

Probing the deuteron at extremely large internal momenta

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Hall A and Hall C Collaborations

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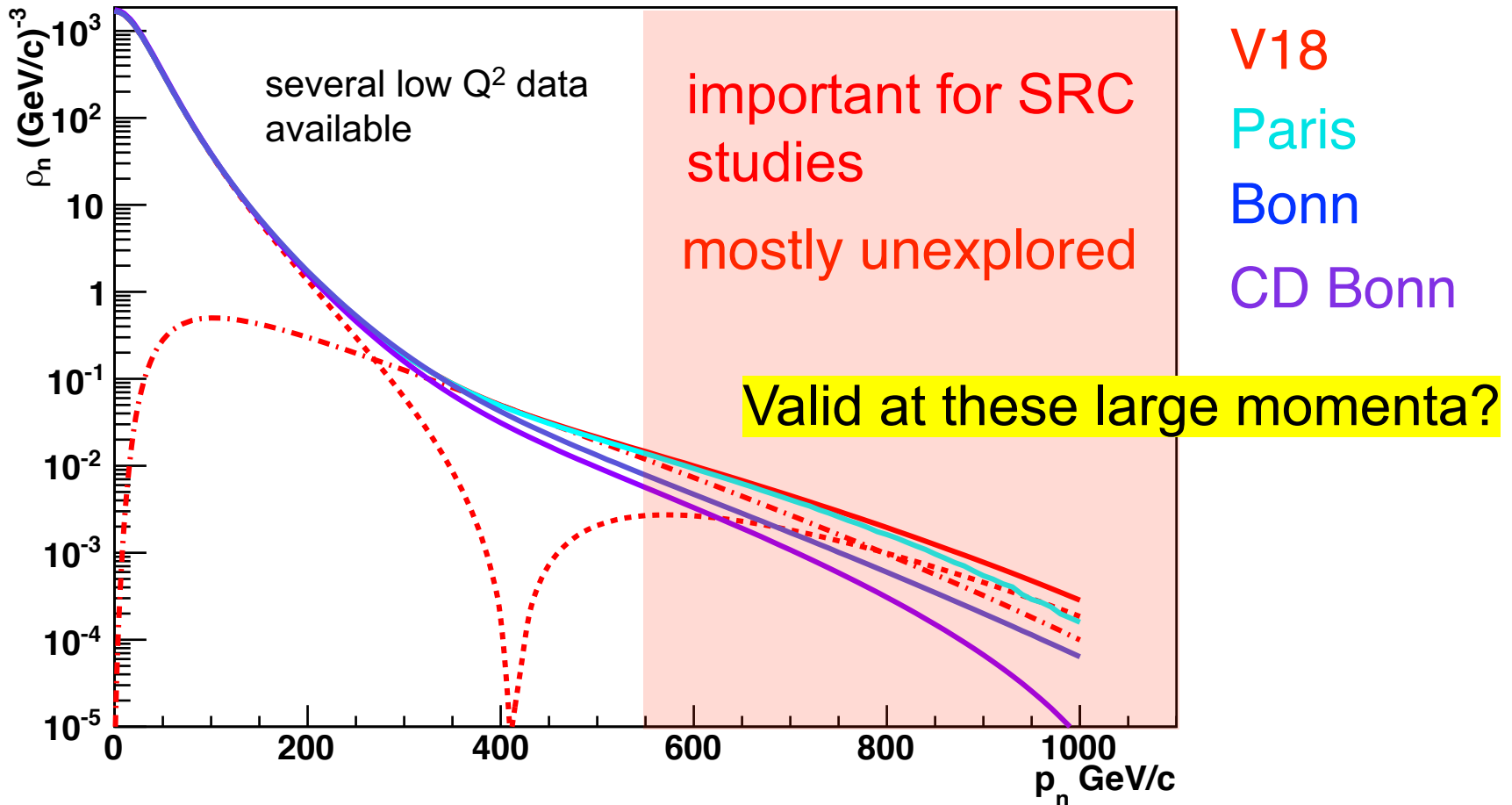
- Introduction
- History and previous Results
- Angular distributions for fixed p_m
- Missing momentum dependences for fixed neutron angle
- New Hall C Results
- Light Cone Momentum Distributions
- Inelastic Channels
- Summary

Introduction: Role of the Deuteron

- Key system to investigate the (repulsive) core of the NN interaction.
- Basis for SRC (structure) studies
- Prime nucleus to test NN models
- Structure needs to be understood in detail at all length scales

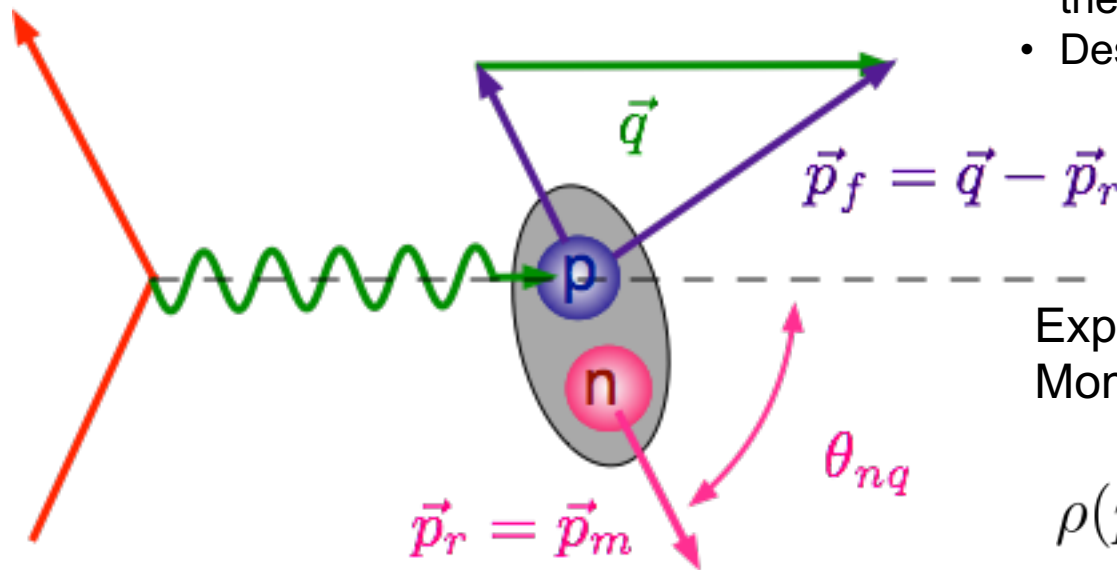
Introduction: Momentum Distribution

virtually no experimental $d(e, ep)n$ data exist for $p_m > 0.5$ GeV/c without large contributions of FSI, MEC and IC



Introduction: D(e,e'p) in PWIA

$$\frac{d^5\sigma}{d\omega d\Omega_e d\Omega_p} = k\sigma_{ep}\rho(p_r)$$



Plane Wave IA:

- Hit nucleon does not interact with the recoiling system
- Described by a plane wave

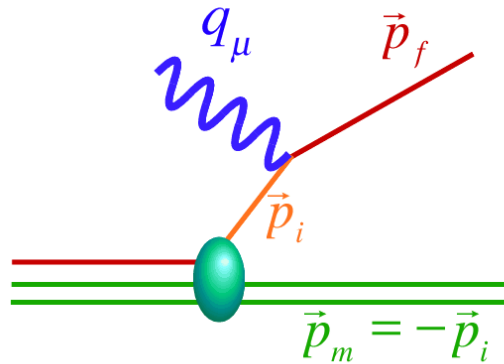
Experimental Momentum distributions:

$$\rho(p_r)_{exp} = \sigma_{red} = \frac{\sigma_{exp}}{k\sigma_{ep}}$$

also called reduced cross sections

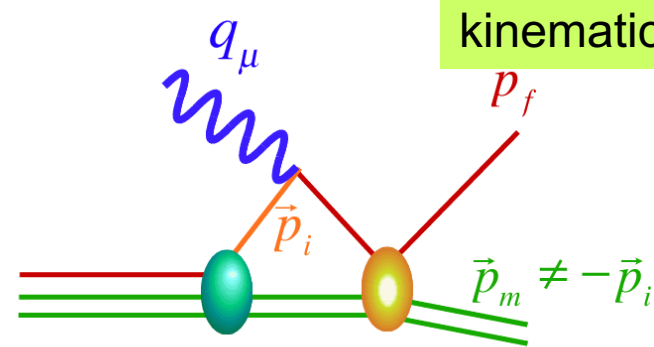
Introduction: D(e,e' p) Reaction Mechanisms

PWIA



$$\frac{d\sigma}{d\omega d\Omega_e d\Omega_N} = k\sigma_{eN} S(E_m, p_m)$$

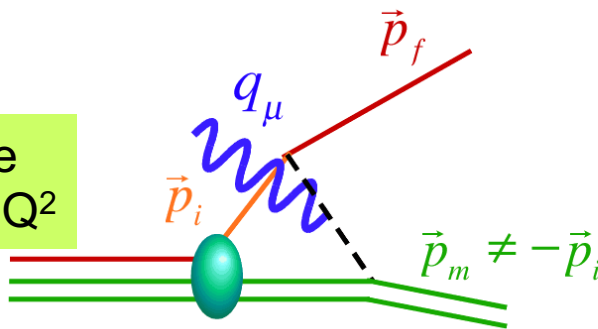
FSI



reduced at certain kinematics ?

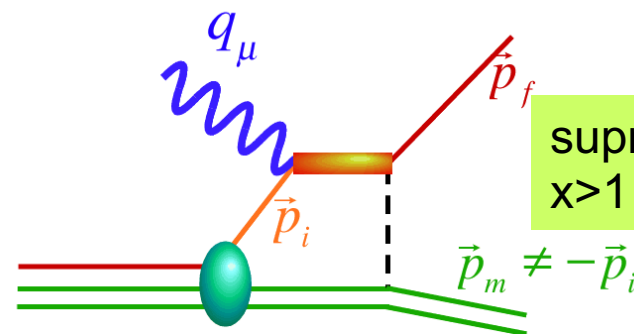
$$\frac{d\sigma}{d\omega d\Omega_e d\Omega_N} = k\sigma_{eN} D(E_m, p_f, p_m)$$

MEC



expected to be small at large Q²

IC

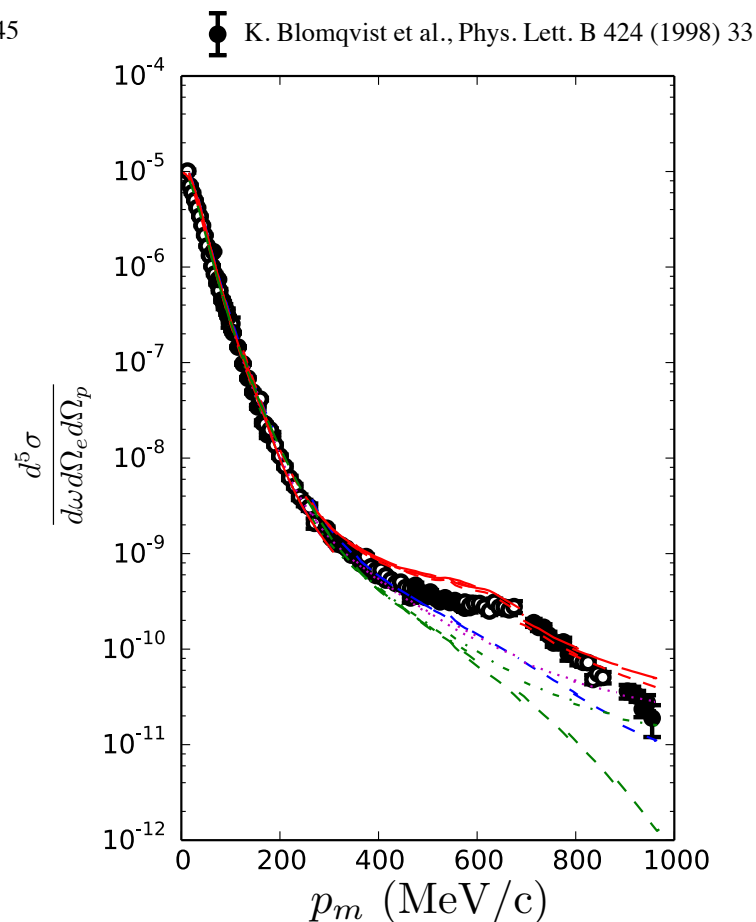
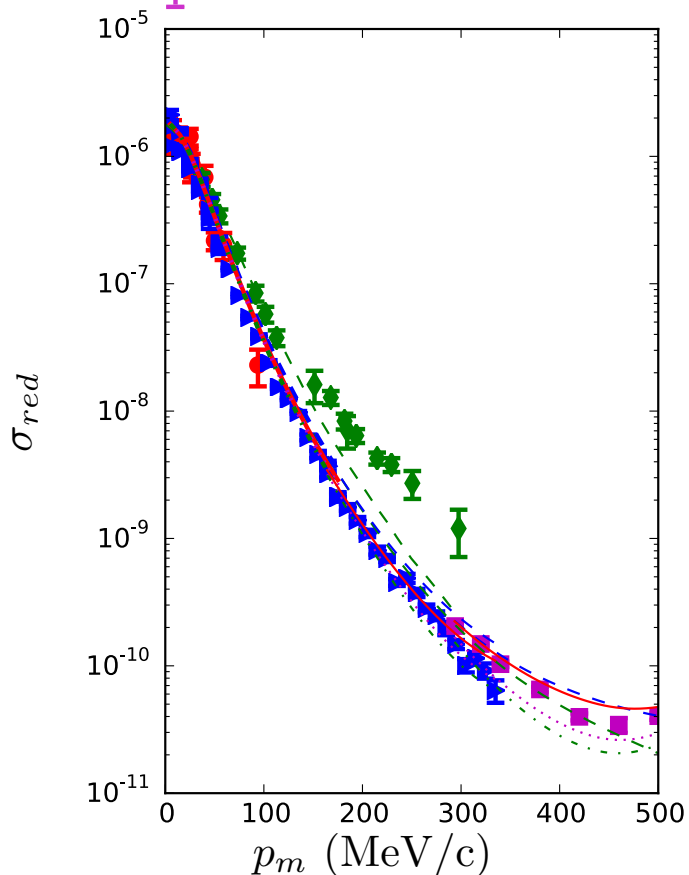


