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VAN AMSTERDAM



AZIMUTHAL DECORRELATION IN V+JET PRODUCTION WITH THE WTA AXIS

Rudi Rahn

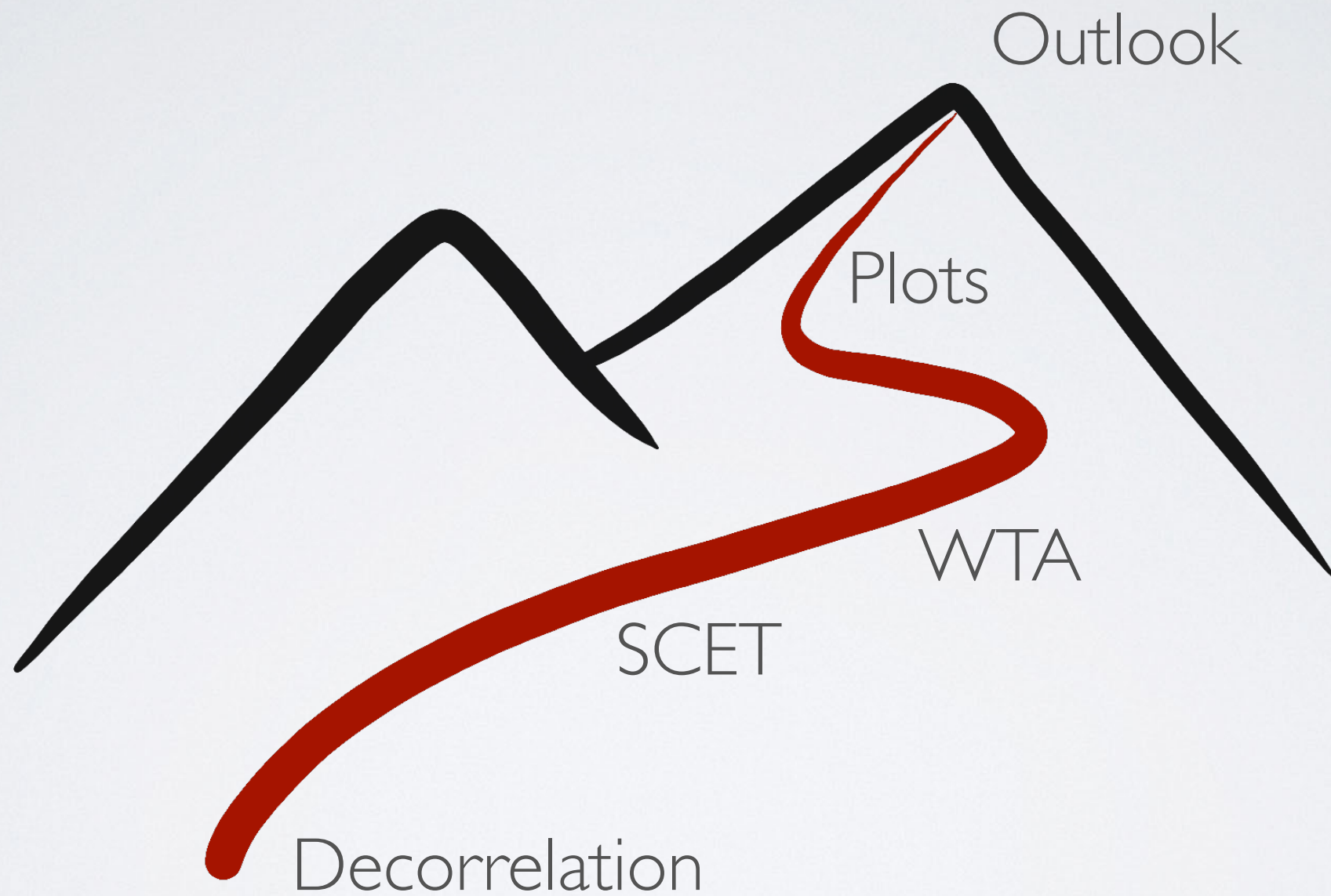
QCD Evolution | 3-05-21,
UCLA

Based on [2005, 12279] and WIP with

Yang-Ting Chien (Stony Brook), Ding Yu Shao (UCLA/Fudan),

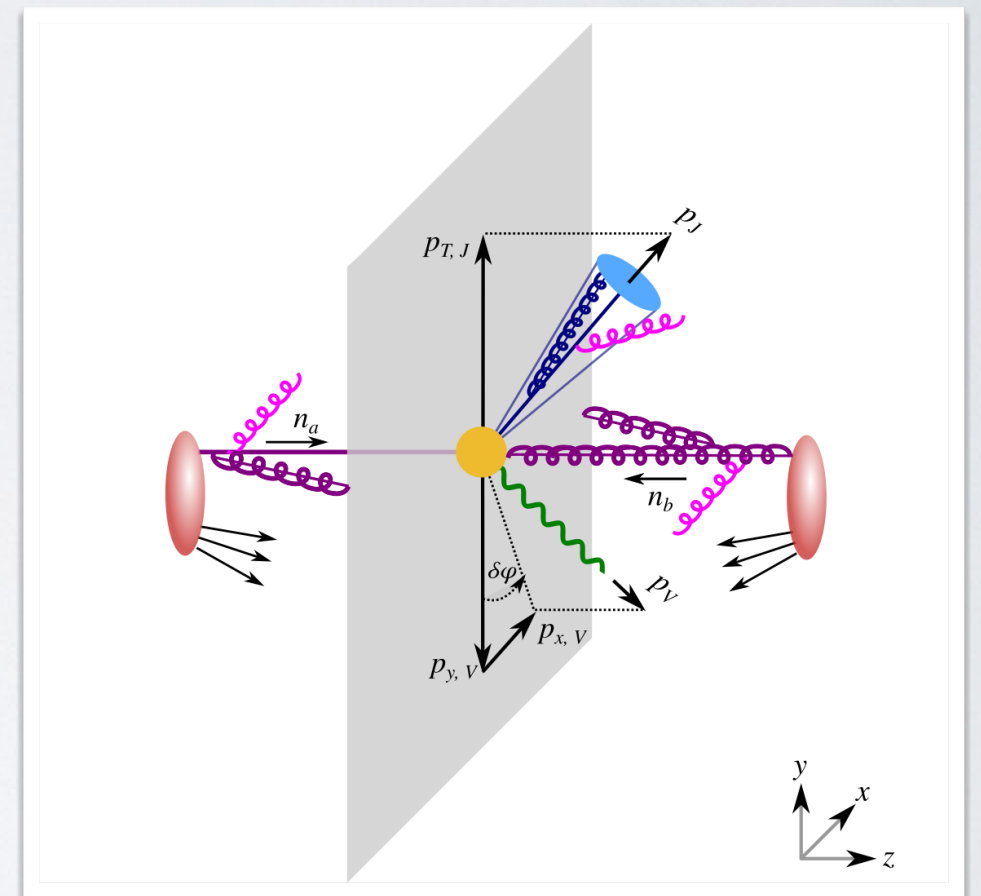
Solange Schrijnder van Velzen, Wouter Waalewijn (Nikhef/UvA), Bin Wu (CERN)

OUTLINE



AZIMUTHAL DECORRELATION I

- V+jet at pp collisions
- Leading partonic: planar
- Soft/Collinear emissions: $\delta\varphi \neq 0$
- Large logarithms $\alpha_s^n \ln^{2n} \delta\varphi$



► Resum using **SCET**

SOFT-COLLINEAR EFFECTIVE THEORY

[Bauer, Fleming, Pirjol, Stewart, Rothstein, '01/'02]

