

# Mixing induced $CP$ asymmetry in semileptonic $B$ -meson decays at $BABAR$

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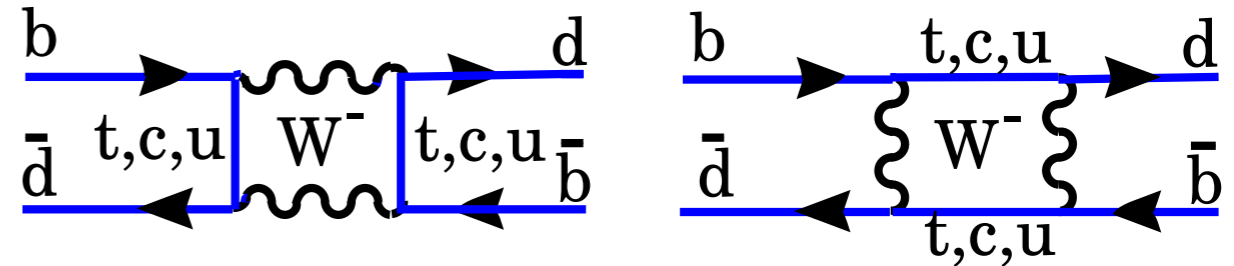
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# Neutral meson mixing

- In neutral meson systems, particle/antiparticle can couple to each other through weak interaction:

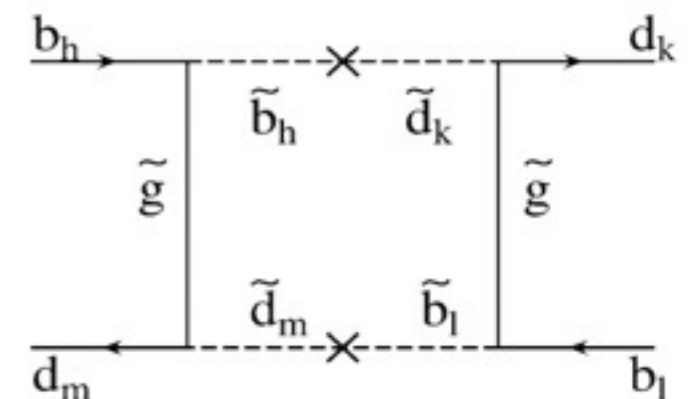
$$B_{d,s}^0 \leftrightarrow \bar{B}_{d,s}^0, D^0 \leftrightarrow \bar{D}^0, K^0 \leftrightarrow \bar{K}^0$$



- Equation of a decaying particle :  $B(t) = \exp[-imt - (\Gamma/2)t]$

$$i \frac{d}{dt} \begin{pmatrix} |B(t)\rangle \\ |\bar{B}(t)\rangle \end{pmatrix} = \left( \hat{M} - \frac{i}{2} \hat{\Gamma} \right) \begin{pmatrix} |B(t)\rangle \\ |\bar{B}(t)\rangle \end{pmatrix}$$

- Off-diagonal elements  $M_{12}, \Gamma_{12}$  in the matrix are non-trivial (and complex).
- Measurements of various aspects of mixing test the theory of weak interaction.
  - Mixing rate, decay rate difference,  $CP$  violation.
- New physics in the loops could alter these observables.



# Mixing induced $CP$ violation

- The eigenstates of the neutral  $B$  system are linear combinations of flavor states  $|B^0\rangle$  and  $|\bar{B}^0\rangle$

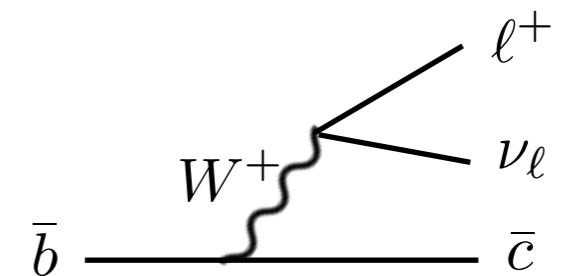
$$|B_{L/H}\rangle = \frac{1}{\sqrt{p^2 + q^2}} (p|B^0\rangle \pm q|\bar{B}^0\rangle)$$

- For flavor specific final states:  $B^0 \rightarrow f, \bar{B}^0 \rightarrow \bar{f}; B^0 \not\rightarrow \bar{f}, \bar{B}^0 \not\rightarrow f$

◆ E.g., semileptonic decays:  $B^0 \rightarrow X\ell^+\nu_\ell$

◆ Assume no direct  $CPV$ :  $|\langle f|B^0\rangle| = |\langle \bar{f}|\bar{B}^0\rangle|$

◆  $CP$  violation due to mixing:



$$A_{sl} = \frac{\Gamma(\bar{B}^0(t=0) \rightarrow f) - \Gamma(B^0(t=0) \rightarrow \bar{f})}{\Gamma(\bar{B}^0(t=0) \rightarrow f) + \Gamma(B^0(t=0) \rightarrow \bar{f})} = \frac{1 - |q/p|^4}{1 + |q/p|^4} \simeq -2 \left( \left| \frac{q}{p} \right| - 1 \right) \simeq \text{Im} \frac{\Gamma_{12}}{M_{12}} \simeq \frac{\Delta\Gamma}{\Delta M} \tan \phi$$

$\bar{B}^0 \rightarrow B^0$ , then decay to the “wrong sign” final state  $f$ .

$CPV$  if  $|q/p| \neq 1$

$$\phi = \arg \left( -\frac{M_{12}}{\Gamma_{12}} \right)$$



# Search/constrain new physics

- Standard Model prediction is small:

