



S. Stone
L. Zhang



Use of $B \rightarrow J/\psi f_0$ decays to discern the $q\bar{q}$ or tetraquark nature of scalar mesons

DPF, Santa Cruz, Cal., Aug. 16, 2013



Scalar meson quandry

- While 0^- and 1^- mesons follow a simple rule that adding an s-quark increases their mass, the 0^+ mesons are difficult to understand in this context

Isospin	1^- state	mass	$q\bar{q}$	0^+ state	mass
1	ρ	776 MeV	$(u\bar{u}+d\bar{d})\sqrt{2}$	$a_0(980)$	980 MeV
0	ω	783 MeV	$(uu-d\bar{d})\sqrt{2}$	$f_0(500)$ or σ	500 MeV
1/2	$K^*(892)$	892 MeV	$(u \text{ or } d) \bar{s}$	$\kappa(800)$	800 MeV
0	ϕ	1020 MeV	$s\bar{s}$	$f_0(980)$	980 MeV

- Suggestions that scalars are tetraquarks



Issues

- What is basic structure of matter?
 - Are the $f_0(500)$ & $f_0(980)$ 2 quark or 4 quark states?
 - We need to know this if we are to understand mesons made only of gluons (glueballs)
- What complications arise in other measurements due to quark substructure?
 - Ex: CP violation studies in $B \rightarrow J/\psi + \text{scalar}$, or any other meson with non- qq structure. See Fleischer, Knegjens & Ricciardi [arXiv:1109.1112]
 - Maybe a concern for other meson states

