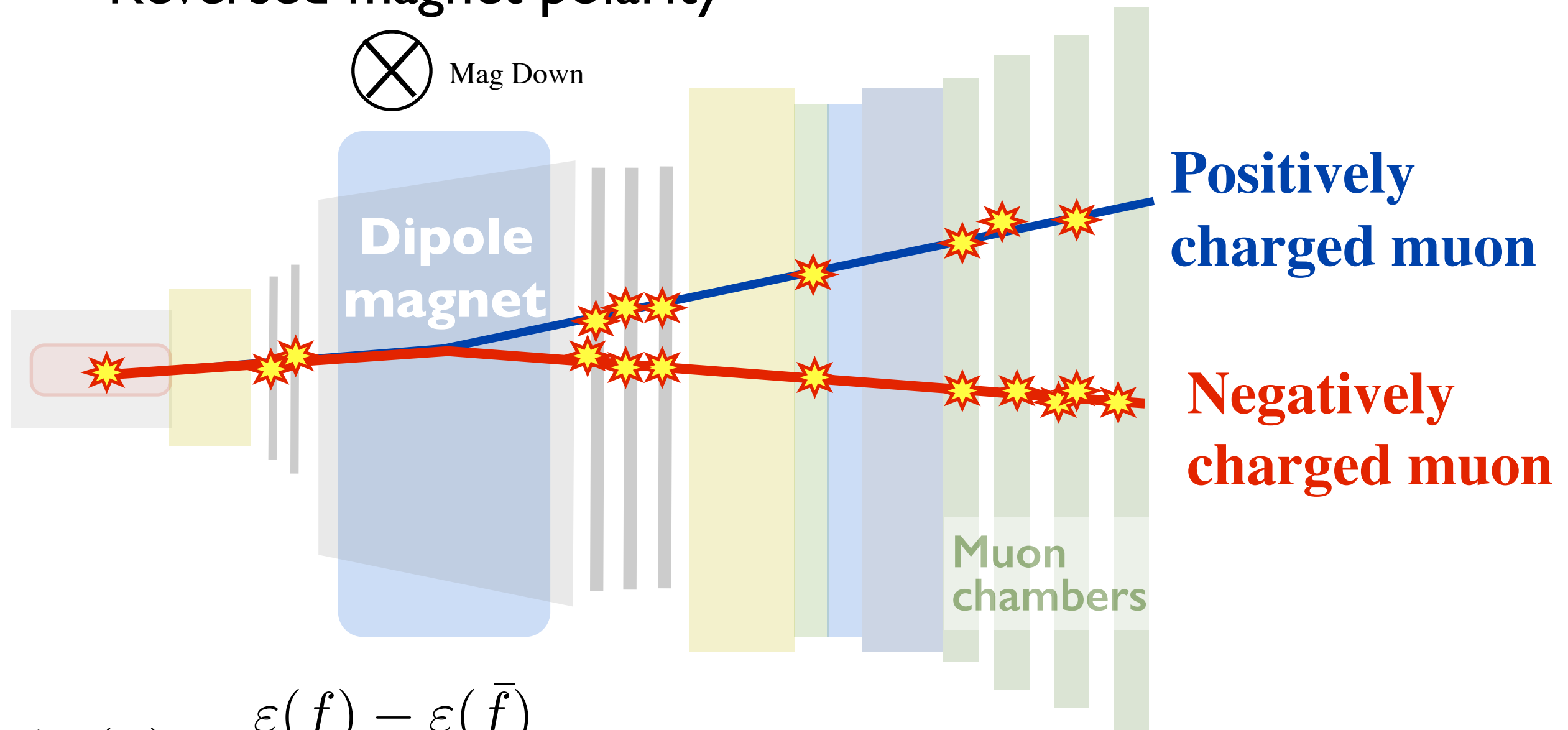


$$A_D(f) = \frac{\varepsilon(f) - \varepsilon(\bar{f})}{\varepsilon(f) + \varepsilon(\bar{f})}$$

Reversing the magnet polarity

[LHCb-PUB-2014-006]

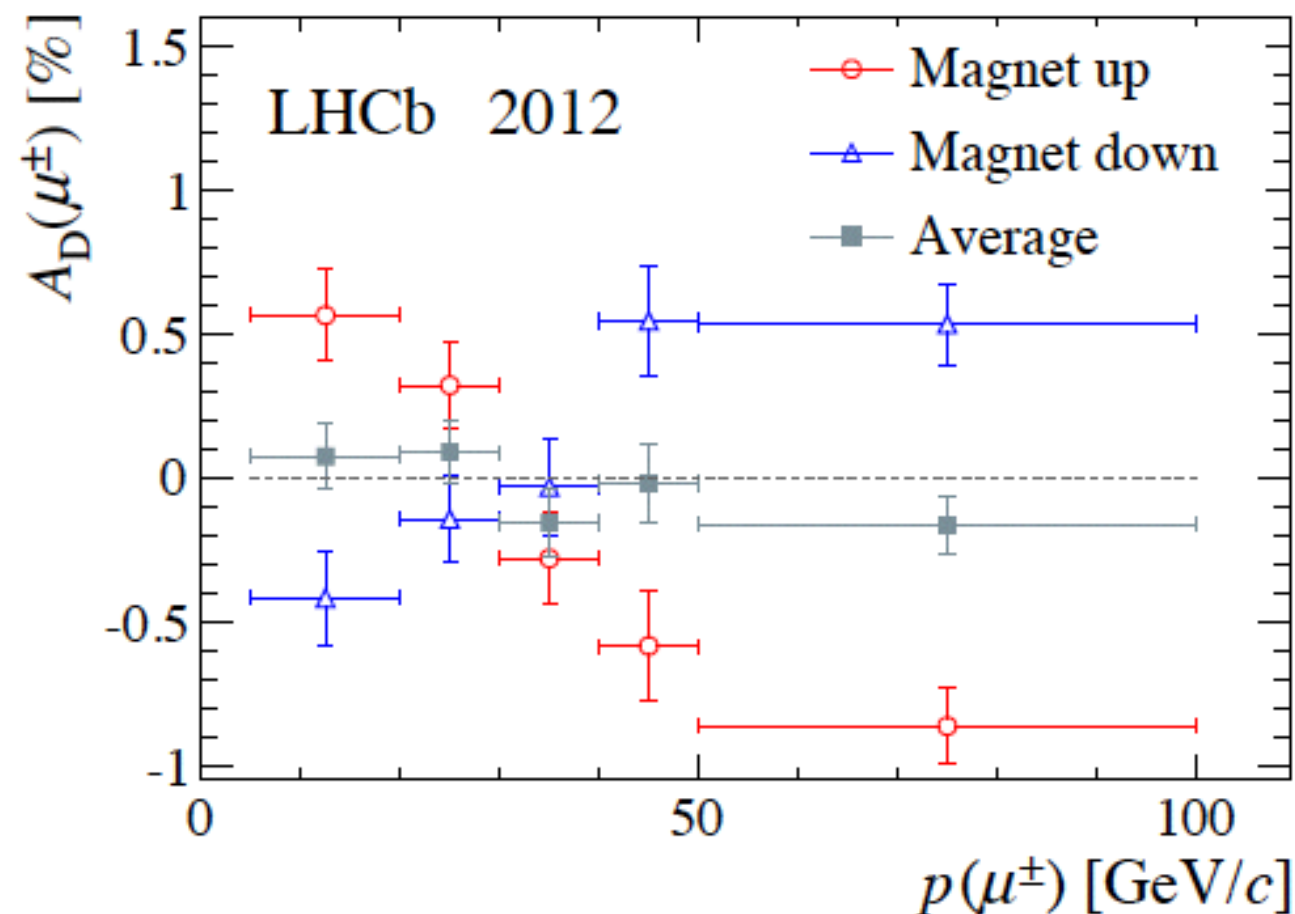
Reversed magnet polarity



$$A_D(f) = \frac{\varepsilon(f) - \varepsilon(\bar{f})}{\varepsilon(f) + \varepsilon(\bar{f})}$$

Reversing the magnet polarity

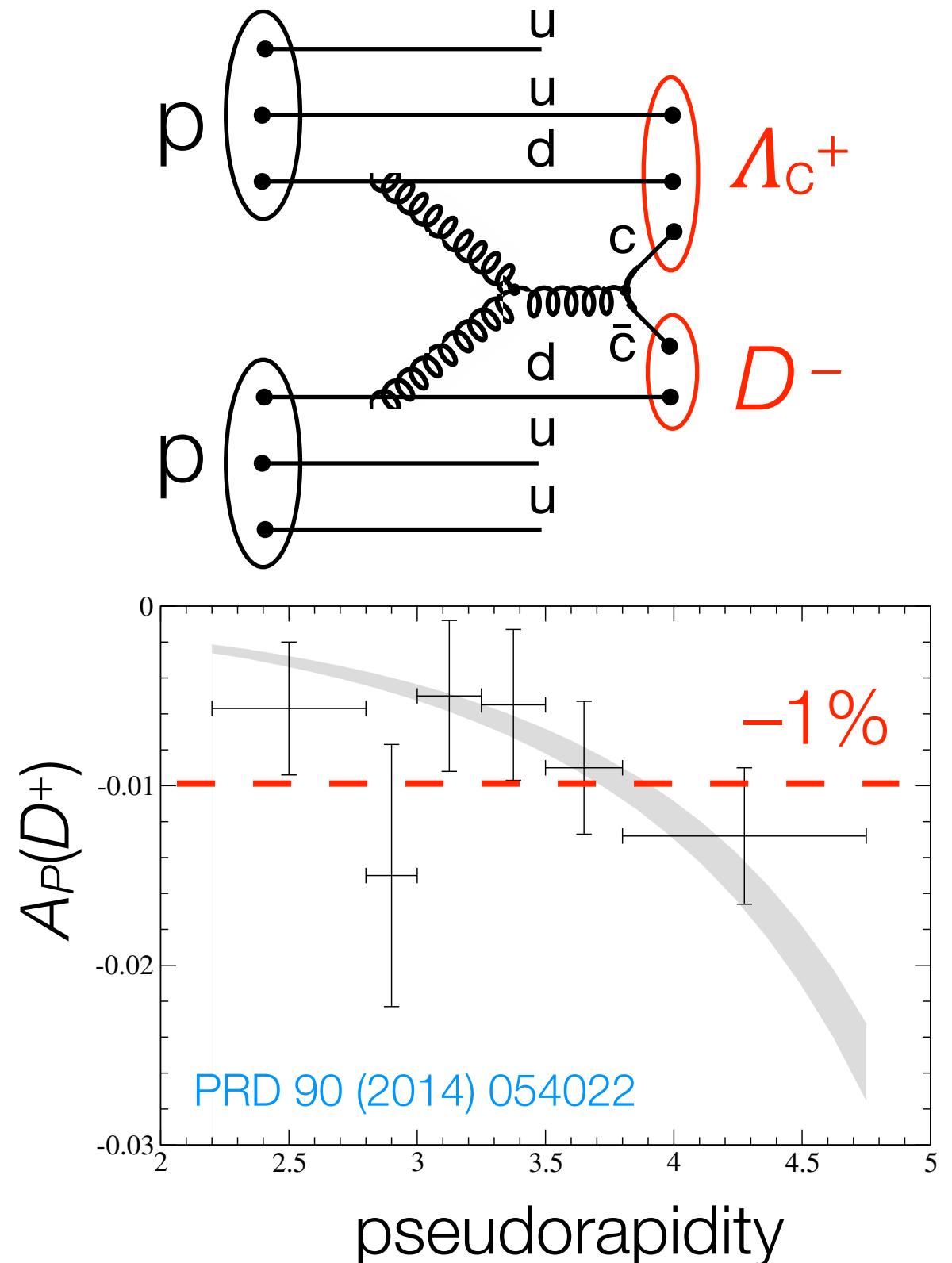
[LHCb-PUB-2014-006]



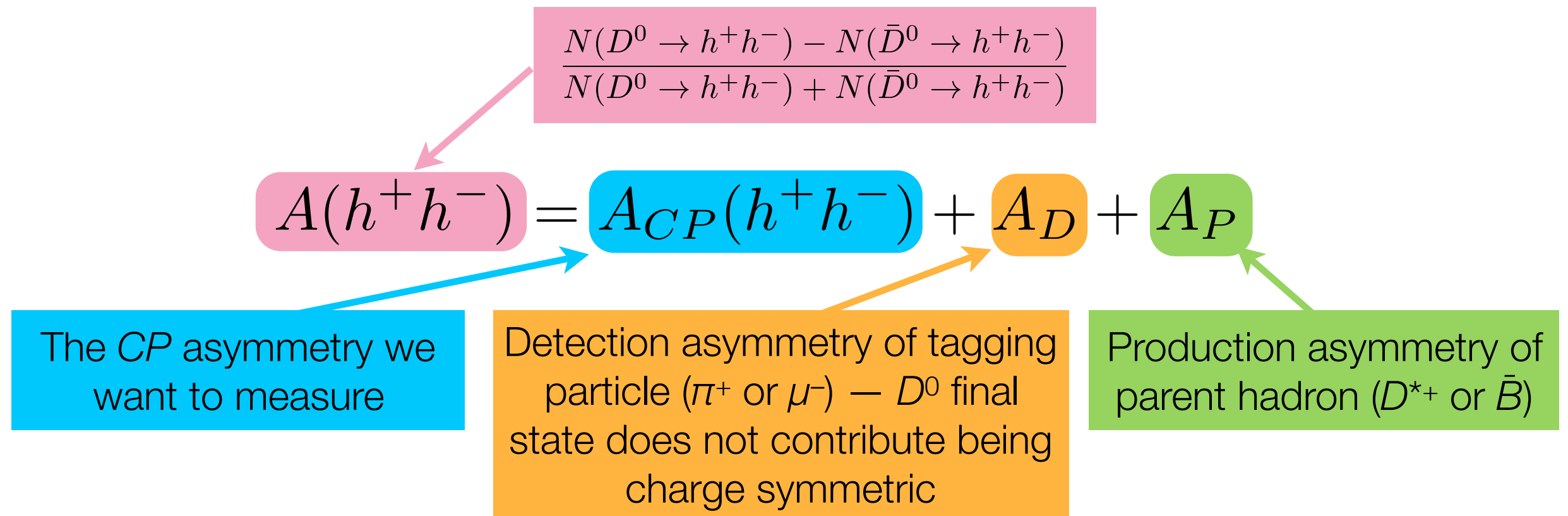
- Effectiveness depends on high frequency (2 weeks) of polarity flips
- Does not cancel asymmetries to 10^{-3} level, but crucial systematic check of results

Production asymmetry

- In pp collision the charm is produced mostly through gluon fusion (which is charge symmetric)
- The quarks may “recombine” with some of the proton valence quarks (beam-drag effect) creating an asymmetric production of charm mesons or baryons
- The effect is enhanced in the forward region, where the charm pair is collinear with the protons' direction



Why do we measure ΔA_{CP} ?

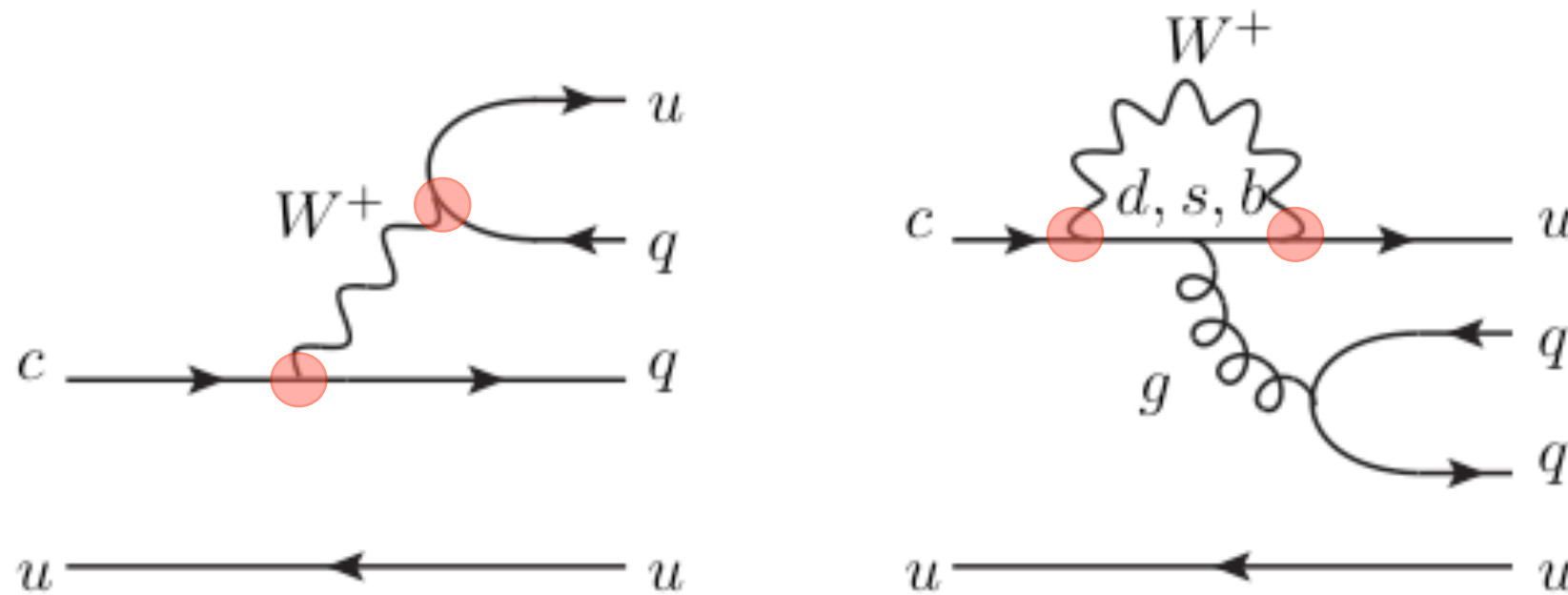


- The difference between raw asymmetries of two kinematically-similar final states removes the biases due to asymmetric detection and production

$$\Delta A_{CP} = A_{CP}(K^+ K^-) - A_{CP}(\pi^+ \pi^-) = A(K^+ K^-) - A(\pi^+ \pi^-)$$

Why do we measure ΔA_{CP} ?

