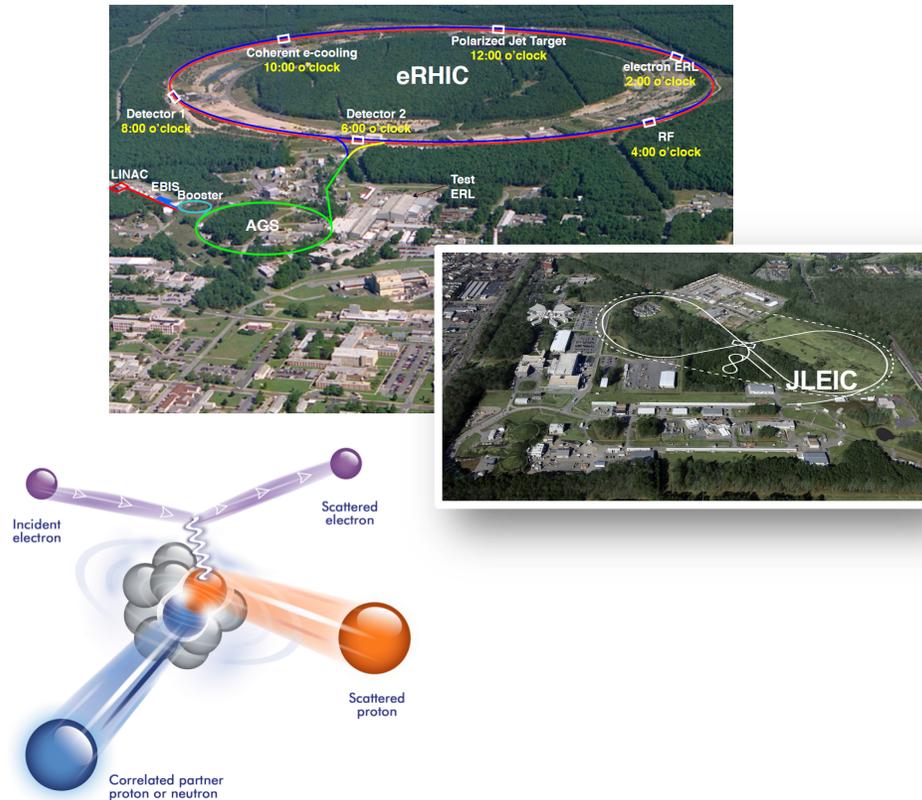


Experimental SRC/EMC Studies: Opportunities and Challenges at EIC

- in a broader context of forward detection requirements



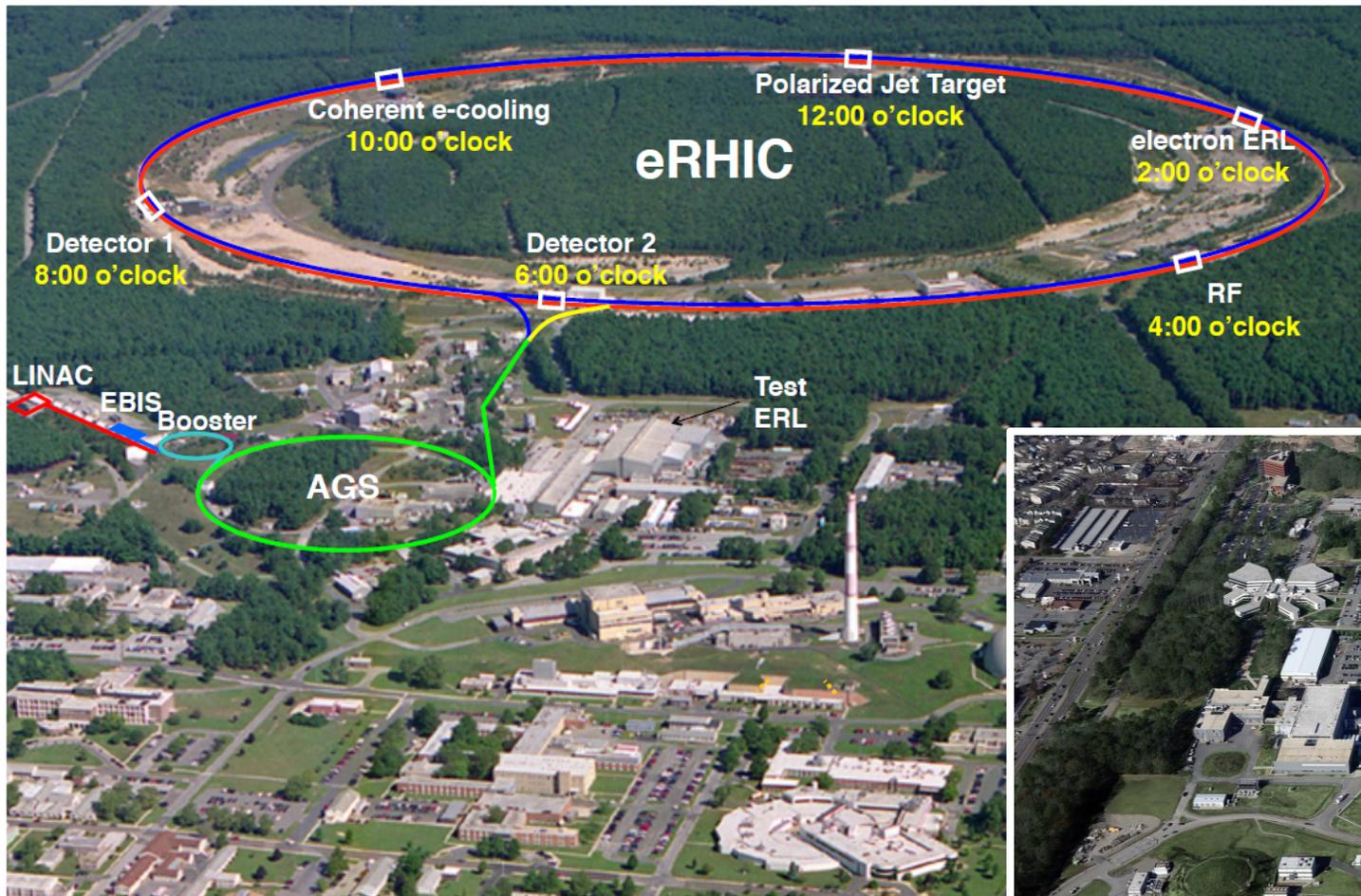
Pawel Nadel-Turonski
Stony Brook University

Joint CFNS/RBRC on Physics and Detector Requirements at Zero-Degree of Colliders,,
Stony Brook, September 24-26, 2019

Outline

- Forward detection requirements at the EIC
- SRC measurements at the EIC
- Key measurements that have similar requirements

Proposed EIC implementations at JLab and BNL



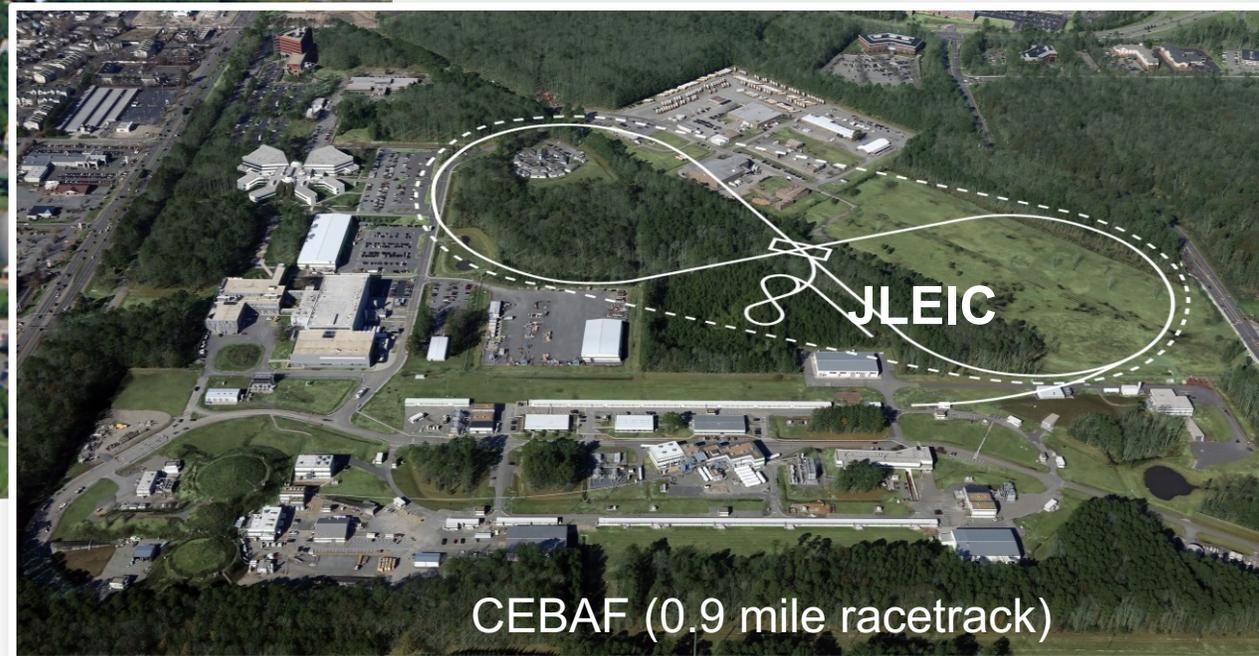
BNL EIC at RHIC (eRHIC)

18 GeV e (10 GeV lumi max) on 275 GeV p

Parameters are similar

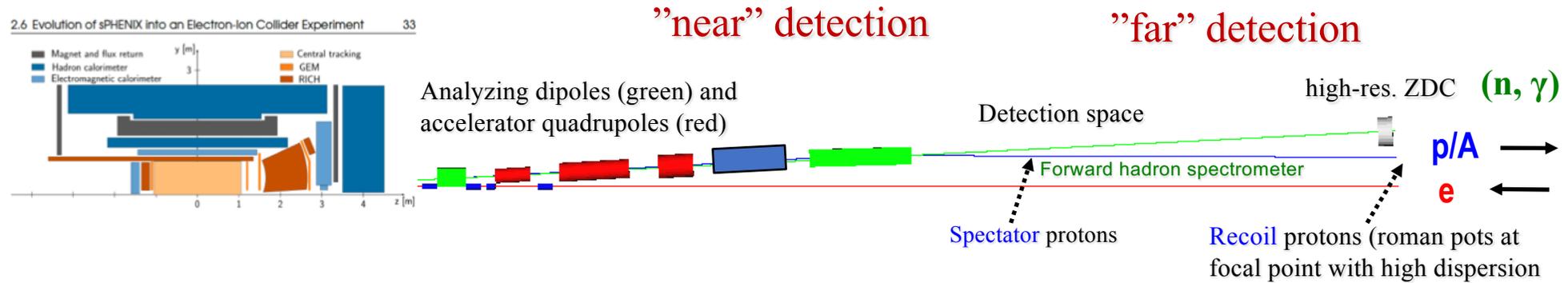
12 GeV e (5 GeV lumi max) on 200 GeV p

JLab EIC (JLEIC)



Near-beam detection could be similar

Detection of target fragments

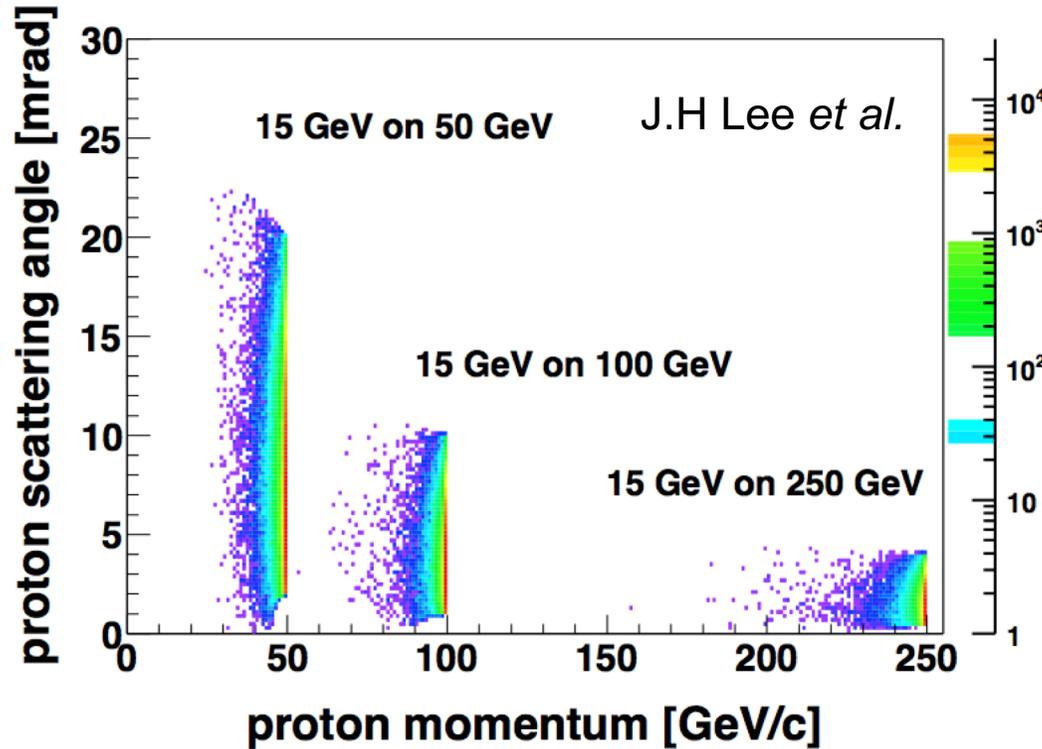


- A large part of the EIC program relies on detecting recoil baryons and target fragments
 - Spatial imaging through *exclusive* meson/photon production on the proton, and in coherent diffraction on nuclei
 - Neutron structure through spectator tagging in light nuclei
 - Various incoherent processes on heavy nuclei, **including SRC and EMC studies**
- Forward detection requirements for the EIC are very demanding
 - Need to detect particles very near the beam (down to $p_T = 0$ for $p/p_0 < 99\%$)
 - Need to detect protons with p_T up to at least 1 GeV/c and spectators with A/Z very different from that of the beam (e.g., spectator protons from deuterium)
 - Need excellent momentum resolution, detection of neutrals, and PID for ions

Forward detection – processes and requirements

- Which processes drive the near-beam (small p_T) acceptance?
 - Coherent diffraction on light nuclei (detecting recoiling ion for clean signal)
 - DIS and (in)coherent diffraction on medium- and heavy nuclei (detecting residual nucleus)
- Which processes drive the “large” angle (p_T) acceptance?
 - Tagging of spectator protons from nuclei
 - Exclusive production of photons (DVCS) and mesons (DVMP) on the proton at large t at *low energies*
 - Detection of photons and neutrons from nuclei (cone with line-of-sight)
- Which processes drive the resolution?
 - Magnetic spectrometer: tagging of protons from nuclei (spectators have $p \sim$ Fermi momentum)
 - Hcal (ZDC): tagging of neutrons from exclusive charged meson production on the proton and spectator neutron tagging (e.g., for reactions on the proton in deuterium)
 - EMcal: Photons from nuclear de-excitations (coherent diffraction and rare isotopes)

Near-beam acceptance using a magnetic spectrometer (large dispersion)

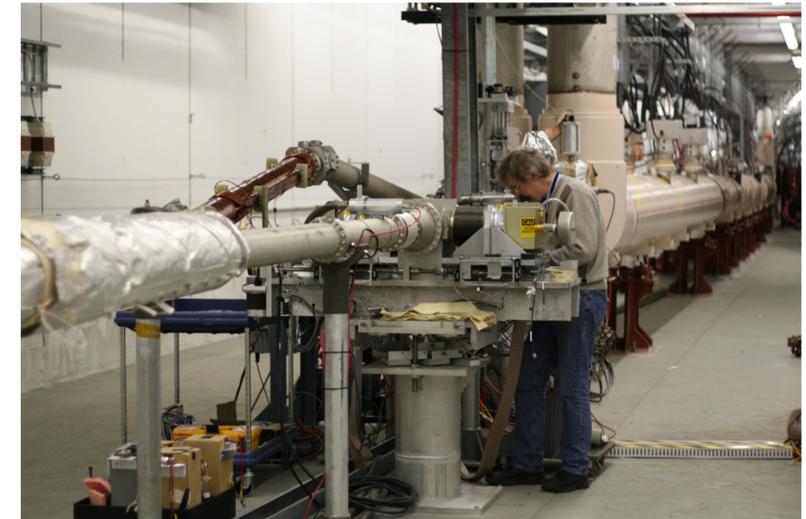


- The intrinsic beam momentum spread is on the order of a few $\times 10^{-4}$.
Insufficient to consider only the beam divergence at the IP
- The target energy loss is typically 1% (DVCS on the proton shown)
- For heavy ions, the momentum change is small, but losing a single proton changes the rigidity ($\sim A/Z$) by about 1%

- With sufficiently *large dispersion* (which separates off-momentum particles from the beam), it is possible to detect even protons scattered at zero degrees or nuclei losing one nucleons.
- Detection of intact nuclei is *much* harder than recoil protons, since at fixed t the larger mass implies a smaller change in both longitudinal and transverse momentum (scattering angle)

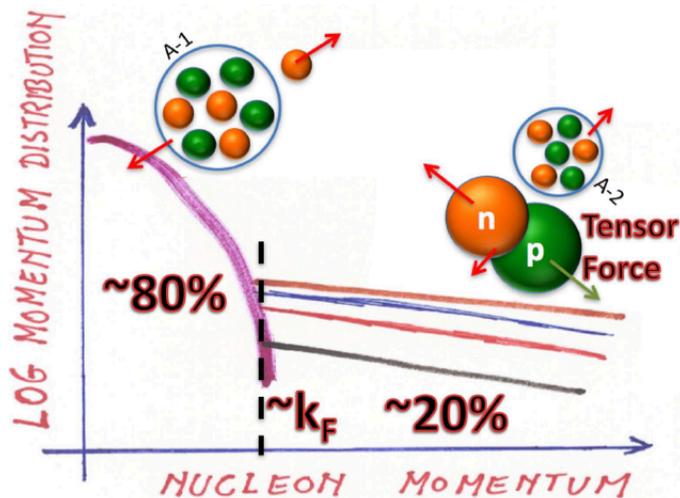
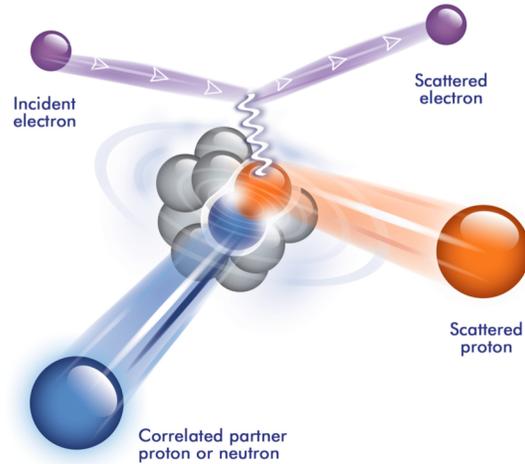
Near-beam detection at JLab and BNL

- At JLab, a full forward spectrometer was always a key part of the MEIC (now JLEIC) concept since its introduction in 2009.
 - proton acceptance of $E/E_{\text{beam}} < 99\%$ for *all* angles and down to 3 mrad for *all* energies
- At BNL, the IR currently lacks a high-dispersion section, limiting its potential for eA.
 - However, nothing prevents BNL to add a forward spectrometer with large dispersion.
 - Implementation easier at BNL since a larger e-ring allows reversing the order of electron and ion quads.
- Input on from the eA community on the detection requirements is thus important and urgent



Roman pots in STAR located ~50 m downstream from the collision point

But SRCs are not a process – so where do they fit in?



