



University of
Zurich^{UZH}

Studies of the electroweak penguin transitions and radiative decays at LHCb

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On behalf of the LHCb collaboration

DPF Santa Cruz – August 2013



Overview

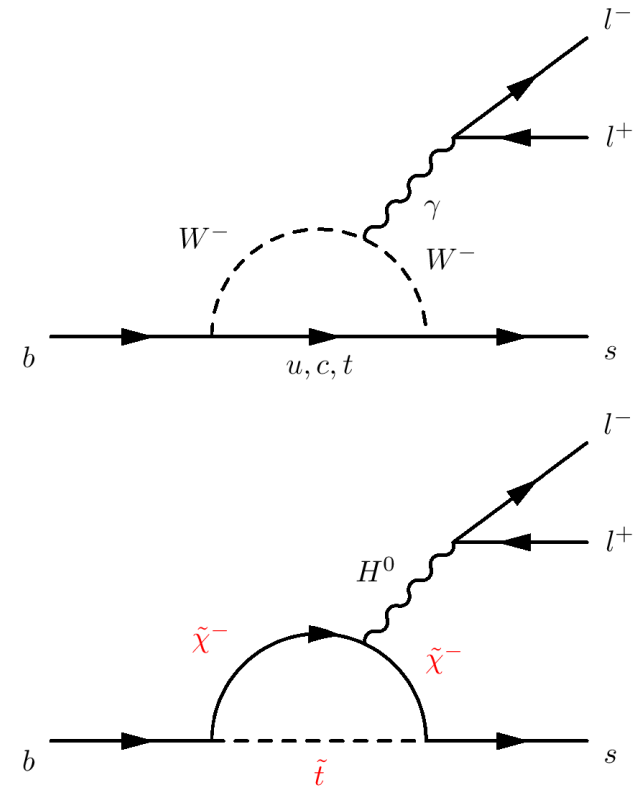
- Flavour changing neutral current (FCNC) decays can only proceed via loop-diagrams in the SM
- New Physics can enter at the same level as SM physics

$$\mathcal{O}_7 = \frac{e}{16\pi^2} m_b (\bar{s} \sigma^{\mu\nu} P_R b) F_{\mu\nu}$$

$$\mathcal{O}_9 = \frac{e^2}{16\pi^2} (\bar{s} \gamma^\mu P_L b) (\bar{\ell} \gamma_\mu \ell)$$

$$\mathcal{O}_{10} = \frac{e^2}{16\pi^2} (\bar{s} \gamma^\mu P_L b) (\bar{\ell} \gamma_\mu \gamma_5 \ell)$$

Operators in SM



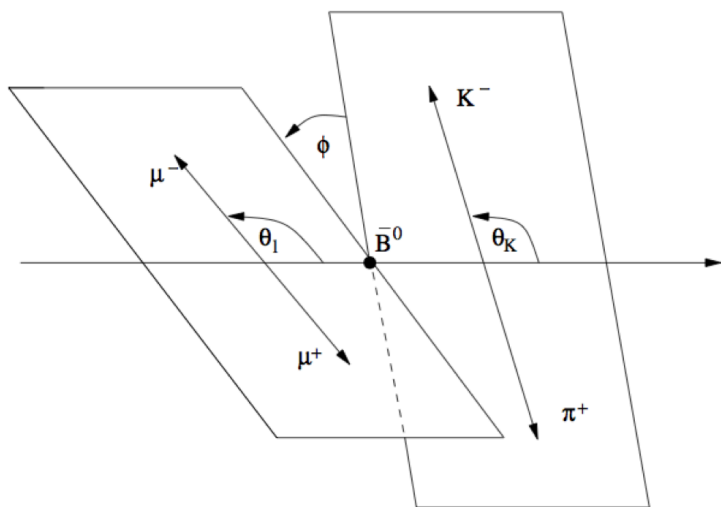
B⁰ → K* (→ K⁺π⁻) μμ

- B⁰ → K* μμ is a FCNC process mediated by penguin and W-box diagrams in the SM
- Its angular distribution is described by three angles (θ_l, θ_K, φ) and the dimuon invariant mass squared (q²)
- Observables F_L and S_i are functions of the Wilson coefficients and form-factors

$$\frac{1}{\Gamma} \frac{d^3(\Gamma + \bar{\Gamma})}{d \cos \theta_\ell d \cos \theta_K d \phi} = \frac{9}{32\pi} \left[\frac{3}{4}(1 - F_L) \sin^2 \theta_K + F_L \cos^2 \theta_K + \frac{1}{4}(1 - F_L) \sin^2 \theta_K \cos 2\theta_\ell \right. \\ \left. - F_L \cos^2 \theta_K \cos 2\theta_\ell + S_3 \sin^2 \theta_K \sin^2 \theta_\ell \cos 2\phi \right. \\ \left. + S_4 \sin 2\theta_K \sin 2\theta_\ell \cos \phi + S_5 \sin 2\theta_K \sin \theta_\ell \cos \phi \right. \\ \left. + S_6 \sin^2 \theta_K \cos \theta_\ell + S_7 \sin 2\theta_K \sin \theta_\ell \sin \phi \right. \\ \left. + S_8 \sin 2\theta_K \sin 2\theta_\ell \sin \phi + S_9 \sin^2 \theta_K \sin^2 \theta_\ell \sin 2\phi \right]$$

- Performing the folding $\phi \rightarrow \phi + \pi$ ($\phi < 0$) leads to:

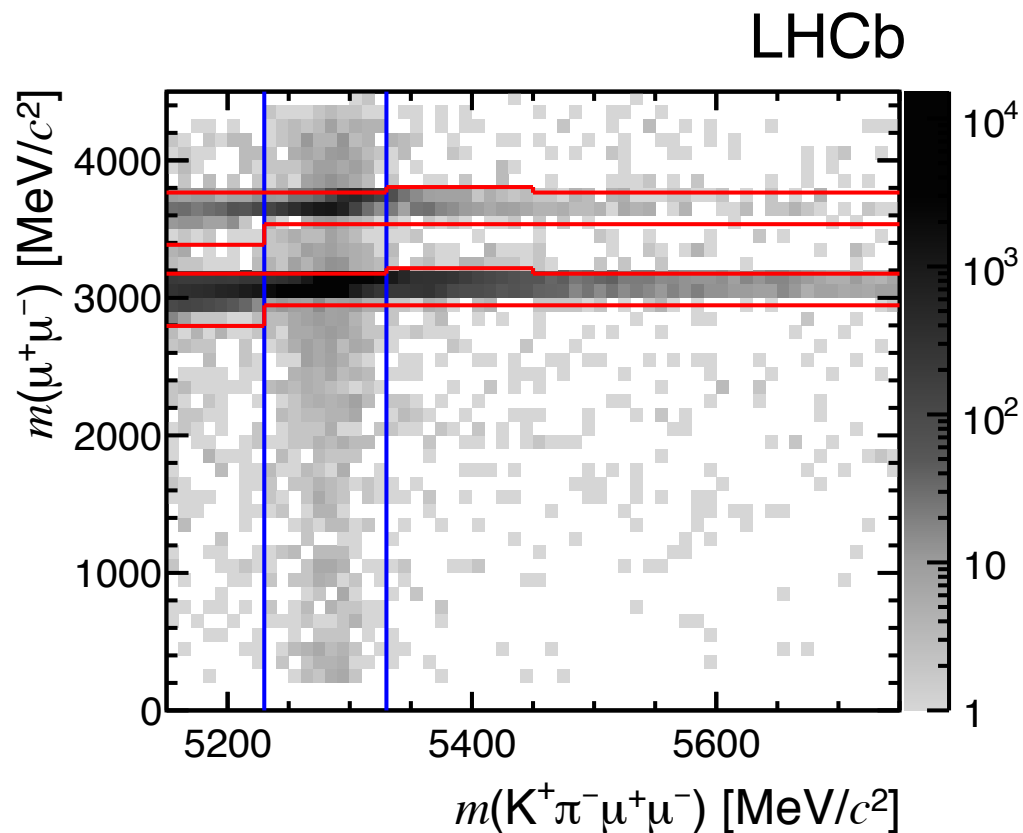
$$\frac{1}{\Gamma} \frac{d^3(\Gamma + \bar{\Gamma})}{d \cos \theta_\ell d \cos \theta_K d \hat{\phi}} = \frac{9}{32\pi} \left[\frac{3}{4} (1 - F_L) \sin^2 \theta_K + F_L \cos^2 \theta_K \right. \\ \left. + \frac{1}{4} (1 - F_L) \sin^2 \theta_K \cos 2\theta_\ell - F_L \cos^2 \theta_K \cos 2\theta_\ell \right. \\ \left. + S_3 \sin^2 \theta_K \sin^2 \theta_\ell \cos 2\hat{\phi} + \frac{4}{3} A_{FB} \sin^2 \theta_K \cos \theta_\ell \right. \\ \left. + 2\theta_K \sin 2\theta_\ell \sin \hat{\phi} + A_9 \sin^2 \theta_K \sin^2 \theta_\ell \sin 2\hat{\phi} \right]$$