$$A_{sl}^s = (1.9 \pm 0.3) \times 10^{-5}$$

$$A_{sl}^d = -(4.1 \pm 0.6) \times 10^{-4}$$

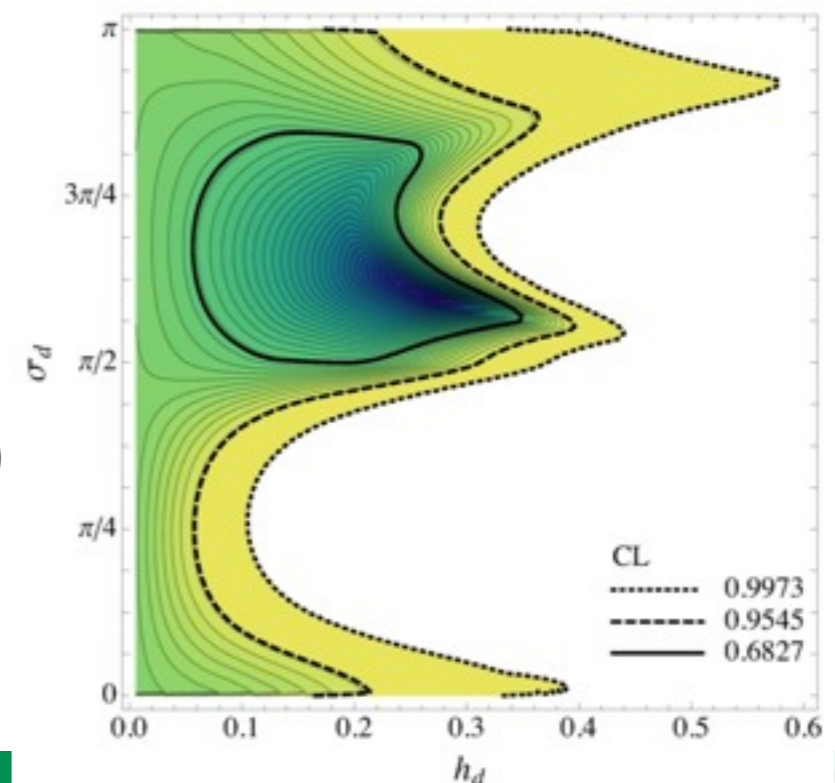
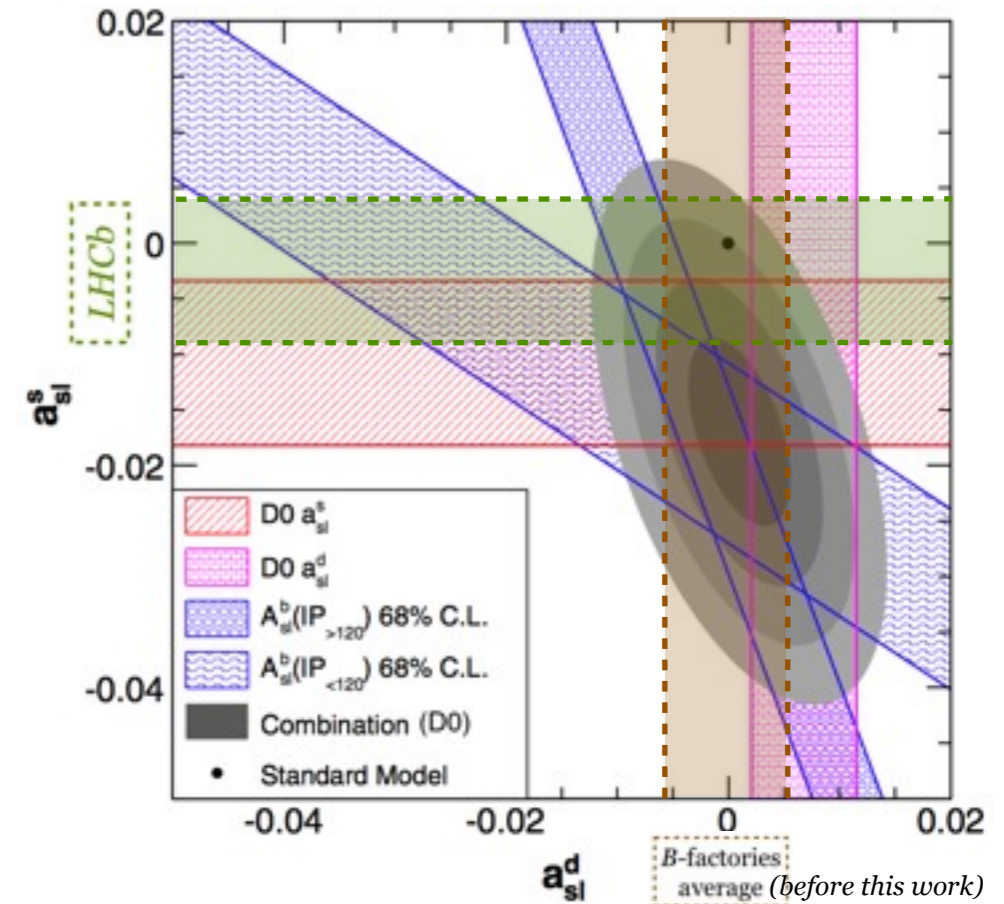
[Lenz, Nierste, arXiv:1102.4274 (2011)]

- ◆ well below current experimental sensitivity ( $\mathcal{O}(10^{-3})$ ).
- ◆ observation of non-zero *CPV* would indicate new physics.
- ◆ Some tension observed in D0 di-muon analysis

- Allowed new physics (model independent) parameter space is still sizable.

$$M_{12}^{d,s} = (M_{12}^{d,s})^{\text{SM}} (1 + h_{d,s} e^{2i\sigma_{d,s}})$$

[Ligeti, et al. PRL 105, 131601 (2010)]

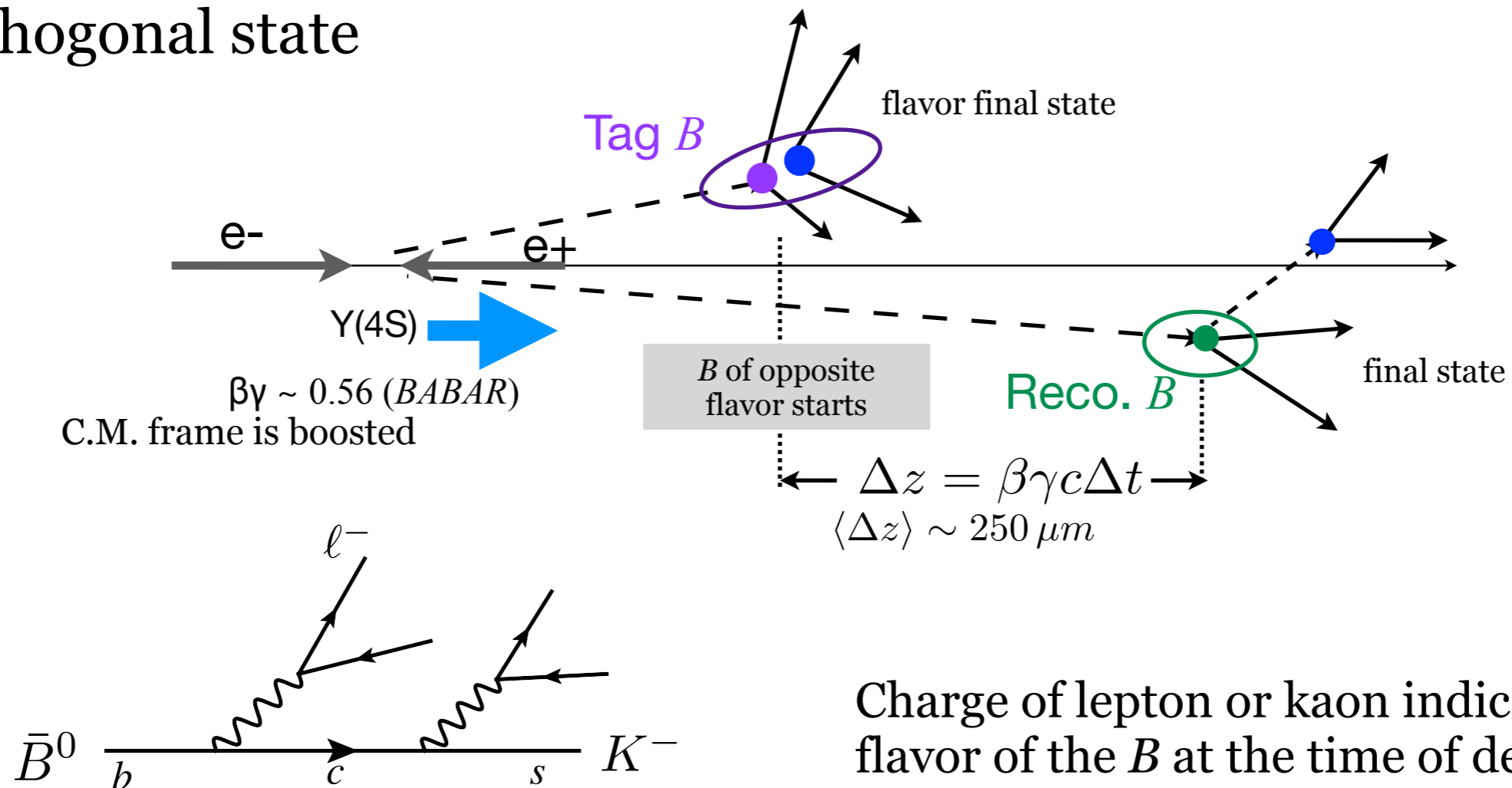


# Experiment

- In  $B$  factories  $\Upsilon(4S)$  are produced and decay to a pair of  $B$  mesons in a coherent  $L=1$  antisymmetric quantum state

$$|i\rangle = 1/\sqrt{2}[B^0(t_1)\bar{B}^0(t_2) - \bar{B}^0(t_1)B^0(t_2)]$$

- Once one  $B$  decays to a basis state, the other projects to the orthogonal state