- In the standard model, CP violation in $D^0 \rightarrow h^+ h^-$ is generated by the interference between the tree and the penguin diagrams

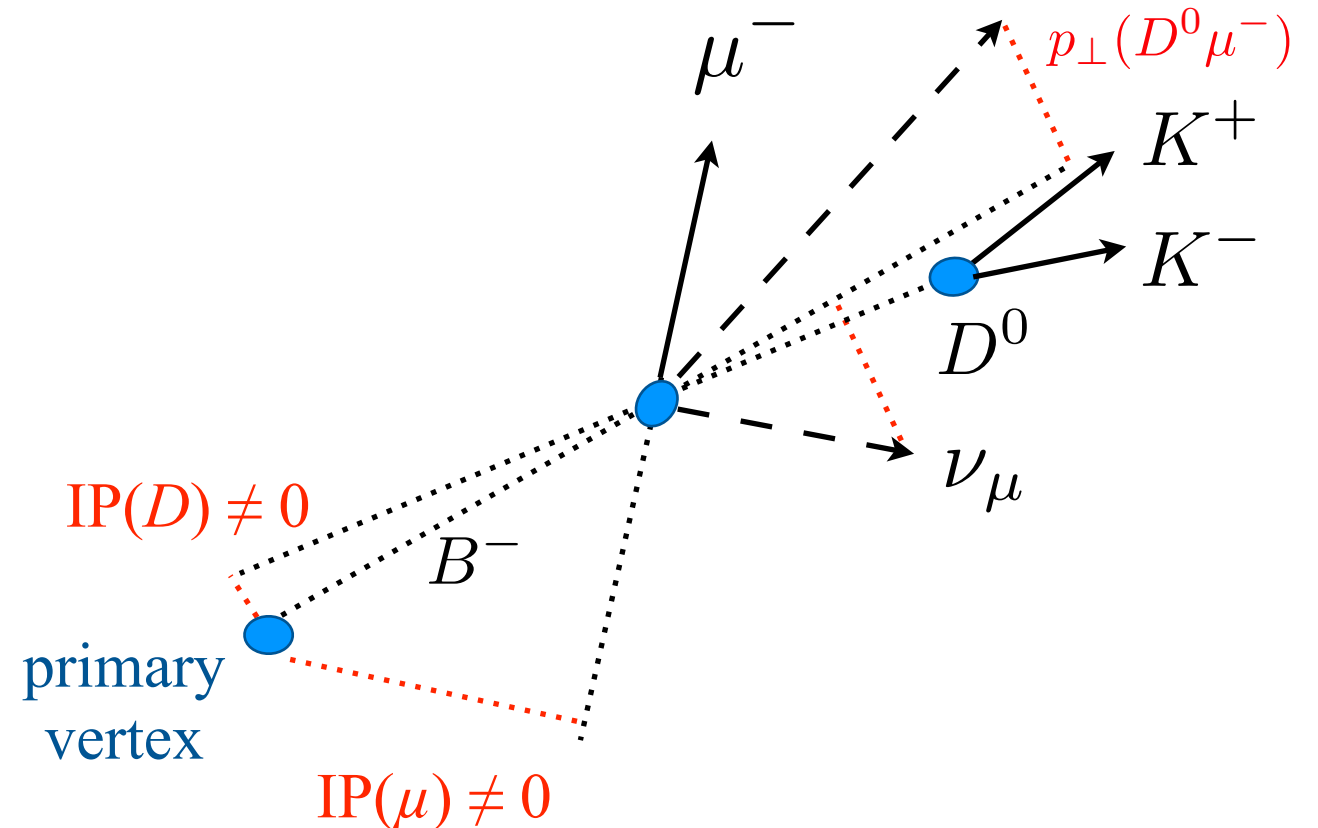
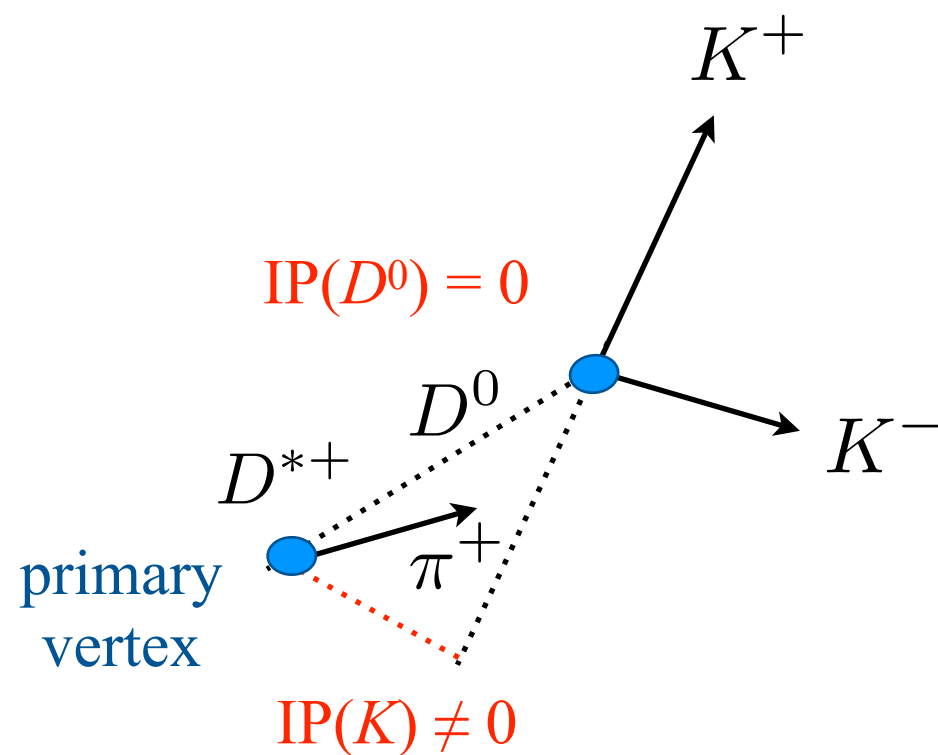


$$V_{ud}V_{cd}^* + V_{us}V_{cs}^* + V_{ub}V_{cb}^* = 0$$

- In the limit of $SU(3)$ symmetry (*i.e.*, ignoring QCD effects), $A_{CP}(K^+K^-)$ and $A_{CP}(\pi^+\pi^-)$ have same magnitude and opposite signs
- In addition to be robust against experimental biases, ΔA_{CP} provides enhanced sensitivity to CP violation

Signal (online) selection/reconstruction

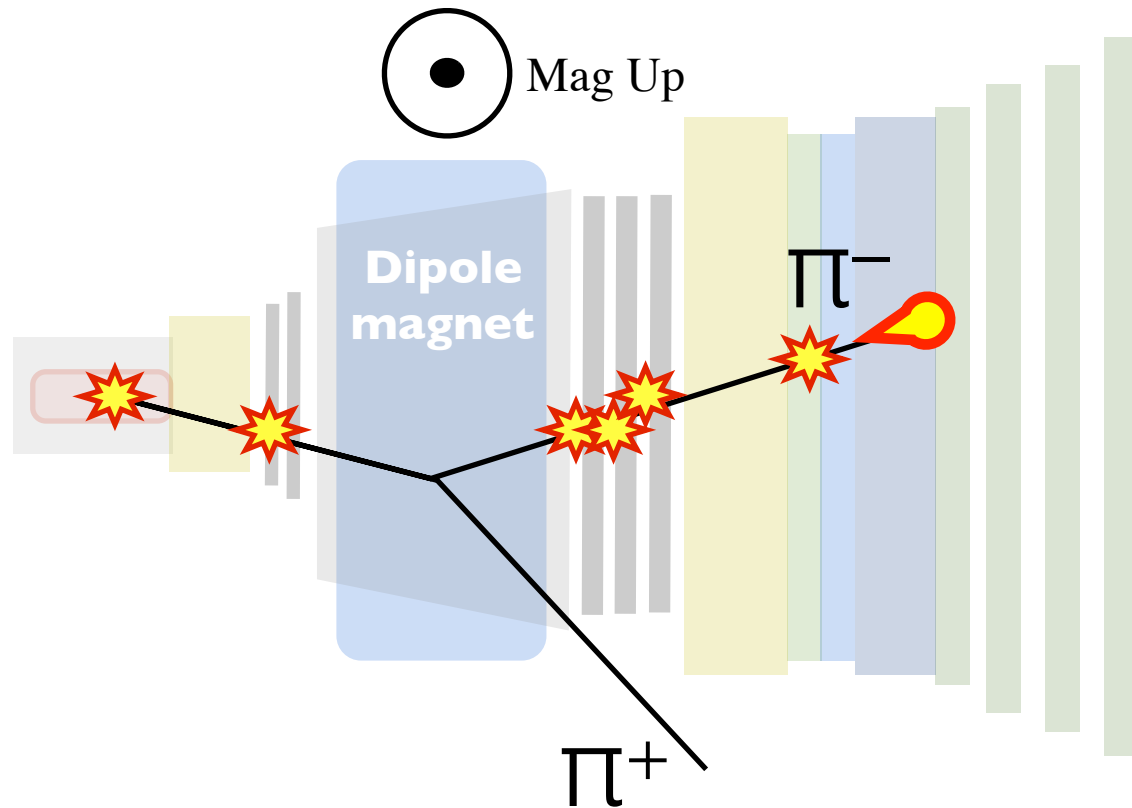
$$m_{\text{corr}} = \sqrt{m^2(D^0\mu^-) + p_{\perp}(D^0\mu^-)^2 + p_{\perp}(D^0\mu^-)}$$



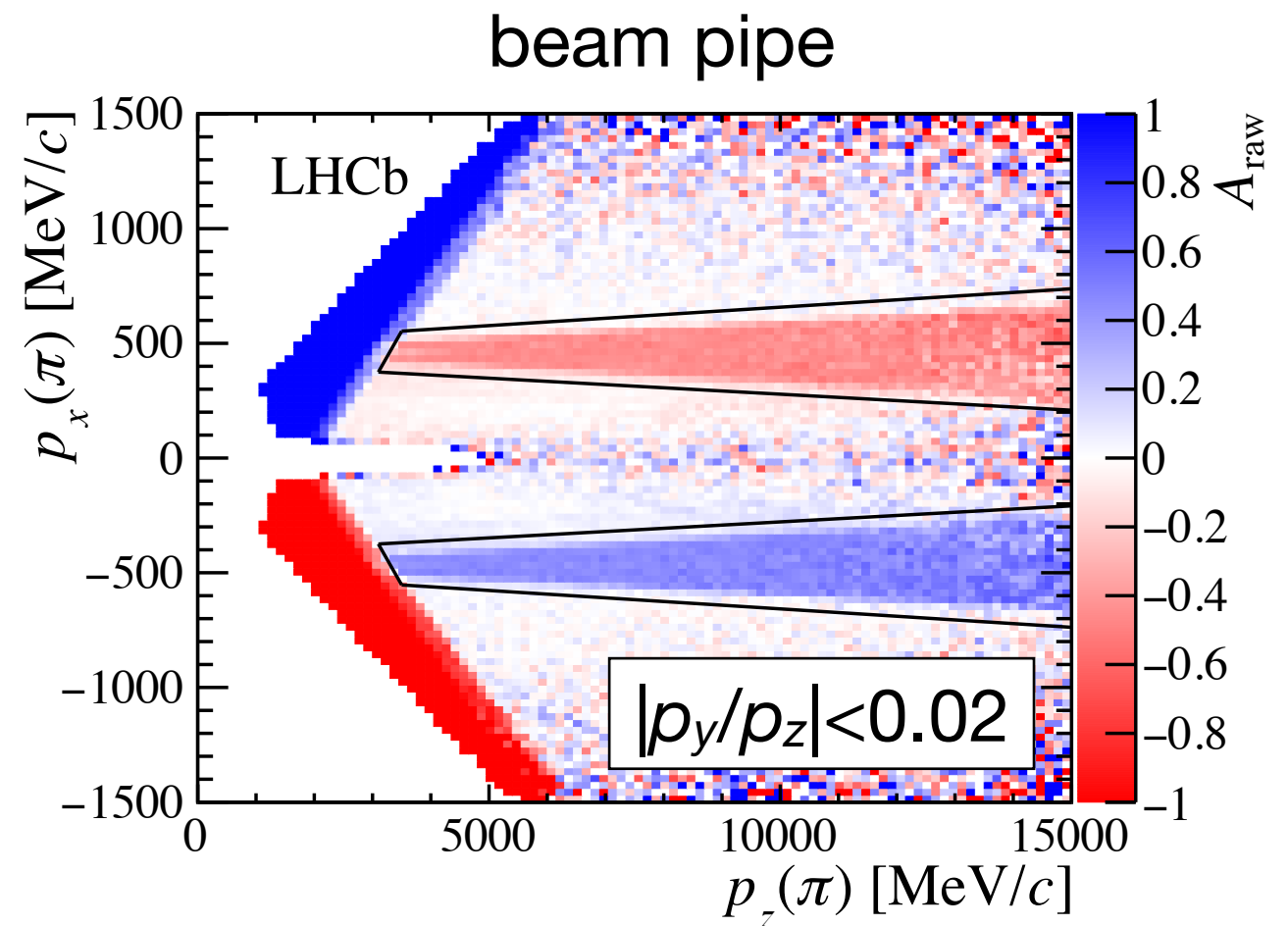
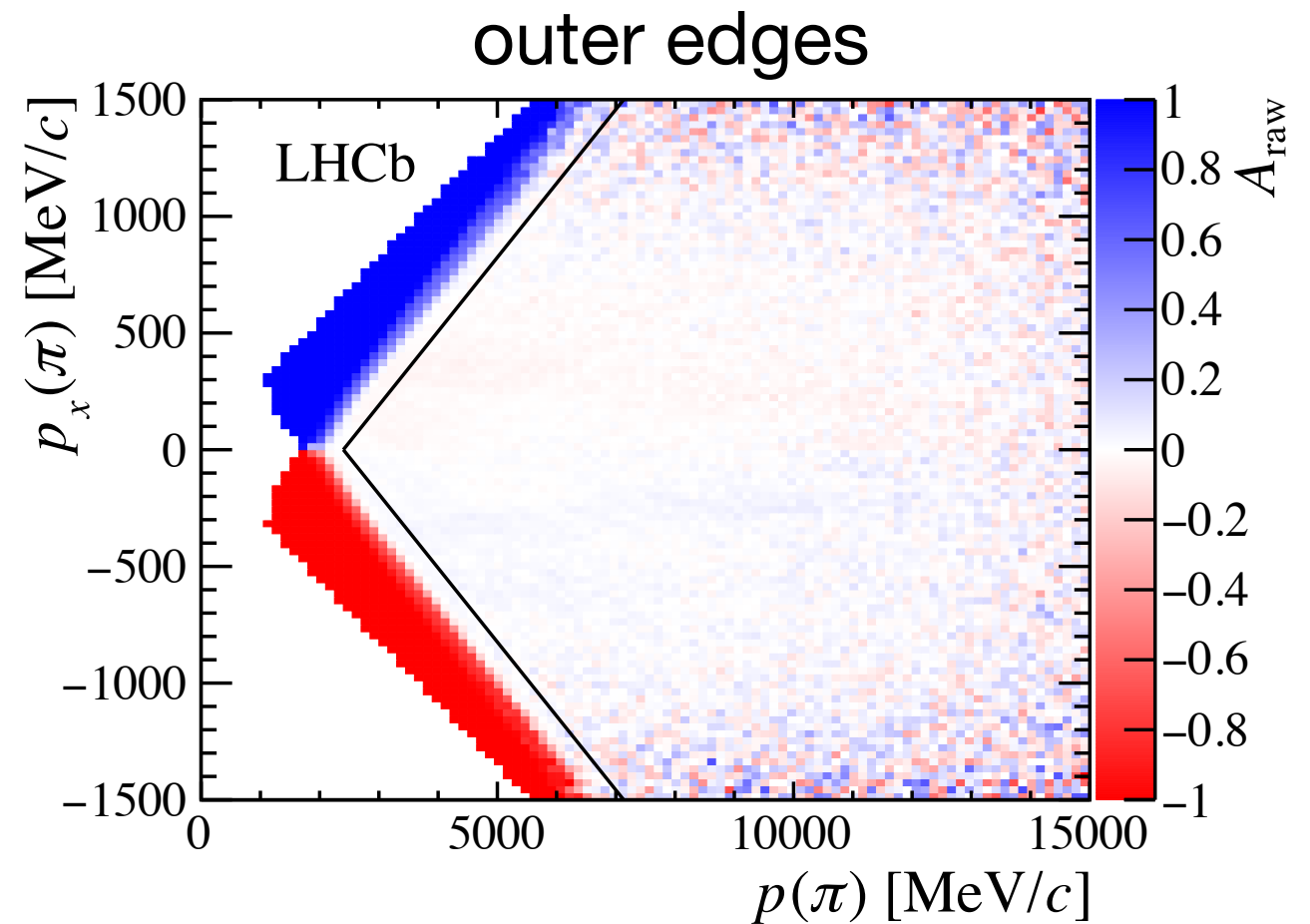
- π tagged: well identified kaons/pions forming a displaced secondary vertex, paired with low-momentum pion to form a D^{*+} vertex that coincides with the PV
- μ tagged: well identified and displaced muon paired with D^0 candidate, requirement on corrected mass to (partly) compensate for undetected neutrino