- Theoretically cleaner observables (reduced form-factor uncertainties) can be calculated from existing ones

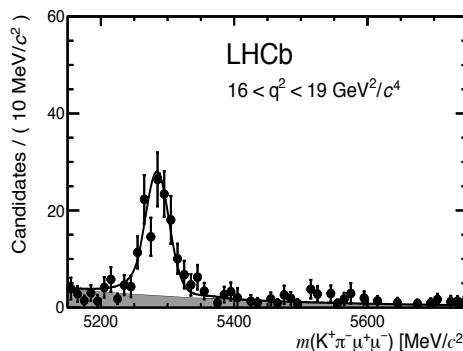
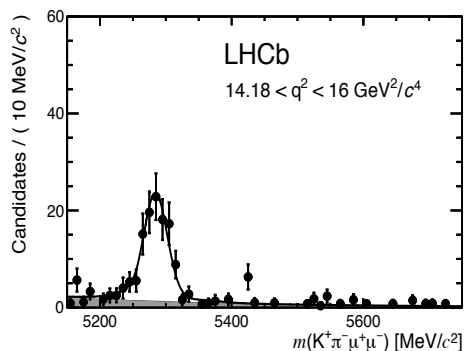
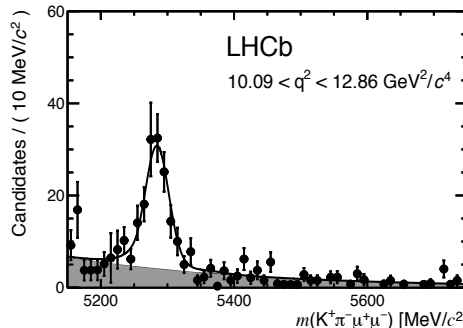
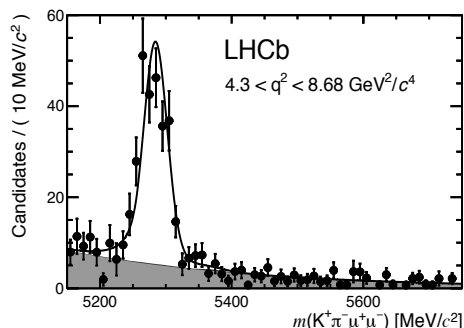
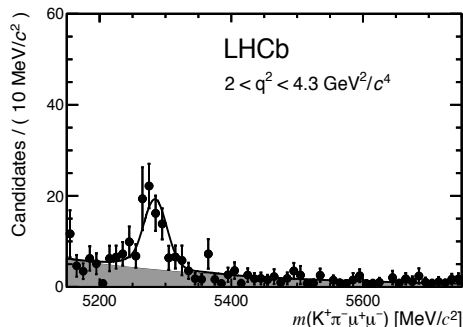
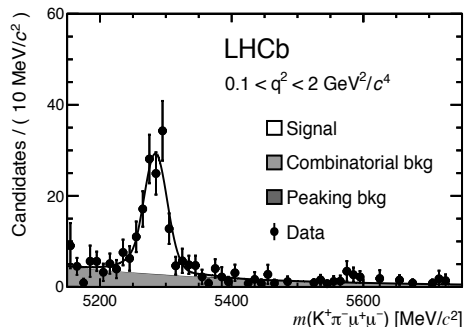
$$A_T^{\text{Re}} = \frac{\frac{4}{3} A_{FB}}{(1 - F_L)} \quad A_T^2 = \frac{2S_3}{(1 - F_L)}$$

Candidate selection



- Candidates selected using a BDT
- (Tree-level) Charmonium resonances vetoed
- Exclusive and partially reconstructed background investigated and reduced to negligible level

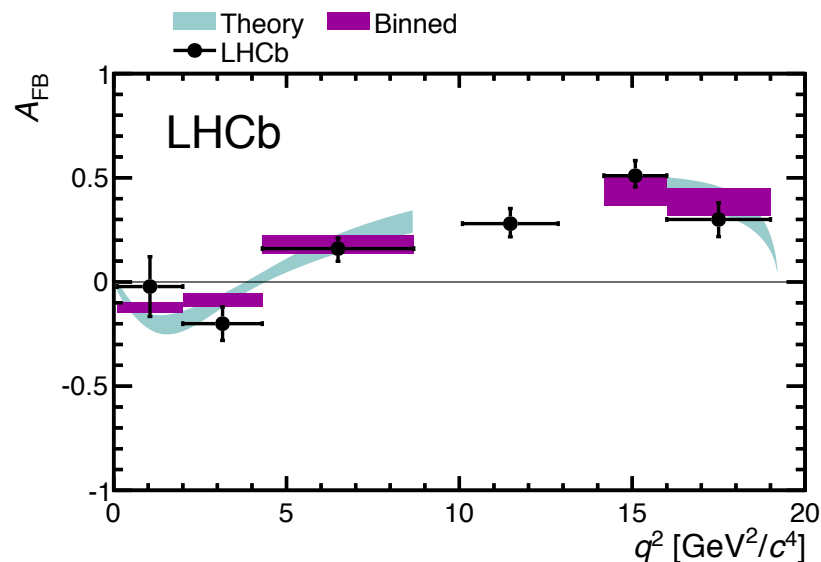
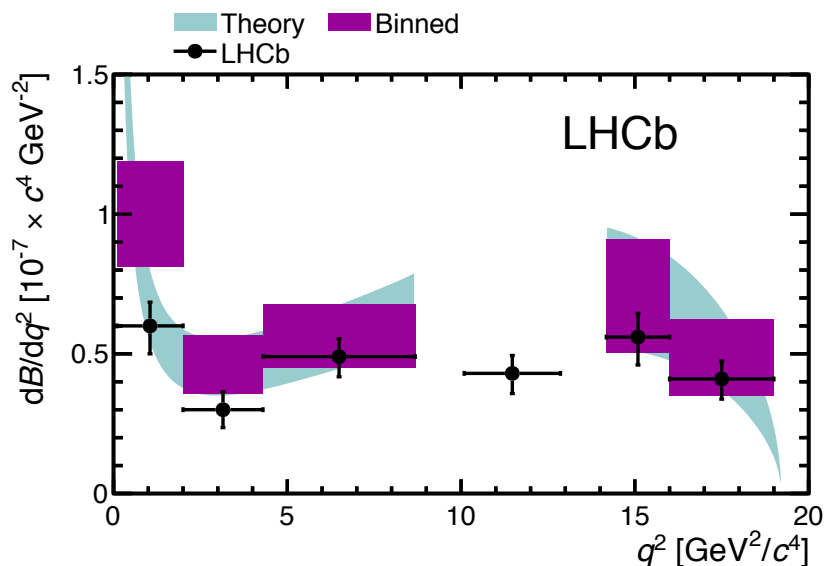
Candidate selection