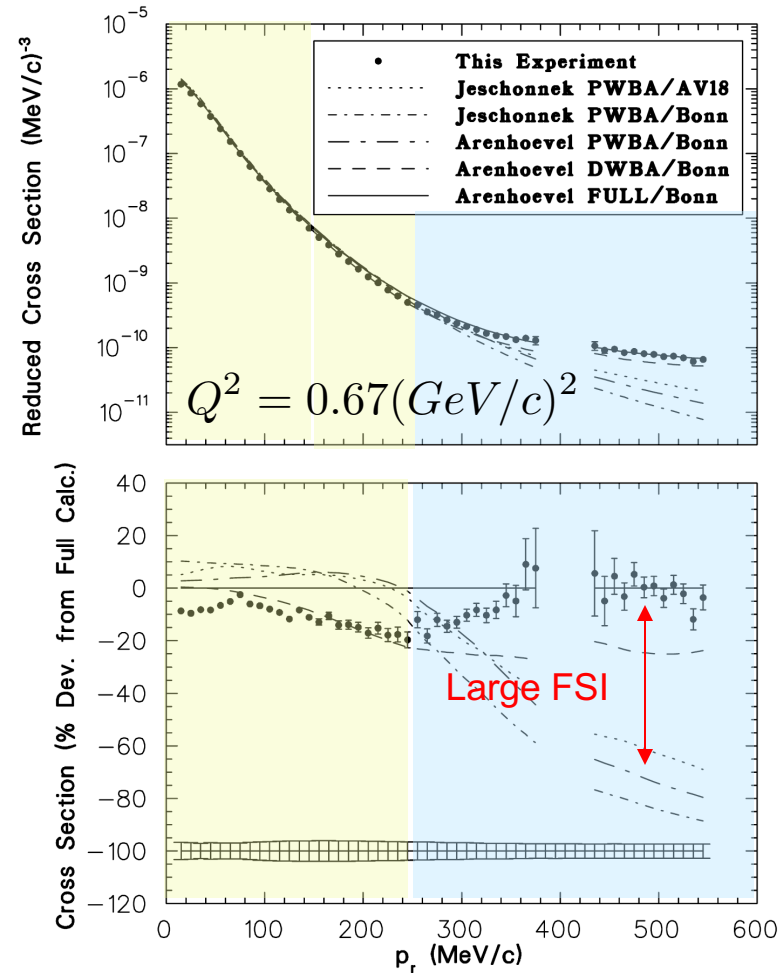
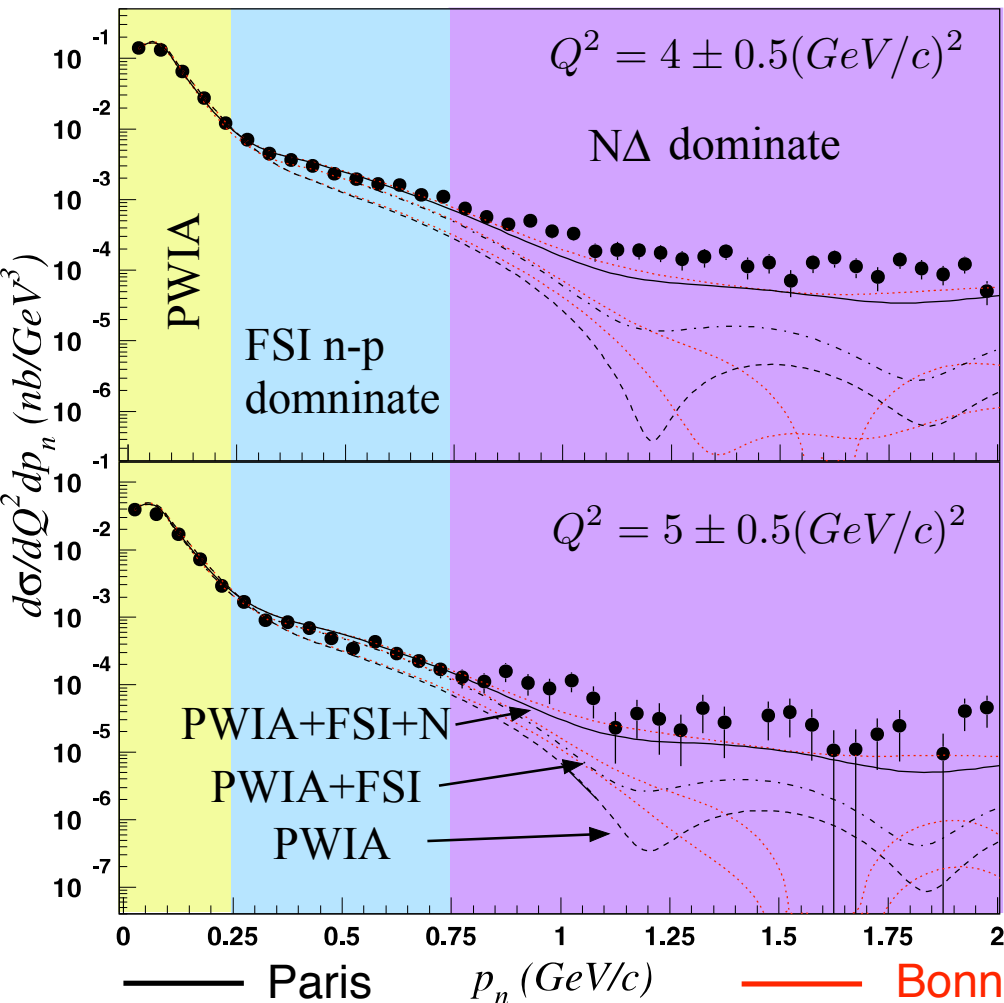
suppressed for x > 1

History of Missing Momentum Dependences

- P. Bounin and M. Croissiaux, Nucl. Phys 70 (1965) 401
- Y. Antufev, et al. Pis'ma Zh. Eksp. Teor. Fiz. 19 (1974) 657
- A. Bussiere et al., Nucl. Phys. A 365 (1981) 349
- S. Turck-Chieze et al., Phys. Lett. B 142 (1984) 145

$Q^2 < 0.33 \text{ (GeV/c)}^2$

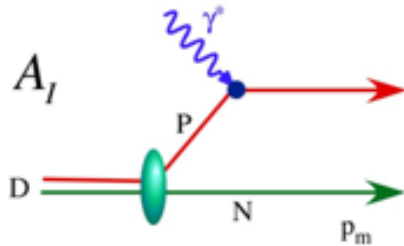




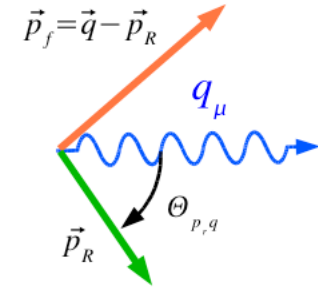
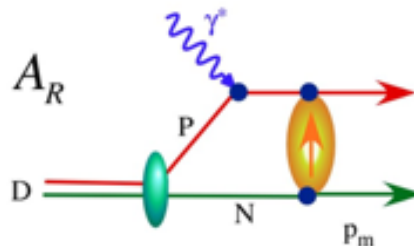
Cross sections integrated over full CLAS acceptance, all recoil angles and wide range of kinematic settings

At high Q^2 FSI as Nucleon Re-Scattering

IA Amplitude (real):



Rescattering Amplitude (at high energy mostly imaginary):



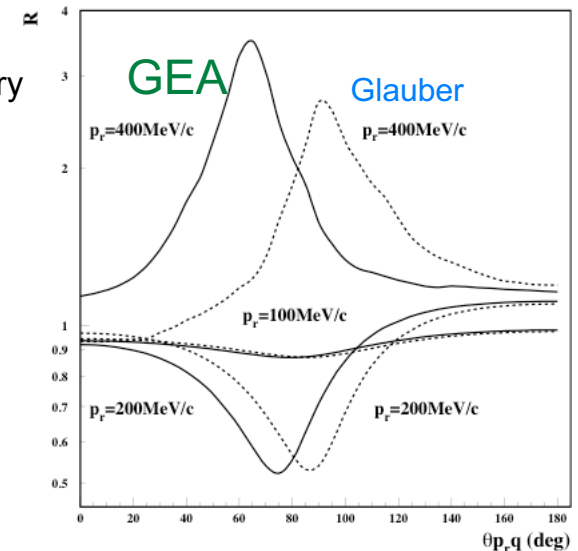
Total Scattering Amplitude: $A = A_I + iA_R$ $A_R \approx i |A_R|$ mostly imaginary

Cross Sections: $\sigma \sim |A|^2 = |A_I + iA_R|^2$
 $\sigma \sim |A_I|^2 - 2|A_I||A_R| + |A_R|^2$

Ratio to PWIA $R = \frac{\sigma}{\sigma_I} = 1 - 2 \frac{|A_I||A_R|}{|A_I|^2} + \frac{|A_R|^2}{|A_I|^2}$

Glauber: spectator nucleons frozen

GEA: spectator nucleon moving

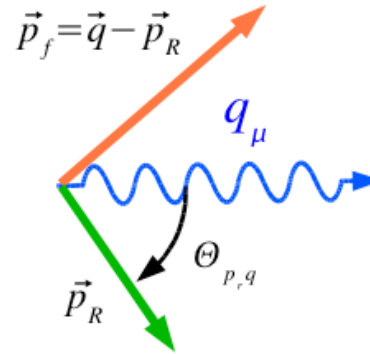


M.Sargsian, PRC **82**, 014612 (2010)

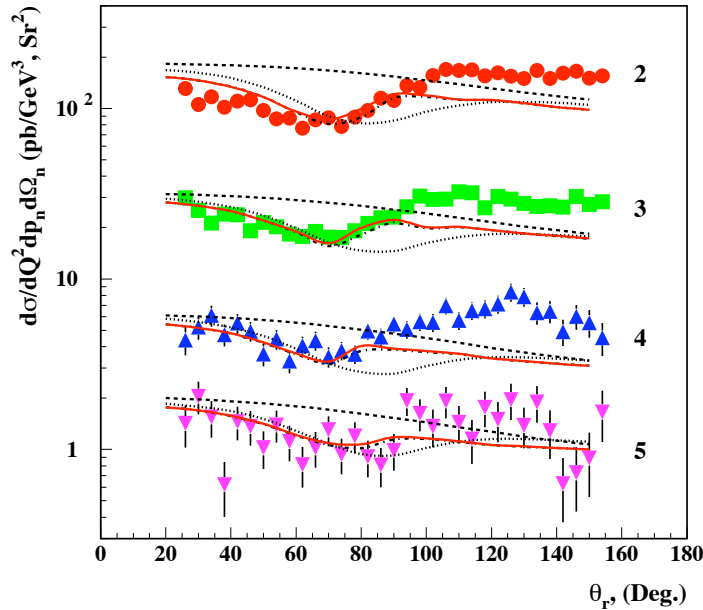
CLAS

Data: Egyian et al. (CLAS) PRL 98 (2007)

Calculation M. Sargsian

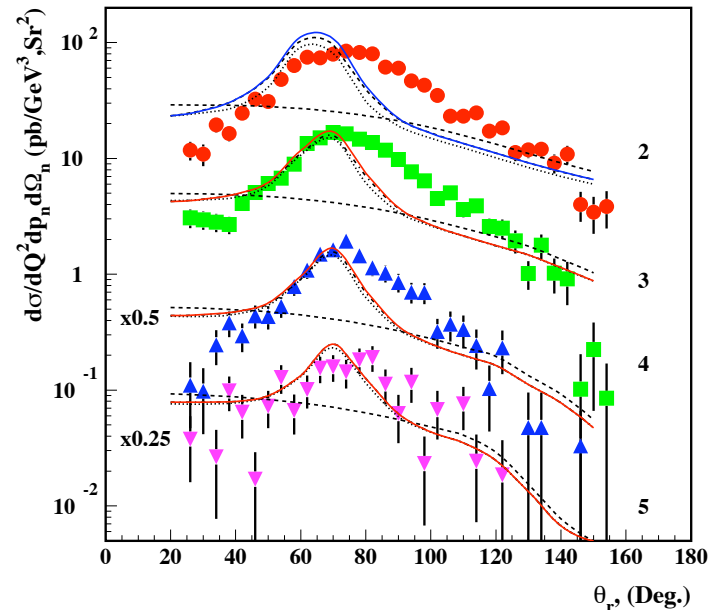


$p_m = 0.25 \pm .05 \text{ GeV}/c$



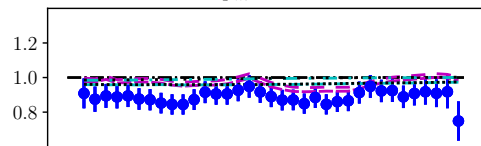
M.Sargsian, PRC **82**, 014612 (2010)

$p_m = 0.5 \pm 0.1 \text{ GeV}/c$

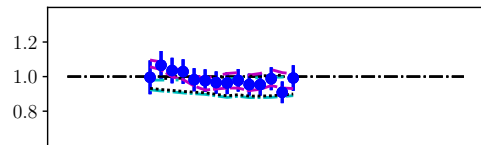


$$R = \frac{\sigma_{EXP}}{\sigma_{PWIA}}$$

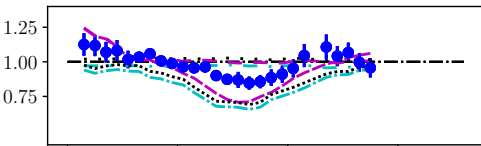
$Q^2 = 0.8$
 $p_m = 0.0$



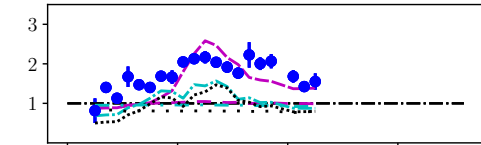
$p_m = 0.1$



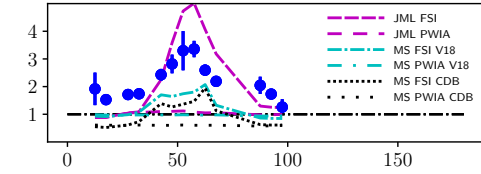
$p_m = 0.2$



$p_m = 0.4$



$p_m = 0.5$



FSI

..... CD Bonn

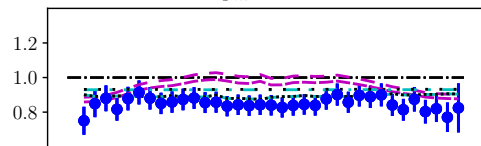
PWIA

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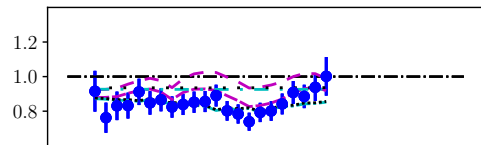
Hall A

WB et al. PRL 107 (2011) 262501

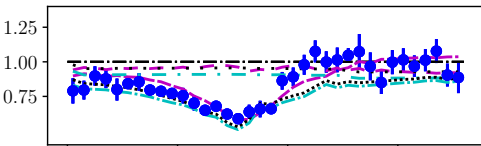
$Q^2 = 2.1$
 $p_m = 0.0$



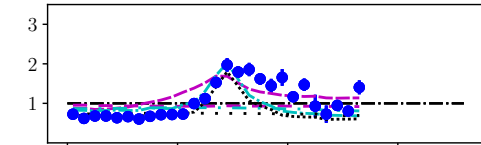
$p_m = 0.1$



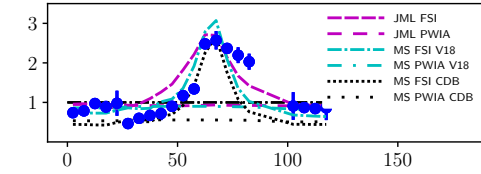
$p_m = 0.2$



$p_m = 0.4$



$p_m = 0.5$



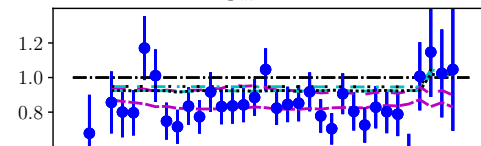
FSI

----- V18

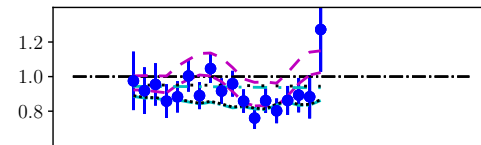
PWIA

Exploring QCD with light nuclei at
EIC

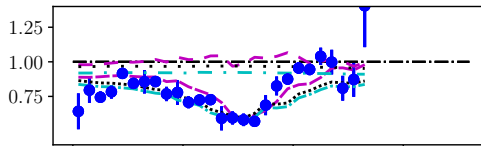
$Q^2 = 3.5$
 $p_m = 0.0$



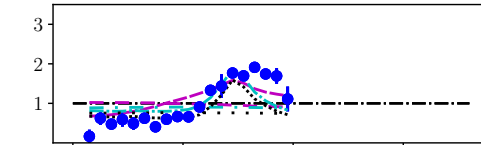
$p_m = 0.1$



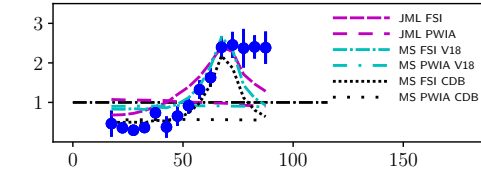
$p_m = 0.2$



$p_m = 0.4$



$p_m = 0.5$



FSI

----- Laget

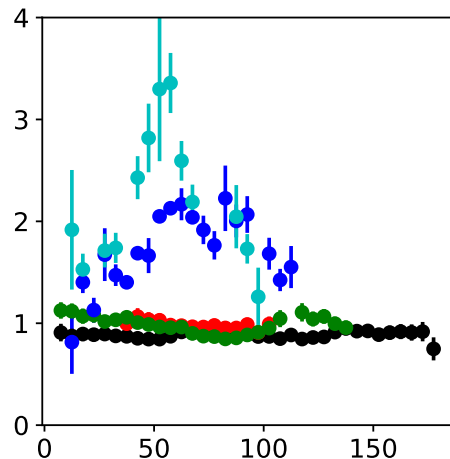
PWIA

$\Delta p_m = \pm 0.02$ GeV/c

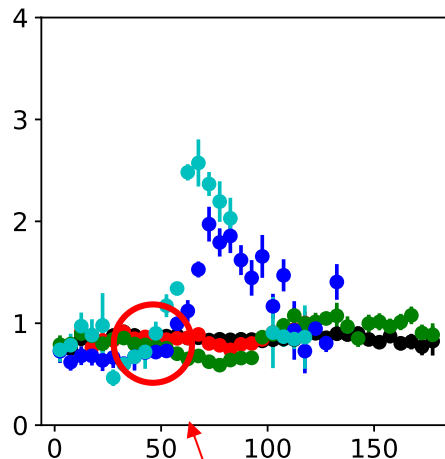
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$$R = \frac{\sigma_{EXP}}{\sigma_{PWIA}}$$

