

eRHIC: THE MACHINE TO IMAGE QUARKS AND GLUONS



Electron Ion Collider: The Next QCD Frontier

Understanding the glue
that binds us all

arXiv: 1212.1701 & 1108.1713

BROOKHAVEN
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a passion for discovery



U.S. DEPARTMENT OF
ENERGY

Office of
Science

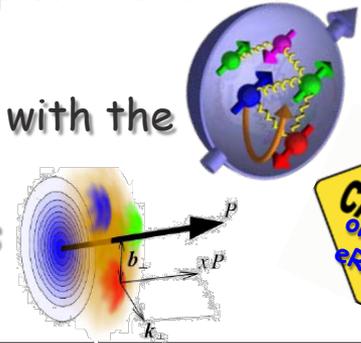
MOST COMPELLING SCIENCE QUESTIONS

How are sea quarks and gluons and their spin distributed in space and momentum inside the nucleon?



How are these quark and gluon distributions correlated with the over all nucleon properties, such as spin direction?

What is the role of the motion of sea quarks and gluons in building the nucleon spin?

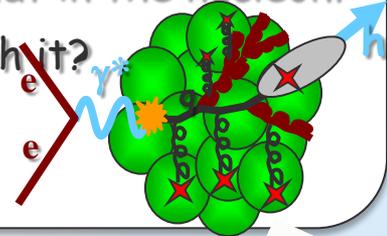


How does the nuclear environment affect the distribution of quarks and gluons and their interaction in nuclei?

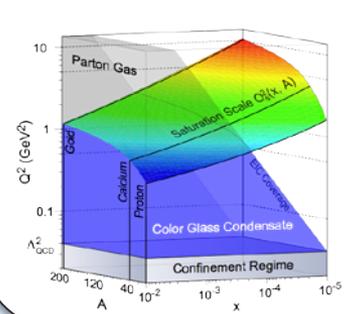


How does the transverse spatial distribution of gluons compare to that in the nucleon?

How does matter respond to fast moving color charge passing through it?
Is this response different for light and heavy quarks?



Where does the saturation of gluon densities set in?

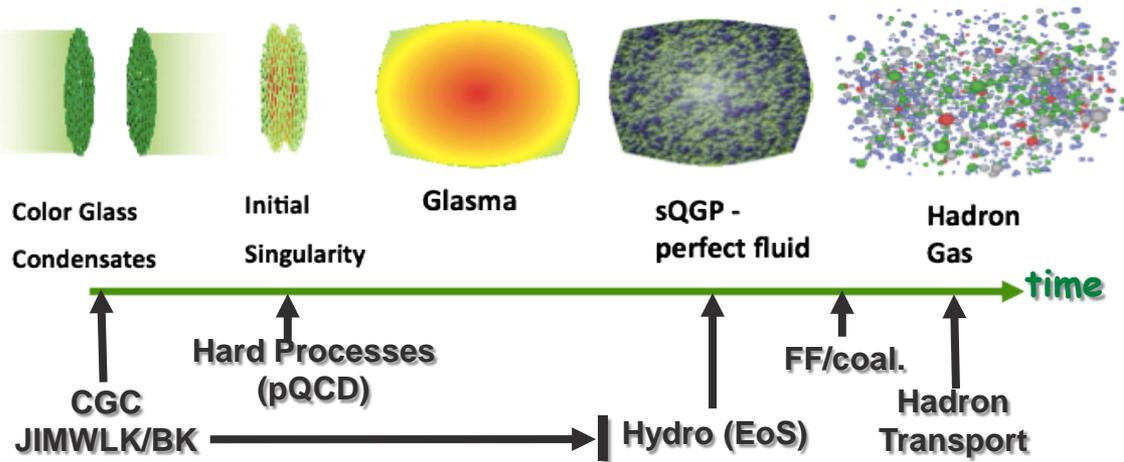


Is there a simple boundary that separates the region from more dilute quark gluon matter? If so how do the distributions of quarks and gluons change as one crosses the boundary?

Does this saturation produce matter of universal properties in the nucleon and all nuclei viewed at nearly the speed of light?



THE eA PHYSICS PROGRAM



Our understanding of some **fundamental** properties of the Glasma, sQGP and Hadron Gas depend **strongly** on our knowledge of the initial state!

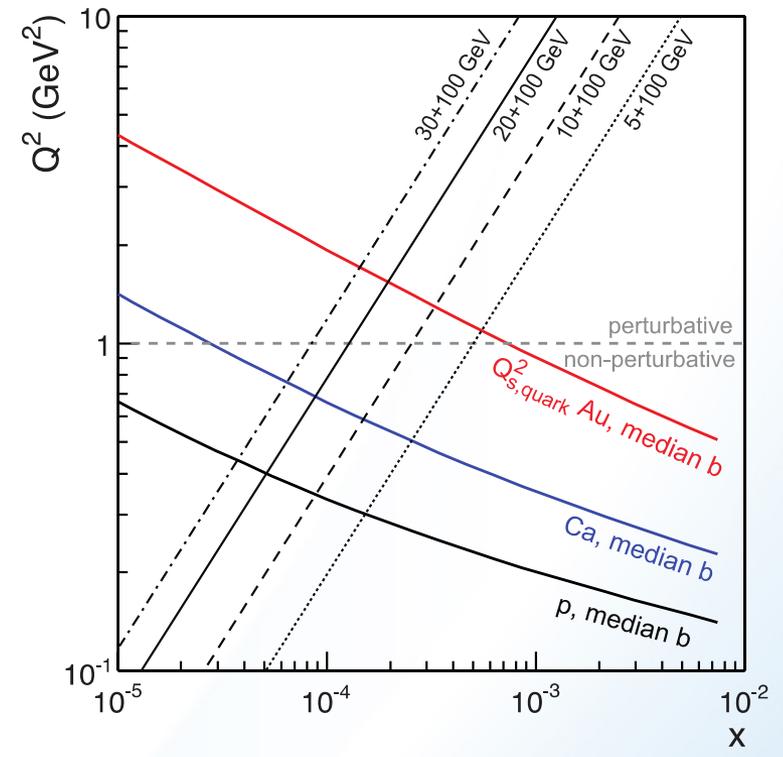
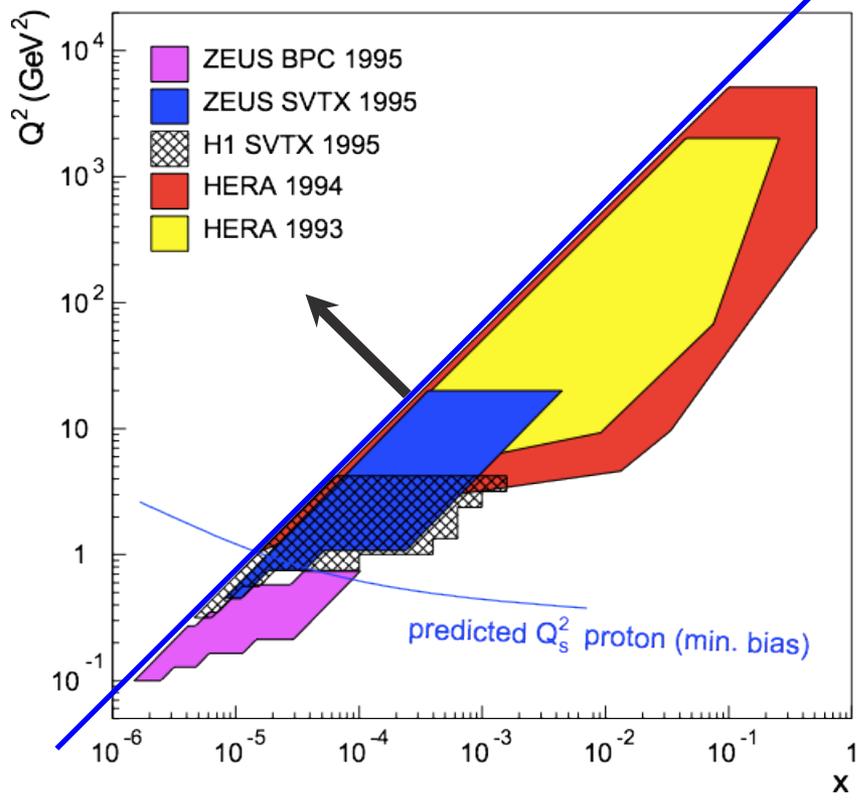
3 conundrums of the initial state:

1. What is the spatial transverse distributions of nucleons and gluons?
2. How much does the spatial distribution fluctuate? Lumpiness, hot-spots etc.
3. How saturated is the initial state of the nucleus?
→ unambiguously see saturation

Advantage over p(d)A:

- eA experimentally much cleaner
 - ✓ no “spectator” background to subtract
- Access to the parton kinematics through scattered lepton (x , Q^2)
- initial and final state effects can be disentangled cleanly
- Saturation:
 - ✓ no alternative explanations, i.e. no hydro in eA

eRHIC: REACHING THE SATURATION REGION

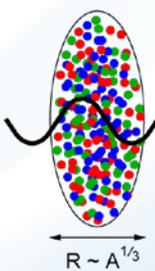


HERA (ep):

Despite high energy range:

- $F_2, G_p(x, Q^2)$ outside the saturation regime
- Need also Q^2 lever arm!
- Only way in ep is to increase \sqrt{s}
- Would require an ep collider at

eRHIC (eA):



$$(Q_s^A)^2 \sim c Q_0^2 \left(\frac{A}{x} \right)^{1/3}$$

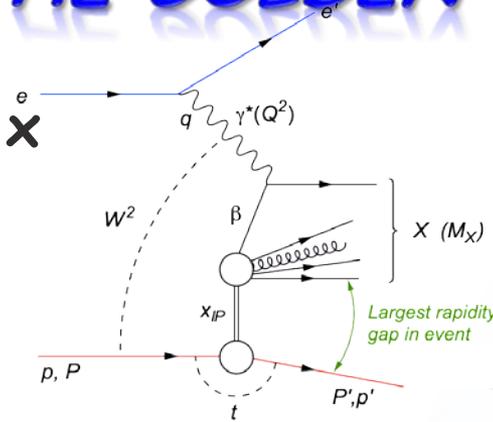
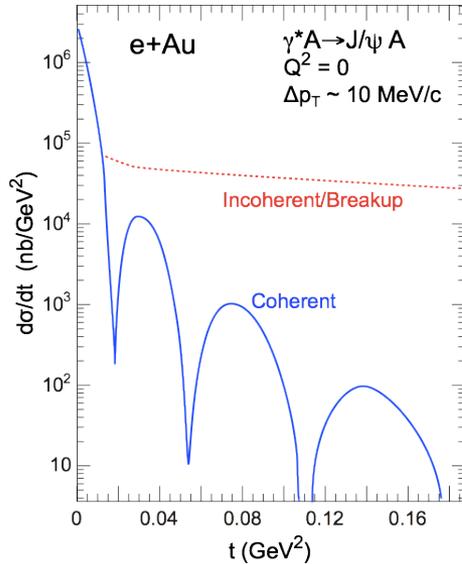
$$L \sim (2m_N x)^{-1} > 2 R_A \sim A^{1/3}$$

Probe interacts
coherently with all nucleons

Gold: 197 times smaller effective x !

SATURATION: THE GOLDEN CHANNEL

Hard diffraction in DIS at small x



Why is diffraction so important

Sensitive to spatial gluon distribution

$$\frac{d\sigma}{dt} \equiv \frac{\text{Fourier Transformation}}{\text{of Source Density } \rho_g(b)}$$

Hot topic:

➤ Lumpiness?

➤ Just Wood-Saxon+nucleon $g(b)$

Incoherent case:

measure fluctuations/lumpiness in $g_A(b)$

VM: Sensitive to gluon momentum distributions

➤ $\sigma \sim g(x, Q^2)^2$

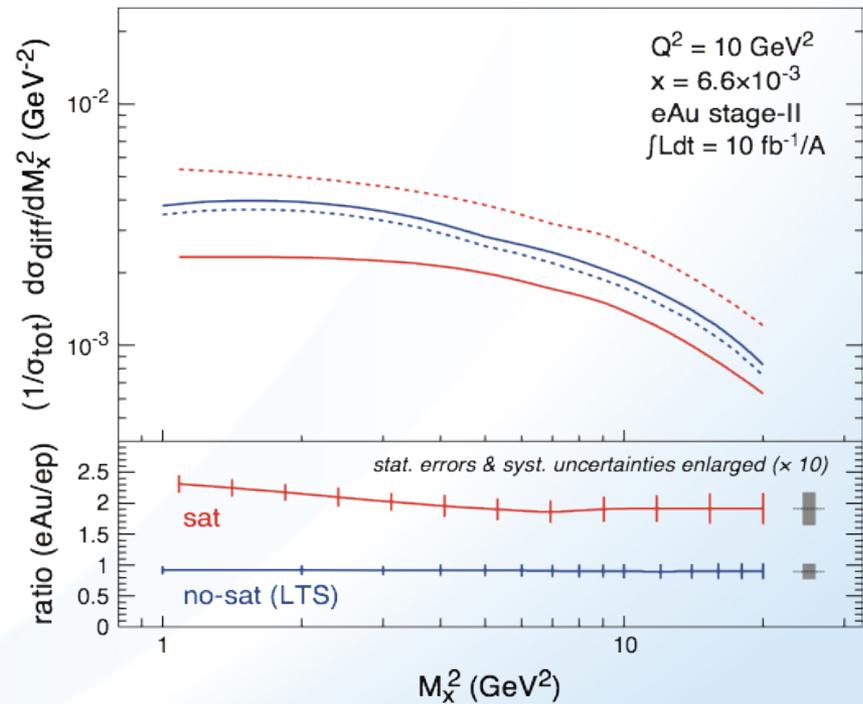
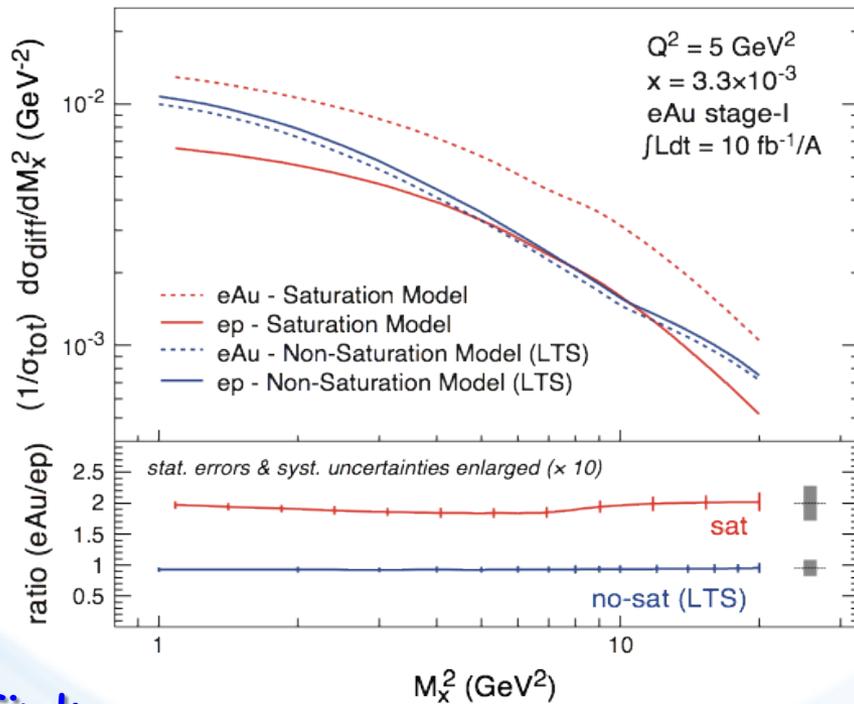
Predictions: $\sigma_{\text{diff}}/\sigma_{\text{tot}}$ in e+A ~25-40%

HERA: 15% of all events are hard diffractive

RATIO OF DIFFRACTIVE TO TOTAL CROSS-SECTION

- Black disc limit characterized by $\sigma_{\text{diff}}/\sigma_{\text{tot}} = 1/2$ (Hera sees 1/7)
- Large fraction of diffractive event is **unambiguous** signature for reaching the **saturated limit**

Fraction of low-mass coherent diffraction in ep and eA at eRHIC:



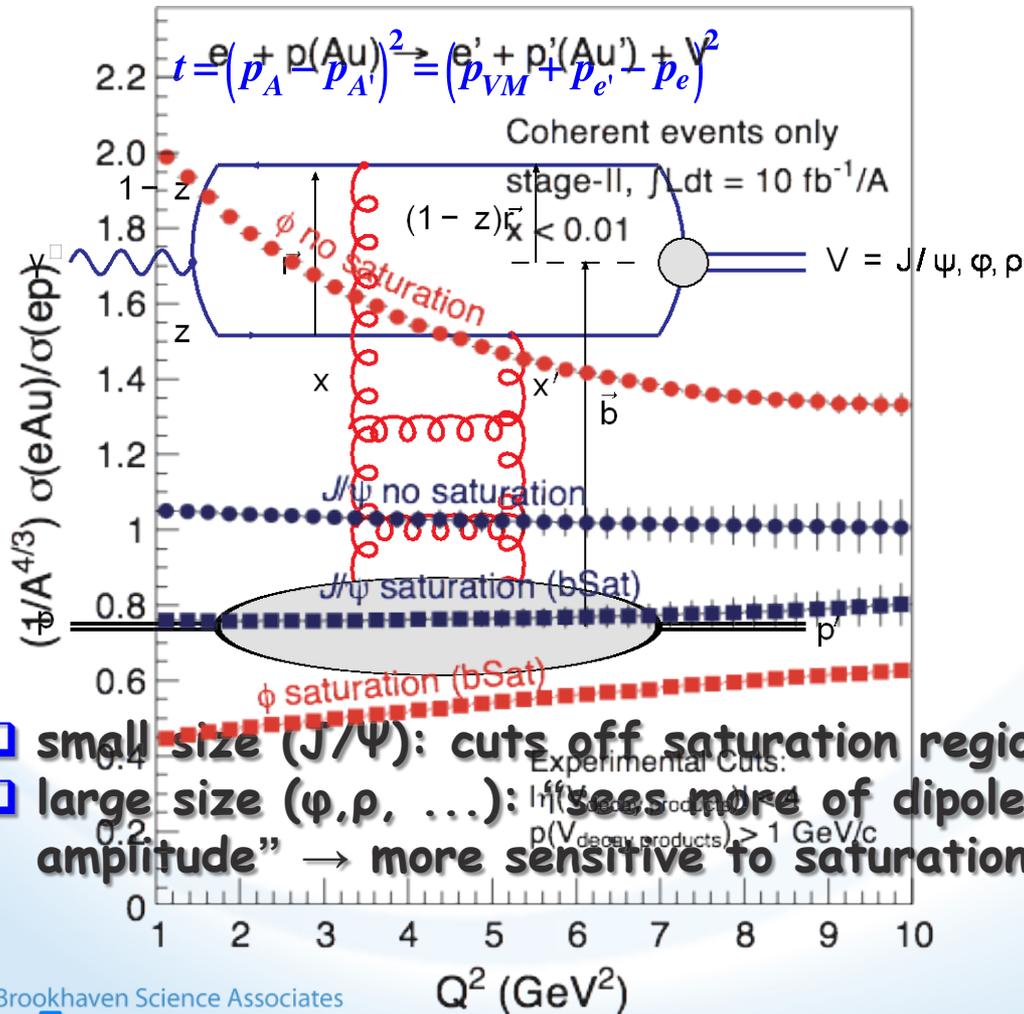
Find:

- w/o non-linear effects eA/ep ratio stays roughly one
- non-linear effects enhance σ_{diff} in eA scattering

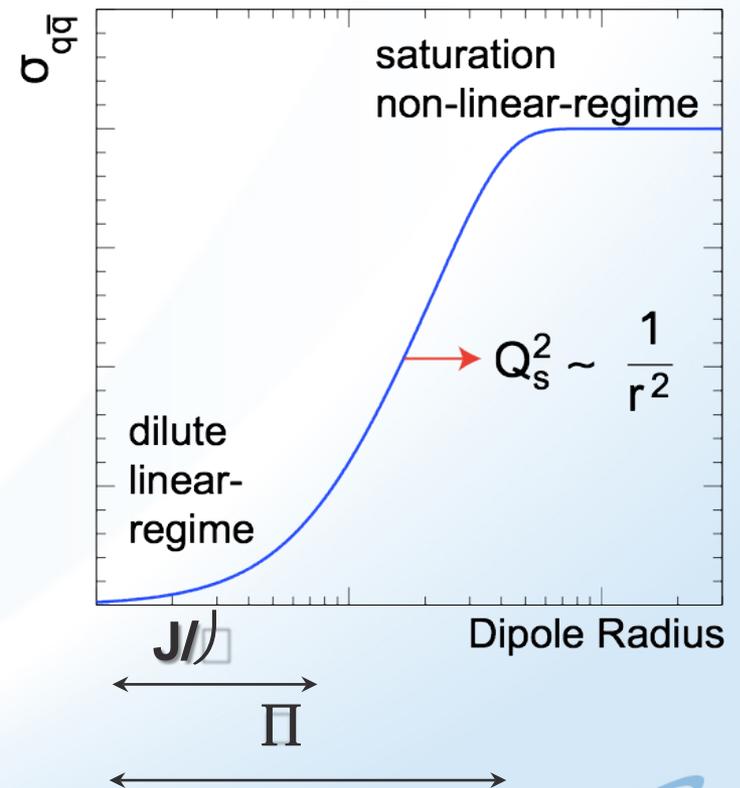
Day-1 signature
for Saturation

EXCLUSIVE VECTOR MESON PRODUCTION

- Unique probe - allows to measure momentum transfer t in eA diffraction
 - ➔ in general, one cannot detect the outgoing nucleus and its momentum



Dipole Cross-Section:

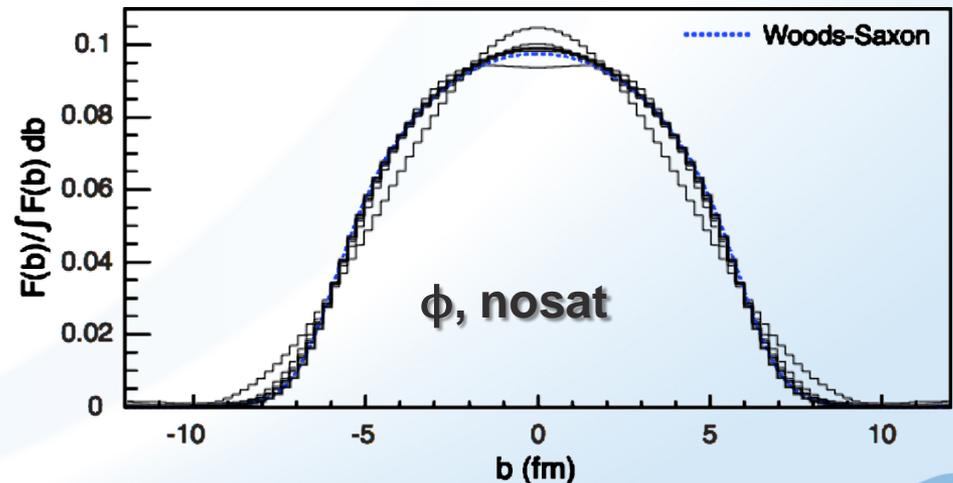
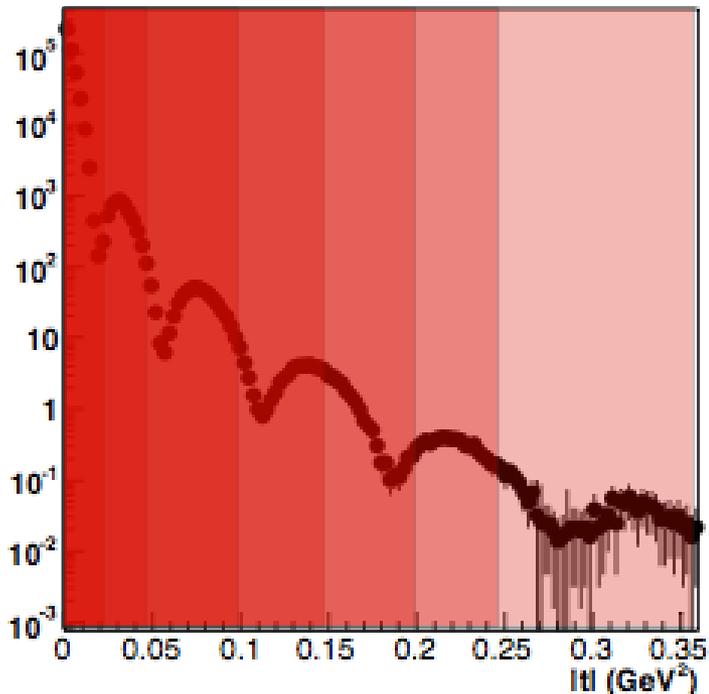


SPATIAL GLUON DISTRIBUTION THROUGH DIFFRACTION

- Idea: momentum transfer t conjugate to transverse position (b_T)
 - coherent part probes “shape of black disc”
 - incoherent part (dominant at large t) sensitive to “lumpiness” of the source (fluctuations, hot spots, ...)

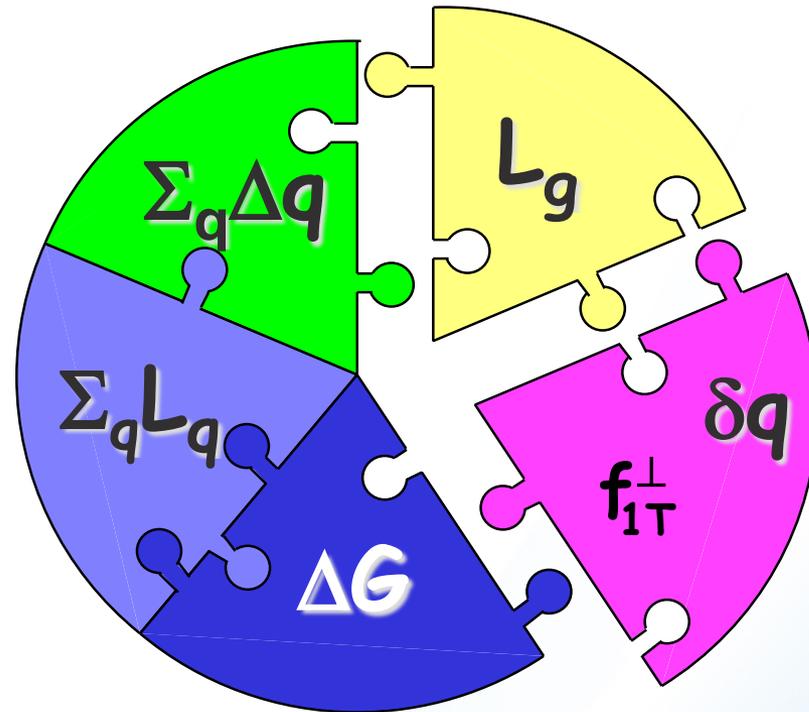
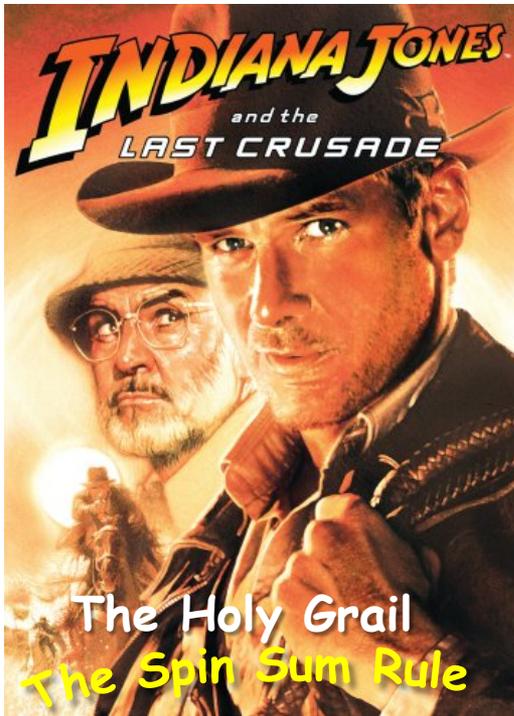
Spatial source distribution: $F(b) \sim \frac{1}{2\pi} \int_0^\infty d\Delta \Delta J_0(\Delta b) \sqrt{\frac{d\sigma}{dt}}$

$$t = \Delta^2 / (1-x) \approx \Delta^2 \quad (\text{for small } x)$$



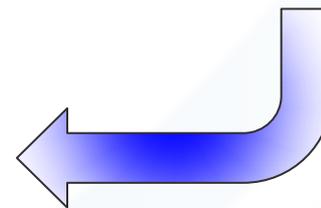
Golden eA measurement for eRHIC

WHAT COMPOSES THE SPIN OF THE PROTON



"Helicity sum rule"

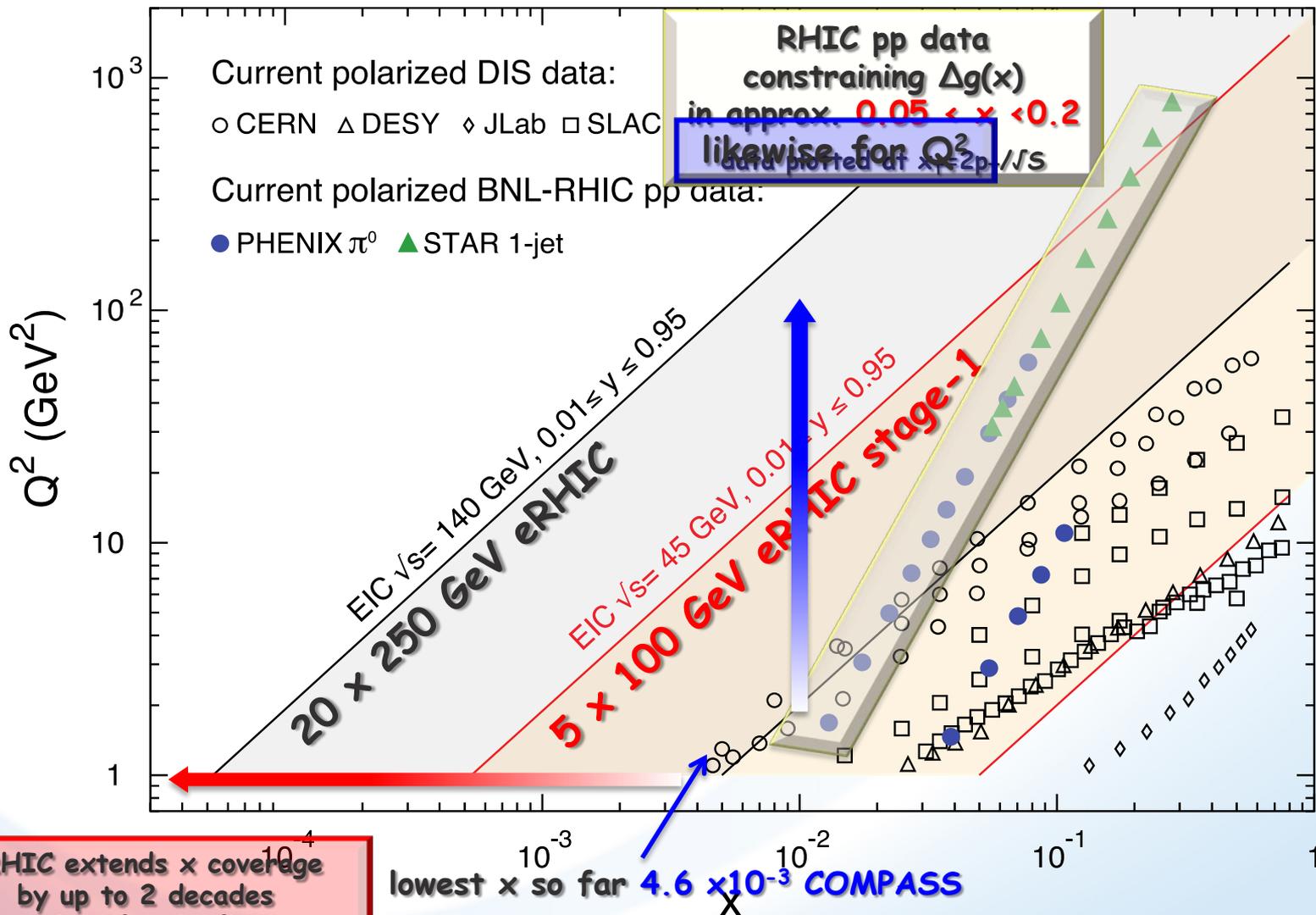
$$\frac{1}{2}\hbar = \left\langle P, \frac{1}{2} \left| J_{QCD}^z \right| P, \frac{1}{2} \right\rangle = \underbrace{\sum_q \frac{1}{2} S_q^z}_{\text{total u+d+s quark spin}} + \overbrace{S_g^z}^{\text{gluon spin}} + \underbrace{\sum_q L_q^z + L_g^z}_{\text{angular momentum}}$$



Can an eRHIC give the final answer?

Contribution to proton spin to date:
 Gluon: 20% (RHIC)
 Quarks: 30% (DIS)
MISS 50% → low x

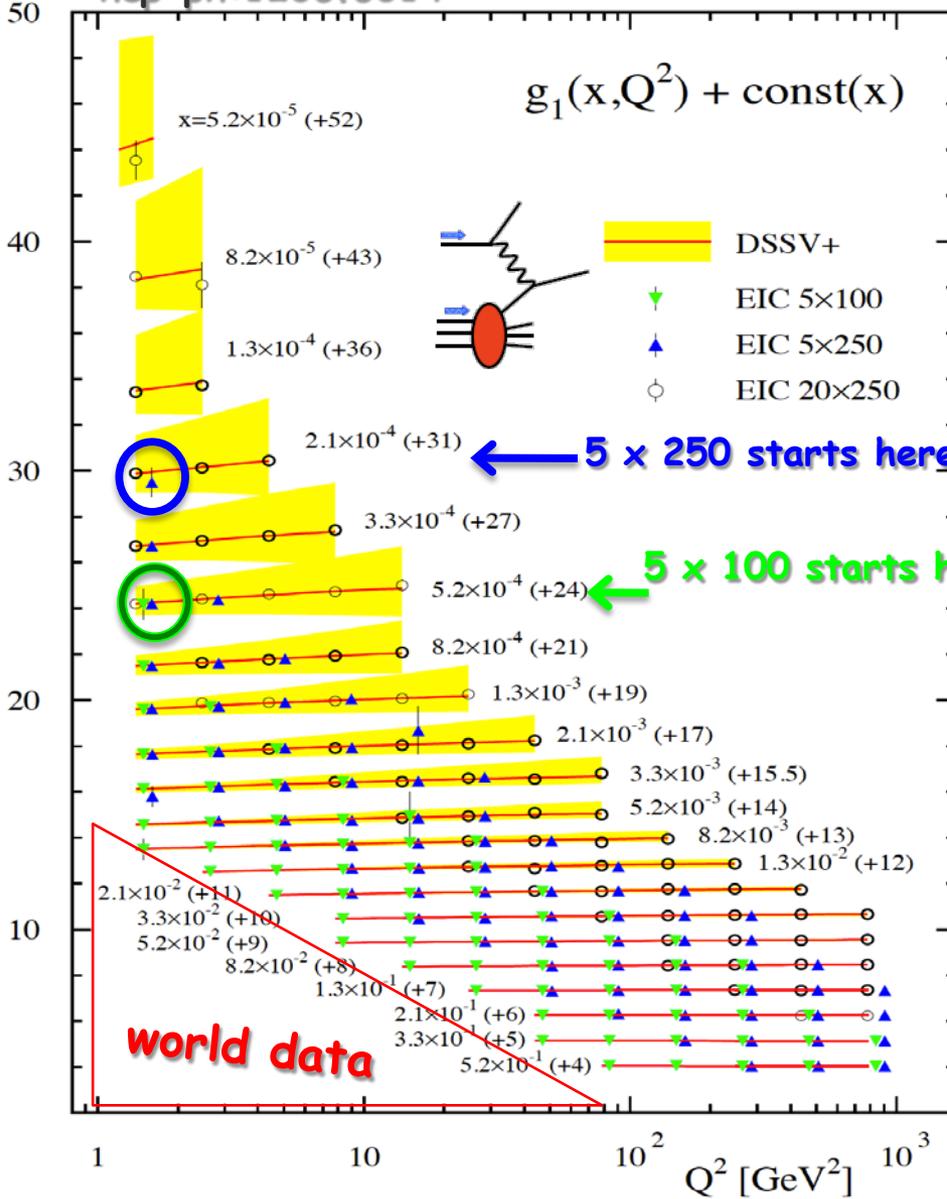
PRESENT VS eRHIC KINEMATIC COVERAGE



eRHIC extends x coverage
 by up to 2 decades
 (at $Q^2 = 1$ GeV²)

g_1^P THE WAY TO FIND THE SPIN

hep-ph:1206.6014



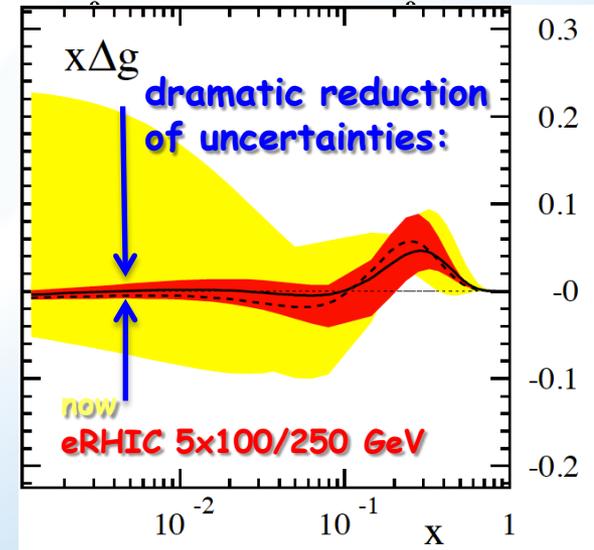
cross section: $\frac{d^2\sigma}{d\Omega dE'} \sim L_{\mu\nu} W^{\mu\nu}$

$$W^{\mu\nu} = -g^{\mu\nu} F_1 - \frac{p^\mu p^\nu}{v} F_2 + \frac{i}{v} \epsilon^{\mu\nu\lambda\sigma} q^\lambda s^\sigma g_1 + \frac{i}{v^2} \epsilon^{\mu\nu\lambda\sigma} q^\lambda (p \cdot q s^\sigma - s \cdot q p^\sigma) g_2$$

pQCD scaling violations

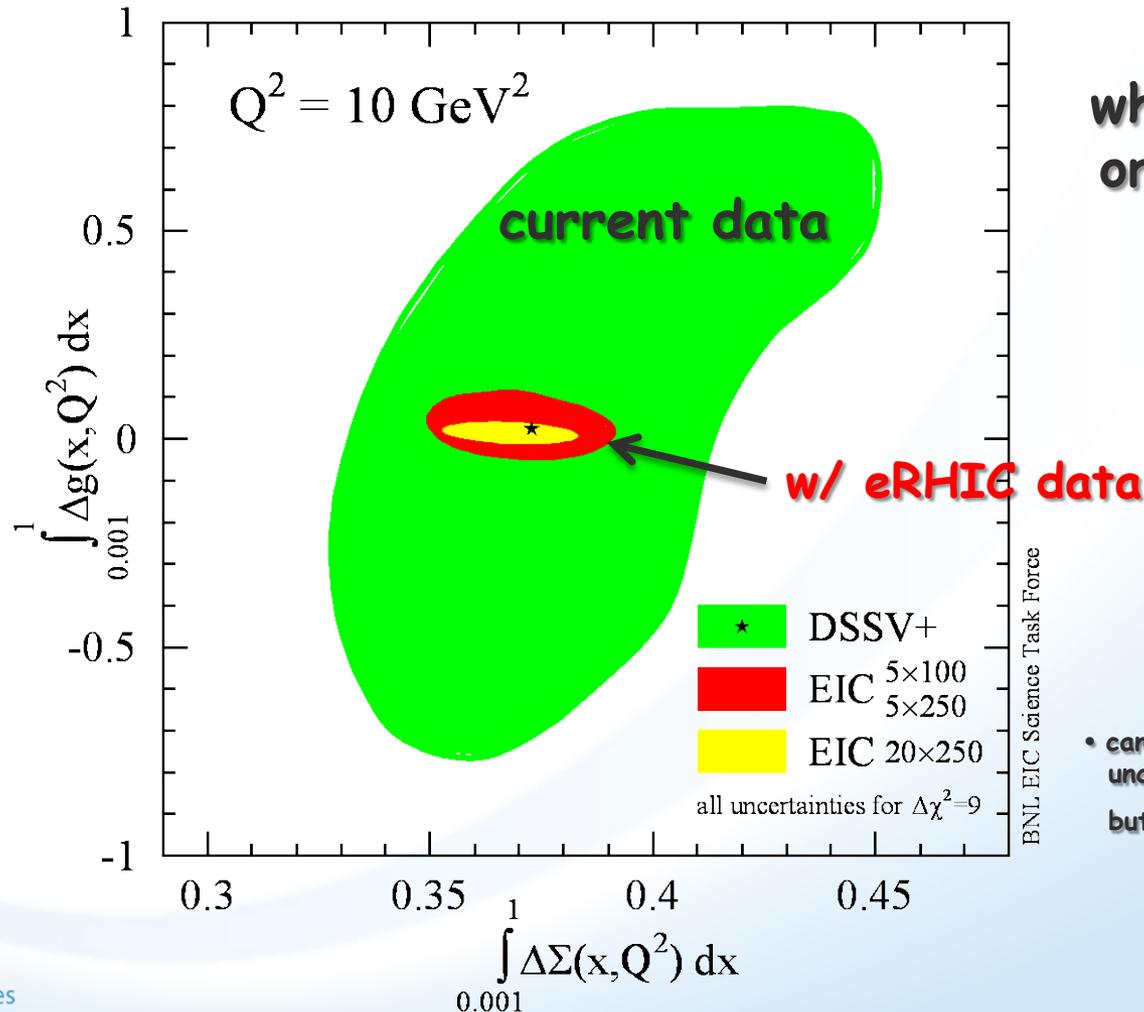
$$\frac{dg_1}{d \log(Q^2)} \sim -\Delta g(x, Q^2)$$

$$\Delta \Sigma(Q^2) = \int g_1(x, Q^2) dx = \int \Delta q_f(x, Q^2) dx$$



CAN WE SOLVE THE SPIN SUM RULE ?

$$\frac{1}{2}\hbar = \left\langle P, \frac{1}{2} \left| J_{QCD}^z \right| P, \frac{1}{2} \right\rangle = \sum_q \underbrace{\frac{1}{2} S_q^z}_{\text{total quark spin } \Delta\Sigma} + \underbrace{S_g^z}_{\text{gluon spin } \Delta g} + \sum_q \underbrace{(L_q^z + L_g^z)}_{\text{orbital angular momentum}}$$

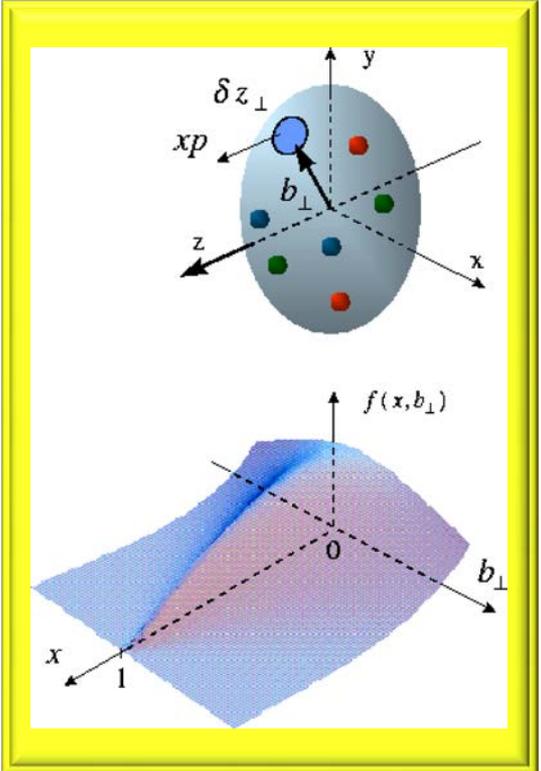


what about the orbital angular momentum?