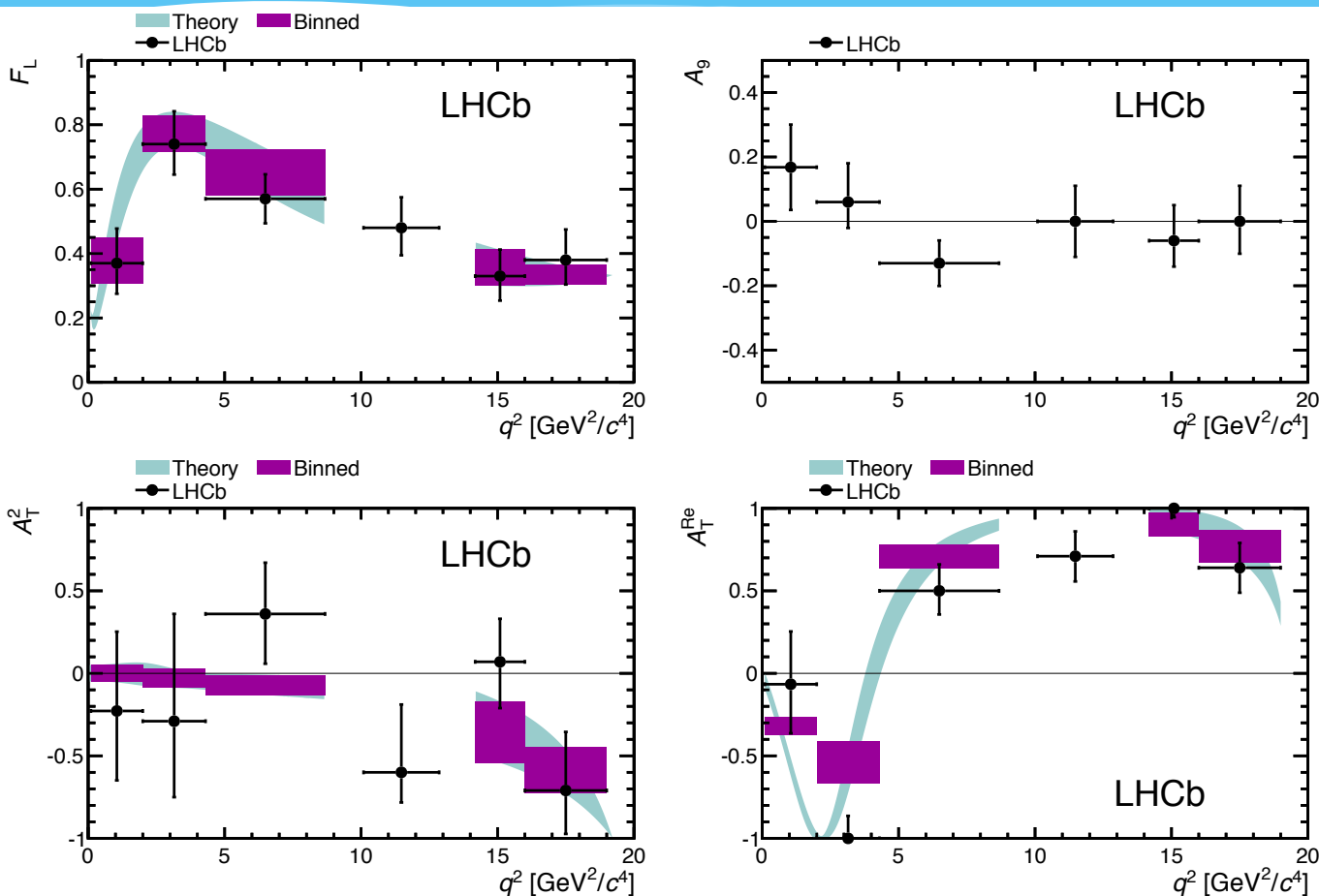


- Data/MC agreement calibrated using control modes e.g. $B^0 \rightarrow J/\psi K^*$, $D^* \rightarrow D(-\rightarrow K\pi)\pi_s$
- Acceptance correction accounted for by weighting candidates with inverse of efficiency determined from simulation
- Total yield of 883 ± 34 signal candidates
- Analysis performed in 6 bins of q^2 ($m_{\mu\mu}^2$) and theoretically clean $1 < q^2 < 6 \text{ GeV}^2/c^4$

- Results accord with the SM predictions (Bobeth et al., JHEP 01 (2012) 107)
- First measurement of the zero-crossing point of A_{FB} (error is stat+sys) (SM predictions in range 3.9-4.4 GeV²/c⁴)

$$q_0^2 = 4.9 \pm 0.9 \text{ GeV}^2/c^4$$





- Consistent with SM predictions (Bobeth et al., JHEP 01 (2012) 107)

New observables

- In low q^2 limit, observables denoted as P_4' , P_5' , P_6' , P_8' are predicted to be largely free from form-factor uncertainties (J. Mathias et al., JHEP 05 (2013) 137)

$$P'_{i=4,5,6,8} = \frac{S_{j=4,5,7,8}}{\sqrt{F_L(1 - F_L)}}$$

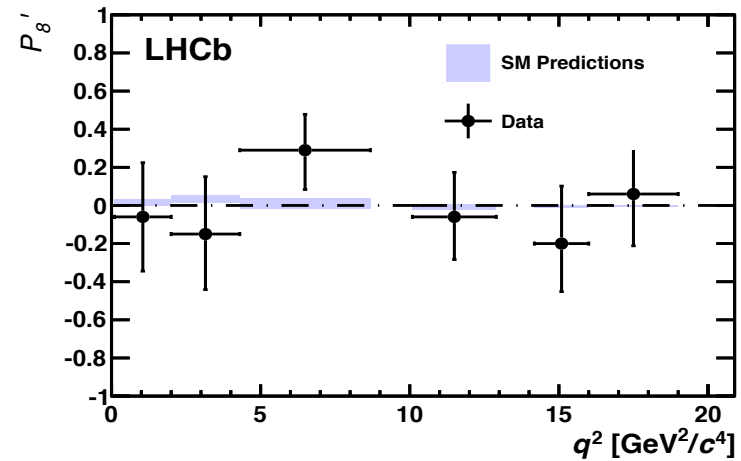
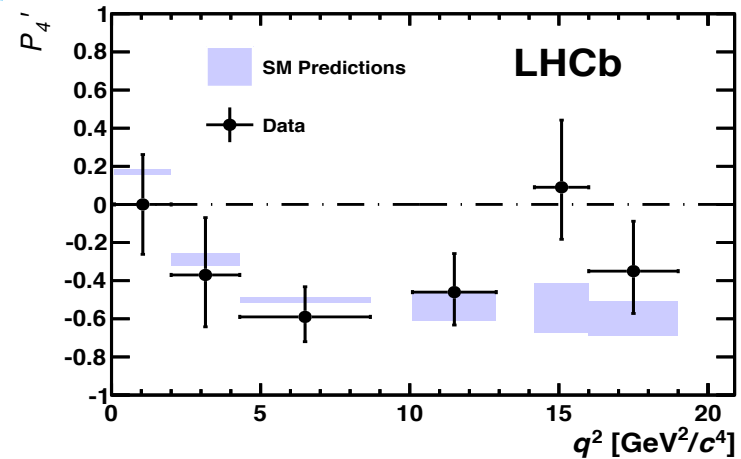
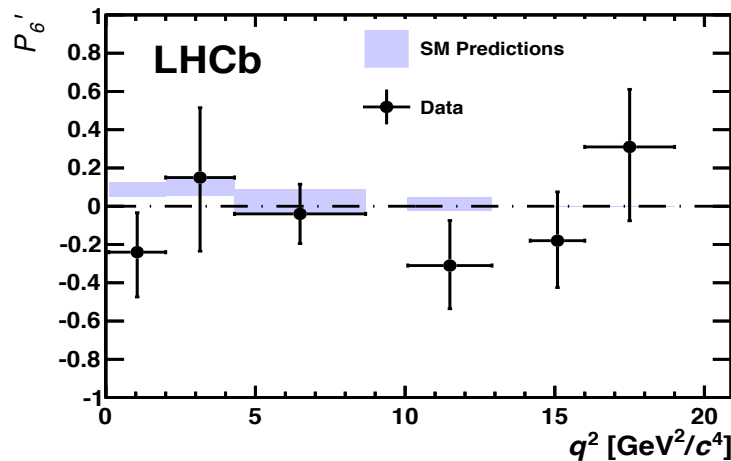
Here a 2D folding is used, for example:

$$P_5', S_5: \begin{cases} \phi \rightarrow -\phi & \text{for } \phi < 0 \\ \theta_\ell \rightarrow \pi - \theta_\ell & \text{for } \theta_\ell > \pi/2 \end{cases}$$