[Beneke, Chapovsky, Diehl, Feldmann, '02]

- Dominant contributions from select phase space regions
- Represent them by dedicated fields
- Power counting of EFT $\Leftrightarrow$ large logarithm argument



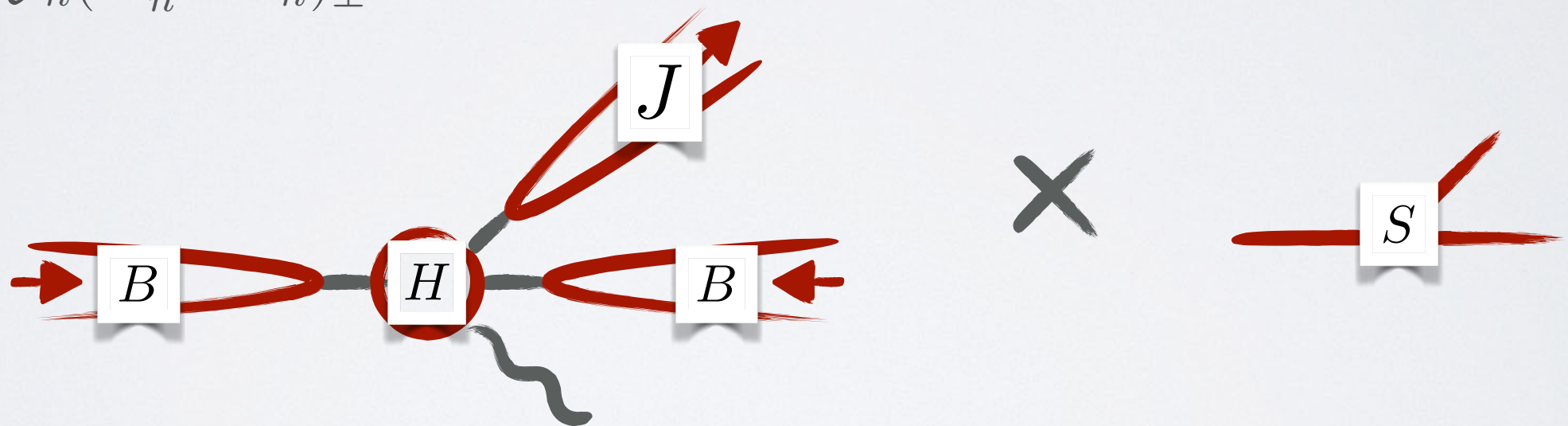
SOFT-COLLINEAR EFFECTIVE THEORY

- Crucial advantage: EFT modes decouple

- Only effective operator joins them

$$\mathcal{O} \sim C_{\pm} \cdot \mathcal{B}_{a,\pm} (\bar{\chi}_J \Gamma \chi_b)_{\pm}$$

$$\mathcal{B}_{n,\pm} \sim \mathcal{Y}_n(W_n^\dagger DW_n)_{\pm}$$



$$d\sigma = H \otimes B_a \otimes B_b \otimes J \otimes S$$

- Resummation via RG equations

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$$\mathcal{O} \sim C_\pm \cdot \mathcal{B}$$

$$)_\pm$$

Only if the observable plays along!

