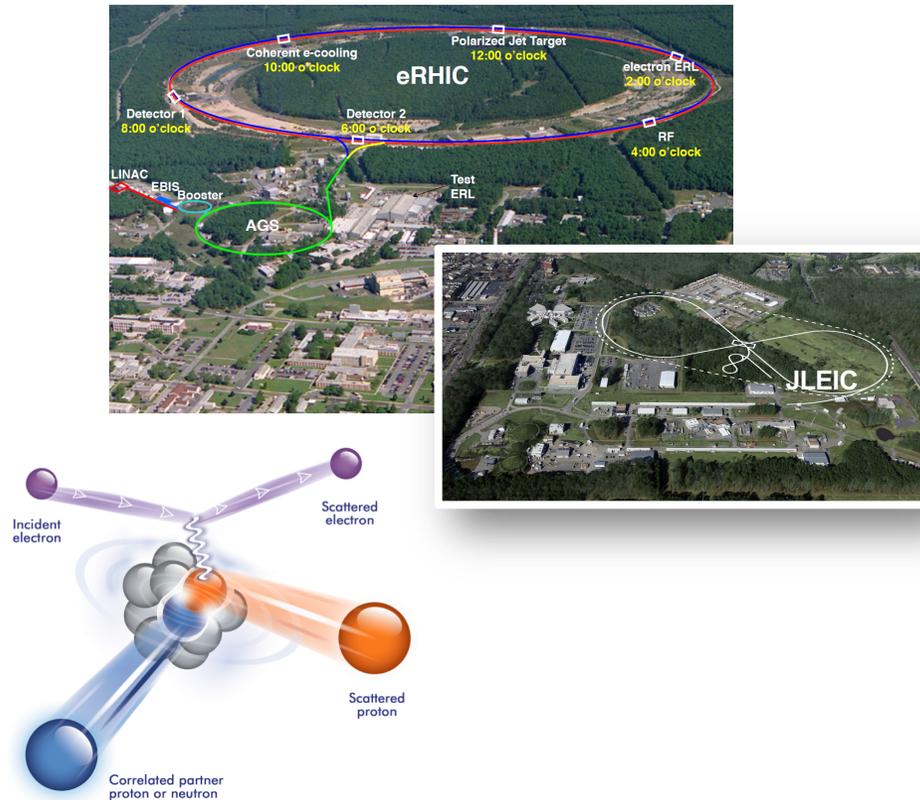


# Experimental SRC/EMC Studies: Opportunities and Challenges at EIC

- in a broader context of forward detection requirements



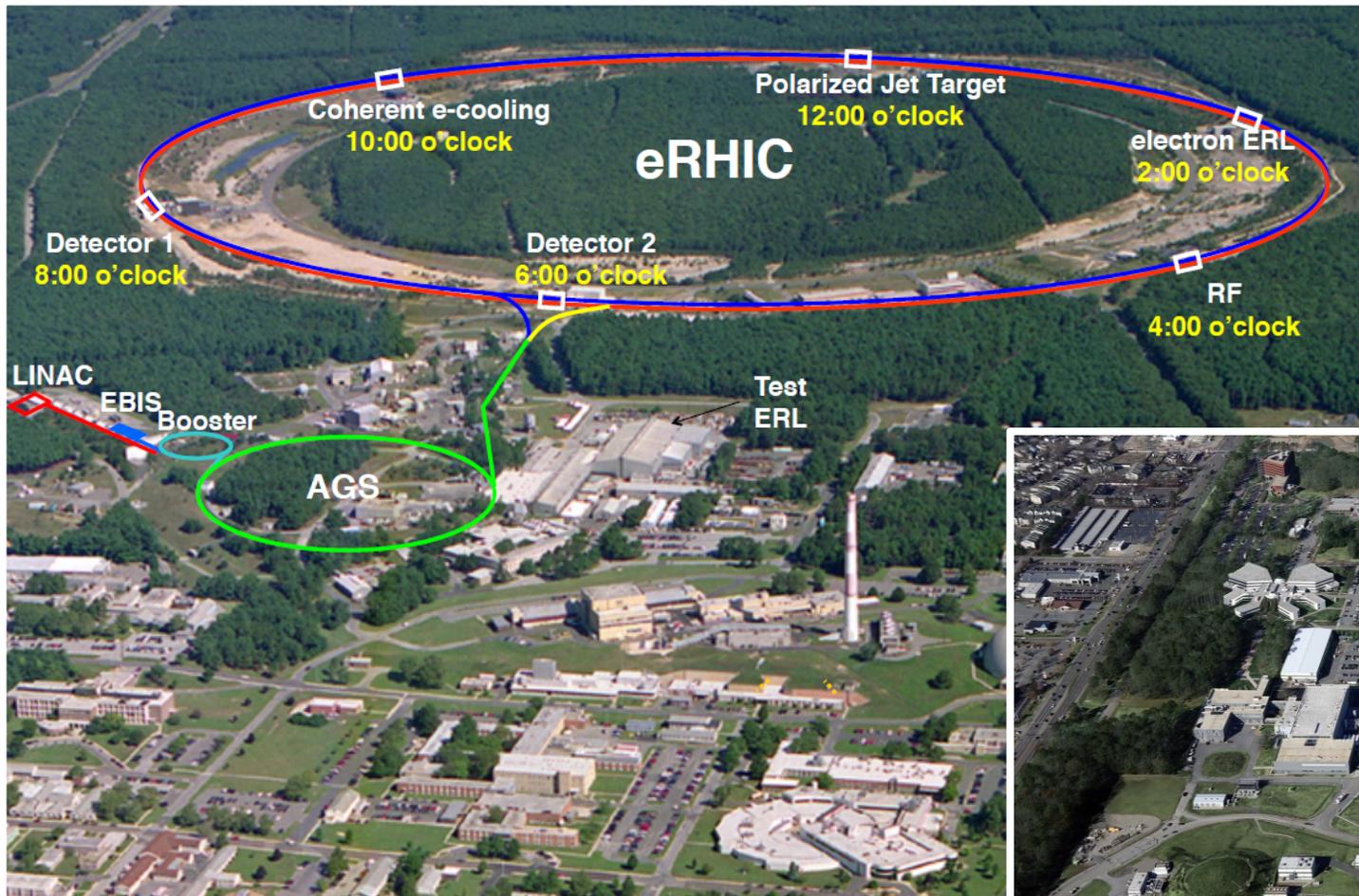
Pawel Nadel-Turonski  
Stony Brook University

Joint CFNS/RBRC on Physics and Detector Requirements at Zero-Degree of Colliders,,  
Stony Brook, September 24-26, 2019

# Outline

- Forward detection requirements at the EIC
- SRC measurements at the EIC
- Key measurements that have similar requirements

# Proposed EIC implementations at JLab and BNL



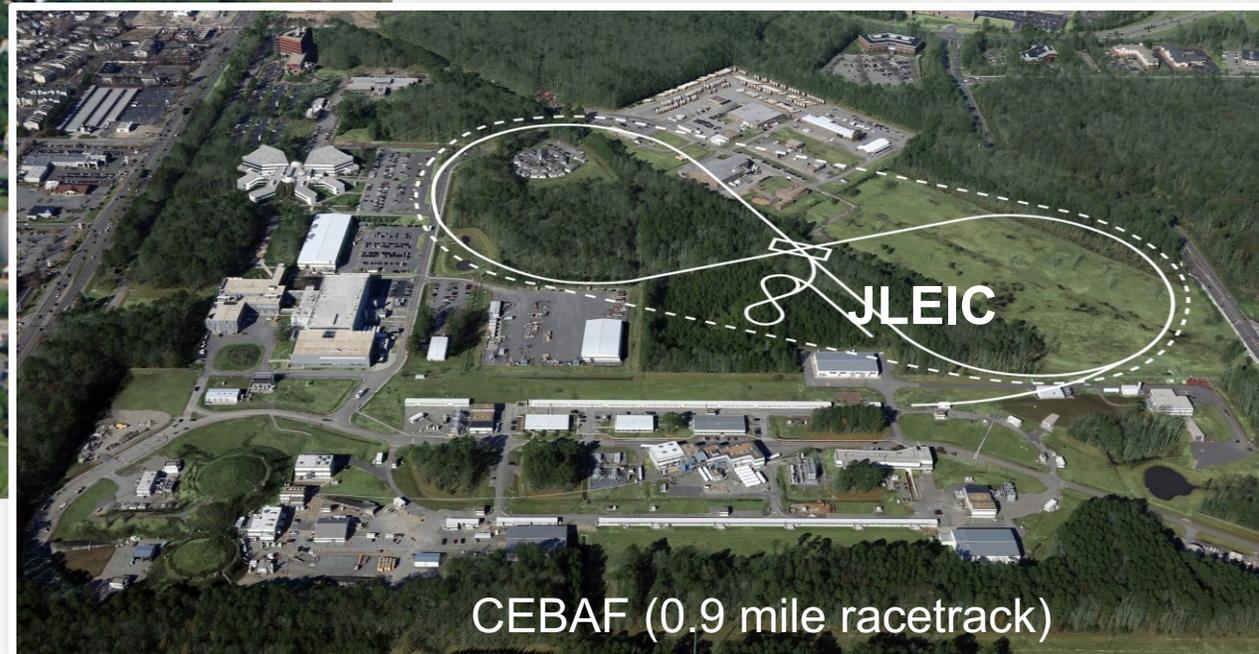
## BNL EIC at RHIC (eRHIC)

18 GeV e (10 GeV lumi max) on 275 GeV p

Parameters are similar

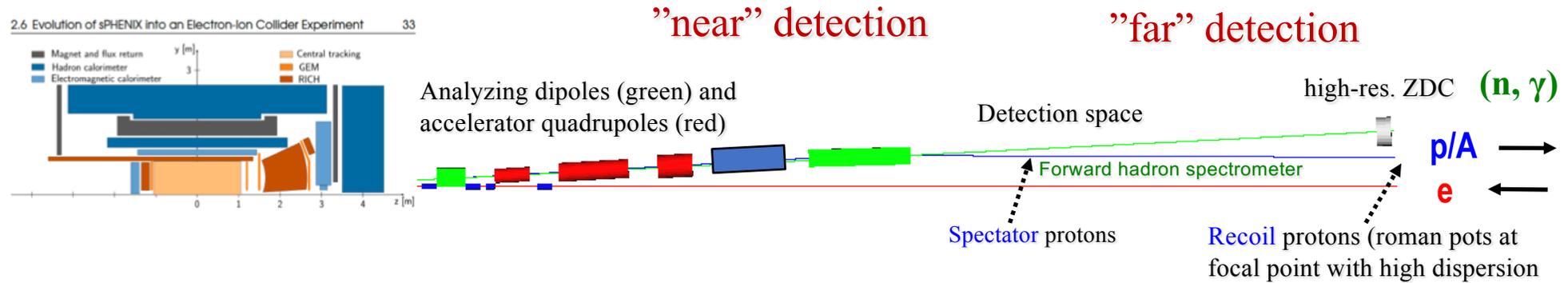
12 GeV e (5 GeV lumi max) on 200 GeV p

## JLab EIC (JLEIC)



Near-beam detection could be similar

# Detection of target fragments

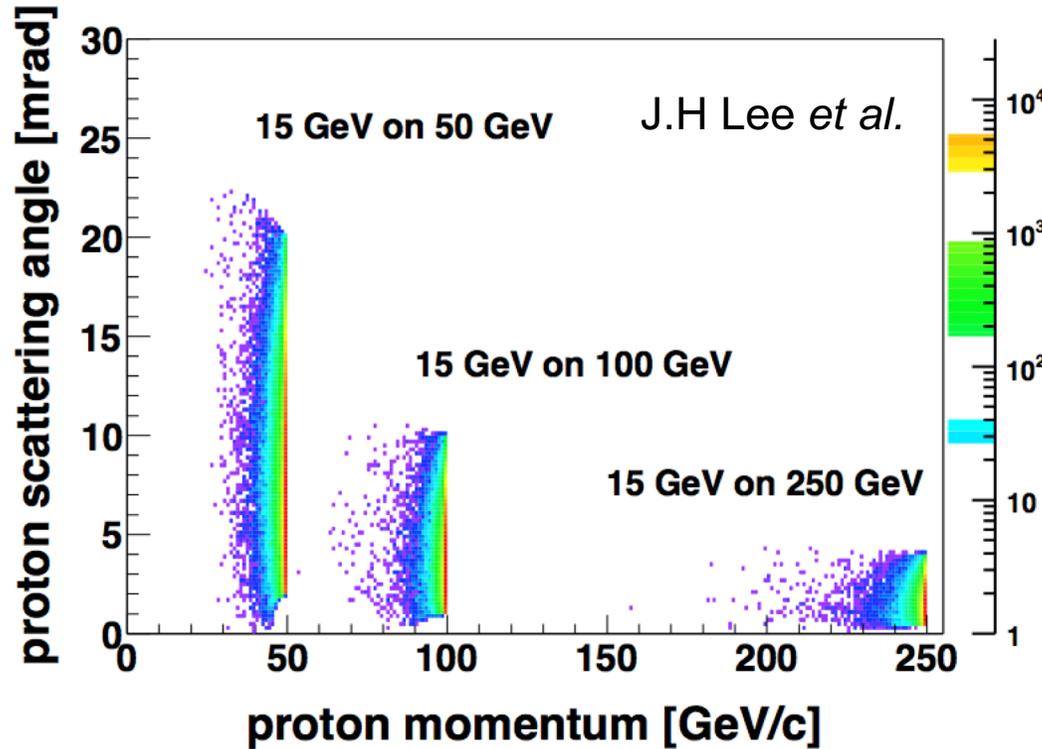


- A large part of the EIC program relies on detecting recoil baryons and target fragments
  - Spatial imaging through *exclusive* meson/photon production on the proton, and in coherent diffraction on nuclei
  - Neutron structure through spectator tagging in light nuclei
  - Various incoherent processes on heavy nuclei, **including SRC and EMC studies**
- Forward detection requirements for the EIC are very demanding
  - Need to detect particles very near the beam (down to  $p_T = 0$  for  $p/p_0 < 99\%$ )
  - Need to detect protons with  $p_T$  up to at least 1 GeV/c and spectators with  $A/Z$  very different from that of the beam (e.g., spectator protons from deuterium)
  - Need excellent momentum resolution, detection of neutrals, and PID for ions

# Forward detection – processes and requirements

- Which processes drive the near-beam (small  $p_T$ ) acceptance?
  - Coherent diffraction on light nuclei (detecting recoiling ion for clean signal)
  - DIS and (in)coherent diffraction on medium- and heavy nuclei (detecting residual nucleus)
- Which processes drive the “large” angle ( $p_T$ ) acceptance?
  - Tagging of spectator protons from nuclei
  - Exclusive production of photons (DVCS) and mesons (DVMP) on the proton at large  $t$  at *low energies*
  - Detection of photons and neutrons from nuclei (cone with line-of-sight)
- Which processes drive the resolution?
  - Magnetic spectrometer: tagging of protons from nuclei (spectators have  $p \sim$  Fermi momentum)
  - Hcal (ZDC): tagging of neutrons from exclusive charged meson production on the proton and spectator neutron tagging (e.g., for reactions on the proton in deuterium)
  - EMcal: Photons from nuclear de-excitations (coherent diffraction and rare isotopes)

# Near-beam acceptance using a magnetic spectrometer (large dispersion)

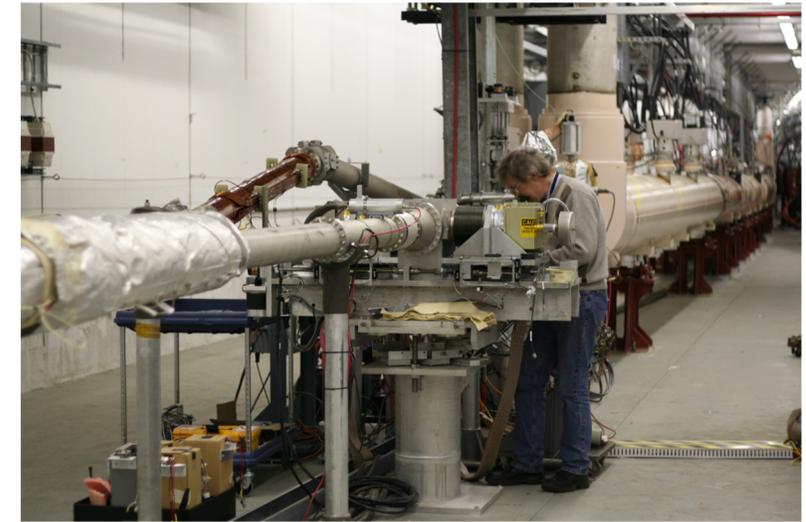


- The intrinsic beam momentum spread is on the order of a few  $\times 10^{-4}$ .  
*Insufficient to consider only the beam divergence at the IP*
- The target energy loss is typically 1% (DVCS on the proton shown)
- For heavy ions, the momentum change is small, but losing a single proton changes the rigidity ( $\sim A/Z$ ) by about 1%

- With sufficiently *large dispersion* (which separates off-momentum particles from the beam), it is possible to detect even protons scattered at zero degrees or nuclei losing one nucleons.
- Detection of intact nuclei is *much* harder than recoil protons, since at fixed  $t$  the larger mass implies a smaller change in both longitudinal and transverse momentum (scattering angle)

