

x-dependent GPDs from Lattice QCD (twist-2 & twist-3)

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QCD Evolution 2021

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Collaborators

Twist-2 GPDs

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Twist-3 GPDs

- ▶ **S. Bhattacharya**
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- ▶ **K. Cichy**
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- ▶ **A. Scapellato**
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- ▶ **F. Steffens**
University of Bonn

Relevant publications

Twist-2 GPDs

- *Unpolarized and helicity generalized parton distributions of the proton within lattice QCD*
C.Alexandrou, K.Cichy, M.C, K. Hadjiyiannakou, K.Jansen, A.Scapellato, F.Steffens, **PRL 125 (2020), [arXiv:2008.10573]**

Twist-3 PDFs

- *New insights on proton structure from lattice QCD: the twist-3 parton distribution function $g_T(x)$*
S.Bhattacharya, K.Cichy, M.C, A.Metz, A.Scapellato, F.Steffens, **PRD Editors' Suggestion, [arXiv:2004.04130]**
- *One-loop matching for the twist-3 parton distribution $g_T(x)$*
S.Bhattacharya, K.Cichy, M.C, A.Metz, A.Scapellato, F.Steffens, **PRD 102 (2020) 3, 034005, [arXiv:2005.10939]**
- *The role of zero-mode contributions in the matching for the twist-3 PDFs $e(x)$ and $h_L(x)$*
S. Bhattacharya, K. Cichy, M.C, A. Metz, A. Scapellato, F. Steffens, **PRD 102 (2020) 114025, [arXiv:2006.12347]**

Why GPDs?

- ★ GPDs provide information on spatial distribution of partons inside the hadron, and its mechanical properties (OAM, pressure, etc.)

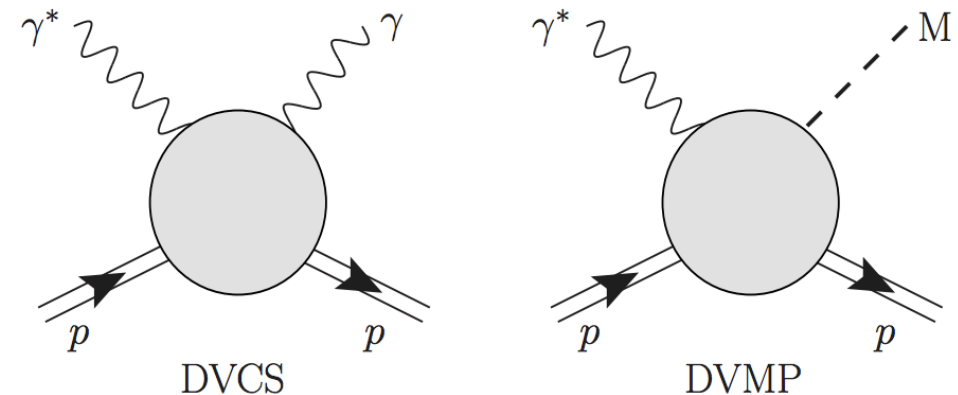
[M. Burkardt, Phys.Rev.D62 071503 (2000), hep-ph/0005108]

[M. V. Polyakov, Phys. Lett. B555 (2003) 57, hep-ph/0210165]

- ★ Experimentally accessed in DVCS and DVMP

[X. D. Ji, Phys. Rev. Lett. 78, 610 (1997), hep-ph/9603249]

(Halls A,B,C (JLab),
PHENIX, STAR, HERMES,
COMPASS, GSI, BELLE, J-PARC)



- ★ Experimentally, GPDs are not well-constrained:

- independent measurements to disentangle GPDs
- limited coverage of kinematic region
- data on certain GPDs
- indirectly related to GPDs through the Compton FFs
- GPDs phenomenology more complicated than PDFs (multi-dimensionality)

Why twist-3 GPDs?

$$f_i = \underset{\substack{\uparrow \\ \text{Twist-2}}}{f_i^{(0)}} + \underset{\substack{\uparrow \\ \text{Twist-3}}}{\frac{f_i^{(1)}}{Q}} + \frac{f_i^{(2)}}{Q^2} \dots$$

Higher-twist distributions:

- ★ Lack density interpretation, but can be sizable
- ★ Sensitive to soft dynamics
- ★ Challenging to probe experimentally and isolate from leading-twist
- ★ Needed for proton tomography
- ★ Related to certain spin-orbit correlations [C. Lorce, PLB 735 (2014) 344, arXiv:1401.7784]
- ★ Estimation on power corrections in hard exclusive processes (DVCS)
- ★ $[\widetilde{H} + \widetilde{G}_2](x, \xi, t)$ related to tomography of $F_\perp$ acting on the active q in DIS off a transversely polarized N right after the virtual photon absorbing

[M. Burkardt, PRD 88 (2013) 114502, arXiv:0810.3589]

- ★ $G_2(x, \xi, t)$ related to L_q :

$$L_q = - \int_{-1}^1 dx x G_2^q(x, \xi, t = 0)$$

[X. D. Ji, Phys. Rev. Lett. 78, 610 (1997), hep-ph/9603249], [M. Penttinen et al., PLB 491 (2000) 96, arXiv:hep-ph/0006321]

GPDs from lattice QCD

(quasi-distributions method)

Useful Reviews:

[K. Cichy, M. Constantinou, Advances in HEP, Volume 2019, Article ID 3036904, [arXiv:1811.07248](#)]

[X. Ji, Y.-S. Liu, Y. Liu, J.-H. Zhang, and Y. Zhao (2020), [2004.03543](#)]

[M. Constantinou (invited review) Eur. Phys. J. A 57 (2021) 2, 77, [arXiv:2010.02445](#)]

Access of GPDs on a Euclidean Lattice

- ★ Relies on Large Momentum Effective Theory (**LaMET**) to reconstruct GPDs
[X. Ji, Phys. Rev. Lett. 110 (2013) 262002; X. Ji, Sci. China Phys. Mech. Astron. 57, 1407 (2014)]