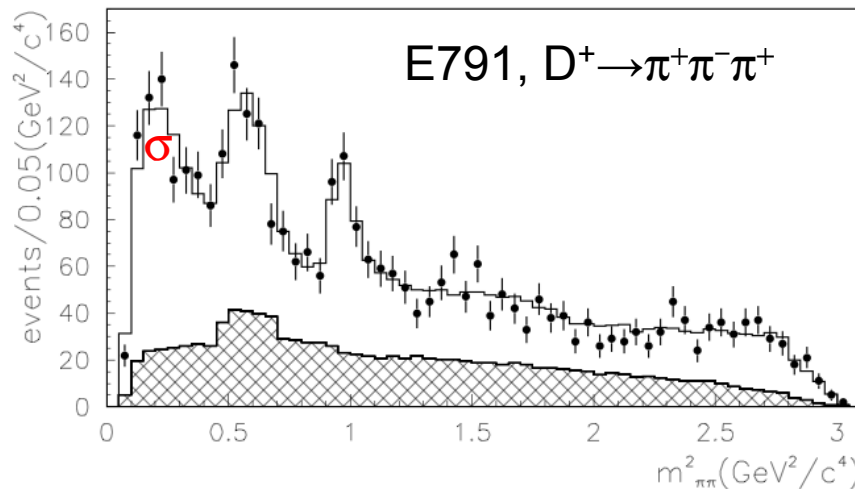


Much Interest

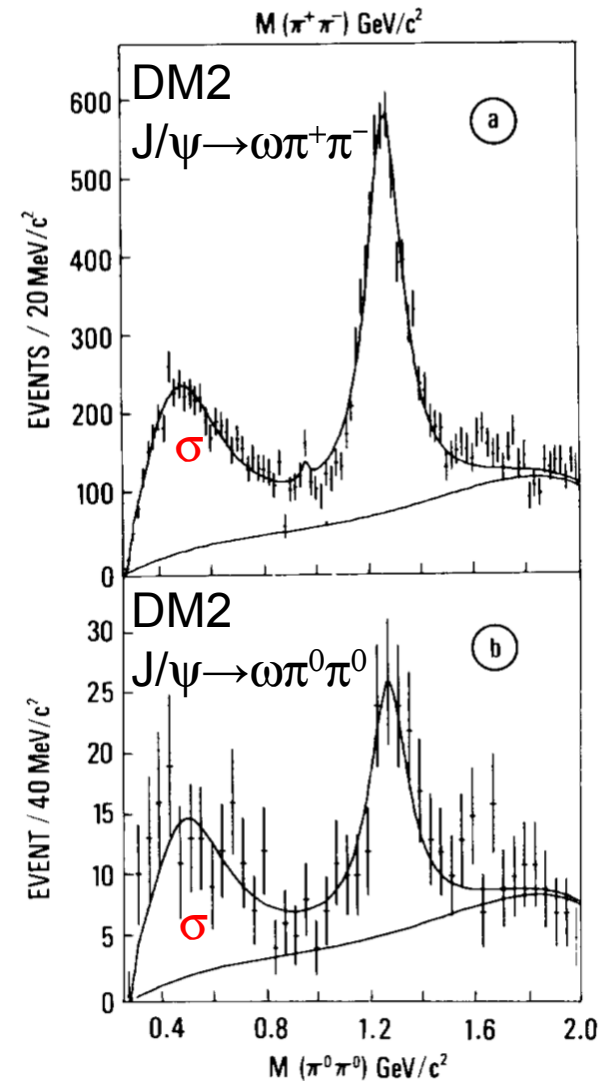
- S. Weinberg, “Tetraquark Mesons in Large N Quantum Chromodynamics,” arXiv:hep-ph/1303.0342
- G. t’Hooft et al., “A Theory of Scalar Mesons,” arXiv:hep-ph/0801.2288
- A. H. Fariborz et al. “Global aspects of the scalar meson puzzle,” arXiv:hep-ph/0902.2825
- R. L. Jaffe, Multi-Quark Hadrons I. The Phenomenology of $Q^2\bar{Q}^2$ Mesons, Phys. Rev. D15 (1977) 267



Knowledge about σ



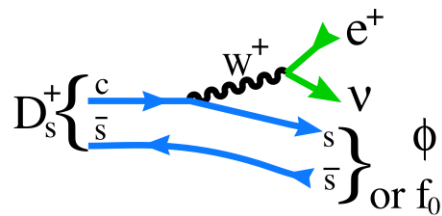
- Also $\pi\pi$ scattering experiments, vector meson radiative decays, etc...
- Mass & Width are uncertain
PDG “estimate” $M=400\text{-}550$ MeV,
 $\Gamma=400\text{-}700$ MeV





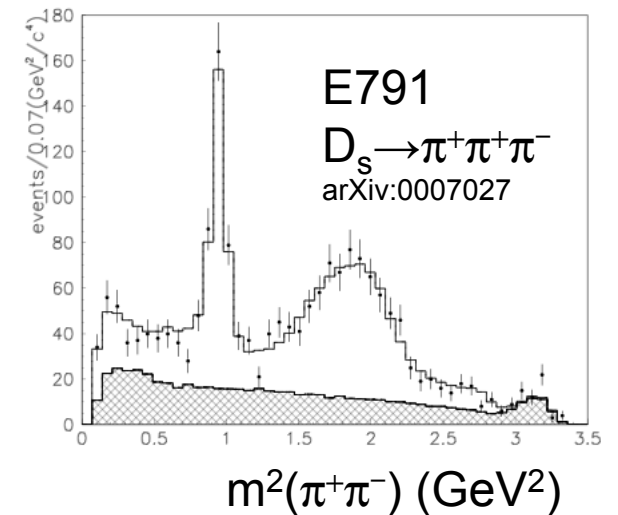
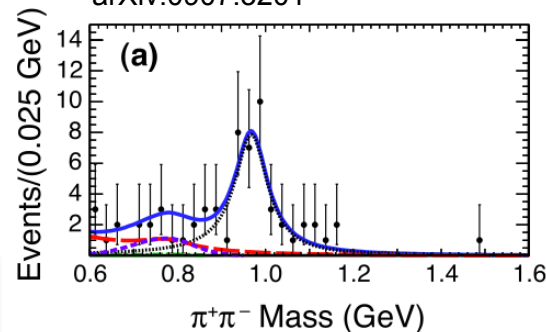
Knowledge about $f_0(980)$

