

Physics with Heavy Ions and Exotic Hadrons at LHCb

Matt Durham

Los Alamos National Laboratory

BNL Nuclear Physics Seminar 7 Jan 2020

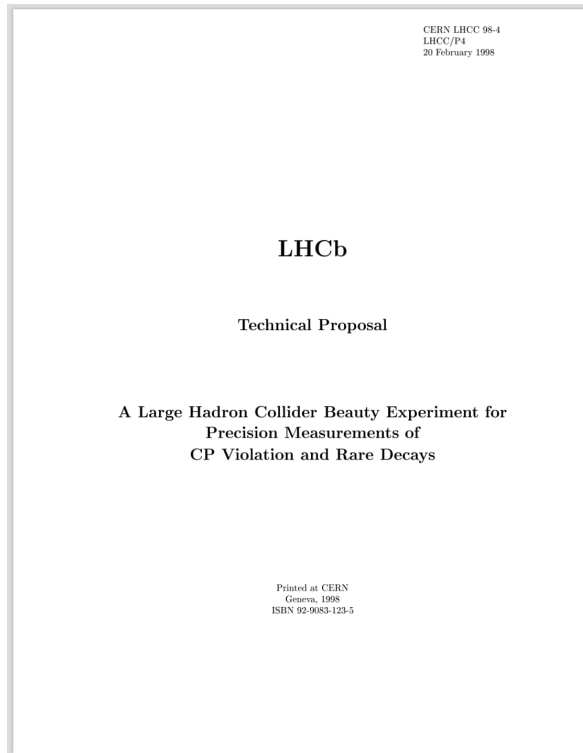


Outline

- LHCb Physics program
- Apparatus
- Heavy Ions at LHCb
 - Collider Mode
 - Fixed Target Mode (SMOG)
- Exotics at LHCb
- Upgrades

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1998 Technical Proposal:
“Beauty Experiment for
CP Violation and Rare Decays”

PHYSICAL REVIEW LETTERS **122**, 211803 (2019)

Editors' Suggestion

Featured in Physics

Observation of *CP* Violation in Charm Decays

R. Aaij *et al.**
(LHCb Collaboration)

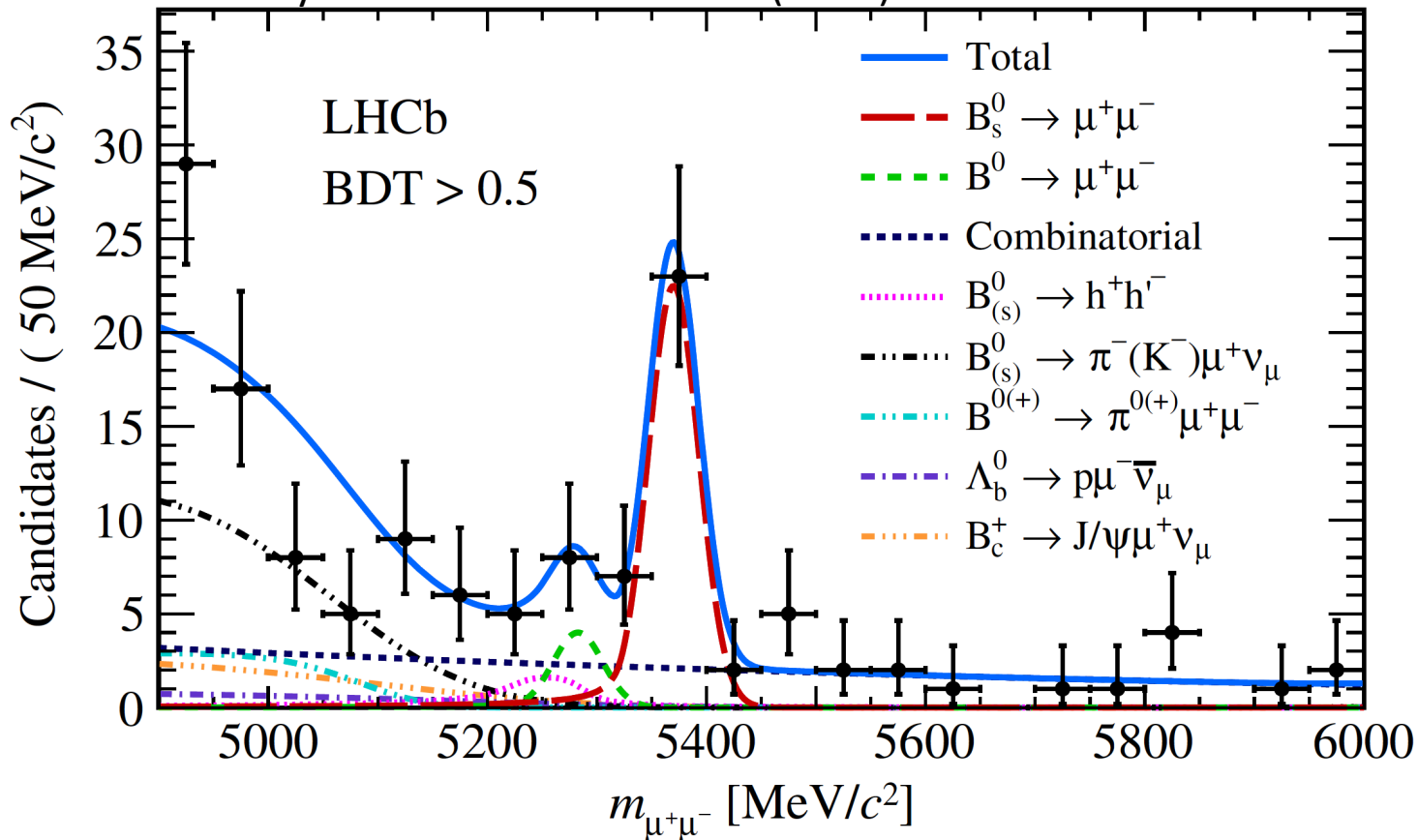
 (Received 21 March 2019; revised manuscript received 2 May 2019; published 29 May 2019)

A search for charge-parity (*CP*) violation in $D^0 \rightarrow K^- K^+$ and $D^0 \rightarrow \pi^- \pi^+$ decays is reported, using pp collision data corresponding to an integrated luminosity of 5.9 fb^{-1} collected at a center-of-mass energy of 13 TeV with the LHCb detector. The flavor of the charm meson is inferred from the charge of the pion in $D^*(2010)^+ \rightarrow D^0 \pi^+$ decays or from the charge of the muon in $\bar{B} \rightarrow D^0 \mu^- \bar{\nu}_\mu X$ decays. The difference between the *CP* asymmetries in $D^0 \rightarrow K^- K^+$ and $D^0 \rightarrow \pi^- \pi^+$ decays is measured to be $\Delta A_{CP} = [-18.2 \pm 3.2(\text{stat}) \pm 0.9(\text{syst})] \times 10^{-4}$ for π -tagged and $\Delta A_{CP} = [-9 \pm 8(\text{stat}) \pm 5(\text{syst})] \times 10^{-4}$ for μ -tagged D^0 mesons. Combining these with previous LHCb results leads to $\Delta A_{CP} = (-15.4 \pm 2.9) \times 10^{-4}$, where the uncertainty includes both statistical and systematic contributions. The measured value differs from zero by more than 5 standard deviations. This is the first observation of *CP* violation in the decay of charm hadrons.

- CKM Mechanism and CP violation
 - First observation of CP violation in the charm sector

LHCb Physics Program

Phys. Rev. Lett. 118 191801 (2017)



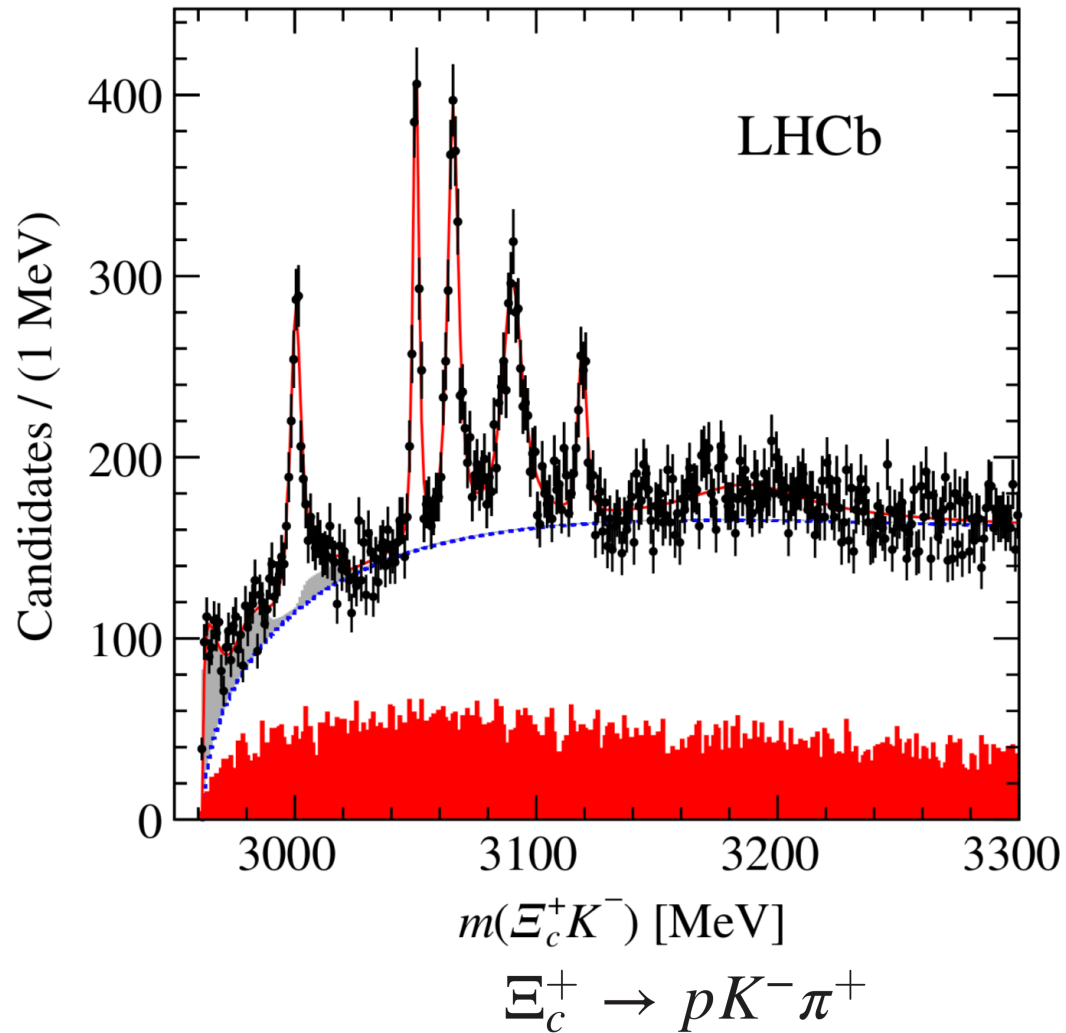
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Rare decays

- $B_s \rightarrow \mu\mu, b \rightarrow s \ell^+\ell^-$

$$\mathcal{B}(B_s^0 \rightarrow \mu^+\mu^-) = (3.0 \pm 0.6_{-0.2}^{+0.3}) \times 10^{-9}$$

LHCb Physics Program

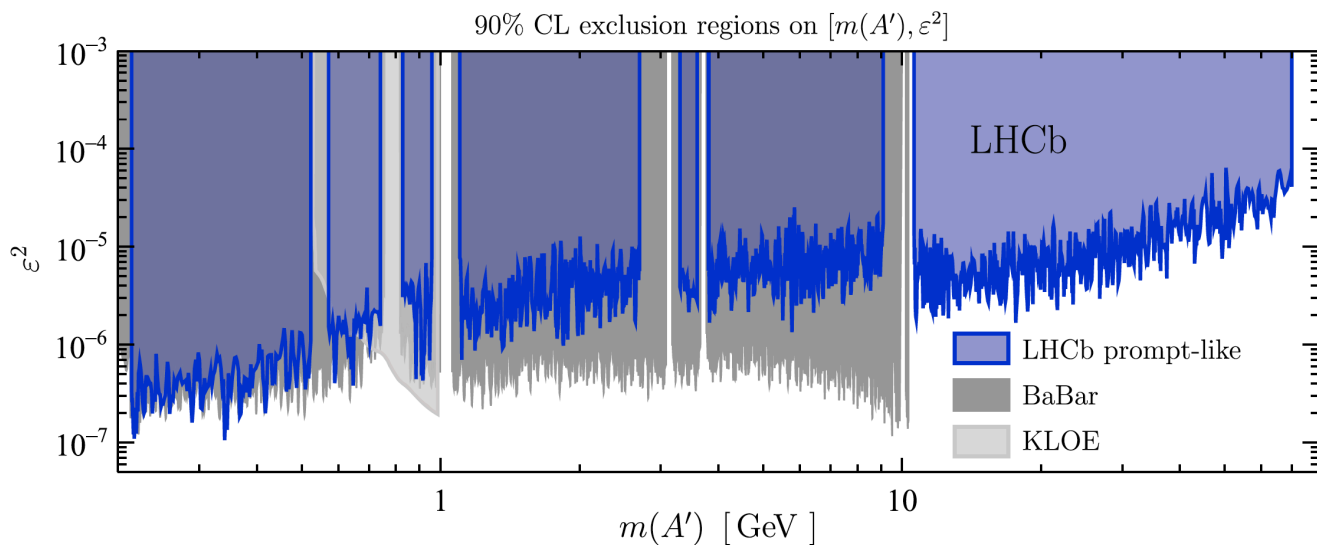
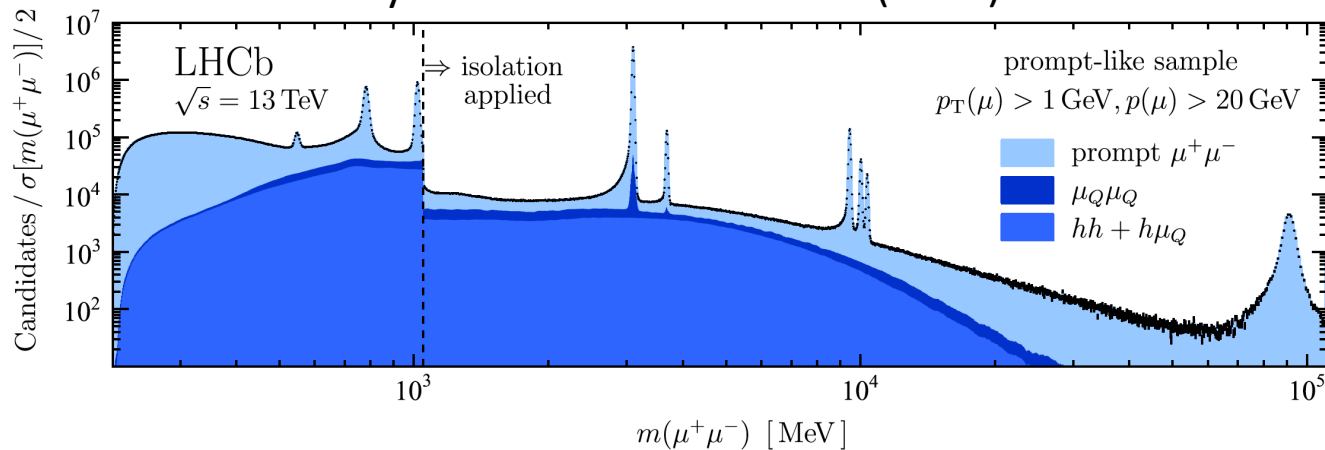


- CKM Mechanism and CP violation
 - First observation of CP violation in the charm sector
- Rare decays
 - $B_s \rightarrow \mu\mu, b \rightarrow s \ell^+ \ell^-$
- Spectroscopy
 - Observation of 5 new Ω_c baryons

Phys. Rev. Lett. 118 182001 (2018)

LHCb Physics Program

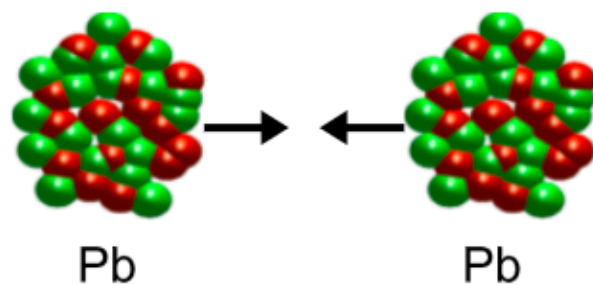
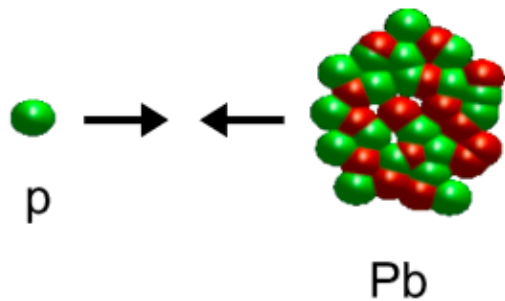
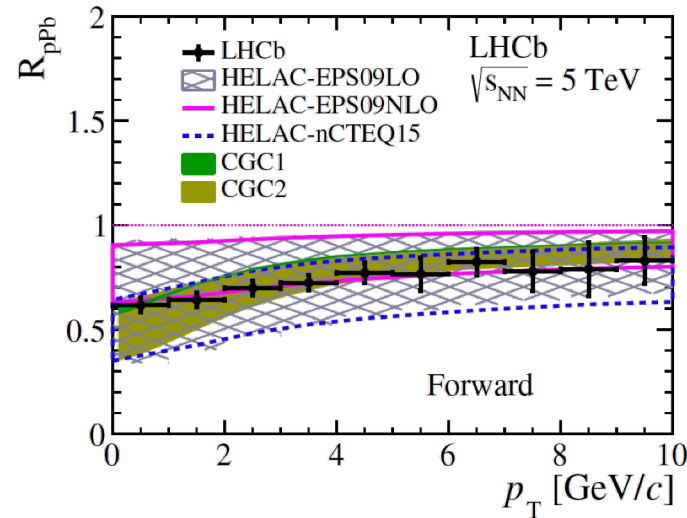
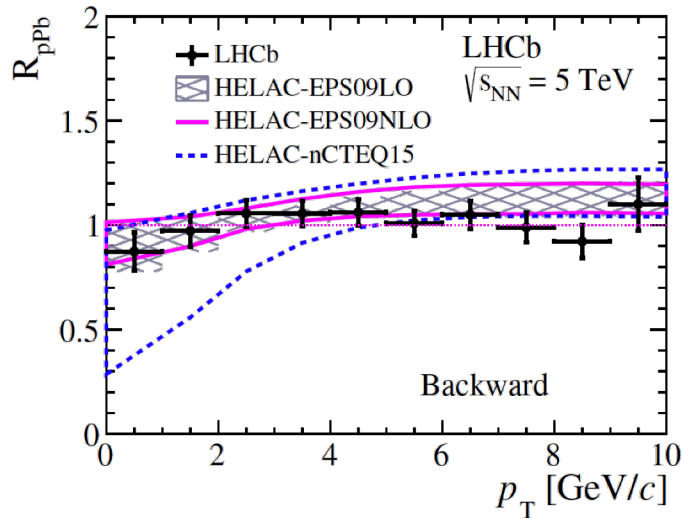
Phys. Rev. Lett. 120 061801 (2018)



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 - First observation of CP violation in the charm sector
- Rare decays
 - $B_s \rightarrow \mu\mu$, $b \rightarrow s \ell^+ \ell^-$
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 - Observation of 5 new Ω_c baryons
- EW, QCD, direct searches for BSM particles
 - Limits on $A' \rightarrow \mu\mu$ decays

LHCb Physics Program

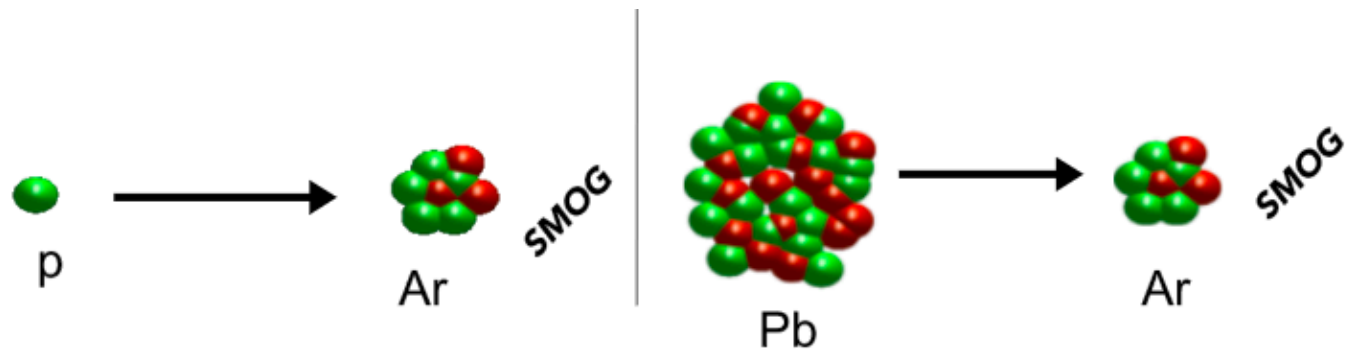
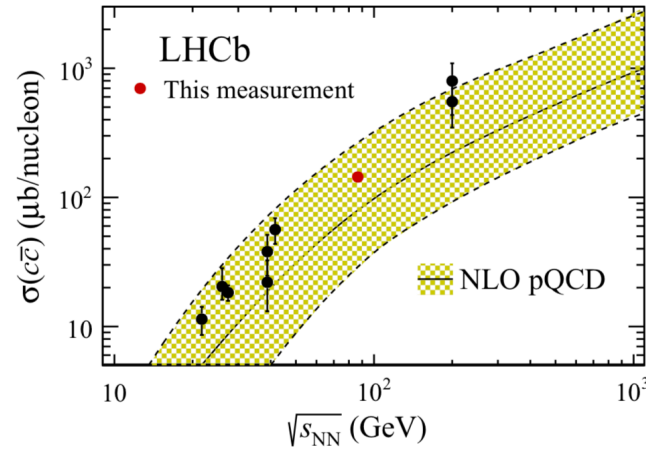
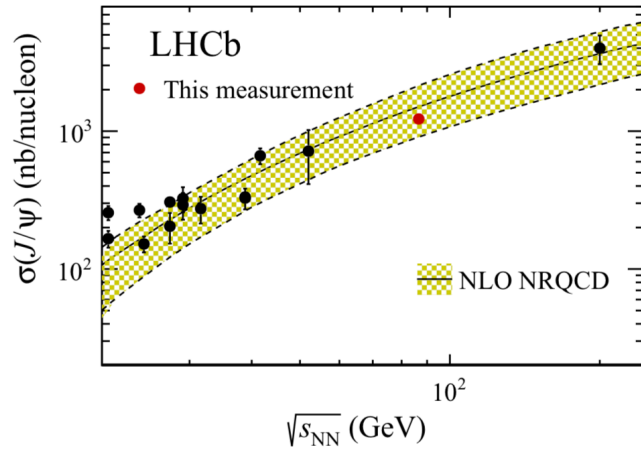
JHEP10 (2017) 090



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- Heavy ion collisions
 - Heavy flavor in pPb, PbPb collisions

LHCb Physics Program

Phys. Rev. Lett. 122, 132002 (2019)



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 - Heavy flavor in pPb, PbPb collisions
- Fixed target collisions
 - pHe, pNe, pAr, PbNe, PbAr, etc accessible by injecting gas into beampipe

Huge evolution in physics program:

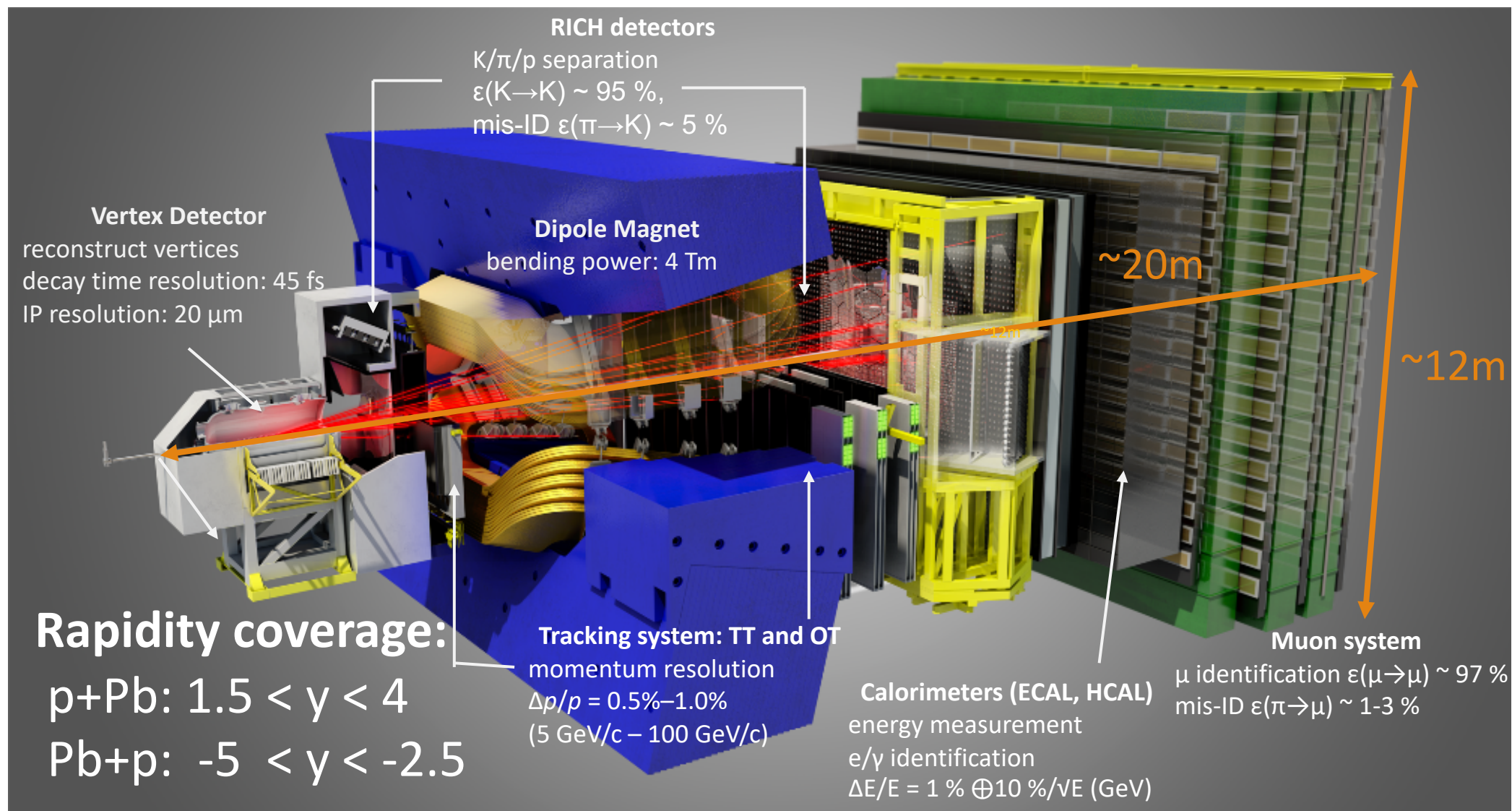
LHCb is a *general purpose* detector covering forward rapidity

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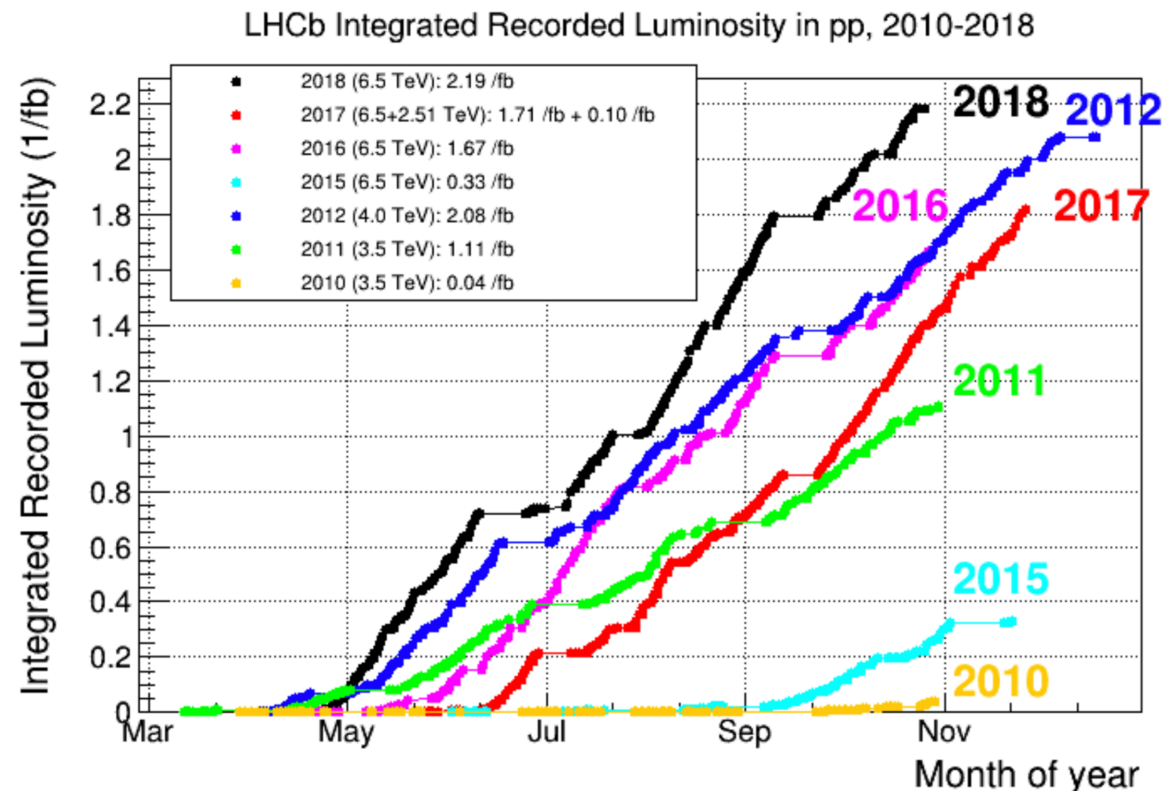
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The LHCb Detector