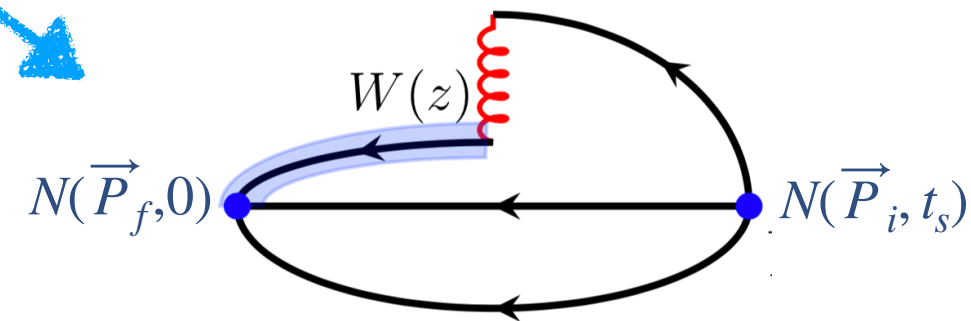
- ★ Matrix elements of spatial operators with **fast moving hadrons**

$$\tilde{q}_{\Gamma}^{\text{GPD}}(x, t, \xi, P_3, \mu) = \int \frac{dz}{4\pi} e^{-i x P_3 z} \langle N(P_f) | \bar{\Psi}(z) \Gamma \mathcal{W}(z, 0) \Psi(0) | N(P_i) \rangle_{\mu}$$

$$\begin{aligned} \Delta &= P_f - P_i \\ t &= \Delta^2 = -Q^2 \\ \xi &= \frac{Q_3}{2P_3} \end{aligned}$$

Variables:

- length of the Wilson line (z)
 - nucleon momentum boost (P_3)
 - momentum transfer (t)
 - skewness (ξ)
- } PDFs & GPDs
- } GPDs



- ★ Challenges of calculation:

- Increased statistical uncertainties due to momentum transfer
- Need for multiple matrix elements to disentangle GPDs
- Frame dependence
- Matching for nonzero skewness, twist-3

**hadronic
matrix elements**

1

**Matching to
light-cone GPDs**

6

**Identification of
ground state**

2

**quasi
distribution
approach**

5

**x-dependence
reconstruction**

Renormalization

3

4

**“form factors”
disentanglement**

High-complexity calculation

Computational and theoretical challenges

each step may introduce systematic uncertainties

Parameters of calculation

(u-d flavor combination)

★ Nf=2+1+1 twisted mass fermions & clover term

★ Ensemble parameters:

Pion mass: 260 MeV
Lattice spacing: 0.093 fm
Volume: $32^3 \times 64$
Spatial extent: 3 fm

★ Kinematical setup:

Twist-2

P_3 [GeV]	$\vec{Q} \times \frac{L}{2\pi}$	$-t$ [GeV ²]	ξ	N_{confs}	N_{meas}
0.83	(0,2,0)	0.69	0	519	4152
1.25	(0,2,0)	0.69	0	1315	42080
1.67	(0,2,0)	0.69	0	1753	112192
1.25	(0,2,2)	1.39	1/3	417	40032
1.25	(0,2,-2)	1.39	-1/3	417	40032



zero and nonzero
skewness

Twist-3

P_3 [GeV]	$\vec{Q} \times \frac{L}{2\pi}$	$-t$ [GeV ²]	ξ	N_{meas}
0.83	(2,0,0)	0.69	0	4288
1.25	(2,0,0)	0.69	0	4288
1.25	(2,2,0)	1.39	0	4288
1.67	(2,0,0)	0.69	0	4288

We utilize

$\pm P_3$

$\pm \vec{Q}$

$\gamma^j \gamma^5, \quad j = 1, 2$

★ Excited states: T_{sink} up to 1.13 fm

GPDs Decomposition

[D. Kiptily and M. Polyakov, Eur. Phys. J. C37 (2004) 105, arXiv:hep-ph/0212372]

[F. Aslan et al., Phys. Rev. D 98, 014038 (2018), arXiv:1802.06243]

$$\begin{aligned}
 F^\mu &= P^\mu \frac{h^+}{P^+} H + P^\mu \frac{e^+}{P^+} E \\
 &+ \Delta_\perp^\mu \frac{b}{2m} G_1 + h_\perp^\mu (H + E + G_2) + \Delta_\perp^\mu \frac{h^+}{P^+} G_3 + \tilde{\Delta}_\perp^\mu \frac{\tilde{h}^+}{P^+} G_4, \\
 \tilde{F}^\mu &= P^\mu \frac{\tilde{h}^+}{P^+} \tilde{H} + P^\mu \frac{\tilde{e}^+}{P^+} \tilde{E} \\
 &+ \Delta_\perp^\mu \frac{\tilde{b}}{2m} (\tilde{E} + \tilde{G}_1) + \tilde{h}_\perp^\mu (\tilde{H} + \tilde{G}_2) + \Delta_\perp^\mu \frac{\tilde{h}^+}{P^+} \tilde{G}_3 + \tilde{\Delta}_\perp^\mu \frac{h^+}{P^+} \tilde{G}_4,
 \end{aligned}$$

$$\begin{aligned}
 h^\mu &= \bar{u}(p') \gamma^\mu u(p), & e^\mu &= \bar{u}(p') \frac{i\sigma^{\mu\nu} \Delta_\nu}{2m} u(p), & b &= \bar{u}(p') u(p), \\
 \tilde{h}^\mu &= \bar{u}(p') \gamma^\mu \gamma_5 u(p), & \tilde{e}^\mu &= \frac{\Delta^\mu}{2m} \tilde{b}, & \tilde{b} &= \bar{u}(p') \gamma_5 u(p)
 \end{aligned}$$

Twist-2 GPDs

Decomposition of matrix elements (twist-2)

