

Bayesian Monte-Carlo Evaluation Framework of Differential and Integral Data

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Presentation Outline

- Extant methods (linear approximation)
- Bayes theorem for generalized data
- Schematic diagram
- Framework Demonstration on U-233 integral and differential data
 - Comparison to linear approximation
- Summary and outlook

Summary of conventional evaluation of diff. and integral data:

- Extant R -matrix resolved resonance range (RRR) evaluations assume normal probability density functions (PDFs) and use Newton-Raphson method to minimize χ^2 of R -matrix and differential data only.
- R -matrix parameters are subsequently further optimized to improve the fit of neutron transport simulations of integral benchmark experiments (IBEs), by a generalized linear least squares (GLLS) of the SAMINT module of SAMMY.
- The 2 methods above are based on Bayes theorem but assume all PDFs are normal (Gaussian) and use a linear approximation for IBEs.
- The proposed framework would remove both approximations
- Similar MC methods exist for optimization of TENDL optical model cross sections
 - P. Helgesson, H. Sjostrang, A.J. Koning, J. Ryden, D. Rochman, E. Alhassan, S. Pomp, Prog. Nucl. Ener. 96 (2017) 76-98
- MC method for fitting RRR to diff. data alone announced in the AZURE R -mat. Code
- SAMPLER module in SCALE randomly perturbs c.s.'s and geometry of IBEs

A general form of Bayes theorem

- Model parameters, data, and any model defect treated on the same footing:
 - G. Arbanas et al., “Bayesian Optimization of Generalized Data”, CW2017, EPJ-N, 4 (2018) 30, <https://doi.org/10.1051/epjn/2018038>
- Uses *posterior* expectation values of constraints relating model, data, and any defect
- Its likelihood function is an exponential function of constraints; prior may be any PDF.
 - Sergio Davis, “Exponential Family Models from Bayes’ Theorem under Expectation Constraints” (2016) <https://arxiv.org/abs/1503.03451>
- Constraints could be imposed selectively on posterior expectation values of
 - 1st moment i.e. mean values of $T(P) - D - \delta$
 - 2nd moment i.e. covariances of the above
 - Conventional Bayes’ theorems is a special case where both constraints are set to 0 (next slide)
- *cf.* conventional Bayes’ theorem: posterior exp. values determined by priors & model
- Conventional form of Bayes’ assumes normal PDFs: χ^2 -minimization
 - Demonstrated to be inferior to using MCMC: G.B. King et al. Phys. Rev. Lett. 122, 232502 (2019)
 - <https://doi.org/10.1103/PhysRevLett.122.232502>

Linear approx. in the RRR is used twice: for $\Delta\sigma$ and Δk_{eff}

- $\Delta\sigma$ from R-matrix resonance parameter cov. matrix in ENDF File 32 (e.g. SAMMY):

$$\langle \delta\sigma_i \delta\sigma_j \rangle = \sum_{n,m} \frac{\partial\sigma_i}{\partial u_n} M_{nm} \frac{\partial\sigma_j}{\partial u_m}$$

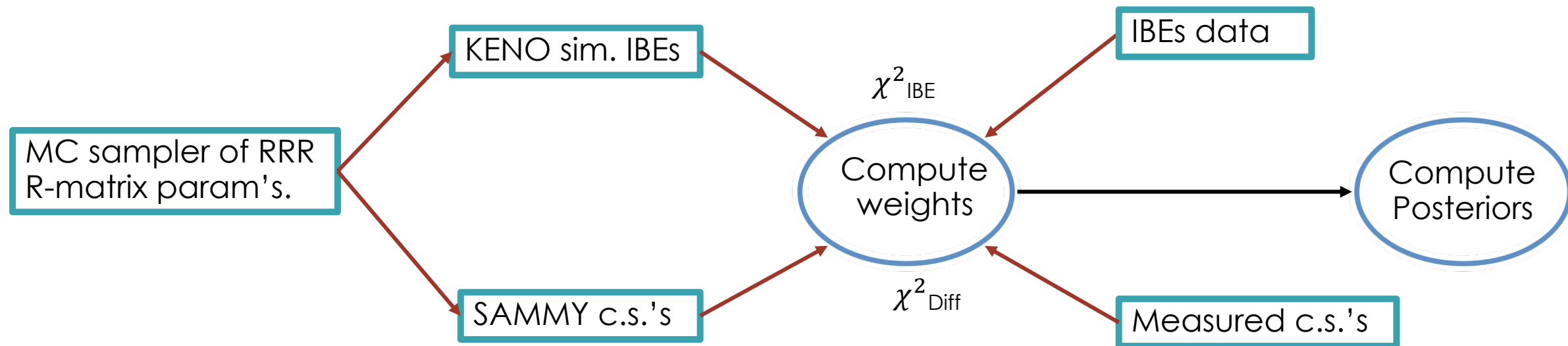
- Δk_{eff} from the covariance matrix of cross sections computed above (e.g. TSURFER)

$$\mathbf{C}_{\text{kk}} = \mathbf{S}_{\text{k}\alpha} \mathbf{C}_{\alpha\alpha} \mathbf{S}_{\text{k}\alpha}^T \quad \mathbf{S}_{\alpha} = \frac{\alpha}{R} \frac{\partial R}{\partial \alpha}$$

- Both linear approximations need to be accurate to obtain accurate results

Bayesian MC optimization framework overview

- Flexible: IBEs or Diff. Data alone could be analyzed independently or simultaneously



- MC sample weight takes into account agreement with differential and IBE data
 - Initially assuming normal PDFs:

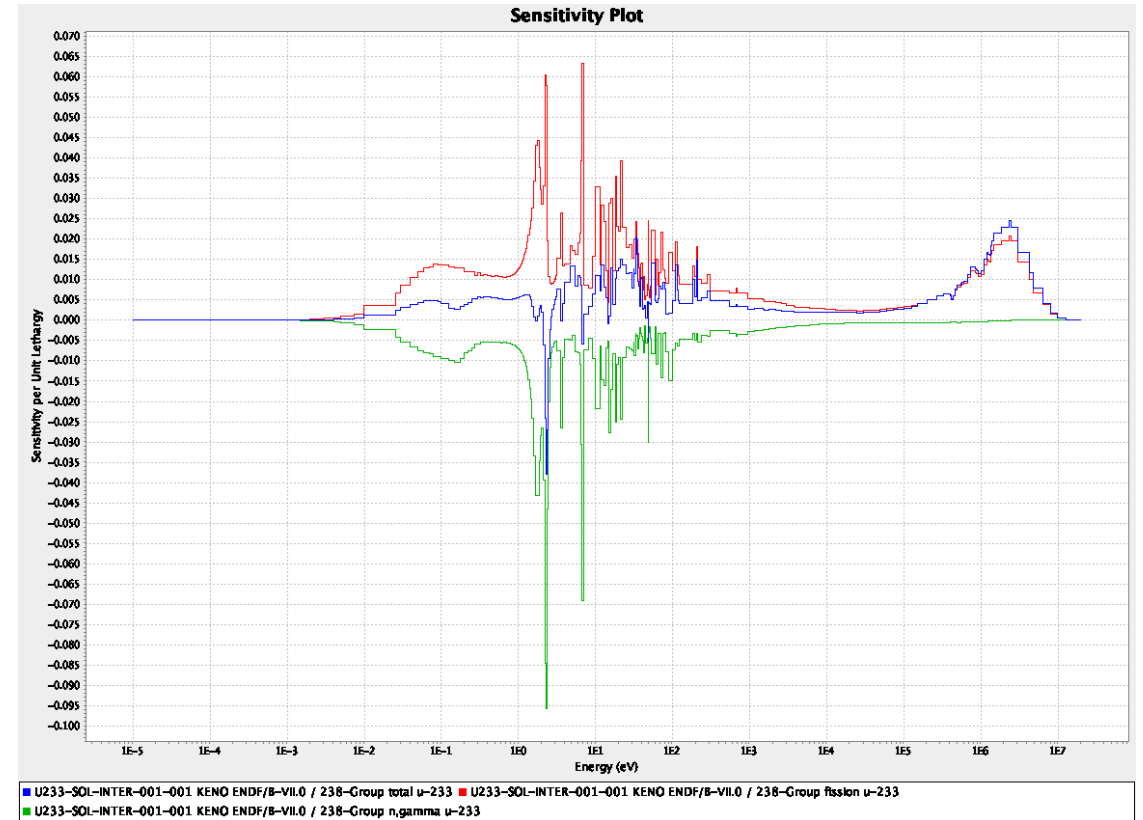
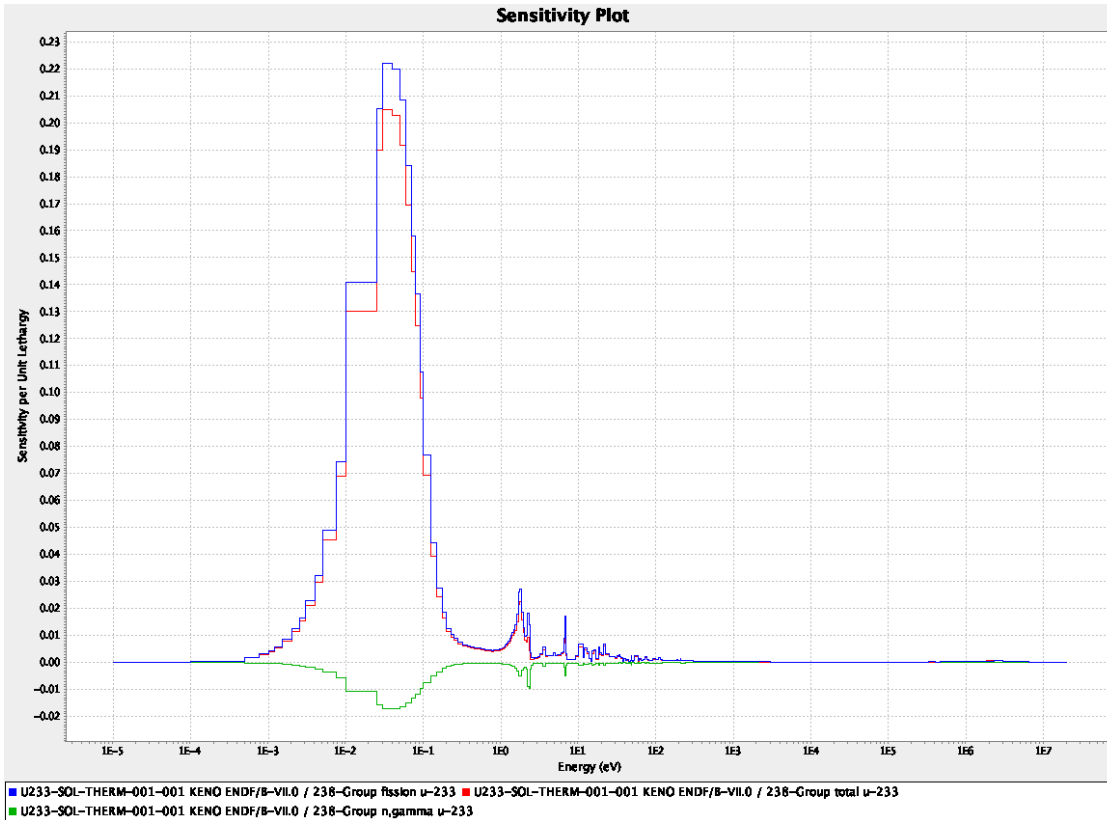
$$w = \exp[-\chi^2/2] \quad \chi^2 = \chi^2_{\text{IBE}} + \chi^2_{\text{Diff}}$$