# Experiment

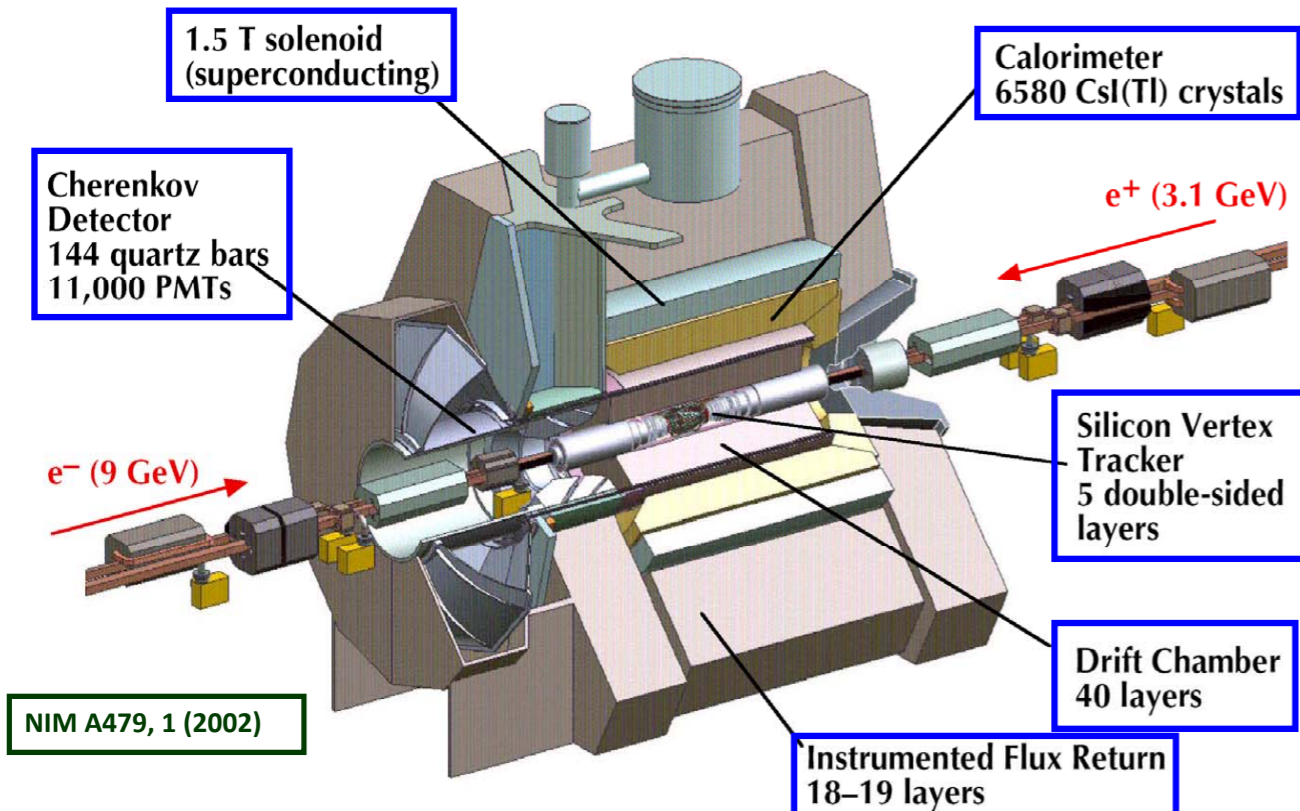
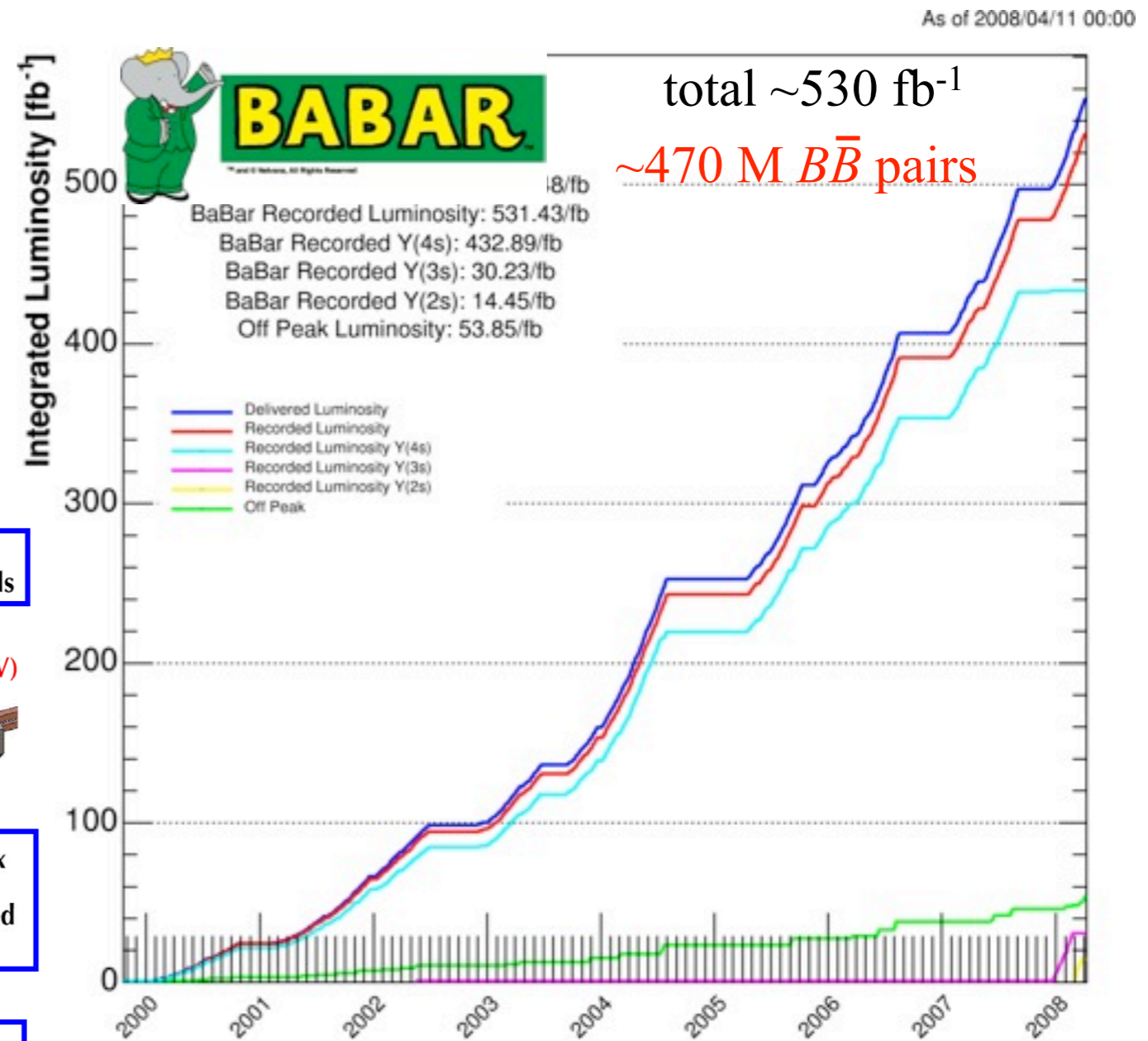
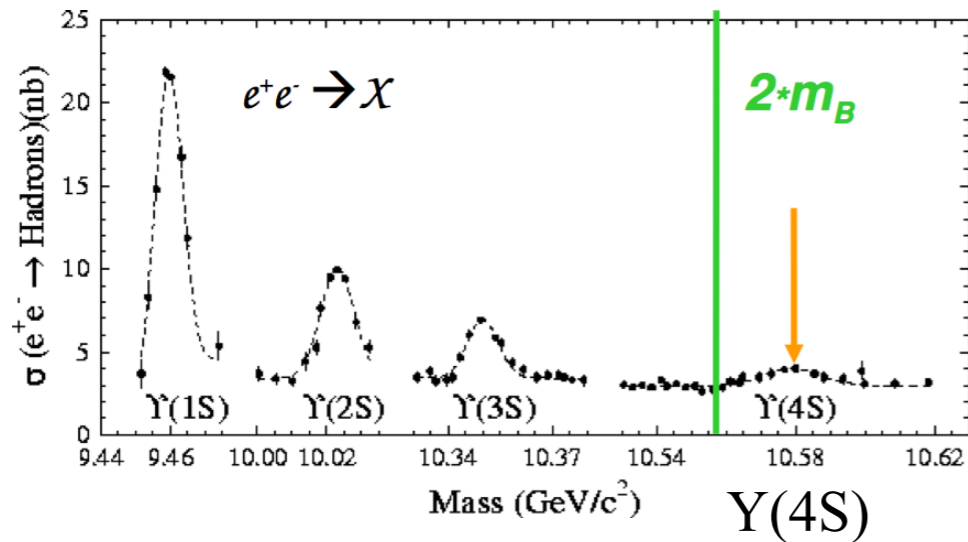
- Reconstruct  $B^0 \rightarrow \ell^+ \nu_\ell D^{*-}$  using a partial reconstruction technique.
- Find the flavor of the “tag  $B$ ” using charged kaons in the rest of the event.
- The “wrong-sign” decays (via mixing) are identified by same-sign lepton-kaon combination:

$$A_{\text{sl}} = \frac{N(B^0 B^0) - N(\bar{B}^0 \bar{B}^0)}{N(B^0 B^0) + N(\bar{B}^0 \bar{B}^0)} = \frac{N(\ell^+ K^+) - N(\ell^- K^-)}{N(\ell^+ K^+) + N(\ell^- K^-)}$$

- Partial reconstruction gives higher efficiency than exclusive reconstruction, and better purity than inclusive lepton reconstruction.
- Kaon tag has a higher efficiency than lepton tag.
  - ◆ Somewhat higher mistag probability
  - ◆ Need to take care of interference due to doubly-Cabibbo-suppressed decay

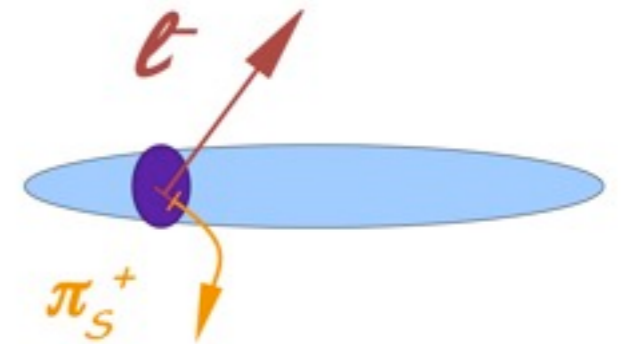


# Dataset and detector

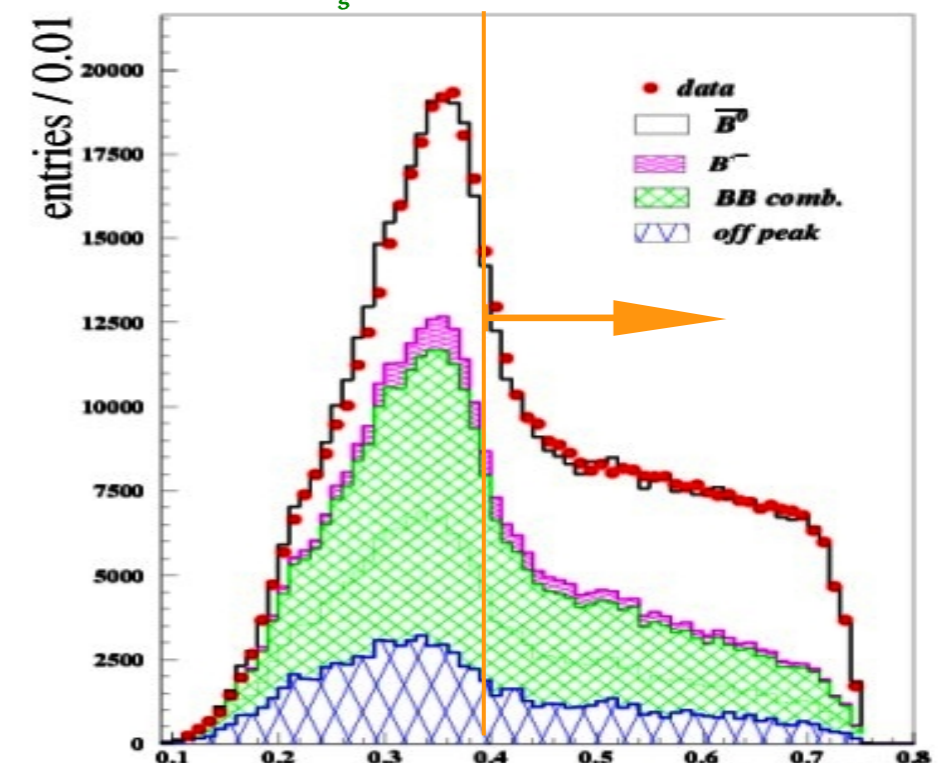


# $B \rightarrow \ell \nu D^*$ partial reconstruction

- Use only lepton ( $e, \mu$ ) and the low momentum  $\pi_s$  of the opposite sign from  $D^{*-} \rightarrow \pi_s^- \bar{D}^0$ , ignoring the remaining products from charm.
- Assume  $B$  at rest in  $\mathcal{V}(4S)$  frame.
- Approximate  $D^*$  direction using  $\pi_s$ :  $\vec{p}_{D^*} = f(\vec{\pi}_s)$
- Reconstruct  $B$  decay vertex using a beamspot constraint:
- Build a likelihood ratio function using lepton and pion momenta, and vertexing probability to reduce background.



$$\mathcal{L}(\mathcal{P}_{\pi_s}, \mathcal{P}_\ell, \text{Prob}(\text{Vtx}))$$





# $B \rightarrow \ell \nu D^*$ partial reconstruction

- Assume only the neutrino is not accounted for, and then calculate the missing neutrino mass

- Good signal

$$B^0 \rightarrow D^{*-} (X) \ell^+ \nu_\ell$$

$$B^0 \rightarrow D^{*-} (X) \tau^+ \nu_\tau; \tau^+ \rightarrow \ell^+ \nu_\ell \bar{\nu}_\tau$$