• Deuterium

- As **spectator tagging on the nucleon** *-> talk by C. Weiss*
but with sufficient acceptance for high- p_T tail
- Additional p_T for the struck nucleon will depend on process and kinematics, but can be substantial *-> talk by F. Hauenstein*

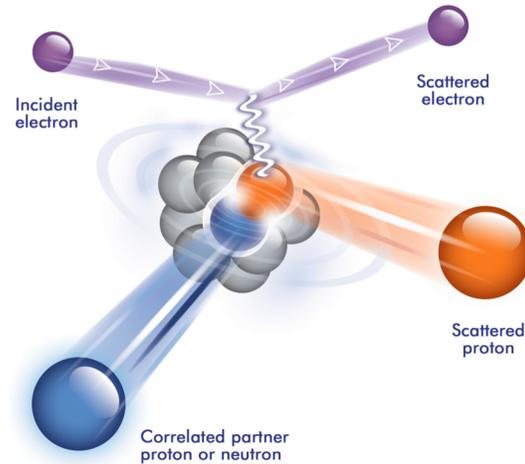
• Light ions

- $A-2$ fragment(s) easy to detect, except for $Z=A/2$
- The latter is, however, a very important case (d, ^{10}B , etc) with detection challenges similar to **coherent processes**

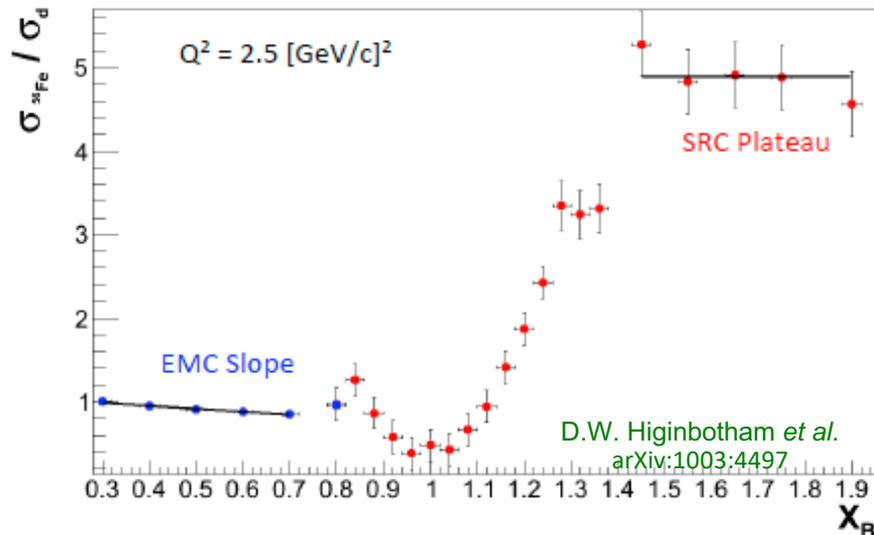
• Heavy ions

- Tagging and identification of heavy fragments needed.
- Similar to measurements of **rare isotopes at the EIC** and helpful for **coherent diffraction** on heavy ions.

Traditional SRC measurements at an EIC



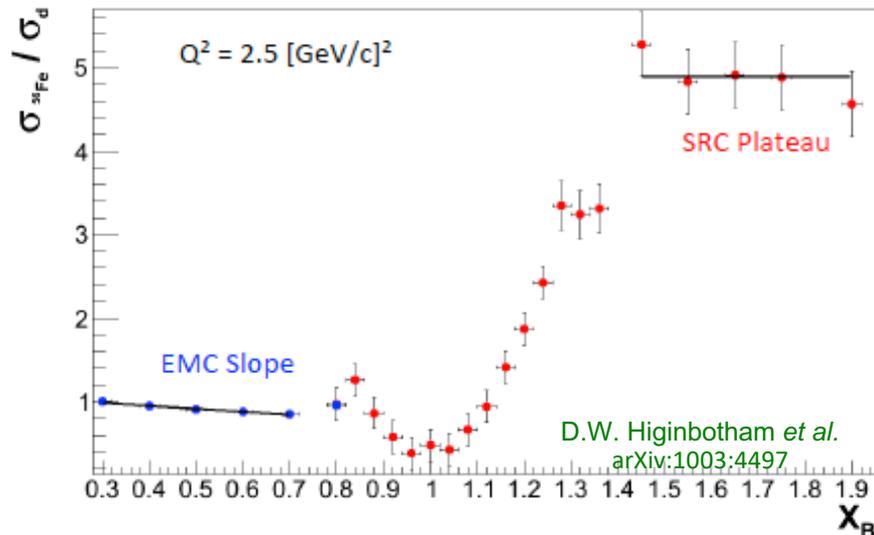
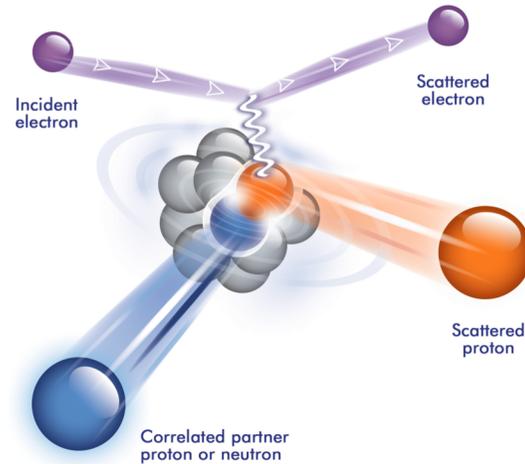
- SRCs have traditionally been studied at high x in quasielastic (QE) reactions in medium- and heavy nuclei.
 - $x > 1$ ensures participation by more than one nucleon
 - QE process provides straightforward measurement and easier-to-understand final-state interactions



- The EIC can contribute by providing a large lever arm in Q^2 at high x and detection of all target fragments – but there are some challenges.
 - Event rates at high x are much lower than for fixed target
 - Resolutions at large x deteriorate rapidly

-> talk by B. Schmookler

SRCs in DIS at lower x?



- The EMC effect ($x \sim 0.5$) is much easier to measure in “typical” EIC kinematics.
- But also for SRC studies at an EIC it would be natural to seek relevant measurements in DIS (or exclusive reactions) at lower x .
 - Can we learn something new about SRCs?
 - Is it important to introduce SRCs into nuclear models (e.g., BeAGLE) for studies of low- x processes?
- Nuclei over the full mass range will provide important input.

JLab LDRD-funded project for SRCs at an EIC

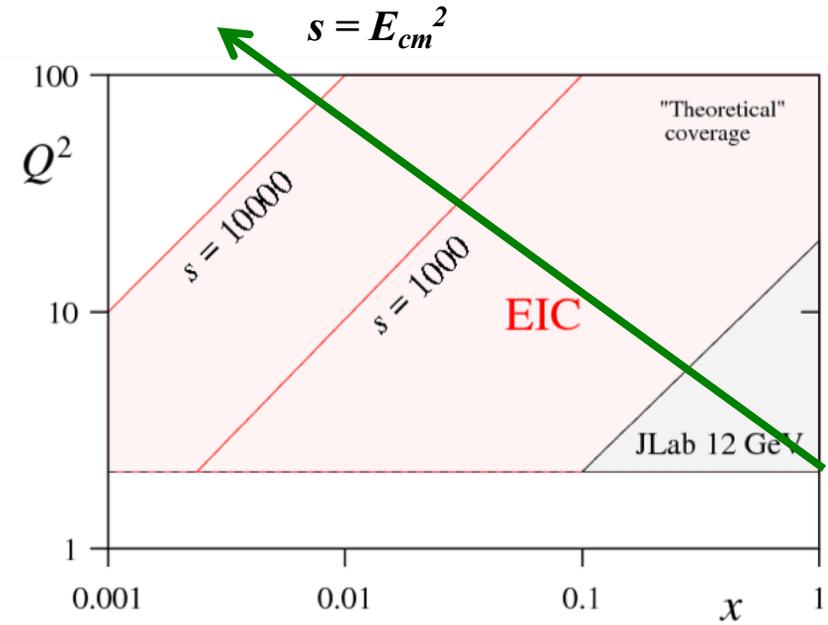
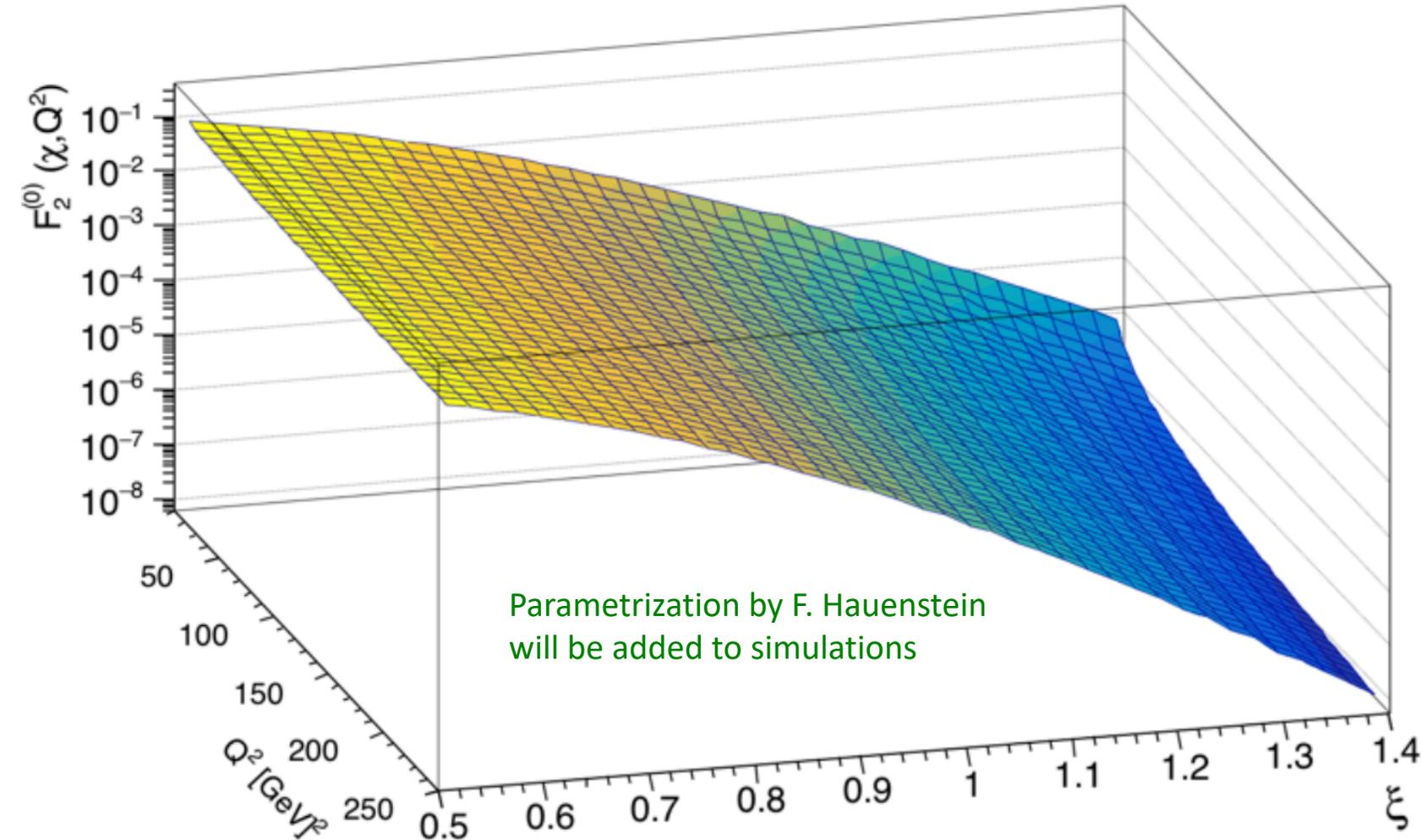
JLab “SRC” LDRD1912: D. Higinbotham, M. Baker, A. Deshpande, F. Hauenstein, O. Hen, C. Hyde, V. Morozov, PNT, A. Schimdt, B. Schmookler, Z. Tu, L. Zheng

- Feasibility studies
 - Rates at high x *F. Hauenstein*
 - Resolution in x at high x *B. Schmookler*
- Simulations, modeling, and detection requirements
 - Incorporating SRCs into BeAGLE *M. Baker*
 - Detection requirements