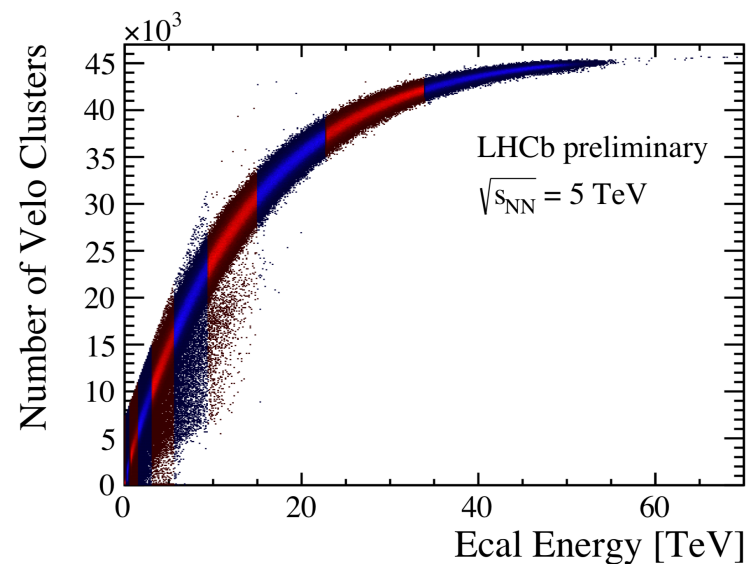
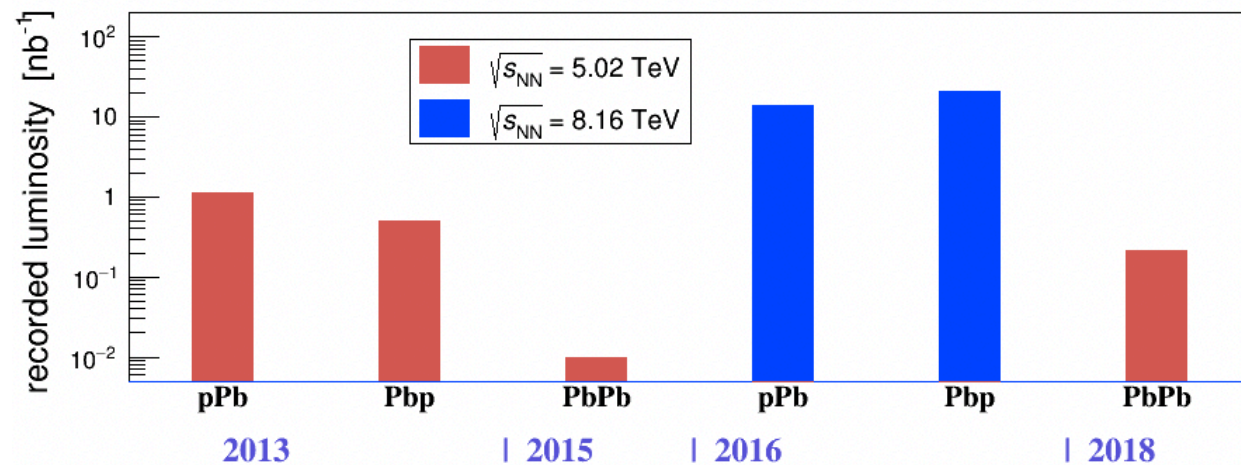
Tracking detector granularity designed for *pp* collisions is not optimal for measurements in central PbPb collisions → upgrade ongoing

Data on Tape



First heavy ion data recorded: 2013 pPb

Tracking works well in full pPb centrality



...limited in PbPb to >60% centrality

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The nuclear PDF

Generic cross section for heavy quark production:

$$d\sigma(Q^2, \sqrt{s})_{pA \rightarrow a+X} = \sum_{i,j=q,\bar{q},g} f_i^P(x_1, Q^2) \otimes A f_i^A(x_2, Q^2) \otimes d\hat{\sigma}(Q^2, x_1, x_2)_{i,j \rightarrow a+X}$$

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Measurable
at experiments

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Measurable at experiments
Calculable by pQCD

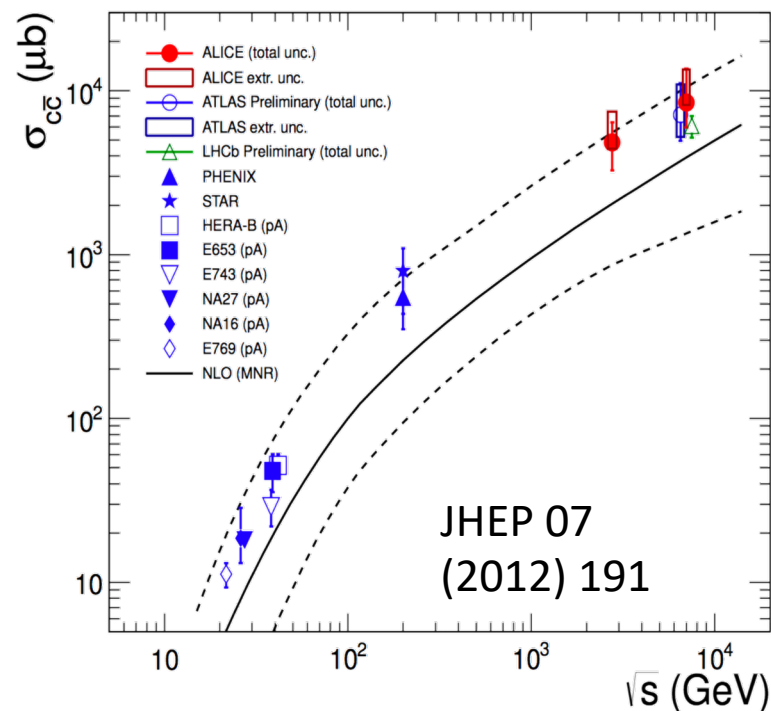
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**Measurable
at experiments**

**Well constrained
HERA and other data**

Calculable by pQCD

The nuclear PDF

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Measurable
at experiments

Well constrained
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Calculable by pQCD

Due to incredible effort, proton PDF is
reasonably well known

A sample of some recent work:

NNPDF3.1: EPJ C77 663 (2017)

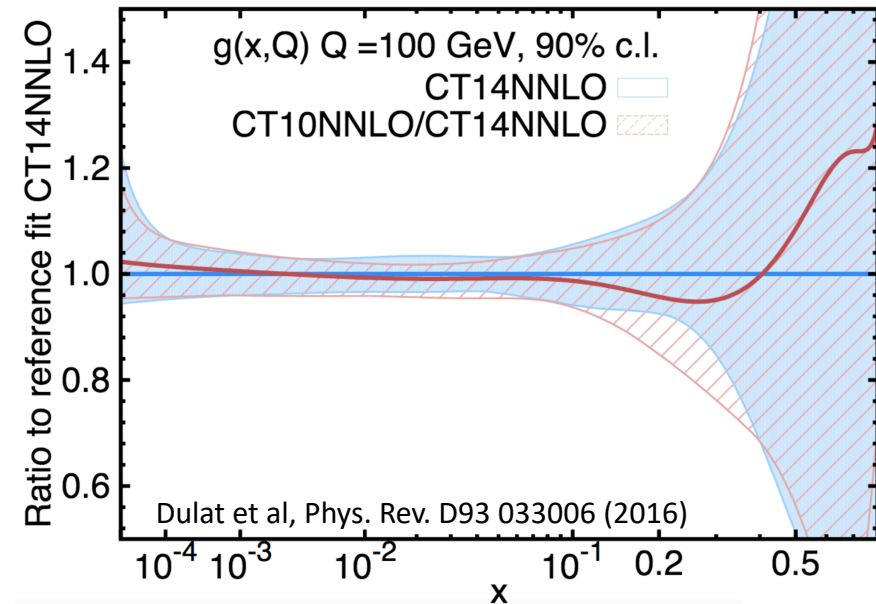
CT14: Phys. Rev. D93 033006 (2016)

MMHT 2014: EPJ C75 204 (2015)

CJ15: Phys. Rev. D93, 114017 (2016)

ABMP16: Phys. Rev. D96, 014011 (2017)

Boughezal et al JHEP (2017) 130



The nuclear PDF

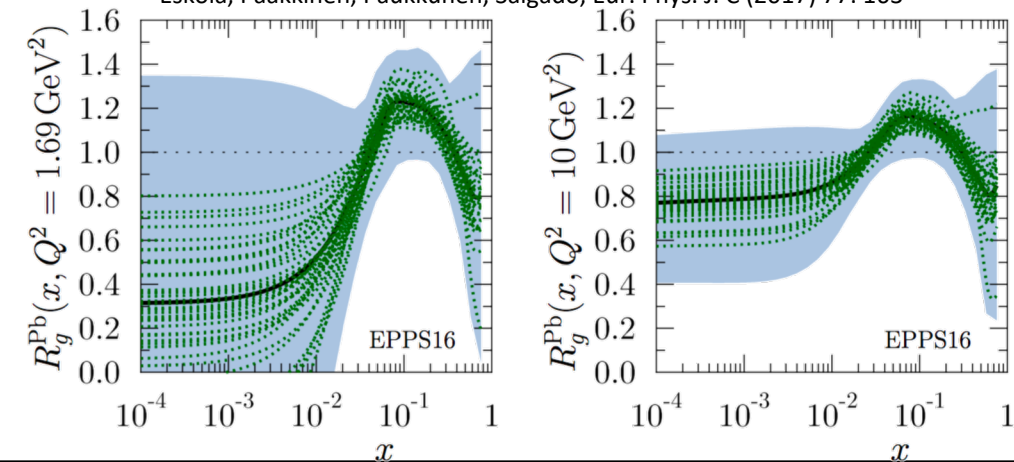
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**Measurable
at experiments**
**Well constrained
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Calculable by pQCD

**Despite incredible effort, nuclear PDF is
not well constrained, esp gluons at low x**

Eskola, Paakkinen, Paukkunen, Salgado, Eur. Phys. J. C (2017) 77: 163



The nuclear PDF

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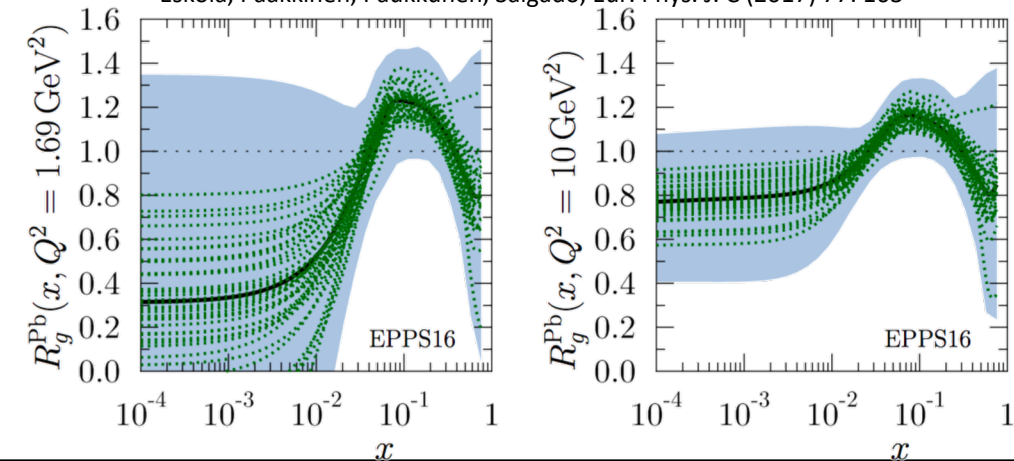
**Measurable
at experiments**
**Well constrained
HERA and other data**
Calculable by pQCD

Despite incredible effort, nuclear PDF is not well constrained, esp gluons at low x

Solution: constrain fits with data at low x with probes that are sensitive to gluon distribution

-> Heavy quarks at forward rapidity

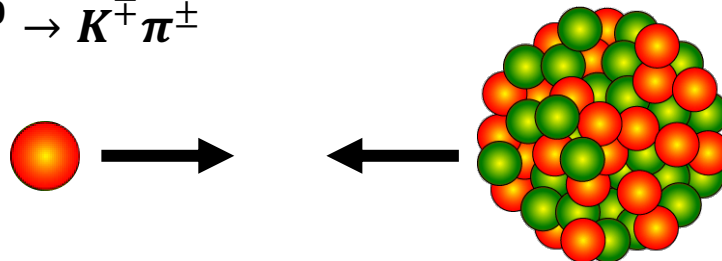
Eskola, Paakkinen, Paukkunen, Salgado, Eur. Phys. J. C (2017) 77: 163



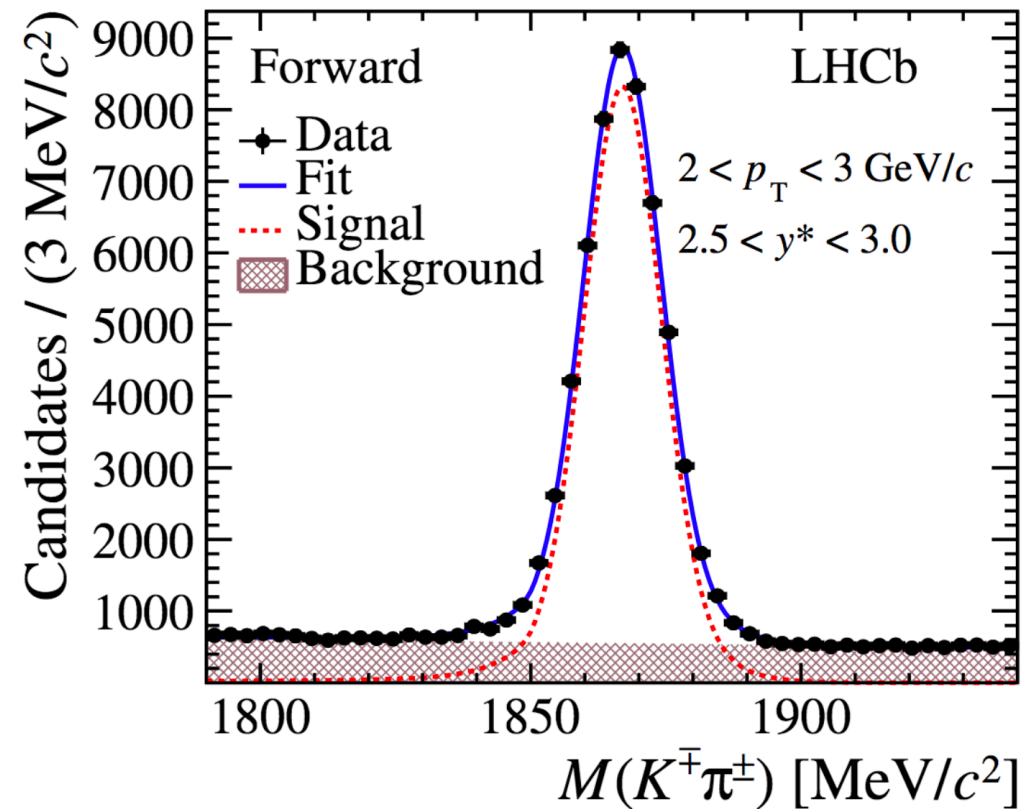
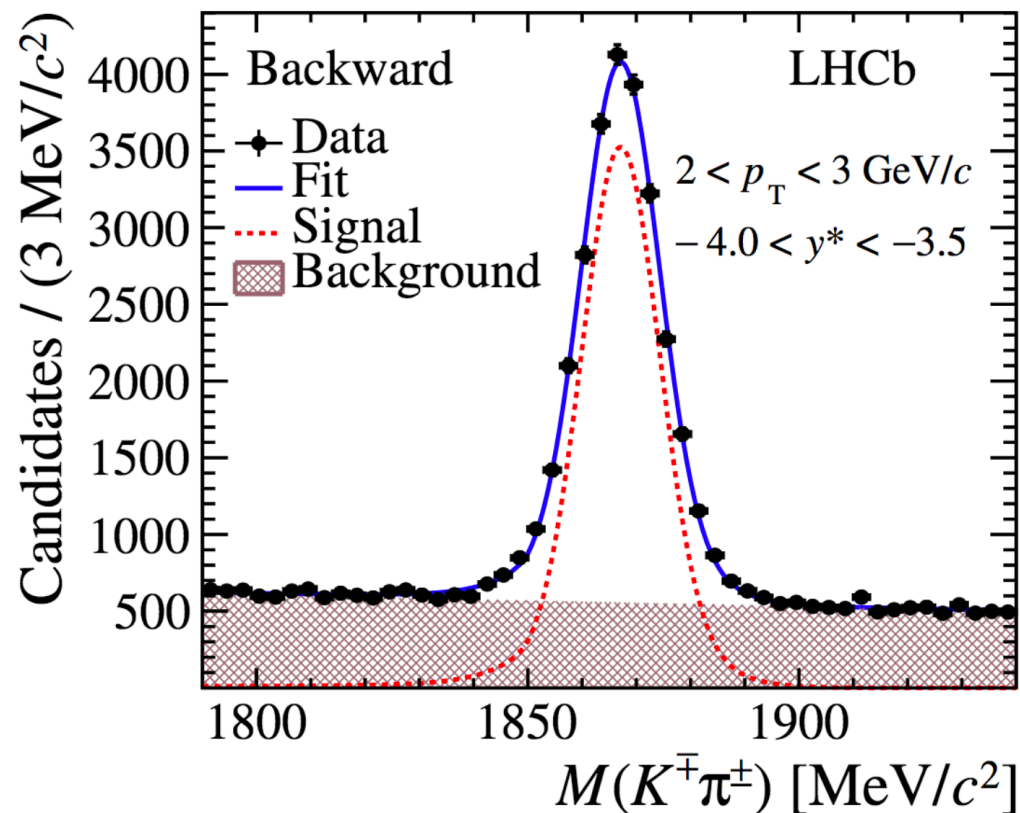
Open Charm Mesons in p Pb collisions: D^0

Fully reconstructed through decay channel $D^0 \rightarrow K^{\mp} \pi^{\pm}$

$$\sqrt{s_{NN}} = 5 \text{ TeV}$$



J. High Energ. Phys. 10 (2017) 90

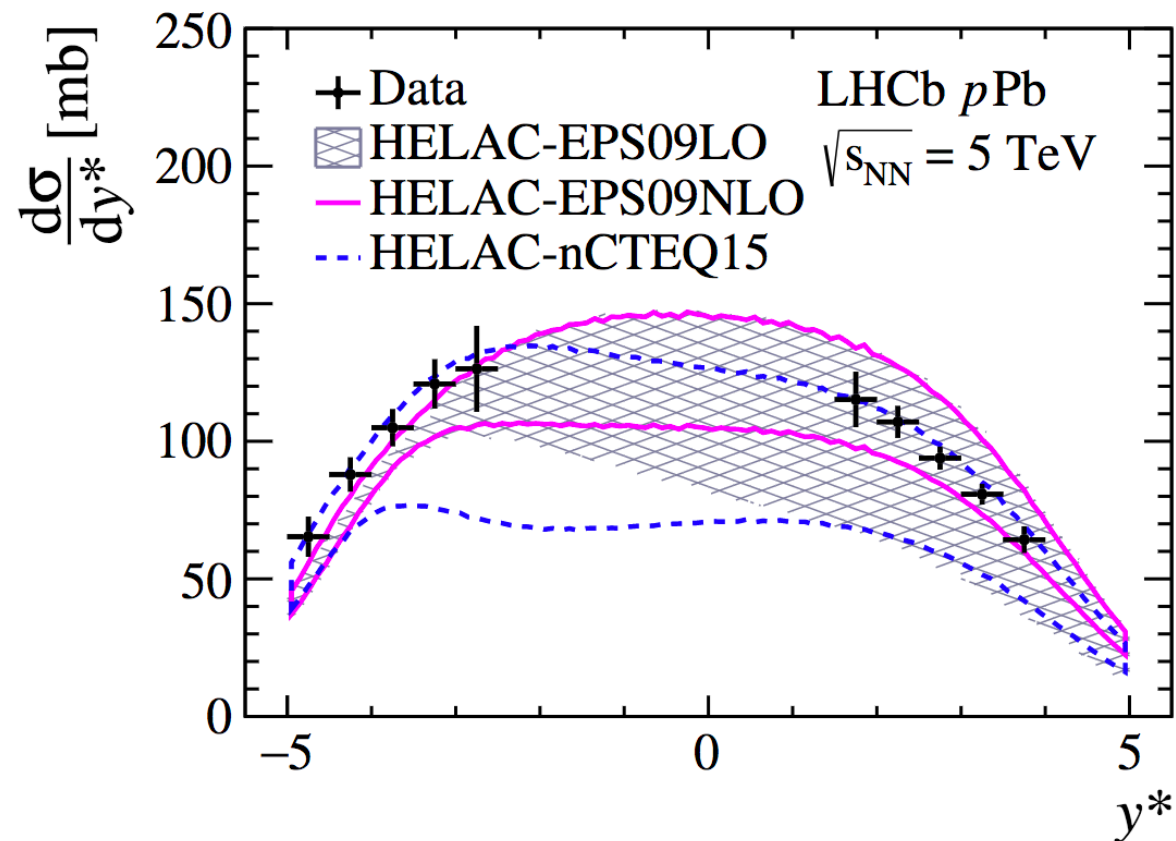
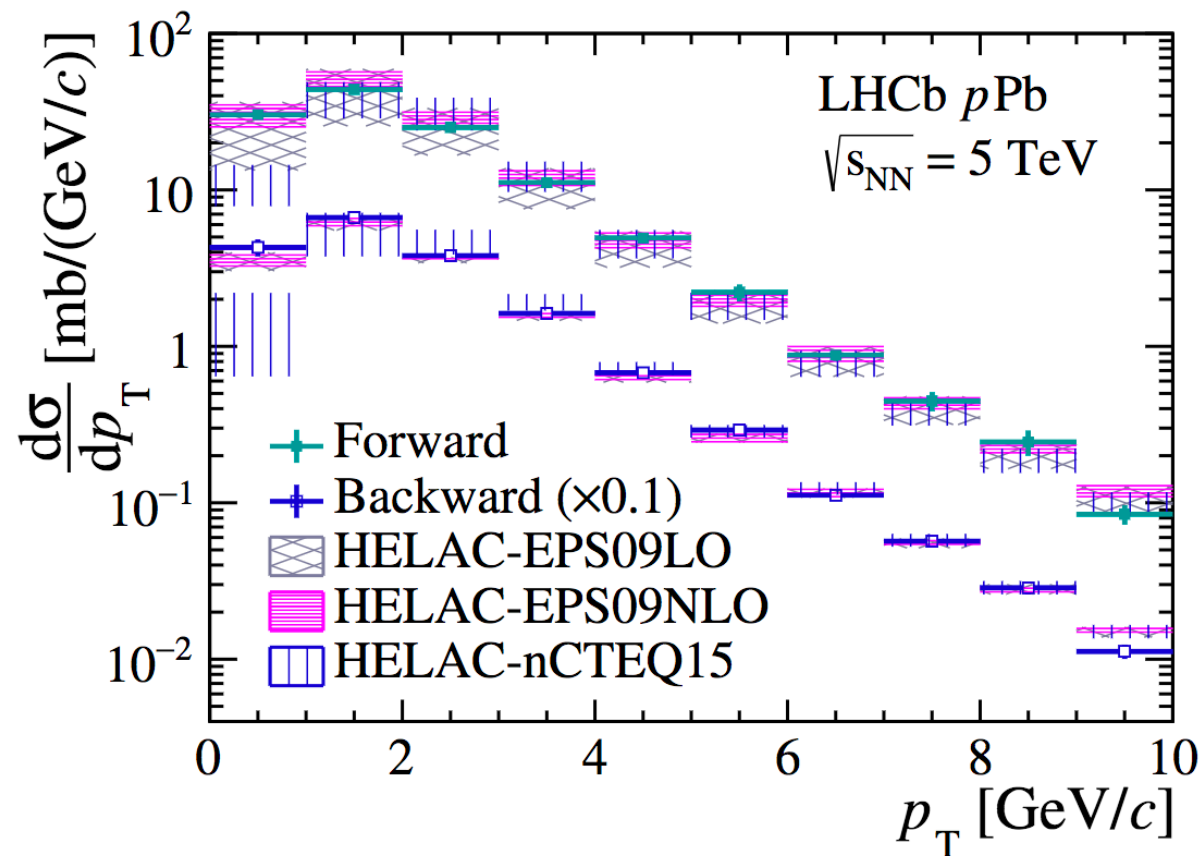


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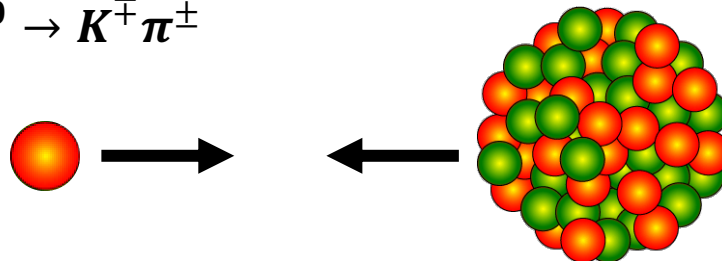
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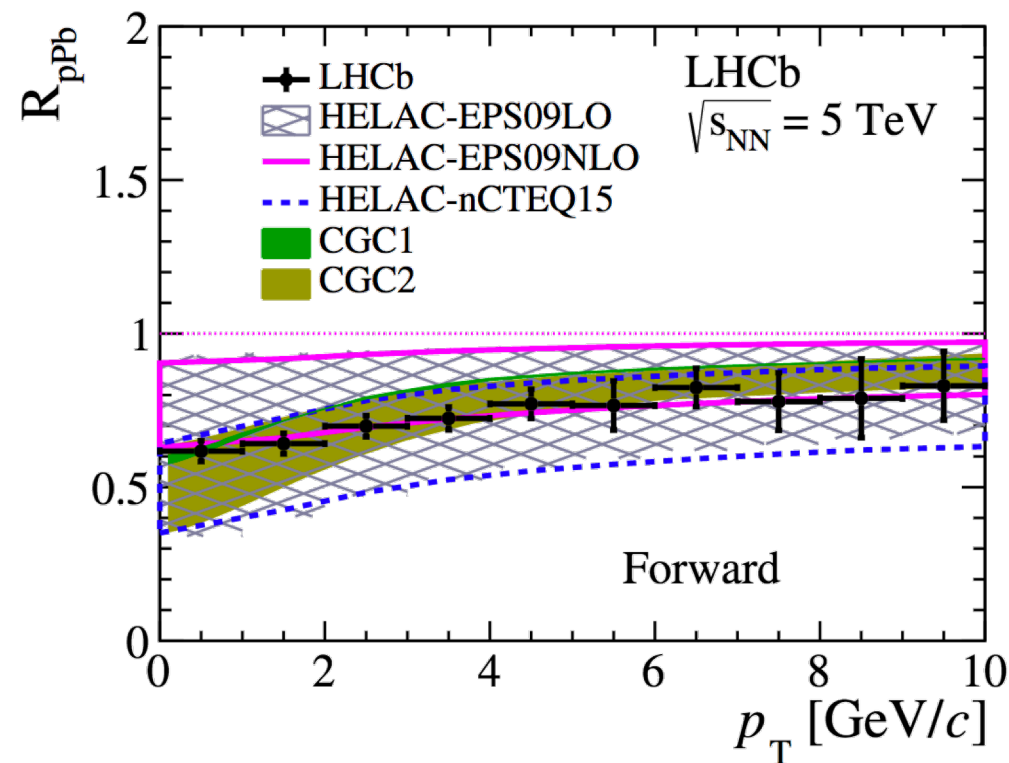
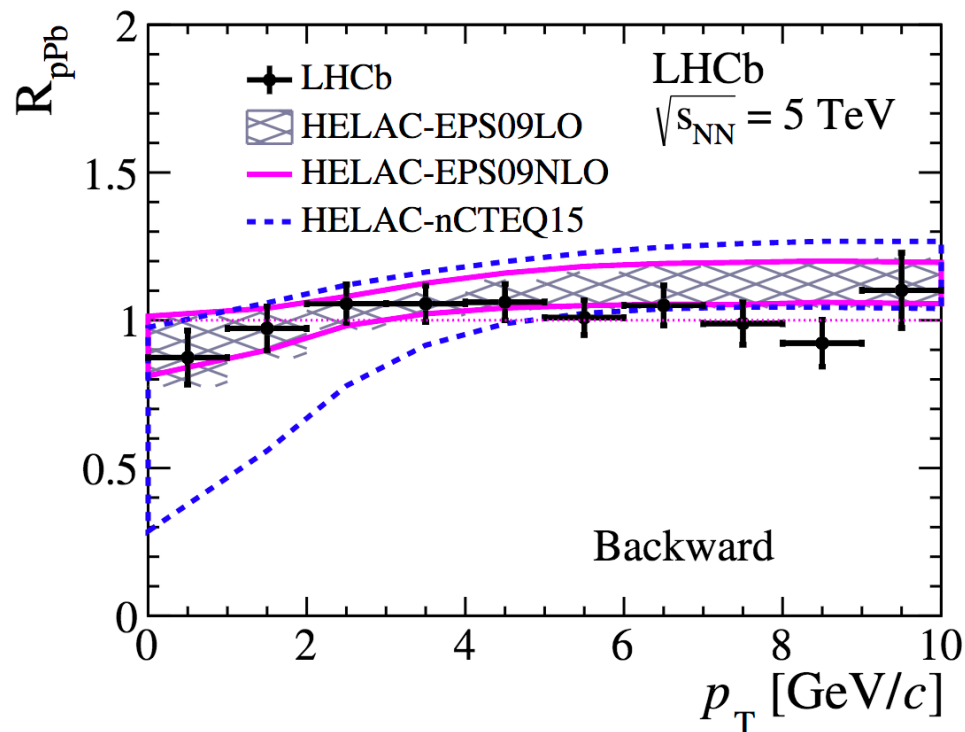
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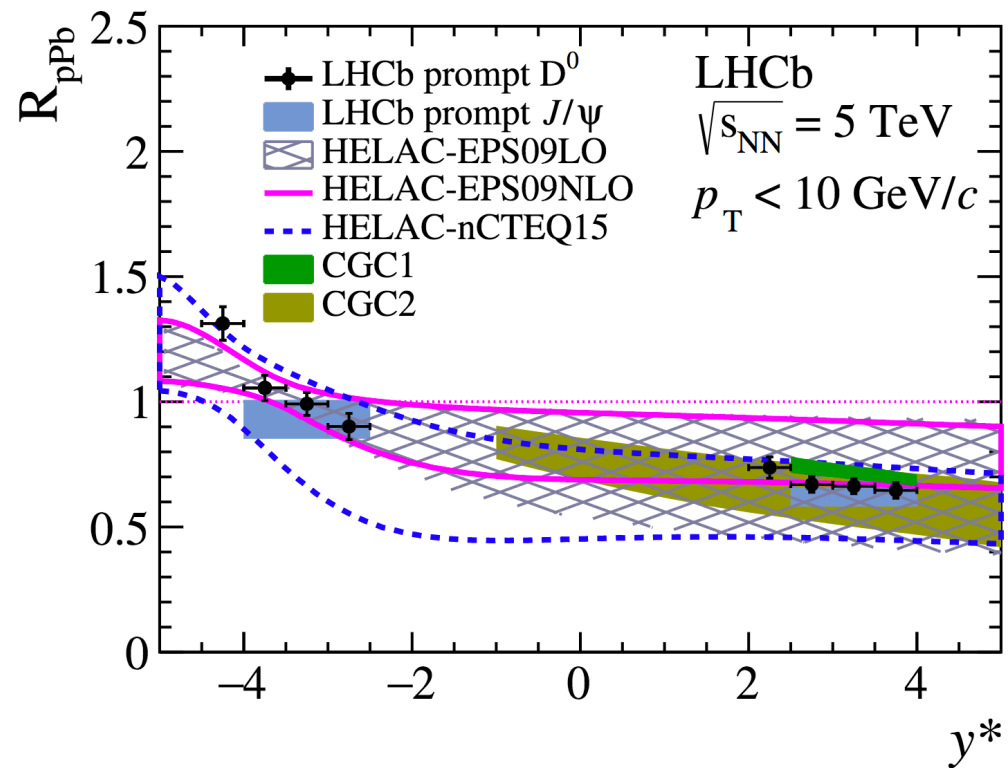