- $f_0(980)$ seen in many experiments
- In weak decays

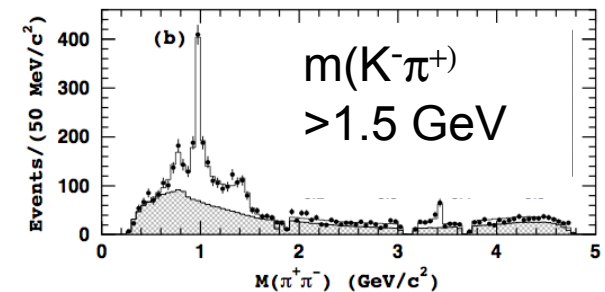


Note $s\bar{s}$ is $I=0$,
e.g. ρ not allowed

CLEO $D_s \rightarrow \pi^+ \pi^- e^- \nu$
arXiv:0907.3201



Belle $B^- \rightarrow K^- \pi^+ \pi^-$
arXiv:0512066

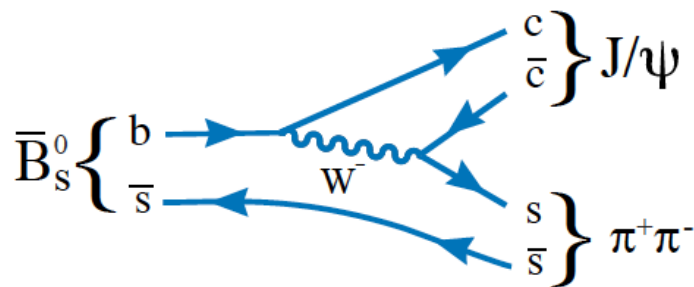


- Also LHCb $B_s \rightarrow J/\psi \pi^+ \pi^-$
(predicted in arXiv:0812.2832)

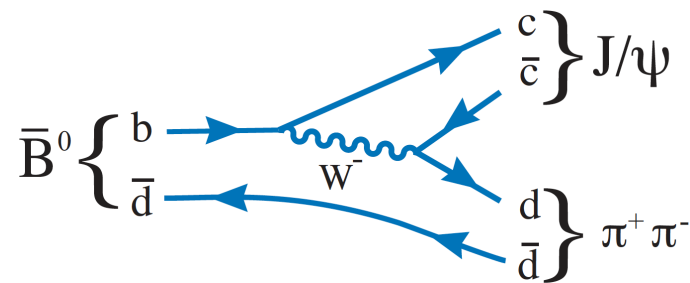


$B_{(s)}^0 \rightarrow J/\psi \pi^+ \pi^-$

- What can these decays tell us about σ & $f_0(980)$?
- These resonances plus others can be made via the following decay diagrams:

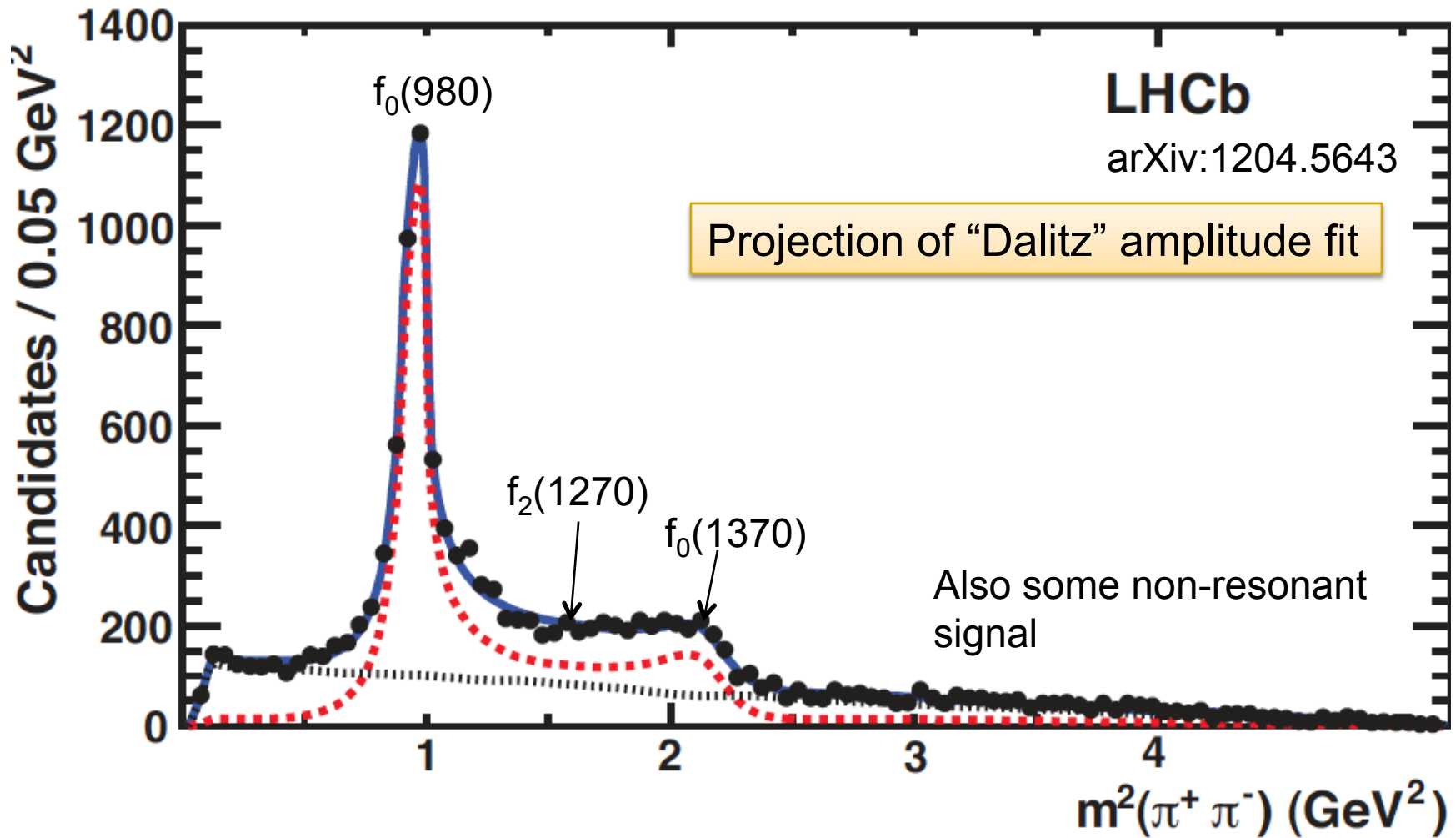


Again, $I=0$





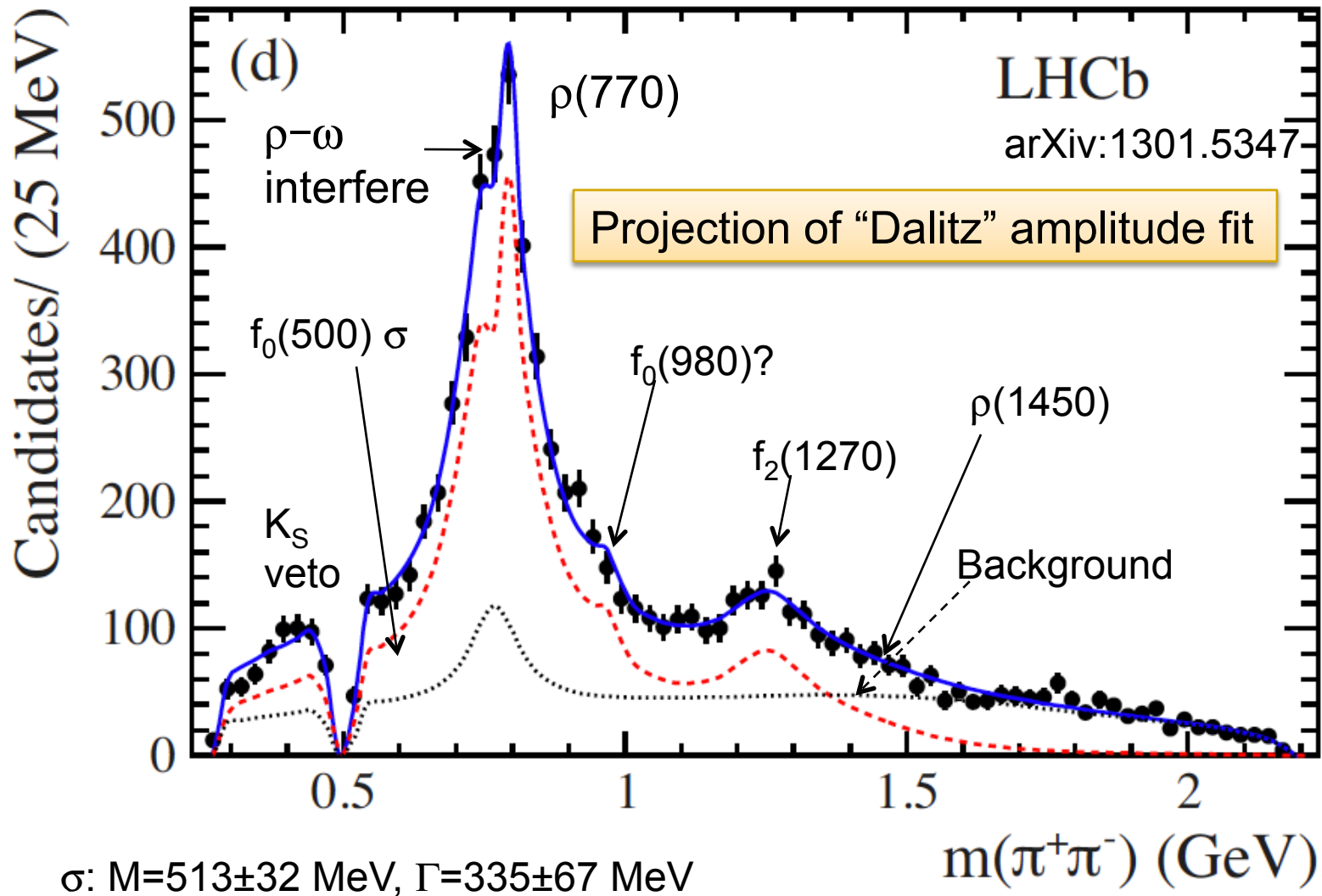
LHCb: $B_s \rightarrow J/\psi \pi^+ \pi^-$



DPF, Santa Cruz, Cal., Aug. 16, 2013



LHCb: $B^0 \rightarrow J/\psi \pi^+ \pi^-$





$\mathcal{B}(B \rightarrow J/\psi f)$

LHCb Measurements

Final state	\bar{B}_s^0	\bar{B}^0
σ	-	$9.60_{-1.70}^{+3.79} \times 10^{-6}$
f_0	$3.40_{-0.16}^{+0.63} \times 10^{-4}$	$< 1.7 \times 10^{-6}$



Quark substructure

- $q\bar{q}$ model

$$|f_0\rangle = \cos \phi |s\bar{s}\rangle + \sin \phi |n\bar{n}\rangle$$

$$|\sigma\rangle = -\sin \phi |s\bar{s}\rangle + \cos \phi |n\bar{n}\rangle,$$

$$\text{where } |n\bar{n}\rangle \equiv \frac{1}{\sqrt{2}} (|u\bar{u}\rangle + |d\bar{d}\rangle).$$