Rate estimate: parametrization of modified F_2

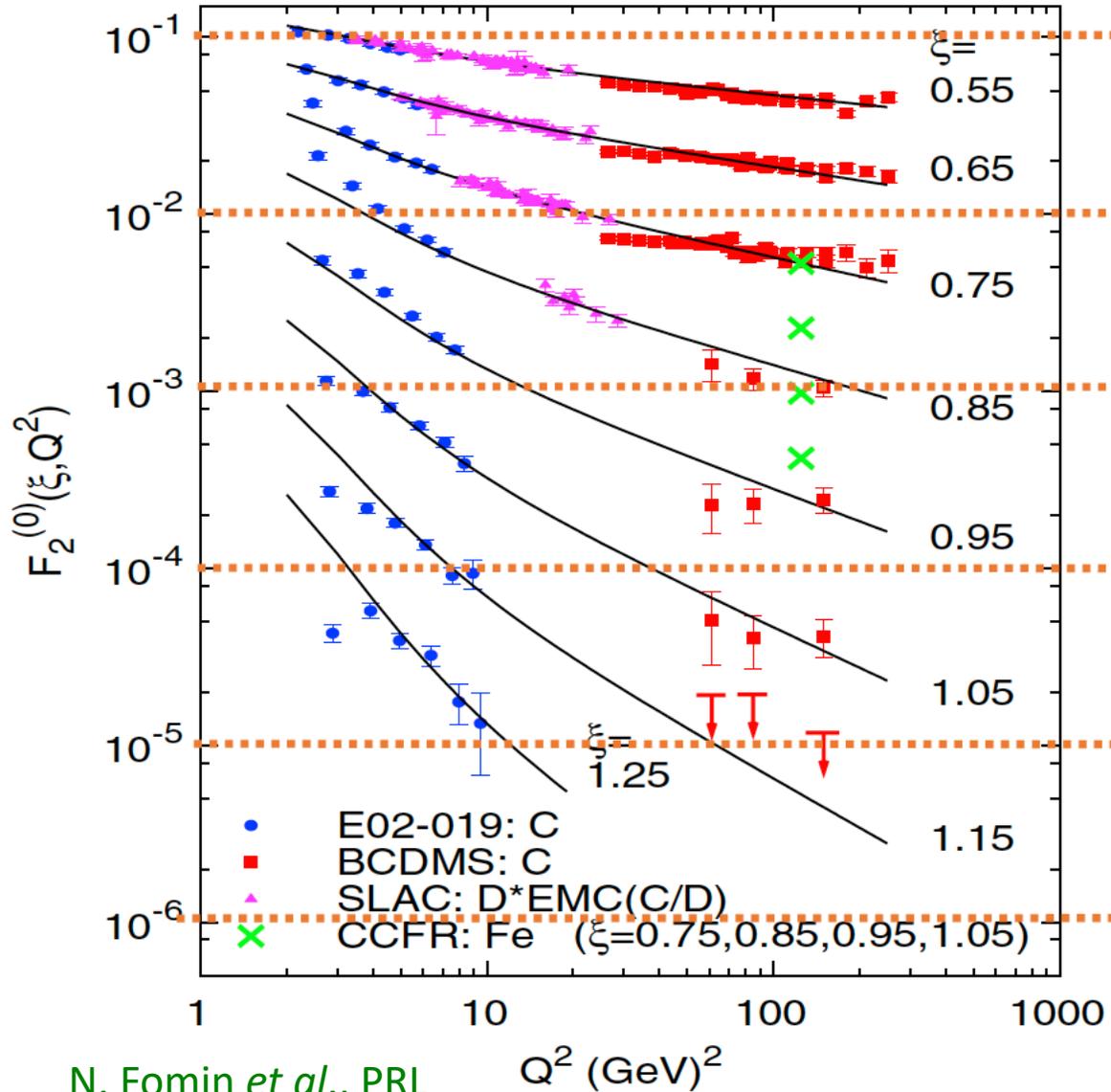
$$F_2^0(\xi, Q^2)$$

$$Q^2 \sim sxy$$

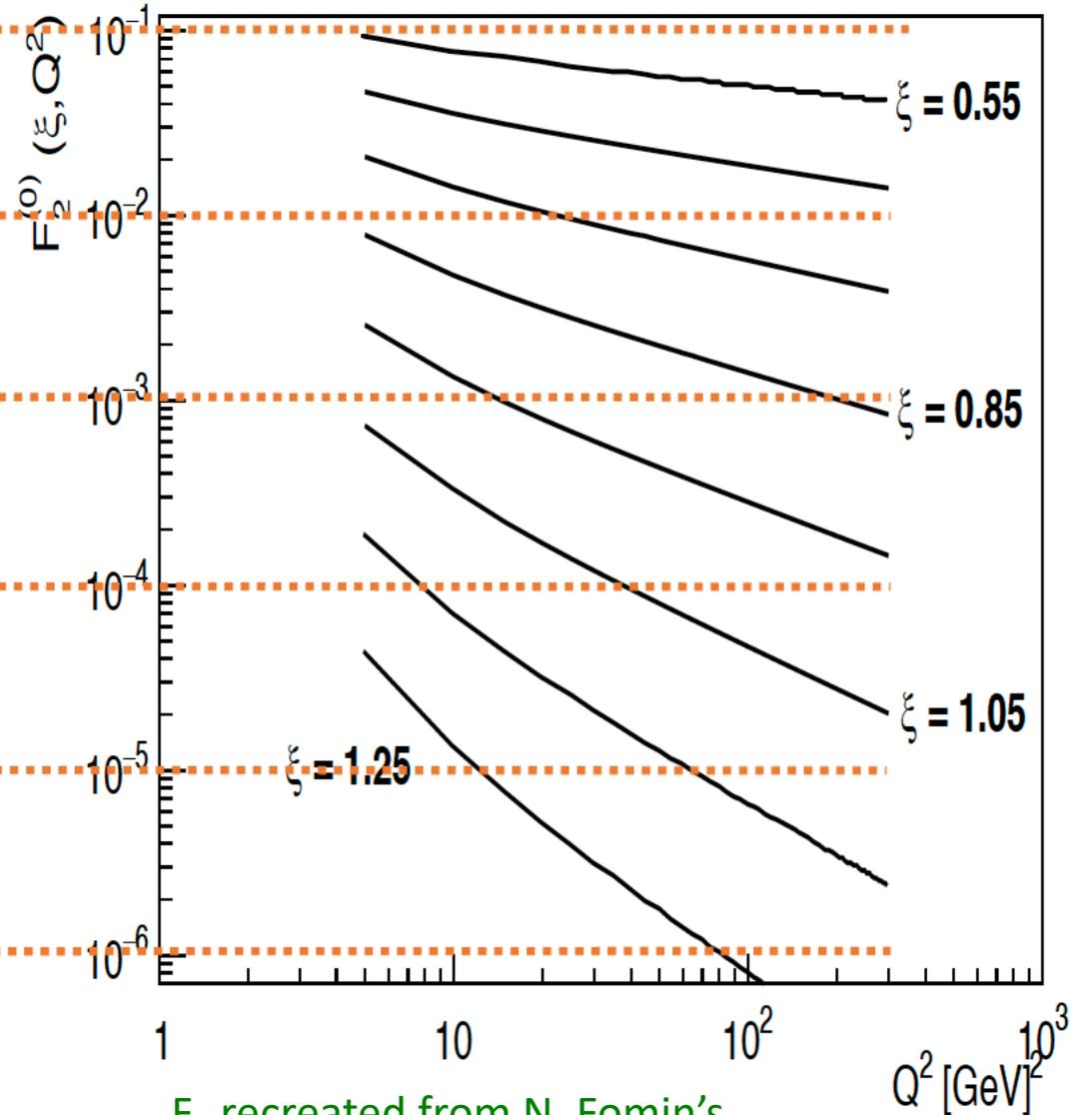


- To reach high x , the EIC would have to operate at low s and low y and/or high Q^2
 - Lower s reduces luminosity
 - Lower y reduces resolution
- At high x and Q^2 , the value of F_2 is also low.
 - Rate is challenging
 - Running conditions need to be optimized

A more detailed look at the modified F_2



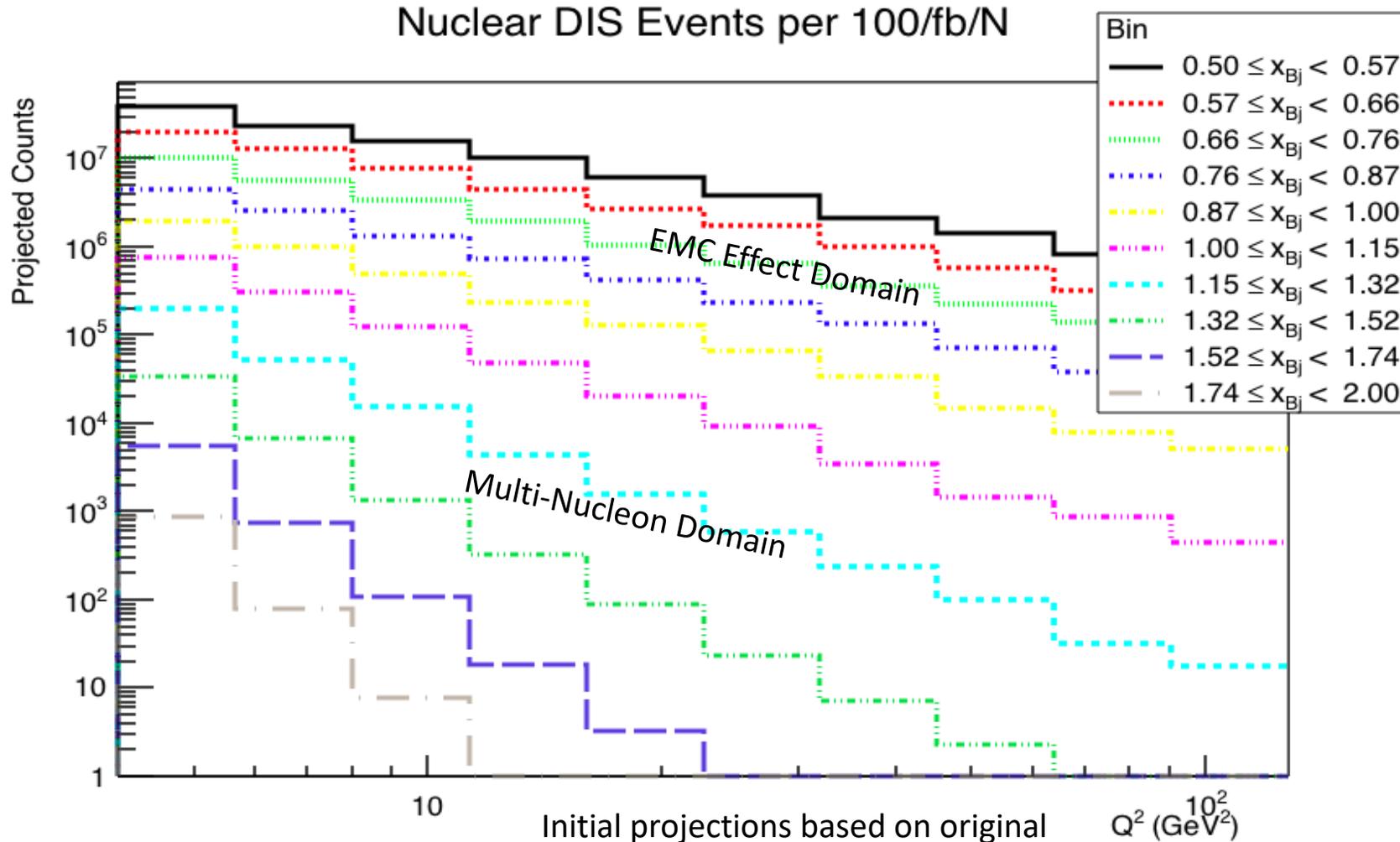
N. Fomin *et al.*, PRL
105, 212502 (2010)



F_2 recreated from N. Fomin's
scripts by F. Hauenstein

Initial rate estimates for high x (and Q^2)

5x50 GeV e+C JLEIC with 6 weeks @ 100% eff. (at lower energy rates would be lower)

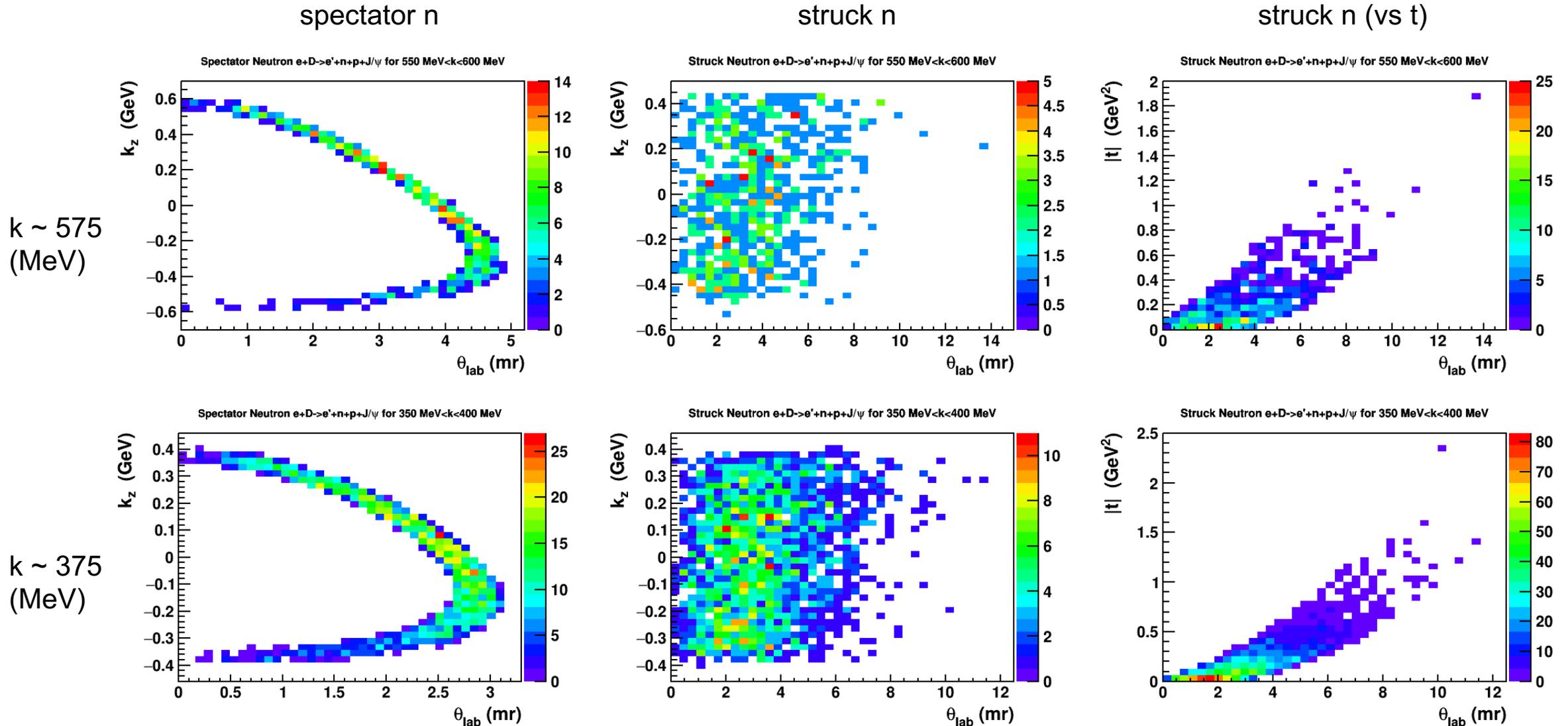


- Measuring the EMC effect at high Q^2 is straightforward

- Measuring SRCs at $x > 1$ is challenging, but there will be some events.

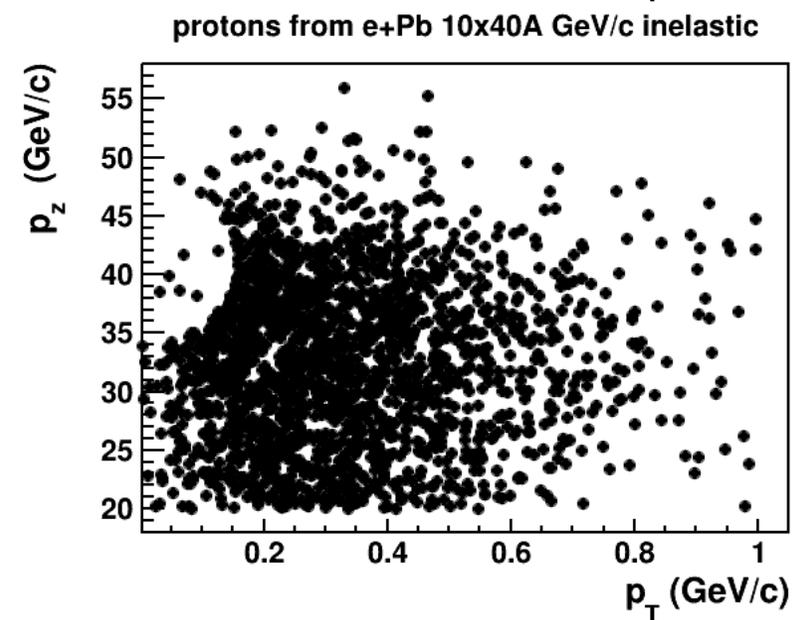
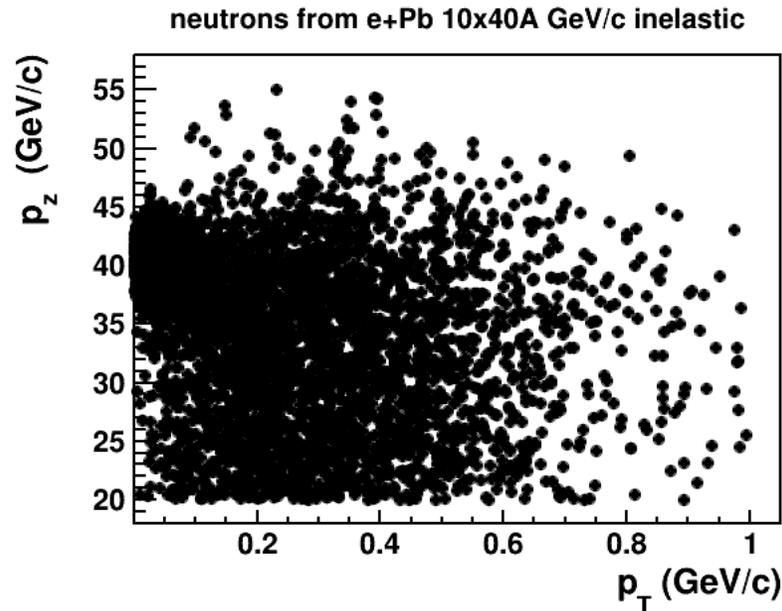
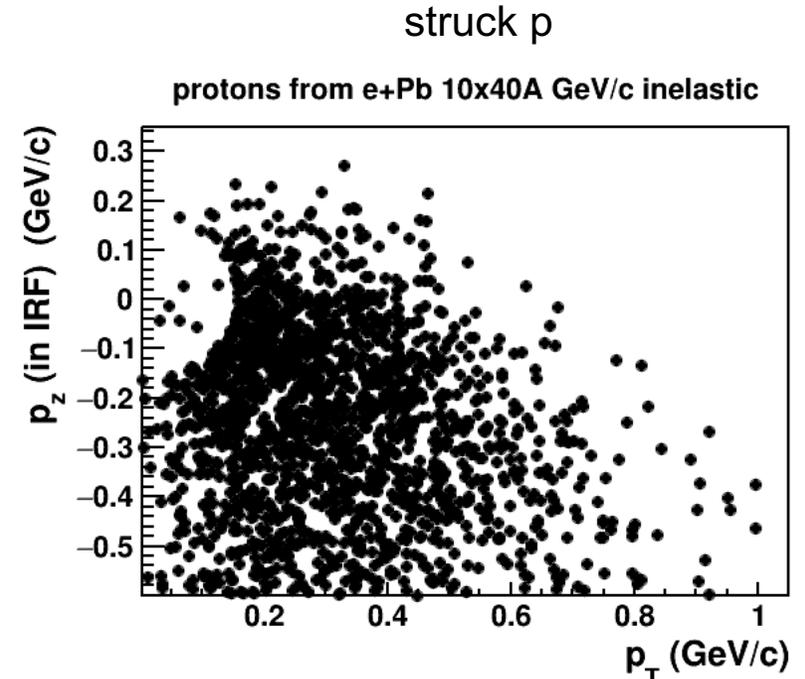
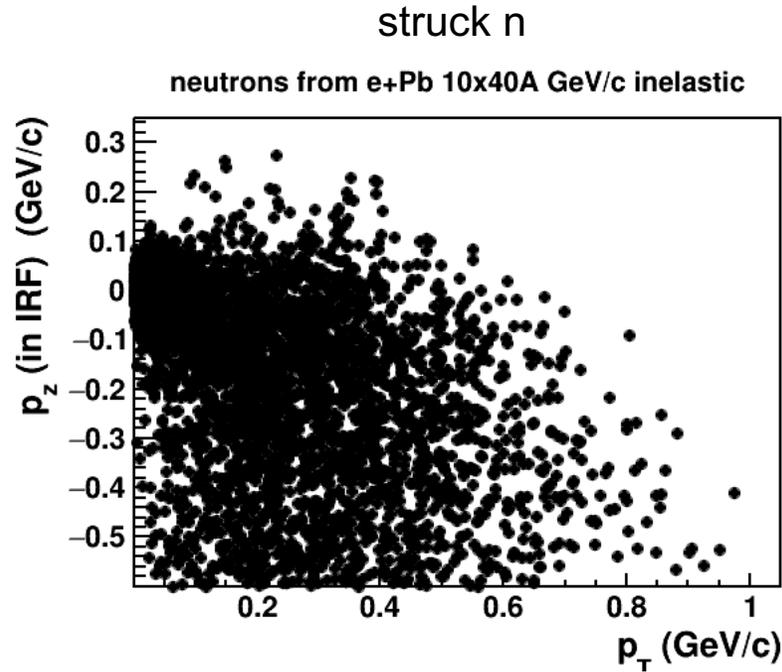
Lab kinematics of spectator and struck neutrons at high energy

M. Baker



And protons...