- Large number of R-matrix param's. requires Metropolis-Hastings MCMC method
 - For MC random sampling to arrive at the posterior PDFs of parameters

Application to U-233

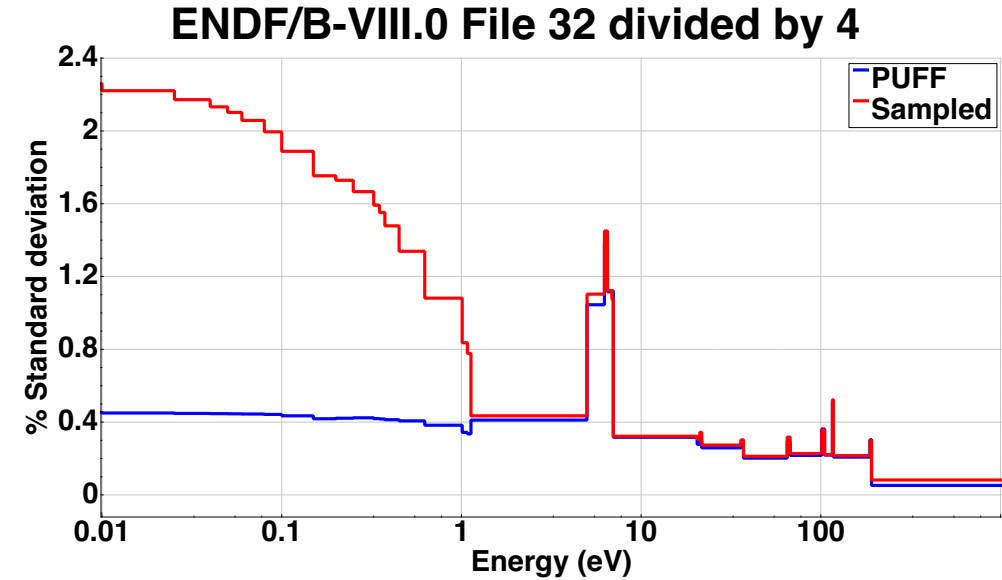
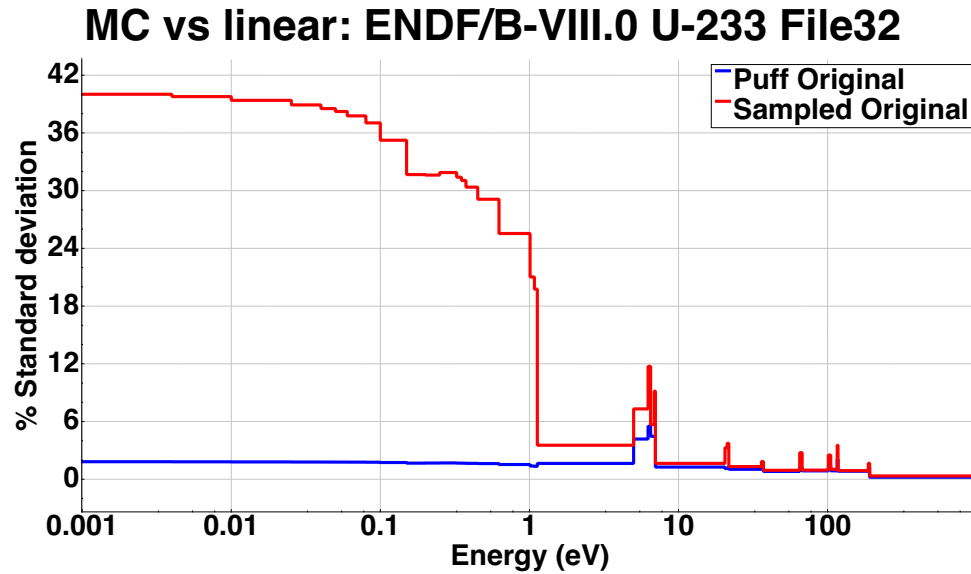
- 1000 randomly perturbed resonance parameter sets created by sampling File 32
- For each set calculate k_{eff} for U233-SOL-{THERM, INTER}-001-001 (KENO code), and then calculate k_{eff} mean values and uncertainties,
 - compare to corresponding TSUNAMI-IP's
 - Compare to measured IBE data
- For each set calculate differential cross sections using the SAMMY and then calculate mean values and uncertainties (transmission, fission)
 - Compare to SAMMY File 32 calculation, assuming it can be done
 - Compare to differential data (transmission, total, fission) by K. Guber (ORNL)