Unpolarized →

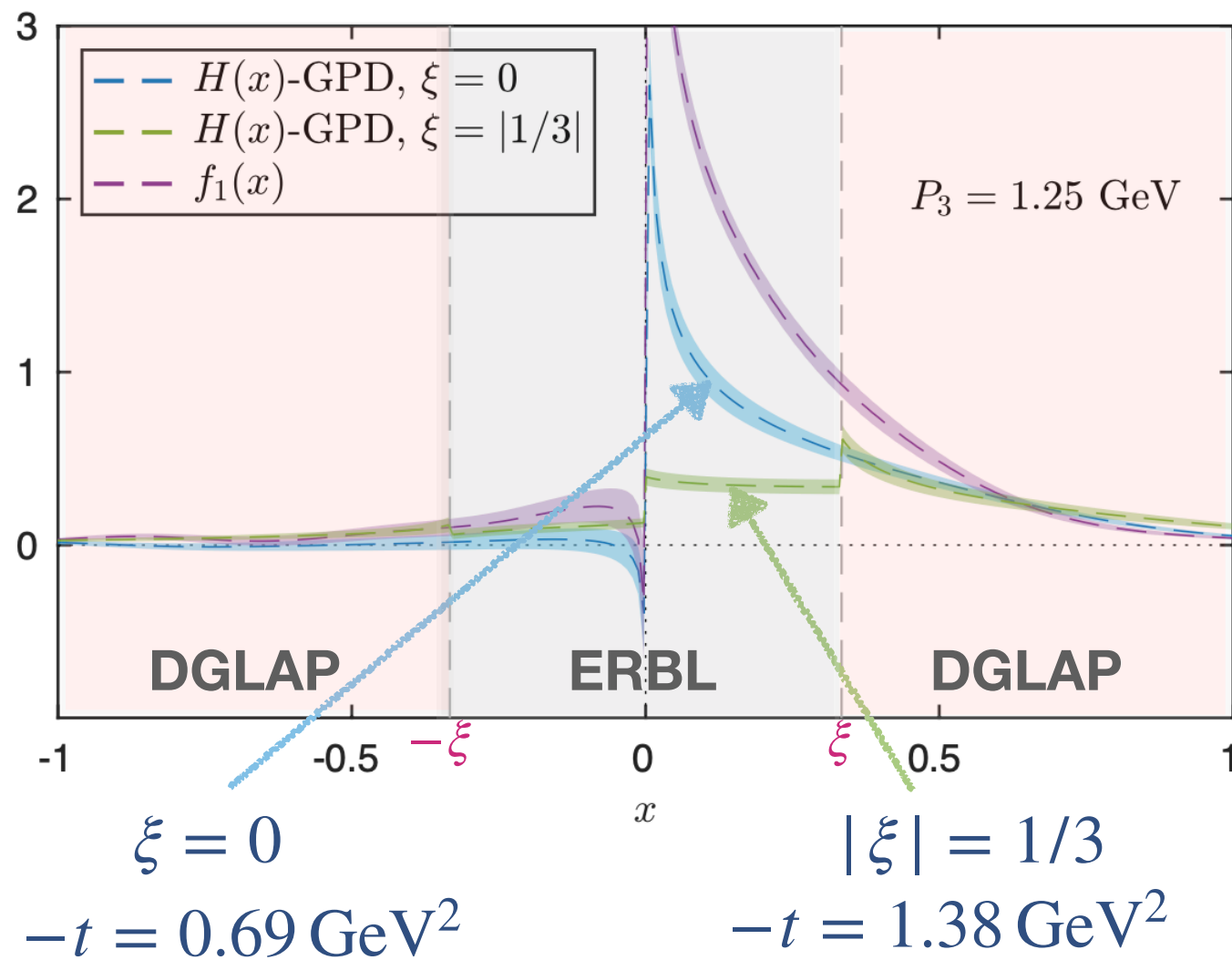
$$\left\{ \begin{aligned} \Pi_{\gamma^0}(\mathcal{P}_0; P_f, P_i) &= C \left[F_H(Q^2) \left(\frac{E_f E_i}{2m^2} + \frac{E_f + E_i}{4m} + \frac{P_{f\rho} P_{i\rho}}{4m^2} + \frac{1}{4} \right) \right. \\ &\quad \left. + F_E(Q^2) \left(P_{f\rho} P_{i\rho} \left(\frac{E_f + E_i}{8m^3} + \frac{1}{4m^2} \right) - \frac{(E_f - E_i)^2}{8m^2} + \frac{E_f + E_i}{8m} + \frac{1}{4} \right) \right], \\ \Pi_{\gamma^0}(\mathcal{P}_j; P_f, P_i) &= i \epsilon_{j0\rho\tau} C \left[F_H(Q^2) \frac{P_{f\rho} P_{i\tau}}{4m^2} + F_E(Q^2) \left(\frac{(E_f + E_i) P_{f\rho} P_{i\tau}}{8m^3} + \frac{P_{f\rho} P_{i\tau} - P_{f\tau} P_{i\rho}}{8m^2} \right) \right], \end{aligned} \right.$$

Helicity →

$$\left\{ \begin{aligned} \Pi_{\gamma^3 \gamma^5}(\mathcal{P}_j; P_f, P_i) &= i C \left[F_{\tilde{H}}(Q^2) \left(\delta_{j3} \left(\frac{E_f + E_i}{4m} - \frac{P_{f\rho} P_{i\rho}}{4m^2} + \frac{1}{4} \right) + \frac{P_{i3} P_{fj} + P_{f3} P_{ij}}{4m^2} \right) \right. \\ &\quad \left. - F_{\tilde{E}}(Q^2) \frac{(P_{f3} - P_{i3})(-E_f P_{ij} + E_i P_{fj} + m(P_{fj} - P_{ij}))}{8m^3} \right], \end{aligned} \right.$$

- Kinematic factors defined by calculation setup
- Decomposition similar to the decomposition of form factors
- Zero skewness: $F_{\tilde{E}}(Q^2) = 0$

Comparison of zero and nonzero skewness



◆ ERBL region ($|x| < \xi$):

- amplitude of removing a quark-antiquark pair with momentum $-\Delta$

◆ DGLAP region ($\xi < |x| < 1$):

- amplitude of removing a quark/antiquark from the hadron and inserting it back with momentum increase Δ

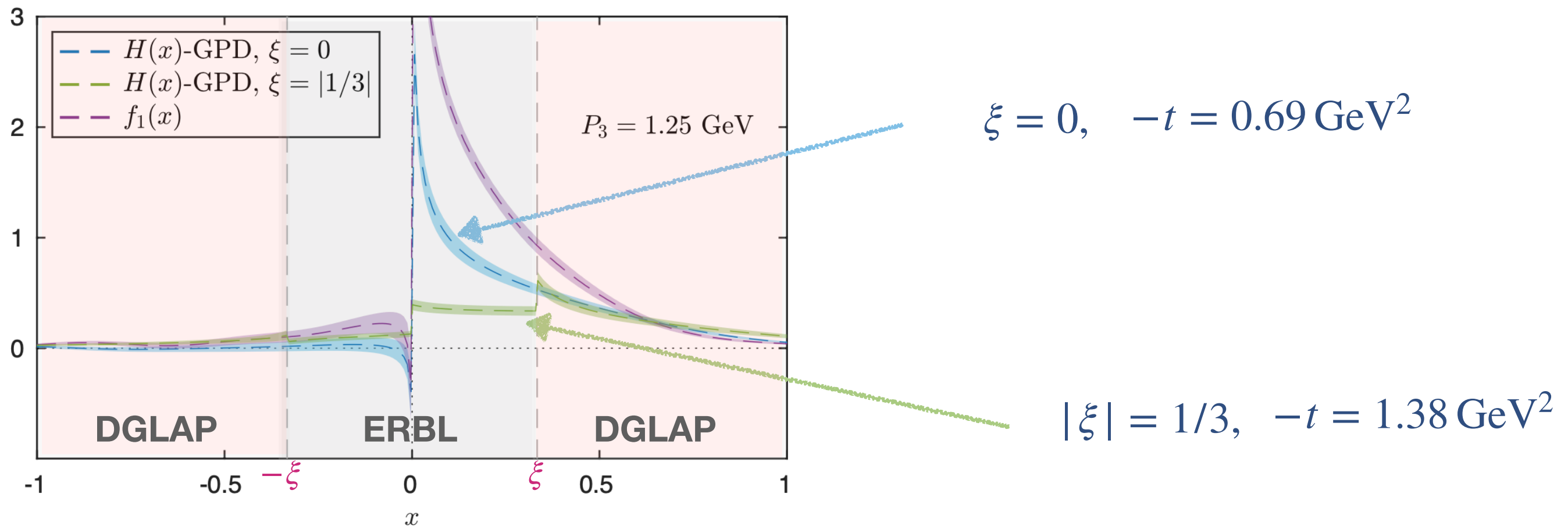
◆ ERBL region:

- $f_1(x)$ dominant
- $H(x, 1/3)$ suppressed compared to $H(x, 0)$

◆ DGLAP region:

- $H(x, 1/3)$ similar to $H(x, 0)$

Comparison of zero and nonzero skewness

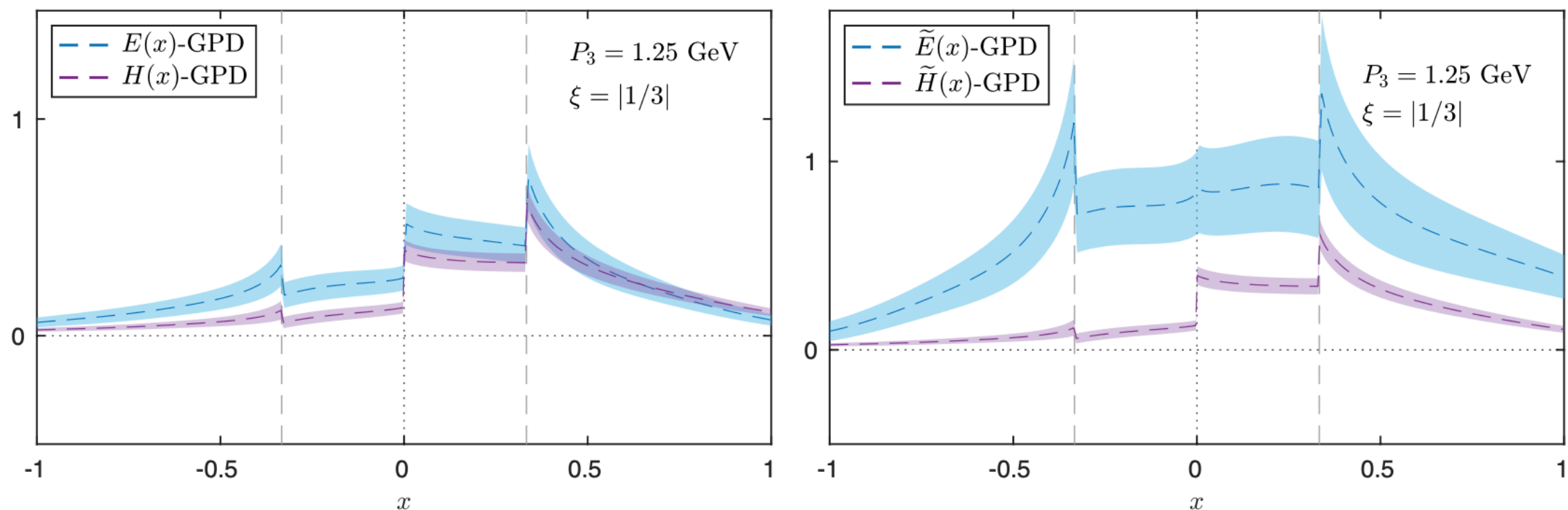


$x \rightarrow 1$ region: qualitatively comparison with power counting analysis

[F. Yuan, Phys.Rev. D69 (2004) 051501, hep-ph/0311288]

- ♦ t -dependence vanishes at large- x
- ♦ $H(x,0)$ asymptotically equal to $f_1(x)$
- ♦ $1/(1 - \xi)^2$ dependence predicts GPDs larger than PDFs

GPDs comparison



- ◆ H -GPD and E -GPD are similar in magnitude

important contribution in the proton spin

$$\int_{-1}^{+1} dx x^2 H^q(x, \xi, t) = A_{20}^q(t) + 4\xi^2 C_{20}^q(t), \quad \int_{-1}^{+1} dx x^2 E^q(x, \xi, t) = B_{20}^q(t) - 4\xi^2 C_{20}^q(t)$$

- ◆ $\widetilde{E}$ -GPD dominant compared to $\widetilde{H}$ -GPD

Twist-3 GPDs

Helicity flip twist-3 GPDs

★ Transverse matrix element of axial operator (proton boost: $\vec{P} = (0,0,P_3)$)

$$\begin{aligned}\tilde{F}^\mu &= P^\mu \frac{\tilde{h}^+}{P^+} \tilde{H} + P^\mu \frac{\tilde{e}^+}{P^+} \tilde{E} \\ &+ \Delta_\perp^\mu \frac{\tilde{b}}{2m} (\tilde{E} + \tilde{G}_1) + \tilde{h}_\perp^\mu (\tilde{H} + \tilde{G}_2) + \Delta_\perp^\mu \frac{\tilde{h}^+}{P^+} \tilde{G}_3 + \tilde{\Delta}_\perp^\mu \frac{h^+}{P^+} \tilde{G}_4\end{aligned}$$

[D. Kiptily and M. Polyakov, Eur. Phys. J. C37 (2004) 105, arXiv:hep-ph/0212372]

[F. Aslan et al., Phys. Rev. D 98, 014038 (2018), arXiv:1802.06243]

★ Sum Rules (generalization of Burkhardt-Cottingham)

$$\int_{-1}^1 dx \tilde{H}(x, \xi, t) = G_A(t), \quad \int_{-1}^1 dx \tilde{E}(x, \xi, t) = G_P(t), \quad \int_{-1}^1 dx \tilde{G}_i(x, \xi, t) = 0$$

[X. D. Ji, Phys. Rev. Lett. 78, 610 (1997), hep-ph/9603249]

★ Sum Rules (generalization of Efremov-Leader-Teryaev)

[A. Efremov, O. Teryaev, E. Leader, PRD 55 (1997) 4307, hep-ph/9607217]

