

# Nucleon EDM & Decay from the Lattice

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for the RBC-UKQCD collaborations

# Topics

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- ▶ Introduction
- ▶ Neutron and proton EDM
- ▶ Proton decay matrix elements
- ▶ Summary and future work

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# Introduction

# Search for NP from lattice QCD

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- ▶ Precise bound of undetected observables
  - ▶ Direct constraint on BSM  
EDM (nucleon (quark), electron, ...), Proton(neutron) decay, NNbar oscillation, LFV, dark matter search, ...
  - ▶ Intensity frontier physics
  - ▶ Hadronic correction should be significant  
Need to take account of nucleon structure, and low energy physics of QCD
  - ▶ Require the model independent method for strong interaction

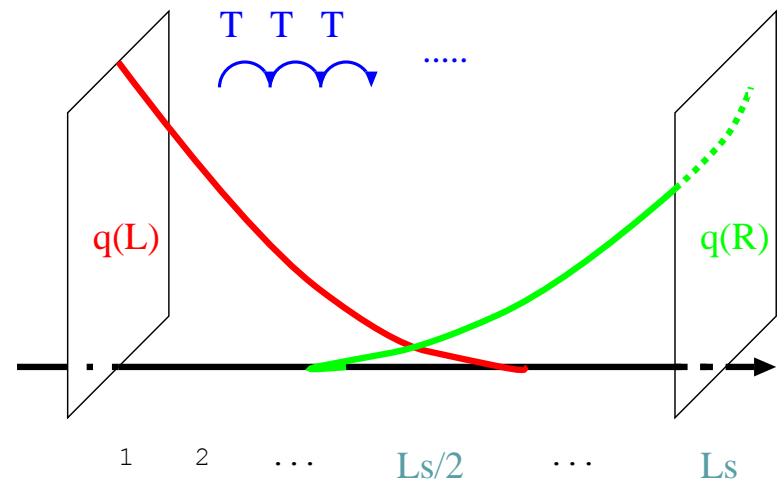
Lattice QCD plays a key role !

# Lattice fermion

- ▶ There are several kinds of fermion definition on the lattice  
Due to Nielsen-Ninomiya no-go theorem
- ▶ Require “realistic” fermion for the **precise calculation**
  - ▶ Wilson-clover (and staggered fermions) may not be appropriate (due to sys. error maybe...)
  - ▶ **Domain-wall fermion** (and also overlap fermion) we use here
- ▶ **Domain-Wall fermion (DWF)**

- L, R fermion are localized on boundaries  
⇒ Chiral symmetry (if  $L_s \rightarrow \infty$ ).
- **Good chiral symmetry**  
Chiral symmetry breaking is suppressed as  $am_{\text{res}} \sim \exp(-L_s)$ .

[Blum Soni, (97), CP-PACS(99), RBC(00),  
RBC/UKQCD. (05 --) ]



# Our strategy

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- ▶ Non-perturbative determination of QCD contribution to nucleon (neutron and proton) EDM & decay from lattice
- ▶ Domain-wall fermion
  - ▶ High precision, but highly computational cost  
spend more than 1 year using powerful supercomputer
  - ▶ New development of algorithm to carry out efficiently
- ▶ All-mode-averaging (AMA) algorithm Blum, Izubichi, ES, I208.4349
  - ▶ more than 5 times faster !
  - ▶ various applications to other observables  
(nucleon form factor)
  - ▶ possible to perform high precision measurement

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# Nucleon EDM from the lattice

# Nucleon EDM in the SM and BSM

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- ▶ Sensitive to P, CP violation
- ▶ Upper limit from experiment:  $< 2.9 \times 10^{-26} \text{ e} \cdot \text{cm}$
- ▶ Contribution from weak boson: CKM phase

Very **tiny**, which is 3-loop :  $d_N^{\text{KM}} \simeq 10^{-30} \text{ -- } 10^{-32} \text{ e} \cdot \text{cm}$

Khriplovich and Zhitnitsky, PLB109, 490 (1982); Czarnecki, Krause, PRL78, 4339 (1997)

- ▶ Contribution from QCD:  $\theta$  term

Unnaturally small (strong CP problem):  $\bar{\theta} < 10^{-9 \pm 1}$

Crewther, et al. (1979), Ellis, Gaillard (1979)