$$d\sigma = H \otimes B_a \otimes B_b \otimes J \otimes S$$

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AZIMUTHAL DECORRELATION II

- Starting point: momentum conservation

$$\vec{p}_{T,a} + \vec{p}_{T,b} + \vec{p}_{T,S} + \vec{p}_{T,c} + \vec{p}_{T,V} = 0$$

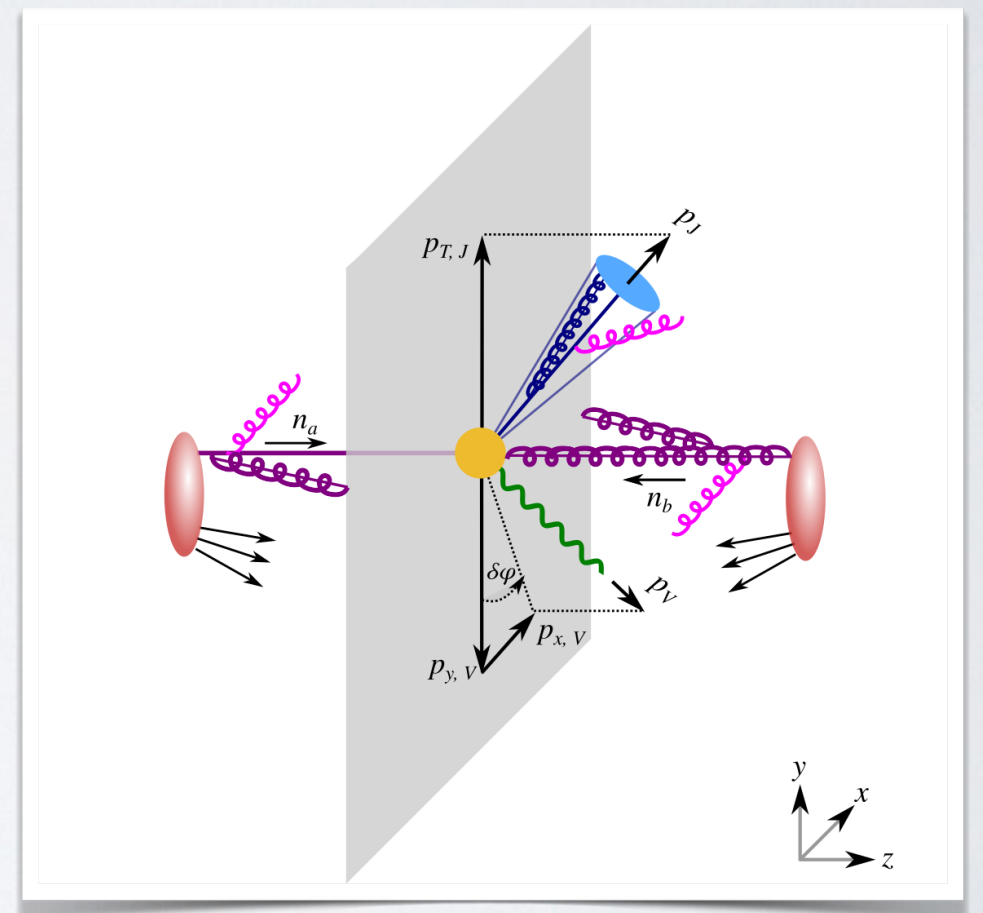
- For WTA axis $\vec{p}_{T,c} \neq \vec{p}_{T,J}$, and rearrange

$$\begin{aligned}\vec{q} &:= \vec{p}_{T,V} + \vec{p}_{T,J} \\ &= \vec{p}_{T,J} - \vec{p}_{T,c} - \vec{p}_{T,a} - \vec{p}_{T,b} - \vec{p}_{T,S}\end{aligned}$$

- For azimuthal decorrelation, pick q_x

$$q_x = -p_{x,c} - p_{x,a} - p_{x,b} - p_{x,S}$$

$$= p_{T,V} \sin \delta\varphi$$



THE WINNER-TAKES-ALL AXIS

[Salam, unpublished]

[Bertolini, Chan, Thaler, '14]

- Recluster a jet, pairwise: pair emissions i, j with $p_{T,i}, \hat{n}_i$ & $p_{T,j}, \hat{n}_j$

$$p_{T,i+j} = p_{T,i} + p_{T,j}$$

$$\hat{n}_{i+j} = \begin{cases} \hat{n}_i & \text{if } p_{T,i} > p_{T,j} \\ \hat{n}_j & \text{if } p_{T,i} < p_{T,j} \end{cases}$$

- Combine with SCET:

$$p_{T,c} \sim Q, \quad p_{T,S} \sim Q\delta\varphi$$

- Direction purely a collinear affair:

$$\max(p_{T,c}, p_{T,S}) = p_{T,c}$$

- Soft subleading in magnitude:

$$p_{T,c} + p_{T,S} \approx p_{T,c}$$

$\Rightarrow$ No soft recoil, no non-global logarithms*

[Larkoski, Neill, Thaler, '14]

WHY SO SIMPLE?