- tetraquark model

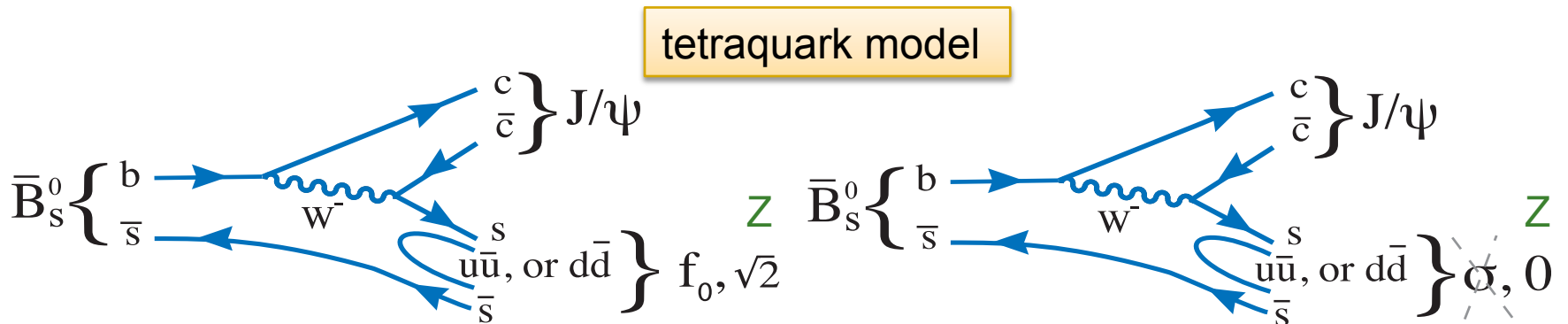
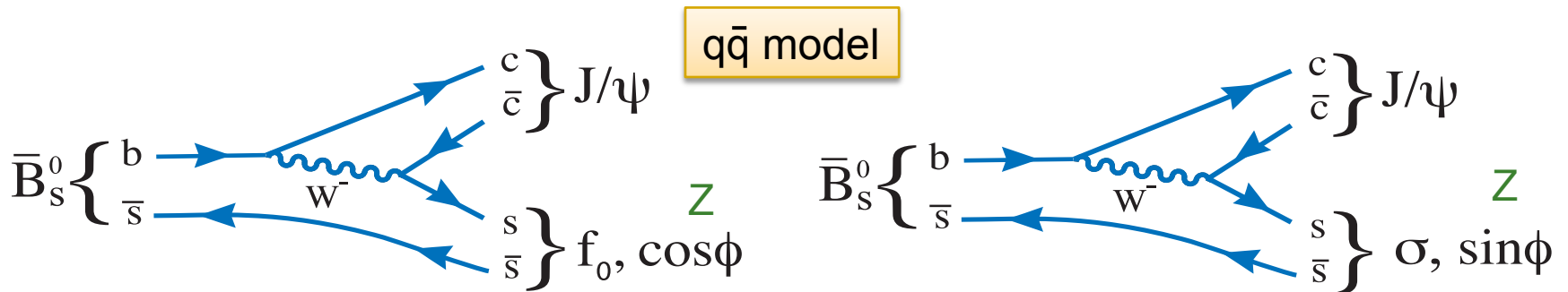
$$|f_0\rangle = \frac{1}{\sqrt{2}} ([su][\bar{s}\bar{u}] + [sd][\bar{s}\bar{d}]), \quad |\sigma\rangle = [ud][\bar{u}\bar{d}]$$



B_s decay diagrams

$$\Gamma(\bar{B}_s^0 \rightarrow J/\psi f) = C \left| F_{B_s}^f \left(m_{J/\psi}^2 \right) \right|^2 |V_{cs}|^2 \Phi Z^2$$