Cross Section Sensitivity of U233-SOL-{THERM, INTER}-001-001

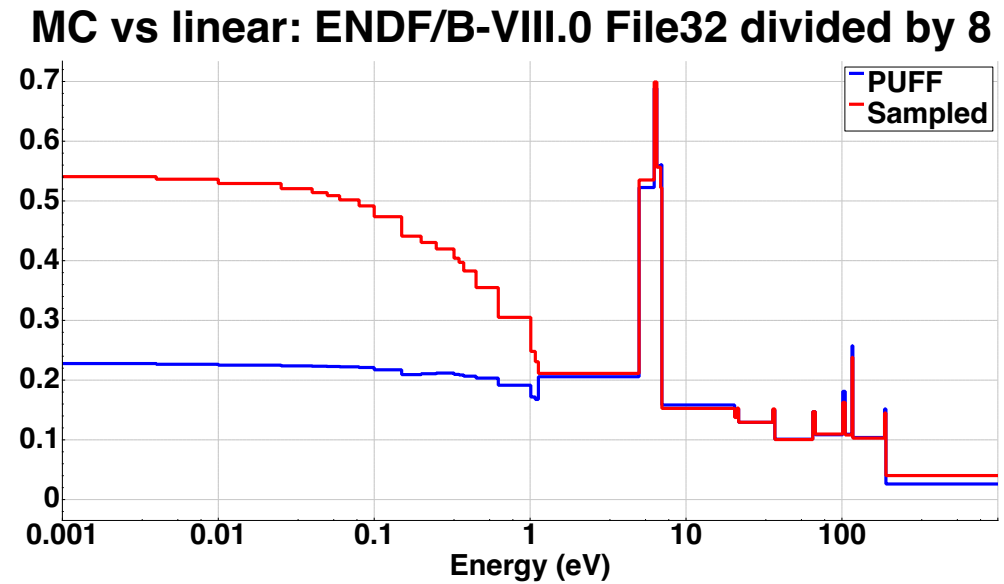


Uncertainties from MC vs. lin. approx. of ENDF/VIII.0 U-233 File 32

- Linear approximation significantly underestimates uncertainties encoded in File 32



- $\rightarrow \Delta k_{\text{eff}} (\text{MC}) \gg \Delta k_{\text{eff}} (\text{linear approx.})$
- MC and linear approx. reach similar uncertainty in the RRR for File 32/8
- The effect of large uncertainty on sub-threshold resonance seen below 1 eV



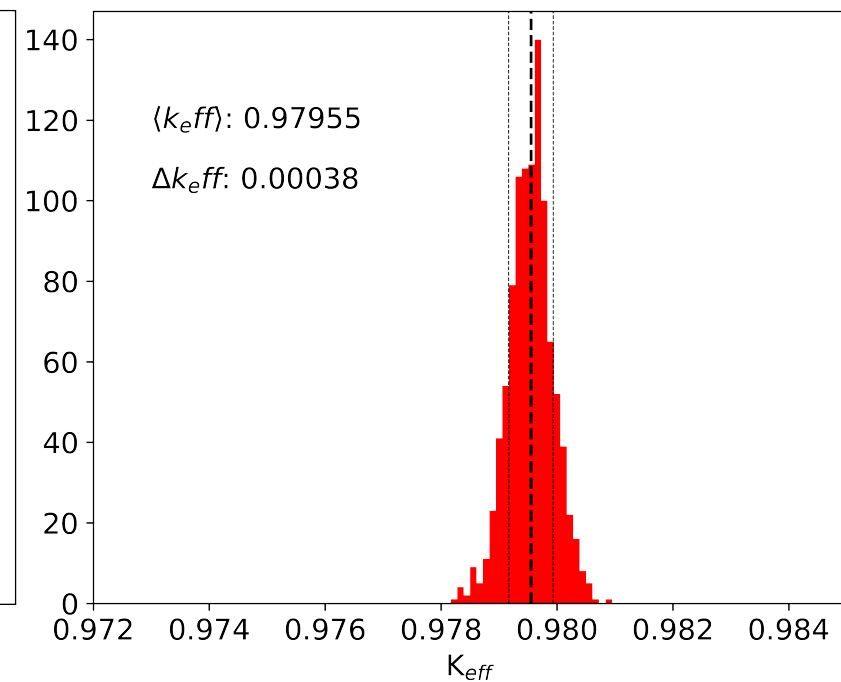
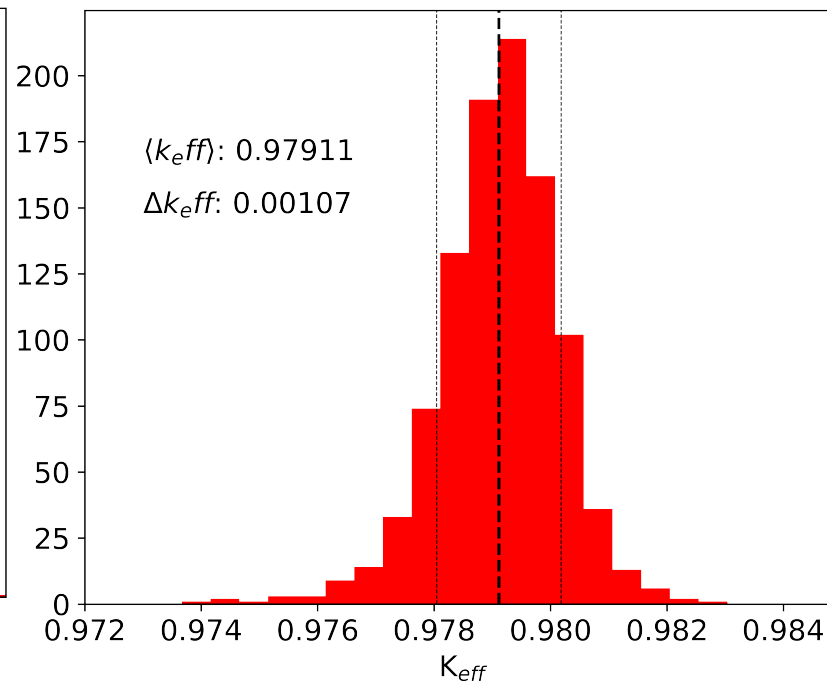
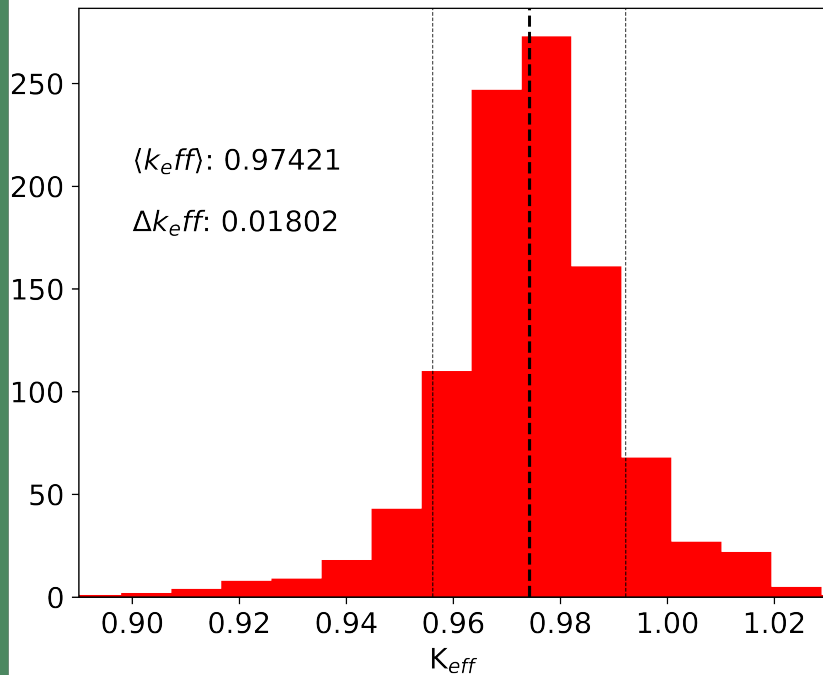
MC vs. linear approx.: Δk_{eff} of U233-SOL-INTER-001