Fiducial requirements

- For specific regions of phase space, the soft pion of a given charge is kicked out of the detector acceptance by the magnetic field

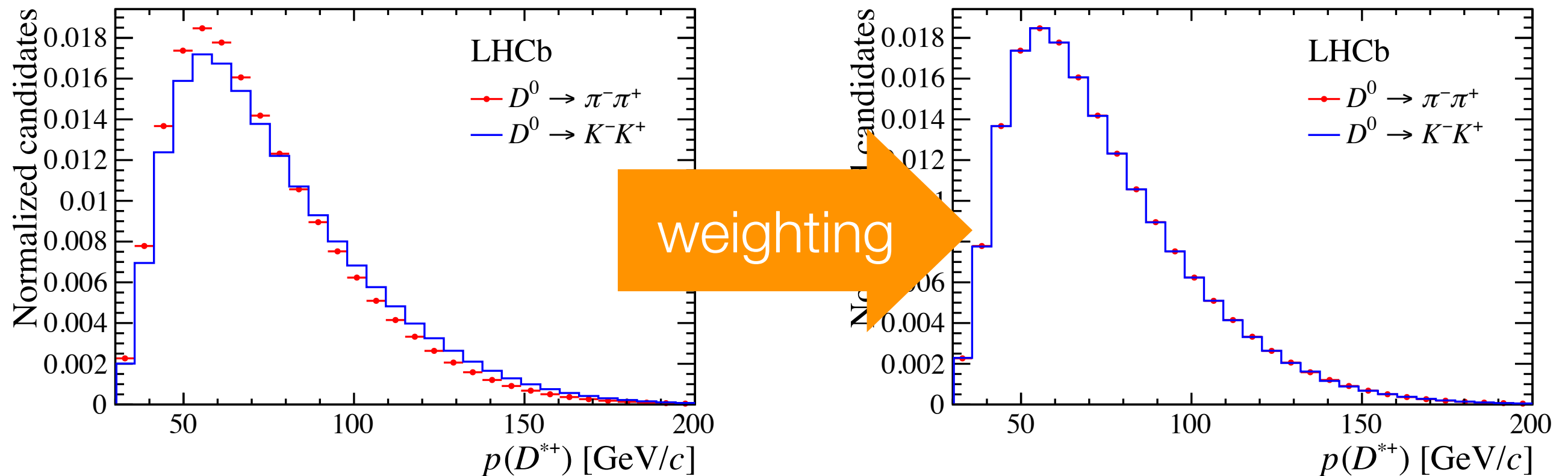


- $O(1)$ detection asymmetries break the linear decomposition of the raw asymmetry

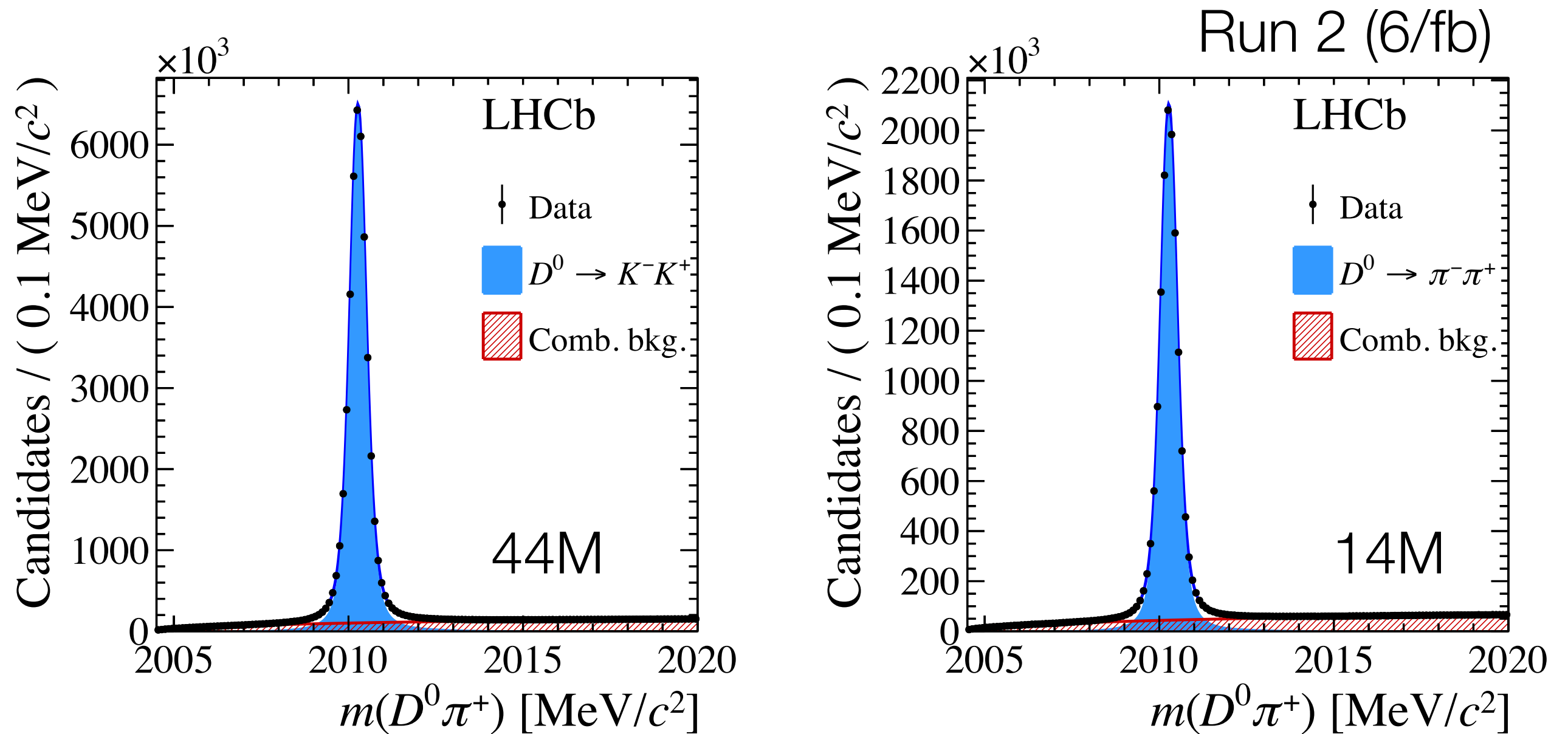


Kinematic weighting

- Detection and production asymmetries are expected to depend on kinematics
- Cancellation of nuisance asymmetries may not be accurate if the kinematic distributions of the reconstructed $D^0 \rightarrow K^+ K^-$ and $D^0 \rightarrow \pi^+ \pi^-$ candidates differ
- The observed differences are removed through a weighting procedure — the resulting net effect on ΔA_{CP} is below 10^{-4}

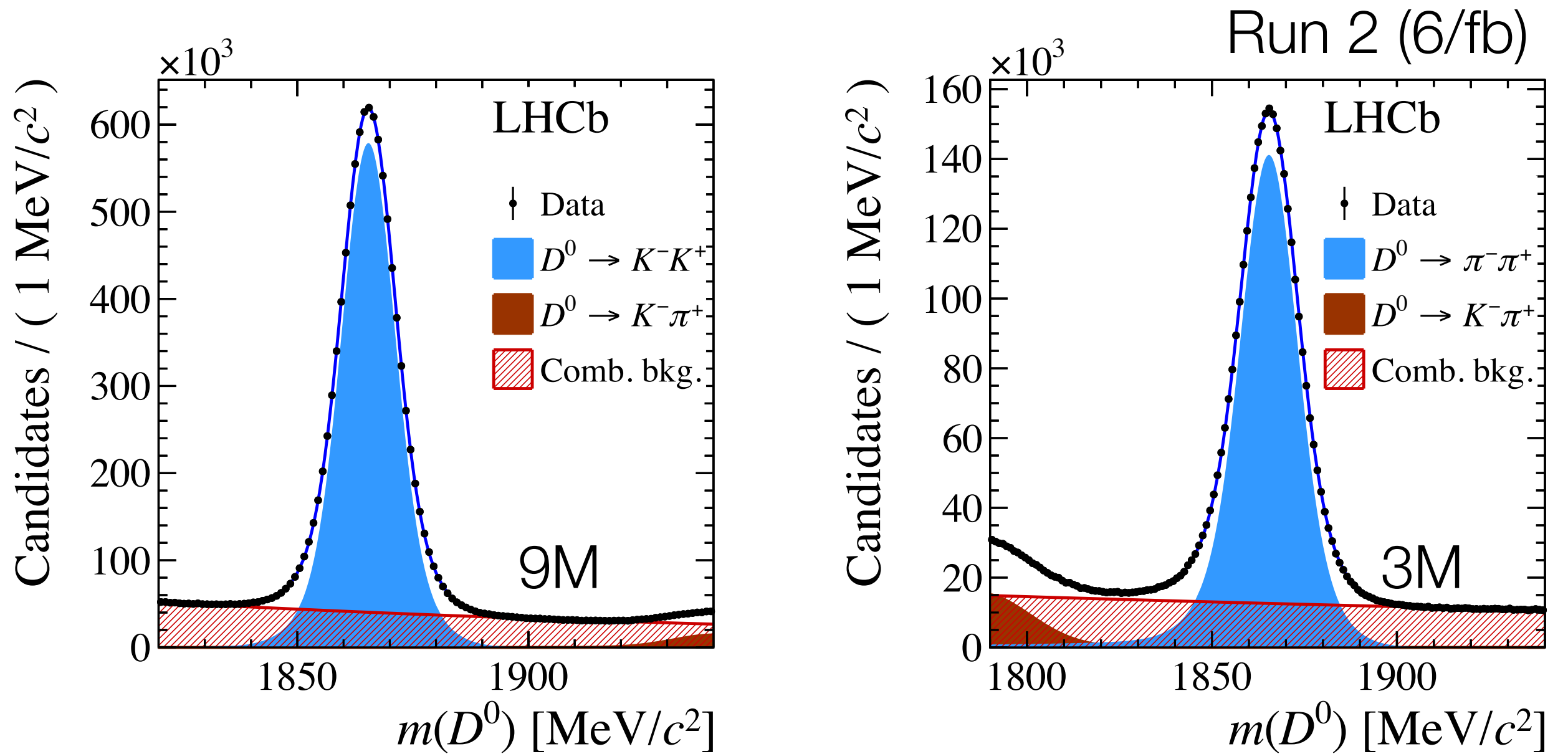


π tagged samples



Raw asymmetries measured with sub-0.1% statistical precision from simultaneous fit to D^0 and $\bar{D}^0$ samples

μ tagged samples



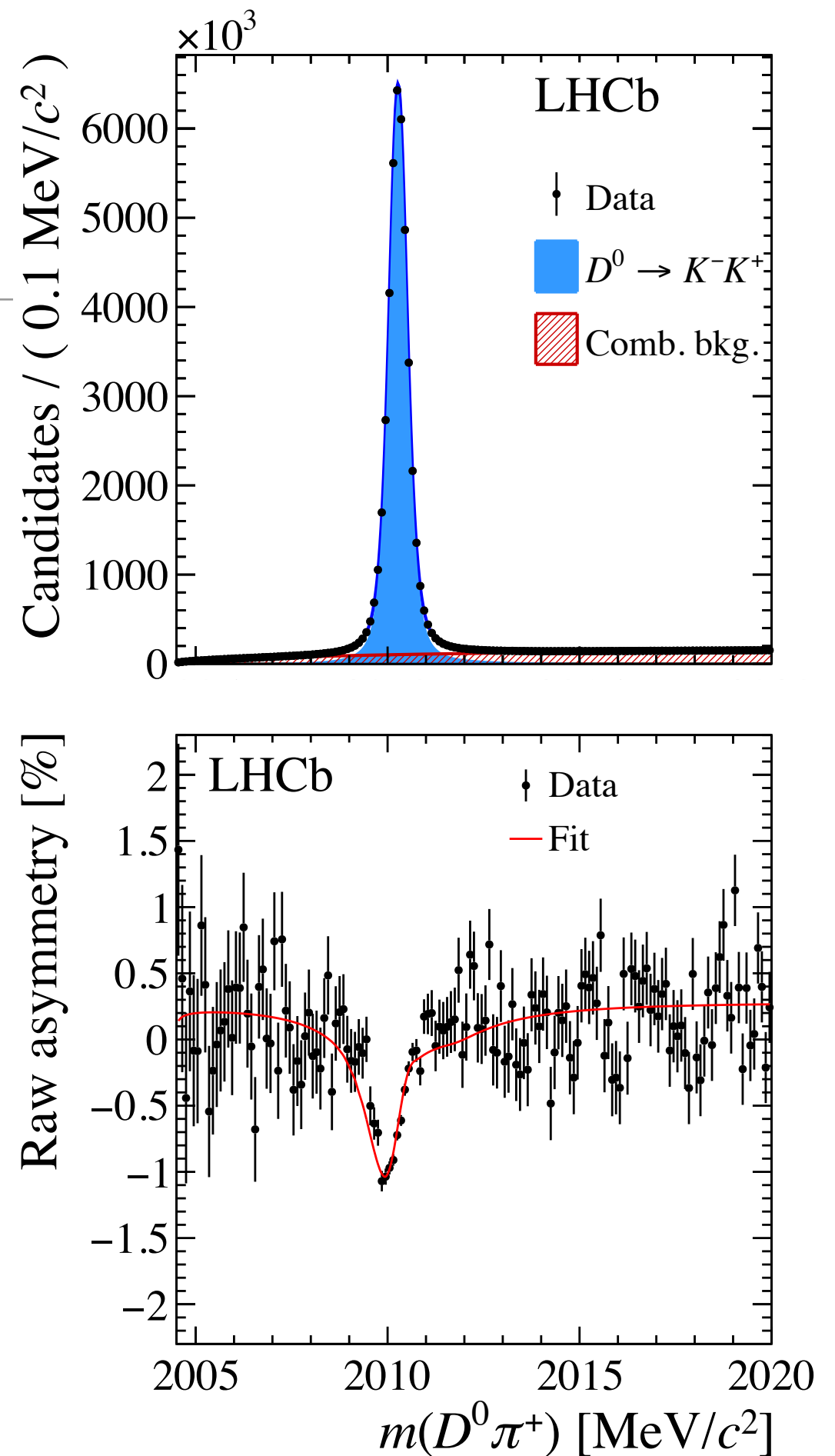
Raw asymmetries measured with sub-0.1% statistical precision from simultaneous fit to D^0 and $\bar{D}^0$ samples