↑ form factor
 ↑ phase space
 ↑ coupling



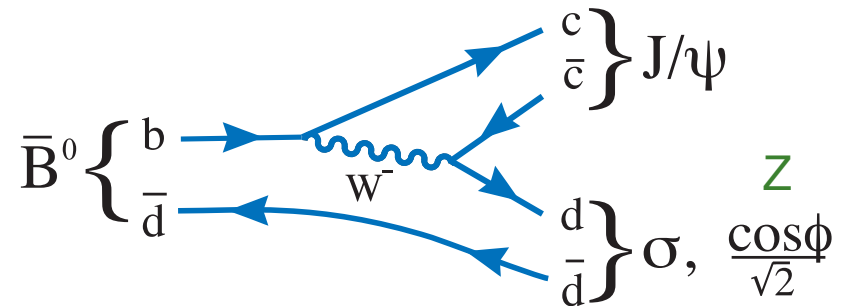
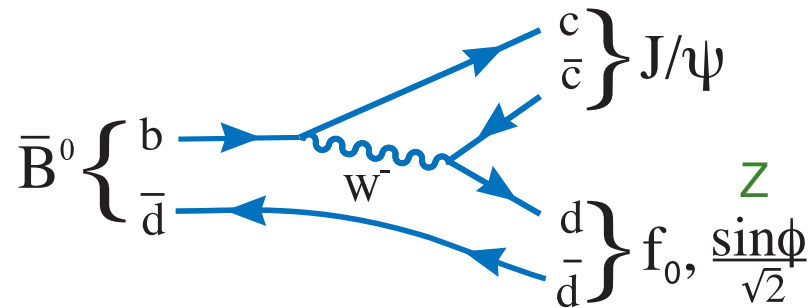
First prediction: If σ is a tetraquark it will not be seen in $B_s \rightarrow J/\psi \sigma$



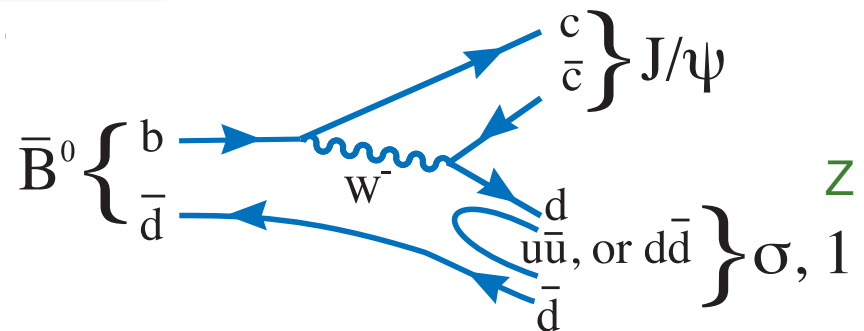
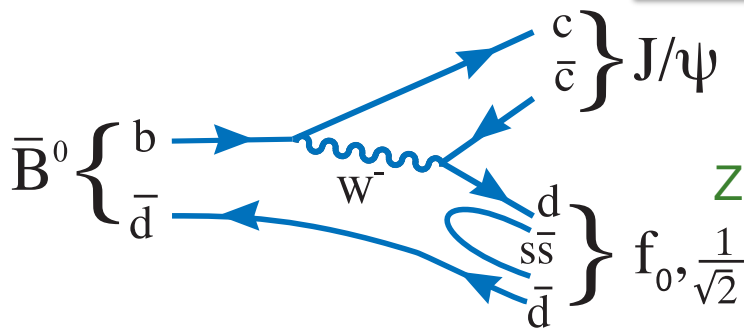
B⁰ decay diagrams

$$\Gamma(\bar{B}^0 \rightarrow J/\psi f) = C \left| F_{B^0}^f \left(m_{J/\psi}^2 \right) \right|^2 |V_{cd}|^2 \Phi Z^2$$