J. High Energ. Phys. 10 (2017) 90

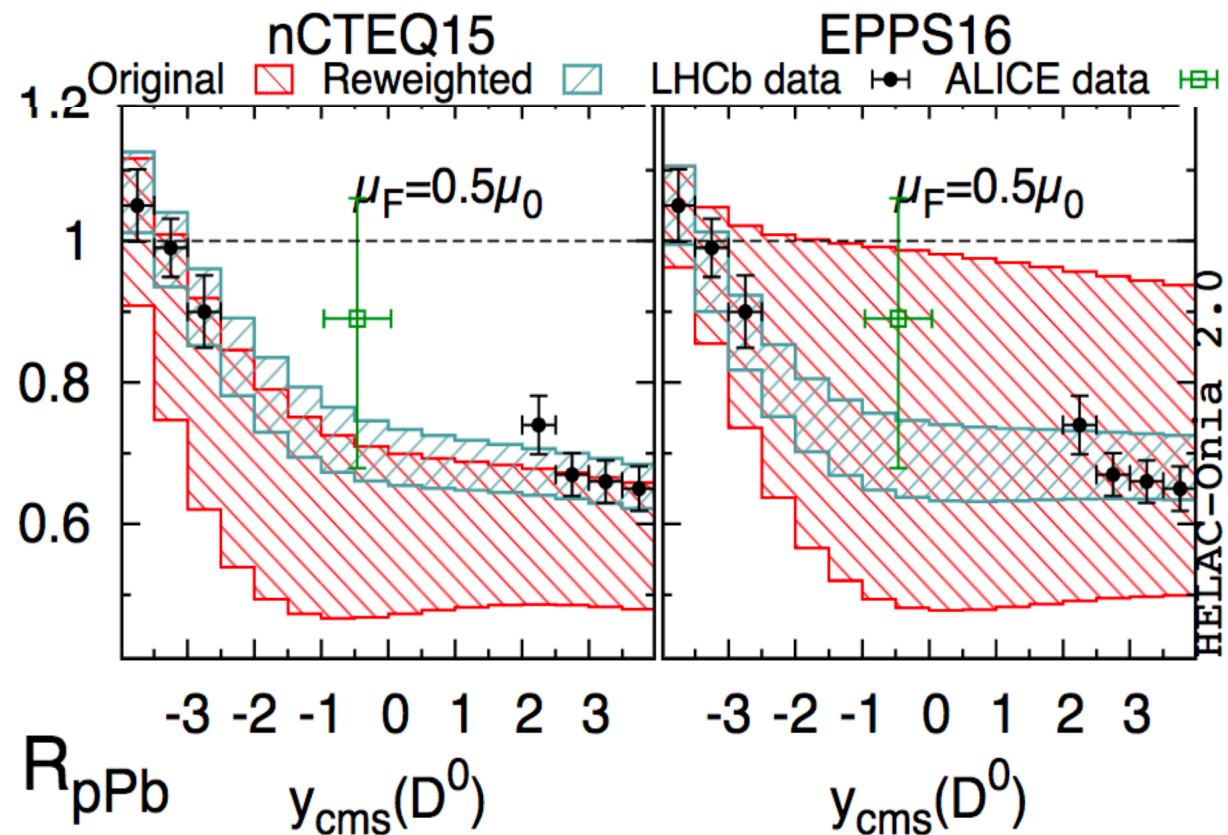
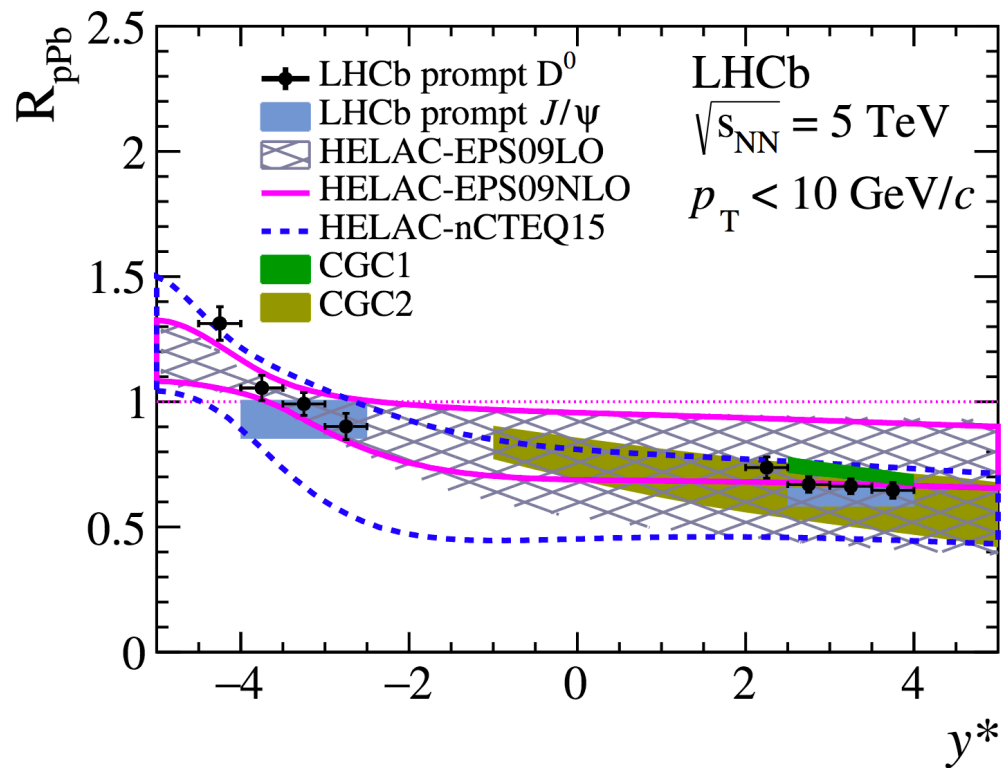


Error bars < calculation uncertainties

Open Charm Mesons in pPb collisions: D^0



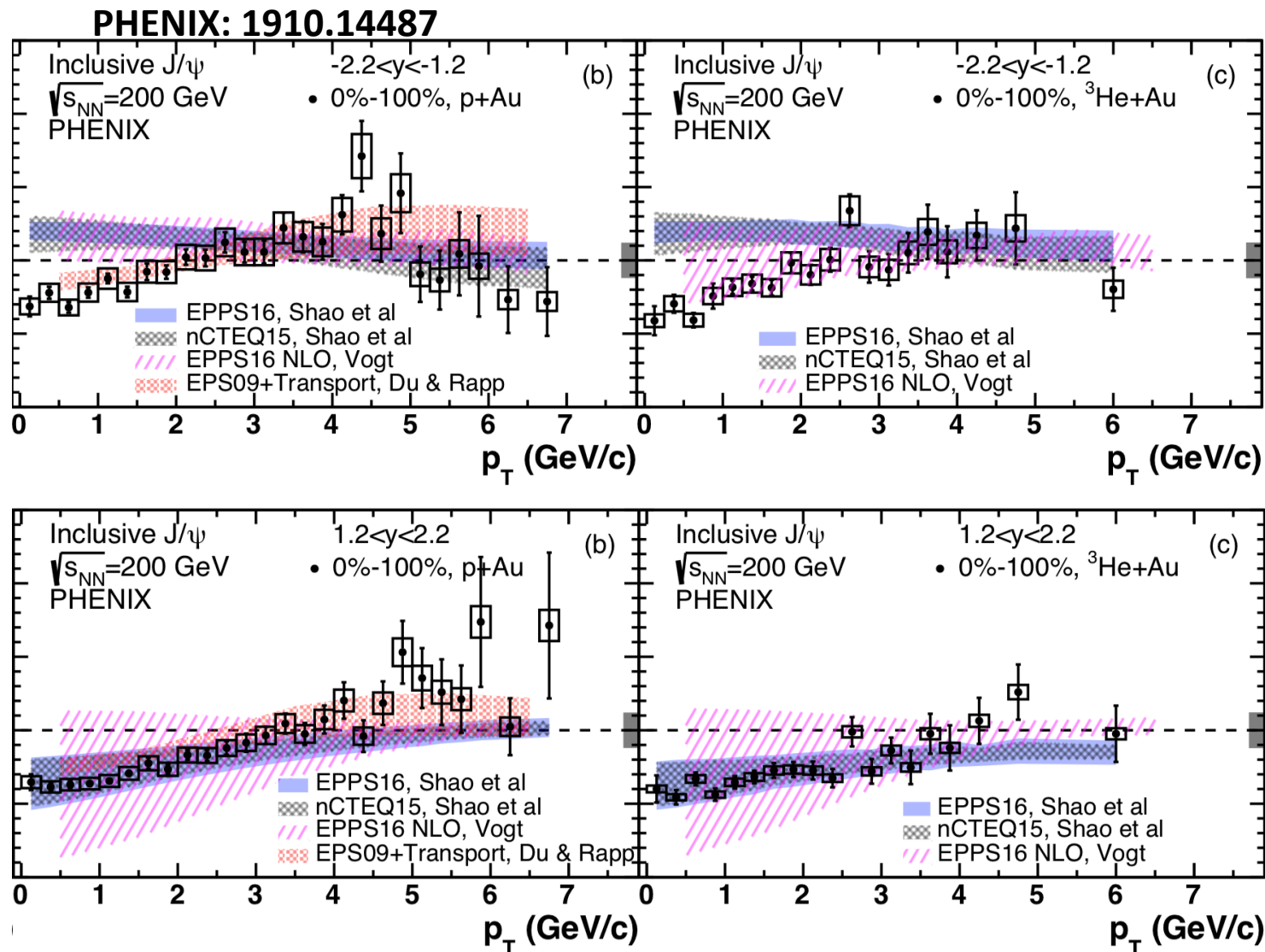
Open Charm Mesons in pPb collisions: D^0



This data is already being used to constrain the gluon nPDF down to $x \sim 5 \times 10^{-6}$

Kusina, Lansberg, Schienbein, Shao,
Gluon shadowing and antishadowing in heavy-flavor production at the LHC
Phys. Rev. Lett. 121, 052004 (2018)

Comparisons with PHENIX p/He+A data

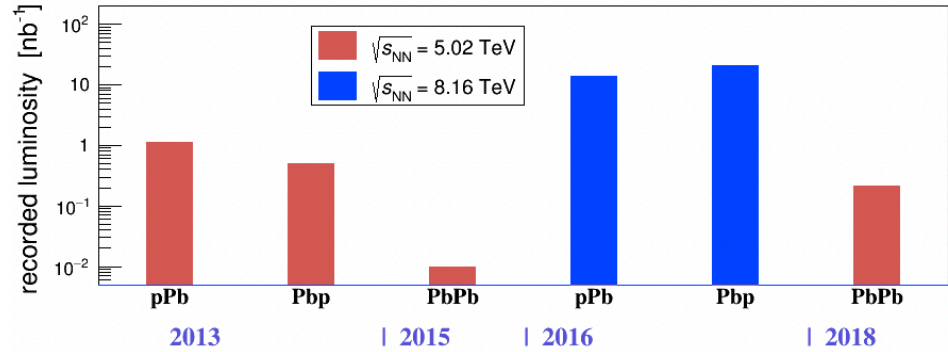


PHENIX recently released p/He+Au Jpsi data at fwd/bkwd rapidity

Comparison with reweighted nPDF calculation shows deviations at low p_T in backwards rapidity

Described by transport calculation

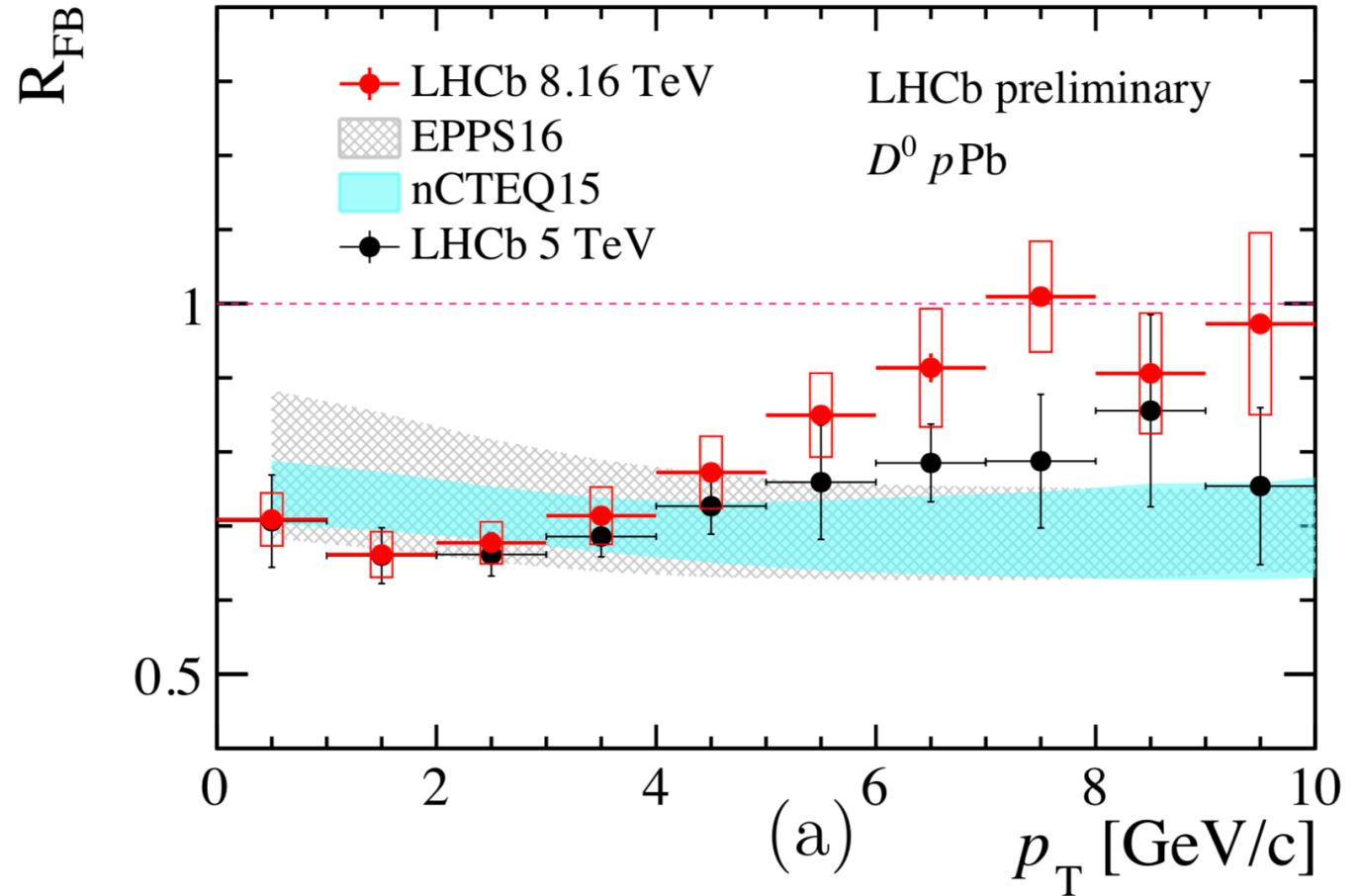
New D results



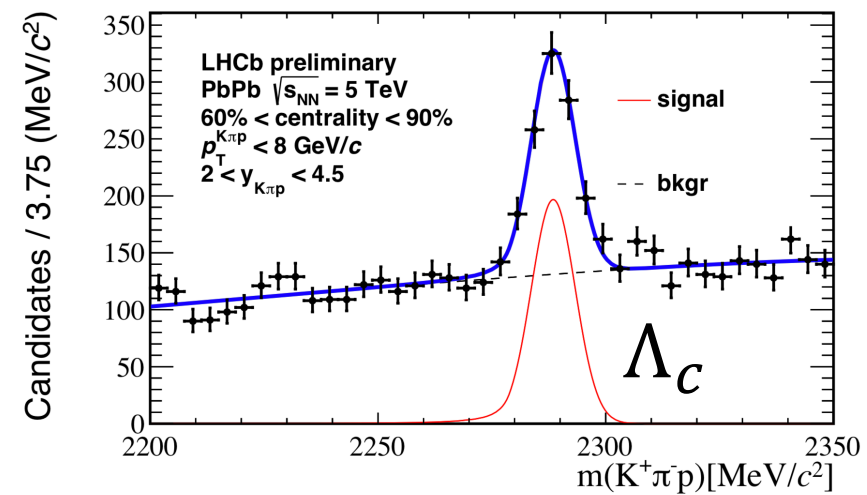
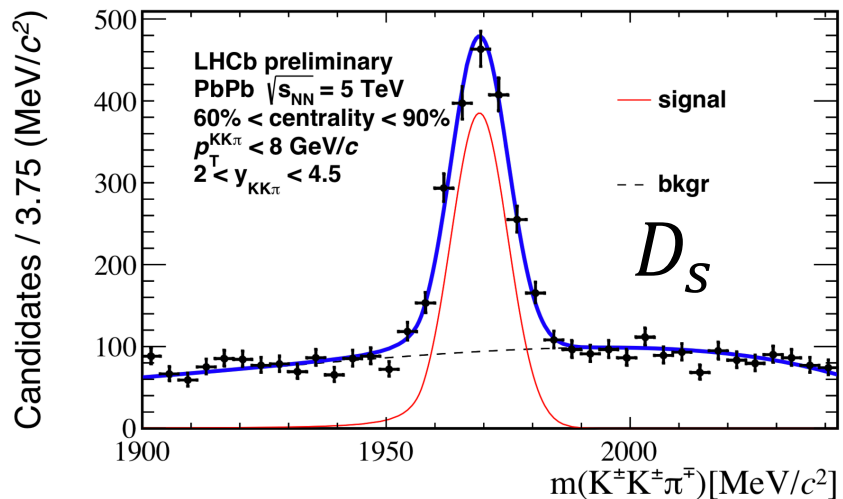
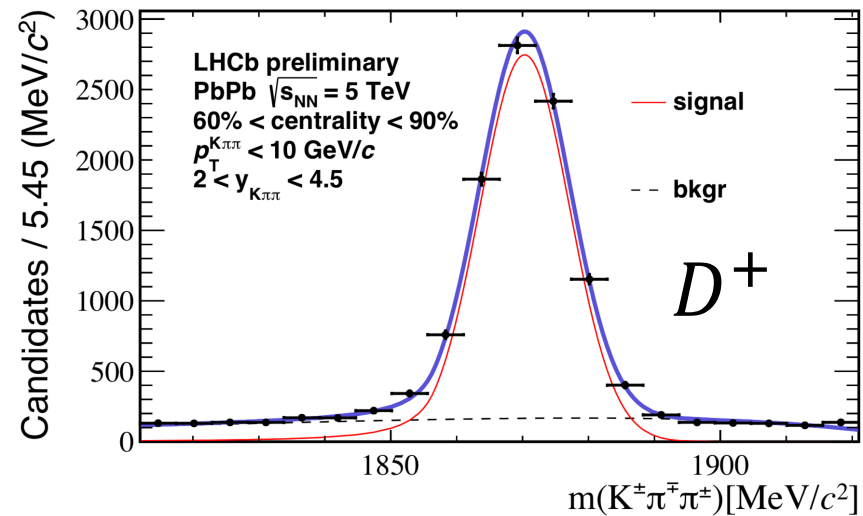
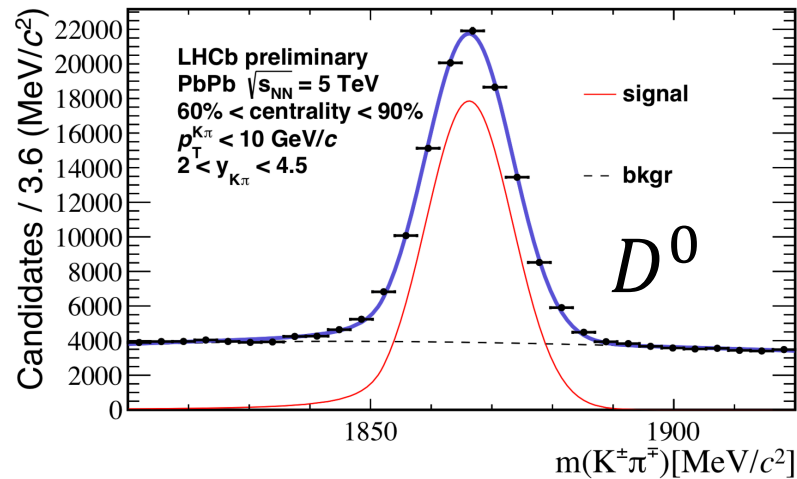
LHCb-CONF-2019-004

New preliminary results from 8 TeV pPb data set:

Significantly better statistical precision
Final evaluation of systematics underway
->More precise constraints on nPDFs



PbPb results to come



Unique results on charm hadrochemistry to come

Fixed target configuration - SMOG

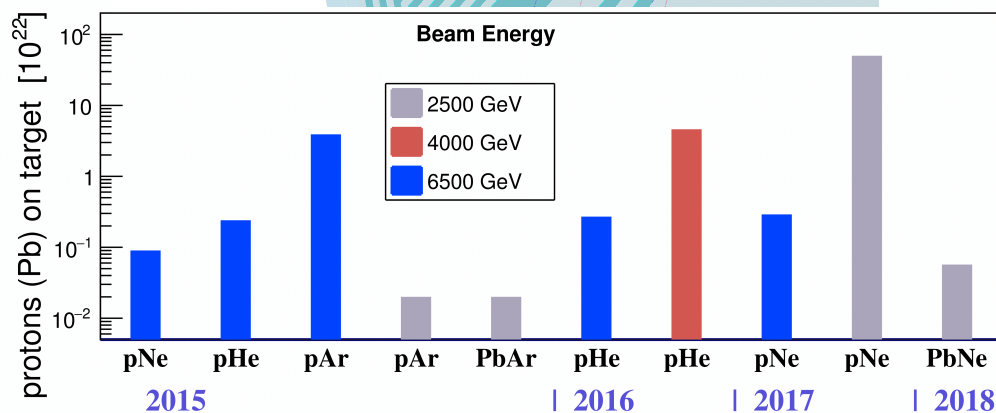
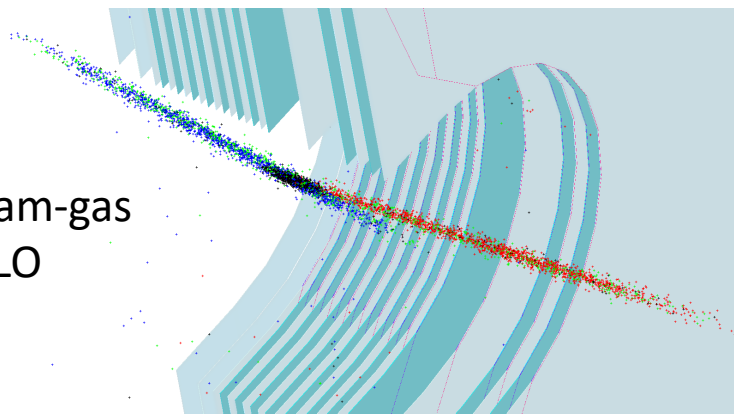
A unique capability at LHCb: inject noble gas into beampipe

Originally intended for precise luminosity measurements:

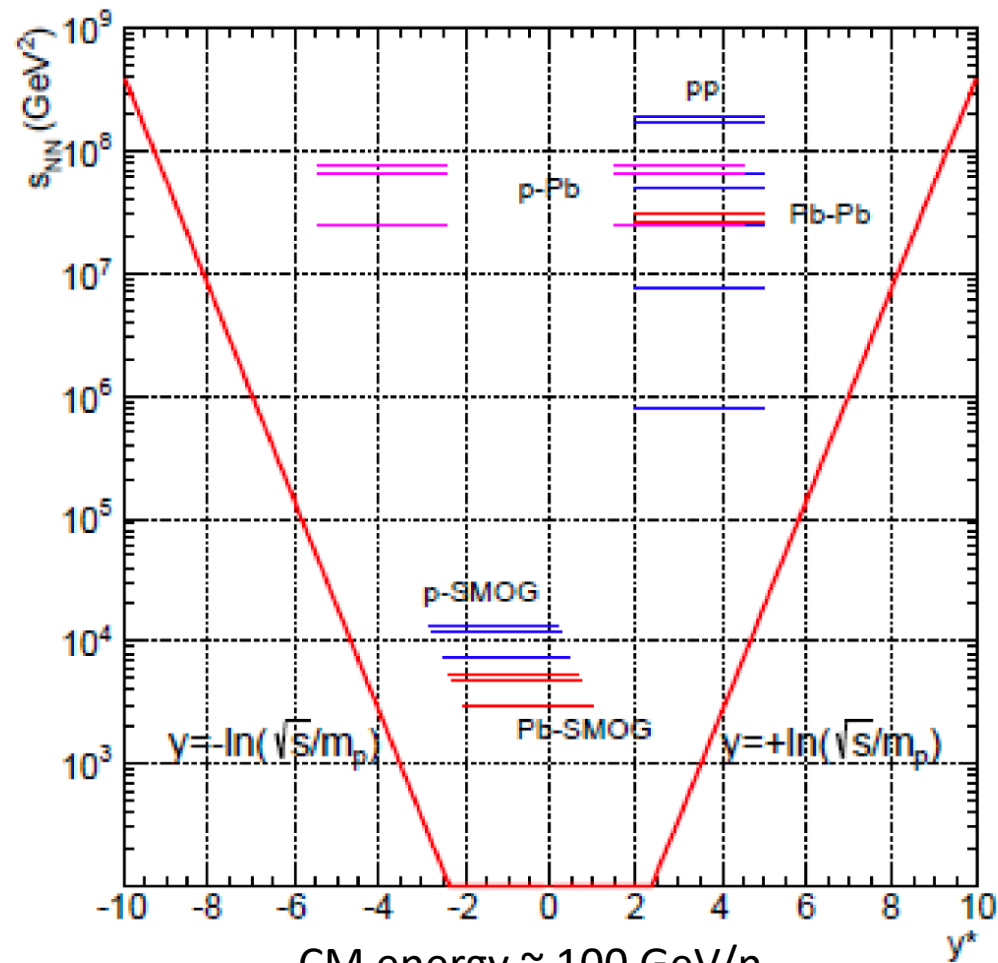
Precision on 2012 pp data is $\pm 1.16\%$, best ever at bunched beam collider