$$B^0 \rightarrow D^{*-} h^+; \text{ (misidentified)}$$

- Peaking but no  $B^0$  flavor info

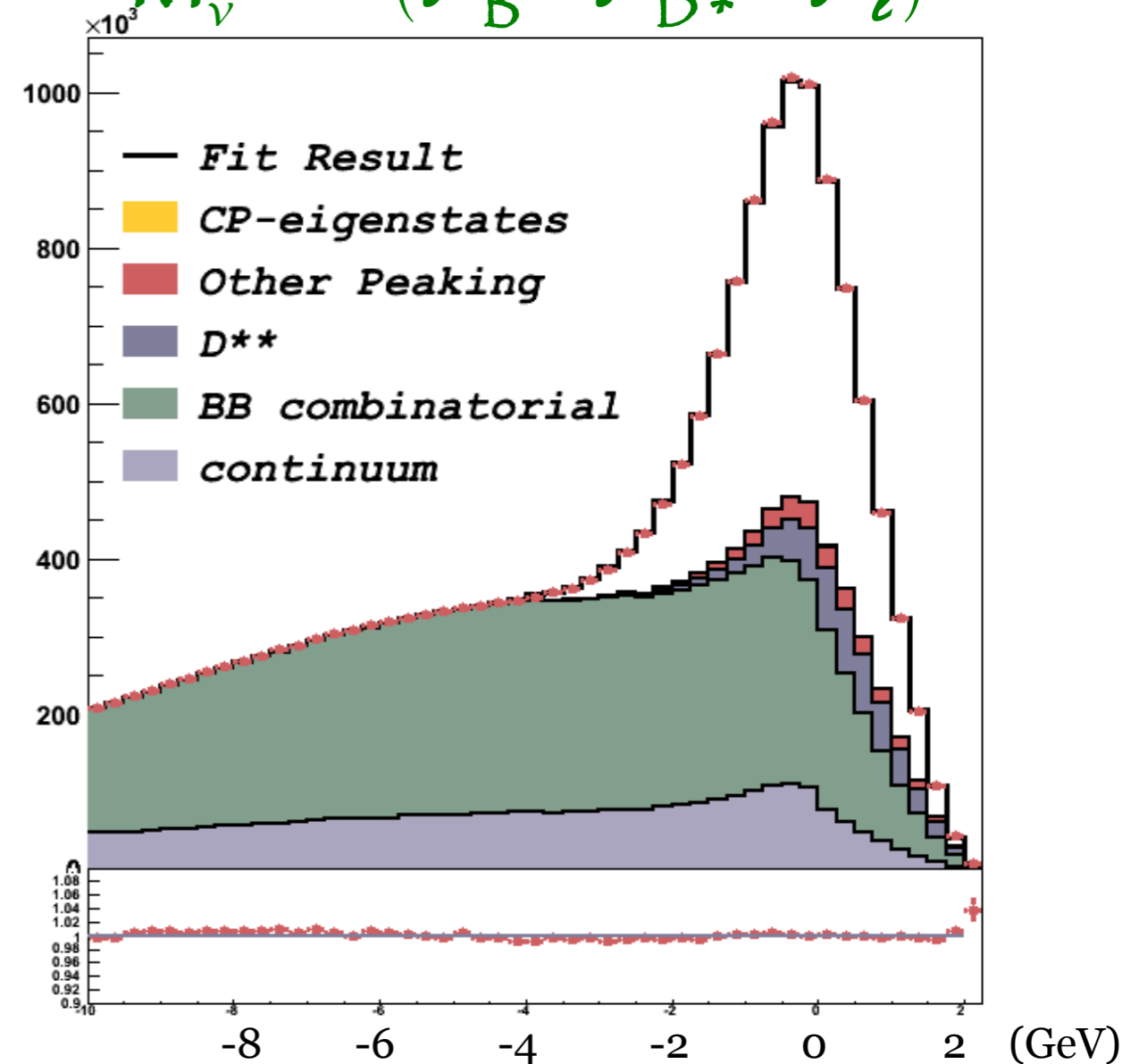
- ◆ double charm followed by charm semileptonic decay

- ◆ charged  $B$ .

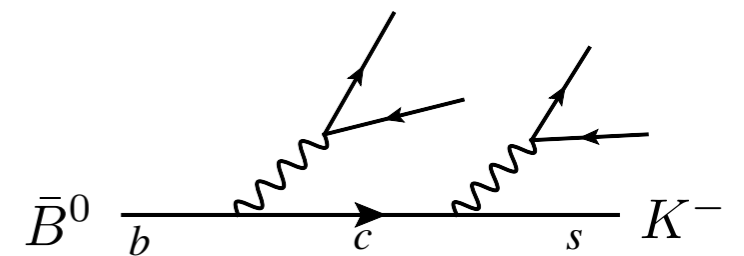
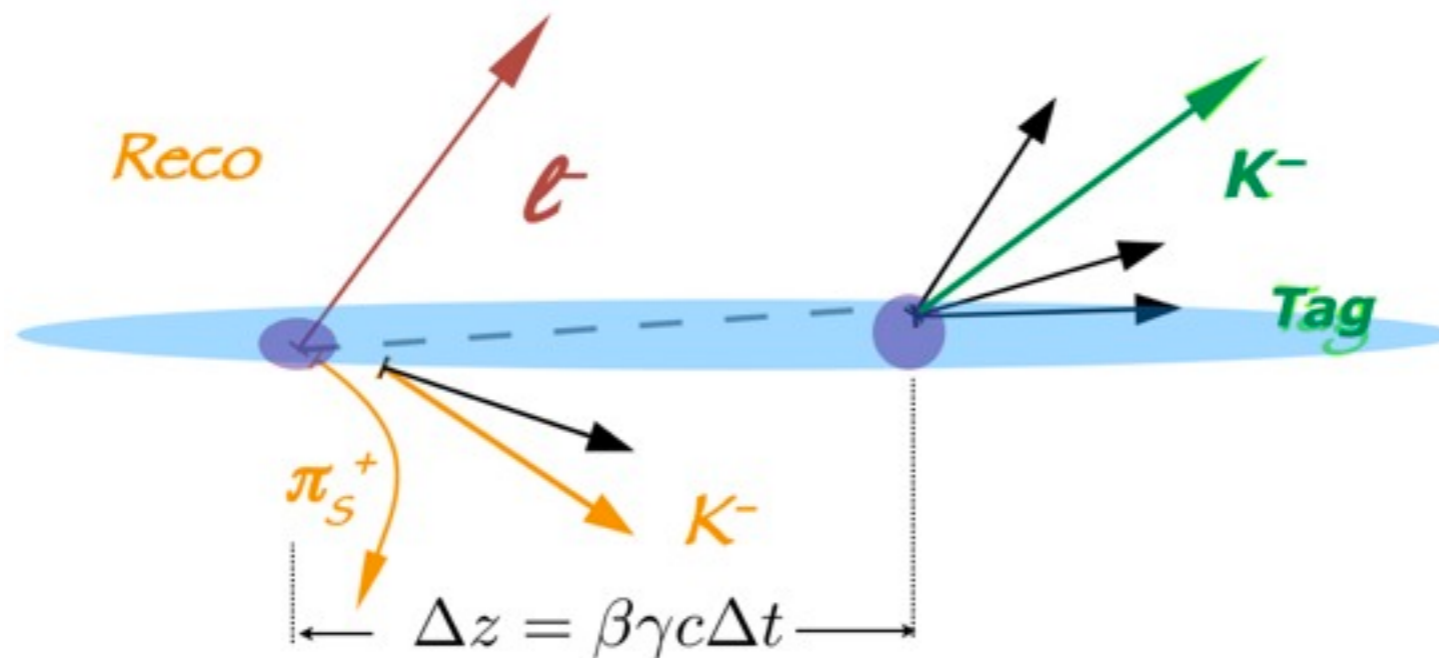
- Yield:

$(5945 \pm 7) \times 10^3$  peaking events

$$M_\nu^2 = (\mathcal{P}_B - \mathcal{P}_{D^*} - \mathcal{P}_\ell)^2$$



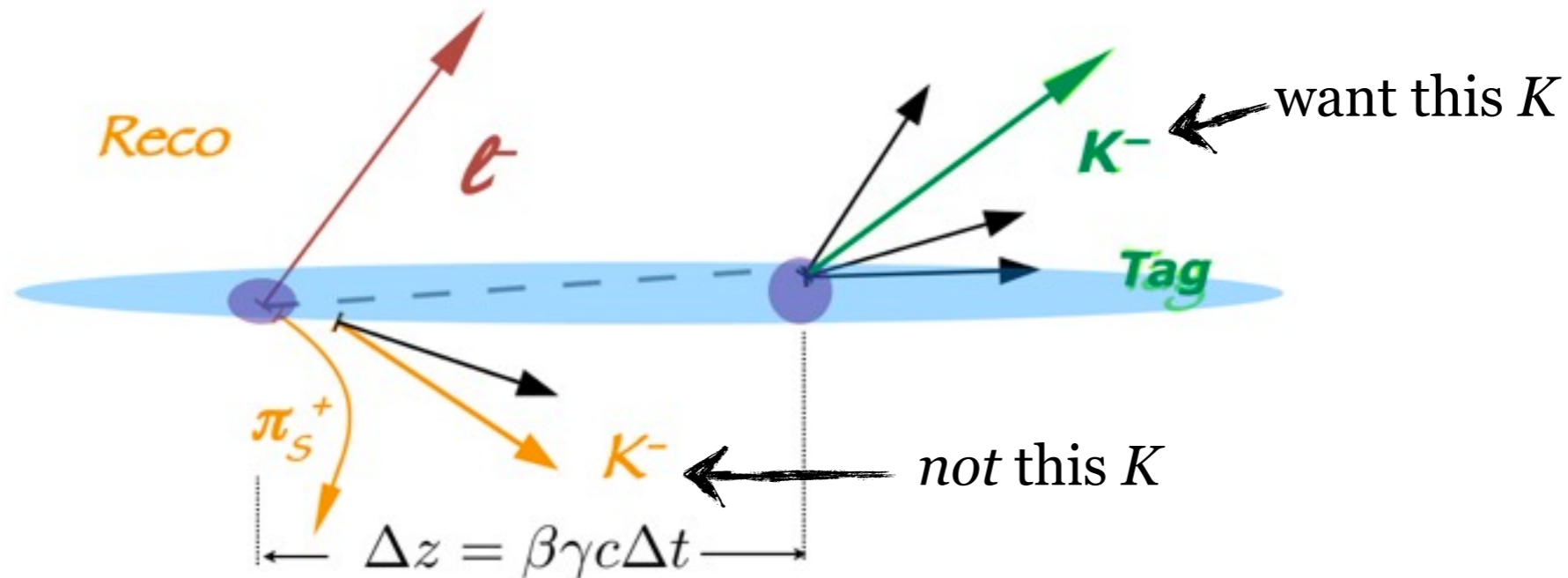
# Kaon tag



- Kaons are identified using ionization energy loss in the tracking devices and Cherenkov angles:  $\sim 85\%$  efficiency,  $\sim 3\%$  pion misidentification.
- Tag- $B$  vertex is identified by intersecting kaon track with the beamspot.
- Define  $\Delta t = \frac{z_{\text{reco.}} - z_{\text{tag}}}{\gamma\beta c}$ , which should follow  $B$  decay distribution (plus some smearing from charm lifetime) for signal events.

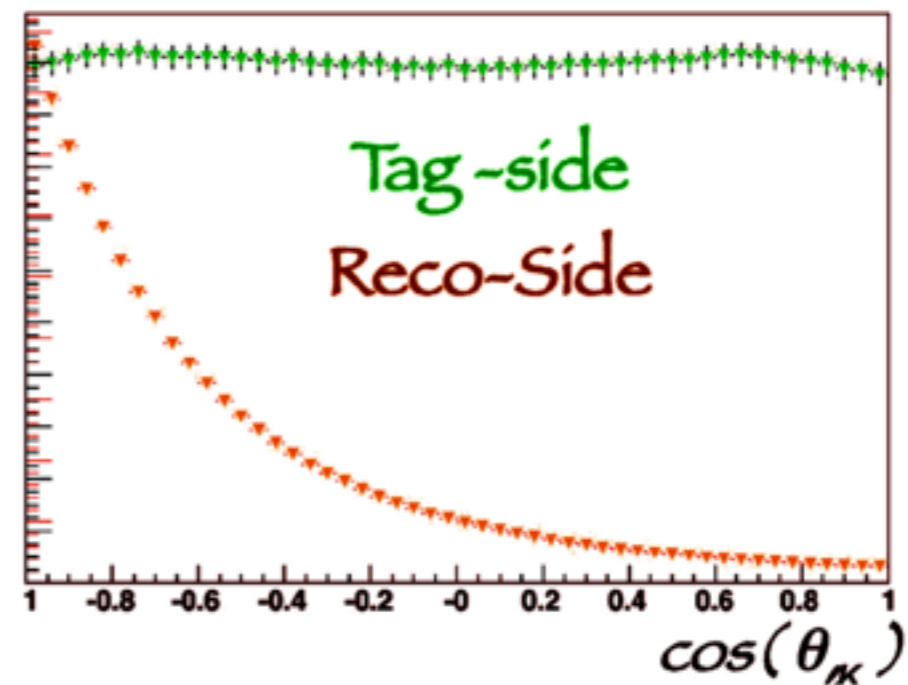
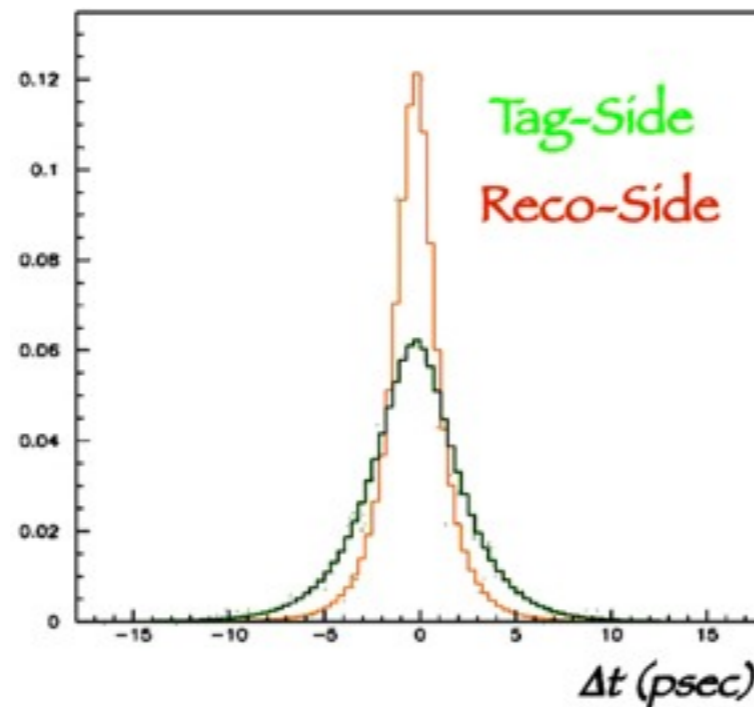


# Kaon tag



- Kaon can also come from the reco. side, mimicking a mixed event.

Use  $\Delta t$  and angular correlation in the PDF to separate the two.



# Contributing asymmetries

- Observed asymmetry has contributions from physics and detector.
- Observed asymmetry for same-sign signal events with tag-side kaon reflects  $(\ell, \pi_s)$  reconstruction charge asymmetry, kaon-id charge asymmetry, and physics  $A_{sl}$ .
  - ◆  $A_{obs;K-tag} \approx A_{\ell\pi} + A_K + A_{sl}$
- If kaon is from reco. side, then physics  $A_{sl}$  contribution is diluted by the mixing probability  $\chi_d$ . (mixed and unmixed are equally likely to be reconstructed.)
  - ◆  $A_{obs;K-rec} \approx A_{\ell\pi} + A_K + \chi_d A_{sl}$
- Asymmetry before kaon tagging  $\frac{N(\ell^+ \pi_s^-) - N(\ell^- \pi_s^+)}{N(\ell^+ \pi_s^-) + N(\ell^- \pi_s^+)}$ 
  - ◆  $A_{obs;rec} \approx A_{\ell\pi} + \chi_d A_{sl}$
- We have enough observables to determine all three  $A_{\ell\pi}$ ,  $A_K$ , and  $A_{sl}$  from data simultaneously.



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- We have enough observables to determine all three  $A_{\ell\pi}$ ,  $A_K$ , and  $A_{sl}$  from data simultaneously.