Systematic uncertainties

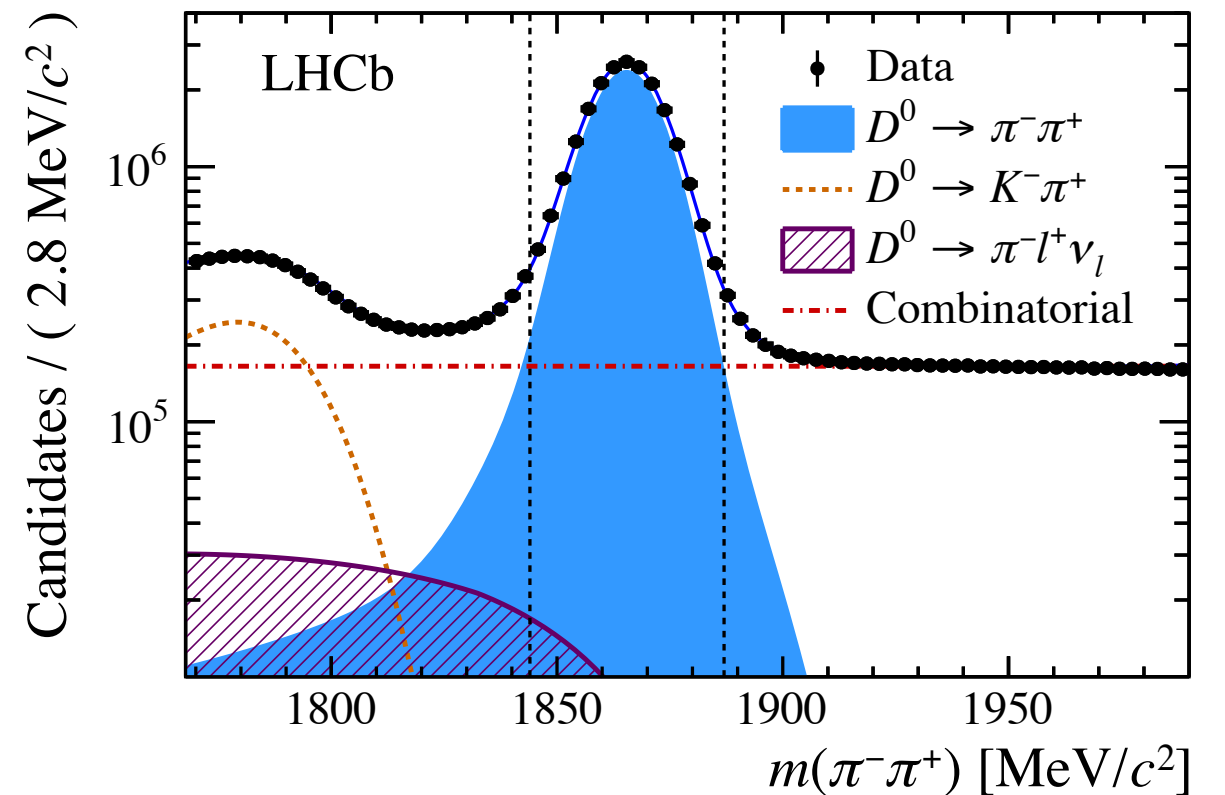
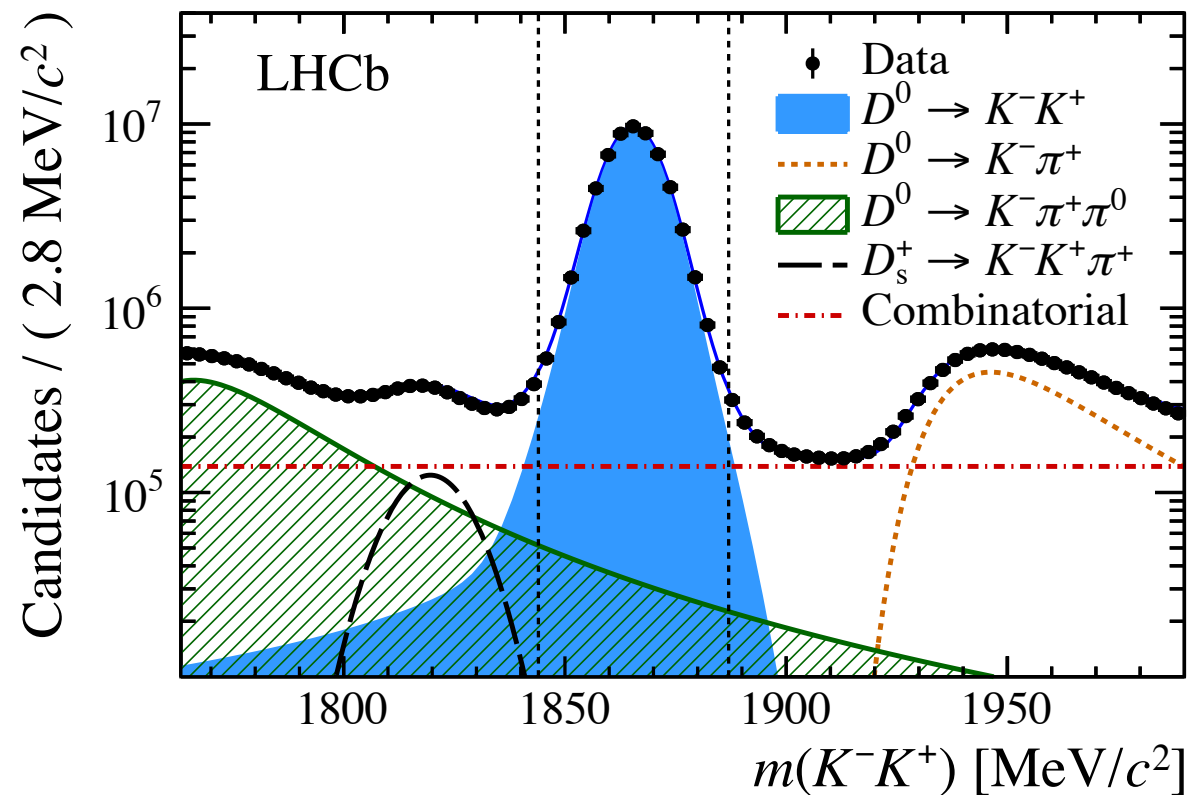
Source	[10 ⁻⁴]	
	π -tagged	μ -tagged
Fit model	0.6	2
Mistag	—	4
Weighting	0.2	1
Secondary decays	0.3	—
B fractions	—	1
B reco. efficiency	—	2
Peaking background	0.5	—
Total	0.9	5

Fit model

- Mass model with $O(10M)$ candidates is never perfect
- Impact on asymmetries is generally marginal
- However, mass shapes may also be charge asymmetric (momentum scale/resolution is slightly asymmetric)
- Tested a variety of alternative models, bias on signal asymmetry estimated using pseudoexperiments



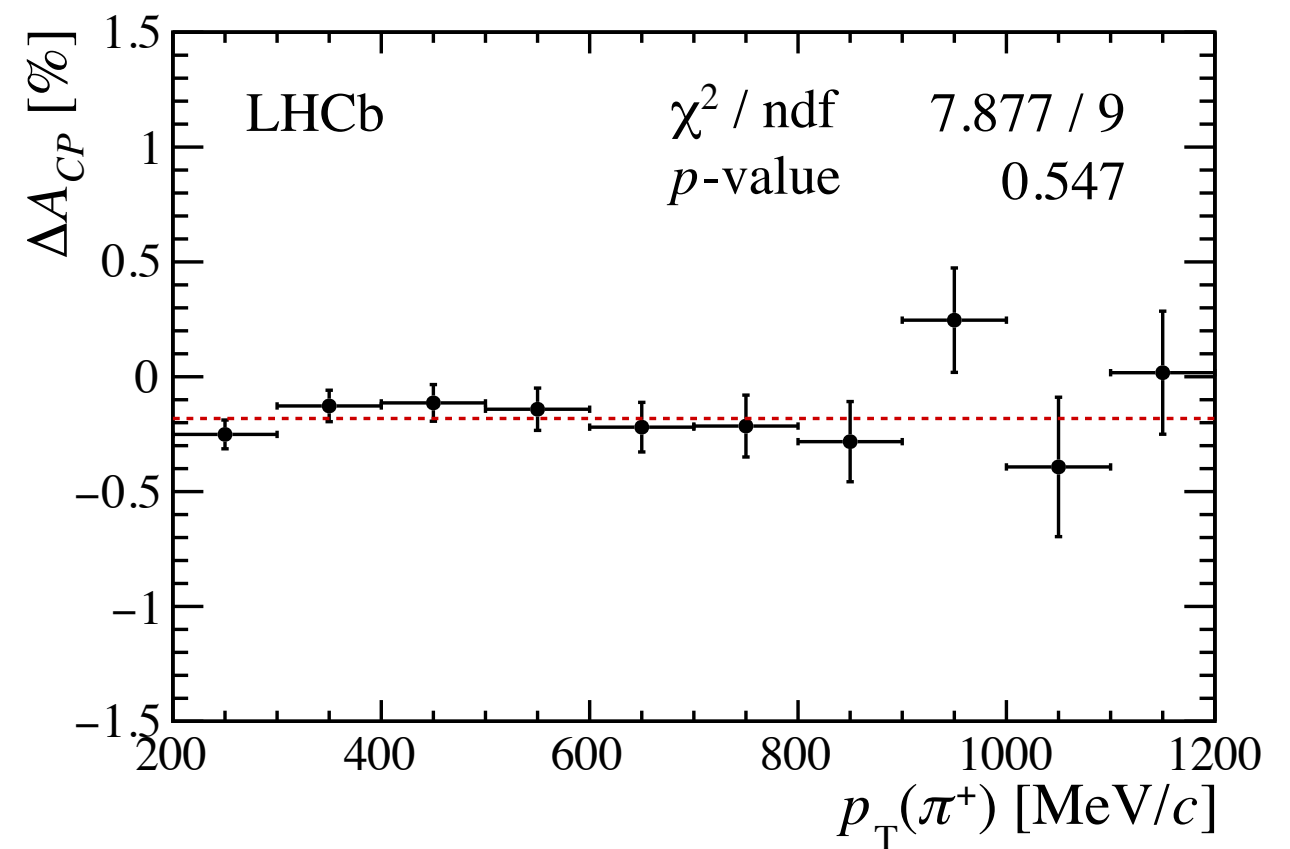
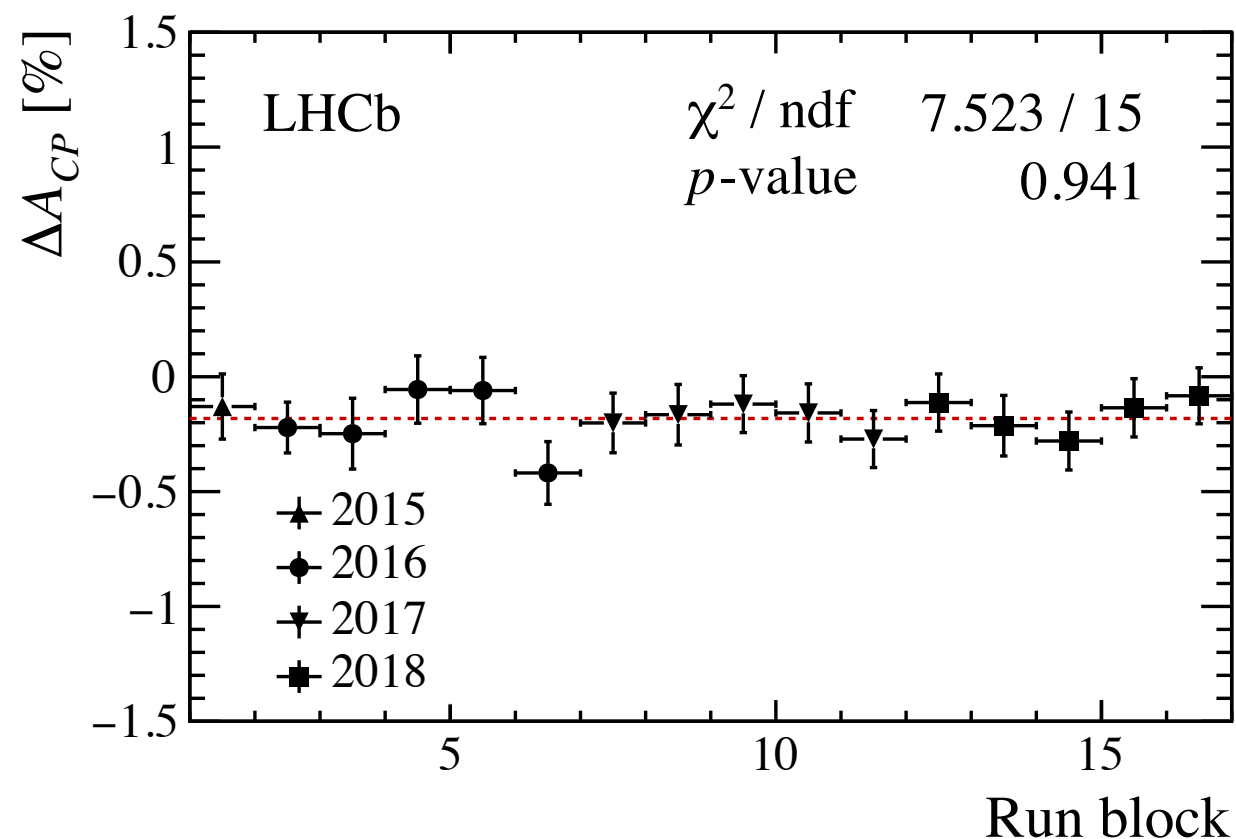
Peaking background



- Backgrounds due to $D^{*+} \rightarrow D^0\pi^+$ decays, where the D^0 is mis- or partially reconstructed would peak in $m(D^0\pi^+)$
- Bias on signal asymmetry estimated by measuring backgrounds asymmetry and contamination in the D^0 -mass signal region

Consistency checks

- Measured value of ΔA_{CP} studied as a function of several variables (e.g., magnet polarity, data-taking period, kinematics)



- No evidence found for unexpected dependences

Results

$$\Delta A_{CP} [10^{-4}]$$

$$-18.2 \pm 3.2 \text{ (stat)} \pm 0.9 \text{ (syst)} \quad \pi \text{ tagged (Run 2)}$$

$$-9 \pm 8 \text{ (stat)} \pm 5 \text{ (syst)} \quad \mu \text{ tagged (Run 2)}$$

$$-10 \pm 8 \text{ (stat)} \pm 3 \text{ (syst)} \quad \pi \text{ tagged (Run 1)}$$

[\[PRL 116 \(2016\) 191601\]](#)

$$-14 \pm 16 \text{ (stat)} \pm 8 \text{ (syst)} \quad \mu \text{ tagged (Run 1)}$$

[\[JHEP 07 \(2014\) 041\]](#)

$$\boxed{-15.4 \pm 2.9 \text{ (stat+syst)}}$$

combined

5.3 σ deviation from zero
first observation of CP violation in charm

Interpretation

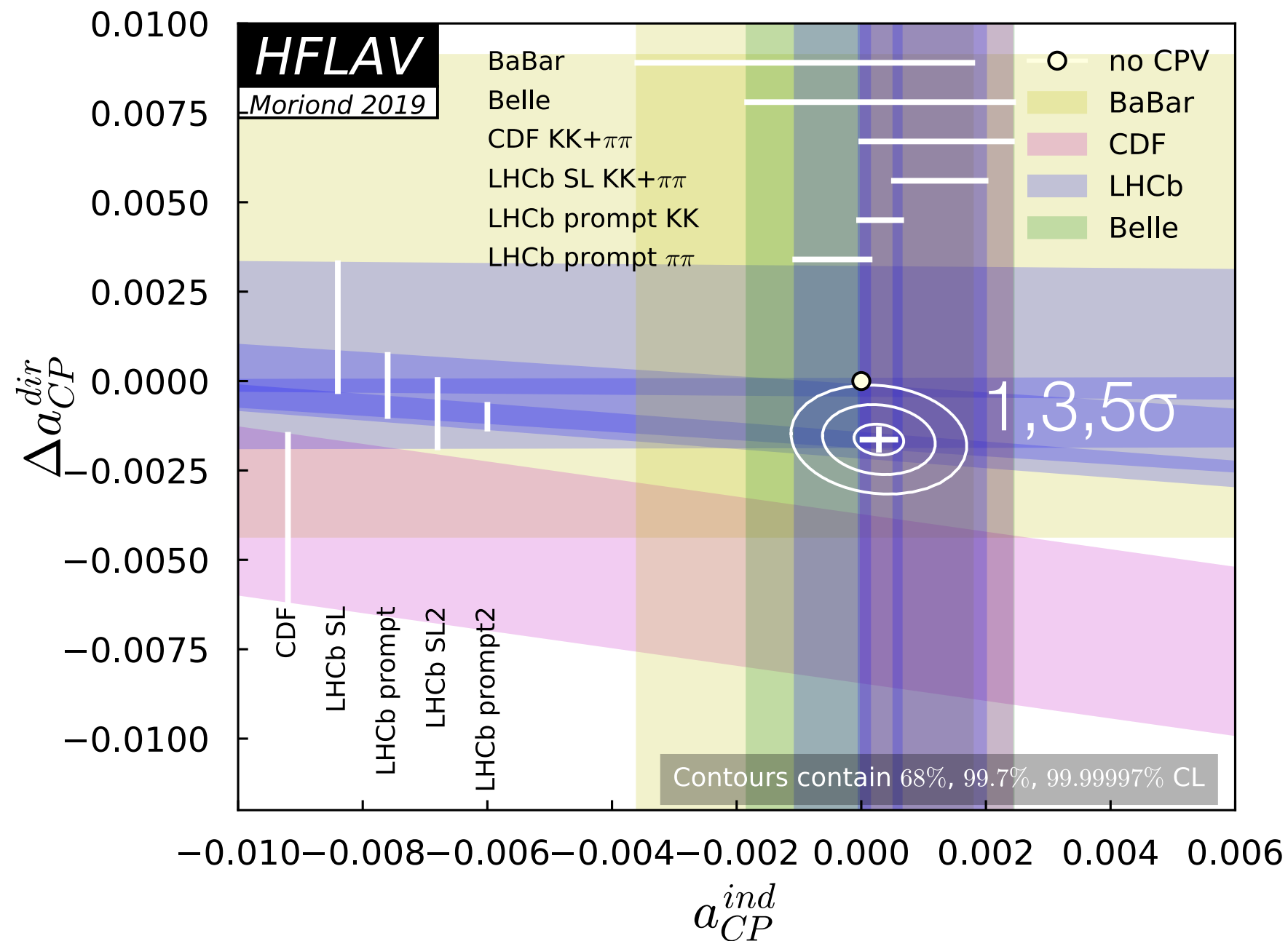
- The time-integrated asymmetry receives contribution from CP violation in the decay amplitudes (direct) and in D^0 - $\bar{D}^0$ mixing (indirect)

$$A_{CP}(h^+h^-) \approx a_{CP}^{\text{dir}}(h^+h^-) + \frac{\langle t \rangle}{\tau} a_{CP}^{\text{ind}}$$

- Indirect contribution is to a good approximation independent of the decay mode
- Since $\Delta\langle t \rangle \sim 0.1\tau$, ΔA_{CP} is mostly a measurement of direct CP violation

$$\Delta A_{CP} \approx \Delta a_{CP}^{\text{dir}} + \frac{\Delta\langle t \rangle}{\tau} a_{CP}^{\text{ind}}$$

Experimental status of CP violation in $D^0 \rightarrow h^+ h^-$



$$\Delta a_{CP}^{dir} = (-16.4 \pm 2.8) \times 10^{-4}$$

$$a_{CP}^{ind} = (2.8 \pm 2.6) \times 10^{-4}$$

so what? shall I
be excited?



yes! it's new
physics!

(just two examples)

it's not, but
we can learn a lot
about QCD

Alexander Lenz



[arXiv:1903.10490]

Yuval Grossman



[arXiv:1903.10952]

yes! it's new
physics!