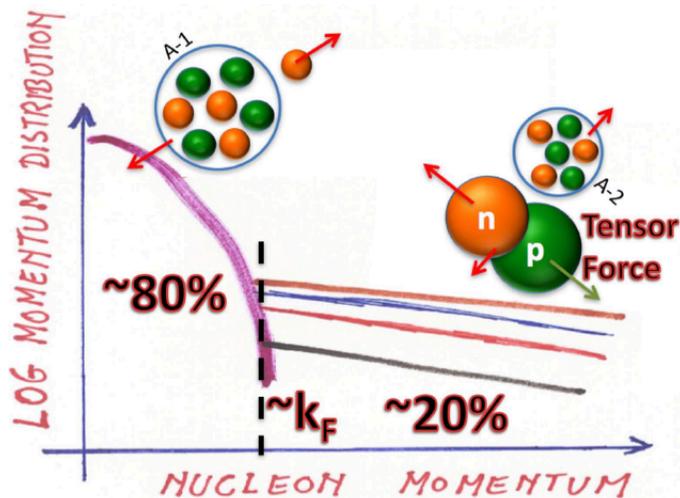
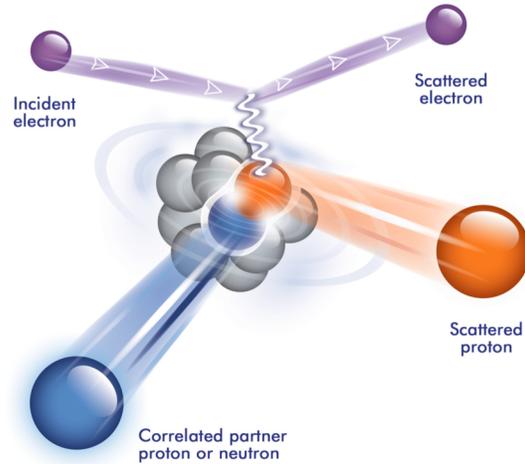
# Near-beam detection at JLab and BNL

- At JLab, a full forward spectrometer was always a key part of the MEIC (now JLEIC) concept since its introduction in 2009.
  - proton acceptance of  $E/E_{\text{beam}} < 99\%$  for *all* angles and down to 3 mrad for *all* energies
- At BNL, the IR currently lacks a high-dispersion section, limiting its potential for eA.
  - However, nothing prevents BNL to add a forward spectrometer with large dispersion.
  - Implementation easier at BNL since a larger e-ring allows reversing the order of electron and ion quads.
- Input on from the eA community on the detection requirements is thus important and urgent



Roman pots in STAR located ~50 m downstream from the collision point

# But SRCs are not a process – so where do they fit in?



## • Deuterium

- As **spectator tagging on the nucleon** -> talk by C. Weiss  
but with sufficient acceptance for high- $p_T$  tail
- Additional  $p_T$  for the struck nucleon will depend on process and kinematics, but can be substantial -> talk by F. Hauenstein

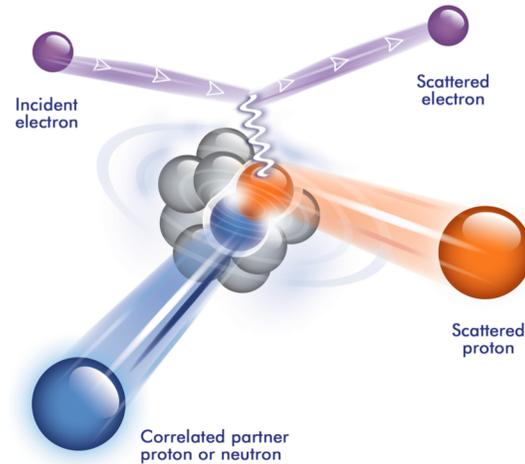
## • Light ions

- $A-2$  fragment(s) easy to detect, except for  $Z=A/2$
- The latter is, however, a very important case (d,  $^{10}\text{B}$ , etc) with detection challenges similar to **coherent processes**

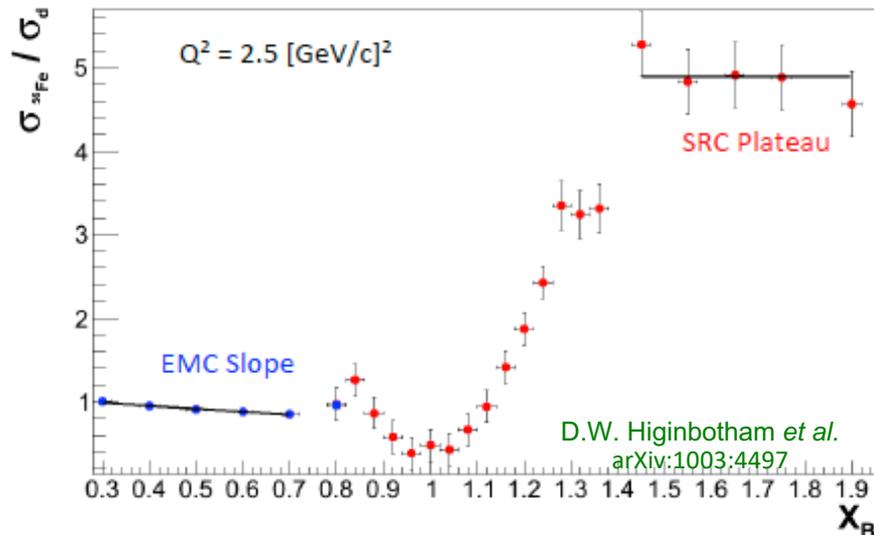
## • Heavy ions

- Tagging and identification of heavy fragments needed.
- Similar to measurements of **rare isotopes at the EIC** and helpful for **coherent diffraction** on heavy ions.

# Traditional SRC measurements at an EIC



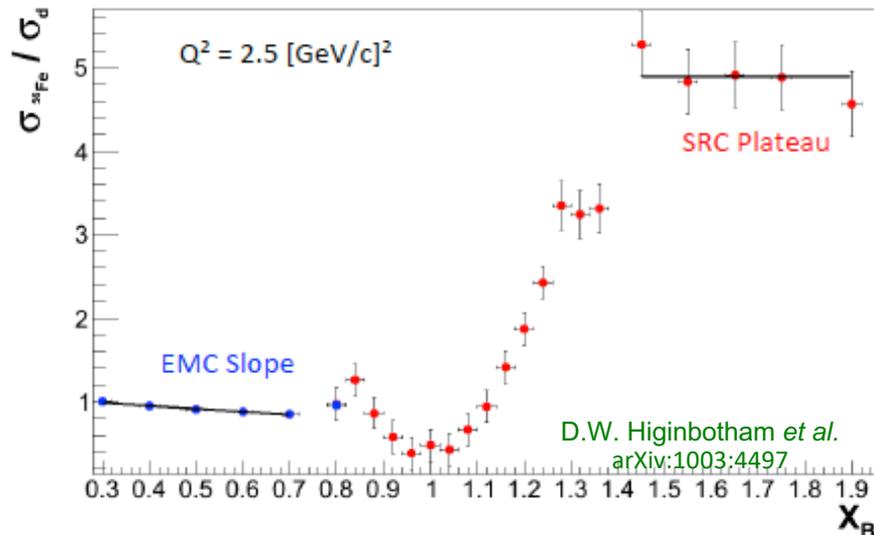
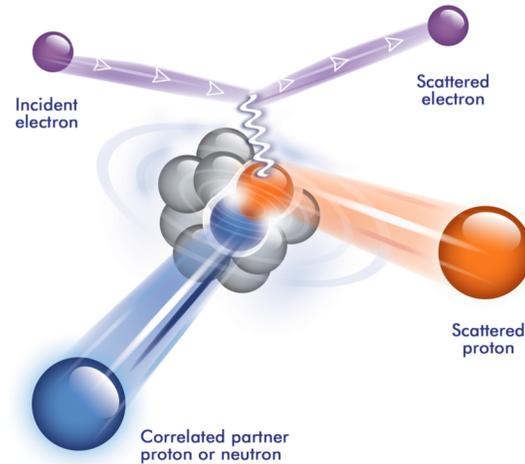
- SRCs have traditionally been studied at high  $x$  in quasielastic (QE) reactions in medium- and heavy nuclei.
  - $x > 1$  ensures participation by more than one nucleon
  - QE process provides straightforward measurement and easier-to-understand final-state interactions



- The EIC can contribute by providing a large lever arm in  $Q^2$  at high  $x$  and detection of all target fragments – but there are some challenges.
  - Event rates at high  $x$  are much lower than for fixed target
  - Resolutions at large  $x$  deteriorate rapidly

-> talk by B. Schmookler

# SRCs in DIS at lower x?



- The EMC effect ( $x \sim 0.5$ ) is much easier to measure in “typical” EIC kinematics.
- But also for SRC studies at an EIC it would be natural to seek relevant measurements in DIS (or exclusive reactions) at lower  $x$ .
  - Can we learn something new about SRCs?
  - Is it important to introduce SRCs into nuclear models (e.g., BeAGLE) for studies of low- $x$  processes?
- Nuclei over the full mass range will provide important input.

# JLab LDRD-funded project for SRCs at an EIC

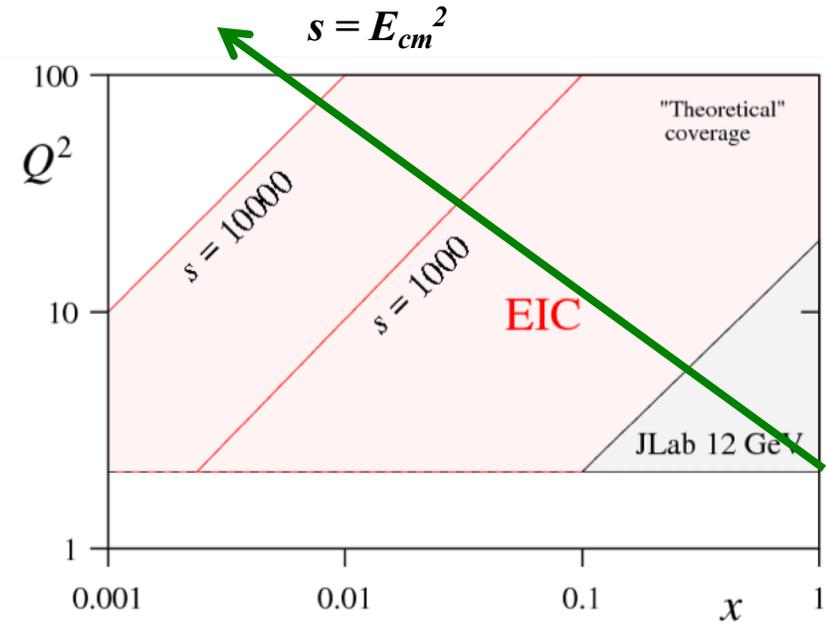
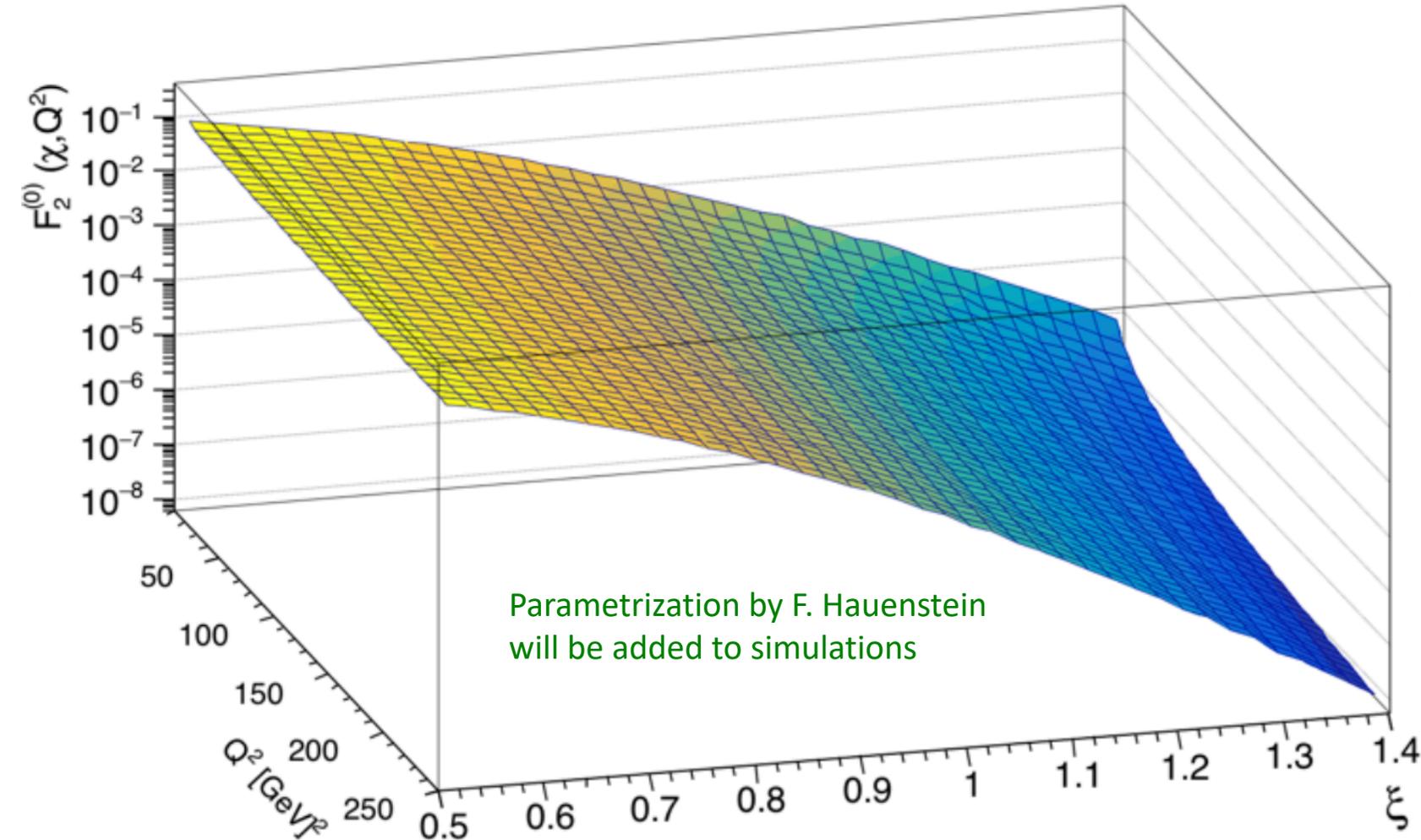
JLab “SRC” LDRD1912: D. Higinbotham, M. Baker, A. Deshpande, F. Hauenstein, O. Hen, C. Hyde, V. Morozov, PNT, A. Schimdt, B. Schmookler, Z. Tu, L. Zheng

- Feasibility studies
  - Rates at high  $x$  *F. Hauenstein*
  - Resolution in  $x$  at high  $x$  *B. Schmookler*
- Simulations, modeling, and detection requirements
  - Incorporating SRCs into BeAGLE *M. Baker*
  - Detection requirements