- Each folding preserves the first 5 terms in the differential angular distribution and the corresponding S_i term

First measurement of new observables

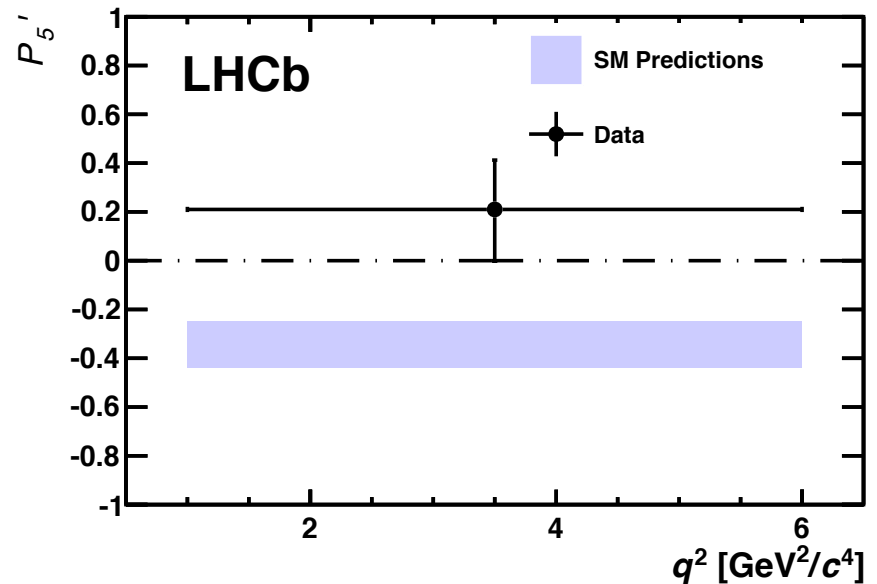
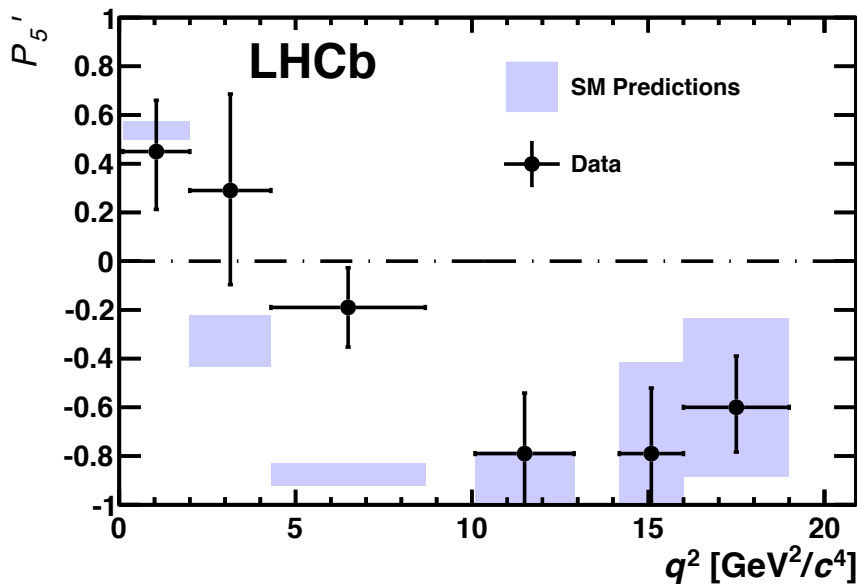
- Results accord with SM predictions (J. Mathias et al., JHEP 05 (2013) 137)



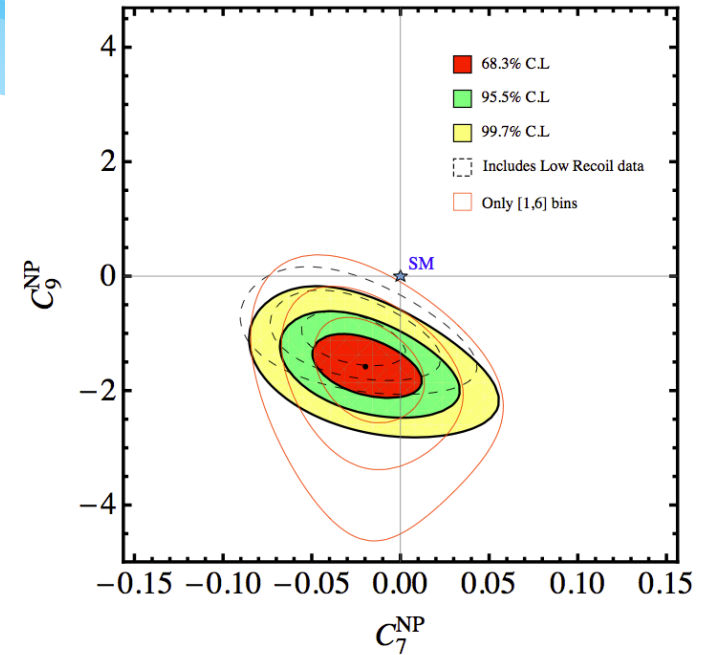
First measurement of new observables

- Local discrepancy of 3.7σ in 3rd bin of P_5' (wrt J. Mathias et al., JHEP 05 (2013) 137)
- 2.5σ discrepancy in theoretically clean $1 < q^2 < 6 \text{ GeV}^2/c^4$ bin

- 0.5% probability to observe 3.7σ deviation in 24 independent measurements

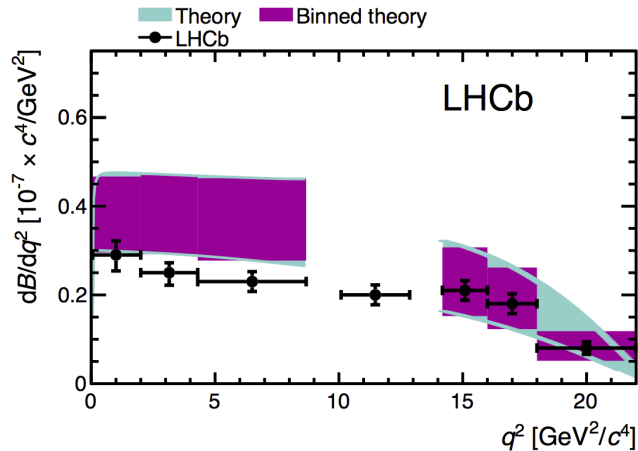


- Descotes-Genon et al. (arXiv:1307.5683) explain the discrepancy in P_5' and other smaller deviations through a large New Physics contribution to the Wilson coefficient of the semileptonic operator O_9