(just two examples)

it's not, but

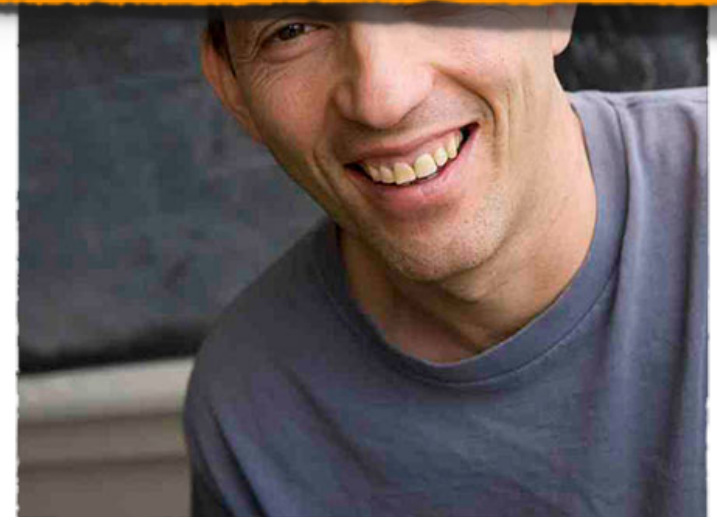
- Experimentally test relations between CP asymmetries of many flavor-SU(3) related decays to help constraining the QCD parameters
- Huge program of measurements. Belle II role will be crucial

Alexander Lenz



[arXiv:1903.10490]

Yuval Grossman



[arXiv:1903.10952]

What about CP violation in mixing?

- Time-evolution described by Schrödinger's equation

$$i \frac{\partial}{\partial t} \begin{pmatrix} |P^0(t)\rangle \\ |\bar{P}^0(t)\rangle \end{pmatrix} = \left[\begin{pmatrix} M_{11} & M_{12} \\ M_{12}^* & M_{22} \end{pmatrix} - \frac{i}{2} \begin{pmatrix} \Gamma_{11} & \Gamma_{12} \\ \Gamma_{12}^* & \Gamma_{22} \end{pmatrix} \right] \begin{pmatrix} |P^0(t)\rangle \\ |\bar{P}^0(t)\rangle \end{pmatrix}$$

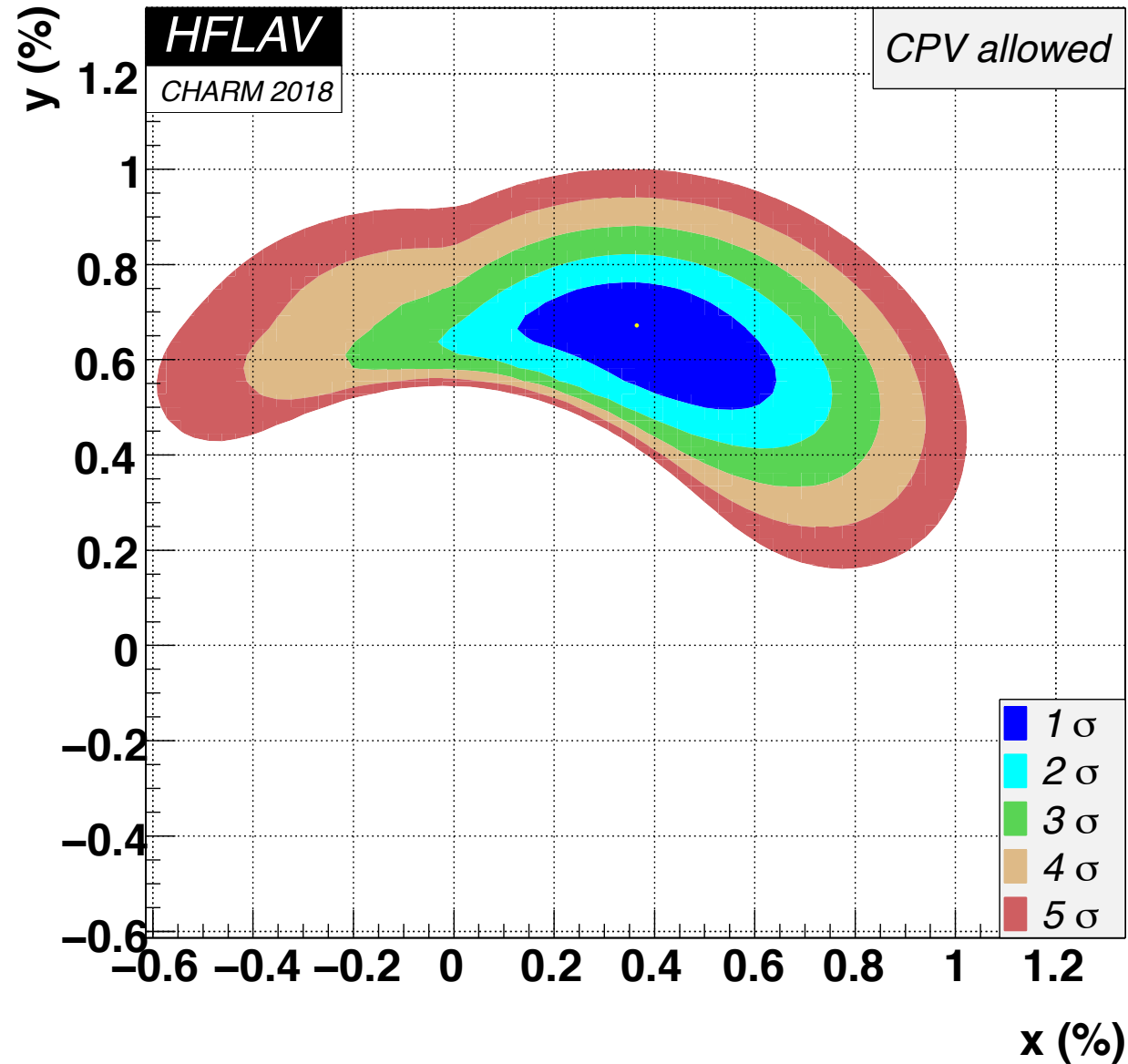
- Eigenstates can have different masses and decay width

$$|P_{1,2}\rangle = p|P^0\rangle \pm q|\bar{P}^0\rangle \quad \frac{q}{p} = \sqrt{\frac{M_{12}^* - \frac{i}{2}\Gamma_{12}^*}{M_{12} - \frac{i}{2}\Gamma_{12}}}$$

$$x = \frac{m_1 - m_2}{\Gamma} \quad y = \frac{\Gamma_1 - \Gamma_2}{2\Gamma} \quad \Gamma = \frac{\Gamma_1 + \Gamma_2}{2}$$

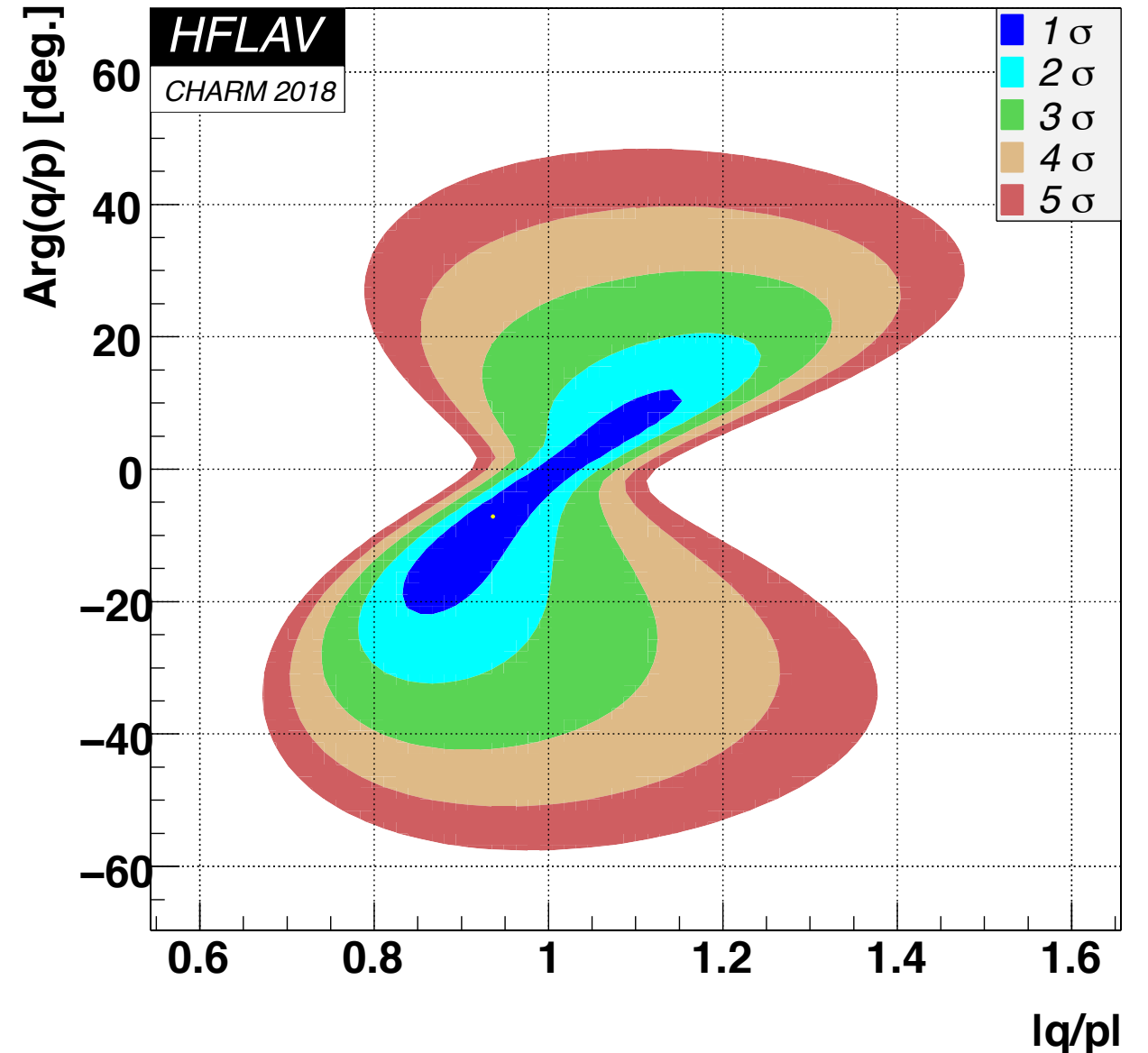
- CPT symmetry implies $|q|^2 + |p|^2 = 1$
- If CP is conserved, q and p are real, and $P_{1,2}$ are CP eigenstates

Pre-Moriond status



$$x = (3.6^{+2.1}_{-1.6}) \times 10^{-3}$$

$$y = (6.7^{+0.6}_{-1.3}) \times 10^{-3}$$



$$|q/p| = 0.94^{+0.17}_{-0.07}$$

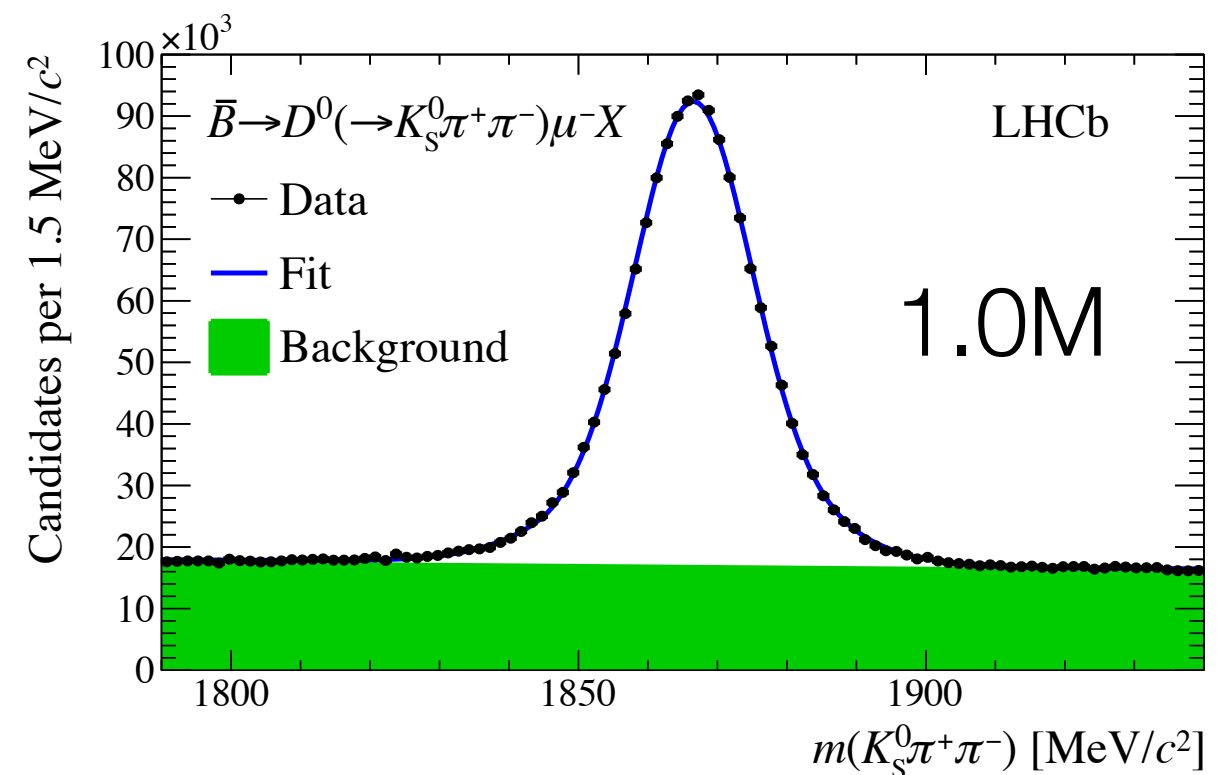
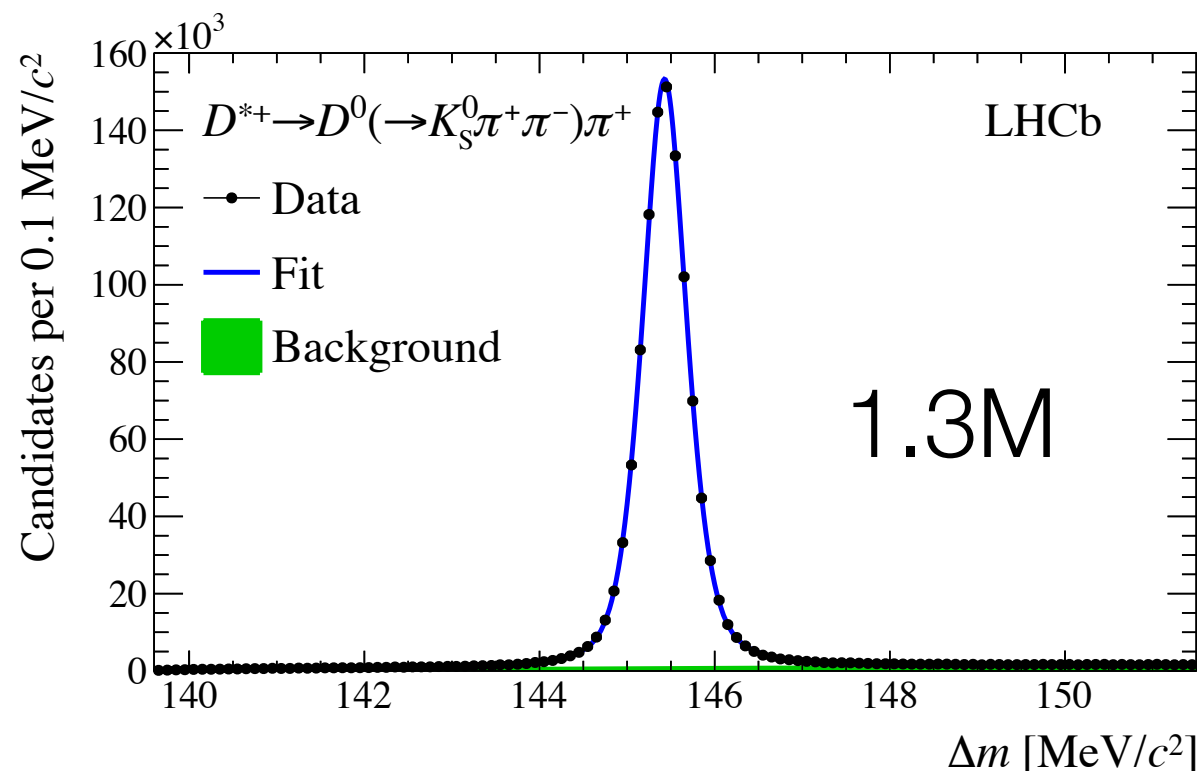
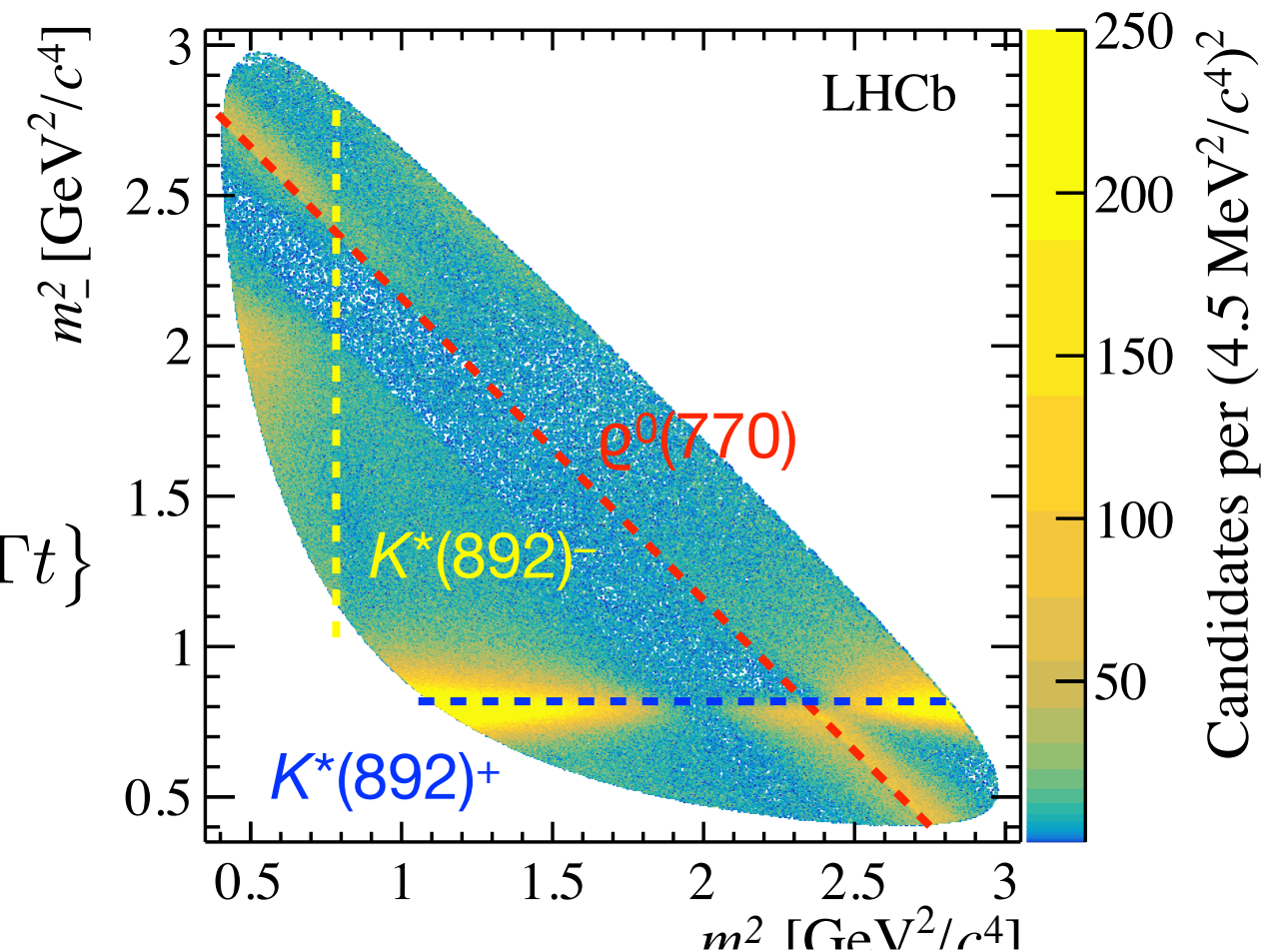
$$\phi = (-7.2^{+15}_{-10})^\circ$$