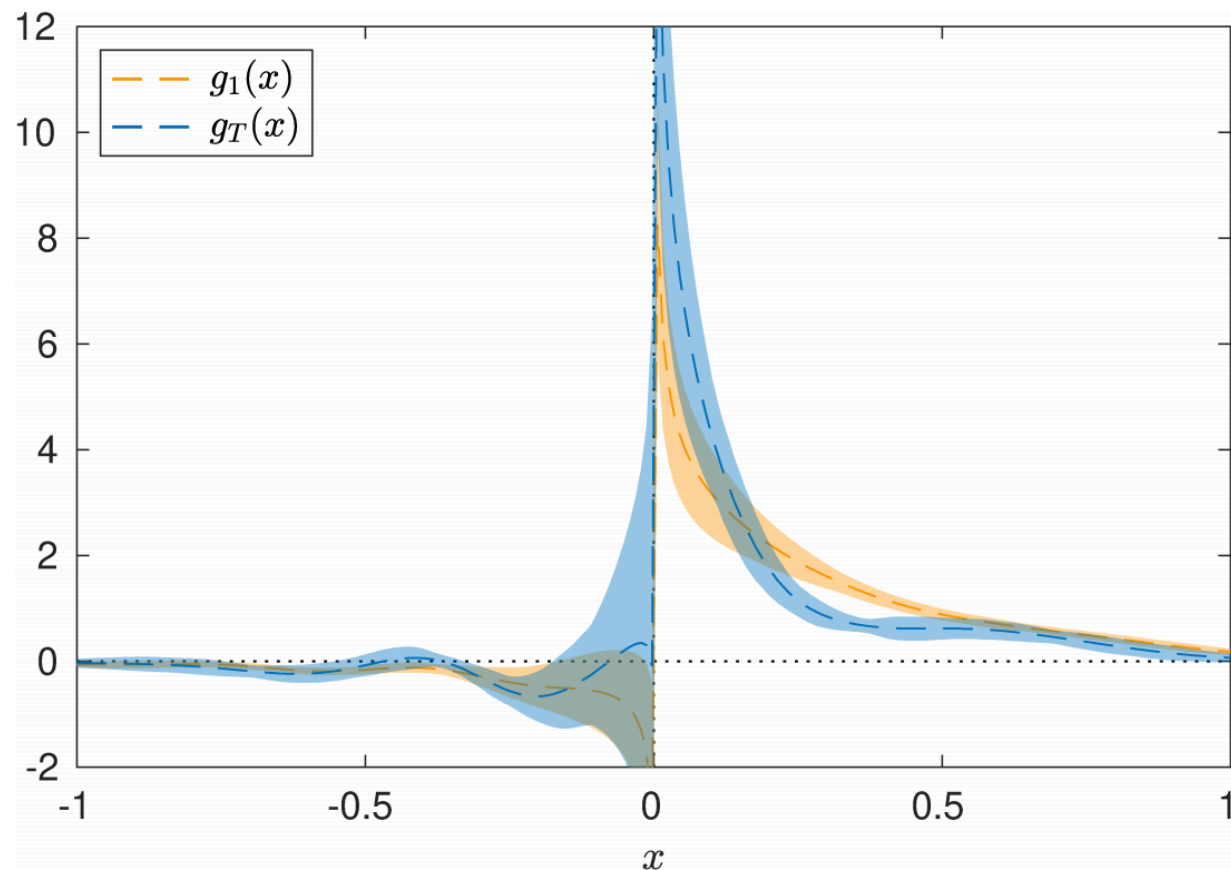
$$\begin{aligned}\int_{-1}^1 dx x \tilde{G}_1(x, \xi, t) &= \frac{1}{2} \left[F_2(t) + \left(\xi \frac{\partial}{\partial \xi} - 1 \right) \int_{-1}^1 dx x \tilde{E}(x, \xi, t) \right], & \int_{-1}^1 dx x \tilde{G}_2(x, \xi, t) &= \frac{1}{2} \left[\xi^2 G_E(t) - \frac{t}{4m^2} F_2(t) - \tilde{A}_{20}(t) \right], \\ \int_{-1}^1 dx x \tilde{G}_3(x, \xi, t) &= \frac{\xi}{4} G_E(t), & \int_{-1}^1 dx x \tilde{G}_4(x, \xi, t) &= \frac{1}{4} G_E(t)\end{aligned}$$

Sum rules simplify for zero skewness

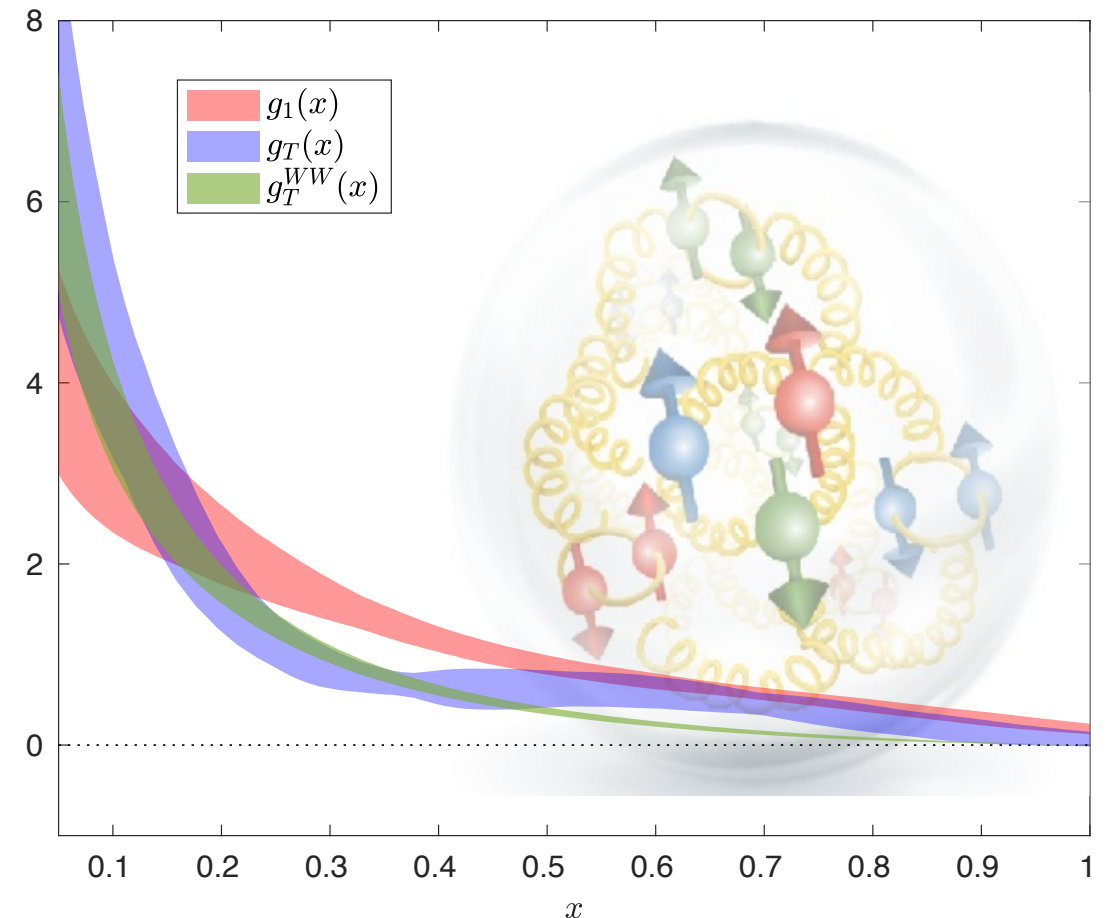
F_2 : Pauli FF
 G_E : electric FF
 $\tilde{A}_{20}, \tilde{B}_{20}$: axial GFFs