- For U233-SOL-INTER-001 consistency between MC and linear approx. is achieved after dividing the U-233 ENDF/B-VIII.0 File 32 by 8

MC samples from: ENDF/B-VIII.0 File 32,

File 32 divided by 4,

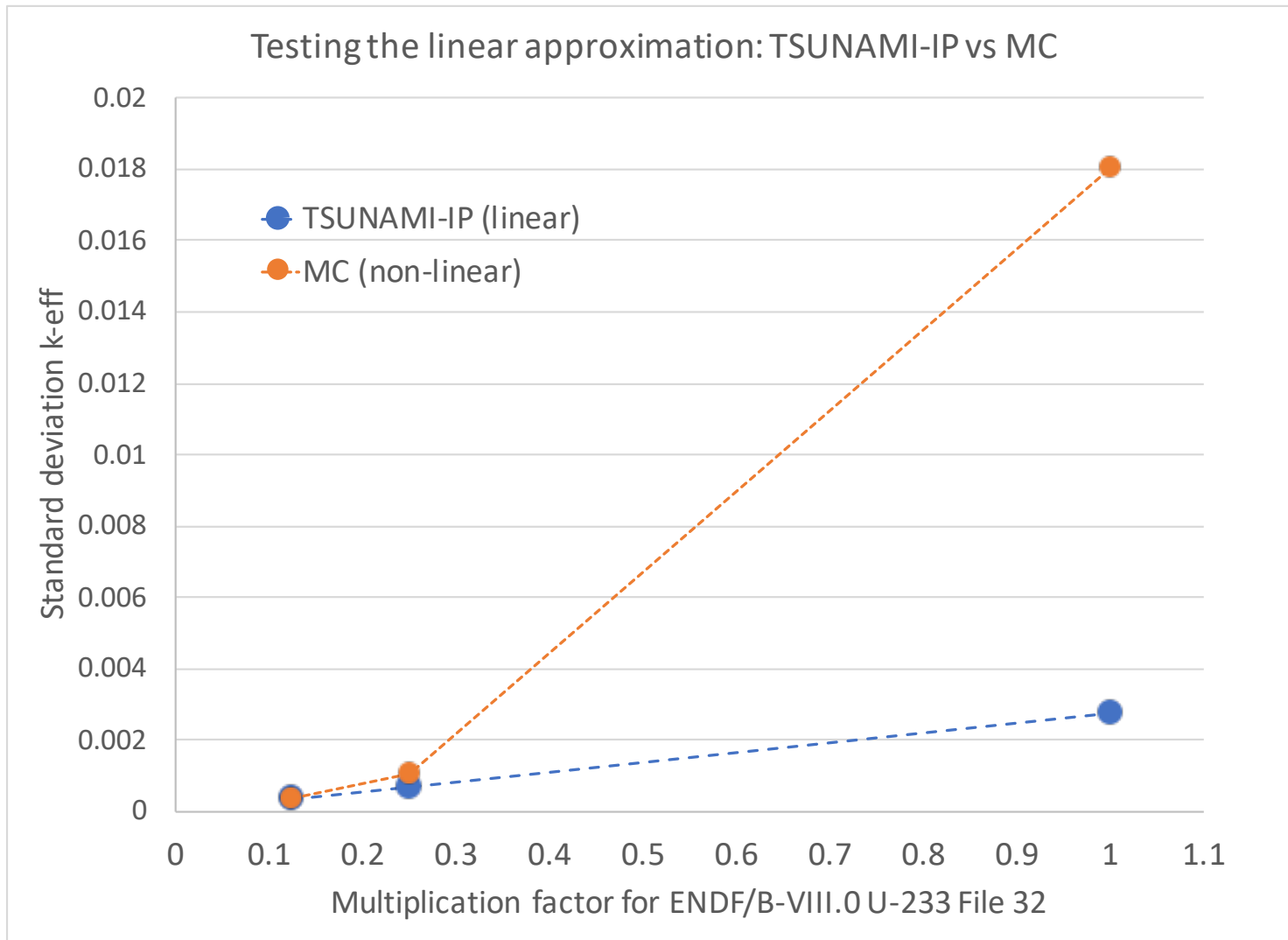
File 32 divided by 8



- k_{eff} uncertainty is decreasing significantly faster than linear scaling would imply