- For protons, Coulomb effects are not negligible.
- We need to merge the primary process with a detailed simulation of the nuclear response (BeAGLE) to understand the measurements



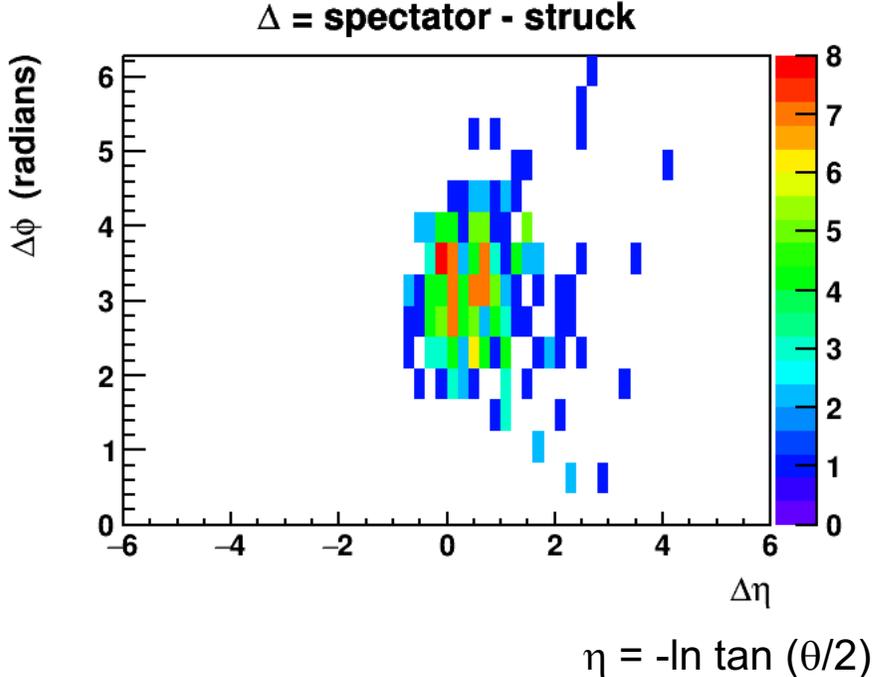
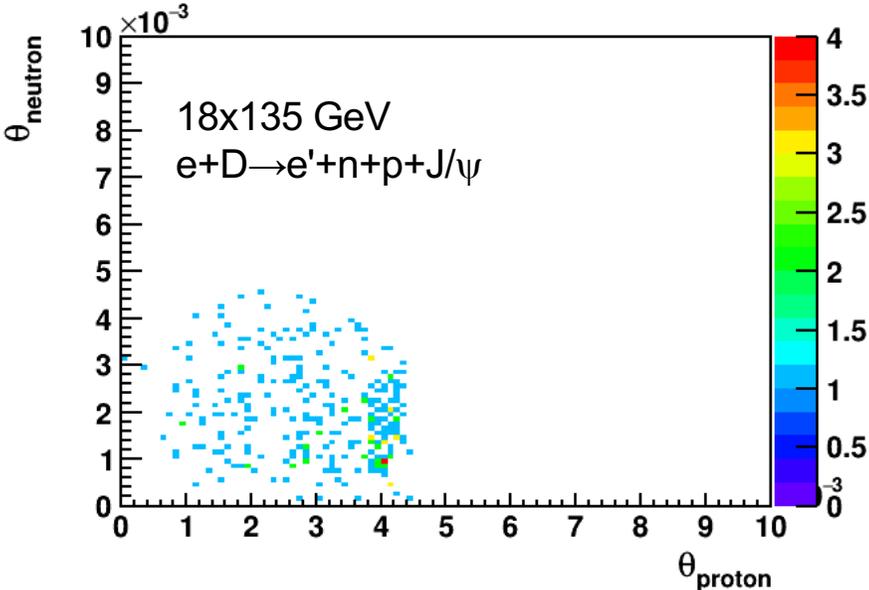
M. Baker

IRF = ion rest frame

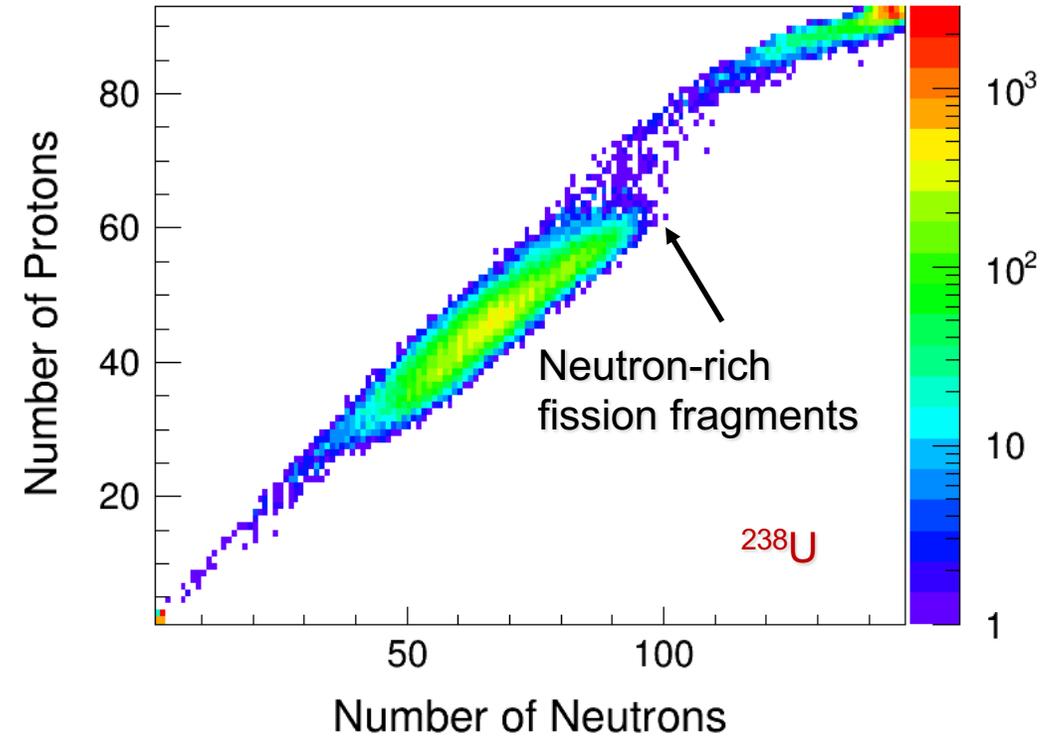
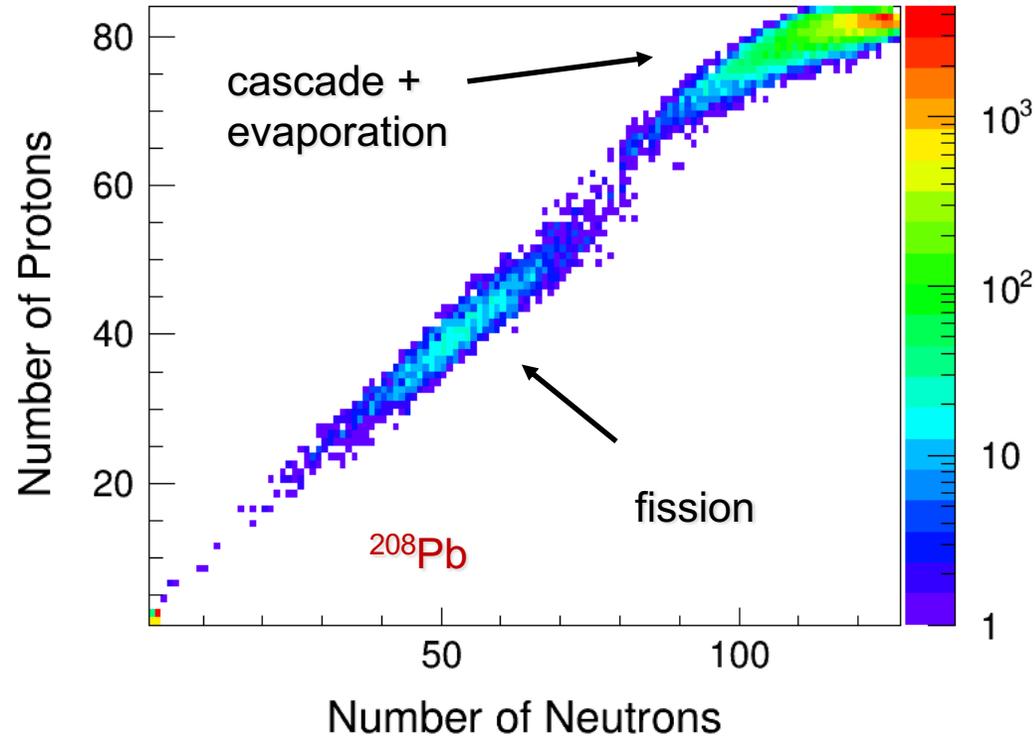
And correlations (deuteron breakup)...

struck neutron, spectator proton
 $|t| < 0.1 \text{ GeV}^2, 550 < k < 600 \text{ MeV}$

M. Baker

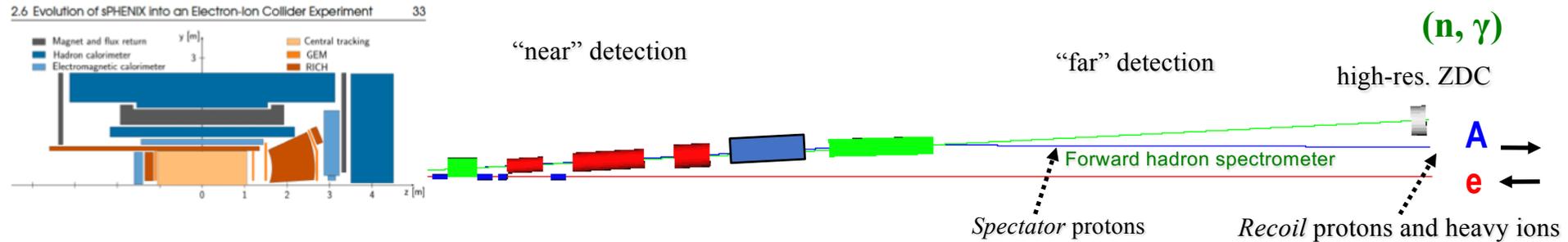


Nuclear fragments produced in eA reactions at the EIC

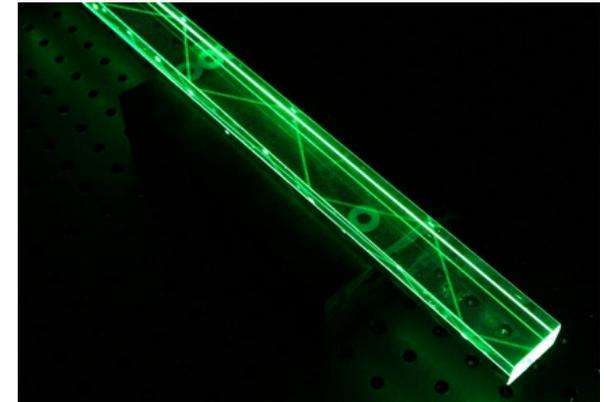


- Fragment detection (“tagging”) is needed to, for instance, reconstruct the Fermi momentum of the struck nucleon

Detection and identification of the produced nuclei

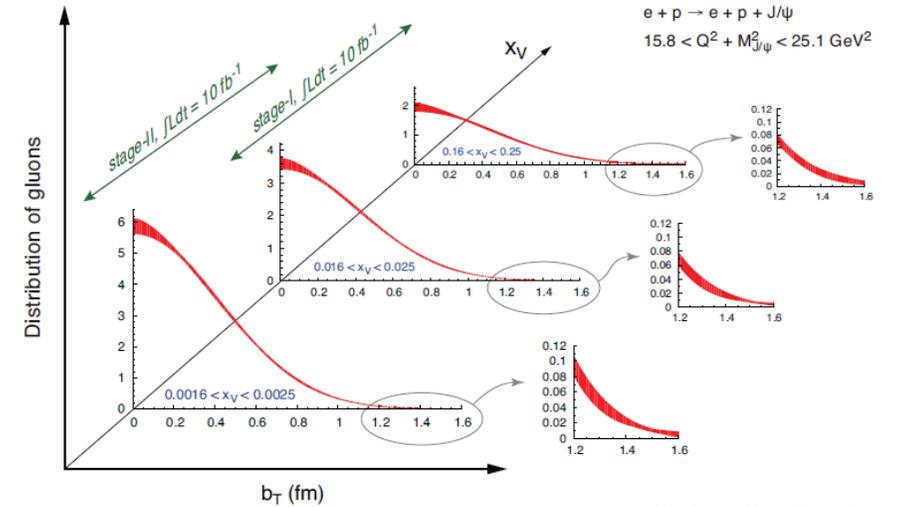
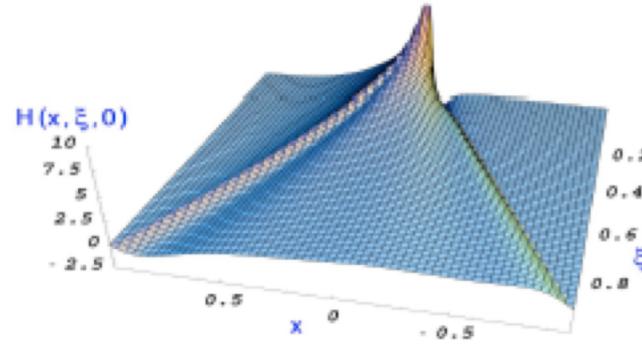
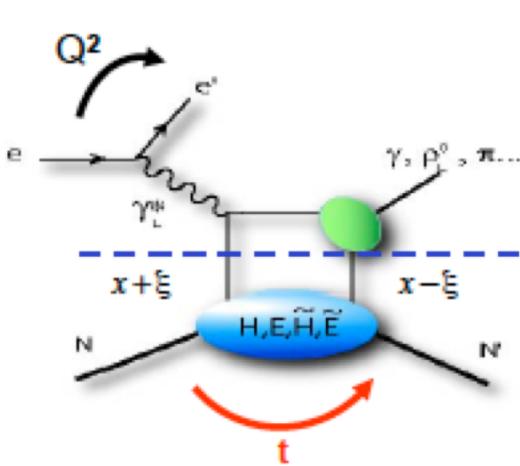


- Detection of nuclei with rigidities (A/Z) close to that the beam require a dedicated “far” spectrometer where the beam is small and dispersion large.
- But to identify the ion we need both A/Z and an independent measurement of Z . The requirement for sensitivity in Z^2 is 2% in order to identify heavy residual nuclei down to $A-1$.
- A “mini-DIRC” can produce 100,000 photons ($\ll 1\%$ error) in a few mm of fused silica.
- R&D in progress (Generic Detector R&D for an EIC program).