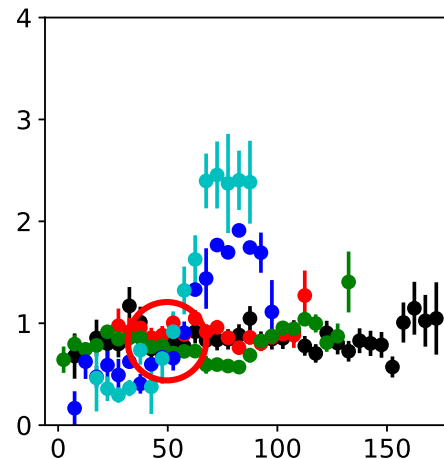
$Q^2 = 0.8$



$Q^2 = 2.1$



$Q^2 = 3.5$



0.1 GeV/c
0.2 GeV/c
0.4 GeV/c
0.5 GeV/c

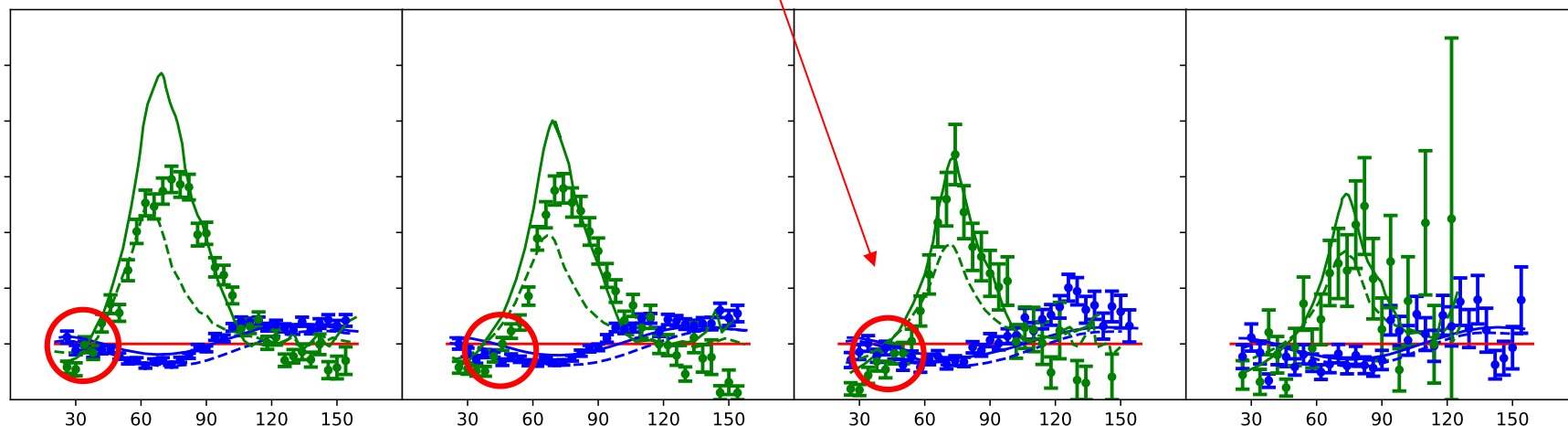
$Q^2 = 2$

$Q^2 = 3$

$Q^2 = 4$

$Q^2 = 5$

$\sim p_m$ independent, small FSI

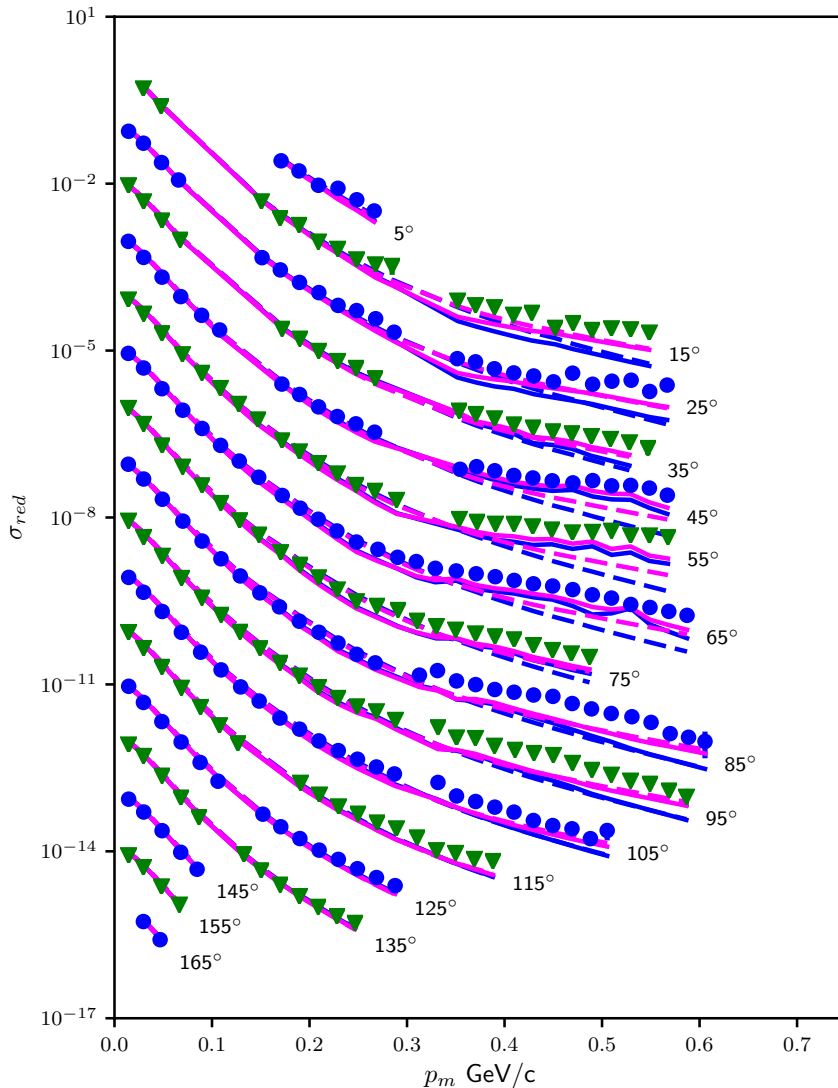


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Exploring QCD with light nuclei at
EIC

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Momentum Distributions



$$Q^2 = 0.8 \text{ (GeV/c)}^2$$

- Large FSI for $p_m > 0.25 \text{ GeV/c}$
- Little sensitivity to W.F.
- Eikonal regime not yet reached

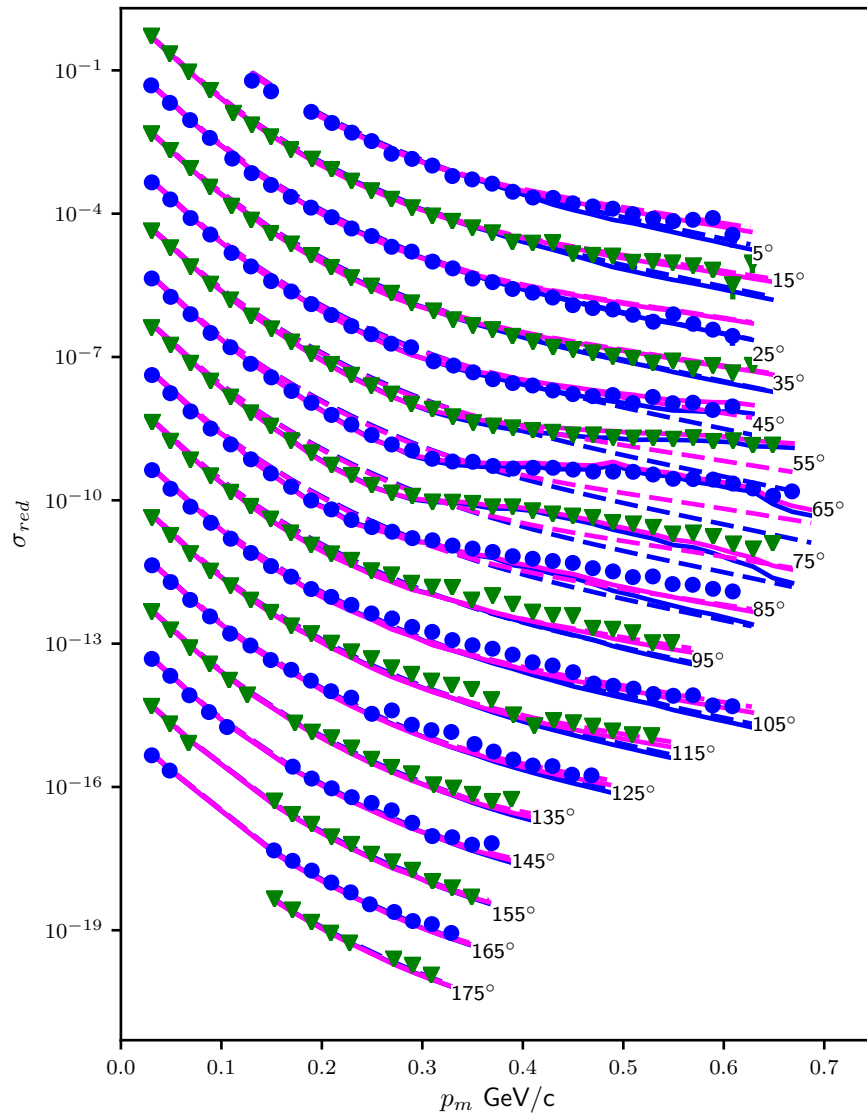
Experimental data are scaled for display purposes

1/20/20

Exploring QCD with light nuclei at
EIC

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Experimental data are scaled for display purposes

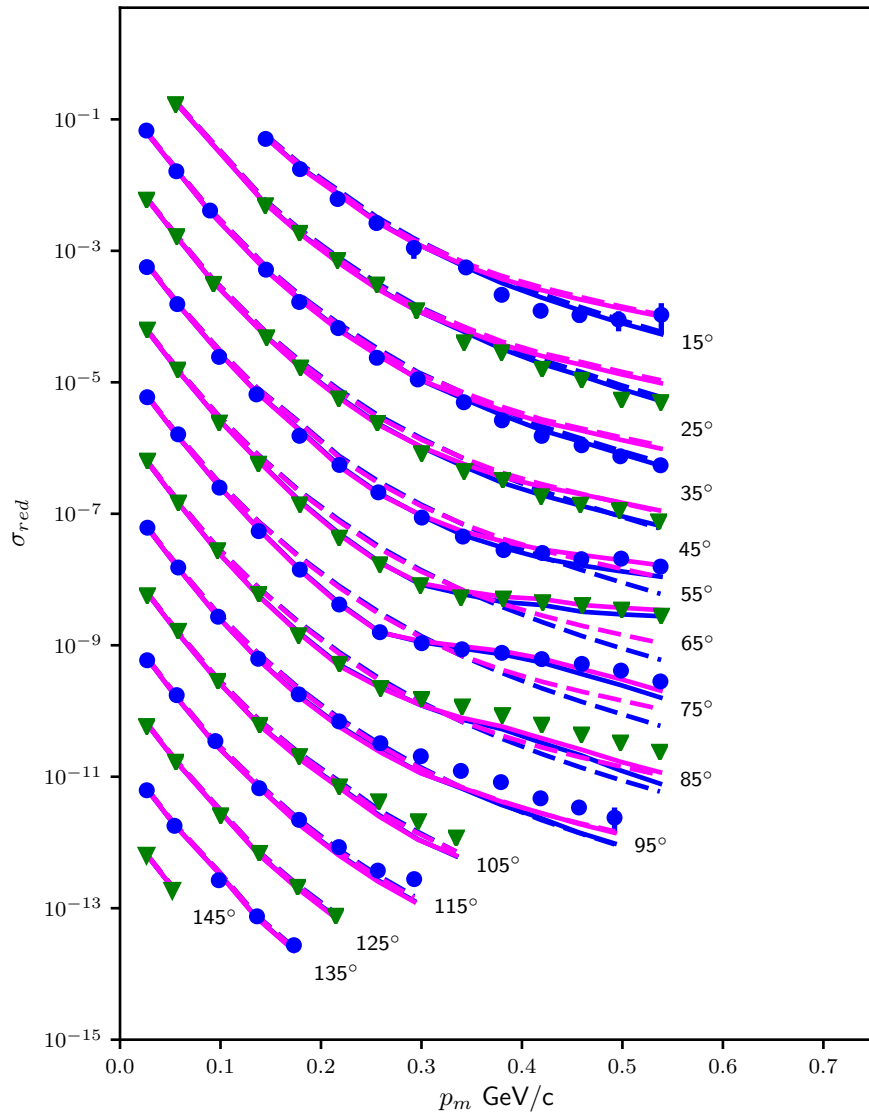


$$Q^2 = 2.1 (\text{GeV}/c)^2$$

Small FSI

Large FSI

Small FSI, Resonances ?



