Threshold Resummation and Determinations of Parton Distribution Functions

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Threshold Resummation and PDFs

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• An accurate knowledge of parton distribution functions (PDFs) is vital to precision phenomenology.

• e.g. A massive state produced at high rapidity would require knowledge of PDFs at high x, where resummation effects are known to be large.

$$\sigma_{H_1+H_2}(x) = \sum_{a,b} \iiint dx_1 dx_2 dz f_{a/H_1}(x_1) f_{b/H_2}(x_2) \hat{\sigma}_{ab}(z) \delta(x - x_1 x_2 z)$$

- Observables for hadron-initiated processes comprise a convolution of two parts:
 - Parton distribution functions (PDFs) $f_{a/H}(x)$
 - Hard-scattering cross section $\hat{\sigma}_{ab}(x)$
- PDFs are not calculable using perturbation theory; their forms are inferred by comparing data to theoretical predictions of observables.

• $\hat{\sigma}_{ab}$ is calculated using perturbation theory.

- "Threshold regions" are encountered in the kinematics of some observables at higher orders of perturbation theory.
 - The leading order (LO) process takes up all available energy of the partonic system
 - Final state gluons are soft

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 - Final state gluons are soft
- When the kinematics is constrained (as occurs in differential cross sections), logarithms associated with soft gluons become large.

• These "threshold logarithms" appear at every order beyond LO in a predictable manner.

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$$\int_0^1 \, \mathrm{d}z \, z^{N-1} \, \int_0^1 \, \mathrm{d}x \, f(x) \, \int_0^1 \, \mathrm{d}y \, g(y) \, \delta(z - xy) = \\ \int_0^1 \, \mathrm{d}x \, x^{N-1} \, f(x) \, \int_0^1 \, \mathrm{d}y \, y^{N-1} \, g(y) = \tilde{f}(N) \, \tilde{g}(N)$$

• Threshold logarithms in Mellin space manifest in powers of ln N.

Organization of Threshold Logantin	Organization	of	Threshold	Logarithm
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	LL	NLL	
NLO	$\alpha_{S} \ln^{2} N$	$\alpha_{\mathcal{S}} \ln N$	
NNLO	$\alpha_S^2 \ln^4 N$	$\alpha_S^2 \ln^3 N$	
:		:	- N.

• Threshold (or soft-gluon) resummation is summing the logarithms to all orders in $\alpha_S^m \ln^n N$ for all n = 2m (LL), $2m - 1 \le n \le 2m$ (NLL), etc.

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• It is known that this sum is an exponential in Mellin space.

- Two processes being considered:
 - Deep inelastic scattering (DIS): $I + H \rightarrow I + X$
 - ▶ Lepton pair production (LPP): $H_1 + H_2 \rightarrow I + I + X$

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- Two processes being considered:
 - Deep inelastic scattering (DIS): $I + H \rightarrow I + X$
 - Lepton pair production (LPP): $H_1 + H_2 \rightarrow I + I + X$
- These two processes are primary sources of information on PDFs.
 - DIS is used to constrain valence PDFs ($F_2 \sim 4u + d$)
 - LPP is used to constrain antiquark PDFs ($\sigma \sim u\bar{u}, dd$)

DIS Kinematics



- The squared mass of the final hadronic state in DIS is given by $W^2 = M^2 + Q^2 (rac{1}{x} 1)$
- Threshold occurs at $W^2 = M^2$.
- This corresponds to x = 1.



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LPP Kinematics



- x₁ and x₂ are the momentum fractions that the partons take from the parent hadrons.
- Threshold occurs when $x_1x_2 = \tau = \frac{Q^2}{S}$.
- LPP data is often in the form of the x_F distribution, where $x_F = \frac{2p_L}{\sqrt{s}}$.

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LPP Kinematics

• At NLO, x₁ and x₂ are integrated, implying that threshold can occur at many values of x_F.



- Threshold kinematics requires that at large x_F, x₁ is large and x₂ is small. (And vice-versa)
- The PDFs fall rapidly at large x, so the largest contribution comes from the threshold region.
- Therefore, threshold corrections dominate at high $|x_F|$.

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$\frac{\rm NLO+NLL}{\rm NLO}$ for Proton-Proton LPP x_F Distribution

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• Others have found similar results:



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Figure 1 from Owens, J. F.; Accardi, A. & Melnitchouk, W. Global parton distributions with nuclear and finite-Q² corrections.

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• Threshold resummation has a large effect on DIS and LPP calculations.

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- Threshold resummation affects DIS and LPP differently because of their differing kinematics.
 - PDFs are sensitive to resummation effects at different values of x for DIS and LPP.
- Preliminary results of the global fit are currently being investigated.