Results on twist-3 PDFs

★ Helicity flip ($g_T(x)$) and transversity ($h_L(x)$) twist-3 PDFs



PRD Editors' Suggestion Highlight



[S. Bhattacharya et al., PRD 102 (2020) 11, arXiv:2004.04130]

See talks @ GHP 2021 by:

S. Bhattacharya (on matching)

A. Scapellato (on lattice results)

Decomposition of matrix elements

Decomposition in Minkowski space

[F. Aslan et al., Phys. Rev. D 98, 014038 (2018), arXiv:1802.06243]

$$\tilde{F}^\mu = P^\mu \frac{\tilde{h}^+}{P^+} \tilde{H} + P^\mu \frac{\tilde{e}^+}{P^+} \tilde{E}$$

$$\begin{aligned} h^\mu &= \bar{u}(p') \gamma^\mu u(p), & e^\mu &= \bar{u}(p') \frac{i\sigma^{\mu\nu} \Delta_\nu}{2m} u(p), & b &= \bar{u}(p') u(p), \\ \tilde{h}^\mu &= \bar{u}(p') \gamma^\mu \gamma_5 u(p), & \tilde{e}^\mu &= \frac{\Delta^\mu}{2m} \tilde{b}, & \tilde{b} &= \bar{u}(p') \gamma_5 u(p) \end{aligned}$$

$$+ \Delta_\perp^\mu \frac{\tilde{b}}{2m} (\tilde{E} + \tilde{G}_1) + \tilde{h}_\perp^\mu (\tilde{H} + \tilde{G}_2) + \Delta_\perp^\mu \frac{\tilde{h}^+}{P^+} \tilde{G}_3 + \tilde{\Delta}_\perp^\mu \frac{h^+}{P^+} \tilde{G}_4$$

Kinematic factors defined by calculation setup

All twist-3 helicity GPDs: 4 x computational cost compared to PDFs !

