

The Covariant, time-dependent Aharonov-Bohm Effect*

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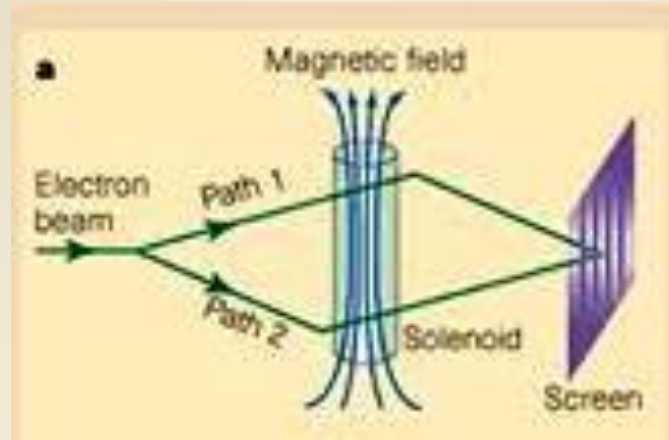
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*D. Singleton and E. Vagenas, *Phys.Lett. B* **723** (2013) 241-244 [[arXiv:1305.1498 \(hep-th\)](#)]

Static Magnetic AB Effect

- The magnetic AB effect is the QM two-slit experiment with an infinite solenoid placed between the slits



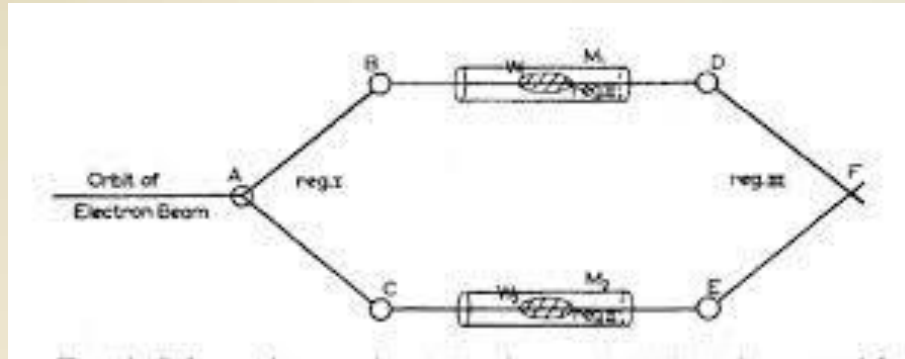
- There is a phase difference between the two paths given by

$$\delta\alpha_B = \alpha_{B_1} - \alpha_{B_2} = \frac{e}{\hbar} \oint_{2-1} \mathbf{A} \cdot d\mathbf{x} = \frac{e}{\hbar} \int \mathbf{B} \cdot d\mathbf{S} = \frac{e}{\hbar} \Phi_0$$

- Leading to a shift in the phase pattern $\Delta x = L\lambda(\delta\alpha_B)/2\pi d$

Static Electric AB Effect

- In the electric AB effect electrons travel along two paths with different potentials and $\mathbf{E}=0$



- When the electrons recombine there is a phase difference

$$\delta\alpha_E = \frac{e}{\hbar} \int_{t_1}^{t_2} \Delta\phi dt = \frac{e}{\hbar} \int_{t_1}^{t_2} \int \mathbf{E} \cdot d\mathbf{x} dt$$

Common description of AB effects

- The common form of the phase shift is

$$\delta(\textit{Phase}) \propto (\textit{Field}) \times (\textit{Area})$$

- For the electric case the “area” has one temporal “leg”

$$\delta\alpha_E = \frac{e}{\hbar} (\mathbf{E} \cdot \Delta\mathbf{x}) \Delta t$$

Covariant versions of AB phase

- In terms of the four-potential the AB phase shift can be written as

$$\delta\alpha_{EB} = \frac{e}{\hbar} \oint A_\mu dx^\mu = \frac{e}{\hbar} \left[\int_{t_1}^{t_2} \Delta\phi dt - \oint \mathbf{A} \cdot d\mathbf{x} \right]$$

- In terms of the **E** and **B** fields the AB phase shift can be written in terms of differential forms, dx^μ , wedge product, \wedge , and the Faraday two form, F .

$$\delta\alpha_{EB} = -\frac{e}{2\hbar} \int F_{\mu\nu} dx^\mu \wedge dx^\nu = \frac{e}{\hbar} \int F$$

Recovery of Static Results

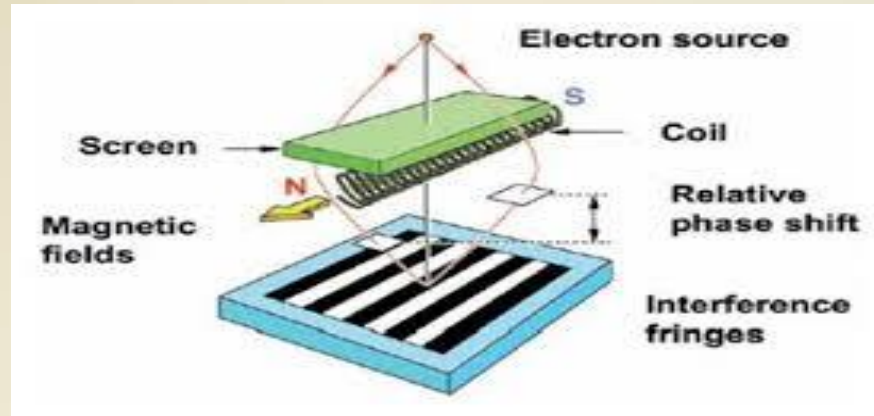
- The static magnetic and electric AB results can be recovered from the Faraday two form expression given above and expanded below.