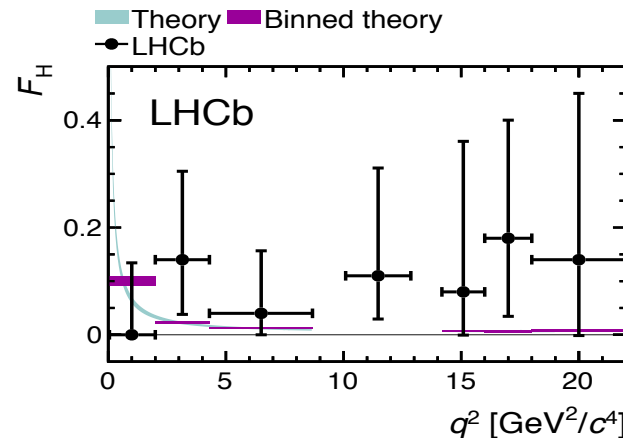
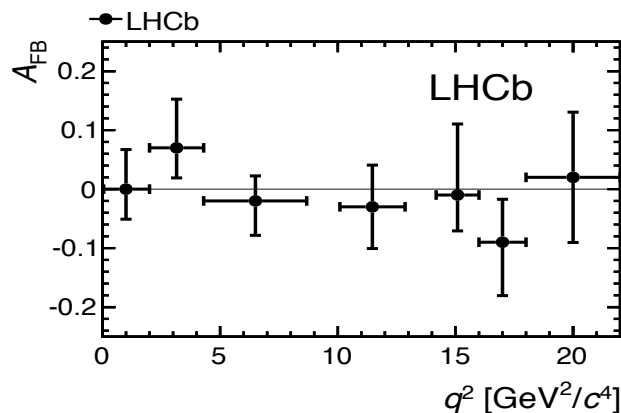
- Predictions for the first two bins and $1 < q^2 < 6 \text{ GeV}^2/c^4$ are also given by Jäger et al. (JHEP 05 (2013) 043)
- Leads to a larger theoretical uncertainty wrt J. Mathias et al.
- Measurements with higher statistics and further theoretical studies are necessary to draw more definitive conclusions

B⁺ -> K⁺ μμ



- Differential angular distribution described by single angle θ_l
- A_{FB} and F_H consistent with SM prediction of zero

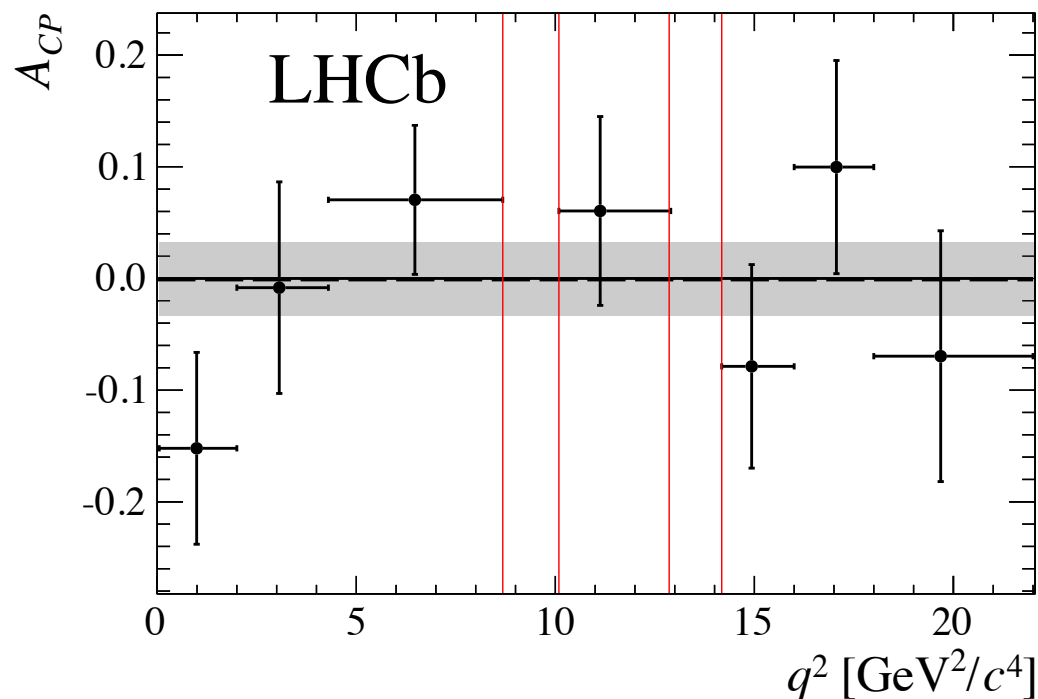
$$\frac{1}{\Gamma} \frac{d\Gamma[B^+ \rightarrow K^+ \mu^+ \mu^-]}{d\cos\theta_l} = \frac{3}{4}(1 - F_H)(1 - \cos^2\theta_l) + \frac{1}{2}F_H + A_{\text{FB}} \cos\theta_l$$



A_{CP} in $B^+ \rightarrow K^+ \mu \mu$

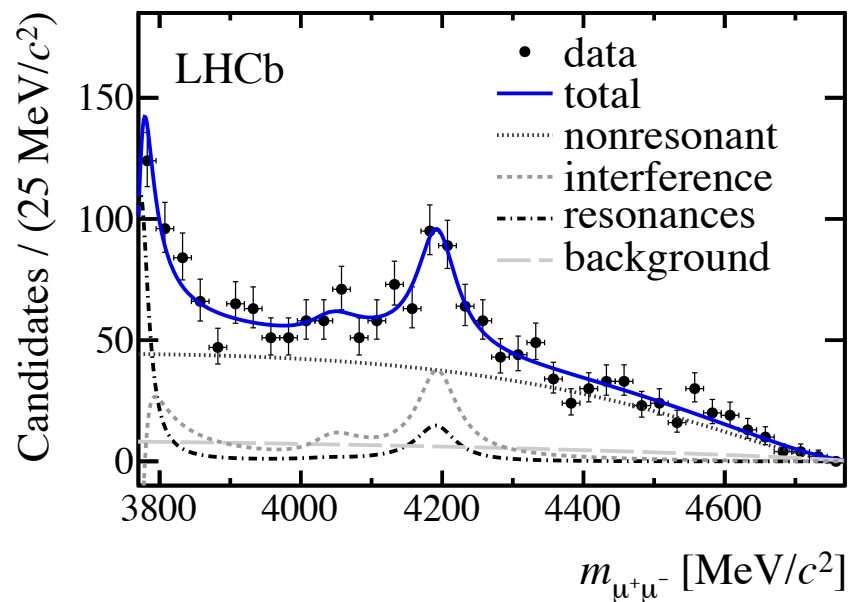
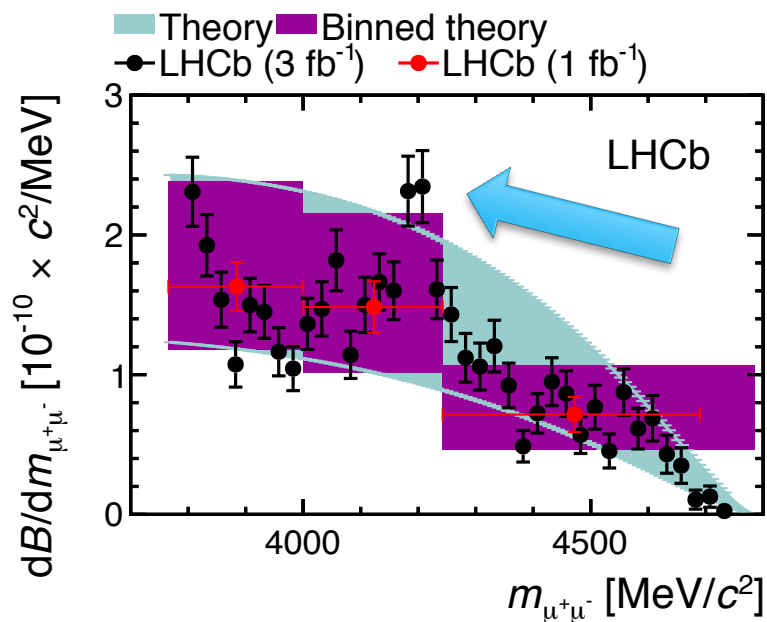
$$A_{CP} = \frac{\Gamma(B^- \rightarrow K^- \mu^+ \mu^-) - \Gamma(B^+ \rightarrow K^+ \mu^+ \mu^-)}{\Gamma(B^- \rightarrow K^- \mu^+ \mu^-) + \Gamma(B^+ \rightarrow K^+ \mu^+ \mu^-)} = -0.0004 \pm 0.033 \pm 0.005 \pm 0.007$$

- Control channel $B^+ \rightarrow J/\psi K^+$ used to account for production and detection asymmetries
- Left-Right asymmetry accounted for by averaging magnet polarity
- Result consistent with SM prediction



