

Extracting the Heavy-Quark Potential from Bottomonium Observables in Heavy-Ion Collisions

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Outline

X. Du, S. Liu, R. Rapp, submitted to PLB

- Our Motivation: Learning QCD Force
 - Why quarkonium? What to expect from it?
- Our Tools: Quarkonium Transport Approach
 - Transport equation
 - Key coefficients, parameters and inputs
 - Probing strong/QCD force/potential with quarkonium
- We've learned: Quarkonium Production in A-A Collisions
 - J/ψ production: $R_{AA}(N_{part})$, $R_{AA}(p_T)$, $v_2(p_T)$, $R_{AA}(\sqrt{s})$, ...
 - $\Upsilon(1S,2S,3S)$ production : $R_{AA}(N_{part})$, $R_{AA}(p_T)$, $v_2(p_T)$, ...
 - Comparison between charmonium and bottomonium: $R_{AA}(\sqrt{s})$, ...
 - $\Upsilon(1S)$ as a probe of in-medium heavy quark-antiquark potential
- We are learning: Extraction of In-Medium Heavy Quark-Antiquark Potential
 - Joint study of heavy quark (open heavy flavor)/quarkonium (hidden heavy flavor)
 - Selecting relevant data points
 - Indicating a strong potential
- Summaries

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Quarkonium as Probe of Potential/Force

Heavy quarkonia ($q\bar{q}$ bound states), J/ψ , $\psi(2S)$, $\Upsilon(1S)$, $\Upsilon(2S)$,as probes:

1. Hard scale ($m \gg T_{\text{QGP}}$): pQCD baseline for primordial production
2. (Ultra-)Soft scale ($v^2 \ll v \ll 1$): nonrelativistic potential picture works
3. Large binding ($E_{\text{BINDING}} \sim mv^2 > T_c$): survive in QGP
4. Various species: bound at different parts of potential

Expect to be ideal for probing:

QCD force / $q\bar{q}$ potential in medium

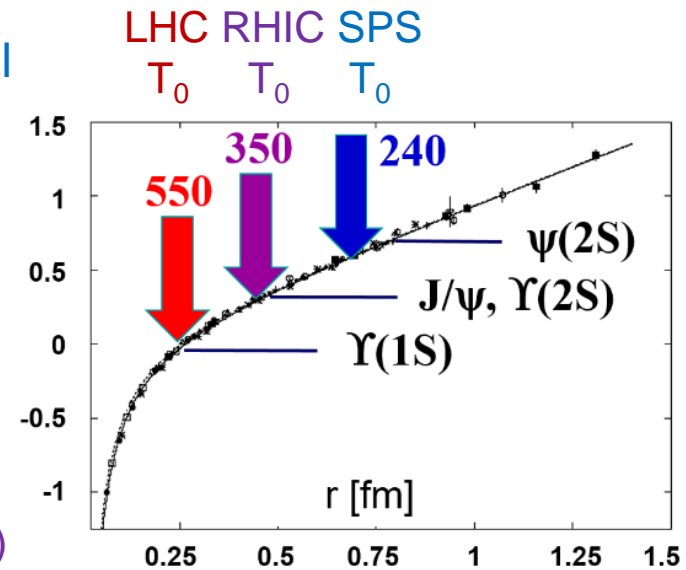
Observables:

Nuclear Modification Factor (R_{AA})

$$R_{AA} = \frac{N_Y^{AA}}{N_{coll} N_Y^{pp}}$$

Elliptic Flow (v_2)

$$\frac{d^2 N}{d^2 p_T} = \frac{1}{2\pi} \frac{dN(p_T)}{p_T dp_T} (1 + 2v_2(p_T) \cos(2\phi) + \dots)$$



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Transport Equation Approach

Kinetic Rate Equation

$$\frac{dN_{\Psi}(\tau)}{d\tau} = -\Gamma_{\Psi}(T(\tau)) [N_{\Psi}(\tau) - N_{\Psi}^{\text{eq}}(T(\tau))]$$

Transport Coefficients

- Reaction Rates

$\Gamma_{\Psi}(T(\tau))$



NLO Quasi-Free
LO Gluo-Diss

- Equilibrium Limits

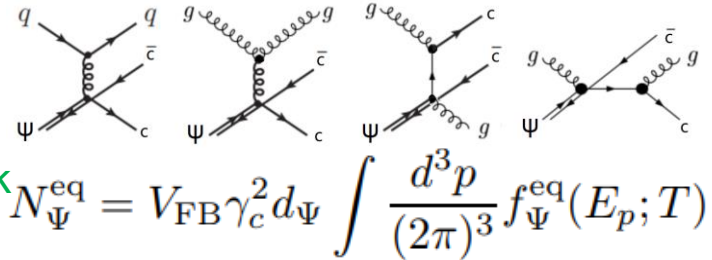
$N_{\Psi}^{\text{eq}}(T(\tau))$



From Heavy quark
conservation

Primordial

Regeneration



Key Parameters

- Coupling α_s



Affects Reaction Rates

Fixed from Previous Calculations compared with data

- Thermal Relaxation Time



Modifies Equilibrium Limit

Extracted from Heavy Quark Diffusion Simulations

Key Inputs

- In-Medium $Q\bar{Q}$ Potential/Binding Energy

- Heavy Quark/-onium pp Cross Section -> fugacity factor γ_c

- Initial State Effects (nPDF)

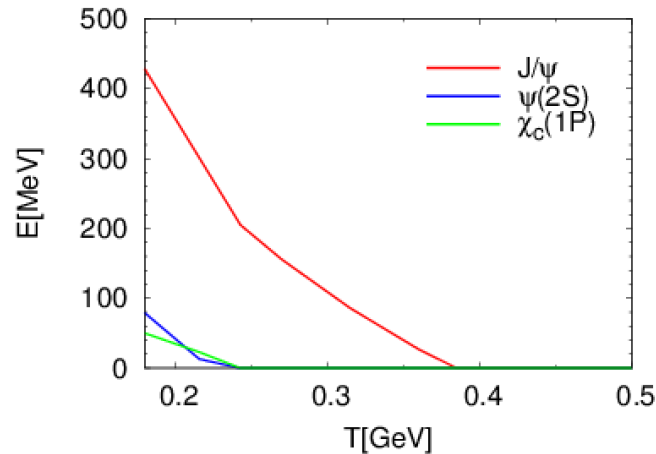
- Fireball Evolution

From Strong Force to Observables

In-Medium $q\bar{q}$ Potential



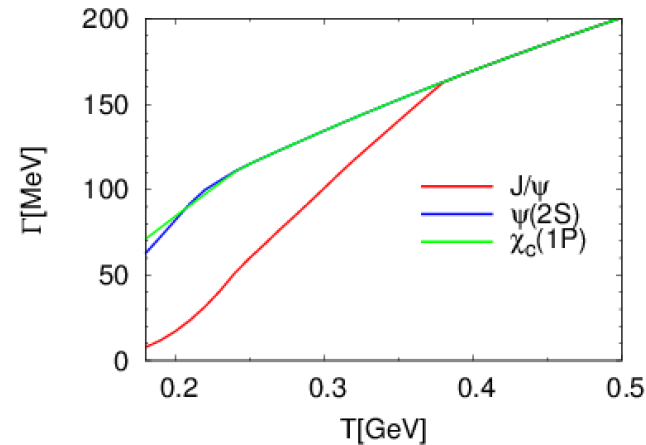
Quarkonium Binding Energies



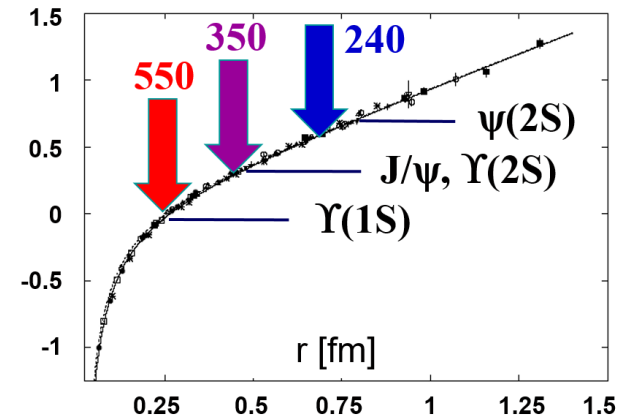
Transport Model \rightarrow Observables



Reaction Rates $\Gamma_{\Psi}(T(\tau))$



- **Hierarchy:**
 - \rightarrow Different Melting temperatures for J/ψ , $\psi(2S)$,
 - Sequential Suppression ...
 - Sequential Regeneration ? ...
- **Probing In-Medium Potential**

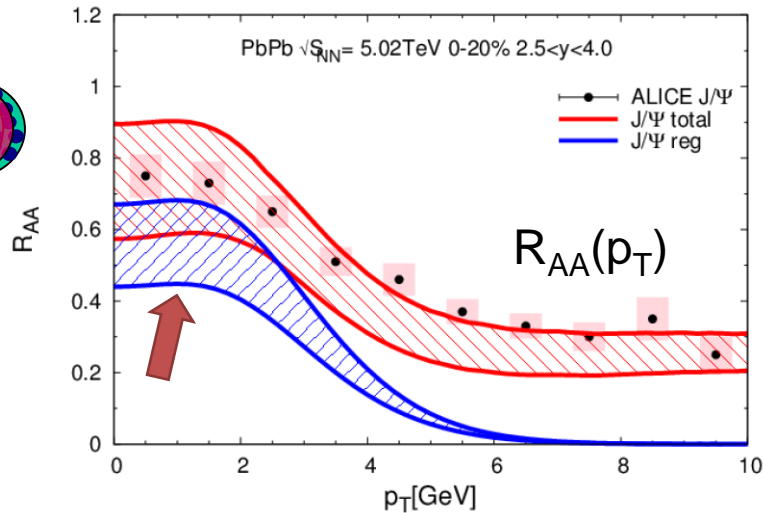
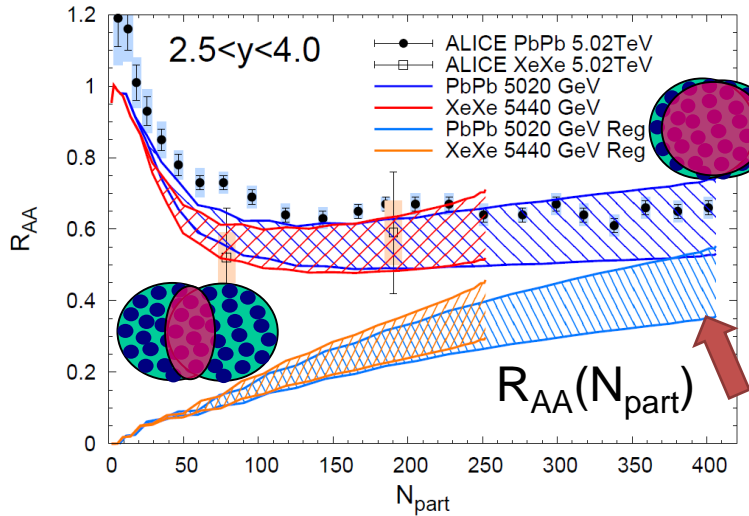