A “mini-DIRC” inside a Roman pot at the downstream focus can identify ions to $\sim 1\%$ in Z^2

Generalized Parton Distributions (GPDs) and 3D spatial imaging



-> talk by S. Fazio

- Exclusive production of a photon or meson is sensitive to GPDs
 - In the limit of small “skewness” ξ , the Fourier transform of t , the four-momentum transfer to the nucleon, can be interpreted as a spatial image in impact-parameter space at different values of x
- To create a spatial image in b , data over a sufficient range in t ($\sim p_T^2$) are needed ($\sim 0-1 \text{ GeV}^2$)
 - Small $t \Leftrightarrow$ large b and vice versa.
 - For nuclei with $A > 1$, that scattering angle is much smaller for a given t , cutting the low t acceptance

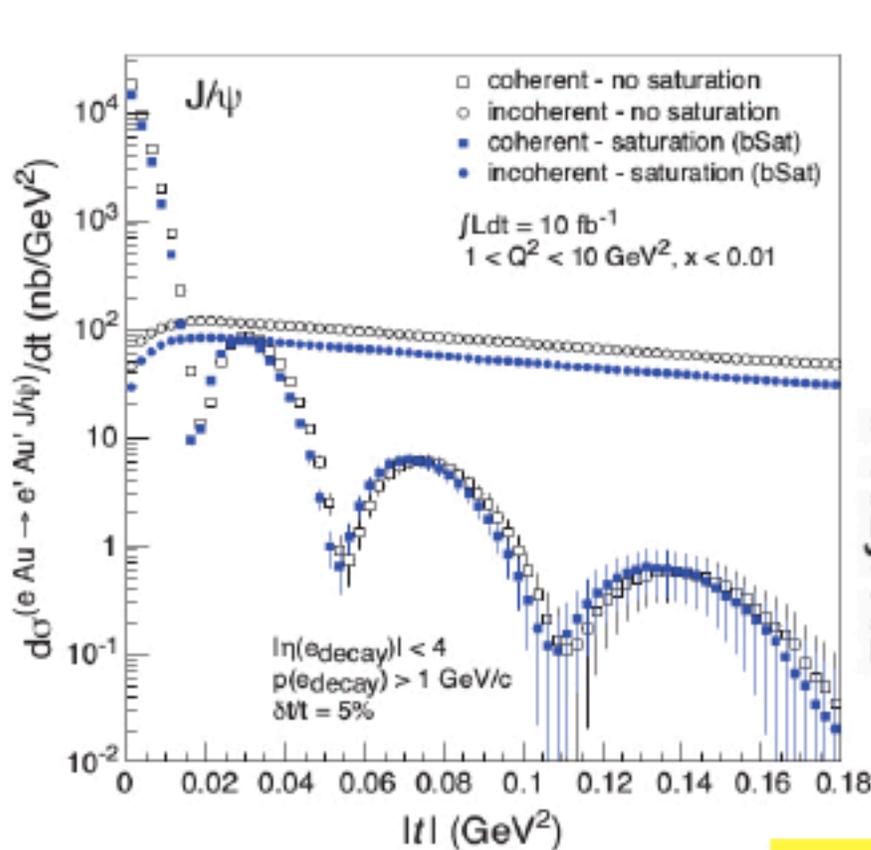
Coherent exclusive (diffractive) reactions on light nuclei



- Coherent exclusive reactions on light ions give *unique* access to the 3D structure of *nuclei*.
- In contrast to heavy nuclei, scattered light nuclei can be detected
 - The momentum transfer t can be determined directly from the ion
- Light ions span the full range of nuclear densities
 - D is the least dense nucleus unbound, while He-4 is comparable to heavy ions, and He-3 falls in the middle
- Polarized He-3 beams will allow for simultaneous measurement of both tagged neutron structure and coherent diffraction on He-3
 - Interesting comparison since spin of He-3 is dominated by the neutron
- He-4 has only one GPD and large asymmetries
 - Easy to measure and interpret

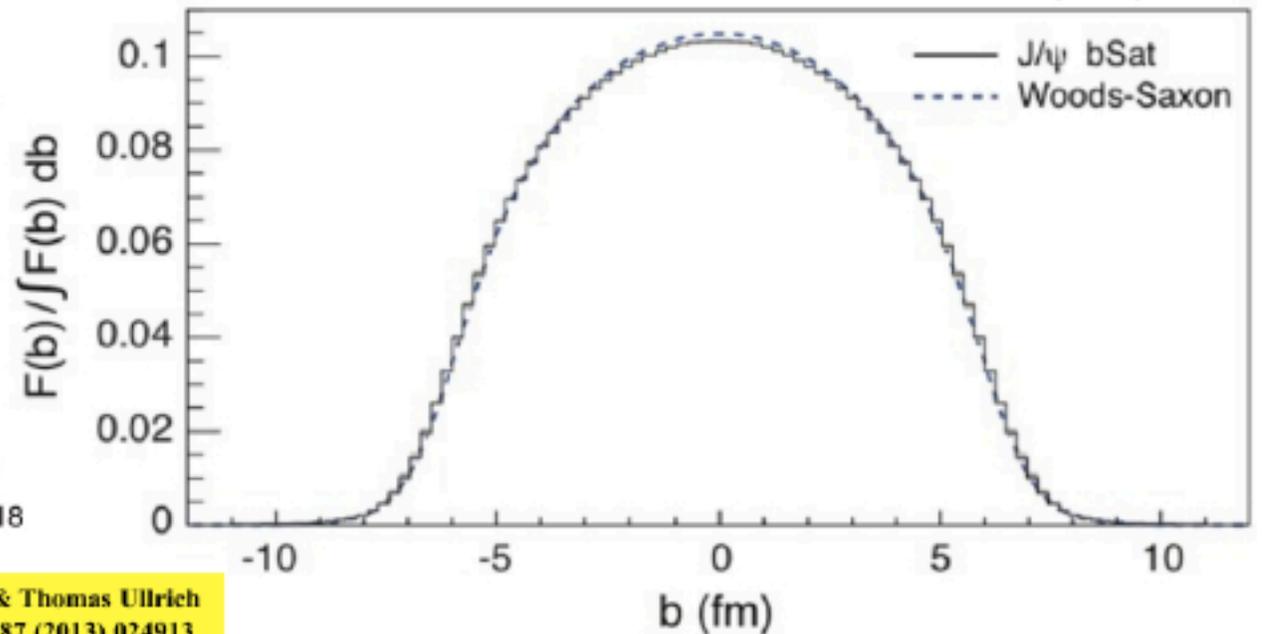
2D spatial gluon imaging in heavy nuclei through coherent diffraction

Momentum transfer t conjugate to transverse coordinate b



$$\rightarrow F(b) \propto \int d\Delta \Delta J_0(\Delta b) \sqrt{\frac{d\sigma}{dt}}$$

$$\downarrow t = \Delta^2/(1-x) \approx \Delta^2$$

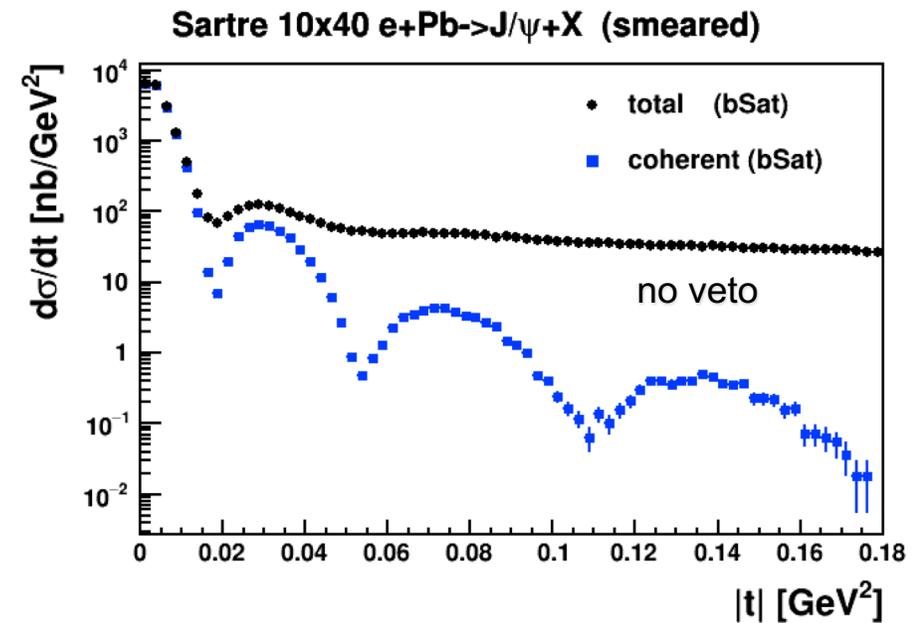


T.T. & Thomas Ullrich
PRC 87 (2013) 024913

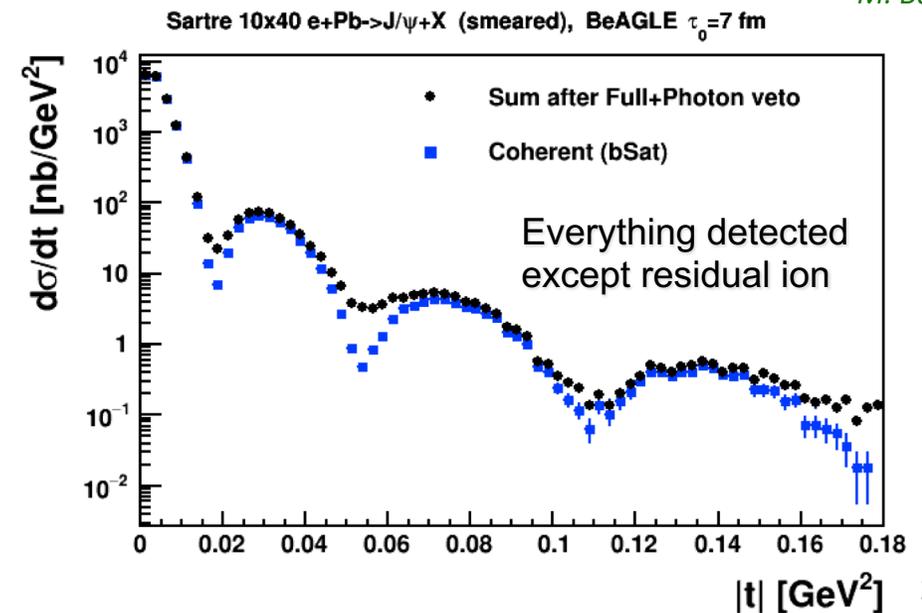
- Experimental challenge: veto the large incoherent background.

Evaluate the veto efficiency of target fragment detection

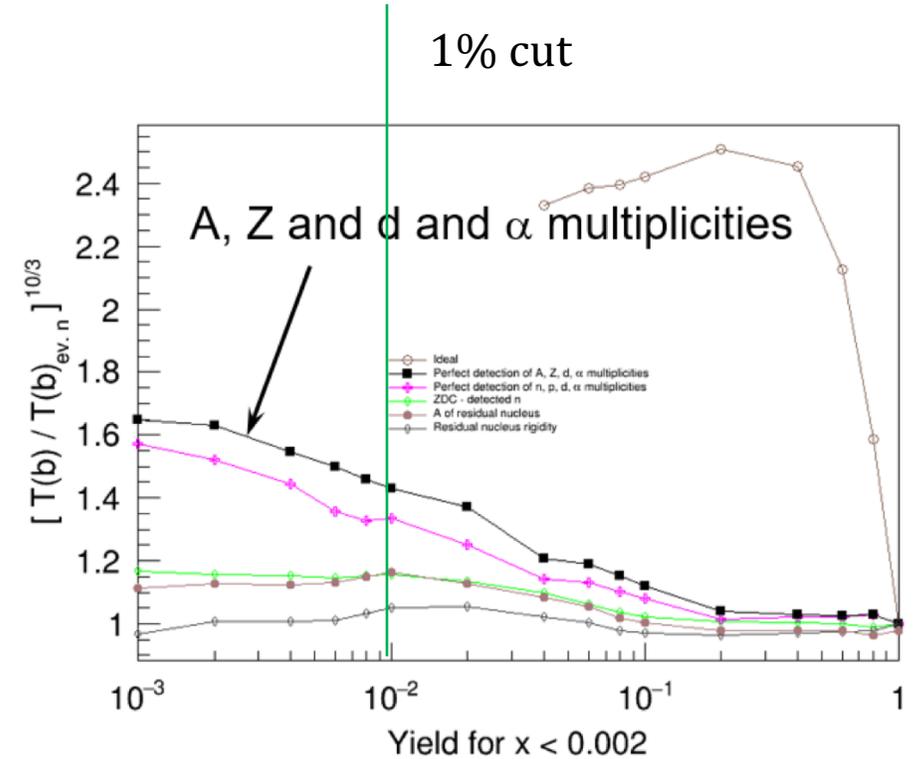
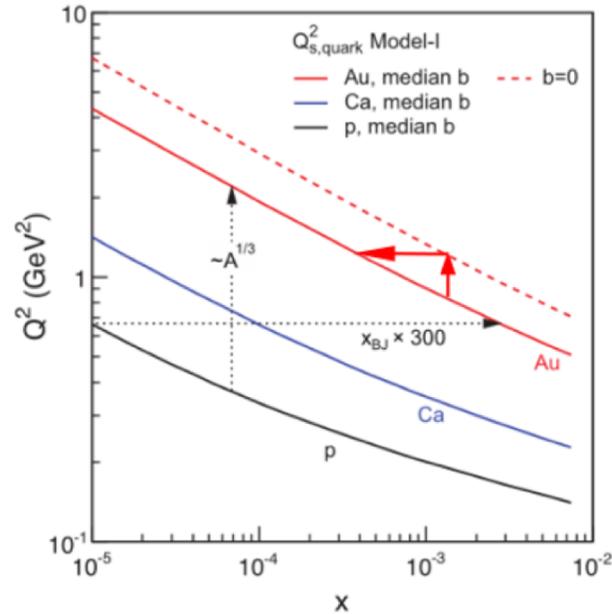
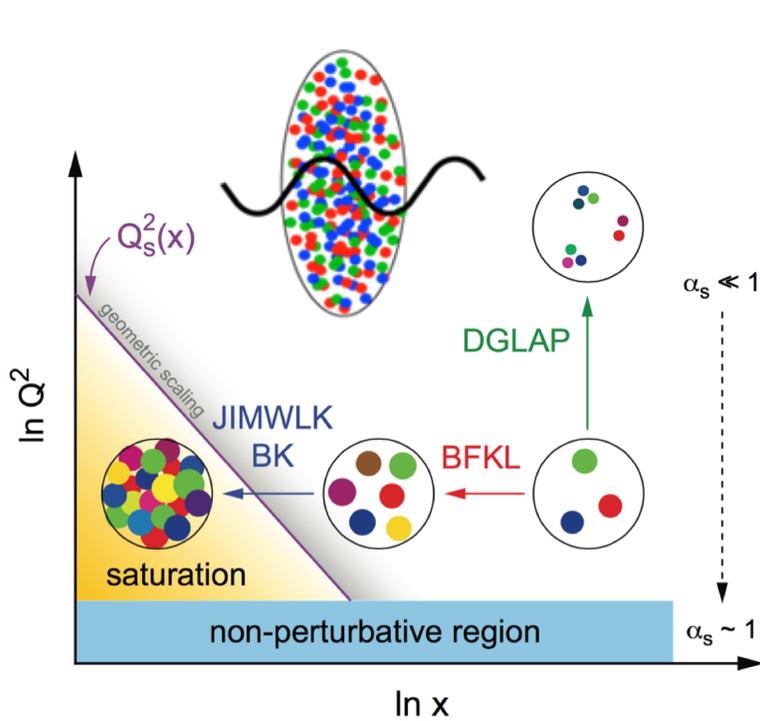
- Intact heavy ions cannot be detected directly
 - The momentum transfer is too small to kick intact heavy ions out of the beam envelope
- Need high-efficiency veto of incoherent events
 - Detection of *all* produced particles is required: protons, neutrons, light nuclei, and photons from nuclear de-excitations *-> talk by C. Hyde*
 - Detection of the residual nucleus (from incoherent diffraction) is also helpful
- BeAGLE was used to simulate the incoherent part – in this case the background
 - The coherent simulation used Sartre



M. Baker



Kinematic reach for probing gluon saturation in DIS



- At low x , the photon interacts coherently with the gluons along its path inside a nucleus, allowing the EIC to reach into the saturation regime.
 - Impact parameter tagging selects events with larger average density $T(b)$
- A larger thickness $T(b)$ is equivalent to a higher beam energy
 - With a evaporation neutron multiplicities only, this factor is 3.2 at 1%
 - Detecting all final-state particles, this increases to 4.8