- ▶ Contribution from BSM: dim-5,6 operator

$$\mathcal{O}_{\text{qEDM}} = d_q \bar{q}(\sigma \cdot F)\gamma_5 q, \quad \mathcal{O}_{\text{cEDM}} = d_q^c \bar{q}(\sigma \cdot G)\gamma_5 q, \quad \mathcal{O}_{\text{Weinberg}} = d^G G G \tilde{G}$$

$$d_N = d_N^{\text{QCD}} \bar{\theta} + d_N(d_q, d_q^c) + d_N(d^G)$$

$$\sim 10^{-17} [\text{e} \cdot \text{cm}] \bar{\theta} + (1.4 - 0.47) d_d - (0.12 - 0.35) d_u + O(10^{-2}) d_q^c$$

$$\sim O(10^{-25} - 10^{-27}) \text{ e} \cdot \text{cm}$$

Hisano, Shimizu (04), Ellis, Lee, Pilaftsis (08),

Hisano, Lee, Nagata, Shimizu (12)



# Nucleon EDM from lattice QCD

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- ▶ Non-perturbative determination of QCD effect

- ▶  $L_\theta$

From lattice QCD we obtain  $d_N$  in  $d_N^{\text{QCD}} = \bar{\theta} d_N$

$\theta$  parameter can be estimated by  $d_N^{\text{exp}}/d_N = \theta$  (if there is no BSM)

- ▶  $L_{\text{qEDM}}, L_{\text{cEDM}}$

From lattice QCD we obtain  $C_{\text{qEDM}}, C_{\text{cEDM}}$

$$d_N^{\text{BSM}} = \sum_q \left[ d_q^{\text{qEDM}} C_{\text{qEDM}}^q + d_q^{\text{cEDM}} C_{\text{cEDM}}^q \right]$$

$d^{\text{qEDM}}, d^{\text{cEDM}}$  depend on BSM parameters.

Result of lattice QCD is an important input value for BSM search

# Methods

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## ▶ Spectrum

$$m_{\uparrow \text{ spin}} - m_{\downarrow \text{ spin}} = 2d_N \theta E \quad R_3 = \frac{\langle N(t) \bar{N}(0) \rangle_{\theta, E}^{\text{up}}}{\langle N(t) \bar{N}(0) \rangle_{\theta, E}^{\text{down}}} \simeq 1 + d_N E \theta t$$

E: External electric field

Aoki and Gocksch, PRL63, 1125 (1989); ES, et al., (CP-PACS) PRD75, 034507 (2007); ES et al., PRD78, 014503 (2008)

## ▶ CPV Form factor

$$\langle n(P_1) | J_{\mu}^{\text{EM}} | n(P_2) \rangle_{\theta} = \bar{u}_N^{\theta} \left[ \underbrace{\frac{F_3^{\theta}(Q^2)}{2m_N} \gamma_5 \sigma_{\mu\nu} Q_{\nu}}_{\text{P, T-odd}} + \underbrace{F_1 \gamma_{\mu} + \frac{F_2}{2m_N} \sigma_{\mu\nu} Q_{\nu} + \dots}_{\text{P, T-even}} \right] u_N^{\theta}$$

$$d_N = \lim_{Q^2 \rightarrow 0} F_3(Q^2) / 2m_N$$

ES, et al., (CP-PACS), PRD72, 014504 (2005); Berruto, et al. (RBC) PRD73, 05409 (2006).