$$Q^2 = 3.5 \text{ (GeV/c)}^2$$

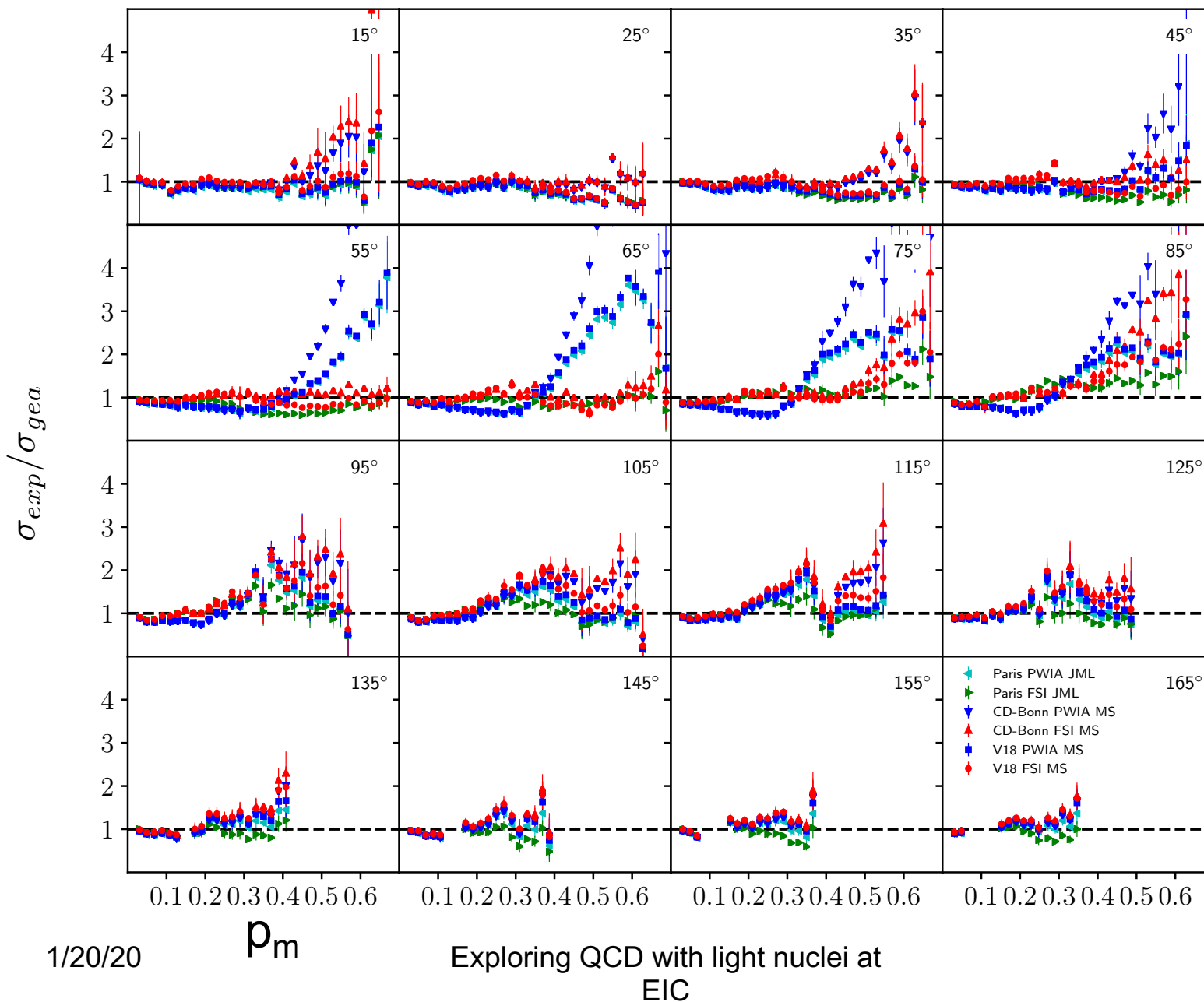
Small FSI

Large FSI

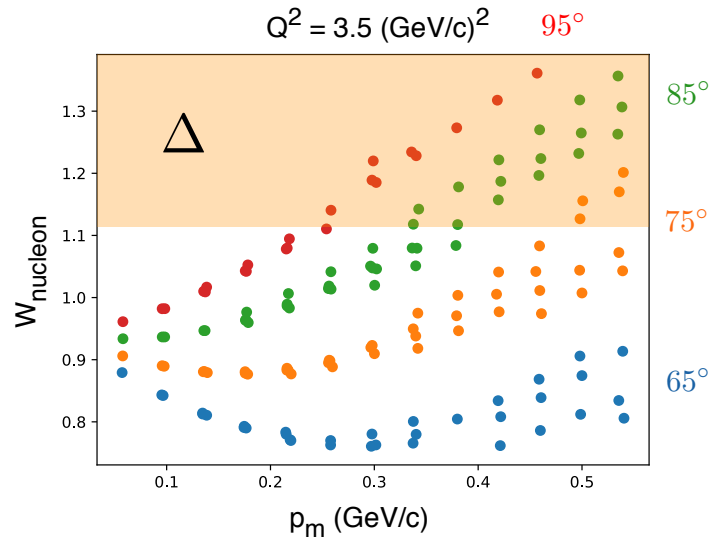
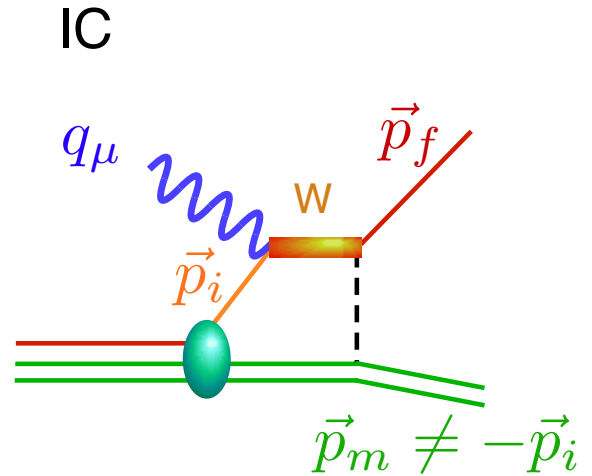
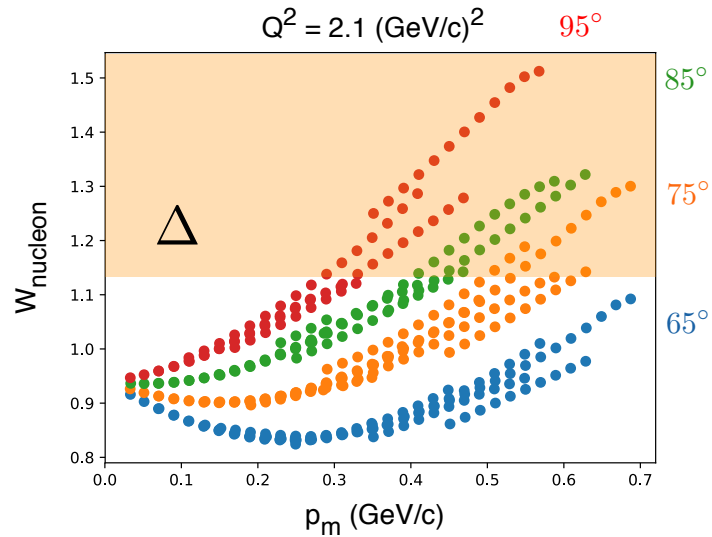
Small FSI, Resonances ?

Experimental data are scaled for display purposes

$Q^2 = 2.1 \text{ (GeV/c)}^2$ GEA M.Sargsian



Resonance Contributions to Re-Scattering ?



Pushing the p_m limit: Experiment at 12 GeV in Hall C

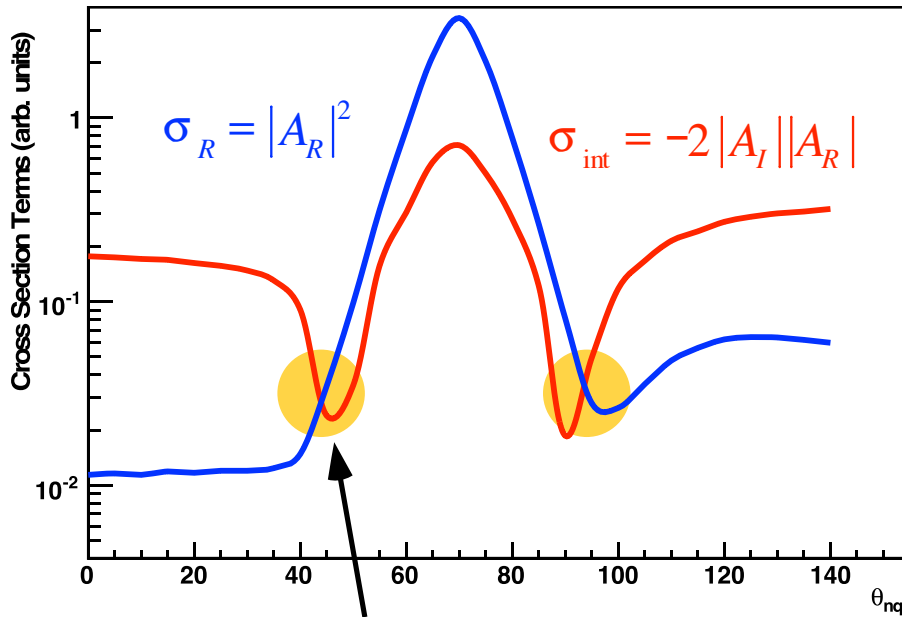
- Determine cross sections at missing momenta up to 1 GeV/c
- Measure at well defined kinematic settings
- Selected kinematics to minimize contributions from FSI
- Selected kinematics to minimize effects of delta excitation

Original Parameters

- Beam: 11 GeV, $\sim 50\mu\text{A}$
- Electron Detector: SHMS at $p_{\text{cen}} = 9.32 \text{ GeV}/c$
 - $\theta_e = 11.68^\circ$, $Q^2 = 4.25 (\text{GeV}/c)^2$, $x = 1.35$
- Proton Detector: HMS $1.96 \leq p_{\text{cen}} \leq 2.3 \text{ GeV}/c$
 - $p_m = 0.5, 0.6, 0.7, 0.8, 0.9, 1.0 \text{ GeV}/c$
 - Angles: $63.5^\circ \geq \theta_p \geq 53.1$
- Target: 10 cm LHD

FSI Reduction

Reduction of FSI: $\sigma : |A_I|^2 - 2|A_I||A_R| + |A_R|^2$



Rescattering determined by slope factor:

$$f_s = e^{-\frac{b}{2}k_t^2}$$

$$k_t = p_m \sin(\theta_{p_m q})$$

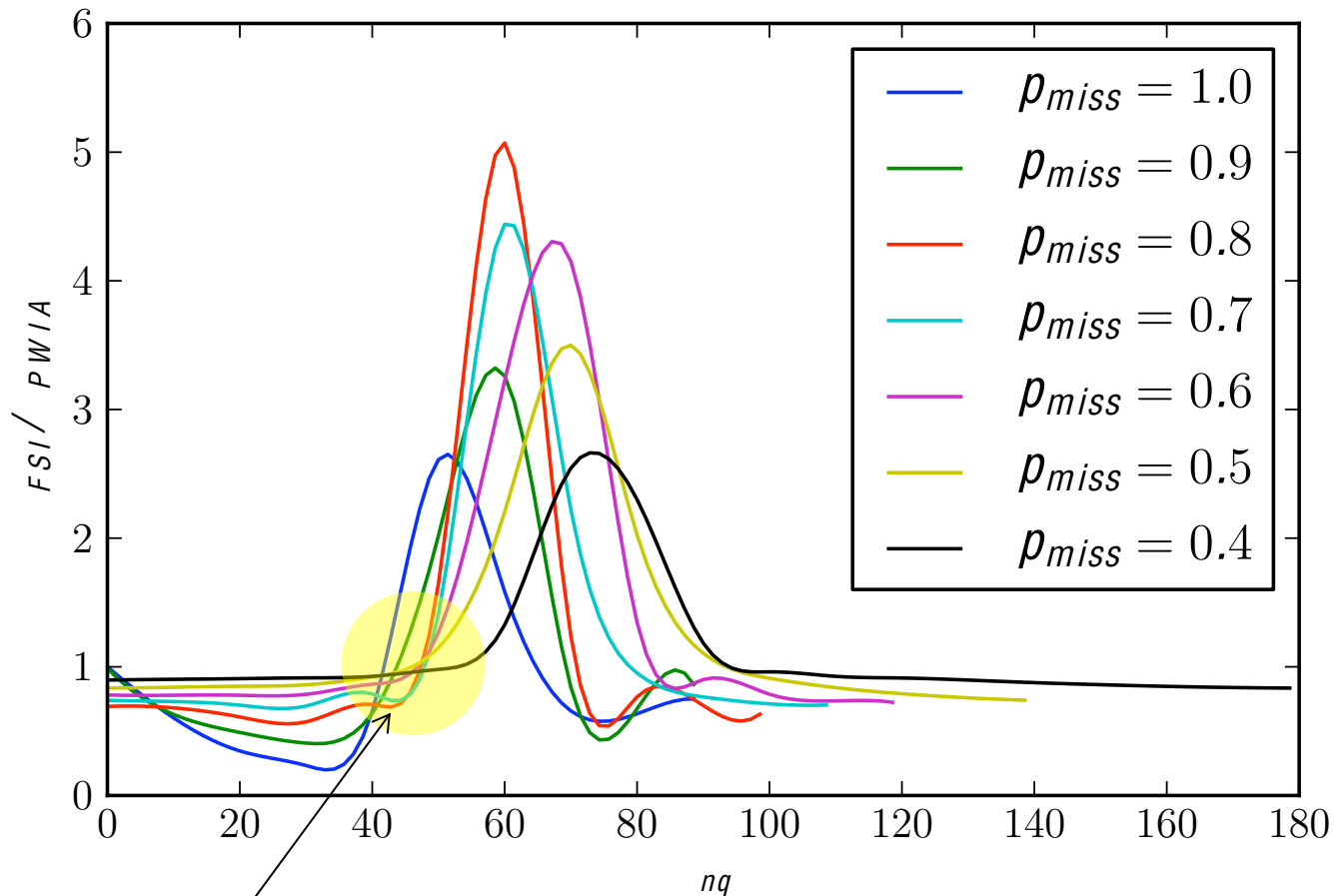
$$b : 6(\text{GeV} / c)^{-2}$$

f_s relatively flat up to $k_t \approx 0.5(\text{GeV} / c)$
 $\Rightarrow p_m \approx 0.8(\text{GeV} / c)$

both terms are equal \Rightarrow **interference** and **rescattering** cancel

- b determined by nucleon size
- cancellation due to imaginary rescattering amplitude
- valid only for high energy (GEA)

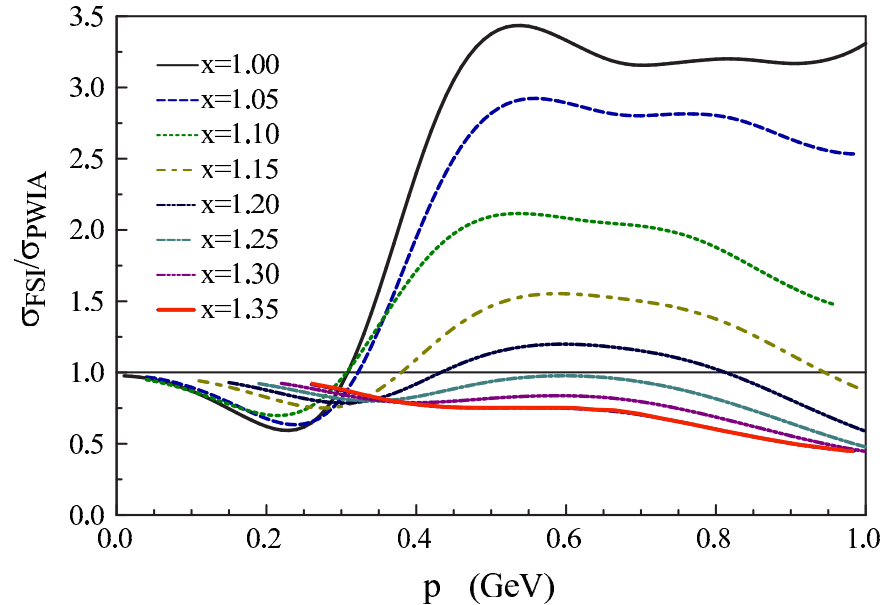
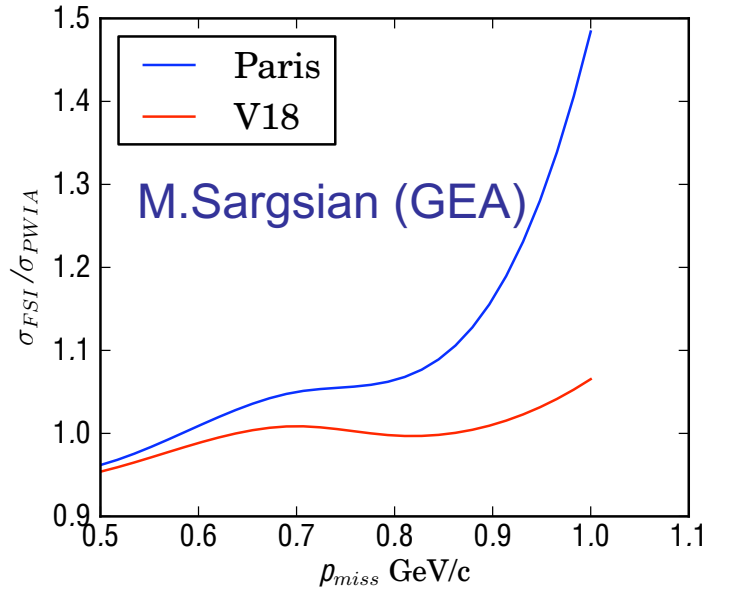
Estimated Angular Distributions up to $p_m = 1\text{ GeV}/c$