JINST 9 P12005 (2014)

Reconstructed beam-gas
vertices inside VELO



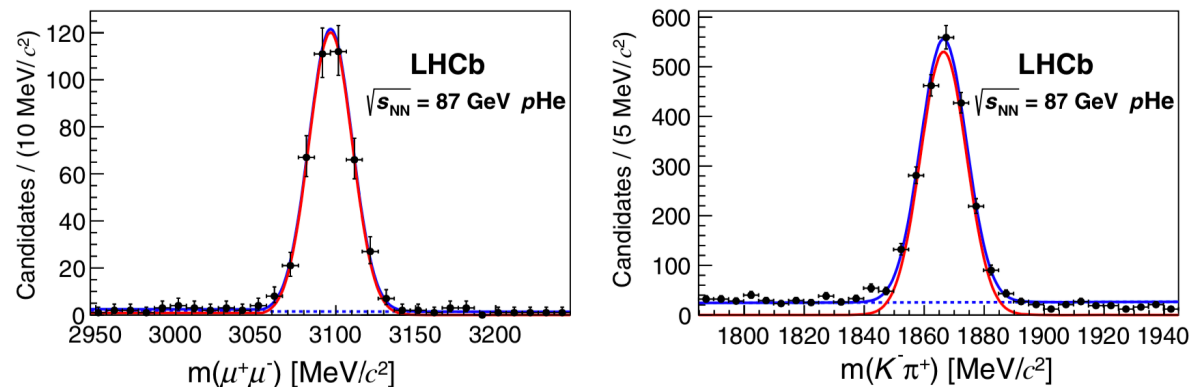
“System for Measurement of Overlap with Gas”



CM energy ~ 100 GeV/n
No limits on centrality

First SMOG Heavy Flavor Result

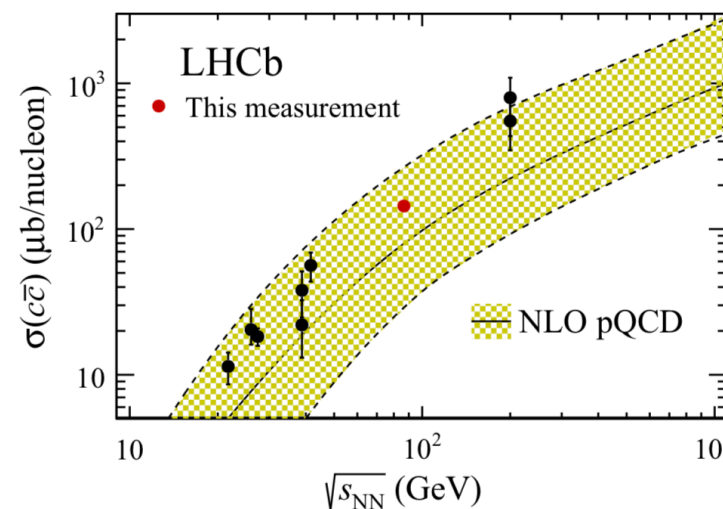
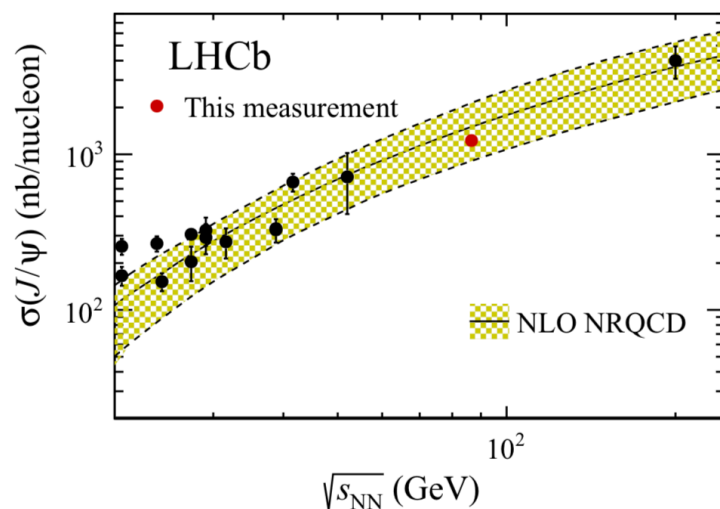
Phys. Rev. Lett. 122, 132002 (2019)



First charm result from SMOG: pHe and pAr data

Unique access to precision heavy flavor probes in energy range between SPS and RHIC

MAJOR upgrade to SMOG system ongoing



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Quark Model of Hadrons

Volume 8, number 3

PHYSICS LETTERS

1 February 1964

AN SU_3 MODEL FOR STRONG INTERACTION SYMMETRY AND ITS BREAKING

A SCHEMATIC MODEL OF BARYONS AND MESONS *

M. GELL-MANN

California Institute of Technology, Pasadena, California

Received 4 January 1964

G. Zweig *)

CERN - Geneva

8182/TH.401

17 January 1964

A simpler and more elegant scheme can be constructed if we allow non-integral values for the charges. We can dispense entirely with the basic baryon b if we assign to the triplet t the following properties: spin $\frac{1}{2}$, $z = -\frac{1}{3}$, and baryon number $\frac{1}{3}$. We then refer to the members $u^{\frac{2}{3}}$, $d^{-\frac{1}{3}}$, and $s^{-\frac{1}{3}}$ of the triplet as "quarks" q and the members of the anti-triplet as anti-quarks \bar{q} . Baryons can now be constructed from quarks by using the combinations (qqq) , $(qqq\bar{q})$, etc., while mesons are made out of $(q\bar{q})$, $(q\bar{q}\bar{q})$, etc. It is assumed that the lowest baryon configuration (qqq) gives just the representations **1**, **8**, and **10** that have been observed, while the lowest meson configuration $(q\bar{q})$ similarly gives just **1** and **8**.

In general, we would expect that baryons are built not only from the product of three quarks, AAA , but also from $\bar{A}AAAA$, $\bar{A}AAAAA$, etc., where \bar{A} denotes an anti-quark. Similarly, mesons could be formed from $\bar{A}A$, $\bar{A}AAA$ etc. For the low mass mesons and baryons we will assume the simplest possibilities, $\bar{A}A$ and AAA , that is, "deuces and treys".

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8182/TH.401

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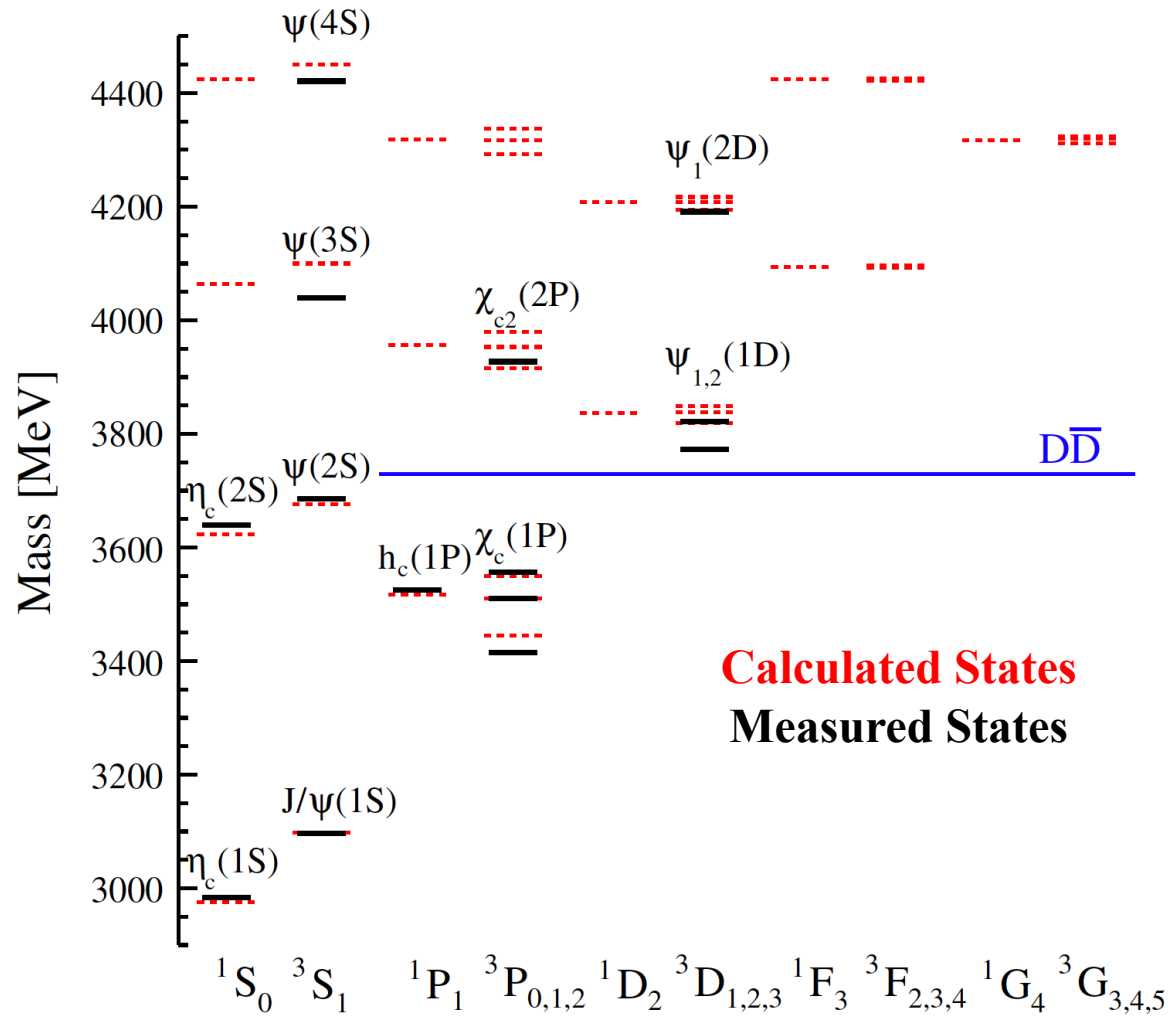
CERN - Geneva

A simpler and more elegant scheme can be constructed if we allow non-integral values for the charges. We can dispense entirely with the basic baryon b if we assign to the triplet t the following properties: spin $\frac{1}{2}$, $z = -\frac{1}{3}$, and baryon number $\frac{1}{3}$. We then refer to the members $u^{\frac{2}{3}}$, $d^{-\frac{1}{3}}$, and $s^{-\frac{1}{3}}$ of the triplet as "quarks" q and the members of the anti-triplet as anti-quarks \bar{q} . Baryons can now be constructed from quarks by using the combinations (qqq) , $(qqq\bar{q})$, etc., while mesons are made out of $(q\bar{q})$, $(qq\bar{q}\bar{q})$, etc. It is assumed that the lowest baryon configuration (qqq) gives just the representations **1**, **8**, and **10** that have been observed, while the lowest meson configuration $(q\bar{q})$ similarly gives just **1** and **8**.

In general, we would expect that baryons are built not only from the product of three quarks, AAA , but also from $\bar{A}AAAA$, $\bar{A}AAAAA$, etc., where \bar{A} denotes an anti-quark. Similarly, mesons could be formed from $\bar{A}A$, $\bar{A}AAA$ etc. For the low mass mesons and baryons we will assume the simplest possibilities, $\bar{A}A$ and AAA , that is, "deuces and treys".

States with >3 quarks have been expected since the beginning of the quark model

Conventional $c\bar{c}$ States



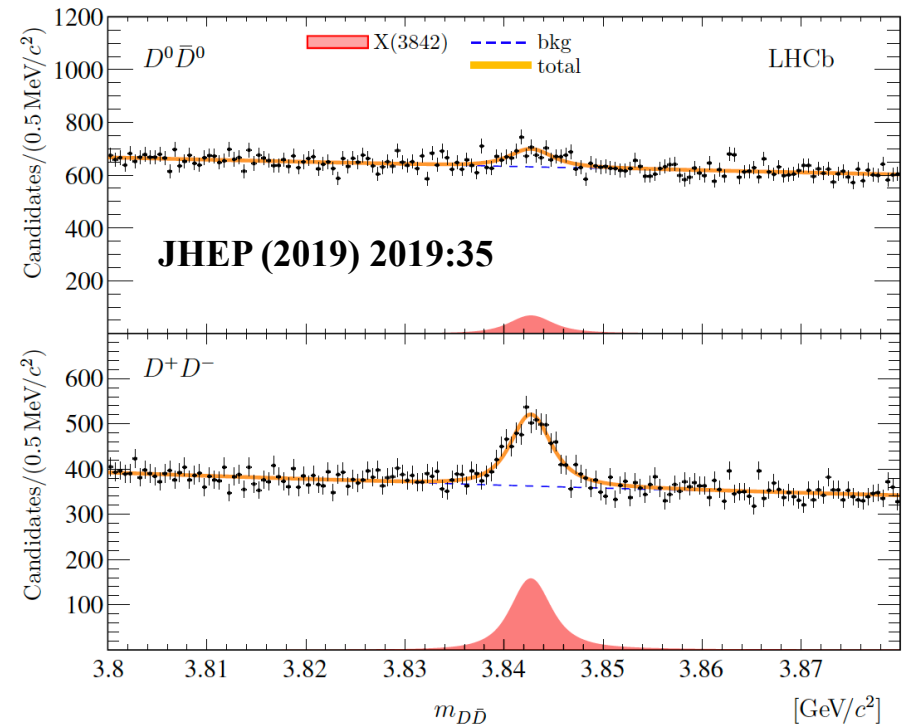
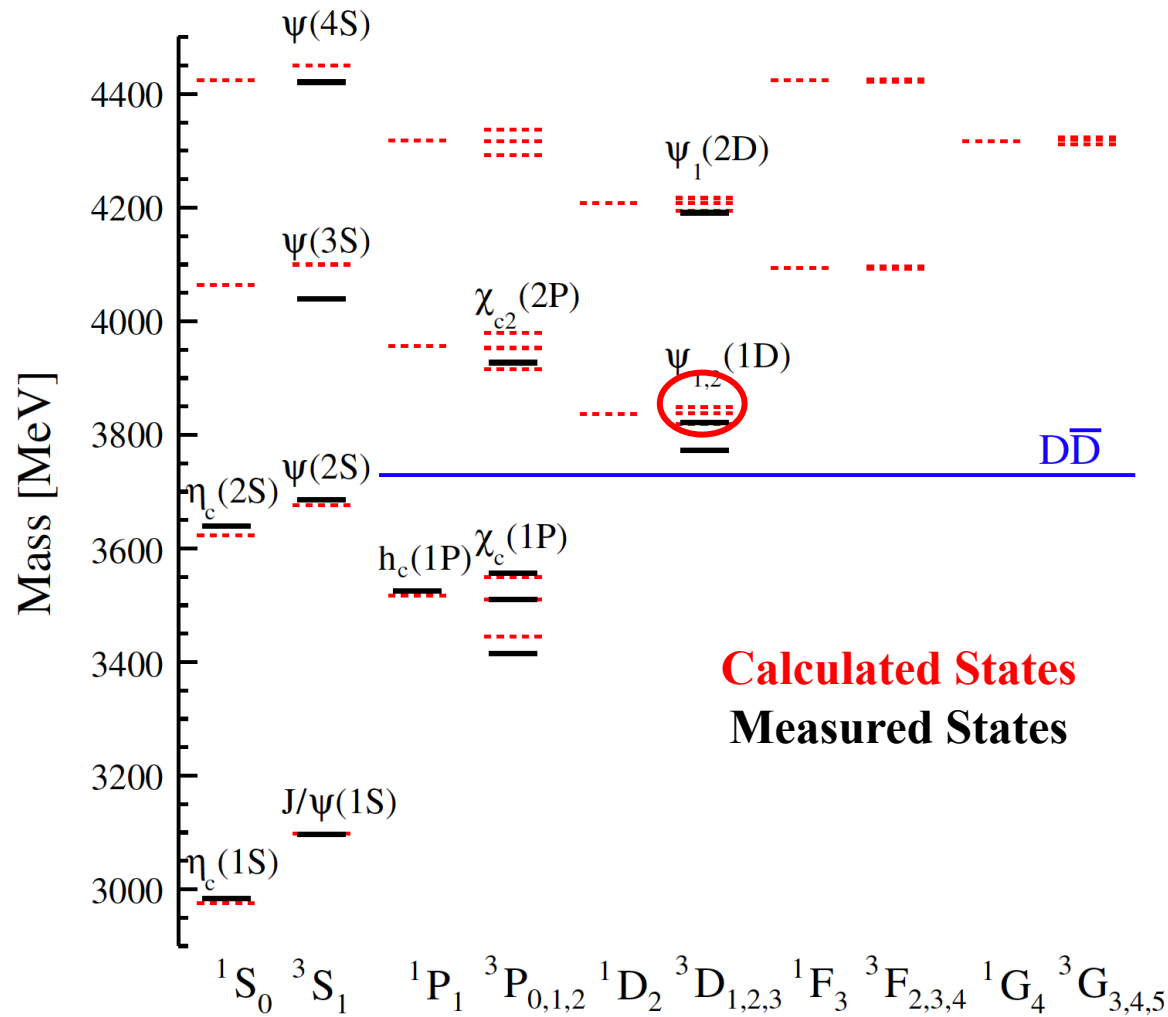
Nonrelativistic potential model: solve Schrodinger equation with the potential

$$V_0^{(c\bar{c})}(r) = -\frac{4}{3} \frac{\alpha_s}{r} + br + \frac{32\pi\alpha_s}{9m_c^2} \tilde{\delta}_\sigma(r) \vec{S}_c \cdot \vec{S}_{\bar{c}}$$

Barnes, Godfrey, Swanson,
Phys. Rev. D 72, 054026 (2005)

Rev. Mod. Phys. 90, 015003 (2018)

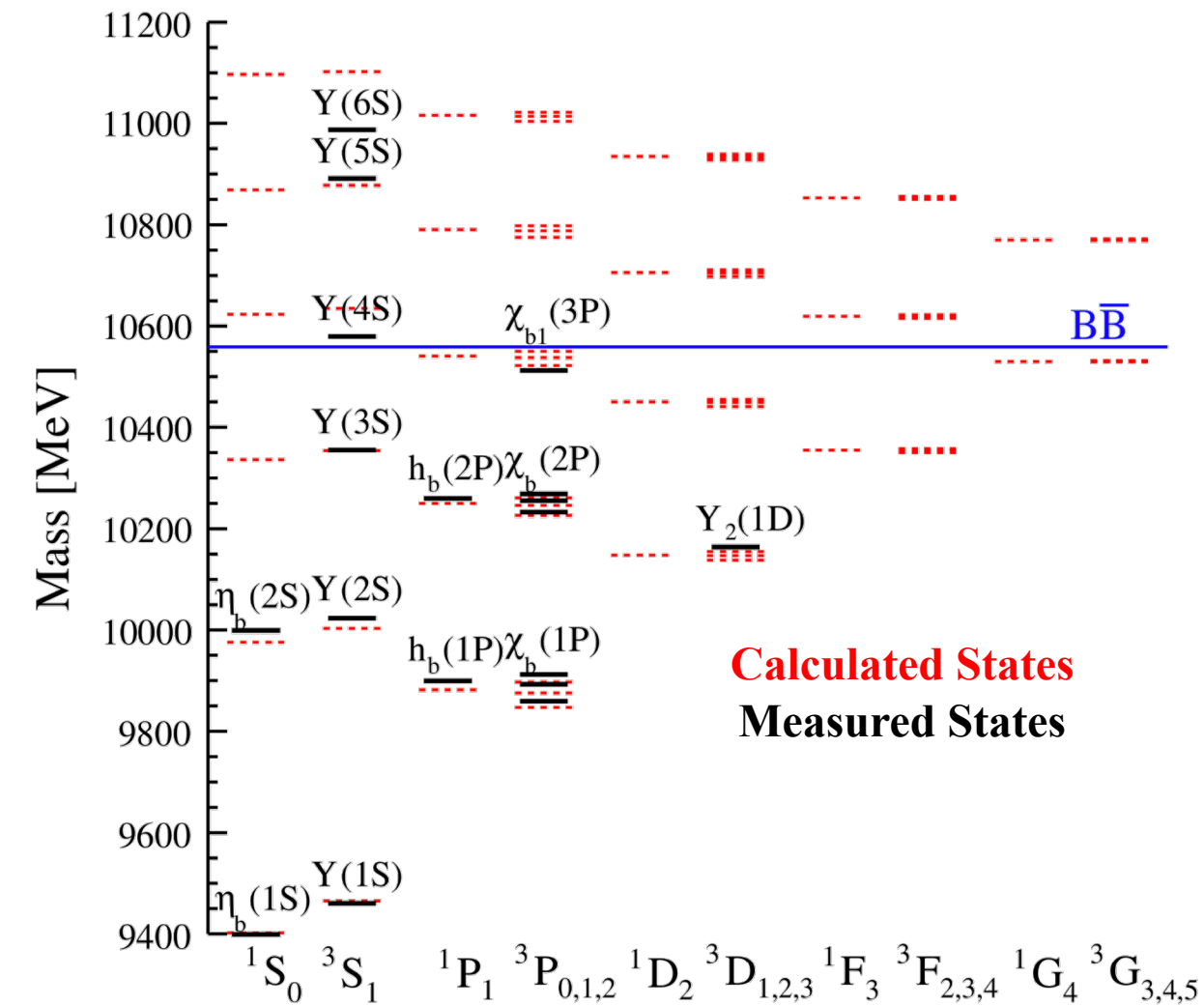
Conventional $c\bar{c}$ States



New charmonium states still being found: LHCb observed state consistent with $\psi_3(1^3D_3)$ found in $D\bar{D}$ and D^+D^- mass spectra in 2019

Measured mass: $3842.71 \pm 0.16 \pm 0.12$ MeV
Predicted mass: 3849 MeV

Conventional $b\bar{b}$ States

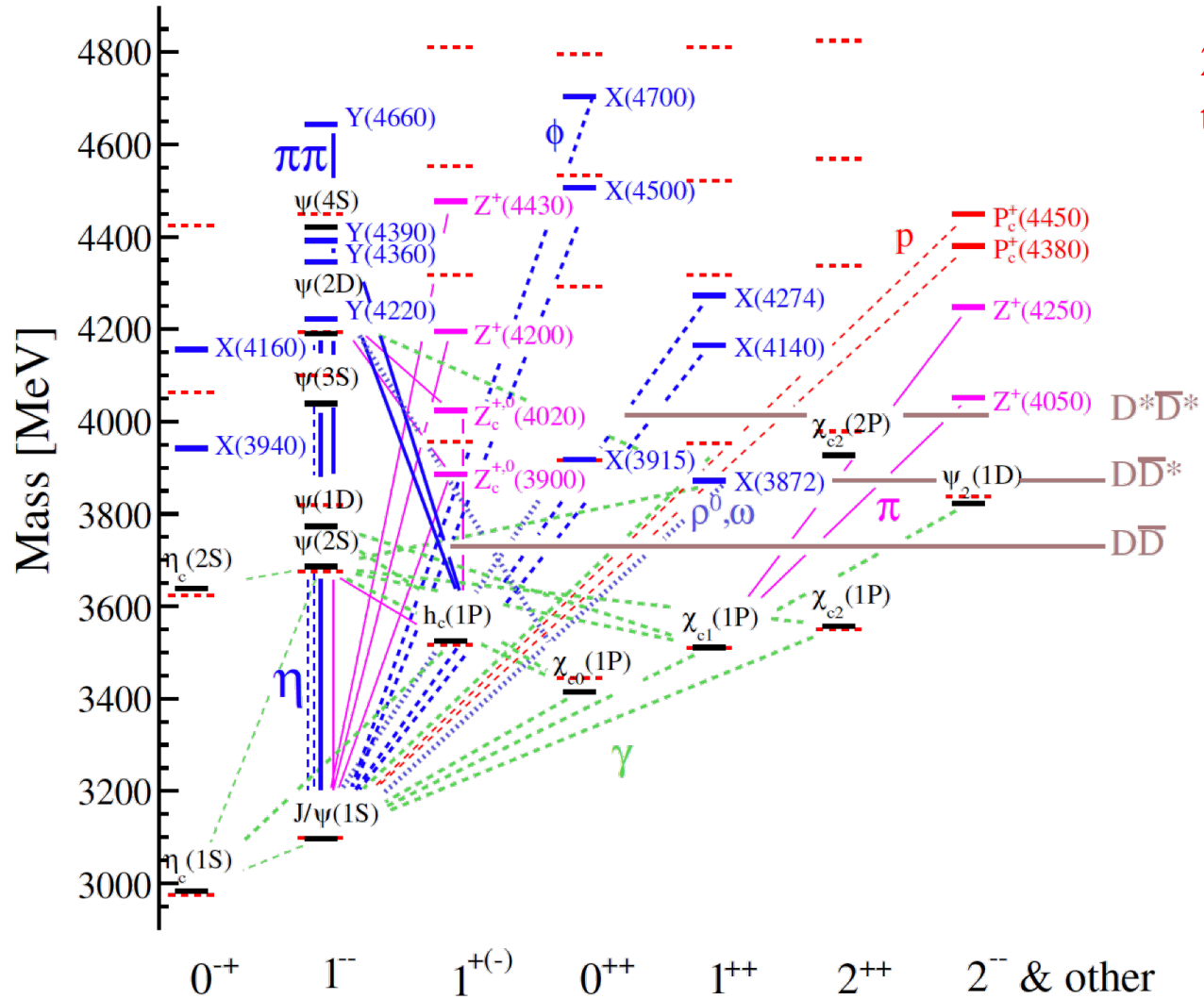


*The allowed $c\bar{c}$ and $b\bar{b}$ bound states are well understood**

Crucial to account for conventional states when searching for exotics

Rev. Mod. Phys. 90, 015003 (2018)

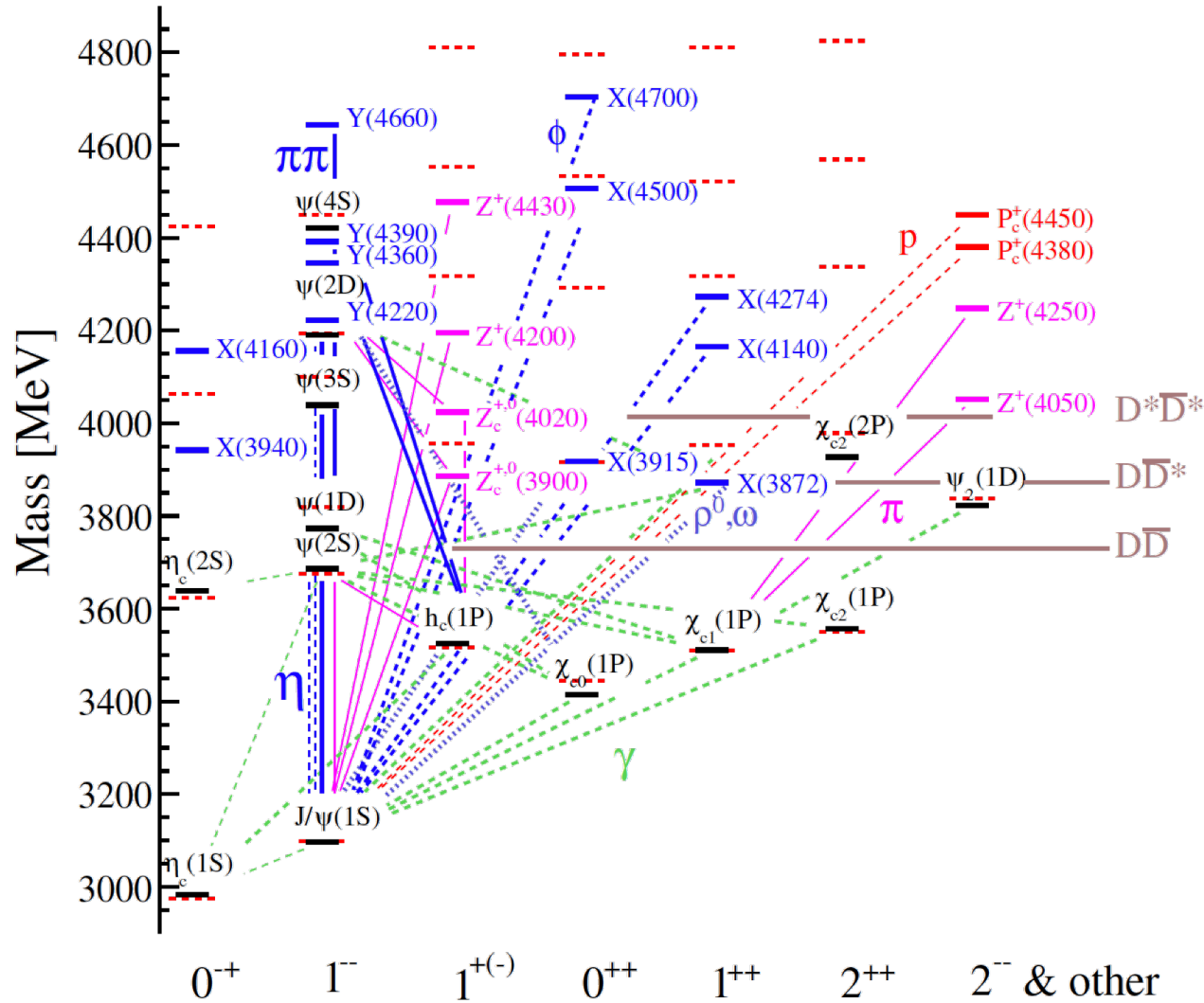
Exotic $c\bar{c}$ States



20+ states containing $c\bar{c}$ have been discovered since 2003 that do not fit in the picture of typical charmonium: Collectively known as “XYZ” particles