# Rate estimate: parametrization of modified $F_2$

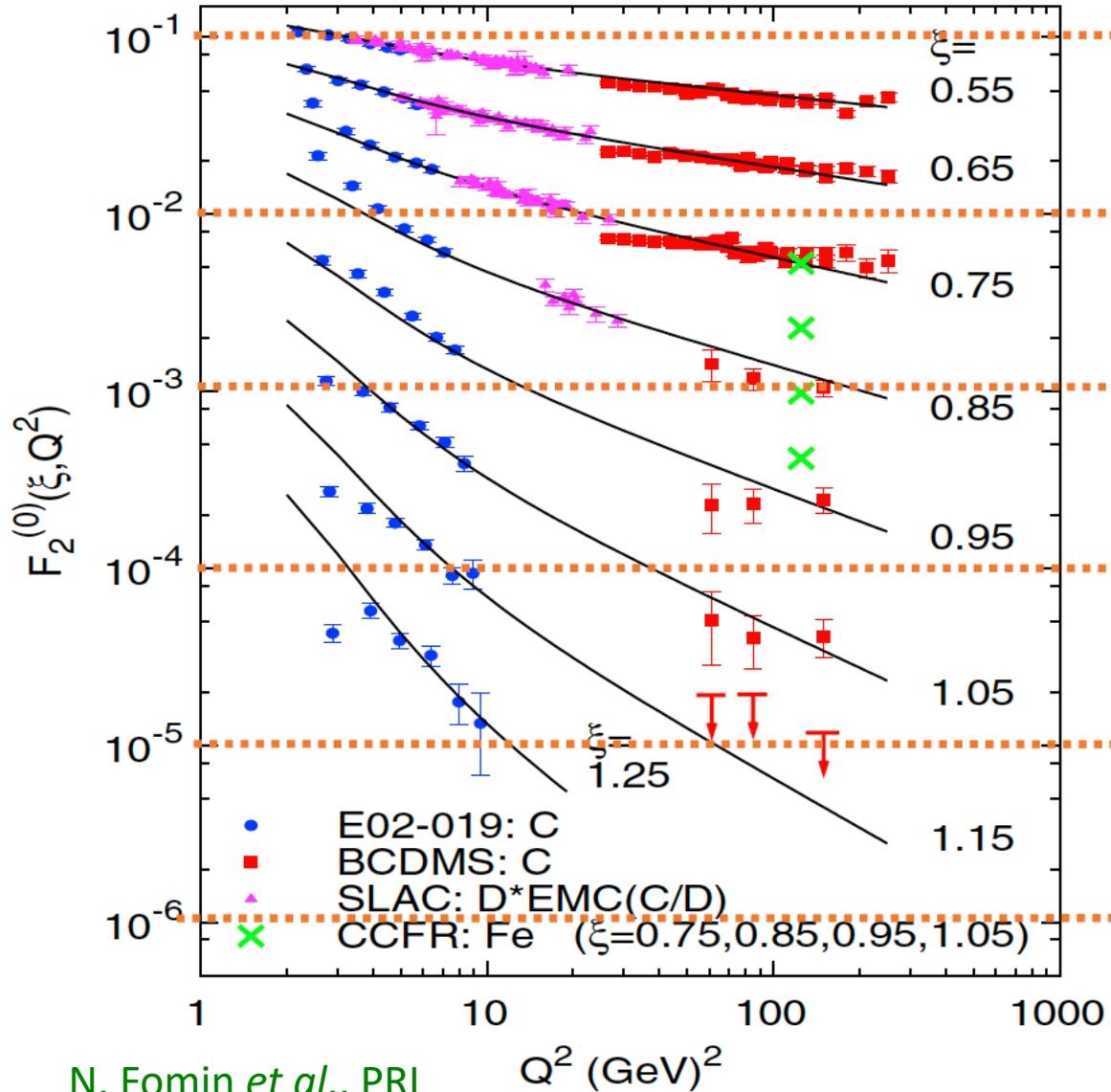
$$F_2^0(\xi, Q^2)$$

$$Q^2 \sim sxy$$

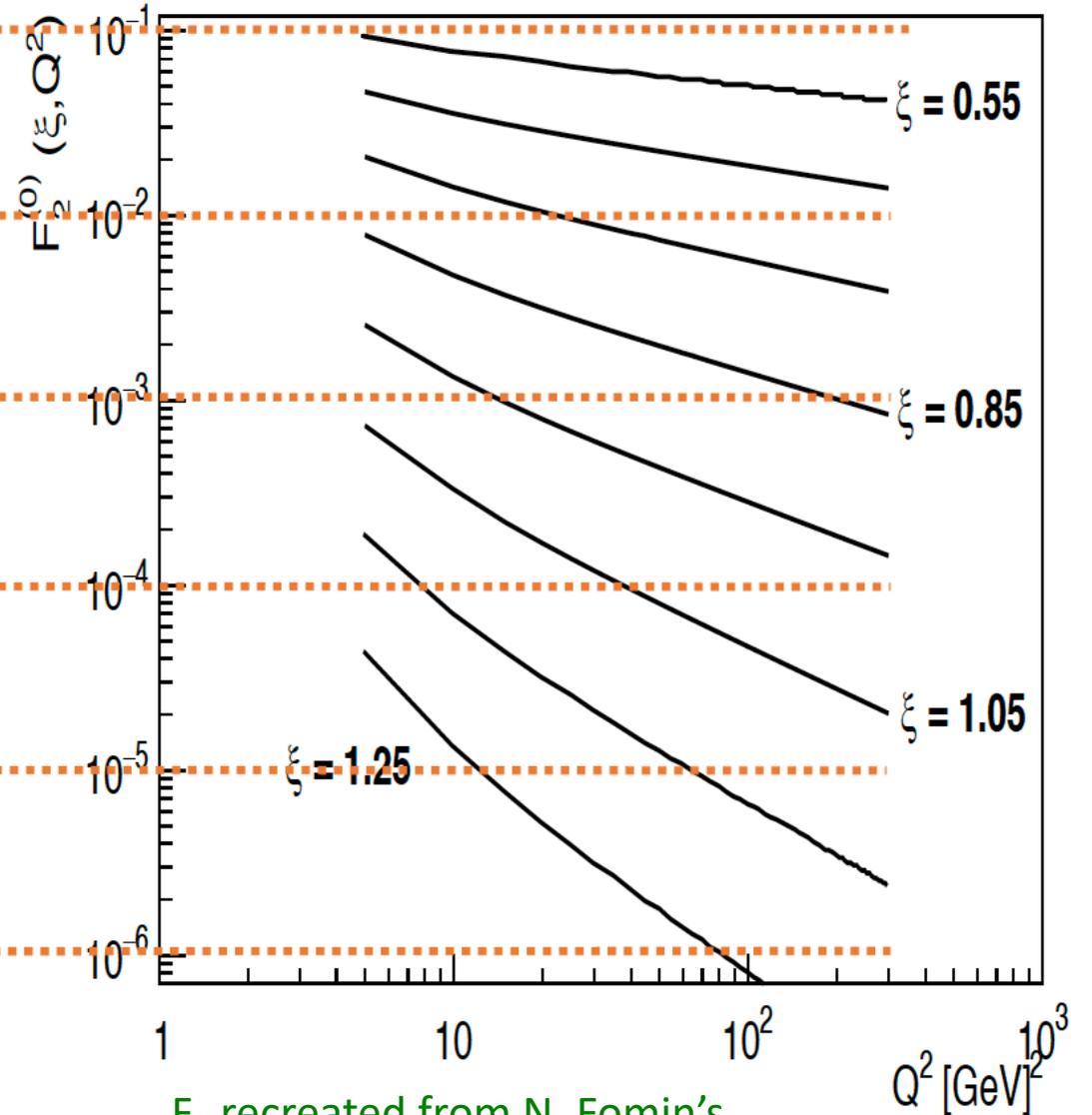


- To reach high  $x$ , the EIC would have to operate at low  $s$  and low  $y$  and/or high  $Q^2$ 
  - Lower  $s$  reduces luminosity
  - Lower  $y$  reduces resolution
- At high  $x$  and  $Q^2$ , the value of  $F_2$  is also low.
  - Rate is challenging
  - Running conditions need to be optimized

# A more detailed look at the modified $F_2$



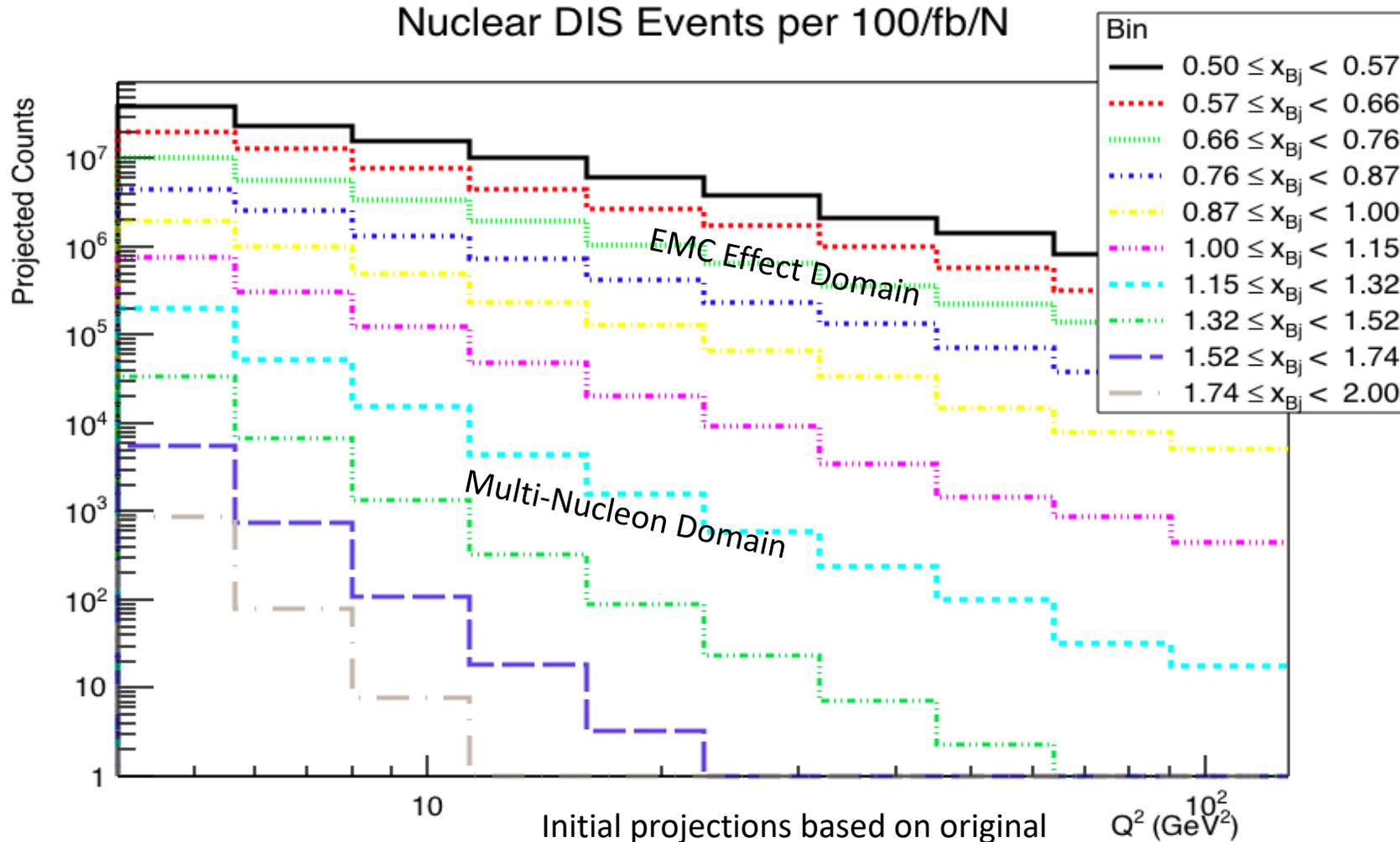
N. Fomin *et al.*, PRL  
105, 212502 (2010)



$F_2$  recreated from N. Fomin's  
scripts by F. Hauenstein

# Initial rate estimates for high x (and $Q^2$ )

5x50 GeV e+C JLEIC with 6 weeks @ 100% eff. (at lower energy rates would be lower)

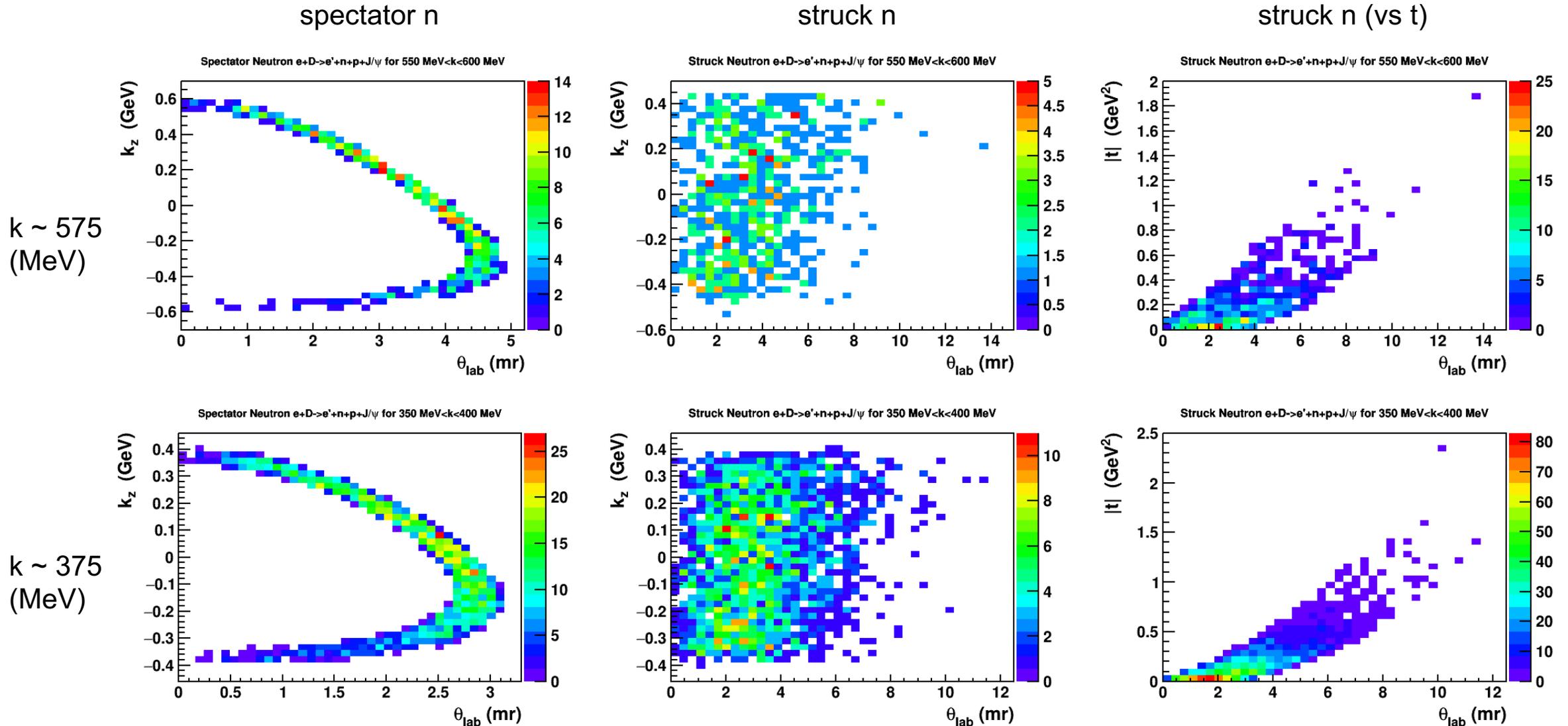


- Measuring the EMC effect at high  $Q^2$  is straightforward

- Measuring SRCs at  $x > 1$  is challenging, but there will be some events.

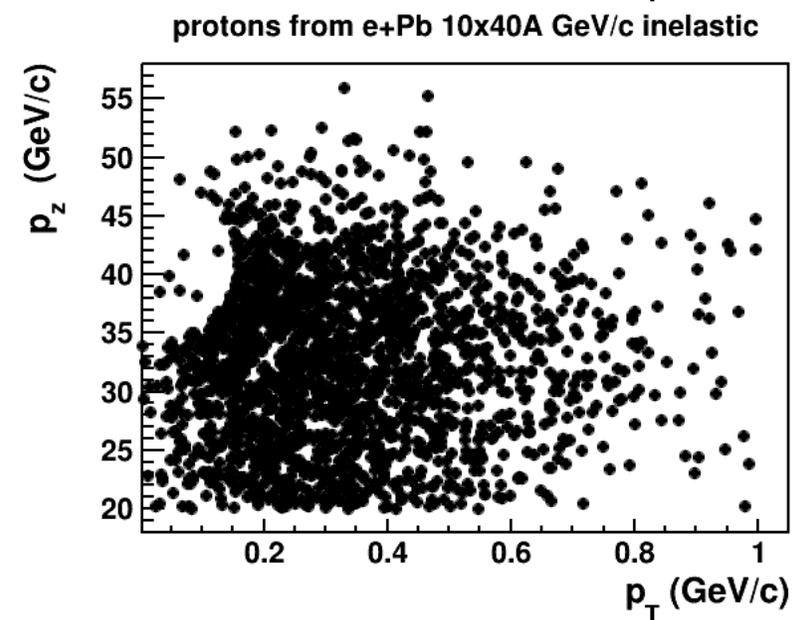
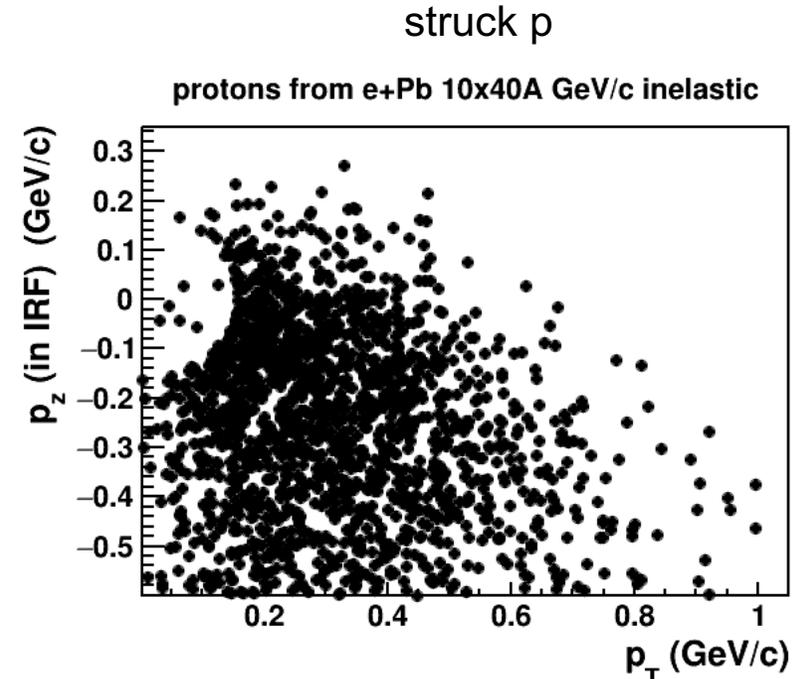
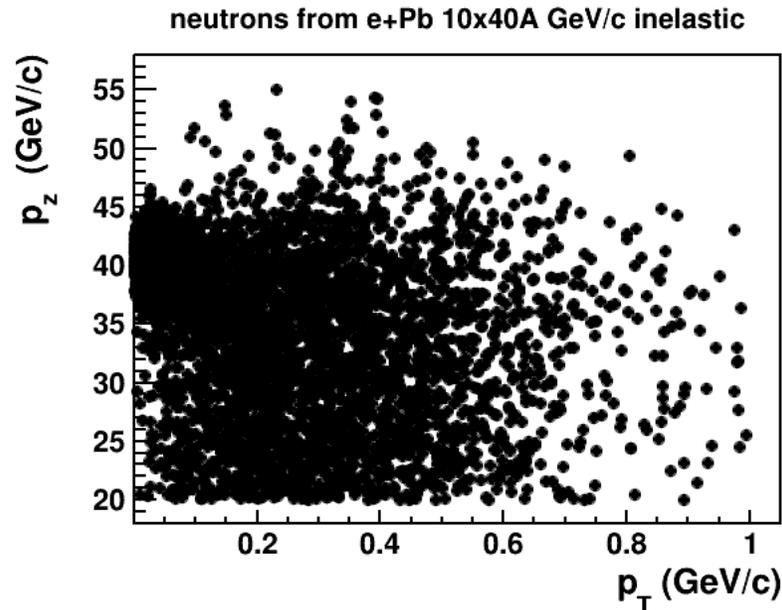
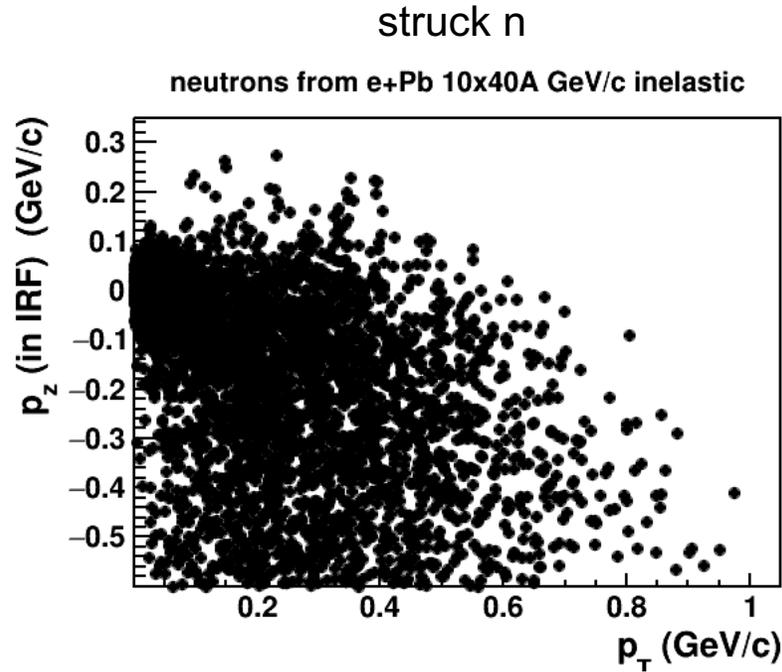
# Lab kinematics of spectator and struck neutrons at high energy