FSI depend weakly on p_m

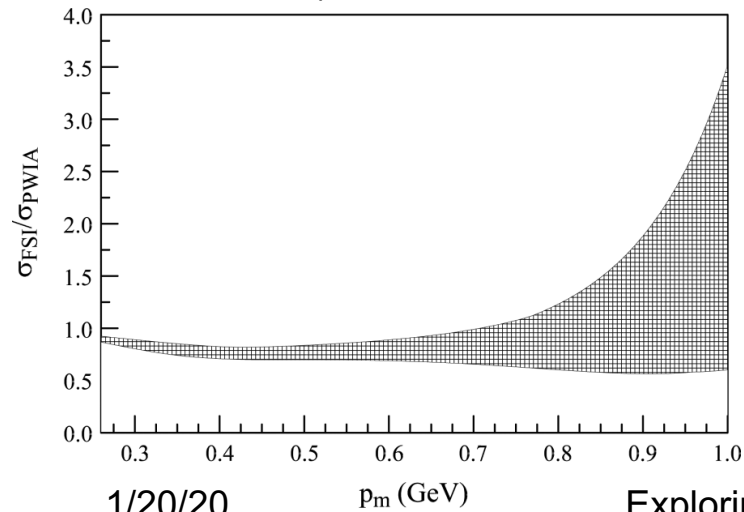
Calculation: M.Sargsian

FSI uncertainties



W.P.Ford et al. PRC 90 064006 (2014)

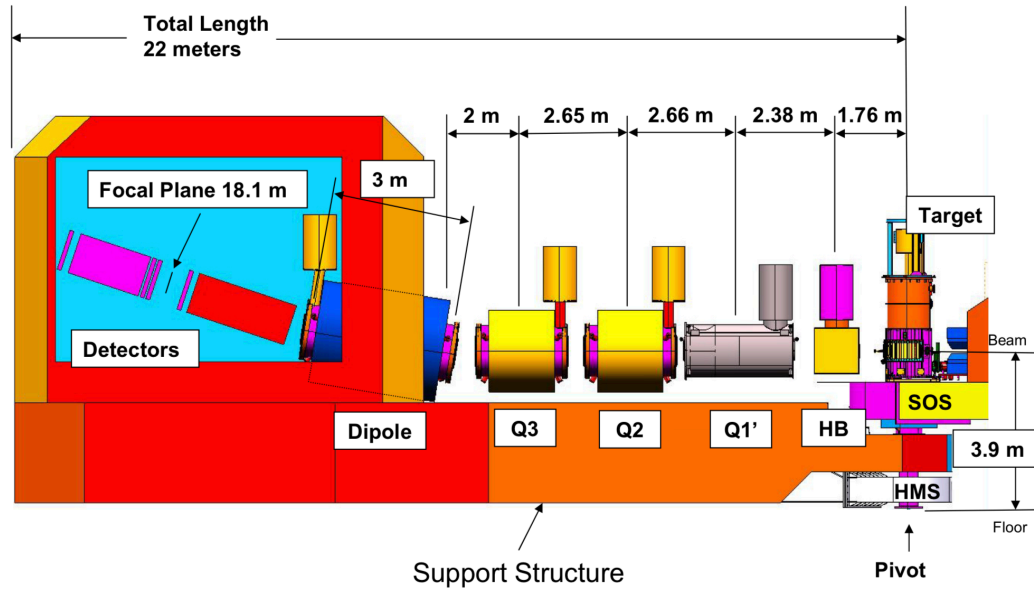
range of possible values of R for all form factor and W.F. used



First Results of Commissioning Experiment



SHMS Overall Dimensions

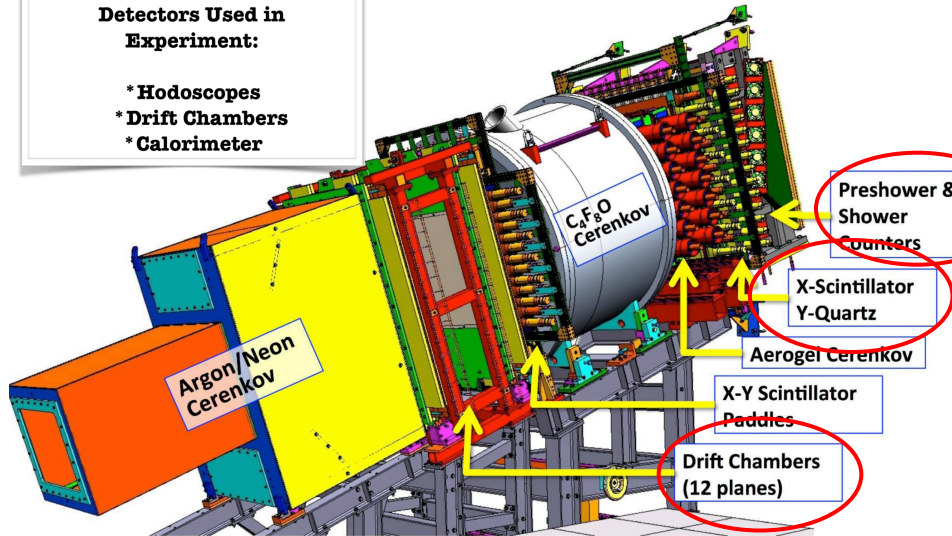


Electrons

Particle Detectors inside the SHMS

Detectors Used in Experiment:

- *Hodoscopes
- *Drift Chambers
- *Calorimeter

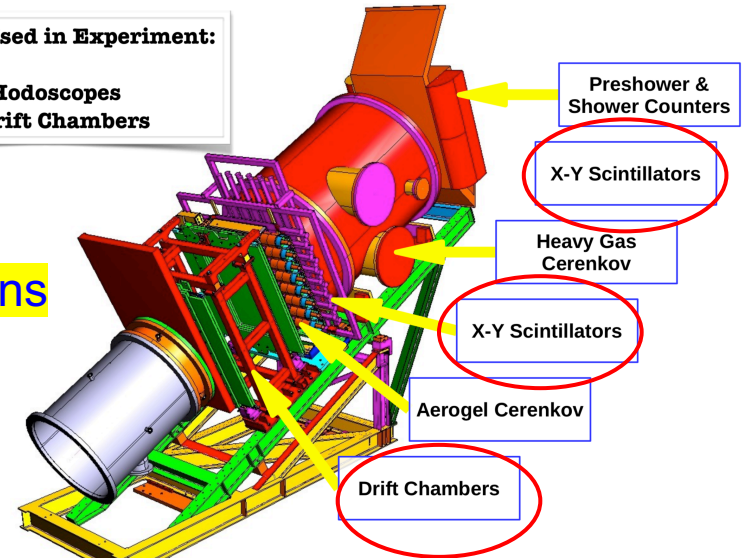


Particle Detectors inside the HMS

Detector Used in Experiment:

- *Hodoscopes
- *Drift Chambers

Protons

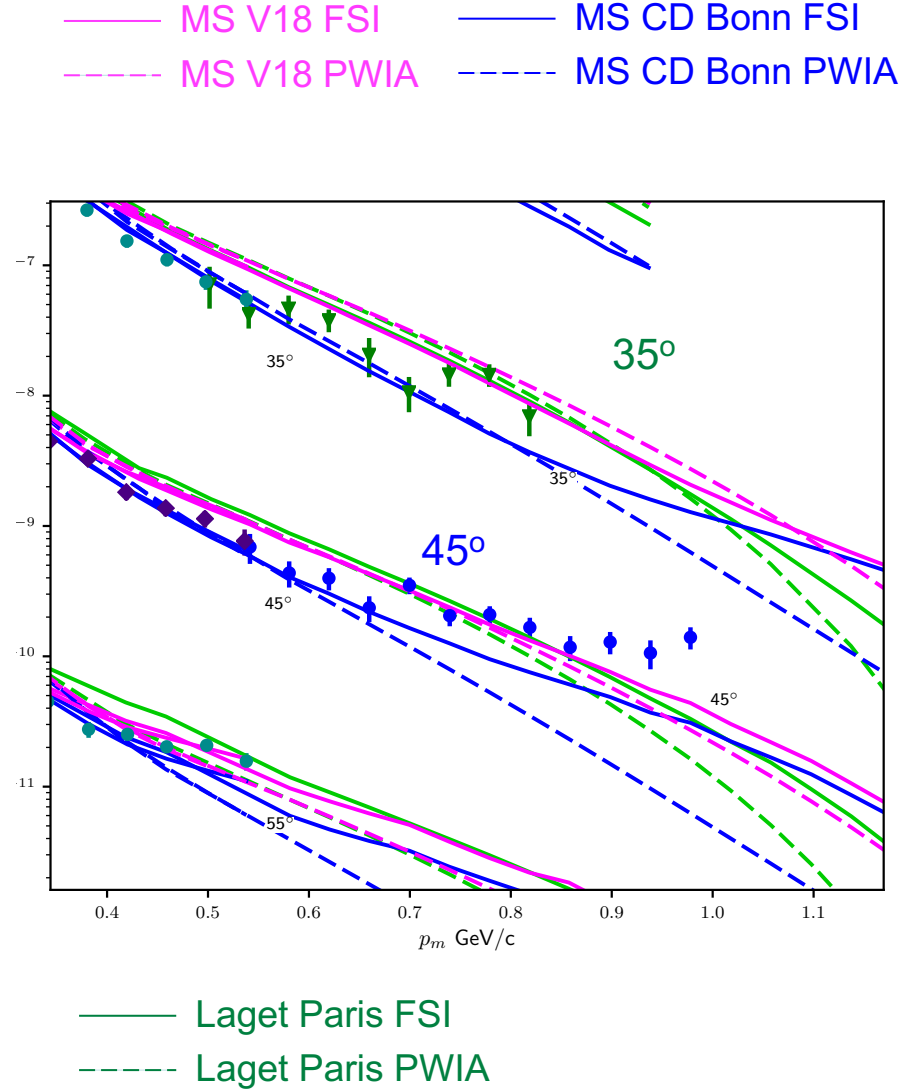
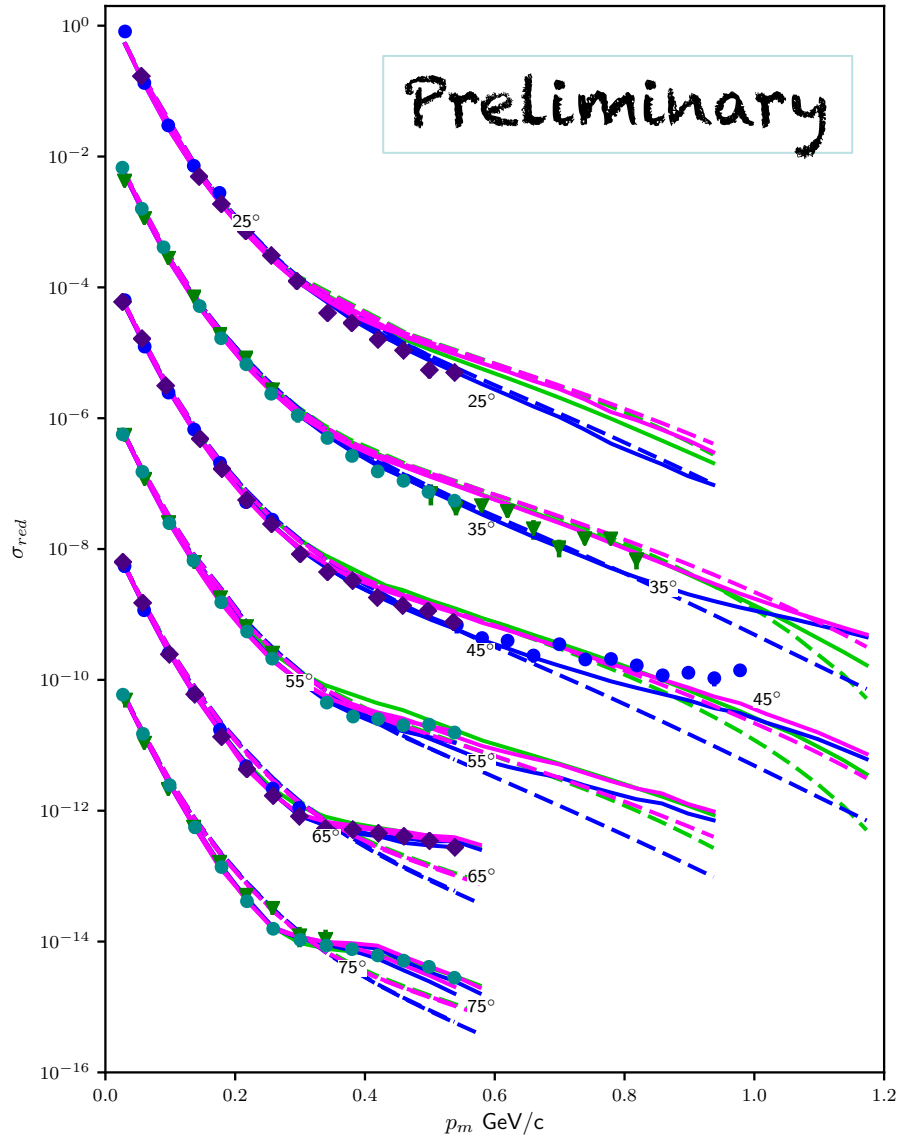