Rev. Mod. Phys. 90, 015003 (2018)

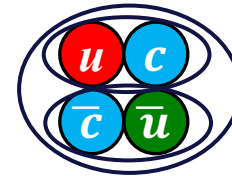
Exotic $c\bar{c}$ States



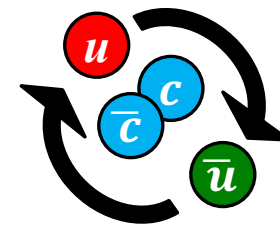
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Multiple explanations explored in literature:

Compact tetraquark/pentaquark



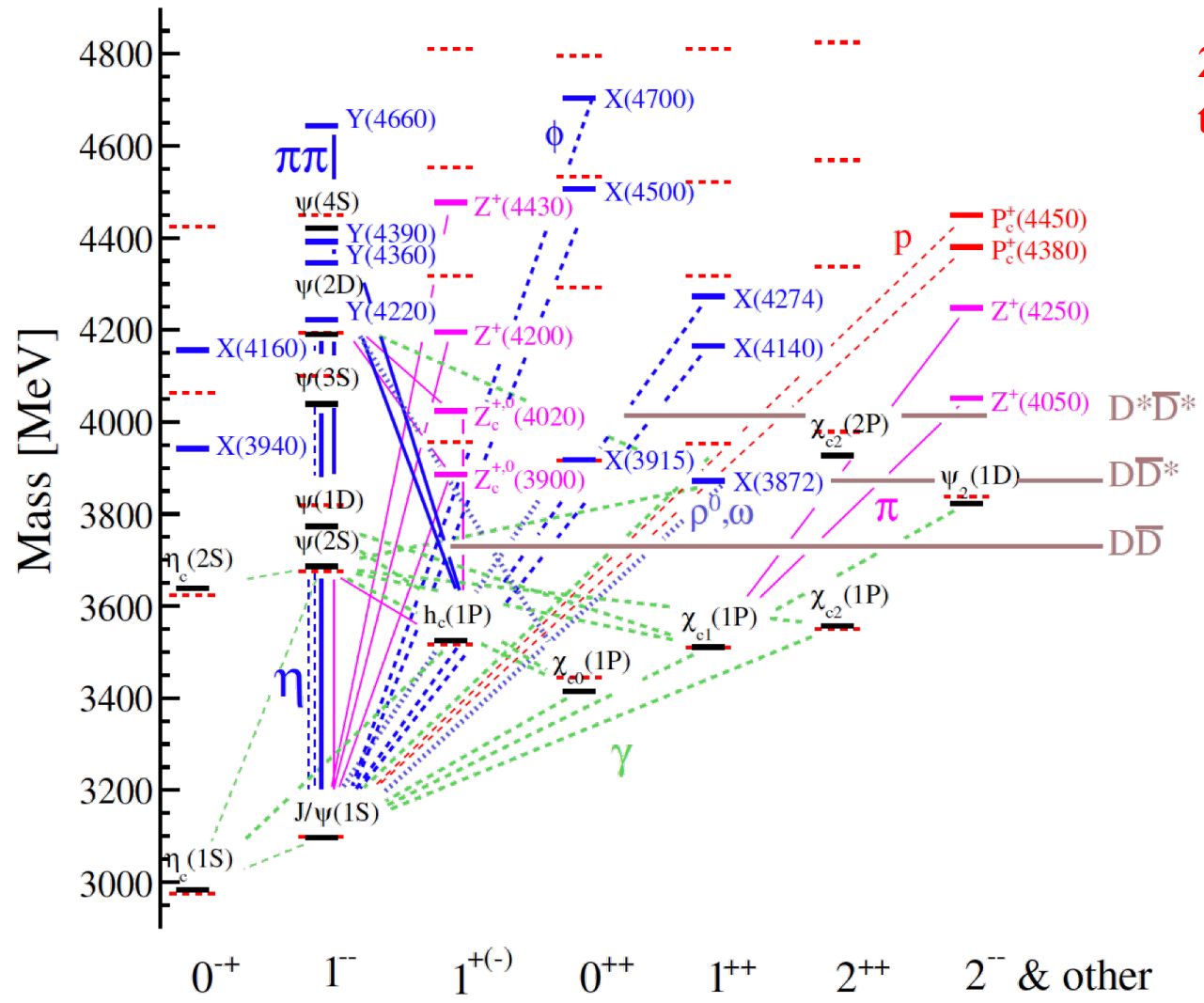
Diquark-diquark
PRD 71, 014028 (2005)
PLB 662 424 (2008)



Hadrocharmonium/
adjoint charmonium
PLB 666 344 (2008)
PLB 671 82 (2009)

Rev. Mod. Phys. 90, 015003 (2018)

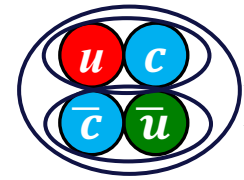
Exotic $c\bar{c}$ States



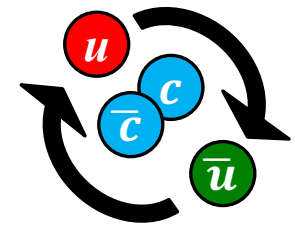
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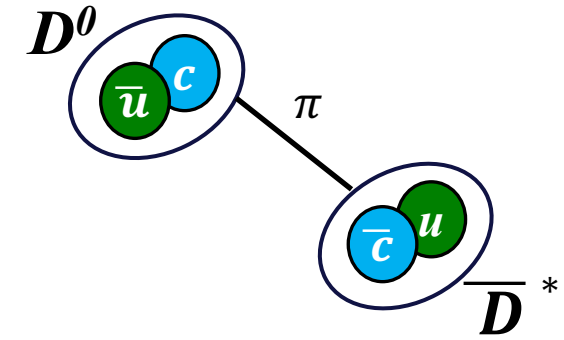
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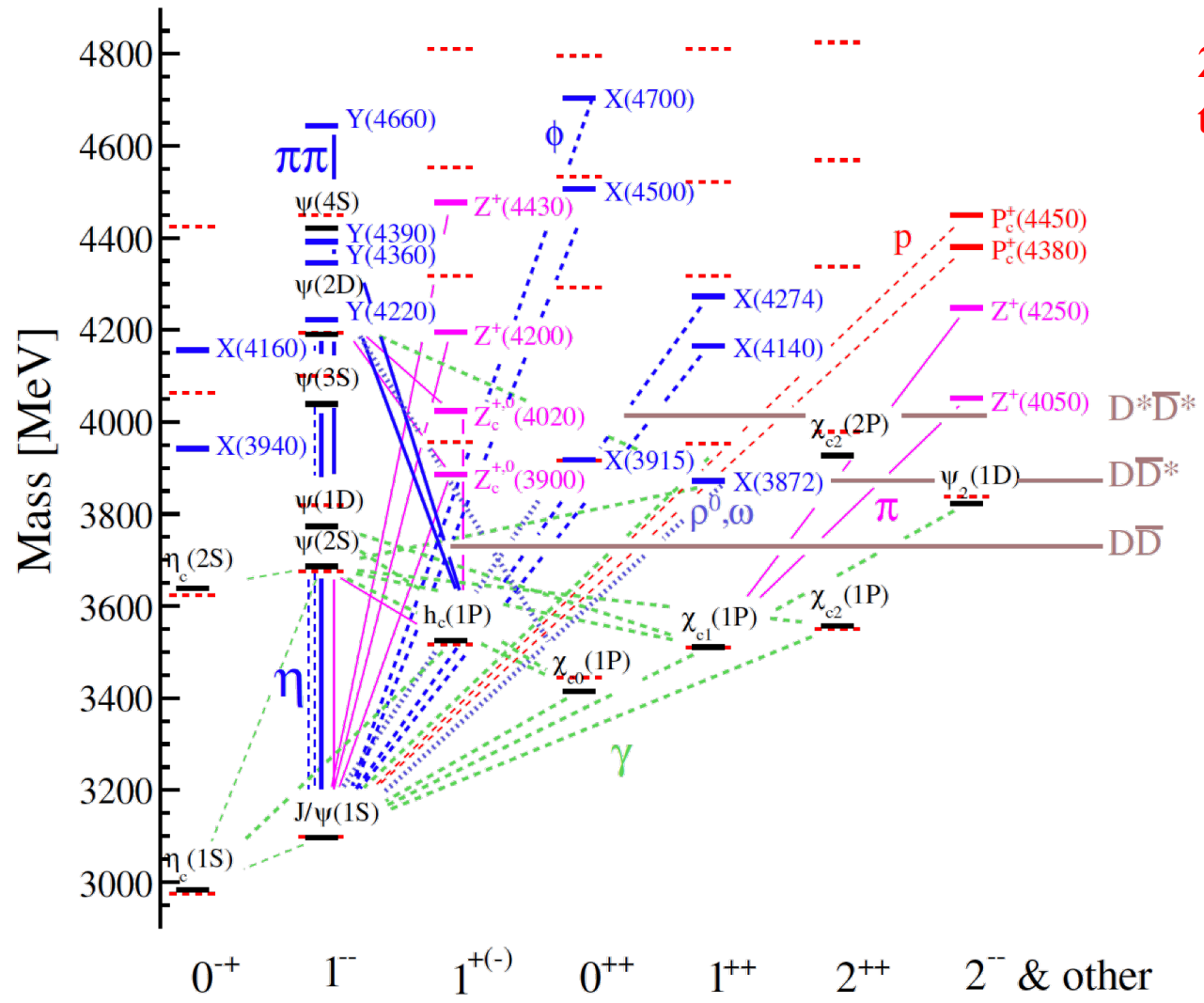
Hadronic Molecules

PLB 590 209 (2004)
PRD 77 014029 (2008)
PRD 100 0115029(R) (2019)



Rev. Mod. Phys. 90, 015003 (2018)

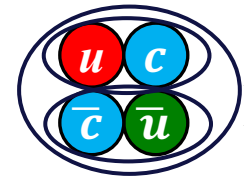
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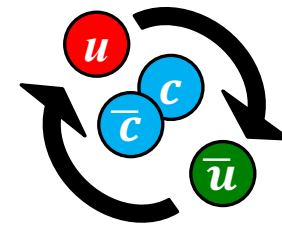
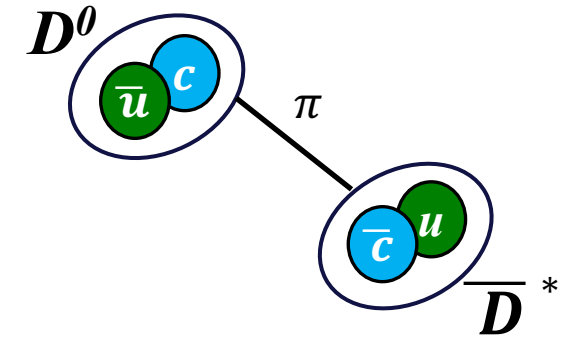
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PRD 100 0115029(R) (2019)



Hadrocharmonium/adjoint charmonium
PLB 666 344 (2008)
PLB 671 82 (2009)

Mixtures of exotic + conventional states

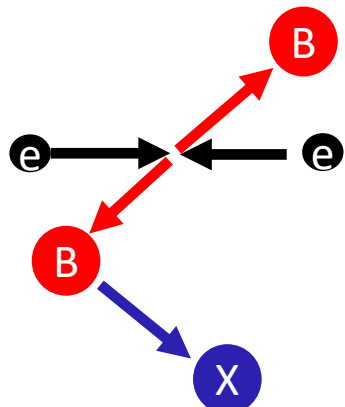
$$X = a |c\bar{c}\rangle + b |c\bar{c}q\bar{q}\rangle$$

PLB 578 365 (2004)
PRD 96 074014 (2017)

Rev. Mod. Phys. 90, 015003 (2018)

Production of Exotic States

Discovery in B decays



B factories: $e^+e^- \rightarrow \Upsilon(4S) \rightarrow B\bar{B}$

Belle, BaBar

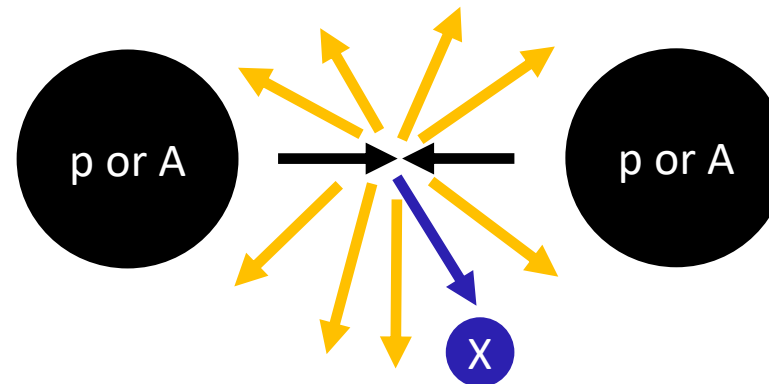
- Well constrained initial state
- Low backgrounds

Hadron Colliders eg LHC: $pp \rightarrow B\bar{B} + X$

LHCb, ATLAS, CMS, D0, CDF, etc

- High rate
- Large total cross section
- Access to wider range of states

Prompt production



AA, pA(?): high particle density

- Production via coalescence or recombination at freezeout

pp, pA: low particle density

- Breakup of weakly bound states in late stages

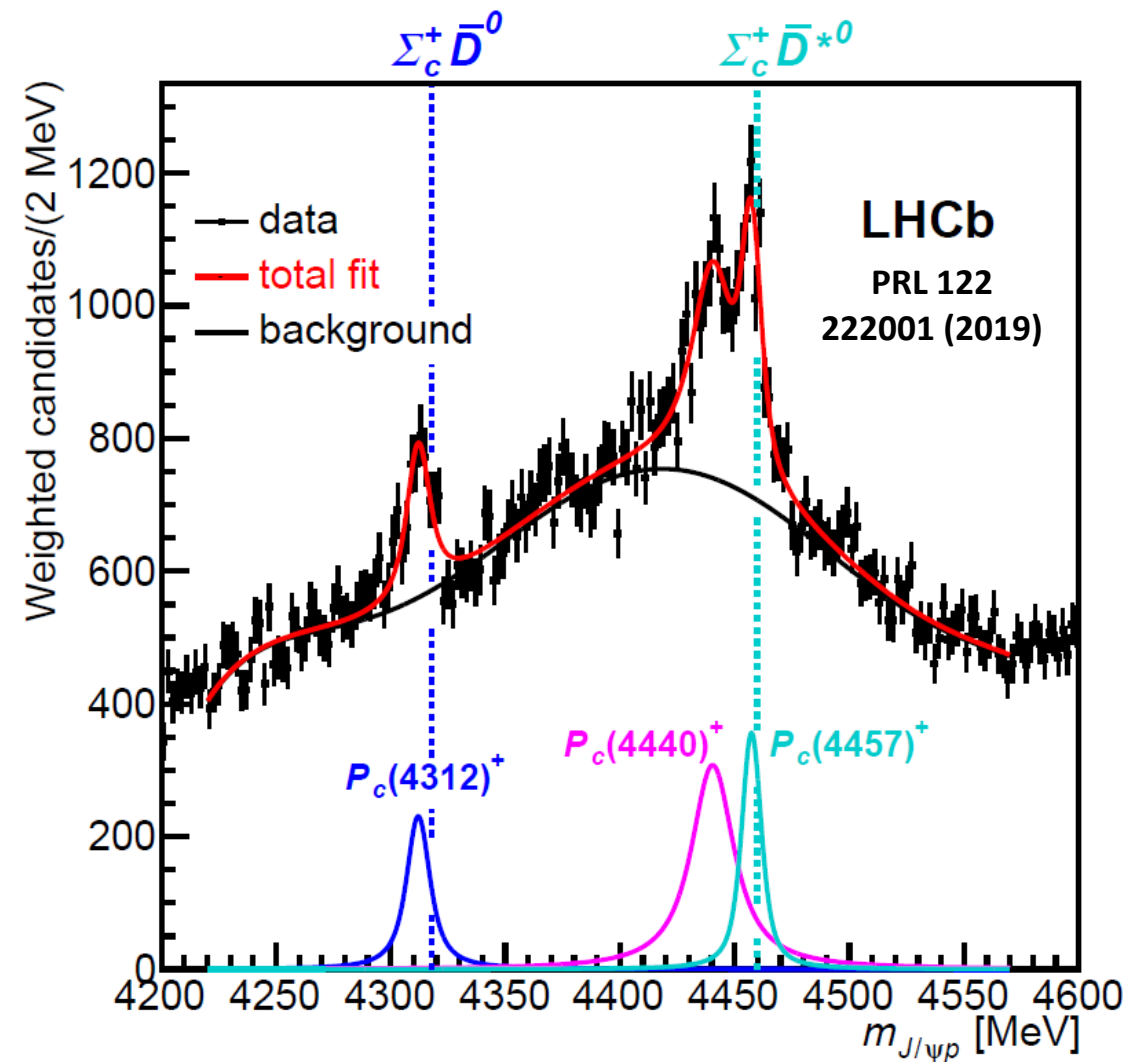
Recent LHCb Discovery – P_c States

Select daughters from the decay

$$\Lambda_b^0 \rightarrow J/\psi p K^-$$

Recent LHCb Discovery – P_c States

Select daughters from the decay

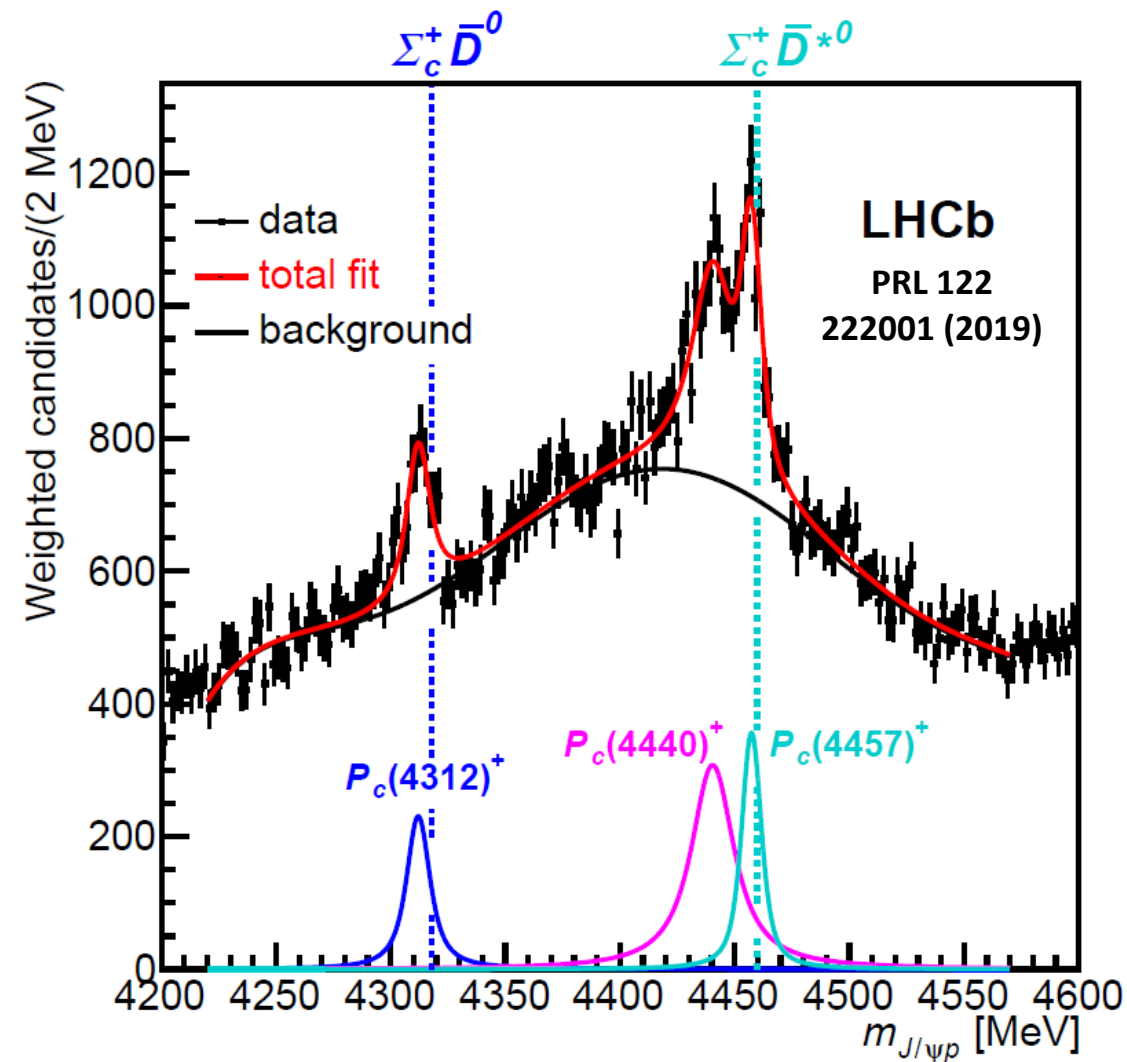
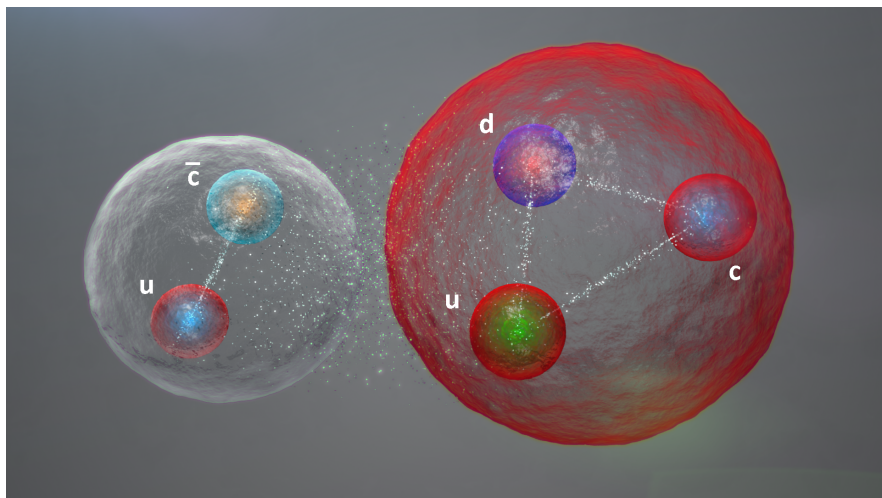


Recent LHCb Discovery – P_c States

Select daughters from the decay



Masses are close to meson+baryon thresholds
– candidate for hadronic molecule



Prompt production has not been studied

Charged Exotic – Z_c State

Select daughters from the decay

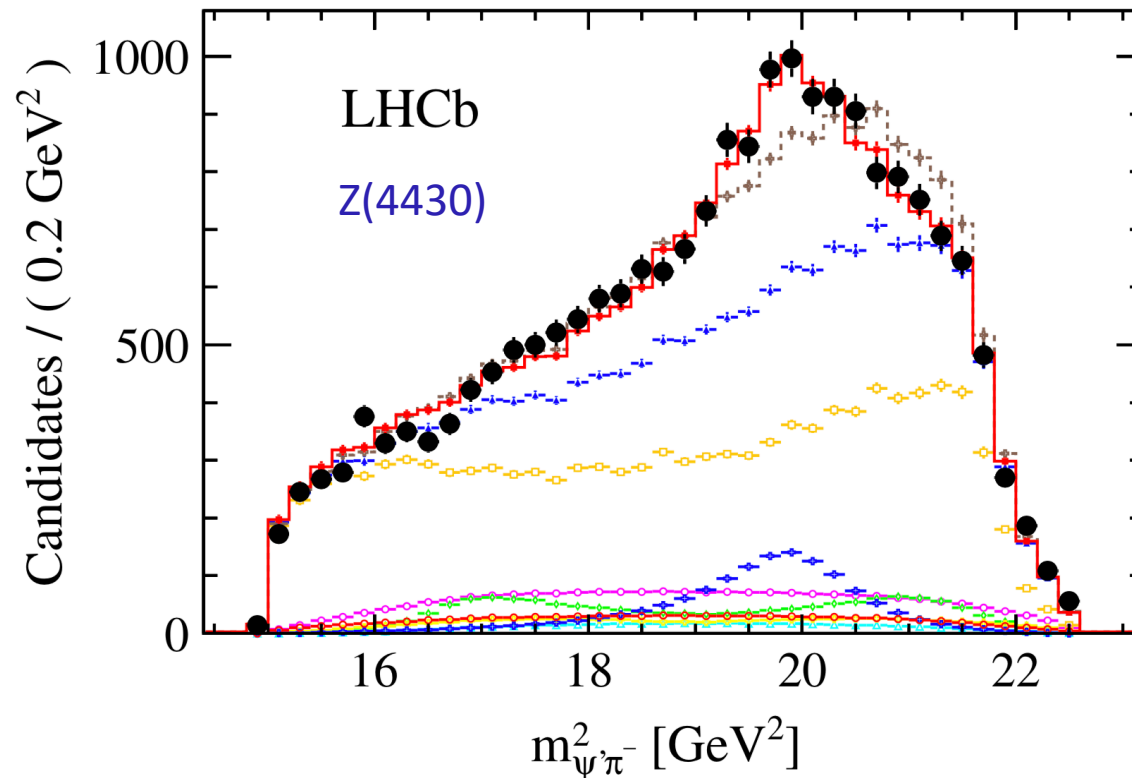
$$B^0 \rightarrow \psi' K^+ \pi^-$$

Charged Exotic – Z_c State

Select daughters from the decay

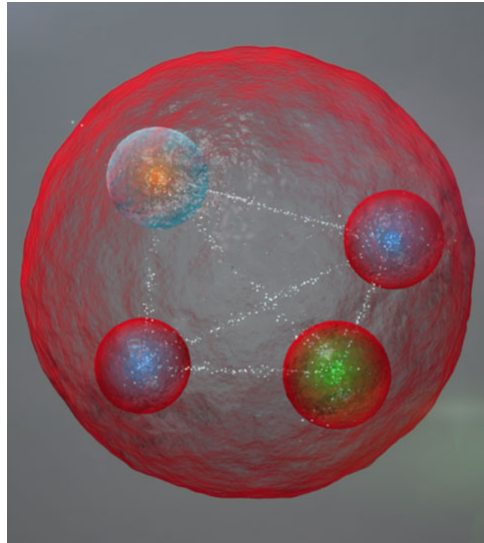


Phys Rev Lett 112 222002 (2014)



Charged Exotic – Z_c State

Select daughters from the decay

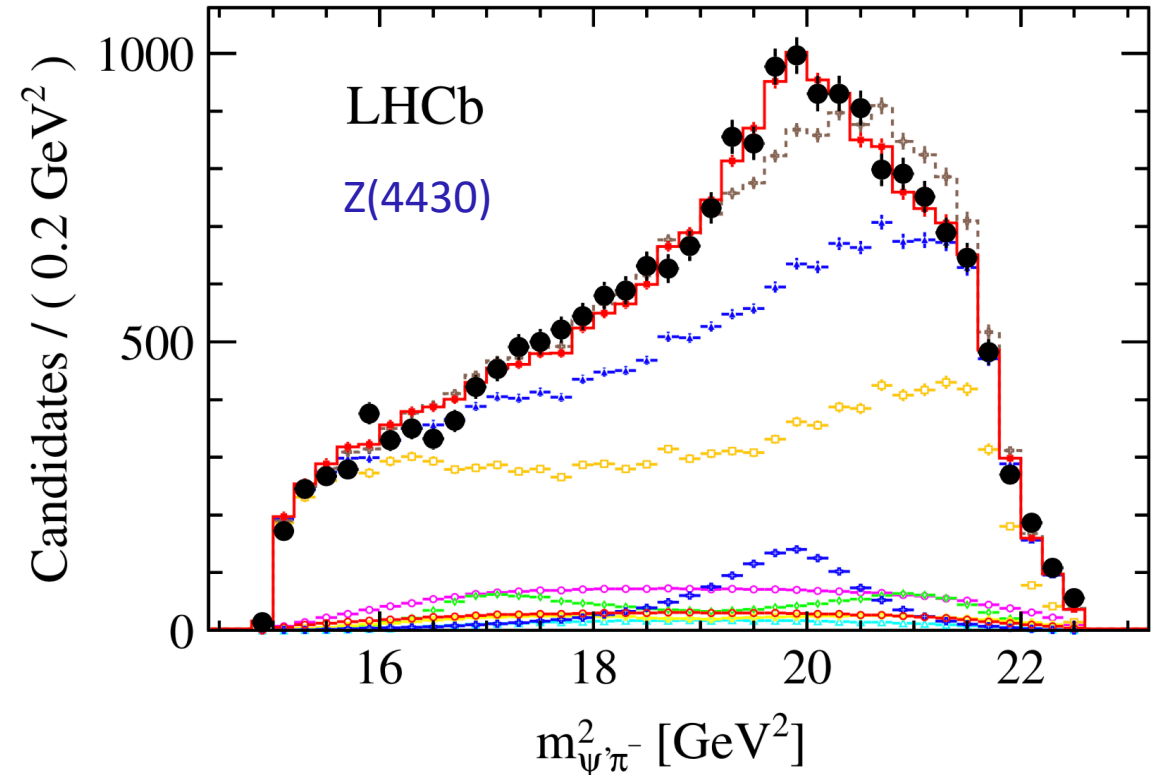


Minimal 4 quark content: $c\bar{c}q\bar{q}$

Masses are NOT close to any hadron+hadron threshold – candidate for compact tetraquark