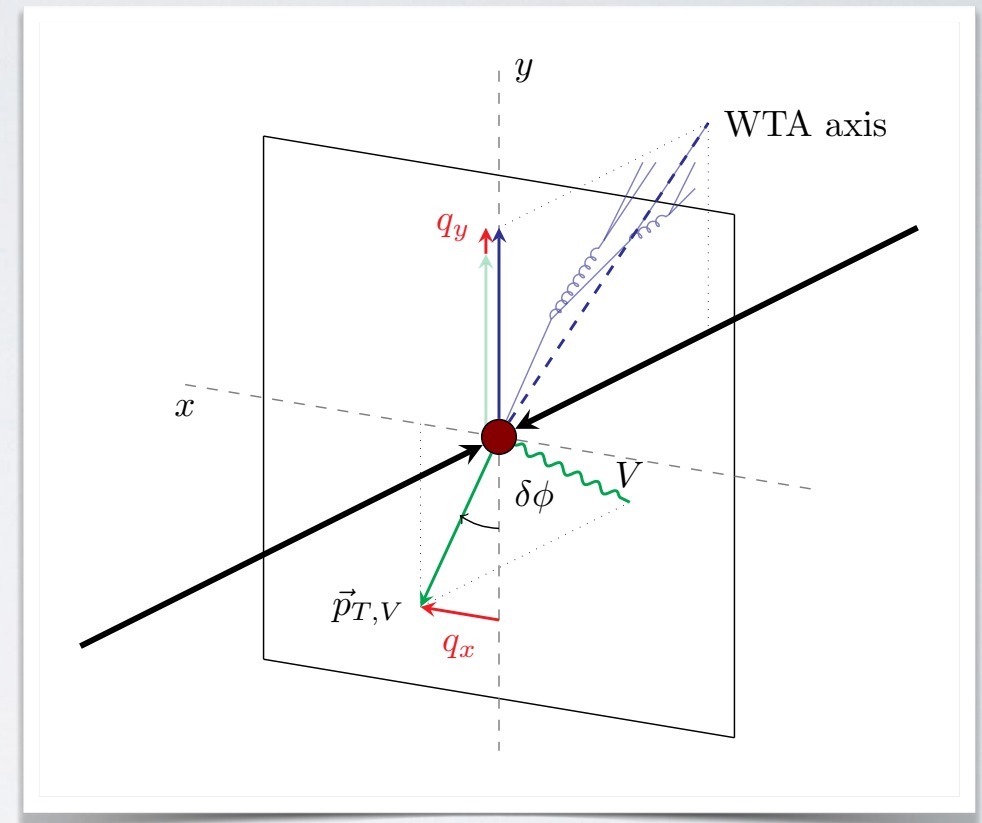
- Compare: *radial* decorrelation

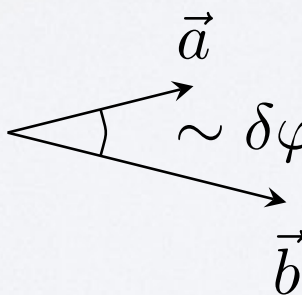
$$q_x = p_{T,V} \sin \delta\varphi$$

$$\sim \delta\varphi$$

$$q_y = p_{T,J} - p_{T,V} \cos \delta\varphi$$

$$\sim (\delta\varphi)^2$$



- Scalar vs. vector sum:  $\Rightarrow (|\vec{a}| + |\vec{b}|) - |(\vec{a} + \vec{b})| \sim \delta\varphi^2$

- Large cancellation:

$$p_{T,c} + p_{T,s} \approx p_{T,c}$$

$\Rightarrow$ Sensitive to in-jet soft $\Rightarrow$ NGL

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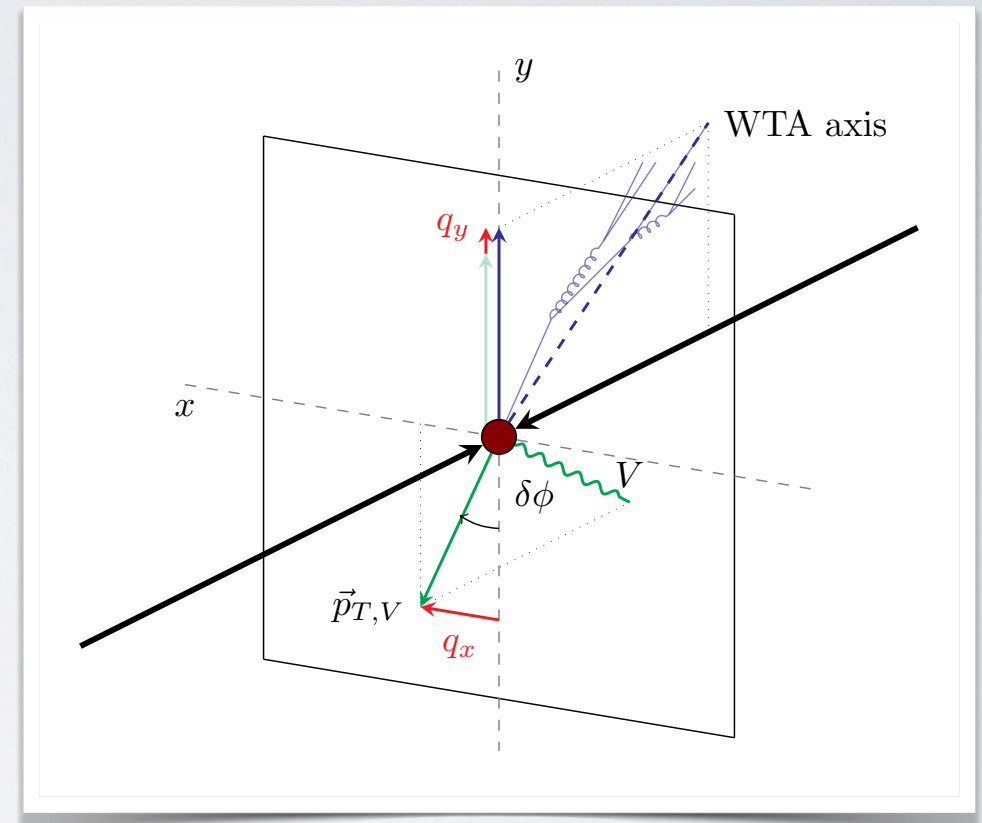
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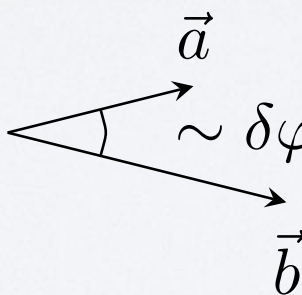
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$$\cancel{p_{T,c}} + p_{T,S} = p_{T,S}$$

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CALCULATION

$$\frac{d\sigma}{dp_{x,V} dp_{T,J} dy_V d\eta_J} = \int \frac{db_x}{2\pi} e^{ip_{x,V} b_x} \sum_{i,j,k} B_i(x_a, b_x) B_j(x_b, b_x) S_{ijk}(b_x, \eta_J) \\ \times \mathcal{H}_{ij \rightarrow V k}(p_{T,V}, y_V - \eta_J) \mathcal{J}_k(b_x) \left[1 + \mathcal{O}\left(\frac{p_{x,V}^2}{p_{T,V}^2}\right) \right]$$

- NNLL resummation

[Arnold, Reno, '89]

[Becher, Lorentzen, Schwartz, '12]

[Moch, Vermaseren, Vogt, '04/'05]

[Becher, Neubert, '09]

[Gehrmann, Luebbert, Yang, '14]

[Echevarria, Scimemi, Vladimirov, '16]

[Luebbert, Oredsson, Stahlhofen, '16]

- 3-loop Γ_{cusp} , 2-loop γ_i , 1-loop finite H,J,S,B

- Most known, only minor recalculations:

$$S_{ij}^{(1)}(b_x, \eta_J, \mu, \nu) = \\ - \sum_{i < j} \mathbf{T}_i \cdot \mathbf{T}_j S^{(1)}\left(b_x, \mu, \nu \sqrt{n_i \cdot n_j / 2}\right)$$

- S at NLO via boost (non-)invariance:

- TMD jet function with η -regulator and large jet radius

[Gutierrez-Reyes, Scimemi,
Waalewijn, Zoppi, '18/'19]

[Chiu, Jain, Neill, Rothstein, '12]

LINEAR POLARISATION

- Two Lorentz structures in gluon beam function

$$B_g^{\mu\nu}(\vec{b}_\perp, \mu, \nu) = \frac{-g_\perp^{\mu\nu}}{d-2} B_g^{(1)}(\vec{b}_\perp, \mu, \nu) + \left(\frac{g_\perp^{\mu\nu}}{d-2} + \frac{b_\perp^\mu b_\perp^\nu}{b_T^2} \right) B_g^{(2)}(\vec{b}_\perp, \mu, \nu)$$