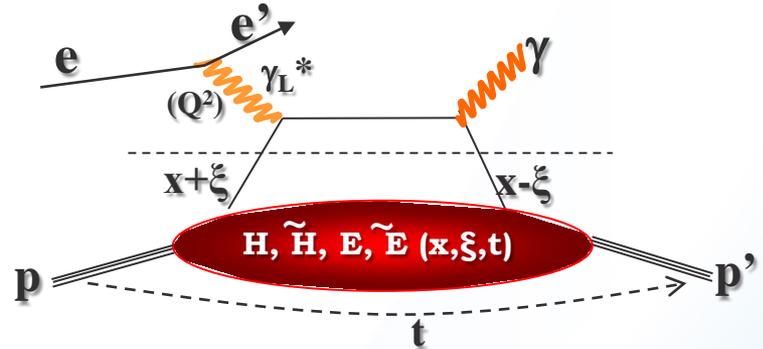
• can expect approx. 5-10% uncertainties on $\Delta\Sigma$ and Δg but need to control systematics

GENERALIZED PARTON DISTRIBUTIONS

the way to 3d imaging of the proton and the orbital angular momentum L_q & L_g



Measure them through exclusive reactions
golden channel: **DVCS**



Spin-Sum-Rule in PRF:

$$\frac{1}{2} = J_q^z + J_g^z = \frac{1}{2} \Delta\Sigma + \sum_q \mathcal{L}_q^z + J_g^z$$

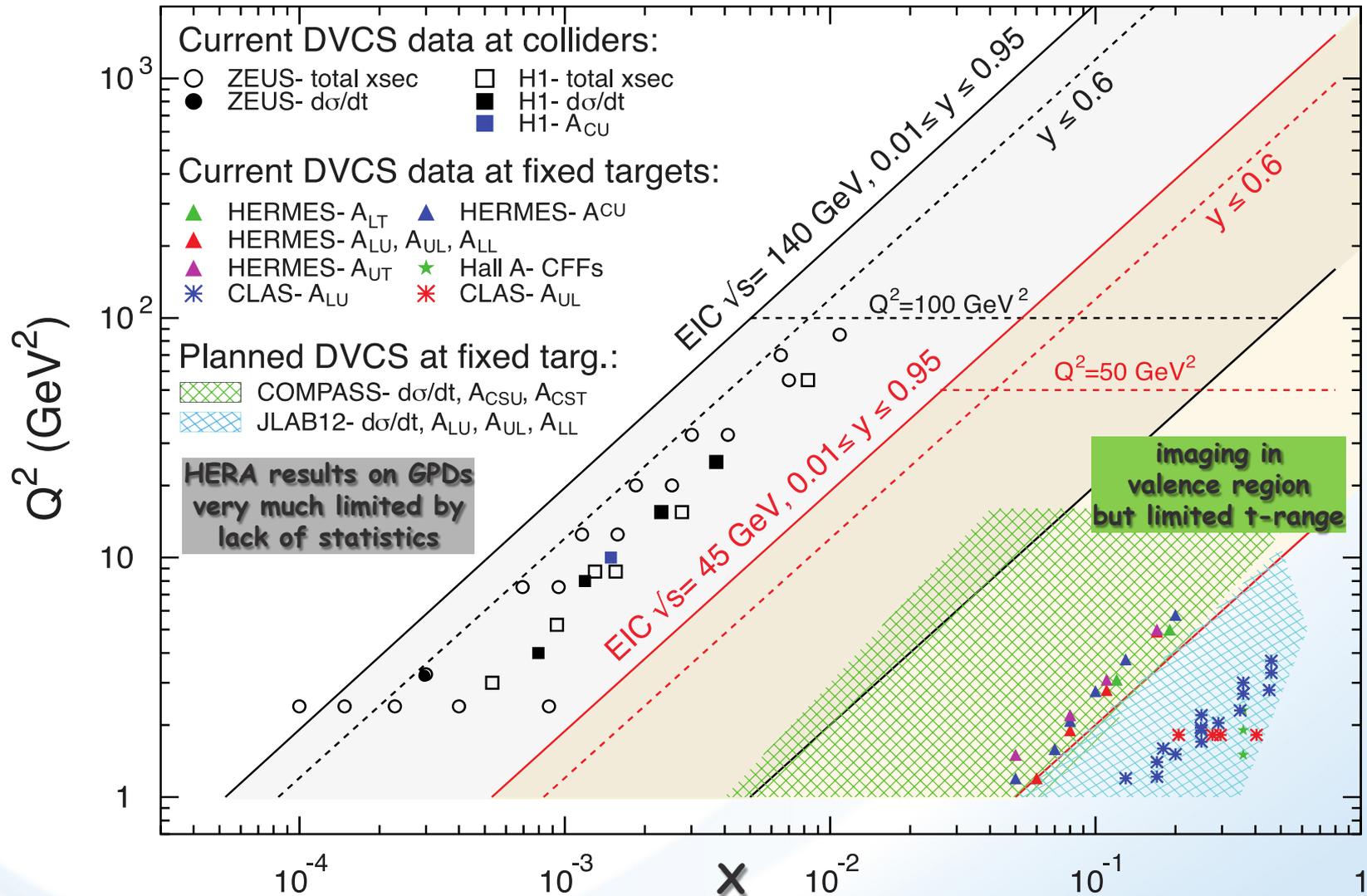
from g_1

$$J_{q,g}^z = \frac{1}{2} \left(\int_{-1}^1 x dx \left(H^{q,g} + E^{q,q} \right) \right)_{t \rightarrow 0}$$

GPDs:
Correlated quark momentum
and helicity distributions in
transverse space

responsible for orbital angular momentum

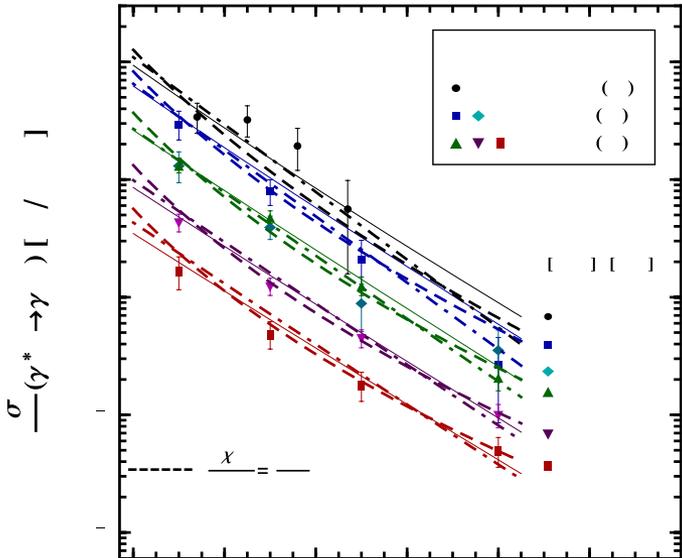
THE DVCS PHASE SPACE



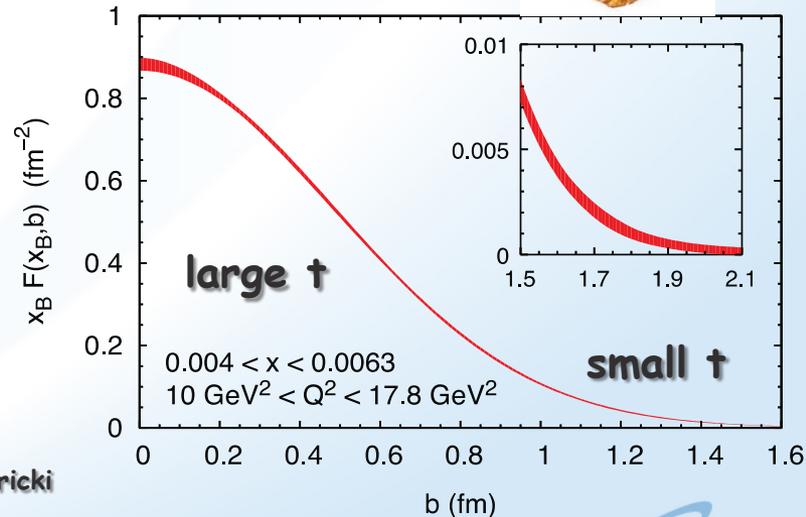
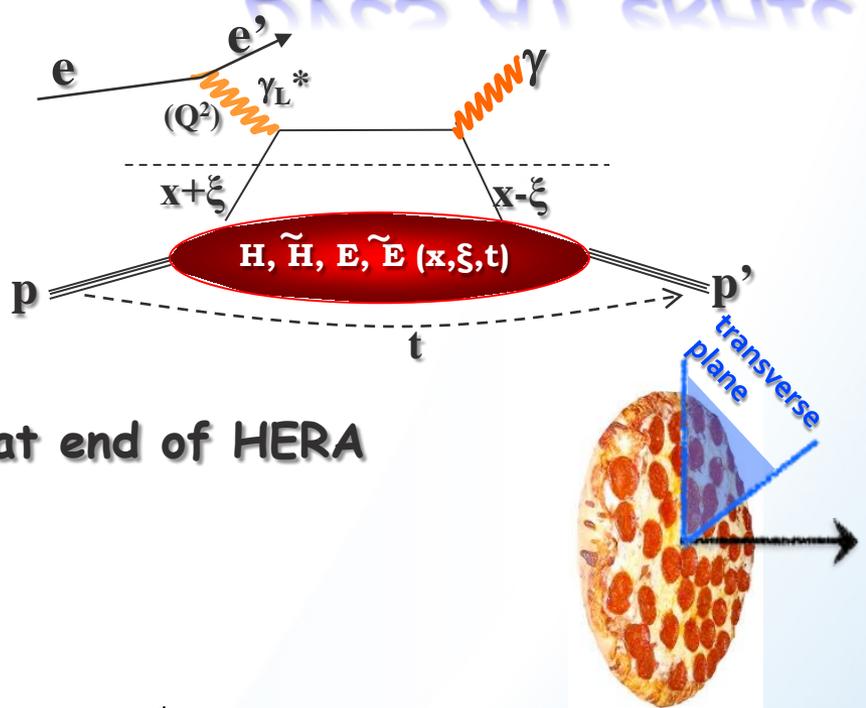
quantum numbers of final state → selects different GPD

DVCS: wide range of observables ($\sigma, A_{UT}, A_{LU}, A_{UL}, A_C$) to disentangle GPDs

DVCS AT eRHIC



DVCS data at end of HERA

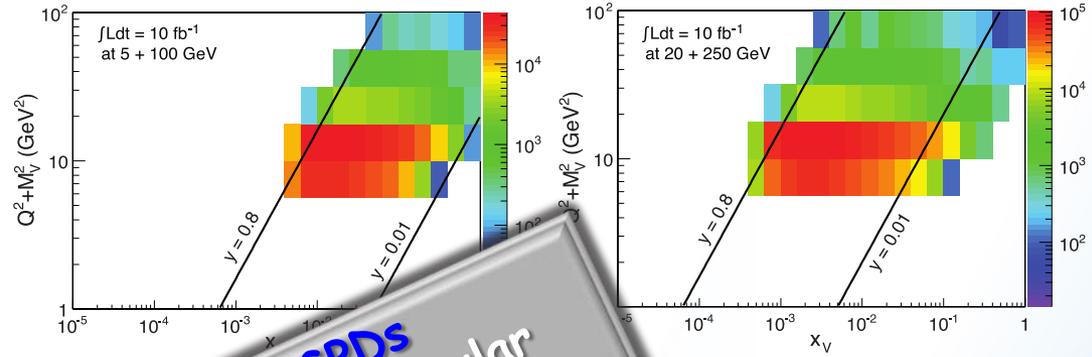


D. Mueller, K. Kumericki
S. Fazio, and ECA
[arXiv:1304.0077](https://arxiv.org/abs/1304.0077)

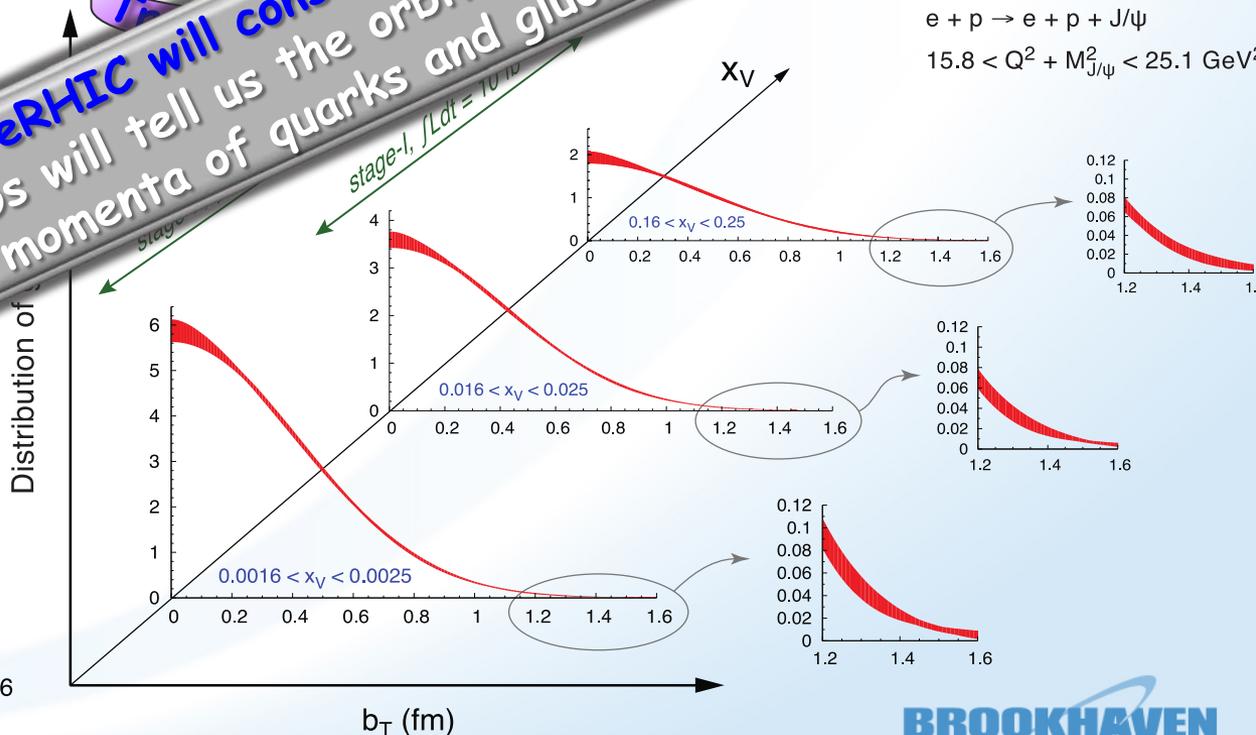
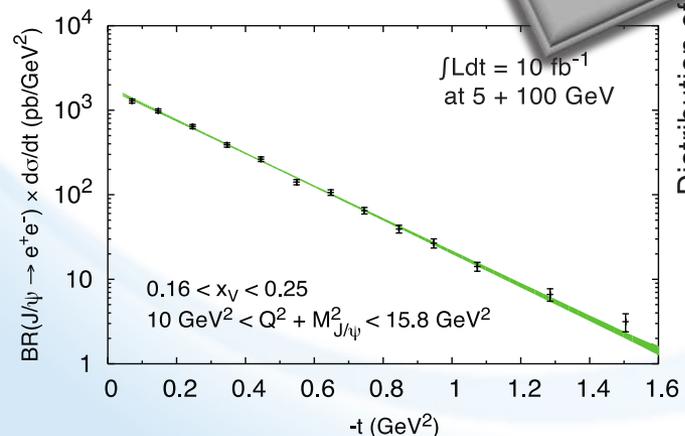
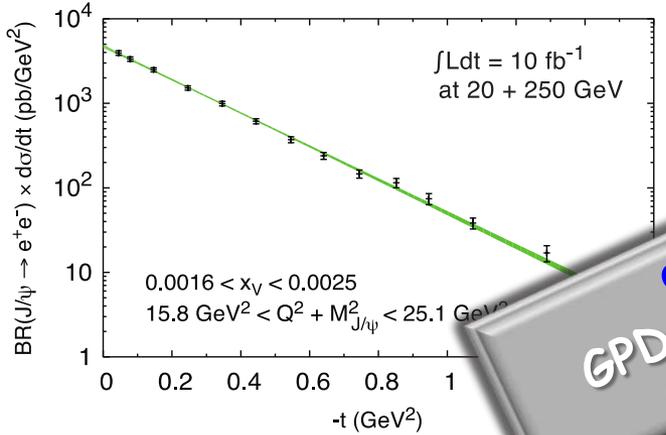
M. Diehl & ECA

To improve imaging on gluons
add J/ψ observables

- cross section
- A_{UT}
-



Fourier
eRHIC will constrain GPDs
GPDs will tell us the orbital angular
momenta of quarks and gluons



$e + p \rightarrow e + p + J/\psi$
 $15.8 < Q^2 + M_{J/\psi}^2 < 25.1 \text{ GeV}^2$

Where do we stand to realize EIC@RHIC



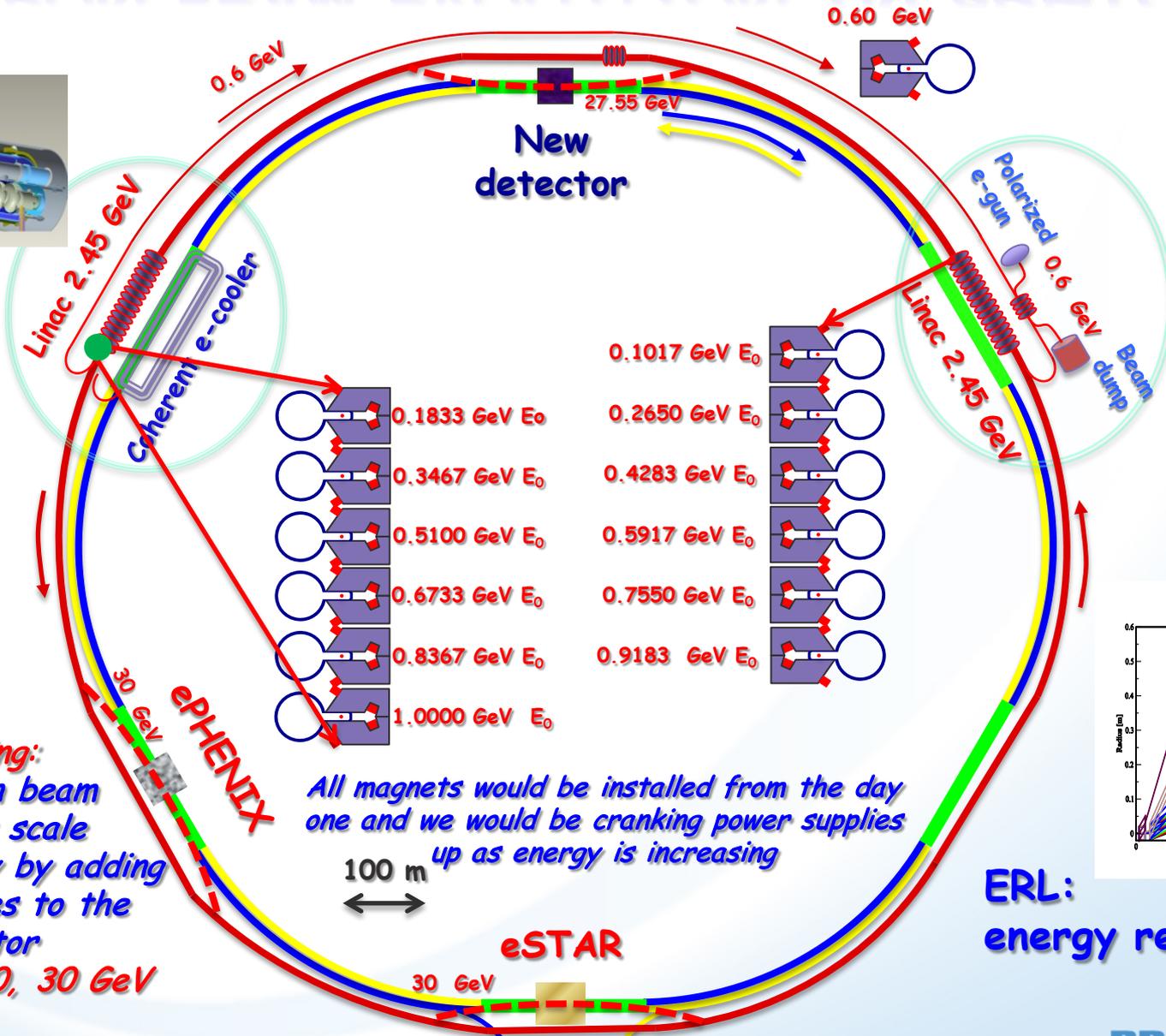
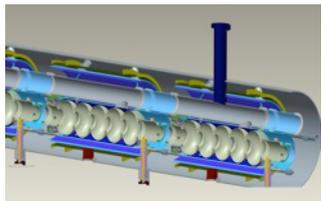
Latest Review:

NSAC 2013 Subcommittee Report on Scientific Facilities:

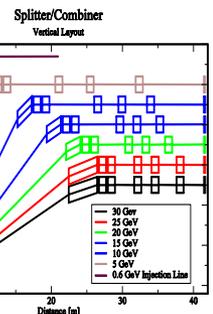
“The Subcommittee ranks an EIC as Absolutely Central in its ability to contribute to world-leading science in the next decade.”

“There are outstanding R&D issues that remain to be addressed in order to achieve performance metrics. Staging approaches to the EIC are also being explored by [BNL and JLab]. Both laboratories are actively addressing R&D issues and are making good progress.”

ELECTRON BEAM EVOLUTION IN eRHIC'S ERL



E/E ₀
0.0200
0.1017
0.1833
0.2650
0.3467
0.4283
0.5100
0.5917
0.6733
0.7550
0.8367
0.9183
1.0000



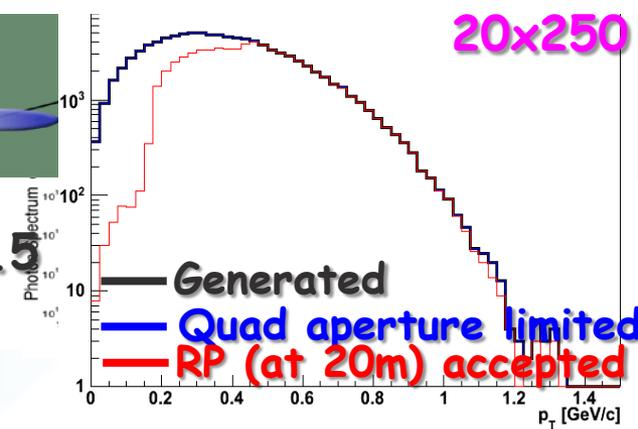
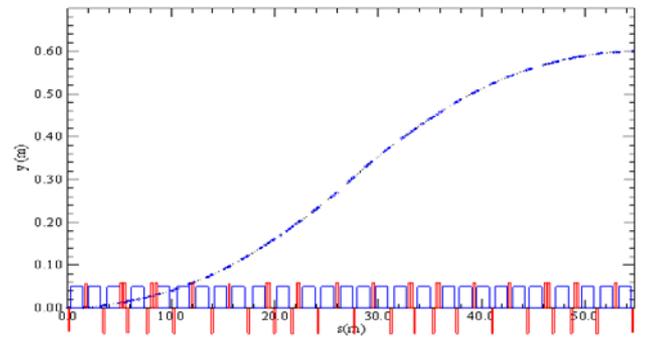
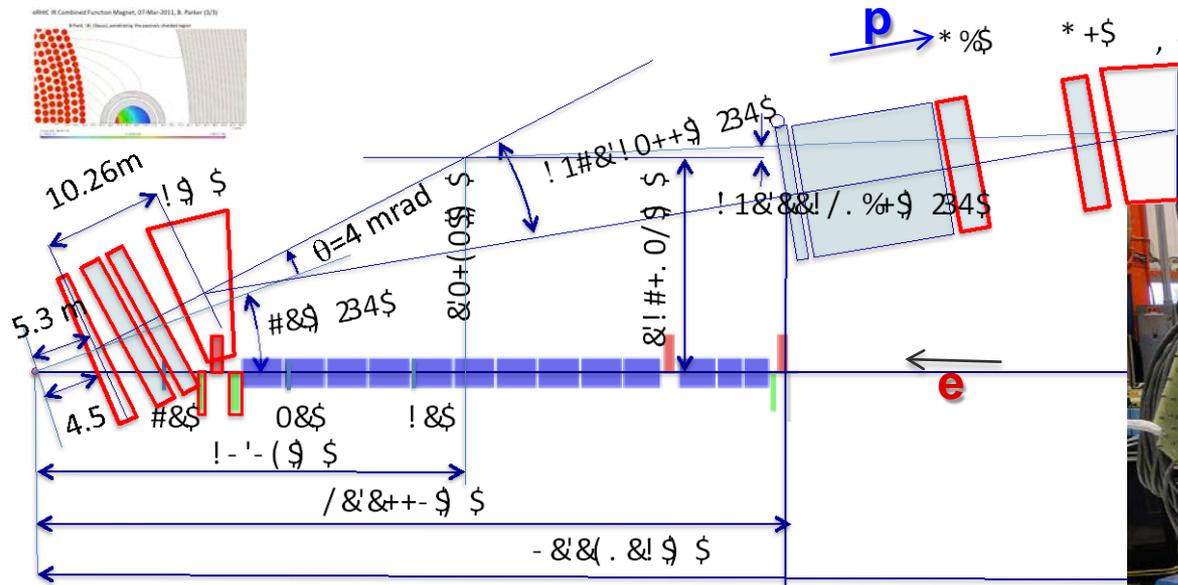
Staging:
All lepton beam energies scale proportionally by adding SRF cavities to the injector
 $E_0 = 5, 10, 20, 30 \text{ GeV}$

All magnets would be installed from the day one and we would be cranking power supplies up as energy is increasing

ERL:
energy recovery linac

Animation is by N. Tsoupas

eRHIC: HIGH-LUMINOSITY IR

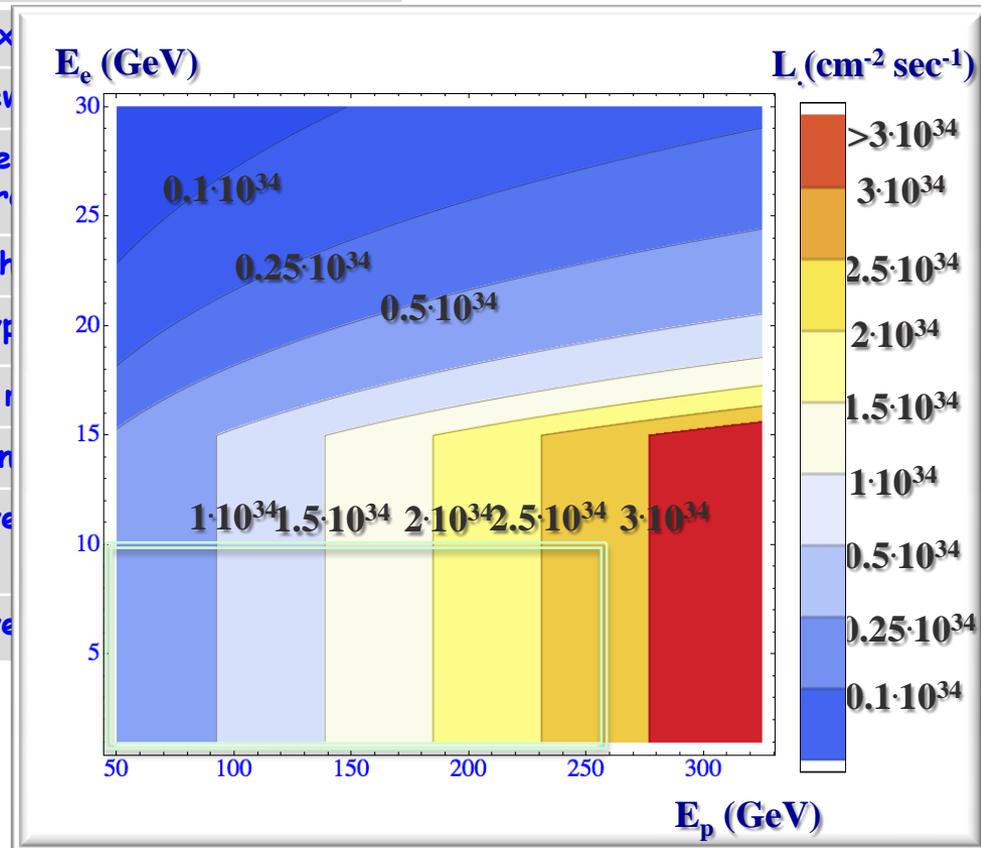


eRHIC - Geometry high-lumi IR with $\beta^*=5$ cm, $l^*=4.5$ and 10 mrad crossing angle $\rightarrow 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$

- 10 mrad crossing angle and crab-crossing
- High gradient (200 T/m) large aperture Nb₃Sn focusing magnets
- Arranged free-field electron pass through the hadron triplet magnets
- Integration with the detector: efficient separation and registration of low angle collision products
- Gentle bending of the electrons to avoid SR impact in the detector

eRHIC R&D HIGHLIGHTS AND LUMINOSITY

Challenge	Increase/reduction beyond the state of the art
Polarized electron gun	10 x
Coherent Electron Cooling	New
Multi-pass SRF ERL	5 x incre 30 x incre
Crab crossing	New for h
Understanding beam-beam effects	New typ
$\beta^*=5$ cm	5x r
Multi-pass SRF ERL	2-3 x in
Feedback for kink instability suppression	Novel
Space charge effect compensation	Novel



- Hourglass the pinch effects are included. Space charge effects are compensated.
- Energy of electrons can be selected at any desirable value at or below 30 GeV
- The luminosity does not depend on the electron beam energy below or at 20 GeV
- The luminosity falls as E_e^{-4} at energies above 20 GeV
- The luminosity is proportional to the hadron beam energy: $L \sim E_h/E_{\text{top}}$

Can we do with less ?

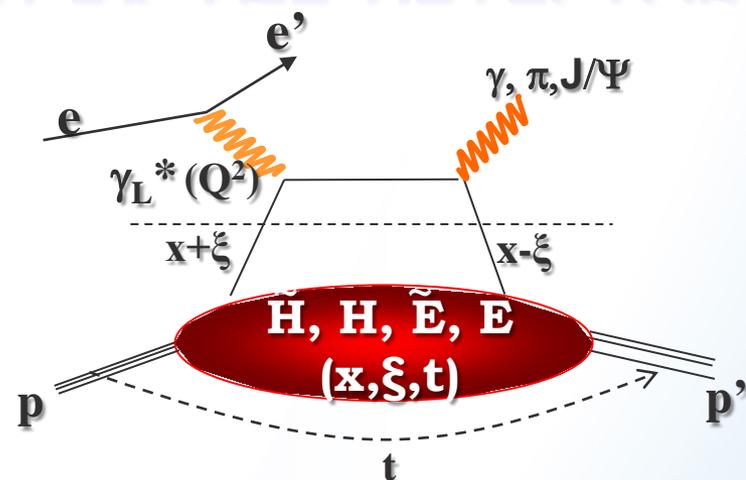
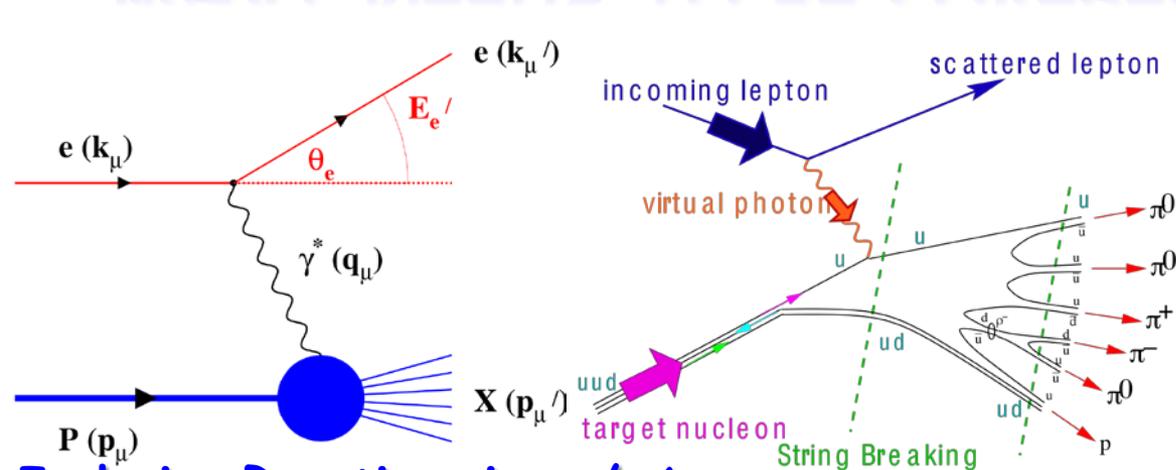
Challenge	Increase/reduction beyond the state of the art
Polarized electron gun	10 x increase
Coherent Electron Cooling	New concept
Multi-pass SRF ERL	5 x increase in current 30 x increase in energy
Crab crossing	New for hadron colliders
Understanding beam-beam effects	New type of collider
$\beta^*=5$ cm	5x reduction
Multi-pass SRF ERL	2-3 x in # of passes
Feedback for kink instability suppression	Novel concept
Space charge effect compensation	Novel concept

no cooling not possible lose all exclusive physics
 → more conventional cooling
 → reduced lumi 10^{33}

need to remove crossing angle and go to head on collisions
 → beams get separated by dipoles close to IR
 → impact on synchrotron rad.