Prompt production has not been studied

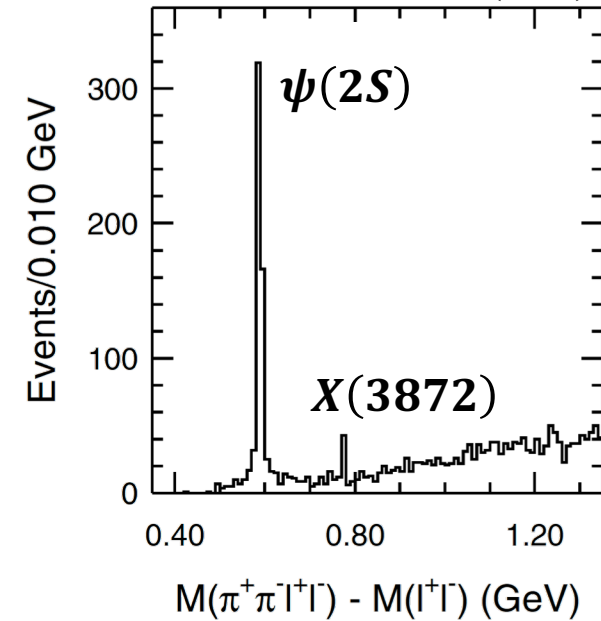
Phys Rev Lett 112 222002 (2014)



X(3872) - a puzzle

Recently renamed
 $\chi_{c1}(3872)$ by PDG

Belle Collaboration
PRL 91 262001 (2003)



The first exotic hadron – discovered in $J/\psi \pi^+ \pi^-$ mass spectrum from B decays by Belle in 2003

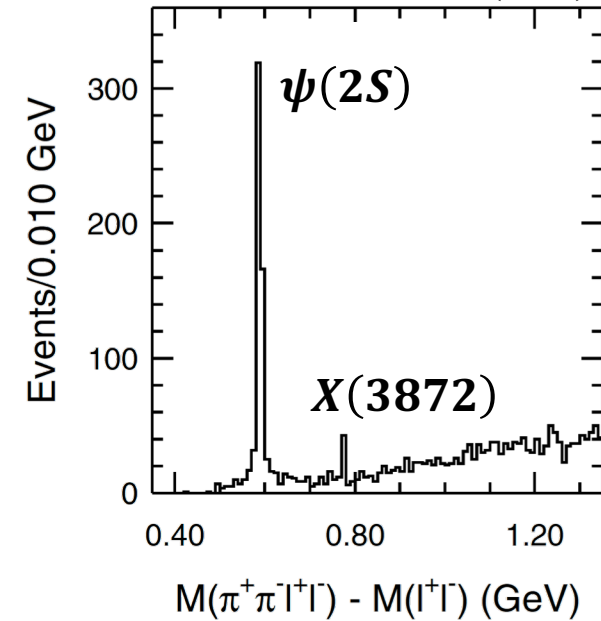
LHCb measured quantum numbers (PRL 110 222001 2013)

- Incompatible with expected charmonium states

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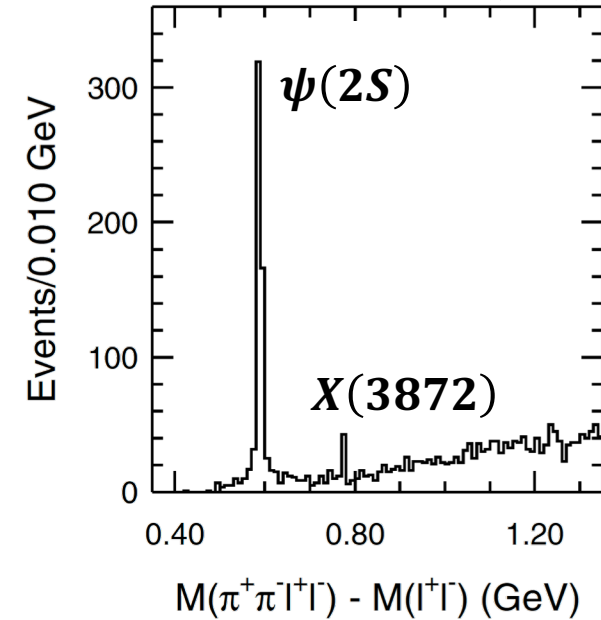
Mass is consistent with sum of D^0 and \bar{D}^{*0} masses:

$$M_{\chi_{c1}(3872)} - (M_{D^0} + M_{\bar{D}^{*0}}) = 0.01 \pm 0.27 \text{ MeV}$$

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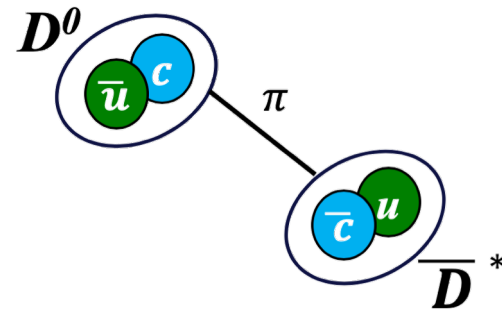
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$D^0\bar{D}^{*0}$ Molecule

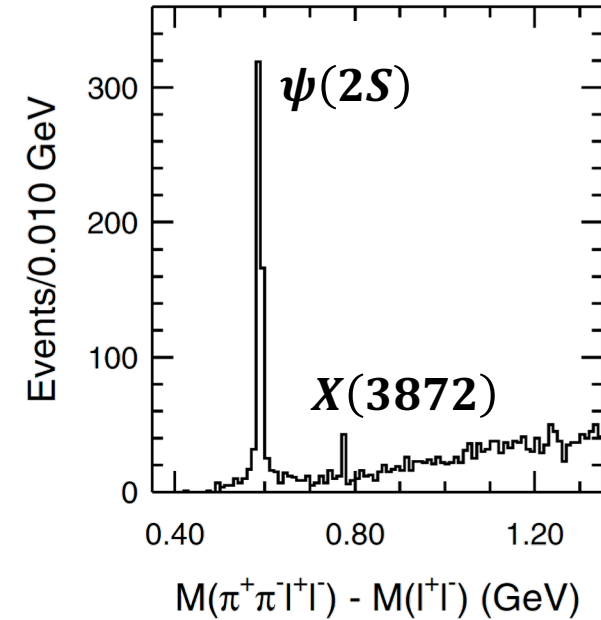


VERY small binding energy
VERY large radius, ~ 7 fm

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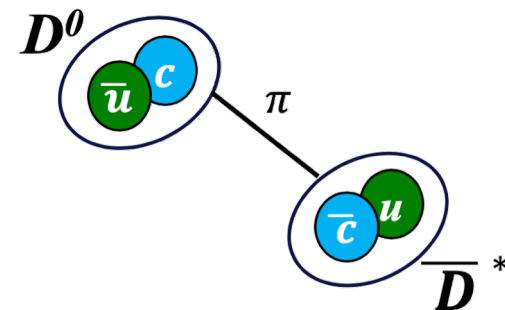
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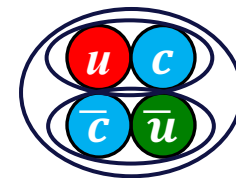
Large prompt production fraction ($\sim 80\%$) – inconsistent with D recombination in pp^*

$D^0 \bar{D}^{*0}$ Molecule



VERY small binding energy
VERY large radius, ~ 7 fm

Compact tetraquark



Tightly bound via color exchange between diquarks
Small radius, ~ 1 fm

*Tension in theoretical literature:

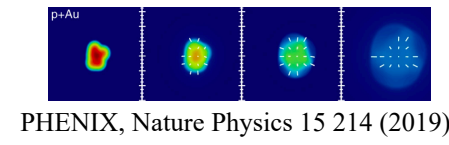
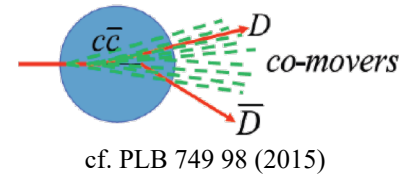
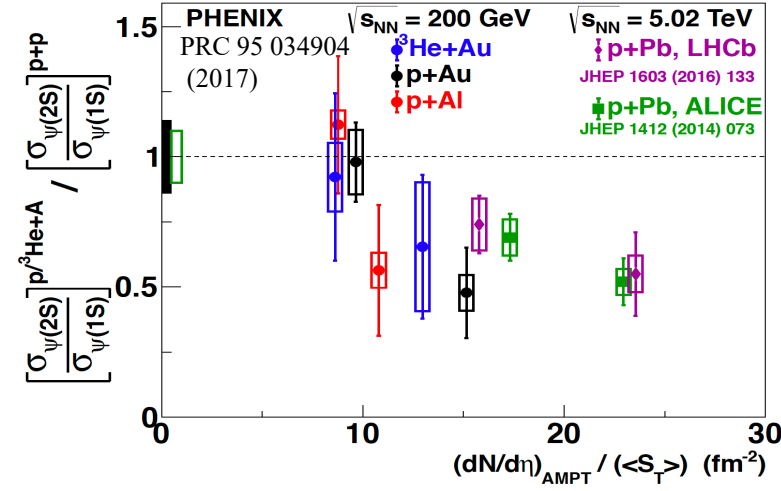
c.f. Bignamini, Grinstein et al
PRL 103 162001 (2009)

VS

Artoisenet, Braaten
PRD 81 114018 (2010)

Effects of Binding Energy

- Suppression of weakly-bound quarkonia states has been studied for decades in pA collisions
 - Ratios of $\psi(2S)/J/\psi$ and $\Upsilon(2S,3S)/\Upsilon(1S)$

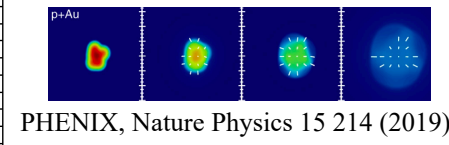
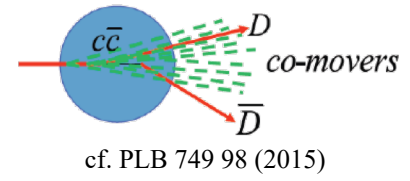
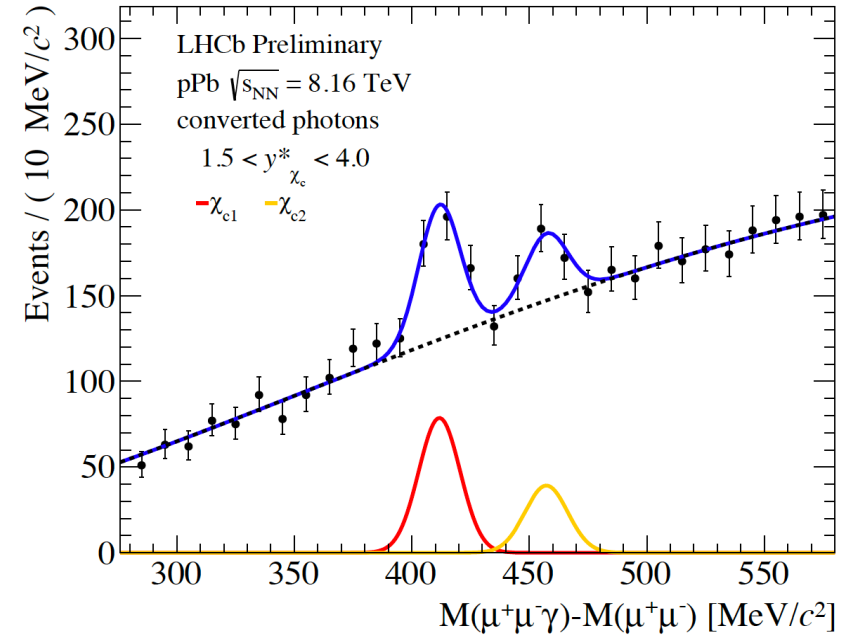


state	η_c	J/ψ	χ_{c0}	χ_{c1}	χ_{c2}	ψ'
mass [GeV]	2.98	3.10	3.42	3.51	3.56	3.69
ΔE [GeV]	0.75	0.64	0.32	0.22	0.18	0.05

Satz, J. Phys. G 32 (3) 2006

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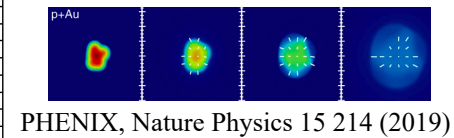
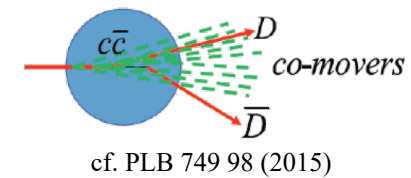
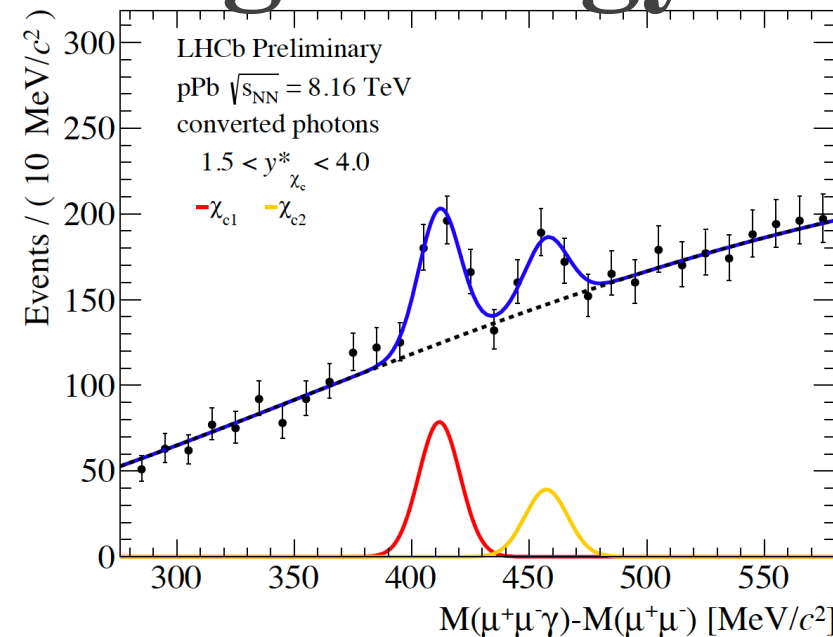


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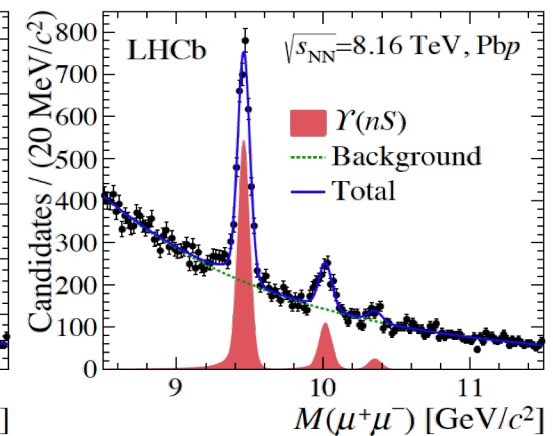
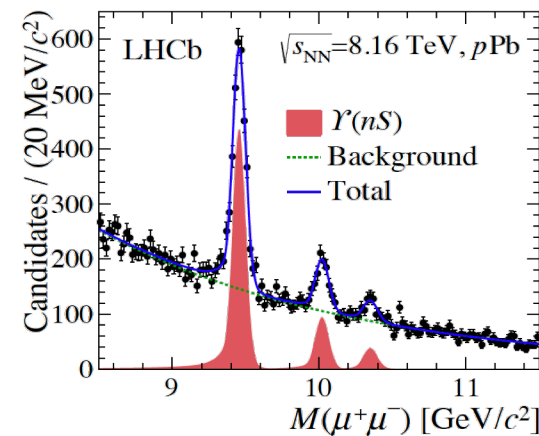
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- In general, final state effects are required to explain difference in suppression between states
- Prevalent in regions with high particle multiplicity



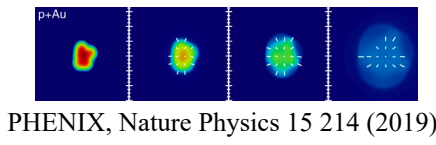
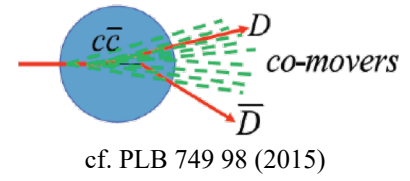
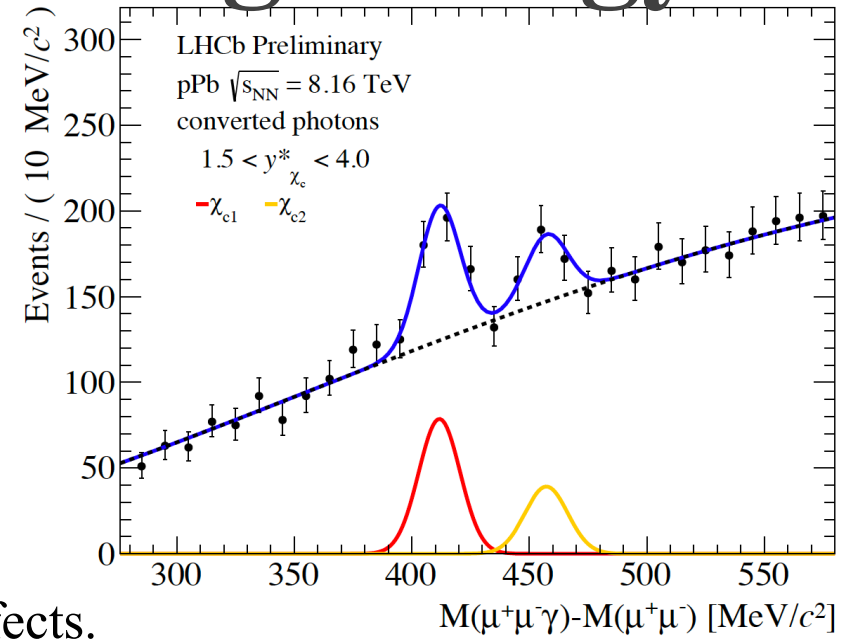
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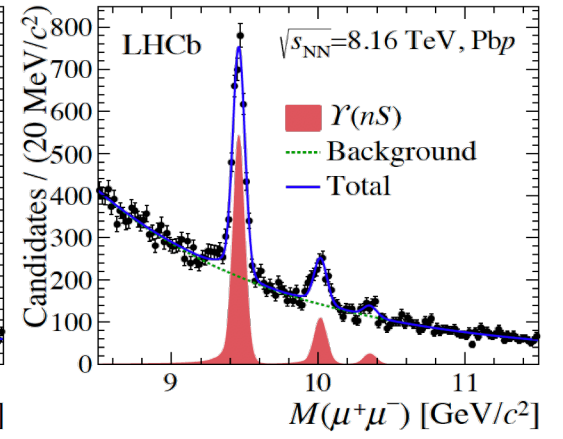
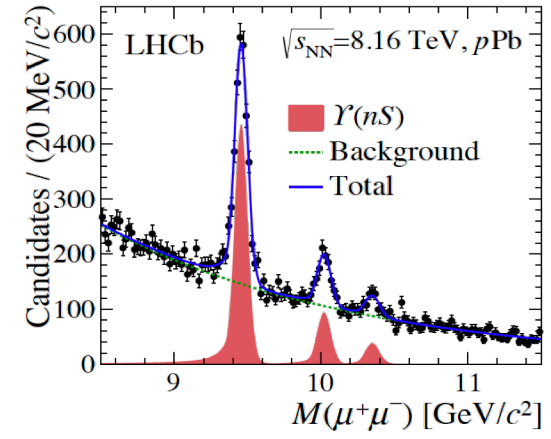
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- Prevalent in regions with high particle multiplicity
- Weakly bound hadronic molecules may show similar effects.



state	η_c	J/ψ	χ_{c0}	χ_{c1}	χ_{c2}	ψ'	$D\bar{D}^*$ Molecule
mass [GeV]	2.98	3.10	3.42	3.51	3.56	3.69	X(3872)
ΔE [GeV]	0.75	0.64	0.32	0.22	0.18	0.05	0.00001 ± 0.00027

Satz, J. Phys. G 32 (3) 2006

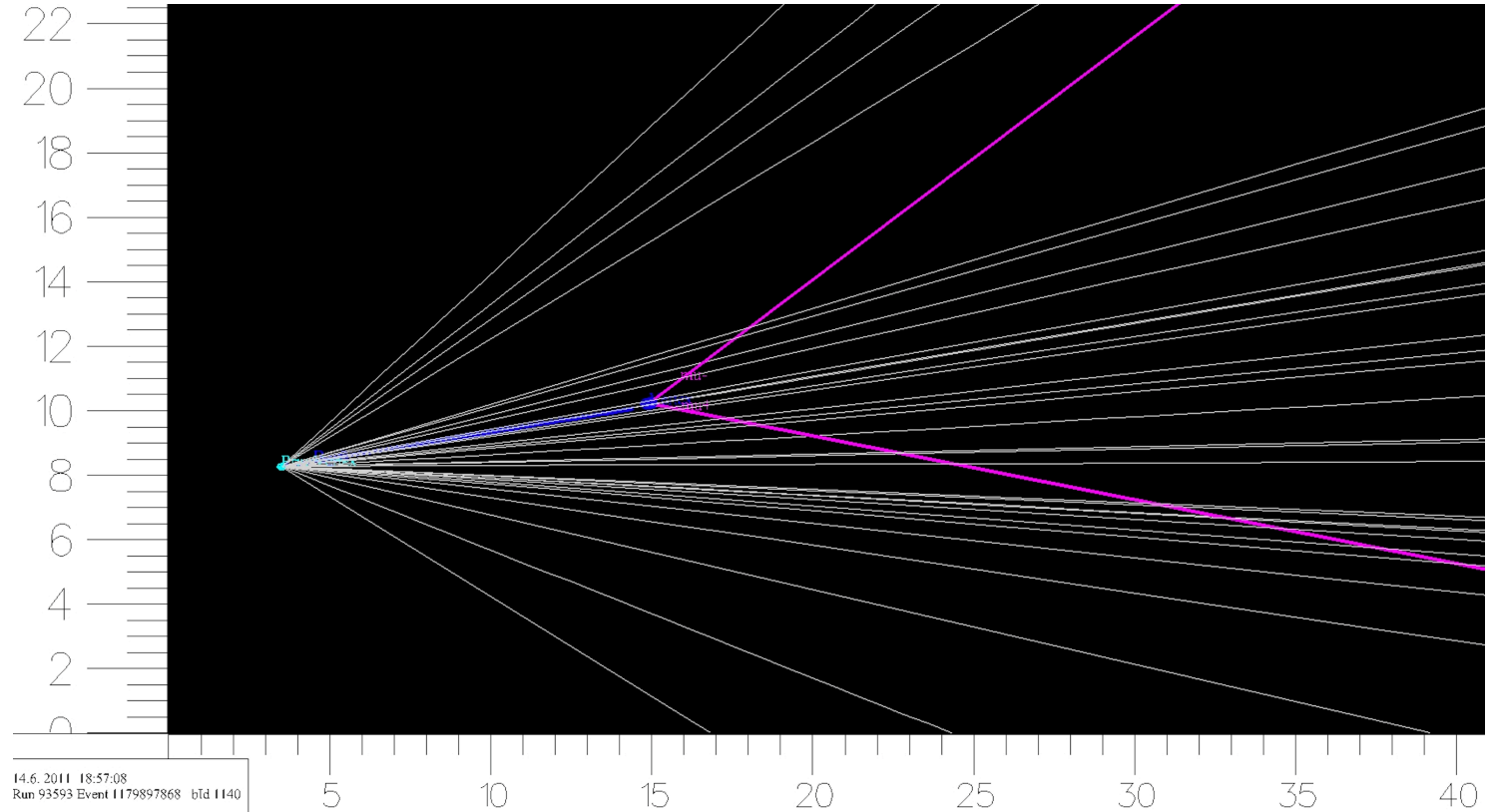


Probing X(3872) structure via interactions with the underlying event

Prompt production:

- X(3872) produced at collision vertex can be subject to further interactions with co-moving particles (medium?) produced in the event
- Potentially subject to breakup effects

Event display of $B_S^0 \rightarrow \mu^+ \mu^-$ candidate, PRL 118 191801 (2017)



Probing X(3872) structure via interactions with the underlying event

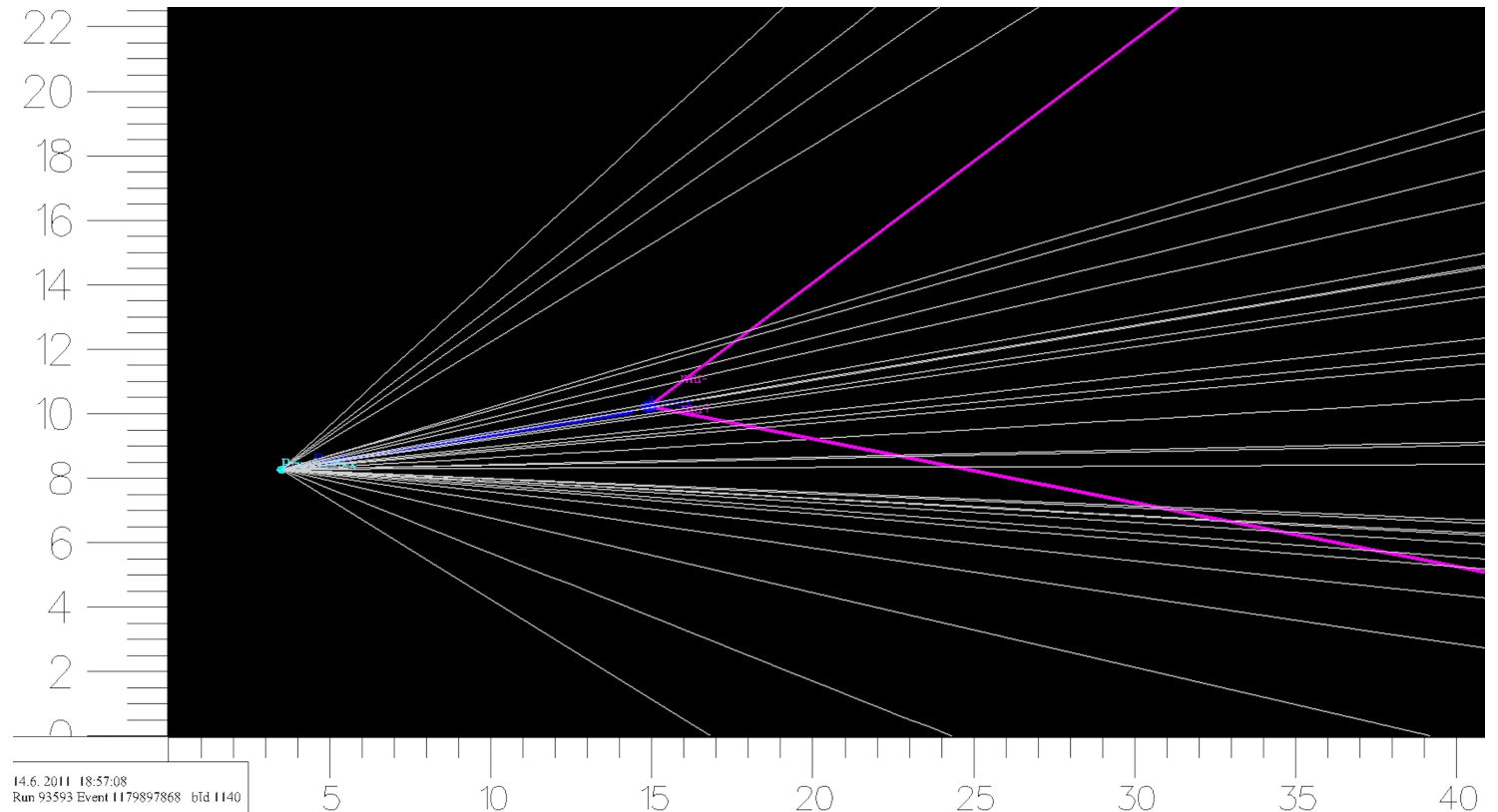
Prompt production:

- X(3872) produced at collision vertex can be subject to further interactions with co-moving particles (medium?) produced in the event
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Production in b -decays:

- Hadrons containing b travel down the beampipe and decay away from the primary vertex and decay in vacuum
- X(3872) from decays not subject to further interactions
- Control sample

Event display of $B_S^0 \rightarrow \mu^+ \mu^-$ candidate, PRL 118 191801 (2017)



X(3872) selection

LHCb-CONF-2019-005

Reconstruct the $\mu^+ \mu^- \pi^+ \pi^-$ final state from the decays:

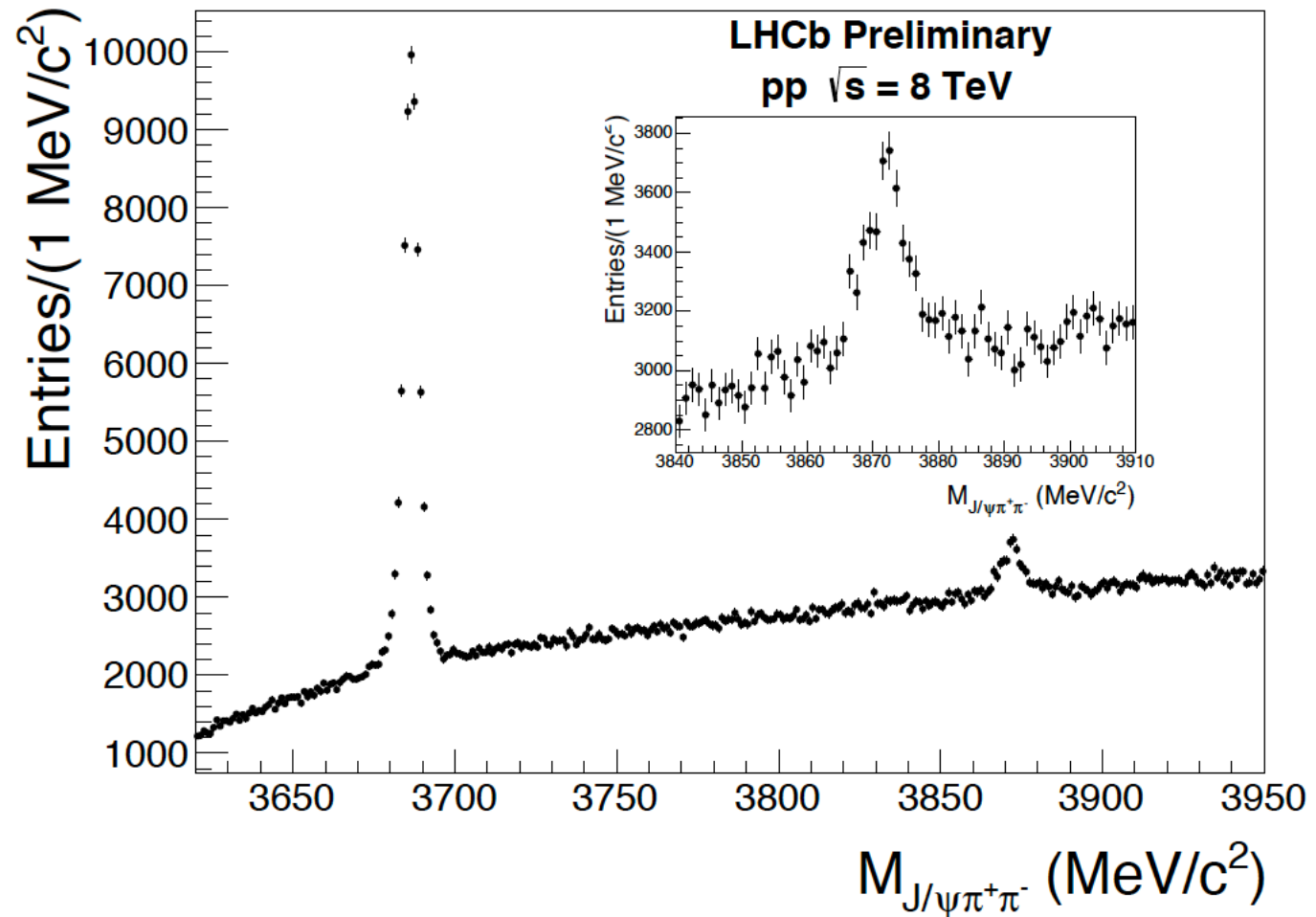
$$X(3872) \rightarrow J/\psi(\rightarrow \mu^+ \mu^-) \rho(\rightarrow \pi^+ \pi^-)$$

$$\psi(2S) \rightarrow J/\psi(\rightarrow \mu^+ \mu^-) \pi^+ \pi^-$$

Select J/ψ from dimuons, combine with two identified pions. Perform kinematic refit, constraining J/ψ mass to known value and all four tracks to identical vertex.