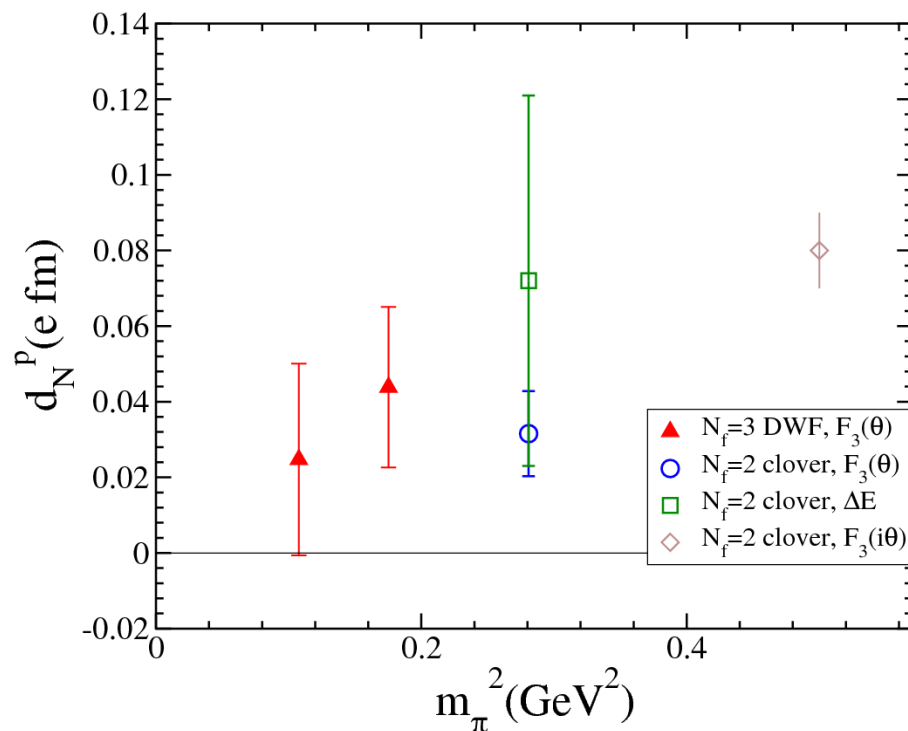
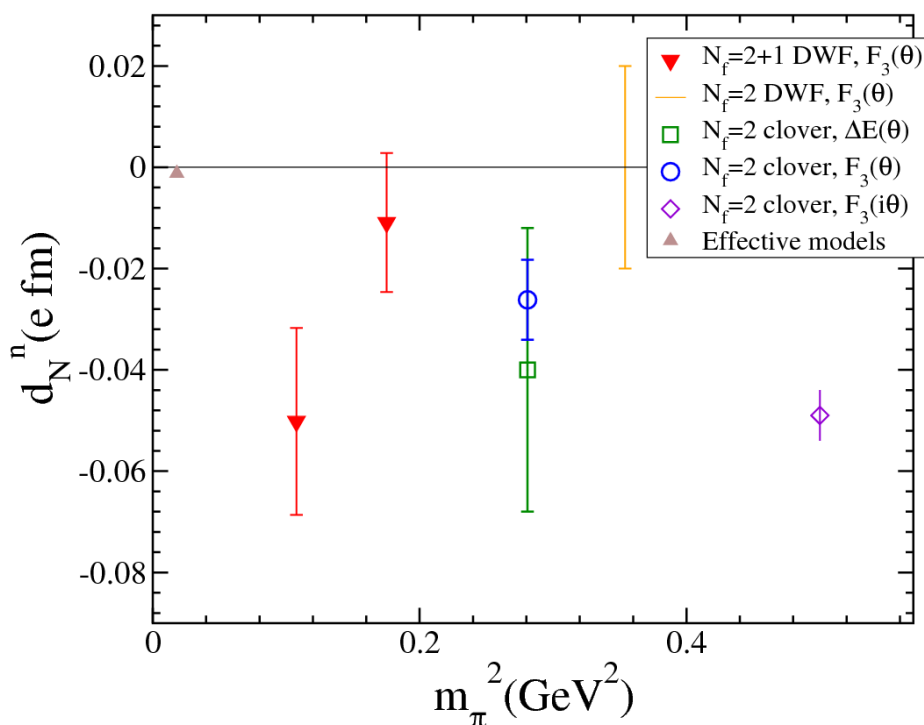
## ▶ Imaginary $\theta$

New generation of imaginary  $\theta$  action:  $\langle O e^{i\theta Q} \rangle \rightarrow \langle O e^{-\theta^I Q} \rangle$

T. Izubuchi, Lattice 2007

# Recent results

- ▶ DWF in  $N_f=2+1$  (RBC/UKQCD)  $m_\pi = 300\text{--}400$  MeV
  - ▶ AMA is very helpful, cost is reduced to 1/5 or less.
  - ▶ Statistical error is more than 40%, and also we need to estimate systematic error (finite size effect).