M. Baker



# And protons...

- For protons, Coulomb effects are not negligible.
- We need to merge the primary process with a detailed simulation of the nuclear response (BeAGLE) to understand the measurements



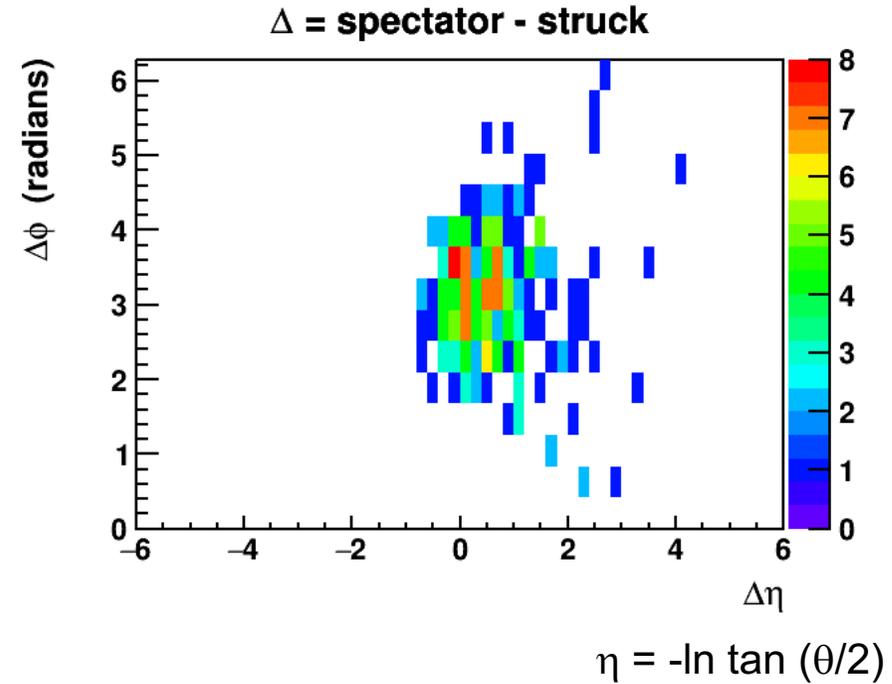
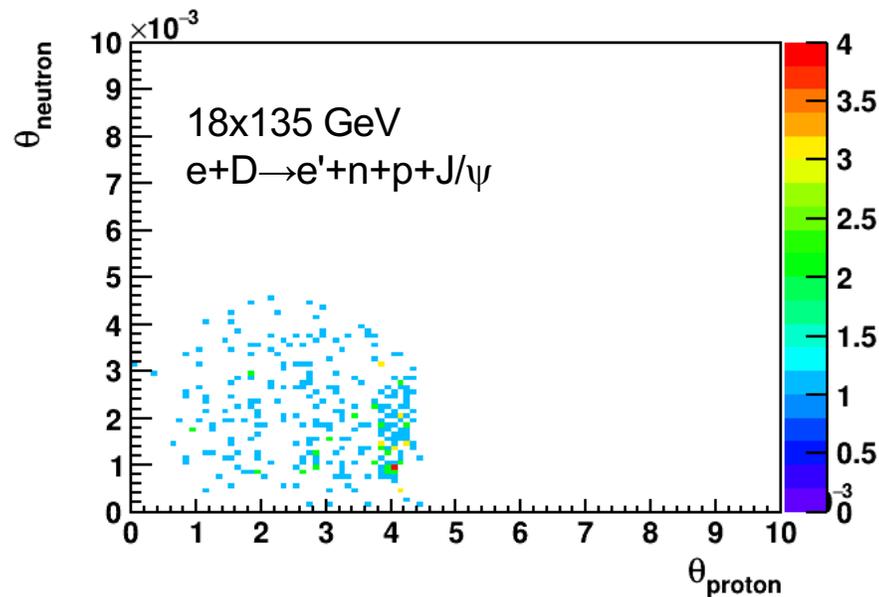
M. Baker

IRF = ion rest frame

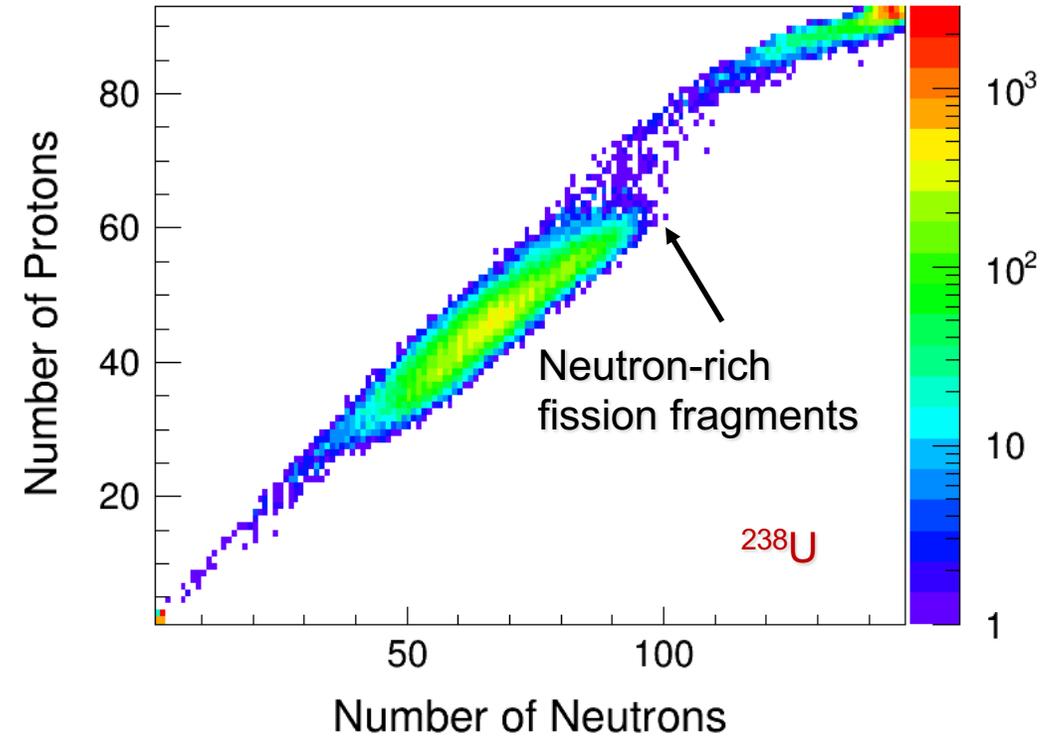
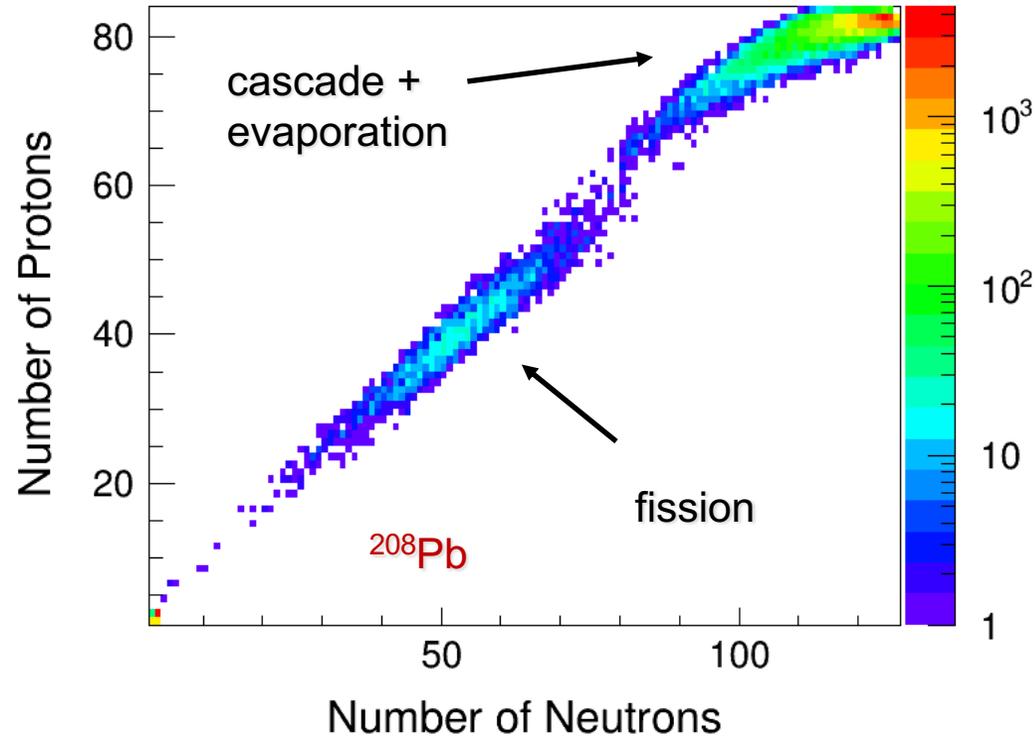
# And correlations (deuteron breakup)...

struck neutron, spectator proton  
 $|t| < 0.1 \text{ GeV}^2$ ,  $550 < k < 600 \text{ MeV}$

*M. Baker*

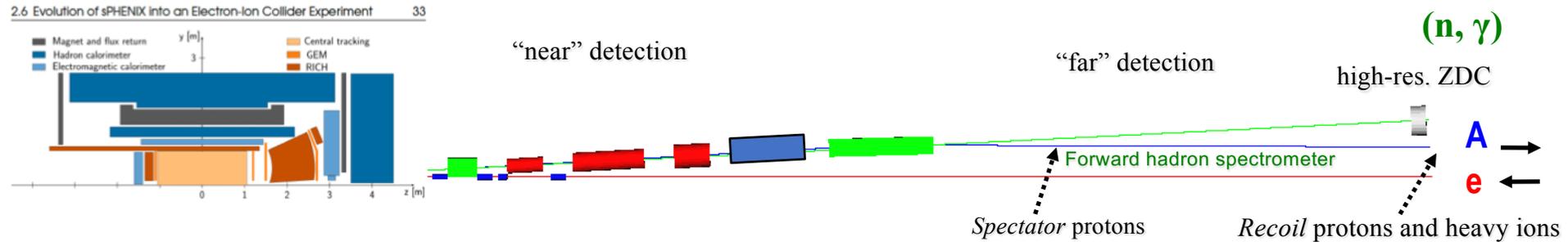


# Nuclear fragments produced in eA reactions at the EIC

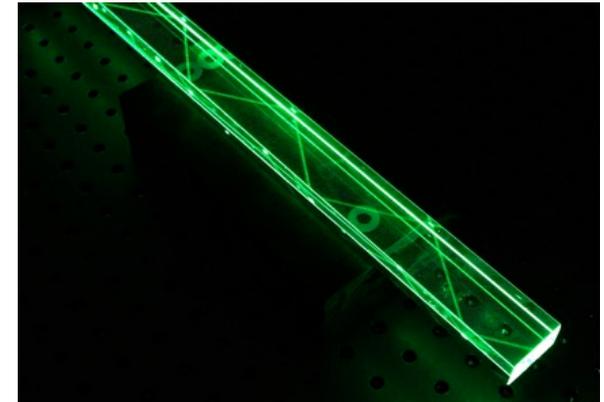


- Fragment detection (“tagging”) is needed to, for instance, reconstruct the Fermi momentum of the struck nucleon

# Detection and identification of the produced nuclei

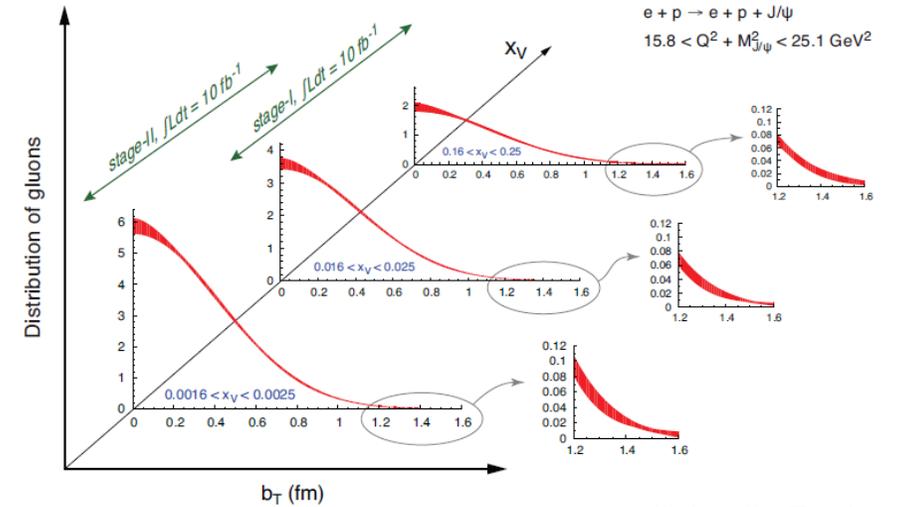
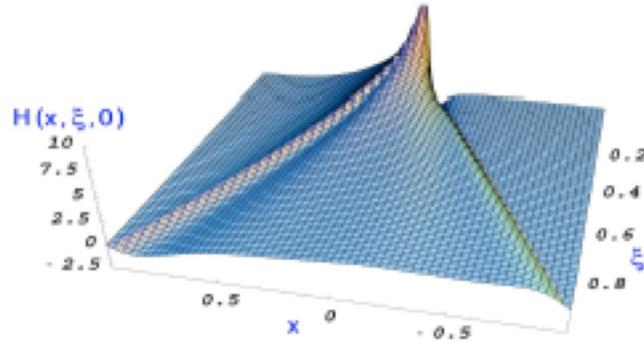
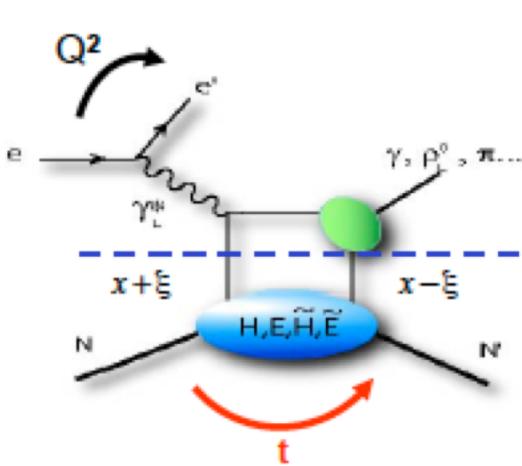


- Detection of nuclei with rigidities ( $A/Z$ ) close to that the beam require a dedicated “far” spectrometer where the beam is small and dispersion large.
- But to identify the ion we need both  $A/Z$  and an independent measurement of  $Z$ . The requirement for sensitivity in  $Z^2$  is 2% in order to identify heavy residual nuclei down to  $A-1$ .
- A “mini-DIRC” can produce 100,000 photons ( $\ll 1\%$  error) in a few mm of fused silica.
- R&D in progress (Generic Detector R&D for an EIC program).