$$\begin{aligned} F &= -\frac{1}{2}F_{\mu\nu}dx^\mu \wedge dx^\nu \\ &= (E_x dx + E_y dy + E_z dz) \wedge dt + B_x dy \wedge dz + B_y dz \wedge dx + B_z dx \wedge dy \end{aligned}$$

- If $\mathbf{E}=0$ the last three terms are just the infinitesimal magnetic phase shift.
- If $\mathbf{B}=0$ and there is no time dependence ($\partial_t \mathbf{A}=0$) the first three terms are the infinitesimal electric phase shift.

Time dependent AB phase 1.

- These two covariant expressions can be used to study the time-dependent AB effect e.g. a solenoid with time varying flux ($\phi=0$ for this situation)



- The Faraday two form version of the phases gives exactly cancelling results i.e.

$$\frac{e}{\hbar} \int [B_x dy \wedge dz + B_y dz \wedge dx + B_z dx \wedge dy] = \frac{e}{\hbar} \int \mathbf{B}(\mathbf{x}, t) \cdot d\mathbf{S}$$

Cancels

$$\frac{e}{\hbar} \int [(E_x dx + E_y dy + E_z dz) \wedge dt] = -\frac{e}{\hbar} \oint \mathbf{A} \cdot d\mathbf{x} = -\frac{e}{\hbar} \int \mathbf{B}(\mathbf{x}, t) \cdot d\mathbf{S}$$

$$\mathbf{E} = -\partial_t \mathbf{A}$$

Time dependent AB phase 2

- The expression four-vector potential expression gives the same result. Consider a time dependent current and vector potential

$$\mathbf{A} = \frac{kI(t)}{r} \hat{\theta}$$

- Taylor expand the current

$$I(t_i + \Delta t/2) \approx I(t_i) + I'(t_i) \frac{\Delta t}{2} + \dots$$

- This gives two contributions to the phase difference

$$\delta\alpha_B = \frac{ekI(t_i)\Delta\theta}{\hbar}$$

$$\delta\alpha_{A(t)} = \frac{ekI'(t)\Delta t\Delta\theta}{2\hbar}$$

Time dependent AB phase 2

- Now due to $\mathbf{E} = -\partial_t \mathbf{A}$ there is an electric field outside the solenoid which contributes to the phase difference

$$\mathbf{E} = -\frac{\partial \mathbf{A}}{\partial t} = -\frac{kI'(t)}{r} \hat{\theta}$$

$$\mathbf{a} = \frac{e\mathbf{E}}{m} = -\frac{ekI'(t)}{mr} \hat{\theta}$$

- This leads to a change in distance of

$$\Delta d = \frac{1}{2} a \Delta t^2 = -\frac{ekI'(t)}{2mr} \frac{r \Delta \theta}{v} \Delta t = -\frac{ekI'(t) \Delta \theta \Delta t}{2mv}$$

- And a phase difference of

$$\delta \alpha_{E-field} = \frac{\Delta d}{\lambda/(2\pi)} = -\frac{ekI'(t) \Delta t \Delta \theta}{2\hbar}$$

- Which cancels the $\delta \alpha_{A(t)}$ contribution from the previous slide.

Prediction for time dependent AB phase

- Both the Faraday two-form expression and the four-vector expression for the phase difference give no *time dependent* phase shift due to the time variation of the flux in the solenoid.
- This is easier to calculate/see using the Faraday two-form expression.
- The four-vector expression shows this result is due to a cancellation between the usual AB phase and a non-AB phase coming from the electric field external to the solenoid.

Experimental Status of Time-Dependent AB Effect

- Not much work has been done on the time-dependent AB effect

B. Lee, E. Yin, T. K. Gustafson, and R. Chiao, "Analysis of Aharonov-Bohm effect due to time-dependent vector potentials", *Phys. Rev. A* **45**, 4319 (1992).

A. N. Ageev, S. Yu. Davydov, and A. G. Chirkov, "Magnetic Aharonov-Bohm Effect under Time-Dependent Vector Potential" *Tech. Phys. Lett.* **26**, 392 (2000).

- These two papers predict a time-changing phase shift but ignore the external electric field.
- The second paper reports on an experiment that did find no time-dependent phase shift but this was attributed to a bad choice of parameter (time rate of change of \mathbf{A} was too rapid).

Summary and Conclusions

- Two covariant expressions for AB phase difference were given.
- These expressions allow one to study the time-dependent AB effect.
- The result for a time varying version of the usual AB set-up is a prediction of no time dependent phase shift.
- This is due to a cancellation between a standard AB phase shift and a non-AB shift due to the external electric field.

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