Mixing with $D^0 \rightarrow K_S \pi^+ \pi^-$

- Multiple interfering amplitudes enhance the sensitivity to mixing

$$\mathcal{P}_{D^0} \propto e^{-\Gamma t} \left\{ |\mathcal{A}_{D^0}|^2 - \text{Re}[\mathcal{A}_{D^0}^* \mathcal{A}_{\bar{D}^0} (y + ix)] \Gamma t \right\}$$

- New LHCb analysis with Run 1 data [\[arXiv:1903.03074\]](https://arxiv.org/abs/1903.03074)



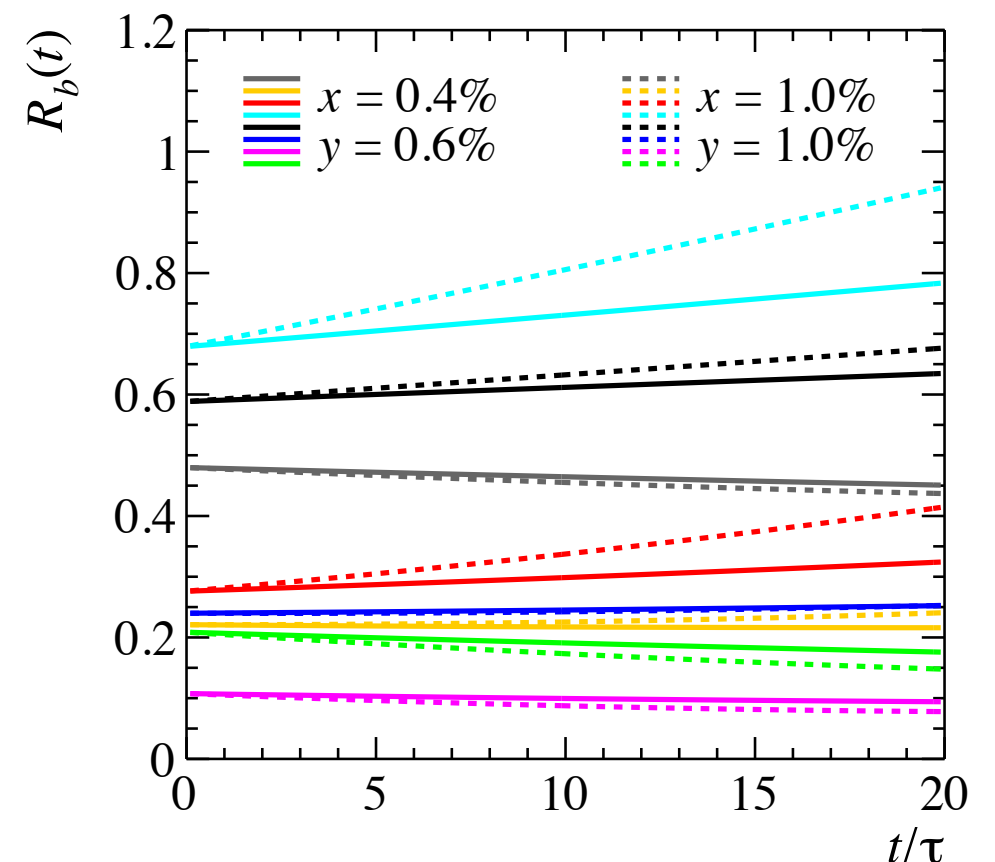
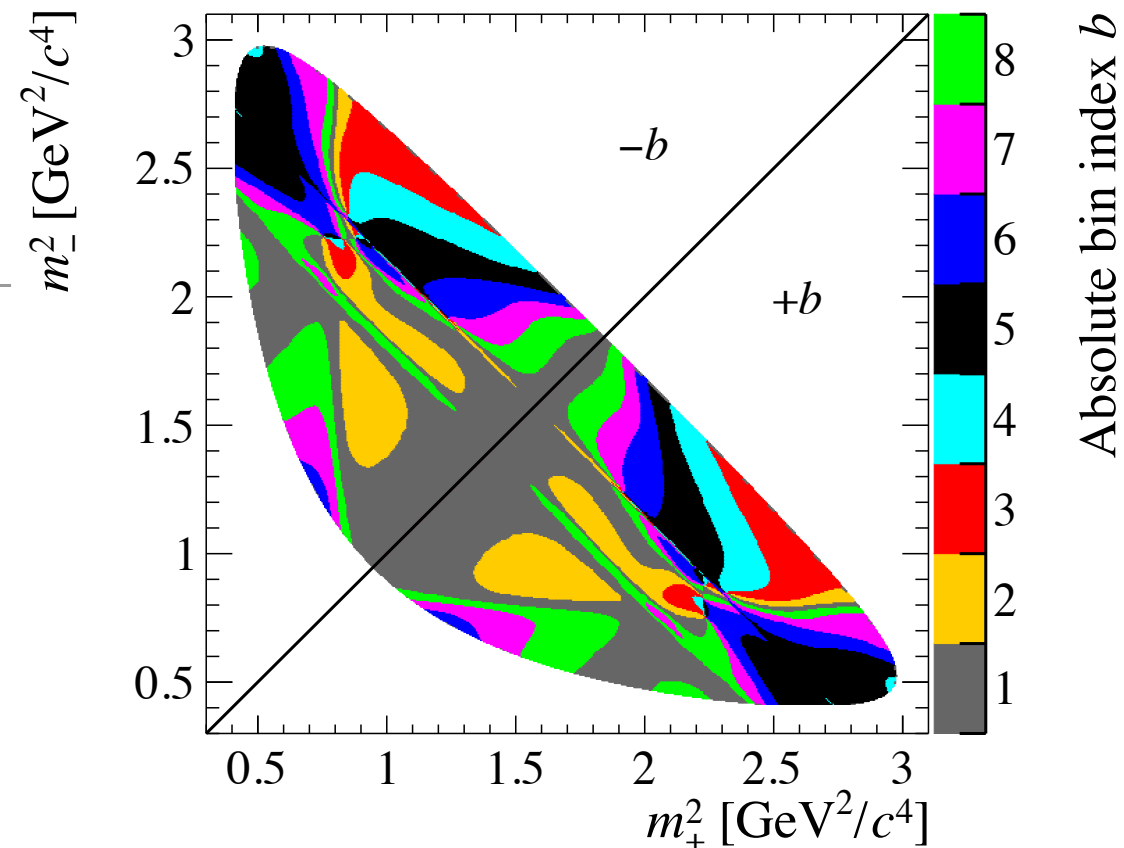
Bin-flip method

- Bin the Dalitz-plot to avoid amplitude model
- Ratio of events in bin $-b$ to events in bin b to suppress effects due to non-uniform efficiency variations

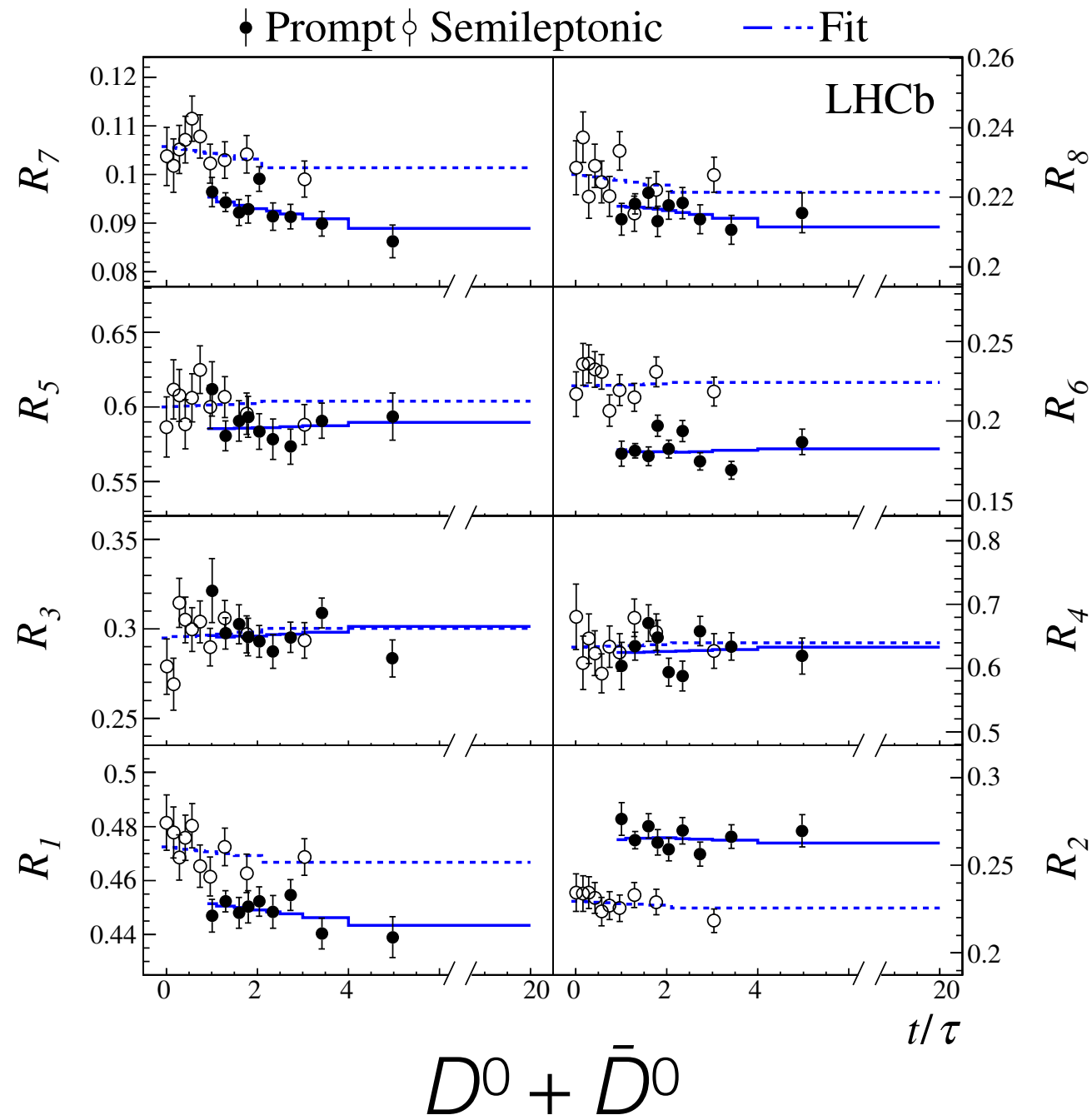
$$R_b \approx r_b - \sqrt{r_b} \left[(1 - r_b) c_b y - (1 + r_b) s_b x \right] \Gamma t$$

- Constrain hadronic parameters (c_b , s_b) using measurements with quantum-correlated $D^0 \bar{D}^0$ pairs, *i.e.* at CLEO
- Model-independent and completely data driven
- Comes with the price of degraded sensitivity to mixing effects from CP -even/odd amplitudes (*i.e.*, to y)

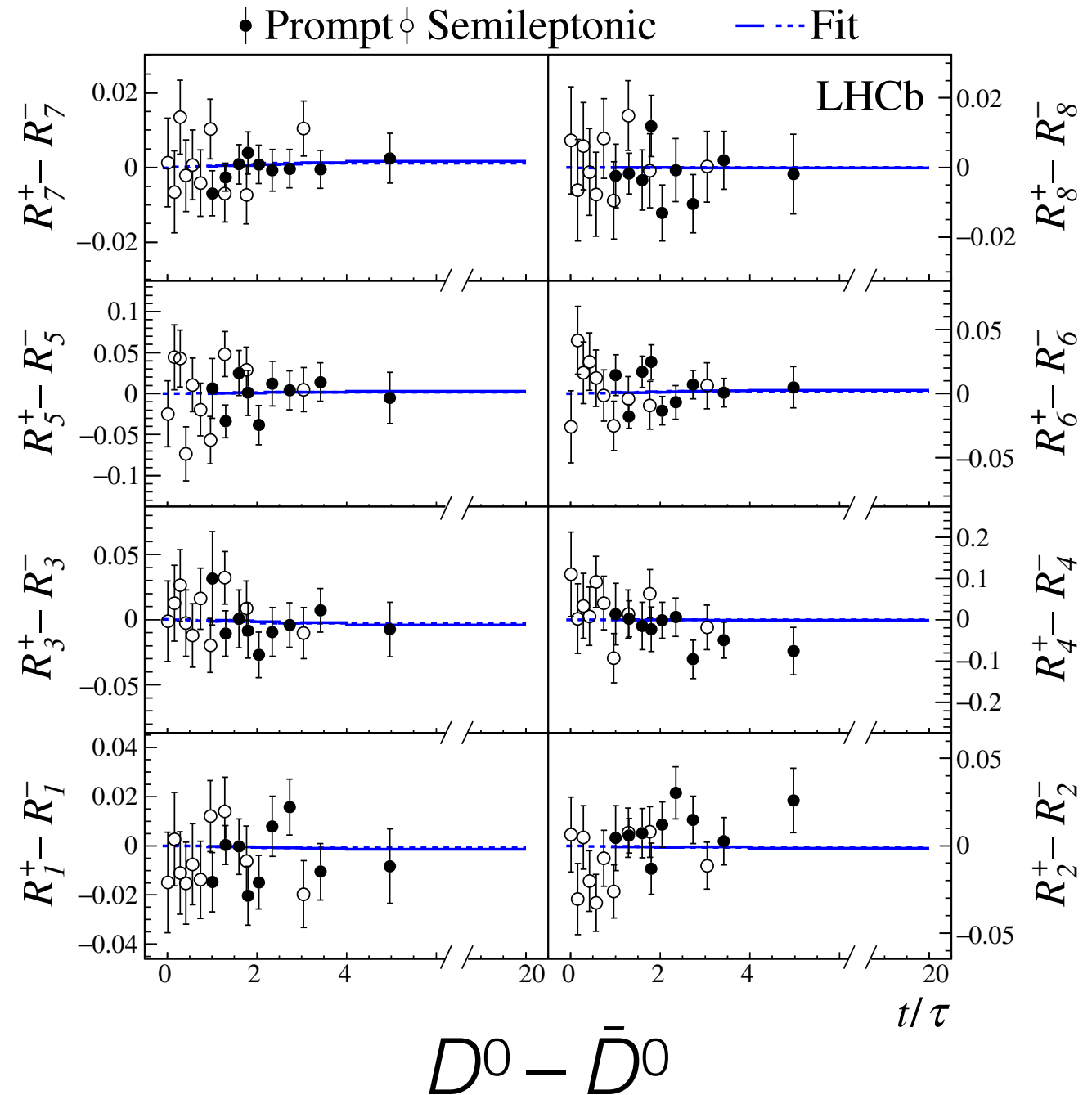
[PRD 99 (2019) 012007]