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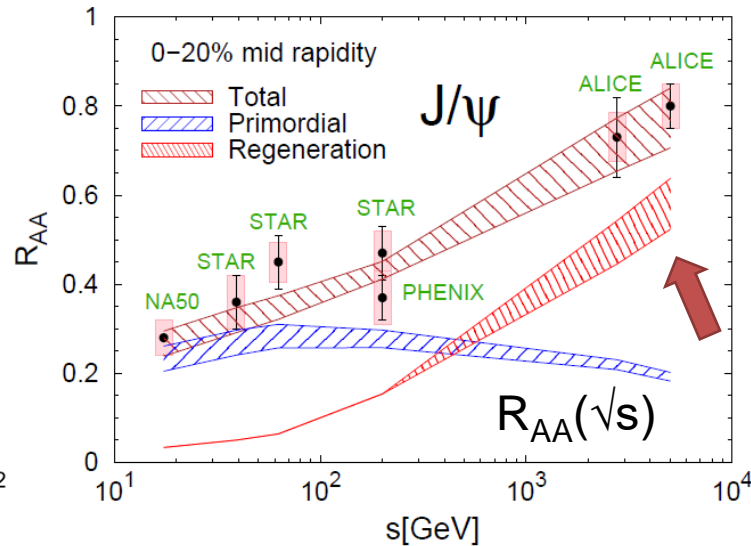
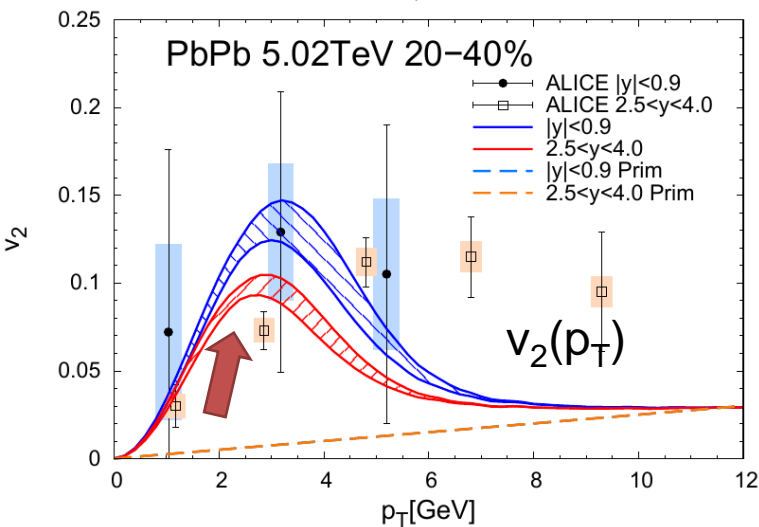
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J/ψ Production in AA

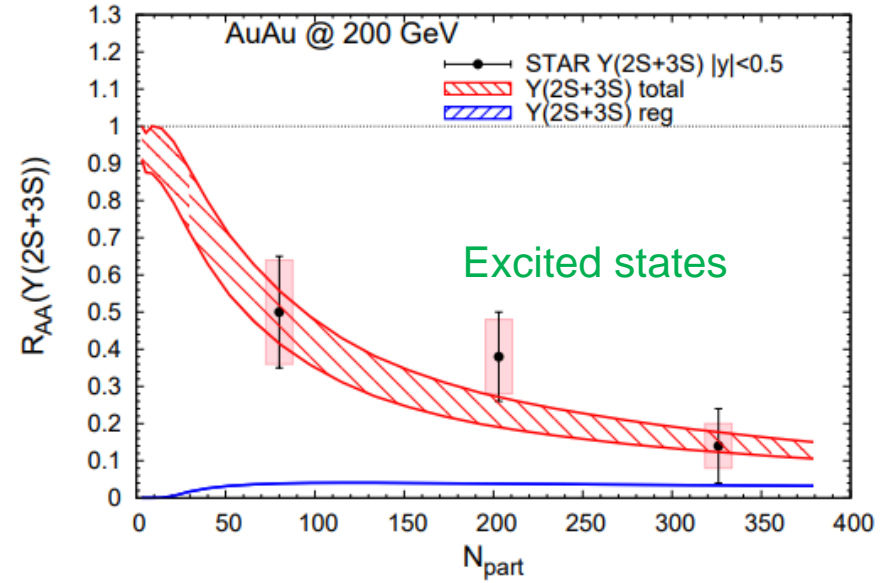
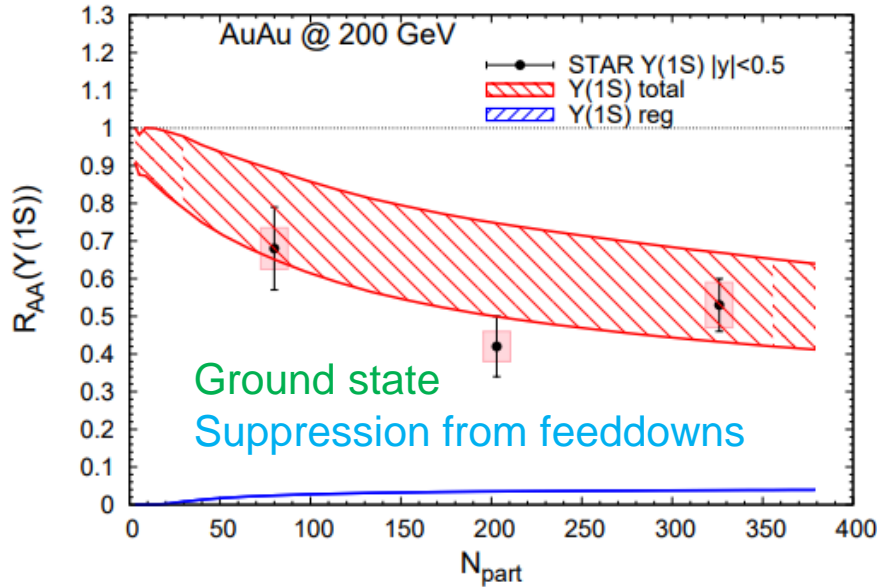


See:
[PLB785 \(2018\) 419](#)
[ALI-PREL-126572](#)
[PRL119 \(2017\) 242301](#)
 ...



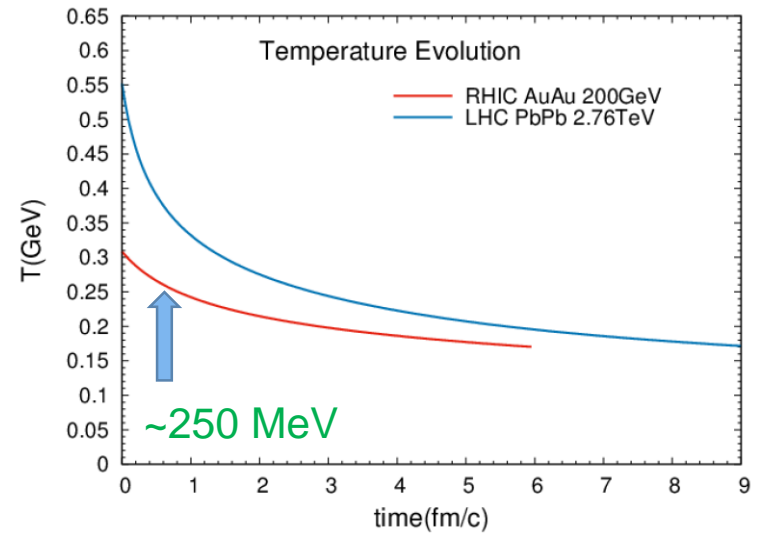
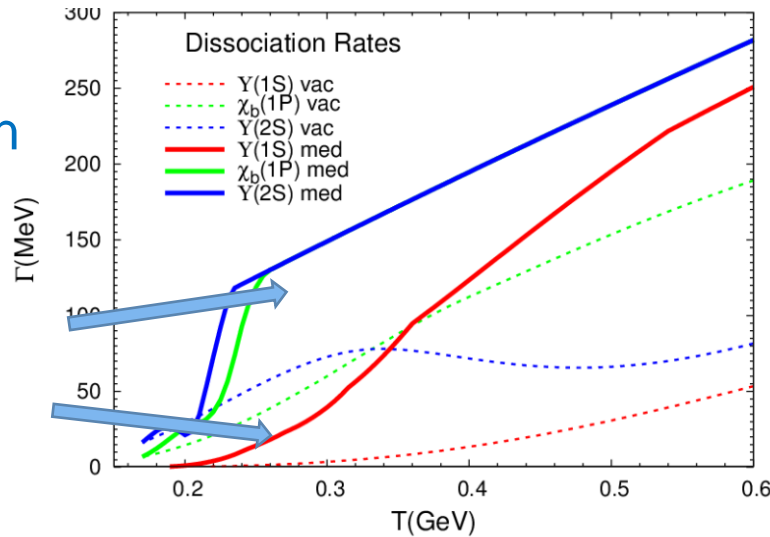
Charmonium
 Large
 Regeneration