A “mini-DIRC” inside a Roman pot at the downstream focus can identify ions to  $\sim 1\%$  in  $Z^2$

# Generalized Parton Distributions (GPDs) and 3D spatial imaging



-> talk by S. Fazio

- Exclusive production of a photon or meson is sensitive to GPDs
  - In the limit of small “skewness”  $\xi$ , the Fourier transform of  $t$ , the four-momentum transfer to the nucleon, can be interpreted as a spatial image in impact-parameter space at different values of  $x$
- To create a spatial image in  $b$ , data over a sufficient range in  $t$  ( $\sim p_T^2$ ) are needed ( $\sim 0-1 \text{ GeV}^2$ )
  - Small  $t \Leftrightarrow$  large  $b$  and vice versa.
  - For nuclei with  $A > 1$ , that scattering angle is much smaller for a given  $t$ , cutting the low  $t$  acceptance

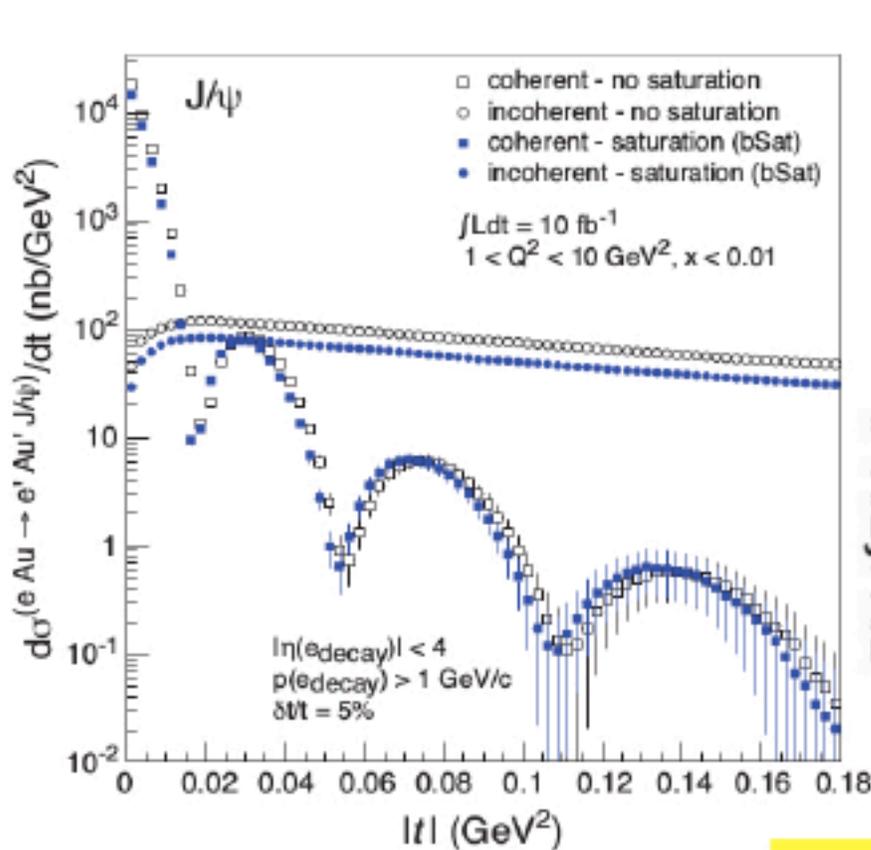
# Coherent exclusive (diffractive) reactions on light nuclei



- Coherent exclusive reactions on light ions give *unique* access to the 3D structure of *nuclei*.
- In contrast to heavy nuclei, scattered light nuclei can be detected
  - The momentum transfer  $t$  can be determined directly from the ion
- Light ions span the full range of nuclear densities
  - D is the least dense nucleus unbound, while He-4 is comparable to heavy ions, and He-3 falls in the middle
- Polarized He-3 beams will allow for simultaneous measurement of both tagged neutron structure and coherent diffraction on He-3
  - Interesting comparison since spin of He-3 is dominated by the neutron
- He-4 has only one GPD and large asymmetries
  - Easy to measure and interpret

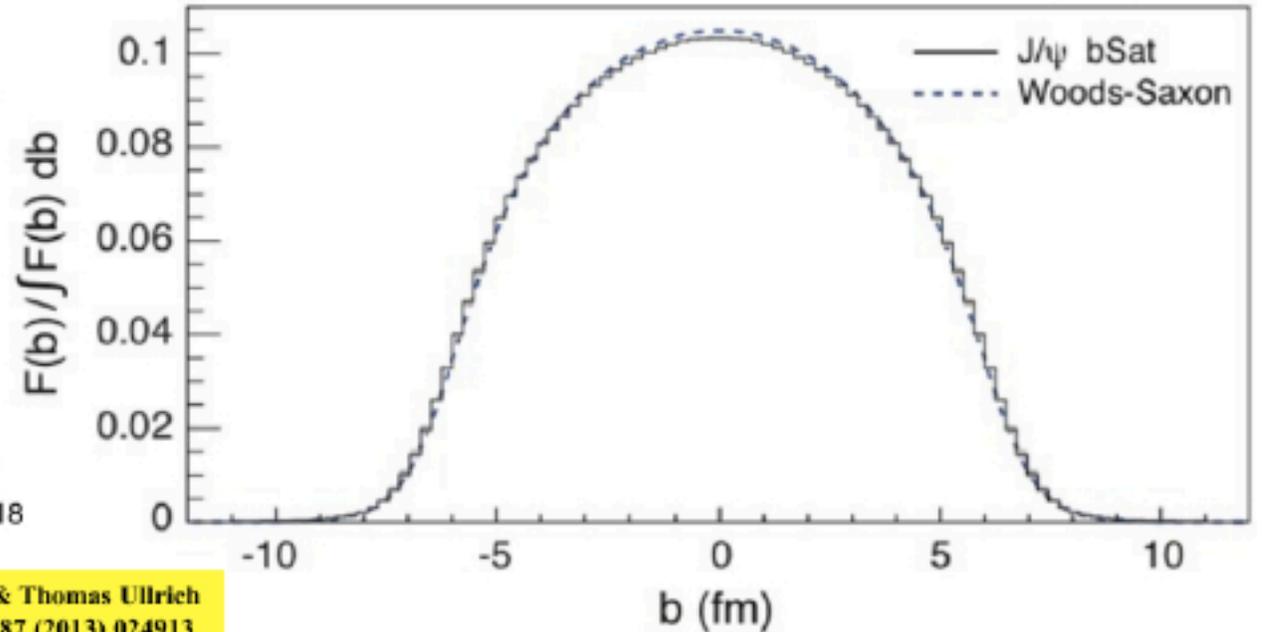
# 2D spatial gluon imaging in heavy nuclei through coherent diffraction

Momentum transfer  $t$  conjugate to transverse coordinate  $b$



$$F(b) \propto \int d\Delta \Delta J_0(\Delta b) \sqrt{\frac{d\sigma}{dt}}$$

$$t = \Delta^2/(1-x) \approx \Delta^2$$

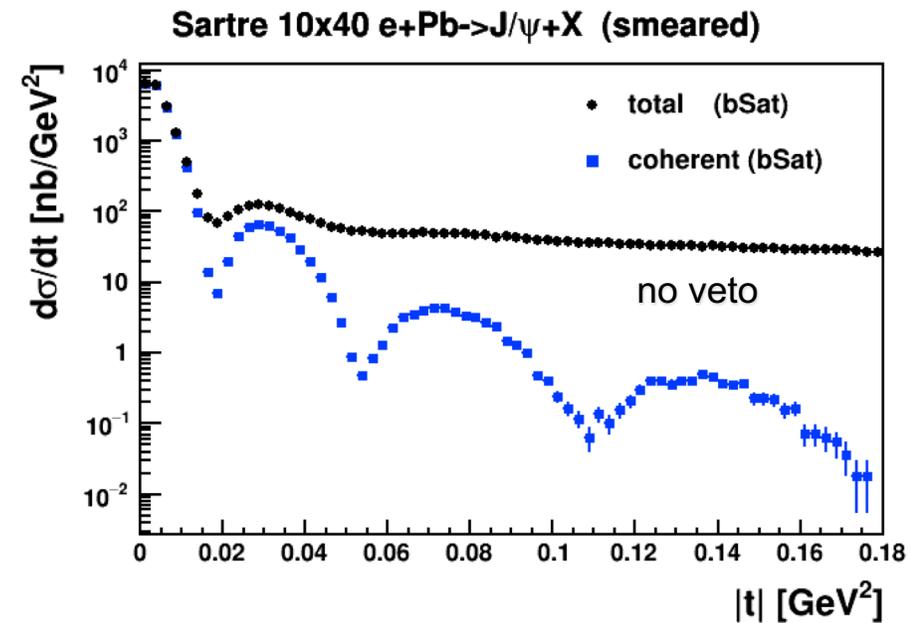


T.T. & Thomas Ullrich  
 PRC 87 (2013) 024913

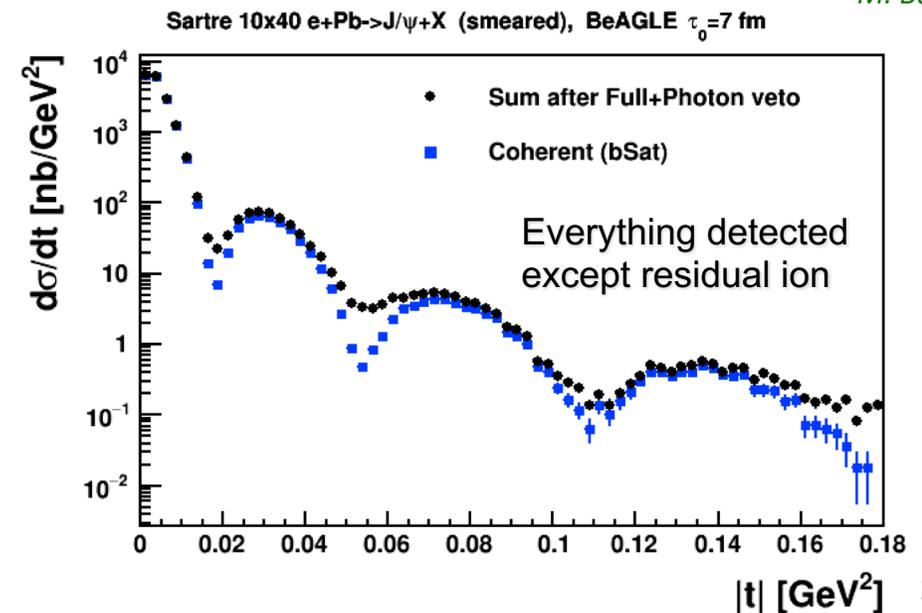
- Experimental challenge: veto the large incoherent background.

# Evaluate the veto efficiency of target fragment detection

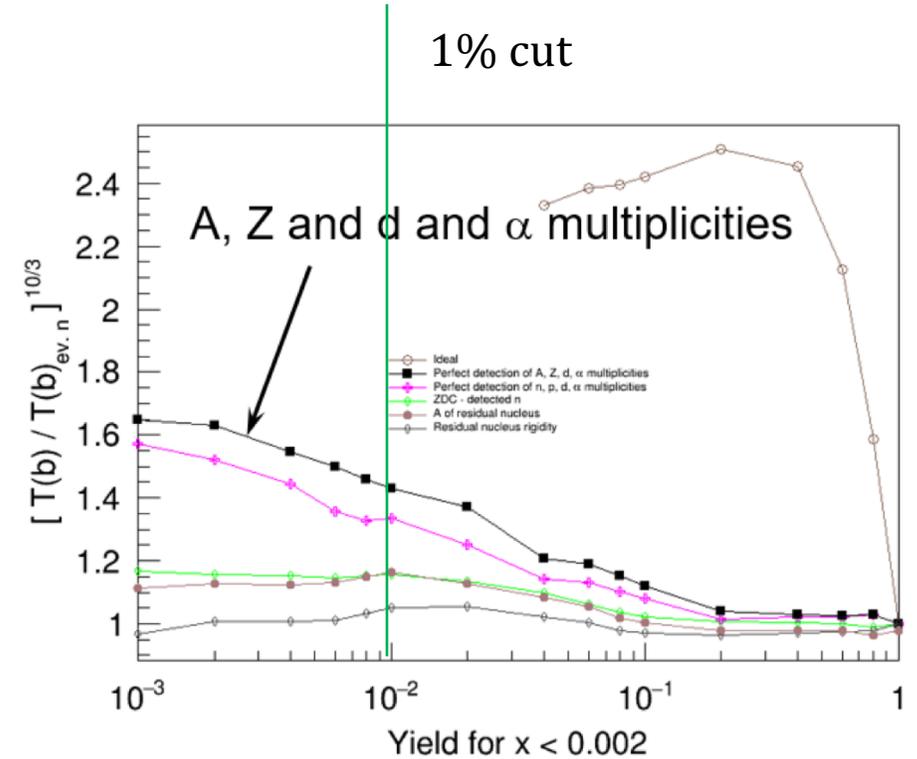
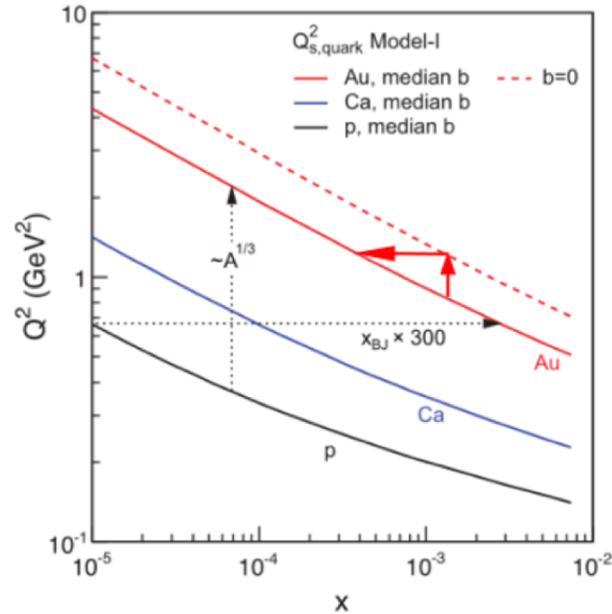
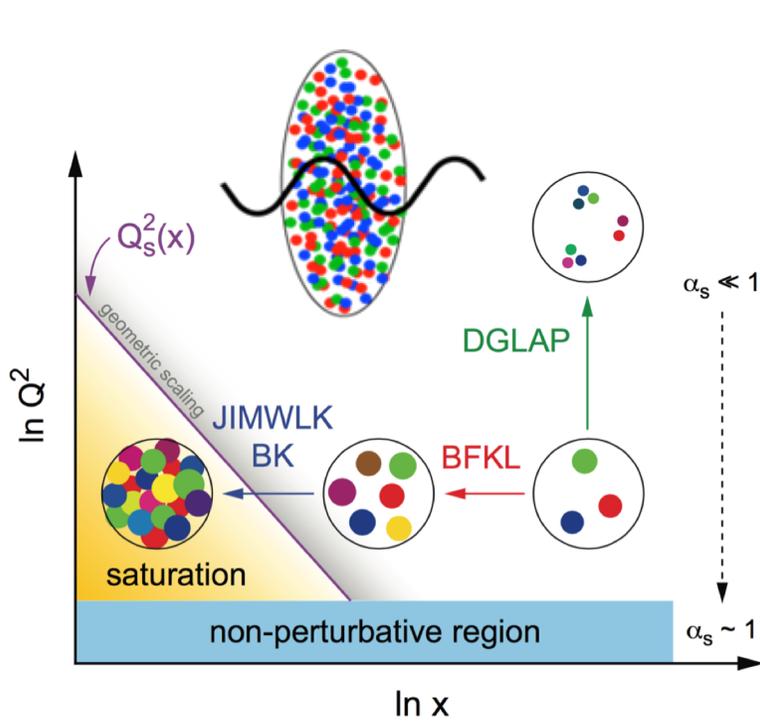
- Intact heavy ions cannot be detected directly
  - The momentum transfer is too small to kick intact heavy ions out of the beam envelope
- Need high-efficiency veto of incoherent events
  - Detection of *all* produced particles is required: protons, neutrons, light nuclei, and photons from nuclear de-excitations -> talk by C. Hyde
  - Detection of the residual nucleus (from incoherent diffraction) is also helpful
- BeAGLE was used to simulate the incoherent part – in this case the background
  - The coherent simulation used Sartre



M. Baker



# Kinematic reach for probing gluon saturation in DIS



- At low  $x$ , the photon interacts coherently with the gluons along its path inside a nucleus, allowing the EIC to reach into the saturation regime.
  - Impact parameter tagging selects events with larger average density  $T(b)$
- A larger thickness  $T(b)$  is equivalent to a higher beam energy
  - With a evaporation neutron multiplicities only, this factor is 3.2 at 1%
  - Detecting all final-state particles, this increases to 4.8