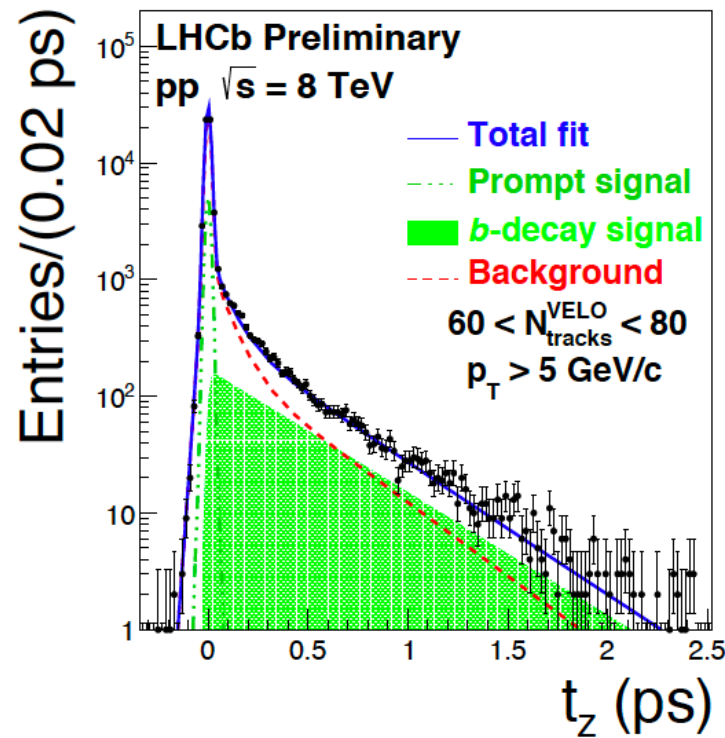
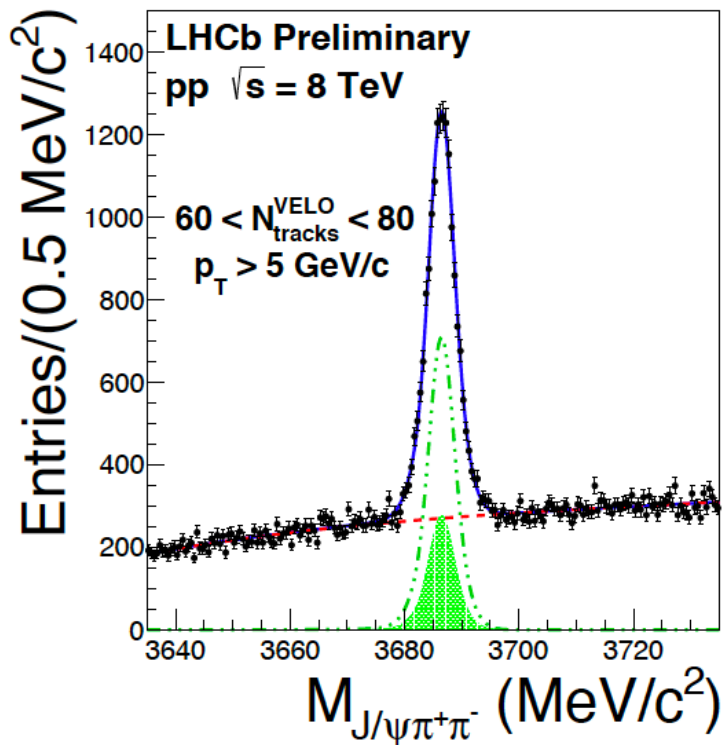
Direct comparison between conventional charmonium $\psi(2S)$ and exotic $X(3872)$ via ratio of cross sections:

$$\frac{\sigma_{\chi_{c1}(3872)}}{\sigma_{\psi(2S)}} \times \frac{\mathcal{B}[\chi_{c1}(3872) \rightarrow J/\psi \pi^+ \pi^-]}{\mathcal{B}[\psi(2S) \rightarrow J/\psi \pi^+ \pi^-]}$$



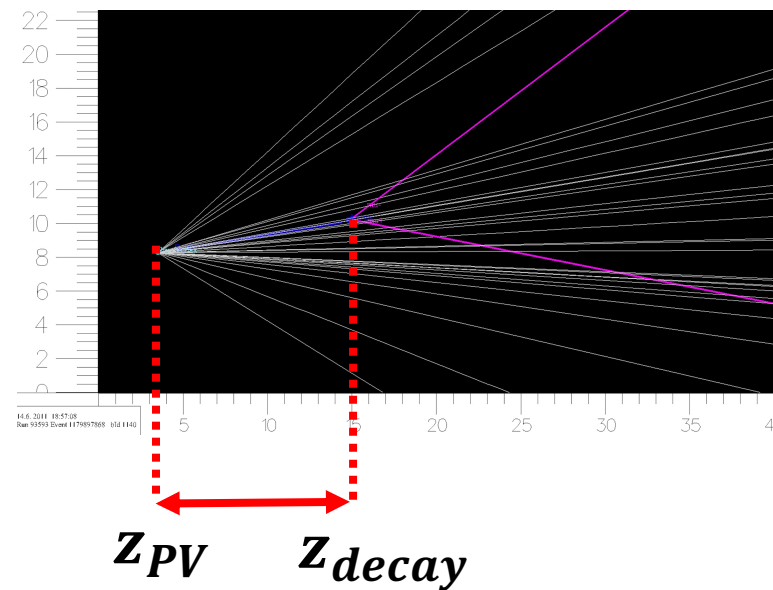
Prompt / b -decay separation

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Simultaneous fit to invariant mass and pseudo proper time spectrum:

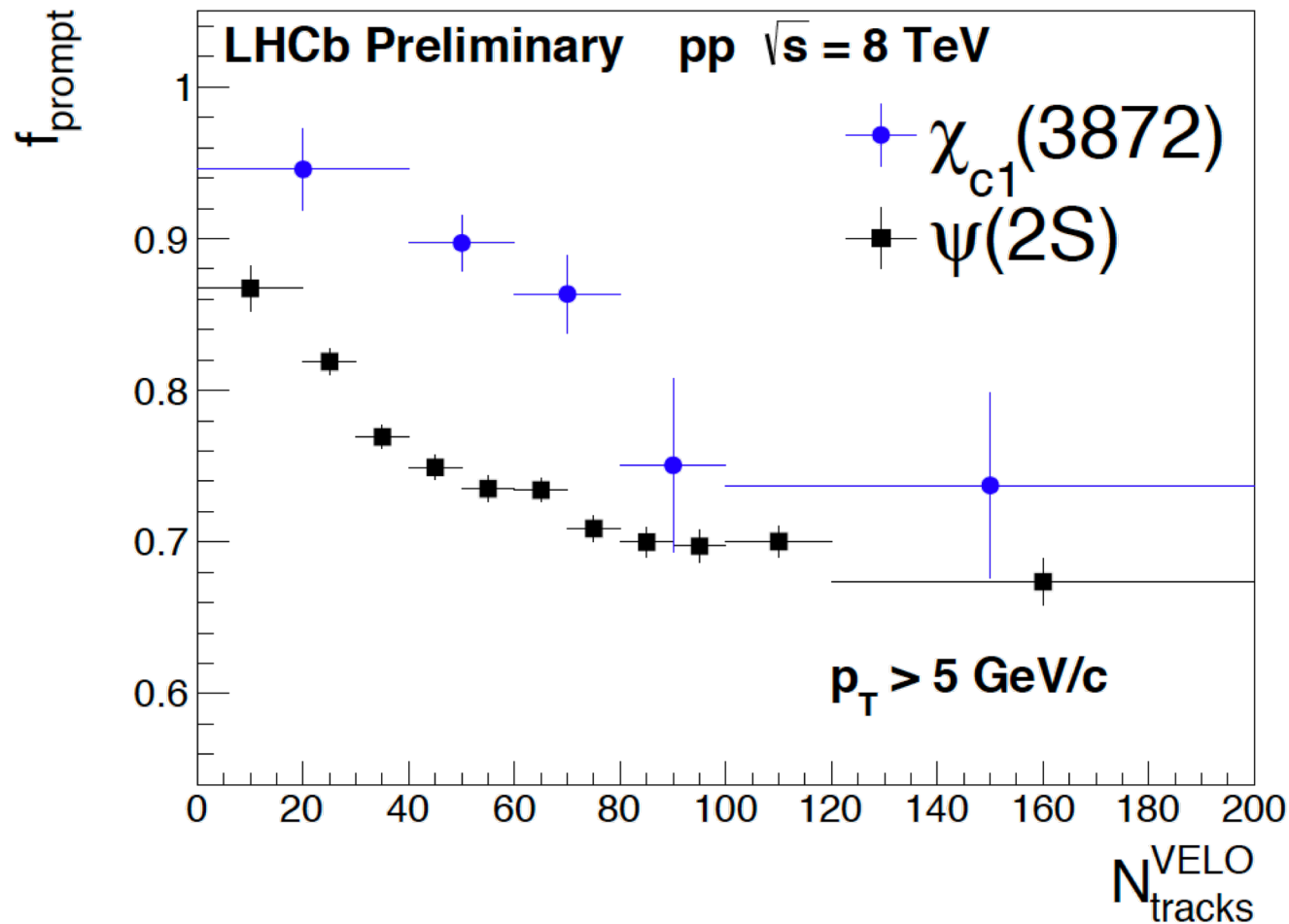
$$t_z = \frac{z_{\text{decay}} - z_{PV}}{p_z} M$$



Fit to mass constrains S/B while fit to t_z constrains prompt fraction

Prompt fraction

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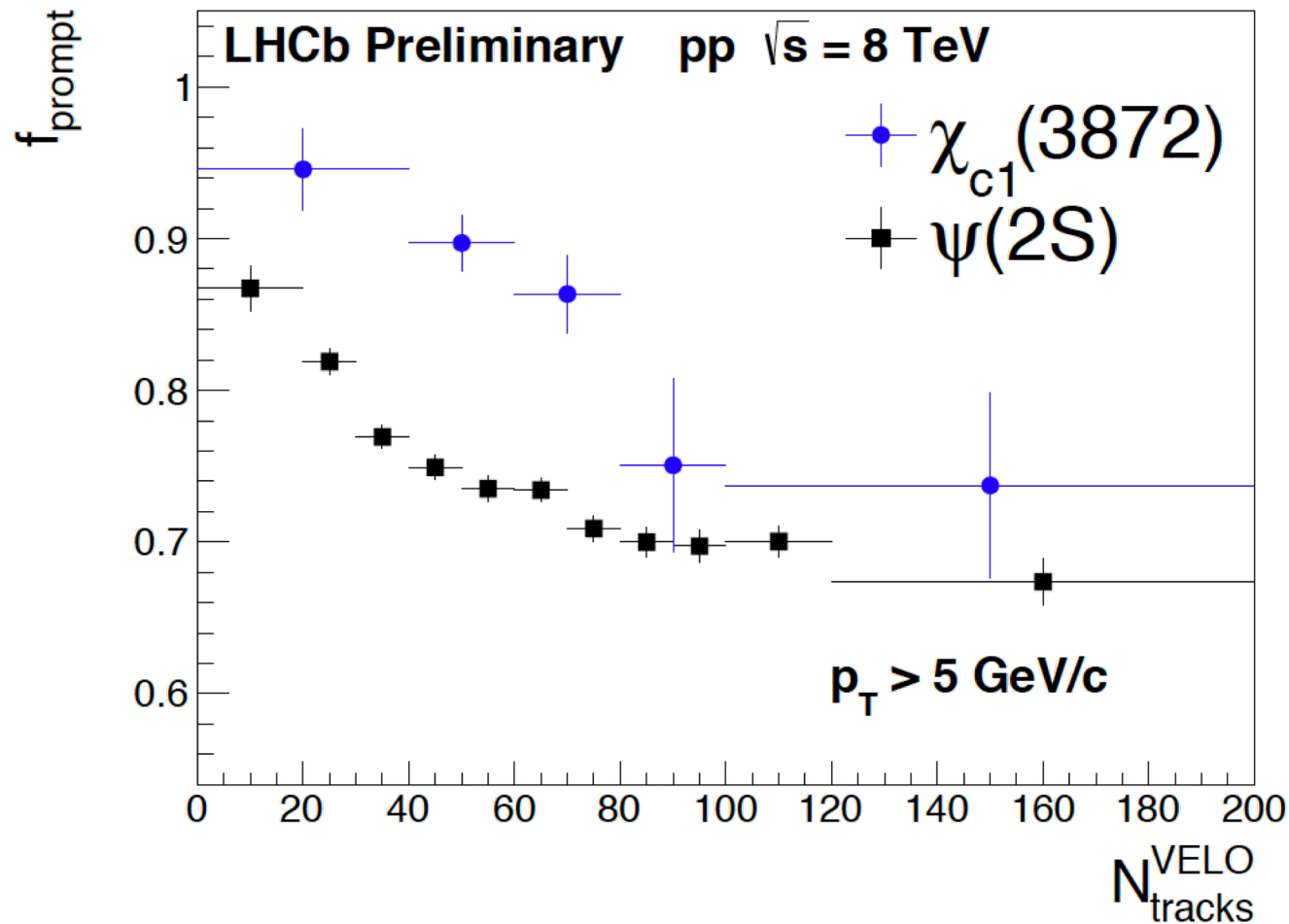


$$f_{prompt} = \frac{N_{prompt}}{N_{prompt} + N_{b-decay}}$$

- Significant decrease in prompt fraction of both $X(3872)$ and $\psi(2S)$ as event activity increases:

Prompt fraction

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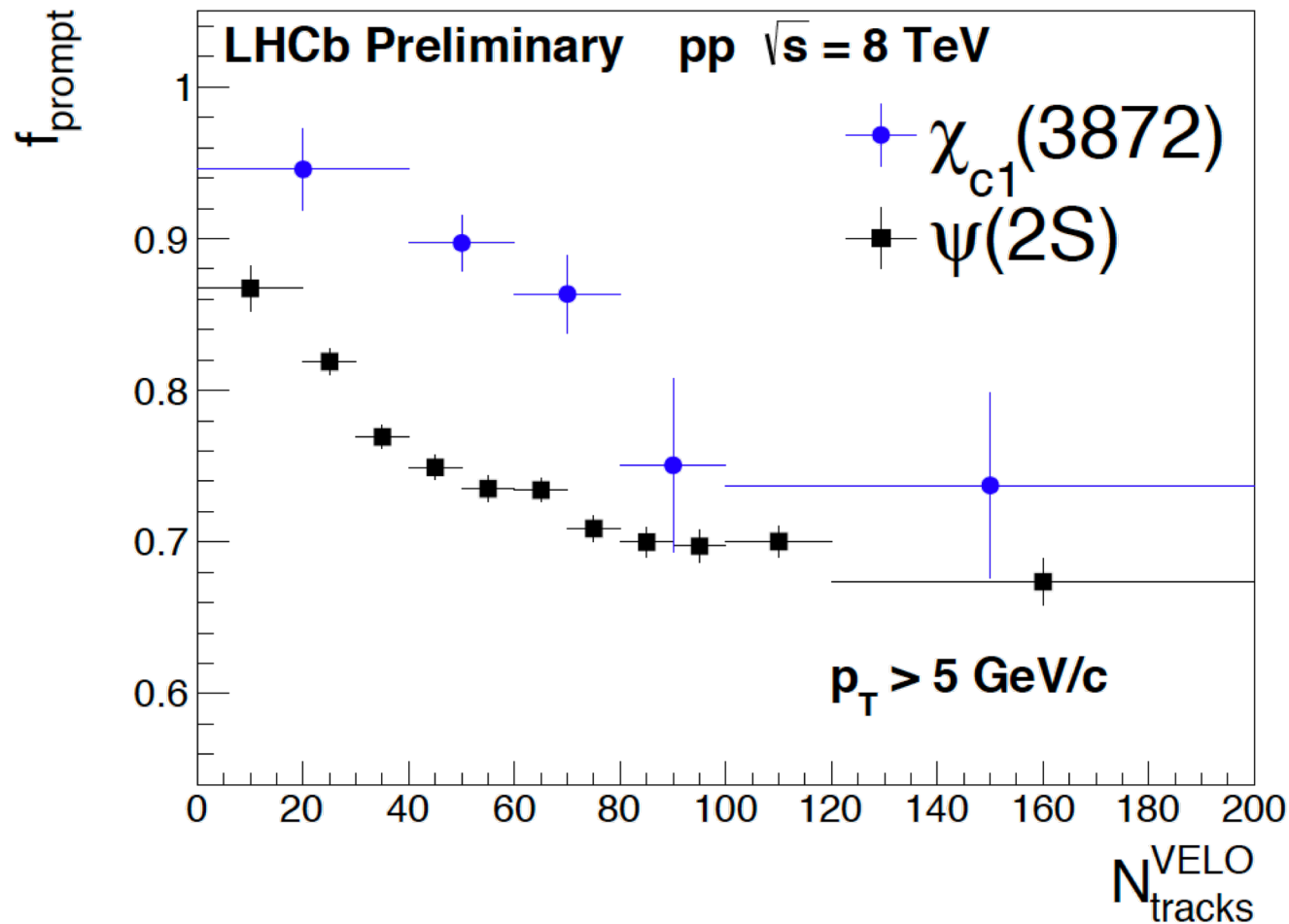


$$f_{\text{prompt}} = \frac{N_{\text{prompt}}}{N_{\text{prompt}} + N_{b\text{-decay}}}$$

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- Events with $b\bar{b}$ production naturally have higher multiplicity, due to fragmentation and decays
 - OPAL, PLB 550 33 (2002)

Prompt fraction

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$$f_{\text{prompt}} = \frac{N_{\text{prompt}}}{N_{\text{prompt}} + N_{b\text{-decay}}}$$

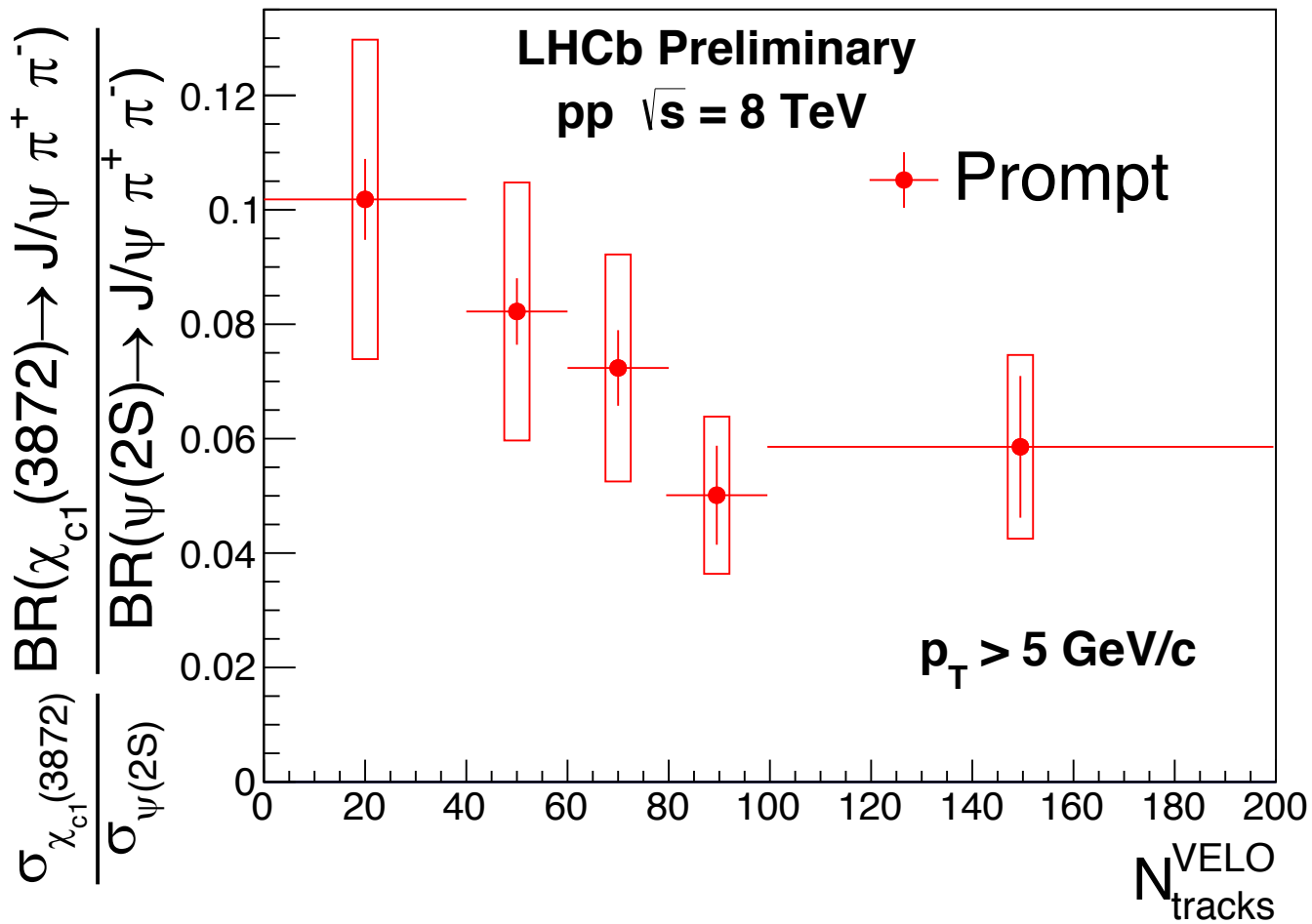
- Significant decrease in prompt fraction of both $X(3872)$ and $\psi(2S)$ as event activity increases:
- Events with $b\bar{b}$ production naturally have higher multiplicity, due to fragmentation and decays
 - OPAL, PLB 550 33 (2002)
- Formation of prompt $X(3872)$ and $\psi(2S)$ may be disrupted at the vertex, which cannot affect production via b decays in vacuum.