Bottomonium $R_{AA}(N_{part})$ at RHIC

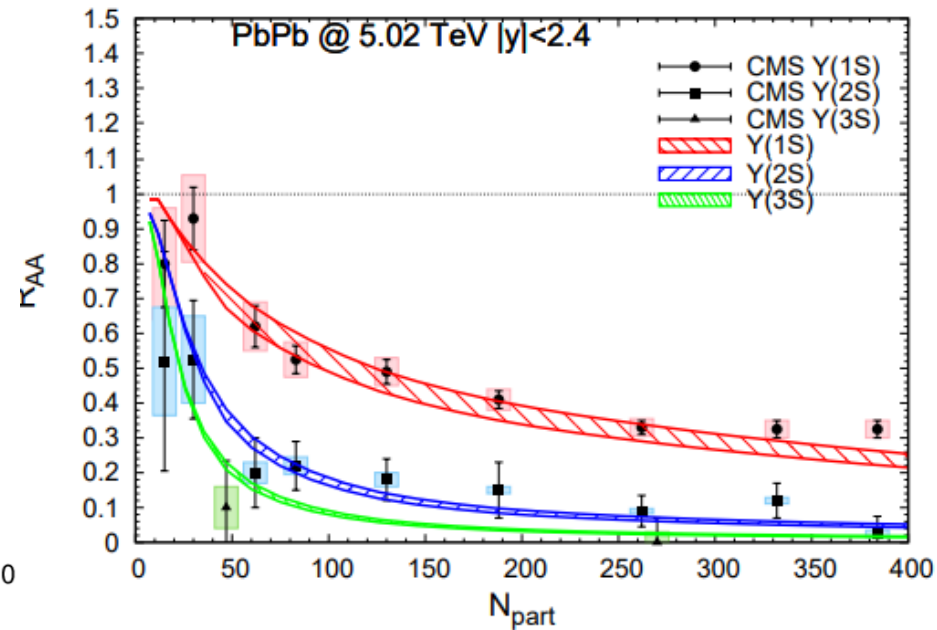
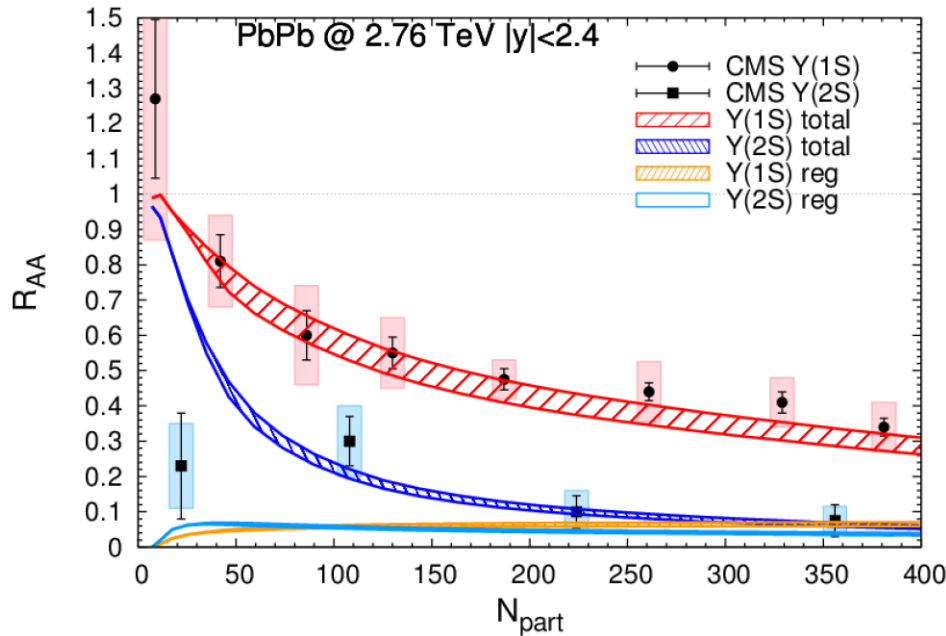


Sequential
Suppression

Excited states
Ground state

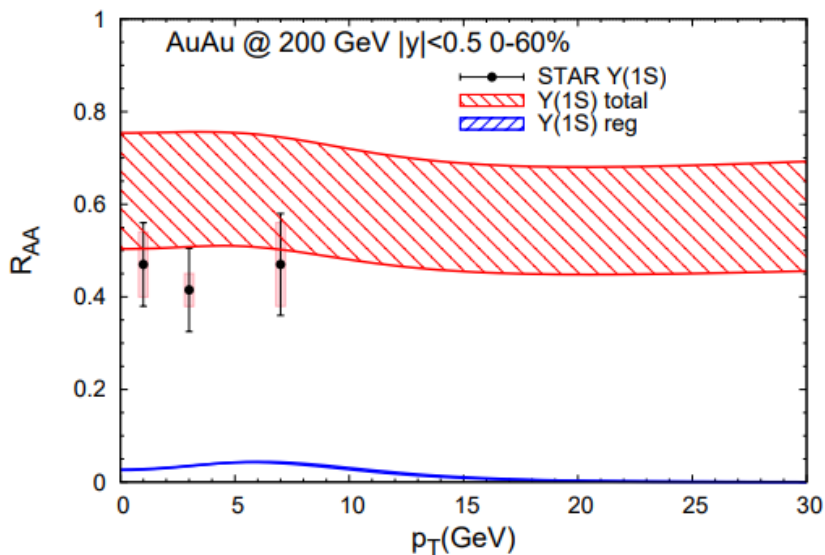
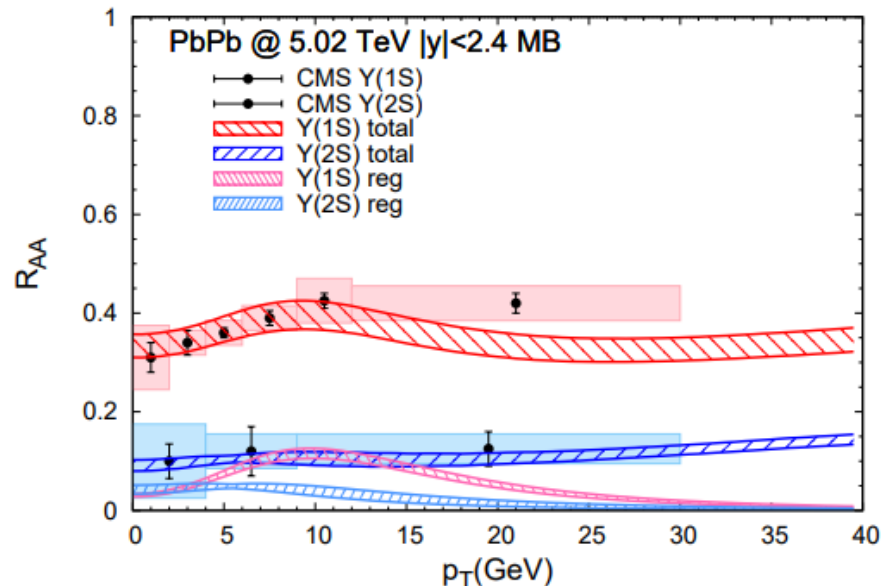
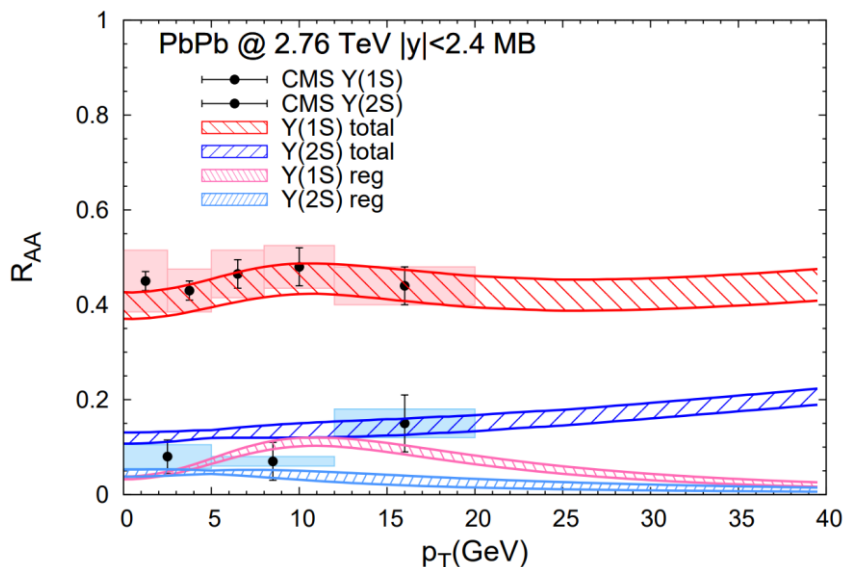


Bottomonium $R_{AA}(N_{part})$ at the LHC



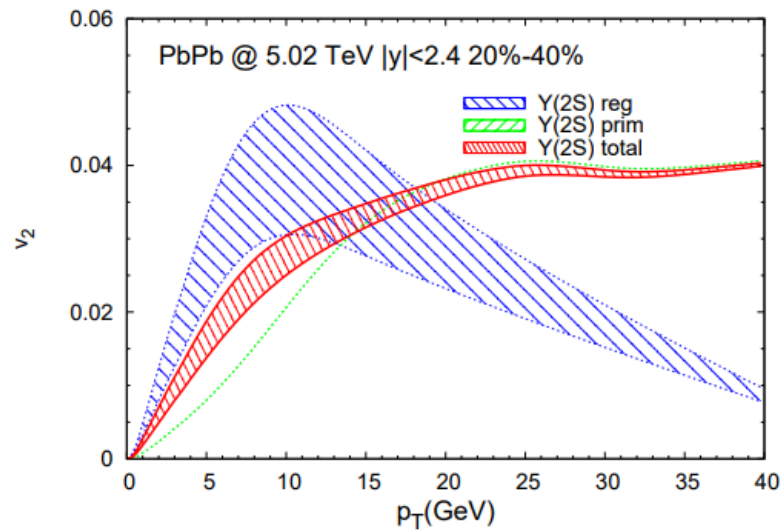
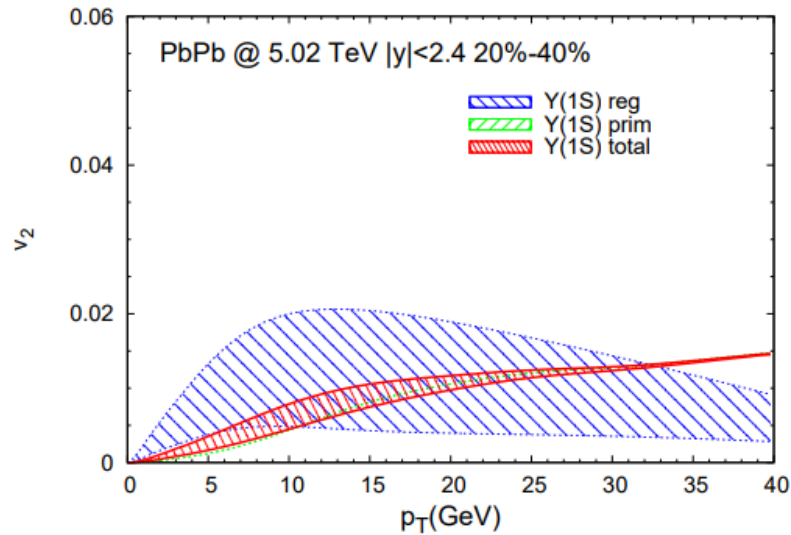
- Sequential suppression
- Direct $\Upsilon(1S)$ suppression, small regeneration
- Regeneration may not be negligible for $\Upsilon(2S)$

Bottomonium $R_{AA}(p_T)$



- No peak at low- p_T
- Non-thermalized b-quarks induces small p_T -dependence for bottomonium

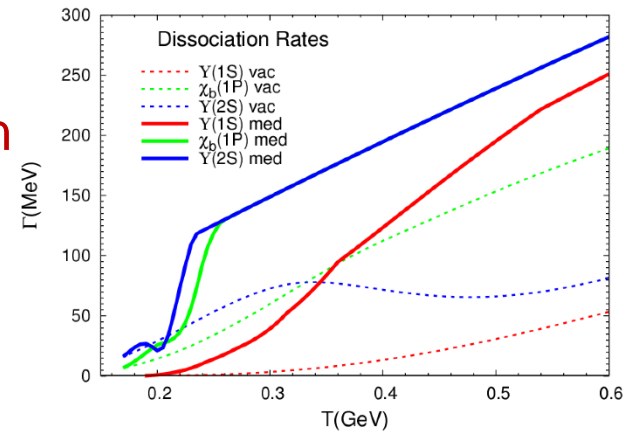
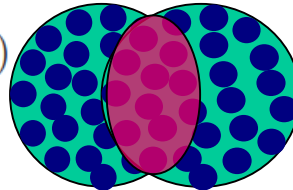
Bottomonium $v_2(p_T)$ at LHC



elliptic flow v_2 :