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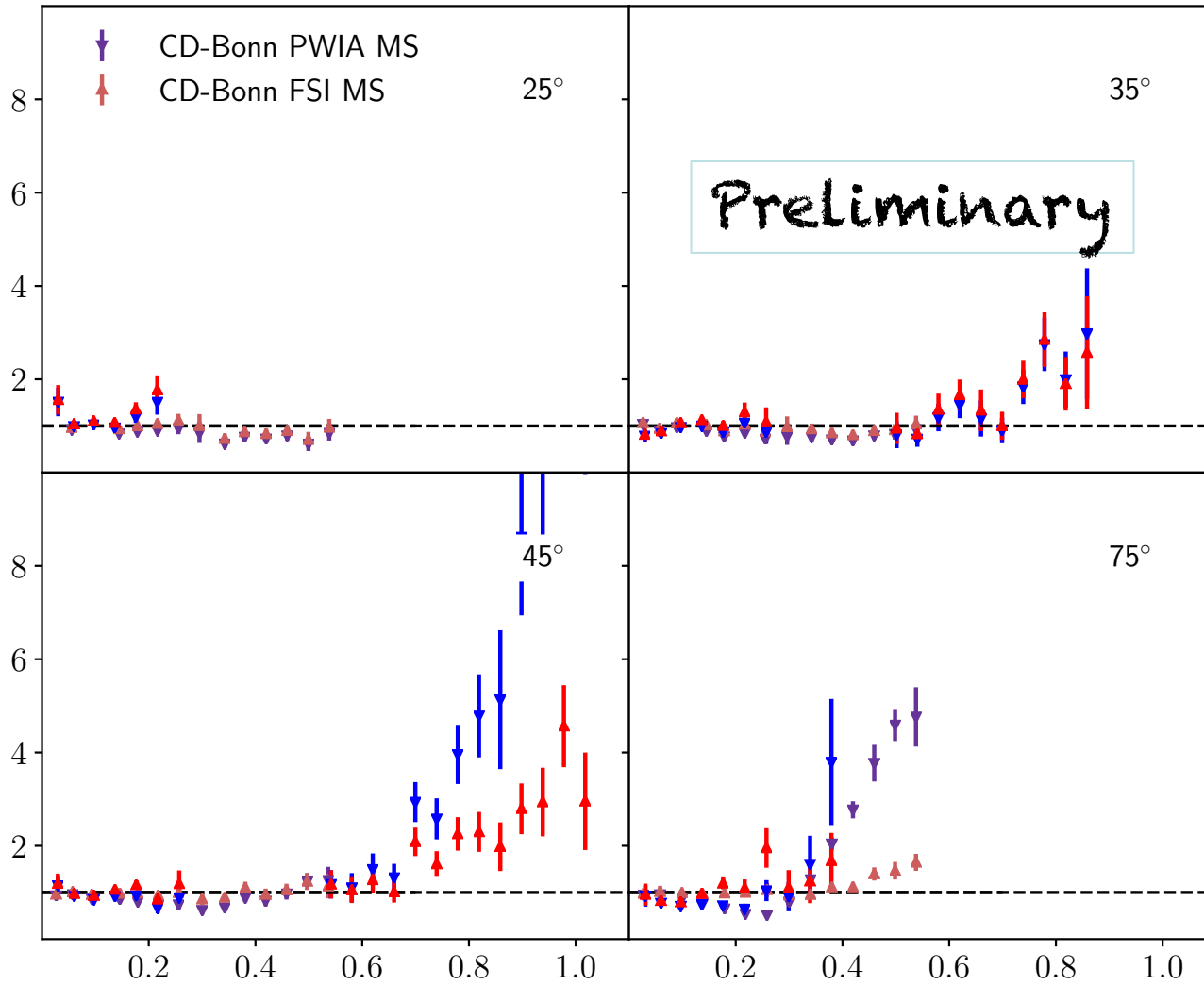
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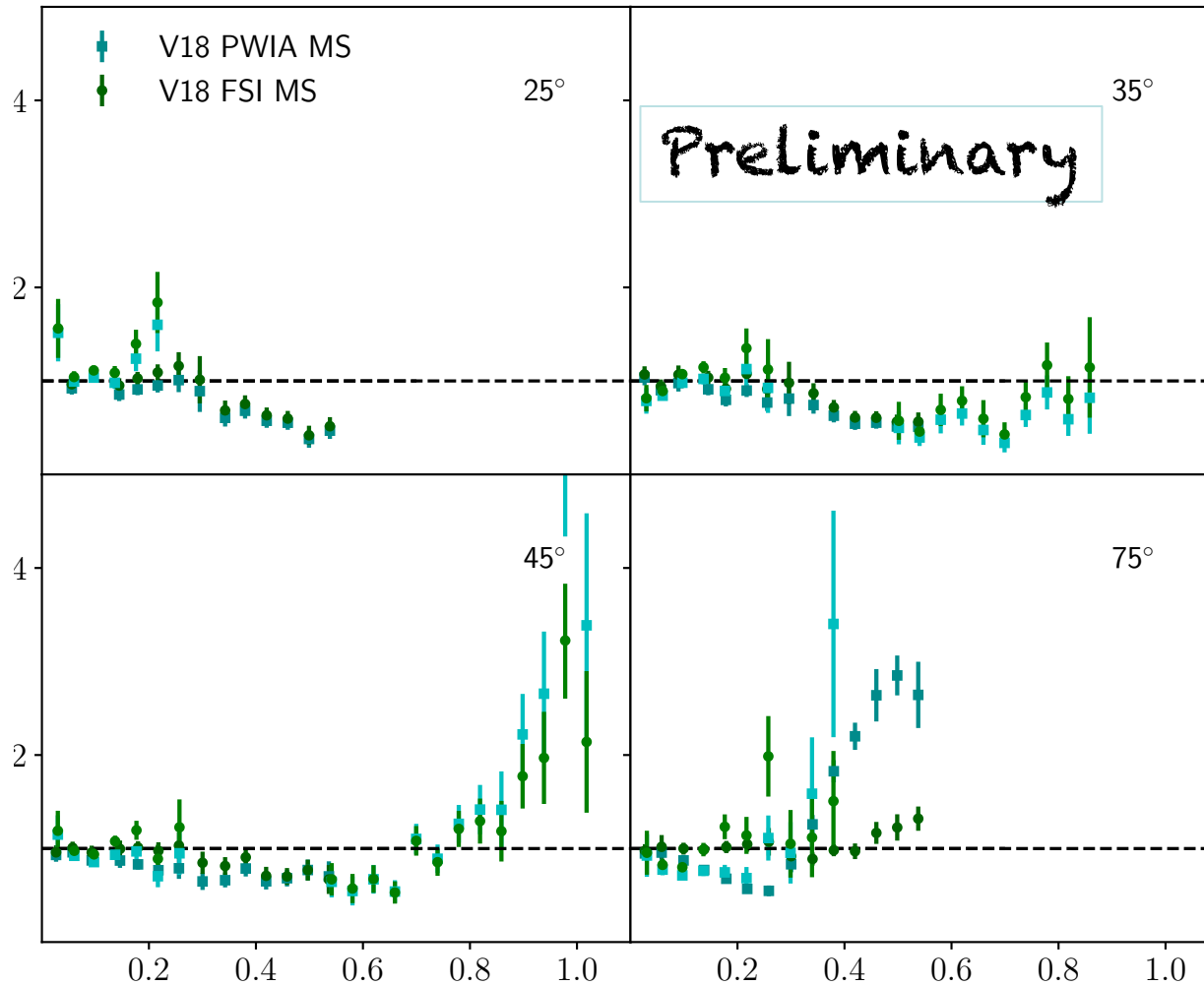
New Results



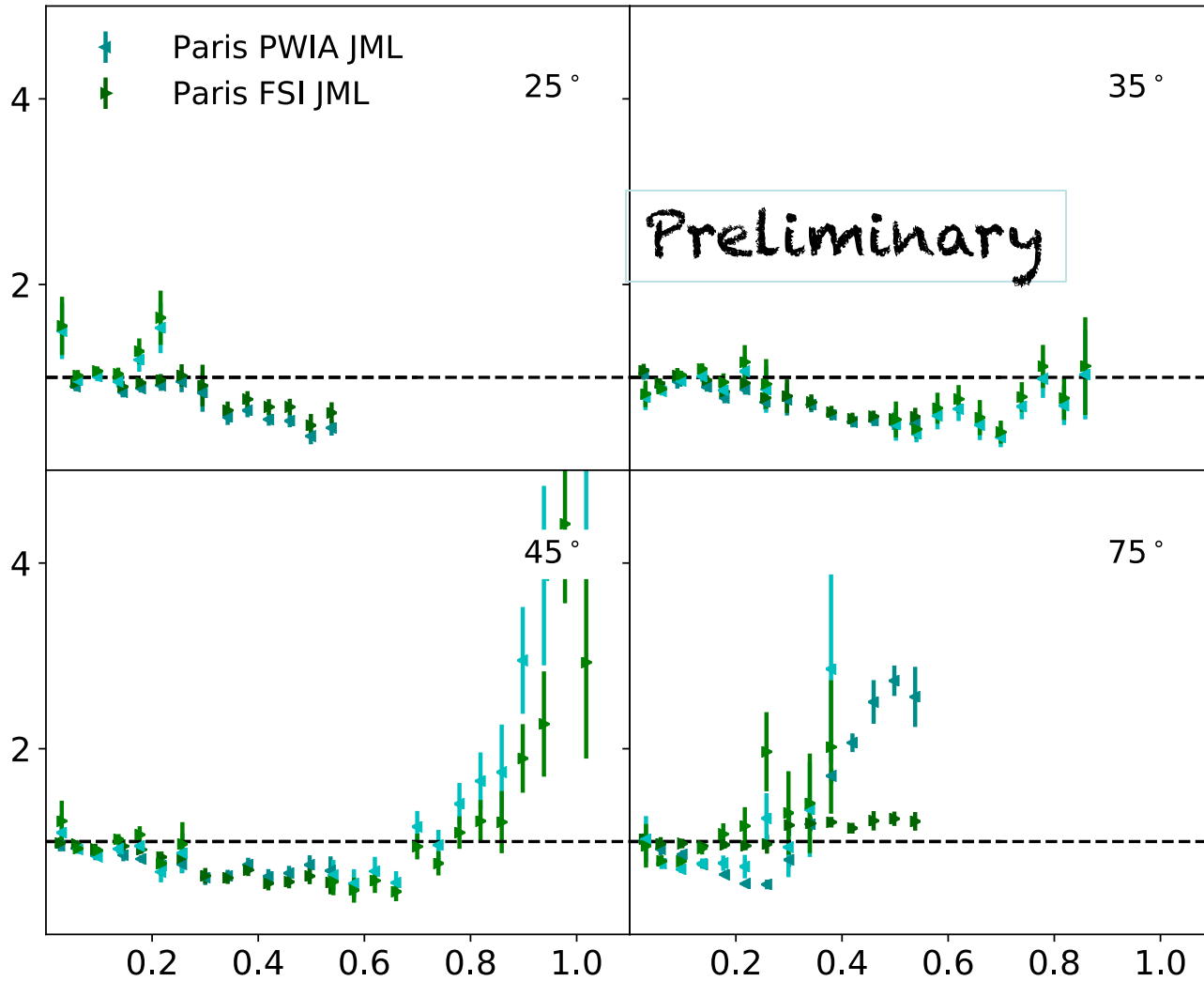
$$\sigma_{exp}/\sigma_{calc}$$



$$\sigma_{exp}/\sigma_{calc}$$



$$\sigma_{exp}/\sigma_{calc}$$



Experimental $\rho(\alpha, p_t)$ distributions

$$\rho_{EXP}(\alpha, p_T) = \frac{\sigma_{EXP}(\alpha, p_T)}{K \sigma_{eN}^{LC}(\alpha, p_T)}$$

- Hall A
- $Q^2 = 0.8, 2.1$ and 3.5 $(\text{GeV}/c)^2$: constant for each set
 - $p_{\text{miss}} = 0.2, 0.4$ and 0.5 GeV/c : angular distribution
 - $20^\circ \leq \theta_{nq} \leq 140^\circ$
 - angular range for each p_{miss} dependent on kinematics

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- Missing information due to finite spectrometer acceptance
- Interpolation necessary for missing data
- Various methods possible
- Large FSI at small α and large P_T for small Q^2

LC PWIA cross section

$$\frac{d\sigma}{dE'_e d\Omega_e d\Omega_p} = K \sigma_{eN}^{LC}(\alpha, p_t) \rho(\alpha, p_t)$$

Nuclear analog to parton distribution

$$f_N(\alpha) = \rho(\alpha) = \int \frac{\rho(\alpha, p_t)}{\alpha} d^2 p_t$$

Normalization:

$$\int \rho(\alpha, p_t) \frac{d\alpha}{\alpha} 2\pi p_t dp_t = 1$$

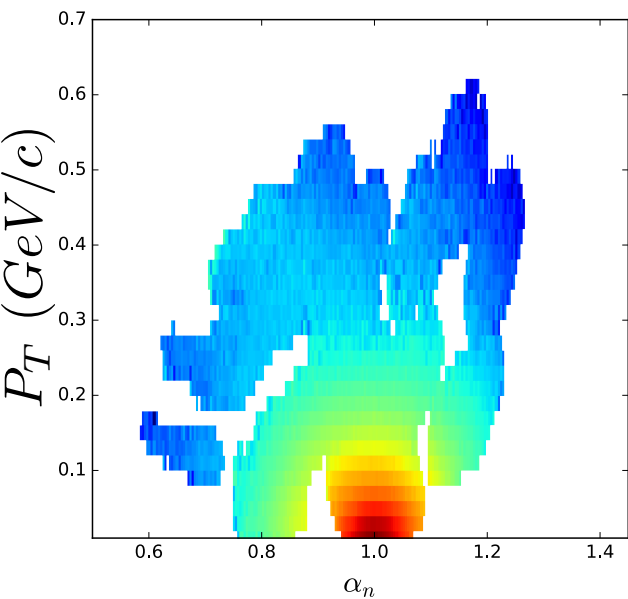
LC Momentum sum rule:

$$\int \alpha \rho(\alpha, p_t) \frac{d\alpha}{\alpha} 2\pi p_t dp_t = 1$$

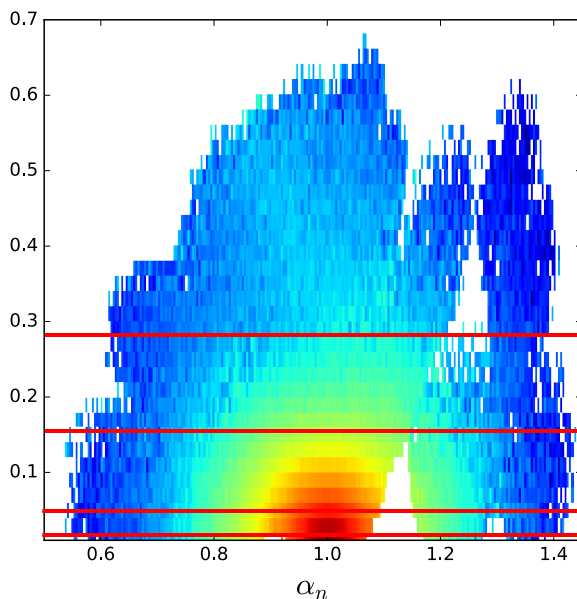
Problem: how do FSI affect $f_N(\alpha)$

First Results

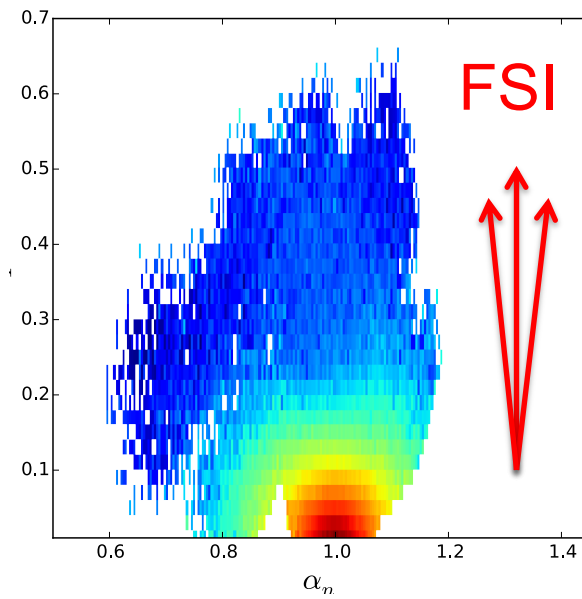
$0.8 (GeV/c)^2$



$2.1 (GeV/c)^2$

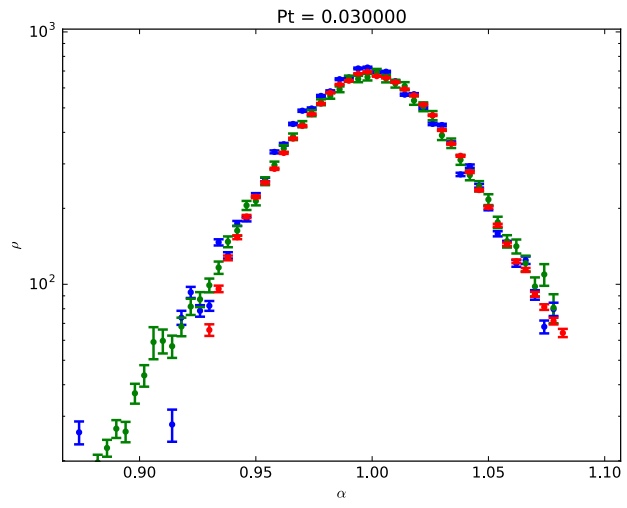
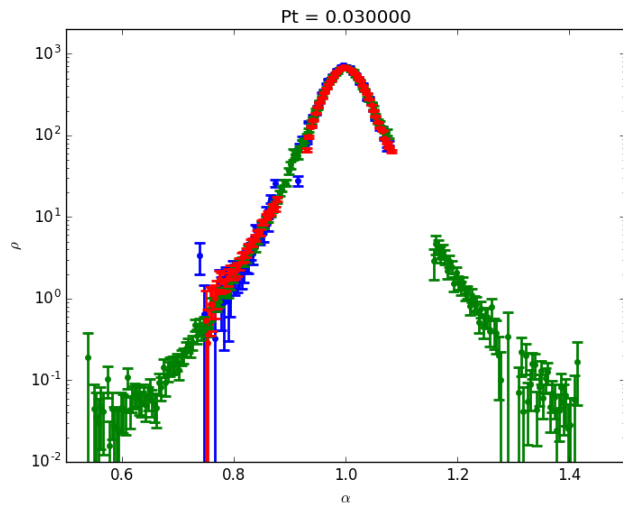
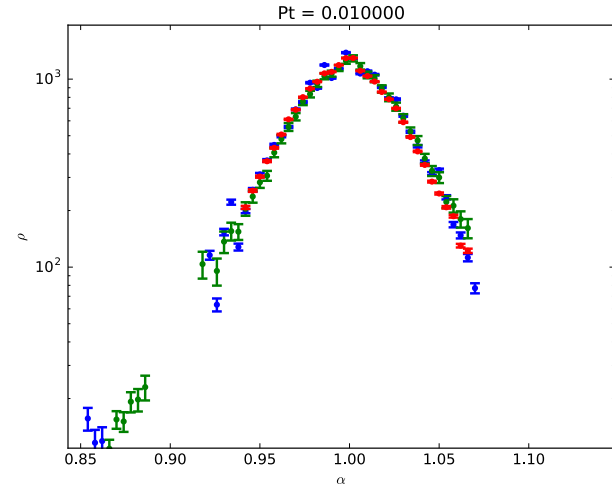
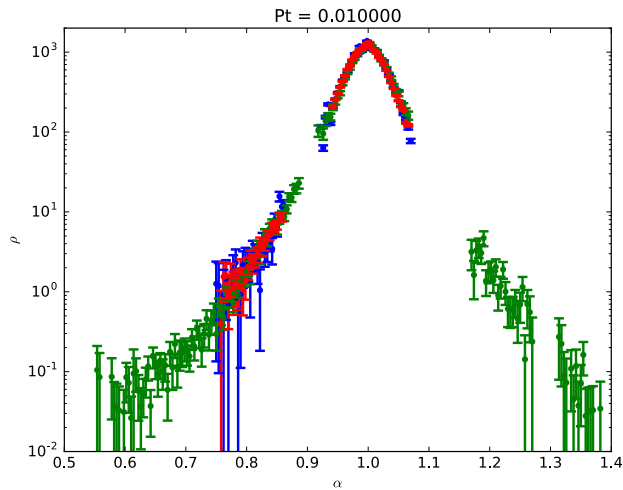