Currently investigate machine and IR design with a less ambiguous R&D to reach $10^{33} \text{ cm}^{-2}\text{s}^{-1}$ with 10 GeV initial beam energy

WHAT NEEDS TO BE COVERED BY THE DETECTOR



Inclusive Reactions in ep/eA :

- Physics: Structure Fcts.: g_1, F_2, F_L
- Very good electron id \rightarrow find scattered lepton
- Momentum/energy and angular resolution of e' critical
- scattered lepton \rightarrow kinematics

Semi-inclusive Reactions in ep/eA :

- Physics: TMDs, Helicity PDFs \rightarrow flavor separation, dihadron-corr., ...
 \rightarrow Kaon asymmetries, cross sections
- Excellent particle ID: π^\pm, K^\pm, p^\pm separation over a wide range in η
- full Φ -coverage around γ^*
- Excellent vertex resolution \rightarrow Charm, Bottom identification

Exclusive Reactions in ep/eA :

- Physics: GPDs, proton/nucleus imaging, DVCS, excl. VM/PS prod.
- Exclusivity \rightarrow large rapidity coverage \rightarrow rapidity gap events
- \searrow reconstruction of all particles in event
- high resolution in $t \rightarrow$ Roman pots

THE eRHIC DETECTOR CONCEPT

Extremely wide physics program puts **stringent requirements** on detector performance

- ❑ high acceptance $-5 < \eta < 5$
- ❑ good PID (π, K, p and lepton) and vertex resolution
- ❑ same rapidity coverage for tracking and calorimeter
 - good momentum resolution, lepton PID
- ❑ low material density because of low scattered lepton p
 - minimal multiple scattering and brems-strahlung
- ❑ very forward electron and proton/neutron detection
- ❑ Fully integrated in machine IR design

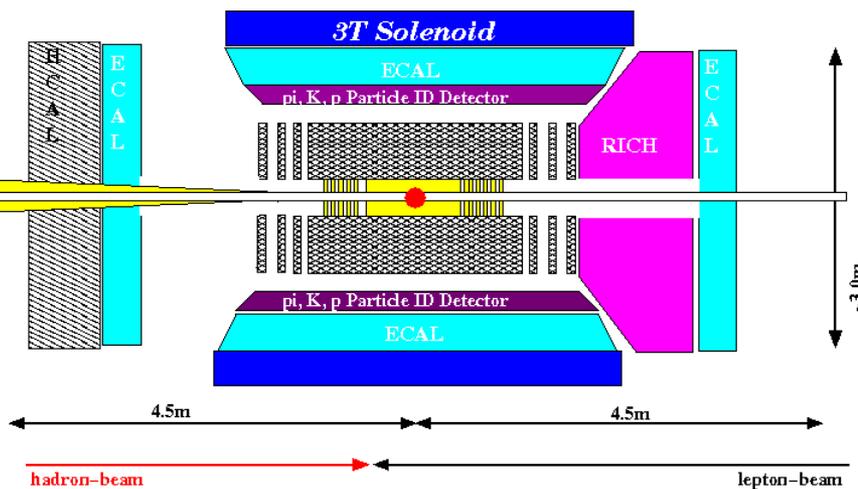
Summary:

Full Geant Model based on Generic EIC R&D detector concepts

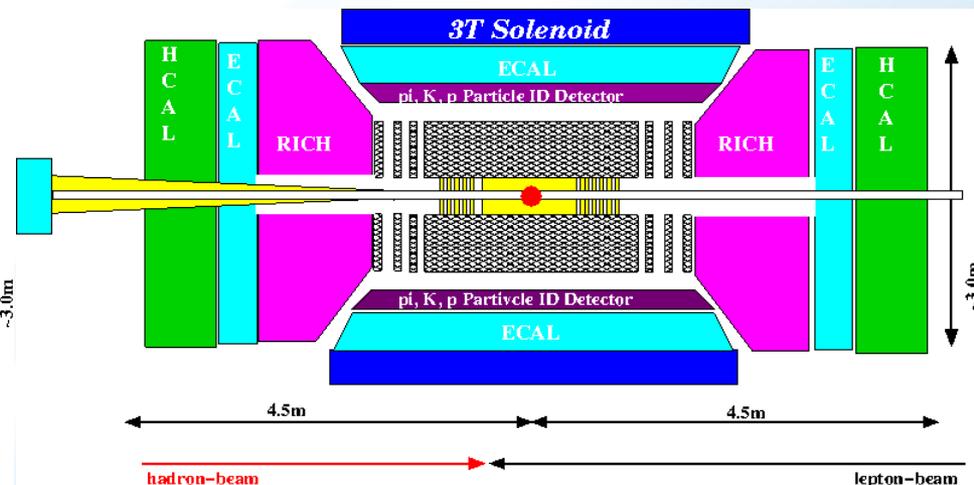
https://wiki.bnl.gov/eic/index.php/DIS: What_is_important

https://wiki.bnl.gov/eic/index.php/ERHIC_Dedicated_Detector_Design

Phase-I (5 - 10 GeV):



Phase-II (>10 GeV):



barrel silicon tracker:

- MAPS technology: $\sim 20 \times 20 \text{ mm}^2$ chips, $\sim 20 \text{ } \mu\text{m}$ 2D pixels
- 6 layers at [30..160] mm radius
- 0.37% X_0 in acceptance per layer simulated precisely;
- digitization: single discrete pixels, one-to-one from MC points

forward/backward silicon trackers:

- 2x7 disks with up to 280 mm radius
- N sectors per disk; 200 μm silicon-equivalent thickness
- digitization: discrete $\sim 20 \times 20 \text{ } \mu\text{m}^2$ pixels

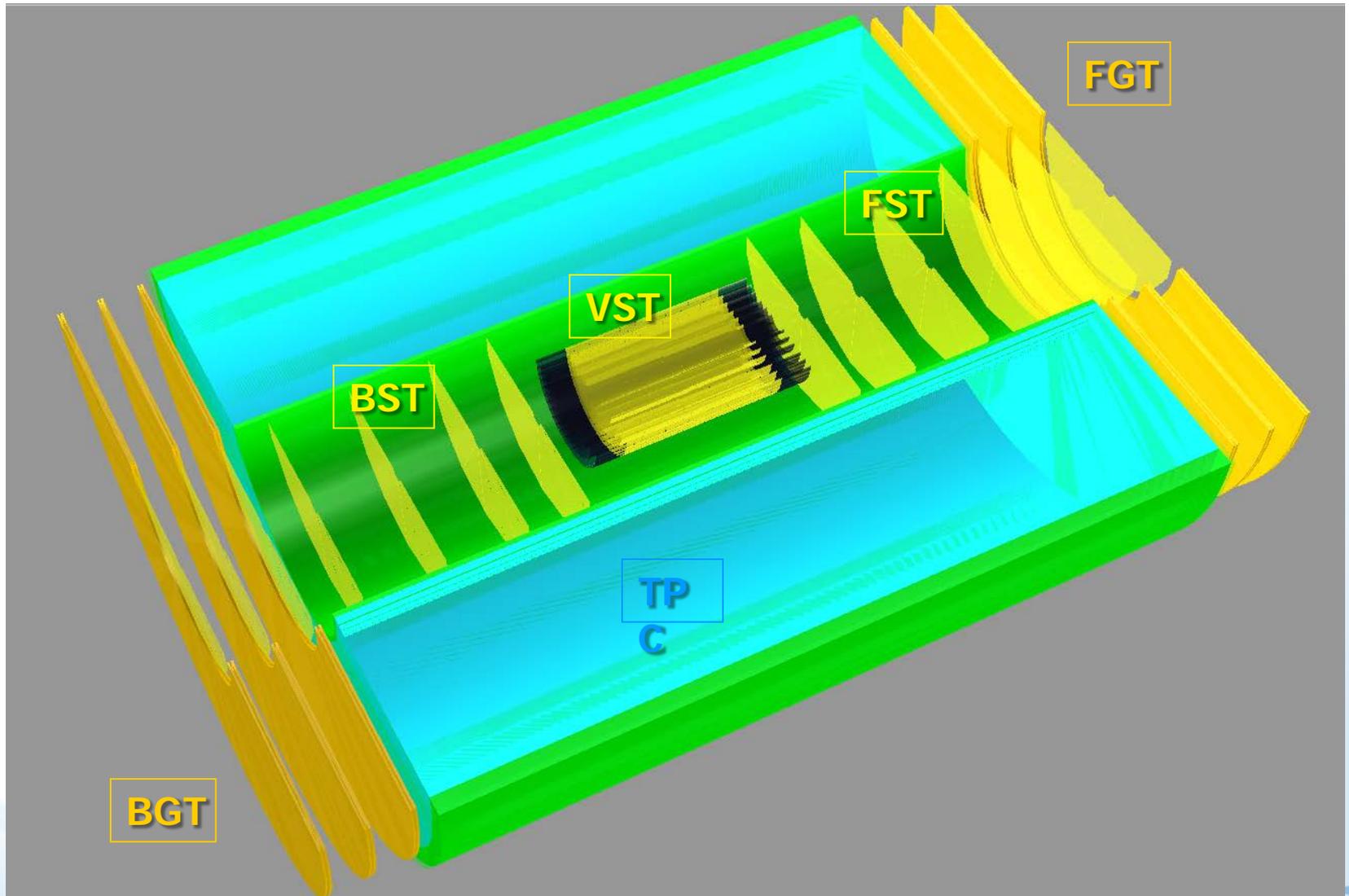
TPC:

- $\sim 2\text{m}$ long; gas volume radius [300..800] mm
- 1.2% X_0 IFC, 4.0% X_0 OFC; 15.0% X_0 aluminum endcaps
- digitization: idealized, assume 1x5 mm GEM pads

GEM trackers:

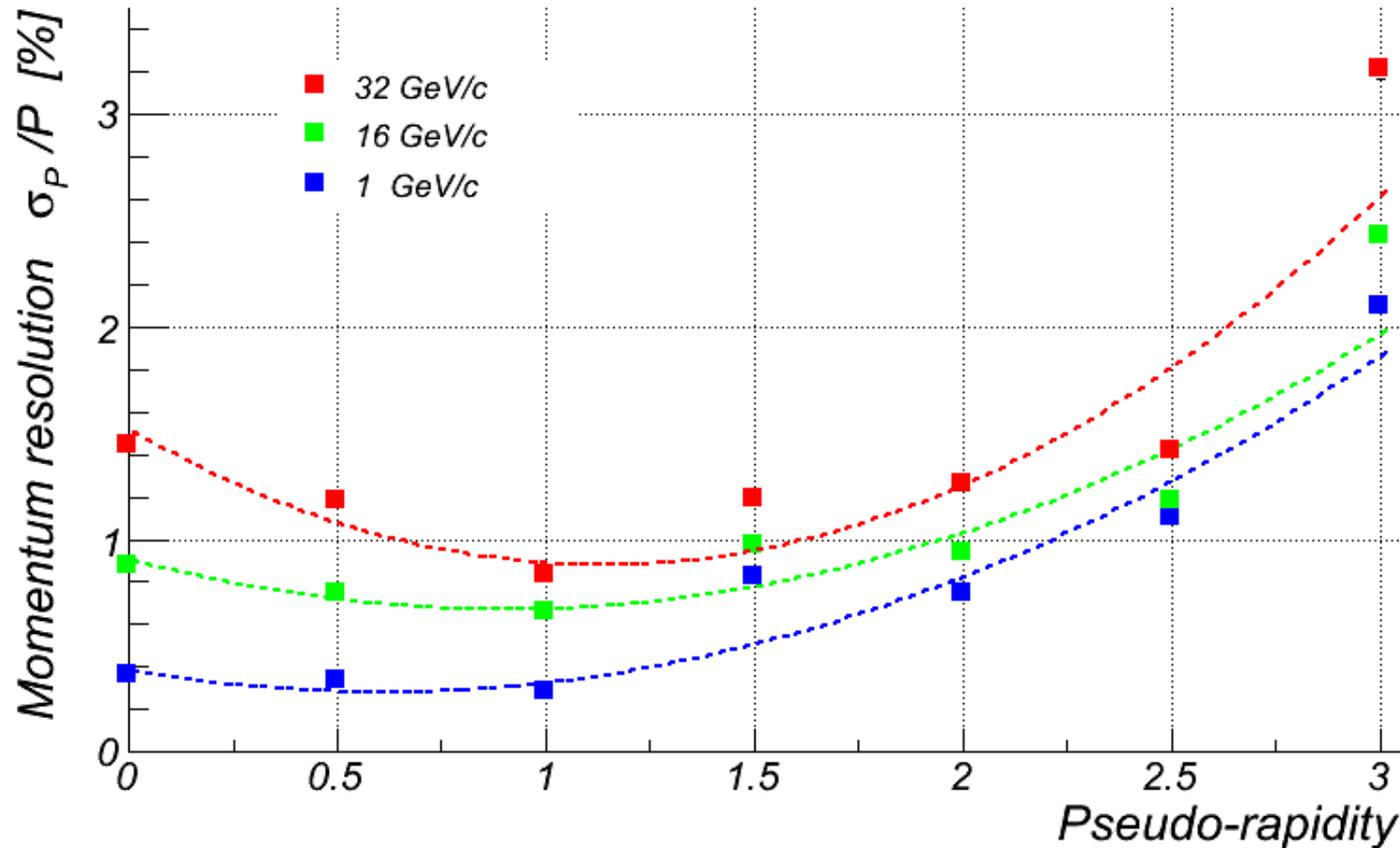
- 3 disks behind the TPC endcap
- STAR FGT design
- digitization: 100 mm resolution in X&Y; gaussian smearing

TRACKER ZOOMED VIEW



MOMENTUM RESOLUTION STUDY (1)

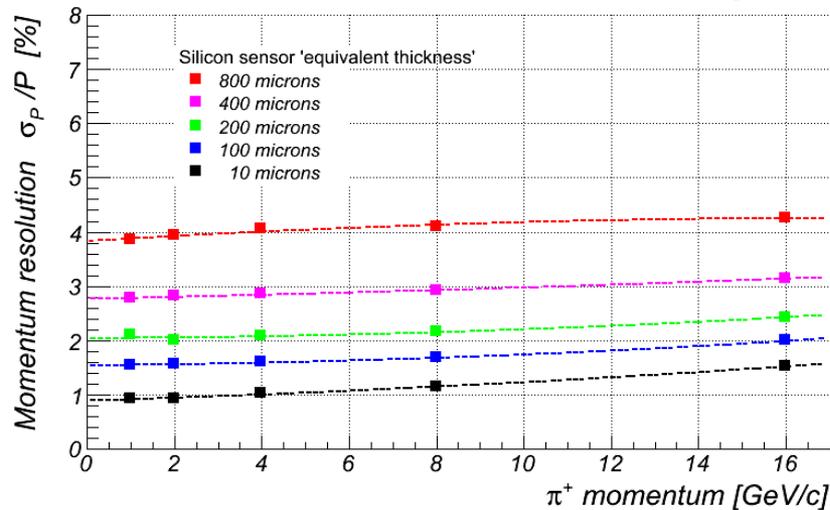
π^+ track momentum resolution vs. pseudo-rapidity



-> expect 2% or better momentum resolution in the whole kinematic range

MOMENTUM RESOLUTION STUDY (2)

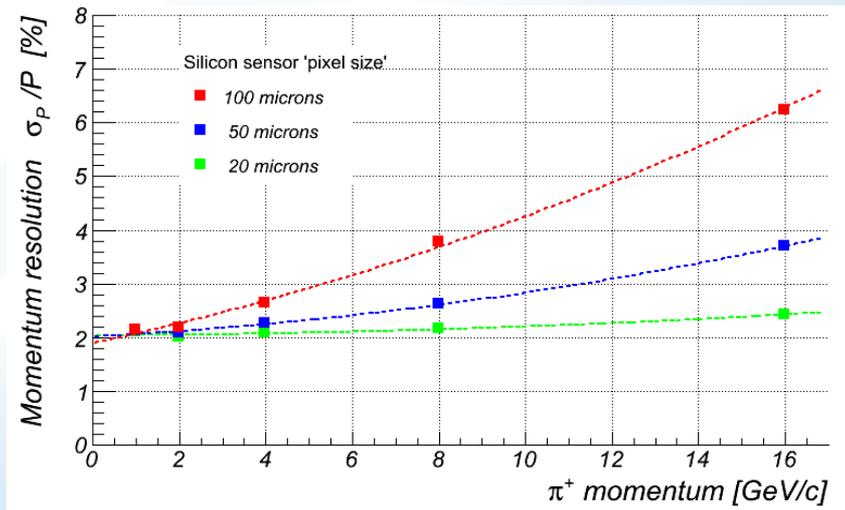
π^+ track momentum resolution at $\eta = 3.0$ vs. Silicon thickness



-> ~flat over inspected momentum range because of very small Si pixel size

π^+ track momentum resolution at $\eta = 3.0$ vs. Silicon pixel size

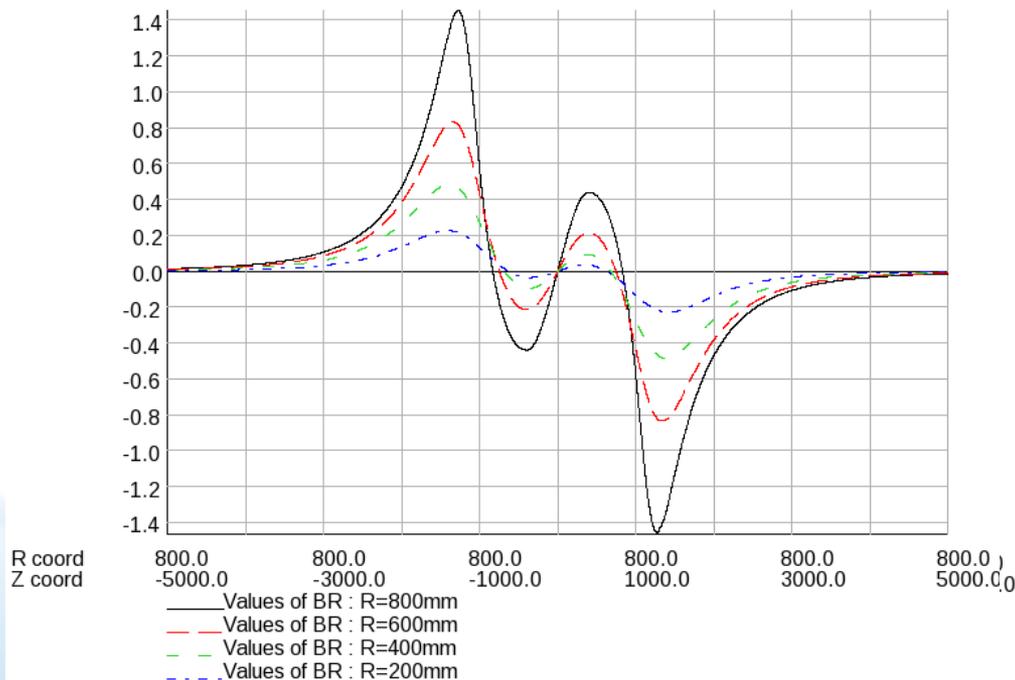
-> 20 micron pixel size is essential to maintain good momentum resolution



EIC SOLENOID MODELING

main requirements:

- Yield large enough bending for charged tracks at large η
- Keep field inside TPC volume as homogeneous as possible
- Keep magnetic field inside RICH volume(s) small



-> use OPERA-3D/2D software

Presently used design: **MRS-B1**

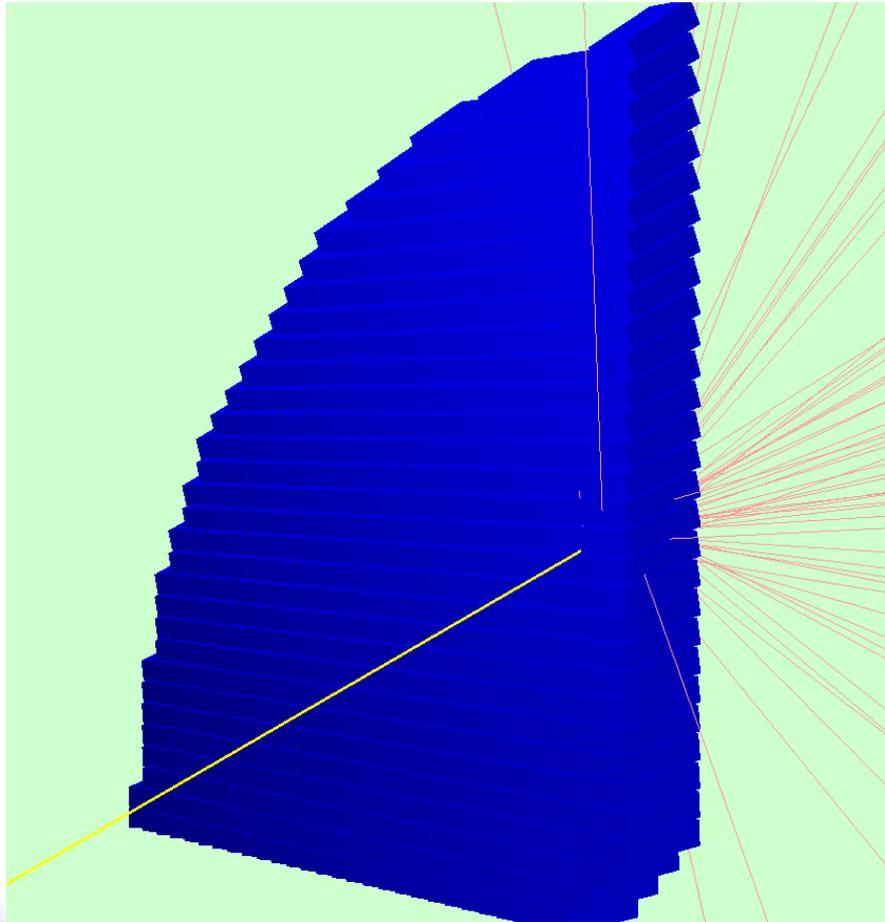
Total Length : 2.4 m

Inner Radius : 1.0 m

Outer Radius : 1.1 m

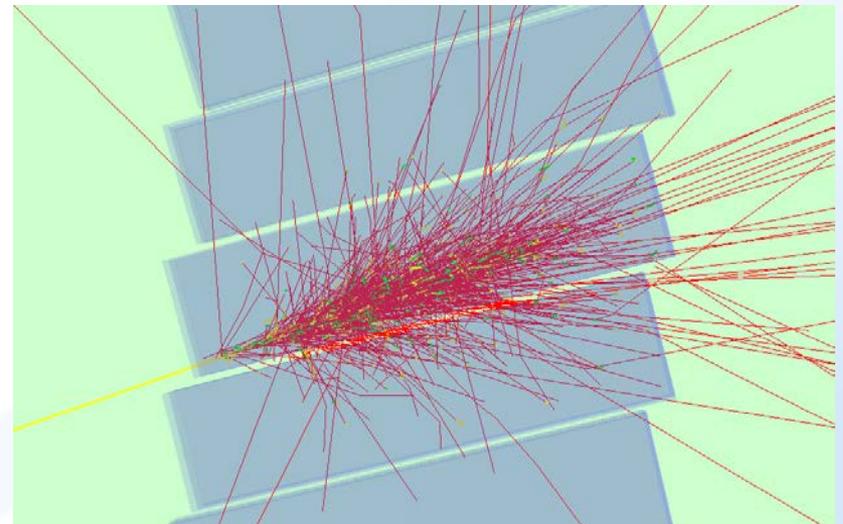
Central B field: 3.0 T

BACKWARD EM CALORIMETER (BEMC)



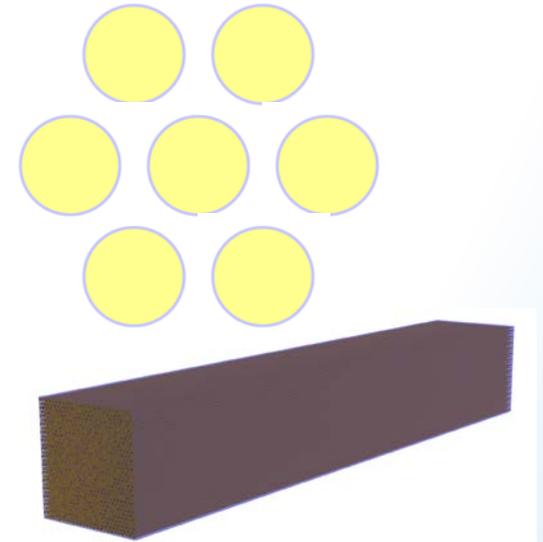
- PWO-II, layout a la CMS & PANDA
- -2500mm from the IP
- both projective and non-projective geometry implemented
- digitization based on PANDA R&D

10 GeV/c electron hitting one of the four BEMC quadrants



Same event (details of shower development)

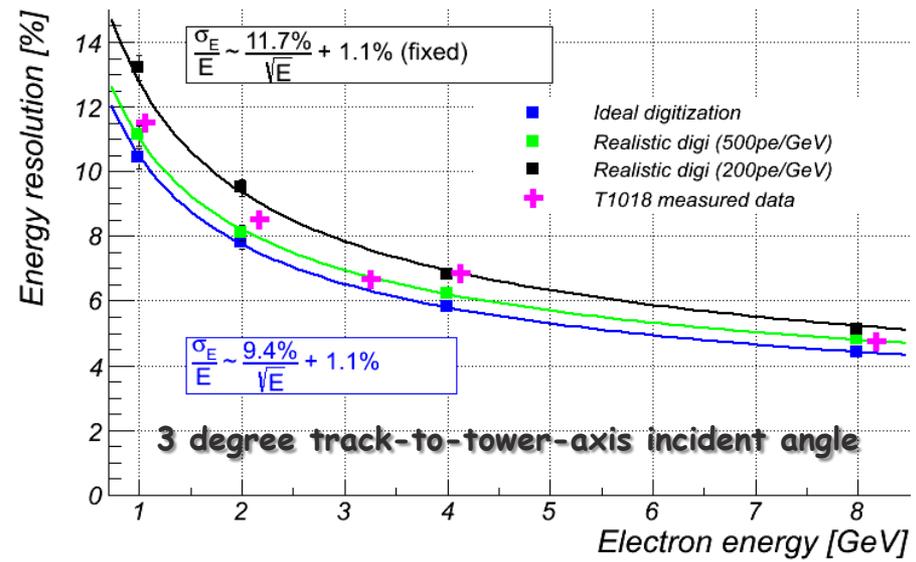
FORWARD EM CALORIMETER (FEMC)



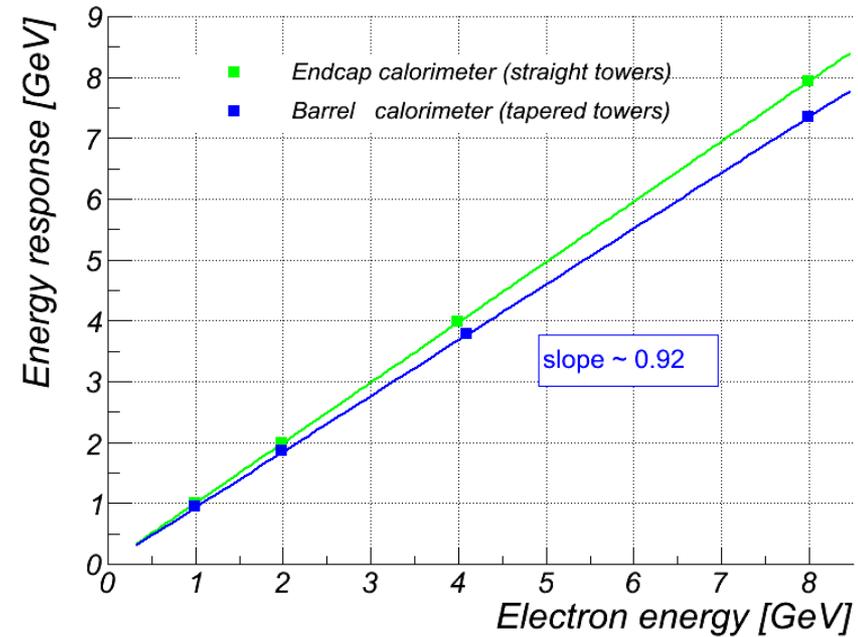
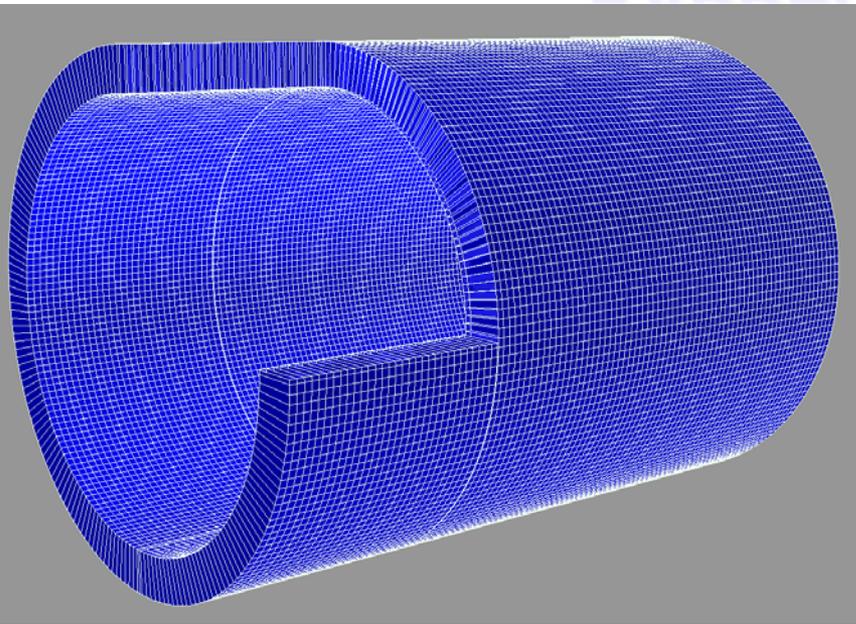
tower (and fiber) geometry described precisely

- ❑ tungsten powder scintillating fiber sampling calorimeter technology
- ❑ +2500mm from the IP; non-projective geometry
- ❑ sampling fraction for e/m showers ~2.6%
- ❑ “medium speed” simulation (up to energy deposit in fiber cores)
- ❑ reasonably detailed digitization: “ideal” clustering code
- ❑ “Realistic” digitization: 40MHz SiPM noise in 50ns gate;
- ❑ 4m attenuation length; 5 pixel single tower threshold;
- ❑ 70% light reflection on upstream fiber end;

-> good agreement with original MC studies and measured data



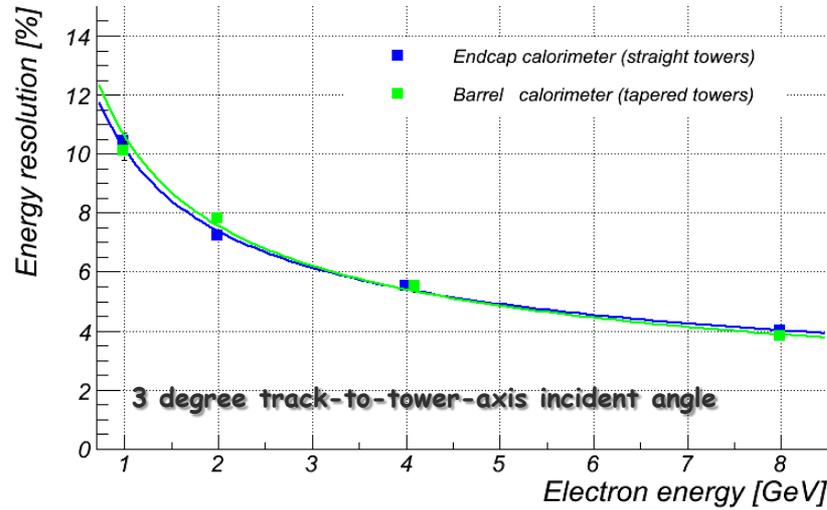
BARREL EM CALORIMETER (CEMC)



- same tungsten powder + fibers technology as FEMC,
- ... but towers are tapered
- non-projective

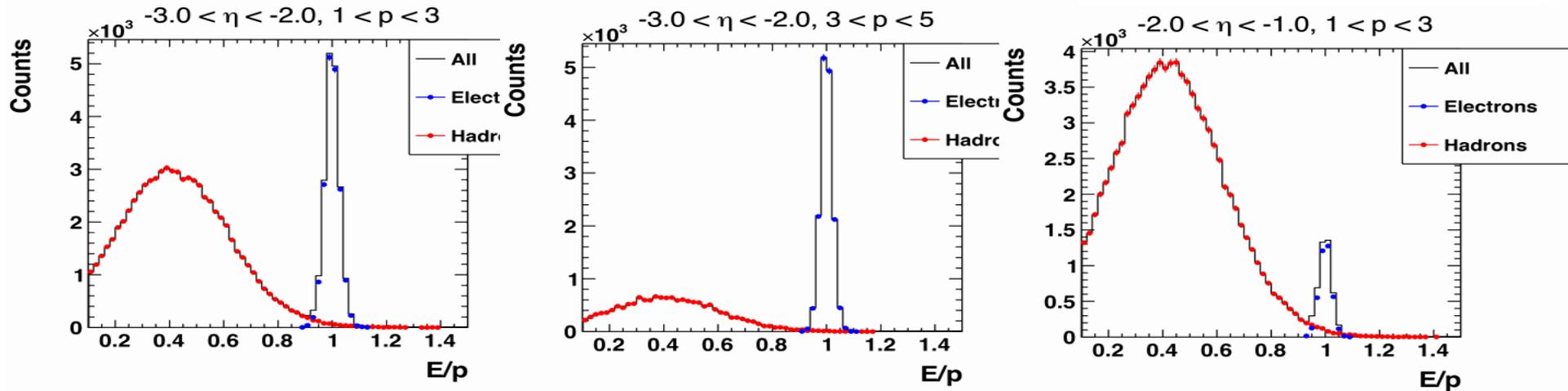
⇒ barrel calorimeter collects less light, but response (at a fixed 3° angle) is perfectly linear

→ simulation does not show any noticeable difference in energy resolution between straight and tapered tower calorimeters



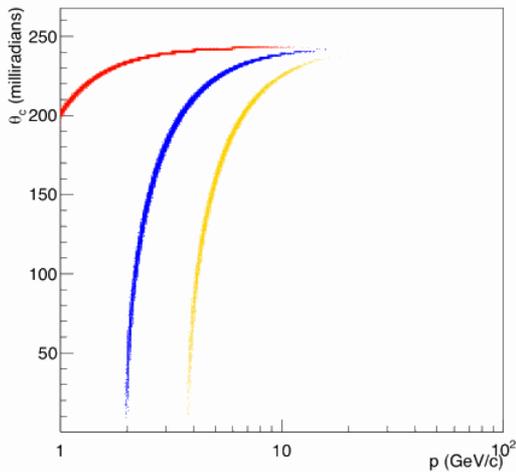
LEPTON-HADRON SEPARATION VIA E/P

all plots: 10GeV x 100GeV beams

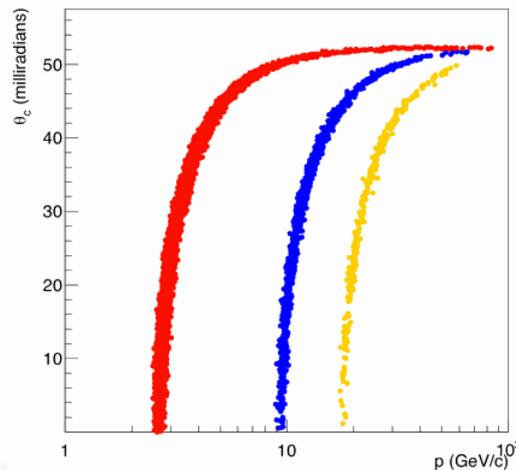


HADRON IDENTIFICATION WITH RICH

Cherenkov angle vs. p for aerogel ($n = 1.0304$)



Cherenkov angle vs. p for C4F10 ($n = 1.00137$)

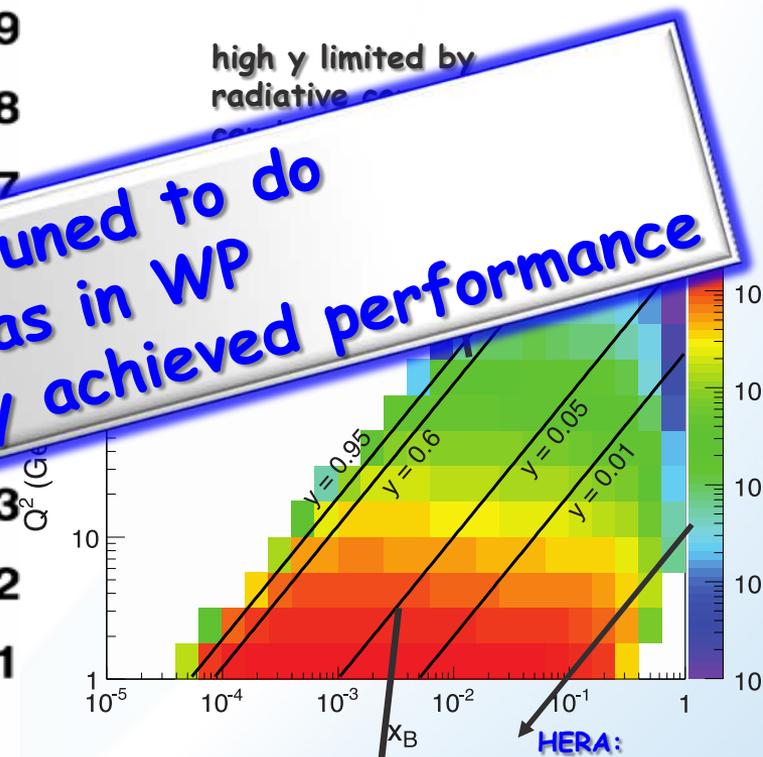
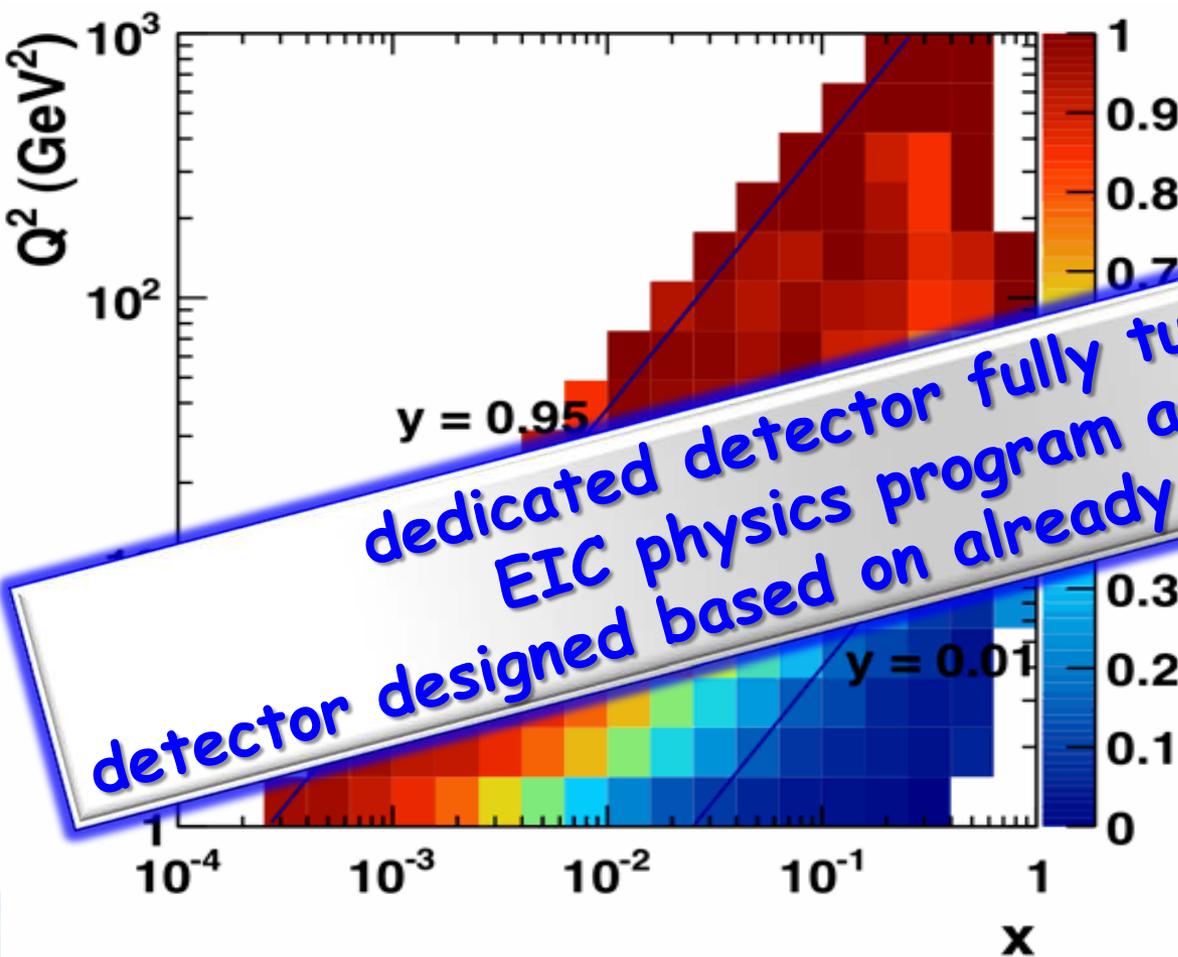


consider hadrons in
pseudo-rapidity range
 $\sim [1.0 \dots 3.0]$

-> pion/kaon/proton identification should be possible up
to momenta ~ 40 GeV/c

MIGRATION IN (X, Q^2) BINS

10 GeV x 100 GeV beams

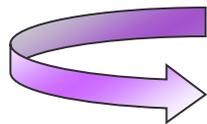


low y -coverage:
limited by E'_e resolution
→ hadron method
→ or change beam energy

-> "survival probability" is above ~80% in the region, where tracking has superior resolution

FROM RHIC TO eRHIC

- ❑ 2013/14: Develop physics case and preliminary design for eRHIC.
- ❑ Run RHIC annually until 2016; thereafter every other year.
- ❑ Use reduced RHIC running to complete accelerator and detector R&D, detailed design and prototyping for eRHIC.
- ❑ Construct sPHENIX.
- ❑ 2017: Install electron cooling
- ❑ 2018/19: Perform beam energy scan II.
- ❑ 2020/21: Perform high luminosity jet measurements.
- ❑ 2022-24: RHIC shutdown and conversion into eRHIC (10 GeV e-beam); upgrade of STAR/sPHENIX to eRHIC detectors.
- ❑ 2024: Commission eRHIC.