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# Nucleon decay from the lattice

# Effective operator

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## ▶ Dimension-6 BV operator

$$\mathcal{L}_{\text{GUT}} = \mathcal{L}_{\text{SM}} + \sum_i C_i(\mu) O_i(\mu) / \Lambda_{\text{GUT}}^2 + \mathcal{O}((O(\mu) / \Lambda_{\text{GUT}}^2)^2)$$

$$O_i(\mu) = (qq)_{\Gamma}(ql)_{\Gamma'} \quad \text{“}i\text{” labels chirality } (\Gamma) \text{ and flavor } (q,l)$$

$C_i$  depends on type of GUTs model

## ▶ Matrix element

Lattice QCD provides each decay channels of  $W_0$  from matrix element;

$$\langle \pi^0 | (ud)_{\Gamma} u_{\Gamma'} | p \rangle = P_{\Gamma'} \left[ W_0^{\Gamma} - \frac{i \not{q}}{m_p} W_1^{\Gamma} \right] u_p \quad \begin{array}{l} \text{Aoki et al. (JLQCD), PRD62, 014506} \\ \text{(2000); Aoki et al. (RBC), PRD75, 014507} \\ \text{(2007)} \end{array}$$

which is extracted from 3-pt function.

$W_0$  : determine from QCD matrix element (model independent)

## ▶ Decay rate

$$\Gamma_{p \rightarrow \pi^0 e^+} = \frac{m_p}{32\pi^2} \left[ 1 - \left( \frac{m_e}{m_p} \right)^2 \right]^2 \left| \sum_i C_i W_0^i(p \rightarrow \pi^0) \right|^2$$

Precision of  $W_0$  is significant, since the decay rate is affected by twice of that.

# How to obtain $W_0$ from lattice QCD

## ▶ The “indirect” method

- ▶ Measurements of low-energy constant via tree level chiral perturbation theory.

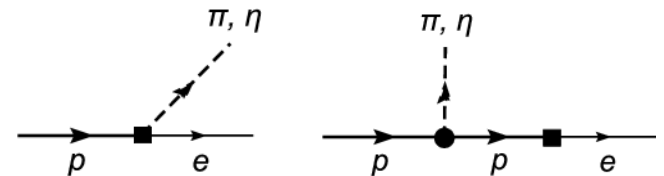
$$W_0^{LR}(p \rightarrow \pi^0) \simeq \alpha(1 + D + F)/\sqrt{2}f_0,$$

where D and F is given by experiment, and  $\alpha$  is given by 2-pt function:

$$\langle 0 | ((ud)_{RuL}) | p \rangle = \alpha P_L u_p$$

Easy calculation, BUT it has systematic error due to being based on ChPT.

Claudson, et al., NPBI95 (1982) 297



S.Aoki et al. (JLQCD), PRD62, 014506 (2000), Y. Aoki et al.(RBC), PRD75, 014507 (2007), Y.Aoki et al. (RBC-UKQCD), PRD78, 054505 (2008)

## ▶ The “direct” method

- ▶ Measurement of matrix element extracted from 3-pt function.
- ▶ Rather expensive, while there is no uncertainty depending on models.
- ▶ Provides each channels of decay mode.

S Aoki et al. (JLQCD), PRD62, 014506 (2000), Y.Aoki et al.(RBC), PRD75, 014507 (2007)