$3.5 (GeV/c)^2$

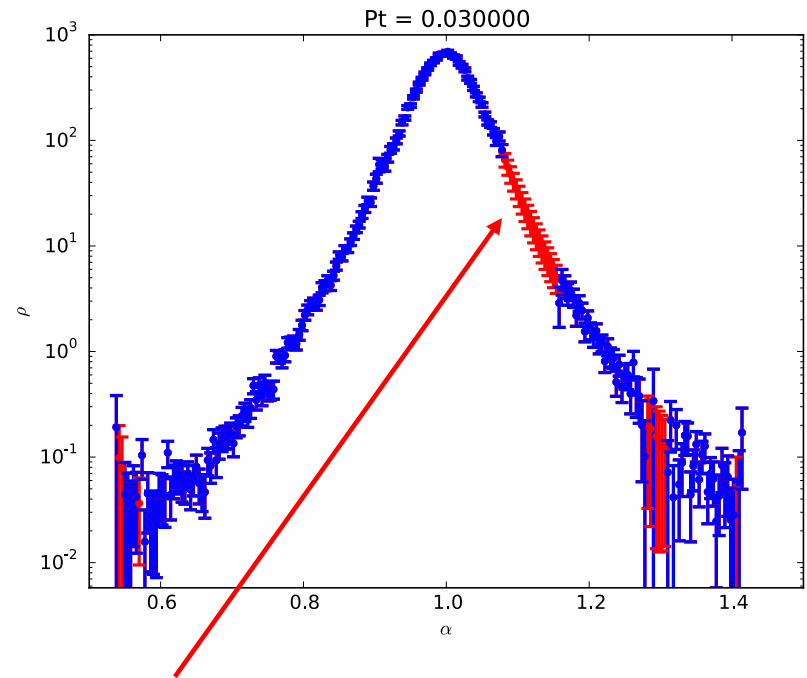
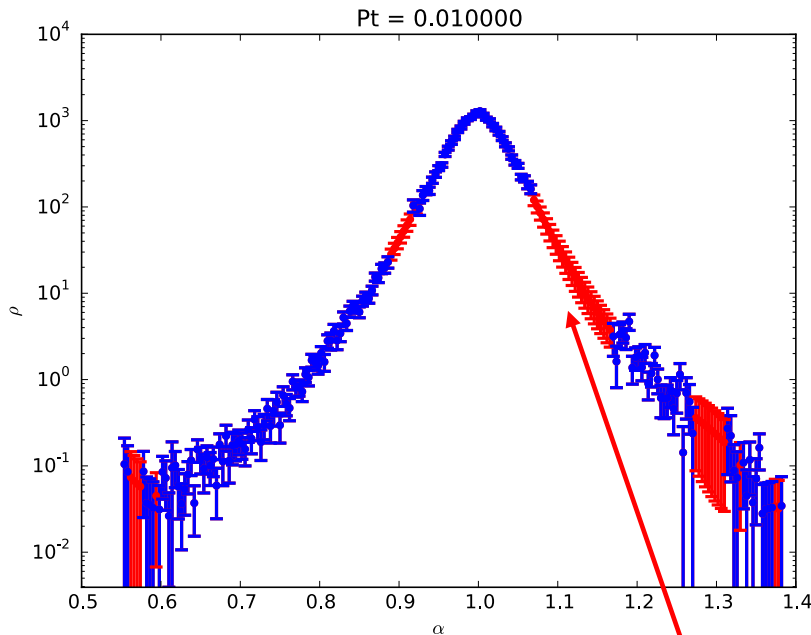


- Systematic error $\sim 7\%$
- Original experiment design to measure angular distributions

Small $P_T < 0.05$ GeV/c



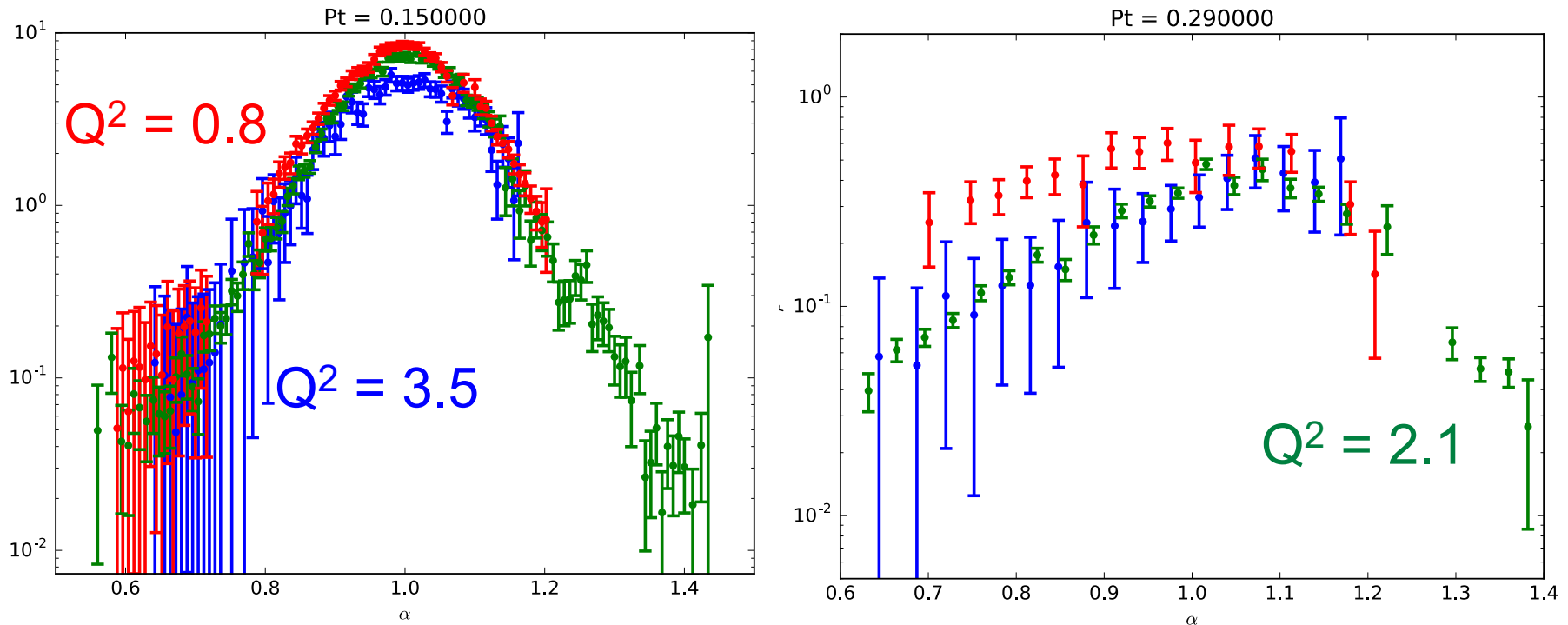
Missing Data



Missing data

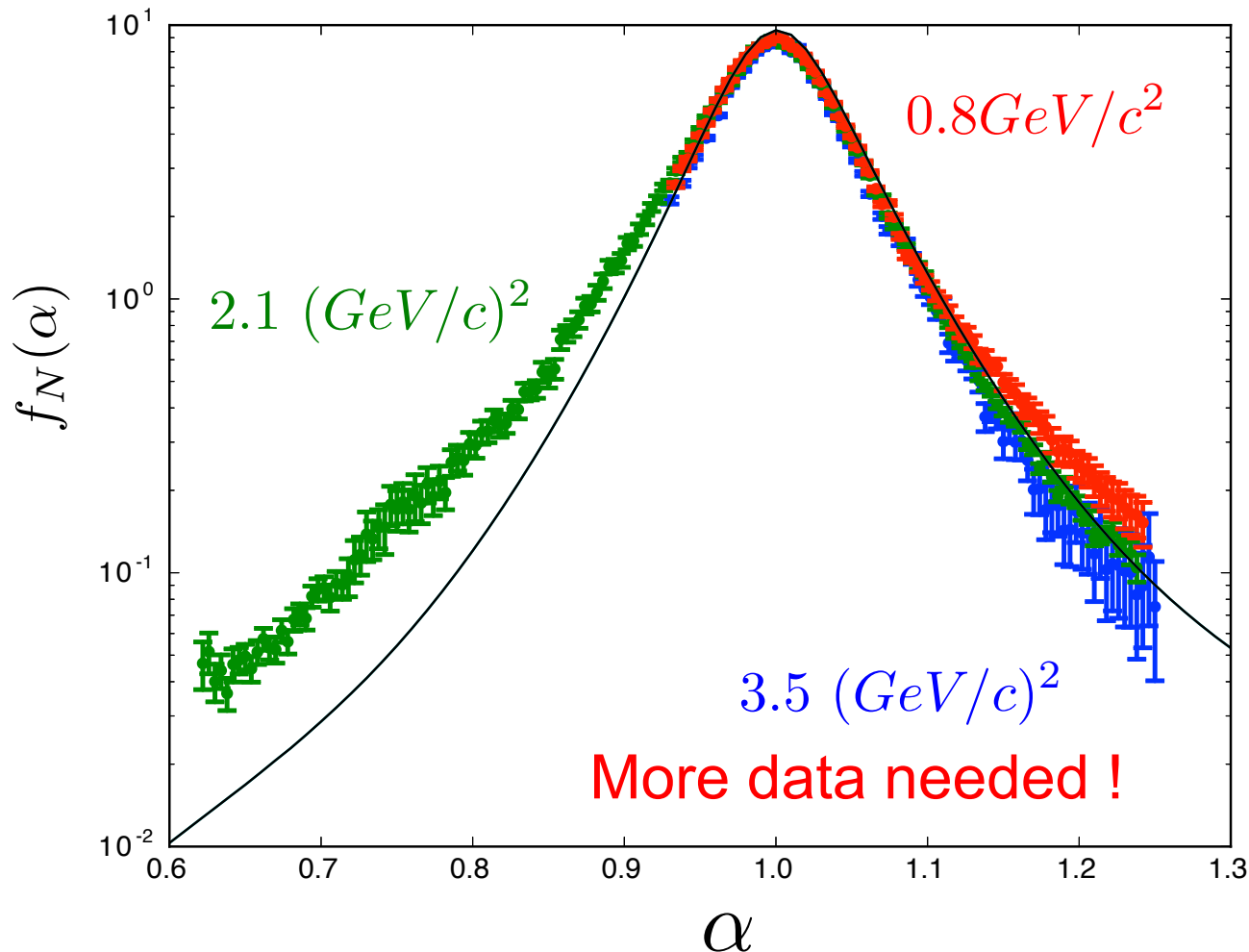
- Filled with fitting procedure
- Experimental data are not changed
- No extrapolation

Larger P_T



- $Q^2 = 0.8$ does not follow higher Q^2 behavior
- Qualitatively different (Large FSI etc.)

$$f_N(\alpha) = \rho(\alpha) = \int \frac{\rho(\alpha, p_t)}{\alpha} 2\pi p_t dp_t$$



$$I = \int f_N(\alpha) d\alpha$$

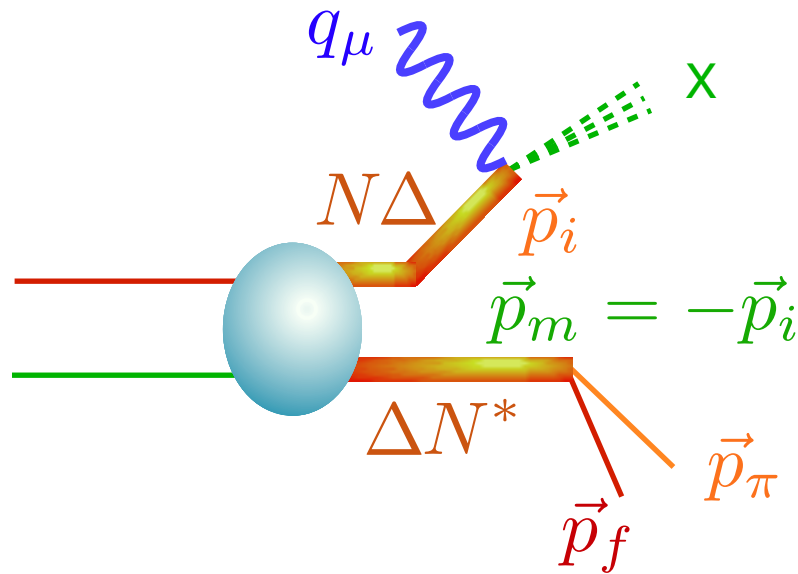
$$\frac{I_{exp}}{I_{PWIA}} \approx 1.3$$

$$\frac{I_{exp}}{I_{PWIA}} \approx 1.$$

$$\frac{I_{exp}}{I_{PWIA}} \approx 0.9$$

$$\Delta I / I \approx 0.1$$

Explore NN interaction core using inelastic channel: EIC



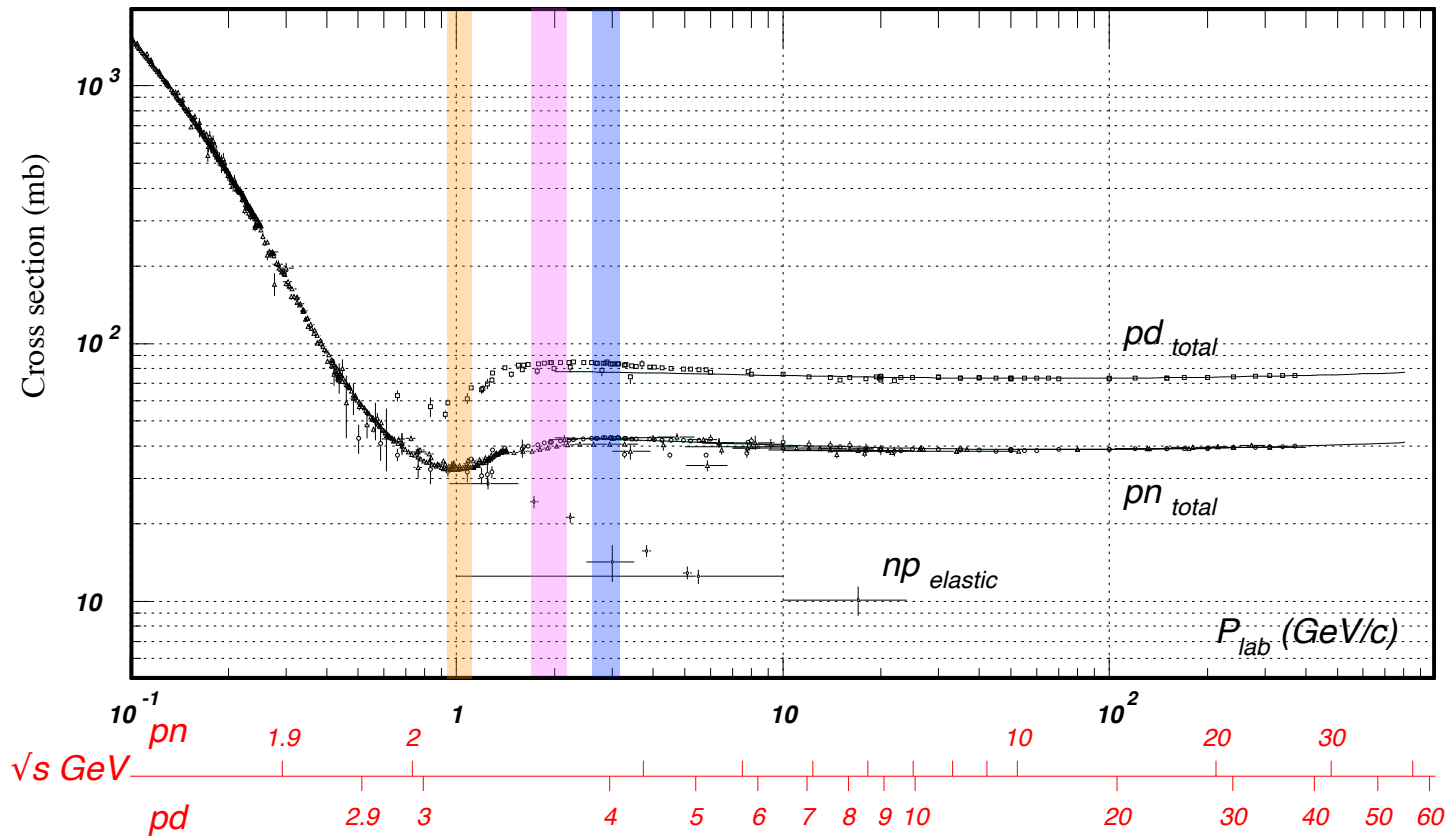
- Backwards Δ or N^*
- Large momentum $\sim 1\text{GeV}/c$