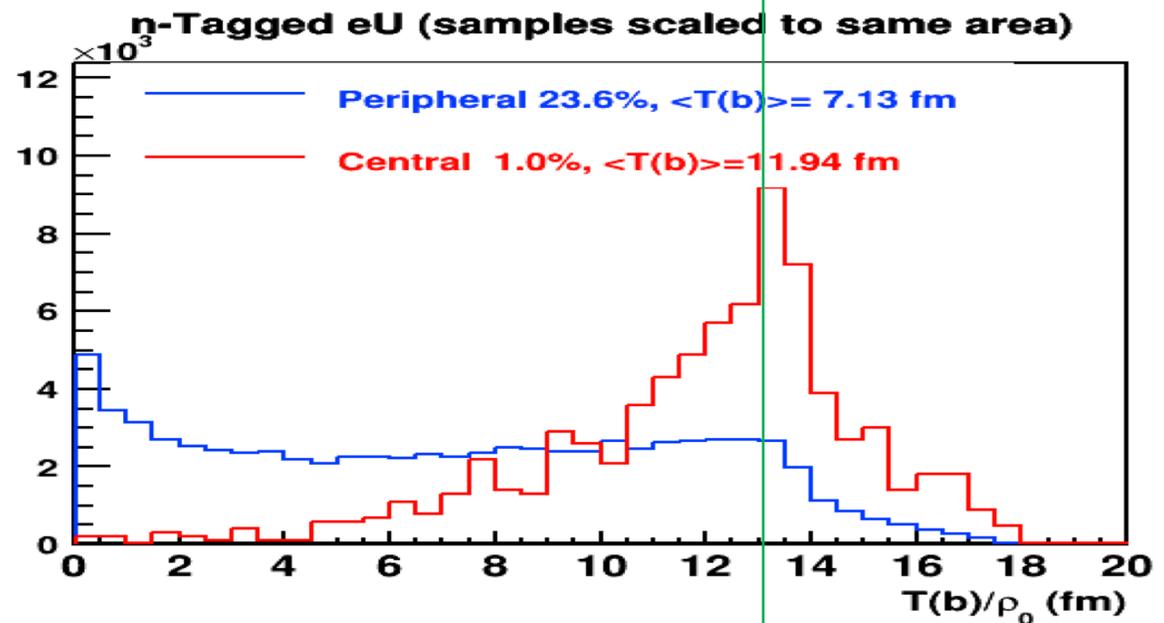
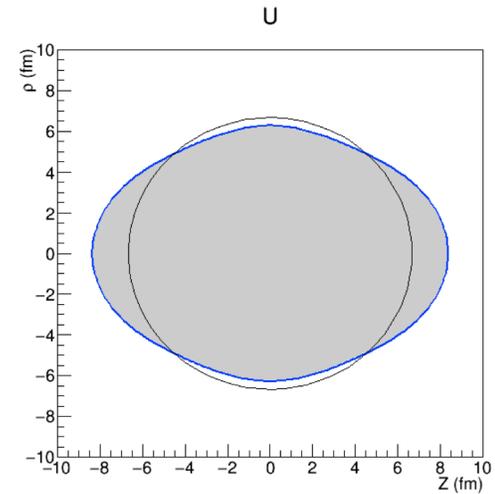
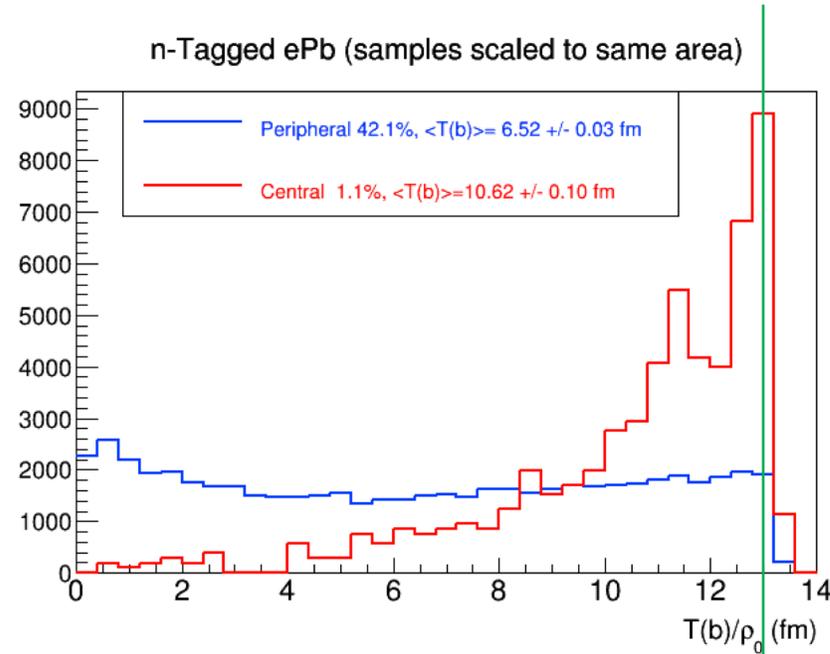
$$Q_s^2 \sim A^{1/3} / x^{0.3} \sim T(b) * (E_e E_A)^{0.3}$$

$$F_E = \left( \frac{\langle T(b) \rangle_{central}}{\langle T(b) \rangle_{minbias}} \right)^{10/3}$$

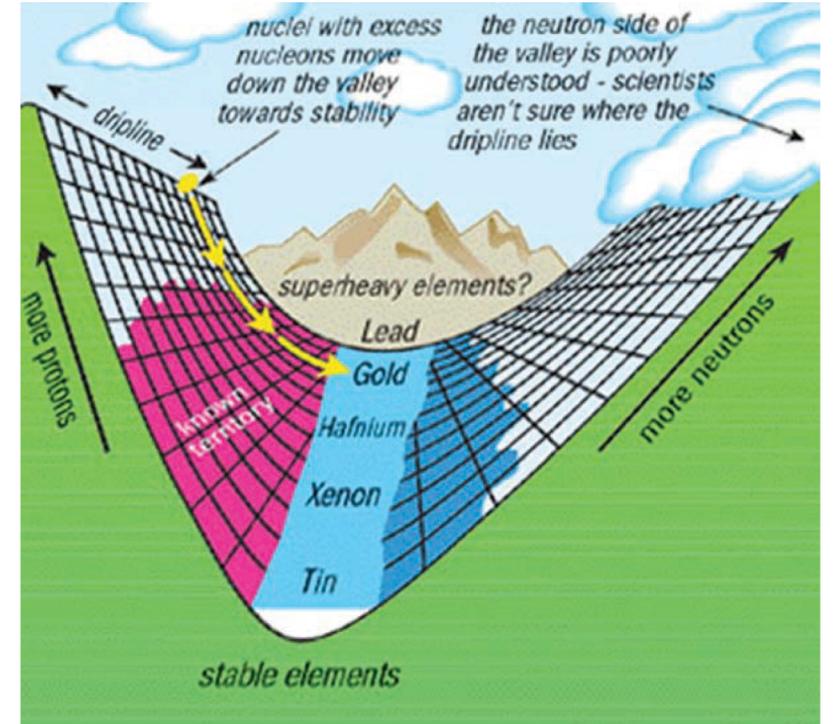
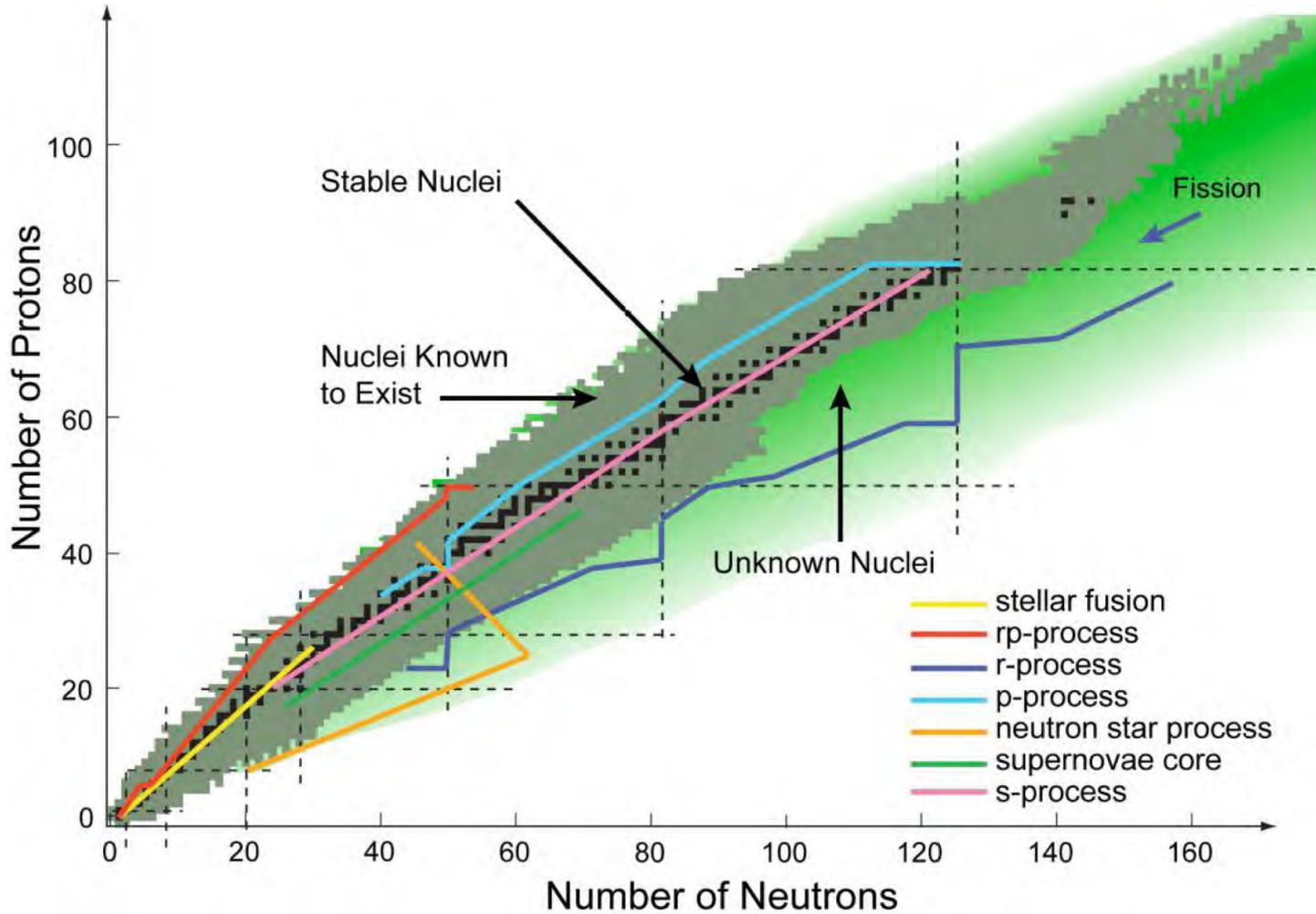
## T(b) for $^{208}\text{Pb}$ and $^{238}\text{U}$

- Using U-238 instead of Au or Pb further increases the equivalent energy, from a factor 4.8 to 6.7 at 1% yield.
  - U has slightly larger A than Pb
  - U is a deformed nucleus. Longer paths along the long axis
  - It is easier to align one axis in eA than two in AA
- A factor 6.7 is equivalent to running untagged Pb at 740 GeV/A
- Note that detection of all nuclear fragments, including spectator protons, is crucial

M. Baker, 2018



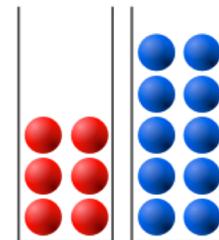
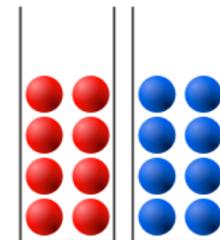
# Rare isotopes at an EIC