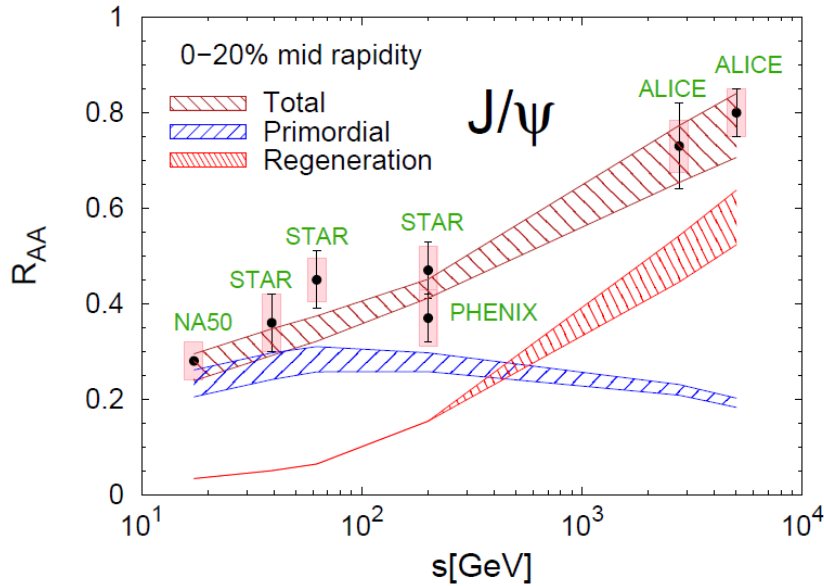
Anisotropy of production

$$\frac{d^2 N}{d^2 p_T} = \frac{1}{2\pi} \frac{dN(p_T)}{p_T dp_T} (1 + 2v_2(p_T) \cos(2\phi) + \dots)$$

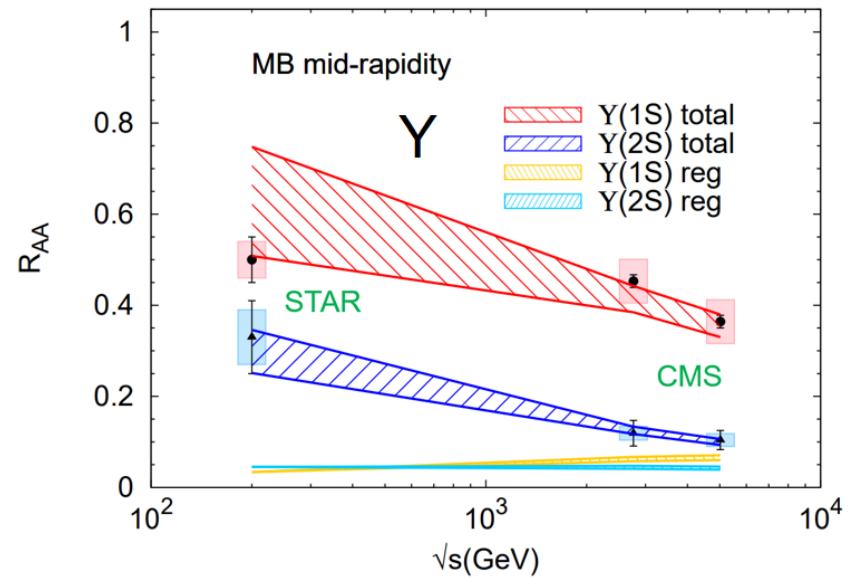


Both destruction and coalescence occur early for $\Upsilon(1S)$, later for $\Upsilon(2S)$

Difference of Charmonium/Bottomonium



R. Rapp, X. Du, NPA967 (2017) 216



X. Du, M. He, R. Rapp, X. Du, PRC96 (2017)

Charm/Bottom-onium difference:

Charmonium:

→ Large regeneration

Bottomonium:

→ Small regeneration

J/ψ and $Y(2S)$:

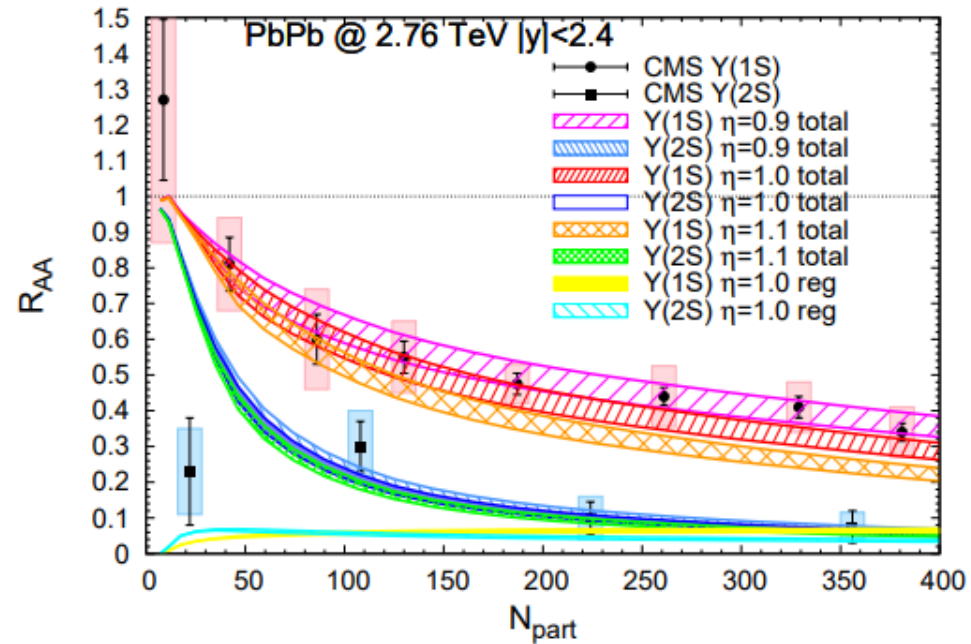
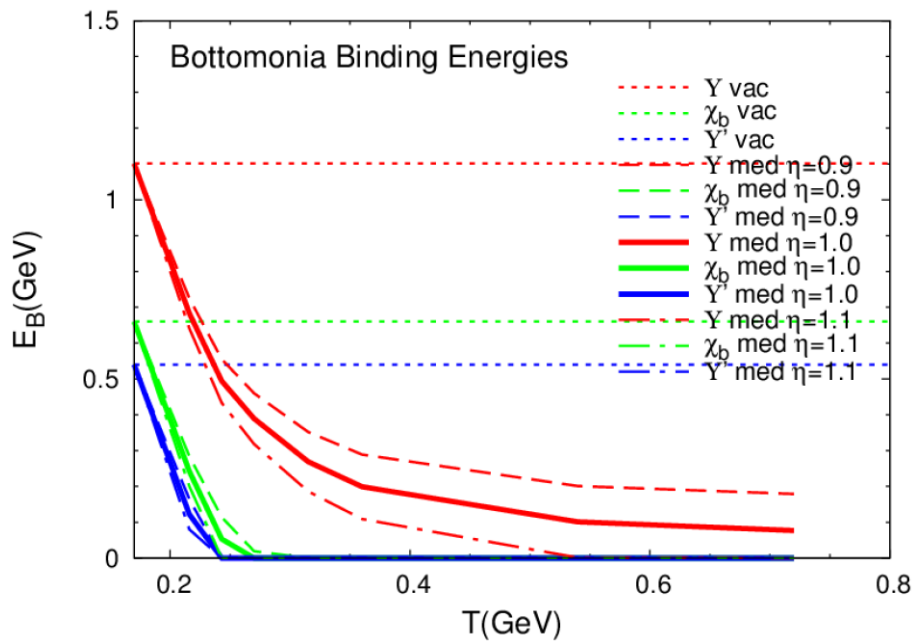
Similar binding energies

BUT

Different excitation functions

→ Due to Large regeneration for J/ψ

$\Upsilon(1S)$ as a Probe of In-Medium Potential



Again: Potential \rightarrow Binding Energies \rightarrow Reaction Rates $\rightarrow R_{AA}$ Observables

- Significant binding energy dependence of the $\Upsilon(1S)$ R_{AA}
- Small regeneration (clean probe)

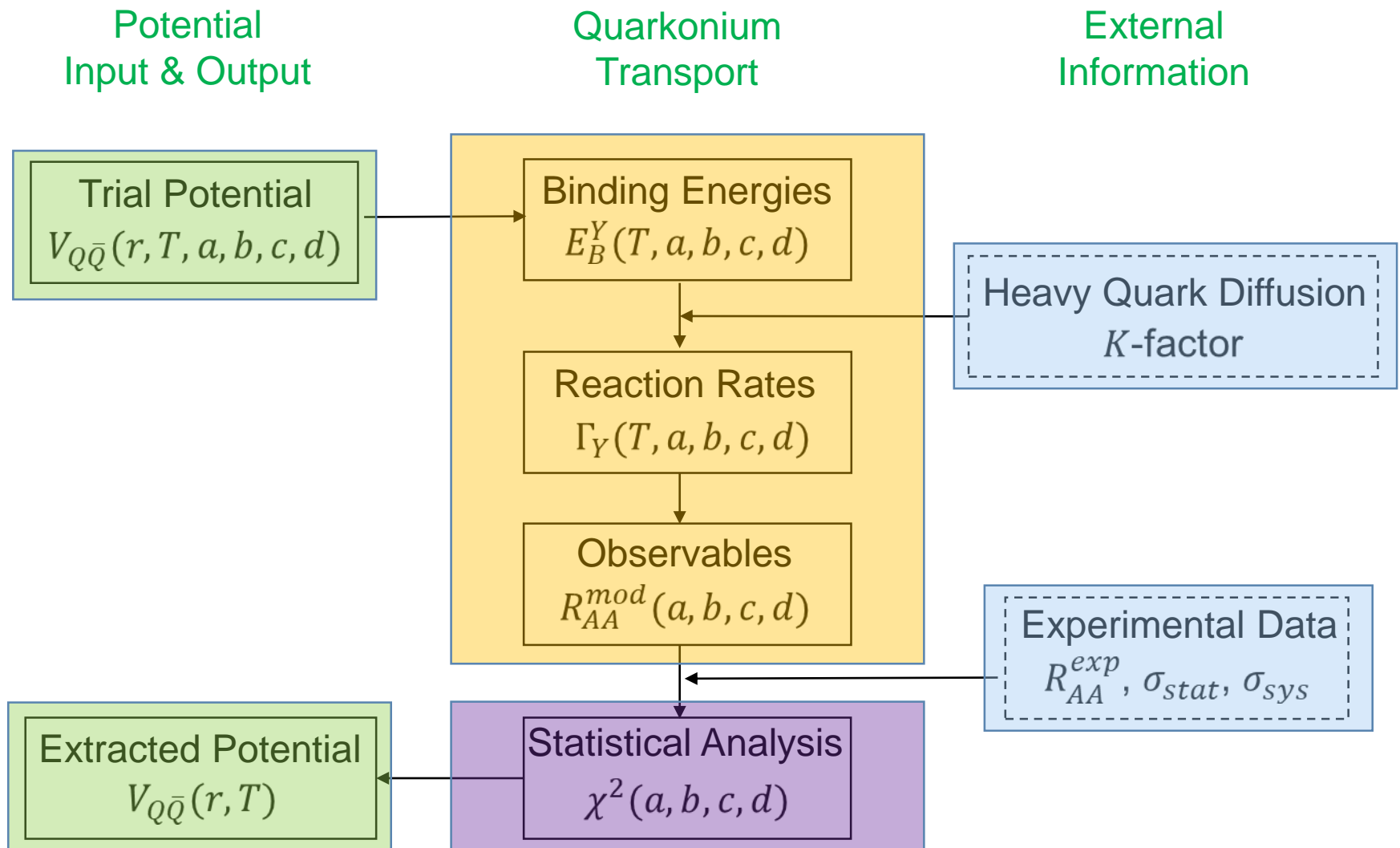
$\Upsilon(1S)$ is ideal for probing in-medium $q\bar{q}$ potential