**of course all subject to
funding by DOE**

TAKE AWAY MESSAGE

The eRHIC will profoundly impact our understanding of QCD with its high energy, high luminosity eA and polarized collisions

ACHIEVED
if
eRHIC is built

entire science program uniquely tied to a future high energy polarized electron ion collider never been measured before & never without

program as and it will be critical to form a eRHIC community beyond RHIC

eRHIC machine design ambitious, but will push collider technologies in several regions

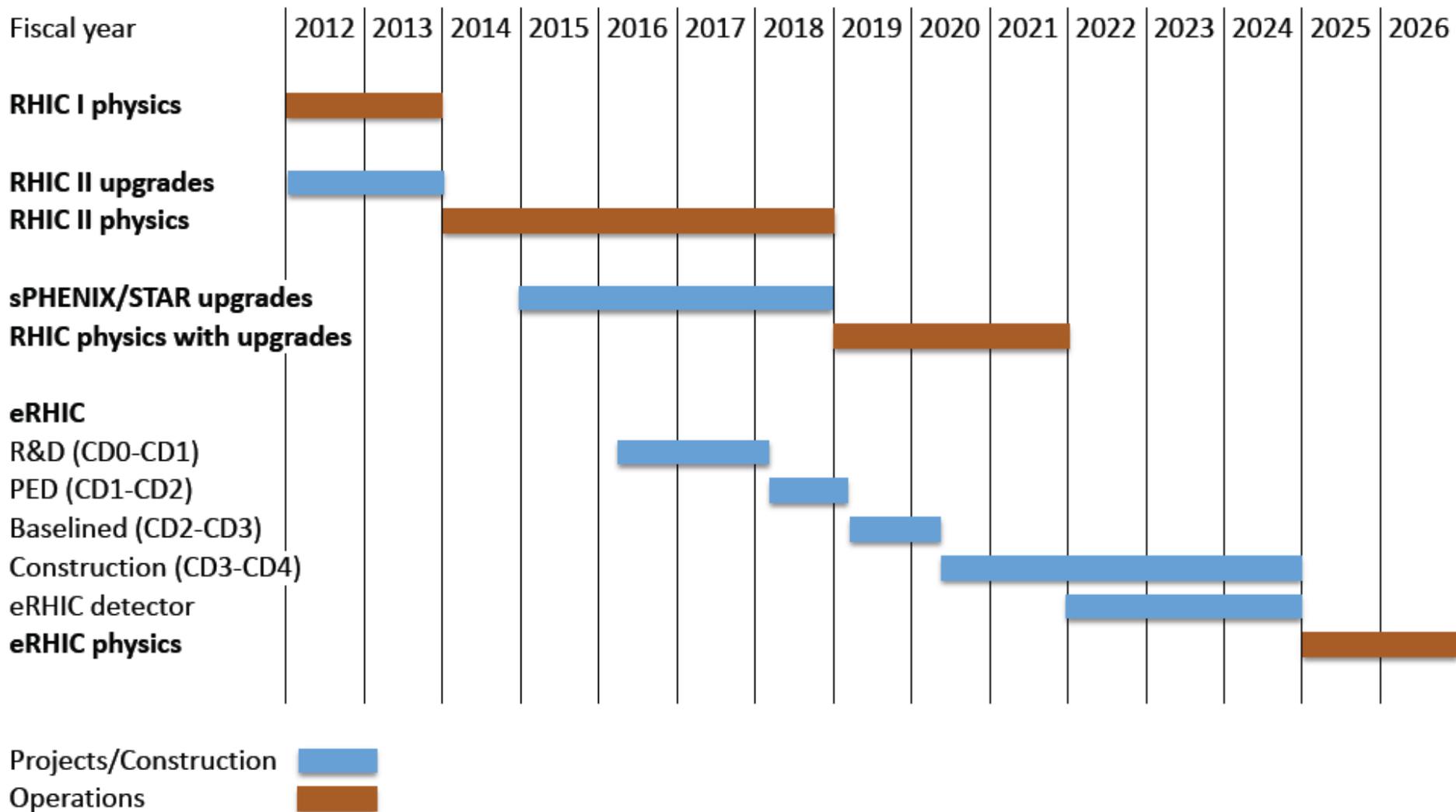
BACKUP

INTERNATIONAL CONTEXT

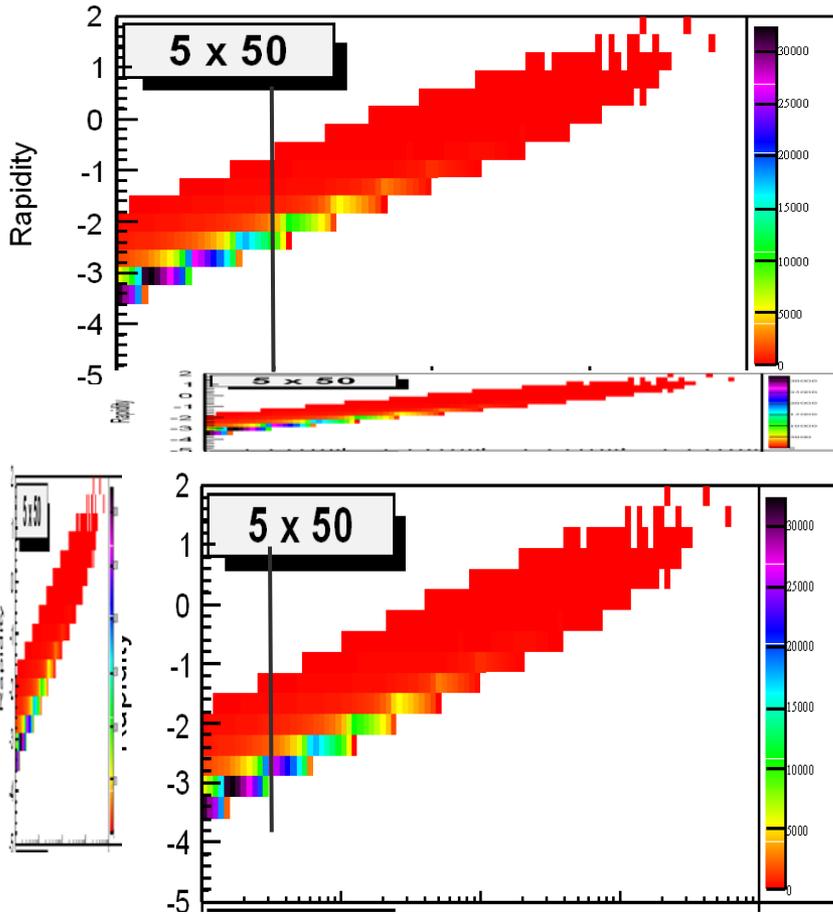
Electron-“Ion” colliders in the past and future:

	HERA@DESY	LHeC@CERN	eRHIC@BNL	MEIC@JLab	HIAF@CAS	ENC@GSI
E_{CM} (GeV)	320	800-1300	45-175	12-140	12 → 65	14
proton x_{min}	1×10^{-5}	5×10^{-7}	3×10^{-5}	5×10^{-5}	7×10^{-3} → 3×10^{-4}	5×10^{-3}
ion	p	p to Pb	p to U	p to Pb	p to U	p to ~ ^{40}Ca
polarization	-	-	p, ^3He	p, d, ^3He (^6Li)	p, d, ^3He	p,d
L [$\text{cm}^{-2} \text{s}^{-1}$]	2×10^{31}	10^{33}	10^{33-34}	10^{33-34}	$10^{32-33} \rightarrow 10^{35}$	10^{32}
IP	2	1	2+	2+	1	1
Year	1992-2007	2022 (?)	2022	Post-12 GeV	2019 → 2030	upgrade to FAIR

POSSIBLE SCHEDULE TO REALIZE eRHIC



LEPTON KINEMATICS

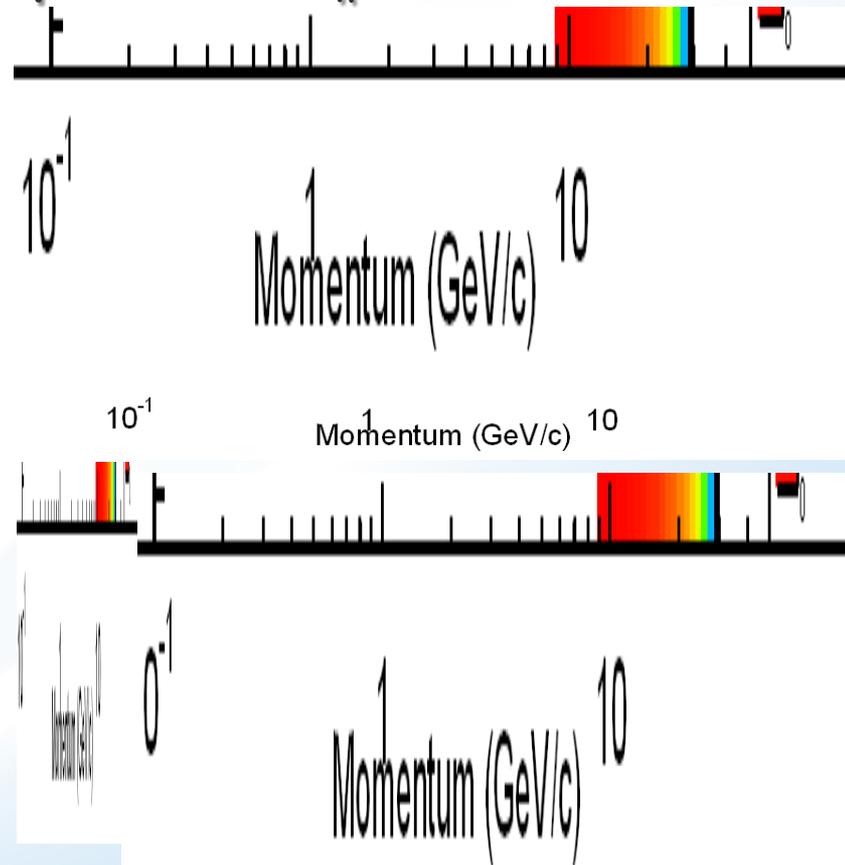


low Q^2 coverage
critical for
saturation physics

\sqrt{s}

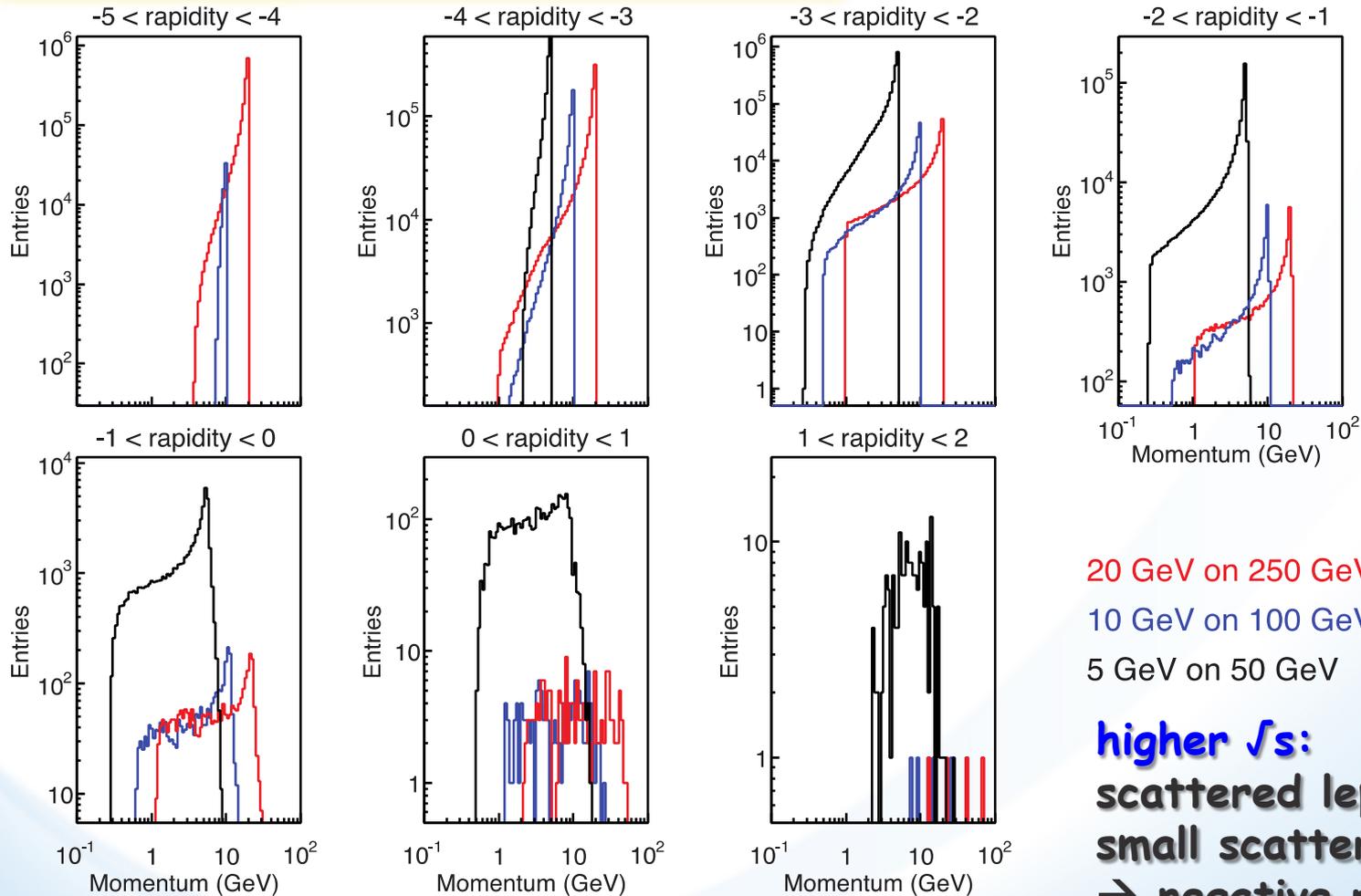
Increasing Lepton Beam Energy:
 5 GeV: $Q^2 \sim 1 \text{ GeV} \rightarrow \eta \sim -2$
 10 GeV: $Q^2 \sim 1 \text{ GeV} \rightarrow \eta \sim -4$

highest E'_e at most negative rapidities
 independent of E_h



SCATTERED LEPTON KINEMATICS

CUTS: $Q^2 > 0.1 \text{ GeV}^2$ && $0.01 < y < 0.95$



20 GeV on 250 GeV

10 GeV on 100 GeV

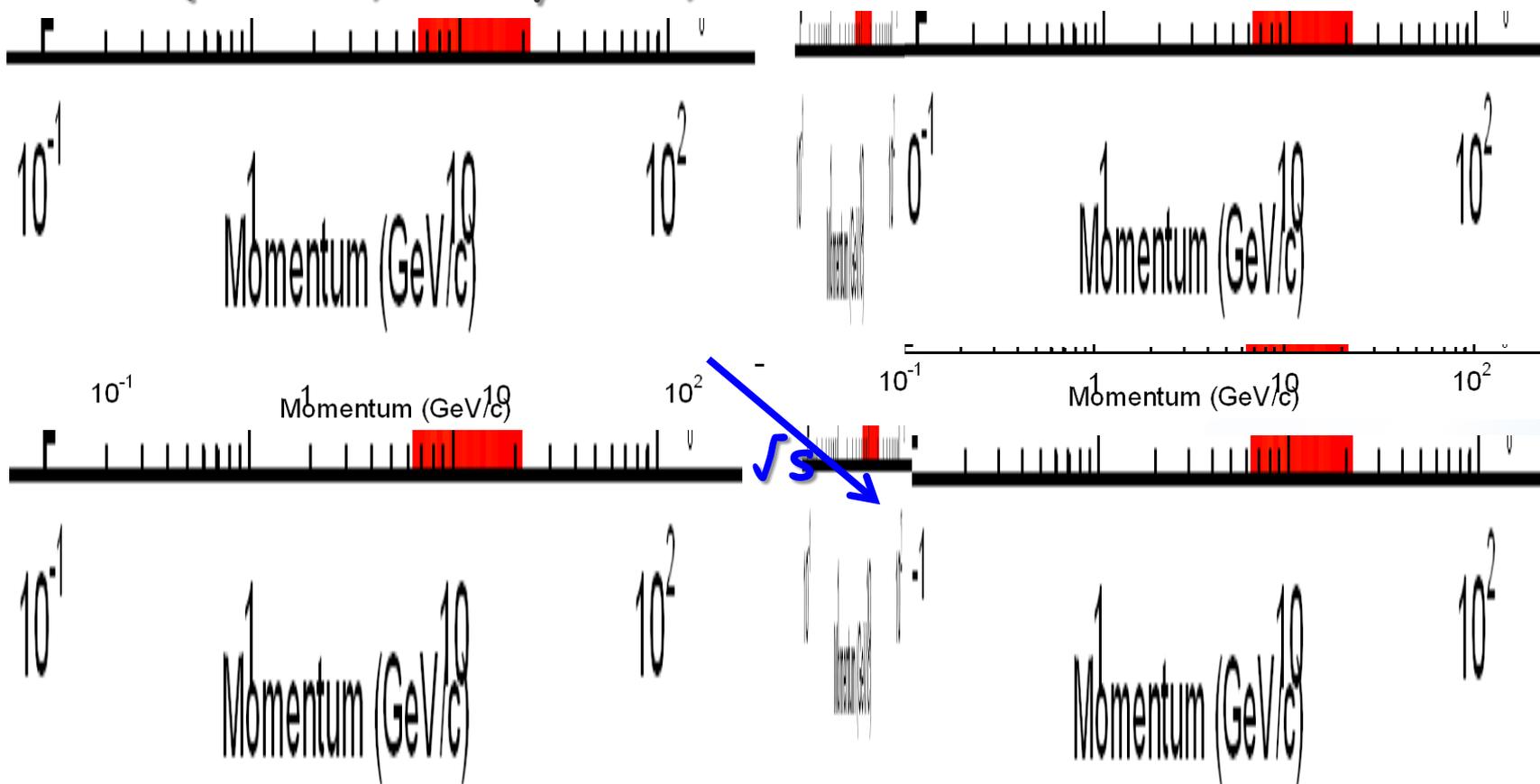
5 GeV on 50 GeV

higher \sqrt{s} :

**scattered lepton has
small scattering angle
→ negative rapidities**

PION KINEMATICS

Cuts: $Q^2 > 1 \text{ GeV}$, $0.01 < y < 0.95$, $z > 0.1$

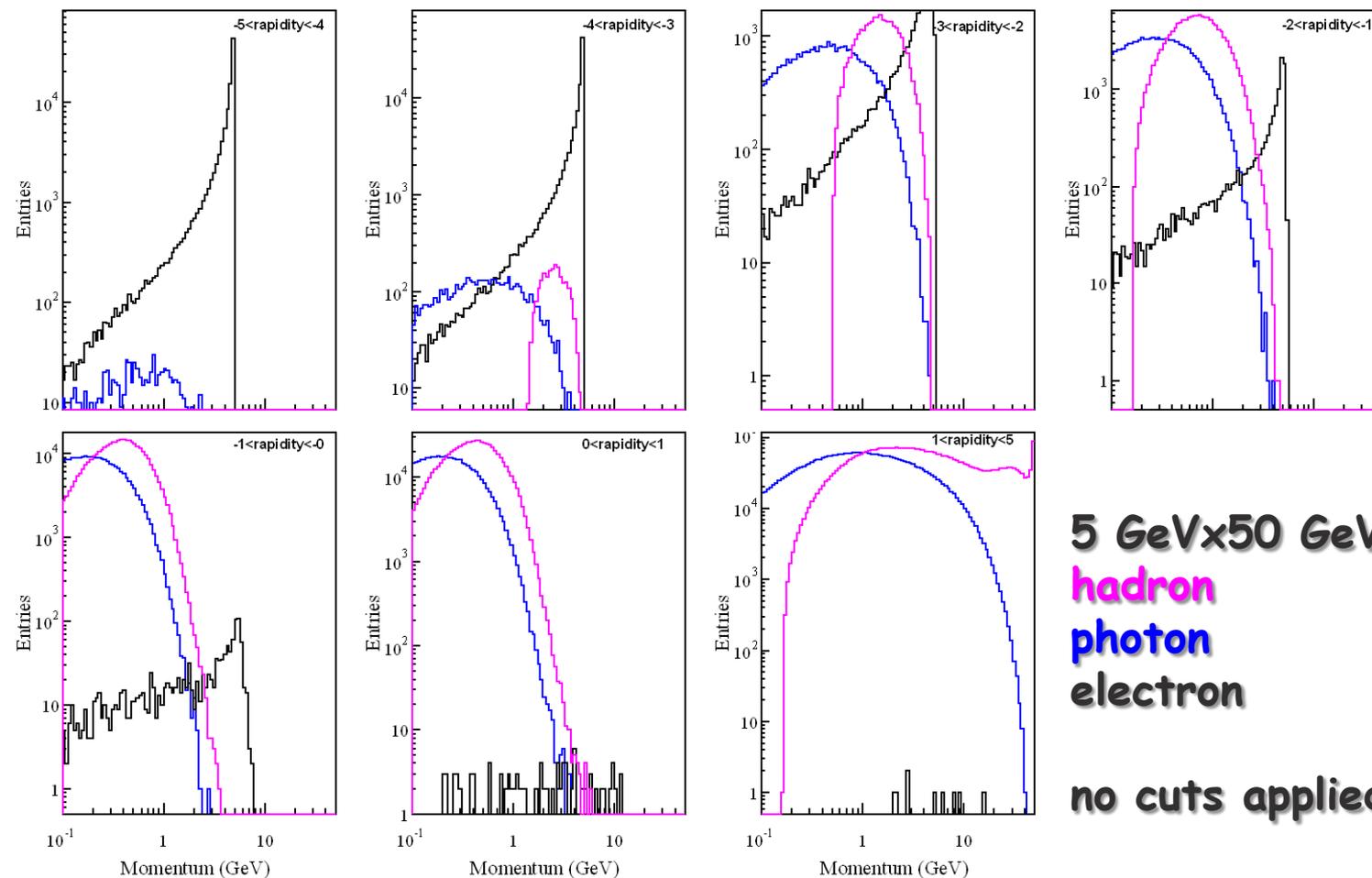


Increasing Hadron Beam Energy: influences max. hadron energy at fixed η
 Increasing $30 \text{ GeV} < \sqrt{s} < 170 \text{ GeV}$

→ hadrons are boosted from forward rapidities to negative rapidities

→ no difference between π^\pm , K^\pm , p^\pm

HADRON, LEPTON, PHOTON SEPARATION



5 GeVx50 GeV

hadron

photon

electron

no cuts applied

hadron/photon suppression factor

needed for $p_e > 1 \text{ GeV}$:

$-3 < \eta < -2$: ~ 10

$-2 < \eta < -1$: > 100

$-1 < \eta < 0$: ~ 1000

p_{max} hadron for PID:

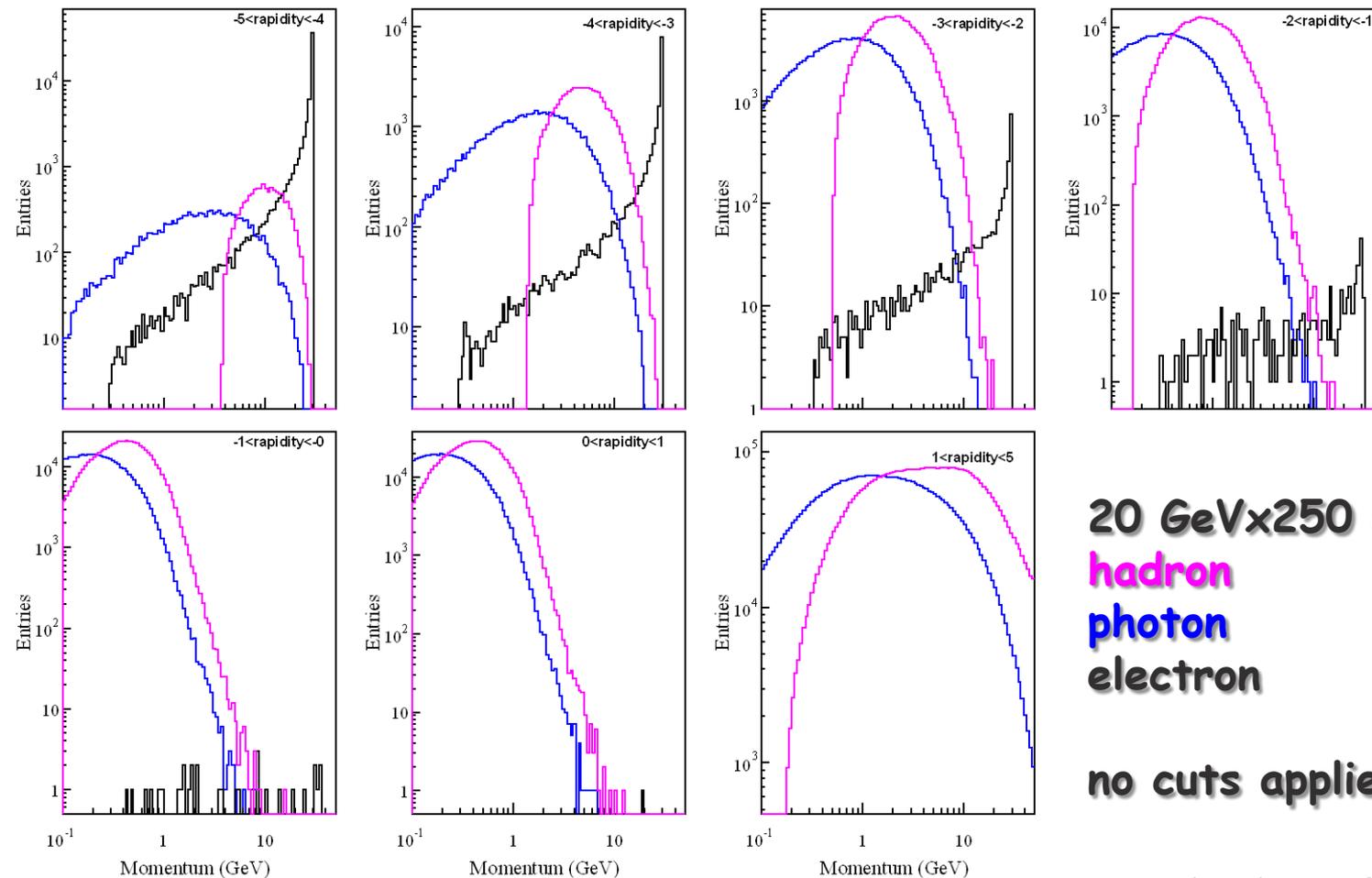
$-5 < \eta < -1$: $< 10 \text{ GeV}$

$-1 < \eta < 1$: $< 5 \text{ GeV}$

$1 < \eta < 5$: $< 50 \text{ GeV}$



LEPTON IDENTIFICATION



20 GeVx250 GeV

hadron

photon

electron

no cuts applied

hadron/photon suppression factor

needed for $p_e > 1\text{ GeV}$:

$-4 < \eta < -3$: > 100

$-3 < \eta < -2$: ~ 1000

$-2 < \eta < -1$: $> 10^4$

p_{max} hadron for PID:

$-5 < \eta < -1$: $< 30\text{ GeV}$

$-1 < \eta < 1$: $< 10\text{ GeV}$

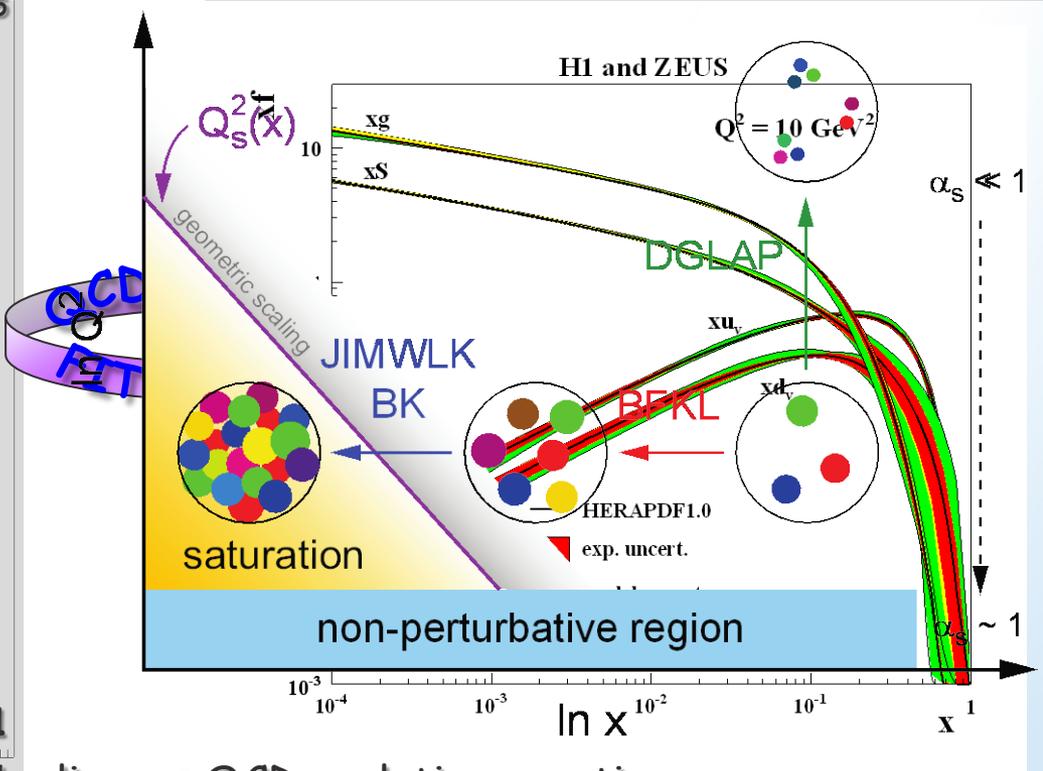
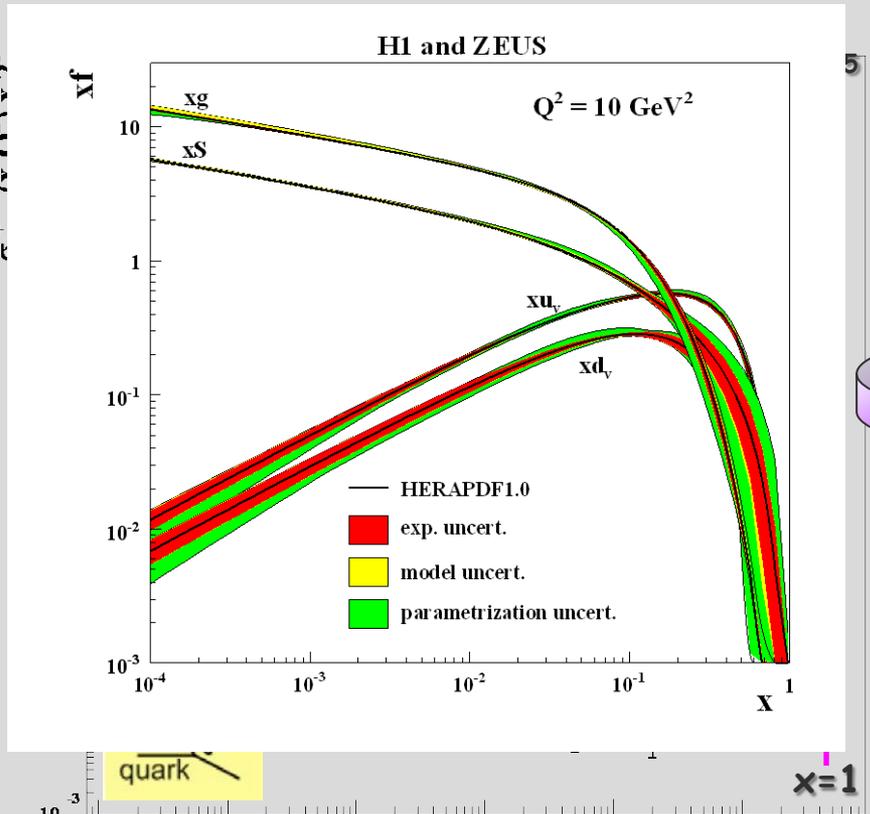
$1 < \eta < 5$: $< 100\text{ GeV}$



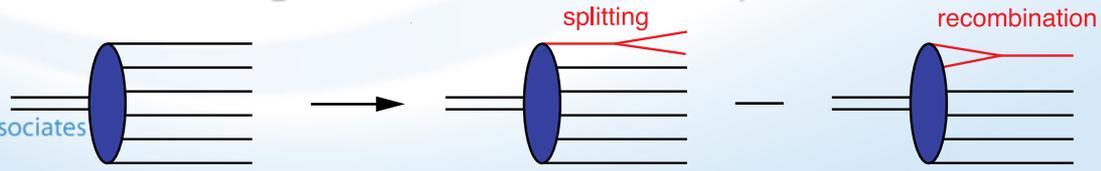
WHAT DO WE KNOW

$d^2\sigma_{\text{ep}}^{NC} \sim 2\pi\alpha^2 Y \frac{v^2}{Y}$
Gluon density dominates at $x < 0.1$

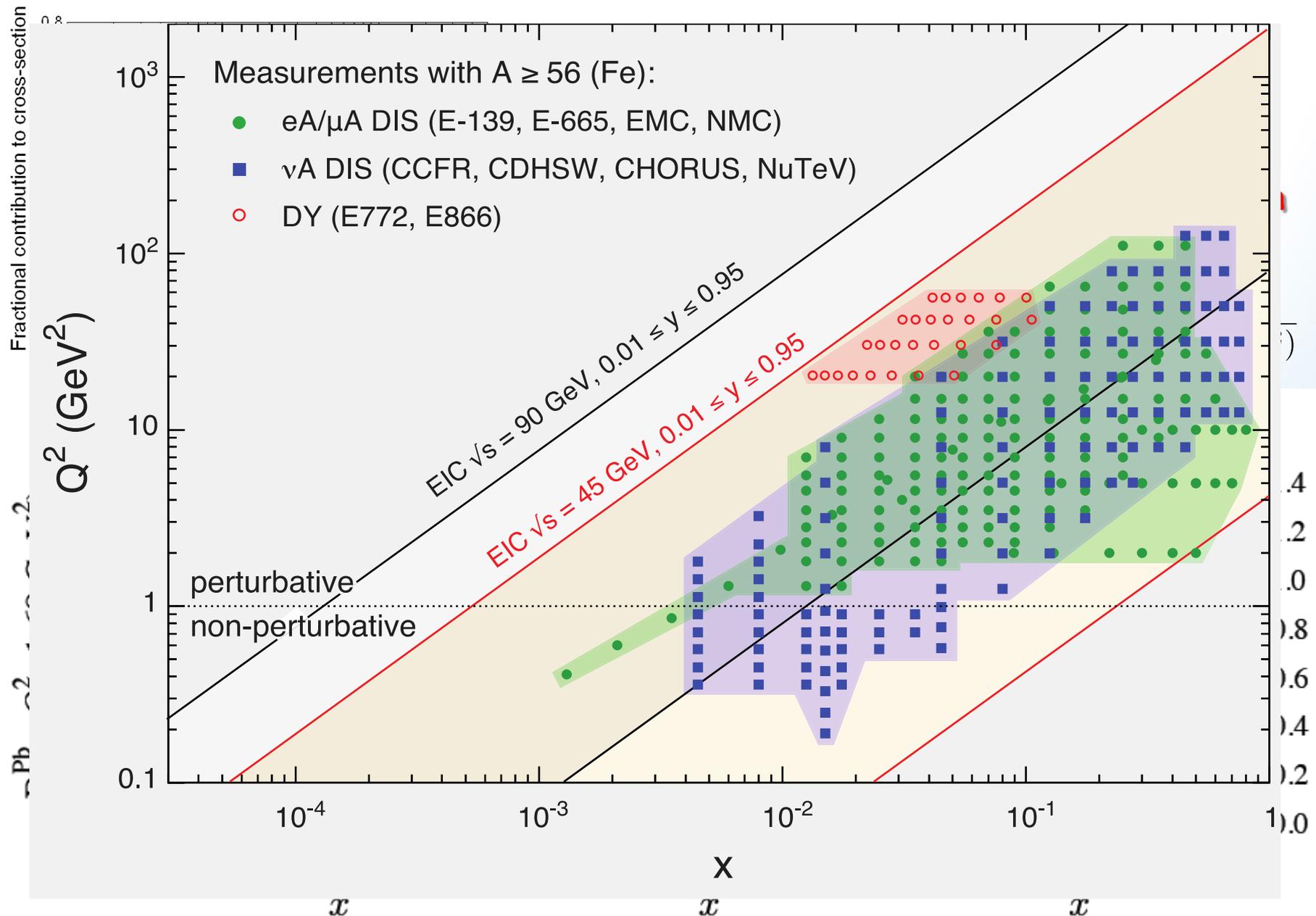
Gluon density dominates at $x < 0.1$



- Rapid rise in gluons described naturally by linear pQCD evolution equations
- This rise cannot increase forever - limits on the cross-section
- non-linear pQCD evolution equations provide a natural way to tame this growth and lead to a saturation of gluons, characterised by the saturation scale $Q_s^2(x)$



WHAT DO WE KNOW ABOUT xG IN NUCLEI?