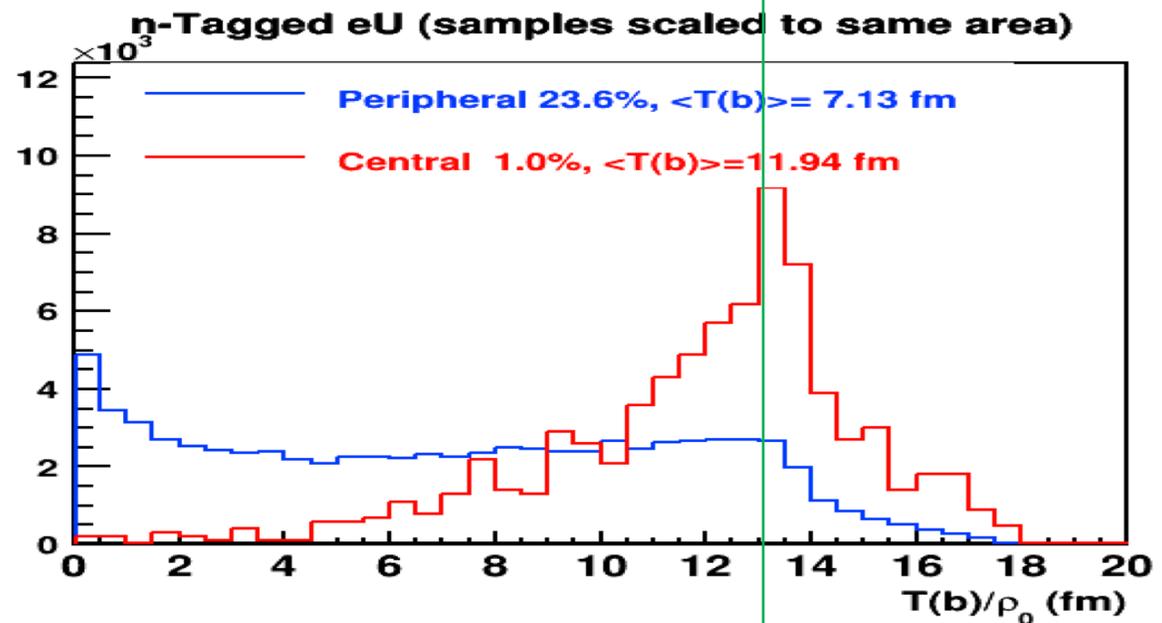
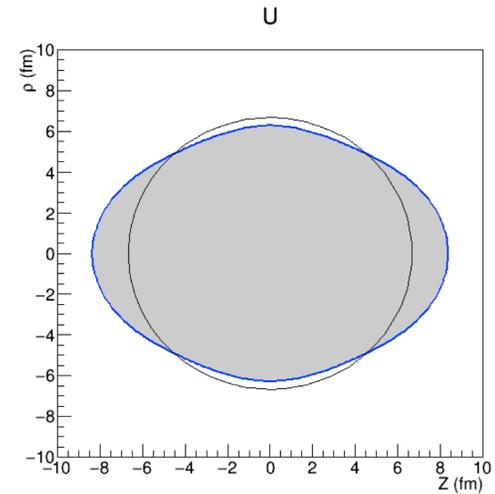
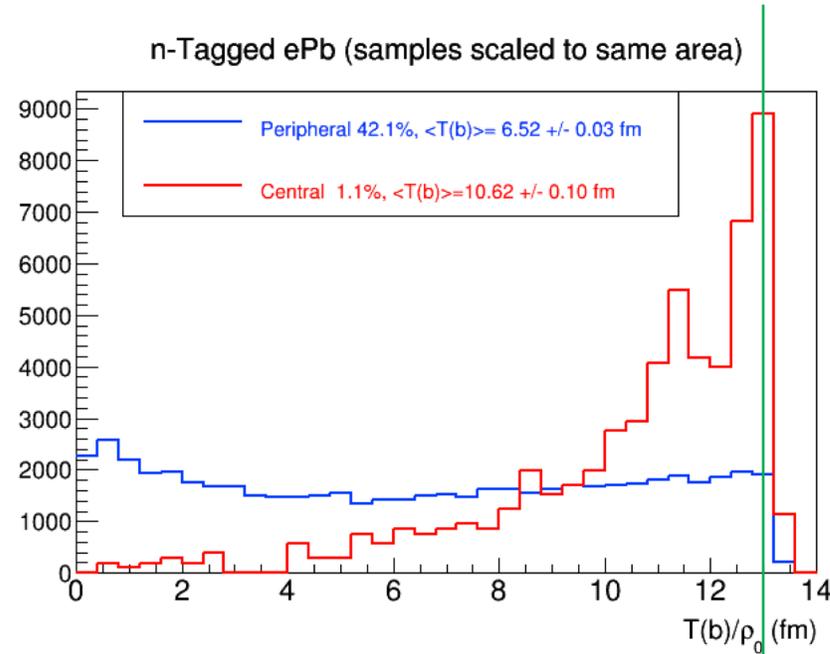
$$Q_s^2 \sim A^{1/3} / x^{0.3} \sim T(b) * (E_e E_A)^{0.3}$$

$$F_E = \left(\frac{\langle T(b) \rangle_{central}}{\langle T(b) \rangle_{minbias}} \right)^{10/3}$$

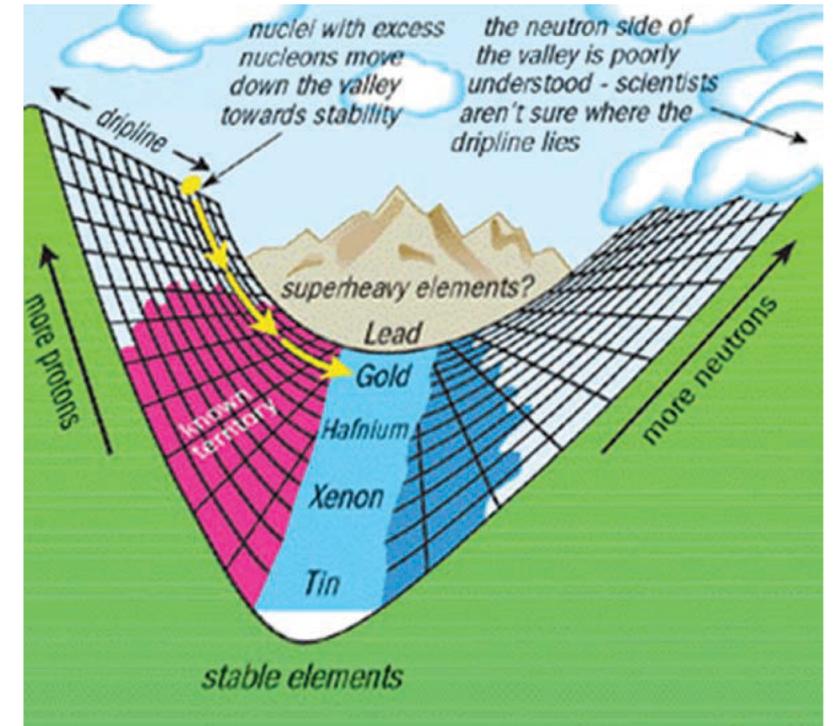
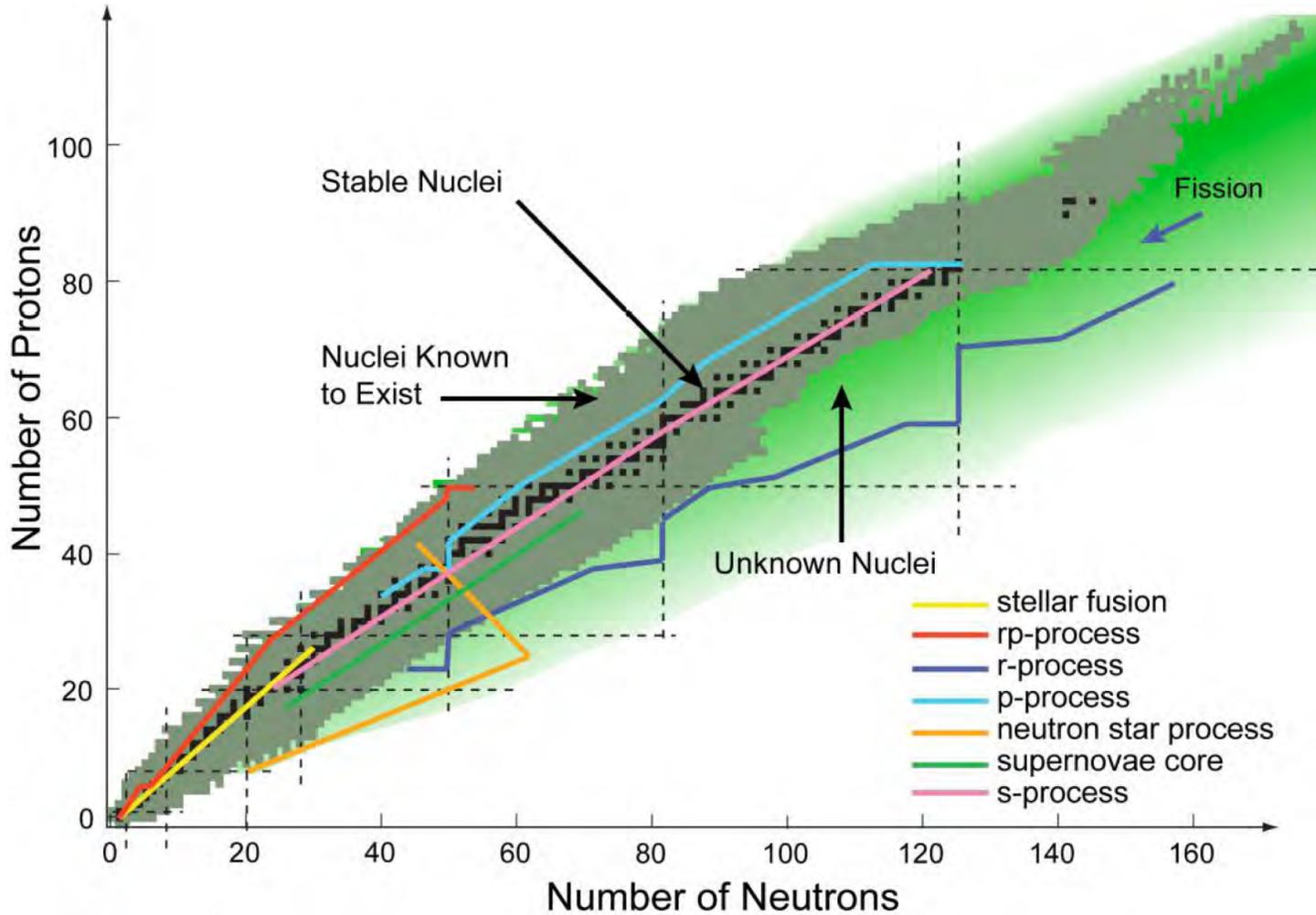
T(b) for ^{208}Pb and ^{238}U

- Using U-238 instead of Au or Pb further increases the equivalent energy, from a factor 4.8 to 6.7 at 1% yield.
 - U has slightly larger A than Pb
 - U is a deformed nucleus. Longer paths along the long axis
 - It is easier to align one axis in eA than two in AA
- A factor 6.7 is equivalent to running untagged Pb at 740 GeV/A
- Note that detection of all nuclear fragments, including spectator protons, is crucial

M. Baker, 2018



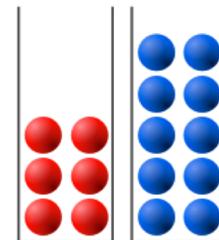
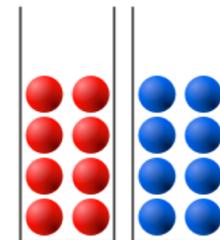
Rare isotopes at an EIC