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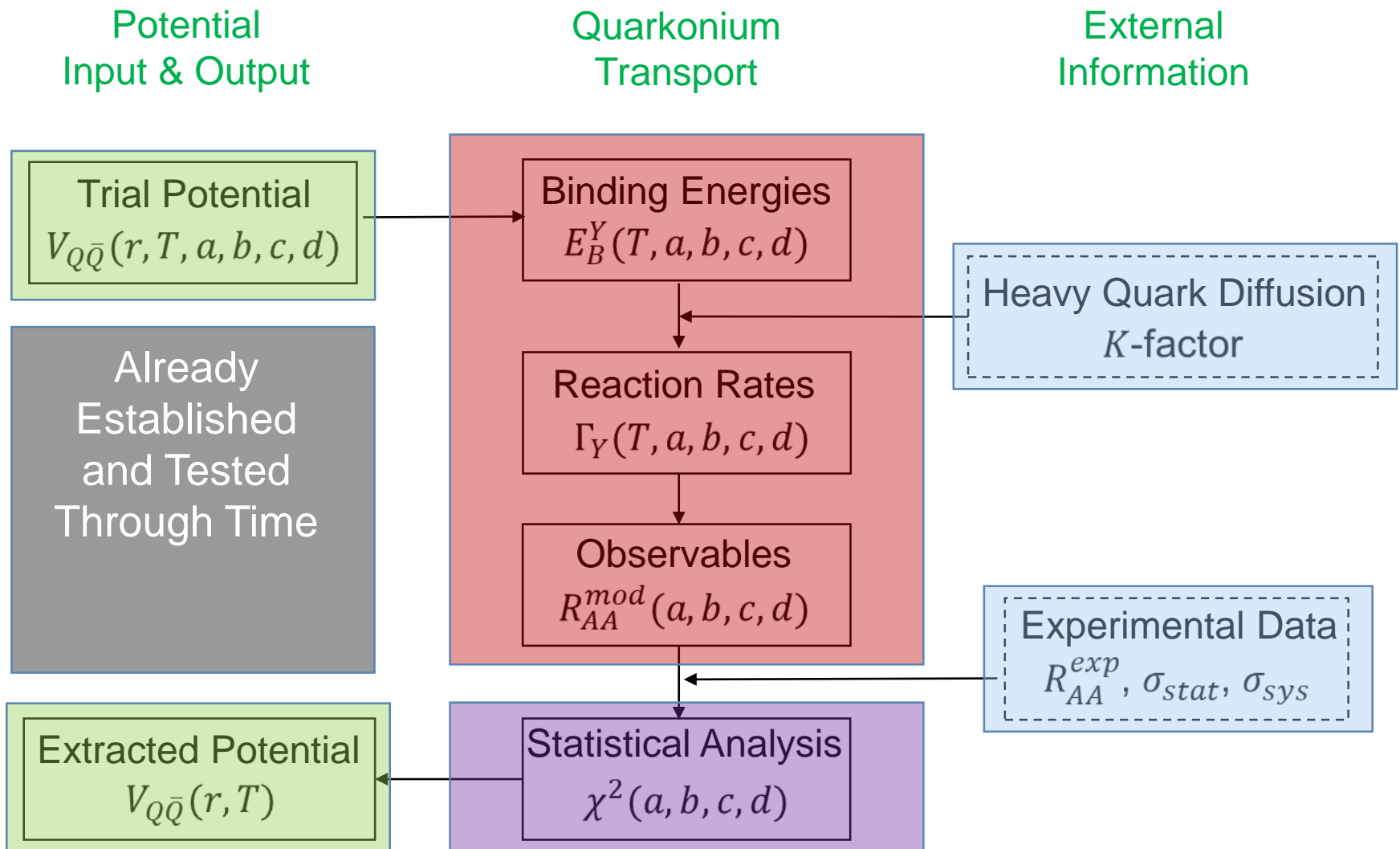
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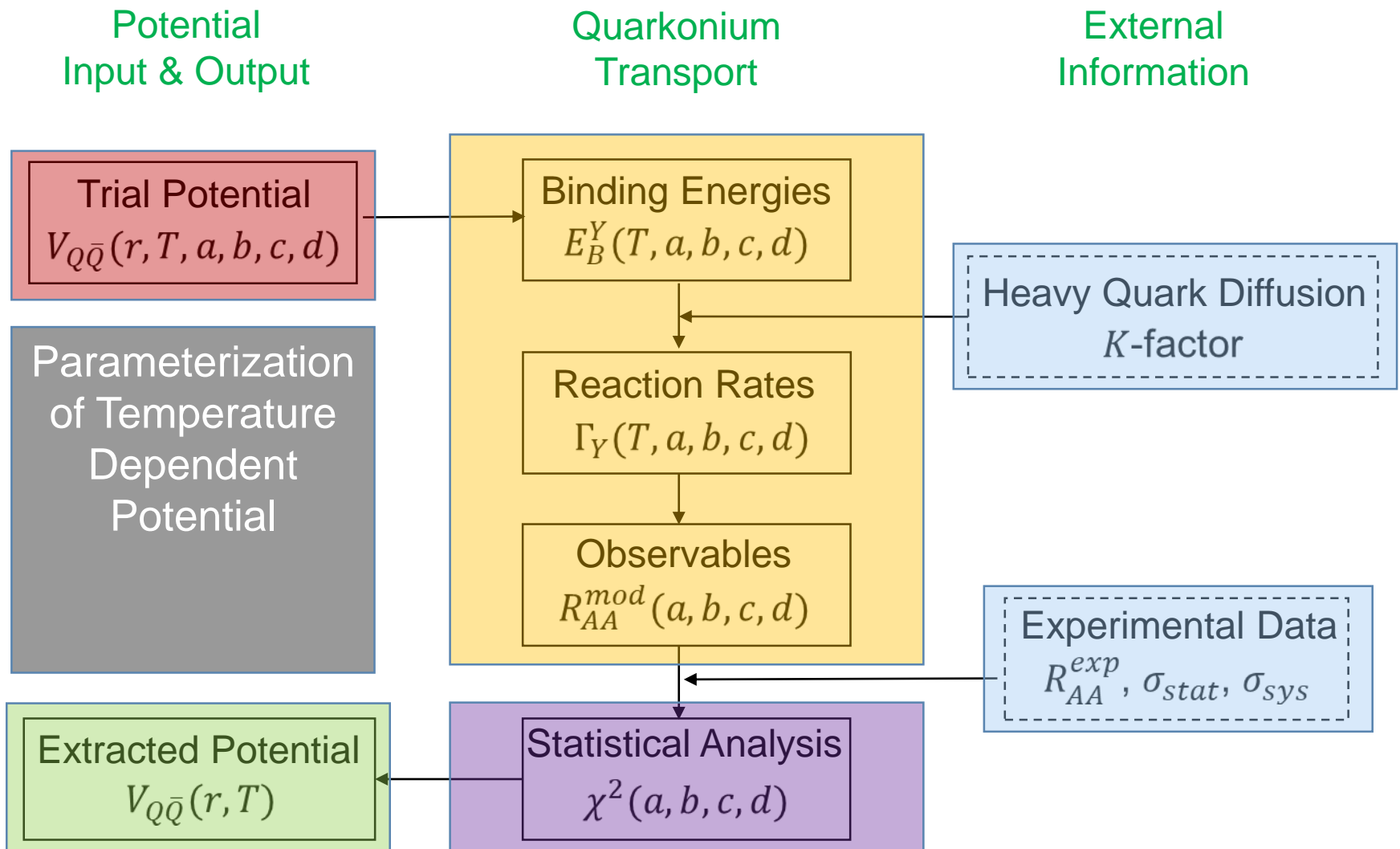
Extracting In-Medium Potential



Extracting In-Medium Potential



Extracting In-Medium Potential



Trial Potential

$$V_{Q\bar{Q}}(r) = -\frac{4}{3}\alpha_s \left(\frac{e^{-\overbrace{m_D} r}}{r} + m_D \right) + \min \left(\sigma r, \frac{\sigma}{\overbrace{m_S}} \right)$$

Coulomb

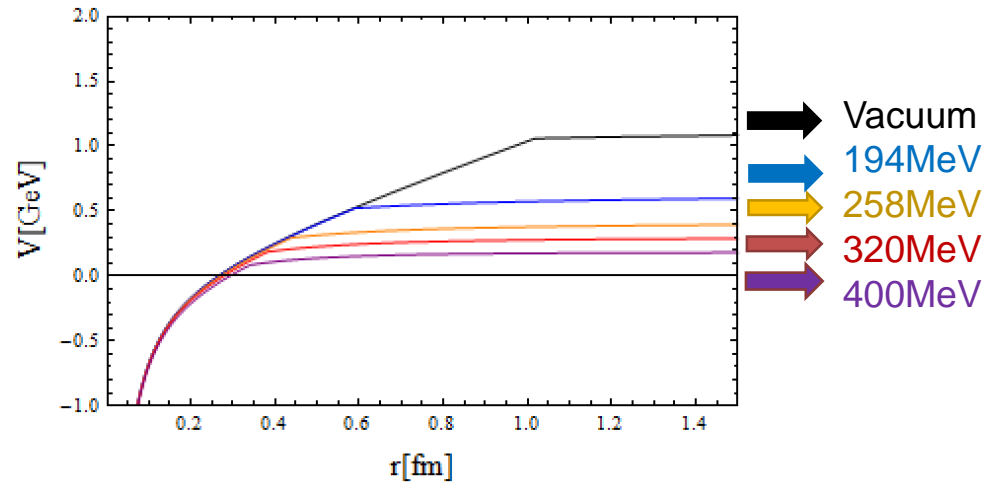
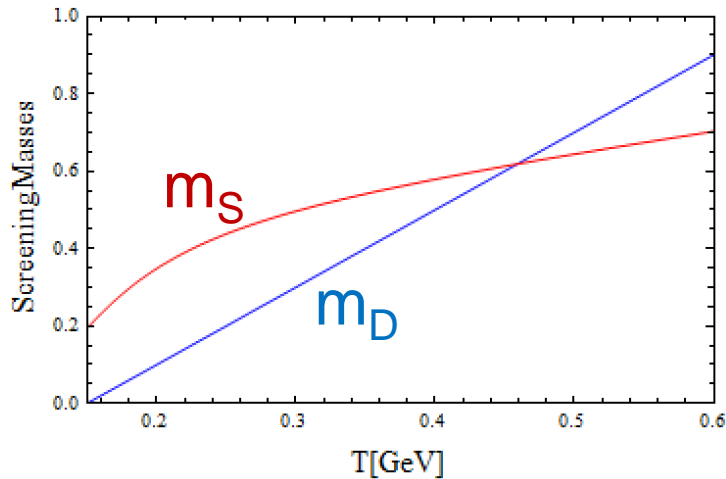
Screening
Masses

$$m_D = a(T - T_0)$$

$$m_S = m_v + b \left(\sqrt{(T - T_0)^2 + (dT_0)^2} - dT_0 \right) - c \left(\sqrt{(T - T_0)^2 + (dT_0)^2} - (T - T_0) - dT_0 \right)$$