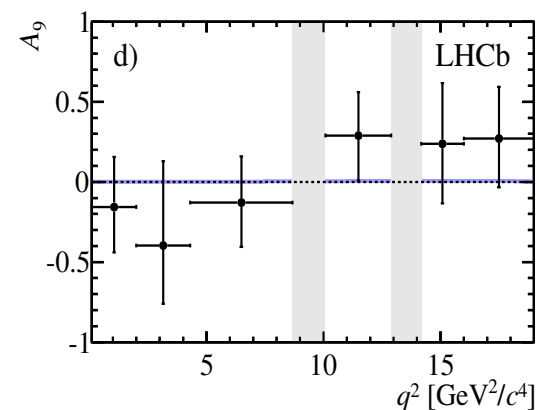
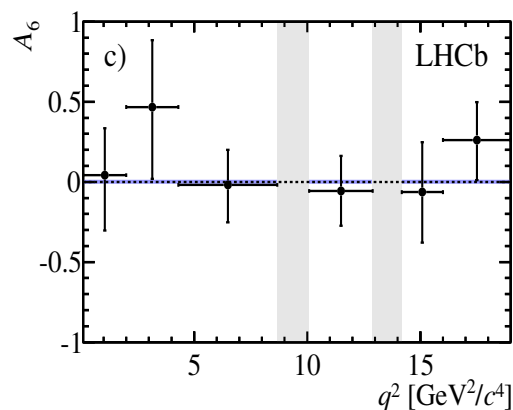
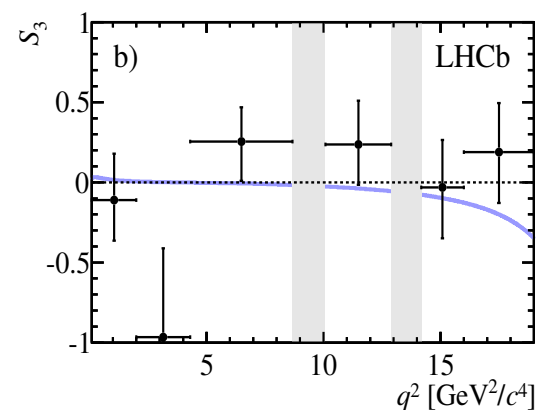
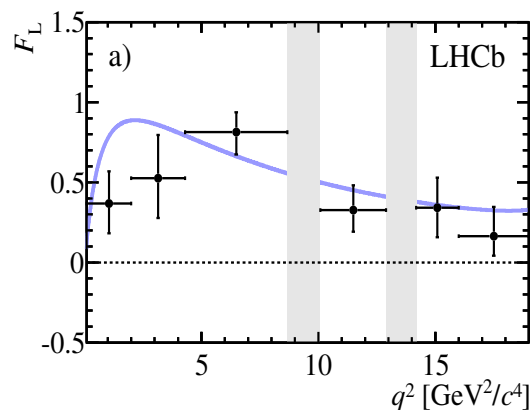
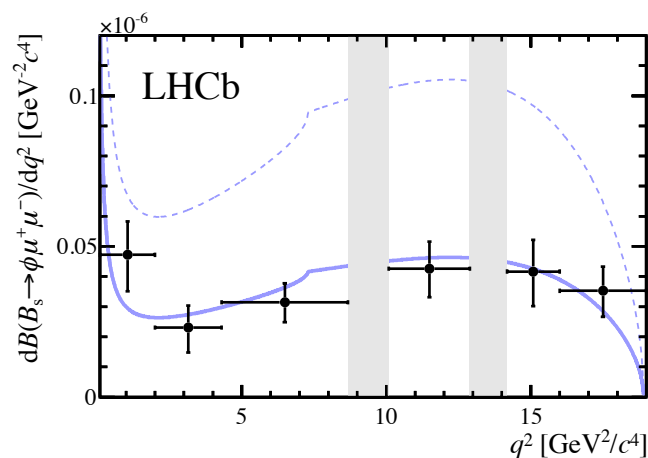
New resonance at in $B^+ \rightarrow K^+ \mu^+ \mu^-$ at high q^2

- Resonance compatible with the $\psi(4160)$ meson
- Larger than theoretical estimates
- First observation of both $B^+ \rightarrow \psi(4160)K^+$ and $\psi(4160) \rightarrow \mu^+ \mu^-$
- Could complicate further $b \rightarrow s \mu^+ \mu^-$ measurements at high q^2



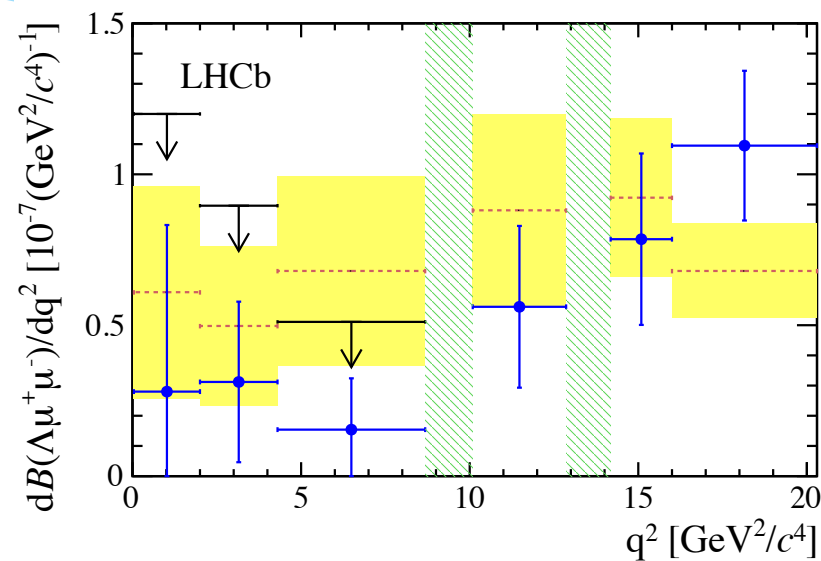
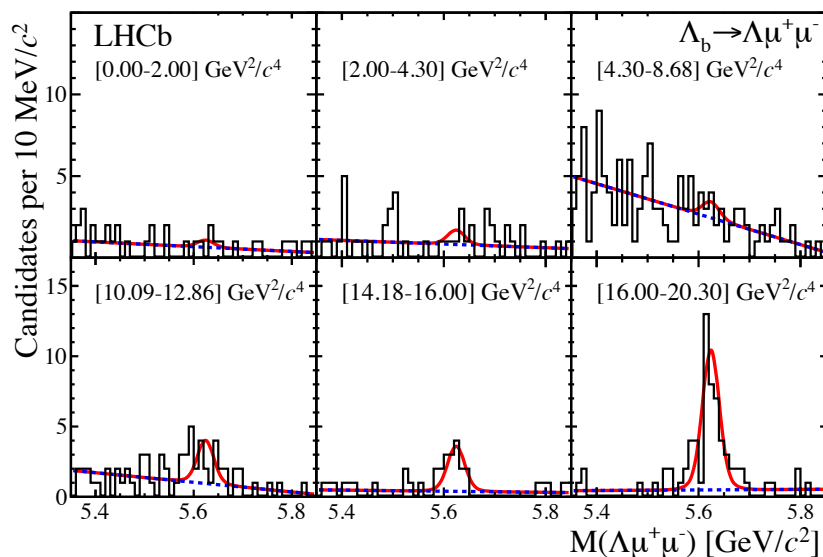
$B_s^0 \rightarrow \phi (-> K^+ K^-) \mu \mu$

- Branching fraction lower than SM theory predictions (blue dotted line)
- First angular analysis of $B_s^0 \rightarrow \phi \mu^+ \mu^-$
- All observables are consistent with SM expectation



$\Lambda_b \rightarrow \Lambda \mu \mu$

- Baryon decays more theoretically complex
- Yield of 78 ± 12 $\Lambda_b \rightarrow \Lambda \mu^+ \mu^-$ decays observed, mostly in q^2 regions above J/ψ mass
- $\Lambda_b \rightarrow J/\psi \Lambda$ normalization mode