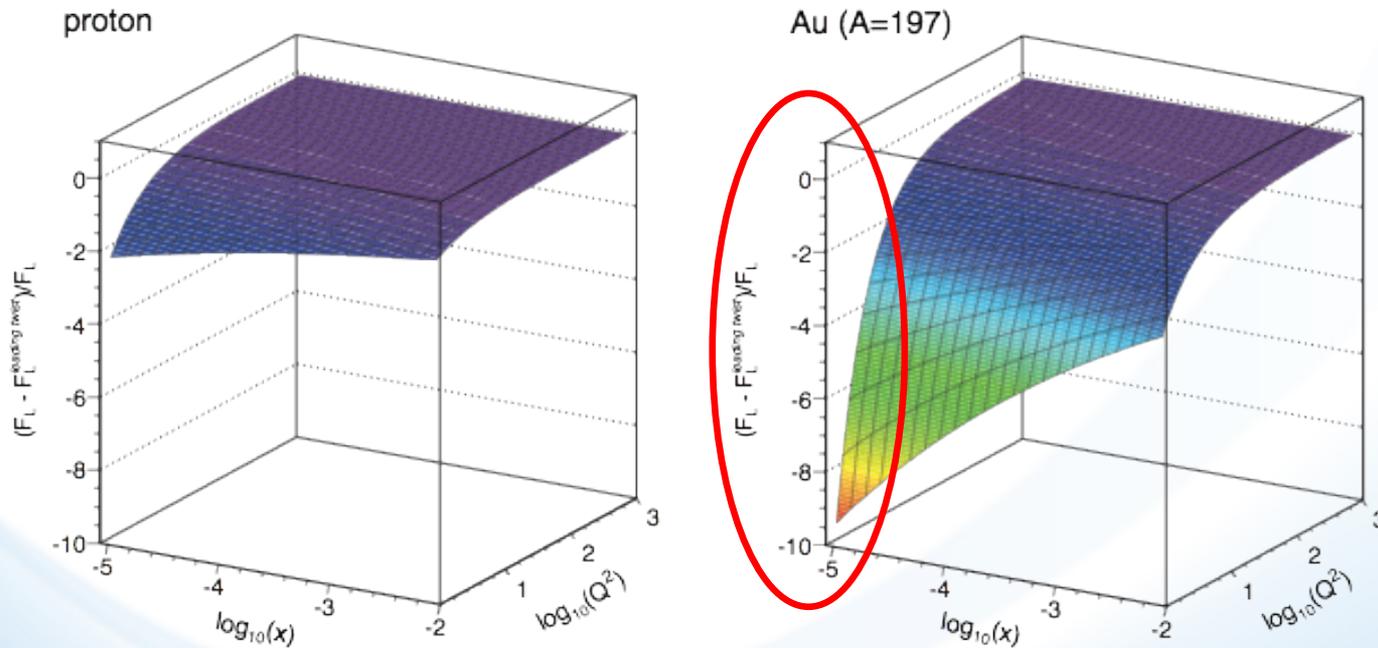


INCLUSIVE DIS IN eA : NUCLEAR PDFS

$$\frac{d^2\sigma^{eA \rightarrow eX}}{dx dQ^2} = \frac{4\pi\alpha^2}{xQ^4} \left[\left(1 - y - \frac{y^2}{2}\right) F_2(x, Q^2) - \frac{y^2}{2} F_L(x, Q^2) \right]$$

quark+anti-quark gluon

➤ Expect strong non-linear effects in F_L

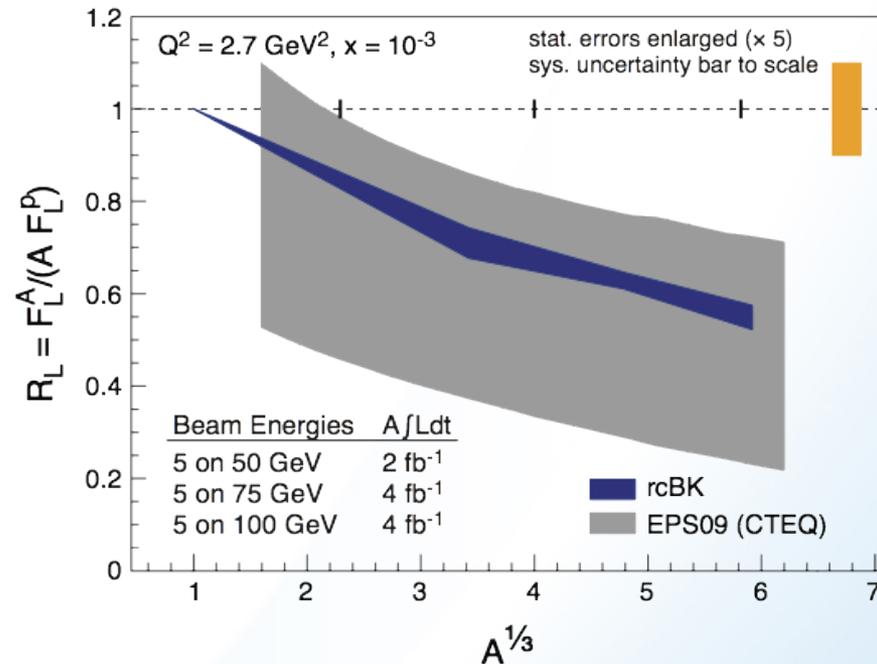
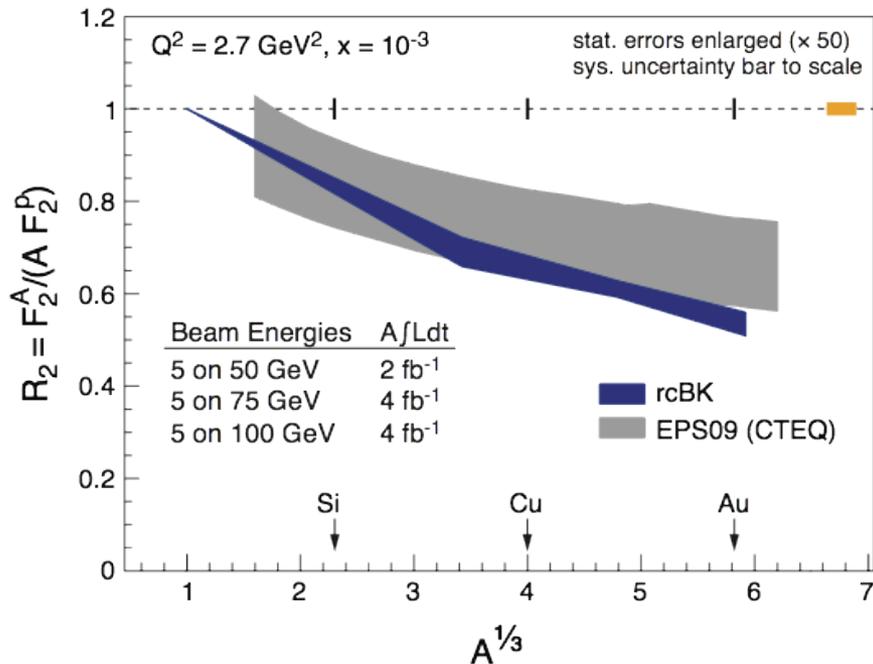


Relative contributions of higher twist effects to F_L amplified in eA

Dipole model (J. Bartels et al.)

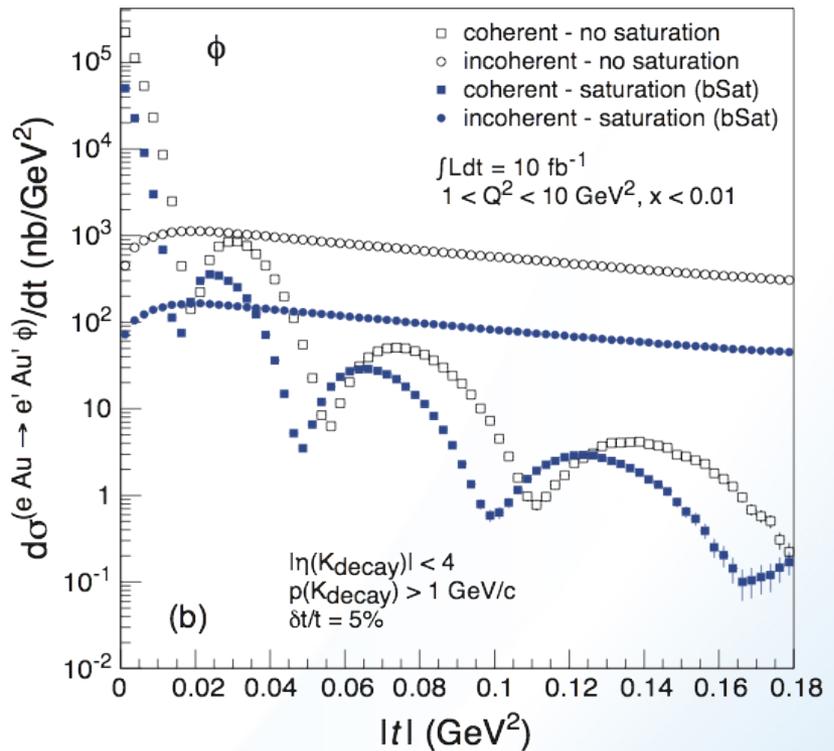
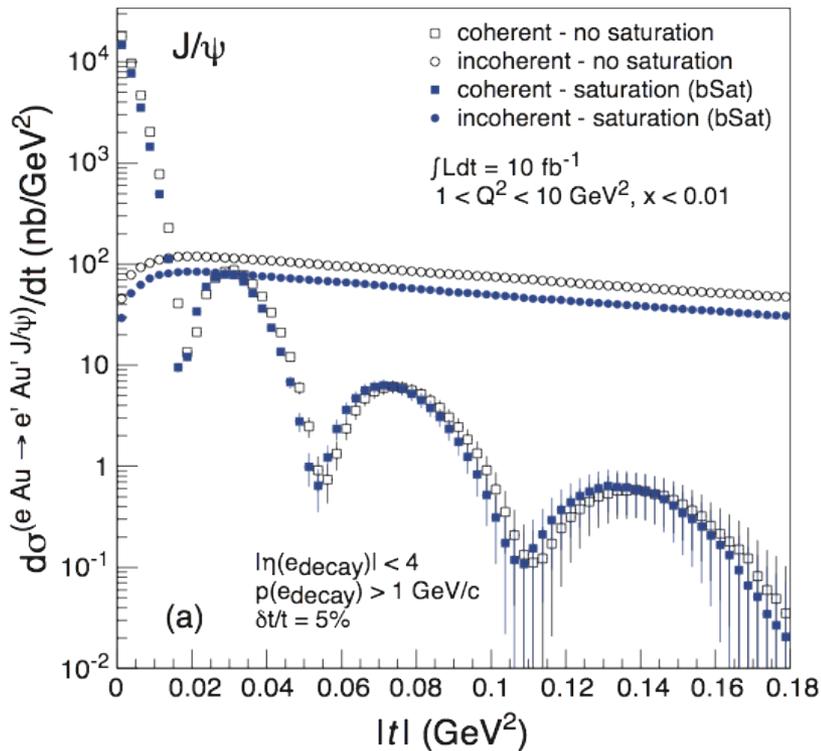
Brookhaven Science Associates

INCLUSIVE DIS IN EA: NUCLEAR PDFS



- measurement of F_L requires running at different \square s
- F_2, F_L : negligible stat. error, **systematics dominated**
- A dependence helps to **discriminate between linear and non-linear (saturation) models**
- Precision nPDF: Huge impact on pA, AA programs

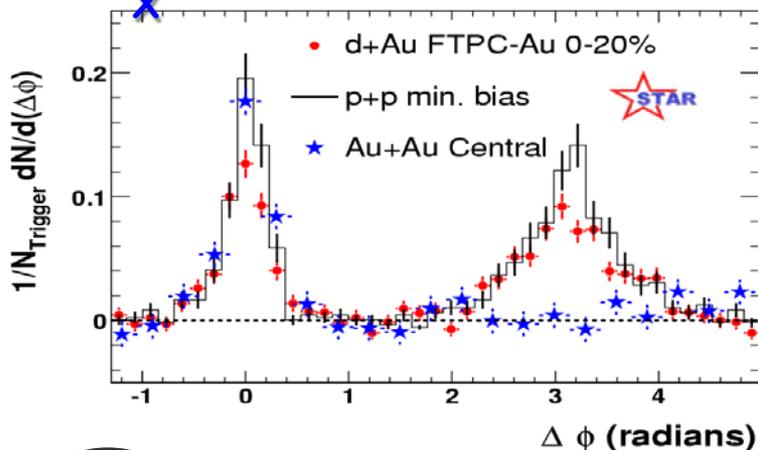
SPATIAL GLUON DISTRIBUTION THROUGH DIFFRACTION



- Goal: going after the **source distribution** of gluons **through Fourier transform of $d\sigma/dt$**
- Find: Typical diffractive pattern for coherent (non-breakup) part
- As expected: J/ψ less sensitive to saturation effects than larger Φ-meson

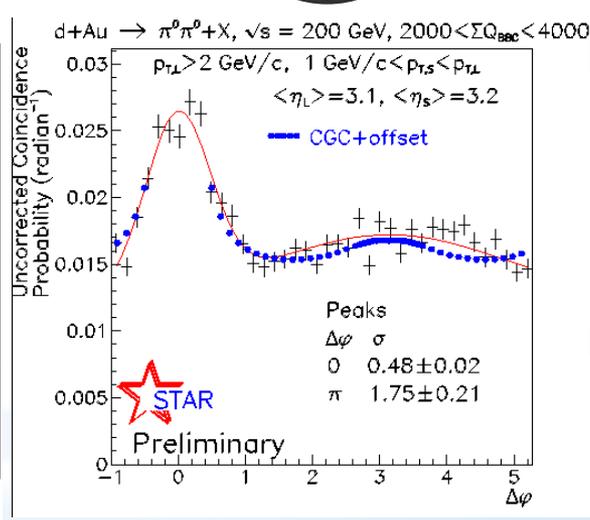
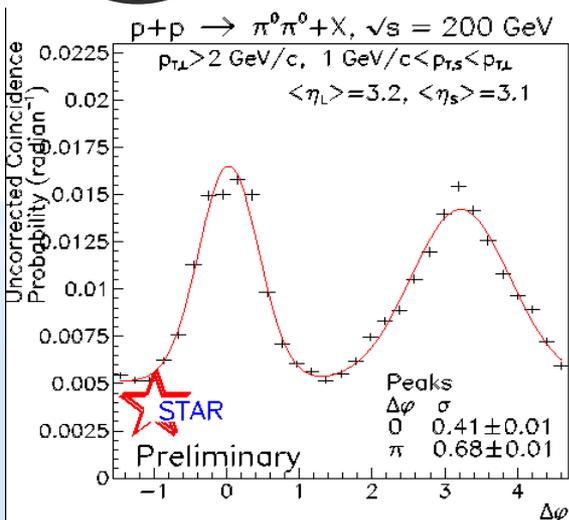
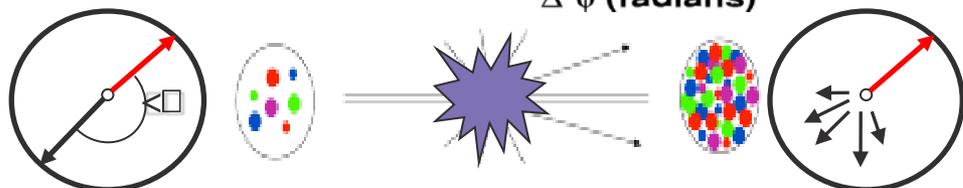
DI-HADRON CORRELATIONS IN dA

comparisons between $d+Au \rightarrow h_1 h_2 X$
(or $p+Au \rightarrow h_1 h_2 X$) and $p+p \rightarrow h_1 h_2$



- At $y=0$, suppression of away-side jet is observed in A+A collisions
- No suppression in p+p or d+A
 $\rightarrow x \sim 10^{-2}$

$$x_A = \frac{k_1 e^{-y_1} + k_2 e^{-y_2}}{\sqrt{s}} \ll 1$$

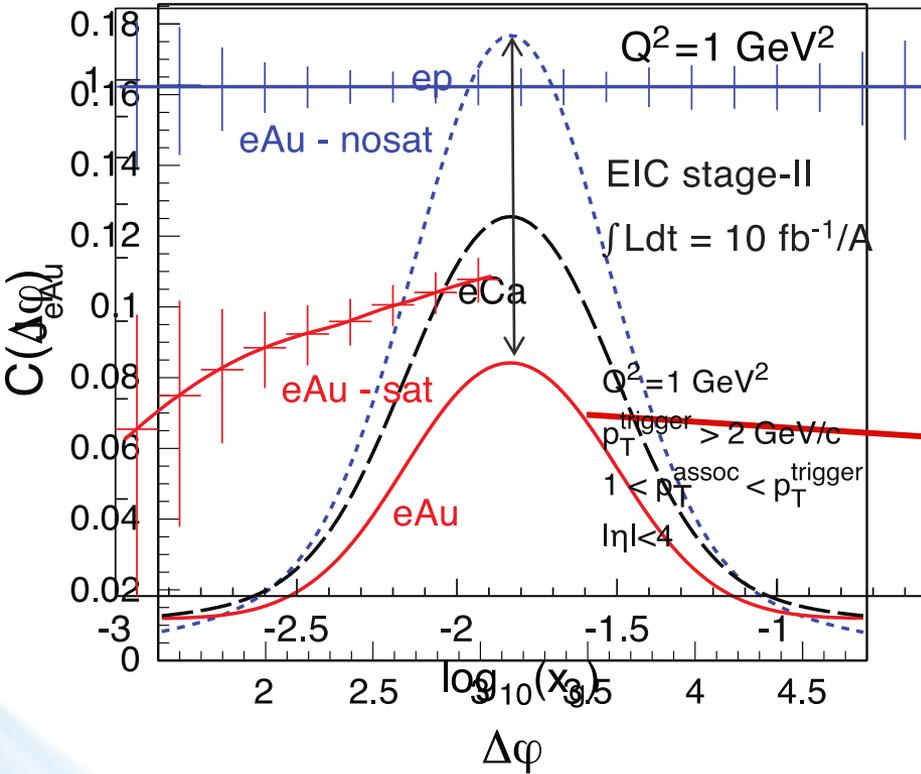


- However, at forward rapidities ($y \sim 3.1$), an away-side suppression is observed in dAu
- Away-side peak also much wider in d+Au compared to pp
 $\rightarrow x \sim 10^{-3}$

eA DIHADRON CORRELATIONS STUDIES

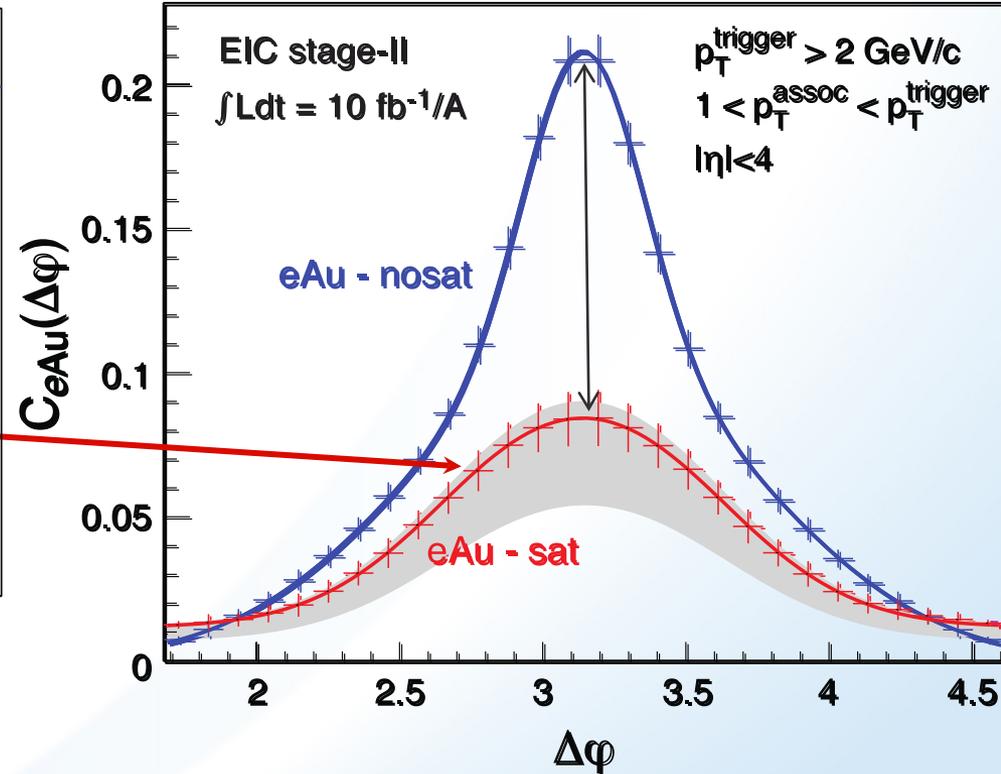
Theory: Saturation

Dominguez, Xiao, Yuan, Lee, Zheng '11/12



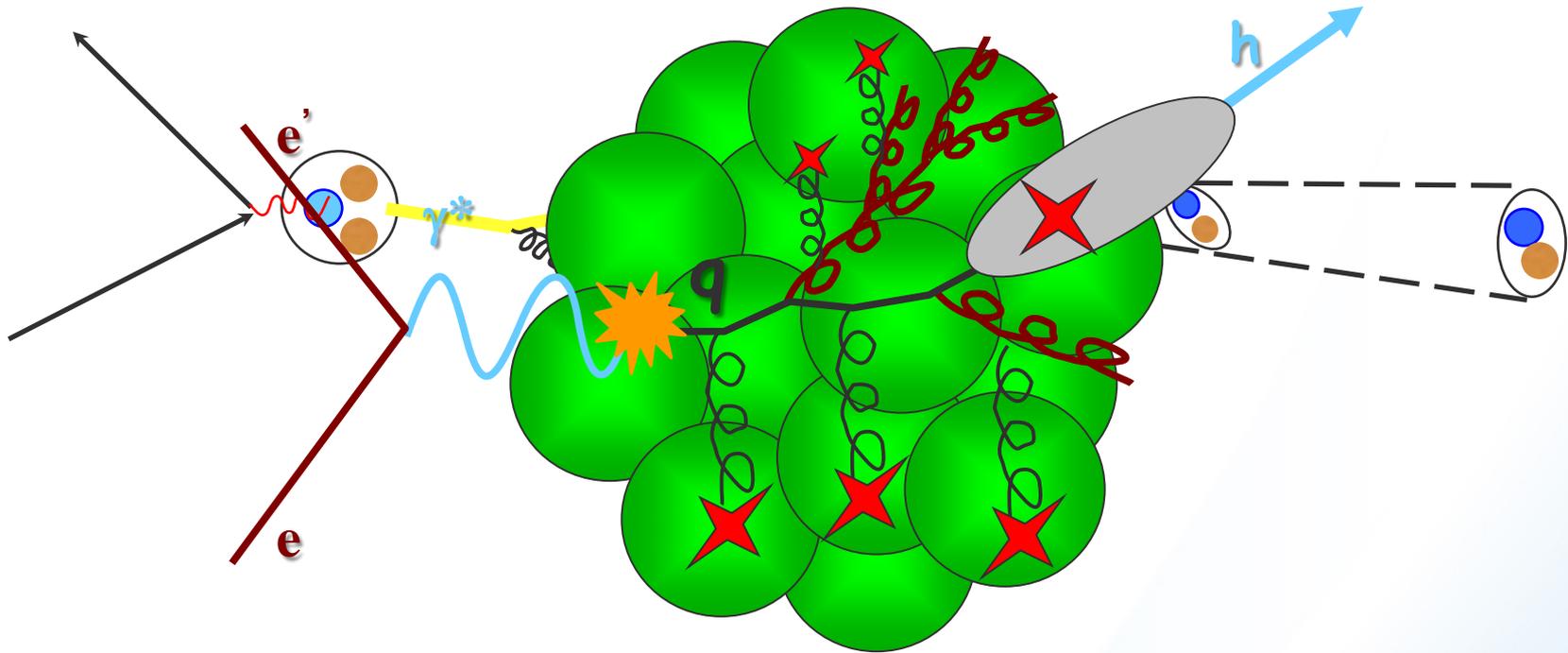
Exp: Saturation versus "conventional" scenario

J.H. Lee, L. Zheng (CCNU) '11/12



- ❑ eA-MC: Pythia6.4 + nPDF (EPS09) + nuclear geometry from DPMJetIII without PS
- ❑ Here for 10 fb⁻¹/A (~ 20 weeks), std. experimental cuts
- ❑ **Clear signal, pronounced differences between sat and no-sat**

DEEP INELASTIC SCATTERING - VACUUM



What happens if we add a nuclear medium

Observables:

Broadening:

Δp_t^2 linked directly with saturation scale

$$\Delta p_t^2 = \langle p_t^2 \rangle_A - \langle p_t^2 \rangle_p$$

Attenuation:

ratio of hadron production in A to d
modifications of nPDF cancel out

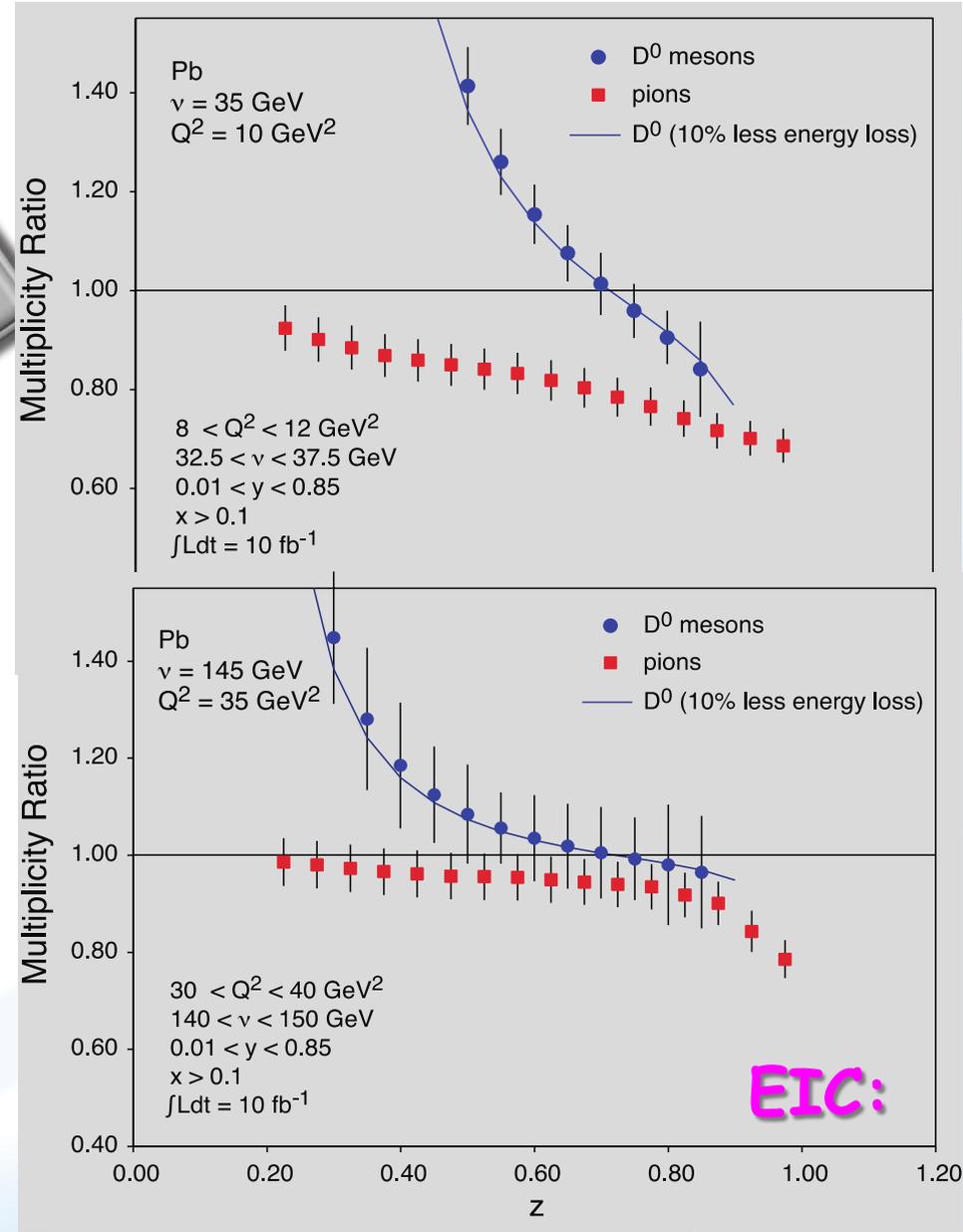
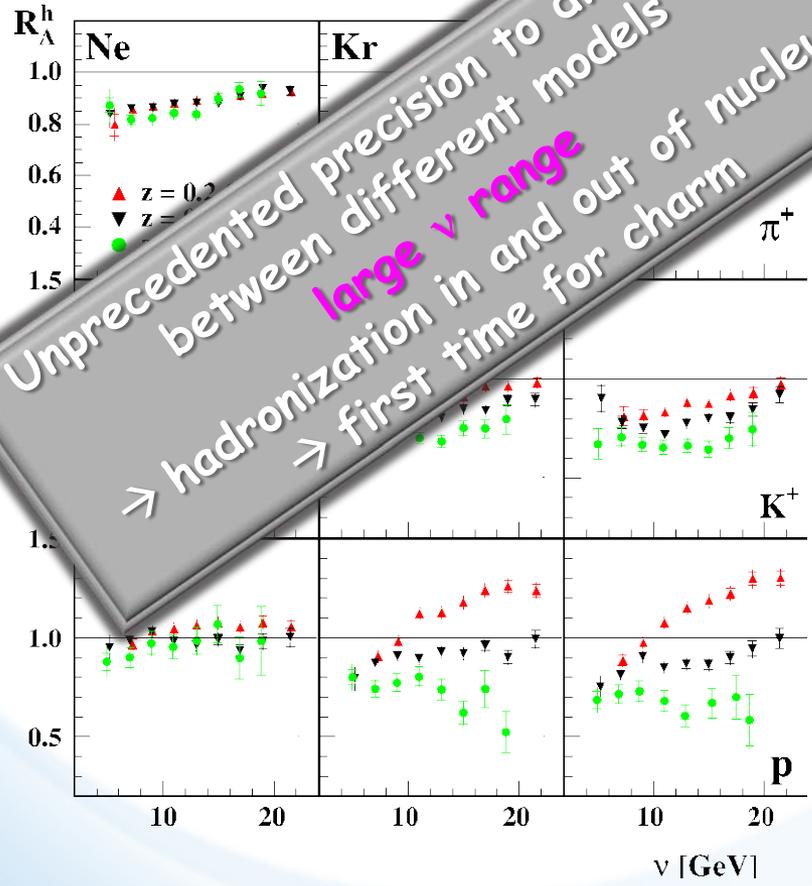
$$R_A^h(Q^2, x, z, p_t, \Theta)$$

WHAT DO WE KNOW AND WHAT CAN EIC DO

Hermes:

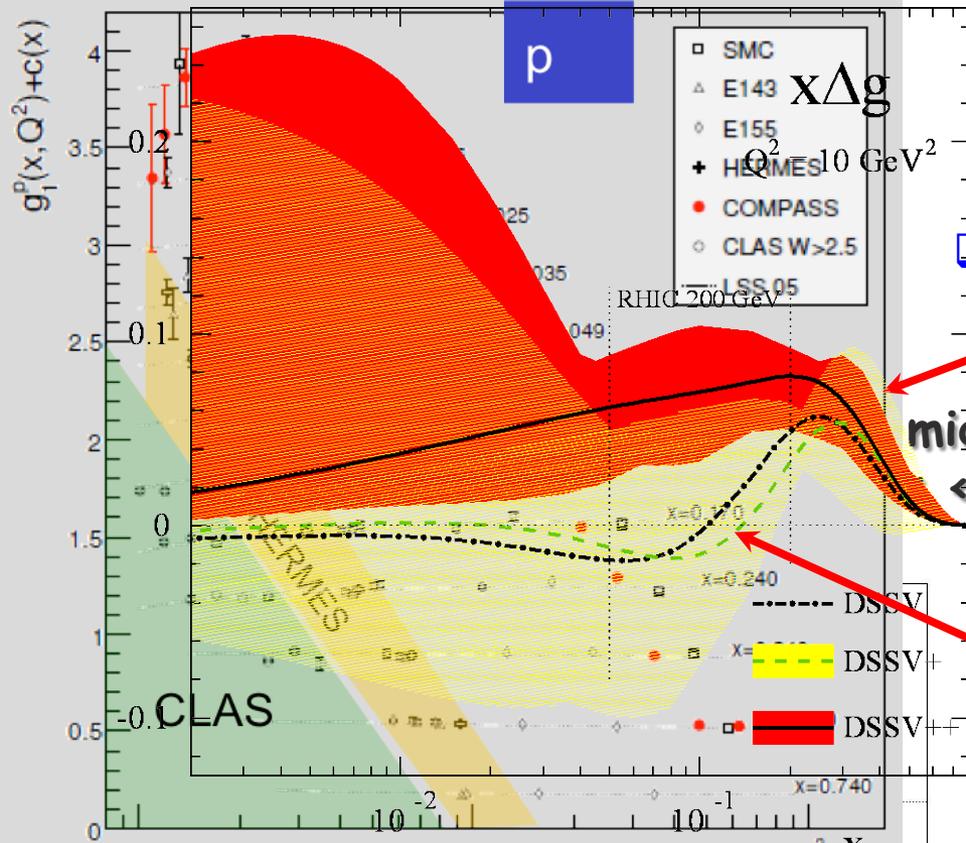
$E_e = 27 \text{ GeV} \rightarrow \sqrt{s} = 7.2 \text{ GeV}$

$E_h = 2-15 \text{ GeV}$



EIC:

WHAT DO WE KNOW NOW ON $\Delta g(x)$



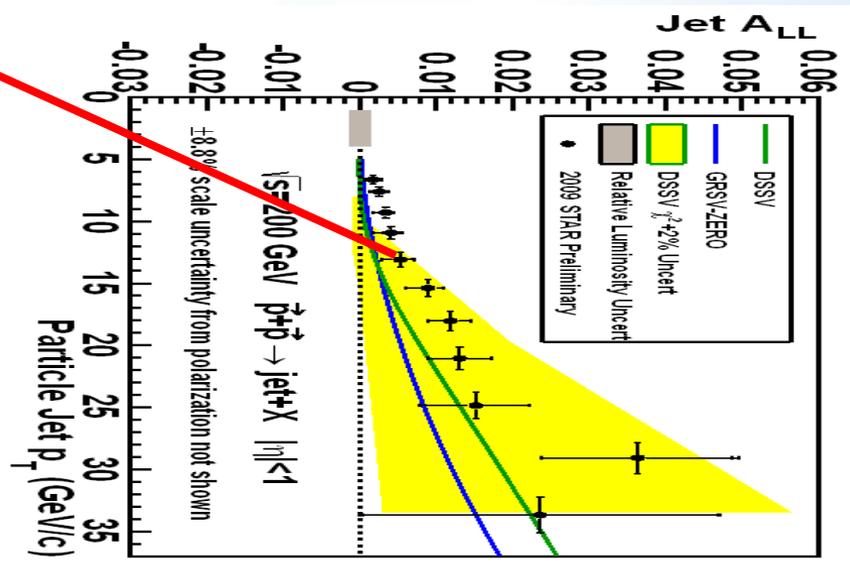
$$\frac{dg_1}{d \log(Q^2)} \sim -\Delta g(x, Q^2)$$

Scaling violations of g_1 (Q^2 -dependence) give indirect access to the gluon distribution via DGLAP midrapidity; direct access to gluons (gg, qg) $x < 0.2$

Integral in RHIC x -range: Q^2 [GeV²]