# Low energy constant (indirect)

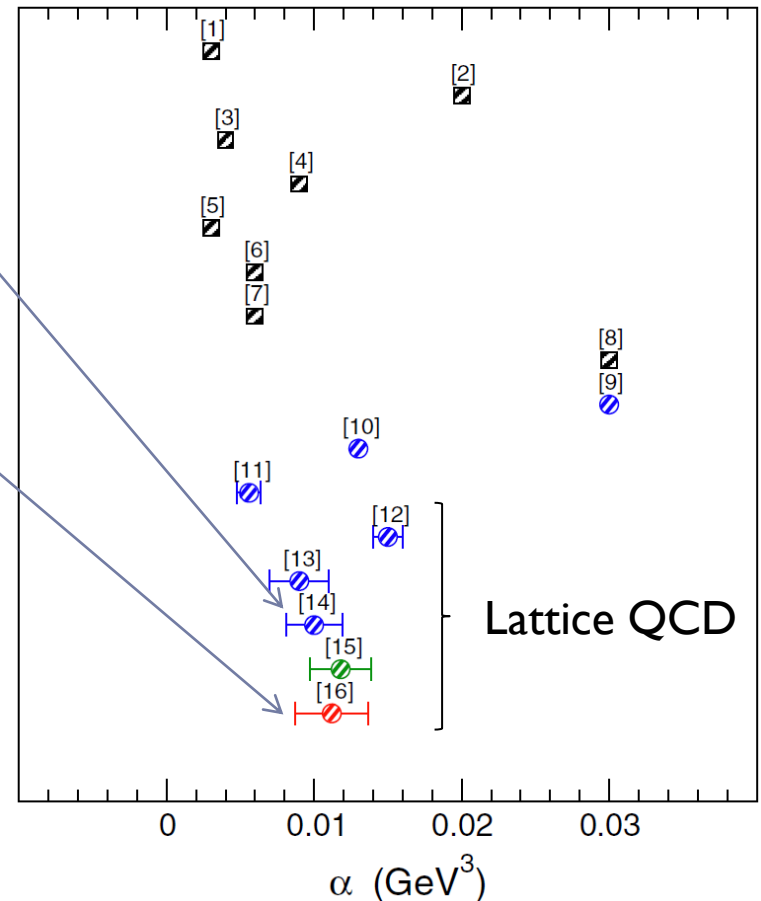
- ▶ Works with DWFs
- ▶ Quenched QCD (direct/indirect)

Y.Aoki, C. Dawson, J. Noaki, and A. Soni,  
Phys. Rev. D75, 014507 (2007)

- ▶  $N_f=2+1$  (indirect)

Y.Aoki et al. (RBC-UKQCD), Phys. Rev.  
D78, 054505 (2008)

- Large model dependence
- Need to subtract ChPT ambiguity  
⇒ direct calculation



# $W_0$ (direct)

Aoki, ES, Soni, arXiv:1304.7424

- DWFs in  $N_f=2+1$  (direct)

$24^3 \times 64$  lattice in RBC/UKQCD collaboration,

Chiral extrapolation with  $m=0.005, 0.01, 0.02, 0.03$  ( $m_\pi = 0.3 \text{ -- } 0.8 \text{ GeV}$ )

- Physical kinematics
- Estimate all systematic errors
- Uncertainty is still large.

(stat + sys error):

30--40% for  $p \rightarrow \pi$

20--40% for  $p \rightarrow K$

$$-\langle \pi^0 | (ud)_R u_L | p \rangle$$

$$\langle \pi^0 | (ud)_L u_L | p \rangle$$

$$\langle K^0 | (us)_R u_L | p \rangle$$

$$\langle K^0 | (us)_L u_L | p \rangle$$

$$-\langle K^+ | (us)_R d_L | p \rangle$$

$$\langle K^+ | (us)_L d_L | p \rangle$$

$$-\langle K^+ | (ud)_R s_L | p \rangle$$

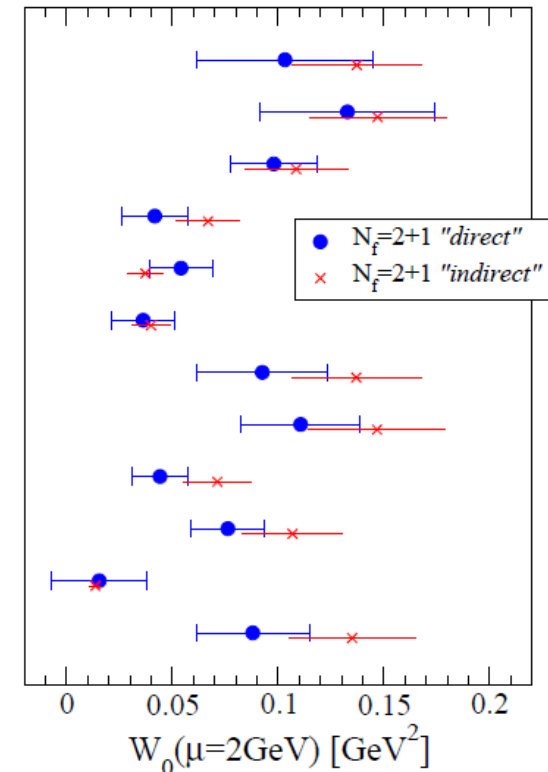
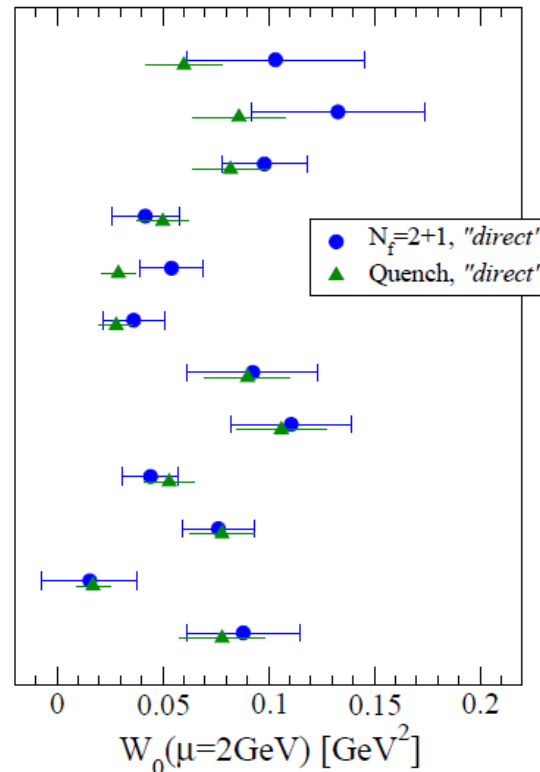
$$\langle K^+ | (ud)_L s_L | p \rangle$$