$A = 16$

Lower energy

Higher energy



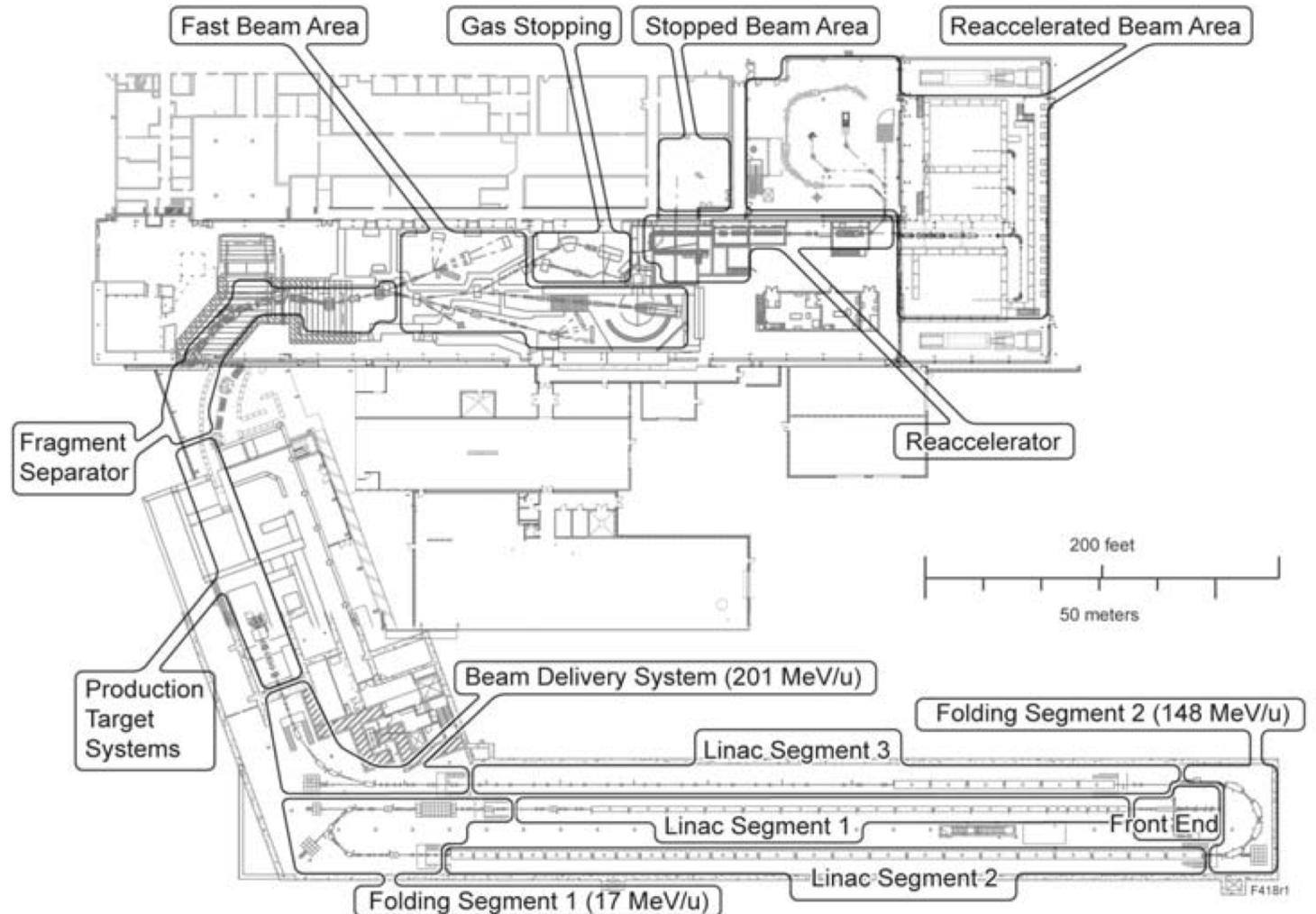
$|N - Z| = 0$

$|N - Z| = 4$

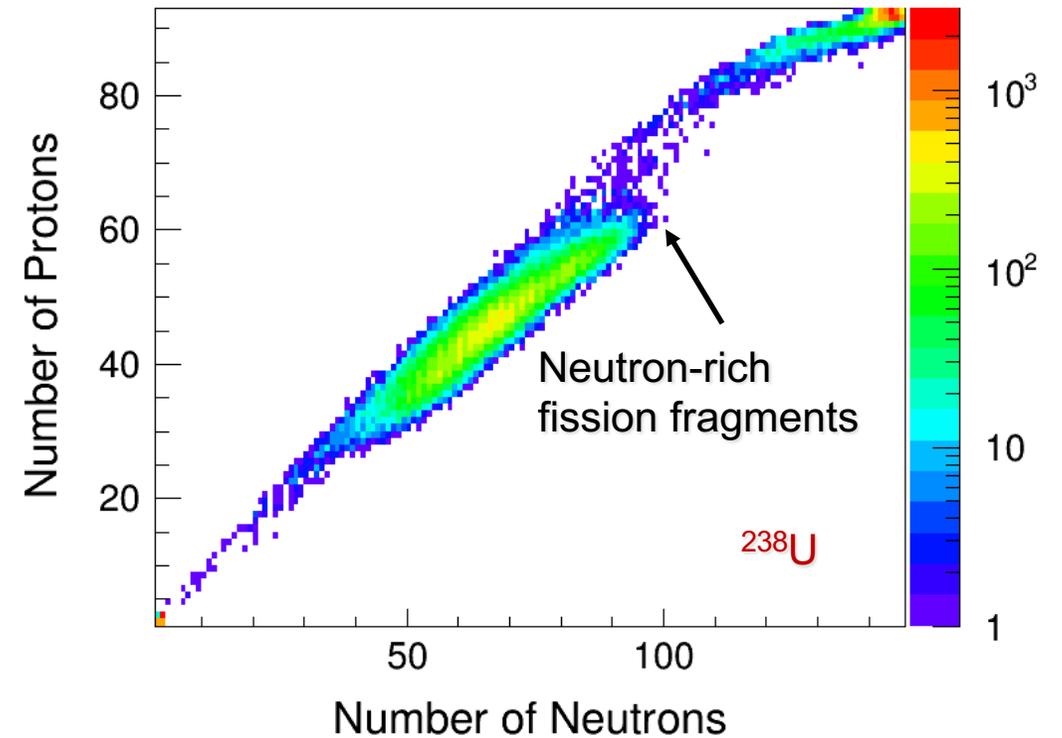
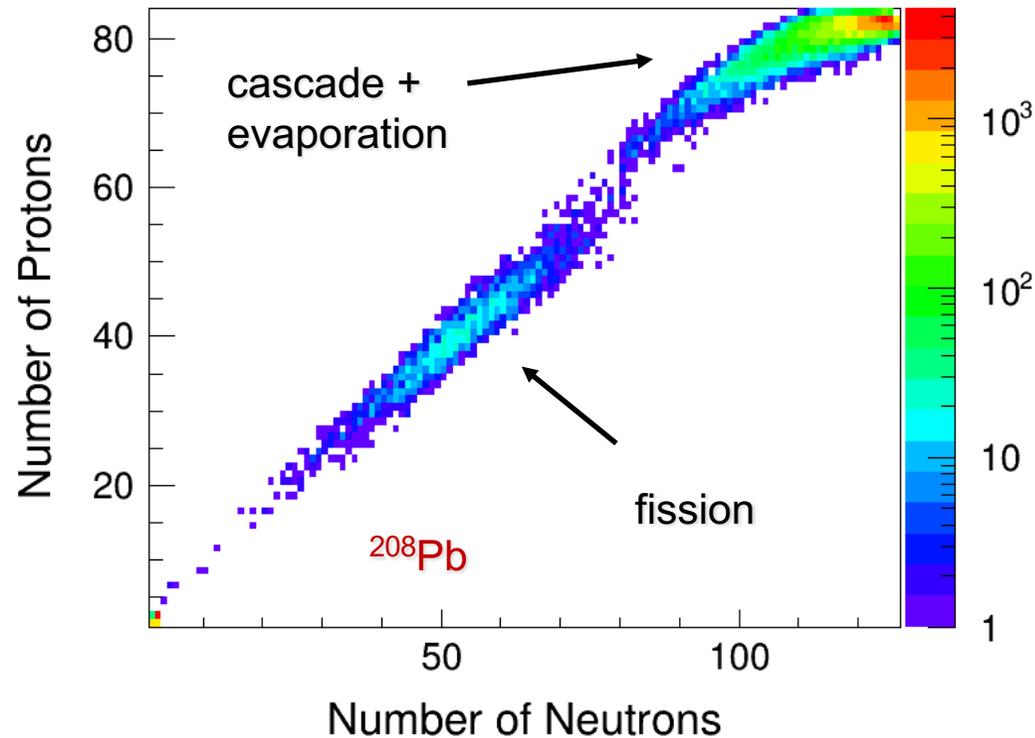
Facility for Rare Isotope Beams (FRIB)



- Close to completion at MSU
- Will produce radioactive beams through in-flight projectile fragmentation followed by fragment separation in a downstream spectrometer
- Focus in on neutron-rich nuclei

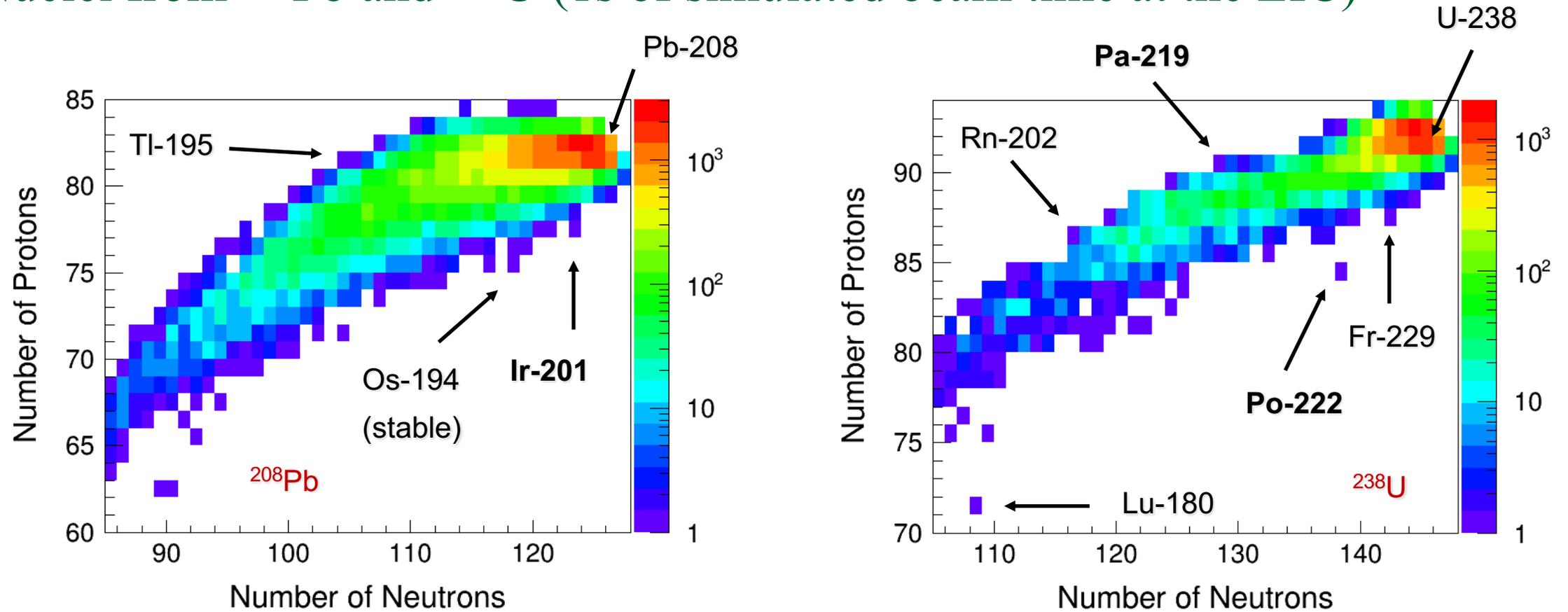


Nuclei from ^{208}Pb and ^{238}U (1s of simulated beam time at the EIC)



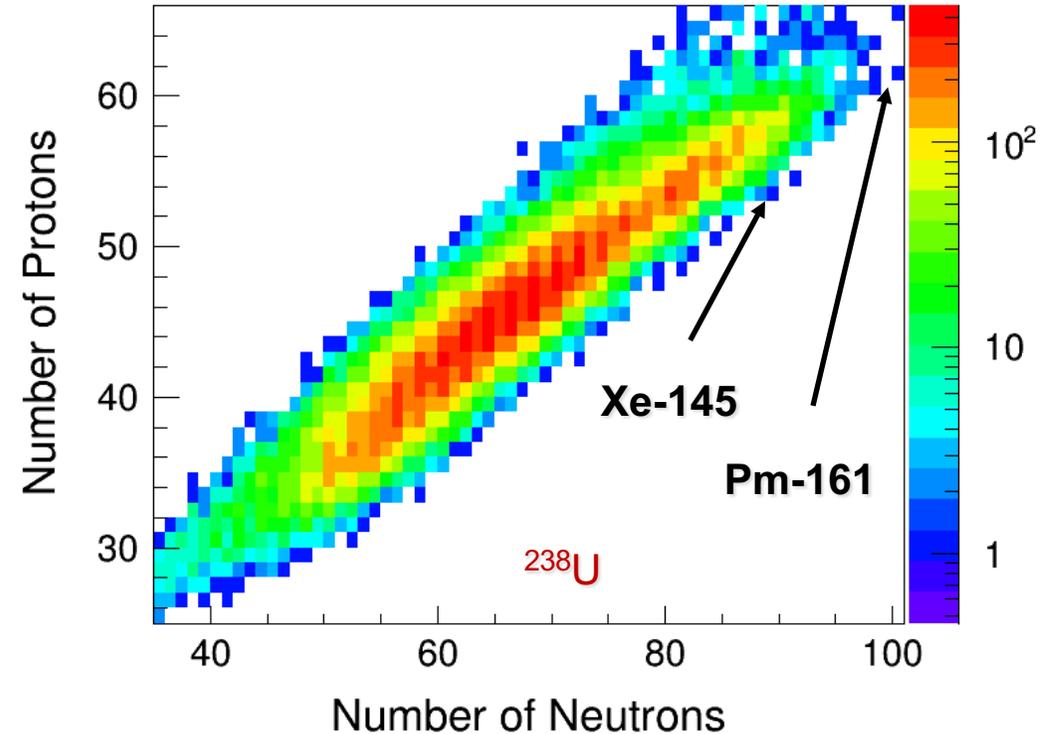
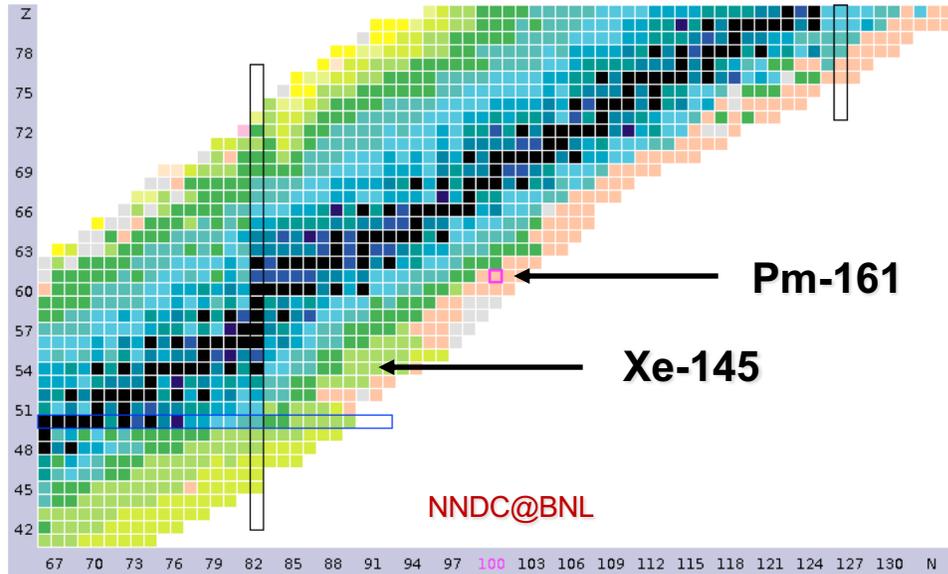
- ^{208}Pb (left) produces mainly heavy isotopes from evaporation
- ^{238}U (right) produces fewer, but heavier isotopes from evaporation. It also produces very neutron-rich fission fragments (medium-mass nuclei have fewer neutrons).

Nuclei from ^{208}Pb and ^{238}U (1s of simulated beam time at the EIC)



- ^{208}Pb (left) produces mainly heavy isotopes from evaporation
- ^{238}U (right) produces fewer, but heavier isotopes from evaporation. It also produces very neutron-rich fission fragments (medium-mass nuclei have fewer neutrons).

How “exotic” are ^{161}Pm and ^{145}Xe (produced in 1s of simulated beam time)?



- Lifetime for ^{161}Pm : >130 ns
- Lifetime for ^{145}Xe : 188 ms
- Both nuclei are very close to the most neutron-rich nuclei currently known – and in a spot where rates at FRIB are relatively low

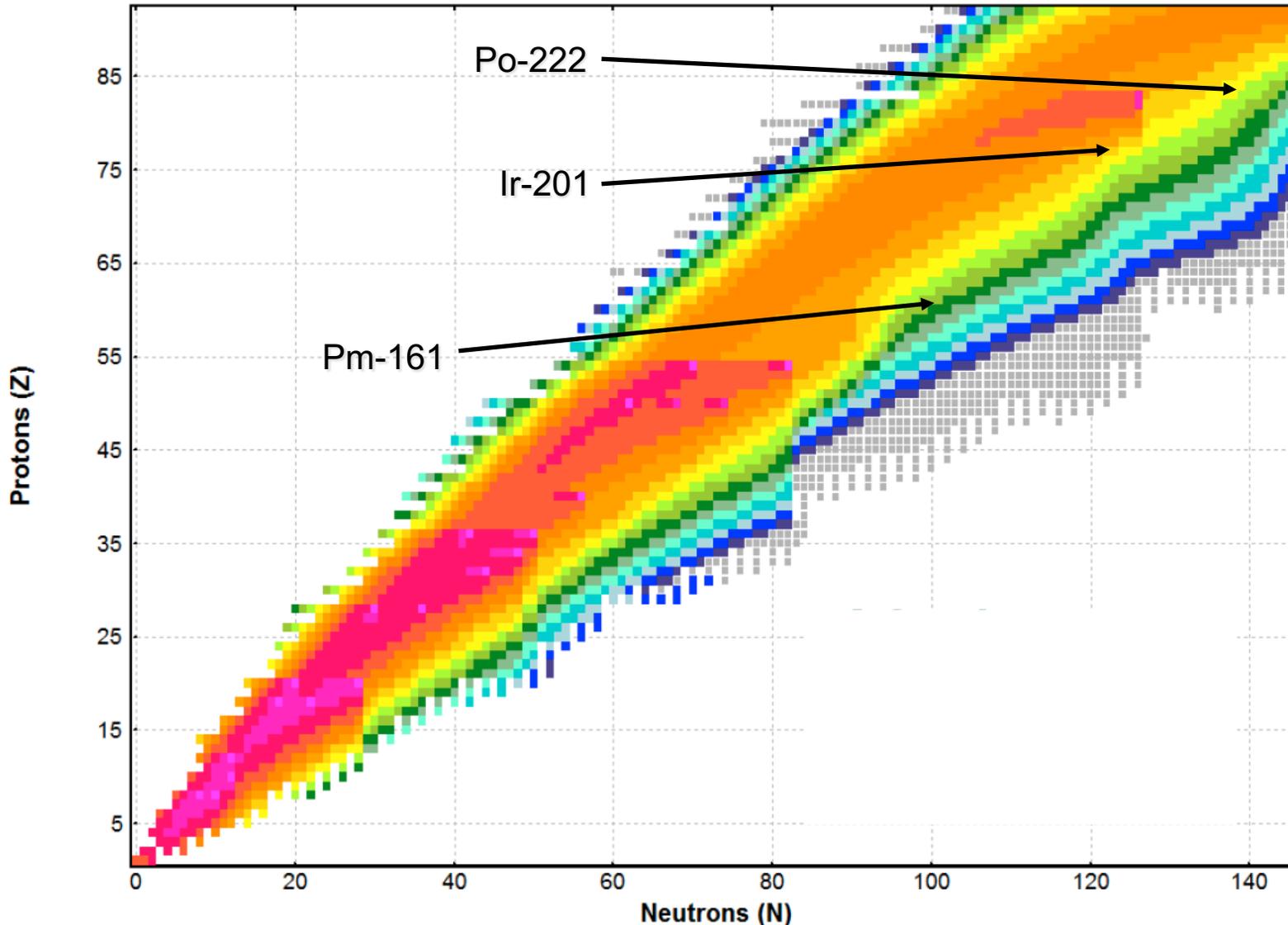
Rates at the EIC?

NSCL PAC35 rates (v.1.03)

https://groups.nsl.ms.edu/frib/rates/nsl_pac35_rates.html The rates are estimated based on

the EPAX 2.15 cross section parameterization for fragmentation and the LISE++ 3EER model for in-flight fission.

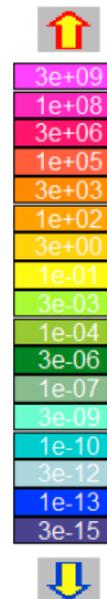
Primary beam intensities and energies have been used from the PAC35 beam list



Needs further studies!

A bold extrapolation

- Let's assume that we have one "yellow-orange" event per second, or 10^8 in a year, distributed over ~ 100 isotopes of interest
- Further, let's assume that we can use the FRIB rate estimates for the extrapolation.