q \bar{q} model



tetraquark model





Rate ratios

Label	Mode ratio	Rate ratio	$\mathcal{Z}^2 q\bar{q}$	\mathcal{Z}^2 tetraquark
$r_{sf_0}^{0f_0}$	$\frac{\Gamma(\bar{B}^0 \rightarrow J/\psi f_0)}{\Gamma(\bar{B}_s^0 \rightarrow J/\psi f_0)}$	$= \frac{ F_{B^0}^{f_0}(m_{J/\psi}^2) ^2 V_{cd} ^2 \Phi_{B^0}^{f_0}}{ F_{B_s^0}^{f_0}(m_{J/\psi}^2) ^2 V_{cs} ^2 \Phi_{B_s^0}^{f_0}}$	$\frac{1}{2} \tan^2 \phi$	$\frac{1}{4}$
$r_{0\sigma}^{0f_0}$	$\frac{\Gamma(\bar{B}^0 \rightarrow J/\psi f_0)}{\Gamma(\bar{B}^0 \rightarrow J/\psi \sigma)}$	$= \frac{ F_{B^0}^{f_0}(m_{J/\psi}^2) ^2 \Phi_{B^0}^{f_0}}{ F_{B^0}^{\sigma}(m_{J/\psi}^2) ^2 \Phi_{B^0}^{\sigma}}$	$\tan^2 \phi$	$\frac{1}{2}$
$r_{sf_0}^{s\sigma}$	$\frac{\Gamma(\bar{B}_s^0 \rightarrow J/\psi \sigma)}{\Gamma(\bar{B}_s^0 \rightarrow J/\psi f_0)}$	$= \frac{ F_{B_s^0}^{\sigma}(m_{J/\psi}^2) ^2 \Phi_{B_s^0}^{\sigma}}{ F_{B_s^0}^{f_0}(m_{J/\psi}^2) ^2 \Phi_{B_s^0}^{f_0}}$	$\tan^2 \phi$	0
$r_{0\sigma}^{sf_0}$	$\frac{\Gamma(\bar{B}_s^0 \rightarrow J/\psi f_0)}{\Gamma(\bar{B}^0 \rightarrow J/\psi \sigma)}$	$= \frac{ F_{B_s^0}^{f_0}(m_{J/\psi}^2) ^2 V_{cs} ^2 \Phi_{B_s^0}^{f_0}}{ F_{B^0}^{\sigma}(m_{J/\psi}^2) ^2 V_{cd} ^2 \Phi_{B^0}^{\sigma}}$	2	2

Last ratio is independent of model, allows measurement of form factor ratio



Form-factors

- $r_{0\sigma}^{sf_0}$ is independent of whether the states are $q\bar{q}$ or tetraquark, so we determine

$$\frac{|F_{B_s^0}^{f_0}(m_{J/\psi}^2)|}{|F_{B^0}^\sigma(m_{J/\psi}^2)|} = 0.99^{+0.13}_{-0.04}$$

- Li et al suggest using $r_{sf_0}^{s\sigma}$ to measure ϕ in the $q\bar{q}$ model, & compute $|F_{B_s^0}^\sigma(m_{J/\psi}^2)|^2 / |F_{B_s^0}^{f_0}(m_{J/\psi}^2)|^2 = 1$
- Fleischer et al. predict in the tetraquark model that $\mathcal{B}(\bar{B}^0 \rightarrow J/\psi f_0, f_0 \rightarrow \pi^+\pi^-) \sim (1 - 3) \times 10^{-6}$



Form-factors II

- They use $|F_{B^0}^{f_0}(m_{J/\psi}^2)|/|F_{B_s^0}^{f_0}(m_{J/\psi}^2)| = 0.69$ from El-Bennich et al. If a unit form-factor is used then their prediction doubles & the tetraquark nature of the $f_0(980)$ becomes inconsistent with the LHCb upper limit of 1.1×10^{-6} .
- Using the limit on the measured ratio $r_{sf_0}^{0f_0}$, we find $\phi < 29^\circ$ @ 90% cl
- LHCb assumes the similar ratio $|F_{B^0}^{f_0}(m_{J/\psi}^2)|/|F_{B^0}^\sigma(m_{J/\psi}^2)| = 1$, & find $\phi < 31^\circ$ @ 90% cl (using $r_{0\sigma}^{0f_0}$)



Predictions

- Our null prediction for $B_s \rightarrow J/\psi \sigma$ is ameliorated somewhat by the possibility of mixing between the tetraquark states and higher order diagrams. However this leads to rate of no more than 1% of the $B_s \rightarrow J/\psi f_0$ rate
- Mixing of other isoscalar mesons can be measured by just using decays of B^0 & $B_s \rightarrow J/\psi F$, i.e. the mixing angle can be measured from one state alone!
- Thus future data using $B_{(s)} \rightarrow J/\psi F$ decays can tell us about the substructure of isoscalar mesons

The End

For more details see S. Stone &
L. Zhang, PRL 111, 062001 (2013)
arXiv: 1305.6554