For $\vec{Q} = (Q_x, 0, 0)$ the following matrix elements contribute

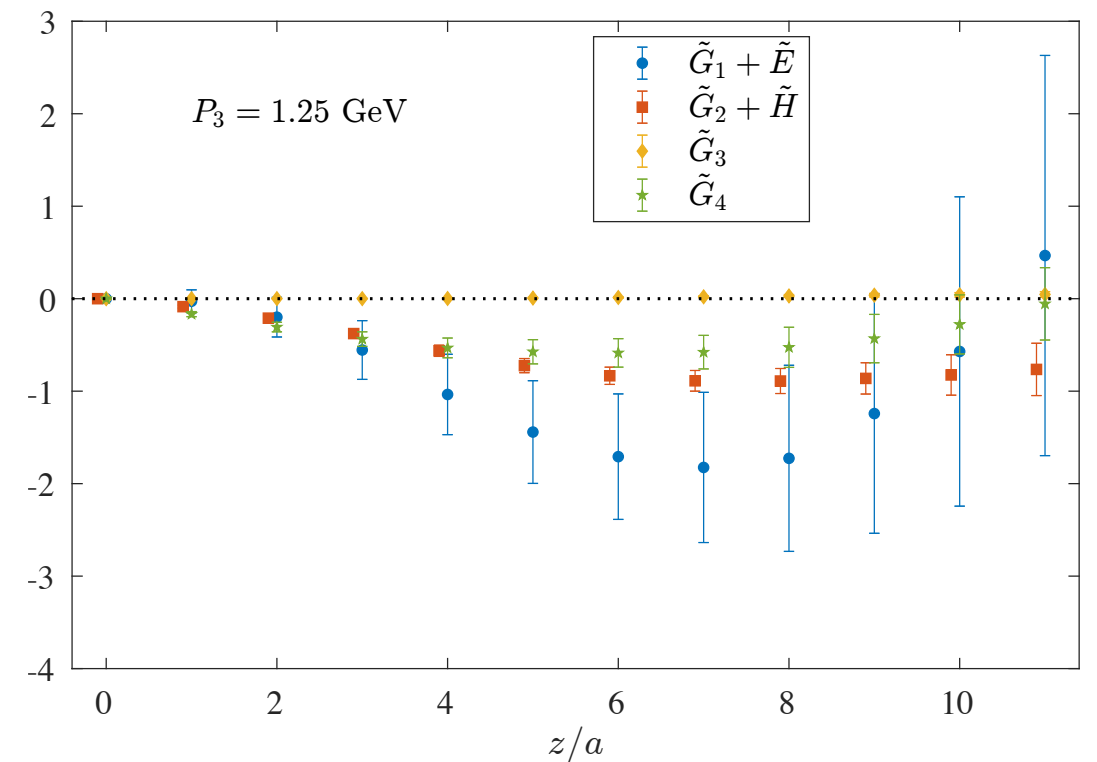
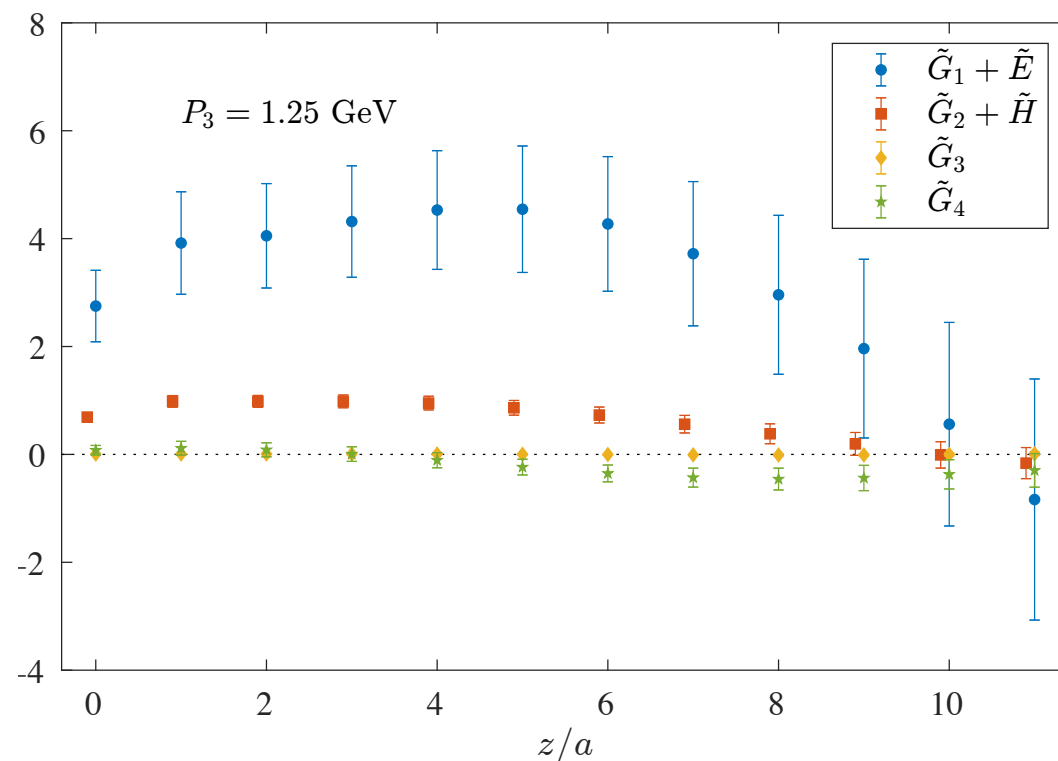
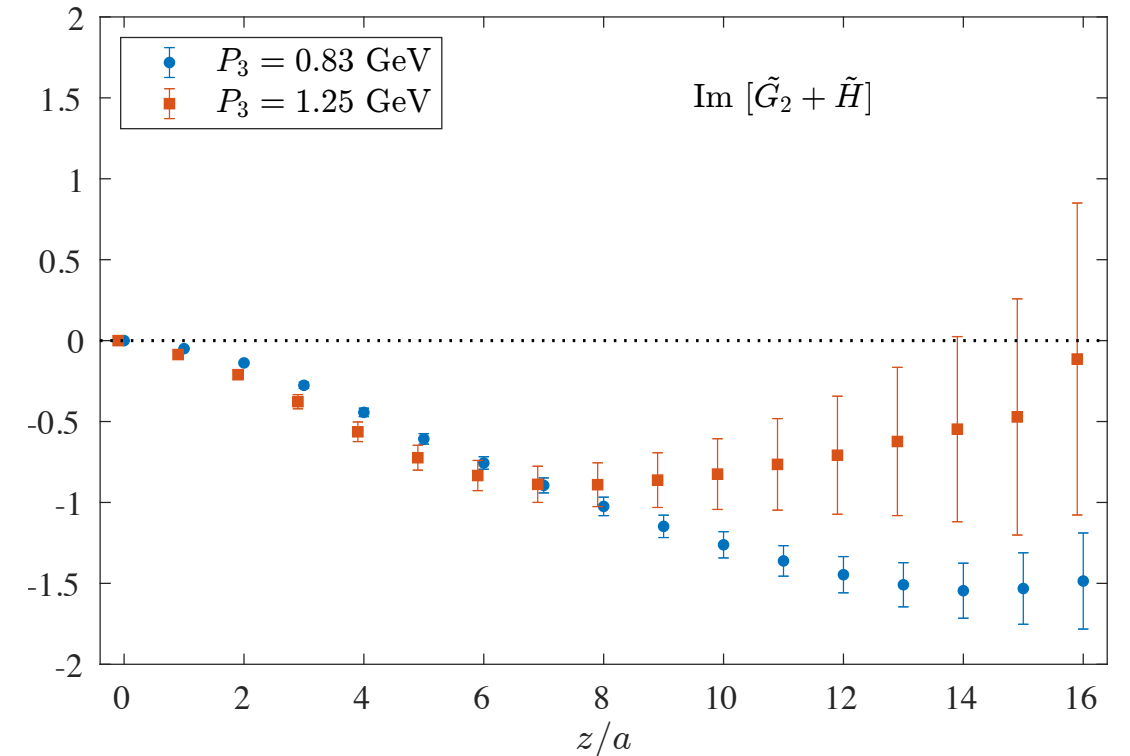
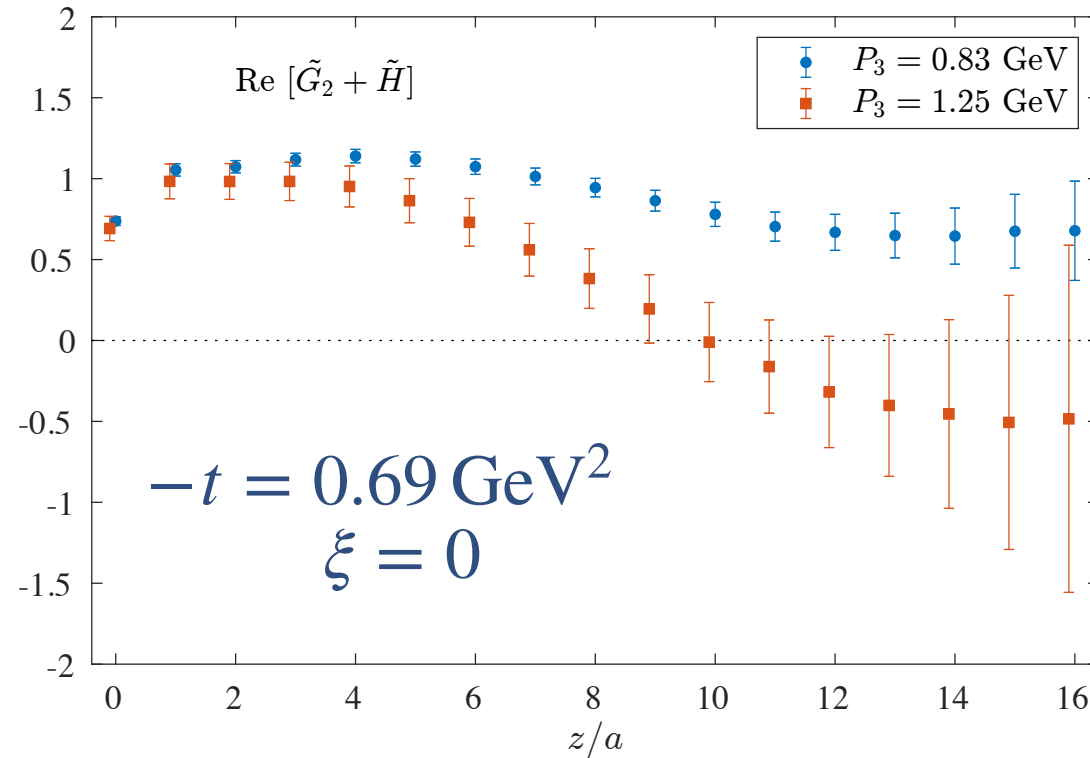
- $\Pi(\gamma^2 \gamma^5, \Gamma_0) : \tilde{H} + \tilde{G}_2, \quad \tilde{G}_4$
- $\Pi(\gamma^2 \gamma^5, \Gamma_2) : \tilde{H} + \tilde{G}_2, \quad \tilde{G}_4$
- $\Pi(\gamma^1 \gamma^5, \Gamma_1) : \tilde{H} + \tilde{G}_2, \quad \tilde{E} + \tilde{G}_1$
- $\Pi(\gamma^1 \gamma^5, \Gamma_3) : \tilde{G}_3$