Summary

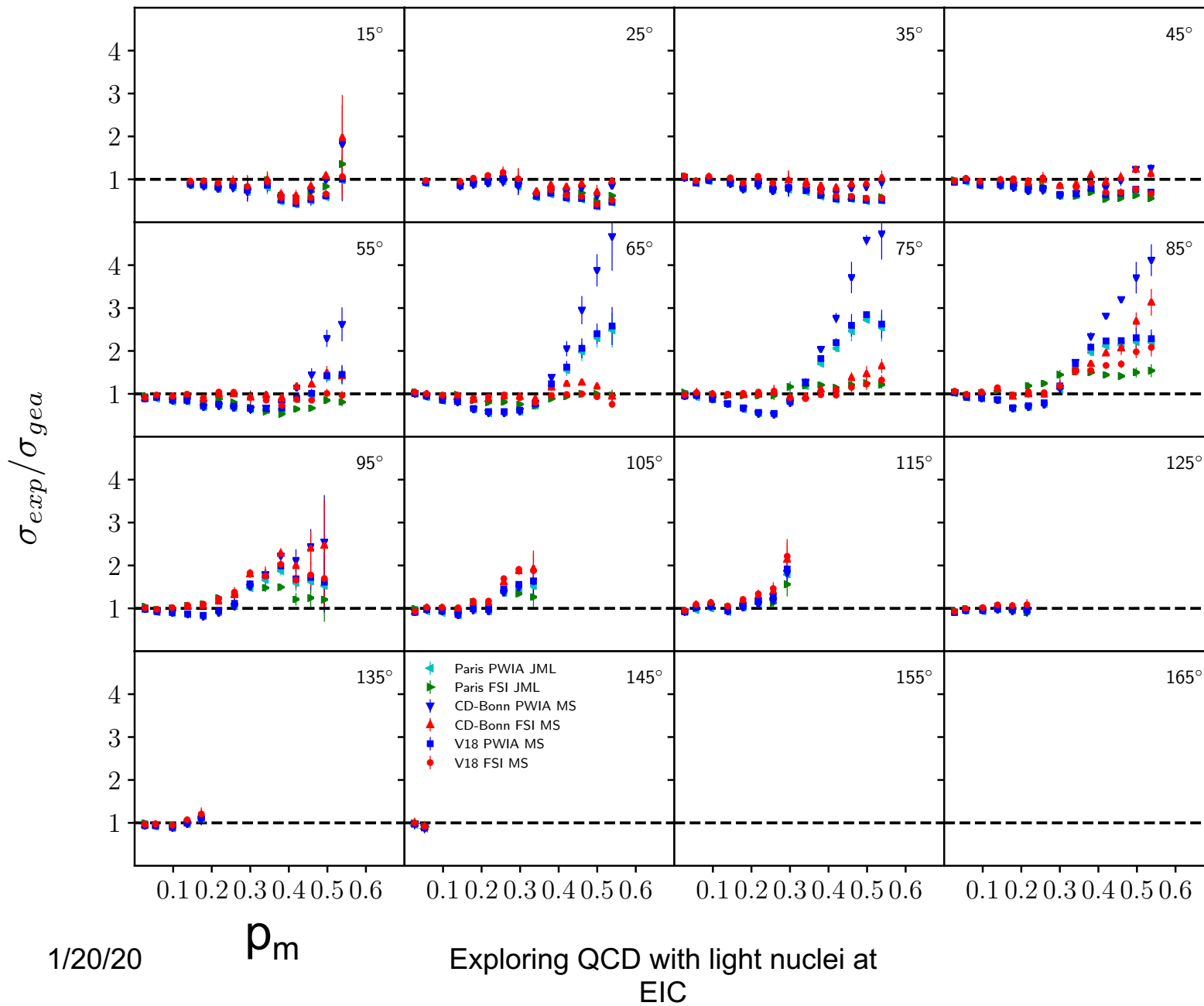
- High Q^2 $d(e, e' p)n$ can be described using generalized eikonal approximation for $Q^2 > 2 \text{ GeV}^2/c^2$
- There is a kinematic window to study the Deuteron momentum distribution.
- New data on very high missing momenta obtained in Hall C at JLAB for $Q^2 > 4 \text{ (GeV}^2/c^2)$ up to $p_m = 0.98 \text{ GeV}/c$
- Reduced cross section behavior different from Paris, V18 or CD-Bonn prediction
- LC momentum distribution determination needs more data
- Inelastic channels should be included (EIC)
Supported by the Department of Energy, contract DESC0013620

Backup

0.8 2.1 3.5 (GeV/c)²



$Q^2 = 3.5 \text{ (GeV/c)}^2$ GEA M.Sargsian



Extracting ρ_{LC}

Relativistic Description of the Deuteron, L.L Frankfurt and M. Strikman, Nuclear Physics **B148** (1979) 107

High-Energy Phenomena, Short-Range Nuclear Structure and QCD, L.L Frankfurt and M. Strikman, Physics Reports **76**, (1981) 215

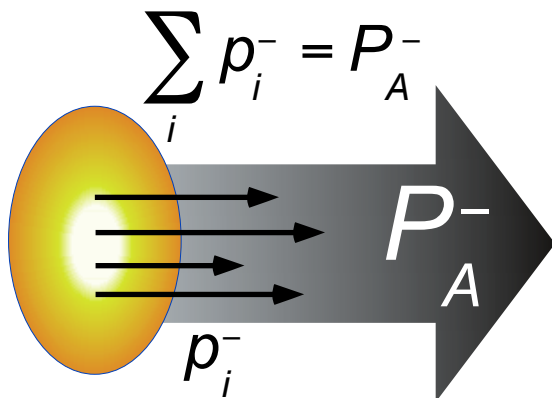
LC momentum

$$p^- = E - p_z$$

LC momentum fraction

$$\alpha = A \frac{p_i^-}{P_A^-}$$

α is frame independent for boosts along the z-axis



analogous to “x” for quark distributions

Spectator (neutron) momentum fraction $\alpha_s = 2 \frac{E_s - p_s^z}{M_D}$
from experiment

remember in lab: $P_D^- = M_D$ and

Proton momentum fraction $\alpha = 2 - \alpha_s$

Small α large p_z

$\alpha_s \rightarrow 2$ $p_{pz} \rightarrow \infty$

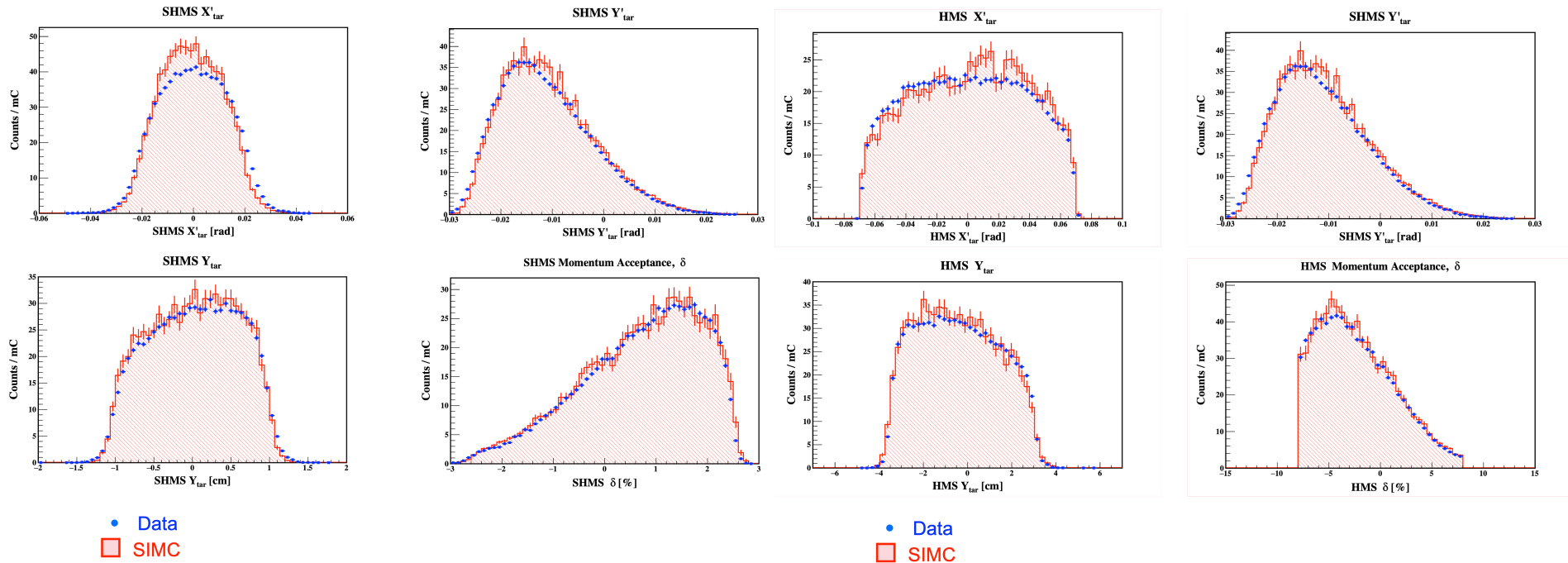
Advantages of working on LC:

- at high Q^2 , FSI is mostly transverse α is approx. conserved by FSI (M.Sargsian Int. J. Mod. Phys. E10 2001)
- $\rho(\alpha)$ is very little affected by re-scattering
- at high energies: $N\bar{N}$ become important but
- unimportant on LC (photon energy is 0)
- $\rho(\alpha)$ necessary for interpretation of DIS data of nuclei

Detector Performance

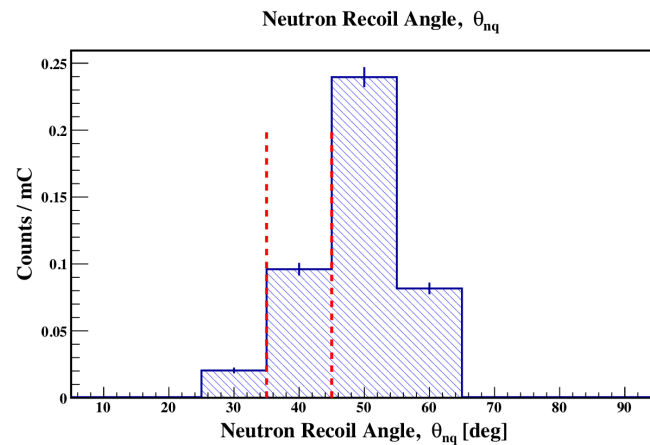
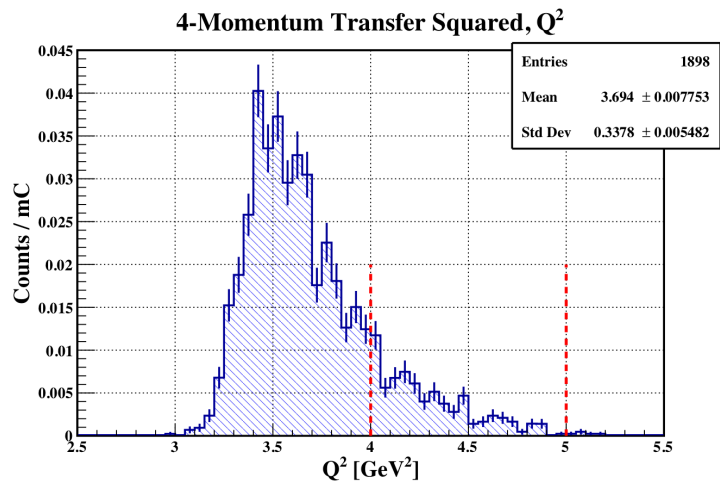
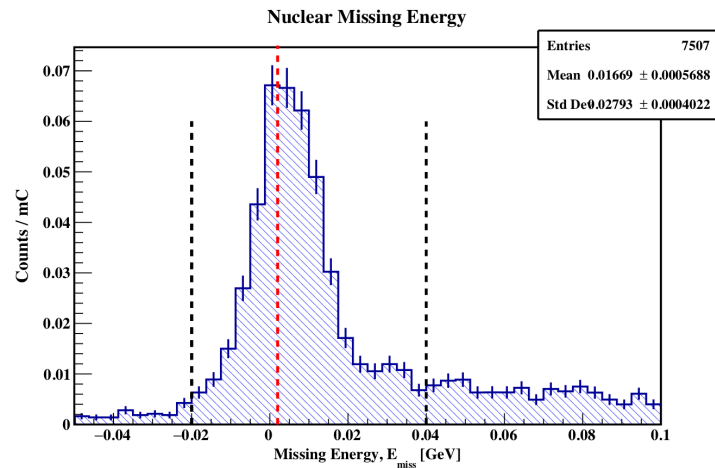
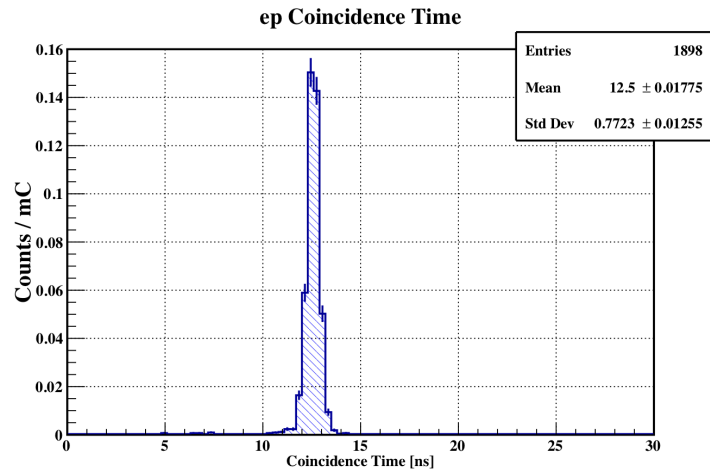
SHMS

HMS



Kinematic Variables

$p_m = 750 \text{ MeV}/c$



- Data

Problems

- **Reaction dynamics:**
 - how does the photon interact with a deeply bound nucleon ?
 - what is the EM current structure ?
- **Final State Interactions**
 - high Q^2 : eikonal approximations valid ?
(Answer: **yes** for $Q^2 > 2$ (GeV/c)²)
- **Deuteron wave function**
 - can one probe NN wave function at small distances ?
 - can one find evidence for new degrees of freedom ?

All these problems are interconnected
High Q^2 data are necessary for progress!