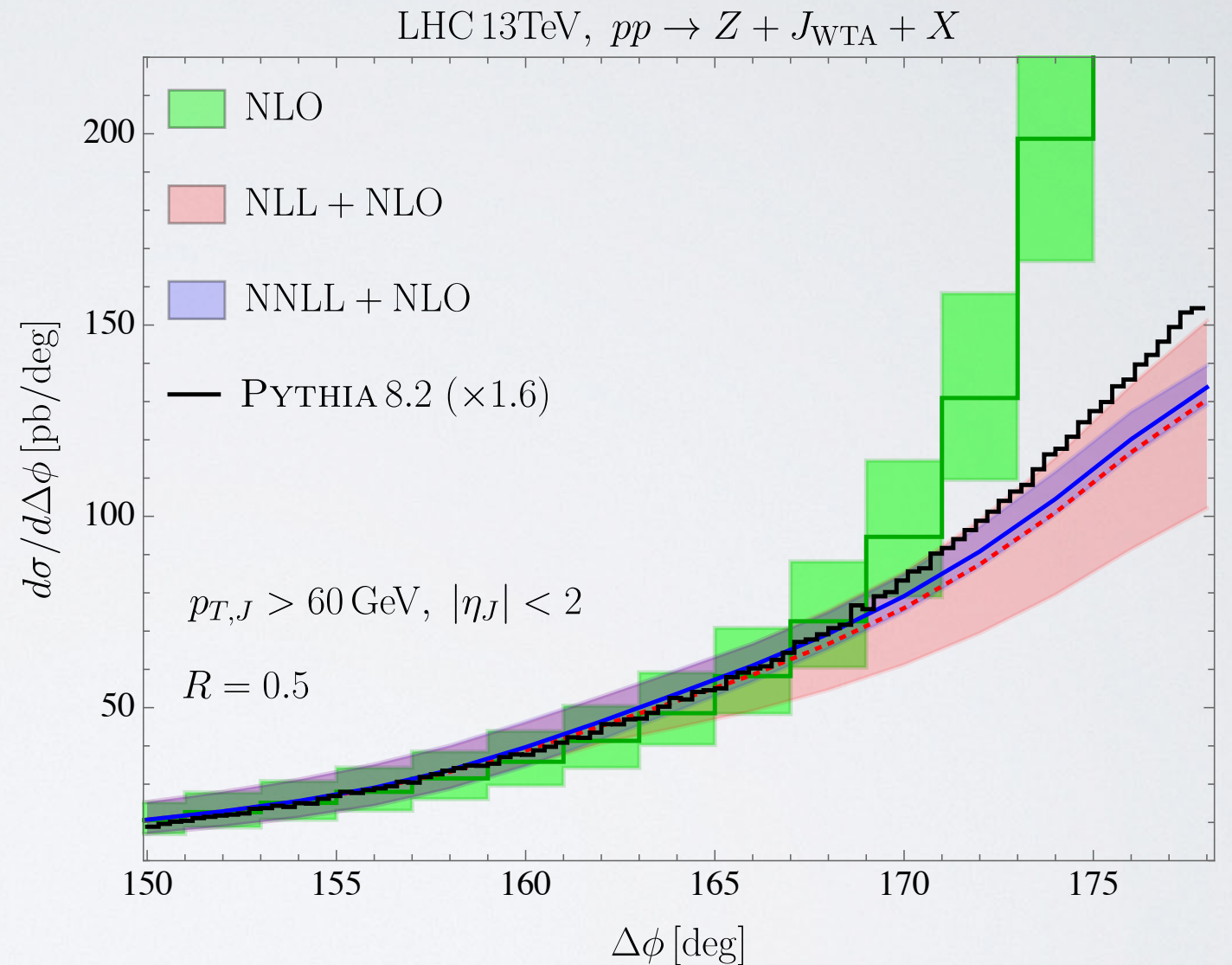
[Catani, Grazzini, '11]

- $B_g^{(2)}$ contributes, with $\mathcal{H}_{qg \rightarrow Vq}^{(L)} \sim \frac{\hat{u} m_V^2}{\hat{s} \hat{t}}$
- Somewhat surprising, only one gluon
- Identically for gluon jet!

$$\mathcal{J}_g^{L(1)}(\vec{b}_\perp, \mu, \nu) = -\frac{1}{3} C_A + \frac{2}{3} T_F n_f$$

PLOTS I

- Using CT14nlo
- NLO from MCFM
- No difference Pythia hadron/parton
- Matching via transition function