String

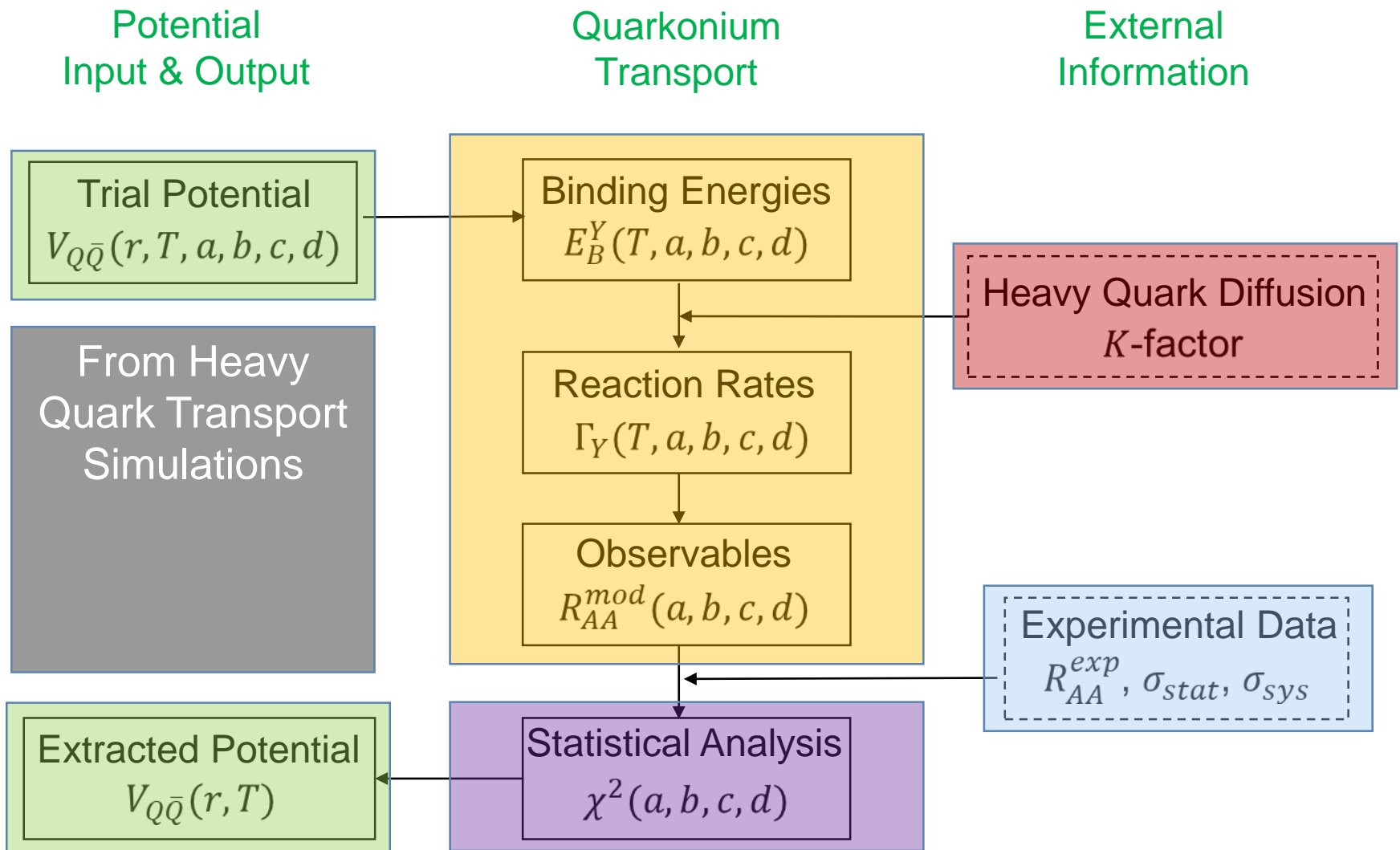
α_s, σ
Fit to vacuum
spectroscopy



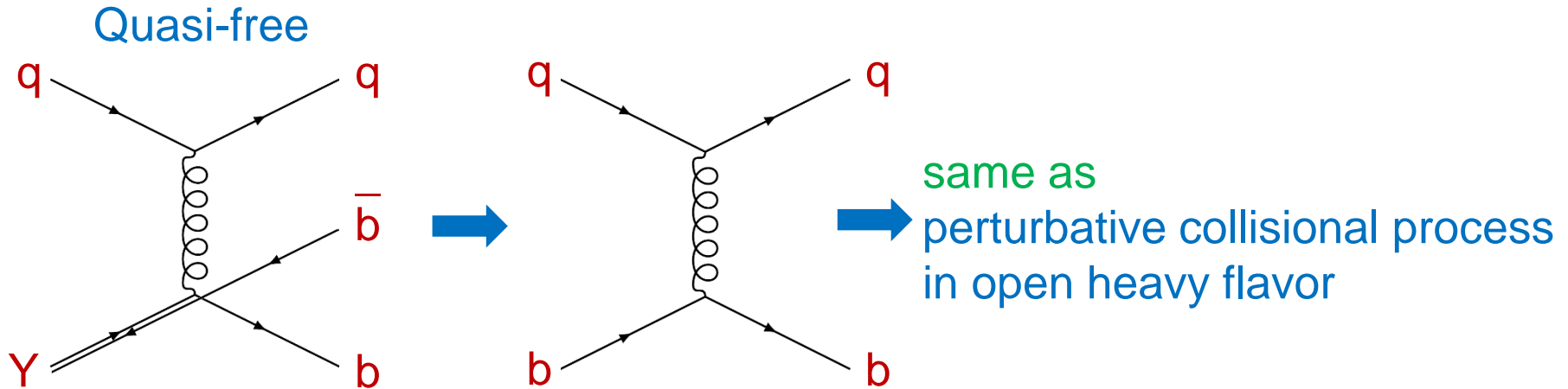
- Vacuum
- 194 MeV
- 258 MeV
- 320 MeV
- 400 MeV

In relevant parameter space, sensitive parameters: a ($\sim m_D/T$) & b ($\sim m_S/T$ at high-T)

Extracting In-Medium Potential



K-factor



K-factor:

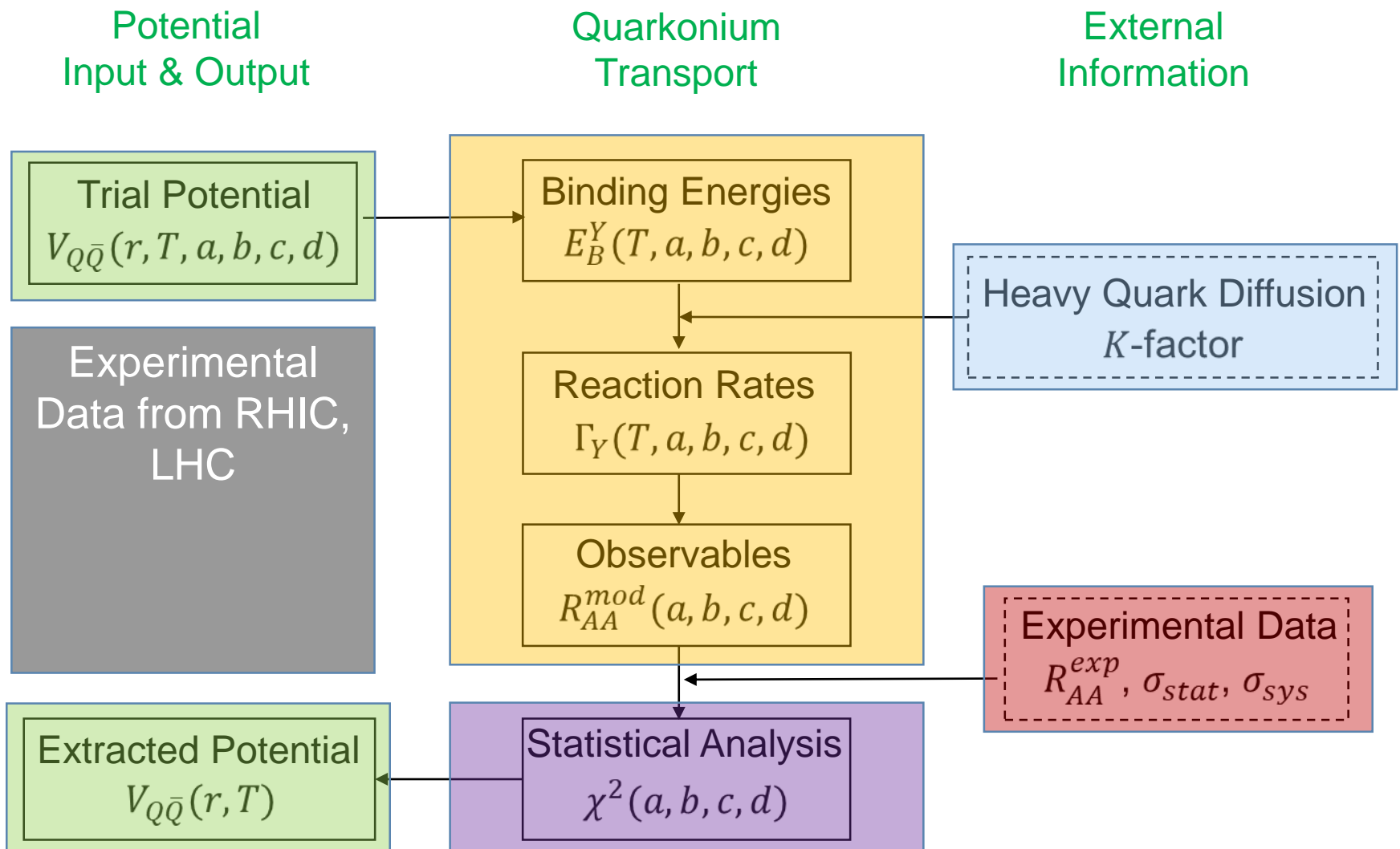
- Multiplication to pQCD calculated Reaction Rates
- Accounts for nonperturbative effects
- Open heavy-flavor suggests **K-factor** > 5

R. Rapp, et al, NPA979 (2018) 21-86
S. Cao, et al, arXiv: 1809.07894
R. Rapp, arXiv: 1901.06440

$D_s(2\pi T)$:

- pQCD: ~ 30
- Required by phenomenology: $2\sim 5$

Extracting In-Medium Potential



Experimental Data

Data selection (Bottomonium data):

In total : 53 R_{AA} data points from RHIC to LHC energies

• 193 GeV U-U	$ y < 1.0$	Y(1S), Y(1S+2S+3S)
• 200 GeV Au-Au	$ y < 0.5$	Y(1S), Y(2S+3S), Y(1S+2S+3S)
• 2.76 TeV Pb-Pb	$ y < 2.4$	Y(1S), Y(2S)
• 2.76 TeV Pb-Pb	$2.5 < y < 4.0$	Y(1S)
• 5.02 TeV Pb-Pb	$ y < 2.4$	Y(1S), Y(2S), Y(3S)
• 5.02 TeV Pb-Pb	$2.5 < y < 4.0$	Y(1S)

Multiple Collision Systems, Multiple Rapidities, Multiple Species, ...