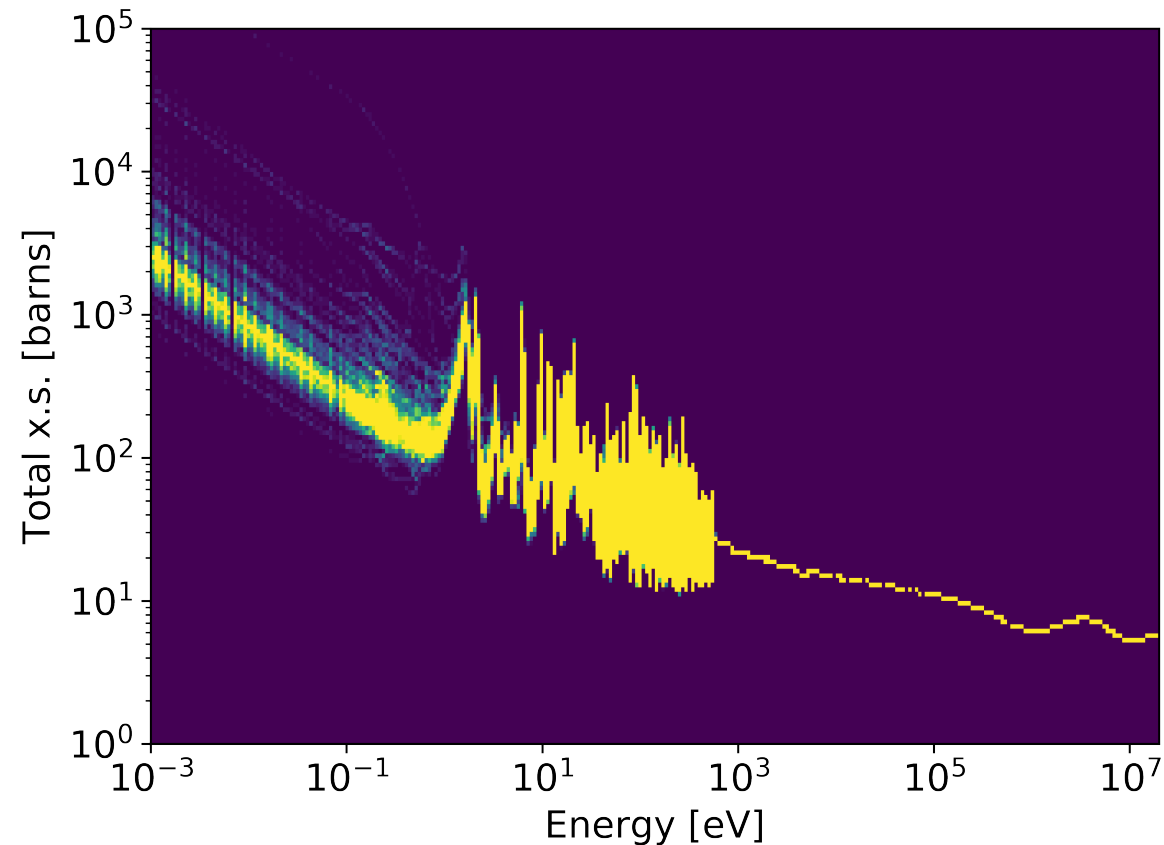
MC vs. linear approx. for Δk_{eff} of U233-SOL-INTER-001-001

- MC reveals large deviation from non-linearity for ENDF/B-VIII.0 U-233 File 32



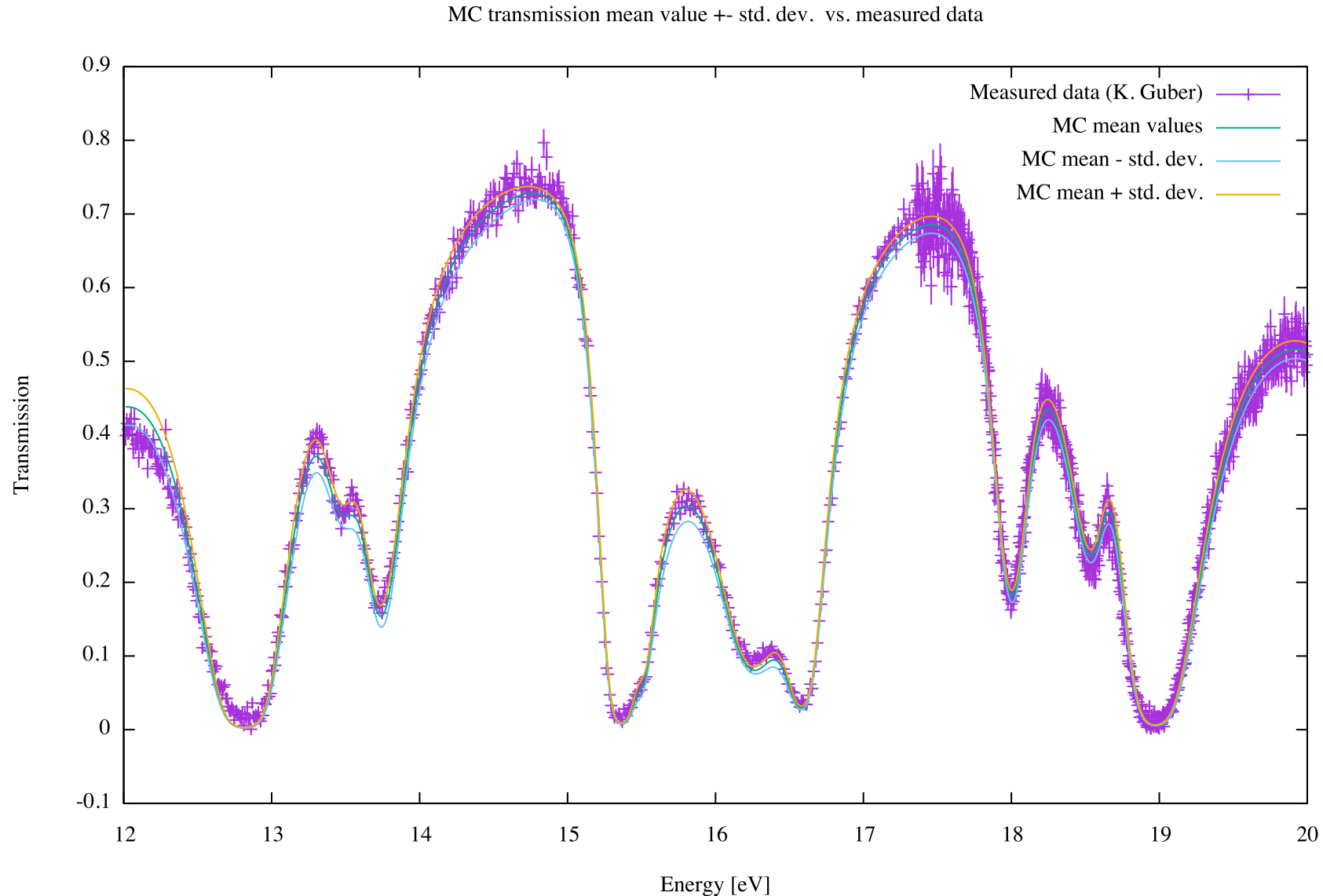
MC total cross section heat map: large variation below 1 eV