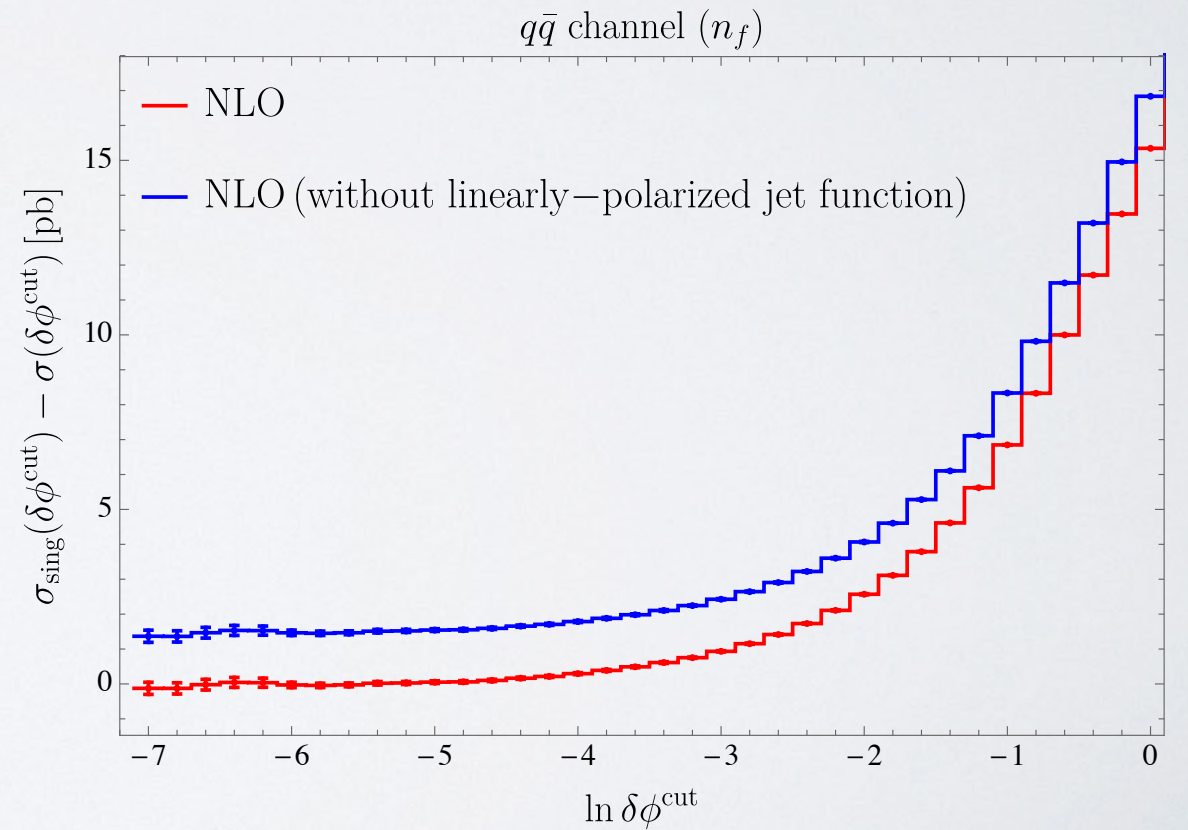
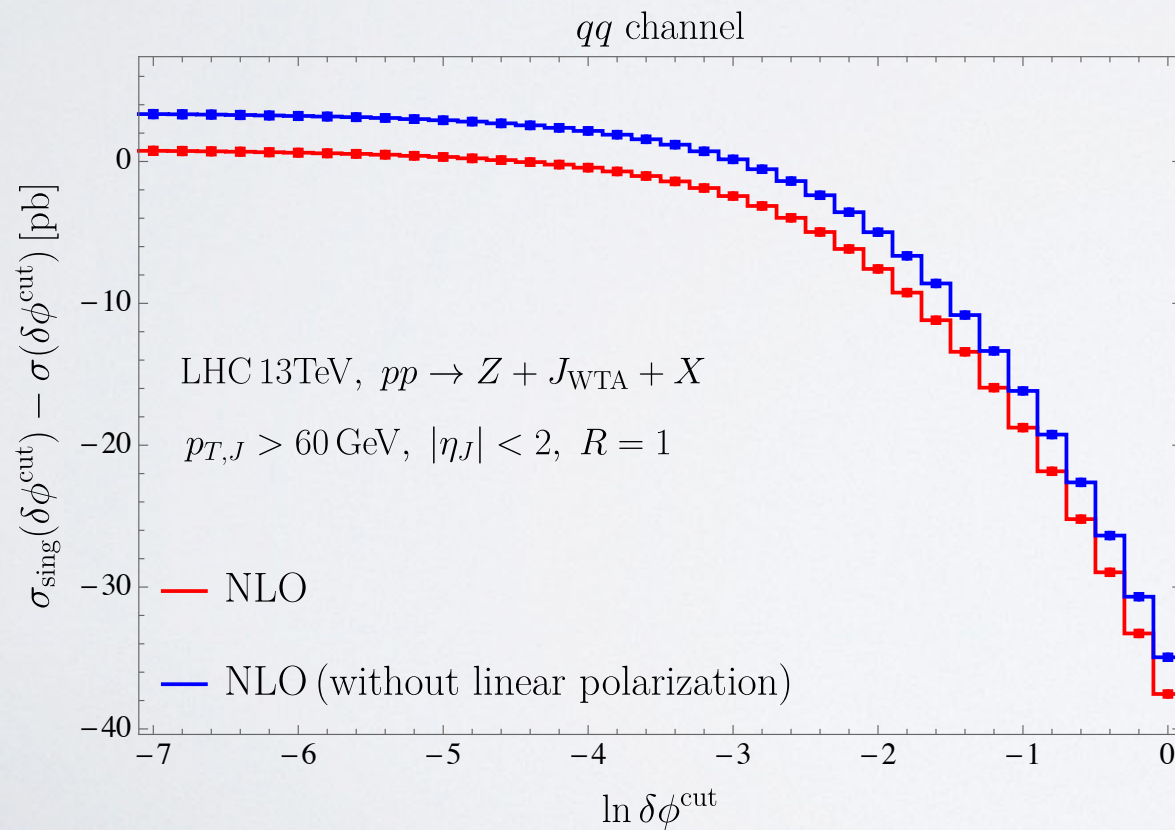
$$\mu_H = \sqrt{p_{T,V}^2 + m_V^2}, \nu_S = \mu_B = 2e^{-\gamma_e}/|b_x|, \nu_{B_{a,b}} = x_{a,b}\sqrt{s}, \nu_J = 2p_{T,J} \cosh \eta_J$$