- Then, if we want to accumulate a total of 10,000 events for the isotopes of interest, in a year we can move from the orange-yellow to light green (although for ^{161}Pm we are there in 1 minute)

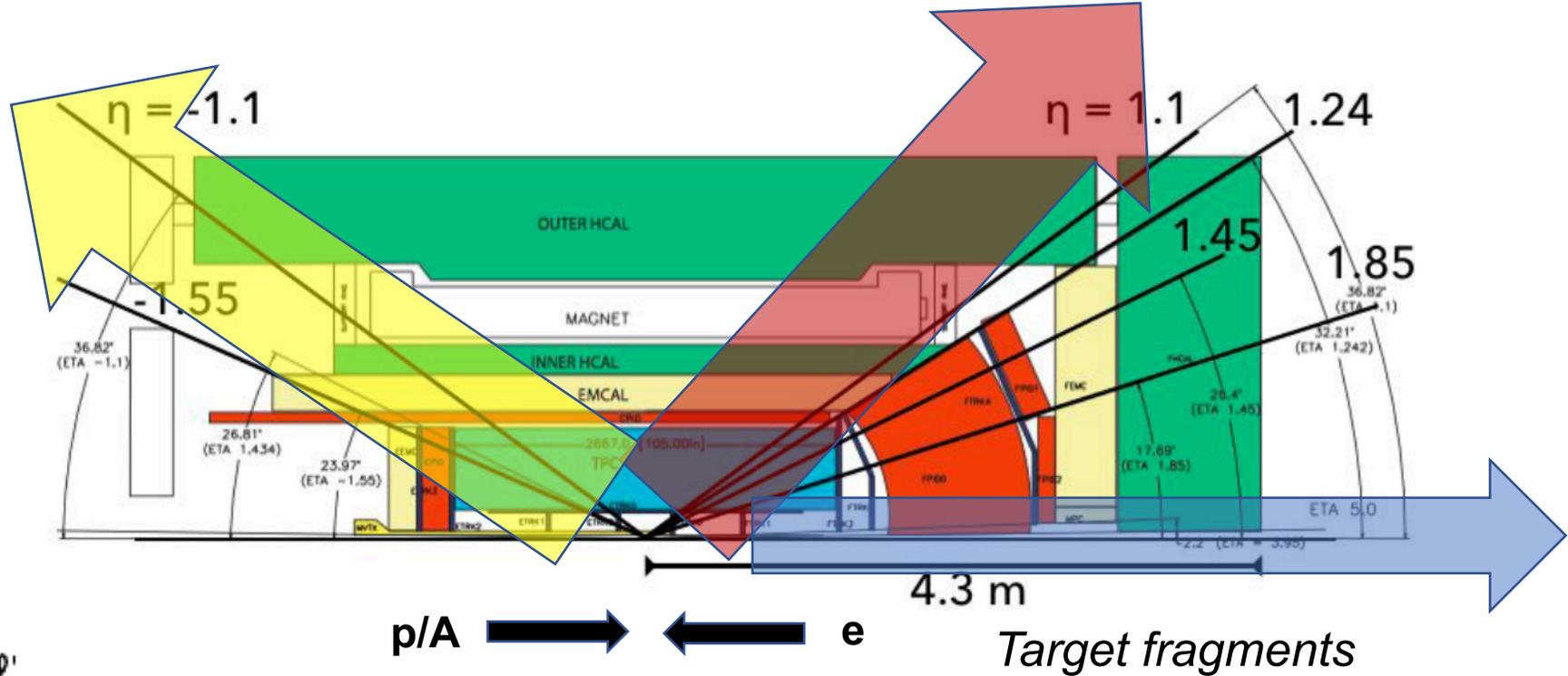
Thank you!

Backup

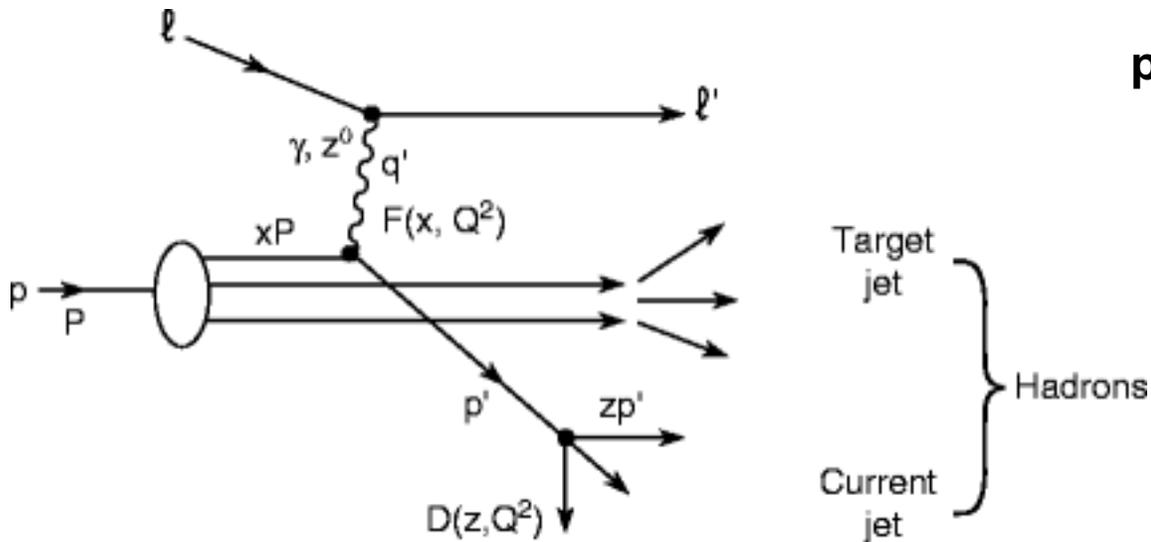
What do we measure?

Current jet (or hadron)

Scattered electron



Lepton scattering on a proton



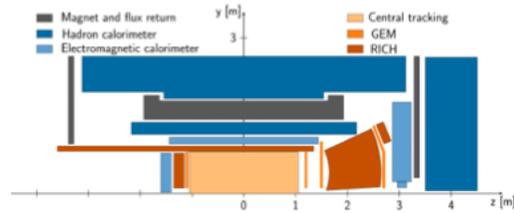
Inclusive DIS: only electron is detected

Semi-Inclusive DIS (SIDIS): electron and current jet (hadron) are detected.

Exclusive reactions: all particles are detected

Detection of target fragments – forward spectrometer

2.6 Evolution of sPHENIX into an Electron-Ion Collider Experiment 33



”near” detection

”far” detection

Analyzing dipoles (green) and accelerator quadrupoles (red)

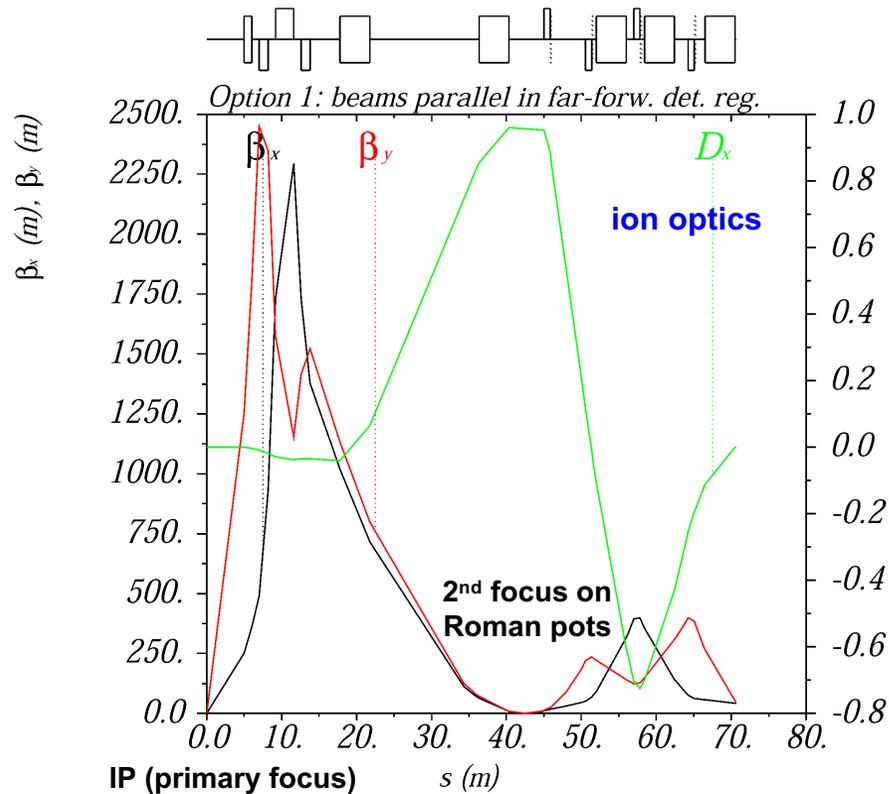
Detection space

high-res. ZDC (n, γ)

Forward hadron spectrometer
Spectator protons

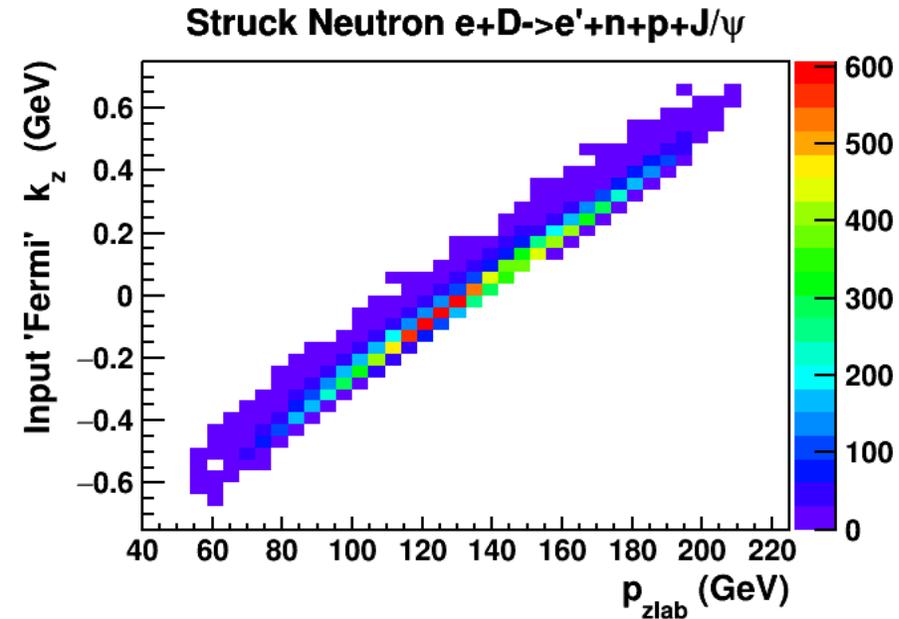
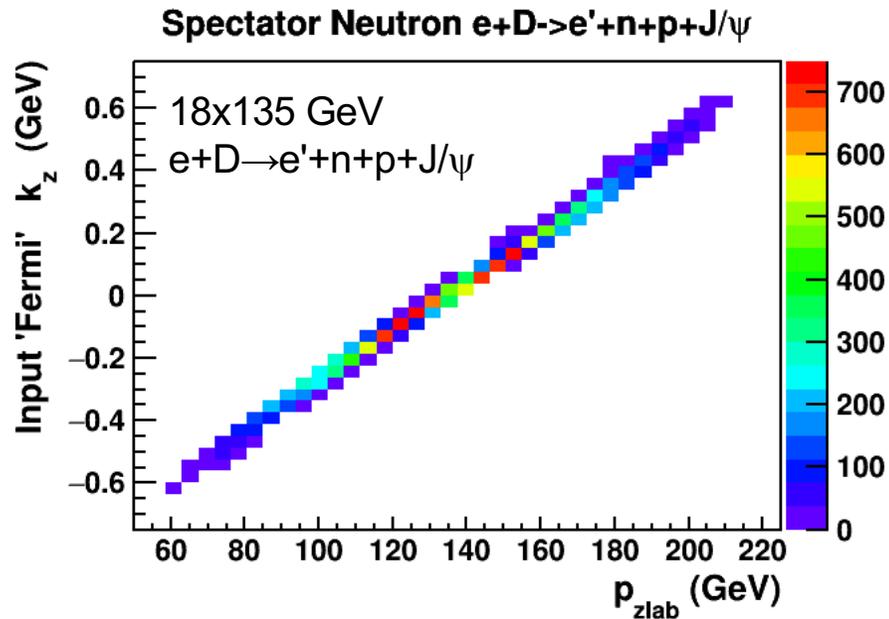
Recoil protons (roman pots at focal point with high dispersion)

p/A →
 e ←



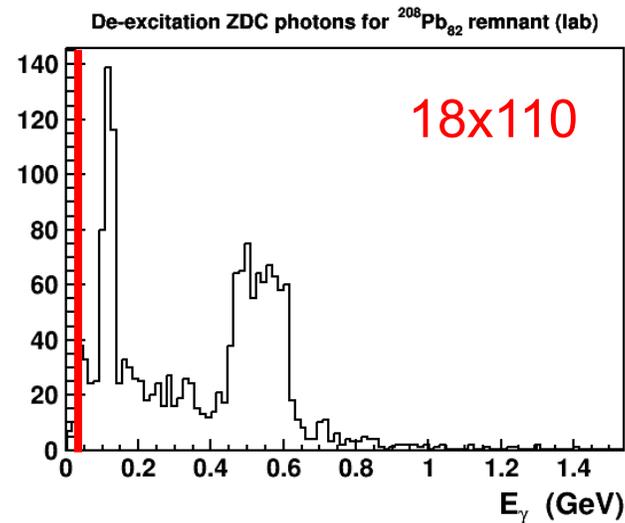
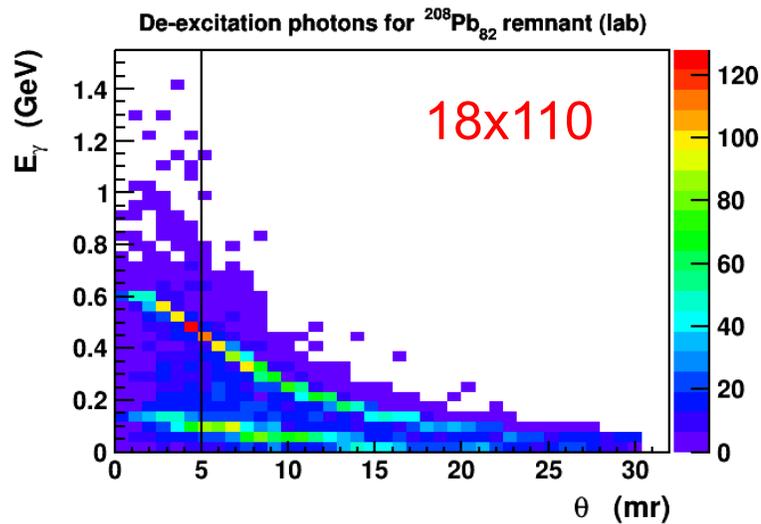
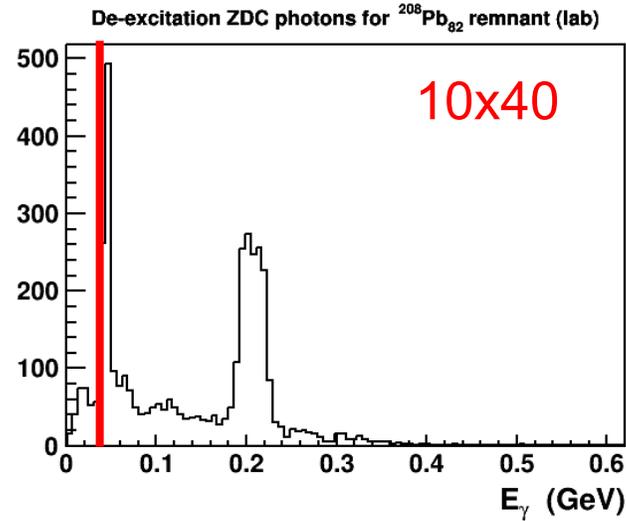
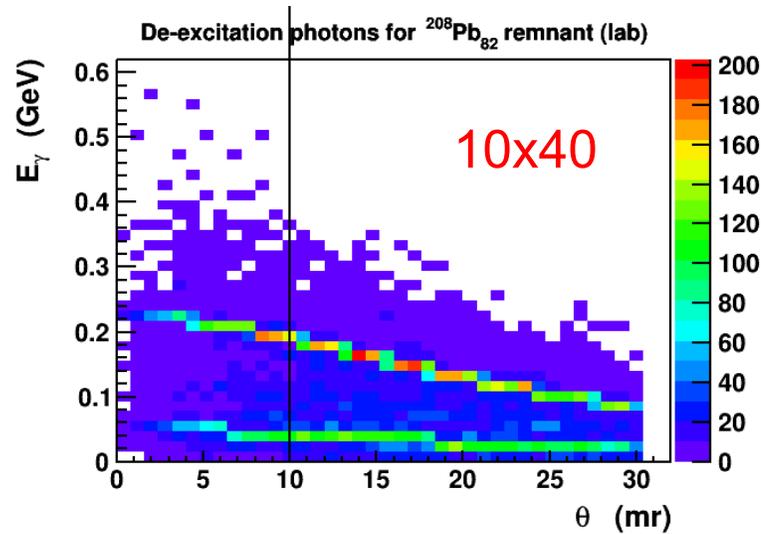
- Functionally, the forward detection is naturally separated into a ”near” and ”far” parts
- ”Near” detection:
 - *Goal*: off-momentum/rigidity particles or ones scattered at ”large” angles (high p_T)
 - *Requirement*: large magnet apertures
- ”Far” detection (can be after a crab cavity):
 - *Goal*: small-angle particles with momentum/rigidity close to that of the beam
 - *Requirement*: large dispersion and small beam size

Effect of boosting k_z to the lab frame



- The boost makes precision measurements of nucleons moving along the z-axis easier
 - It also makes precision gamma spectroscopy of photons from nuclear de-excitations in heavier nuclei possible.
- However, as noted earlier, this is also the reason why hadronic reconstruction methods do not work as well for eA as they do for pA.

Photons from $^{208}\text{Pb}_{82}$ in lab frame



Detailed studies ongoing.

It is clear that γ 's will be needed for low $|t|$!

w/ Morozov, Hyde, Turonski et al.

Density fluctuations from incoherent diffraction

- Incoherent diffraction probes the variance of the density
 - Calculations at $x \sim 10^{-3}$ show "lumpiness" in the gluon distribution and suggest sensitivity to model assumptions
- Could we in the future see the imprint of correlations from nucleonic degrees of freedom at the partonic level?
 - Experimentally straightforward, but a clear interpretation needs much more progress on the theory side