# Signal physics model

- Without interference from doubly-Cabibbo-suppressed decays (DCSD) on the tag side, the underlying physics would be relatively simple:

$$\begin{array}{l}
 \text{opposite-sign/} \\
 \text{unmixed} \\
 \\
 \text{same-sign/} \\
 \text{mixed}
 \end{array}
 \left\{ \begin{array}{l}
 \mathcal{F}_{\bar{B}^0 B^0}(\Delta t) = \frac{\Gamma_0 e^{-\Gamma_0 |\Delta t|}}{2} [\cosh(\Delta\Gamma \Delta t/2) + \cos(\Delta m_d \Delta t)] \\
 \mathcal{F}_{B^0 \bar{B}^0}(\Delta t) = \frac{\Gamma_0 e^{-\Gamma_0 |\Delta t|}}{2} [\cosh(\Delta\Gamma \Delta t/2) + \cos(\Delta m_d \Delta t)] \\
 \\
 \mathcal{F}_{\bar{B}^0 \bar{B}^0}(\Delta t) = \frac{\Gamma_0 e^{-\Gamma_0 |\Delta t|}}{2} [\cosh(\Delta\Gamma \Delta t/2) - \cos(\Delta m_d \Delta t)] \left| \frac{q}{p} \right|^2 \\
 \mathcal{F}_{B^0 B^0}(\Delta t) = \frac{\Gamma_0 e^{-\Gamma_0 |\Delta t|}}{2} [\cosh(\Delta\Gamma \Delta t/2) - \cos(\Delta m_d \Delta t)] \left| \frac{p}{q} \right|^2
 \end{array} \right.$$





# Signal physics model

- Including interference from doubly-Cabibbo-suppressed decays (DCSD) on the tag side.

$$\begin{aligned} \mathcal{F}_{\bar{B}^0 B^0}(\Delta t) &= \frac{\Gamma_0 e^{-\Gamma_0 |\Delta t|}}{2(1+r'^2)} \left[ \left(1 + \left|\frac{q}{p}\right|^2 r'^2\right) \cosh(\Delta\Gamma\Delta t/2) + \left(1 - \left|\frac{q}{p}\right|^2 r'^2\right) \cos(\Delta m_d \Delta t) - \left|\frac{q}{p}\right| (b+c) \sin(\Delta m_d \Delta t) \right], \\ \mathcal{F}_{B^0 \bar{B}^0}(\Delta t) &= \frac{\Gamma_0 e^{-\Gamma_0 |\Delta t|}}{2(1+r'^2)} \left[ \left(1 + \left|\frac{p}{q}\right|^2 r'^2\right) \cosh(\Delta\Gamma\Delta t/2) + \left(1 - \left|\frac{p}{q}\right|^2 r'^2\right) \cos(\Delta m_d \Delta t) + \left|\frac{p}{q}\right| (b-c) \sin(\Delta m_d \Delta t) \right], \\ \mathcal{F}_{\bar{B}^0 \bar{B}^0}(\Delta t) &= \frac{\Gamma_0 e^{-\Gamma_0 |\Delta t|}}{2(1+r'^2)} \left[ \left(1 + \left|\frac{p}{q}\right|^2 r'^2\right) \cosh(\Delta\Gamma\Delta t/2) - \left(1 - \left|\frac{p}{q}\right|^2 r'^2\right) \cos(\Delta m_d \Delta t) - \left|\frac{p}{q}\right| (b-c) \sin(\Delta m_d \Delta t) \right] \left|\frac{q}{p}\right|^2, \\ \mathcal{F}_{B^0 B^0}(\Delta t) &= \frac{\Gamma_0 e^{-\Gamma_0 |\Delta t|}}{2(1+r'^2)} \left[ \left(1 + \left|\frac{q}{p}\right|^2 r'^2\right) \cosh(\Delta\Gamma\Delta t/2) - \left(1 - \left|\frac{q}{p}\right|^2 r'^2\right) \cos(\Delta m_d \Delta t) + \left|\frac{q}{p}\right| (b+c) \sin(\Delta m_d \Delta t) \right] \left|\frac{p}{q}\right|^2, \end{aligned}$$

$$r' \sim \frac{\mathcal{A}(\bar{b} \rightarrow \bar{u}c\bar{d})}{\mathcal{A}(b \rightarrow c\bar{u}d)} \sim \mathcal{O}(\%)$$