$$-\langle K^+ | (ds)_R u_L | p \rangle$$

$$-\langle K^+ | (ds)_L u_L | p \rangle$$

$$\langle \eta | (ud)_R u_L | p \rangle$$

$$\langle \eta | (ud)_L u_L | p \rangle$$





# Summary and future work

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- ▶ Lattice study of nucleon EDM & decay
- ▶ Precise estimate of non-perturbative contribution
- ▶ Need to reduce the statistical + systematic error
  - ▶ Nucleon EDM : > 40% stat. error and large sys. error expected
  - ▶ Nucleon decay :  $p \rightarrow \pi$  channel : 30--40% total error  
 $p \rightarrow K$  channel: 20--40% total error
- ▶ **AMA algorithm is very helpful.** Blum, Izubichi, ES, I208.4349
- ▶ Larger lattice size ( $\sim 5 \text{ fm}^3$ ) at physical point will be available soon. We expect 10–20% error level.

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Thank you

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# Backup

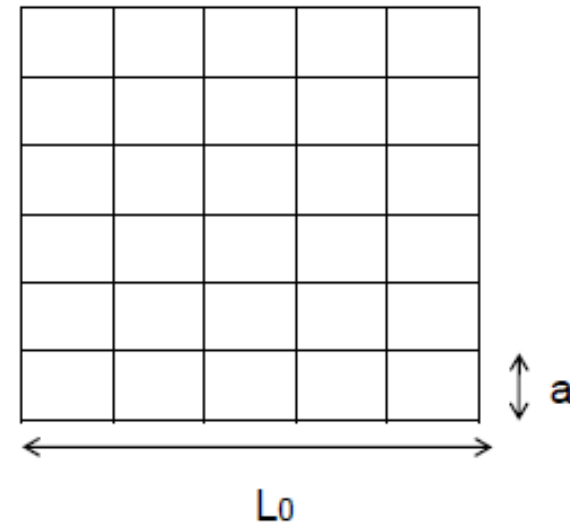
# Lattice QCD

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In lattice regularization, the path integral of  $\langle O \rangle$  is computed by Monte-Carlo integral:

$$\langle O \rangle = Z^{-1} \int D\Psi O(\Psi) e^{-S(\Psi)} \simeq \frac{1}{N} \sum_i O(\Psi_i)$$

- ▶ **Exact** QCD calculation (enough large number of sampling N)
- ▶ Gauge invariant
- ▶ Translational invariant
- ▶ Ultraviolet cut-off  $a$  (lattice spacing)
- ▶ Infrared cut-off  $V=L_0^D$  (lattice volume)
- ▶ Taking continuum limit, and infinite volume



# Lattice QCD

- ▶ Hadron spectrum in  $N_f=2+1$  QCD
  - ▶ Good agreement with various lattice action and fermion with experimental results !

Kronfeld, 1209.3468