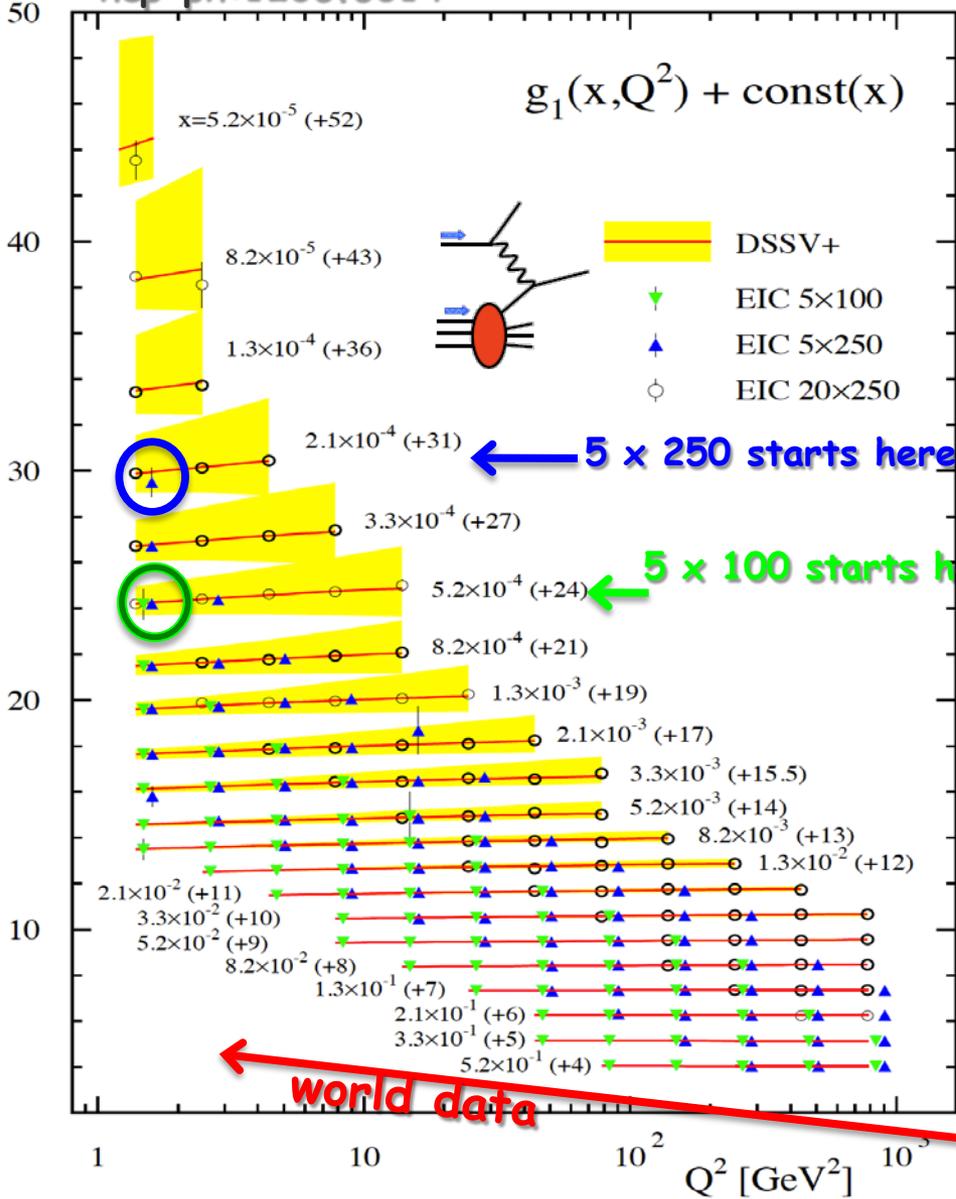
$\int_{0.05}^{0.2} dx \Delta g \sim 0.1 \pm 0.06 \pm 0.07$ @ 10 GeV²

Contribution to proton spin to date:
 Gluon: 20% Quarks: 30%



g_1^p THE WAY TO FIND THE SPIN

hep-ph:1206.6014



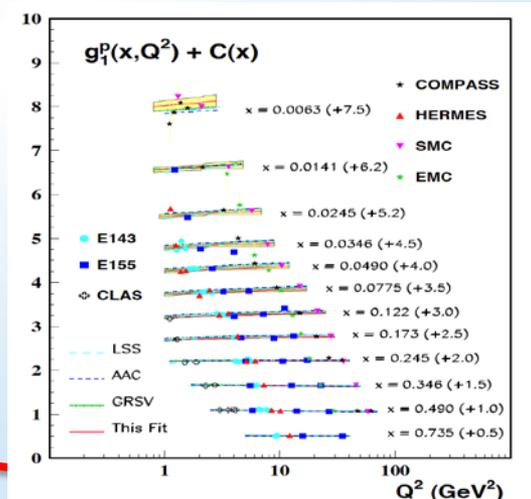
cross section: $\frac{d^2\sigma}{d\Omega dE'} \sim L_{\mu\nu} W^{\mu\nu}$

$$W^{\mu\nu} = -g^{\mu\nu} F_1 - \frac{p^\mu p^\nu}{v} F_2 + \frac{i}{v} \epsilon^{\mu\nu\lambda\sigma} q^\lambda s^\sigma g_1 + \frac{i}{v^2} \epsilon^{\mu\nu\lambda\sigma} q^\lambda (p \cdot q s^\sigma - s \cdot q p^\sigma) g_2$$

pQCD scaling violations

$$\frac{dg_1}{d \log(Q^2)} \sim -\Delta g(x, Q^2)$$

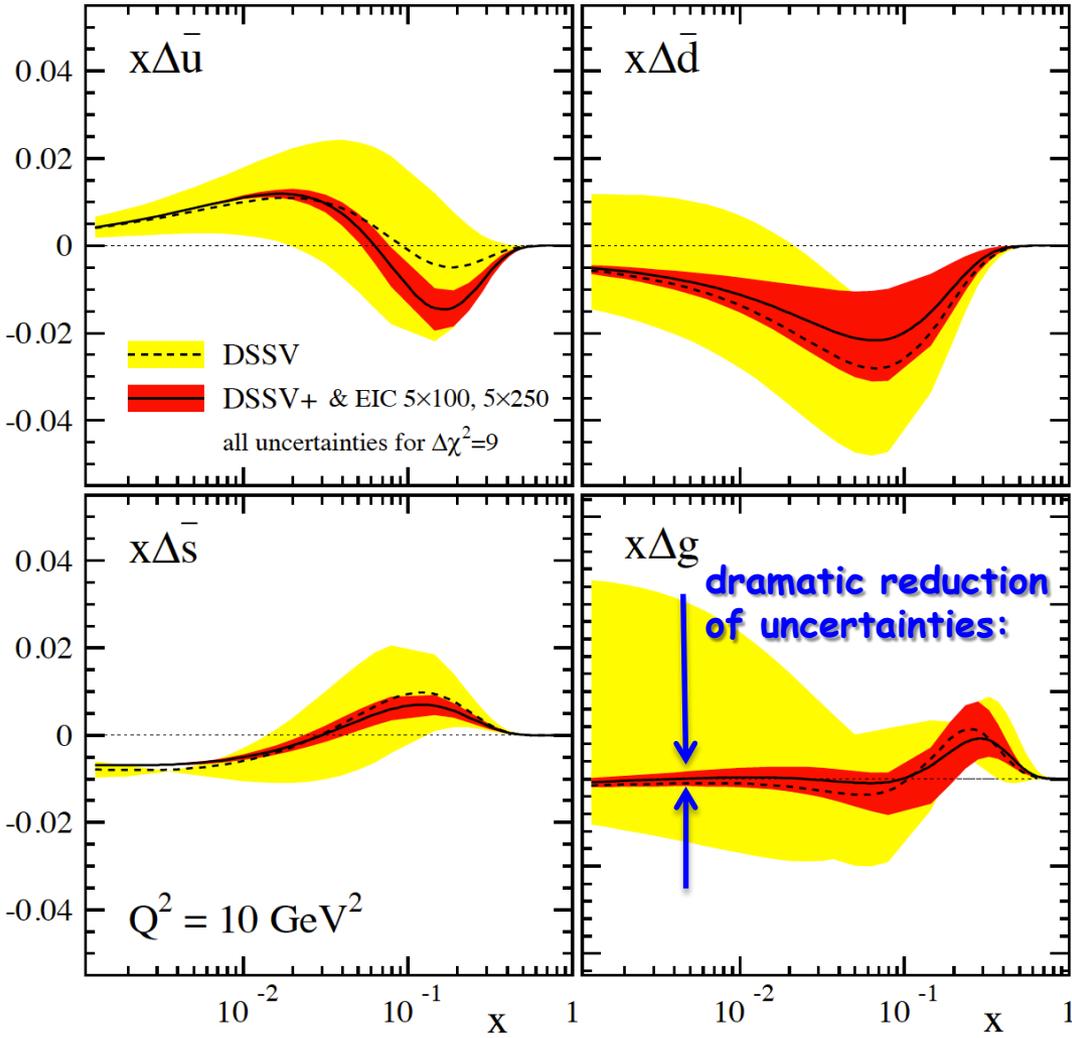
$$\Delta \Sigma(Q^2) = \int_0^1 g_1(x, Q^2) dx = \int_0^1 \Delta q_f(x, Q^2) dx$$



IMPACT OF eRHIC DATA ON HELICITY PDFs

DIS scaling violations mainly determine Δg at small x

in addition, SIDIS data provide detailed **flavor separation of quark sea**



yet, small x behavior completely unconstrained
 \rightarrow determines x -integral, which enters proton spin sum

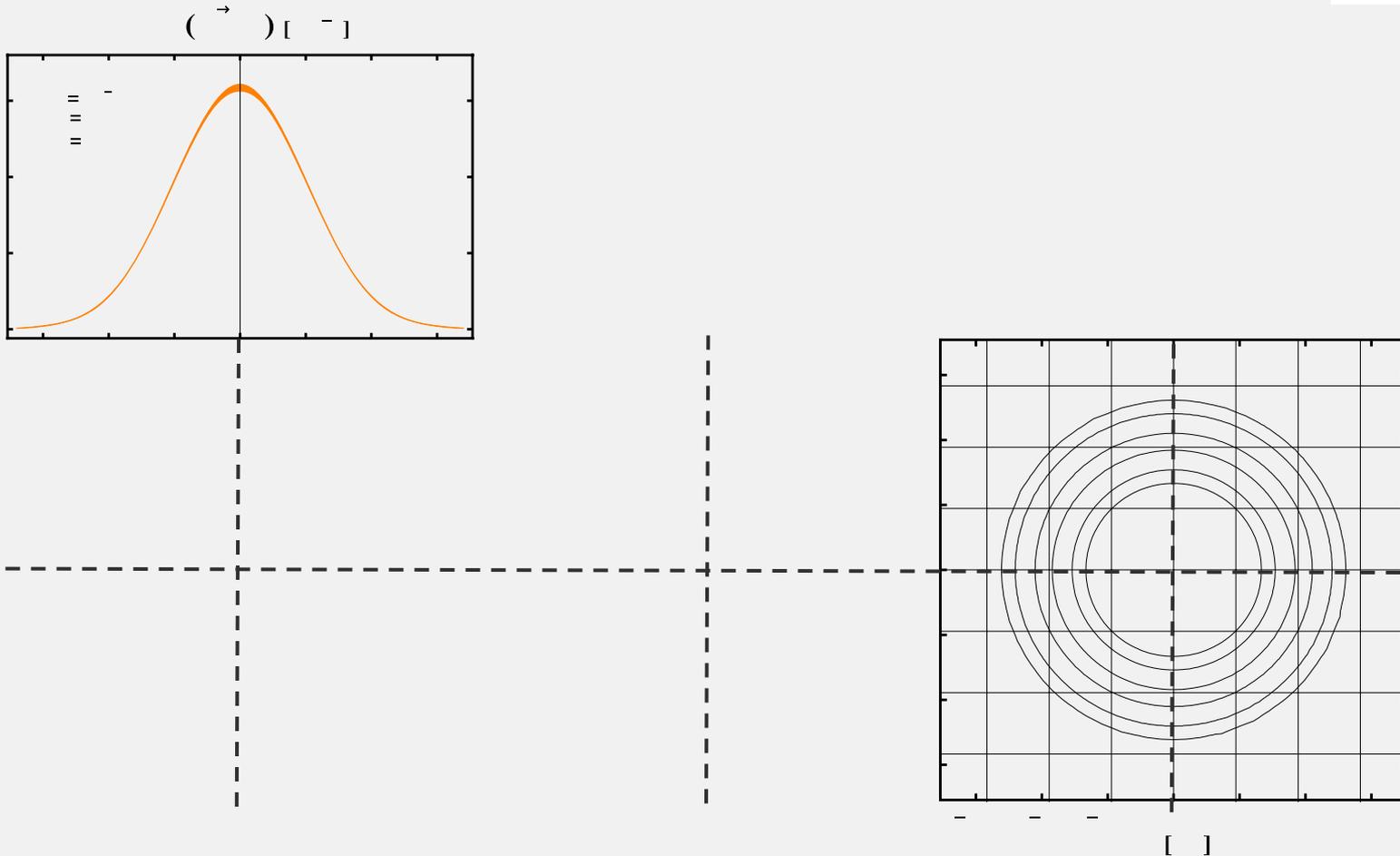
- includes only "stage-1 data"
- can be pushed to $x=10^{-4}$ with $20 \times 250 \text{ GeV}$ data

"issues":

- (SI)DIS @ eRHIC limited by systematic uncertainties need to control rel. lumi, polarimetry, detector performance, ... very well

WHAT WILL WE LEARN ABOUT 2D+1 STRUCTURE OF THE PROTON

GPD H and E as function of t , x and Q^2



GPD H and E 2d+1 structure for sea-quarks and gluons

→ Excellent reconstruction of H^A, H^B
and good reconstruction of H^C (from $d\sigma/dt$)

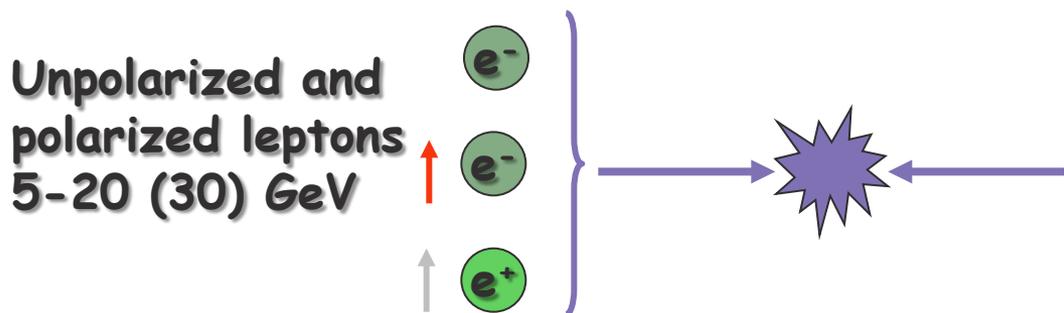
AN ELECTRON ION COLLIDER IN THE US

Requirements:

- ❑ High Luminosity $> 10^{33} \text{ cm}^{-2}\text{s}^{-1}$
- ❑ Flexible center of mass energies
- ❑ Electrons and protons/light nuclei polarised
- ❑ Wide range of nuclear beams
- ❑ a wide acceptance detector with good PID (e/h and π , K, p)
- ❑ wide acceptance for protons from elastic reactions and neutrons from nuclear breakup

Electron accelerator

to be build



70% e^- beam polarization goal
polarized positrons?

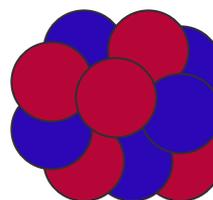


RHIC

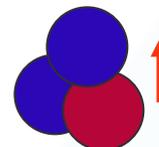
Existing = \$2B



Polarized protons
50-250 GeV

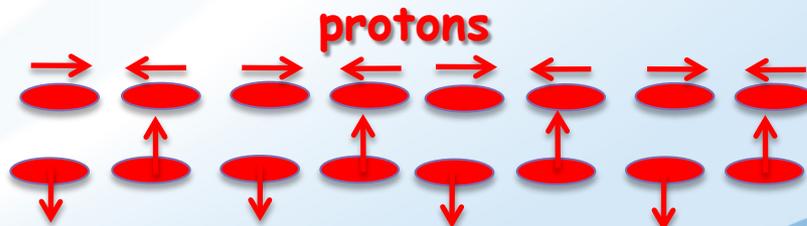
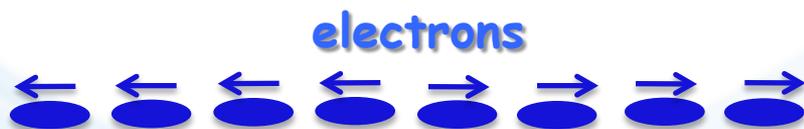


Light ions (d, Si, Cu)
Heavy ions (Au, U)
50-100 GeV/u



Polarized light ions
 He^3 166 GeV/u

Center mass energy range: $\sqrt{s}=30-200$ GeV; $L \sim 100-1000 \times$ Hera
longitudinal and transverse polarization for p/ He^3 possible



eRHIC: design luminosity

	e	p	${}^2\text{He}^3$	${}^{79}\text{Au}^{197}$	${}^{92}\text{U}^{238}$
Energy, GeV	20	250	167	100	100
CM energy, GeV		100	82	63	63
Number of bunches/distance between bunches	107 nsec	111	111	111	111
Bunch intensity (nucleons) , 10^{11}	0.36	4	6	6	6
Bunch charge, nC	5.8	64	60	39	40
Beam current, mA	50	556	556	335	338
Normalized emittance of hadrons , 95% , mm mrad		1.2	1.2	1.2	1.2
Normalized emittance of electrons, rms, mm mrad		16	24	40	40
Polarization, %	80	70	70	none	none
rms bunch length, cm	0.2	5	5	5	5
β^* , cm	5	5	5	5	5
Luminosity per nucleon, $\times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$		2.7	2.7	1.6	1.7

Hourglass the pinch effects are included. Space charge effects are compensated.

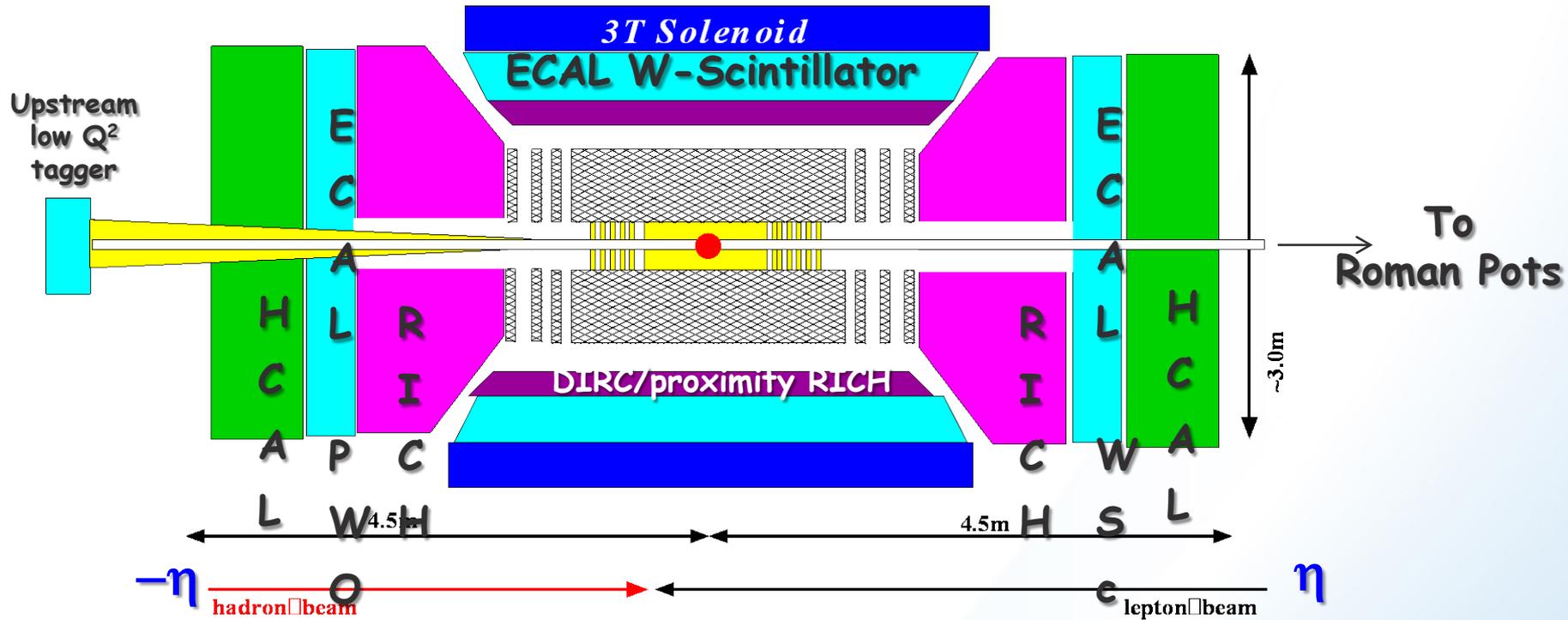
Energy of electrons can be selected at any desirable value at or below 30 GeV

The luminosity does not depend on the electron beam energy below or at 20 GeV

The luminosity falls as E_e^{-4} at energies above 20 GeV

The luminosity is proportional to the hadron beam energy: $L \sim E_h/E_{\text{top}}$

BNL: 1ST DETECTOR DESIGN CONCEPT



PID:

$-1 < \eta < 1$: DIRC or proximity focusing Aerogel-RICH

$1 < |\eta| < 3$: RICH

Lepton-ID:

$-3 < \eta < 3$: e/p

$1 < |\eta| < 3$: in addition Hcal response & γ suppression via tracking

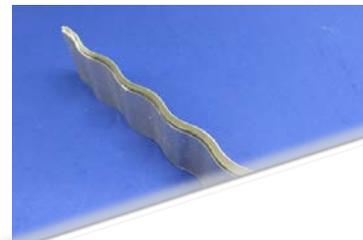
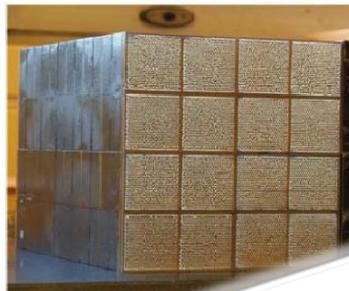
$|\eta| > 3$: ECal+Hcal response & γ suppression via tracking

$-5 < \eta < 5$: Tracking (TPC+GEM+MAPS)

VIBRANT DETECTOR R&D PROGRAM

Calorimetry

- W-Scintillator & W-Si
 - compact and high resolution
- Crystal calorimeters PbW & BGO



BNL, Indiana University, Penn State Univ., UCLA, USTC, TAMU

Pre-Shower

- W-Si
- LYSO pixel array with readout via X-Y WLS fibers



Univ. Tecnica Valparaiso

PID via Cerenkov

- DTD

Univ. of South Carolina, JLab, GSI

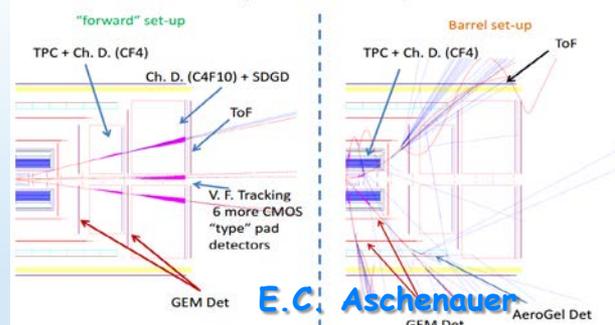
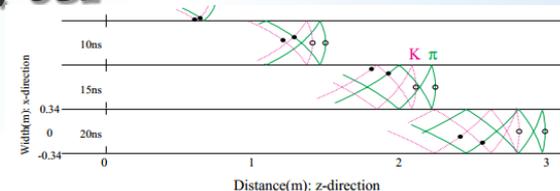
- Readout based TRD → eSTAR

Indiana Univ., USTC, VECC, ANL

Tracking

BNL, Florida Inst. Of Technology, Iowa State, LBNL, MIT, Stony Brook, Temple, Jlab, Virginia, Yale

- **μ -Vertex:** central and forward based on MAPS
- **Central:** TPC/HBD provides low mass, good momentum, dE/dx , eID
Fast Layer: μ -Megas or PIMMS
- **Forward:** Planar GEM detectors



Brookhaven Science Associates

More Info on EIC-Detector R&D:
https://wiki.bnl.gov/conferences/index.php/EIC_R%25D

MODELING THE DETECTOR IN GEANT

μ -vertex detector:

- ❑ 6 layers with [30..160] mm radius
- ❑ 0.37% X_0 in acceptance per layer simulated precisely;
- ❑ digitization: single discrete pixels one-to-one from MC points

Forward/b

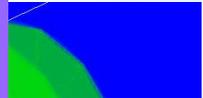
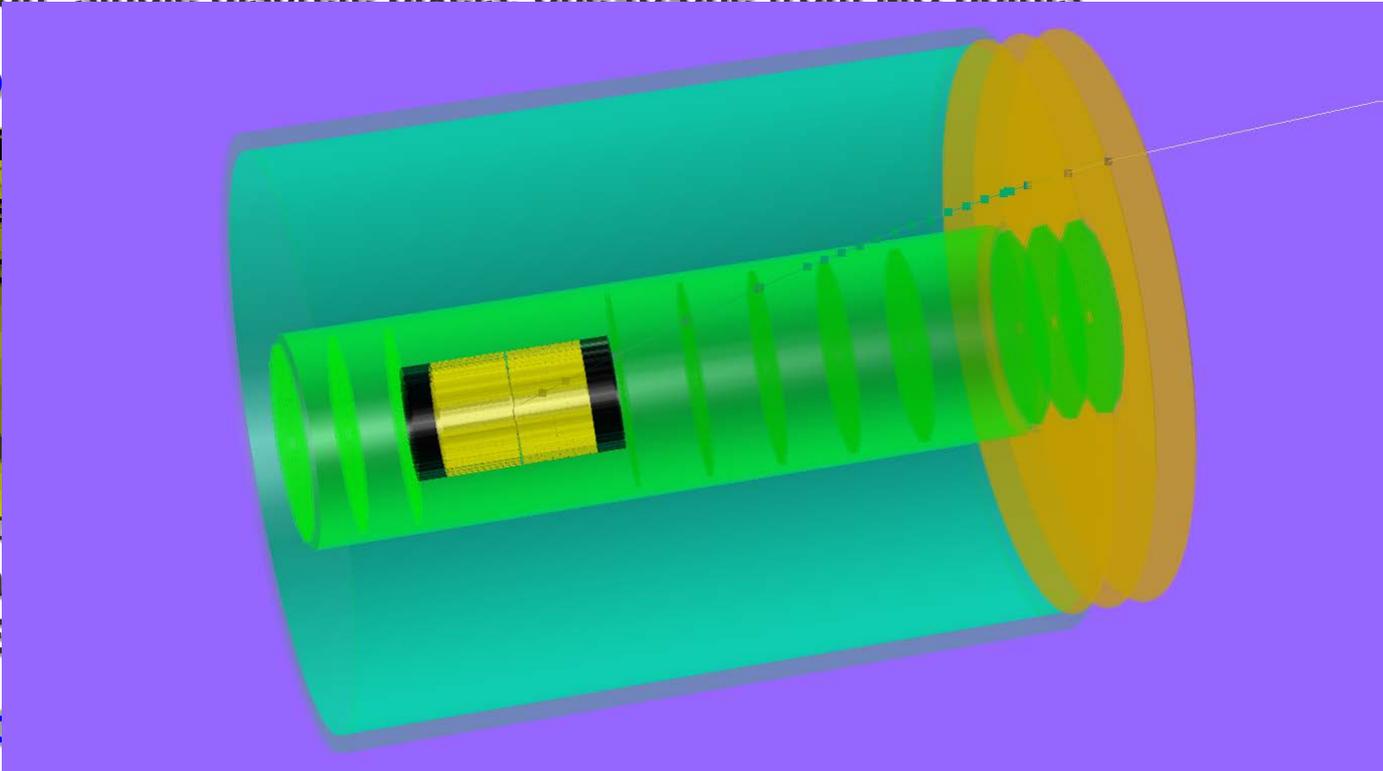
- ❑ 3+5+3 sil
- ❑ N sectors
- ❑ digitization

TPC

- ❑ ~2m lon
- ❑ 1.2% X_0
- ❑ digitiza
- mm GE

Forward t

- ❑ 3 disks behind the TPC endcap
- ❑ rather precise START FGT design implemented
- ❑ digitization: 100 μ m resolution in X&Y; gaussi



KINEMATICS OF BREAKUP NEUTRONS

Results from GEMINI++ for 50 GeV Au

theta distribution of neutrons at $E^* = 10$ MeV

histoTheta10
Entries 9143

theta distribution of neutrons at $E^* = 50$ MeV

histoTheta50

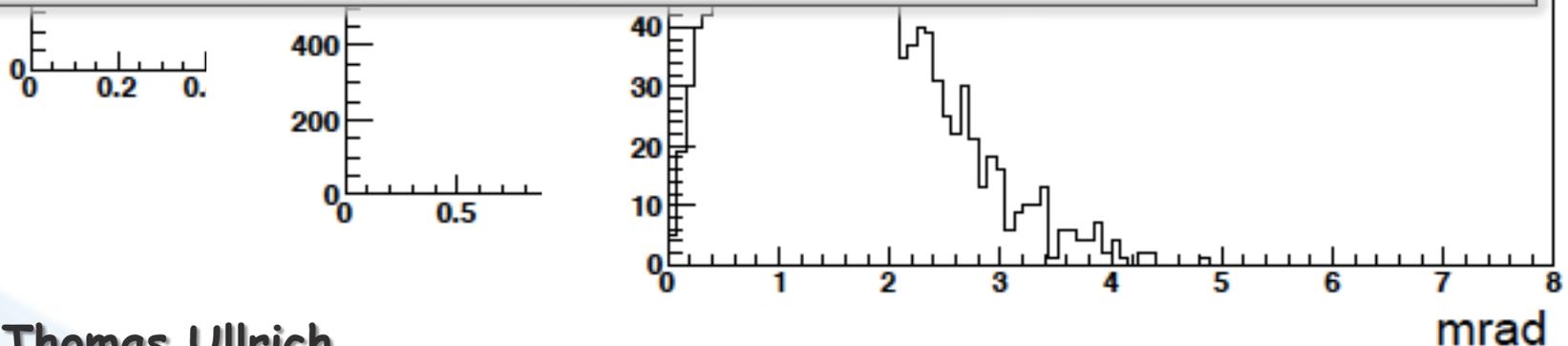
Results:

- With an aperture of ± 3 mrad we are in relative good shape
- enough "detection" power for $t > 0.025$ GeV²
 - below $t \sim 0.02$ GeV² we have to look into photon detection
 - Is it needed?

Question:

- For some physics rejection power for incoherent is needed $\sim 10^4$
- How efficient can the ZDCs be made?

eta500
2098
1.445
0.8048

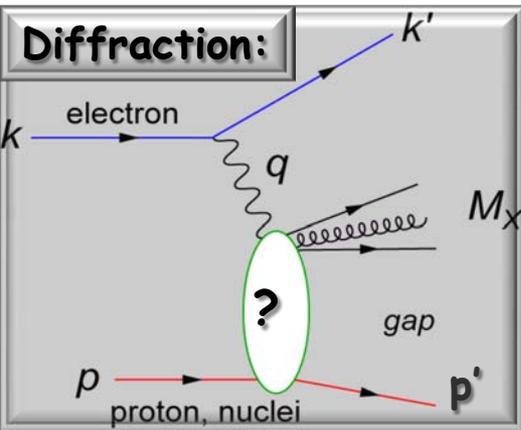


by Thomas Ullrich



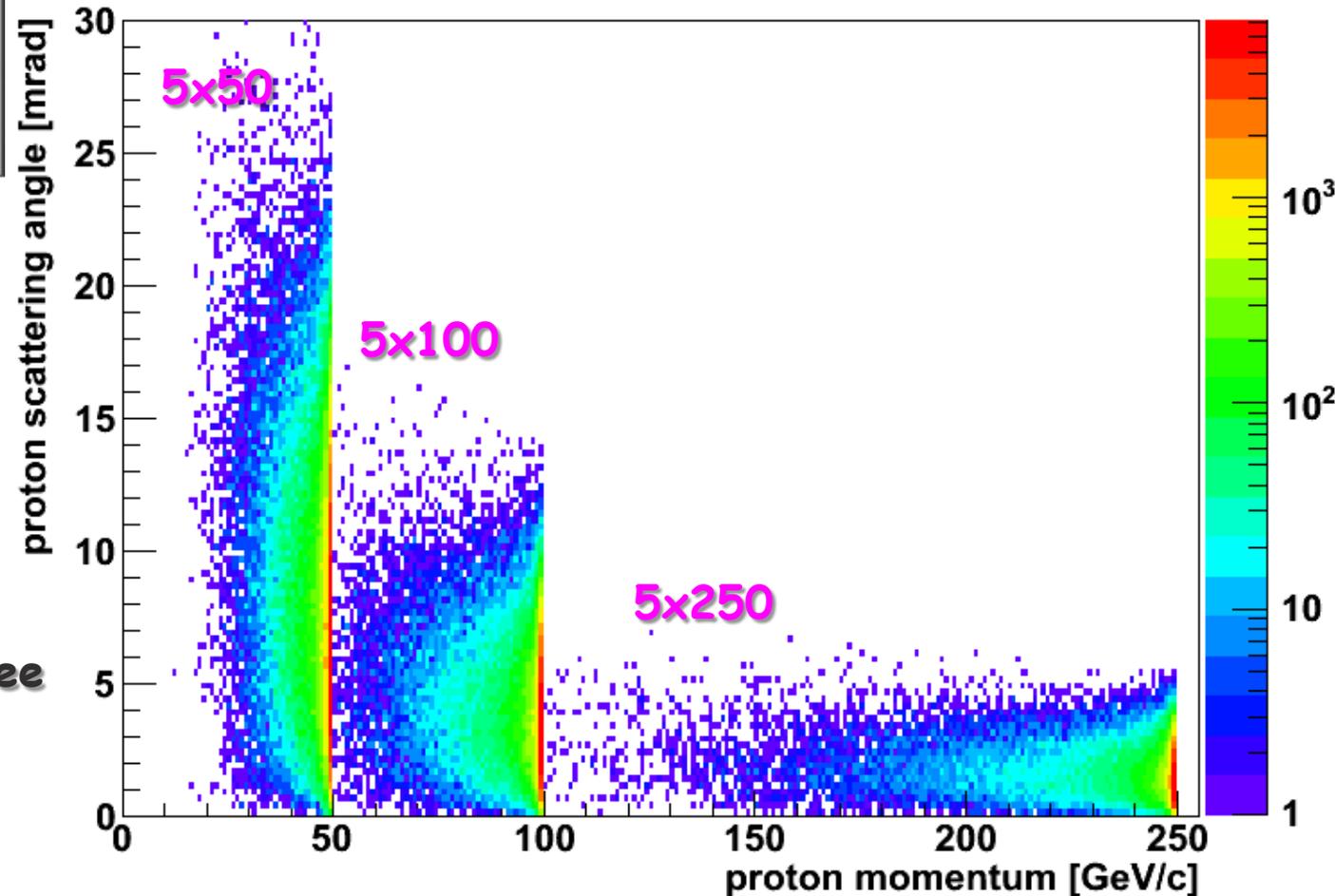
± 5 mrad acceptance seems sufficient

DIFFRACTIVE PHYSICS: P' KINEMATICS



$$t = (p_4 - p_2)^2 = 2[(m_p^{\text{in}} \cdot m_p^{\text{out}}) - (E^{\text{in}} E^{\text{out}} - p_z^{\text{in}} p_z^{\text{out}})]$$

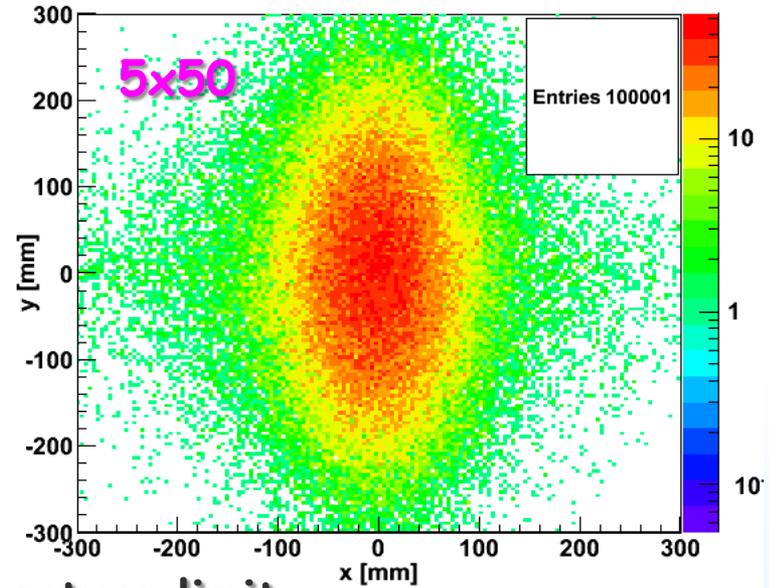
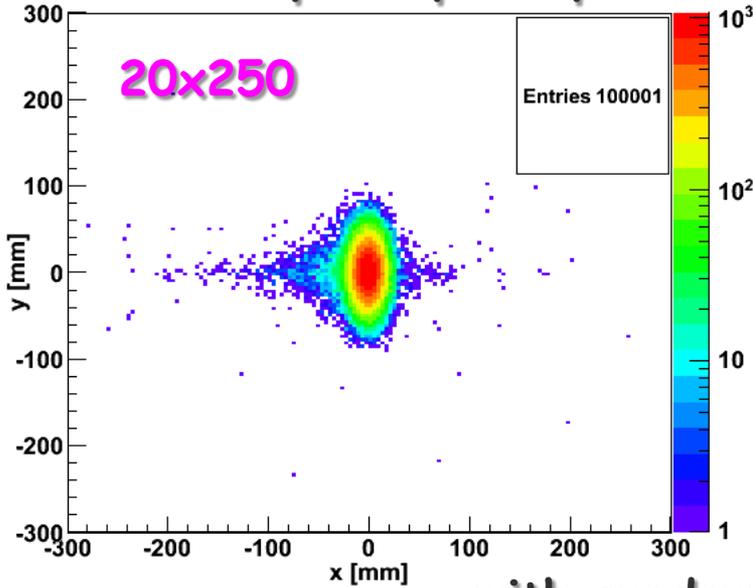
→ "Roman Pots" acceptance studies see later



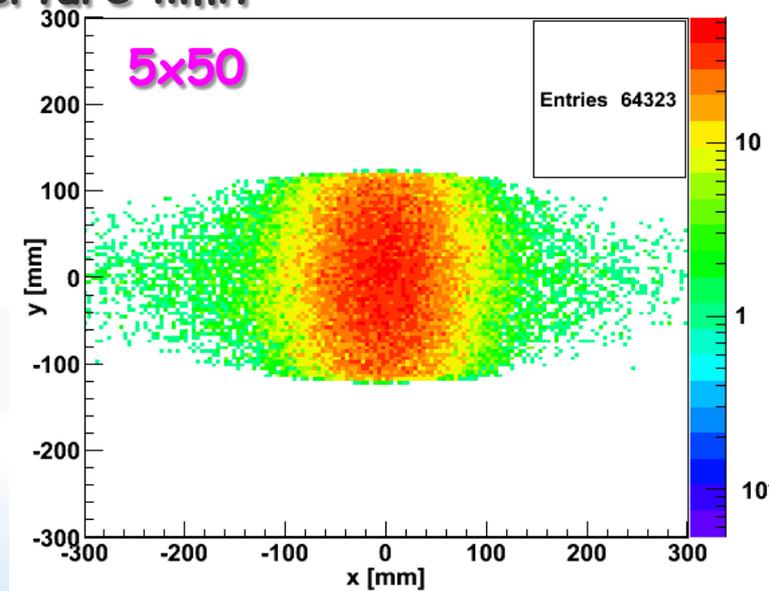
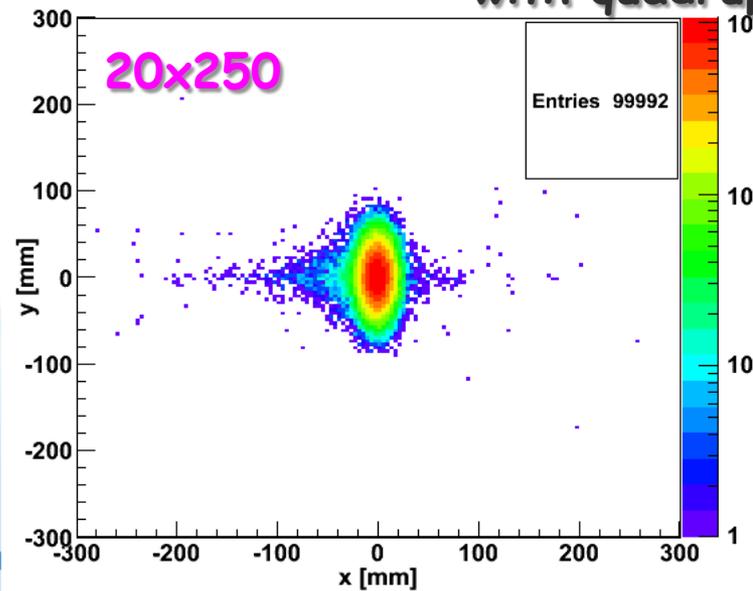
Simulations by J.H Lee