$$b = 2r' \sin(2\beta + \gamma) \cos \delta'$$

$$c = -2r' \cos(2\beta + \gamma) \sin \delta'$$



# Fitting

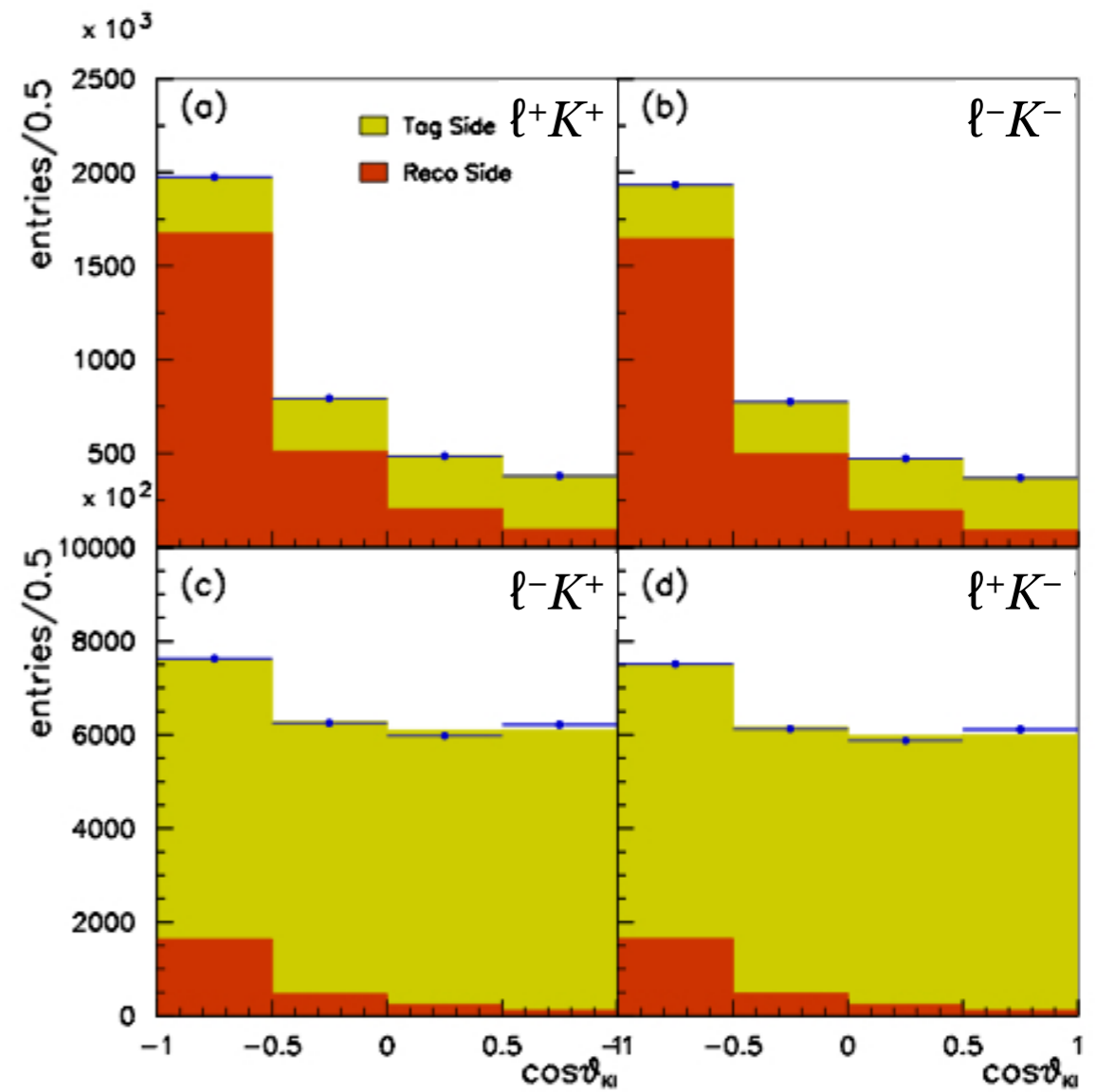
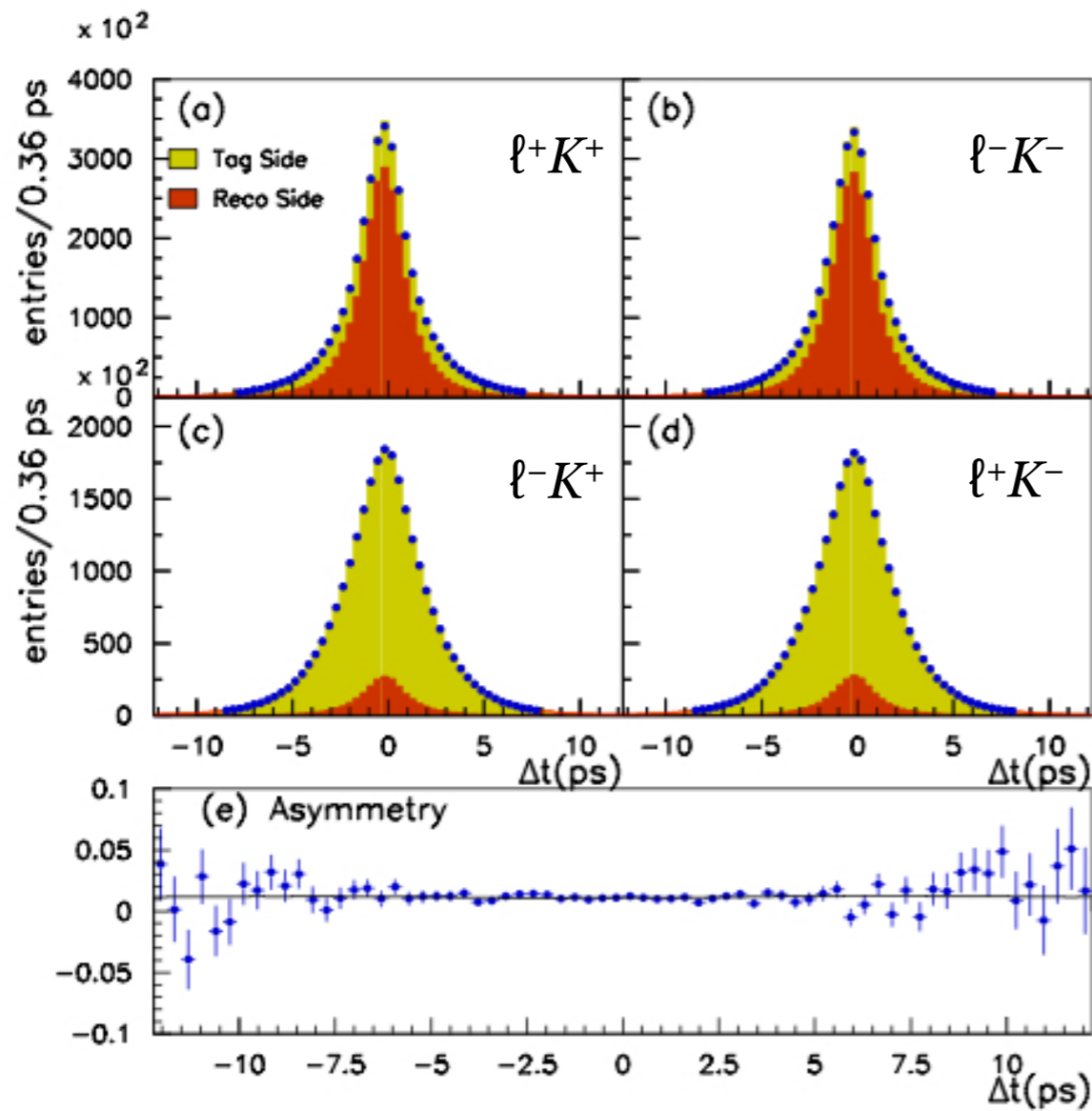
- Five variables (  $\Delta t$ ,  $\sigma_{\Delta t}$ ,  $\cos\theta_{\ell K}$ ,  $M_{\nu}^2$ ,  $p_K$  ) binned fit to separate signal to background, tag-side kaon from reco-side.
- Also include opposite-sign (unmixed) events  $\ell^+K^-/\ell^-K^+$  to improve resolution and mis-tag parameters, etc.
- More than 100 free parameters:
  - ◆  $A_{sl}$ ,  $A_{\ell\pi}$ ,  $A_K$ ,
  - ◆ reco-side  $K$  fraction,
  - ◆ wrong tag fractions (charge dependent),
  - ◆ Doubly-Cabibbo-suppressed decay parameters,
  - ◆  $\Delta t$  resolution parameters,
  - ◆  $B$  lifetime, mixing rate, etc.



# Result

$$A_{sl} = (0.6 \pm 1.6_{-3.2}^{+3.6}) \times 10^{-3}$$

BABAR Preliminary



# Result summary

- Lifetime and mixing rate consistent with the world average.
- A few per mille reconstruction asymmetry.
- Order of a percent kaon tag asymmetry.
- Fit to MC reproduces the generated value (also checked varying the generated value).

	Fit to Data	Fit to MC	MC value
$A_{sl}$	$(0.6 \pm 1.6) \times 10^{-3}$	$(0.4 \pm 0.5) \times 10^{-3}$	0
$A_{e\pi_s}$	$(3.0 \pm 0.4) \times 10^{-3}$	$(9.7 \pm 0.2) \times 10^{-3}$	-
$A_{\mu\pi_s}$	$(3.1 \pm 0.5) \times 10^{-3}$	$(8.4 \pm 0.3) \times 10^{-3}$	-
$A_K$	$(13.7 \pm 0.3) \times 10^{-3}$	$(14.7 \pm 0.1) \times 10^{-3}$	-
$\tau_B^0$	$1.5535 \pm 0.0019$	$1.5668 \pm 0.0012$	1.540
$\Delta m_d$	$0.5085 \pm 0.0009$	$0.4826 \pm 0.0006$	0.489

$$\Delta_{CP} = 1 - \left| \frac{q}{p} \right| \simeq \frac{1}{2} A_{sl}$$

Source	$\sigma(\Delta_{CP})$
Peaking Sample Composition	$+1.50$ $-1.17 \times 10^{-3}$
Combinatorial Sample Composition	$\pm 0.39 \times 10^{-3}$
$\Delta t$ Resolution Model	$\pm 0.60 \times 10^{-3}$
$K_R$ Fraction	$\pm 0.11 \times 10^{-3}$
$K_R$ $\Delta t$ Distribution	$\pm 0.65 \times 10^{-3}$
Fit Bias	$+0.58$ $-0.46 \times 10^{-3}$
$CP$ eigenstate Description	$\pm 0$
Physical Parameters	$+0$ $-0.28 \times 10^{-3}$
Total	$+1.88$ $-1.61 \times 10^{-3}$



# Conclusions

$$A_{sl} = (0.6 \pm 1.6_{-3.2}^{+3.6}) \times 10^{-3}$$

BABAR Preliminary  
[accepted by PRL]

- Consistent with and more precise than previous  $B$  factories average  $(-0.5 \pm 5.6) \times 10^{-3}$ .
- Competitive and complementary to similar measurements at hadron colliders.
- Result consistent with the Standard Model, no sign of  $CP$  violation in  $B_d$  mixing.
  - ◆ tension still exists in  $B_s$  mixing;  $B$  factories are not able to contribute to that measurement.