$A = 16$

Lower energy

Higher energy



Protons Neutrons

Protons Neutrons

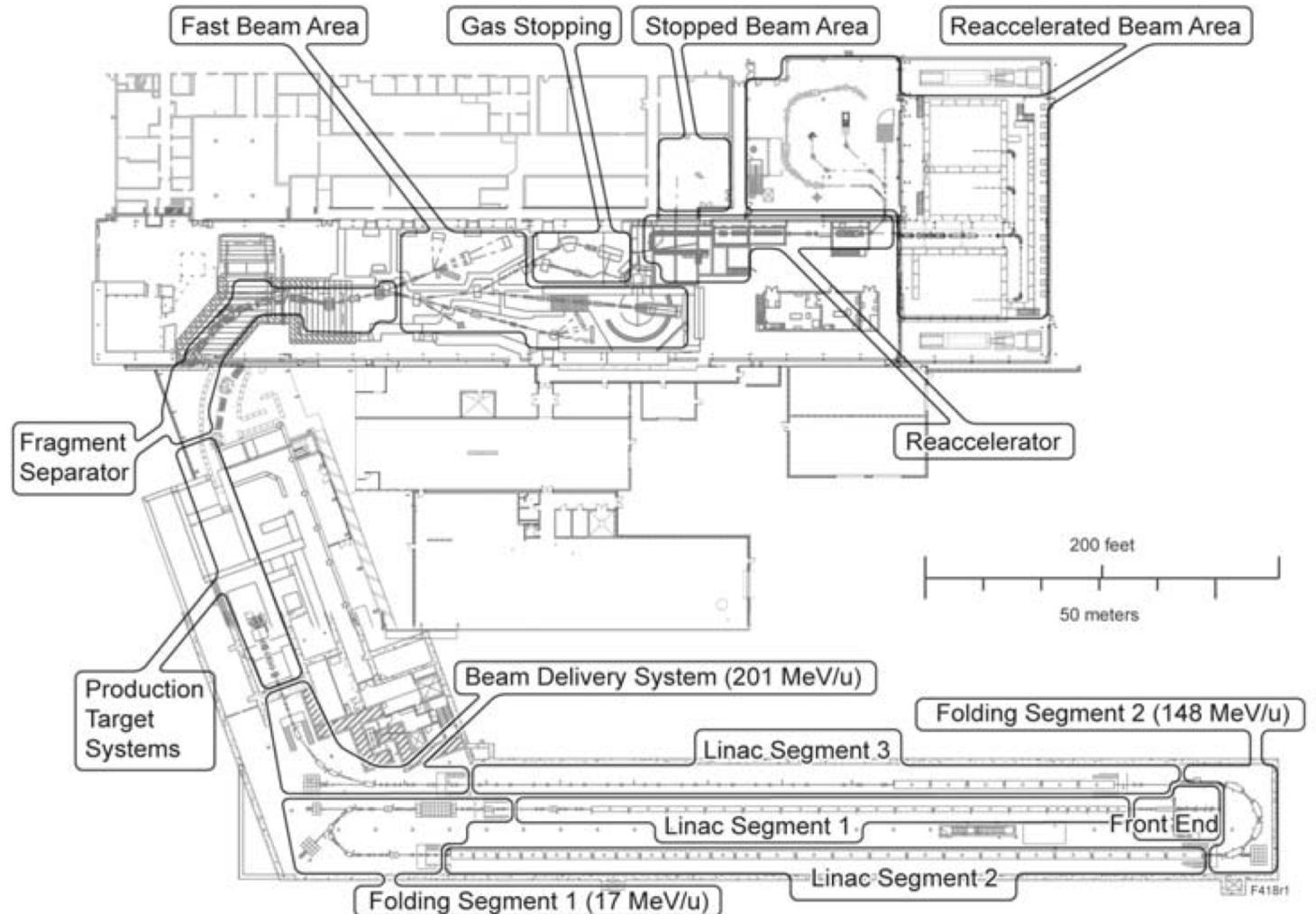
$$|N - Z| = 0$$

$$|N - Z| = 4$$

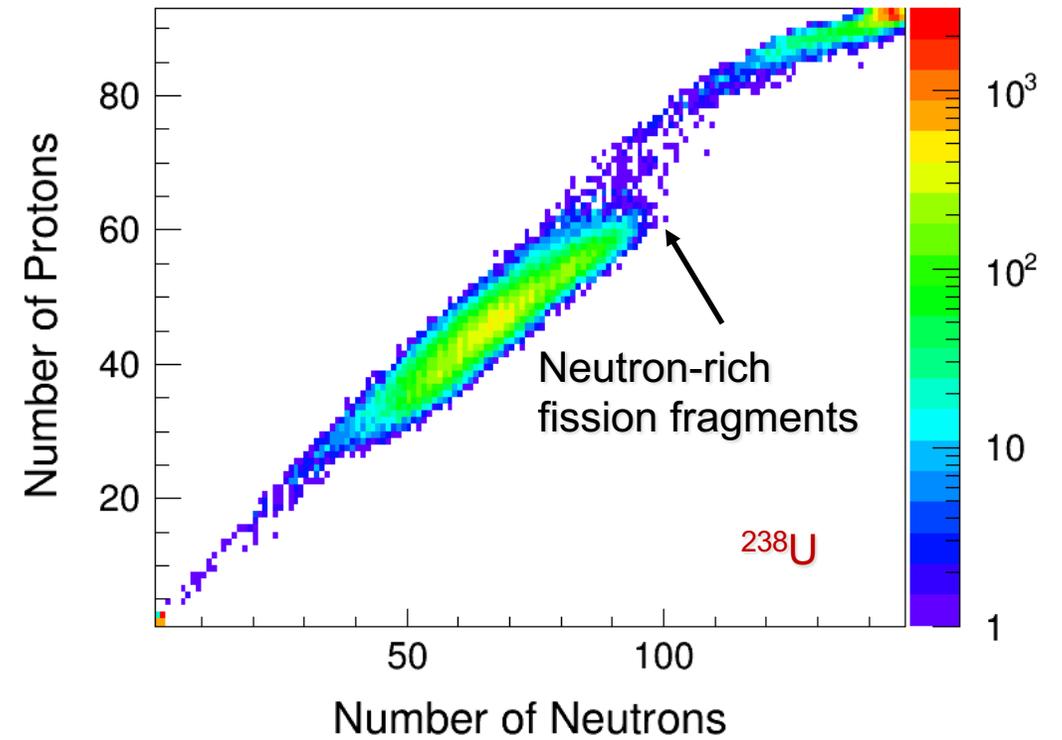
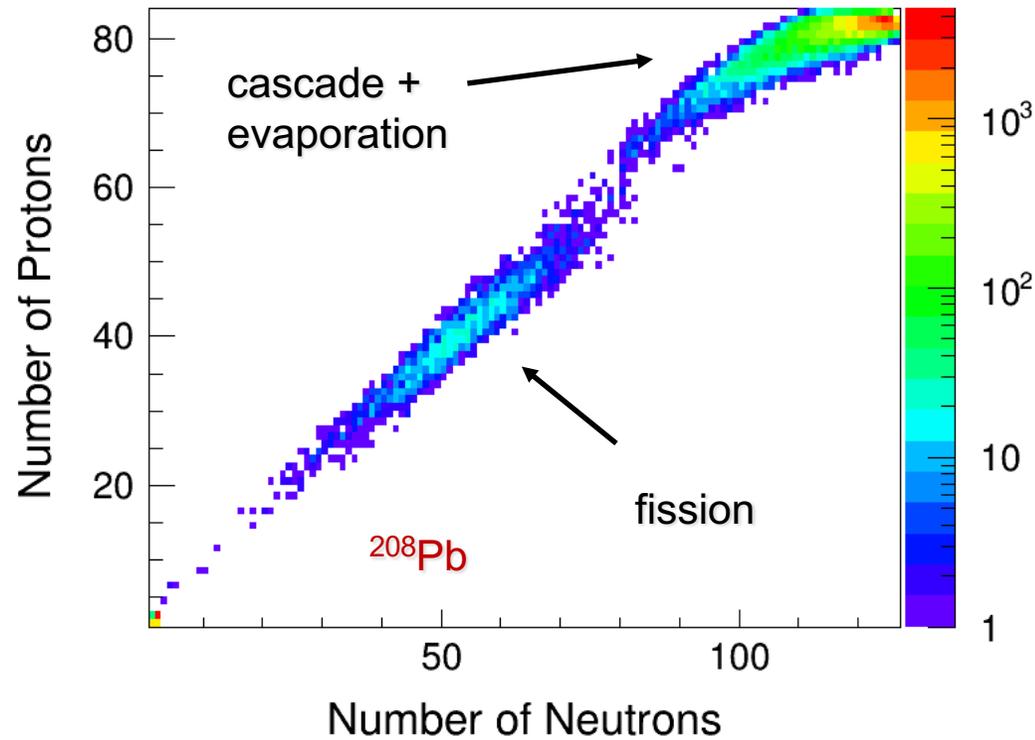
# Facility for Rare Isotope Beams (FRIB)



- Close to completion at MSU
- Will produce radioactive beams through in-flight projectile fragmentation followed by fragment separation in a downstream spectrometer
- Focus in on neutron-rich nuclei

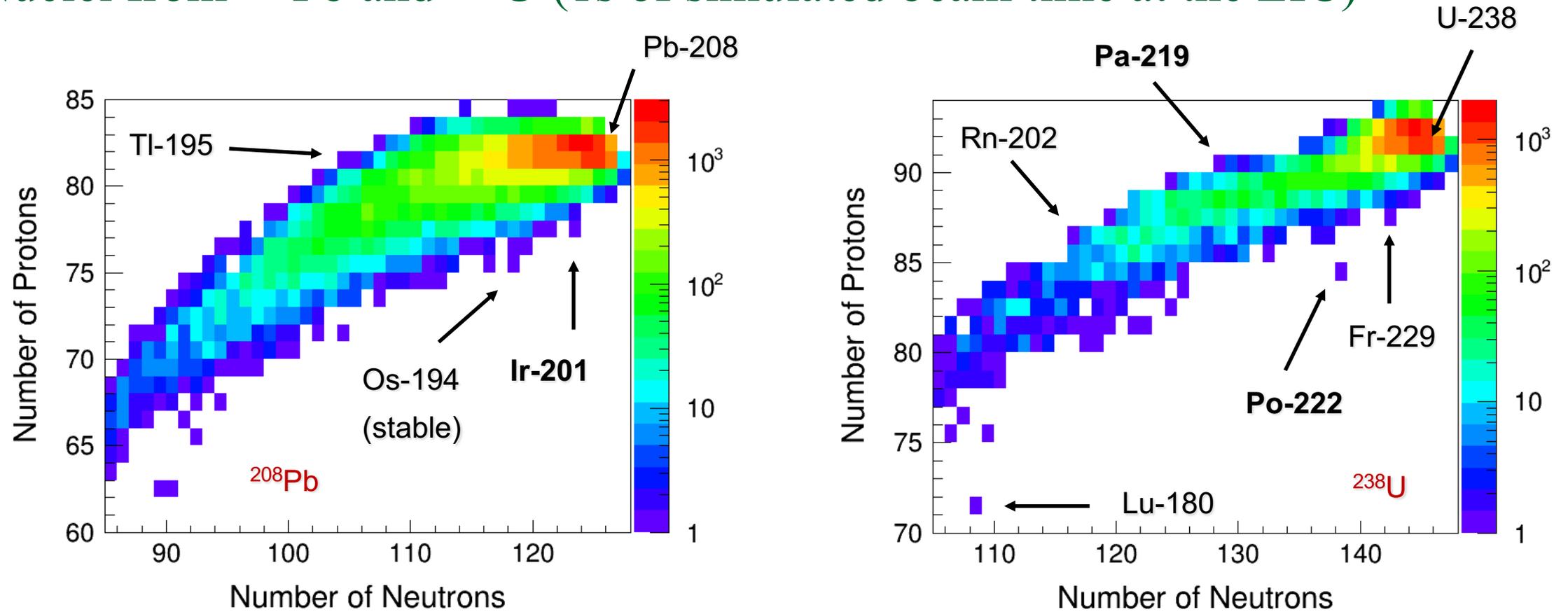


# Nuclei from $^{208}\text{Pb}$ and $^{238}\text{U}$ (1s of simulated beam time at the EIC)



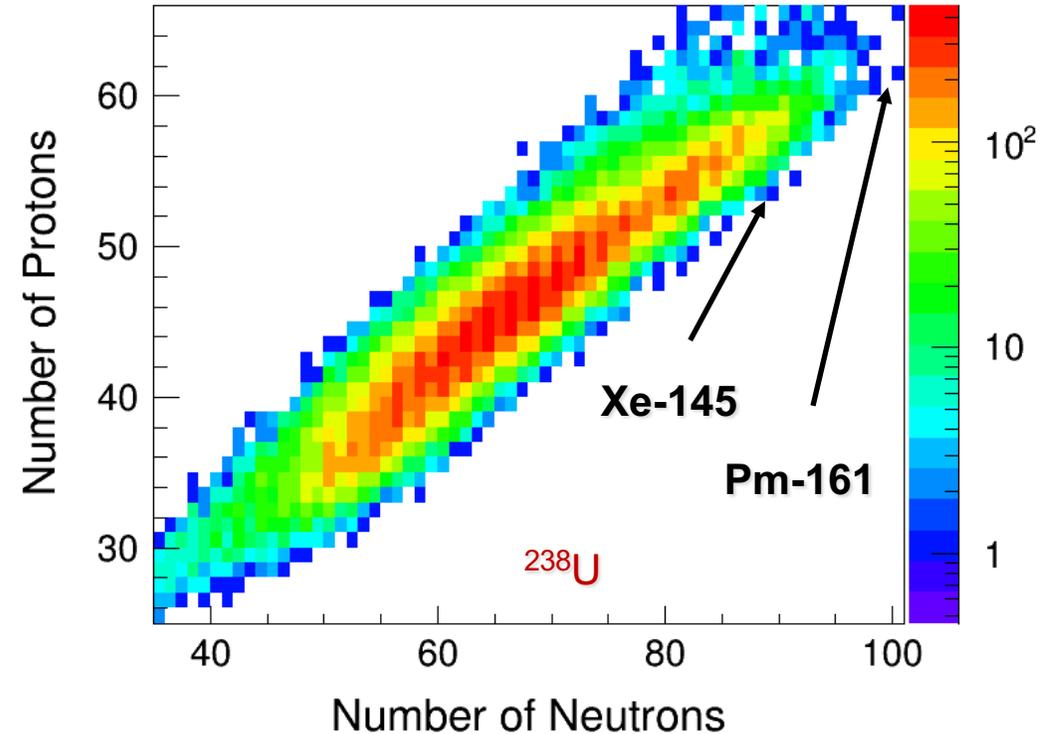
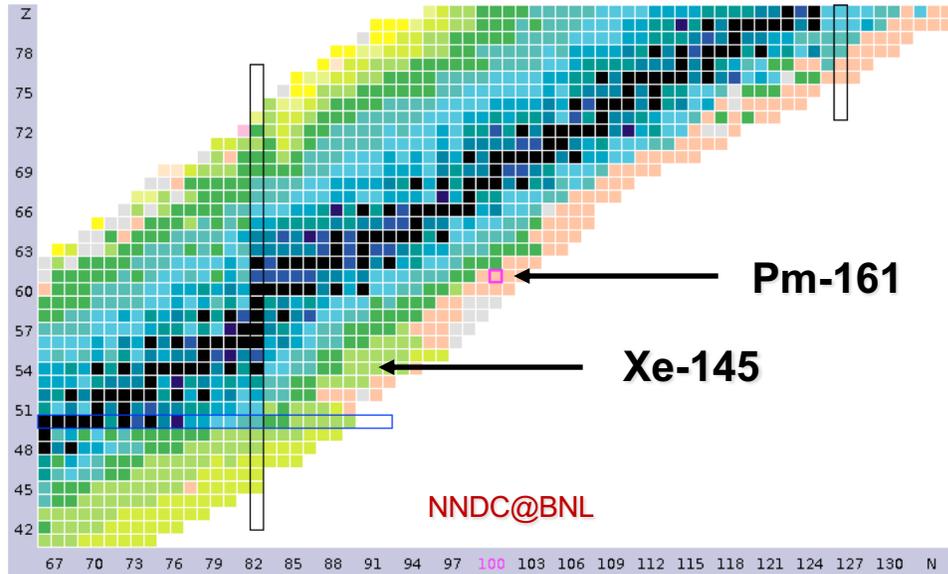
- $^{208}\text{Pb}$  (left) produces mainly heavy isotopes from evaporation
- $^{238}\text{U}$  (right) produces fewer, but heavier isotopes from evaporation. It also produces very neutron-rich fission fragments (medium-mass nuclei have fewer neutrons).

# Nuclei from $^{208}\text{Pb}$ and $^{238}\text{U}$ (1s of simulated beam time at the EIC)



- $^{208}\text{Pb}$  (left) produces mainly heavy isotopes from evaporation
- $^{238}\text{U}$  (right) produces fewer, but heavier isotopes from evaporation. It also produces very neutron-rich fission fragments (medium-mass nuclei have fewer neutrons).

How “exotic” are  $^{161}\text{Pm}$  and  $^{145}\text{Xe}$  (produced in 1s of simulated beam time)?



- Lifetime for  $^{161}\text{Pm}$ :  $>130$  ns
- Lifetime for  $^{145}\text{Xe}$ : 188 ms
- Both nuclei are very close to the most neutron-rich nuclei currently known – and in a spot where rates at FRIB are relatively low

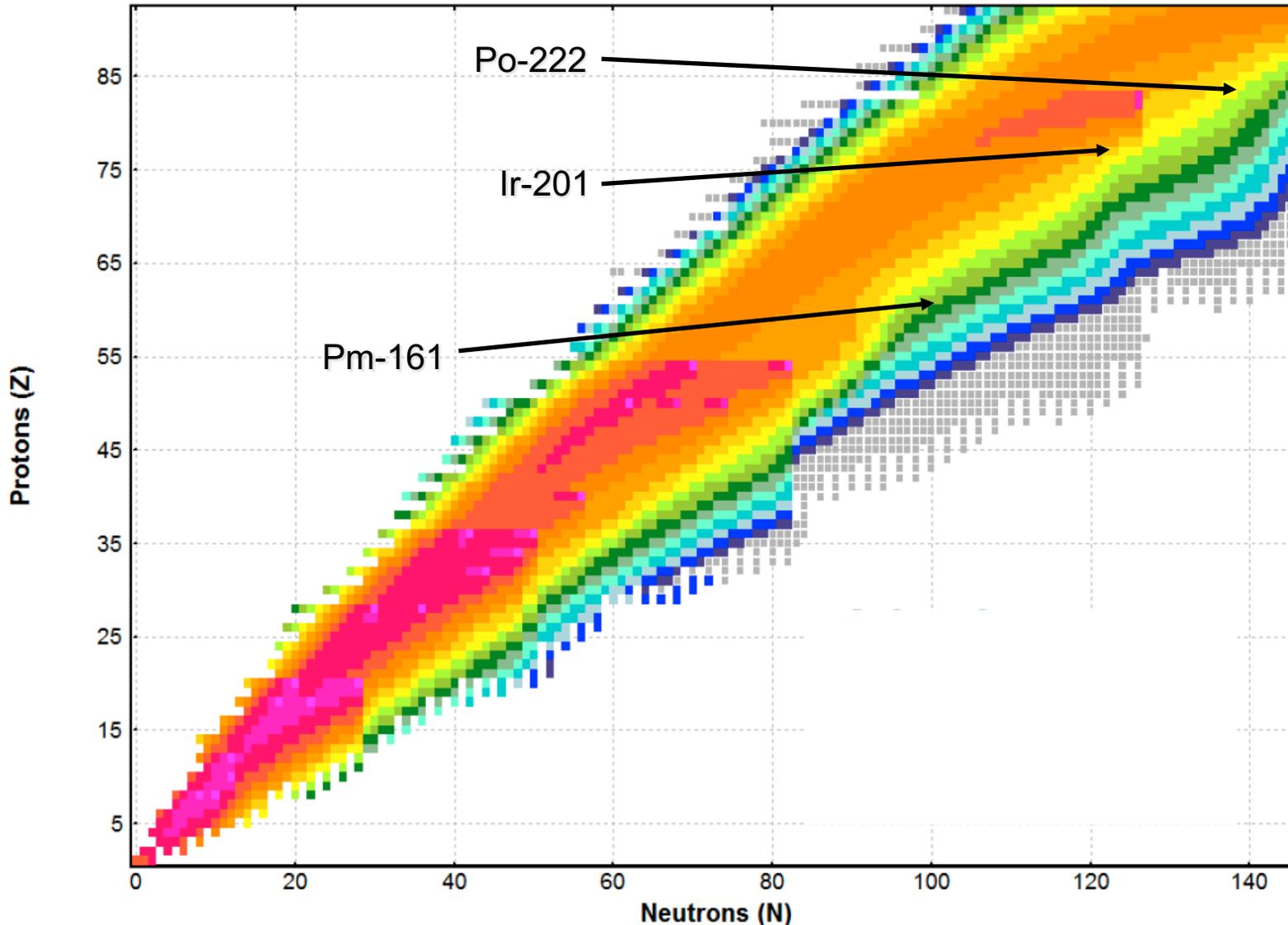
# Rates at the EIC?

## NSCL PAC35 rates (v.1.03)

[https://groups.nsl.msu.edu/frib/rates/nsl\\_pac35\\_rates.html](https://groups.nsl.msu.edu/frib/rates/nsl_pac35_rates.html) The rates are estimated based on

the EPAX 2.15 cross section parameterization for fragmentation and the LISE++ 3EER model for in-flight fission.

Primary beam intensities and energies have been used from the PAC35 beam list



## Needs further studies!

### A bold extrapolation

- Let's assume that we have one "yellow-orange" event per second, or  $10^8$  in a year, distributed over  $\sim 100$  isotopes of interest
- Further, let's assume that we can use the FRIB rate estimates for the extrapolation.