Results

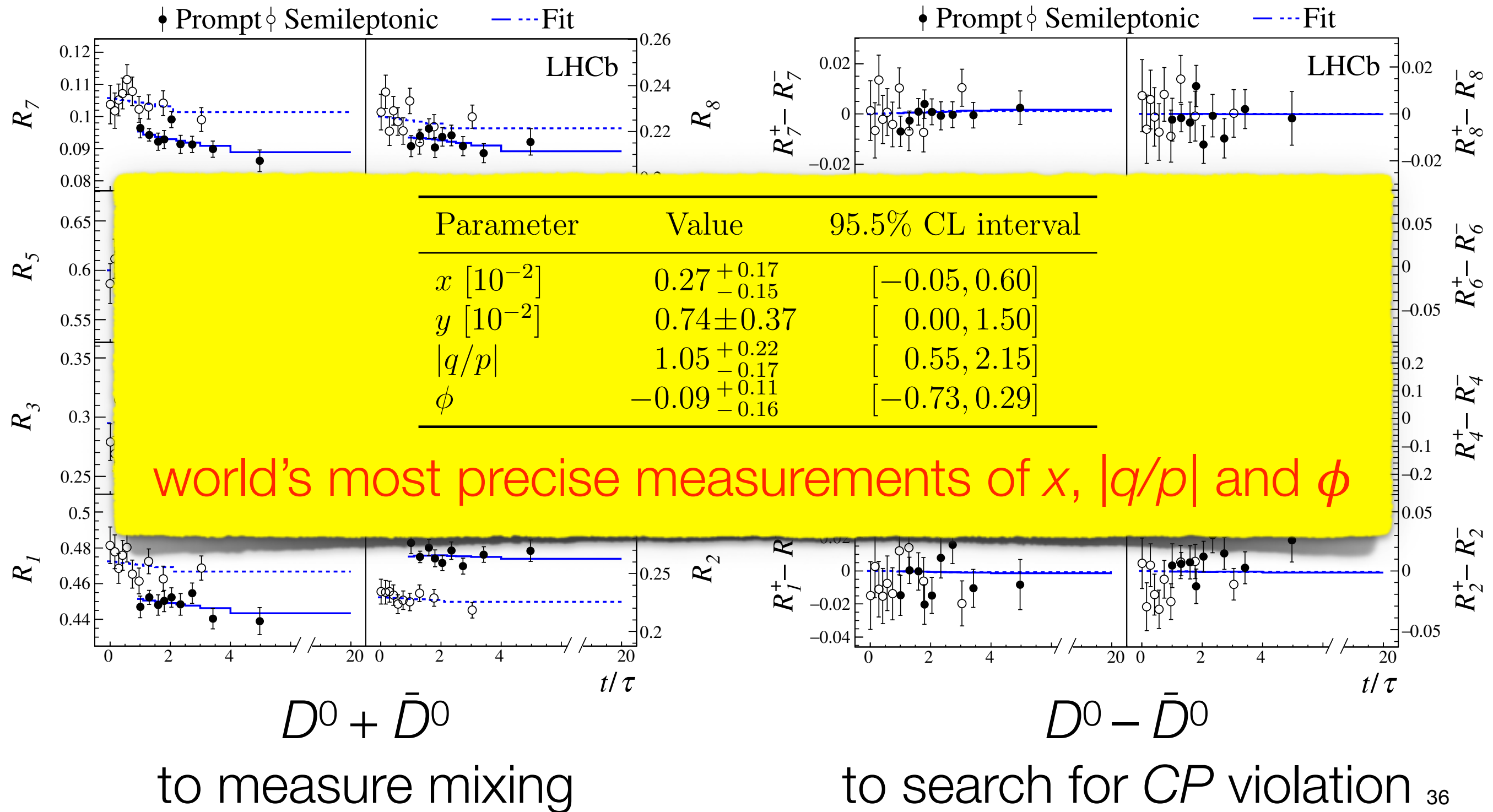


to measure mixing



to search for CP violation ₃₆

Results



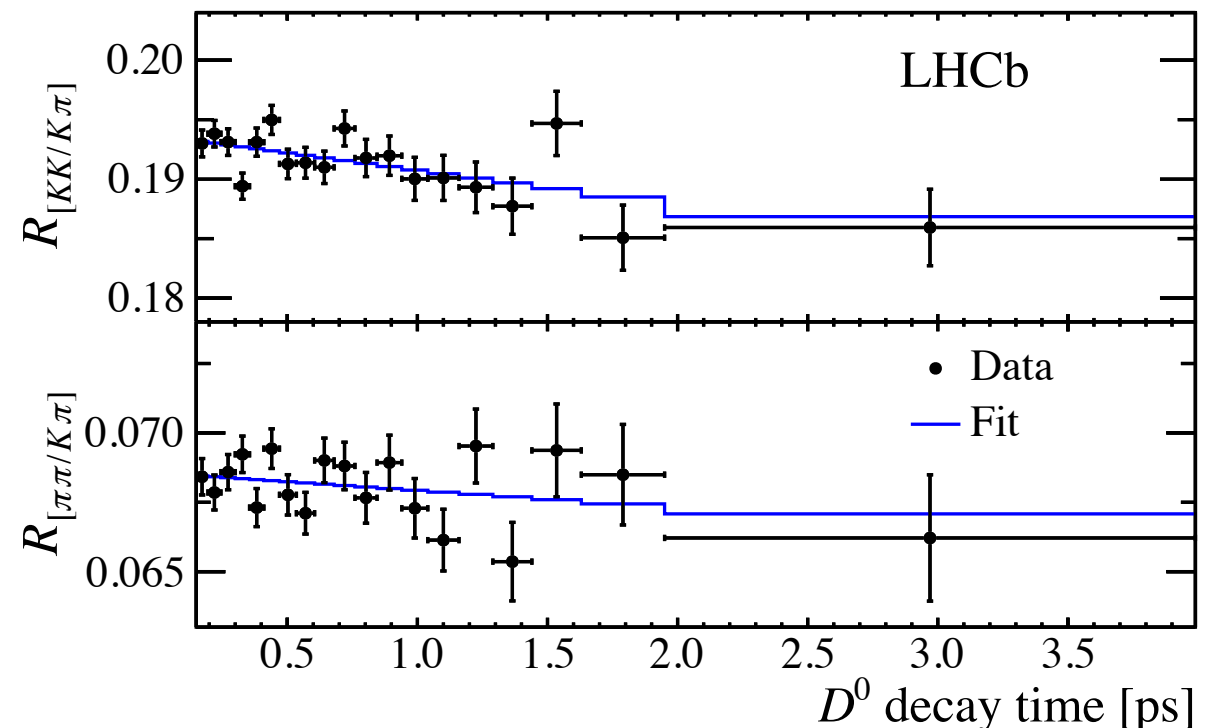
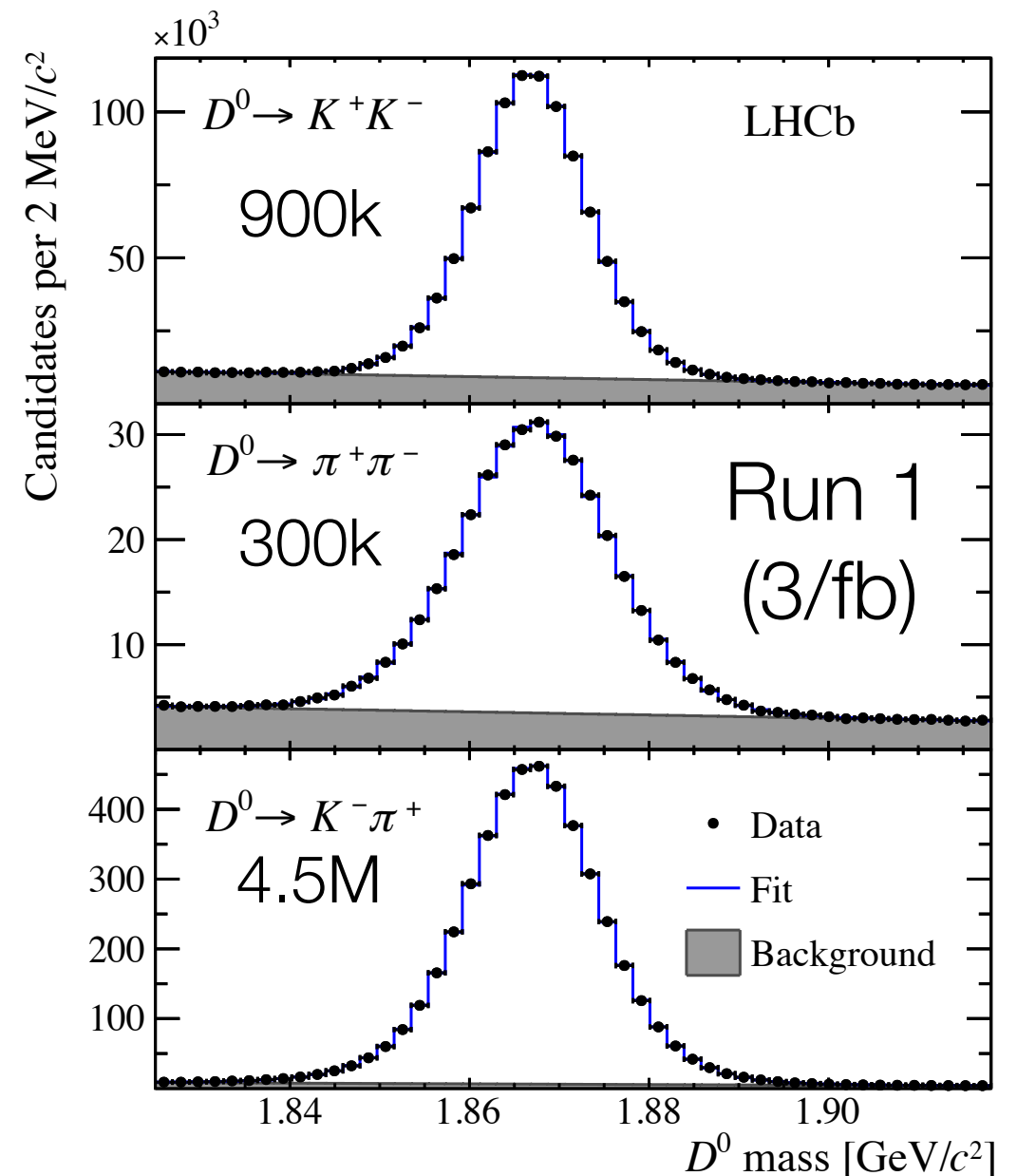
Measurement of y_{CP}

- Effective lifetime of decays to CP -even final state (relative to CP -mixed)

$$y_{CP} \equiv \frac{\tau(K^- \pi^+)}{\tau(h^+ h^-)} - 1$$

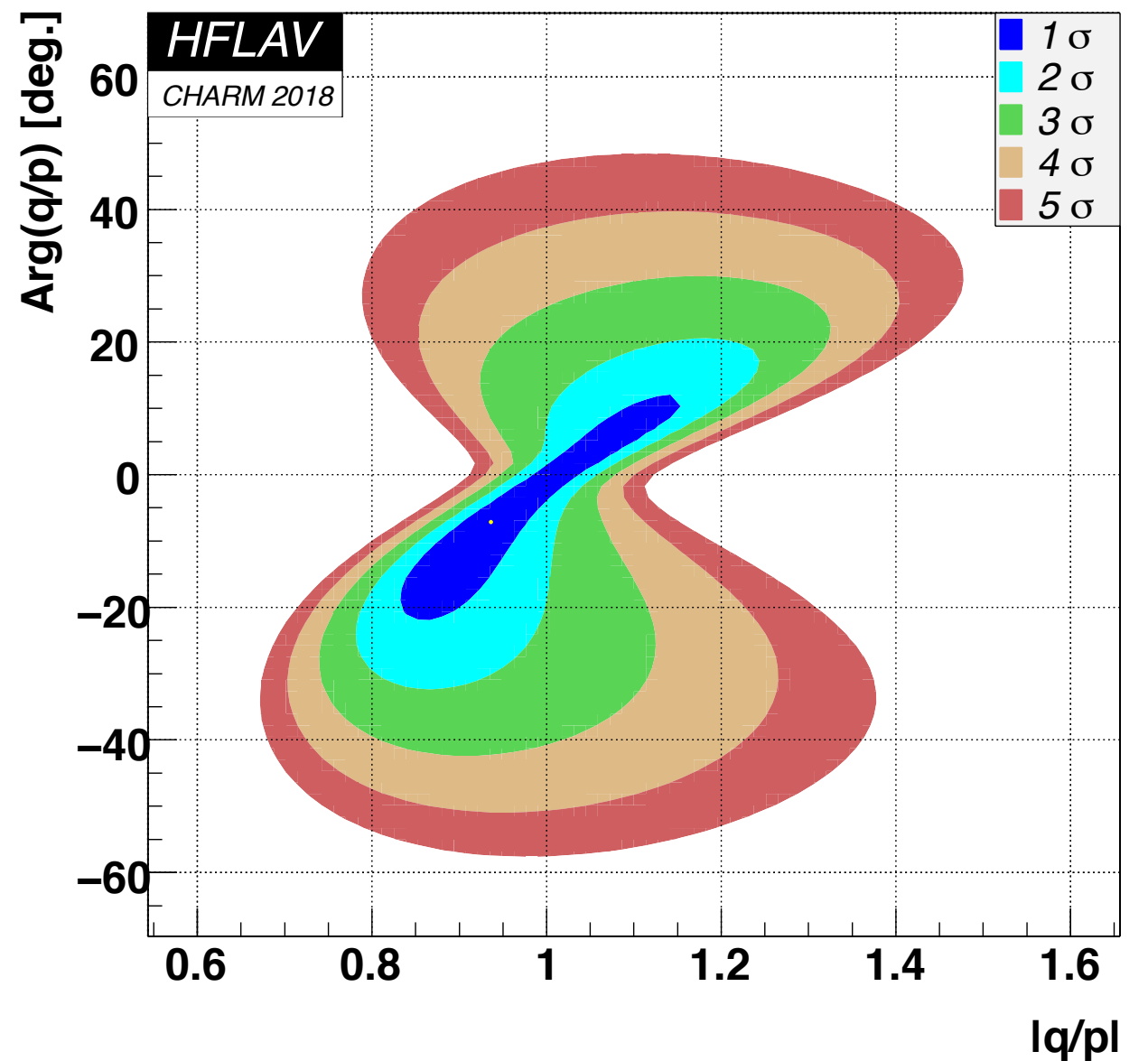
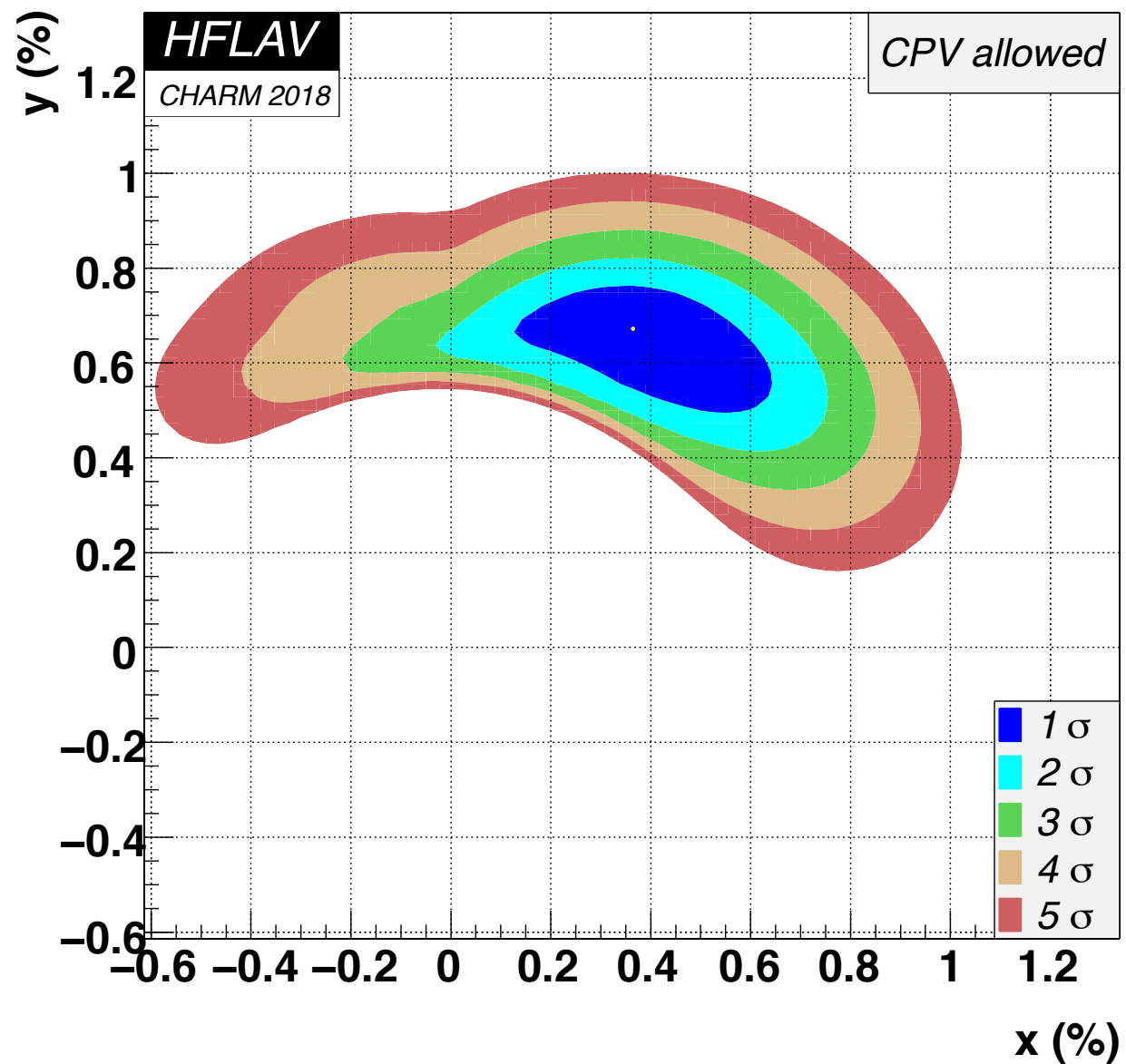
- Equal to y in case of CP symmetry
- Use μ -tagged decays to reduce biases on D^0 decay time
- Result with Run 1 data is as precise as the world average