Parity projectors

$$\Gamma_0 = \frac{1}{4}(1 + \gamma^0)$$

$$\Gamma_i = \frac{1}{4}(1 + \gamma^0) \gamma^5 \gamma^i$$

Matrix elements decomposition

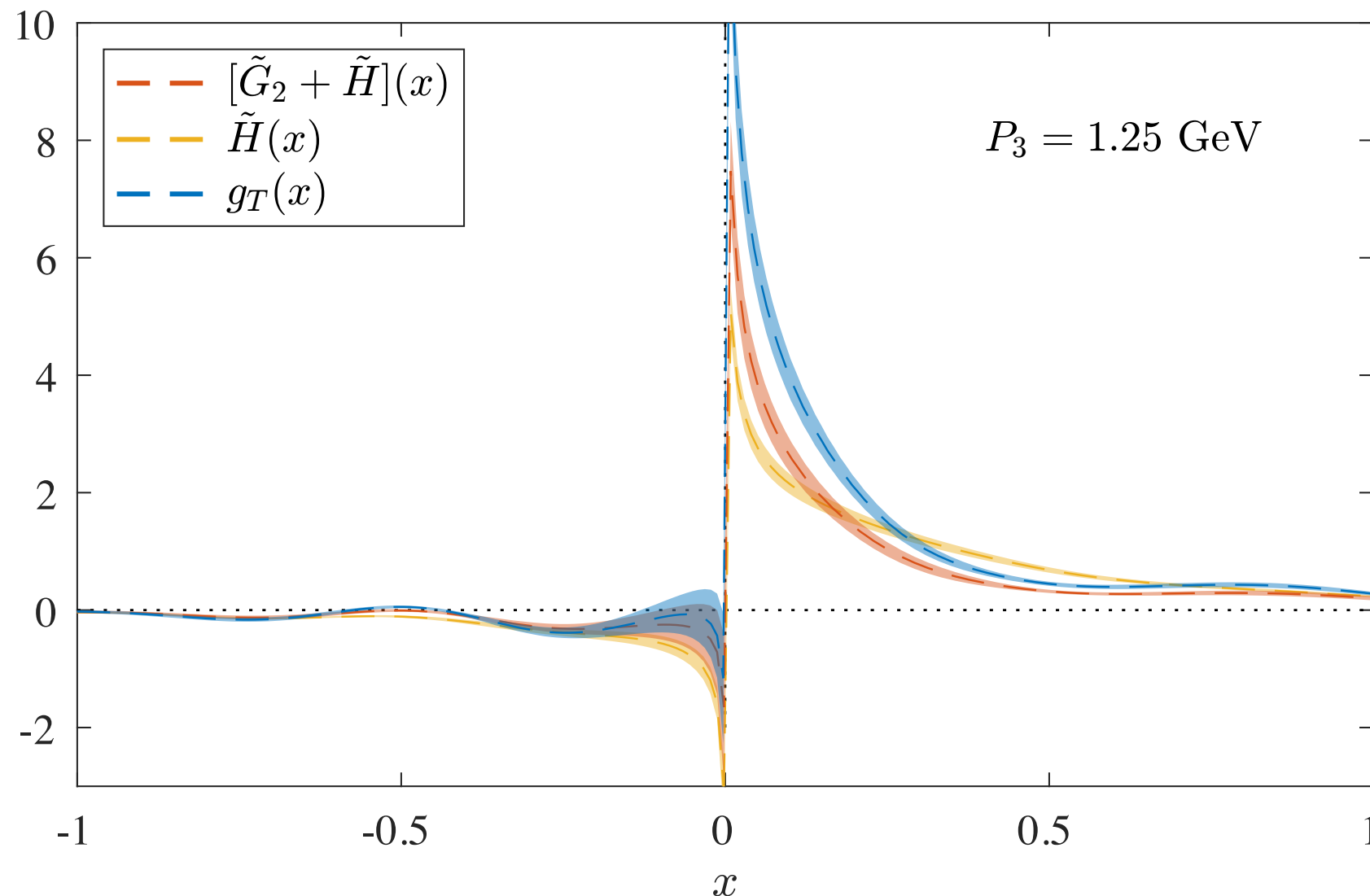


★ $\tilde{E} + \tilde{G}_1$: highest contribution

★ $\tilde{G}_3, \tilde{G}_4$: suppressed

x-dependence of twist-3 GPDs

$$-t = 0.69 \text{ GeV}^2, \quad \xi = 0$$



- ★ $g_T(x)$: dominant distribution
- ★ $\tilde{H} + \tilde{G}_2$ similar in magnitude to $\tilde{H}$
- ★ $\tilde{G}_2$ is expected to be small

Concluding Remarks

Concluding Remarks

- ★ GPDs **multi-dimensionality** poses computational challenges
- ★ Current status: **exploratory and qualitative**
- ★ At twist-3: 2-parton and 3-parton correlations, such as qgq
 - **0th moments of twist-3 GPDs is zero** [D. Kiptily and M. Polyakov, Eur. Phys. J. C37 (2004) 105]
 - **1st moments of twist-3 GPDs have zero qgq contribution**
 - **alternative matching proposed in Braun, Ji, Vladimirov, arXiv: 2103.12105**
- ★ Extraction of twist-3 GPDs is promising with several interesting investigations (WW-approximation, sum rules)
- ★ Nonzero skewness of particular interest: twist-3 GPDs (in models) exhibit discontinuities at $x = \pm \xi$

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Thank you



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