- Then, if we want to accumulate a total of 10,000 events for the isotopes of interest, in a year we can move from the orange-yellow to light green (although for  $^{161}\text{Pm}$  we are there in 1 minute)

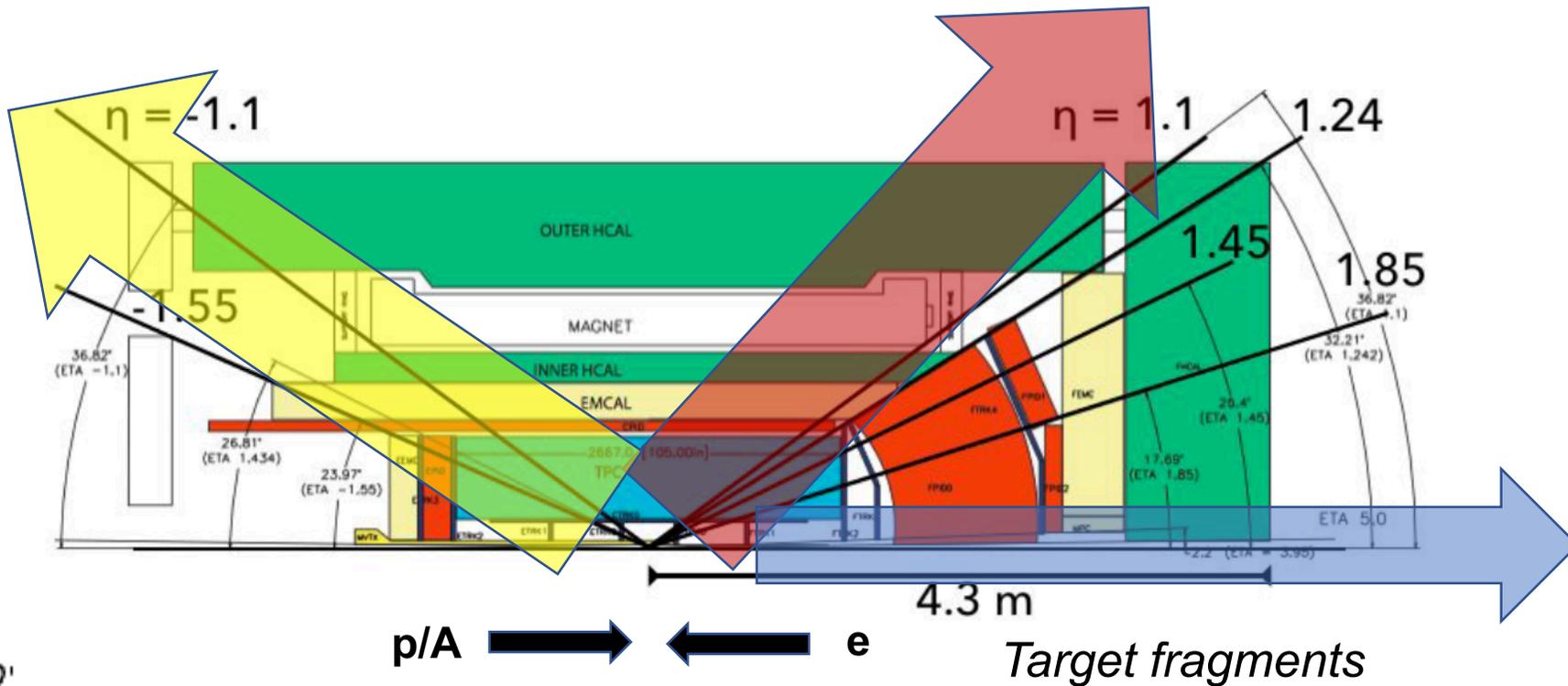
Thank you!

Backup

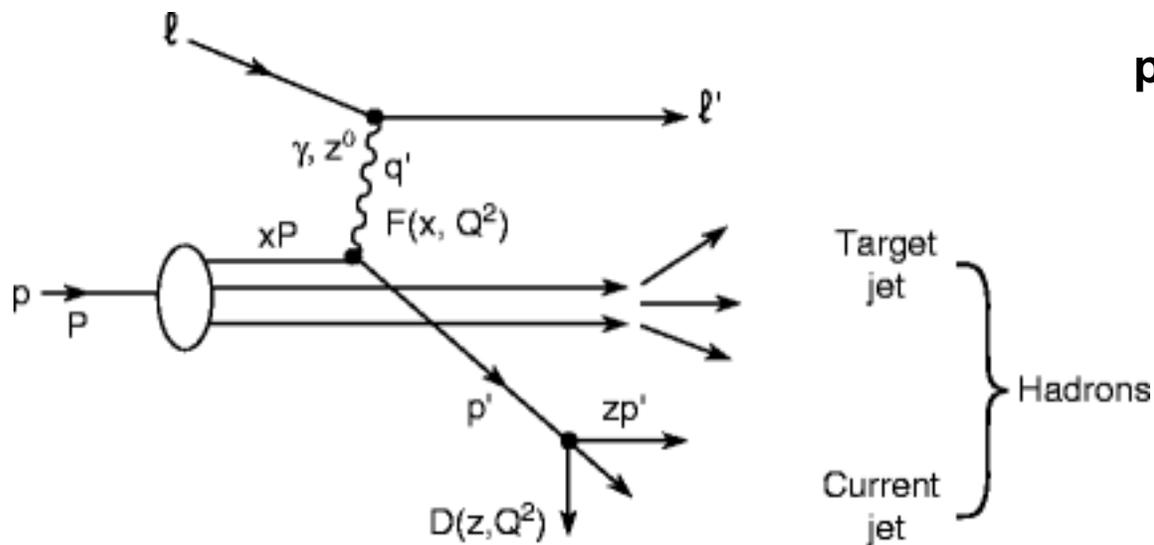
# What do we measure?

Current jet (or hadron)

Scattered electron



Lepton scattering on a proton



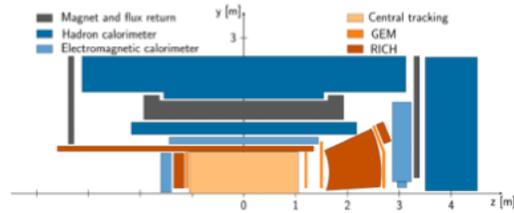
**Inclusive DIS:** only electron is detected

**Semi-Inclusive DIS (SIDIS):** electron and current jet (hadron) are detected.

**Exclusive reactions:** all particles are detected

# Detection of target fragments – forward spectrometer

2.6 Evolution of sPHENIX into an Electron-Ion Collider Experiment 33



”near” detection

”far” detection

Analyzing dipoles (green) and accelerator quadrupoles (red)

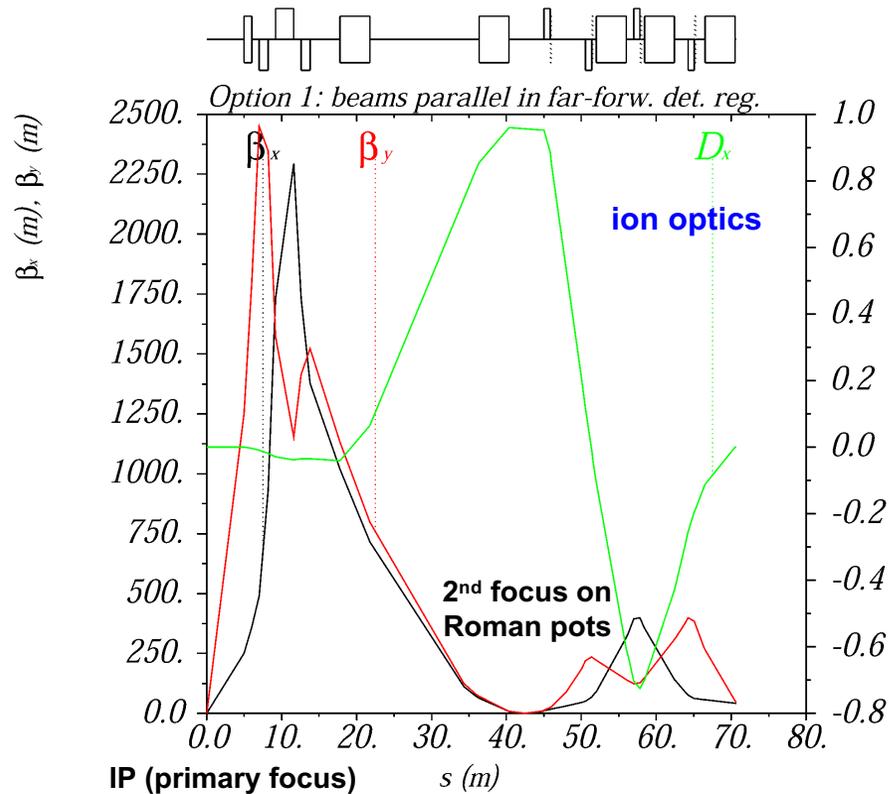
Detection space

high-res. ZDC ( $n, \gamma$ )

Forward hadron spectrometer  
Spectator protons

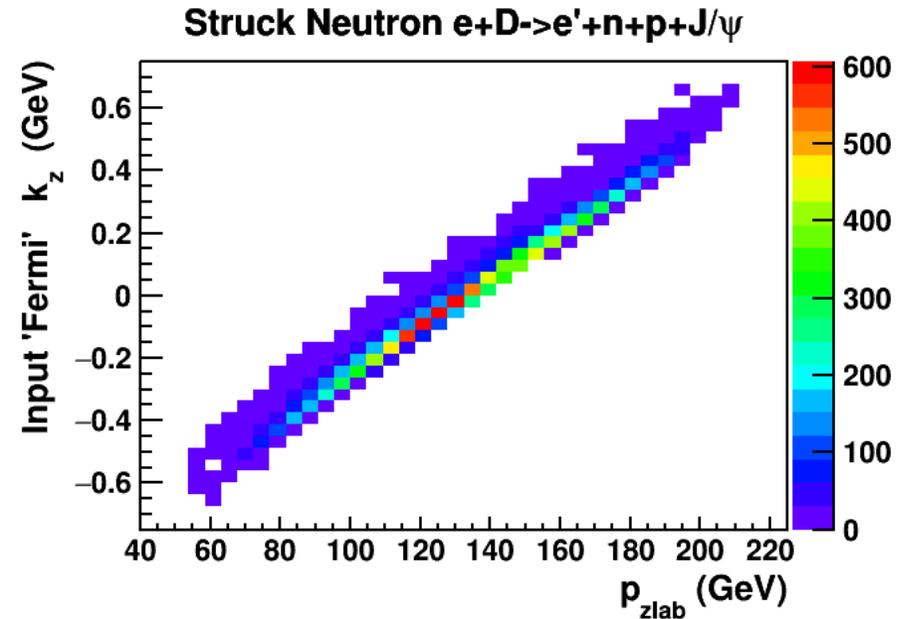
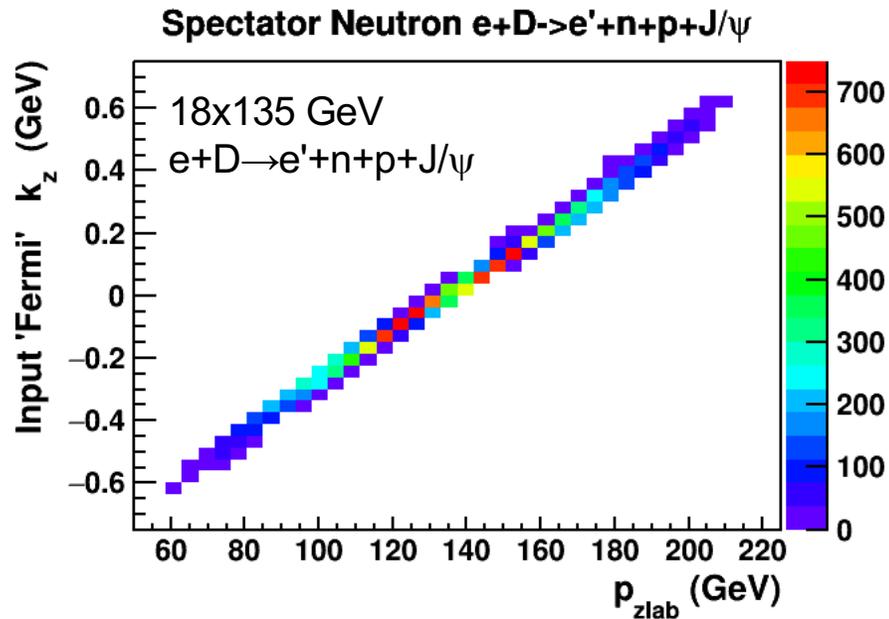
Recoil protons (roman pots at focal point with high dispersion)

$p/A$  →  
 $e$  ←



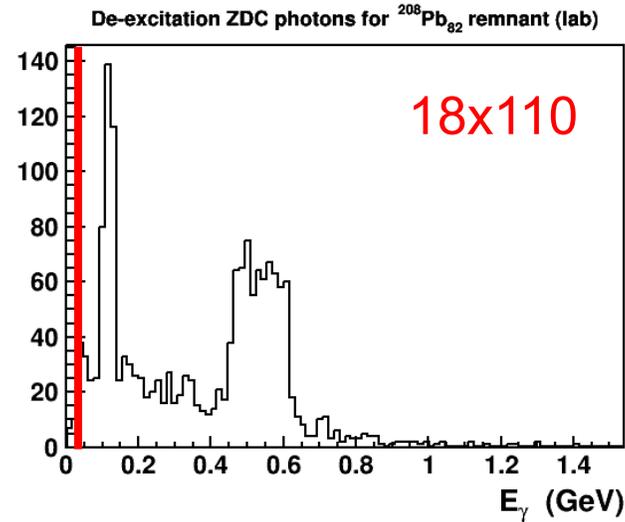
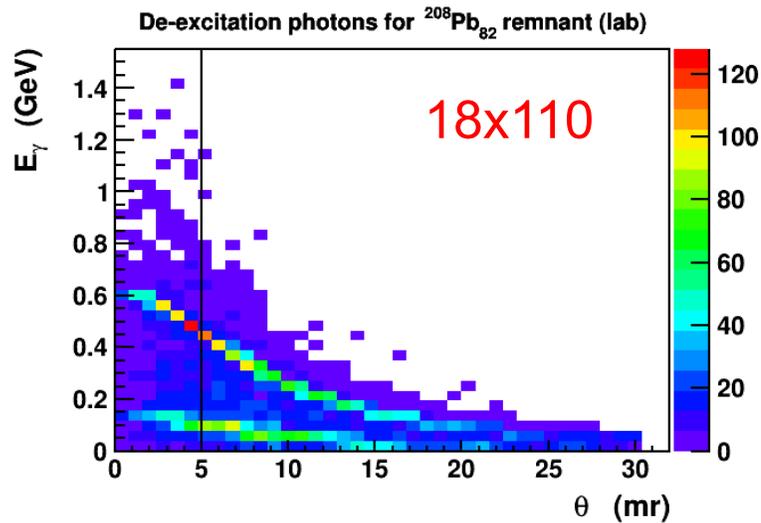
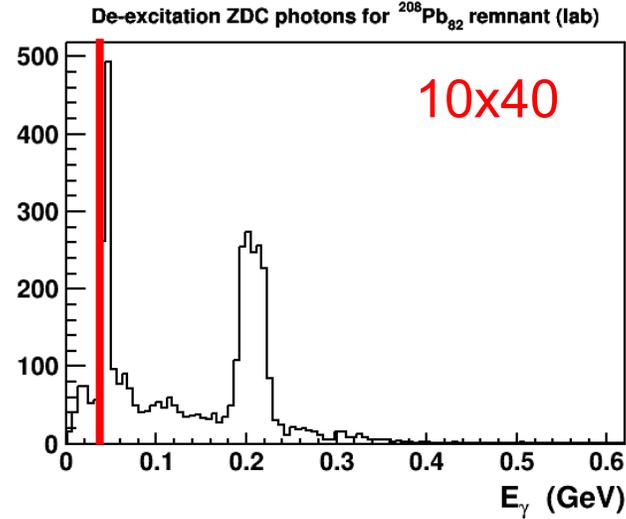
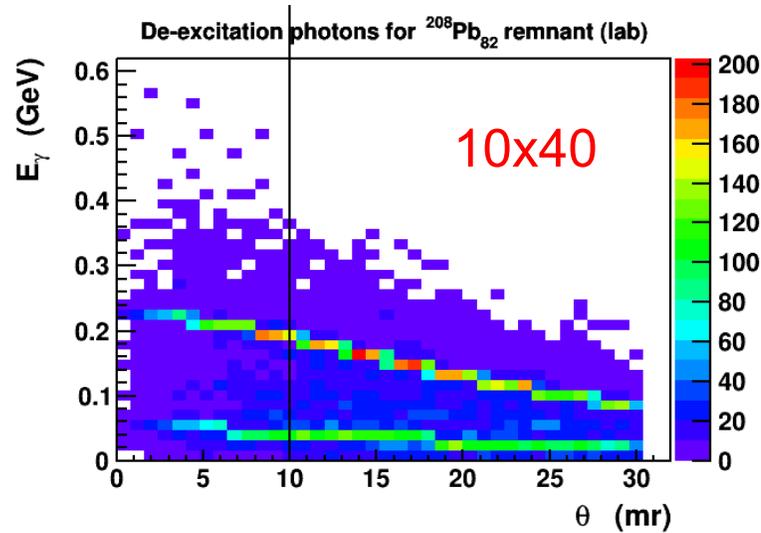
- Functionally, the forward detection is naturally separated into a ”near” and ”far” parts
- ”Near” detection:
  - *Goal*: off-momentum/rigidity particles or ones scattered at ”large” angles (high  $p_T$ )
  - *Requirement*: large magnet apertures
- ”Far” detection (can be after a crab cavity):
  - *Goal*: small-angle particles with momentum/rigidity close to that of the beam
  - *Requirement*: large dispersion and small beam size

# Effect of boosting $k_z$ to the lab frame



- The boost makes precision measurements of nucleons moving along the z-axis easier
  - It also makes precision gamma spectroscopy of photons from nuclear de-excitations in heavier nuclei possible.
- However, as noted earlier, this is also the reason why hadronic reconstruction methods do not work as well for eA as they do for pA.

# Photons from $^{208}\text{Pb}_{82}$ in lab frame



Detailed studies ongoing.

It is clear that  $\gamma$ 's will be needed for low  $|t|$ !

w/ Morozov, Hyde, Turonski et al.

# Density fluctuations from incoherent diffraction

- Incoherent diffraction probes the variance of the density
  - Calculations at  $x \sim 10^{-3}$  show "lumpiness" in the gluon distribution and suggest sensitivity to model assumptions
- Could we in the future see the imprint of correlations from nucleonic degrees of freedom at the partonic level?
  - Experimentally straightforward, but a clear interpretation needs much more progress on the theory side