Ratio of cross sections

$$\frac{\sigma_{\chi_{c1}(3872)}}{\sigma_{\psi(2S)}} \times \frac{\mathcal{B}[\chi_{c1}(3872) \rightarrow J/\psi \pi^+ \pi^-]}{\mathcal{B}[\psi(2S) \rightarrow J/\psi \pi^+ \pi^-]} = \frac{N_{\chi_{c1}(3872)} f_{\text{prompt}}^{\chi_{c1}(3872)}}{N_{\psi(2S)} f_{\text{prompt}}^{\psi(2S)}} \times \frac{\epsilon_{\psi(2S)}}{\epsilon_{\chi_{c1}(3872)}}$$

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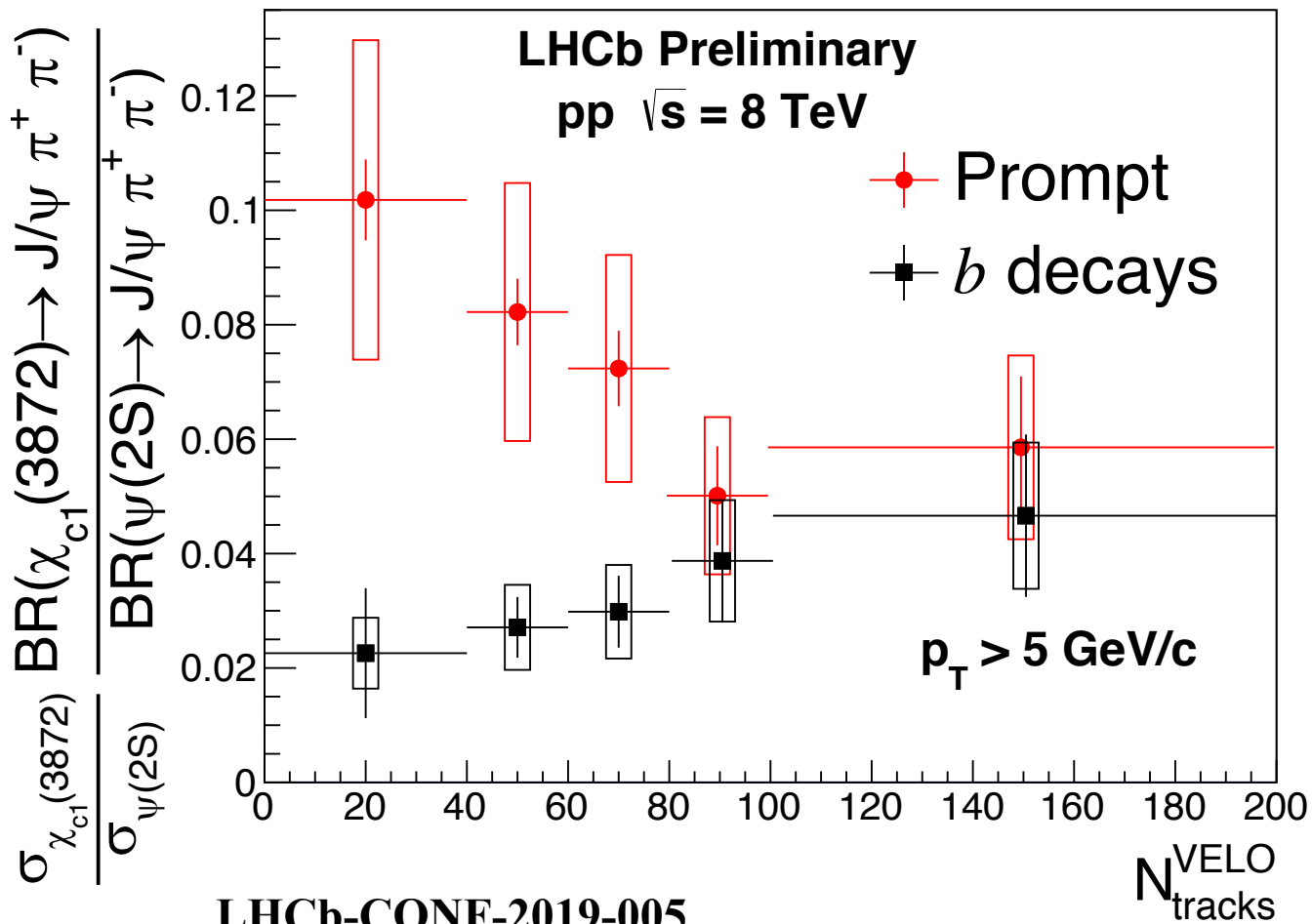


Prompt Component:
Increasing suppression of
 $X(3872)$ production relative to
 $\psi(2S)$ as event activity increases

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Ratio of cross sections

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Prompt Component:
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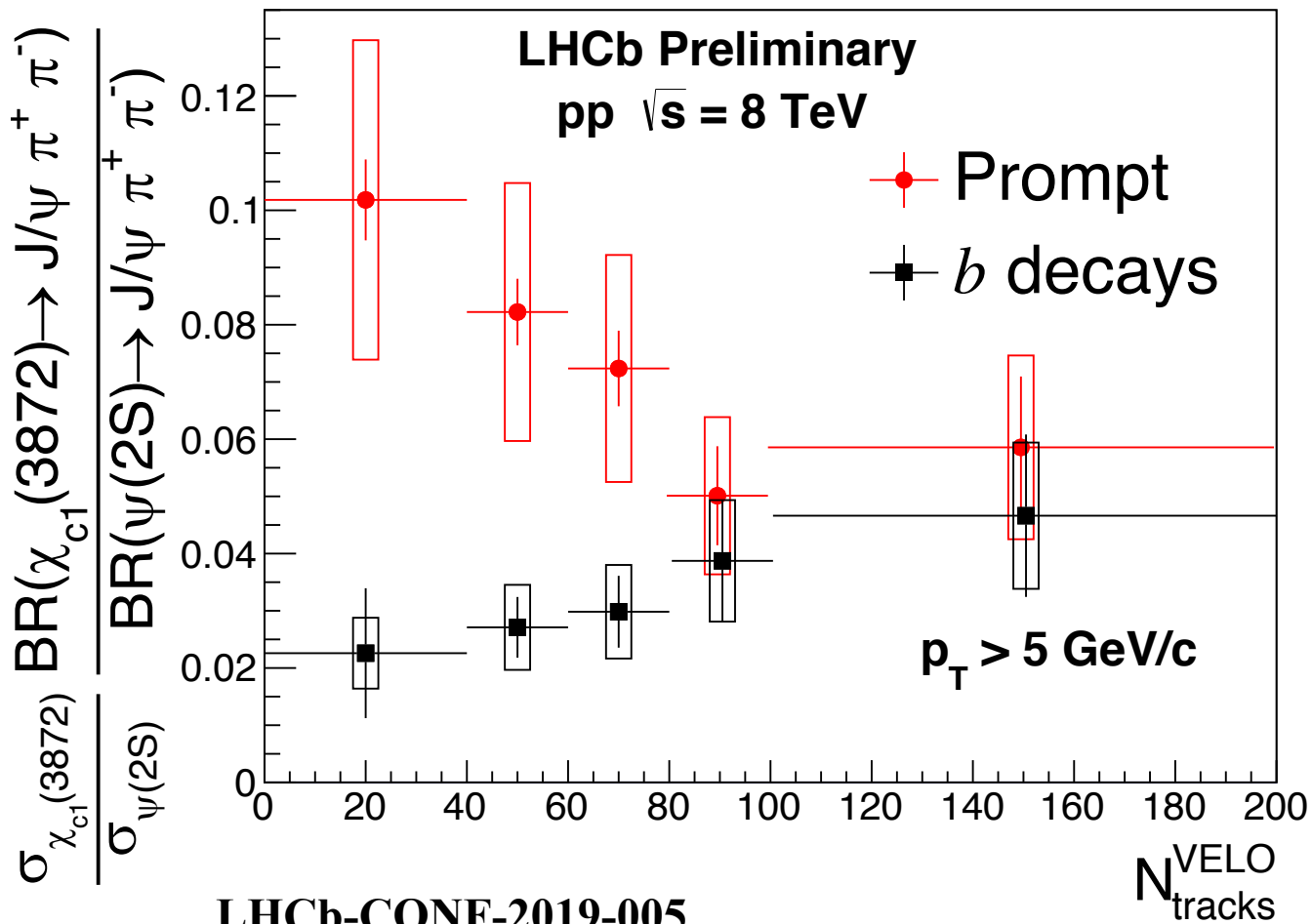
***b*-decay component:**
 No significant change in relative production, as expected for decays in vacuum. Ratio is set by ***b*** decay branching fractions.

Consistent with ATLAS measurement
 $R = 0.0395 \pm 0.0032 \pm 0.0008$ ($p_T > 10$ GeV/c)

JHEP 2017:117 (2017)

Ratio of cross sections

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LHCb-CONF-2019-005

Prompt Component:
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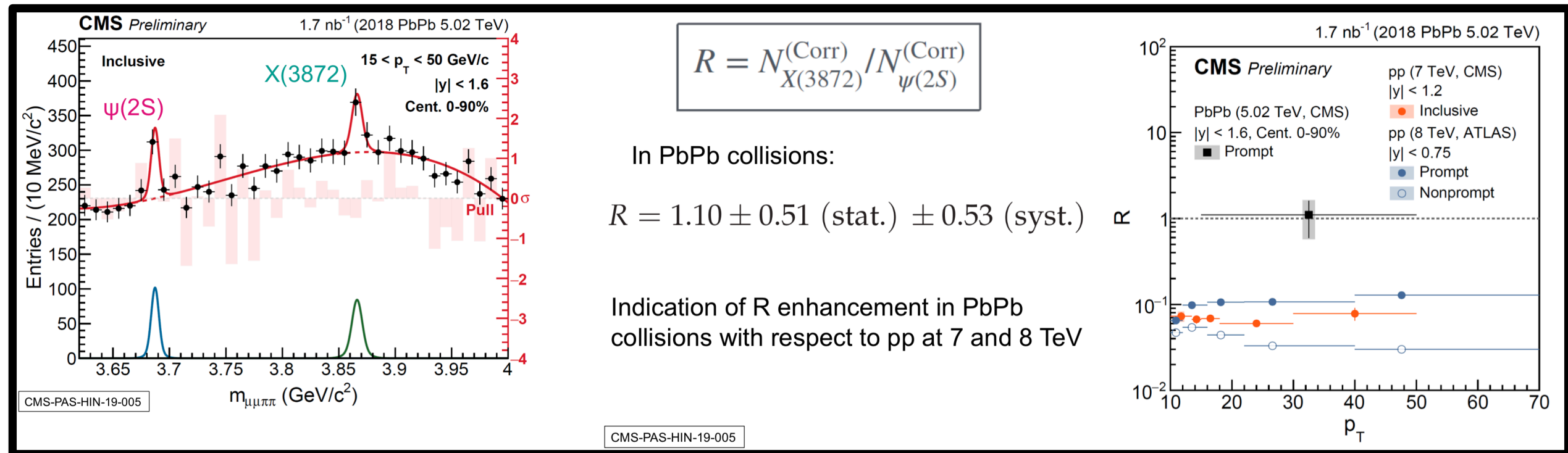
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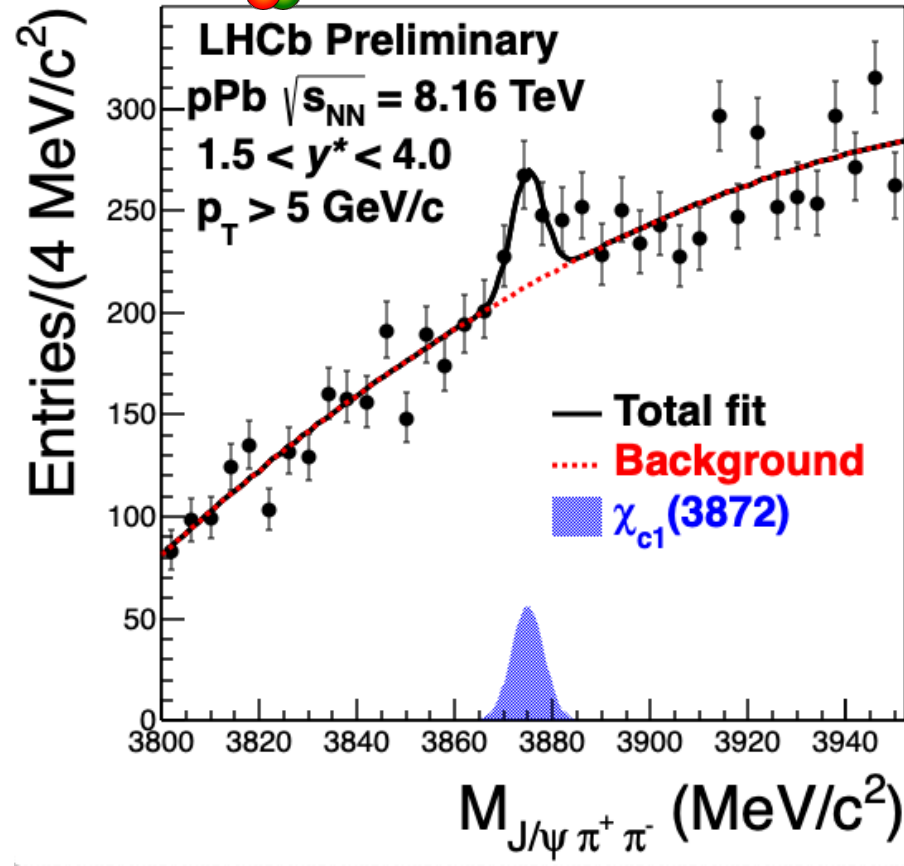
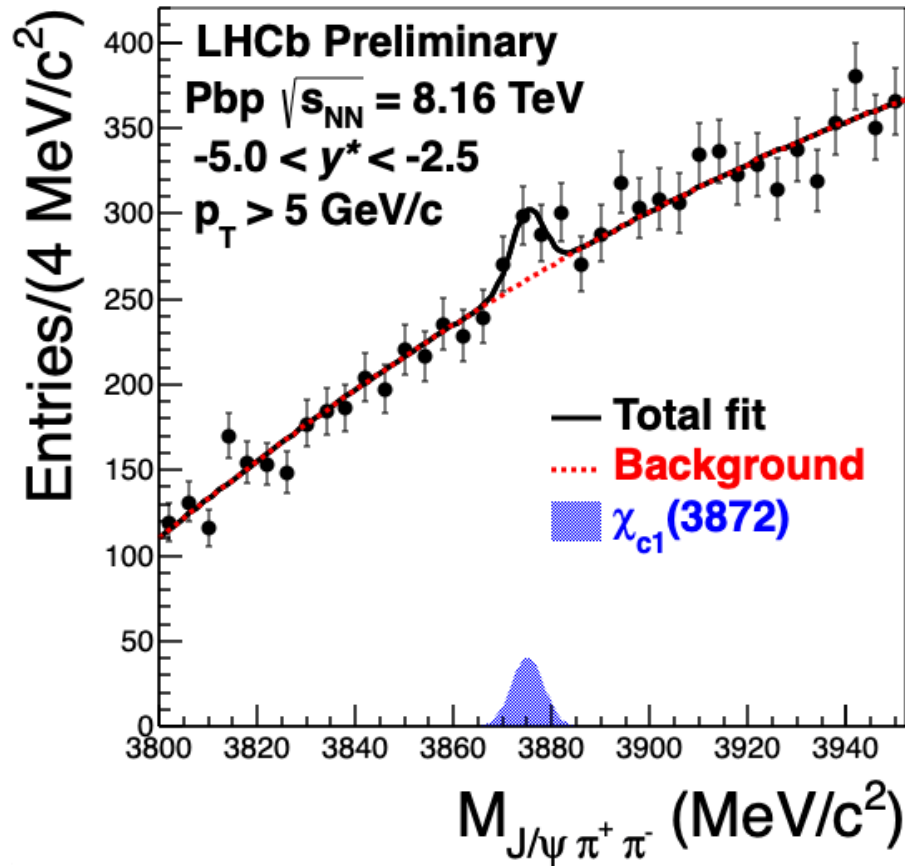
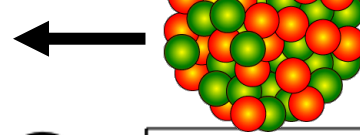
CMS Result – X(3872) in PbPb

See slides from YJ Lee, Quark Matter 2019



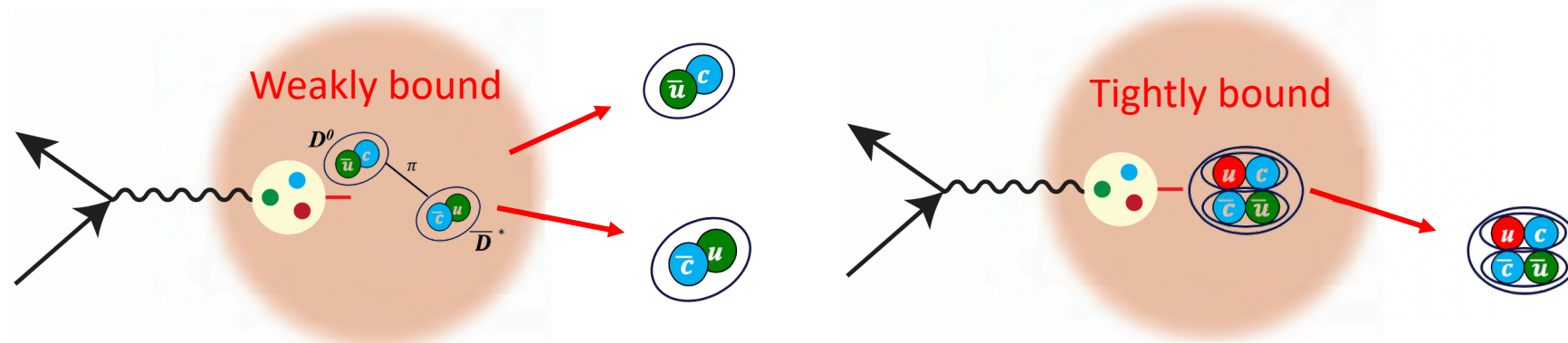
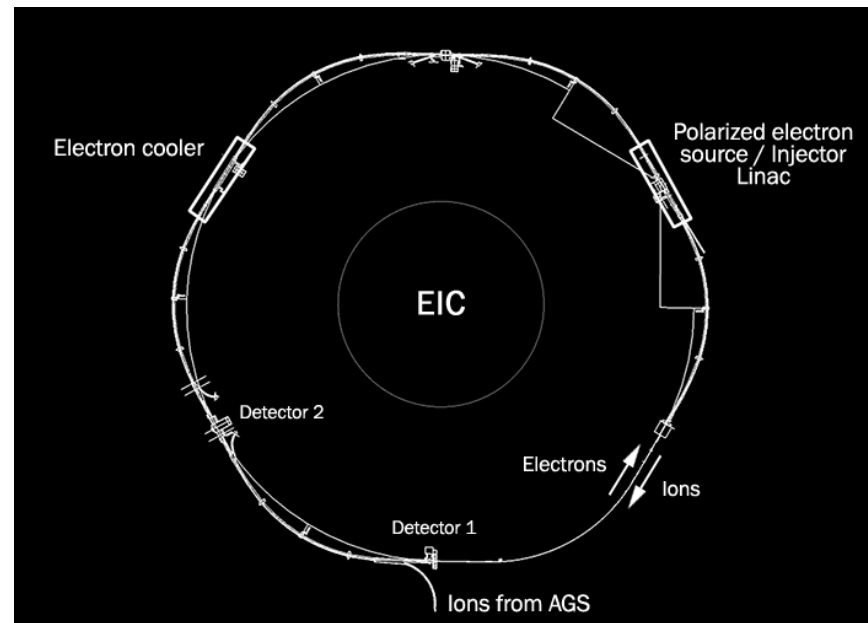
Hint of recombination in PbPb ($p_T > 15 \text{ GeV}$) ??
 Large error bars preclude firm conclusions

X(3872) in pPb collisions



Studies at a future EIC

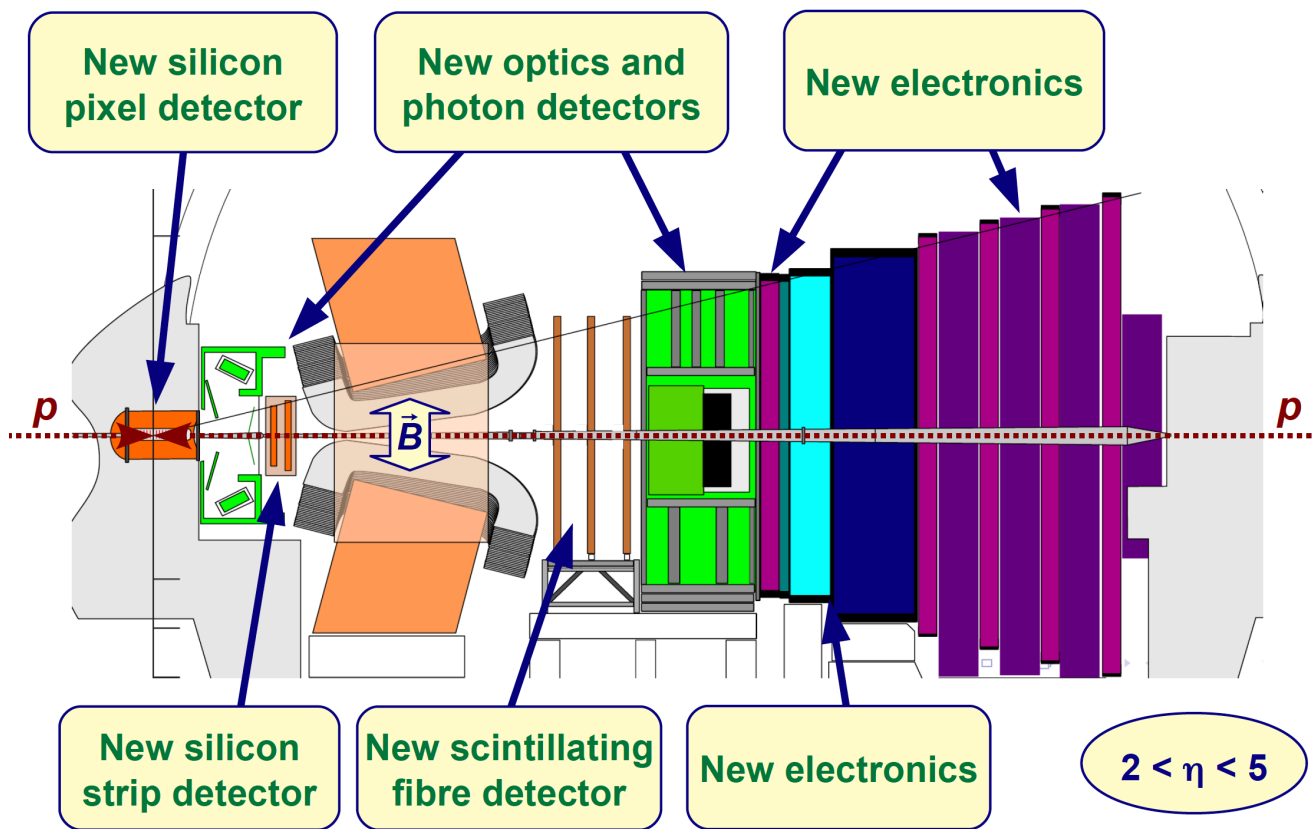
- Electron ion collisions – clean environment for complicated reconstruction
- High luminosity, high rate -> high statistics for rare probes
- Production inside the nucleus exposes exotics to a dense QCD environment – potentially disrupting formation
 - Discrimination between molecular and compact tetraquark pictures
- Requires tracking and reconstruction in forward (nucleus-going) direction – research at LANL is ongoing w/LDRD funds



Outline

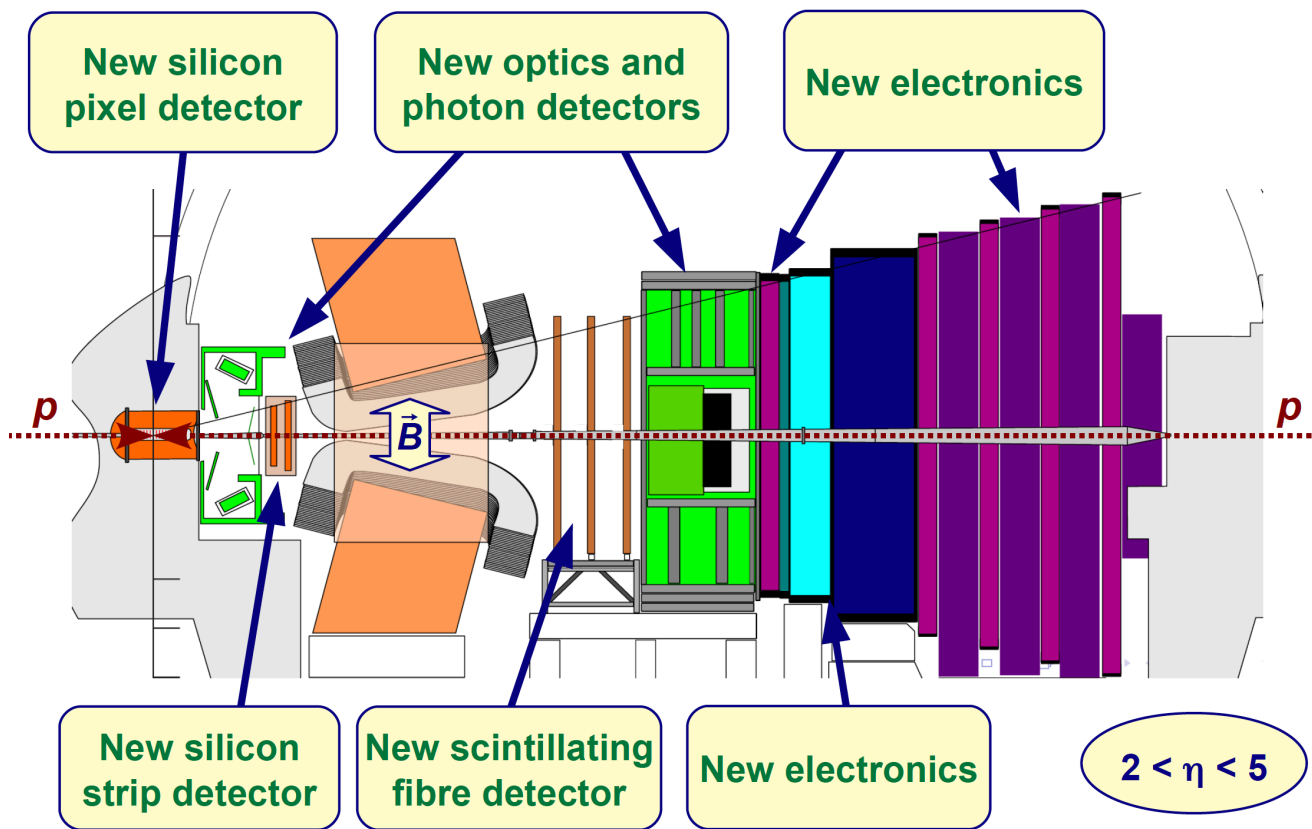
- LHCb Physics program
- Apparatus
- Heavy Ions at LHCb
 - Collider Mode
 - Fixed Target Mode (SMOG)
- Exotics at LHCb
- Upgrades

Massive Upgrades Underway

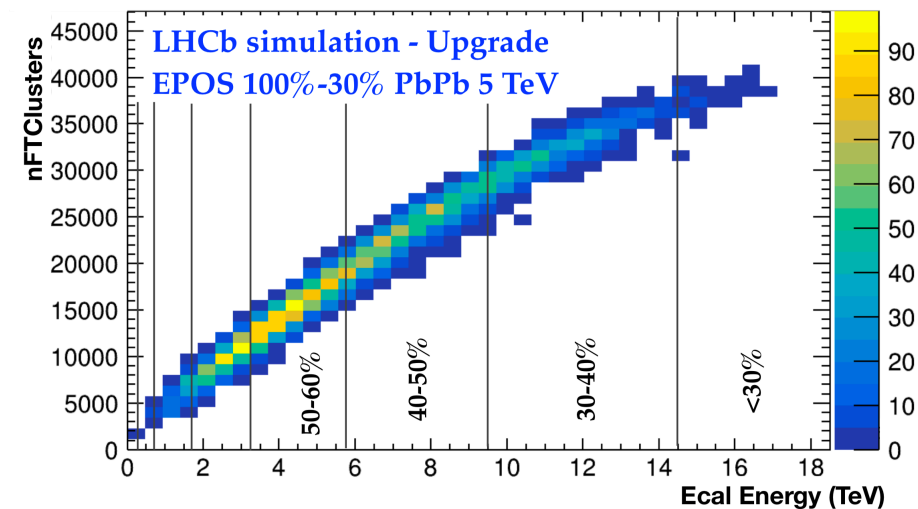
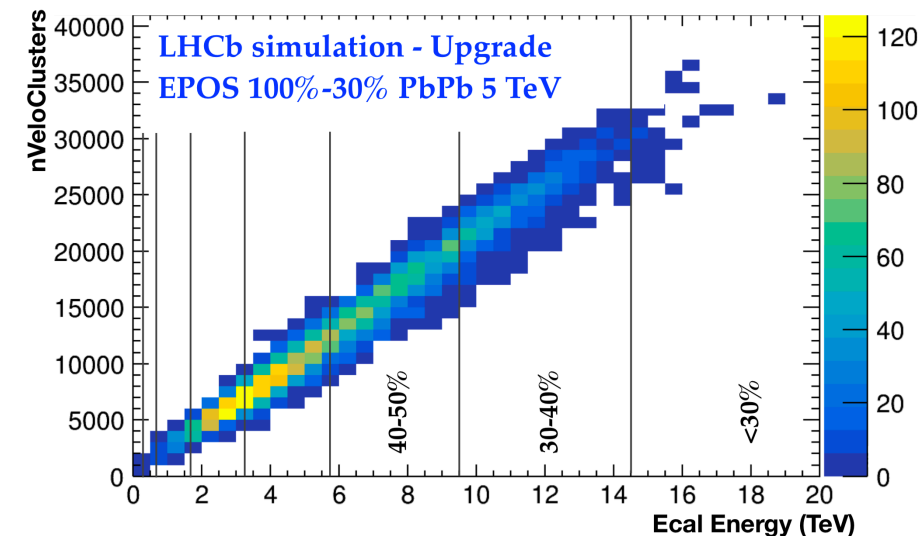


**Trigger-less readout of full
40MHz LHC collision rate**

Massive Upgrades Underway

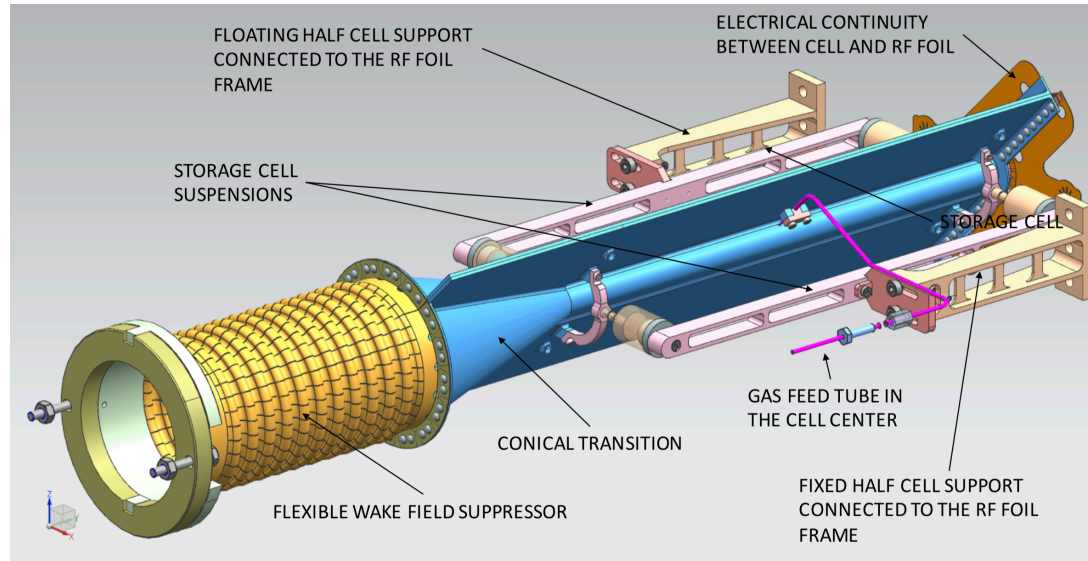


Trigger-less readout of full
40MHz LHC collision rate



Tracking up to at least 30% in PbPb collisions

SMOG II



1 yr of pAr data at 115 GeV:

Int. Lumi.		80 pb ⁻¹
Sys.error of J/Ψ xsection		~3%
J/Ψ yield		28 M
D^0 yield		280 M
Λ_c yield		2.8 M
Ψ' yield		280 k
$\Upsilon(1S)$ yield		24 k
$DY \mu^+ \mu^-$ yield		24 k

<https://cds.cern.ch/record/2673690/files/LHCb-TDR-020.pdf>

+ flow in various small systems?
Net proton fluctuations?

SMOG II can potentially run simultaneously with pp physics
Factor of ~100 increase in rates over current SMOG
Installation is happening NOW

Summary

- LHCb has an active, growing program in heavy ion physics, in fixed target and collider modes.
- The LHCb detector has unique capabilities for the study of exotics:
 - Full particle ID
 - Access to low p_T
 - Prompt production of exotics is new, largely unexplored territory
- Significant upgrades are well underway that will significantly expand LHCb reach in heavy ions
 - Full software reconstruction and triggering
 - Increased reach in centrality
 - SMOG II



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