- Extremely large thermal cross sections can occur when MC random perturbations of subthreshold resonance energy fall near 0 eV
 - Due to large uncertainty of subthreshold resonance in the ENDF/B-VIII.0 U233 File 32.



Transmission measured data vs. MC ensemble

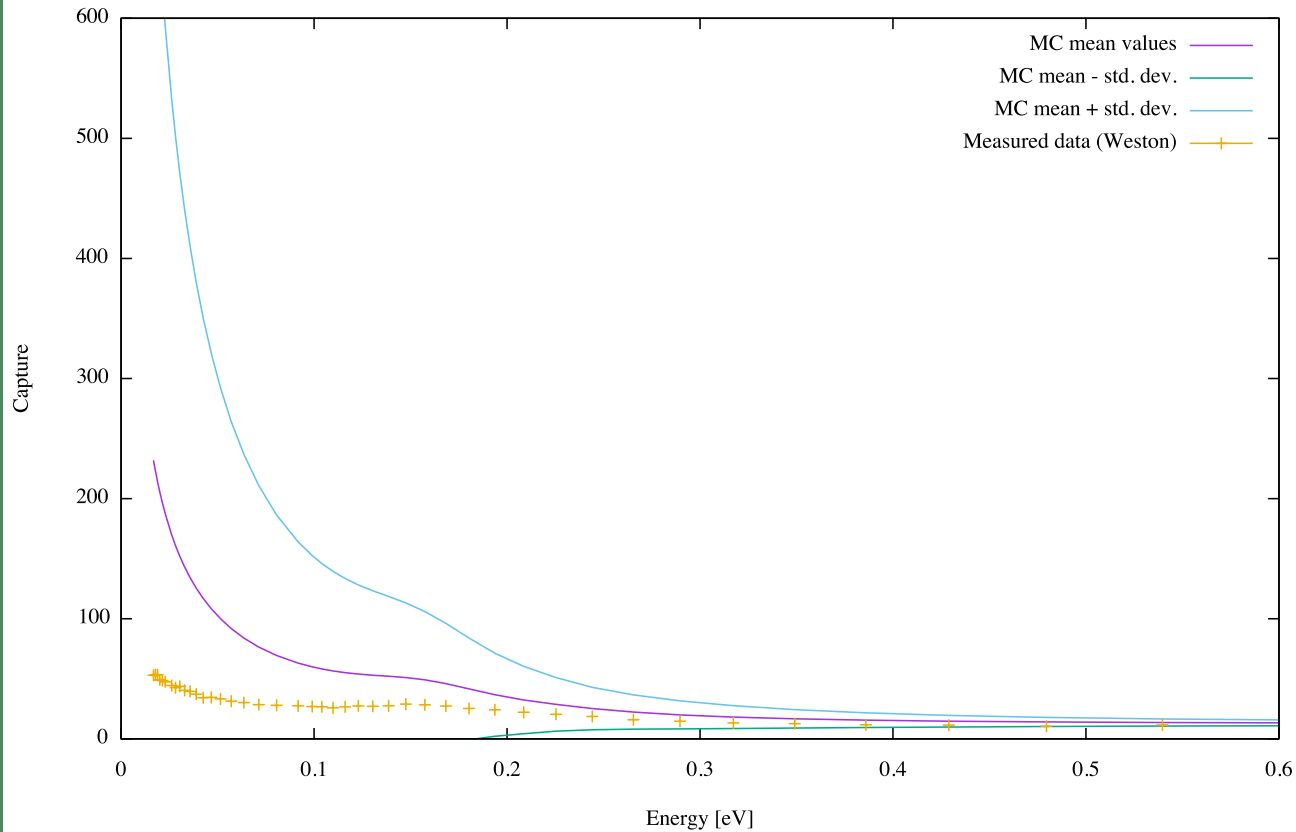
- Agreement above 12 eV is consistent with agreement seen with PUFF on slide 10