$R_{AA}(N_{part})$:

Most relevant data
for probing force

STAR Collaboration, L. Adamczyk et al., Phys. Rev. C 94(2016) 064904.

STAR Collaboration, Z. Ye, Nucl. Phys. A 967 (2017) 600.

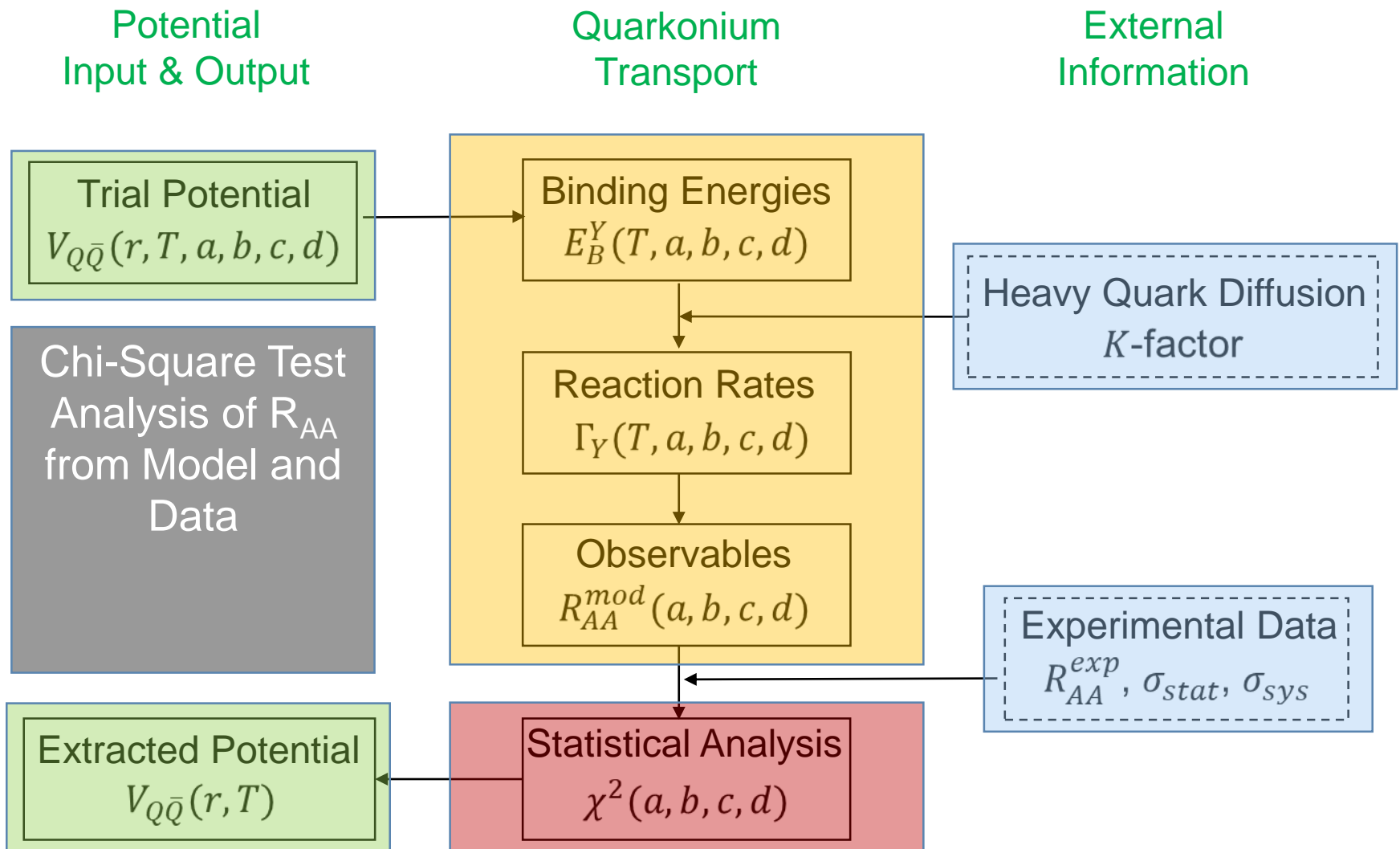
CMS Collaboration, V. Khachatryan et al., Phys. Lett. B 770(2017) 357.

ALICE Collaboration, B.B. Abelev et al., Phys. Lett. B 738(2014) 361.

CMS Collaboration, A.M. Sirunyan et al., Phys. Lett. B 790(2019) 270.

ALICE Collaboration, S. Acharya et al., Phys. Lett. B 790(2019) 89.

Extracting In-Medium Potential



Statistical Analysis: χ^2 Test

Model calculations
with trial potential

$$R_{AA}^{mod}(a, b, c, d)$$

53 R_{AA} experimental
data points

$$R_{AA}^{exp}, \sigma_{stat}, \sigma_{sys}$$

$$\chi^2 = \sum_{i=1}^N \left(\frac{R_{AA}^{mod}(a, b, c, d) - R_{AA}^{exp}}{\sigma_{exp}} \right)^2$$

$$\sigma_{exp} = \sqrt{\sigma_{stat}^2 + \sigma_{sys}^2}$$

p-value

Assuming the model is correct, there is p-value chance that the data is at least this discrepancy from the model