PROTON DISTRIBUTION IN Y VS X AT S=20 M

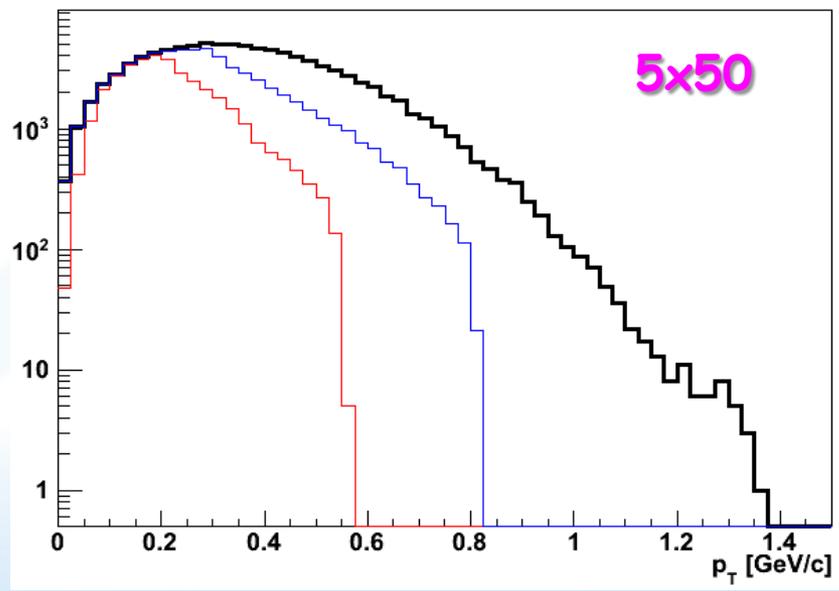
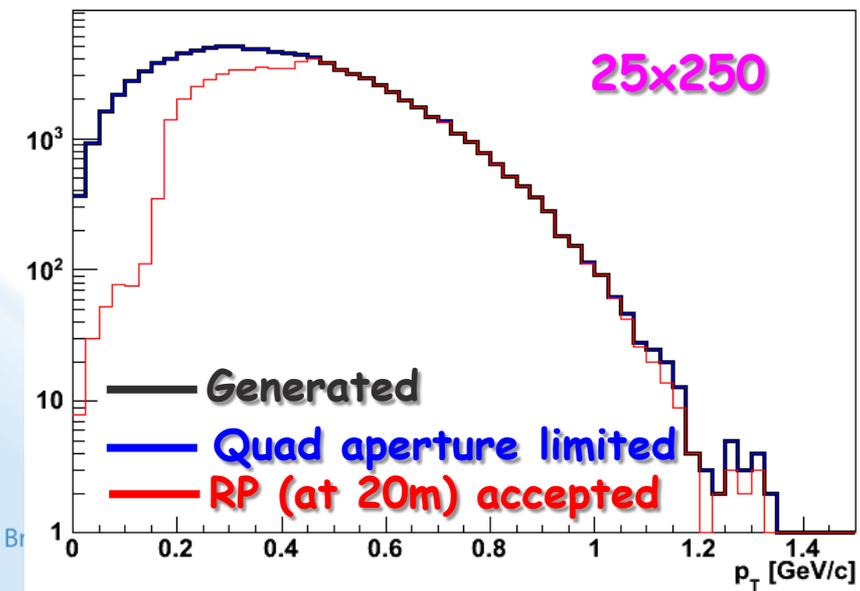
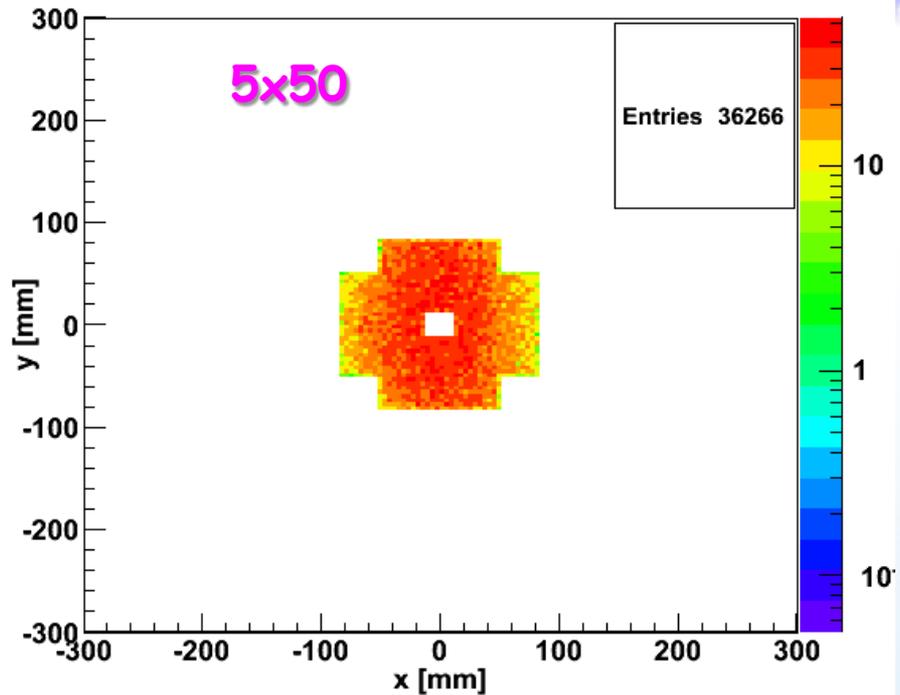
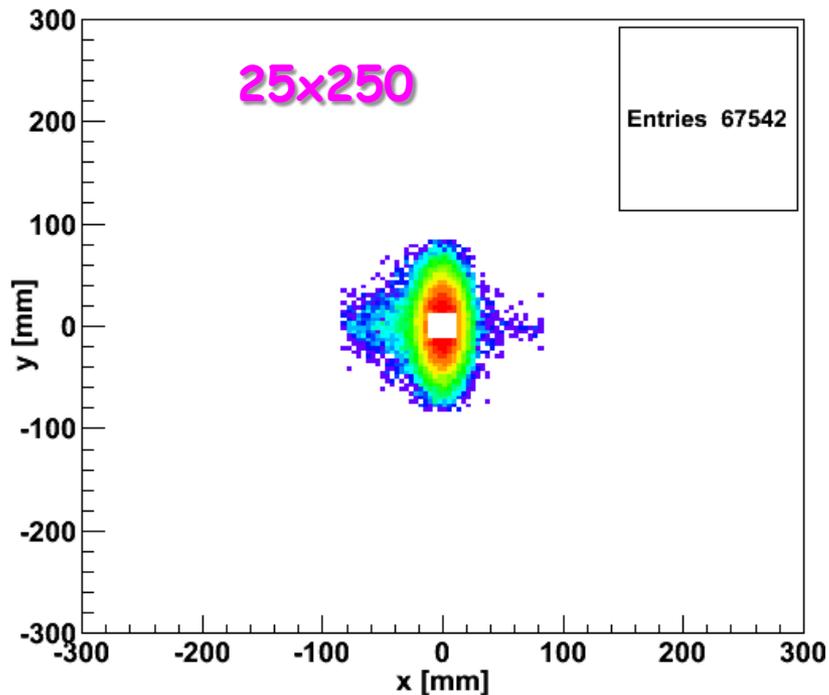
without quadrupole aperture limit



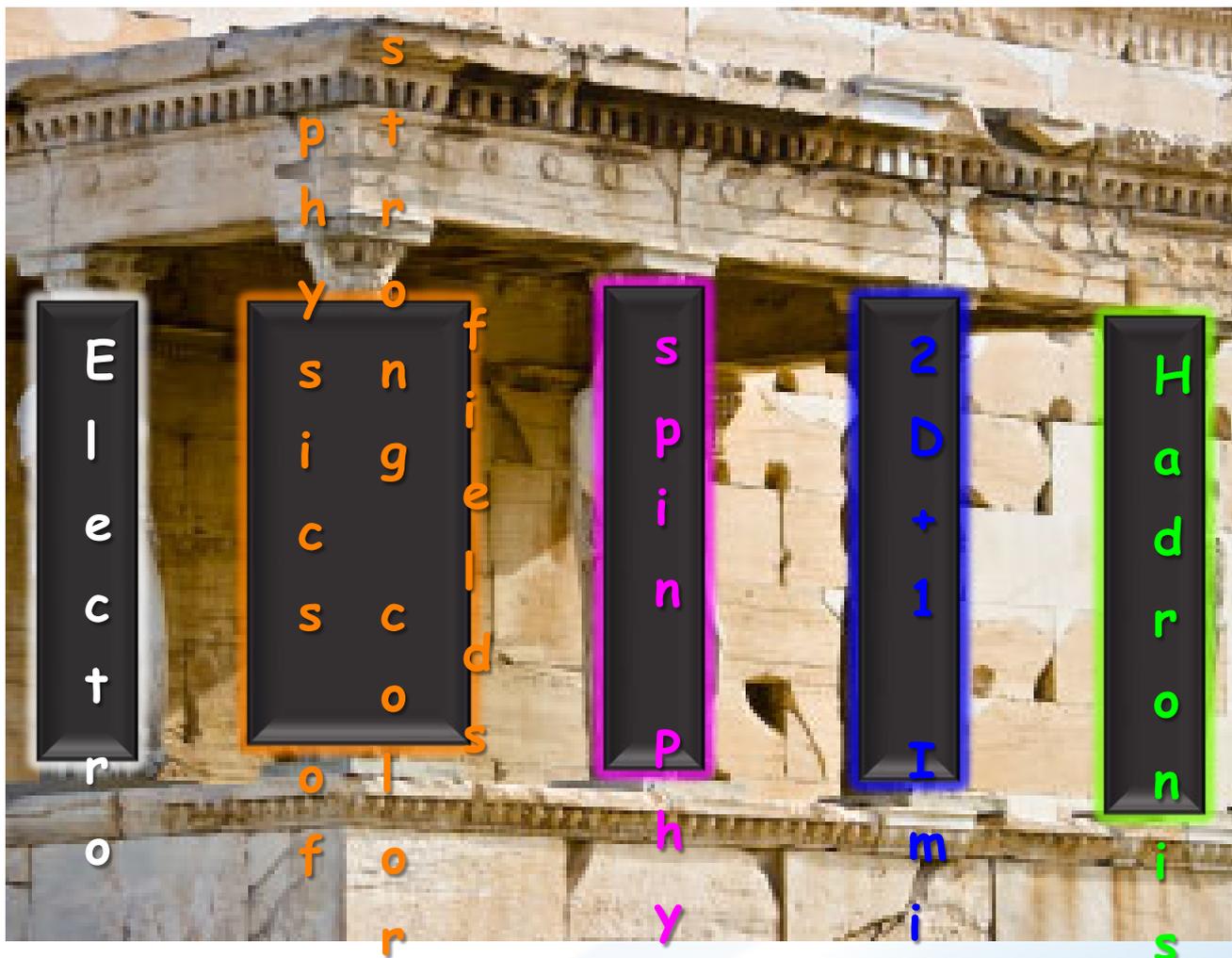
with quadrupole aperture limit



ACCEPTED IN "ROMAN POT" (EXAMPLE) AT $S=20M$



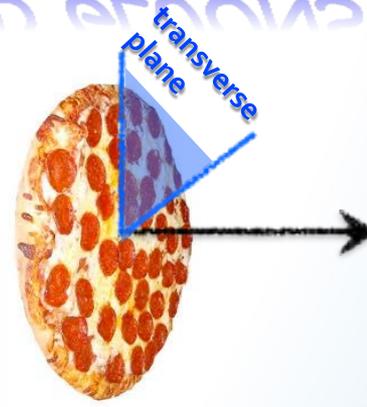
THE PILLARS OF THE eRHIC PHYSICS PROGRAM



Wide physics program with high requirements on detector and machine performance

THE PATH TO IMAGING QUARKS AND GLUONS

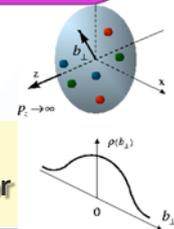
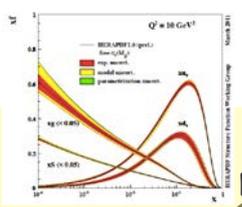
- PDFs do not resolve transverse momenta or positions in the nucleon
- fast moving nucleon turns into a 'pizza' but transverse size remains about 1 fm



compelling questions

- how are quarks and gluons spatially distributed
- how do they move in the transverse plane
- do they orbit and do we have access to spin-orbit correlations

→ required set of measurements & theoretical concepts



1-D

$f(x)$
parton densities

$F(t)$
form factor

$f(x, k_T)$

$f(x, b_T)$

$H(x, 0, t)$

$H(x, \xi, t)$

2+1-D

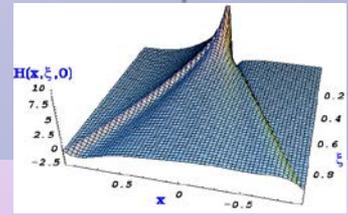
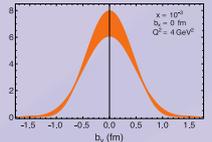
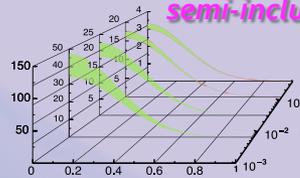
transv. mom. dep. PDF

impact par. dep. PDF

generalized PDF

semi-inclusive DIS

exclusive processes



4+1-D

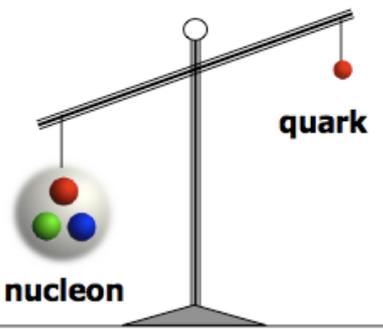
$W(x, k_T, b_T)$

high-level connection

measurable?

important in other branches of Physics

THE HADRONIC MASS PUZZLE

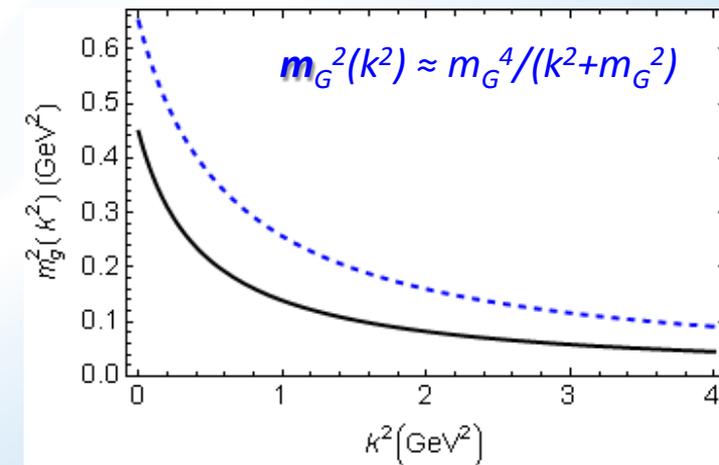
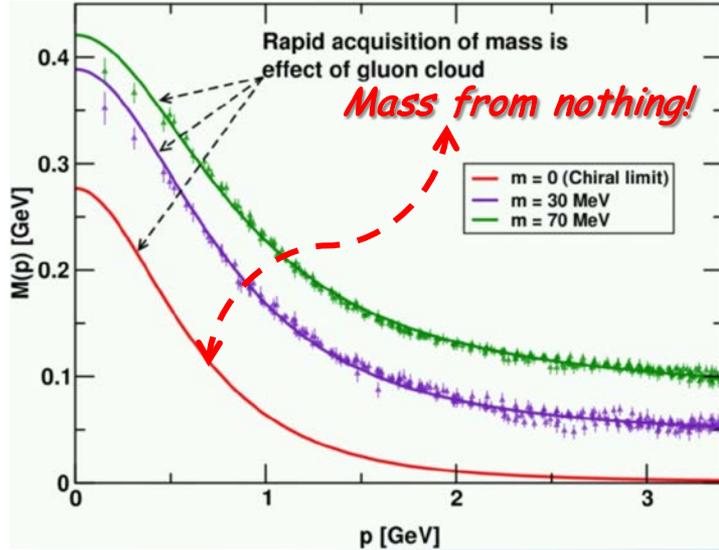


- ❑ In QCD, all “constants” of quantum mechanics are actually **strongly** momentum dependent: couplings, number density, mass, etc.
- ❑ So, a quark’s mass depends on its momentum.
- ❑ Mass function can be calculated and is depicted here.

in agreement: the vast bulk of the light-quark mass comes from a cloud of gluons, dragged along by the quark as it propagates.

- ❑ Continuum- and Lattice-QCD
- ❑ Running gluon mass
 - ❑ Gluon is massless in UV, in agreement with pQCD
 - ❑ Massive in infrared
 - ❑ $m_G(0) = 0.67\text{-}0.81 \text{ GeV}$
 - ❑ $m_G(m_G^2) = 0.53\text{-}0.64 \text{ GeV}$
- ❑ DSE prediction confirmed by numerical simulations of lattice-regularised QCD

C.D. Roberts, *Prog. Part. Nucl. Phys.* 61 (2008) 50
 M. Bhagwat & P.C. Tandy, *ATP Conf. Proc.* 842 (2006) 225-227



Qin et al., *Phys. Rev. C* 84 042202 (Rapid Comm.)

E.C. Aschenauer

HOW MANY GLUONS HAVE SPACE IN A PROTON?

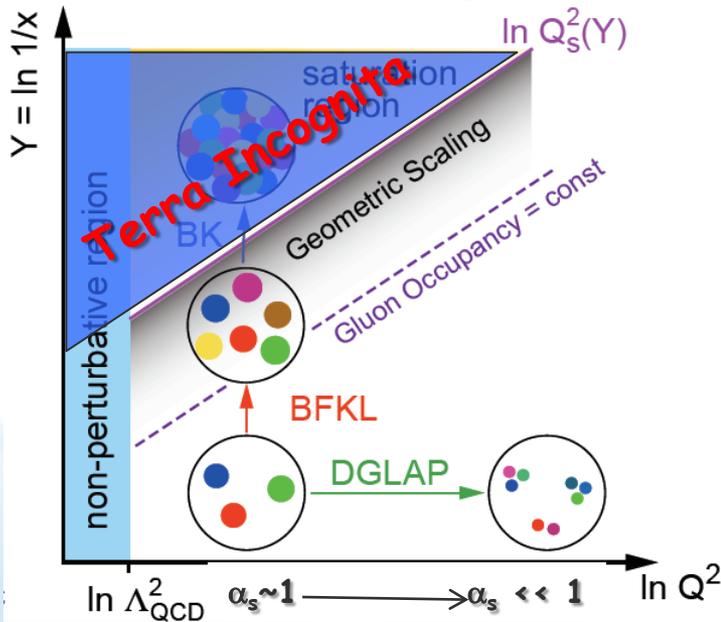
- at small x linear evolution gives strongly rising $g(x)$
 - cannot go on forever

- BK/JIMWLK **non-linear** evolution includes **recombination effects** \rightarrow **saturation**

- Dynamically generated scale

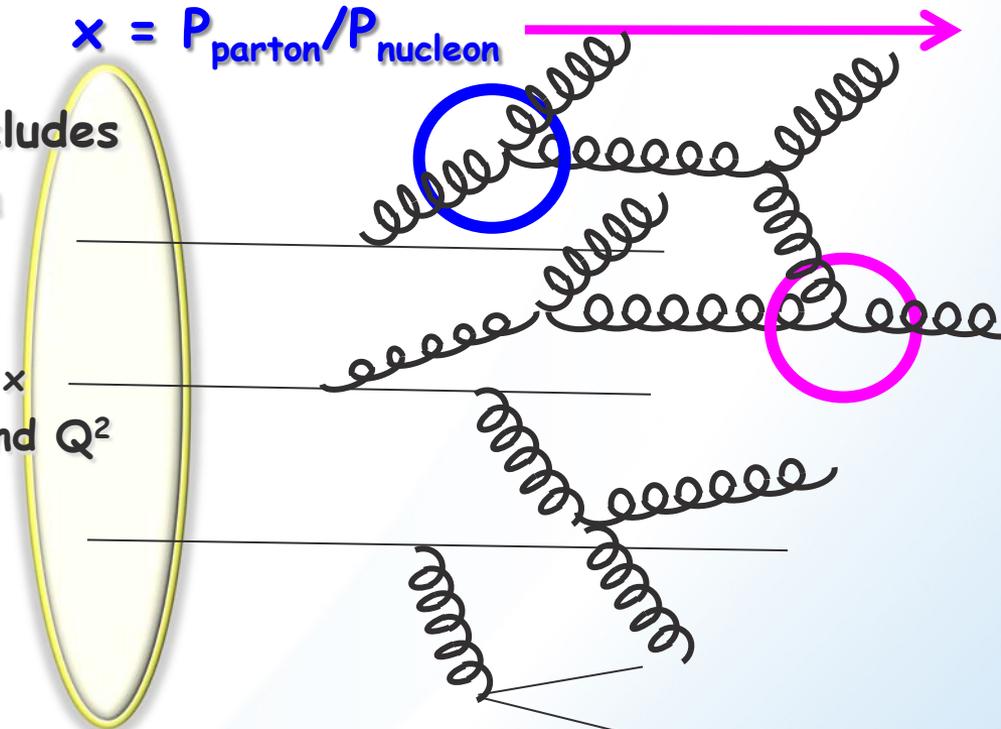
Saturation Scale: $Q_s^2(x)$

- Increases with energy or decreasing x
- Scale with $Q^2/Q_s^2(x)$ instead of x and Q^2



small x / higher energy

$x = P_{\text{parton}}/P_{\text{nucleon}}$



Bremsstrahlung
 $\sim \alpha_s \ln(1/x)$

Recombination
 $\sim \alpha_s \rho$

Saturation must set in at low $x \rightarrow$ high occupancy

DIHADRON CORRELATIONS IN eA AT EIC

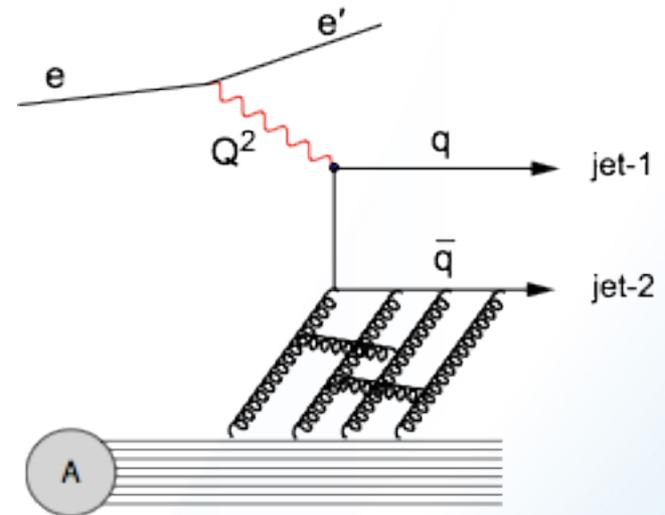
EIC:

- Extract the spatial multi-gluon correlations and study their non-linear evolution
 - essential for understanding the transition from a deconfined into a confined state.

Advantage over $p(d)A$:

- eA experimentally much cleaner
 - no “spectator” background to subtract
 - Access to the exact kinematics of the DIS process (x , Q^2)

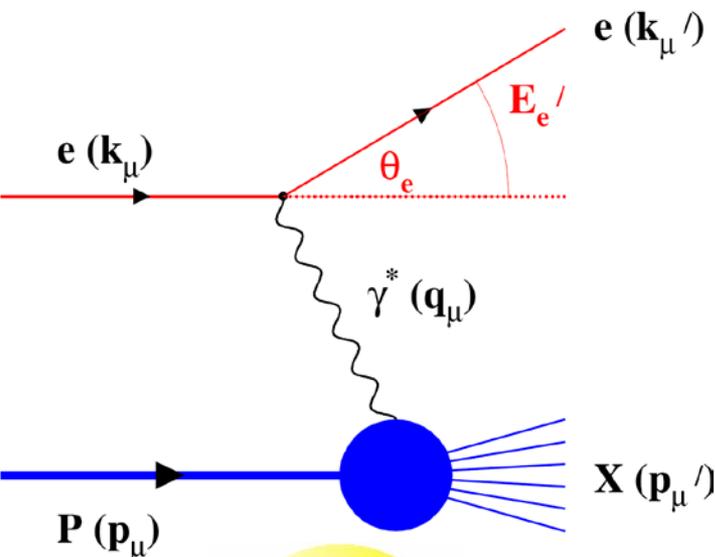
Perfect saturation signature:



Either jets or use leading hadrons from jets (dihadrons)

DEEP INELASTIC SCATTERING

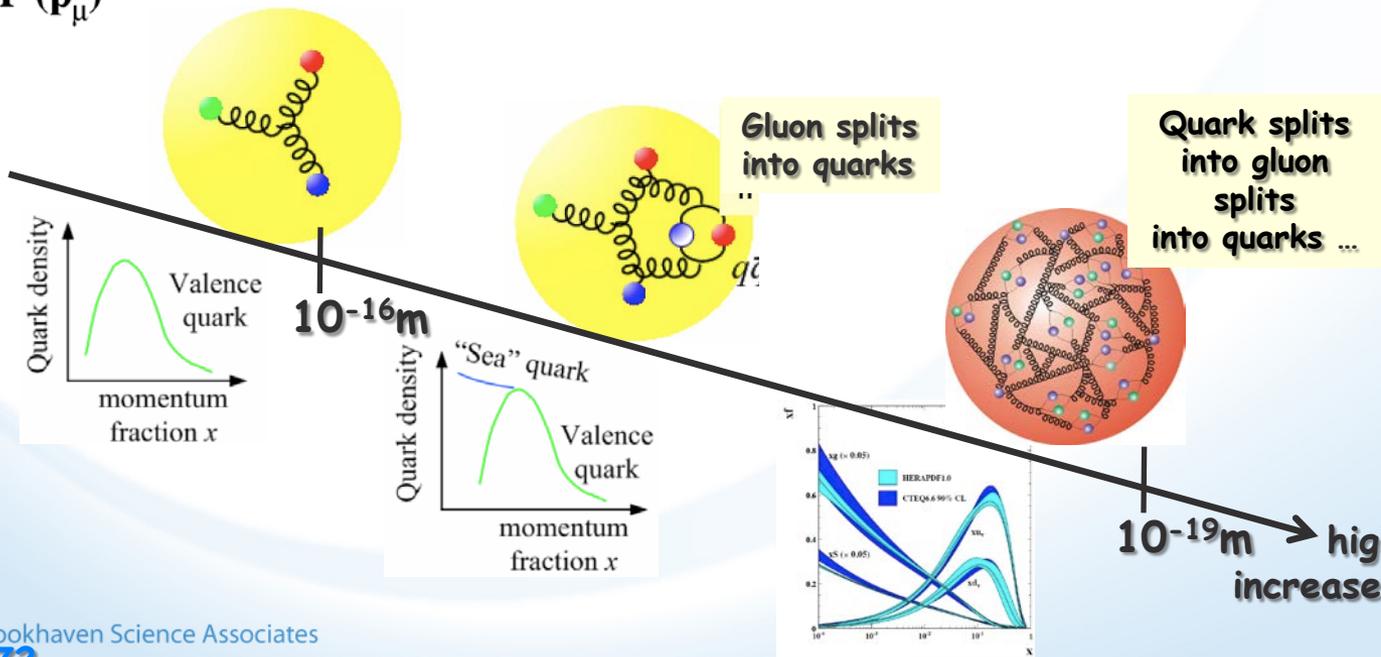
Kinematics:



$Q^2 = -q^2 = -(k_\mu - k'_\mu)^2$ Measure of resolution power
 $Q^2 = 2E_e E'_e (1 - \cos \Theta_{e'})$

$y = \frac{pq}{pk} = 1 - \frac{E'_e}{E_e} \cos^2 \left(\frac{\theta'_e}{2} \right)$ Measure of inelasticity

$x = \frac{Q^2}{2pq} = \frac{Q^2}{sy}$ Measure of momentum fraction of struck quark



$10^{-19}m \rightarrow$ higher \sqrt{s}
 increases resolution

HELICITY STRUCTURE - OPEN QUESTIONS



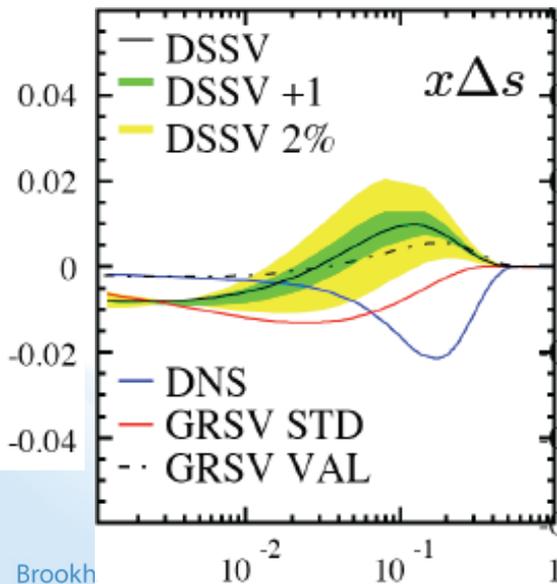
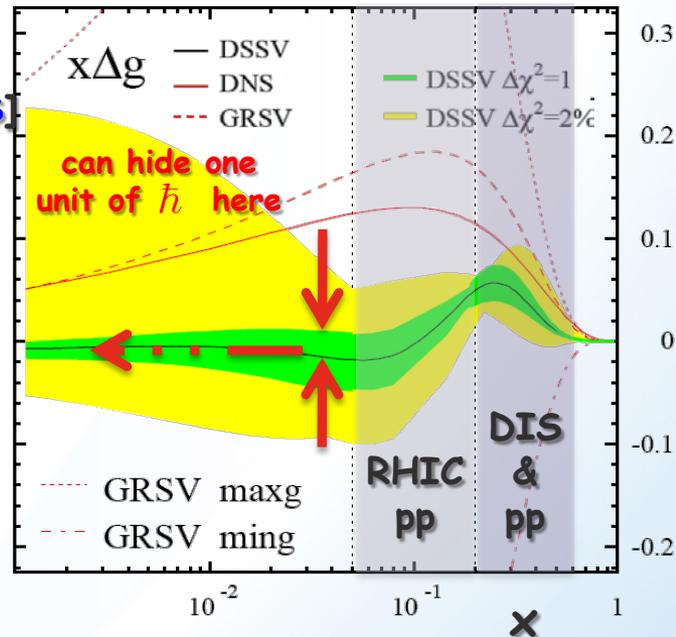
significant experimental and theoretical progress in past 25+ years, yet many unknowns ...



$\Delta g(x, Q^2)$

- found to be small at $0.05 < x < 0.2$ [RHIC, COMPASS, HERMES]
- RHIC can slightly extend x range & reduce uncertainties [500 GeV running & particle correlations]

yet, full 1st moment [proton spin sum] will remain to have significant uncertainties from unmeasured small x region



Δq 's (x, Q^2)

- known: quarks contribute much less to proton spin than expected from quark models
large uncertainties in $\Delta\Sigma$ from unmeasured small x
- surprisingly small/positive Δs from SIDIS: large SU(3) breaking?
- flavor separation not well known, e.g., $\Delta\bar{u} - \Delta\bar{d}$

SCALING VIOLATIONS AT SMALL X

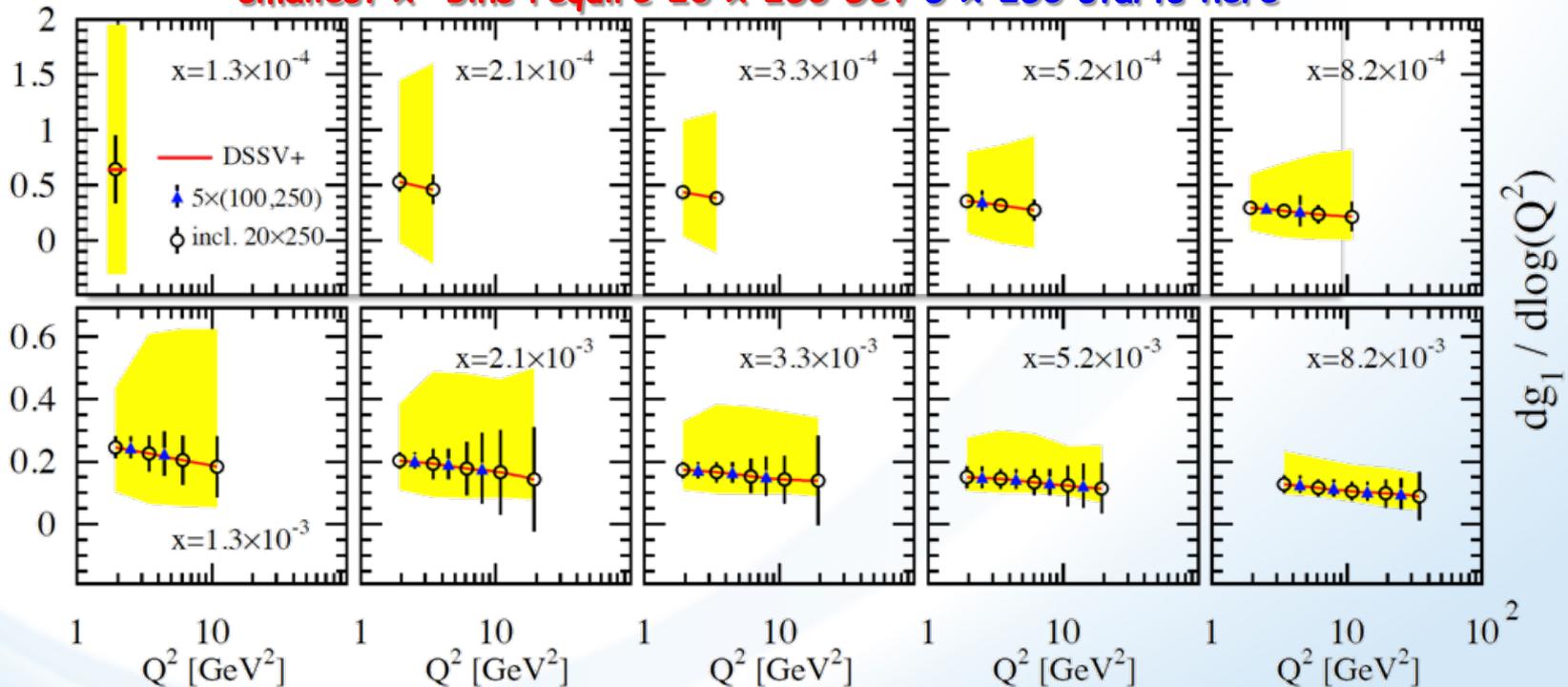
rough small-x approximation to Q^2 -evolution:

$$\frac{dg_1}{d \log(Q^2)} \propto -\Delta g(x, Q^2)$$

spread in $\Delta g(x, Q^2)$ translates into spread of scaling violations for $g_1(x, Q^2)$

- need x-bins with a least two Q^2 values to compute derivative (limits x reach somewhat)

smallest x bins require 20 x 250 GeV 5 x 250 starts here



- error bars for moderate 10fb^{-1} per c.m.s. energy; bands parameterize current DSSV+ uncertainties

CONSTRAIN J_q VIA GPD E

observables sensitive to E :

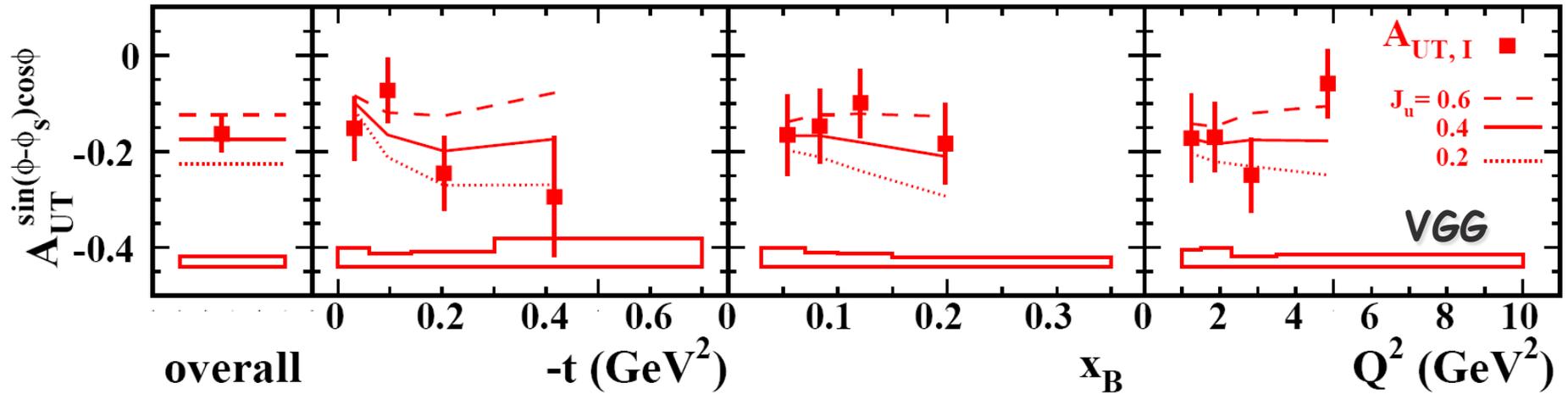
(J_q input parameter in ansatz for E)

□ DVCS A_{UT} : HERMES

□ nDVCS A_{LU} : Hall A

Hermes DVCS-TTSA [arXiv: 0802.2499]:

$$A_{UT}^{\sin(\phi-\phi_T)\cos\phi} \sim \text{Im}(F_2 H - F_1 E)$$



eRHIC:

HERMES like A_{UT}

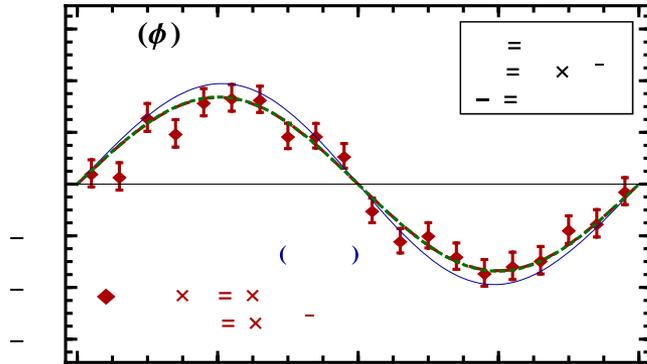
20 GeV x 250 GeV

Lumi: $2 \times 50 \text{fb}^{-1}$

— $\kappa = +$
 - - $\kappa = =$
 ··· $\kappa = -$

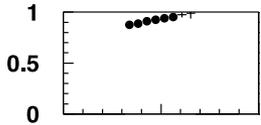
DIFFERENT DVCS ASYMMETRIES

[arXiv:1304.0077](https://arxiv.org/abs/1304.0077)



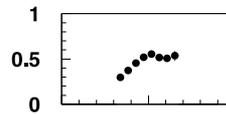
- - -
—

BH FRACTION



5 GeV x 100 GeV

BH subtraction will be relevant in stage 1, at large y , depending on the x - Q^2 bin



<2.51,3.98>

<3.98,6.31>

<6.31,10>

<10,15.8>

<15.8,25.1>

<25.1,39.8>

<39.8,63.1>

<63.1,100>

BUT...