Use b^* prescription [Collins, Soper, Sterman, '85]

PLOTS II

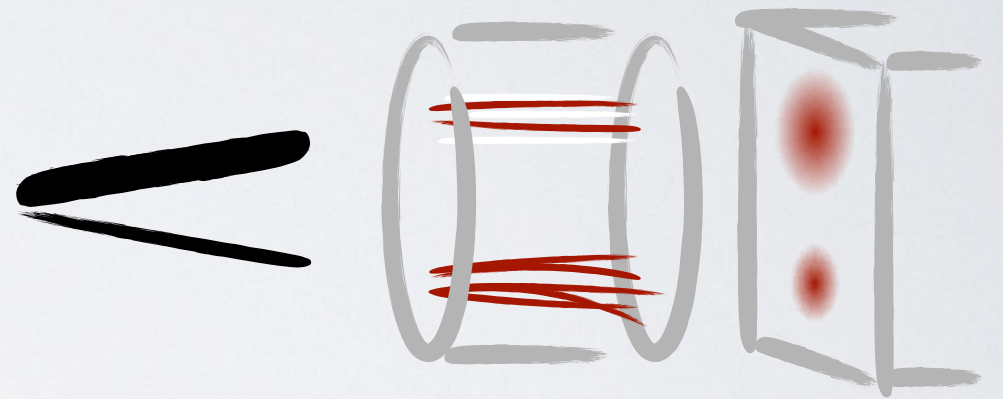
- Linearly polarised beam and jet functions needed
- MCFM vs. Our singular:



TRACKS

- Angular resolution of calorimetry could be a problem

⇒ Use charged particle tracks

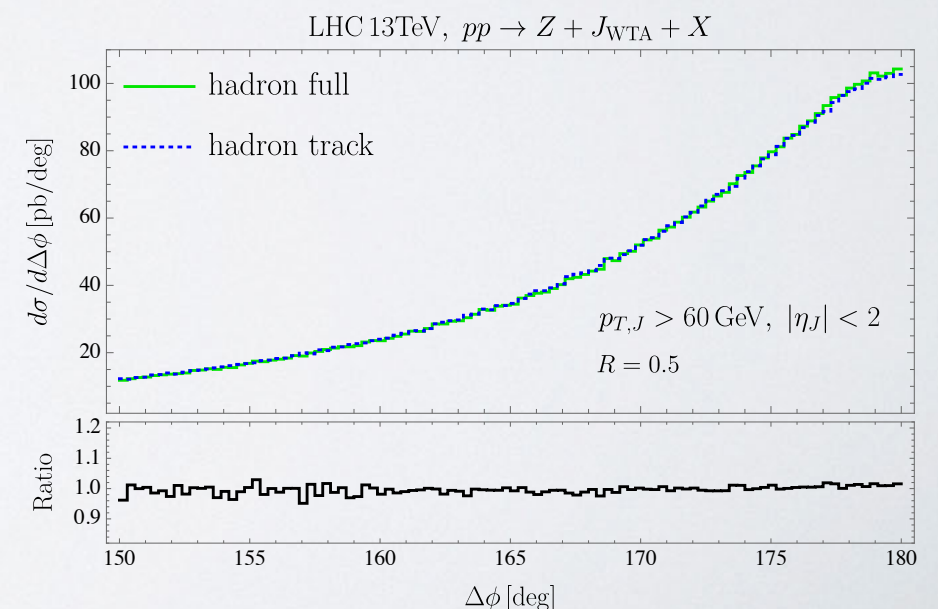


- Only affects Jet function finite term

$$\begin{aligned} \bar{\mathcal{J}}_q^{(1)} = & \mathcal{J}_q^{(1)} + 4C_F \int_0^1 dx \frac{1+x^2}{1-x} \ln \frac{x}{1-x} \int_0^1 dz_1 T_q(z_1, \mu) \\ & \times \int_0^1 dz_2 T_g(z_2, \mu) [\theta(z_1 x - z_2(1-x)) - \theta(x - \frac{1}{2})] \end{aligned}$$

- Track function T ⇒ after the break

[Chang, Procura, Thaler, Waalewijn, '13]



OUTLOOK / CONCLUSION

- Resummation for azimuthal decorrelation in $V+\text{jet}$
 - Surprisingly simple (WTA+SCET)
 - Surprisingly rich (linear polarisation)
 - Surprisingly flexible (Tracks)
- Still some work to do (Glauber contribution,...)