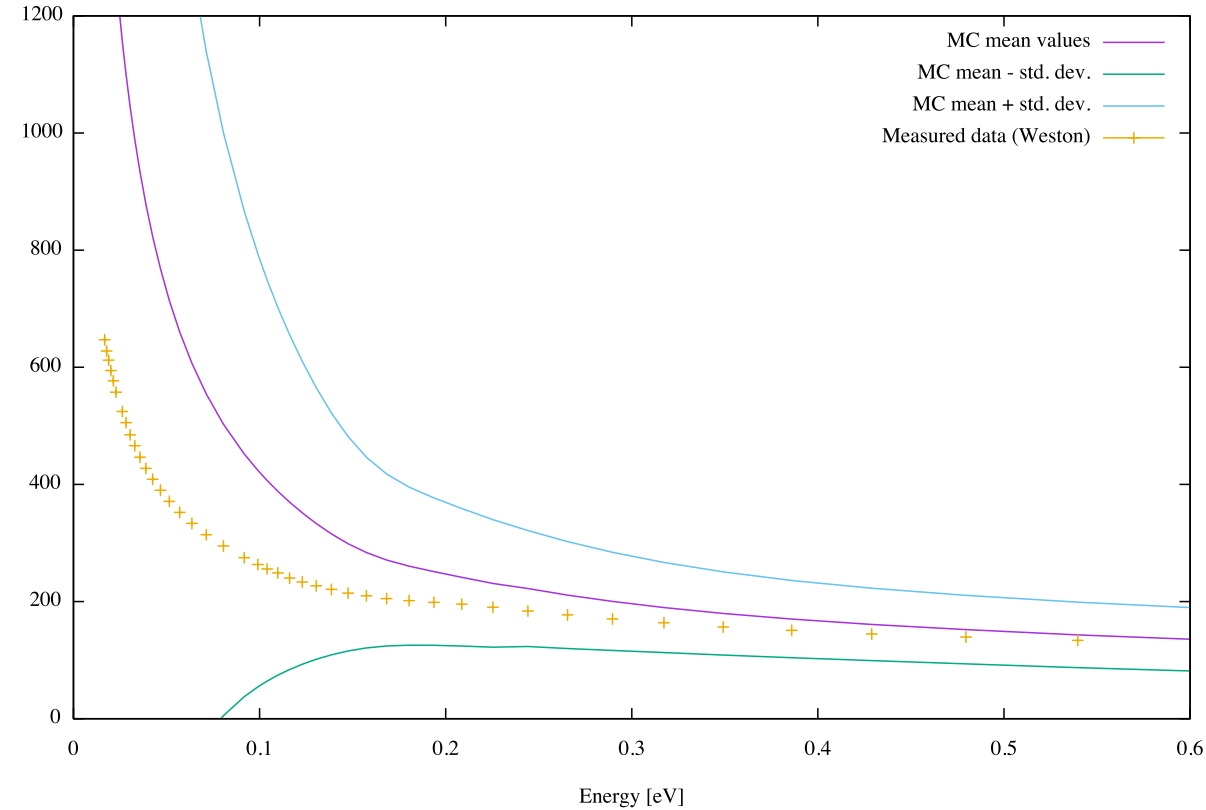
Large variance of MC cross sections below 0.6 eV

- Comparison to Weston capture and fission data:

MC capture mean value +/- std. dev. vs. measured data



MC fission mean value +/- std. dev. vs. measured data



Conclusions and outlook

- Basic components of the MC evaluation framework of differential and integral data
 - Computation of random MC ensemble from ENDF File 32
 - Simulation of IBEs and R-matrix cross section compared to experimental data
 - Computation of weighted averages
- Application to U-233 indicates deviation from the conventional linear approximation
 - IBEs: U233-SOL-{INTER,THERM}-001-001
 - Diff. data: transmission and fission
- Evaluation framework will require MCMC method e.g. Metropolis-Hastings (M.-H.)
 - Computational burden of IBEs makes this more realistic for differential data evaluation
 - Currently surveying parallelized generalizations of the M.-H. method

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- Marco T. Pigni:
 - U-233 SAMMY input files
- Klaus Guber:
 - U-233 transmission and fission data
- B.J. Marshall:
 - Guidance with IBEs and SAMPLER
- Vladimir Sobes:
 - Analytically solvable neutron transport problem for testing

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Auxiliary slides

Overview of generalized form of Bayes' Theorem

- Generalized data = (parameters, data, model defect): $z \equiv (P, D, \delta)$
- Generalized data covariance matrix $\mathbf{C} \equiv \langle (z - \langle z \rangle)(z - \langle z \rangle)^T \rangle$
- Constraint on the posterior expectation values define the evaluation:

$$\omega \equiv T(P) - D - \delta \qquad \langle \omega \rangle' = \omega'_f$$

$$\mathbf{\Omega}' \equiv \langle (\omega - \langle \omega \rangle')(\omega - \langle \omega \rangle')^T \rangle' = \mathbf{\Omega}'_f,$$

- Posterior PDF; model T(P) appears only in the the likelihood function via constraints:

$$p(z|\langle z \rangle, \mathbf{C}, f) \propto p(z|\langle z \rangle, \mathbf{C}) \times p(f|z, \langle z \rangle, \mathbf{C})$$

- Exponential likelihood function

$$p(f|z, \langle z \rangle, \mathbf{C}) = e^{-\sum_i \lambda_i \omega_i - \sum_{ij} \Lambda_{ij} (\omega - \langle \omega \rangle')_i (\omega - \langle \omega \rangle')_j}$$

- Extant evaluations impose constraints $\omega'_f = 0 \quad \mathbf{\Omega}'_f = 0$ leading to:

$$p(f_0|z, \langle z \rangle, \mathbf{C}) = \delta_{\text{Dirac}}(\omega)$$