$$y_{CP} = (5.7 \pm 1.3 \pm 0.9) \times 10^{-3}$$

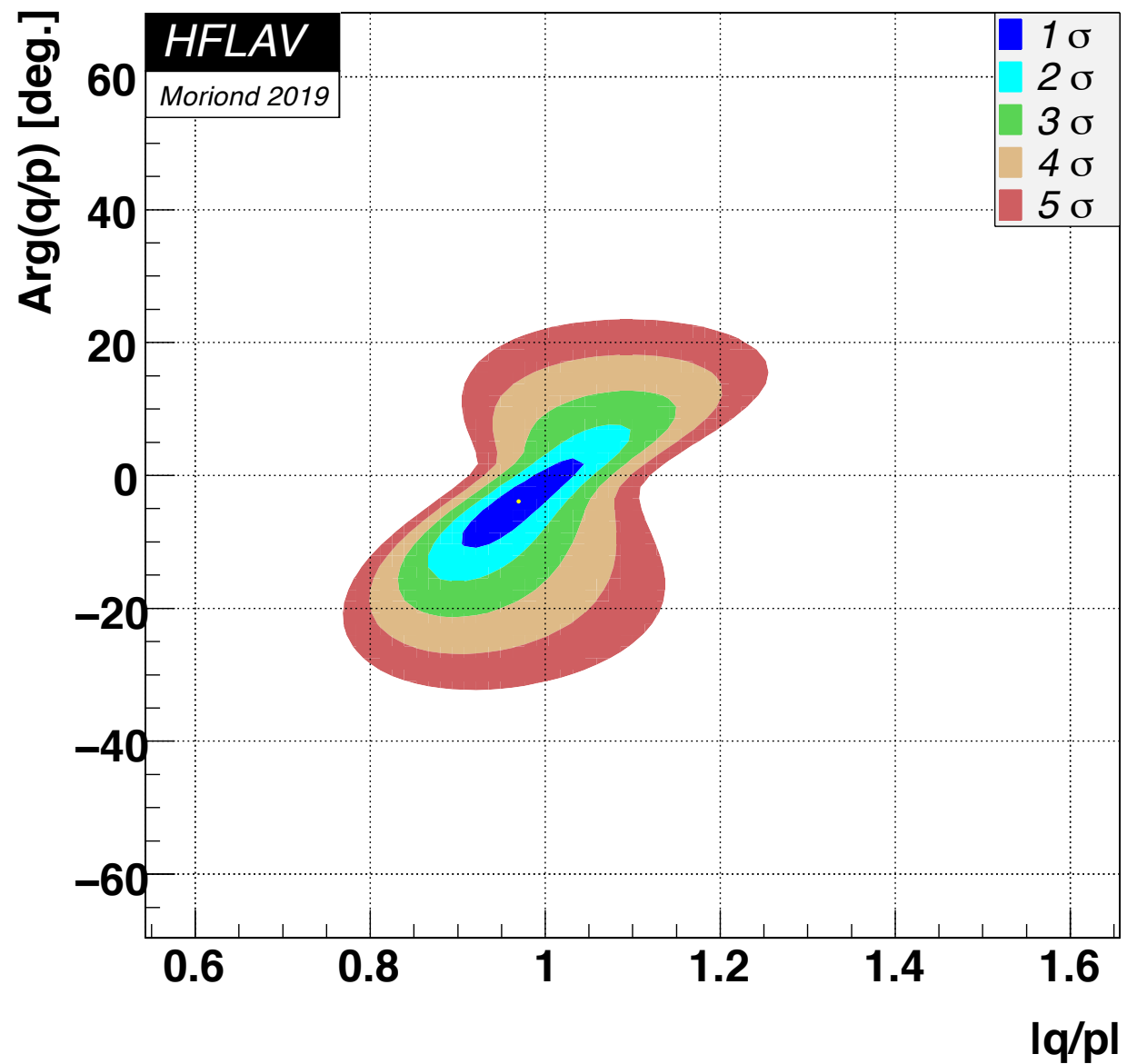
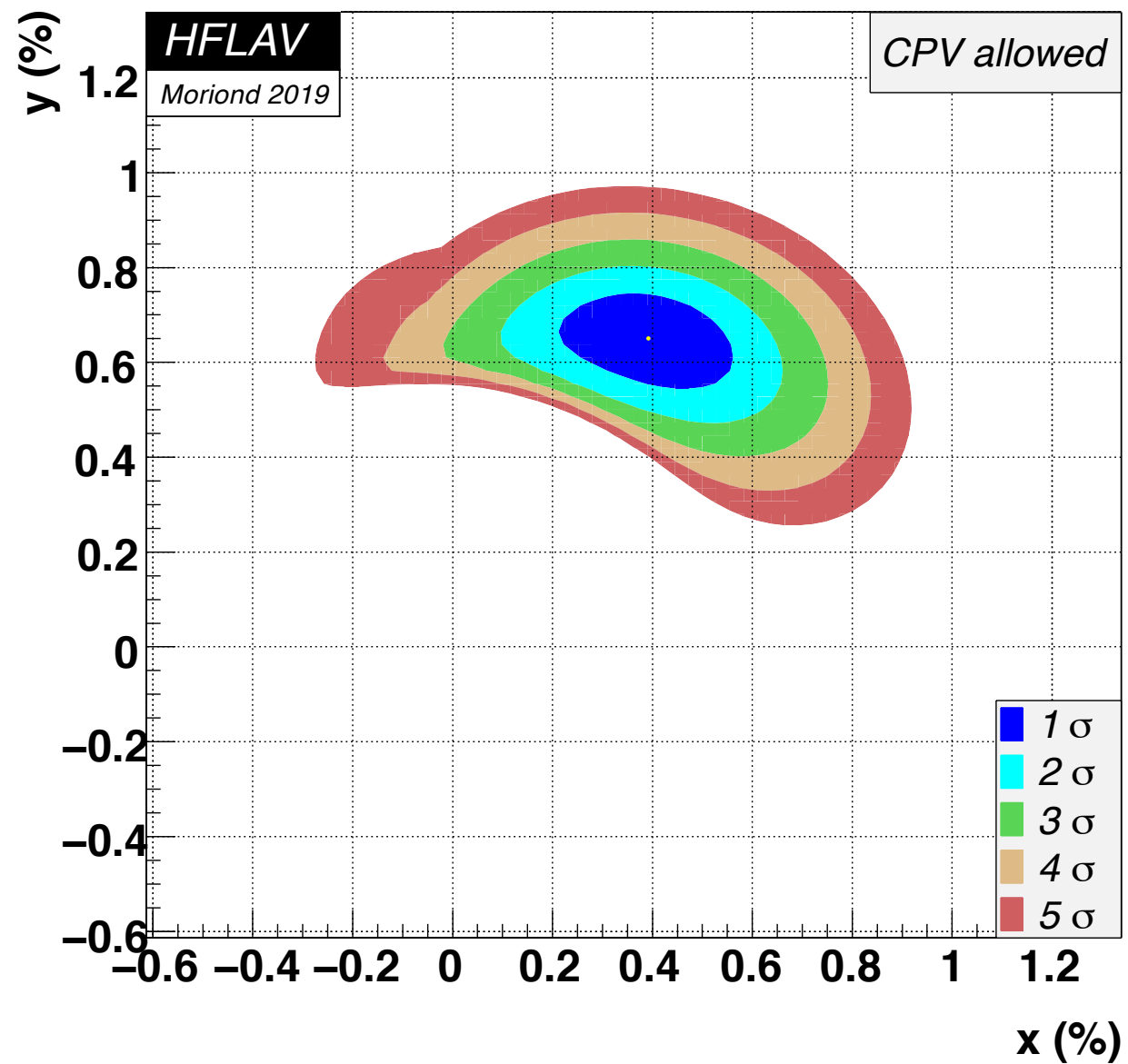


[PRL 122 (2019) 011802]

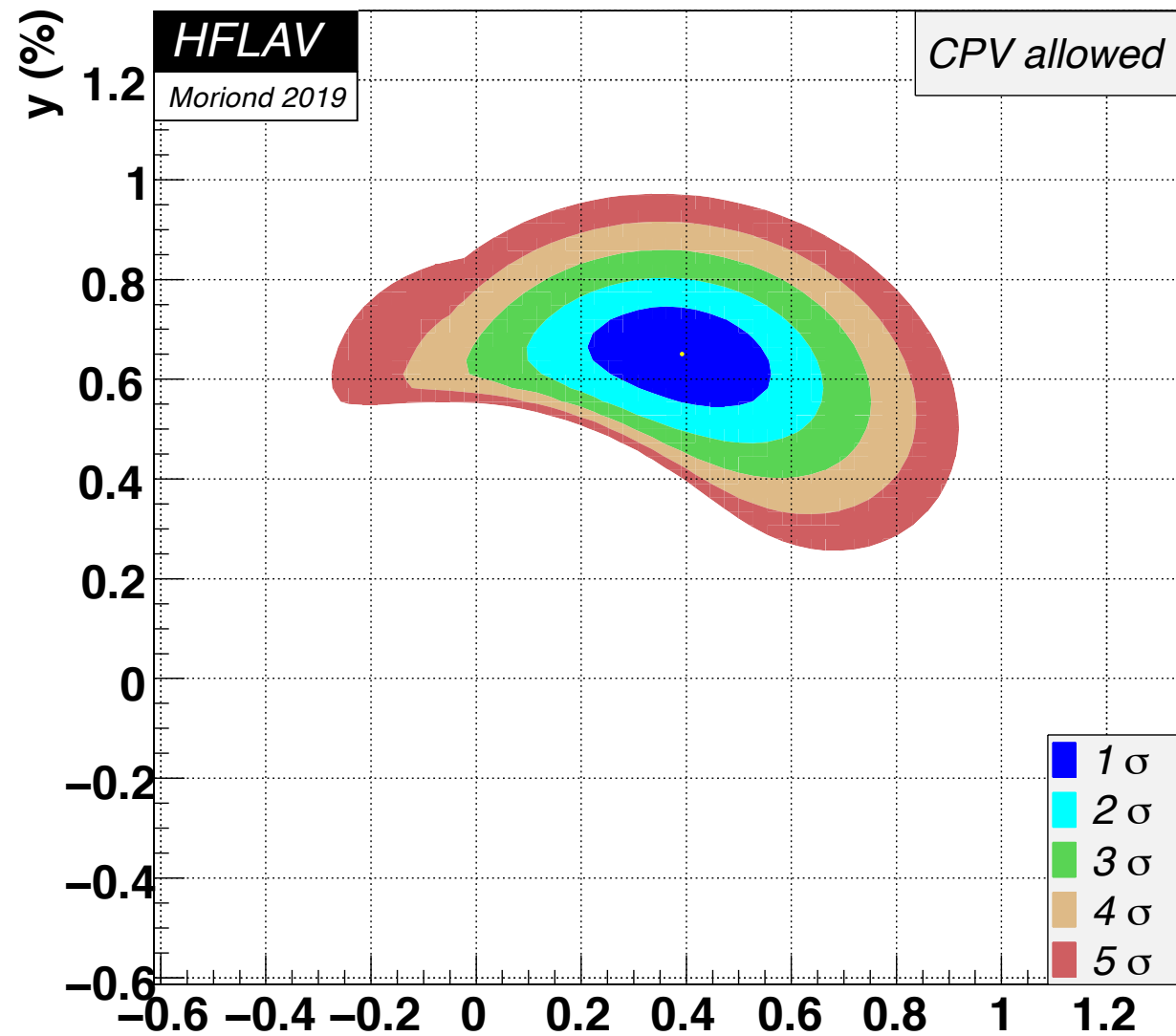
Impact on world average



Impact on world average

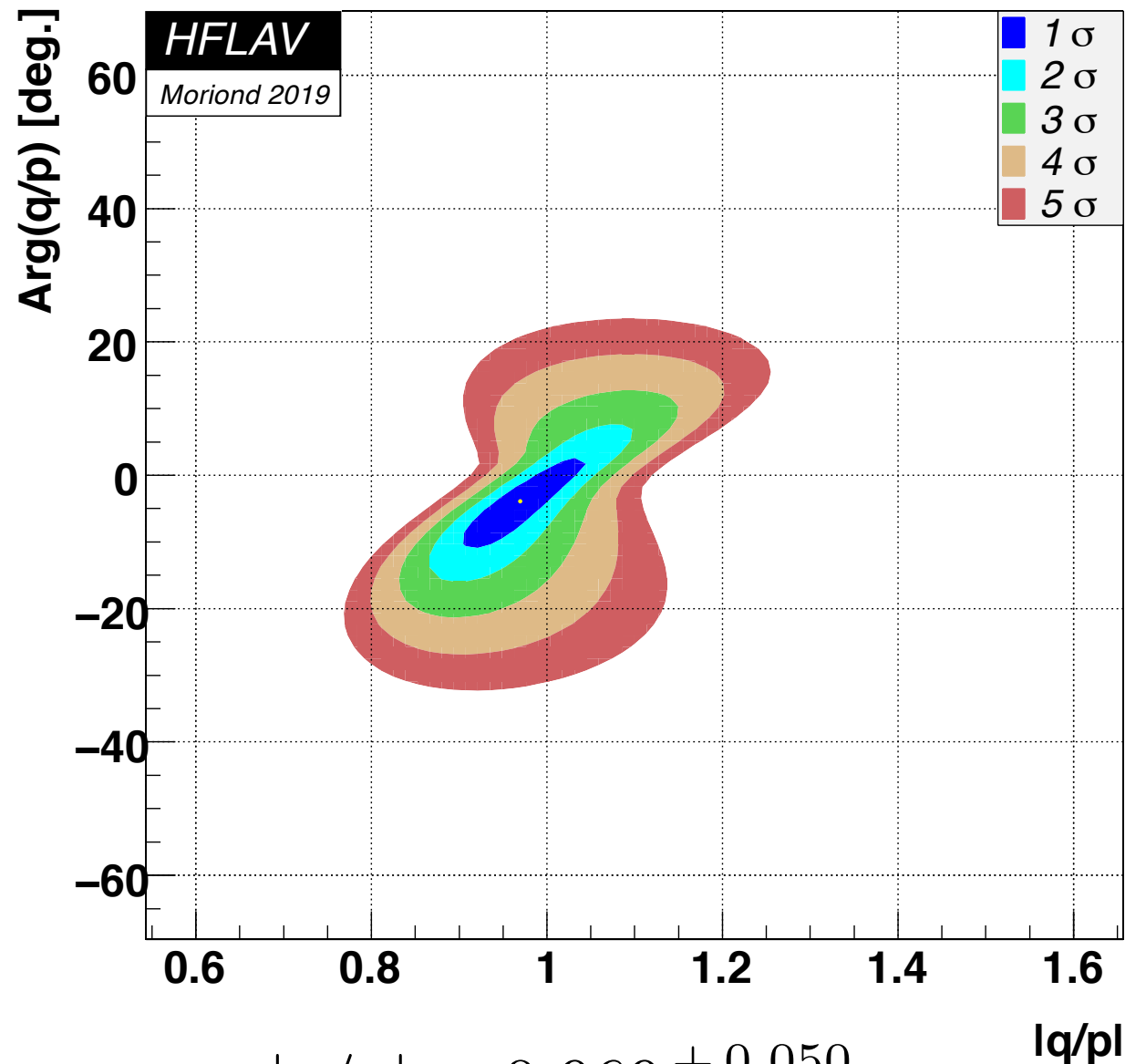


Impact on world average



$$x = (3.9 \pm_{-1.2}^{+1.1}) \times 10^{-3} \text{ } x \text{ } (\%)$$

$$y = (6.51 \pm_{-0.69}^{+0.63}) \times 10^{-3}$$



$$|q/p| = 0.969 \pm_{-0.045}^{+0.050}$$

$$\phi = (-3.9 \pm_{-4.6}^{+4.5})^\circ$$

first evidence for nonzero (positive) value of x

Summary

- Observed (direct) CP violation in charm decays

$$\Delta A_{CP} = (-15.4 \pm 2.9) \times 10^{-4}$$

- Measured asymmetry seems consistent with standard model, although predictions suffer from large uncertainties due to strong-interaction effects
- Additional searches for CP violation in different decay modes can help to clarify the picture
- First evidence for nonzero difference between neutral charm-meson eigenstate. Yet no signs for mixing-induced CP violation

Any 1.1 Questions