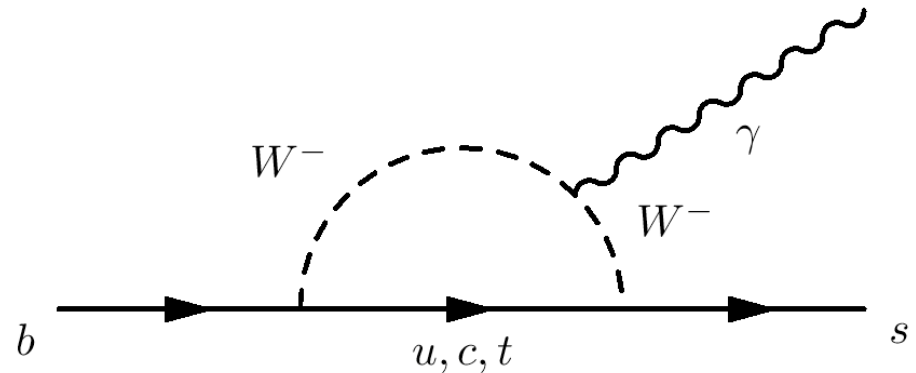


- Limits set at low q^2
- Good agreement with SM expectation

(Detmold et al., Phys.Rev. D87 (2013) 074502)

Radiative B decays

- $b \rightarrow s \gamma$ is a FCNC process
- In SM emitted photons in such decays are predicted to be predominately left handed since the recoil s quark that couples to the W boson is left handed

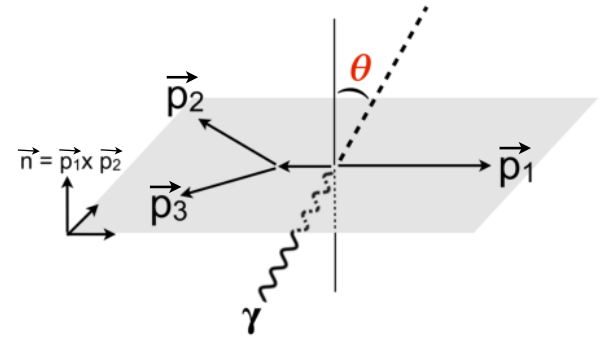


The photon polarization is defined as:

$$\lambda_{\gamma} \equiv \frac{|c_{\text{R}}|^2 - |c_{\text{L}}|^2}{|c_{\text{R}}|^2 + |c_{\text{L}}|^2}$$

B → K_{res} γ → K π π γ

- For decays of the type B → K_{res} γ → K π π γ, the photon polarization can be studied through the angular correlations of the daughters of the K_{res}
- For a single resonance, A_{ud} is proportional to λ_γ
- If the helicity amplitude J is known, λ_γ can be determined from a measurement of A_{ud}

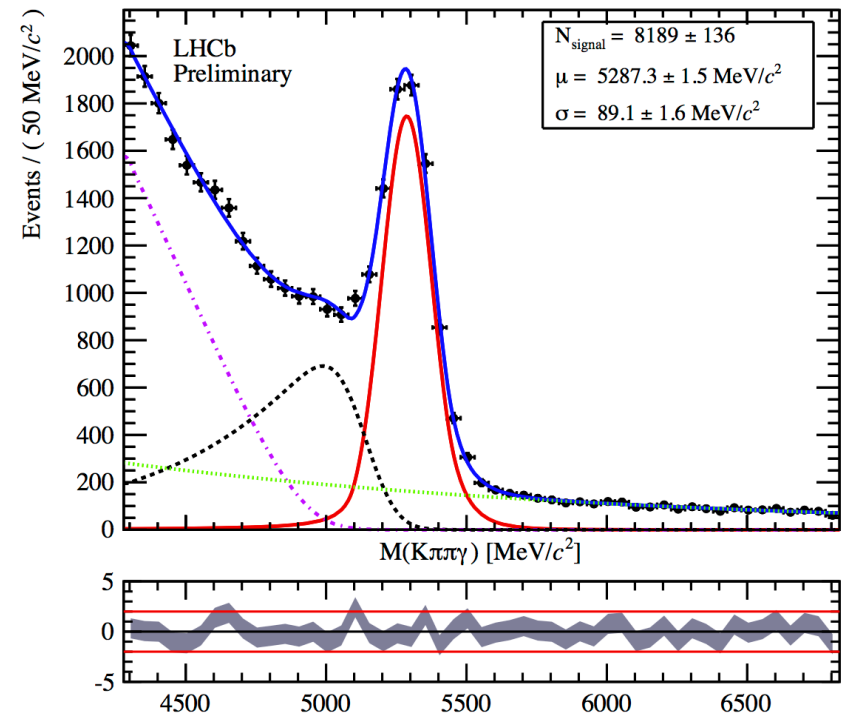


$$\mathcal{A}_{ud} \equiv \frac{\int_0^1 d\cos\tilde{\theta} \frac{d\Gamma}{d\cos\tilde{\theta}} - \int_{-1}^0 d\cos\tilde{\theta} \frac{d\Gamma}{d\cos\tilde{\theta}}}{\int_{-1}^1 d\cos\tilde{\theta} \frac{d\Gamma}{d\cos\tilde{\theta}}} = \frac{3}{4} \lambda_\gamma \frac{\int ds ds_{13} ds_{23} \text{Im} [\vec{n} \cdot (\vec{\mathcal{J}} \times \vec{\mathcal{J}}^*)]}{\int ds ds_{13} ds_{23} |\mathcal{J}|^2}$$

Theory input needed!

B → K⁺π⁻π⁺γ

- Inclusive CP asymmetry measured for first time
- A_{CP} compatible with zero
- A_{ud} measured for the first time
- Result 4.5σ from zero
- Evidence of photon polarization in b → sγ decays
- If theoretical predictions existed, first measurement of λ_γ possible



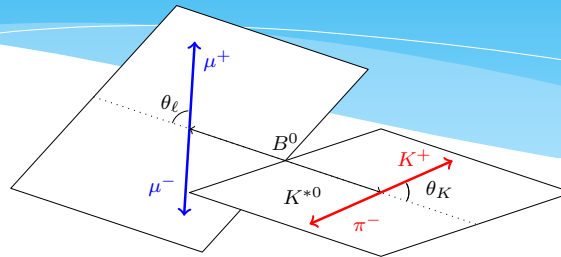
$$\mathcal{A}_{CP} = -0.007 \pm 0.015 (\text{stat})_{-0.011}^{+0.012} (\text{syst}) \quad \mathcal{A}_{ud} = -0.085 \pm 0.019 (\text{stat}) \pm 0.004 (\text{syst})$$

Summary

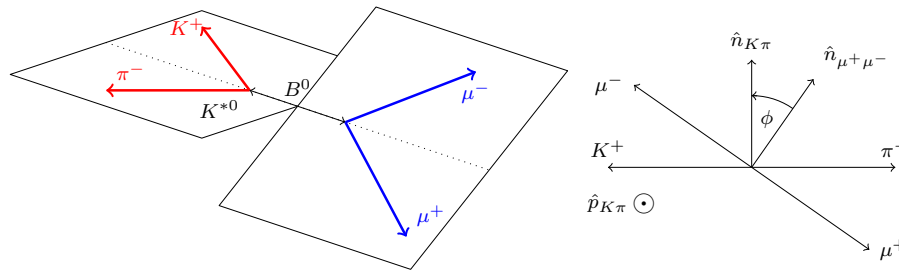
- LHCb has many recent precision measurements in electroweak penguin and radiative decay modes
 - New angular observables measured in $B^0 \rightarrow K^* \mu \mu$
 - First measurement of zero crossing point of A_{FB} in $B^0 \rightarrow K^* \mu \mu$
 - New resonance observed at high q^2 in $B^+ \rightarrow K^+ \mu \mu$
 - First angular analysis of $B_s^0 \rightarrow \phi \mu \mu$
 - First measurements of A_{CP} and A_{ud} in $B \rightarrow K \pi \pi \gamma$
- In general, there is good agreement between results and SM predictions except for a deviation in low q^2 region of P_5' in $B^0 \rightarrow K^* \mu \mu$
- Many results shown only using 1 fb^{-1} , updates with full 3 fb^{-1} to follow!