- Larger χ^2 smaller p-value
- Smaller χ^2 larger p-value

Chi-square distribution
degree of freedom:
53 data – 4 parameters = 49

Accepted band

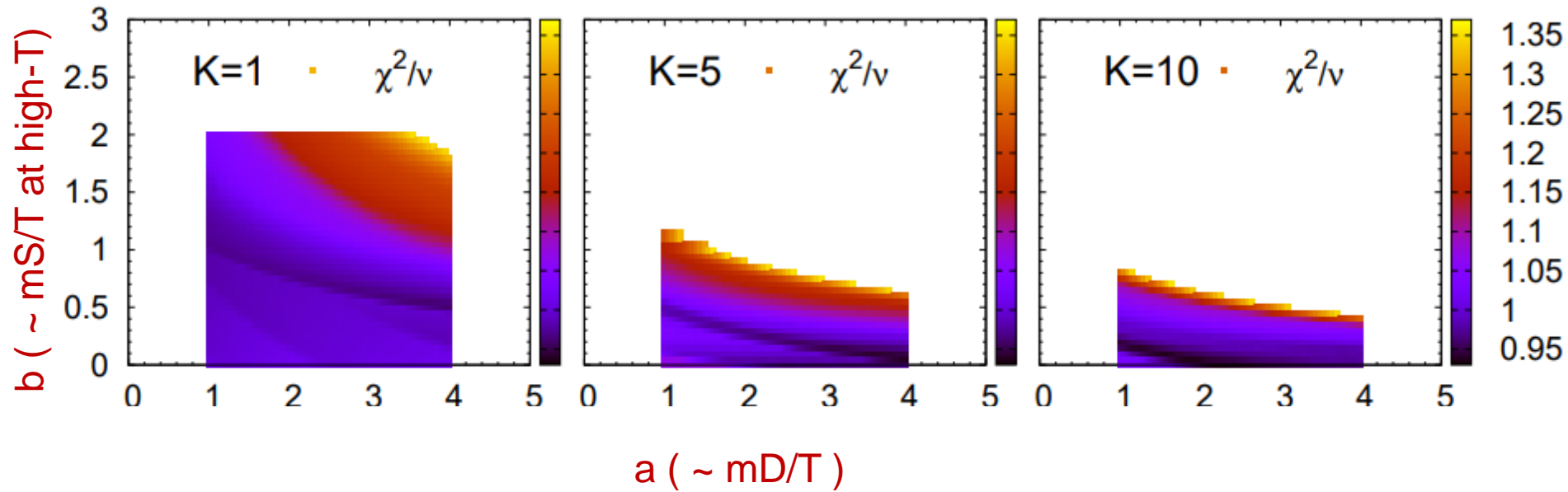
95% confidence: p-value > 0.05



Best fit

Smallest χ^2 , Largest p-value

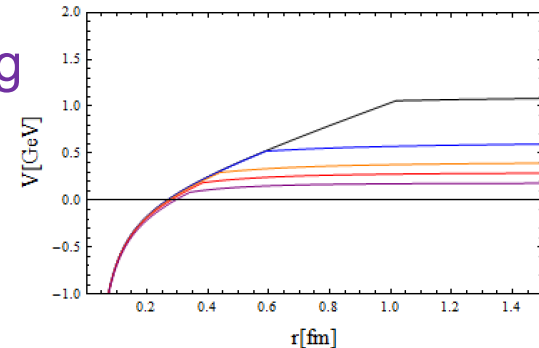
Resultant Parameter Space



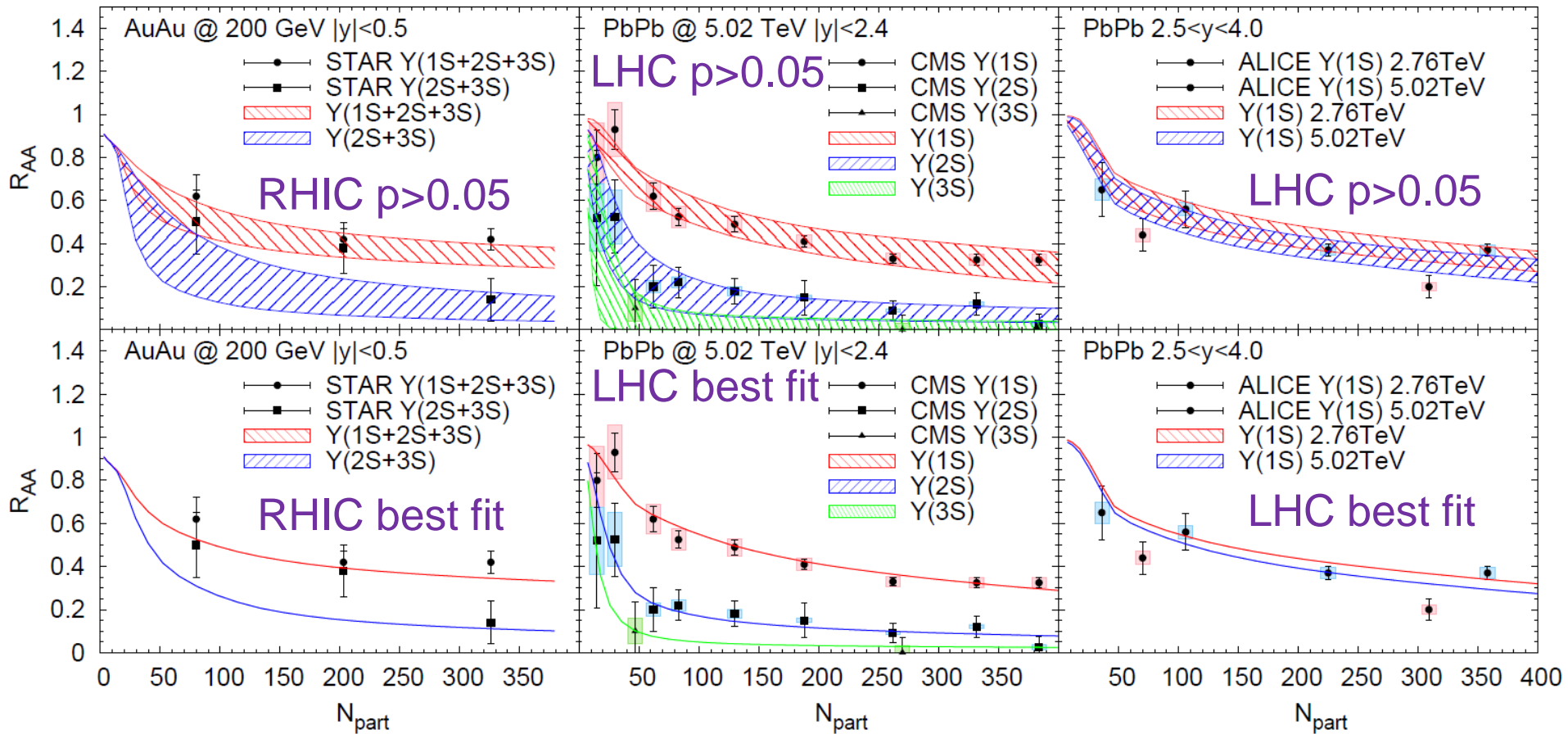
Smaller b \rightarrow smaller m_S \rightarrow larger r_s \rightarrow stronger binding
 Larger b \rightarrow larger m_S \rightarrow smaller r_s \rightarrow weaker binding

In order to describe data, rates are approximately constant at relevant temperature ($<400\text{MeV}$)

Thus Increasing K-factor requires stronger binding



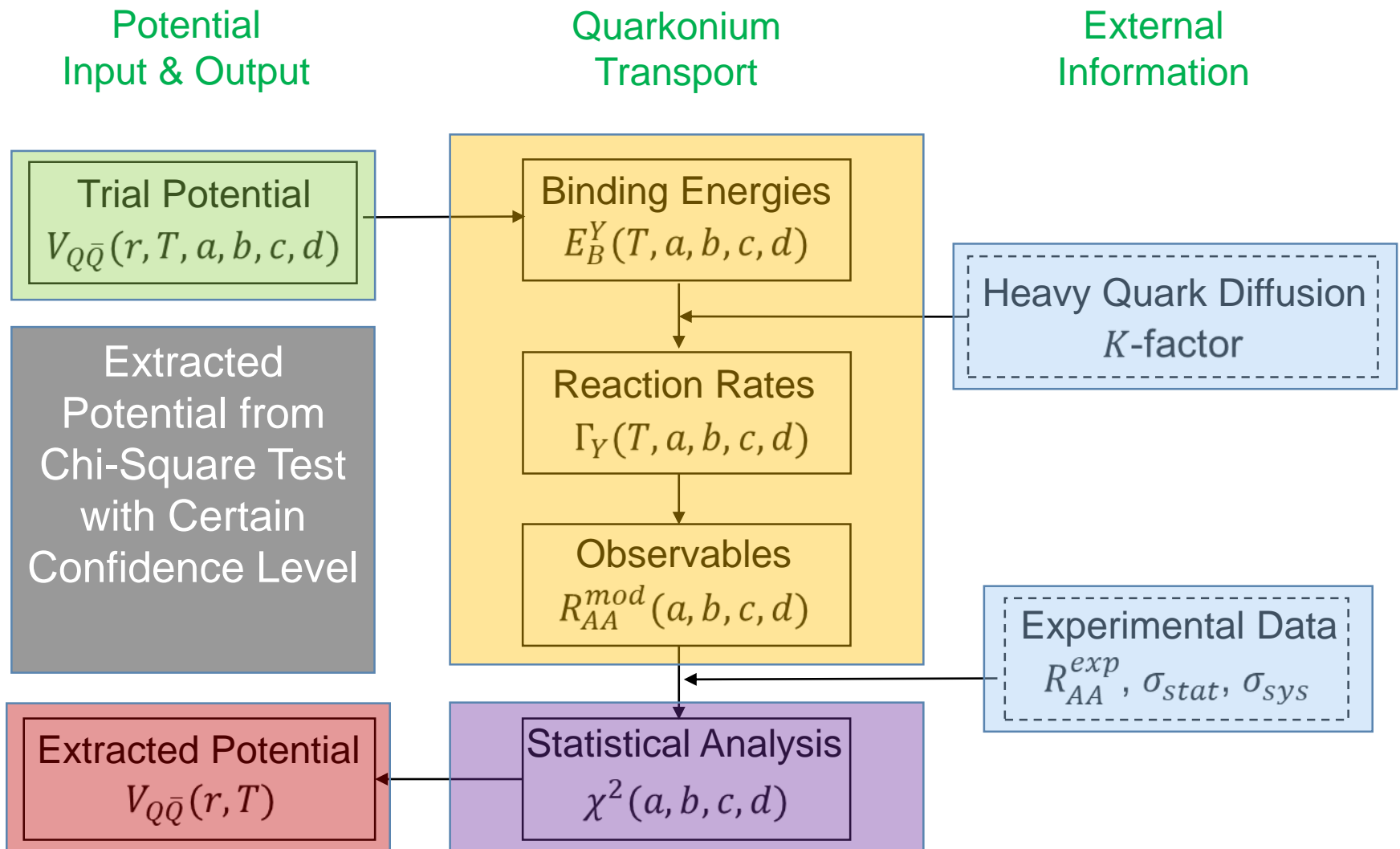
Part of K=5 Resultant Bands of R_{AA} s



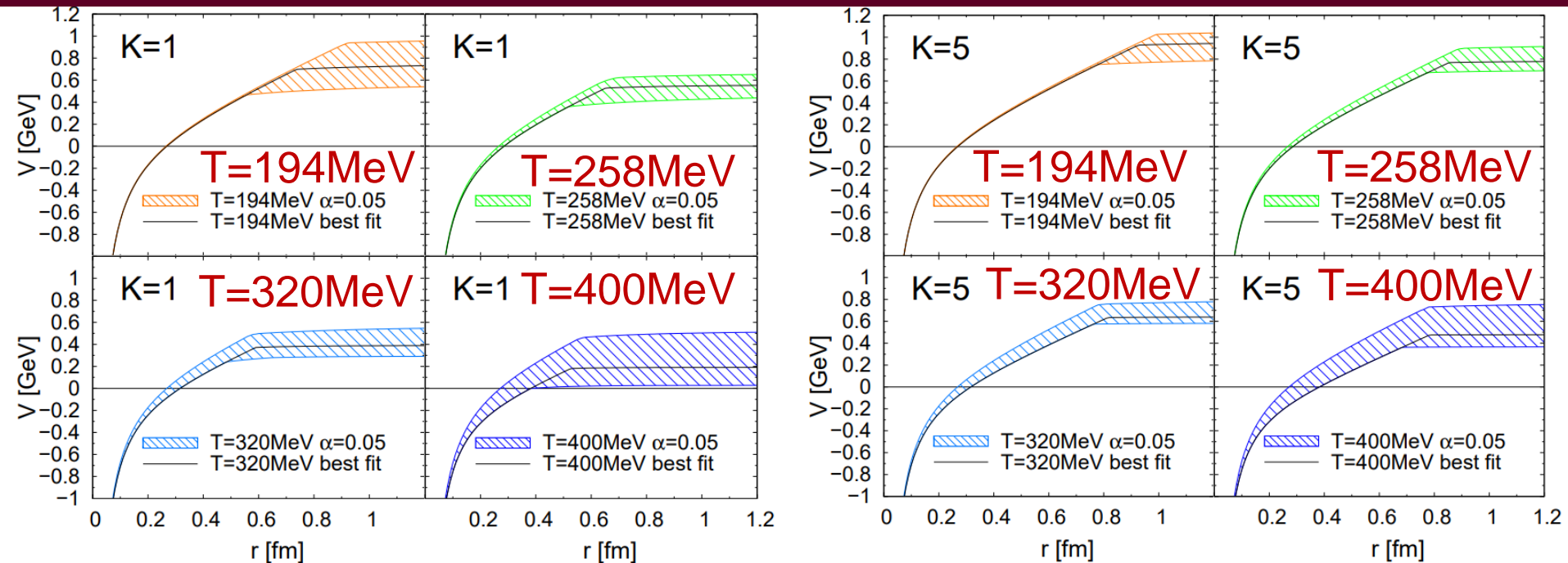
Overall good description of data with parameters selected by statistics

Few data points do not have constraint power

Extracting In-Medium Potential



Resultant Bands of In-Medium Potential



$V^{\text{best-fit}}(r=1\text{fm})$	$K=1$	$K=5$
$T=194\text{MeV}$	$\sim 730\text{MeV}$	$\sim 940\text{MeV}$
$T=320\text{MeV}$	$\sim 400\text{MeV}$	$\sim 630\text{MeV}$

$K=1$: allows for both weak/strong binding

$K=5$: suggests strong binding

$K=10$: even stronger..

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Summaries

- Transport model is previously well tested in multiple aspects in heavy-ion collisions
 - J/ψ , Υ , ...
 - R_{AA} , v_2 , ...
- Regeneration is importance for J/ψ , but not so important for Υ
 - $R_{AA}(N_{part})$, $R_{AA}(p_T)$, $v_2(p_T)$, $R_{AA}(\sqrt{s})$, ...
- $\Upsilon(1S)$ is ideal for probing in-medium heavy quark-antiquark potential/QCD force
 - $\Upsilon(1S)$'s observable is sensitive to in-medium potential
 - $\Upsilon(1S)$ has small regeneration, is cleaner than charmonium
- QGP is a strongly coupled medium
 - Joint study of open/hidden heavy flavors is important
 - We should select the most relevant data points for probing the force

Thank you!