5 GeV x 100 GeV

and

5 GeV x 250 GeV

are overlapping:

x -sec. measurements at 5x250 at low- y can crosscheck the BH subtrac. made for 5x100

,6.31>

<6.31,10>

<10,15.8>

<15.8,25.1>

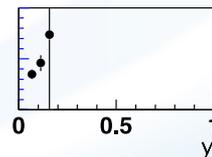
<25.1,39.8>

<39.8,63.1>

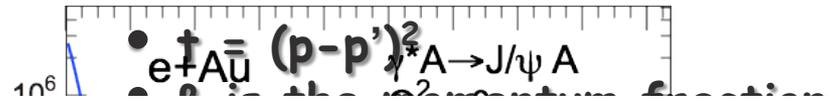
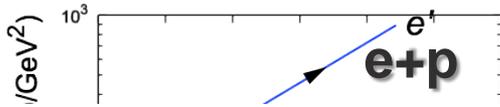
<63.1,100>

20 GeV x 250 GeV
BH subtraction will be not an issue for $y < 0.6$

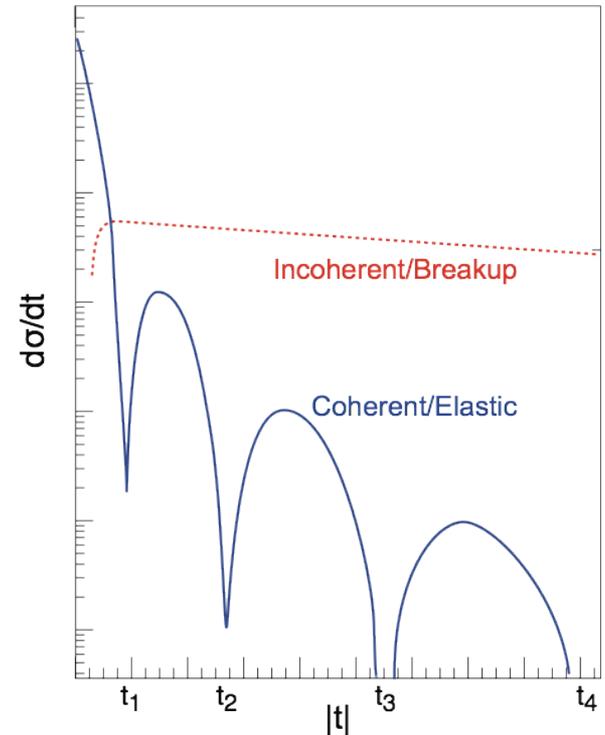
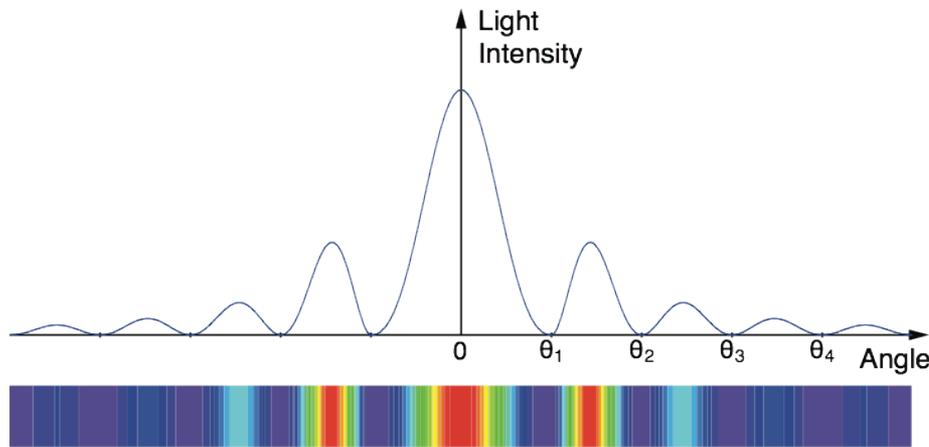
BNL EIC Science Task Force



HARD DIFFRACTION IN DIS AT SMALL X



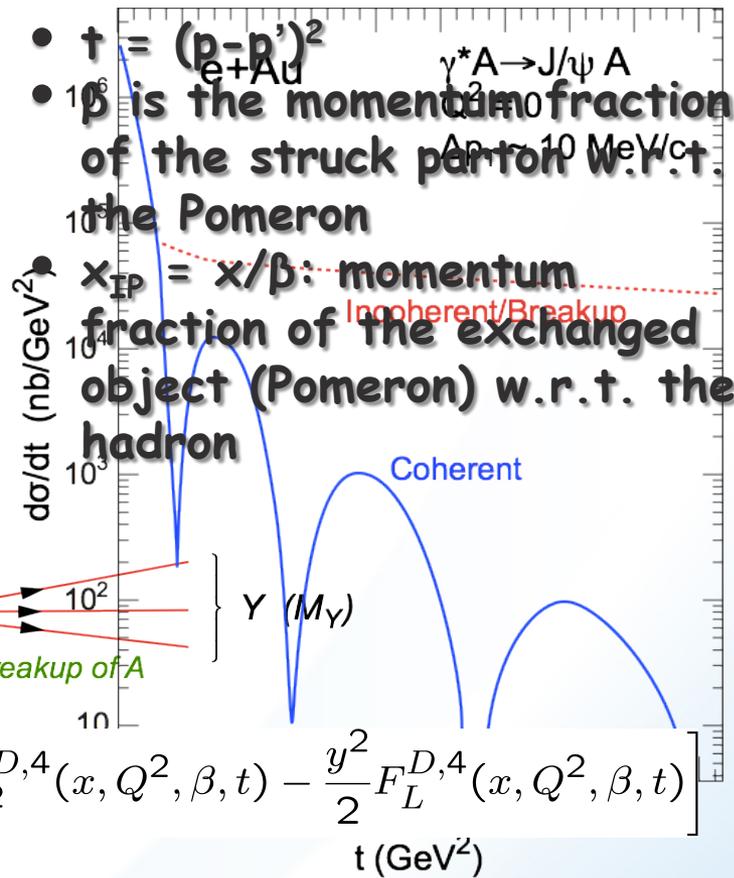
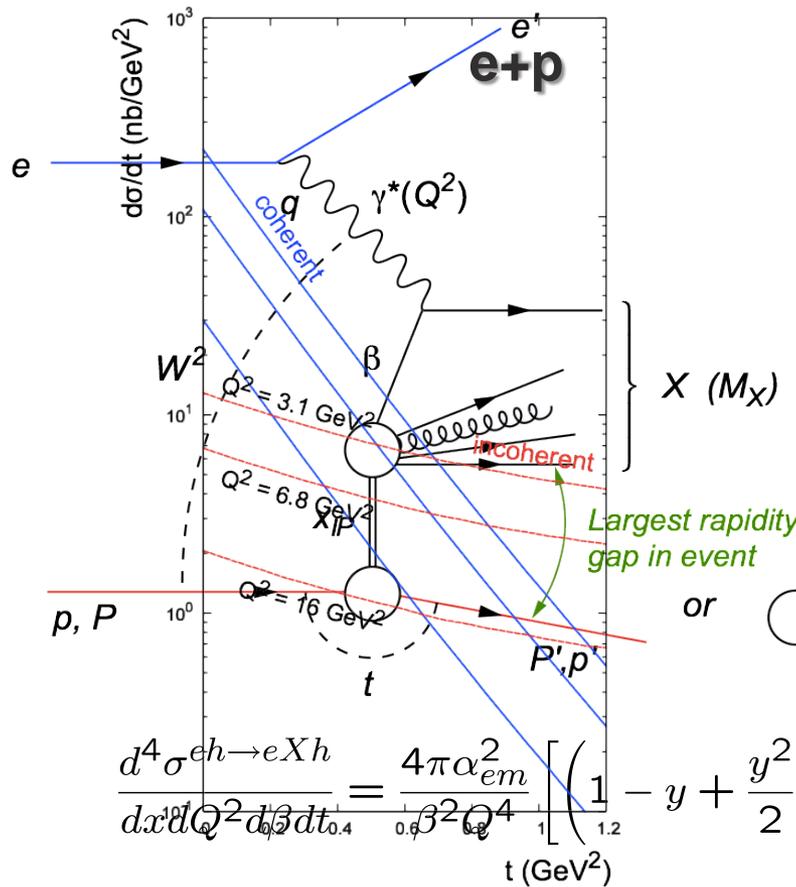
Diffraction Analogy: plane monochromatic wave incident on a circular screen of radius R



- ▶ incoherent \Leftrightarrow breakup of p
- ▶ HERA: 15% of all events are hard diffractive

- ▶ breakup into nucleons (nucleons intact)
- ▶ incoherent diffraction
- ▶ Predictions: $\sigma_{\text{diff}}/\sigma_{\text{tot}}$ in $e+A \sim 25-40\%$

HARD DIFFRACTION IN DIS AT SMALL X



- Diffraction in e+p:
 - ▶ coherent \Leftrightarrow p intact
 - ▶ incoherent \Leftrightarrow breakup of p
 - ▶ HERA: 15% of all events are hard diffractive

- Diffraction in e+A:
 - ▶ coherent diffraction (nuclei intact)
 - ▶ breakup into nucleons (nucleons intact)
 - ▶ incoherent diffraction
 - ▶ Predictions: $\sigma_{diff}/\sigma_{tot}$ in e+A \sim 25-40%

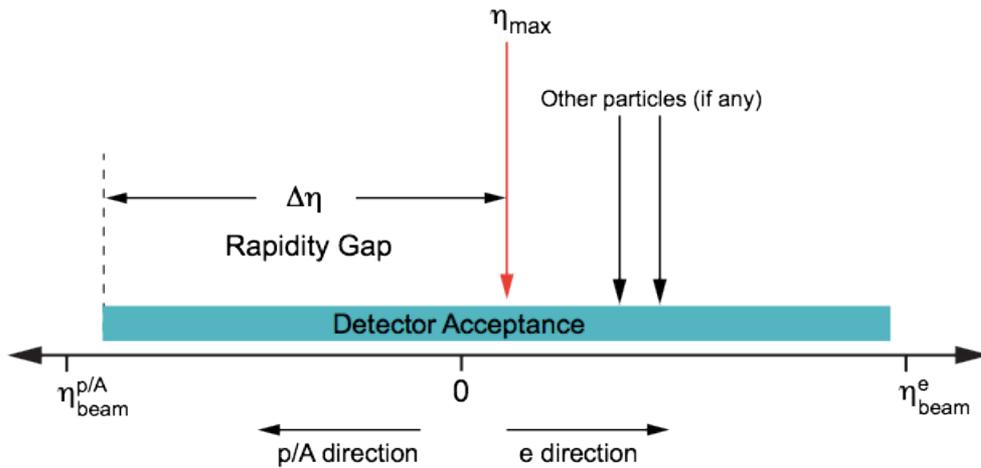
LARGE RAPIDITY GAP METHOD (LRG)

Identify Most Forward Going Particle (MFP)

- Works at HERA but at higher \sqrt{s}
- EIC smaller beam rapidities

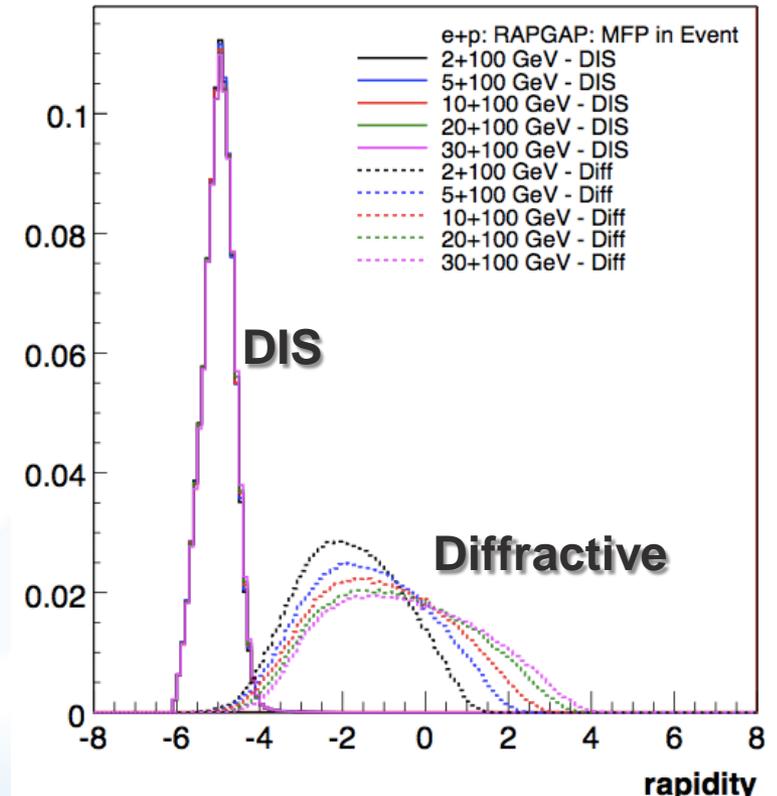
Diffractive ρ^0 production at EIC:
 η of MFP

M. Lamont '10

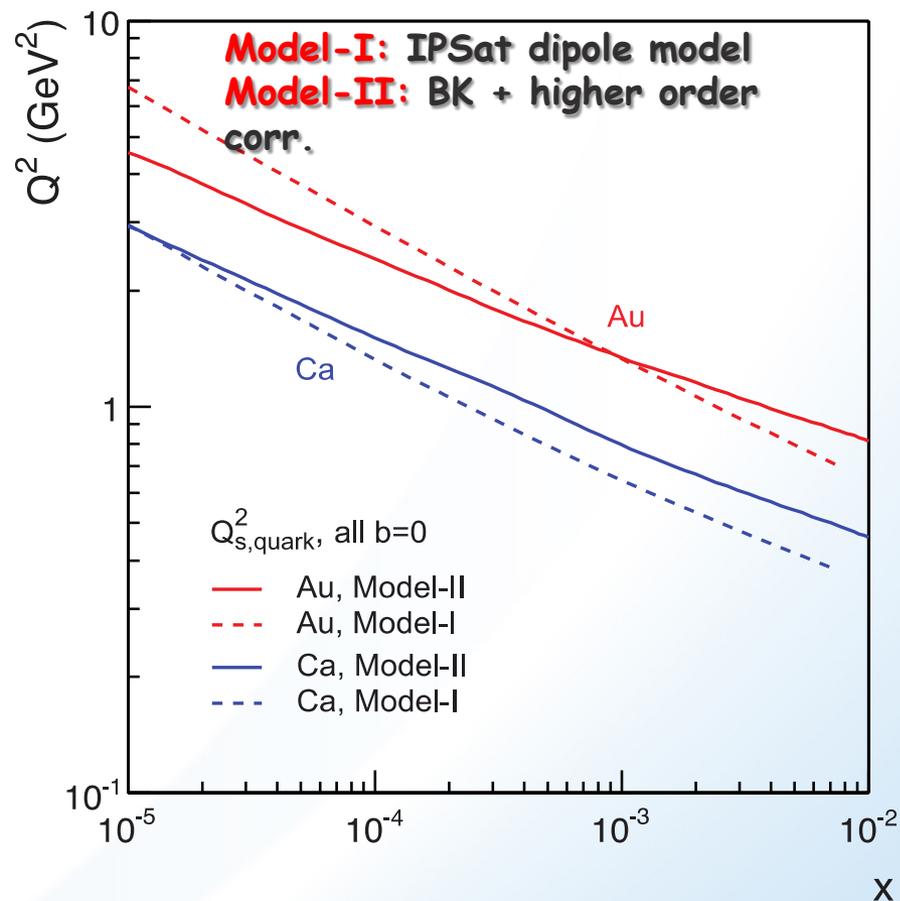
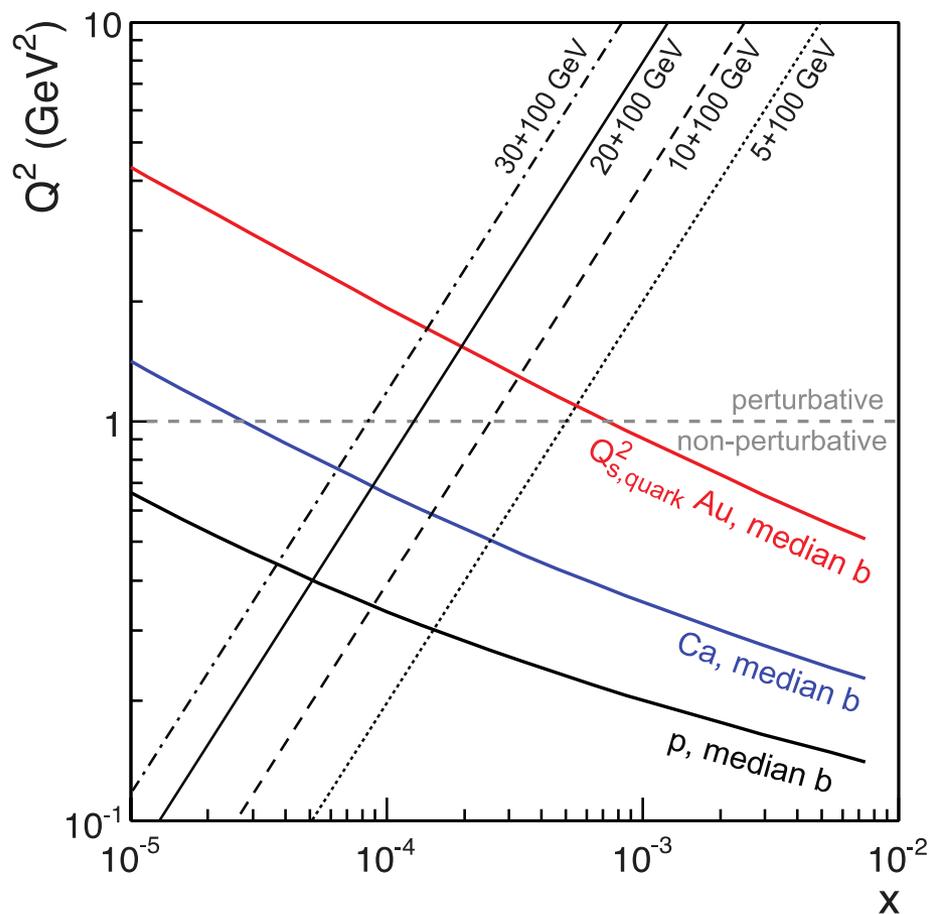


Hermeticity requirement:

- needs just to detector presence
- does not need momentum or PID
- simulations: \sqrt{s} not a show stopper for EIC (can achieve 1% contamination, 80% efficiency)



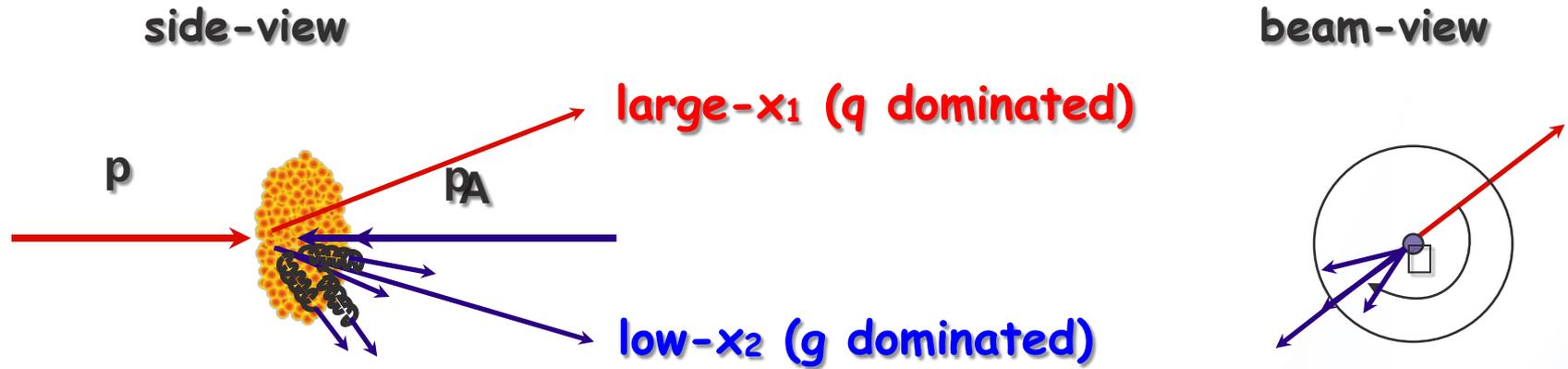
BEST KNOWLEDGE OF Q_s IN eA



Enhancement of Q_s with A

⇒ saturation regime reached at significantly lower energy in nuclei

h-h FORWARD CORRELATION IN p(d)A AT RHIC



Low gluon density (pp):
pQCD predicts $2 \rightarrow 2$ process
 \Rightarrow back-to-back di-jet

High gluon density (pA):
 $2 \rightarrow$ many process
 \Rightarrow expect broadening of away-side

- Small- x evolution \leftrightarrow multiple emissions
- Multiple emissions \rightarrow broadening
- Back-to-back jets (here leading hadrons) may get broadening in p_T with a spread of the order of Q_s

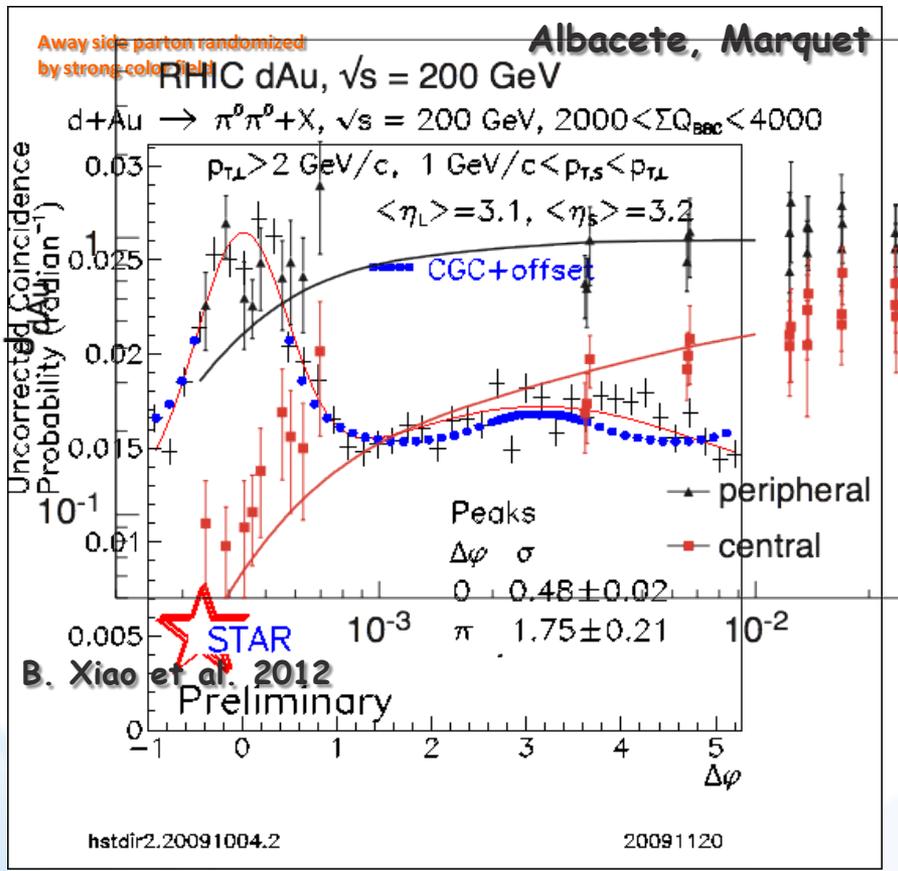
First prediction by: C. Marquet ('07)

Latest review: Stasto, Xiao, Yuan arXiv:1109.1817 (Sep. '11)

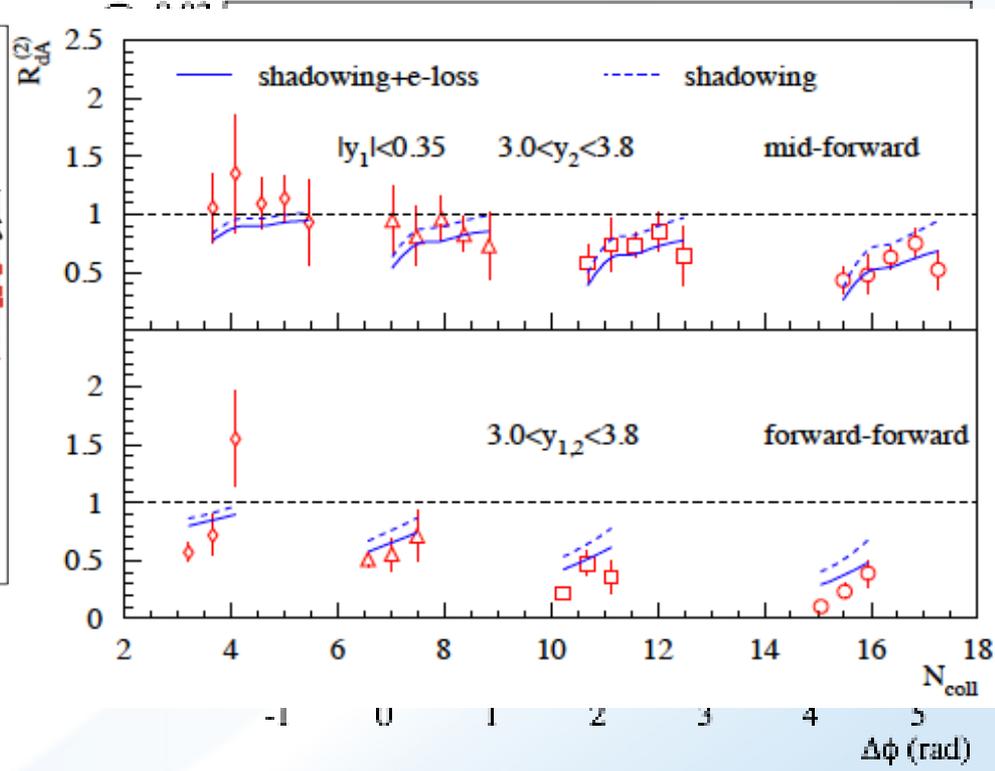
FORWARD CORRELATIONS IN dA AT RHIC

1 question, 2 answers

Initial state saturation model



"Non-initial state" shadowing model



$$\langle q_{\perp}^2 \rangle_{dAu} = \langle q_{\perp}^2 \rangle_{pp} + \Delta \langle q_{\perp}^2 \rangle$$

How saturated is the initial state?

WHY IS DIFFRACTION SO IMPORTANT

- Sensitive to **spatial** gluon distribution

$$\frac{d\sigma}{dt} \equiv \frac{\text{Fourier Transformation}}{\text{of Source Density } \rho_g(b)}$$

- Hot topic:

- Lumpiness?
- Just Wood-Saxon+nucleon $g(b)$

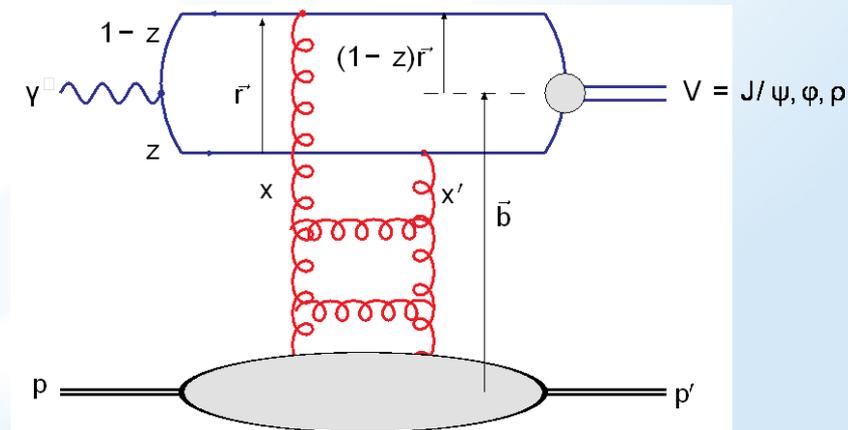
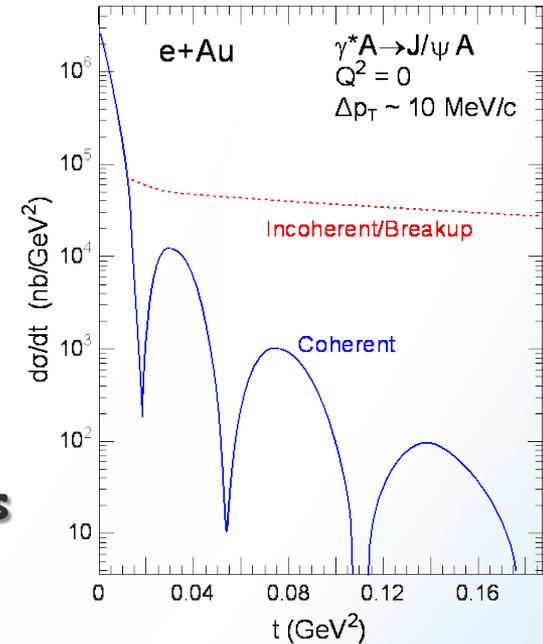
- Incoherent case: measure fluctuations/lumpiness in $g_A(b)$

- Sensitive to gluon momentum distributions

- $\sigma \sim g(x, Q^2)^2$

$$\frac{d\sigma^{\gamma^* p \rightarrow pV}}{dt} \sim \left| \int \Psi_V^* \frac{d\sigma_{q\bar{q}}}{d^2b} \Psi e^{-ib\Delta} \right|^2$$

$$\frac{d\sigma_{q\bar{q}}}{d^2b} \sim r^2 \alpha_s x g(x, \mu^2) T(b)$$



MEASURING F_L WITH THE EIC (I)

$$\frac{d^2\sigma^{ep \rightarrow eX}}{dx dQ^2} = \frac{4\pi\alpha_{e.m.}^2}{xQ^4} \left[\left(1 - y + \frac{y^2}{2}\right) F_2(x, Q^2) - \frac{y^2}{2} F_L(x, Q^2) \right]$$

quark+anti-quark
momentum distributions

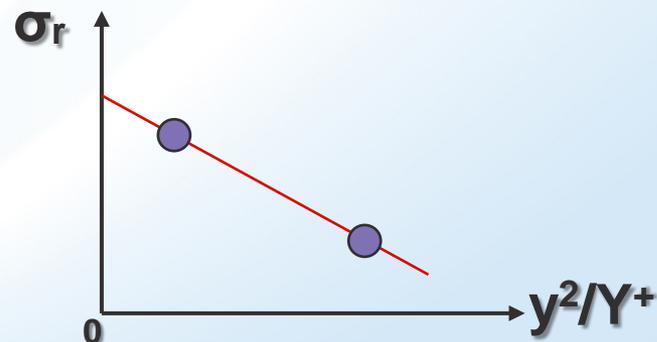
gluon momentum
distribution

In practice use reduced cross-section:

$$\begin{aligned} \sigma_r &= \left(\frac{d^2\sigma}{dx dQ^2} \right) \frac{xQ^4}{2\pi\alpha^2 [1 + (1 - y)^2]} = F_2(x, Q^2) - \frac{y^2}{1 + (1 - y)^2} F_L(x, Q^2) \\ &= F_2(x, Q^2) - \frac{y^2}{Y^+} F_L(x, Q^2) \end{aligned}$$

□ How to extract F_L

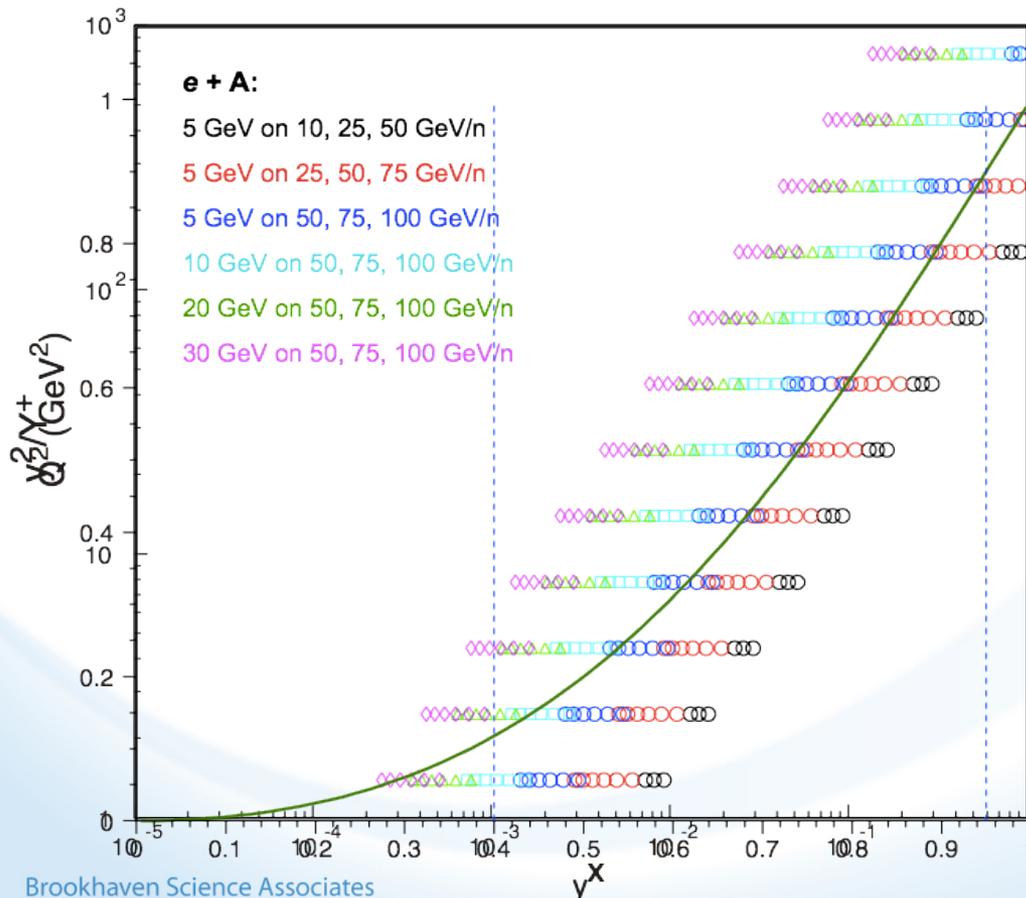
- Need different values of y^2/Y^+
- F_L slope of σ_r vs y^2/Y^+
- F_2 intercept of σ_r vs y^2/Y^+ with y -axis



MEASURING F_L WITH THE EIC (II)

In order to extract F_L one needs **at least two measurements** of the inclusive cross section with “wide” span in inelasticity parameter y ($Q^2 = sxy$)

F_L requires runs at various \sqrt{s} \Rightarrow longer program



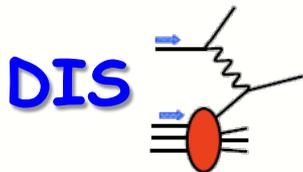
EIC studies:

- Statistical error is negligible in essentially whole range
- Systematical Error
 - Calibration
 - Normalization
 - Experiment
 - Radiative Corrections

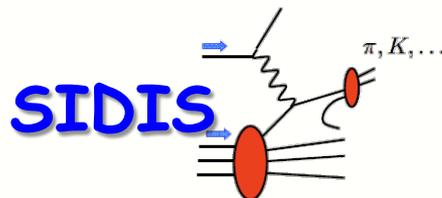
PREPARATION OF DIS AND SIDIS MOCK DATA



- PEPSI MC to generate σ^{++} and σ^{+-} with LO GRSV PDFs



inclusive final-state



identified charged pions and kaons

assume modest 10 fb^{-1} for each energy, 70% beam polarizations

$Q^2 > 1 \text{ GeV}^2$, $0.01 < y < 0.95$, invariant mass $W^2 > 10 \text{ GeV}^2$

depolarization factor of virtual photon $D(y, Q^2) > 0.1$ (cuts on small y)

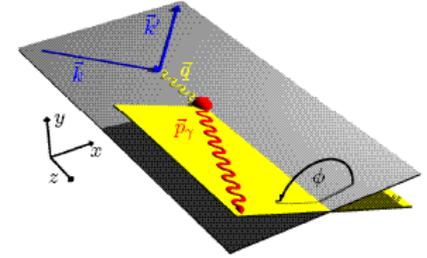
scattered lepton: $1^\circ < \theta_{\text{elec}} < 179^\circ$ and $p_{\text{elec}} > 0.5 \text{ GeV}$

hadron: $p_{\text{hadr}} > 1 \text{ GeV}$, $0.2 < z < 0.9$,
 $1^\circ < \theta_{\text{hadr}} < 179^\circ$

- use rel. uncertainties of data to generate mock data by randomizing around **DSSV+** by $1-\sigma$
- **SIDIS**: incl. typical 5% (10%) uncertainty for pion (kaon) frag. fcts (from **DSS** analysis)

DVCS ASYMMETRIES

$$d\sigma \sim \left(\tau_{BH}^* \tau_{DVCS} + \tau_{DVCS}^* \tau_{BH} \right) + |\tau_{BH}|^2 + |\tau_{DVCS}|^2$$



→ different charges: $e^+ e^-$:

$$\Delta\sigma_C \sim \cos\phi \cdot \text{Re}\{ H + \cancel{\xi \tilde{H}} + \dots \} \quad \Rightarrow \quad H$$

→ polarization observables:

$$\Delta\sigma_{LU} \sim \sin\phi \cdot \text{Im}\{ H + \cancel{\xi \tilde{H}} + \cancel{kE} \} \quad \Rightarrow \quad \Delta\sigma_{UT}$$

$$\Delta\sigma_{UL} \sim \sin\phi \cdot \text{Im}\{ \tilde{H} + \cancel{\xi H} + \dots \} \quad \Rightarrow \quad \begin{matrix} \tilde{H} \\ \swarrow \quad \searrow \\ \text{beam} \quad \text{target} \end{matrix}$$

$$\Delta\sigma_{UT} \sim \sin\phi \cdot \text{Im}\{ k(H - E) + \dots \} \quad \Rightarrow \quad H, E$$

$$\xi = x_B / (2 - x_B) \quad k = t / 4M^2 \quad \text{kinematically suppressed}$$