Backup

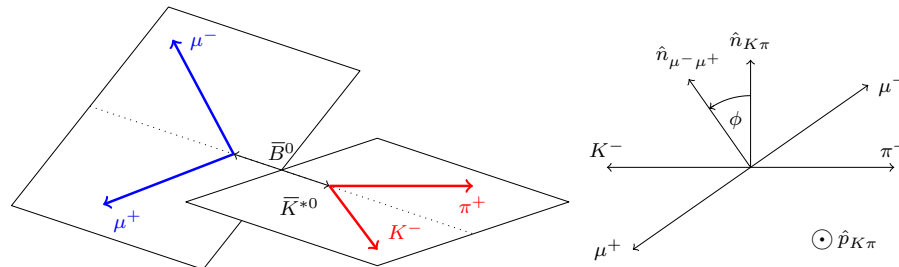
Angles in $B^0 \rightarrow K^* \mu \mu$



(a) θ_K and θ_ℓ definitions for the B^0 decay



(b) ϕ definition for the B^0 decay



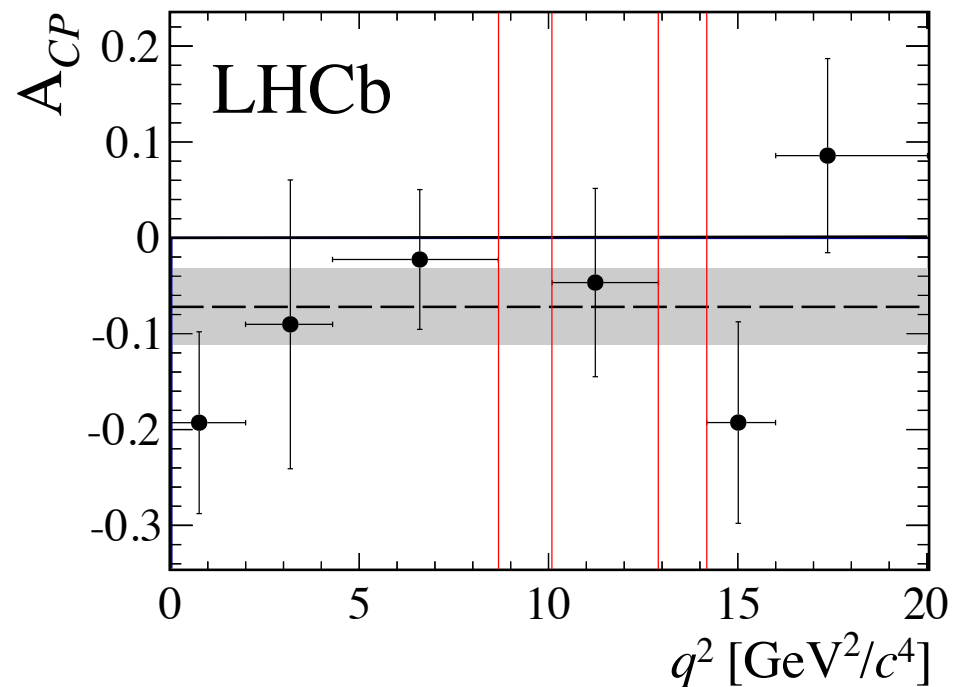
(c) ϕ definition for the \bar{B}^0 decay

A_{CP} in $B^0 \rightarrow K^* \mu \mu$

$$A_{CP} = \frac{\Gamma(B^- \rightarrow K^- \mu^+ \mu^-) - \Gamma(B^+ \rightarrow K^+ \mu^+ \mu^-)}{\Gamma(B^- \rightarrow K^- \mu^+ \mu^-) + \Gamma(B^+ \rightarrow K^+ \mu^+ \mu^-)}$$

$$-0.072 \pm 0.040 \text{ (stat.)} \pm 0.005 \text{ (syst.)}$$

- Control channel $B^+ \rightarrow J/\psi K^+$ used to account for production and detection asymmetries
- Left-Right asymmetry accounted for by averaging magnet polarity
- Result consistent with SM prediction



New observables in $B^0 \rightarrow K^* \mu \mu$

GeV^2/c^4	P'_4	P'_5	P'_6	P'_8
0.10 – 2.00	$0.00^{+0.26}_{-0.26} \pm 0.03$	$0.45^{+0.19}_{-0.22} \pm 0.09$	$-0.24^{+0.19}_{-0.22} \pm 0.05$	$-0.06^{+0.28}_{-0.28} \pm 0.02$
2.00 – 4.30	$-0.37^{+0.29}_{-0.26} \pm 0.08$	$0.29^{+0.39}_{-0.38} \pm 0.07$	$0.15^{+0.36}_{-0.38} \pm 0.05$	$-0.15^{+0.29}_{-0.28} \pm 0.07$
4.30 – 8.68	$-0.59^{+0.15}_{-0.12} \pm 0.05$	$-0.19^{+0.16}_{-0.16} \pm 0.03$	$-0.04^{+0.15}_{-0.15} \pm 0.05$	$0.29^{+0.17}_{-0.19} \pm 0.03$
10.09 – 12.90	$-0.46^{+0.20}_{-0.17} \pm 0.03$	$-0.79^{+0.16}_{-0.19} \pm 0.19$	$-0.31^{+0.23}_{-0.22} \pm 0.05$	$-0.06^{+0.23}_{-0.22} \pm 0.02$
14.18 – 16.00	$0.09^{+0.35}_{-0.27} \pm 0.04$	$-0.79^{+0.20}_{-0.13} \pm 0.18$	$-0.18^{+0.25}_{-0.24} \pm 0.03$	$-0.20^{+0.30}_{-0.25} \pm 0.03$
16.00 – 19.00	$-0.35^{+0.26}_{-0.22} \pm 0.03$	$-0.60^{+0.19}_{-0.16} \pm 0.09$	$0.31^{+0.38}_{-0.37} \pm 0.10$	$0.06^{+0.26}_{-0.27} \pm 0.03$
1.00 – 6.00	$-0.29^{+0.18}_{-0.16} \pm 0.03$	$0.21^{+0.20}_{-0.21} \pm 0.03$	$0.18^{+0.21}_{-0.21} \pm 0.03$	$0.23^{+0.18}_{-0.19} \pm 0.02$
	S_4	S_5	S_7	S_8
0.10 – 2.00	$-0.01^{+0.12}_{-0.12} \pm 0.03$	$0.22^{+0.09}_{-0.10} \pm 0.04$	$-0.12^{+0.11}_{-0.11} \pm 0.03$	$-0.04^{+0.12}_{-0.12} \pm 0.01$
2.00 – 4.30	$-0.14^{+0.13}_{-0.12} \pm 0.03$	$0.11^{+0.14}_{-0.13} \pm 0.03$	$0.06^{+0.15}_{-0.15} \pm 0.02$	$-0.05^{+0.12}_{-0.12} \pm 0.02$
4.30 – 8.68	$-0.29^{+0.06}_{-0.06} \pm 0.02$	$-0.09^{+0.08}_{-0.08} \pm 0.01$	$-0.03^{+0.07}_{-0.08} \pm 0.04$	$0.13^{+0.08}_{-0.08} \pm 0.01$
10.09 – 12.90	$-0.22^{+0.09}_{-0.08} \pm 0.02$	$-0.40^{+0.08}_{-0.10} \pm 0.10$	$-0.17^{+0.11}_{-0.11} \pm 0.03$	$-0.03^{+0.10}_{-0.10} \pm 0.01$
14.18 – 16.00	$0.05^{+0.14}_{-0.08} \pm 0.01$	$-0.38^{+0.10}_{-0.09} \pm 0.09$	$-0.08^{+0.13}_{-0.14} \pm 0.01$	$-0.10^{+0.13}_{-0.12} \pm 0.02$
16.00 – 19.00	$-0.16^{+0.11}_{-0.09} \pm 0.01$	$-0.29^{+0.09}_{-0.08} \pm 0.04$	$0.15^{+0.16}_{-0.15} \pm 0.03$	$0.03^{+0.12}_{-0.12} \pm 0.02$